

Production of molecular structure hadron in Pb-Pb collisions



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Contents

1. 2-dimensional coalescence model

- 2-body coalescence
- 3-body coalescence

2. Result

- p_T distribution of deuteron and helium-3 in Pb-Pb collisions at 2.76TeV
- Prediction of $X(3872)$ ($D\bar{D}$) and T_{cc} (DD^* or 4-quark) p_T distribution in Pb-Pb collisions at 5.02TeV

3. Summary

Coalescence model

2-dimension coalescence model

- The yields of produced hadron

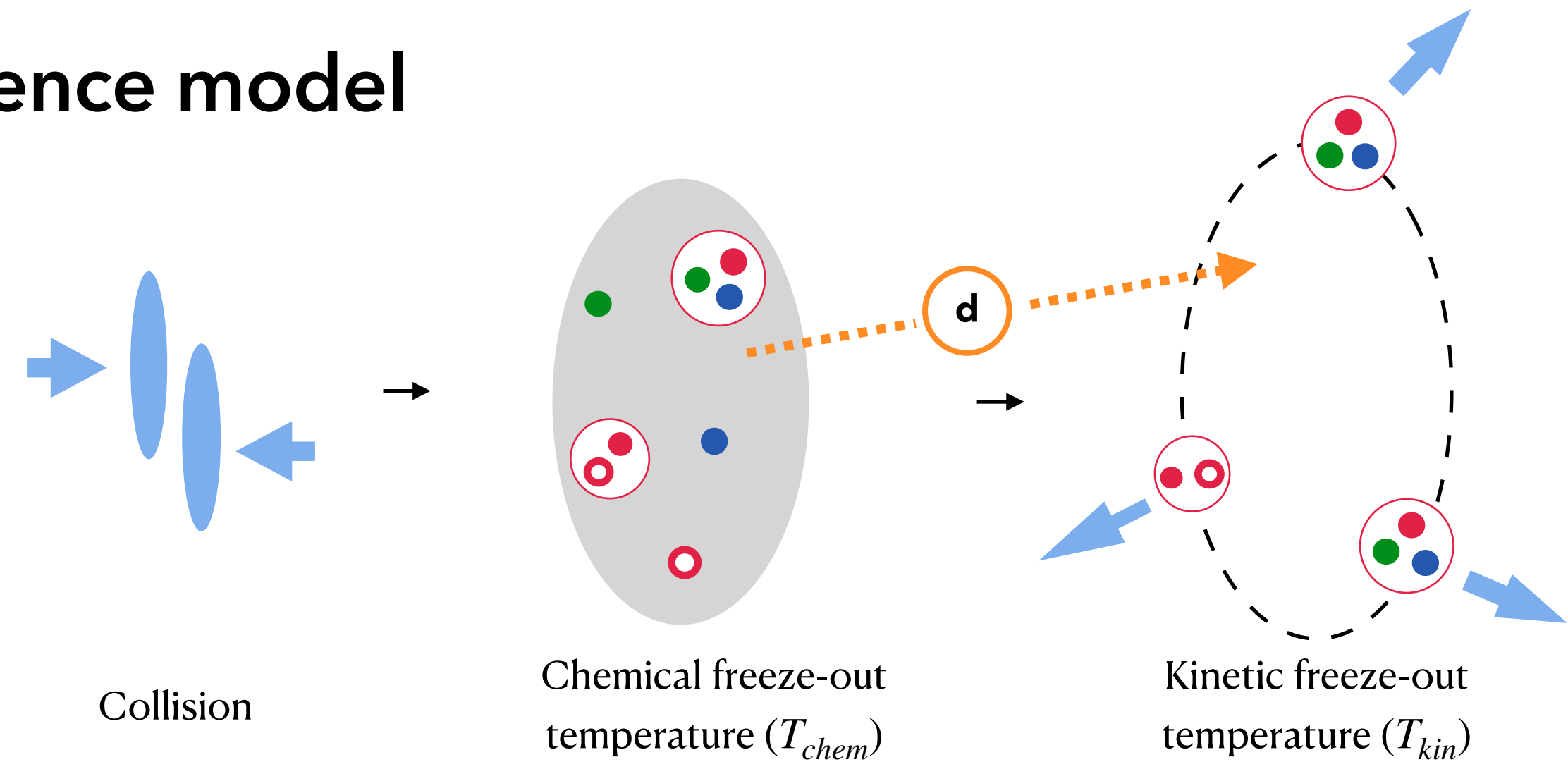
$$N_h = g_h \int \prod_{i=1}^N d^2x_i d^2p_i f_i(x_i, p_i) W(r_1, \dots, r_{N-1}, k_1, \dots, k_{N-1})$$

- Hadron transverse momentum distribution

$$\frac{d^2N_h}{d^2P_T} = g_h \int \prod_{i=1}^N d^2x_i d^2p_i f_i(x_i, p_i) W(r_1, \dots, r_{N-1}, k_1, \dots, k_{N-1}) \delta^{(2)} \left(P_T - \sum_{j=1}^N p_j \right)$$

- Normalization condition

$$\int d^2x_i d^2p_i f_i(\vec{x}_i, \vec{p}_i) = N_i, \quad \int \prod_{i=1}^{N-1} d^2r_i d^2k_i W_H(\vec{r}_1, \dots, \vec{r}_{N-1}, \vec{k}_1, \dots, \vec{k}_{N-1}) = (2\pi)^{2(N-1)}$$



Coalescence model

2-body coalescence

$$d : p + n$$

$$X(3872) : D^* + \bar{D}^0$$

$$T_{cc} : D + D^*$$

S. Cho, K. J. Sun, C. M. Ko, S. H. Lee, Y. Oh, Phys. Rev. C 101, 024909 (2020)

• Wigner function : $W(\vec{r}, \vec{k}) = 4 \exp \left[-\frac{(r')^2}{\sigma^2} - \sigma^2(k')^2 \right]$

\vec{r}' (\vec{k}') : relative distance (momentum) of constituent in center of mass frame of produced hadron

• Constituent distribution : $f(x_i, p_i) = \frac{d^2 N_i}{A_L d^2 p_{iT}}$

A_L : Coalescence area at kinetic freeze-out point in Lab frame

Lorentz transformation

$$\Delta t' = \gamma(\Delta t - \beta r_x), \quad r'_x = \gamma(r_x - \beta \Delta t)$$

$$\Delta E' = \gamma(\Delta E - \beta k_x), \quad k'_x = \gamma(k_x - \beta \Delta E)$$

• Coordinate :

