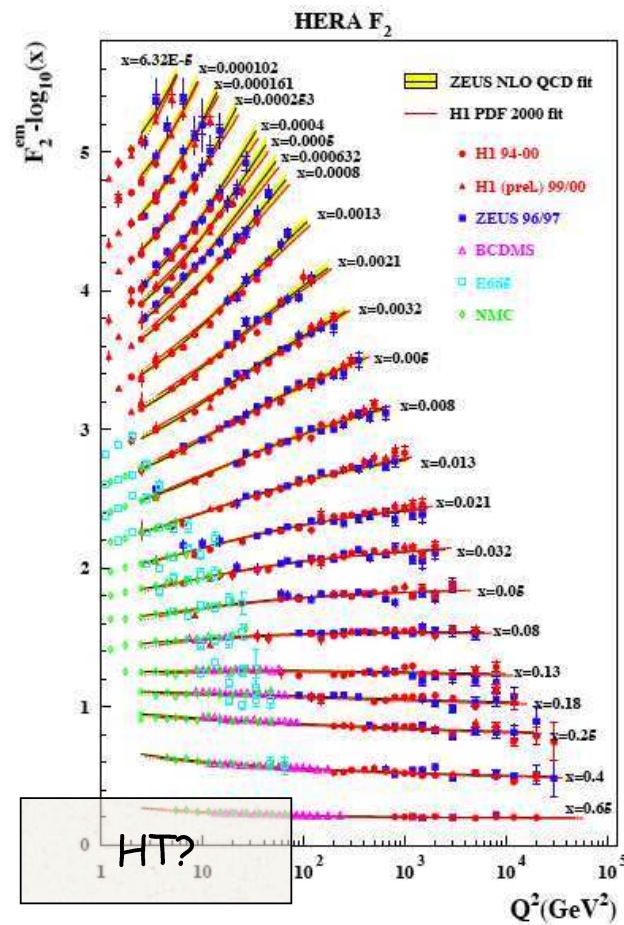
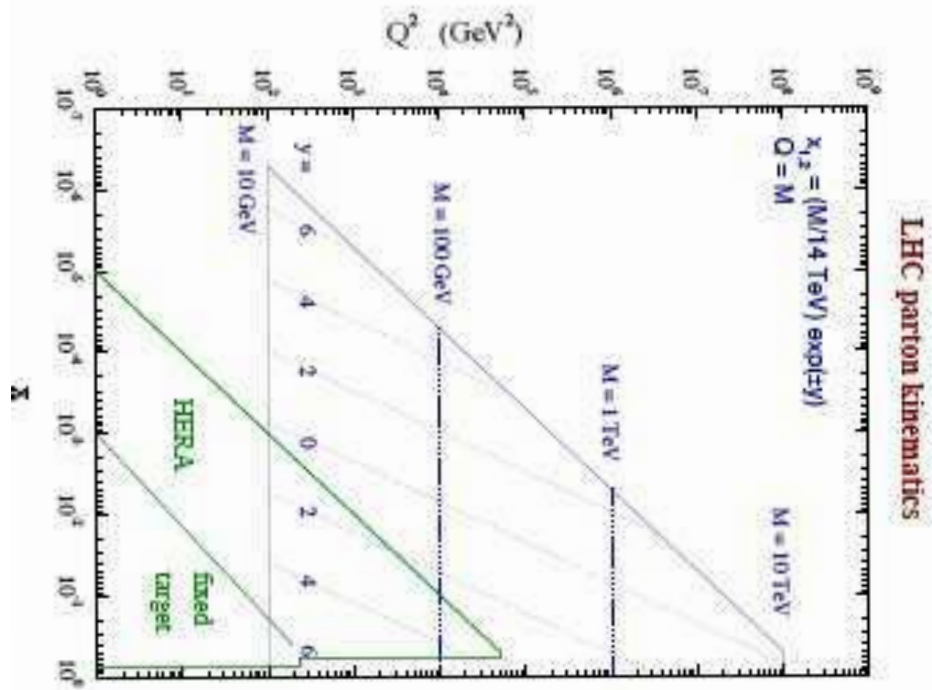


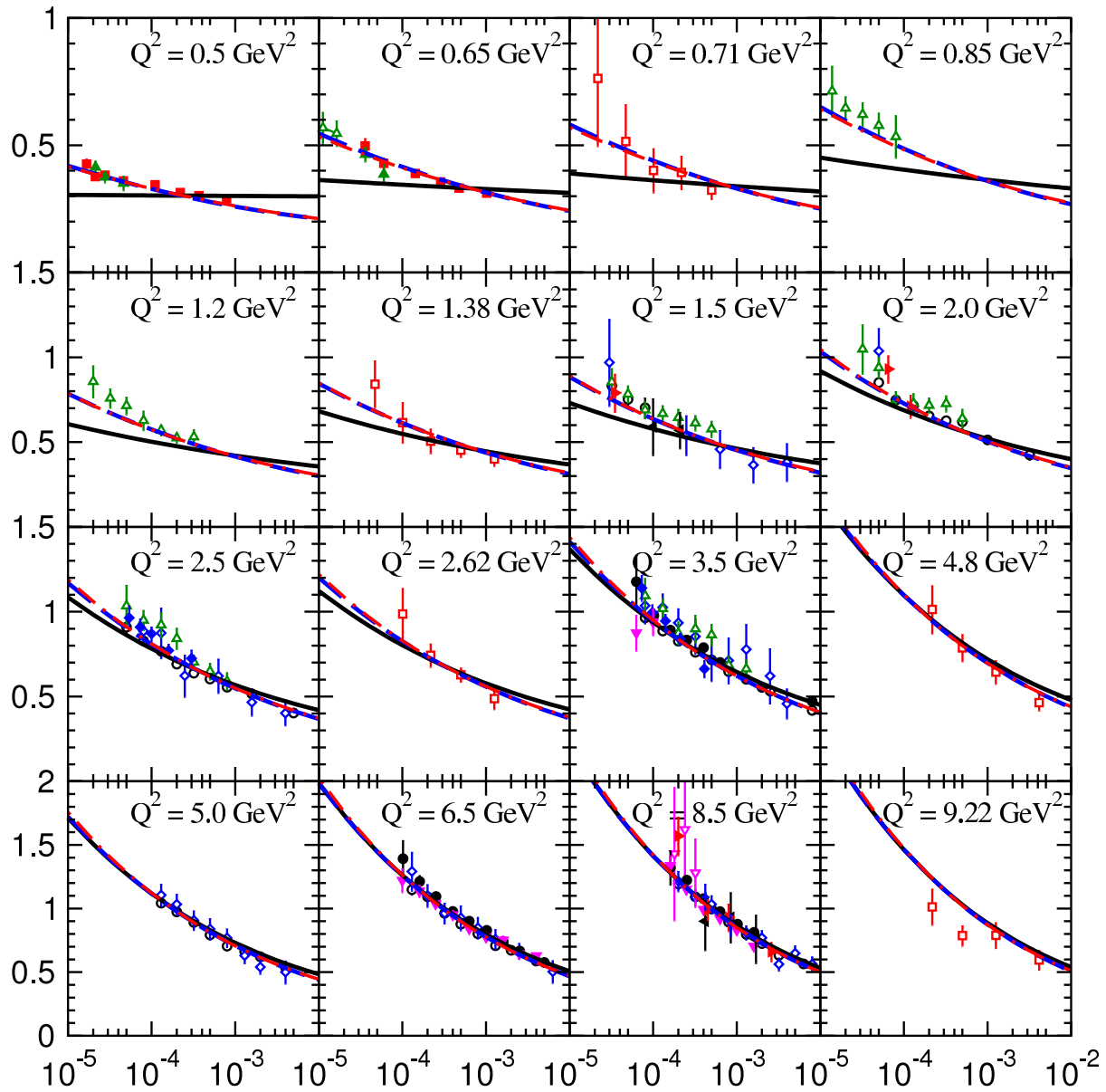
## II: FUTURE PROSPECTS: THEORY AND EXPERIMENT

# FROM HIGHER ORDERS TO RESUMMATION

# PERTURBATIVE SUCCESS!



# FITS INCLUDING HIGHER TWISTS...



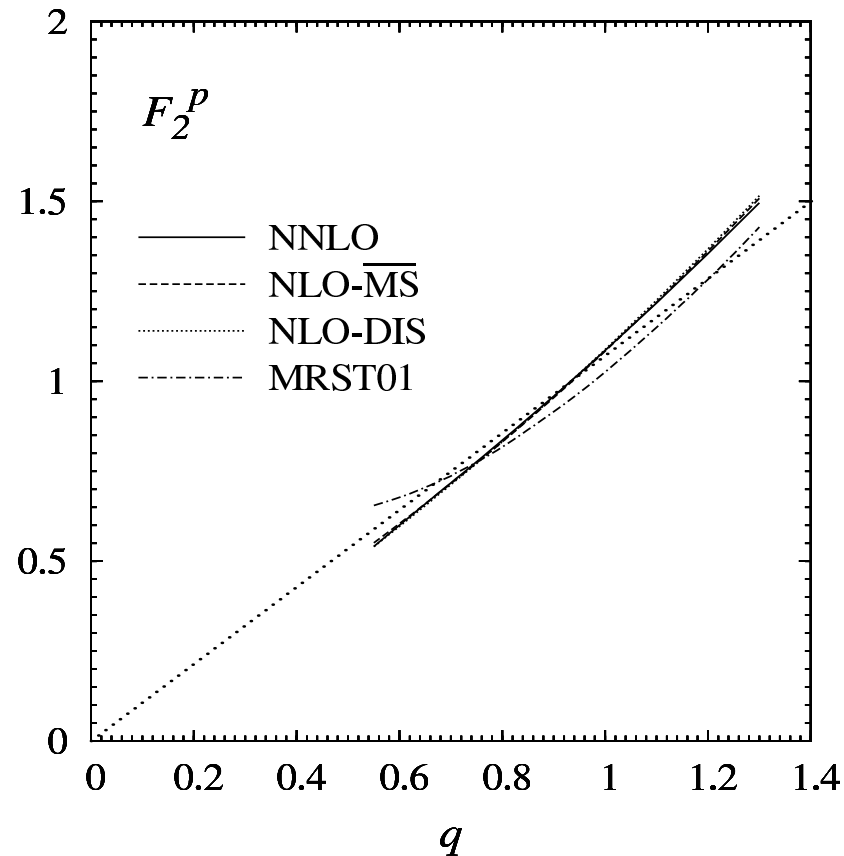
**A. KOTIKOV**

## Curvature test of $F_2$

- At  $x = 10^{-4}$  most measurements lie along a straight (dotted) line, if plotted versus

$$q = \log_{10} \left( 1 + \frac{Q^2}{0.5 \text{ GeV}^2} \right)$$

D. Haidt, EPJ C35, 519 (2004)



...CAN REPRODUCE PECULIAR SCALING OF THE DATA...  
the valence-like input gluon distribution at  $Q_0 = 1 \text{ GeV}$

**C. PISANO**

# ...NNLO MOMENTS OF PDFS COMPARED TO LATTICE...

RESULTS -  $\alpha_s$ ,  $\Lambda_{QCD}$  and PDF moments

## $\alpha_s$ determination

	$\alpha_s(M_Z^2)$	expt	theory
NNLO			
MRST03	0.1153	$\pm 0.0020$	$\pm 0.0030$
A02	0.1143	$\pm 0.0014$	$\pm 0.0009$
SY01(ep)	0.1166	$\pm 0.0013$	
SY01( $\nu N$ )	0.1153	$\pm 0.0063$	
<b>BBG</b>	<b>0.1134</b>	$+0.0019$ $-0.0021$	
<b>World Average</b>	<b>0.1182</b>	<b><math>\pm 0.0027</math></b>	

## PDF moments

$f$	$n$	BBG(NNLO)	MRST04	A02
$u_v$	2	$0.2986 \pm 0.0029$	0.285	0.304
	3	$0.0871 \pm 0.0011$	0.082	0.087
	4	$0.0333 \pm 0.0005$	0.032	0.033
$d_v$	2	$0.1239 \pm 0.0026$	0.115	0.120
	3	$0.0315 \pm 0.0008$	0.028	0.028
	4	$0.0105 \pm 0.0004$	0.009	0.010
$u_v - d_v$	2	$0.1747 \pm 0.0039$	0.171	0.184
	3	$0.0556 \pm 0.0014$	0.055	0.059
	4	$0.0228 \pm 0.0007$	0.022	0.024

## Comparison with lattice results

BBG	Lattice
$N^3LO - \Lambda_{QCD}^{(4)}$ MeV	Alpha Collaboration - $\Lambda_{QCD}^{(2)}$ MeV
$231 \pm 26$	$245 \pm 16 \pm 16$

[M. Della Morte, *et al.*, Nucl.Phys.B713,(2005),378]

