

PY411 / 506  
Computational Physics 2

**Salvatore Rappoccio**

# Fluid Dynamics

- Now turn to fluid dynamics
  - Many particles together acting under conservation of mass and momentum
  - Simple cases are incompressible fluids without friction (“Newtonian fluids”)
- Consider a volume  $V$  inside a fluid. The mass  $m$  inside that volume is the integral of the density  $\rho$  :

$$m = \int \rho dV$$

# Fluid Dynamics

- The change in mass is determined by the rate of change at the surface

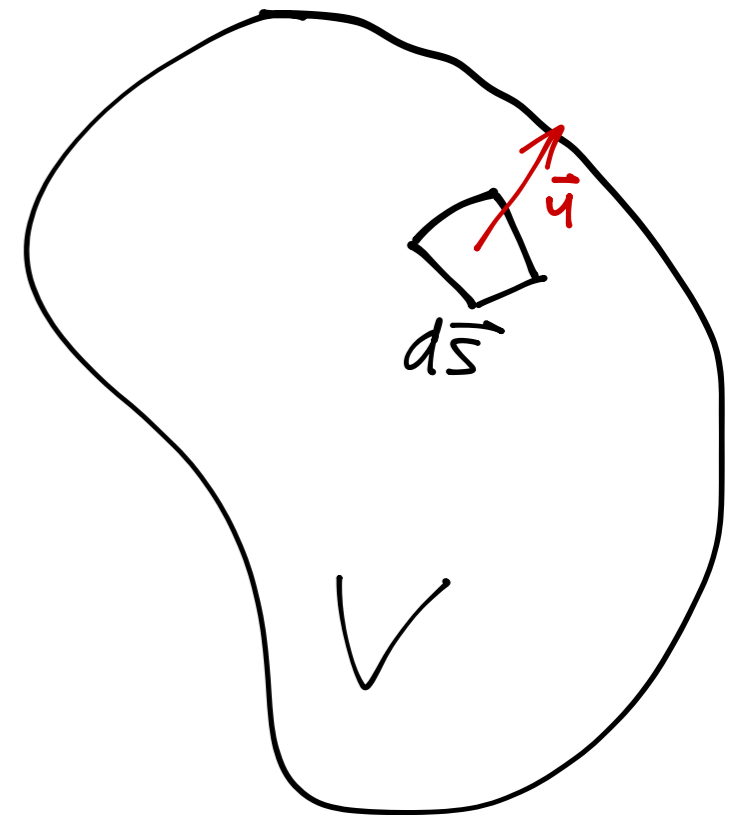
$$\frac{d}{dt} \int \rho dV = - \int d\vec{S} \cdot \rho \vec{u}$$

- Then apply the divergence theorem:

$$\int d\vec{S} \cdot \rho \vec{u} = \frac{d}{dt} \int \rho dV = \int dV \nabla \cdot \rho \vec{u}$$

- So we then obtain the continuity equation:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \vec{u} = 0$$



# Fluid Dynamics

- Now we can apply Newton's second law (conservation of momentum):

$$\rho \frac{d\vec{u}}{dt} = \vec{F}$$

- and then we have:

$$\frac{d\vec{u}}{dt} = \frac{\partial \vec{u}}{\partial t} + (\vec{u} \cdot \vec{\nabla}) \vec{u}$$

Change in fluid velocity  
at a fixed point in space

Change in fluid velocity  
due to motion of fluid  
from neighboring  
points (Advection)

# Fluid Dynamics

- What about friction?
  - Add force of gravity  $\vec{g}$
- Introduce viscous forces:
  - Dynamic viscosity coefficient  $\mu$
  - Bulk viscosity coefficient  $\xi$
  - Pressure  $p$
- Then total force is:

$$\vec{F} = \mu \nabla^2 \vec{u} + (\mu + \xi) \vec{\nabla} (\vec{\nabla} \cdot \vec{u})$$

- Then we get the Navier-Stokes equation with dynamic viscosity  $\nu = \mu/\rho$

$$\frac{\partial \vec{u}}{\partial t} + (\vec{u} \cdot \vec{\nabla}) \vec{u} = \vec{g} - \frac{1}{\rho} \vec{\nabla} p + \nu \nabla^2 \vec{u}$$

# Fluid Dynamics

- If we look at this in 1-d, this is :

$$\frac{\partial u}{\partial t} + \left( u \frac{\partial u}{\partial x} \right) = g - \frac{1}{\rho} \frac{\partial p}{\partial t} + \nu \frac{\partial^2 u}{\partial x^2}$$

- If  $g = 0$  and pressure  $p = 0$  (or gradient = 0), this is Burgers' equation!

$$\frac{\partial u}{\partial t} + \left( u \frac{\partial u}{\partial x} \right) = \nu \frac{\partial^2 u}{\partial x^2}$$

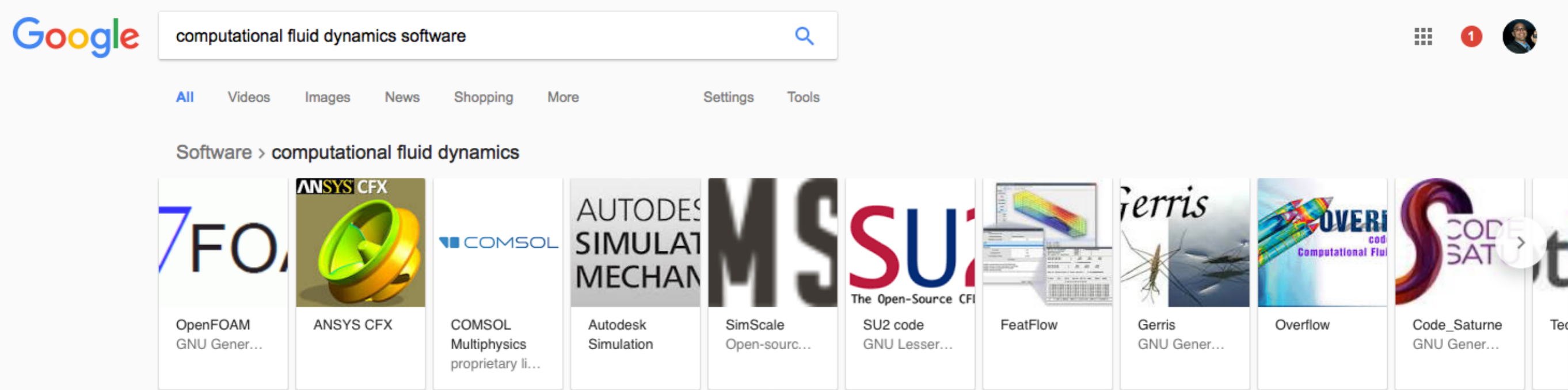
- So shock propagation is a special case of the Navier-Stokes dynamics

