# **High-precision measurement of the W boson mass with the CDF II detector**



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**LHCb student seminar 23 June, 2022**





### Boson masses

### Higgs field potential



Potential with a minimum at a non-zero value of the scalar<br>Potential potential with a minimum at a minim **Gauge field potential** 

$$
V = -\frac{g^2 v^2}{8} [(W_{\mu}^+)^2 + (W_{\mu}^-)^2]
$$
  
The **n**  
the **va**  

$$
-\frac{v^2 (g^2 + g^2)}{8} Z^{\mu} Z_{\mu}
$$

 $\mathcal{L}$   $\overline{\phantom{125}Z}$   $\overline{Z}$   $\overline{Z}$   $\overline{Z}$   $\overline{Z}$   $\overline{Z}$   $\overline{Z}$   $\overline{Z}$  $m_Z = \frac{1}{2} \sqrt{g^2 + g^2}$  $\sqrt{2\pi}$  $m_W^{} =$ *v* 2 *g*  $m_Z =$ *v*  $\frac{y}{2}\sqrt{g^2+g^2}$ 

> *e*<sup>2</sup>*µ* 2 Ge T. and  $g = 0$ *e*2  $v = 246$  GeV and  $g = 0.64$ :

There are an in this remarkable phenomena in this  $m_W =$  . In the term of the term of the term of the term of term  $m_W = 78.7$  GeV

 $m_H - v \sqrt{2\pi - 125}$  GOV  $m_H = v\sqrt{2\lambda} = 125$  GeV

#### **Quantum corrections** Guantum cor **parameters**  $\overline{\bullet}$  trees-level predictions of the theory, with subsequently measurements probing the problem of  $\bf Q$ uantum ! ions

#### crease the W boson mass by terms proportional to  $\mathcal{N}$  $\mathcal{L}_{\text{G}}$  $\sim$  $\sim$  $\sim$  $\sim$  $\sim$  $\mathbf{z}$  $\overline{\phantom{a}}$ of California Santa Cruz, Santa Cruz, CA 95064, <sup>g</sup>Cornell University, Ithaca, NY 14853, <sup>h</sup>University of Cyprus, Nicosia CY-1678, Cyprus, <sup>i</sup>University College Dublin, Dublin 4, Ireland, <sup>j</sup>University of Edinburgh, Edinburgh EH9 3JZ, United Kingdom, <sup>k</sup>University of Heidelberg, D-69120 Heidelberg, Ger $m_{\pi}^2$  $m<sub>W</sub>$ <sup>n</sup>Nagasaki Institute of Applied Science, Nagasaki, Japan, *W<sup>+</sup> W<sup>+</sup> t b* FIG. 1: The one-loop contribution to the W boson mass  $m_W^2 = \frac{\hbar^3}{c} \frac{\pi \alpha_{EM}}{\sqrt{2}C (1 - m^2/\sqrt{m^2}) (1 - m^2/\sqrt{m^2})}$  $\mathcal{L}$  is the precision of the precision of the W boson mass  $\mathcal{L}$ m , making improved measurements of making including in making in making in making in priority in priority in probing the masses and electronic states  $\widehat{t}$ plies of new hypothetical particles. We describe in  $\mathcal{P}$  $\mathcal{M}$  dragger must precise  $\bigcup$  $\begin{array}{cc} \bullet & \bullet \\ \bullet & \bullet \end{array}$ other precisely measured parameters in the "on-shell" scheme as [4]:  $\mathcal{C}_{0}$  $\pi\alpha_{EM}$  $\sqrt{2}G_F(1-m_W^2/m_Z^2)(1-\Delta r)$ , (1)  $W^+$ measured the properties of new unit function  $\mathcal{C}$  $\int_0^w$  for  $\sqrt{2}G_F(1-m_W^2/m_Z^2)(1-\Delta r)$

