



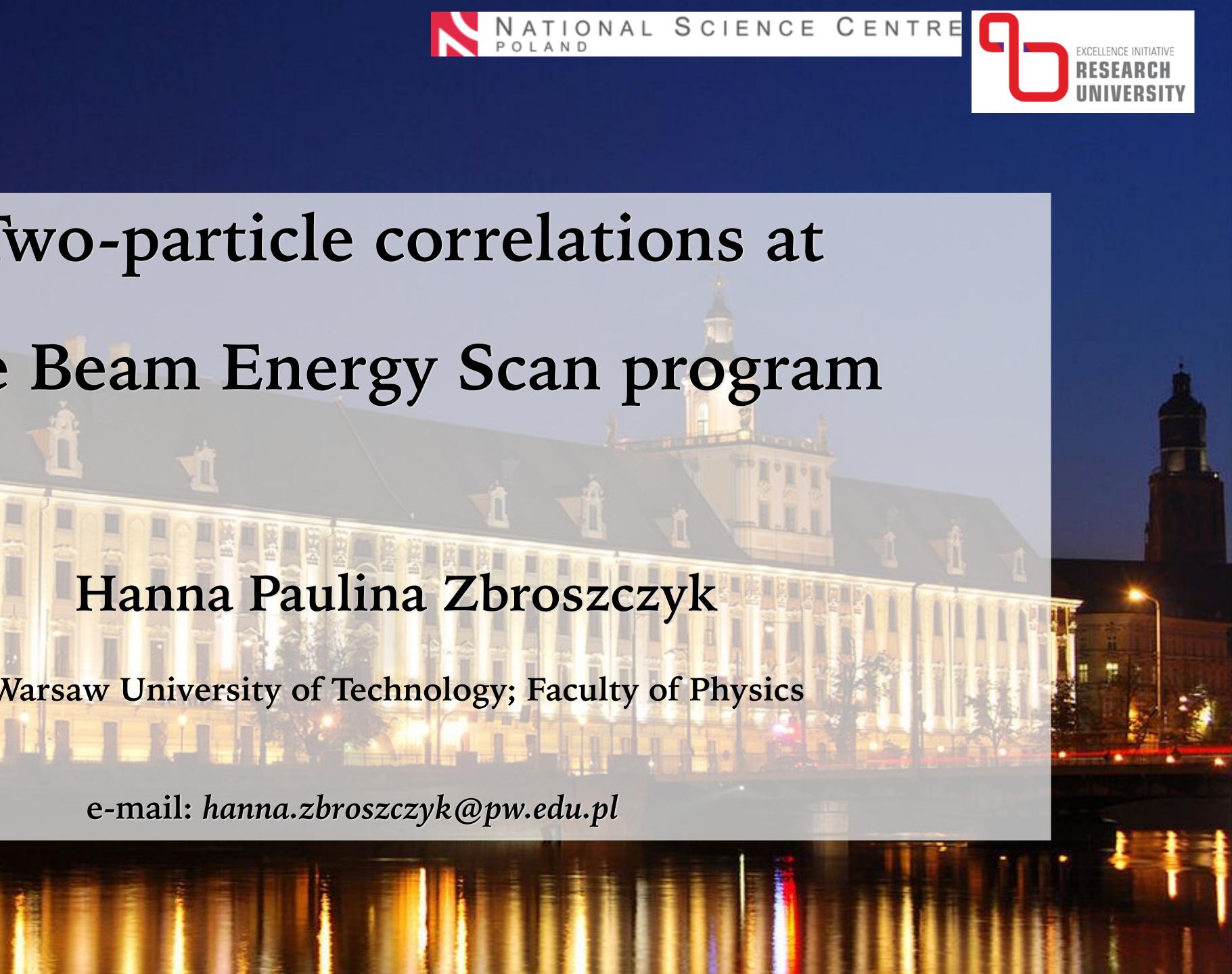
Faculty of Physics

Two-particle correlations at the Beam Energy Scan program

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Outline

Introduction

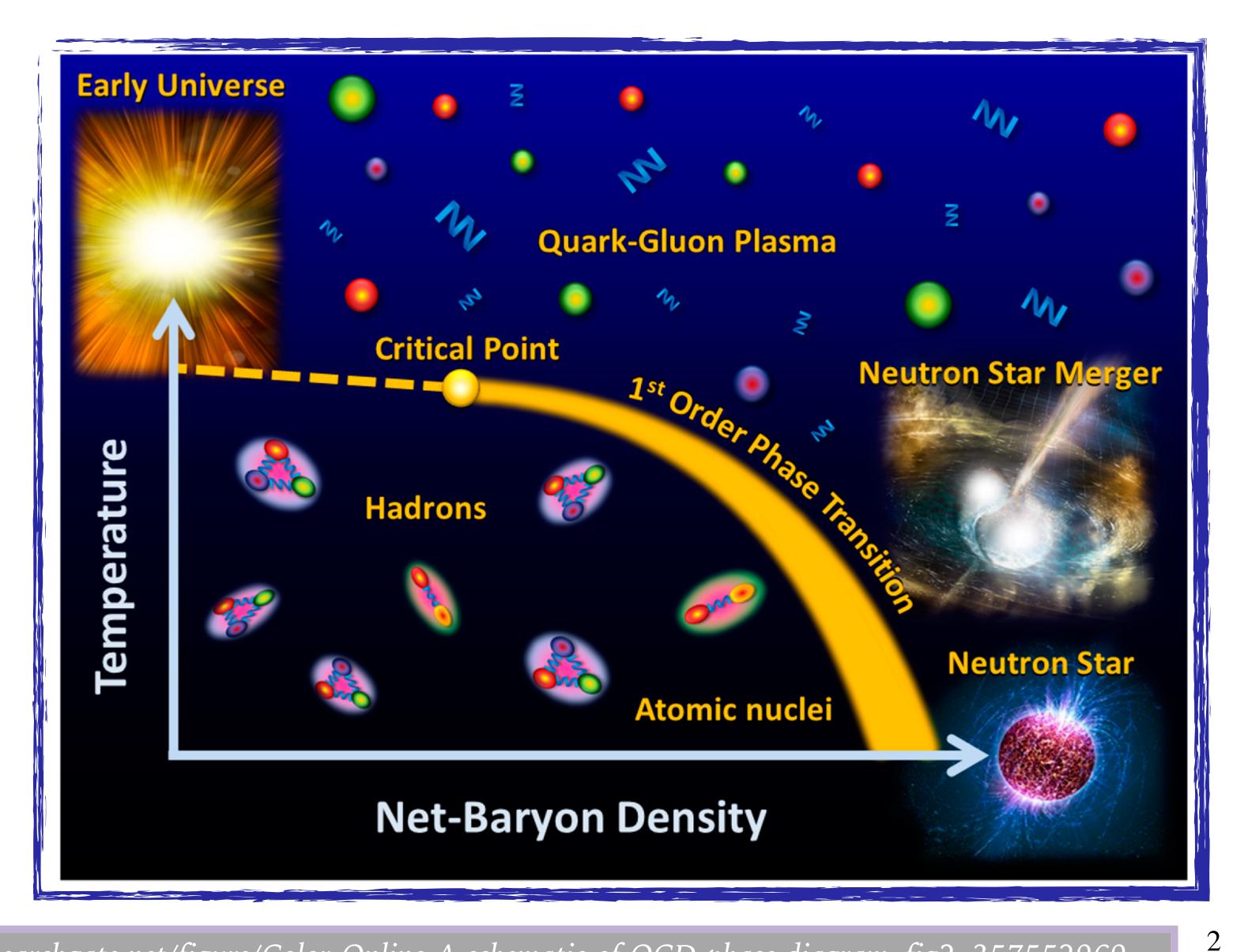
- QCD phase diagram
- Correlation femtoscopy

Results

- Final State Interactions
- Phase transition

Future

Conclusions



https://www.researchgate.net/figure/Color-Online-A-schematic-of-QCD-phase-diagram_fig2_357552969

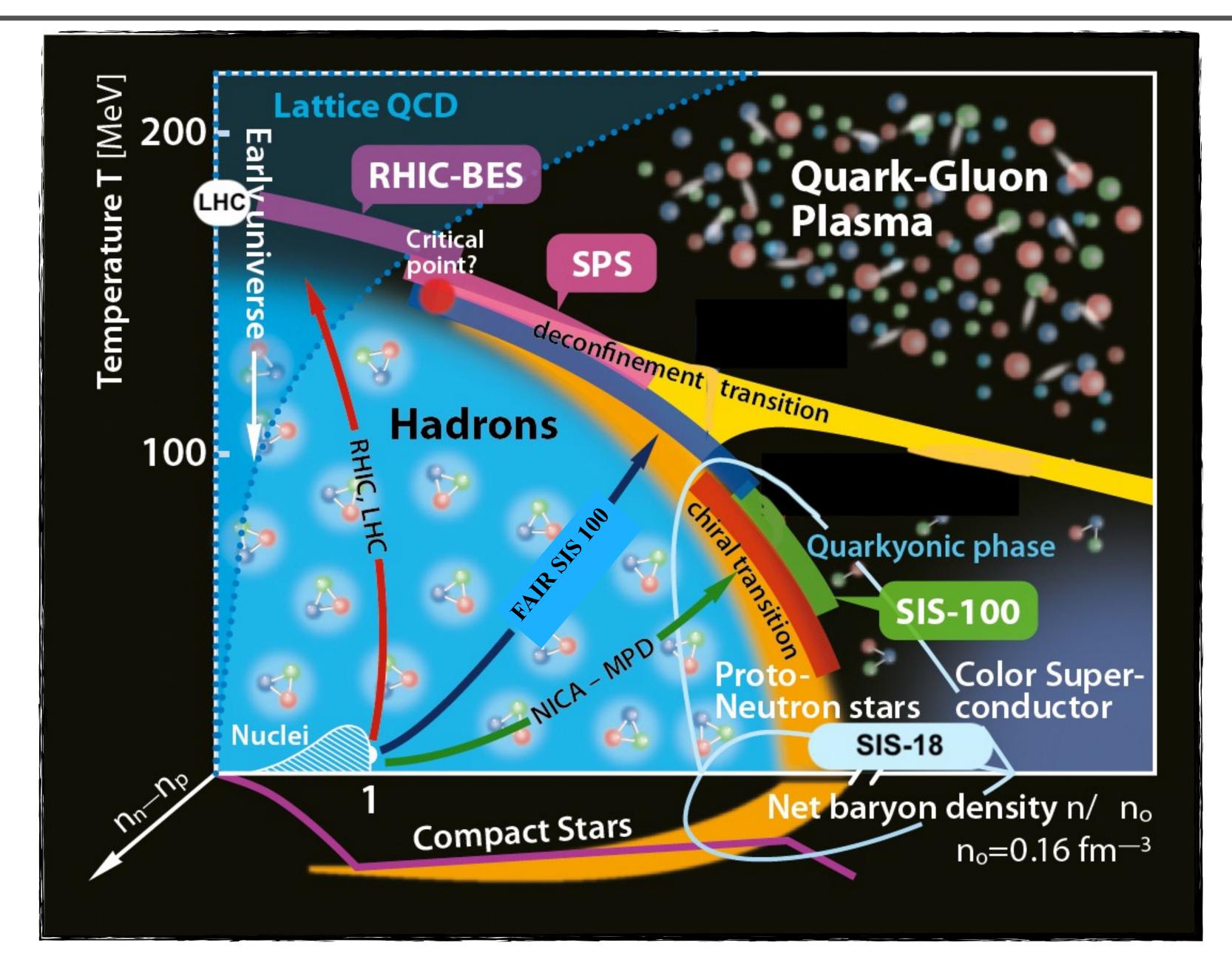
Introduction

QCD phase diagram

Correlation femtoscopy



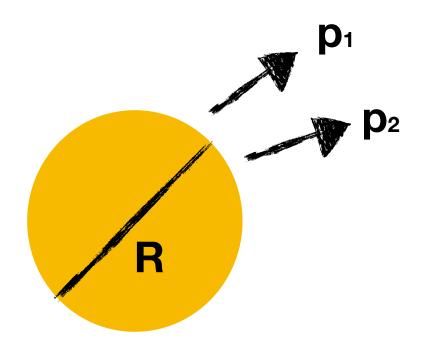
QCD Phase diagram of strongly interacting matter





Traditional and non-traditional femtoscopy

Femtoscopy (originating from HBT): the method to probe geometric and dynamic properties of the source



Space-time properties $(10^{-15}m, 10^{-23}s)$ can be determined due to two-particle correlations that arise due to: Quantum Statistics (Fermi-Dirac, Bose-Einstein); Final State Interactions (Coulomb, strong)

determined assumed

$$C(k^*, r^*) = \int S(r^*) |\Psi(k^*, r^*)|^2 d^3 r^* = \frac{1}{4}$$

emission function $S(r^*)$

 $\Psi(k^*, r^*)$ – two-particle wave function (includes e.g. FSI interactions)

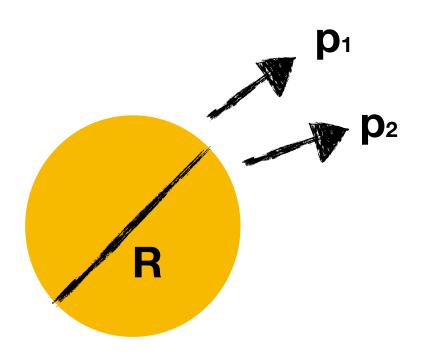
 $Sgnl(k^*)$ - correlation function $Bckg(k^*)$

measured $Sgnl(k^*)$ $Bckg(k^*)$



Traditional and non-traditional femtoscopy

If we assume we know the emission function, measured correlation function can be used to determine parameters of Final State Interactions



two-particle correlations that arise due to: Final State Interactions (Coulomb, strong)

assumed determined

$$C(k^*, r^*) = \int S(r^*) |\Psi(k^*, r^*)|^2 d^3r^* = S(r^*)$$

$$S(r^*) - \text{emission function}$$

 $\Psi(k^*, r^*)$ – two-particle wave function (includes e.g. FSI interactions)

 $\frac{Sgnl(k^*)}{Bckg(k^*)}$ - correlation function

- Space-time properties $(10^{-15}m, 10^{-23}s)$ can be determined due to
- Quantum Statistics (Fermi-Dirac, Bose-Einstein);

measured $Sgnl(k^*)$ $Bckg(k^*)$

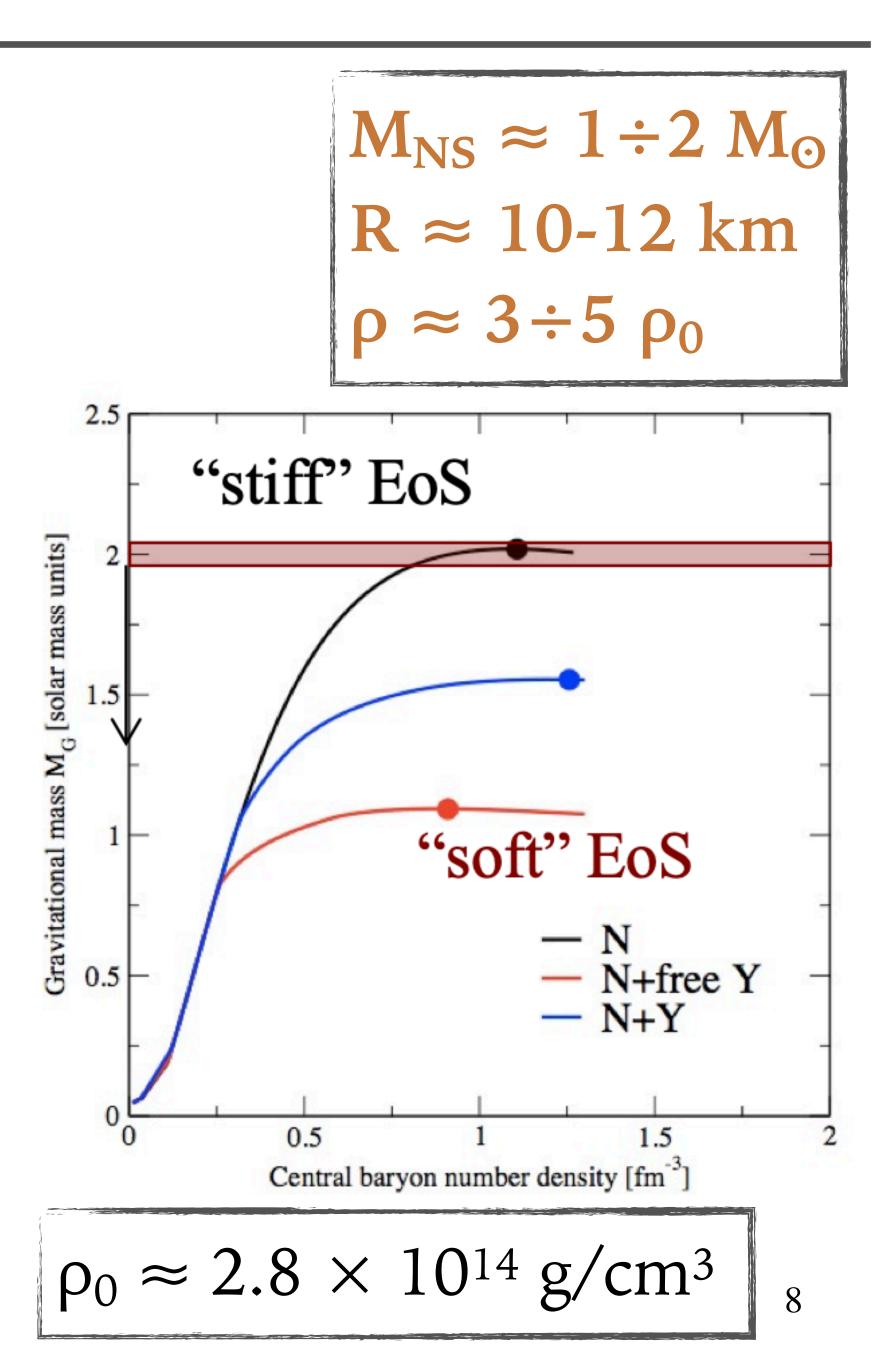


Results Final State Interactions Phase transition



Neutron star puzzle

- Hyperons: expected in the core of neutron stars; conversion of N into Y energetically favorable.
- Appearance of Y: The relieve of Fermi pressure \rightarrow softer EoS \rightarrow mass reduction (incompatible with observation)



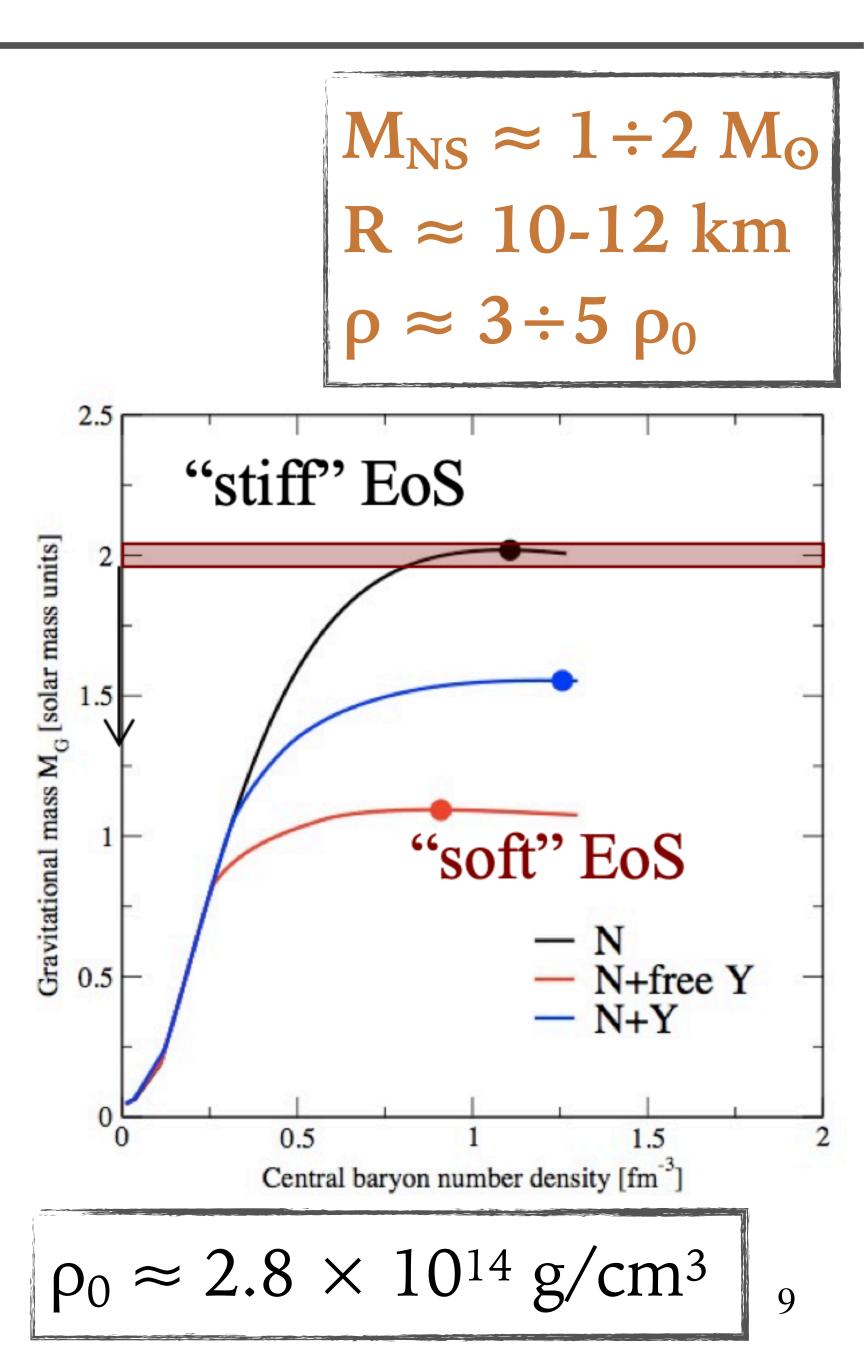
Neutron star puzzle

- Hyperons: expected in the core of neutron stars; conversion of N into Y energetically favorable.
- Appearance of Y: The relieve of Fermi pressure \rightarrow softer EoS \rightarrow mass reduction (incompatible with observation).

