New constraints on extended Higgs sectors from the trilinear Higgs coupling

Based on

arXiv:2202.03453 (accepted in PRL) in collaboration with Henning Bahl and Georg Weiglein,

(as well as arXiv:1903.05417 (PLB), 1911.11507 (EPJC) in collaboration with Shinya Kanemura)

Johannes Braathen

The 19th Workshop of LHC Higgs Working Group

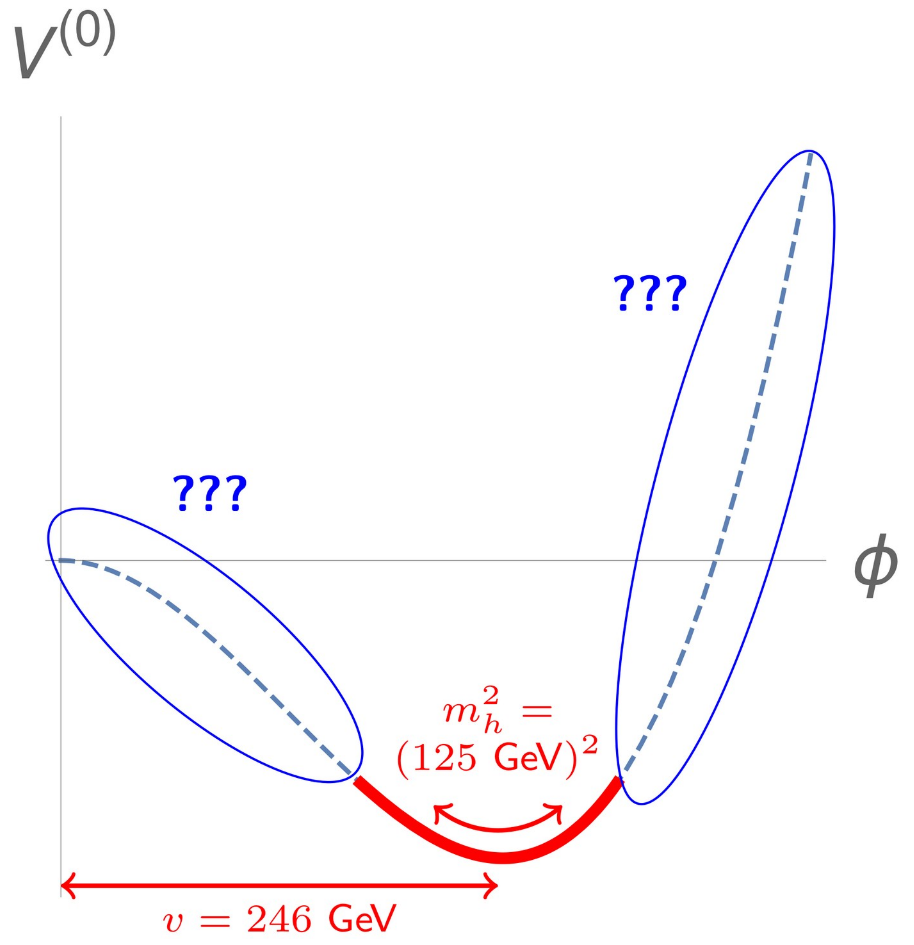
CERN | November 28-30, 2022

HELMHOLTZ RESEARCH FOR GRAND CHALLENGES

DESY.

Why study the trilinear Higgs coupling?

- **Probing the Higgs potential:**
  Since the Higgs discovery, the existence of the Higgs potential is confirmed, but at the moment we only know:
  → the location of the EW minimum:
  v = 246 GeV
  → the curvature of the potential around the EW minimum:
  m_h = 125 GeV
  However we still don't know the **shape** of the potential, away from EW minimum → **depends on λ_{hhh}**



- **λ_{hhh} determines the nature of the EWPT!**
  ⇒ O(20%) deviation of λ_{hhh} from its SM prediction needed to have a strongly first-order EWPT → necessary for EWBG [Grojean, Servant, Wells '04], [Kanemura, Okada, Senaha '04]

- ***New in this talk***: **studying λ_{hhh} can also serve to constrain the parameter space of BSM models!**

BSM contributions to λ_{hhh}

The Two-Higgs-Doublet Model

- 2 $SU(2)_L$ doublets $\Phi_{1,2}$ of hypercharge ½
- CP-conserving 2HDM, with softly-broken Z_2 symmetry ($\Phi_1 \to \Phi_1$, $\Phi_2 \to -\Phi_2$) to avoid tree-level FCNCs

$$V^{(0)}_{\rm 2HDM} = m_1^2|\Phi_1|^2 + m_2^2|\Phi_2|^2 - m_3^2(\Phi_2^\dagger\Phi_1 + \Phi_1^\dagger\Phi_2)$$
$$+ \frac{\lambda_1}{2}|\Phi_1|^4 + \frac{\lambda_2}{2}|\Phi_2|^4 + \lambda_3|\Phi_1|^2|\Phi_2|^2 + \lambda_4|\Phi_2^\dagger\Phi_1|^2 + \frac{\lambda_5}{2}\Big((\Phi_2^\dagger\Phi_1)^2 + \text{h.c.}\Big)$$
$$v_1^2 + v_2^2 = v^2 = (246\ \text{GeV})^2$$

- **Mass eigenstates**:
  h, H: CP-even Higgs bosons *(h → 125-GeV SM-like state);* A: CP-odd Higgs boson;
  $H^\pm$: charged Higgs boson

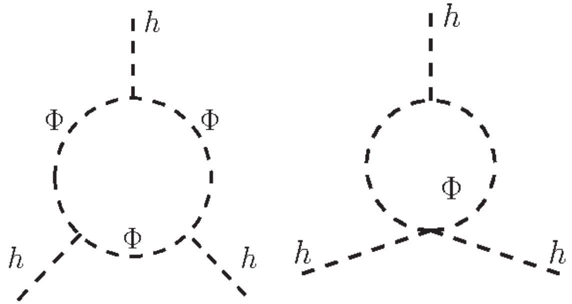
- **BSM parameters**: 3 BSM masses m_H, m_A, $m_{H\pm}$, BSM mass scale M (defined by $M^2 \equiv 2m_3^2/s_{2\beta}$), angles α (CP-even Higgs mixing angle) and β (defined by $\tan\beta = v_2/v_1$)

- **BSM-scalar masses** take form $m_\Phi^2 = M^2 + \tilde\lambda_\Phi v^2\,, \quad \Phi \in \{H, A, H^\pm\}$

- We take the **alignment limit α=β-π/2** → all Higgs couplings are SM-like at tree level
  → compatible with current experimental data!

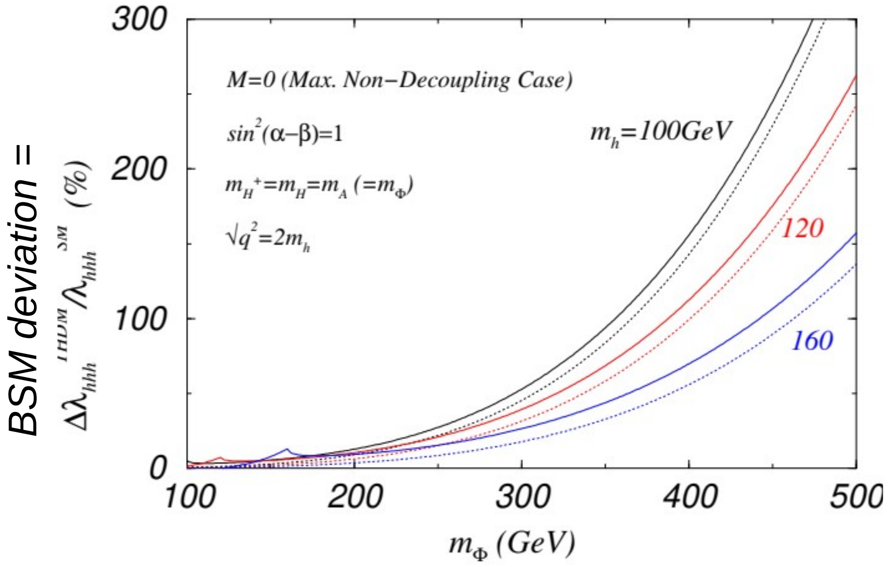
Non-decoupling effects in λ_{hhh}

- First investigation of 1L BSM contributions to λ_{hhh} in 2HDM: [Kanemura, (Kiyoura), Okada, Senaha, Yuan '02, '04]



$$g_{hh\Phi\Phi} = -\frac{2(M^2 - m_\Phi^2)}{v^2}$$

$$(\Phi \in \{H, A, H^\pm\})$$



- **Deviations of tens/hundreds of % from SM possible,** for **large $g_{h\Phi\Phi}$ or $g_{hh\Phi\Phi}$ couplings**
- **Non-decoupling effects**, now found in various models (2HDM, inert doublet model, singlet extensions, etc.)

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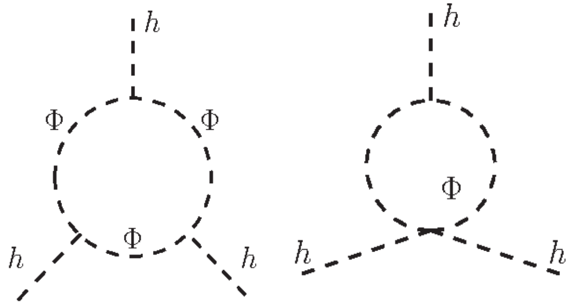

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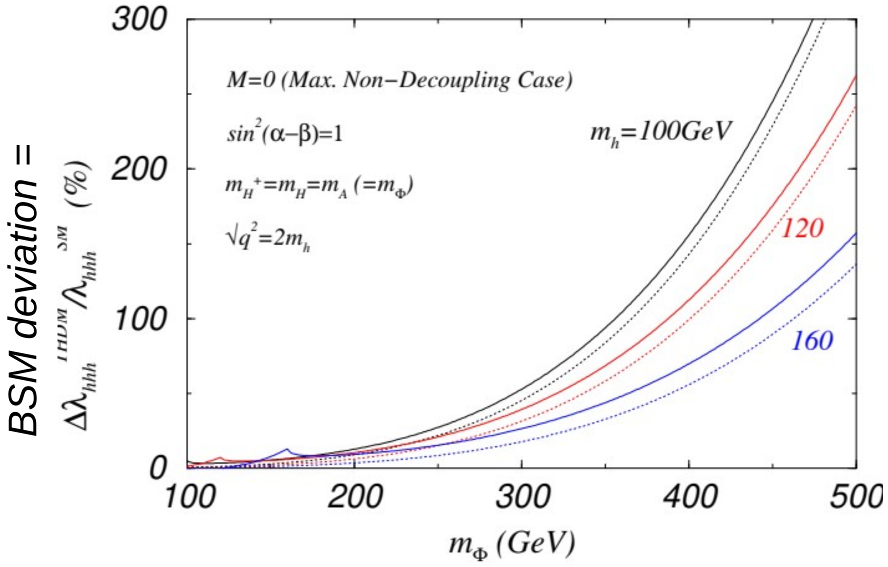
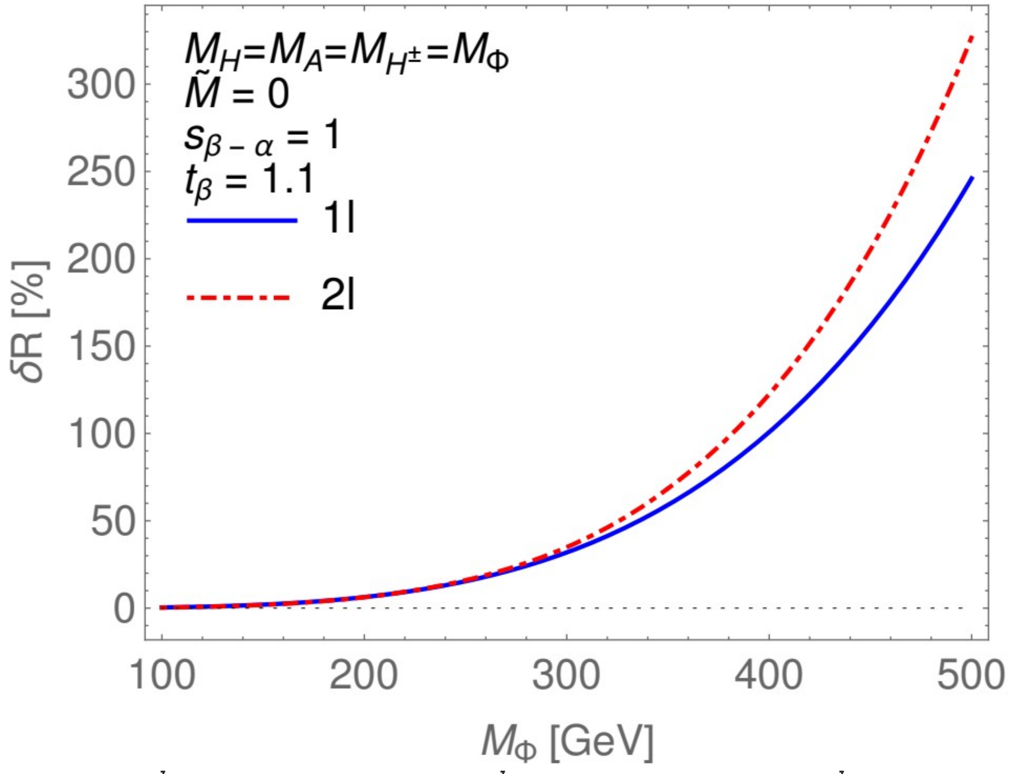
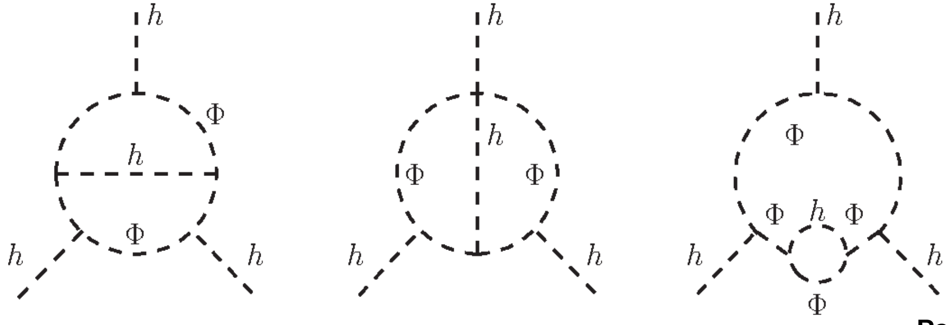


- **Deviations of tens/hundreds of % from SM possible,** for **large $g_{h\Phi\Phi}$ or $g_{hh\Phi\Phi}$ couplings**
- **Non-decoupling effects**, now found in various models (2HDM, inert doublet model, singlet extensions, etc.)

- Non-decoupling effects **confirmed at 2L** in [JB, Kanemura '19] → **leading 2L corrections involving BSM scalars (H,A,H±) and top quark**, computed in effective potential approximation





Constraining the 2HDM with λ_{hhh}

i. *Can we apply the limits on κ_λ, extracted from experimental searches for double-Higgs production, for BSM models?*

ii. *Can large BSM deviations occur for points still allowed in light of theoretical and experimental constraints? If so, how large can they become?*

Can we apply hh-production results for the aligned 2HDM?

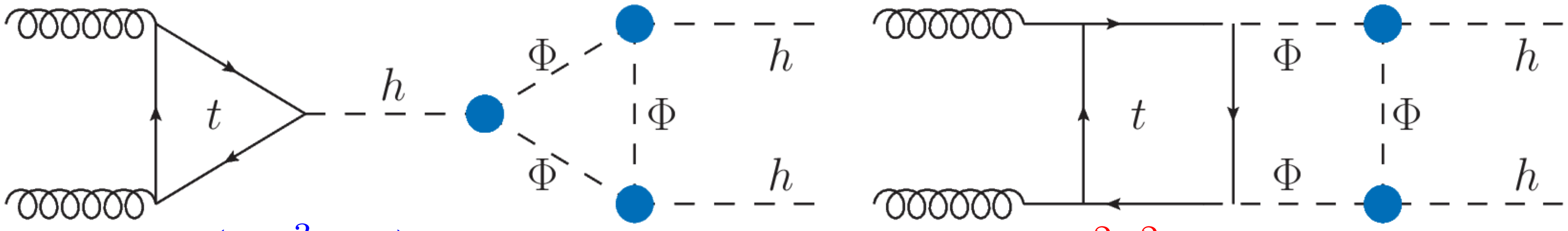
- Current strongest limit on κ_λ are from ATLAS double- (+ single-) Higgs searches

  $-0.4 < \kappa_\lambda < 6.3$ [ATLAS-CONF-2022-050]

  [where $\kappa_\lambda \equiv \lambda_{hhh}/(\lambda_{hhh}^{(0)})^{SM}$]

- What are the *assumptions* for the ATLAS limits?
  - All other Higgs couplings (to fermions, gauge bosons) are SM-like

    → this is **ensured by the alignment** ✓
  - The modification of λ_{hhh} is the only source of deviation of the *non-resonant Higgs-pair production cross section* from the SM



$\propto \mathcal{O}(y_t g_{hh\Phi\Phi}^3)$ ***included*** $\propto \mathcal{O}(y_t^2 g_{hh\Phi\Phi}^2)$ *not included*

→ We **correctly include all leading BSM effects to double-Higgs production, in powers of $g_{hh\Phi\Phi}$, up to NNLO!** ✓

- **We can apply the ATLAS limits to our setting!**

  (Note: BSM resonant Higgs-pair production cross section also suppressed at LO, thanks to alignment)

A parameter scan in the aligned 2HDM

[Bahl, JB, Weiglein `2202.03453`]

- Our strategy:
  1. **Scan BSM parameter space**, keeping only points passing various theoretical and experimental constraints (*see below*)
  2. Identify regions with **large BSM deviations in $\boldsymbol{\lambda_{hhh}}$**
  3. Devise a **benchmark scenario** allowing large deviations and investigate impact of experimental limit on λ_{hhh}
- *Here*: we consider an **aligned 2HDM of type-I**, but similar results expected for other 2HDM types, or other BSM models with extended Higgs sectors
- Constraints in our parameter scan:

experimental

  - SM-like Higgs measurements with `HiggsSignals`
  - Direct searches for BSM scalars with `HiggsBounds`
  - b-physics constraints, using results from [Gfitter group 1803.01853]

  Checked with `ScannerS` [Mühlleitner et al. `2007.02985`]

