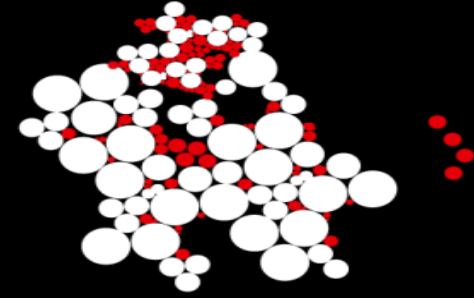


UNIVERSITY OF TWENTE.



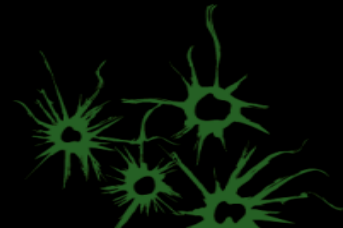
European Course of Cryogenics 2022



Thursday 25/08

Superconductivity & Cryogenics

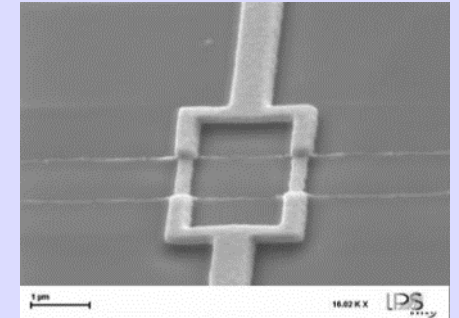
Dr. Marc Dhallé, EMS UTwente



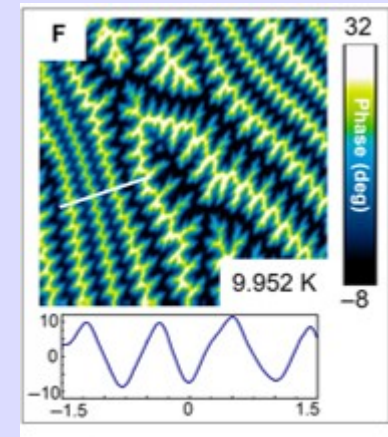
1. Phenomenology & understanding
2. Materials & applications
3. Cooling requirements & - strategies



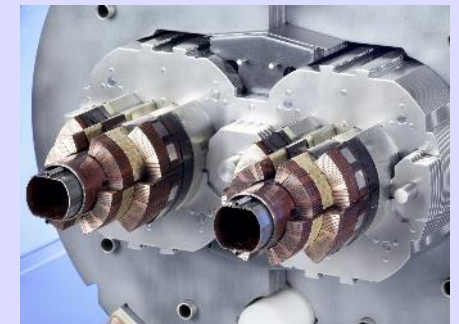
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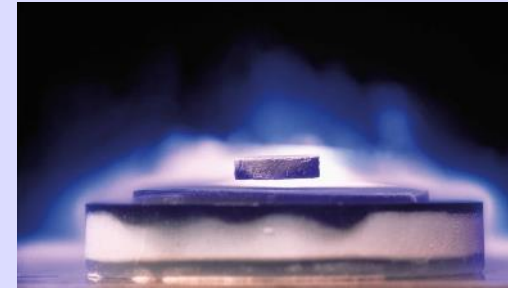
Science Advances



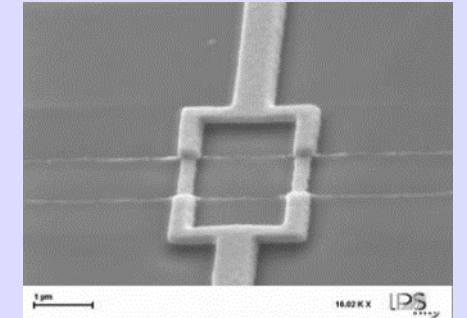
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Want to know more?

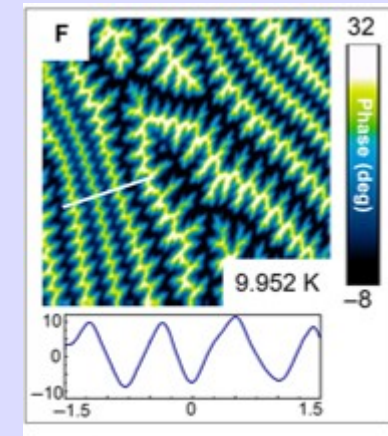
- V.L. Ginzburg and E.A. Andryushin, "*Superconductivity*" (1994)
Excellent layman's primer.
- V.V. Schmidt "*The Physics of Superconductors*" (1997)
- M. Cyrot and D. Pavuna, "*Introduction to Superconductivity and high-T_c materials*" (1992)
- M. Tinkham, "*Introduction to Superconductivity*" (1975)
- P.G. De Gennes, "*Superconductivity of metals and alloys*" (1966)
... and many more books of varying level of detail.



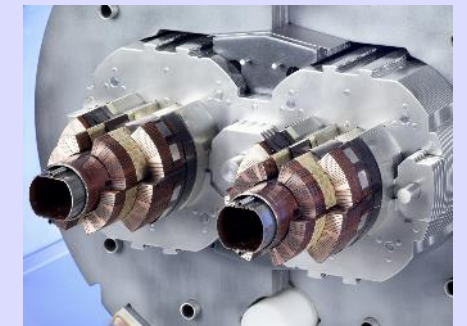
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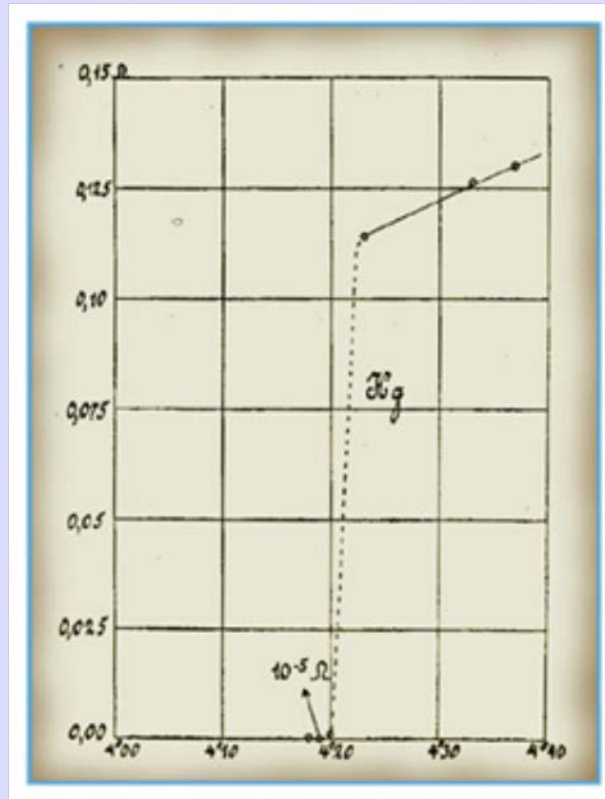
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... or contact us: <https://www.utwente.nl/en/tnw/ems/>

Disappearance of electrical resistivity

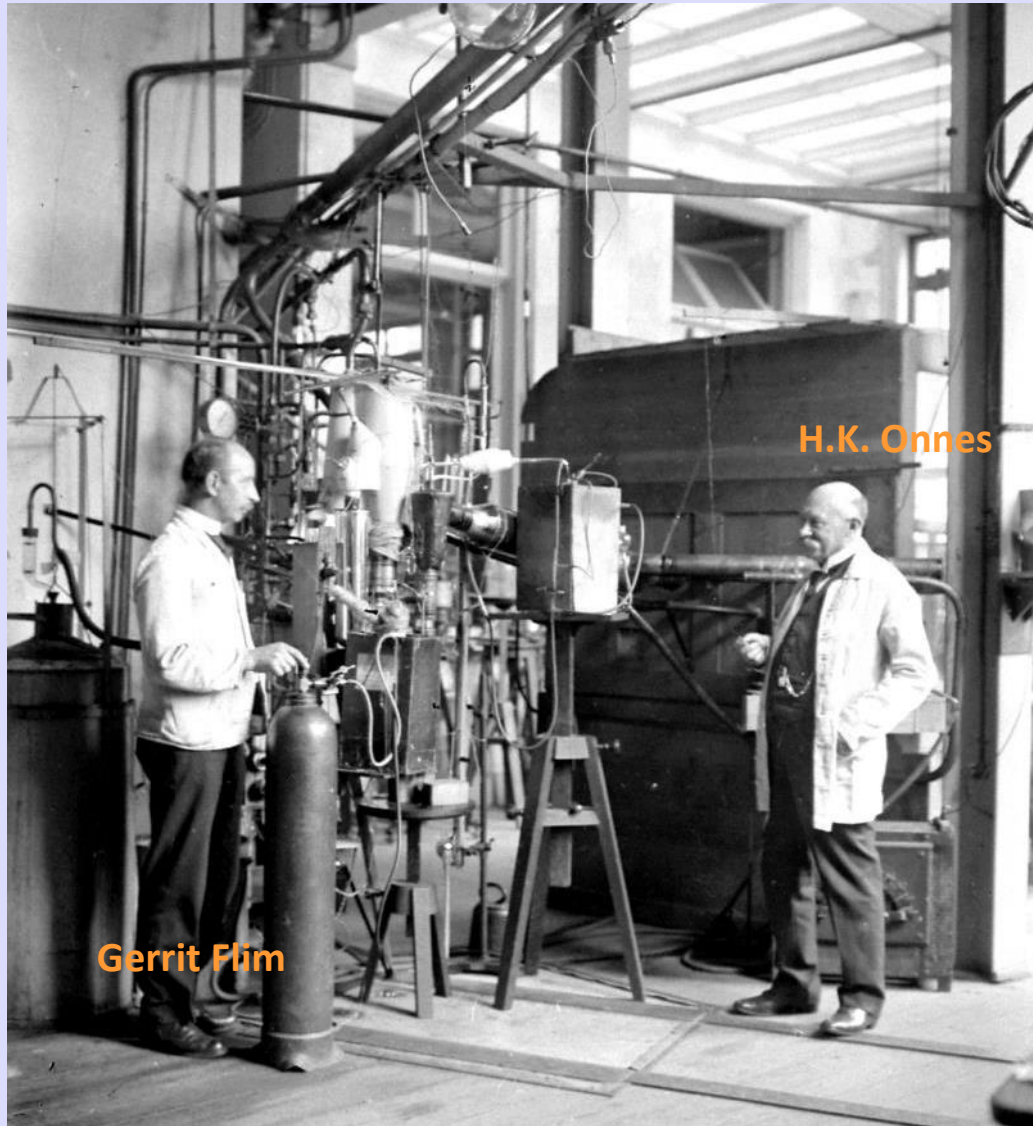


Onnes, 1911

Below a critical temperature T_c :

$$\rho = 0$$

No scattering of charge carriers at lattice imperfections



H.K. Onnes

Gerrit Flim

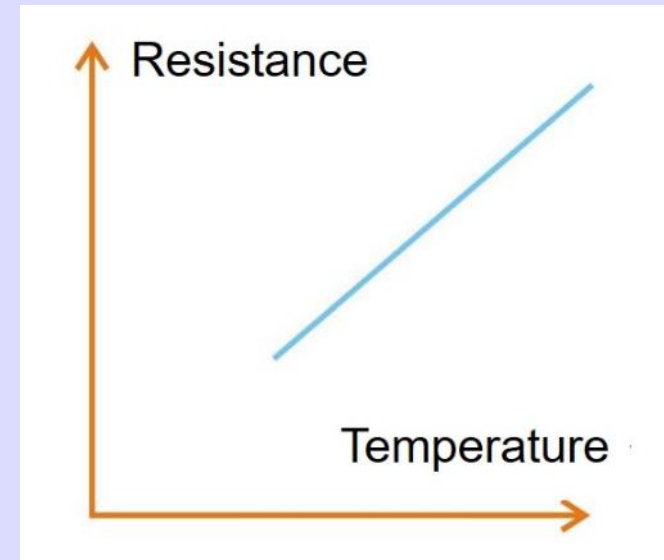
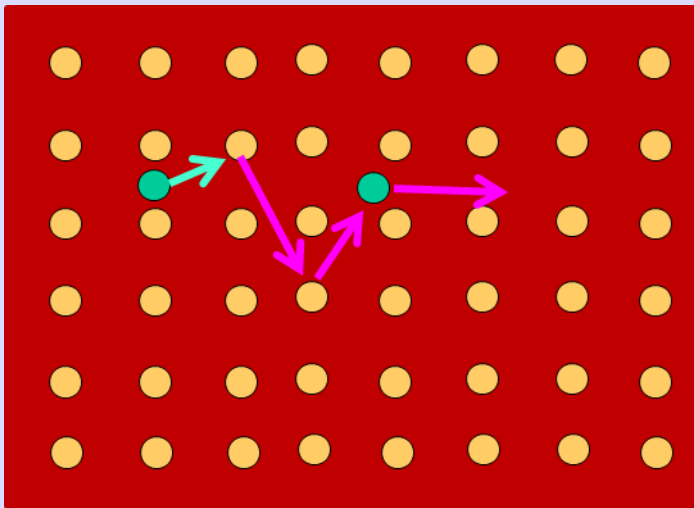
Liquid helium
4.2 Kelvin (1908)



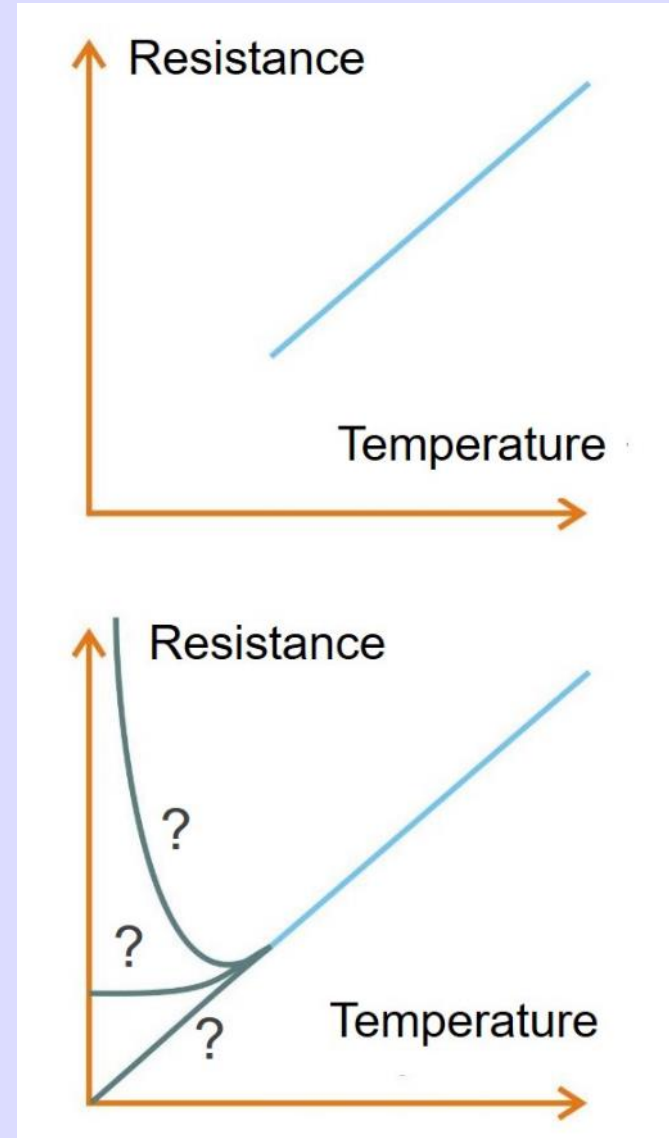
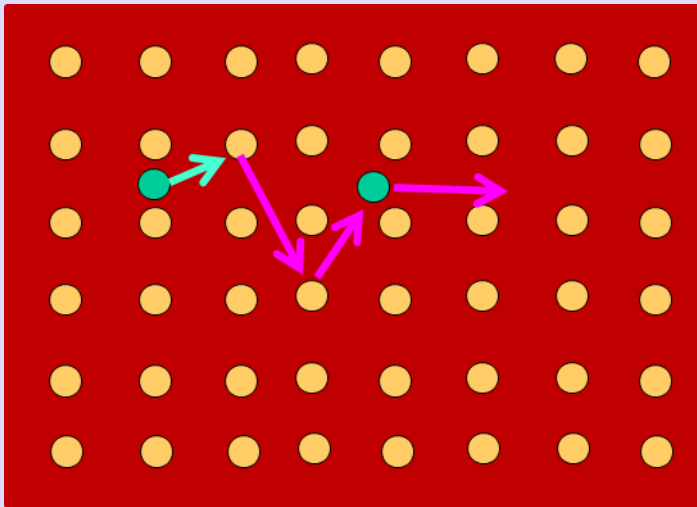
<https://www.youtube.com/watch?v=9FudzqfpLLs>

Lorentz.leidenuniv.nl

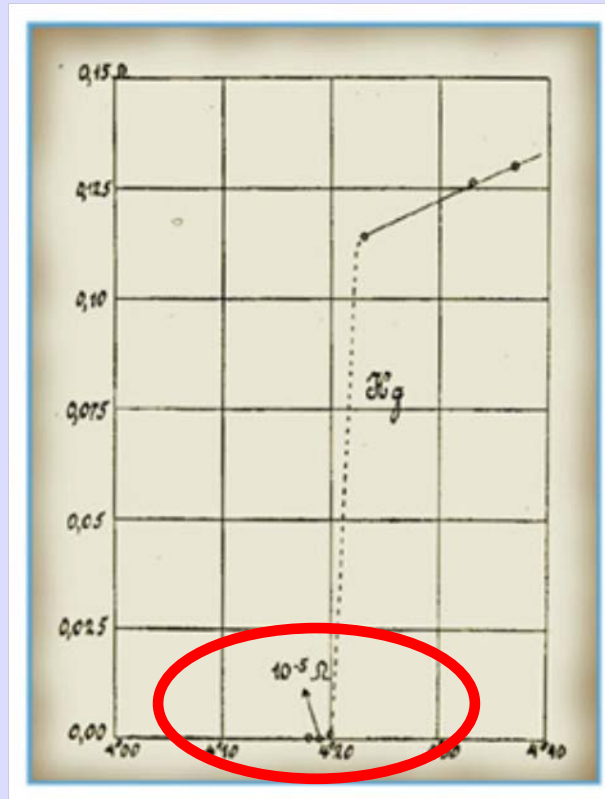
Electrical conductivity & resistivity



Electrical conductivity & resistivity



Disappearance of electrical resistivity



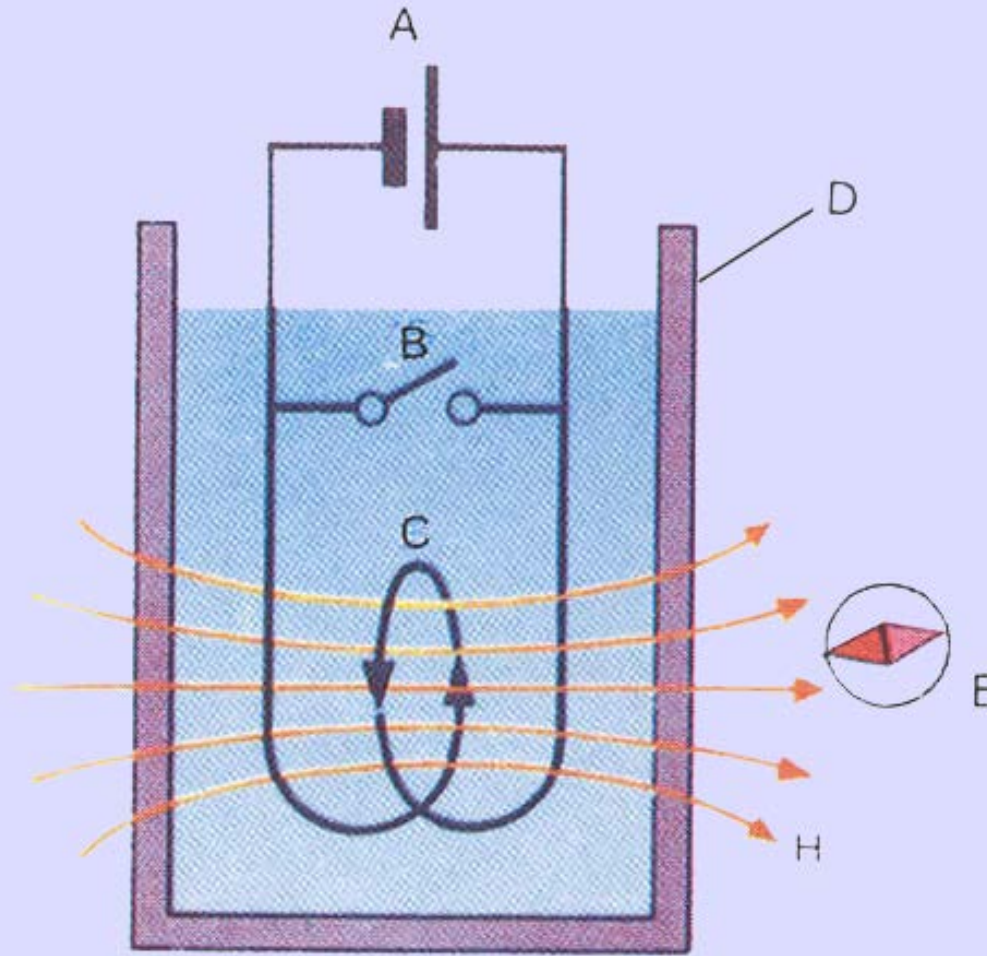
Onnes, 1911

Below a critical temperature T_c :

