

Introduction to Accelerator Physics

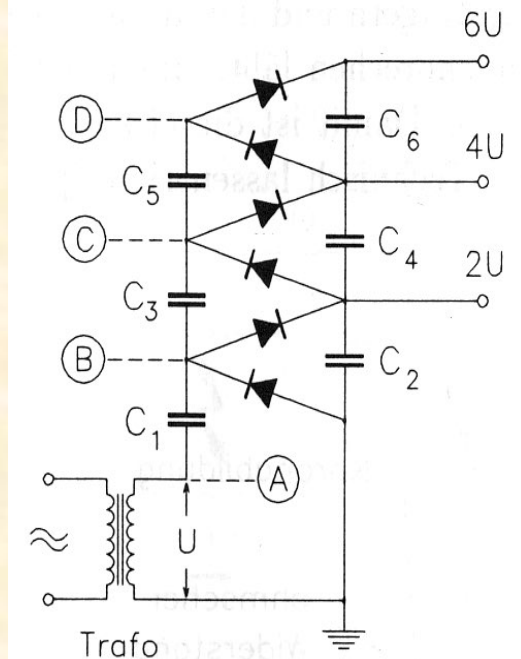
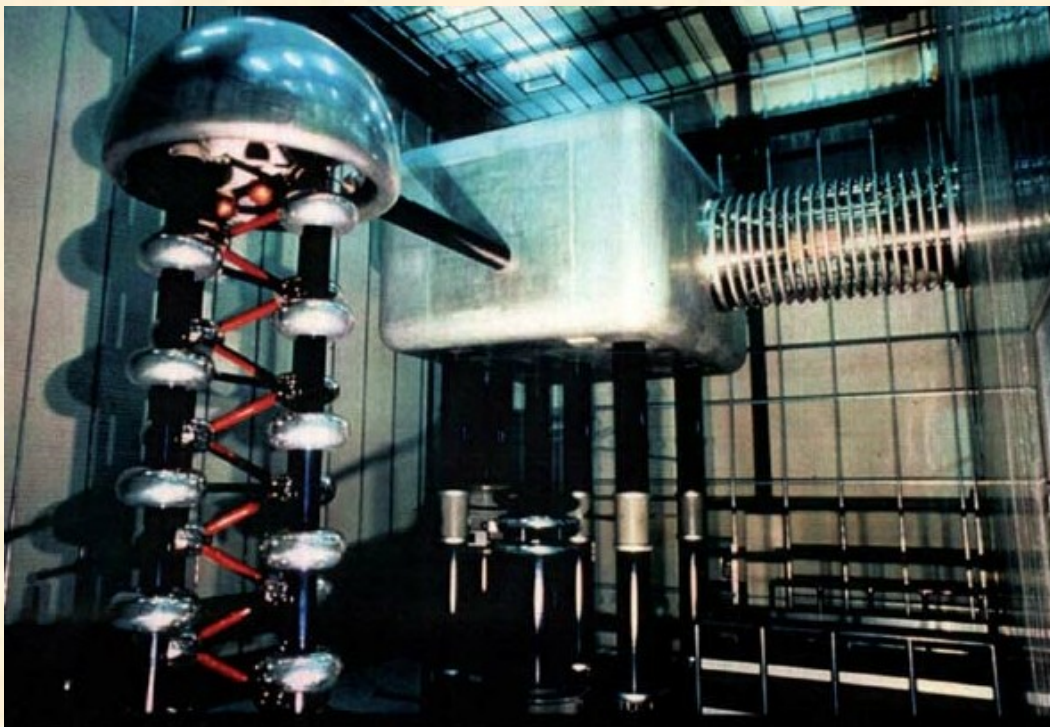
A Short Introduction ...

Bernhard Holzer
CERN, ABP

1.) Electrostatic Machines: The Cockcroft-Walton Generator

1928: Encouraged by Rutherford Cockcroft and Walton start the design & construction of a high voltage generator to accelerate a proton beam

1932: First particle beam (protons) produced for nuclear reactions: splitting of Li-nuclei with a proton beam of 400 keV



Particle source: Hydrogen discharge tube on 400 kV level

Accelerator: evacuated glass tube

Target: Li-Foil on earth potential

Technically: rectifier circuit, built of capacitors and diodes (Greinacher)

Problem:

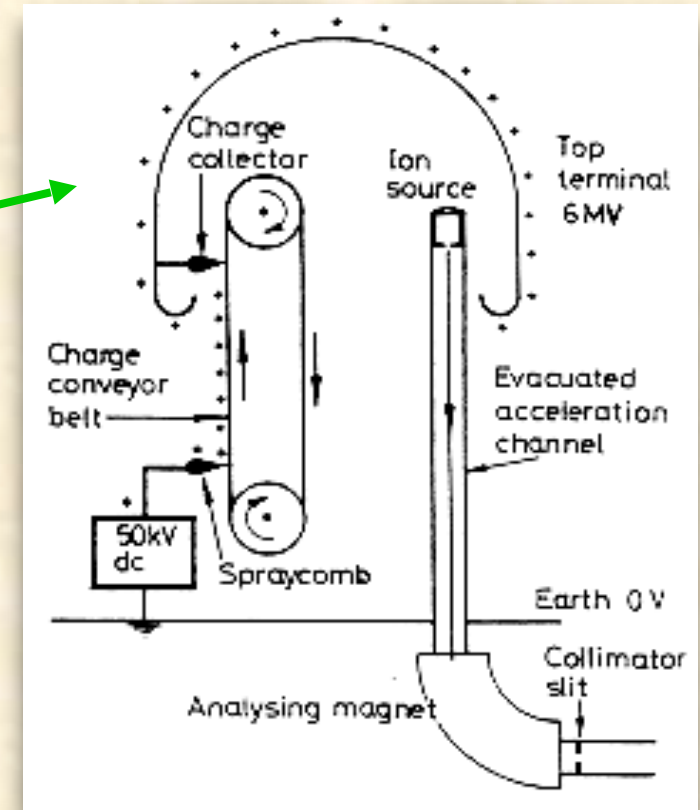
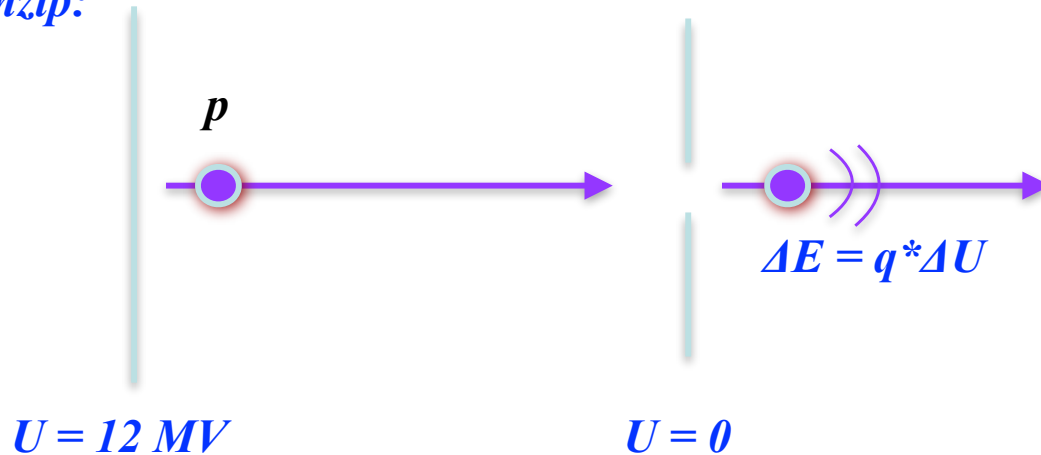
DC Voltage can only be used once

2.) Electrostatic Machines: van de Graaff Accelerator (1930 ...)

*creating high voltages by mechanical
transport of charges*

* *Terminal Potential: $U \approx 12 \dots 28 \text{ MV}$
using high pressure gas to suppress discharge (SF_6)*

Das Prinzip:



Problems:

- * Particle energy limited by high voltage discharges
- * high voltage **can only be applied once per particle ...**
... or twice ?

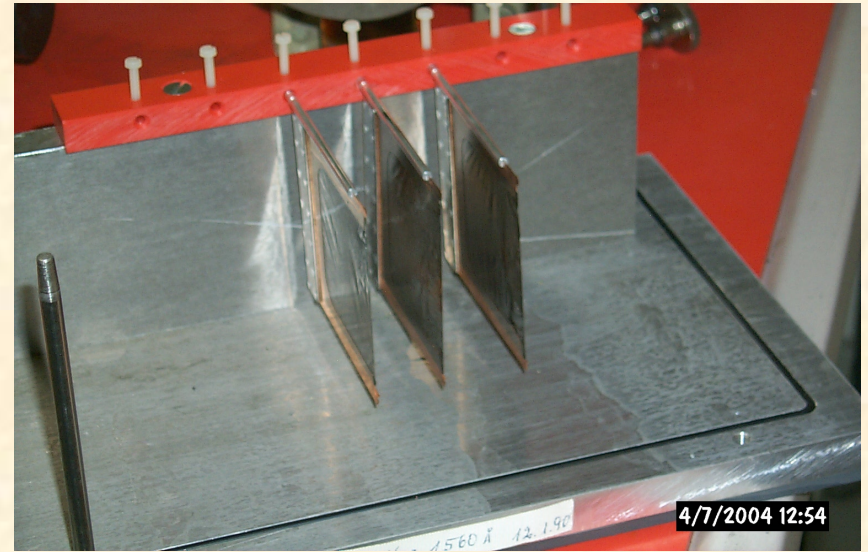
*The „Tandem principle“: Apply the accelerating voltage twice ...
... by working with **negative ions (e.g. H⁻)** and
stripping the electrons in the centre of the structure*

*Example for such a „steam engine“: 12 MV-Tandem van de Graaff
Accelerator at MPI Heidelberg*



*Innen-Ansicht einer solchen
„Dampfmaschine“*

*stripping foils: 1500 Å
... so nebenbei: das sind ca 1000
Atomlagen.*



Poggensee, DESY

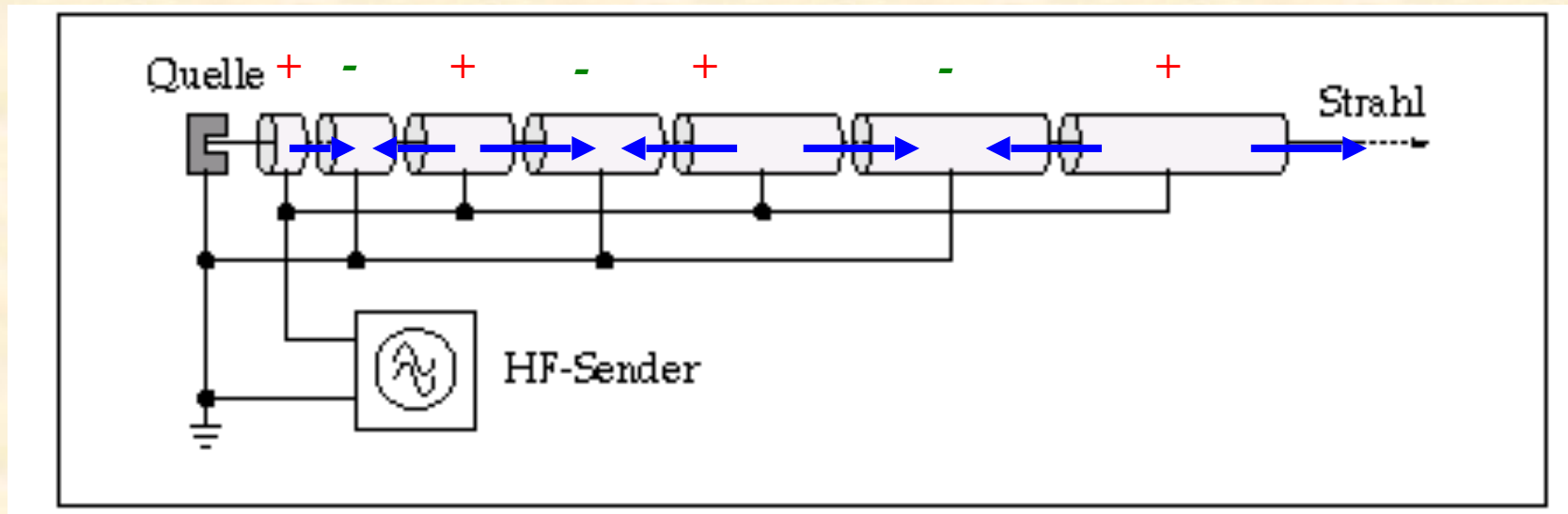


*mechanischer Transport der
Ladung mittels Glasfaserband*

3.) The first RF-Accelerator: „Linac“

1928, Wideroe: how can the acceleration voltage be applied several times to the particle beam

schematic Layout:



Energy gained after n acceleration gaps

$$E_n = n \cdot q \cdot U_0 \cdot \sin \psi_s$$

n number of gaps between the drift tubes

q charge of the particle

U_0 Peak voltage of the RF System

Ψ_s synchronous phase of the particle

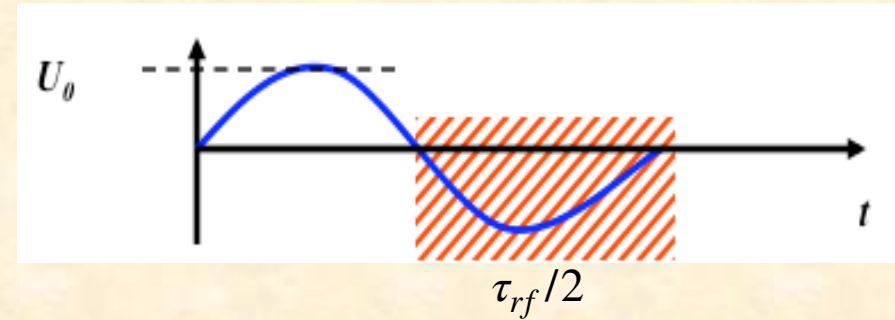
** acceleration of the proton in the first gap*

** voltage has to be „flipped“ to get the right sign in the second gap → RF voltage*

→ shield the particle in drift tubes during the negative half wave of the RF voltage

Wideroe-Structure: the drift tubes

shielding of the particles during the negative half wave of the RF



Time span of the negative half wave: $\tau_{rf}/2$

Length of the Drift Tube: $l_n = v_n \cdot \frac{\tau_{rf}}{2}$

Kinetic Energy of the Particles $E_n = \frac{1}{2}mv^2$

$$v_n = \sqrt{2E_n/m}$$

$$l_n = \frac{1}{f_{rf}} \cdot \sqrt{\frac{n \cdot q \cdot U_0 \cdot \sin\psi_s}{2m}}$$

Achtung !!! valid for **non relativistic** particles ...

