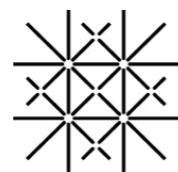


# LFUV at MuC

Admir Greljo

*u*<sup>b</sup> UNIVERSITÄT  
BERN



University  
of Basel

**FNSNF**

SWISS NATIONAL SCIENCE FOUNDATION

[Eccellenza, Project-186866](#)

# Motivation

- The LHCb anomalies in  $b \rightarrow s\mu^+\mu^-$  decays remind us that — New Physics might take an exotic form — an option we should embrace given the present status of the field.
- Several *anomalous* observables: BRs, angular distributions, LFUV ratios.  
see [LHCb Implications](#) next week
- Coherent explanation by a short-distance  $bs\mu\mu$  contact interaction —  $\mathcal{O}(10^{-5})G_F$  — the violation of perturbative unitarity  $\lesssim 100 \text{ TeV}$
- New mass threshold in the vicinity of colliders?
- Today: [Azatov, Garosi, AG, Marzocca, Salko, Trifinopoulos; 2205.13552](#)

Complementary high- $p_T$  searches at future colliders:  
**FCC-hh versus MuC**

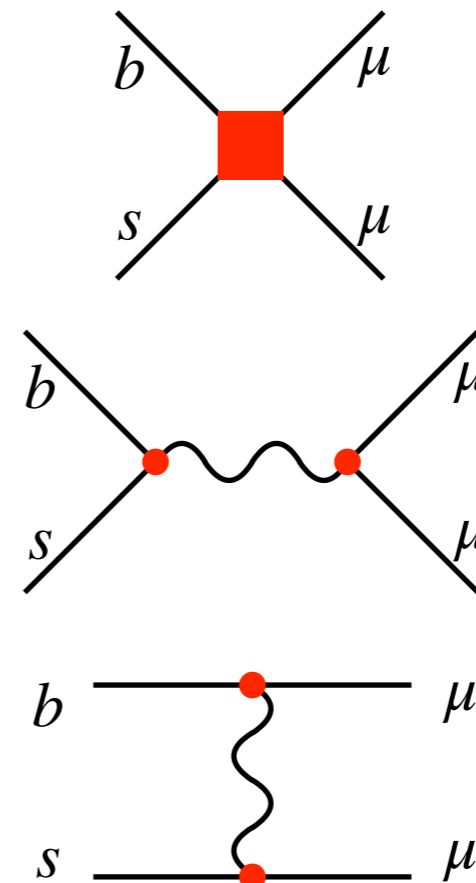
# The scope

## Competitors

Collider	C.o.m. Energy	Luminosity	Label
LHC Run-2	13 TeV	140 fb <sup>-1</sup>	LHC
HL-LHC	14 TeV	6 ab <sup>-1</sup>	HL-LHC
FCC-hh	100 TeV	30 ab <sup>-1</sup>	FCC-hh
Muon Collider	3 TeV	1 ab <sup>-1</sup>	MuC3
Muon Collider	10 TeV	10 ab <sup>-1</sup>	MuC10
Muon Collider	14 TeV	20 ab <sup>-1</sup>	MuC14

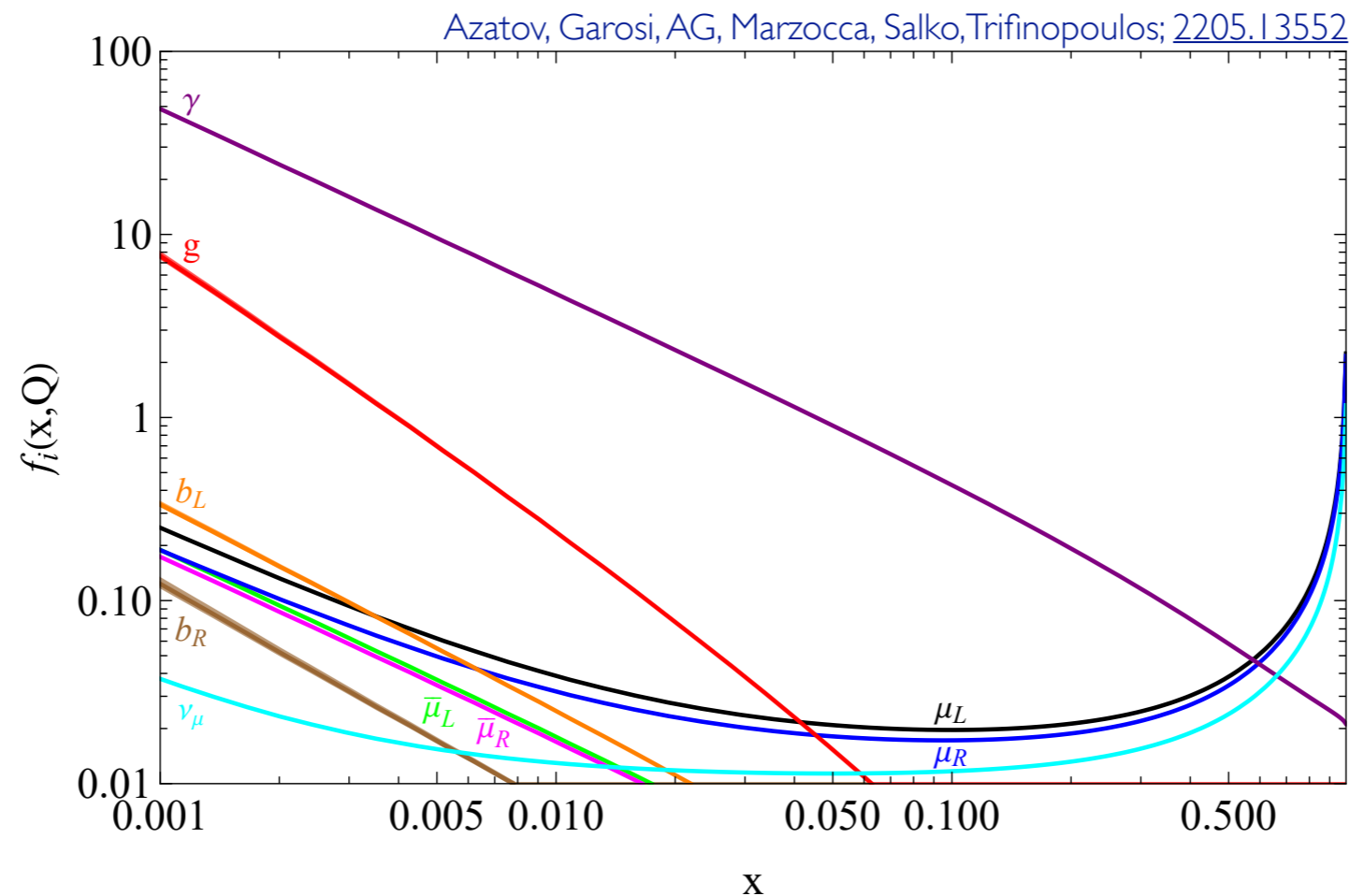
## New Physics benchmarks:

1. Semileptonic 4F interactions
2. Gauged  $U(1)$  extensions
3. Leptoquarks



# The Muon Beam

- Collinear radiation: Spreads the muon energy to lower values and generates different initial states  $\implies$  Parton Distribution Functions
- We cross-check and numerically solve the DGLAP equations from (Han et al, 2007.14300, 2103.09844) with appropriate initial conditions at the LL accuracy
- Selected PDFs at  $Q = 3 \text{ TeV}$ :

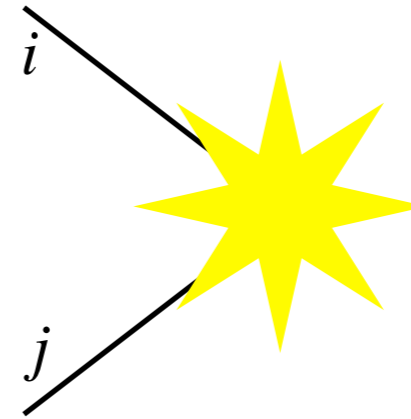




# The Muon Beam

- Parton luminosities

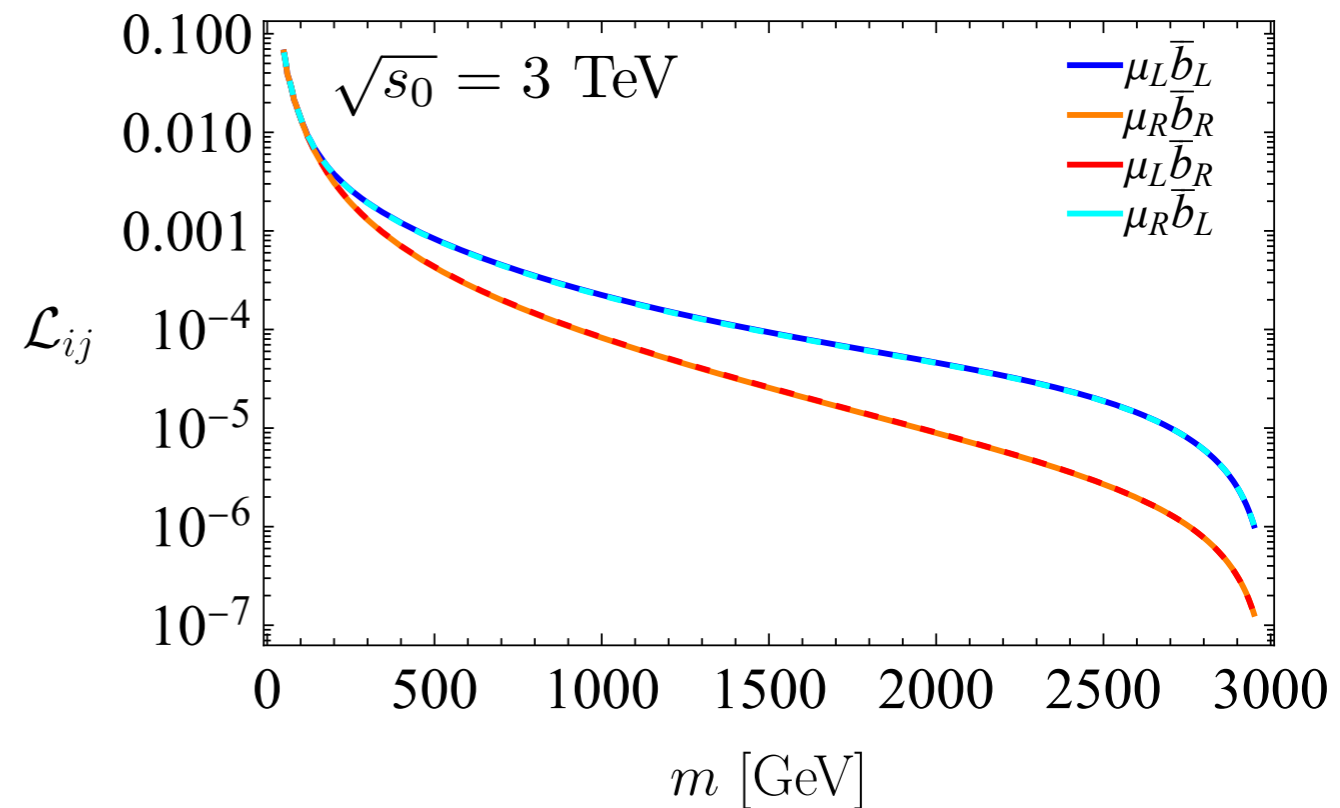
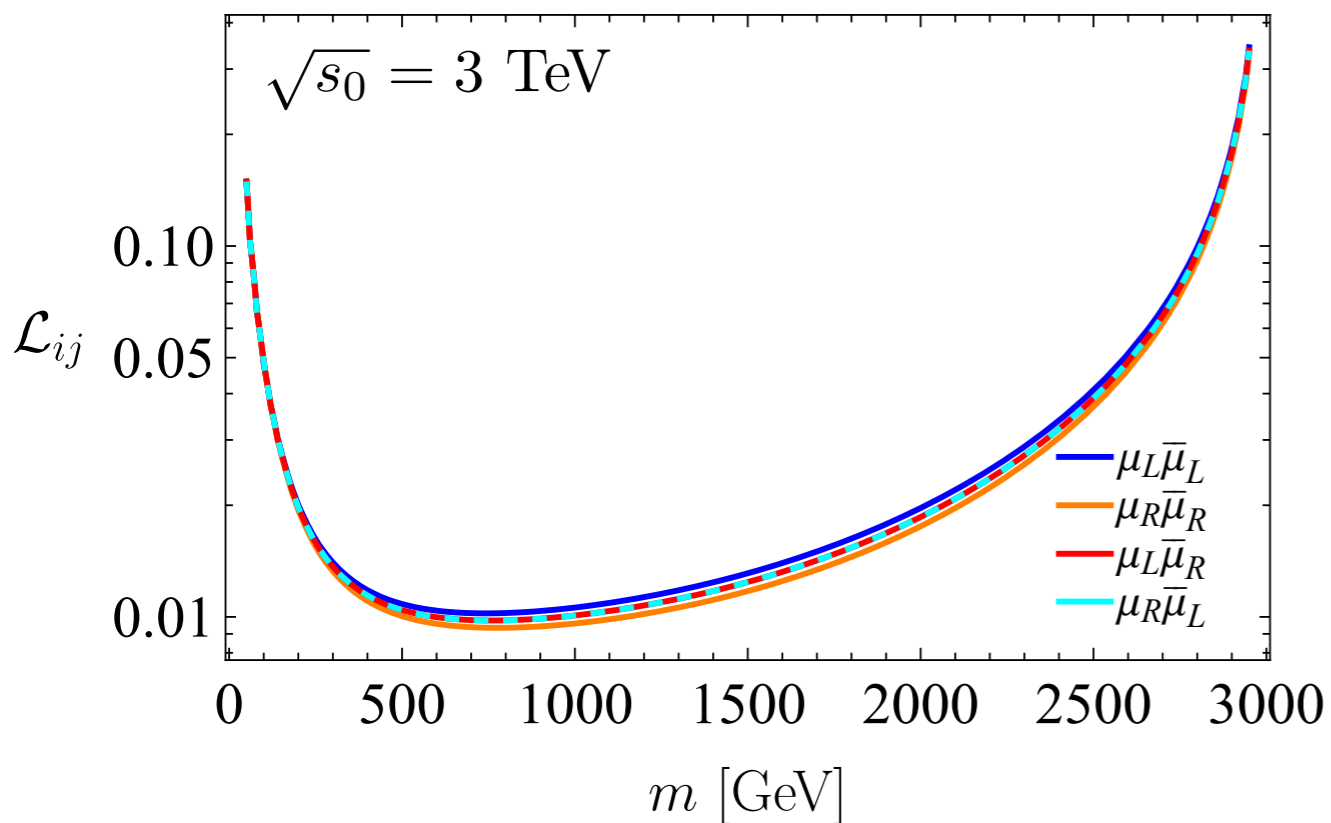
$$\mathcal{L}_{ij}(\tau) = \int_{\tau}^1 \frac{dx}{x} f_i(x, m) f_j\left(\frac{\tau}{x}, m\right)$$



