

# P-even, CP-violating Signals in Scalar Mediated Processes

Venus Keus

# DIAS

Institiúid Ard-Léinn | Dublin Institute for  
Bhaile Átha Cliath | Advanced Studies



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In collaboration with Howard Haber and Rui Santos

LHC Higgs WG2 and WG3 – CP violation and Higgs Sector  
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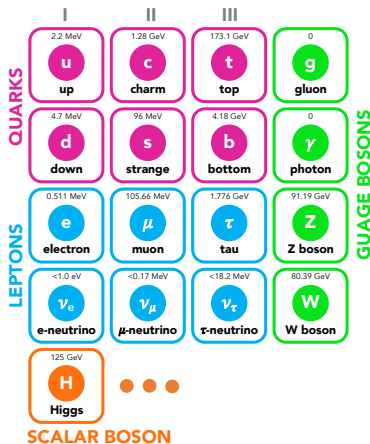
# The Standard Model of Particle Physics

## What is missing:

- a suitable Dark Matter candidate
- a successful baryogenesis mechanism
  - sufficient amount of CP violation
- a natural inflation framework
- an explanation for the fermion mass hierarchy
- a stable electroweak vacuum

⇒ beyond the Standard Model

⇒ scalar extensions of the SM



# The 2-Higgs doublet model (2HDM)

The most general renormalizable 2HDM scalar potential:

$$\begin{aligned} \mathcal{V} = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 - [m_{12}^2 \Phi_1^\dagger \Phi_2 + \text{h.c.}] \\ & + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) \\ & + \left\{ \frac{1}{2} \lambda_5 (\Phi_1^\dagger \Phi_2)^2 + [\lambda_6 (\Phi_1^\dagger \Phi_1) + \lambda_7 (\Phi_2^\dagger \Phi_2)] \Phi_1^\dagger \Phi_2 + \text{h.c.} \right\} \end{aligned}$$

With the fields and vevs:

$$\Phi_a = \begin{pmatrix} \Phi_a^+ \\ \Phi_a^0 \end{pmatrix} \quad \text{and} \quad \langle \Phi_a \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_a \end{pmatrix}$$

where

$$v_1 = v c_\beta, \quad v_2 = v s_\beta e^{i\xi}$$

with  $v = 246$  GeV.

# The Higgs basis

Linear combinations of  $\Phi_1$  and  $\Phi_2$

$$\mathcal{H}_1 \equiv c_\beta \Phi_1 + s_\beta \Phi_2, \quad \mathcal{H}_2 = e^{i\eta} [-s_\beta \Phi_1 + c_\beta \Phi_2]$$

such that

$$\langle \mathcal{H}_1^0 \rangle = v/\sqrt{2} \quad \text{and} \quad \langle \mathcal{H}_2^0 \rangle = 0$$

The scalar potential is now given by

$$\begin{aligned} \mathcal{V} = & Y_1 \mathcal{H}_1^\dagger \mathcal{H}_1 + Y_2 \mathcal{H}_2^\dagger \mathcal{H}_2 + [Y_3 e^{-i\eta} \mathcal{H}_1^\dagger \mathcal{H}_2 + \text{h.c.}] \\ & + \frac{1}{2} Z_1 (\mathcal{H}_1^\dagger \mathcal{H}_1)^2 + \frac{1}{2} Z_2 (\mathcal{H}_2^\dagger \mathcal{H}_2)^2 + Z_3 (\mathcal{H}_1^\dagger \mathcal{H}_1) (\mathcal{H}_2^\dagger \mathcal{H}_2) + Z_4 (\mathcal{H}_1^\dagger \mathcal{H}_2) (\mathcal{H}_2^\dagger \mathcal{H}_1) \\ & + \left\{ \frac{1}{2} Z_5 e^{-2i\eta} (\mathcal{H}_1^\dagger \mathcal{H}_2)^2 + [Z_6 e^{-i\eta} (\mathcal{H}_1^\dagger \mathcal{H}_1) + Z_7 e^{-i\eta} (\mathcal{H}_2^\dagger \mathcal{H}_2)] \mathcal{H}_1^\dagger \mathcal{H}_2 + \text{h.c.} \right\} \end{aligned}$$

# The physical mass eigenstates

The charged scalar mass eigenstates:  $G^\pm = \mathcal{H}_1^\pm$ ,  $H^\pm \equiv \mathcal{H}_2^\pm$