$f$	$n$	BBG	Lattice
		NNLO	QCDSF
$u_v - d_v$	2	$0.1747 \pm 0.0039$	$0.191 \pm 0.012$

[G. Schierholz, *private communication*]



# DOES IT WORK ALSO AT SMALL x?

HERA and the LHC- transporting PDFs to hadron-hadron cross-sections

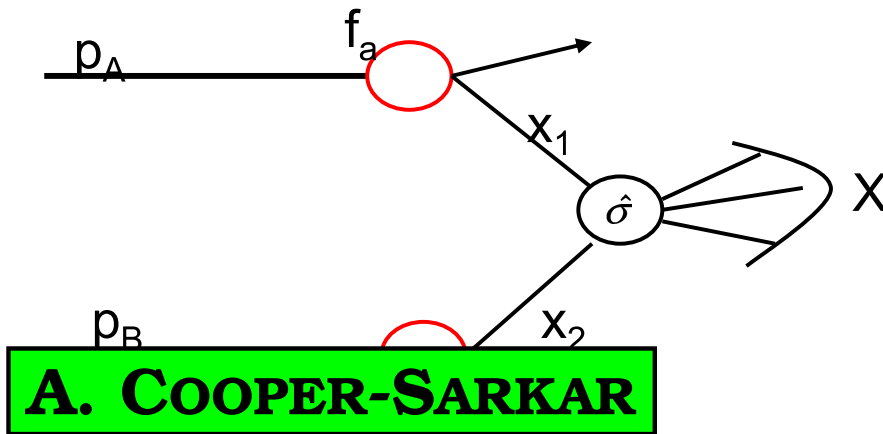
QCD factorization theorem for short-distance **inclusive** processes

$$\sigma_X = \sum_{a,b} \int_0^1 dx_1 dx_2 f_a(x_1, \mu_F^2) f_b(x_2, \mu_F^2) \times \hat{\sigma}_{ab \rightarrow X} \left( x_1, x_2, \{p_i^\mu\}; \alpha_S(\mu_R^2), \alpha(\mu_R^2), \frac{Q^2}{\mu_R^2}, \frac{Q^2}{\mu_F^2} \right)$$

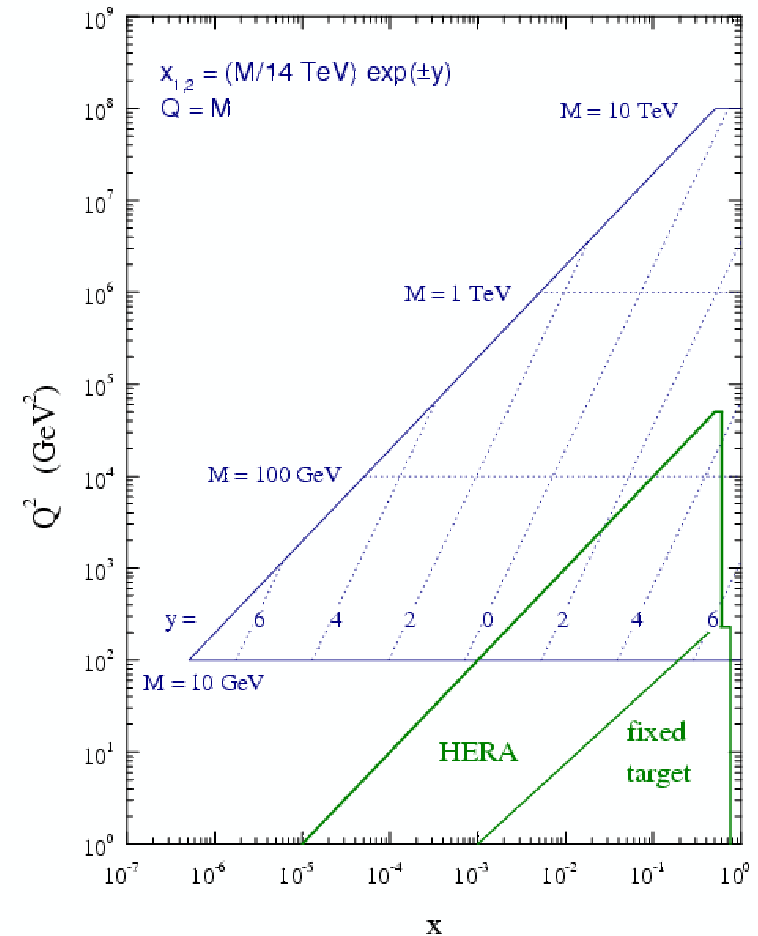
where  $X=W, Z, D\text{-}Y, H, \text{high-}E_T \text{ jets, } \hat{p} \text{ prompt-}\gamma$

and  $\sigma$  is known

- to some fixed order in pQCD and EW
- in some leading logarithm approximation (LL, NLL, ...) to all orders via resummation



LHC parton kinematics

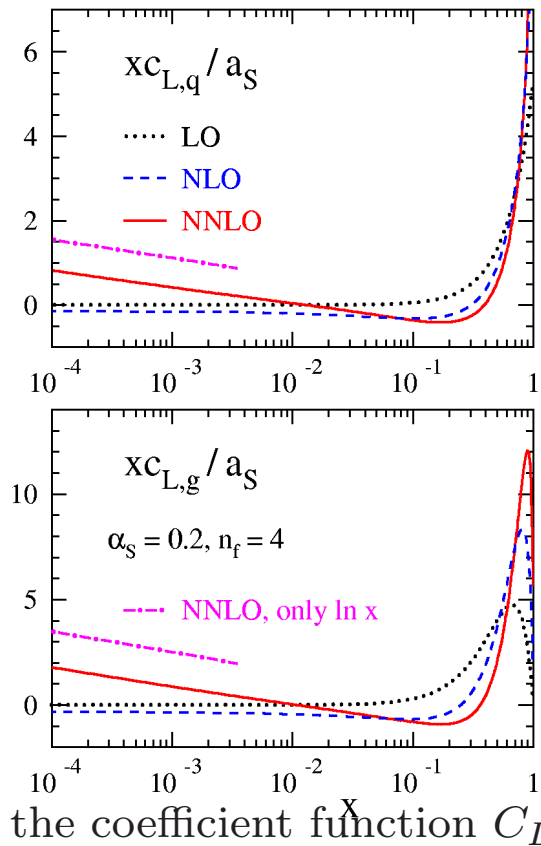


The central rapidity range for W/Z production AT LHC is at low-x ( $5 \times 10^{-4}$  to  $5 \times 10^{-2}$ )

Knowledge of the PDFs is vital

# A PROBLEM: NNLO CORRECTIONS PERTURBATIVE INSTABILITY AT SMALL $x$ THEORY

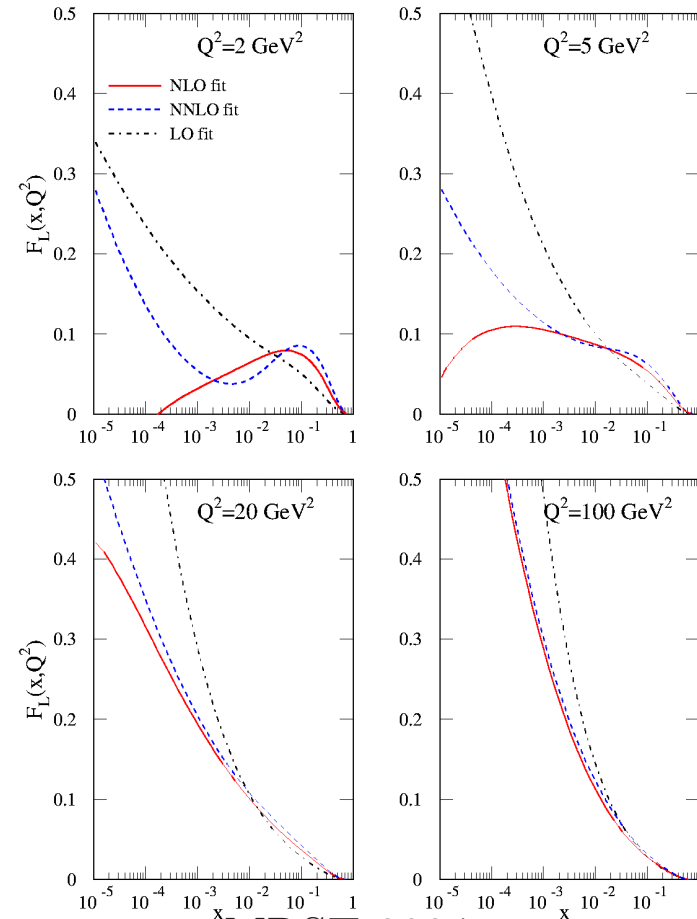
- Perturbation theory unstable
- leading log approx no good



Moch, Vermaseren, Vogt 2005

# PHENOMENOLOGY $F_L$ FIT

$F_L$  LO, NLO and NNLO



MRST 2004

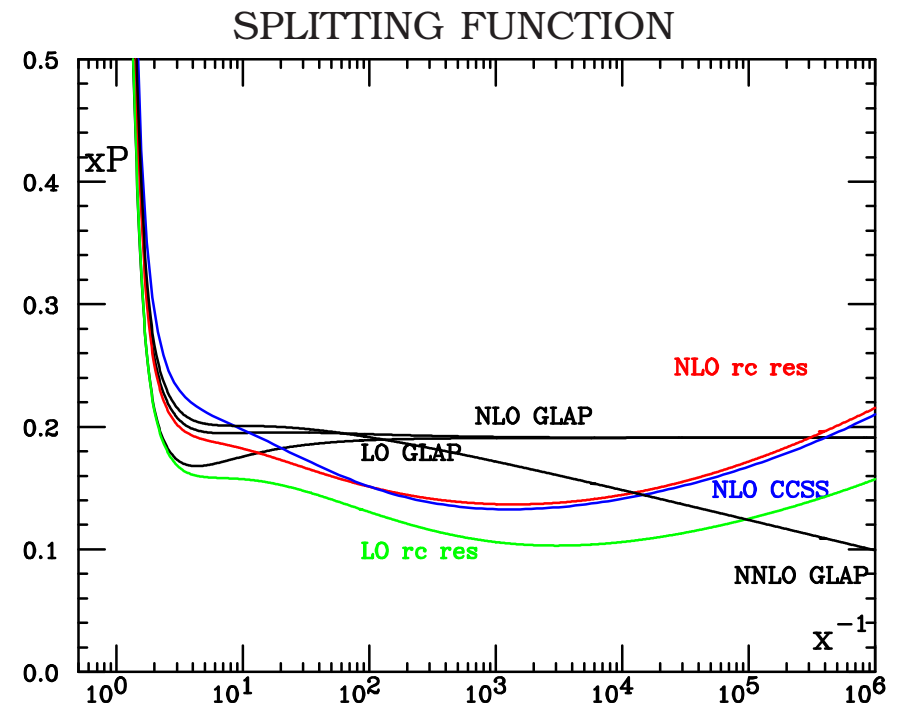
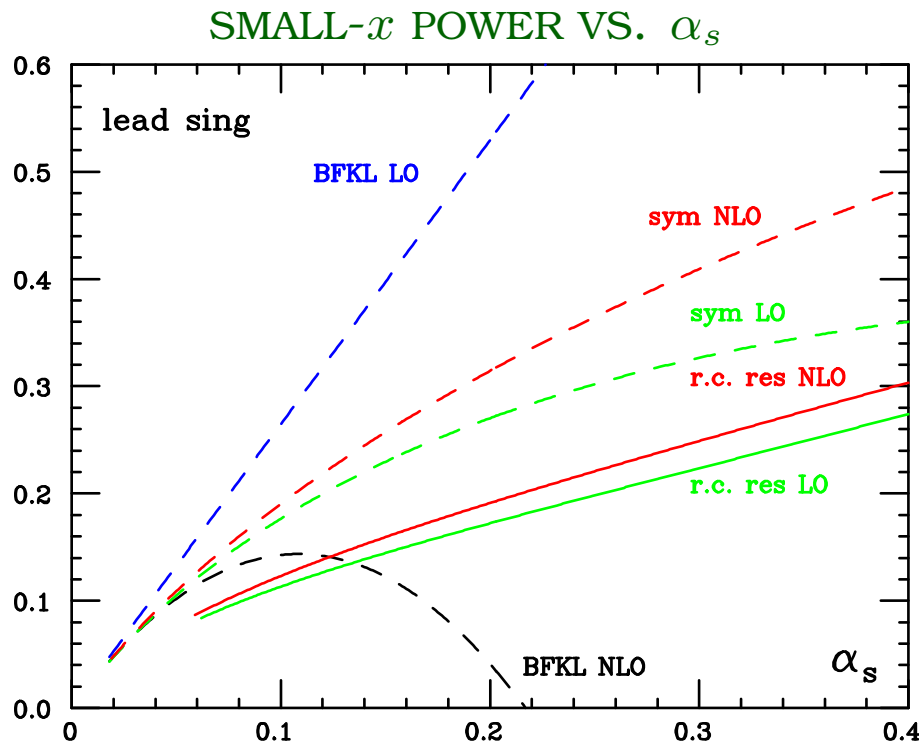