$$\rho = 0$$

No scattering of charge carriers at lattice imperfections

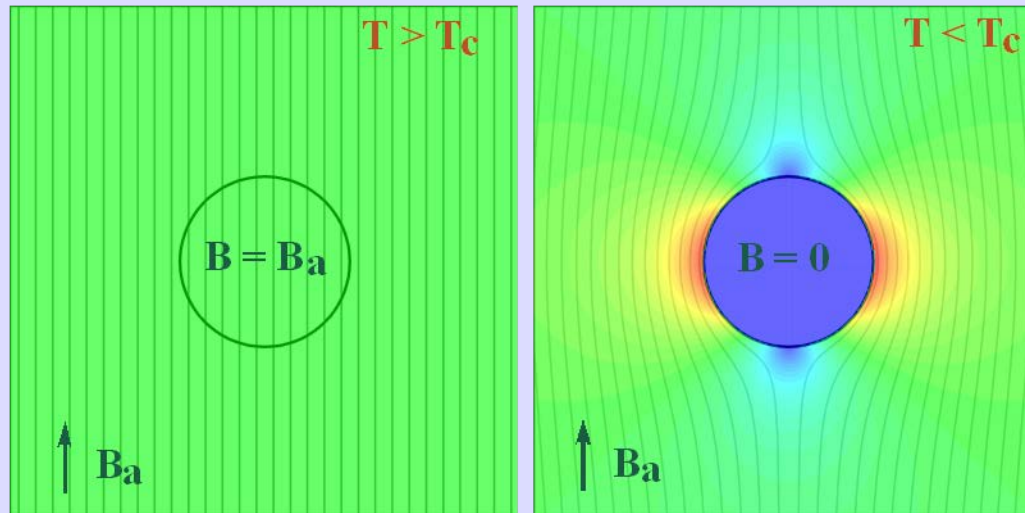
Disappearance of electrical resistivity



Resistance is at least 10^{17} times smaller than that of copper

Ginzburg & Andryushin, 1994

Meissner effect

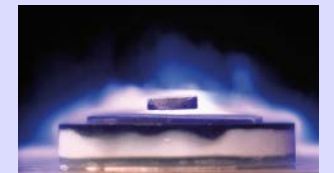


Perfect diamagnetism

$$B = 0$$

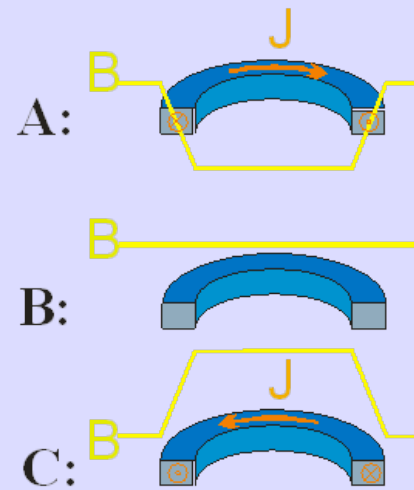
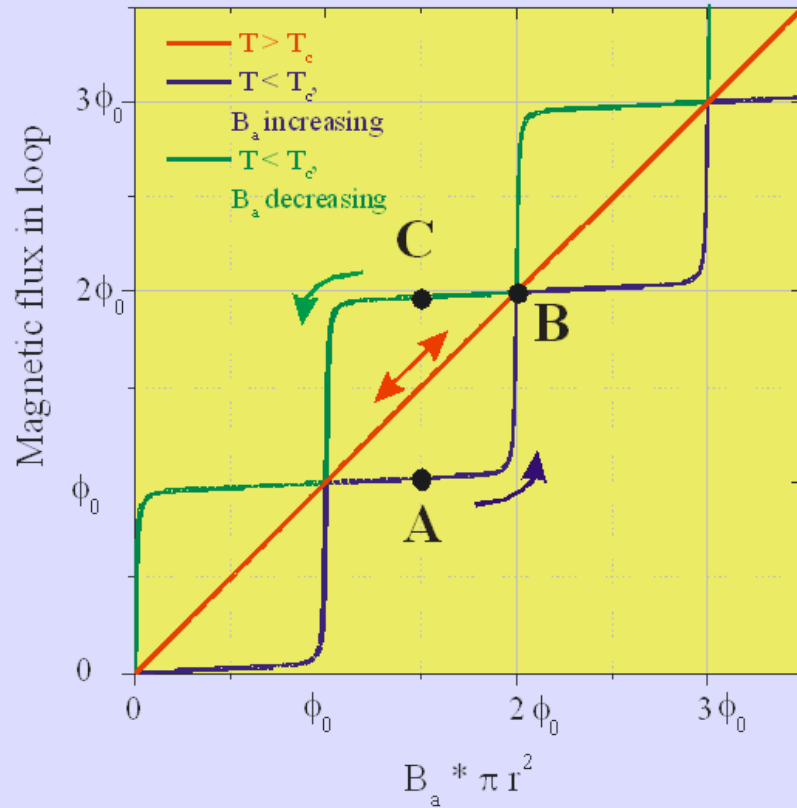
Magnetic fields are screened out by surface currents (*)

W. Meissner and R. Ochsenfeld, 1933



(*) up to a critical field H_{c1}

Flux Quantization



Only integer multiples

$$\Phi = n\phi_0$$

of the flux quantum

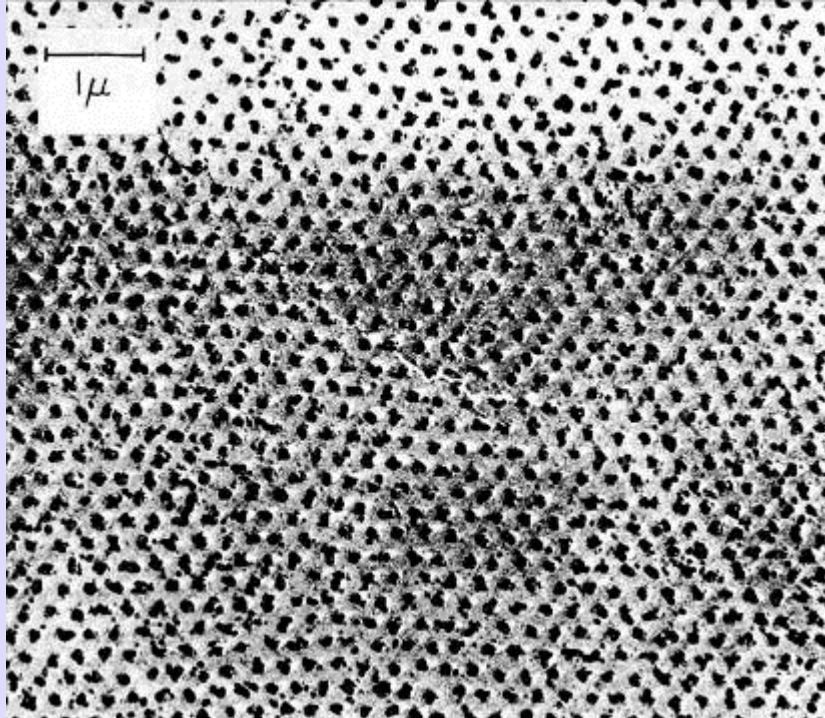
$$\phi_0 = \frac{h}{2e}$$

are allowed inside a superconducting ring.

F. London, 1950

W. A. Little and R.D. Parks, 1962

Flux Quantization



A.A. Abrikosov, 1957
Essman and Traube, 1967

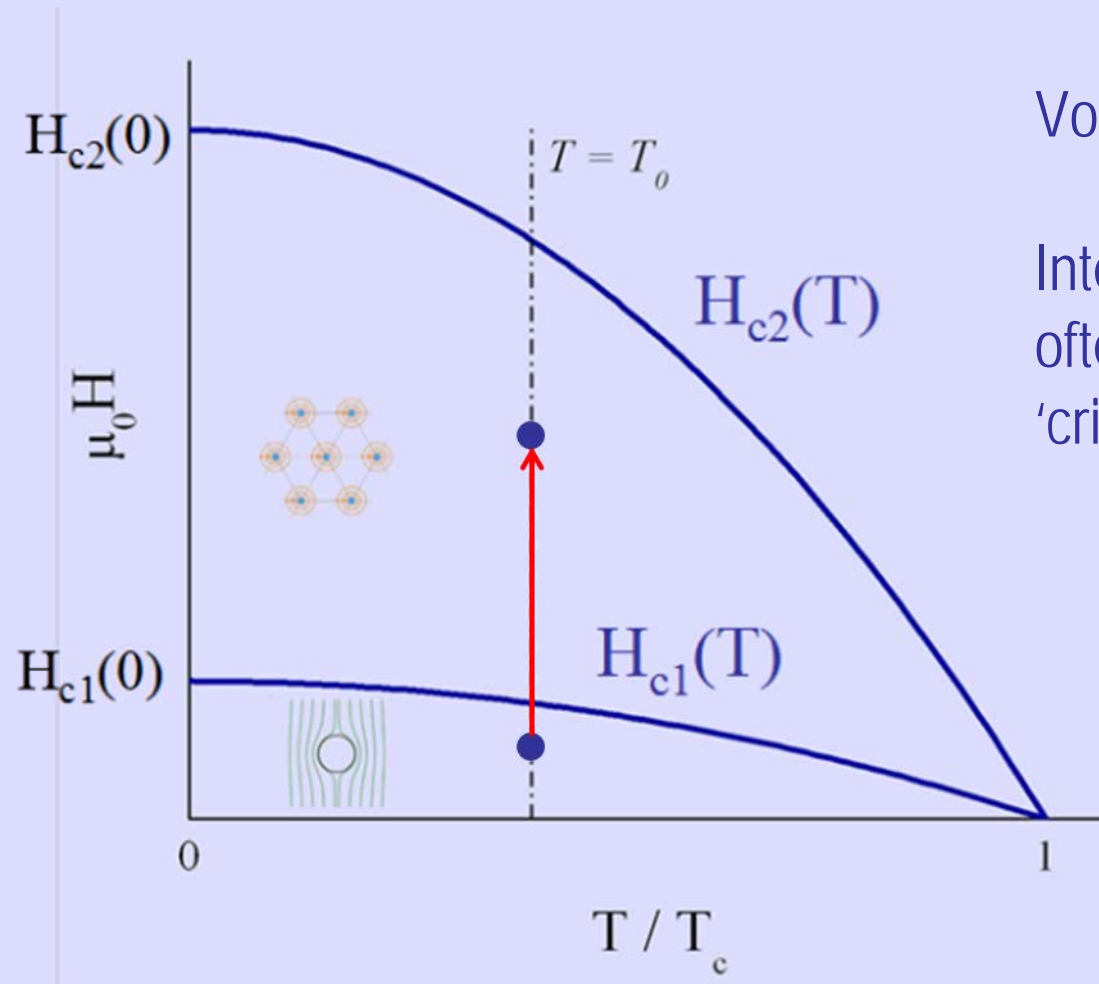
For “type II” superconducting materials:

$$B = n A \phi_0$$

Fields higher than H_{c1} are admitted in the form of mutually repulsive flux tubes ^(*), generated by current vortices

(*) up to a second *critical field* H_{c2}

Flux Quantization



Vortex- or 'Abrikosov'- state.

Interactions between vortices and current often determine maximal or 'critical' current density J_c

Josephson effect

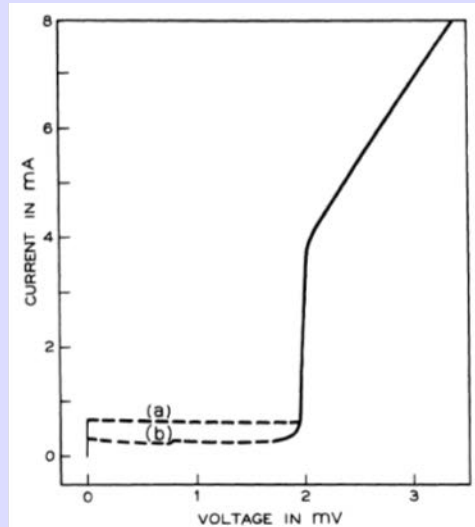
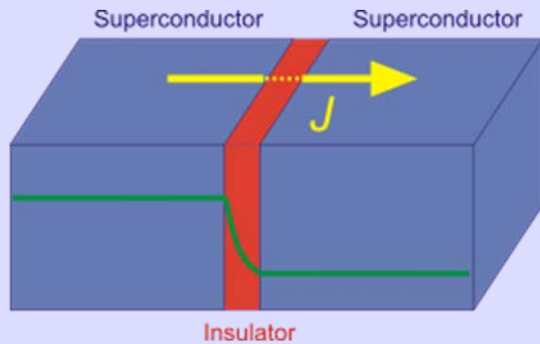


FIG. 1. Current-voltage characteristic for a tin-lead tunnel structure at ~1.5°K, (a) for a field of 6×10^{-3} gauss and (b) for a field 0.4 gauss.

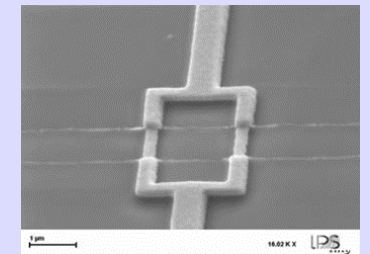
Current can 'leak' through a barrier:

$$J_c = J_0 \sin \delta$$

$$\delta(t) = \delta(0) + \frac{Vt}{\Phi_0}$$

The maximum tunneling current is modulated by a "phase difference" across the barrier that depends on magnetic field.

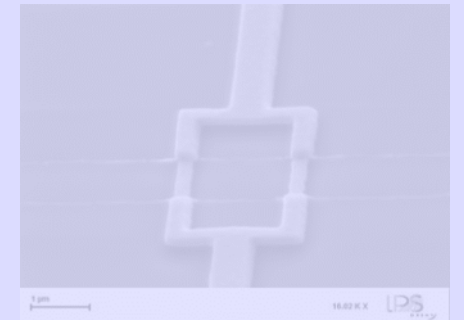
B.D. Josephson, 1962
P.W. Anderson and J.M.Rowell, 1963



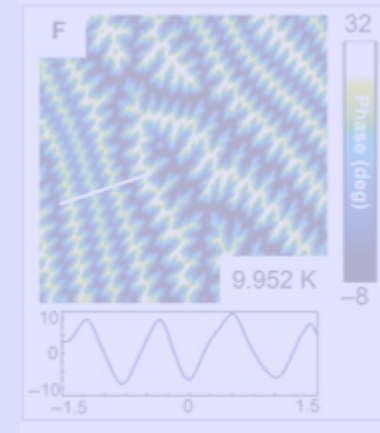
1. Phenomenology & *understanding*
2. Materials & applications
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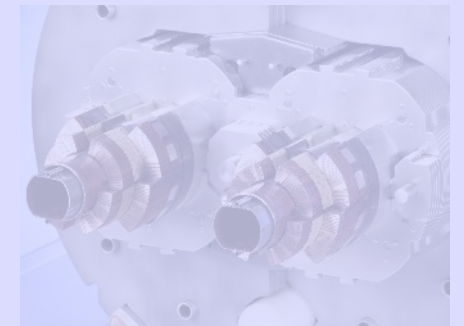
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The London theory

H. London and F. London, 1935

$$\lambda^2 (\nabla \times J_s) + H = 0$$

Equation of motion for “superelectrons”

→ penetration depth λ

$$n_s = n_0 \left(1 - T/T_c\right) ; \quad n_n = n_0 T/T_c$$

“2 fluid”-model

→ kinetic inductance, surface impedance

Ginzburg-Landau theory

V.L. Ginzburg and L.D. Landau, 1950

$$\psi(\vec{r}) = |\psi(\vec{r})| e^{i\varphi(\vec{r})} \quad \text{with} \quad |\psi|^2 = n_s/2$$

complex "order parameter"

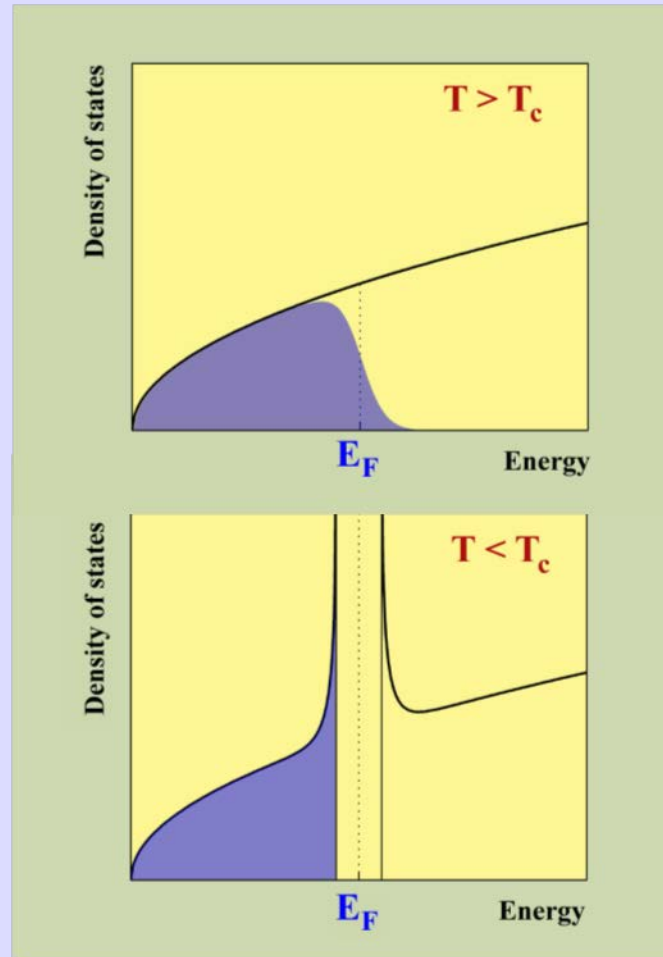
$$\Delta F = F_S - F_N = \alpha |\psi|^2 + \beta |\psi|^4 + \frac{1}{2m} |(i\hbar\nabla + 2e\vec{A})\psi|^2 + \frac{\mu_0 H^2}{2}$$

Minimization of free energy $F \rightarrow$

flux quantization; coherence length ξ ; penetration depth; vortex state

BCS theory (1)

L.N. Cooper, 1956

**Bound Electron Pairs in a Degenerate Fermi Gas***

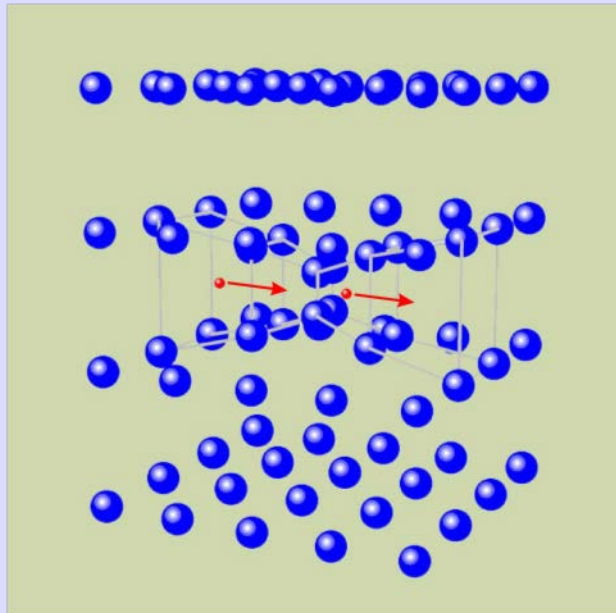
LEON N. COOPER

Consider a pair of electrons which interact above a quiescent Fermi sphere with an interaction of the kind that might be expected due to the phonon and the screened Coulomb fields. If there is a net attraction between the electrons, it turns out that they can form a bound state, though their total energy is larger than

In the presence of an attractive interaction, electrons can 'team up', they form "Cooper pairs".

Cooper pairs are bosons.

To break them up, one needs to overcome an energy gap Δ .

BCS theory (2)*J. Bardeen, L.N. Cooper, J.R. Schrieffer 1957***Theory of Superconductivity***J. BARDEEN, L. N. COOPER,[†] AND J. R. SCHRIEFFER[‡]

We shall call the interaction, H_2 , between electrons resulting from the electron-phonon interaction the "phonon interaction." This interaction is attractive when the energy difference, $\Delta\epsilon$, between the electron states involved is less than $\hbar\omega$. Diagonal or self-energy

The attractive interaction comes about through interactions with lattice vibrations ("*phonons*").

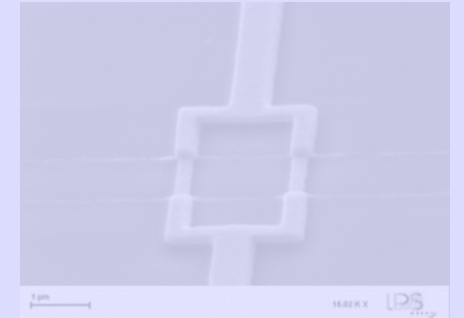
- T_c (isotope effect), critical field H_c
- Gap Δ
- Cooper pair density
- penetration depth λ , coherence length ξ



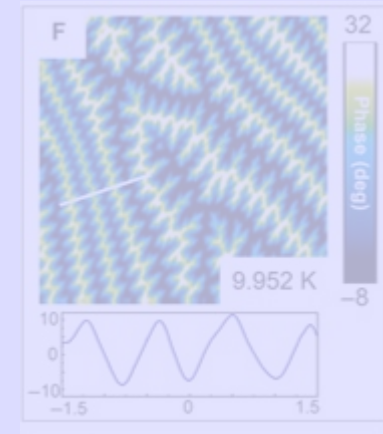
1. Phenomenology & understanding
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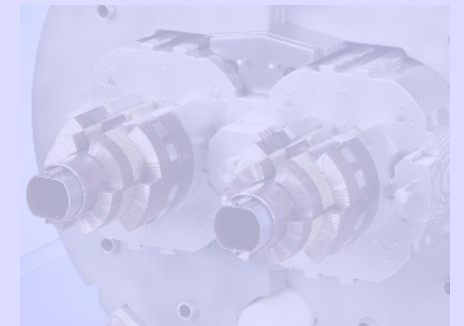
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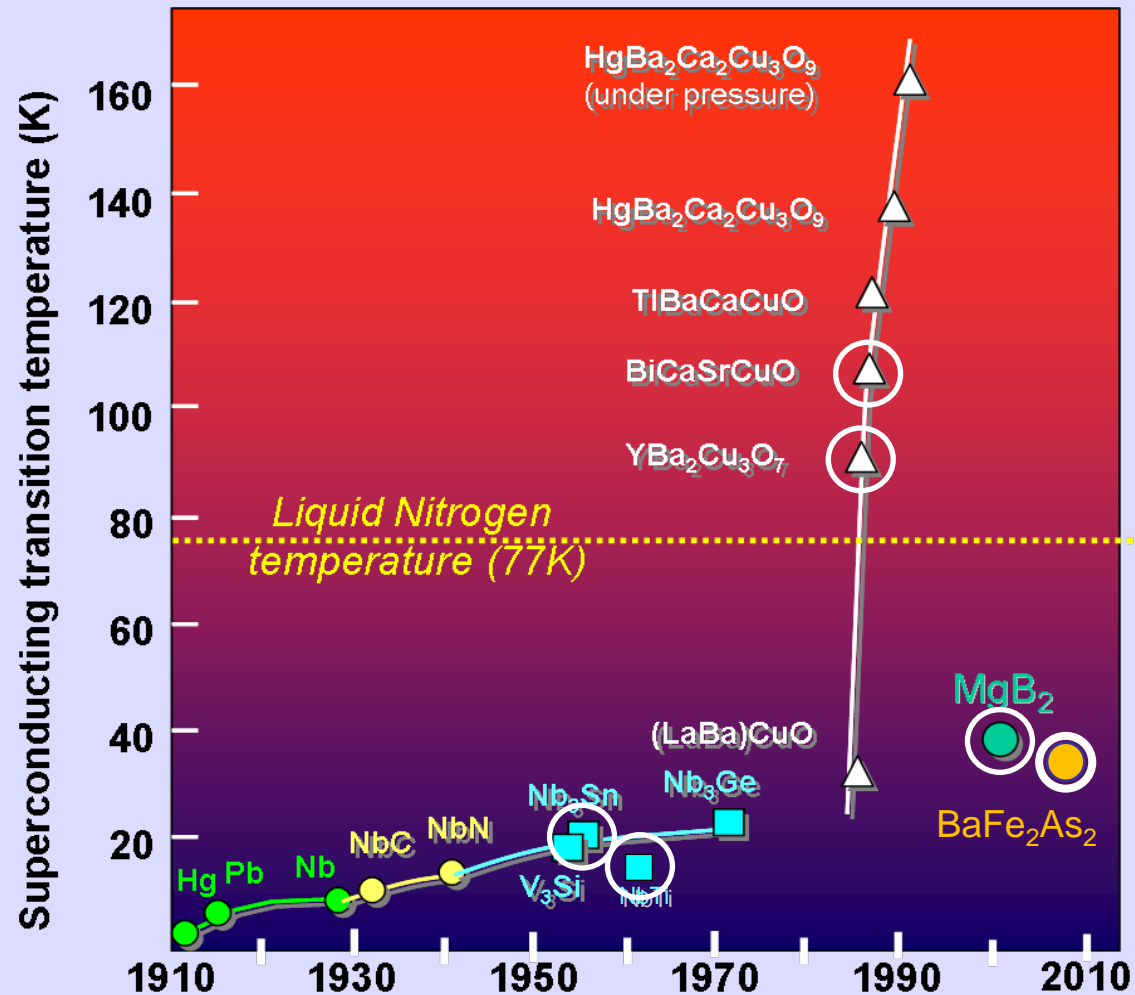