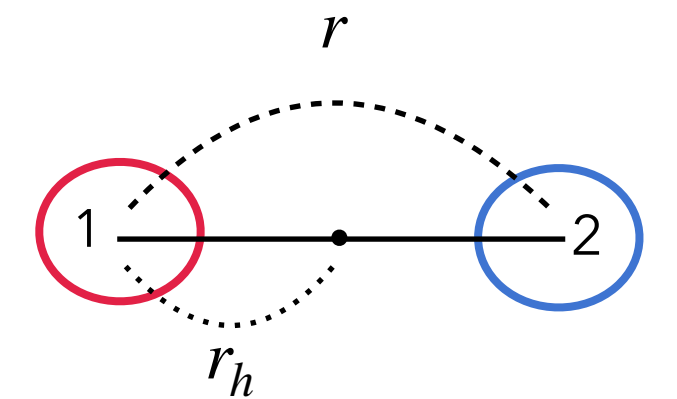
$$R^\mu = \frac{x_1^\mu + x_2^\mu}{2}, \quad r^\mu = x_1^\mu - x_2^\mu,$$

$$P^\mu = p_1^\mu + p_2^\mu, \quad k^\mu = \frac{p_1^\mu - p_2^\mu}{2}$$

$$\frac{d^2 N_h}{d^2 P_T} = \frac{g_h}{g_1 g_2} (2\sqrt{\pi})^2 \sigma^2 \frac{1}{A} \int d^2 p_1 d^2 p_2 \frac{d^2 N_1}{d^2 p_{1T}} \frac{d^2 N_2}{d^2 p_{2T}} \exp[-\sigma^2(k')^2] \delta^{(2)}(P_T - p_{1T} - p_{2T}), \quad \langle r^2 \rangle = \frac{3}{2} \sigma^2 \rightarrow \langle r_h^2 \rangle = \frac{3}{8} \sigma^2$$

• In $\sigma \rightarrow \infty$ limit, Wigner function : $W(\vec{r}, \vec{k}) = 4 \left(\frac{\pi}{\sigma^2} \right) e^{-\frac{r^2}{\sigma^2}} \times \delta^{(2)}(\vec{k}')$

$$\frac{d^2 N_h}{d^2 P_T} = \frac{g_h}{g_1 g_2} (2\pi)^2 \left(\frac{\gamma}{A} \right) \frac{d^2 N_1}{d^2 p_{1T}} \Big|_{\vec{p}_{1T} = \frac{\vec{P}_T}{2}} \frac{d^2 N_2}{d^2 p_{2T}} \Big|_{\vec{p}_{2T} = \frac{\vec{P}_T}{2}}$$



Coalescence model

3-body coalescence

$${}^3\text{He} : p + p + n$$

· Wigner function : $W(\vec{r}, \vec{k}) = 4^2 \exp \left[-\frac{(r'_1)^2}{\sigma_1^2} - \sigma_1^2 (k'_1)^2 \right] \exp \left[-\frac{(r'_2)^2}{\sigma_2^2} - \sigma_2^2 (k'_2)^2 \right]$ · Coordinate :

$$R^\mu = \frac{x_1^\mu + x_2^\mu + x_3^\mu}{3}, \quad r_1^\mu = x_1^\mu - x_2^\mu, \quad r_2^\mu = \frac{x_1^\mu + x_2^\mu}{2} - x_3^\mu$$

$$P^\mu = p_1^\mu + p_2^\mu, \quad k_1^\mu = \frac{p_1^\mu - p_2^\mu}{2}, \quad k_2^\mu = \frac{p_1^\mu + p_2^\mu - 2p_3^\mu}{3}$$

Lorentz transformation :

$$\Delta t'_{1,2} = \gamma(\Delta t_{1,2} - \beta r_{1,2}) \quad \Delta E'_{1,2} = \gamma(\Delta E_{1,2} - \beta k_{1,2})$$

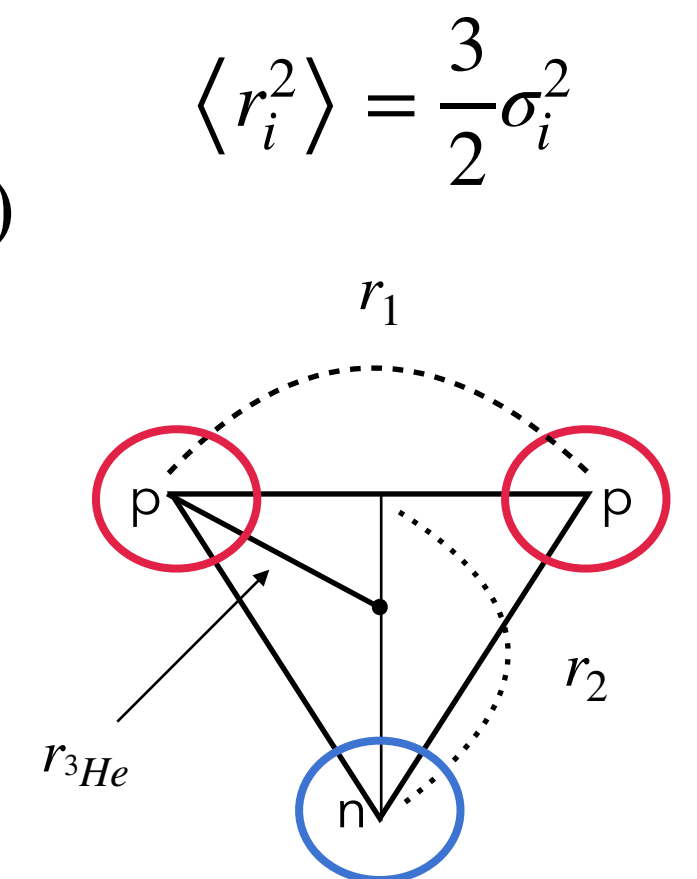
$$r'_{1,2} = \gamma(r_{1,2} - \beta \Delta t_{1,2}) \quad k'_{1,2} = \gamma(k_{1,2} - \beta \Delta E_{1,2})$$

$$\frac{d^2 N_h}{d^2 P_T} = \frac{g_h}{g_1 g_2 g_3} (2\sqrt{\pi})^4 (\sigma_1 \sigma_2)^2 \frac{1}{A^2} \int d^2 p_1^2 p_2 d^2 p_3 \frac{d^2 N_1}{d^2 p_{1T}} \frac{d^2 N_2}{d^2 p_{2T}} \frac{d^2 N_3}{d^2 p_{3T}} \exp \left[-\sigma_1^2 (k'_1)^2 - \sigma_2^2 (k'_2)^2 \right] \delta^{(2)}(P_T - p_{1T} - p_{2T} - p_{3T})$$

$$\langle r_i^2 \rangle = \frac{3}{2} \sigma_i^2$$

In $\sigma \rightarrow \infty$ limit,

$$\frac{d^2 N_h}{d^2 P_T} = \frac{g_h}{g_1 g_2 g_3} (2\pi)^4 \left(\frac{\gamma}{A} \right)^2 \frac{d^2 N_1}{d^2 p_{1T}} \Big|_{\vec{p}_{1T} = \frac{\vec{P}_T}{3}} \frac{d^2 N_2}{d^2 p_{2T}} \Big|_{\vec{p}_{2T} = \frac{\vec{P}_T}{3}} \frac{d^2 N_3}{d^2 p_{3T}} \Big|_{\vec{p}_{3T} = \frac{\vec{P}_T}{3}}$$



