



# HIGHER-ORDER QCD CALCULATIONS FOR THE LHC

---

Alexander Huss



# THE PLAN.

THE PLAN.

## 1. NNLO predictions for the LHC

- jets & interpolations grids
- identified photons & fragmentation

## 2. Differential N<sup>3</sup>LO

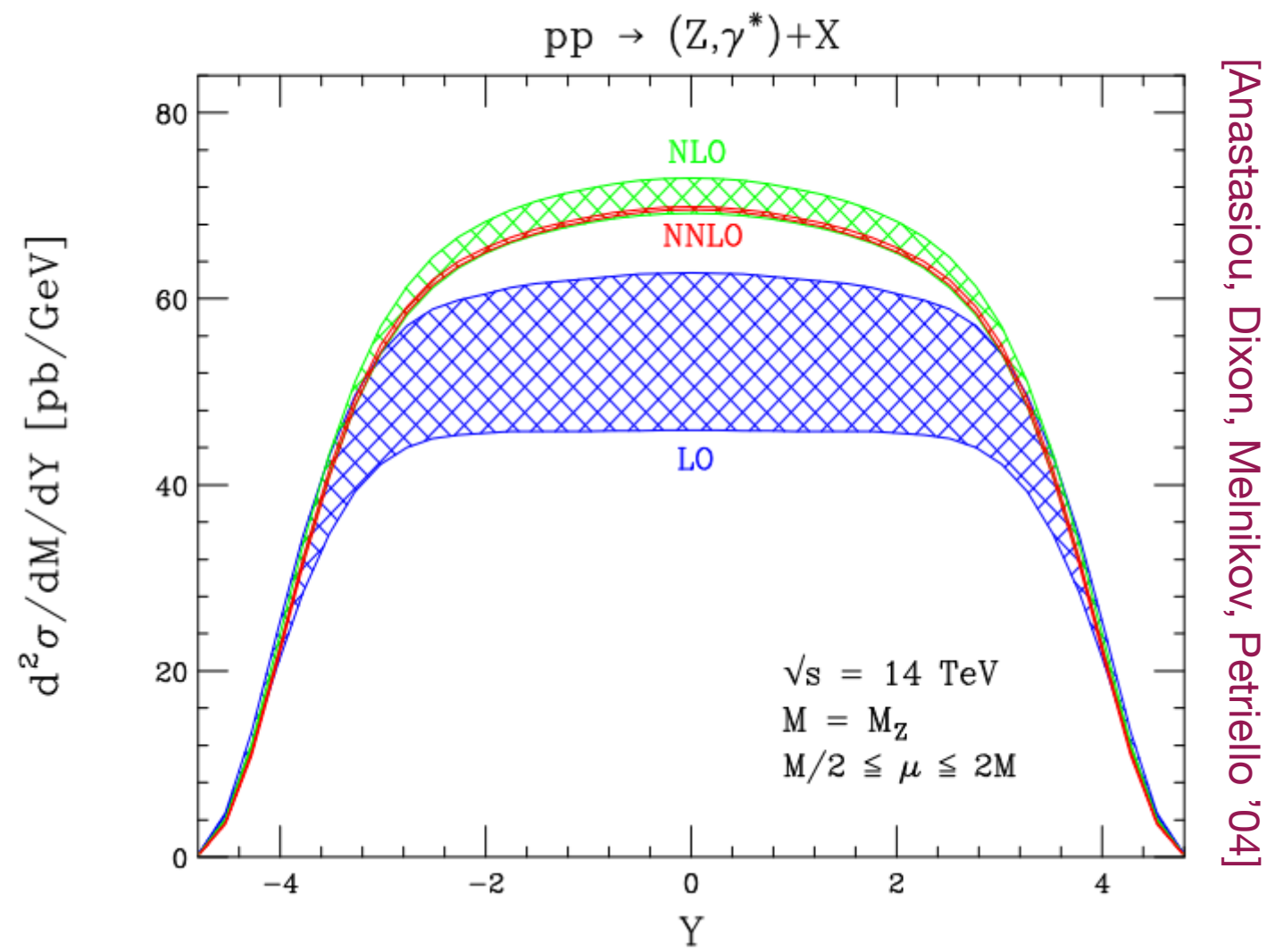
- Higgs & fiducial power corrections
- Drell-Yan & PDFs

## 3. Bayesian approach to MHO

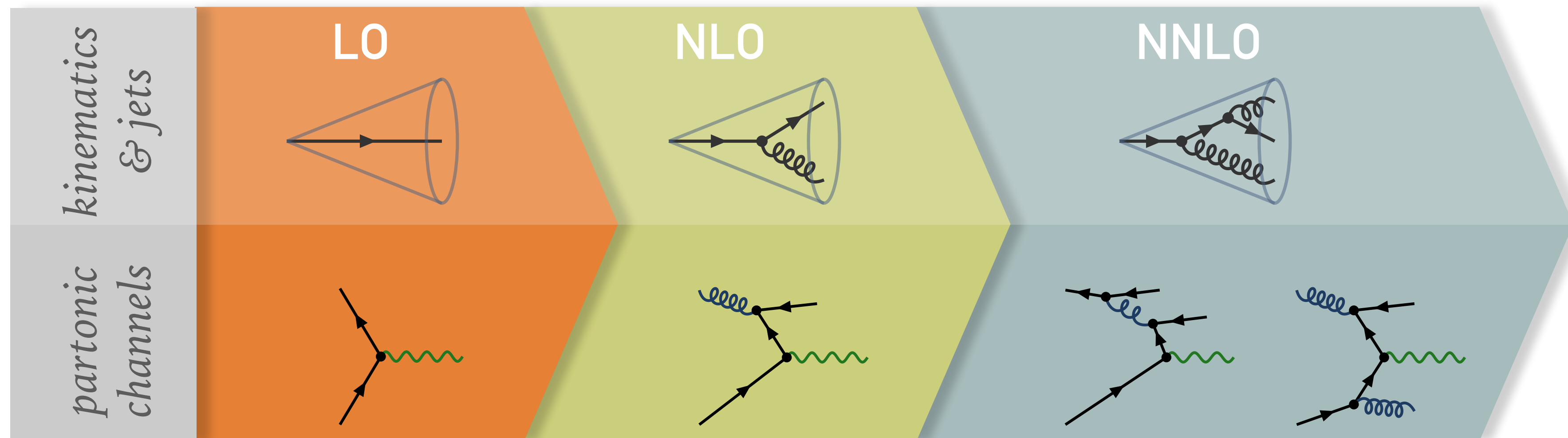
- the abc model & correlations

## 4. Summary & Outlook

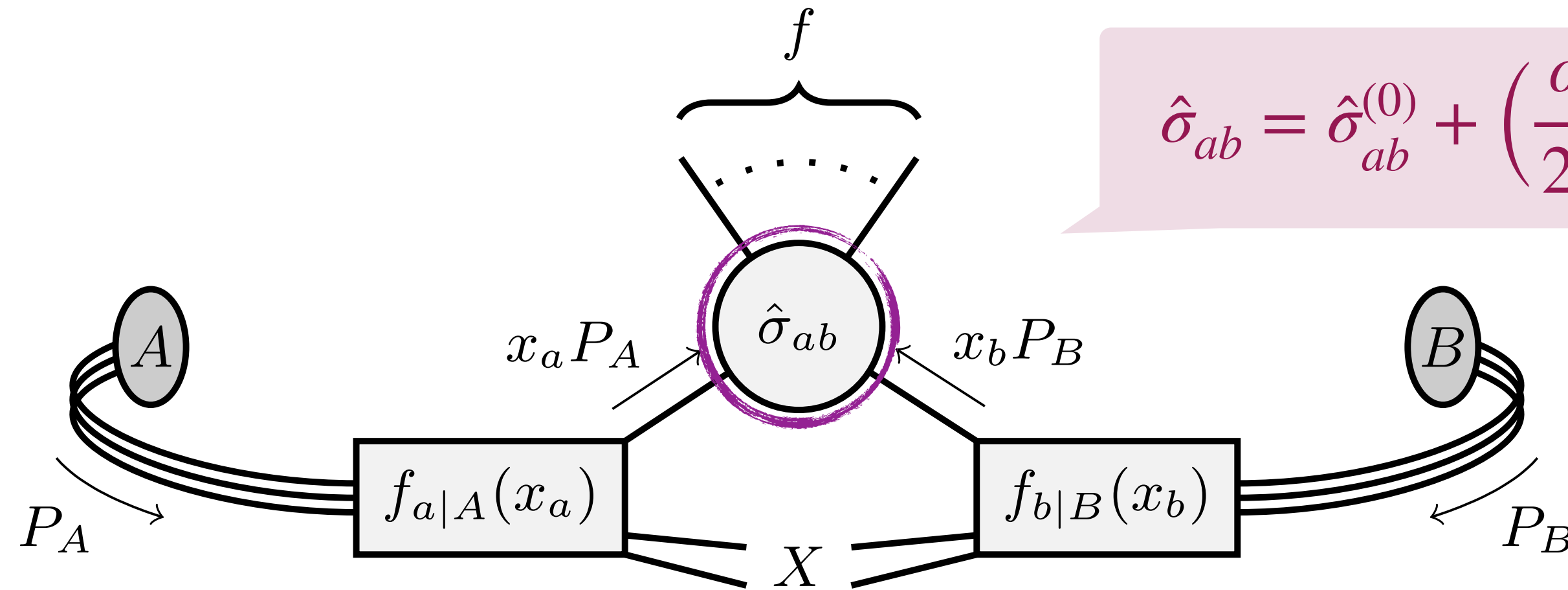
# WHAT WE HOPE NNLO WILL GIVE US



- reduced uncertainties ( $\leftrightarrow$  missing higher orders)
- guaranteed that all partonic channels open up at NNLO
- better modelling of final-state kinematics & jets



# THE MASTER FORMULA



$$\hat{\sigma}_{ab} = \hat{\sigma}_{ab}^{(0)} + \left(\frac{\alpha_s}{2\pi}\right) \hat{\sigma}_{ab}^{(1)} + \left(\frac{\alpha_s}{2\pi}\right)^2 \hat{\sigma}_{ab}^{(2)} + \dots$$

$$\sigma_{AB} = \sum_{ab} \int_0^1 dx_a \int_0^1 dx_b f_{a|A}(x_a) f_{b|B}(x_b) \hat{\sigma}_{ab}(x_a, x_b) (1 + \mathcal{O}(\Lambda_{\text{QCD}}/Q))$$

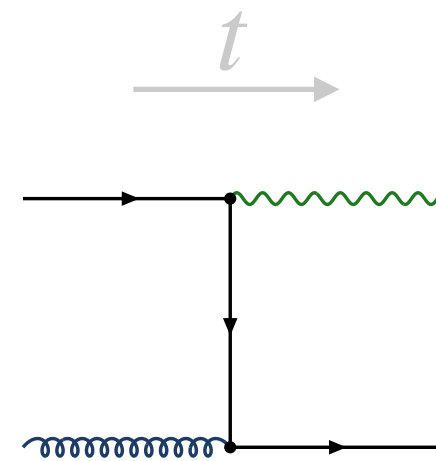
parton distribution functions  
(non-perturbative, universal)  
in principle, improvable

hard scattering  
(perturbation theory)  
systematically improvable

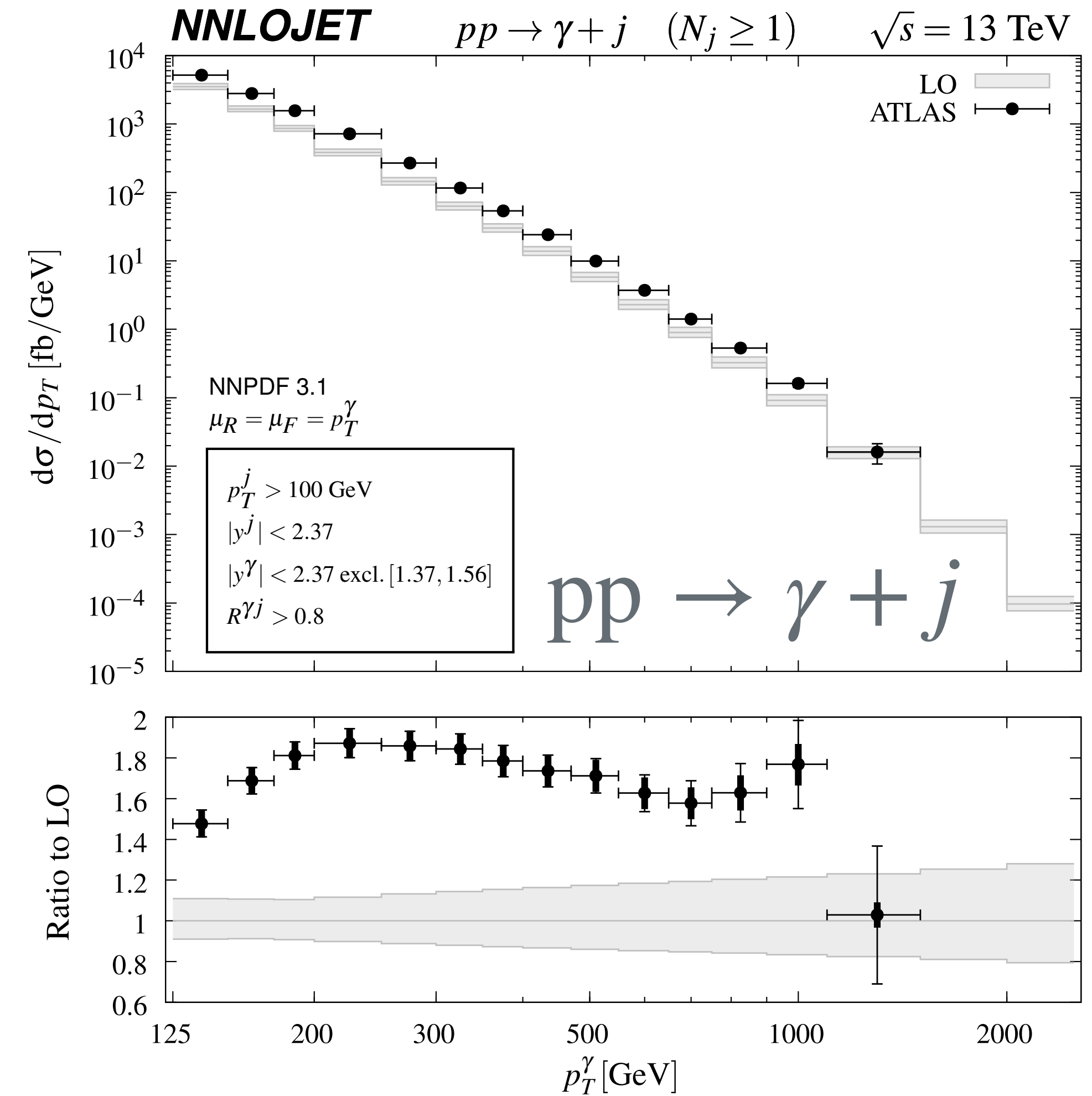
non-perturbative effects  
(power suppressed)  
ultimately, limiting factor?

# PERTURBATION THEORY @ LEADING ORDER

01



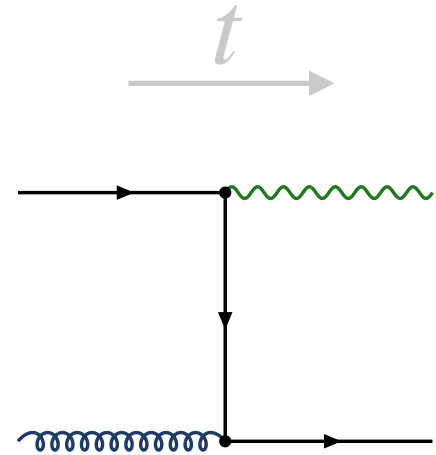
Only captures gross features & unreliable uncertainty estimates



[Chen, Gehrmann, Glover, Höfer, A.Huss '19]

# PERTURBATION THEORY @ NEXT-TO-LEADING ORDER

LO



NLO

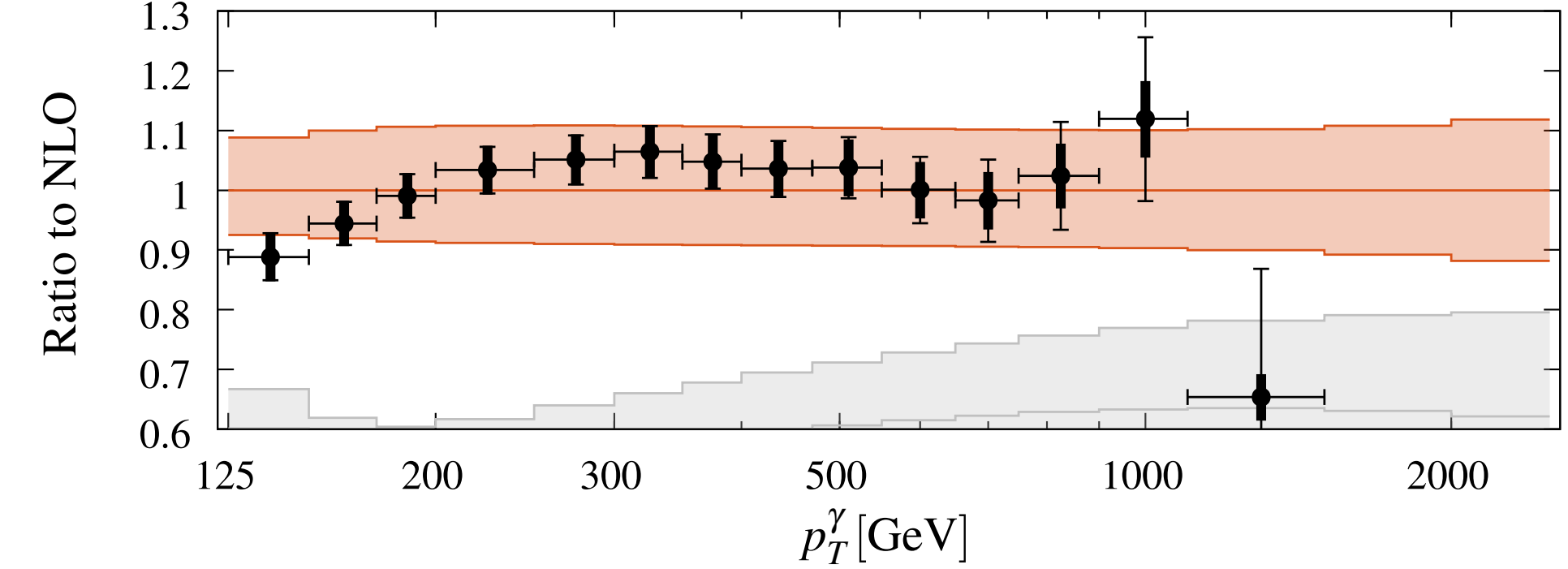
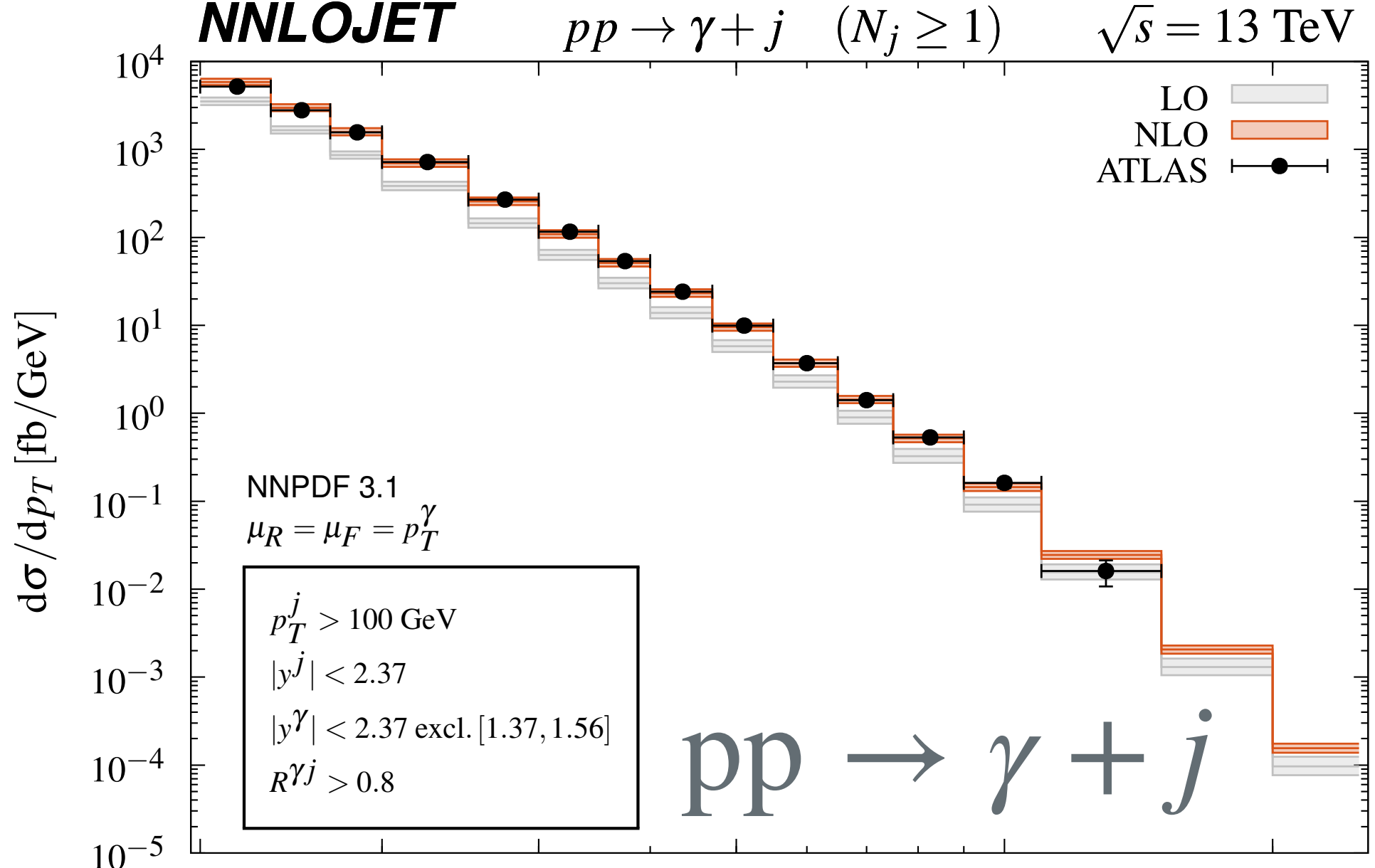


can become unresolved

soft:  $E_g \rightarrow 0$       collinear:  $\theta_{qg} \rightarrow 0$

higher order: more loops & legs

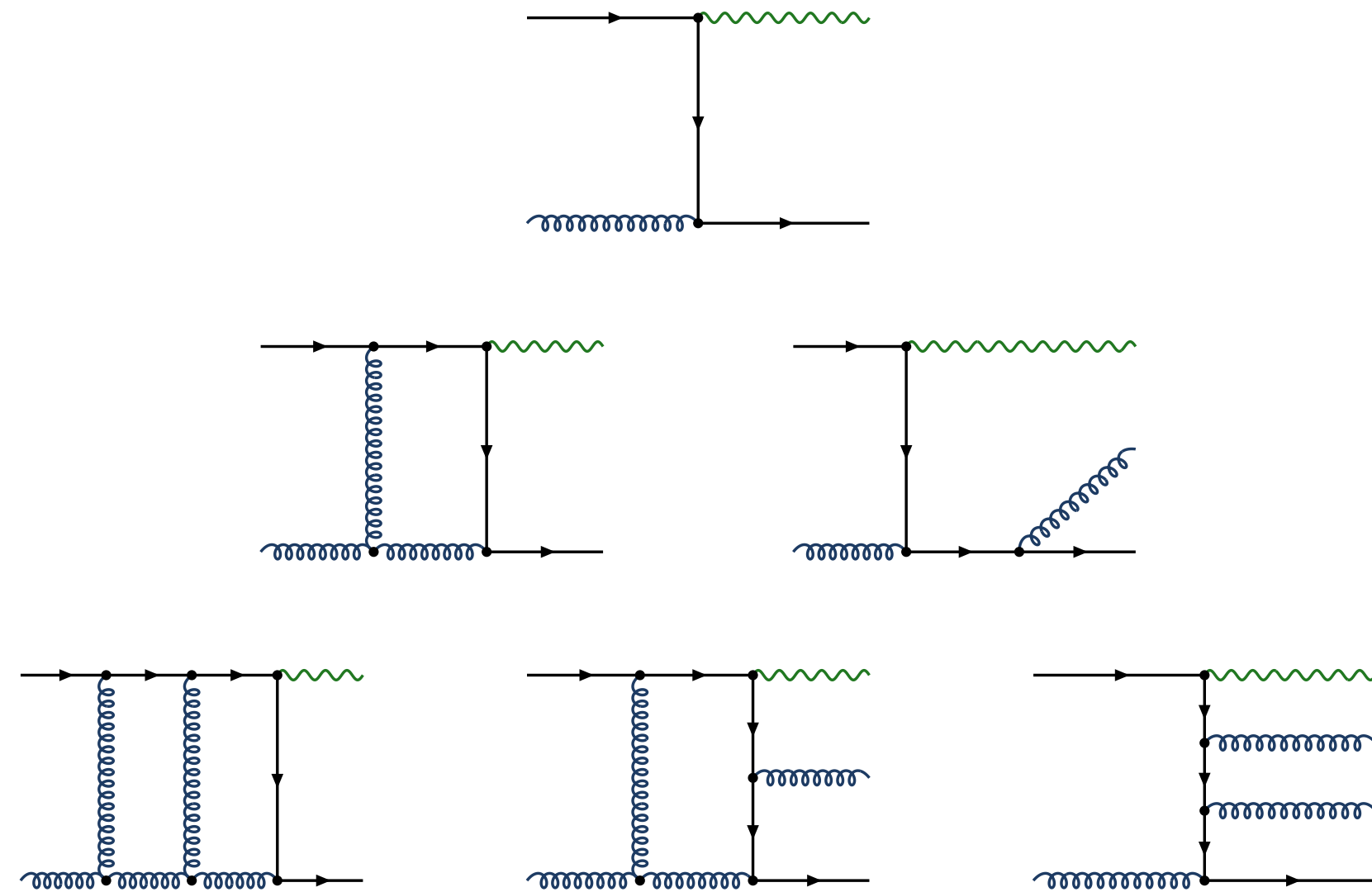
typically  $\mathcal{O}(10 - 30\%)$  precision  
here: limits the *interpretation* of data!



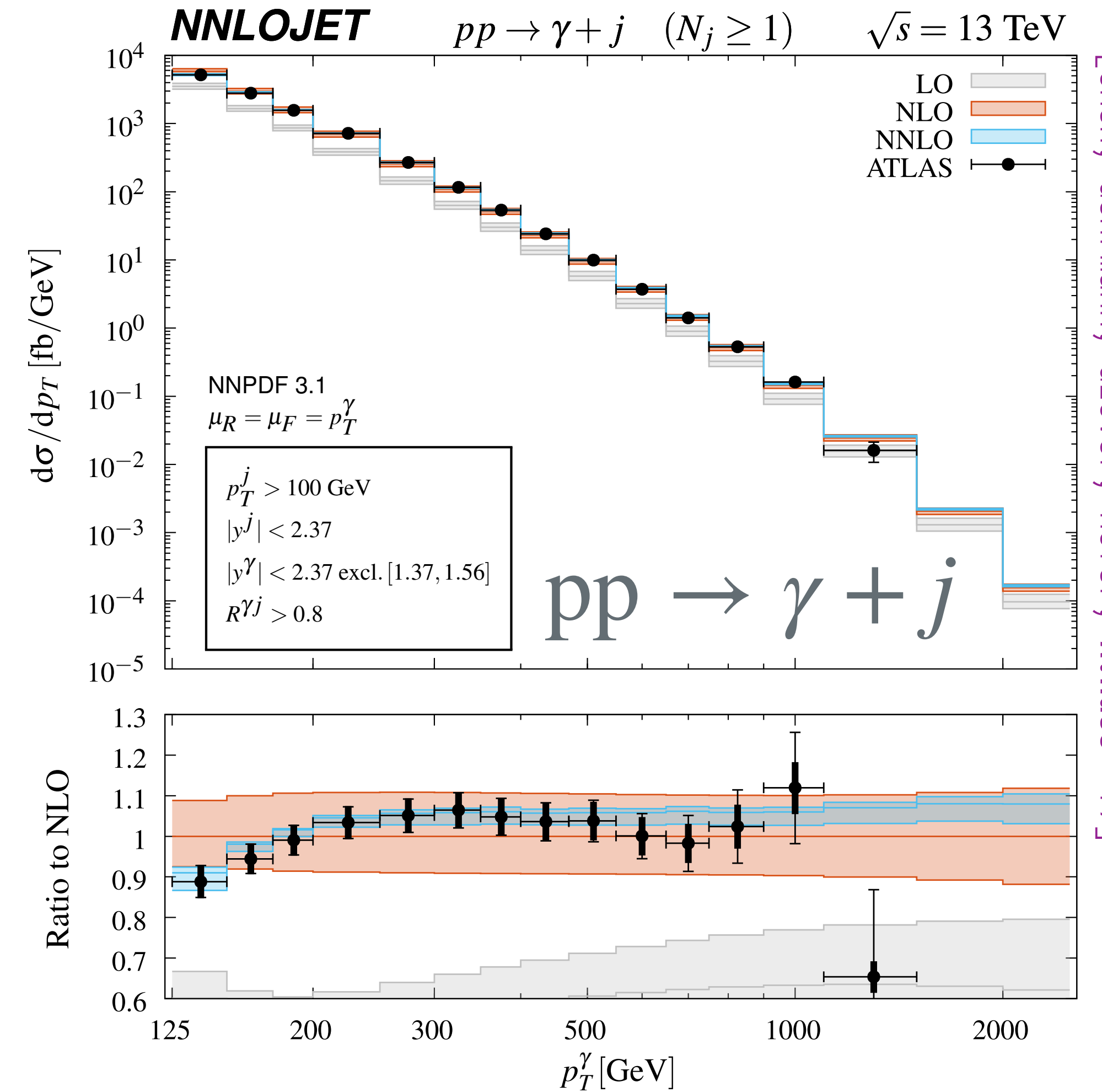
[Chen, Gehrmann, Glover, Höfer, A.Huss '19]

# PERTURBATION THEORY @ NEXT-TO-NEXT-TO-LEADING ORDER

LO  
NLO  
NNLO



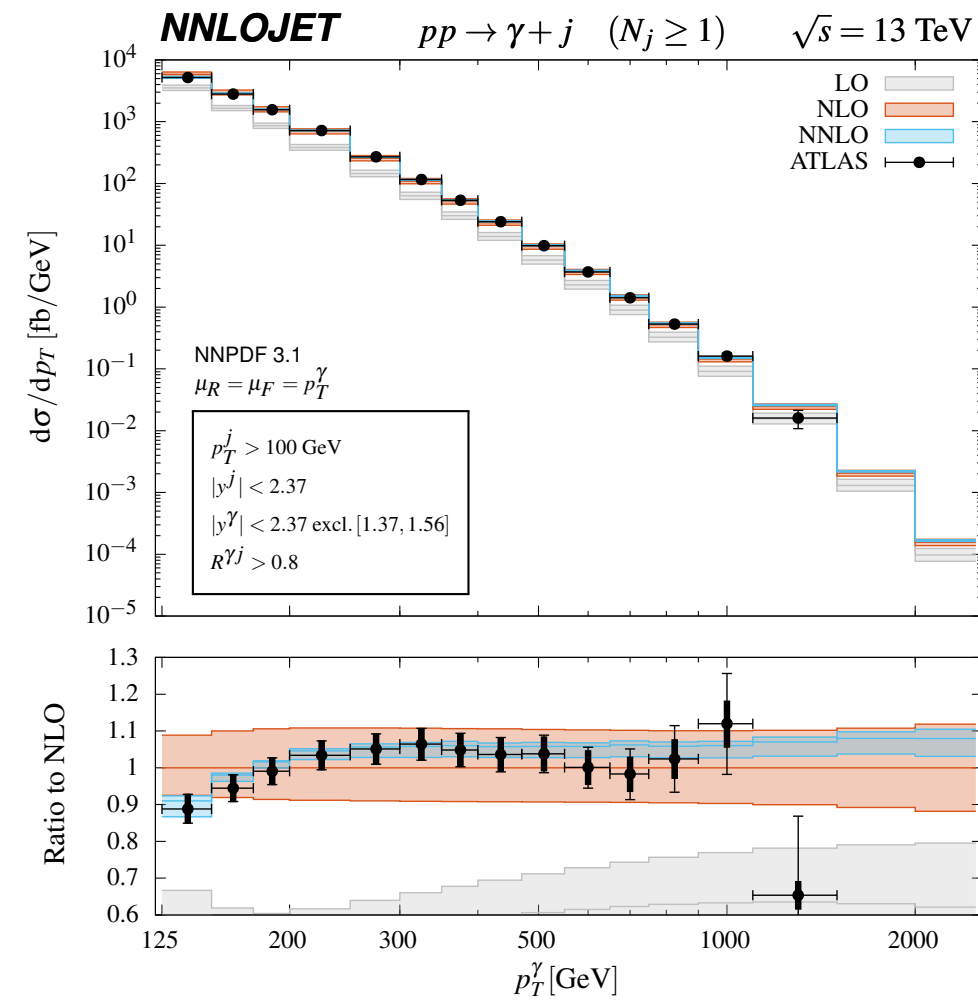
mandatory to achieve *single digit* of relative precision



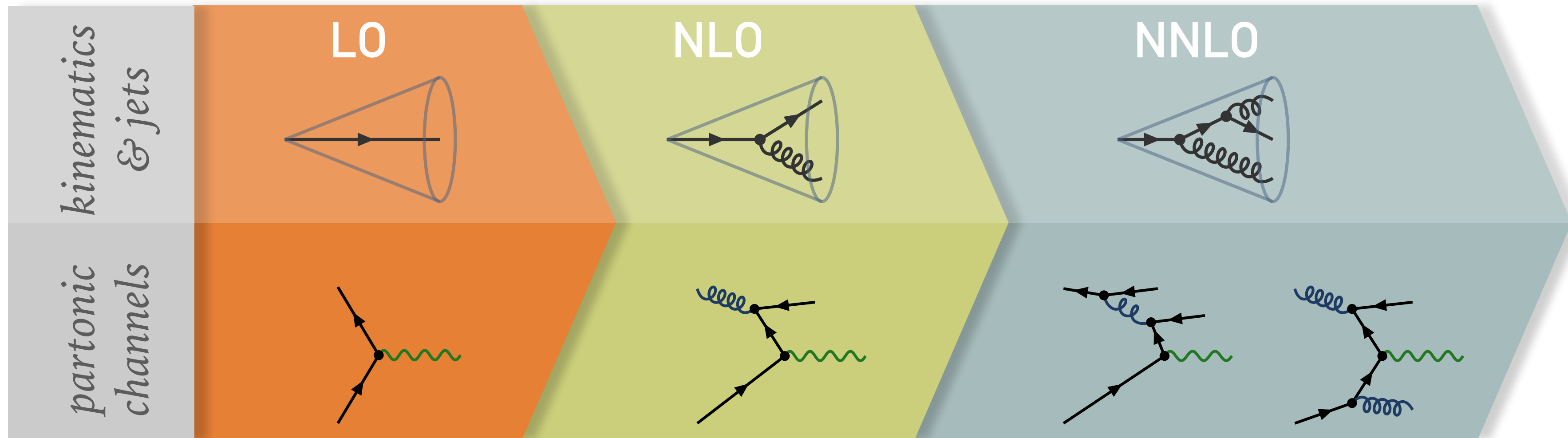
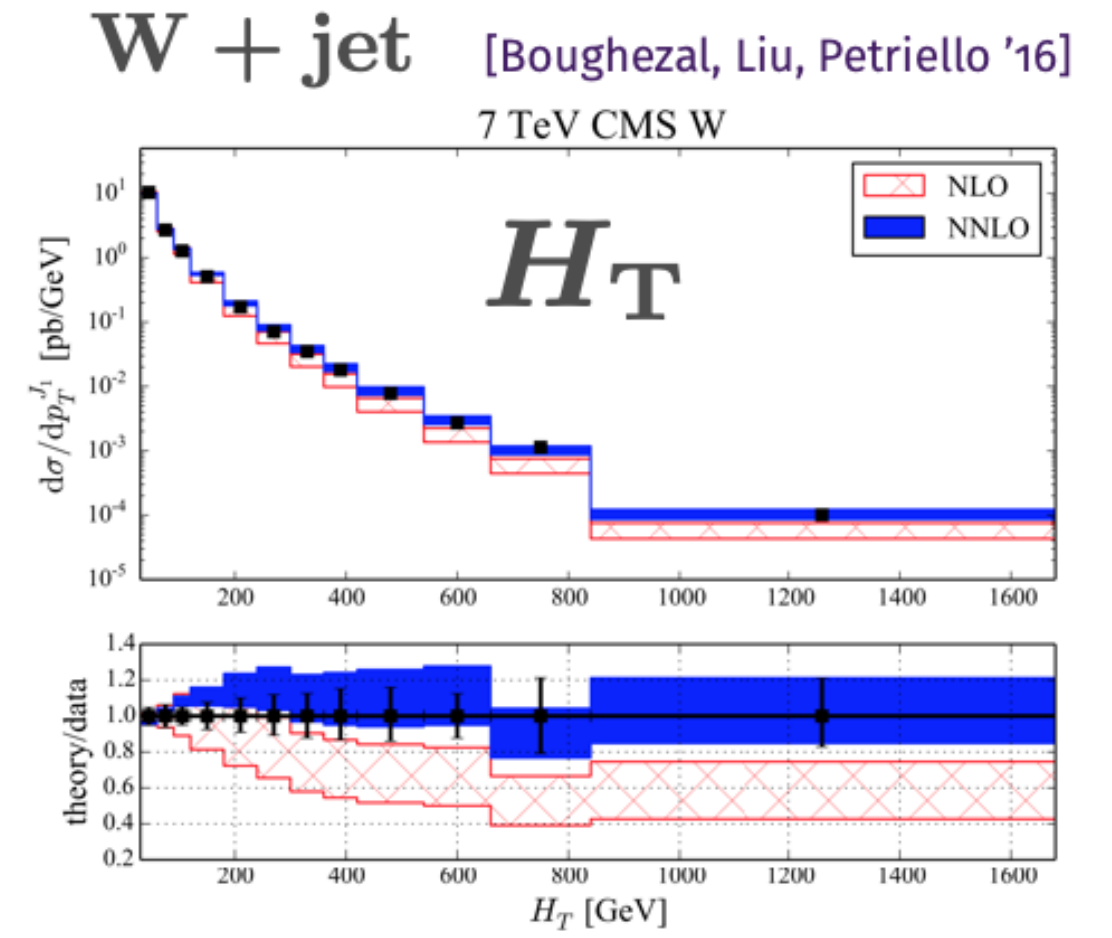
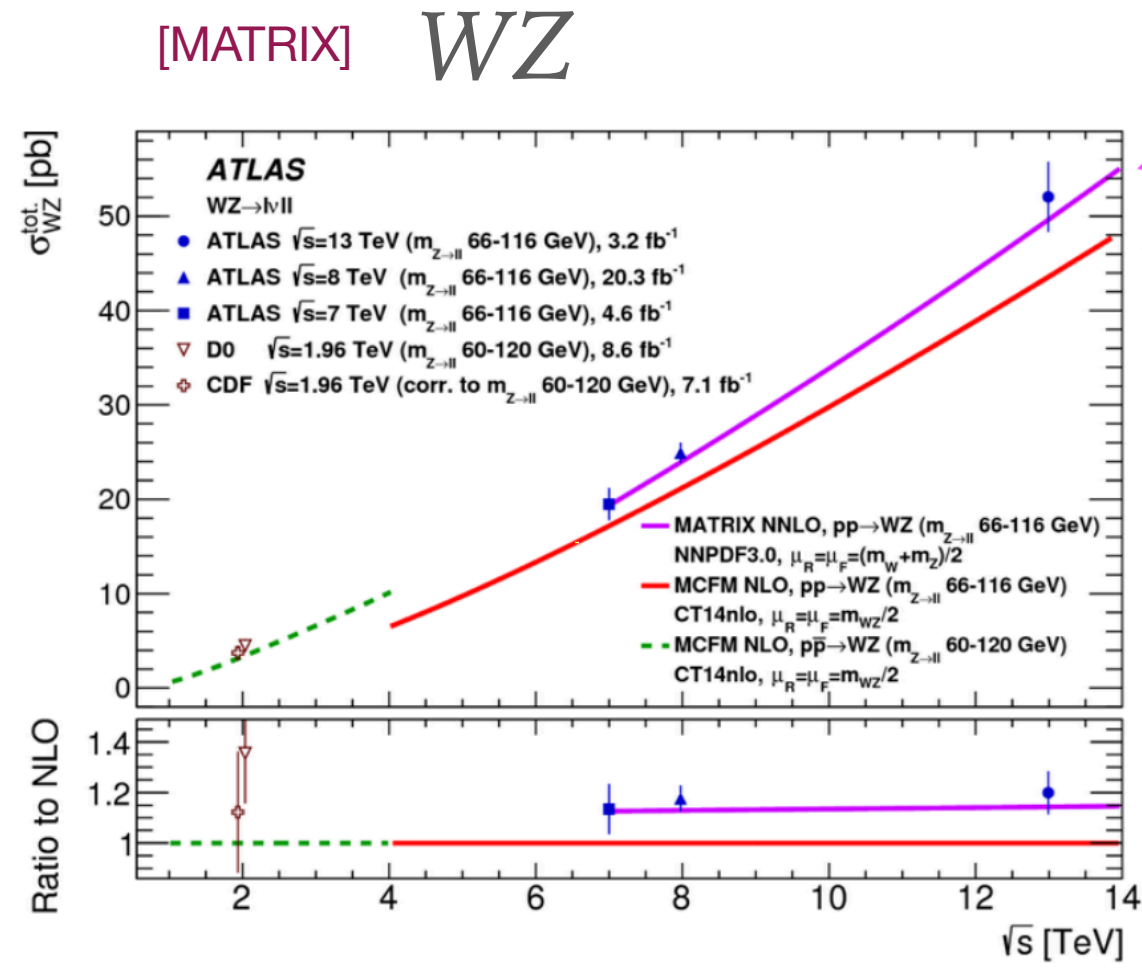
[Chen, Gehrmann, Glover, Höfer, A.Huss '19]

# EXPECT

# WHAT WE ~~HOPE~~ NNLO WILL GIVE US



[Chen, Gehrmann, Glover, Höfer, AH '19]

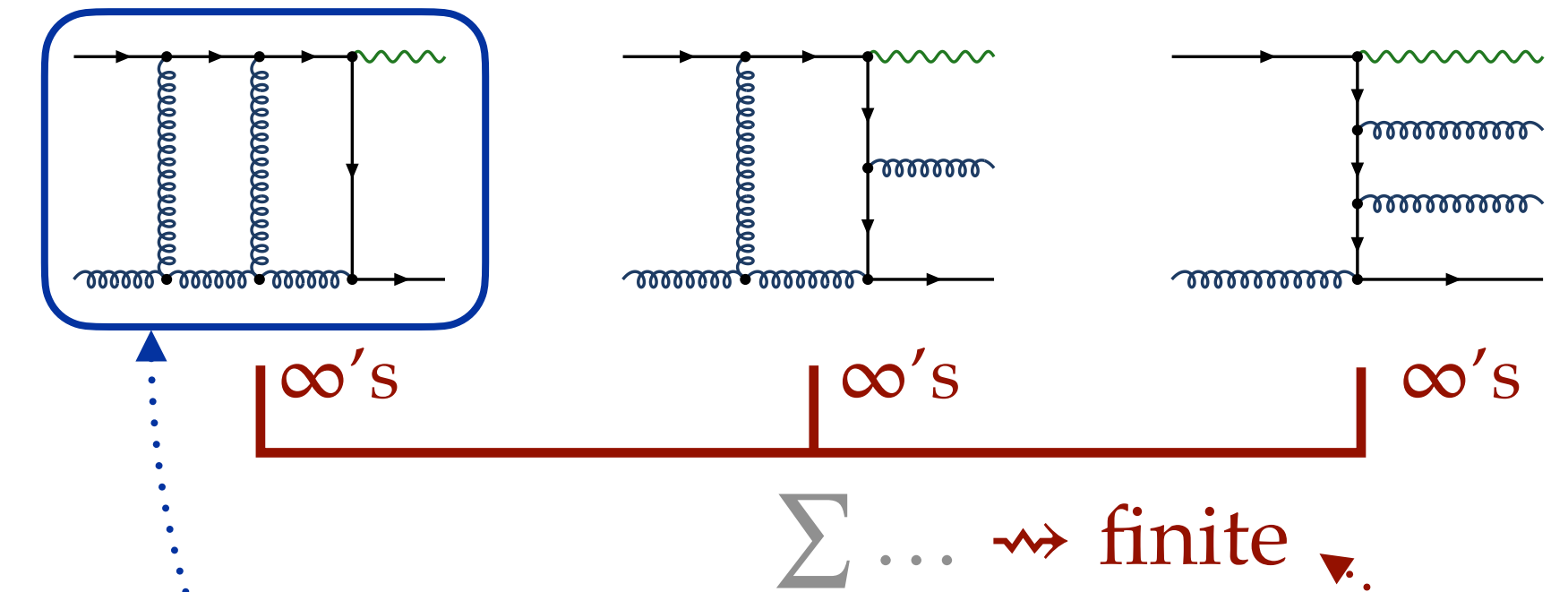




# WHAT CAN WE DO TODAY? – THE NNLO TIMELINE

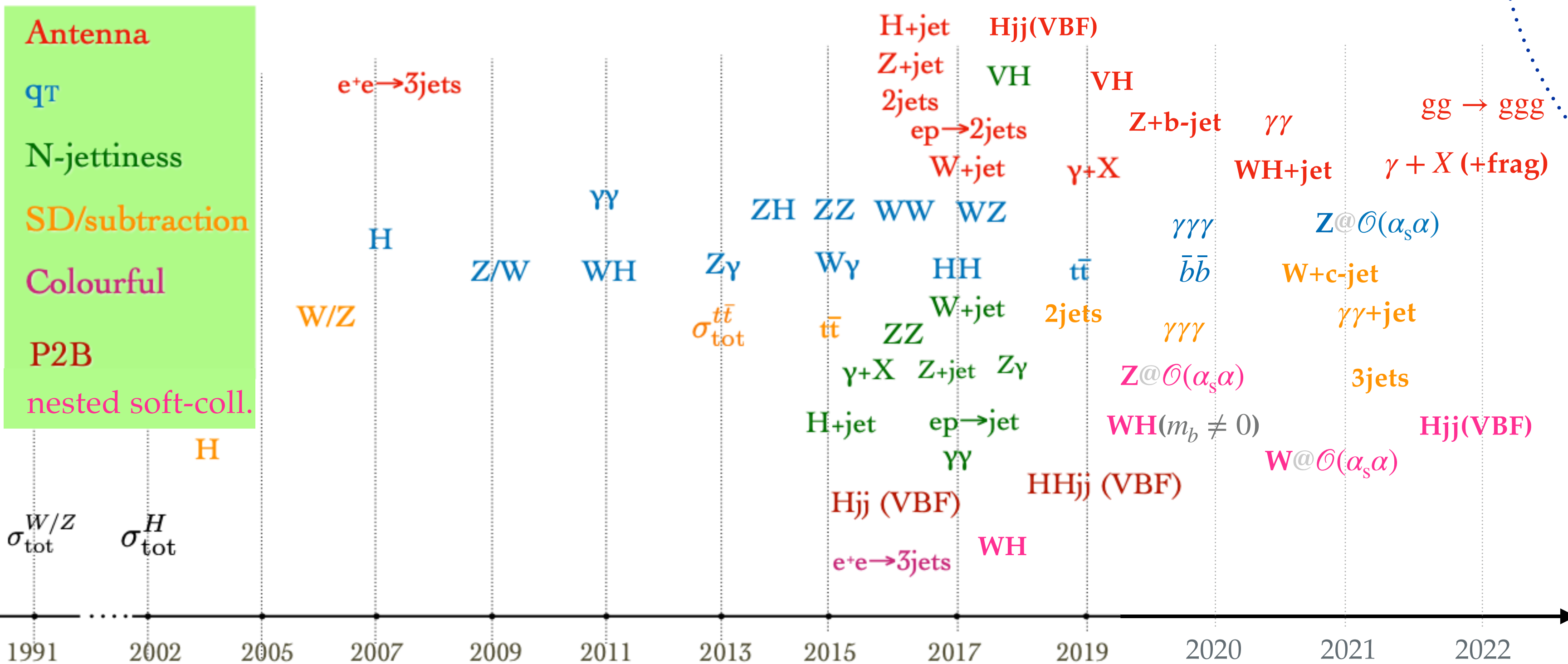
Tremendous progress in the past  $\sim 5-10$  years!