  - EW precision observables, computed at two loops with `THDM_EWPOS` [Hessenberger, Hollik '16, '22]

theoretical

  - Vacuum stability
  - Boundedness-from-below of the potential

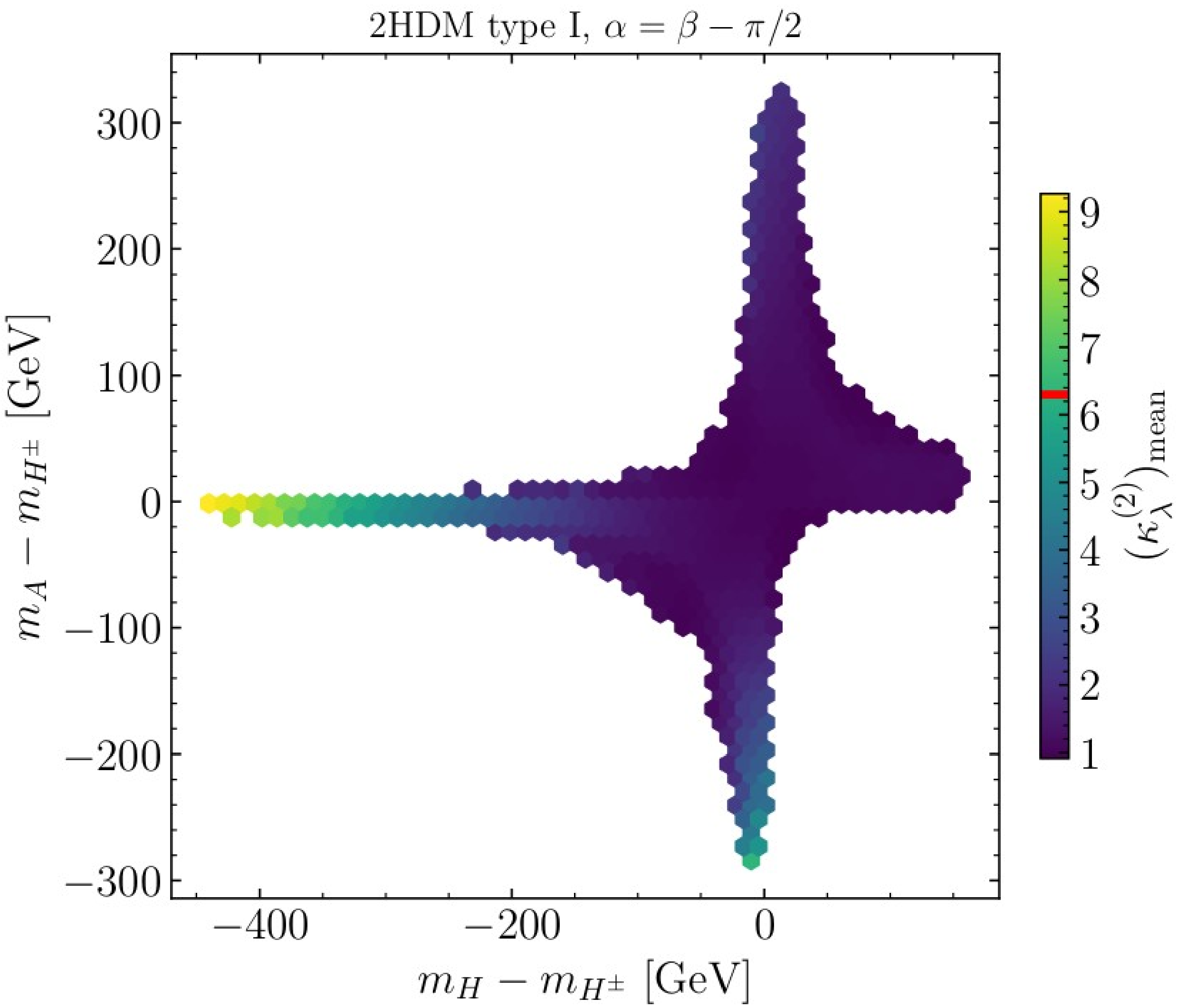
  Checked with `ScannerS`

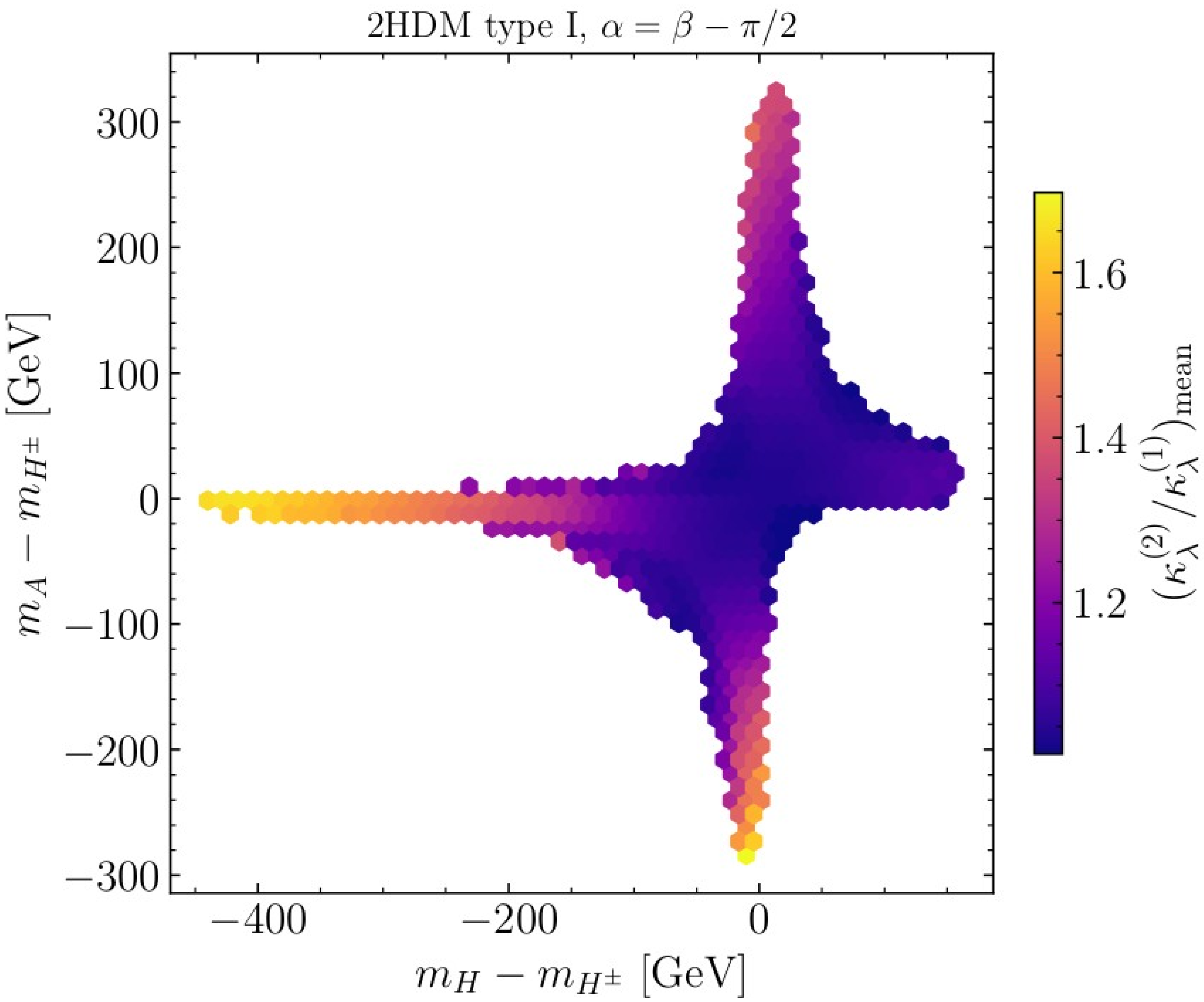
  - NLO perturbative unitarity, using results from [Grinstein et al. `1512.04567`], [Cacchio et al. `1609.01290`]
- For points passing these constraints, we **compute $\boldsymbol{\kappa_\lambda}$ at 1L and 2L**, using results from [JB, Kanemura '19]

Parameter scan results

[Bahl, JB, Weiglein 2202.03453]

<u>Mean value</u> for $\kappa_\lambda^{(2)} = (\lambda_{hhh}^{(2)})^{2HDM} / (\lambda_{hhh}^{(0)})^{SM}$ [left] and $\kappa_\lambda^{(2)}/\kappa_\lambda^{(1)} = (\lambda_{hhh}^{(2)})^{2HDM} / (\lambda_{hhh}^{(1)})^{2HDM}$ [right] in $(m_H - m_{H\pm}, m_A - m_{H\pm})$ plane





NB: all previously mentioned constraints are fulfilled by the points shown here

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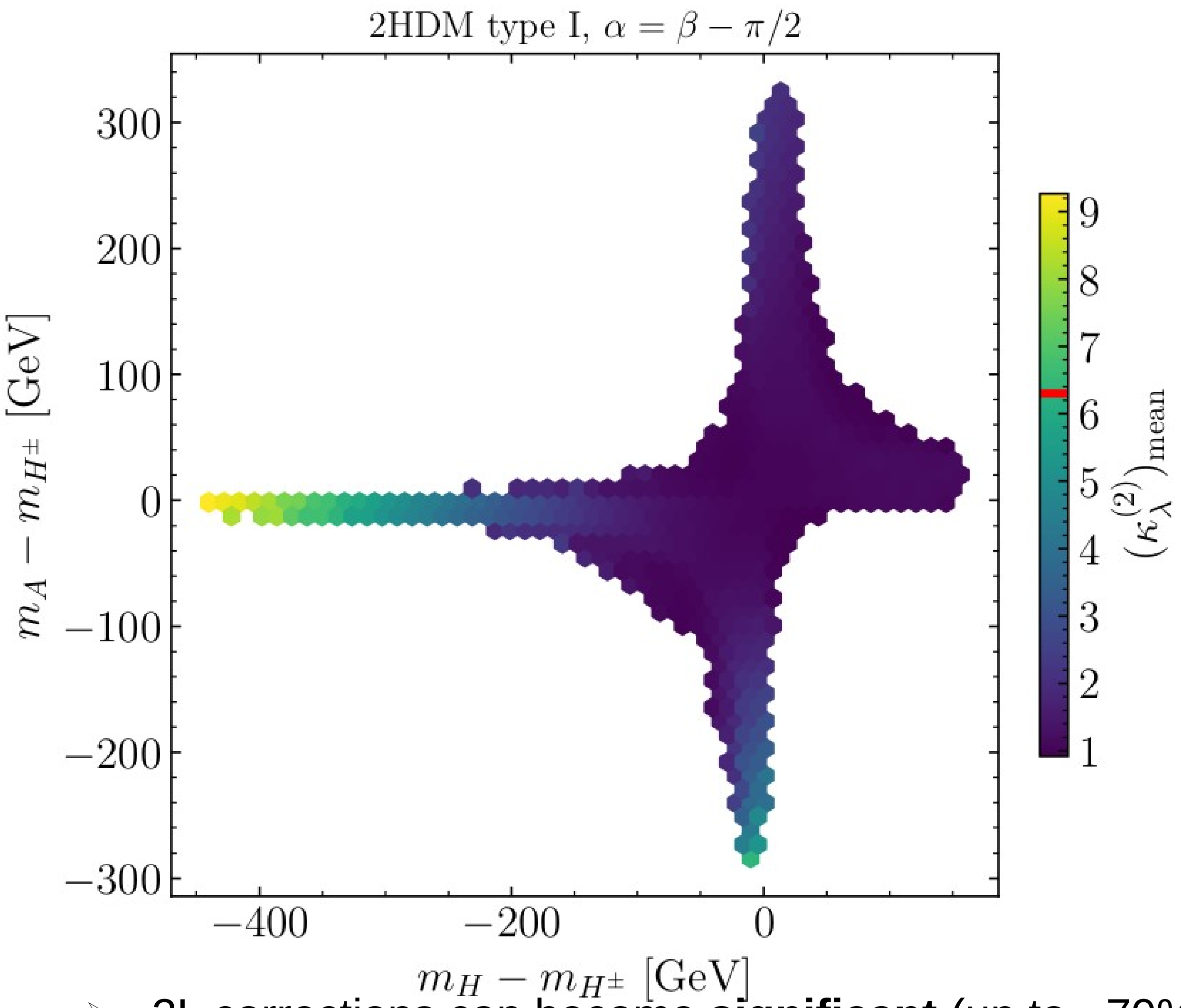

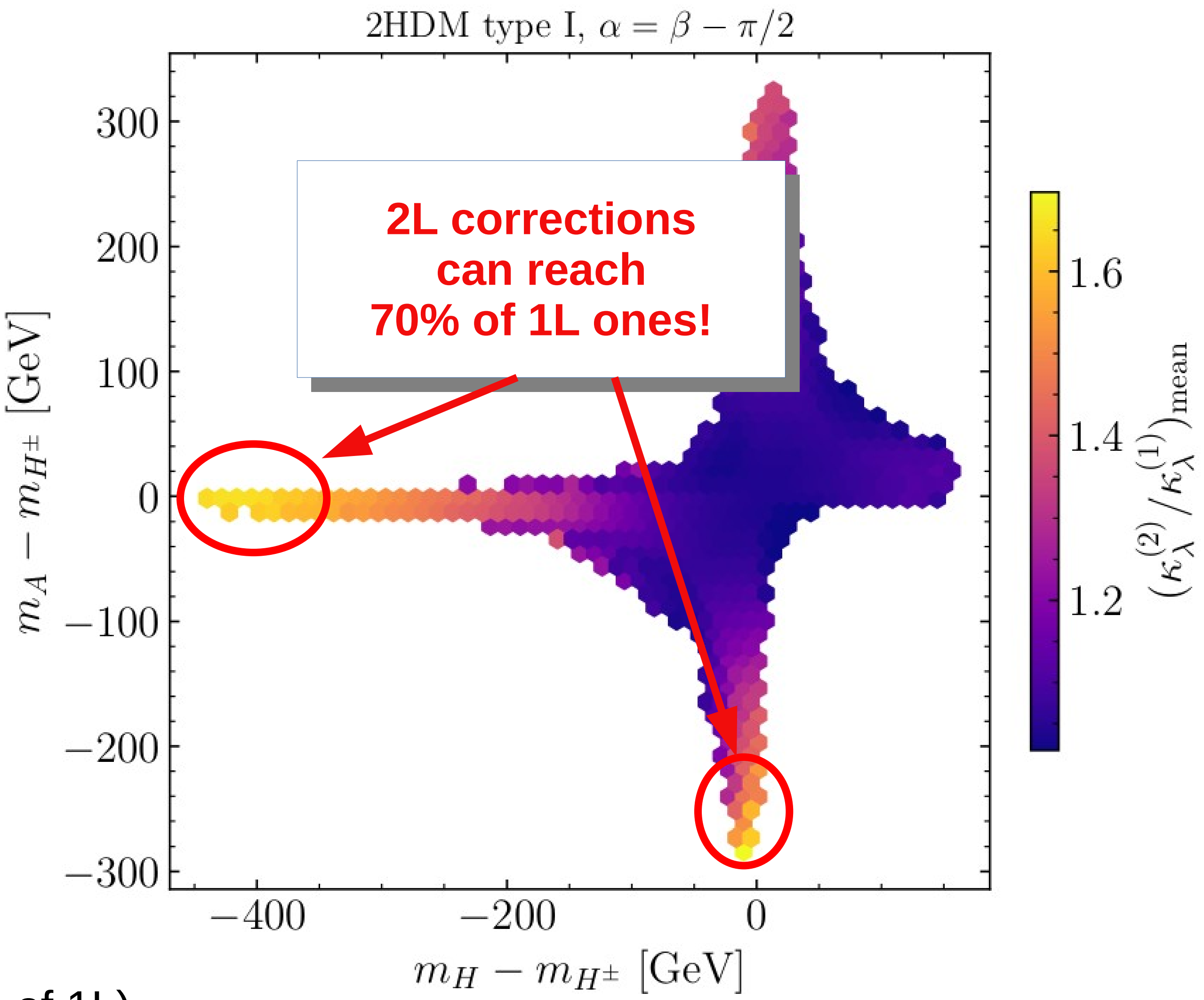

- 2L corrections can become **significant** (up to ~70% of 1L)

