



Radiation studies for the target area

Daniele Calzolari SY-STI-BMI (CERN), with inputs from A. Lechner, C. Rogers, R. Franqueira Ximenes, F. J. Saura Esteban and Marco Calviani



Outline

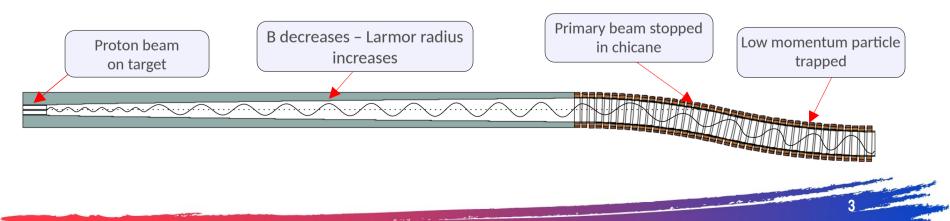
Introduction

- General target overview
- Target options for the radiation studies: graphite and liquid lead
- Radiation load to the target
 - Energy deposition and long-term radiation damage in the target
- Radiation load to the target solenoid
 - Heat load and long-term radiation damage in the superconducting coils
 - Total power to the cold mass
- Adiabatic tapering and chicane (only for graphite target)
 - Magnetic field modeling
 - Pions/muons energy spectrum
 - Particle time distribution
 - Energy deposition in the chicane
- Conclusions



Introduction

- The MC under current investigation is proton driven. **Protons** impact on a solid or liquid target **generating pions** by inelastic collisions. [1]
- The generated pions travels through a tapering region where the magnetic field is adiabatically decreasing. The effect of this section is to decrease the angular divergence of the produced pions.
 [2,3]
- Finally, the beam enters a chicane where the high energy component of the beam is intercepted. Low momentum components (muons and pions) are forced to follow the field lines generated by a series of solenoids. [4]
- The scope of these studies is to assess the radiation load to the equipment in the target area (target and magnets) and develop a shielding design. The studies consider a generic setup as a first step using the MAP design as a starting point. All the simulation are conducted using FLUKA.
- All the results will be normalized per unit MW proton beam intensity with 200 days of operation per year.





Parameters considered for these radiation studies

Table 1: Parameters table.		
Material	Graphite $(1.8 \mathrm{g/cm^3})$	Liquid lead $(10.5\mathrm{g/cm^3})$
Inelastic scattering length	44.94 cm	$17.34\mathrm{cm}$
Target radius	$15\mathrm{mm}$	$15 \mathrm{mm} \ (+ 5 \mathrm{mm} \ \mathrm{vessel})$
Target length	80 cm [3]	$29.7 \mathrm{cm} (+2 \mathrm{cm} \mathrm{vessel})$
Beam size (round)	$5\mathrm{mm}$	
Beam power (normalization purposes)	$(1 \mathrm{MW})$	
Beam energy	$5{ m GeV}$	
Shielding thickness	$42.2\mathrm{cm}$	
Magnet aperture (radius)	$60\mathrm{cm}$	
Peak magnetic field		20 T
Realistic values under consideration are higher (1.5-4 MW)		

