BSM Higgs Flavoured Correlations

by Arturo de Giorgi (Madrid, IFT)

based on **2109.07490** - *JHEP* 05 (2022)

in collaboration with J. Alonso-Gonzalez (Madrid, IFT), L. Merlo (Madrid, IFT), S. Pokorski (Warsaw U.)

Contribution to: LHC Higgs Working Group - Common WG2 and WG3 - CP violation and Higgs Sector



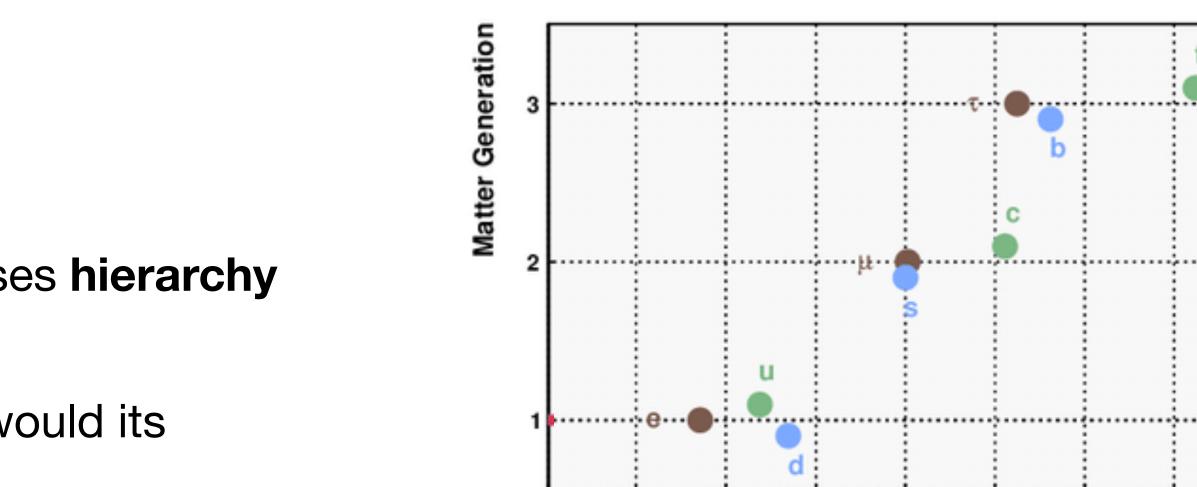
June 23, 2022





1. **SM Flavour-puzzle**: lack of explanation for masses **hierarchy** in the SM

⇒ is there a deeper **underlying theory**? What would its **consequences** imply?



10⁶

10⁷

10⁸

10

10⁵

10⁴



10¹⁰

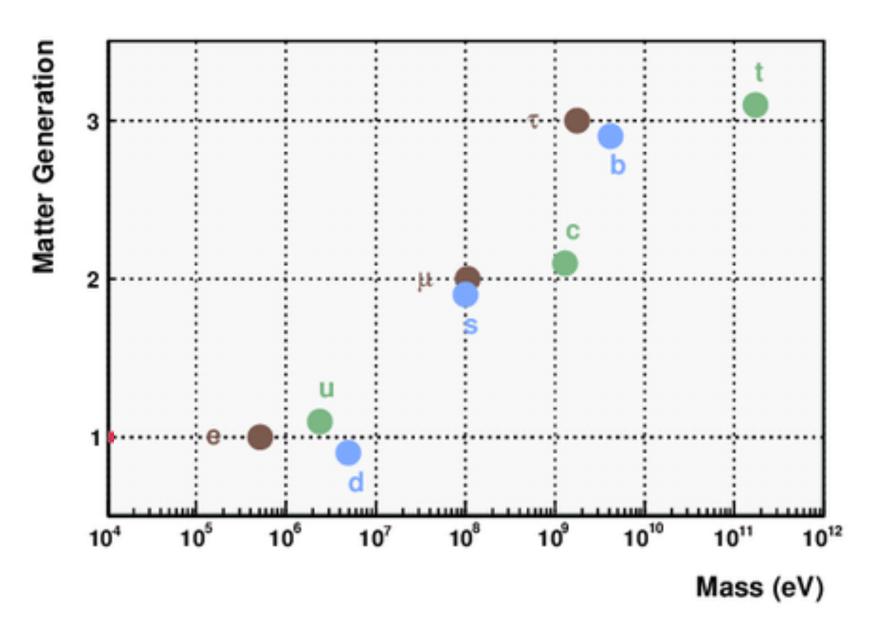


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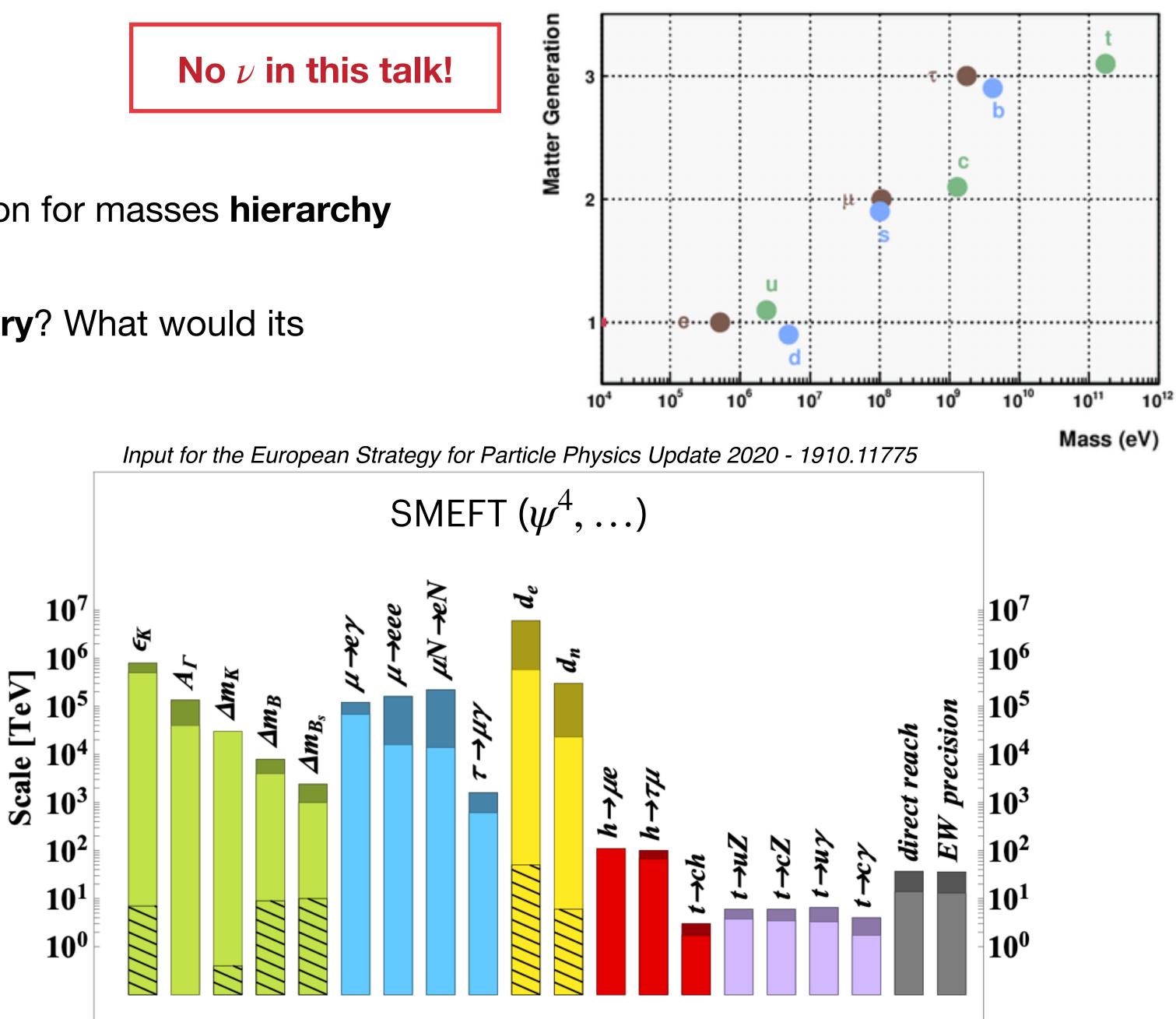




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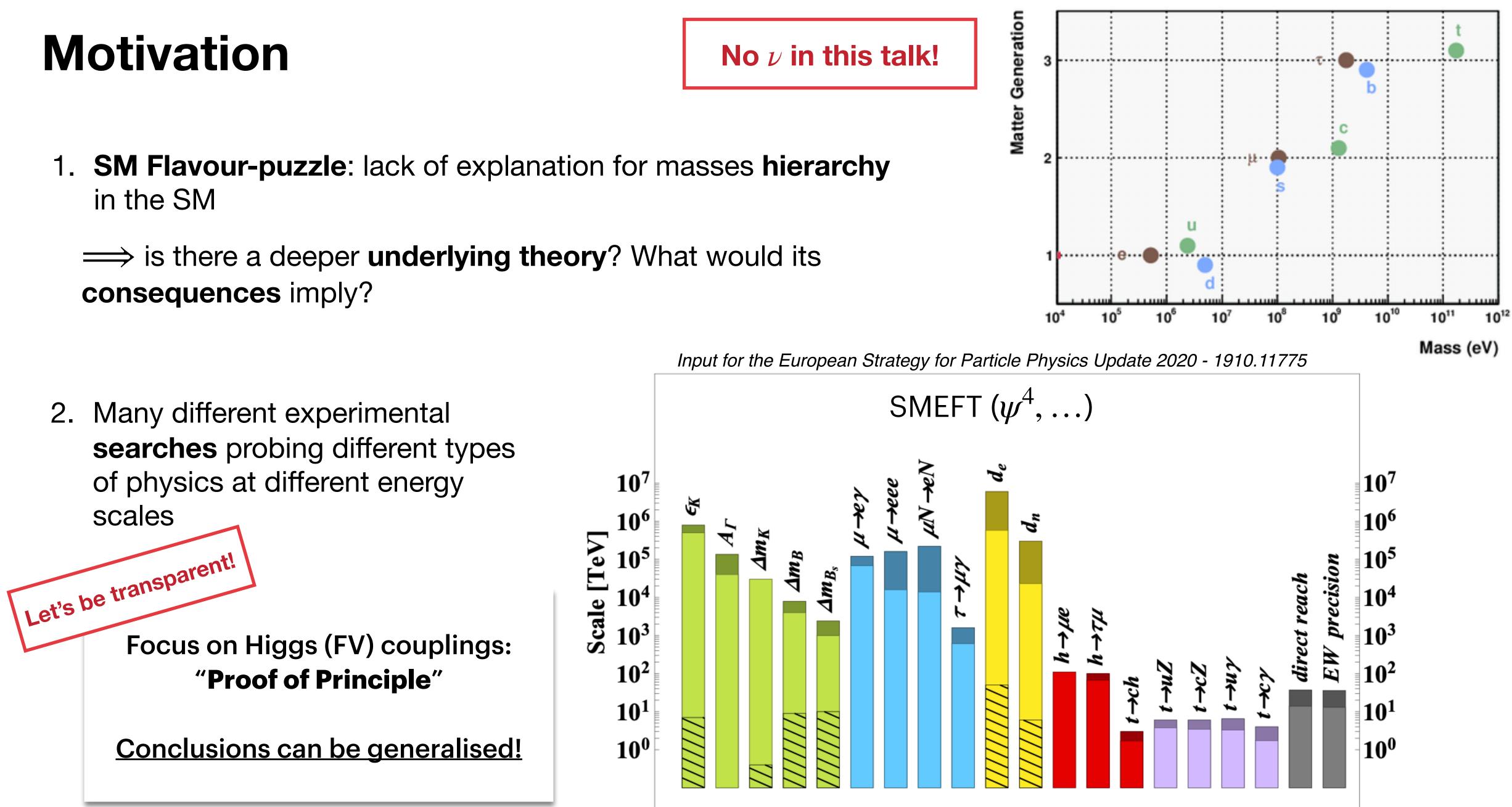
2. Many different experimental searches probing different types of physics at different energy scales



Observable



in the SM



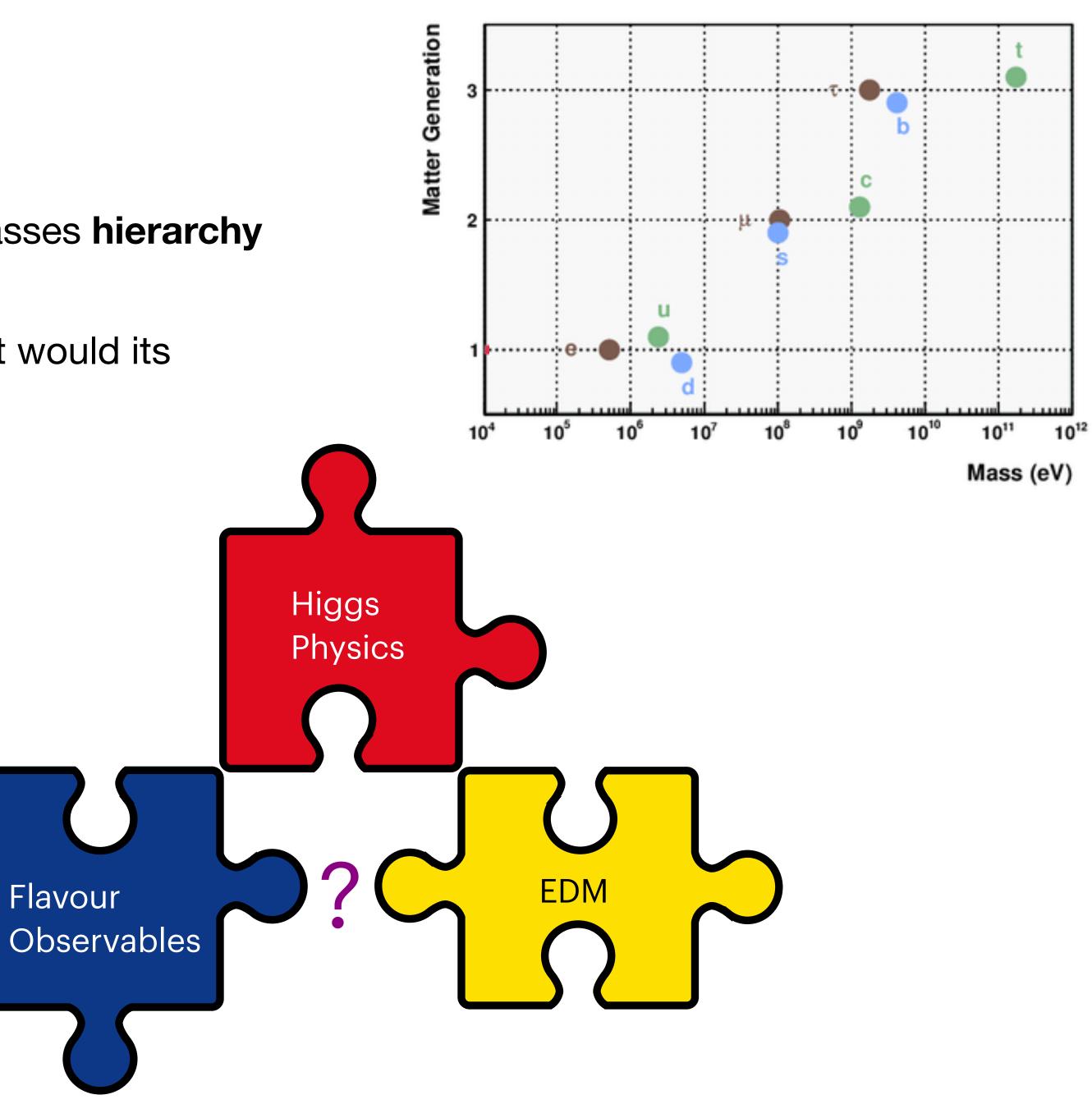
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 \implies how far can we push the **synergy** between **Higgs physics**, **EDMs** and **flavour observables**?

