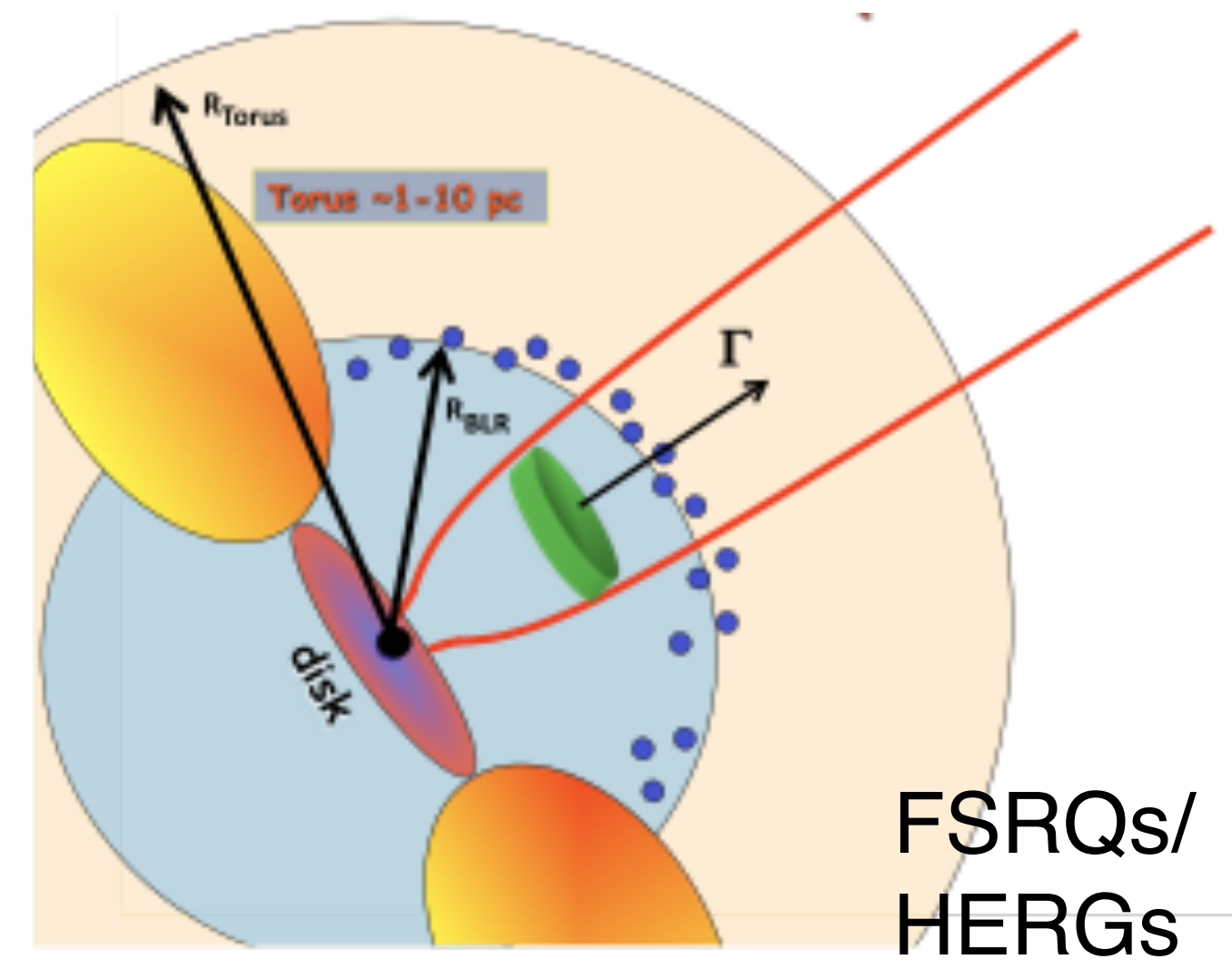
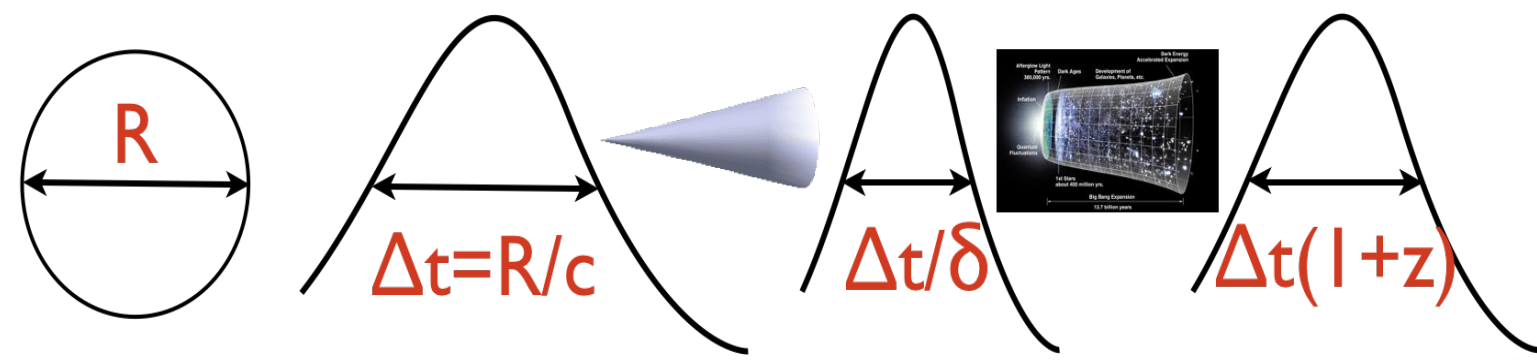
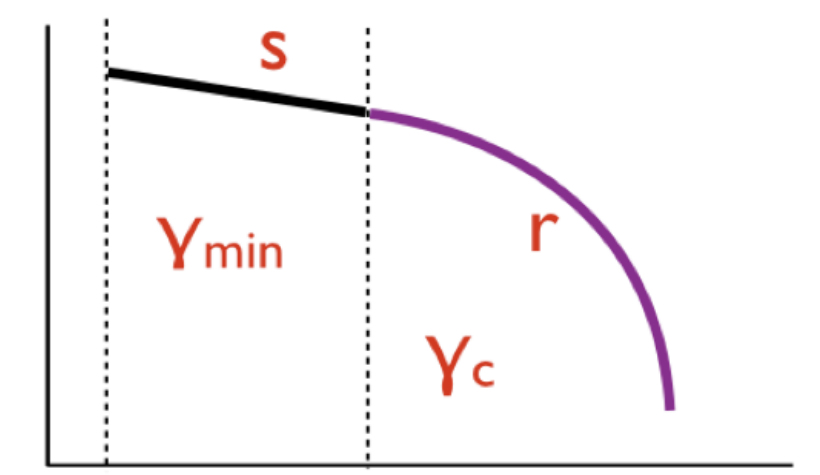
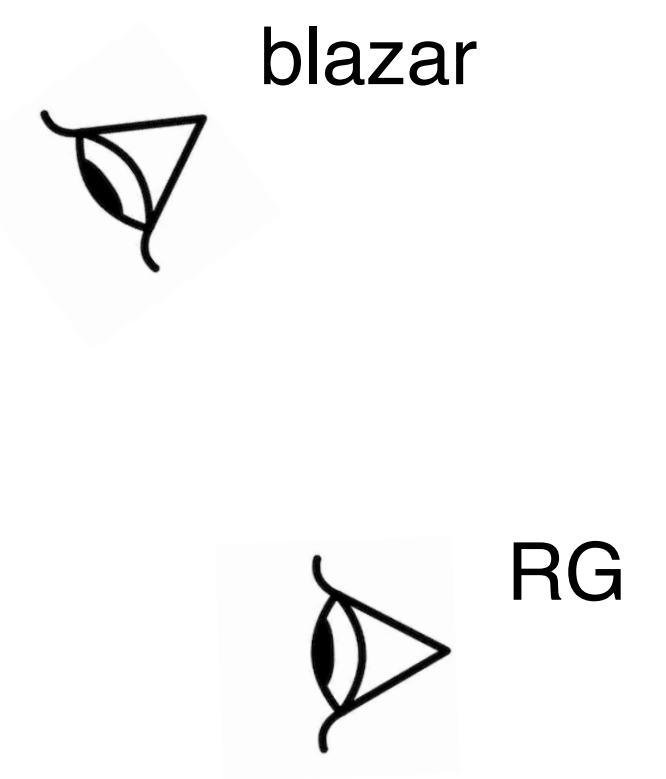
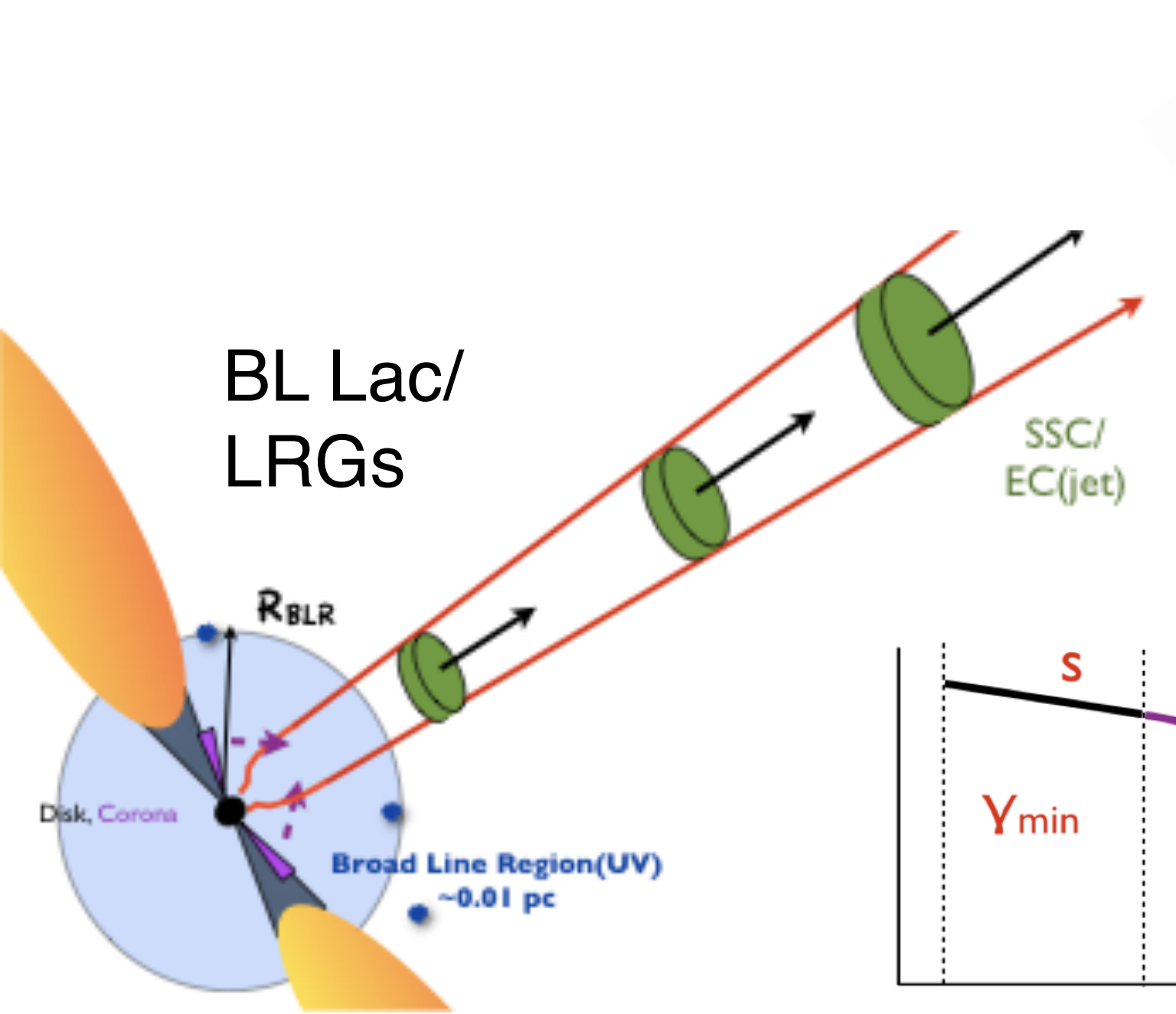


**Broad band modelling of Blazars:
a tool to understand the physics of relativistic jets.**

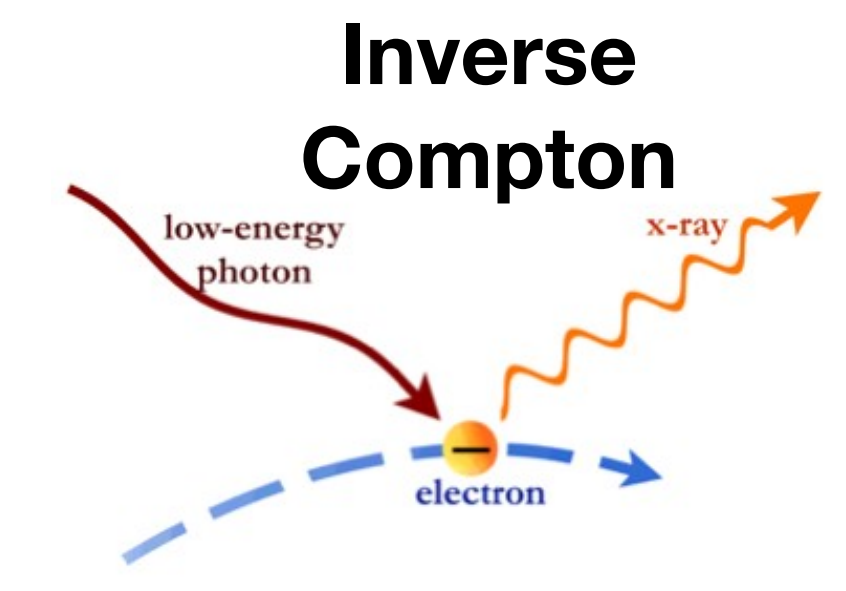
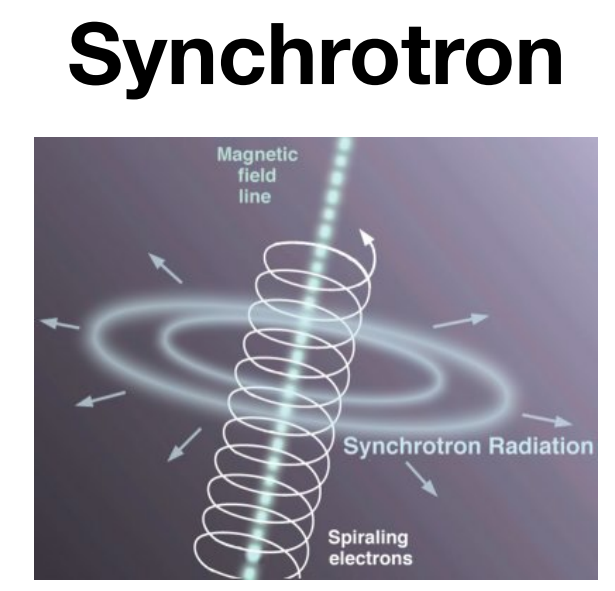
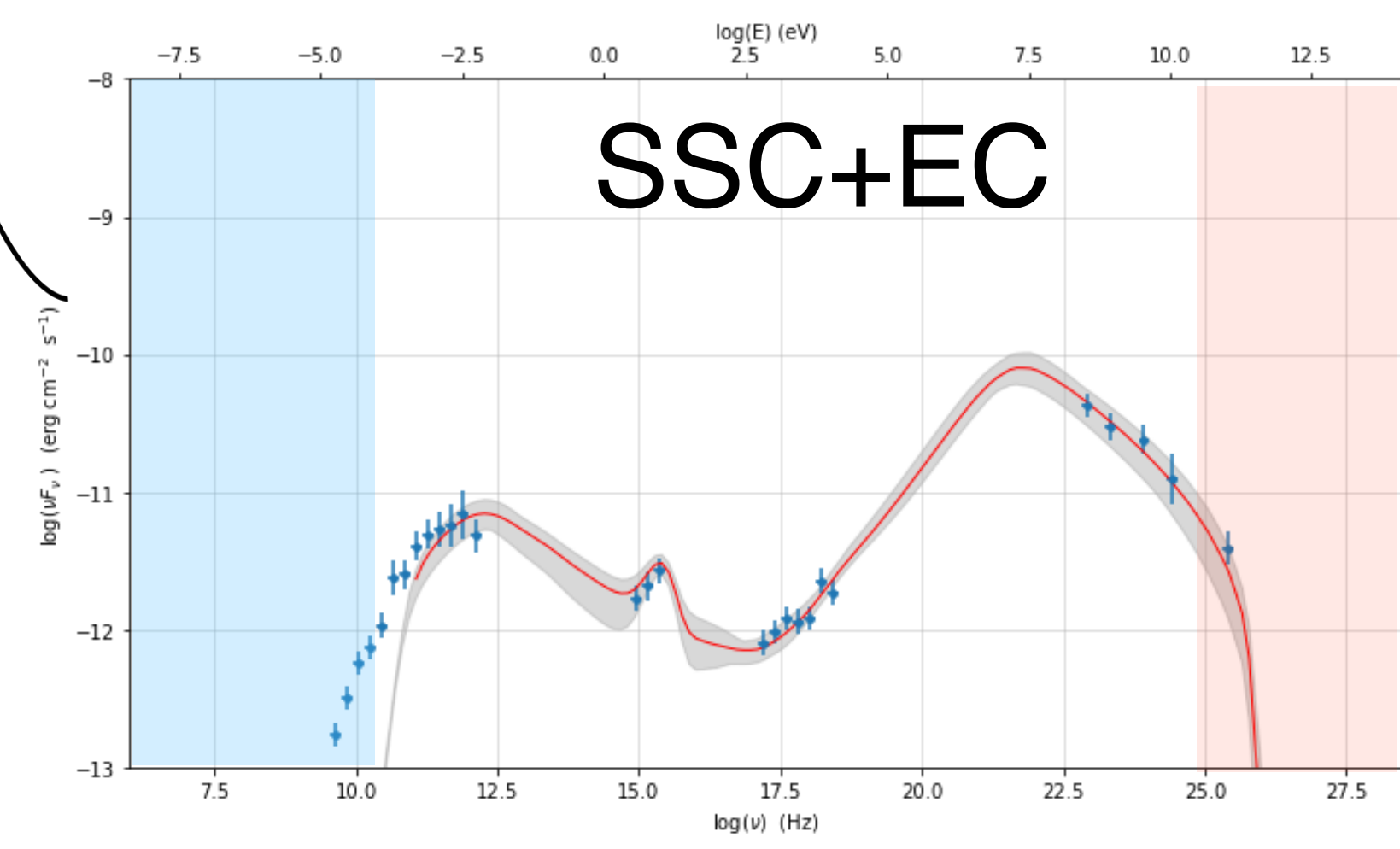
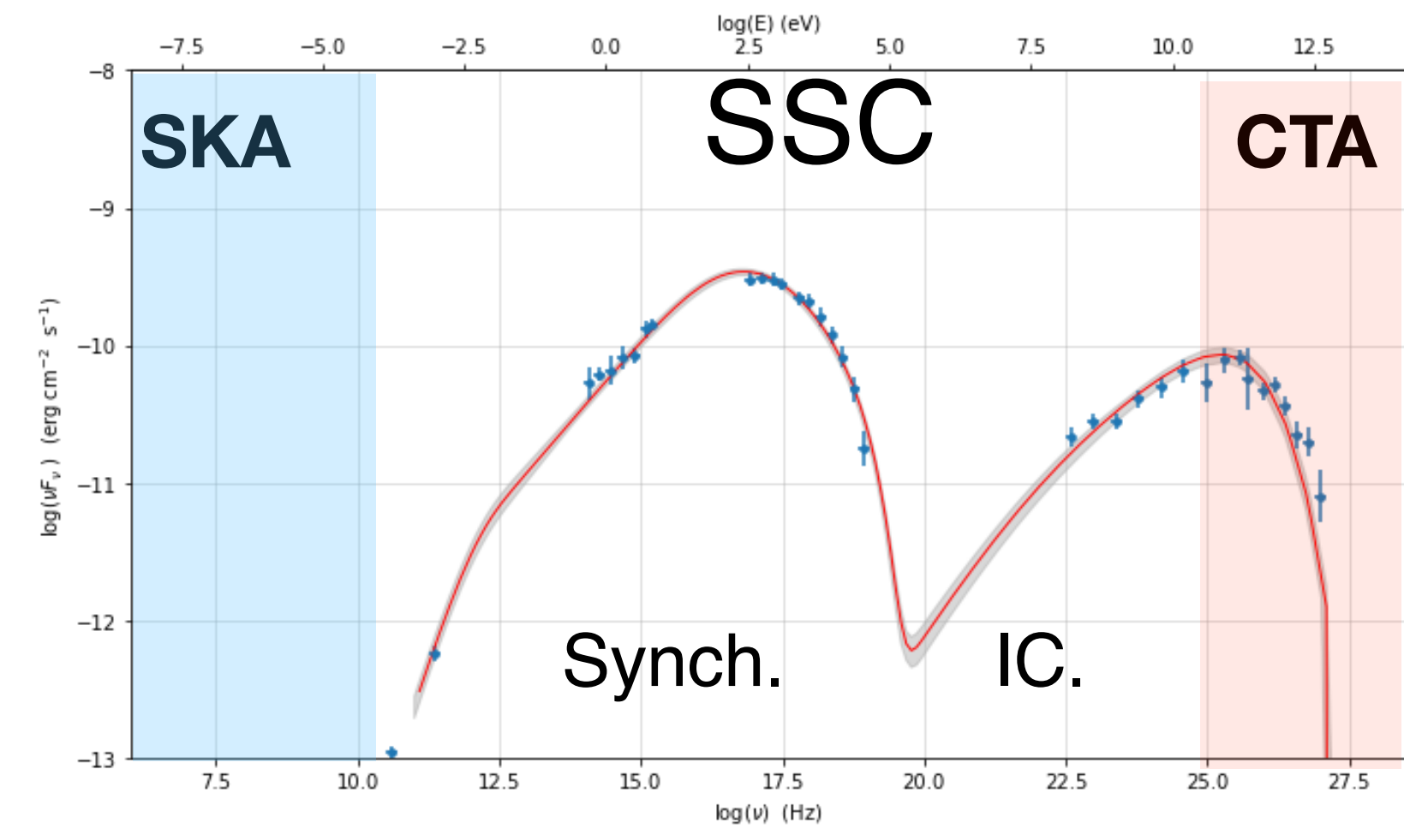
Andrea Tramacere

astronomy department university of Geneva

- current status of the modeling
- phenomenological signatures hinting for complex scenarios
- self-consistent modeling and degeneracy
- connection of micro to macro physics



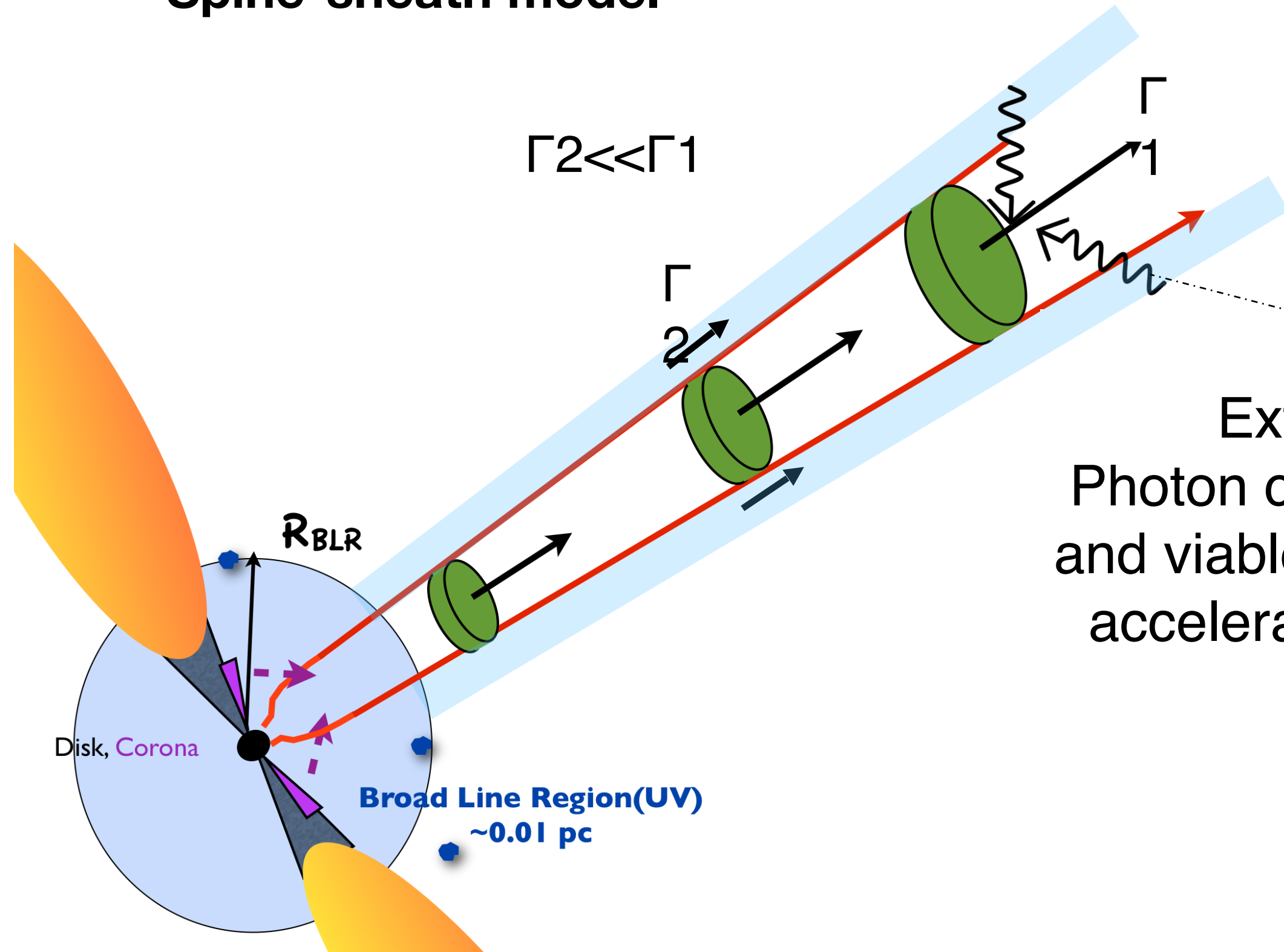
- rel. jets are powered by accr. on BH (10^6-9M_{sun})
- rel. plasma ($\Gamma > 10$) acc. by non-thermal processes
- jet axis vs l.o.s. \rightarrow beaming (blazar/RG)
- variable emission over the entire EM spectrum
- Low energy bump S, high energy bump IC



FSRQs/
HERGs

Hadronic models

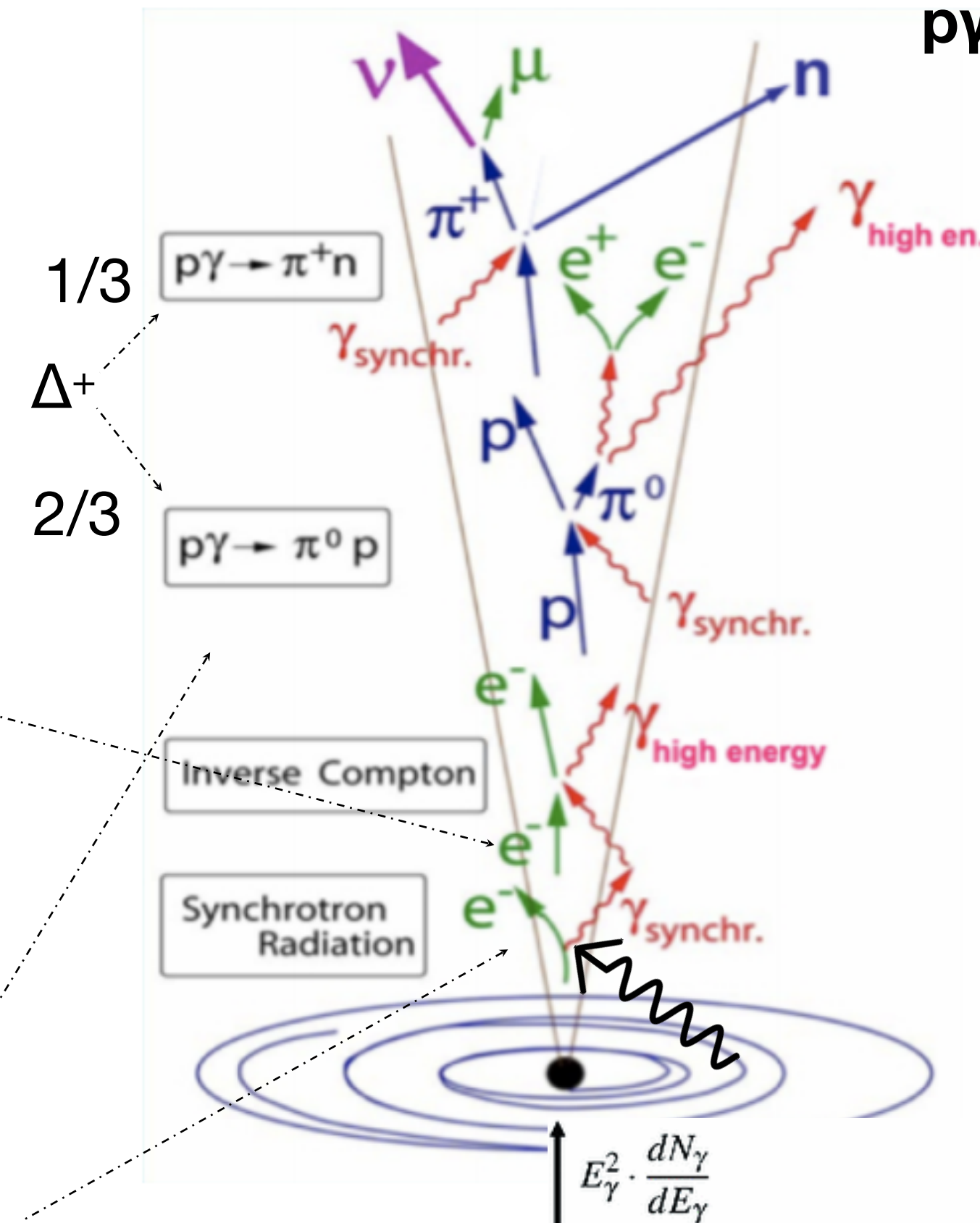
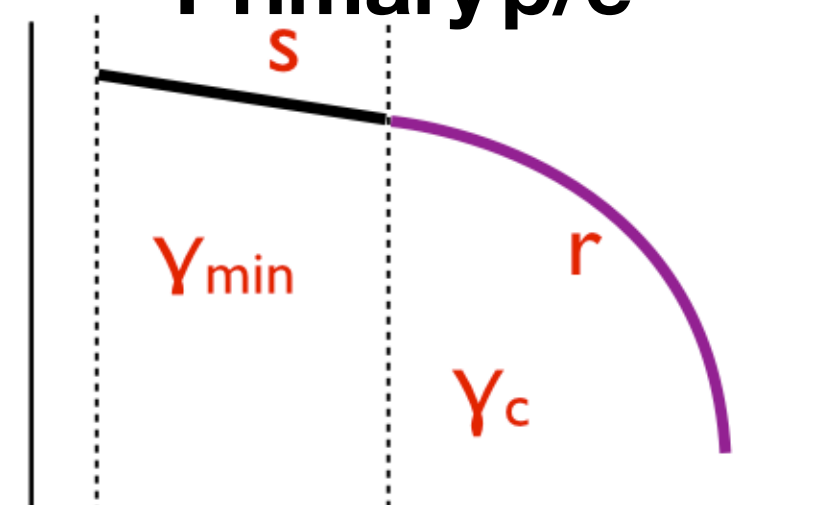
Structured jets Spine-sheath model



