$\Delta \alpha_{had}$: the KNT evaluation and the big picture

Alex Keshavarzi University of Manchester

ECFA WG1 Higgs Factory Parametric uncertainties at future e^+e^- colliders: a_{EM}

 14^{th} July 2022



The University of Manchester





 Dyson summation of Real part of one-particle irreducible blobs Π into the effective, real running coupling α_{QED}:

 $\Pi = \underset{q}{\overset{\gamma^{*}}{\underset{q}{\longrightarrow}}}$

Full photon propagator ~ 1 + Π + $\Pi \cdot \Pi$ + $\Pi \cdot \Pi \cdot \Pi \cdot \Pi$ + ...

$$\rightsquigarrow \qquad \alpha(q^2) = \frac{\alpha}{1 - \operatorname{Re}\Pi(q^2)} = \alpha / \left(1 - \Delta \alpha_{\operatorname{lep}}(q^2) - \Delta \alpha_{\operatorname{had}}(q^2)\right)$$

• The Real part of the VP, $\text{Re}\Pi$, is obtained from the Imaginary part, which via the *Optical* Theorem is directly related to the cross section, $\text{Im}\Pi \sim \sigma(e^+e^- \rightarrow hadrons)$:

$$\Delta lpha_{
m had}^{(5)}(q^2) = -rac{q^2}{4\pi^2 lpha} \operatorname{P} \int_{m_{\pi}^2}^{\infty} rac{\sigma_{
m had}^0(s) \, \mathrm{d}s}{s-q^2} , \quad \sigma_{
m had}(s) = rac{\sigma_{
m had}^0(s)}{|1-\Pi|^2}$$

Identical input as for a_{μ}^{HVP} !



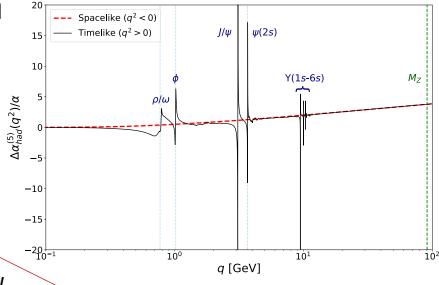


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Full photon propagator ~ 1 + Π + $\Pi \cdot \Pi$ + $\Pi \cdot \Pi \cdot \Pi \cdot \Pi$ + ...

$$\rightarrow \alpha(q^2) = \frac{\alpha}{1 - \operatorname{Re}\Pi(q^2)} = \alpha / \left(1 - \Delta \alpha_{\operatorname{lep}}(q^2) - \Delta \alpha_{\operatorname{had}}(q^2)\right)_{\operatorname{res}}$$

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 q^2 can be:

- Timelike $(q^2 = s > 0)$, rapidly changing function. - Data-driven only.
- Spacelike (q² = t < 0), smooth function.
 Data-driven & Lattice.



 \rightarrow



• Dyson summation of Real part of one-particle irreducible blobs Π into the effective, real running coupling α_{QED} :

 $\Pi = \tilde{\gamma}$

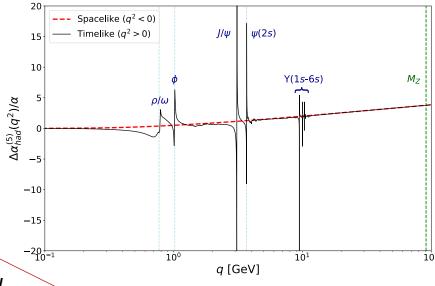
Full photon propagator $\sim 1 + \Pi + \Pi \cdot \Pi + \Pi \cdot \Pi \cdot \Pi + \dots$

$$\rightarrow \alpha(q^2) = \frac{\alpha}{1 - \operatorname{Re}\Pi(q^2)} = \alpha / \left(1 - \Delta \alpha_{\operatorname{lep}}(q^2) - \Delta \alpha_{\operatorname{had}}(q^2)\right)_{\operatorname{res}}$$