# SOLUTION: RESUMMATION:

## $n_f = 0$ QUALITATIVE FEATURES

### HERALHC 2005

SINGULARITY IN ANOM. DIM. AT  $N = \alpha \Rightarrow$  ASYMPT. SMALL- $x$  POWER  $G \sim x^{-\alpha}$



- RESUMMED EXPANSION CONVERGES RAPIDLY
- SMALL  $x$  INTERCEPT & CURVATURE DETERMINE RESUMMED BEHAVIOUR
- RESUMMED RESULTS OBTAINED BY DIFFERENT GROUPS AGREE WELL

# PROGRESS: EXTENSION TO QUARKS (CCSS)

## The Matrix Kernel

$$\mathcal{K}_0 = \begin{pmatrix} \Gamma_{qq}^0(\omega)\chi_c^\omega(\gamma) & \Gamma_{qg}^0(\omega)\chi_c^\omega(\gamma) + \Delta_{qg}(\gamma, \omega) \\ \Gamma_{gq}^0(\omega)\chi_c^\omega(\gamma) & [\Gamma_{gg}^0(\omega) - \frac{1}{\omega}]\chi_c^\omega(\gamma) + \frac{1}{\omega}\chi_0^\omega(\gamma) + \Delta_{gg}(\gamma, \omega) \end{pmatrix}$$

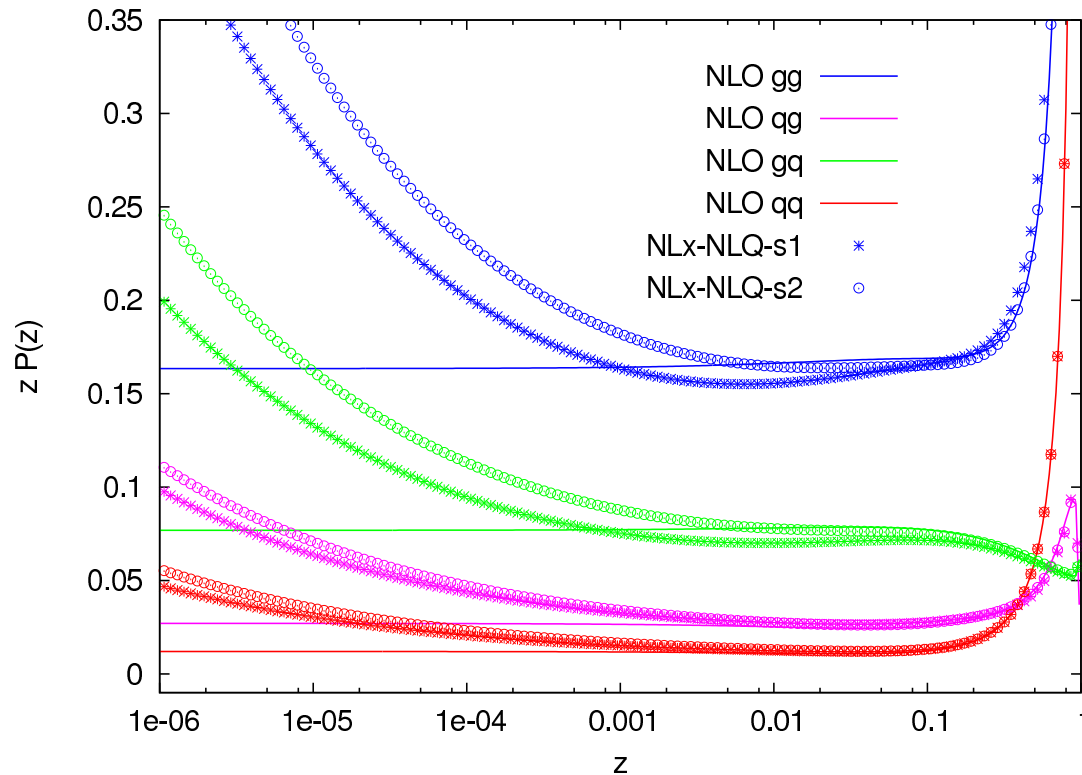
- $\mathcal{K}_0$  has **at most simple poles** in  $\gamma$  and simple poles in  $\omega$  in the gluon row
- **No  $\omega$ -poles** are present **in the quark row**, consistently with LO DGLAP and reggeization of the quark at  $\omega = -1$ ;  
We'll keep this structure also in  $\mathcal{K}_1$
- At NLO  $\Gamma_{qq}^1$  and  $\Gamma_{qg}^1$  contain  $\frac{\bar{\alpha}_s^2}{\omega}$ . Instead of adding such terms in  $\mathcal{K}_1$  (see above) we add a proper non-singular  $\Delta_{qg}(\gamma, \omega)$  term
- **Momentum Sum Rule**: restored by adding a non-singular subleading  $\Delta_{gg}(\gamma, \omega)$  term
- $\mathcal{K}_1$ : we add **NLO** DGLAP matrix  $\Gamma_1$  and **NLL $x$**  BFKL kernel  $\chi_1$  in  $\mathcal{K}_{1,gg}$  with subtractions to avoid double-counting
- **Running coupling**: introduced in  $(\mathbf{k}, x)$  space (analytic double inverse Mellin transf.)

$$\mathcal{K}(\mathbf{k}, \mathbf{k}'; x) = \bar{\alpha}_s(\mathbf{k}_>^2)\mathcal{K}_0(\mathbf{k}, \mathbf{k}'; x) + \bar{\alpha}_s^2(\mathbf{k}_>^2)\mathcal{K}_1(\mathbf{k}, \mathbf{k}'; x)$$

( $\mathcal{K}_1$  depends on the choice of run.coupl. scale  $\mathbf{k}_> \equiv \max(\mathbf{k}, \mathbf{k}')$ )

# Frozen Coupling Features

Resummed  $\overline{\text{MS}}$  splitting functions  $z P(\alpha_s, z)$  for  $\alpha_s = 0.2$



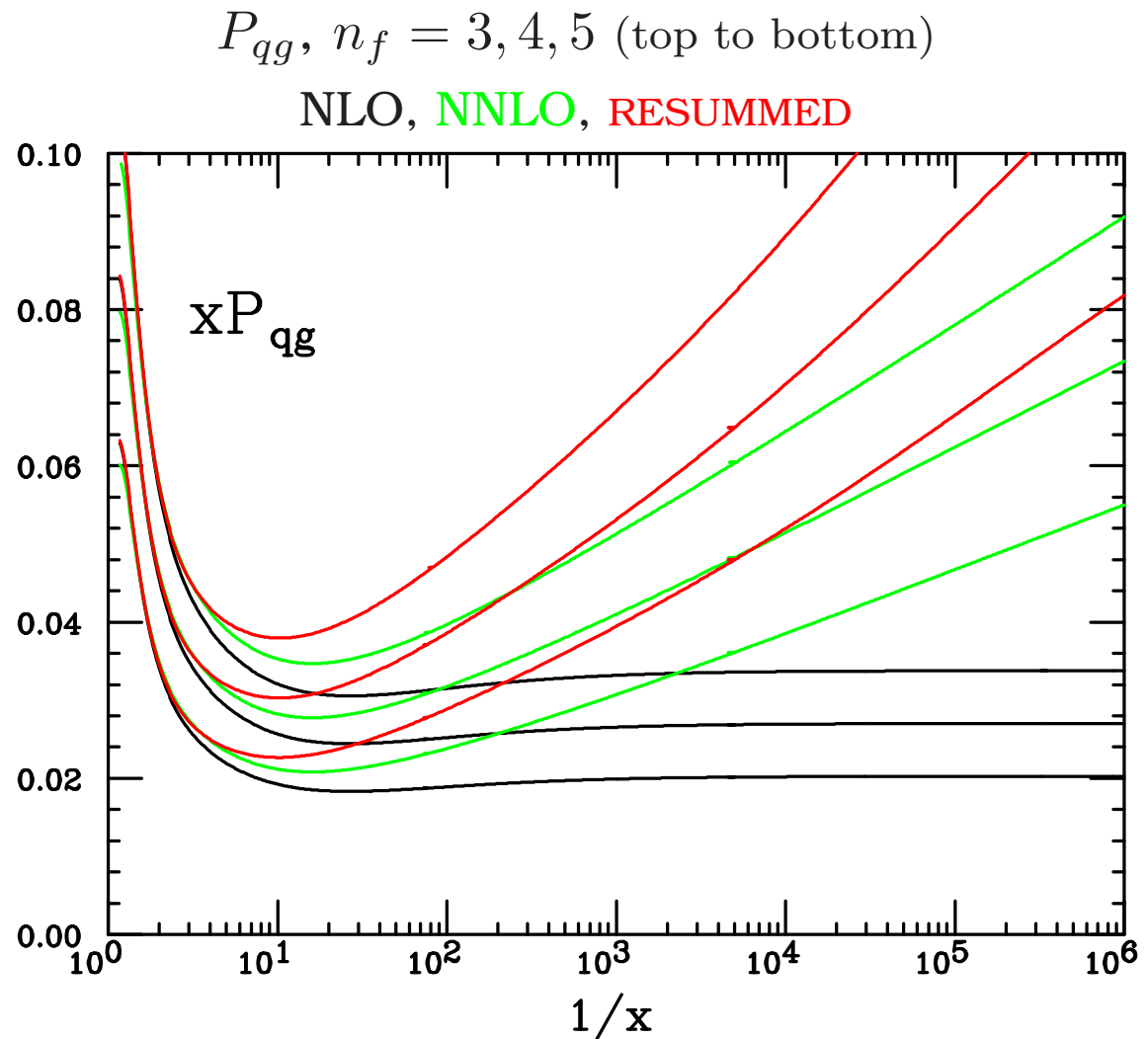
- at large- $x$  fixed order and resummed splitting functions overlap
- at moderate- $x$  resummed splitting functions show a small dip
- final rise sets in at very small- $x$
- resummation scheme uncertainty is small

# EXTENSION TO QUARKS (ABF)

USE  $\mathcal{Q}_0\overline{\text{MS}}$  SCHEME: (Ball, Forte, 2005)