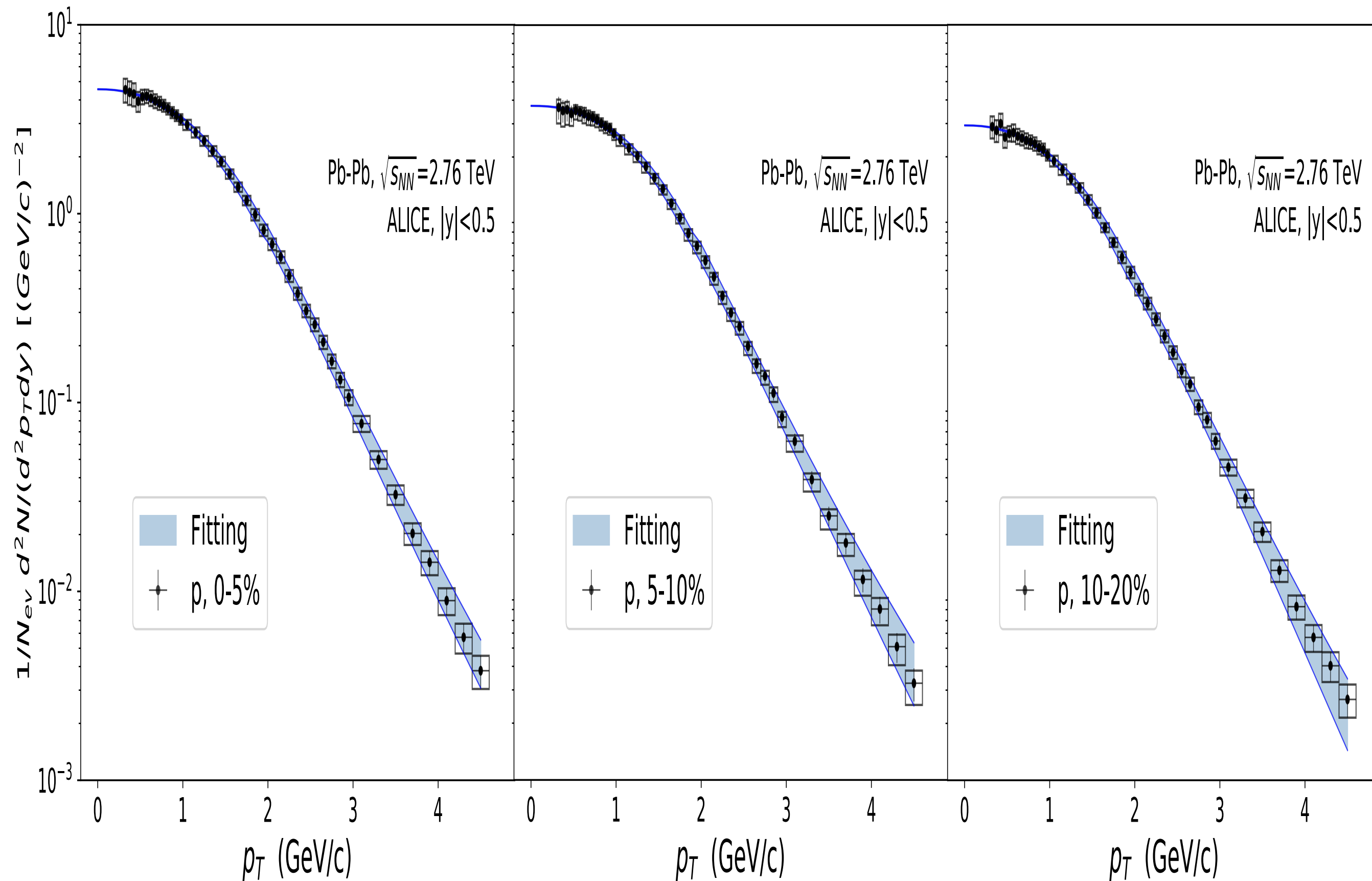
d and ${}^3\text{He}$

in Pb-Pb at $\sqrt{s_{NN}}=2.76\text{TeV}$

Proton p_T distribution

Pb-Pb collisions at 2.76 TeV

○ Fitting (ALICE Collaboration, Phys. Rev. C 88, 044910)



○ Feed-down (Pb-Pb collisions, 2.76 TeV)

Measured proton data includes the feed-down contribution (N , Δ , *etc.*)

$$\bullet \frac{d^2 N_p}{d^2 p_T} \Big|_{t=t_k} = R_b \frac{d^2 N_p}{d^2 p_T} \Big|_{Exp}$$

$$\left(R_b = \frac{\text{the number of proton at formation point}}{\text{the number of final proton}} \right)$$

R_b can be determined using deuteron and helium-3 data ($R_b = 0.368$)