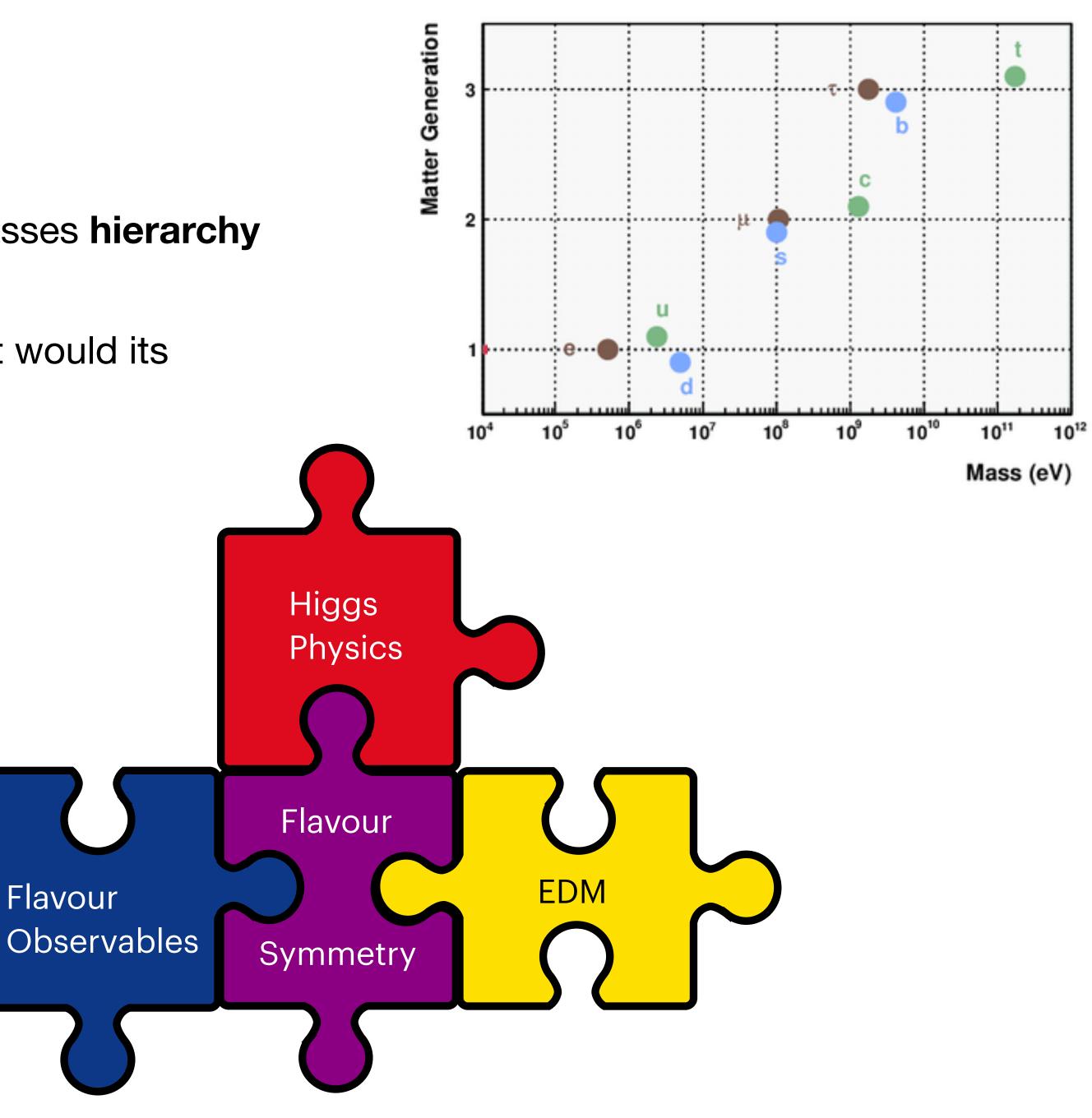


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Outlook

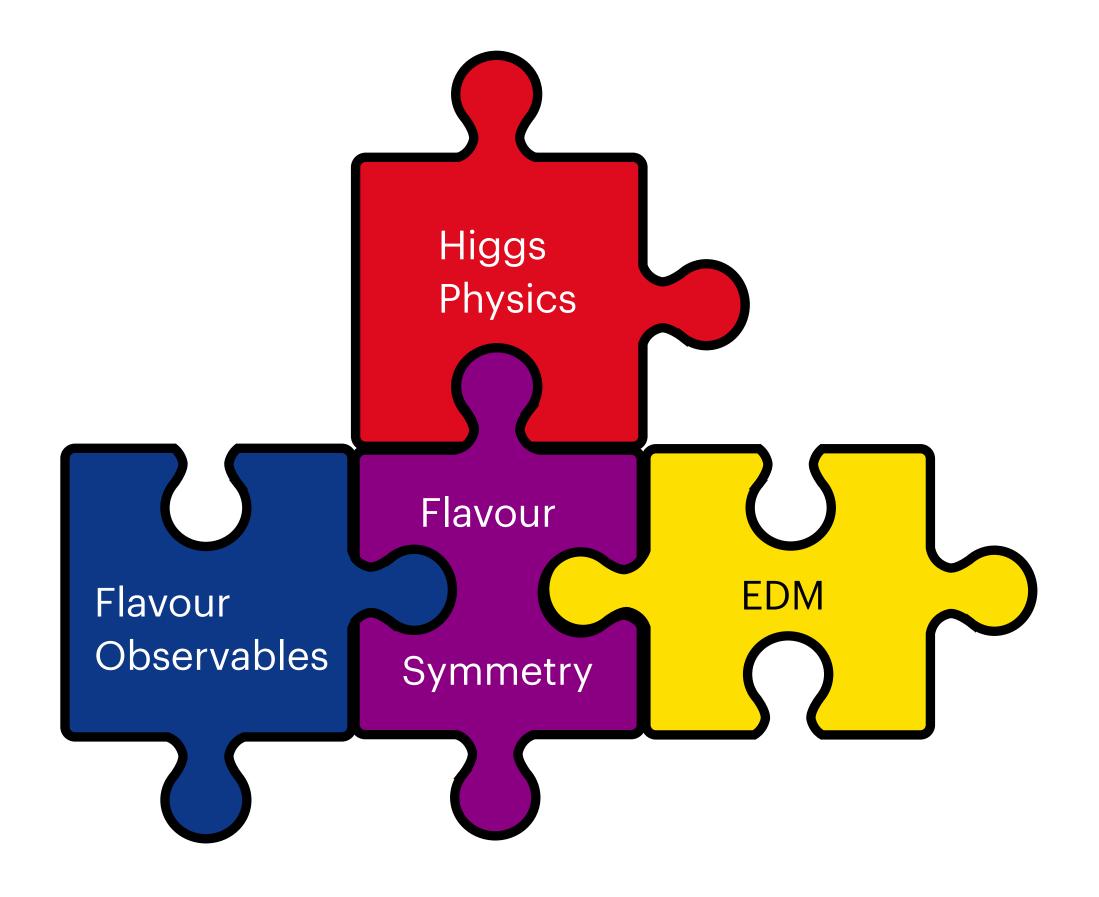
A. Theoretical Framework

B. Flavour-Symmetries

- 1. Minimal Flavour Violation
- 2. Froggatt-Nielsen

C. Searches & Consequences

- 1. Higgs Physics
- 2. EDM
- 3. Flavour Observables
- D. Synergy & Prospects



Theoretical Framework

• Assuming NP can be described within the **SMEFT**, there is just one type of **dimension-6 operator** contributing to the modification of the Yukawa interactions:

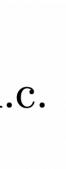
EV $\mathcal{L} \subset -\bar{F}'_L Y'_f \overset{(\sim)}{H} f'_R - \bar{F}'_L C'_f \overset{(\sim)}{H} f'_R \frac{H^{\dagger} H}{\Lambda_f^2} + \text{h.c.}$ + Mass

Notice that in principle the **NP scale** Λ can be **different** in the quark and lepton sector!

NSB

$$\mathcal{L} \subset -\bar{f}_L Y_f f_R \frac{v}{\sqrt{2}} - \bar{f}_L \left(Y_f + \frac{v^2}{\Lambda_f^2} C_f\right) f_R \frac{h}{\sqrt{2}} + h$$

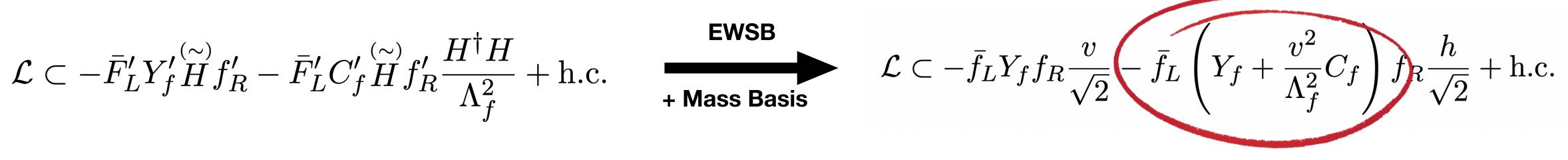
s Basis



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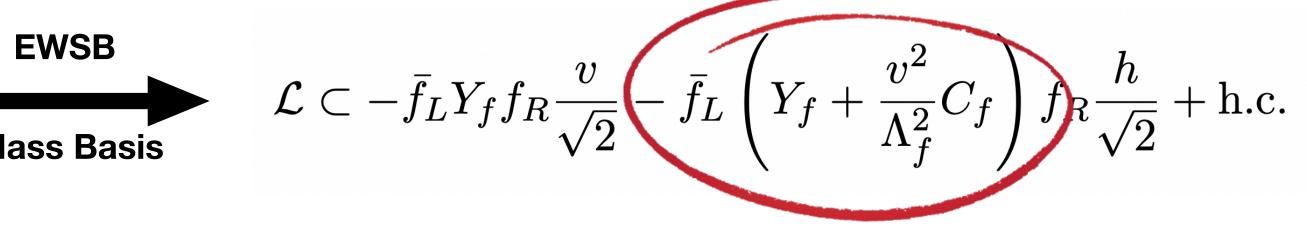
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• Parametrize deviations:

$$\begin{array}{ll} \mbox{Effective}\\ \mbox{Yukawa:} & \hat{Y}_f \equiv Y_f + \frac{v^2}{\Lambda^2} C_f \\ \mbox{Phenomenological}\\ \mbox{Lagrangian:} & \mathcal{L}_{\rm eff.} = \frac{y_{ffh}^{\rm SM}}{\sqrt{2}} \left(\kappa_f \bar{f} f + \tilde{\kappa} \bar{f} i \gamma_5 f\right) h & \mbox{SM:} \ \kappa_f = 1 \ , \ \tilde{\kappa}_f = 0 \\ \mbox{Deviation}\\ \mbox{Parameter:} & r_f^2 \equiv \frac{|\hat{y}_{ffh}|^2}{|y_{ffh}^{\rm SM}|^2} = \frac{v^2 |\hat{y}_{ffh}|^2}{2m_f^2} = \kappa_f^2 + \tilde{\kappa}_f^2 \end{array}$$



M. E. Peskin (1312.4974) Handbook of LHC Higgs Cross Sections: 3. Higgs Properties (1307.1347)



A Flavour Model (1/2) **Minimal Flavour Violation**

The Model

- SM: accidental symmetry $U(3)^5 = U(3)_q^3 \times U(3)_l^2$ broken solely by Yukawas
- MFV: the only source of flavour and CP-violation in the SM comes from the Yukawas
- The Yukawas are promoted to **spurion fields** transforming as bi-triplets of the flavour symmetry
- \implies all higher dimensional flavour-violating operators must be **controlled by Yukawas**!

Consequences



G. D'Ambrosio, G.F. Giudice, G. Isidori, A. Strumia (0207036)

No flavour-violating terms! Only one c'_{f} for each fermion sector!





A Flavour Model (2/2)

Froggatt-Nielsen

The Model

- New U(1) symmetry and SM-singlet scalar field ϕ (conventionally, with charge $n_{\phi} = -1$)
- of ϕ/Λ_F

$$\mathcal{L} \subset -\left[y_{f,ij}'\bar{F}_{i,L}'\overset{(\sim)}{H}f_{j,R}' + c_{f,ij}'\bar{F}_{L}'\overset{(\sim)}{H}f_{R}'\frac{H^{\dagger}H}{\Lambda_{f}^{2}}\right]\left(\frac{\phi}{\Lambda_{F}}\right)^{n_{F_{i}}+n_{f_{j}}} + \text{h.c.}$$

Consequences

• Once the ϕ takes VEV, each term is suppressed by powers of $\epsilon \equiv \langle \phi \rangle / \Lambda_F$

$$Y_f = \text{diag}\left(y_{f_1} \epsilon^{n_{F_1} + n_{f_1}}, y_{f_2} \epsilon^{n_{F_2} + n_{f_2}}, y_{f_3} \epsilon^{n_{F_2} + n_{f_2}}, y_{f_3} \epsilon^{n_{F_2} + n_{f_2}}, y_{f_3} \epsilon^{n_{F_3} + n_{f_3}}\right)$$

• Fermions and ϕ transform under the new symmetry and the Yukawa terms are made invariant adding powers

 $C_{f,ij} \approx \mathcal{O}(1) \epsilon^{n_{F_i} + n_{f_j}} e^{i\theta_{f,ij}}$ $\varepsilon^{n_{F_3}+n_{f_3}})$



FN: Benchmarks

• Free parameters: singlet vev, the fermion charges and the phases

 \implies we can **trade** the **right-field charges** to the **masses** but we are still left with quite some **freedom**!

- The conclusions are fairly independent of the cha assignation, while **phases** can play a **mayor role**!
- We consider a specific scenario for quarks and thre possibilities for leptons

ATTENTION!

The **phases** are **not fixed** at all and can have major impact on the constraints!

Incomplete list of authors: G. Altarelli, F. Feruglio, N. Haba, I. Masina, D. Meloni, L. Merlo, H. Murayama, ...

rges

$$\epsilon \equiv \langle \phi \rangle / \Lambda_F = 0.23$$

 ee
 Quarks
 Q'_L
 u'_R
 d'_R

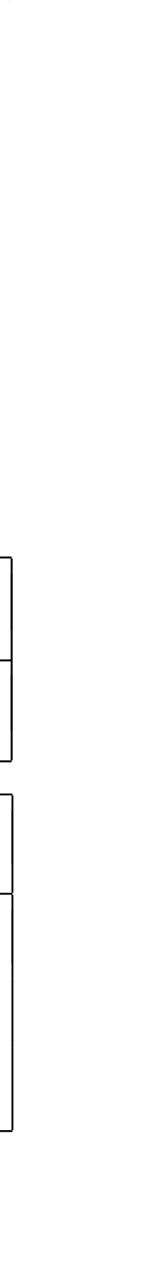
 (2,1,0)
 (5,2,0)
 (5,4,2)

 Leptons
 L'_L
 e'_R

 Anarchy (A)
 (0,0,0)
 (10,5,3)

 $\mu\tau$ -Anarchy (A)
 (1,0,0)
 (9,5,3)

 Hierarchy (H)
 (2,1,0)
 (8,4,3)



Flavour Symmetry at Work

• **Deviations** in each sector are **related**!

$$Y_f K_f = Y_f + rac{v^2}{\Lambda^2} \mathrm{diag}(\mathrm{Re}C_f)\,, \qquad Y_f \tilde{K}_f = rac{v^2}{\Lambda^2} \mathrm{diag}(\mathrm{Im}C_f)$$

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Minimal Flavour Violation

$$C_f' = c_f' Y_f'$$

$$\begin{aligned} \kappa_{f_1} &= \kappa_{f_2} = \kappa_{f_3} = 1 + \frac{v^2}{\Lambda_f^2} \operatorname{Re} c'_f \\ \tilde{\kappa}_{f_1} &= \tilde{\kappa}_{f_2} = \tilde{\kappa}_{f_3} = \frac{v^2}{\Lambda_f^2} \operatorname{Im} c'_f \\ r_f^2 &= 1 + \frac{v^4}{\Lambda_f^4} |c'_f|^2 + 2\frac{v^2}{\Lambda_f^2} \operatorname{Re} c'_f \end{aligned}$$

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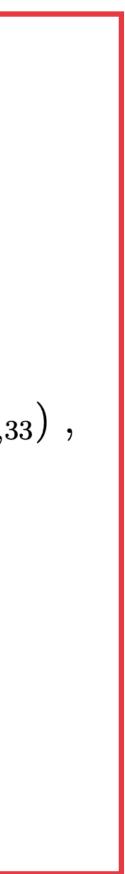
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Froggatt-Nielsen

$$C_{f,ij} \approx \mathcal{O}(1) \epsilon^{n_{F_i} + n_{f_j}} e^{i\theta_{f,ij}}$$

