Interpretation of MINOS data in terms of non-standard neutrino interactions

Stephen Parke
Fermilab

arXiv:1009.0014 with Joachim Kopp and Pedro Machado

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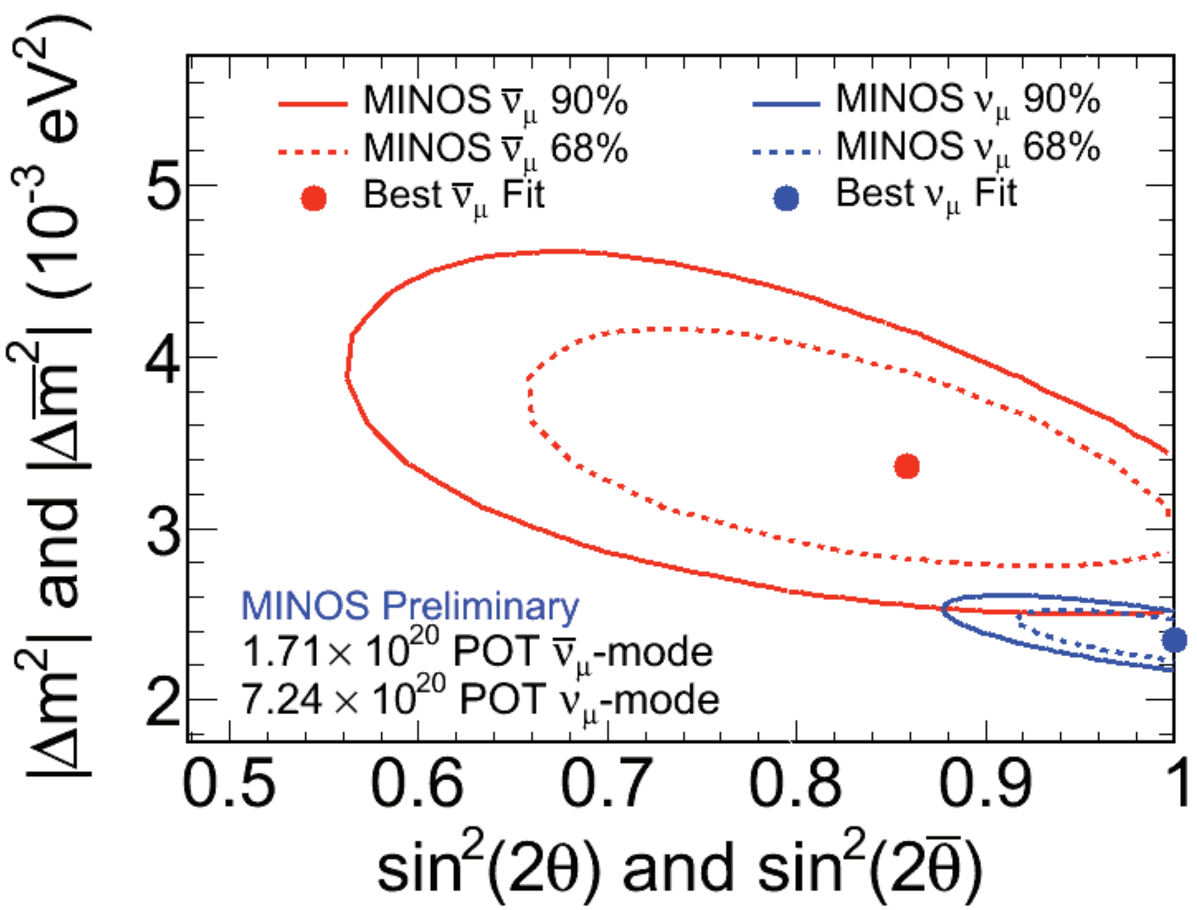
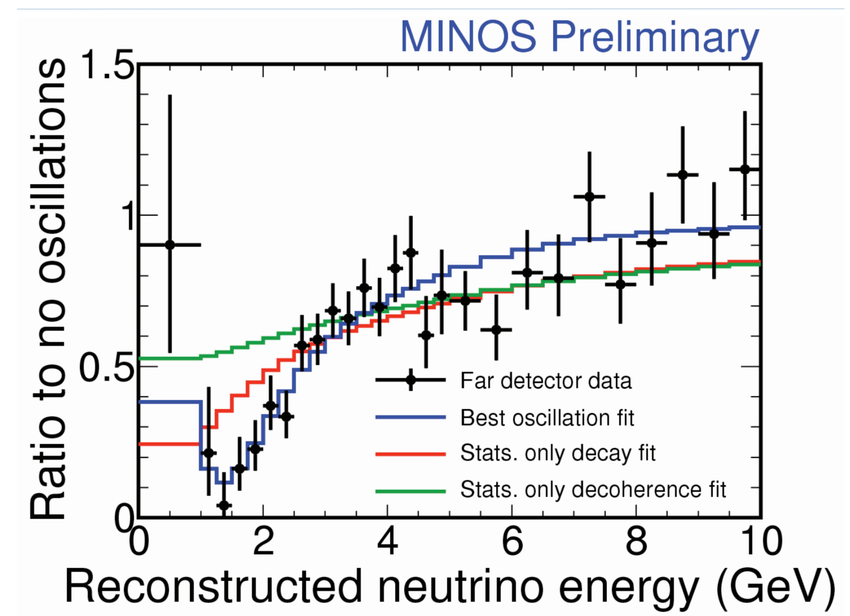
MINOS Results