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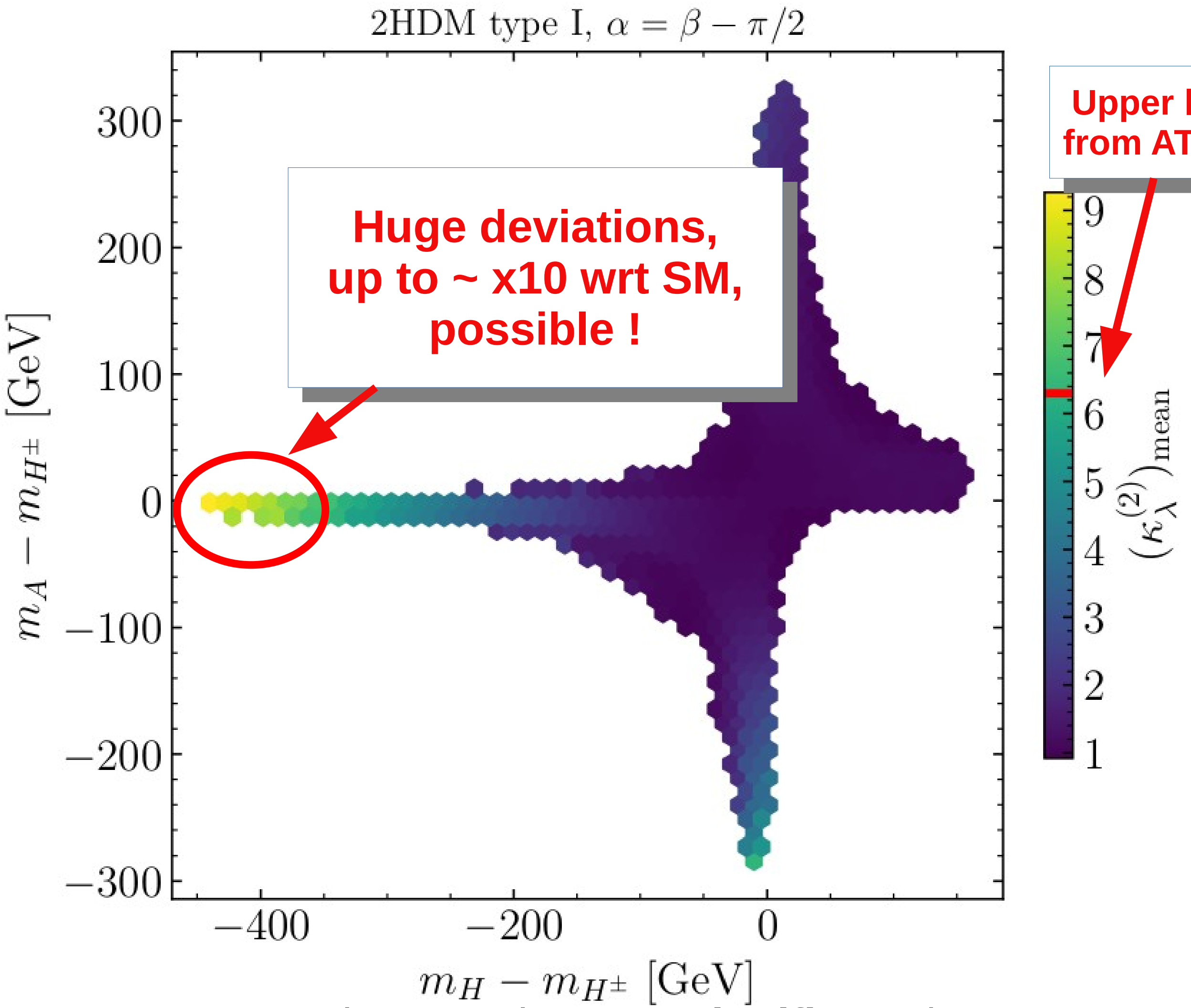
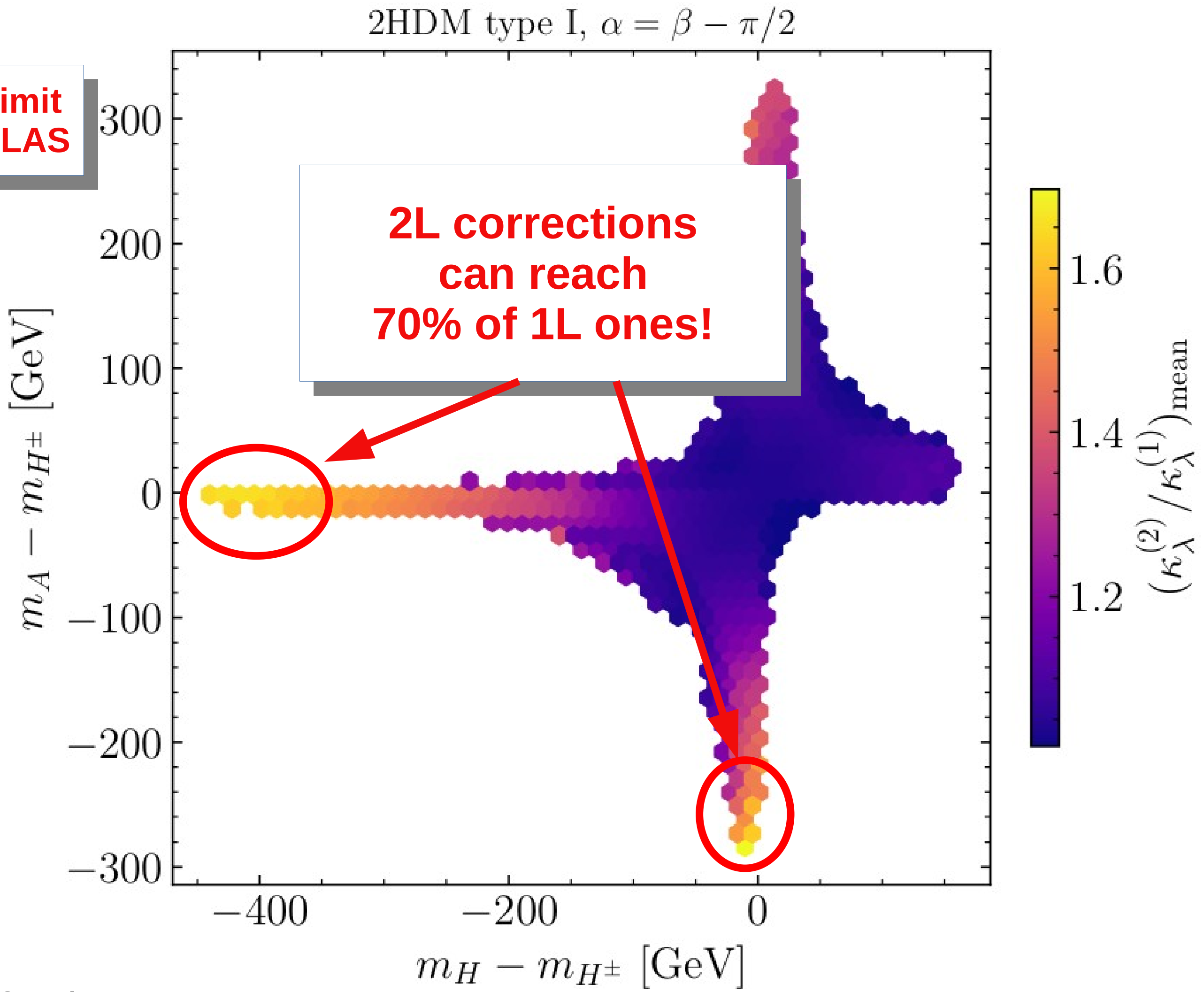




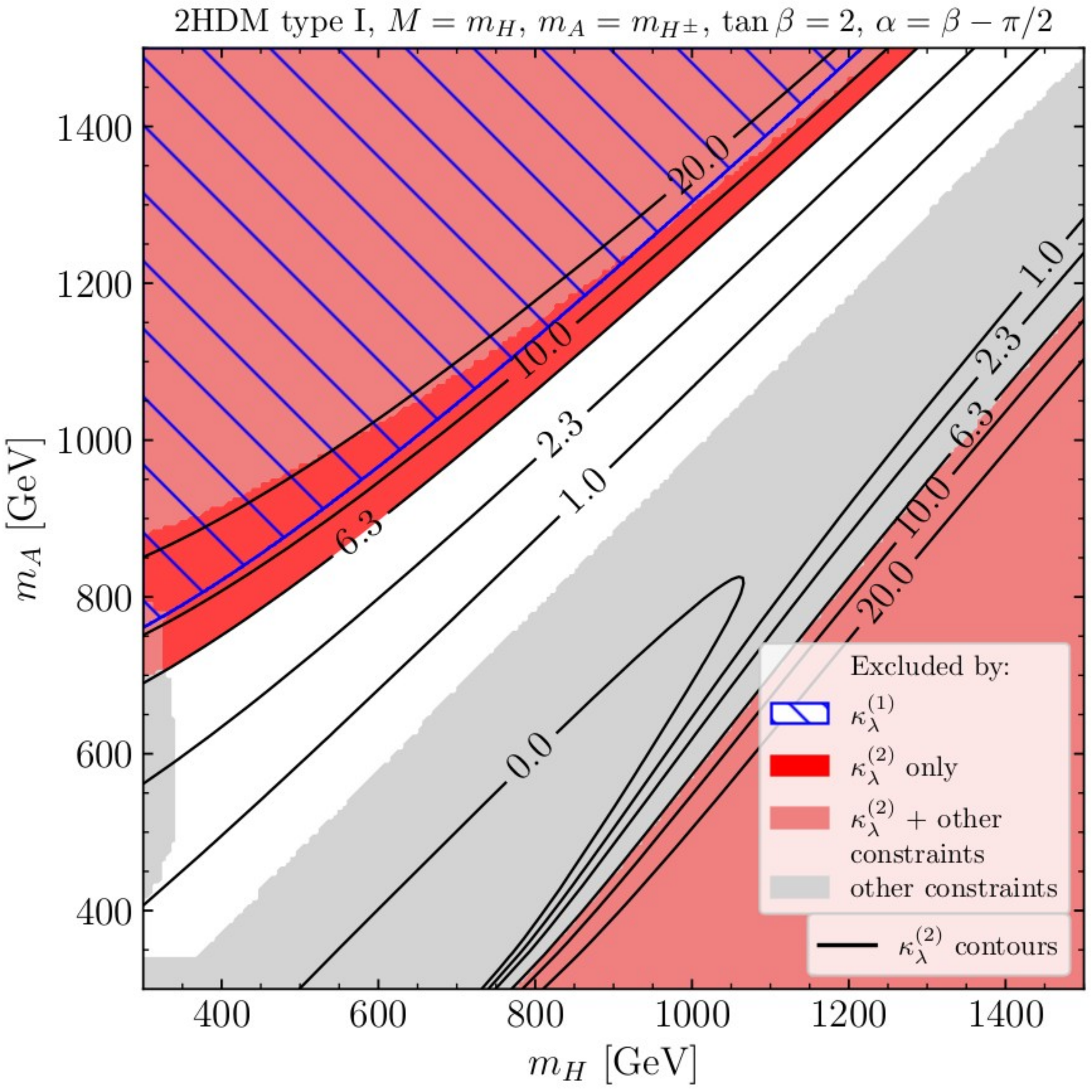
- 2L corrections can become **significant** (up to ~70% of 1L)
- **Huge enhancements** (by a factor ~10) of λ_{hhh} possible for $m_A \sim m_{H\pm}$ and $m_H \sim M$

A benchmark scenario in the aligned 2HDM

[Bahl, JB, Weiglein 2202.03453]

Results shown for aligned 2HDM of type-I, similar for other types (*available in backup*)

We take $m_A=m_{H\pm}$, $M=m_H$, $\tan\beta=2$



- ***Grey area:*** area excluded by other constraints, in particular Higgs physics, boundedness-from-below (BFB), perturbative unitarity
- ***Light red area:*** area excluded both by other constraints (BFB, perturbative unitarity) and by $\kappa_\lambda^{(2)} > 6.3$ [in region where $\kappa_\lambda^{(2)} < -0.4$ the calculation isn't reliable]
- ***Dark red area:*** new area that is **excluded ONLY by $\kappa_\lambda^{(2)}$ > 6.3**. Would otherwise not be excluded!
- ***Blue hatches:*** area excluded by $\kappa_\lambda^{(1)} > 6.3$ → impact of including 2L corrections is significant!

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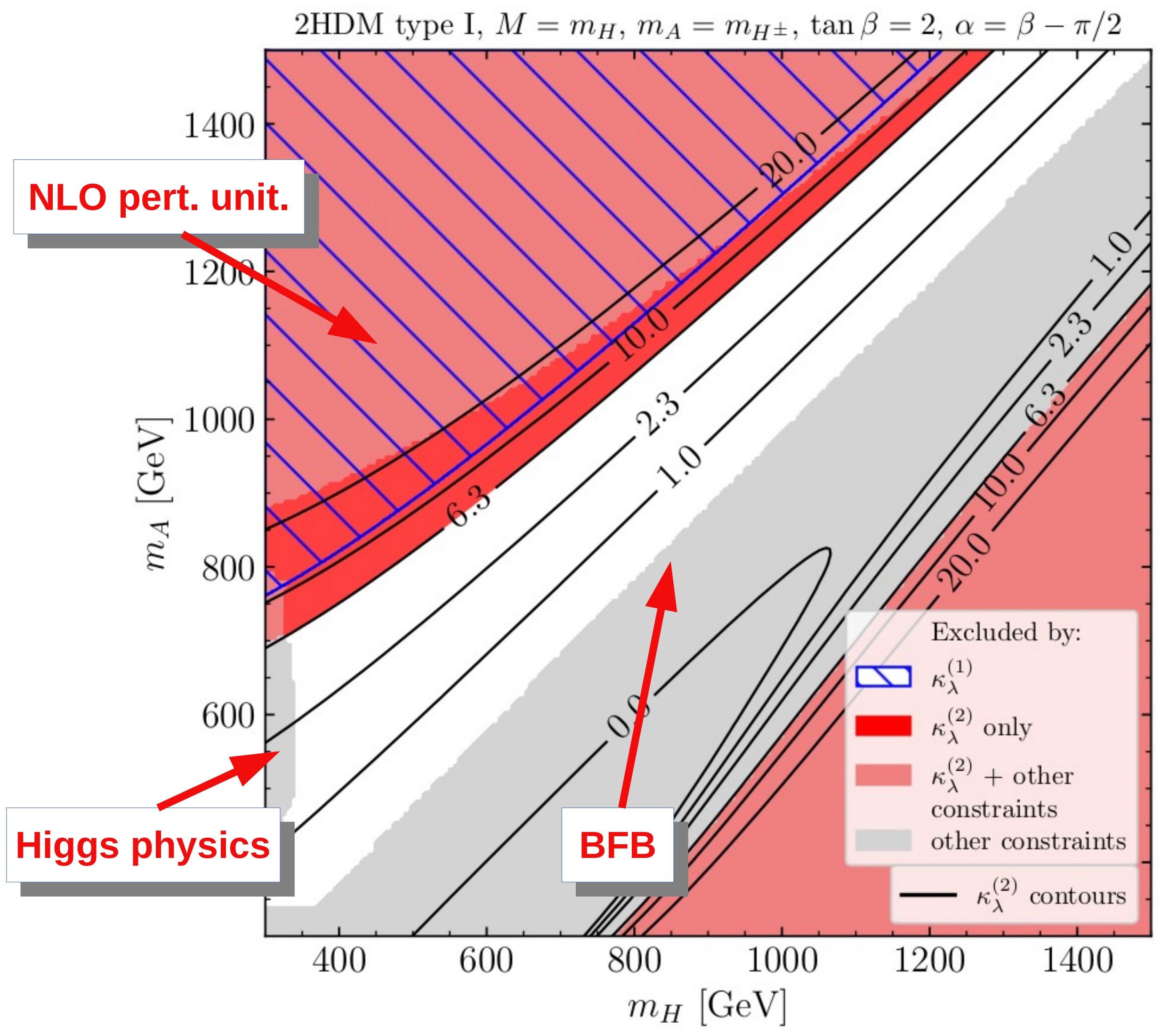

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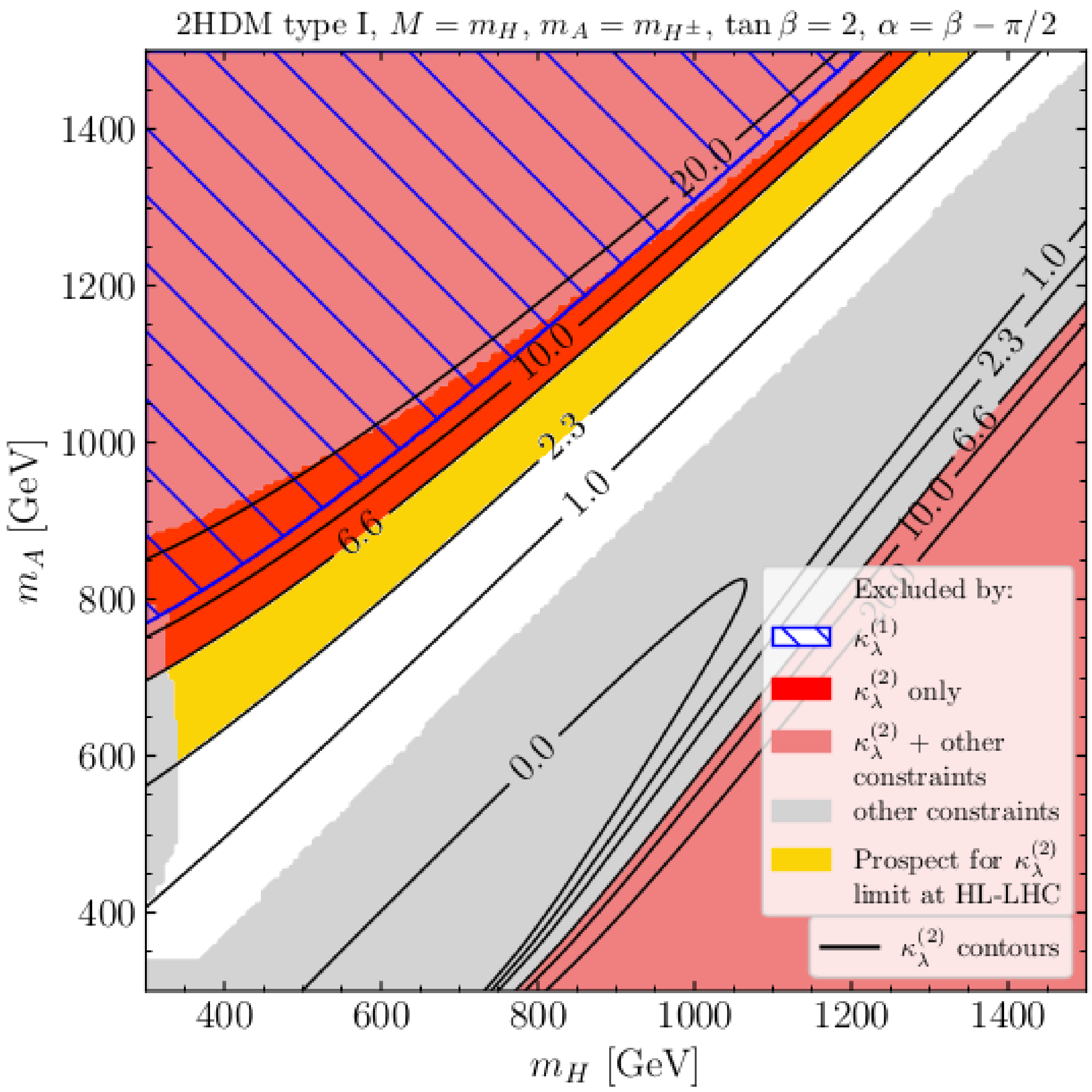
A benchmark scenario in the aligned 2HDM – future prospects

Suppose for instance the upper bound on κ_λ becomes $\kappa_\lambda < 2.3$



- ***Golden area:*** additional exclusion if the limit on κ_λ becomes $\kappa_\lambda^{(2)} < 2.3$ (achievable at HL-LHC)
- Of course, **prospects even better with an e^+e^- collider!**
- Experimental constraints, such as Higgs physics, may also become more stringent, however **not** theoretical constraints (like BFB or perturbative unitarity)