Energy: ≈ 20 MeV per Nucleon $\beta \approx 0.04 \dots 0.6$, Particles: Protons/Ions

Und jetzt a bissi Einstein:

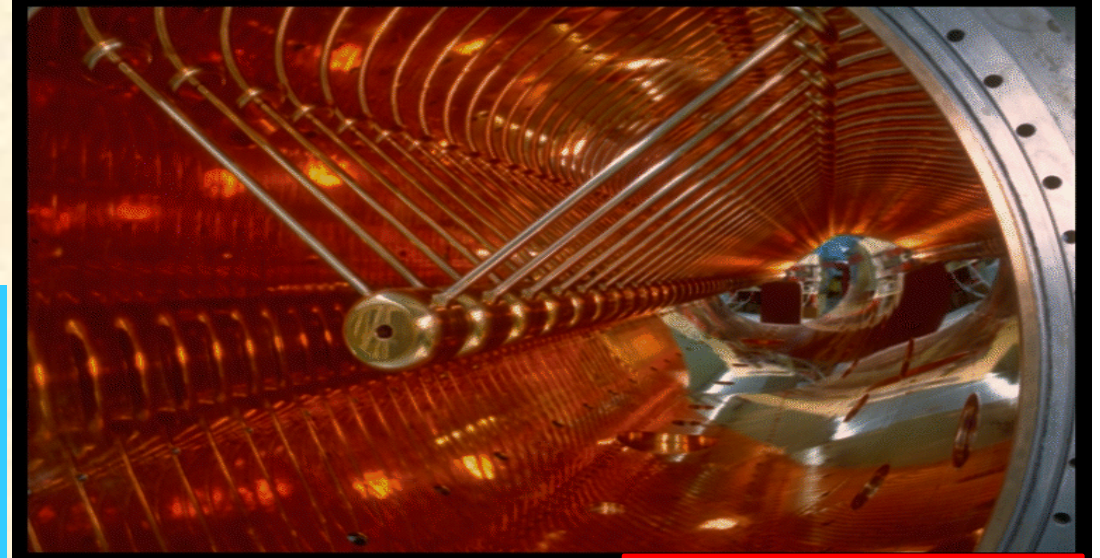
Beam energies

1.) reminder of some relativistic formula

rest energy $E_0 = m_0 c^2$

total energy $E = \gamma \cdot E_0 = \gamma \cdot m_0 c^2$

kinetic energy $E_{kin} = E_{total} - m_0 c^2$



Energy Gain per „Gap“:

$$W = q U_0 \sin \omega_{RF} t$$

Linac III:

$$E_{total} = 988 \text{ MeV}$$

$$m_0 c^2 = 938 \text{ MeV}$$

$$E_{kin} = 50 \text{ MeV}$$

$$\gamma = \frac{E_{ges}}{E_0} = 988/938 = 1.05$$

—> *im klassischen Bereich*

3.) The Cyclotron: (Livingston / Lawrence ~1930)

Problem:

Linacs werden bei $v=c$ sehr schnell sehr langgggg.

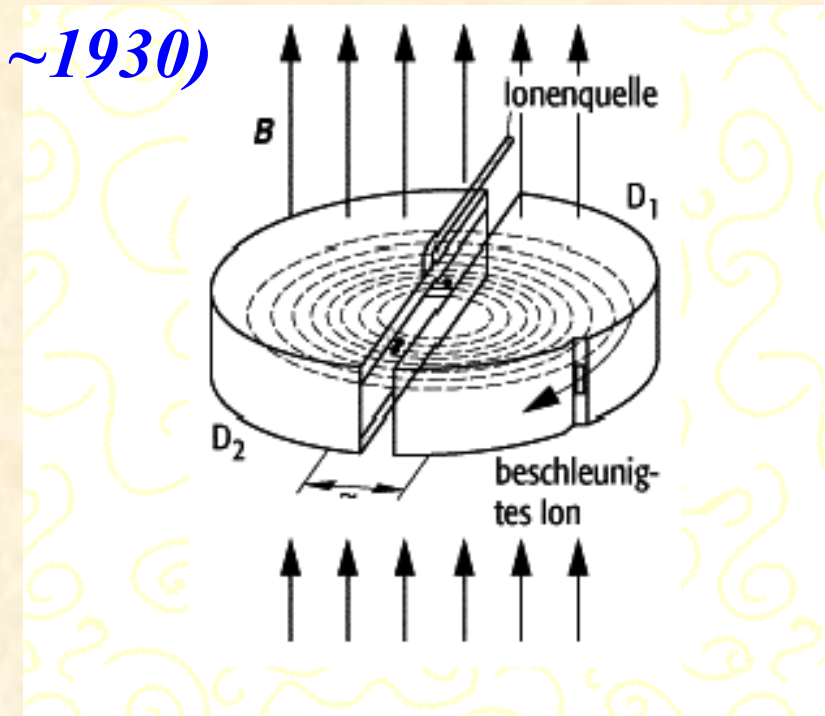
—> Man erhaelt ne kompakte (d.h. billigere) Maschine, wenn man den Orbit der Teilchen aufwickelt.

Idea: Apply a magnetic field: $B = \text{const}$

Lorentzforce

$$F = q \cdot v \cdot B$$

geladene Teilchen in Bewegung werden im Magnetfeld abgelenkt.



Kreisbahn-Bedingung:

Zentrifugalkraft wird durch die entgegengesetzte Lorentz-Kraft aufgehoben.

$$F_{\text{Lorentz}} = F_{\text{zentrifugal}}$$

$$q \cdot v \cdot B = \frac{mv^2}{r}$$

$$B \cdot R = \frac{mv}{q} \quad \longrightarrow \quad B \cdot R = \frac{p}{q}$$

3.) The Cyclotron: (Livingston / Lawrence ~1930)

revolution frequency

$$\omega_{\text{revol}} = \frac{v}{r} = \frac{q}{m} \cdot B = \text{const!!!}$$

Die Umlauf-frequenz im Cyclotron ist konstant.

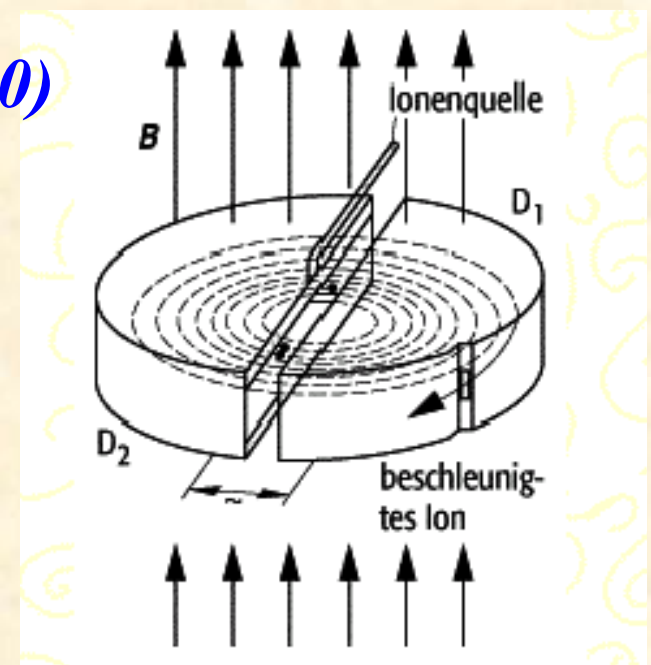
Wir lassen eine gleich-grosse konstante RF frequenz auf die Teilchen los und die Kiste funktioniert.

$$\omega_{\text{rf}} = \omega_{\text{revolution}} \quad \text{oder} \quad \omega_{\text{rf}} = h \cdot \omega_{\text{revolution}}$$

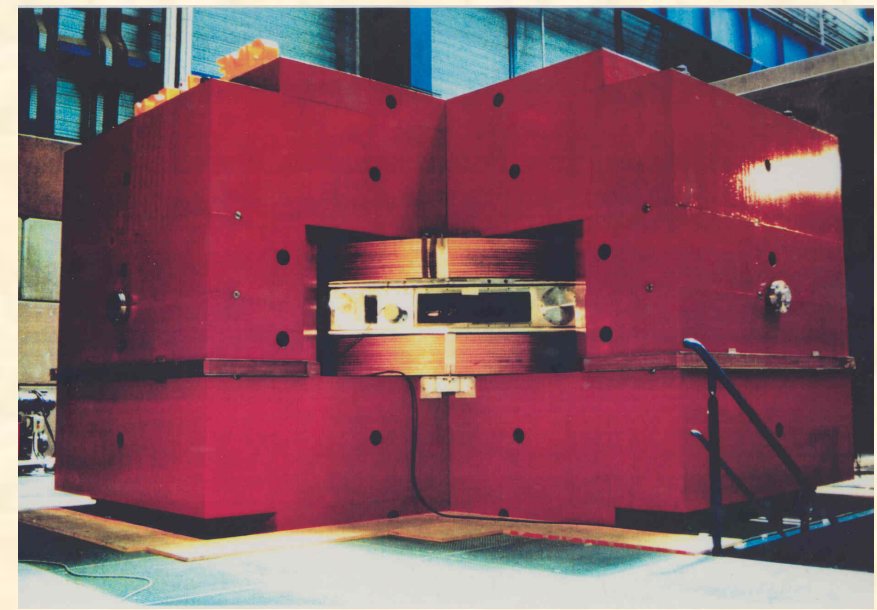
*rf-frequency = h * revolution frequency,
h = "harmonic number"*

Problem: Albert !!!

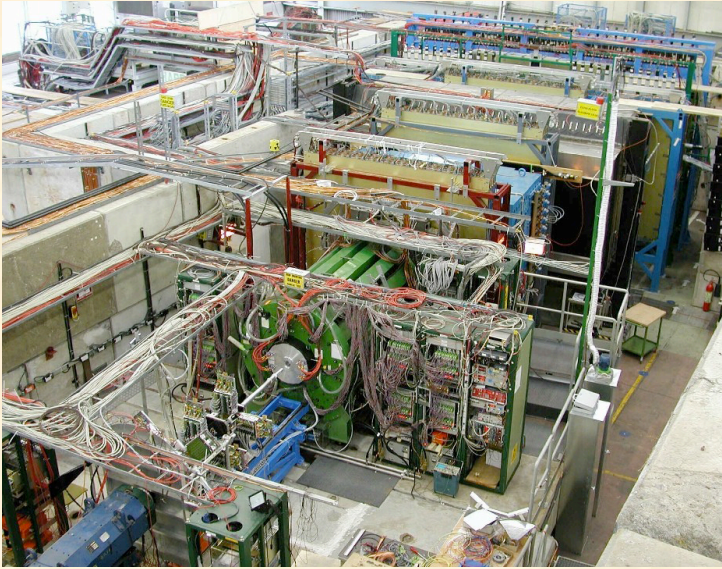
$$\left. \begin{aligned} \omega_{\text{revol}} &= \frac{q}{\gamma m} \neq \text{const} \\ \omega_s(t) &= \omega_{\text{rf}}(t) = \frac{q}{\gamma(t) \cdot m_0} B \end{aligned} \right\} \begin{aligned} &\text{Synchro-Cyclotron} \\ &\text{Nachfahren der RF} \\ &\text{Frequenz} \end{aligned}$$



increasing radius for increasing momentum → Spiral Trajectory

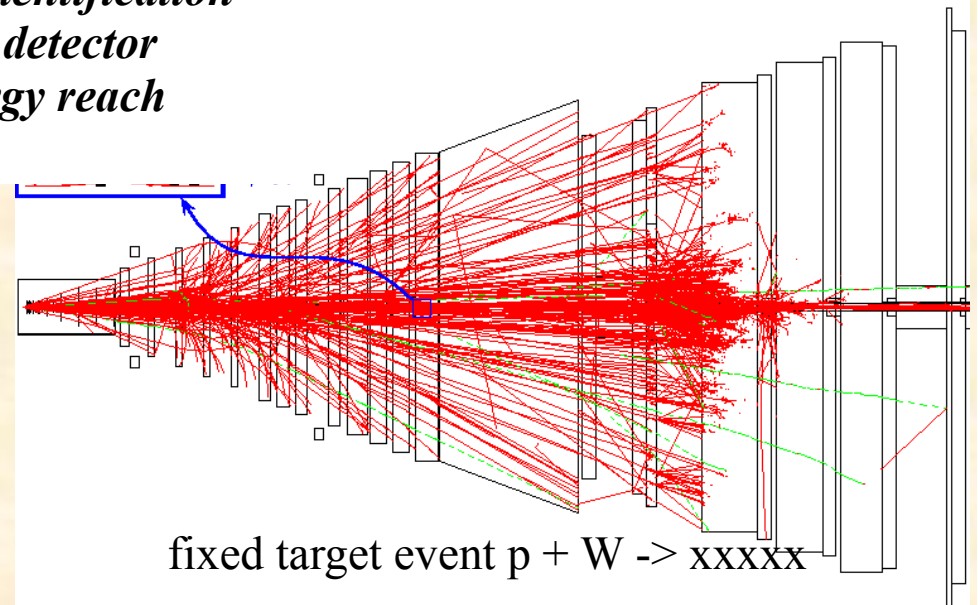


Fixed target experiments:



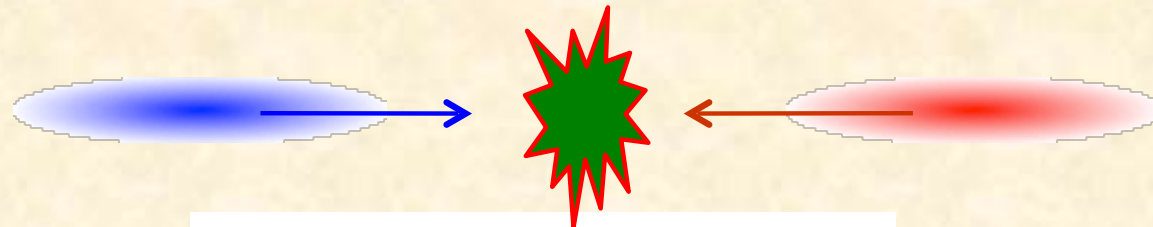
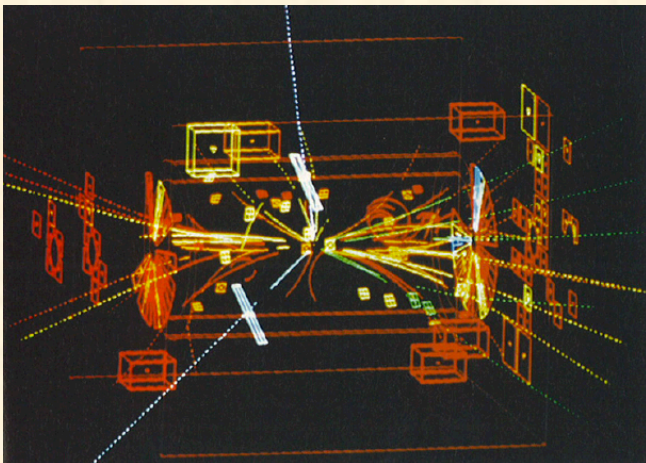
HARP Detector, CERN

high event rate
easy track identification
asymmetric detector
limited energy reach



Collider experiments:

$$E=mc^2$$



low event rate (luminosity)
challenging track identification
symmetric detector

$$E_{lab} = E_{cm}$$

Z_0 boson discovery at the UA2 experiment (CERN).