from P. Vahle, Neutrino 2010

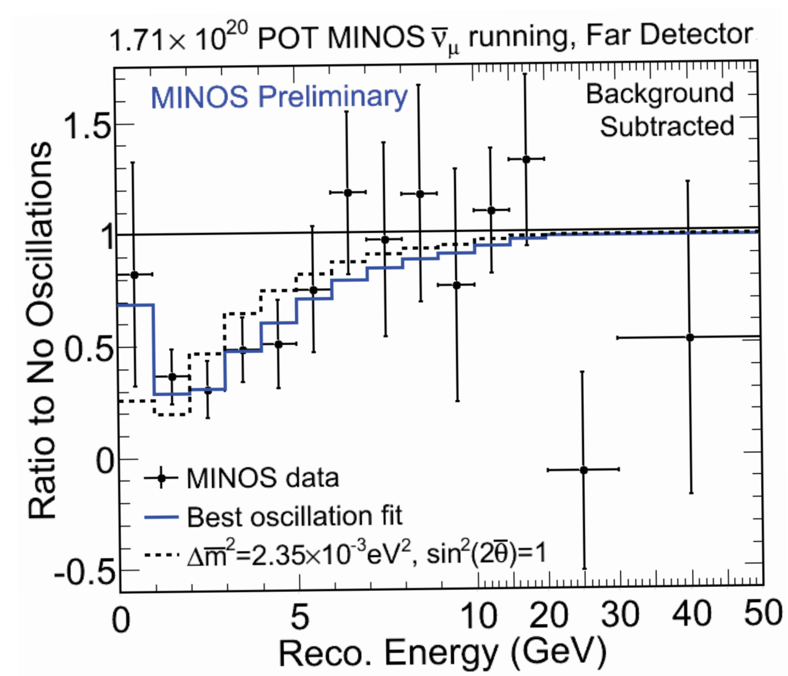


P. Vahle, Neutrino 2010

~2000 events



~100 events



Some Possibilities:

Is this. . .

- . . .Non-Standard Interaction (Mann, Cherdack, Musial, Kafka, 1006.5720; Kopp, Machado, Parke, 1009.0014)?
- . . .sterile neutrino (plus gauged Z' from $U(1)$ according to $B - L$) (Engelhardt, Nelson, Walsh, 1002.4452)?
- . . .gauged ultra-light Z' from $U(1)$ according to $L_\mu - L_\tau$ (Heeck, W.R., 1007.2655)?
- . . .CPT violation? (Barenboim, Lykken, 0908.2993; Choudhury, Datta, Kundu, 1007.2923)?
- . . .nothing and will go away (common sense)?

Outline:

Non-Standard Interactions:

(I) Neutral Current
(II) Charge Current

(I) Neutral Current Non-Standard Interactions:

$$H = \frac{1}{2E}\left[U\begin{pmatrix}0 & \\ & \Delta m^2_{32}\end{pmatrix}U^\dagger + A\begin{pmatrix}\epsilon^m_{\mu\mu} & \epsilon^m_{\mu\tau} \\ \epsilon^{m*}_{\mu\tau} & \epsilon^m_{\tau\tau}\end{pmatrix}\right],$$

$$\epsilon^m_{\tau\tau} - \epsilon^m_{\mu\mu}.$$

$$U = \begin{pmatrix}\cos\theta_{23} & \sin\theta_{23} \\ -\sin\theta_{23} & \cos\theta_{23}\end{pmatrix}$$

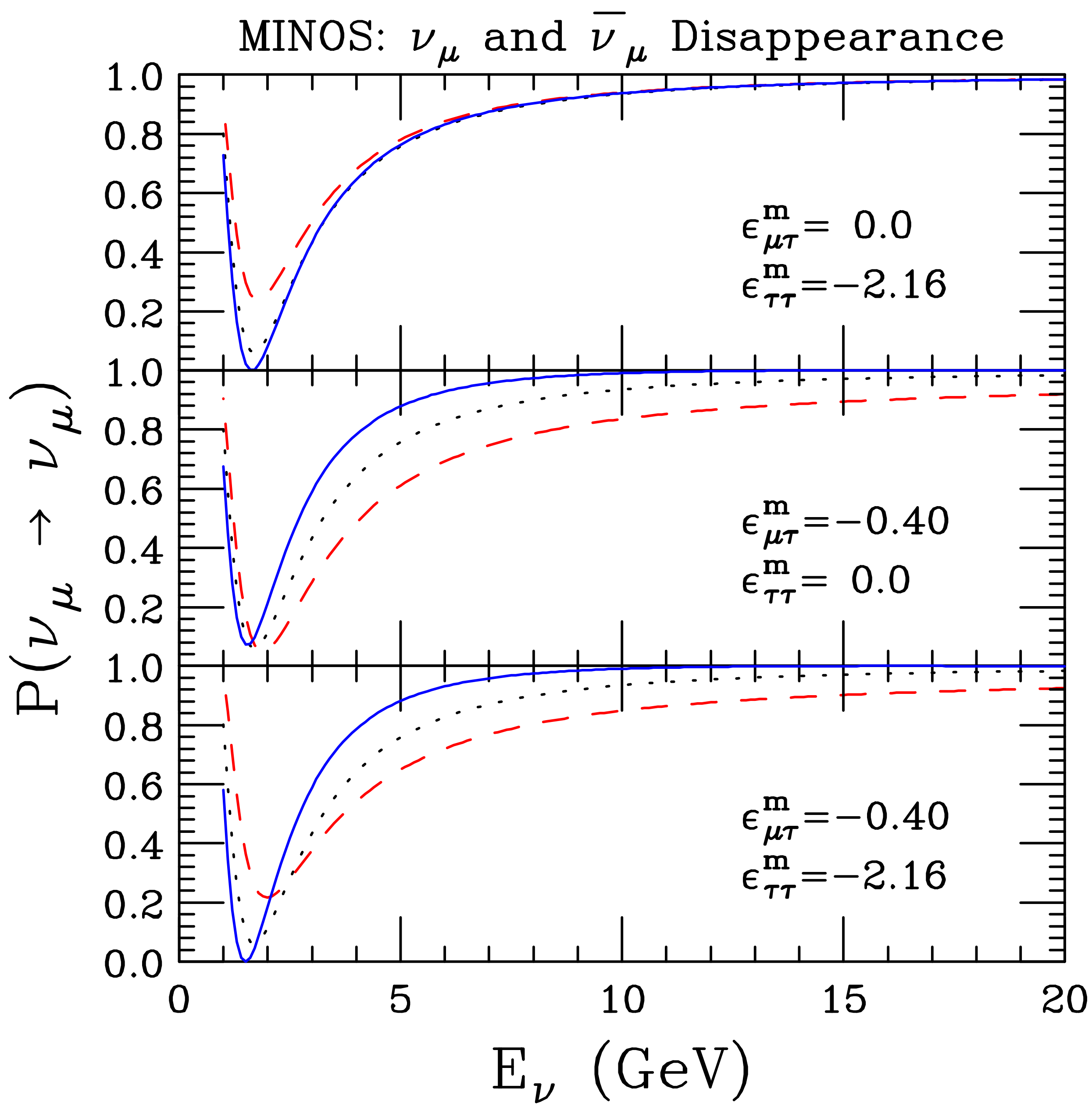
$$P(\nu_\mu \to \nu_\mu) \;=\; 1 - \frac{|\Delta m^2_{32}\sin 2\theta_{23} + 2\epsilon^m_{\mu\tau}A|^2}{\Delta m^4_N}\sin^2\left(\frac{\Delta m^2_N L}{4E}\right)$$