# Fluid Dynamics

- What about the full case?
  - It's a \$1M question
  - (No, really, it's a Millennium Prize question).
  - <http://www.claymath.org/millennium-problems>

# Fluid Dynamics

- OK, well, what about simple-ish solutions?
  - Even then, still pretty hard
  - Commercial software in abundance:



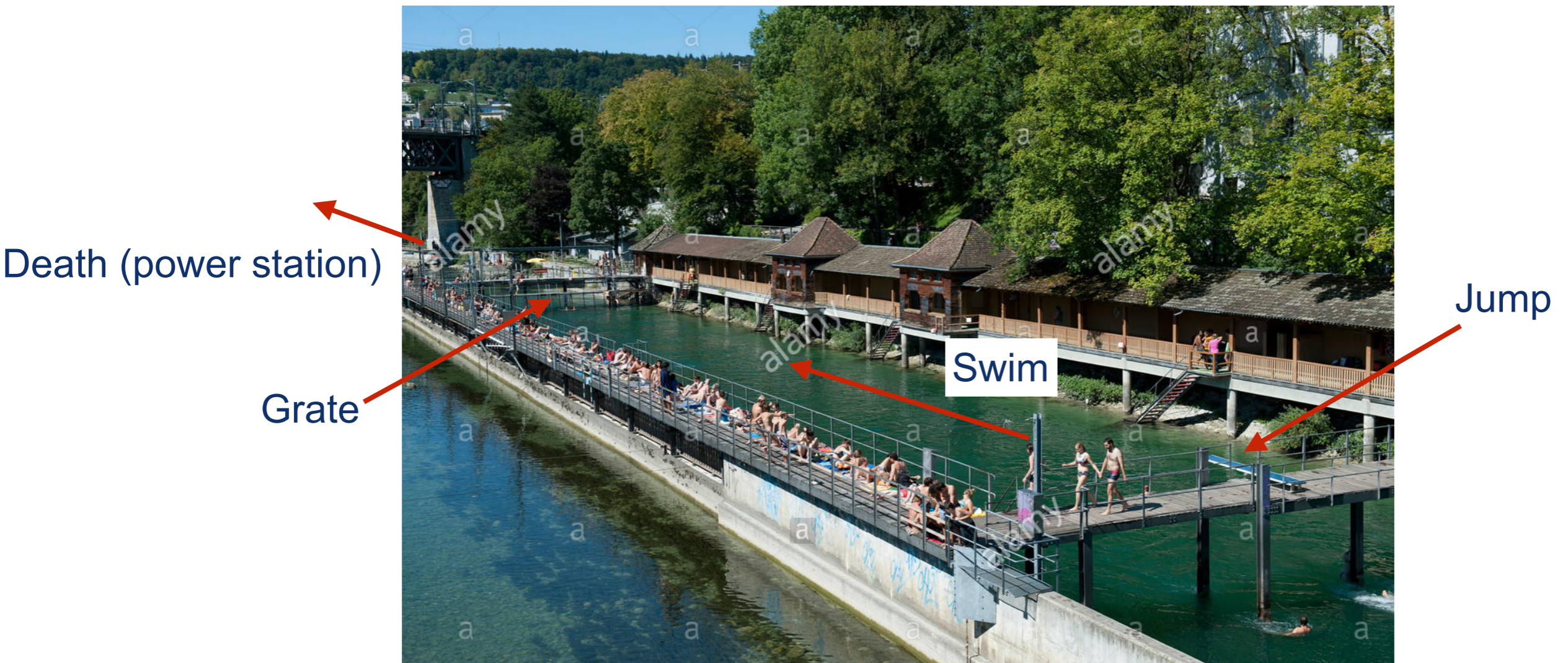
## – Open source examples:

- [https://people.sc.fsu.edu/~jburkardt/py\\_src/navier\\_stokes\\_2d\\_exact/navier\\_stokes\\_2d\\_exact.html](https://people.sc.fsu.edu/~jburkardt/py_src/navier_stokes_2d_exact/navier_stokes_2d_exact.html)
- <http://lorenabarba.com/blog/cfd-python-12-steps-to-navier-stokes/>



# Fluid Dynamics

- Let's look at one special case of flow through a long tube with constant pressure, like the flow in a river.