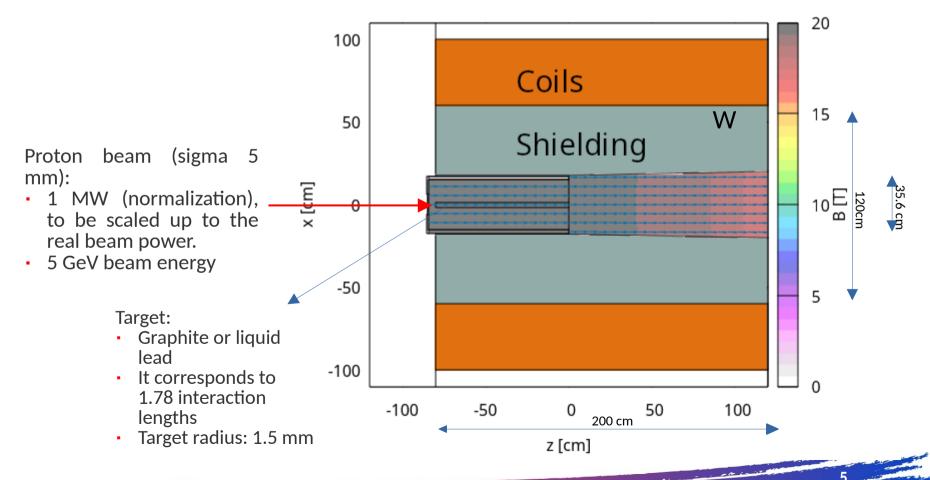
A Contraction of the second second



Target geometry

Generic shielding and magnet geometry:

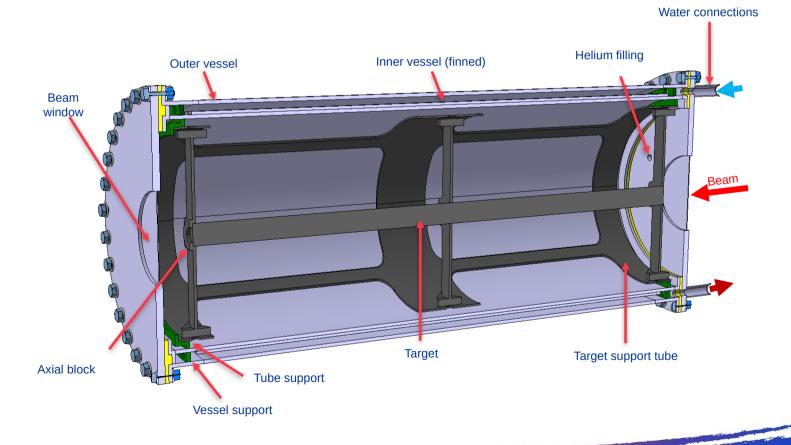
- Tungsten considered for the shielding (engineering and material aspects to be studied)
- No neutron absorber yet included





Target geometry

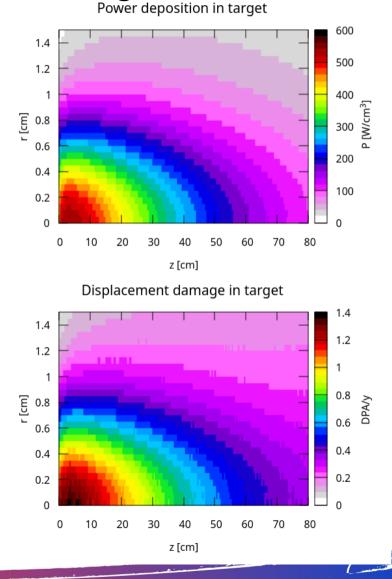
 The graphite target and vessel geometry has been studied by means of thermomechanical simulations (see the presentation of Rui Franqueira Ximenes)





Graphitic target: heat deposition and displacement damage

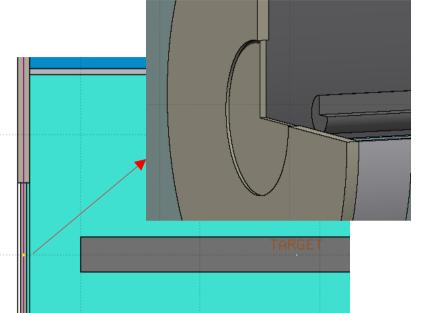
- In a graphitic target, the energy deposition is mostly concentrated in the first section of the target.
- Values up to 600 W/cm³ are observed per MW proton beam.
- The displacement damage follows the same profile, with values up to
 1.5 DPA per year per MW beam.
- The total power deposition in the target is 5.5% of the proton beam power. (34% goes in the shielding surrounding the target and 0.2% in the coils)