$$\Delta m^2_N \;=\; \sqrt{(\Delta m^2_{32}\cos 2\theta_{23} + \epsilon^m_{\tau\tau}A)^2 + |\Delta m^2_{32}\sin 2\theta_{23} + 2\epsilon^m_{\mu\tau}A|^2}$$

For anti-neutrinos, $\epsilon^m_{\mu\tau} \to \epsilon^{m*}_{\mu\tau}$ and $A \to -A$, so that in matter

$P(\nu_\mu \to \nu_\mu) \neq P(\bar{\nu}_\mu \to \bar{\nu}_\mu)$ without CPT violation.

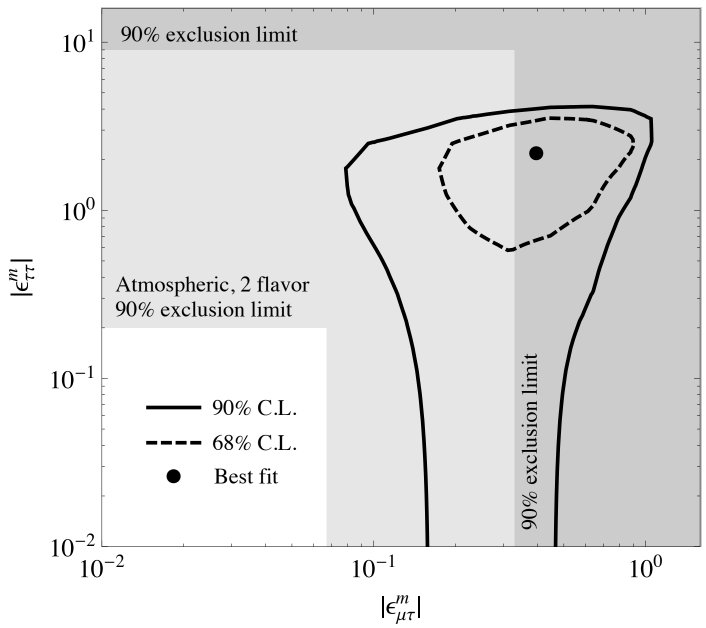
NC-NSI (conti)

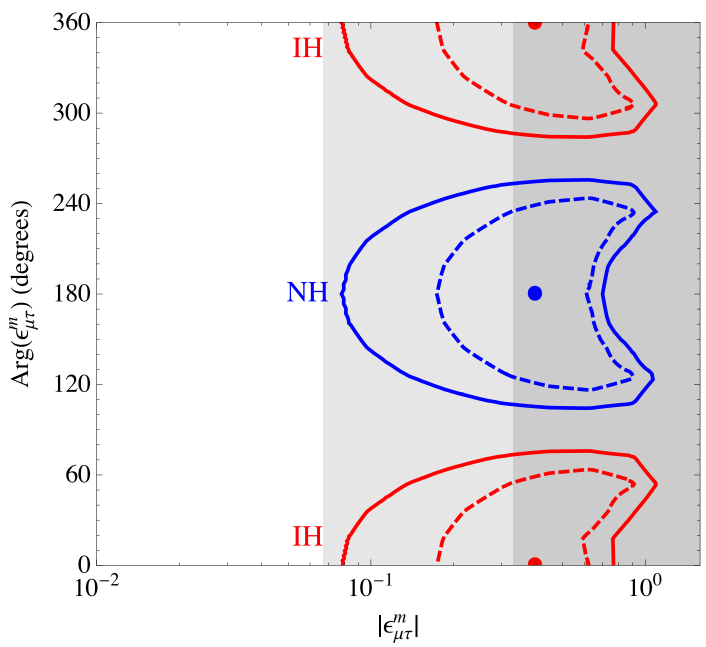


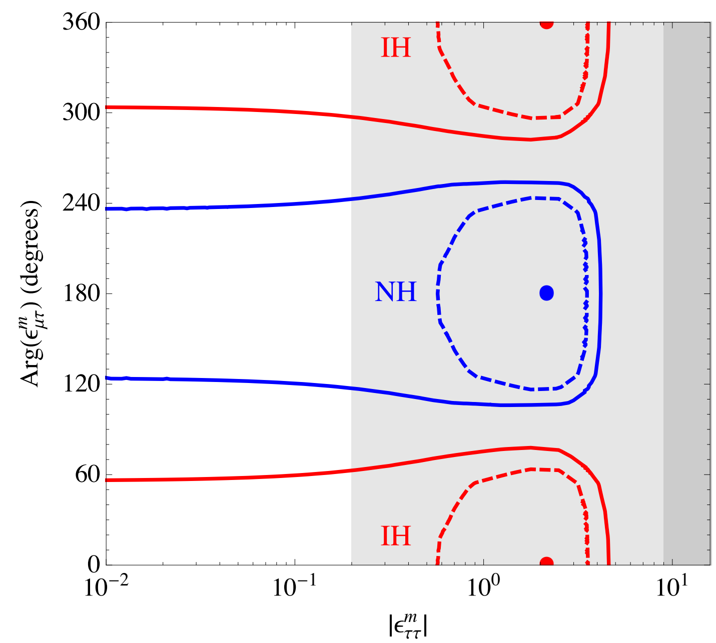
MINOS: ν_μ and $\overline{\nu}_\mu$ Disappearance
$P(\nu_\mu \to \nu_\mu)$
E_ν (GeV)
$\epsilon^{m}_{\mu\tau}$= 0.0
$\epsilon^{m}_{\tau\tau}$=−2.16
$\epsilon^{m}_{\mu\tau}$=−0.40
$\epsilon^{m}_{\tau\tau}$= 0.0
$\epsilon^{m}_{\mu\tau}$=−0.40
$\epsilon^{m}_{\tau\tau}$=−2.16
0.0
0.2
0.4
0.6
0.8
1.0
0
5
10
15
20