CTA/SKA/IceCube strong synergy

External Photon contribution, and viable process to accelerate protons

Primary p/e-

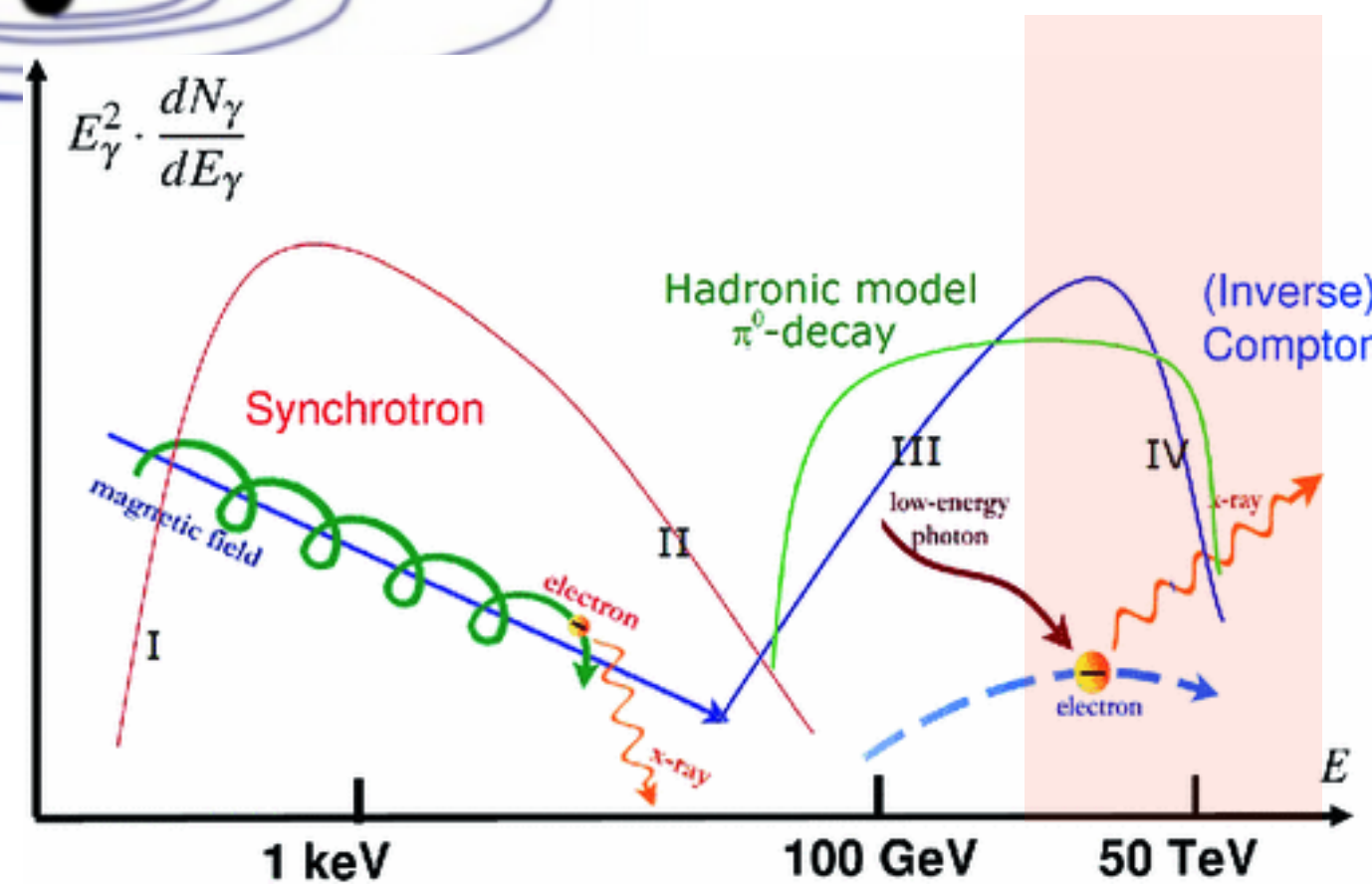


pγ, p-synch

$$E_p E_\gamma \sim E_{\Delta^+}^2$$

$$E_\nu \sim E_p / 20$$

$$E_p \sim 10^{17} \text{ (eV/E}_\gamma\text{)}$$

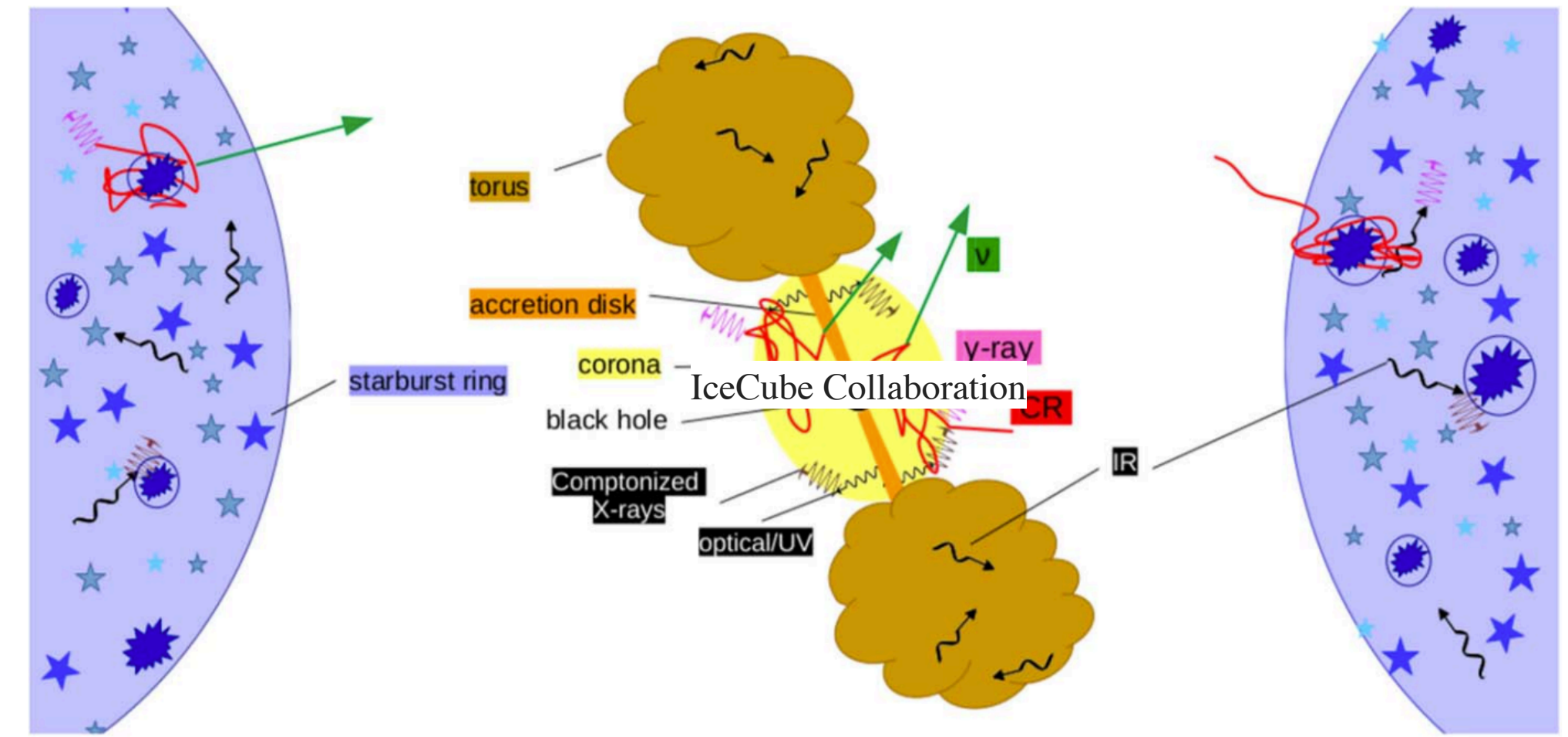


more complex mode with pp corona/outflows

can explain IceCube results for Sy2 NGC 1068

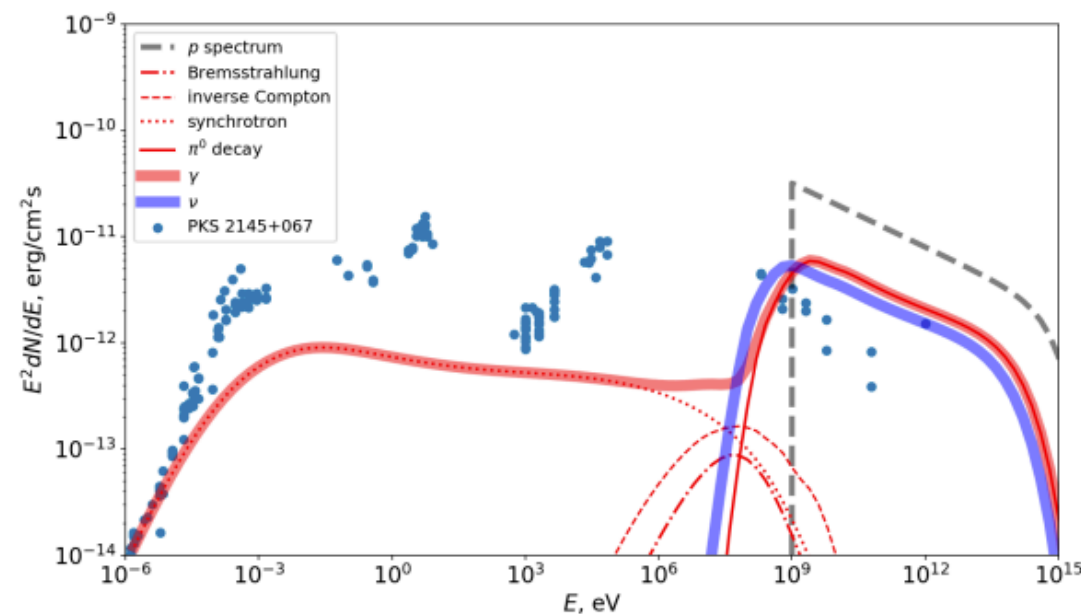
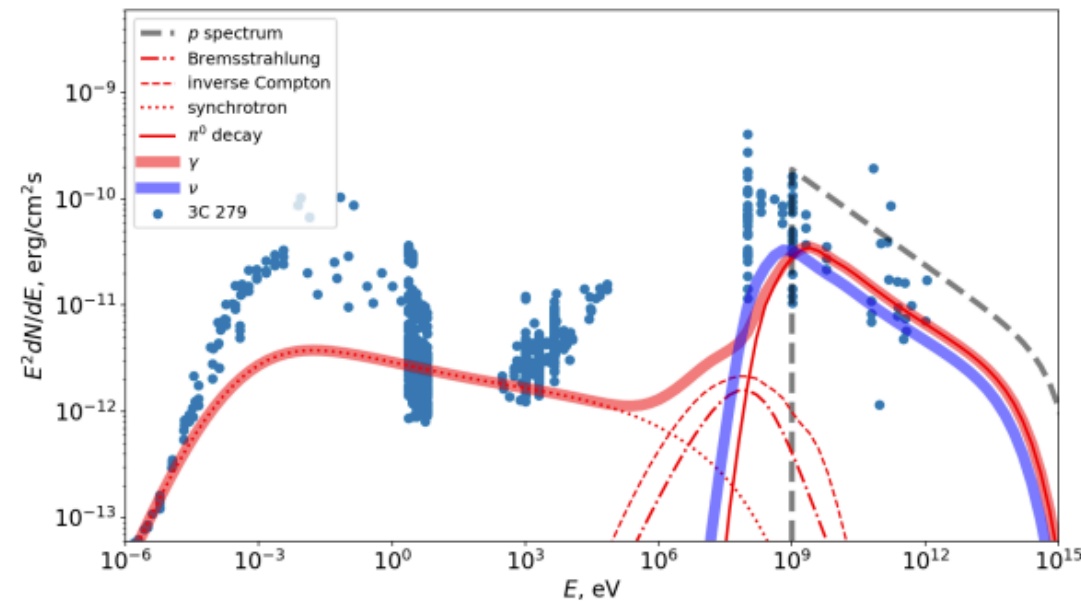
but also γ -NLS1/CSS have a corona and possible signatures of winds/outflows, hence such a model might produce neutrinos in the gamma-ray opaque regions of jetted AGNs

see Teresa Talk/ IceCube results 2022

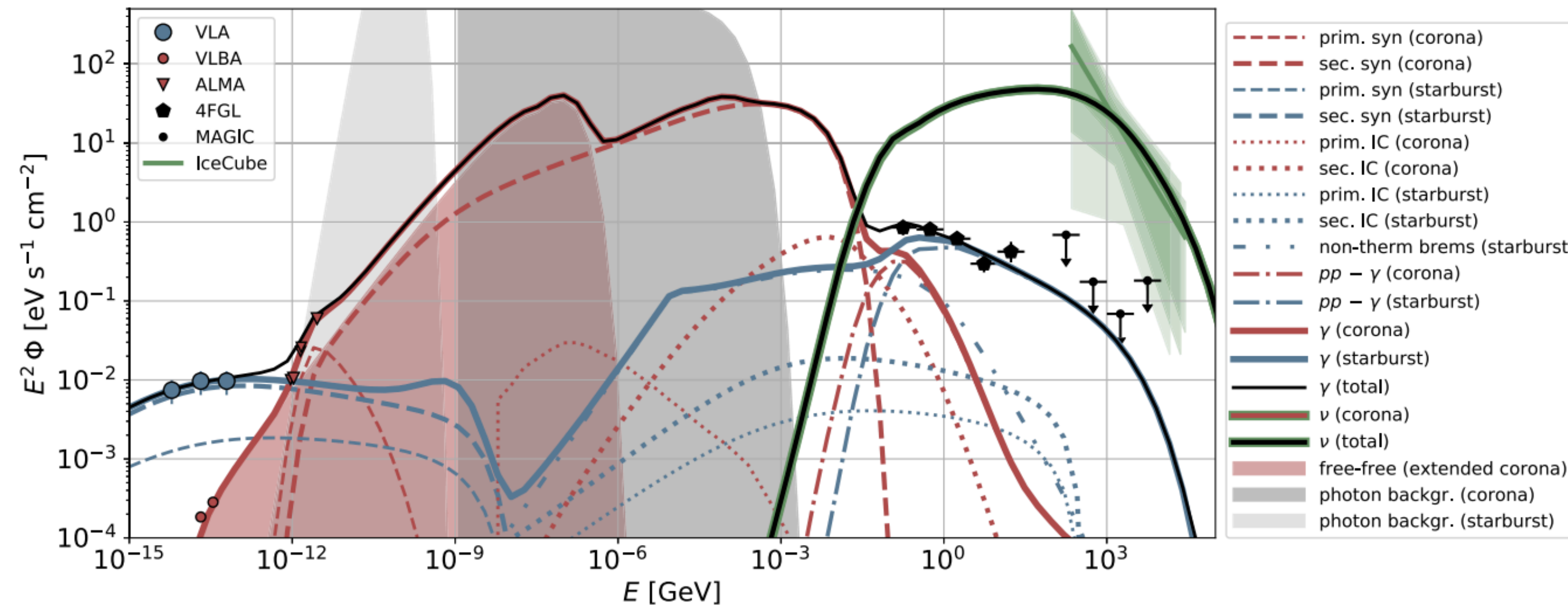


Neronov&Semikoz 2020

for pp applied to pc scale of relativistic jets propagating through circumnuclear medium of the AGN

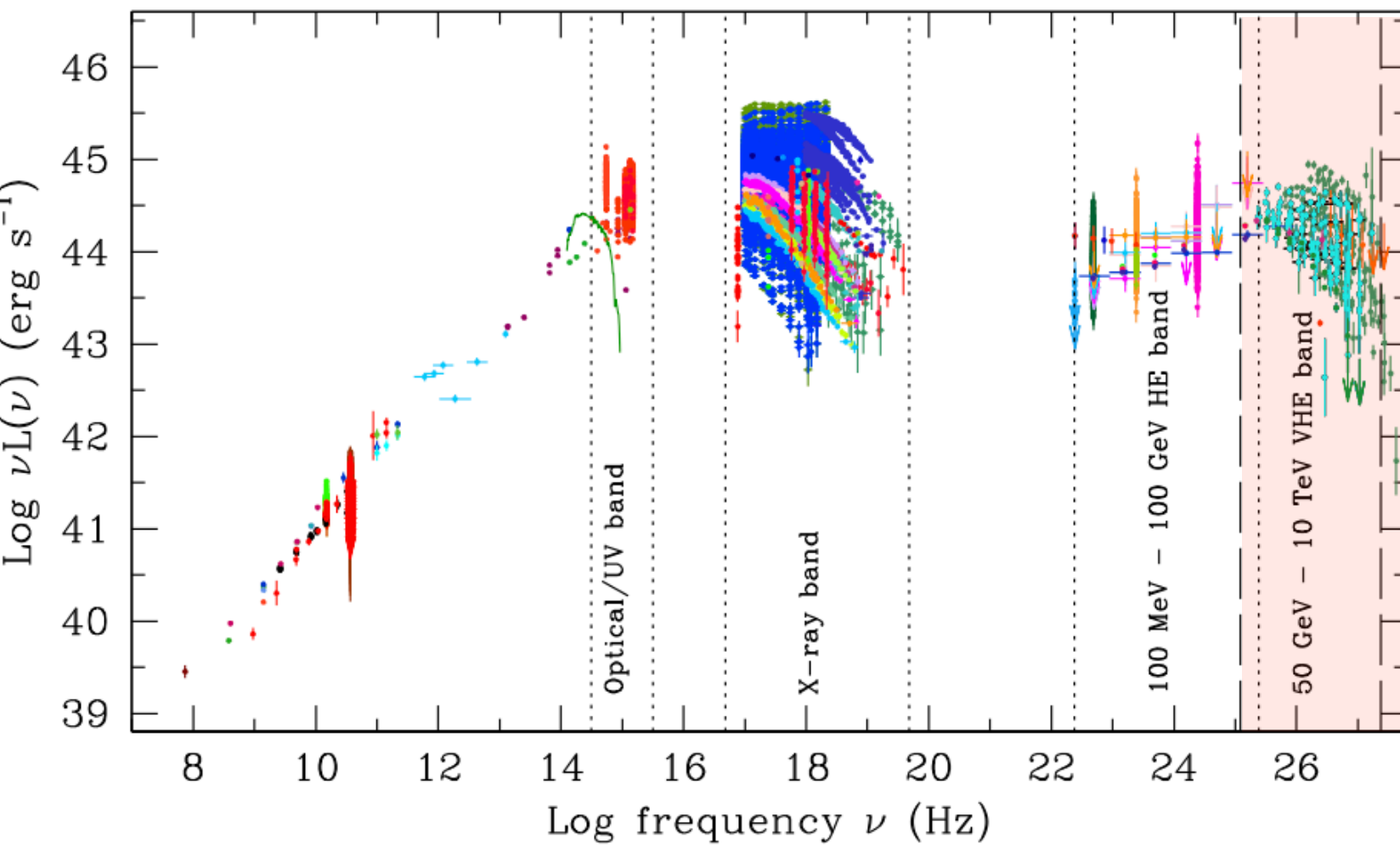


Eichmann+ 2022, see also Murase papers

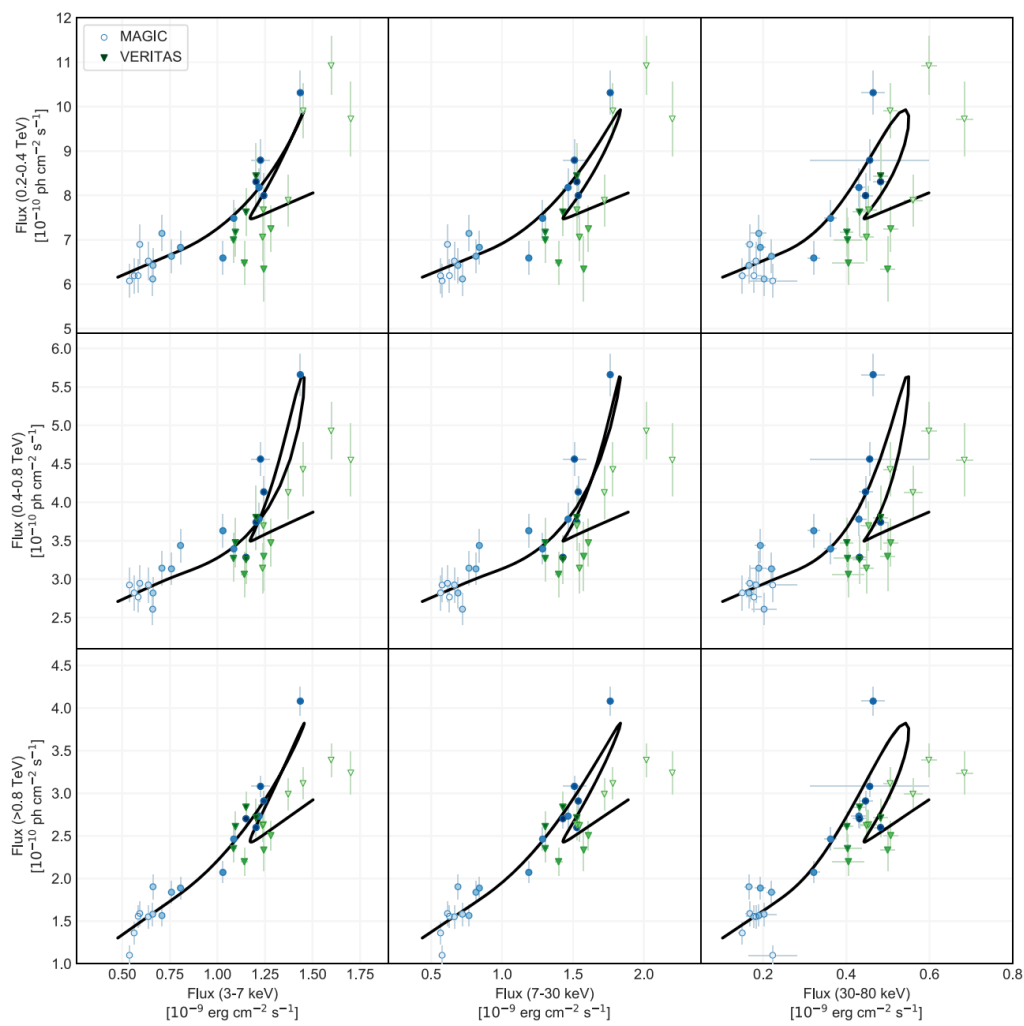


Padovani+ 2017

Mrk 421



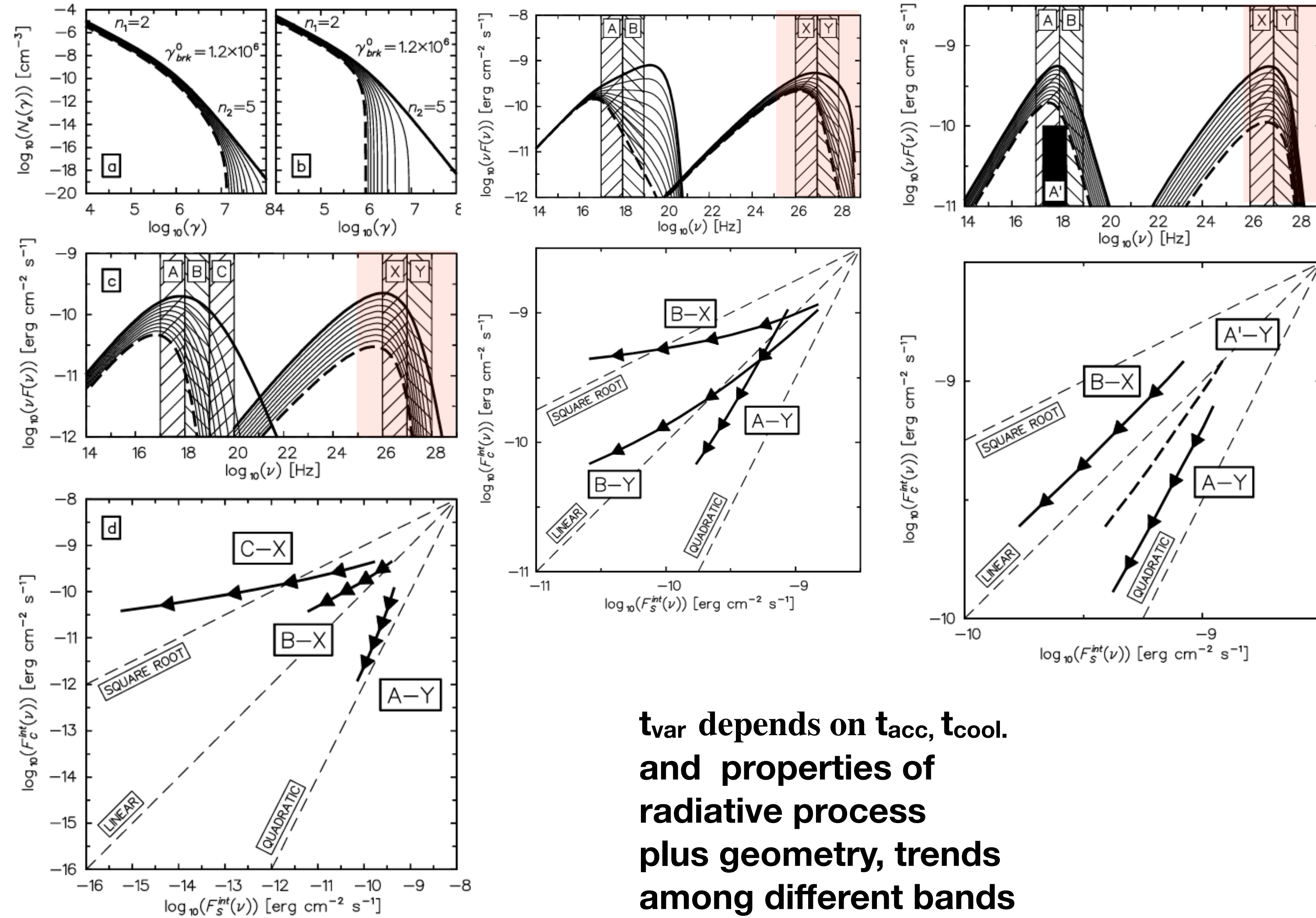
Mrk 421 2013 Acciari+ 2020



CW: soft lag, when the flare is observed at frequencies where the higher energy variability occurs more rapidly than at lower energy

CCW: In contrast, when observed at frequencies for which the acceleration and cooling timescale are almost equal, the loops are expected to be counter-clockwise with a possible hard lag

single episode, single process Katarzyński+ 2005

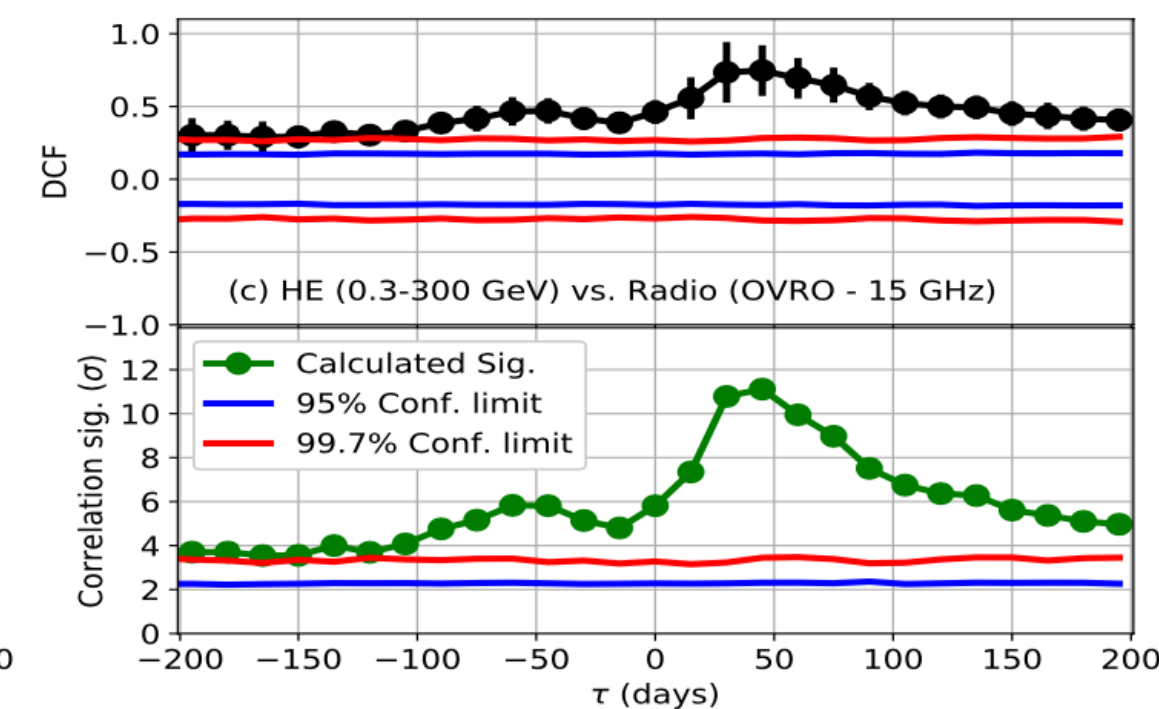
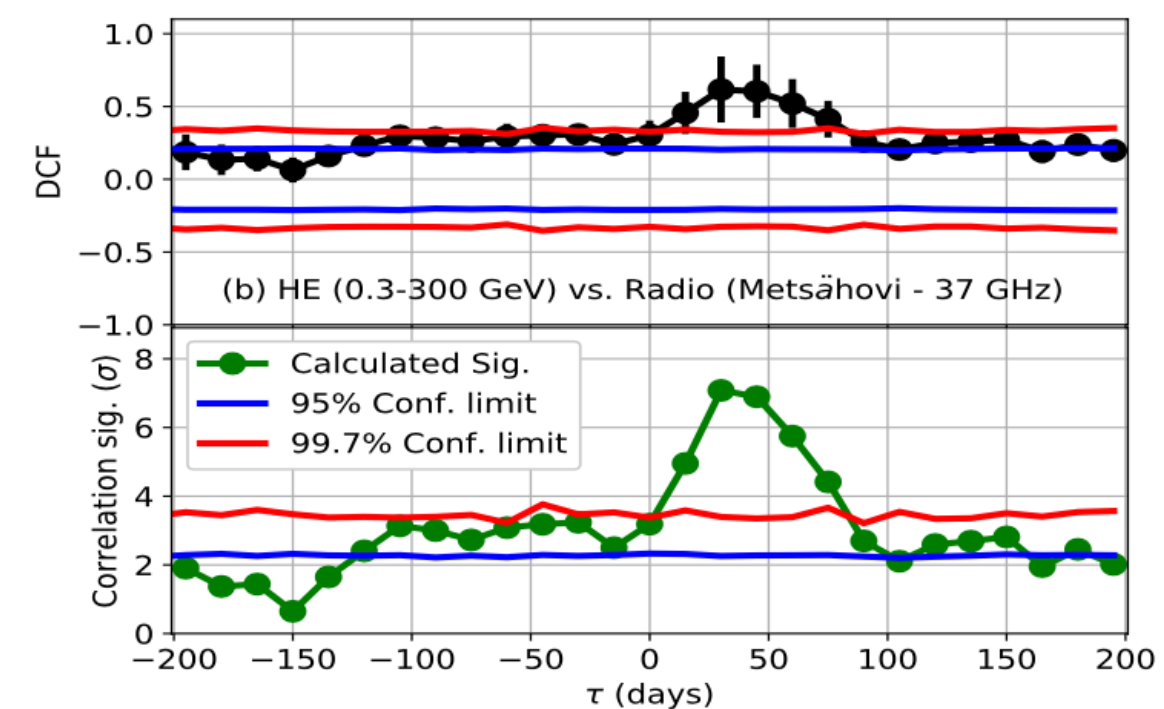
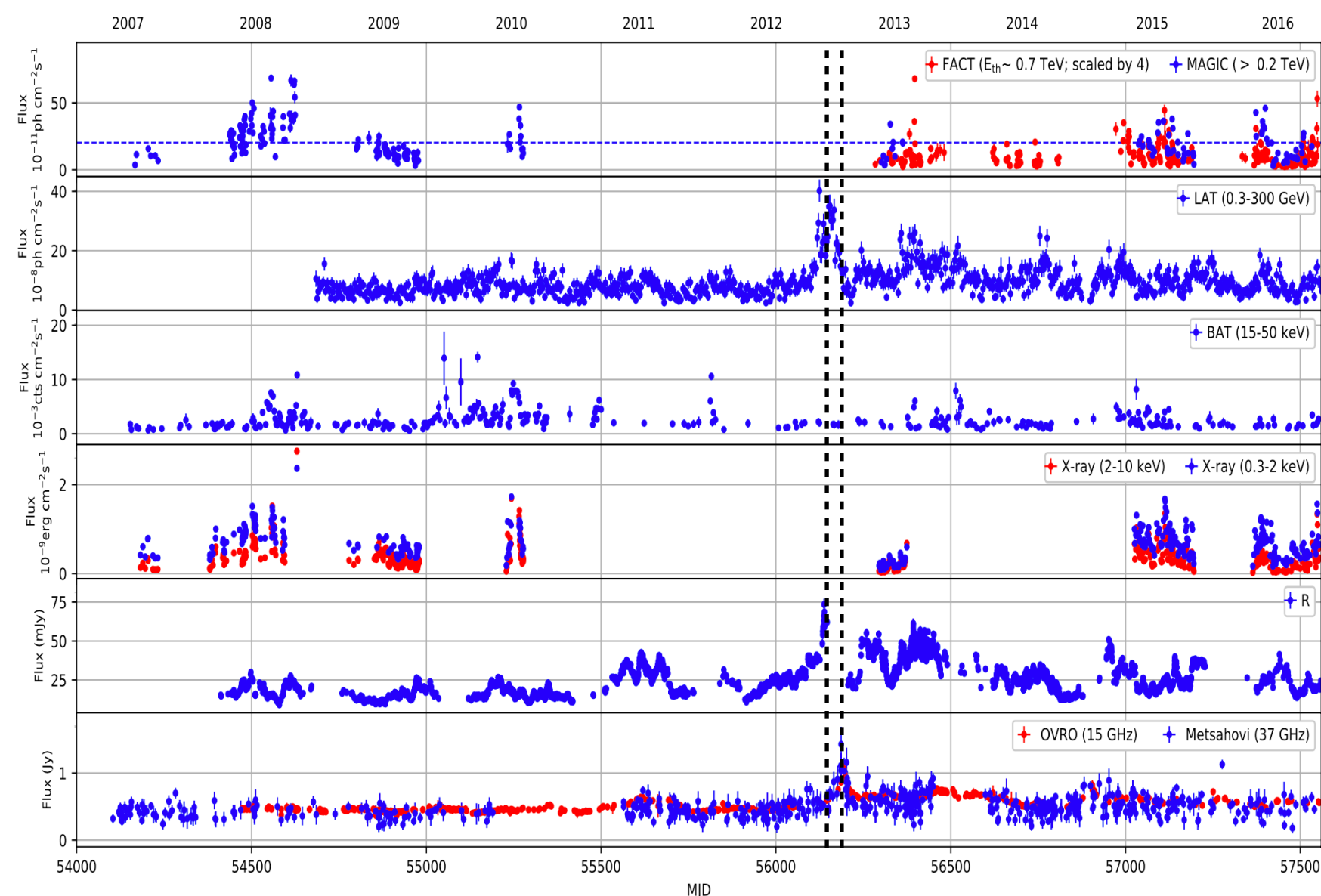


t_{var} depends on t_{acc} , t_{cool} .
and properties of radiative process plus geometry, trends among different bands bring signatures of different processes.