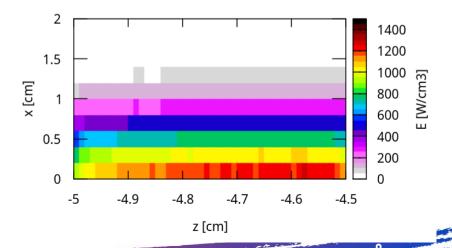


Window for the graphite target

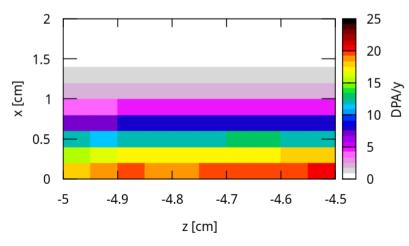
- Various windows option are under exploration. Current possibilities are Ti based materials with thickness ranging from 1 to 5 mm (with a C target, smaller targets can be foreseen). (Ti DPA energy threshold: 30 eV)
- They will need to withstand high power deposition values and intense radiation damage.
- Based on the current configuration, the displacement damage might be too excessive. To reduce this value, larger beam sizes might be needed.
- The possibility of a windowless target could be explored.







5 mm DPA in window per year



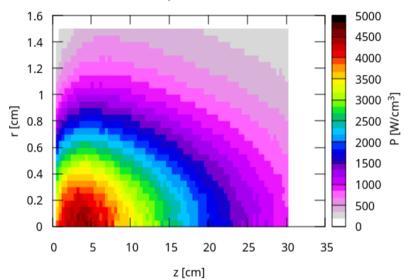


Liquid lead target and vessel: radiation effects

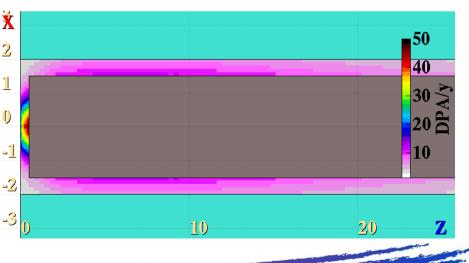
- In a liquid lead target, the power densities are higher in comparison with the graphite target due to the high Z and material density.
- Values up to 4500 W/cm³ are observed per MW proton beam.
- The total power deposition in the target is 20% of the proton beam power. (32% goes in the shielding and 0.28% in the coils)

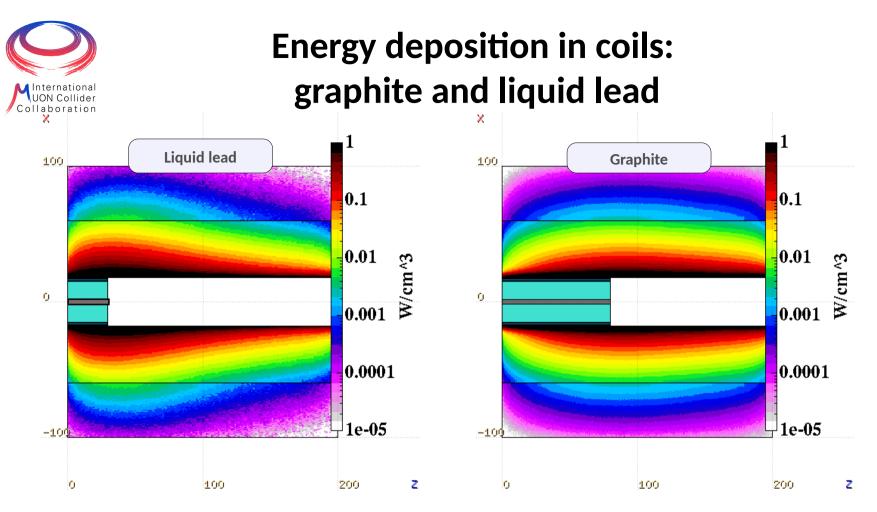
The vessel is assumed to be made of stainless steel.