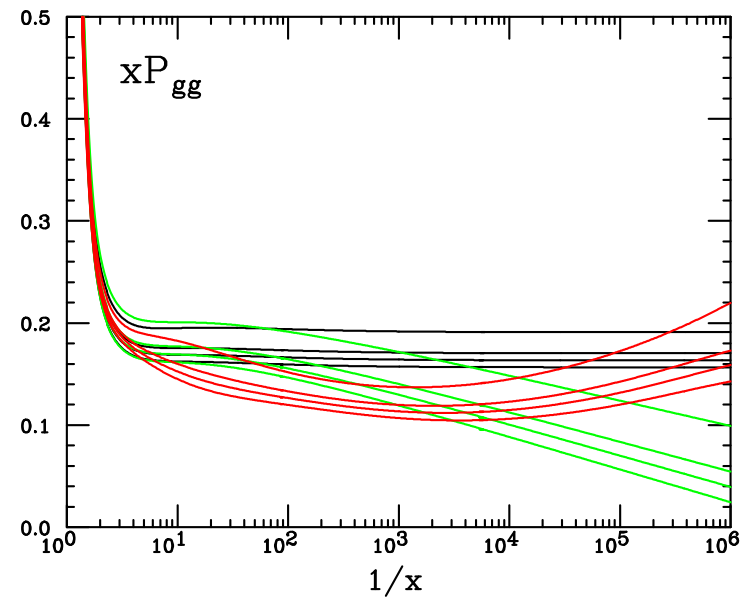
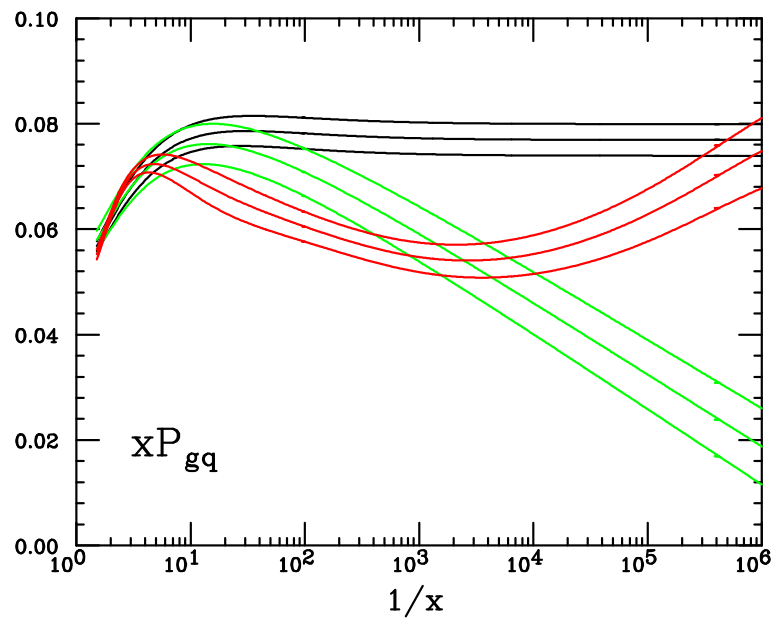
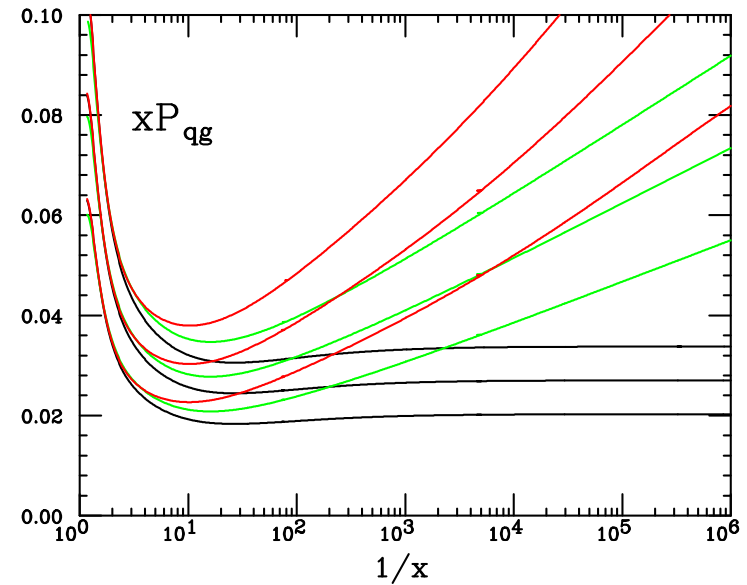
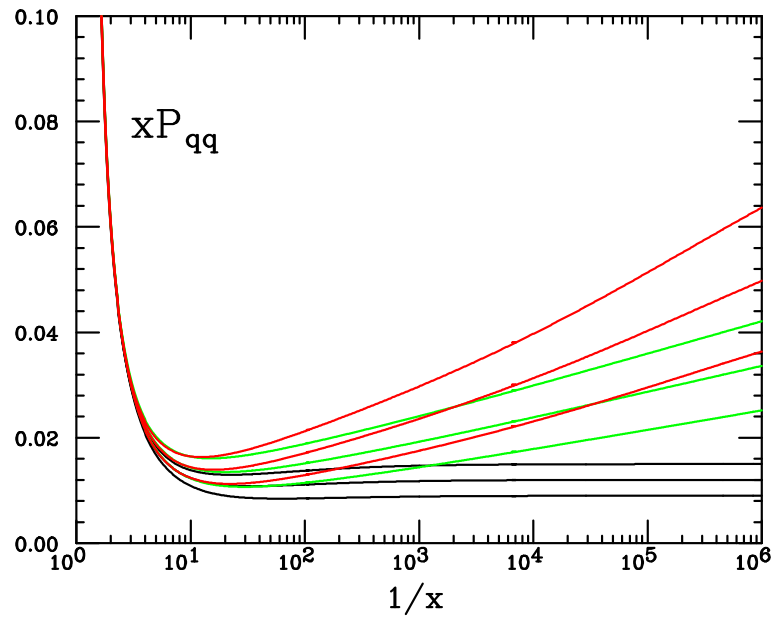
COINCIDES WITH  $\overline{\text{MS}}$  AT LARGE  $x$  (NLO) BUT SMALL  $x$  R.C. SINGULARITIES IN GLUON SECTOR

- $\gamma_{qg}$  SAME AS IN  $\overline{\text{MS}}$
- $\gamma_+$  SAME AS IN  $\mathcal{Q}_0$
- CAN DETERMINE RESUMMED  $\gamma_{qg}$  (Catani & Hautmann)



# THE SPLITTING FUNCTION MATRIX (ABF)

NLO, NNLO, RESUMMED



# QUARK AND GLUON EVOLUTION (ABF)

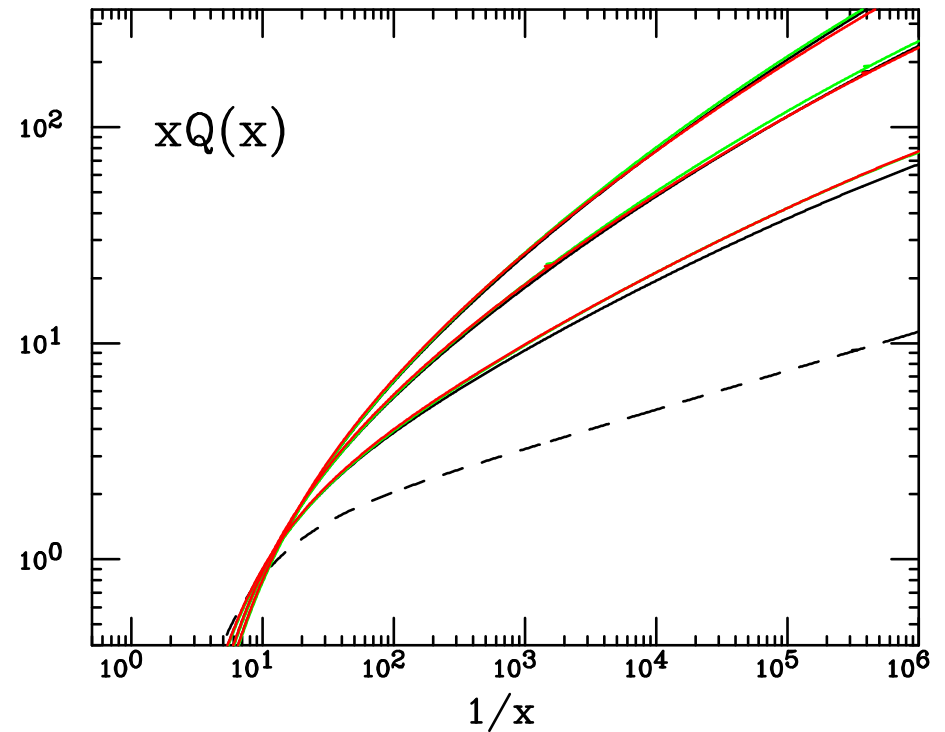
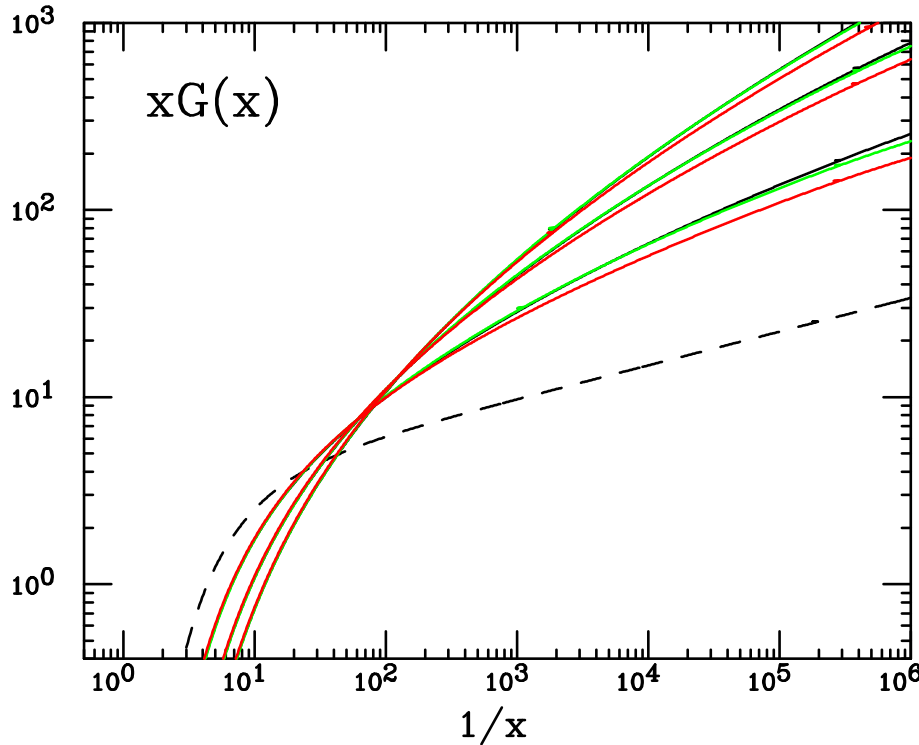
evolve toy  $G = (x, Q_0) = x^{-0.18}(1-x)^5$ ,  $Q(x, Q_0) = \frac{1}{3}G(x, Q_0)$ ,  $Q_0 = 2$  GeV

GLUON

QUARK

$Q = 2, 10, 100, 1000$  GeV (bottom to top)

NLO, NNLO, RESUMMED

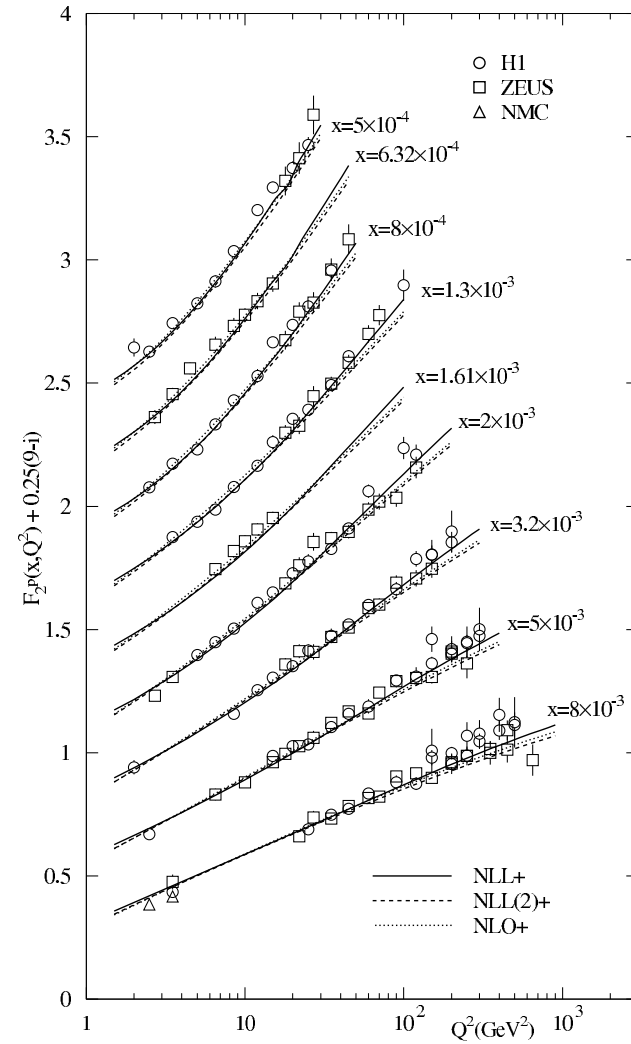
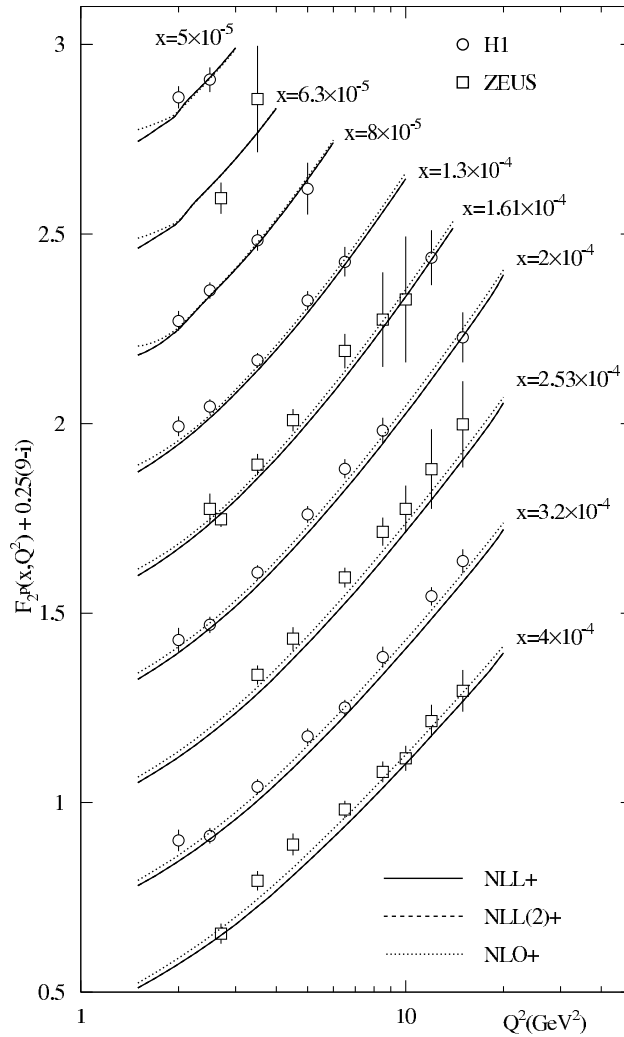