Gauge quantum corrections

<sup>o</sup>University de Oviedo, E-33007 Oviedo, Spain, <sup>p</sup>Queen Mary's

$$
\Delta r_{tb} = \frac{c}{\hbar^3} \frac{-3G_F m_W^2}{8\sqrt{2}\pi^2 (m_Z^2 - m_W^2)} \times \left[ m_t^2 + m_b^2 - \frac{2m_t^2 m_b^2}{m_t^2 - m_b^2} \ln(m_t^2 / m_b^2) \right]
$$

 $\mathbf{f}_\mathrm{max} = \mathbf{f}_\mathrm{max} + \mathbf{$ 

 $81358 \pm 4$  MeV  $GISSO = I$  into  $V$ SM calculation of W boson mass vields  $\frac{1}{2}$  measurement of 11 500011 made yields SM calculation of W boson mass yields

,



Higgs quantum corrections

 $L$ iggs quentum sorreation with a contribution proportion proportions a contribution of  $\mathbf{r}_i$ 

from supersymmetric particles can be as large as large as  $\mathbf{t}$ 

eral hundred MeV/c<sup>2</sup> [7].

Naively integrating to a cutoff scale  $\Lambda$ :

$$
\Delta m_H = \frac{3g^2m_t^2}{16\pi^2m_W^2}\Lambda^2
$$

If there is no new physics up to scale  $\Lambda$ then we need 'fine-tuning' to cancel the quantum corrections oretical corrections beyond second order, which have if there is no new physics up to sca

**1% fine tuning:**  $\Lambda = 6.6$  **TeV** 

**Motivates TeV-scale new physics** *W<sup>+</sup> W<sup>+</sup>*

**duniversity Library Brussels, Brussels, Brussels, Brussels, Brussels, Brussels, Brussels, Brussels, Brussels, B** euniversity of California, Irvine, Irv ally, the lighter the squark masses and the larger the

culture site of Bristol, Bris

of California Santa Cruz, Santa Cruz, CA 95064, <sup>g</sup>Cornell

#### **W** boson mass 8 √ 2 π<sup>2</sup>

The W boson mass is the most sensitive observable to sources of 'naturalness'  $\cos$  of 'n <u>b is the moot condition operatives to codicted</u> or materialities , 0)  $\alpha$ 

### Classic example: **Supersymmetry**



is in deutron by Bigging of District is dependent of the set of th<br>In the set of the set o Mass splittings in supersymmetric isospin doublets: different mass shifts for W & Z bosons

#### **W boson mass** vertx coupling of the <sup>Z</sup> to leptons (l). If this vertex is written as <sup>i</sup>¯lγ<sup>µ</sup>(g<sup>V</sup> <sup>−</sup> <sup>g</sup>Aγ5)lZ<sup>µ</sup> <sup>1</sup> <sup>−</sup> Reg<sup>V</sup> "  $\mathbf W$  boson mass mostly concentrate on the effective leptonic weak mixing and  $\sim$ for m<sup>t</sup> # m<sup>b</sup> eq. (2.61) reduces to the well known quadratic correction

Difference in corrections to W and Z propagators encapsulated by *ρ* parameter At the the the the this single this single this interest the sine of the sine of the sine of the single, single Difference in corrections to w and Z propagators encaps / and Z prop  $\overline{\phantom{a}}$ ce in corrections to W a



#### **W boson mass** *5.2. The standard model effective field theory* The SMEFT is a constructed out of the SM constructed out of a series of SUL(3)  $\frac{1}{2}$   $\frac$

More generally the SM effective field theory parameterizes high-scale effects Sulorally the Sir

$$
\mathcal{L}_{SMEFT} = \mathcal{L}_{SM} + \mathcal{L}^{(5)} + \mathcal{L}^{(6)} + \mathcal{L}^{(7)} + \cdots, \qquad \mathcal{L}^{(d)} = \sum_{i=1}^{n_d} \frac{C_i^{(d)}}{\Lambda^{d-4}} Q_i^{(d)}
$$
 for  $d > 4$ .  
\n**1. Brivio and M. Trott, Phys. Rep. 793 (2019) 1  
\nPhys. Rep. 793 (2019) 1  
\n
$$
\frac{\delta m_W}{m_W} = (0.34c_{HD} + 0.72c_{HWB} + 0.37c_{Hl3} - 0.19c_{ll1}) \frac{v^2}{\Lambda^2}
$$
\n
$$
V_i
$$**