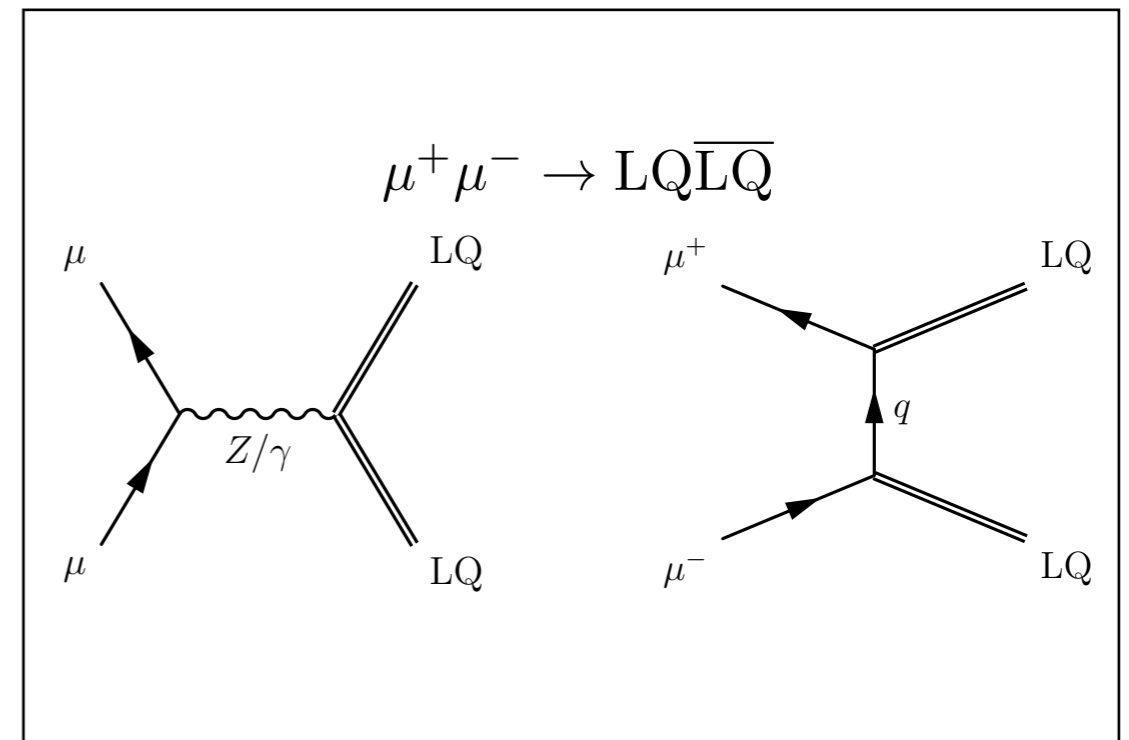
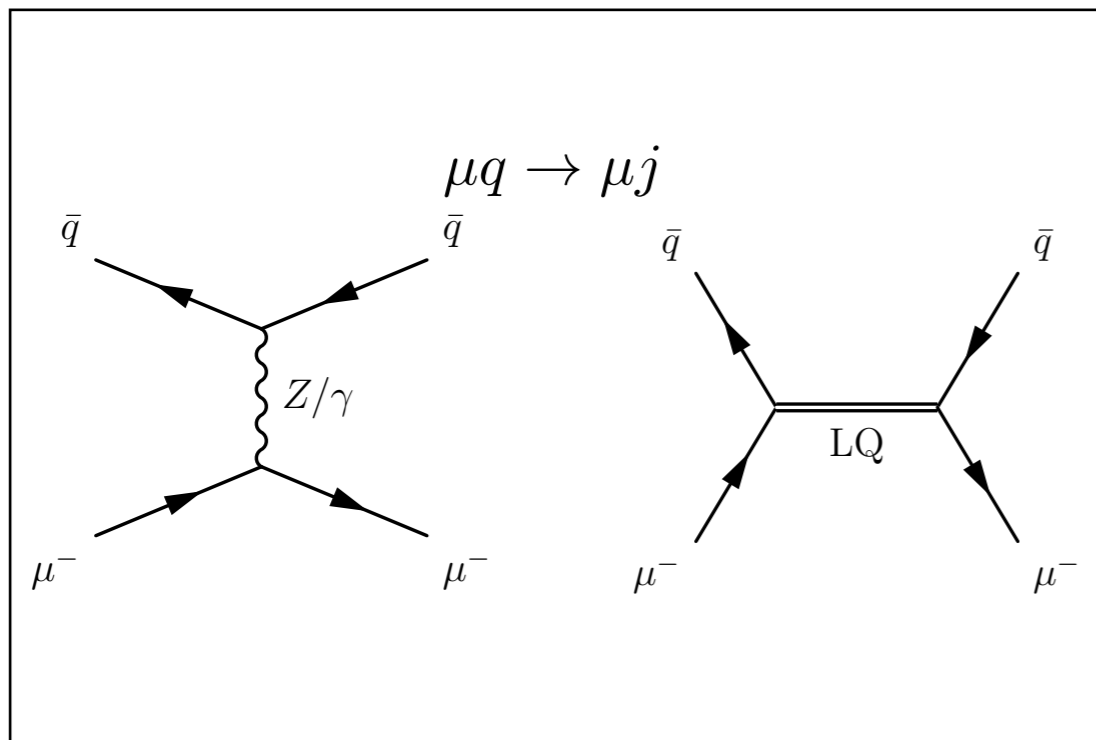
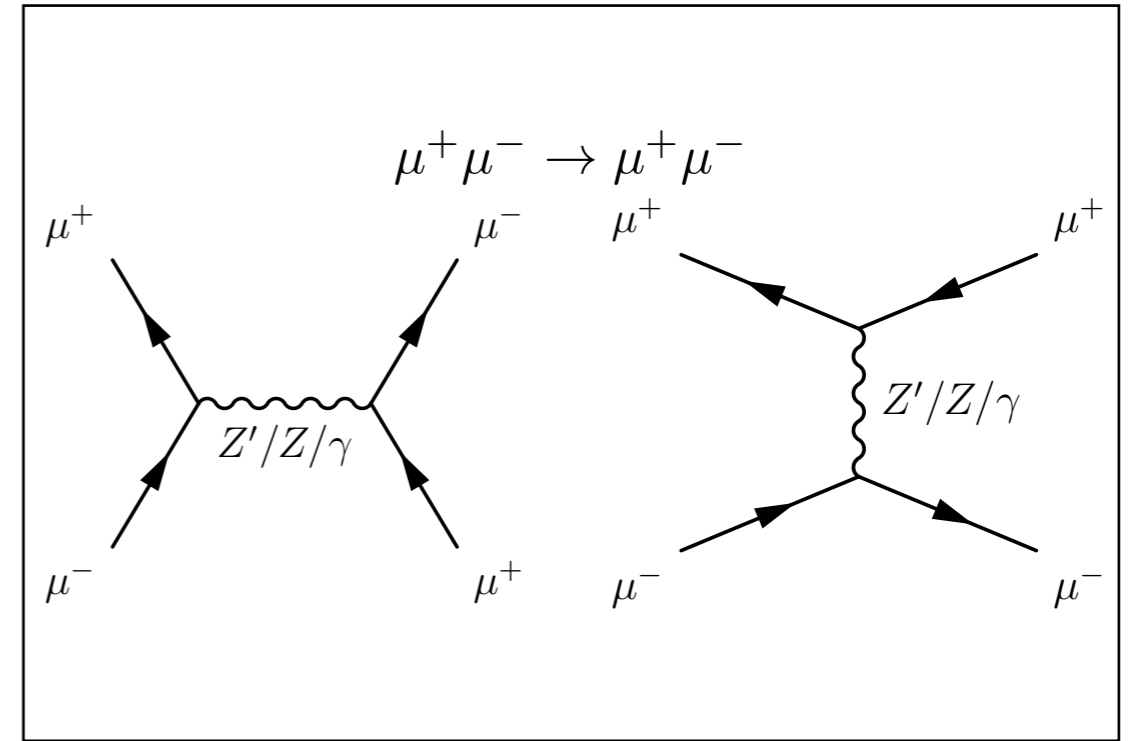
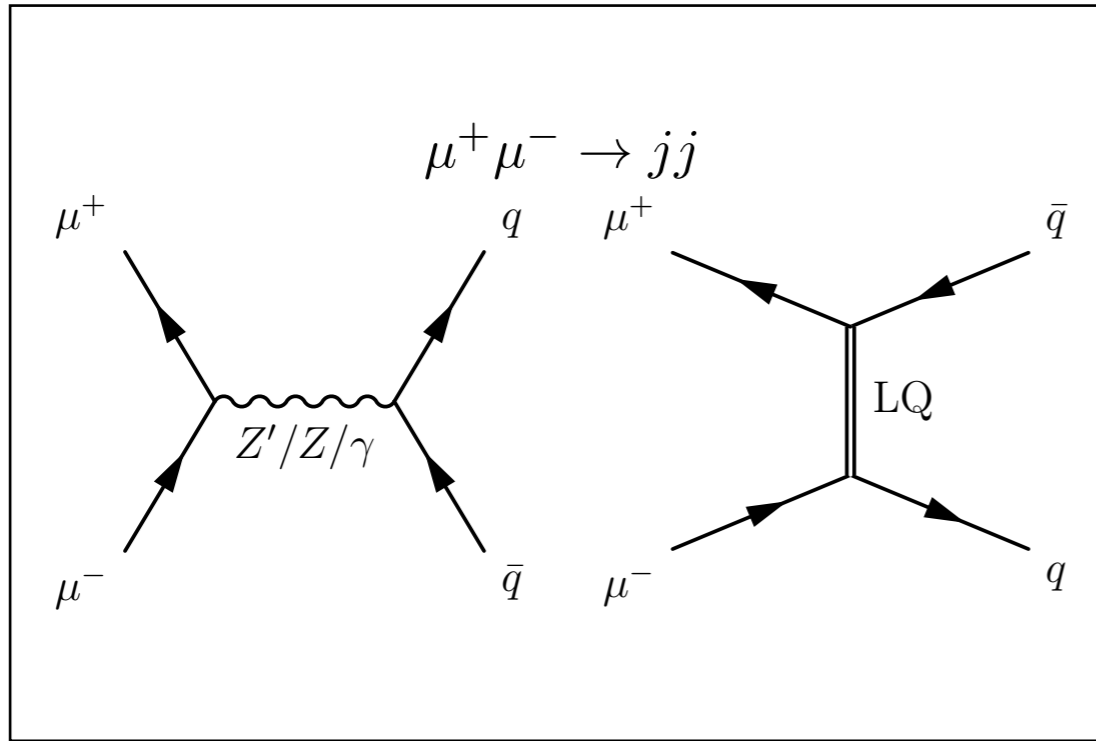
$$m^2 = (p_i + p_j)^2$$