$\hookrightarrow 2 \rightarrow 2$  under good control;  $2 \rightarrow 3$  next frontier



## Main challenges:

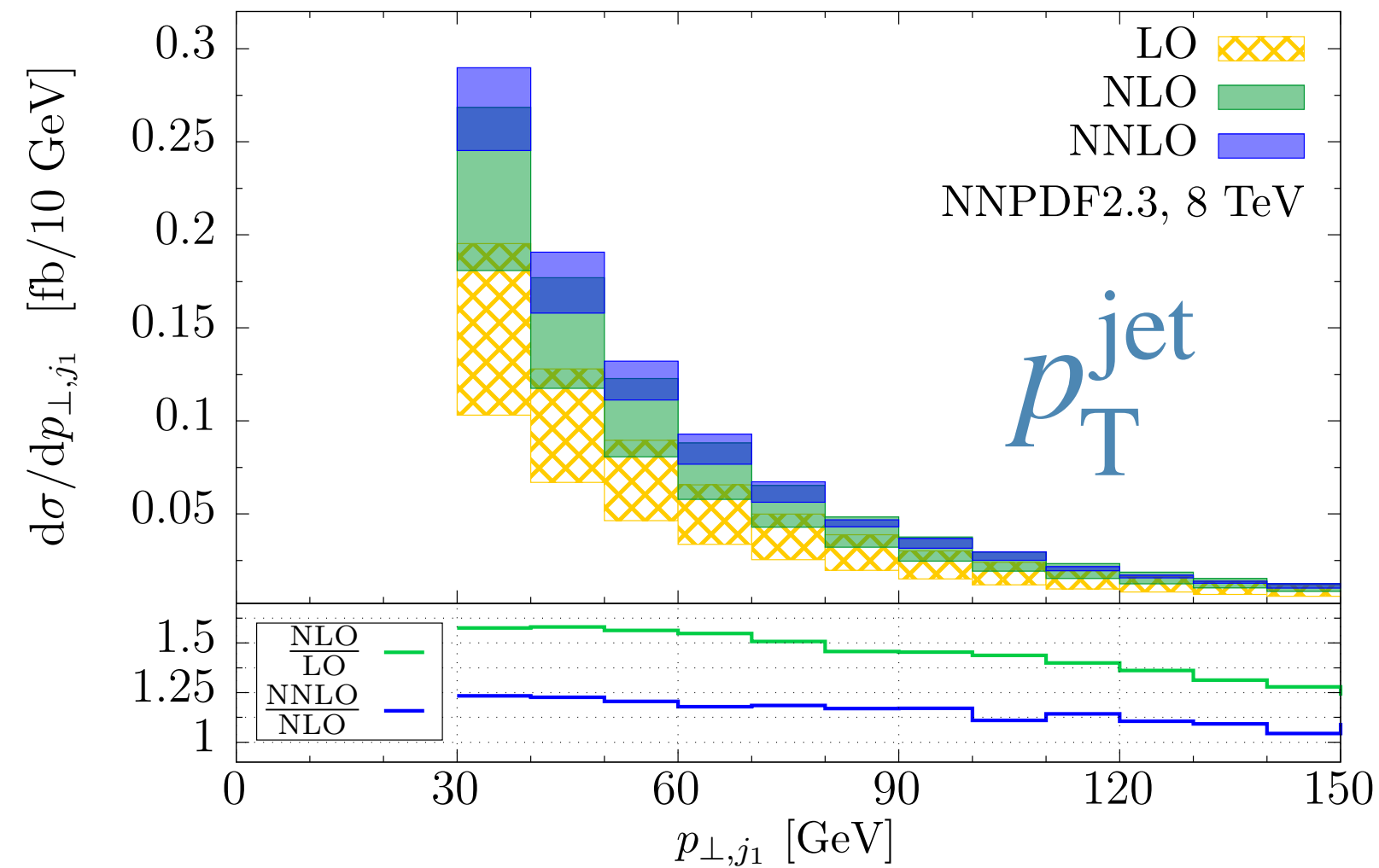
- amplitudes & multi-loop integrals
- infrared subtractions



# INDEPENDENT CALCULATIONS — H + jet $\times 3!$

## residue subtraction

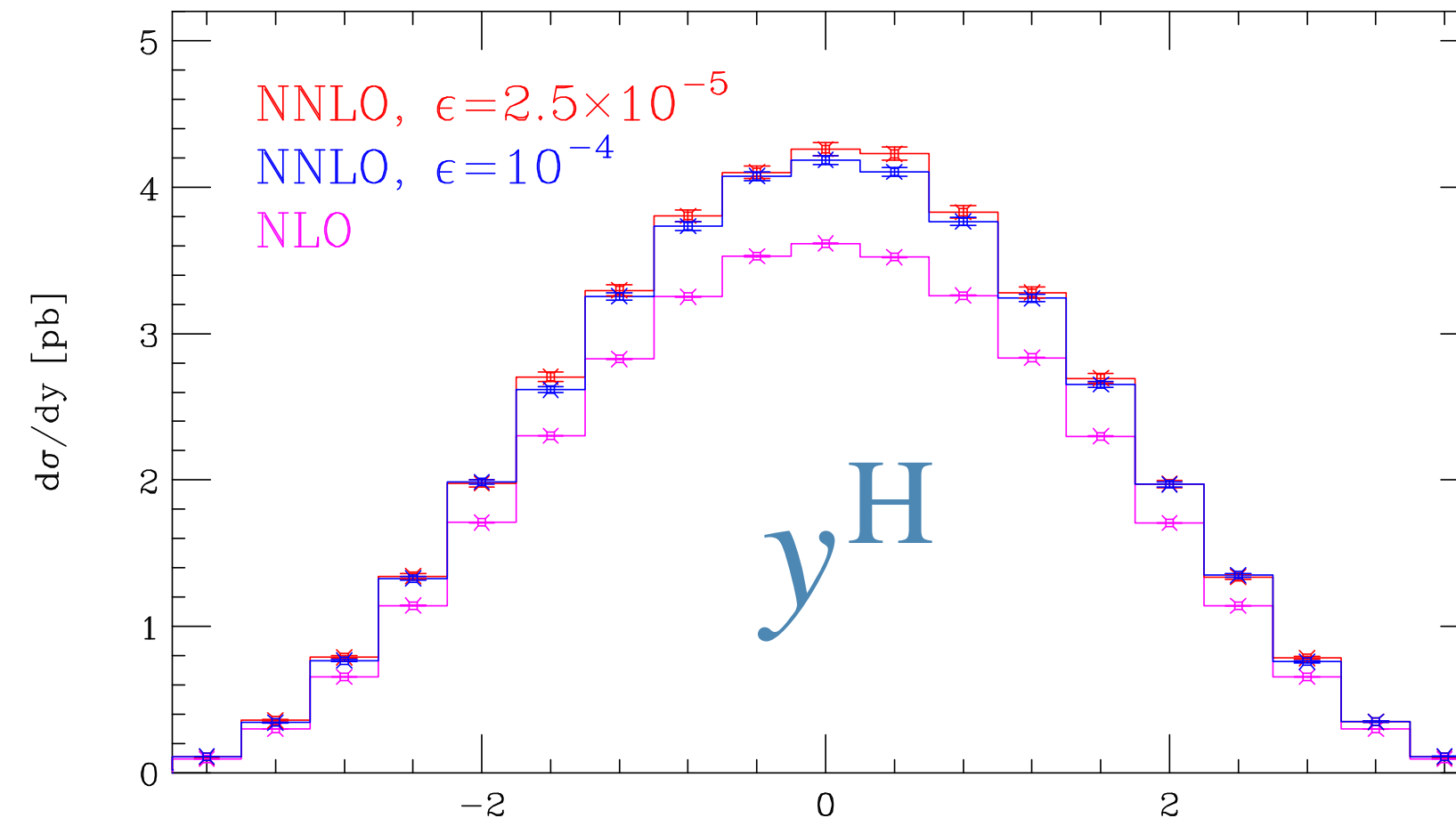
[Caola, Melnikov, Schulze '15]



## $\tau_1$ jettiness subtraction

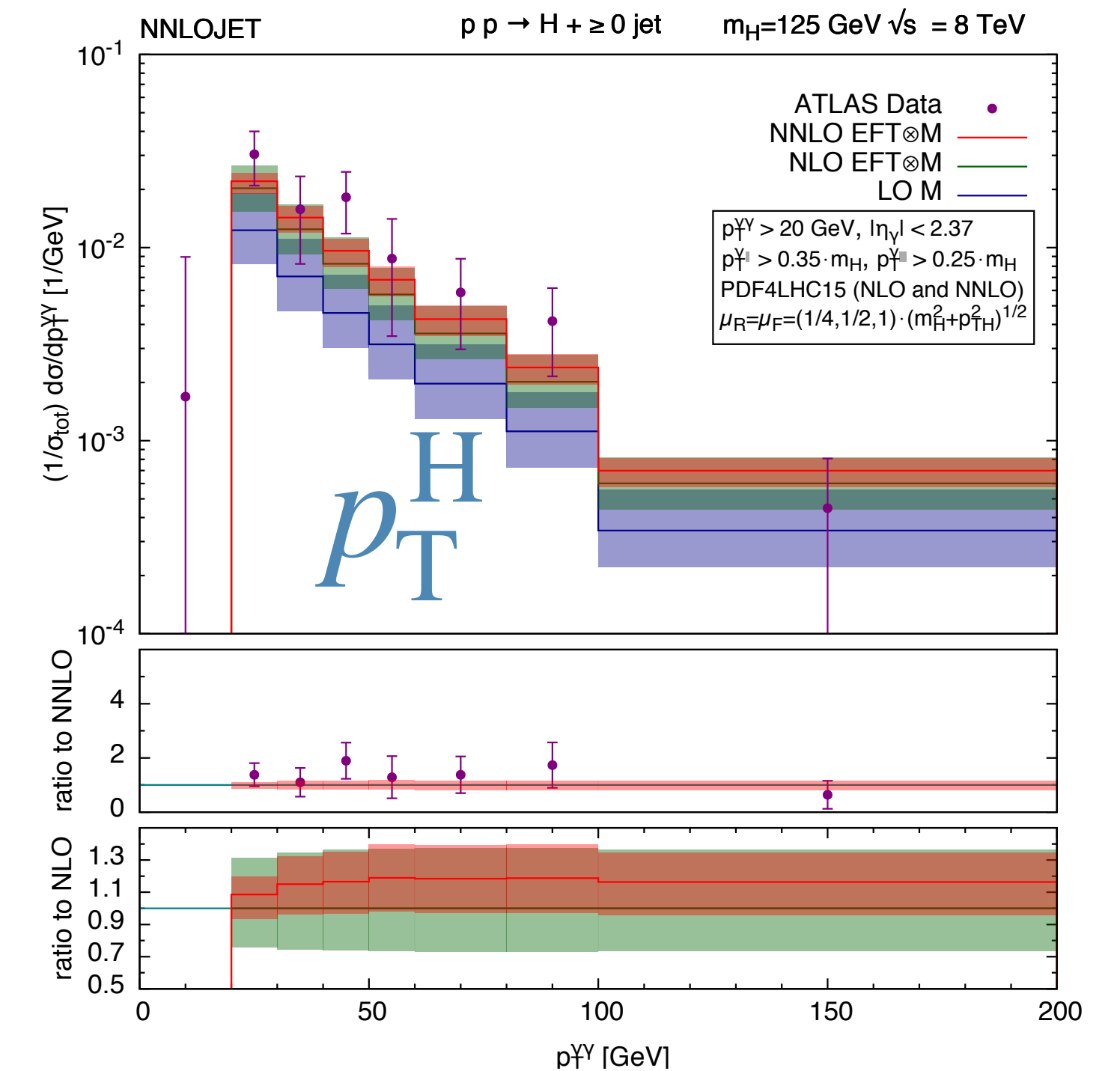
[Boughezal, Focke, Giele, Liu, Petriello '15]

[Campbell, Ellis, Seth '19]



## antenna subtraction

[Chen, Cruz-Martinez, Gehrmann, Glover, Jaquier '16]

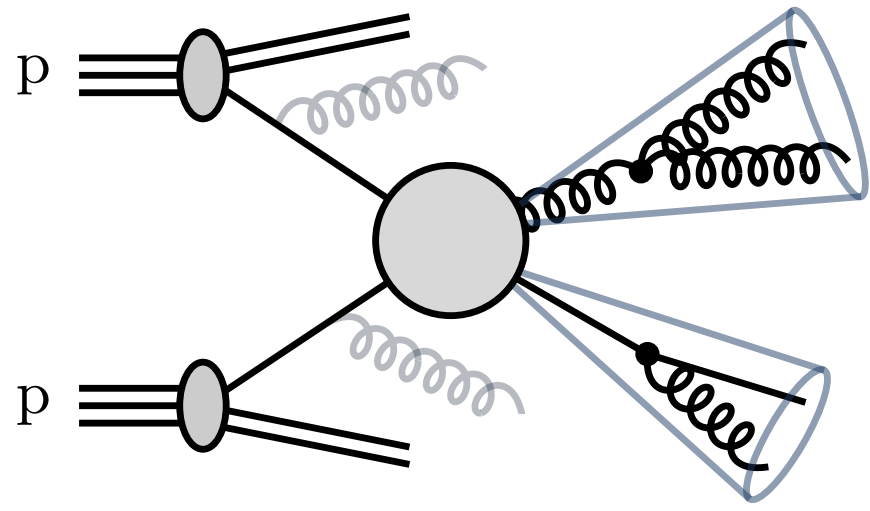


very complex calculations  $\leftrightarrow$  validation!

- long-standing [ $\sim$ '15] discrepancy in H + jet  
 $\hookrightarrow$  only resolved in ['19]

benchmark approaches

# JETS ARE...



the ideal pQCD laboratory  
simple  $2 \rightarrow 2$  parton scattering

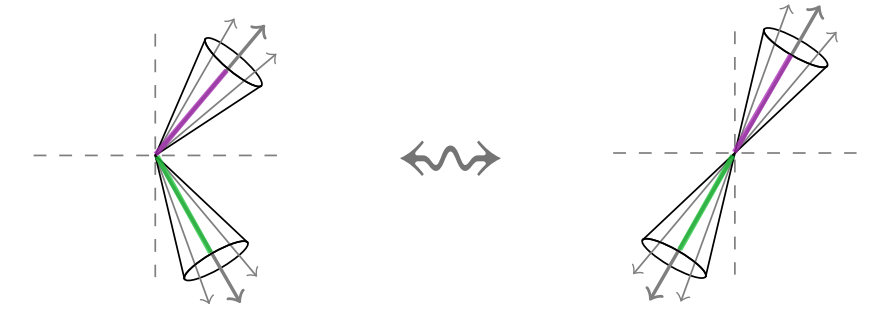
produced in *abundance*

wide kinematical range

Constrain PDFs  
(high- $x$  gluon)

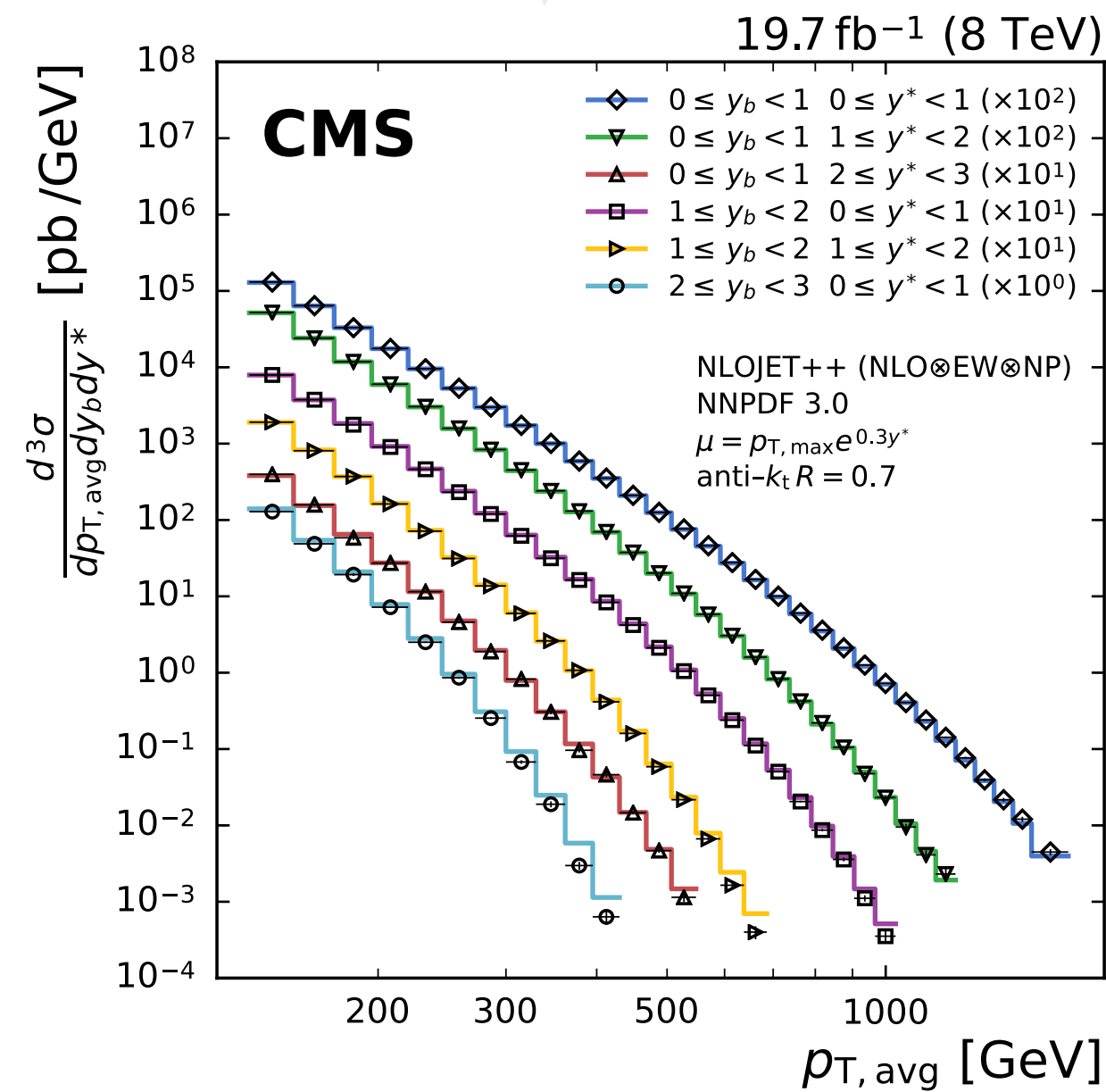
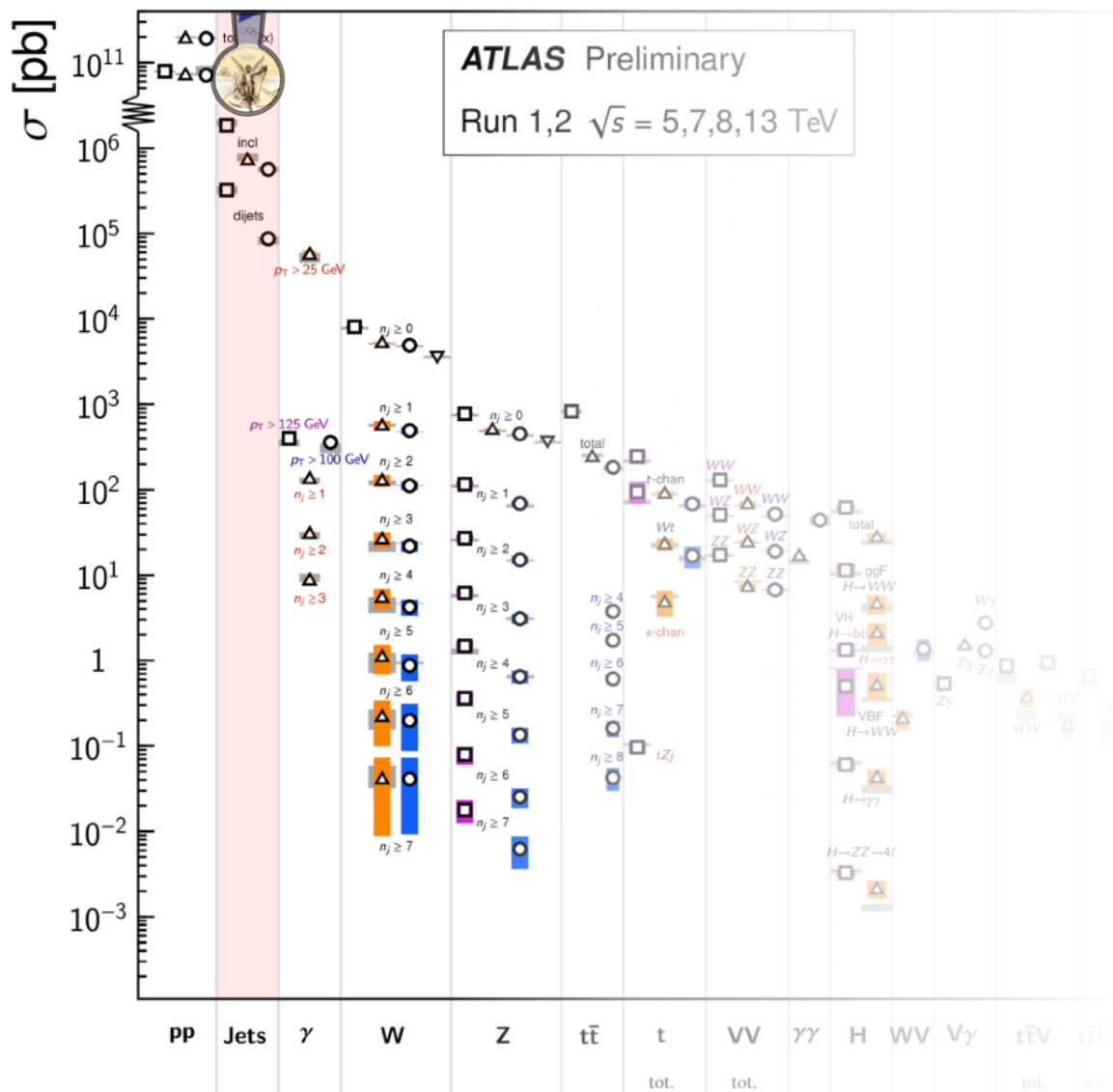
$$x = \frac{p_T}{\sqrt{s}} (e^{\pm y_j} + e^{\pm y_j'})$$

- @ LO: 3 variables ( $p_T, y_j, y_j'$ )
- inclusive jet[2] (some smearing)

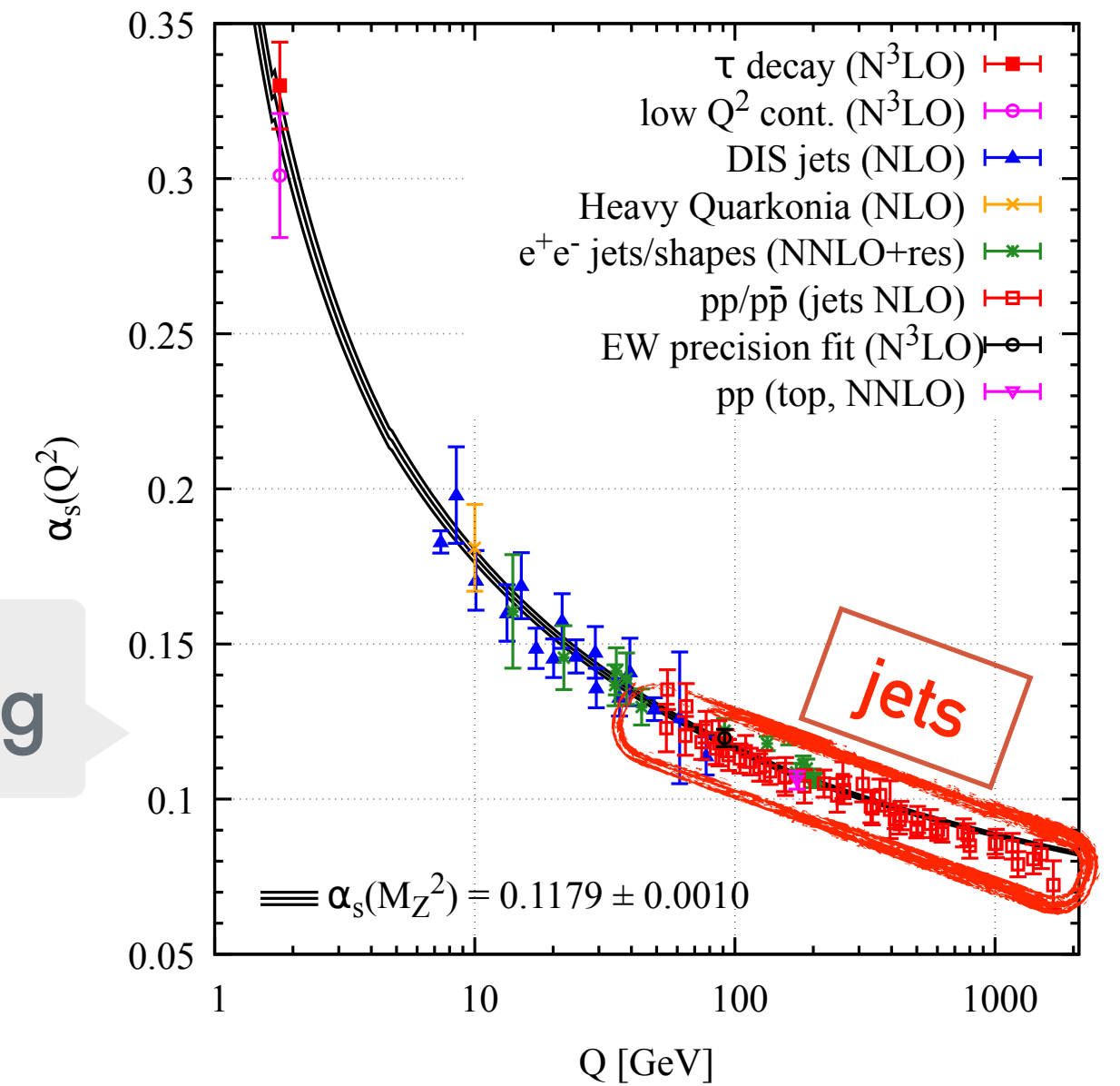


- di-jet[3] (reconstructible: 3-D)


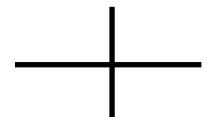
Standard Model Production Cross Section Measurements



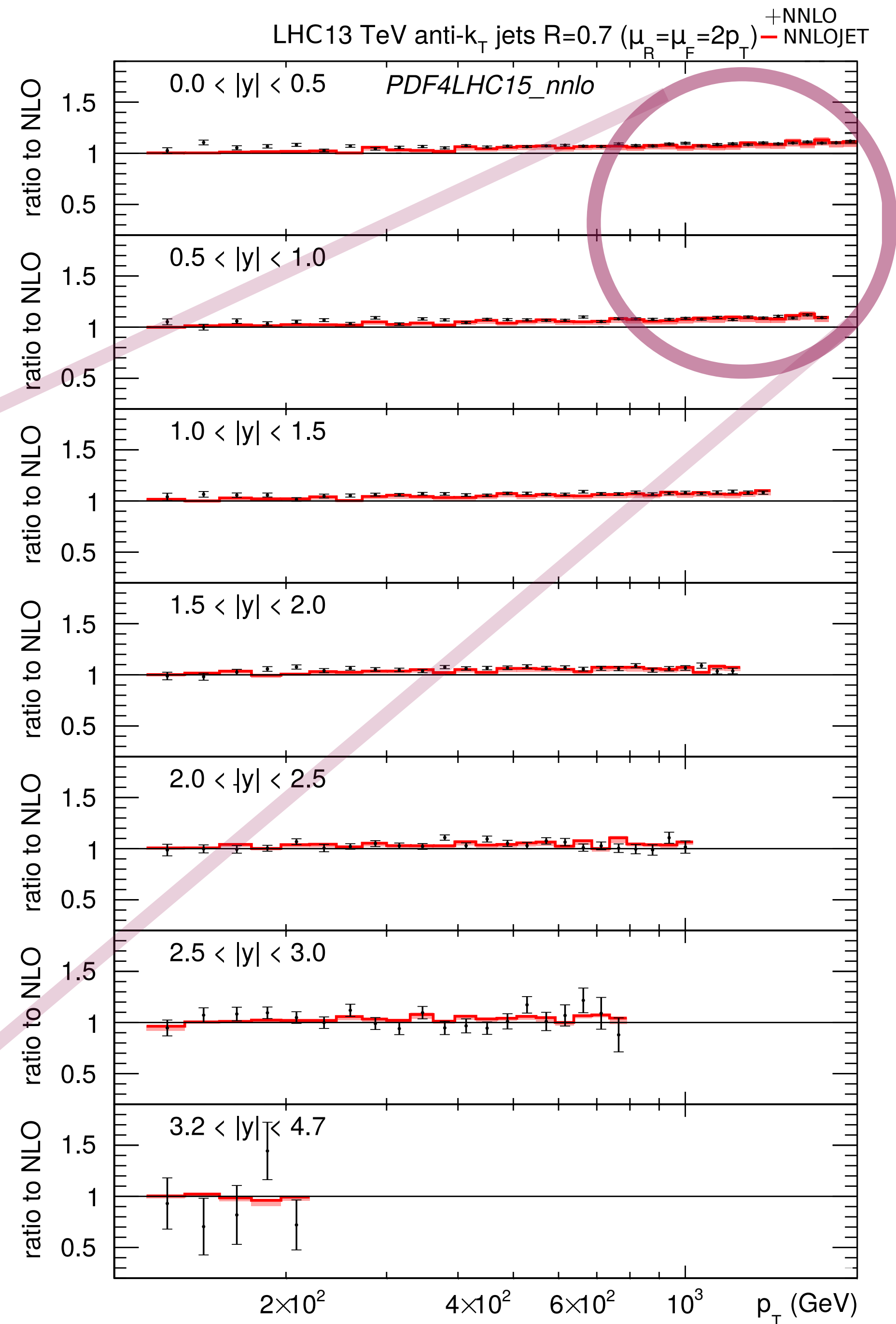
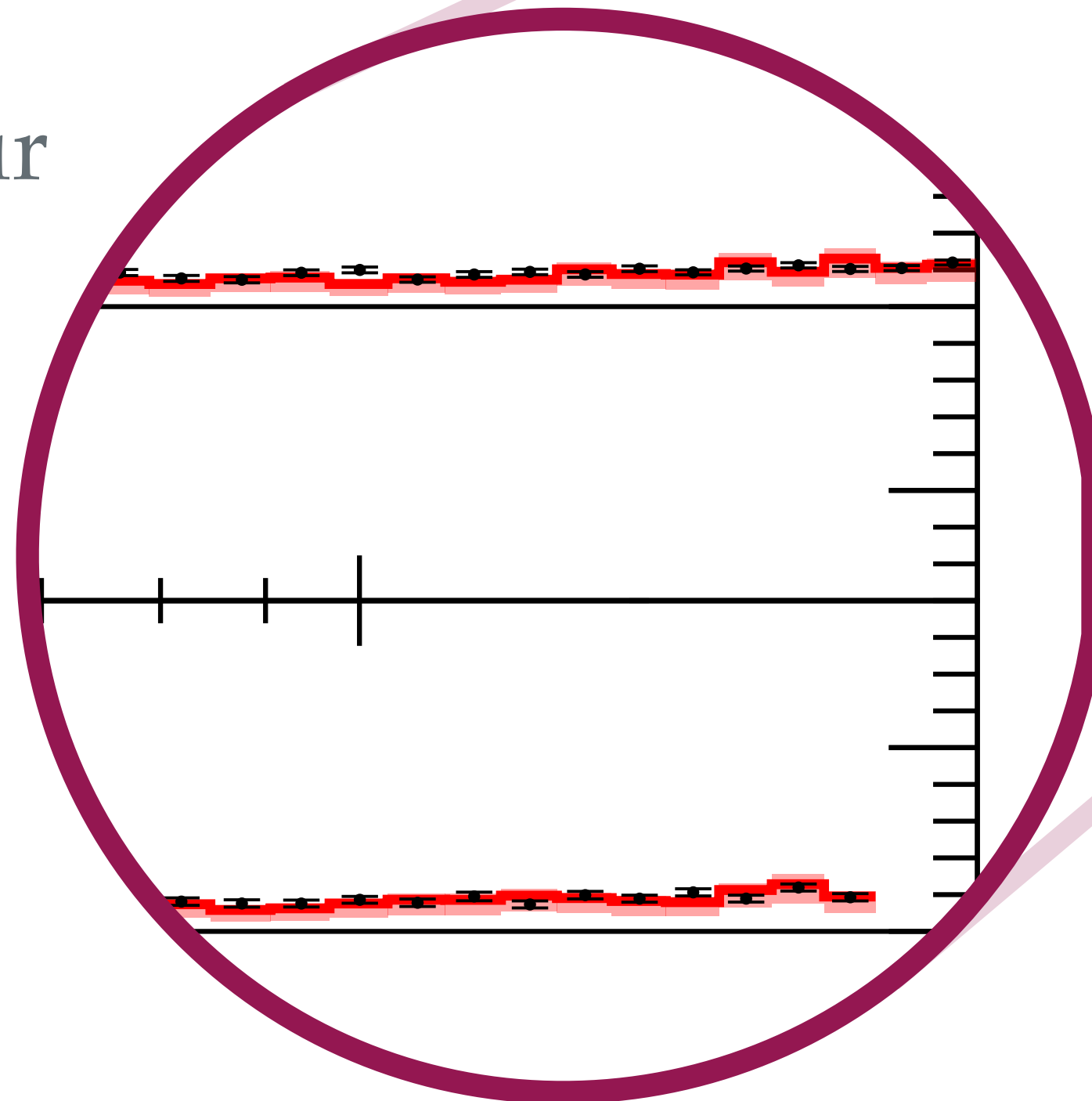
$\alpha_s$  & running



# INCLUSIVE JETS – 2 CALCULATIONS!

-  **NNLOJET** [Currie, Glover, Pires '16]
-  **STRIPPER** [Czakon, van Hameren, Mitov, Poncelet '19]

- in very good agreement!
- sub-leading colour negligible!(?)



# FAST INTERPOLATION GRIDS — APPLFAST

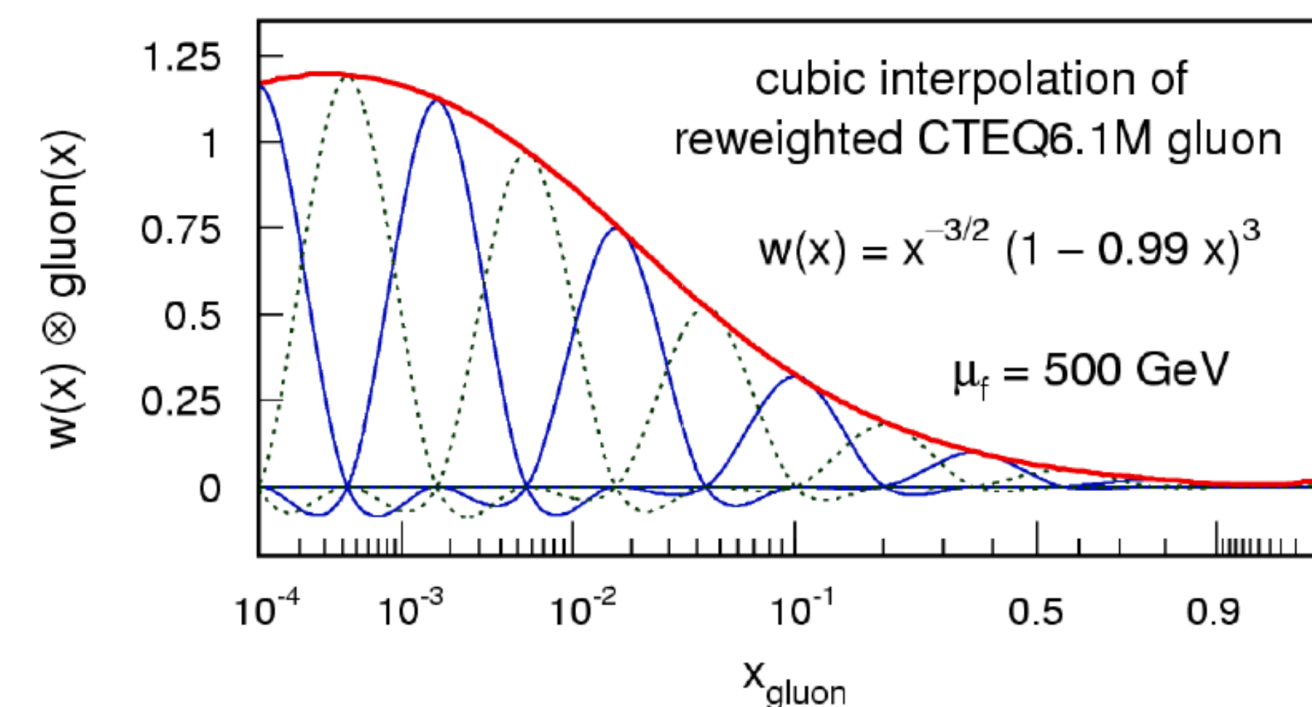
[APPLgrid, fastNLO, NNLOJET '19, '22]

- NNLO calculations  $\mathcal{O}(100k)$  CPU hours  $\rightsquigarrow$  prohibitive in PDF &  $\alpha_s$  fits!

$\hookrightarrow$  approximate the costly convolution using a grid:

$$\sigma = \int_0^1 dx f_a(x) \alpha_s^n \hat{\sigma}_a(x) \simeq \sum_i f_a(x^{(i)}) \alpha_s^n \underbrace{\left[ \int_0^1 dx \hat{\sigma}_a(x) E^{(i)}(x) \right]}_{\hat{\sigma}_a^{(i)}}$$

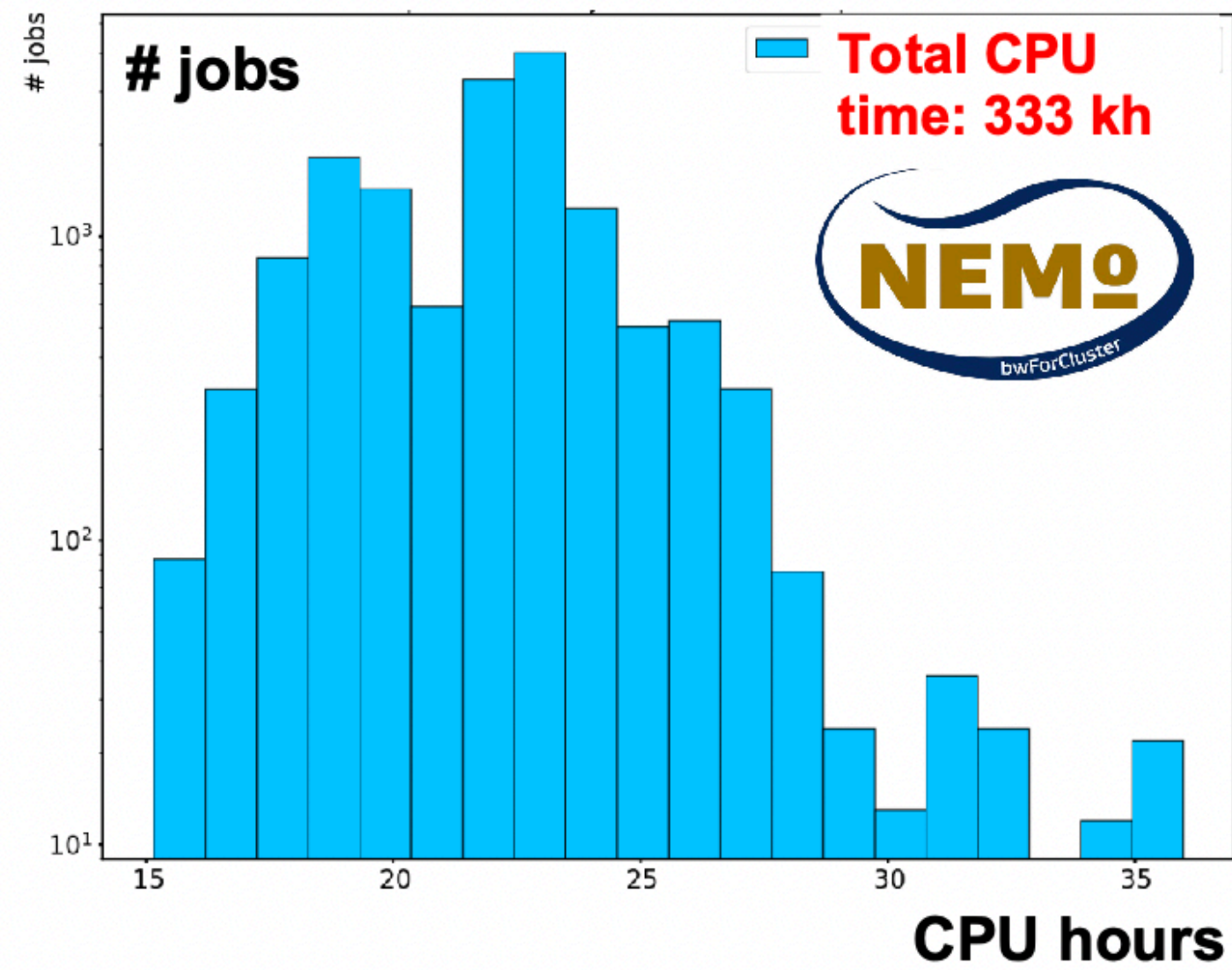
*sum: cheap!*



$$f_a(x) \simeq \sum_i f_a(x^{(i)}) E^{(i)}(x)$$

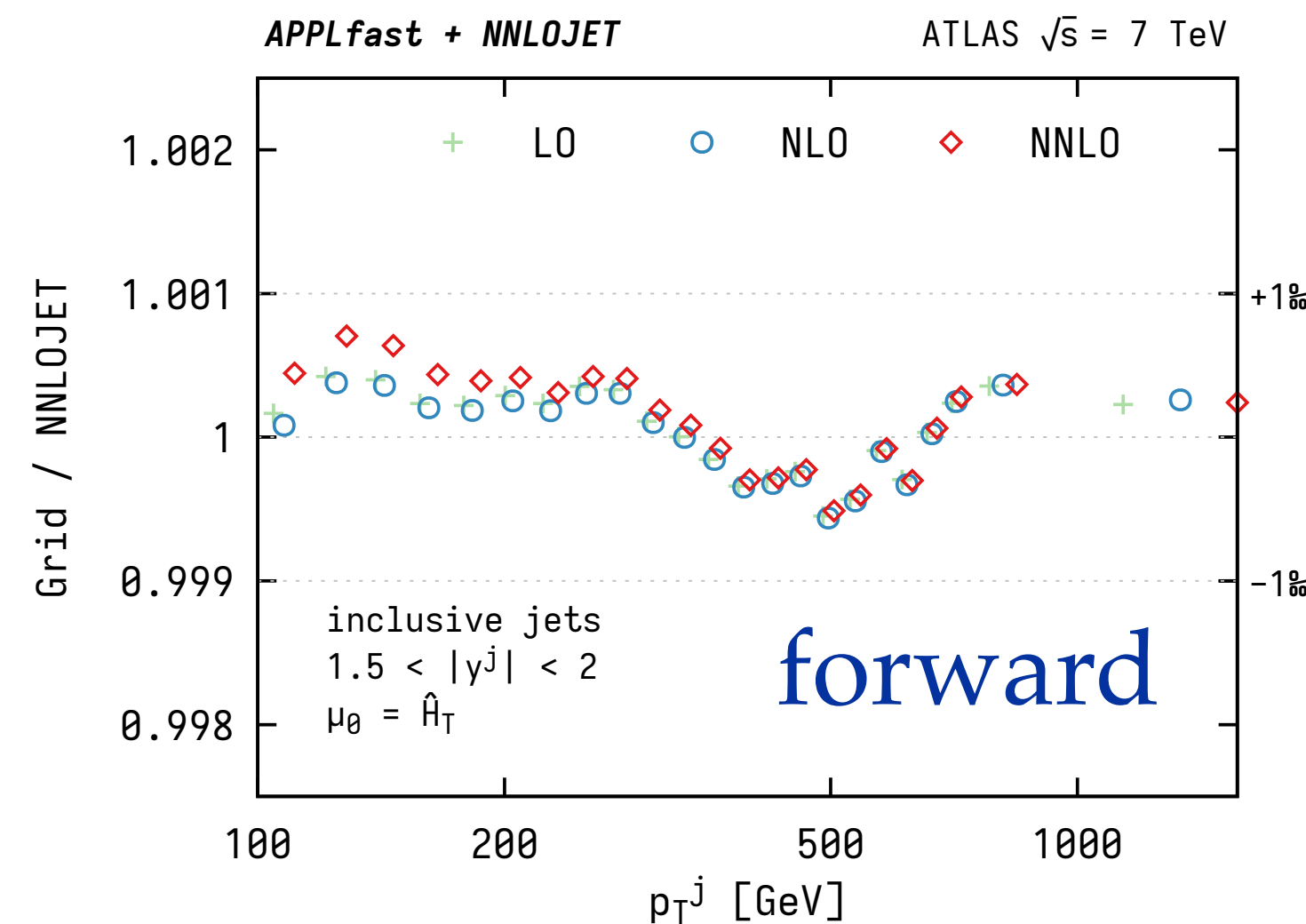
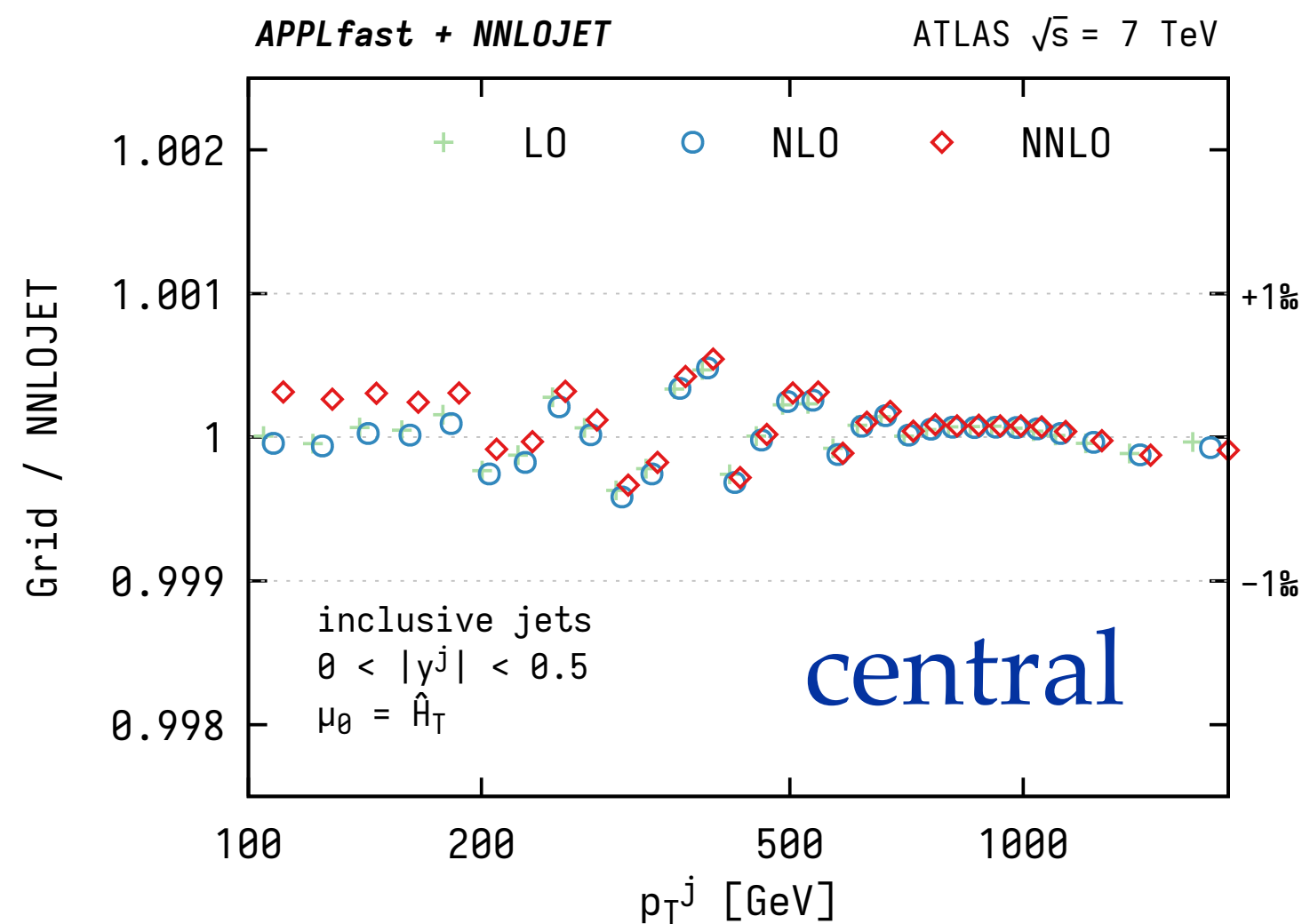
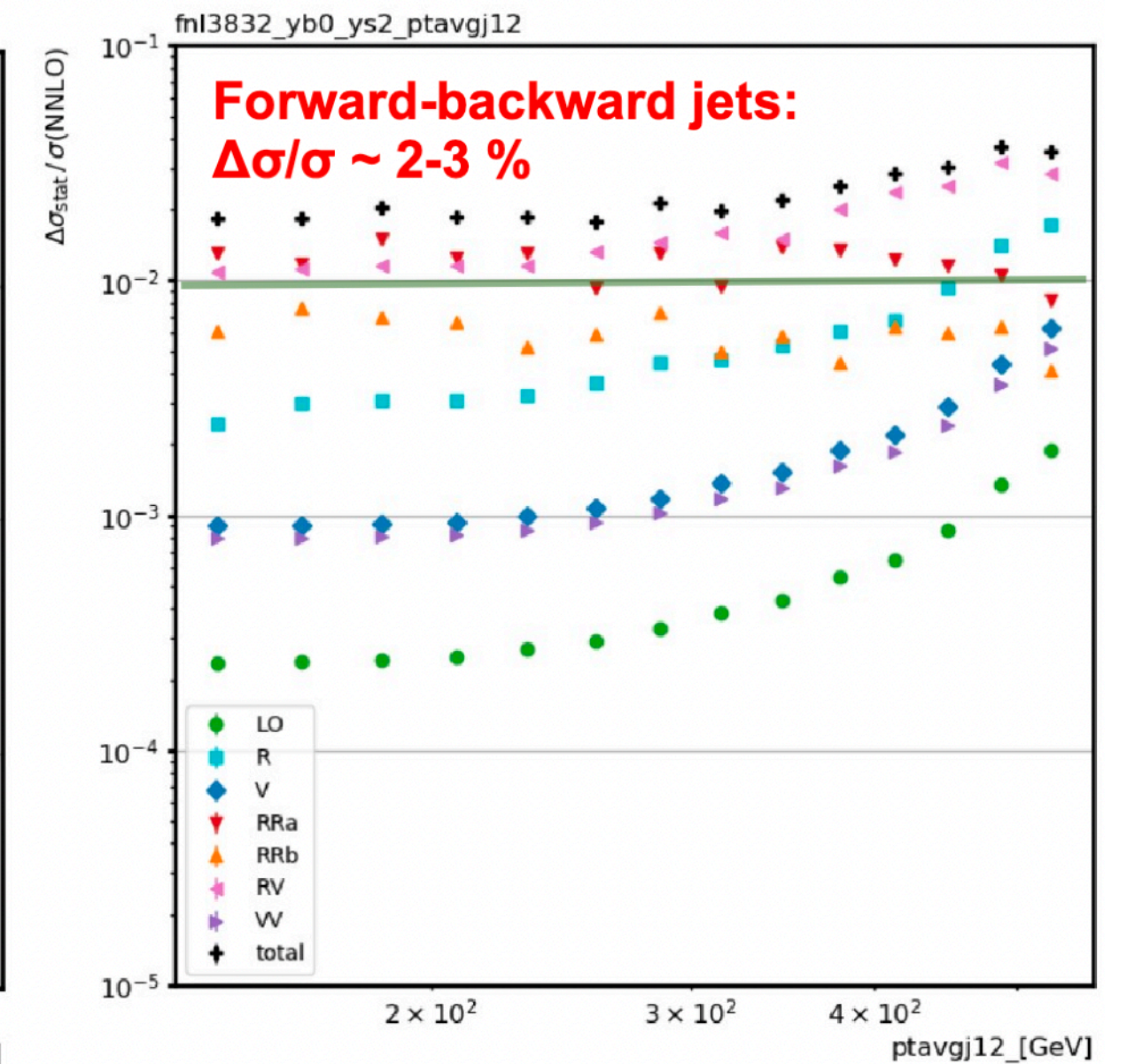
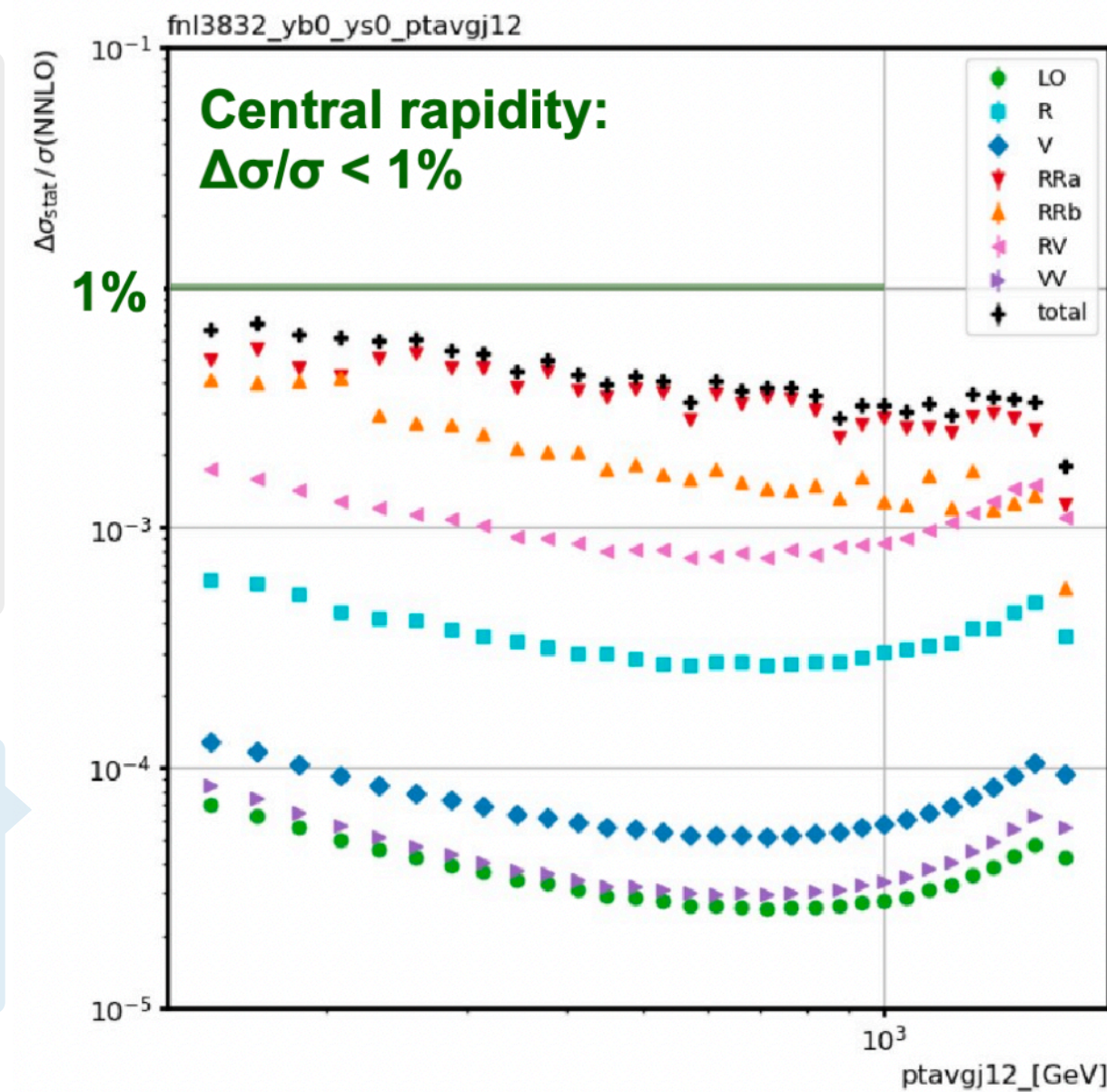
# THE INVESTMENT