*V*(*H*†  $\overline{n}$  $W$ <sup> $/m$ </sup> $W$  = **•**<br>2 *CH H*  $\frac{1}{2}$ For  $\delta m_W^{} / m_W^{} = 0.1$   $\%$  and c<sub>HD</sub>=1,  $\Lambda = 4.5$  TeV e.g. Z' boson

**For** *v*2  $\overline{a}$  $\frac{1}{4}$ 3*CH v*<sup>2</sup>  $\sqrt{2}$ ◆  $\mathbf{I}$ *For δ* $m_W/m_W = 0.1\%$  *and c<sub>HWB</sub>=1,*  $\Lambda = 6.6$  *TeV* on the exact solution (  $\mathbf{e}$ ,  $\mathbf{g}$ ,  $\mathbf{c}$ ,  $\mathbf{c}$ <sup>2</sup> <sup>3</sup>*CH v*2)*/*(3*CH* ) to first order in *CH* . This expansion assumes a mass gap to the e.g. compositeness

 $\alpha$  was suppressed in the previous equations. We absorb the  $\alpha$ Smaller  $c_i \rightarrow s$ maller  $\Lambda$ 

### **W boson mass measurements**



# **CDF II measurement of the W boson mass**



CDF II detector consists of

- silicon vertex detector
- large drift chamber
- coarse calorimeter towers
- outer muon chambers

 $s=1.96$  TeV proton-antiproton collisions from the Fermilab Tevatron





#### **DAMAAANIKA** Manufade Ester **CONE II MAQQUIYAMANT AF THA IN hACAN M** MAY TOWER ENDING U meascurement of the W hees **CDF II measurement of the W boson mass**

10



W bosons identified in their decays to  $e\nu$  and  $\mu\nu$  $\mathbf{r}$ 

Mass measured by fitting template distributions of transverse momentum and mass

$$
m_T = \sqrt{2 p_T^{\;l} \rlap{\,/} \
$$





#### **Calibrations** DAIS E transverse Eta-Phi LEGO Plot Max tower E= 40.7 Min tower E= 0.50 N clusters= University of the Continued of the



Measurement requires precise calibrations and momentum scale and resoution

 $$ 





#### **Calibrations** pairs. Decays to quark pairs are not observable direction and neglecting particle masses. The event assuming massless particles using calorimeter tower DAIS E transverse Eta-Phi LEGO Plot University of the Continued of the



Measurement requires precise calibrations and momentum scale and resoution mentum vector  $\mathcal{L}$  . The component of recoil pro-Measurement requires precise calibrations

$$
\vec{p}_{T} = -(\vec{p}_{T}^{l} + \vec{u}_{T}) \qquad p_{T}^{l}
$$
\nRecoil scale\n
$$
p_{T}^{W}
$$
\n
$$
\downarrow
$$
\n
$$
p_{T}^{W}
$$
\n
$$
\downarrow
$$
\n
$$
p_{T}^{V}
$$





## **Detector simulation**

### **Developed custom simulation for analysis**

Models ionization energy loss, multiple scattering, bremsstrahlung, photon conversion, Compton scattering

Acceptance map for muon detectors

Parameterized GEANT4 model of electromagnetic calorimeter showers **0.08** Includes shower losses due to finite calorimeter thickness **0.09 6 6** Includes shower losses due to finite calorime

Hit-level model of central outer tracker **0.05** *Layer-by-layer resolution functions and efficiencies*  **0.03 0.04 2 2**

Material map of inner silicon detector **1** *Includes radiation lengths and Bethe-Bloch terms* 





**0.06** Kotwal & CH, NIMA 729, 25 (2013)

### **First step is to align the drift chamber** (the "central outer tracker" or COT)

Two degrees of freedom (shift & rotation) for each of 2520 cells made up of twelve sense wires constrained using hit residuals from cosmic-ray tracks