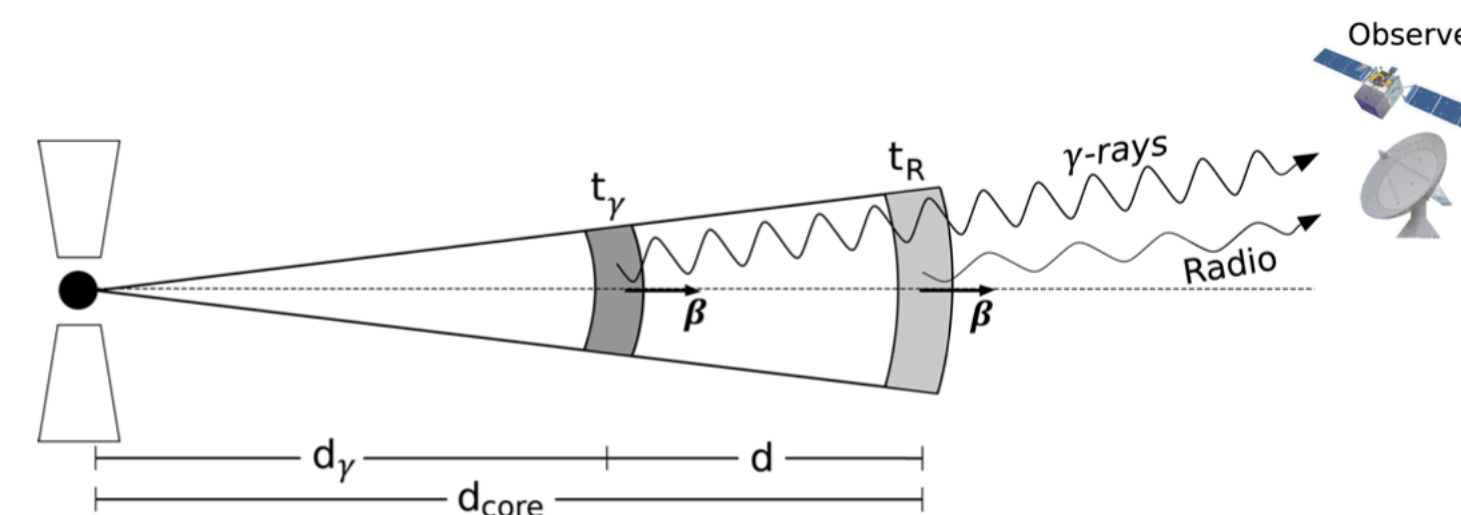
MW variability and correlation studies of Mrk 421 during historically low X-ray and γ -ray activity in 2015-2016

Radio- γ delay in Mrk 421 (months)

Magic coll. 2020



Radio- γ delay $\sim 1/\nu$



W. Max-Moerbeck+ 2014
 B. Pushkarev+ 2010
 Ghisellini+1985
 McCray,R. 1968

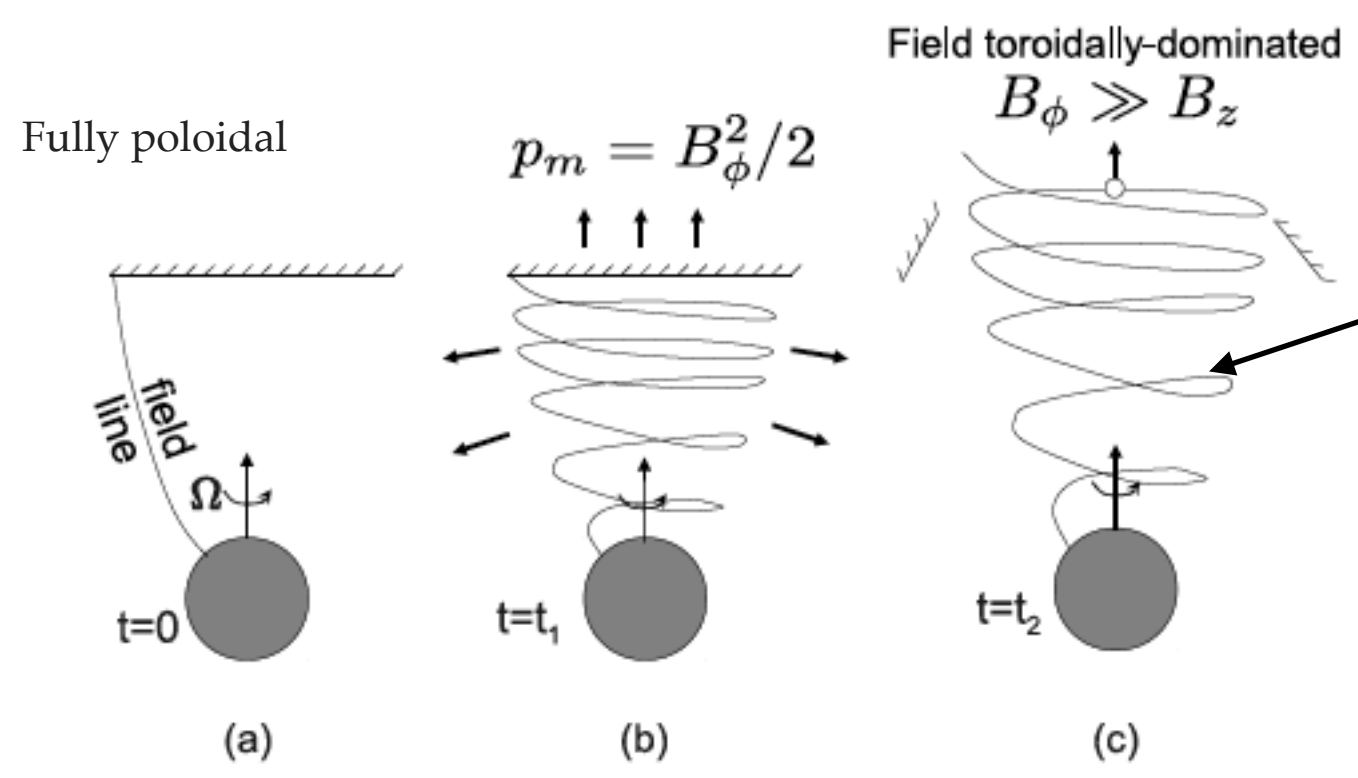
see Vitalii Talk

long terms depends on jet

feeding processes and jet structure:

AD instabilities ,BH spin, jet geometry

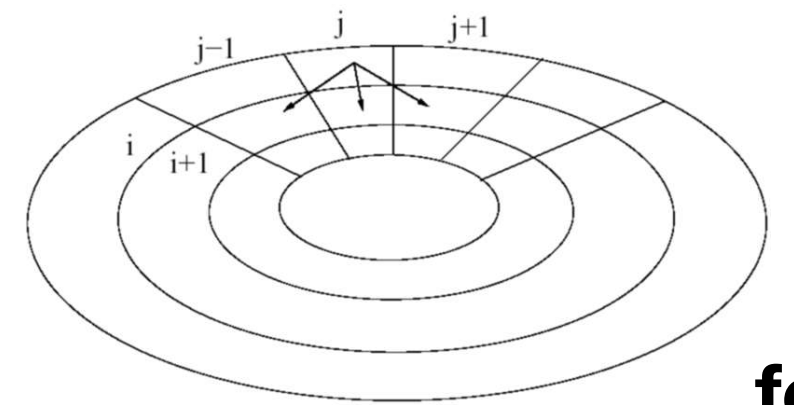
credit picture Sera Markov



BZ process
mechanism to work, the system requires two ingredients:
- a central spinning black hole and
- a surrounding disk of plasma with a strong poloidal magnetic field. Once accretion onto the black hole has begun, the magnetic field lines embedded in the disk, due to the frame dragging, twists the field lines into following the rotation of the black hole.

provides the B gradient to accelerate the jet

Frame dragging and differential rotation
disk variability



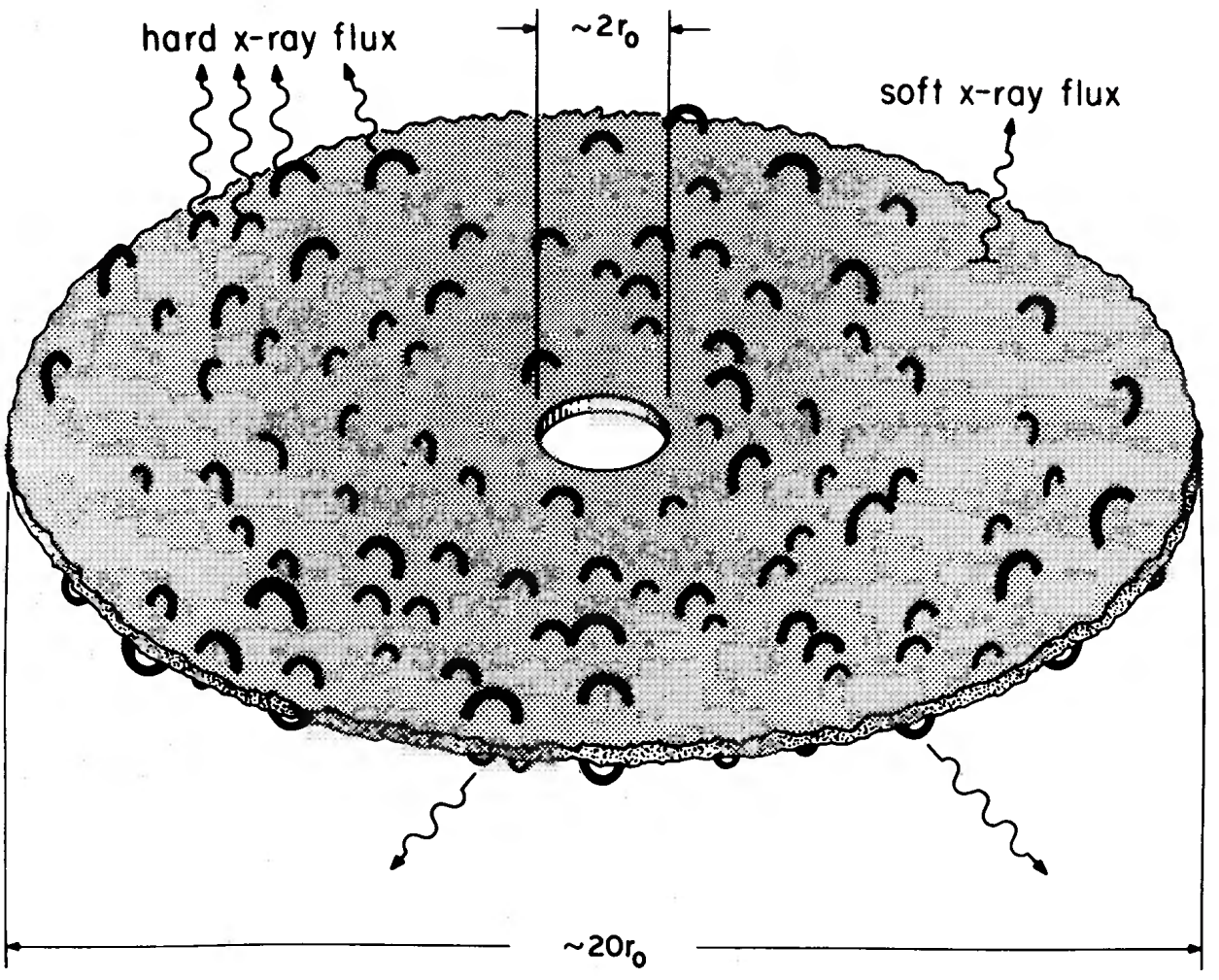
feeding

jet powering

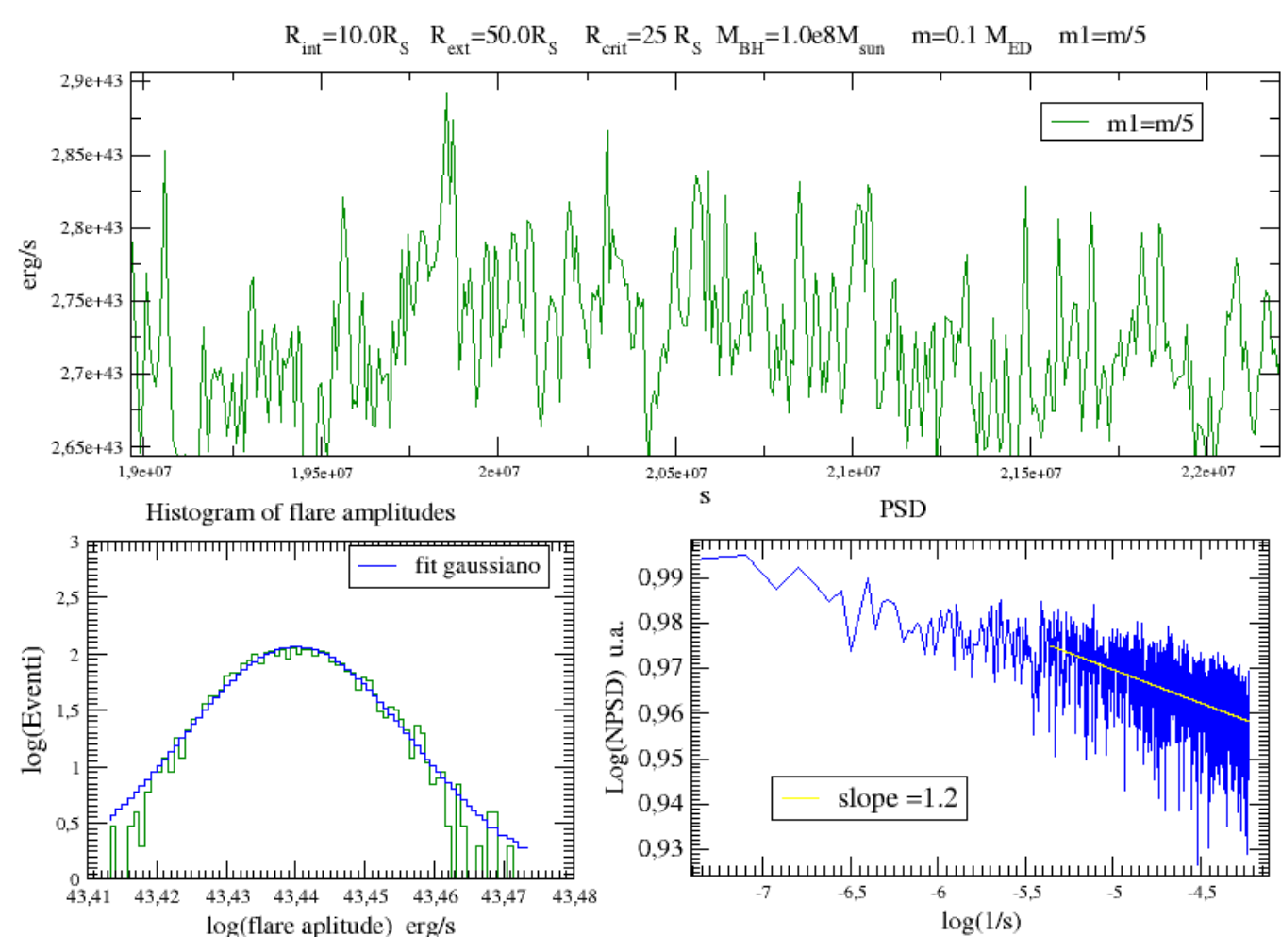
jet variability

provides content of pairs during magneti reconnection plus pre-acceleration of particles

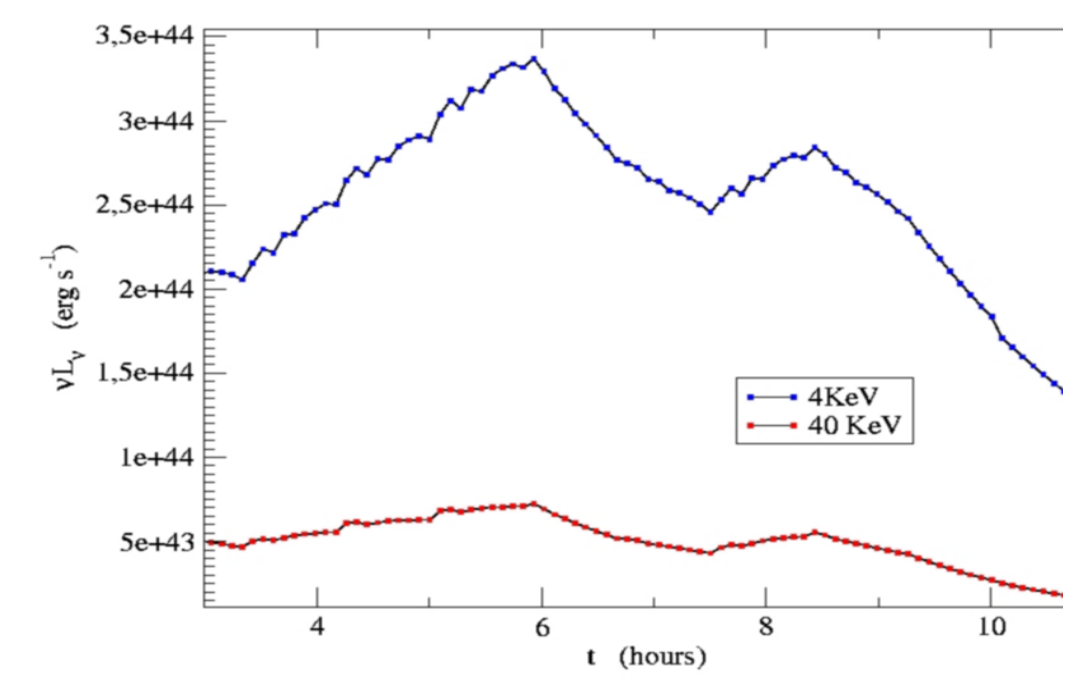
Disk, differential rotation, B&P wind plus plasmoids



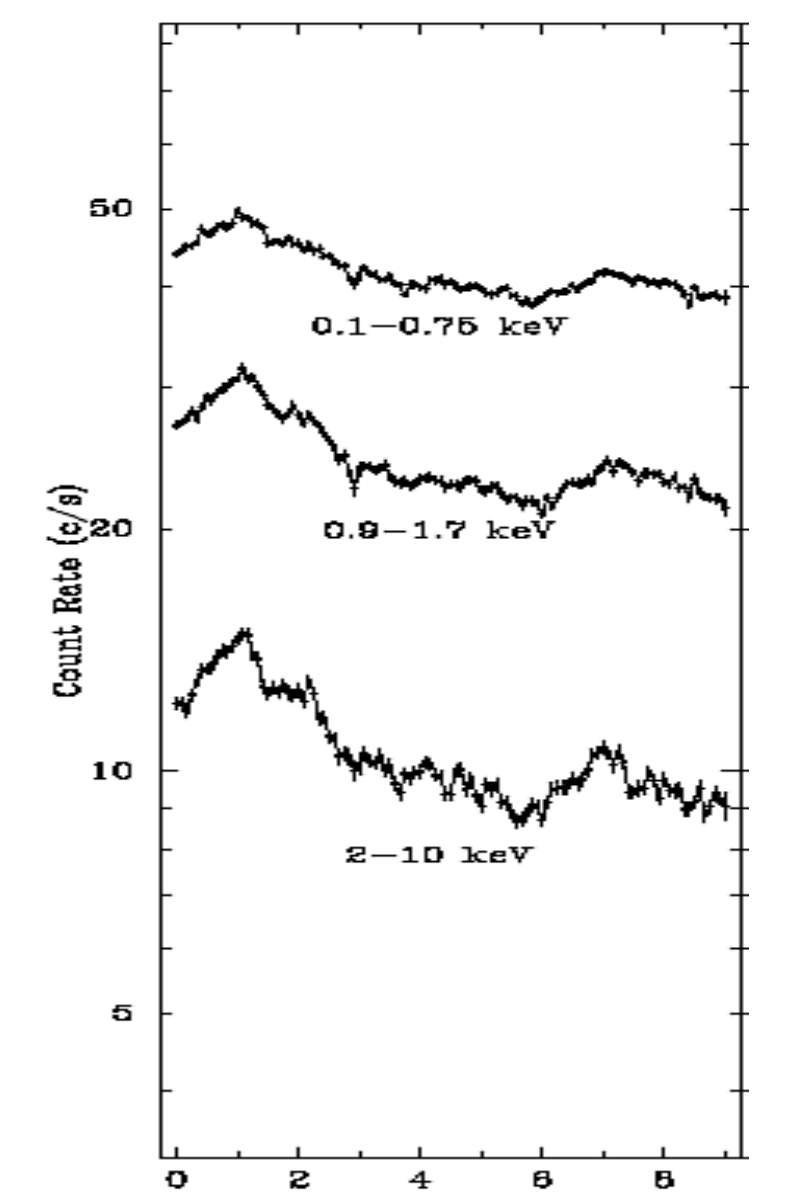
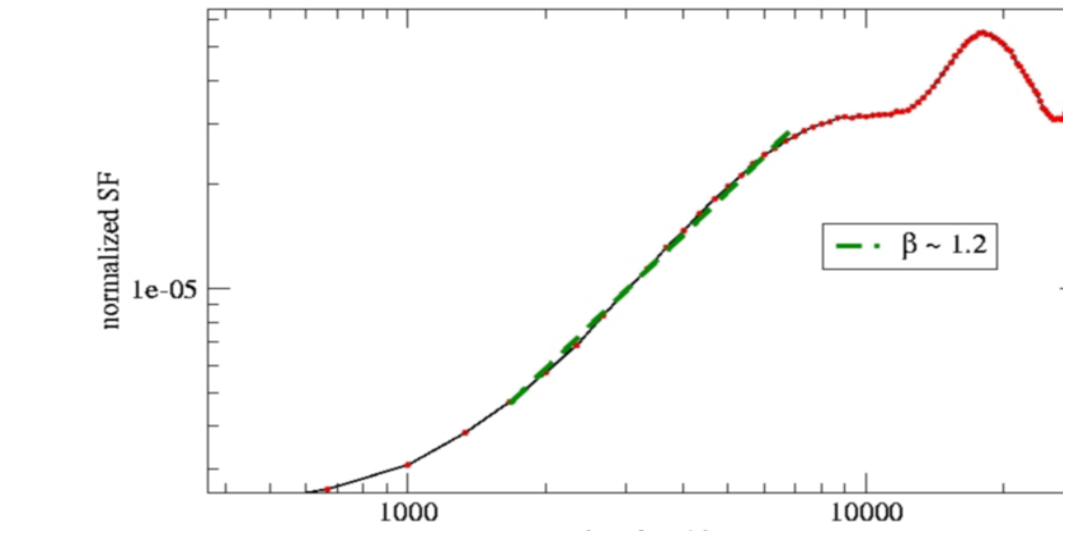
Galev&Rosen 1979



tramacere

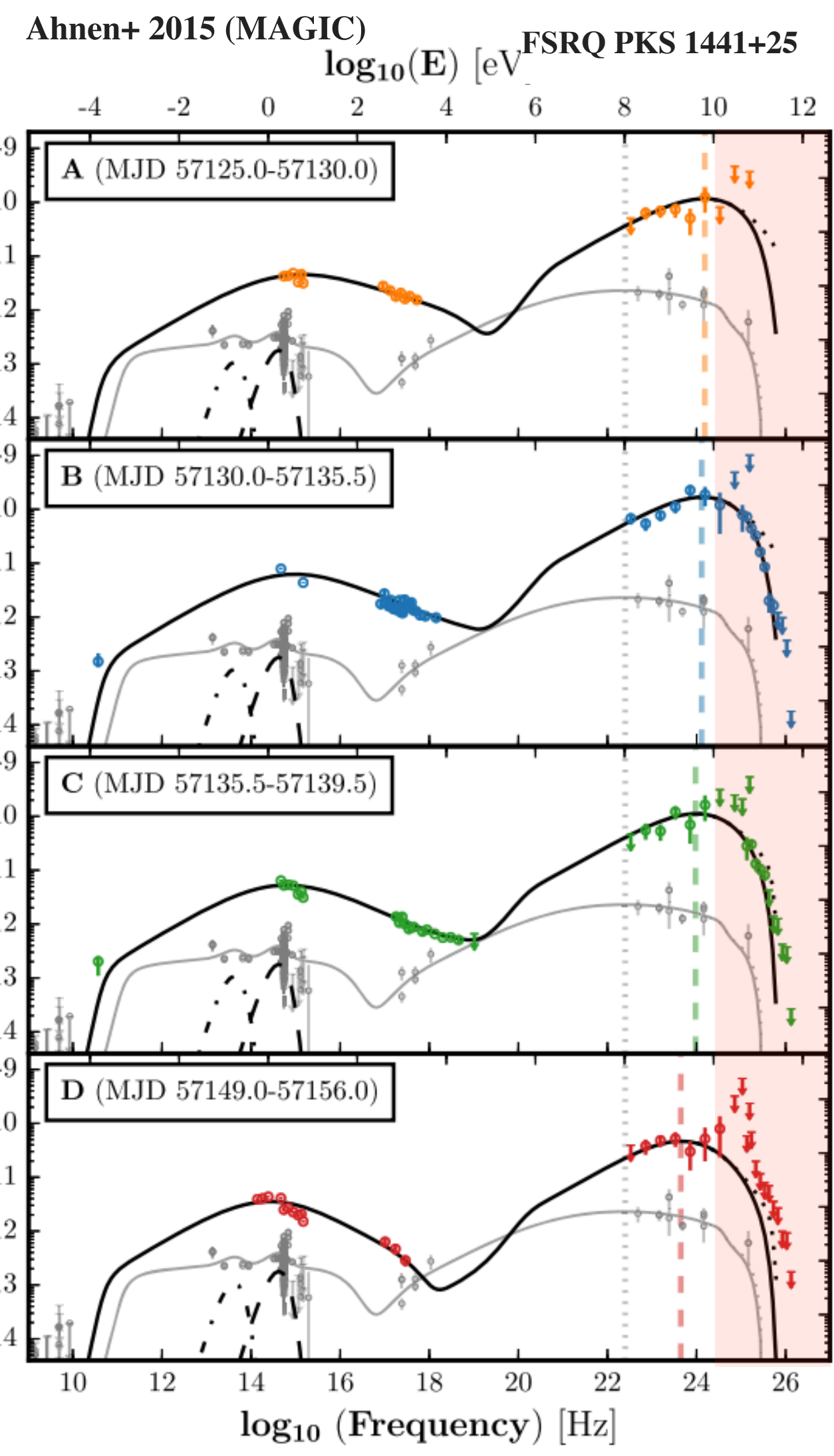


tramacere

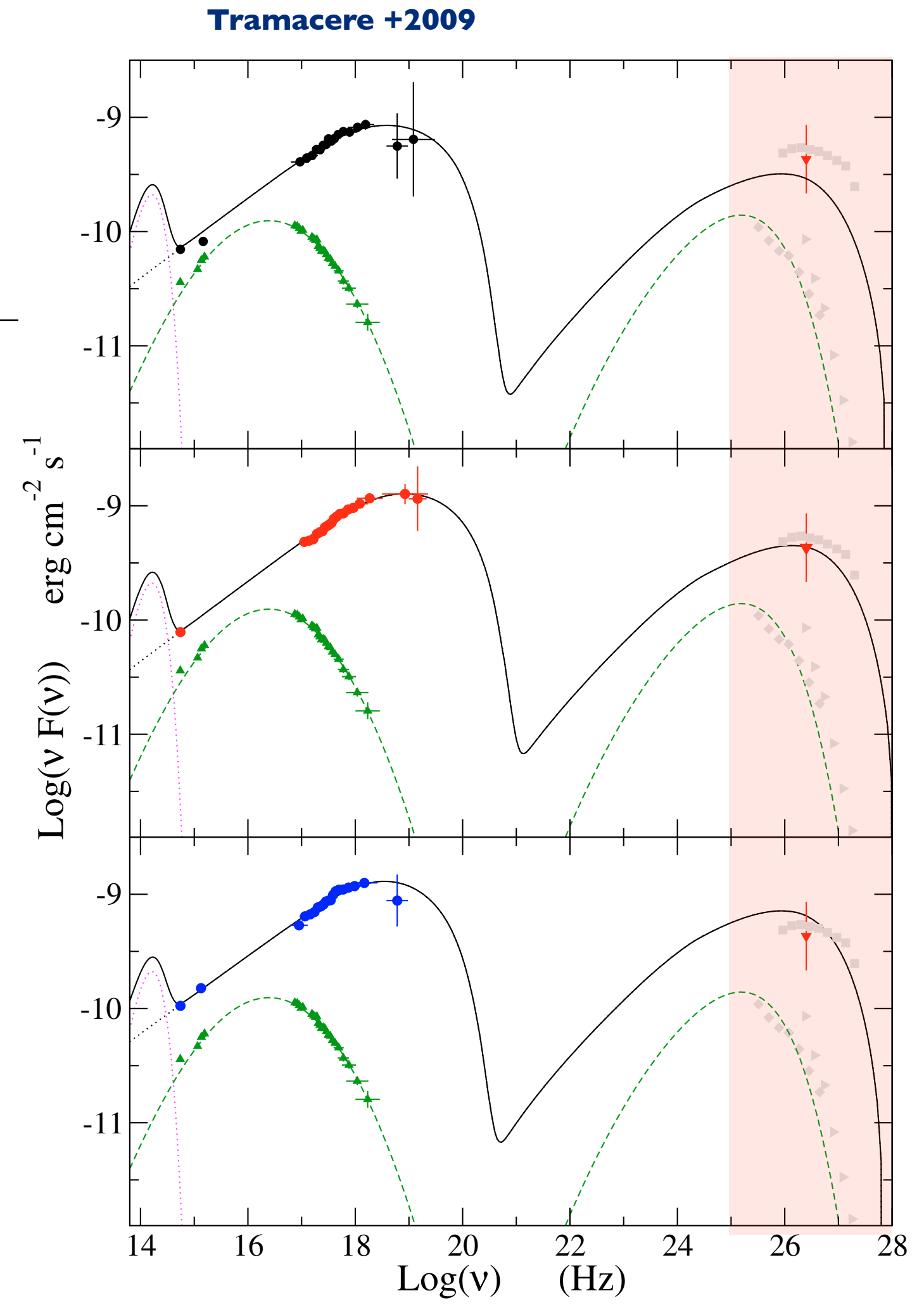
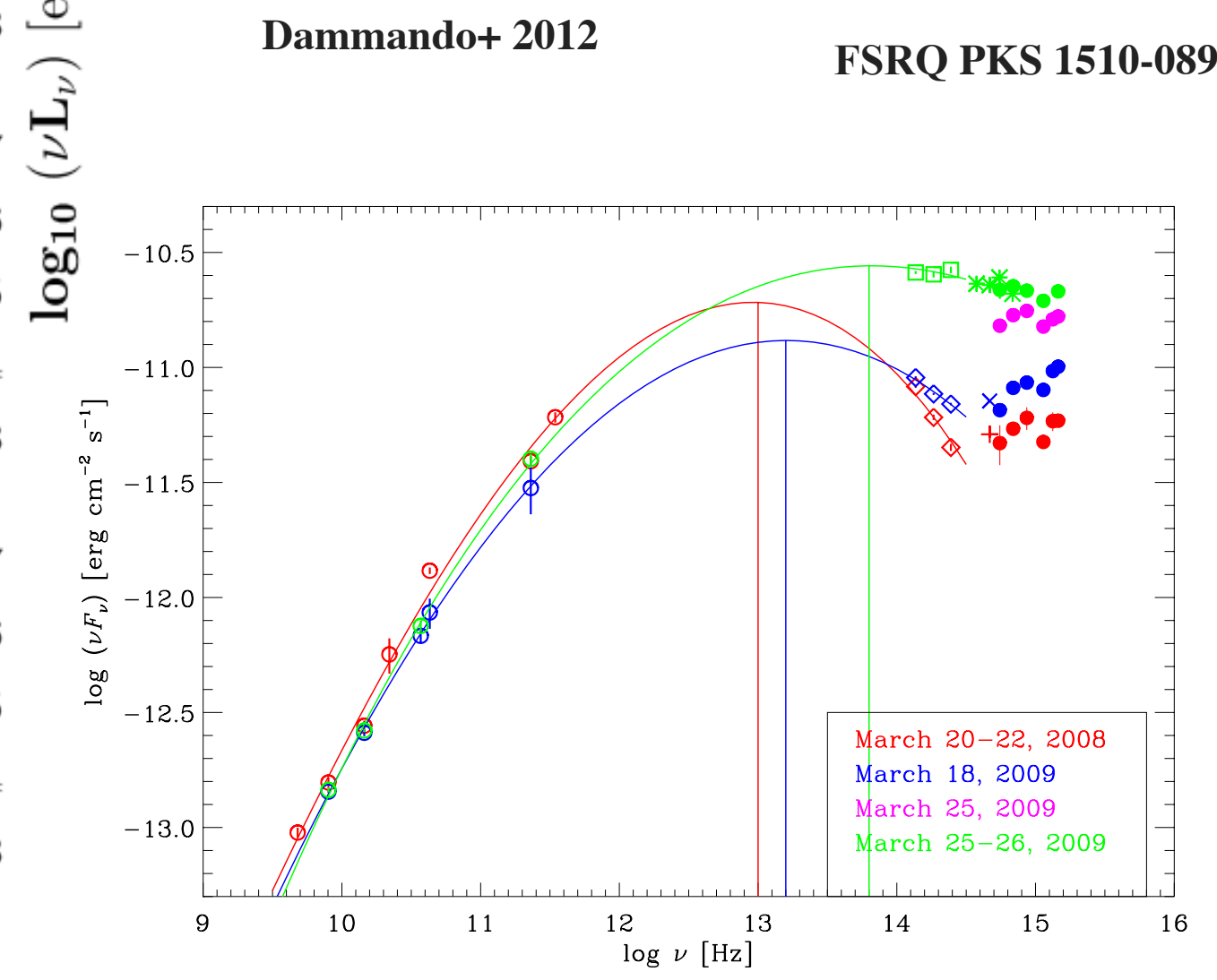
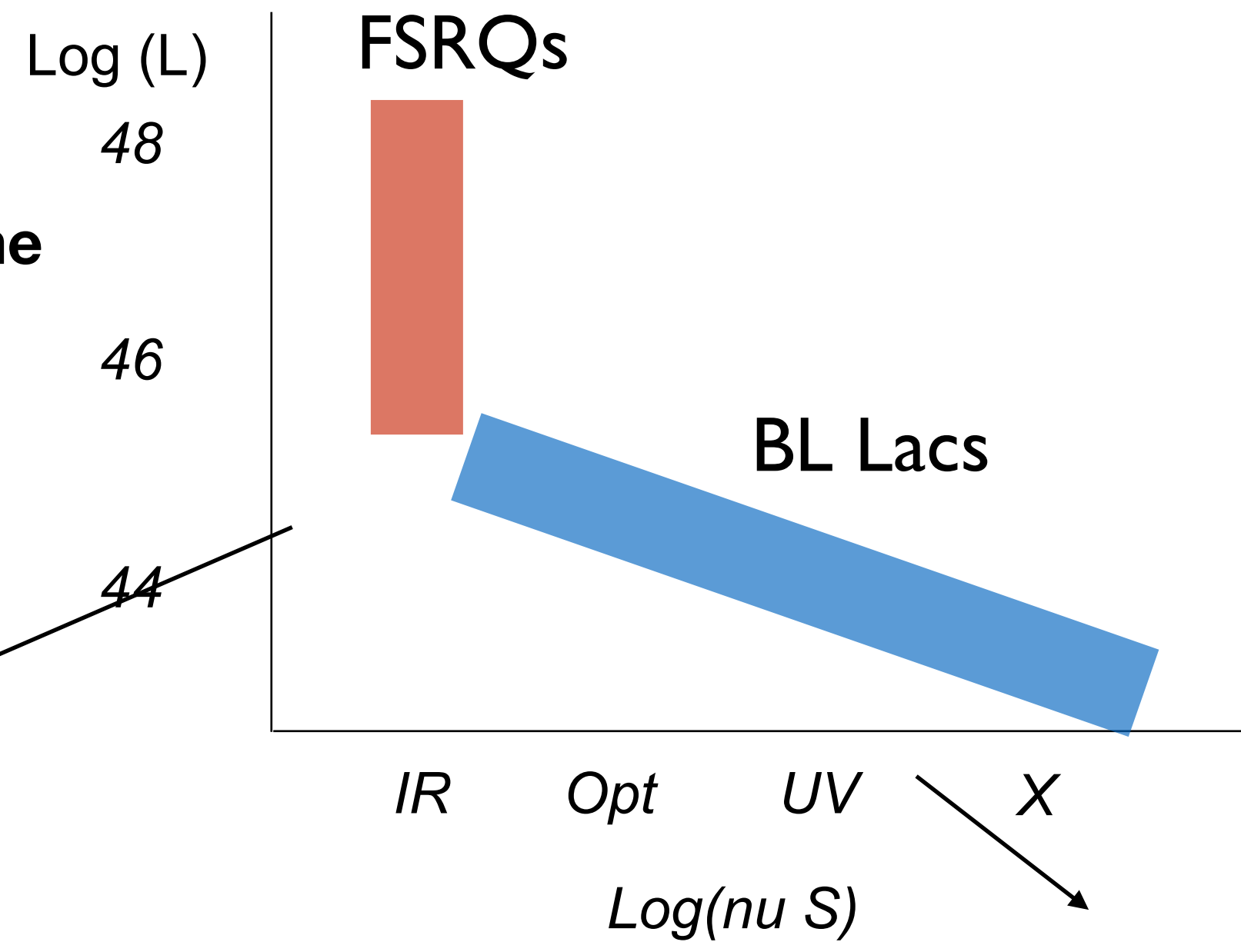