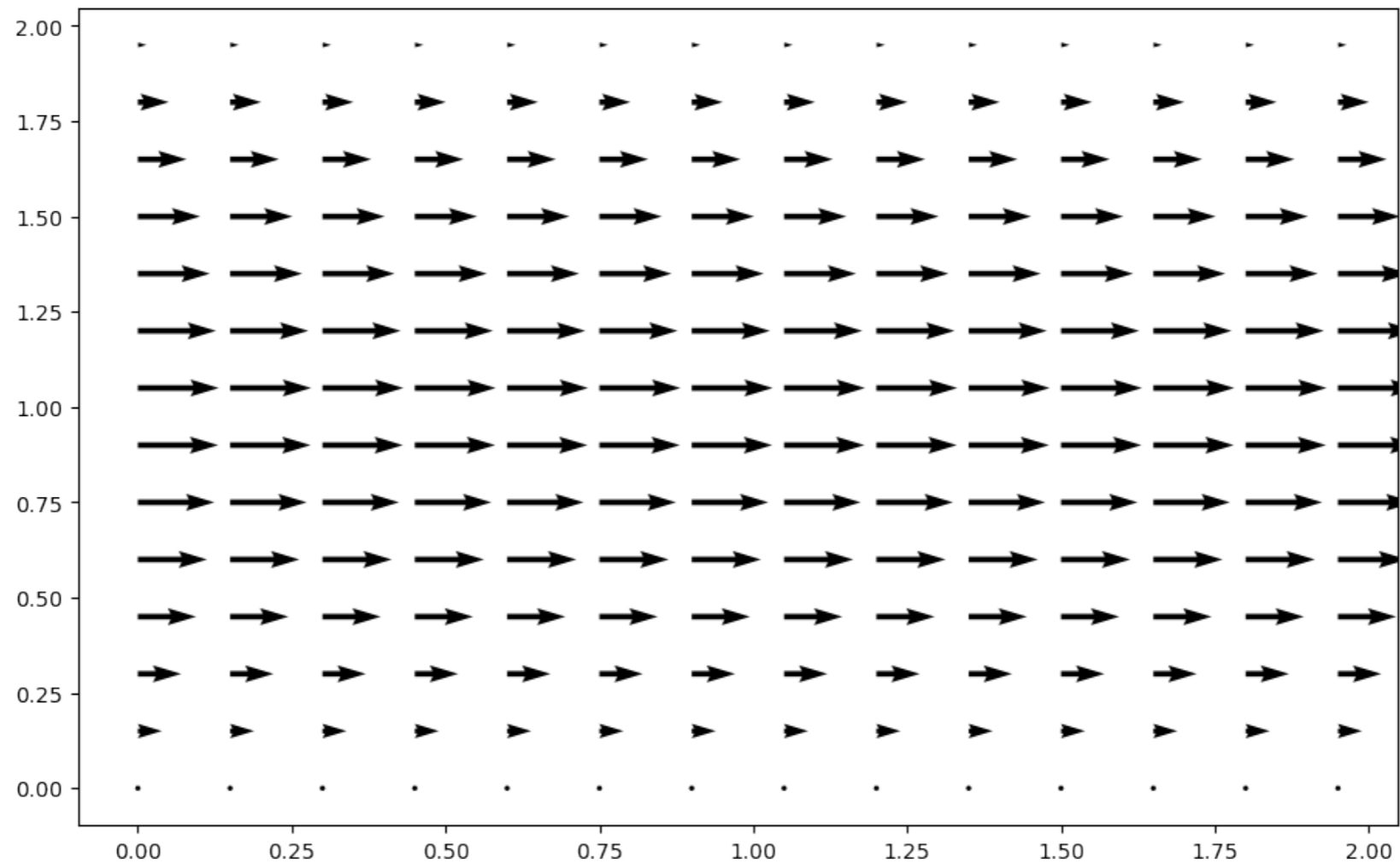
Limmat river in Zurich, CH

# Jupyter!

- Let's play with some Jupyter now.
- We will use the example from here :
  - <http://lorenabarba.com/blog/cfd-python-12-steps-to-navier-stokes/>
- We will walk through it.

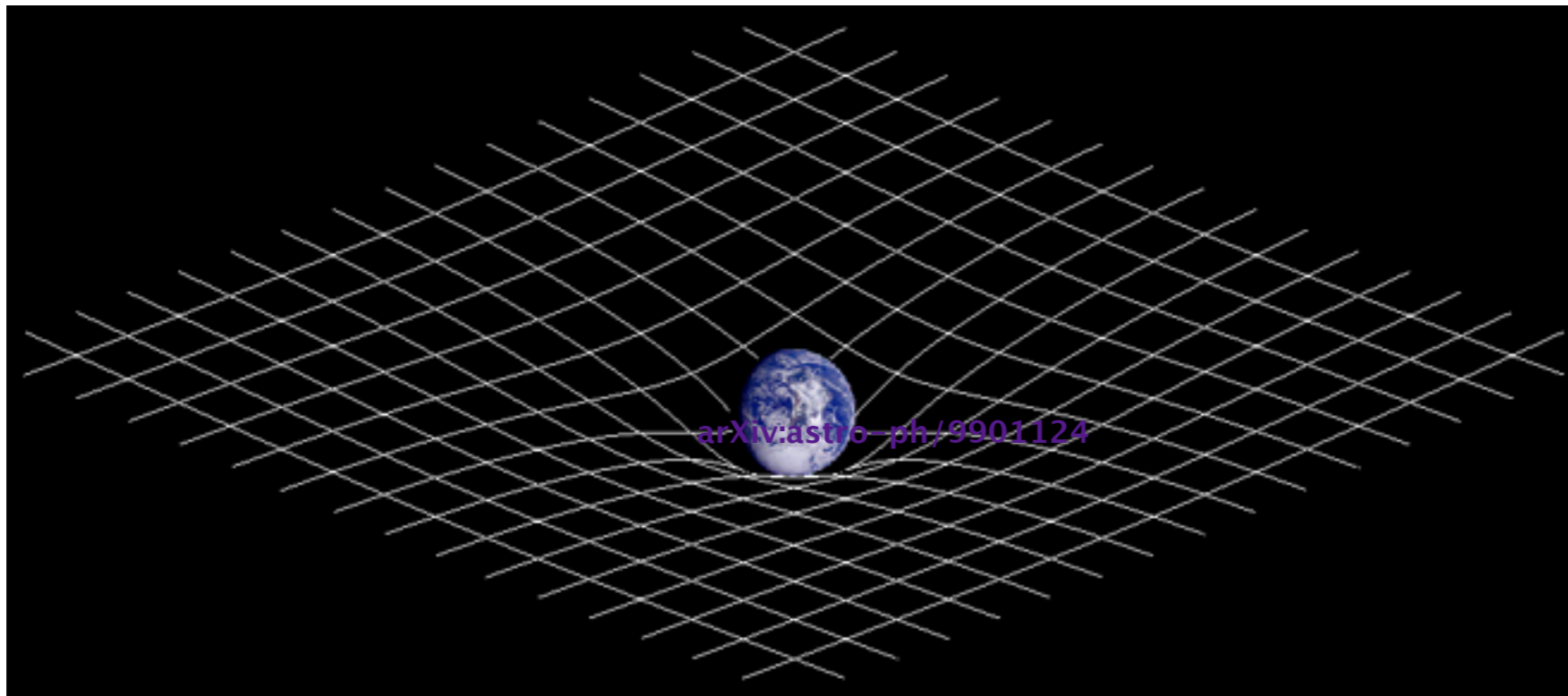
# Fluid Dynamics

- The flow in the center is highest, and the flow at the edges is lowest:



# Recall: General Relativity

- [http://en.wikipedia.org/wiki/General\\_relativity](http://en.wikipedia.org/wiki/General_relativity)
- Relates gravity to the curvature of space-time!



- Objects with mass or energy distort space-time, and this induces a gravitational field

# Recall: General Relativity

- Space-time is a tensor
- So gravity is a tensor
- Einstein's equations :

$$\mathcal{R}_{\mu\nu} - \frac{1}{2}g_{\mu\nu}\mathcal{R} = \frac{8\pi G_N}{c^4}T_{\mu\nu} - g_{\mu\nu}\Lambda$$

The diagram shows the Einstein field equations with several terms labeled and connected by arrows:

- Ricci tensor (gravitational force)**: Points to  $\mathcal{R}_{\mu\nu}$ .
- Metric tensor**: Points to  $g_{\mu\nu}$ .
- Curvature scalar**: Points to  $\mathcal{R}$ .
- Newton's constant**: Points to  $G_N$ .
- Speed of light**: Points to  $c^4$ .
- Stress-energy tensor (energy and momentum density of matter + radiation)**: Points to  $T_{\mu\nu}$ .
- Cosmological constant**: Points to  $\Lambda$ .

# Recall: General Relativity

- That's a huge set of nonlinear partial differential equations, and can be arbitrarily complicated ( $T_{\mu\nu}$  has no constraint to its format)
- A few simple cases can be derived :
  - If spacetime is homogeneous and isotropic, this is the Robertson-Walker metric :

Cosmological scale factor

$$ds^2 \equiv \sum_{\mu=0}^3 \sum_{\nu=0}^3 g_{\mu\nu} dx^\mu dx^\nu = c^2 dt^2 - R^2(t) \left[ \frac{dr^2}{1 - kr^2} + r^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right],$$

- Assuming that the matter+radiation behave like a uniform perfect fluid with density  $\rho$  and pressure  $p$ , this is the Friedmann-Lamaitre equations:

$$H^2 \equiv \left( \frac{\dot{R}}{R} \right)^2 = \frac{8\pi G_N \rho}{3} - \frac{kc^2}{R^2} + \frac{\Lambda c^2}{3}, \quad \frac{\ddot{R}}{R} = -\frac{4\pi G_N}{3} (\rho + 3p) + \frac{\Lambda c^2}{3}.$$

# Recall: Hubble's Law

- Hubble used this equation to determine a linear relationship :

$$rK + X \cos \alpha \cos \delta + Y \sin \alpha \cos \delta + Z \sin \delta = v ,$$

Distance      Longitude      Latitude      Constants      Velocity

- Plotting the data :

