

Dark matter at the muon colliders

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Based on arXiv:[22xx.xxxxx](#)

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- Dark matter is firmly established based on astrophysics observations and cosmology.
- Beyond the standard model
- Various candidates
 - Weakly interacting massive particles(WIMPs)
 - Axions
 - Sterile neutrinos
 - ...

Detecting DM

- Direct detection: $DM+SM \rightarrow DM+SM$,
- Indirect detection: $DM+DM \rightarrow SM+SM$
- Collider experiments: $SM+SM \rightarrow DM+SM$

Real scalar or Majorana fermion:

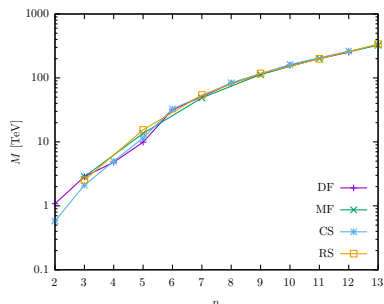
- Additional $SU(2)$ weak n -plet, which can be real scalar or Majorana fermion

$$\mathcal{L}_S = \frac{1}{2}(D_\mu\chi)^2 - \frac{1}{2}M_\chi^2\chi^2 \quad (1)$$

$$\mathcal{L}_F = \frac{1}{2}\chi(i\bar{\sigma}^\mu D_\mu - M_\chi)\chi \quad (2)$$

- Electroweak corrections introduce mass splitting
- The neutral one is also the lightest, hence DM candidates

Thermal masses including Sommerfeld enhancement and bound-state formation:



Bottaro et.al.2107.09688&2205.04486

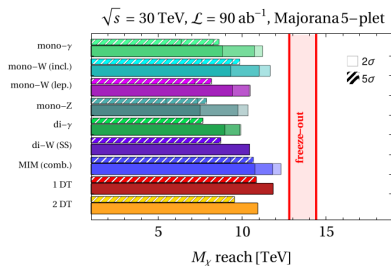
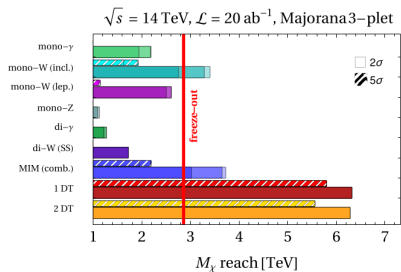
- MF and RS: odd n , $Y = 0$
- DF and CS: both even and odd n , non-zero Y
- Masses are determined by n .
- $n = 4, 5$ is at $\mathcal{O}(10)$ TeV, beyond the reach of LHC!

direct production at the muon collider

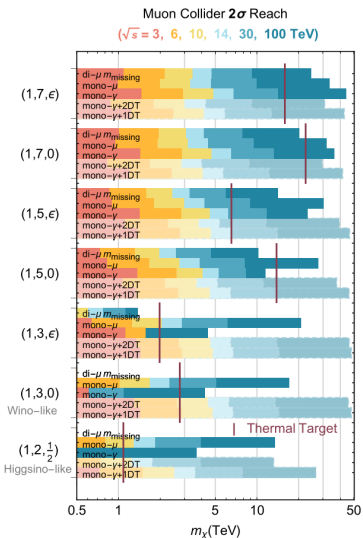
- mono- V : $\mu^+\mu^- \rightarrow \chi^i\chi^j + V, V = \gamma, Z, W$
- mono- μ : $\mu^+\mu^- \rightarrow \chi^i\chi^j + \mu^\pm\nu, \gamma\mu^\pm \rightarrow \chi^i\chi^j + \mu^\pm$
- VBF di- μ : $\mu^+\mu^- \rightarrow \chi^i\chi^j + \mu^+\mu^-$
- di- V : $\mu^+\mu^- \rightarrow \chi^i\chi^j + VV$
- Disappearing tracks due to decay of $\chi^i\chi^j$