MRK 421 XMM-NEWTON SEMBAY et al

The peak frequency dynamic range, for LSP, is low, this might hide the emitters evolution

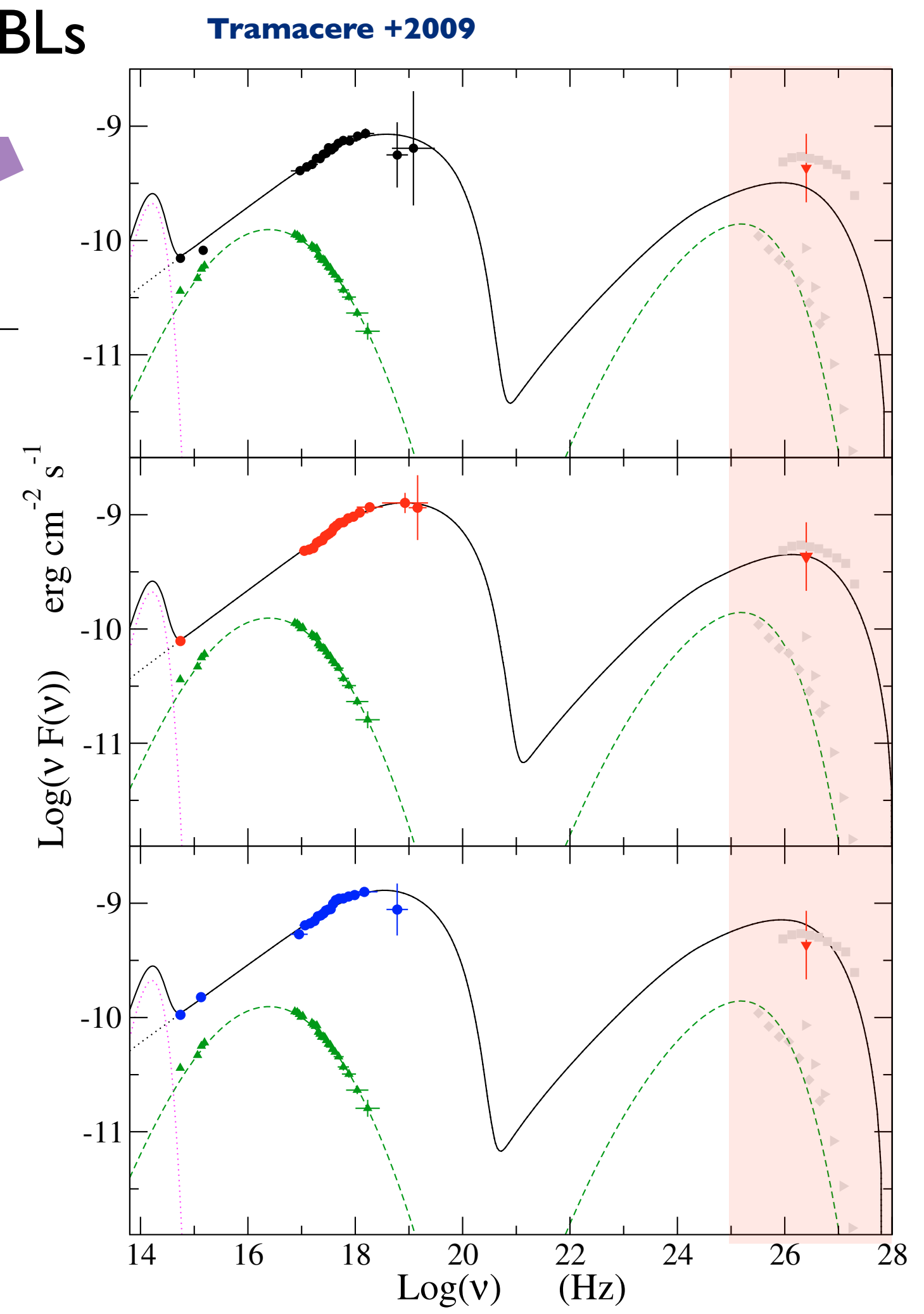
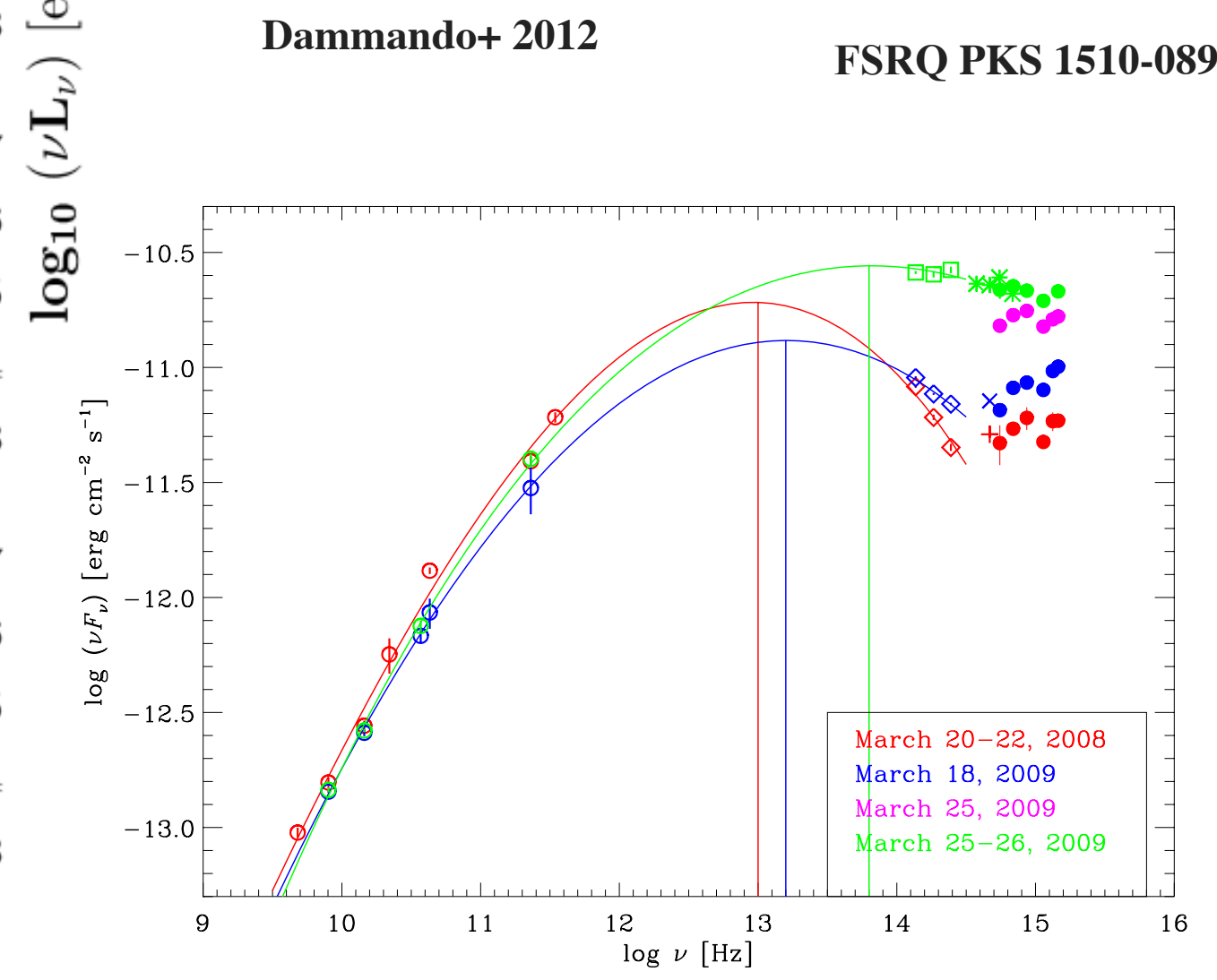
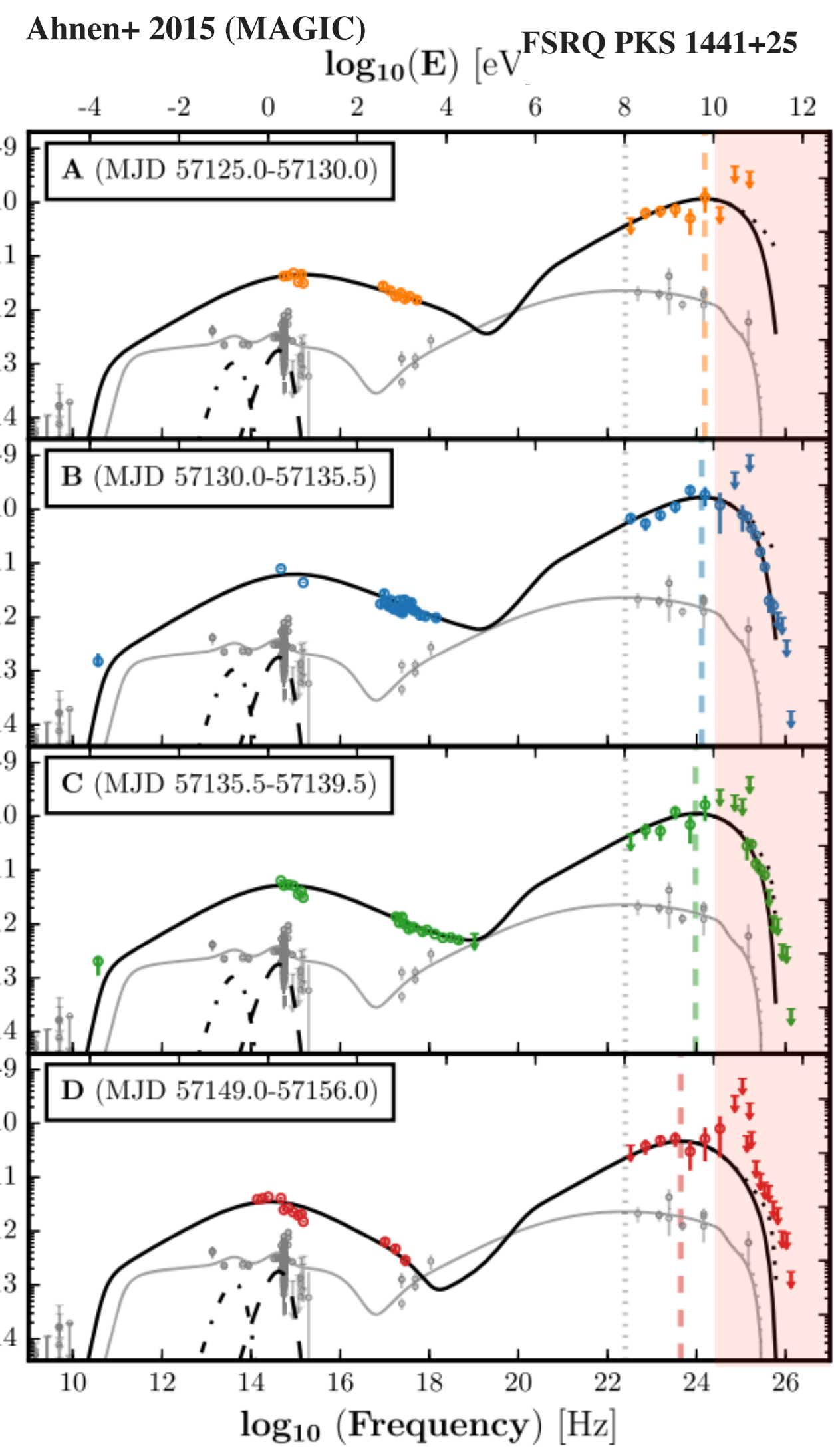
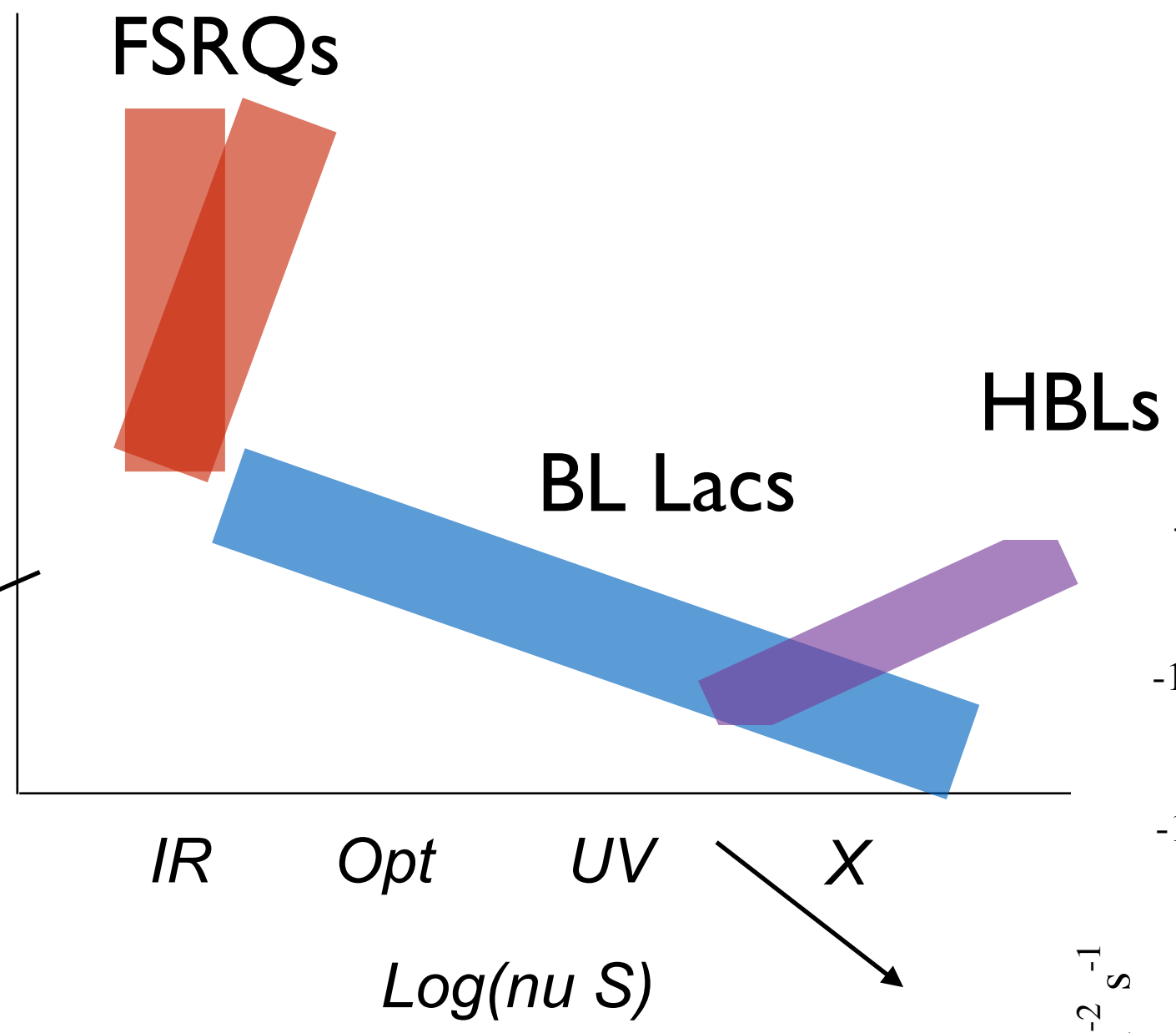


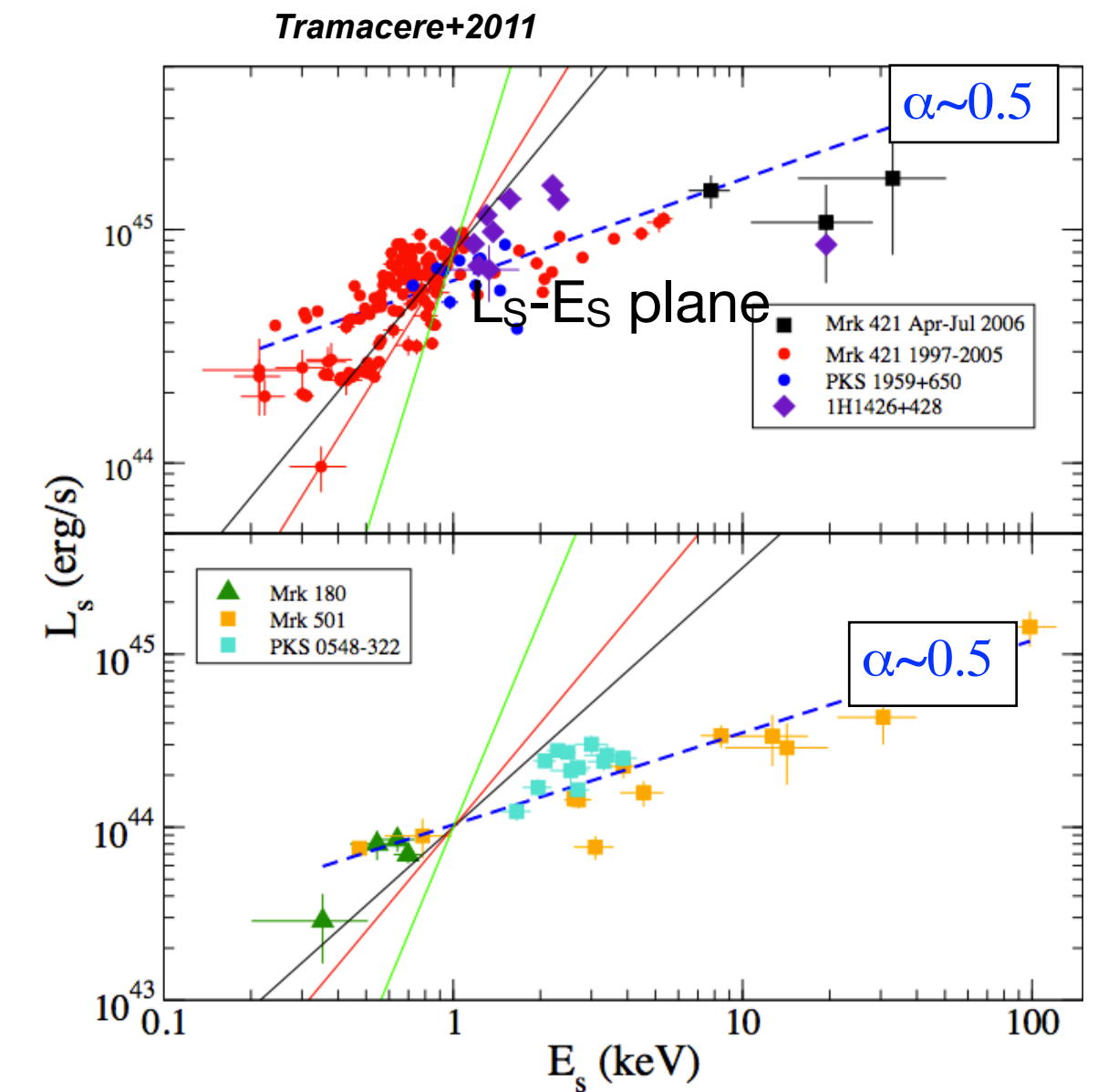
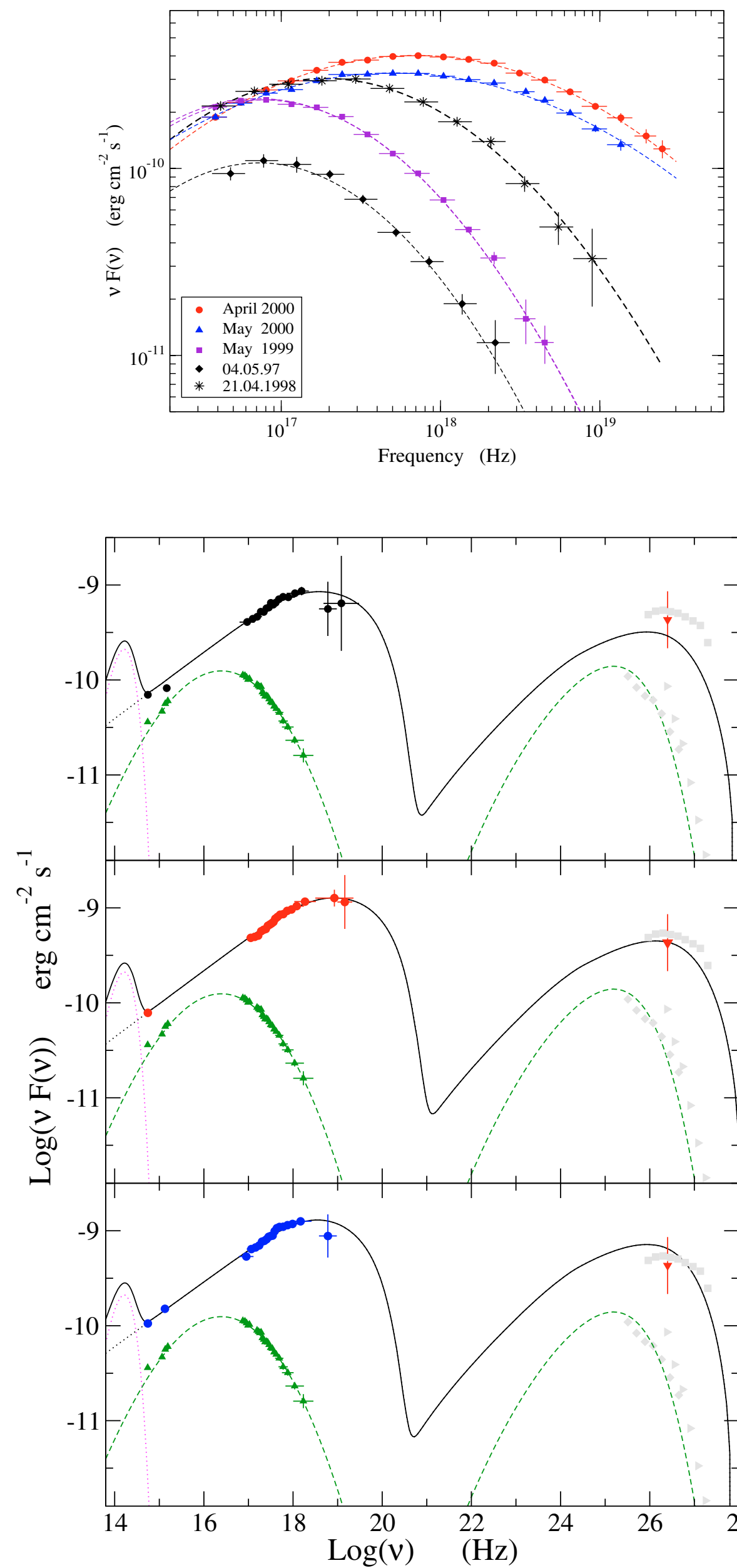
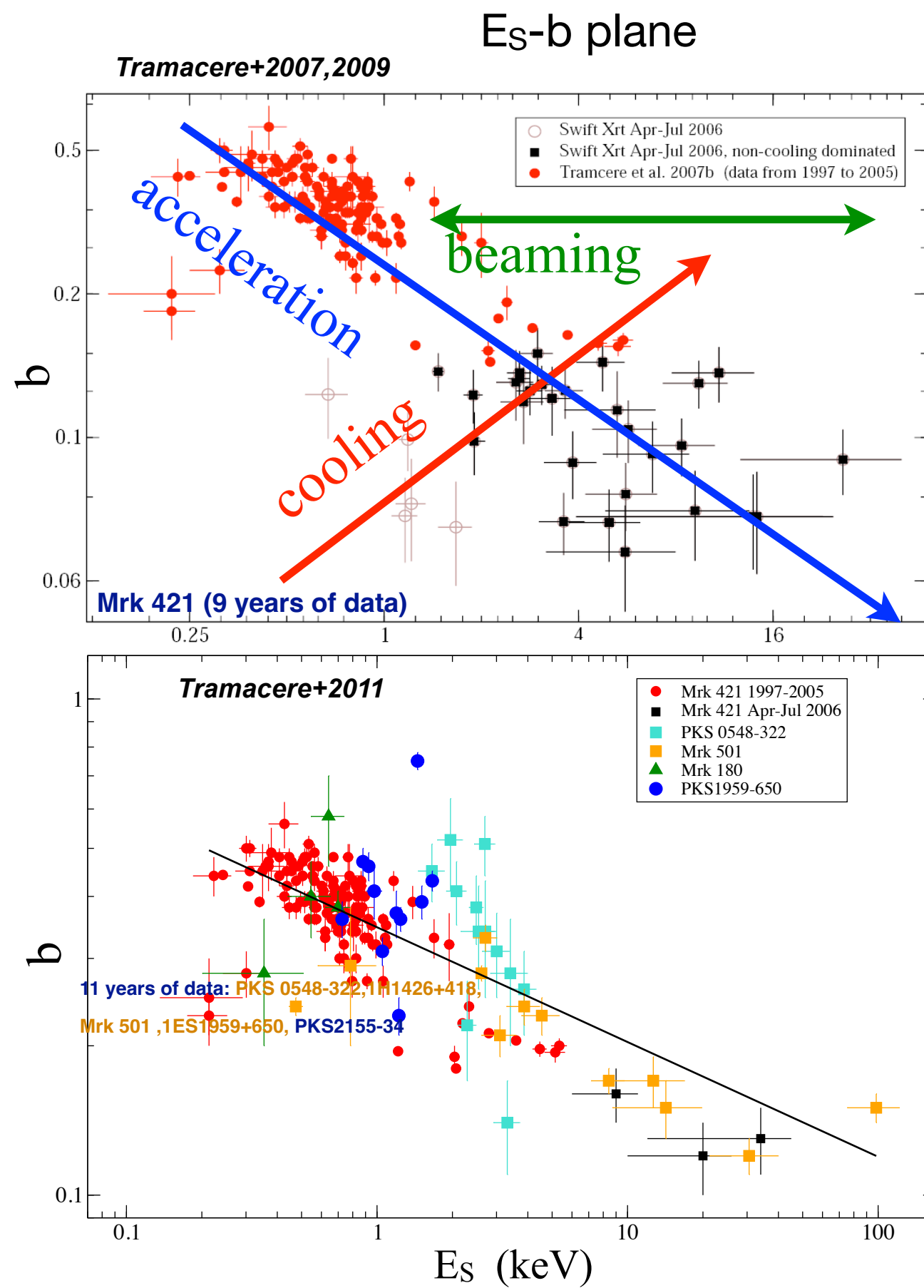
even though some peak shift are observed...



The peak frequency dynamic range, for LSP, is low, this might hide the emitters evolution

even though some peak shift are observed...



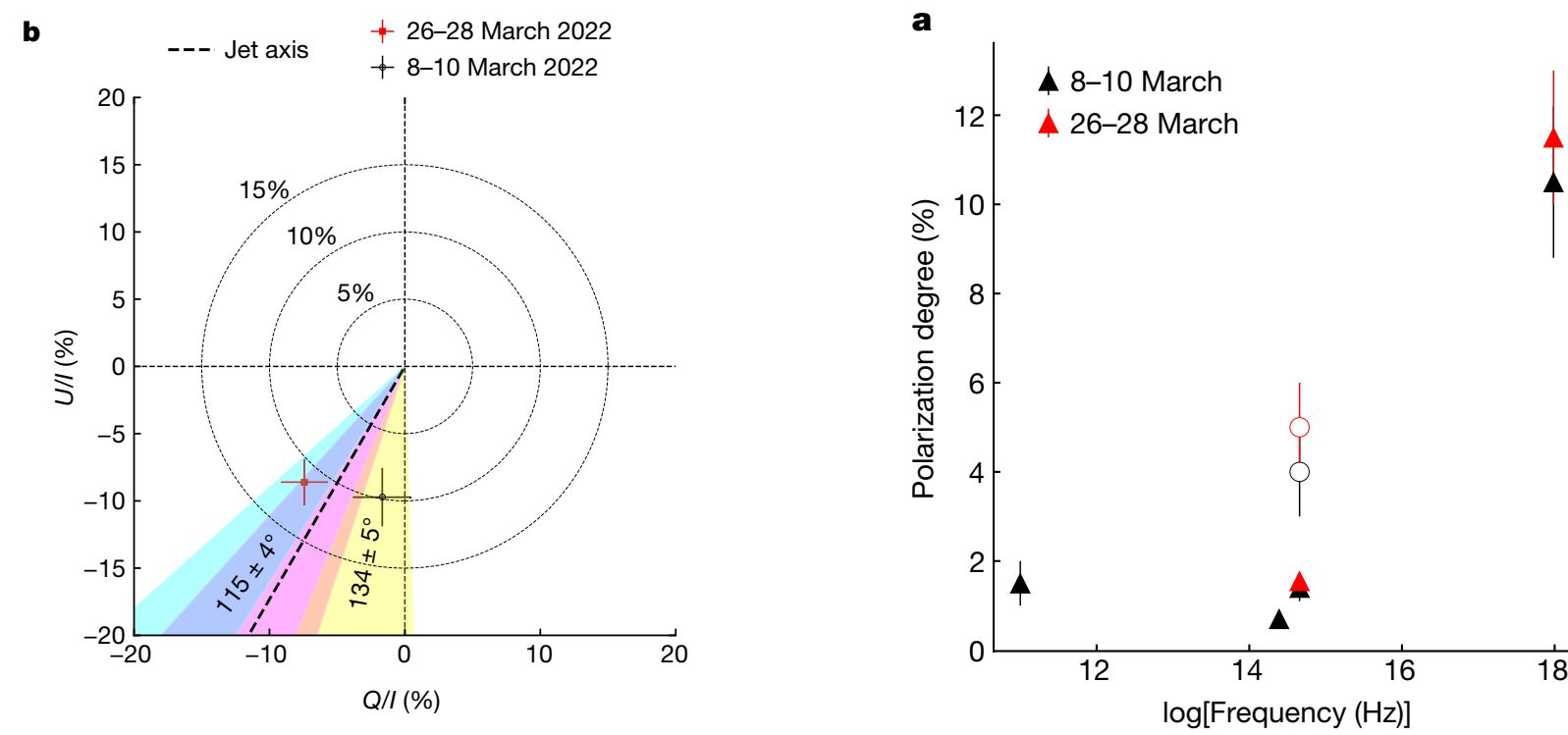


Does HBLs provide the most efficient class to study the acceleration?

Is it the blazar sequence: the larger the peak luminosity the lower the luminosity a signature of the physics, or is there a bias from considering only the cooling?