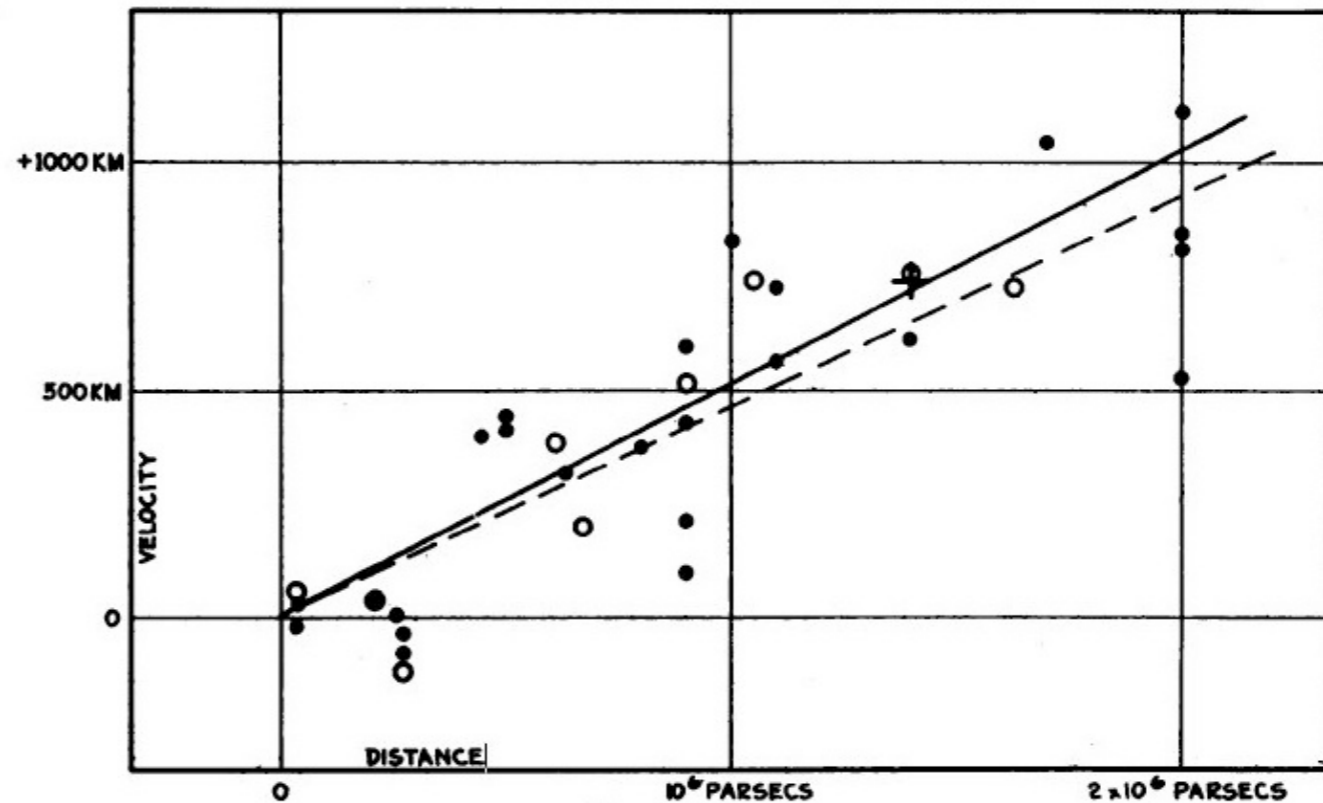


FIGURE 1  
Velocity-Distance Relation among Extra-Galactic Nebulae.

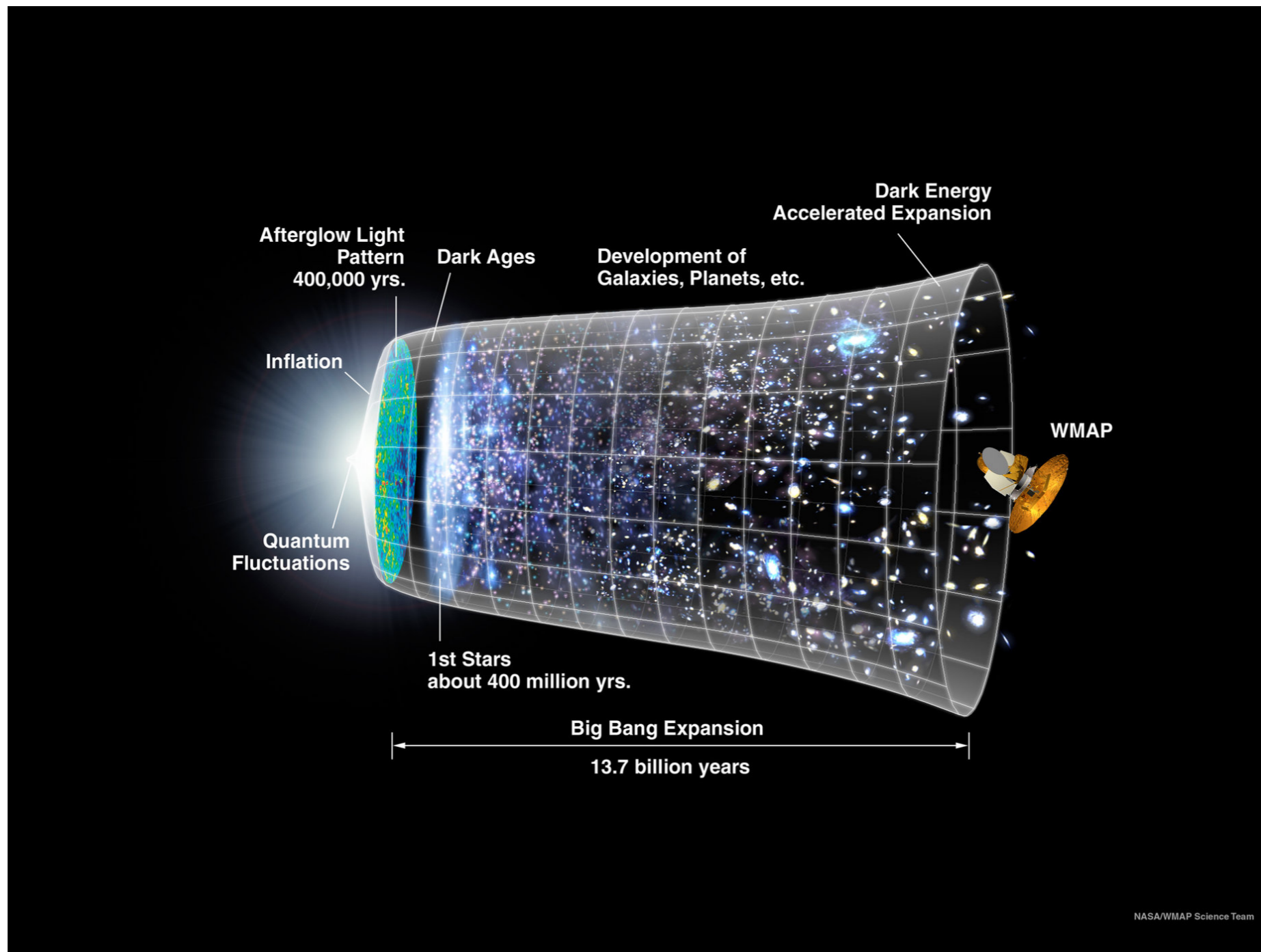
# General Relativity

- Now in a position to calculate some GR numerically instead of just analyzing the data
- Some approaches approximate matter as perfect fluid evolving according to GR
  - Morally equivalent equations to Navier-Stokes, so similar techniques can be used for solutions
- We will investigate inflation
  - For an overview : [arXiv:astro-ph/9901124](https://arxiv.org/abs/astro-ph/9901124)



# General Relativity

- Cosmic inflation



# General Relativity

- Go back to our equations:

$$H^2 \equiv \left(\frac{\dot{R}}{R}\right)^2 = \frac{8\pi G_N \rho}{3} - \frac{kc^2}{R^2} + \frac{\Lambda c^2}{3}, \quad \frac{\ddot{R}}{R} = -\frac{4\pi G_N}{3}(\rho + 3p) + \frac{\Lambda c^2}{3}.$$

Diagram annotations:

- Arrows from "Density" point to  $\rho$  in both equations.
- Arrows from "Constants" point to  $G_N$ ,  $k$ , and  $\Lambda$  in both equations.
- An arrow from "Pressure" points to  $p$  in the second equation.
- An arrow from "Cosmological scale factor" points to  $R$  in the first equation.
- An arrow from  $H^2$  points to the left.

