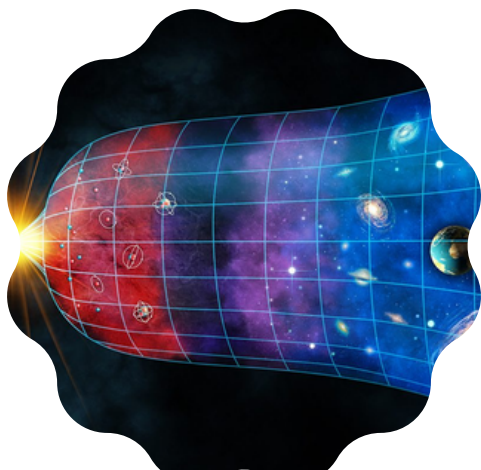
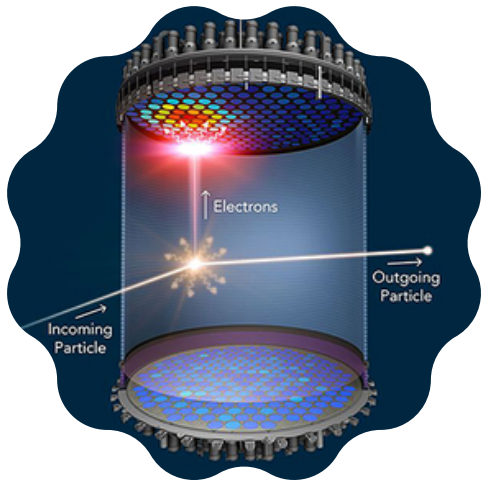


Complementarity of direct detection, indirect detection and collider searches (LHC)

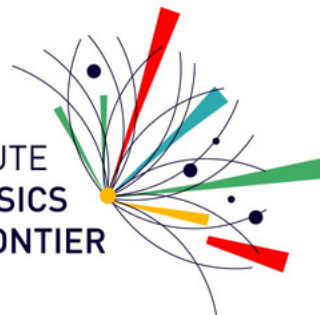
Farinaldo Queiroz

UFRN

The 2nd DMNet International Symposium, Heidelberg
2022



MILLENNIUM INSTITUTE
FOR SUBATOMIC PHYSICS
AT HIGH-ENERGY FRONTIER
SAPHIR



What you've seen so far...

DM models, properties and particle physics candidates, Takashi Toma

Searching for Dark Matter with HyperK, Nicole Bell

Dark matter in galaxy clusters from X-ray & SZ effect, Stefano Ettori

Some gamma-ray talks

Astrophysical anomalies

Direct and Indirect Detection of Dark Matter

The 2nd DMNet International Symposium

Sep 13 – 15, 2022
Max Planck Institute for Nuclear Physics, Heidelberg, Germany
Europe/Zurich timezone

Emphasize the importance of complementarity

Dark matter sensitivity of neutrino detectors

Take away messages



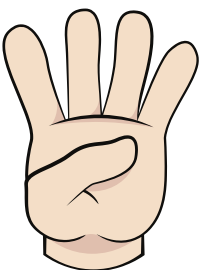
Without interdisciplinary searches one cannot claim the discovery of dark matter particles



The next generation of experiments will exclude most simplified models, but not the WIMP paradigm



The progress on the direct, indirect and collider experiments will move at very different paces in the near future



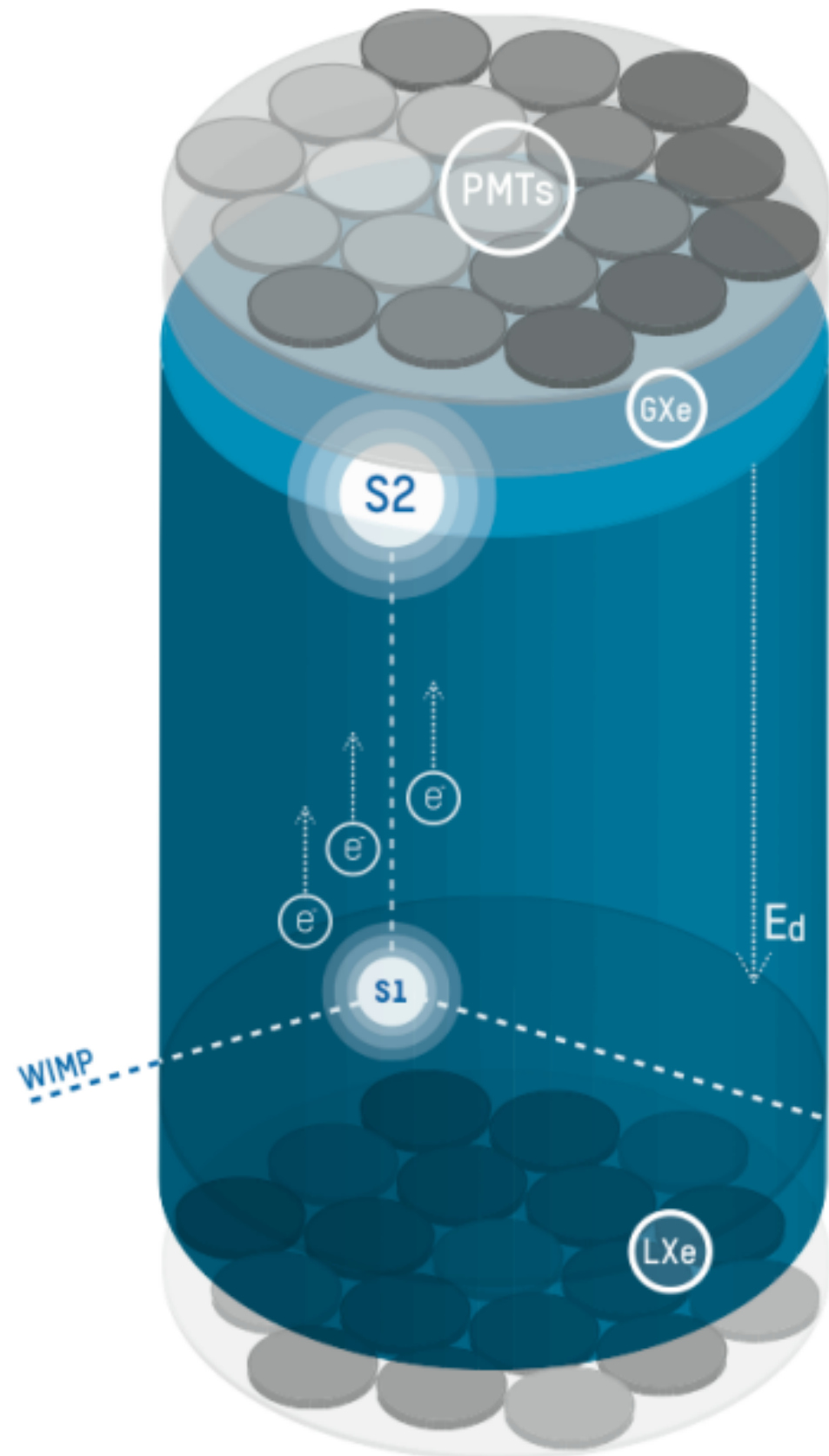
**Part of the field is moving to a multifold dark matter search:
Direct, indirect, collider, *neutron stars, *gravitational waves**



Part of the field is moving to alternative dark matter production mechanisms motivated by "Why not? "



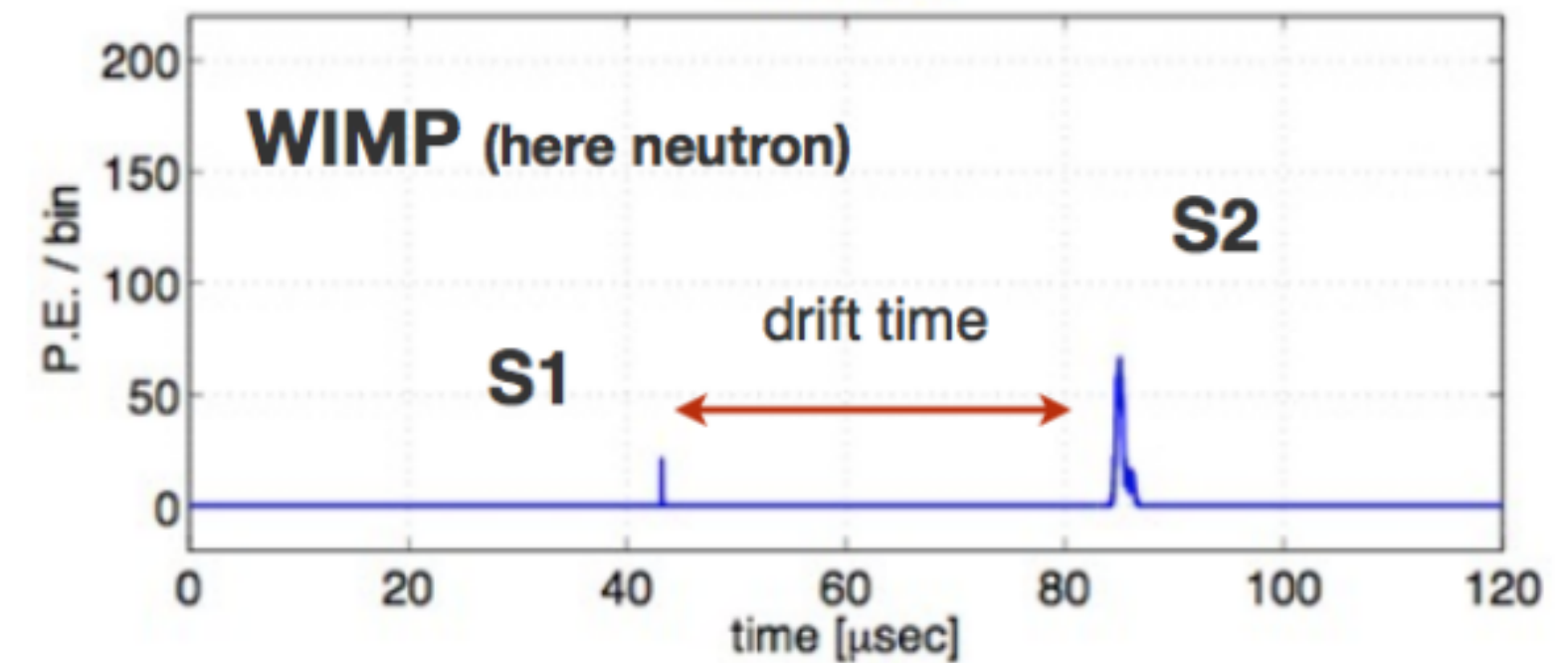
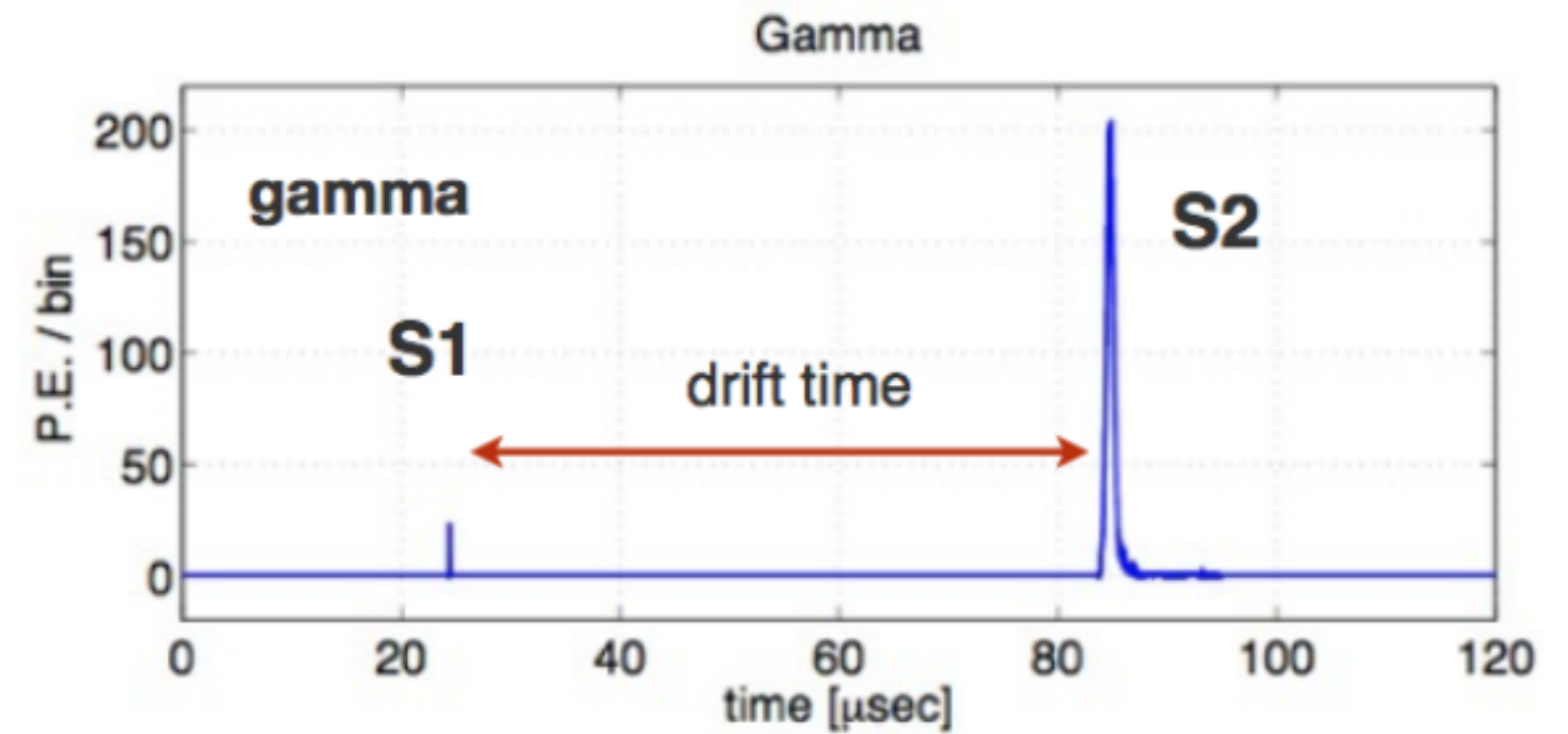
Direct Detection

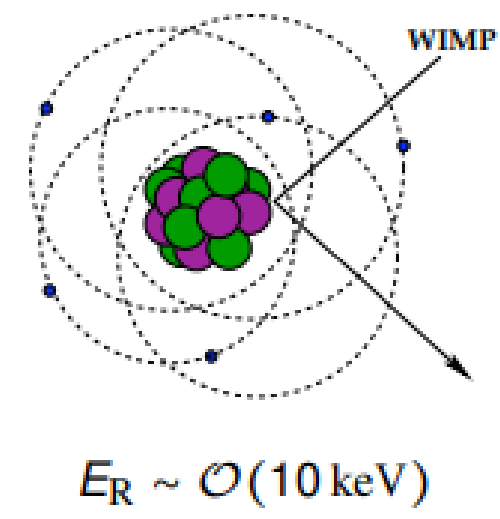
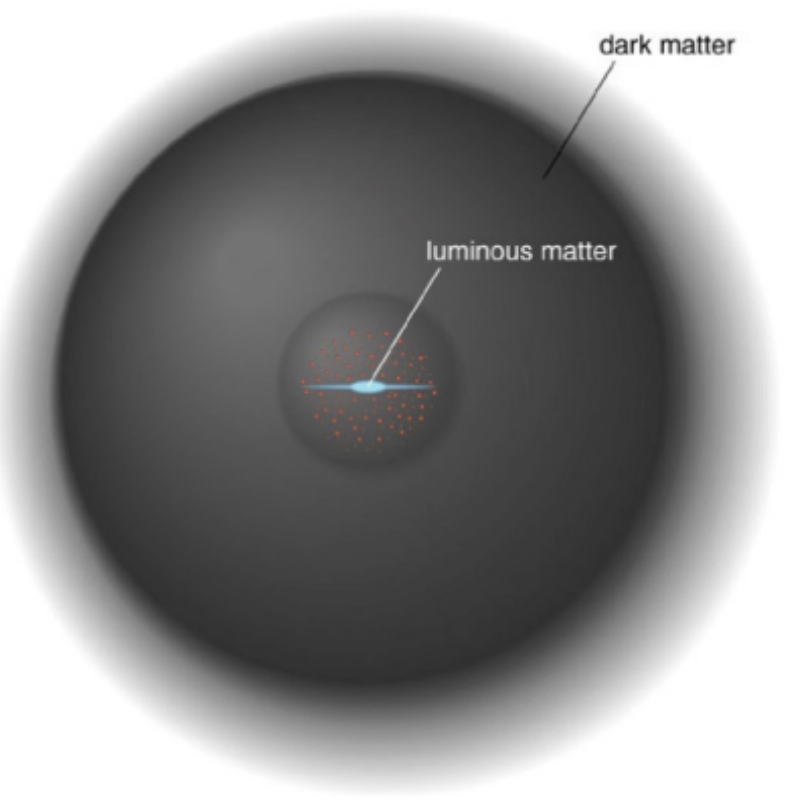