$$\tau = m^2/s_0$$

Azatov, Garosi, AG, Marzocca, Salko, Trifinopoulos; [2205.13552](#)



# The signatures at MuC



# The signatures at MuC

$$m_X < \sqrt{s_0}$$

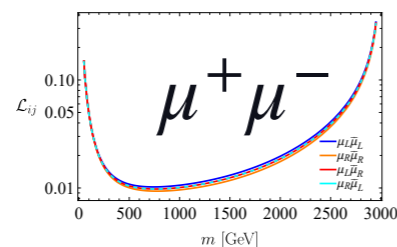
$$m_X > \sqrt{s_0}$$

- Kinematical features at  $m_{\mu\mu} \sim m_X$   
e.g. a resonance peak
- Corrections to the bins  $m_{\mu\mu} \approx \sqrt{s_0}$   
“fifth force searches”

- Corrections to the bins  $m_{\mu\mu} \approx \sqrt{s_0}$   
“EFT searches”

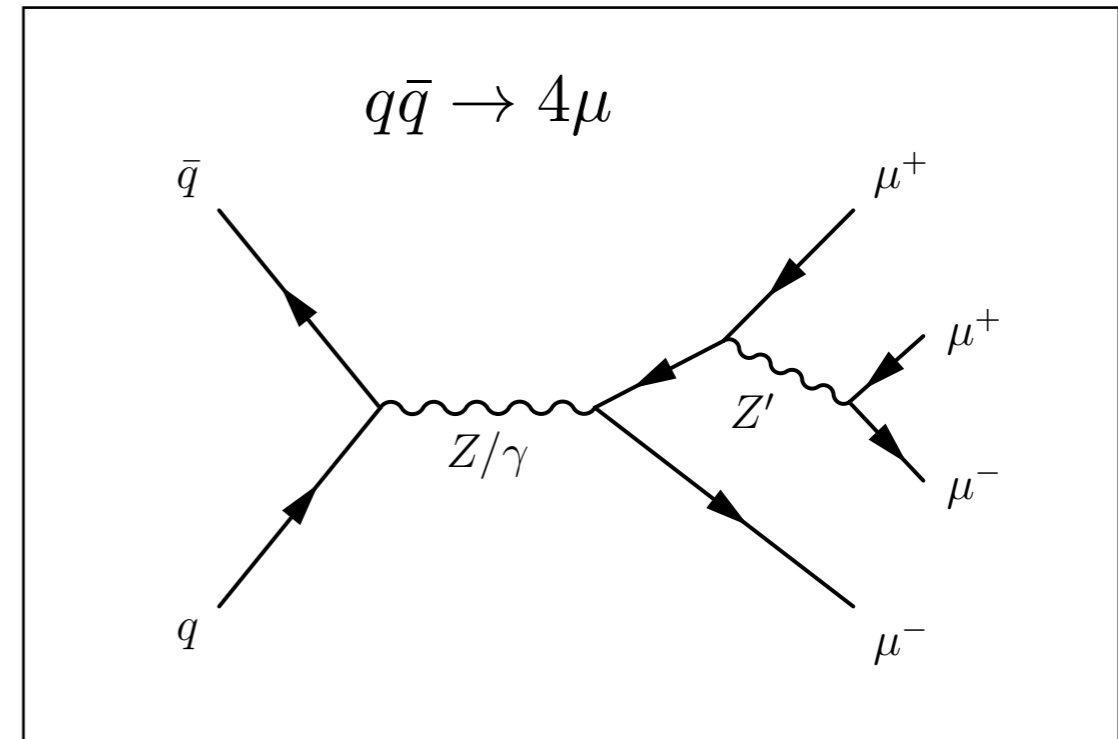
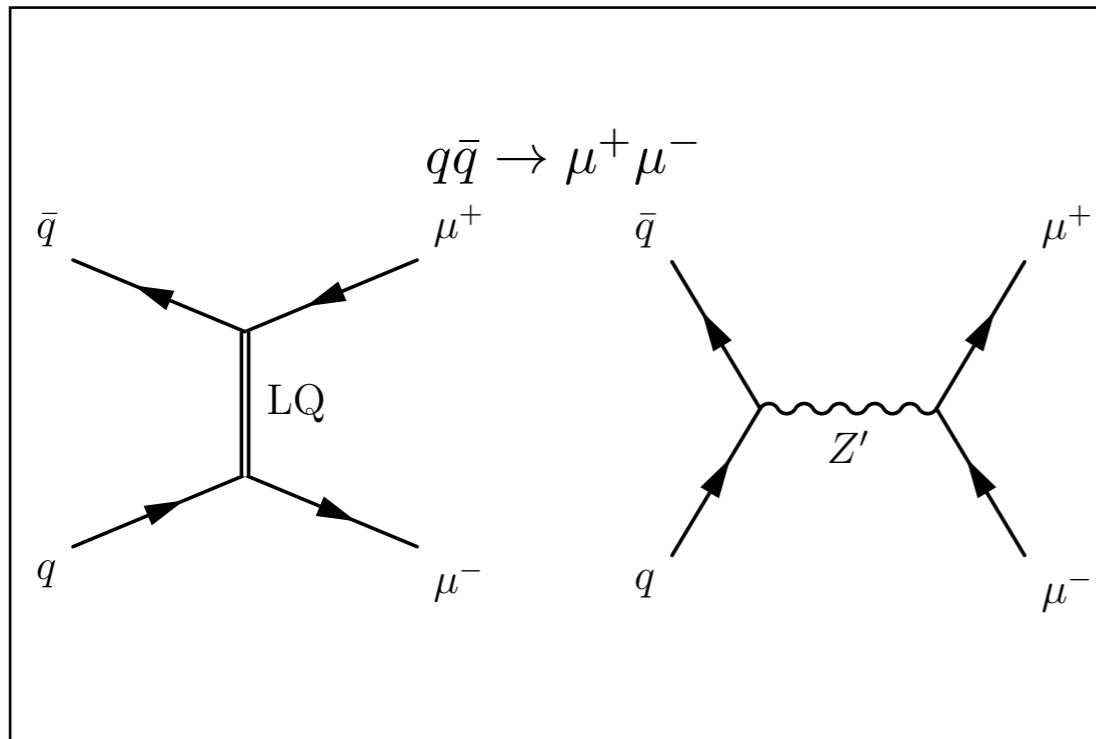


- Only effective  
at MuC due to



- Monotonously decreasing  
luminosities in proton colliders

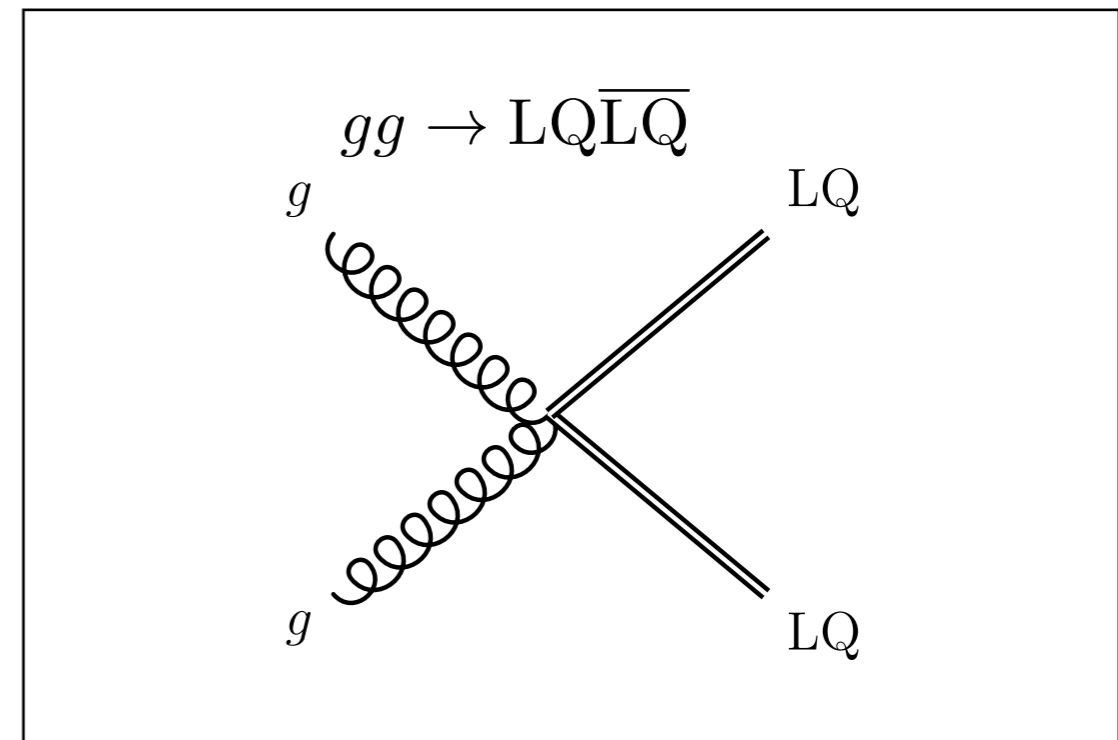
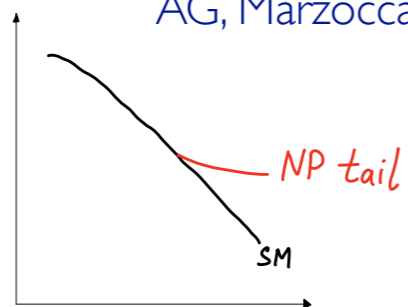
# The signatures at hadron colliders



- Corrections to the high-mass Drell-Yan
- LHC data already useful to constrain models for  $b \rightarrow s\mu\mu$