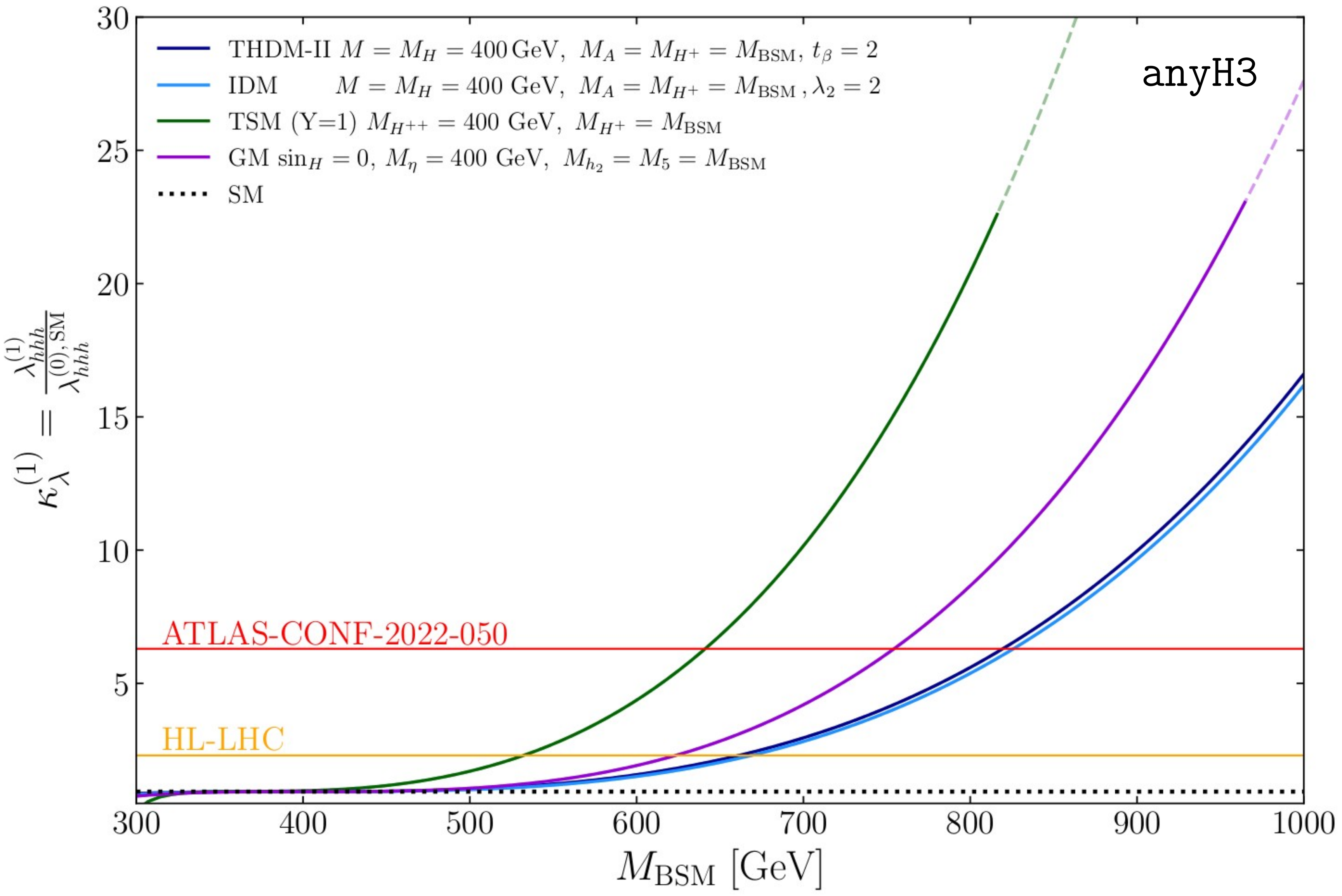
Similar results for other BSM models

Using here predictions for κ_λ computed to full 1L with Python package `anyH3`

[Bahl, JB, Gabelmann, Weiglein *to appear*]

- Example results shown for benchmark scenarios in
  - **2HDM type-II**
  - **Inert Doublet Model**
  - **Complex triplet extension of SM**
  - **Georgi-Machacek model**

- Perturbative unitarity checks
  - Solid: OK ✓
  - Dashed: not OK ✗



Summary

- λ_{hhh} plays a crucial role to understand the shape of the Higgs potential, and probe indirectly signs of New Physics

- λ_{hhh} can **deviate significantly from SM** prediction (by up to a **factor ~10**), for otherwise theoretically and experimentally **allowed points**, due to non-decoupling effects in radiative corrections involving BSM scalars

- Current experimental bounds on λ_{hhh} can **already exclude significant parts of otherwise unconstrained BSM parameter space**, and future prospects even better! Inclusion of 2L corrections [JB, Kanemura ‘19] has significant impact.

- For most of this talk, 2HDM taken as an *example*, but similar results occur for a wider range of BSM models with extended scalar sectors
  → motivates **automating calculations of λ_{hhh}** → **public tool `anyH3` in development**

Thank you for your attention!

Contact

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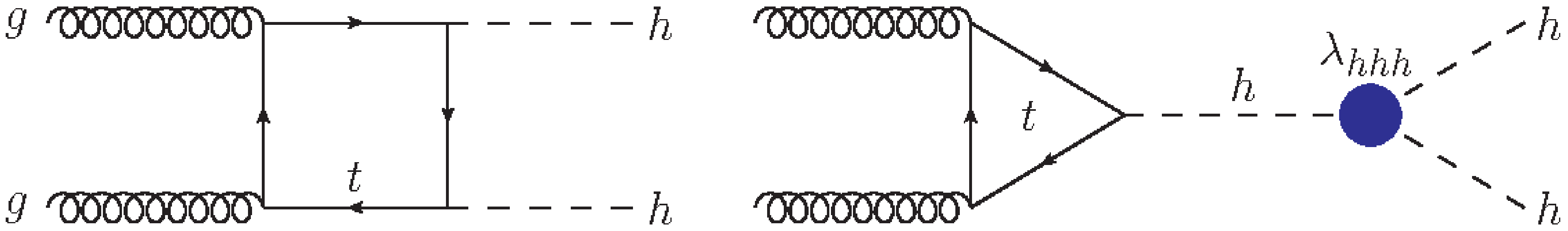
Backup

Accessing λ_{hhh} via double-Higgs production

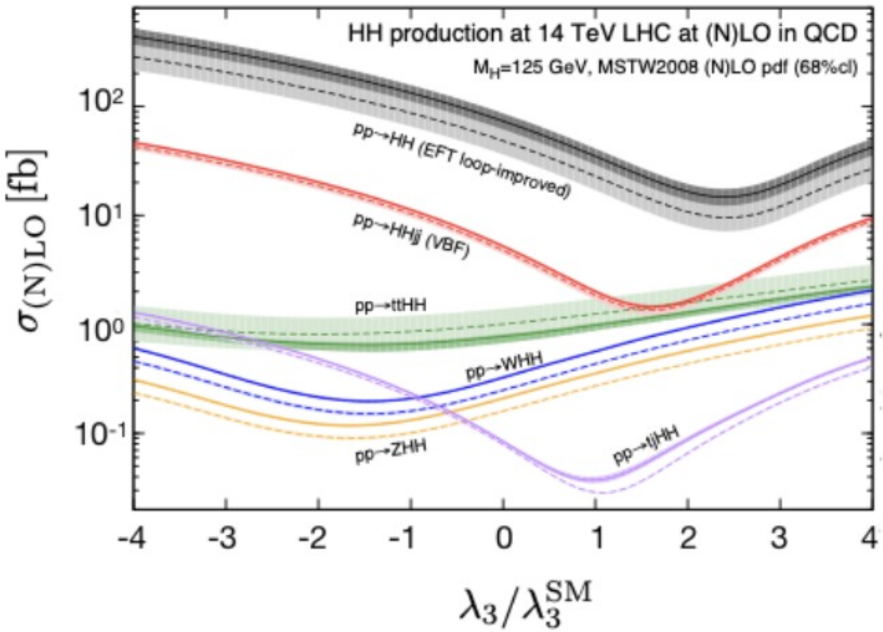
- **Double-Higgs production** → λ_{hhh} enters at LO → **most direct probe of λ_{hhh}**



[Note: Single-Higgs production (EW precision observables) → λ_{hhh} enters at NLO (NNLO)]

- Box and triangle diagrams **interfere destructively**
  → small prediction in SM
  → BSM deviation in λ_{hhh} can **significantly enhance hh-production!**

- Upper limit on hh-production cross-section → **limits on $\kappa_\lambda \equiv \lambda_{hhh}/(\lambda_{hhh}^{(0)})^{SM}$**



[Frederix et al., '14]

Accessing λ_{hhh} via double-Higgs production

- **Dou**

[Not

- Box
  → S
  → B

hh-p

- Uppe
  $\kappa_\lambda \equiv \lambda_{hhh}/(\lambda_{hhh}^{(0)})^{SM}$

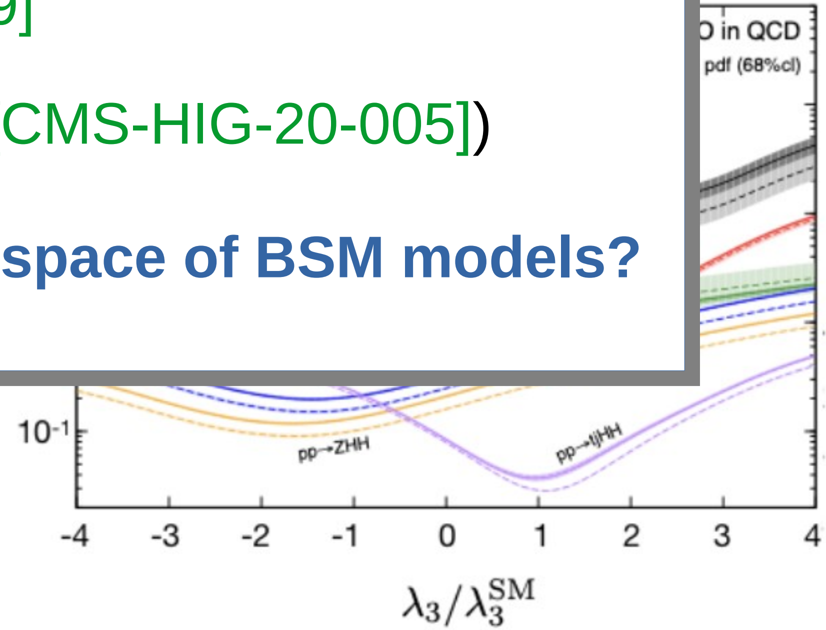
Recent results from ATLAS hh-searches [ATLAS-CONF-2021-052]
yield the limits:

-1.0 < κ_λ < 6.6 at 95% C.L.

→ factor ~2 improvement compared to
previously best ATLAS limits (from single h prod.)
[ATLAS-PHYS-PUB-2019-009]

(CMS recently gave **-2.3 < κ_λ < 9.4** at 95% C.L. [CMS-HIG-20-005])

→ **Can κ_λ now be used to constrain the parameter space of BSM models?**

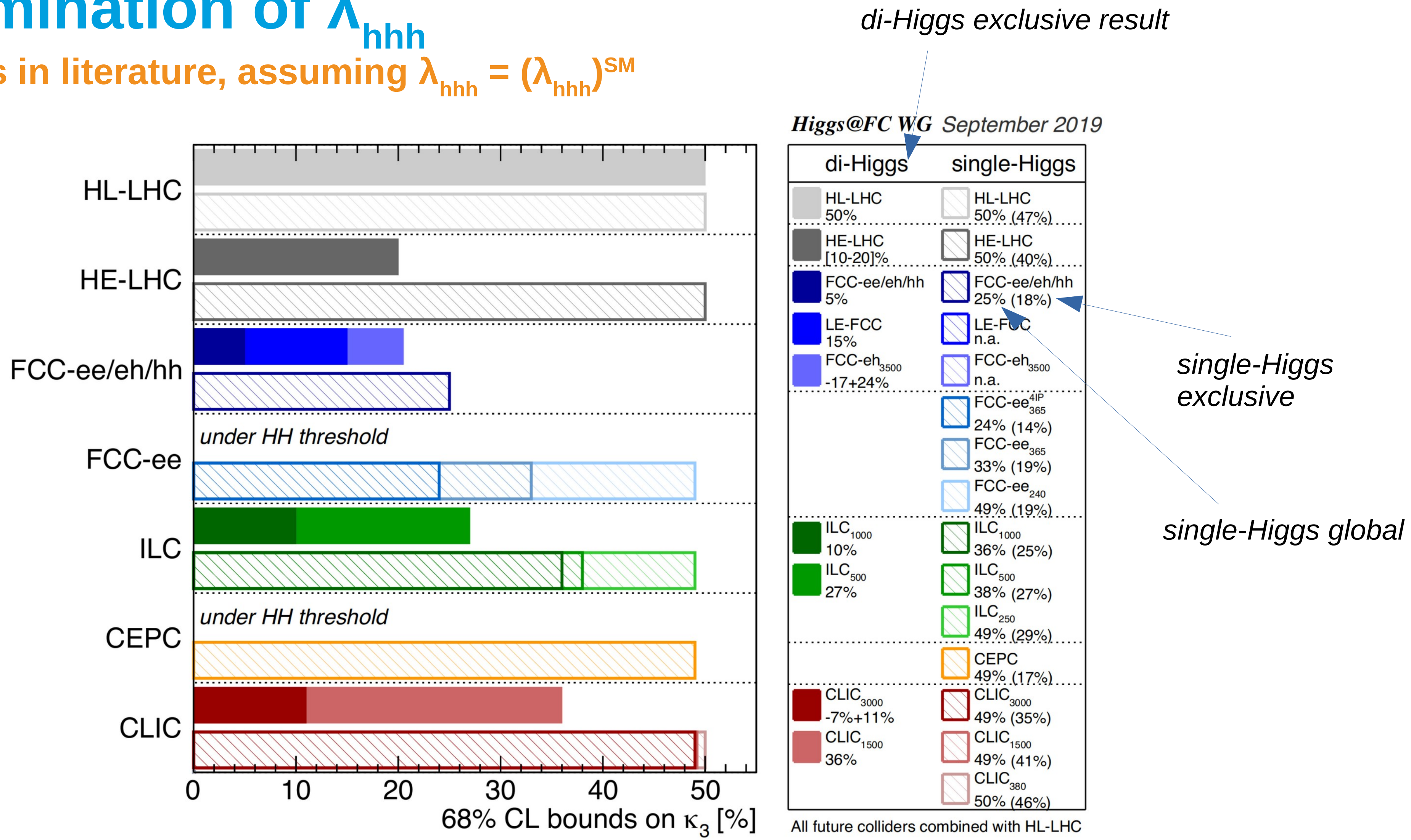


[Frederix et al., '14]

Future determination of λ_{hhh}

Expected sensitivities in literature, assuming $\lambda_{hhh} = (\lambda_{hhh})^{SM}$

Plot taken from
[de Blas et al., 1905.03764]



see also [Cepeda et al., 1902.00134], [Di Vita et al.1711.03978], [Fujii et al. 1506.05992, 1710.07621, 1908.11299], [Roloff et al., 1901.05897], [Chang et al. 1804.07130,1908.00753], *etc.*

Future determination of λ_{hhh}

Higgs production cross-sections (here double Higgs production) depend on λ_{hhh}

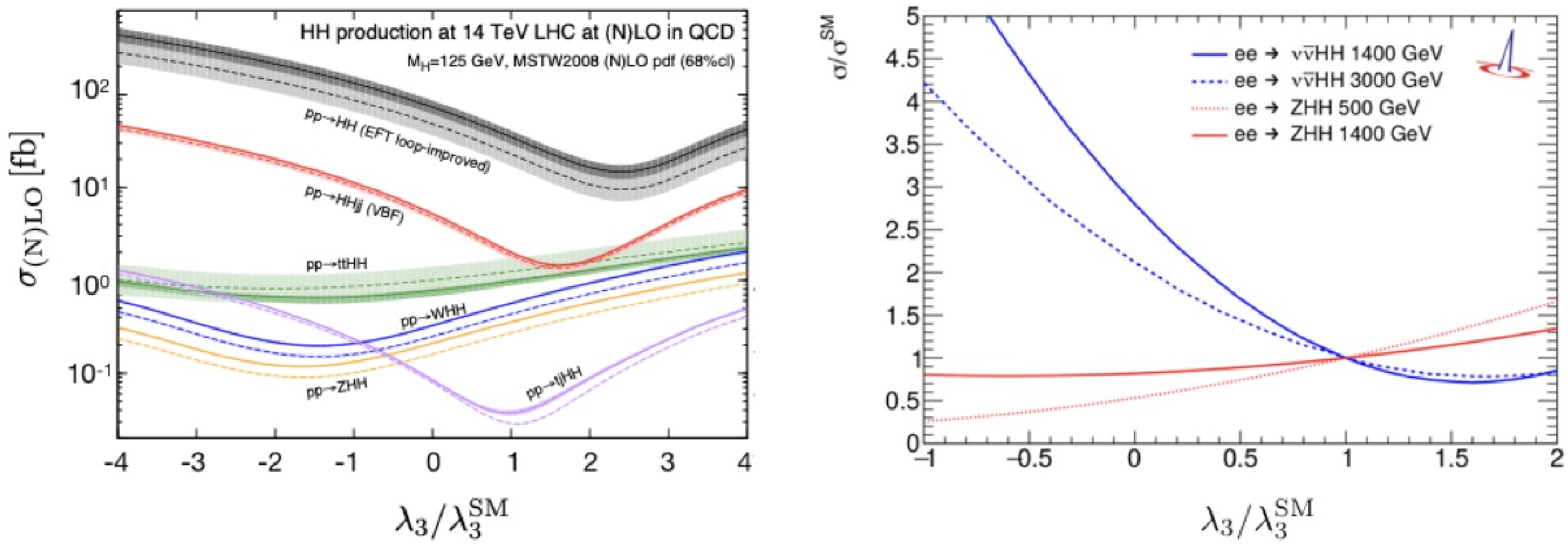


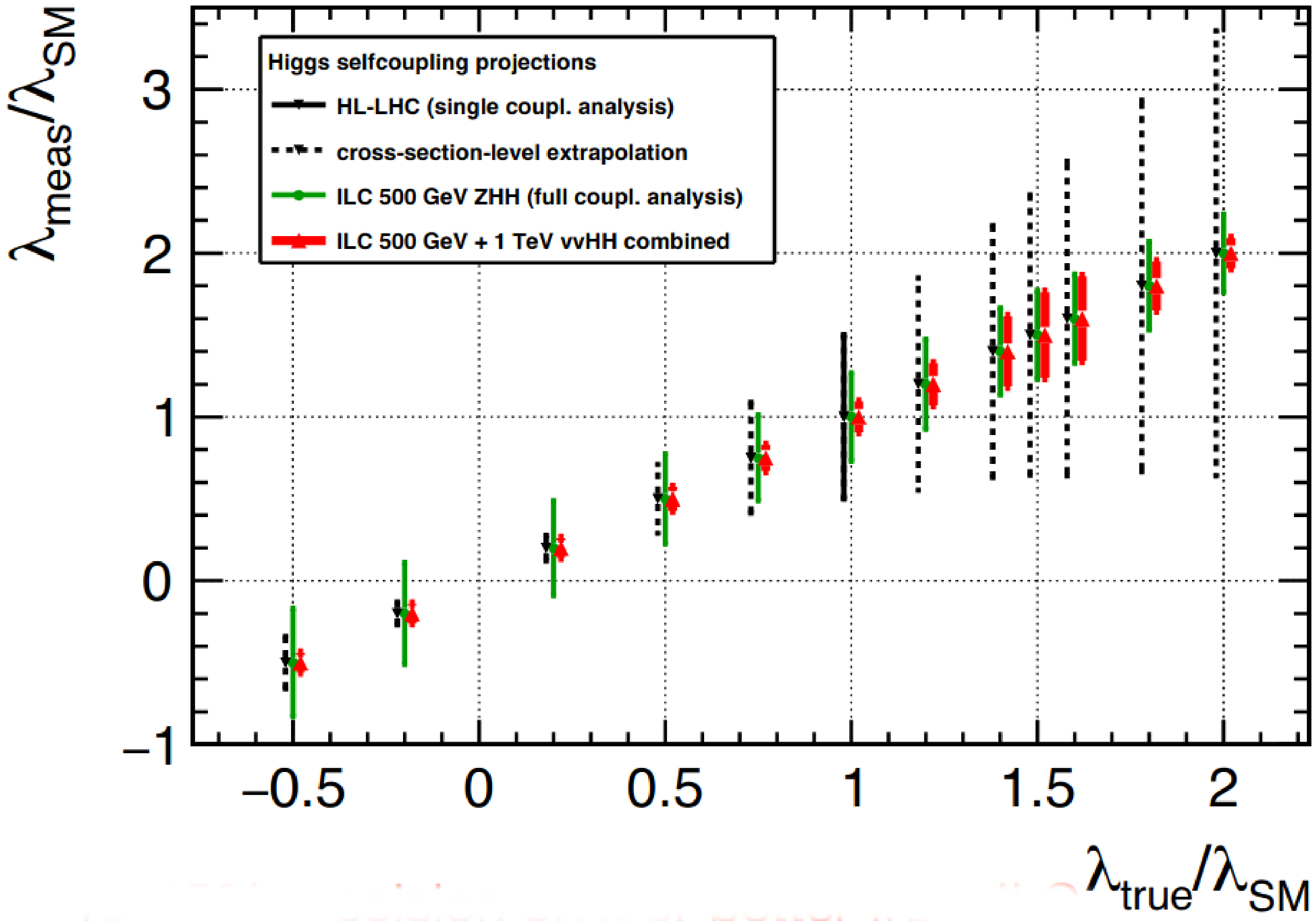
Figure 10. Double Higgs production at hadron (left) [65] and lepton (right) [66] colliders as a function of the modified Higgs cubic self-coupling. See Table 18 for the SM rates. At lepton colliders, the production cross sections do depend on the polarisation but this dependence drops out in the ratios to the SM rates (beam spectrum and QED ISR effects have been included).