- Due to the small beam size, DPA peaks in the front window are up to 30-40 per year. Larger beam sizes might be needed.
- The lateral side of the vessel is exposed to the shower products. We observe DPA values up to 8-10, which would potentially require a larger vessel
- The downstream window shows values of ~1s DPA per year per MW beam

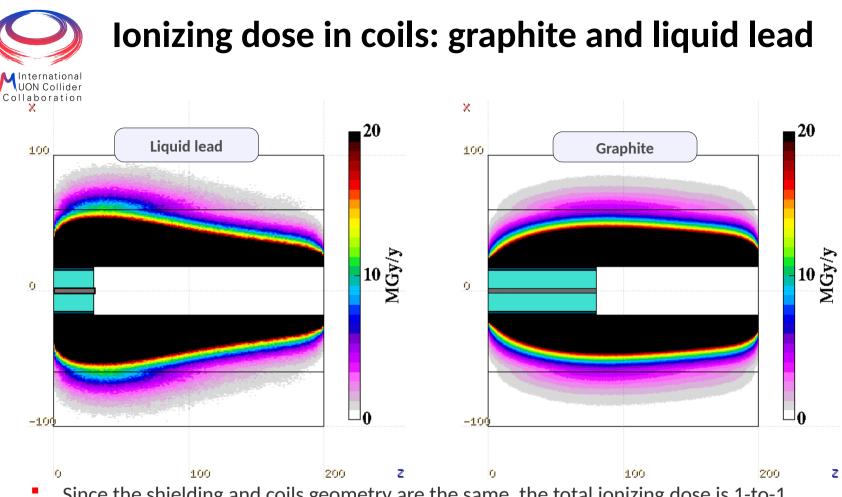


Power deposition in lead

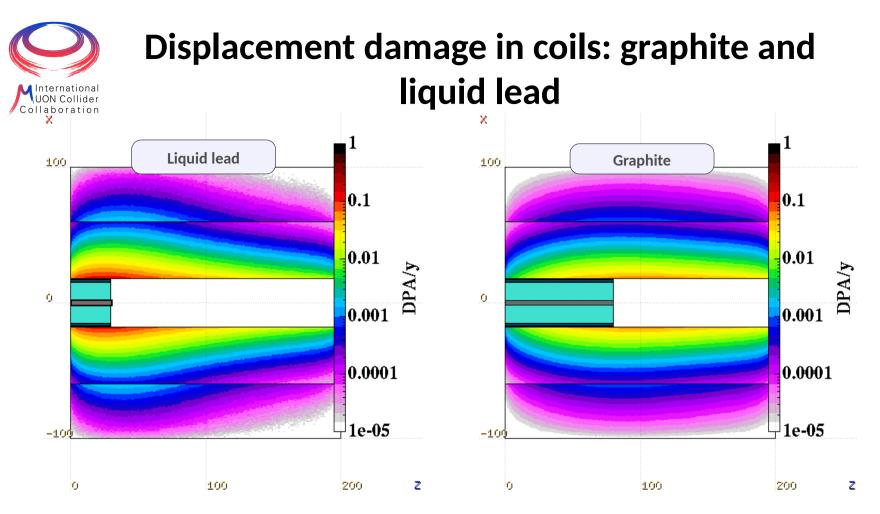




- Considering the same number of inelastic scattering lengths (1.78) the energy deposition distribution is different when considering a graphite and a liquid lead target.
- The lead target produces a shorter shower development, therefore the hotspot in the coils is closer to the target.



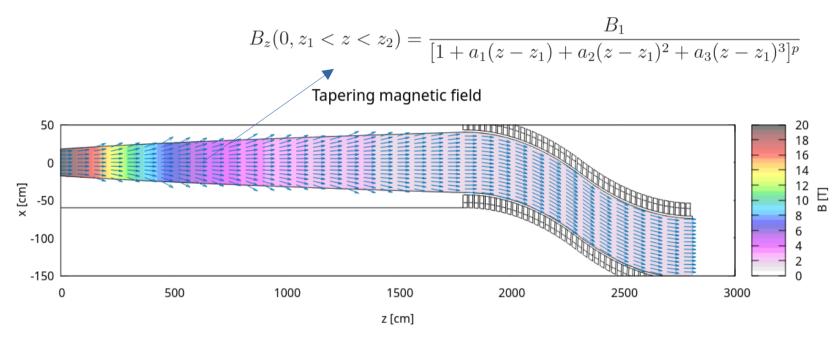
- Since the shielding and coils geometry are the same, the total ionizing dose is 1-to-1 function to the power deposition.
- The scale chosen for these plots **highlight** the **effect on the coils**.
- After 1 year with 1 MW power, the coils receive ~10 MGy in a case of a liquid lead target. This might be not sustainable with the material degradation.
- Additional shielding can be positioned in the liquid lead target case, since the cooling capability is provided by the liquid lead circuit itself.



