Phenomenology of DDVCS in the era of new experiments

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Oct, 28th 2022



- → Starting point: GPD
- Double deeply virtual Compton scattering (DDVCS)
 - Goal & motivation
 - > Formulation à *la* Kleiss & Stirling
 - > Tests of our KS-based formulation
- → Summary

Starting point: GPD

 GPD = Generalized Parton Distribution ≈ "3D version of a PDF (Parton Distribution Function)." With x the fraction of the hadron's longitudinal momentum carried by a quark:

$$GPD(x,\eta,t) = \frac{1}{2} \int \frac{dz^{-}}{2\pi} e^{ix\bar{p}^{+}z^{-}} \langle P_{2} | \bar{\mathfrak{q}}_{f}(-z/2)\gamma^{+} \mathcal{W}[-z/2,z/2] \mathfrak{q}_{f}(z/2) | P_{1} \rangle |_{z_{\perp}=z^{+}=0}$$
$$= \Delta^{2} = (p_{2}-p_{1})^{2}, \quad \eta = -\frac{q\Delta}{pq}, \quad q = \frac{q_{1}+q_{2}}{2}, \quad p = p_{1}+p_{2}$$

- Importance:
 - Connected to QCD energy-momentum tensor, and so to spin. GPDs are a way to address the hadron's spin puzzle
 - Tomography: distribution of longitudinal momentum on the transverse (to hadron's motion) plane

$$q(x, \mathbf{b}_{\perp}) = \int \frac{\mathrm{d}^{2} \boldsymbol{\Delta}}{4\pi^{2}} e^{-i\mathbf{b}_{\perp} \cdot \boldsymbol{\Delta}} \frac{H^{q}(x, 0, t = -\boldsymbol{\Delta}^{2})}{\mathbf{A} \operatorname{particular GPD}}$$

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Our goal

- **Goal:** assessment of feasibility of double deeply virtual Compton scattering (DDVCS) at JLab12, EIC and JLab20+. Phenomenology: xsec, asymmetries, etc
- What is DDVCS? Electroproduction of a lepton pair



$$t = \Delta^2 = (p_2 - p_1)^2, \quad \eta = -\frac{q\Delta}{pq},$$

 $q = \frac{q_1 + q_2}{2}, \quad p = p_1 + p_2$

GPD = Generalized Parton Distribution. Factorized amplitude into GPD and hard contribution (perturbative term)

Why DDVCS?

- Goal: assessment of feasibility of double deeply virtual Compton scattering (DDVCS) at JLab12, EIC and JLab20+
- **Problem:** currently, GPDs are accessible experimentally via deeply virtual Compton scattering (DVCS), timelike Compton scattering (TCS) and deeply virtual meson production (DVMP) only



Why DDVCS?

- **Goal:** assessment of feasibility of double deeply virtual Compton scattering (DDVCS) at JLab12, EIC and JLab20+
- **Problem:** DVCS amplitude allows for measurement of the GPD with restriction to $x = \eta$. Similar situation happens with TCS for $x = -\eta$

DVCS



Sketch of DVCS amplitude (LO)

$$\begin{aligned} \mathcal{A}_{\text{DVCS}} &\sim \int_{-1}^{1} dx \; \frac{1}{x - \eta + i0} \text{GPD}(x, \eta, t) + \cdots \\ &= \text{PV}\left(\int_{-1}^{1} dx \frac{1}{x - \eta} \text{GPD}(x, \eta, t)\right) - \int_{-1}^{1} dx \; i\pi \delta(x - \eta) \text{GPD}(x, \eta, t) + \cdots \end{aligned}$$

So we can measure GPDs at $x = \eta$ only, i.e., we can access $\text{GPD}(\eta, \eta, t)$

Real part may be expressed by the imaginary part by means of dispersion relations

Why DDVCS?

- **Problem:** DVCS amplitude allows for measurement of the GPD with restriction to $x = \eta$. Similar situation happens with TCS for $x = -\eta$
- Solution by DDVCS: the extra virtuality allows for the introduction of a new (generalized) Bjorken variable ξ so that we can access GPDs for $x = \xi \neq \eta$



Formulation à la Kleiss & Stirling

- 1st proposed by Belitsky, Mueller, Guidal and Vanderhaeghen in:
 - Exclusive Electroproduction of Lepton Pairs as a Probe of Nucleon Structure, PRL 90, 022001 (2003)
 - > Double Deeply Virtual Compton Scattering off the Nucleon, PRL 90, 012001 (2003)
- Xsec by Belitsky and Mueller in *Probing generalized parton distributions with electroproduction of lepton pairs off the nucleon*, Phys. Rev. D 68, 116005 (2003)
- That work seems to present some typos or mismatches because we cannot reproduce appropriate limits with it: taking a virtuality of DDVCS to be a reality you get either DVCS or TCS
- Consequently, we have performed a rederivation of DDVCS' formulae via Kleiss & Stirling's methods