The solution: a mechanism providing the additional pressure to make the EoS stiffer.

Possible mechanisms:

- Two-body Y-N & Y-Y interactions
- Chiral forces
- Hyperonic Three Body Forces
- Quark Matter Core Phase transition at densities lower than hyperon threshold



Neutron star puzzle

- Hyperons: expected in the core of neutron stars; conversion of N into Y energetically favorable.
- Appearance of Y: The relieve of Fermi pressure \rightarrow softer EoS \rightarrow mass reduction (incompatible with observation).

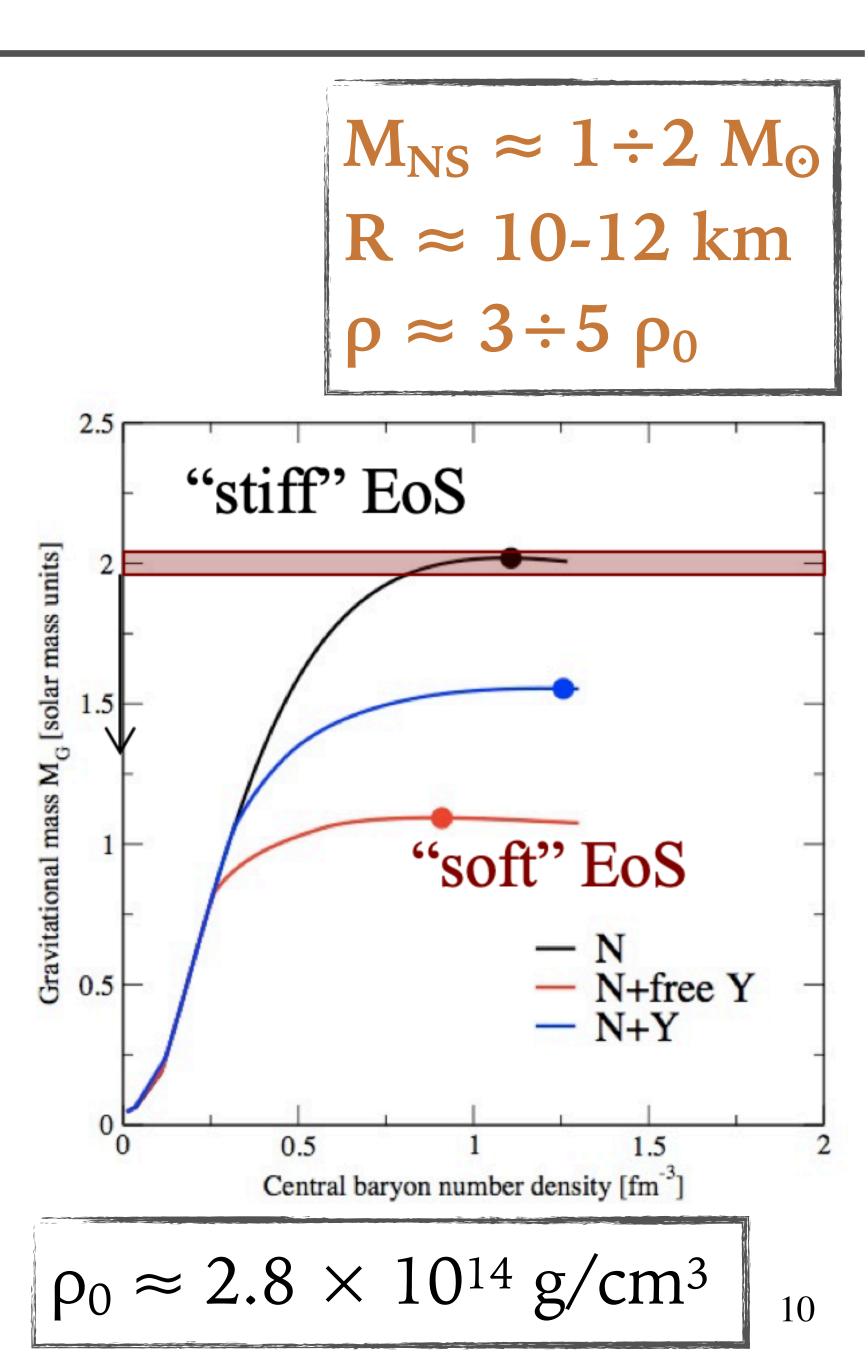
The solution requires a mechanism that could provide the additional pressure at high densities needed to make the EoS stiffer.

Possible mechanisms:

- Two-body Y-N & Y-Y interactions
- Chiral forces
- Hyperonic Three Body Forces
- Quark Matter Core Phase transition at densities lower than hyperon threshold

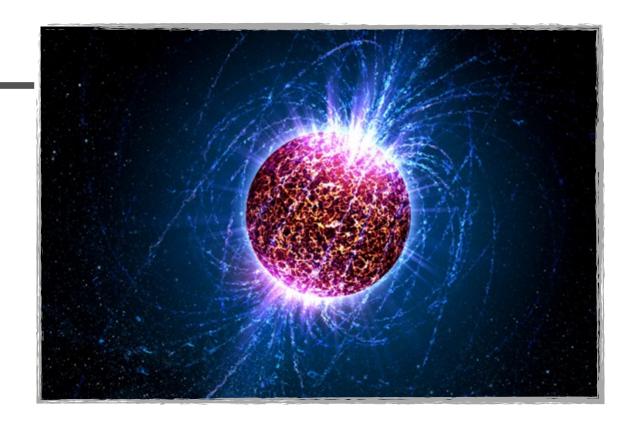
A lot of experimental and theoretical effort to understand:

- The K-N interaction, governed by the presence of $\Lambda(1405)$
- The nature of $\Lambda(1405)$, the consequences of KNN formation
- K and \overline{K} investigated to understand kaon condensation



- Experiment: More interest about Y–N and Y–Y interactions.
- **Theory**: Major steps forward have been made (Lattice QCD).
- Numerous theoretical predictions exist, but no clear evidence for any such bound states, despite many experimental searches.
- The existence of hypernuclei (confirmed by attractive Y–N interaction) \rightarrow indicates the possibility to bind Y to N.
- The measurement of the Y–N and Y–Y interactions leads to important implications for the possible formation of **Y–N** or **Y–Y bound states**.
- A precise knowledge of these interactions help to explore unknown structure of neutron stars.

Y-N and Y-Y interactions



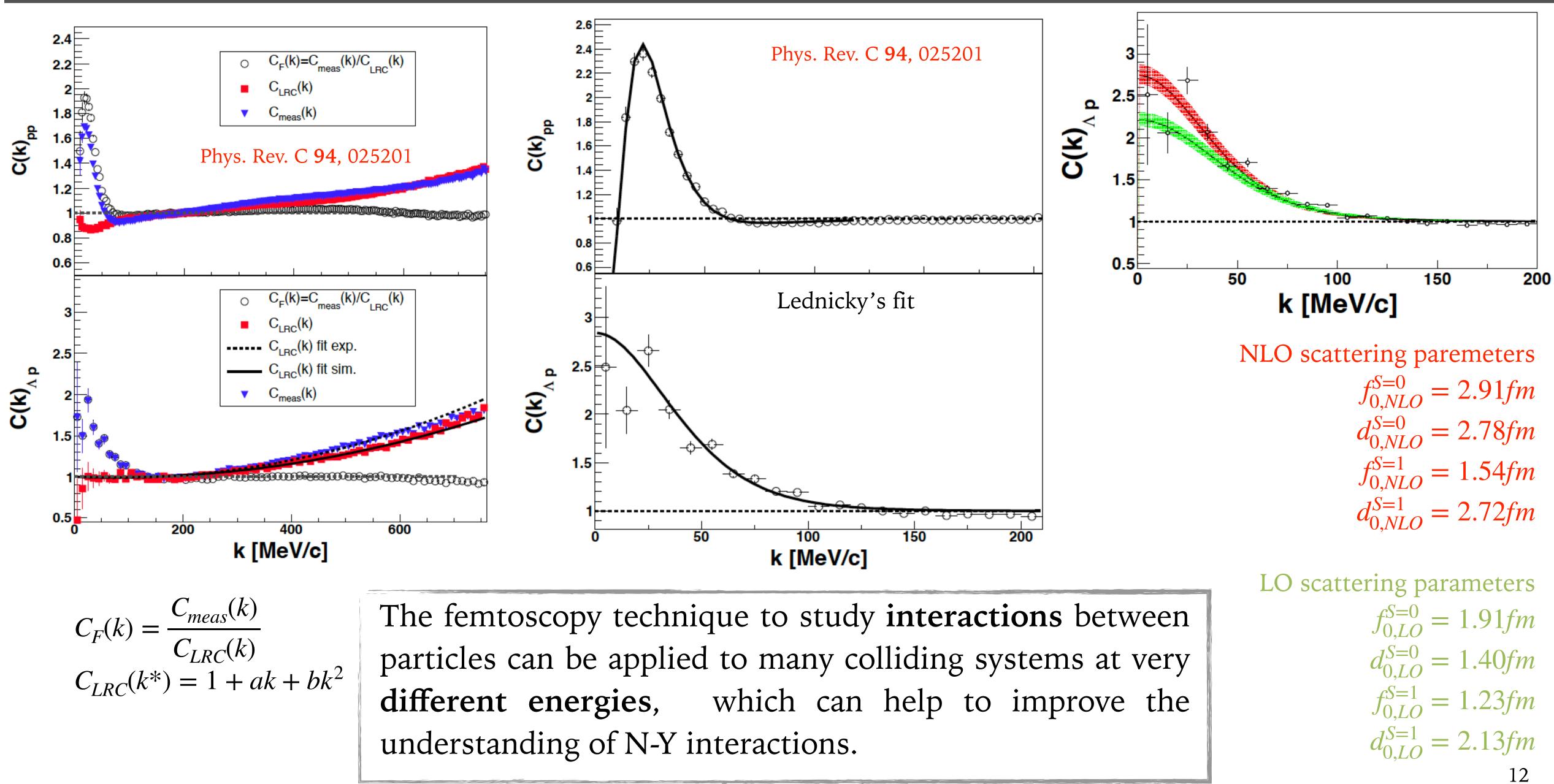






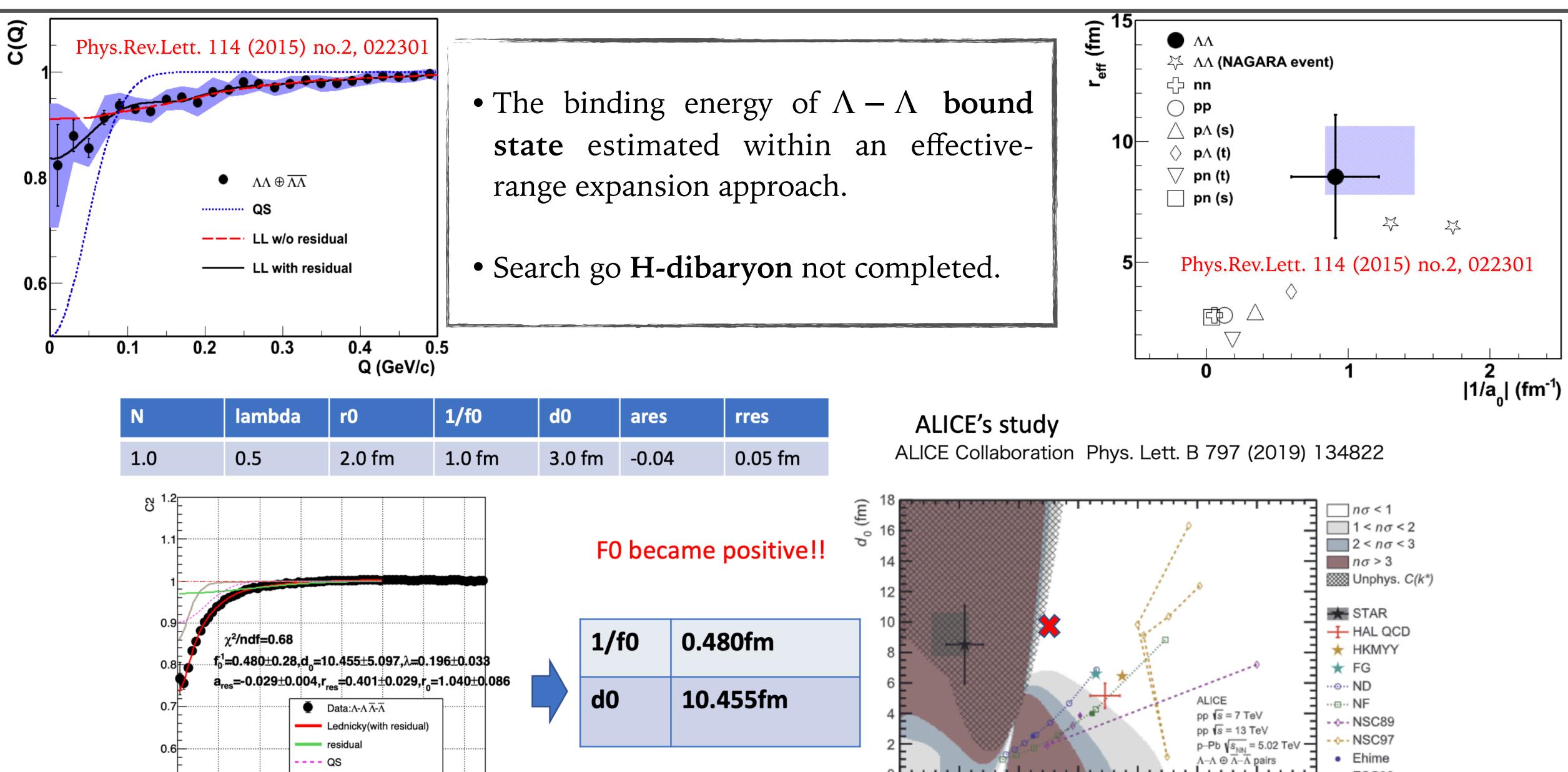
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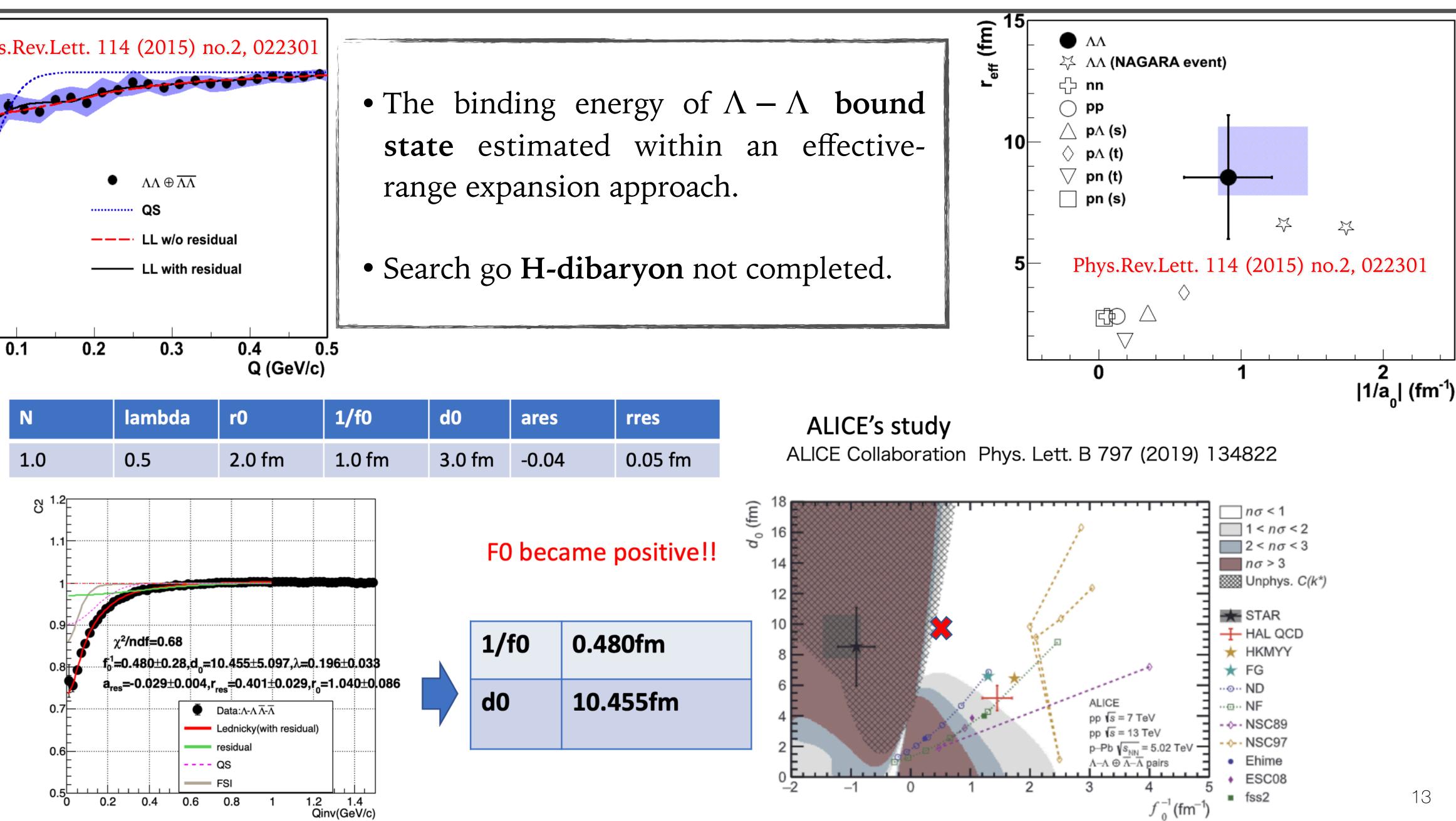
Y-N interactions at HADES