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The history of superconducting materials (only a selected materials are shown)

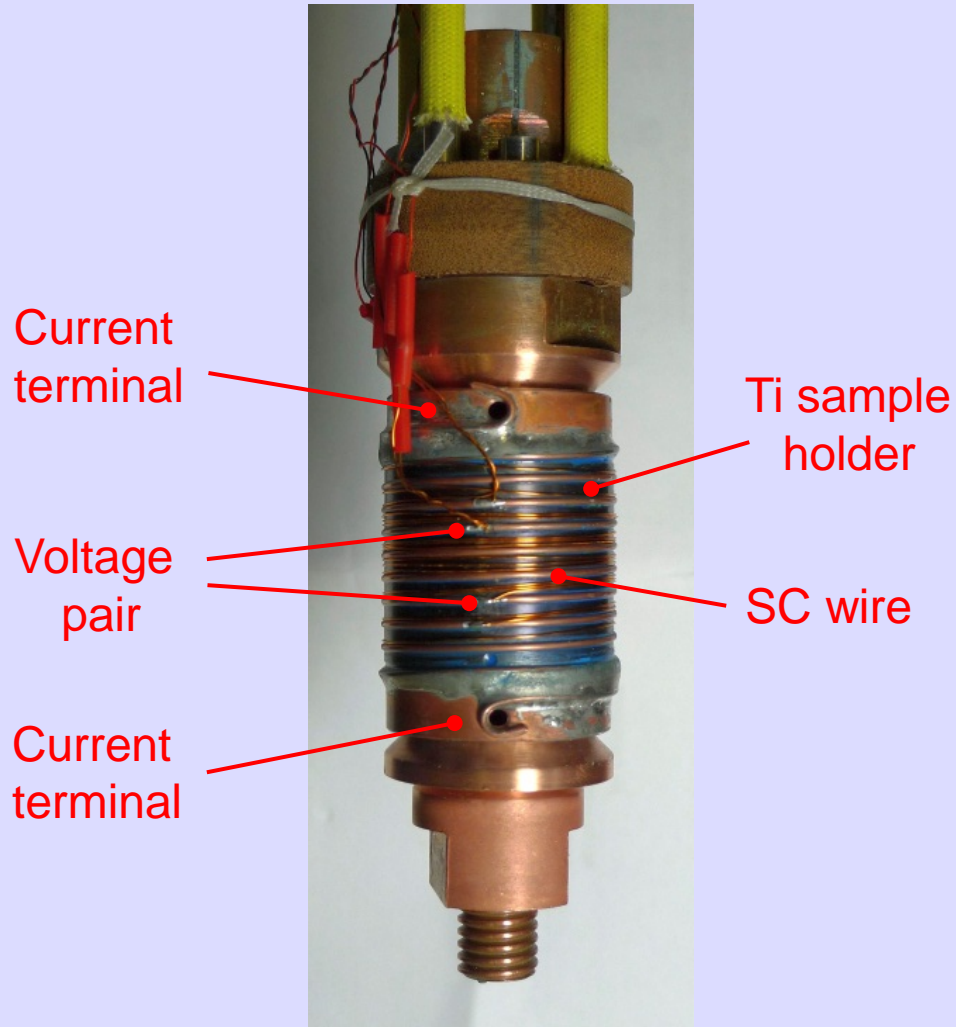


○ : Materials that are developed / commercialized into **practical** superconductors

"Practical" = >1 km long wires (*) with stable and uniform properties that can be used on coil windings

(*) see lecture 2

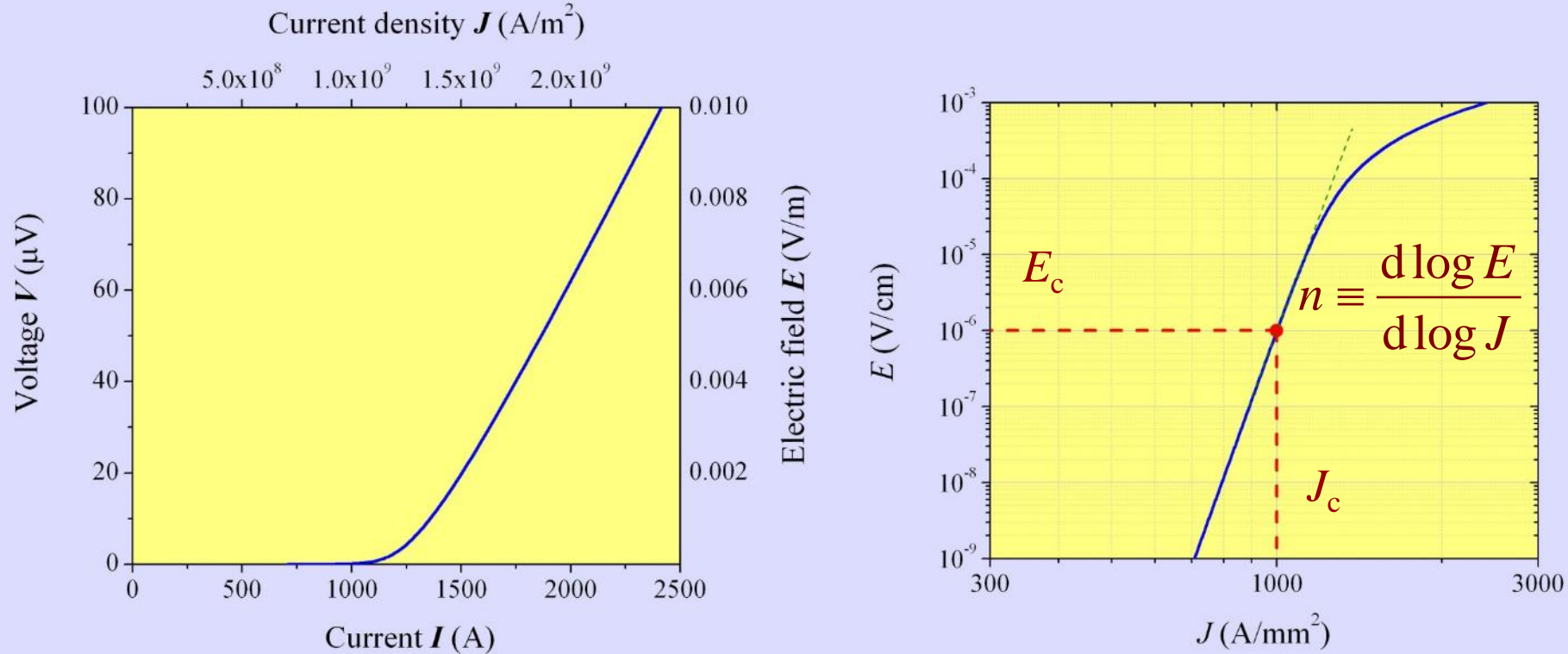
Intermezzo



- Probing the limit of loss-less charge transport: the critical current I_c ;
the critical current density J_c
- Place the sample holder in a controlled T - and H - environment (often a SC solenoid in liquid or gaseous He)
- At each desired T - and H - value, slowly increase the current I and record the voltage V across the test section (IV-measurement)

Voltage-current relation of a superconductor: "sudden" appearance of voltage

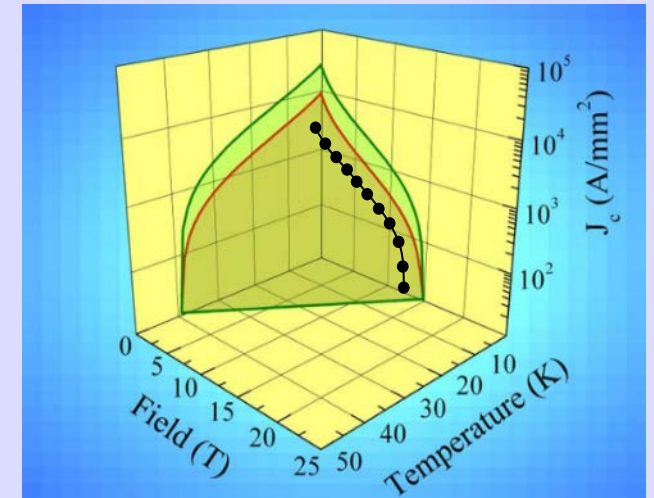
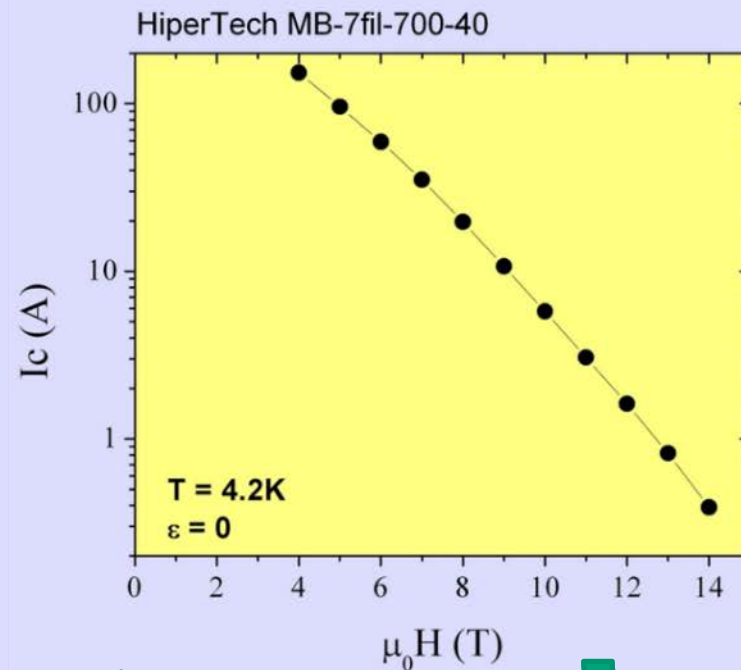
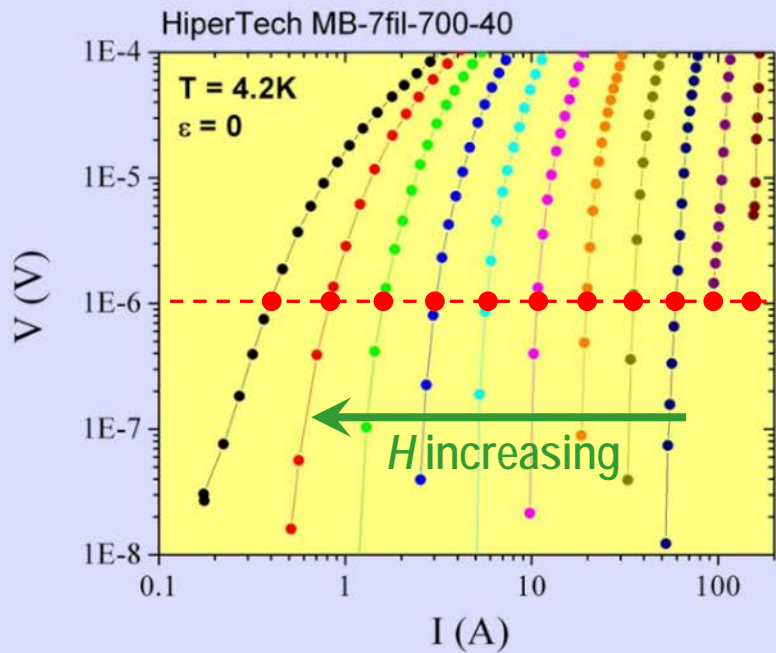
~ power-law relation (curved on lin-lin VI scale, but linear on log-log VI scale)



Note:

1. Not infinitely steep ($n < \infty$); needs **critierion** (E_c or ρ_c)
2. Use of E and J instead of V and I

Measure V - I curves across a range of magnetic fields H and temperatures T , extract $I_c(B, T)$ and from this the critical surface $J_c(T, H)$ of a superconductor



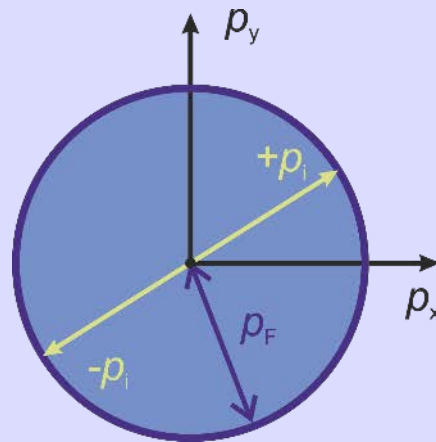
Depairing current:

- $10^4 - 10^6$ A/mm² ;
- Fundamental upper limit;
- Never reached in practical SC

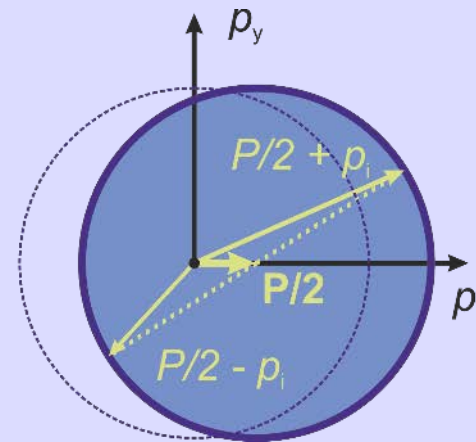
$$J_{c,depairing} = 2n_s e \Delta / p_F$$

Loss-less transport = bodily shifting of Fermi surface in k-space

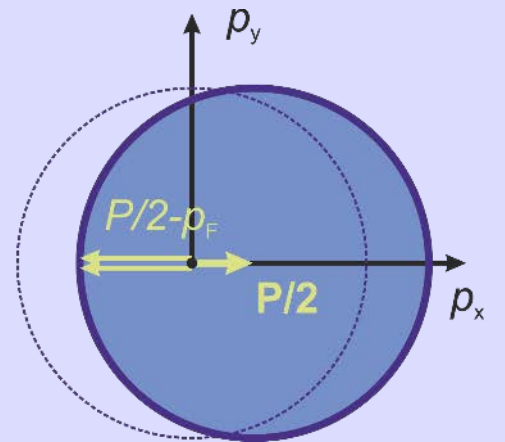
$$E_0 = \frac{1}{2m} \left[\left(\mathbf{P}/2 + \mathbf{p}_i \right)^2 + \left(\mathbf{P}/2 - \mathbf{p}_i \right)^2 \right] = \frac{1}{2m} \left(P^2 / 2 + 2 p_F^2 \right)$$



Non current-carrying



Paired : $E = E_0$



Unpaired : $E = E_0 + 2\Delta - \mathbf{P} \cdot \mathbf{v}_F$

after Rose-Innes & Rhoderick, 1974

Scattering of a pair to the state with lowest available E lowers E_{kin} , but costs condensation energy 2Δ

$$E_{scattered} = 2 \frac{(\mathbf{P}/2 - \mathbf{p}_F)^2}{2m} + 2\Delta(T) = E_0 - \frac{P p_F}{m} + 2\Delta(T)$$

Depinning current:

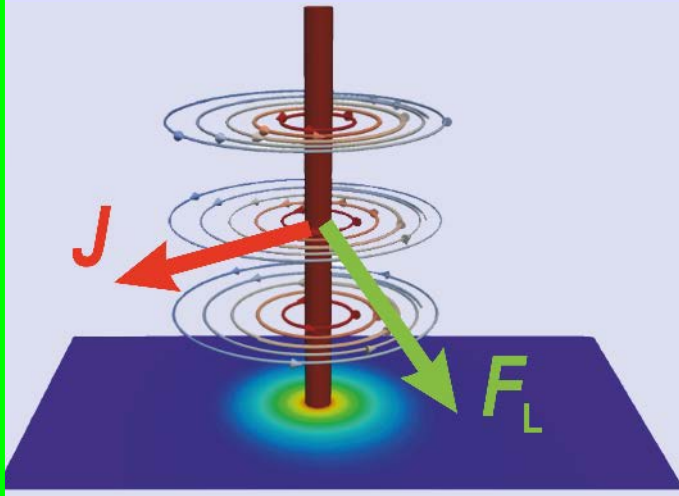
- $10^2 - 10^4$ A/mm² ;
- Sample-dependent;
- Usual limit
in practical SC

Current J + vortices \rightarrow Lorentz force F_L

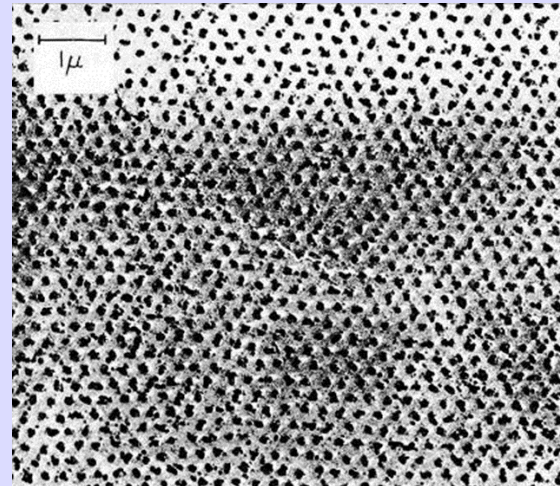
Moving vortex \rightarrow Dissipation

See Maxwell's equations: $\nabla \times \mathbf{B} = \mu \mathbf{J}$, $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$ and $P = \mathbf{E} \cdot \mathbf{J}$

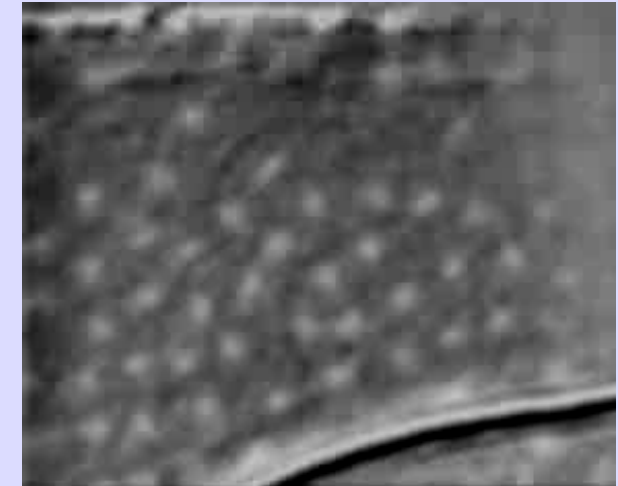
High current density can only be maintained loss-less
when vortices are adequately pinned



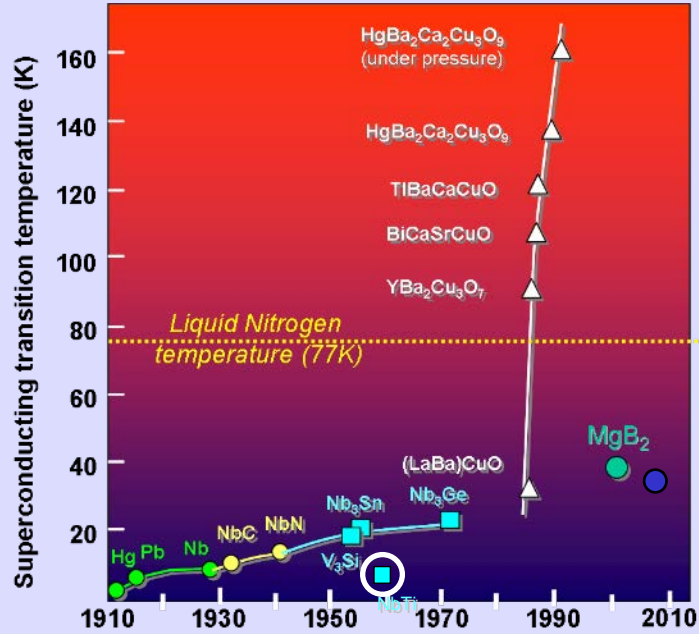
after Glatz et al, 2019



Essman and Traube, 1967



Olsen et al, 2004



NbTi : niobium-titanium alloy

$$T_c = 9\text{ K}$$

$$B_{c2} = 11\text{ T at } 4.2\text{K and } 14\text{ T at } 0\text{K}$$

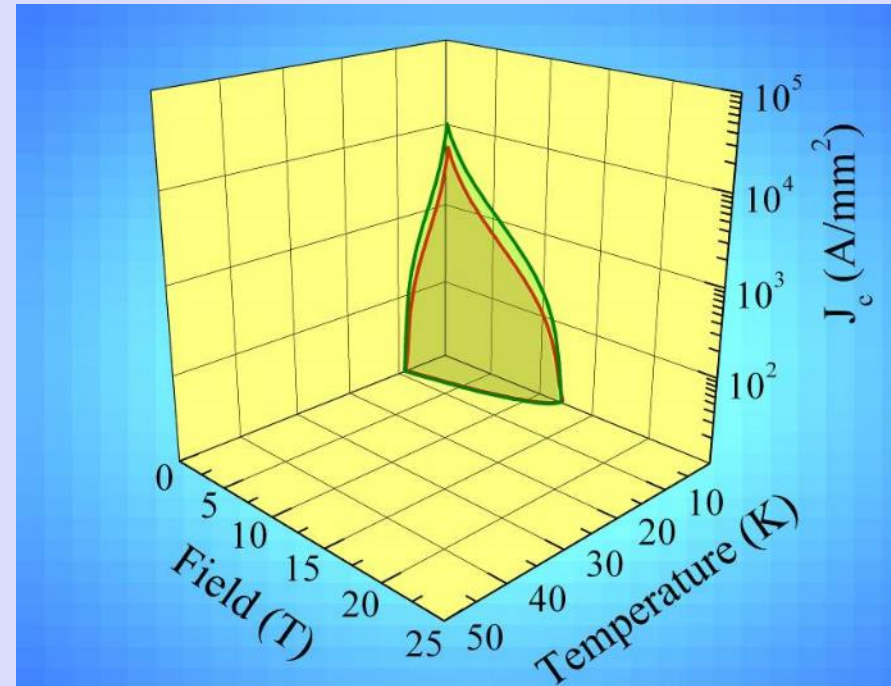
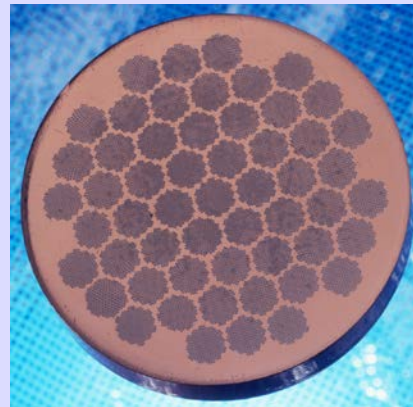
well-developed

costs ~1 €/kA.m at 4.2K, 5T (*)