NC-NSI (conti)

3 flavor simulation:



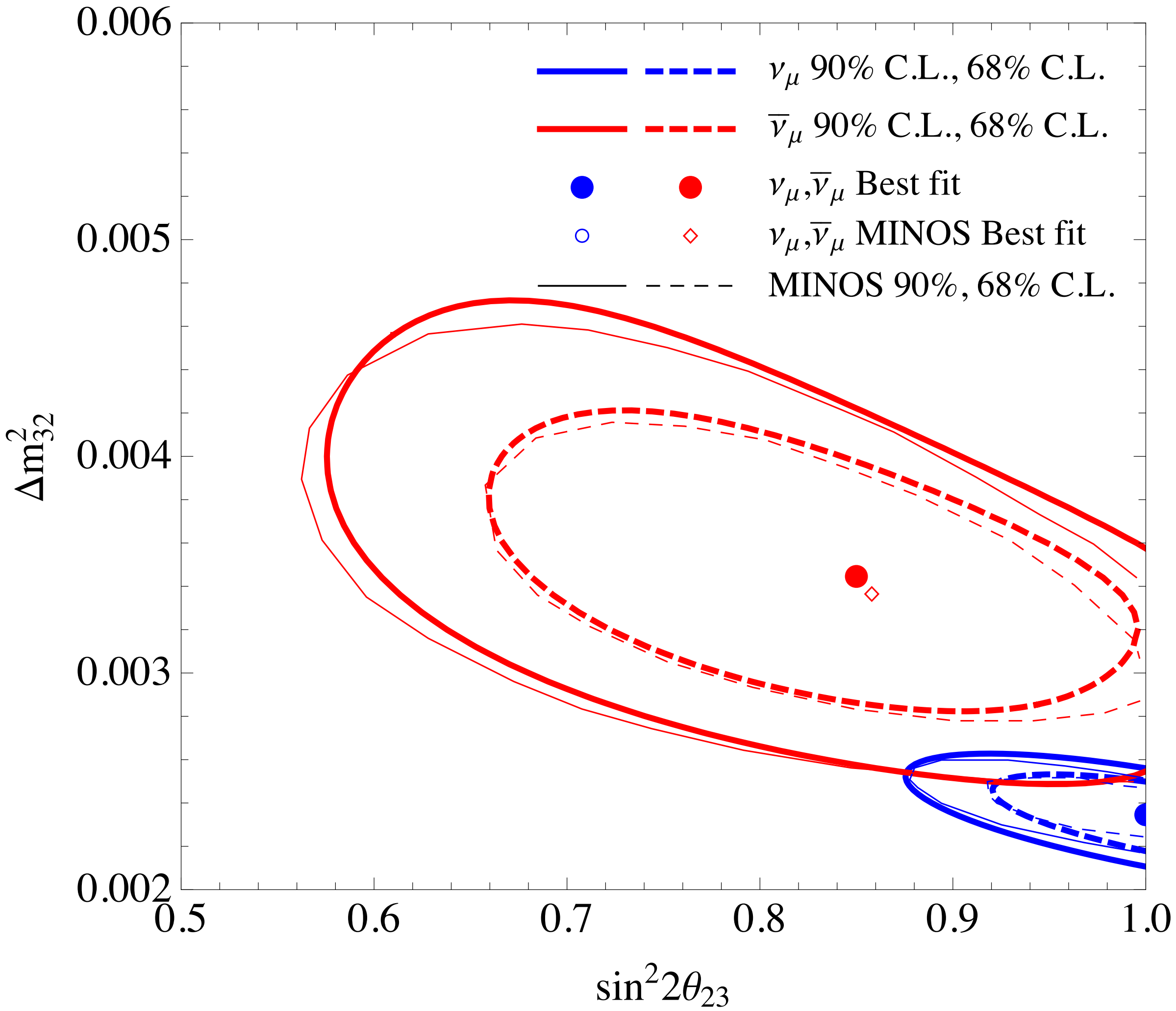




3 times larger than given in

Mann-Cherdack-Musial-Kafka, arXiv:1006.5720 [hep-ph]

NC-NSI (conti)


0.006
0.005
0.004
0.003
0.002
Δm^2_{32}
0.5
0.6
0.7
0.8
0.9
1.0
$\sin^2 2\theta_{23}$
ν_μ 90% C.L., 68% C.L.
$\overline{\nu}_\mu$ 90% C.L., 68% C.L.
$\nu_\mu,\overline{\nu}_\mu$ Best fit
$\nu_\mu,\overline{\nu}_\mu$ MINOS Best fit
MINOS 90%, 68% C.L.