SIGNAL x Background

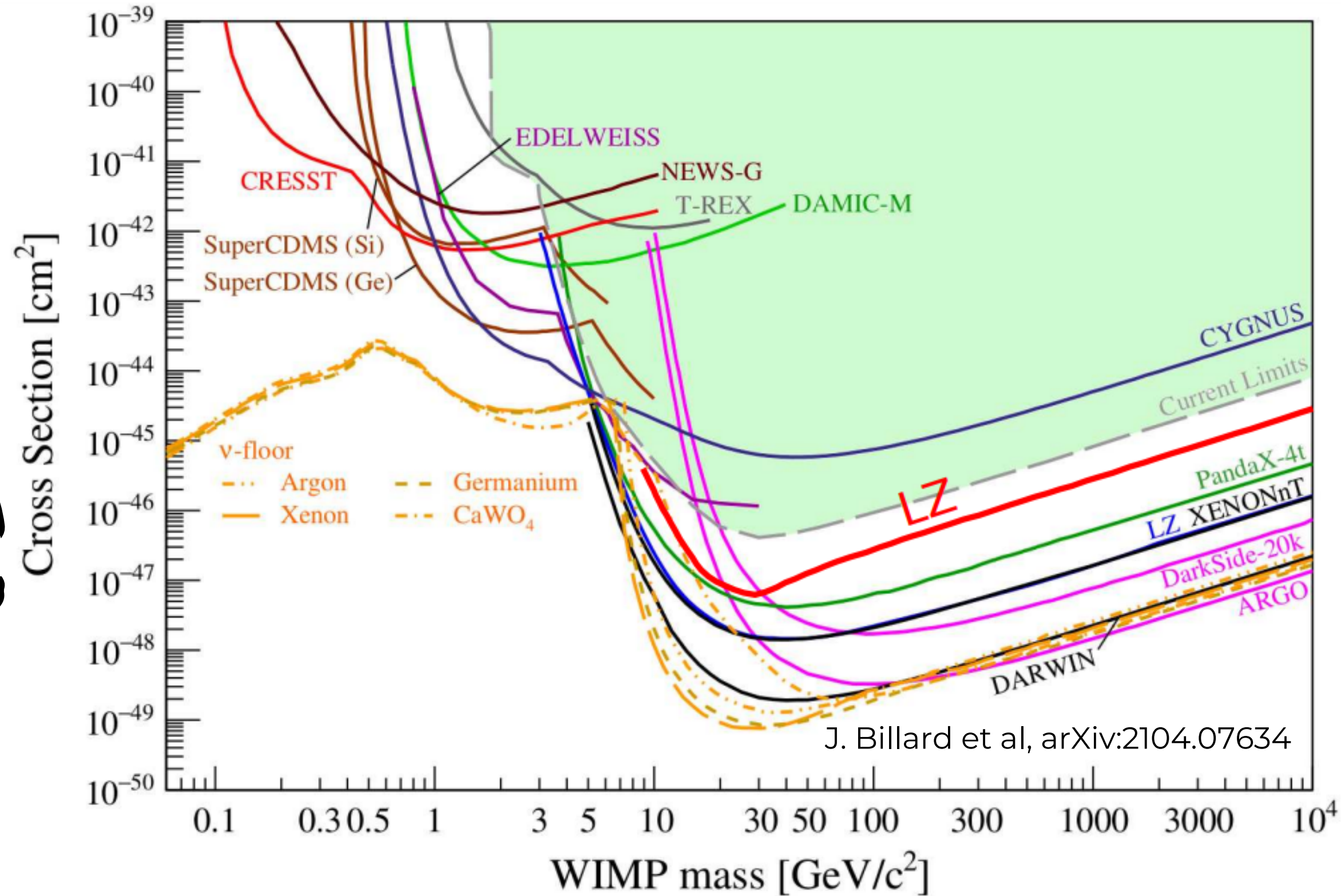
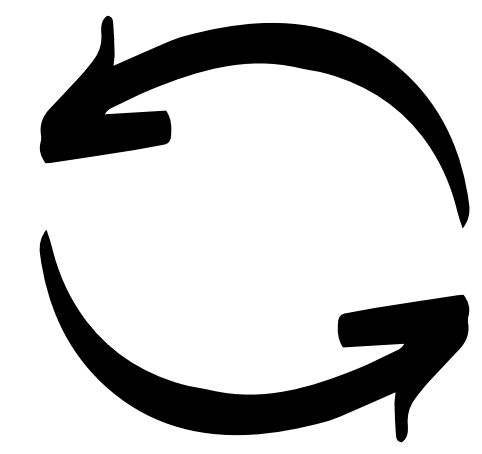
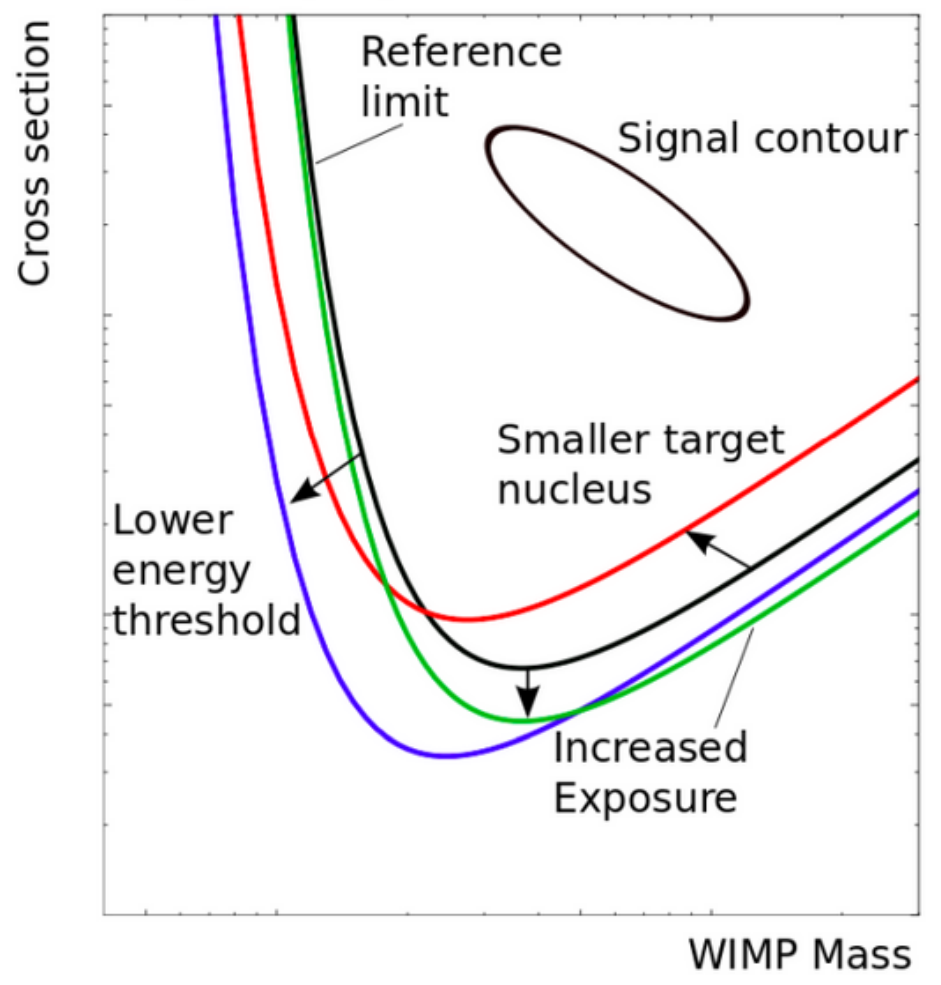
It is related to the dark matter local density.

I don't want to run from it

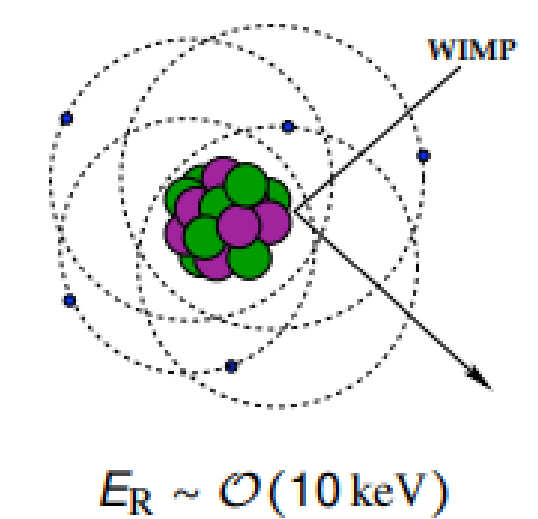
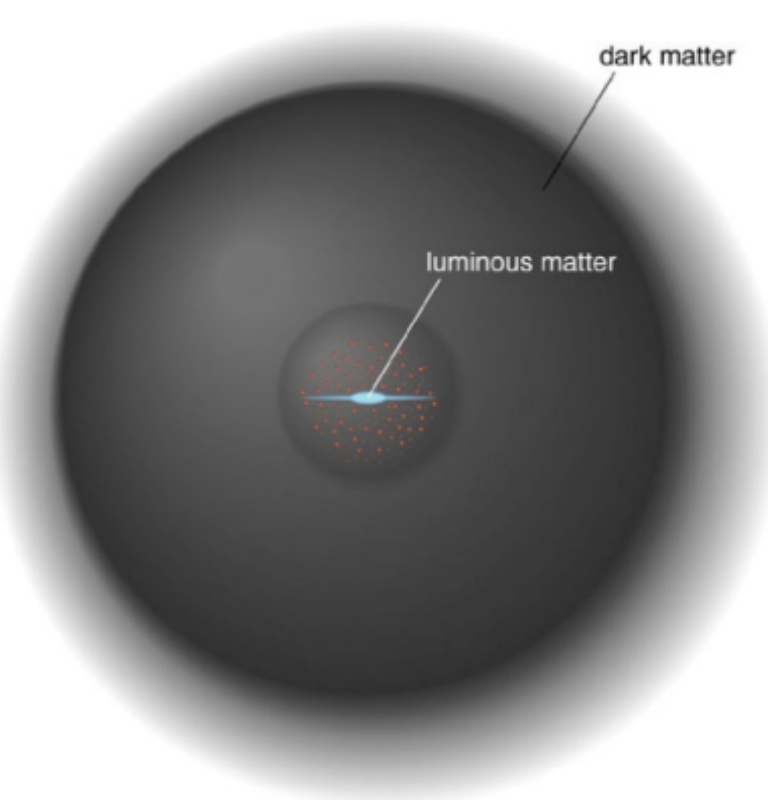




$$\frac{dR}{dE}(E, t) = \frac{\rho_0}{m_\chi \cdot m_A} \cdot \int \mathbf{v} \cdot \mathbf{f}(\mathbf{v}, t) \cdot \frac{d\sigma}{dE}(E, \mathbf{v}) d^3v$$

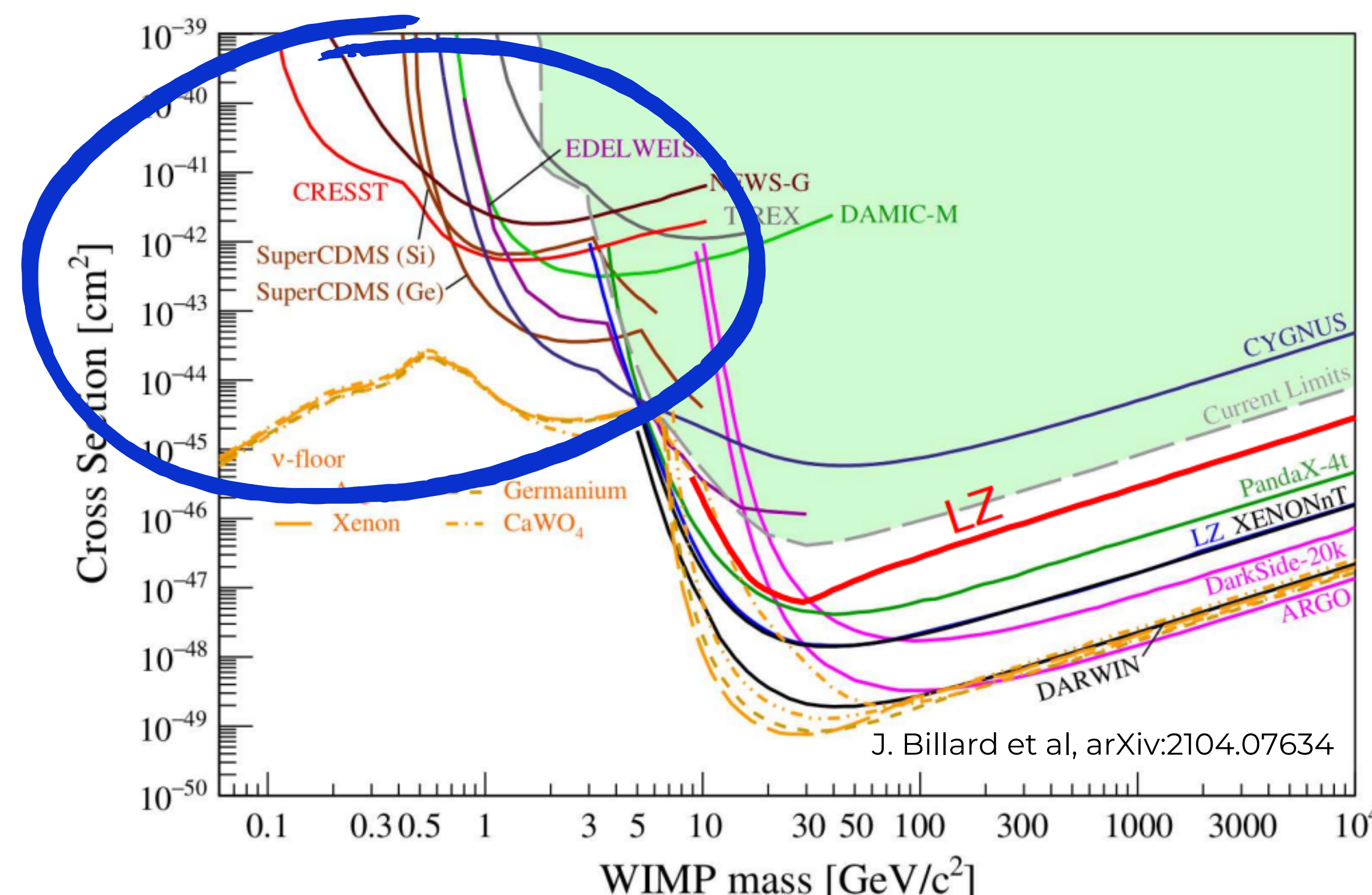


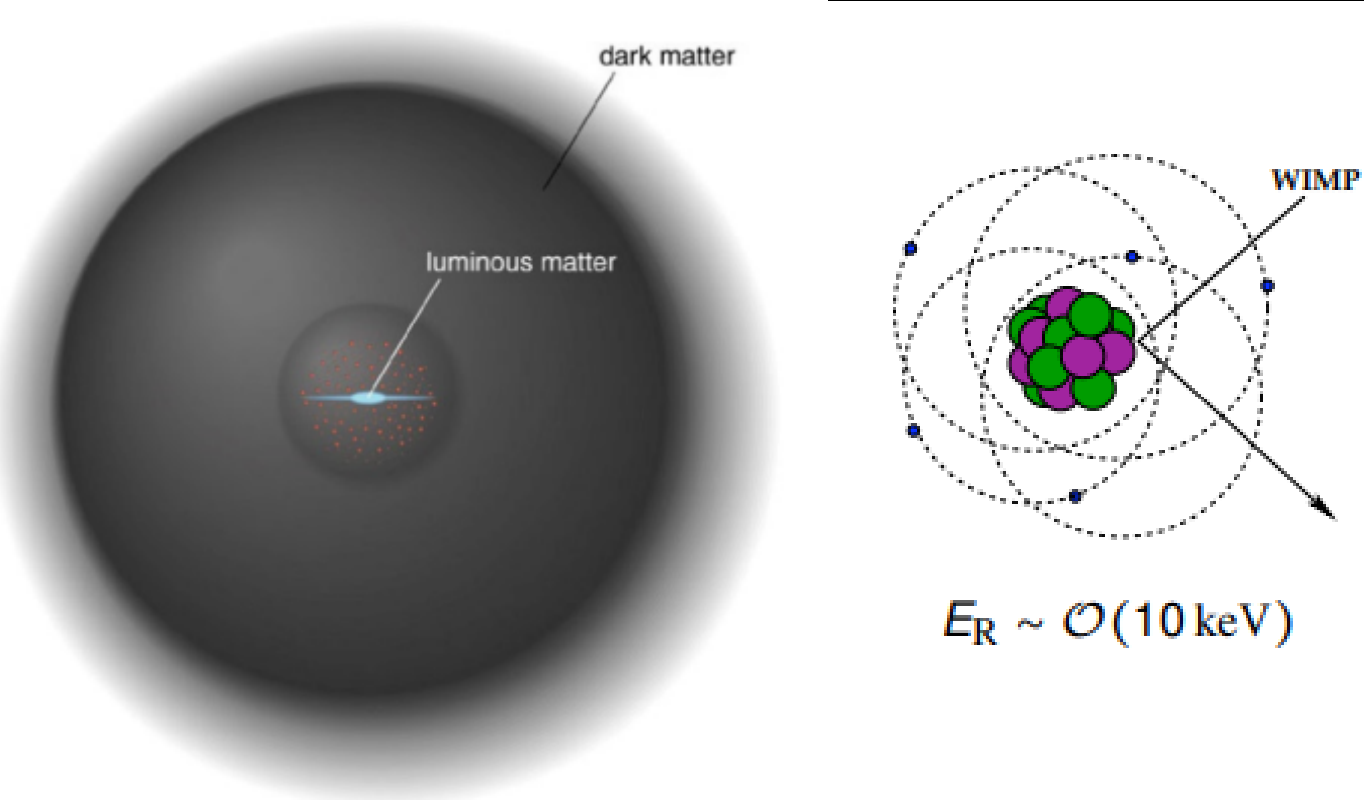
J. Billard et al, arXiv:2104.07634



$$\frac{dR}{dE}(E, t) = \frac{\rho_0}{m_\chi \cdot m_A} \cdot \int \mathbf{v} \cdot \mathbf{f}(\mathbf{v}, t) \cdot \frac{d\sigma}{dE}(E, \mathbf{v}) d^3v$$