[APPLgrid, fastNLO, NNLOJET '22]



grid generation  
 $\leftrightarrow$  runtime  $\sim \times 2-3$   
 (fastNLO)  
 $\leftrightarrow$  storage ;-)

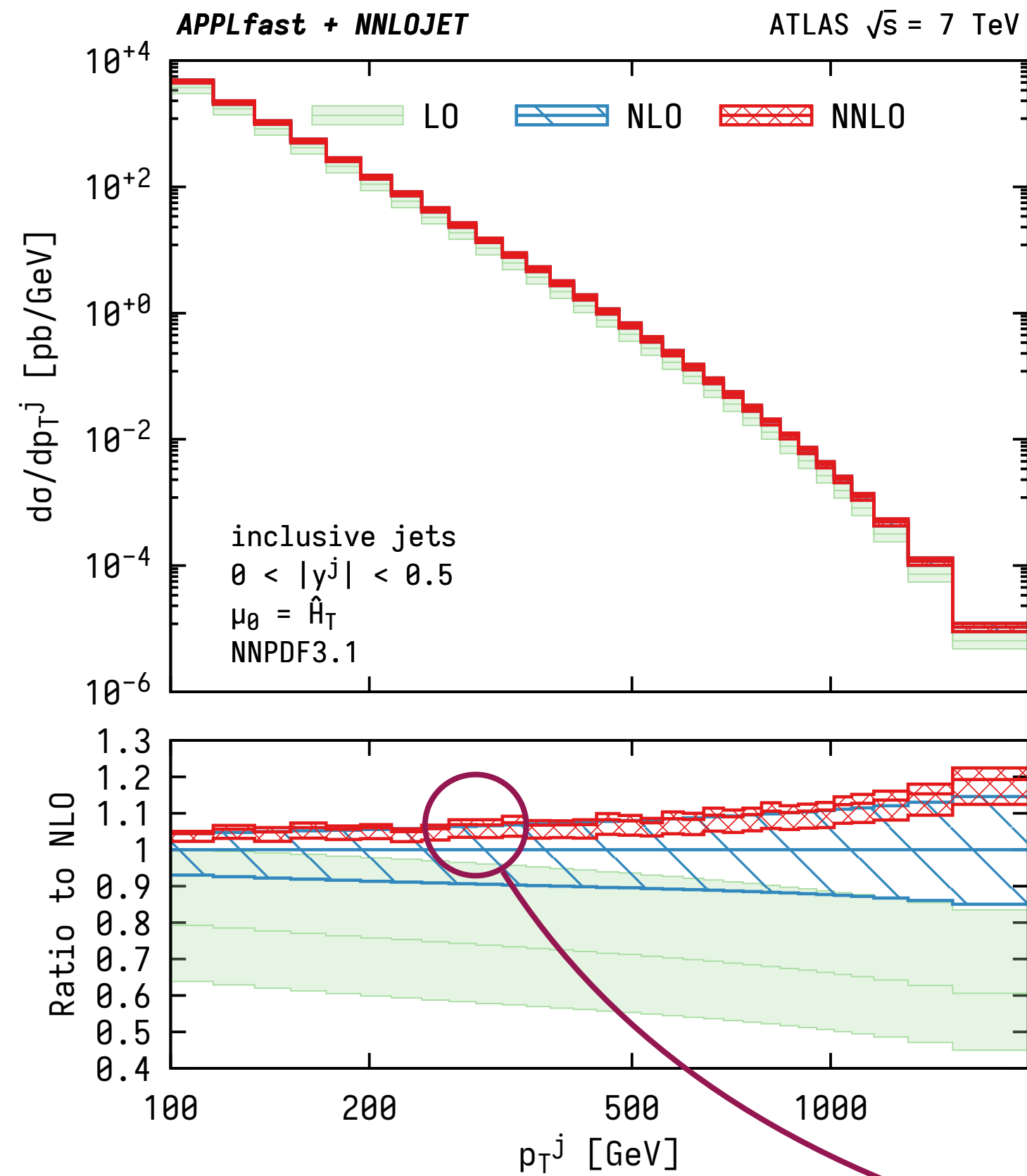
MC uncert.  
 $\lesssim$  few %



closure tests  
 grid evaluation vs.  
 “vanilla” NNLOJET  
 $\hookrightarrow$  interp. bias  $\lesssim 0.1\%$

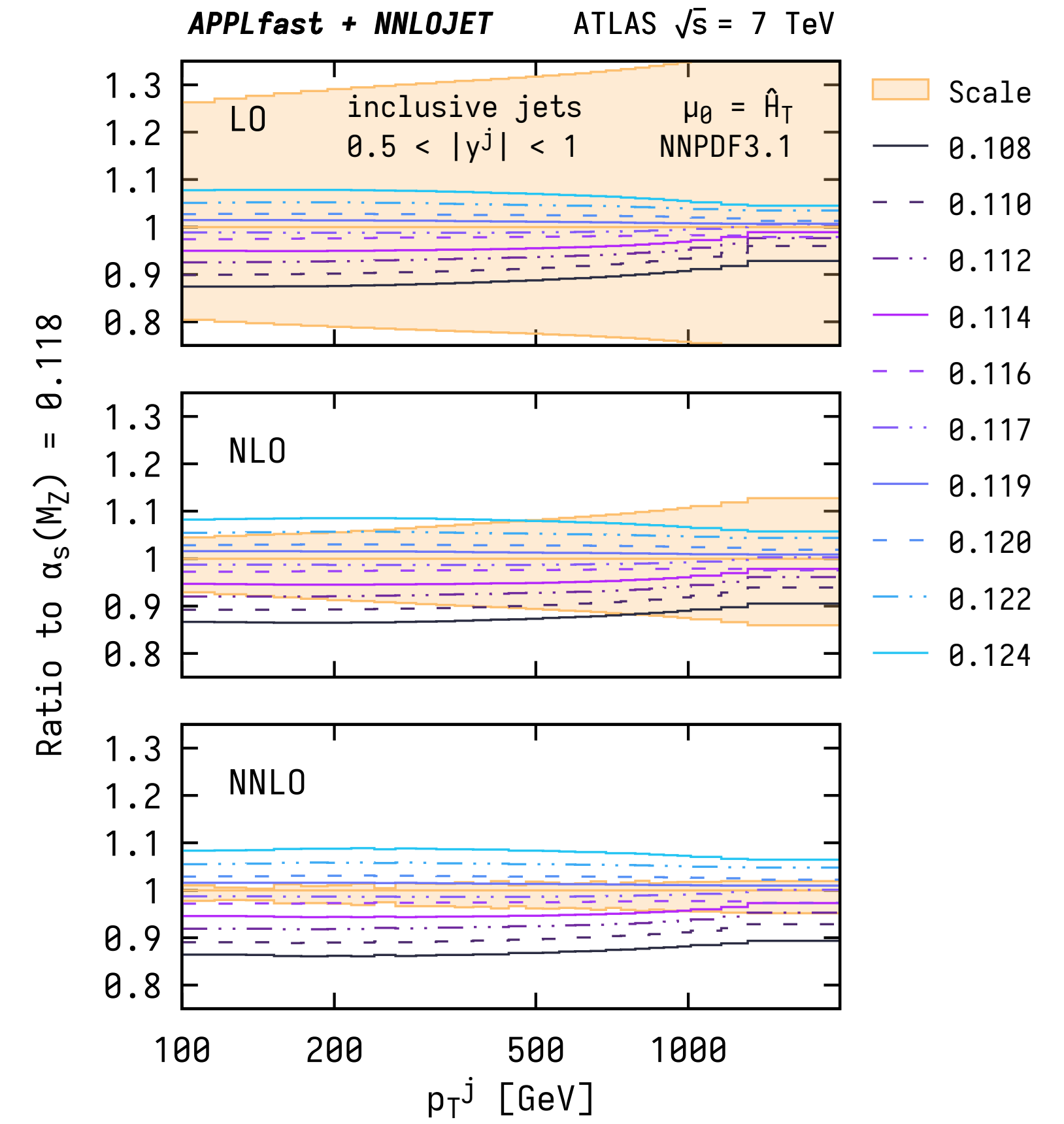
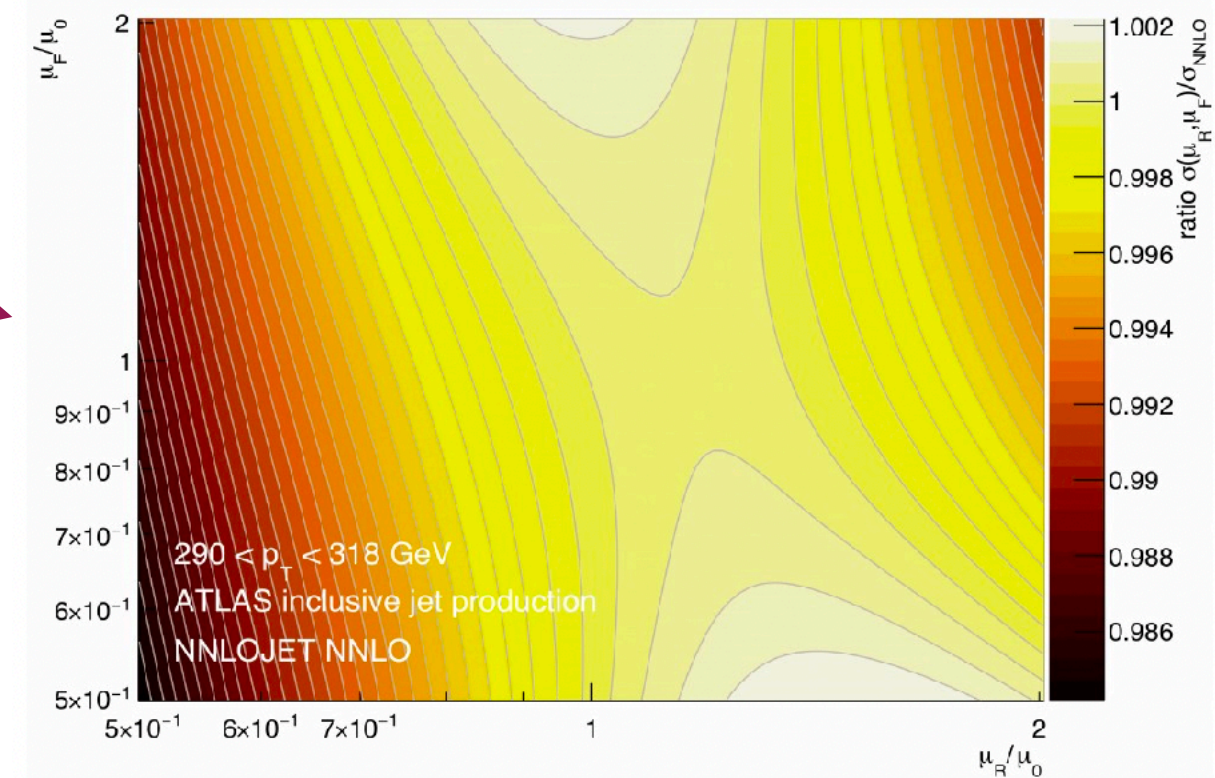
# THE INVESTMENT & RETURN

[APPLgrid, fastNLO, NNLOJET '22]



scale dependence  
 vs.  
 $\alpha_s \in [0.108, 0.124]$

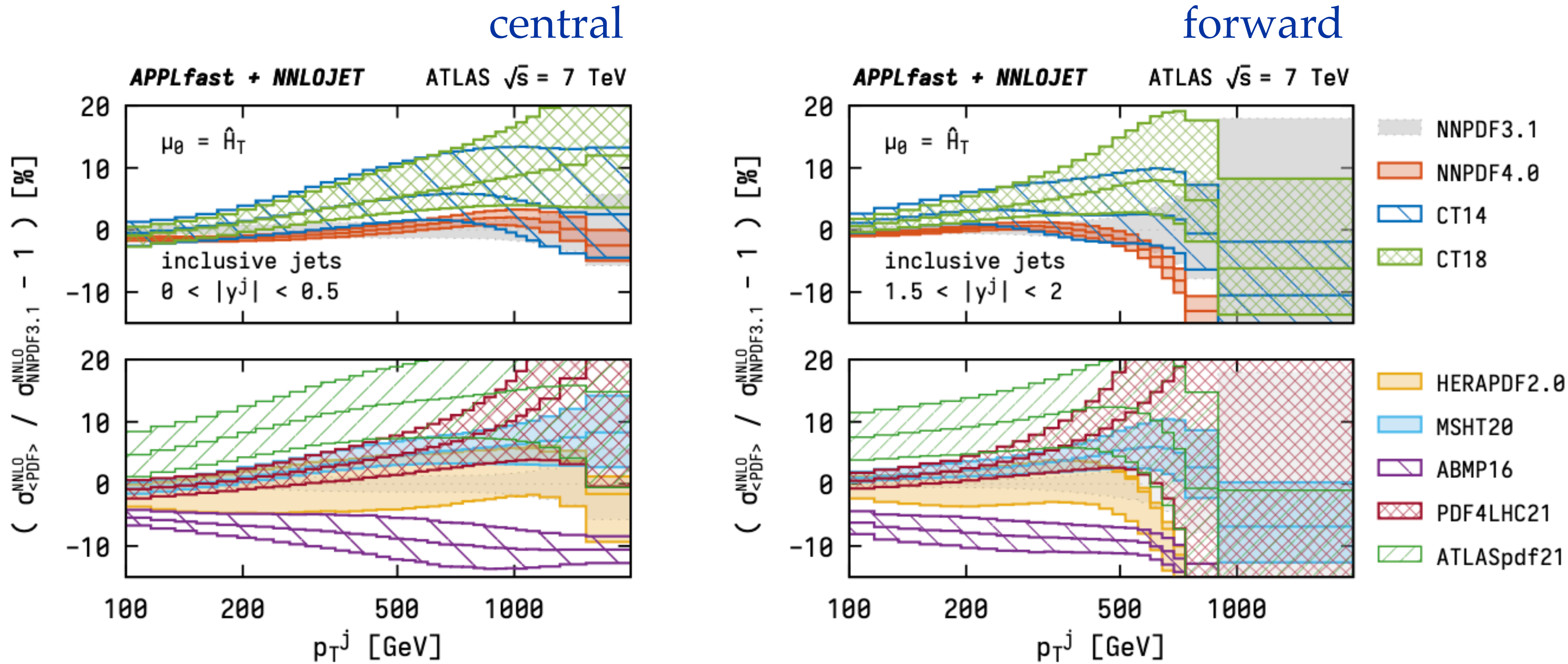
2D  $(\mu_R, \mu_F)$   
 dependence



predictions obtained  
 in seconds!

# PDF DEPENDENCE & UNCERTAINTIES

[APPLgrid, fastNLO, NNLOJET '22]



- ABMP16 & ATLASpdf21 largest excursion from the rest of the “pack”
- extremely small NNPDF4.0 PDF errors



# VALIDITY OF $K$ -FACTORS

[APPLgrid, fastNLO, NNLOJET '22]

$$K^{\text{NNLO}}(\mu) \equiv \frac{d\sigma^{\text{NNLO}}(\mu)/dp_T}{d\sigma^{\text{NLO}}(\mu)/dp_T}$$

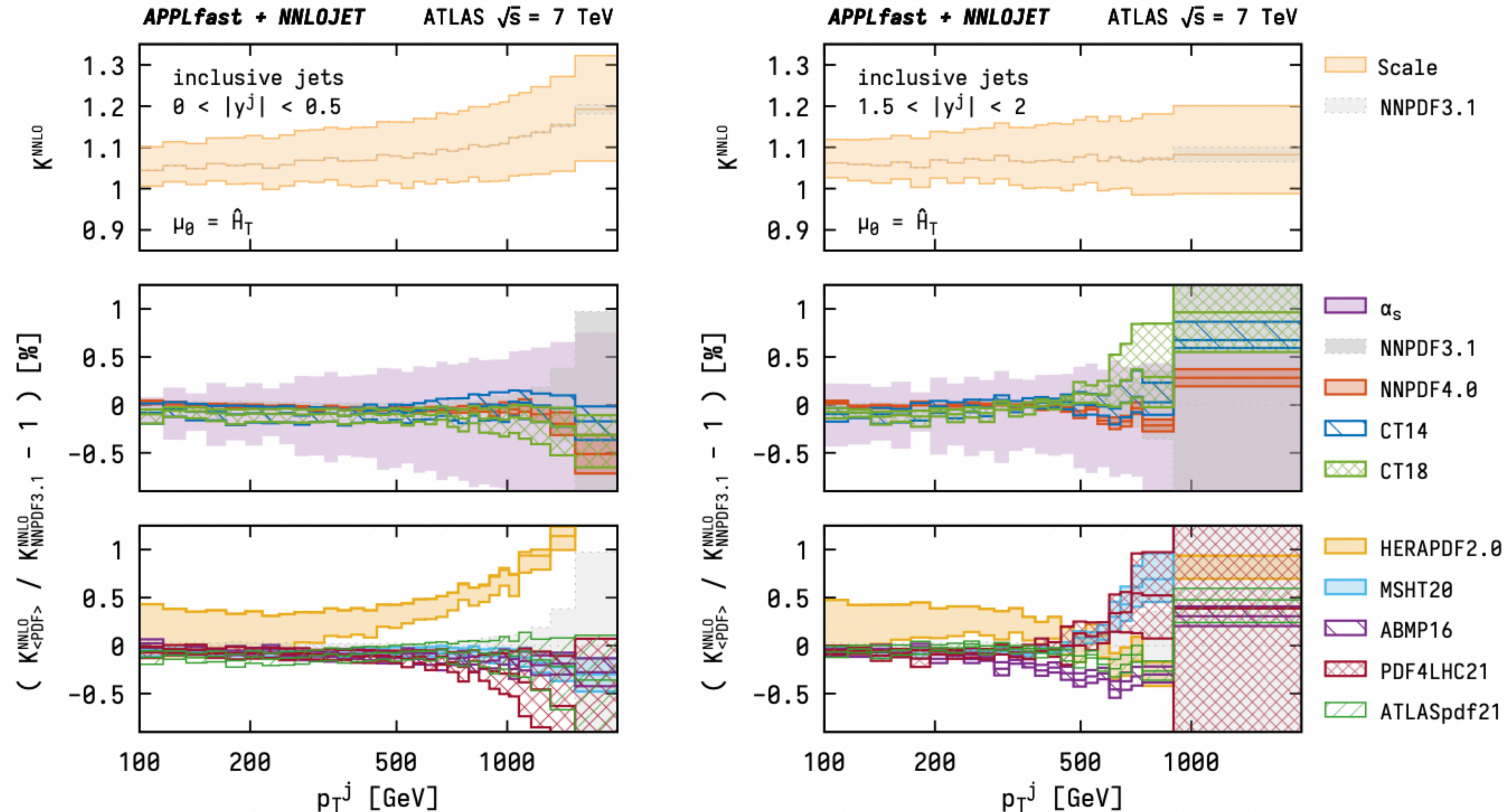
$$\sigma_{\text{approx. 1}}^{\text{NNLO}}(\mu) = \sigma^{\text{NLO}}(\mu) \times K^{\text{NNLO}}(\mu_{\text{ref}})$$

$$\sigma_{\text{approx. 2}}^{\text{NNLO}}(\mu) = \sigma^{\text{NLO}}(\mu) \times K^{\text{NNLO}}(\mu),$$

- $K$ -factor must be applied with correlated scales to avoid  $\mathcal{O}(10\%)$  scale unc.
- extremely robust ( $\lesssim 0.5\%$ ) w.r.t. PDF choice!  
(exception: HERAPDF2.0)

central

forward



# AVAILABLE GRIDS TABLES

[APPLgrid, fastNLO, NNLOJET '22]

## inclusive jets

Data	$\sqrt{s}$ [TeV]	$\mathcal{L}$ [fb <sup>-1</sup> ]	no. of points	anti- $k_T$ $R$	kinematic range [GeV]	fiducial cuts	$\mu_{R/F}$ -choice
CMS [30]	2.76	0.00543	81	0.7	$p_T^{\text{jet}} \in [74, 592]$	$ y  < 3.0$	$p_T^{\text{jet}}, \hat{H}_T$
ATLAS [28]	7.0	4.5	140	0.6	$p_T^{\text{jet}} \in [100, 1992]$	$ y  < 3.0$	$p_T^{\text{jet}}, \hat{H}_T$
CMS [31]	7.0	5.0	133	0.7	$p_T^{\text{jet}} \in [114, 2116]$	$ y  < 3.0$	$p_T^{\text{jet}}, \hat{H}_T$
ATLAS [32]	8.0	20.3	171	0.6	$p_T^{\text{jet}} \in [70, 2500]$	$ y  < 3.0$	$p_T^{\text{jet}}, \hat{H}_T$
CMS [33]	8.0	5.6 19.7	248	0.7	$p_T^{\text{jet}} \in [21, 74]$ $p_T^{\text{jet}} \in [74, 2500]$	$ y  < 4.7$	$p_T^{\text{jet}}, \hat{H}_T$
ATLAS [34]	13.0	3.2	177	0.4	$p_T^{\text{jet}} \in [100, 3937]$	$ y  < 3.0$	$p_T^{\text{jet}}, \hat{H}_T$
CMS [35]	13.0	36.3 33.5	2 × 78	0.4 0.7	$p_T^{\text{jet}} \in [97, 3103]$	$ y  < 2.0$	$p_T^{\text{jet}}, \hat{H}_T$

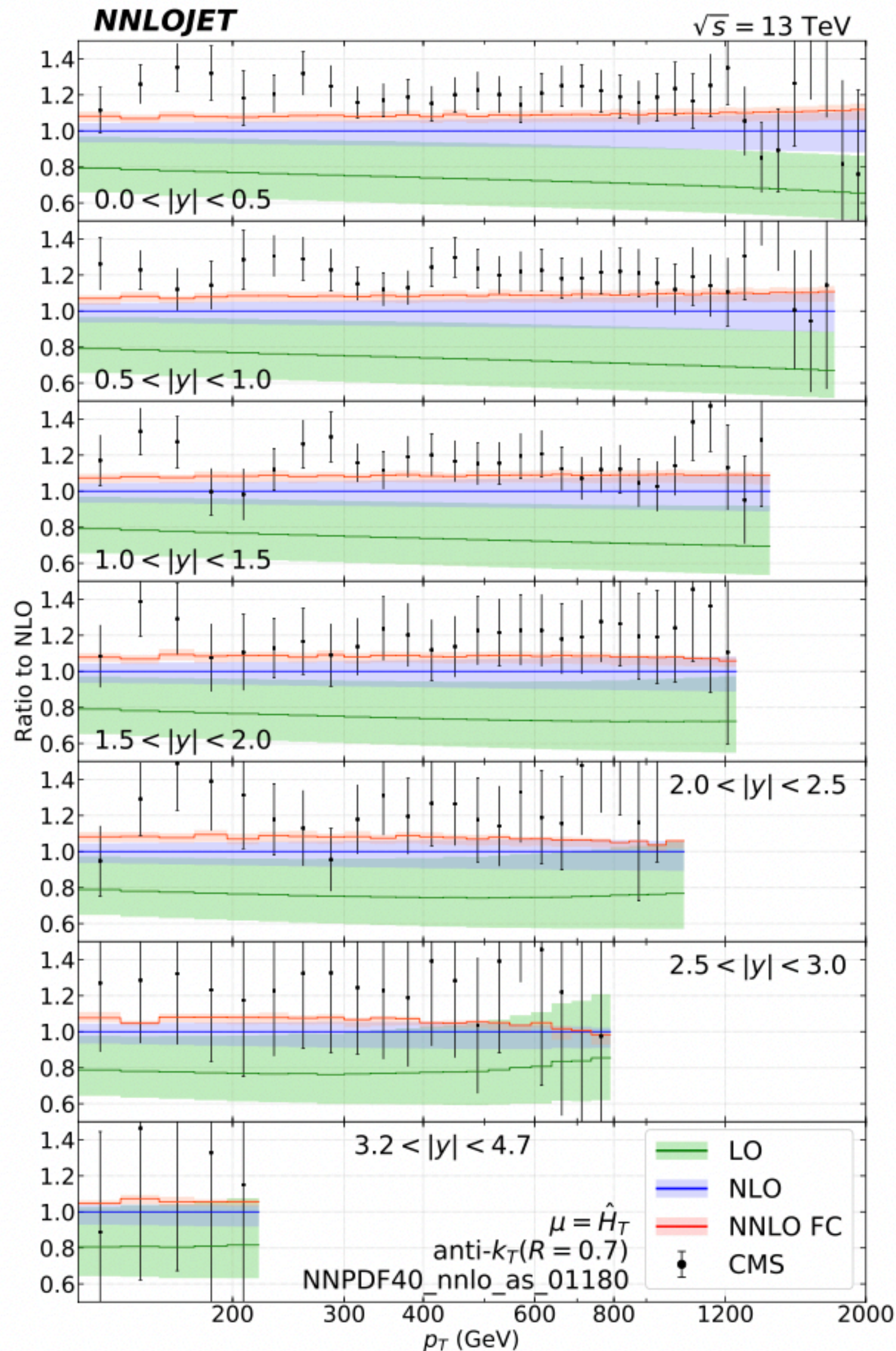
## di-jets

Data	$\sqrt{s}$ [TeV]	$\mathcal{L}$ [fb <sup>-1</sup> ]	no. of points	anti- $k_T$ $R$	kinematic range [GeV]	fiducial cuts	$\mu_{R/F}$ -choice
ATLAS [55]	7.0	4.5	90	0.6	$m_{12} \in [260, 5040]$	$ y_1 ,  y_2  < 3.0$ $[p_{T,1}, p_{T,2}] > [100, 50]\text{GeV}$ $y^* < 3.0$	$m_{12}$
CMS [31]	7.0	5.0	54	0.7	$m_{12} \in [197, 5058]$	$ y  < 5.0$ $[p_{T,1}, p_{T,2}] > [60, 30]\text{GeV}$ $ y_{\text{max}}  < 2.5$	$m_{12}$
CMS [49]	8.0	19.7	122	0.7	$\langle p_{T1,2} \rangle \in [133, 1784]$	$ y  < 5.0$ $p_{T,1}, p_{T,2} > 50\text{GeV}$ $ y_1 ,  y_2  < 3.0$	$p_{T,1} \exp(0.3 y^*)$ $m_{12}$
ATLAS [34]	13.0	3.2	136	0.4	$m_{12} \in [260, 9066]$	$ y_1 ,  y_2  < 3.0$ $p_{T,1}, p_{T,2} > 75\text{GeV}$ $\langle p_{T1,2} \rangle > 100\text{GeV}$ $y^* < 3.0$	$m_{12}$

- all grids available on: [ploughshare.web.cern.ch](http://ploughshare.web.cern.ch)
- caveat*: calculation based on **leading-colour** approximation in NNLO parts  
 $\Leftrightarrow$  leading:  $N_c^2, N_c n_f, n_f^2$  (sub-leading:  $\times 1/N_c^2$ )

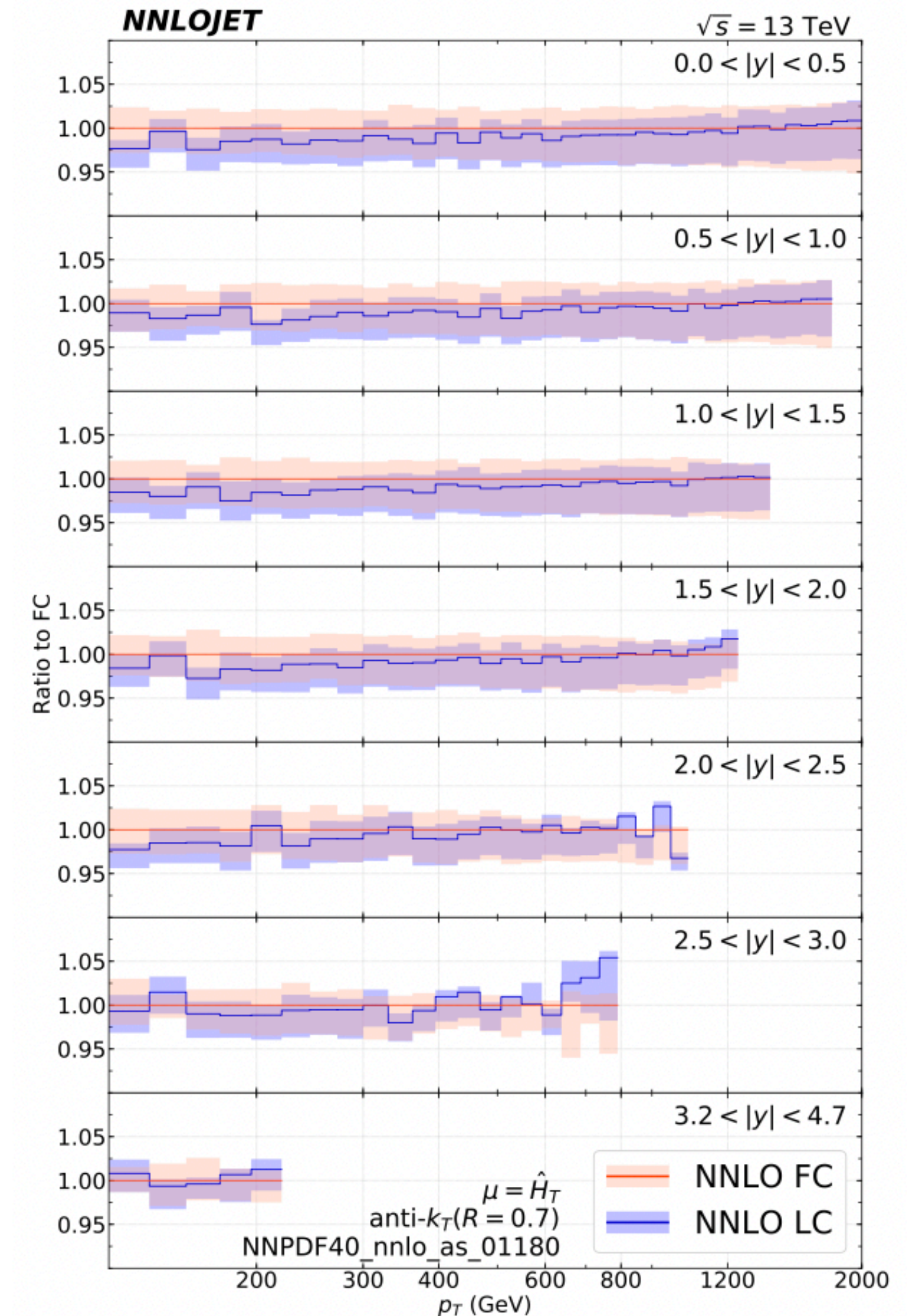
# How Good is LC?

sub-leading colour: SLC  
 leading colour: LC  
 full colour (LC+SLC): FC



- improved agreement with the data
- +ve SLC contribution
  - ↪ up to 20% on  $\delta\sigma^{\text{NNLO}}$
  - ↪ largest @ low- $p_T$
  - ↪ diminishes @ high- $p_T$
- impact on NNLO:
  - ↪  $\lesssim 2\%$
  - ↪ within  $\Delta_{\text{scl}}$

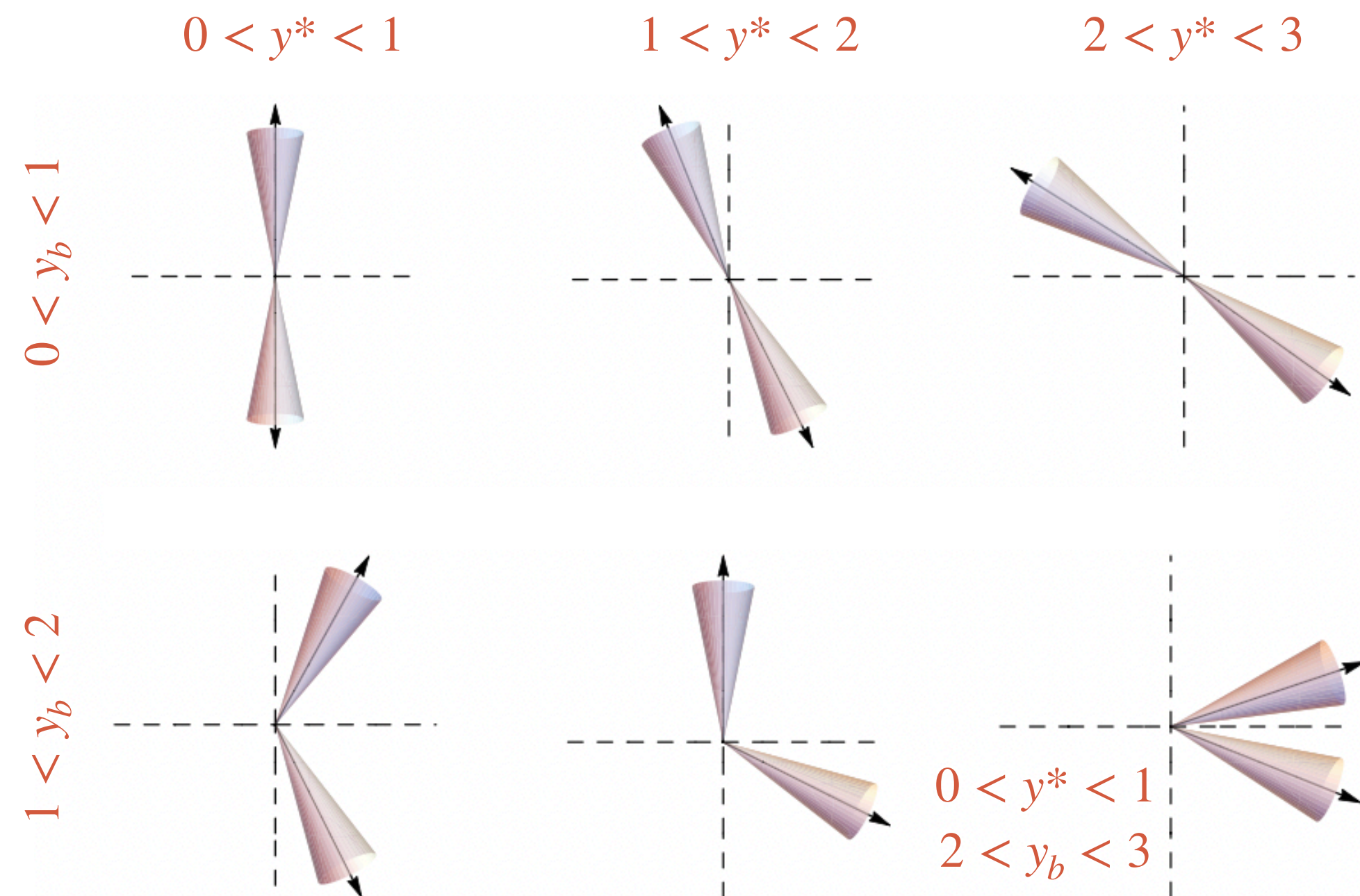
SLC small in incl. jets ( $R=0.4, 0.7$ )  
 still small on di-jet  $d\sigma/dm_{jj}$  ( $R=0.4$ )  
**substantial in 3D di-jet ( $R=0.7$ )**



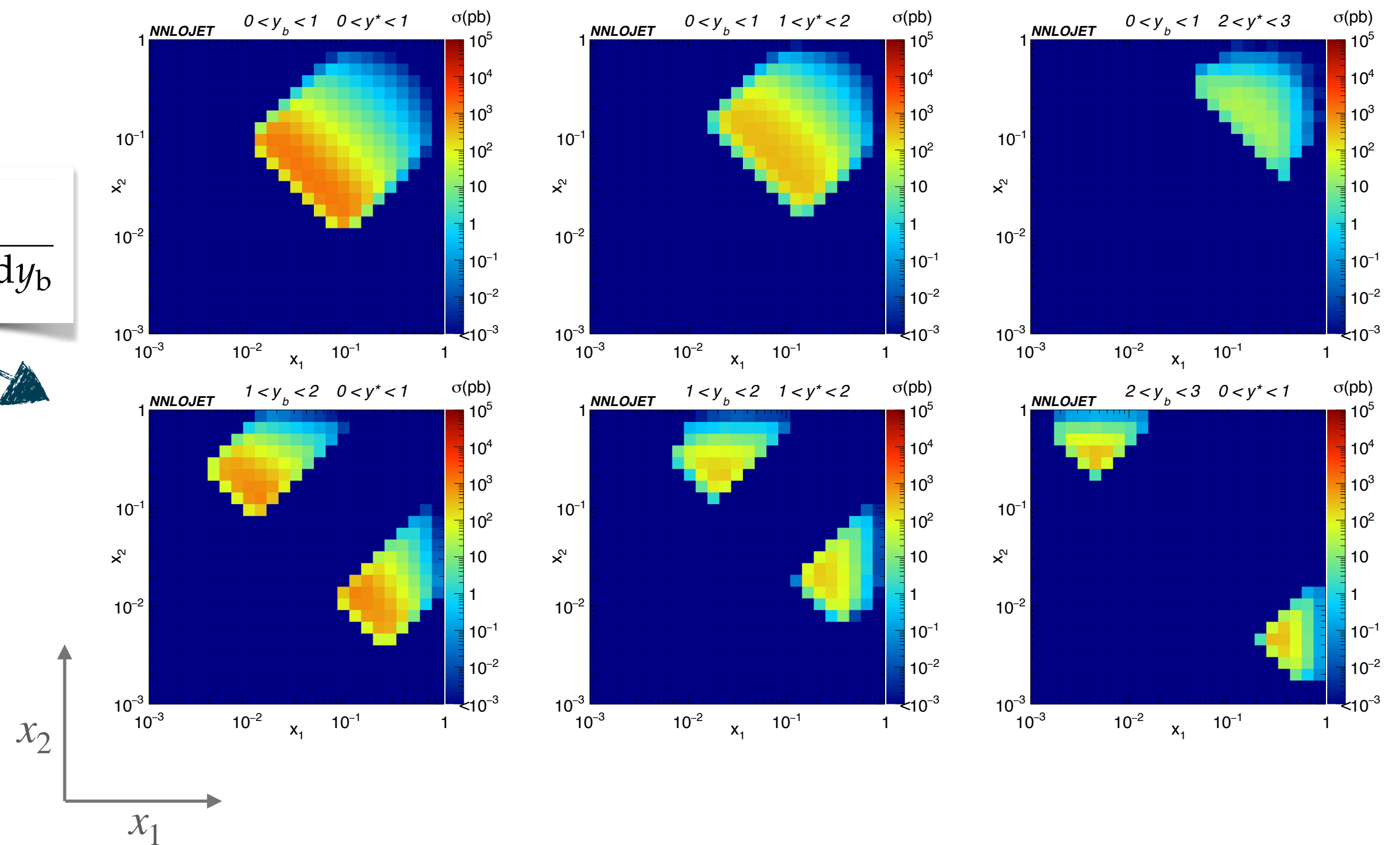
# TRIPLY-DIFFERENTIAL DI-JET PRODUCTION

- different event topologies  
 $\rightsquigarrow$  disentangle mom. fractions  $(x_1, x_2)$