Hubble parameter :  $H(t_0) = 72 \text{ km/s/Mpc}$  at present time

- Values of  $k$ :
  - $-k = 1$ : closed 3-sphere
  - $-k = 0$ : flat
  - $-k = -1$ : open 3-hyperboloid

# General Relativity

- Eliminate Lambda, use conservation of mass+energy, and combine the two equations, you get:

$$\dot{\rho} = -3H(\rho + 3p)$$

- Further simplify the math by scaling mass density and pressure to include all of the constants in the system:

$$H^2 = \left( \frac{\dot{R}}{R} \right)^2 = \frac{8\pi G}{3} \rho - \frac{kc^2}{R^2}$$

$$\dot{H} + H^2 = \left( \frac{\ddot{R}}{R} \right)^2 = -\frac{4\pi G}{3} \left( \rho + \frac{3p}{c^2} \right)$$

# General Relativity

- If we consider a perfect fluid, the pressure is linearly dependent on the density, so you get  $p = w\rho$  for some constant  $w$
- If space is completely flat, and we set  $c=1$ , we have  $k = 0$ , so

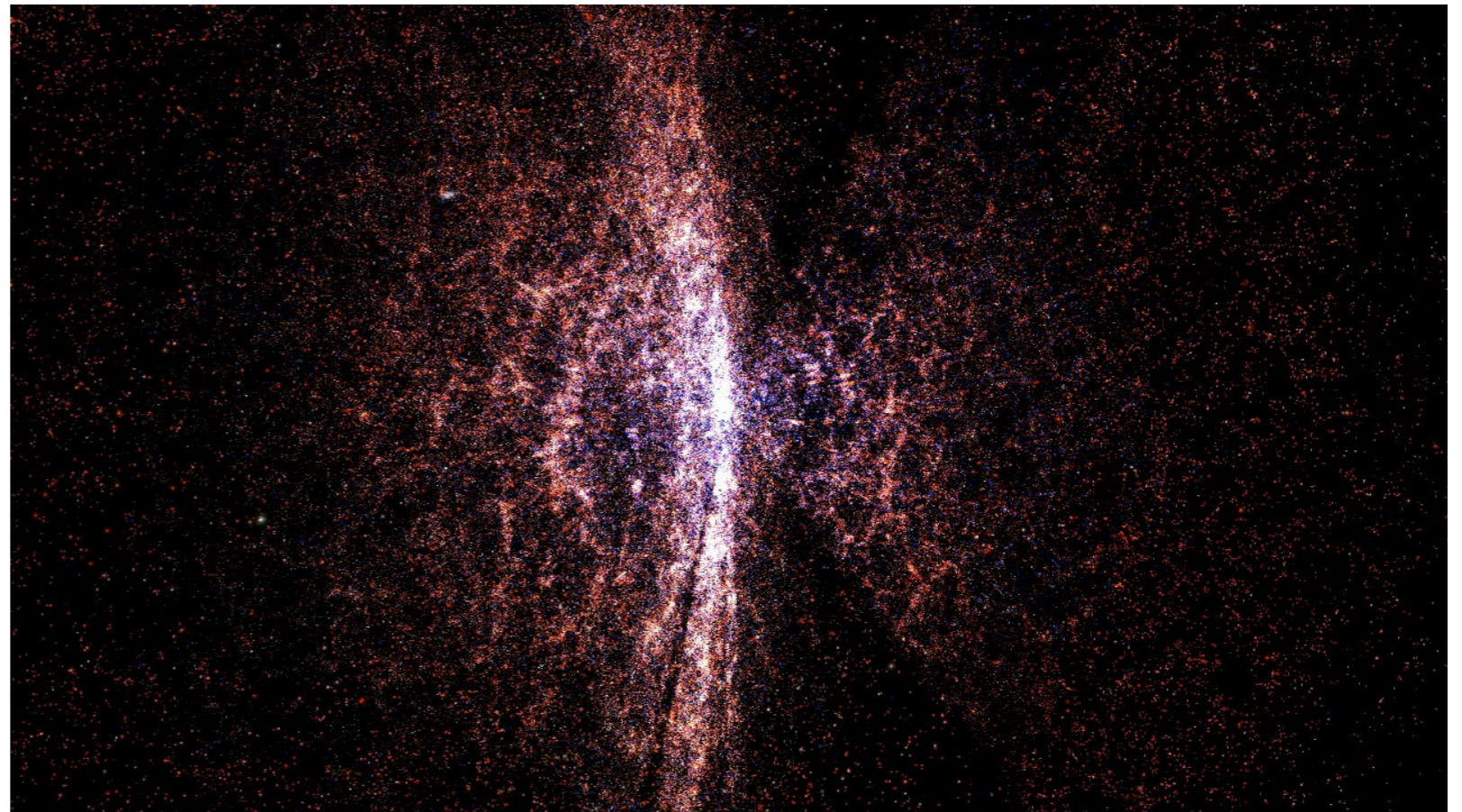
$$H^2 = \left( \frac{\dot{R}}{R} \right)^2 = \frac{8\pi G}{3} \rho$$

$$\dot{H} + H^2 = \frac{\ddot{R}}{R} = -\frac{4\pi G}{3} (1 + 3w) \rho$$

- You will solve this for your homework problem!