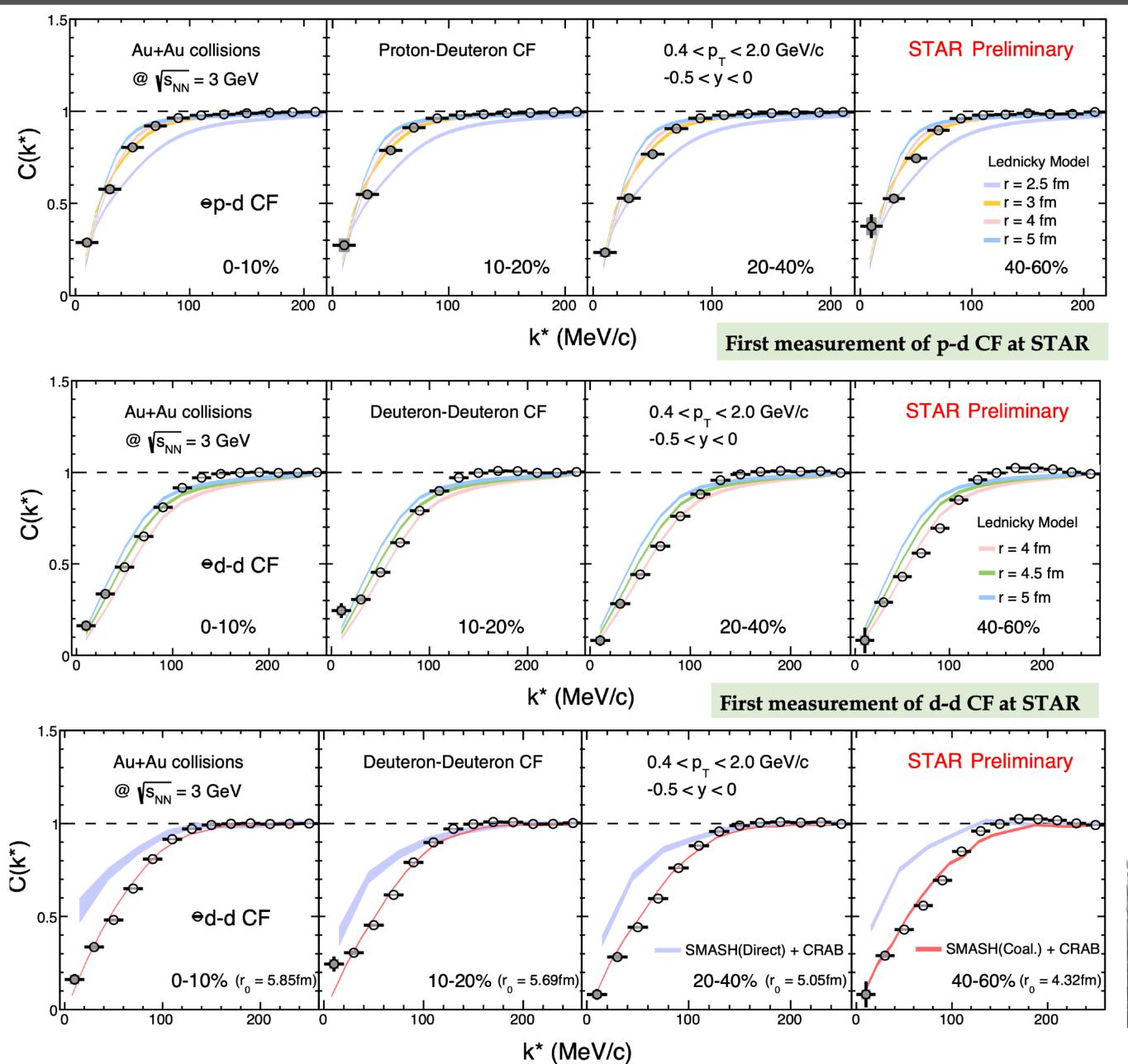
$$C_F(k) = \frac{C_{meas}(k)}{C_{LRC}(k)}$$
$$C_{LRC}(k^*) = 1 + ak + bk^2$$

Strange Baryon Correlations (Including Λ Hyperons)





Light nuclei formation at STAR

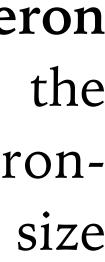


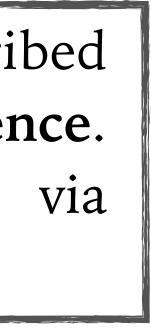
• First measurement of proton-deuteron and deuteron-deuteron correlation functions from STAR

• Proton-deuteron and deuteron-deuteron correlations qualitatively described by the Lednicky-Lyuboshits model; deuterondeuteron has larger emission source size than proton-deuteron

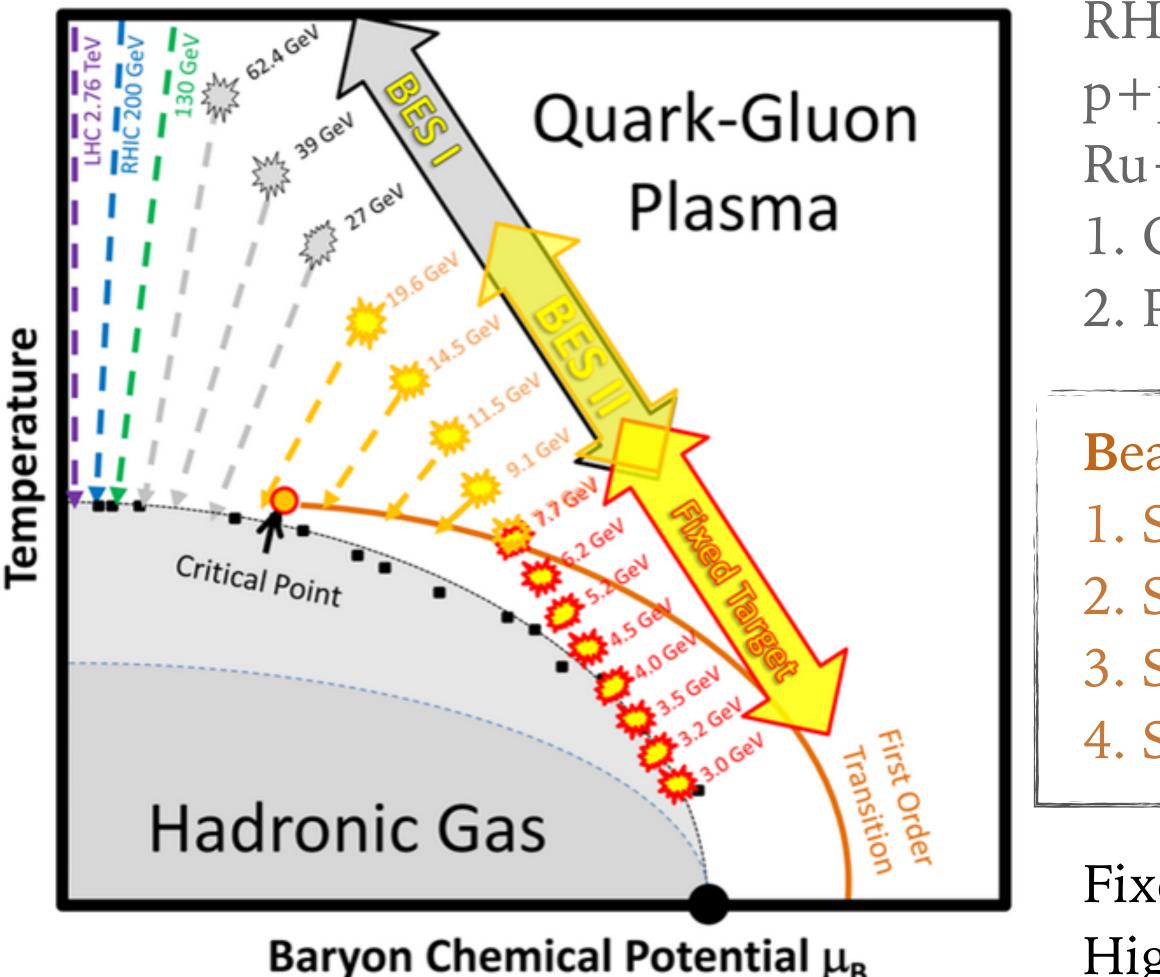
•Deuteron-deuteron correlations described better by the model including coalescence. Light nuclei are likely to be formed via coalescence.







Program Beam Energy Scan



RHIC Top Energy: 200 GeV p+p, p+Al, p+Au, d+Au, 3He+Au, Cu+Cu, Cu+Au, Ru+Ru, Zr+Zr, Au+Au, U+U1. QCD at high energy density/temperature 2. Properties of QGP, EoS

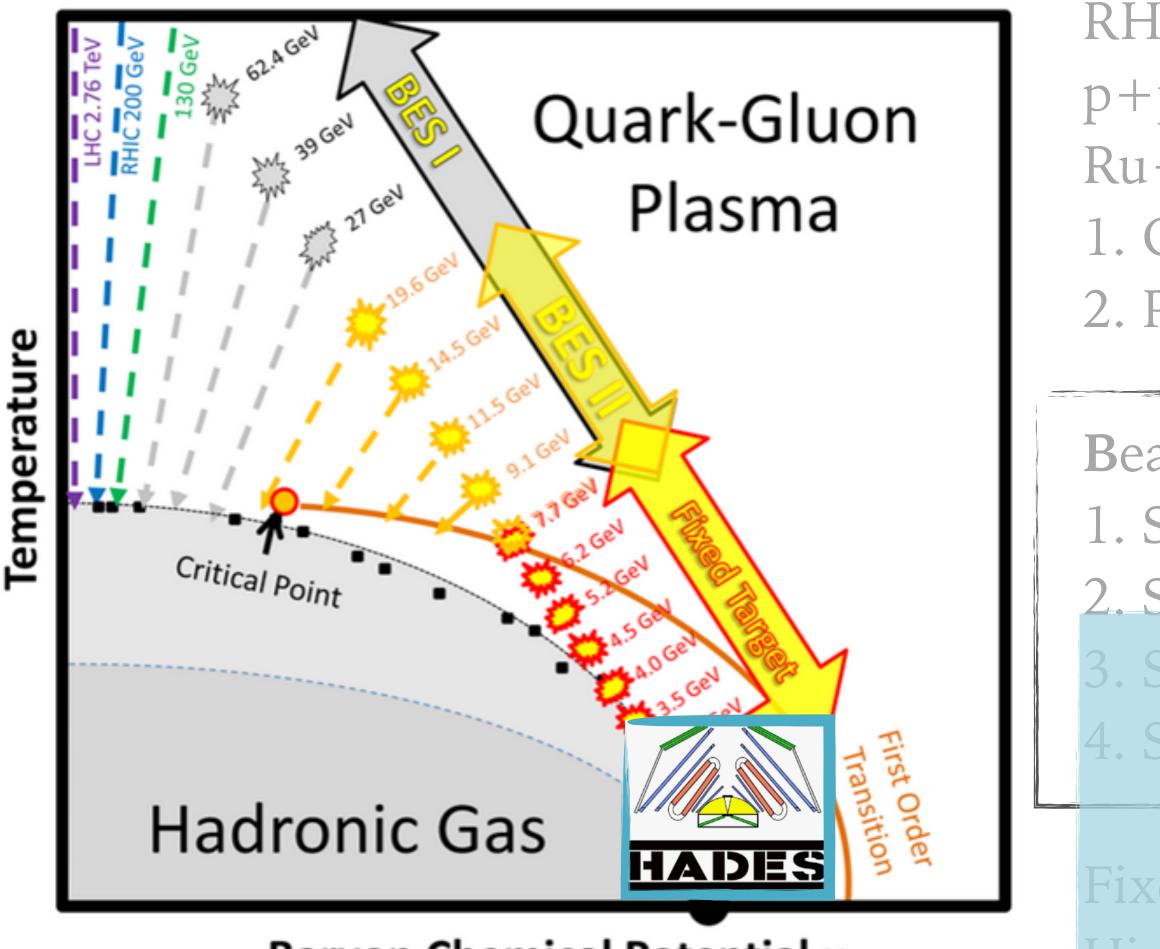
Beam Energy Scan: Au+Au 3-62 GeV 1. Search for turn-off of QGP signatures 2. Search for signals of the first-order phase transition 3. Search for QCD critical point 4. Search for signals of Chiral symmetry restoration

Fixed-Target Program: Au + Au = 3-7.7 GeV High baryon density regime with 420-720 MeV





Program Beam Energy Scan



Baryon Chemical Potential μ_{R}

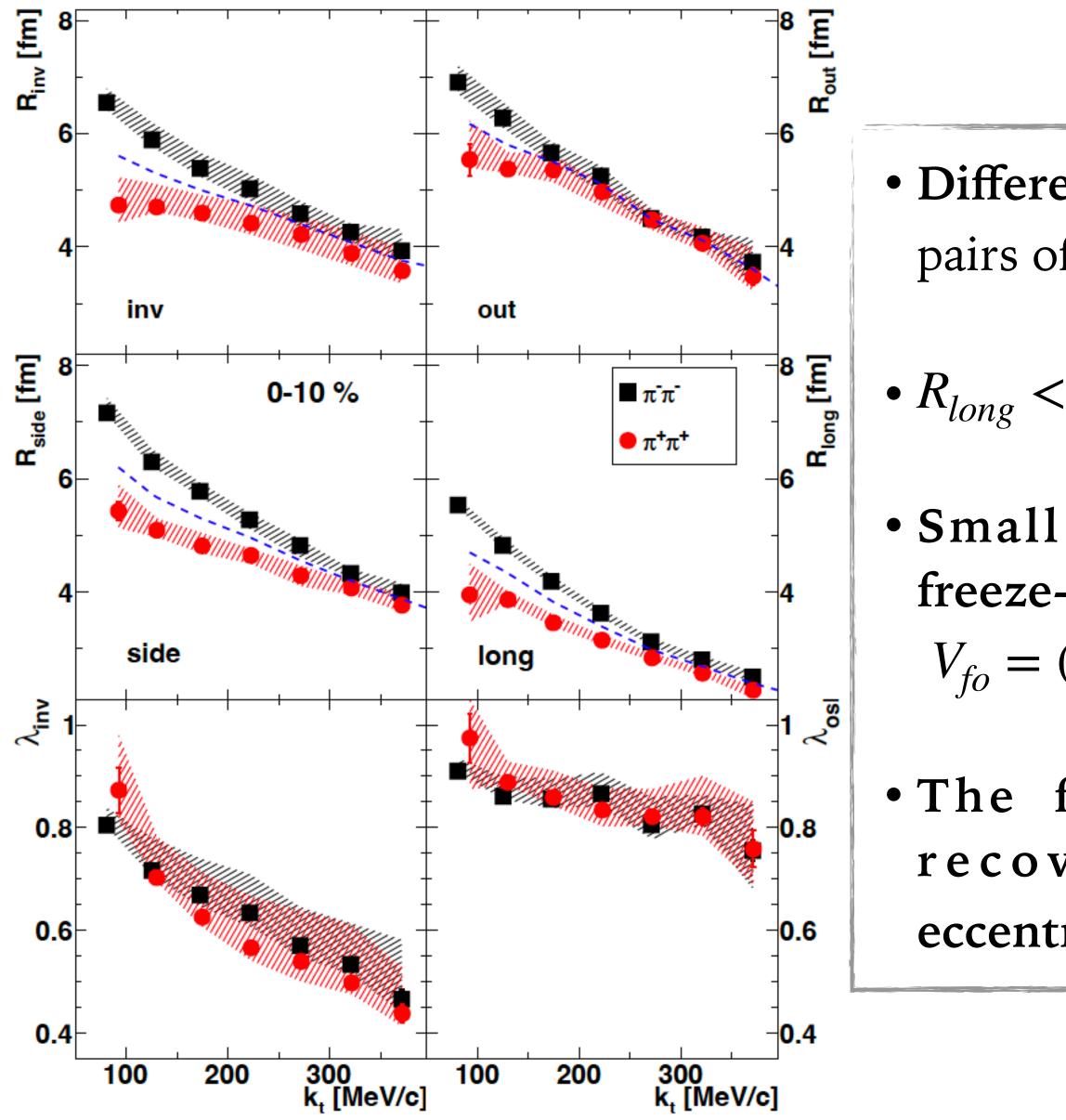
RHIC Top Energy: 200 GeV p+p, p+Al, p+Au, d+Au, 3He+Au, Cu+Cu, Cu+Au, Ru+Ru, Zr+Zr, Au+Au, U+U1. QCD at high energy density/temperature 2. Properties of QGP, EoS

Beam Energy Scan: Au+Au 3-62 GeV 1. Search for turn-off of QGP signatures 2. Search for signals of the first-order phase transition 1-2 GeV/nucleon: hiral symmetry restoration C+C, Ar+KCl, p+p, d+p and p+Nb3-5 GeV/nucleon: p+p 420-720 MeV