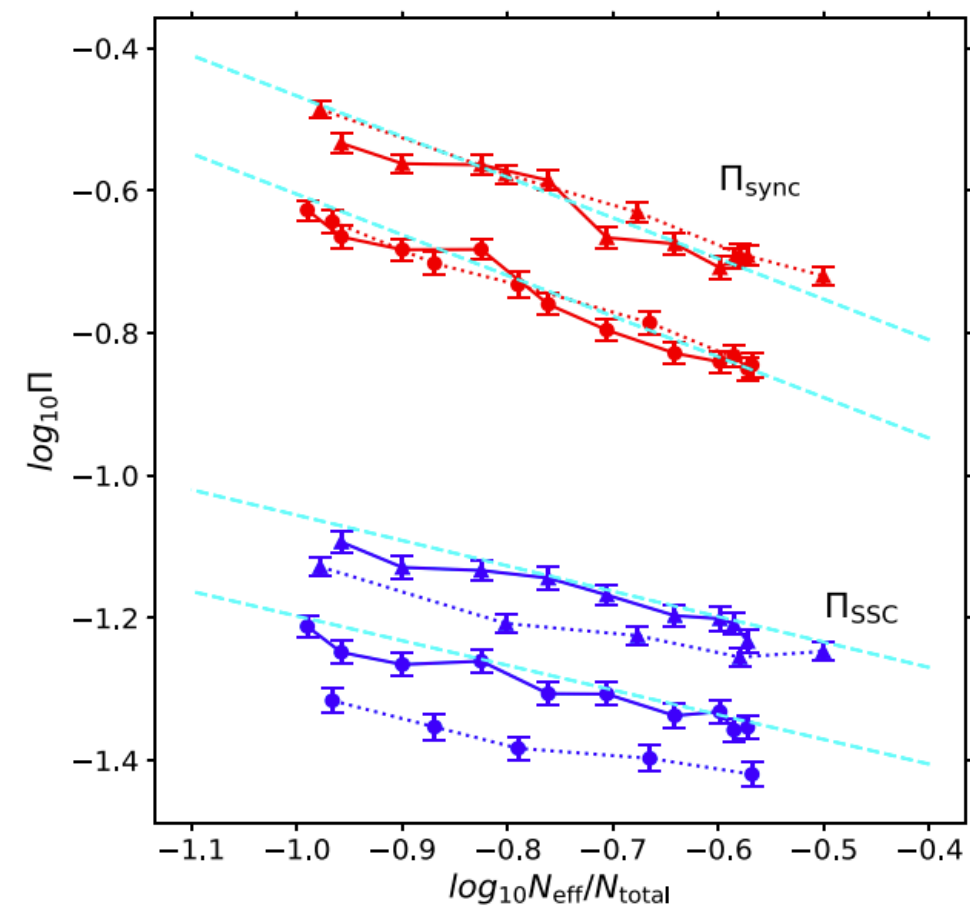
mrk 501, quiescent state iXPE recent results (nature 2022)

see Hancheng Talk



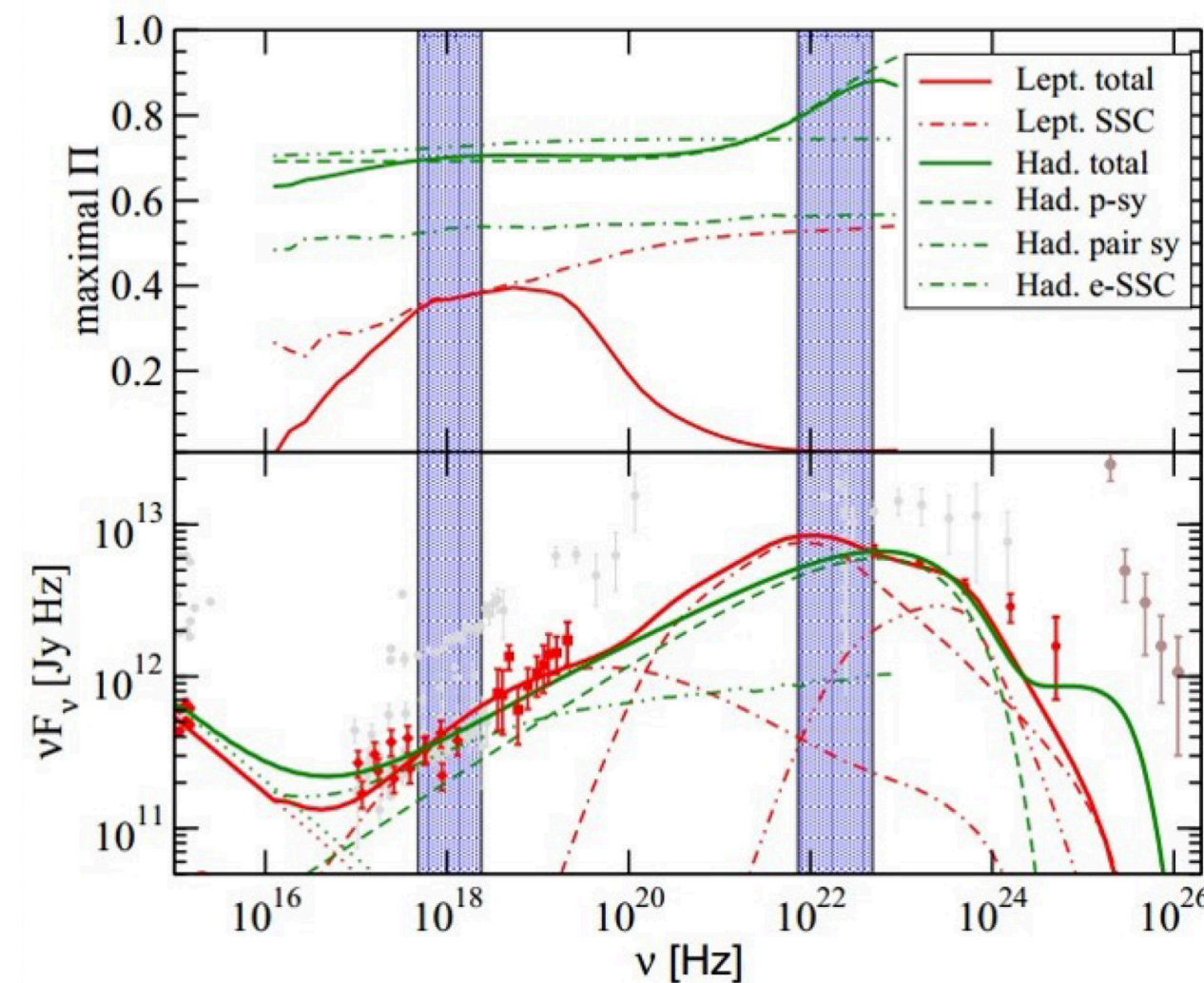
- polarization $\sim 10\% \ll \sim 70\%$ from ordered fields
- hints for turbulence impacting on **B** structure
- multi-zone vs single zone
- **B** orthogonal to the jet axis
- hadronic-leptonic discrimination (but VHE spectral signatures are mandatory to break degeneracy)
- intrinsic difference between classes

IC pol \ll S pol
(Bonometto+ 70, Pairson&Romani 2019)

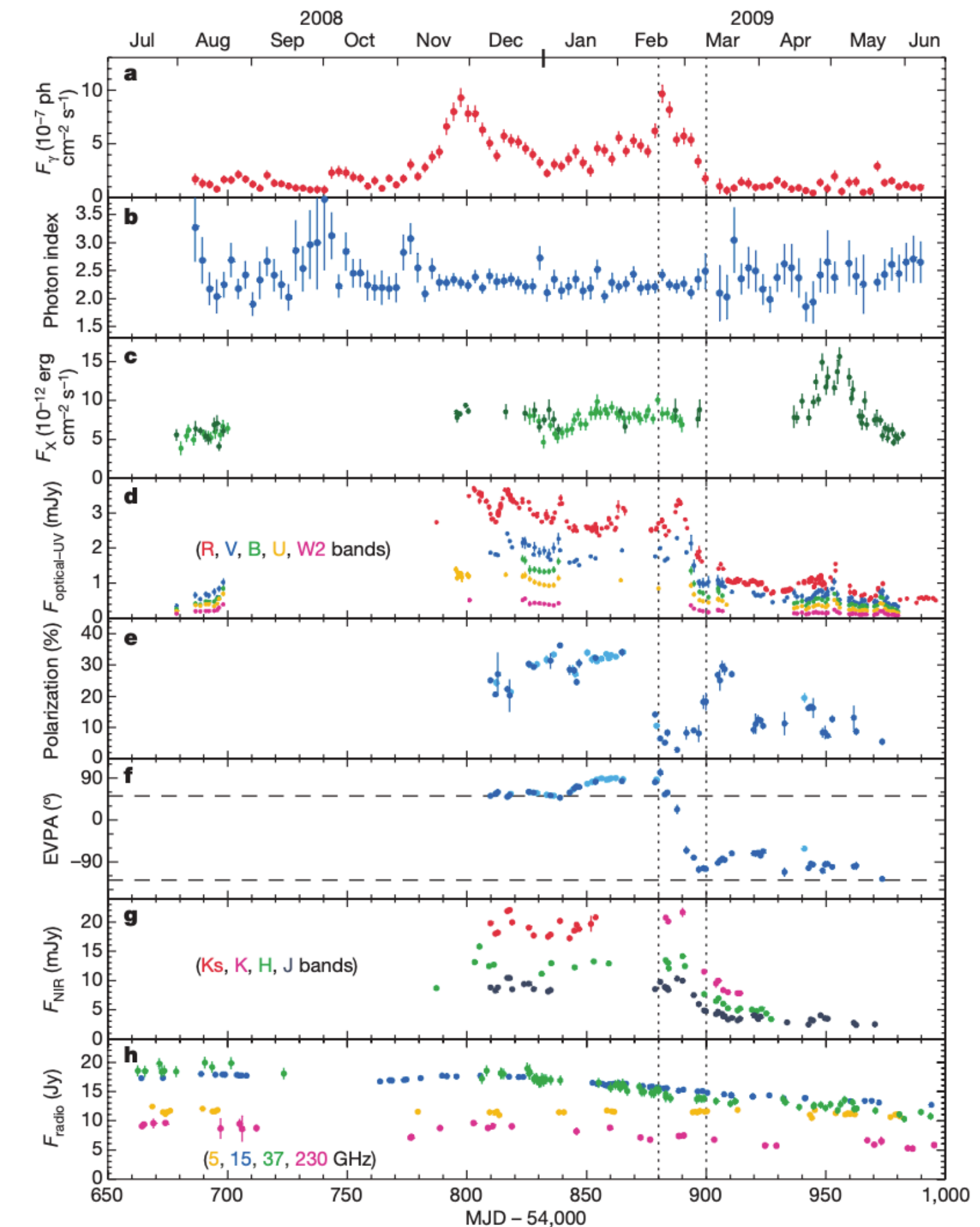


3C279 (Zhang 2017)

3C279



3C279 (nature Fermi-LAT 2010)

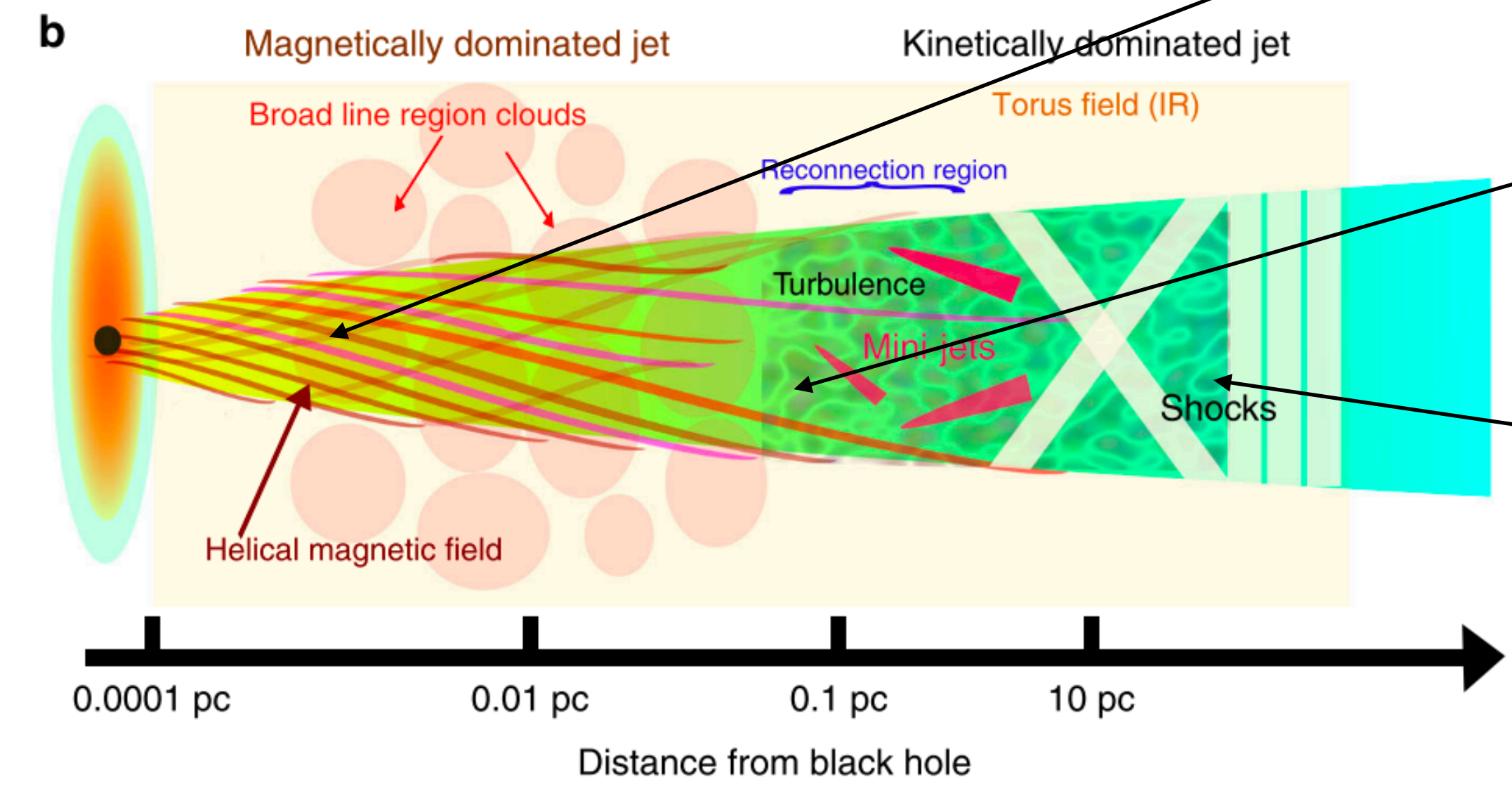




EHT Nature 2021

interaction region between an accretion-powered outflow and the fast jet spine, which is potentially powered by the black hole spin

Shukla and Mannheim,
Nature
Communications 2020



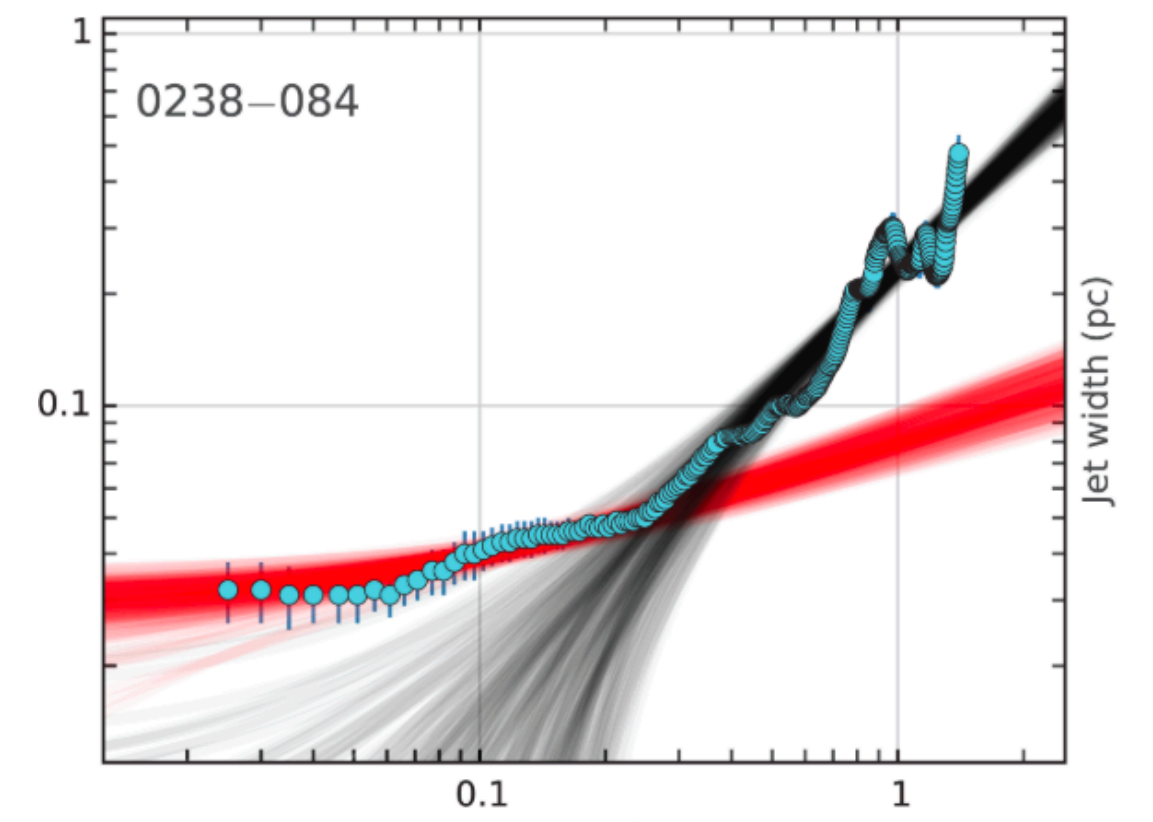
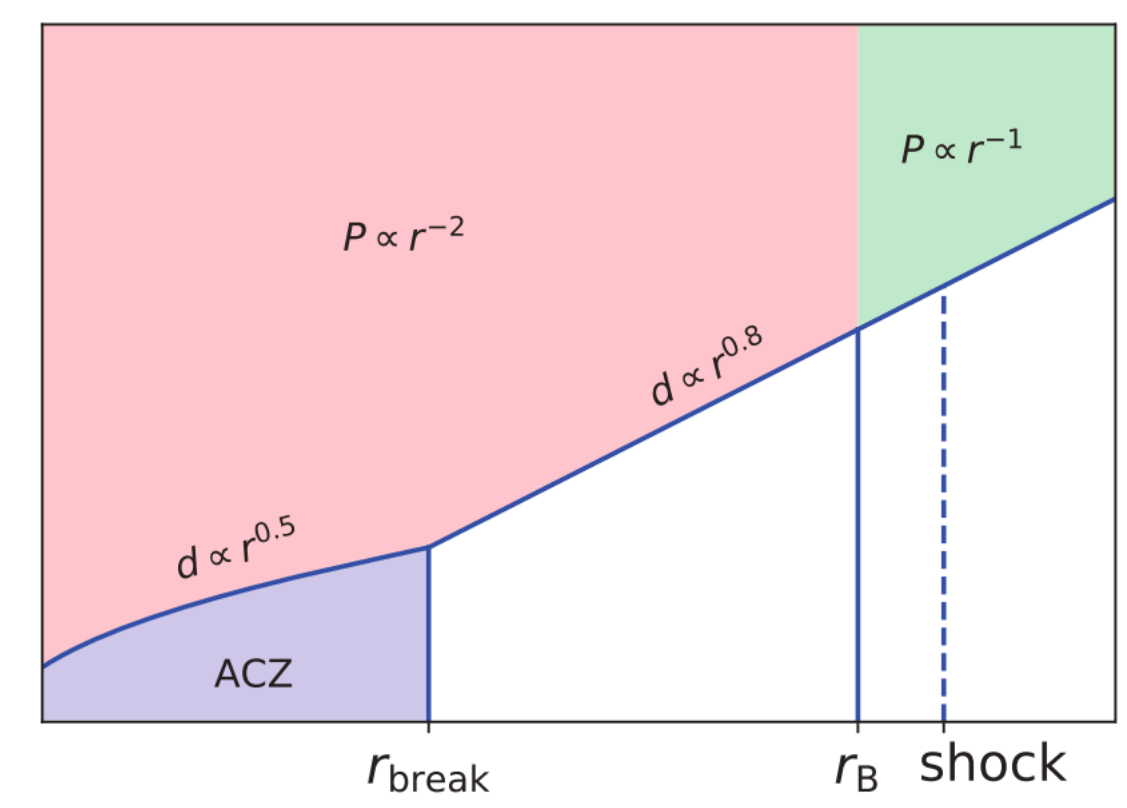
blazar zone
(x/γ) flares
turb. stochastic
acceleration

expansion
radio delay

Radio morphology and radio-gamma delays

Kovalev+ 2020

break occurs at MD~KD, Γ saturates



here we add **adiabatic cooling**
(t=time elapsed from the expansion)

Parker 2006, Stawartz&Petrosian 2008 Tramacere+2011

$$|\dot{\gamma}_{ad}| = \frac{1}{3} \frac{\dot{V}}{V} \gamma = \frac{\dot{R}(t)}{R(t)} \gamma = \frac{\beta_{exp} c}{R(t)} \gamma$$

JetSeT

injection term

$$L_{inj} = V_{acc} \int \gamma m_e c^2 Q(\gamma, t) d\gamma \quad (erg/s)$$

systematic term

$$S(\gamma, t) = -C(\gamma, t) + A(\gamma, t)$$

cooling term

$$C(\gamma) = |\dot{\gamma}_{synch}| + |\dot{\gamma}_{IC}| + |\dot{\gamma}_{ad}|$$

syst. acc. term

$$A(\gamma) = A_{p0} \gamma, \quad t_A = \frac{1}{A_0}$$

$$\frac{\partial n(\gamma, t)}{\partial t} = \frac{\partial}{\partial \gamma} \left\{ - [S(\gamma, t) + D_A(\gamma, t)] n(\gamma, t) \right\} + \frac{\partial}{\partial \gamma} \left\{ D_p(\gamma, t) \frac{\partial n(\gamma, t)}{\partial \gamma} \right\} - \frac{n(\gamma, t)}{T_{esc}(\gamma)} - \frac{n(\gamma, t)}{T_{ad}} + Q(\gamma, t)$$

$$T_{ad} = \frac{1}{3} \frac{R(t)}{\beta_{exp} c} \quad (\text{Gould 1975})$$

Turbulent magnetic field

$$W(k) = \frac{\delta B(k_0^2)}{8\pi} \left(\frac{k}{k_0} \right)^{-q}$$



momentum diffusion term

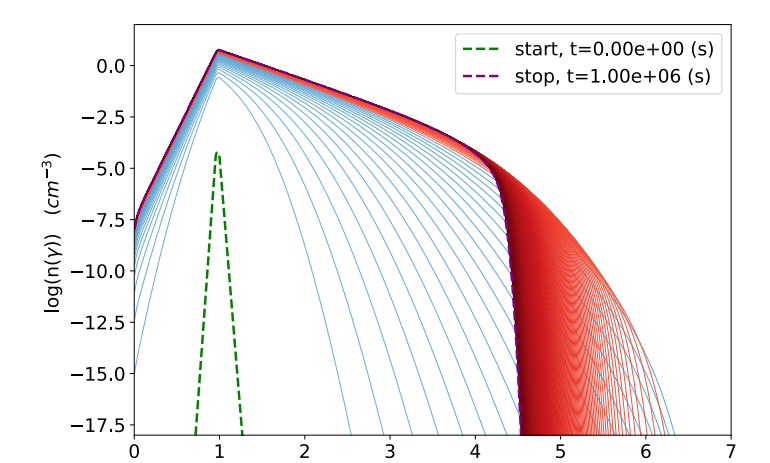
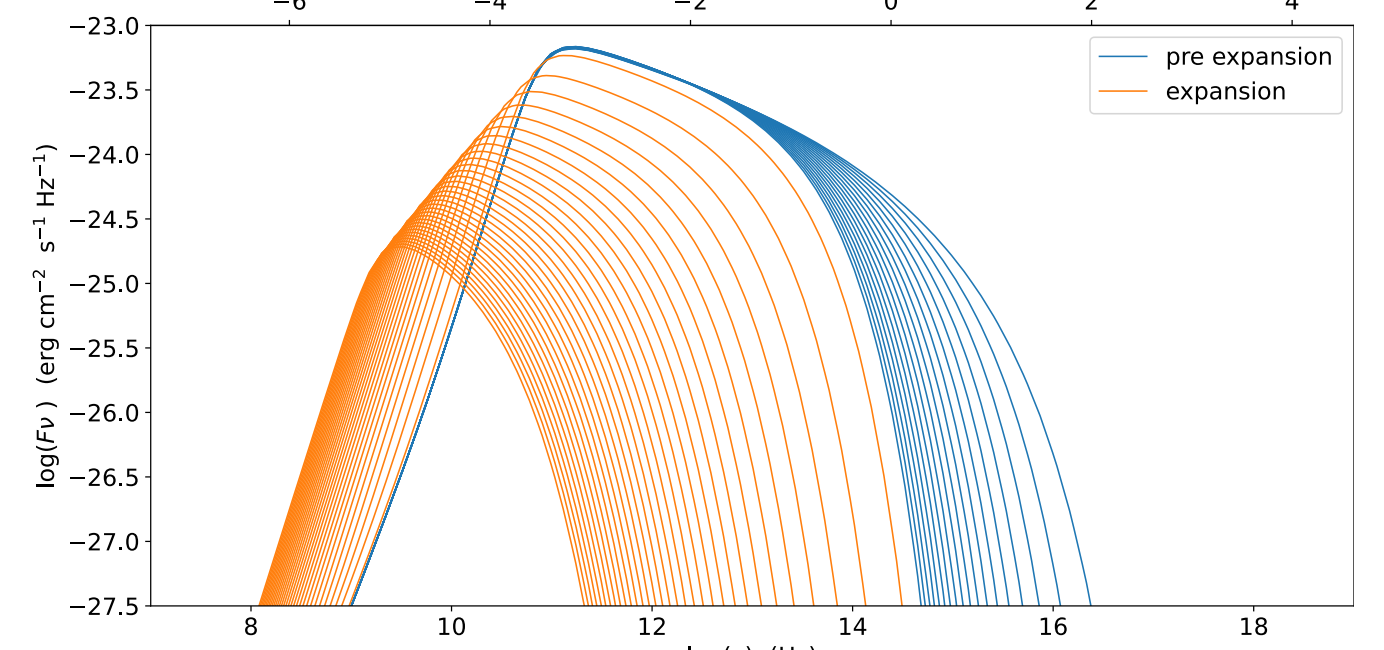
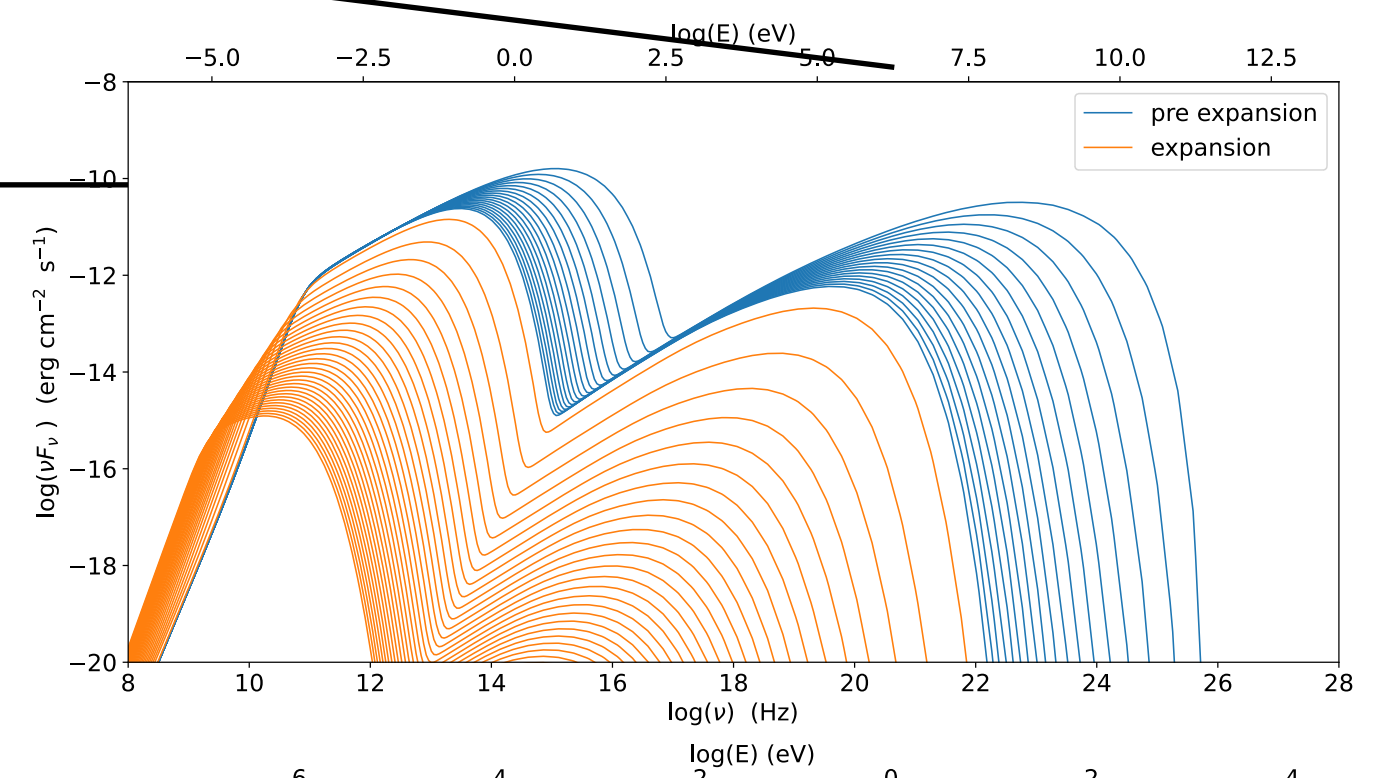
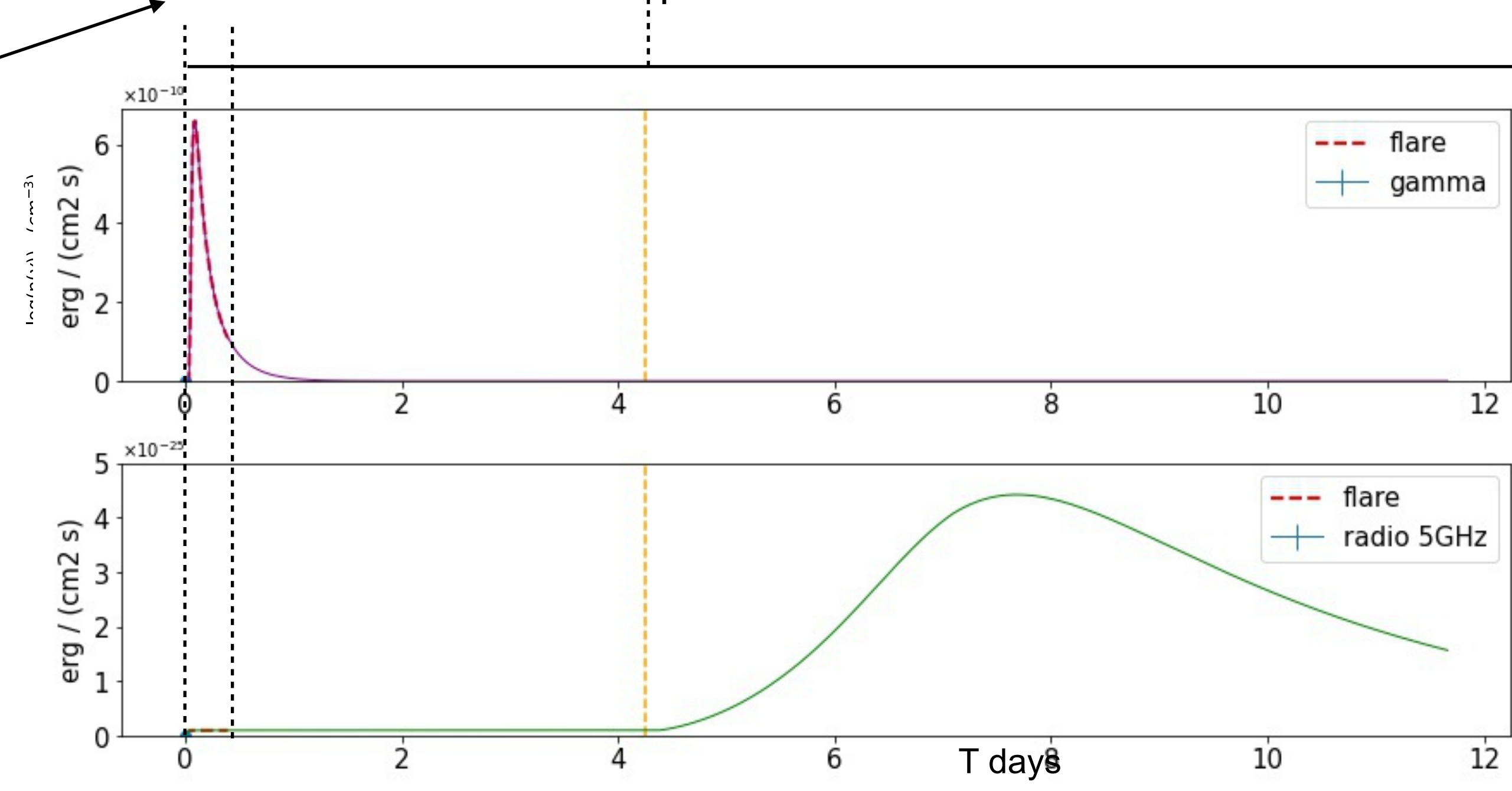
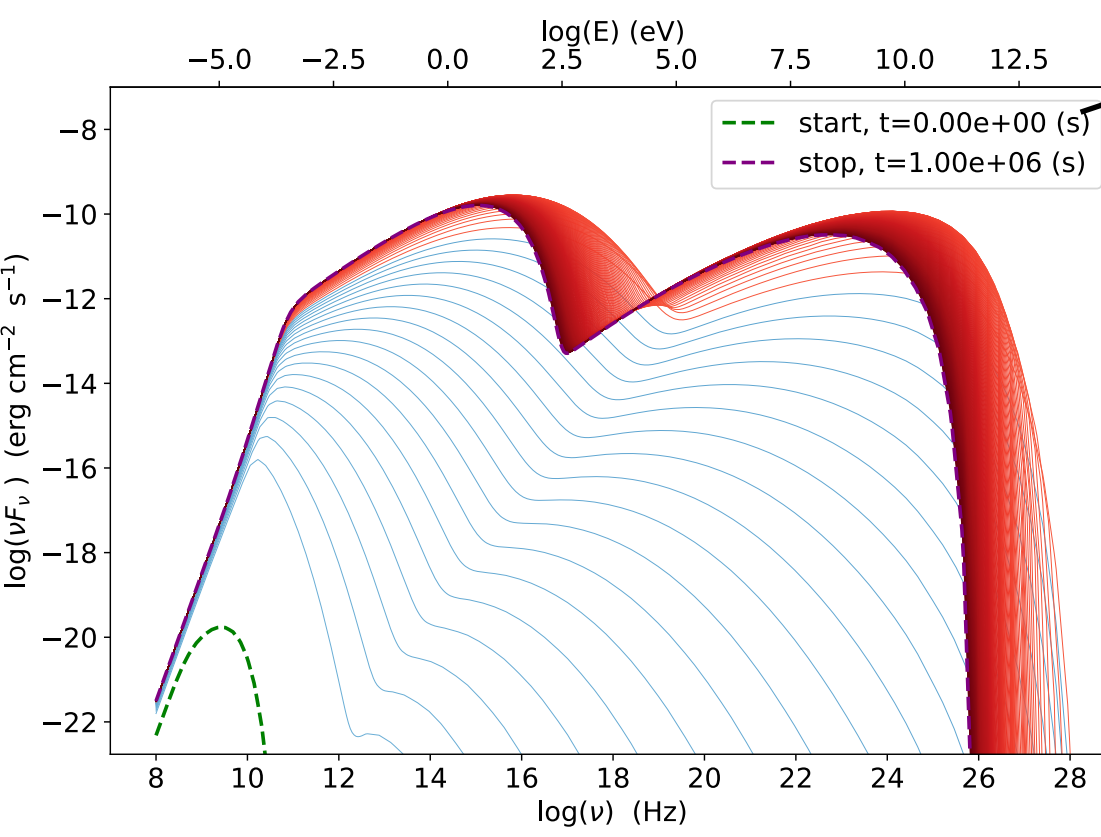
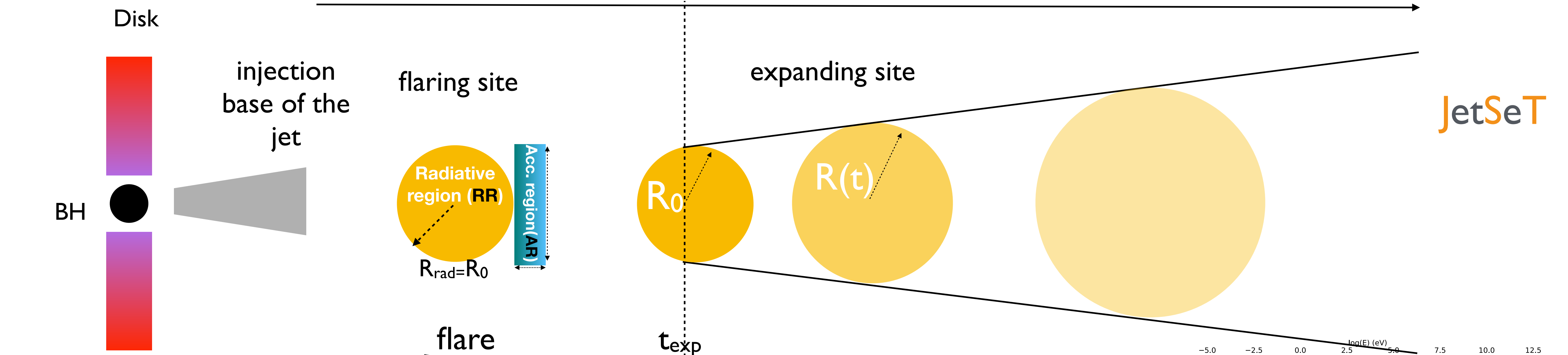
$$D_p \approx \beta_A^2 \left(\frac{\delta B}{B_0} \right)^2 \left(\frac{\rho_g}{\lambda_{max}} \right)^{q-1} \frac{p^2 c^2}{\rho_g c}$$