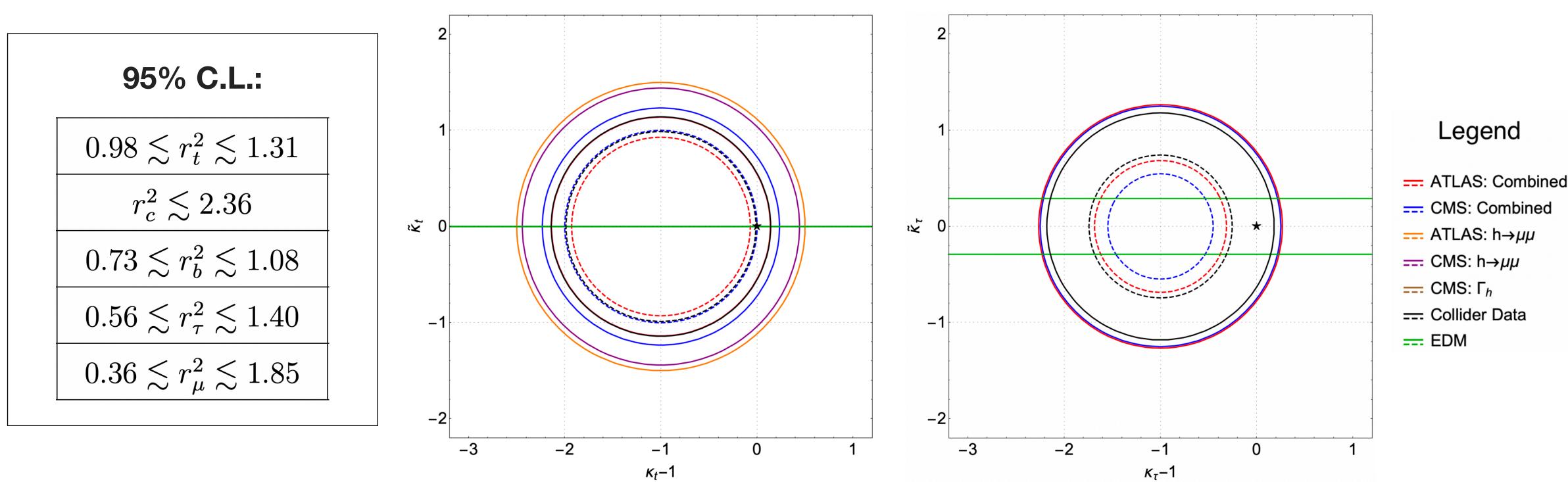
$$\begin{split} K_f = & 1 + \frac{v^2}{\Lambda^2} \operatorname{diag}\left(\mathcal{O}(1)\cos\theta_{f,11}, \, \mathcal{O}(1)\cos\theta_{f,22}, \, \mathcal{O}(1)\cos\theta_{f,33}\right) \,, \\ \tilde{K}_f = & \frac{v^2}{\Lambda^2} \operatorname{diag}\left(\mathcal{O}(1)\sin\theta_{f,11}, \, \mathcal{O}(1)\sin\theta_{f,22}, \, \mathcal{O}(1)\sin\theta_{f,33}\right) \,, \end{split}$$

$$r_f^2 \simeq 1 + \mathcal{O}(1)^2 \frac{v^4}{\Lambda^4} + 2\mathcal{O}(1) \frac{v^2}{\Lambda^2} \cos \theta_f$$



Higgs Physics (1/2)

• Global χ^2 fit:

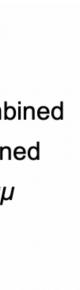


ATLAS (1909.02845) ATLAS (2007.07830) CMS (1809.10733) CMS (2009.04363)

 \mathcal{T}

Strong bounds only for third generations Relatively weak bounds for second generations **Null** bounds for the **first** generations!

Data from:

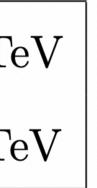


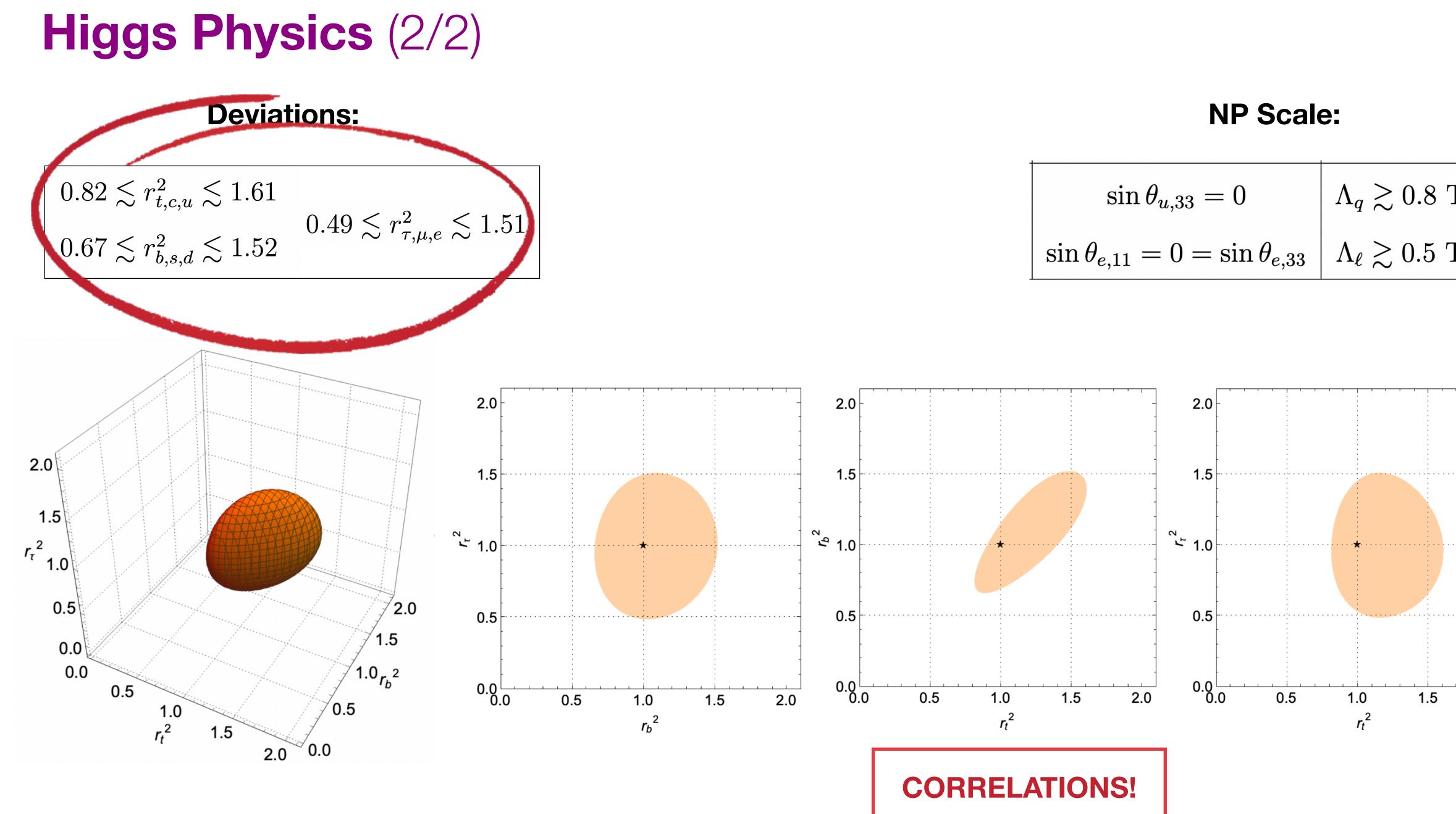


Deviations:

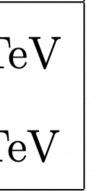
 $\begin{array}{l} 0.82 \lesssim r_{t,c,u}^2 \lesssim 1.61 \\ 0.67 \lesssim r_{b,s,d}^2 \lesssim 1.52 \end{array} \quad 0.49 \lesssim r_{\tau,\mu,e}^2 \lesssim 1.51 \end{array}$

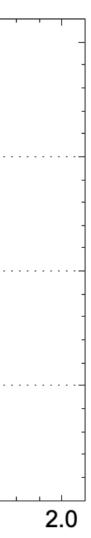
$\sin\theta_{u,33} = 0$	$\Lambda_q\gtrsim 0.8~{ m Te}$
$\sin \theta_{e,11} = 0 = \sin \theta_{e,33}$	$ig \Lambda_\ell \gtrsim 0.5~{ m Te}$

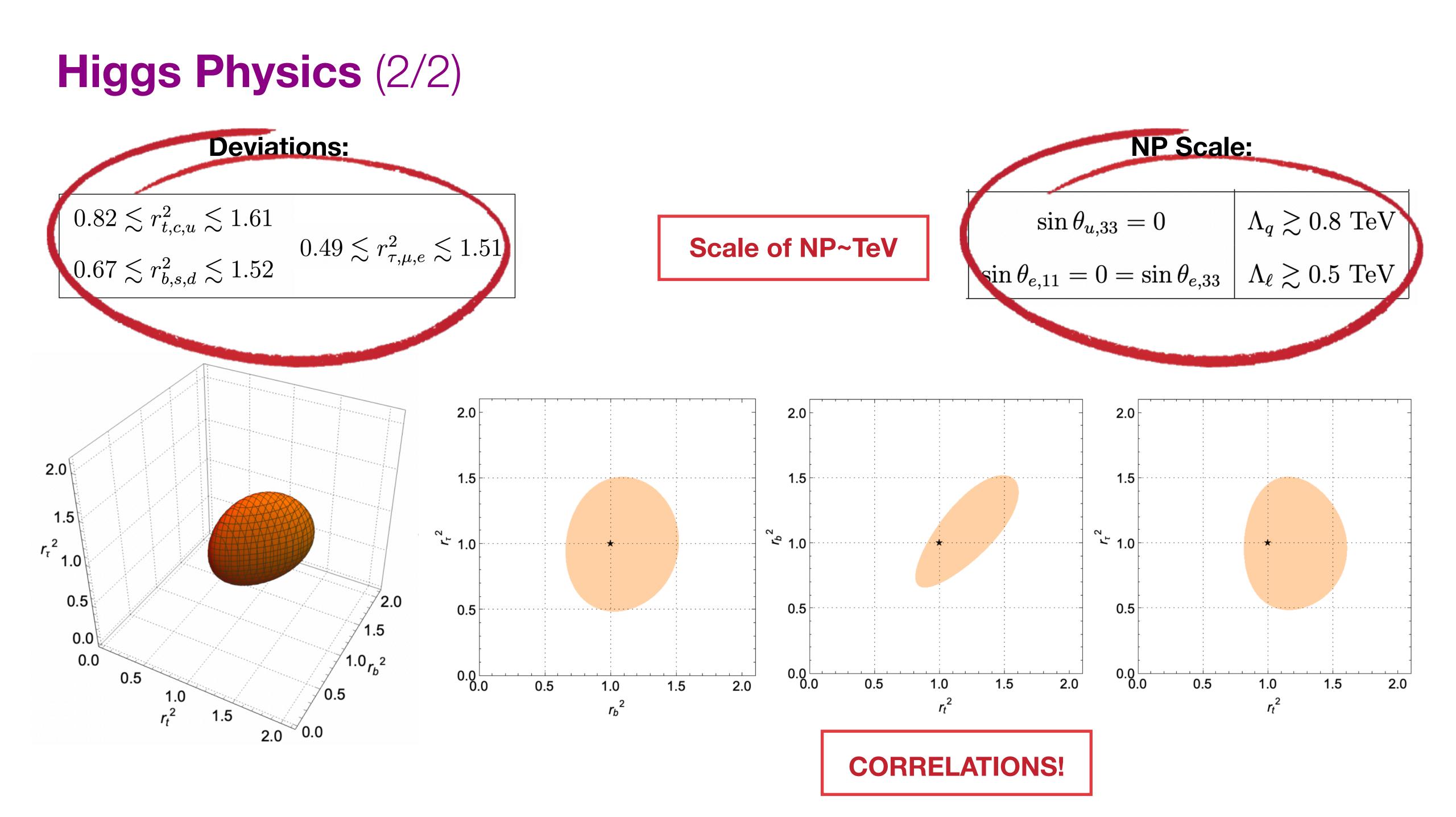




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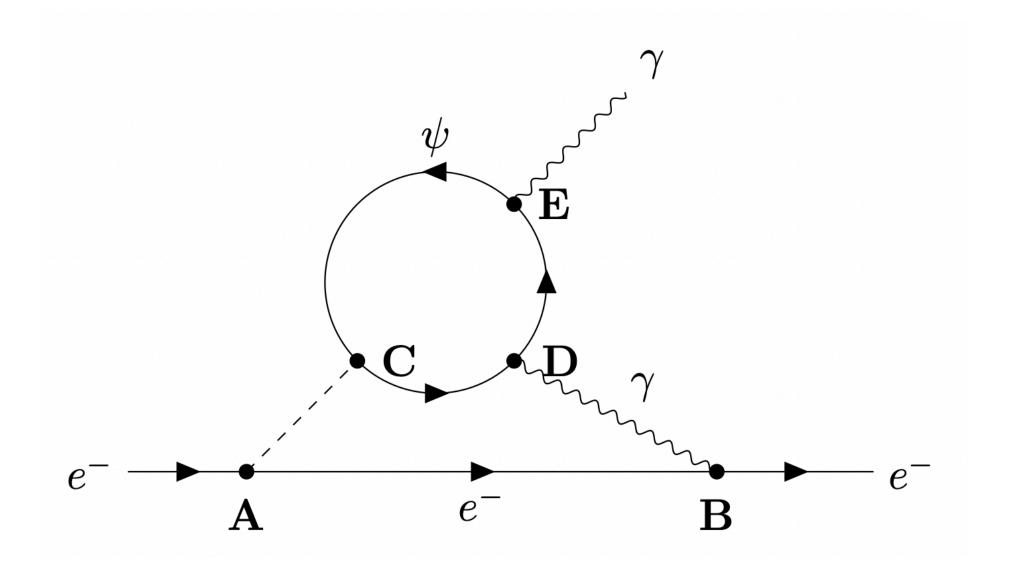
• CP-odd process

SM:

$$d_e^{CKM} \leq 10^{-44} e \text{ cm}$$
 Experimentally:
 $|d_e| < 1.1 \times 10^{-29} e \text{ cm}$, at 90% C.L.

 M. Pospelov, and A. Ritz (1311.5537)

• FV-NP through **Bar-Zee** diagram:



$$\frac{d_e}{e} = 4 N_C Q_{\psi}^2 \frac{\alpha_{\rm em}}{(4\pi)^3} \sqrt{2} G_F m_e \left[\kappa_e \tilde{\kappa}_{\psi} f_1(x_{\psi/h}) + \tilde{\kappa}_e \kappa_{\psi} f_2(x_{\psi/h}) \right]$$
J. Brod, U. Haisch, and J. Zupan (13)

90% C.L.

$$\begin{split} |\tilde{\kappa}_t| &\lesssim 0.0012 \,, \qquad |\tilde{\kappa}_b| \lesssim 0.24 \,, \qquad |\tilde{\kappa}_\tau| \lesssim 0.29 \\ |\tilde{\kappa}_c| &\lesssim 0.37 \,, \qquad |\tilde{\kappa}_\mu| \lesssim 31 \qquad |\tilde{\kappa}_e| \lesssim 0.0017 \end{split}$$