- LO VS NLO DIFFERENCE LARGER THAN FIXED VS RESUMMED
- RESUMMED GLUON BELOW UNRESUMMED, QUARK JUST BELOW

# A FIRST RESUMMED FIT (THORNE-WHITE)

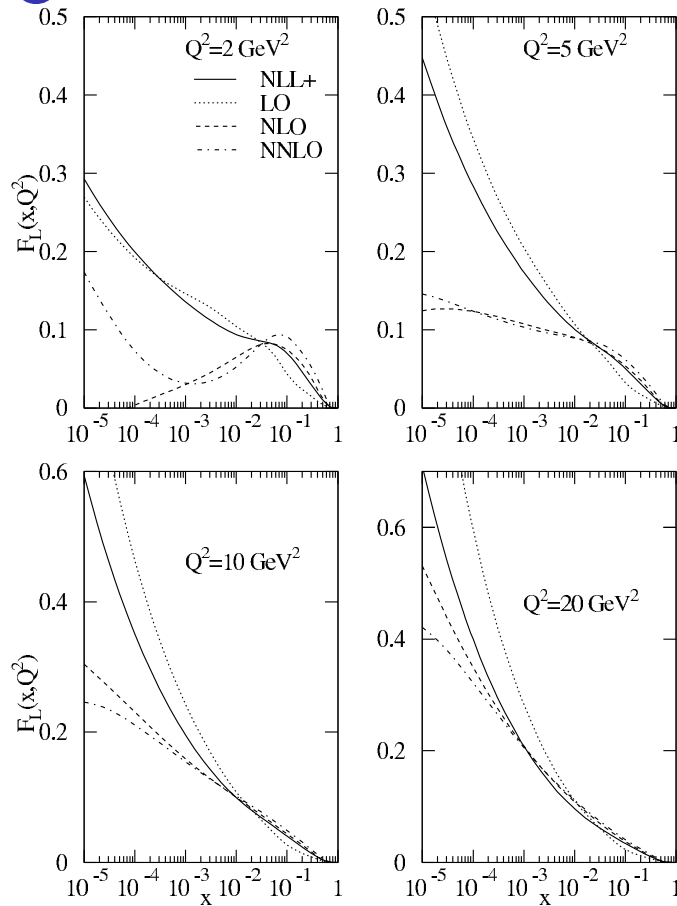
## Global Fit - Results

### Results - $F_2$



**C. WHITE**

## Longitudinal Structure Function

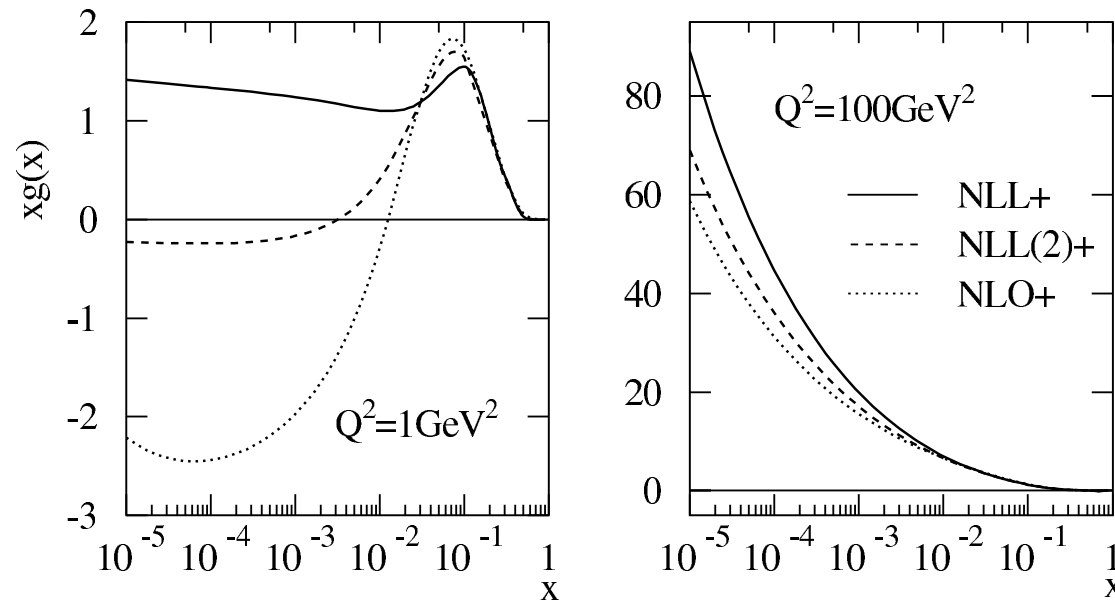


- ▶ Clearly see perturbative instability in fixed order results.
- ▶ This is cured by the resummation.

**NNLO INSTABILITY REMOVED FROM  $F_L$**



## Glucn Distribution



- ▶ Gluons differ for  $x \lesssim 10^{-2}$ .
- ▶ NLL resummed gluon positive and growing at small  $x$ !
- ▶ Not true at fixed order

**HIGHER INITIAL GLUON (LESS EVOLUTION)**

- ▶ See this in e.g.  $F_L$ ...

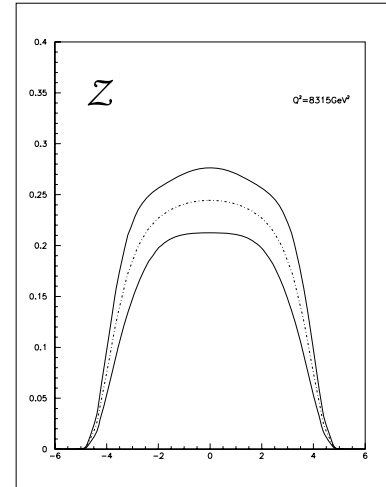
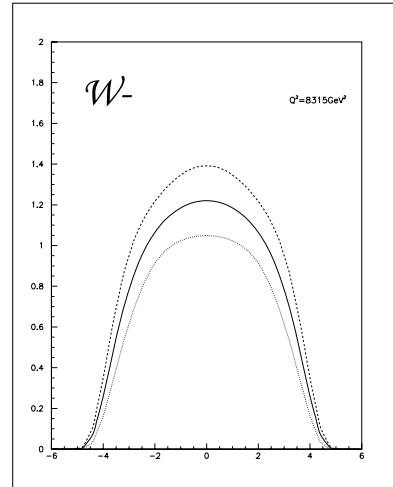
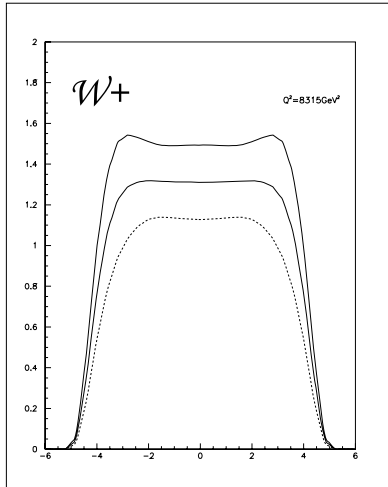
# RESUMMATION: SUMMARY

- FULL RESUMMED EVOLUTION SOON AVAILABLE IN VARIOUS APPROACHES
- DIFFERENT APPROACHES MUST BE BENCHMARKED
- NNLO INSTABILITY REMOVED
- SOME CRITICAL ASPECTS OF FITS MIGHT IMPROVE (NEGATIVE GLUON?)

TOWARDS LHC

# EXPERIMENTAL SUCCESS!

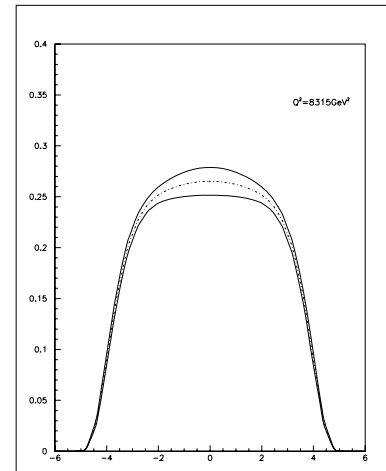
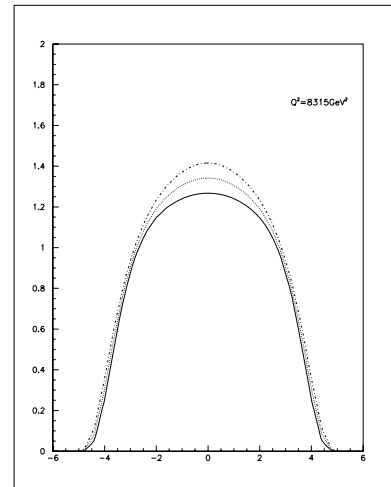
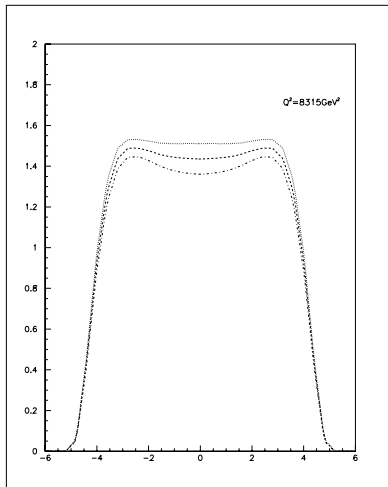
What has HERA data ever done for us?



Pre-HERA  $W^+/W^-/Z$  rapidity spectra  $\sim \pm 15\%$  uncertainties become!

**NO WAY** to use these cross-sections as a good luminosity monitor

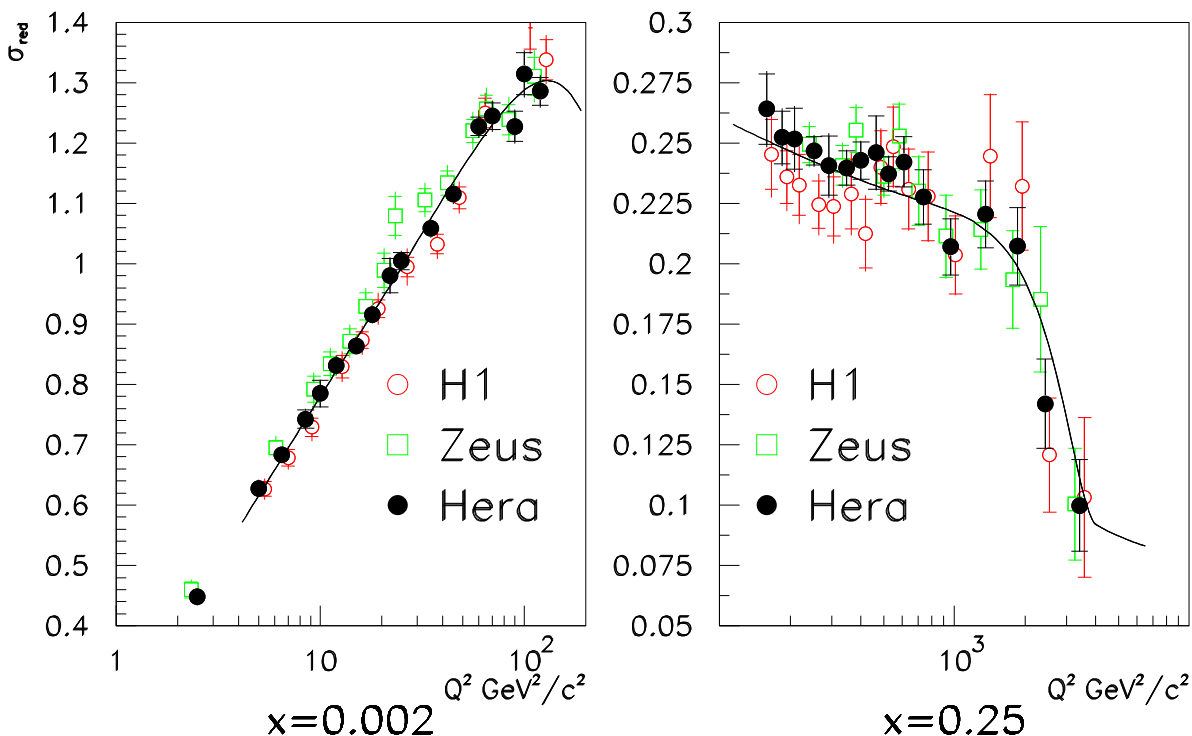
Post-HERA  $W^+/W^-/Z$  rapidity spectra  $\sim \pm 5\%$  uncertainties



**A. COOPER-SARKAR**

# ...STILL NOT OVER!

## AVERAGING OF HERA DATA



Cross calibration of systematic uncertainties leads to better than  $1/\sqrt{2}$  improvement for systematic errors dominated regions. Simple  $1/\sqrt{2}$  improvement for stat. error dominated domain.