(*) L. Cooley et al., 2005



Cubic alloy, isotropic



Pinning in NbTi: very finely distributed Ti-precipitates

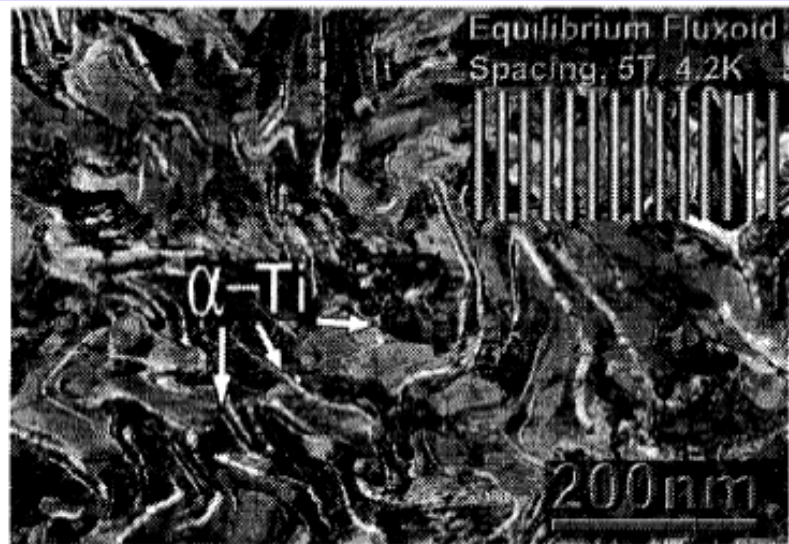
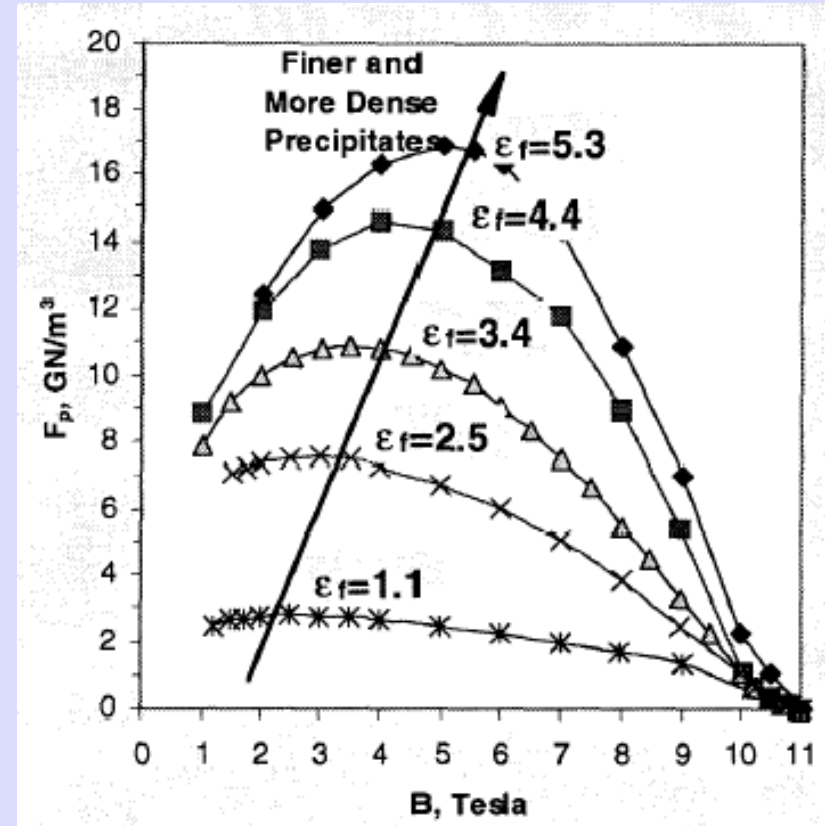
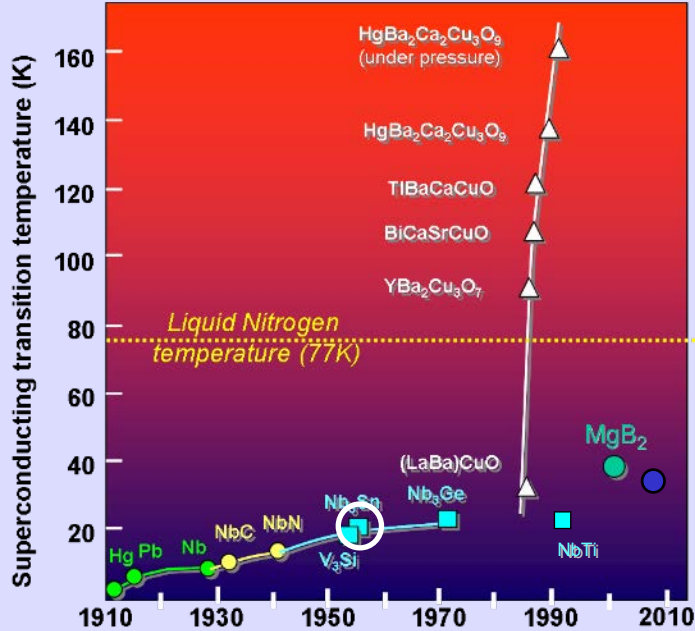


Figure 2 Typical Nb-47wt.%Ti high critical current microstructure (in transverse cross-section) showing the densely folded sheets of α -Ti pinning centers dispersed within the superconducting β -Nb-Ti matrix. Superimposed is a schematic illustration of the equilibrium fluxoid spacing and dimensions appropriate to Nb-47wt.%Ti at 5T, 4.2K.

Larbalestier et al, 1995



$$F_p = J_c \times B = \text{“maximum pinning force”}$$



Nb₃Sn : niobium three tin

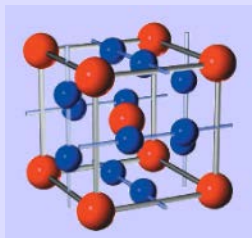
$T_c = 18\text{ K}$

$B_{c2} = 23\text{-}26\text{ T at } 4.2\text{K and } 30\text{T at } 0\text{K}$

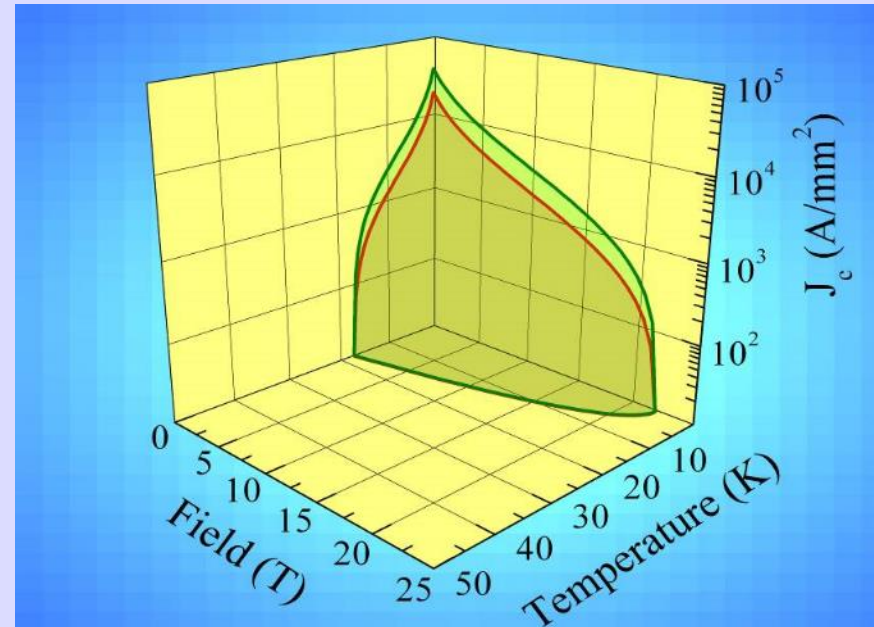
well developed, still in progress

cost ~5-25 €/kA.m at 4.2K, 5T (*)

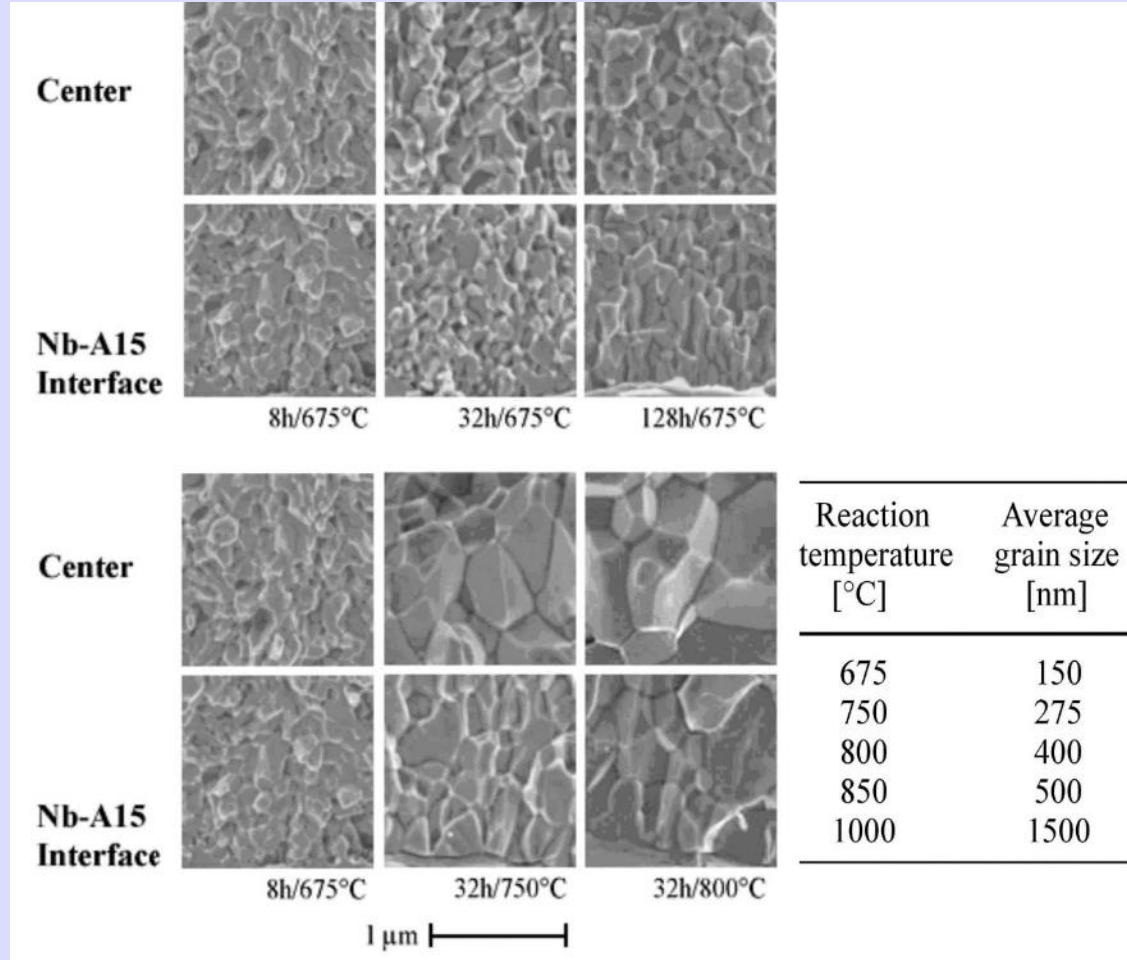
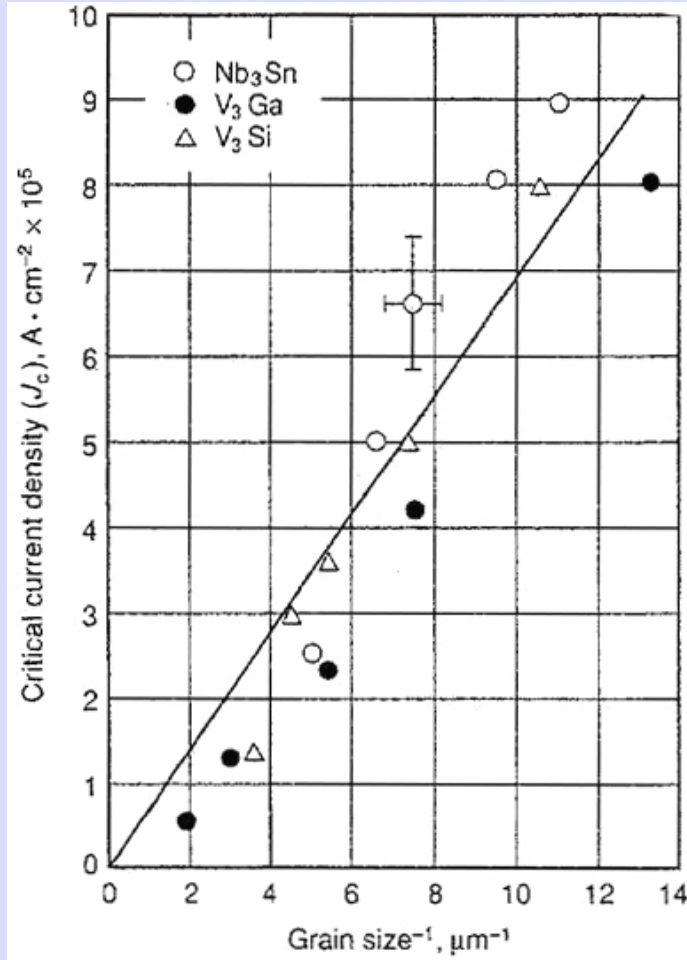
(*) L. Cooley et al., 2005



Cubic inter-metallic, isotropic

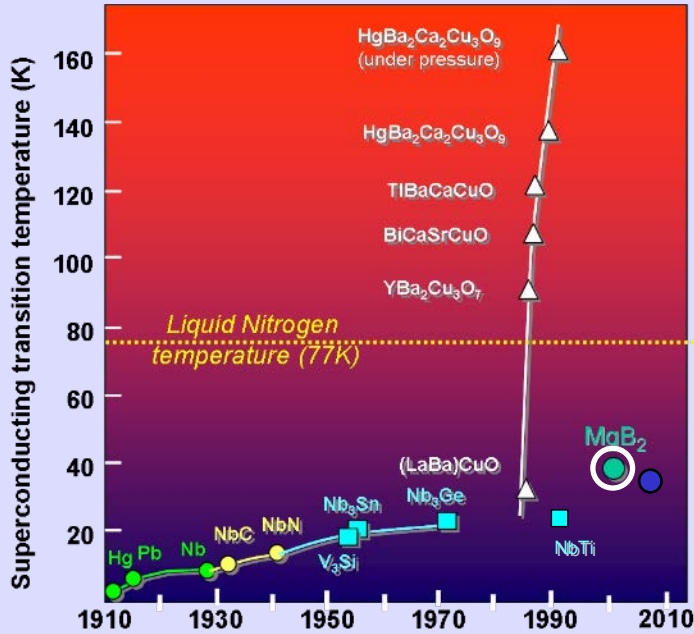


Pinning in Nb₃Sn: grain boundaries



Suenaga et al, 1981

Fischer et al, 2002



MgB₂ : magnesium di-boride

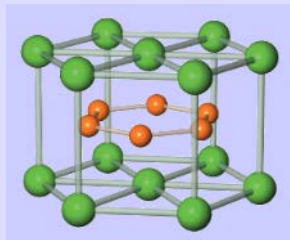
$T_c = 38\text{ K}$

$B_{c2} = 20\text{-}25\text{ T at } 4.2\text{K and } 25\text{-}30\text{ T at } 0\text{K}$

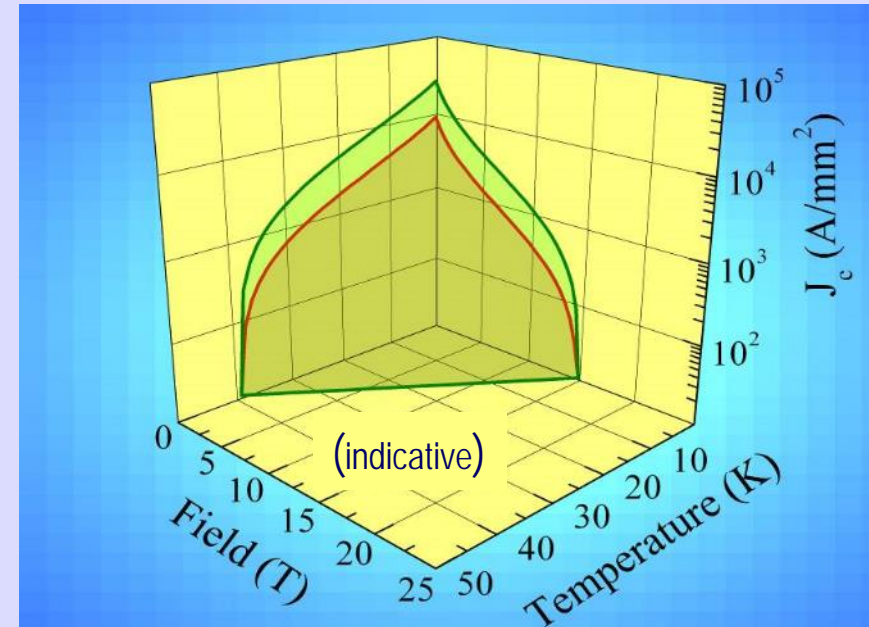
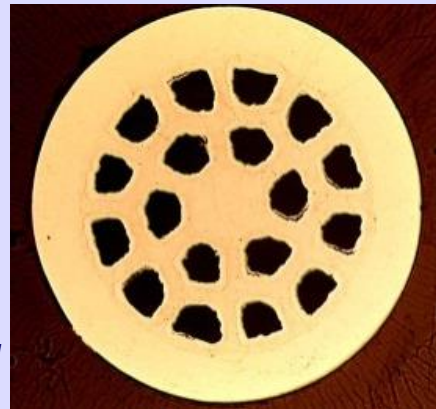
Development in progress

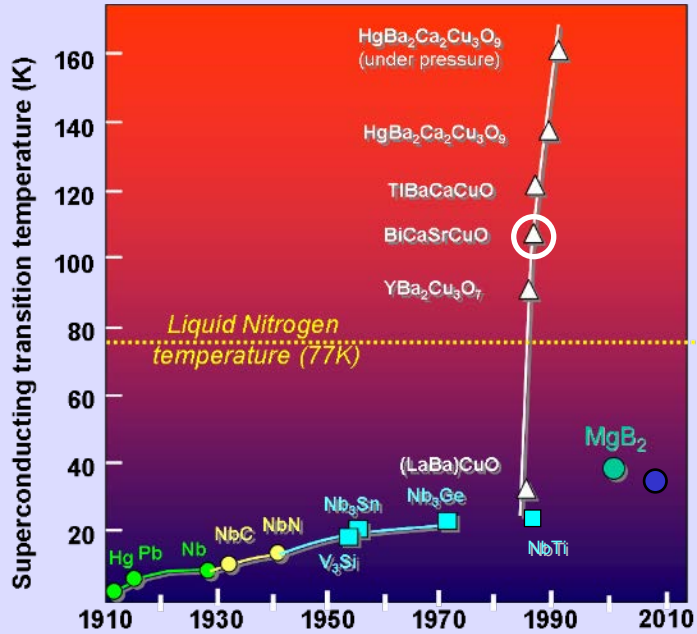
cost ~1 €/kA.m at 4.2K, 5T (*)

(*) L. Cooley et al., 2005



Hexagonal inter-metallic, anisotropic



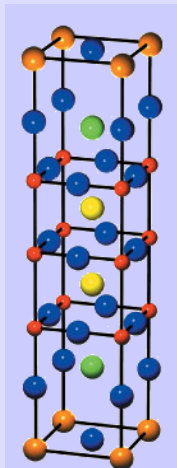


“BSCCO” : bismuth-strontium-calcium-copper-oxide “Bi-2212” and “Bi-2223” (*)

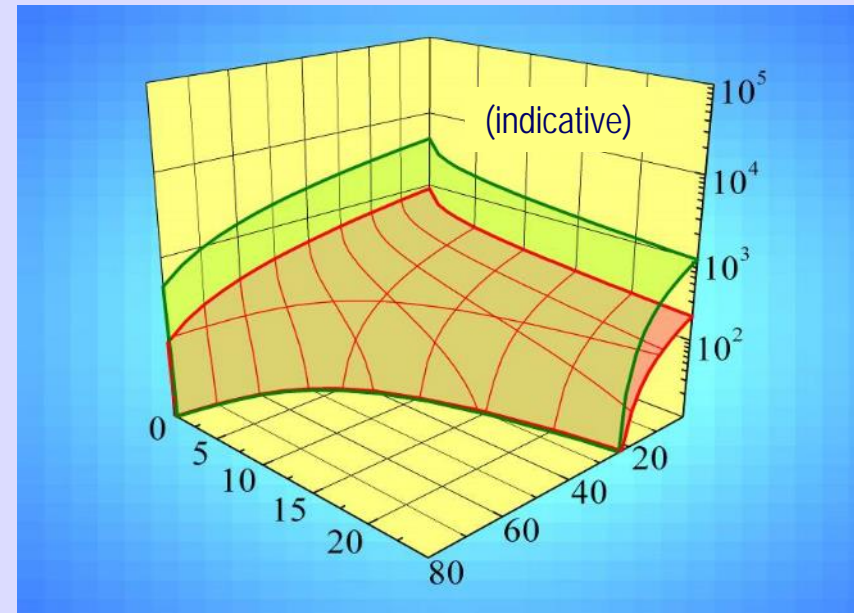
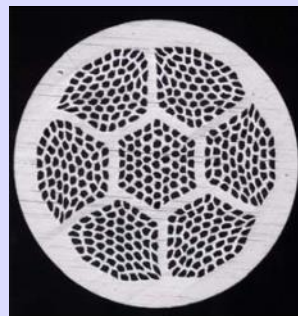
$$T_c = 80 \text{ K (2212) or } 110 \text{ K (2223)}$$

$$B_{c2} = > 100 \text{ T}$$

under development , cost ~ 50 €/kAm (**)