The Z_0 boson decays into a e^+e^- pair, shown as white dashed lines.

II.) A Bit of Theory

The big storage rings: „Synchrotrons“

1.) Introduction and Basic Ideas

„ ... in the end and after all it should be a kind of circular machine“
→ need transverse deflecting force

Lorentz force $\vec{F} = q * (\cancel{\vec{E}} + \vec{v} \times \vec{B})$

typical velocity in high energy machines: $v \approx c \approx 3 * 10^8 \text{ m/s}$

Example:

$$B = 1 \text{ T} \quad \rightarrow \quad F = q * 3 * 10^8 \frac{\text{m}}{\text{s}} * 1 \frac{\text{Vs}}{\text{m}^2}$$

$$F = q * 300 \frac{\text{MV}}{\text{m}}$$

equivalent E
electrical field:

Technical limit for electrical fields:

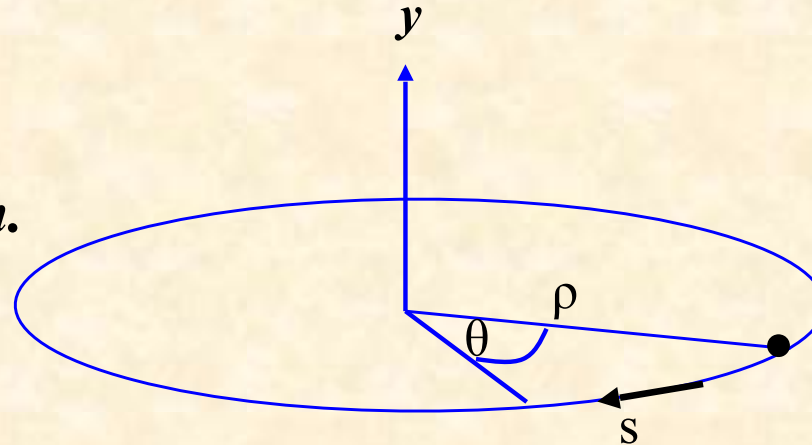
$$E \leq 1 \frac{\text{MV}}{\text{m}}$$

Ein Speicherring besteht aus Magneten, Magneten und Magneten

und ein wenig Vakuum-Kammern, Strahldiagnose, und RF Systemen

The ideal circular orbit

... das hatten wir schon.



circular coordinate system

condition for circular orbit:

Lorentz force

$$F_L = e v B$$

centrifugal force

$$F_{\text{centr}} = \frac{\gamma m_0 v^2}{\rho}$$

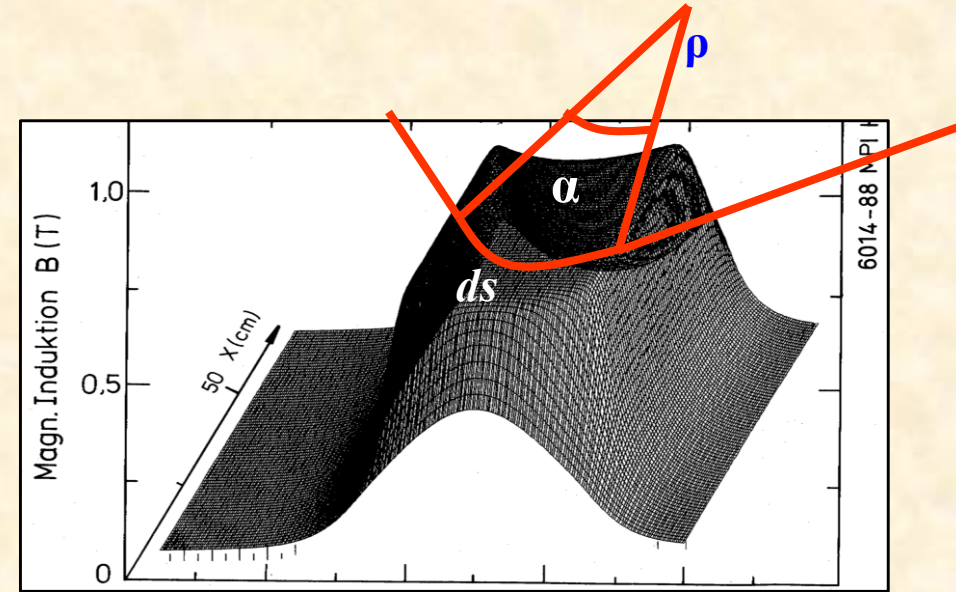
$$\frac{\gamma m_0 v^2}{\rho} = e v B$$

$$\frac{p}{e} = B \rho$$

$B \rho =$ "beam rigidity"

... und jetzt isses sogar relativistisch korrekt.

The Magnetic Guide Field



field map of a storage ring dipole magnet

Dipole erzeugen mit zwei parallelen Polschuhen ein konstantes (!) Magnetfeld

$$B \approx 1 \dots 8 \text{ T}$$

*Achtung: um zum Pluto zu kommen
muessen wir höchste Präzision
fordern.*

$$\frac{\Delta B}{B} \approx 10^{-4}$$

Ablenkradius:

$$\rho = \frac{p}{e B} = \frac{7000 \cdot 10^9 \text{ eV}}{3 \cdot 10^8 \text{ m/s} \cdot 8 \text{ Vs/m}^2}$$

*nota bene:
die allgemeinste Ausdruck fuer
die Energie ist*

$$\rho = 2.8 \text{ km}$$

$$E^2 = p^2 c^2 + m^2 c^4 \rightarrow p \approx \frac{E}{c}$$

*Achtung:
um Energie unabhängige
Gleichungen zu erhalten teilen
wir die Felder durch "p"*

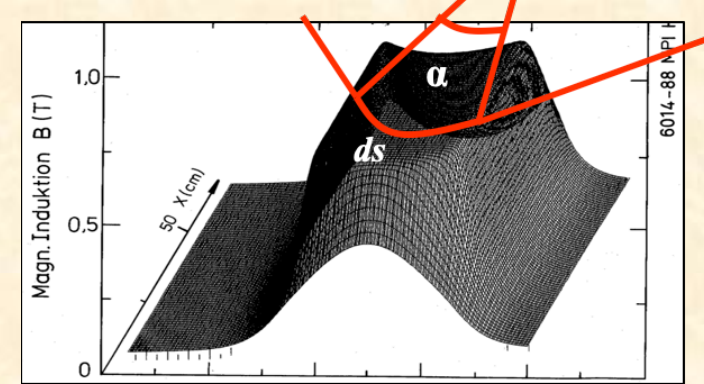
*„normalised
bending strength“*

$$1/\rho = \frac{B}{p/e}$$

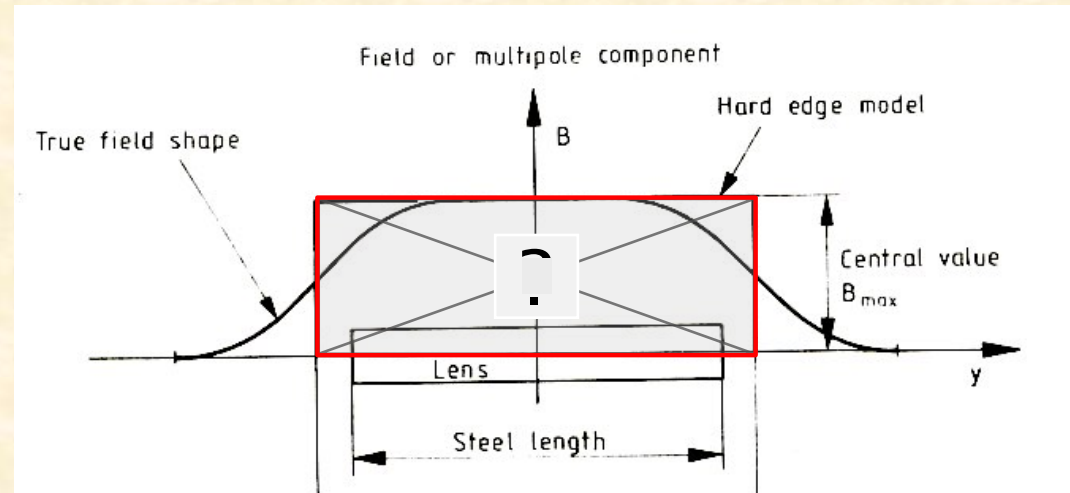
Bending Angle

„integrated field strength”

$$\alpha = \frac{B^* dl}{B^* \rho}$$



$$B l_{eff} = \int_0^{l_{mag}} B ds$$



The angle swept out in one revolution must be 2π , so

$$\alpha = \frac{\int B dl}{B^* \rho} = 2\pi \quad \rightarrow \quad \int B dl = 2\pi * \frac{p}{q} \quad \dots \text{for a full circle}$$

2.) *Focusing Forces: Hook's law*

*... keeping the flocs together:
In addition to the pure bending of the beam
we have to keep 10^{11} particles close together*



focusing force

Um auf unserem Weg zu Pluto unsere 10^{11} Teilchen zusammen zu halten wenden wir den Trick der linearen rücktreibenden Kraft an:

D.h. unsere Teilchen werden wie 10^{11} Pendelchen hin und her schwingen

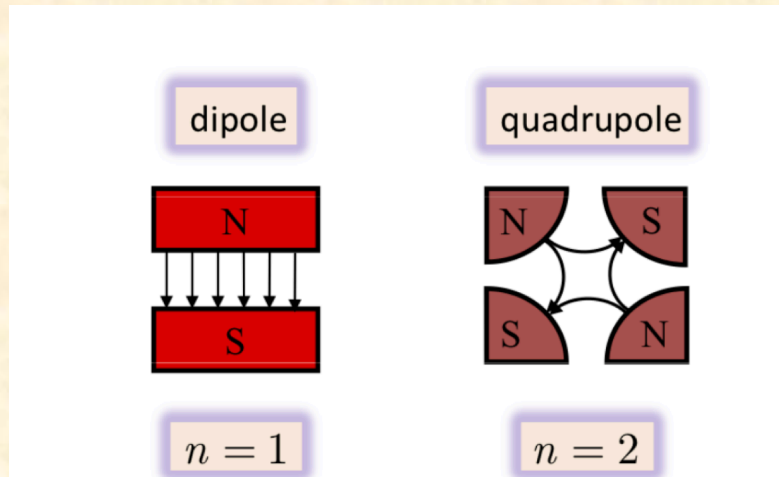
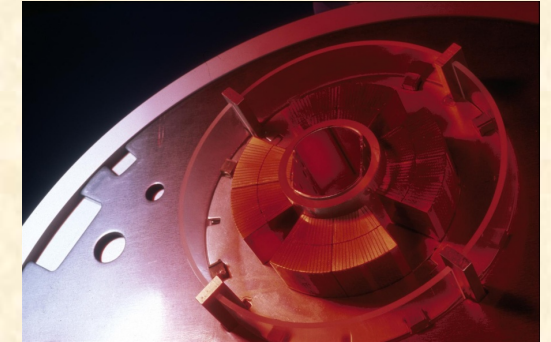
Wir müssen es nur schaffen, ein Magnetfeld zu erzeugen, das linear ansteigt.

2.) Focusing Forces: Quadrupole Fields

Apply this concept to magnetic forces: we need a Lorentz force that rises as a function of the distance to ...