[Frederix et al., 1401.7340]

Plots taken from
[de Blas et al., 1905.03764]

Future determination of λ_{hhh}

Achieved accuracy actually depends on the value of λ_{hhh}



[J. List et al. '21]

See also [Dürig, DESY-THESIS-2016-027]

The Two-Higgs-Doublet Model

- 2 $SU(2)_L$ doublets $\Phi_{1,2}$ of hypercharge ½
- CP-conserving 2HDM, with softly-broken Z_2 symmetry ($\Phi_1 \to \Phi_1$, $\Phi_2 \to -\Phi_2$) to avoid tree-level FCNCs

$$V^{(0)}_{\text{2HDM}} = m_1^2|\Phi_1|^2 + m_2^2|\Phi_2|^2 - m_3^2(\Phi_2^\dagger\Phi_1 + \Phi_1^\dagger\Phi_2) + \frac{\lambda_1}{2}|\Phi_1|^4 + \frac{\lambda_2}{2}|\Phi_2|^4 + \lambda_3|\Phi_1|^2|\Phi_2|^2 + \lambda_4|\Phi_2^\dagger\Phi_1|^2 + \frac{\lambda_5}{2}\Big((\Phi_2^\dagger\Phi_1)^2 + \text{h.c.}\Big)$$

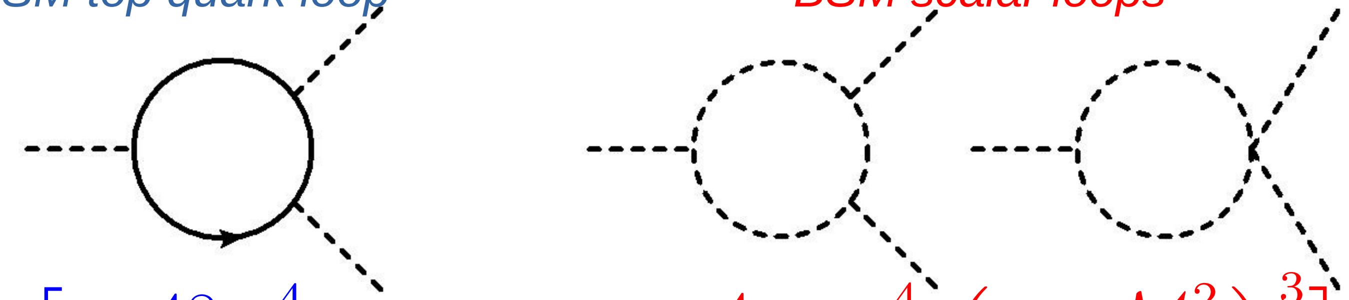
- m_1,m_2 eliminated with tadpole equations, and $v_1^2 + v_2^2 = v^2 = (246\ \text{GeV})^2$
- 7 free parameters in scalar sector: m_3, λ_i (i=1,..,5), $\tan\beta \equiv v_2/v_1$
- Mass eigenstates: h, H: CP-even Higgses, A: CP-odd Higgs, $H^\pm$: charged Higgs, α: CP-even Higgs mixing angle
- λ_i (i=1,..,5) traded for mass eigenvalues m_h, m_H, m_A, $m_{H\pm}$ and angle α
- m_3 replaced by a Z_2 soft-breaking mass scale

$$M^2 = \frac{2m_3^2}{s_{2\beta}}$$

One-loop non-decoupling effects

- **Leading one-loop** corrections to λ_{hhh} in models with extended sectors (like 2HDM):

SM top quark loop *BSM scalar loops*



$$\delta^{(1)}\lambda_{hhh} \supset \frac{1}{16\pi^2}\left[-\frac{48m_t^4}{v^3} + \sum_\Phi \frac{4n_\Phi m_\Phi^4}{v^3}\left(1-\frac{\mathcal{M}^2}{m_\Phi^2}\right)^3\right]$$

First found in 2HDM:
[Kanemura, Kiyoura, Okada, Senaha, Yuan ‘02]

$\mathcal{M}$: **BSM mass scale**, e.g. soft breaking scale M of Z_2 symmetry in 2HDM
n_Φ : # of d.o.f of field Φ

- Size of new effects depends on how the BSM scalars acquire their mass: $m_\Phi^2 \sim \mathcal{M}^2 + \tilde{\lambda}v^2$

$$\left(1-\frac{\mathcal{M}^2}{m_\Phi^2}\right)^3 \longrightarrow \begin{cases} 0, \text{ for } \mathcal{M}^2 \gg \tilde{\lambda}v^2 \\ 1, \text{ for } \mathcal{M}^2 \ll \tilde{\lambda}v^2 \end{cases} \longrightarrow \textbf{Huge BSM effects possible!}$$

One-loop non-decoupling effects

- Leading one-loop c

$\delta^{(1)}\lambda_{hhh} \supset$

$\mathcal{M}$: **BSM mass**
n_Φ : # of d.o.f of

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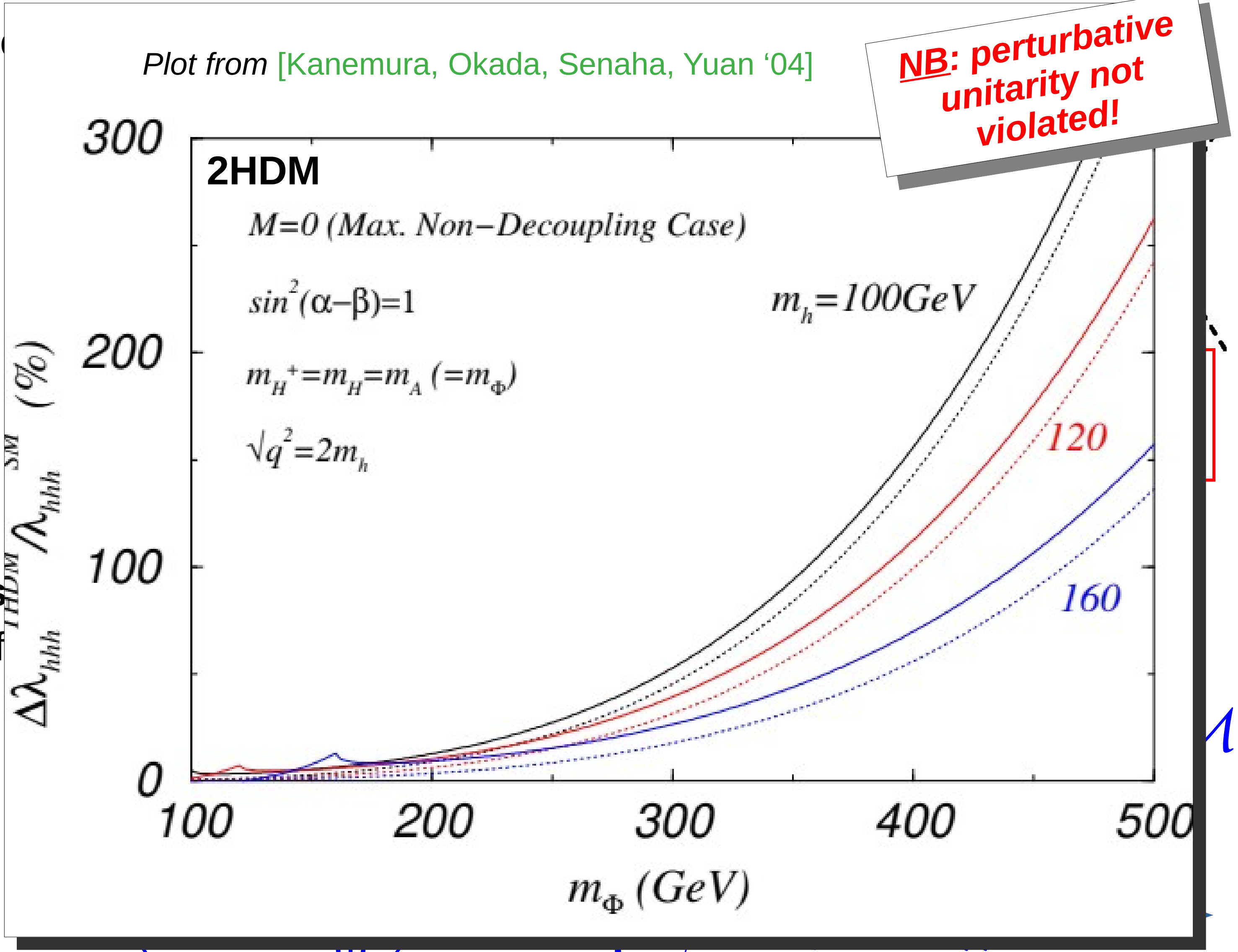


NB: perturbative unitarity not violated!

First found in 2HDM: [Kanemura, Kiyoura, Okada, Senaha, Yuan '02]

$\mathcal{M}^2 + \tilde{\lambda} v^2$

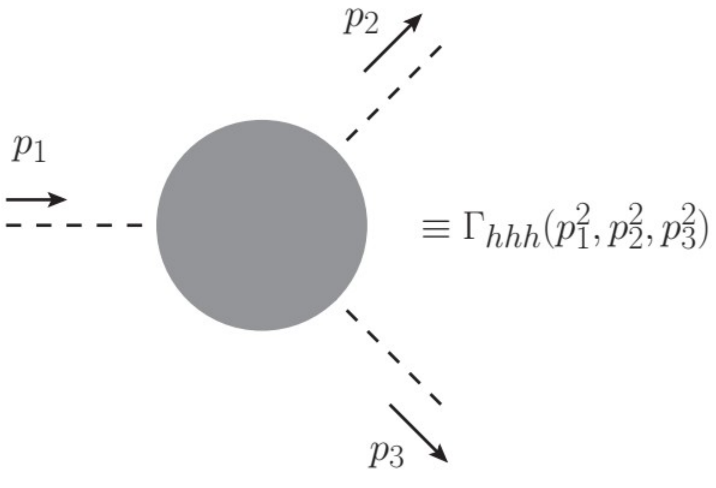
Huge BSM effects possible!

Our calculation

Goal: How large can the two-loop corrections to λ_{hhh} become?

An effective Higgs trilinear coupling

- In principle: consider 3-point function Γ_{hhh}
  but this is momentum dependent → very difficult beyond one loop



- Instead, consider an **effective trilinear coupling**

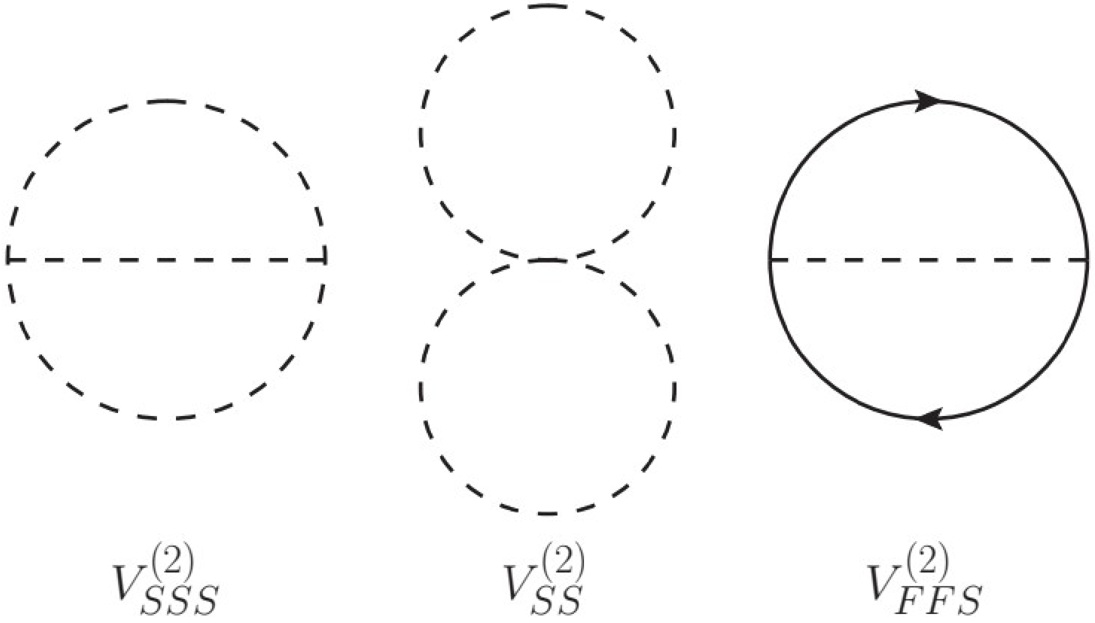
$$\lambda_{hhh} \equiv \frac{\partial^3 V_{\text{eff}}}{\partial h^3}\bigg|_{\text{min}}$$

- Momentum effects are neglected, but are expected to be *sub-leading* anyway
  - At one loop [Kanemura, Okada, Senaha, Yuan '04]: effects of a few % (away from thresholds)
  - At two loops, no study for 3-pt. functions but experience from Higgs mass calculations

Our effective-potential calculation

[JB, Kanemura '19]

- **Step 1**: compute $V_{\text{eff}} = V^{(0)} + \frac{1}{16\pi^2} V^{(1)} + \frac{1}{(16\pi^2)^2} V^{(2)}$ ($\overline{\text{MS}}$ result)
  - $V^{(2)}$: 1PI vacuum bubbles
  - *Dominant BSM contributions to $V^{(2)}$* = diagrams involving **heavy BSM scalars and top quark**



  - **Aligned scenarios** → no mixing + compatible with experimental results
  - **Neglect masses of light states** (SM-like Higgs, light fermions, ...)

Our effective-potential calculation

[JB, Kanemura '19]

- **Step 1**: compute $V_{\text{eff}} = V^{(0)} + \frac{1}{16\pi^2}V^{(1)} + \frac{1}{(16\pi^2)^2}V^{(2)}$ ($\overline{\text{MS}}$ result)
  - $V^{(2)}$: 1PI vacuum bubbles
  - *Dominant BSM contributions to* $V^{(2)}$ = diagrams involving **heavy BSM scalars and top quark**
  - *Aligned scenarios + neglect light masses*

- **Step 2**:

$$\lambda_{hhh} \equiv \frac{\partial^3 V_{\text{eff}}}{\partial h^3}\bigg|_{\text{min.}} = \frac{3[M_h^2]_{V_{\text{eff}}}}{v} + \left[\frac{\partial^3}{\partial h^3} - \frac{3}{v}\left(\frac{\partial^2}{\partial h^2} - \frac{1}{v}\frac{\partial}{\partial h}\right)\right]\Delta V\bigg|_{\text{min.}}$$

($\overline{\text{MS}}$ result too)

Express tree-level result in terms of effective-potential Higgs mass

Our effective-potential calculation

[JB, Kanemura '19]

- **Step 1**: compute $V_{\rm eff} = V^{(0)} + \frac{1}{16\pi^2}V^{(1)} + \frac{1}{(16\pi^2)^2}V^{(2)}$ ($\overline{\rm MS}$ result)
  - $V^{(2)}$: 1PI vacuum bubbles
  - *Dominant BSM contributions to* $V^{(2)}$ = diagrams involving **heavy BSM scalars and top quark**
  - *Aligned scenarios + neglect light masses*

- **Step 2**: ($\overline{\rm MS}$ result too)

$$\lambda_{hhh} \equiv \frac{\partial^3 V_{\rm eff}}{\partial h^3}\bigg|_{\rm min.} = \frac{3[M_h^2]_{V_{\rm eff}}}{v} + \left[\frac{\partial^3}{\partial h^3} - \frac{3}{v}\left(\frac{\partial^2}{\partial h^2} - \frac{1}{v}\frac{\partial}{\partial h}\right)\right]\Delta V\bigg|_{\rm min.}$$

- **Step 3**: conversion from $\overline{\rm MS}$ to OS scheme (*details in the following*)
  - Express result in terms of **pole masses**: M_t, M_h, M_Φ (Φ=H,A,H$^\pm$); OS Higgs VEV $v_{\rm phys} = \frac{1}{\sqrt{\sqrt{2}G_F}}$
  - Include **finite WFR**: $\hat{\lambda}_{hhh} = (Z_h^{\rm OS}/Z_h^{\overline{\rm MS}})^{3/2}\lambda_{hhh}$
  - Prescription for M to ensure **proper decoupling** with $M_\Phi^2 = \tilde{M}^2 + \tilde{\lambda}_\Phi v^2$ and $\tilde{M} \to \infty$

Our effective-potential calculation – scheme conversion

- OS result is obtained as

$$\hat{\lambda}_{hhh} = \underbrace{\left(\frac{Z_h^{\text{OS}}}{Z_h^{\overline{\text{MS}}}}\right)^{3/2}}_{\text{inclusion of WFR}} \times \underbrace{\lambda_{hhh}}_{\substack{\overline{\text{MS}}\text{ parameters} \\ \text{translated to OS ones}}}$$

- Let's suppose (for simplicity) that λ_{hhh} only depends on one parameter x, as

$$\lambda_{hhh} = f^{(0)}(x^{\overline{\text{MS}}}) + \kappa f^{(1)}(x^{\overline{\text{MS}}}) + \kappa^2 f^{(2)}(x^{\overline{\text{MS}}}) \qquad \left(\kappa = \frac{1}{16\pi^2}\right)$$

and

$$x^{\overline{\text{MS}}} = X^{\text{OS}} + \kappa\delta^{(1)}x + \kappa^2\delta^{(2)}x$$

then in terms of OS parameters

$$\begin{aligned}\lambda_{hhh} = f^{(0)}(X^{\text{OS}}) &+ \kappa\left[f^{(1)}(X^{\text{OS}}) + \frac{\partial f^{(0)}}{\partial x}(X^{\text{OS}})\delta^{(1)}x\right] \\ &+ \kappa^2\left[f^{(2)}(X^{\text{OS}}) + \frac{\partial f^{(1)}}{\partial x}(X^{\text{OS}})\delta^{(1)}x + \frac{\partial f^{(0)}}{\partial x}(X^{\text{OS}})\delta^{(2)}x + \frac{\partial^2 f^{(0)}}{\partial x^2}(X^{\text{OS}})(\delta^{(1)}x)^2\right]\end{aligned}$$