Main limitation: $M_\chi < \frac{\sqrt{s}}{2}$

direct production at muon colliders



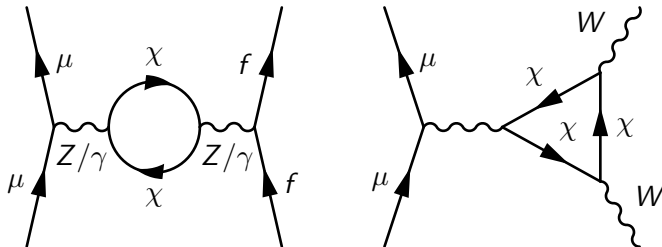
Bottaro et.al.2107.09688



Han,Liu,Wang,Wang,2009.11287

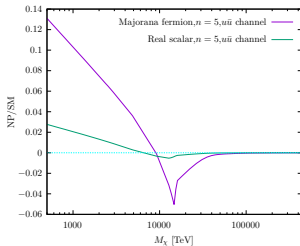
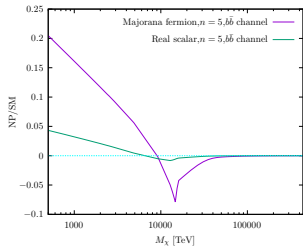
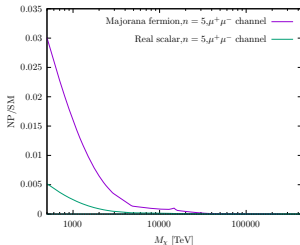
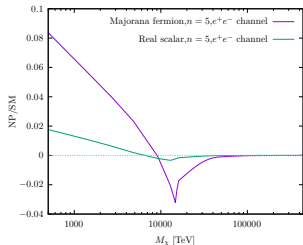
Indirect probes

- Direct production is limited by $M < \frac{\sqrt{s}}{2}$.
- Indirect probes through loop corrections: no such limit!



Neutral current contribution

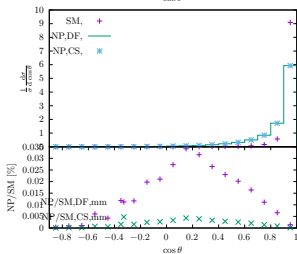
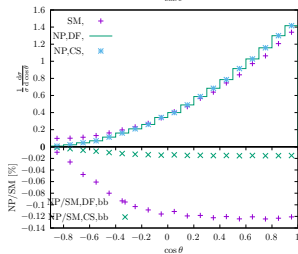
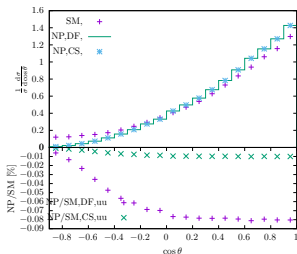
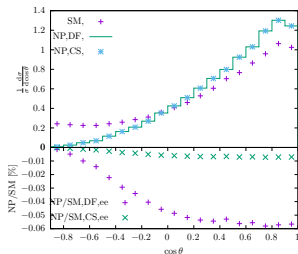
30 TeV muon collider:



- Effects of fermions are much larger than scalars
- Large effects for low mass
- small effects for large mass (EFT regime)
- local peak for the threshold

$$M_\chi = \frac{\sqrt{s}}{2}$$

Differential distribution



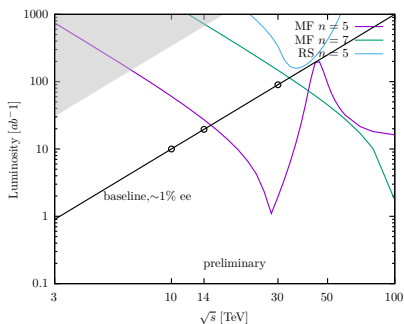
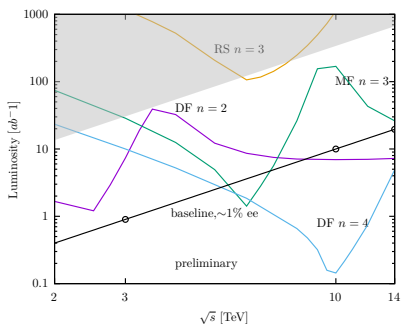
$e^+e^-, u\bar{u}, b\bar{b}$