$$x_{1,2} = \frac{2p_{T,avg}}{\sqrt{s}} e^{\pm y_b} \cosh(y^*)$$

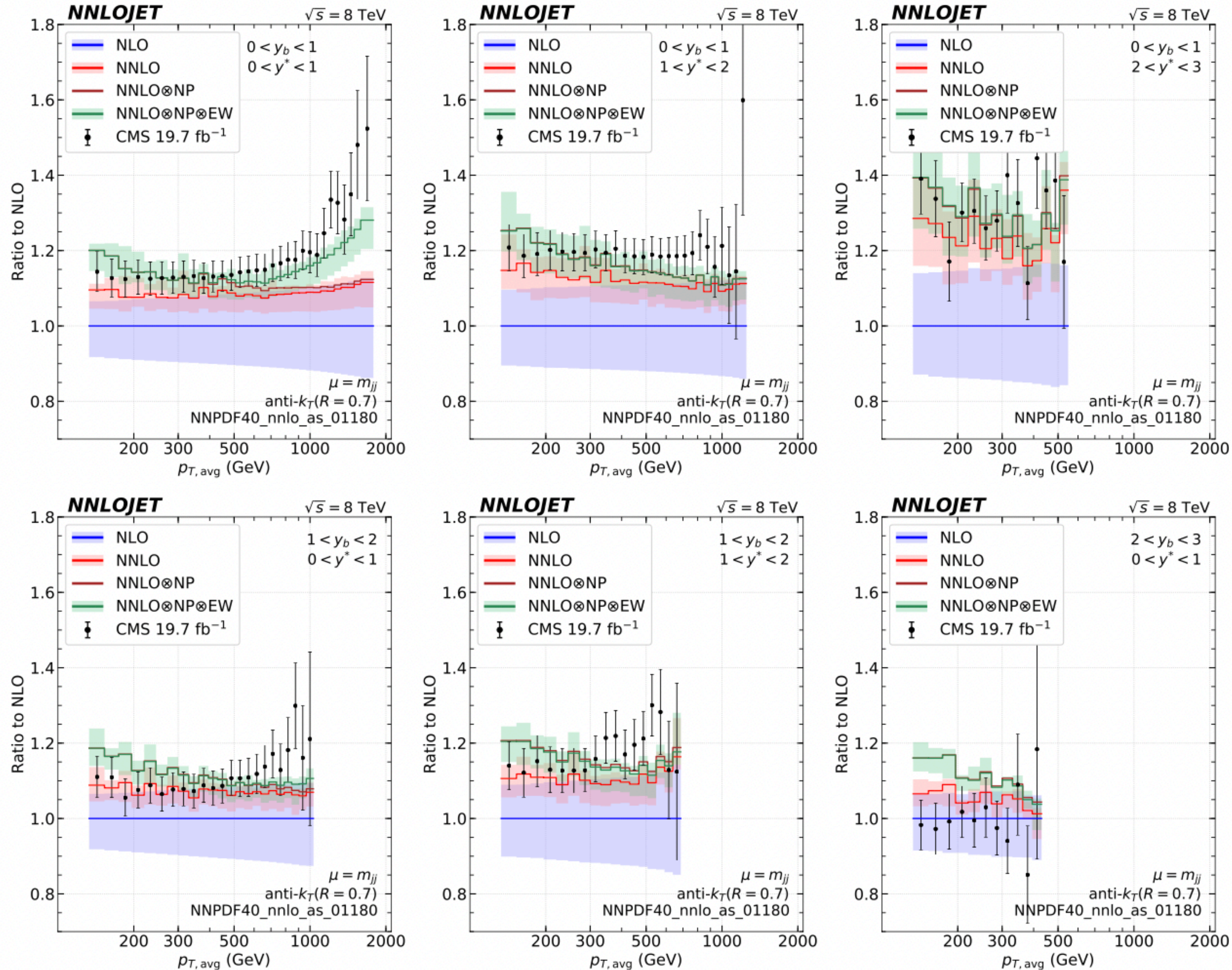


$$\frac{d^3\sigma}{dp_{T,avg} dy^* dy_b}$$



# TRIPLY-DIFFERENTIAL DI-JET PRODUCTION — TH vs. DATA

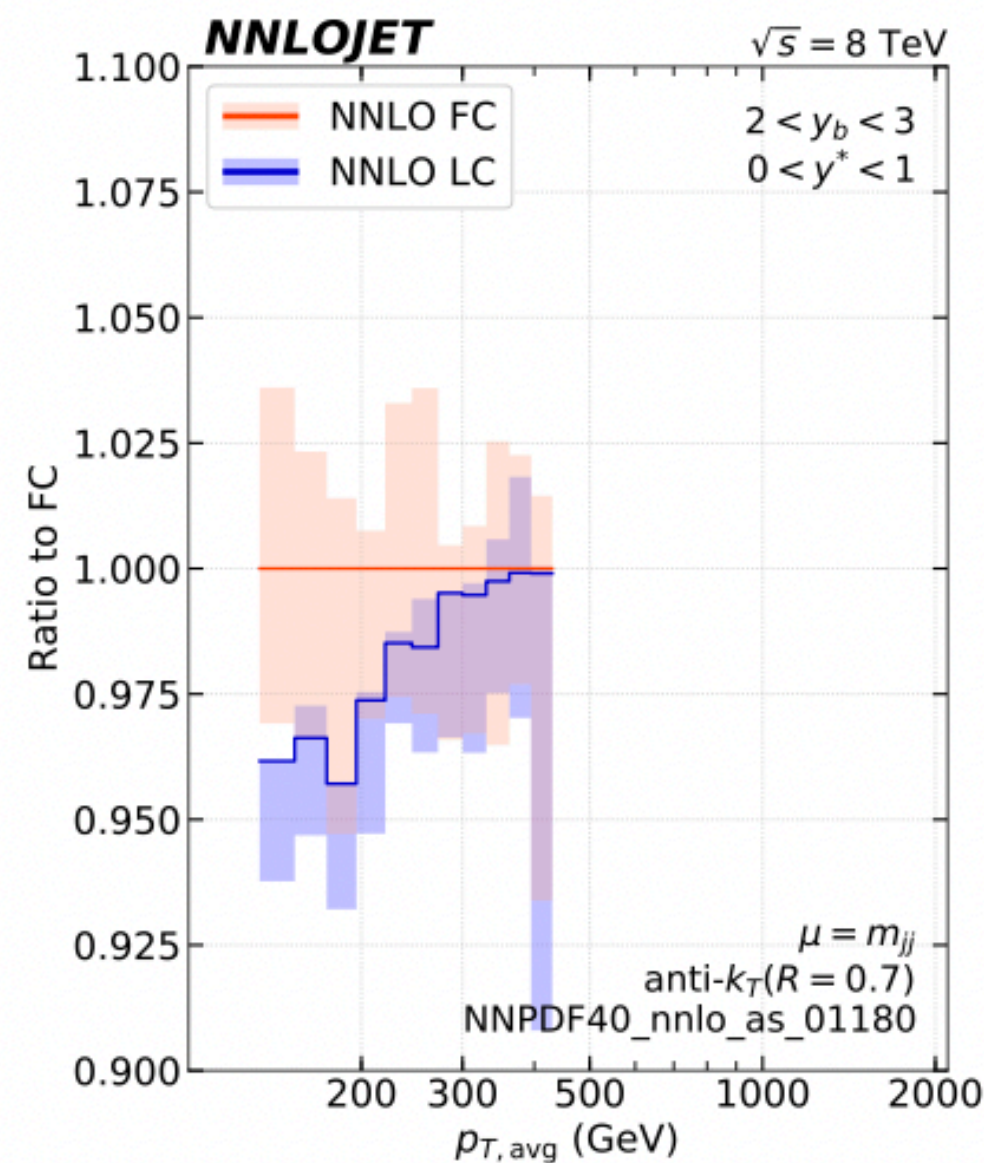
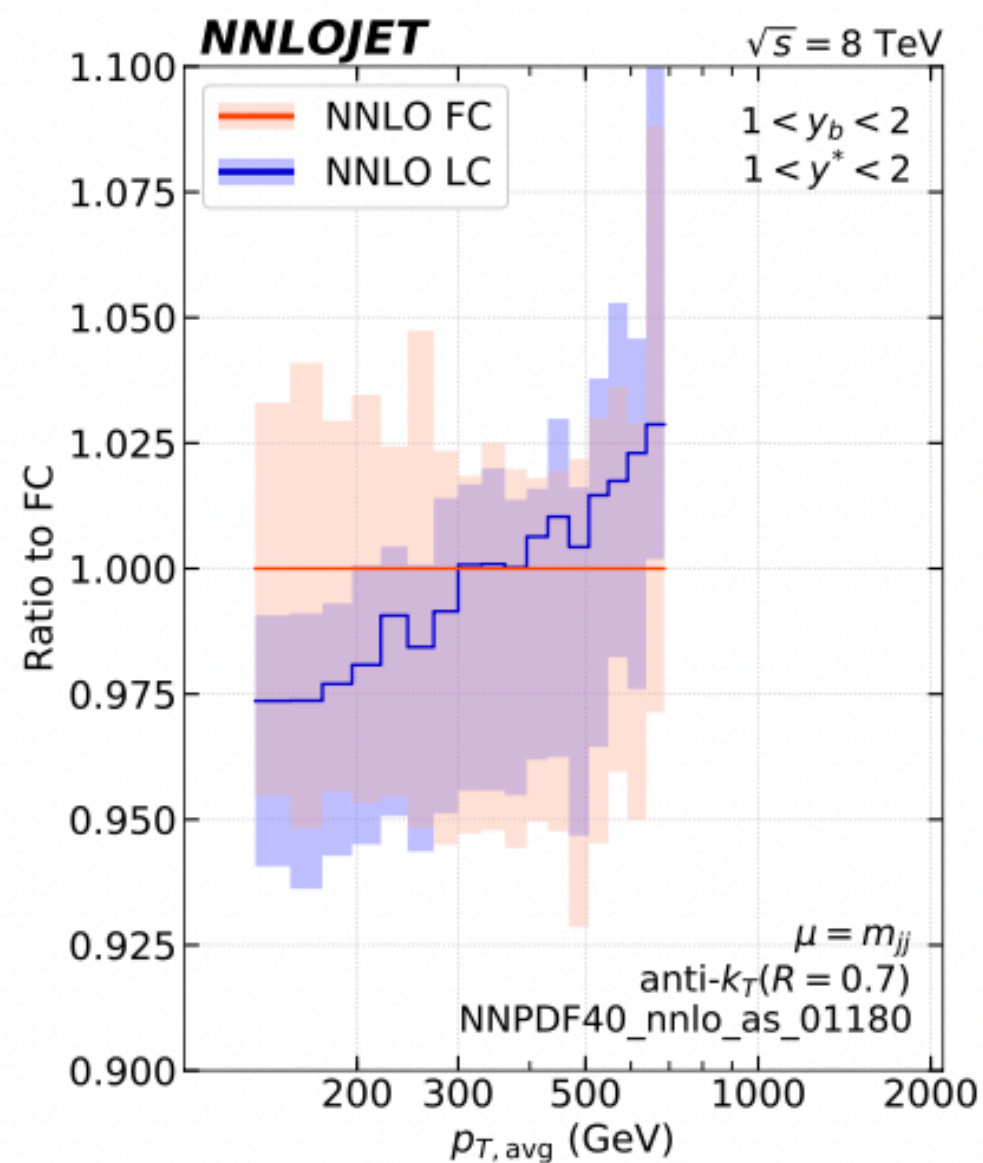
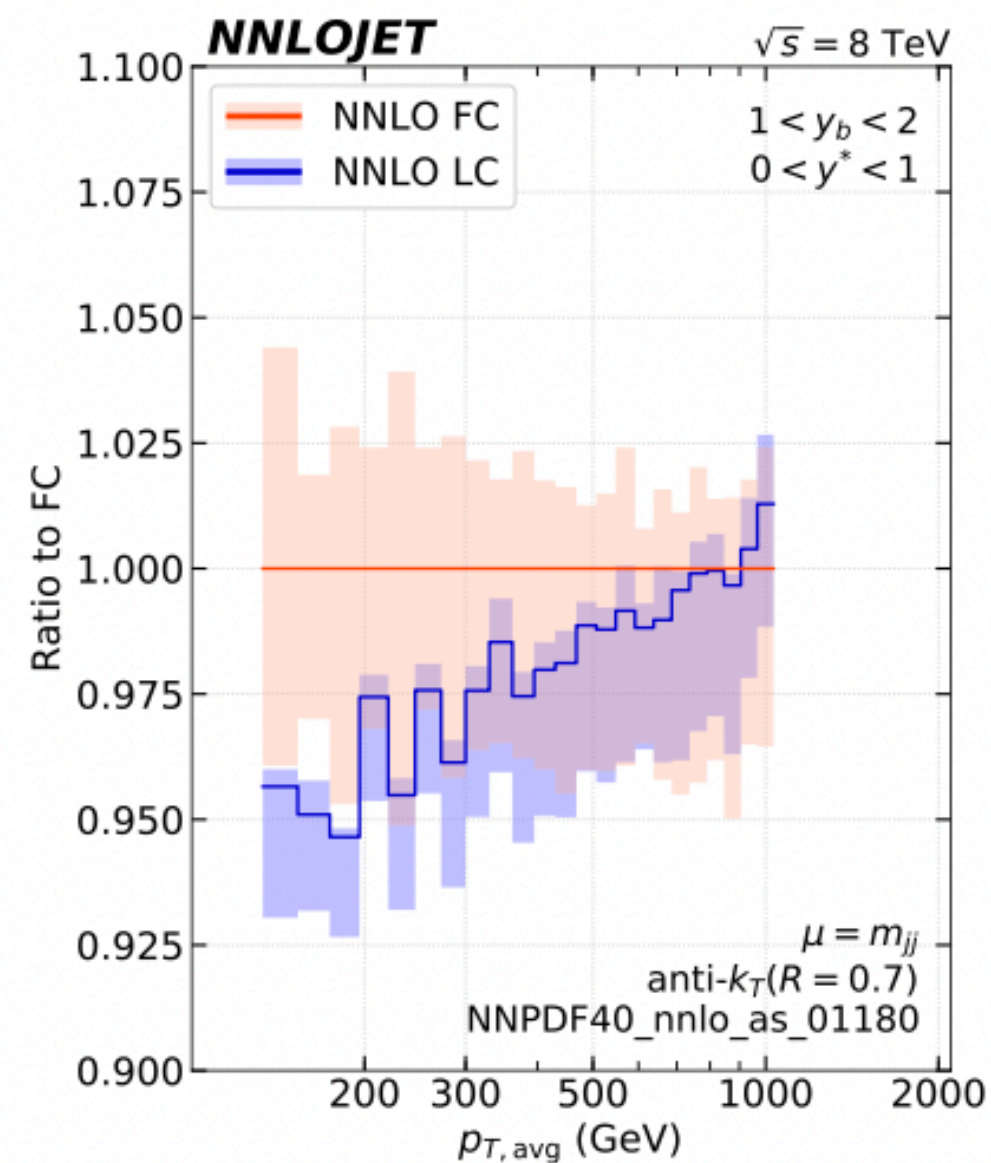
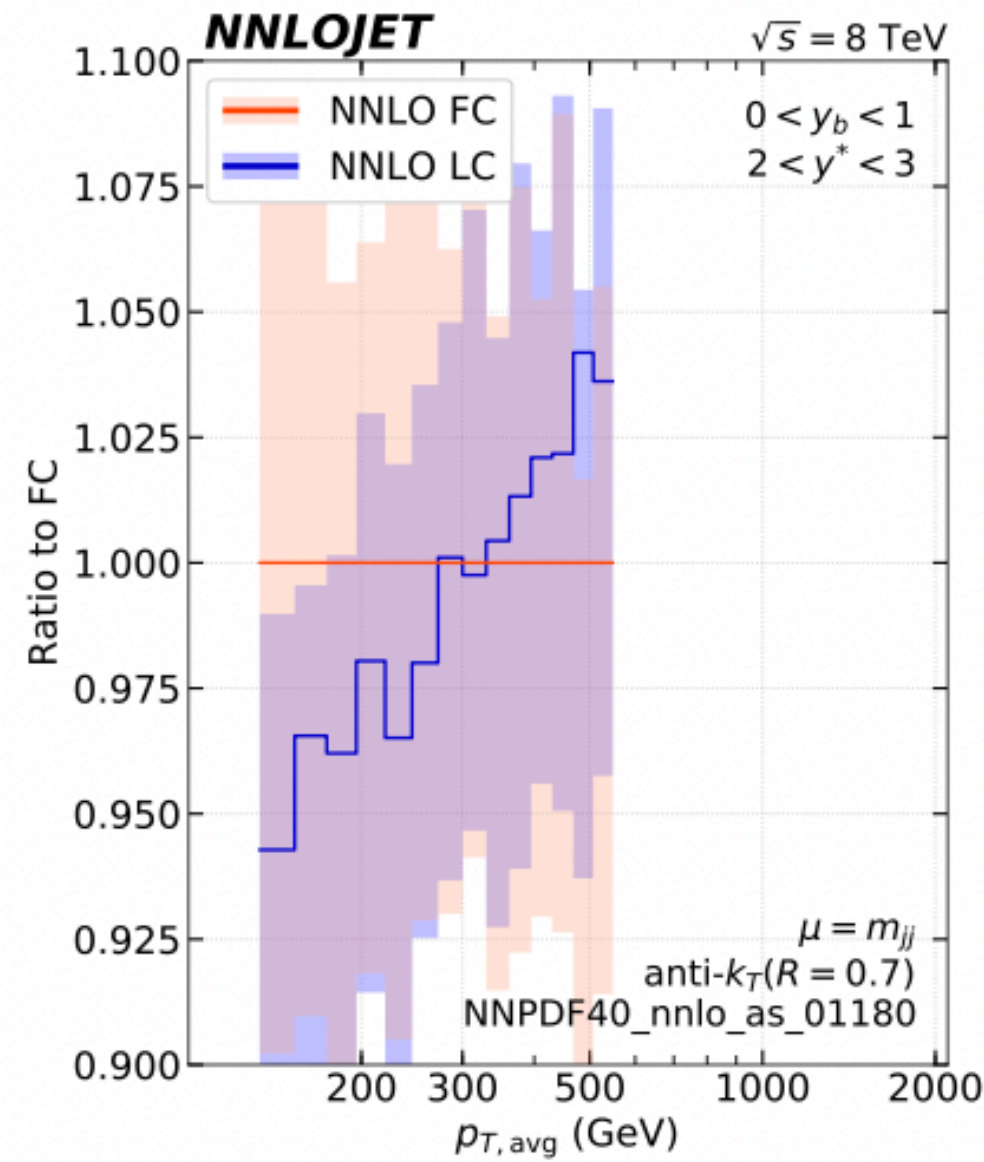
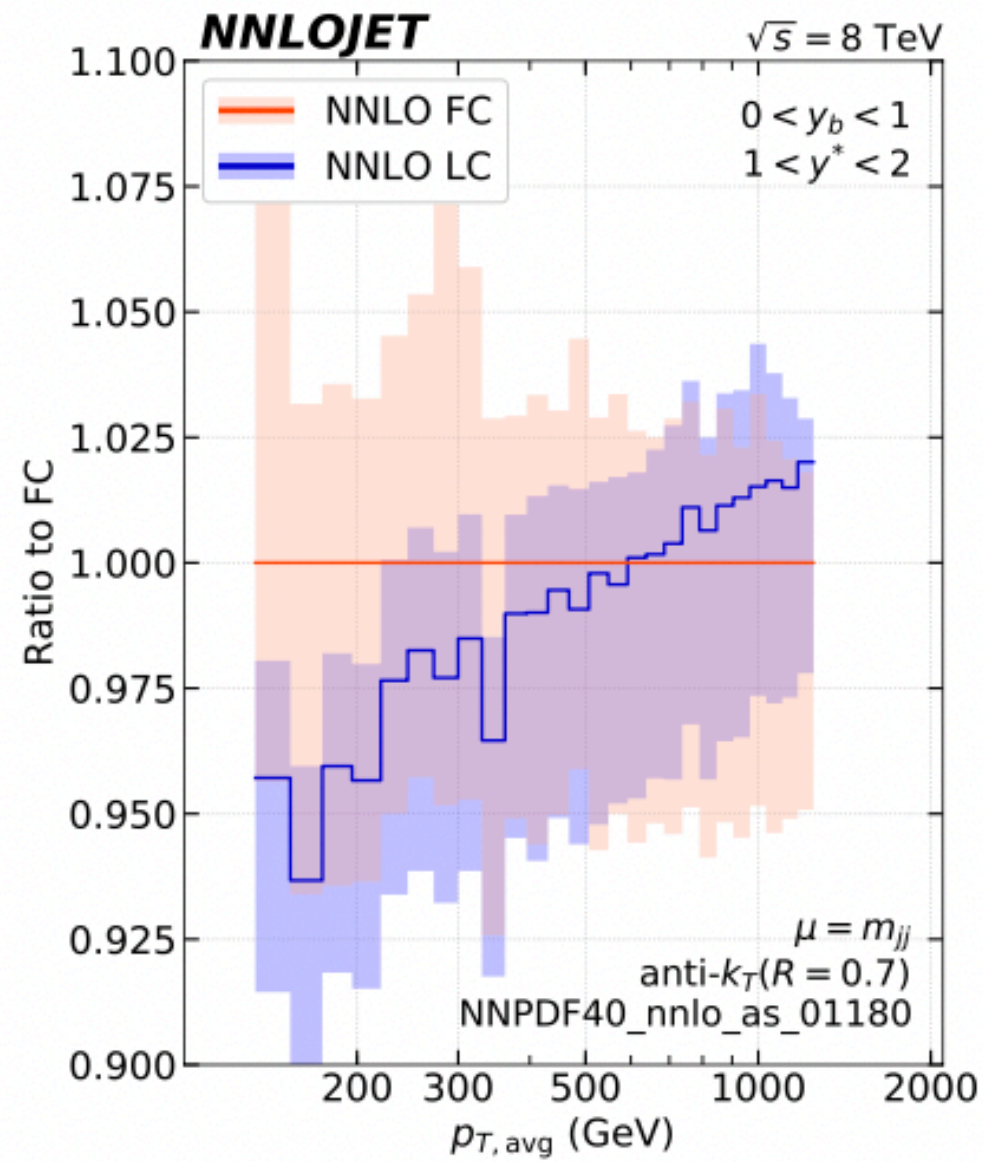
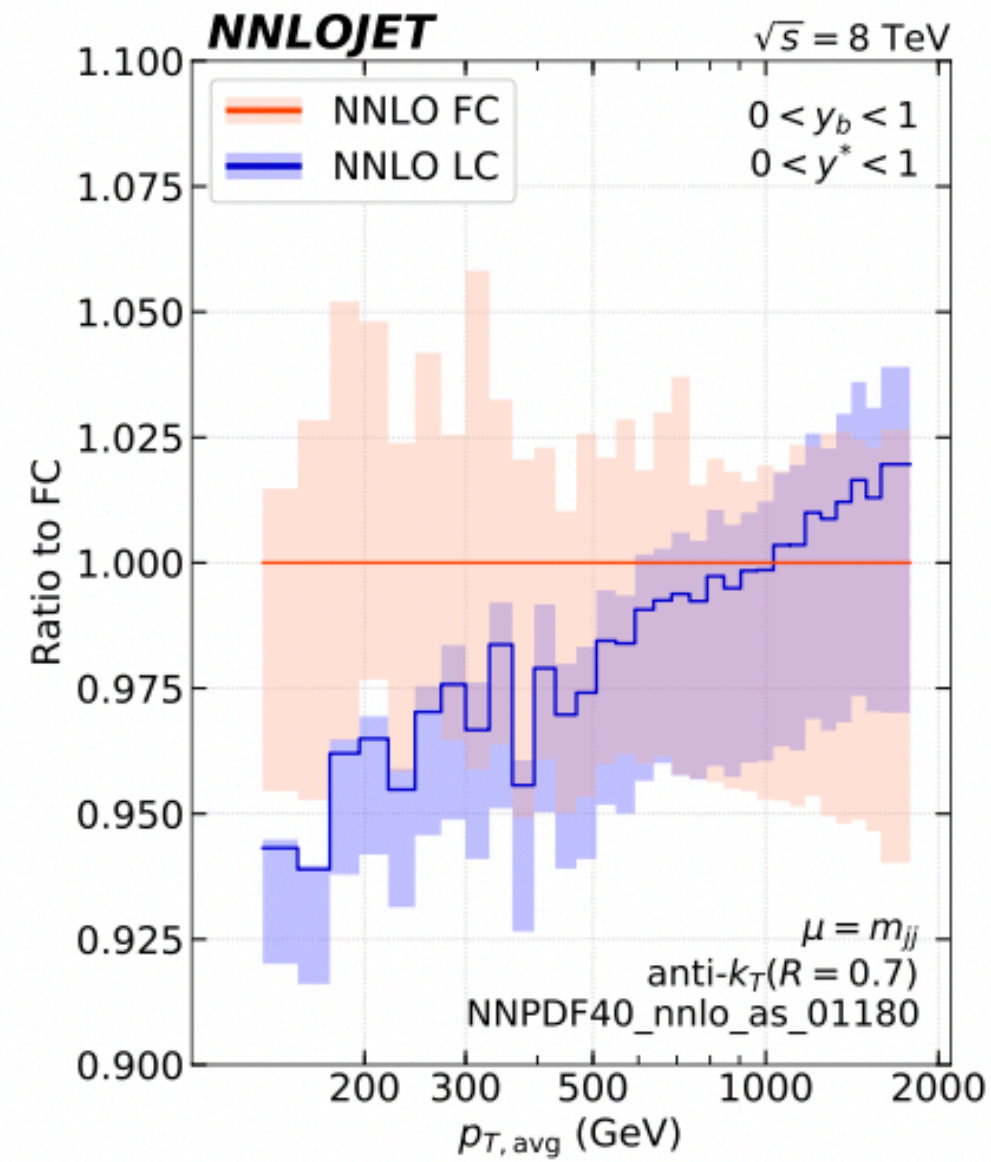
[Chen, Gehrmann, Glover, AH, Mo '22]



- large **NP** corrections  
@ low- $p_{T,avg}$
- **EW** corrections only impacts  
↔ high- $p_{T,avg}$   
&  $y_b, y^* < 1$
- improved description of data  
& reduced uncertainties

# TRIPLY-DIFFERENTIAL DI-JET PRODUCTION — FC vs. LC

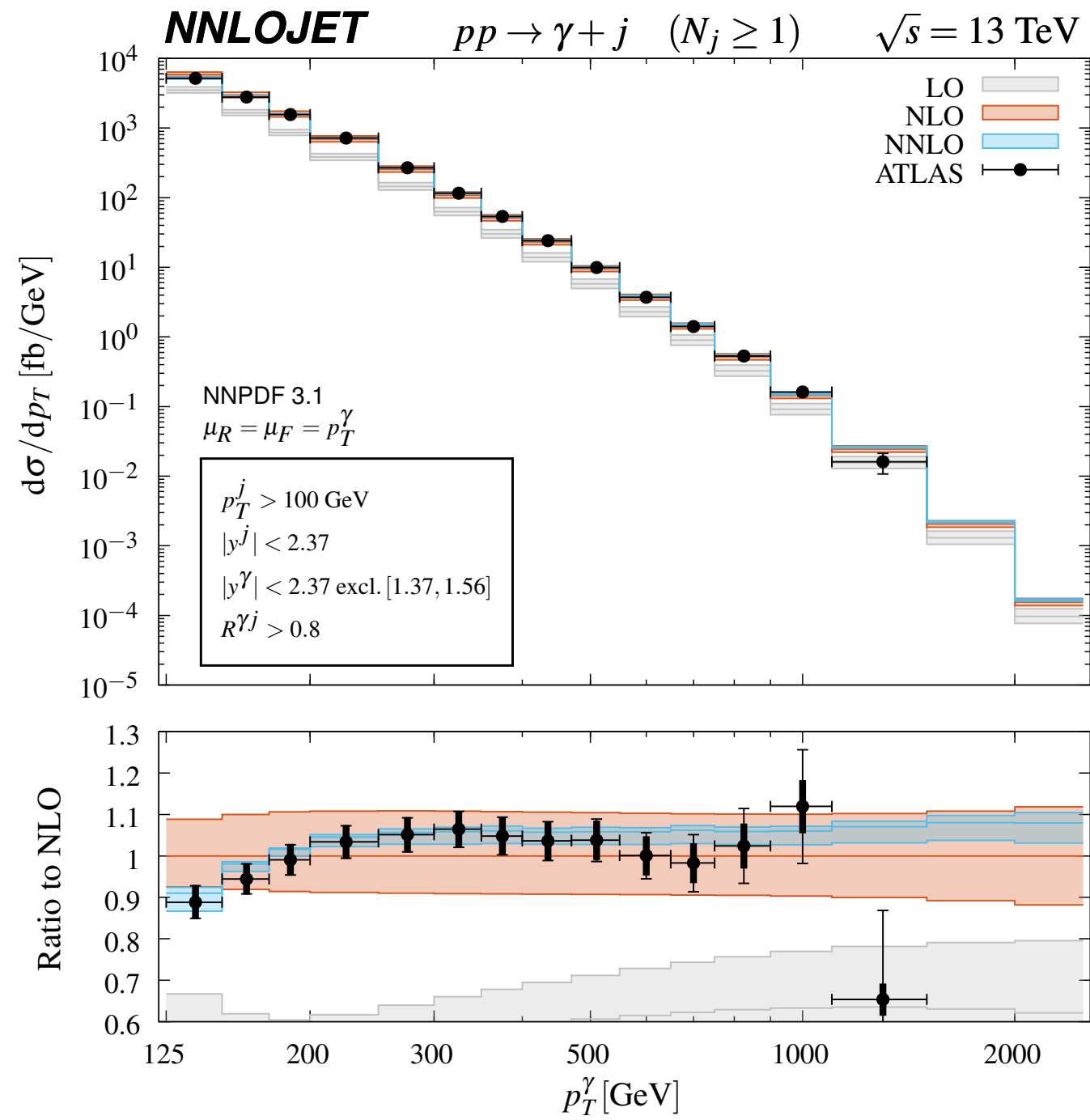
[Chen, Gehrmann, Glover, AH, Mo '22]



- large SLC contributions
  - ↪ low- $p_{T,avg}$   $\leftrightarrow$  30–60%
  - ↪ med- $p_{T,avg}$   $\leftrightarrow$  small | · |
  - ↪ high- $p_{T,avg}$   $\leftrightarrow$  –20 %
- **LC** → **FC**
  - ↪ +5 % enhancement
- grids with FC very desirable!
  - ↪ resolve tension with other datasets? [NNPDF4.0]

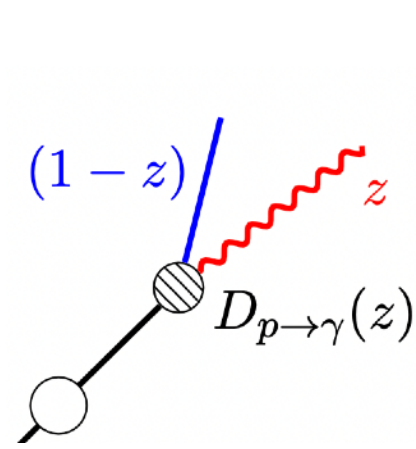
# IDENTIFIED OBJECTS — CHALLENGES IN TH VS. EXP

## ISOLATED PHOTONS $\gamma + \text{jet}$

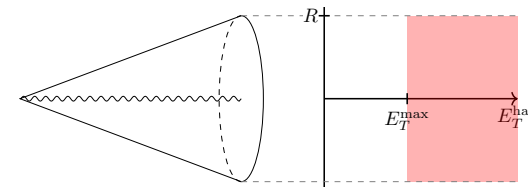


[Chen, Gehrmann, Glover, Höfer, AH '19]

“ok” for tight isol.  
tune @ NLO?  
per-cent mismatch

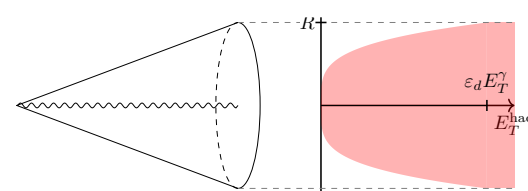


Exp:



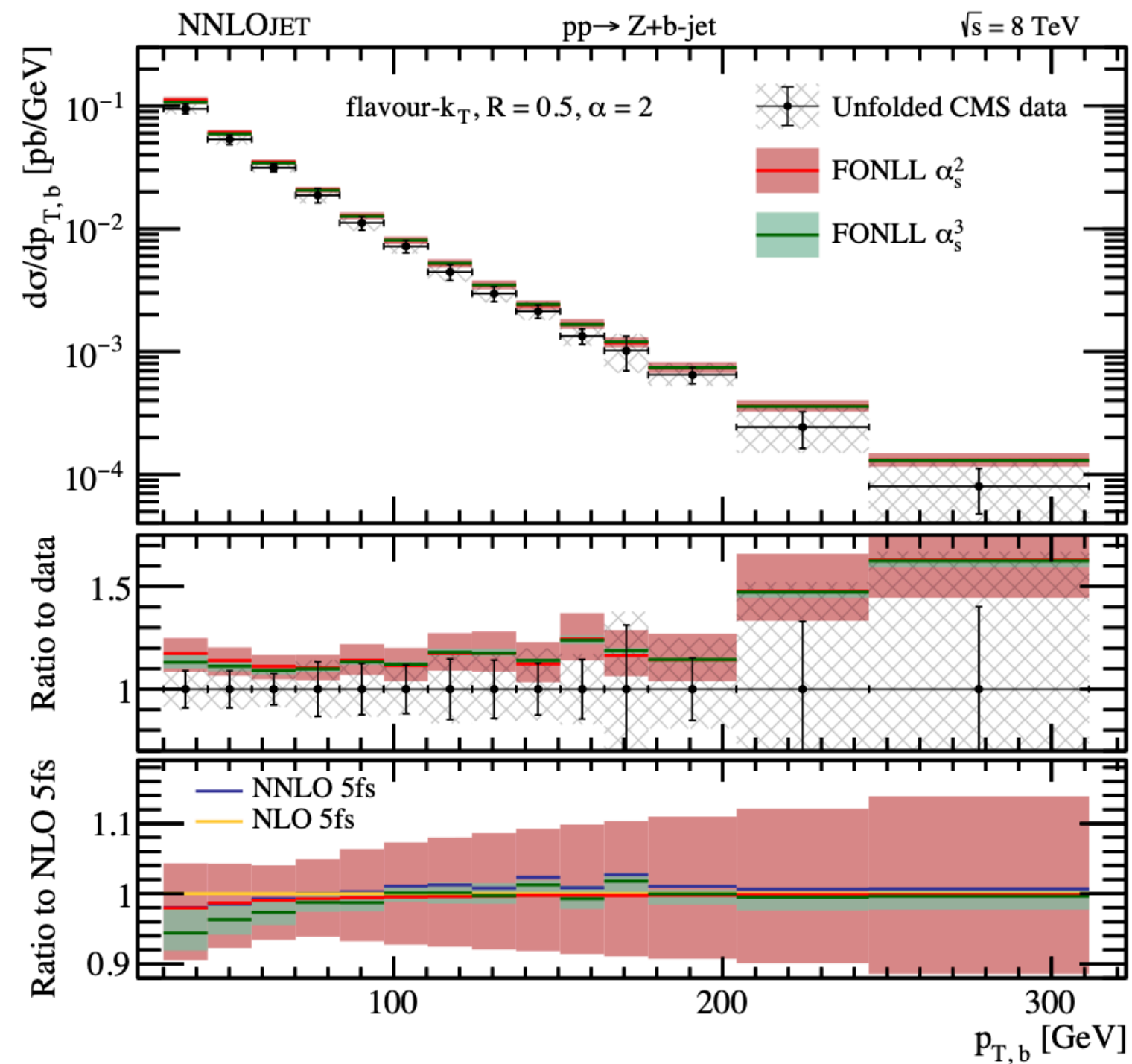
fixed cone

TH:

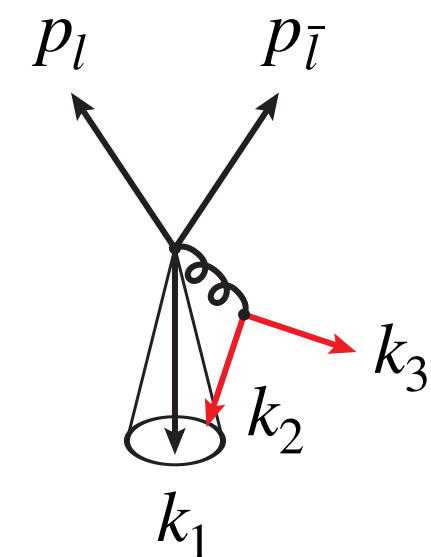


smooth cone

## JET FLAVOUR $Z + \text{b-jet}$



[Gauld, Gehrmann-De Ridder, Glover, AH, Majer '20]



Exp: anti- $k_T$

TH: flavour- $k_T$

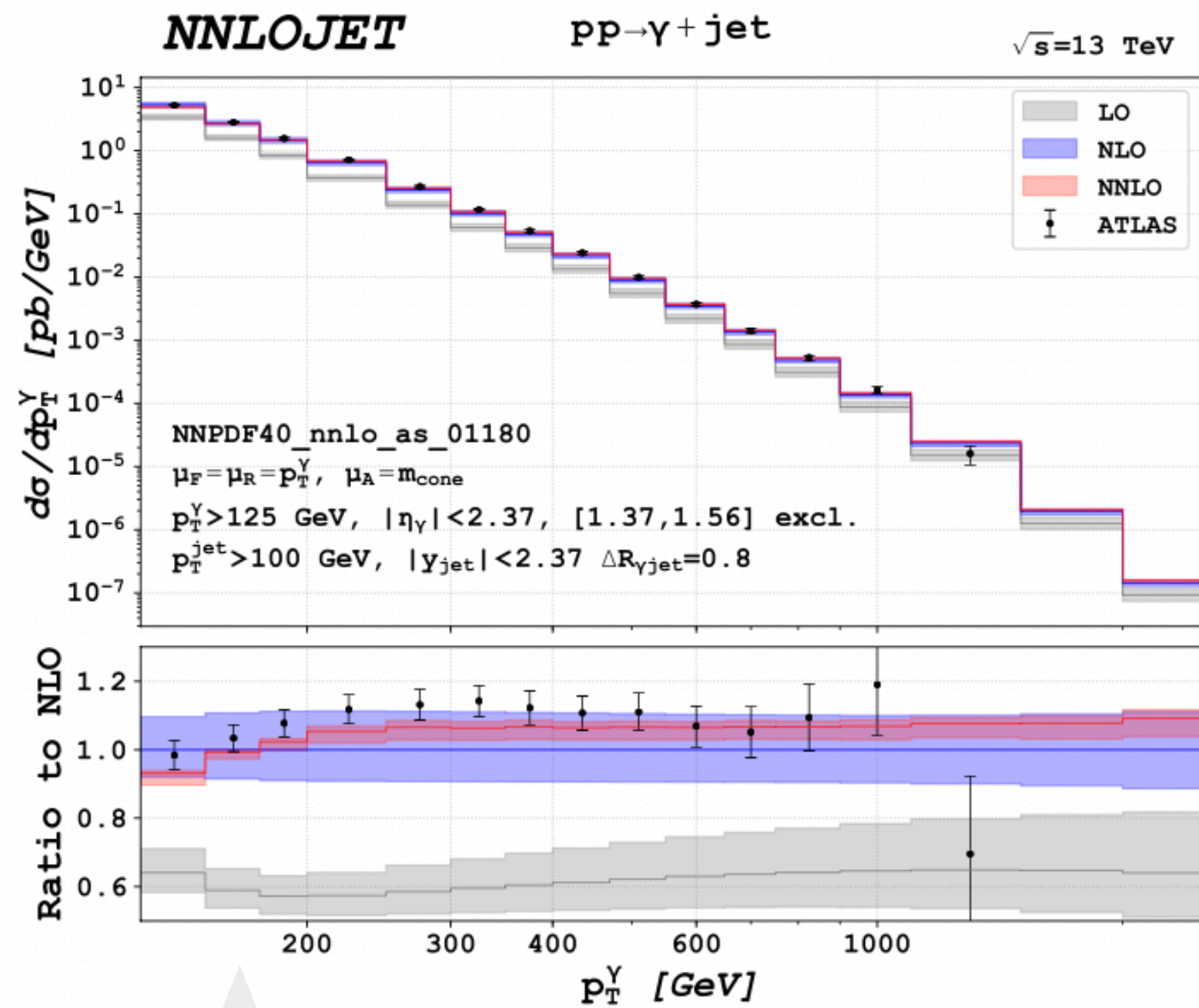
## HADRONS

- B-hadron Stripper [Czakon, Generet, Mitov, Poncelet '21]
- antenna final-final [Gehrmann, Stagnitto '22]

...

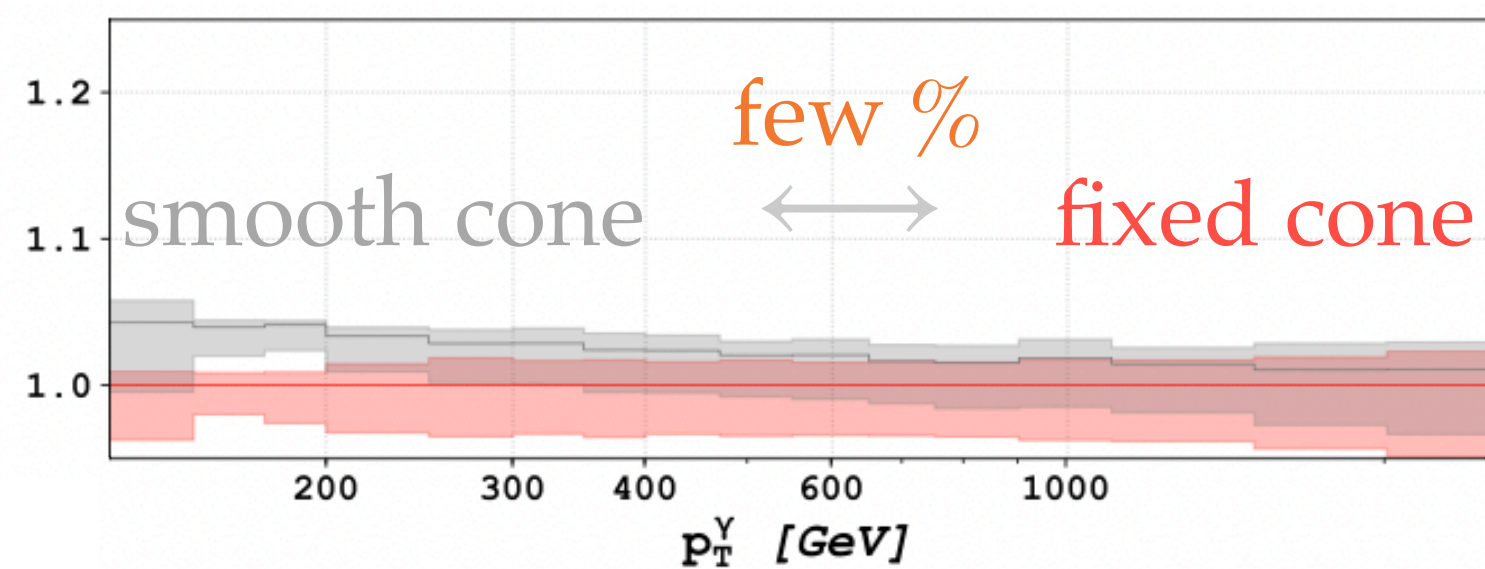
unfolding  
 $\mathcal{O}(10\%)$

# $\gamma + \text{jet}$ @ NNLO WITH FRAGMENTATION



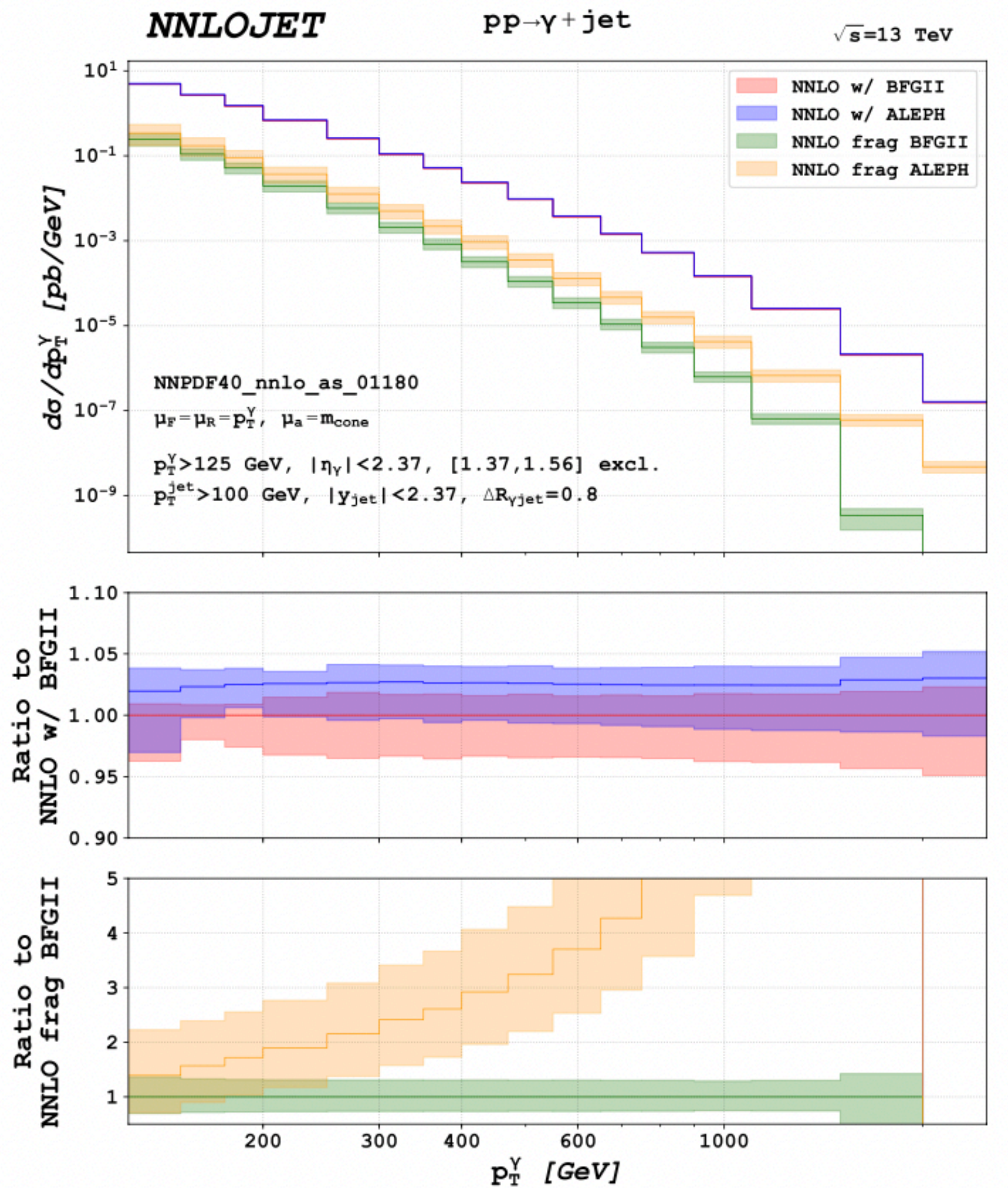
slight deterioration of agreement

maybe ok at NLO; **matters** at NNLO



## DEPENDENCE ON $D_{a \rightarrow \gamma}$

- BFG II vs. ALEPH
  - [Bourhis, Fontannaz, Guillet '98]
  - [ALEPH collab. '96]
- differences on  $d\sigma/dp_T^Y$ 
  - $\rightsquigarrow$  2 – 4%
- frag. contrib.  $\times 10^{-1}$ 
  - $\rightsquigarrow$   $\mathcal{O}(1)$  differences
- access to  $D_{a \rightarrow \gamma}$  @ LHC
  - $\hookrightarrow$  new observables?
  - $\hookrightarrow$  NNFRag?
  - $\hookrightarrow$  ...





# CONCLUSIONS & OUTLOOK PART 1

---

- NNLO QCD calculations in good shape
  - $2 \rightarrow 2$  essentially solved
  - $2 \rightarrow 3$  new frontier  $\leftrightarrow$  methods reaching maturity
  - *loop amplitudes* becoming a bottleneck again
  - in the quest for percent-level theory  $\leftrightarrow$  mixed QCD $\times$ EW important
- dissemination of results
  - public codes (MCFM, Matrix), nTuples, ...
  - fast interpolation grids  $\leftrightarrow$  APPLgrid fastNLO PineAPPL (anyway needed in fitting)
- identified objects  $\leftrightarrow$  mismatch in TH vs. Exp/NNLO
  - photon isolation, flavour tagging, hadron fragmentation, ...

# THE PLAN.

THE PLAN.

## 1. NNLO predictions for the LHC

- jets & interpolations grids
- identified photons & fragmentation

## 2. Differential N<sup>3</sup>LO

- Higgs & fiducial power corrections
- Drell-Yan & PDFs

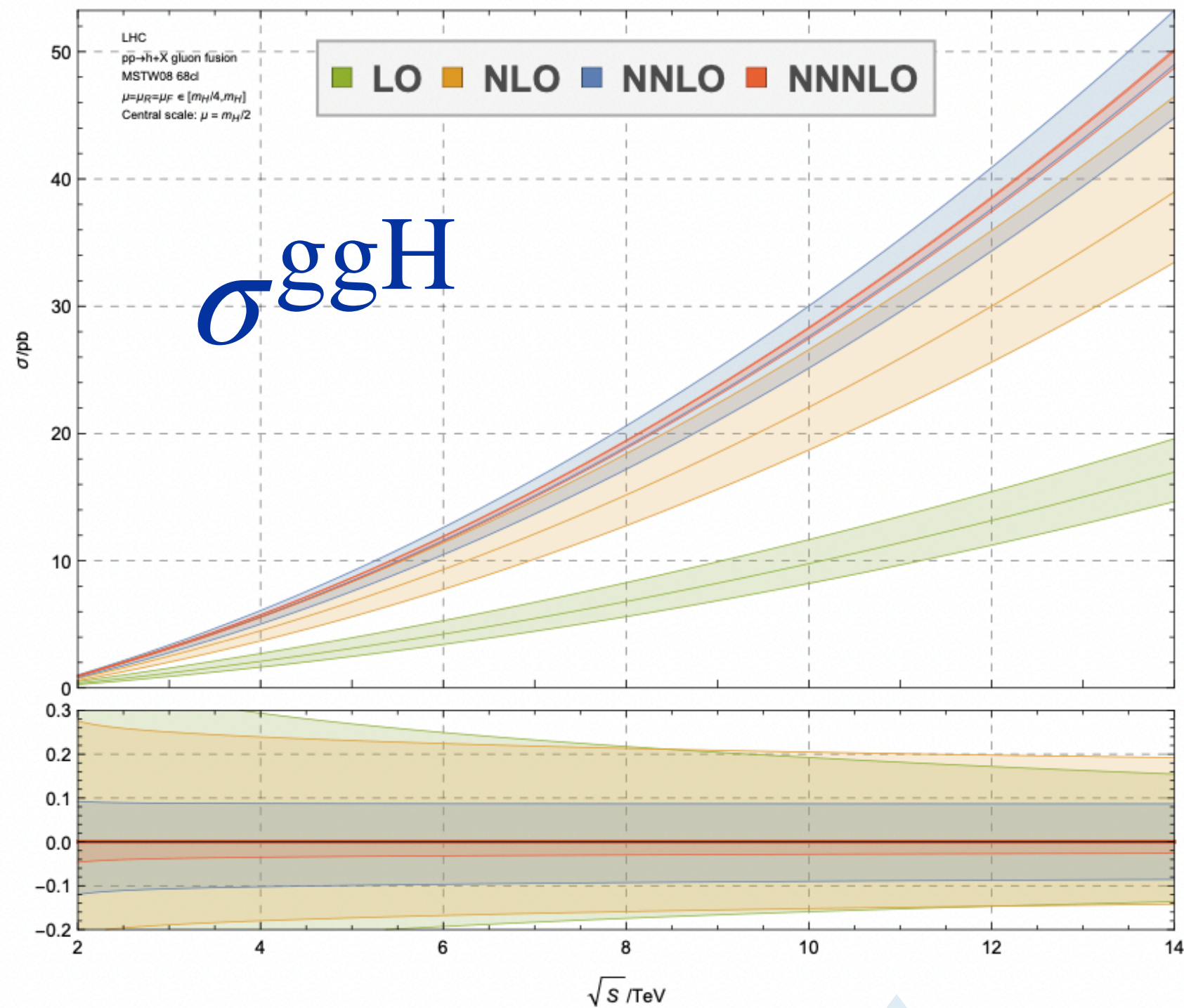
## 3. Bayesian approach to MHO

- the abc model & correlations

## 4. Summary & Outlook

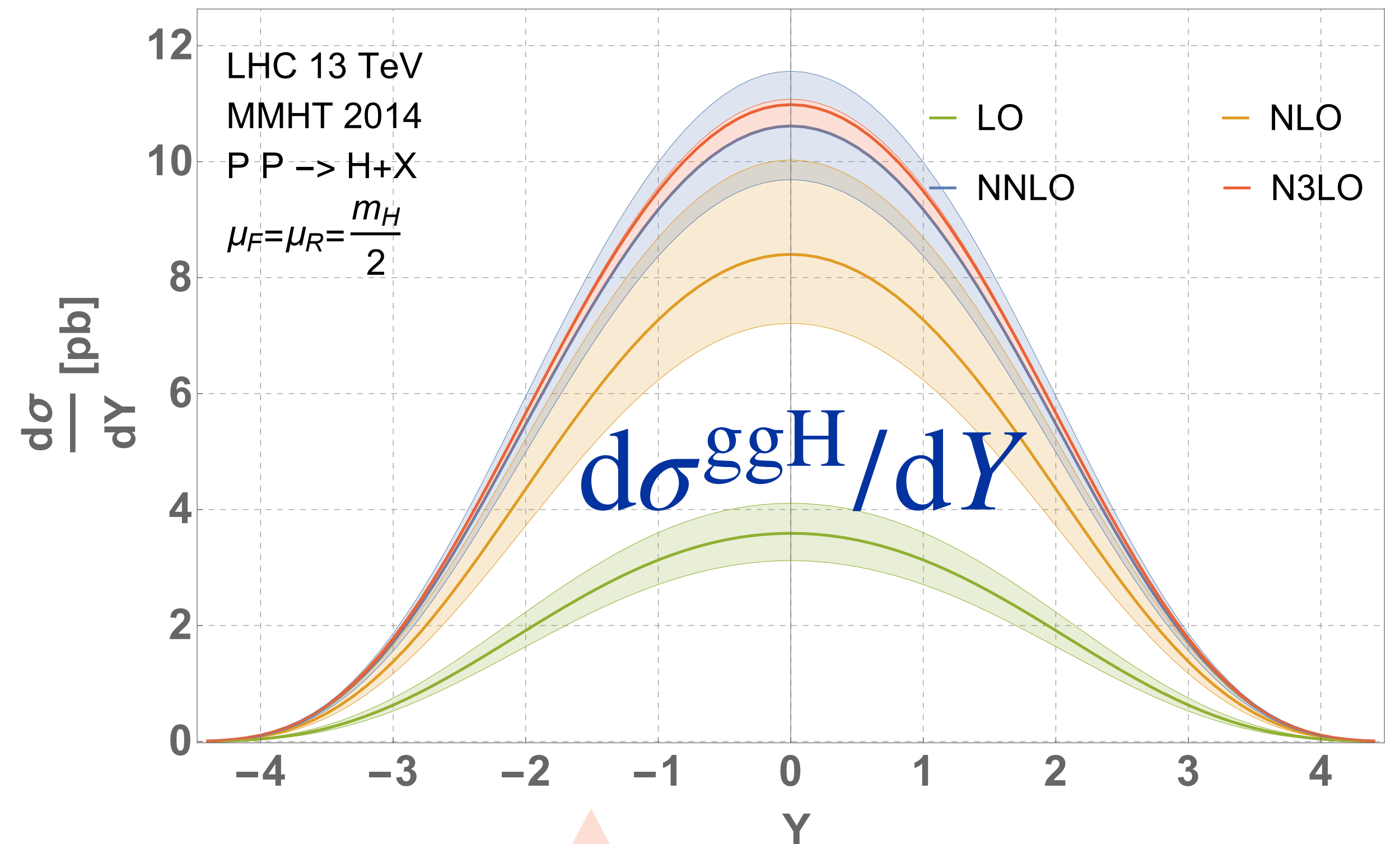
# HIGGS $ggH$ @ N3LO – INCLUSIVE\* PREDICTIONS

[Anastasiou, Duhr, Dulat, Herzog, Mistlberger '15]



nice convergence of perturbative expansion

[Dulat, Mistlberger, Pelloni '18]



differential info lost:  
→ Higgs kinematics, QCD radiation, ...