J. Alonso-Gonzalez, L. Merlo, and S. Pokorski (2103.16569)







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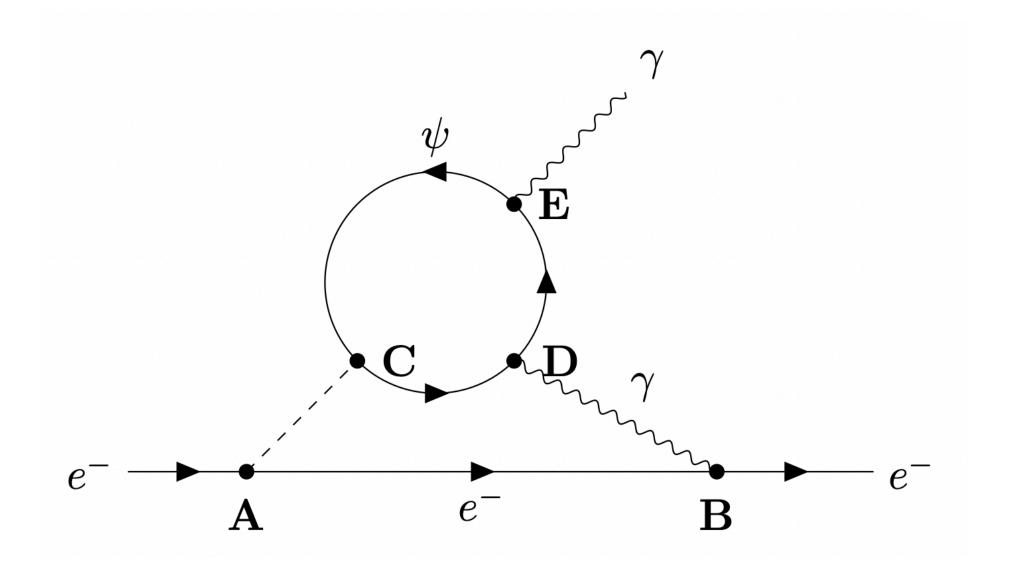
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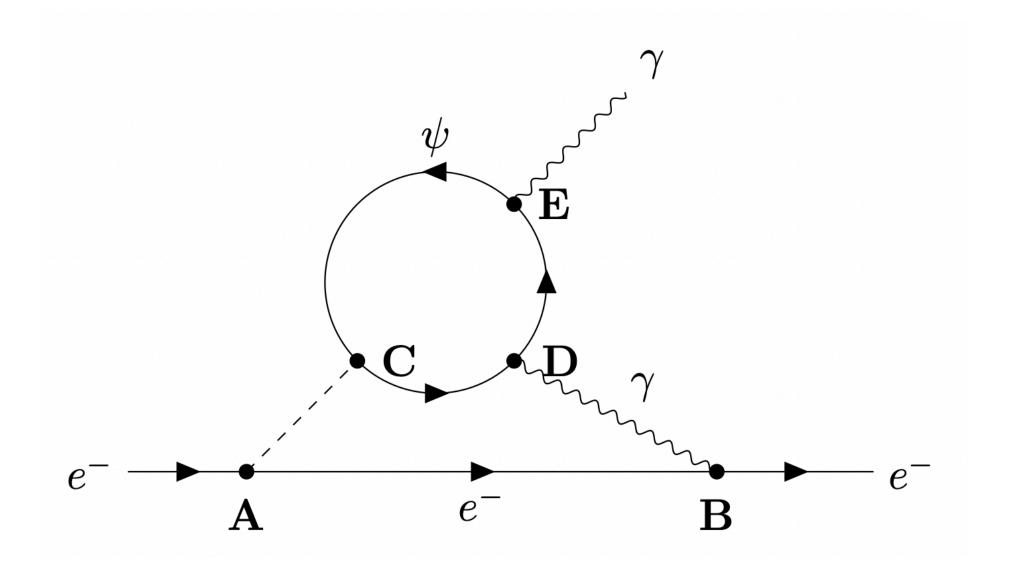
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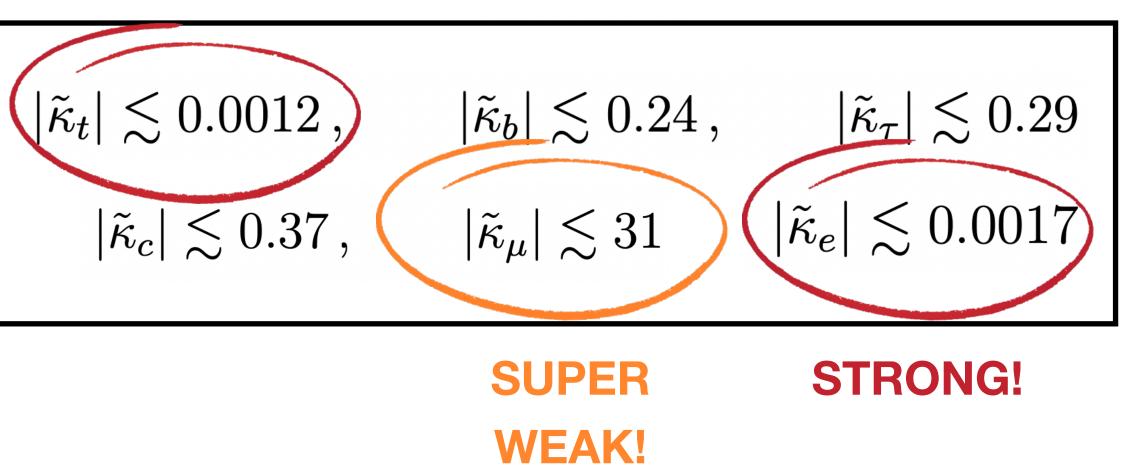
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• Stronger bounds on other generations!

$$|\tilde{\kappa}_{t,c,u}| \lesssim 0.0012$$
,

Deviations:

 $\begin{array}{l} 0.998 \lesssim r_q^2 \lesssim 1.002 \\ 0.997 \lesssim r_\ell^2 \lesssim 1.003 \end{array}$

J. Alonso-Gonzalez, L. Merlo, and S. Pokorski (2103.16569)

$$|\tilde{\kappa}_{b,s,d}| \lesssim 0.24$$
, $|\tilde{\kappa}_{\tau,e,\mu}| \lesssim 0.0017$

$$\begin{aligned} \sin \theta_{u,33} &= 1 & \Lambda_q \gtrsim 7.4 \text{ TeV} \\ \sin \theta_{e,11} &= 1 = \sin \theta_{e,33} & \Lambda_\ell \gtrsim 6.0 \text{ TeV} \end{aligned}$$



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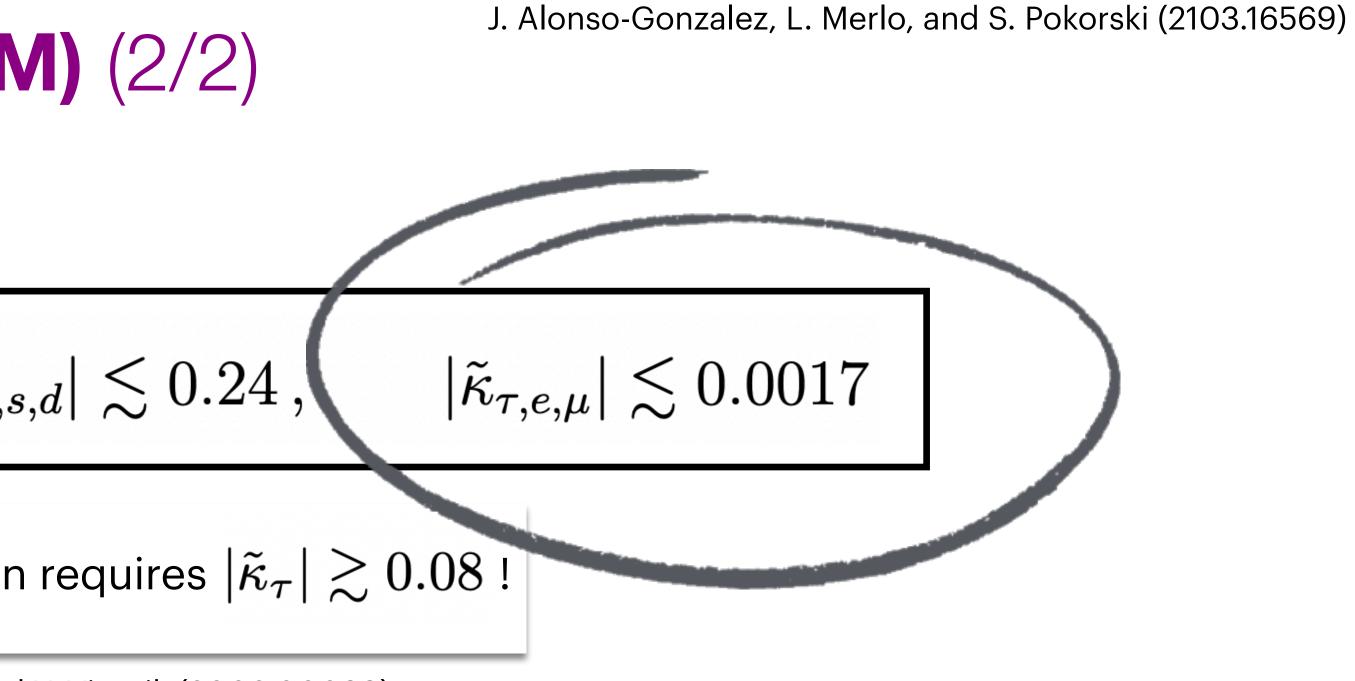
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Elina Fuchs, M. Losada, Y. Nir, and Y. Viernik (2003.00099)

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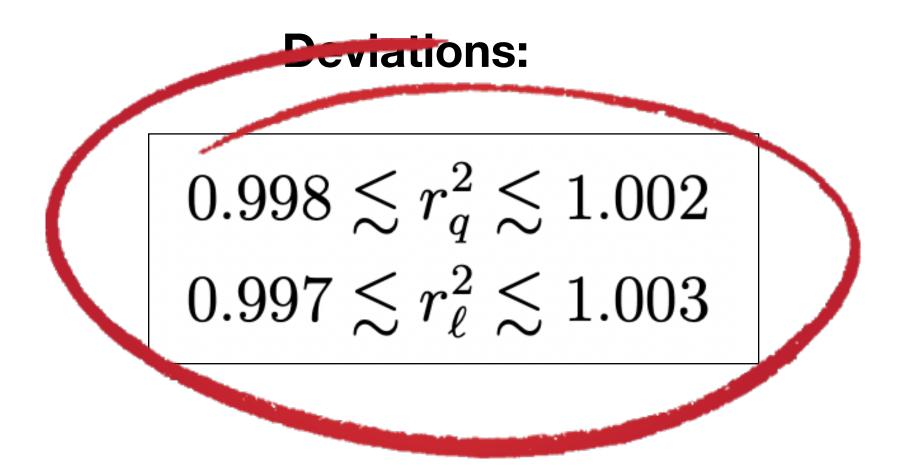
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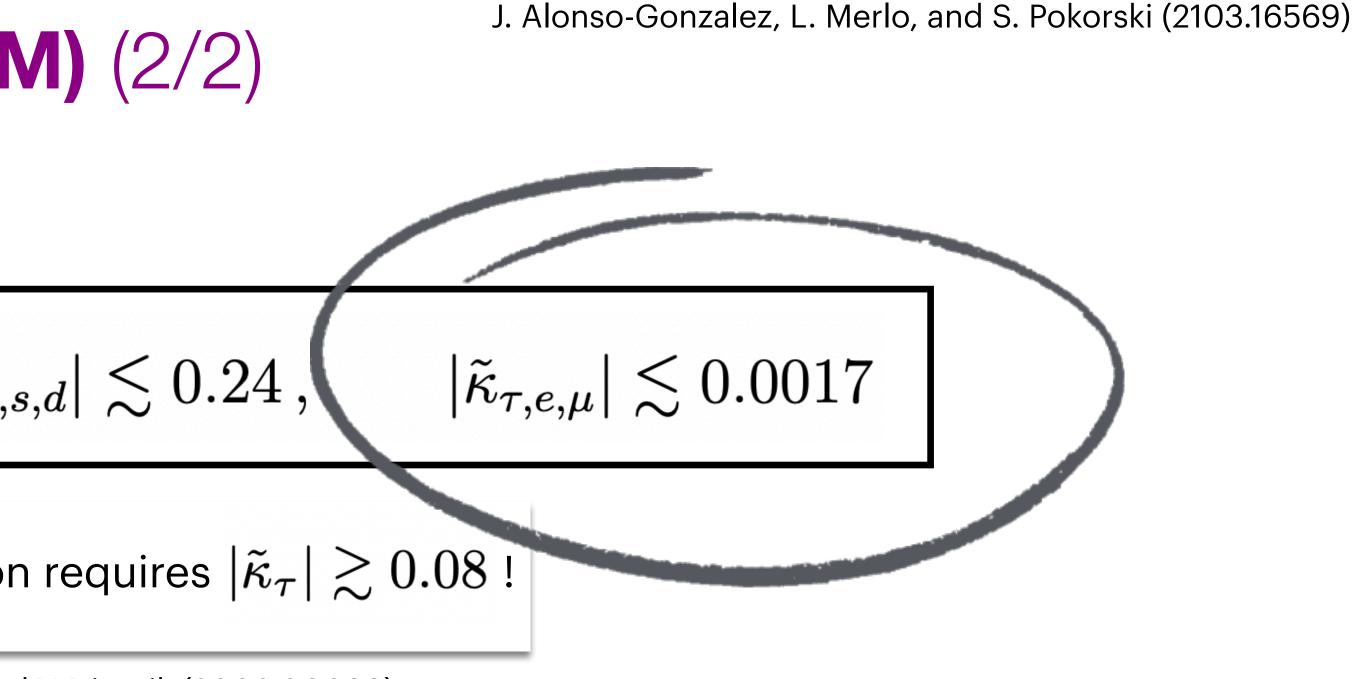
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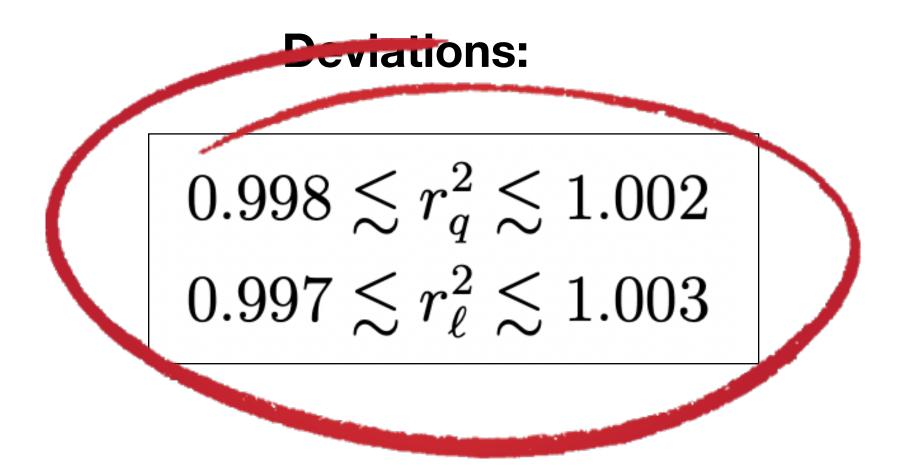
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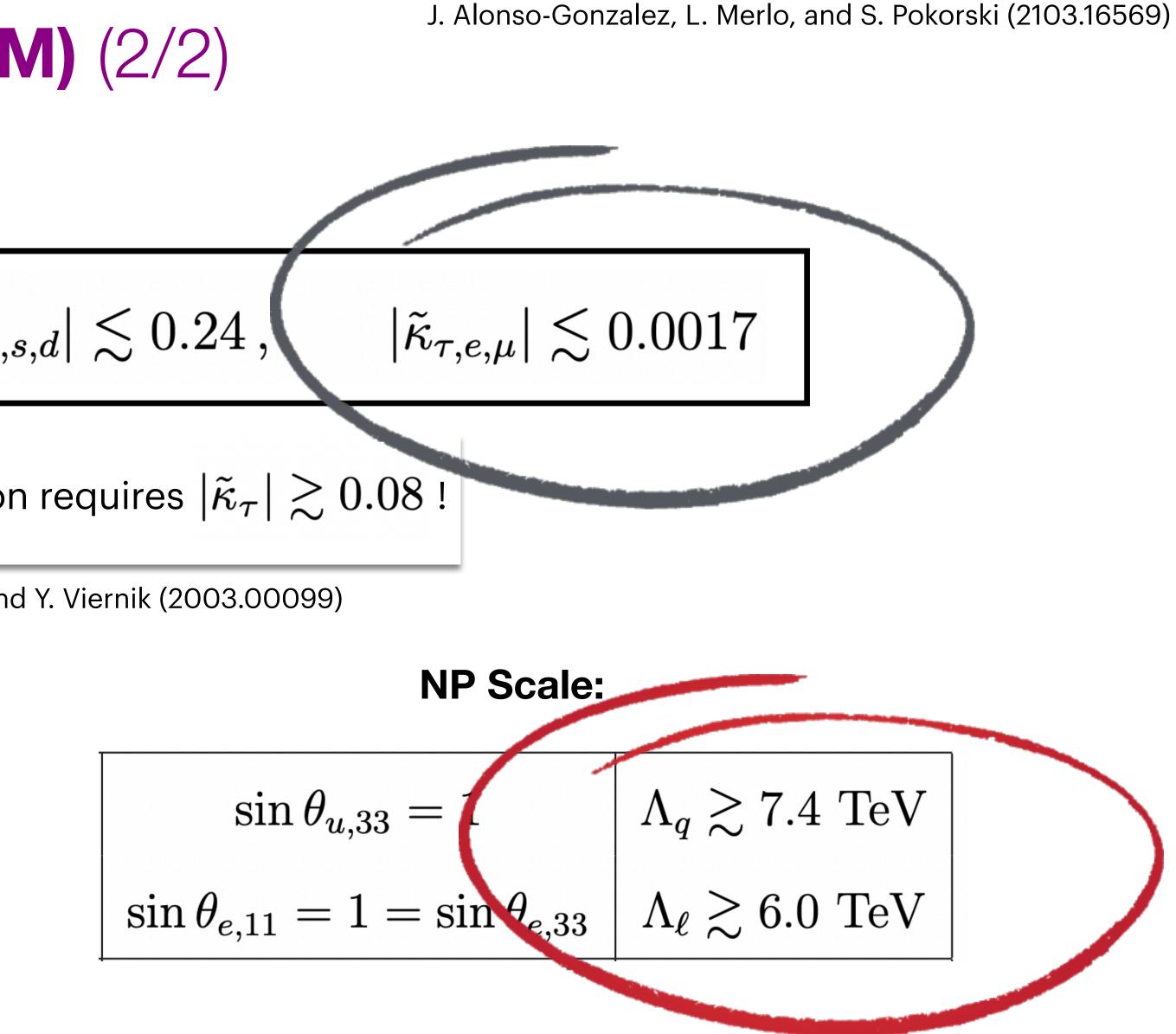
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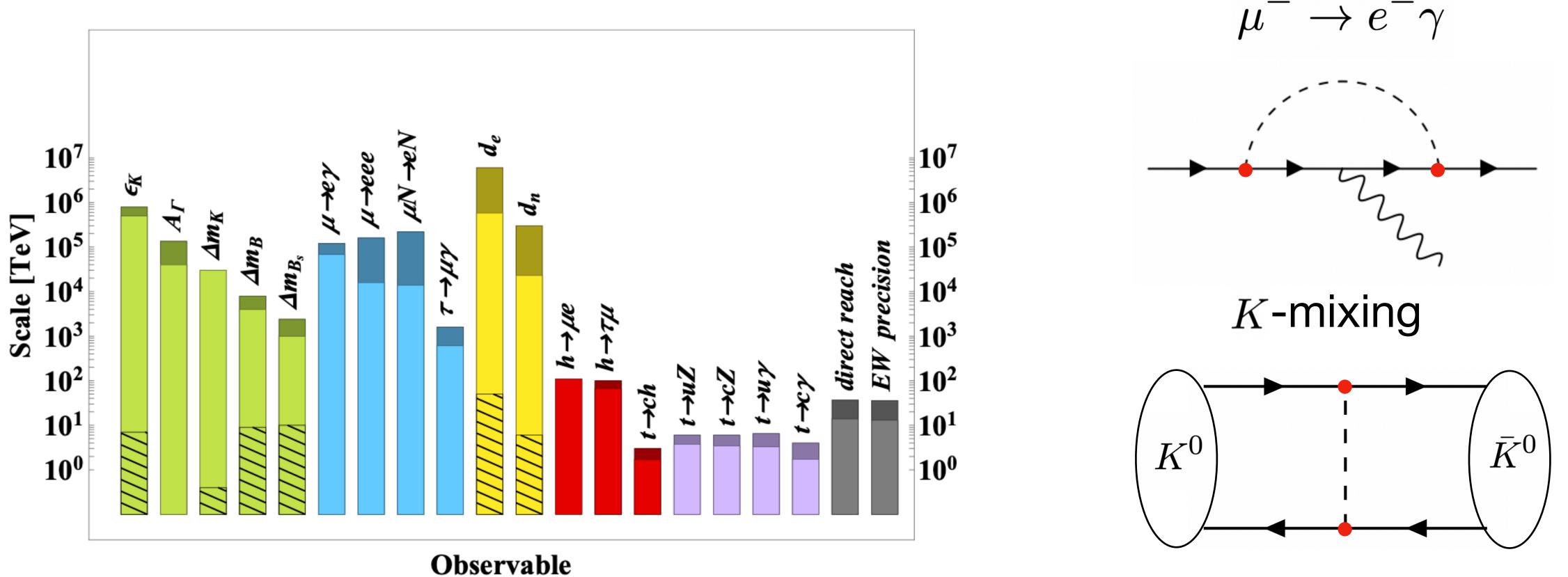