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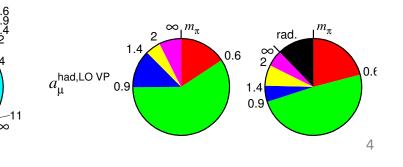
$$\Delta \alpha_{had}^{(5)}(q^2) = -\frac{q^2}{4\pi^2 \alpha} P \int_{m_{\pi}^2}^{\infty} \frac{\sigma_{had}^0(s) \, ds}{s - q^2} , \quad \sigma_{had}(s) = \frac{\sigma_{had}^0(s)}{|1 - \Pi|^2}$$

$$Identical input as for a_{\mu}^{HVP}!$$
Most commonly evaluated at Z-pole: $q^2 = \pm M_Z^2$.
 \Rightarrow Principal component for global EW precision
fits and limits their precision.
KNT19: $\Delta \alpha_{had}^{(5)}(M_Z^2) = 0.02761(12)$
(Note that timelike $\Delta \alpha_{had}^{(5)}(q^2)$ is smooth at $s = M_Z^2$)



 a^2 can be:

- Timelike $(q^2 = s > 0)$, rapidly changing function. - Data-driven only.
- Spacelike $(q^2 = t < 0)$, smooth function. - Data-driven & Lattice.





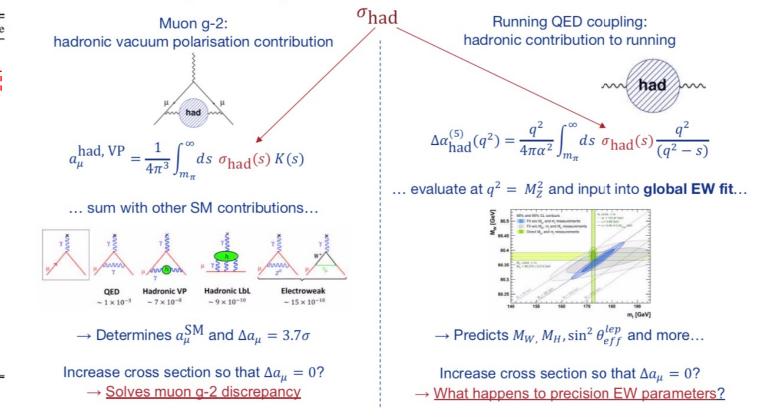
Why is $\Delta \alpha_{had}$ important?

- Limits precision of EW precision fits and so the effectiveness of high-precision EW measurements.
- Can draw a direct parallel with evaluation of the Muon g-2 and probe the muon g-2 discrepancy.
- Is a test of low-energy hadronic theory, e.g. Lattice QCD vs dispersive e^+e^- data.

r V) 7)	Input value 80.379(12)	Fit result 80.359(3)	Result w/o input value
7)		80 250/2)	
	125.10(14)	125.10(14)	$80.357(4)(5) \\ 94^{+20+6}_{-18-6}$
$_{Z}^{2}) \times 10^{4}$	276.1(1.1)	275.8(1.1)	272.2(3.9)(1.2)
<u>)</u>	172.9(4) 0.1179(10)	173.0(4) 0.1180(7)	
7))	91.1876(21) 2.4952(23)	91.1883(20) 2.4940(4)	
)	2.085(42) 41.541(37)	2.0903(4) 41.490(4)	•••
	20.767(25) 0.1721(30)	20.732(4) 0.17222(8)	•••
)	0.21629(66)	0.21581(8) 1.27(2)	•••
Ó	$4.18^{+0.03}_{-0.02}$	$4.18_{-0.02}^{+0.03}$	•••
	0.0171(10) 0.0707(35)	0.01622(7) 0.0737(2)	•••
	0.0992(16)	0.1031(2) 0.1471(3)	•••
	0.670(27)	0.6679(2)	
Q _{FB})	0.2324(12)	0.23152(4)	0.23152(4)(4) 0.23152(4)(4)
	Q _{FB}) ład Coll)	0.0992(16) 0.1499(18) 0.670(27) 0.923(20) $2_{FB})$ 0.2324(12)	$\begin{array}{cccc} 0.0707(35) & 0.0737(2) \\ 0.0992(16) & 0.1031(2) \\ 0.1499(18) & 0.1471(3) \\ 0.670(27) & 0.6679(2) \\ 0.923(20) & 0.93462(7) \\ 0.2324(12) & 0.23152(4) \end{array}$