# N-body simulations

- Consider  $N$  interacting particles with long range force acting between them (i.e. gravity)
  - Structure formation of galaxies, stars, superclusters, etc
- Recall our previous attempts at the 3-body problem
  - Will now extend to  $N$  bodies



Sloan Digital Sky Survey (NASA)

# N-body simulations

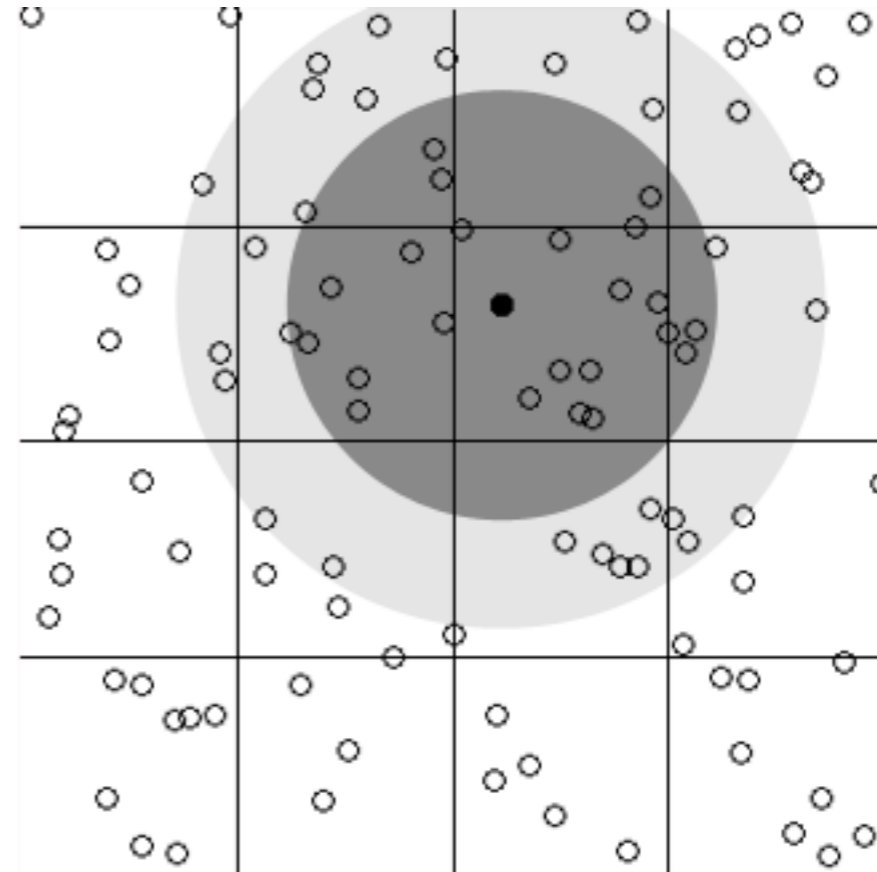
- Long-range interaction very inconvenient here
- Force goes like

$$F \sim \frac{1}{r^2}$$

- But surface area goes like

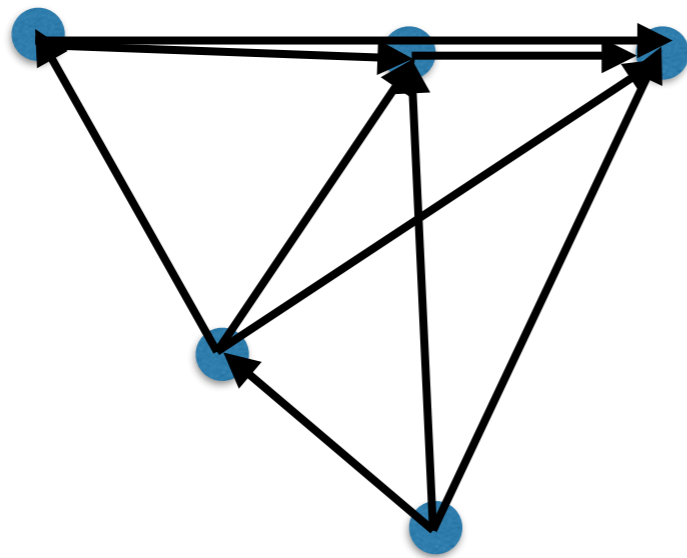
$$A \sim 4\pi r^2$$

- Thus the number of particles at distance  $r$  times the strength of the force is basically constant
  - Cannot truncate!



# N-body simulations

- If we calculate the forces between N bodies, we therefore need to compute  $N(N-1)/2$  forces each iteration



$$\vec{F}_i = \sum_j \vec{F}_{ij} = - \sum_j \frac{Gm_i m_j}{r^3} \vec{r}$$

Simple, but computationally intractable

# N-body simulations

- Try the next level of sophistication in our ODE solvers
  - Currently evolve position and velocity (according to velocity and acceleration, respectively)
  - How about evolving acceleration!?
  - Use the JERK!



# N-body simulations

- Try the next level of sophistication in our ODE solvers
  - Currently evolve position and velocity (according to velocity and acceleration, respectively)
  - How about evolving acceleration!?

– Use the JERK!

Not this jerk



This jerk

$$\vec{J}_{ij} = \dot{\vec{a}}_{ij} = \frac{M_i}{r_{ij}^3} \left[ \vec{v}_{ij} - 3 \frac{\vec{v}_{ij} \cdot \vec{r}_{ij}}{r_{ij}^2} \vec{r}_{ij} \right]$$

# N-body simulations

- First let's look at simple systems where we actually solve the ODEs for all particles
- Going to use third-party software:
  - <https://www.ids.ias.edu/~piet/act/comp/algorithms/starter>
  - Instructions:
  - <https://github.com/rappoccio/PHY410/blob/master/Lecture35/README.md>
- Can't animate and solve the ODEs in real time for any large-ish number of  $N$  (like, 30).