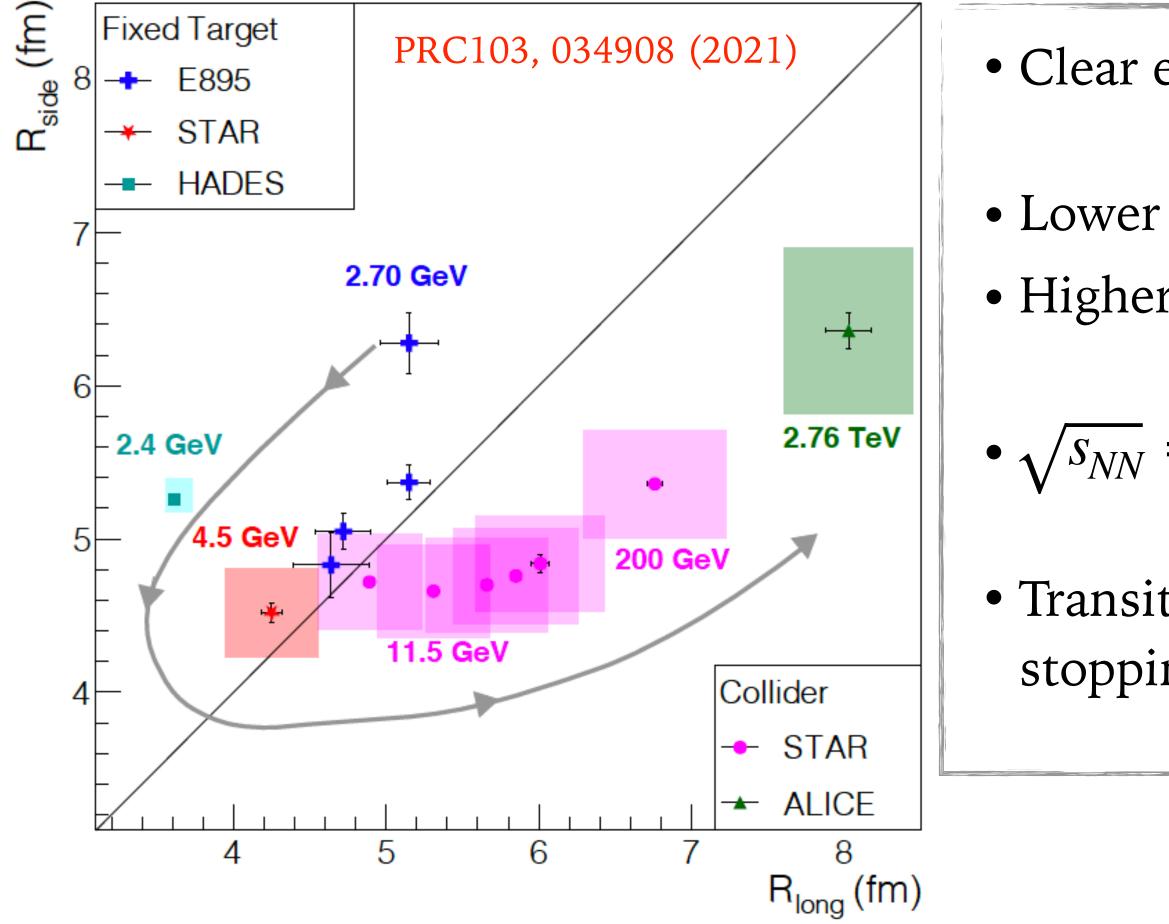
Identical pion correlations at HADES



[fm³] HADES $\pi^{-}\pi^{-}$ 5² 6000 HADES $\pi^+\pi^+$ E895 π⁻π E866 π⁻π E866 π⁺π⁺ • Differences of the radii for ΝΑ49 π'π' CERES $\pi \pi + \pi^+ \pi^-$ 4000 STAR $\pi^{+}\pi^{+}\pi^{+}\pi^{+}$ pairs of π^- and π^+ ALICE π⁻π⁻+π⁺π⁺ 2000 • $R_{long} < R_{out} \simeq R_{side}$ $m_{t} = 260 \text{ MeV}$ • Small variation of the 10² 10³ 10 ∕ √s_{NN} [GeV] freeze-out volume ^Efinal $V_{fo} = (2\pi)^{3/2} R_{side}^2 R_{long}$ + p₁₁₂ = 160 MeV/c $\pi \pi$ 0.4 + p_{t.12} = 250 MeV/c + p_{t,12} = 440 MeV/c • The final eccentricity 0.3 recovers the initial 0.2 eccentricity for large p_T . 0.1 0.1 0.2 0.3 0.4 0

 $\epsilon_{initial}$





• Clear evolution in the freeze-out shape indicated

• Lower energies: system more oblate ($R_{side} > R_{long}$) • Higher energies: system more prolate ($R_{side} < R_{long}$)

• $\sqrt{s_{NN}} = 4.5$ GeV: round system ($R_{side} \simeq R_{long}$)

• Transition region between dynamics dominated by stopping and boost-invariant dynamics.

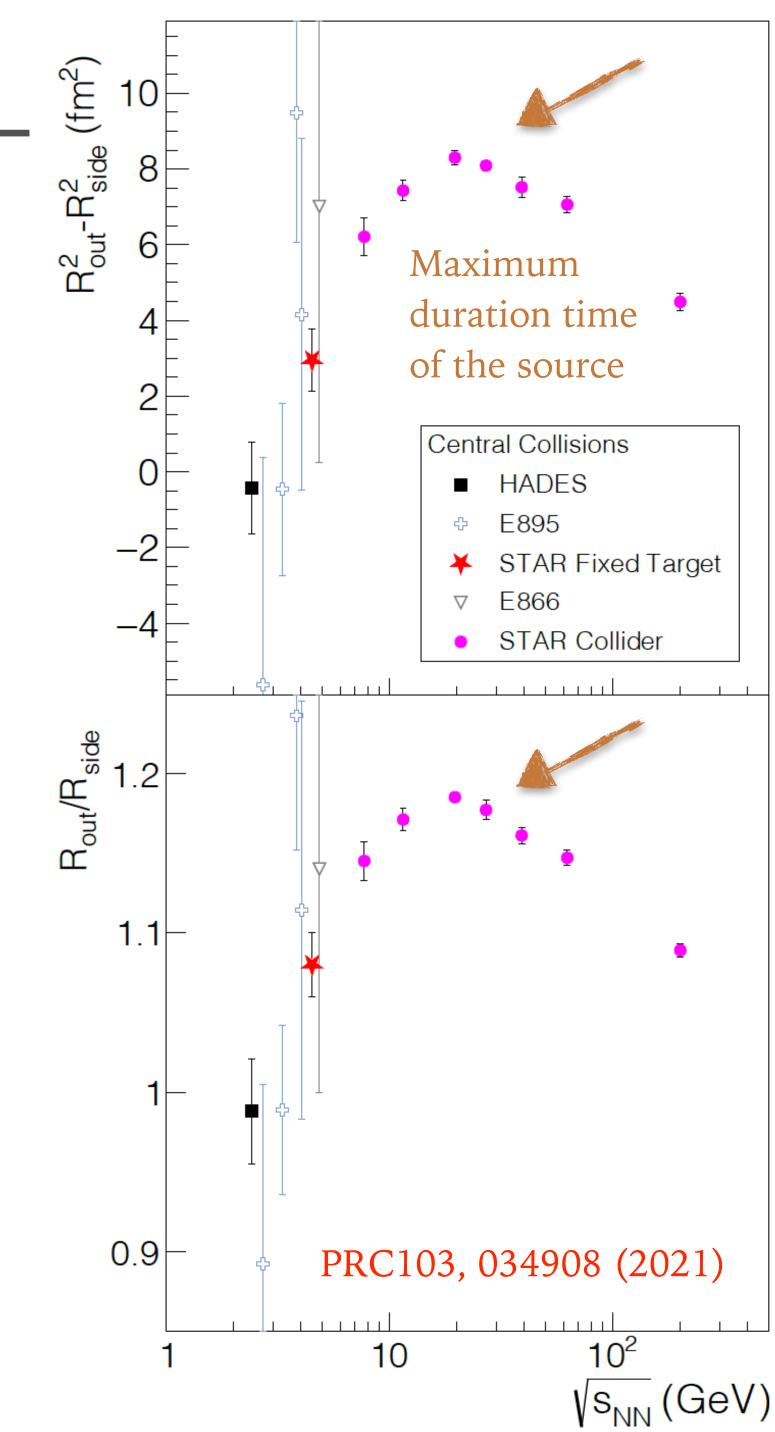
$$R_{out}^2 - R_{side}^2 = \beta_t^2 \Delta \tau^2$$

Visible peak in
$$\frac{R_{out}}{R_{side}}(\sqrt{s_{NN}})$$
 near the $\sqrt{s_{NN}} \simeq 20$

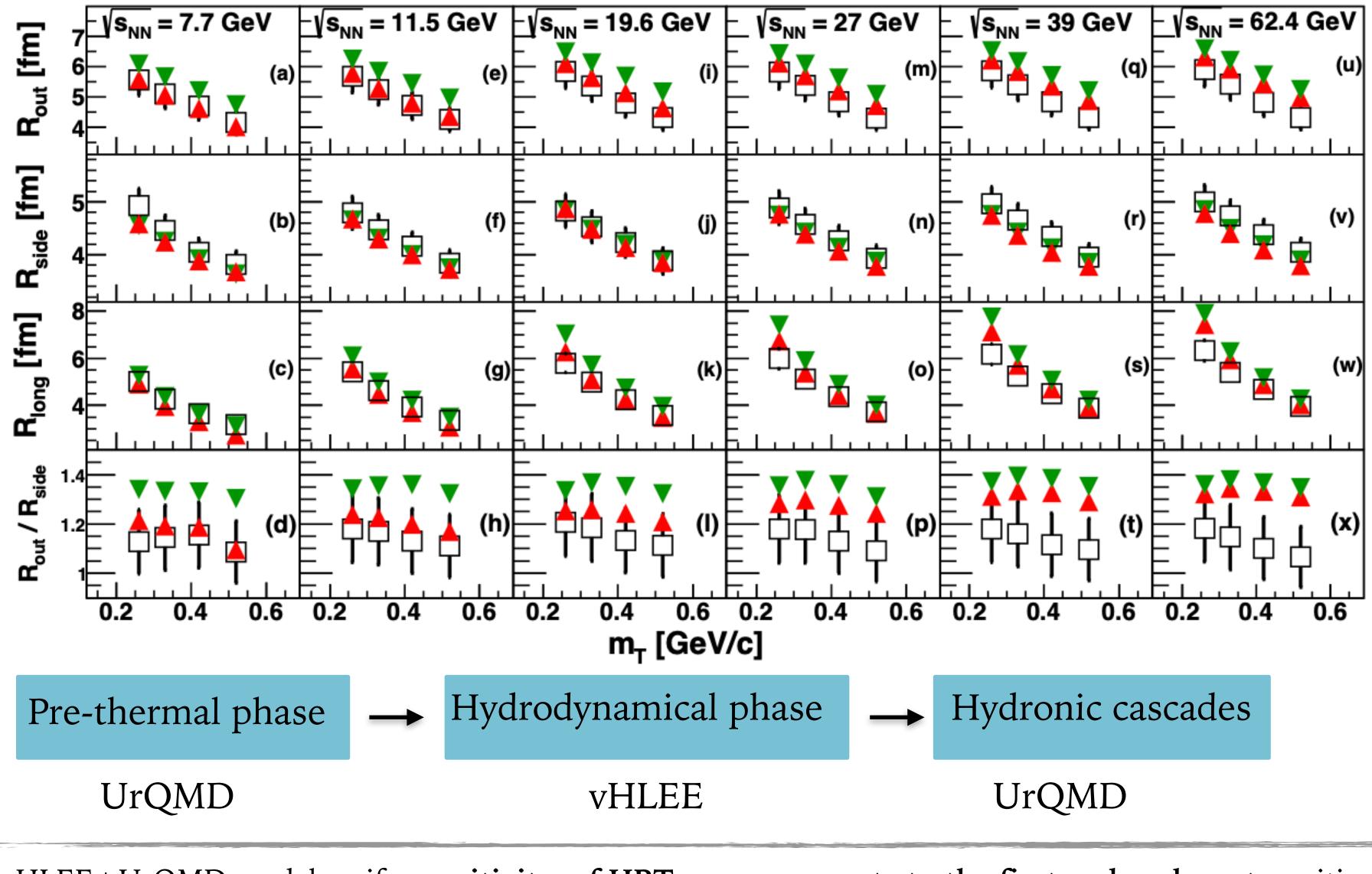
QCD calculations predict a peak near to the QGP transition threshold - signature of first-order phase transition?

Theoretical attention from hydro and transport models needed

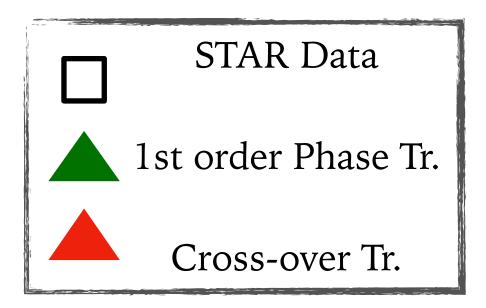
) GeV



How to measure a phase transition?



vHLEE+UrQMD model verify sensitivity of HBT measurements to the first-order phase transition

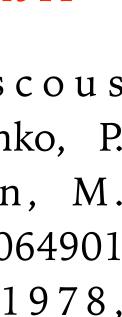


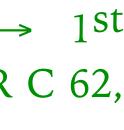
Phys.Rev. C96 (2017) no.2, 024911

vHLLE (3+1)-D viscous hydrodynamics: Iu. Karpenko, P. Huovinen, H. Petersen, M. Bleicher; Phys.Rev. C 91, 064901 (2015), arXiv:1502.01978, 1509.3751

HadronGas + Bag Model $\rightarrow 1^{st}$ order PT; P.F. Kolb, et al, PR C 62, 054909 (2000)

Chiral EoS \rightarrow crossover PT (XPT); J. Steinheimer, et al, J. Phys. G 38, 035001 (2011)









Future



CBM (and HADES) challenges

arXiv:1607.01487v3 [nucl-ex] 29 Mar 2017

The scientific programme of the Compressed Baryonic Matter experiment at FAIR

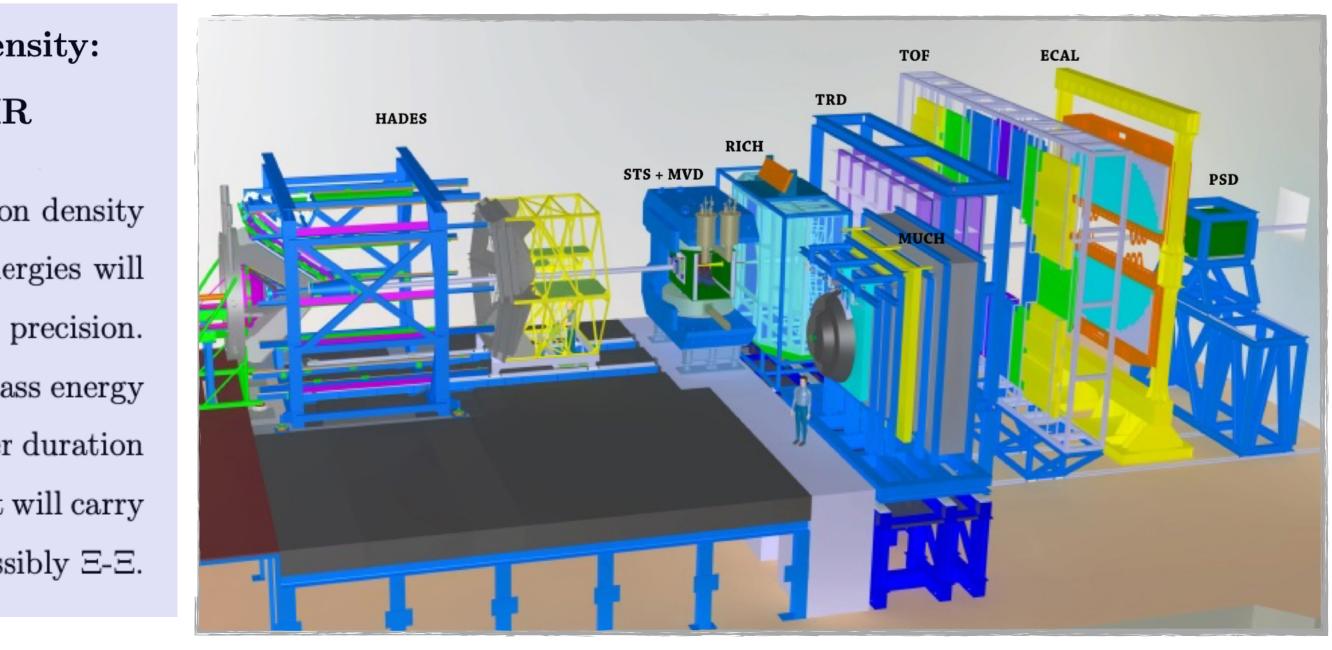
diagnostic probes which are sensitive to the dense phase of the nuclear fireball. The goal of the CBM experiment at SIS100 ($\sqrt{s_{NN}} = 2.7 - 4.9$ GeV) is to discover fundamental properties of QCD matter: the phase structure at large baryon-chemical potentials ($\mu_B > 500$ MeV), effects of chiral symmetry, and the equation-of-state at high density as it is expected to occur in the core of neutron stars. In this article, we review the motivation for and the physics programme of CBM, including activities before the start of data taking in 2024, in the context of the worldwide efforts to explore high-density QCD matter.

QCD Phase Structure and Interactions at High Baryon Density: Completion of BES Physics Program with CBM at FAIR

the nature (attractive or repulsive) of some hyperon interactions. The high baryon density and large number of hyperons produced in nucleus-nucleus collisions at FAIR energies will put CBM in a unique position to measure those interactions with unprecedented precision. Unlike at RHIC and LHC, the collision dynamics at the relatively low center-of-mass energy of FAIR are less explosive, allowing hyperons and nucleons to interact over a longer duration during the collision evolution, yielding stronger correlations. The CBM experiment will carry out a program to measure the correlations of Λ -N, Ξ -N, Ω -N, Λ - Λ , Λ - Ξ , and possibly Ξ - Ξ .