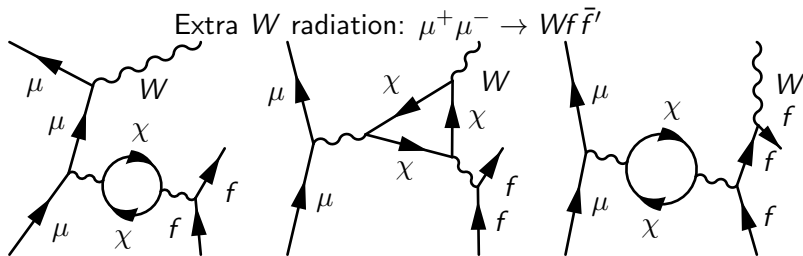
- peaks in forward region
- θ coverage of detectors
- Inclusive

$\mu^+\mu^-$:

- SM peak in forward region
- NP effects peak in center region
- Need differential

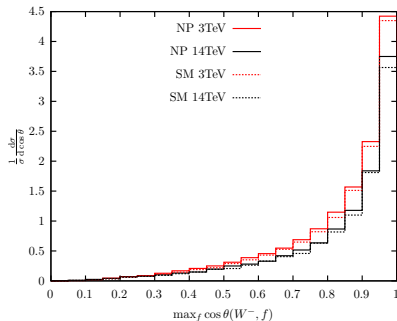
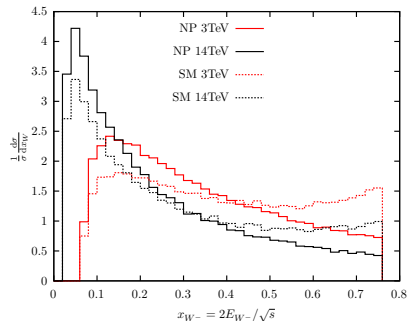
- Inclusive level: e^+e^- , $\tau^+\tau^-$, jj , $t\bar{t}$, HZ , W^+W^-
- Differential level: $\mu^+\mu^-$





Differential distributions

$$\mu^+ \mu^- \rightarrow W^- u \bar{d}$$



Dominated by soft and collinear W bosons.

Soft and collinear radiation

Double logarithm $\ln^2 \frac{s}{m_W^2}$ and single logarithm are large: NC at fixed order \rightarrow NC+CC with resummation

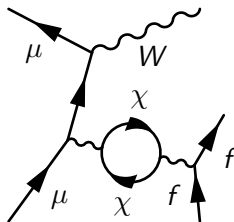
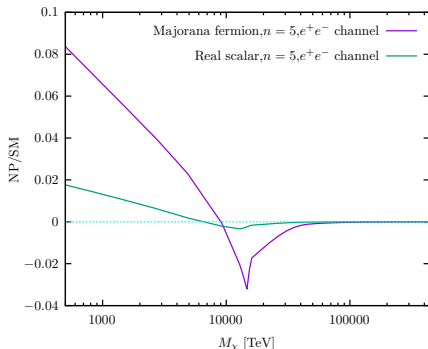
	3 TeV			10 TeV			30 TeV		
	DL	$e^{\text{DL}-1}$	$\text{SL}(\frac{\pi}{2})$	DL	$e^{\text{DL}-1}$	$\text{SL}(\frac{\pi}{2})$	DL	$e^{\text{DL}-1}$	$\text{SL}(\frac{\pi}{2})$
$\ell_L \rightarrow \ell'_L$	-0.46	-0.37	0.25	-0.82	-0.56	0.33	-1.23	-0.71	0.41
$\ell_L \rightarrow q_L$	-0.44	-0.36	0.25	-0.78	-0.54	0.34	-1.18	-0.69	0.42
$\ell_L \rightarrow e_R$	-0.32	-0.27	0.13	-0.56	-0.43	0.17	-0.85	-0.57	0.21
$\ell_L \rightarrow u_R$	-0.27	-0.24	0.11	-0.48	-0.38	0.15	-0.72	-0.51	0.18
$\ell_L \rightarrow d_R$	-0.24	-0.21	0.10	-0.43	-0.35	0.13	-0.64	-0.47	0.16
$\ell_R \rightarrow \ell'_L$	-0.32	-0.27	0.13	-0.56	-0.43	0.17	-0.85	-0.57	0.21
$\ell_R \rightarrow q_L$	-0.30	-0.26	0.12	-0.53	-0.41	0.16	-0.79	-0.55	0.21
$\ell_R \rightarrow \ell'_R$	-0.17	-0.16	0.07	-0.30	-0.26	0.09	-0.46	-0.37	0.12
$\ell_R \rightarrow u_R$	-0.12	-0.12	0.05	-0.22	-0.20	0.07	-0.33	-0.28	0.08
$\ell_R \rightarrow d_R$	-0.09	-0.09	0.04	-0.17	-0.16	0.05	-0.25	-0.22	0.06