# CAN WE USE HERA/TEVATRON WISDOM @ LHC?

- W/Z cross sections serve as precision physics monitors

- ◆ all cross sections at Tevatron/LHC could be normalized to W/Z
- ◆ Tevatron is a *W factory*

W Factory	Mode	Events/Week/Exp. (after trigger & cuts)
	$W \rightarrow ev$	~ 15,000
	$Z \rightarrow ee$	~ 1,500

- Both experimental and theoretical errors are under control

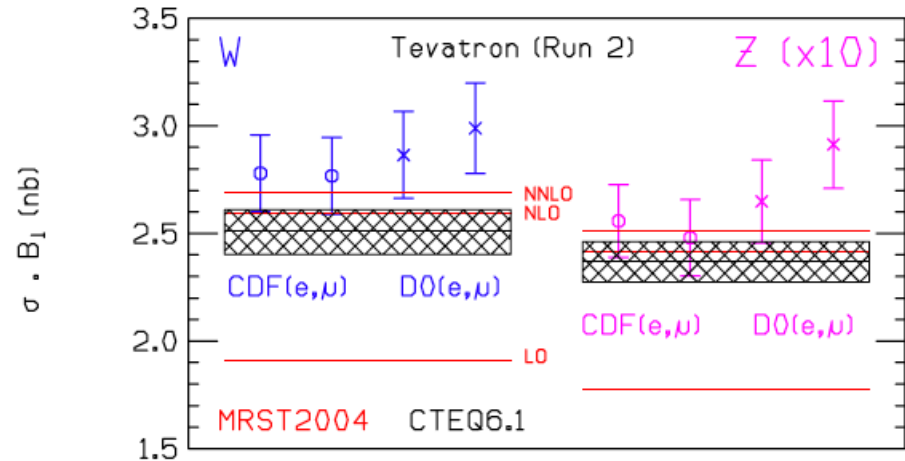
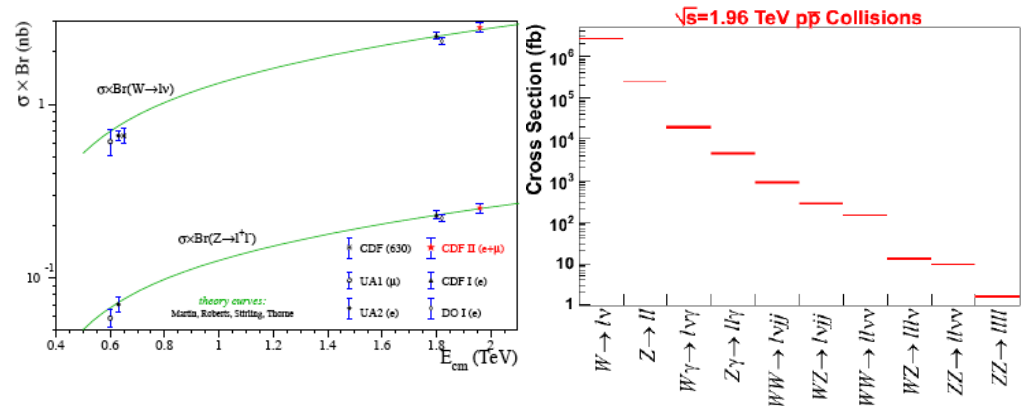
- ◆ NNLO a small (positive) correction to NLO

$$\frac{\sigma(L)}{L} = 2.5\% \oplus 5.5\%$$

$\sigma_{TOT}(p\bar{p})$

$\sigma_{EXP}$  (lumi detector acceptance)

dominated by experimental uncertainty



- Note that CTEQ and MRST NLO predictions agree within CTEQ6.1 pdf errors (but MRST at edge of CTEQ6.1 error band)

**J. HUSTON**

# Precision benchmarks: W/Z cross sections at the LHC

- CTEQ and MRST NLO predictions in good agreement with each other
- NNLO corrections are small and negative
- NNLO mostly a K-factor; NLO predictions adequate for most predictions at the LHC

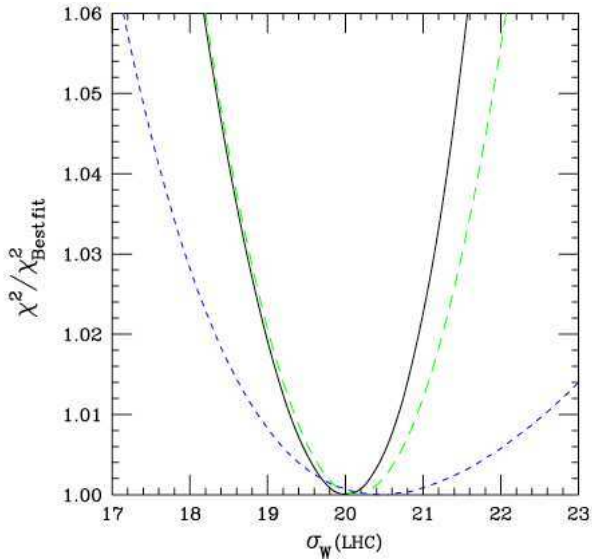


Figure 82. Lagrange multiplier results for the  $W$  cross section (in nb) at the LHC using a positive-definite gluon. The three curves, in order of decreasing steepness, correspond to three sets of kinematic cuts, standard/intermediate/strong.

removing low x data from global fits increases uncertainty but does not significantly move central answer

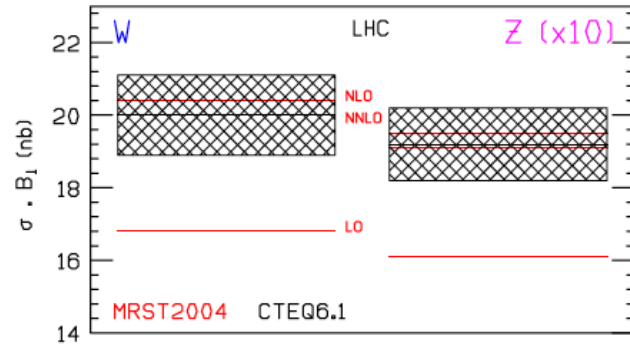
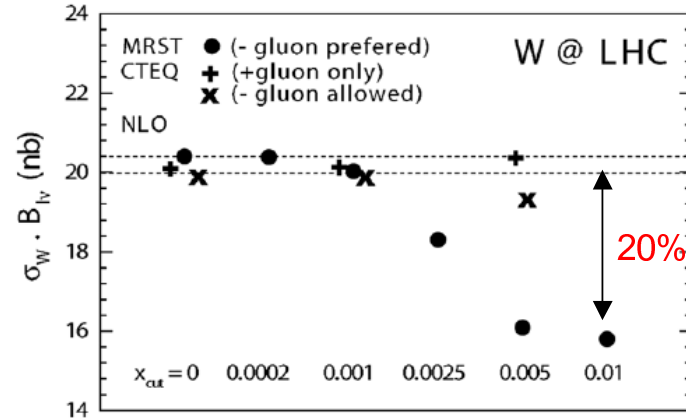


Figure 80. Predicted cross sections for  $W$  and  $Z$  production at the LHC using MRST2004 and CTEQ6.1 pdfs. The overall pdf uncertainty of the NLO CTEQ6.1 prediction is approximately 5%, consistent with figure 77.



MRST found a tension between low x and high x data; not present in CTEQ analysis

Figure 81. Predicted total cross section of  $W^+ + W^-$  production at the LHC for the fits obtained in the CTEQ stability study, compared with the MRST results. The overall pdf uncertainty of the prediction is  $\sim 5\%$ , as observed in figure 77.

**J. HUSTON**

# EXAMPLE OF APPLICATION

## LUMINOSITY MEASUREMENT

$$pp \rightarrow W \rightarrow \ell \nu$$

$$pp \rightarrow Z \rightarrow \ell^+ \ell^-$$

Clean signature of leptonic final state.  
 High rate: O (10 Hz) at  $L = 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ .  
 $\Delta L/L = 1\%$  in 20 min.

The increasing precision of the QCD calculation of W and Z production cross section makes the counting of the W and Z leptonic decays an attractive means to measure the luminosity at the LHC (ref. G. Polesello)

$$L = \frac{N_{W/Z} - N_{background}}{A \cdot \varepsilon \cdot \sigma_{W/Z}}$$

$N_{W/Z}$	Number of identified bosons
$N_{background}$	Number of background events
A	Geometrical acceptance
$\varepsilon$	Reconstruction efficiency
$\sigma_{W/Z}$	Theoretical cross section

Largest uncertainties: theoretical estimate of  $\sigma_{W,Z}$  and detector model

**A. SBRIZZI**



# Uncertainty on luminosity from W/Z decays

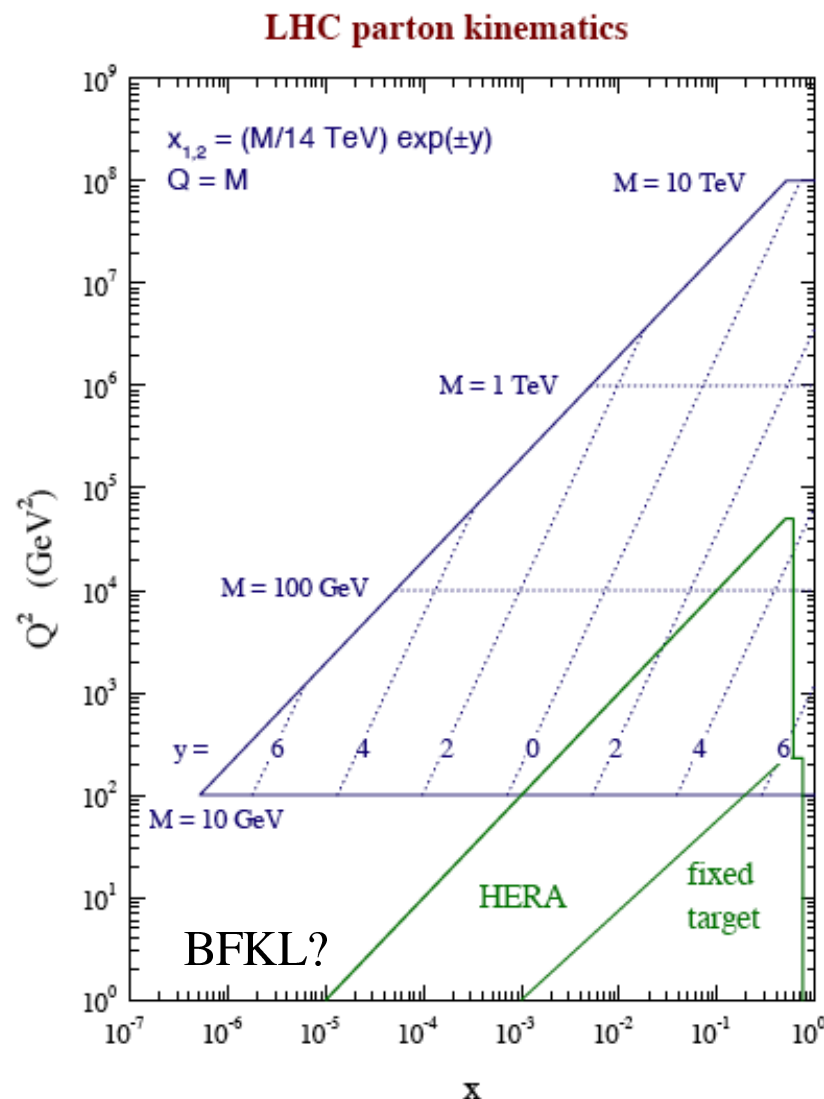
In  $pp$  collision,  $\sigma_{W,Z}$  is the convolution of the PDF and the partonic cross section

$$\sigma_{W/Z} = \sum_{i,j} \int dx_i dx_j f_1(x_i) f_2(x_j) \hat{\sigma}_{ij \rightarrow W/Z} (S_{\max} x_i x_j)$$

- The partonic cross section is known at NNLO at 1% level (PRD 094008)
- The PDF contribution controversial (commonly set at 3%) (hep-ph/0307219)
- Currently available tools are NLO + parton shower (see also CMS 2004/056)
- Uncertainty on background simulation to be studied
  - QCD and heavy quarks (difficult to simulate, real data needed)
  - Leptonic decays of top quark pairs
  - $W \rightarrow \tau \rightarrow l$  decays
- Uncertainty on acceptance determination at 2% level (hep-ph/0405130)
  - Variation of NLO QCD scale 1%
  - PDF contribution 1%
  - Improvement are expected by better fixing PDF by looking at W/Z data
  - EW effect of multiple photon radiation to be studied