... the design orbit

$$F(x) = q \cdot v \cdot B(x)$$



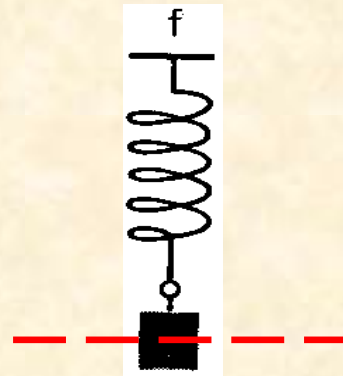
Dipoles: Create a constant field

$$B_y = \text{const}$$

Quadrupoles: Create a linear increasing magnetic field:

$$B_y(x) = g \cdot x, \quad B_x(y) = g \cdot y$$

Federpendel im Physik Buch



there is a **restoring force**, proportional to the elongation x :

$$F = m * a = - \text{const} * x$$

$$F = m * \frac{d^2x}{dt^2} = - \text{const} * x$$

Hook's Federgesetz: $F = - k * x$

*Integration liefert uns eine cos- artige Lösung
oder eine sinus artige*

$$x(t) = A \cdot \cos(\omega t)$$

$$x(t) = B \cdot \sin(\omega t)$$

oder eine Kombination aus beiden

$$x_{\text{allg}}(t) = A \cdot \cos(\omega t) + B \cdot \sin(\omega t)$$

Vorteil:

harmonische Schwingungen sind sehr (!!) stabil,
haben eine wohldefinierte Frequenz
sind in der Natur (i.,e. Physik) weitverbreitet

Focusing forces and particle trajectories:

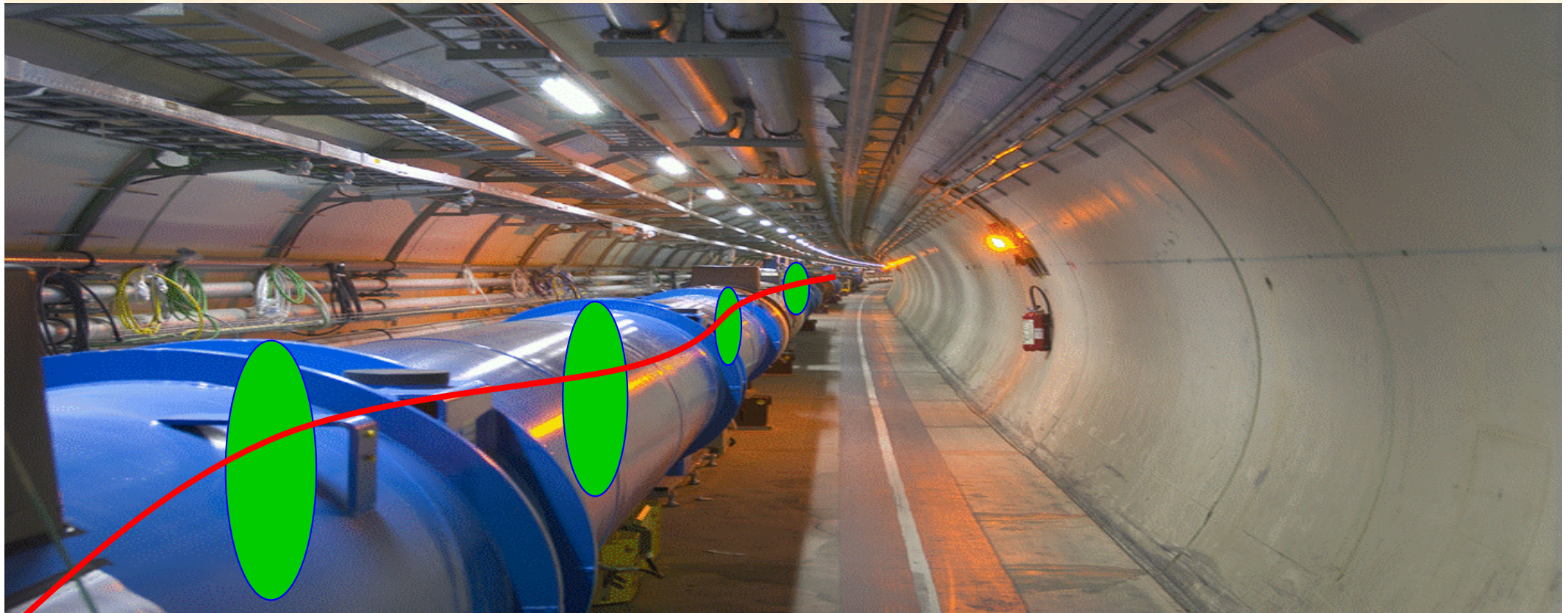
*normalise magnet fields to momentum
(remember: $\mathbf{B} * \rho = p / q$)*

Dipole Magnet

$$\frac{B}{p/q} = \frac{B}{B\rho} = \frac{1}{\rho}$$

Quadrupole Magnet

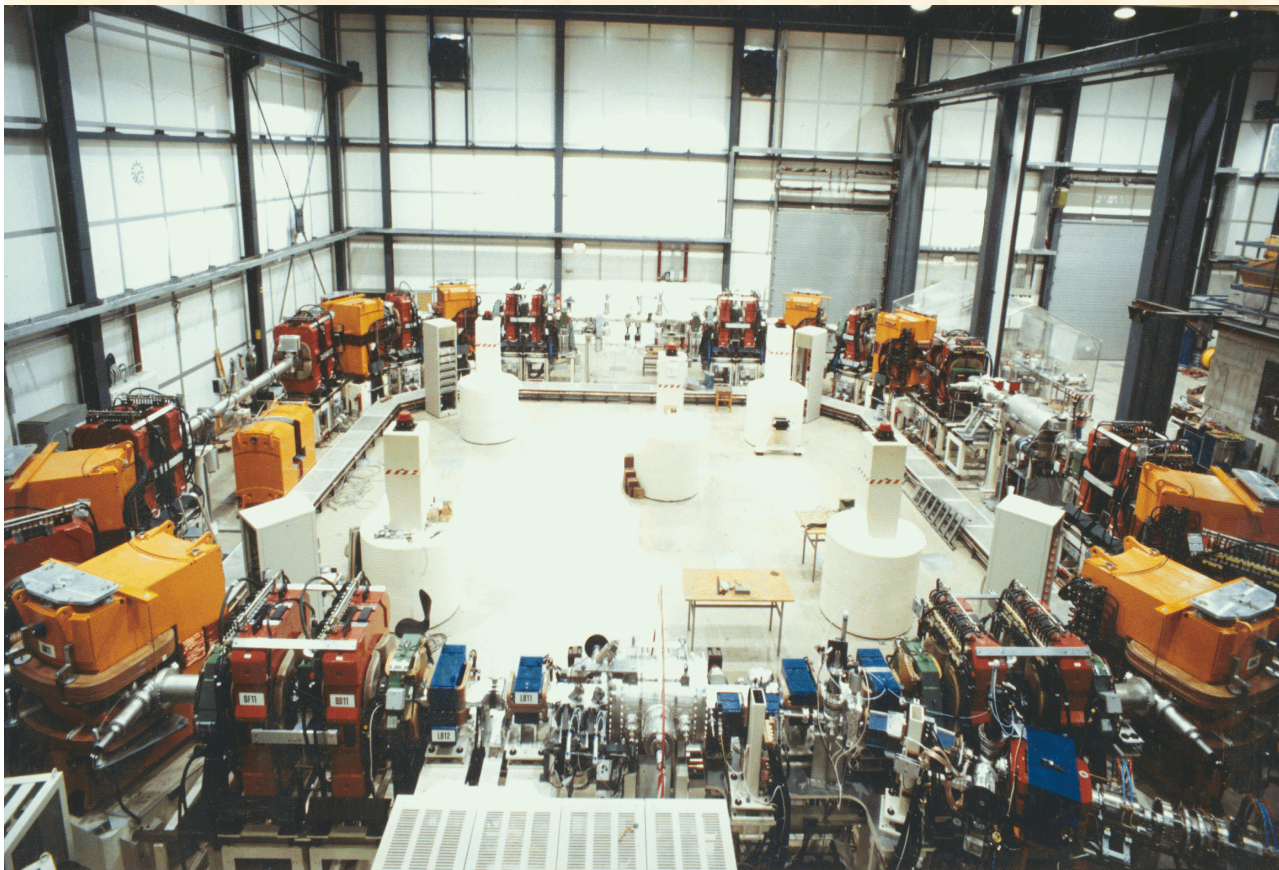
$$k := \frac{g}{p/q}$$



3.) The Equation of Motion:

$$\frac{B(x)}{p/e} = \frac{1}{\rho} + kx + \frac{1}{2!} \cancel{m} x^2 + \frac{1}{3!} \cancel{n} x^3 + \dots$$

only terms linear in x, y taken into account **dipole fields**
quadrupole fields



Separate Function Machines:

Split the magnets and optimise them according to their job:

bending, focusing etc

*Example:
heavy ion storage ring TSR*

* *man sieht nur
dipole und quads → linear*

The Equation of Motion:

* Equation for the *horizontal motion*:

$$x'' + x \cdot \left(\frac{1}{\rho^2} + k \right) = 0$$

x = *particle amplitude*

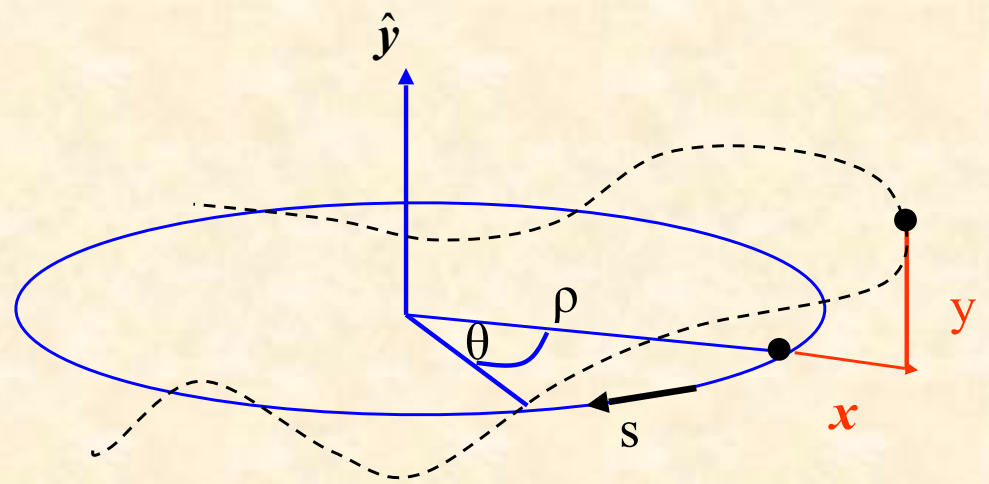
x' = *angle of particle trajectory (wrt ideal path line)*

$$x'' = - x \cdot \underbrace{\left(\frac{1}{\rho^2} + k \right)}$$

$$x'' = - K \cdot x$$

Hook's Gesetz fuer Speicherringe

... es gibt da nur ein kleines Problem:



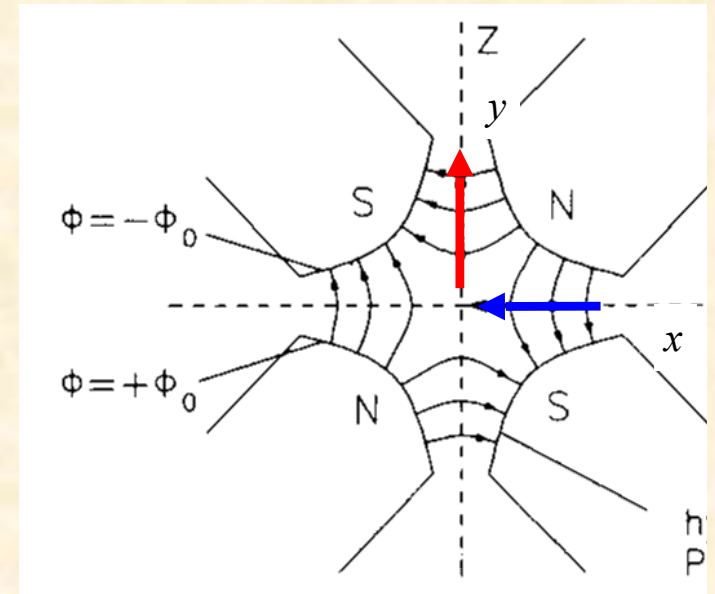
In der vertikalen Ebene drehen sich die Magnetfeld-Linien um

* *Equation for the vertical motion:*

$$\frac{1}{\rho^2} = 0 \quad \text{no dipoles ... in general ...}$$

$$k \leftrightarrow -k \quad \text{quadrupole field changes sign}$$

$$y'' - k \cdot y = 0$$



*... und Teilchen, die in der horizontalen Ebene fokussiert werden,
werden im gleichen Atemzug in der vertikalen Ebene aus der Maschine befördert.*

4.) Solution of Trajectory Equations

Define ... hor. plane: $K = 1/\rho^2 + k$

... vert. Plane: $K = -k$

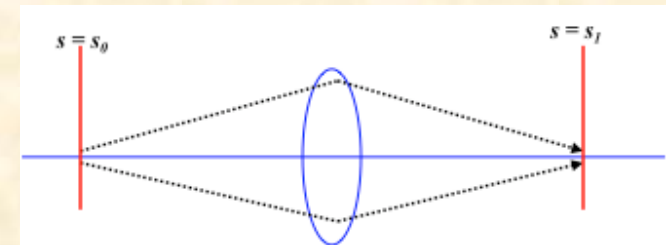
$$x'' + K x = 0$$

Differential Equation of harmonic oscillator ... with spring constant K

Ansatz: **Hor. Focusing Quadrupole $K > 0$:**

$$x(s) = x_0 \cdot \cos(\sqrt{|K|}s) + x'_0 \cdot \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|}s)$$

$$x'(s) = -x_0 \cdot \sqrt{|K|} \cdot \sin(\sqrt{|K|}s) + x'_0 \cdot \cos(\sqrt{|K|}s)$$



... da ist wieder unsere Kuckucksuhr.

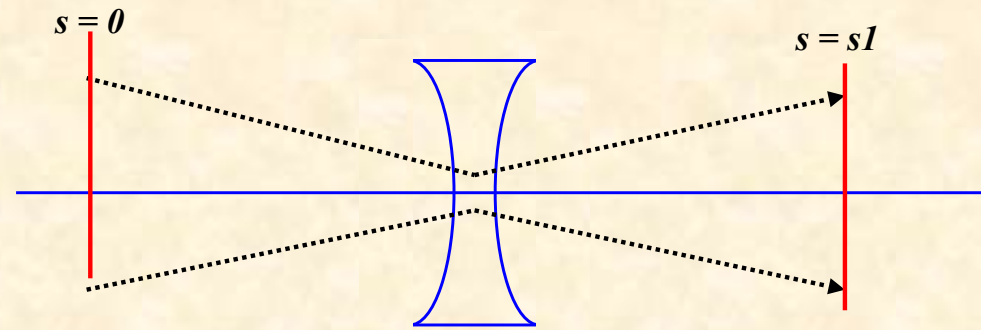
For convenience expressed in matrix formalism:

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s1} = M_{foc} * \begin{pmatrix} x \\ x' \end{pmatrix}_{s0}$$

$$M_{foc} = \begin{pmatrix} \cos(\sqrt{|K|}l) & \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|}l) \\ -\sqrt{|K|} \sin(\sqrt{|K|}l) & \cos(\sqrt{|K|}l) \end{pmatrix}$$

hor. defocusing quadrupole:

$$x'' - K x = 0$$



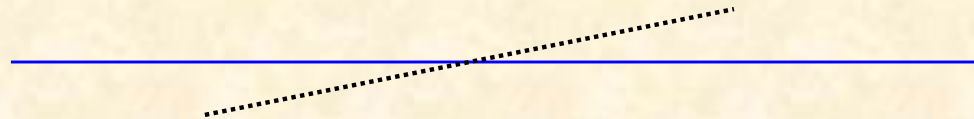
Ansatz: Remember from school

$$x(s) = a_1 \cdot \cosh(\omega s) + a_2 \cdot \sinh(\omega s)$$

$$M_{defoc} = \begin{pmatrix} \cosh \sqrt{|K|} l & \frac{1}{\sqrt{|K|}} \sinh \sqrt{|K|} l \\ \sqrt{|K|} \sinh \sqrt{|K|} l & \cosh \sqrt{|K|} l \end{pmatrix}$$

drift space:

$$K = 0$$



$$x(s) = x'_0 \cdot s$$

$$M_{drift} = \begin{pmatrix} 1 & l \\ 0 & 1 \end{pmatrix}$$

! *with the assumptions made, the motion in the horizontal and vertical planes are independent „ ... the particle motion in x & y is uncoupled“*

... zur Erinnerung:

hyperbolische Funktionen führen leicht zu Panik Attacken !