AG, Marzocca; 1704.09015

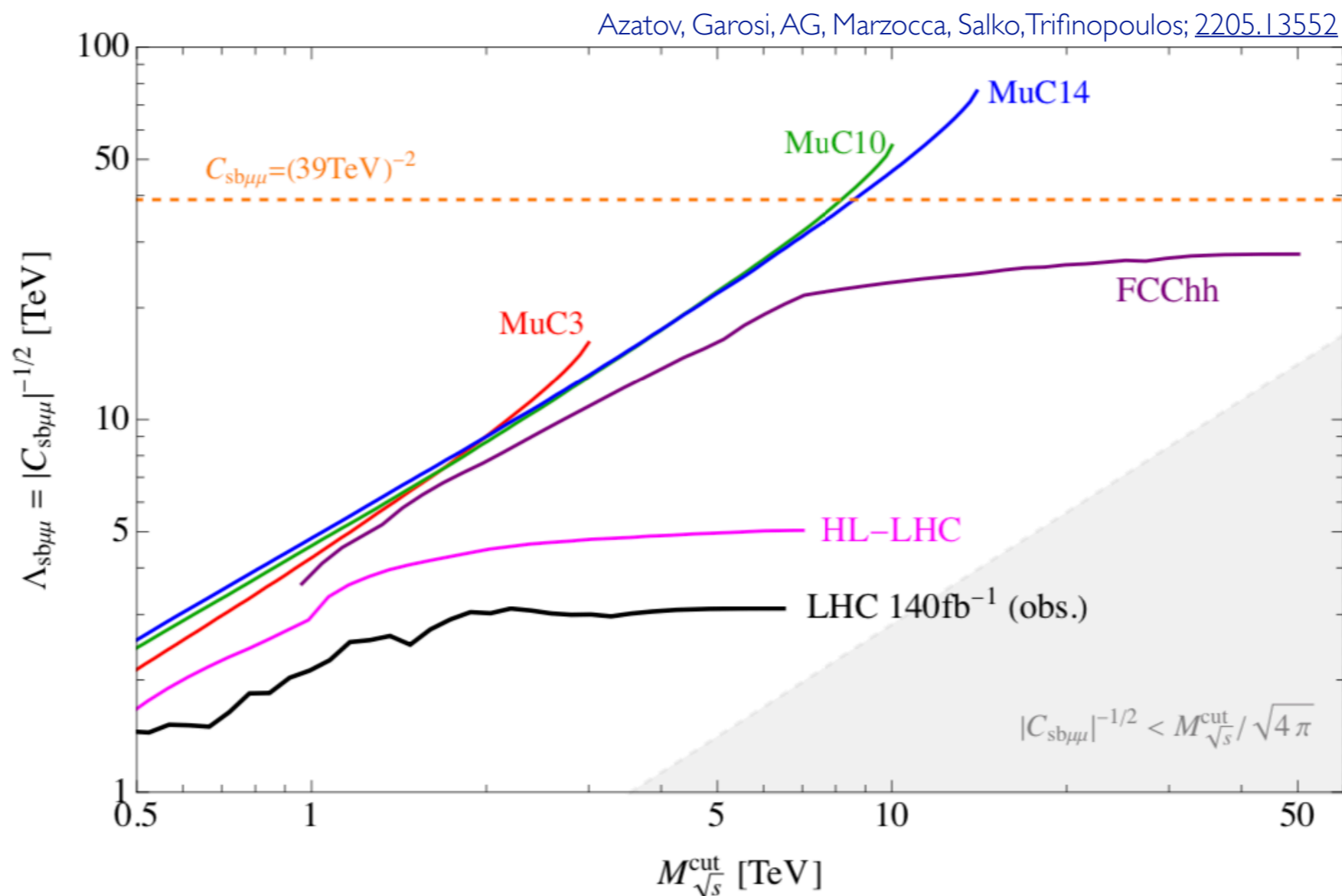
EFT limit



# ***Results***

# Contact interactions

The pessimistic case

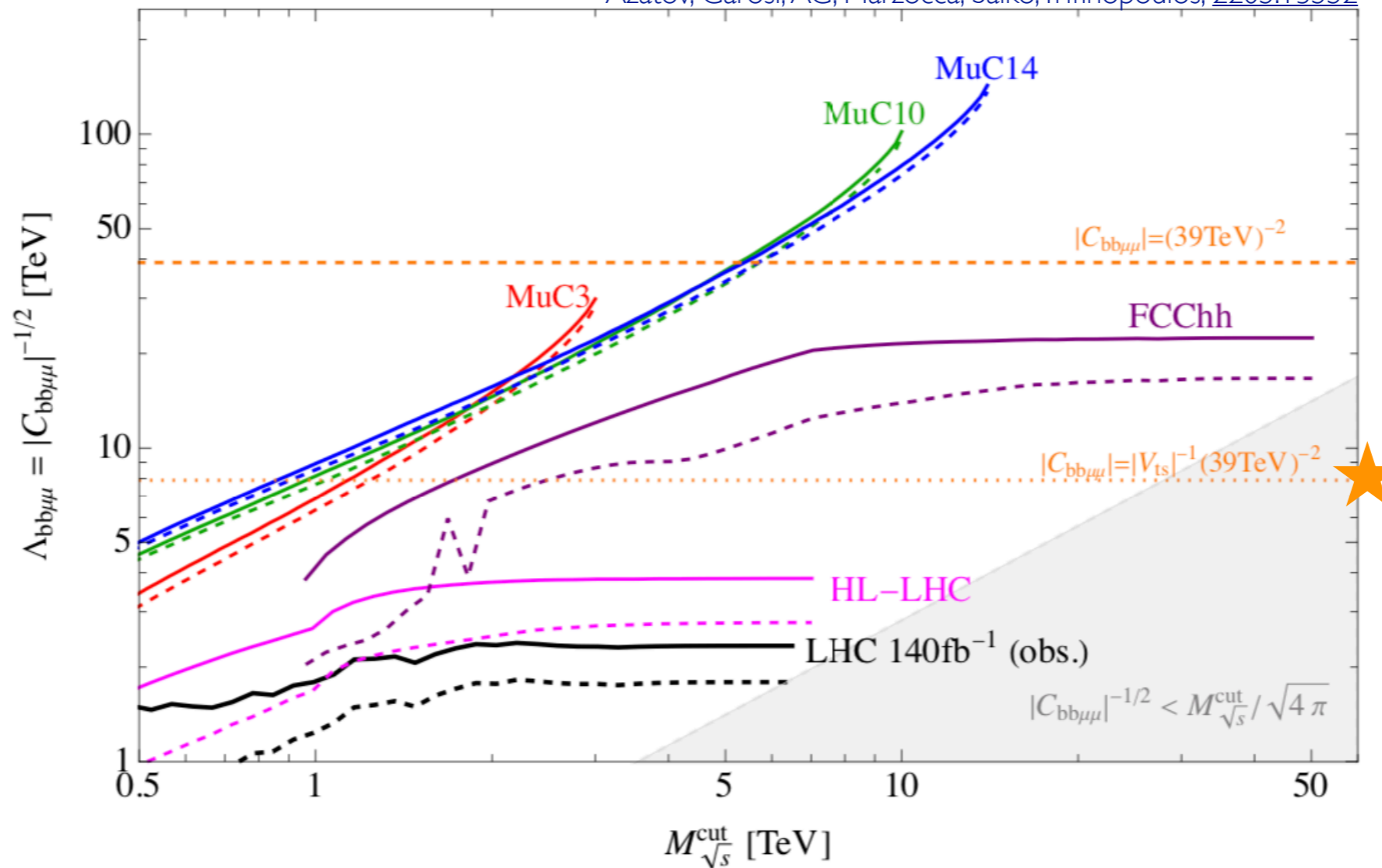


**Figure 7.** Sensitivity reach (95%CL) for the  $(\bar{s}_L \gamma_\alpha b_L)(\bar{\mu}_L \gamma^\alpha \mu_L)$  contact interaction as function of the upper cut on the final-state invariant mass, compared to the value required to fit  $bs\mu\mu$  anomalies (dashed orange line).

# Contact interactions

The realistic case ★

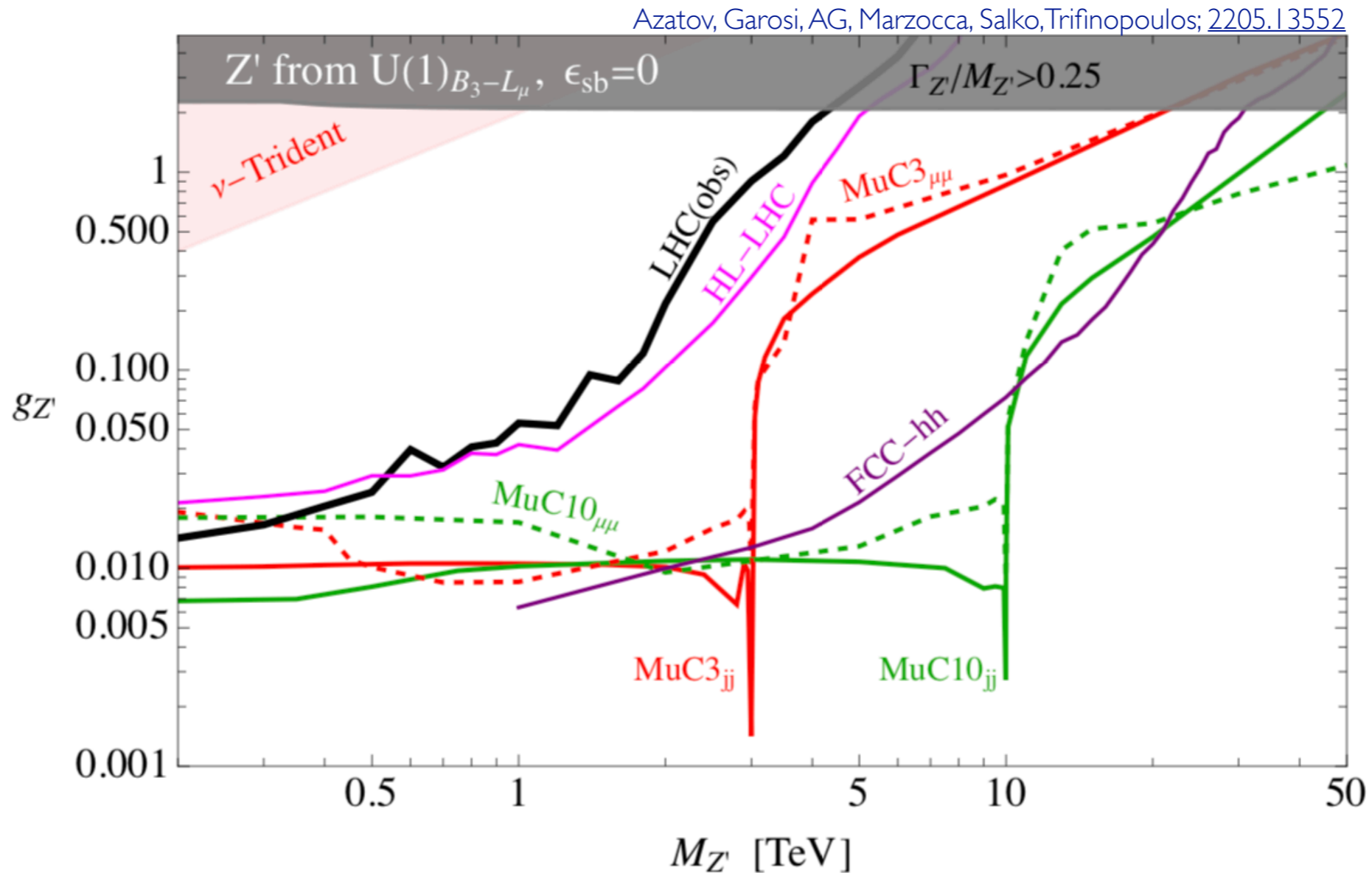
Azatov, Garosi, AG, Marzocca, Salko, Trifinopoulos; [2205.13552](#)



**Figure 8.** Sensitivity reach (95%CL) for the  $(\bar{b}_L \gamma_\alpha b_L)(\bar{\mu}_L \gamma^\alpha \mu_L)$  contact interaction as function of the upper cut on the final-state invariant mass. Solid (dashed) lines represent the limit for positive (negative) values of  $C_{bb\mu\mu}$ . The orange dotted and dashed lines shows reference values in relation to the  $bs\mu\mu$  anomalies fit, with or without a  $1/V_{ts}$  enhancement of the  $bb$  operator compared to the  $bs$  one, respectively.