# A WORD OF CAUTION

- Experience at the Tevatron is very useful, but scattering at the LHC is not necessarily just “rescaled” scattering at the Tevatron
- Small typical momentum fractions  $x$  in many key searches
  - ◆ dominance of gluon and sea quark scattering
    - ▲ where HERA experience comes in handy
  - ◆ large phase space for gluon emission and thus for production of extra jets
  - ◆ intensive QCD backgrounds
  - ◆ or to summarize,...lots of Standard Model to wade through to find the BSM pony



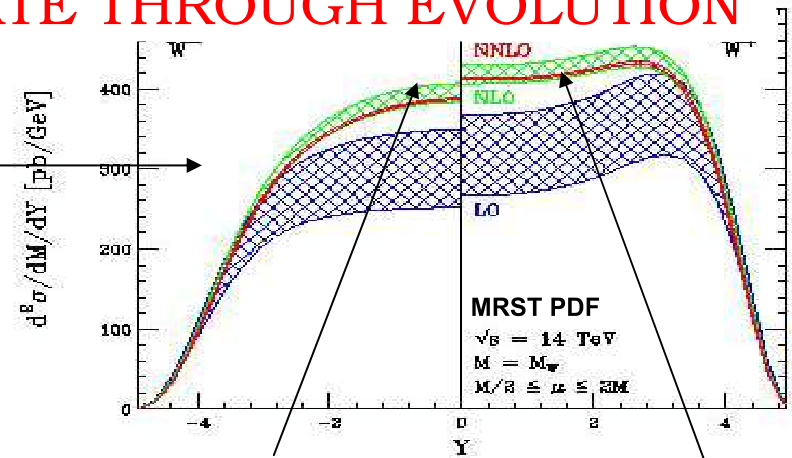
**J. HUSTON**

# PARTON UNCERTAINTIES PROPAGATE THROUGH EVOLUTION

W/Z production have been considered as good standard candle processes with small theoretical uncertainty.

BUT- there are also QED effects to be considered of a similar size to NNLO QCD

PDF uncertainty has been considered as a dominant contribution and most PDF groups quote uncertainties  $< \sim 5\%$

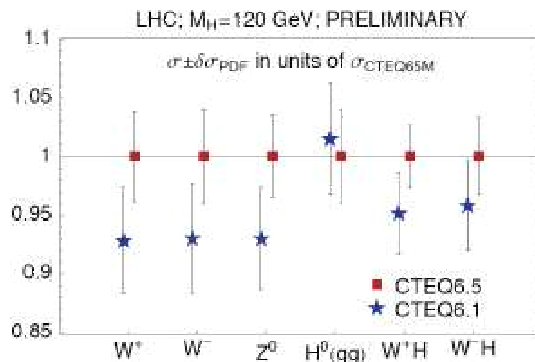


NNLO corrections small  $\sim$  few%  
 NNLO residual scale dependence  $< 1\%$

PDF Set	$\sigma_{W^+} \cdot B_{W \rightarrow \nu}$ (nb)	$\sigma_W \cdot B_{W \rightarrow \nu}$ (nb)	$\sigma_Z \cdot B_{Z \rightarrow ll}$ (nb)
ZEUS-S	$12.07 \pm 0.41$	$8.76 \pm 0.30$	$1.89 \pm 0.06$
CTEQ6.1	$11.66 \pm 0.56$	$8.58 \pm 0.43$	$1.92 \pm 0.08$
MRST01	$11.72 \pm 0.23$	$8.72 \pm 0.16$	$1.96 \pm 0.03$

BUT the central values differ by more than some of the uncertainty estimates.

AND the situation just got dramatically worse. The new CTEQ6.5 estimate is 8% higher



MRST-CTEQ QUARK AGREEMENT AT LOW SCALE

$\Rightarrow$  DISAGREEMENT AT HIGH SCALE

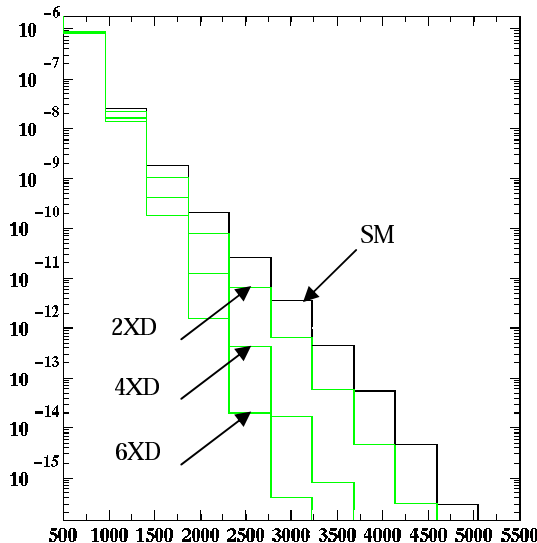
monitor

**A. COOPER-SARKAR**

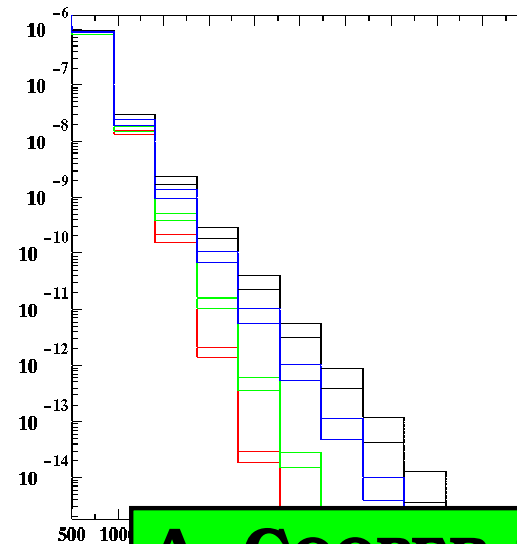
# LHC LIMITED BY PDFS (J. HUSTON)

Such PDF uncertainties the jet cross sections compromise the LHC potential for discovery.  
E.G. Dijet cross section potential sensitivity to compactification scale of extra dimensions ( $M_c$ ) reduced from  $\sim 6$  TeV to 2 TeV. (Ferrag et al)

$M_c = 2$  TeV,  
no PDF error



$M_c = 2$  TeV,  
with PDF error



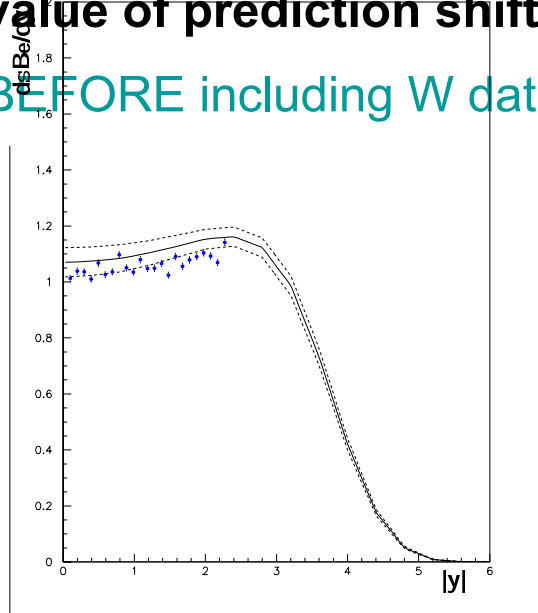
**A. COOPER-SARKAR**

## TURNING THE QUESTION AROUND...

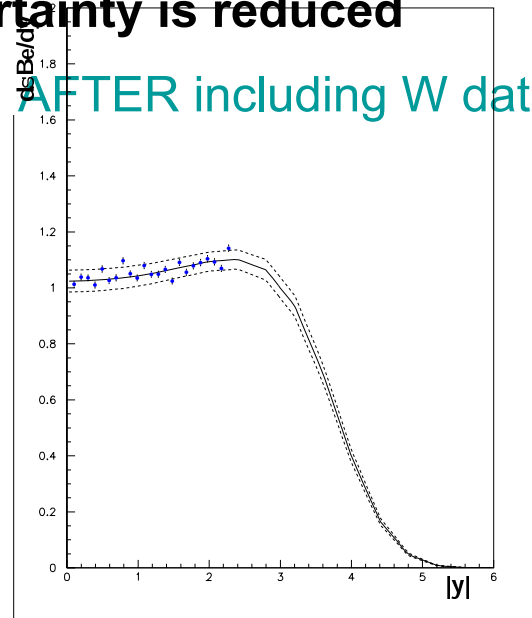
by how much can we reduce the PDF errors with early LHC data?

Generate data with 4%  $\epsilon$  **EXAMPLE I: SMALL  $x$  GLUON** through ATLFast detector simulation and then include this pseudo-data in the global ZEUS PDF fit **Central value of prediction shifts and uncertainty is reduced**

BEFORE including W data



AFTER including W data



Lepton+ rapidity spectrum data generated with CTEQ6.1 PDF compared to predictions from ZEUS PDF

Lepton+ rapidity spectrum data generated with CTEQ6.1 PDF compared to predictions from ZEUS PDF **AFTER these data are included in the fit**

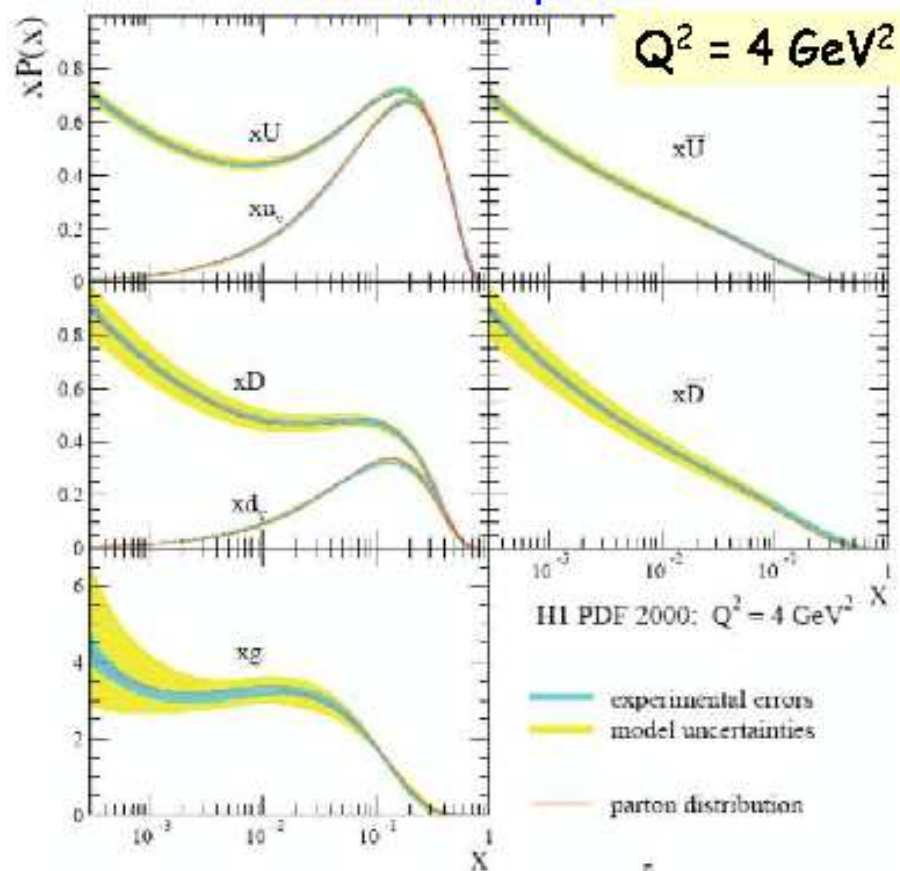
Specifically the low- $x$  gluon shape parameter  $\lambda$ ,  $xg(x) = x^{-\lambda}$ , was  $\lambda = -.199 \pm .046$  for the ZEUS PDF before including this pseudo-data. It becomes  $\lambda = -.181 \pm .030$  after including the pseudo-data.