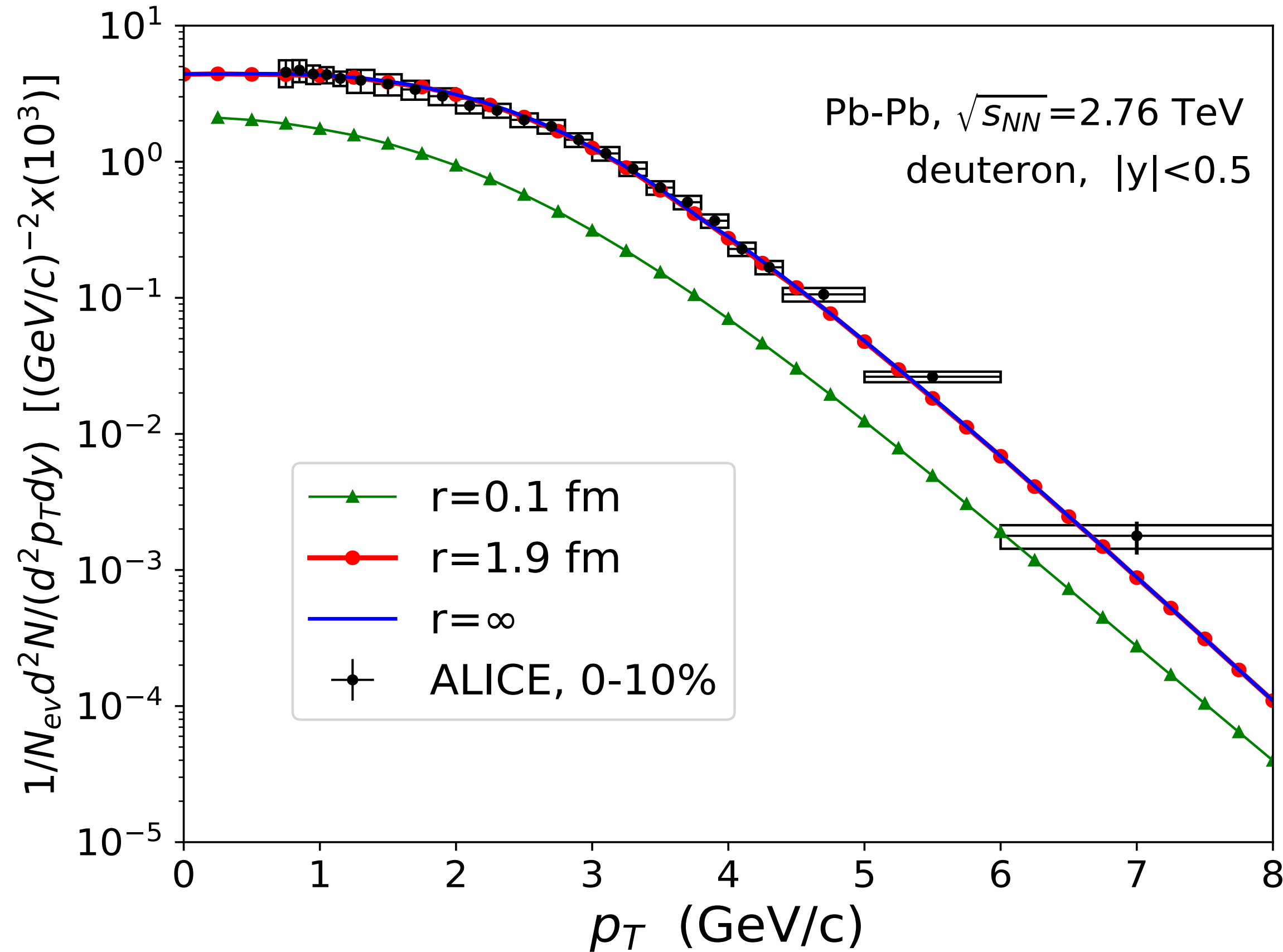
→ 36.8% of the measured p participate in coalescence

R_b is close to the result of statistical hadronization model

Deuteron p_T distribution

Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76\text{TeV}$

- Experimental Data - ALICE Collaboration, Eur. Phys. J. C (2017) 77:658



- 2-body formula

$$1. \quad \frac{d^2 N_d}{d^2 P_T} = g_d (2\sqrt{\pi})^2 \sigma^2 \left(\frac{R_b^2}{A} \right) \int d^2 p_p d^2 p_n \frac{d^2 N_p^{Exp}}{d^2 p_{pT}} \frac{d^2 N_n^{Exp}}{d^2 p_{nT}} \times \exp[-\sigma^2 (k')^2] \delta^{(2)}(P_T - p_{pT} - p_{nT})$$

$$2. \quad \frac{d^2 N_d}{d^2 P_T} (\sigma \rightarrow \infty) = g_d (2\pi)^2 \gamma \left(\frac{R_b^2}{A} \right) \frac{d^2 N_p}{d^2 p_{pT}} \Big|_{\vec{p}_{pT}=\frac{\vec{P}_T}{2}}^{Exp} \frac{d^2 N_n}{d^2 p_{nT}} \Big|_{\vec{p}_{nT}=\frac{\vec{P}_T}{2}}^{Exp}$$

- Parameter

G. Röpke, Phys. Rev. C 79 (2009) 014002

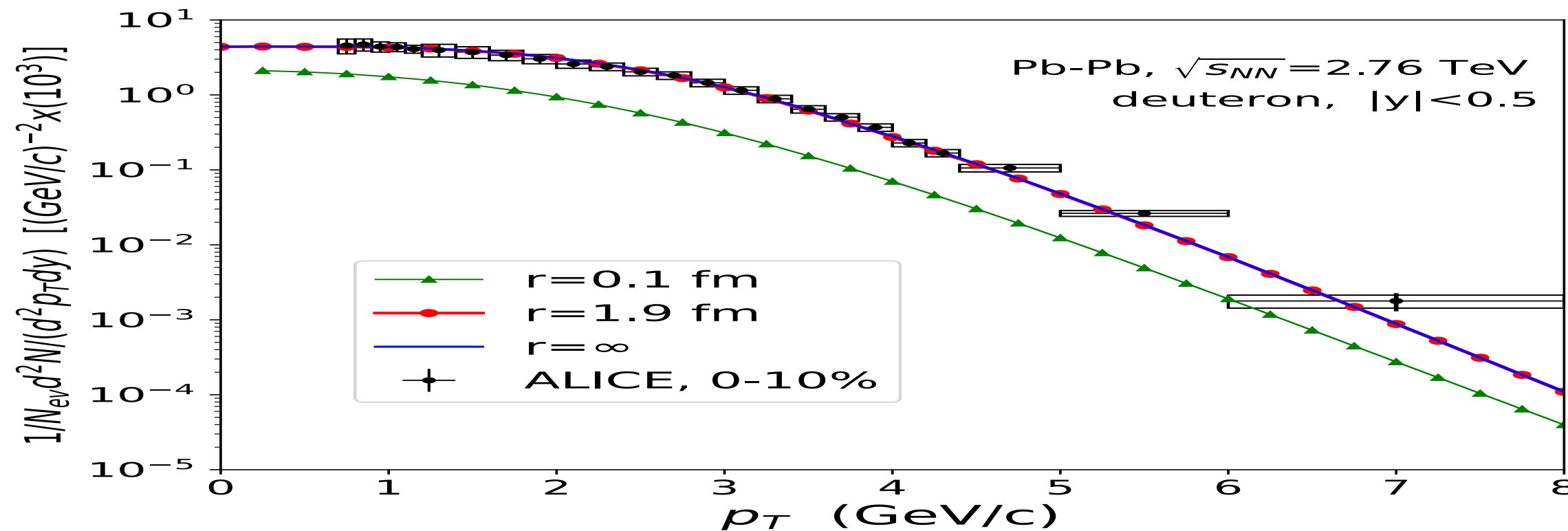
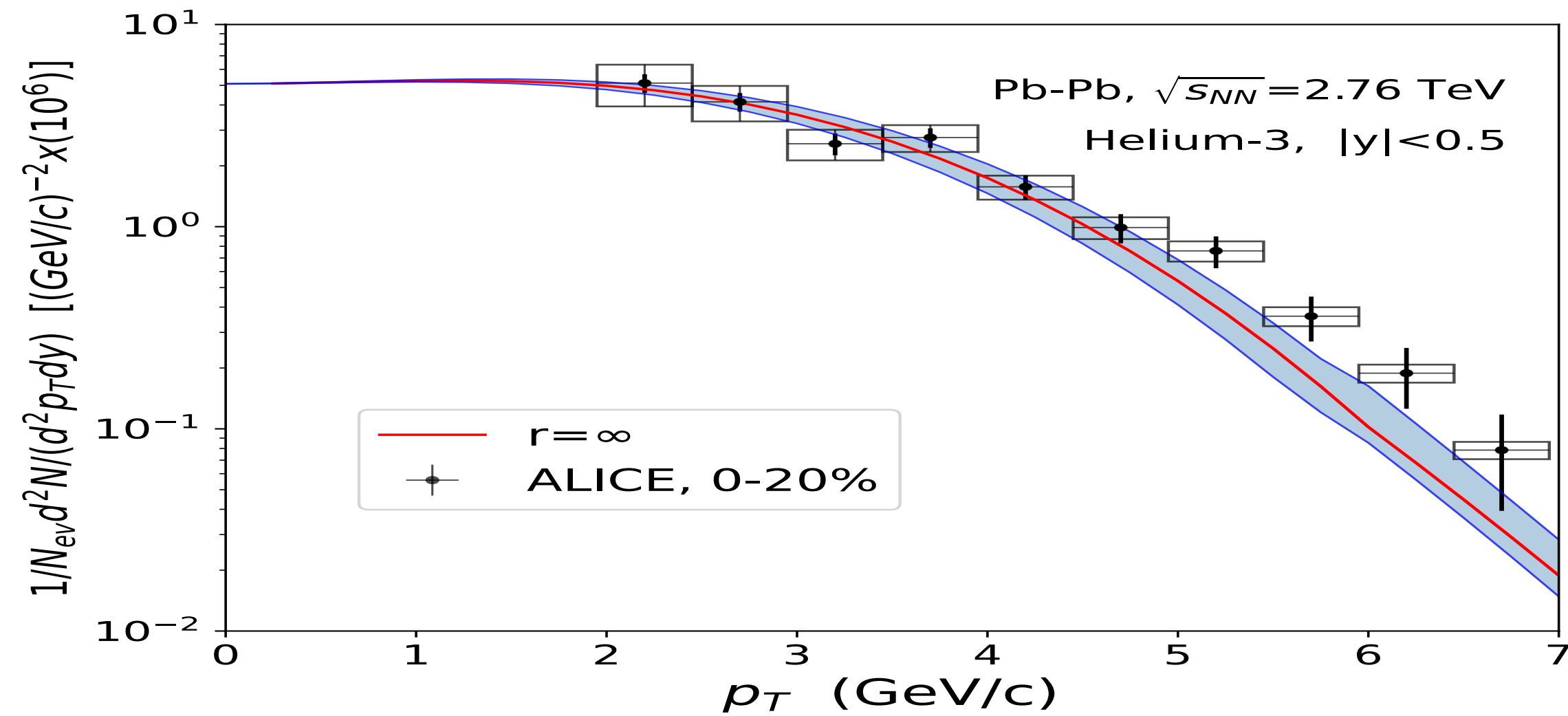
$$\sigma = \sqrt{8/3} r_d, \quad r_d \sim 1.9 \text{ fm}, \quad R_b = 0.368, \quad A_{0-10\%} = 608 \text{ fm}^2$$