Uncertainty from

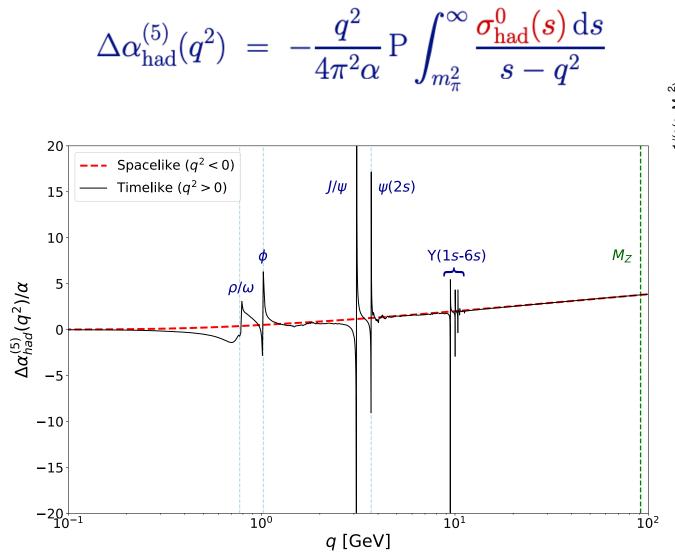
Experimentally measured hadronic cross section:

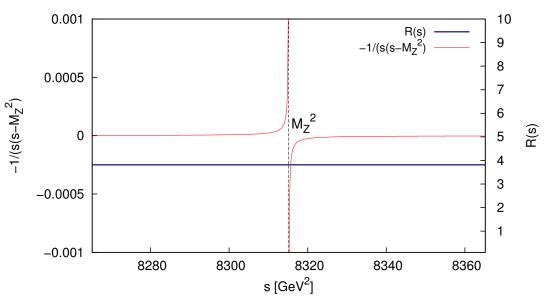


Keshavarzi, Marciano, Passera and Sirlin, *Phys.Rev.D* 102 (2020) 3, 033002



 $\Delta \alpha_{\rm had}$: M_Z^2 vs. low Q^2





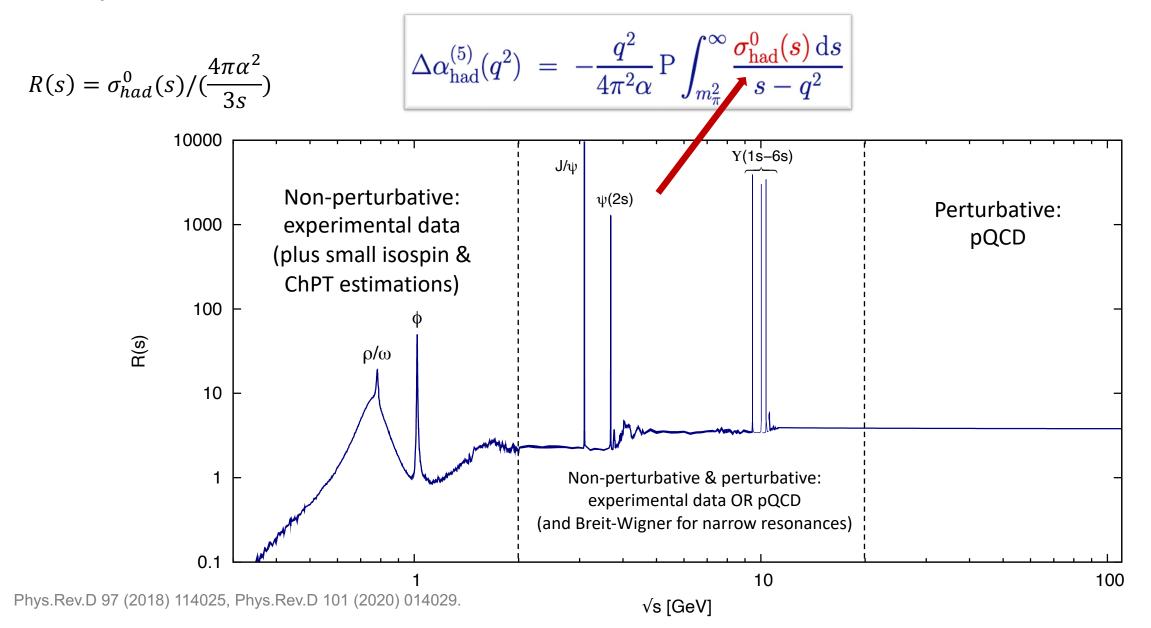
- Both functions smooth at $Q^2 = \pm M_Z^2$. Not true for lowenergy timelike resonances.
- Timelike principal value integral over combined (nonsmooth) data evaluated at low Q^2 can result in fluctuations.