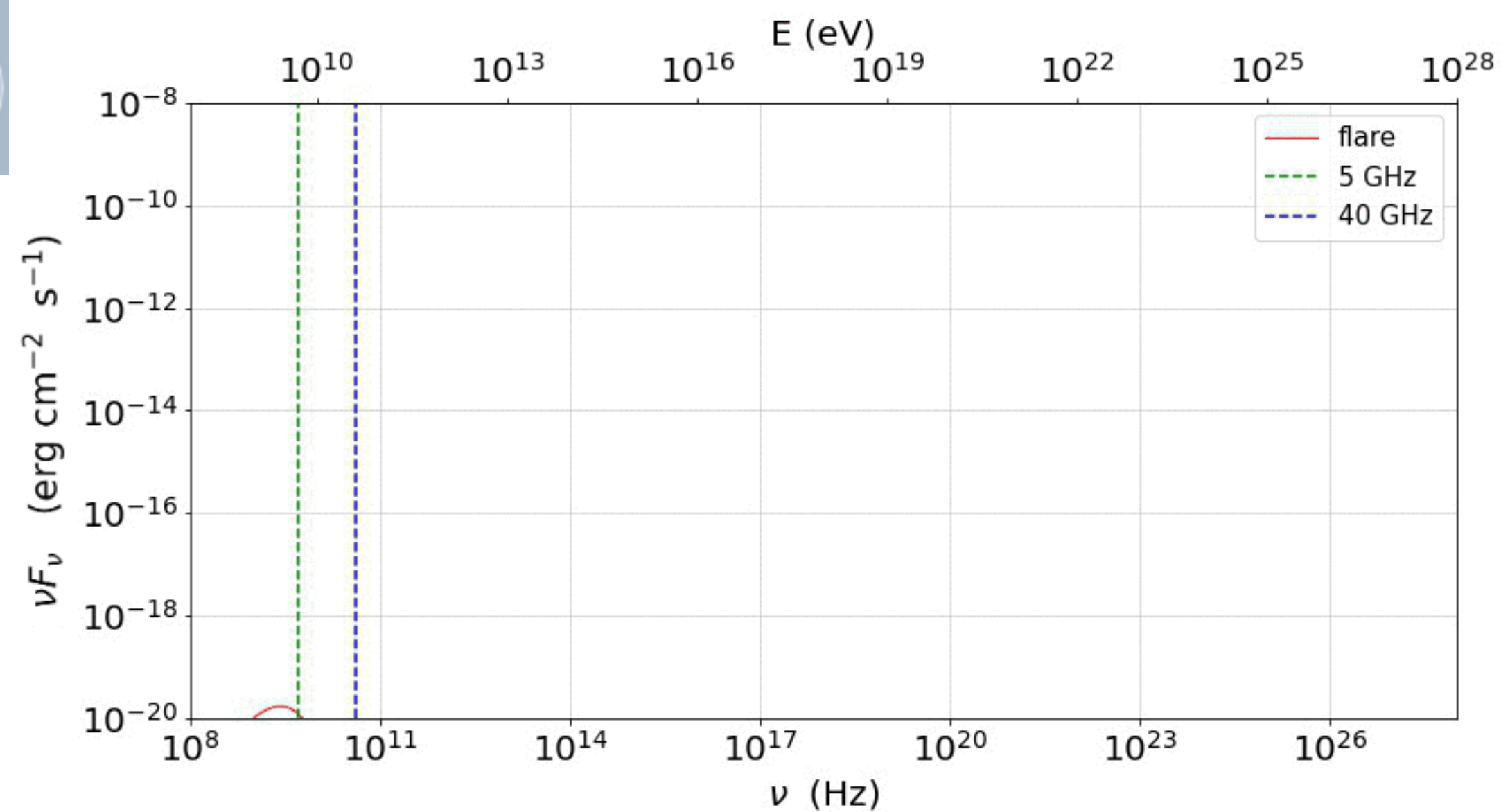
Tramacere+ 2022 (see also Boula 2022)

$$\Delta r = t_{\text{exp}} \beta c \Gamma \quad (\text{obs rest frame})$$

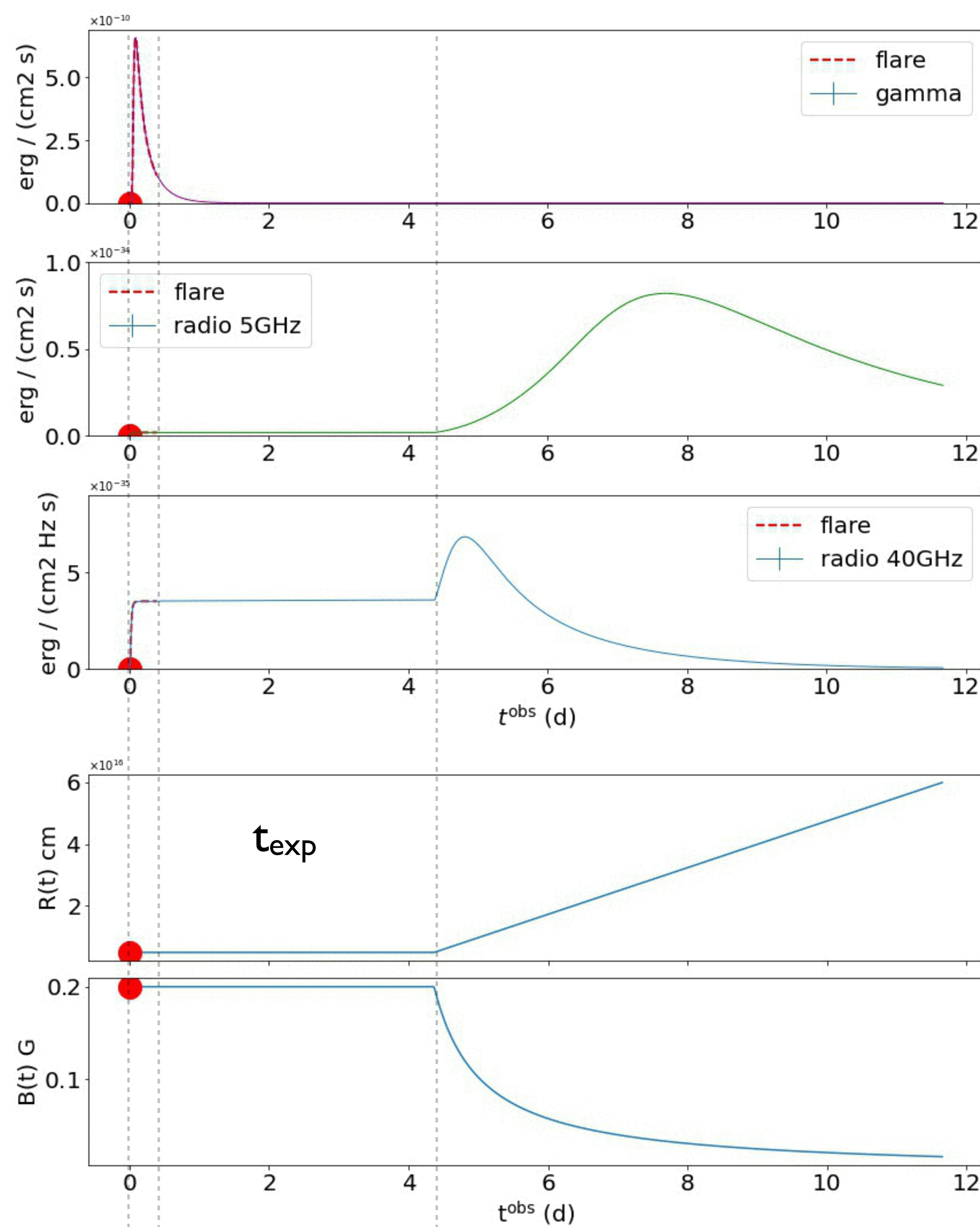
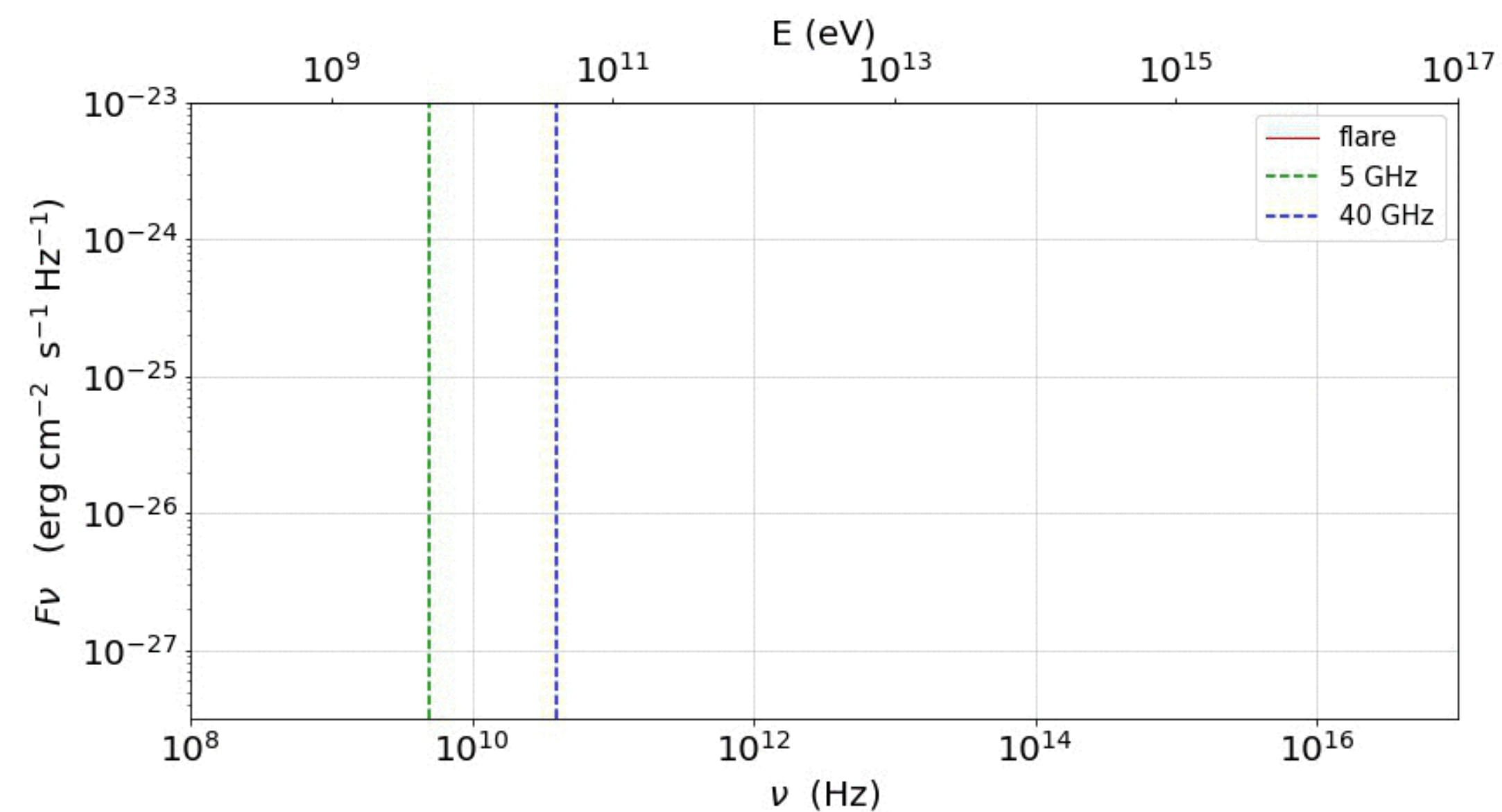
R_H

JetSeT





- duration $\sim 3 \times 10^7$ s (blob frame) ~ 11 d obs
- $t_{\text{exp}} = 1 \times 10^7$ s
- $\beta_{\text{exp}} = 0.1c$

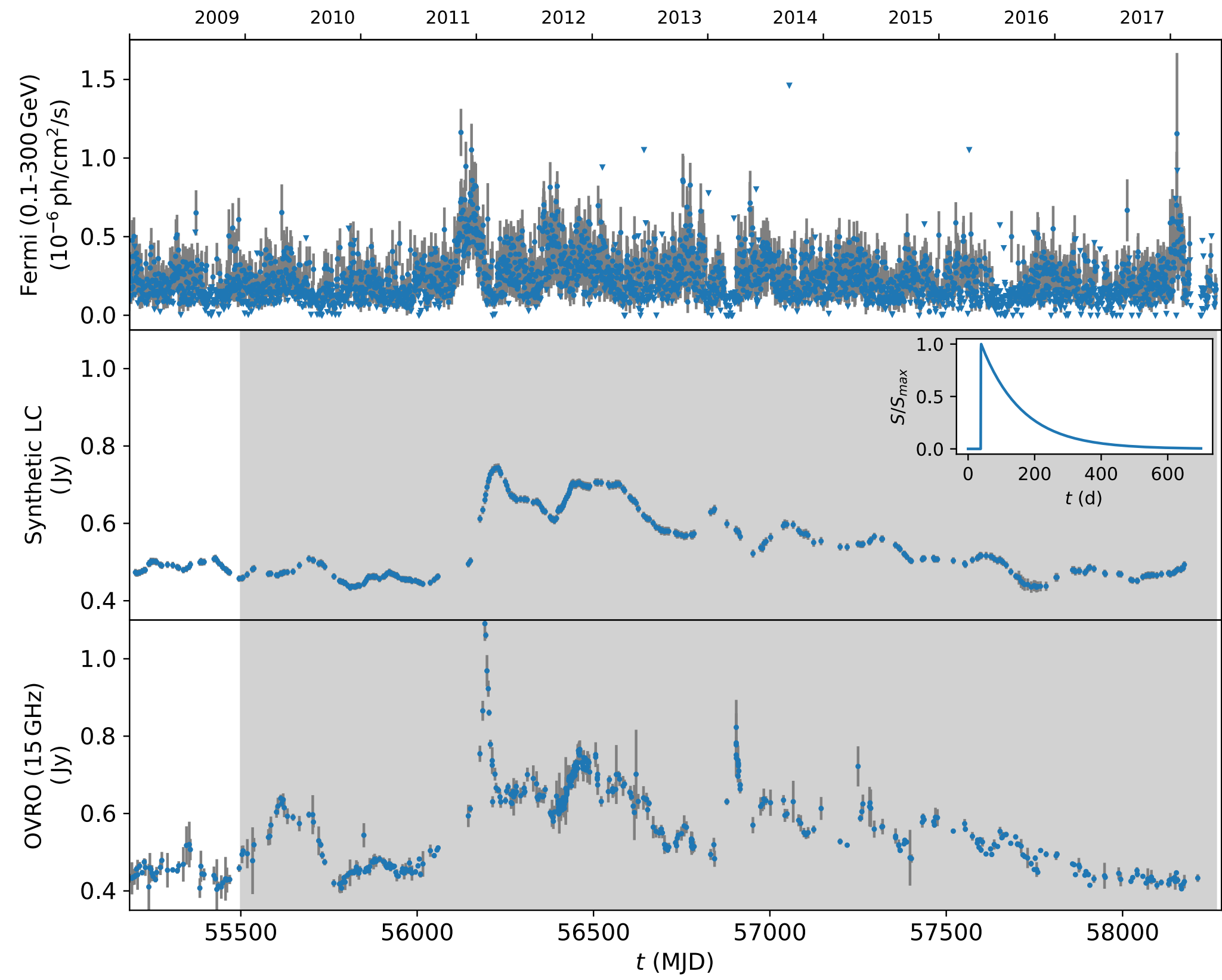


flare

expansion

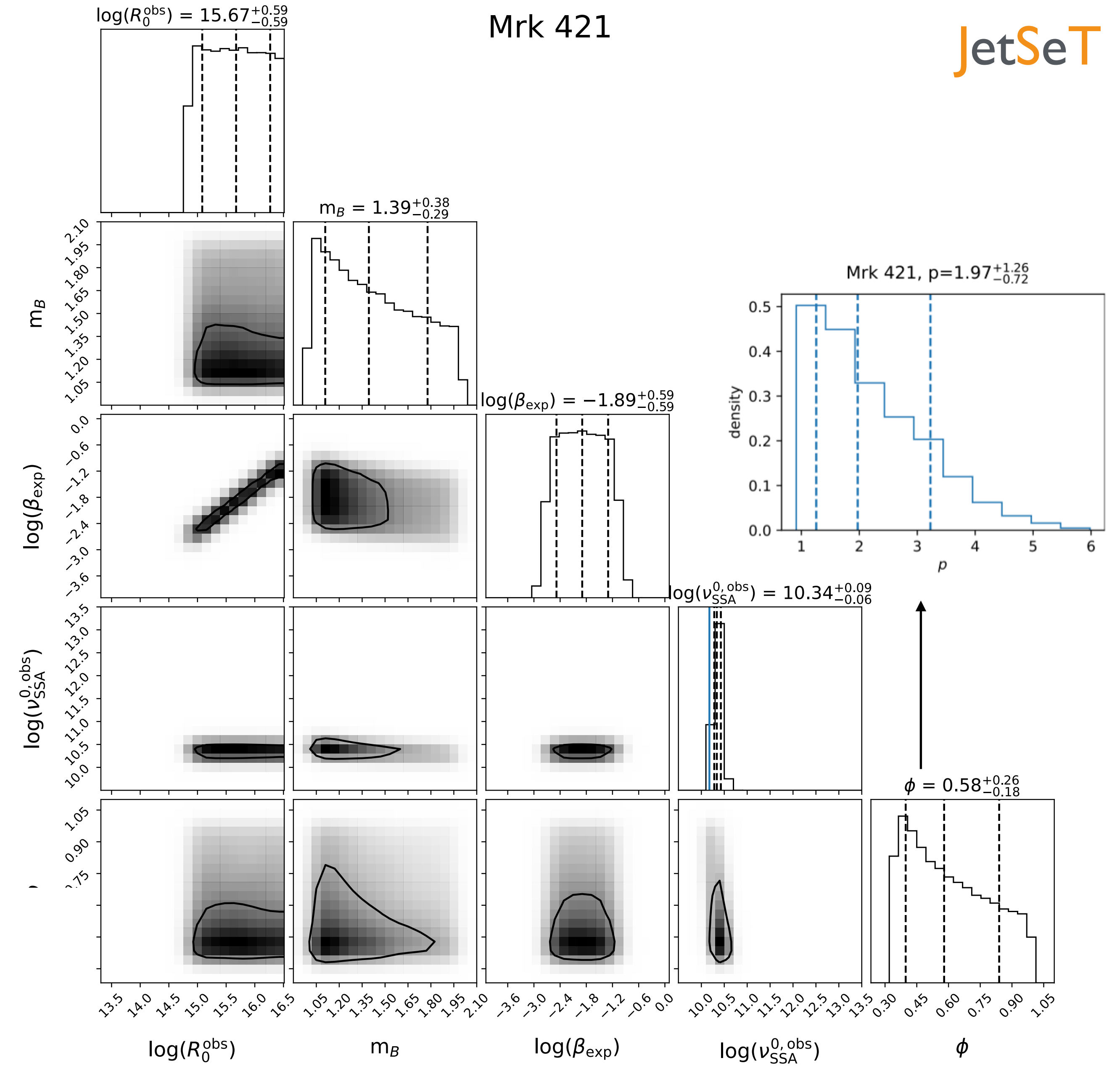
Tramacere+ 2022

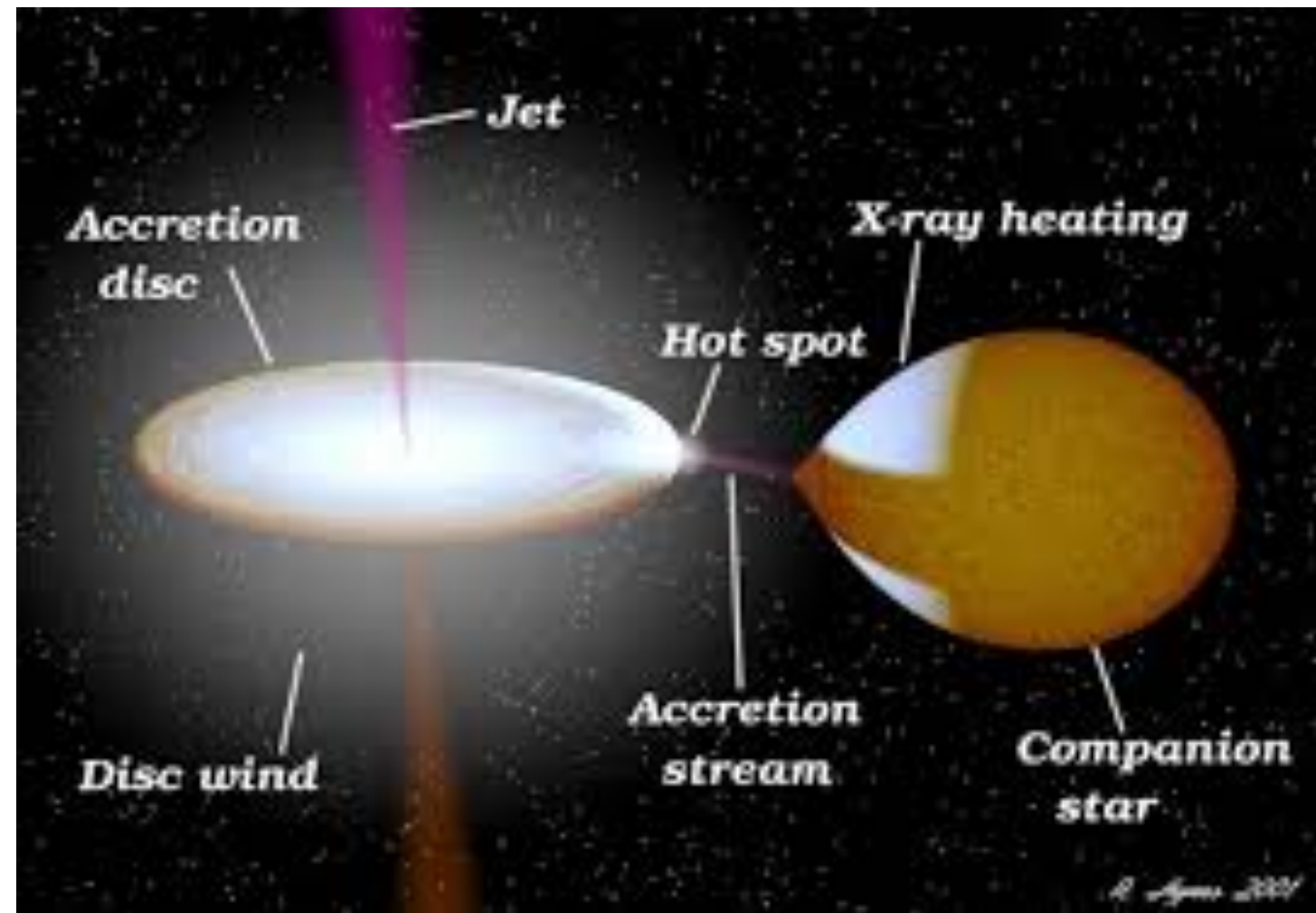
Tramacere+ 2022



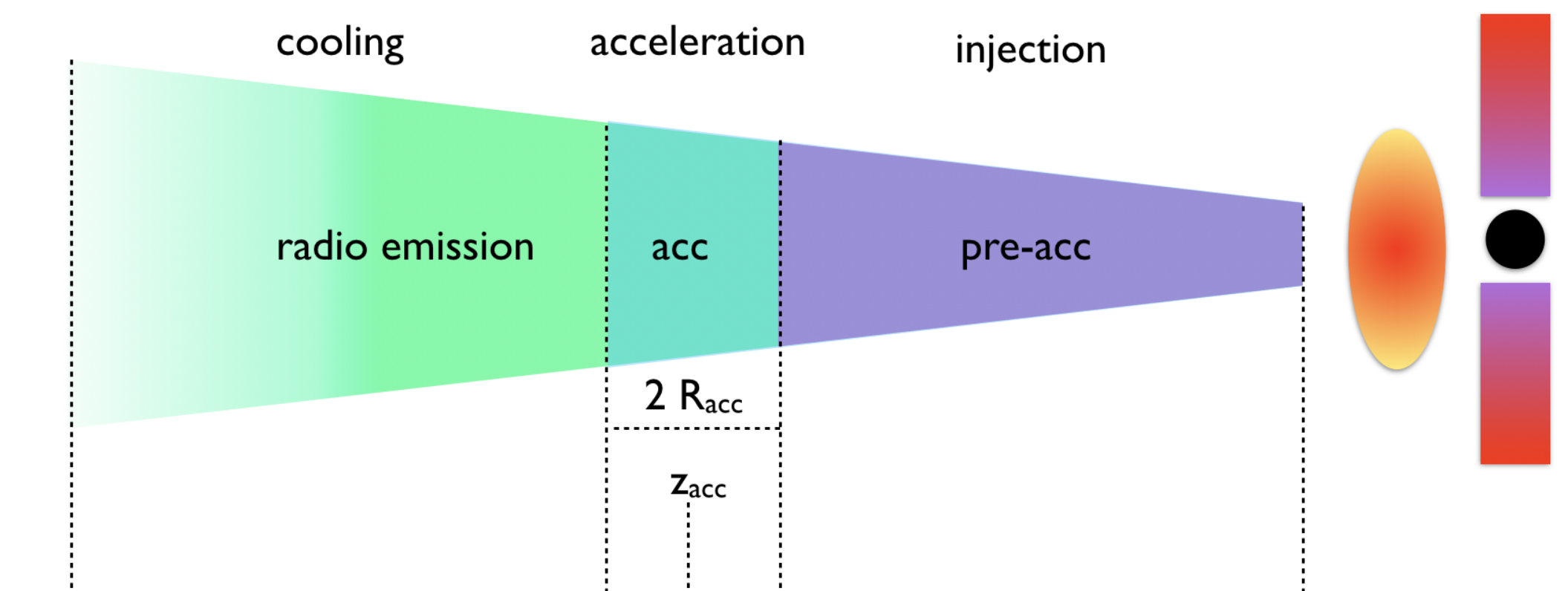
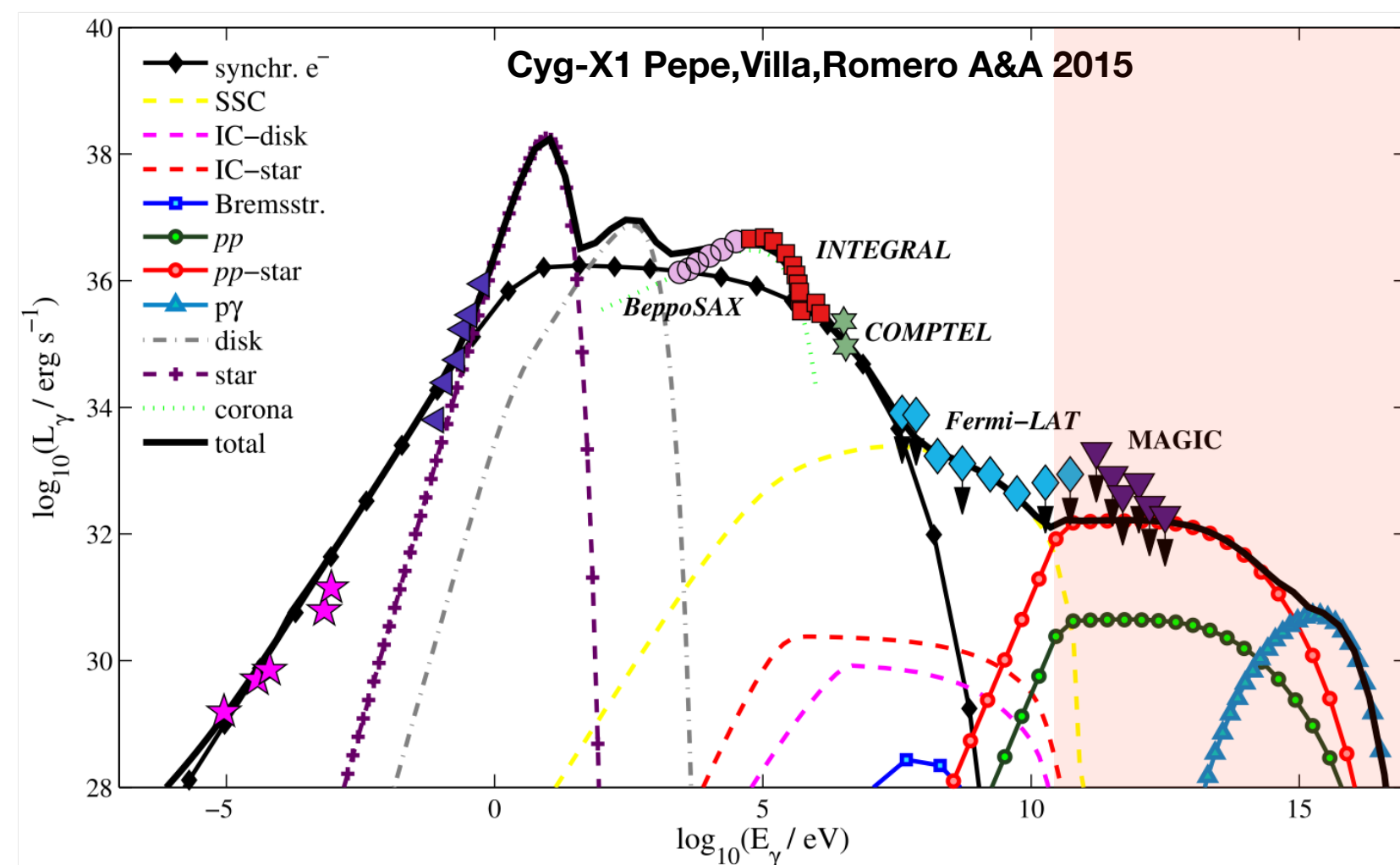
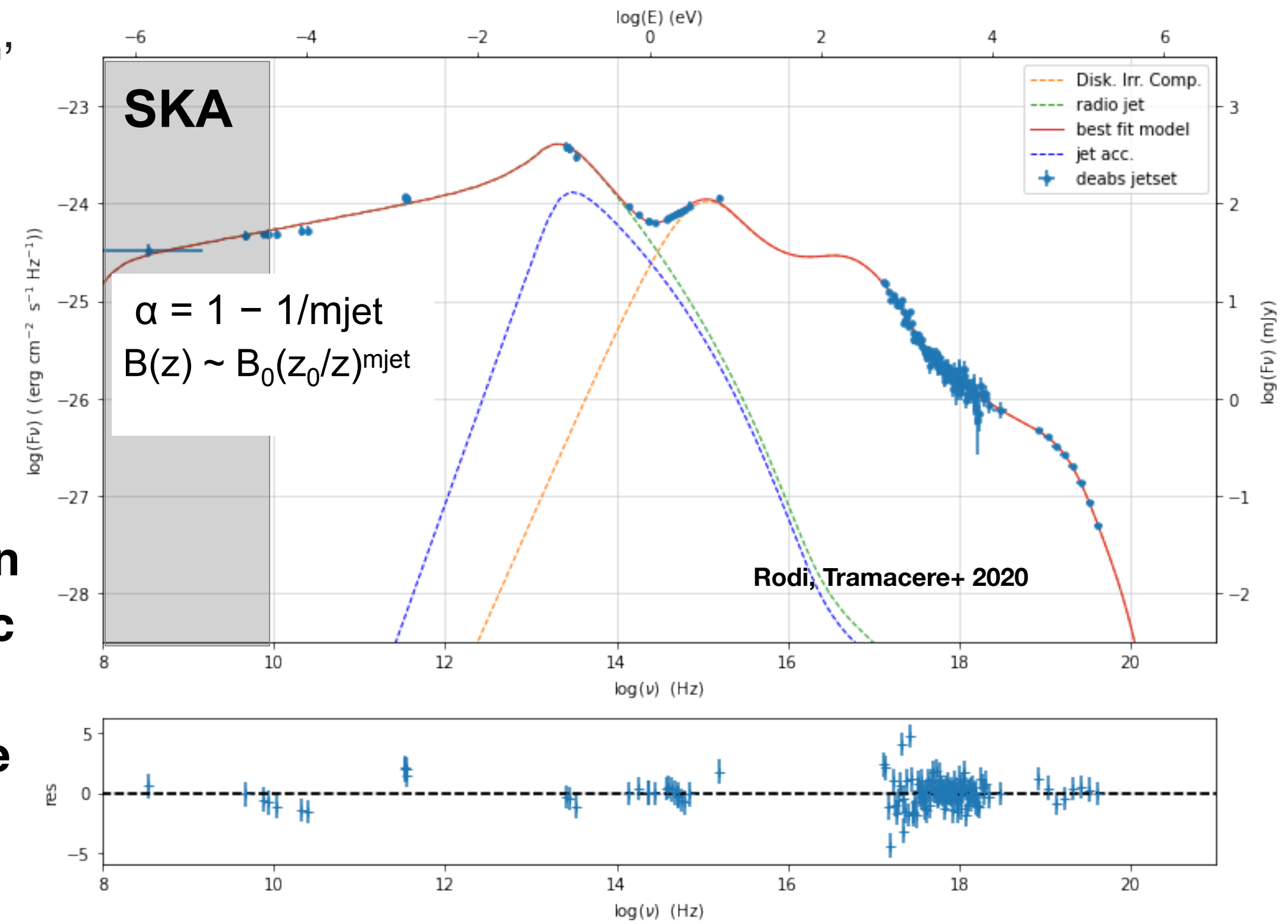
Parameter	Value
A	$12.5^{+0.5}_{-0.013} \times 10^3 \text{ Jy cm}^2 \text{ s/ph}$
t_{rise}	$\lesssim 1 \text{ day}$
t_{decay}	$126.5^{+1.3}_{-1.3} \text{ days}$
Δt	$37.58^{+0.13}_{-0.13} \text{ days}$
$F_{\text{background}}$	$0.18^{+0.008}_{-0.0004} \text{ Jy}$

$$S(t) = A \frac{\exp^{-(t-\Delta t)/t_{\text{decay}}^*}}{1 + \exp^{-(t-\Delta t)/t_{\text{rise}}}}$$