Kleiss & Stirling's technique (KS): the basics

- The idea of KS: compute amplitudes, not the modulus squared of them
- Transform spinor products into new scalars s and t (prevents the use of traces of Dirac gamma matrices):

$$s(p_1, p_2) := \bar{u}_+(p_1)u_-(p_2) = -s(p_2, p_1)$$
$$t(p_1, p_2) := \bar{u}_-(p_1)u_+(p_2) = [s(p_2, p_1)]^*$$
$$s(p_1, p_2) = (p_1^y + ip_1^z)\sqrt{\frac{p_2^0 - p_2^x}{p_1^0 - p_1^x}} - (p_2 \leftrightarrow p_1)$$

KS' paper: Spinor Techniques for Calculating p anti $p \rightarrow W+-/Z0 + Jets$. Nuclear Physics B262 (1985) 235-262

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Not only VCS but also BH (pure QED)



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Example: BH1 à la KS

• KS application to BH1 diagram of DDVCS: $i\widetilde{\mathcal{M}}_{BH1} = \left(\frac{ie^4}{(q_2^2 + i0)(\Delta^2 + i0)((k - \Delta)^2 + i0)}\right)^{-1} i\mathcal{M}_{BH1}$ amplitude

$$\begin{split} i\widetilde{\mathcal{M}}_{\rm BH1} = & (F_1 + F_2) \sum_L \left(Y_{s_2s_1} f(s_\ell, \ell_-, \ell_+; s, k', L) f(s, L, k; +, r'_{s_2}, r_{s_1}) + Z_{s_2s_1} f(s_\ell, \ell_-, \ell_+; s, k', L) f(s, L, k; -, r'_{-s_2}, r_{-s_1}) \right) - \frac{F_2}{2M} J_{s_2, s_1}^{(2)} \sum_{L, R} f(s_\ell, \ell_-, \ell_+; s, k', L) g(s, L, R, k) \end{split}$$

• For example, *f* function is defined as

 $f(s, k_0, k_1; s', k_2, k_3) = u_s(k_0)\gamma^{\mu}u_s(k_1)u_{s'}(k_2)\gamma_{\mu}u_{s'}(k_3)$

that can be expressed by means of *s* and *t* KS scalars

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Dedicated softwares

→ PARTONS platform: open-source C++ program

- Contains several GPD models
- Leading twist... but higher twist corrections will be included in near future
- Useful for theorists and experimentalists
- Provides xsecs, Compton Form Factors, etc
- DVCS, TCS and DVMP are already included

To download and for tutorials: http://partons.cea.fr

Description of architecture: Eur. Phys. J. C78 (2018), 478

Software



Dedicated softwares

→ EpIC Monte Carlo event generator in C++

- Uses PARTONS framework
- Includes radiative corrections
- Generates the kinematic configurations following the probability distributions given by PARTONS
- DVCS, TCS and DVMP are already included



Access EpIC via GitHub:

https://github.com/pawelsznajder/epic

Detail description and architecture: arXiv:2205.01762 [hep-ph]

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DDVCS to DVCS



DDVCS to TCS

Evaluate energy of incoming virtual photon to be used as energy of TCS photon beam

$$\nu = \frac{Q^2}{2Mx_B}$$

Divide by flux Γ and get rid of x_B and Q^2 differentiation

$$\Gamma = \frac{\alpha_{em}}{2\pi Q^2} \left(1 + \frac{(1-y)^2}{y} - \frac{2(1-y)Q_{\min}^2}{yQ^2} \right) \frac{\nu}{Ex_B}, \qquad Q_{\min}^2 = \frac{(ym_e)^2}{1-y}$$

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DVCS limit (BH1 + crossed)



 $x_B = 0.04$, t = -0.1 GeV², $-q_1^2 = 10$ GeV², $q_2^2 \approx 0.001$ GeV², $E_{beam} = 160$ GeV

DVCS formulae:

Belitsky et al., *Theory of deeply virtual Compton scattering on the nucleon*, Nuclear Physics B629 (2002)

Belitsky et al., *Compton scattering: from deeply virtual to quasi-real*, Nuclear Physics B878 (2014)

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TCS limit (BH2 + crossed)



 $x_B = 2 \cdot 10^{-4}$, t = -0.5 GeV², -q₁² = 2 \cdot 10^{-3} GeV², q₂² = 1 GeV², E_{beam} = 12 GeV

TCS formulae:

Berger et al., *Timelike Compton scattering: exclusive photoproduction of lepton pairs*, The European Physics Journal C 23 (2002)

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DVCS limit (VCS)



 $x_B = 0.04, t = -0.01 \text{ GeV}^2, -q_1^2 = 10$ GeV², E_{beam} = 160 GeV, $\phi = 0$

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- Xsec goes as 1/Q² meaning that there is a delicate cancellation of such a dependence on the amplitude
- As shown, these limits are not trivial and require a careful analysis



- New analytical formulae has been derived •
- DDVCS is already implemented in PARTONS (Lo and LT)
- We are interested in observables such as the beam spin asymmetry proportional to . Im(BH x VCS*) **CONC**
- Code will be included in EpIC MC generator to study feasibility of DDVCS •
- We are in contact with experimentalists that will use it and we invite everybody to it •