- Yield

$$\frac{N_{d, 0-10\%}^{coal}}{N_{d, 0-10\%}^{stat}} = 0.97$$

Helium-3 p_T distribution and consistency of the parameters (R_b , A)

• Phys. Rev. C 93, 024917 (2016)



$$\frac{d^2N}{d^2p_T} \Big|_{t=t_k} = R_b \frac{d^2N}{d^2p_T} \Big|_{Exp}$$

$$R_b = \frac{\text{the number of proton at formation point}}{\text{the number of final proton}}$$

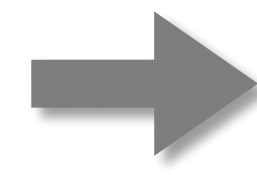
$$\frac{d^2N_d}{d^2P_T} = g_d (2\pi)^2 \gamma \left(\frac{R_b^2}{A} \right) \frac{d^2N_p}{d^2p_{pT}} \Big|_{\vec{p}_{pT}=\frac{\vec{P}_T}{2}} \Big|_{Exp} \frac{d^2N_n}{d^2p_{nT}} \Big|_{\vec{p}_{nT}=\frac{\vec{P}_T}{2}} \Big|_{Exp}$$

$$\frac{d^2N_{^3He}}{d^2P_T} = g_{^3He} (2\pi)^4 \gamma^2 \left(\frac{R_b^3}{A^2} \right) \frac{d^2N_p}{d^2p_{pT}} \Big|_{\vec{p}_{pT}=\frac{\vec{P}_T}{3}} \Big|_{Exp} \frac{d^2N_p}{d^2p_{pT}} \Big|_{\vec{p}_{pT}=\frac{\vec{P}_T}{3}} \Big|_{Exp} \frac{d^2N_n}{d^2p_{nT}} \Big|_{\vec{p}_{nT}=\frac{\vec{P}_T}{3}} \Big|_{Exp}$$

When $R_b \rightarrow \alpha R_b$,

$$\frac{d^2N_d}{d^2p_T} \sim \frac{R_b^2}{A} \rightarrow \frac{\alpha^2 R_b^2}{\alpha^2 A}, \quad \frac{d^2N_{^3He}}{d^2p_T} \sim \frac{R_b^3}{A^2} \rightarrow \frac{\alpha^3 R_b^3}{\alpha^4 A^2}$$

Determine A and R_b correctly
($R_b = 0.368$, $A_{0-10\%} = 608 \text{ fm}^2$)