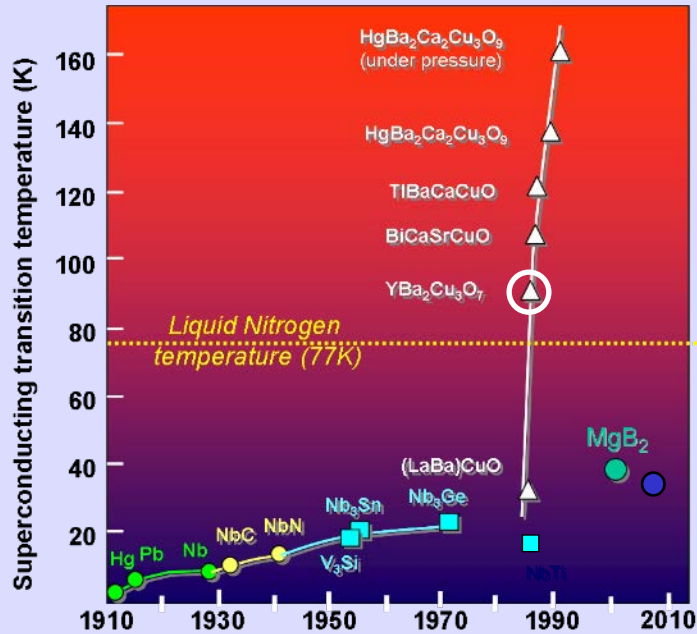


Orthorhombic oxide;
strongly anisotropic



(*) $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_8$
& $\text{Bi}_2\text{Sr}_2\text{Ca}_2\text{Cu}_3\text{O}_{10}$

(**) L. Cooley et al.,
2005



“ReBCO” : yttrium-barium-copper-oxide (*)

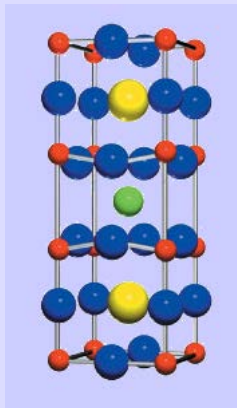
$T_c = 90\text{ K}$

$B_{c2} = > 100\text{ T}$

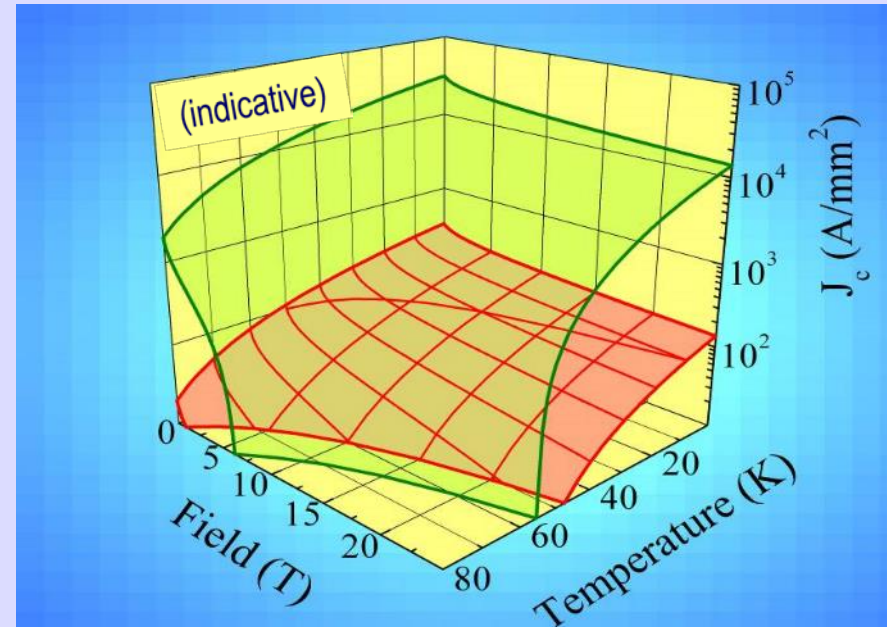
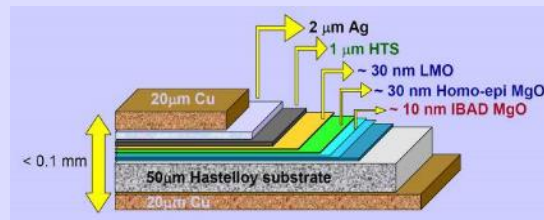
Under development

cost ~100-200 €/kAm (**)

(*) “Re” = rare earth,
used to be $\text{YBa}_2\text{Cu}_3\text{O}_7$,
now often $\text{GdBa}_2\text{Cu}_3\text{O}_7$

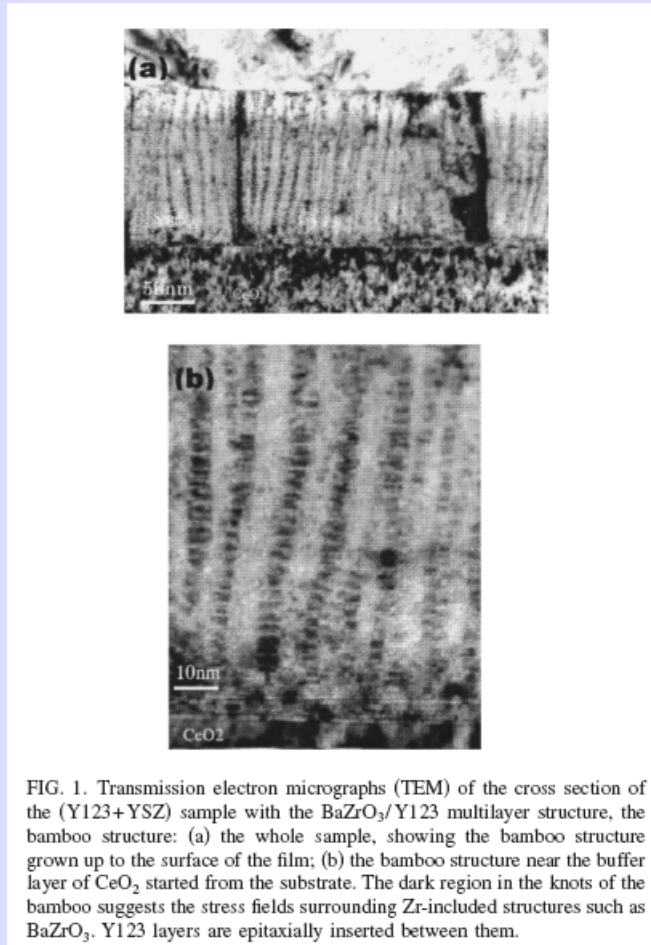


Orthorhombic oxide; anisotropic

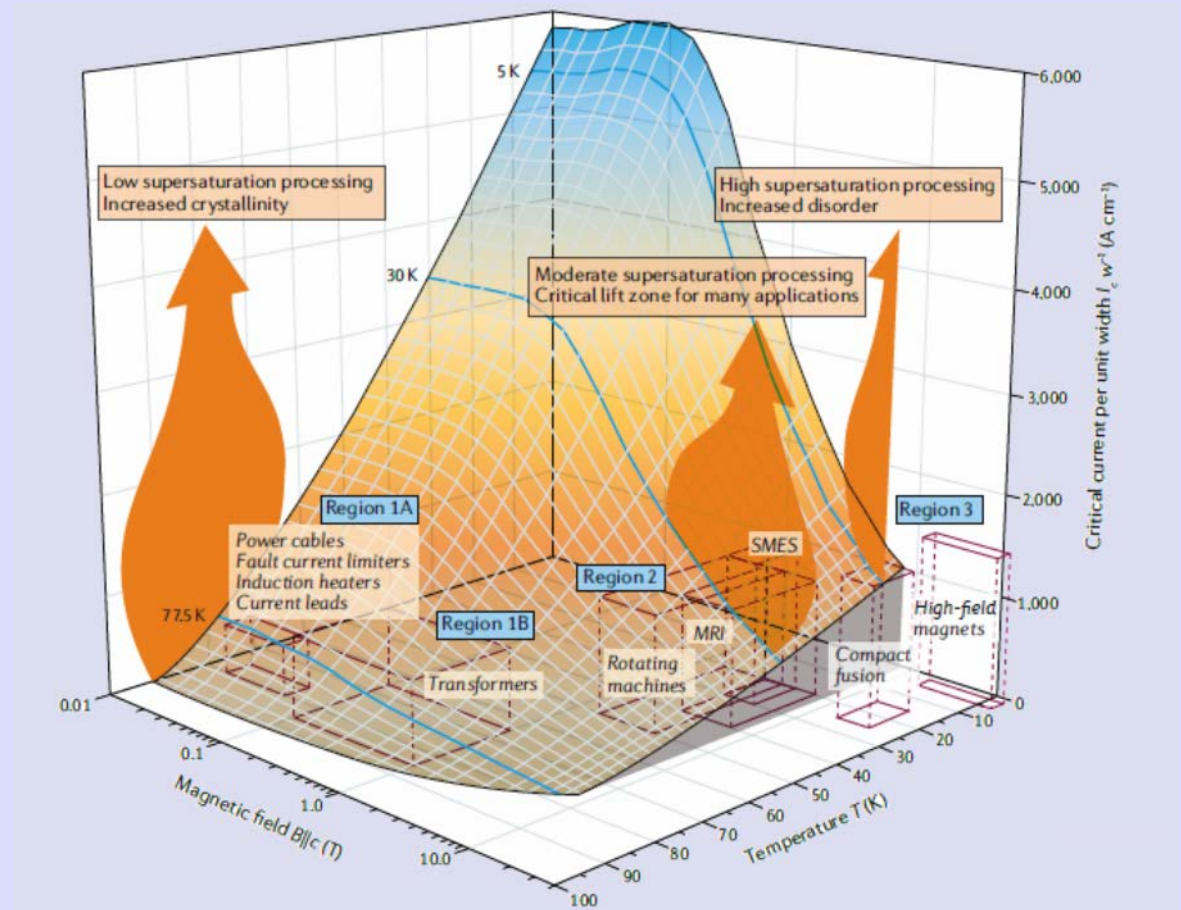


(**) at 30K, 2T;
N. Bykowsky 2016,
J.H. Kim 2016

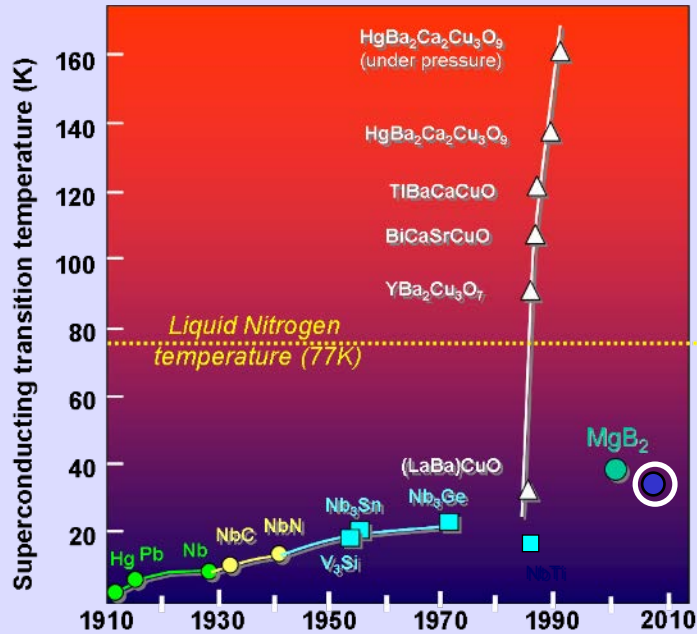
Pinning in $YBa_2Cu_3O_7$: "correlated" disorder (anisotropy!)



Yamada et al *Appl. Phys. Lett.* 2005



MacManus-Driscoll, 2021



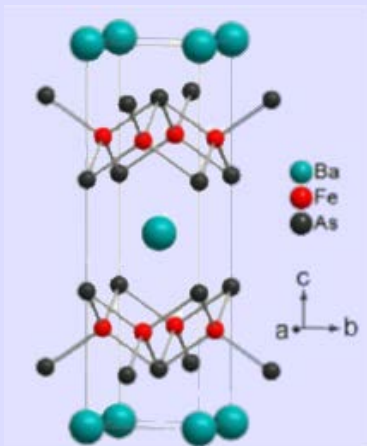
“Ba122” : barium-iron-arsenic (*)

$T_c = 39\text{ K}$

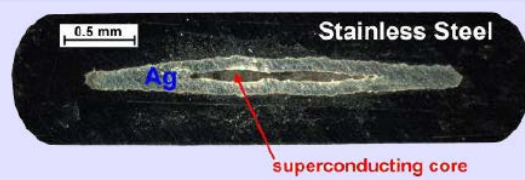
$B_{c2} \approx 100\text{ T}$

(*) $\text{Ba}_{0.55}\text{K}_{0.45}\text{Fe}_2\text{As}_2$

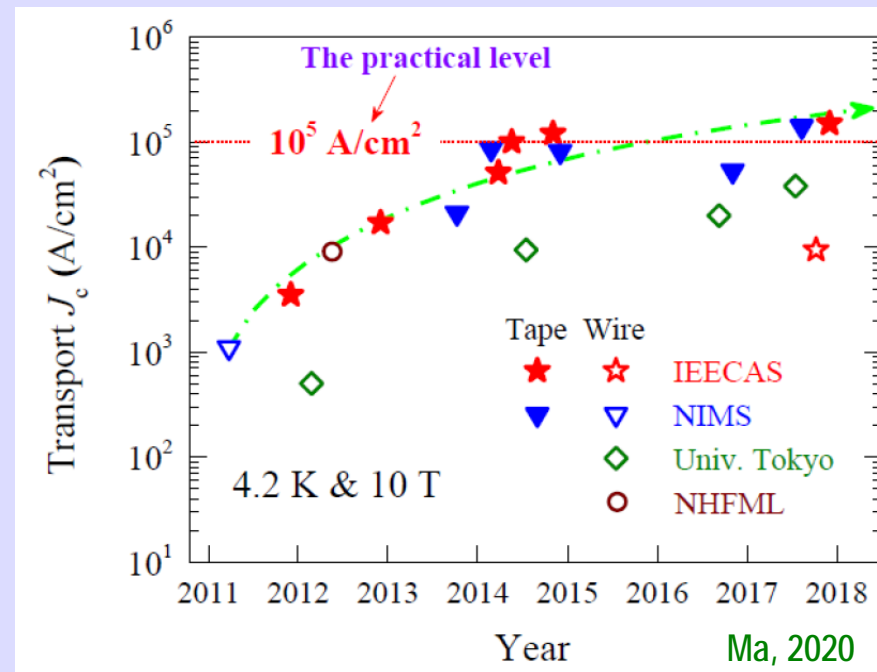
Early development
not yet commercialized



Orthorhombic pnictide;
near-isotropic (!)



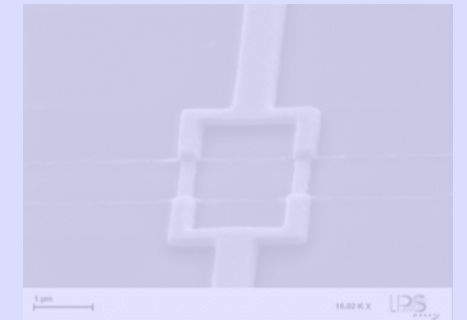
Ma, 2020



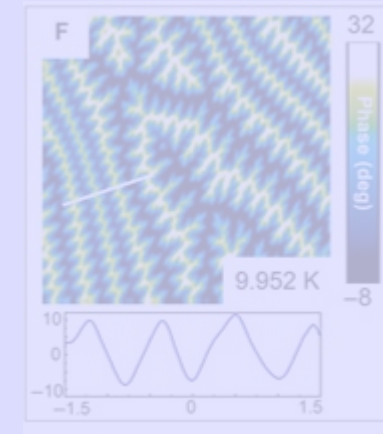
1. Phenomenology & understanding
2. Materials & **applications**
3. Cooling requirements & - strategies



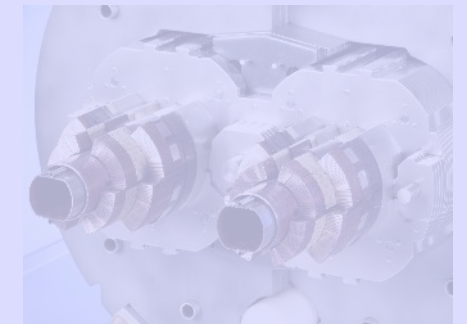
Physics World



Supraconductivité.fr



Science Advances

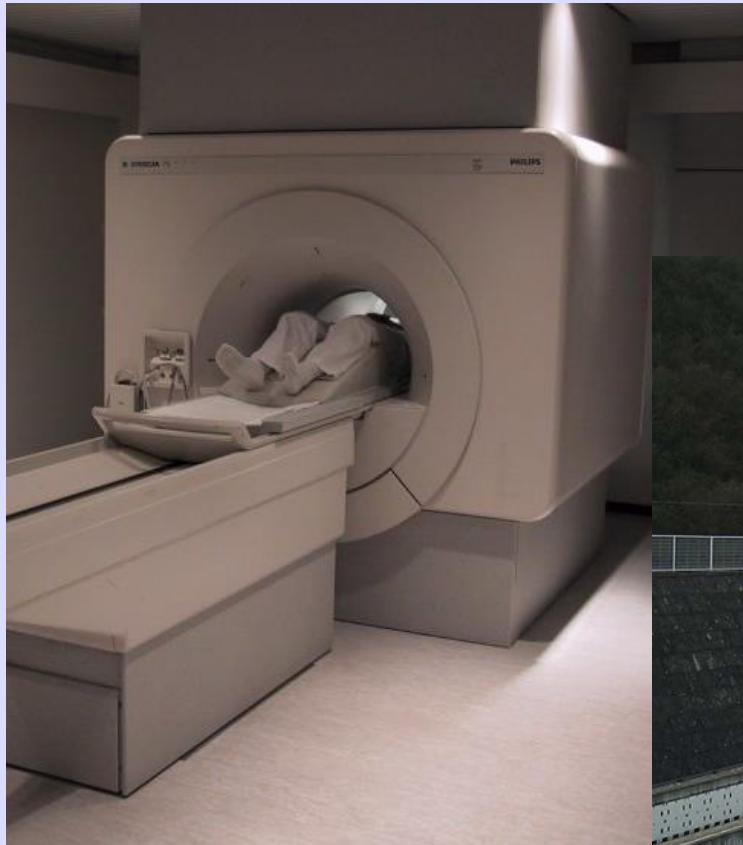


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Large scale (1)

High current density without ohmic heating

($J \sim 100 - 1000 \text{ A/mm}^2$ instead of $\sim 1 \text{ A/mm}^2$)



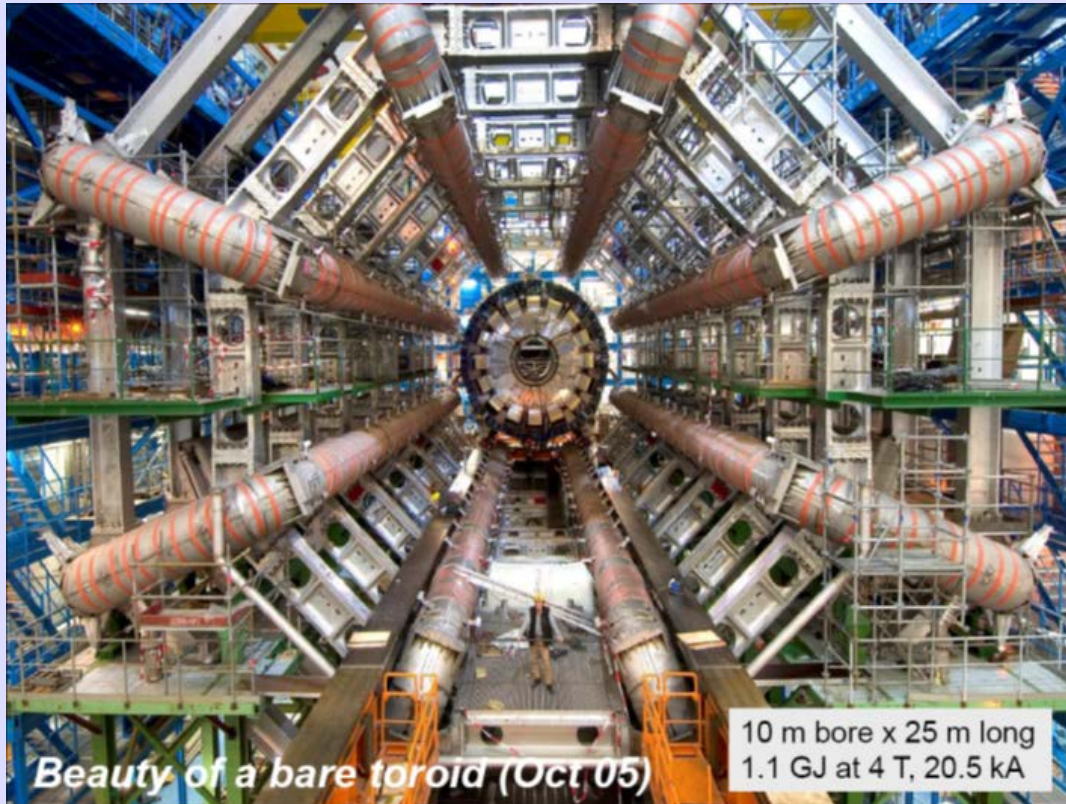
Large scale (2)

High current density without ohmic heating

($J \sim 100 - 1000 \text{ A/mm}^2$ instead of $\sim 1 \text{ A/mm}^2$)

⇒ Powerful **electro-magnets**
(MRI, NMR, HEP, fusion,...)

⇒ Compact electro-technical devices
(cables, FCL, generators, motors...)



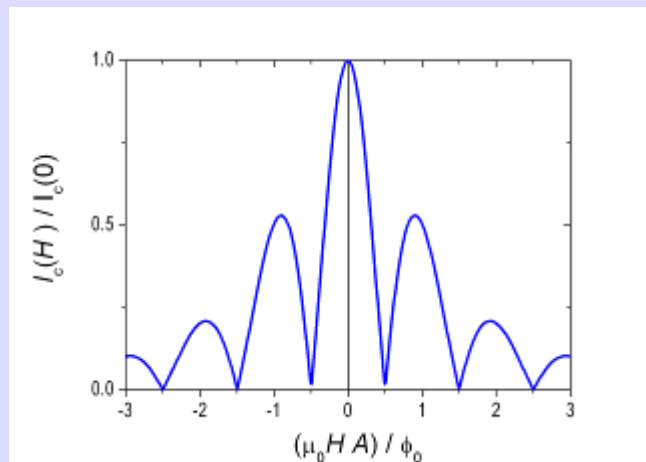
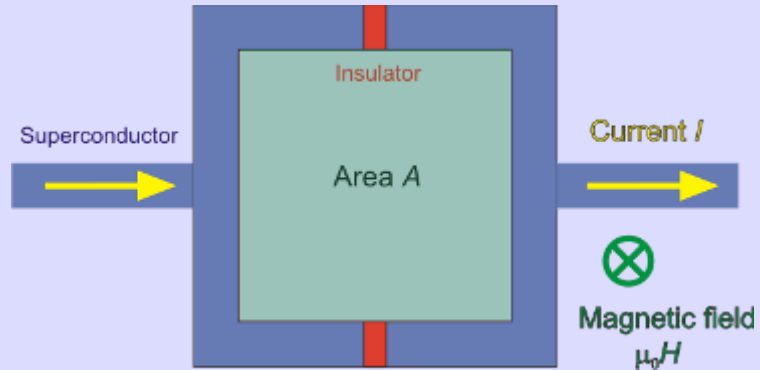
Large scale (3)

High current density without ohmic heating

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- ⇒ Powerful electro-magnets
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- ⇒ Compact **electro-technical devices**
(cables, FCL, generators, motors...)