 \rightarrow Improved KNT VP evaluation planned to address this.

Differences for $Q^2 = \pm M_Z^2$ are: KNT19: $\Delta \alpha_{had}^{(5)}(+M_Z^2) - \Delta \alpha_{had}^{(5)}(-M_Z^2) = 0.40(6) \times 10^{-4}$. F. Jegerlehner, 1905.05078 - pQCD: = 0.45(2)×10^{-4}.



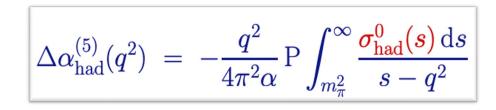
The KNT evaluation of $\sigma_{had}^0(s)$

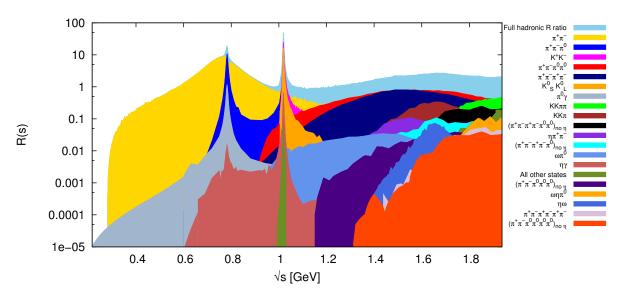


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Must calculate $\sigma_{had}^0(s)$ for each hadronic in the low-energy region, and combine many measurements for each final state:

- Radiative corrections to account for VP and FSR contributions, with corresponding systematic uncertainties.
- Re-binning of data, with scans to optimise e.g., resonance regions.
- Combine measurements via linear χ^2 minimisation in given channel, avoiding fitting biases and incorporating all available experimental uncertainty and correlation information.
- Inflation of uncertainties in local regions of tension in data.
- Sum all channels (+ isospin/ChPT estimates + BW narrow resonances + pQCD) to get the full $\sigma_{had}^0(s)$.
- Integrate over $\sigma_{had}^0(s)$ for all $\pm q^2$ to get $\Delta \alpha_{had}^{(5)}(\pm q^2)$:





To arrive at the final KNT software routine containing full VP function $\Pi(q^2)$:

• Note that $\sigma_{had}^0(s)$ requires a VP correction from $\Pi(q^2)$ which contains the calculated $\Delta \alpha_{had}^{(5)}(q^2)...$