arXiiv:2209.05009 [nucl-ex]12 Sept 2022

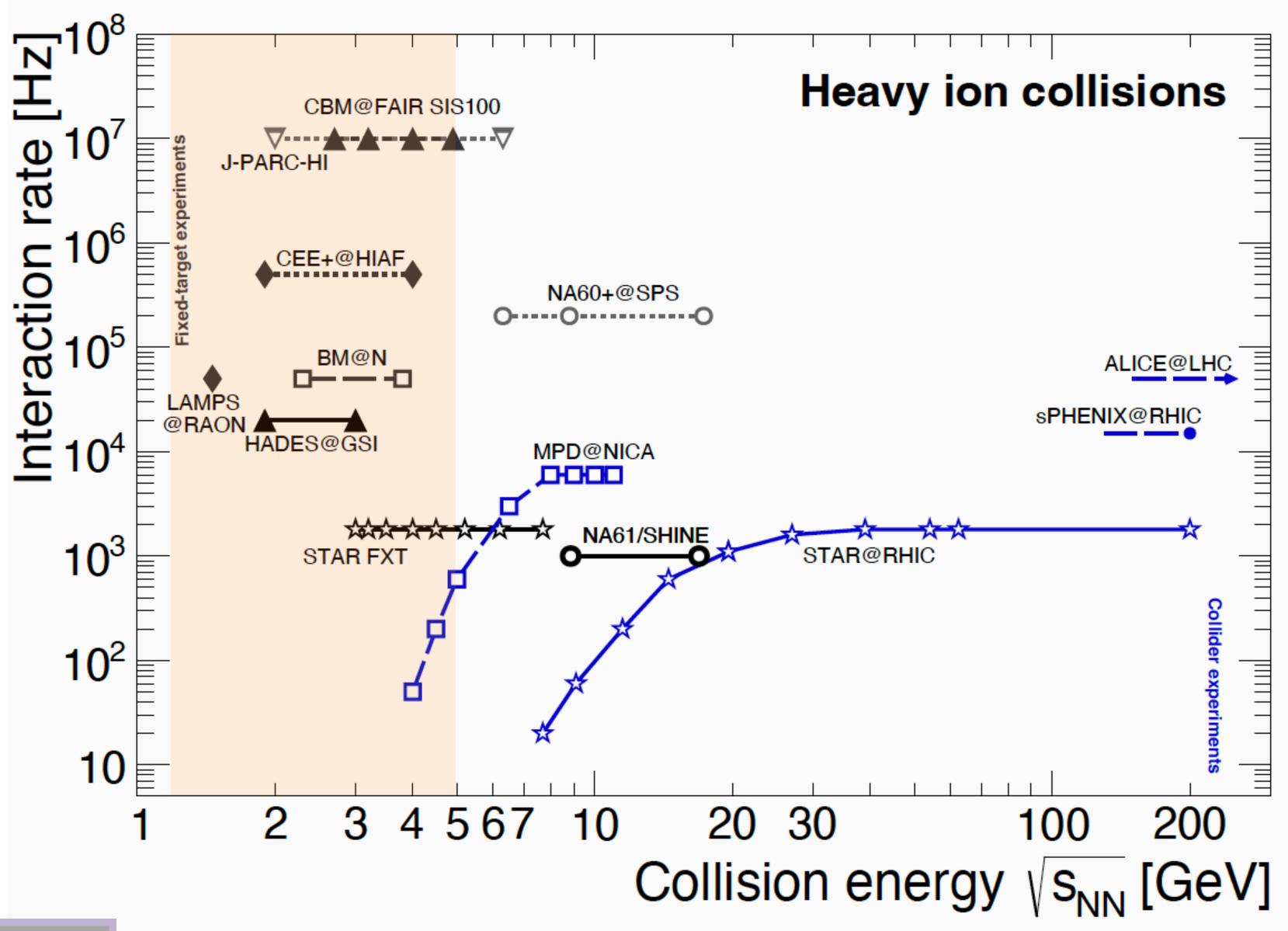
Challenges in QCD matter physics –



The CBM experimental setup together with the HADES detector 22







Bright future



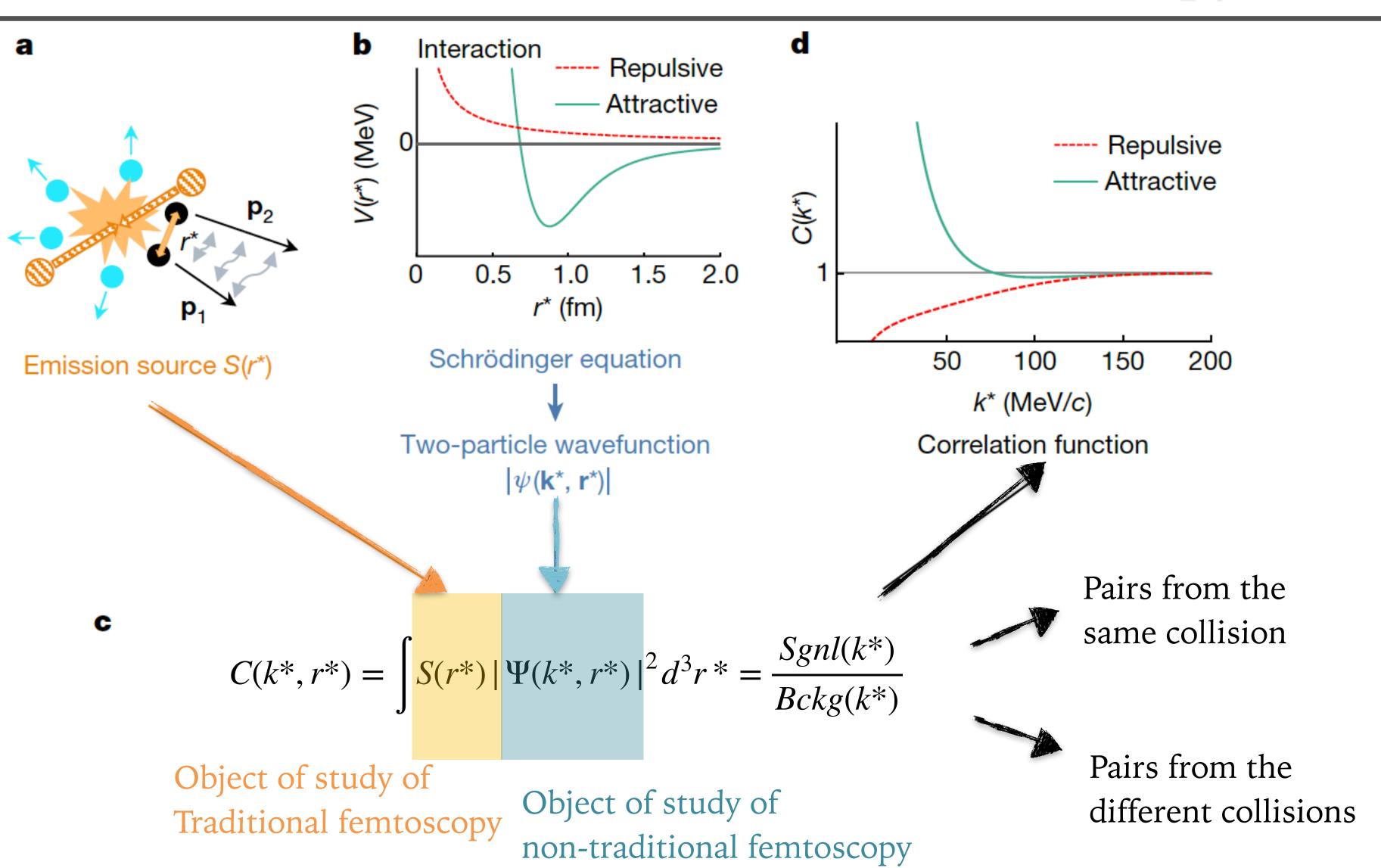
Summary

Energy Scan Programs give un access to:

- Low energy gives a real access to the area of Neutron Stars
- Intermediate energies can help to search for the signatures of the phase transition between hadron and quark matter
- All energies help a lot to study strong interactions at the final state

Back-up slides

Traditional and non-traditional femtoscopy



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nature

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Published: 04 November 2015

Measurement of interaction between antiprotons

The STAR Collaboration

Nature527, 345–348 (2015)Cite this article9961Accesses47Citations368AltmetricMetrics

1 This article has been updated

Abstract

One of the primary goals of nuclear physics is to understand the force between nucleons, which is a necessary step for understanding the structure of nuclei and how nuclei interact with each other. Rutherford discovered the atomic nucleus in 1911, and the large body of knowledge about the nuclear force that has since been acquired was derived from studies made on nucleons or nuclei. Although antinuclei up to antihelium-4 have been discovered¹ and their masses measured, little is known directly about the nuclear force between antinucleons. Here, we study antiproton pair correlations among data collected by the STAR experiment² at the Relativistic Heavy Ion Collider (RHIC)³, where gold ions are collided with a centre-of-mass energy of 200 gigaelectronvolts per nucleon pair. Antiprotons are abundantly produced in such collisions, thus making it feasible to study details of the antiproton-antiproton interaction. By applying a technique similar to Hanbury Brown and Twiss intensity interferometry⁴, we show that the force between two antiprotons is attractive. In addition, we report two key parameters that characterize the corresponding strong interaction: the scattering length and the effective range of the interaction. Our measured parameters are consistent within errors with the corresponding values for proton-proton interactions. Our results provide direct information on the interaction between two antiprotons, one of the simplest systems of antinucleons, and so are fundamental to understanding the structure of more-complex antinuclei and their properties.

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Article | Open Access | Published: 09 December 2020

Unveiling the strong interaction among hadrons at the LHC

ALICE Collaboration

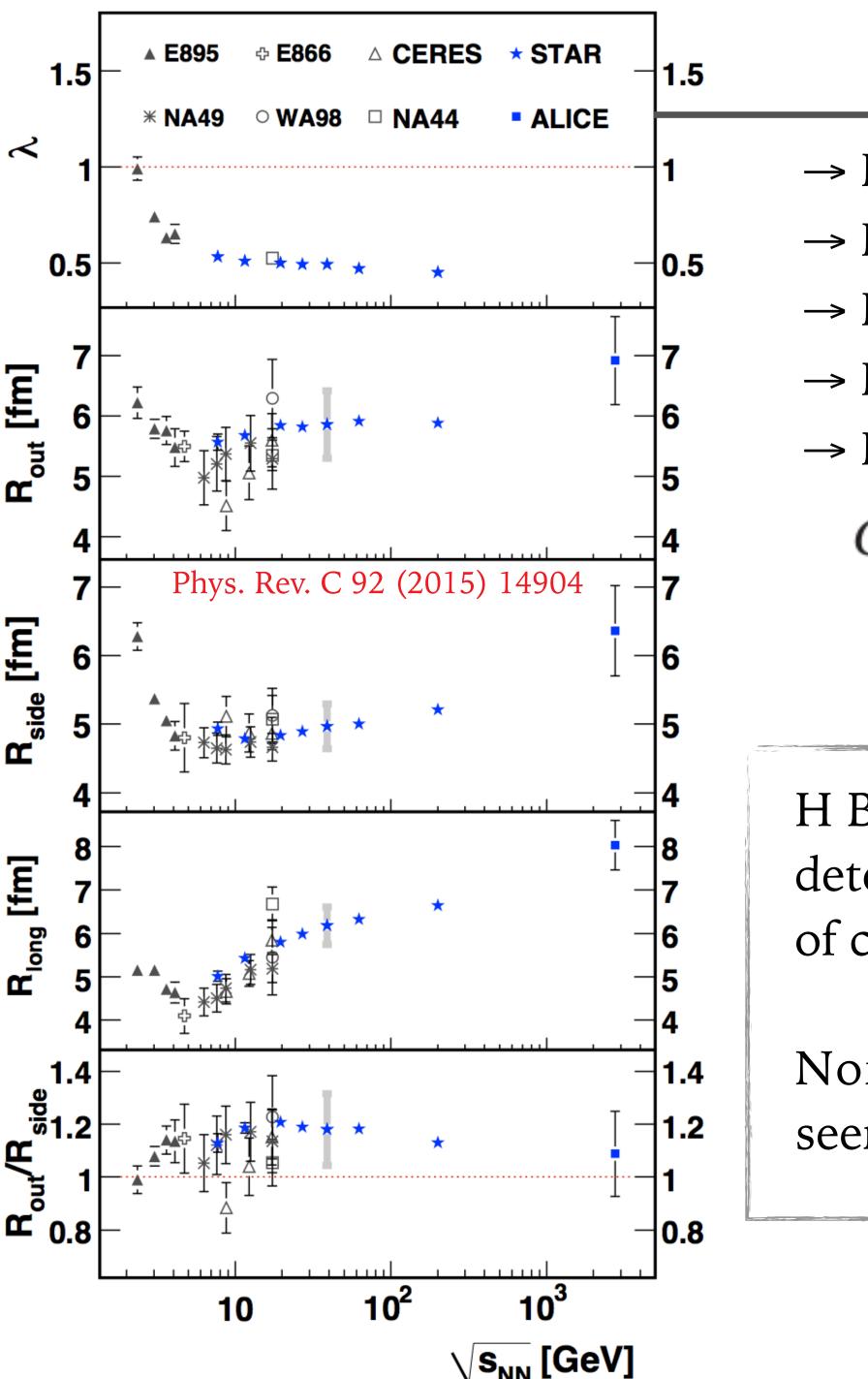
Nature588, 232–238 (2020)Cite this article9258Accesses6Citations231AltmetricMetrics

• A Publisher Correction to this article was published on 15 January 2021

• This article has been updated

Abstract

One of the key challenges for nuclear physics today is to understand from first principles the effective interaction between hadrons with different quark content. First successes have been achieved using techniques that solve the dynamics of quarks and gluons on discrete space-time lattices^{1,2}. Experimentally, the dynamics of the strong interaction have been studied by scattering hadrons off each other. Such scattering experiments are difficult or impossible for unstable hadrons^{3,4,5,6} and so high-quality measurements exist only for hadrons containing up and down quarks⁷. Here we demonstrate that measuring correlations in the momentum space between hadron pairs^{8,9,10,11,12} produced in ultrarelativistic proton-proton collisions at the CERN Large Hadron Collider (LHC) provides a precise method with which to obtain the missing information on the interaction dynamics between any pair of unstable hadrons. Specifically, we discuss the case of the interaction of baryons containing strange quarks (hyperons). We demonstrate how, using precision measurements of proton-omega baryon correlations, the effect of the strong interaction for this hadron-hadron pair can be studied with precision similar to, and compared with, predictions from lattice calculations^{13,14}. The large number of hyperons identified in proton-proton collisions at the LHC, together with accurate modelling 15 of the small (approximately one femtometre) inter-particle distance and exact predictions for the correlation functions, enables a detailed determination of the short-range part of the nucleon-hyperon interaction.



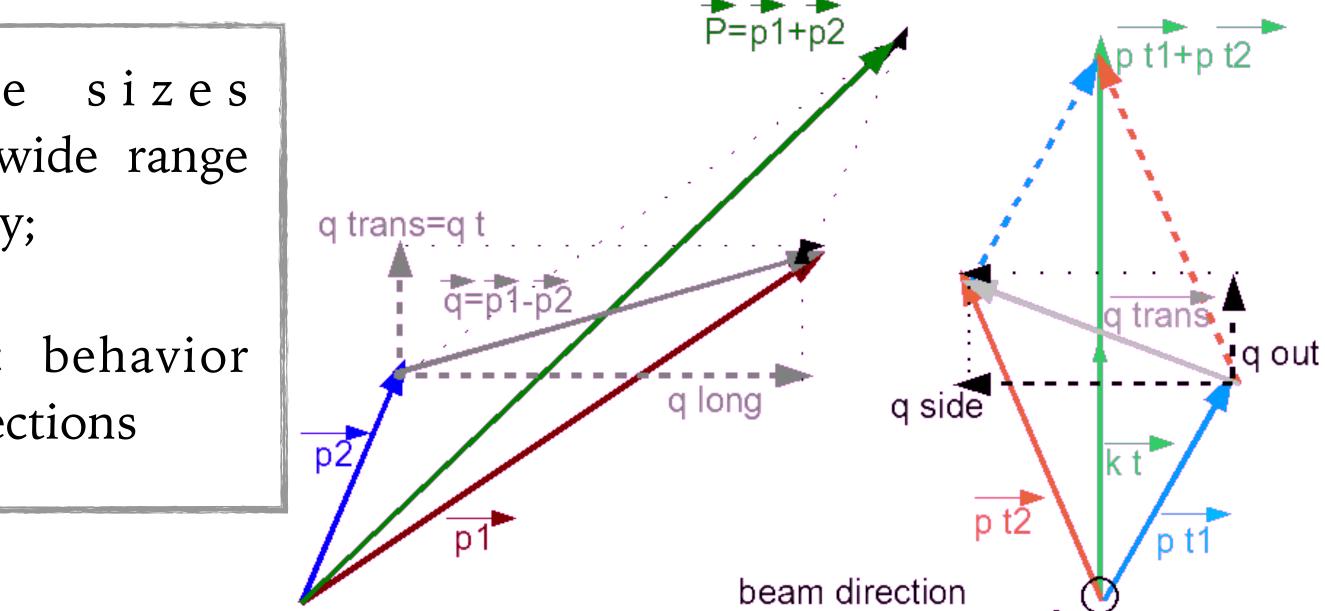
Identical pion femtoscopy

- \rightarrow R_{side} spatial source evolution in the transverse direction
- $\rightarrow R_{out}$ related to spatial and time components
- $\rightarrow R_{out}/R_{side}$ signature of phase transition
- \rightarrow R_{out}²- R_{side}² = $\Delta \tau^2 \beta_t^2$; $\Delta \tau$ emission time
- \rightarrow R_{long} temperature of kinetic freeze-out and source lifetime
 - $C(\vec{q}) = (1 \lambda)$ $\times \exp\left(-q_o^2 R\right)$

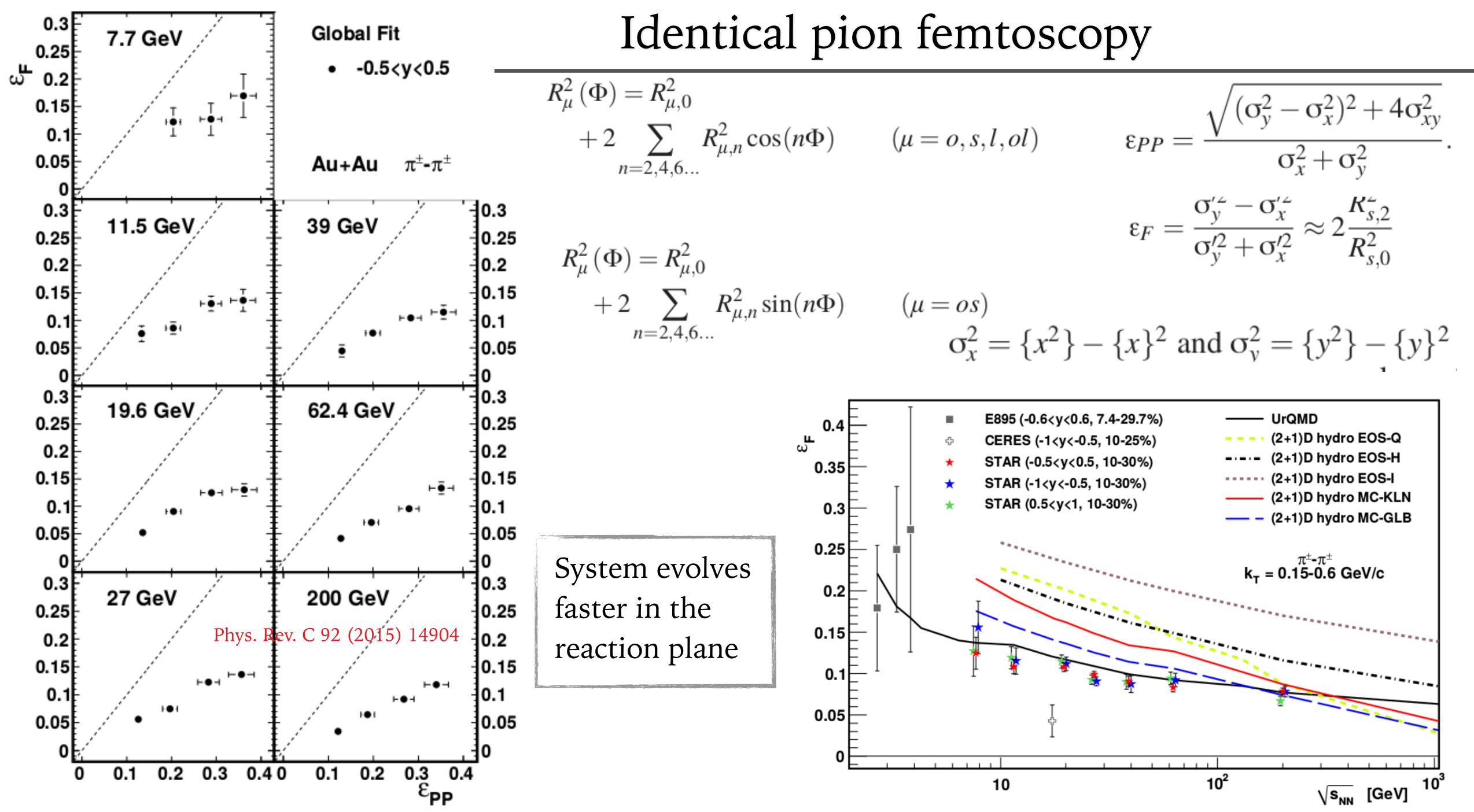
HBT source sizes determined for wide range of collision energy;

Non-monotonic behavior seen in three directions

$$k_{\rm Coul}(q_{\rm inv})\lambda R_o^2 - q_s^2 R_s^2 - q_l^2 R_l^2 - 2q_o q_s R_{os}^2 - 2q_o q_l R_{ol}^2$$