Flavour Observables (1/2)

Data from flavour-violating processes:



Input for the European Strategy for Particle Physics Update 2020 - 1910.11775

Higgs Flavour Violating couplings:

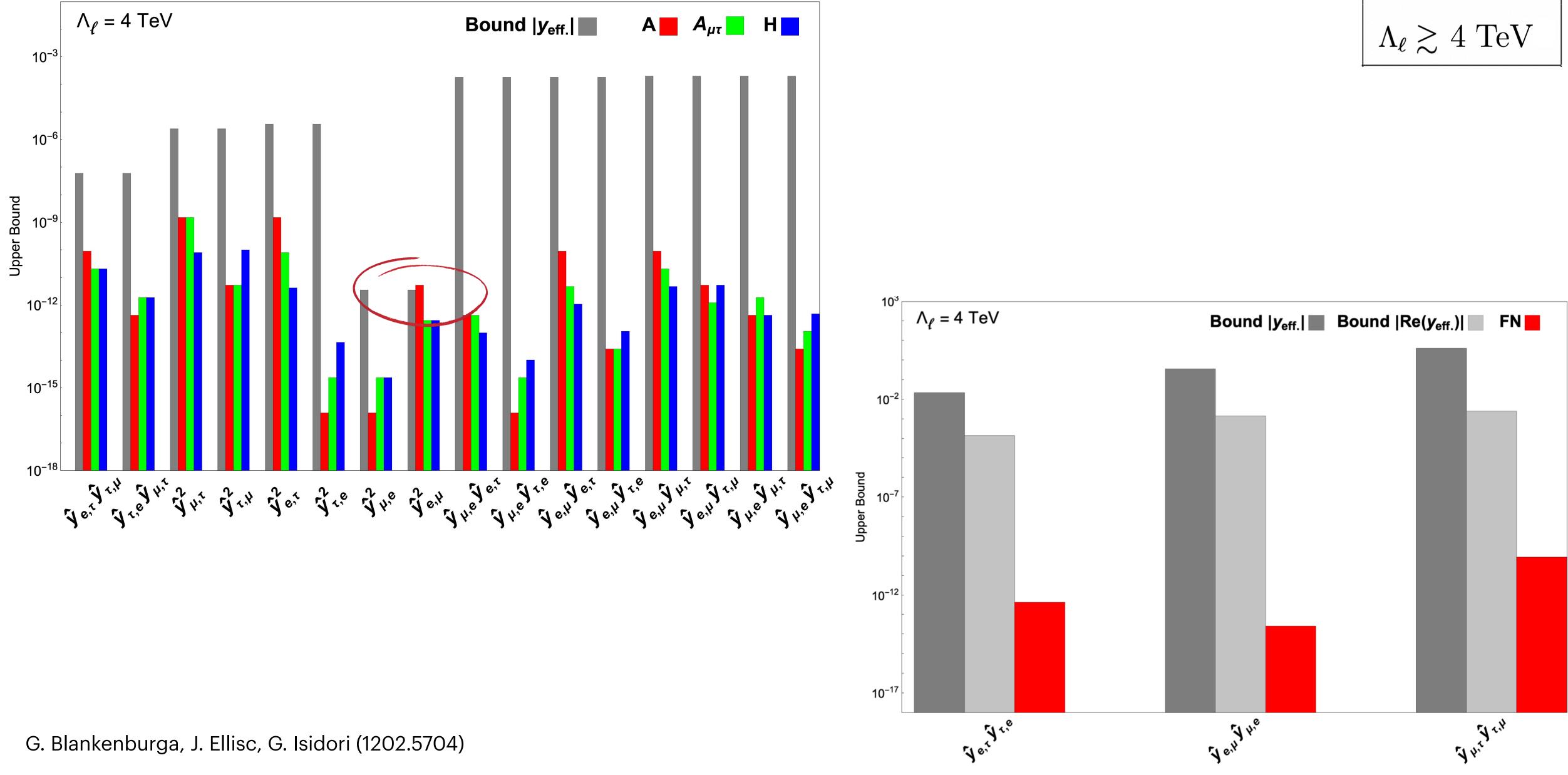
• without flavour symmetry: $|\hat{Y}|^2 \sim$

$$\frac{v^4}{\Lambda_q^4} \Longrightarrow \Lambda_q \approx (60 - 300)$$

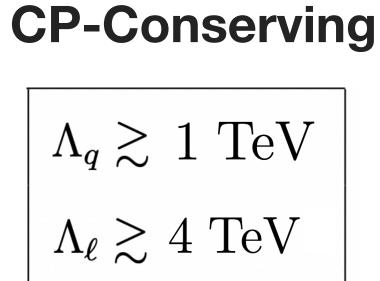
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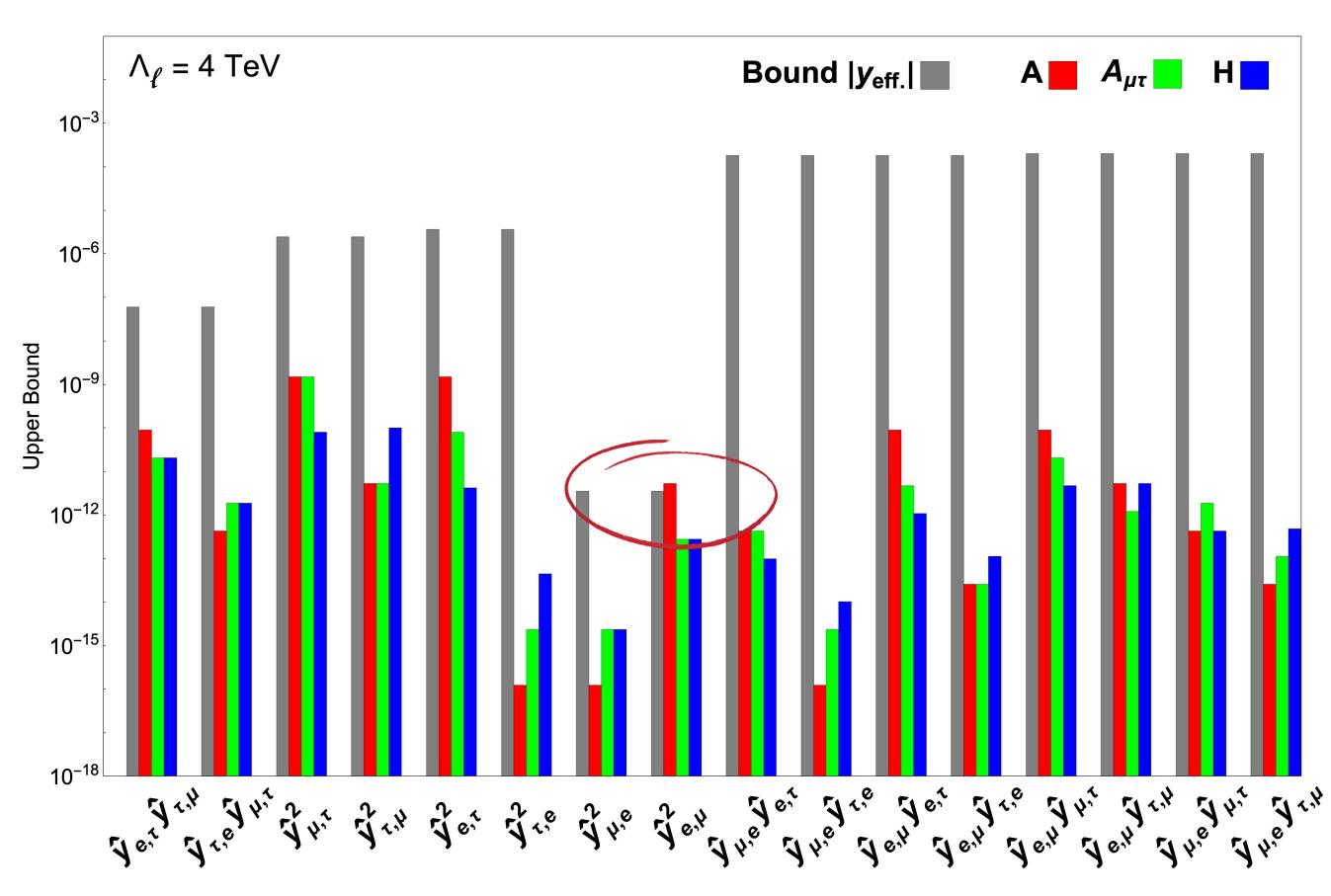
Flavour Observables (2/2)







Flavour Observables (2/2)

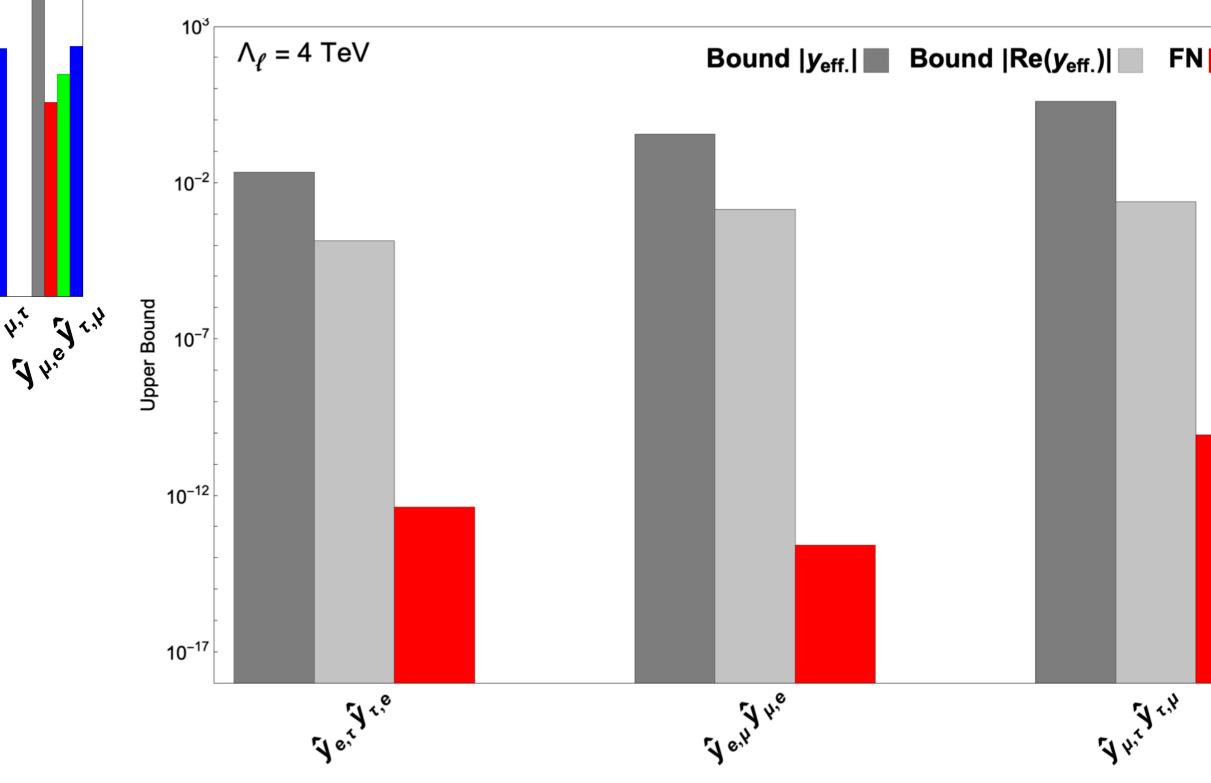


G. Blankenburga, J. Ellisc, G. Isidori (1202.5704)

CP-Conserving

 $\Lambda_q \gtrsim 1 \text{ TeV}$ $\Lambda_\ell \gtrsim 4 \text{ TeV}$

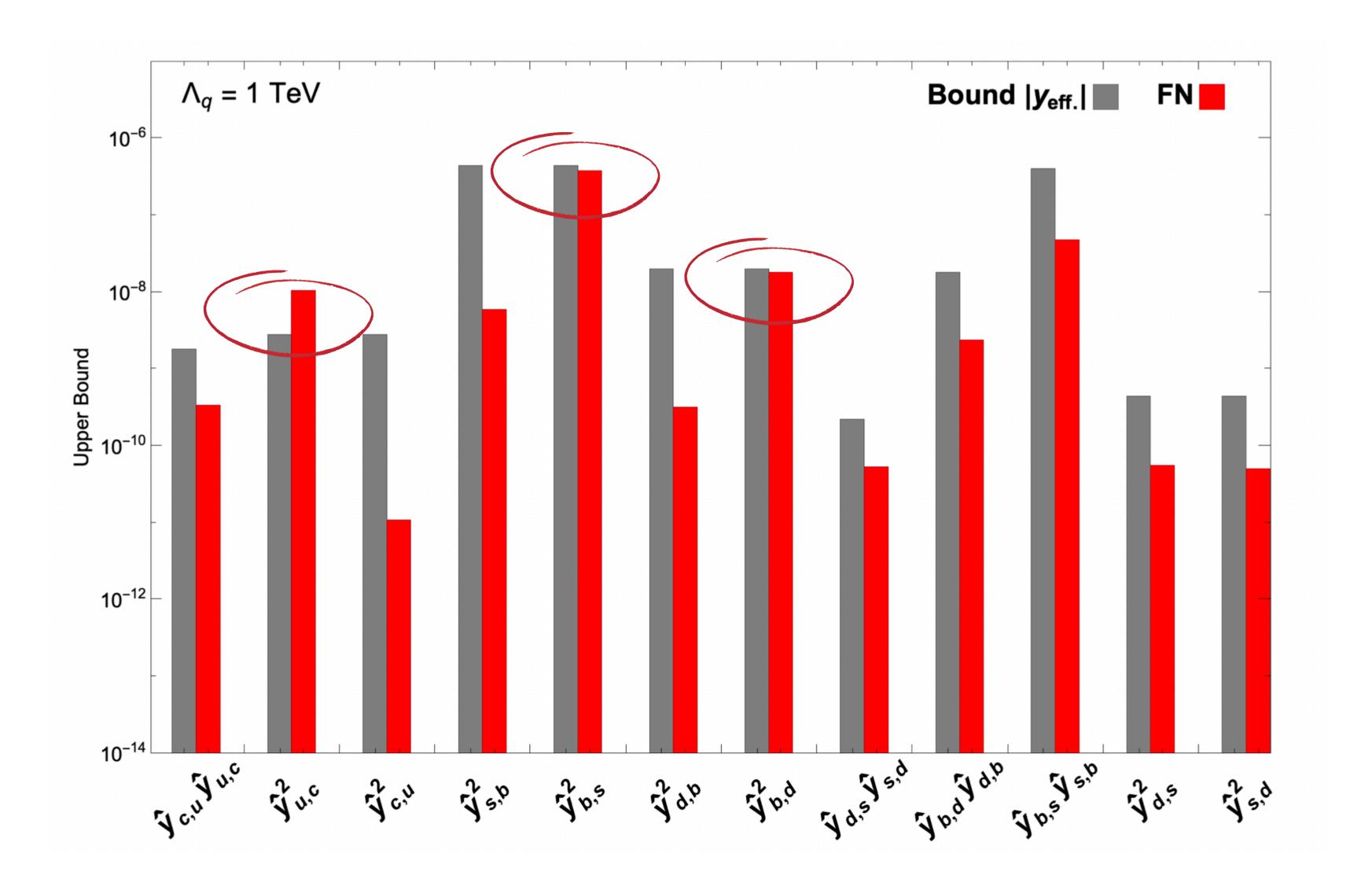
Several orders of magnitude stronger!