Before After After After After After After After  $\mathcal{L}(\mathcal{L})$  the transverse (x, y) coordinates of the center of the center of each center of each cell, at the longi-Kotwal & CH, NIMA 762, 85 (2014)

#### **Muon momentum calibration** correction to the sinusoidal amplitude a(R) is also introduced to reduce pull sain ba sl1 45 -78

#### **First step is to align the drift chamber** (the "central outer tracker" or COT) resulting in a(R) varying between 45 µm and 49 µm for the different superight and arm onar

**Sing directive between in parameters -10** coming and outgoing cosmic-ray tra **m)** µ $\ln$ on of the wire within the chamber using difference between fit parameters of incoming and outgoing cosmic-ray tracks Two parameters for the electrostatic deflection of the wire within the chamber constrained Dr the electrostatic deflection of the wire within the chamber constrained is illustrated in Fig. 15.  $\overline{A}$ Parameters describing the electrostatic deflection varying with azimuth as given by









 ${\bf Second \ step \ is \ to \ calibrate \ the \ momentum \ scale \ using \ } J/\psi$  decays to muons

### **Simulation:**

Adjust kinematics to match the data Model resonance shape using hit-level simulation and NLO form factor for QED radiation



### ${\bf Second \ step \ is \ to \ calibrate \ the \ momentum \ scale \ using \ } J/\psi$  decays to muons

#### **Simulation corrections:**

Correct the length scale of the tracker with mass measurement as a function of  $\Delta \cot\theta$ Correct the amount of upstream material with mass measurement as a function of  $p_T^{-1}$ 



### Third step is to calibrate the scale using  $\Upsilon$  decays to muons

Compare fit results with and without constraining the track to the collision point



Fig. 12: Distribution of the best-fit temperature mass for the best-fit temperature (histograms) and the properties (histograms) and the M  $\mu$ <sub>N</sub>  $\mu$ sample used to call the muon tracks are reconstructed with tracks are reconstructed with  $\alpha$ with constraint without constraint

### **Muon momentum calibration** detectable collision products accompanying the W or Z boson is defined as the hadronic recoil is defined as the

**Final step is to measure the Z boson mass**  Final stan is to moneura the 7 hopen mose Final Step is to measure the  $\boldsymbol{\Sigma}$  boson mass

 $M_Z = 91\,192.0 \pm 6.4_{stat} \pm 4.0_{sys}$  MeV  $\mathcal{L}$  because  $\mathcal{L}$  because  $\mathcal{L}$  and  $\mathcal{L}$   $\math$ 

Result blinded with [-50,50] MeV offset until previous steps were complete Combine all measurements into a final charged-track momentum scale Boqult blinded with [50.50] MeV offect until previous steps were complete ficall simula with [ 00,00] we vanded and previous steps were complete.<br>Combine all meeoguremente into a final eberged treak mementum eeele **WE WE ARE AN INCREDITED** 



## **Electron momentum calibration**

**First step is to transfer the track calibration to the calorimeter (E/p) using W & Z decays** 

#### **Data corrections:**

Use mean E/p to remove time dependence & response variations in tower Fit ratio of calorimeter energy to track momentum to correct each tower in *η*  $\overline{\phantom{a}}$  combination of scintillator and aluminum. Table 1 shows the materials in  $\overline{\phantom{a}}$ 





### **Electron momentum calibration** CEM 2.175 × 3.175 × 3.175 × 3.175 × 3.175 × 3.175 × 3.175 × 3.175 × 3.175 × 3.175 × 3.175 × 3.175 × 3.175 × 3.

#### **First step is to transfer the track calibration to the calorimeter (E/p) using W & Z decays**  CES aluminum 6.0 0.07 Solenoid complete the solenoid and the solenoid of the solenoid control in the solenoid of the

Parameterize calorimeter shower deposition and leakage based on GEANT4 Determine small calorimeter thickness corrections using region of low E/p in data Fit calorimeter scale as a function of  $E_T$  to correct for any remaining energy dependence CEM front-plate aluminum 14.0 0.157  $T$  are all total  $T$  and  $T$  are all the  $T$  and  $T$  are all the  $T$ 





for this extra plastic material.