# $Z'$ models: $B_3 - L_\mu$

$$\mathcal{L}_{Z'_{B_3-L_\mu}}^{\text{int}} = -g_{Z'} Z'_\alpha \left[ \frac{1}{3} \bar{Q}_L^3 \gamma^\alpha Q_L^3 + \frac{1}{3} \bar{b}_R \gamma^\alpha b_R + \frac{1}{3} \bar{t}_R \gamma^\alpha t_R - \bar{L}_L^2 \gamma^\alpha L_L^2 - \bar{\mu}_R \gamma^\alpha \mu_R + \left( \frac{1}{3} \epsilon_{sb} \bar{Q}_L^2 \gamma^\alpha Q_L^3 + \text{h.c.} \right) + \mathcal{O}(\epsilon_{sb}^2) \right]$$



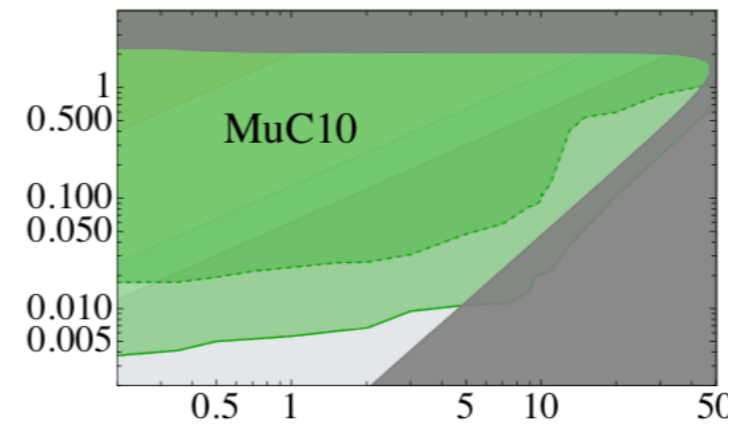
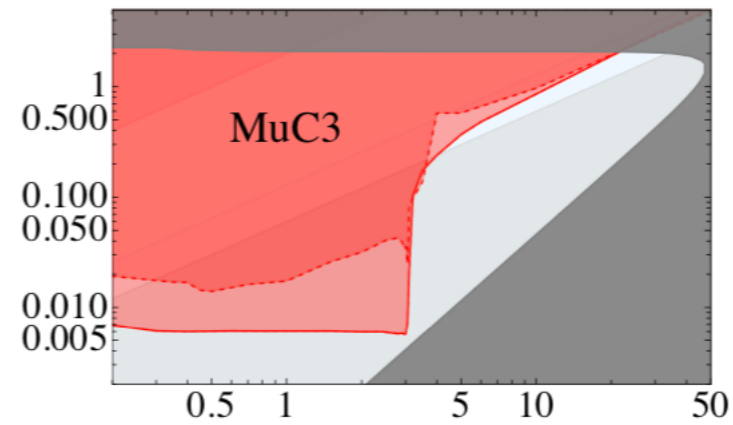
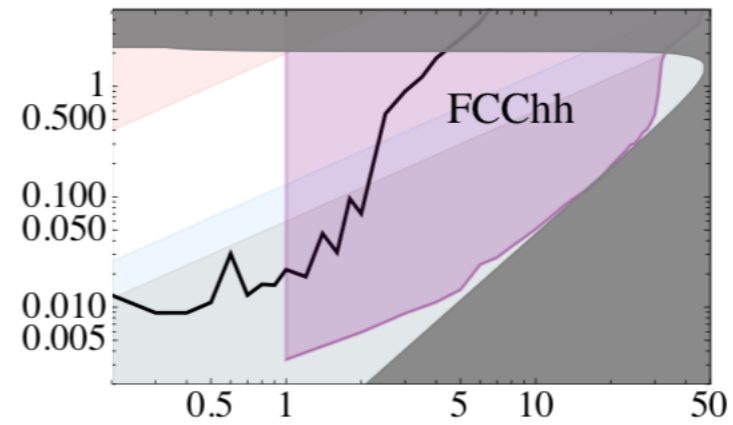
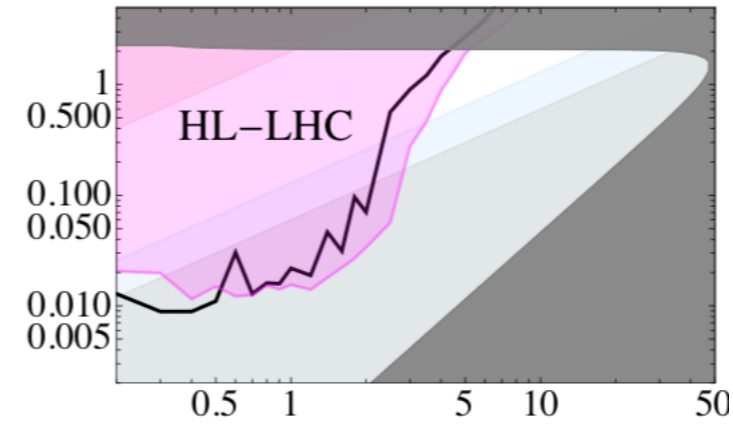
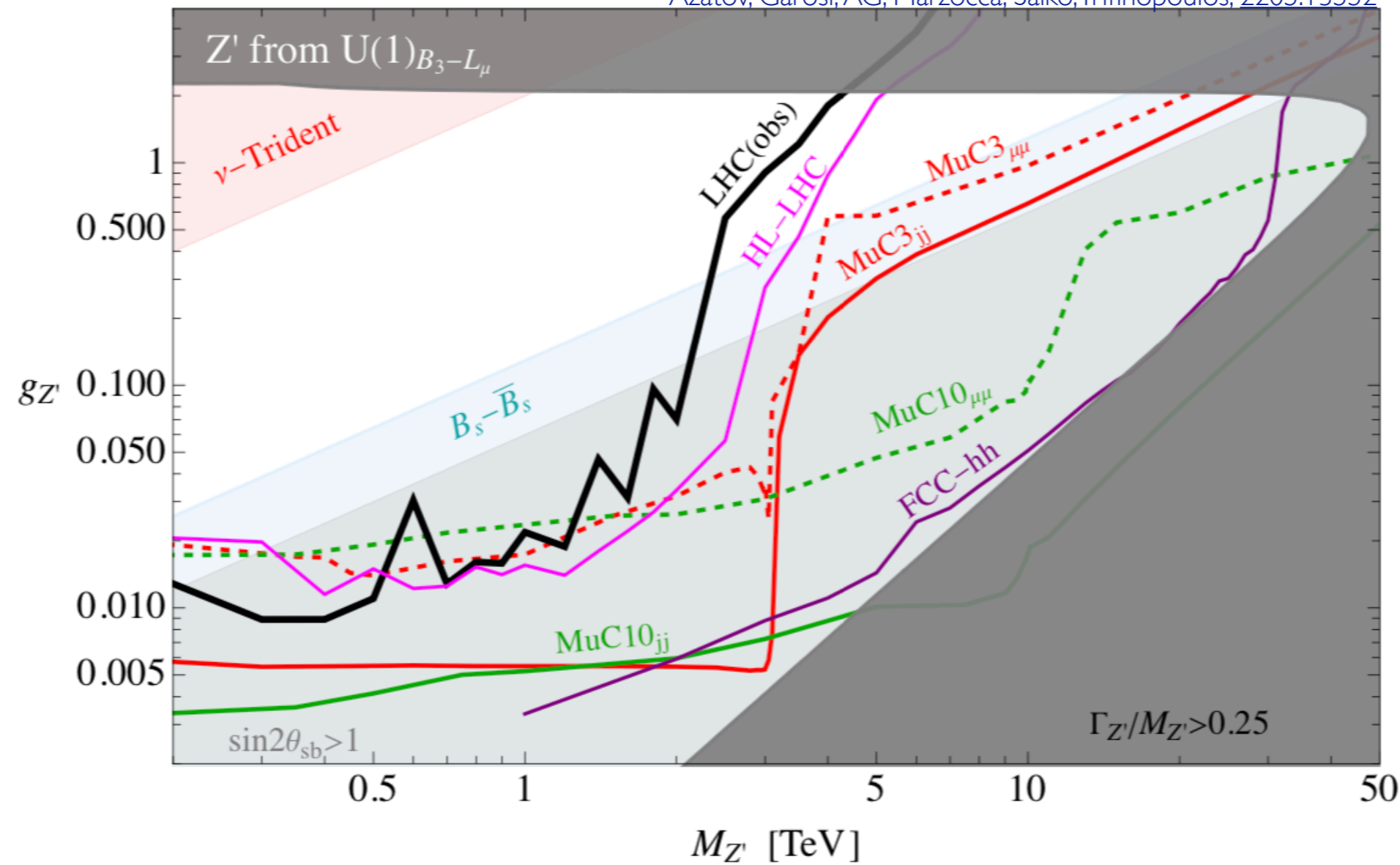
**Figure 9.** Discovery reach at  $5\sigma$  for the  $B_3 - L_\mu$  model with  $\epsilon_{sb} = 0$ , for different final states at each collider (as indicated by the labels). The region excluded at 95% CL by LHC [111] is above the black line while in the dark gray region the  $Z'$  has a large width, signaling a loss of perturbativity.



# $Z'$ models: $B_3 - L_\mu$

$$bs\mu\mu : \epsilon_{sb} = -1.7 \times 10^{-3} \left( \frac{M_{Z'}}{g_{Z'} \text{TeV}} \right)^2$$

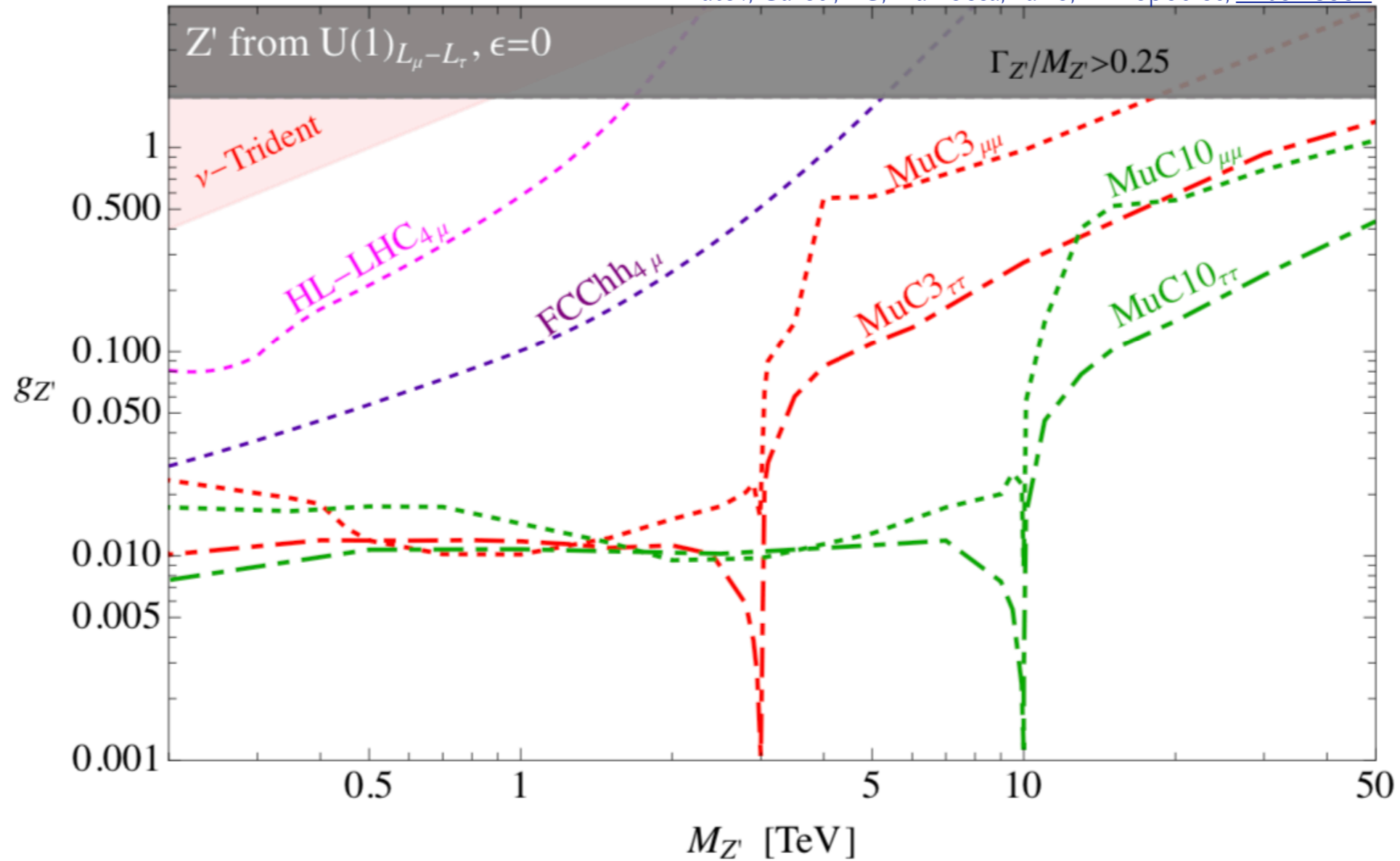
Azatov, Garosi, AG, Marzocca, Salko, Trifinopoulos; 2205.13552