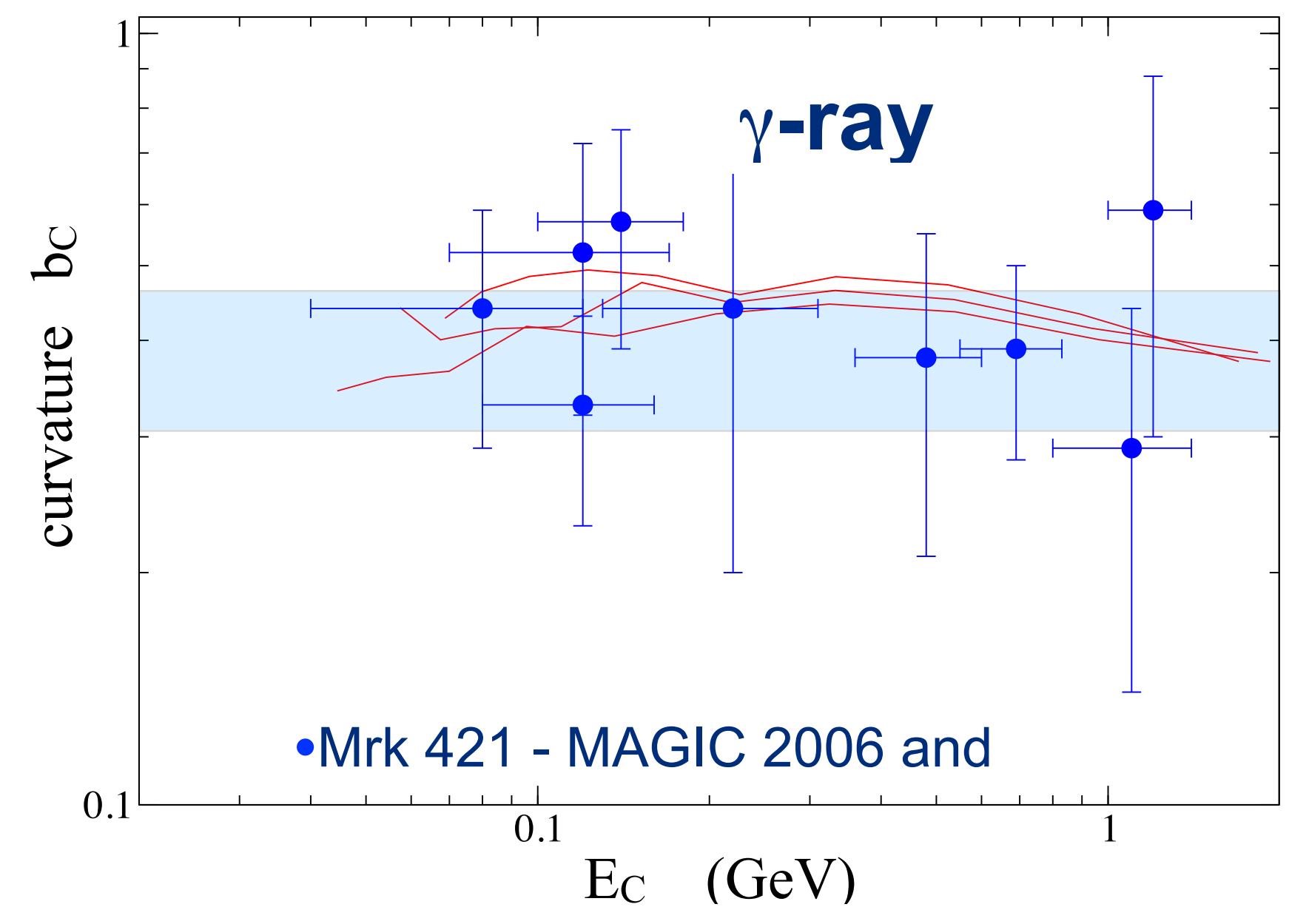
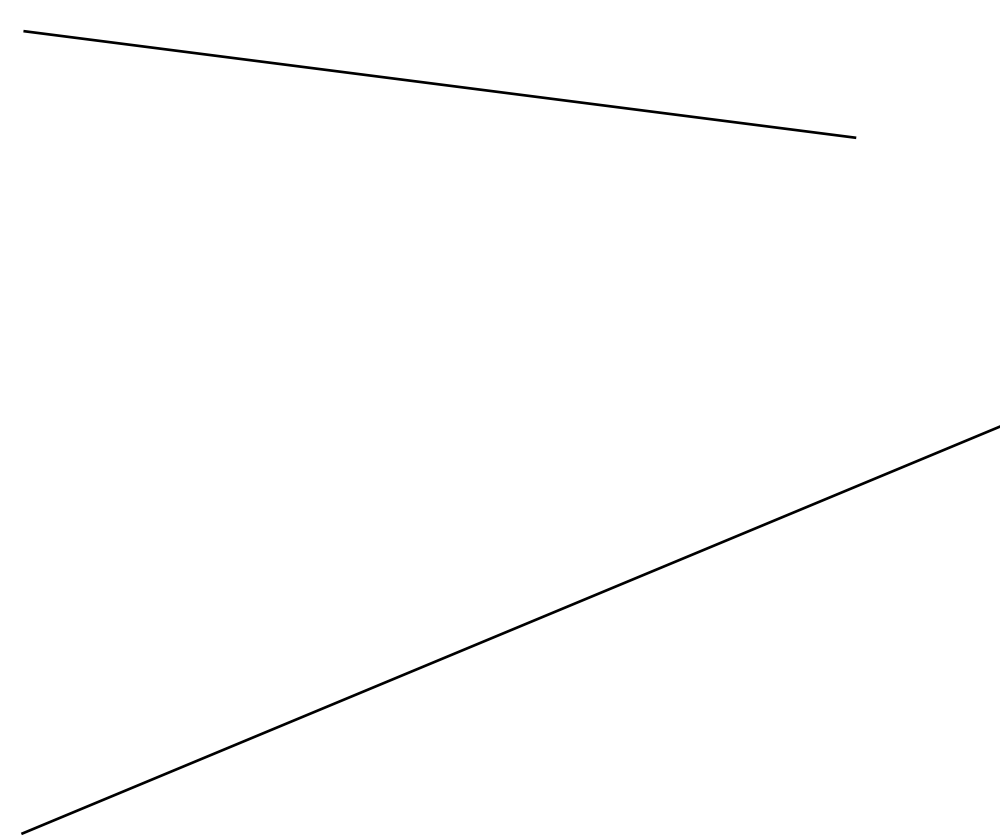
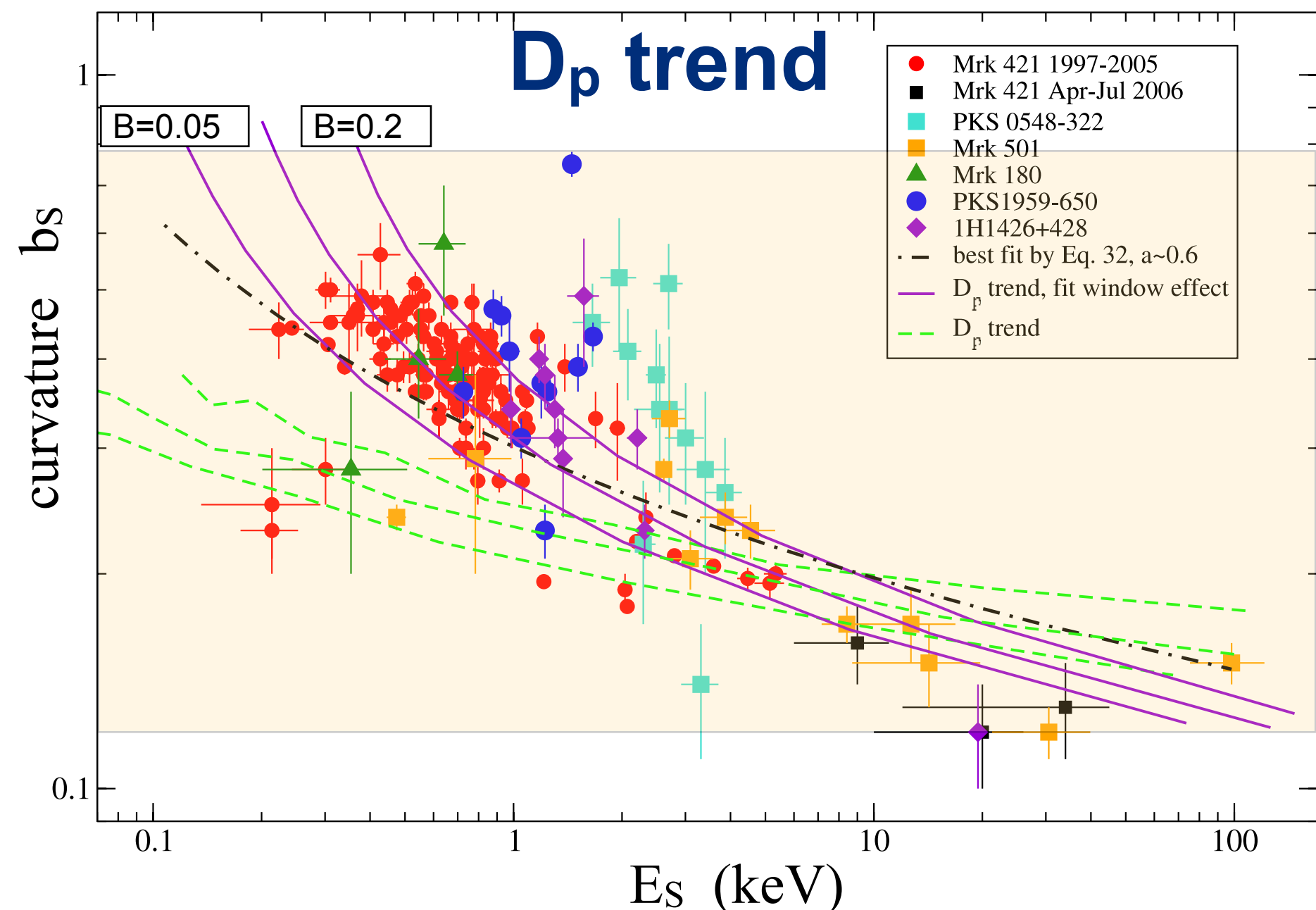
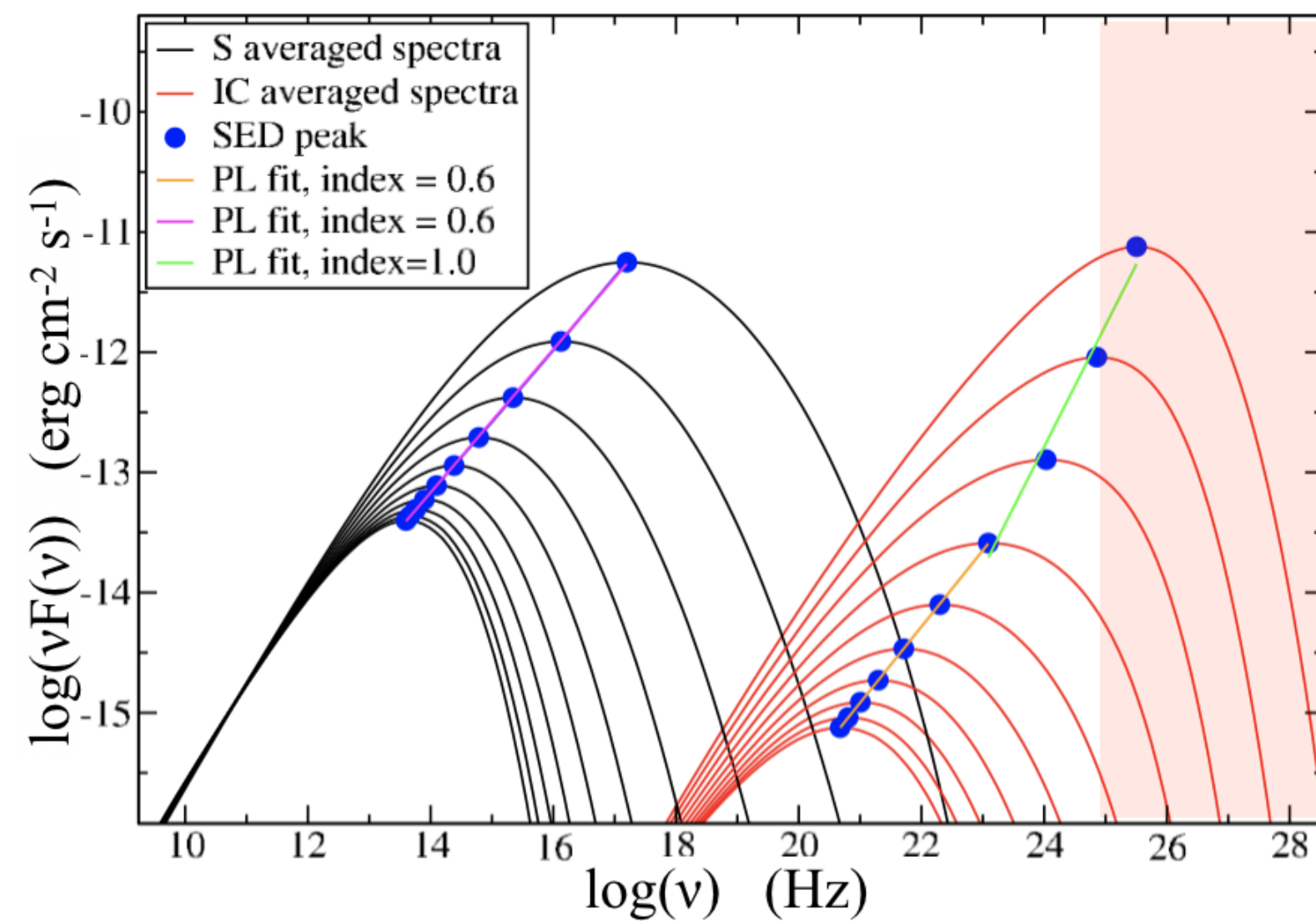
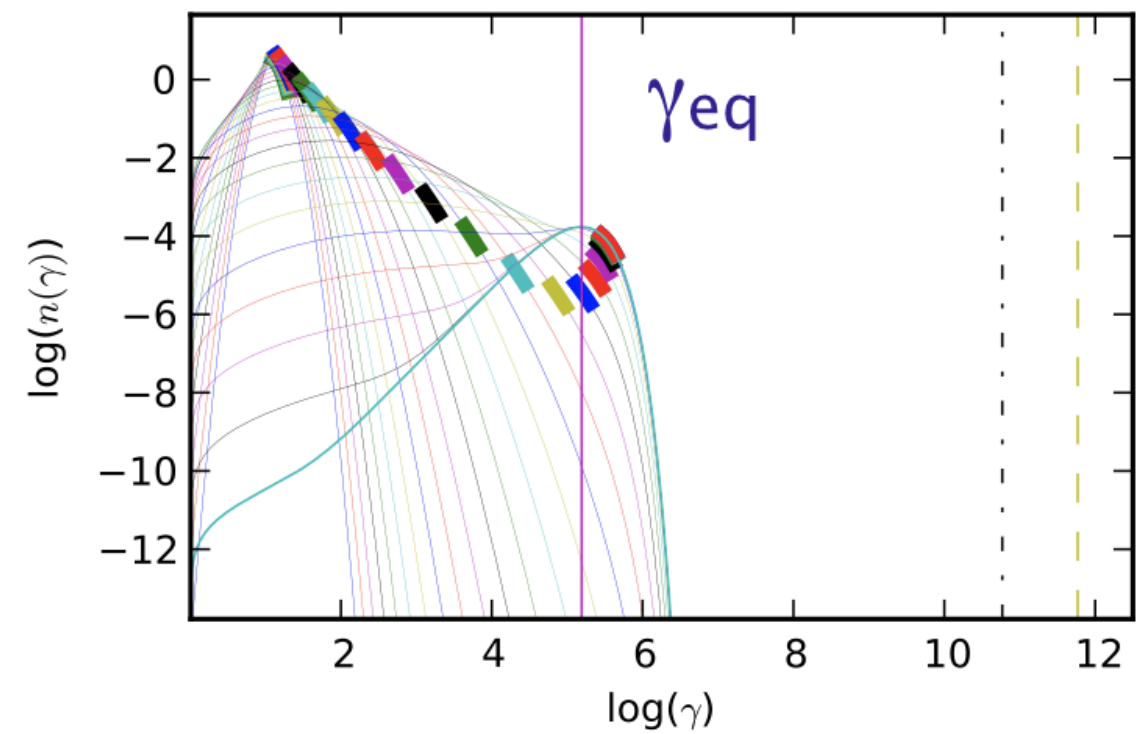
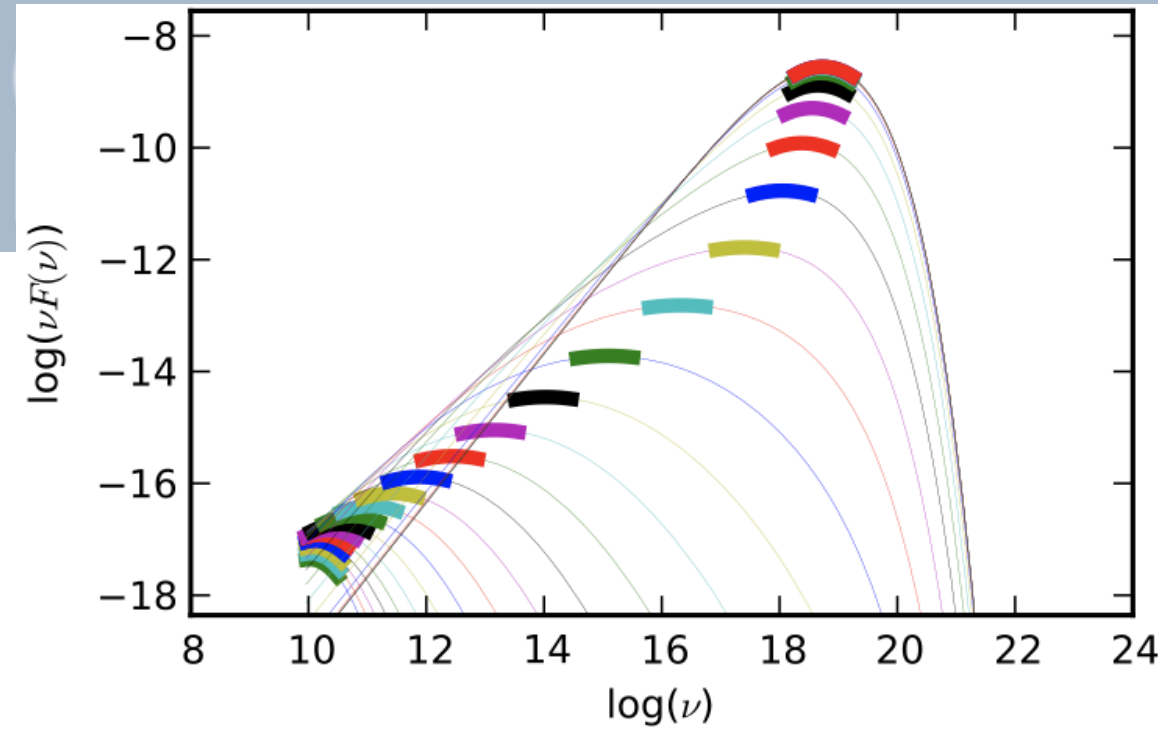
- Transient jets with $BH \sim M_{Sun}$, and complex interaction between accretion and **jet formation**. Crucial to understand processes of jets at AGN scales, and for understanding mechanisms such as BZ
- **CTA sensitivity will allow discriminate better between leptonic and lepto-hadronic models**
- **Synergy with SKA will have strong implication on jets energetic and Magnetic field topology**
- **CTA+SKA Jet feeding (BZ,etc)**
- Recent iXPE measurements confirm pol. angle parallel to the jet axis as in the case of blazars!



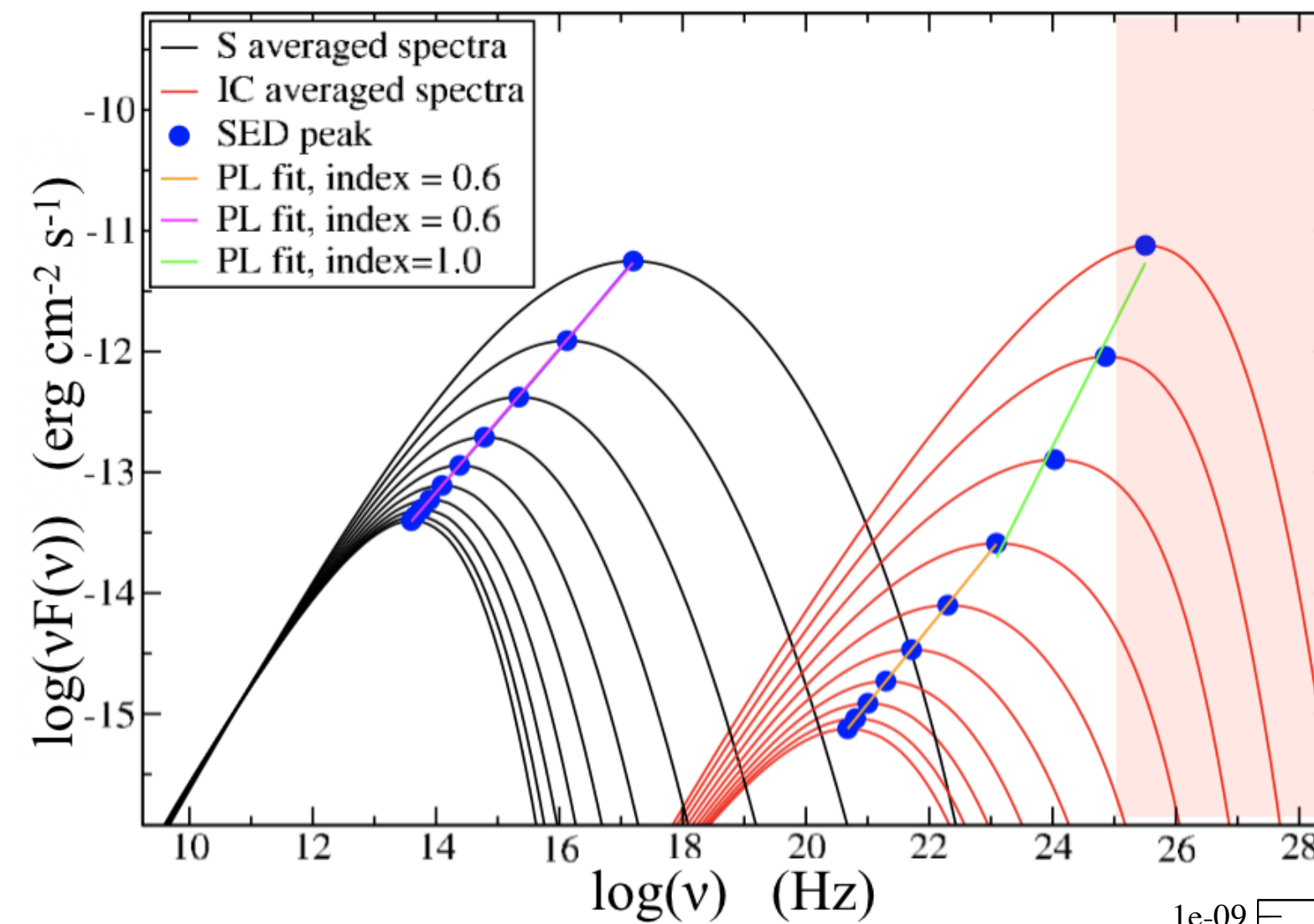
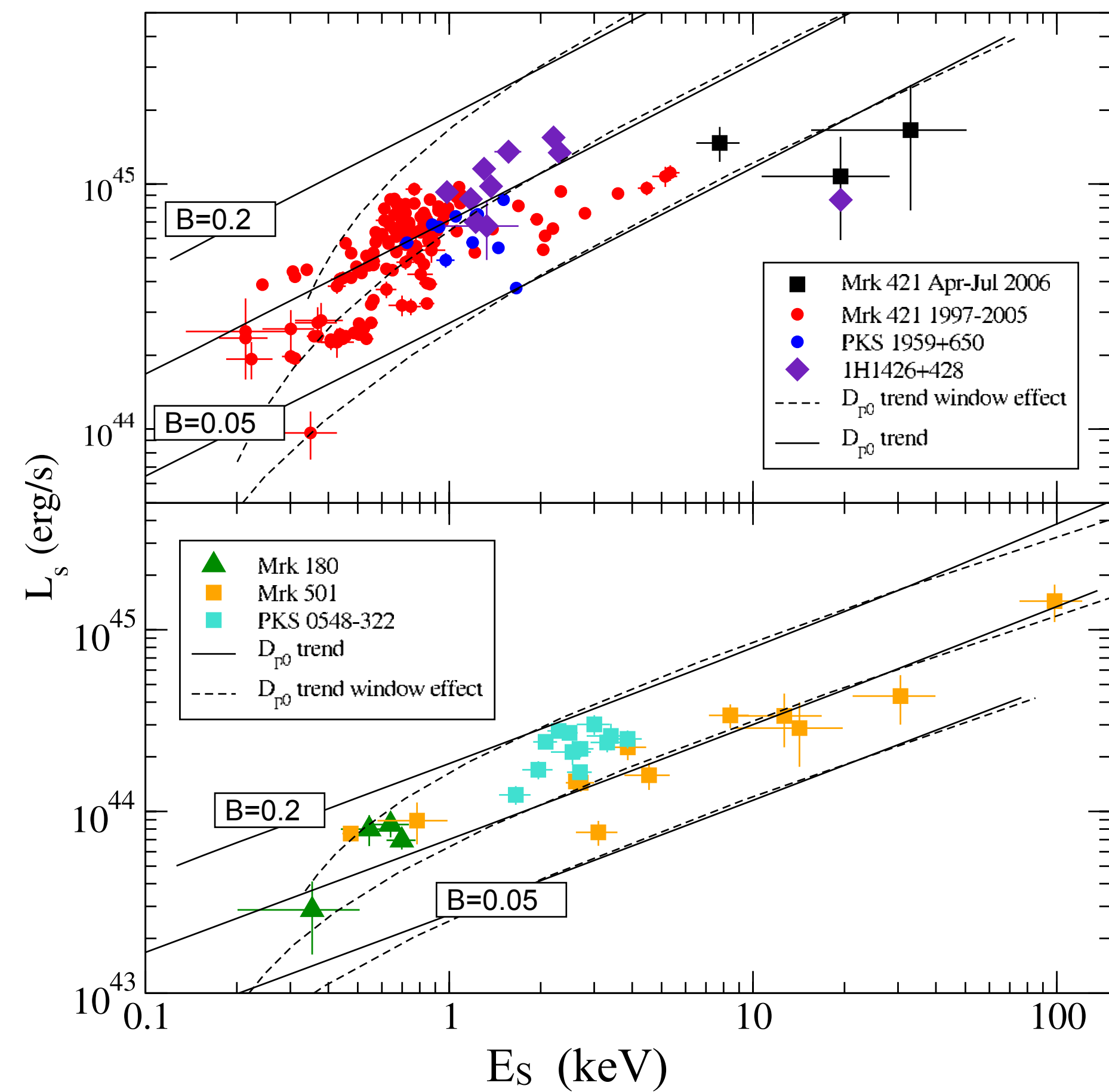
Ep-vs-curvature

Tramacere+2011

Curvature is at acceleration dominated regime!