- The damage profile follows the energy distribution one. Again, the liquid lead target cause an higher displacement damage in the coils.
- **No neutron absorber yet included**! This can significantly change the picture.



Tapering and chicane

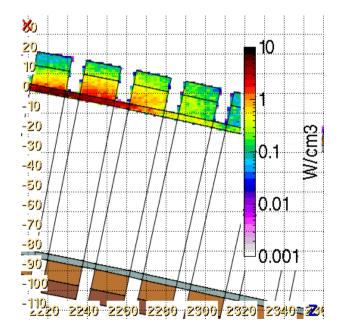


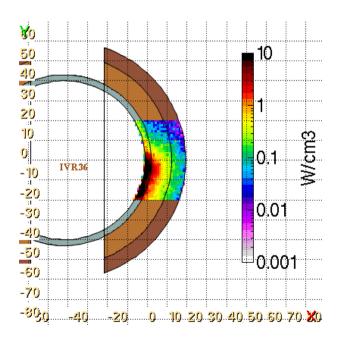
- All the results are simulated considering the **graphite target**.
- The magnetic field chosen for adiabatic section is an inverse cubic tapering.
- The chicane coils are 18 cm long, placed at 0.625° intervals (25 cm in s). The inner radius of the coils is 43 cm. A tungsten shielding of 3 cm is placed uniformly inside the chicane.



Energy deposition in the chicane

- More than 50% of the power is going in the shielding before the chicane!
- The total power deposited in the chicane region is 12.4% of the one from the primary beam. This corresponds to a few 100 kW. A dedicated beam dump might be needed

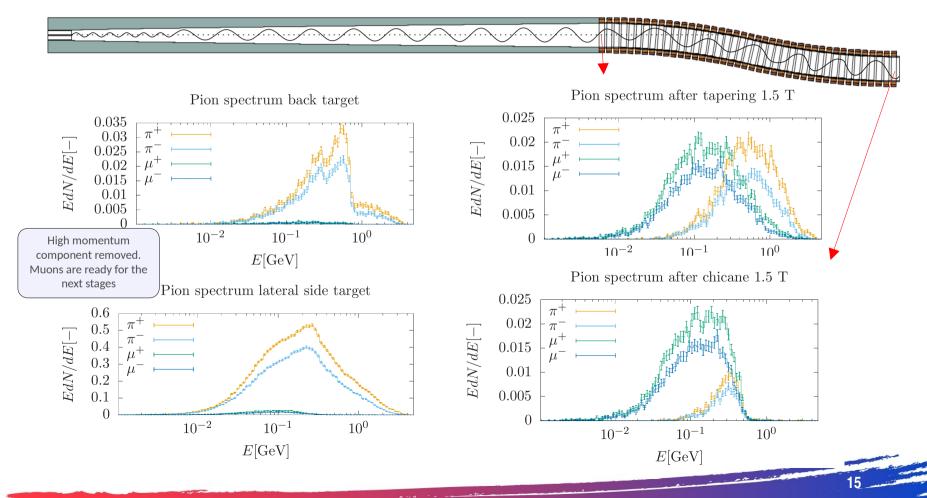






Tapering effect: transverse momentum correction.

- Here there are reported the secondary spectra per primary particle of the proton beam. (normalization factor: 1.25E15 particle per second per MW of beam power with 5 GeV protons)
- The physic optimization should be perform together with the shielding design.