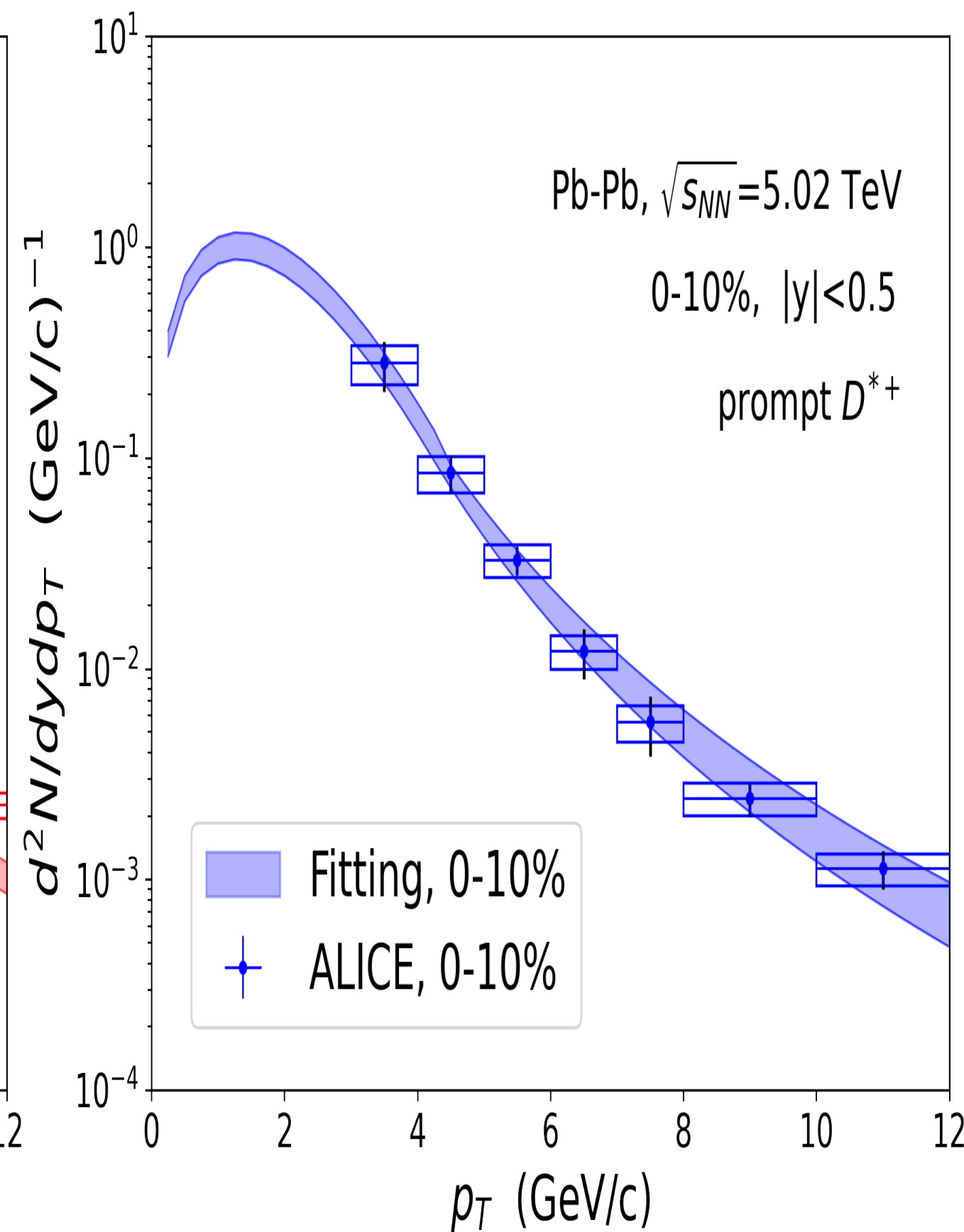
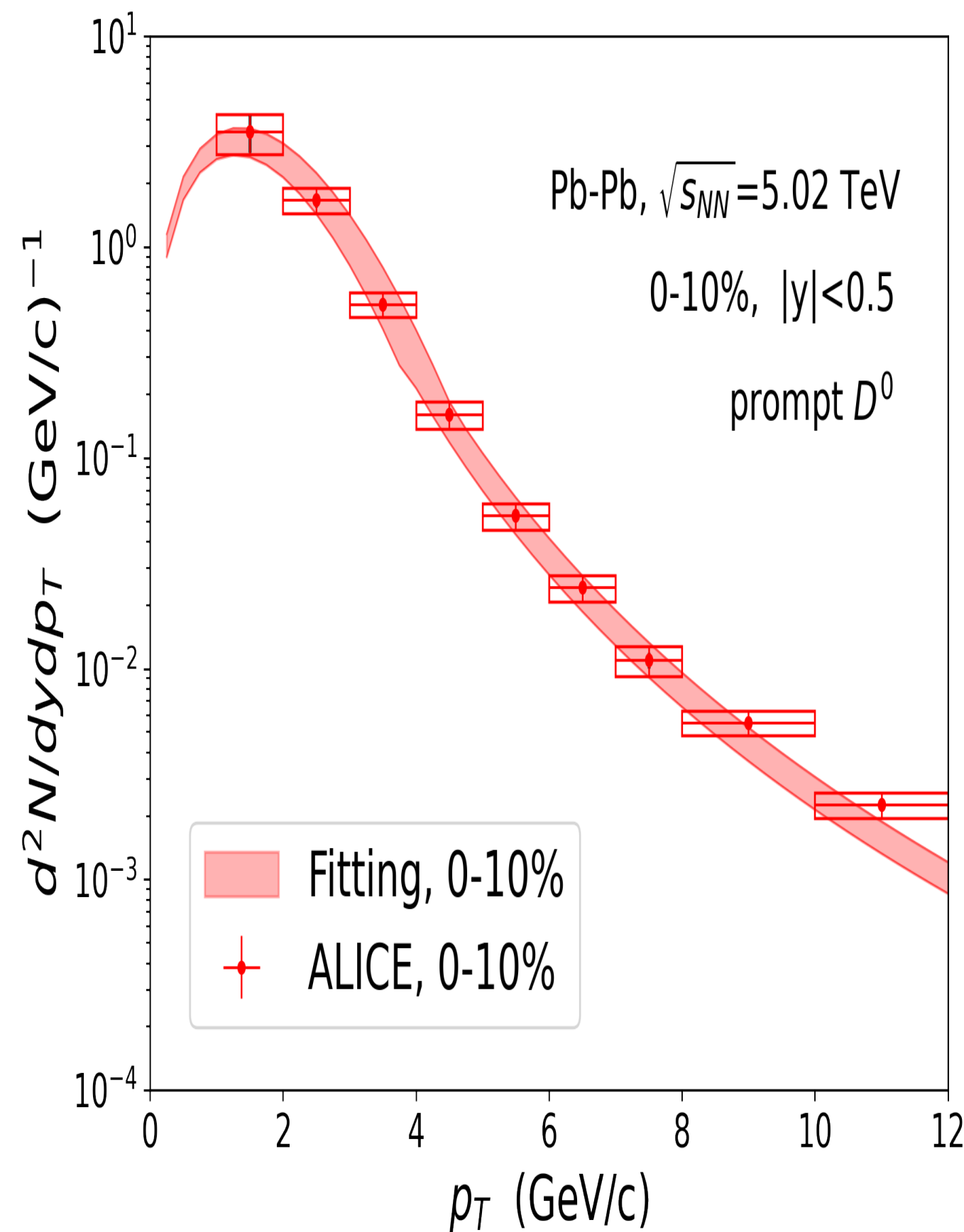
d and 3He can be explained simultaneously

$X(3872)$ and T_{cc}
in Pb-Pb at $\sqrt{s_{NN}}=5.02\text{TeV}$

D meson distribution and Feed-down

Pb-Pb collisions at 5.02TeV

○ Fitting (ALICE Collaboration, JHEP 01 (2022) 174)