(II) Charge Current NSI

$$\nu_\tau + N \to X + \mu\,, \qquad \mathcal{L}_{\rm NSI} \supset -2\sqrt{2} G_F \epsilon^d_{\tau\mu} V_{ud} \left[\bar{u}\gamma^\rho d\right] \left[\bar{\mu}\gamma_\rho P_L \nu_\tau\right] + h.c.\,,$$

Vector only; not axial vector!

gauge invariance? later

$$\nu_\mu + N \to X + \mu$$

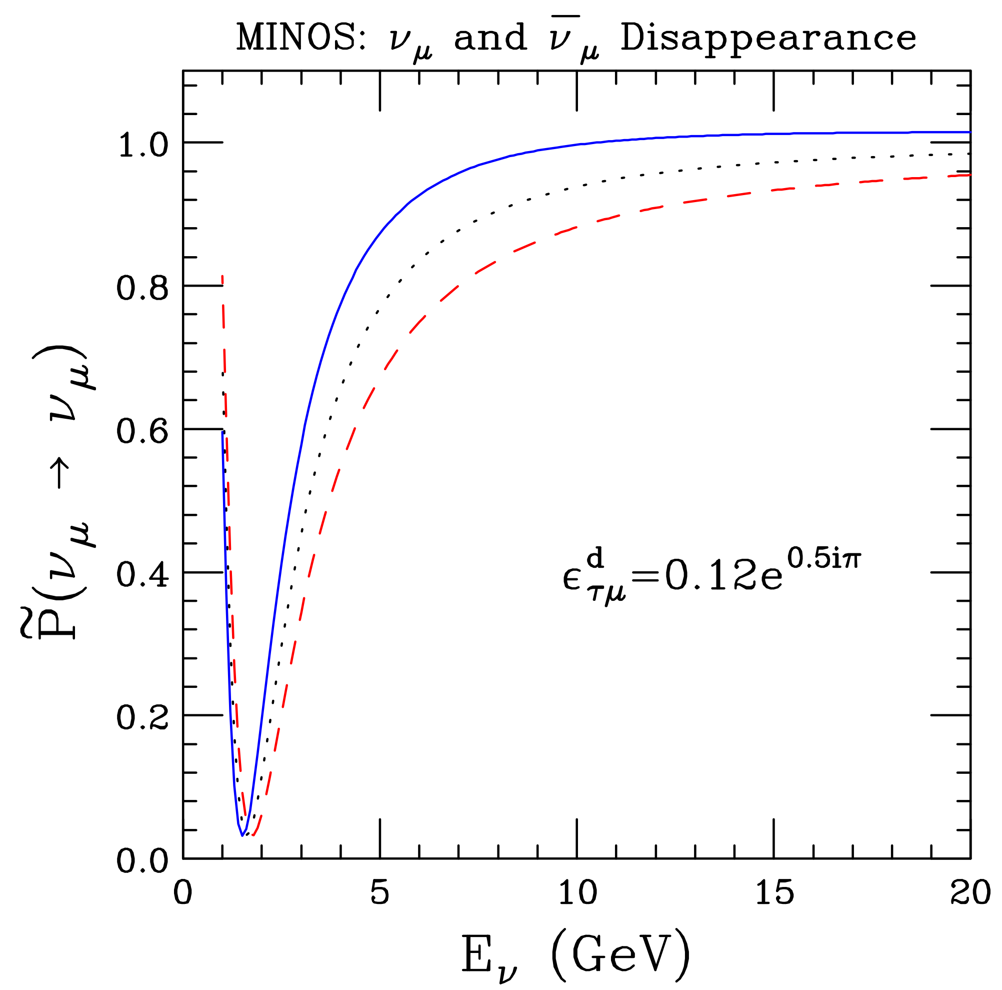
$$\nu_\mu \xrightarrow{\text{osc.}} \nu_\tau + N \to X + \mu$$

$$\tilde{P}(\nu_\mu \to \nu_\mu) = 1 - \left[1 + 2\,|\epsilon^d_{\tau\mu}| \cot 2\theta_{23} \cos\left[\arg(\epsilon^d_{\tau\mu})\right] - |\epsilon^d_{\tau\mu}|^2\right] \sin^2 2\theta_{23} \sin^2\left(\frac{\Delta m^2_{32} L}{4E}\right)$$
$$+ 2\,|\epsilon^d_{\tau\mu}| \sin 2\theta_{23} \sin\left[\arg(\epsilon^d_{\tau\mu})\right] \sin\left(\frac{\Delta m^2_{32} L}{4E}\right) \cos\left(\frac{\Delta m^2_{32} L}{4E}\right)\,.$$

For anti-neutrinos, the sign of $\arg(\epsilon^d_{\tau\mu})$ has to be reversed, and thus

$$\tilde{P}(\nu_\mu \to \nu_\mu) \neq \tilde{P}(\bar{\nu}_\mu \to \bar{\nu}_\mu)\,.$$