The neutral scalar mass eigenstates:  $G^0 = \sqrt{2} \operatorname{Im} \mathcal{H}_1^0$

and three states  $h_1$ ,  $h_2$ ,  $h_3$  with the neutral squared-mass matrix

$$\mathcal{M}^2 = v^2 \begin{pmatrix} Z_1 & \operatorname{Re}(\bar{Z}_6) & -\operatorname{Im}(\bar{Z}_6) \\ \operatorname{Re}(\bar{Z}_6) & \frac{1}{2} [Z_{34} + \operatorname{Re}(\bar{Z}_5)] + Y_2/v^2 & -\frac{1}{2} \operatorname{Im}(\bar{Z}_5) \\ -\operatorname{Im}(\bar{Z}_6) & -\frac{1}{2} \operatorname{Im}(\bar{Z}_5) & \frac{1}{2} [Z_{34} - \operatorname{Re}(\bar{Z}_5)] + Y_2/v^2 \end{pmatrix}$$

in the basis  $\{\sqrt{2} \operatorname{Re} \mathcal{H}_1^0 - v, \sqrt{2} \operatorname{Re} \mathcal{H}_2^0, \sqrt{2} \operatorname{Im} \mathcal{H}_2^0\}$  diagonalized by

$$R\mathcal{M}^2 R^T = \operatorname{diag}(m_1^2, m_2^2, m_3^2)$$

where  $R \equiv R_{12}R_{13}R_{23}$  is the product of three rotation matrices parametrized by  $\theta_{12}$ ,  $\theta_{13}$  and  $\theta_{23}$ , respectively, and

$$Z_{34} \equiv Z_3 + Z_4, \quad \bar{Z}_5 \equiv Z_5 e^{-2i\eta}, \quad \bar{Z}_6 \equiv Z_6 e^{-i\eta}, \quad \bar{Z}_7 \equiv Z_7 e^{-i\eta}$$

# The Higgs alignment limit

Providing  $Z_6 = 0$ , the mass-squared matrix becomes diagonal

$$\mathcal{M}_0^2 = \begin{pmatrix} Z_1 v^2 & 0 & 0 \\ 0 & Y_2 + \frac{1}{2} v^2 [Z_{34} + \text{Re} \bar{Z}_5] & -\frac{1}{2} v^2 \text{Im} \bar{Z}_5 \\ 0 & -\frac{1}{2} v^2 \text{Im} \bar{Z}_5 & Y_2 + \frac{1}{2} v^2 [Z_{34} - \text{Re} \bar{Z}_5] \end{pmatrix}$$

With the rephasing freedom we can assume  $Z_5$  is real without loss of generality. Thus, the only complex parameter in the scalar potential is  $Z_7$ .

$$m_h^2 = Z_1 v^2, \quad m_{H,A}^2 = Y_2 + \frac{1}{2} (Z_3 + Z_4 \pm Z_5)$$

The neutral scalar mass eigenstates are states of definite CP with  $h_1$  identified as the SM Higgs boson.

When  $\text{Re}(\bar{Z}_7) \text{Im}(\bar{Z}_7) \neq 0$  the only potential source of CP violation resides in the scalar self-interactions.

# The condition $\text{Re}(\bar{Z}_7) \text{Im}(\bar{Z}_7) \neq 0$

neutral scalar	$\text{Re}(\bar{Z}_7) \neq 0$ and $\text{Im}(\bar{Z}_7) = 0$	$\text{Re}(\bar{Z}_7) = 0$ and $\text{Im}(\bar{Z}_7) \neq 0$
$h_1$	CP-even	CP-even
$h_2$	CP-even	CP-odd
$h_3$	CP-odd	CP-even
$G$	CP-odd	CP-odd

The interactions of  $h_1$  remain CP-conserving while cubic and quartic scalar interactions involving  $h_2$  and  $h_3$  violate CP:

$$h_3 h_3 h_3 \quad h_3 h_2 h_2 \quad h_3 H^+ H^-, \quad h_3 h_3 h_3 h_1, \quad h_3 h_1 h_2 h_2, \quad h_3 h_1 H^+ H^-$$

since each interaction vertex involves an odd number of would-be CP-odd scalar fields.

# P-conserving, CP-violating processes in the 2HDM

We impose these requirements:

- all processes survive the Higgs alignment limit
- no quartic scalar couplings are involved (due to coupling and phase space suppressions)
- the dominant contribution to the CP-violating signal is P-even (C-even, CP-violating contributions due to scalar-fermion couplings are absent or suppressed)

Simultaneous observation of processes involving:

1.  $h_2 H^+ H^-$ ,  $h_3 H^+ H^-$ ,  $Z h_2 h_3$ ,
2.  $h_2 h_k h_k$ ,  $h_3 H^+ H^-$ ,  $Z h_2 h_3$ , (for  $k = 2$  or  $3$ ),
3.  $h_3 h_k h_k$ ,  $h_2 H^+ H^-$ ,  $Z h_2 h_3$ , (for  $k = 2$  or  $3$ ),
4.  $h_2 h_k h_k$ ,  $h_3 h_\ell h_\ell$ ,  $Z h_2 h_3$ , (for  $k, \ell = 2$  or  $3$ ).

would signal CP violation.