○ Feed-down (Pb-Pb collisions, 5.02TeV)

- Experimental data :

$$dN_{0-10\%}^{D^0}/dy = 6.819 \pm 0.457(stat.)_{-0.936}^{+0.912}(syst.)$$

- Decay channel

$$Br(D^*(2007)^0 \rightarrow D^0\pi^0) = (64.7 \pm 0.9) \%$$

$$Br(D^*(2007)^0 \rightarrow D^0\gamma) = (35.3 \pm 0.9) \%$$

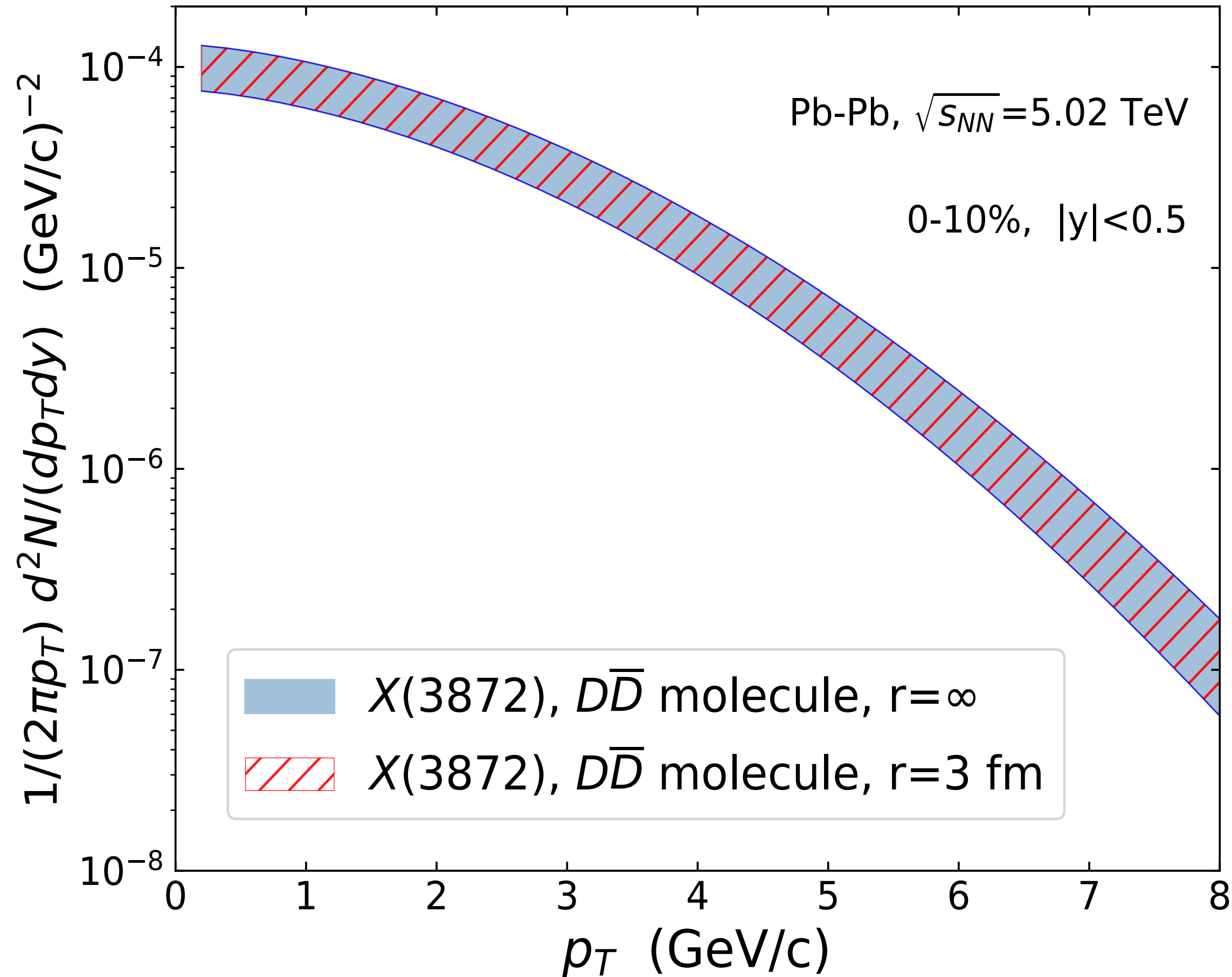
$$Br(D^*(2010)^+ \rightarrow D^0\pi^+) = (67.7 \pm 0.5) \%$$

- From Statistical hadronization model, 31% of measured D^0 participate in coalescence

X(3872) p_T distribution

Pb-Pb collisions at 5.02 TeV

$$X(3872) : D^* + \bar{D}^0$$



$$\frac{d^2 N_{X(3872)}}{d^2 P_T} = g_X (2\sqrt{\pi})^2 \sigma^2 \left(\frac{R_b^2}{A} \right) \int d^2 p_{D^0} d^2 p_{\bar{D}^*} \frac{d^2 N_{D^0}^{Exp}}{d^2 p_{D^0 T}} \frac{d^2 N_{\bar{D}^*}^{Exp}}{d^2 p_{\bar{D}^* T}} \exp[-\sigma^2(k')^2] \times \delta^{(2)}(P_T - p_{D^0 T} - p_{\bar{D}^* T})$$

Radius	X(3872)
Molecule	3fm
Compact 4-quark	Not possible

- Yields

$$\frac{N_{coal}^{X(3872)}}{N_{SHMc}^{X(3872)}} = 2.47 \pm 0.716, \quad \frac{N_{coal}^{X(3872)}}{N_{SHMc}^{\psi(2S)}} = 0.806 \pm 0.234$$