\* analytically integrated over emissions: ⊕ extremely fast; ⊖ idealised setup

# FULLY DIFFERENTIAL ggH @ N3LO

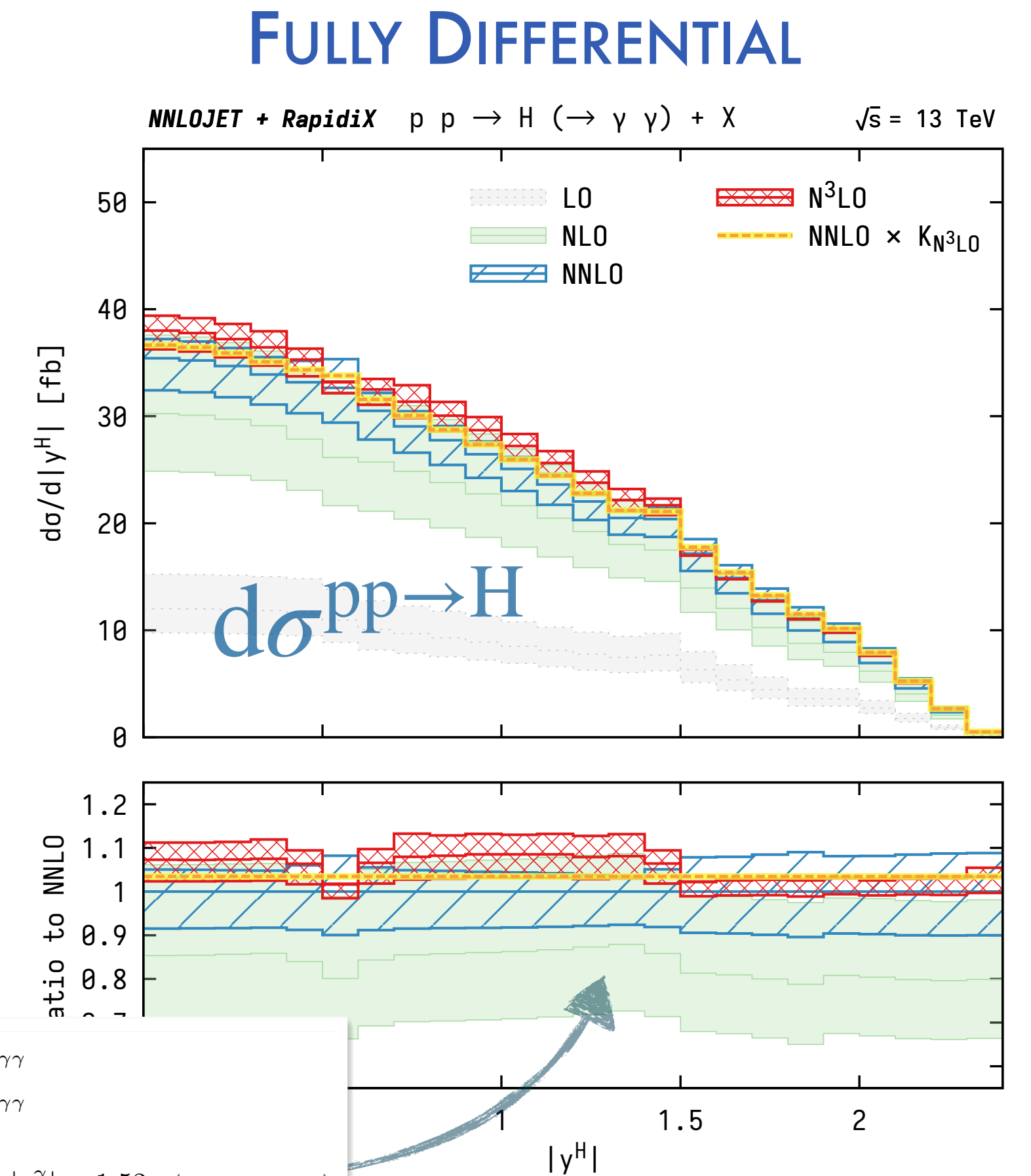
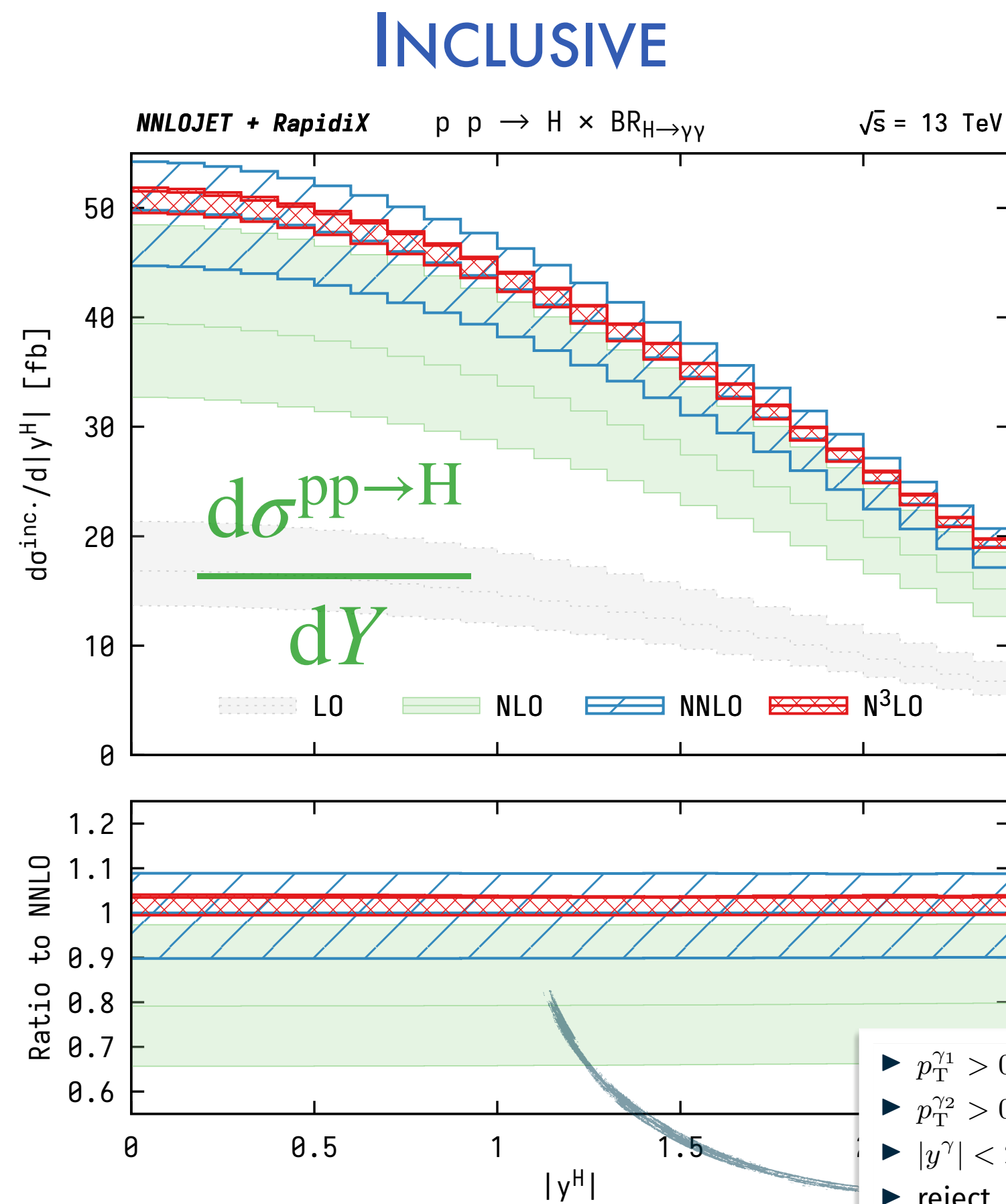
[Chen, Gehrmann, Glover, AH, Mistlberger, Pelloni '21]

## Projection-to-Born:

[VBF @ NNLO: Cacciari et al. '15]

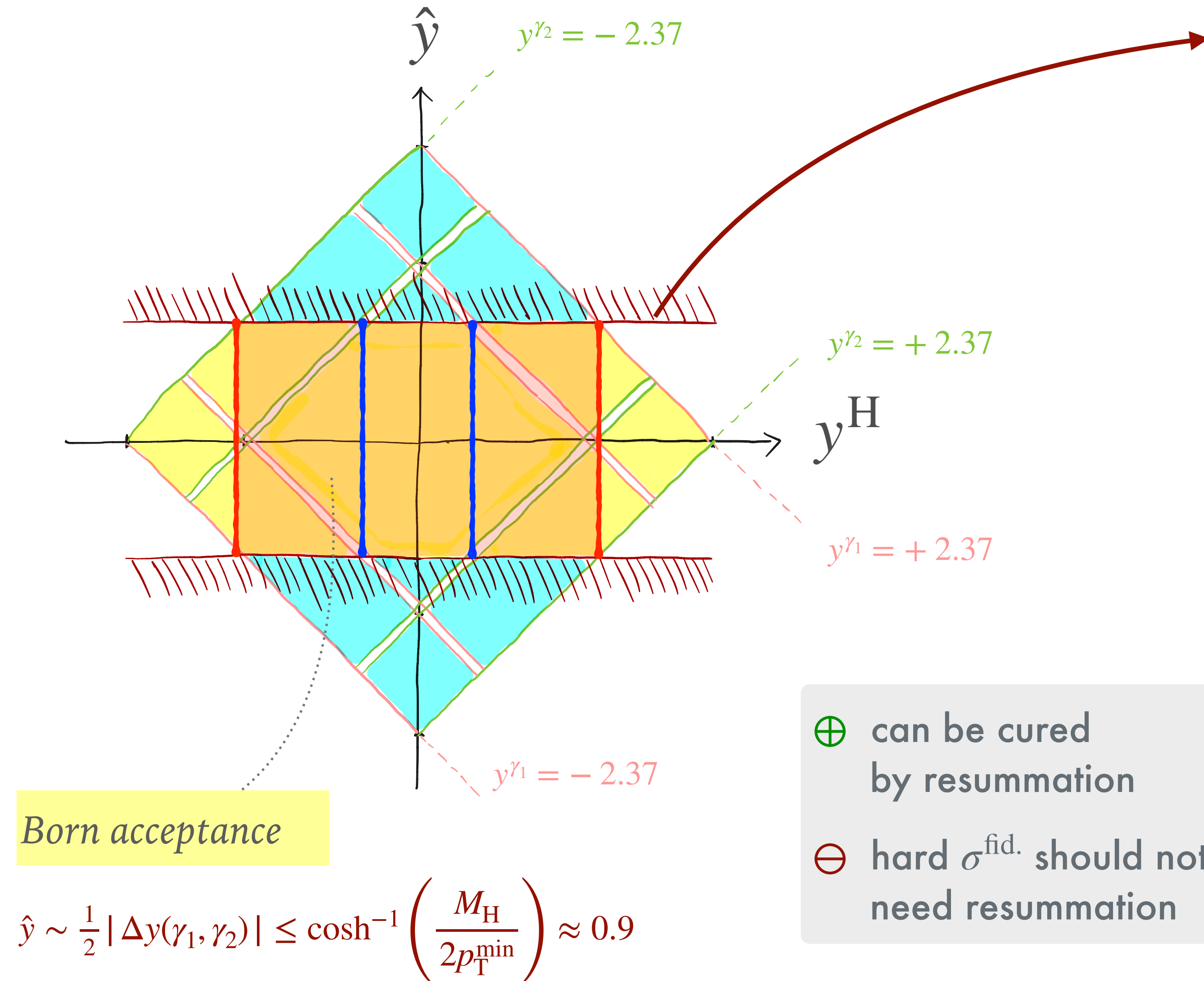
$$\frac{d\sigma_F^{N^k\text{LO}}}{d\mathcal{O}} = \frac{d\sigma_{F, \text{inc.}}^{N^k\text{LO}}}{d\mathcal{O}_B} + \left\{ \frac{d\sigma_{F+\text{jet}}^{N^{k-1}\text{LO}}}{d\mathcal{O}} - \frac{d\sigma_{F+\text{jet}}^{N^{k-1}\text{LO}}}{d\mathcal{O}} \Big|_{\mathcal{O} \rightarrow \mathcal{O}_B} \right\}$$

idea: restore *differential* info of an *inclusive* calculation



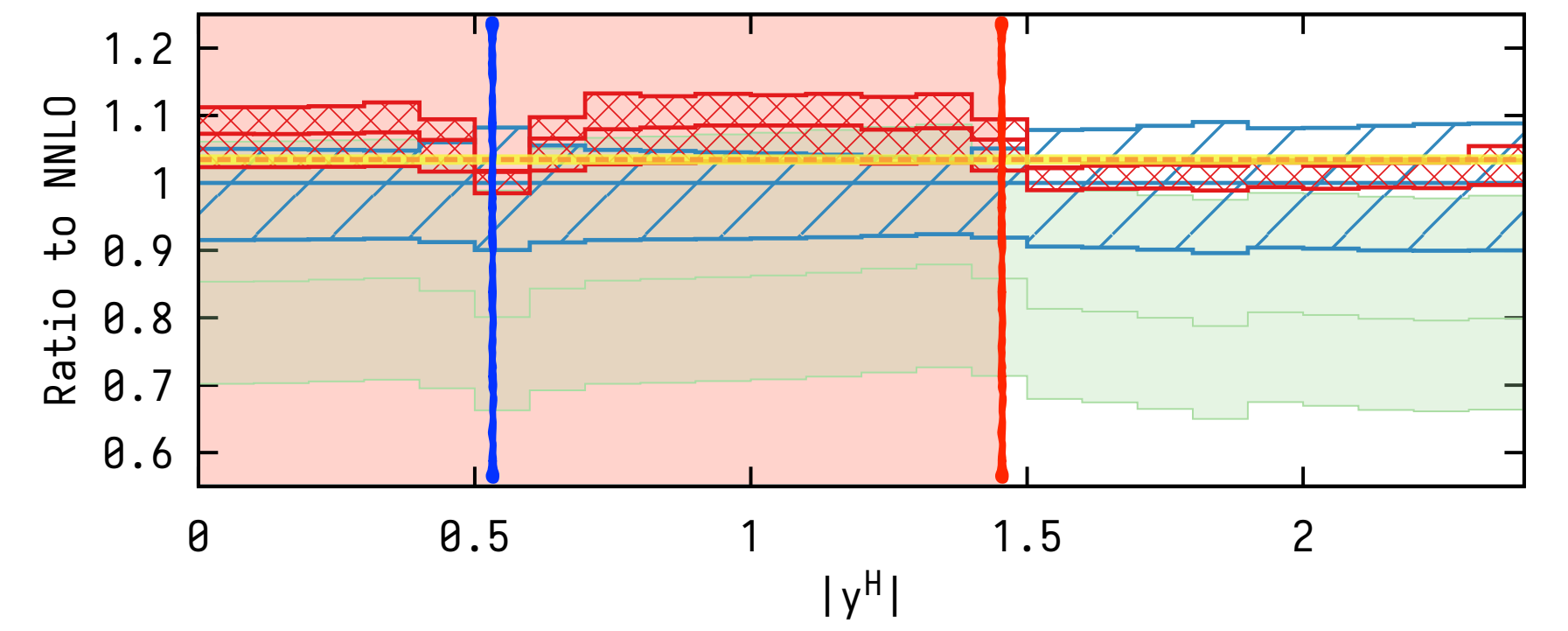
- ▶  $p_T^{\gamma 1} > 0.35 \cdot m_{\gamma\gamma}$
  - ▶  $p_T^{\gamma 2} > 0.25 \cdot m_{\gamma\gamma}$
  - ▶  $|y^\gamma| < 2.37$
  - ▶ reject  $1.37 < |y^\gamma| < 1.52$  (barrel-endcap)
  - ▶ photon isolation in  $\Delta R < 0.2$
- $$\Leftrightarrow \sum_{\Delta R_{i\gamma} < 0.2} p_{T,i} < 0.05 \cdot E_T^\gamma$$

# FIDUCIAL ACCEPTANCES & $y_H$

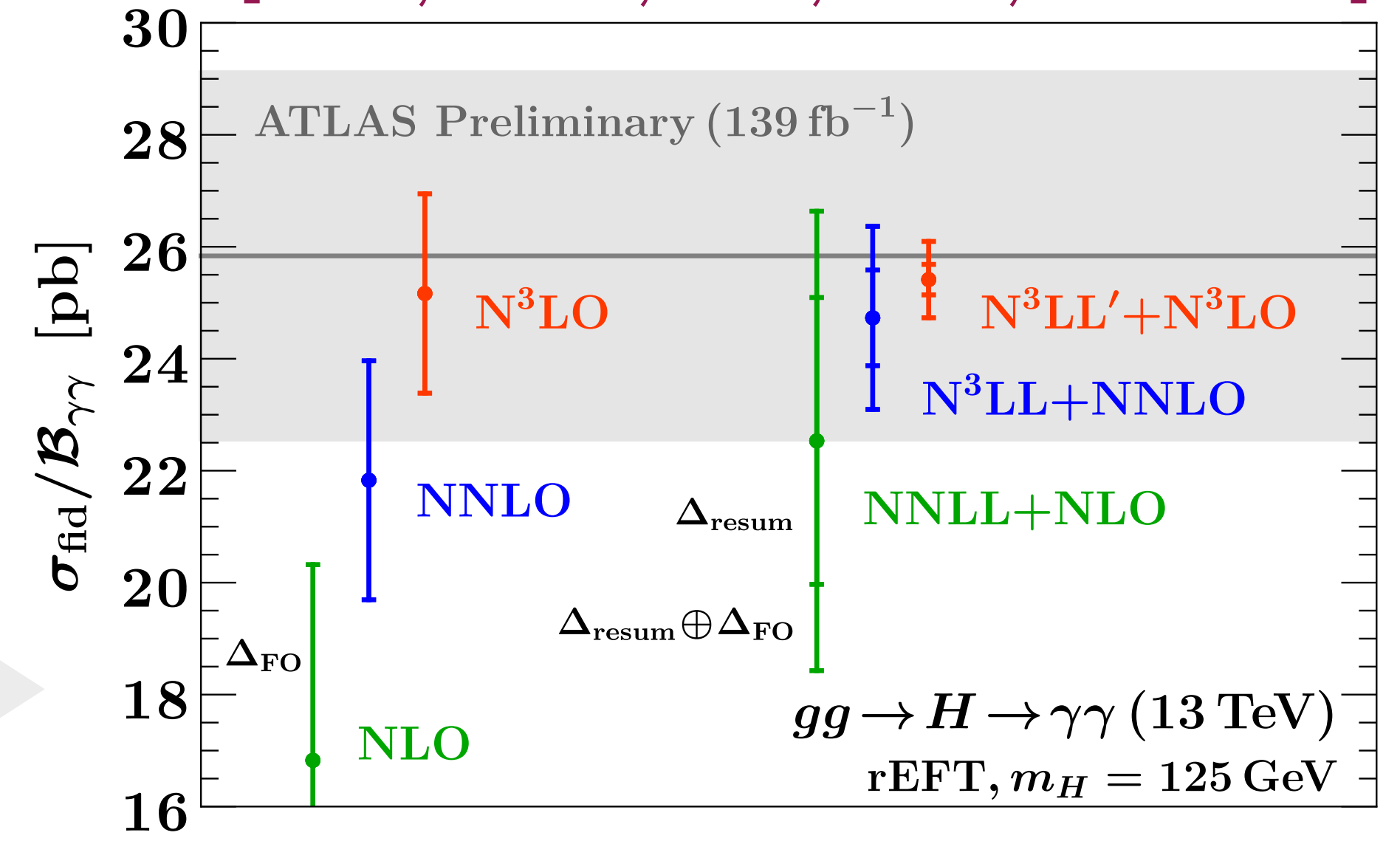


- $\oplus$  can be cured by resummation
- $\ominus$  hard  $\sigma^{\text{fid.}}$  should not need resummation

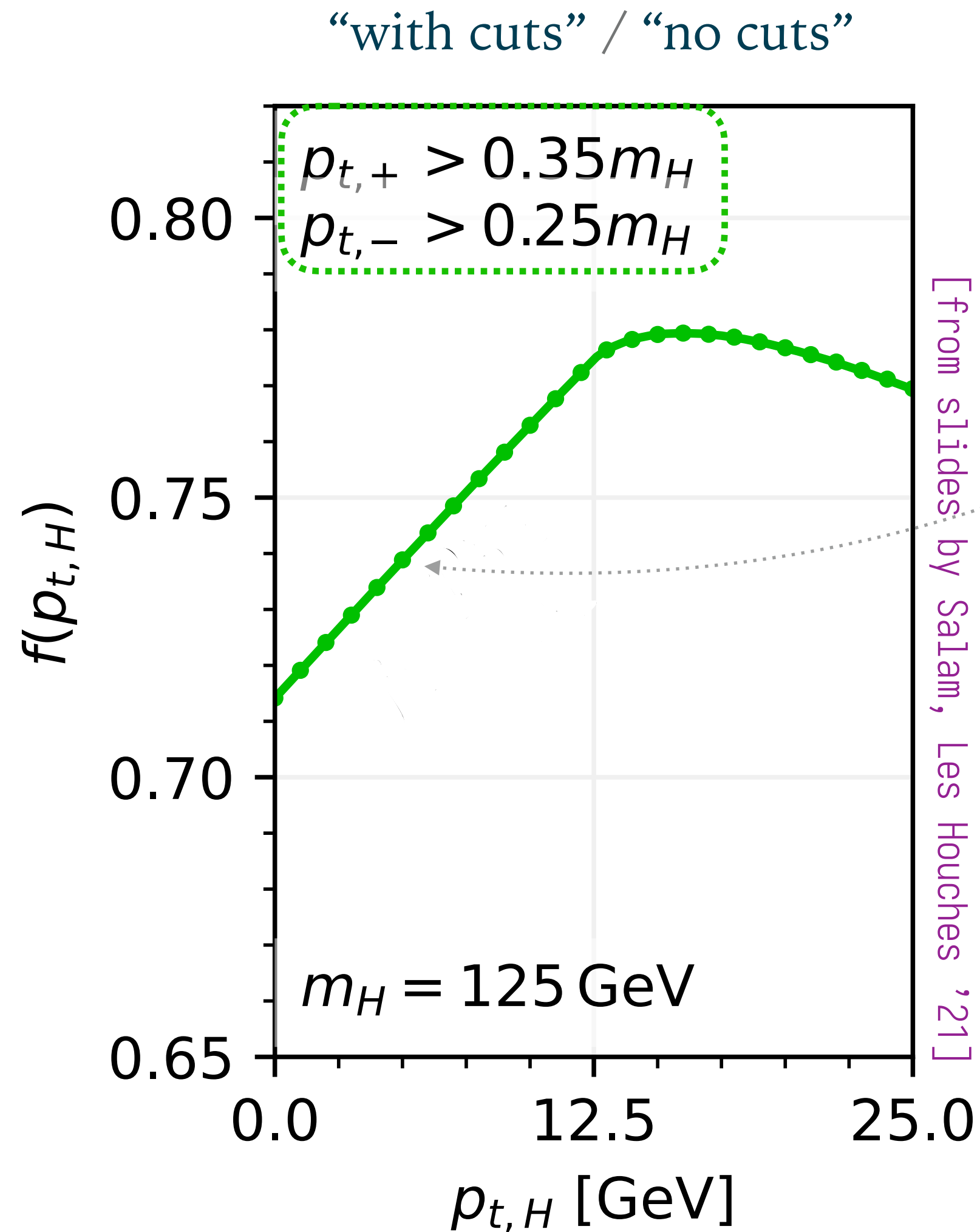
linear fiducial power corrections!



[Billis, Dehnadi, Ebert, Michel, Tackmann '21]



# ACCEPTANCE $f(p_T^H)$



$$f(p_T^H) = f_0 + f_1 \cdot p_T^H + \mathcal{O}((p_T^H)^2)$$

[Frixione, Ridolfi '97; Ebert, Tackmann '19 + Michel, Stewart '21; Alekhin et al. '21]

• Linear  $p_T^H$  dependence

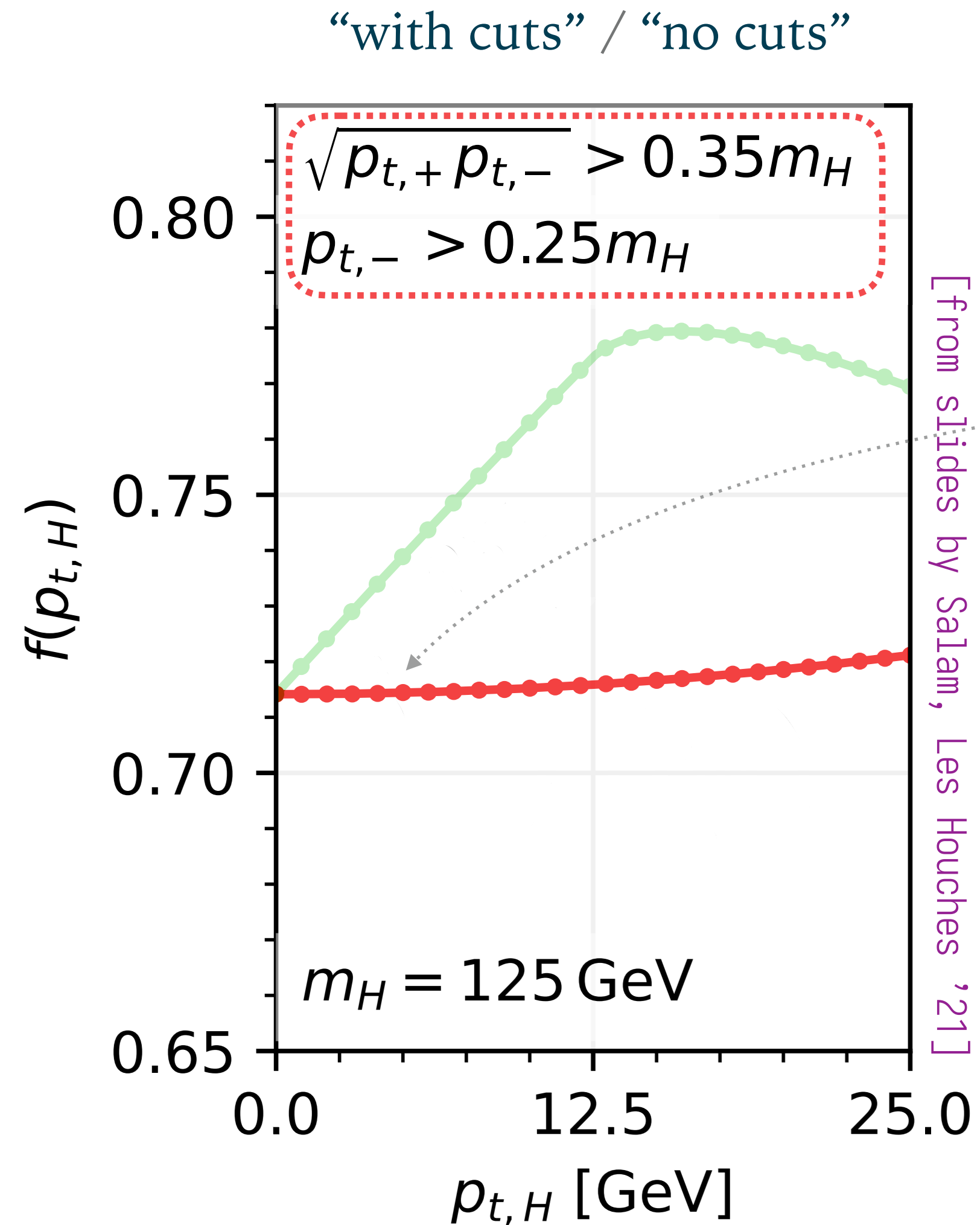
- **factorial growth** for fixed-order
- *sensitivity* to very low  $p_T^H$

$$\frac{\sigma_{\text{asym}} - f_0 \sigma_{\text{inc.}}}{\sigma_0 f_0} \simeq 0.18 \alpha_s - 0.15 \alpha_s^2 + 0.31 \alpha_s^3 + \dots$$

$$\simeq 0.12 @ \text{N}^3\text{LL}$$

[Salam, Slade '21]

# ACCEPTANCE $f(p_T^H)$



$$f(p_T^H) = f_0 + f_1 \cdot p_T^H + f_2 \cdot (p_T^H)^2 + \mathcal{O}((p_T^H)^3)$$

● Quadratic  $p_T^H$  dependence

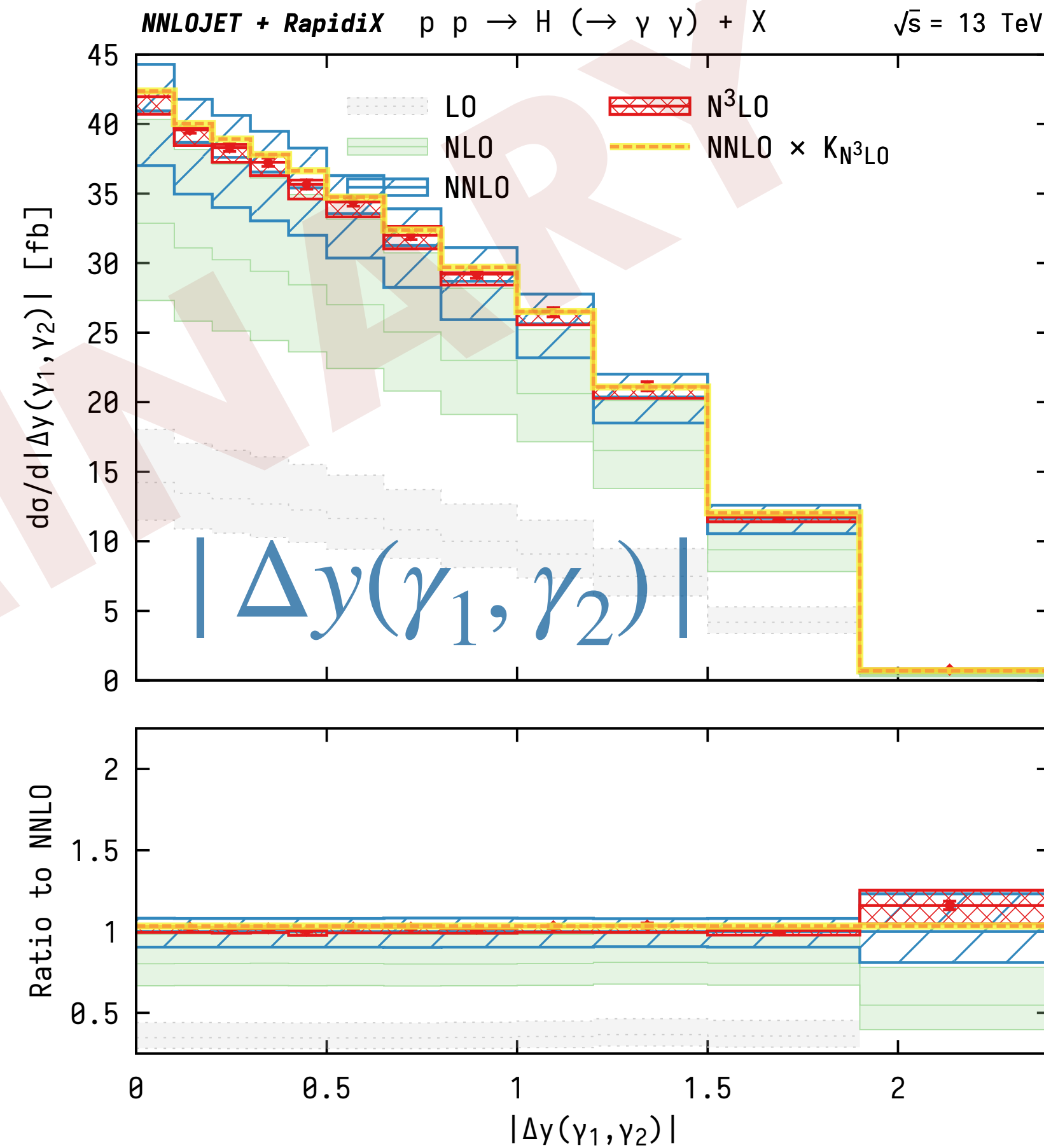
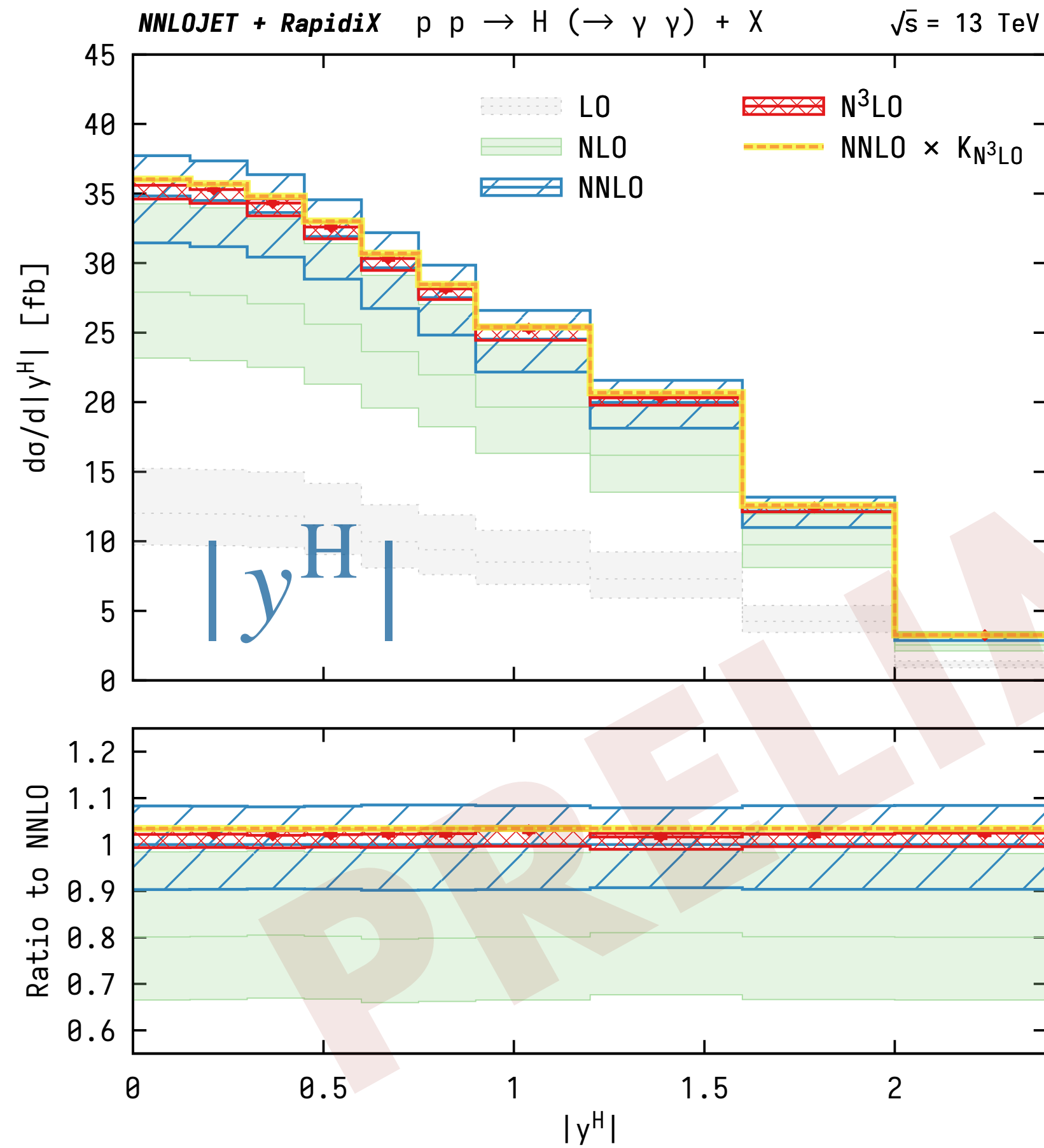
- *suppress* factorial growth
- fixed order  $\simeq$  resummation ✓

$$\frac{\sigma_{\text{prod}} - f_0 \sigma_{\text{inc.}}}{\sigma_0 f_0} \simeq 0.005_{\alpha_s} + 0.002_{\alpha_s^2} - 0.001_{\alpha_s^3} + \dots$$

$$\simeq 0.006 @ \text{N}^3\text{LL}$$

[Salam, Slade '21]

# HIGGS @ N3LO WITH PRODUCT CUTS



$$\sqrt{p_T^{\gamma_1} p_T^{\gamma_2}} \geq 0.35 \cdot M_H$$

$$p_T^{\gamma_2} \geq 0.25 \cdot M_H$$

- no visible instabilities
- ↔ flat  $K$ -factor

- $N^3LO \simeq NNLO \times K_{N^3LO}$

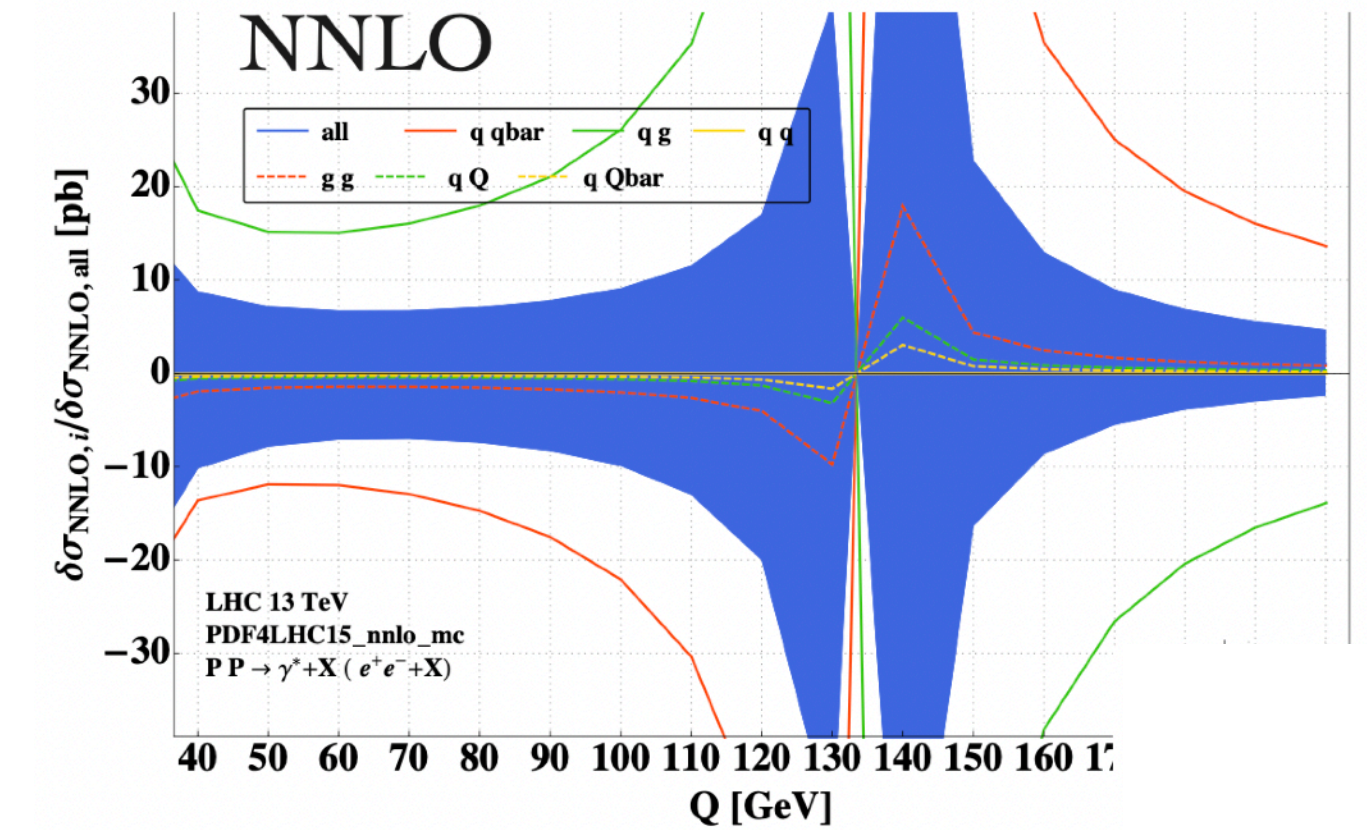
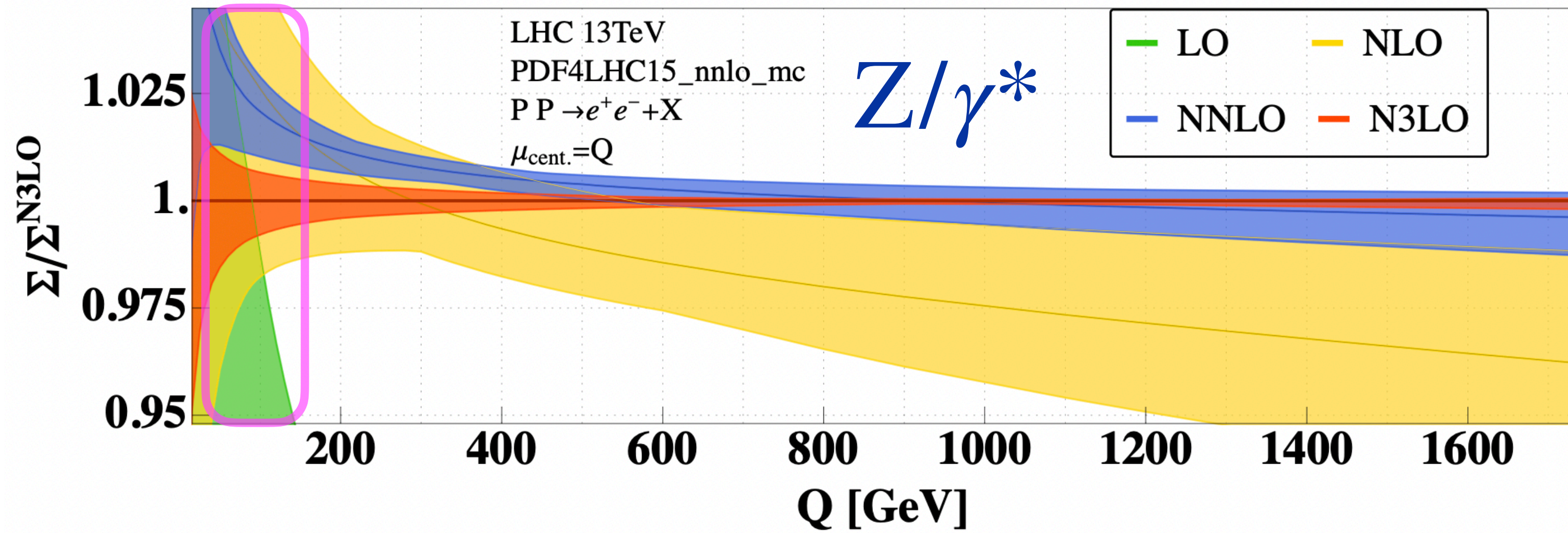
- no visible instabilities; no “features” in the corrections; very flat  $K$ -factors

- $N^3LO \sim NNLO \times K_{N^3LO}$

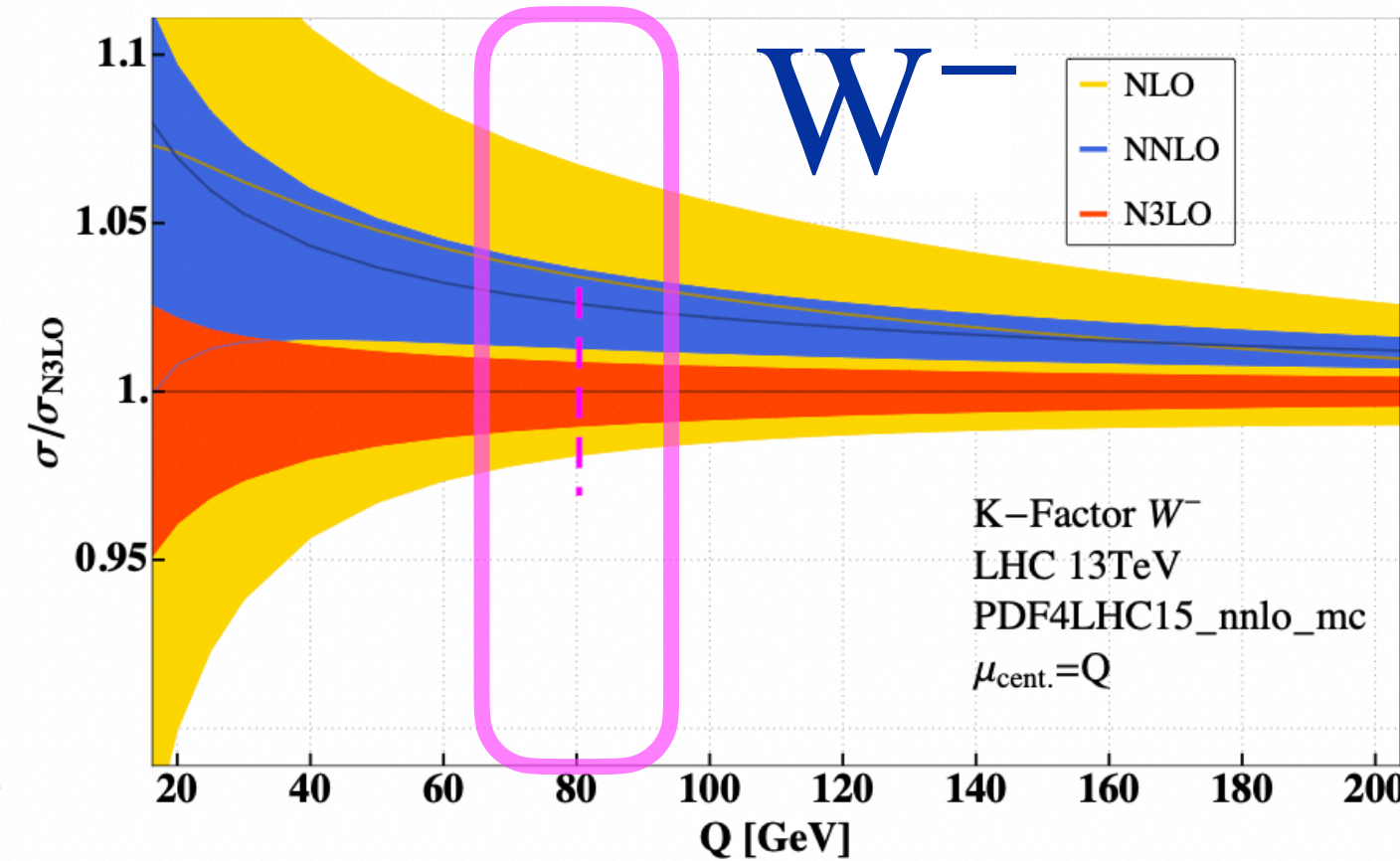
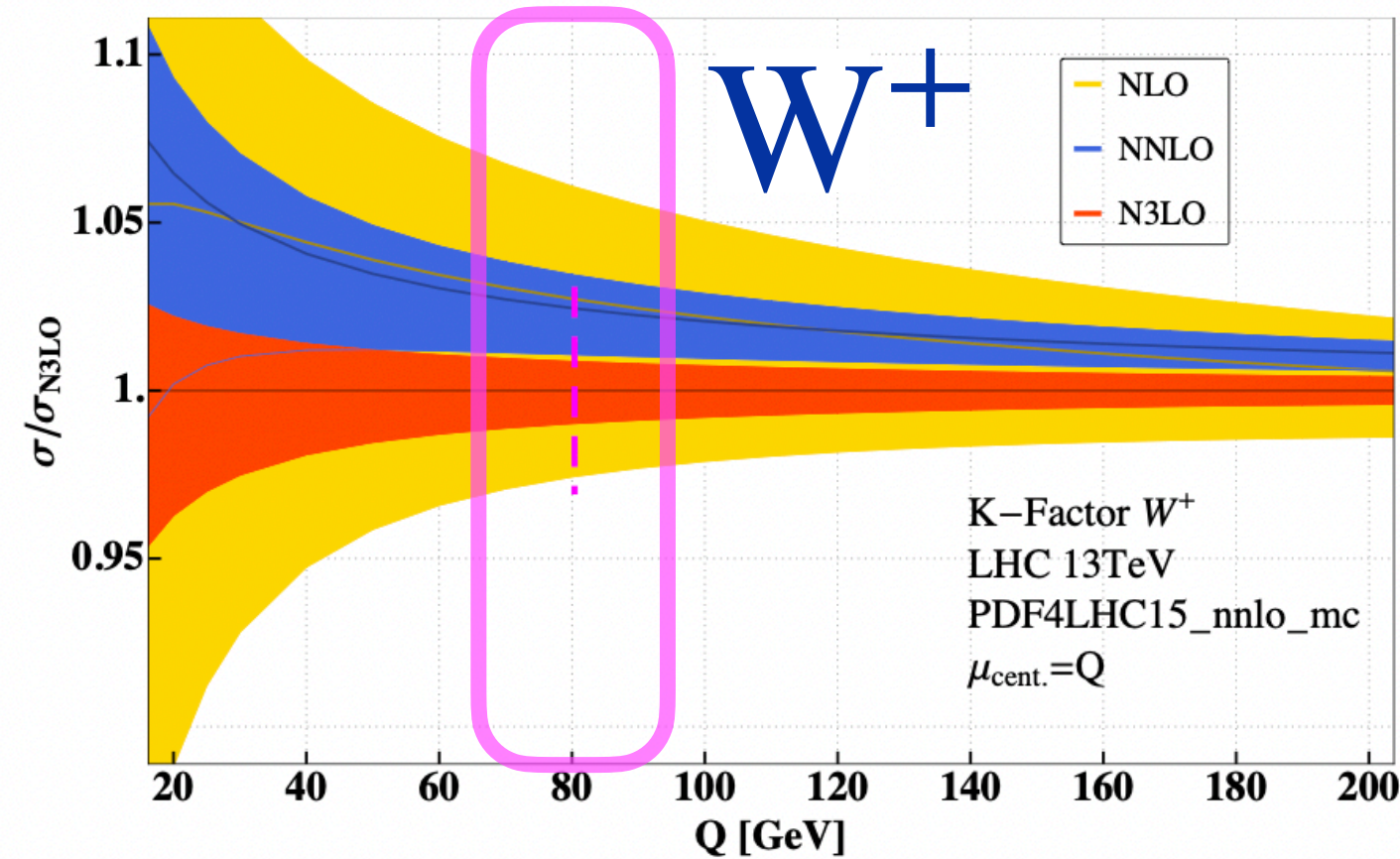


# DRELL-YAN @ N3LO — $Q$ DEPENDENCE

[Dulat, Duhr, Mistlberger '20 '21]



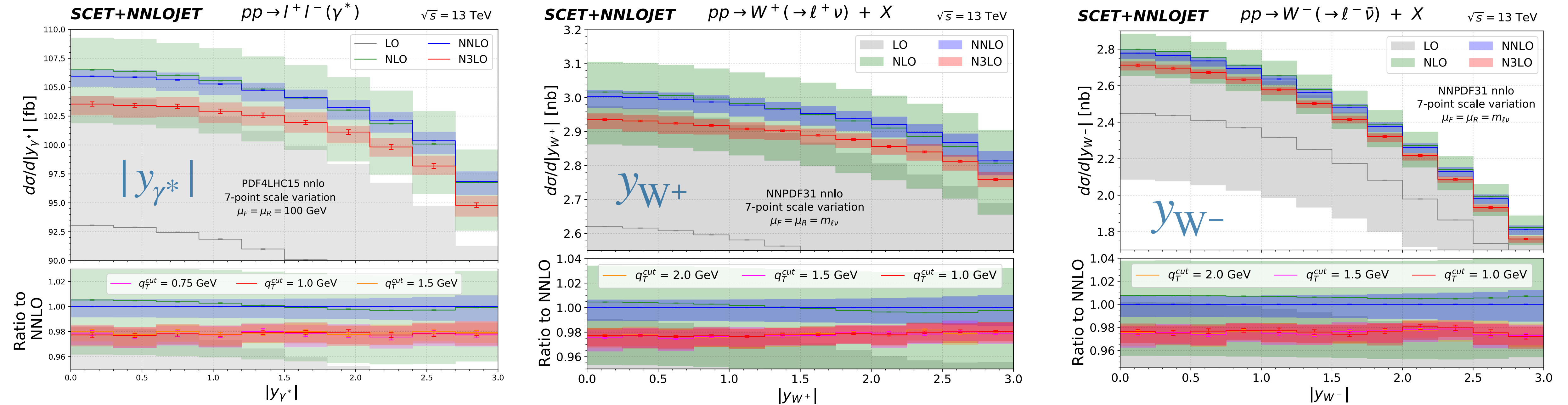
**NNLO:**  $1 \sim \pm 20$   
(large cancellations)  
↪ artificially small?  
**N3LO:**  $1 \sim \pm 2$



resonance region  $\leftrightarrow$  non-overlapping bands;  $\Delta_{\text{scl}}^{\text{NNLO}} \simeq \Delta_{\text{scl}}^{\text{N}^3\text{LO}} ?!$

# DRELL-YAN @ N3LO — $Y_V$ DISTRIBUTIONS

[Chen, Gehrmann, Glover, AH, Yang, Zhu '21, '22]

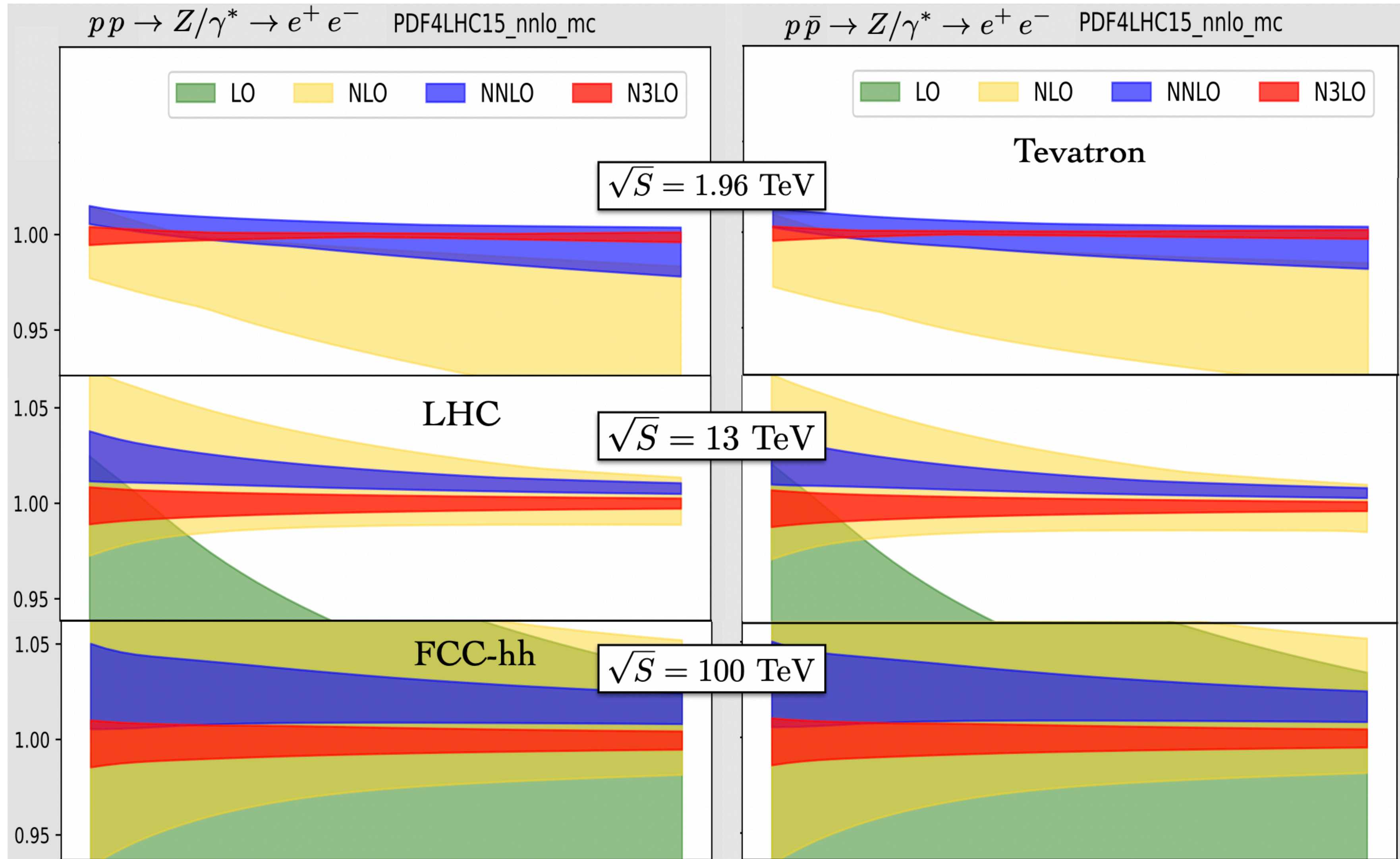


⦿ same collider @ 13 TeV  $\rightsquigarrow$  almost universal NNLO  $\rightarrow$  N<sup>3</sup>LO corrections!