Flavour Observables (2/2)



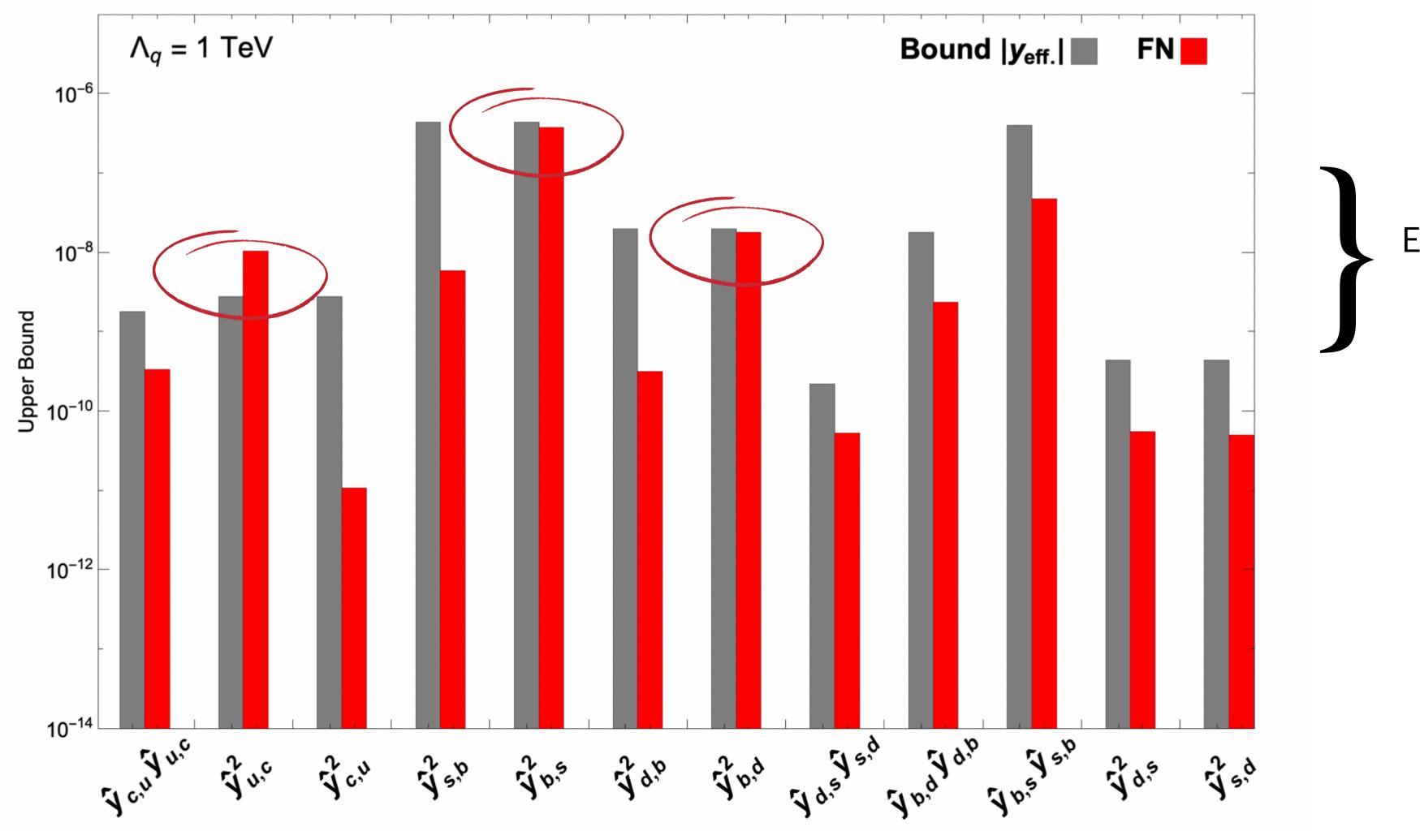
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Experimental searches not that far away!



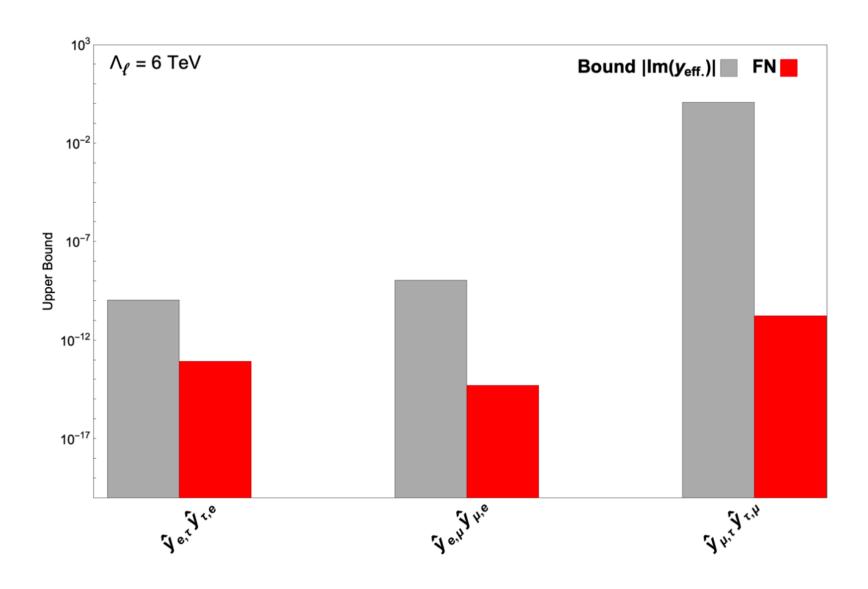
What can we learn?

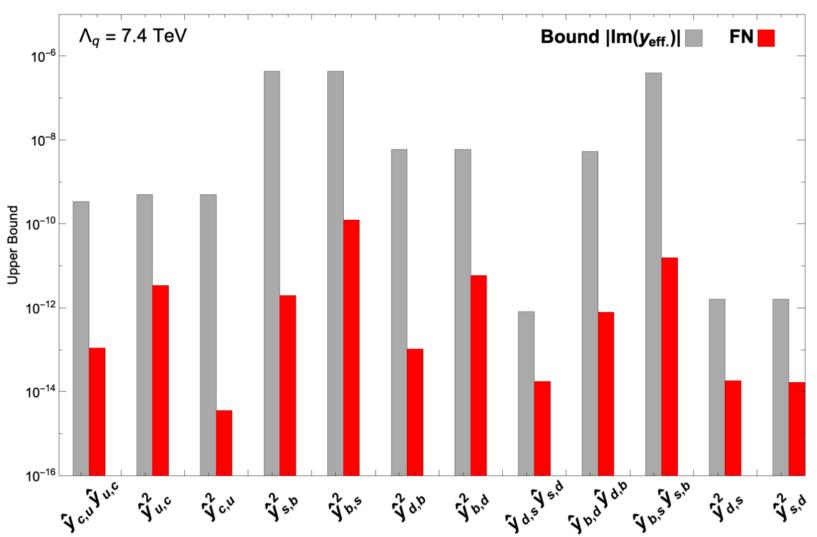
• Flavour Observables + EDM

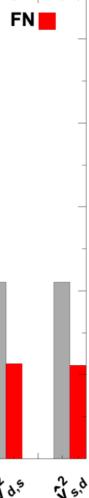
CP-Violating Flavour Observables EDM

 $\Lambda_q \gtrsim 3 \text{ TeV}$ $\Lambda_{\ell} \gtrsim 1 \text{ TeV}$

 $\Lambda_q \gtrsim 7.4 \text{ TeV}$ $\Lambda_\ell \gtrsim 6.0 \text{ TeV}$





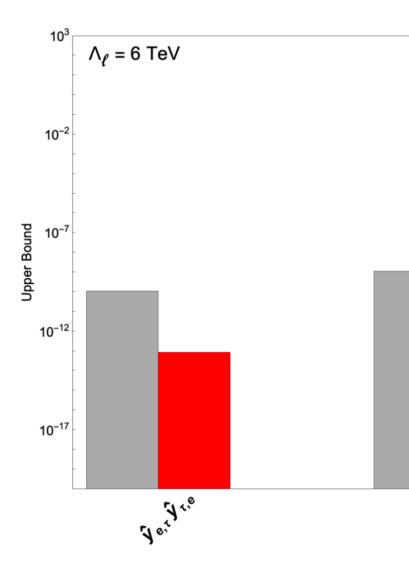


What can we learn?

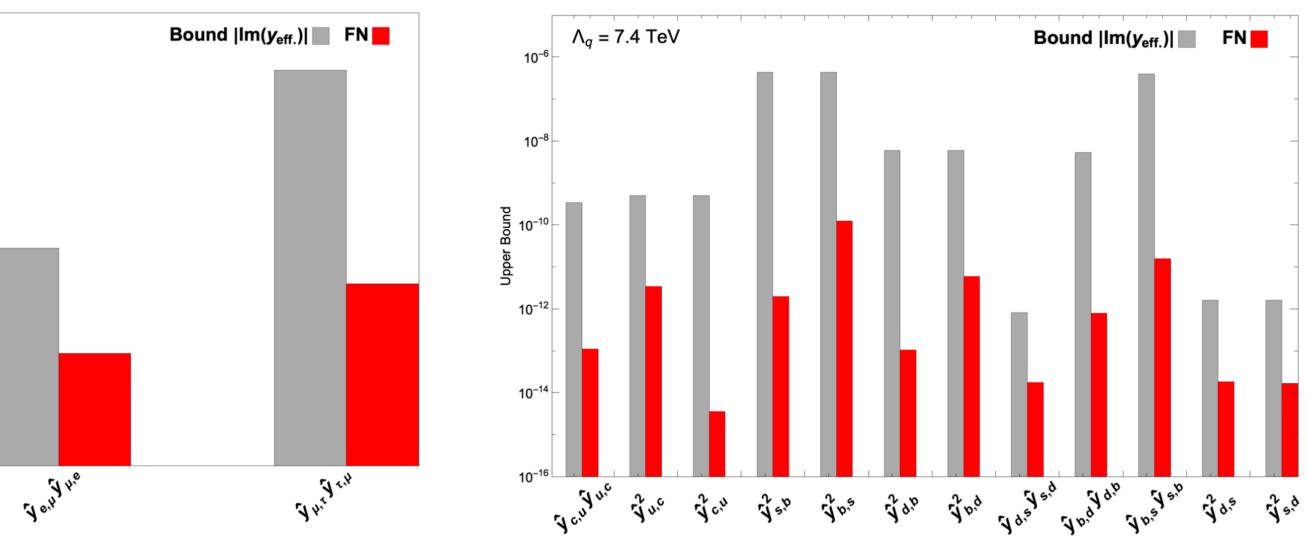
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CP-Violating Flavour Observables EDM

 $\Lambda_q \gtrsim 3 \text{ TeV}$ $\Lambda_\ell \gtrsim 1 \text{ TeV}$ $\Lambda_q \gtrsim 7.4 \text{ TeV}$ $\Lambda_\ell \gtrsim 6.0 \text{ TeV}$

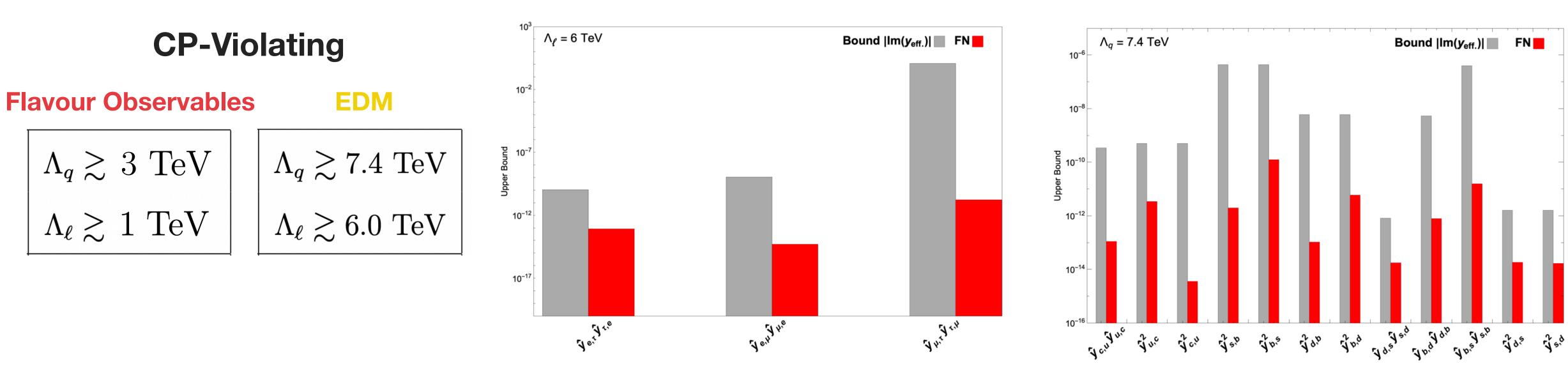


Avoiding all current experimental bounds!

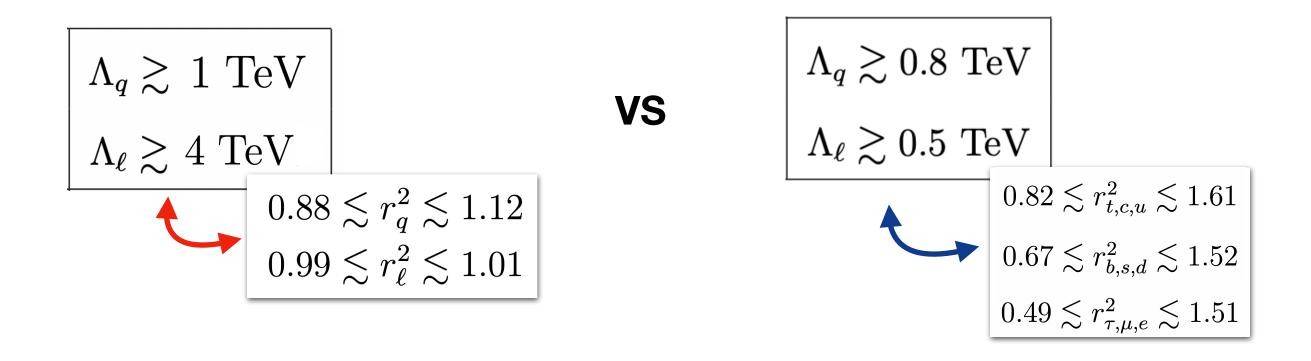


What can we learn?

 Flavour Observables EDM +



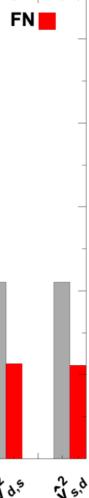
• Flavour Observables + **Higgs Physics**



Avoiding all current experimental bounds!

Bounds are of the **same order**! \implies

⇒ Future LHC searches could become more sensitive to NP than flavour observables!



Prediction Higgs & Top Decays

• We can make **predictions** on processes that have not been measured yet!



 \implies **Top-decay** into higgs+quark!