Kotwal & CH, NIMA 729, 25 (2013)



Tune energy loss due to material upstream of the tracker (high E/p)

Sampling resolution given by  $\sigma_E/E = \sqrt{\frac{E}{E}}$   $\frac{1}{E}$   $\frac{1}{E}$   $\frac{1}{E}$  with  $\kappa = 0.7 - 1.1$  % increasing with tower 12.6 %  $E_T\,$  $+\kappa^2$  with  $\kappa = 0.7 - 1.1$  % increasing with tower  $\eta$ FIG. S13: (Left) Measured calorimeter energy scale in bins of electron tower in W → eν data after corrections are  $\frac{12.6\%}{\pi}$ Sampling resolution given by  $\sigma_E/E = \sqrt{\frac{1-2\pi}{D}} + \kappa^2$  with  $\kappa = 0.7 - 1.1$  % increasing with tower  $\eta$ the fit region is the fit region is  $E_T$ 



### **Electron momentum calibration** design of the detector, such products have small transverse momentum. The transverse momentum vector such a su

Second step is the measurement of the Z boson mass

 $M_Z = 91\;194.3 \pm 13.8_{stat} \pm 7.6_{sys}$  MeV

As a consistency check measure mass using only track information e.g.  $M_Z^{} = 91\,~215.2 \pm 22.4$  MeV for non-radiative electrons (E/p<1.1)  $\mathsf{A}$ 

#### Same blinding as for muon channel



# **Recoil momentum calibration**

*ϕu*

#### **First step is the alignment of the calorimeters**

Misalignments relative to the beam axis cause a modulation in the recoil direction Alignment performed separately for each run period using minimum-bias data

### **Second step is the reconstruction of the recoil**

Remove towers traversed by identified leptons Remove corresponding recoil energy in simulation using towers rotated by 90° *validate using towers rotated by 180o*









rec  $\mathbf r$ 

ξ

 $\vec{p}_{T}^{\,l}$ 

### **Recoil momentum calibration PC 33 %**

### **Fourth step is the calibration of the recoil resolution**  0

Includes jet-like energy and angular resolution, additional dijet fraction term, and pileup ional dijet fraction term





#### **Recoil momentum validation** this χ2, we obtain NW. <sup>s</sup>hard = [0.<sup>828</sup> <sup>±</sup> <sup>0</sup>.028(stat)] GeV1/2. The tuning is The u|| distribution is directly affected by the meascoil momentum validati (Figs. 14 and 16) and the modeling of lepton tower

### W boson recoil distributions validate the model

Most important is the recoil projected along the charged-lepton's momentum  $(u_{||})$ Most important is the recoil projected along the charged-lepton's momentum can be written as:

$$
m_T \approx 2 p_T \sqrt{1+u_{||}/p_T} \approx 2 p_T + u_{||}
$$

 $\vec{p\,}^l_T$ 

 $|u_{\parallel}|$ 

 $\vec{u}_T$ 

ν *T p*

 $\boldsymbol{u}$ 





Triggers with low momentum thresholds (18 GeV) and very loose lepton id

Offline id also loose, efficiencies vary by 2% as hadronic recoil direction changes

**No lepton isolation requirement in trigger or offline selection**





Largest background is  $Z\to\mu\mu$  with one unreconstructed muon: **7.4% of data sample**  $W\to\,\tau\nu$  background is ~1% in each channel: largest background in electron sample  $\mathcal{L}$  in data and includes backgrounds (shaded). The likelihood is computed using events (shaded). The likelihood is computed using events (shaded). The likelihood is computed using events (shaded). The likelihood is c

Background from hadrons misreconstructed as leptons estimated using data: 0.2-0.3% **Deckargund** from be



## **W boson transverse momentum**

#### Boson  $p_T$  impacts the  $p_T$  distributions of the decay leptons

Resbos used to generate events with non-perturbative parameters and NNLL resummation to model the region of low boson  $p_T$ 