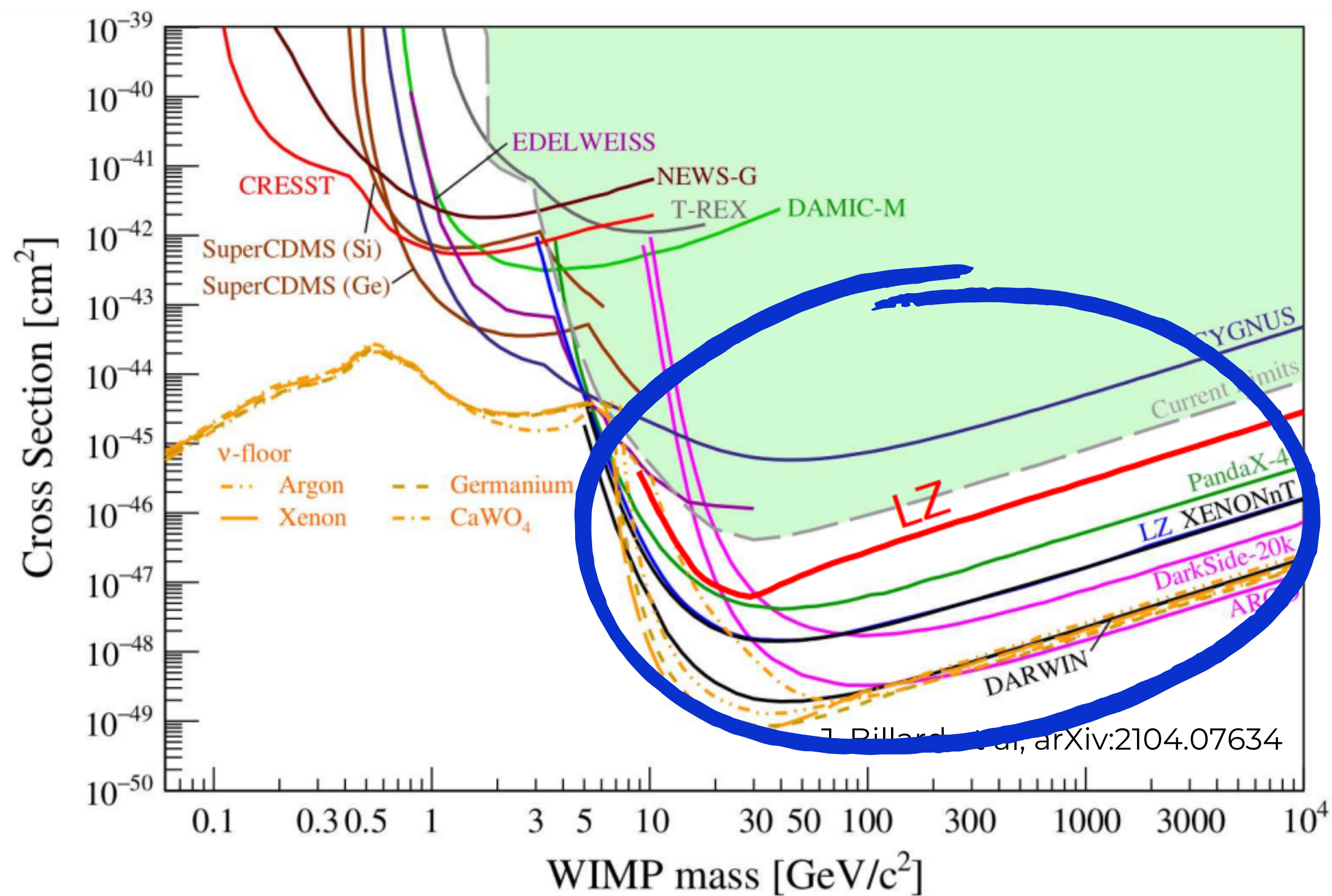
Sub-GeV dark matter?
pseudoscalar mediators
Non-standard Cosmologies



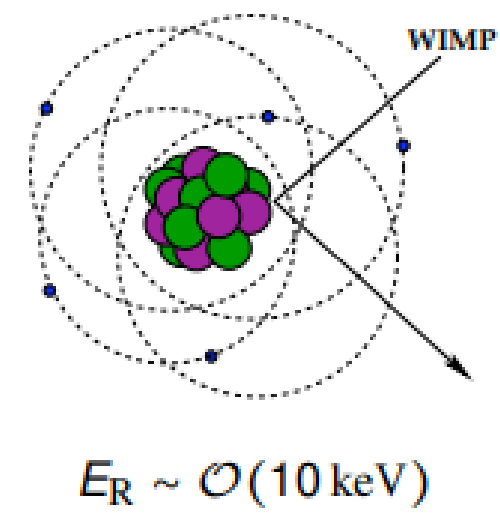
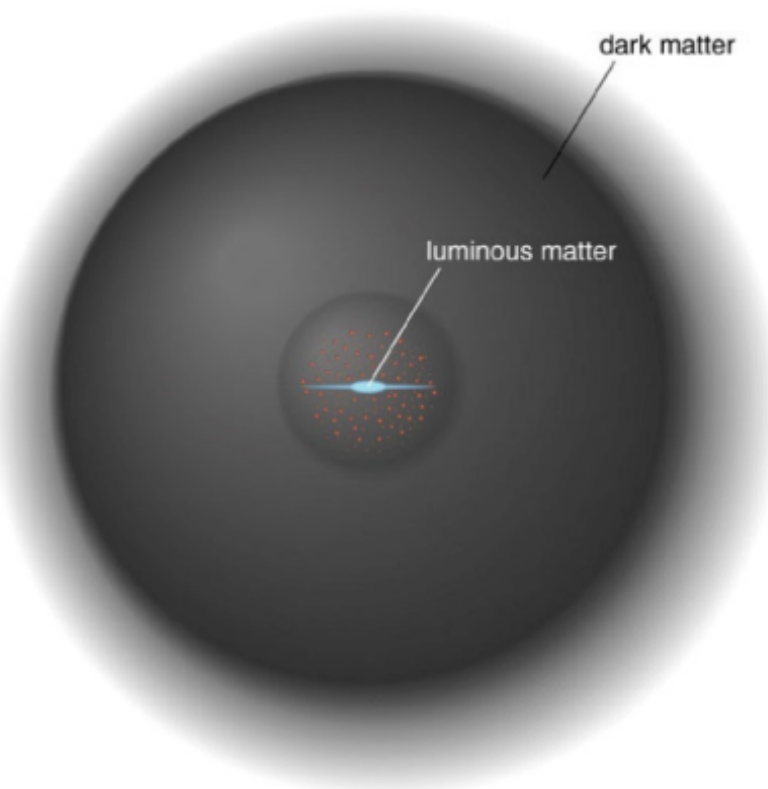


$$\frac{dR}{dE}(E, t) = \frac{\rho_0}{m_\chi \cdot m_A} \cdot \int \mathbf{v} \cdot \mathbf{f}(\mathbf{v}, t) \cdot \frac{d\sigma}{dE}(E, \mathbf{v}) d^3v$$

Sub-GeV dark matter?
pseudoscalar mediators
Non-standard Cosmologies

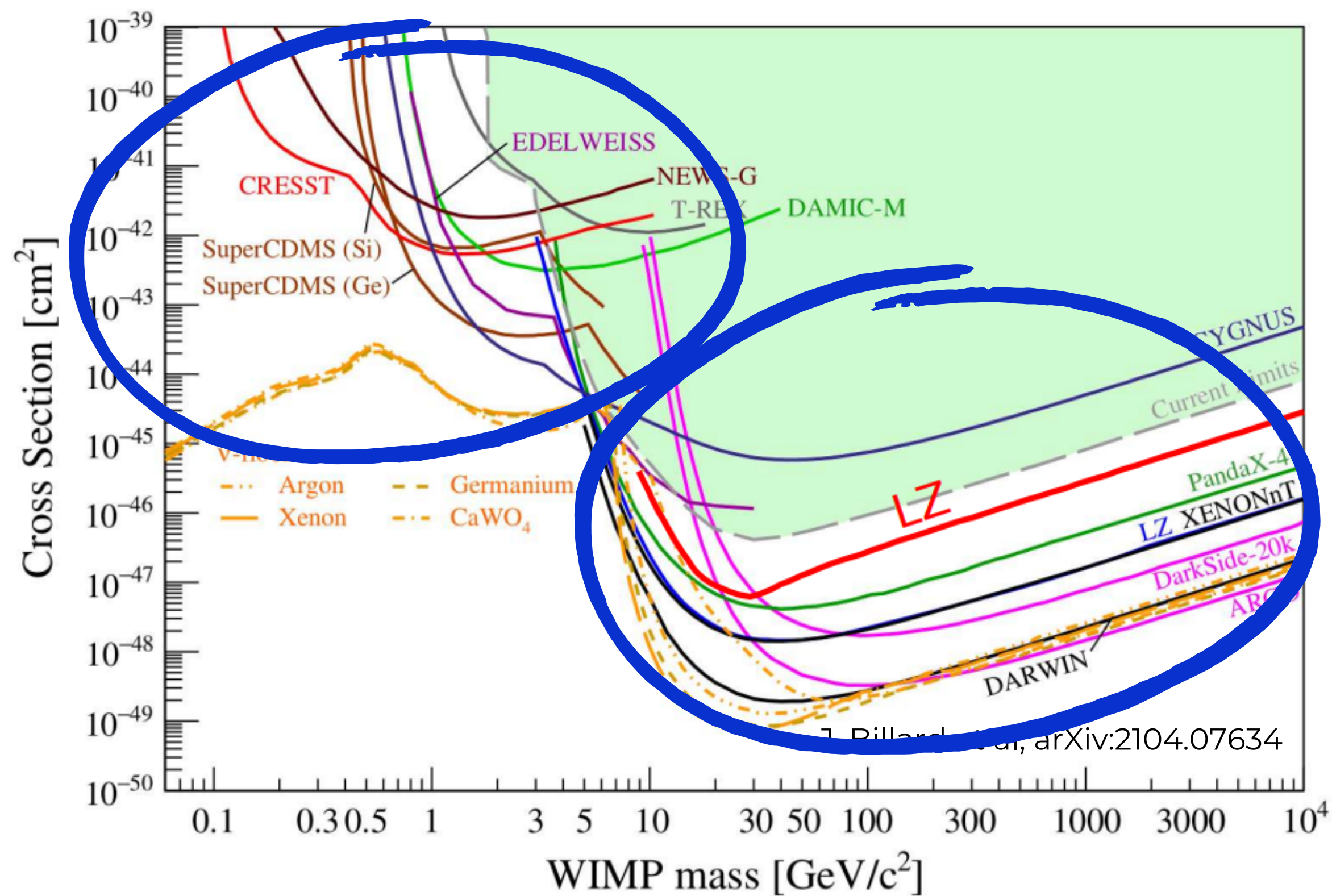


J. Billard et al, arXiv:2104.07634



$$\frac{dR}{dE}(E, t) = \frac{\rho_0}{m_\chi \cdot m_A} \cdot \int \mathbf{v} \cdot \mathbf{f}(\mathbf{v}, t) \cdot \frac{d\sigma}{dE}(E, \mathbf{v}) d^3v$$

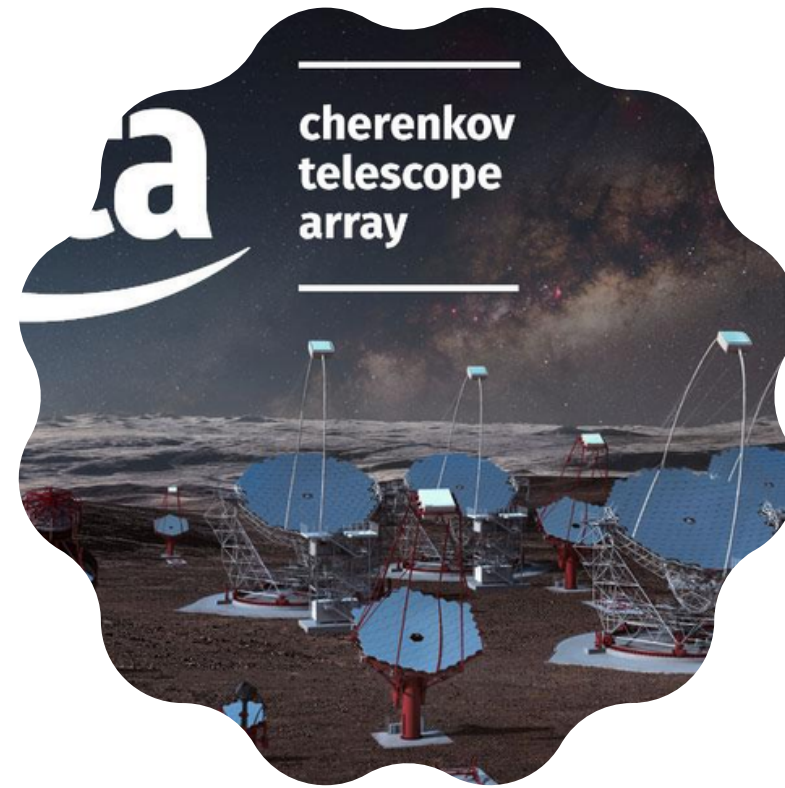
Sub-GeV dark matter?
pseudoscalar mediators
Non-standard Cosmologies



Indirect Detection



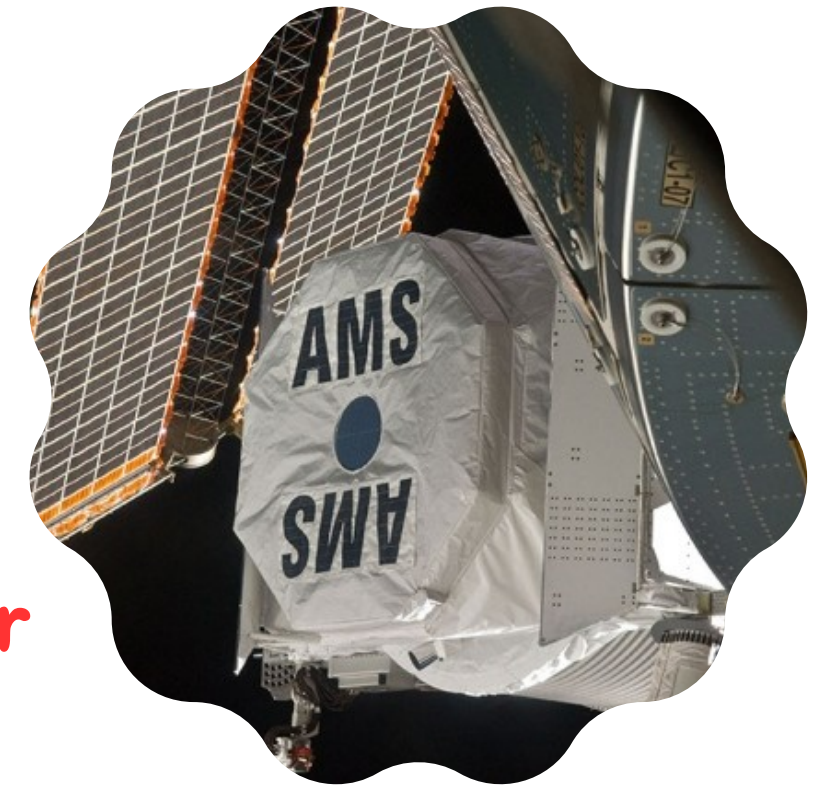
*Fermi-LAT



CTA

It is related
to the dark matter
local density.

I don't want to
run from it



AMS

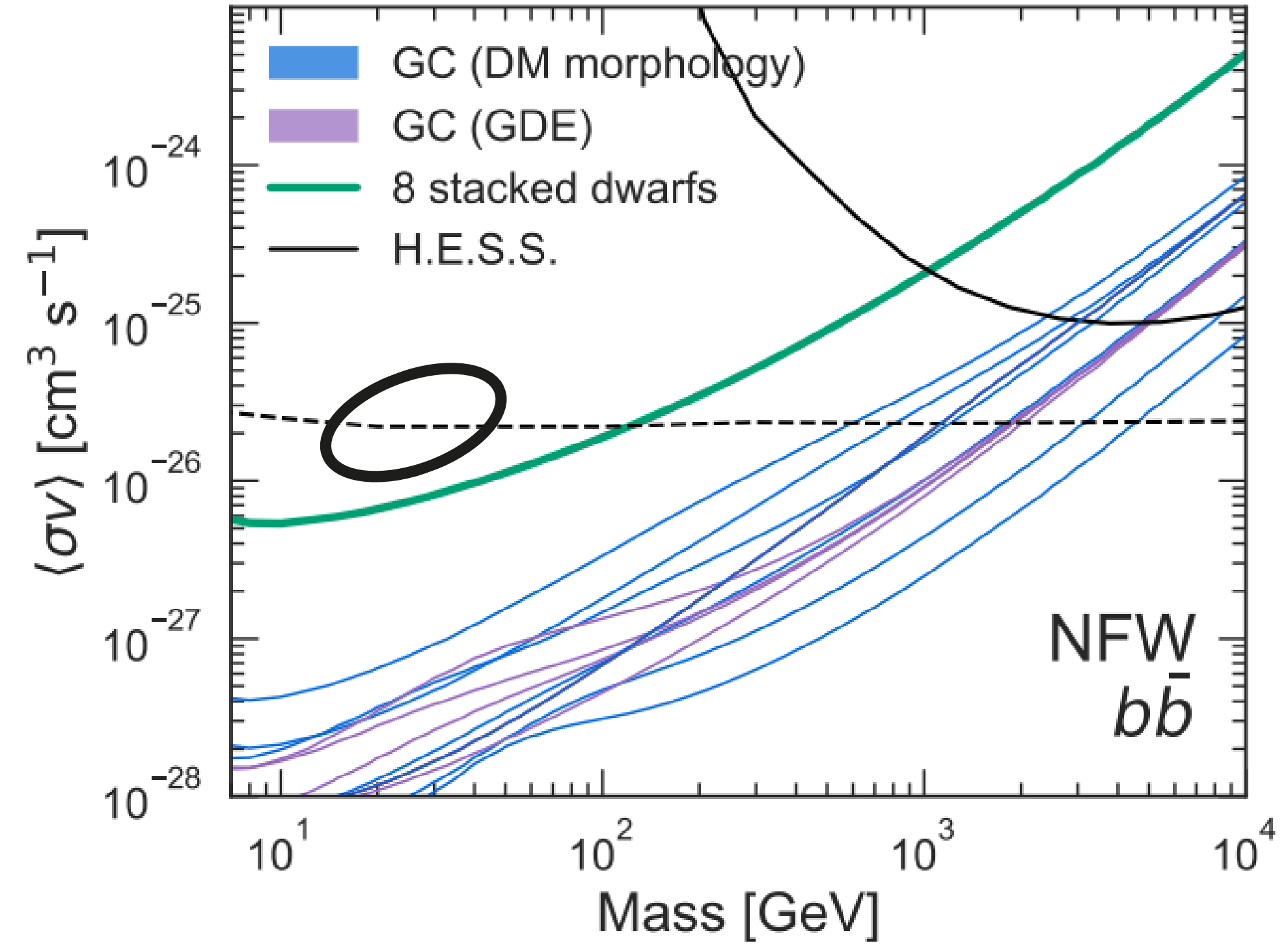
*Neutrino telescopes

$$\underbrace{\frac{d\Phi}{d\Omega dE}}_{\text{Diff. Flux}} = \frac{\underbrace{\sigma v}_{\text{Anni. Cross Section}}}{8\pi m_\chi^2} \times \underbrace{\frac{dN}{dE}}_{\text{Energy Spectrum}} \times \int_{\text{l.o.s}} ds \underbrace{\rho^2(\vec{r}(s, \Omega))}_{\text{Dark Matter Distribution}}$$

Where are we?