Our effective-potential calculation – scheme conversion

- OS result is obtained as

$$\hat{\lambda}_{hhh} = \underbrace{\left(\frac{Z_h^{\mathrm{OS}}}{Z_h^{\overline{\mathrm{MS}}}}\right)^{3/2}}_{\text{inclusion of WFR}} \times \underbrace{\lambda_{hhh}}_{\substack{\overline{\mathrm{MS}}\text{ parameters}\\ \text{translated to OS ones}}}$$

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and

$$x^{\overline{\mathrm{MS}}} = X^{\mathrm{OS}} + \kappa\delta^{(1)}x + \kappa^2\delta^{(2)}x$$

then in terms of OS parameters

$$\lambda_{hhh} = f^{(0)}(X^{\mathrm{OS}}) + \kappa\left[f^{(1)}(X^{\mathrm{OS}}) + \cancel{\frac{\partial f^{(0)}}{\partial x}(X^{\mathrm{OS}})\delta^{(1)}x}\right]$$

$$+\kappa^2\left[f^{(2)}(X^{\mathrm{OS}}) + \frac{\partial f^{(1)}}{\partial x}(X^{\mathrm{OS}})\delta^{(1)}x + \cancel{\frac{\partial f^{(0)}}{\partial x}(X^{\mathrm{OS}})\delta^{(2)}x} + \cancel{\frac{\partial^2 f^{(0)}}{\partial x^2}(X^{\mathrm{OS}})(\delta^{(1)}x)^2}\right]$$

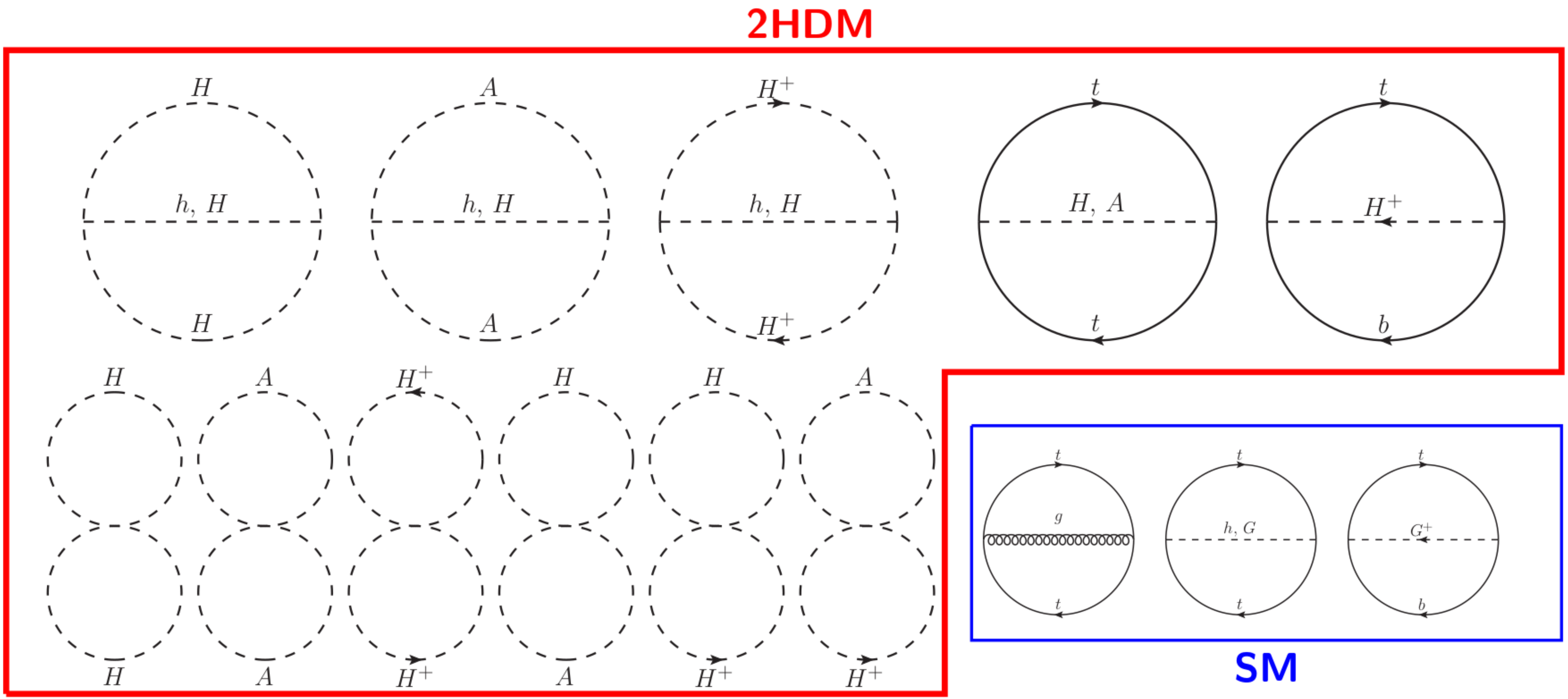
because we neglect m_h in the loop corrections and $\lambda_{hhh}^{(0)} = 3m_h^2/v$ (in absence of mixing)

Effective potential in the 2HDM

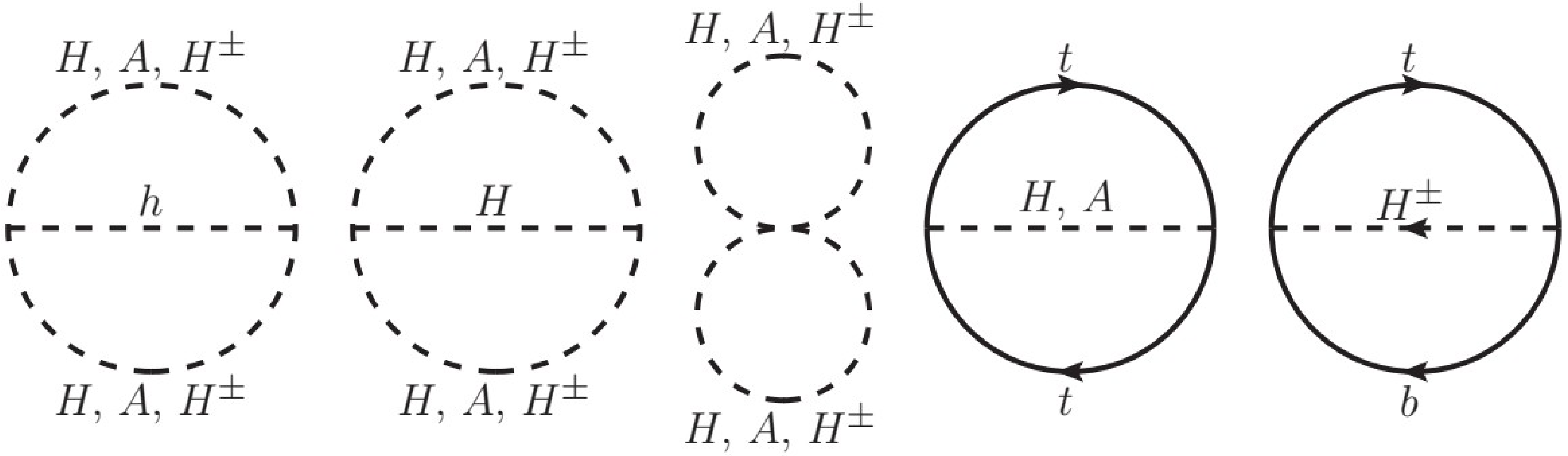
$$\lambda_{hhh} \equiv \left.\frac{\partial^3 V_{\text{eff}}}{\partial h^3}\right|_{\text{min.}}$$

[JB, Kanemura '19]

- 2HDM → **15 new BSM diagrams** appearing in $V^{(2)}$ w.r.t. the SM case



$\overline{\text{MS}}$ result



- Taking BSM scalars to be degenerate $M_\Phi = M_H = M_A = M_{H^\pm}$ we obtain in the $\overline{\text{MS}}$ scheme:
  (expressions for non-degenerate masses → see [JB, Kanemura 1911.11507])

$$\delta^{(2)}\lambda_{hhh} = \frac{16m_\Phi^4}{v^5}\left(4+9\cot^2 2\beta\right)\left(1-\frac{M^2}{m_\Phi^2}\right)^4\left[-2M^2-m_\Phi^2+(M^2+2m_\Phi^2)\,\overline{\log}\,m_\Phi^2\right]$$
$$+\frac{192m_\Phi^6\cot^2 2\beta}{v^5}\left(1-\frac{M^2}{m_\Phi^2}\right)^4\left[1+2\,\overline{\log}\,m_\Phi^2\right]$$
$$+\frac{96m_\Phi^4 m_t^2\cot^2\beta}{v^5}\left(1-\frac{M^2}{m_\Phi^2}\right)^3\left[-1+2\,\overline{\log}\,m_\Phi^2\right]+\mathcal{O}\left(\frac{m_\Phi^2 m_t^4}{v^5}\right)$$

Decoupling property in $\overline{\text{MS}}$ scheme

- Seeing whether corrections from additional BSM states decouple if said state is taken to be very massive is a good way to check the consistency of the calculation

$$\delta^{(2)}\lambda_{hhh} = \frac{16 m_\Phi^4}{v^5}\left(4 + 9\cot^2 2\beta\right)\left(1 - \frac{M^2}{m_\Phi^2}\right)^4\left[-2M^2 - m_\Phi^2 + (M^2 + 2m_\Phi^2)\,\overline{\log}\, m_\Phi^2\right]$$

$$\delta^{(1)}\lambda_{hhh} = \frac{16 m_\Phi^4}{v^3}\left(1 - \frac{M^2}{m_\Phi^2}\right)^3 \qquad + \frac{192 m_\Phi^6 \cot^2 2\beta}{v^5}\left(1 - \frac{M^2}{m_\Phi^2}\right)^4\left[1 + 2\,\overline{\log}\, m_\Phi^2\right]$$
$$+ \frac{96 m_\Phi^4 m_t^2 \cot^2\beta}{v^5}\left(1 - \frac{M^2}{m_\Phi^2}\right)^3\left[-1 + 2\,\overline{\log}\, m_\Phi^2\right] + \mathcal{O}\left(\frac{m_\Phi^2 m_t^4}{v^5}\right)$$

where $m_\Phi^2 = M^2 + \tilde{\lambda}_\Phi v^2$

- To have $m_\Phi \to \infty$, then we must take $M \to \infty$, otherwise the quartic couplings grow out of control
- Fortunately all of these terms go like

$$(m_\Phi^2)^{n-1}\left(1 - \frac{M^2}{m_\Phi^2}\right)^n \underset{m_\Phi^2 = M^2 + \tilde{\lambda}_\Phi v^2}{=} \frac{(\tilde{\lambda}_\Phi v^2)^n}{M^2 + \tilde{\lambda}_\Phi v^2} \xrightarrow[\tilde{\lambda}_\Phi v^2 \text{ fixed}]{M\to\infty} 0$$

$\overline{\text{MS}} \to$ OS scheme conversion

- To express $\delta^{(2)}\lambda_{hhh}$ in terms of physical parameters (v_phys, M_t, $M_A = M_H = M_{H^\pm} = M_\Phi$), we replace

$$m_A^2 \to M_A^2 - \Pi_{AA}(M_A^2), \quad m_H^2 \to M_H^2 - \Pi_{HH}(M_H^2), \quad m_{H^\pm}^2 \to M_{H^\pm}^2 - \Pi_{H^+H^-}(M_{H^\pm}^2),$$

$$v \to v_\text{phys} - \delta v, \quad m_t^2 \to M_t^2 - \Pi_{tt}(M_t^2)$$

- A priori, M is still renormalised in $\overline{\text{MS}}$ scheme, because it is difficult to relate to physical observable

  ... but then, expressions do not decouple for $M_\Phi^2 = M^2 + \tilde{\lambda}_\Phi v^2$ and $M \to \infty$!
- This is because we should relate M_Φ, renormalised in OS scheme, and M, renormalised in $\overline{\text{MS}}$ scheme, with a **one-loop relation** $\to$ then the two-loop corrections decouple properly
- We give a new "OS" prescription for the finite part of the counterterm for M be requiring that
  1. the decoupling of $\delta^{(2)}\hat{\lambda}_{hhh}$ (in OS scheme) is apparent using a relation $M_\Phi^2 = \tilde{M}^2 + \tilde{\lambda}_\Phi v^2$
  2. all the log terms in $\delta^{(2)}\hat{\lambda}_{hhh}$ are absorbed in δM^2

$$\begin{aligned}\delta^{(2)}\hat{\lambda}_{hhh} = &\frac{48M_\Phi^6}{v_\text{phys}^5}\left(1-\frac{\tilde{M}^2}{M_\Phi^2}\right)^4\left\{4+3\cot^2 2\beta\left[3-\frac{\pi}{\sqrt{3}}\left(\frac{\tilde{M}^2}{M_\Phi^2}+2\right)\right]\right\}+\frac{576M_\Phi^6\cot^2 2\beta}{v_\text{phys}^5}\left(1-\frac{\tilde{M}^2}{M_\Phi^2}\right)^4\\ &+\frac{288M_\Phi^4M_t^2\cot^2\beta}{v_\text{phys}^5}\left(1-\frac{\tilde{M}^2}{M_\Phi^2}\right)^3+\frac{168M_\Phi^4M_t^2}{v_\text{phys}^5}\left(1-\frac{\tilde{M}^2}{M_\Phi^2}\right)^3-\frac{48M_\Phi^6}{v_\text{phys}^5}\left(1-\frac{\tilde{M}^2}{M_\Phi^2}\right)^5+\mathcal{O}\left(\frac{M_\Phi^2M_t^4}{v_\text{phys}^5}\right)\end{aligned}$$