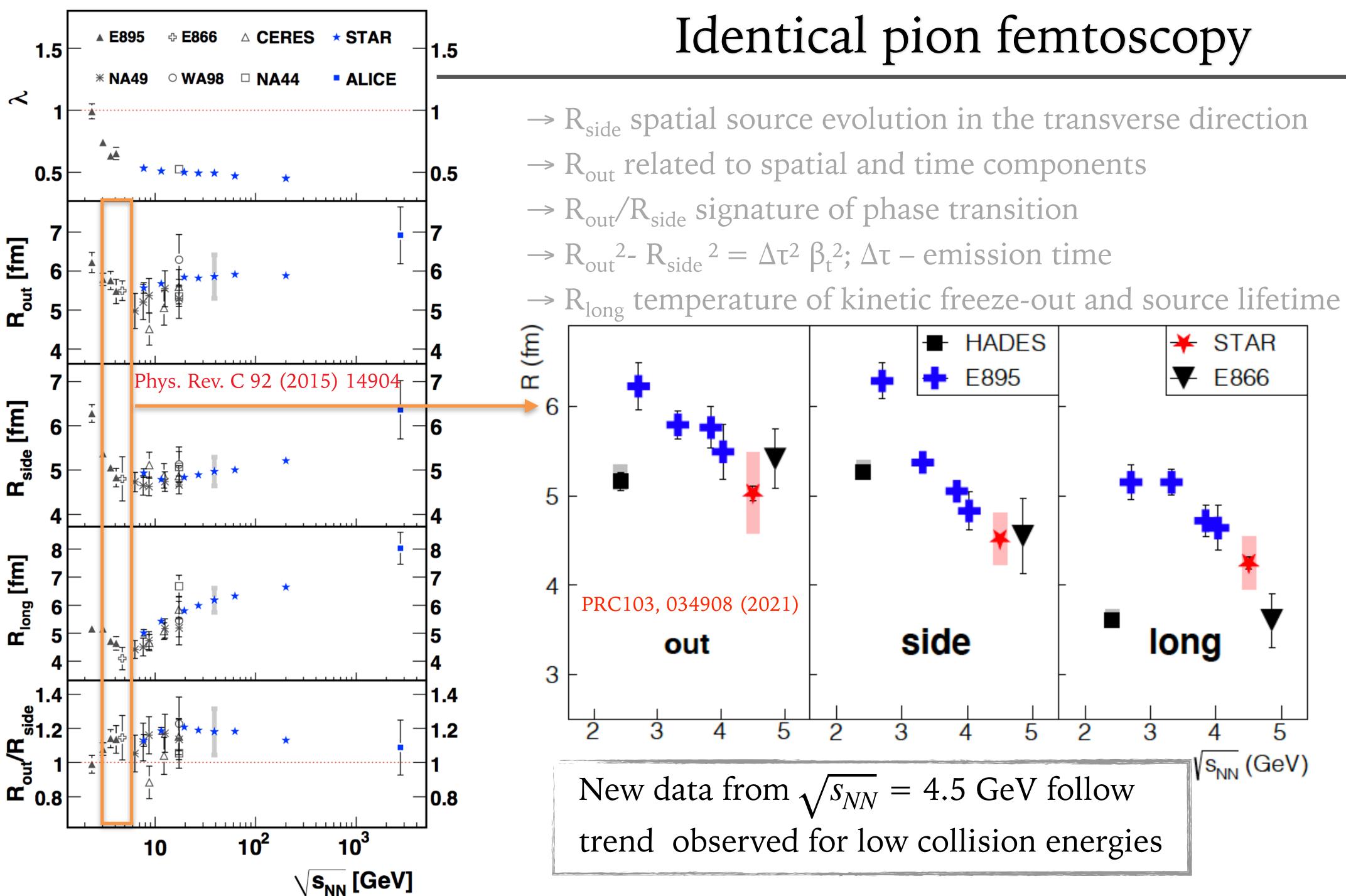






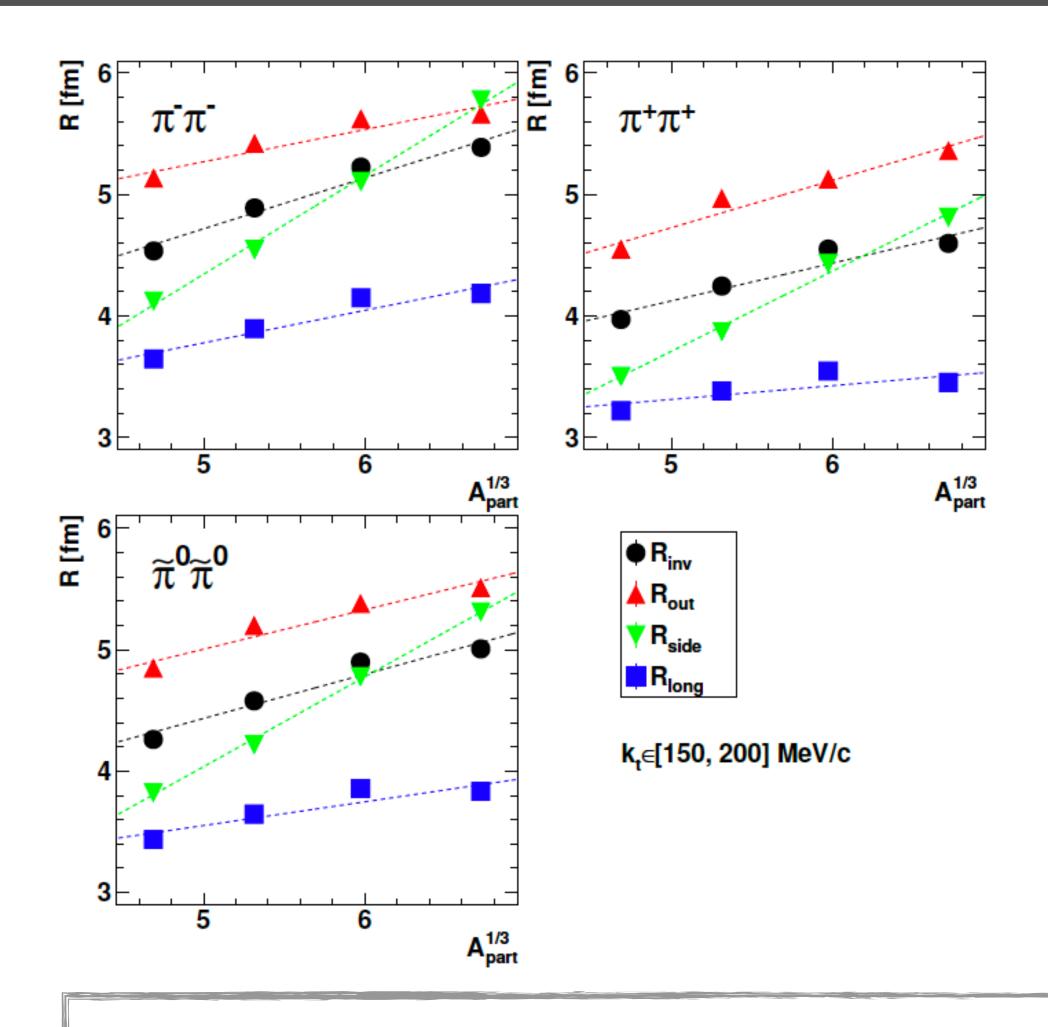
$$R_{\mu,n}^{2}\cos(n\Phi) \qquad (\mu = o, s, l, ol) \qquad \epsilon_{PP} = \frac{\sqrt{(\sigma_{y}^{2} - \sigma_{x}^{2})^{2} + \sigma_{y}^{2}}}{\sigma_{x}^{2} + \sigma_{y}^{2}}$$
$$\epsilon_{F} = \frac{\sigma_{y}^{\prime 2} - \sigma_{x}^{\prime 2}}{\sigma_{y}^{\prime 2} + \sigma_{x}^{\prime 2}} \approx 2\frac{R_{s,0}^{\prime 2}}{R_{s,0}^{2}}$$

$$R_{\mu,n}^{2} \sin(n\Phi) \qquad (\mu = os) \\ \sigma_{x}^{2} = \{x^{2}\} - \{x\}^{2} \text{ and } \sigma_{y}^{2} = \{y^{2}\} - \{$$

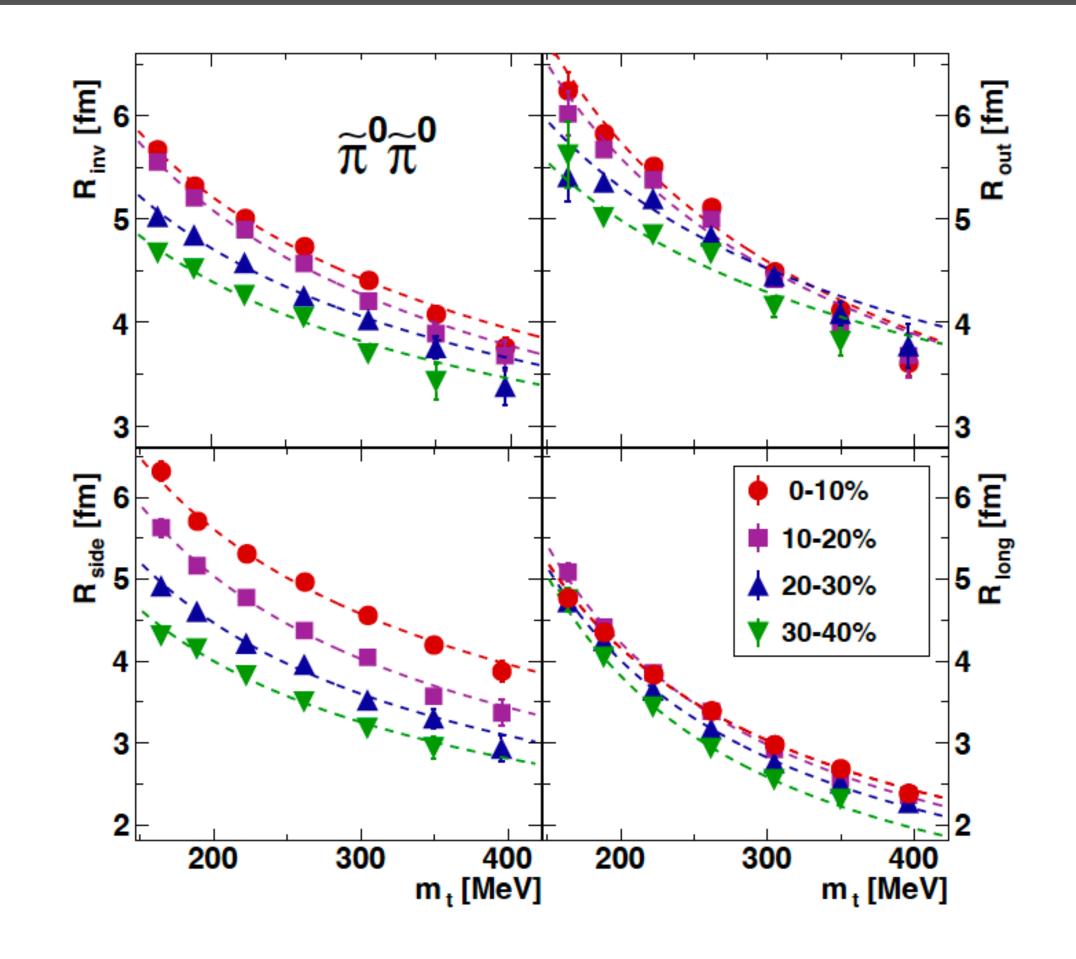




Correlations at HADES at $\sqrt{s_{NN}} = 2.4 \text{ GeV}$



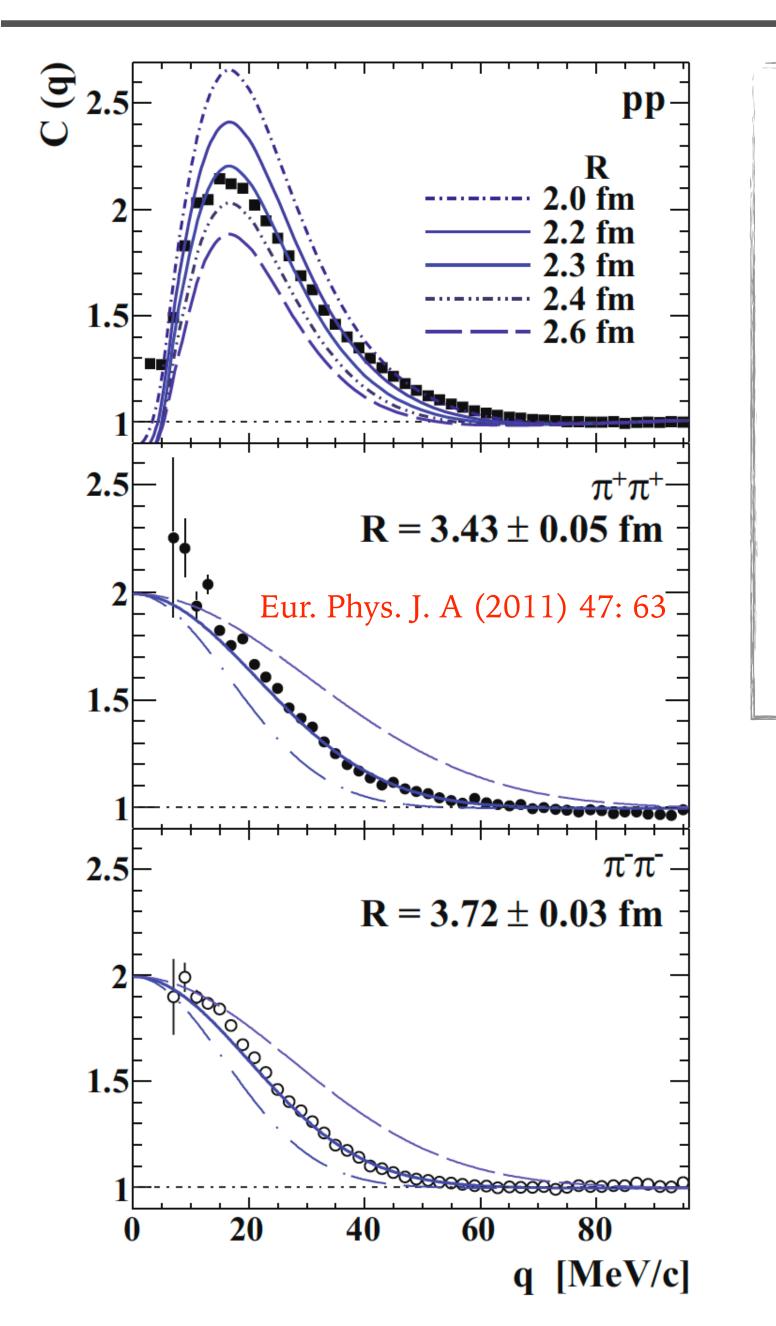
- $R_{long} < R_{out} \simeq R_{side}$
- Data fit with $R = R_0 (m_t/m_\pi)^{\alpha}$



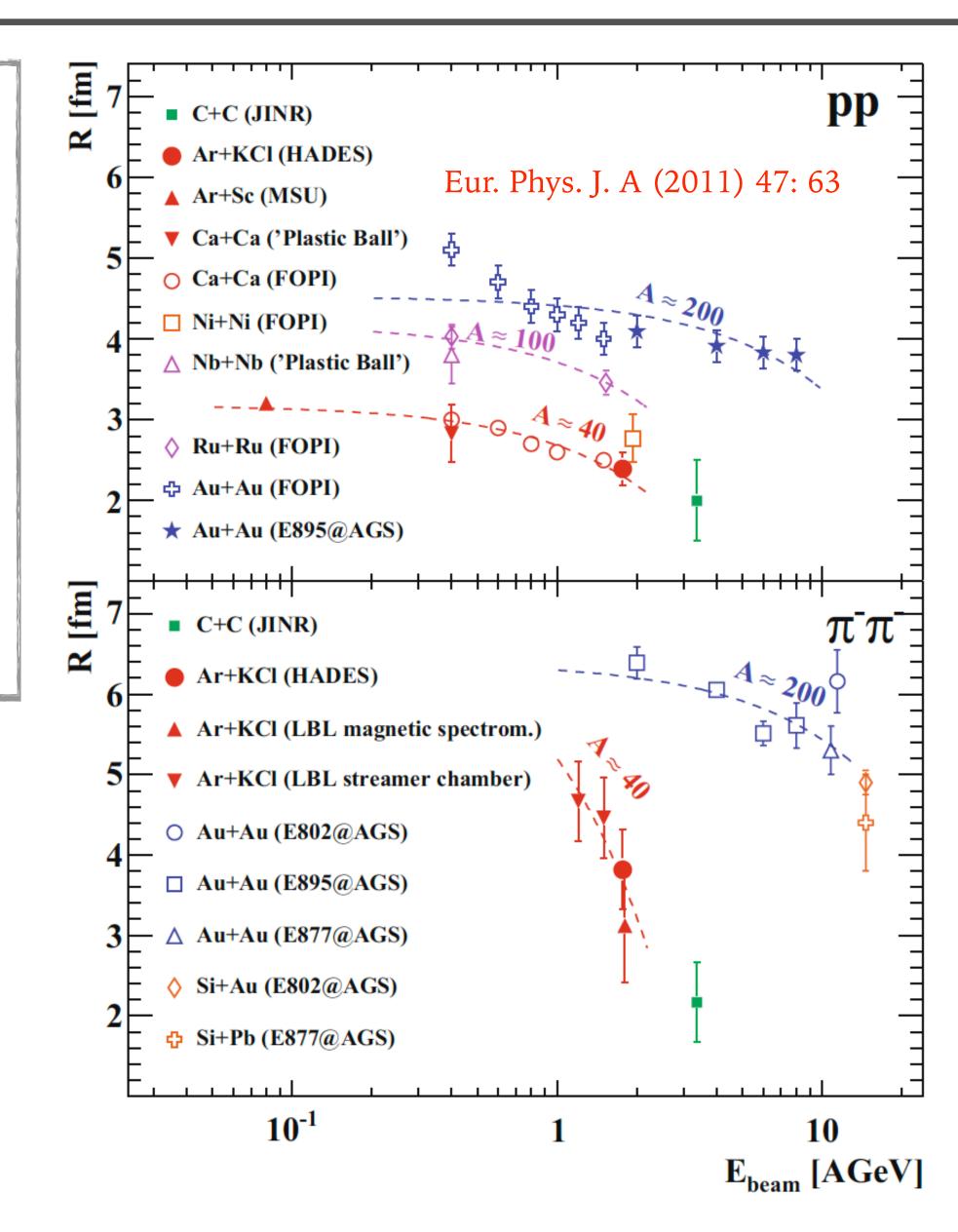
• Differences of the radii for pairs of π^- and π^+ are found (low transverse momenta).