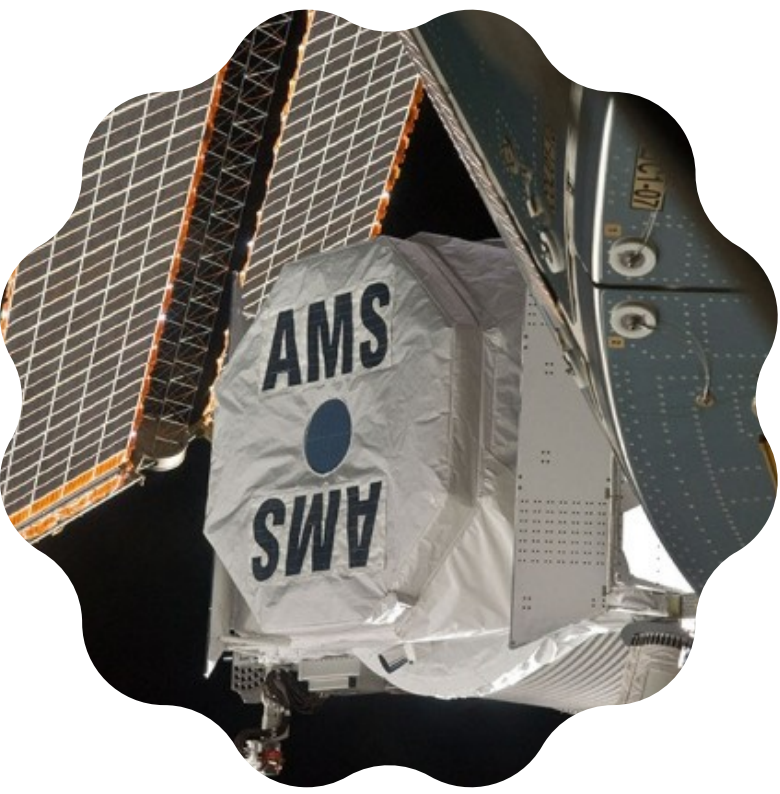


***Fermi-LAT**
through 2022,
possibly will
continue



$$\underbrace{\frac{d\Phi}{d\Omega dE}}_{\text{Diff. Flux}} = \frac{\underbrace{\sigma v}_{\text{Anni. Cross Section}}}{8\pi m_\chi^2} \times \underbrace{\frac{dN}{dE}}_{\text{Energy Spectrum}} \times \int_{\text{l.o.s}} ds \underbrace{\rho^2(\vec{r}(s, \Omega))}_{\text{Dark Matter Distribution}}$$

Where are we?

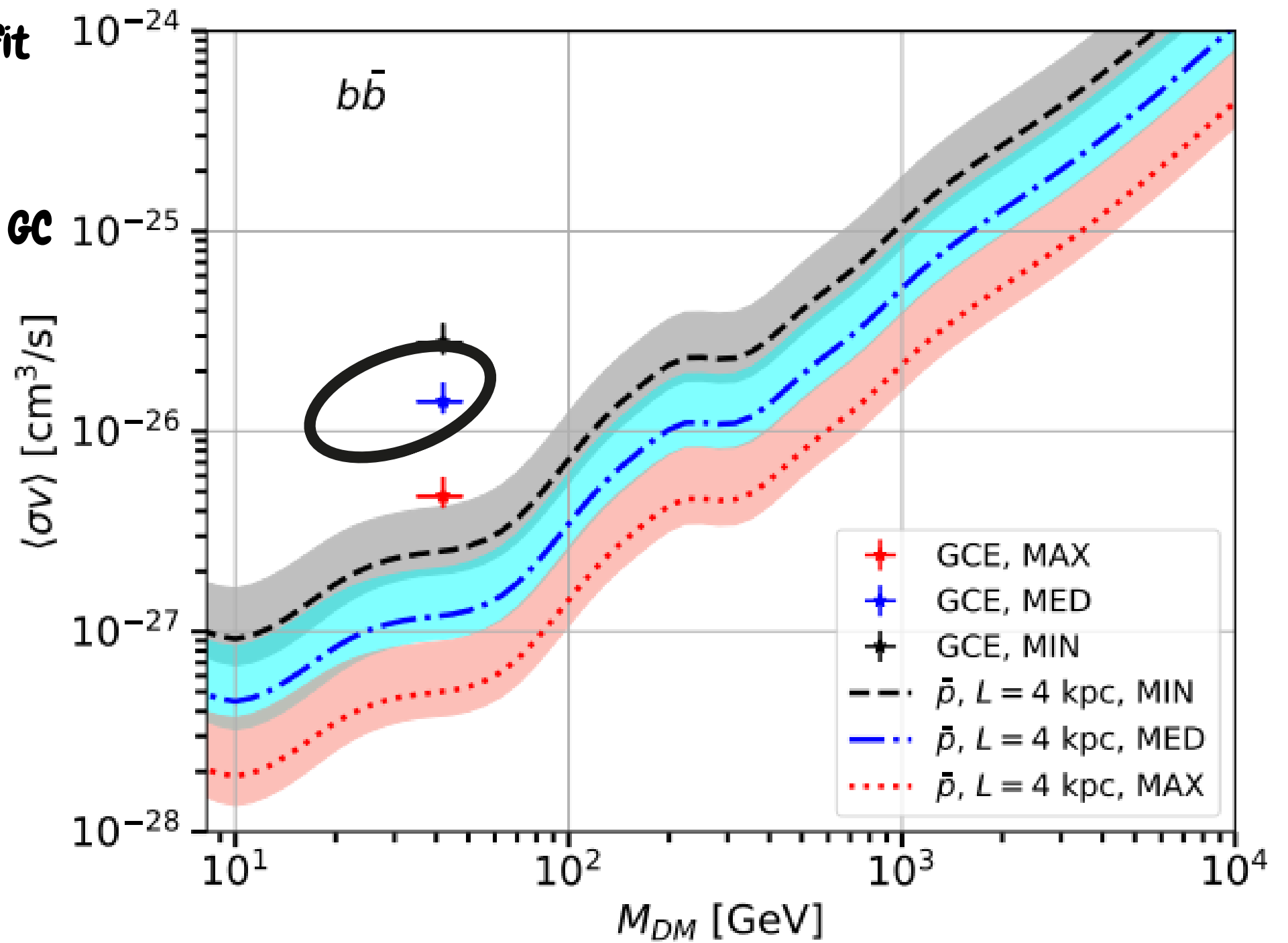


AMS

Our plans are to keep AMS operating on the International Space Station as long as there is a space station.

It is fair to say that the DM interpretation does provide a good fit to the Fermi-LAT data but is disfavored by other probes. The neutron stars interpretation for the GC excess is debatable

Leane and Slatyer, arXiv:1904.08430



Mauro, Winkler arxiv: 2101.11027

$$\underbrace{\frac{d\Phi}{d\Omega dE}}_{\text{Diff. Flux}} = \frac{\underbrace{\sigma v}_{\text{Anni. Cross Section}}}{8\pi m_\chi^2} \times \underbrace{\frac{dN}{dE}}_{\text{Energy Spectrum}} \times \int_{l.o.s} ds \underbrace{\rho^2(\vec{r}(s, \Omega))}_{\text{Dark Matter Distribution}}$$

Where are we?

Forecast

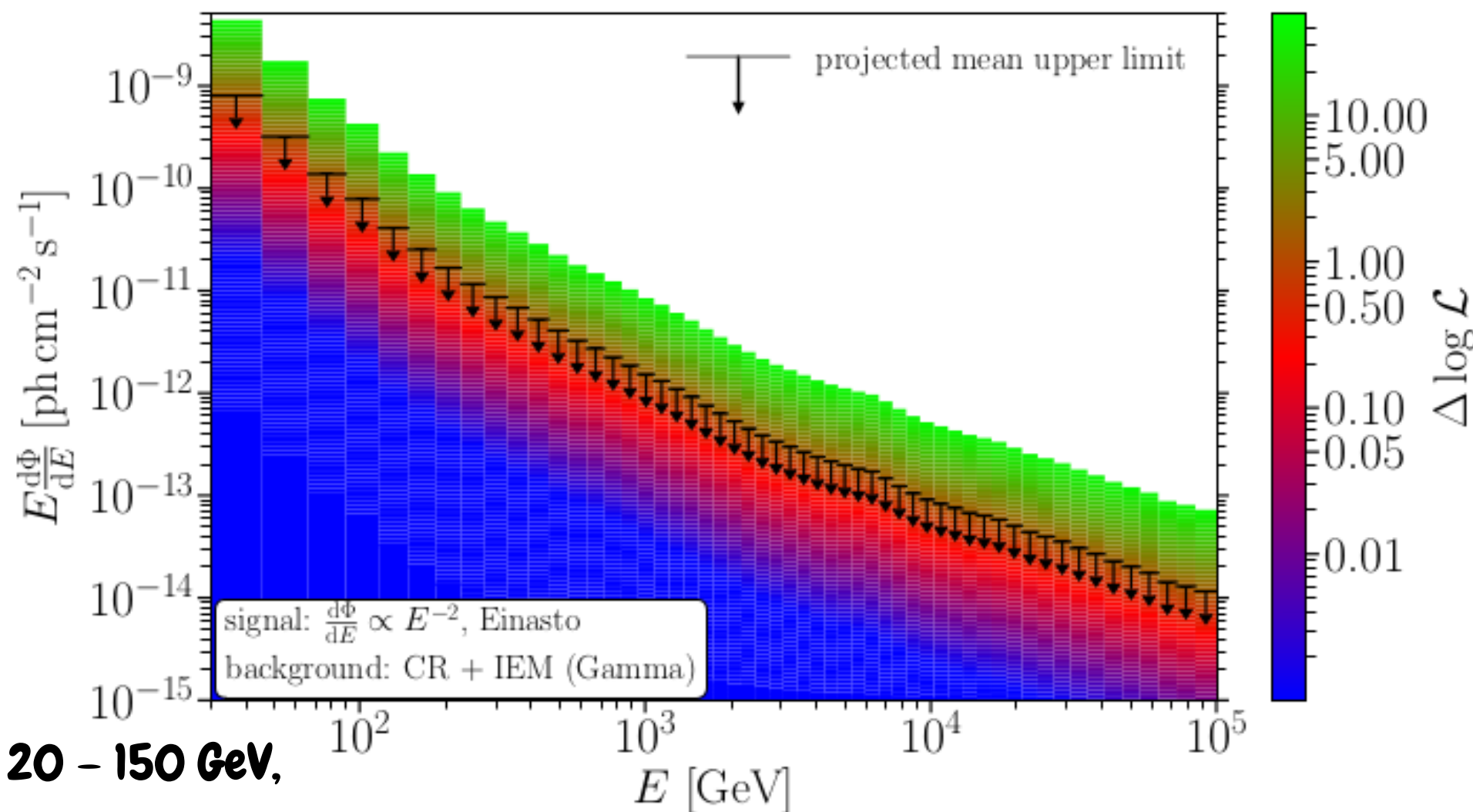


CTA

Hopefully in the coming years

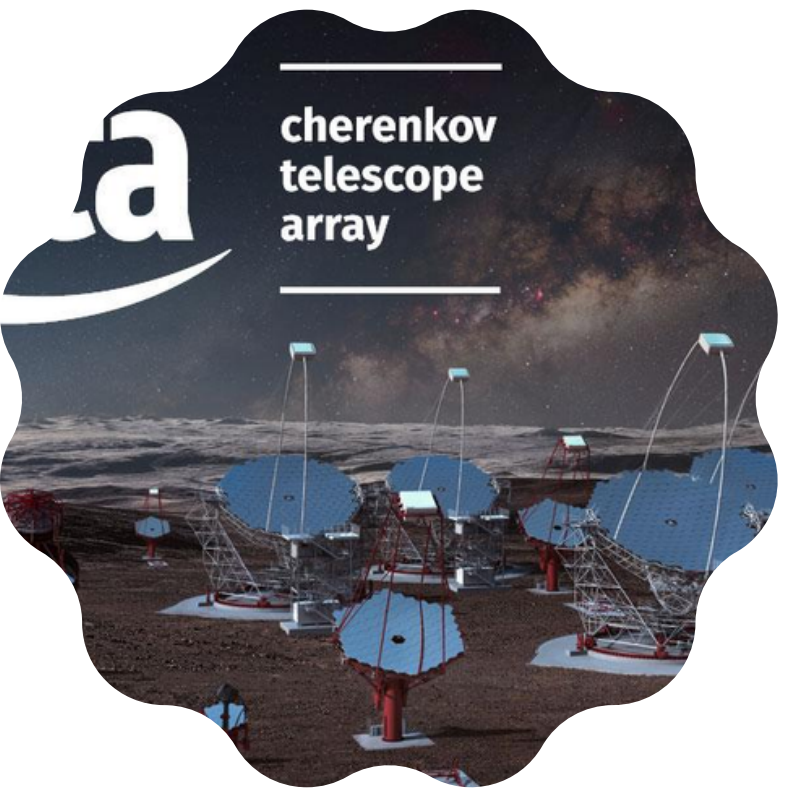
The CTA design concept :

- i) LSTs (Large-Sized Telescopes, 23 m in diameter) E= 20 – 150 GeV,**
- ii) MSTs (Medium-Sized Telescopes, 11.5 m) E=150 GeV–5 TeV**
- iii) a large number of SSTs (Small-Sized Telescope, 4 m) E> 5TeV**



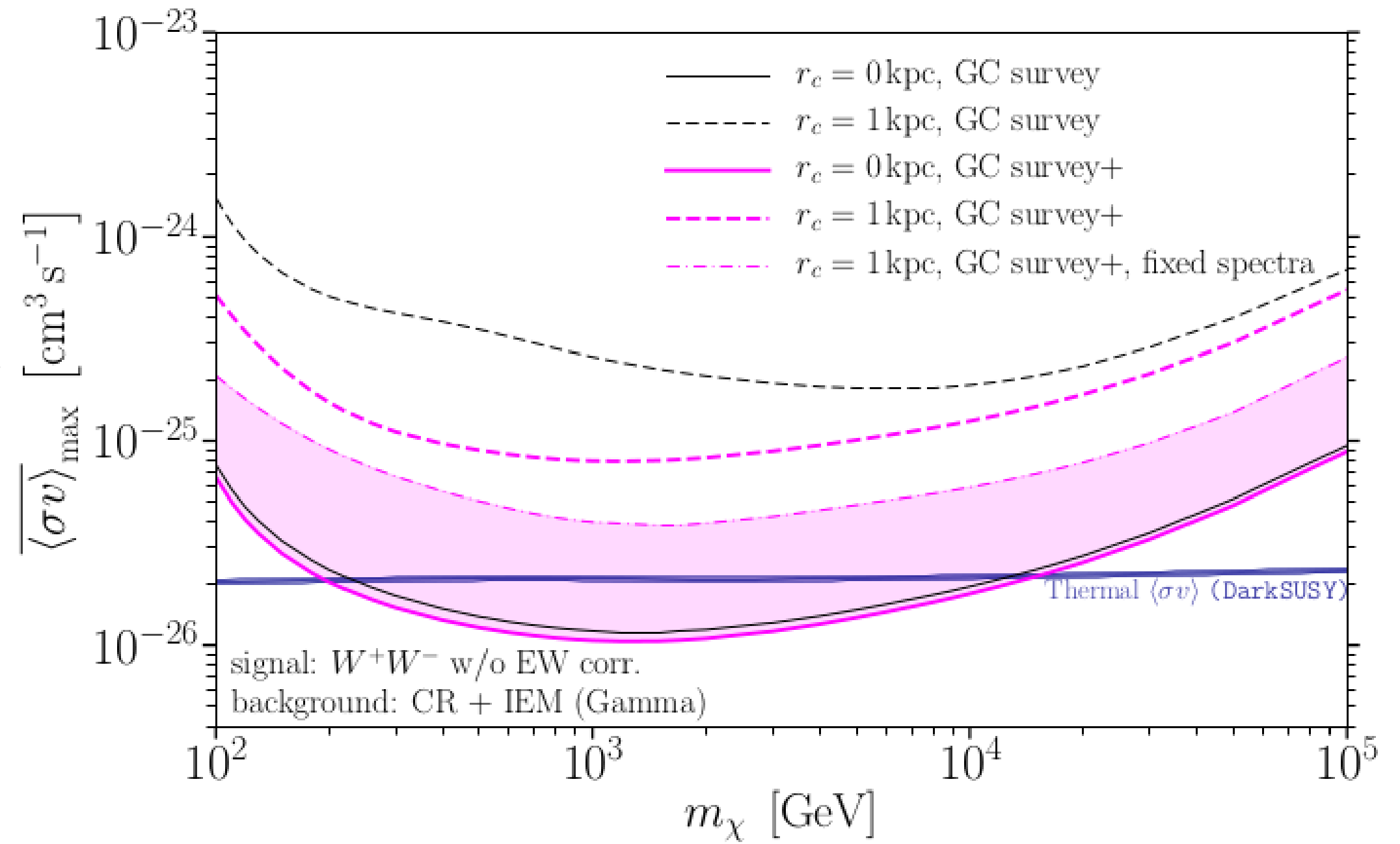
$$\underbrace{\frac{d\Phi}{d\Omega dE}}_{\text{Diff. Flux}} = \frac{\underbrace{\sigma v}_{\text{Anni. Cross Section}}}{8\pi m_\chi^2} \times \underbrace{\frac{dN}{dE}}_{\text{Energy Spectrum}} \times \int_{\text{l.o.s}} ds \underbrace{\rho^2(\vec{r}(s, \Omega))}_{\text{Dark Matter Distribution}}$$

Where are we?