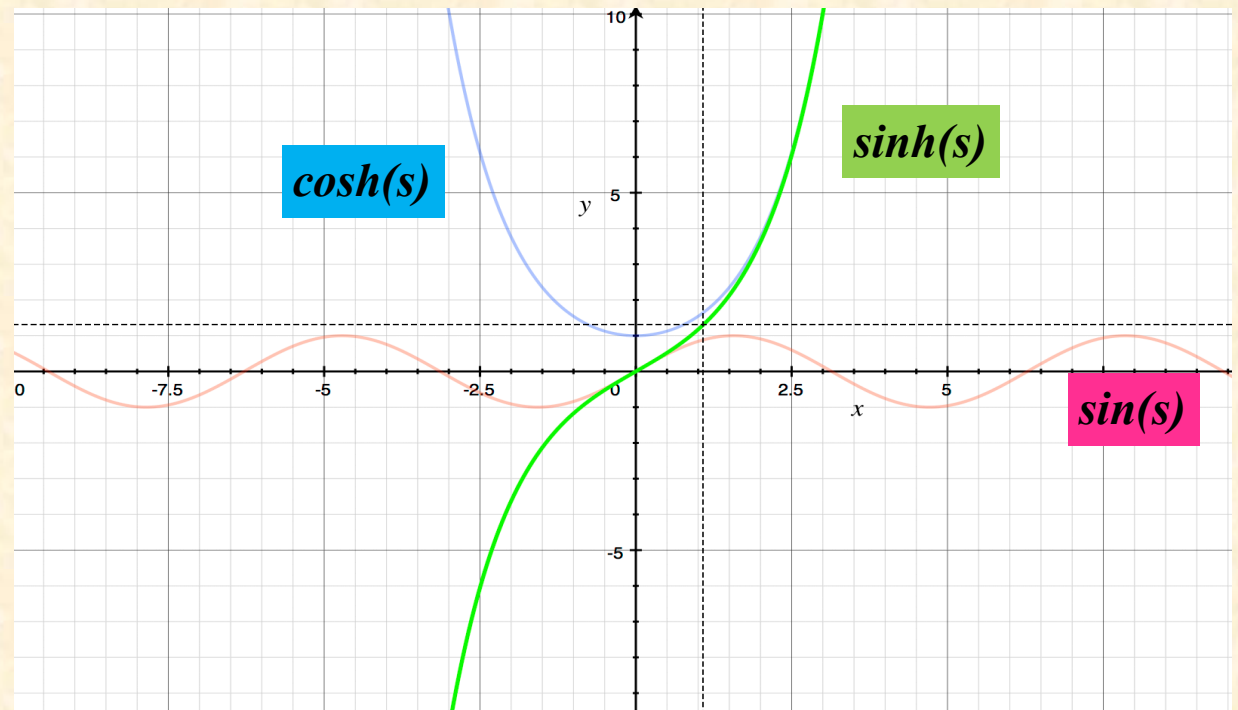
$$M_{defoc} = \begin{pmatrix} \cosh \sqrt{|K|}l & \frac{1}{\sqrt{|K|}} \sinh \sqrt{|K|}l \\ \sqrt{|K|} \sinh \sqrt{|K|}l & \cosh \sqrt{|K|}l \end{pmatrix}$$

$$f(s) = \sin(s) \quad f(s) = \cos(s)$$

$$f(s) = \sinh(s) \quad f(s) = \cosh(s)$$

Ansatz für die Teilchenbewegung im defokussierenden Fall:

$$x(s) = a_1 \cdot \cosh(\omega s) + a_2 \cdot \sinh(\omega s)$$

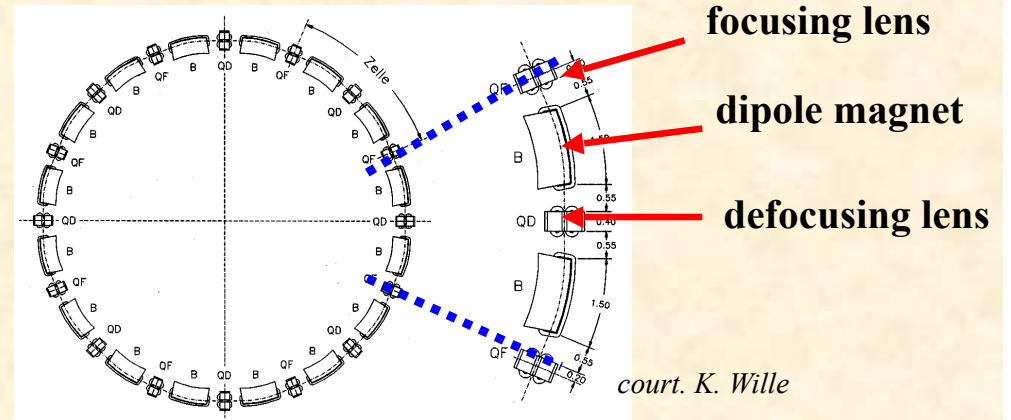


Transformation through a system of lattice elements

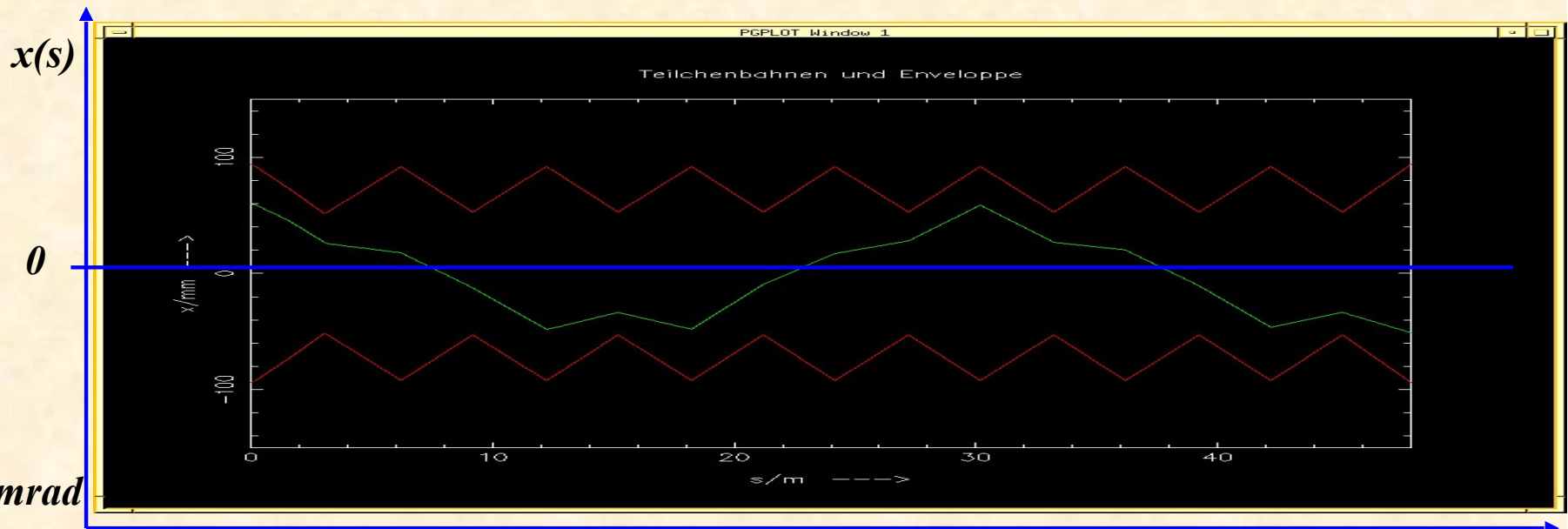
combine the single element solutions by multiplication of the matrices

$$M_{total} = M_{QF} * M_D * M_{QD} * M_{Bend} * M_{D*} \dots$$

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s_2} = M(s_2, s_1) \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_{s_1}$$



in each accelerator element the particle trajectory corresponds to the movement of a harmonic oscillator ,,



typical values
in a strong
foc. machine:
 $x \approx mm, x' \leq mrad$

*“Once more unto the breach, dear friends, once more”
(W. Shakespeare, Henry 5)*

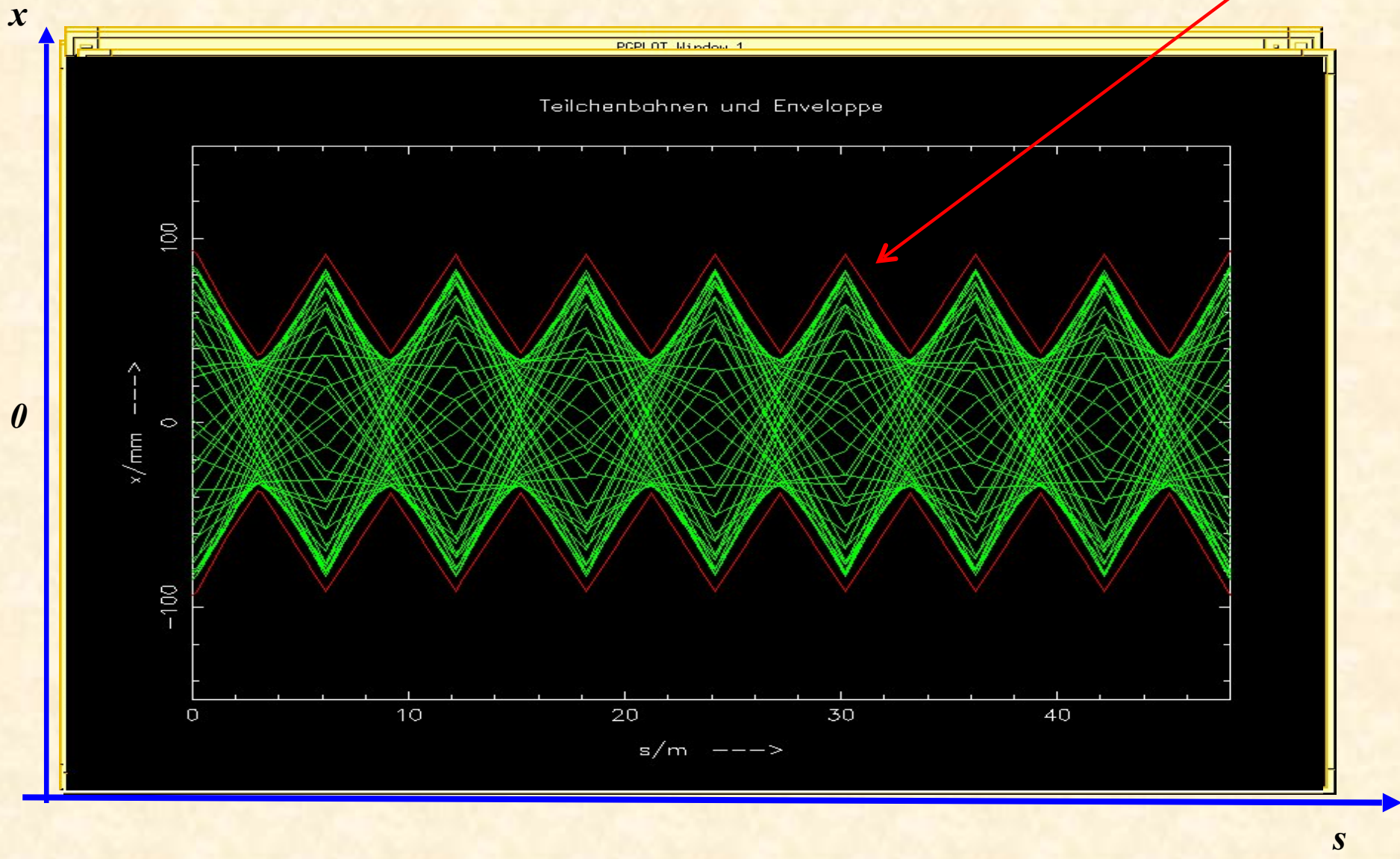
“Do they actually drop ?”

Answer: No

Question: what will happen, if the particle performs a second turn ?