Conclusions

- Target:
 - Radiation load studies for the target area have been started simultaneously with the thermomechanical calculations to assess various target options.
 - Radiation load on graphite and liquid lead target has been studied with a fixed beam size.
 - In both cases, the containment structure (a window to enclose an helium athmosphere for graphitic materials and the vessel for the liquid lead option) are subject to intense displacement damage. This might be a possible lifetime limitation of the target assembly.
- Coils of target solenoid:
 - The energy deposition and long term damage in the shielding and the superconducting coils has been assessed, assuming a generic geometry.
 - The liquid lead target induces shorter shower development and the radiation load to the coils is a factor 1.5 higher.
- Chicane:
 - With a simple geometry, the particle yield after the chicane has been obtained. A dedicated beam dump might be needed





Thank you!

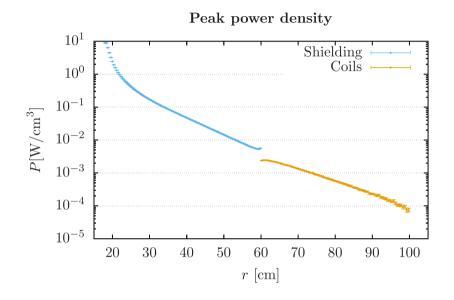


References

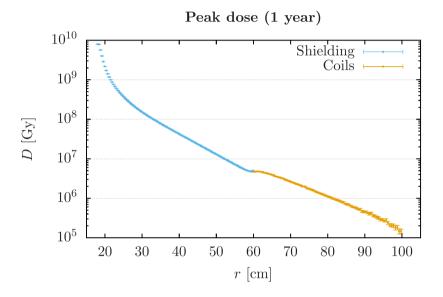
- [1] Considerations on Muon Collider Targetry, Marco Calviani.
- [2] Analytic Forms for an Adiabatic Tapered Solenoid, McDonald. https://www.hep.princeton.edu/mumu/target/taper.pdf
- [3] Preliminary Optimization of the Pion Capture and Decay Channe, Paul et al. https://doi.org/10.1063/1.1818427
- [4] The Target System Baseline, Kirk er al. https://inspirehep.net/files/8020cfe9d0e3474c15ecbf2a5bb29f1a
- [5] CARBON AND MERCURY TARGET SYSTEMS FOR MUON COLLIDERS AND NEUTRINO FACTORIES, McDonald et al. doi:10.18429



Main solenoid: energy deposition and ionizing dose peak



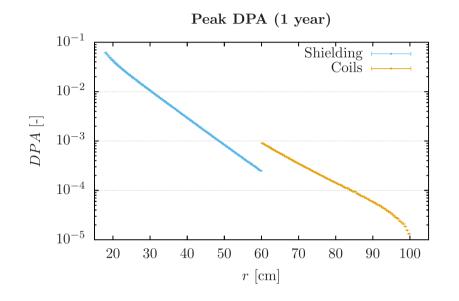
 The small bump in the power deposition in the interface between shielding and coils is mostly due to the difference in density.



 With 40 cm W shielding, the peak dose reaches 4-5 MGy in the coils after 1 year of operation (assuming 1 MW beams)



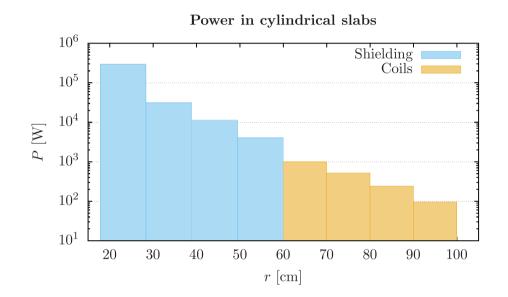
Radial profile of the peak displacement damage



- The considered energy threshold for inducing a displacement is 40 eV for the superconductive coils and 90 eV for the tungsten shielding.
- The gap in the DPA values at the interface is due to:
 - 1) The different values in the energy thresholds (you need to deposit less energy to induce a displacement in the coils).
 - 2) The neutron elastic cross section is larger for light elements.
- With 40 cm W shielding, the peak DPA reaches 10⁻³ in the coils after 1 year of operation (assuming 1 MW beams)