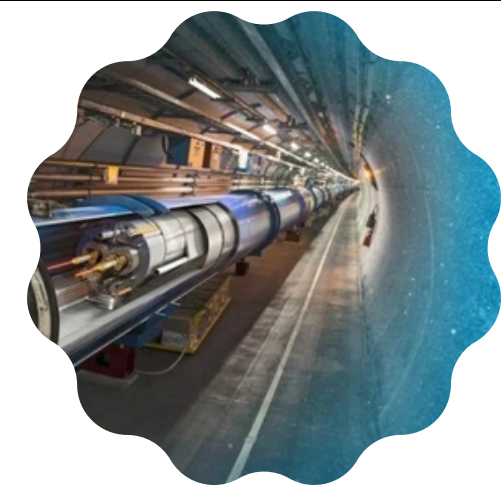
CTA
Hopefully in the coming years

CTA is expected to improve by 1-2 orders of magnitude the limits (masses > 300GeV)



$$\underbrace{\frac{d\Phi}{d\Omega dE}}_{\text{Diff. Flux}} = \frac{\underbrace{\sigma v}_{\text{Anni. Cross Section}}}{8\pi m_\chi^2} \times \underbrace{\frac{dN}{dE}}_{\text{Energy Spectrum}} \times \int_{\text{l.o.s}} ds \underbrace{\rho^2(\vec{r}(s, \Omega))}_{\text{Dark Matter Distribution}}$$

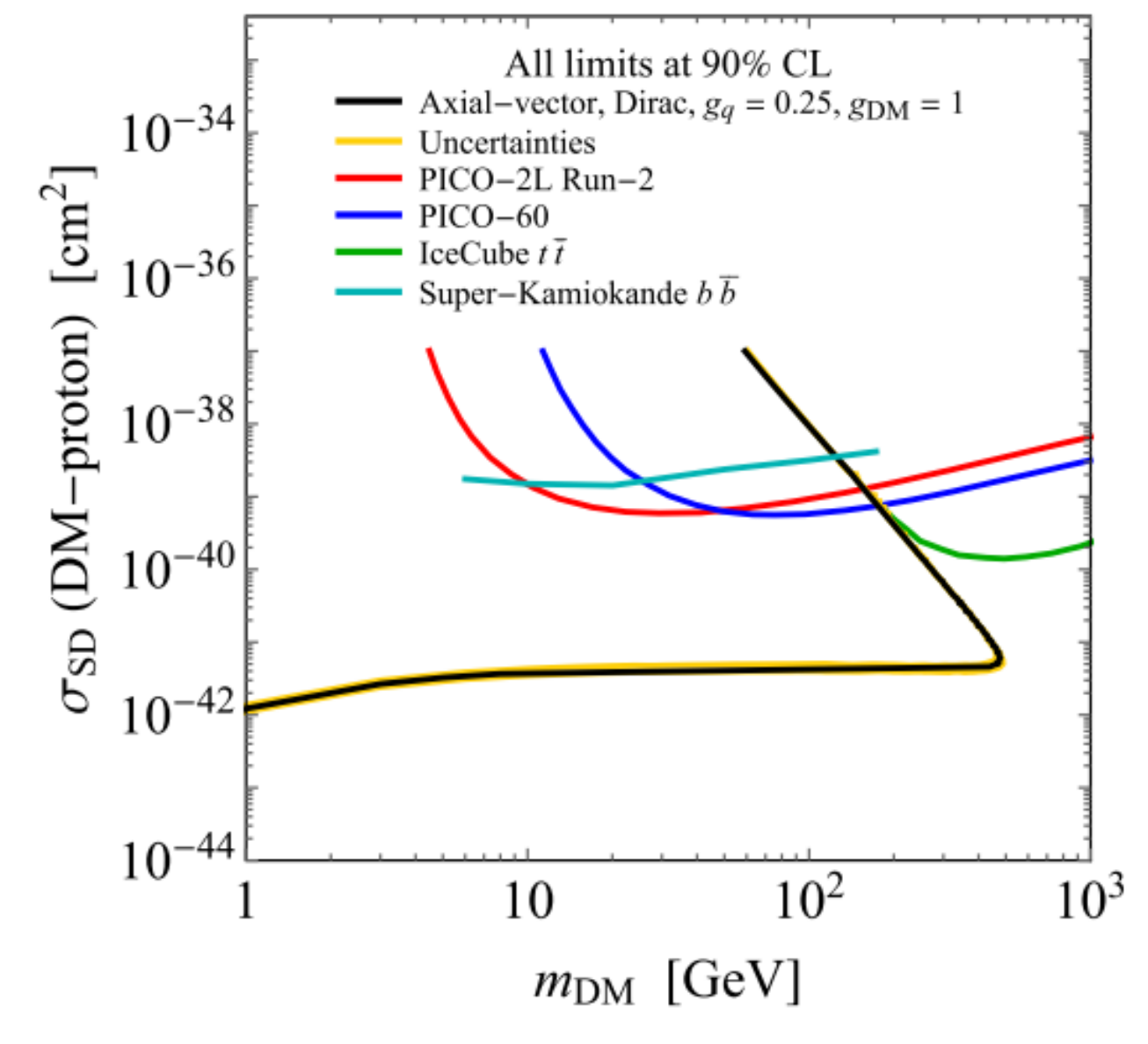
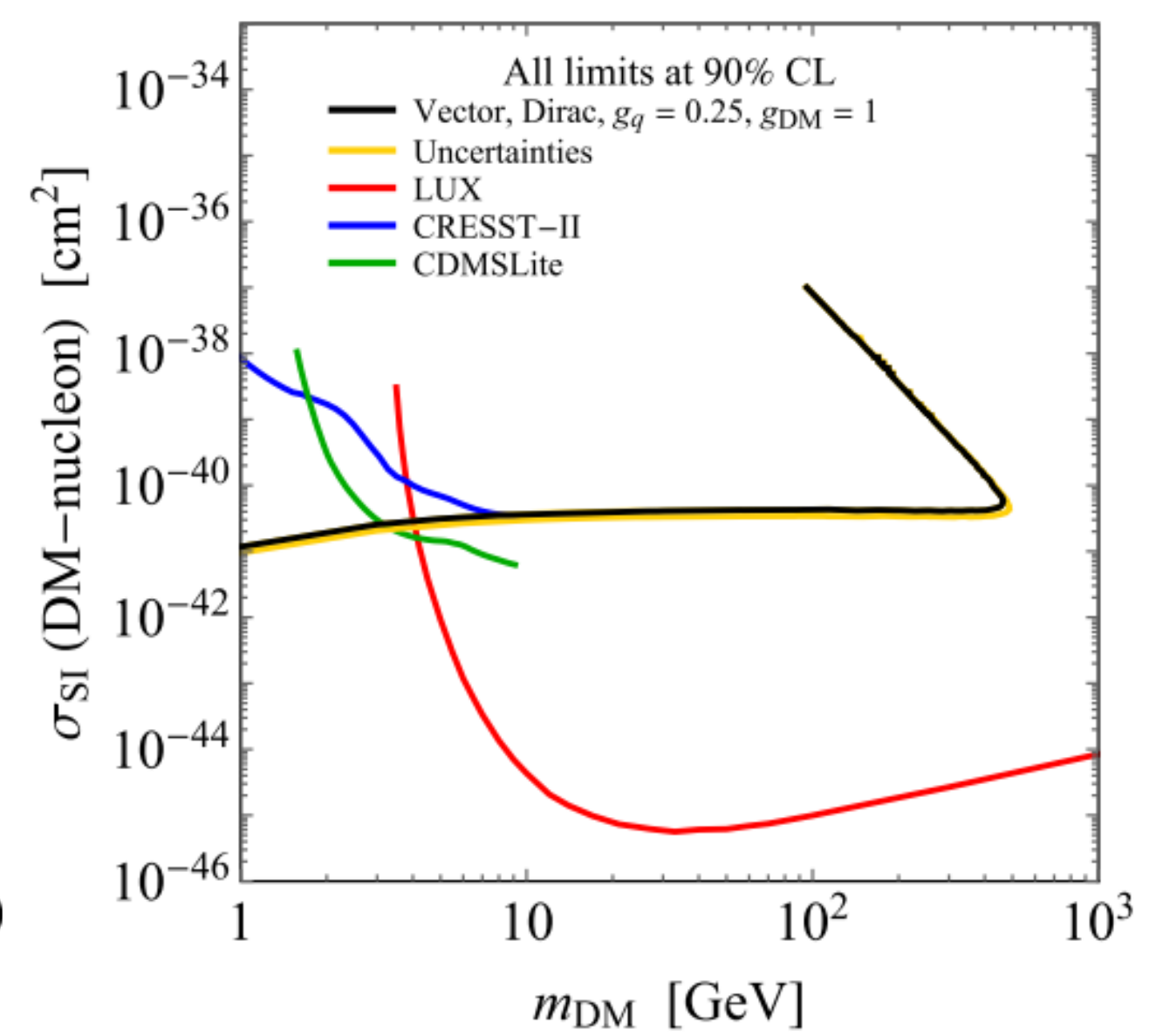
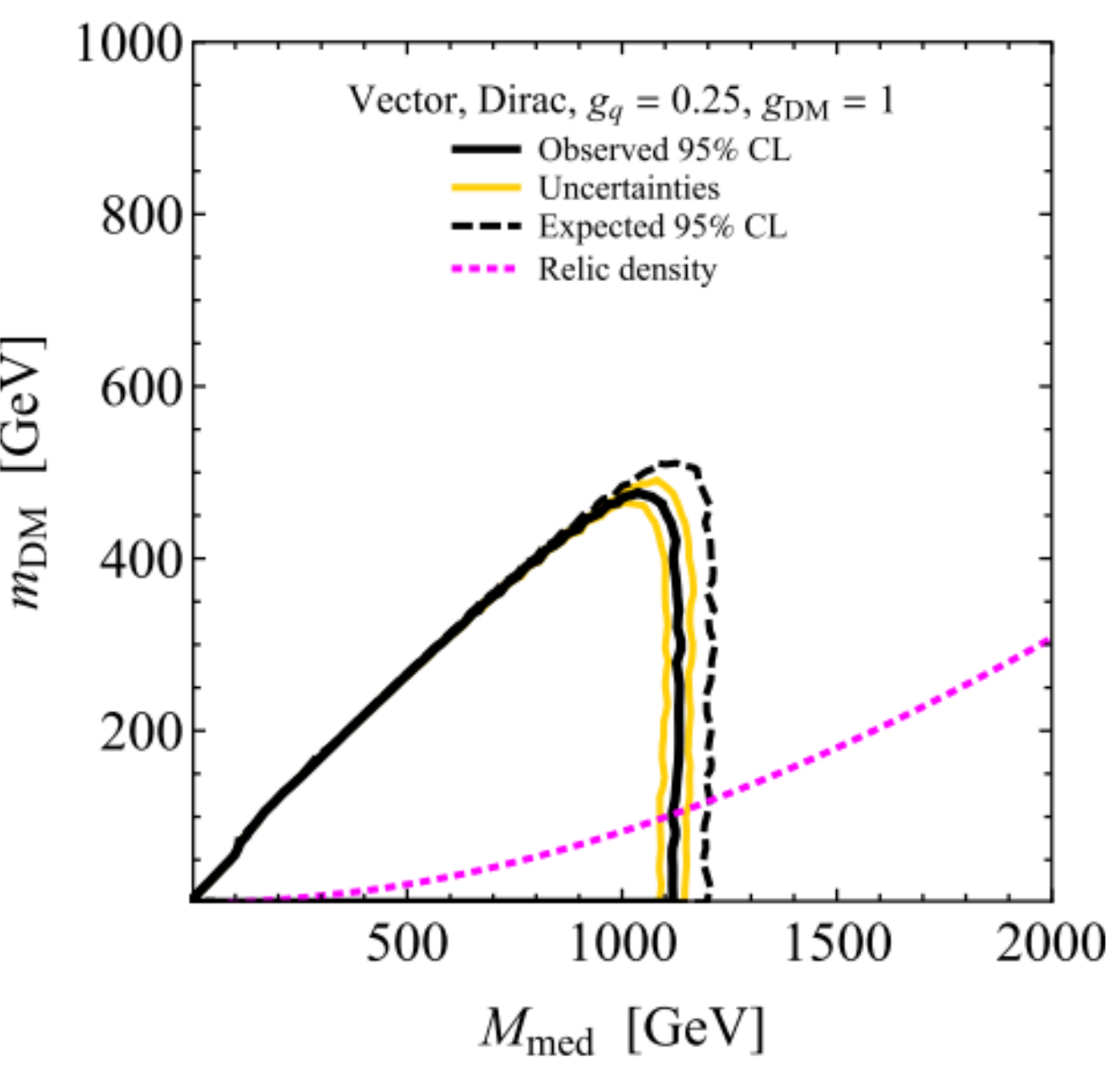
Colliders