Z boson  $\mathsf{p}_\mathsf{T}$  used to constrain the non-perturbative parameter  $\mathsf{g}_2$  and the perturbative coupling  $\alpha_{_S}$ 

### Resbos models W boson  $p_T$  well *uncertainty estimated using DYQT and constrained with data*



#### W boson production and decay on proquetion and s production and decays, Z → ll, are used for calibration.

Parton distributions impact the measurement through lepton acceptance Restriction in *η* reduces the fraction of low-p<sub>T</sub> leptons do modela anoment and cagnitapient de cop

**Small correction applied to update to NNPDF3.1 NNLO PDF**  The set with the most W charge asymmetry measurements at the time  $\overline{a}$ described in this paper and is the most precise single sing paale to innepers. I innel PD



#### **Uncertainty determined using a principal component analysis on the replica set**  description of the CDF II detector is presented in the CDF II detector is presented in the CDF II detector is p ) a principal component analysis on the  $\mathbf{F}$  $\mathsf{p}\mathsf{n}\mathsf{c}\mathsf{a}$  such as initial-states such as initial-states such as in

Measurement sensitive to ~15 eigenvectors Leading 25 eigenvectors used to estimate uncertainty (3.9 MeV) Three general NNLO PDF sets (NNPDF3.1, CT18, and MMHT14) have a range of  $\pm 2.1$  MeV from mean

**Photos resummation with ME corrections used to model final-state photon radiation**  validated by studying the average radiation in EM towers around the charged lepton, *and with the Z mass measurement*











### **W boson mass measurement**





### **The W boson mass is an important parameter in particle physics**

### **Measurement of W boson mass with <10 MeV precision achieved with complete CDF data set**

Result of >20 years of experience with the CDF II detector

0.01% precision required flexibility: all experimental aspects controlled by the analysis team *Reconstruction, alignment, calibration, simulation, analysis* 

Analysis procedures approved pre-unblinding and frozen

Surprising 0.1% deviation from SM motivates expanded study of mw measurements and procedures

### **Backup**









#### **Background fractions TE BACKGround 1** normalization and shape (in parentheses). Where  $\blacksquare$ in the W  $\alpha$  -corresponding to  $\alpha$  $\mathbf{b}$  resolution is taken in the d $\mathbf{c}$ fraction is estimated to be comparisoned that the systematic uncertainties are estimated by comparison by contract uncertainty of  $\mathbf{B}$ made from different impact-parameter regions and from different requirements on the hit residual patterns.



luminosity to obtain the background fraction of (0.01 ± 0.01)% in the current sample.

and 6.8 MeV on MW for the mT , put the mT , p

Fraction δM<sup>W</sup> (MeV)

<sup>∗</sup> → µµ 7.37 ± 0.10 1.6 (0.7) 3.6 (0.3) 0.1 (1.5)

W + 1.0004 + 1.000 +

Hadronic jets 0.01 ± 0.04 0.1 (0.8) -0.6 (0.8) 2.4 (0.5)

Decays in flight 0.20 ± 0.14 1.3 (3.1) 1.3 (5.0) -5.2 (3.2)

Cosmic rays 0.01 ± 0.01 0.3 (0.0) 0.5 (0.0) 0.3 (0.3)

Total 8.47  $\pm$  0.18  $\pm$  0.18  $\pm$  0.18  $\pm$  0.18  $\pm$  0.18 (5.1)  $\pm$  0.18  $\pm$  0.18  $\pm$  0.18  $\pm$  0.18  $\pm$ 

Source (%) may be a set of the set of the



to indicate a negative correlation between fits.