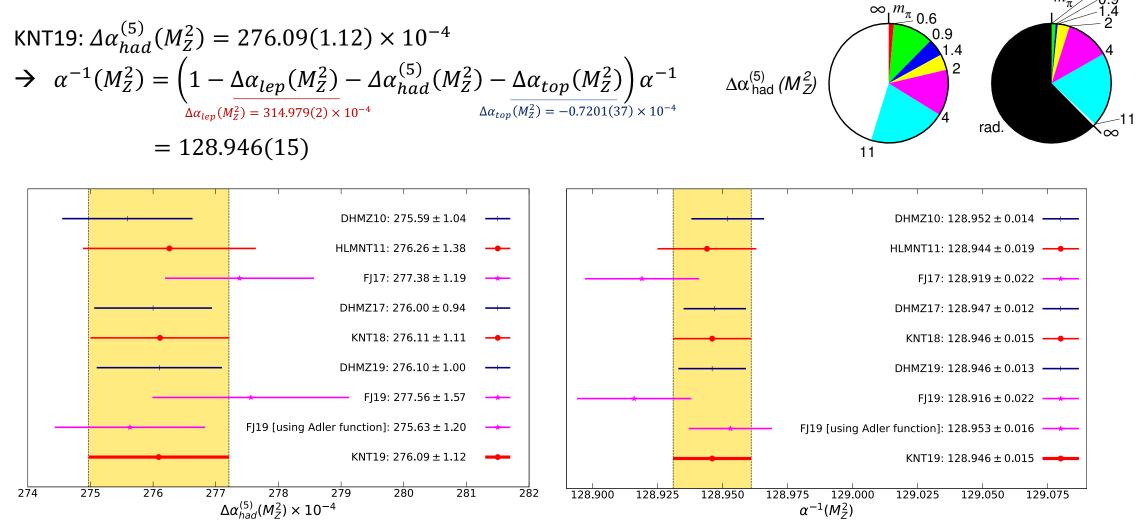
→ Iterate between producing $\Pi(q^2)$ and correcting data until convergence on results.

The self-consistent KNT vacuum polarisation software routine, vp_knt_v3_1, is available by contacting us directly

(and is due to be updated shortly...)



Results for data-driven evaluations of $\Delta \alpha_{had}^{(5)}(M_Z^2)$ and $\alpha(M_Z^2)$



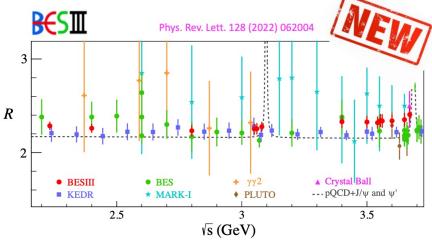
Phys.Rev.D 97 (2018) 114025, Phys.Rev.D 101 (2020) 014029.



Prospects and motivation for improvement

Areas/plans for improvement from KNT:

- New data
 - New cross section measurements are currently being included in preparation for a new update, e.g. BESIII:



- More cross section measurements due to be released.
- Updated data analysis from KNT in the next year(s), including updated VP routine.
- Future plans include a new evaluation of VP with significant improvements and a specific VP-dedicated publication.

Motivation for improvement: Future measurements

 FCC/FCC-ee (for example) would probe new physics at the precision of non-perturbative hadronic corrections to the running coupling for the first time.

→ Order(s) of magnitude improvement expected in e.g., $\sin^2 \theta_{eff}$ and M_W .

World average: $\sin^2 \theta_{eff} = 0.23151(14)$

Erler and Schott, Prog. Part. Nucl. Phys. 2019

EW fit prediction: $\sin^2 \theta_{eff} = 0.23152(4)_{parametric}(4)_{th}$

Keshavarzi, Marciano, Passera and Sirlin, Phys.Rev.D 102 (2020) 033002, using Gfitter

Parametric error 4×10^{-5} on $\sin^2 \theta_{eff}$ is dominated by $\Delta \alpha_{had}^{(5)}$ (M_Z^2) uncertainty.

- Without an improvement in the precision of $\Delta \alpha_{had}^{(5)}$, the precision of the EW fit prediction will become more precise than the current best determination!
- Need an improvement ~×3 in $\Delta \alpha_{had}^{(5)}$ precision to make it compatible with such measurements (e.g. $\sin^2 \theta_{eff}$ precision $\leq 1 \times 10^{-5}$).