Represent an important cross-check

Cover regions of parameter unexplored by direct (indirect) detection experiments

A. Boveia, 2022 Snowmass Summer Study, 2206.03456

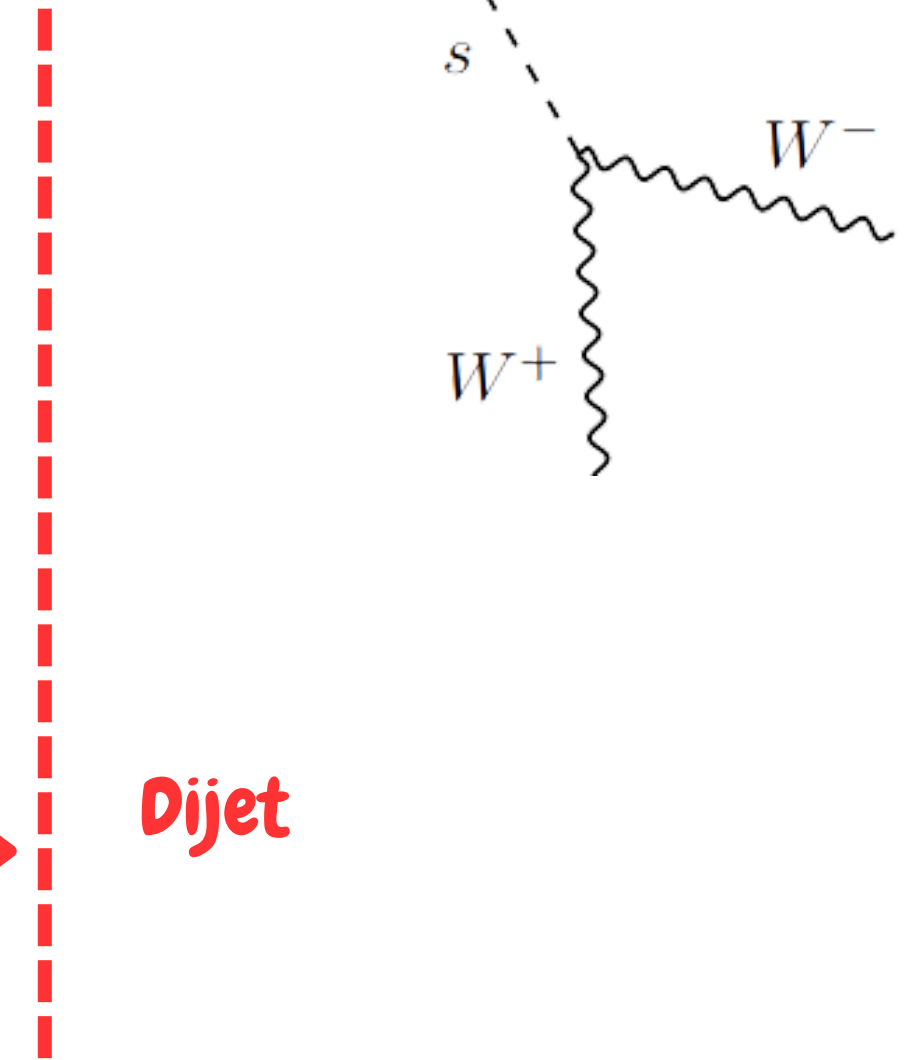
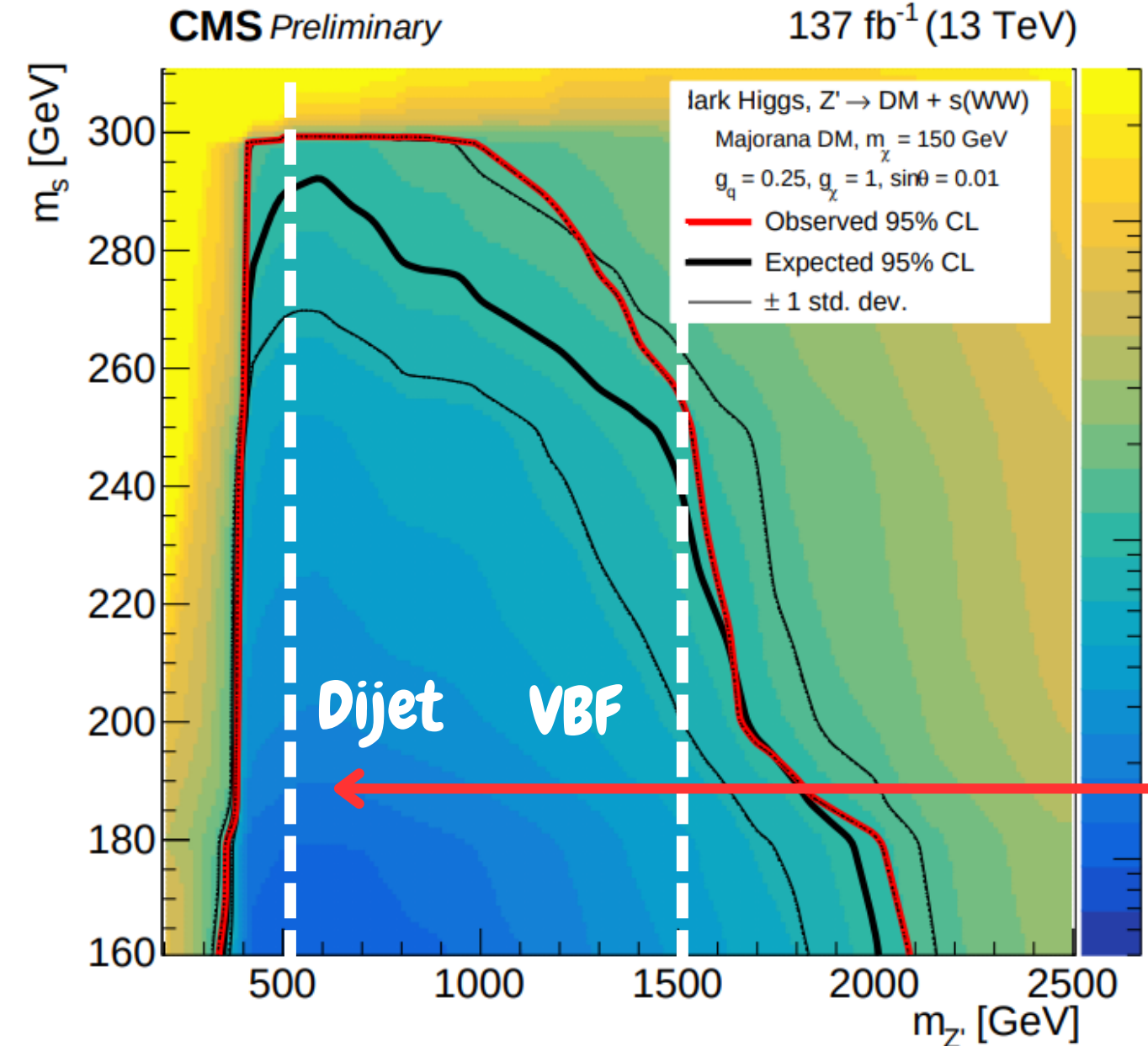
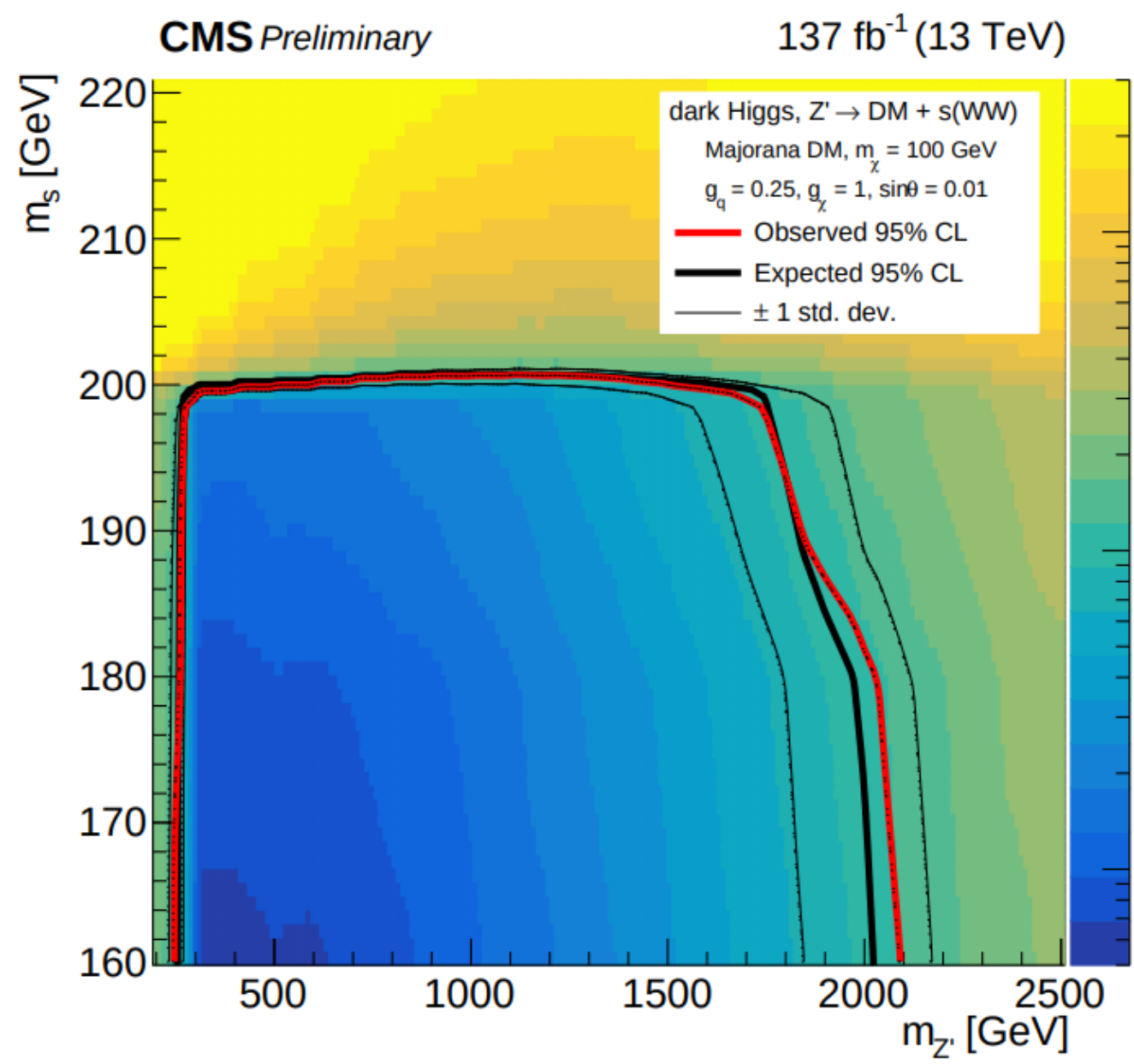
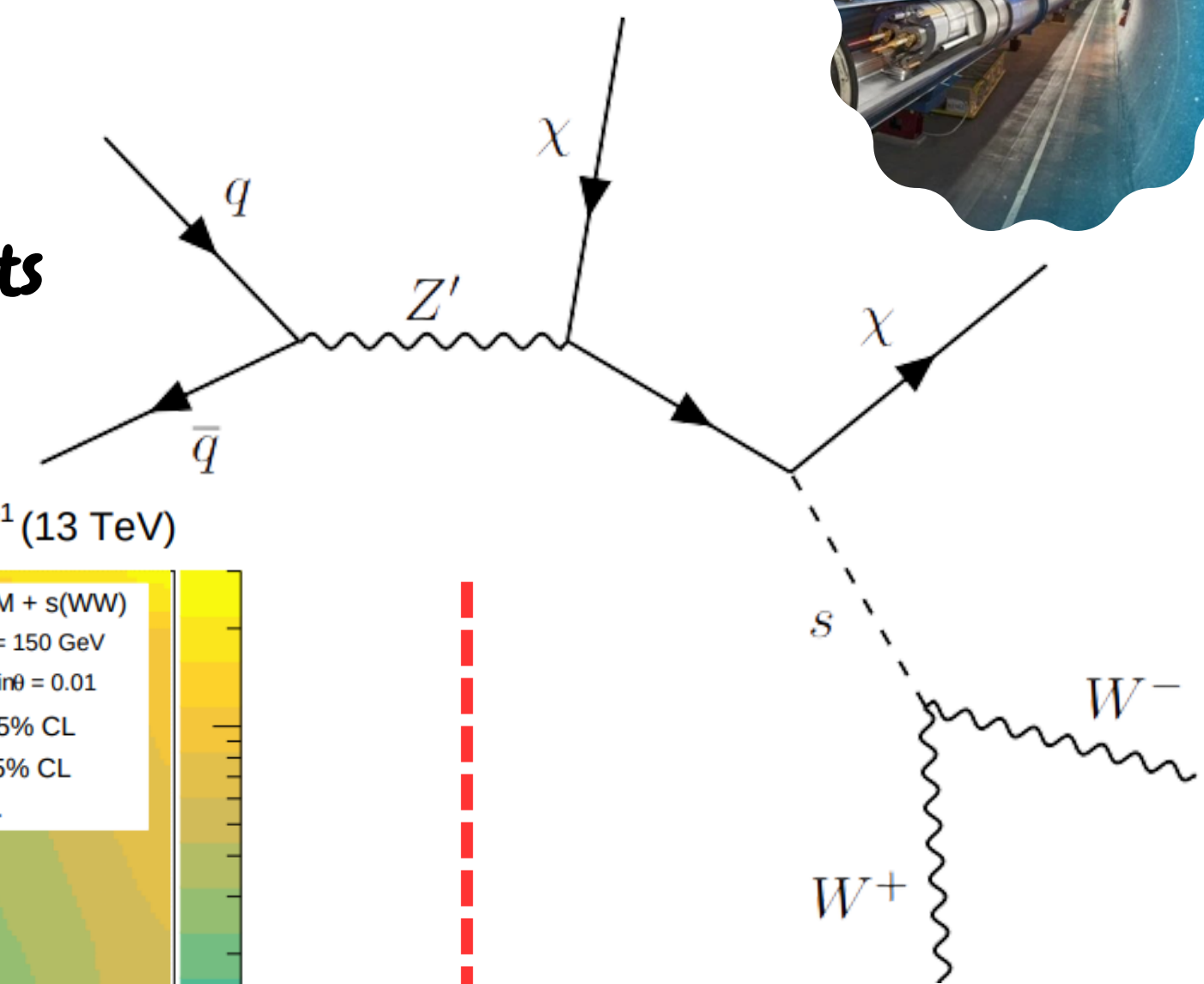
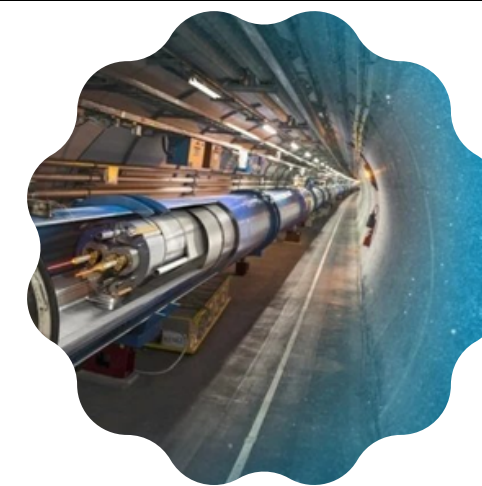


Colliders

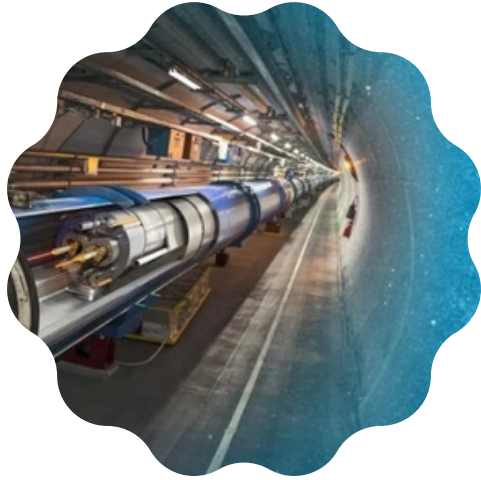
Represent an important cross-check

Cover regions of parameter unexplored by direct detection experiments

CMS, arxiv:1908.01713



Colliders



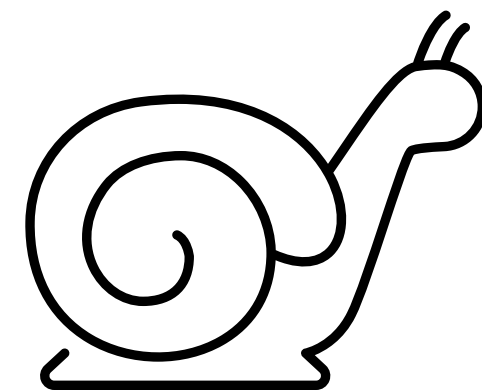
What's the future?

HL-LHC in 2030: it will not improve by one order of magnitude the sensitivity on dark sector

HE-LHC in 2050 ?

Colliders will continue to be important searches for dark sectors but the progress will occur at a very different pace

Muon collider? LHeC?



CoGENT
CDMS anomaly
DAMA
The GeV Fermi excess
AMS02 data
DAMPE
PAMELA
XenonIT anomaly
x-ray signals.....
Shunsaku Horiuchi's tak



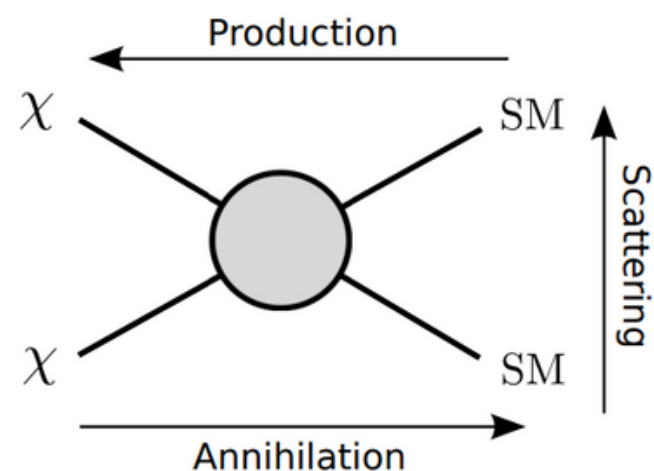
Pandora box

**If it is hard to observe a clear dark matter signal
within the thermal relic paradigm, imagine if we
depart from it**

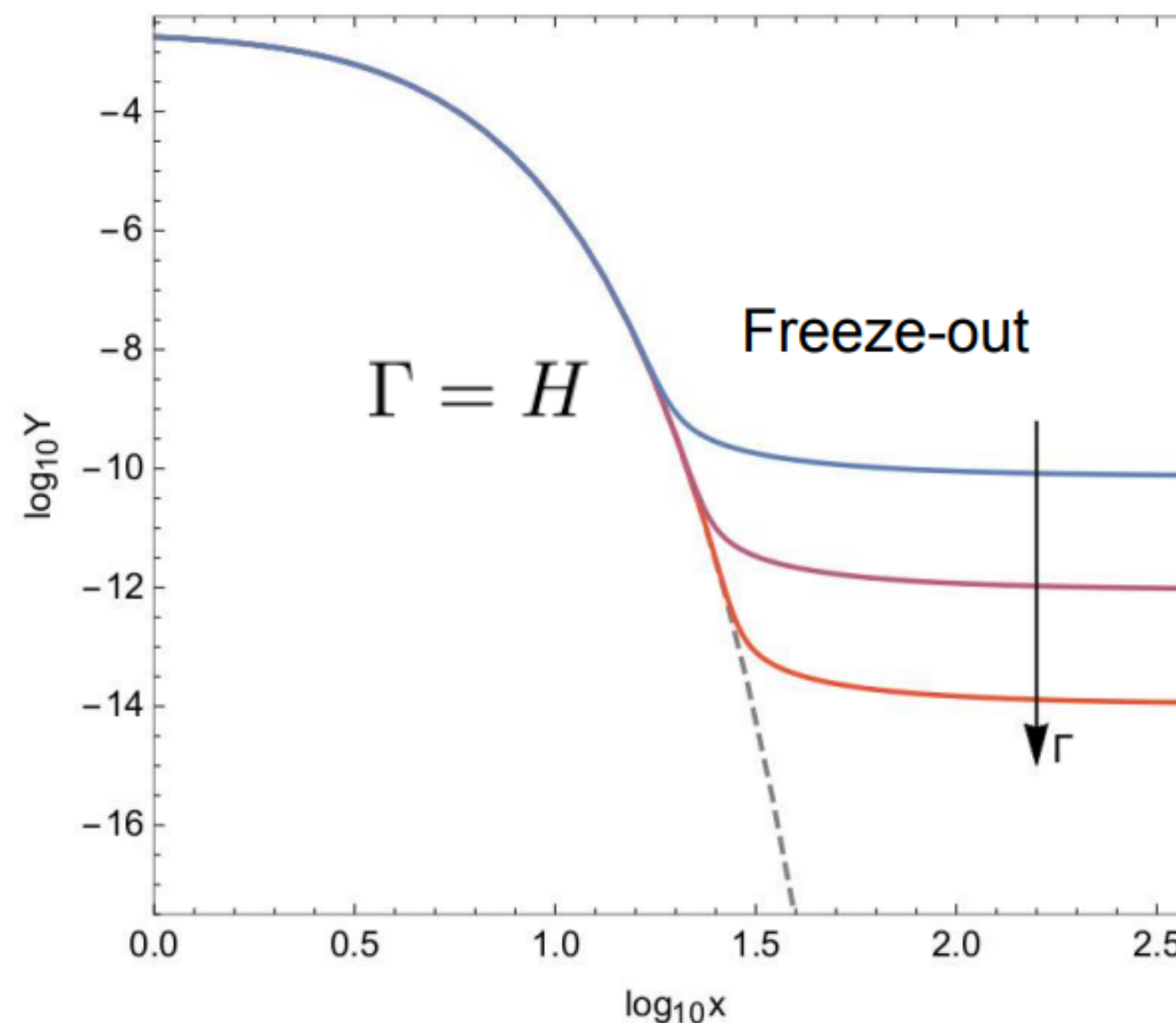
Standard Darkness

We often use as evidence for dark matter

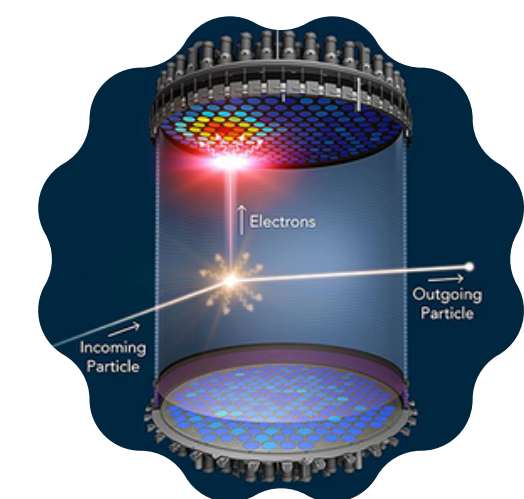
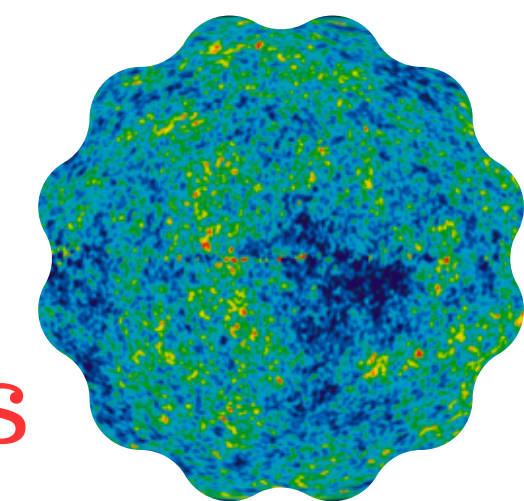
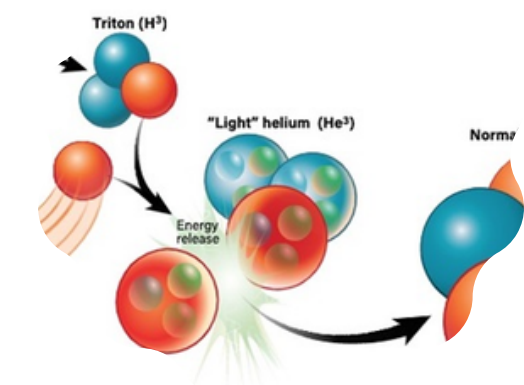
- Galaxy rotation curves
- Cosmic Microwave Background
- Collision of Clusters
- Baryon Acoustic Oscillations
- Gravitational Lensing
- Cosmic Shear
- Structure Formation
- *Fermi GeV excess
- *AMS-02 results