⦿ NC & CC<sup>±</sup> processes probe different parton content across  $Y_V$  (valence u vs. d, ...)

# DRELL-YAN @ N3LO — COLLIDER DEPENDENCE

[slides from C. Duhr: TH colloquium '22]



no "odd" scale behaviour @ 2 TeV

main **difference** from:  
collider energy  
very **similar** between:  
 $pp$  vs.  $p\bar{p}$

# N3LO PARTON DISTRIBUTION FUNCTIONS

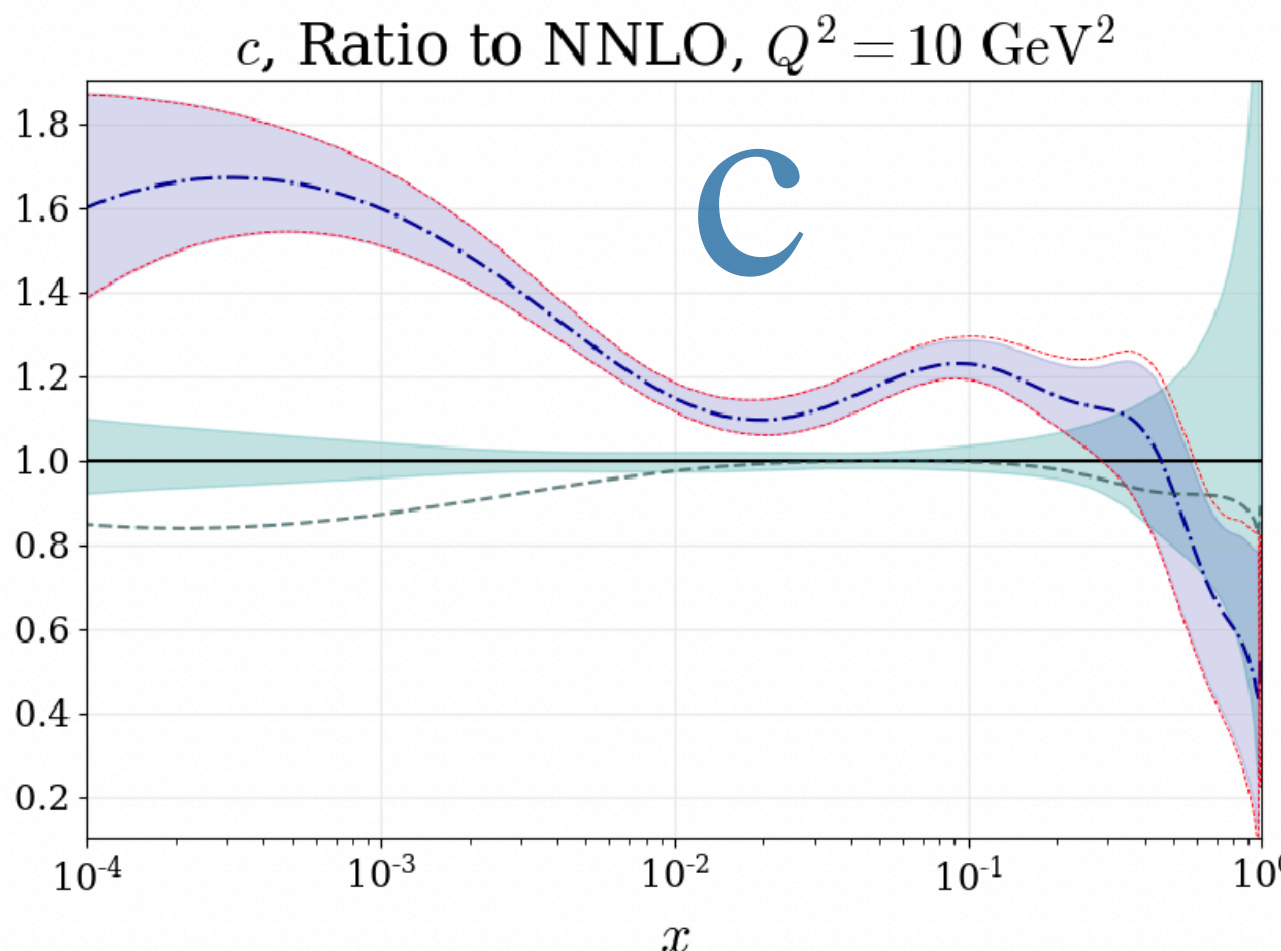
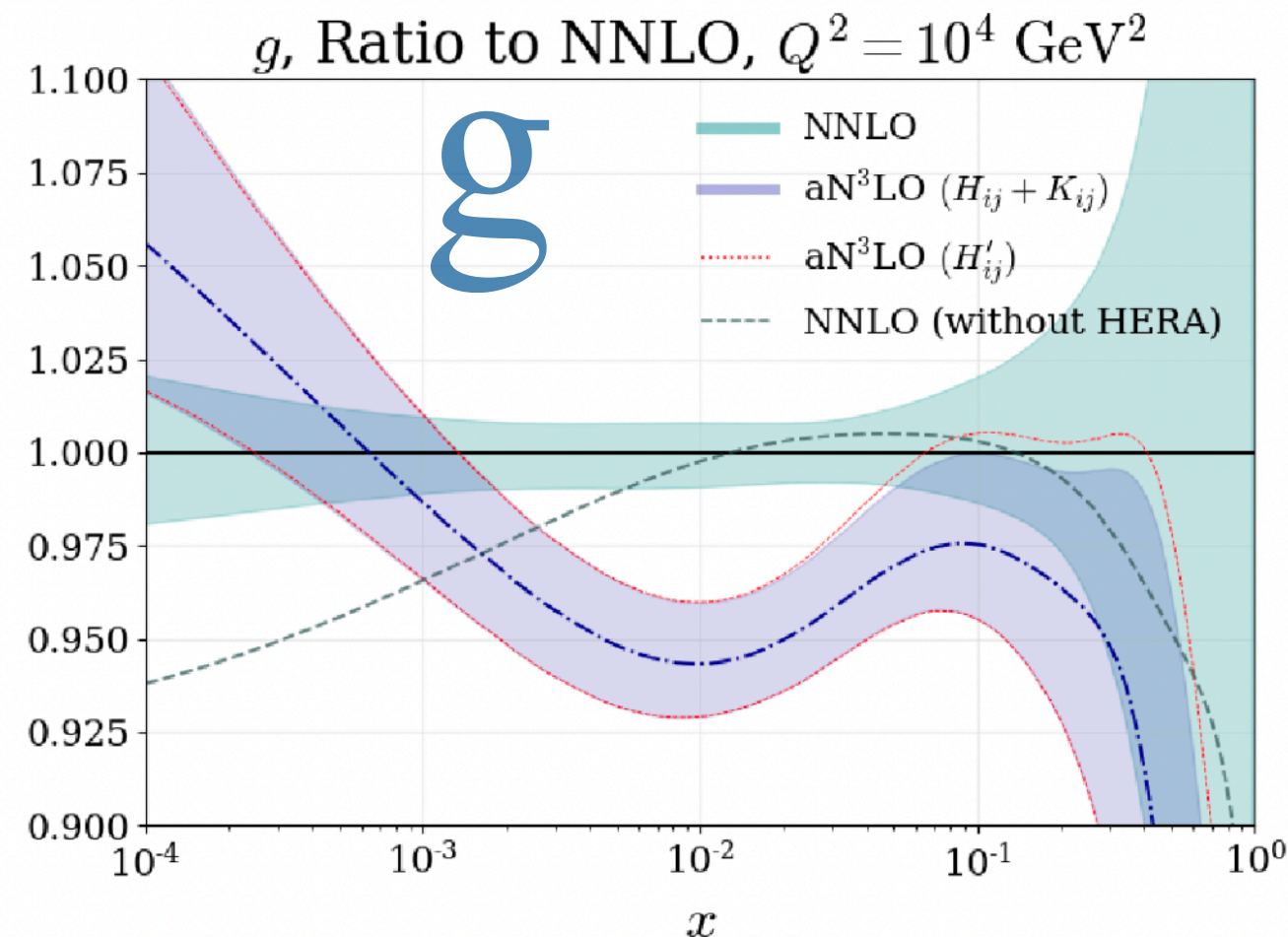
- N3LO evolution

↔ 4-loop splitting functions

[Moch, Ruijl, Ueda, Vermaseren, Vogt '17, '18, '22, in progress]

- aN3LO PDFs (MSHT)

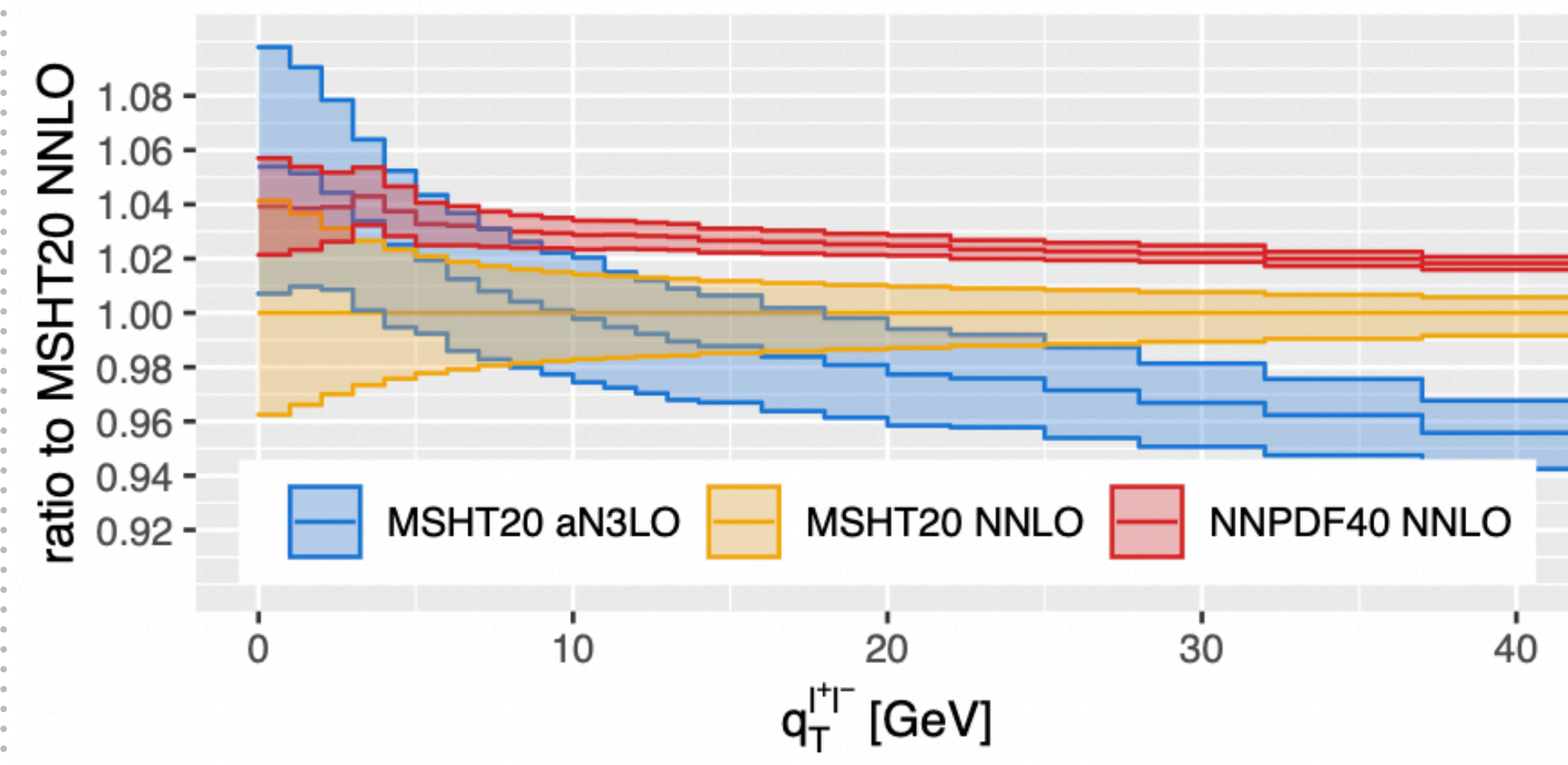
[McGowan, Cridge, Harland-Lang, Thorne '22]



- purely resummed  $p_T^Z$  spectrum

PDF uncertainties

[Neumann, Campbell '22]



sys. differences between PDFs

PDF(NNLO  $\rightarrow$  N<sup>3</sup>LO)  $\delta\sigma^{N^3LO}$   $\nearrow$  (?)

ggH:  $\delta\sigma^{N^3LO}$   $\searrow$

VBF:  $\delta\sigma^{N^3LO}$   $\nearrow$

# CONCLUSIONS & OUTLOOK PART 2

---

- N<sup>3</sup>LO predictions are key to reach percent-level accuracy
  - computation of *inclusive*  $2 \rightarrow 1$  processes very mature  $\leftrightarrow$  ggH, DY, VBF, VH, ...
  - differential predictions for pp  $\rightarrow$  "colour neutral" appearing  $\leftrightarrow$  relies on very stable NNLO "+jet" calculation
  - but: performance of slicing methods very poor  $\leftrightarrow$   $\mathcal{O}(10\text{M})$  CPU core hours
- Fiducial cuts  $\leftrightarrow$  linear power corrections (other processes?)
  - $\hookrightarrow$  crucial for practicability of slicing approaches
- Inadequacies in traditional scale variations  $\leftrightarrow$  DY @ N<sup>3</sup>LO
  - $\hookrightarrow$  effect from missing N<sup>3</sup>LO PDFs?
  - $\hookrightarrow$  more robust TH uncertainties desirable  
(Padé approximant, Bayesian models, PMC, series transforms, ...)

# THE PLAN.

THE PLAN.

## 1. NNLO predictions for the LHC

- jets & interpolations grids
- identified photons & fragmentation

## 2. Differential N<sup>3</sup>LO

- Higgs & fiducial power corrections
- Drell-Yan & PDFs

## 3. Bayesian approach to MHO

- the abc model & correlations

## 4. Summary & Outlook

# WHAT IS THE **UNCERTAINTY** $\Delta_{\text{TH}}$ OF MY RESULT?

---

- increasingly urgent to address with  $\Delta_{\text{EXP}} \searrow$  ( $\leftrightarrow$  HL-LHC)
  - ▶ what does  $5\sigma$  mean if  $\Delta_{\text{TH}}$  non-negligible?
  - ▶ interpretation of data in need for robust  $\Delta_{\text{TH}}$ : PDF fits,  $\chi^2$  in ATLAS jets, ...
- various sources that contribute to  $\Delta_{\text{TH}}$ :
  - ▶  $\Delta_{\alpha_s}$ ,  $\Delta_{\text{param}}$ : parametric uncertainties  $\leftrightarrow$  exp. extraction
  - ▶  $\Delta_{\text{PDF}}$ : parton distribution functions (PDFs)  $\leftrightarrow$  fits
  - ▶  $\Delta_{\text{non pert.}}$ : hadronisation, UE, ...  $\leftrightarrow$  parton showers [e.g. HERWIG vs. PYTHIA]
  - ▶  $\Delta_{\text{MHO}}$ : *missing higher-order (MHO)* corrections Focus here

# CONVENTIONAL APPROACH FOR $\Delta_{\text{MHO}}$ — SCALE VARIATION

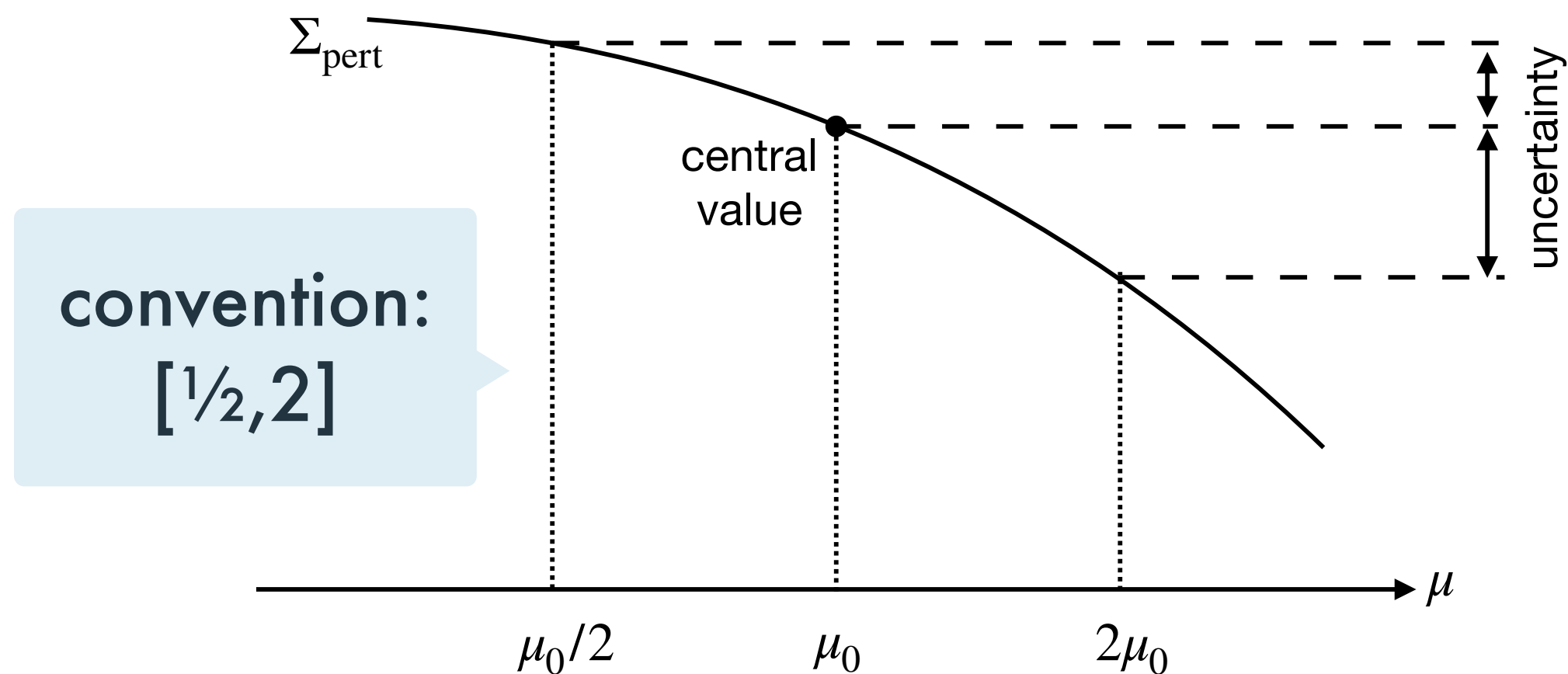
- approximation for an observable @ (next-to-)<sup>n</sup> leading order:

$\triangleright$  N<sup>n</sup>LO:  $\Sigma \simeq \Sigma_n(\mu) = \sum_{k=0}^n \Sigma^{(k)}(\mu) \propto \alpha_s^{n_0+k}$

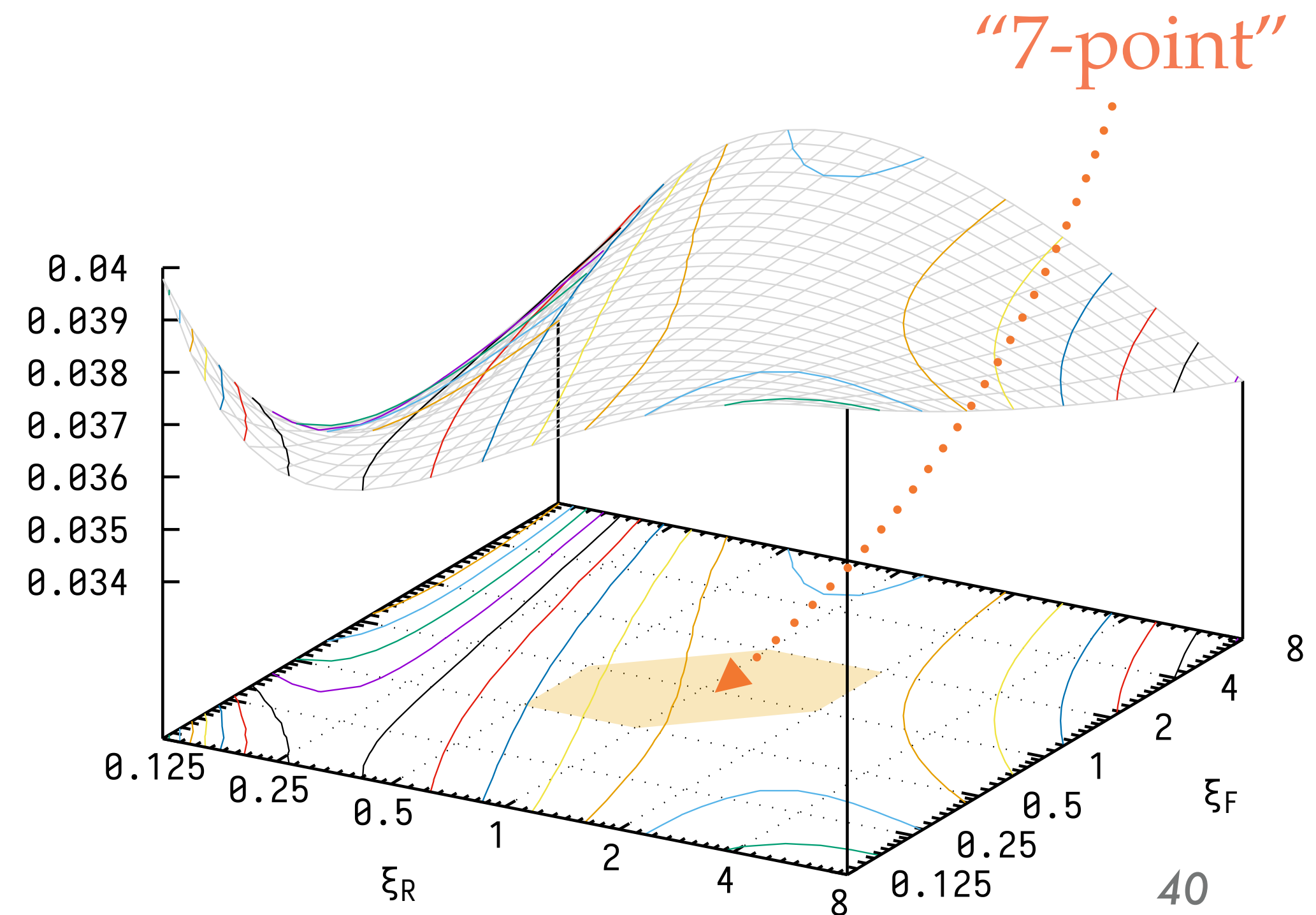
electroweak (EW):  
 $\hookrightarrow$  scheme dependence  
 $\hookrightarrow \alpha \ll \alpha_s$

- truncation of series induces a sensitivity to terms of the next order

$$\mu \frac{d}{d\mu} \Sigma_n(\mu) = \mathcal{O}(\alpha_s^{n_0+n+1}) = \mathcal{O}(\Delta_{\text{MHO}})$$



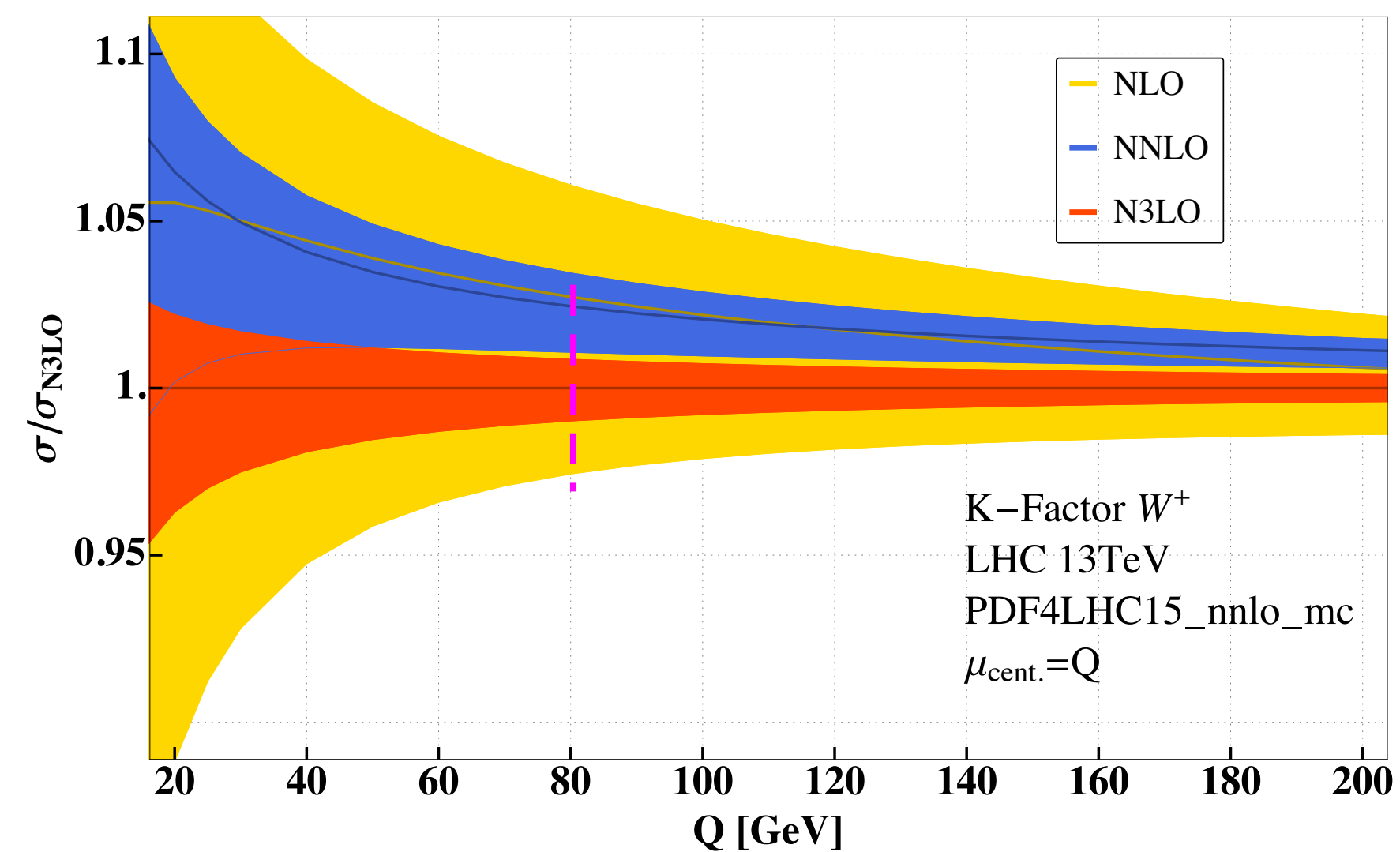
$(\mu_R, \mu_F)$





# ISSUES WITH STANDARD SCALE VARIATIONS

- known to be insufficient:
  - exclusive jet(s) (veto)
  - ratios (correlation?)
  - cancellations (e.g.  $q\bar{q}$  vs.  $qg$  in DY)



[Duhr, Dulat, Mistlberger '20]

- choice of the central scale
  - fastest apparent convergence (FAC)
    - $\hookrightarrow \Sigma^{(n)}(\mu_{\text{FAC}}) = 0$
  - principle of minimal sensitivity (PMS)
    - $\hookrightarrow \left. \frac{\partial}{\partial \mu} \Sigma^{(n)}(\mu) \right|_{\mu_{\text{PMS}}} = 0$
  - BLM/PMC
    - [Brodsky, Lepage, Mackenzie '83]; [Brodsky, Di Giustino '12]
    - ...
- crucially:** *no statistical interpretation!*
  - $\rightsquigarrow$  need to do better

# PROBABILITY DISTRIBUTIONS FOR $\Delta_{\text{MHO}}$

[Cacciari, Houdeau '11]

- Sequence of perturbative *corrections*  $\delta_k$  normalised w.r.t. LO (dimensionless)

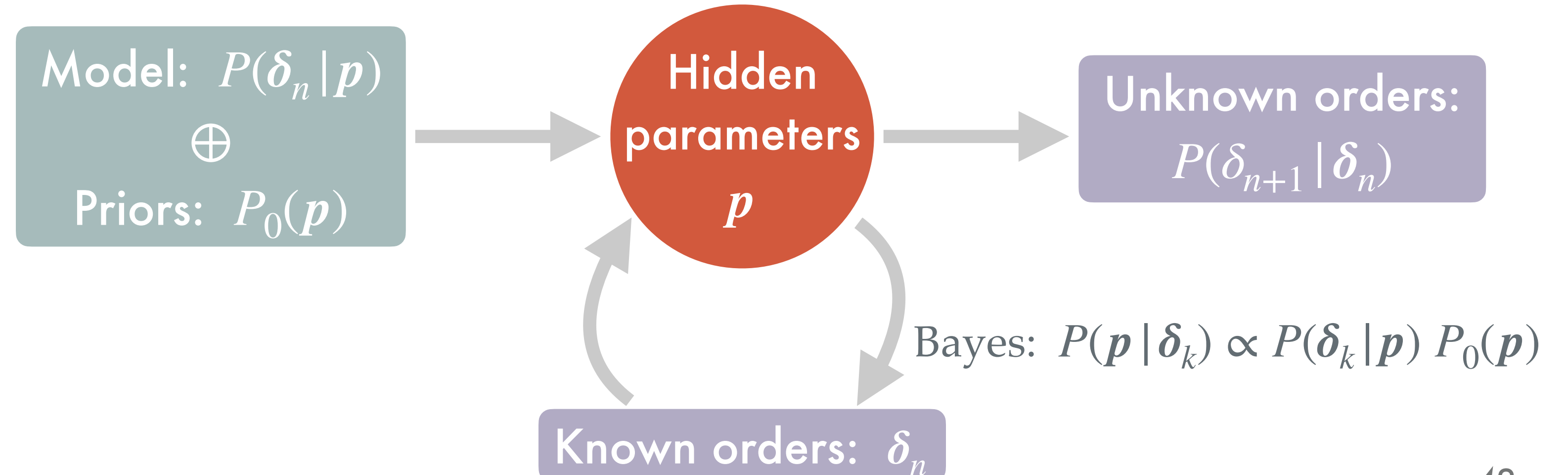
$$\Sigma_n = \Sigma^{(0)} (1 + \delta_1 + \dots + \delta_n) \quad \rightsquigarrow \quad \delta_k = \mathcal{O}(\alpha_s^k)$$

- Probability distribution for  $\delta_{n+1}$ , given  $\delta_n = (\delta_0, \delta_1, \dots, \delta_n)$

$$P(\delta_{n+1} | \delta_n) = \frac{P(\delta_{n+1})}{P(\delta_n)} = \frac{\int d^m \mathbf{p} P(\delta_{n+1} | \mathbf{p}) P_0(\mathbf{p})}{\int d^m \mathbf{p} P(\delta_n | \mathbf{p}) P_0(\mathbf{p})}$$

$$P(A, B) = P(A | B) P(B)$$

$$P(A) = \int dB P(A, B)$$



# THE CH MODEL

[Cacciari, Houdeau '11]

- perturbative expansion  $\delta_k = c_k \alpha_s^k$  bounded by a geometric series:  $|c_k| \leq \bar{c} \quad \forall k$

$$\left| \sum_k \delta_k \right| \leq \sum_k |c_k| \alpha_s^k \leq \sum_k \bar{c} \alpha_s^k$$

- ▶ one hidden parameter:  $\bar{c}$
- ▶ constrain upper bound  $\bar{c}$  from known orders  
     $\rightsquigarrow$  constraint on unknown coefficients  $c_{n+1}$
- limitations:
  - ▶  $\alpha_s$  at what scale? why not:  $\frac{\alpha_s}{\pi}, \frac{\alpha_s}{2\pi}, \alpha_s \ln^2(v), \alpha_s \ln(v), \dots$ ?
- why not let the model figure out the expansion parameter itself?

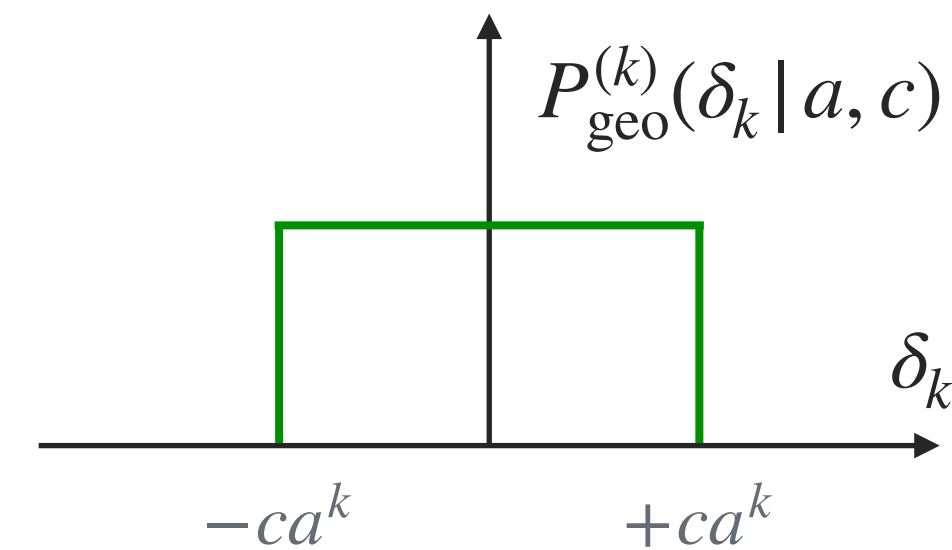
# THE GEOMETRIC MODEL

[Bonvini '20]

- bounded by a geometric series with expansion parameter  $a$ :

$$|\delta_k| \leq c a^k \quad \forall k \quad \Leftrightarrow \text{two model parameters: } a, c$$

- model:** 
$$P_{\text{geo}}^{(k)}(\delta_k | a, c) = \frac{1}{2c a^k} \Theta\left(c - \frac{|\delta_k|}{a^k}\right)$$

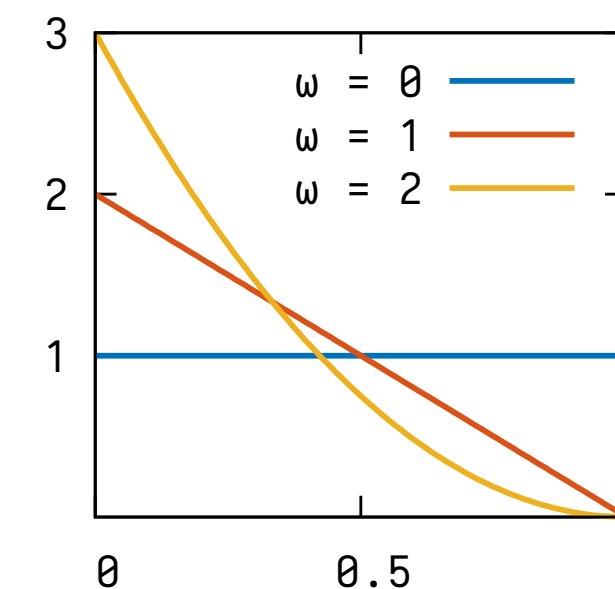


- priors:** 
$$P_0(a, c) = P_0(a) P_0(c)$$

$$P_0(a) = (1 + \omega) (1 - a)^\omega \Theta(a) \Theta(1 - a)$$

$$P_0(c) = \frac{\varepsilon}{c^{1+\varepsilon}} \Theta(c - 1)$$

$$\Leftrightarrow dc/c \sim d \ln(c) \quad (\varepsilon: \text{regulator})$$

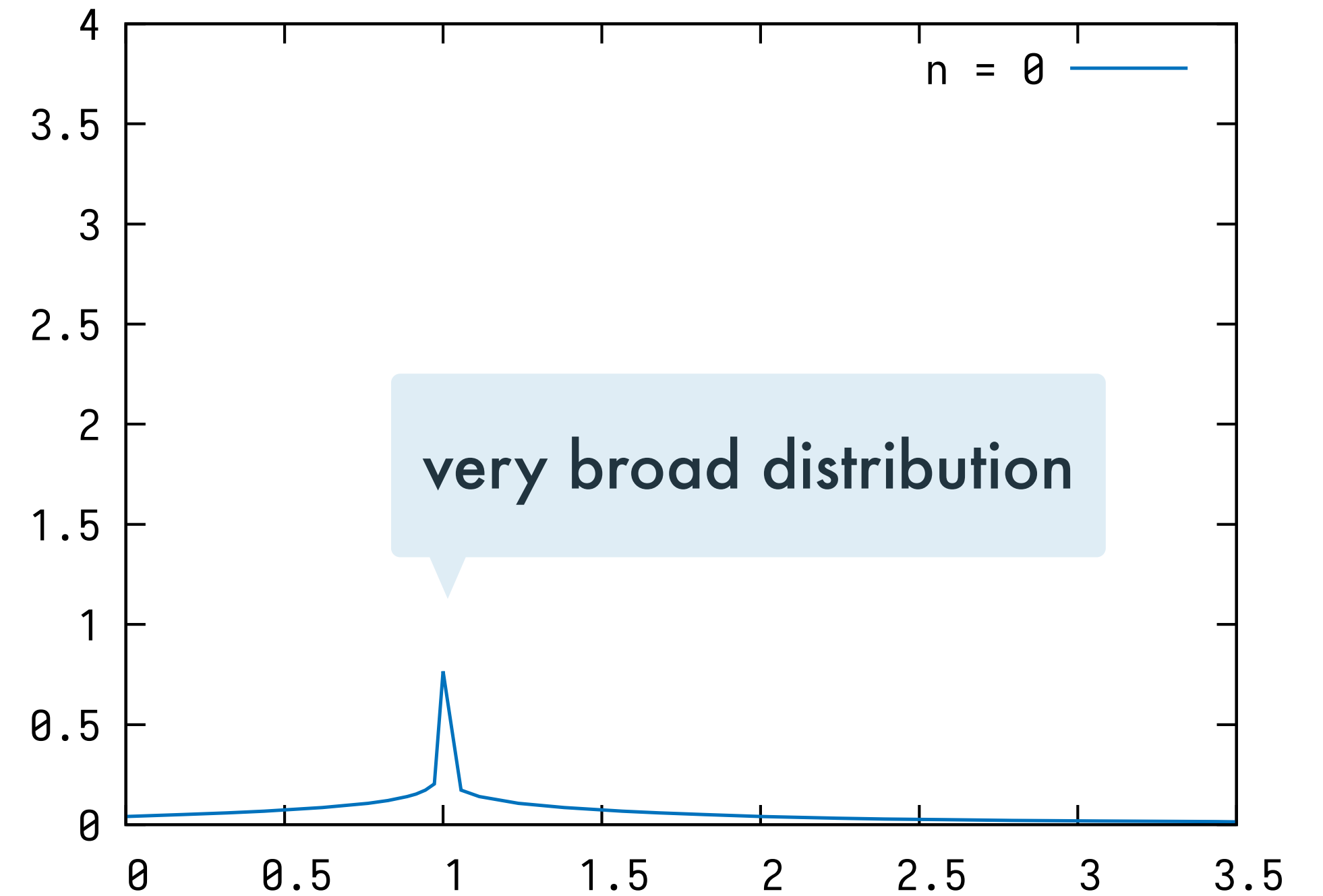
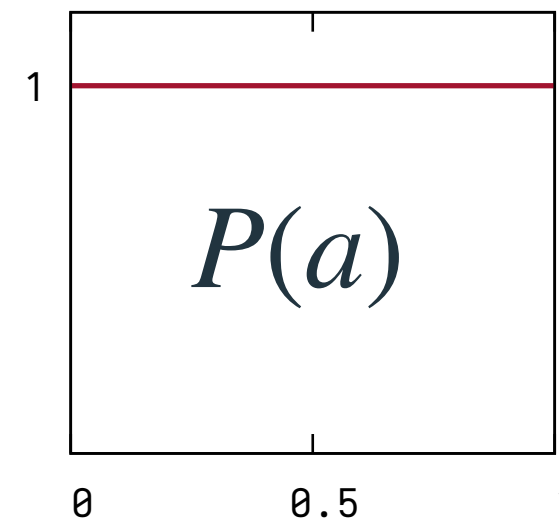


# THE INFERENCE STEP — GEOMETRIC SERIES: $\delta_k = (0.7)^k$

LO  $\delta_0 \equiv 1$

$$P_0(a, c) = \Theta(a) \Theta(1 - a) P_0(c)$$

chose  $\omega = 0$  for flat prior in  $a$



no inference yet!  
 $P(\delta_1)$  entirely determined by the model & priors

$$P(\delta_1) = \int da \int dc P_{\text{geo}}^{(1)}(\delta_1 | a, c) P_0(a, c)$$

# THE INFERENCE STEP — GEOMETRIC SERIES: $\delta_k = (0.7)^k$

- LO  $\delta_0 \equiv 1$

$$P_0(a, c) = \Theta(a) \Theta(1 - a) P_0(c)$$

- NLO  $\delta_1 = 0.7$

$$P(a, c | \delta_1) \propto P_{\text{geo}}^{(1)}(\delta_1 | a, c) P_0(a, c)$$

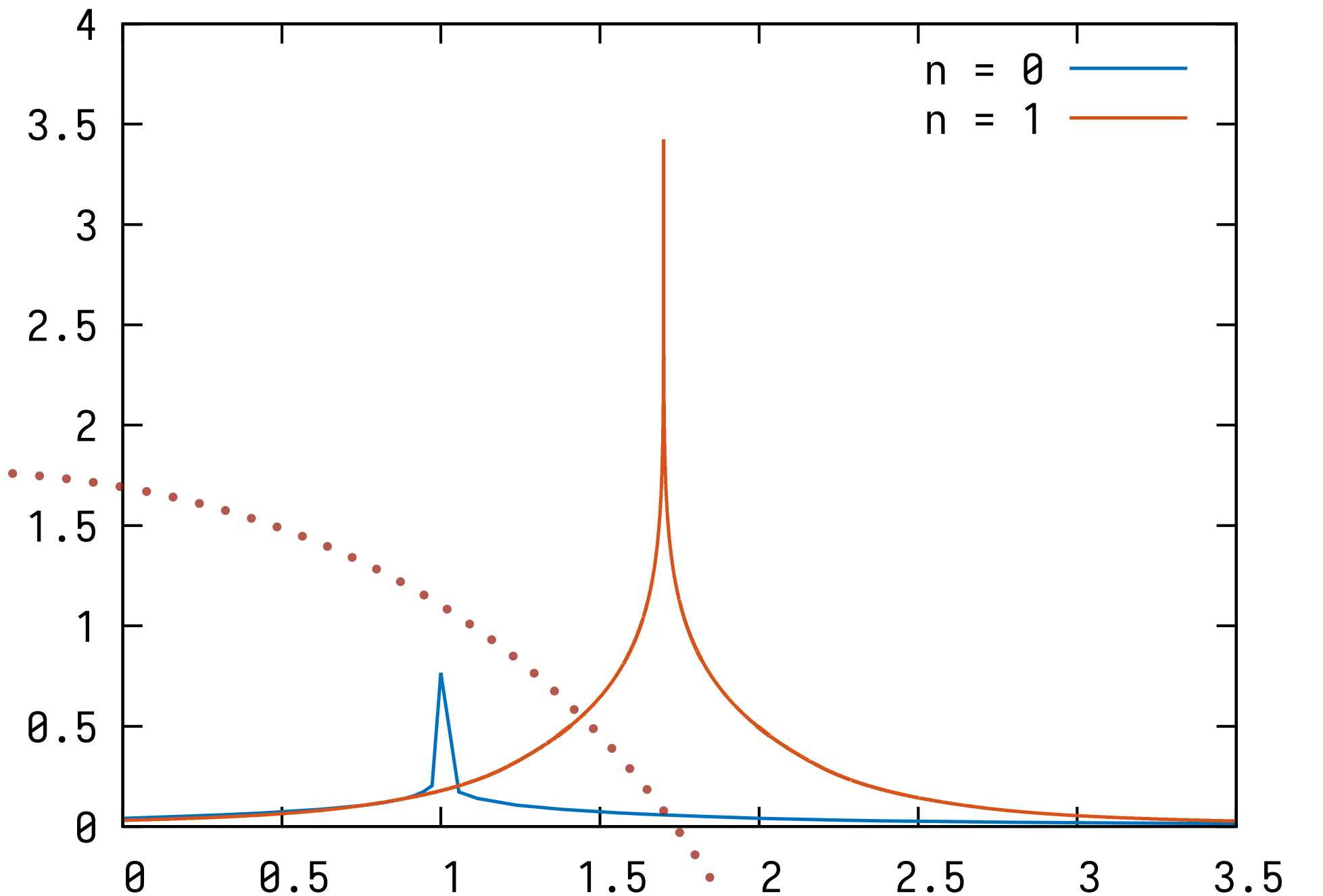
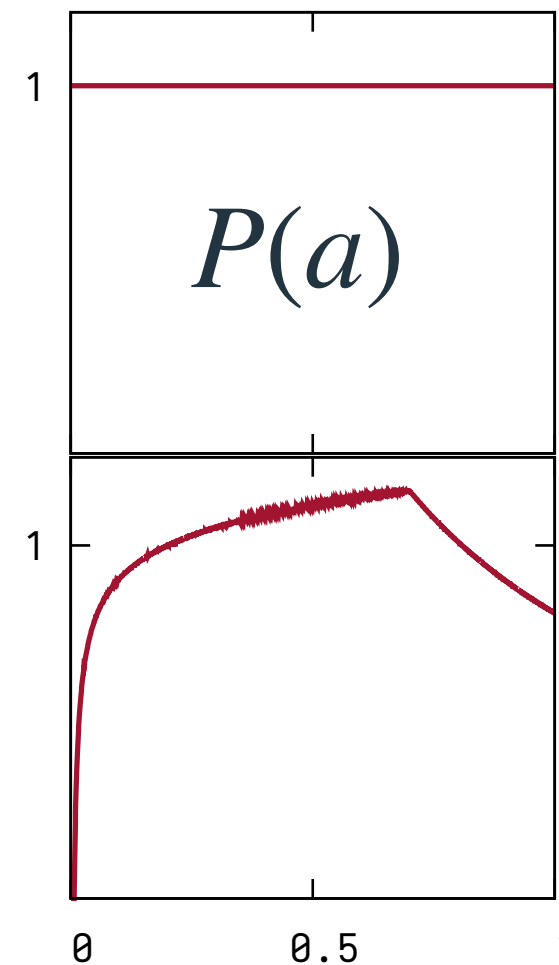
}
}
}
  
 posterior                  likelihood                  prior

Bayes' theorem:

$$P(A | B) = \frac{P(B | A) P(A)}{P(B)}$$

$\delta_k$  independent:

$$P(\delta_2 | \delta_1) = P(\delta_2)$$



$$P(\delta_2 | \delta_1) = \int da \int dc P(\delta_2 | \delta_1, a, c) P(a, c | \delta_1)$$

$$\propto \int da \int dc P_{\text{geo}}^{(2)}(\delta_2 | a, c) P_{\text{geo}}^{(1)}(\delta_1 | a, c) P_0(a, c)$$