Fraction δM<sup>W</sup> (MeV)

<sup>∗</sup> → ee 0.134 ± 0.003 0.2 (0.3) 0.3 (0.0) 0.0 (0.6)

 $W = \{x_0, y_0, \ldots, y_n\}$  ,  $W = \{x_0, y_0, \ldots, y_n\}$  ,  $W = \{x_0, y_0, \ldots, y_n\}$ 

Total 1.41 ± 0.08 2.3 (1.2) 1.1 (6.5) 6.2 (1.3)

Hadronic jets 0.34 ± 0.08 2.2 (1.2) 0.9 (6.5) 6.2 (−1.1)

<sup>T</sup> fit p<sup>ν</sup>

# **Initial state LO & NLO**



## **Recoil in W & Z events**



## **Recoil projections in W events**



# **Recoil model parameters**



*ATLAS, arXiv:1906.02025*





1

**P** <sup>χ</sup>**2 = 82 %**





## **Recoil reconstruction in muon channel**





#### **Electron momentum calibration** Electron Inonientum Calibration.



### **Electron momentum calibration**





## **Track momentum calibration**

### **Residual tracker misalignments studied using difference in E/p between electrons and positrons**

Correction as a function of polar angle applied to measured tracks from W and Z decays

Linear dependence on cot theta would cause a bias in the mw mass fit

No linear correction required, statistical precision from E/p constrains the bias to <0.8 MeV



# **Measurement updates**



#### **Electroweak observables at dimension 6** Example a complete at dimensio Here is www.gar viver: Yalliere di lilliere fields in the SMEFT. Electrow  $\overline{\bm{A}}$ 2Mˆ <sup>2</sup> **e** n V∙ servables at rt dimens! Electroweak observables at dimensio where 'ˆ *<sup>W</sup>* = 9 <sup>√</sup>2*G*<sup>ˆ</sup> *<sup>F</sup> <sup>m</sup>*<sup>ˆ</sup> <sup>3</sup> *<sup>W</sup> /(*12π*)*. Here <sup>δ</sup>*G*<sup>ˆ</sup> *<sup>F</sup>* <sup>=</sup> *<sup>C</sup>(*3*) <sup>H</sup>*&*/G*ˆ *<sup>F</sup>* − *C*& &*/*2*G*ˆ *<sup>F</sup>*

C<sup>H</sup> → 0. Many expressions that follow have explicit dependence on v¯<sup>T</sup> , which is related  $\frac{\partial \mathcal{M}_{HKVB}^2}{\partial \mathcal{M}_{HKB}^2}$ ,  $Q_{HD}$ ,  $Q_{HP}^{(1)}$ ,  $Q_{HH}^{(2)}$ ,  $Q_{HH}^{(3)}$  $\delta M^2$  $\overline{a}$ ∑agrá<br>Tagrá ້<br>ວ**່**  $\frac{1}{\sqrt{2}}$ n<br>F  $r$ arg $\widehat{\theta}$ ens $\overline{a}$ ,  $\overline{a}$ , to the  $\frac{1}{q}$  in port !√  $\frac{2}{3}$  $\frac{\partial^2 u}{\partial x^2} = \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} + \frac{\partial^2 u}{\partial$  $\frac{1}{2}$   $\frac{2}{\ell!}\sqrt{\mathcal{Q}}\hat{\mathcal{L}}_H$  $\overline{d_1}$  $\frac{1}{\sqrt{2H}\hat{\theta}}$ \<br>**∖** ceed∫  $\frac{1}{2}$  +4 c  $H^2$ 2.2 Gauge couplings in the SMEFT: ¯g1, g¯<sup>2</sup>  $\mathcal{H}_{\mathcal{L}}^{\mathcal{U}}$  the Lagrangian parameters  $\bar{p}$  ,  $\bar{q}_1$  to the input parameters  $\delta \frac{\delta \sqrt{4z}}{2\hat{M}^2}$  $\not\!\!Z$  $\frac{1}{2M^2} \approx \frac{2\pi\omega}{\omega} \left( \frac{2\mu}{\omega} \frac{\partial^2 \phi}{\partial \omega} \$  $\frac{B \cdot Q_{HD} Q_{H\ell}^{(1)} Q_{H\ell}^{(2)}}{1 + B Q_{H\ell}^{(1)}}$  $\frac{d}{d}H\frac{d}{dS}\hat{\theta}^{\mathbf{Q}}\hat{\theta}^{\mathbf{R}}$  $\frac{d^4}{d^4}\overline{C} \overline{\mathcal{H}}_H^H\hat{q}^0, Q_H$ e Cl $Hq$   $Q$  Ad $\frac{3}{4}Q$ e C  $\mathop{pr}$  $\mathbf{p}$  $\frac{\delta M^2}{2M^2}$  $\boldsymbol{\widetilde{4}}$  $\gamma$   $\hat{M}$  2 <sup>−</sup> <sup>1</sup>  $4\sqrt{2}\hat{C}$  $\sqrt{ }$  $\mathfrak{P}$  $\frac{s}{3}$ Cêccs $\bar{p}$  $Hq$ pr  $\operatorname{imp}\!\mathfrak{M}_{Hq}^{(3)}$ pr aram  $m = -2$  mapping of the mapping of the mapping of the experimental re- $\partial \mathbb{N}$ tagrangian parameters.  $\partial \mathbb{R}$  $\frac{1}{\sqrt{2}}$  where  $\frac{1}{\sqrt{2}}$  on  $\frac{1}{\sqrt{2}}$  on the measured measured when  $\frac{1}{\sqrt{2}}$  $\mathbf{q}$  at  $\mathbf{g}$   $\mathbf{g}$   $\mathbf{g}$   $\mathbf{g}$  are  $\mathbf{g}$  and  $\delta M^2 = A$ te the Lagrangian para  $\delta \hat{u}$  is an analyzed in detail for the Tevatron measurement in Second  $\begin{pmatrix} 1 & -232 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$  $\mathcal{N}$ ang $\mathcal{L}$ tida $\mathcal{N}$ liga $\mathcal{N}$  $\alpha$  is a  $\mu$  in the effect of  $\alpha$  measurements in the measurement of  $\alpha$  measurements in the measurement in the  $\alpha$  $\mathcal{L}_{HVV}$   $\mathcal{L}_{Hq}$  the analytic method of  $H\bar{q}$ .