BR	Experimental Bound 95%C.L.	FN Prediction				
$h \rightarrow uc$		6×10^{-8}				
$h \rightarrow ds$			6×10^{-10}			
$h \rightarrow db$	_		4×10^{-8}			
$h \rightarrow sb$	_		8×10^{-7}			
BR	Experimental Bound 95%C.L.	A	$A_{\mu au}$	H		
$h \rightarrow e \mu$	6.1×10^{-5}	3×10^{-9}	10^{-10}	1×10^{-1}		
$h \rightarrow e \tau$	2.2×10^{-3}	8×10^{-7}	4×10^{-8}	2×10^{-9}		
$h ightarrow \mu au$	$1.5 imes 10^{-3}$	8×10^{-7}	8×10^{-7}	$9 imes 10^{-3}$		

BR	Experimental Bound 95%C.L.	FN Prediction
$t \rightarrow hu$	$4.5 imes 10^{-3}$	2×10^{-6}
$t \rightarrow hc$	4.6×10^{-3} , 5.6×10^{-3}	4×10^{-5}



Prediction Higgs & Top Decays

• We can make **predictions** on processes that have not been measured yet!



 \implies **Top-decay** into higgs+quark!

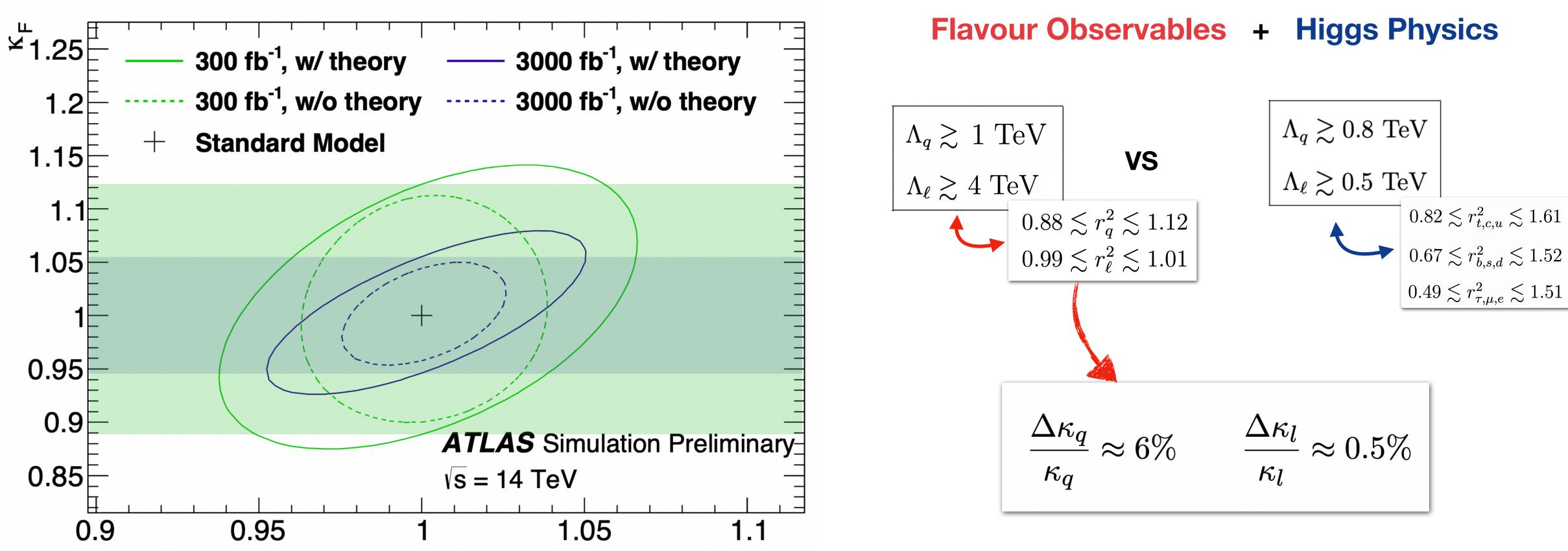
Experimental sensitivity still far away :(

BR	Experimental Bound 95%C.L.	FN Prediction				
$h \rightarrow uc$		6×10^{-8}				
$h \rightarrow ds$			6×10^{-10}			
$h \rightarrow db$	_		4×10^{-8}			
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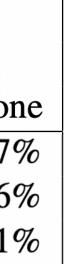
Waiting for HL-LHC - Run III



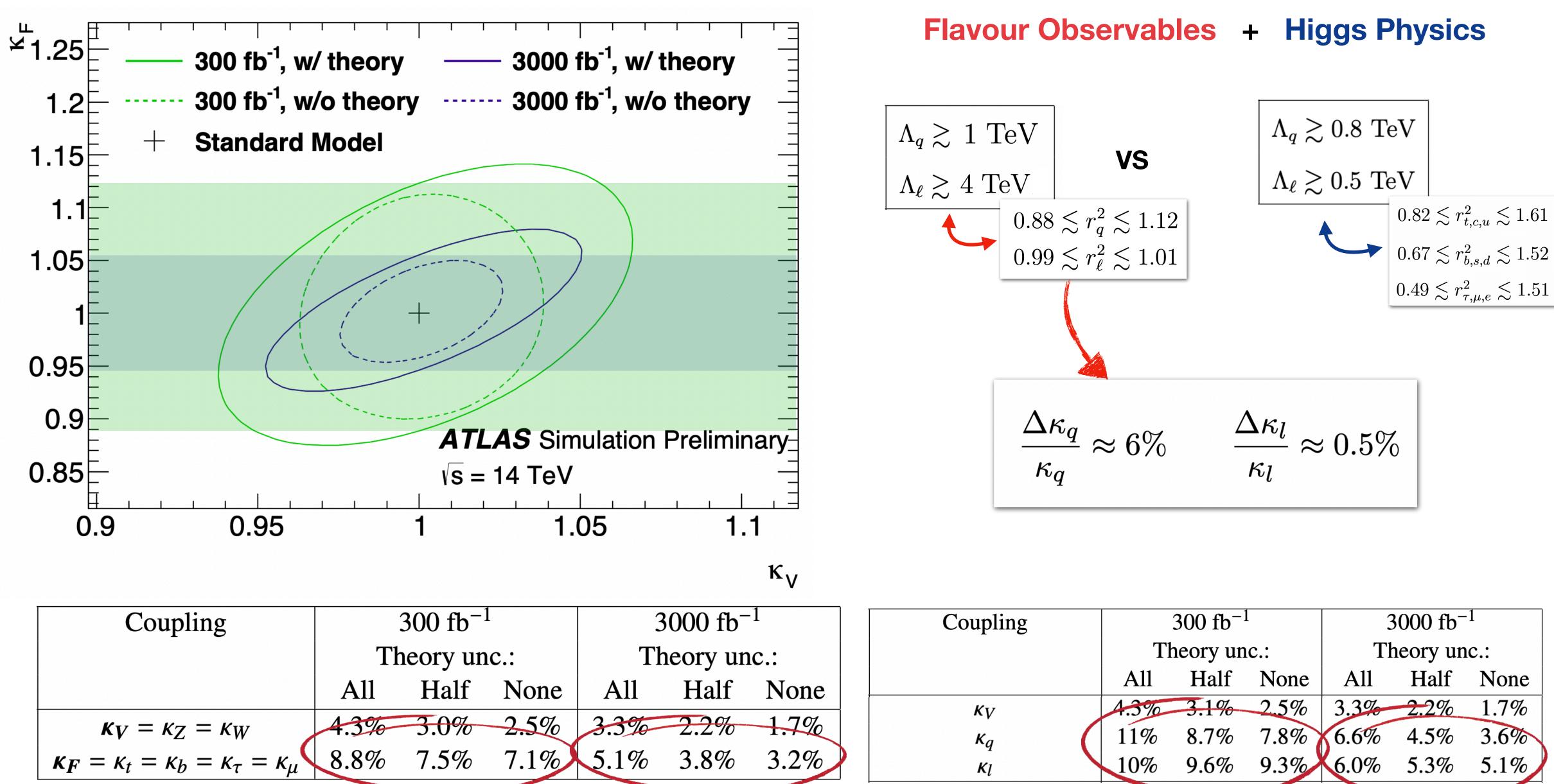
Coupling		300 fb ⁻	1		3000 fb ⁻	-1] [Coupling		300 fb ⁻	1		3000 fb ⁻	-1
		heory un	nc.:	T	neory ur	nc.:				heory ur	nc.:	Tł	heory un	nc.:
	All	Half	None	All	Half	None			All	Half	None	All	Half	No
							[ĸv	4.3%	3.1%	2.5%	3.3%	2.2%	1.7
$\kappa_V = \kappa_Z = \kappa_W$	4.3%	3.0%	2.5%	3.3%	2.2%	1.7%		Kq	11%	8.7%	7.8%		4.5%	
$\boldsymbol{\kappa_F} = \boldsymbol{\kappa}_t = \boldsymbol{\kappa}_b = \boldsymbol{\kappa}_\tau = \boldsymbol{\kappa}_\mu$	8.8%	7.5%	7.1%	5.1%	3.8%	3.2%		γ K _l	10%	9.6%	9.3%	6.0%	5.3%	

κ_v





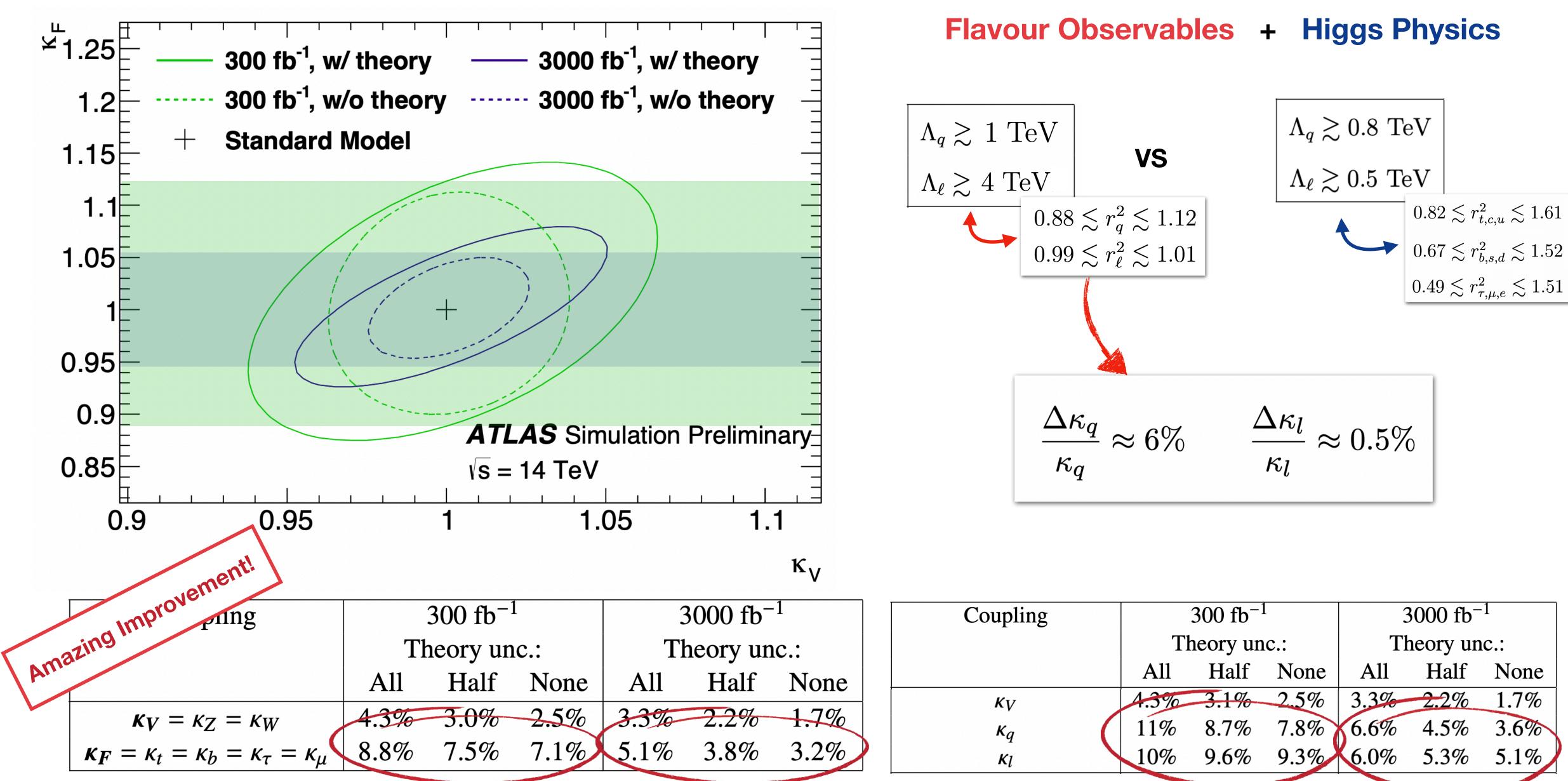
Waiting for HL-LHC - Run III



) ⁻¹	Coupling		300 fb^{-1}	[2	3000 fb ⁻	1
inc.:	1 0		neory un			neory un	
None		All	Half	None	All	Half	Nor
	κ _V	4.3%	3.1%	2.5%	3.3%	2.2%	1.79
1.7%	κ_q	11%	8.7%	7.8%	6.6%	4.5%	3.6
3.2%	κ _l	10%	9.6%	9.3%	6.0%	5.3%	5.19



Waiting for HL-LHC - Run III



) ⁻¹	Coupling		300 fb^{-1}	[2	3000 fb ⁻	1
inc.:	1 0		neory un			neory un	
None		All	Half	None	All	Half	Nor
	κ _V	4.3%	3.1%	2.5%	3.3%	2.2%	1.79
1.7%	κ_q	11%	8.7%	7.8%	6.6%	4.5%	3.6
3.2%	κ _l	10%	9.6%	9.3%	6.0%	5.3%	5.19



Take Home Messages

Flavour symmetries:

- are well motivated both to explore BSM physics and to understand the SM Yukawa structure!
- allow to connect and correlate very different experimental searches and constraints and can lower considerably the **NP scale** to **1-10 TeV**!!
- CP violating \rightarrow unless the phases of the diagonal couplings are zero, the EDM put the strongest bounds, followed by flavour observables.
- CP conserving \rightarrow most constrained by flavour observables, but Higgs physics' bounds are of the same

order of magnitude \rightarrow optimistic for the **quark** scale!

- Bounds from **Collider Searches** could become **competitive** with
 - **Flavour Observables**!





contact at: <u>arturo.degiorgi@uam.es</u>



SIGES

Higgs Physics - Parametrization

Data on signal strength from CMS and ATLAS:

$$\mu_P^F = \frac{\sigma_P}{\sigma_P^{\rm SM}} \left[\frac{\Gamma(h \to F)}{\Gamma^{\rm SM}(h \to F)} \left[\left(\frac{\Gamma_{h, \text{tot}}}{\Gamma_{h, \text{tot}}^{\rm SM}} \right)^{-1} \right] \right]$$

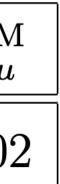
• Each ratio can be parametrised as a function of the **deviation parameters**:

$$\frac{\Gamma_{h,\text{tot.}}}{\Gamma_{h,\text{tot.}}^{\text{SM}}} = 1 + \text{BR}_{bb}^{\text{SM}} \left(r_b^2 - 1\right) + \left(\text{BR}_{gg}^{\text{SM}} + \text{BR}_{cc}^{\text{SM}}\right) \left(r_t^2 - 1\right) + \text{BR}_{\gamma\gamma}^{\text{SM}} \left(0.639 - 0.718 \, r_t^*\right) + \text{BR}_{\tau\tau}^{\text{SM}} \left(r_{\tau}^2 - 1\right)$$

h - Branching Fractions

BR_{bb}^{SM}	BR_{gg}^{SM}	$\mathrm{BR}^{\mathrm{SM}}_{ au au}$	BR_{cc}^{SM}	${ m BR}^{ m SM}_{\gamma\gamma}$	$\mathrm{BR}^{\mathrm{SN}}_{\mu\mu}$
0.58	0.082	0.062	0.029	0.002	0.000