# THE INFERENCE STEP — GEOMETRIC SERIES: $\delta_k = (0.7)^k$

- LO  $\delta_0 \equiv 1$

$$P_0(a, c) = \Theta(a) \Theta(1 - a) P_0(c)$$

- NLO  $\delta_1 = 0.7$

$$P(a, c | \delta_1) \propto P_{\text{geo}}^{(1)}(\delta_1 | a, c) P_0(a, c)$$

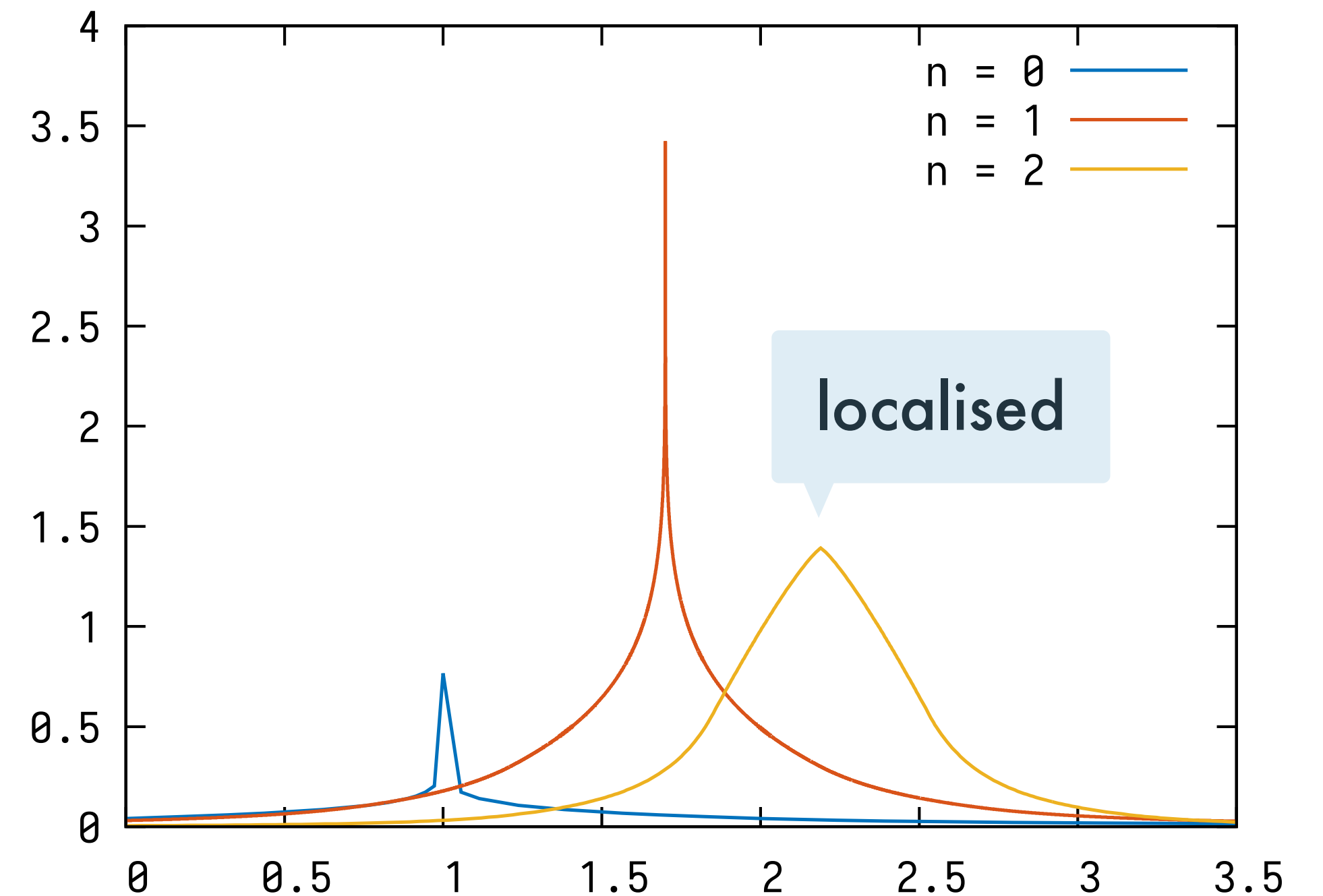
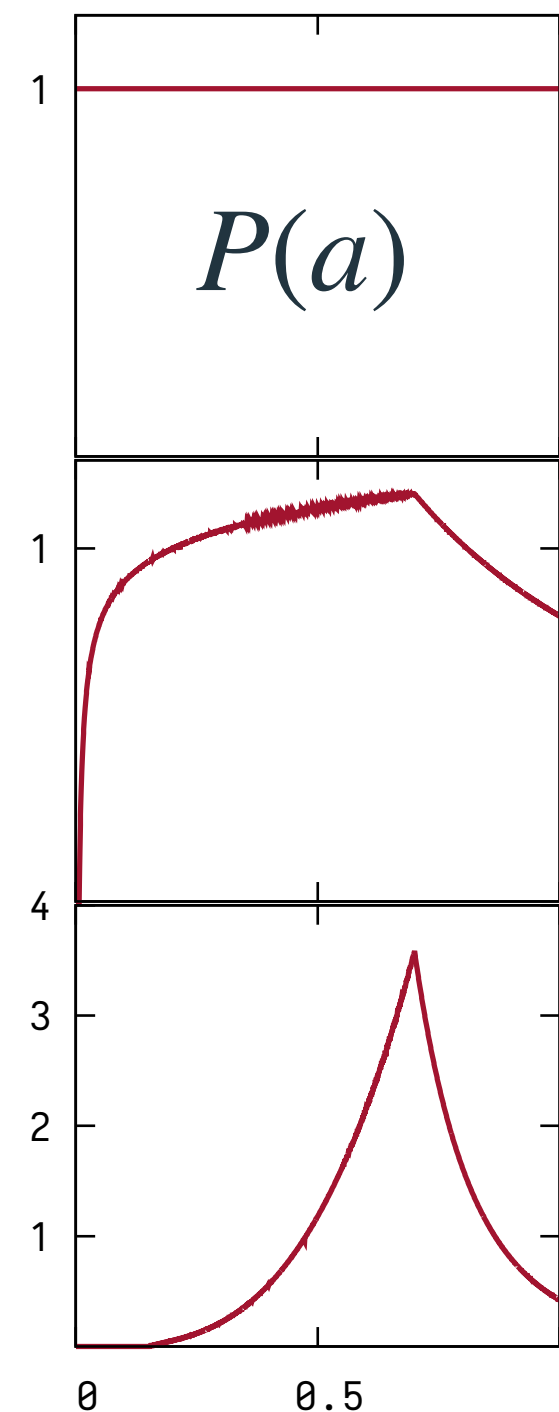
- N<sup>2</sup>LO  $\delta_2 = 0.7^2$

$$P(a, c | \delta_1, \delta_2) \propto P(\delta_2 | \delta_1, a, c) P(a, c | \delta_1)$$

$$\propto P_{\text{geo}}^{(2)}(\delta_2 | a, c) P_{\text{geo}}^{(1)}(\delta_1 | a, c) P_0(a, c)$$

Bayes' theorem  
& independence

$a \sim 0.7$   
also:  $c \sim 1$



$$P(\delta_3 | \delta_1, \delta_2) \propto \int da \int dc \prod_{k=1}^3 \left[ P_{\text{geo}}^{(k)}(\delta_k | a, c) \right] P_0(a, c)$$

# THE INFERENCE STEP — GEOMETRIC SERIES: $\delta_k = (0.7)^k$

- LO  $\delta_0 \equiv 1$

$$P_0(a, c) = \Theta(a) \Theta(1 - a) P_0(c)$$

- NLO  $\delta_1 = 0.7$

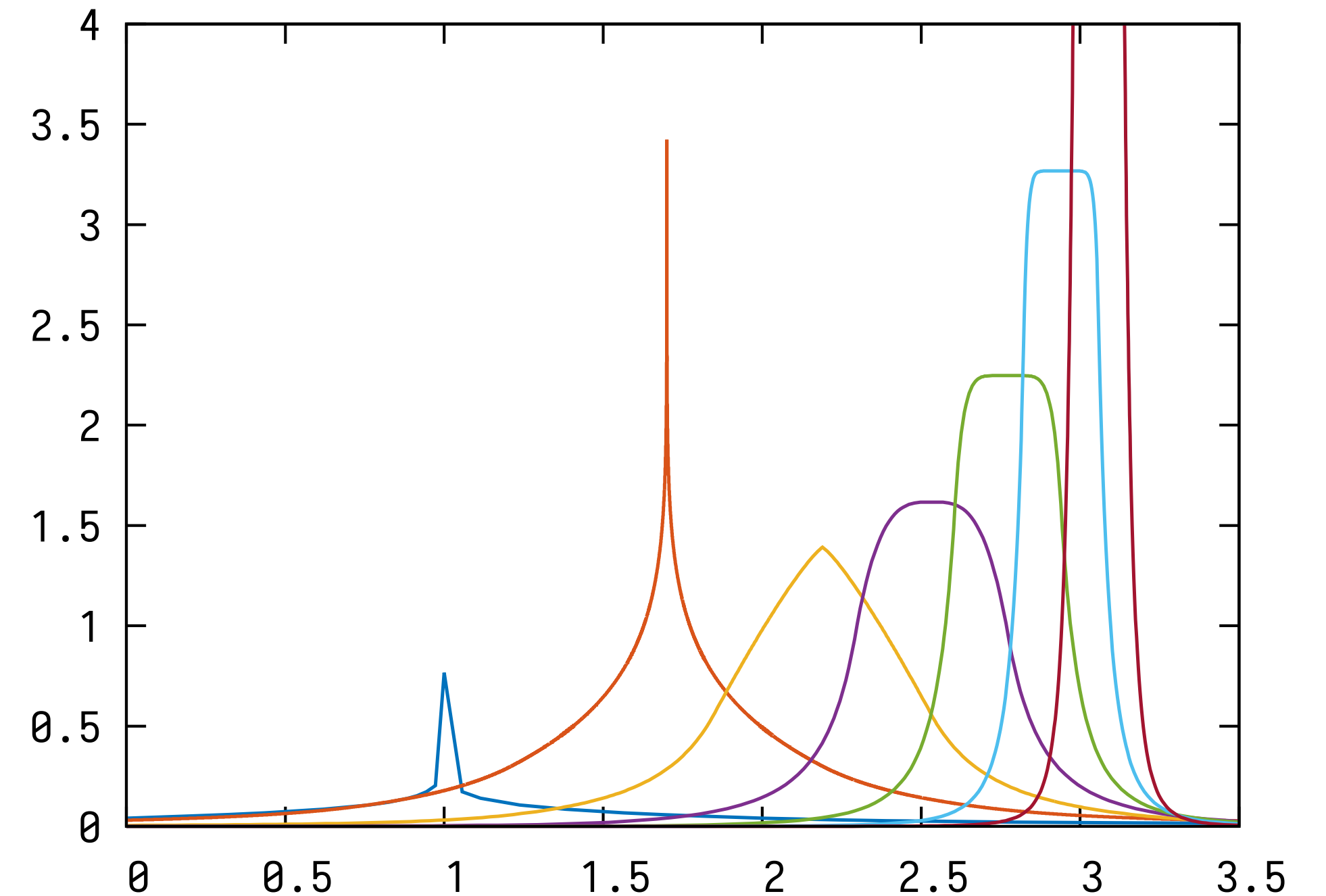
$$P(a, c | \delta_1) \propto P_{\text{geo}}^{(1)}(\delta_1 | a, c) P_0(a, c)$$

- N<sup>2</sup>LO  $\delta_2 = 0.7^2$

$$P(a, c | \delta_1, \delta_2) \propto P(\delta_2 | \delta_1, a, c) P(a, c | \delta_1)$$

$$\propto P_{\text{geo}}^{(2)}(\delta_2 | a, c) P_{\text{geo}}^{(1)}(\delta_1 | a, c) P_0(a, c)$$

- ...



$$P(\delta_{n+1} | \delta_n) \propto \int da \int dc \prod_{k=1}^n \left[ P_{\text{geo}}^{(k)}(\delta_k | a, c) \right] P_0(a, c)$$

can be solved analytically



# THE *abc* MODEL — ASYMMETRIC GEOMETRIC MODEL

[Duhr, AH, Mazeliauskas, Szafron '21]

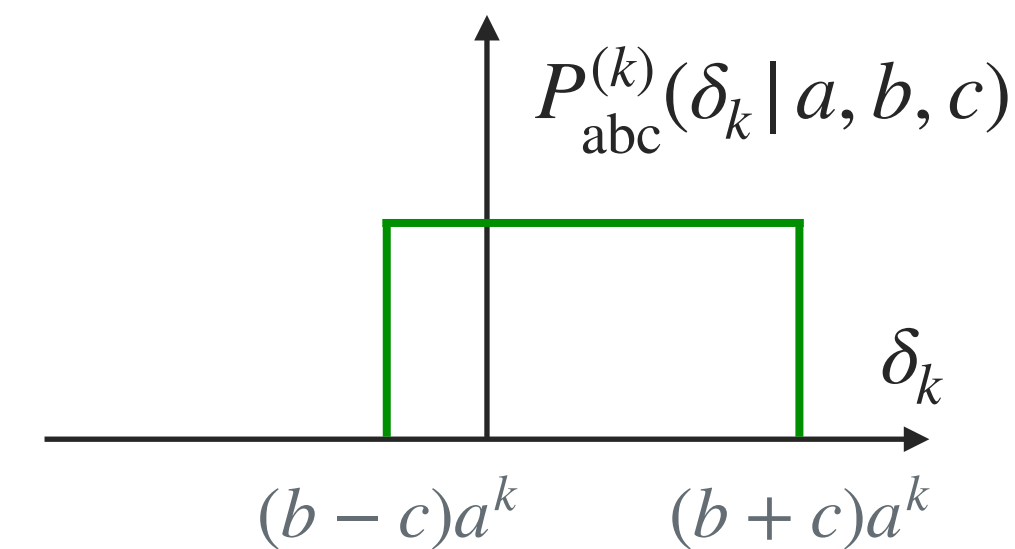
geometric model is symmetric:  $P(\delta_0, \dots, \delta_n) = P(|\delta_0|, \dots, |\delta_n|) \rightsquigarrow \langle \delta_{n+1} \rangle_{\text{geo}} = 0$

allow for different lower & upper bound:

$$b - c \leq \frac{\delta_k}{a^k} \leq b + c \quad \forall k \quad \rightsquigarrow \text{three model parameters: } a, b, c$$

bias/offset

**model:**  $P_{abc}^{(k)}(\delta_k | a, b, c) = \frac{1}{2c|a|^k} \Theta\left(c - \left|\frac{\delta_k}{a^k} - b\right|\right)$



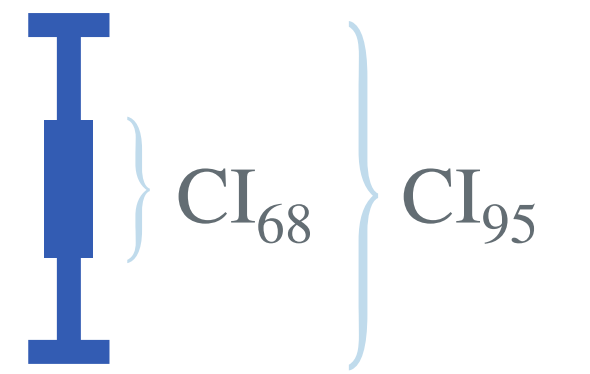
**priors:**  $P_0(a, b, c) = P_0(a) P_0(b, c)$

$$P_0(a) = \frac{1}{2} (1 + \omega) (1 - |a|)^\omega \Theta(1 - |a|)$$

$\rightsquigarrow$  support:  $[-1, +1]$  (alternating ✓)

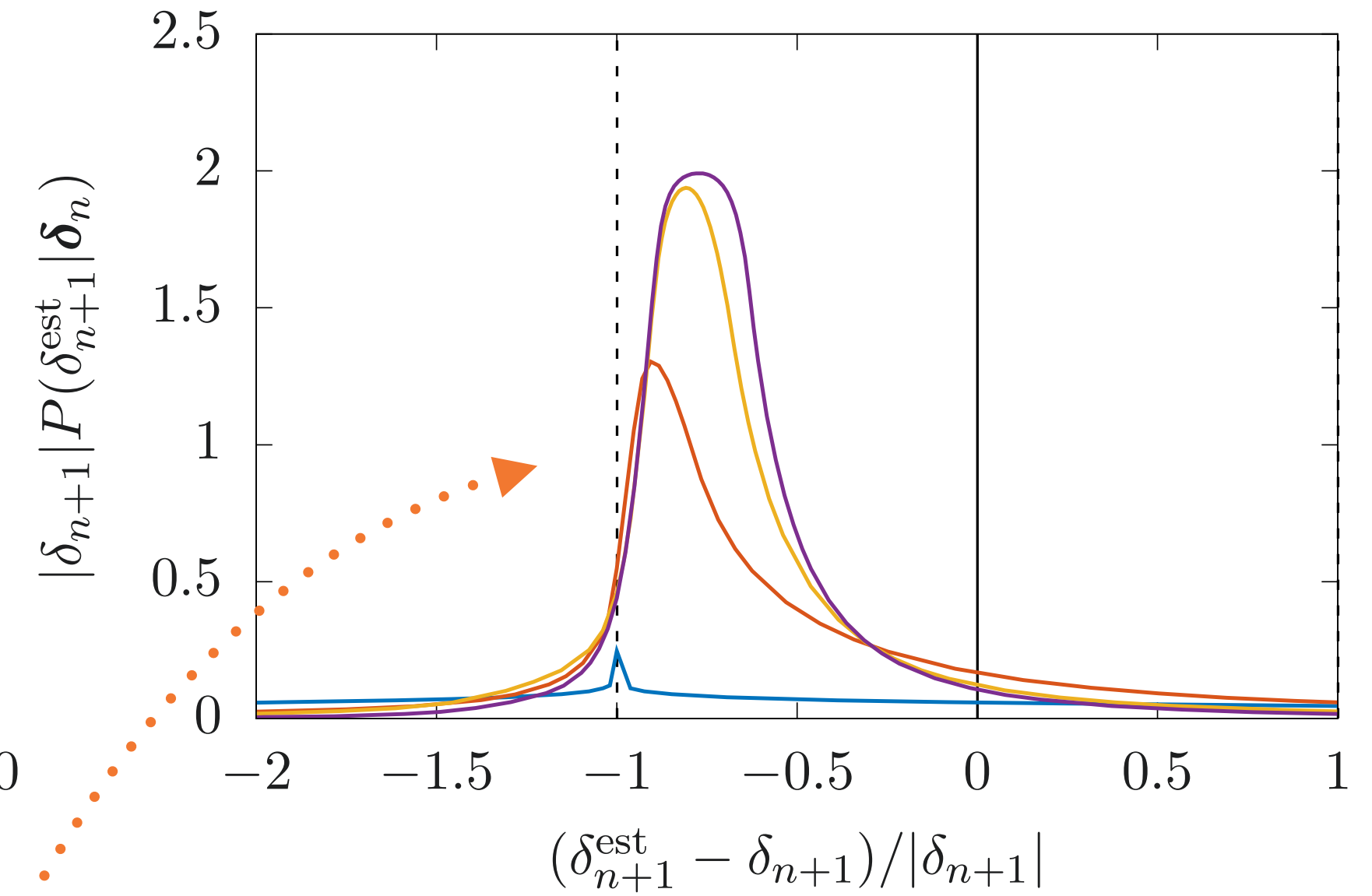
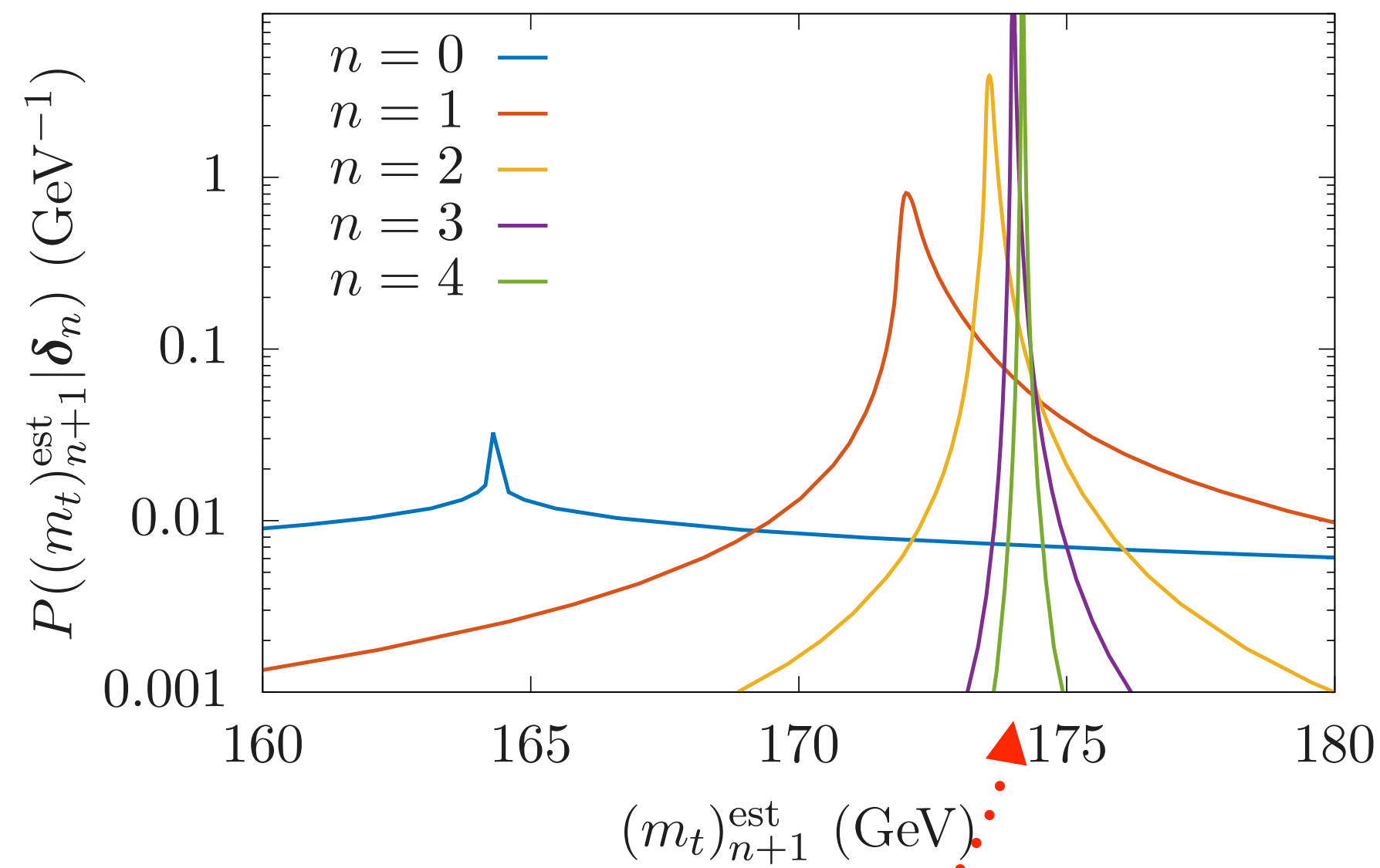
$$P_0(b, c) = \frac{\varepsilon \eta^\varepsilon}{c^{1+\varepsilon}} \Theta(c - \eta) \frac{1}{2\xi c} \Theta(\xi c - b)$$

# A REAL-WORLD EXAMPLE — $m_t$ (OS $\leftrightarrow$ $\overline{\text{MS}}$ )



$$m_t = \frac{Z_m^{\text{OS}}}{Z_m^{\overline{\text{MS}}}(\mu_R)} \bar{m}_t(\mu_R)$$

$$= \sum_k m_t^{(k)}$$

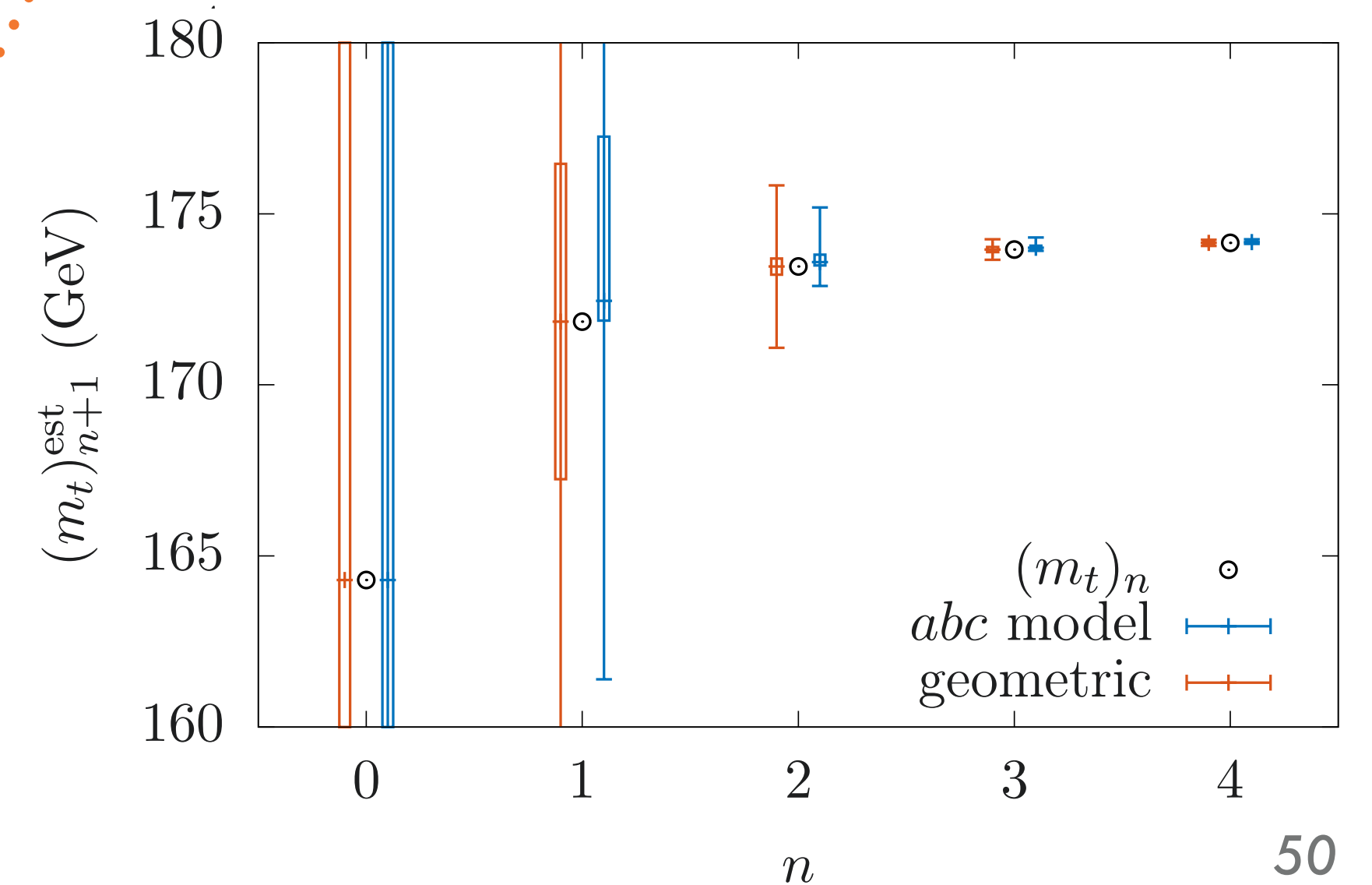


strongly peaked  $n \nearrow$

positive corrections  
anticipated

estimate for  $m_t - (m_t)_4$

- $\text{CI}_{68} = [0.008, 0.046] \text{ GeV}$
- $\text{CI}_{95} = [-0.027, 0.112] \text{ GeV}$



# WHAT TO DO WITH THE THE SCALE $\mu$ ?

•  $\forall \mu \rightsquigarrow P(\delta_3 | \delta_0, \delta_1, \delta_2; \mu)$

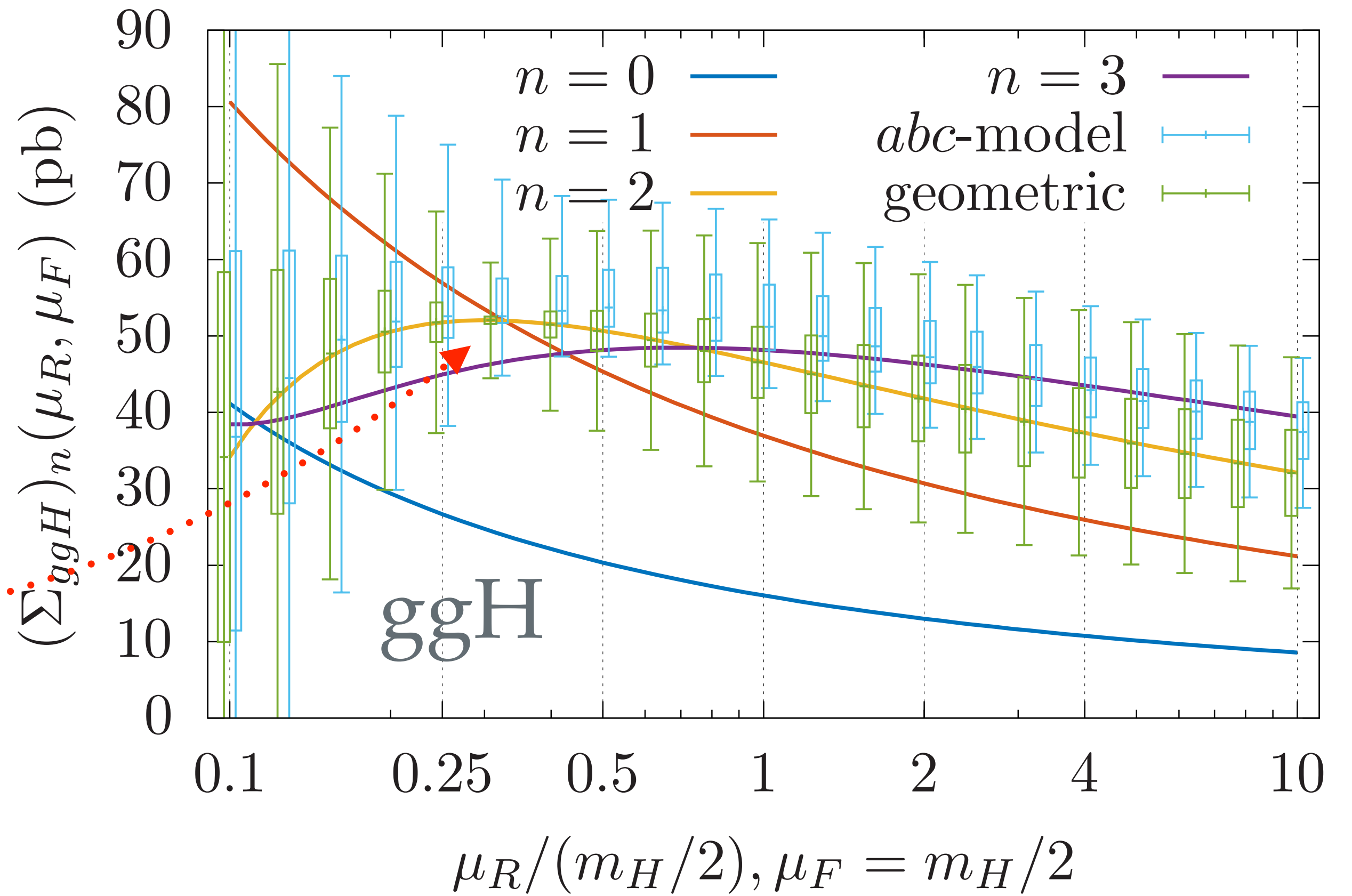
▶  $CI_{68/95}$   (geo)  (*abc*)

• geo

- ▶ always entered around **NNLO**
- ▶ very narrow peak 

• *abc*

- ▶  $\mu/\mu_0 \gtrsim 1 \rightsquigarrow$  anticipate pos. **N3LO**
- ▶  $\mu/\mu_0 \lesssim 1 \rightsquigarrow$  bias slowly disappears



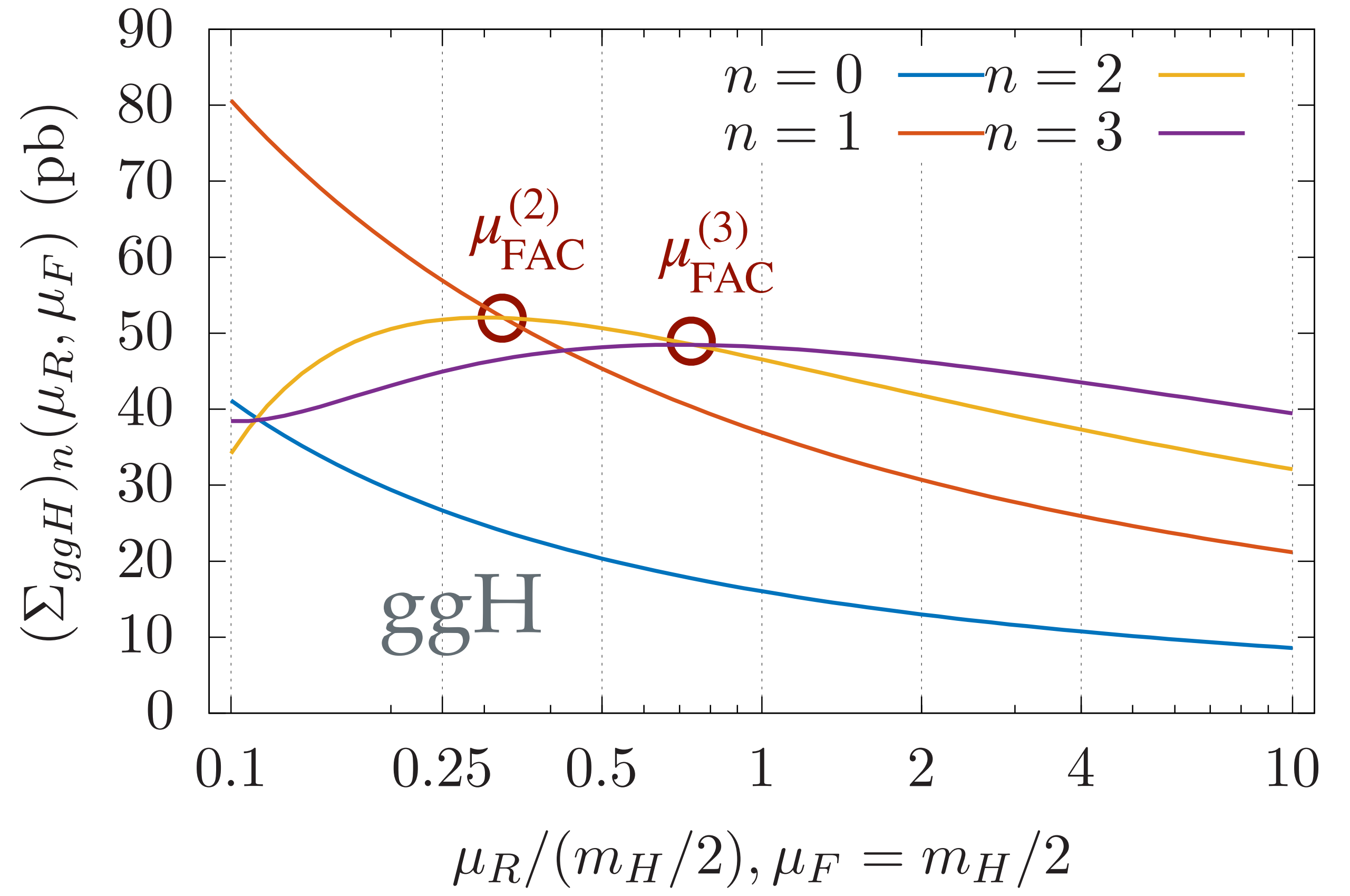
# WHAT TO DO WITH THE THE SCALE $\mu$ ?

•  $\forall \mu \rightsquigarrow P(\delta_3 | \delta_0, \delta_1, \delta_2; \mu)$

▸  $CI_{68/95}$   (geo)  (abc)

• two options:

1. invoke some *principle* to pick the “*optimal*” scale
  - FAC, PMS, PMC, ...



**Fastest Apparent Convergence**  
 $\Sigma_n(\mu_{FAC}) = \Sigma_{n-1}(\mu_{FAC})$

depends on order  
 might not be unique

# WHAT TO DO WITH THE THE SCALE $\mu$ ?

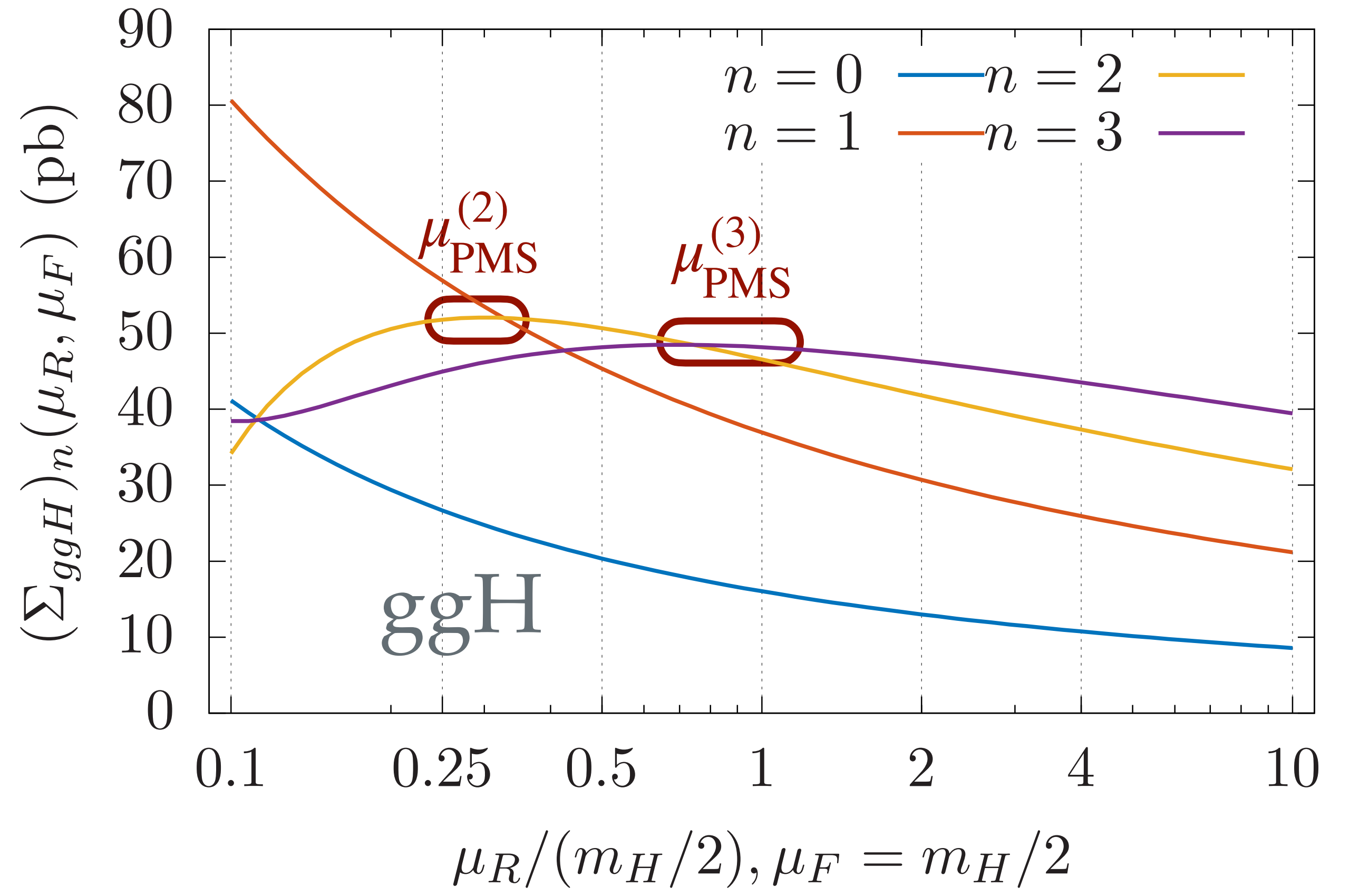
•  $\forall \mu \rightsquigarrow P(\delta_3 | \delta_0, \delta_1, \delta_2; \mu)$

▶  $CI_{68/95}$   (geo)  (abc)

• two options:

1. invoke some *principle* to pick the “*optimal*” scale

▶ FAC, PMS, PMC, ...



**Principle of Minimal Sensitivity**  
 $\frac{\partial}{\partial \mu} \Sigma_n(\mu) \Big|_{\mu_{\text{PMS}}} = 0$

depends on order  
 might not be unique

# WHAT TO DO WITH THE THE SCALE $\mu$ ?

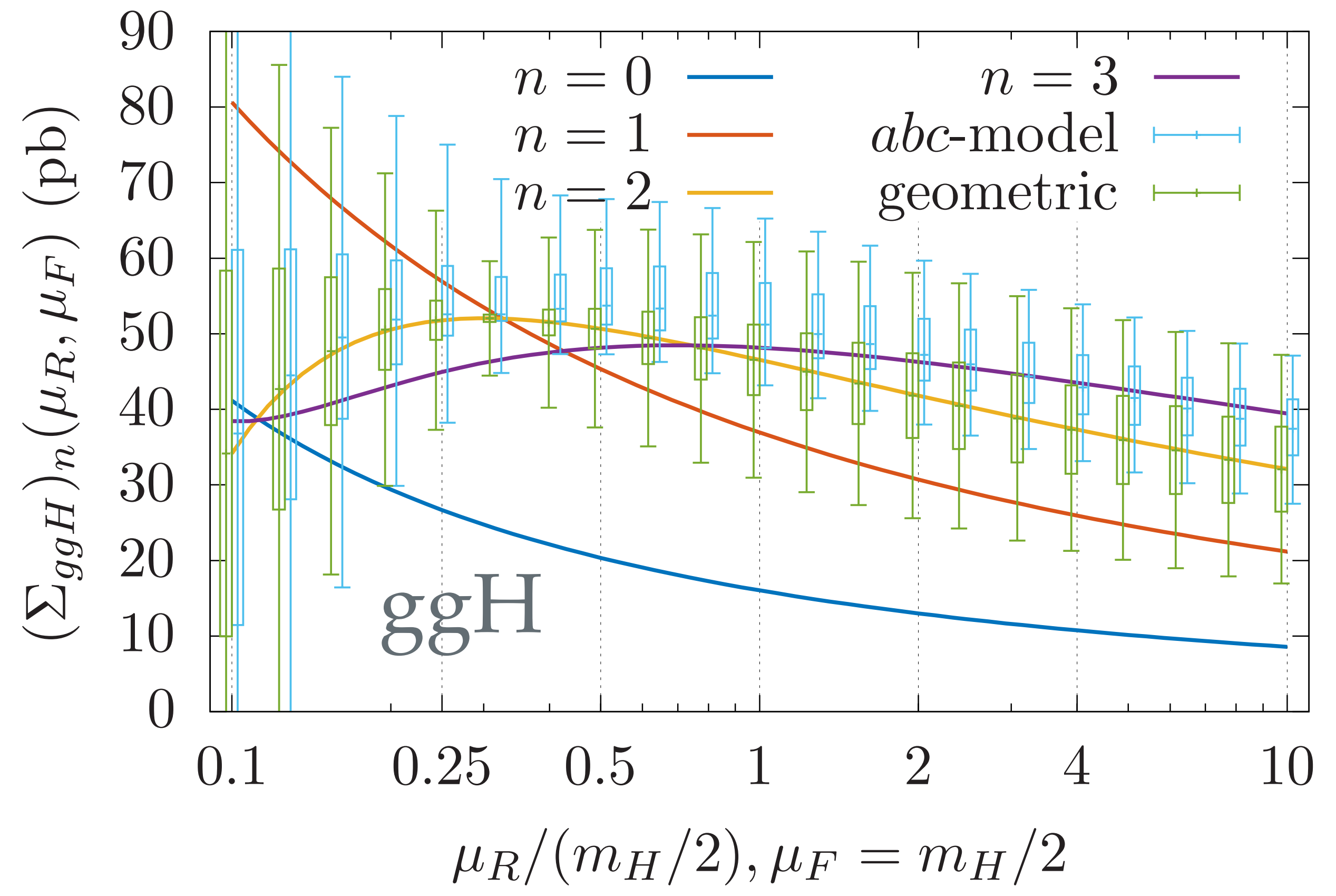
•  $\forall \mu \rightsquigarrow P(\delta_3 | \delta_0, \delta_1, \delta_2; \mu)$

▶  $CI_{68/95}$   (geo)  (abc)

• two options:

1. invoke some principle to pick the “optimal” scale
  - ▶ FAC, PMS, PMC, ...
2. combine different  $P(\delta_{n+1} | \delta_n; \mu)$

pursued in the following



# PRESCRIPTIONS FOR SCALES

---

## Scale Marginalisation (sm):

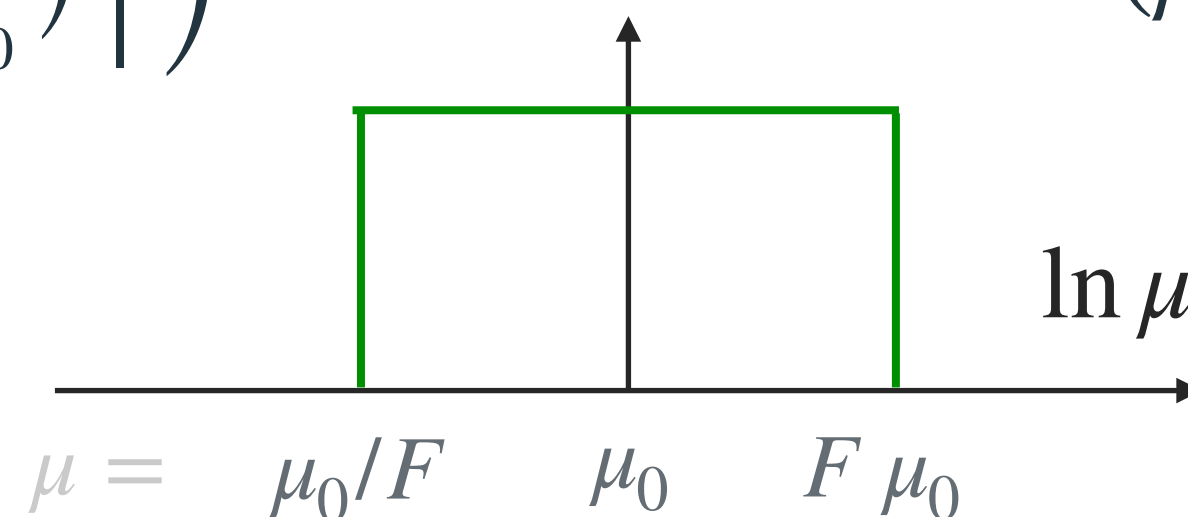
[Bonvini '20]

- treat  $\mu$  as a hidden model parameter & marginalise over it:

$$\begin{aligned} P_{\text{sm}}(\delta_{n+1} | \delta_n) &= \int d\mu P(\delta_{n+1}, \mu | \delta_n) \\ &= \int d\mu P(\delta_{n+1} | \delta_n; \mu) P(\mu | \delta_n) \end{aligned}$$

- $P(\mu | \delta_n) \propto P(\delta_n; \mu) P_0(\mu)$  with prior:

$$P_0(\mu) = \frac{1}{2\mu \ln F} \Theta\left(\ln F - \left|\ln\left(\frac{\mu}{\mu_0}\right)\right|\right)$$



## Scale Average (sa):

[Duhr, AH, Mazeliauskas, Szafron '21]

- $\mu$  has no probabilistic interpretation  $\rightsquigarrow$  average over it:

$$P_{\text{sa}}(\delta_{n+1} | \delta_n) = \int d\mu w(\mu) P(\delta_{n+1} | \delta_n; \mu)$$

- weight function:

$$w(\mu) = \frac{1}{2\mu \ln F} \Theta\left(\ln F - \left|\ln\left(\frac{\mu}{\mu_0}\right)\right|\right)$$

# PEAK OF THE DISTRIBUTIONS\*

[Duhr, AH, Mazeliauskas, Szafron '21]

## Scale Marginalisation (sm):

- if  $\mu_{\text{FAC}} \in [\mu_0/F, F\mu_0]$  then  $P_{\text{sm}}(\delta_{n+1} | \delta_n)$  peaks at  $\Sigma_n(\mu_{\text{FAC}})$ 
  - $P(\delta_n | \mu)$  dominated by ( $k = n$ ) term
  - symmetric model
    - ↪  $\delta_n(\mu) = 0$  enhanced

## Scale Average (sa):

- if  $\mu_{\text{PMS}} \in [\mu_0/F, F\mu_0]$  then  $P_{\text{sa}}(\delta_{n+1} | \delta_n)$  peaks at  $\Sigma_n(\mu_{\text{PMS}})$ 
  - overlap between  $P(\delta_{n+1} | \delta_n; \mu)$  enhanced at stationary point
    - ↪  $\Sigma'_n(\mu_{\text{PMS}}) \approx 0$

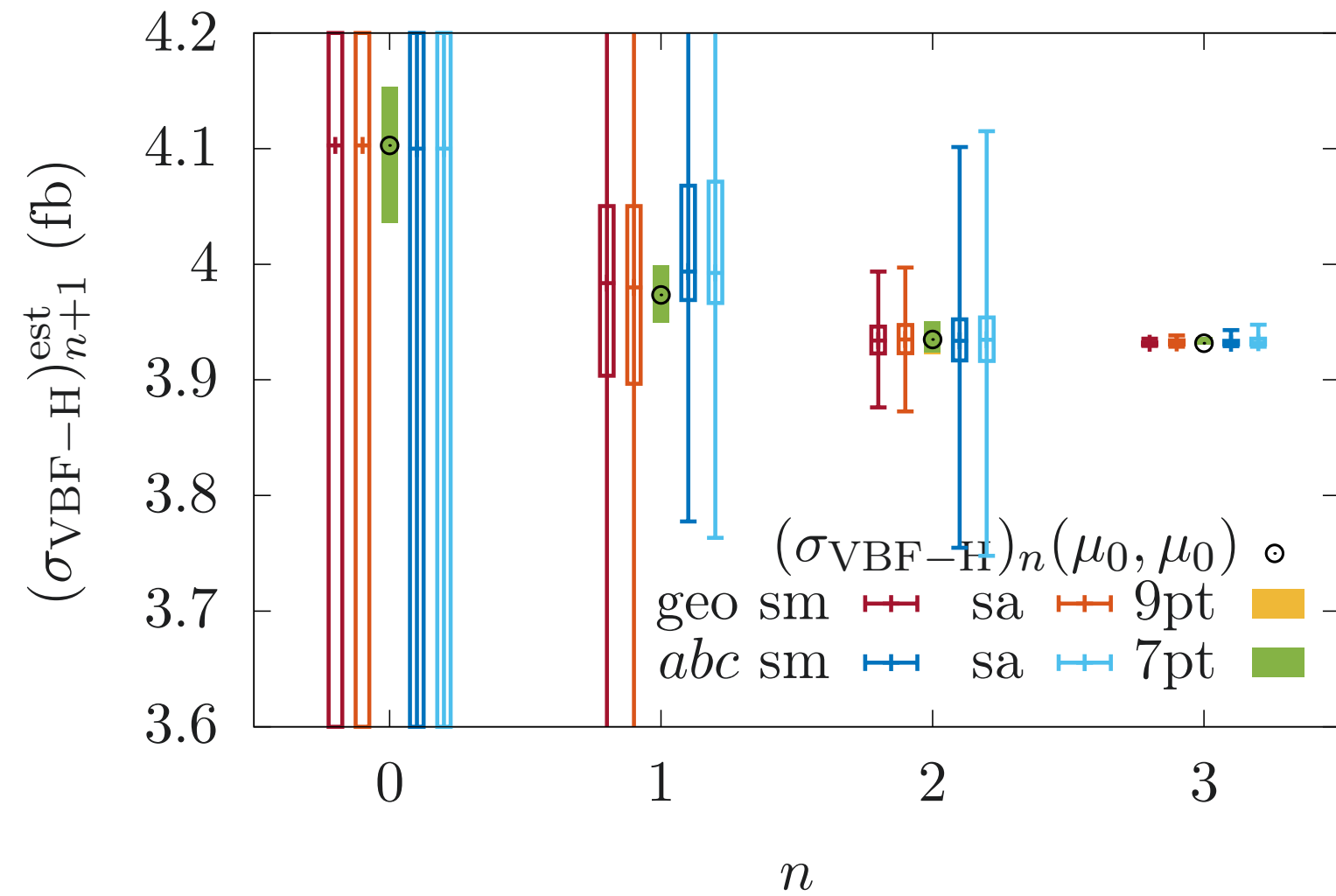
Choice of how to interpret the scale has consequences for predictions!