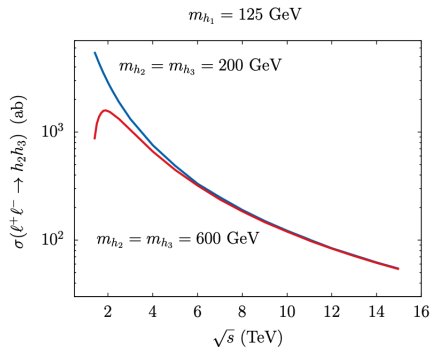


# Practical considerations and discovery potential

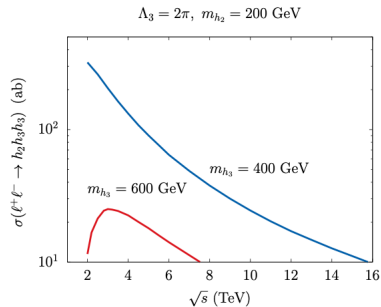
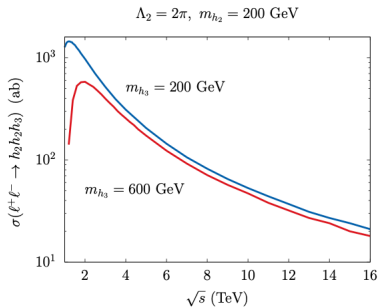
- Some scenarios involve only the discovery of just two neutral states while some also require the observation of a charged Higgs boson
- We omit cases that involve the cubic self interaction of a given scalar
- We focus on processes where  $h_2$ ,  $h_3$  can be observed via an  $s$ -channel  $Z$  exchange
- We also study processes in which  $h_2$ ,  $h_3$  and  $H^\pm$  can be observed
- Once  $h_2$  and  $h_3$  are established, one can search for  $h_2 h_2 h_3$  and  $h_2 h_3 h_3$  production (which are also mediated via an  $s$ -channel  $Z$  exchange)
- Simultaneous observation of both production channels would signal the presence of P-even CP violation

# Discovery potential at future colliders

Accelerator	$\sqrt{s}$ (TeV)	Integrated luminosity ( $ab^{-1}$ )
CLIC	1.5	2.5
CLIC	3	5
Muon Collider	3	1
Muon Collider	7	10
Muon Collider	14	20

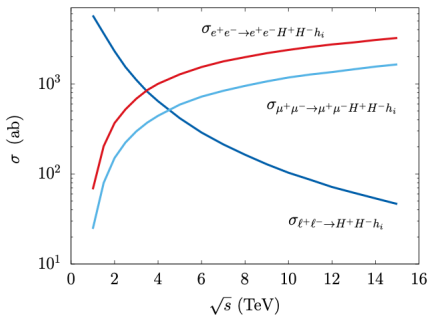


# Discovery potential at future colliders

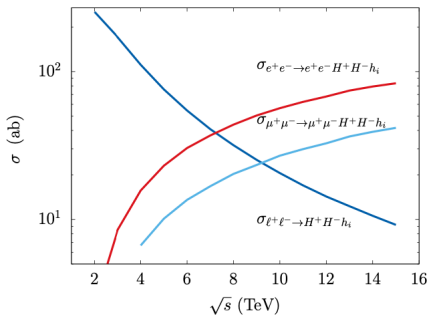


# Discovery potential at future colliders

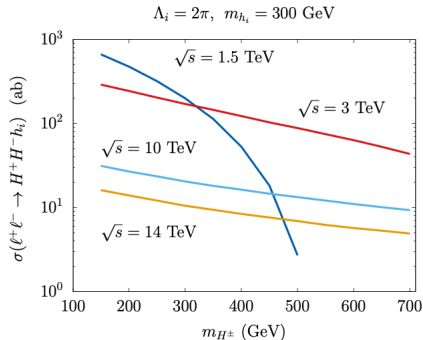
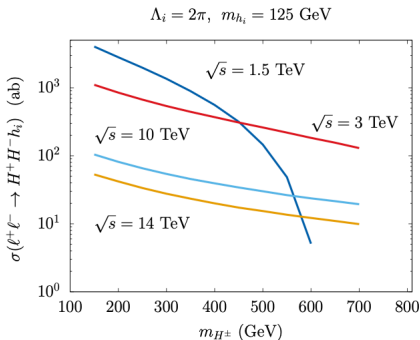
$\Lambda_i = 2\pi$ ,  $m_{h_i} = 125$  GeV,  $m_{H^\pm} = 150$  GeV