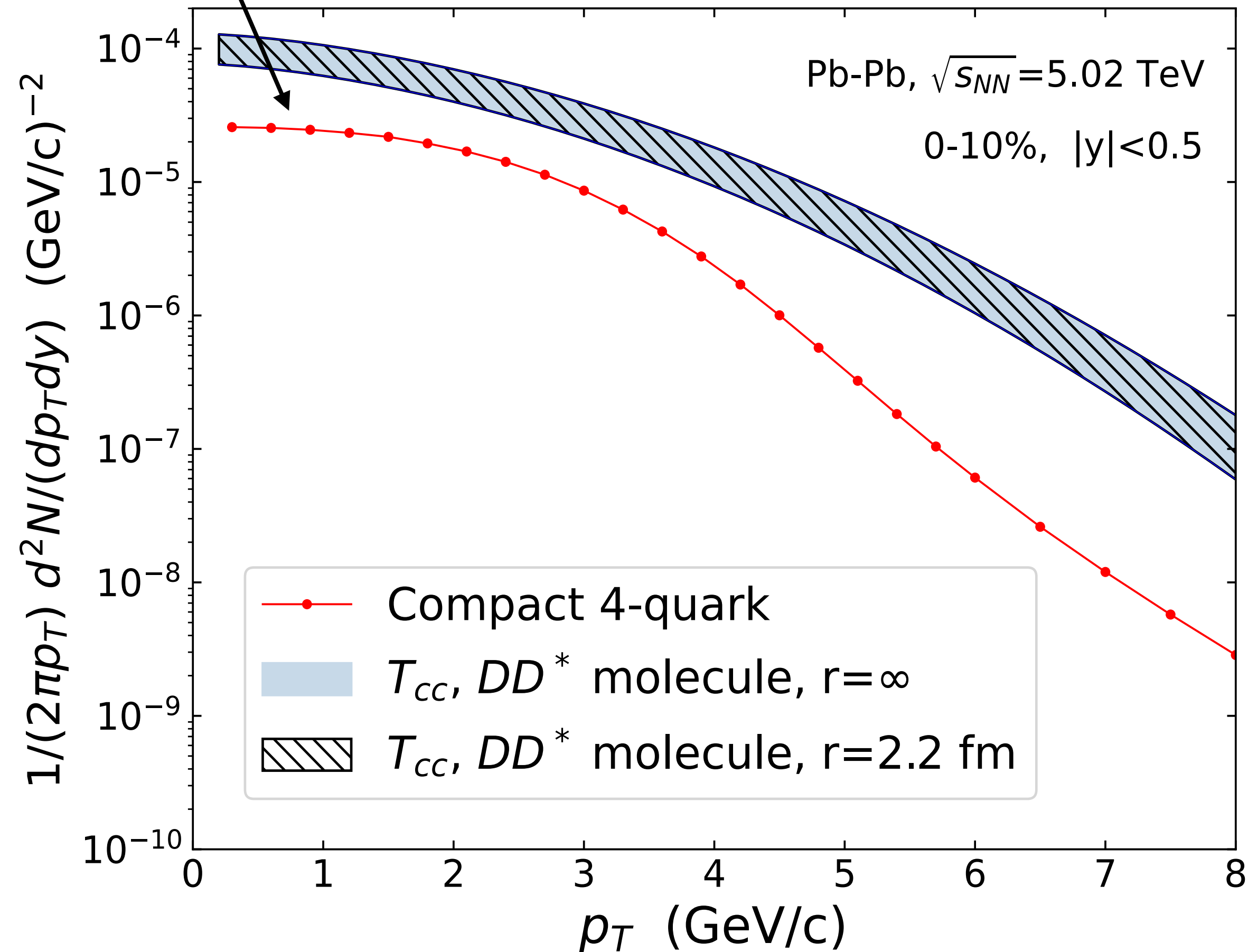
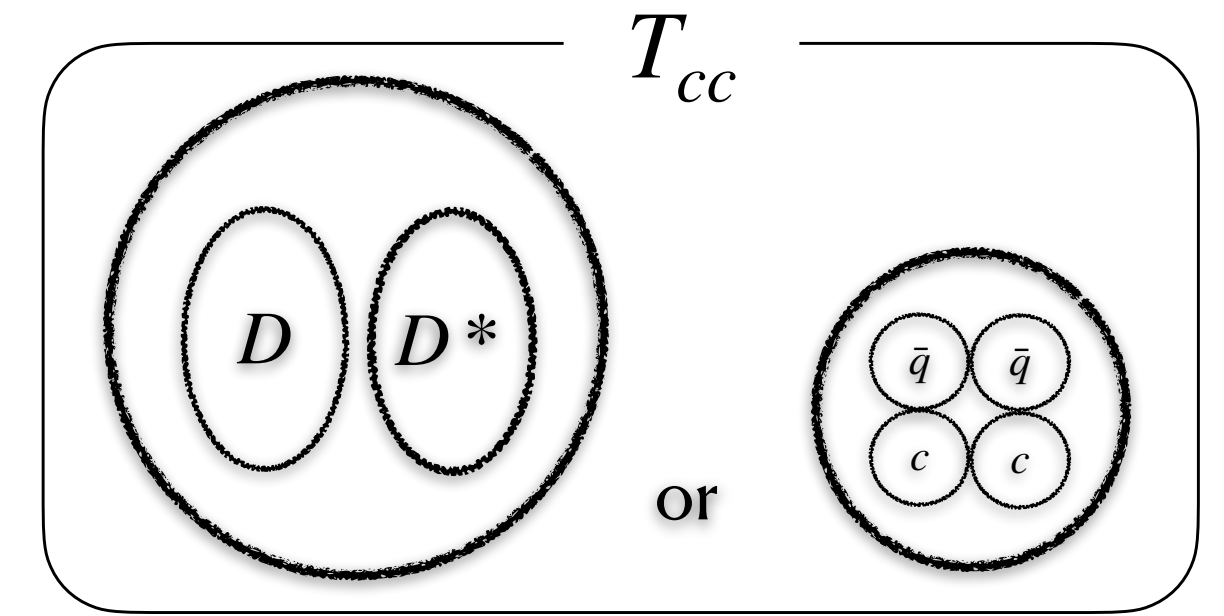
$$(dN_{SHMc}^{\psi(2S)}/dy = 3.04 \times 10^{-3}, \quad N_{SHMc}^{X(3872)}/N_{SHMc}^{\psi(2S)} = 0.326)$$

A. Andronic *et al.* JHEP 07, 035 (2021)

T_{cc} p_T distribution

Pb-Pb collisions at 5.02 TeV

S. H. Lee and S. Cho, Phys. Rev. C 101, 024902
+ Scaling ($\times 1.63, 2.76\text{TeV} \rightarrow 5.02\text{TeV}$)



Radius	T_{cc}
Molecule	2.2fm
Compact 4-quark	0.433fm

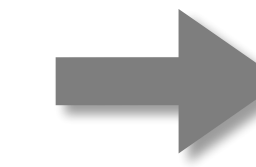
S. Noh, W. Park and S.H. Lee, Phys.Rev. D 103, 114009 (2021)

- Yields

$$dN_{coal}^{DD^*}/dy = (2.47 \pm 0.71) \times 10^{-3}, \quad dN_{coal}^{4q}/dy = 6.2 \times 10^{-4}$$

Two possible configurations of T_{cc} are markedly different

Measurement of the
pt distribution of T_{cc}
in heavy-ion collisions



Confirmation of the
structure of T_{cc}

Summary

- We study the transverse momentum distribution of loosely-bound molecular configuration hadron.
- The $\sigma \rightarrow \infty$ limit coalescence model explained deuteron and helium-3 distribution well.
- We assume that $X(3872)$ is $D\bar{D}$ molecular structures and estimate the transverse momentum distributions using the same formula as deuteron.
- T_{cc} - compact 4-quark state or DD^* molecule. By measuring p_T distribution from heavy ion collisions, the structure of T_{cc} can be confirmed.