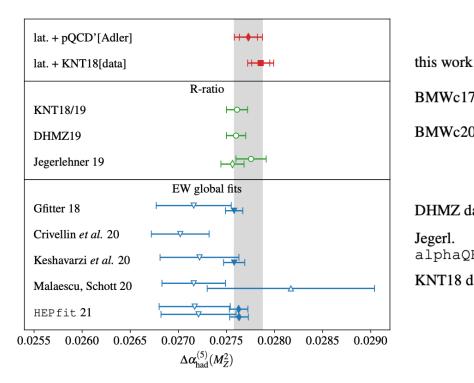


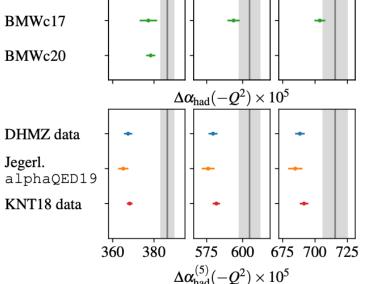
Prospects and motivation for improvement

 $3.0 \,\mathrm{GeV^2}$

 $5.0 \,\mathrm{GeV^2}$

Motivation for improvement: tensions with lattice QCD





 $1.0 \, \text{GeV}^2$

Tension with data-driven results washed out at the Z pole.

Up to 3.5σ tension with data-driven results between 1 and 7 GeV² (comparable to g-2 discrepancy...). Other prospects for improvement:

- New low-energy data for $\sigma_{had}^0(s)$ (CMD-3, SND, KEDR, BESIII, Belle-2, ...).
- Direct determination of $\Delta \alpha_{had}^{(5)}(M_Z^2)$ measuring the muon asymmetry $A_{FB}^{\mu\mu}(s)$ in the vicinity of the *Z*-pole (see Patrick Janot's talk in this workshop).
- Euclidean split method (Adler function). Needs spacelike offset $\Delta \alpha_{had}^{(5)}$ ($-M_0^2$) with $-M_0^2 \sim 2$ GeV and pQCD (see Fred Jegerlehner's talk in this workshop).
- Direct measurement of $\Delta \alpha_{had}^{(5)}(q^2)$ from MUonE muon-electron scattering experiment.
- More lattice QCD evaluations..



The muon g-2 and $\Delta \alpha$ connection

Keshavarzi, Marciano, Passera and Sirlin, Phys.Rev.D 102 (2020) 3, 033002

- Shift KNT hadronic cross section in fully energy-dependent (point-like and binned) analysis to account for Δa_{μ} .
- Input new values of $\Delta \alpha$ into Gfitter to predict EW observables.
- Analysis greatly constrained from more precise EW observables measurements and more comprehensive hadronic cross section.
 - Can Δa_{μ} be due to hypothetical mistakes in the hadronic $\sigma(s)$?
 - An upward shift of $\sigma(s)$ also induces an increase of $\Delta \alpha_{had}^{(5)}(M_Z)$.
 - Consider:

$$\begin{aligned} \mathbf{a}_{\mu}^{\text{HLO}} &\to \\ \mathbf{a} &= \int_{4m_{\pi}^2}^{s_u} ds \, f(s) \, \sigma(s), \qquad f(s) = \frac{K(s)}{4\pi^3}, \ s_u < M_Z^2, \\ \mathbf{\Delta}\alpha_{\text{had}}^{(5)} &\to \end{aligned} \\ b &= \int_{4m_{\pi}^2}^{s_u} ds \, g(s) \, \sigma(s), \qquad g(s) = \frac{M_Z^2}{(M_Z^2 - s)(4\alpha\pi^2)}, \end{aligned}$$

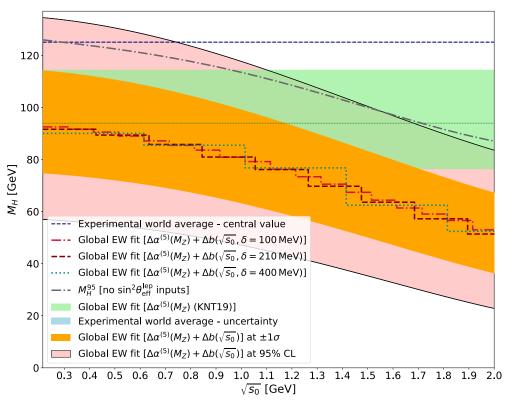
and the increase

Note the very different energydependent weighting of the integrands... $\Delta\sigma(s) = \epsilon\sigma(s)$