Power in cylindrical slabs

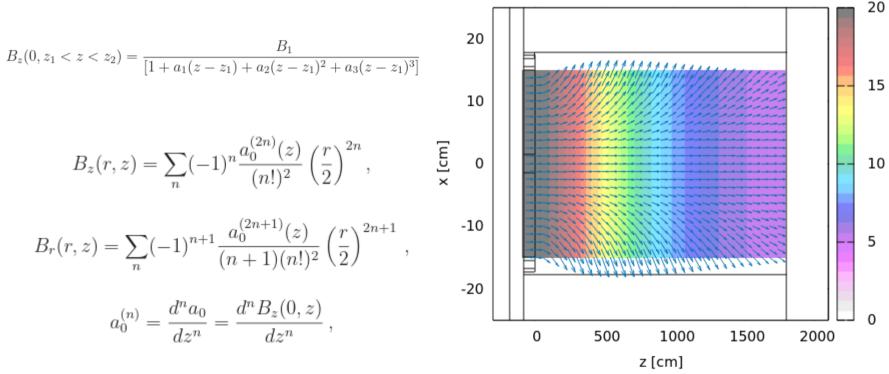


The (integrated) **power deposition** for the 2 m (longitudinal coordinate) **coils** is reported for each cylindrical slab of the magnet.

For example: a cylindrical coil with r_{min} =60cm and r_{max} =70cm (length of 2m) would absorb 1 kW.



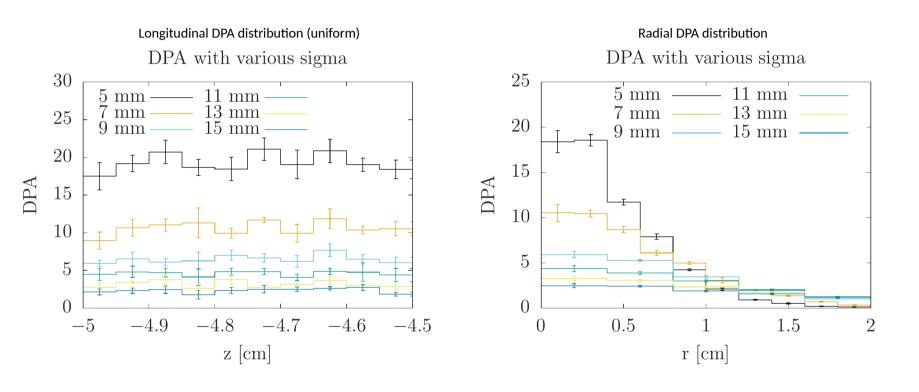
Tapering field: inverse cubic taper



Tapering magnetic field



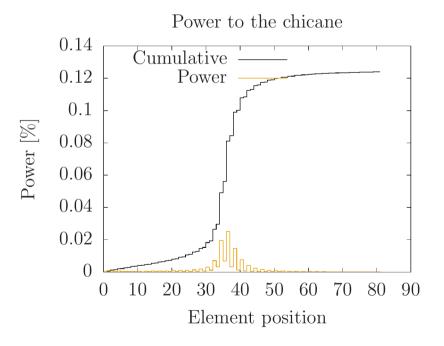
DPA in window (graphite target)