Table 1: Double and single logarithmic corrections to the exclusive processes $\ell^+ \ell^- \rightarrow \bar{f} f$. The single-logarithmic corrections are evaluated at $\theta_* = \pi/2$.

Chen et.al.2202.10509

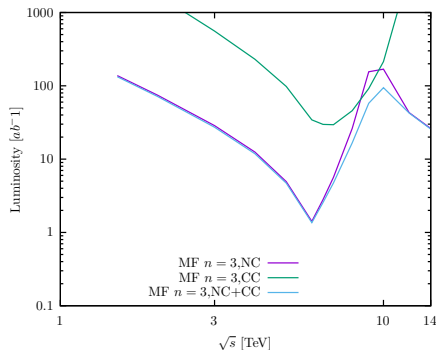
Hard radiation

- NP effects is about 0 when $\sqrt{s} \sim 3.2M_\chi$ for NC
- Initial state hard radiation can reduce $m(f\bar{f})$ from \sqrt{s} to $\sim 2M_\chi$ (threshold region)



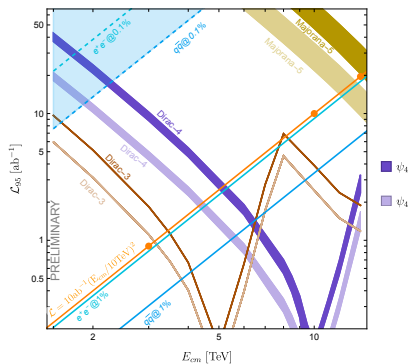
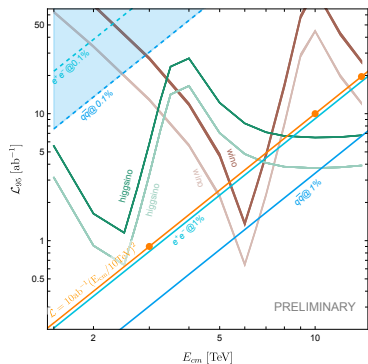
Combination of NC and CC

With hard W radiation:



- Neutral current provides most of the constraints
- Hard W radiation helps to resolve the region $\sqrt{s} \sim 3.2M_\chi$
- Need further improvements on theoretical prediction
- Need to combine with direct productions

Combined NC+CC



Conclusion

- mono- X , di- X and DT for low mass region $M_\chi < \frac{\sqrt{s}}{2}$
- Indirect probes are good at thresholds $M_\chi \sim \frac{\sqrt{s}}{2}$
- and can probe high mass $M_\chi > \frac{\sqrt{s}}{2}$
- Soft/collinear radiations shift NC to NC+CC
- Hard radiations affect the dynamics and sensitivities
- Statistic uncertainties in $\mathcal{O}(0.1 \sim 1\%)$ level: need further improvements on theoretical predictions(NLO+NLL or higher?)

Backup slides