Thermal relic



smaller CS
standard CS
large CS



$$\frac{dn_\chi}{dt} + 3Hn_\chi = -\langle\sigma v\rangle (n_\chi^2 - n_\chi^{eq2})$$

Standard Darkness

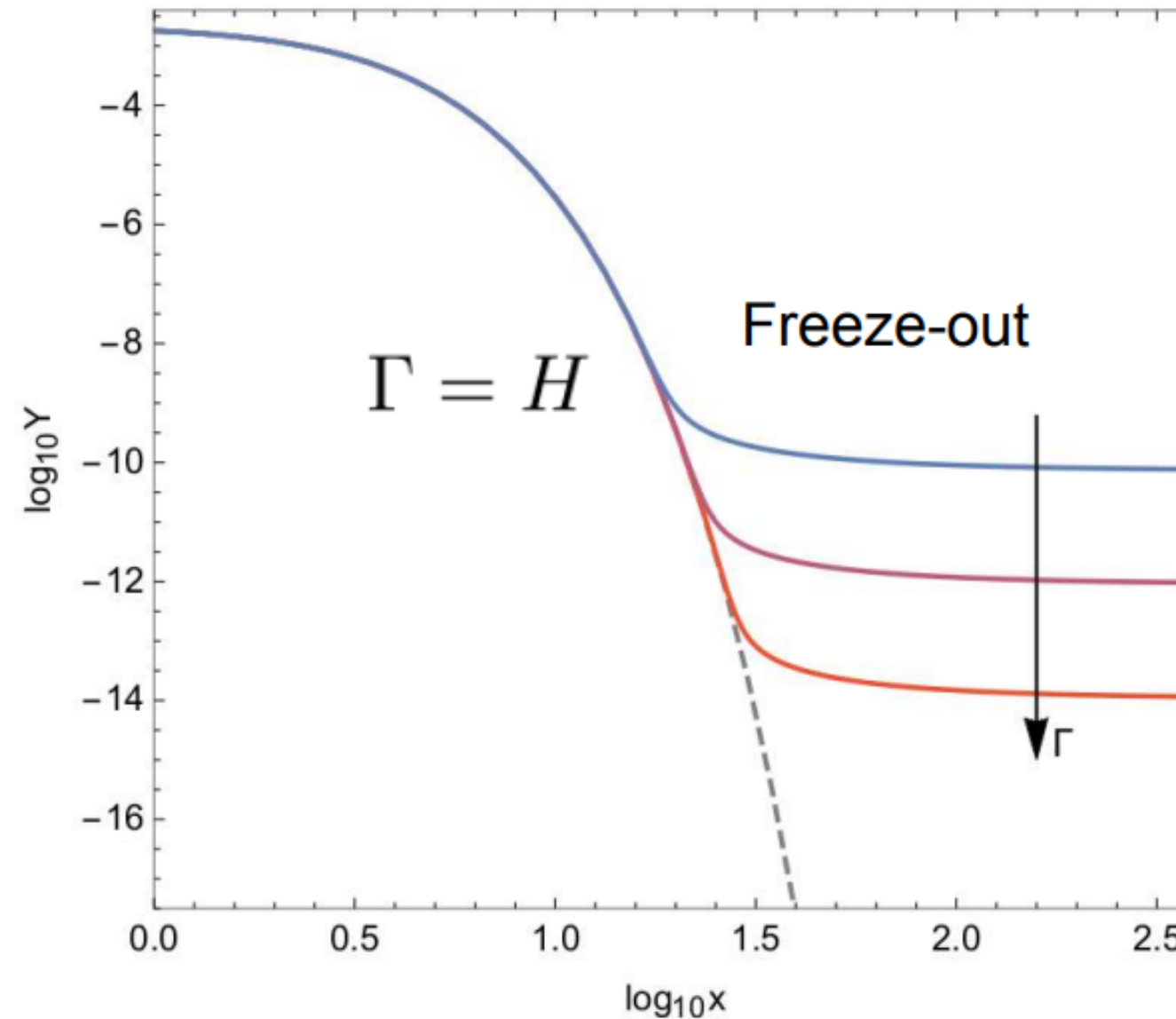
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- Galaxy rotation curves
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- Structure Formation
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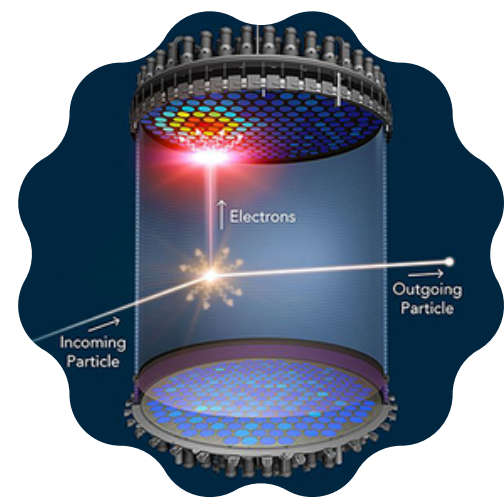
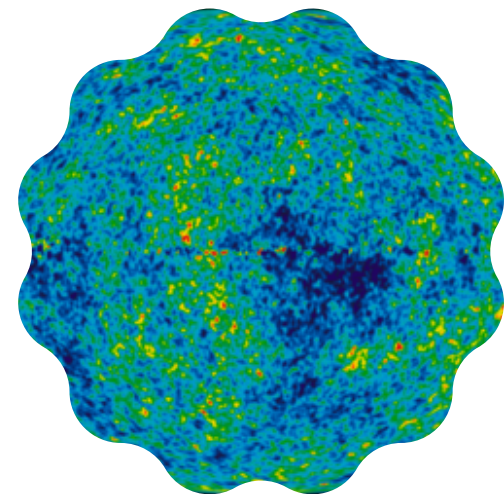
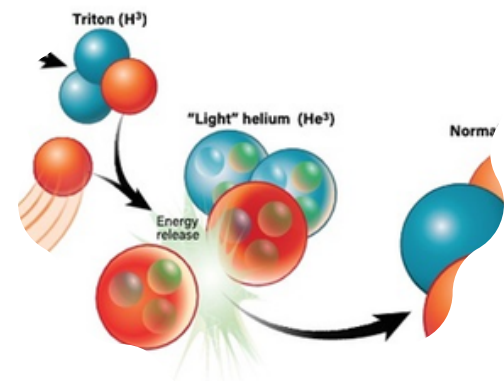
- Neutrino Masses
- Lepton Flavor Violation
- Hierarchy Problem
- Grand Unification

Extra

Thermal relic



$$\frac{dn_\chi}{dt} + 3Hn_\chi = -\langle\sigma v\rangle (n_\chi^2 - n_\chi^{eq2})$$



Bear in mind that I will play the devil's advocate

Standard Darkness

Stranger Darkness

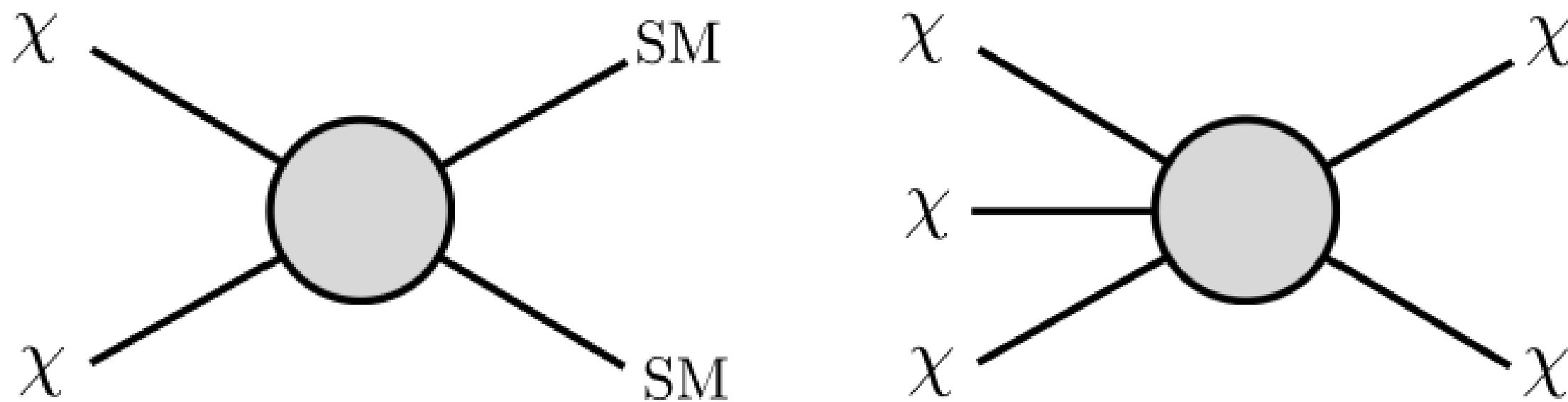
Stranger Darkness



Before we were motivated by "signals" now by why not?

SIMP (Strongly Interacting Massive Particle)

DM abundance is determined by number changing interactions



Large self-interactions are required for observed abundance

Takashi's talk
on Tuesday

Helps to solve the small
scale problem (cusp x core)

How can I prove to my
experimental colleague
that a SIMP was observed?

Stranger Darkness

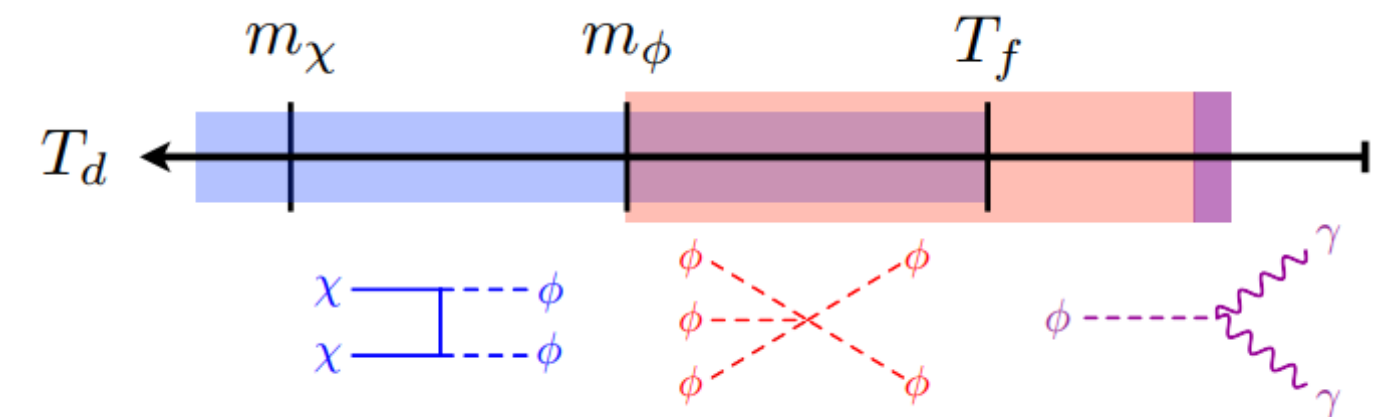


Before we were motivated by "signals" now by why not?

Cannibal DM

1. The dark sector is kinetically decoupled from the SM sector.
2. The dark sector has a mass gap.
3. The dark sector remains in chemical equilibrium, through number changing interactions, at temperatures below the mass of the LDP.
- 4*. The scalar field ϕ decays into radiation
- 5* The scalar field dominated the energy density
- 6* Assume that DM annihilations decouple during cannibalism

Dark fields are kinematically decoupled from SM
The dark sector remains in chemical equilibrium,
through number changing interactions,



A boosted rate for
indirect detection

How can I prove to my
experimental colleague
that a Cannibal DM was observed?

Stranger Darkness

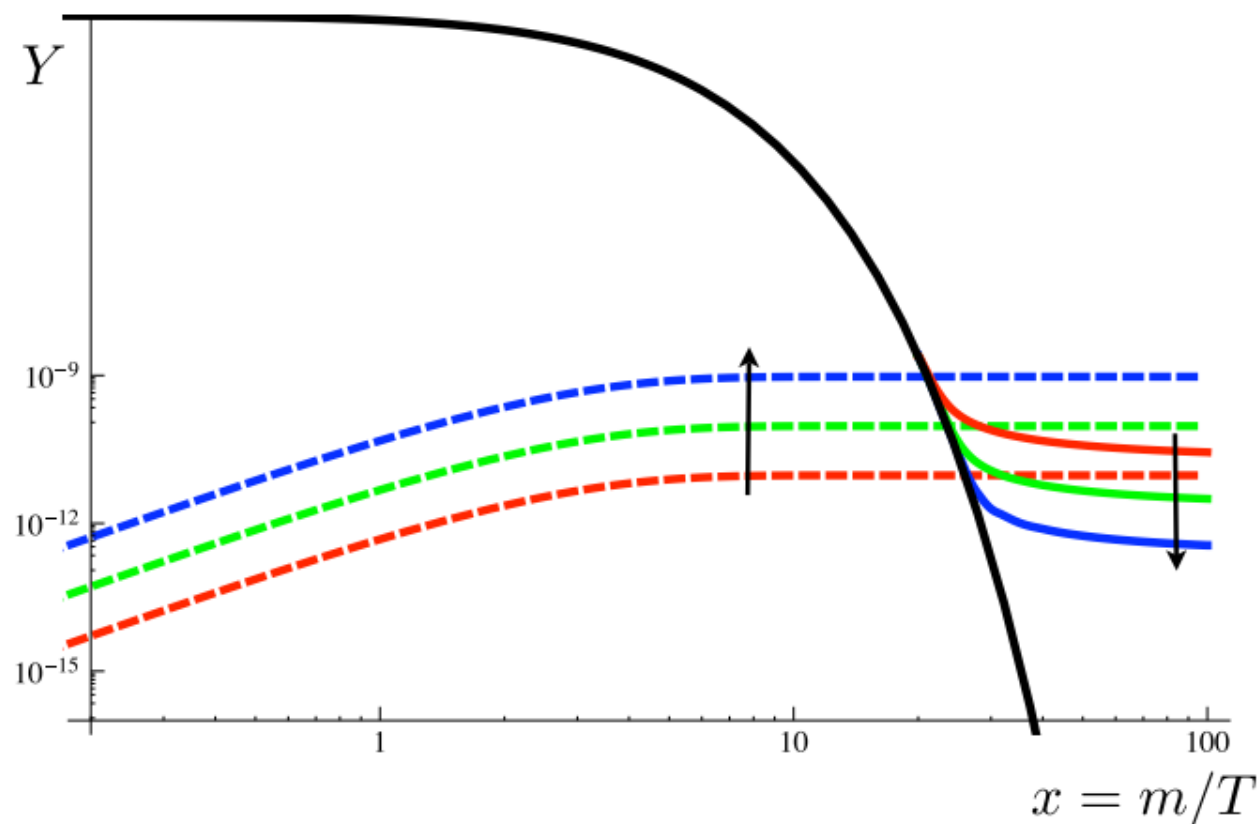


Before we were motivated by "signals" now by why not?