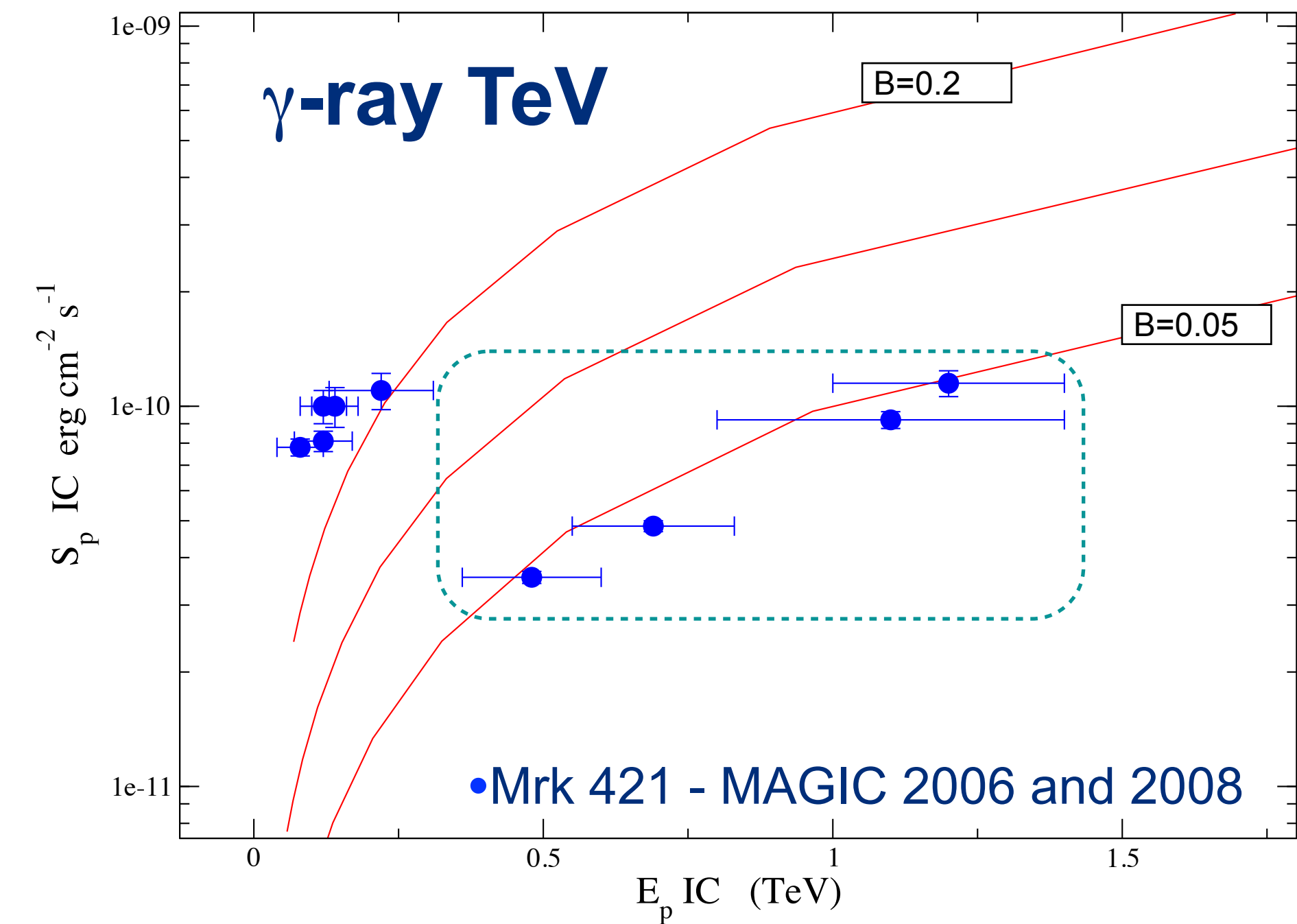


Lp-vs-curvature



Tramacere+2011

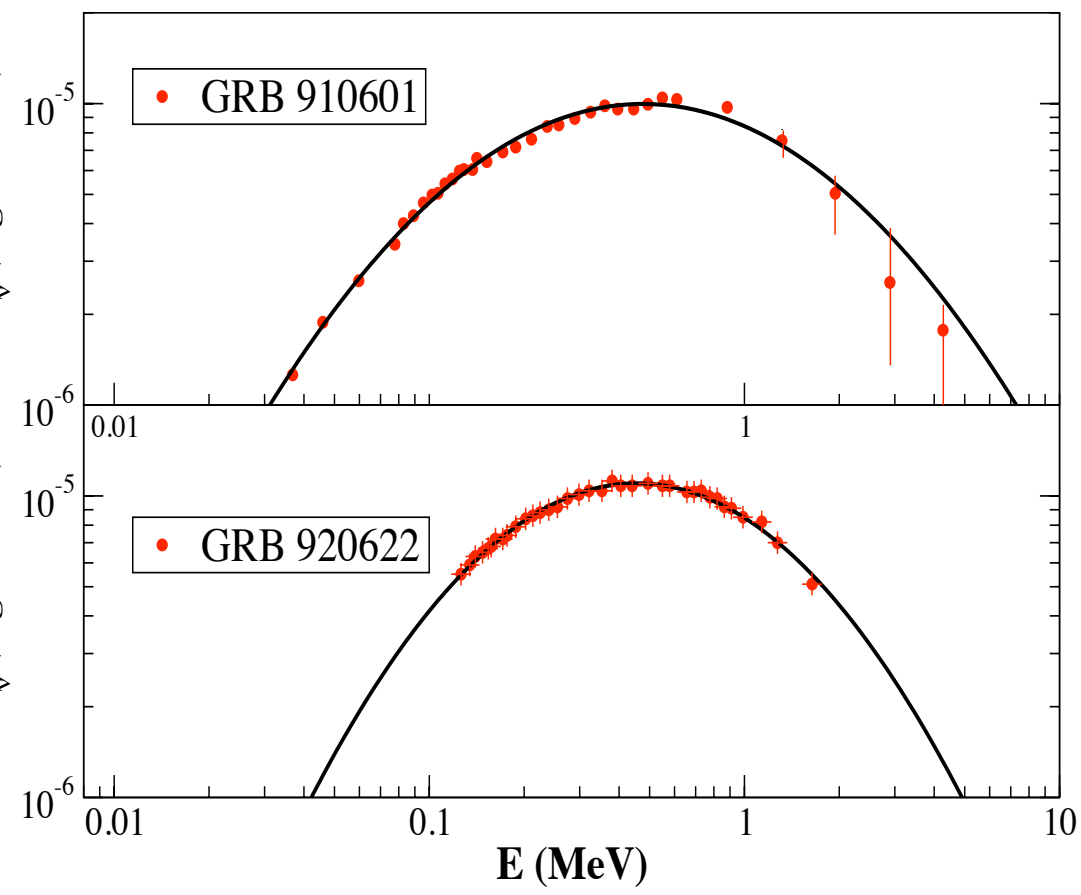
JetSeT



see satoshi

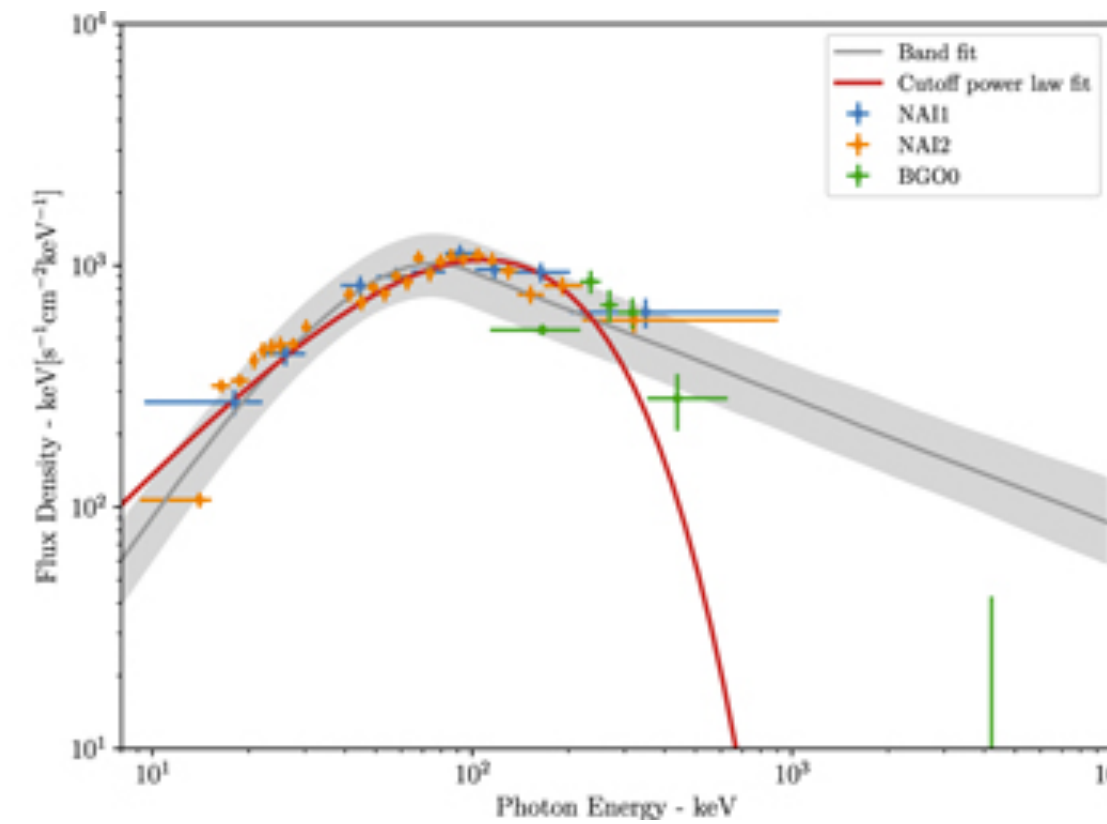
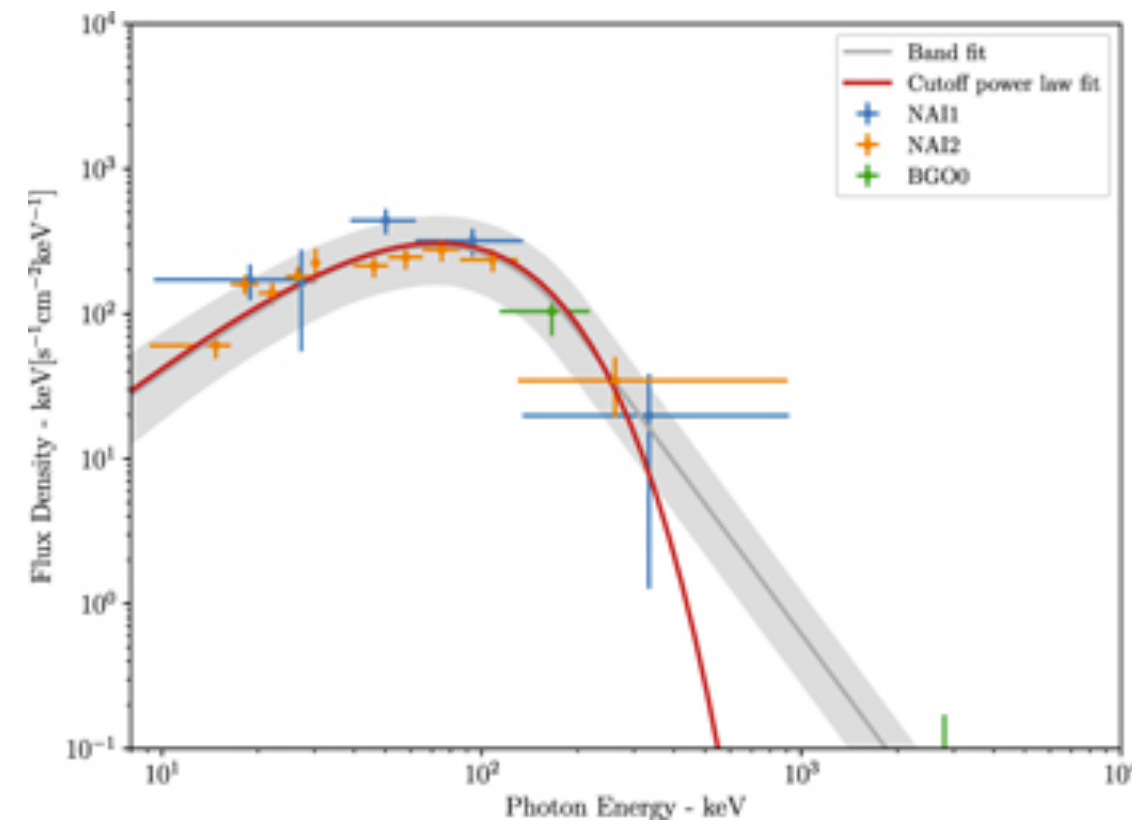
prompt spectral features

Massaro, Grindlay, Paggi 2010

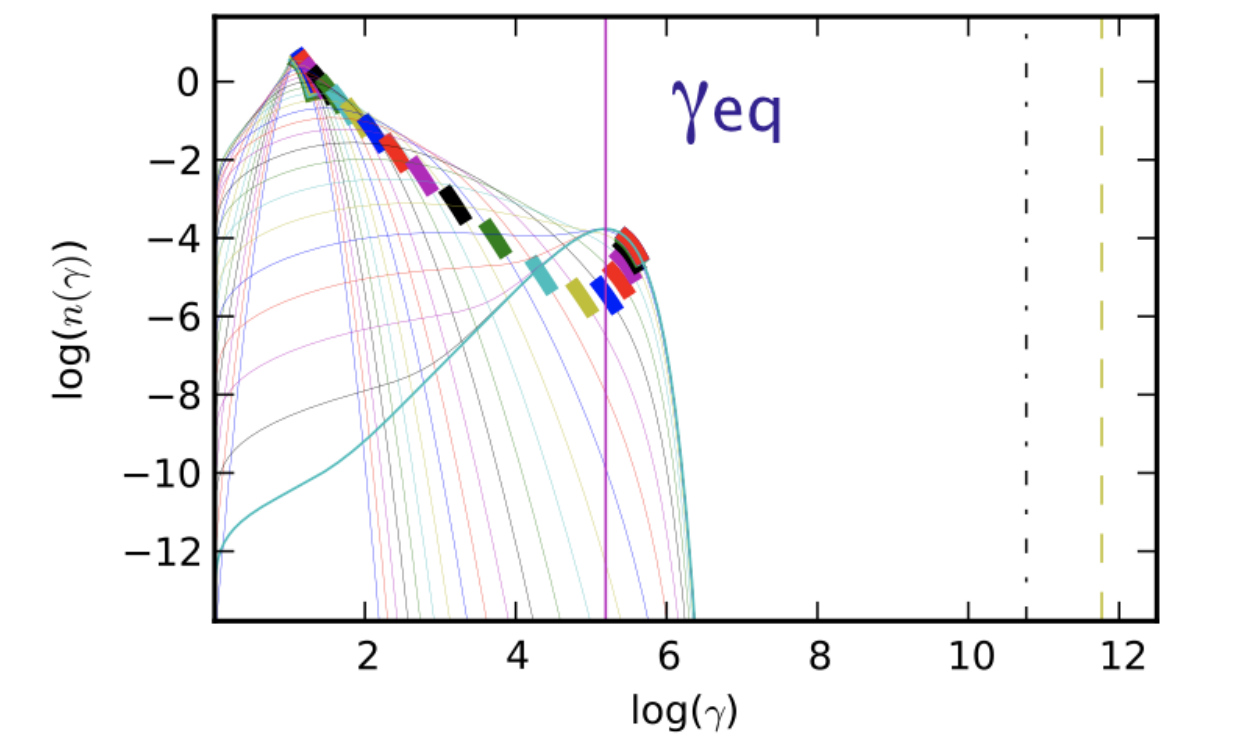
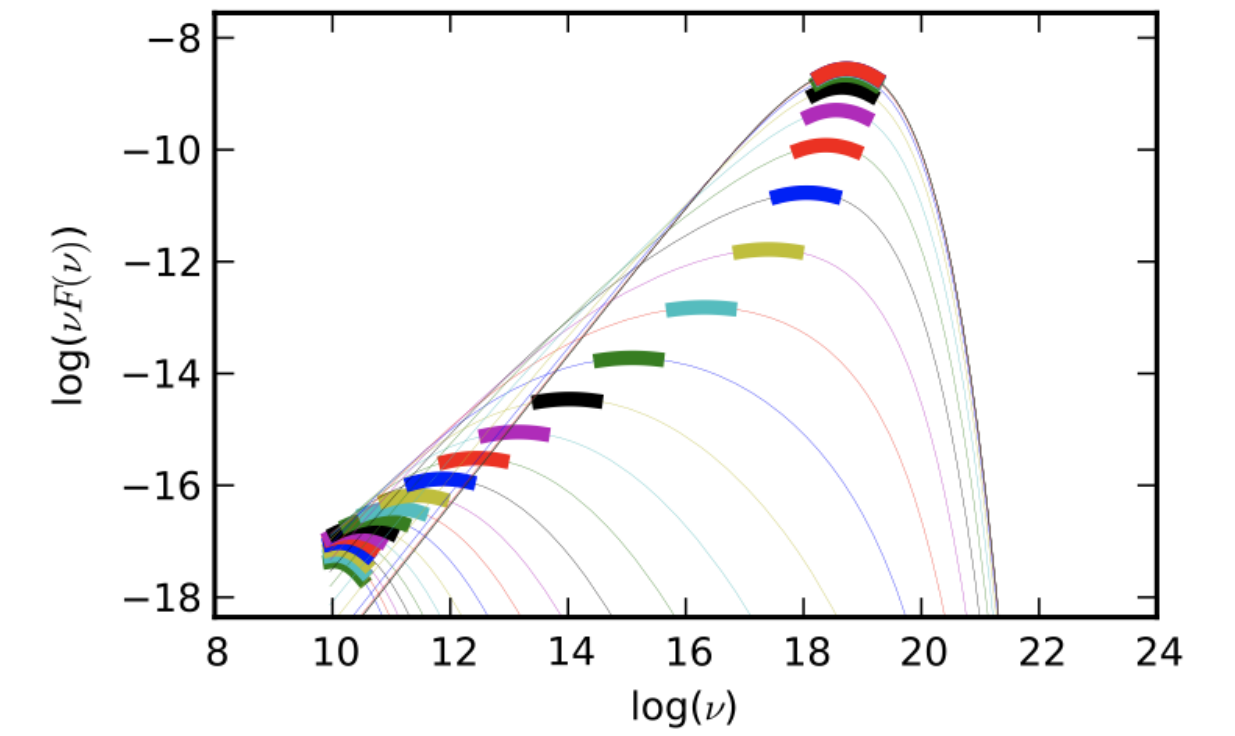


- $b \sim [0.6-0.8]$ compatible with equilibrium, hinting for larger magnetic fields and shorter t_{acc} compared to blazars

Li 2022



Tramacere 2011



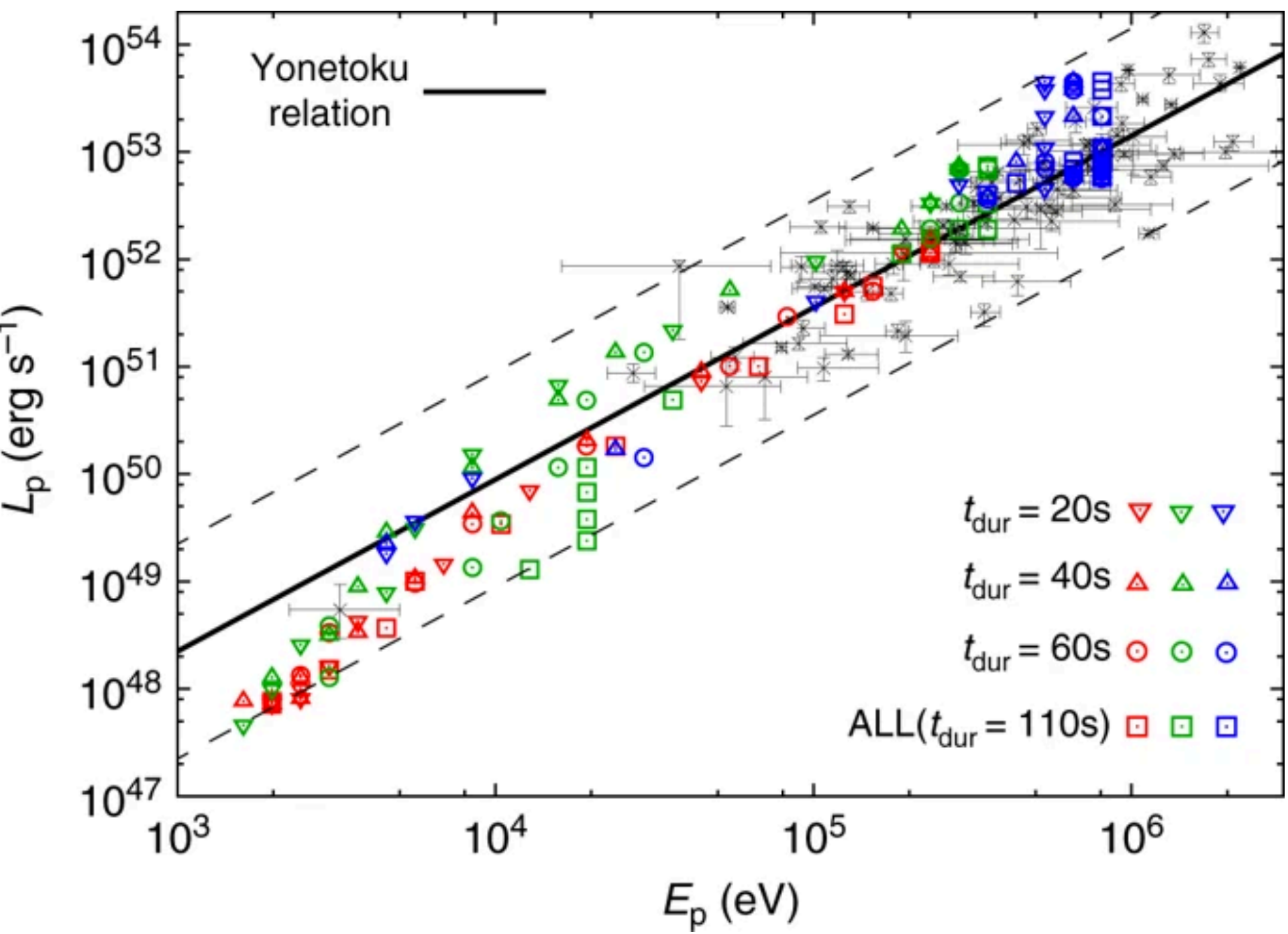
recent *iXPE* polarization measurements with $PD \sim 10\%$ similar to Mrk 501 (see Hancheng)

see satoshi

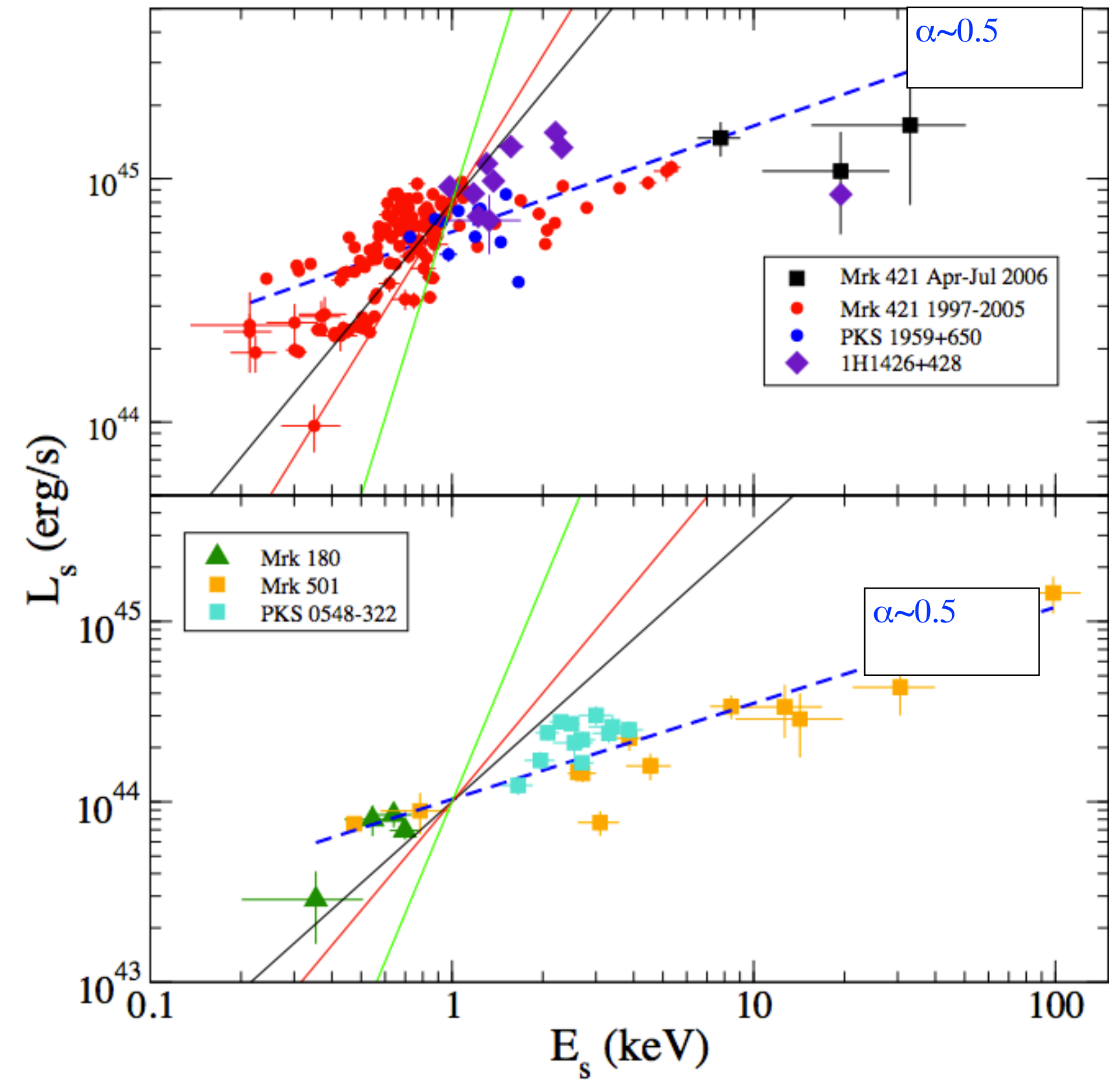
prompt $L_{p,iso} \sim E_p^\alpha$

- index of $\alpha \sim 2$ is compatible with B as main driver
- index of $\alpha \sim 1.5$ is compatible with γ_{3p} increasing keeping $N(\gamma_{3p})$ constant

Ito+ 2019 $\alpha \sim [1.5-2.0]$



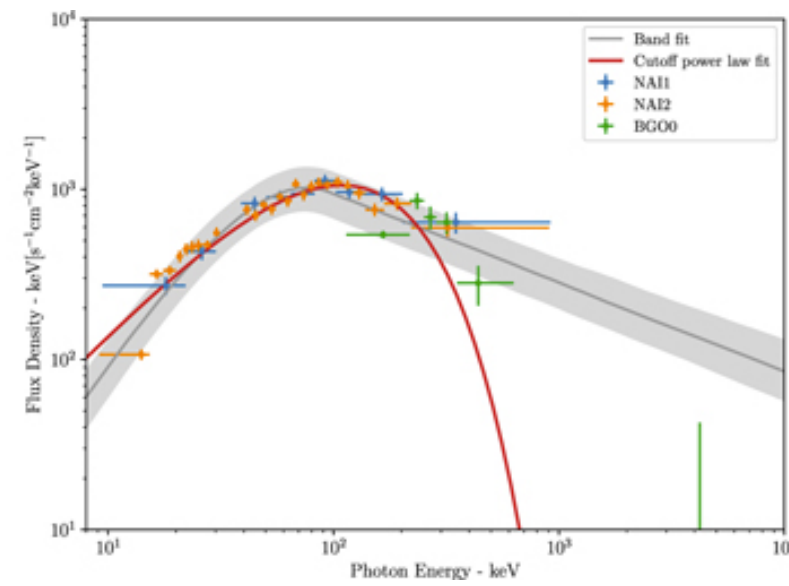
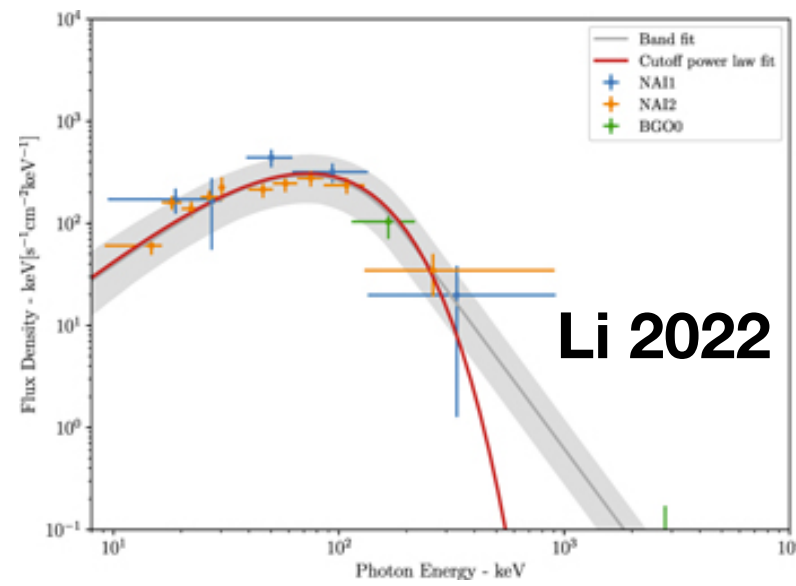
Tramacere 2011



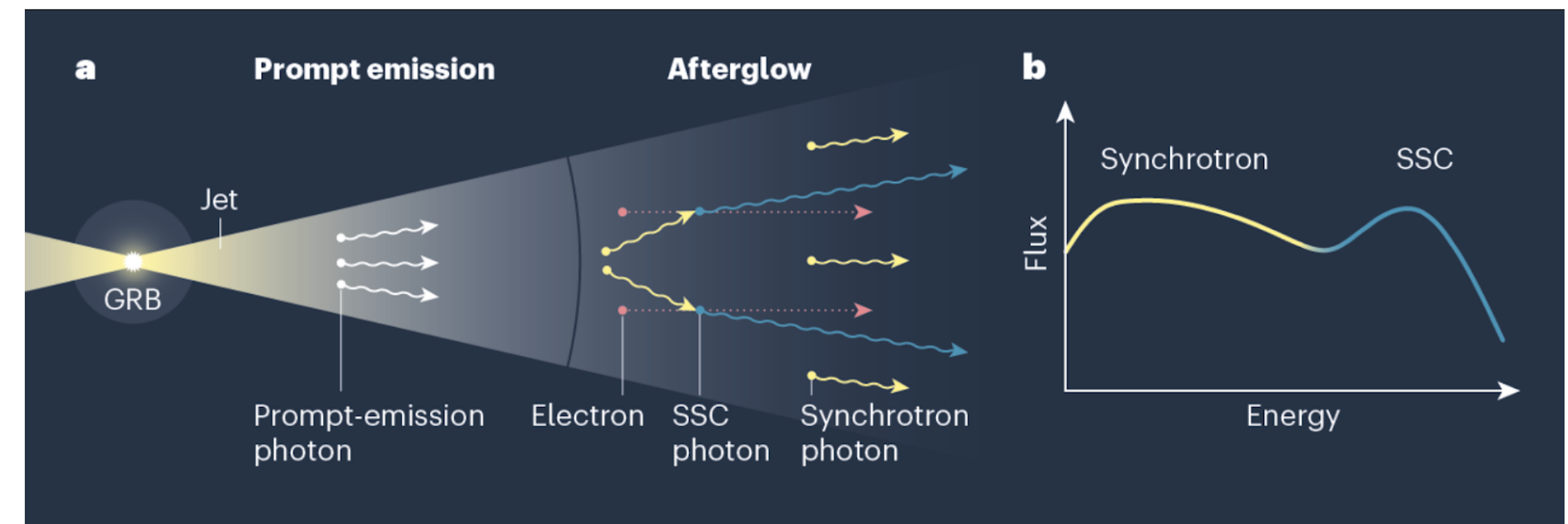
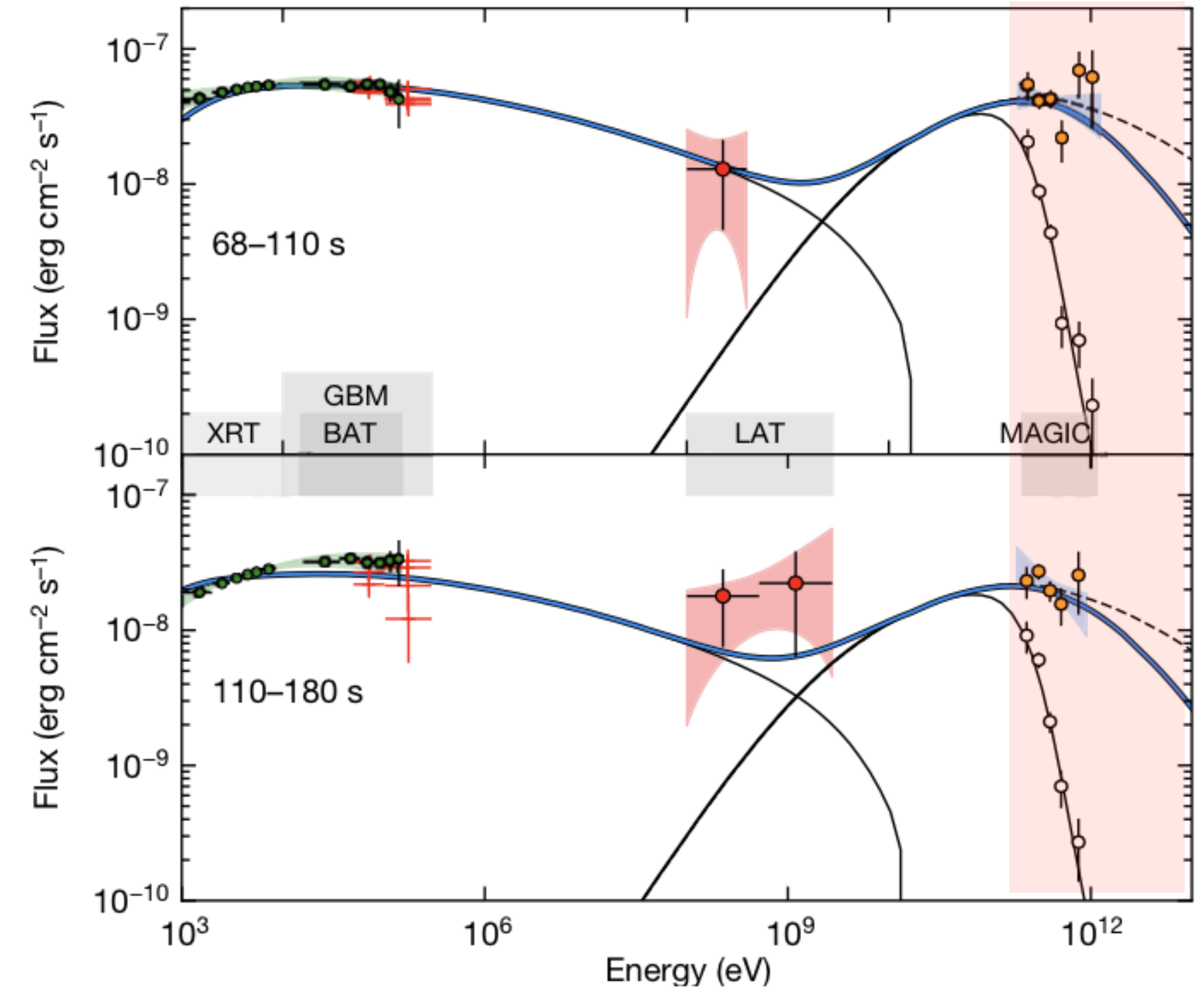
see satoshi

prompt $L_{p,iso} \sim E_p^\alpha$

- $b \sim [0.6-0.8]$ compatible with equilibrium, hinting for larger magnetic fields and shorter t_{acc} compared to blazars
- index of $\alpha \sim 2$ is compatible with B as main driver
- index of $\alpha \sim 1.5$ is compatible with γ_{3p} increasing keeping $N(\gamma_{3p})$ constant



Magic Nature 2019 GRB 190114C **afterglow**



**CTA can help in filling the gap between macro
and micro physics**

**But, we need to have MW/MM simultaneous observations
in particular X-ray (possibly with polarimetry)
we need to look at jets in different environments**