**Works, but computationally intractable!**

# N-body simulations

- Code uses a “predictor-corrector” algorithm:
  - Predicts position and velocity at next time steps

$$\vec{r}_p = \vec{r} + \vec{v}\delta t + \frac{1}{2}\vec{a}\delta t^2 + \frac{1}{6}\vec{J}\delta t^3$$

$$\vec{v}_p = \vec{v} + \vec{a}\delta t + \frac{1}{2}\vec{j}\delta t^2$$

- Computes acceleration and jerk for those predictions using Taylor series

$$\vec{k} \equiv \frac{1}{2}\vec{a}''\delta t^2 = 2(\vec{a} - \vec{a}_p) + \delta t(\vec{J} - \vec{J}_p)$$

$$\vec{l} \equiv \frac{1}{2}\vec{a}'''\delta t^3 = -3(\vec{a} - \vec{a}_p) - \delta t(2\vec{J} - \vec{J}_p)$$

# N-body simulation

- Then corrects position and velocity with the acceleration and jerk:

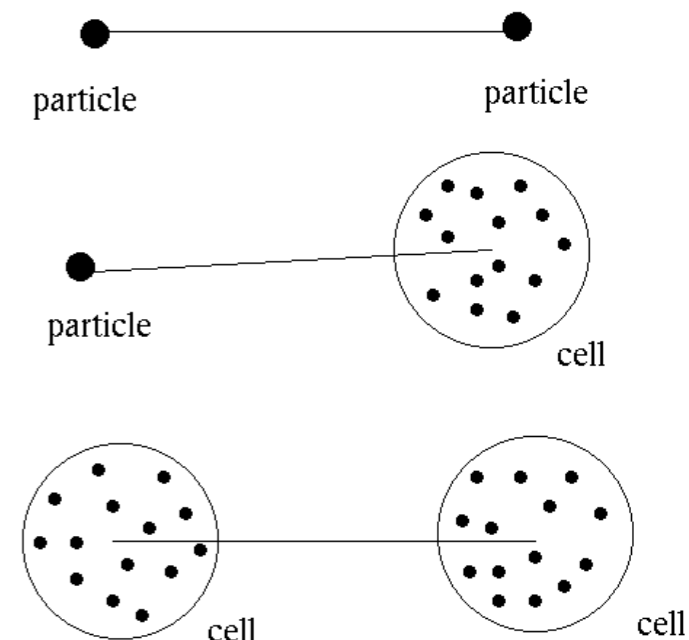
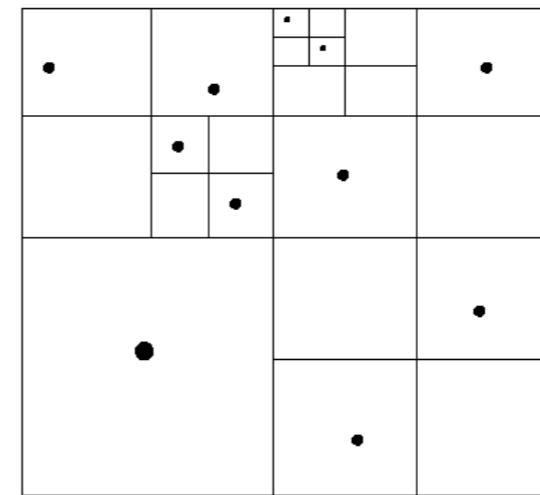
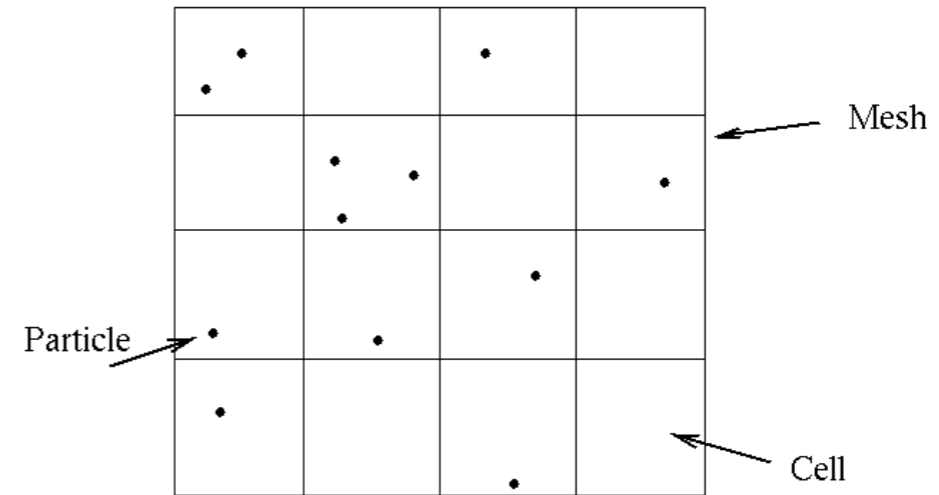
$$\vec{r}_c = \vec{r}_p + \left( \frac{1}{12} \vec{k} + \frac{1}{20} \vec{l} \right) \delta t^2$$

$$\vec{v}_c = \vec{v}_p + \left( \frac{1}{3} \vec{k} + \frac{1}{4} \vec{l} \right) \delta t$$

- Their code also calculates the “least collision time”, i.e. the smallest time between interactions of two objects.
  - Used to adjust the time step, similar to ARK4

# N-body simulations

- We need some kind of approximation then:
  - Particle mesh
    - Create a 3-d lattice, approximate forces from them
    - Good for uniform configurations
    - $O(M \log(M))$  ( $M = \text{num grid points}$ )
  - Trees
    - Partition space into a hierarchy of cubes
    - Compute particle-particle interactions for close interactions
    - Compute particle-cell or cell-cell interactions for far interactions
    - $O(N \log(N))$  ( $N = \text{num particles}$ )



Figures from

<http://www.new-npac.org/projects/cdroms/cewes-1999-06-vol2/cps615course/nbody-materials/nbody-simulations.html>

# N-body simulations

- Fast multipole methods
  - Expands Green's function in a multipole
  - $O(N)$  ( $N$  = num particles)
- Fluid dynamic approximations
  - Approximate by PDEs, not individual particles
  - Only applicable in certain situations

-

# N-body simulations

- Lots of codes out there!
  - AREPO:
    - <http://wwwmpa.mpa-garching.mpg.de/~volker/arepo/>
  - GADGET:
    - <http://wwwmpa.mpa-garching.mpg.de/gadget/>
- Example of simulation from GADGET from a former Comp. Phys. student (Leigh Korbel) for his Master's project