LG<sup>F</sup>

metric version of the SMEFT is given by  $\mathbb{R}^n$ 

<sup>θ</sup><sup>ˆ</sup> <sup>−</sup> *<sup>s</sup>*<sup>2</sup>

√2

where %<sup>ˆ</sup> <sup>=</sup> *<sup>c</sup>*θ<sup>ˆ</sup> *<sup>s</sup>*θˆ*/(c*<sup>2</sup>

$$
\frac{\delta m_W^2}{\hat{m}_W^2} = \hat{\Delta} \left[ 4 C_{HWB} + \frac{c_{\hat{\theta}}}{s_{\hat{\theta}}} C_{HD} + 4 \frac{s_{\hat{\theta}}}{c_{\hat{\theta}}} C_{H\ell}^{(3)} - 2 \frac{s_{\hat{\theta}}}{c_{\hat{\theta}}} C_{\ell \ell} \right]
$$

Expressing M¯ <sup>2</sup>

*)* 2

θ

<sup>2</sup>*G*ˆ *Fm*ˆ <sup>2</sup>

v¯2 T

g2

 $\mathbb{R}^n$ 

<sup>2</sup> + g<sup>1</sup>

Z



Hl

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 $\frac{1}{2}$ 

 $\frac{1}{2}$  ,  $\frac{1}{2}$  ,  $\frac{1}{2}$  ,  $\frac{1}{2}$ 

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In the SMEFT we note that in the flavour symmetric limit  $\sqrt{ }$ 

 $\sqrt{a}$  although the small shown preserve the SMEFT shown preserve the SMEFT shown preserve the SMEFT shown preserve the left handed shown preserve the small shown preserve the small shown preserve the small shown preserv

<sup>2</sup> + g<sup>2</sup>

Hl

2

#### **W boson mass fit results** The combined momentum calibration is used to measure the Z-boson mass in the dimuon channel (see Fig. 3), w boson mass fit results and random of random of range  $\sim$ The variations of the fitted mass values relative to the nominal results, as the fit regions are varied, are consistent

![](_page_51_Picture_424.jpeg)

![](_page_51_Picture_425.jpeg)