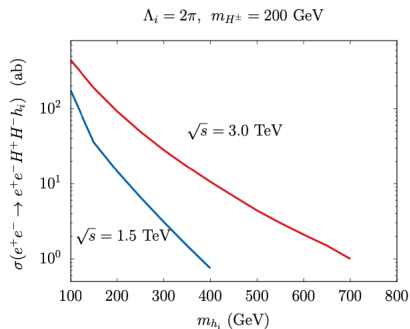
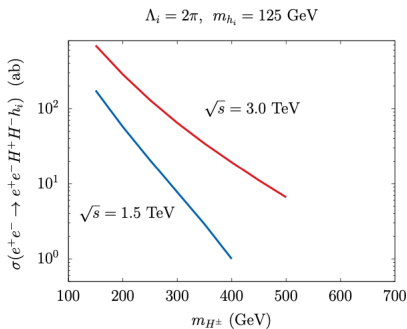
$\Lambda_i = 2\pi$ ,  $m_{h_i} = 300$  GeV,  $m_{H^\pm} = 300$  GeV



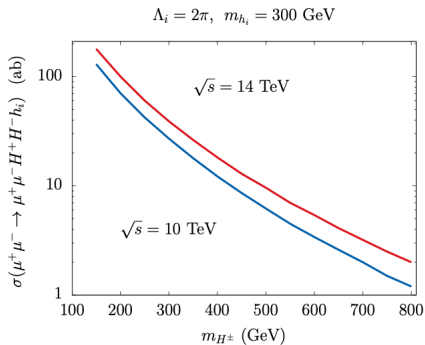
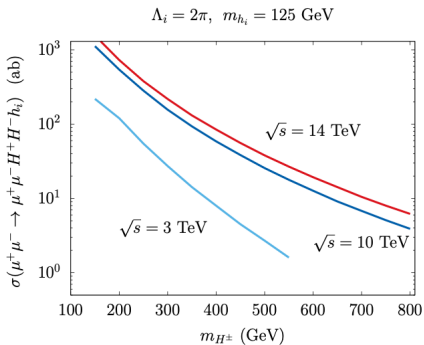
# Discovery potential at future colliders



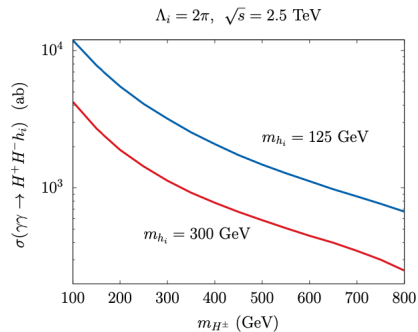
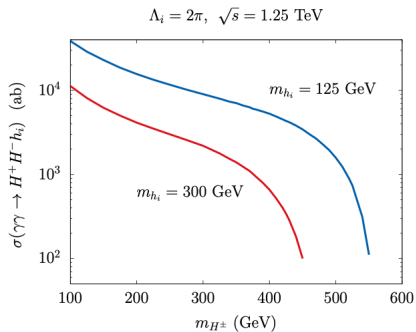
# Discovery potential at future colliders



# Discovery potential at future colliders



# Discovery potential at future colliders





# Summary

- Most studies of Higgs sector CP violation focus on the detection of CP-violating neutral Higgs-fermion Yukawa couplings, which yield P-odd, CP-violating phenomena.
- There is some literature on purely bosonic signatures of Higgs sector CP violation, where the simultaneous observation of three processes (suitably chosen) constitutes a signal of P-even CP violation.
- However, in the examples previously analyzed, some of the processes are strongly suppressed in the approximate Higgs alignment limit (corresponding to the existence of a Standard Model like Higgs boson as suggested by LHC data), in which case the proposed CP-violating signals are difficult to observe in practice.
- We examine processes that do not vanish in the Higgs alignment limit and whose simultaneous observation would provide unambiguous evidence for scalar mediated P-even CP violation.
- We assess the discovery potential of such signals at various future multi-TeV lepton (and  $\gamma\gamma$ ) colliders.
- For loop-induced P-even, CP-violating phenomena, I refer you to our paper ;)