# $Z'$ models: $L_\mu - L_\tau$

$$\mathcal{L}_{Z'L_\mu-L_\tau}^{\text{int}} = -g_{Z'} Z'_\alpha \left[ \bar{L}_L^2 \gamma^\alpha L_L^2 + \bar{\mu}_R \gamma^\alpha \mu_R - \bar{L}_L^3 \gamma^\alpha L_L^3 - \bar{\tau}_R \gamma^\alpha \tau_R + \right. \\ \left. + |\epsilon_b|^2 \bar{Q}_L^3 \gamma^\alpha Q_L^3 + |\epsilon_s|^2 \bar{Q}_L^2 \gamma^\alpha Q_L^2 + (\epsilon_b \epsilon_s^* \bar{Q}_L^2 \gamma^\alpha Q_L^3 + \text{h.c.}) + \dots \right]$$

Azatov, Garosi, AG, Marzocca, Salko, Trifinopoulos; [2205.13552](#)

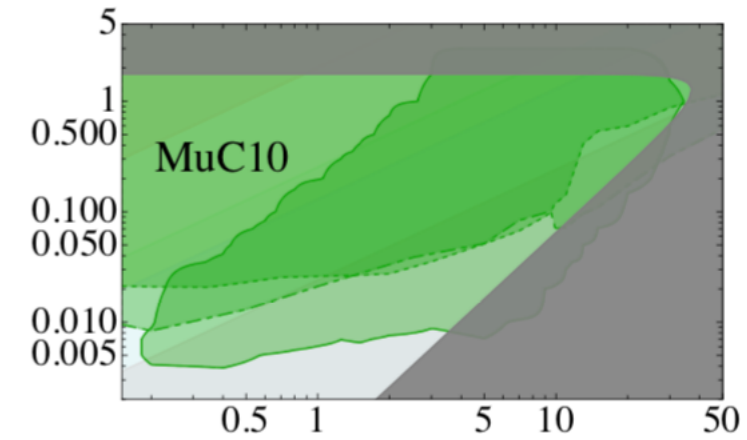
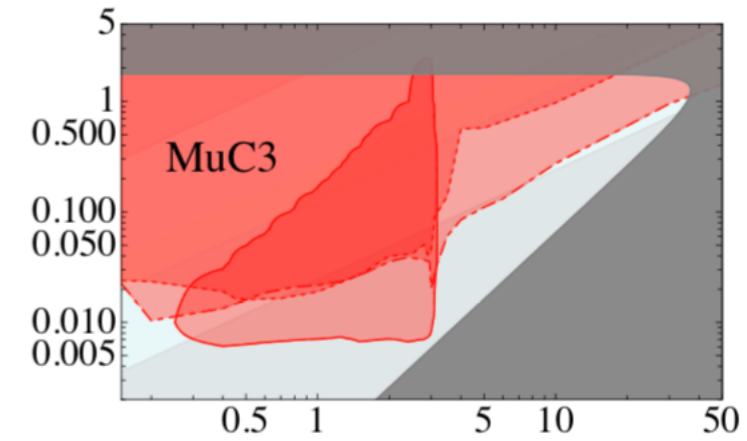
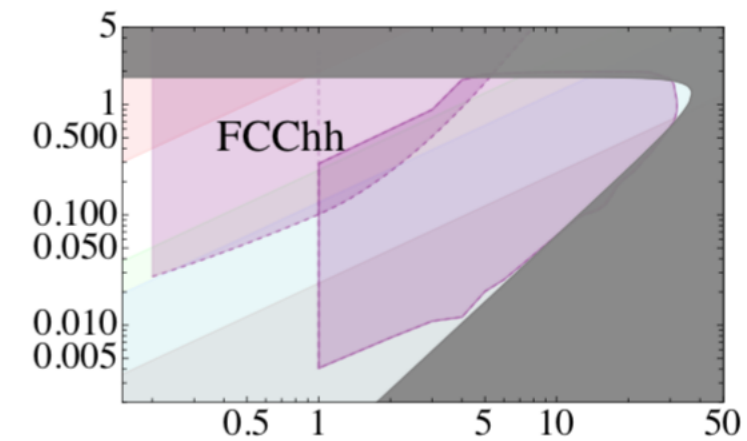
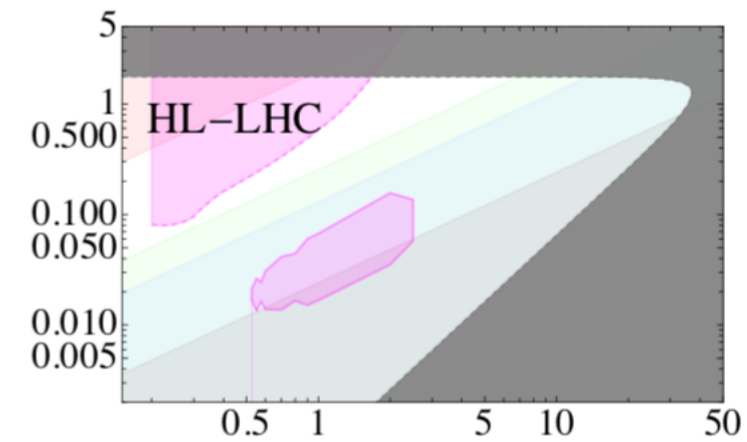
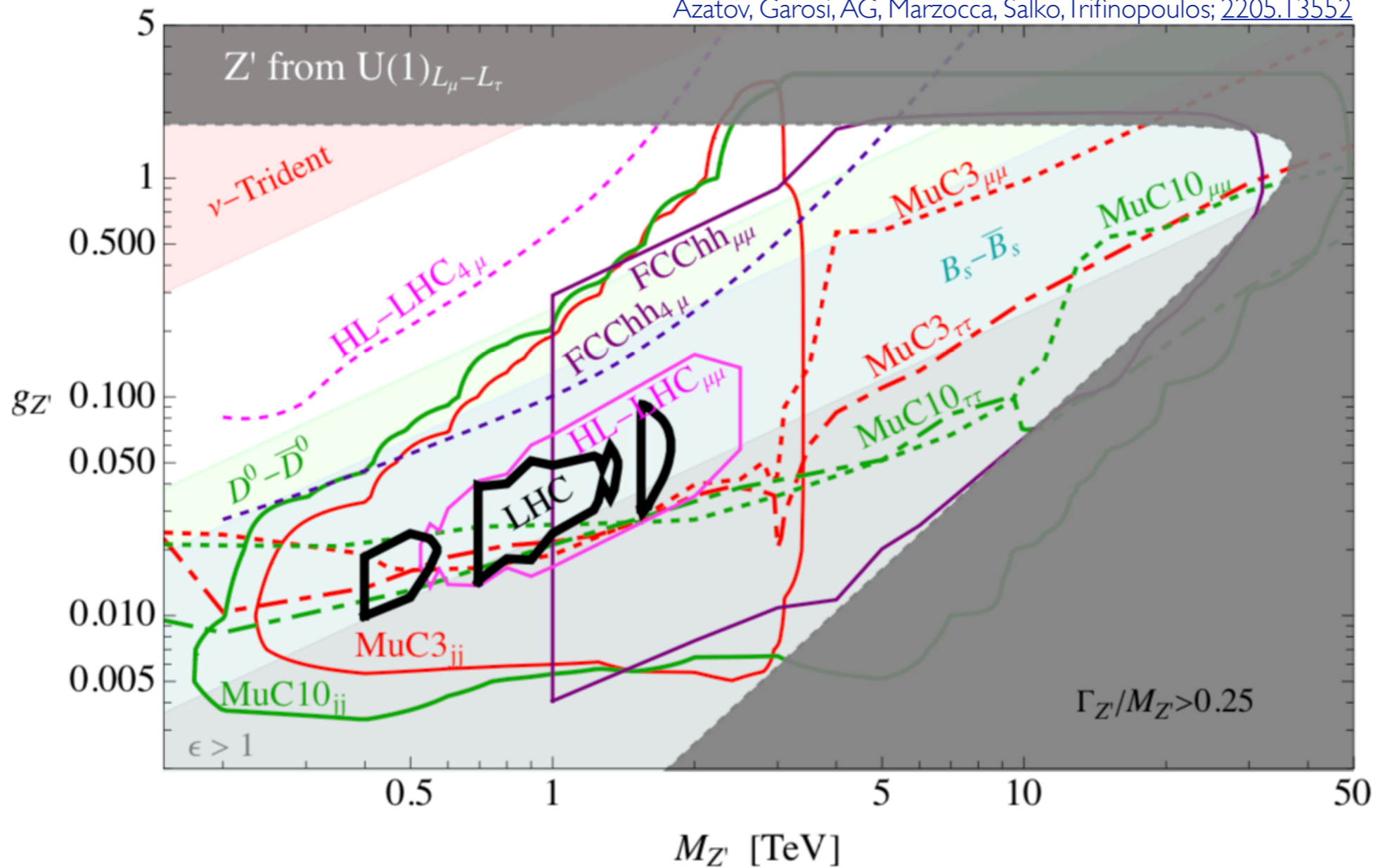


**Figure 11.** Discovery reach at  $5\sigma$  for the  $L_\mu - L_\tau$  model with  $\epsilon_s = \epsilon_b = 0$  in Eq. (5.6). In the dark gray region the  $Z'$  has a large width, signaling a loss of perturbativity.