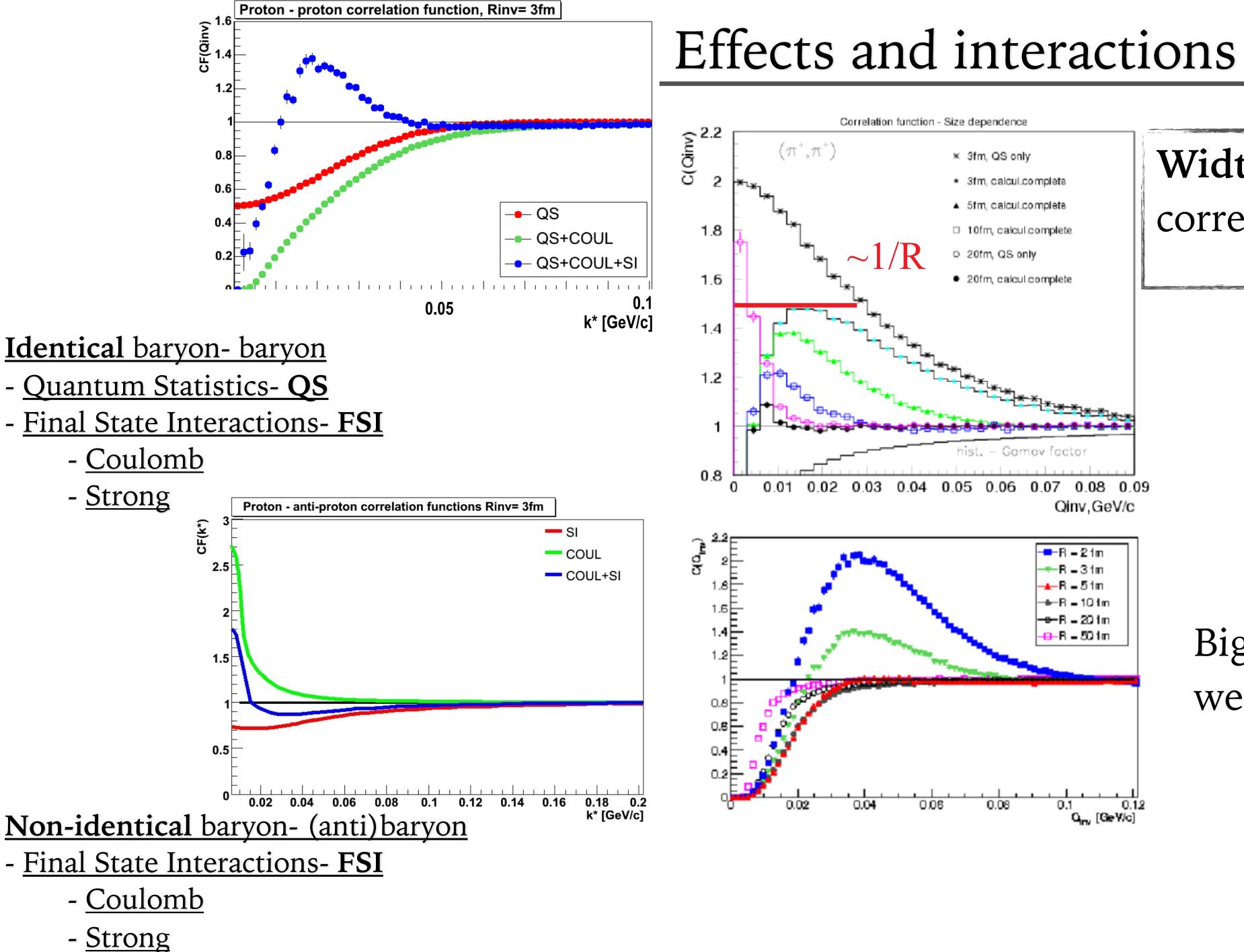
Correlations at HADES at Ar+KCl



Pions emitted from bigger sources than protons The present p - p and $\pi^- - \pi^-$ r a d i i well complement data trends in the SIS energy range.







Width of correlation function $\sim 1/R$

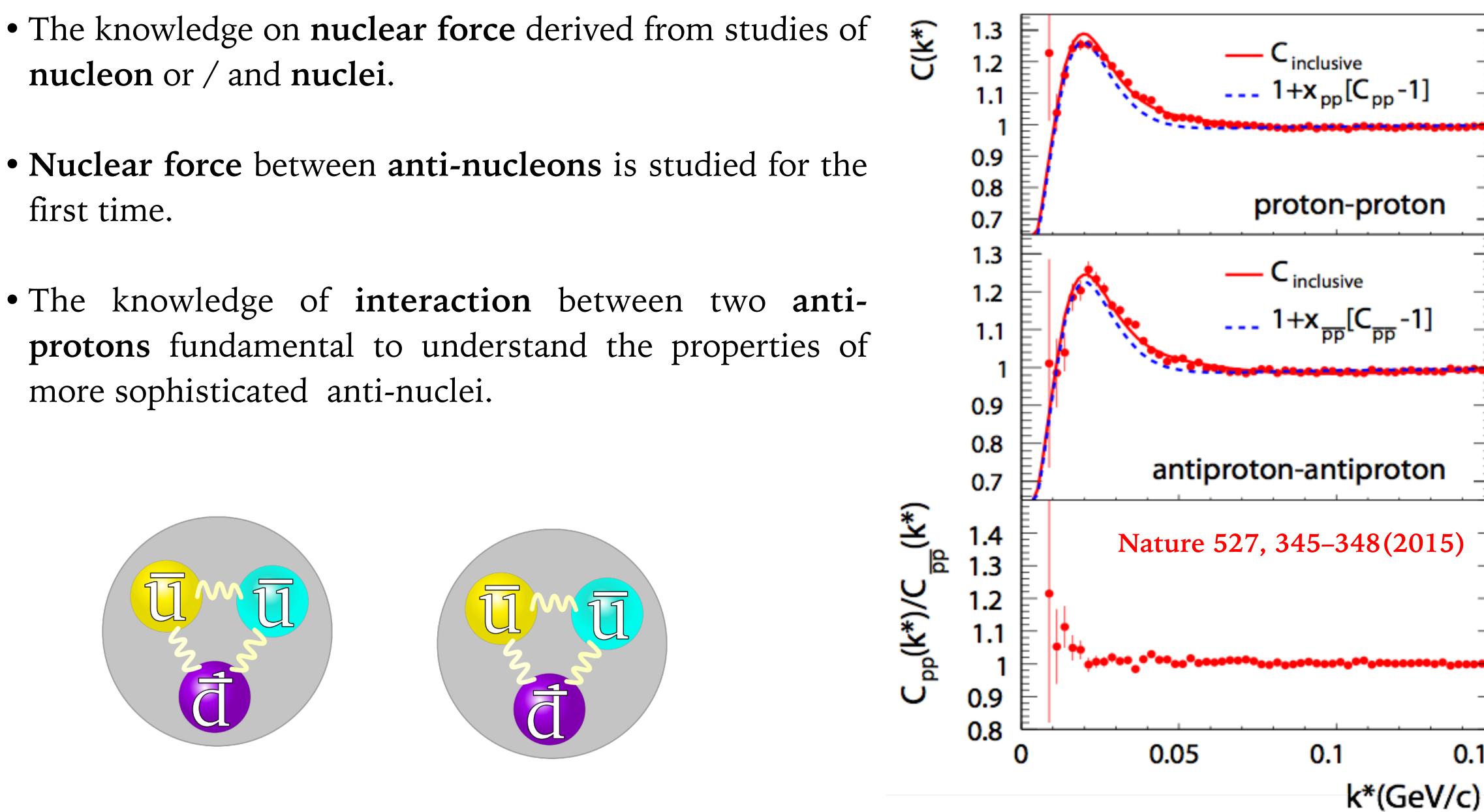
Bigger source and weaker correlation





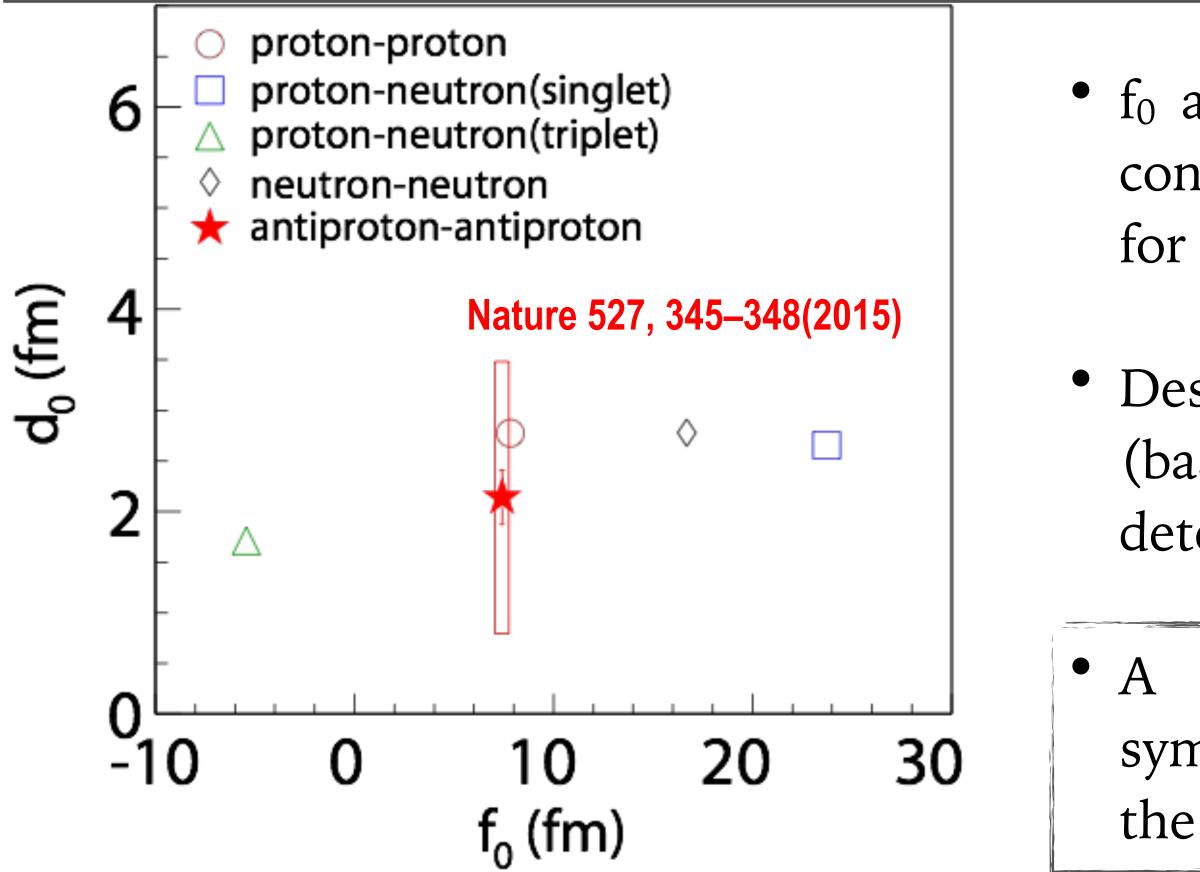
1) Strong interactions between anti-nucleons

- nucleon or / and nuclei.
- first time.
- more sophisticated anti-nuclei.



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.1	5.8

1) Strong interactions between anti-nucleons



The scattering length f₀: determines low-energy scattering.

The elastic cross section, σ_e , (at low energies) determined solely by the scattering length, $\lim_{k \to 0} \sigma_e = 4\pi f_0^2$

 d_0 - the effective range of strong interaction between two particles. It corresponds to the range of the potential in an extremely simplified scenario - the square well potential. f_0 and d_0 - two important parameters of strong interaction between two particles. •

- Theoretical correlation function depends on: source size, k^* , f_0 and d_0 . ۲

• f_0 and d_0 for the antiproton-antiproton interaction consistent with parameters

for the proton-proton interaction.

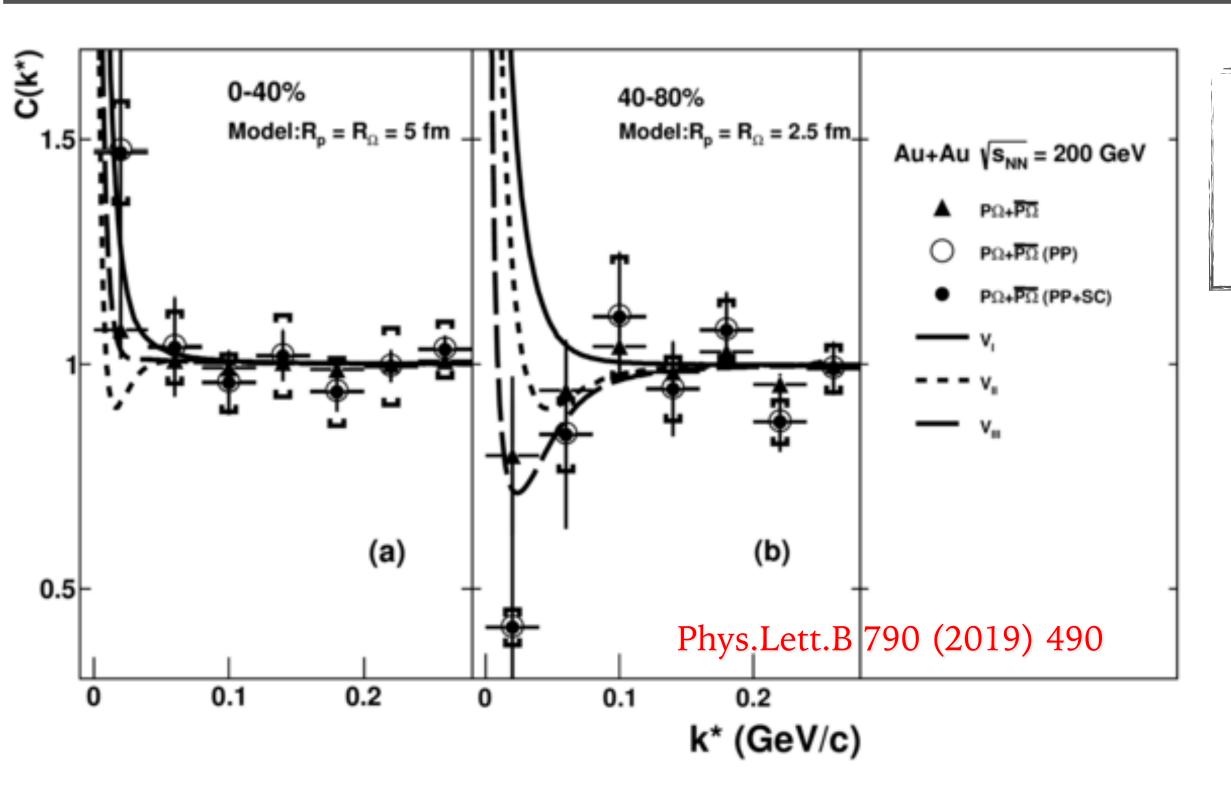
• Descriptions of the interaction among antimatter (based on the simplest systems of anti-nucleons) determined.

• A quantitative verification of matter-antimatter symmetry in context of the forces responsible for the binding of (anti)nuclei.



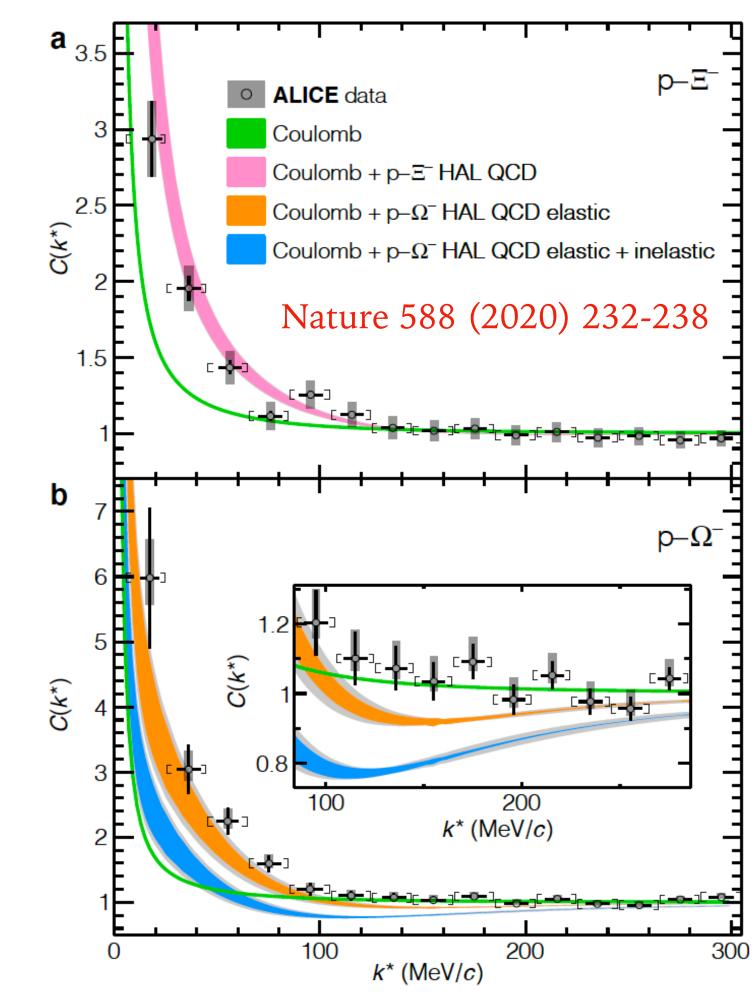


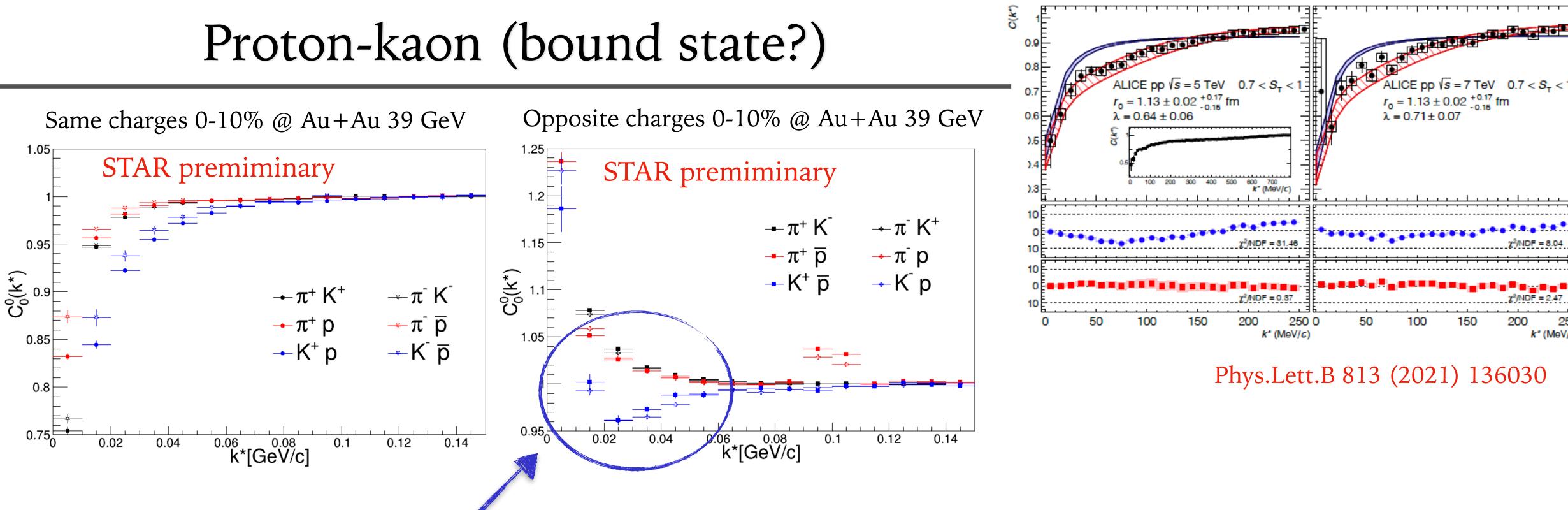
Strange Baryon Correlations (including p- Ω)



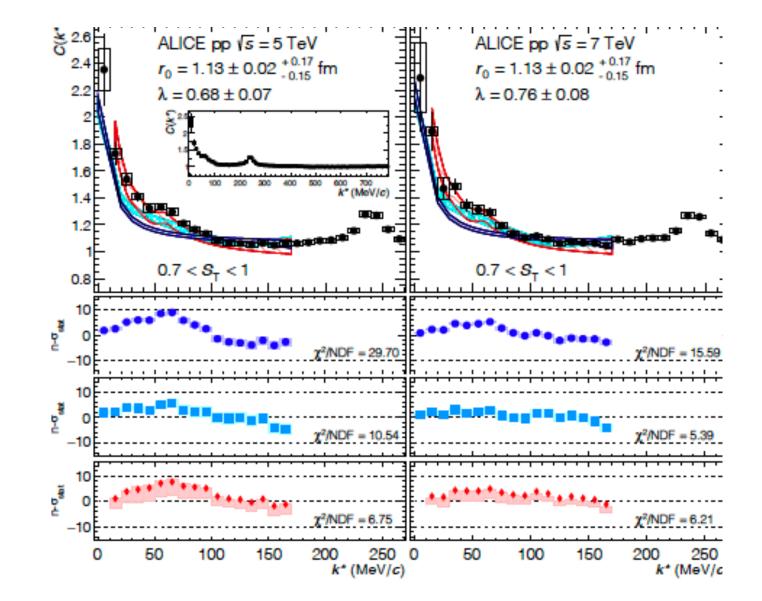
	V1	V2	V3
Ebin [MeV]	-	6.3	26.9
a ⁰ [MeV]	-1.12	5.79	1.29
reff [MeV]	-1.16	0.96	0.65

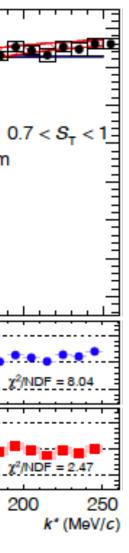
Scattering lenght positive, favor the hypothesis of $p\Omega$ bound state





- High-precision measurement of the strong interaction (anti-correlation) between kaons and protons.
- A structure (ALICE in p+p collisions) observed around a relative momentum of 58 MeV/c in the measured correlation function of opposite charges in p+p collisions.