*full expression with extra $+ 0.079 r_t^2$

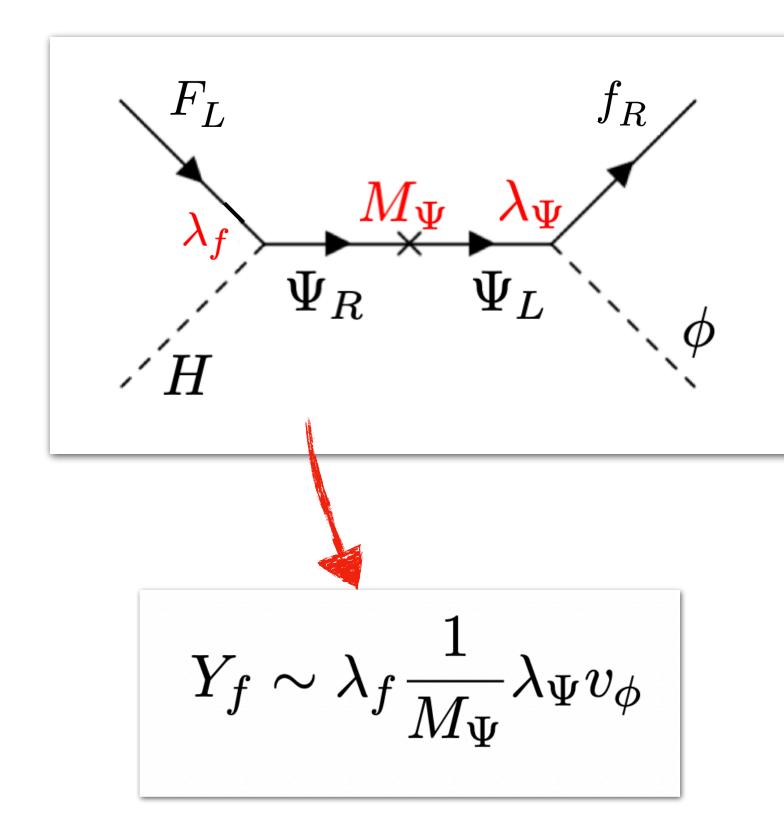


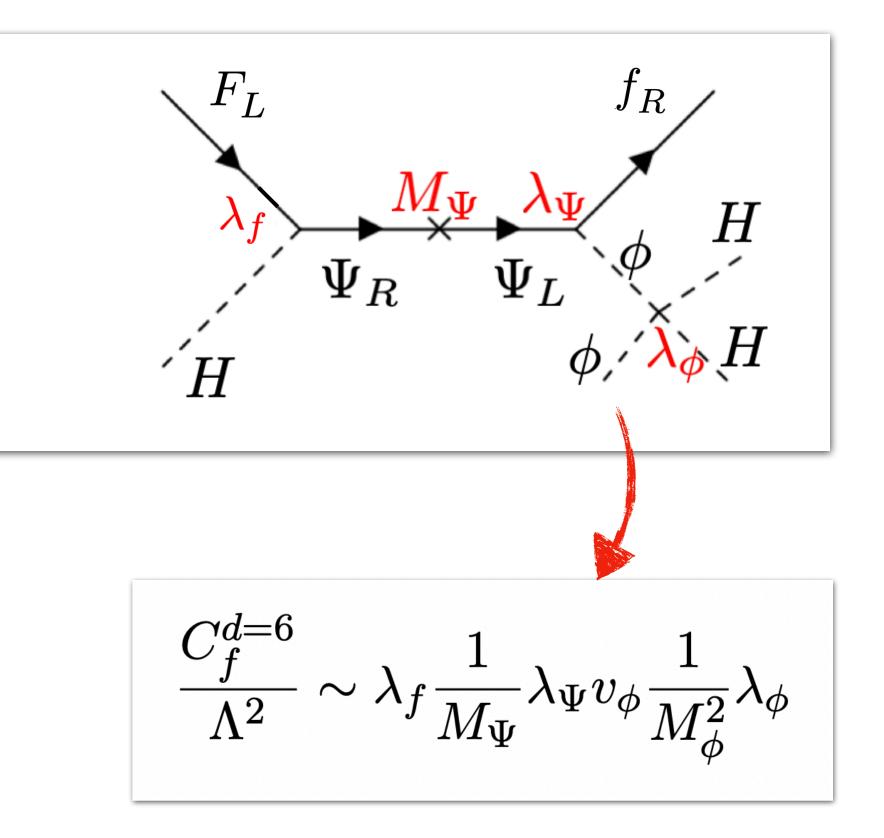




Toy-Model UV Completion

- **Vector-like** exotic fermions: $\Psi_{L,R}$ •
- MFV-inspired symmetry: $F_L, \Psi_{L,R} \sim 3 \subset SU(3)_L$ and $f_R \sim 3 \subset SU(3)_R$ •
- New bi-triplet scalar: $\phi \sim (3,\overline{3})$

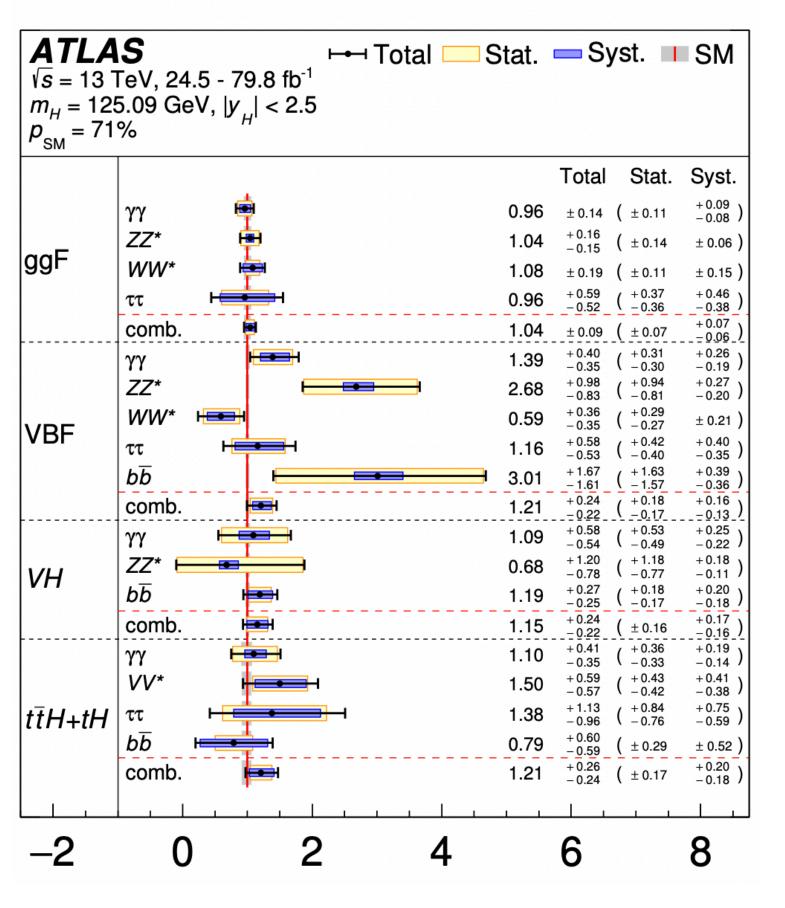


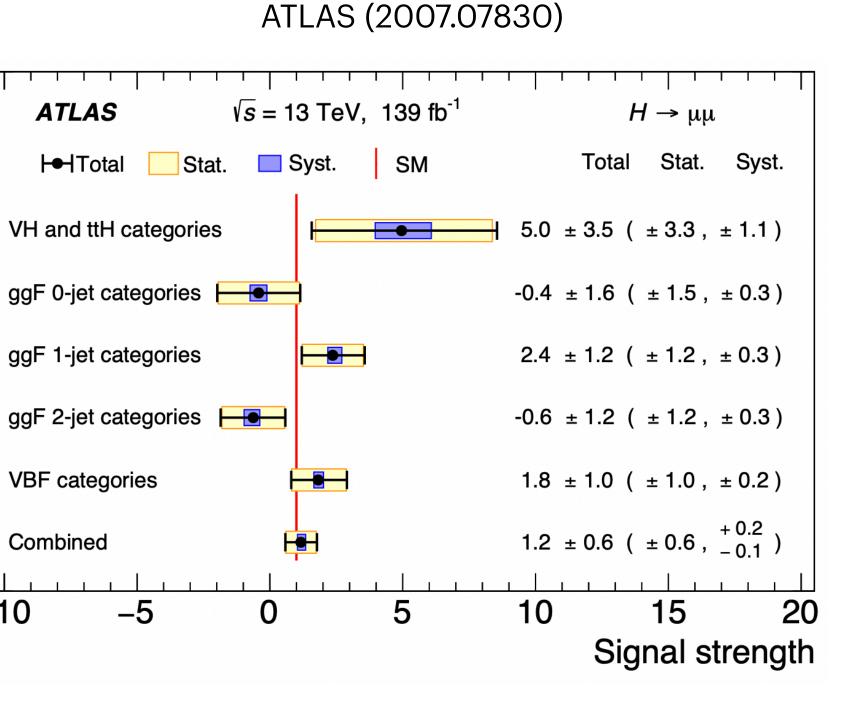


Waiting for HL-LHC - Run III: ATLAS

Combined

-10

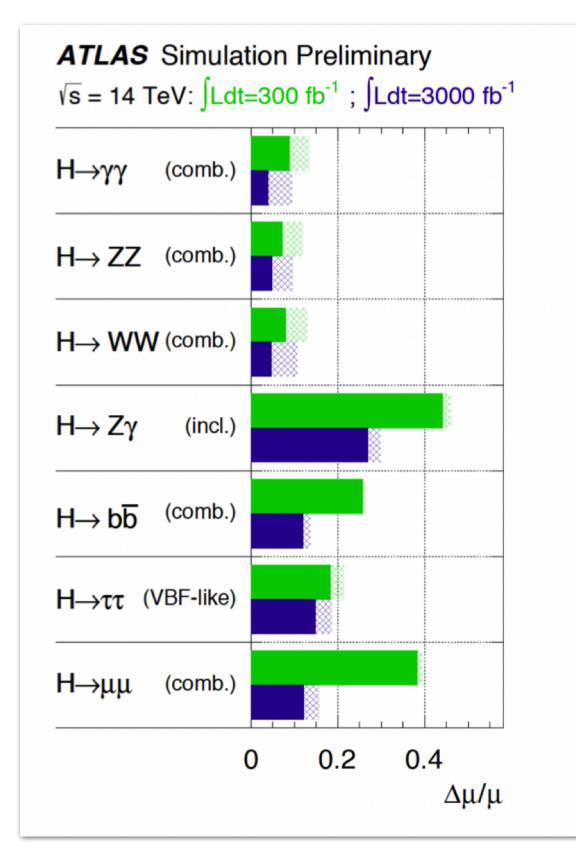




ATLAS (1909.02845)

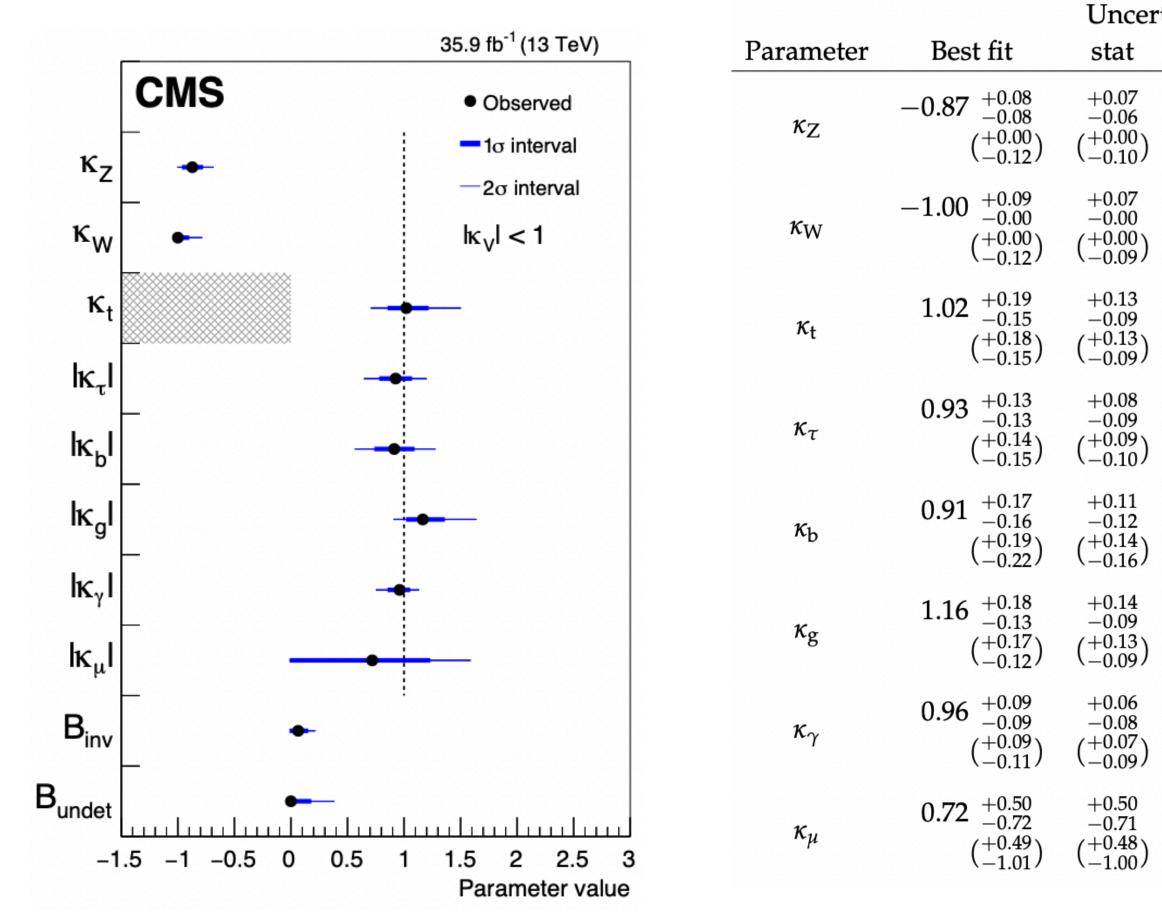
Prospects of LHC Higgs Physics at the end of Run III (1703.07689)

ATLAS Collaboration, ATL-PHYS-PUB-2014-016



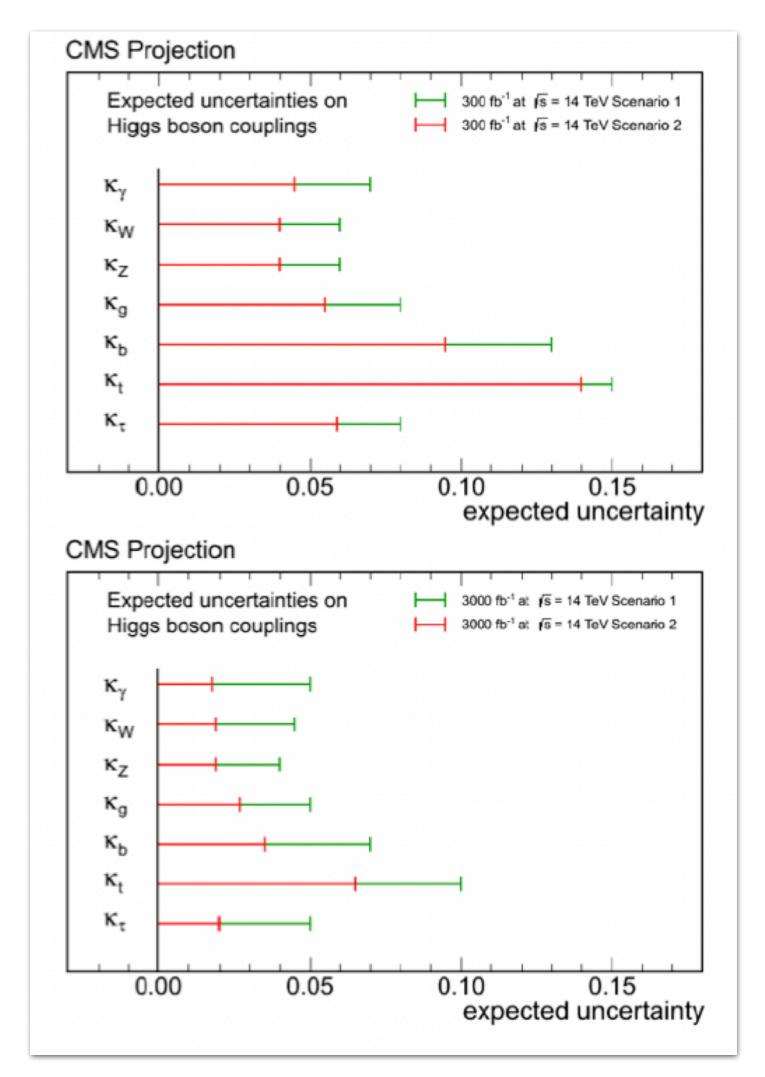


Waiting for HL-LHC - Run III: CMS



CMS (1809.10733)

Prospects of LHC Higgs Physics at the end of Run III (1703.07689)



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$^{+0.11}_{-0.10} \\ (^{+0.11}_{-0.11})$	
$^{+0.13}_{-0.11} \\ (^{+0.13}_{-0.15})$	
$^{+0.12}_{-0.10} \\ (^{+0.11}_{-0.09})$	
$^{+0.06}_{-0.06} \\ (^{+0.05}_{-0.07})$	
$^{+0.00}_{-0.07} \\ (^{+0.06}_{-0.08})$	