# Z' models: $L_\mu - L_\tau$

$bs\mu\mu : \epsilon_b\epsilon_s^* = -5.7 \times 10^{-4} \left( \frac{M_{Z'}}{g_{Z'} \text{TeV}} \right)^2$   
 e.g.  $\epsilon_b = -\epsilon_s$

Azatov, Garosi, AG, Marzocca, Salko, Trifinopoulos; 2205.13552



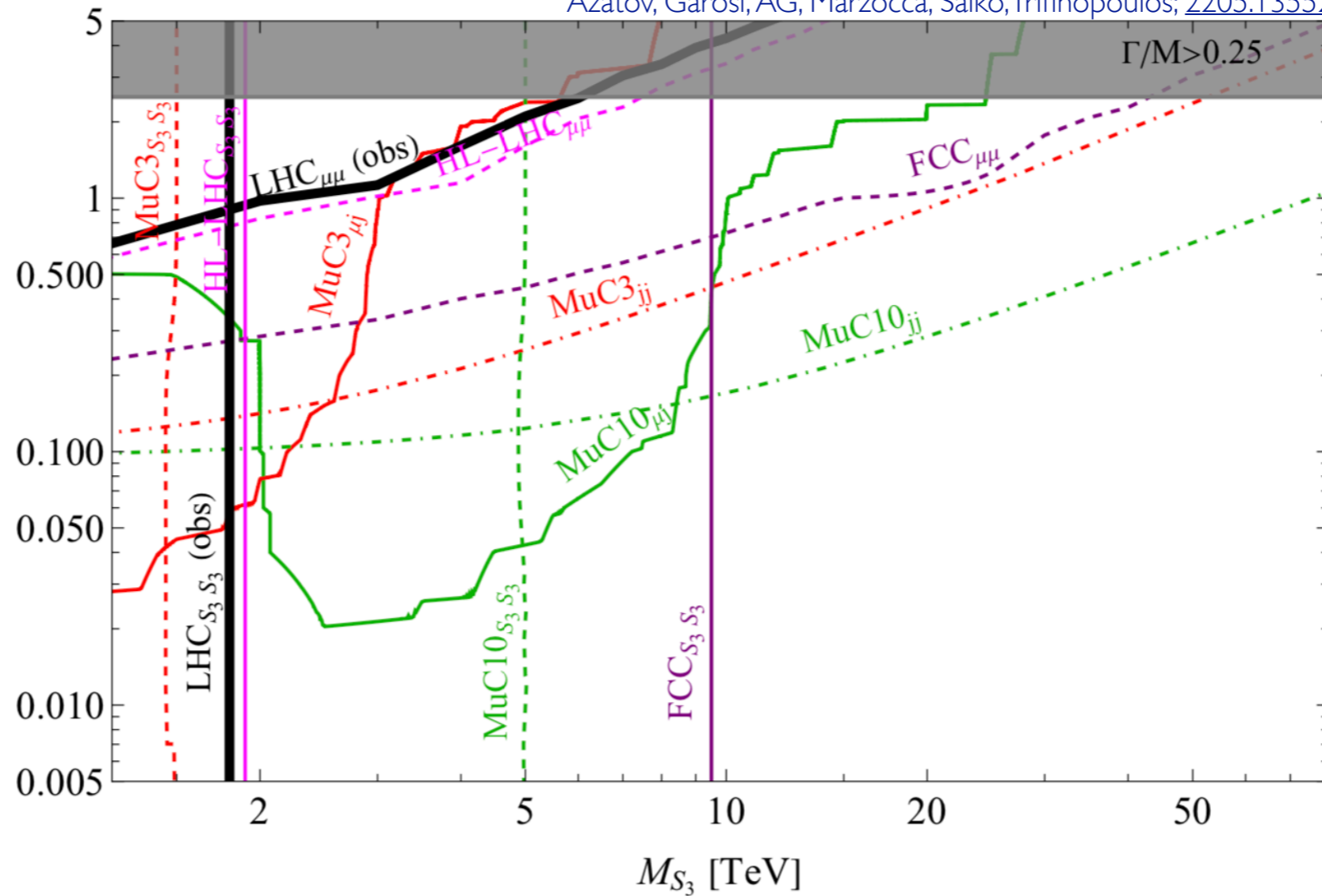
# Scalar Leptoquark

$$S_3 \sim (\bar{\mathbf{3}}, \mathbf{3}, 1/3)$$

$$\begin{aligned} \mathcal{L}_{S_3}^{\text{int}} &= \lambda_{i\mu} \overline{Q}_L^{ic} \epsilon \sigma^I L_L^2 S_3^I + \text{h.c.} , \\ &= -\lambda_{i\mu} S_3^{(1/3)} (V_{ji}^* \overline{u}_L^{jc} \mu_L + \overline{d}_L^{ic} \nu_\mu) + \sqrt{2} \lambda_{i\mu} \left( V_{ji}^* S_3^{(-2/3)} \overline{u}_L^{jc} \nu_\mu - S_3^{(4/3)} \overline{d}_L^{ic} \mu_L \right) + \text{h.c.} \end{aligned}$$

Azatov, Garosi, AG, Marzocca, Salko, Trifinopoulos; [2205.13552](#)

The only coupling  $\rightarrow \lambda_{b\mu}$



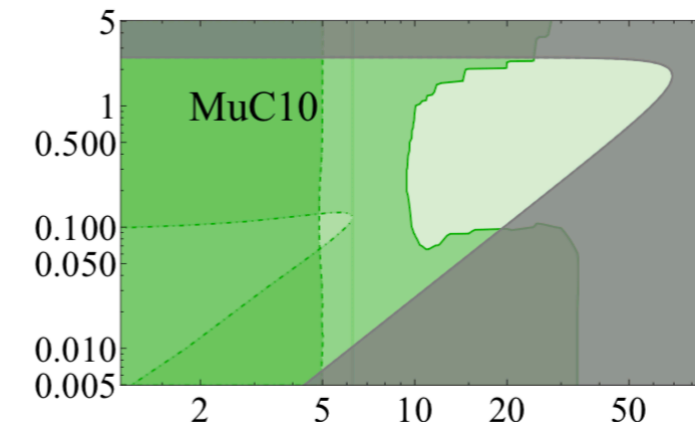
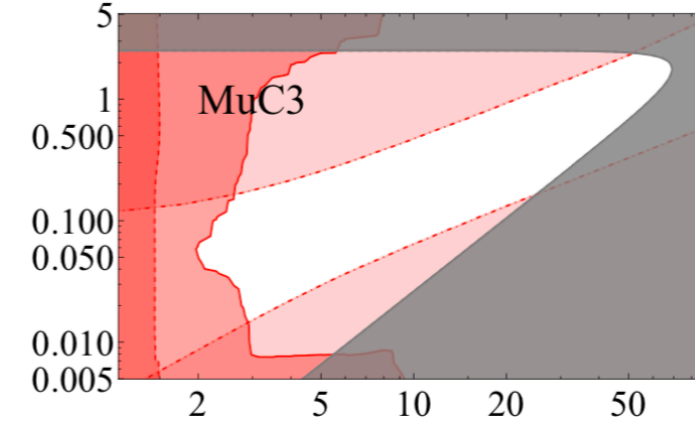
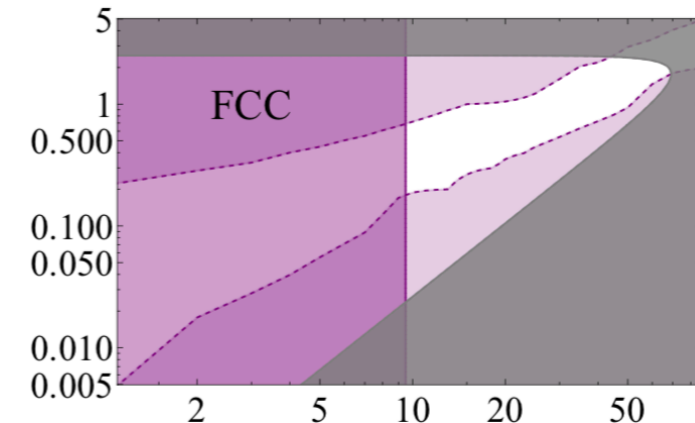
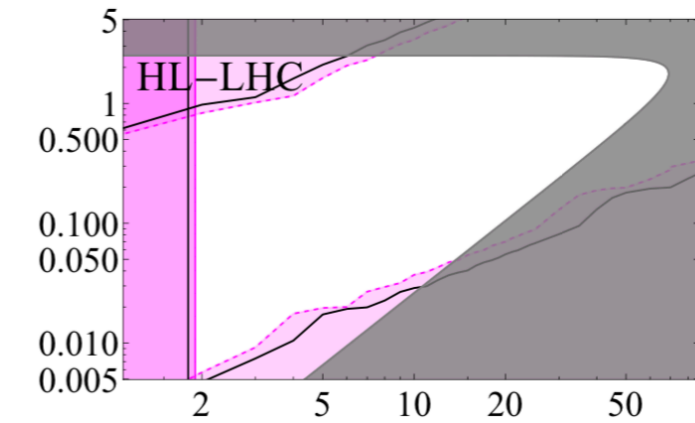
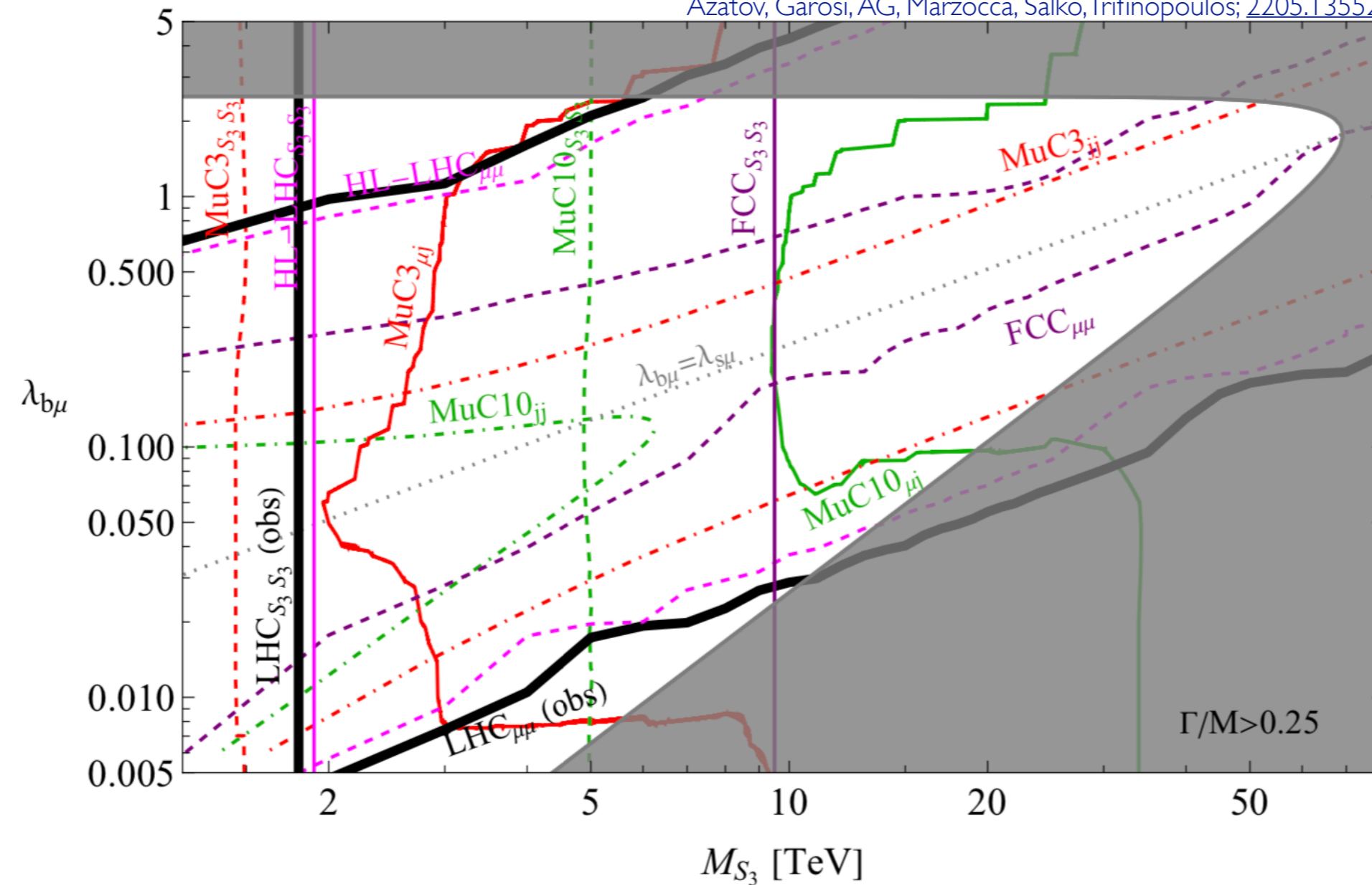
**Figure 13.** The  $5\sigma$  discovery prospects at future colliders for the  $S_3$  leptoquark assuming the  $U(2)^3$  quark flavour symmetry and the exclusive leptoquark coupling to muons (see Section 6.1).