**A. COOPER-SARKAR**

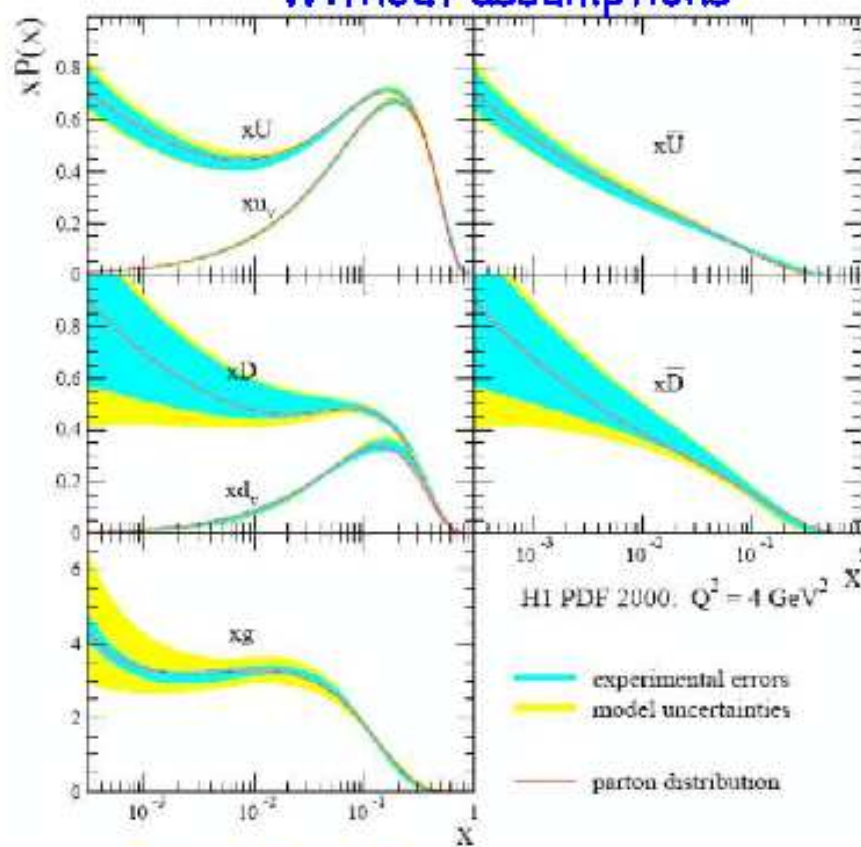
MK/BR,  
March 2005

EXAMPLE II:  $\bar{u} - \bar{d}$  AT SMALL  $x$   
POSSIBLE NON-REGGE BEHAVIOUR

With assumptions



Without assumptions



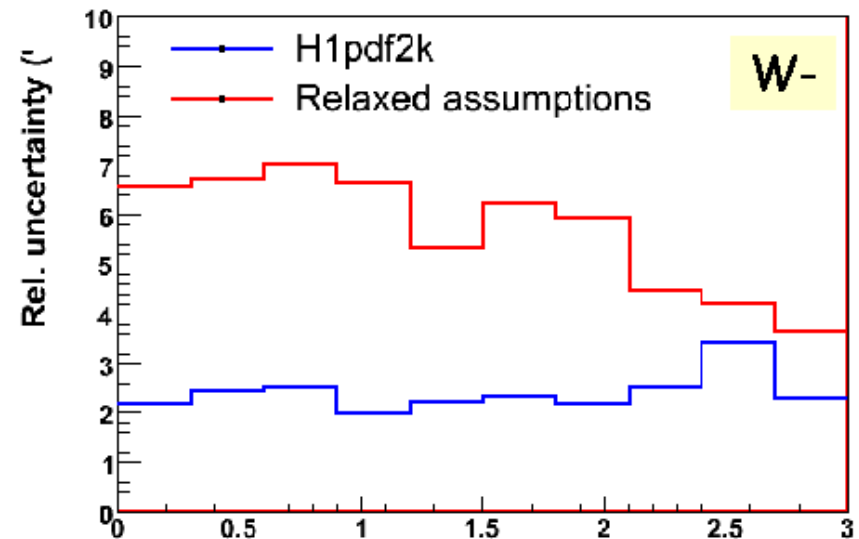
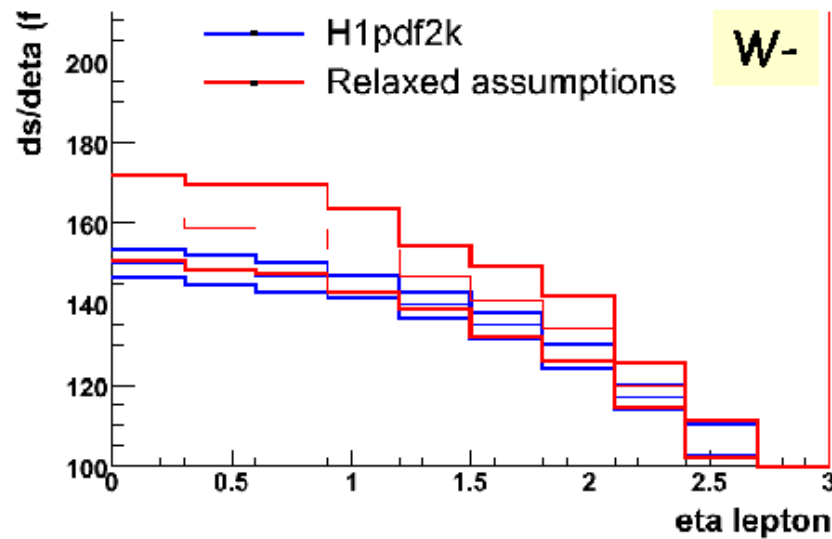
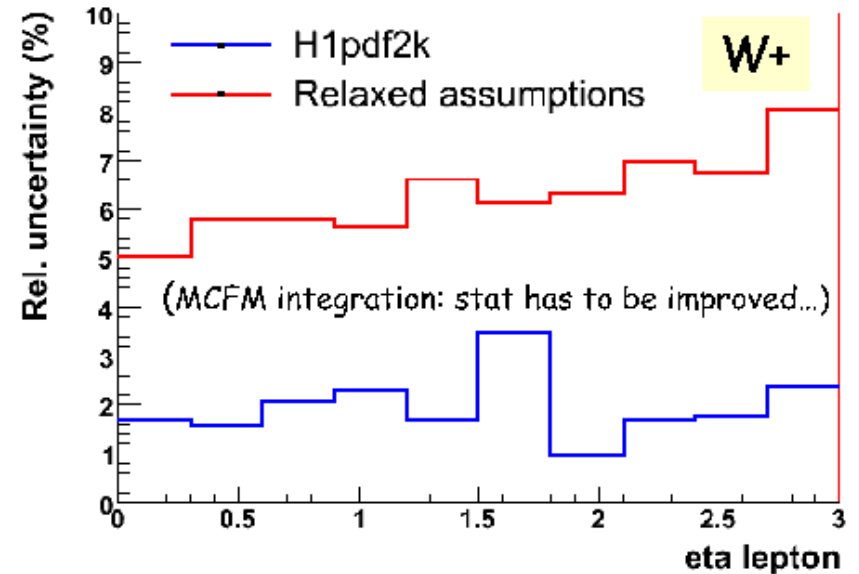
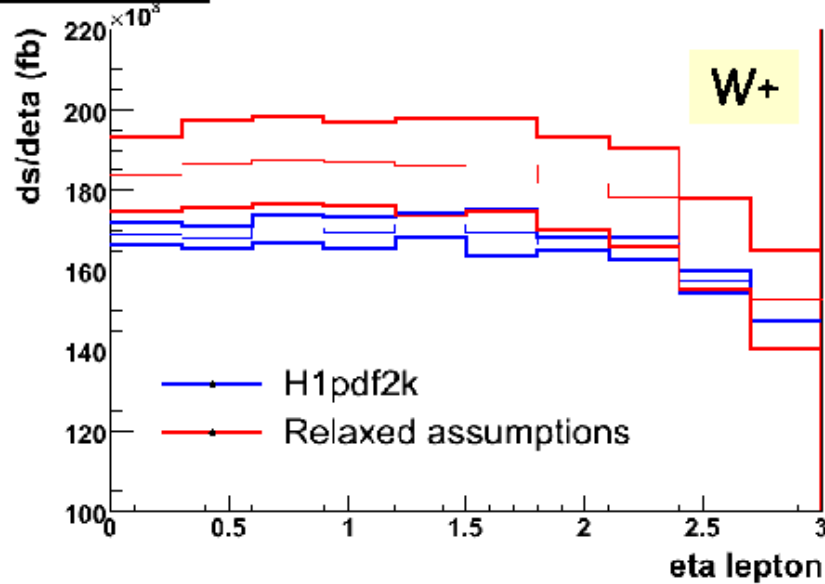
eD at HERA would have helped a lot...

Relax  $B_U = B_D$   
& relation between  $A_U$  and  $A_D$

i.e. still assume  $u = \bar{u}$ ,  $d = \bar{d}$ ,  
relax  $\bar{u} = \bar{d}$

# Uncertainties on W x-sections

W+ cross section



**E. PEREZ**

## EXAMPLE III: GLUONS AND SEA AT SMALL $x$

LHC is a low- $x$  machine (at least for the early years of running)

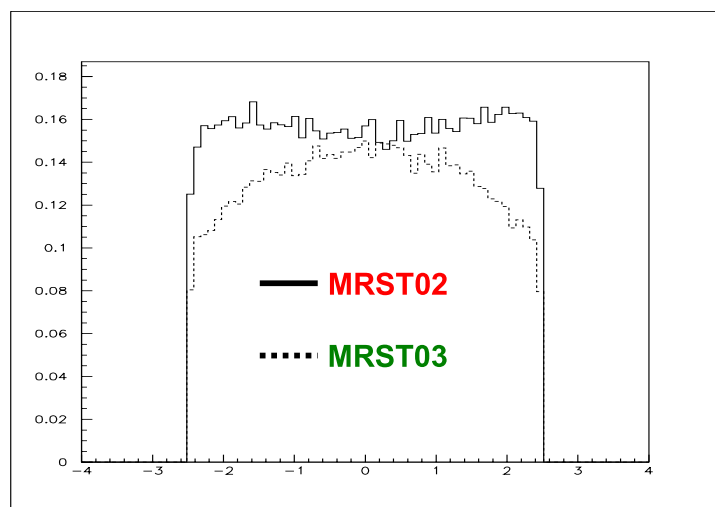
Is NLO (or even NNLO) DGLAP good enough?

The QCD formalism may need extending at small- $x$

MRST03 is a toy PDF set produced without low- $x$  data

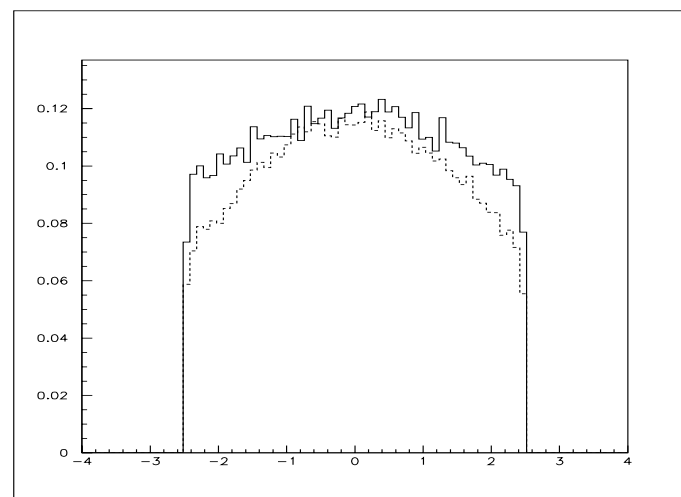
200k events of  $W^{+-} \rightarrow e^{+-}$  generated with MC@NLO using MRST03 and MRST02

Reconstructed Electron Pseudo-Rapidity Distributions (ATLAS fast simulation)



Reconstructed  $e^+$

6 hours  
running



Reconstructed  $e^-$

If something is very different about low- $x$  behaviour it will show up in the our measurable rapidity range

**A. COOPER-SARKAR**



# OUTLOOK

# FITPDF?

NEED A JOINT EFFORT OF THEORISTS AND LHC EXPERIMENTALISTS:

- WHICH PRECISION MEASUREMENTS ARE LIMITED BY PDFS?
- WHEN DOES LACK OF PDF KNOWLEDGE HIDE/SIMULATE NEW PHYSICS?
- HOW CAN LHC MEASUREMENTS IMPROVE PDF DETERMINATION?

# EXPERIMENTAL INPUT

## SOME EXPECTATIONS

- HERA:  $F_L$  & COMBINED  $F_2$
- NEUTRINO DATA: INCOMPATIBILITIES?
- JETS: WHAT IS THE IDEAL DATASET?

## OPEN ISSUES

### SOME EXAMPLES

- W/Z CROSS-SECTION AS STANDARD CANDLE:  
ARE UNCERTAINTIES UNDER CONTROL? (EG HEAVY QUARKS)
- $\bar{u} - \bar{d}$  FROM  $W^\pm$  ASYMMETRY:  
REGGE BEHAVIOUR AT SMALL  $x$ ?
- OBSERVABLES SENSITIVE SMALL  $x$  PARTONS  
RESUMMATION REQUIRED AT NNLO?

# THE PDF CONTROL ROOM AT LHC

Control room is operational and used during the cosmic commissioning runs integrating gradually more detector components.

Cosmic ray data is collected through segments of the full final Event Building and DAQ system