23



Power to the chicane

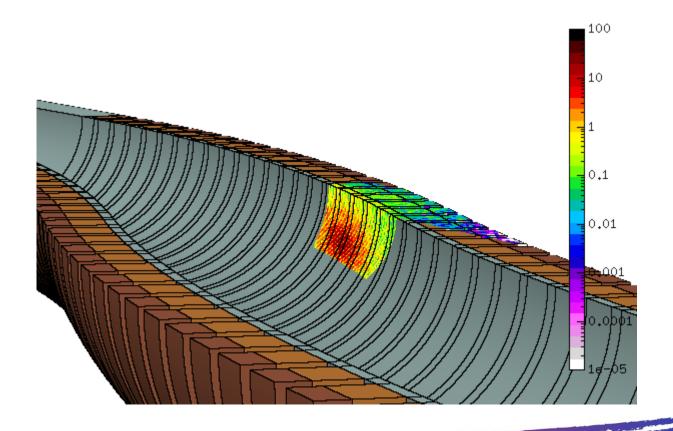


- A C



Heat deposition in the chicane

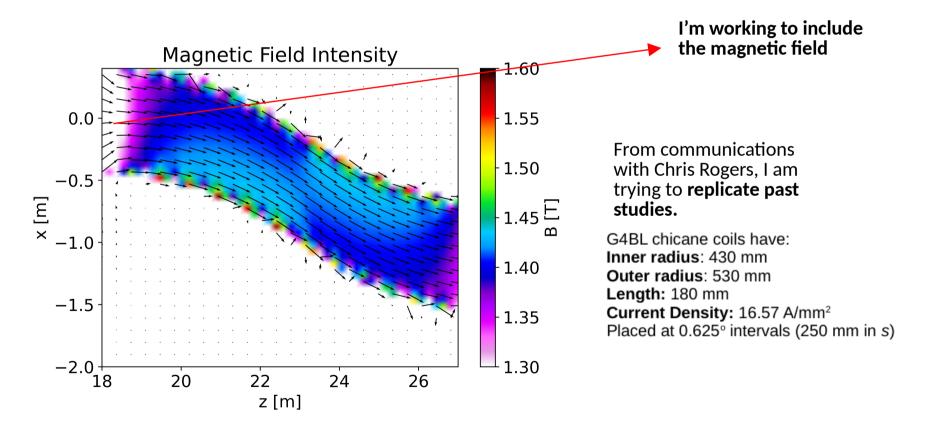
Considering a 1 MW beam, the heat deposition is concerning, a much thicker shielding is required. Since the structure is quite localized, we can implement an asymmetrical shielding if needed.





Tentative magnetic field implementation

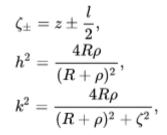
I tried to implement all the different magnets in a numeric way. The source code can be shared upon request. This shouldn't change anything for the heat deposition!





Realistic solenoid (infinitesimal cylinder)

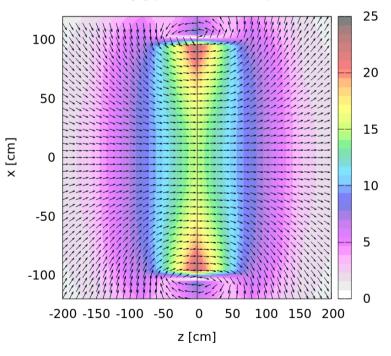
$$\begin{split} B_{\rho} &= \frac{\mu_0 I}{4\pi} \frac{2}{l} \sqrt{\frac{R}{\rho}} \bigg[\frac{k^2 - 2}{k} K(k^2) + \frac{2}{k} E(k^2) \bigg]_{\zeta_-}^{\zeta_+}, \\ B_z &= \frac{\mu_0 I}{4\pi} \frac{1}{l} \frac{1}{\sqrt{R\rho}} \bigg[\zeta k \left(K(k^2) + \frac{R - \rho}{R + \rho} \Pi(h^2, k^2) \right) \bigg]_{\zeta_-}^{\zeta_+} \end{split}$$



Good: they are exact. We can numerically integrate to get a thick solenoid.

Bad: we need to know how the magnet looks like

B [T] (15 T in the centre)

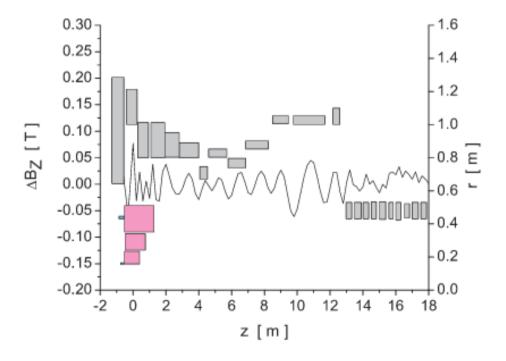


27



Tapering field: achievable?

 Past studies affirm that, with a proper solenoid configuration, the difference between the actual magnetic field and the needed one is minimal. https://journals.aps.org/pr ab/pdf/10.1103/PhysRevS TAB.9.011001





Tapering effect: particle time of arrival

With 1.5 T chicane magnetic field, the time of arrival distribution (0 RMS proton beam)