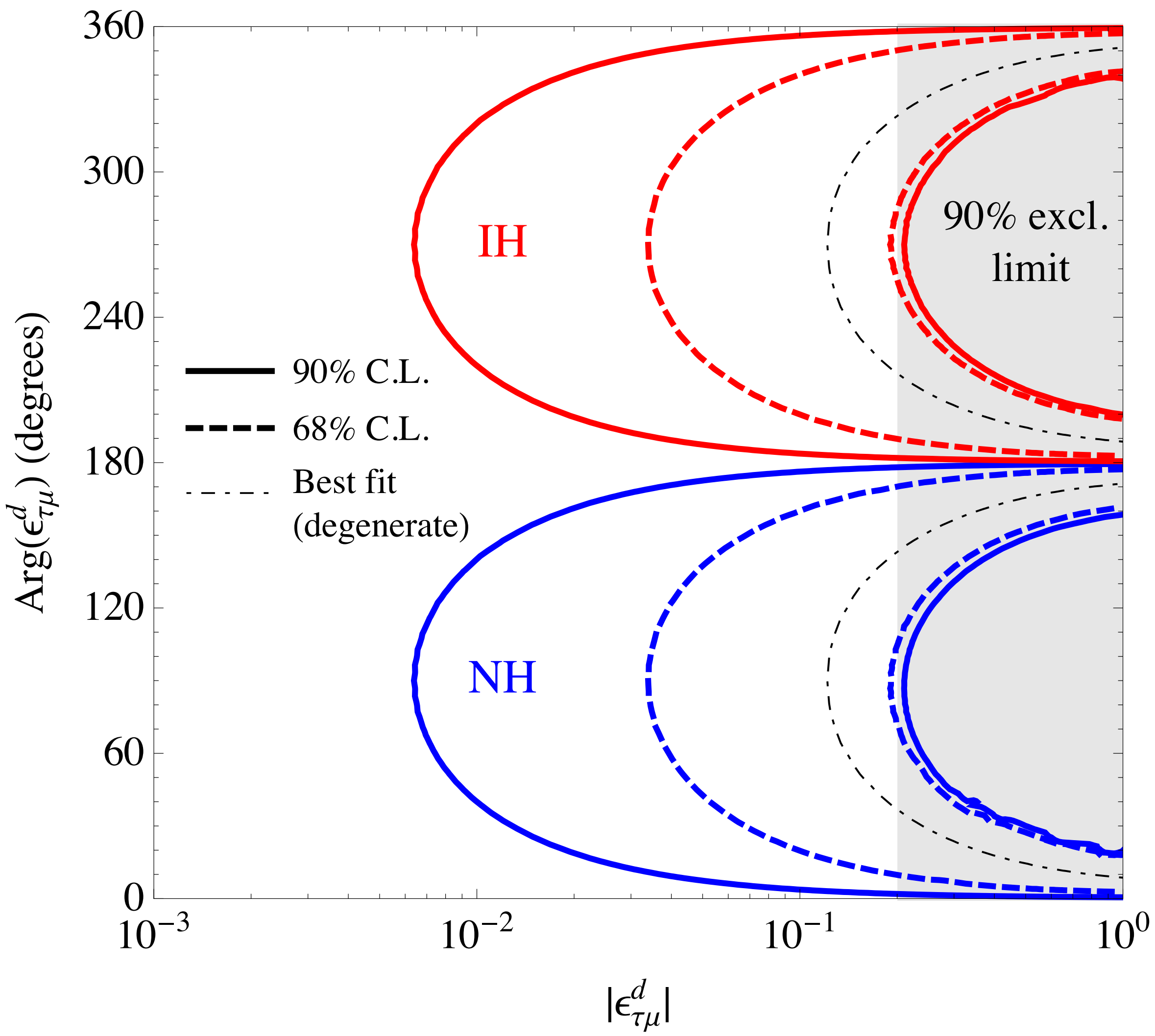
CC-NSI (conti)



MINOS: ν_μ and $\bar{\nu}_\mu$ Disappearance
$\tilde{P}(\nu_\mu \to \nu_\mu)$
1.0
0.8
0.6
0.4
0.2
0.0
0
5
10
15
20
E_ν (GeV)
$\epsilon^{d}_{\tau\mu}=0.12e^{0.5i\pi}$

CC-NSI (conti)

3 flavor simulation:



New Bound on CC-NSI

$$\tau^\pm \to \mu^\pm \pi^0$$

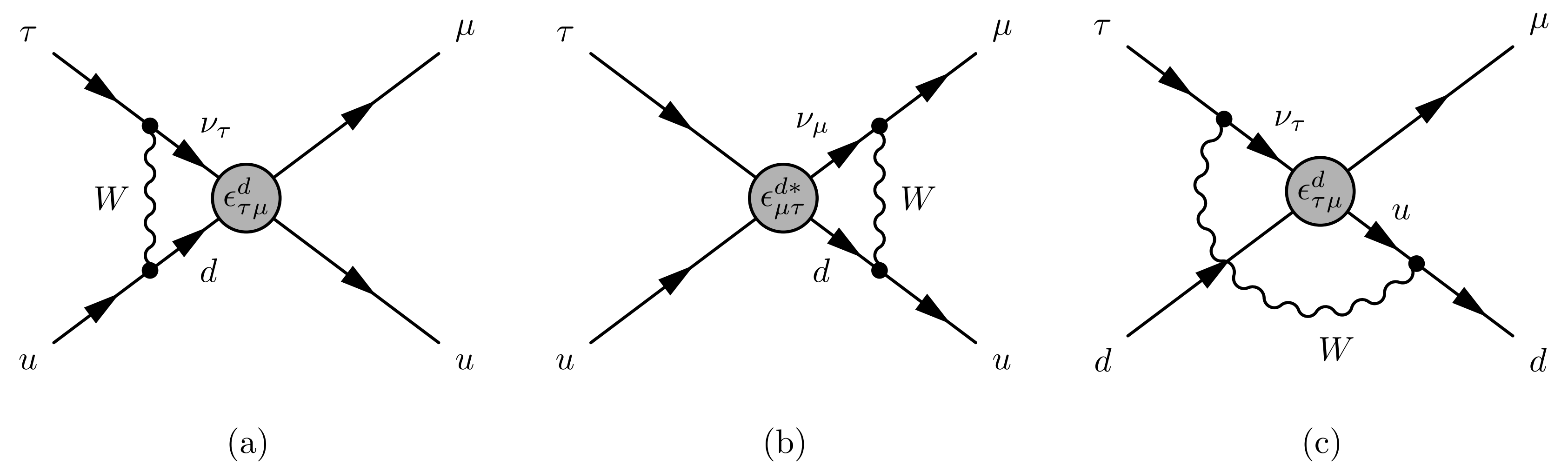


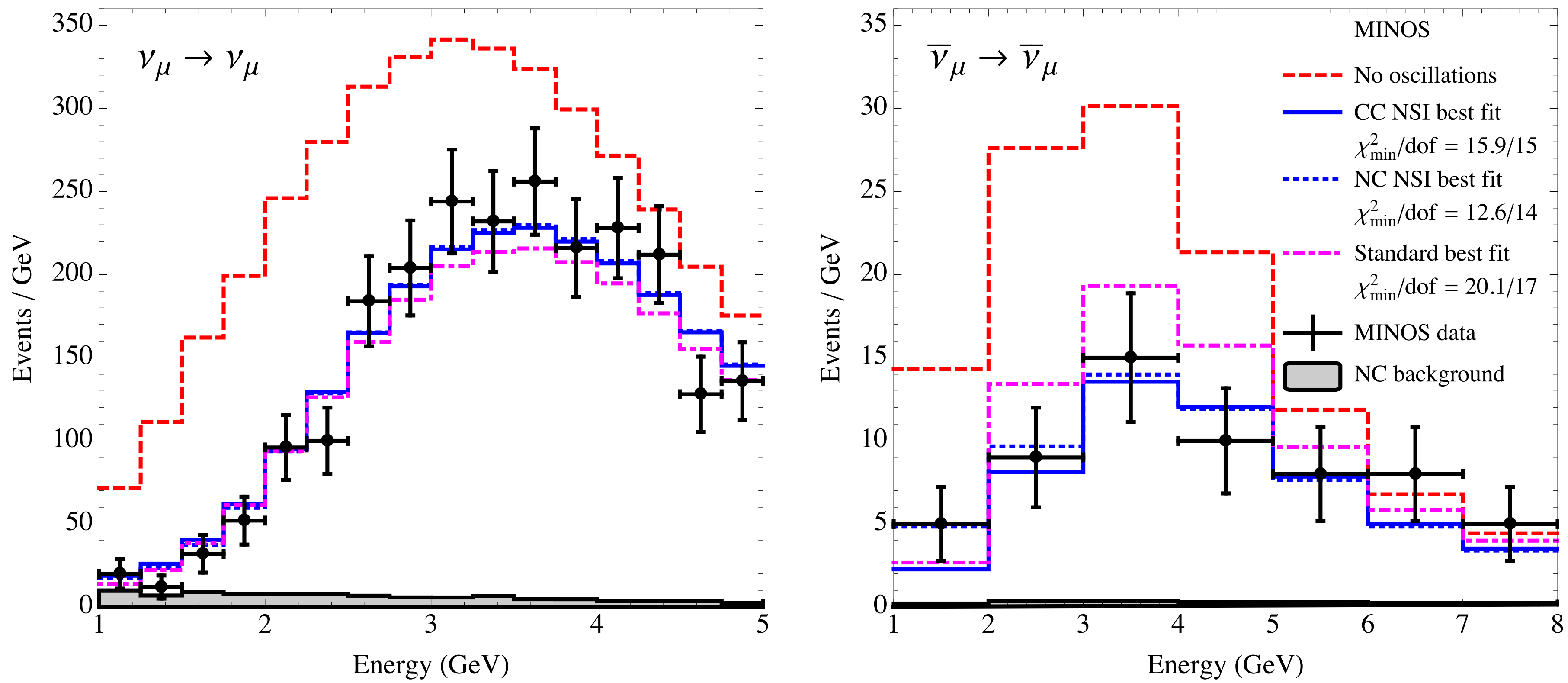
Figure 9: (a) and (b) The loop diagrams used to constrain $\epsilon^d_{\tau\mu}$ and $\epsilon^d_{\mu\tau}$, respectively. (c) A similar diagram that does not have a logarithmic divergence.

$$|\epsilon^d_{\tau\mu}| \simeq \sqrt{\frac{2\,\mathrm{BR}(\tau^\pm \to \mu^\pm\pi^0)}{\mathrm{BR}(\tau^\pm \to \pi^\pm\nu_\tau)}\,\frac{4\pi s_w^2}{3\alpha}}\,.$$

$$\mathrm{BR}(\tau^\pm \to \mu^\pm\pi^0) < 1.1\times 10^{-7}$$

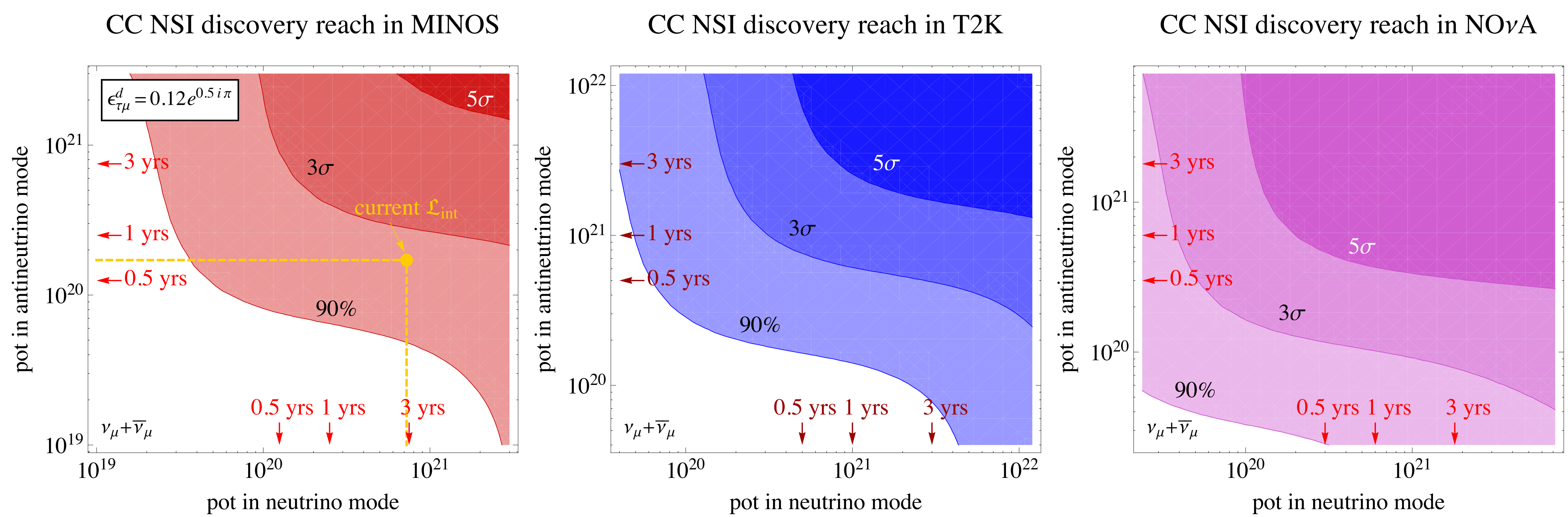
$$|\epsilon^d_{\tau\mu}| \lesssim 0.20\,.$$