# Scalar Leptoquark

■  $bs\mu\mu$  :  $\lambda_{b\mu}\lambda_{s\mu} = -8.4 \times 10^{-4} \left(\frac{M_{S_3}}{\text{TeV}}\right)^2$

Azatov, Garosi, AG, Marzocca, Salko, Trifinopoulos; 2205.13552

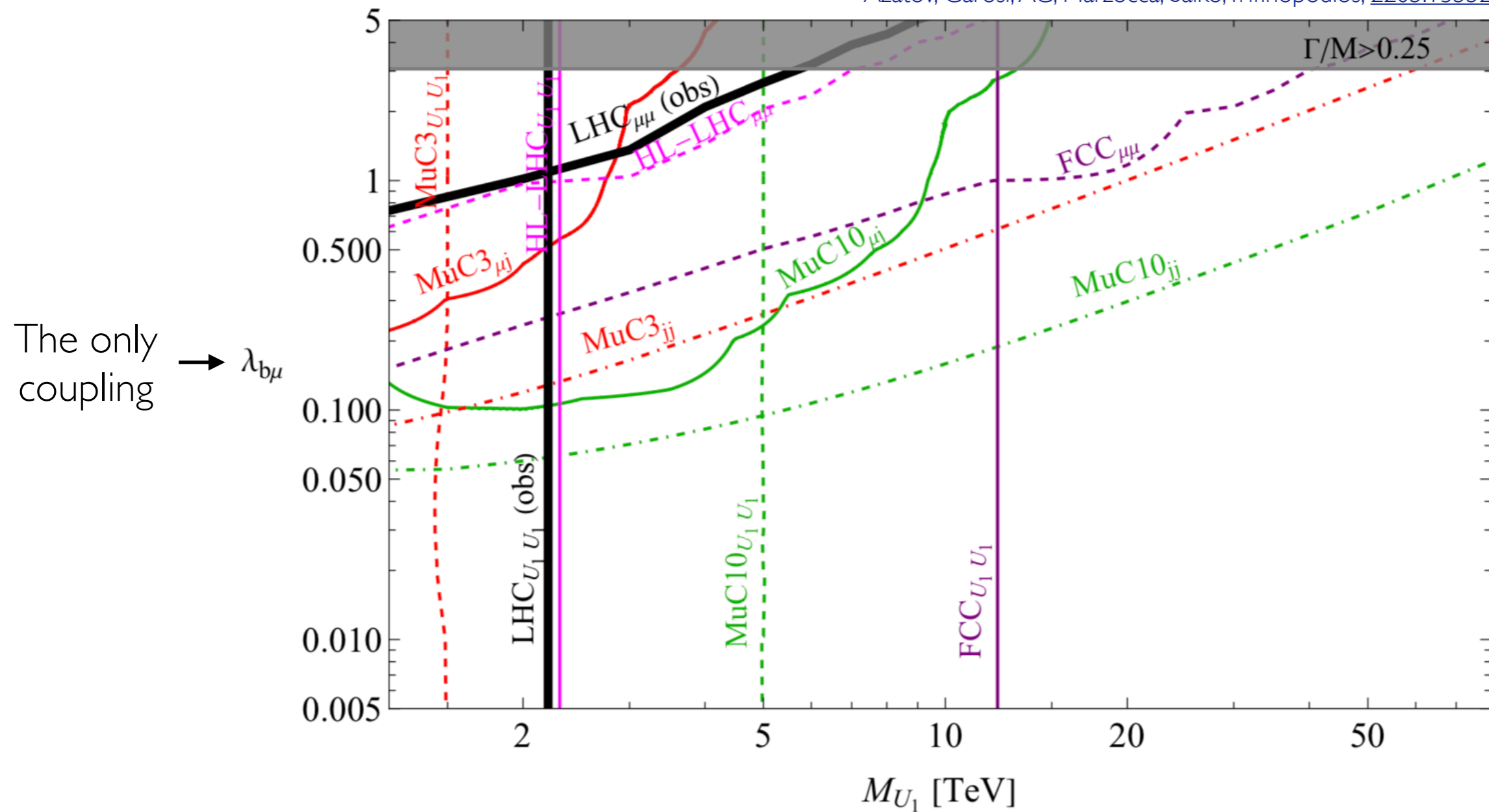


# Vector Leptoquark

$$U_1 \sim (\mathbf{3}, \mathbf{1}, 2/3)$$

$$\mathcal{L}_{U_1}^{\text{int}} = \lambda_{i\mu} \overline{Q}_L^i \gamma_\alpha L_L^2 U_1^\alpha + \text{h.c.} = \lambda_{i\mu} U_1^\alpha \left( V_{ji} \bar{u}_L^j \gamma_\alpha \nu_\mu + \bar{d}_L^i \gamma_\alpha \mu_L \right) + \text{h.c.}$$

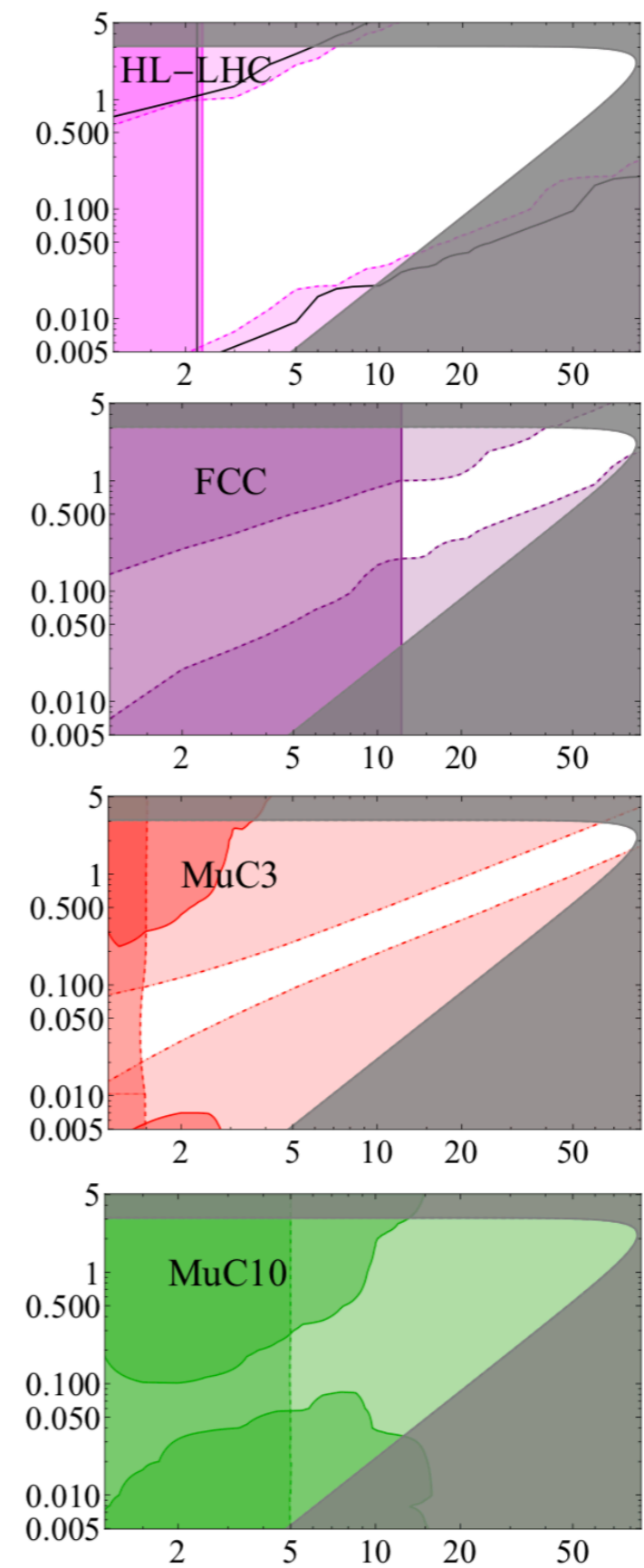
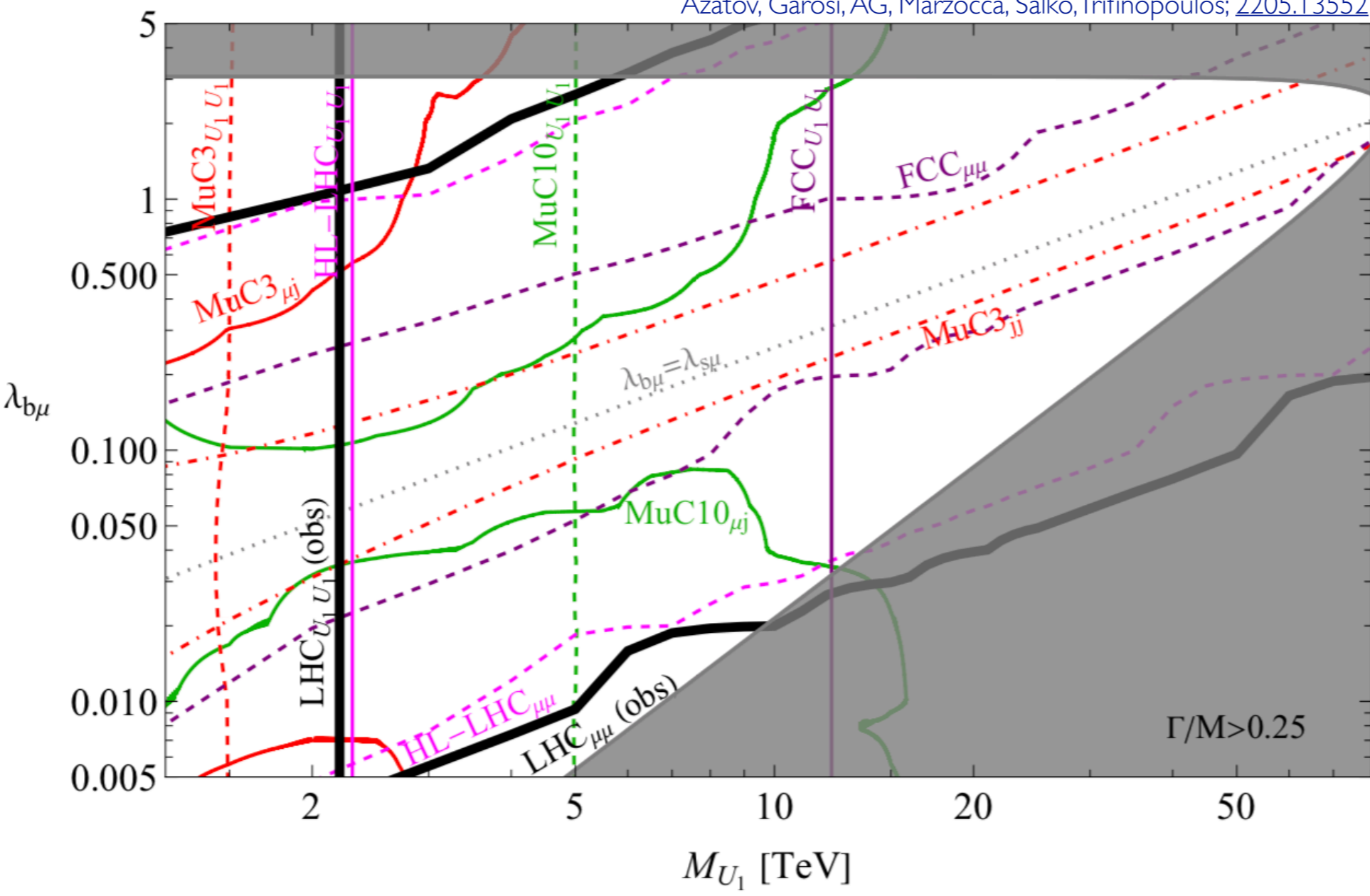
Azatov, Garosi, AG, Marzocca, Salko, Trifinopoulos; 2205.13552



# Vector Leptoquark

■  $bs\mu\mu : \lambda_{b\mu}\lambda_{s\mu} = -8.4 \times 10^{-4} \left(\frac{M_{U_1}}{\text{TeV}}\right)^2$

Azatov, Garosi, AG, Marzocca, Salko, Trifinopoulos; [2205.13552](#)



***Backup***



# Resonant Leptoquark at FCC-hh