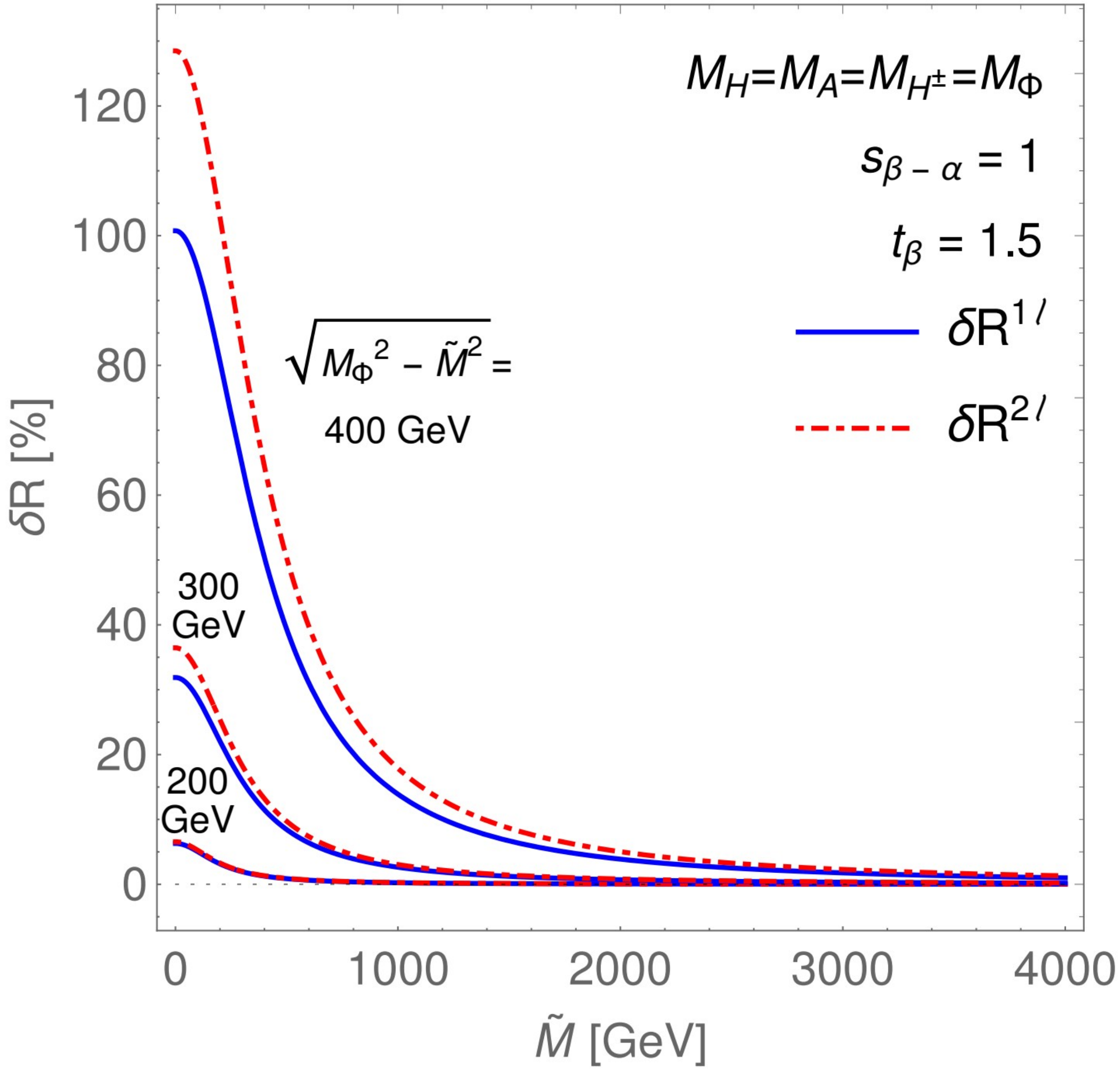
Numerical results in an aligned 2HDM

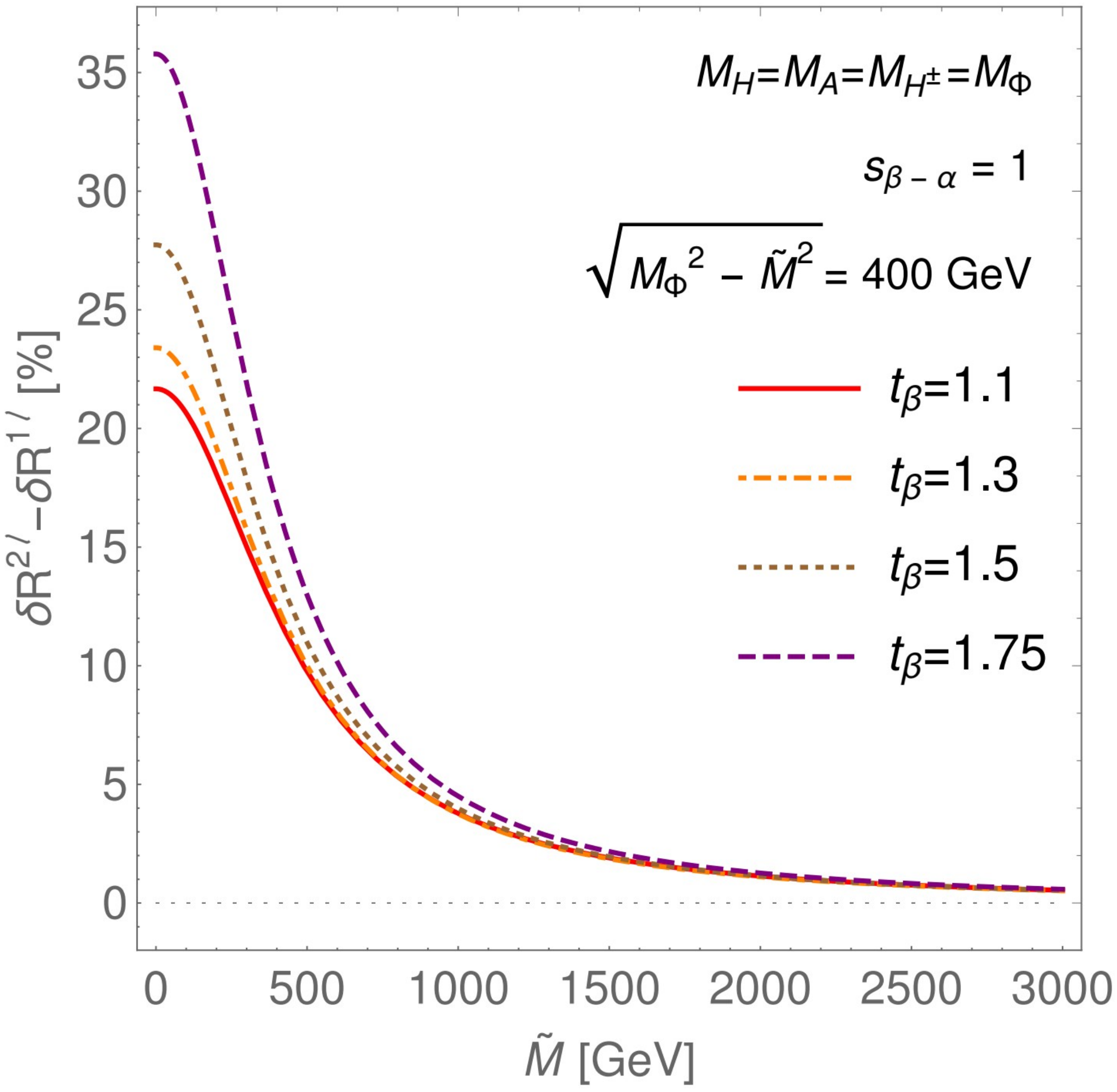
$$\delta R \equiv \frac{\hat{\lambda}_{hhh}^{\text{2HDM}} - \hat{\lambda}_{hhh}^{\text{SM}}}{\hat{\lambda}_{hhh}^{\text{SM}}}$$

Decoupling of BSM effects

$\tilde{M}$: modified "OS" version of Z_2 breaking scale

[JB, Kanemura '19]

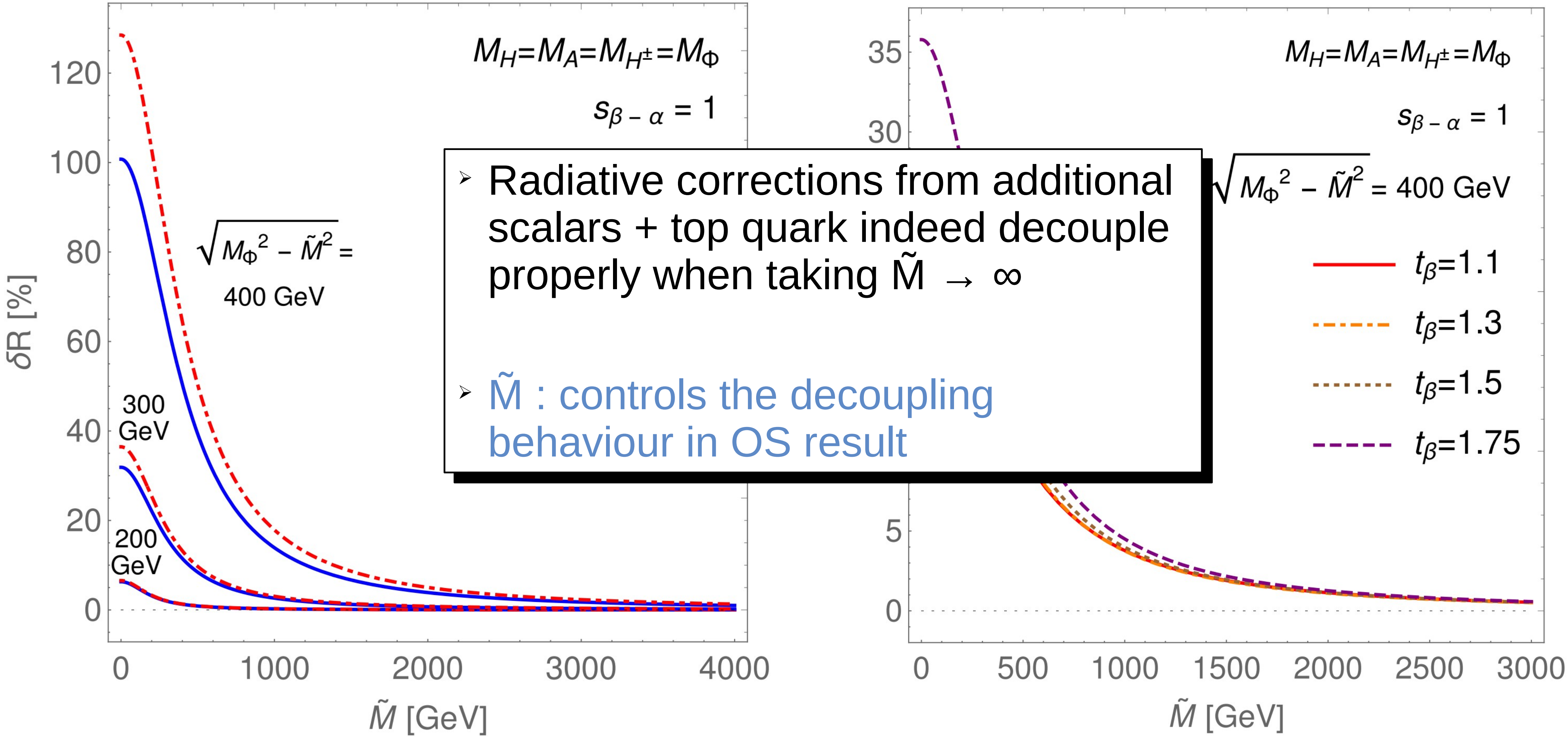




Decoupling of BSM effects

$\tilde{M}$: modified "OS" version of Z_2 breaking scale

[JB, Kanemura '19]

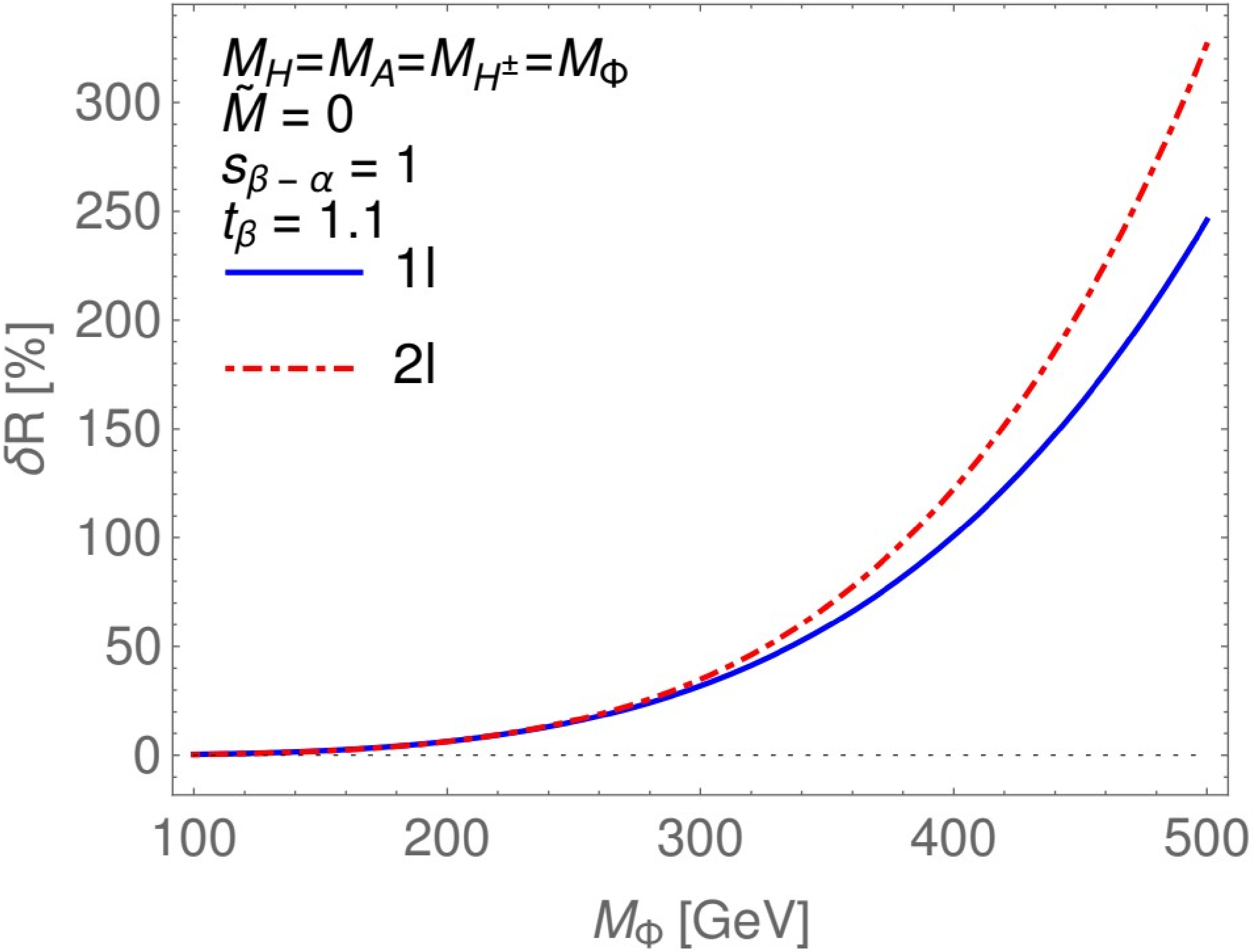


- Radiative corrections from additional scalars + top quark indeed decouple properly when taking $\tilde{M} \to \infty$
- $\tilde{M}$: controls the decoupling behaviour in OS result

BSM deviation of λ_{hhh} in non-decoupling limit

Taking degenerate BSM scalar masses: $M_\Phi = M_H = M_A = M_{H^\pm}$

[JB, Kanemura 1903.05417]



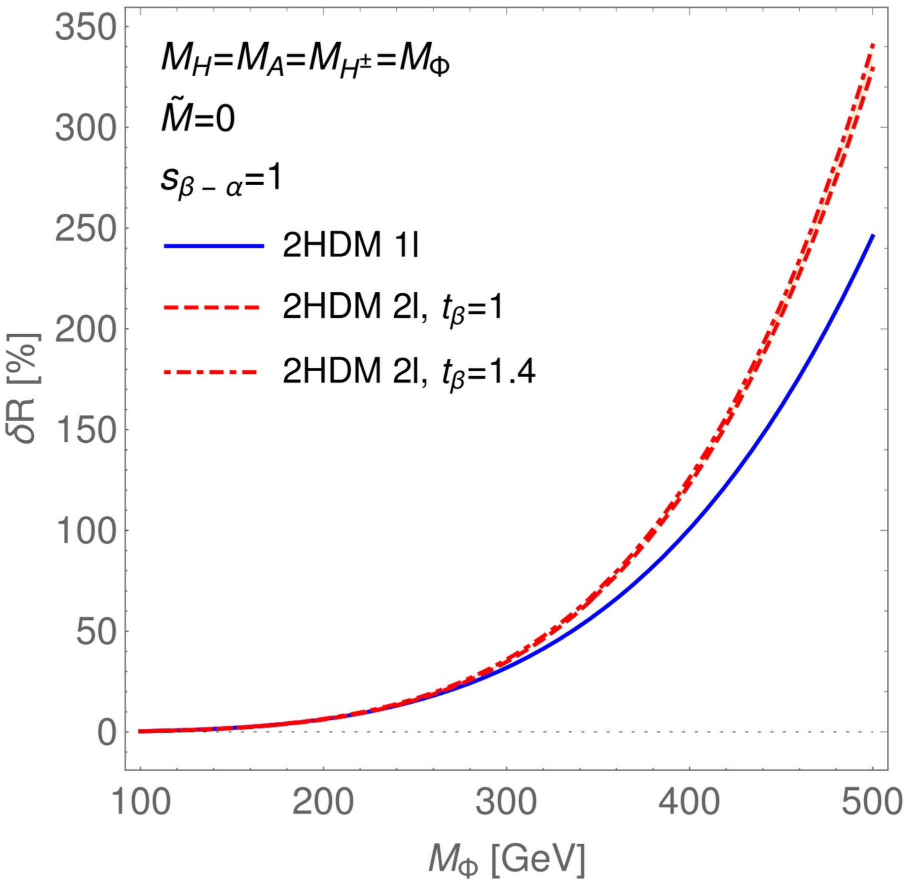
- $\tilde{M} = 0$ → maximal non-decoupling effects
- 1 loop: $\propto M_\Phi^4$
- 2 loops: $\propto M_\Phi^6$
- $\delta^{(2)}\lambda_{hhh}$ typically 10-20% of $\delta^{(1)}\lambda_{hhh}$ for most of mass range, at most 30%

λ_{hhh} at two loops in more models

[JB, Kanemura 1911.11507]

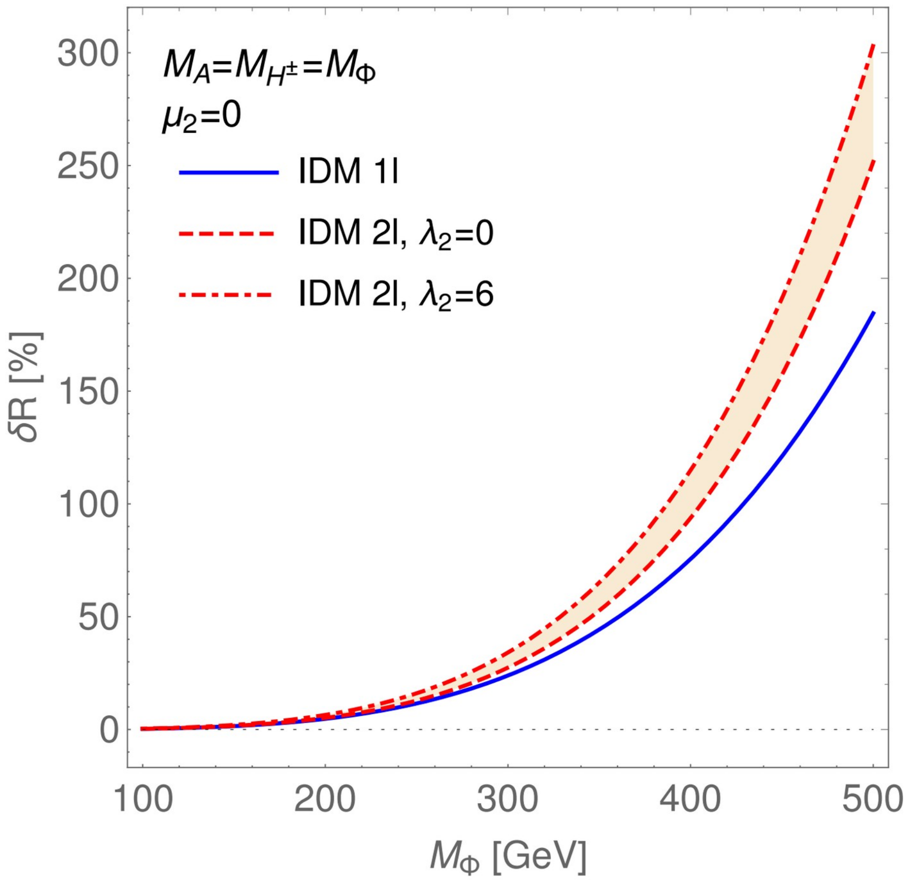
- Calculations in several other models: *IDM*, *singlet extension of SM*
- Each model contains a **new parameter appearing from two loops**:

Aligned 2HDM → tanβ



tanβ constrained by perturbative unitarity
→ only small effects

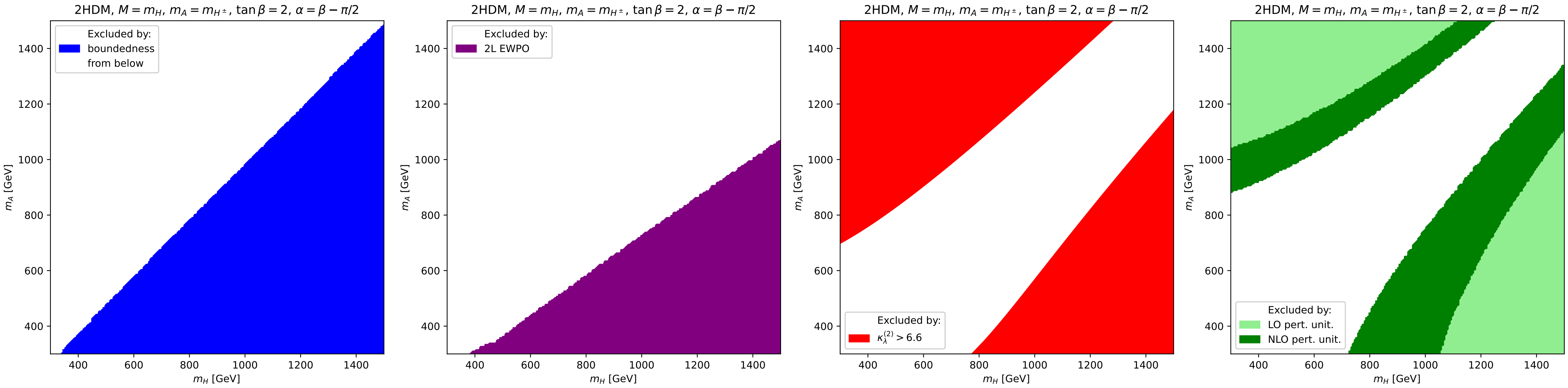
IDM → λ_2 (quartic coupling of inert doublet)