 ϵ >0, in the range:

$$\sqrt{s} \in [\sqrt{s_0} - \delta/2, \sqrt{s_0} + \delta/2]$$

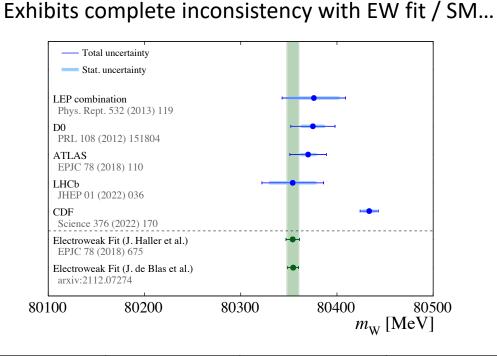
Use Gfitter and precise and up-to-date compilation of total hadronic cross section from KNT, Keshavarzi, Nomura and Teubner, Phys.Rev.D 101 (2020) 014029.



Shifting $\Delta \sigma(s)$ to fix Δa_{μ} is possible but excluded above ~ 1 GeV.

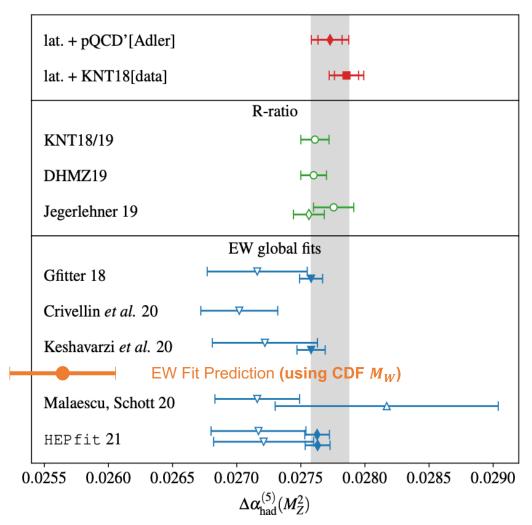


Aside: What about the new CDF M_W ?



Parameter	Input Value (Experiment)	EW Fit Prediction (No CDF <i>M_W</i>)	EW Fit Prediction (using CDF <i>M</i> _W)
$\Delta \alpha_{had}^{(5)}(M_Z^2) \times 10^4$	276.1(1.1)	272.2(4.0) [0.9σ]	256.4(3.9) [4.9σ]
M _H (GeV)	125.10(14)	94(21) [1.5σ]	40(9) [9.5σ]

Taken largest asymmetric $M_{\!H}$ error as symmetric error



And bear in mind that comparisons of the lattice vs data-driven $\Delta \alpha_{had}^{(5)}(q^2)$ are worse evaluated at low-energies... 13



Conclusions

- $\Delta \alpha_{had}$ can be evaluated dispersively from $e^+e^- \rightarrow$ hadrons cross section data.
- $\Delta \alpha_{had}$ limits precision of EW precision fits and so the effectiveness of high-precision EW measurements.
- Can draw a direct parallel with evaluation of the Muon g-2 and probe the muon g-2 discrepancy.
- Is a test of low-energy hadronic theory, e.g. Lattice QCD vs $e^+e^- \rightarrow$ hadrons data used in dispersive approach.
- $\Delta \alpha_{had}(q^2)$ is different for spacelike (smooth) and timelike (resonant) q^2 .
- KNT and other data-driven evaluations (see talks in the workshop) are in good agreement.
- Prospects for improvement in data-driven evaluation and from direct measurement or Lattice QCD.
- Improvements are crucial for effectiveness of future precision EW measurements and for understanding tension with Lattice QCD.
- $\Delta \alpha_{had}(M_Z^2)$ has been used to constrain new physics in hadronic muon g-2 to < 1 GeV.
- CDF M_W results in low prediction of $\Delta \alpha_{had}(M_Z^2)$ that is inconsistent with the SM.

Thank you.