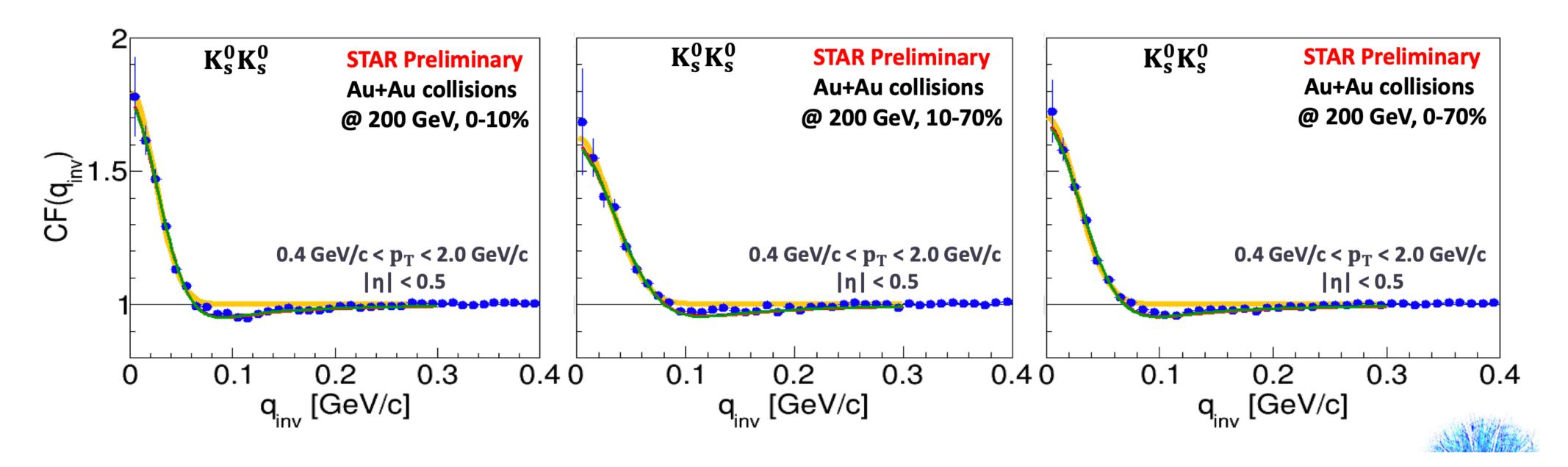








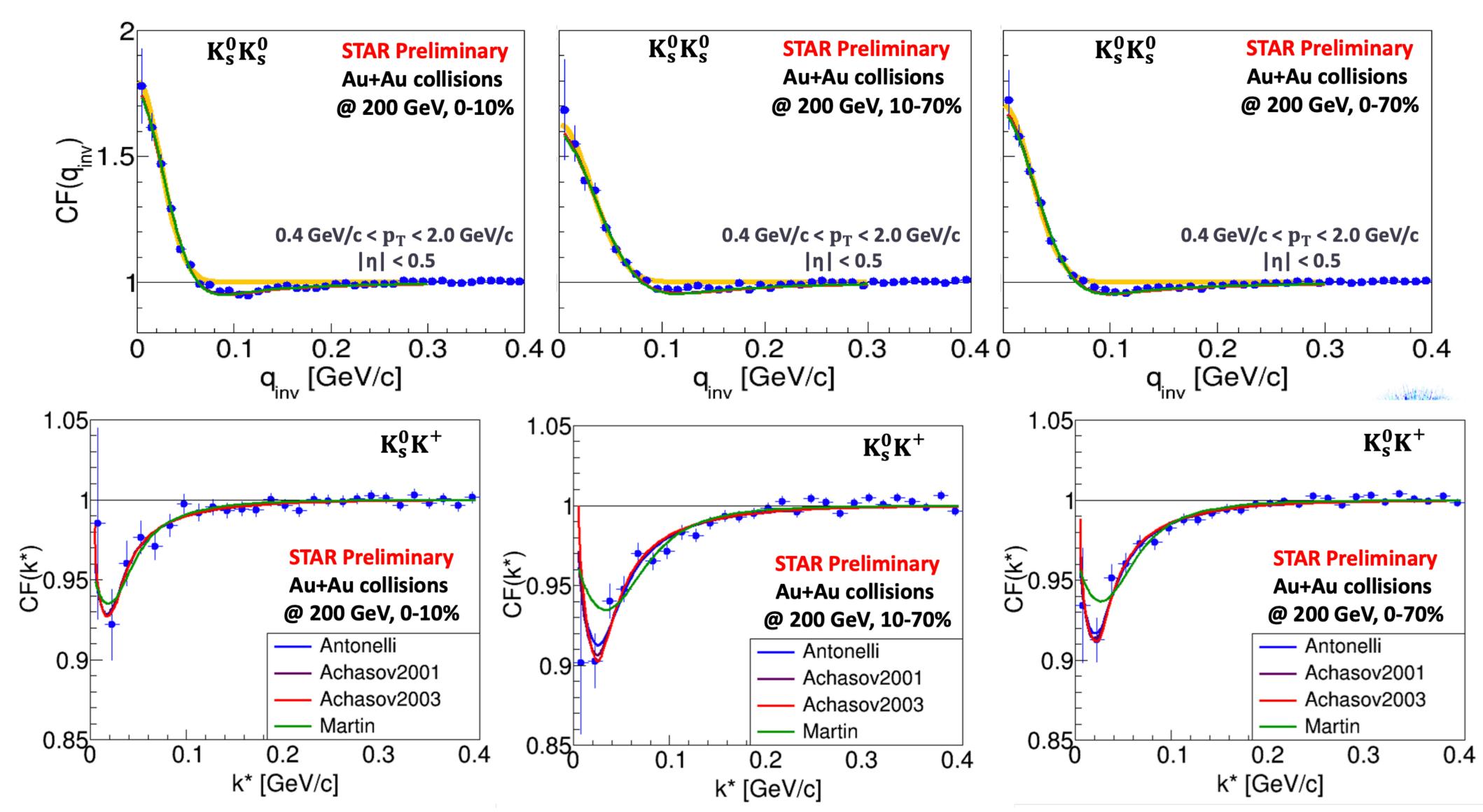
Neutral kaons



The strong final-state interaction has a significant effect on the neutral kaons correlation due to the near-threshold $f_0(980)$ and $a_0(980)$ resonances



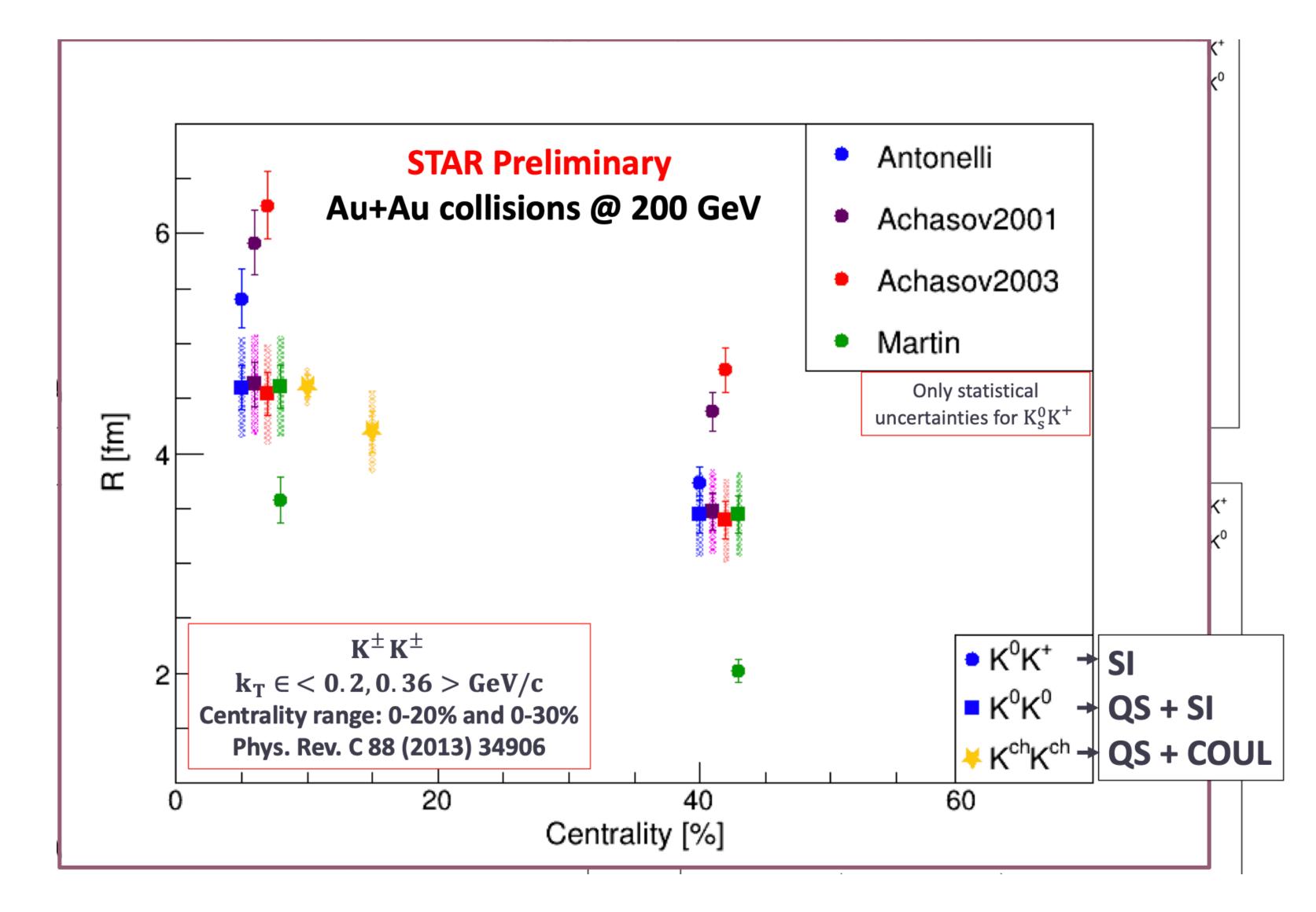
Neutral kaons



The $a_0(980)$ FSI parametrization gives very good representation of the shape of the signal region in CF; the parametrization with the larger $a_0(980)$ mass and decay coupling gives larger size of the source; Antonelli parametrization favors $a_0(980)$ resonance as a tetraquark



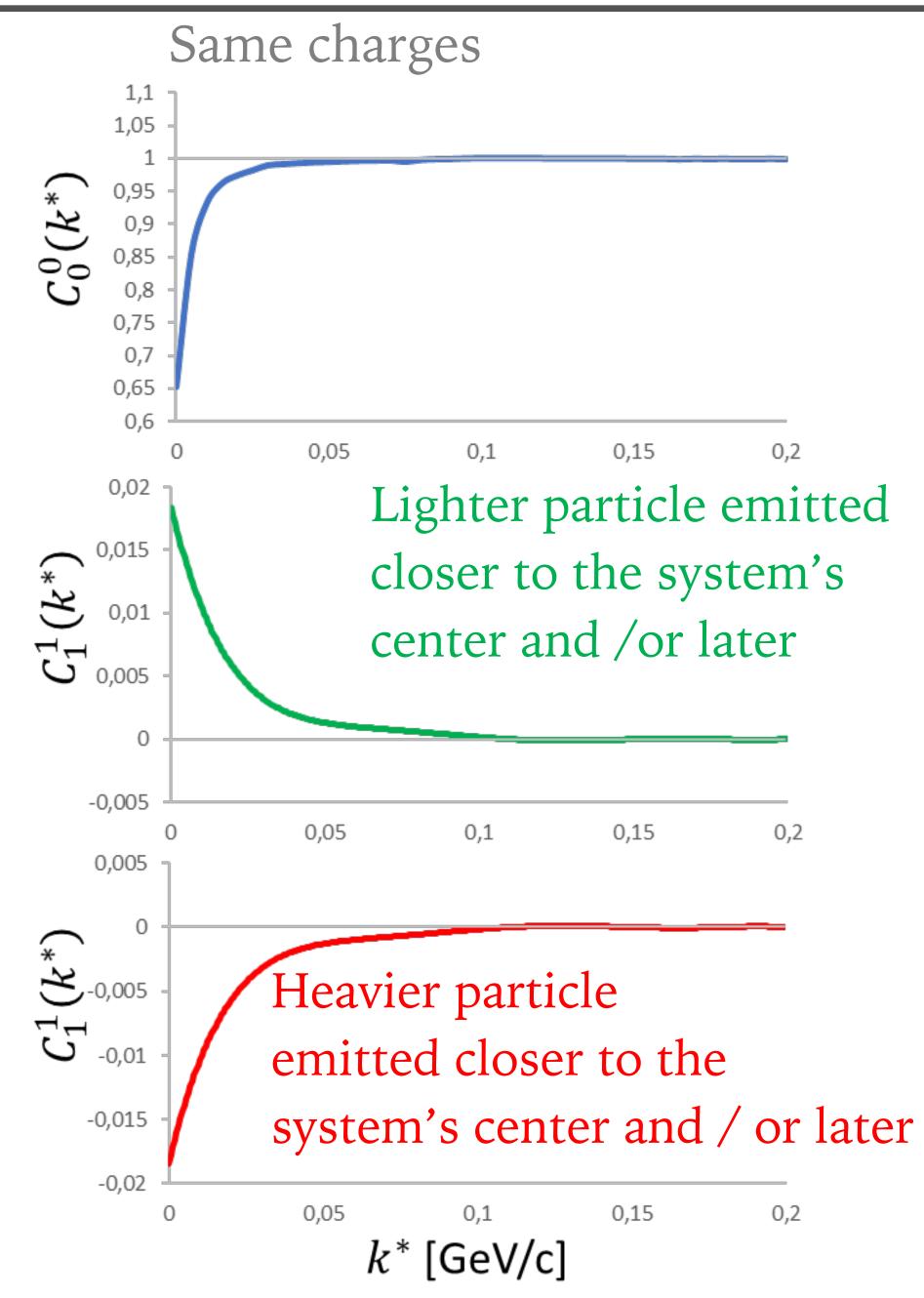
Neutral kaons

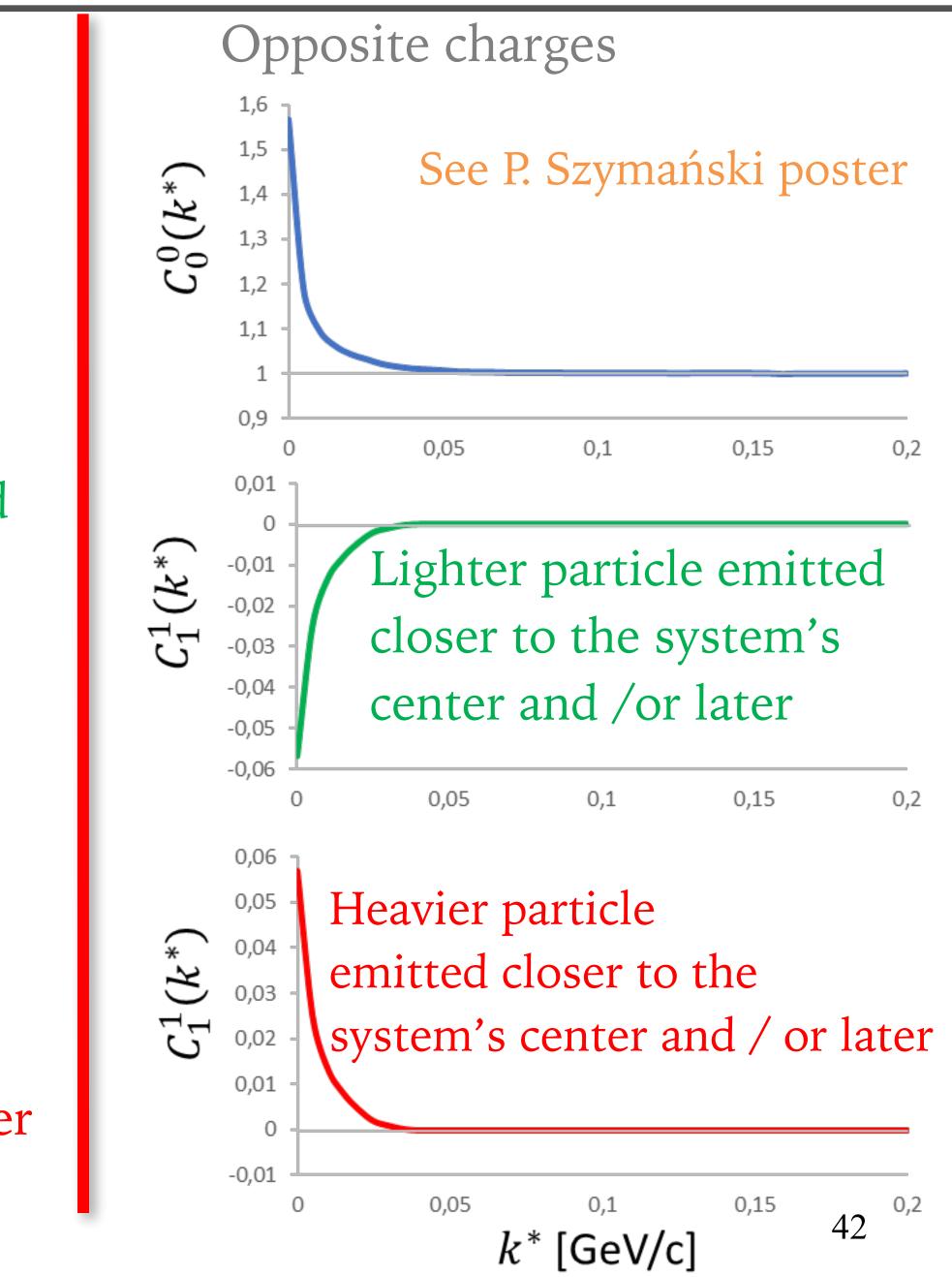


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Nonidentical particles - emission asymmetry





Nonidentical particle correlations