λ_2 is less contrained → **enhancement is possible**
(but 2l effects remains <u>well smaller</u> than 1l ones)

2HDM benchmark plane – individual theoretical constraints

Constraints shown below are independent of 2HDM type



Boundedness from below

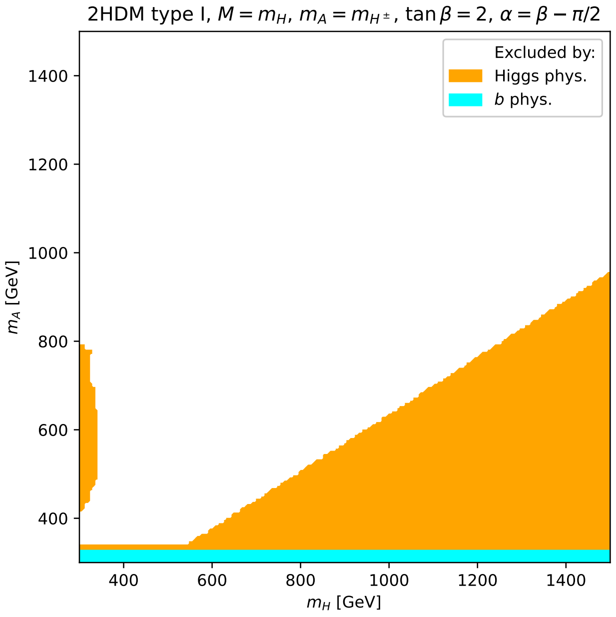
EW precision observables computed at 2L

$\kappa_\lambda^{(2)} > 6.6$

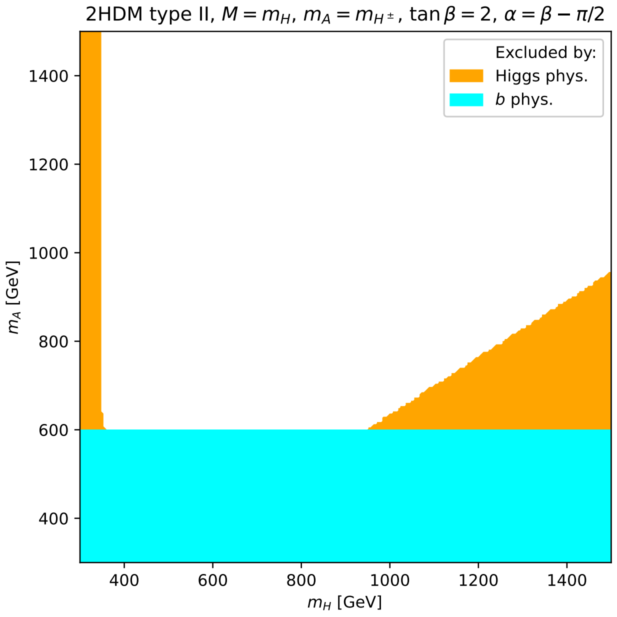
Perturbative unitarity at (N)LO

2HDM benchmark plane – experimental constraints

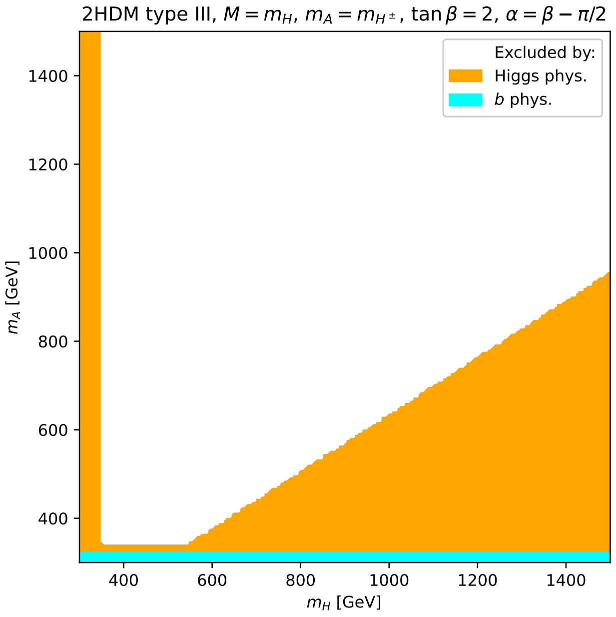
i.e. Higgs physics (via HiggsBounds and HiggsSignals) and *b* physics (from [Gfitter group 1803.01853])



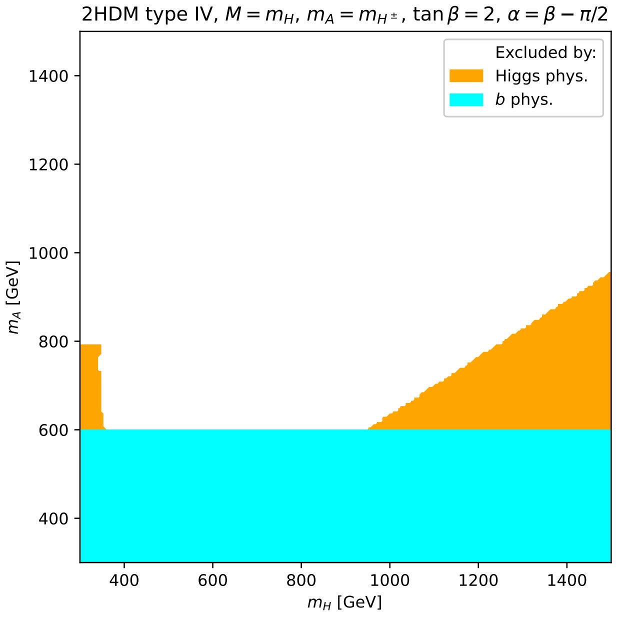
Type-I



Type-II



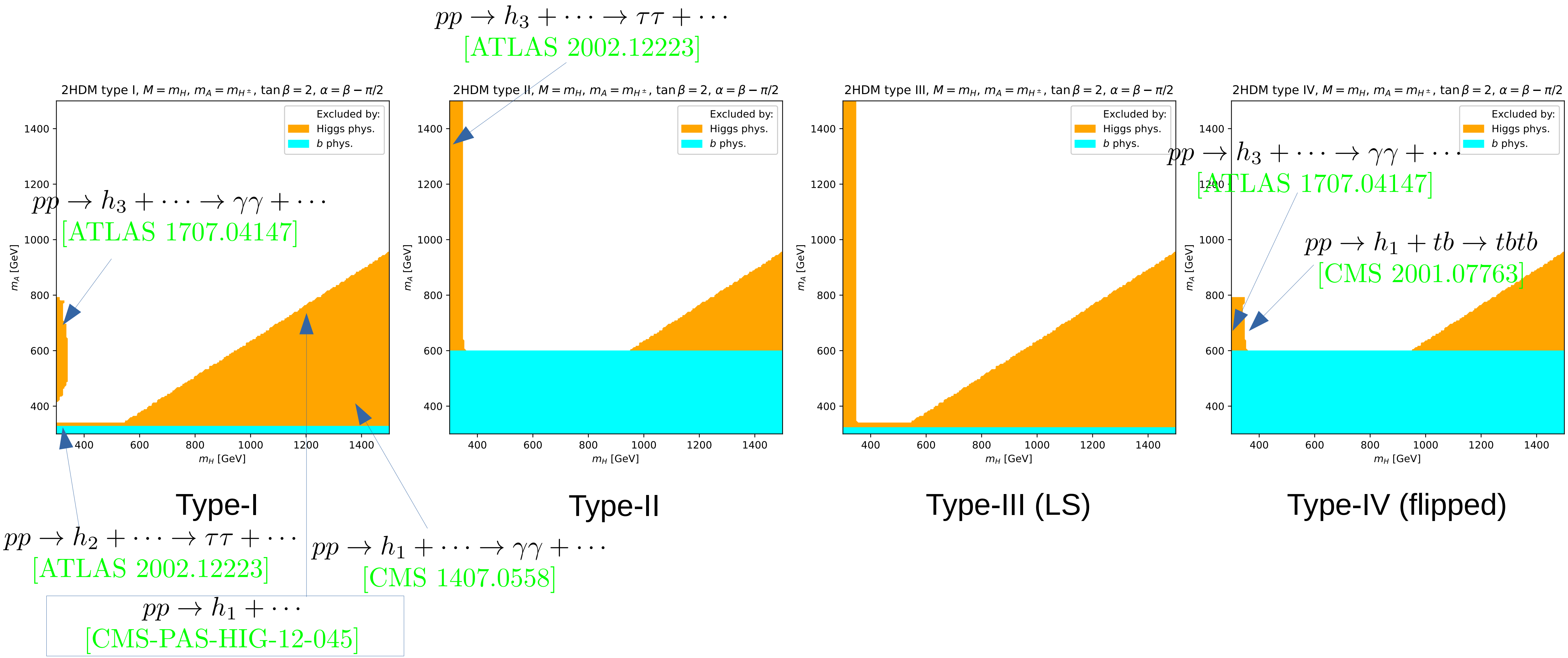
Type-III (LS)



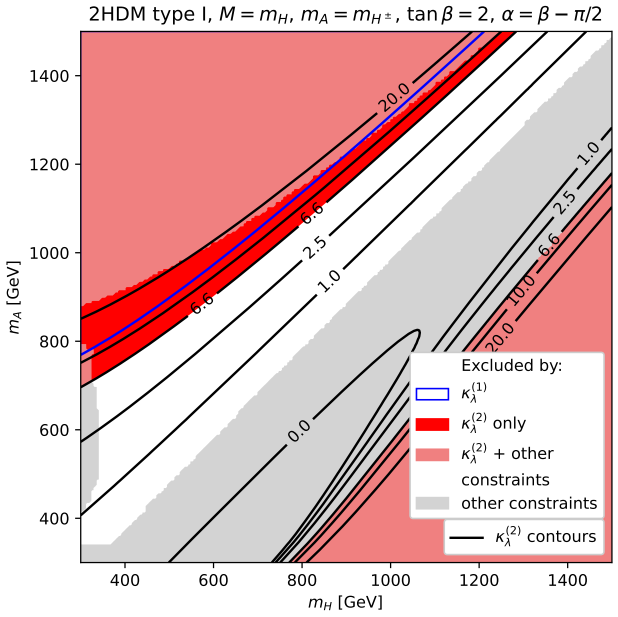
Type-IV (flipped)

2HDM benchmark plane – experimental constraints

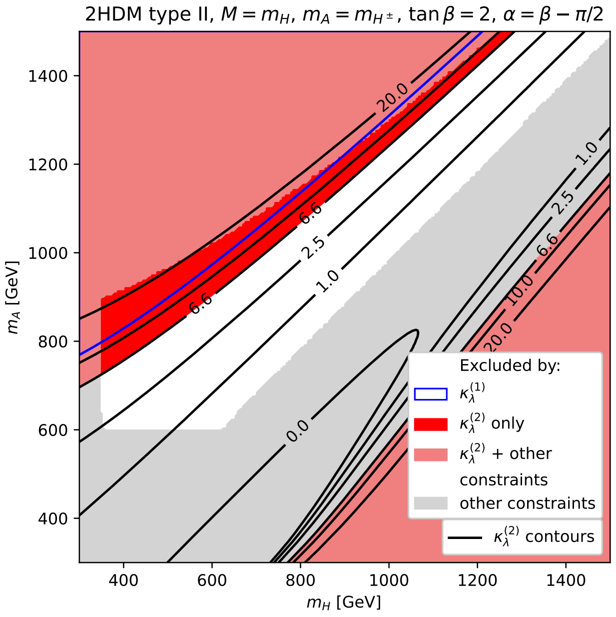
i.e. Higgs physics (via HiggsBounds and HiggsSignals) and *b* physics (from [Gfitter group 1803.01853])



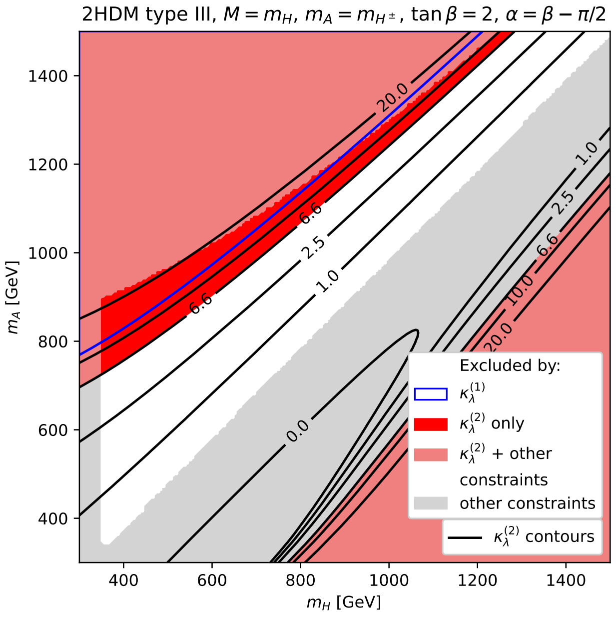
2HDM benchmark plane – results for all types



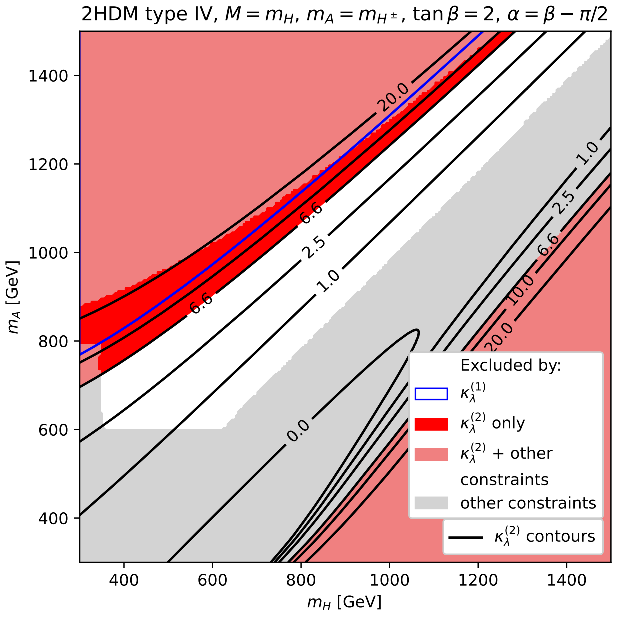
Type-I



Type-II



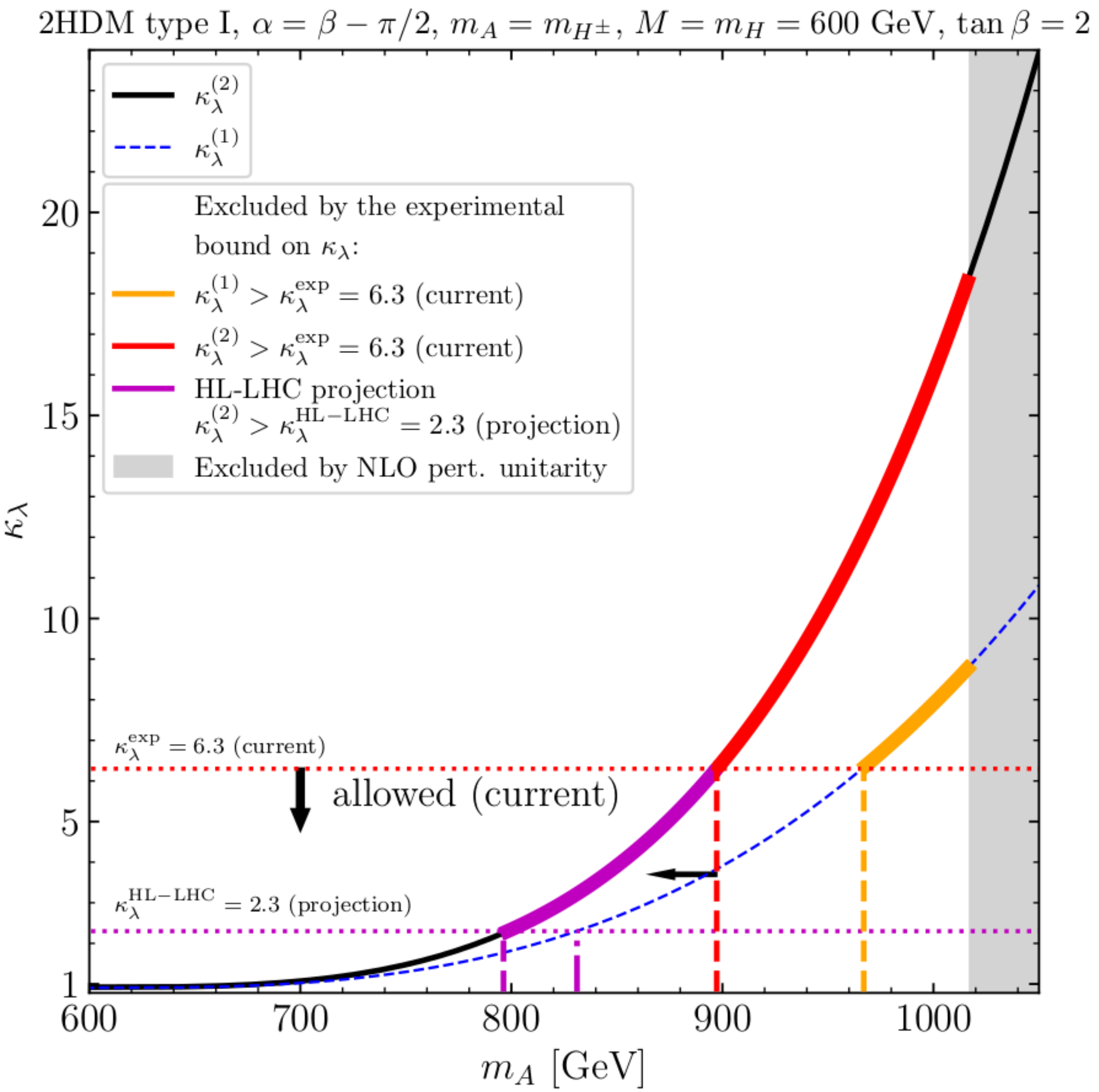
Type-III (LS)



Type-IV (flipped)

A benchmark scenario in the aligned 2HDM – 1D scan

Within the previously shown plane, we fix $M=m_H=600$ GeV, and vary $m_A=m_{H\pm}$



- Illustrates the significantly improved reach of the experimental limit when including **2L corrections** in calculation of κ_λ

Similar results for other BSM models

Using here predictions for κ_λ computed to full 1L with Python package `anyH3`

[Bahl, JB, Gabelmann, Weiglein *to appear*]

- Example results shown for benchmark scenarios in
  - **2HDM type-II**
  - **Inert Doublet Model**
  - **Complex triplet extension of SM**
  - **Georgi-Machacek model**

- Perturbative unitarity checks:
  - Solid: $|a_0| < 0.5$
  - Dashed: $0.5 < |a_0| < 1$
  - Dashed transparent $|a_0| > 1$