Small scale (1)Quantum interference with EM fields

⇒ Superconducting Quantum Interference Device
(**SQUID**)

more precise magnetic field measurements

⇒ Measurement standards

(SI 'volt' definition)

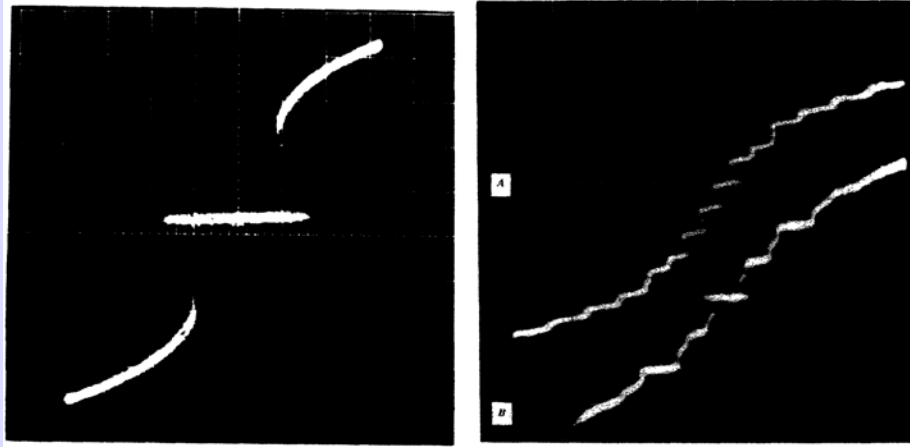
⇒ Faster and compacter electronics

⇒ Qubits, quantum computing

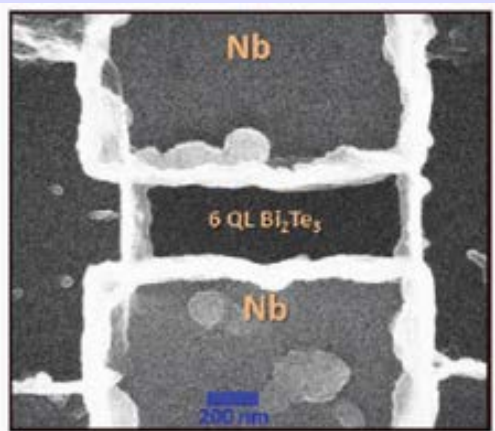


Small scale (2)

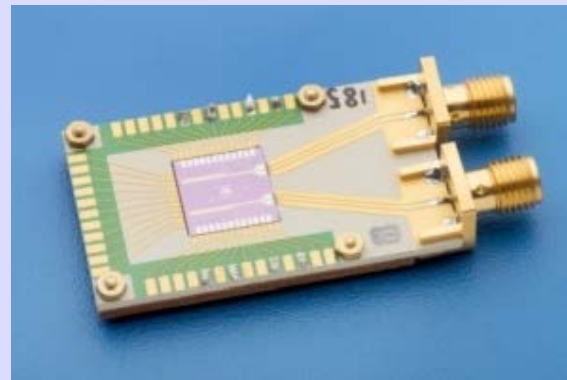
Quantum interference with EM fields



S. Shapiro, 1963



M.P. Stehno, Adv. Mater. 2020



www.PTB.de

⇒ Superconducting Quantum Interference Device
('SQUID')

more precise magnetic field measurements

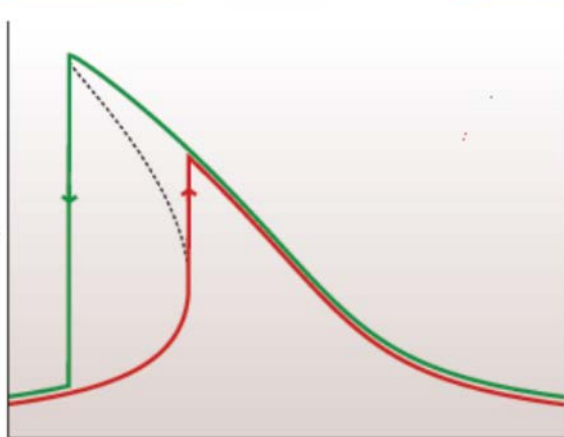
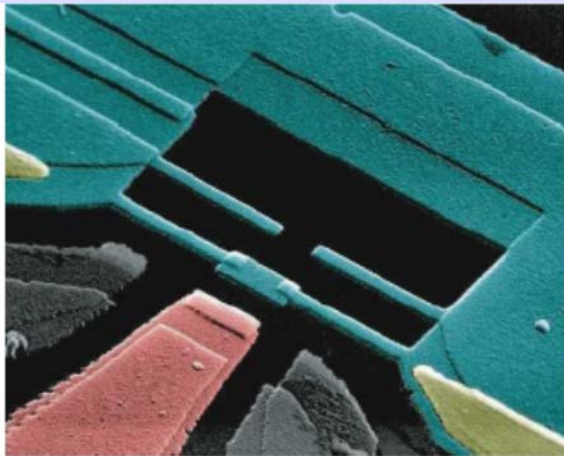
⇒ Measurement **standards**

(SI 'volt' definition)

⇒ Faster and compacter electronics

⇒ Qubits, quantum computing

Small scale (3)



Clarke & Wilhelm, Nature 2008

Quantum interference with EM fields

⇒ Superconducting Quantum Interference Device ('SQUID')

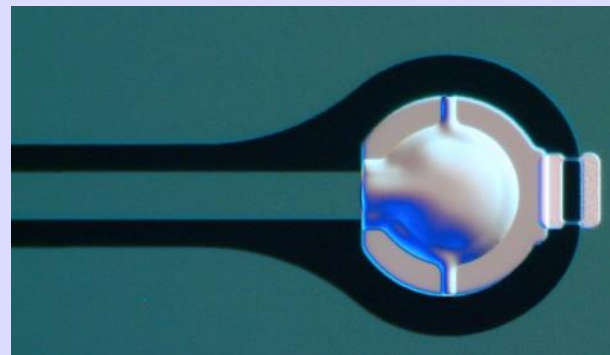
more precise magnetic field measurements

⇒ Measurement standards

(SI 'volt' definition)

⇒ Faster and compacter **electronics**

⇒ Qubits, quantum computing

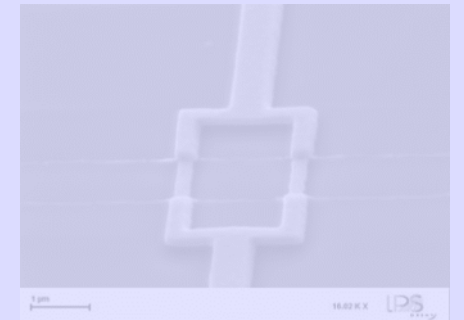


Bosman et al, Nature 2017

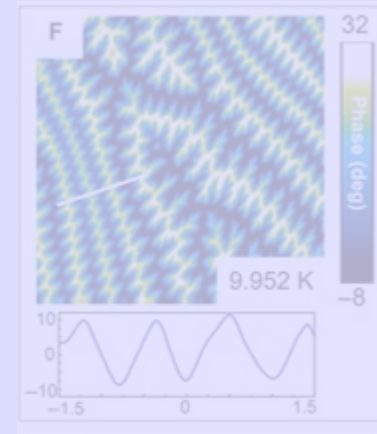
1. Phenomenology & understanding
2. Materials & applications
3. Cooling requirements & - strategies



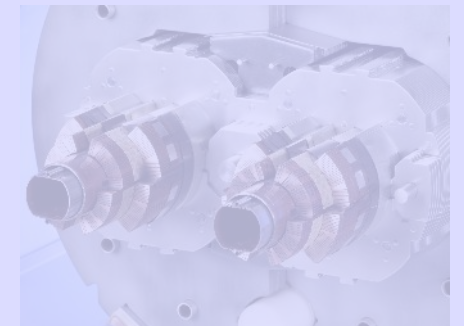
Physics World



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Science Advances

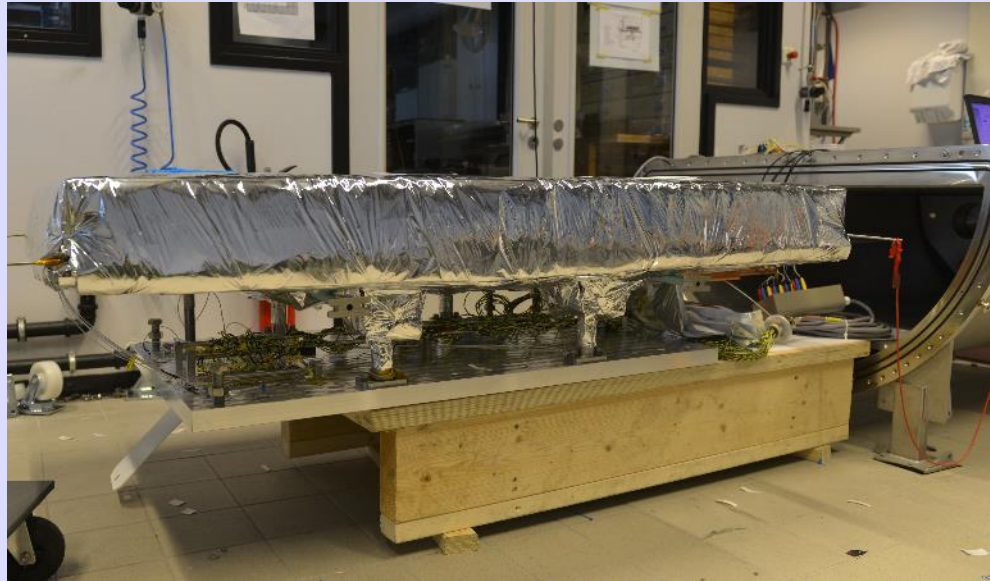


CERN.ch

Heat loads:

- in-leak through cryostat enclosure ➤ in common with other cryogenic systems
- ramp- or AC losses ➤ unavoidable, but strongly application-dependent
- external current leads ➤ perennial worry
- 'warm-cold' structural elements ➤ application-dependent
- 'accidental' ➤ "quench"-detection and - protection

In-leak through cryostat enclosure

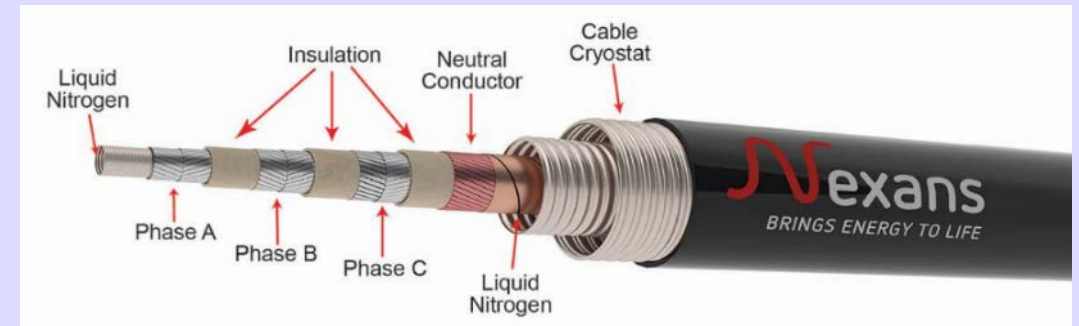


MDS separator



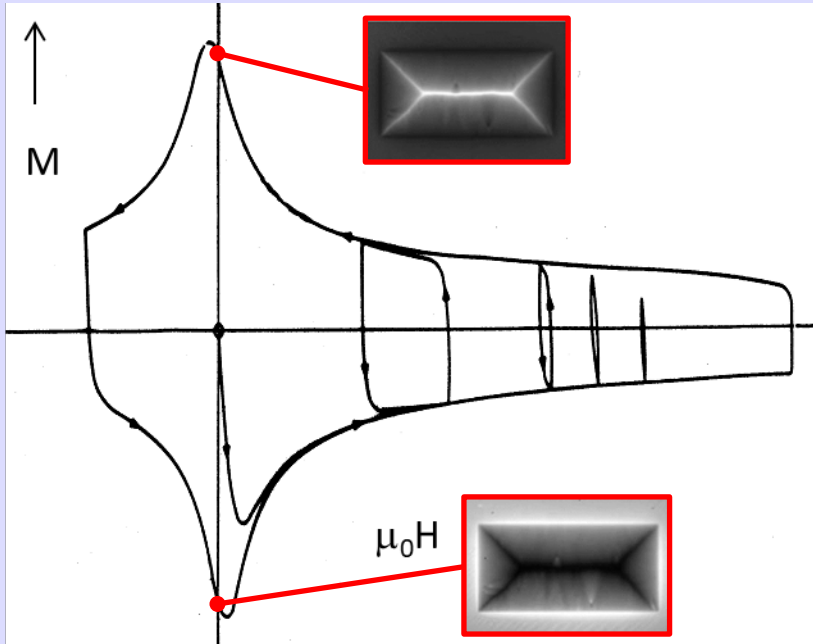
Ecoswing rotor

- Typically vacuum + MLI;
- Challenge: minimize the 'air' gap
- ~ 0.1 - 1 W/m² between 293 & 4.2 K (chambers);
- ~ 0.1 - 1 W/m between 293 & 77 K (flex tubes)



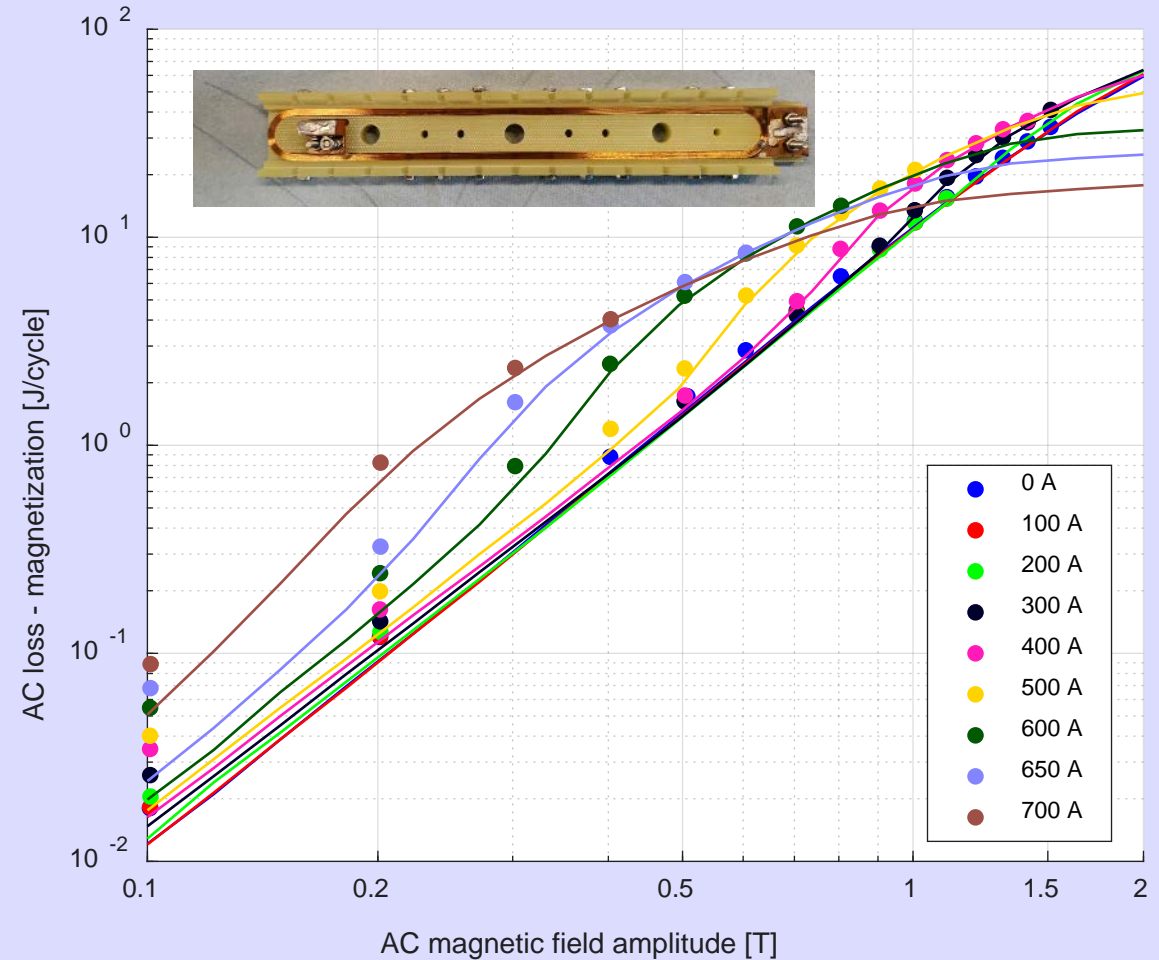
www.nexans.com

Ramp- or AC losses



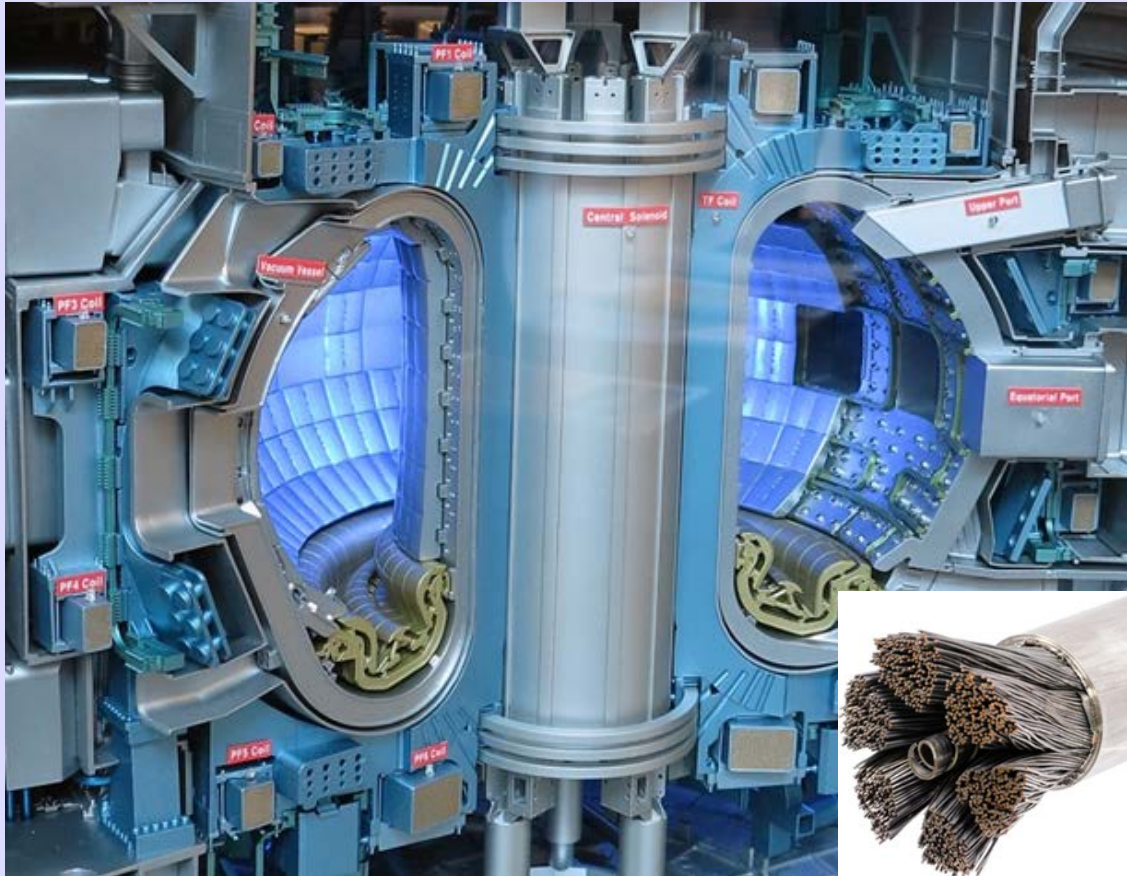
after Jooss et al, 2001

- Time-varying magnetic fields or currents lead to hysteretic internal flux profiles
- Motion of vortices \rightarrow dissipation ...
- **Hysteresis loss**