\* for symmetric models, a convergent series, and reasonable assumptions



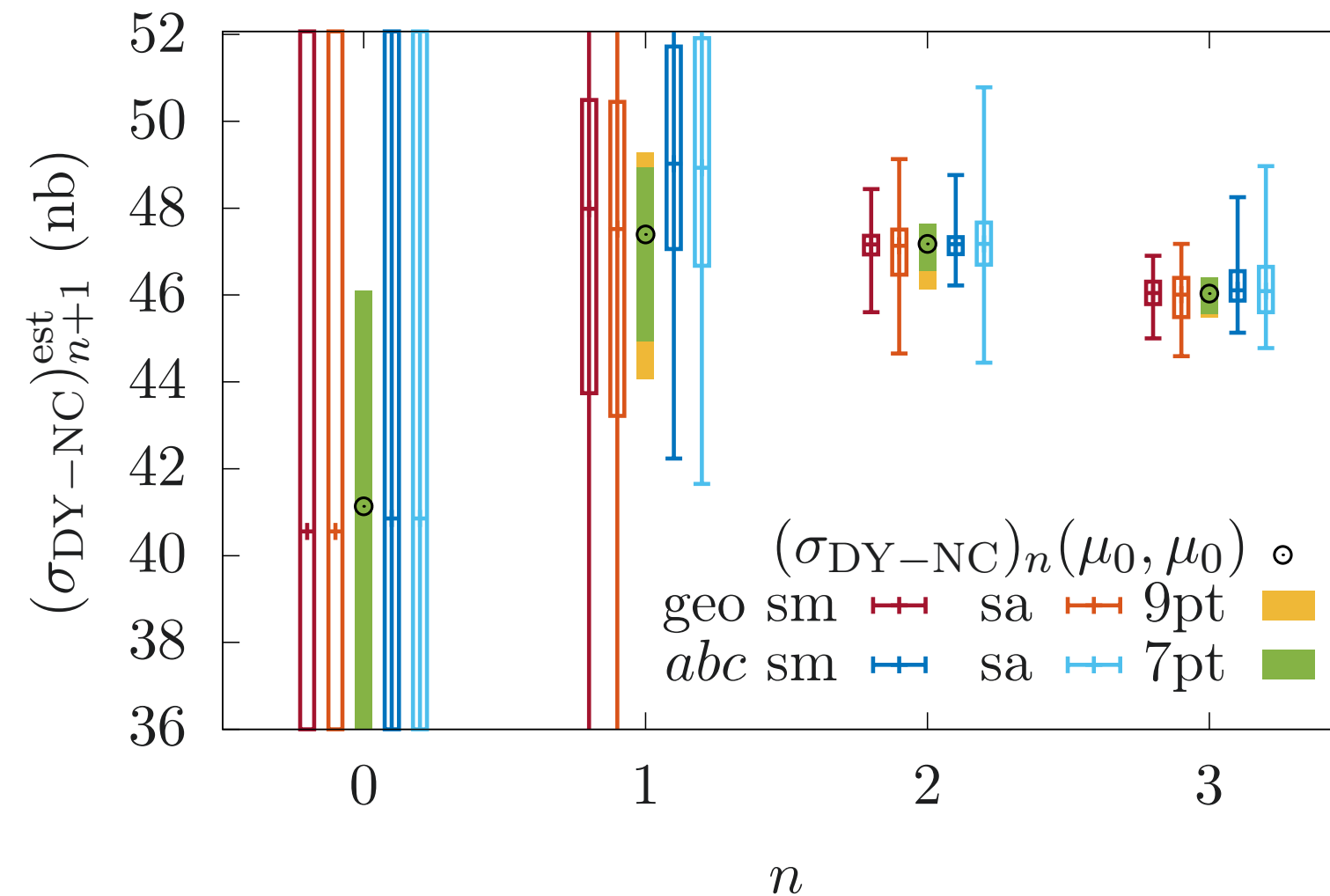
# INCLUSIVE CROSS SECTIONS UP TO N3LO

## VBF-H



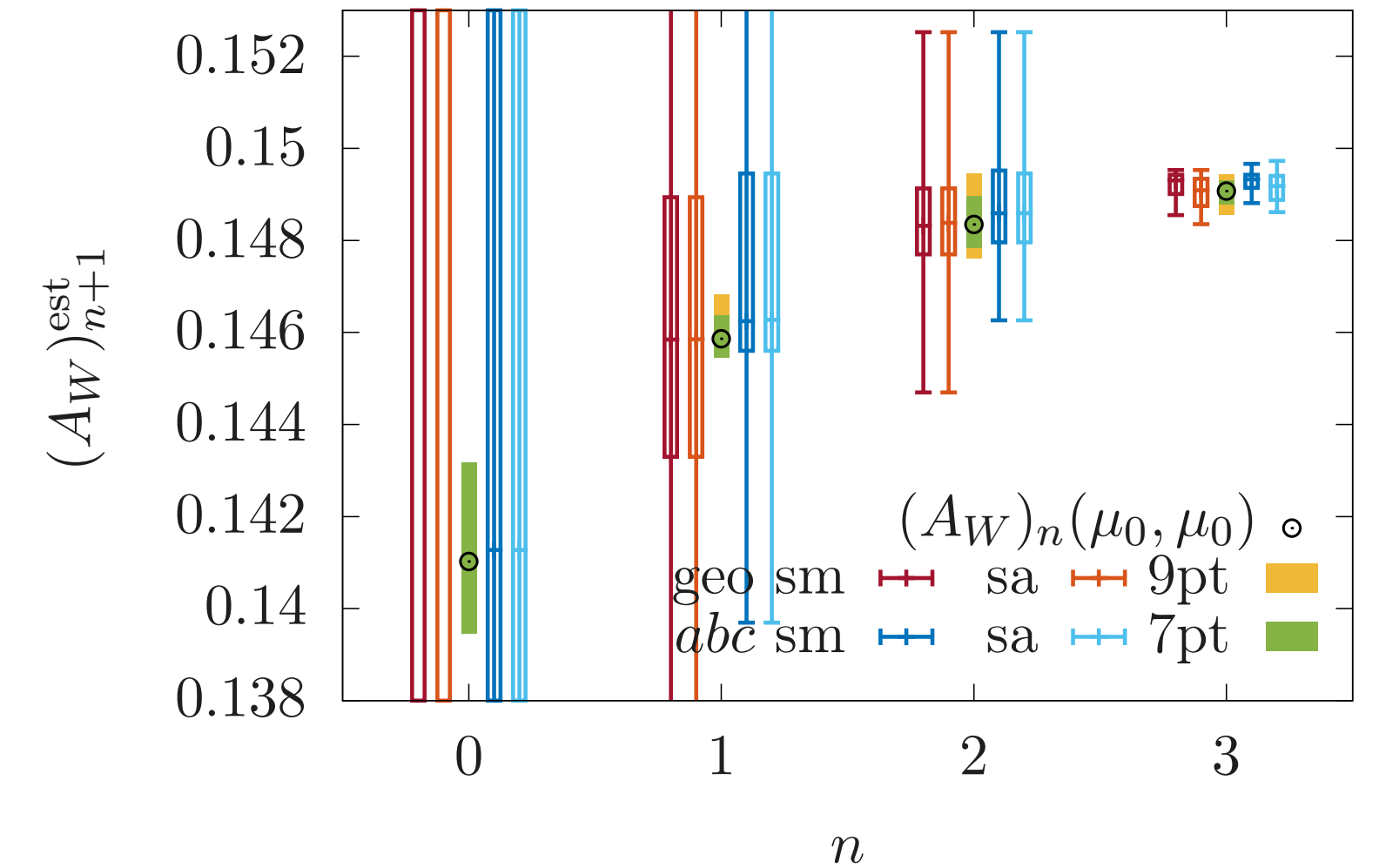
- $n < 2$ :  $\text{CI}_{68}$  bigger than 9pt
- $\delta_1 < 0 \rightsquigarrow abc$  alternating
- $n > 2$ : all prescriptions similar

## DY-NC



- $\delta_3$  is large and outside of 9pt!
- similar unc.:  $sa \simeq 9pt$
- $n = 2$ :  $sm \ll$  others ( $\mu_{\text{FAC}}$ )
- $n = 3$ : all prescriptions similar

$$A_W = \frac{W^+ - W^-}{W^+ + W^-}$$



- large cancellations in the ratio
- $n < 2$ : 9pt performs poorly
- $(A_W)_n \nearrow$  (anticipated by  $abc$ )
- size:  $abc \lesssim$  others

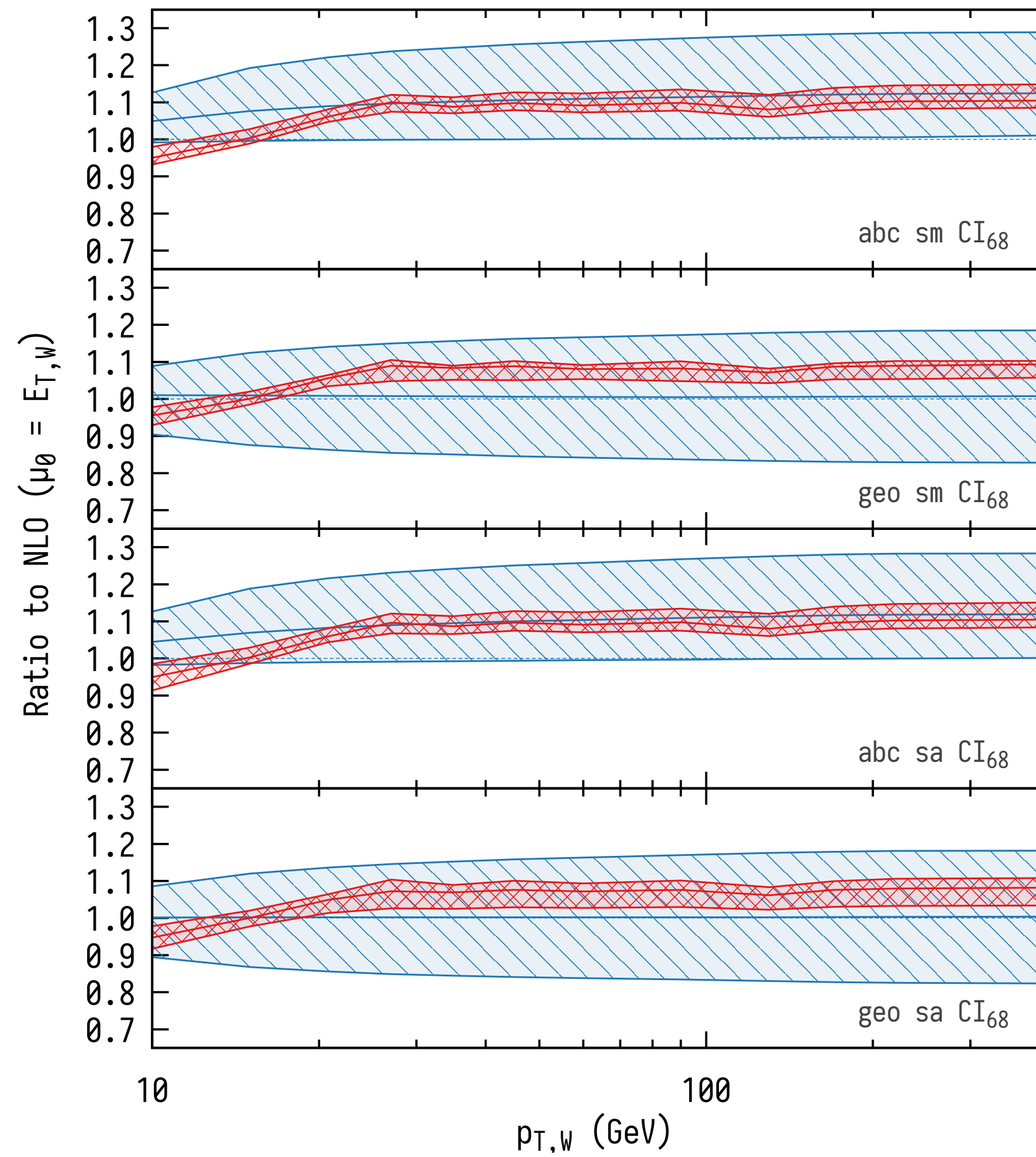
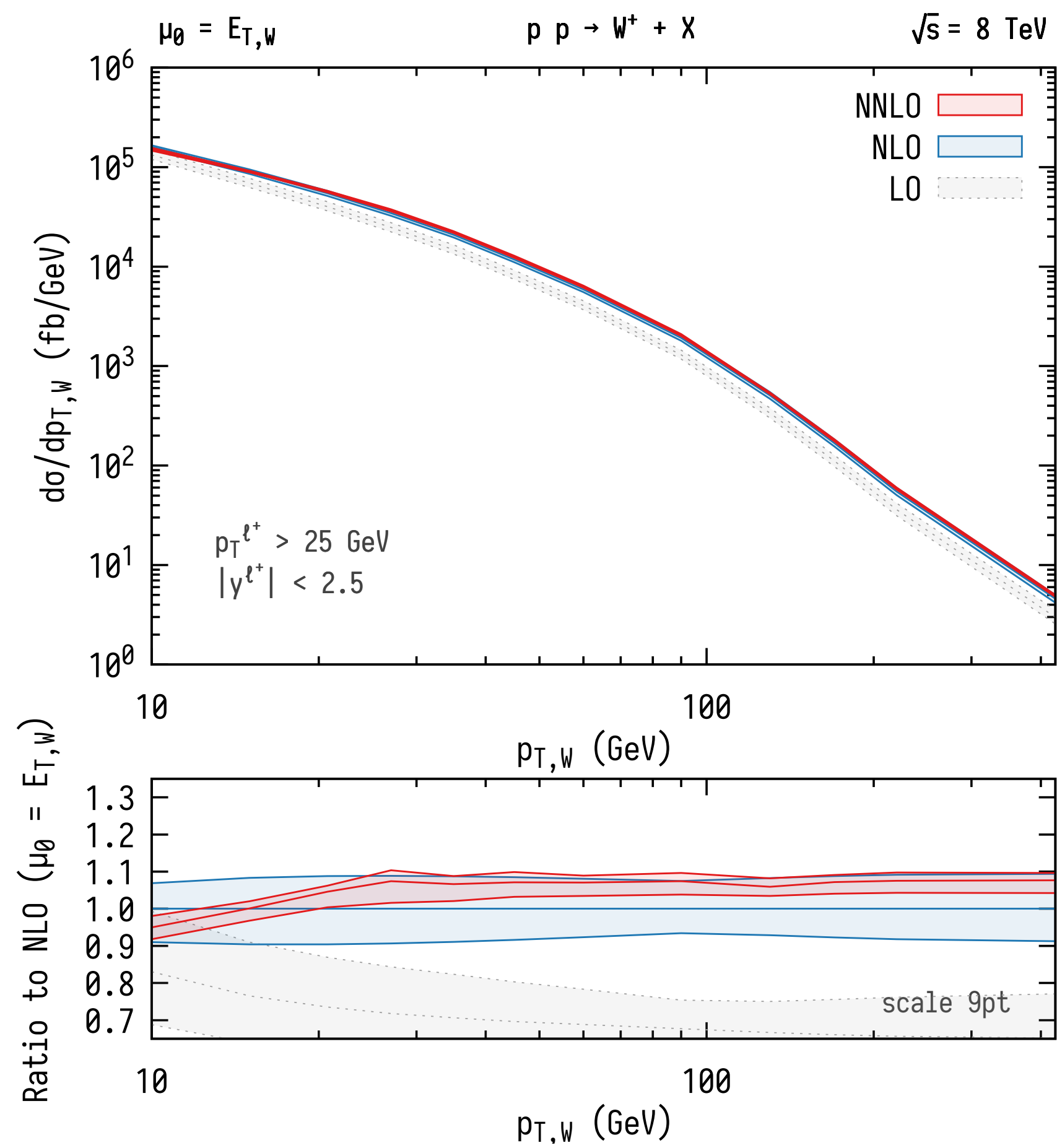
overall: not radically different estimates for  $\Delta_{\text{MHO}}$

# DIFFERENTIAL DISTRIBUTIONS

---

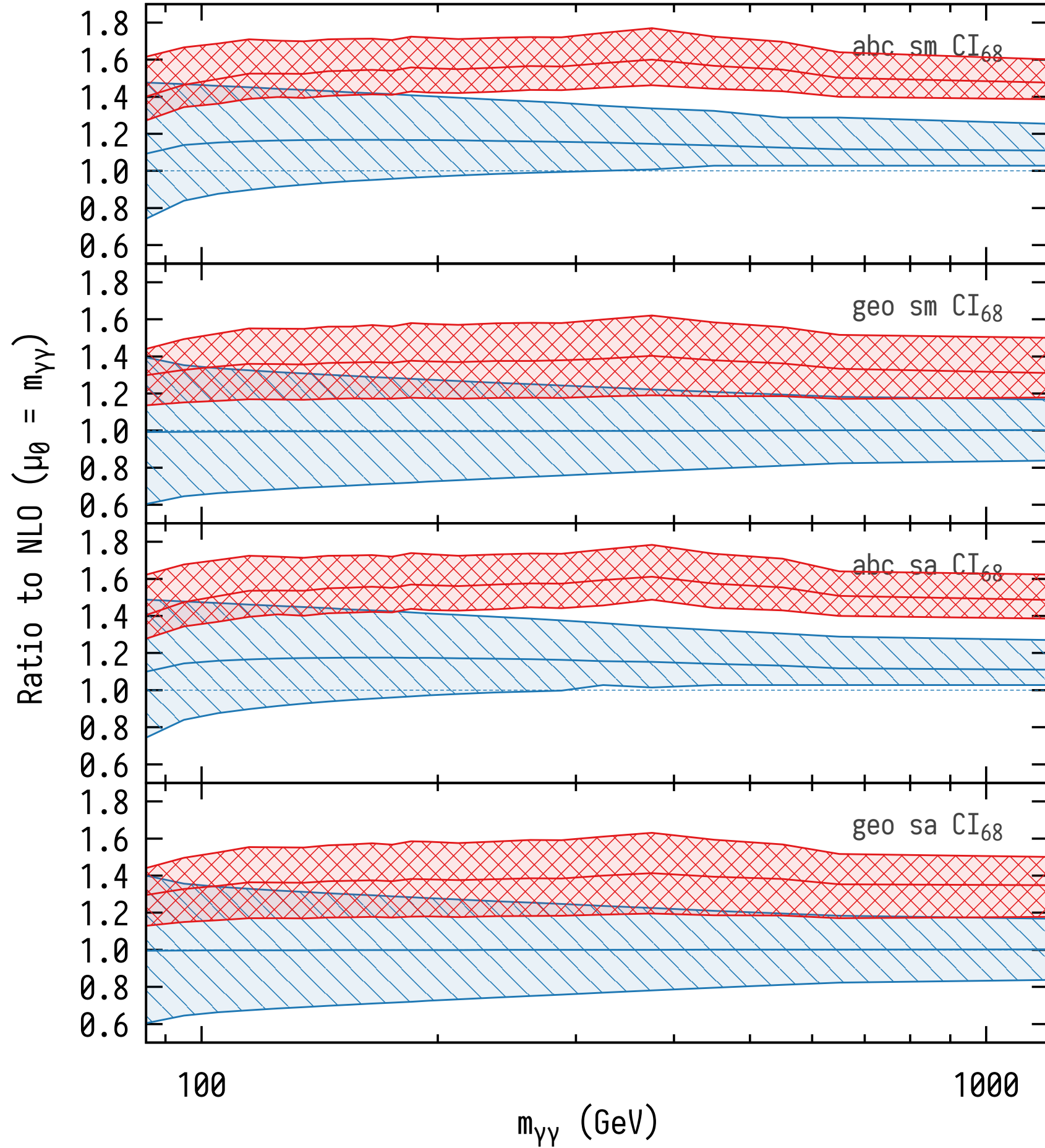
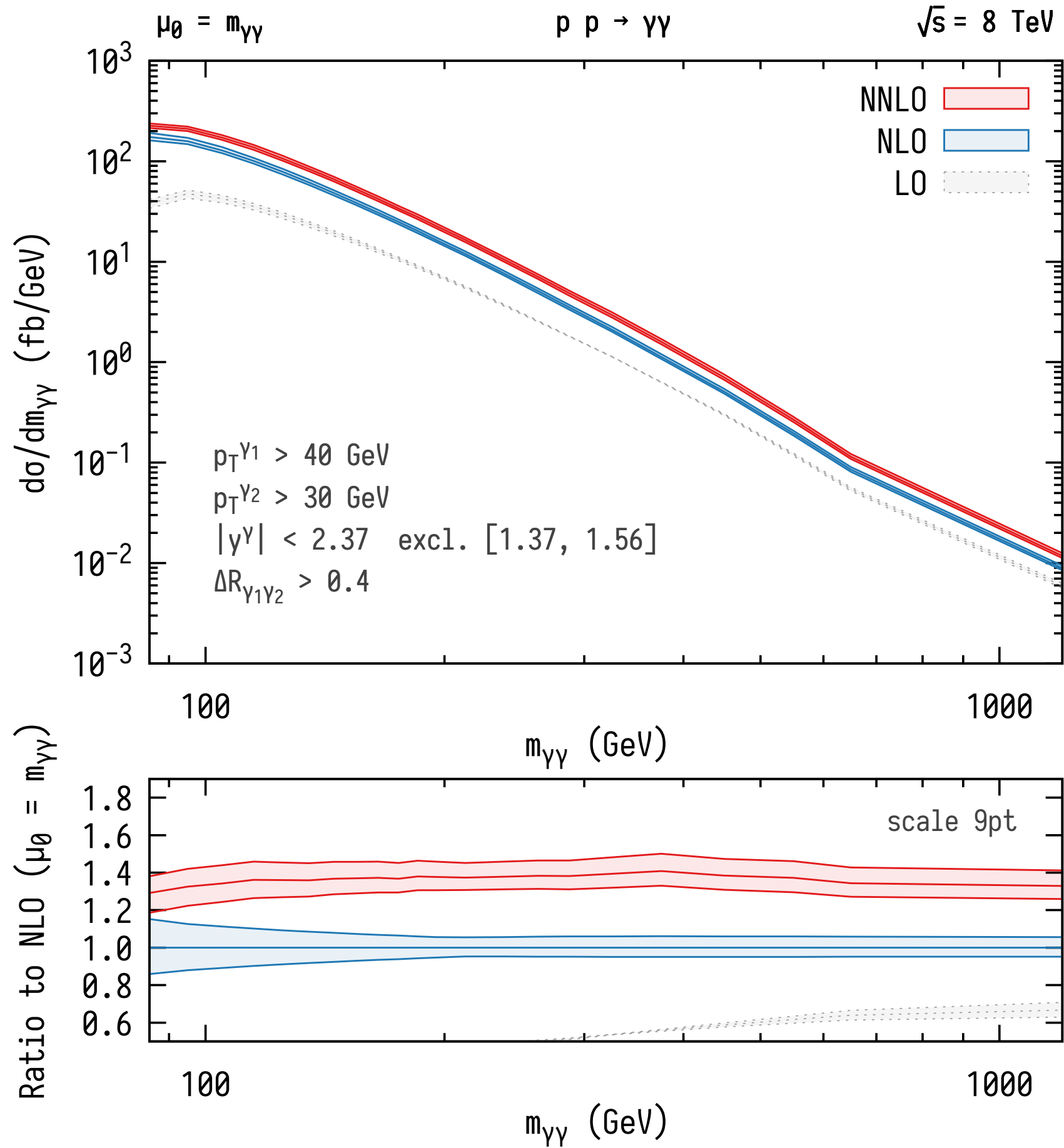
- Bayesian approach also applicable to distributions
  - ↔ treat each bin individually ↔ **will not include correlations!**
- new challenges
  - no longer “easy” to identify an appropriate hard scale  $\mu_0$  (up to rescaling)
    - ↔ inclusive ggH:  $M_H$  vs.  $\frac{1}{2} M_H$ ? Just let the model figure it out.
  - differential distributions can probe different kinematic regimes
    - ↔ dynamical scale choice ↔ *many choices!*
    - ↔ e.g. in jet production:  $p_T^j$ ,  $p_T^{j_1}$ ,  $\langle p_T^j \rangle_{\text{avg}}$ ,  $H_T \equiv \sum_{i \in \text{jets}} p_T^i$ ,  $\hat{H}_T \equiv \sum_{i \in \text{partons}} p_T^i$ , ...
  - re-cycling via quadrature limited ↔ **ideally interpolation grids**

# W-BOSON + JET PRODUCTION



- $n < 2$ :
  - $CI_{68}$  bigger than 9pt
  - *abc* captures pos. shift
- $n = 2$ :
  - almost identical bands
  - $\Delta_{MHO}$  very robust
- sm vs. sa
  - almost identical CI

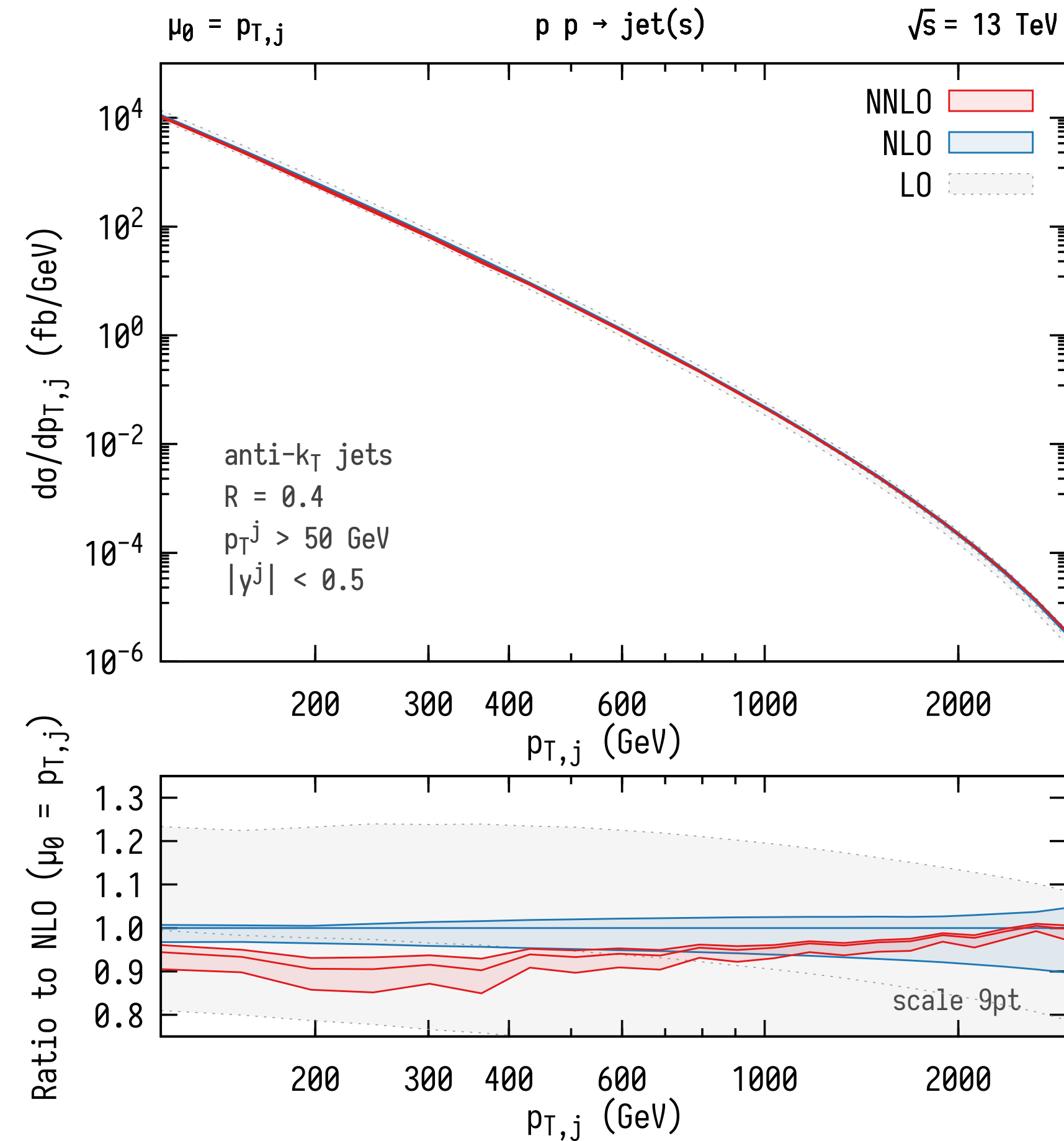
# DI-PHOTON PRODUCTION



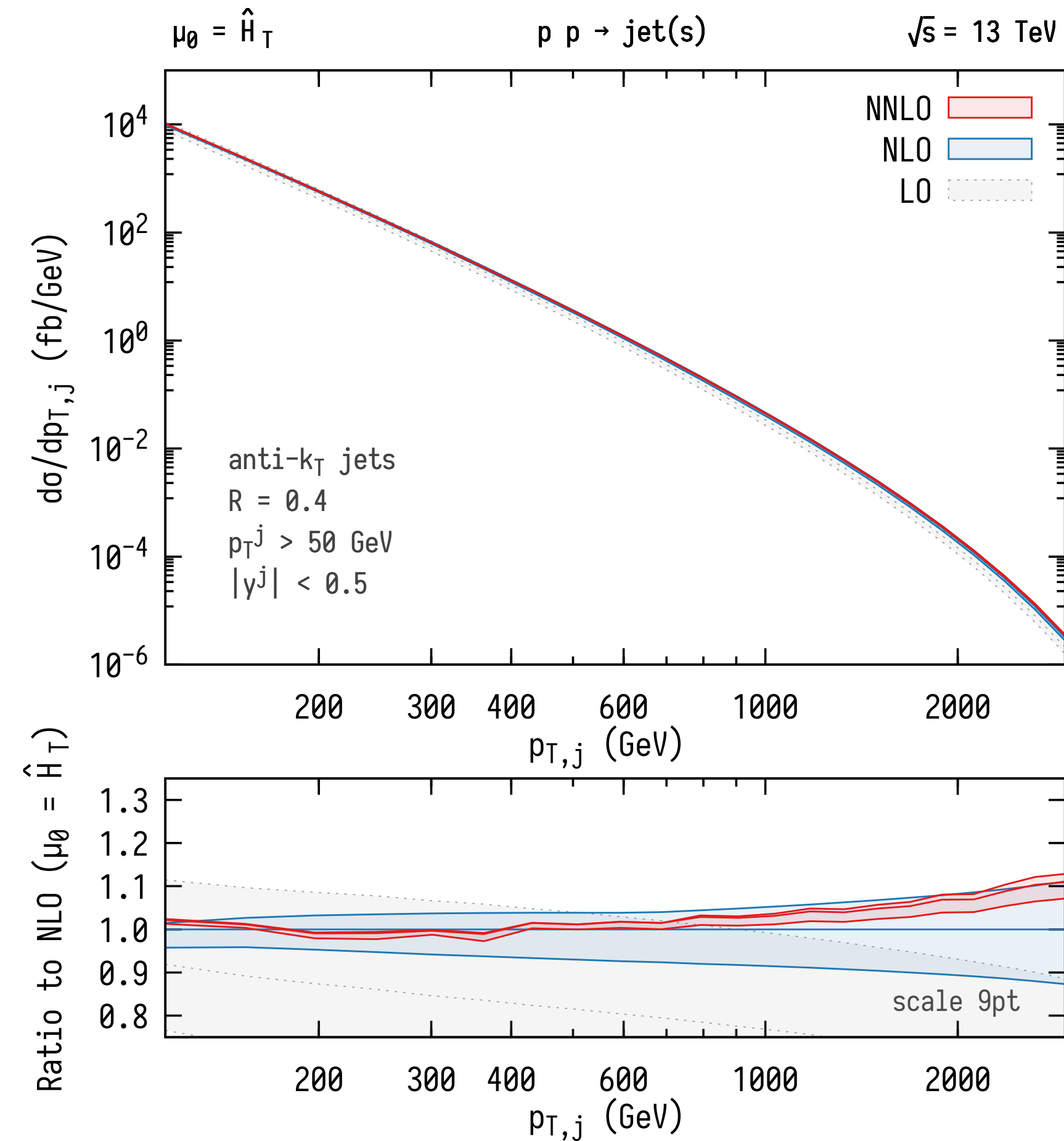
- example where 9pt fails
  - large corrections
  - $\Delta_{\text{MHO}}^{\text{NNLO}} \gtrsim \Delta_{\text{MHO}}^{\text{NLO}}$
  - no sign of convergence
- $n < 2$ :
  - $CI_{68} \sim 2-3 \times 9\text{pt}$
- $n = 2$ :
  - marginal overlap for geo
  - differences in *size & position*
  - ideally N3LO for robust  $\Delta_{\text{MHO}}$
- $\text{sm} \simeq \text{sa}$ 
  - large corrections prohibit FAC points

# THE PROBLEM WITH JETS...

$\mu_0 = p_T^j$ : infrared sensitivity



$\mu_0 = \hat{H}_T$ : recommendation [Currie et al. '18]

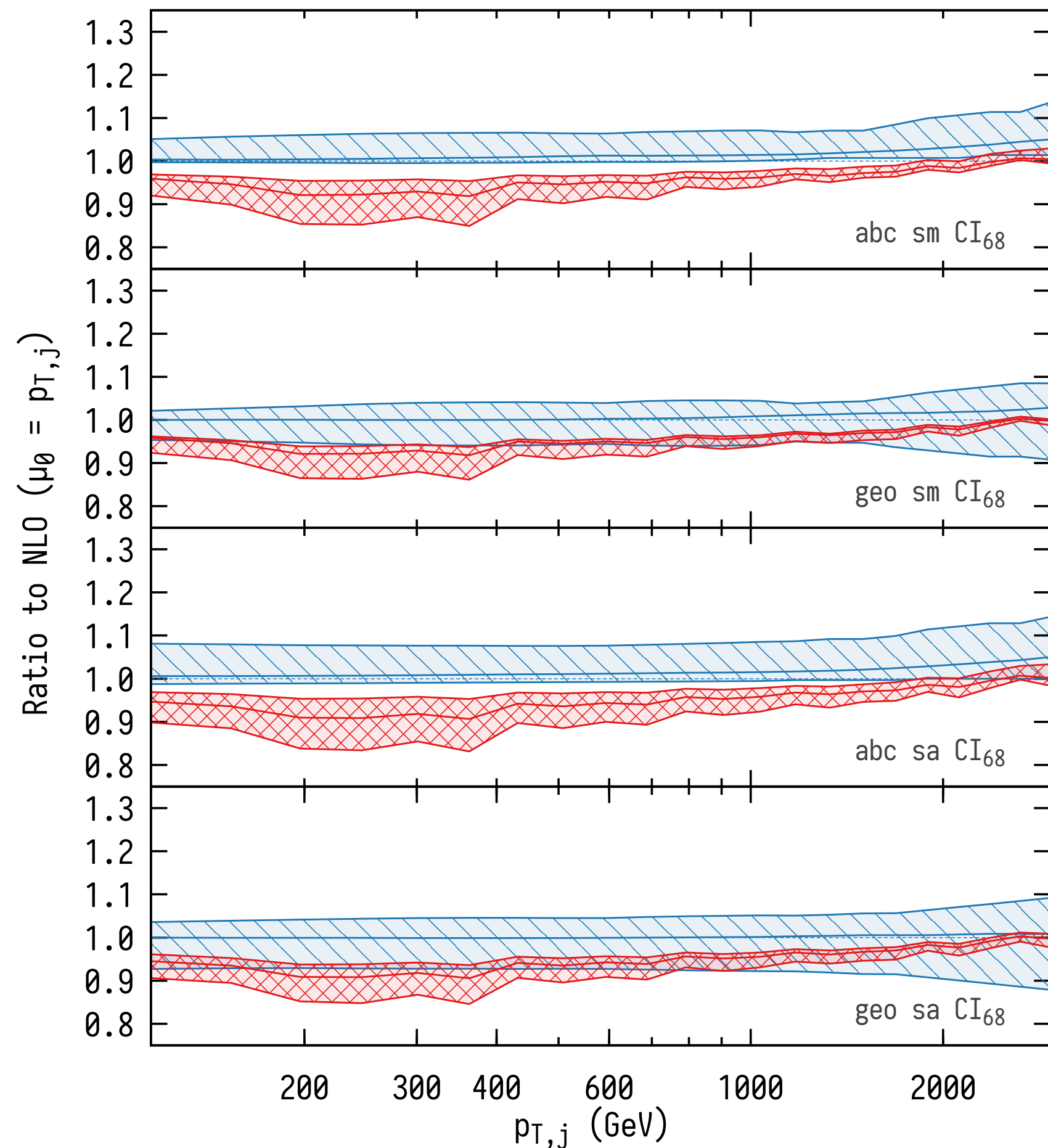


# THE PROBLEM WITH JETS... PERSISTS

non-trivial change of dynamical scales  
cannot be captured by a simple re-scaling

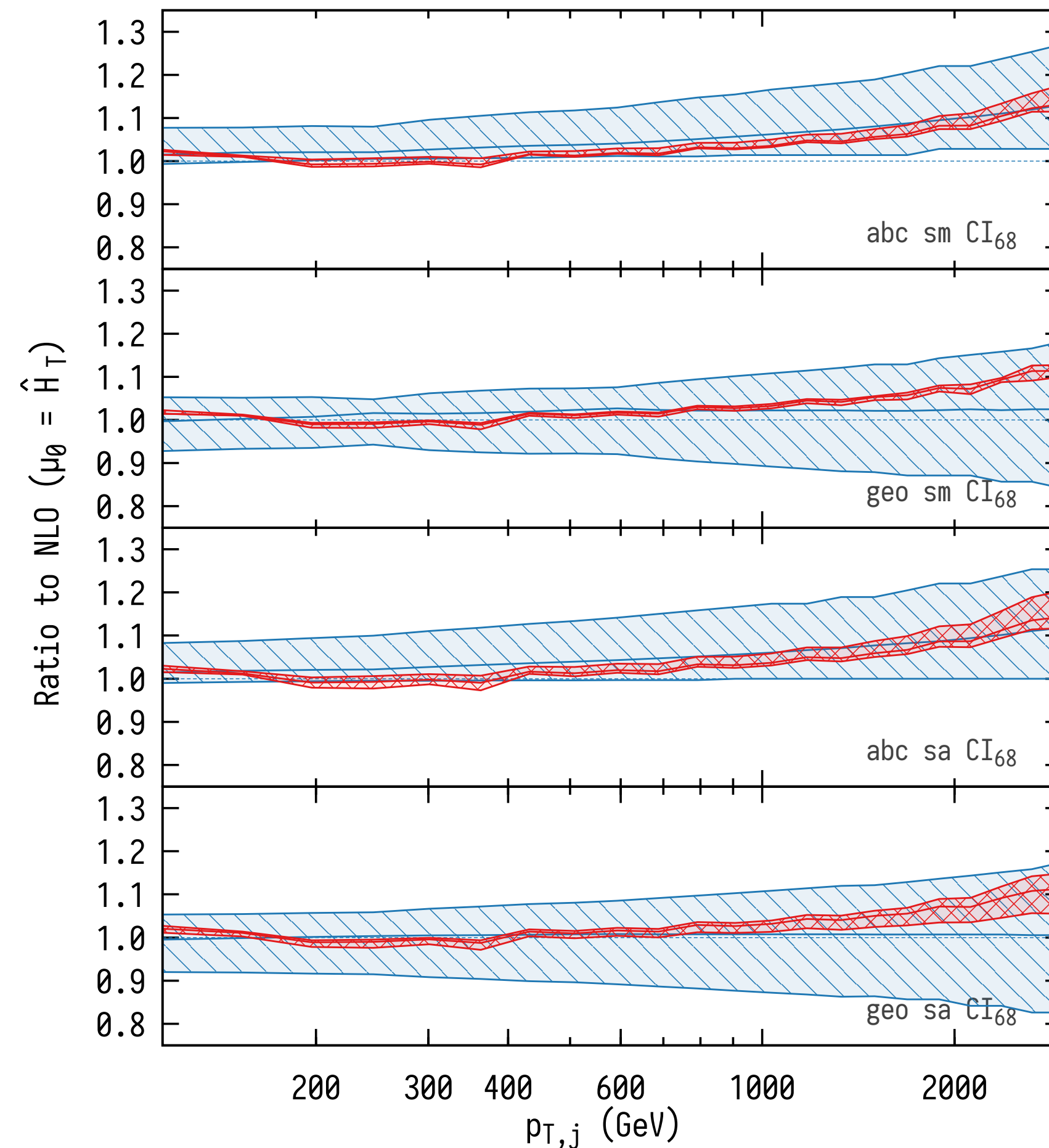
$\mu_0 = p_T^j$ : infrared sensitivity

$\mu_0 = \hat{H}_T$ : recommendation [Currie et al. '18]



larger NLO  
 $\Delta_{\text{MHO}}$

barely any  
difference  
at NNLO



*abc* captures  
positive  
corrections

# WORK IN PROGRESS — CORRELATIONS

---

- idea: if two bins show similar (opposite) perturbative behaviour
  - ↪ two bins should be partially (anti-)correlated.

- we want: joint probability distribution  $P(x, y)$  for two bins  $x$  &  $y$ 
  - ↪ preserve projections for compatibility:

$$P(x) = \int dy P(x, y) = \int dz P(x, z)$$

- ↪ hidden parameter  $-1 < c < +1$  to smoothly implement the correlation
- possibilities: algorithmic “earth movers distance”; map  $P(x)$  onto  $P(y)$ , ...
  - ↪ can be done much simpler

# WORK IN PROGRESS – CORRELATION MODEL IN miho

- projections of multi-dim. Gaussians (+ correlation matrix) are again Gaussian  
 $\hookrightarrow$  map  $P_i$  onto Gaussians, implement correlations, map back

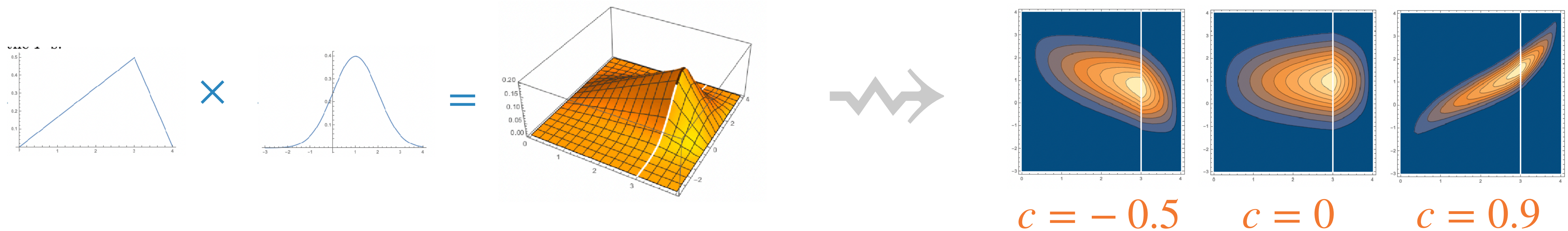
$$\begin{aligned}
 P(x, y) &= P_1(x)P_2(y) \\
 &\times \left. \frac{d\Phi^{-1}(\alpha)}{d\alpha} \right|_{\alpha=\Sigma_1(x)} \left. \frac{d\Phi^{-1}(\beta)}{d\beta} \right|_{\beta=\Sigma_2(y)} \\
 &\times \frac{1}{2\pi\sqrt{1-c^2}} \exp\left(-\frac{1}{2(1-c^2)}[\xi(x)^2 + \eta(y)^2 - c2\xi(x)\eta(y)]\right)
 \end{aligned}$$

$$\Sigma_i(x) = \int_{-\infty}^x dx' P_i(x')$$

$$\Phi^{-1}(p) = \sqrt{2}\text{Erf}^{-1}(-1 + 2p)$$

$$\xi(x) = \Phi^{-1}(\Sigma_1(x))$$

$$\eta(y) = \Phi^{-1}(\Sigma_2(y))$$



use inference to constrain  $c$



# CONCLUSIONS & OUTLOOK PART 3

---

- Bayesian inference is a powerful framework to estimate  $\Delta_{\text{MHO}}$ 
  - statistical interpretation  $\Leftrightarrow P(\delta_{n+1} | \delta_n)$
  - exposes our *assumptions & biases* clearly  $\Leftrightarrow$  model & priors
  - but: it is not more reliable than scale variation  $\rightsquigarrow$  careful analysis required
- typically for  $n < 2$ :  $\text{CI}_{68} > 9\text{pt}$ ;  $n \geq 2$ :  $\text{CI}_{68} \simeq 9\text{pt}$
- *public code*: ミホ (miho)  $\rightsquigarrow$  <https://github.com/aykhuss/miho>
- future directions
  - correlations (PDF fits & data interpretation)
  - marginalisation over models, ...

# CONCLUSIONS & OUTLOOK PART 3

---

- Bayesian inference is a powerful framework to estimate  $\Delta_{\text{MHO}}$ 
  - statistical interpretation  $\Leftrightarrow P(\delta_{n+1} | \delta_n)$
  - exposes our *assumptions & biases* clearly  $\Leftrightarrow$  model & priors
  - but: it is not more reliable than scale variation  $\rightsquigarrow$  careful analysis required
- typically for  $n < 2$ :  $\text{CI}_{68} > 9\text{pt}$ ;  $n \geq 2$ :  $\text{CI}_{68} \simeq 9\text{pt}$
- *public code*: ミホ (miho)  $\rightsquigarrow$  <https://github.com/aykhuss/miho>
- future directions
  - correlations (PDF fits & data interpretation)
  - marginalisation over models, ...

Thank you!