... or a third one or ... 10^{10} turns

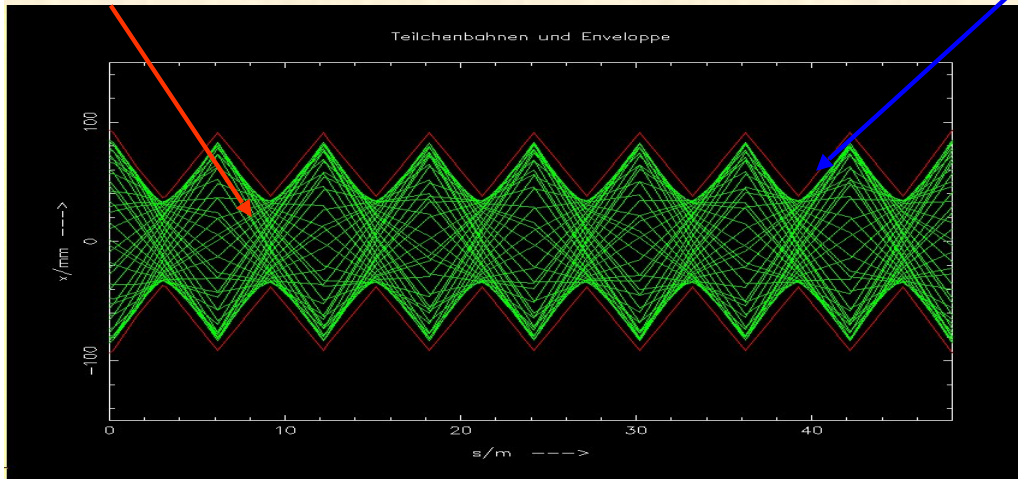
$$\sigma = \sqrt{\epsilon\beta}$$



Emittance of the Particle Ensemble:

$$x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cdot \cos(\Psi(s) + \phi)$$

$$\hat{x}(s) = \sqrt{\varepsilon} \sqrt{\beta(s)}$$



single particle trajectories, $N \approx 10^{11}$ per bunch

*Gauß
Particle Distribution:*

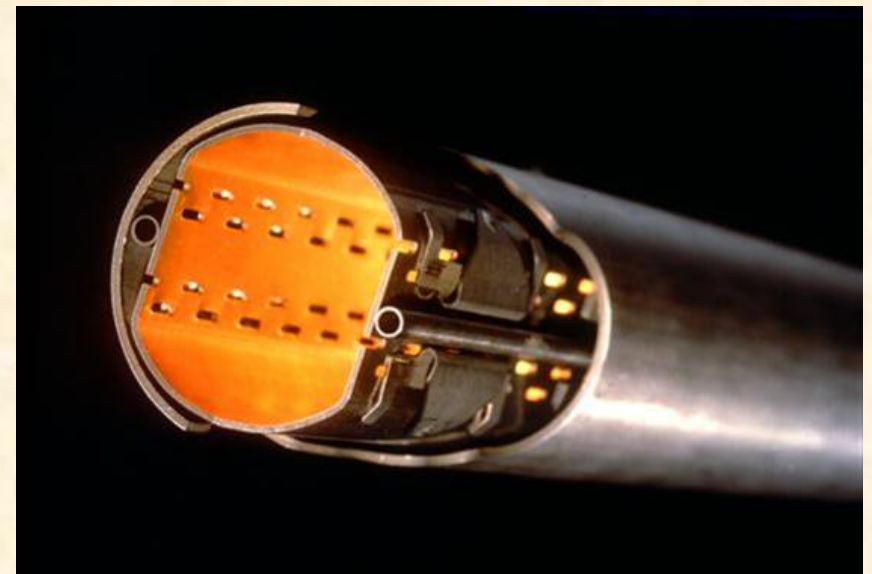
$$\rho(x) = \frac{N \cdot e}{\sqrt{2\pi} \sigma_x} \cdot e^{-\frac{1}{2} \frac{x^2}{\sigma_x^2}}$$

particle at distance 1σ from centre

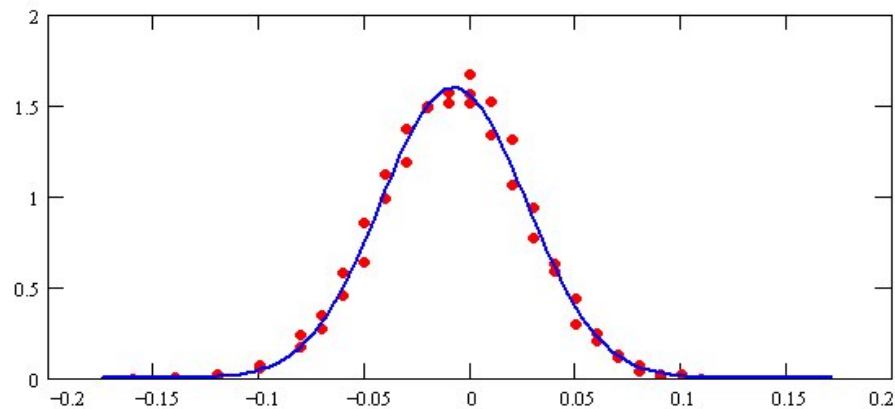
\leftrightarrow 68.3 % of all beam particles

LHC:

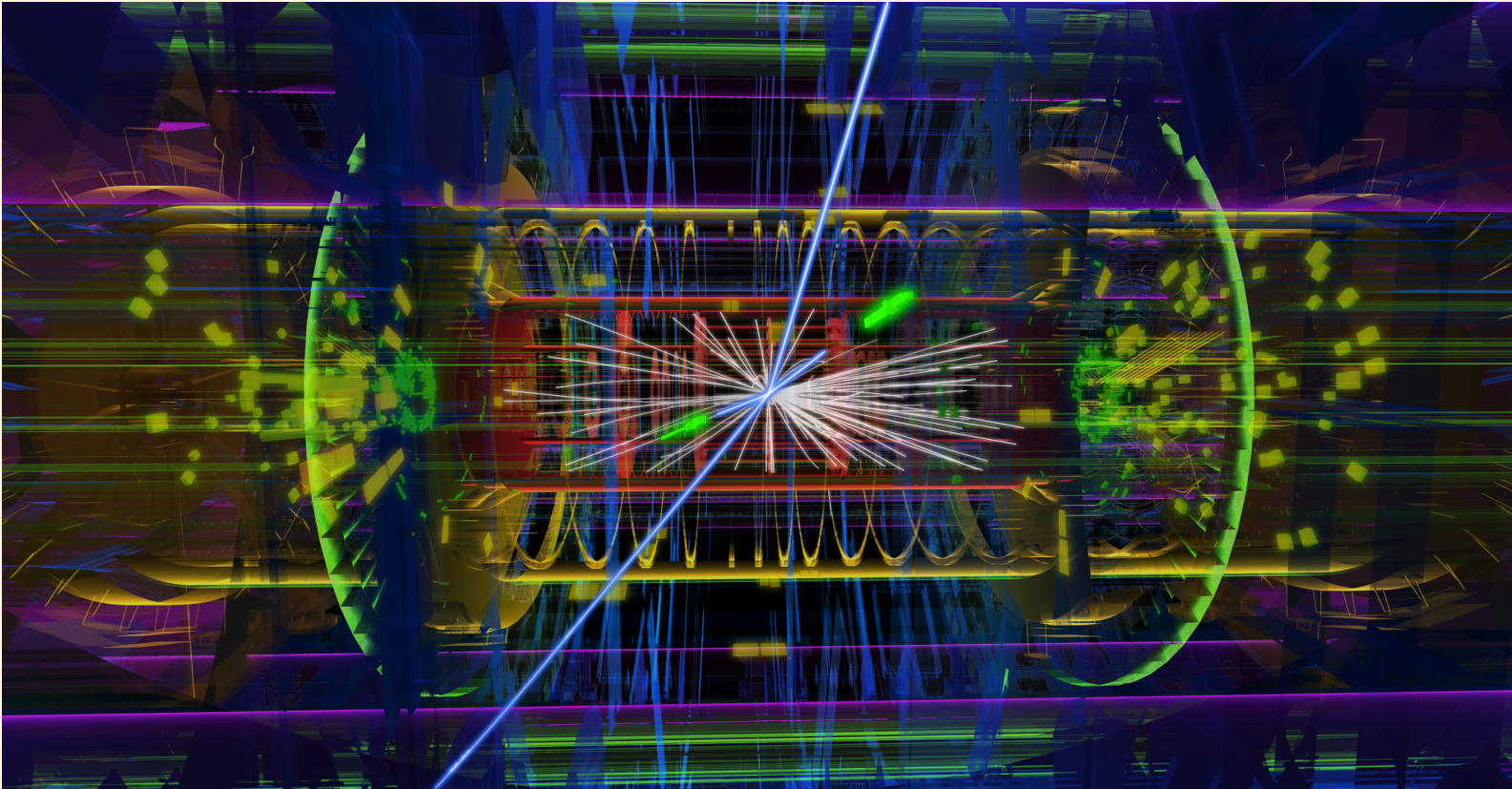
$$\sigma = \sqrt{\varepsilon * \beta} = \sqrt{5 * 10^{-10} m * 180 m} = 0.3 mm$$



*aperture requirements: $r_0 = 17 * \sigma$*



Collisions



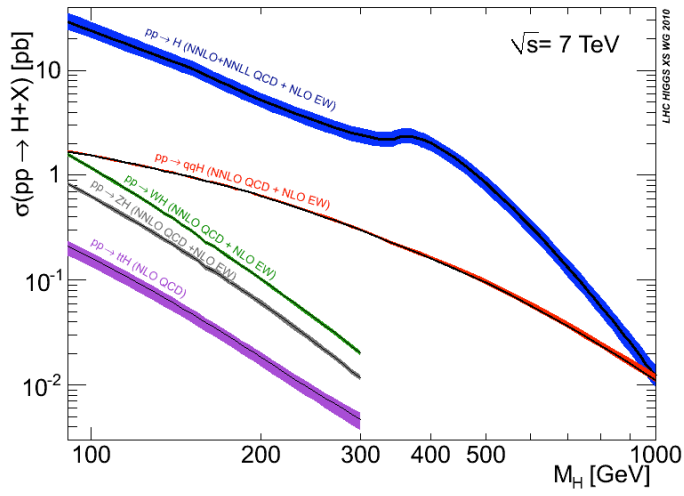
ATLAS event display: Higgs => two electrons & two muons

$$E = m_0c^2 = m_{e1} + m_{e2} + m_{\mu1} + m_{\mu2} = 125.4 \text{ GeV}$$

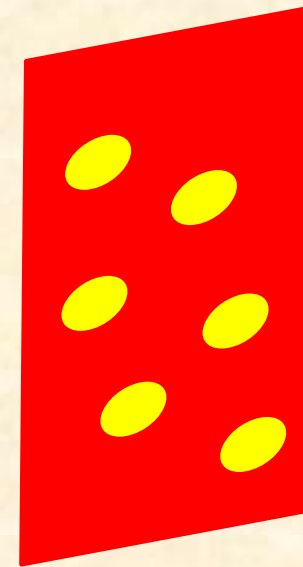
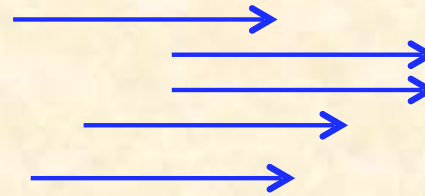
Problem: Our particles are VERY small !!

man trifft nicht so häufig.

Overall cross section of the Higgs:



$$\Sigma_{react} \approx 1pb$$



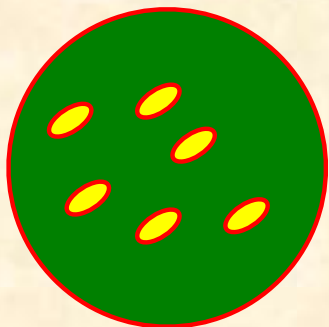
$$1b = 10^{-24} cm^2 = \frac{1}{mio} \cdot \frac{1}{mio} \cdot \frac{1}{mio} \cdot \frac{1}{10000} mm^2$$

$$1pb = 10^{-12}b \approx ZERO$$

The particles are "very small"

The only chance we have:

compress the transverse beam size ... at the IP

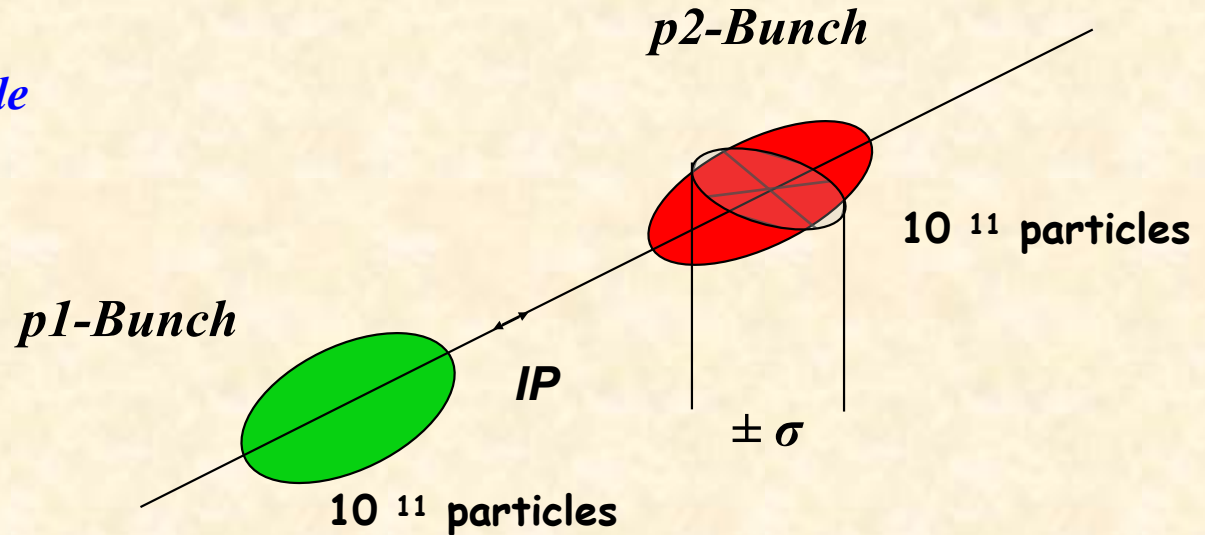


LHC typical \rightarrow 16 μm

5.) Luminosity

Ereignis Rate: "Physik" pro Sekunde

$$R = L \cdot \Sigma_{react}$$



Example: Luminosity run at LHC

$$\sigma_x = \sigma_y = 16 \mu\text{m}$$

Strahlgröße am IP

$$f_0 = 11.245 \text{ kHz}$$

Umlauffrequenz

$$n_b = 2808$$

Zahl der Bunche

$$N_p = 1.2 \cdot 10^{11}$$

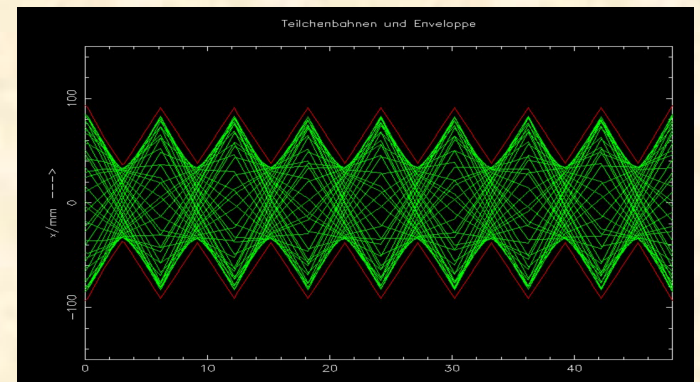
Teilchen in einem Bunch

$$I_p = 584 \text{ mA}$$

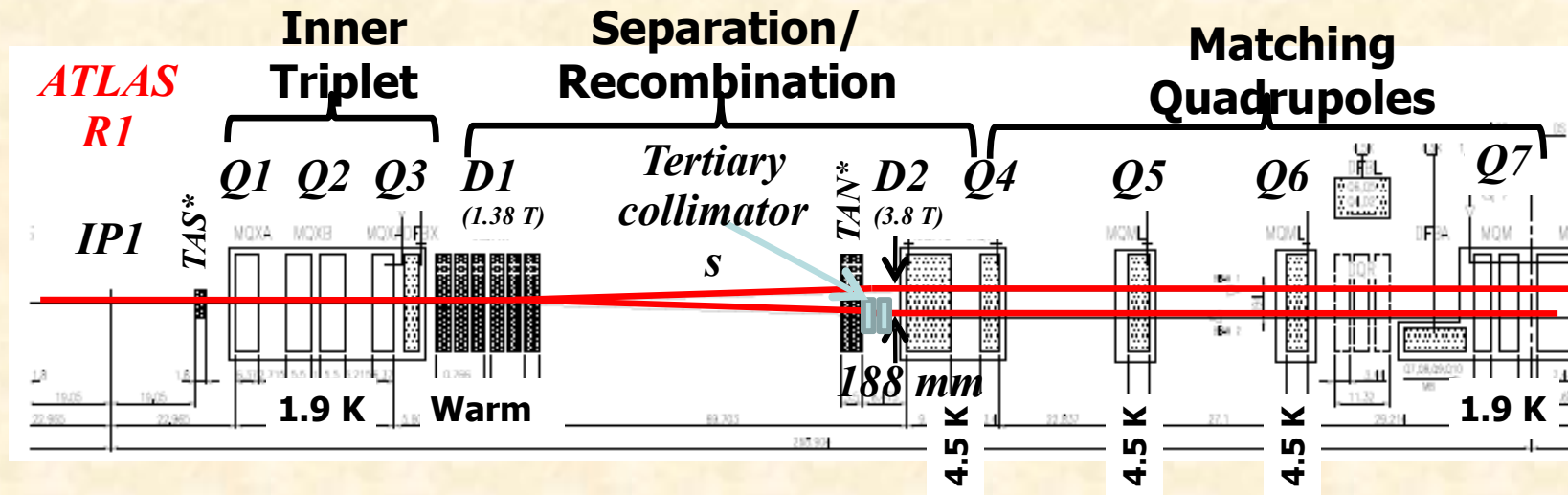
Strahlstrom

$$L = \frac{1}{4\pi e^2 f_0 n_b} * \frac{I_{p1} I_{p2}}{\sigma_x \sigma_y}$$

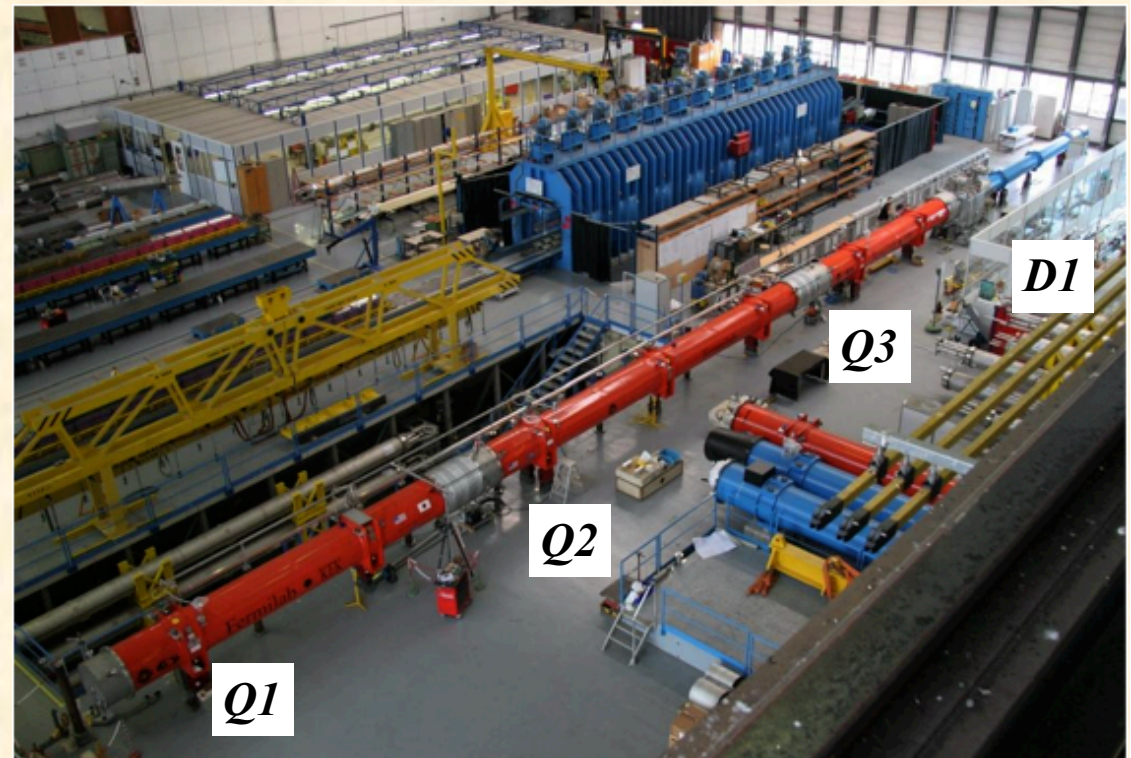
$$L = 1.0 * 10^{34} \frac{1}{\text{cm}^2 \text{s}}$$



The LHC Mini-Beta-Insertions

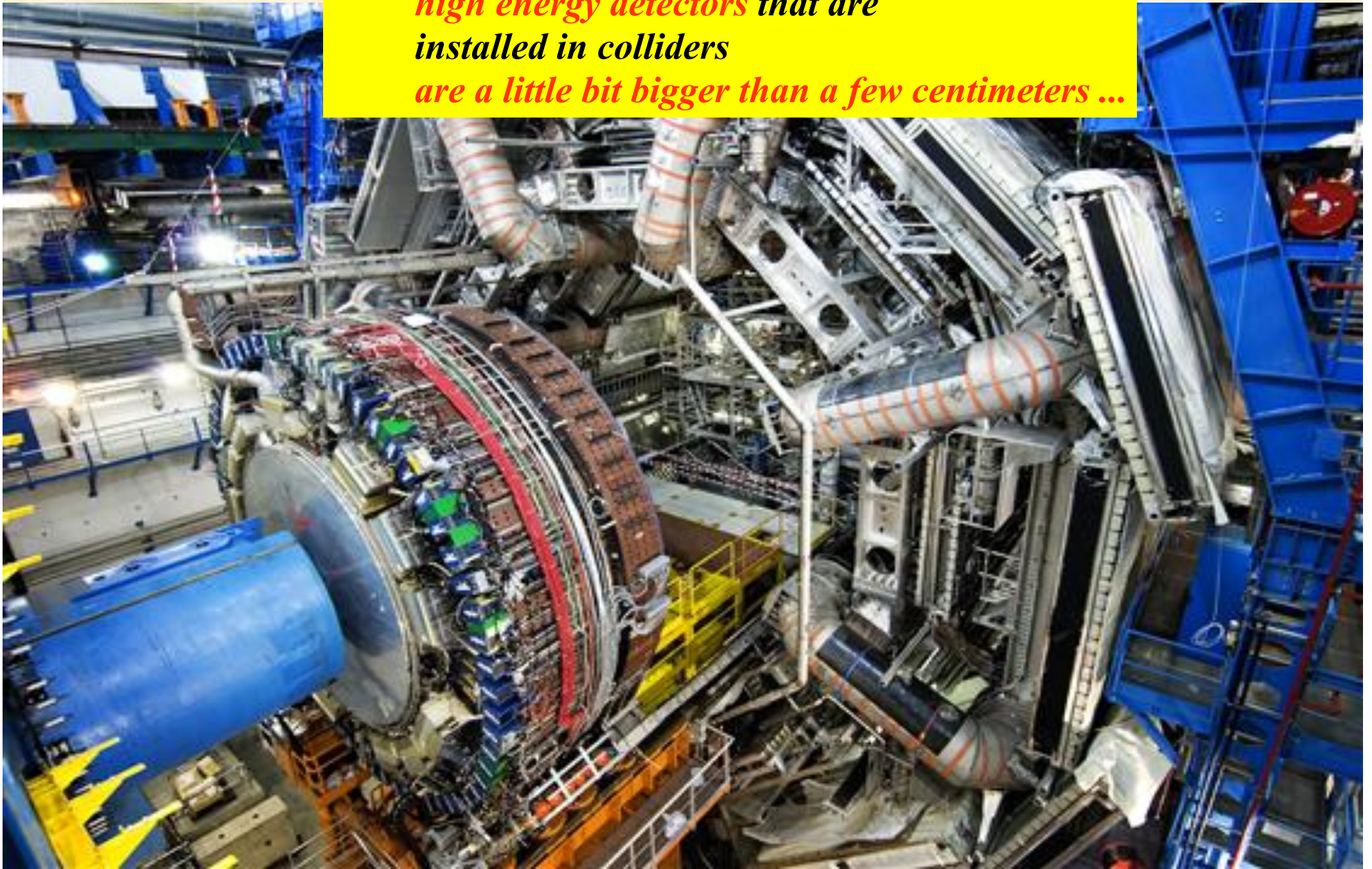


*Extrem starke Fokussierung
(in beiden Ebenen) für beide Strahlen, um
die Trajektorien der 10^{11} Teilchen auf
micro Meter zu komprimieren.*



... clearly there is another problem !!!

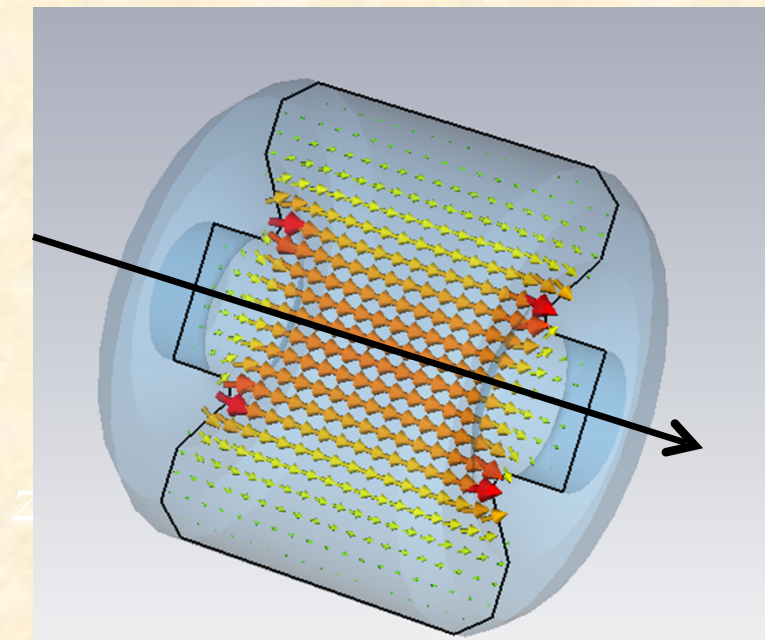
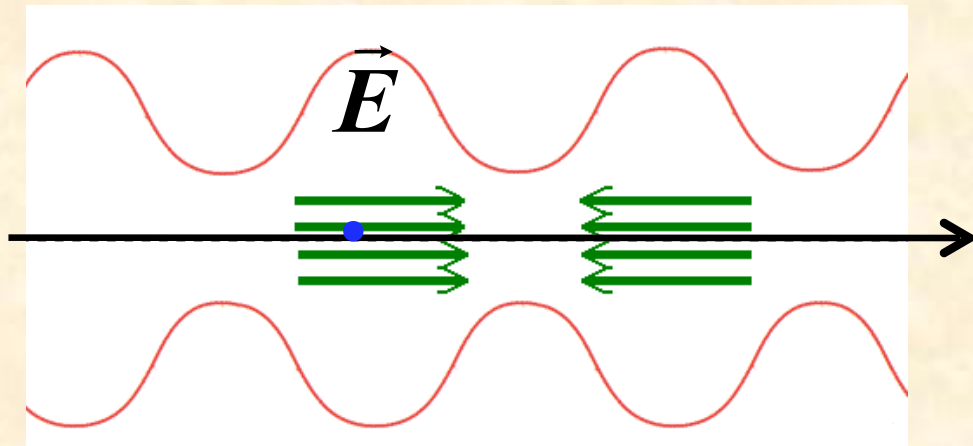
... unfortunately ... in general
high energy detectors that are
installed in colliders
are a little bit bigger than a few centimeters ...



III. The Acceleration

Where is the acceleration?