Hyseresis losses in linear actuator coil

Ramp- or AC losses



www.anl.gov

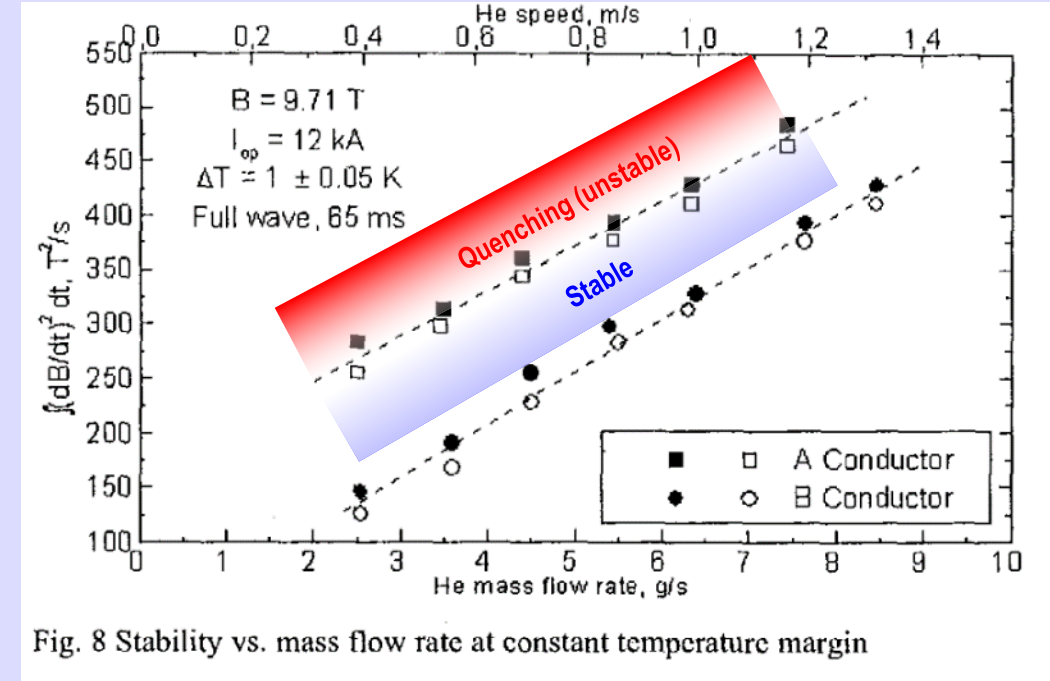
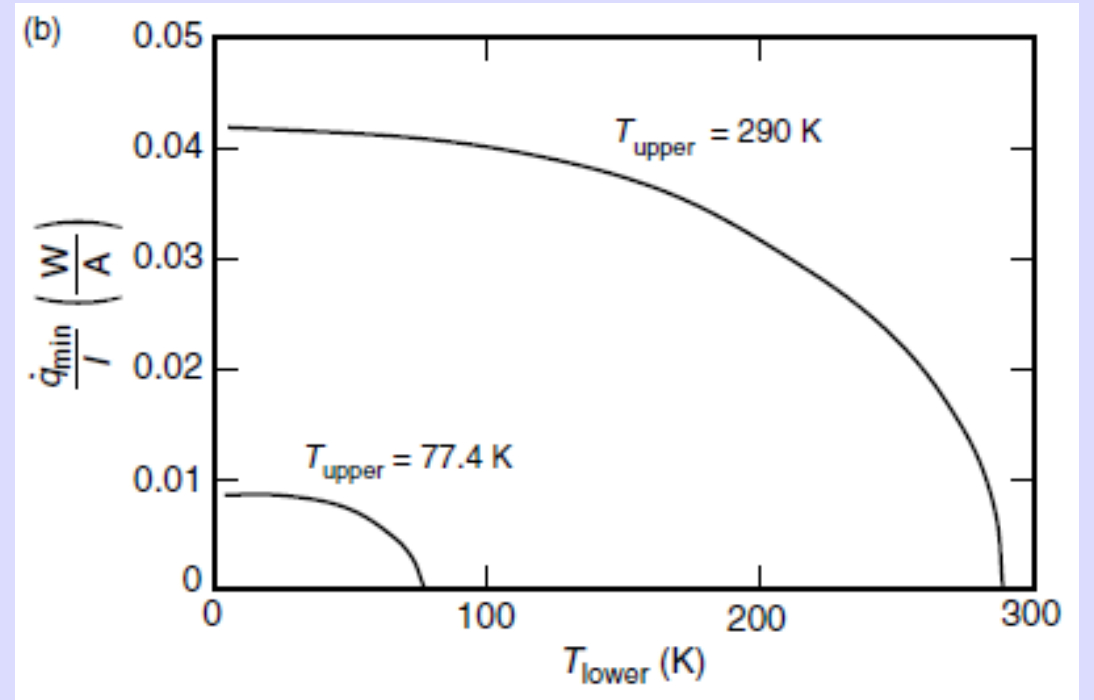
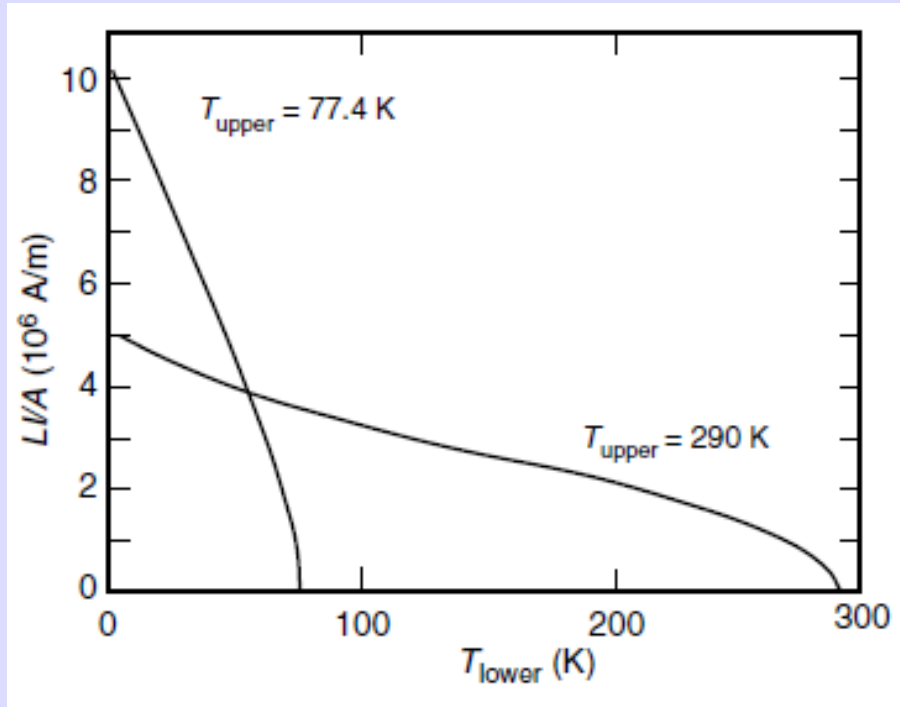


Fig. 8 Stability vs. mass flow rate at constant temperature margin

after P. Bruzzone et al, 2001

- ... in addition, induced currents between strands in cabled structures lead to ohmic loss
- **'Coupling' loss**

External current leads between RT power supply and cryogenic SC device



J. Ekin, "Experimental techniques for low-temperature measurements" 2006

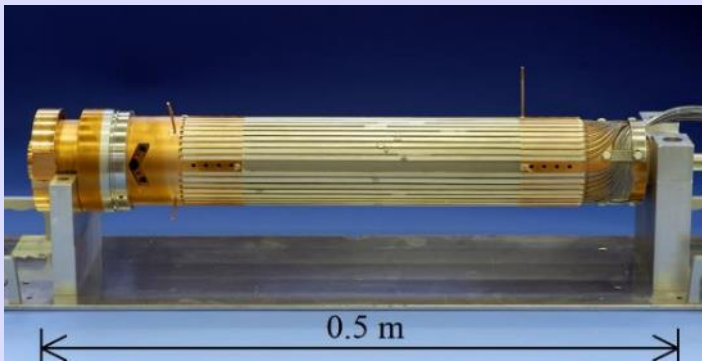
- Minimizing heat conduction calls for long length L and small cross-section A;
 - Minimizing ohmic heating calls for short length L and large cross-section A
- } {
- Optimal geometry L/A
 - Unavoidable loss level

External current leads

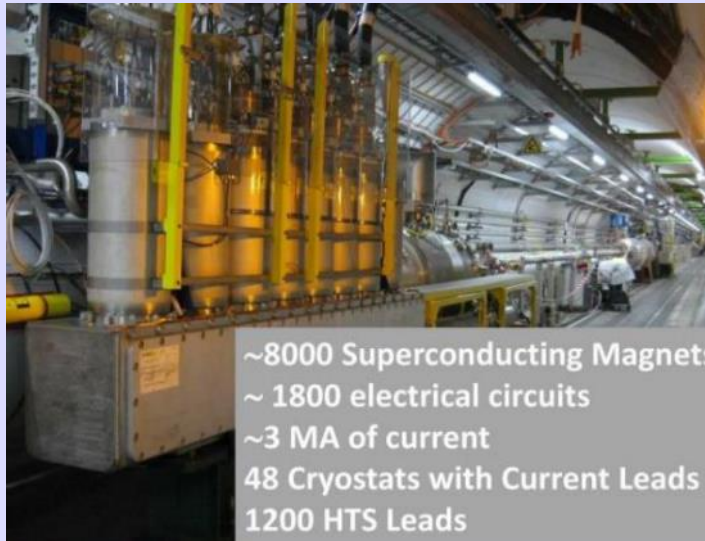
Possible escape (1) :

intercept in-leak a.s.a.p. &
continue with HTS SC leads to actual device

CERN LHC current feedthroughs



A. Ballarino, 2004



~8000 Superconducting Magnets
~ 1800 electrical circuits
~3 MA of current
48 Cryostats with Current Leads
1200 HTS Leads

A. Ballarino, 2013

GM cooler

290 K
terminal

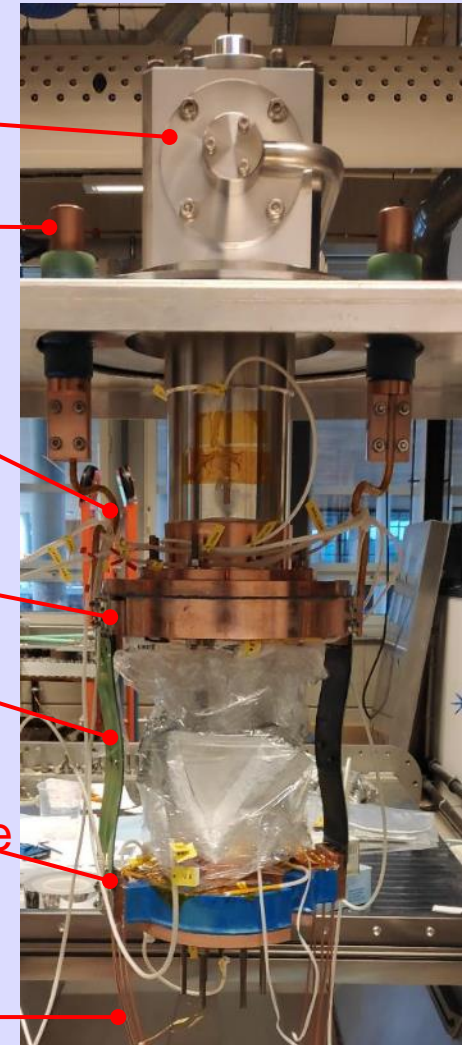
Optimal L/A
Cu lead

60 K 1st stage

BSCCO
SC lead

4.5 K 2nd stage

NbTi
SC lead



MDS current feedthroughs

External current leads

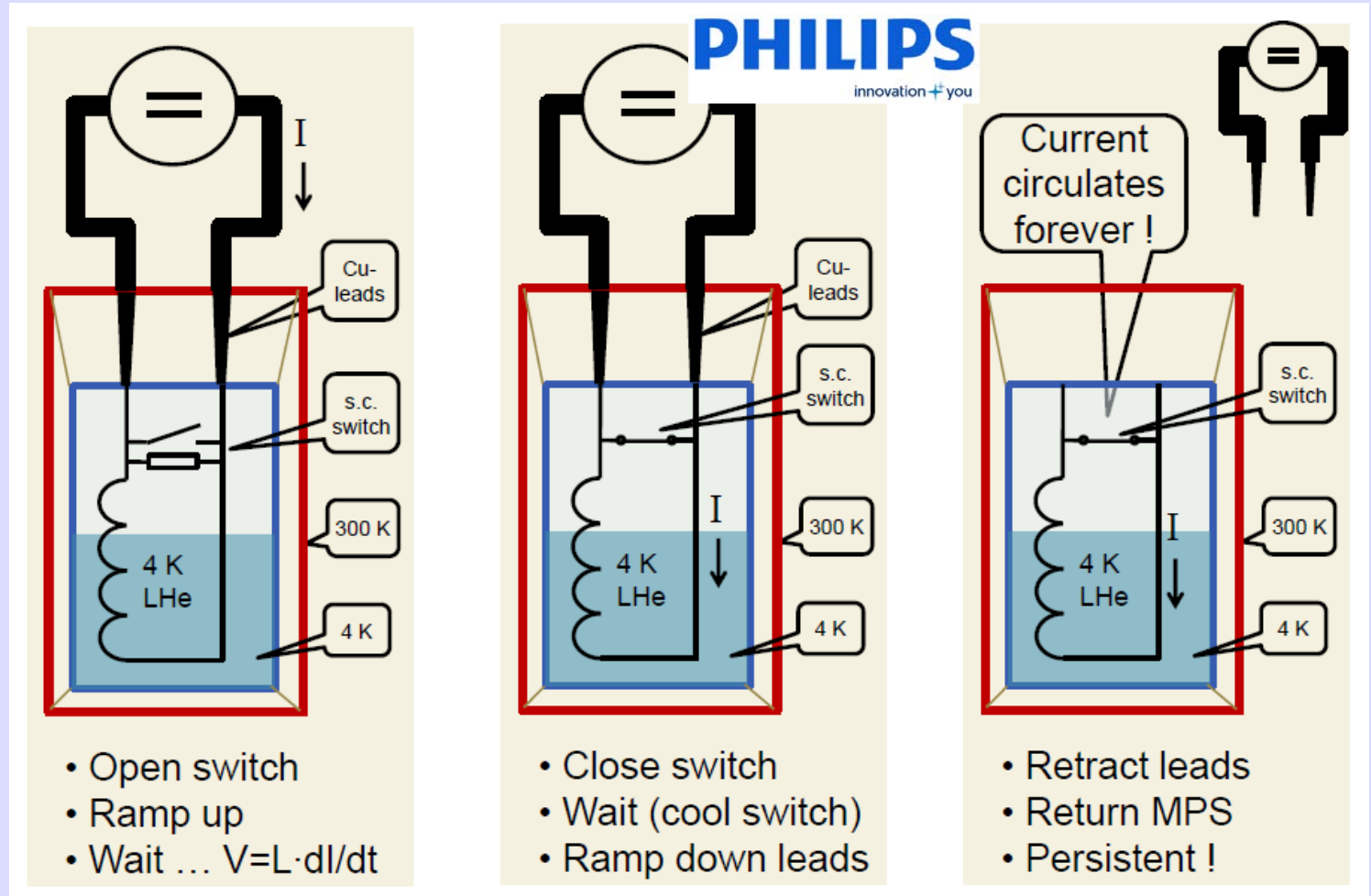
Possible escape (2) :

retract RT connection after charging of the device

persistent mode



G. Mulder, 2018

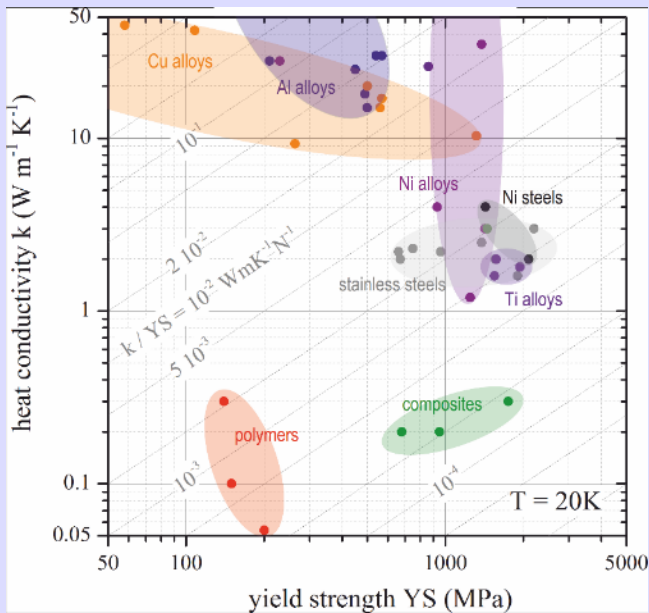


G. Mulder, 2018

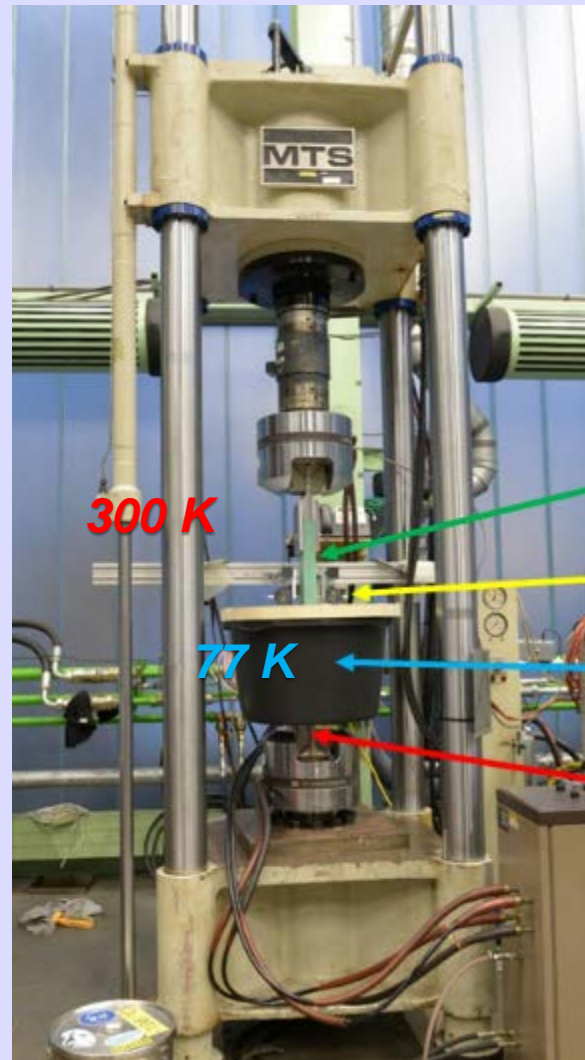
3. Cooling requirements & strategies

Warm-cold supports

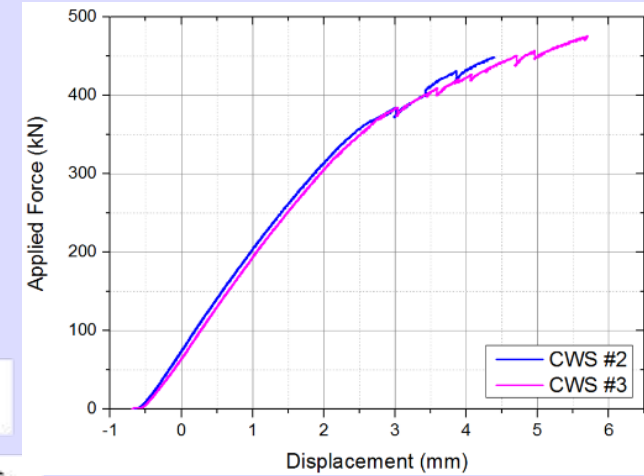
- Especially in electrical machinery (high torque)
- high E_Y – low κ material
- **reinforced polymers**



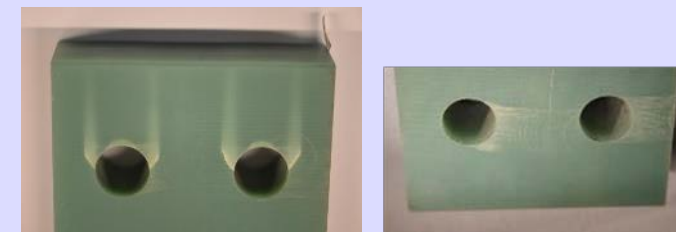
Cryogenic strength vs heat conductivity



Tensile / compressive loading test coupons @ TNO Delft



Tensile yield test



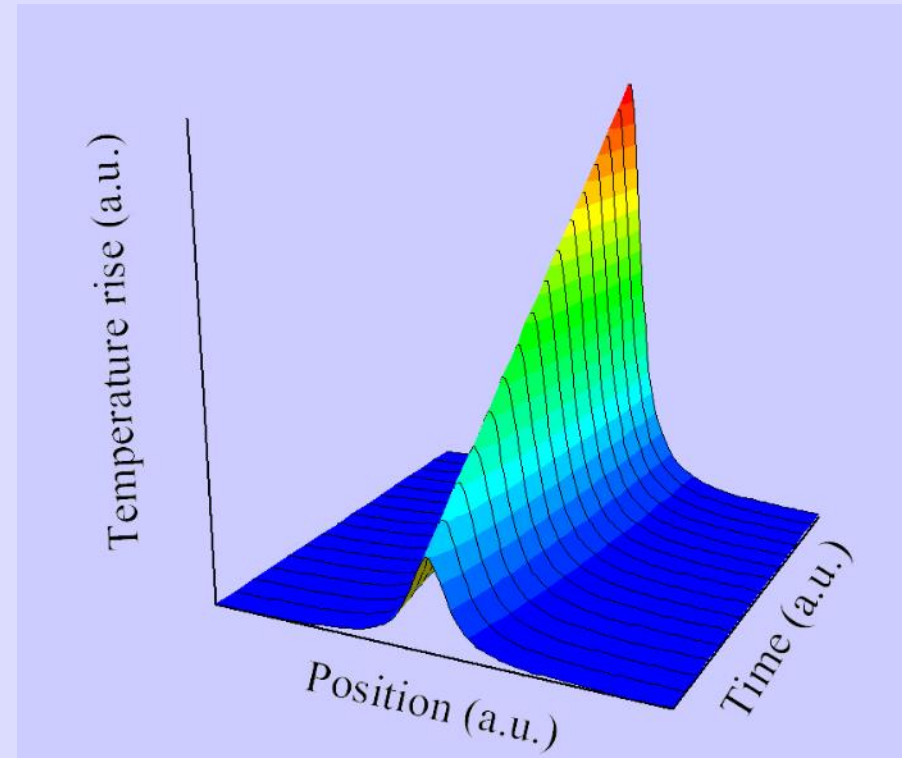
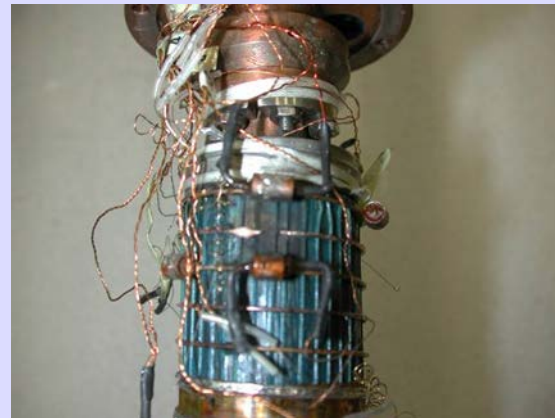
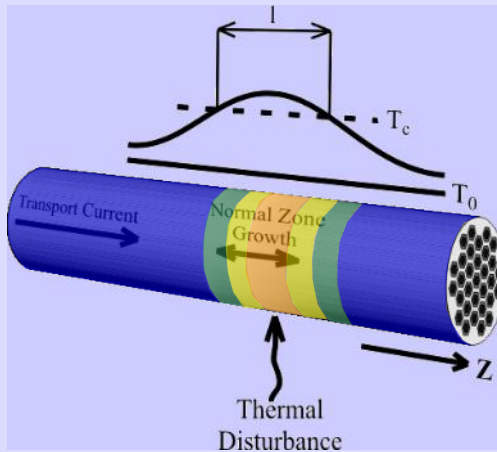
Yielding on warm side, not cold side

- Cold-warm-support test coupon
- Lateral Extensometers
- Thermal insulation
- Heater