CC-NSI (conti)



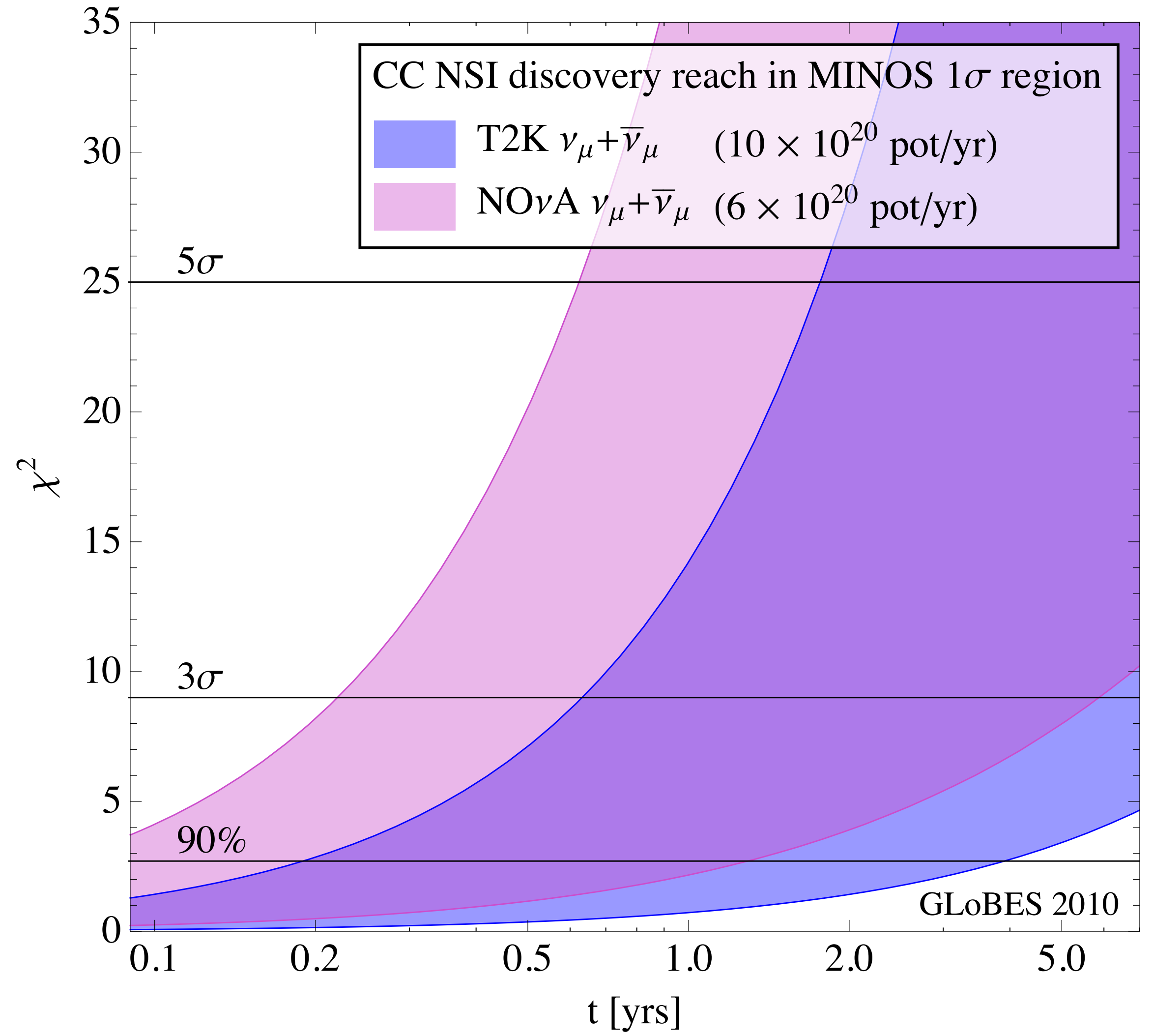
$\nu_\mu \to \nu_\mu$
Events / GeV
Energy (GeV)
$\bar{\nu}_\mu \to \bar{\nu}_\mu$
MINOS
No oscillations
CC NSI best fit
χ^2_{min}/dof = 15.9/15
NC NSI best fit
χ^2_{min}/dof = 12.6/14
Standard best fit
χ^2_{min}/dof = 20.1/17
MINOS data
NC background

CC-NSI in future exp.



Neutrino Factory !!!

CC-NSI in future exp.



Gauge Invariance:

If one imposes SU(2)×U(1) gauge invariance then

$$\mathcal{L}_{\rm NSI} \supset -2\sqrt{2}\, G_F\, \epsilon^d_{\tau\mu}\, V_{ud}\, [\bar{U}\gamma^\mu D]\, [\bar{L}_\mu \gamma_\mu L_\tau]$$

gives tree-level diagram for $\tau \to \mu\, \pi^0$: $|\epsilon^d_{\tau\mu}| \leq 10^{-4}$

Gavela, Talk@NOW2010

see arXiv:1005.0756

The only escape is if our operator is mediated by light particles ($< M_W$)!!!

"Neutrino NSI mediated by light (<< MW), weakly coupled particles—is less well explored in the literature, so a scenario of this type could be responsible for the effects seen by MINOS. This is particularly interesting as models containing light new particles have recently received a lot of attention in the context of Dark Matter searches."

What's interesting MINOS/T2K/NOvA ....

can constrain such operators explicitly!

Conclusions

MINOS anomaly:

- Neutral Current NSI excluded by Atmospheric ν's
- Charge Current NSI can explain the anomaly provide one can evade $SU(2)\times U(1)$ gauge invariance. Explicit model needed!
- MINOS/T2K/NOvA... can directly constrain such CC NSI's.