Install an RF accelerating structure in the ring:

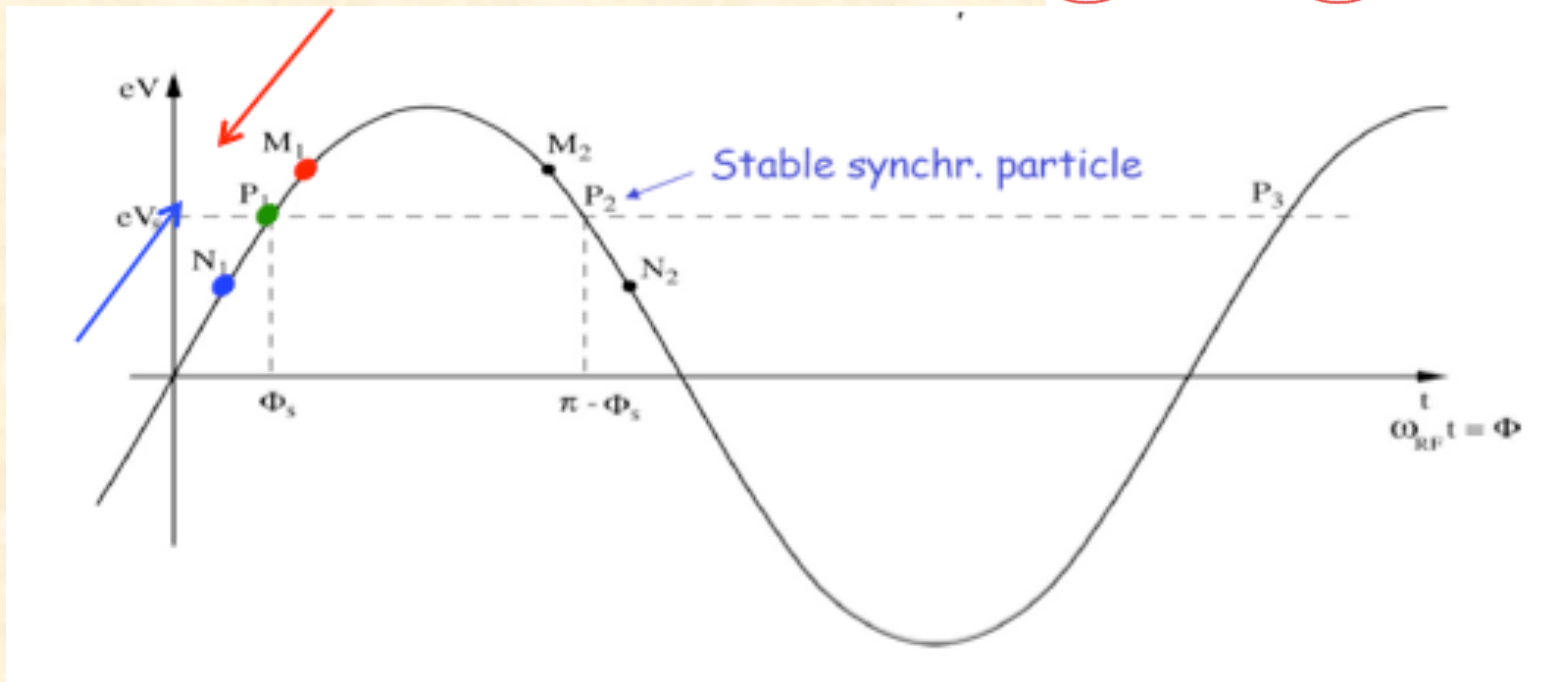
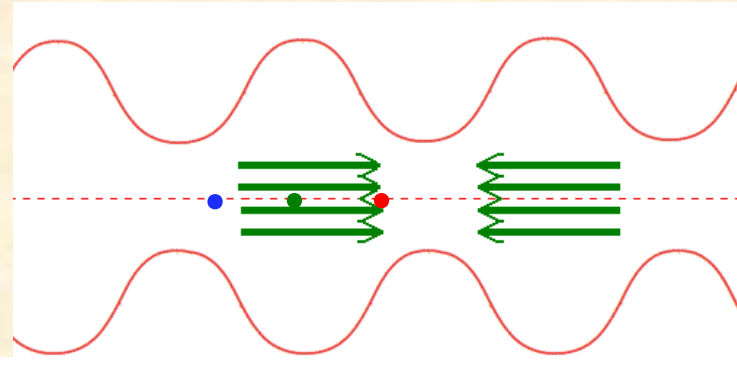


*B. Salvant
N. Biancacci*

The Acceleration & "Phase Focusing"

$\Delta p/p \neq 0$ below transition

- ideal particle* •
- particle with $\Delta p/p > 0$* • *faster*
- particle with $\Delta p/p < 0$* • *slower*



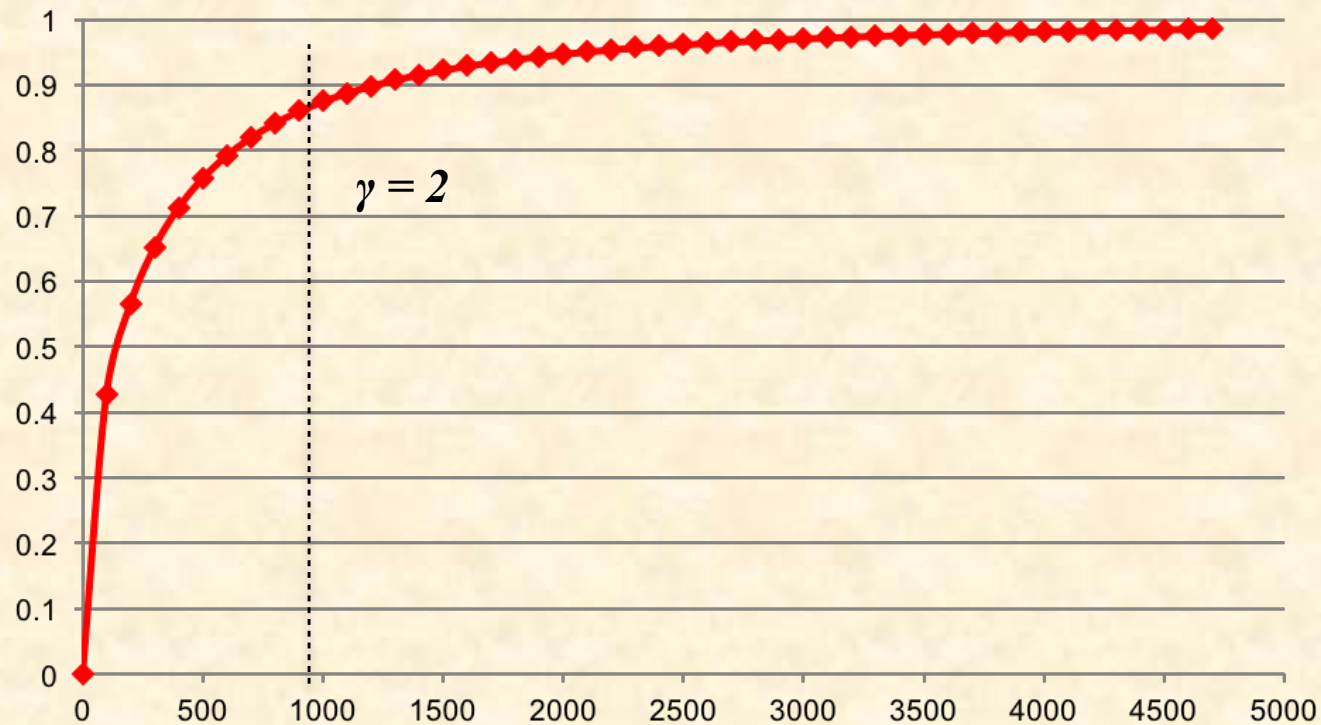
Focussing effect in the longitudinal direction keeping the particles close together ... forming a "bunch"

oscillation frequency: $f_s = f_{rev} \sqrt{-\frac{h\alpha_s * qU_0 \cos\phi_s}{2\pi E_s}} \approx \text{some Hz}$

... so sorry, here we need help from Albert:

$$\gamma = \frac{E_{total}}{mc^2} = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} \rightarrow \frac{v}{c} = \sqrt{1 - \frac{mc^2}{E^2}}$$

v/c

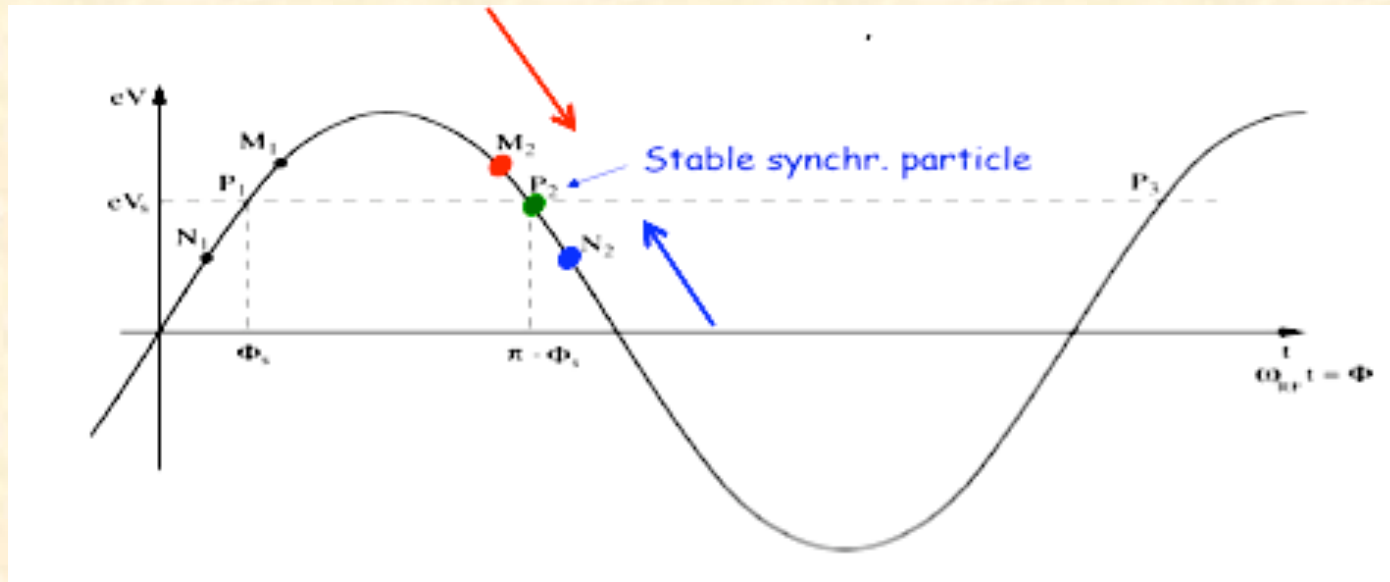
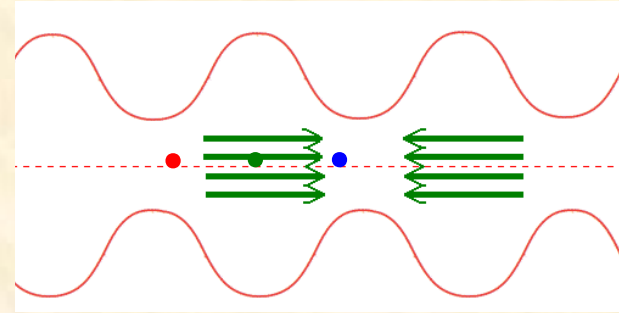


... some when the particles do not get faster anymore

.... but heavier !

The Acceleration above transition

- *ideal particle*
- *particle with $\Delta p/p > 0$* heavier
- *particle with $\Delta p/p < 0$* lighter

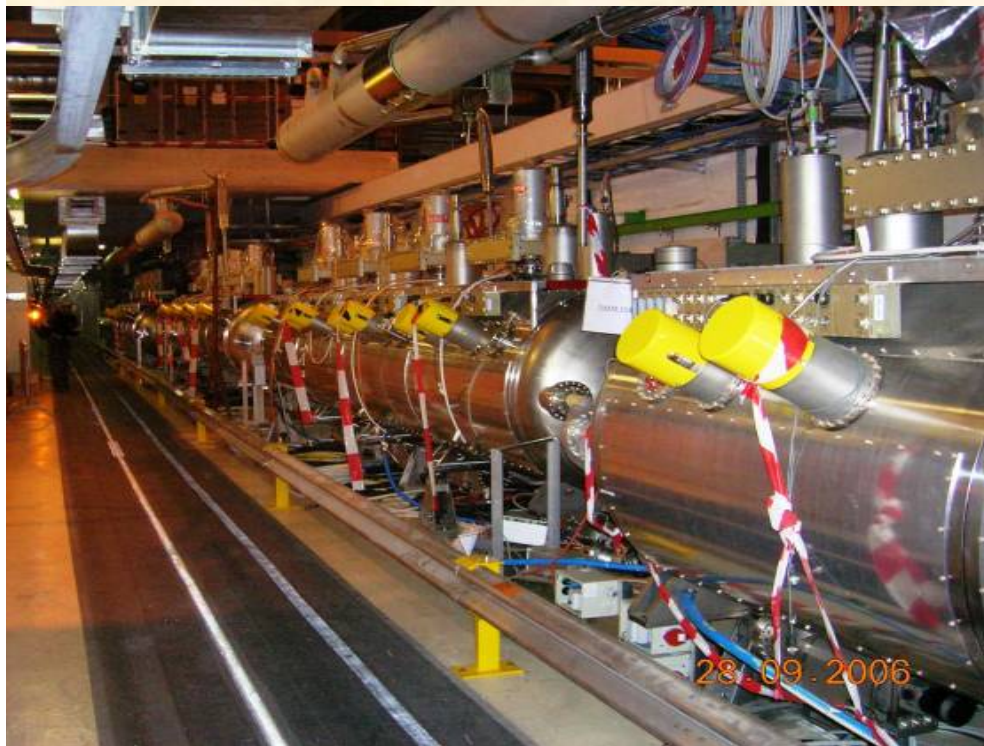
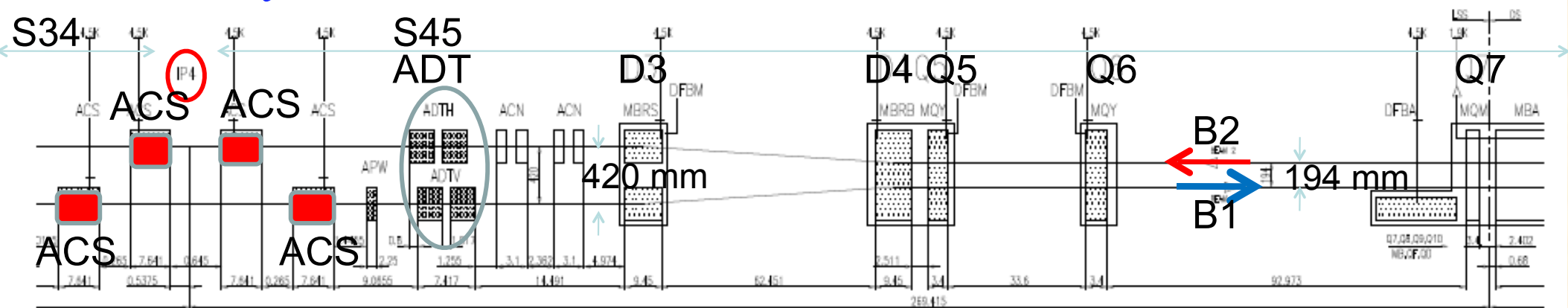


*Focussing effect in the longitudinal direction
keeping the particles close together ... forming a “bunch”*

... and how do we accelerate now ???

with the dipole magnets !

The RF system: IR4



*Nb on Cu cavities @4.5 K (=LEP2)
Beam pipe diam.=300mm*

Bunch length (4σ) *ns* *1.06*

Energy spread (2σ) *10^{-3}* *0.22*

Synchr. rad. loss/turn *keV* *7*

Synchr. rad. power *kW* *3.6*

RF frequency *MHz* *400*

Harmonic number *35640*

RF voltage/beam *MV* *16*

Energy gain/turn *keV* *485*

Synchrotron frequency *Hz* *23.0*

1.) Where are we ?

- * Standard Model of HEP*
- * Higgs discovery*

Merci