Quench-detection and protection

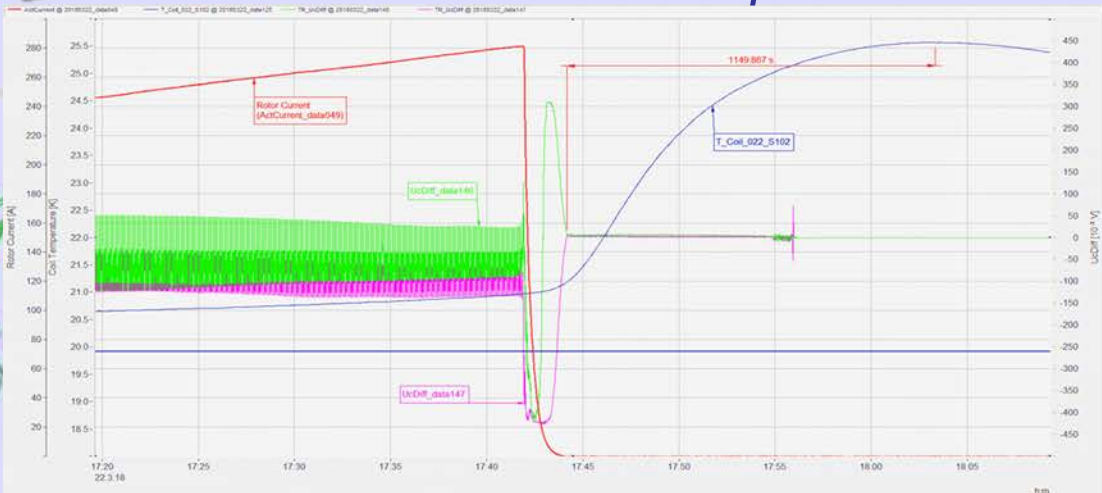
- An unexpected local disturbance can initiate a 'thermal avalanche', a **quench**
- Such a disturbance **must** be detected a.s.a.p. and the stored energy safely extracted

$$C \frac{\partial T}{\partial t} = \frac{\partial}{\partial t} \left(\kappa \frac{\partial T}{\partial x} \right) + \rho J^2 + p_{initial}$$



3. Cooling requirements & strategies

Quench-detection and protection



Temperature rise sub-standard coil after **quench**

- Sub-standard coil had passed (accelerated) acceptance test ...
- ... and failed during power-up ramp;
- Inadequate 'quench-detection' (EM interference)
 - **Required coil replacement & protection upgrade**

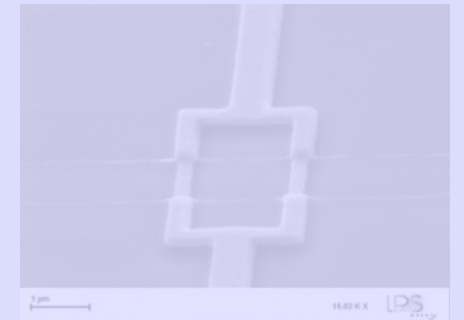


Repair action at Boessenkool Almelo

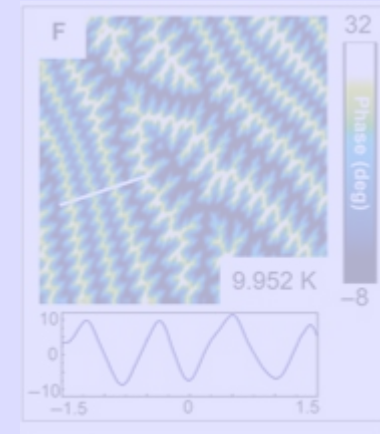
1. Phenomenology & understanding
2. Materials & applications
3. Cooling requirements & - *strategies*



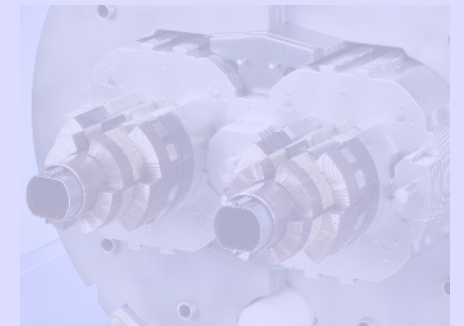
Physics World



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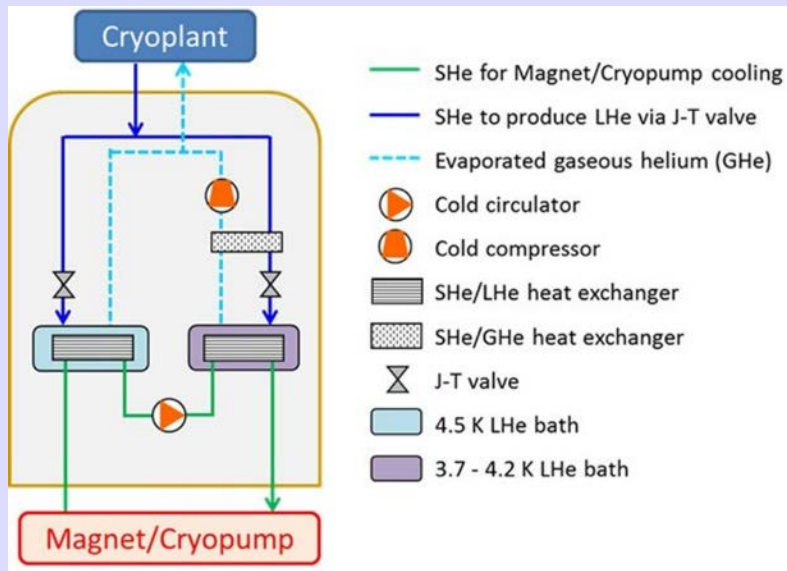
Science Advances



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Large systems ('Big Science')

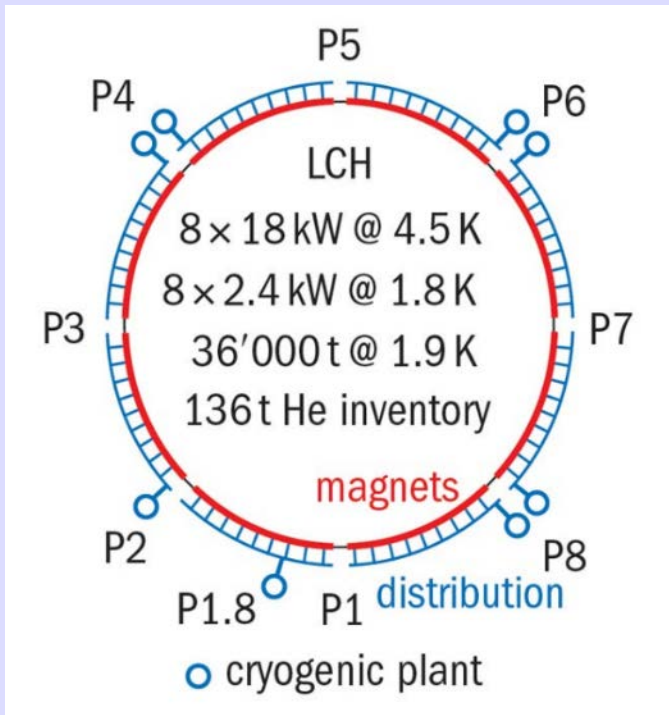
- Forced-flow of (possibly superfluid) He
- Extended cryo-plants for re-liquification



www.iter.org

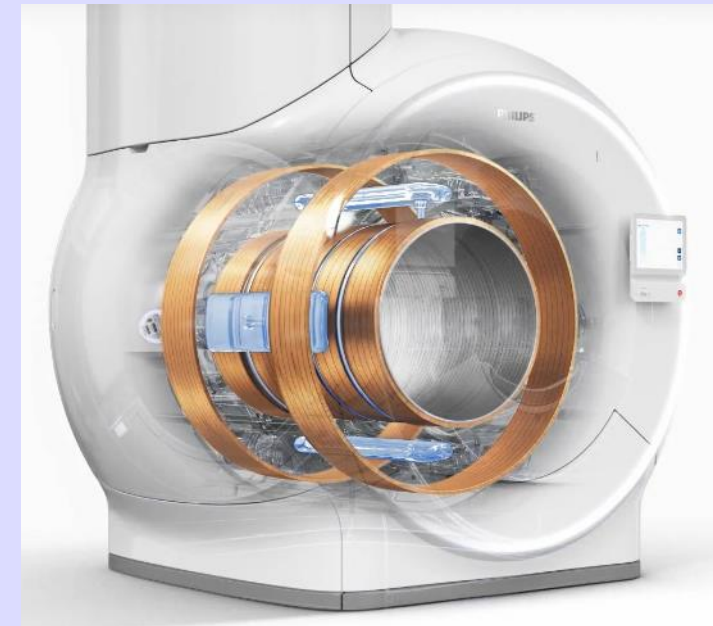
Large systems ('Big Science')

- Forced-flow of (possibly superfluid) He
- Extended cryo-plants for re-liquification



CERN COURIER, 2013

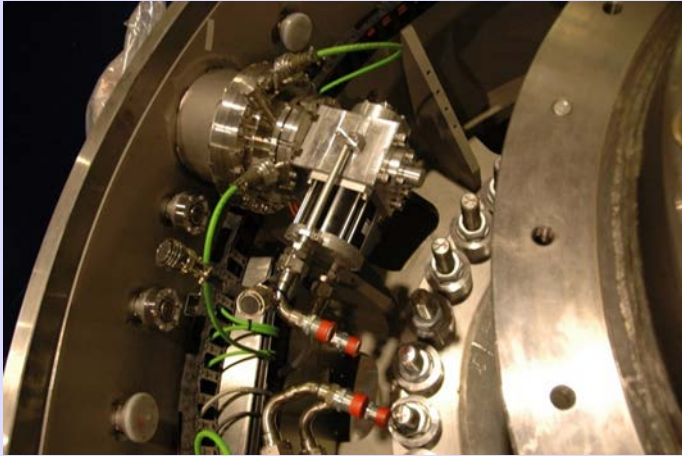
Medical systems (stand-alone)



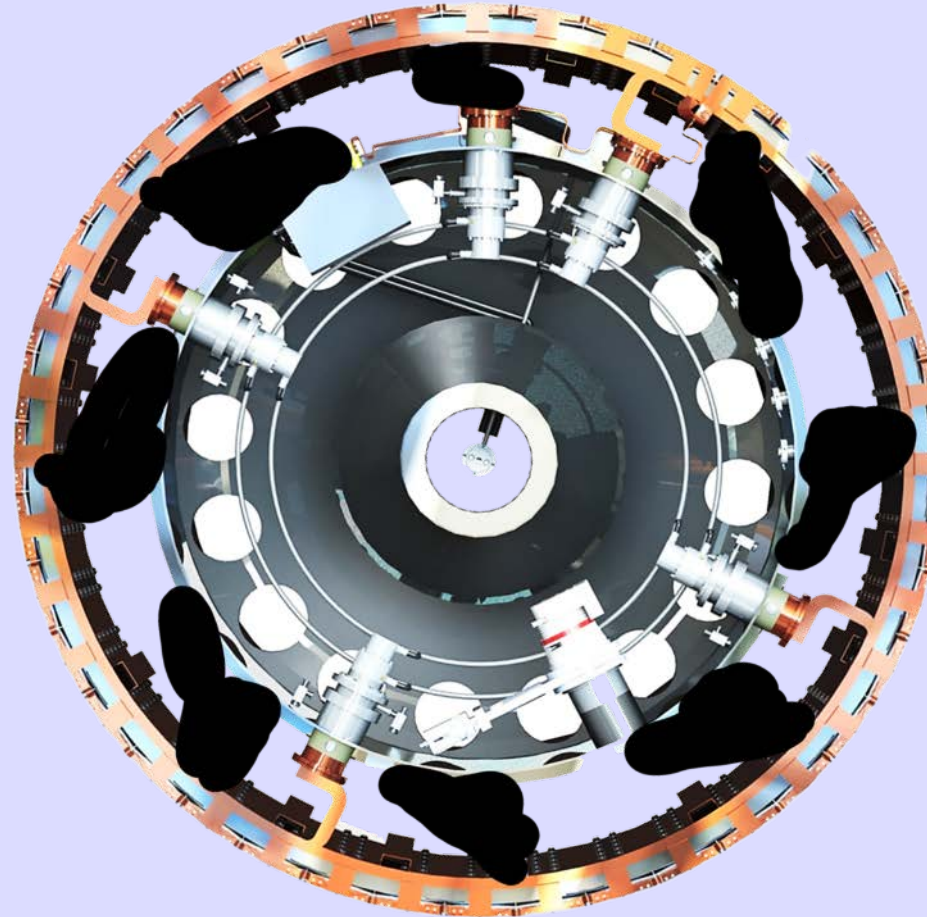
www.philips.com

- Liquid He bath cooling
- “Zero boil-off” technology (cryocooler-based re-condenser)
- Trend towards less & less He on-board (e.g. Philips ‘BlueSeal’ MRI system)

'Industrial' machines (stand-alone) : cryocoolers + conduction cooling



GM cold-heads & compressors



Cryogenic rotor design
(Cu cold-bus, distributed coolers)



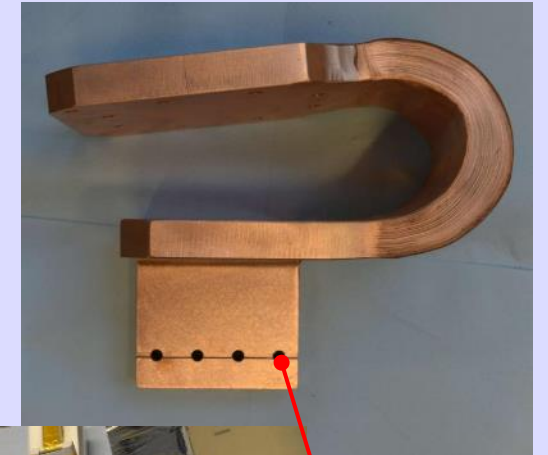
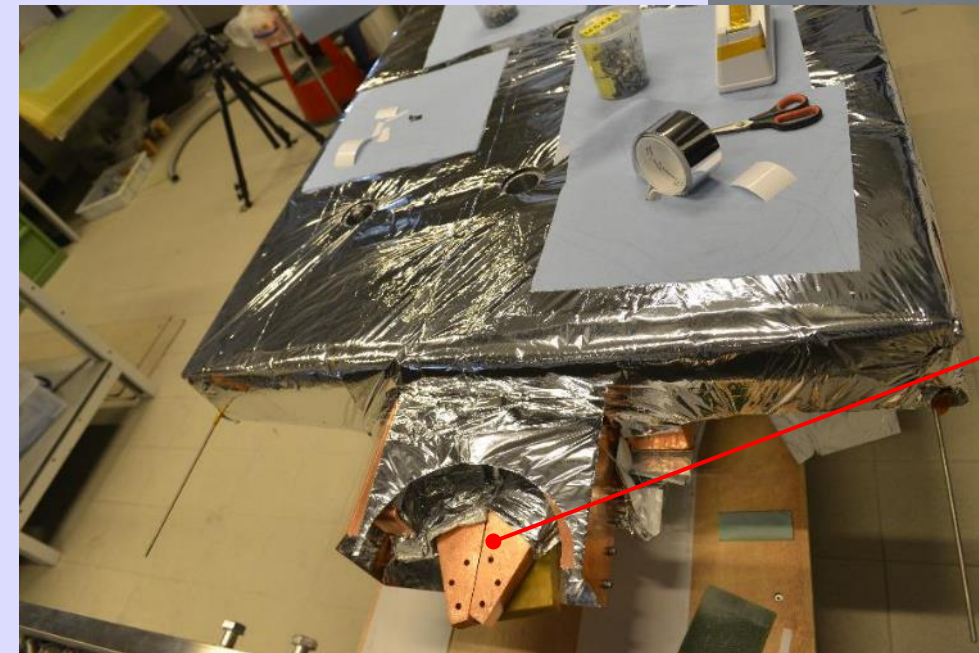
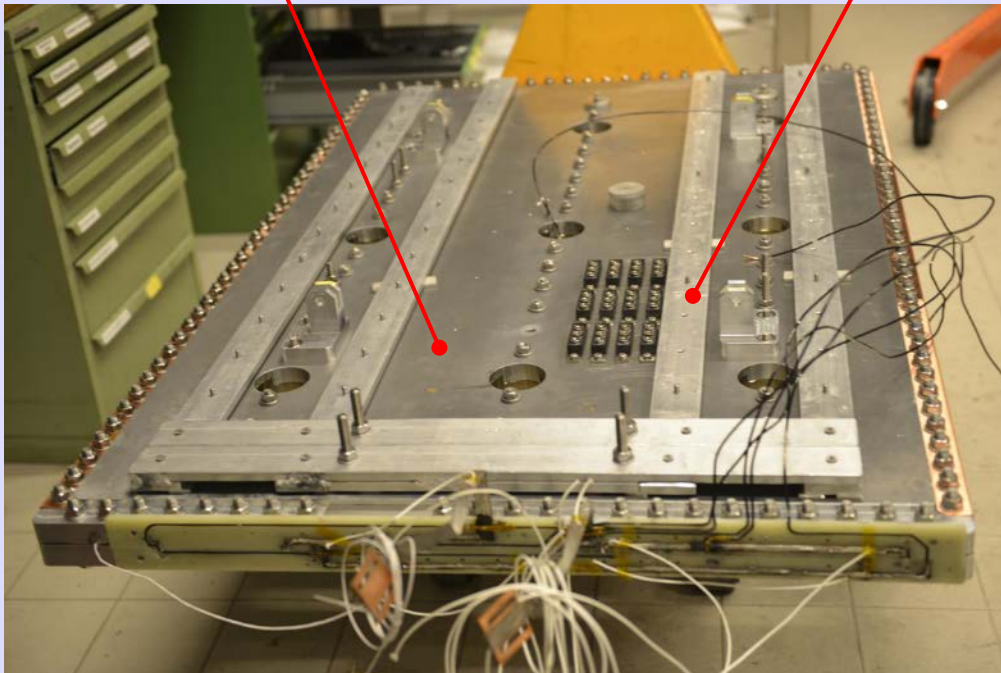
- Cu 'cold-bus' monted on cold back-iron
- 'On-board' rotating GM coolers
- Static compressors
- **Rotating He gas coupling!** (Sumitomo)

'Industrial' machines (stand-alone) : cryocoolers + conduction cooling

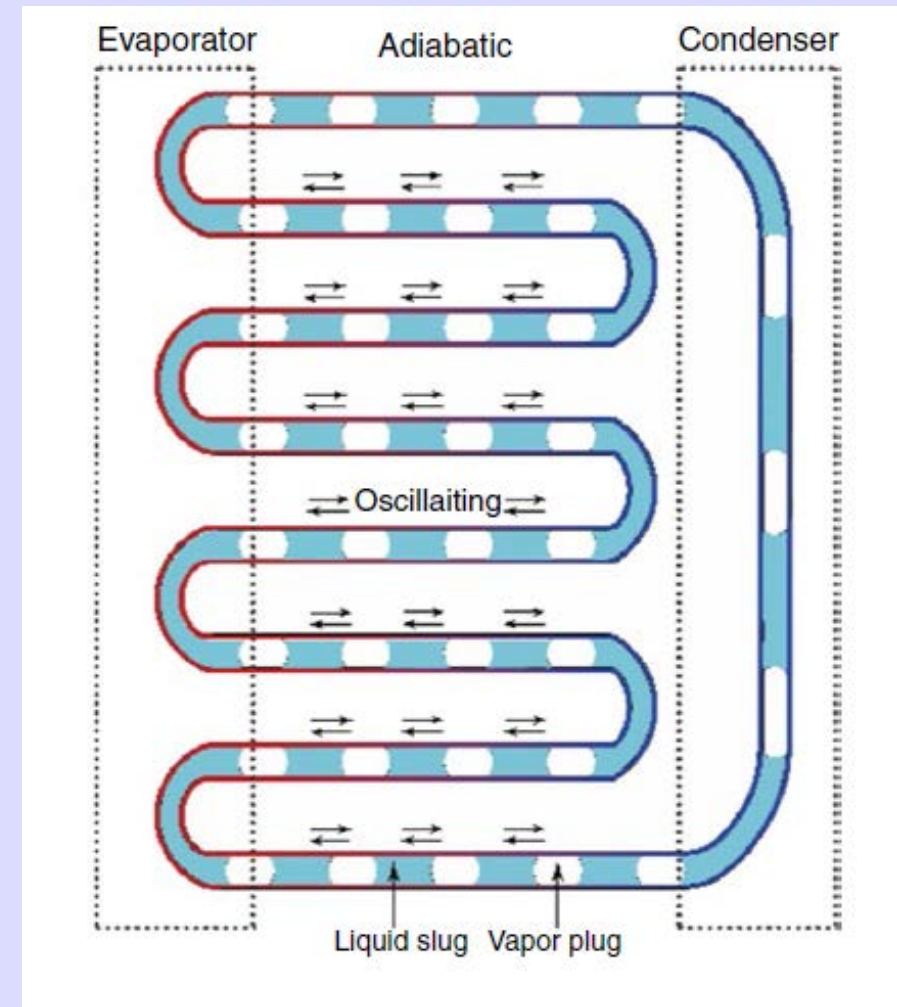
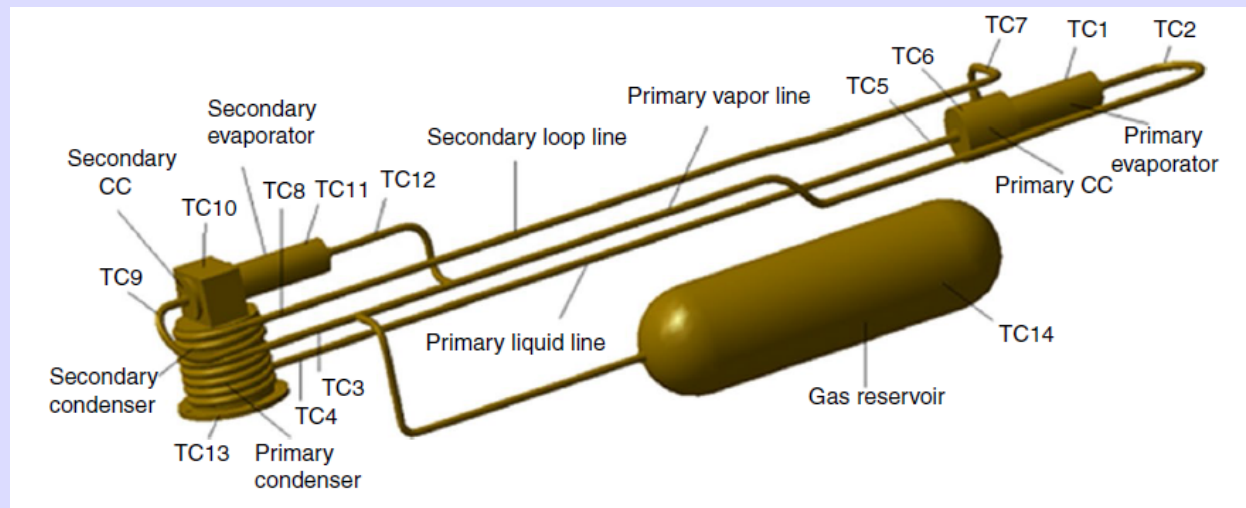
- Thermal gradients need to be kept as low as possible
- High-purity metal flex-links & thermal shunts ...

Al alloy cassette (pre-compression)

high-purity Al shunts



Cryogenic heat-pipes as thermal links for the most demanding SC applications?



A. Haghighi et al, 2022