Hope to observe
the companions

FIMP (Feebly Interacting Massive Particle)

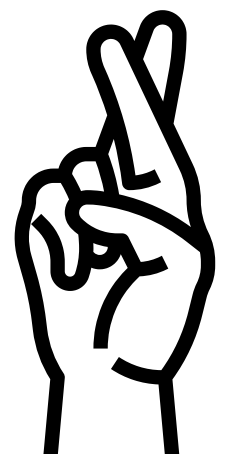
1. Never thermalized with SM sector,
2. Produced slowly by decays or scatterings
3. tiny couplings are needed

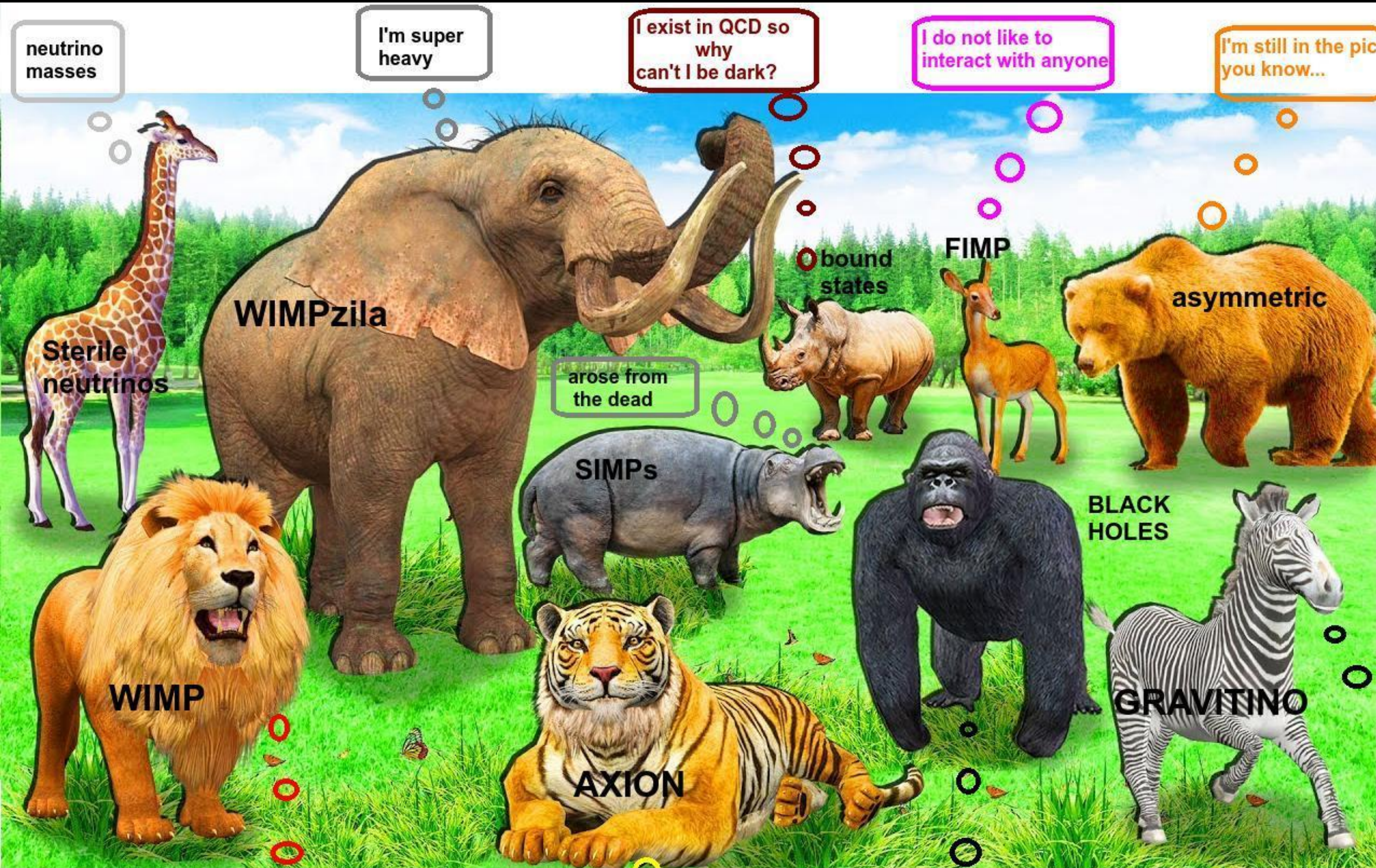


$$\Omega_\chi h^2 \simeq 4.48 \times 10^8 \frac{g_\sigma}{g_{*s} \sqrt{g_*}} \frac{m_\chi}{\text{GeV}} \frac{M_{\text{P}} \Gamma_{\sigma \rightarrow \chi\chi}}{m_\sigma^2}$$

$$y \simeq 10^{-12} \left(\frac{\Omega_\chi h^2}{0.12} \right)^{1/2} \left(\frac{g_*}{100} \right)^{3/4} \left(\frac{m_\sigma}{m_\chi} \right)^{1/2}$$

Good luck proving
it has anything to do
with dark matter





**DARK ZOO is not enough.
We created a genetically
modified dark zoo**

I am still the king

If he dies it's my turn

Well..we never know..

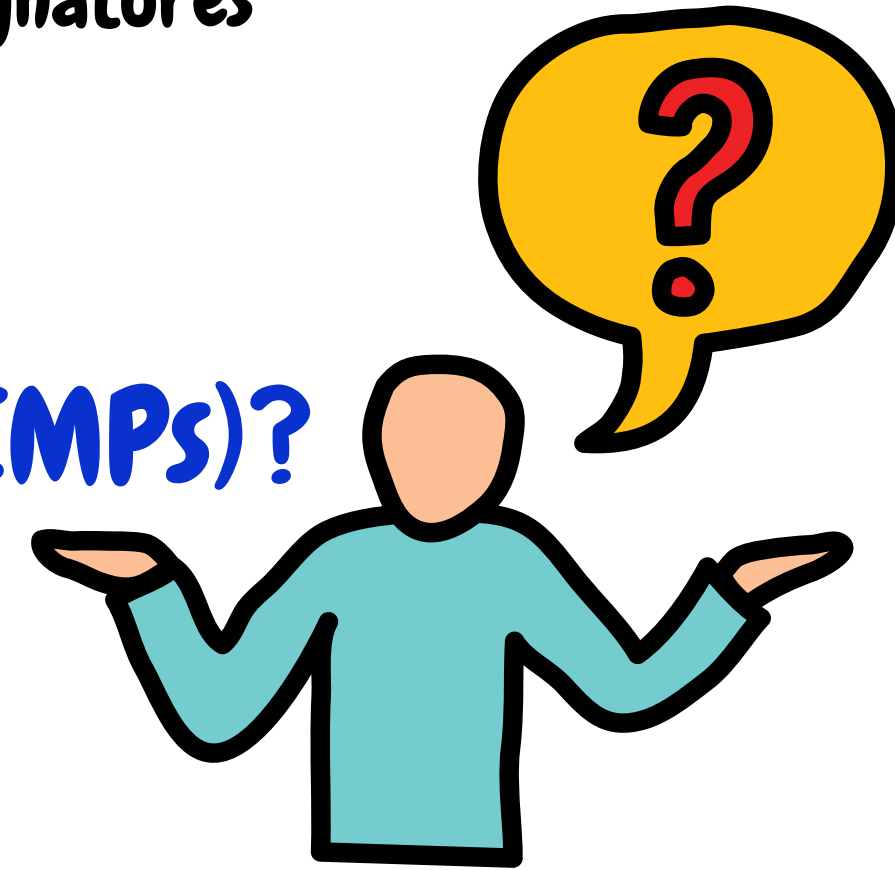
bad luck... nightmare scenario

I have to freeze-in to survive



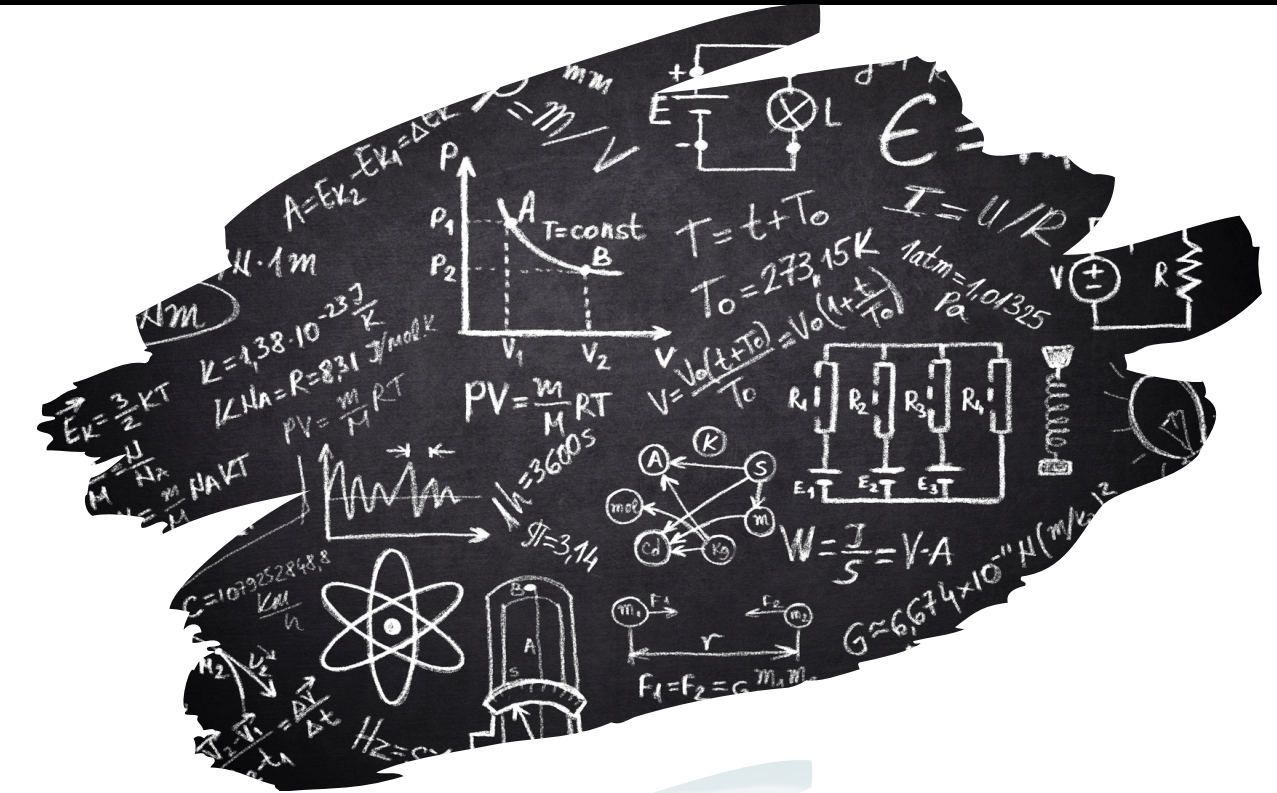
**If you think is plausible to assume tiny couplings, tuned mass differences
several assumptions about complex dark sectors, and lack of smoking gun signatures**

What can you say against thermal relics (or WIMPs)?



**If are near resonances, assume a tiny mass gap
(coannihilation,inelastic), matter domination
periods, we do the same game above, while keeping the
IMPORTANT multifaceted search going**

For theorists this is fun though



For experimentalists, the question is:



which road should I take?

Alternative dark sectors

Thermal relics



Back to Standard Darkness

Thermal Particles

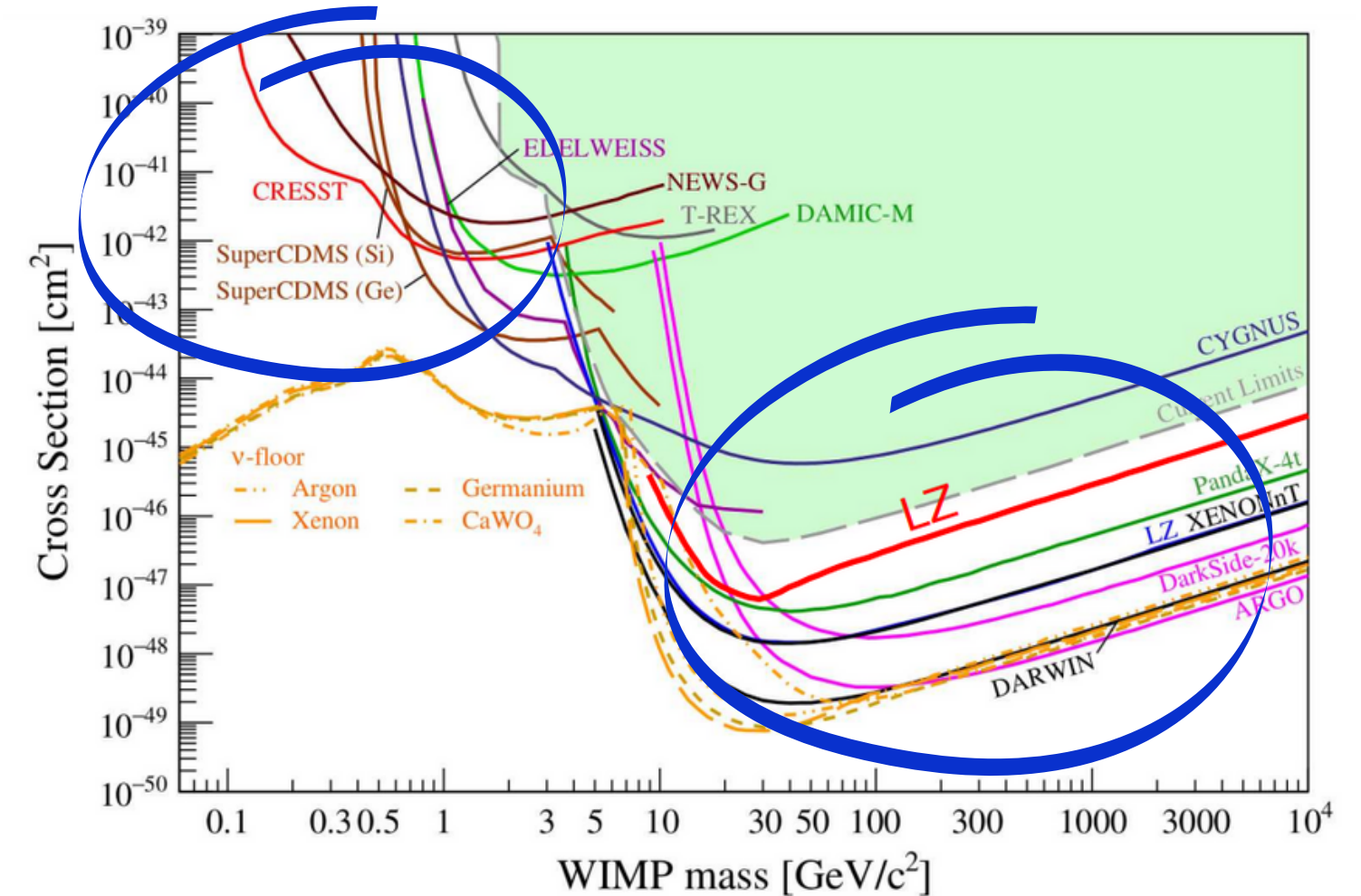
<i>DM bilinear</i>	<i>SM fermion bilinear</i>			
<i>fermion DM</i>	$\bar{f}f$	$\bar{f}\gamma^5 f$	$\bar{f}\gamma^\mu f$	$\bar{f}\gamma^\mu\gamma^5 f$
$\bar{\chi}\chi$	$\sigma v \sim v^2, \sigma_{SI} \sim 1$	$\sigma v \sim v^2, \sigma_{SD} \sim q^2$	—	—
$\bar{\chi}\gamma^5\chi$	$\sigma v \sim 1, \sigma_{SI} \sim q^2$	$\sigma v \sim 1, \sigma_{SD} \sim q^4$	—	—
$\bar{\chi}\gamma^\mu\chi$ (Dirac only)	—	—	$\sigma v \sim 1, \sigma_{SI} \sim 1$	$\sigma v \sim 1, \sigma_{SD} \sim v_\perp^2$
$\bar{\chi}\gamma^\mu\gamma^5\chi$	—	—	$\sigma v \sim v^2, \sigma_{SI} \sim v_\perp^2$	$\sigma v \sim 1, \sigma_{SD} \sim 1$



This translates to the rates at direct detection experiments being suppressed by several orders of magnitude

Sure there are ways to circumvent bounds using particle physics but let's

take very constrained models and see what happens when we play with cosmology



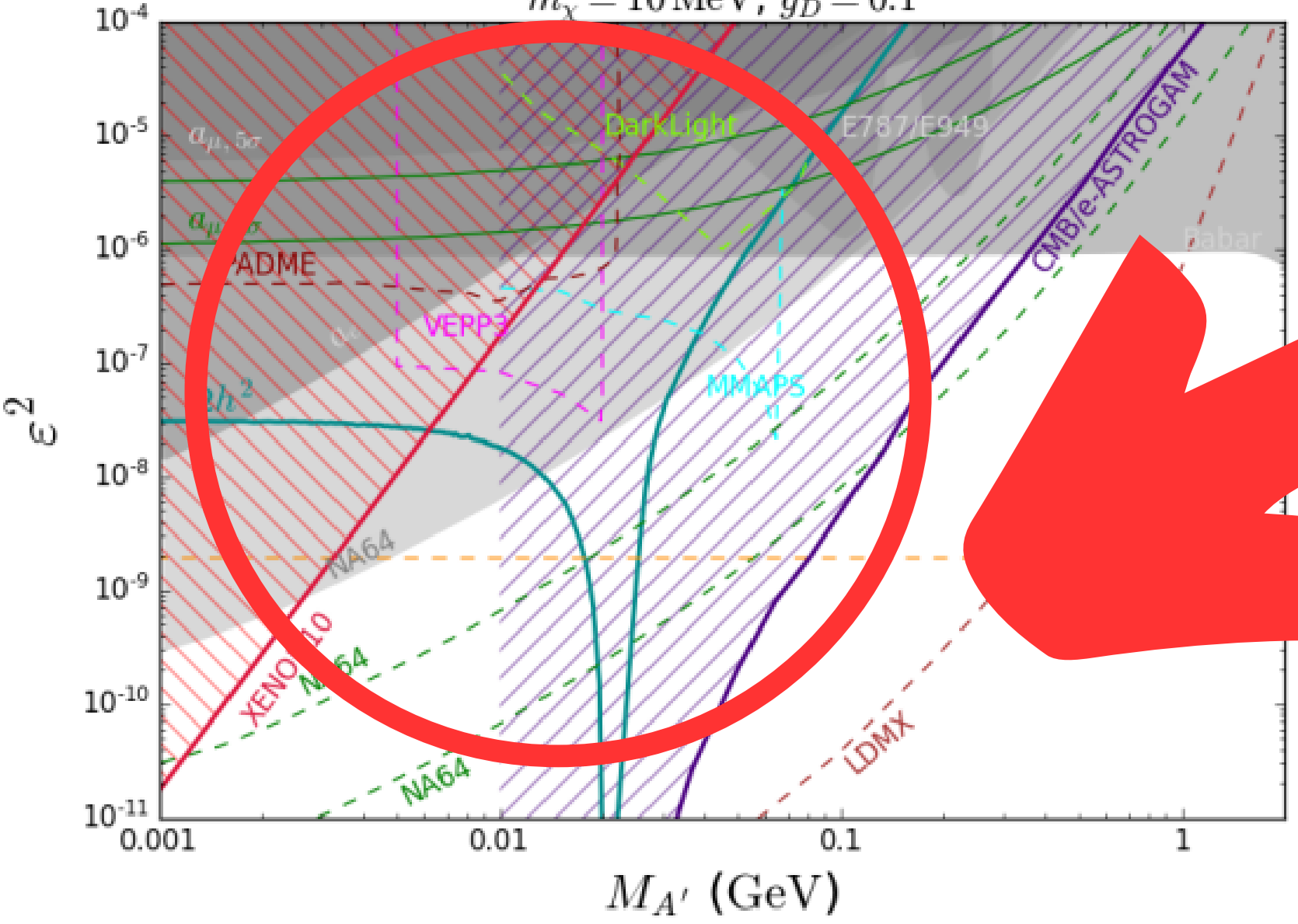
SURE!

Back to Standard Darkness

Taking a very constrained scenario: dark photon

$$\bar{\chi}(-g_D A' - m_\chi)\chi + \sum_i \bar{f}_i(-eq_{f_i} A - \epsilon eq_{f_i} A' - m_{f_i})f_i$$

$m_\chi = 10 \text{ MeV}, g_D = 0.1$



- Thermal relic
- +
- Direct detection
- +
- Indirect Detection
- +
- Accelerator

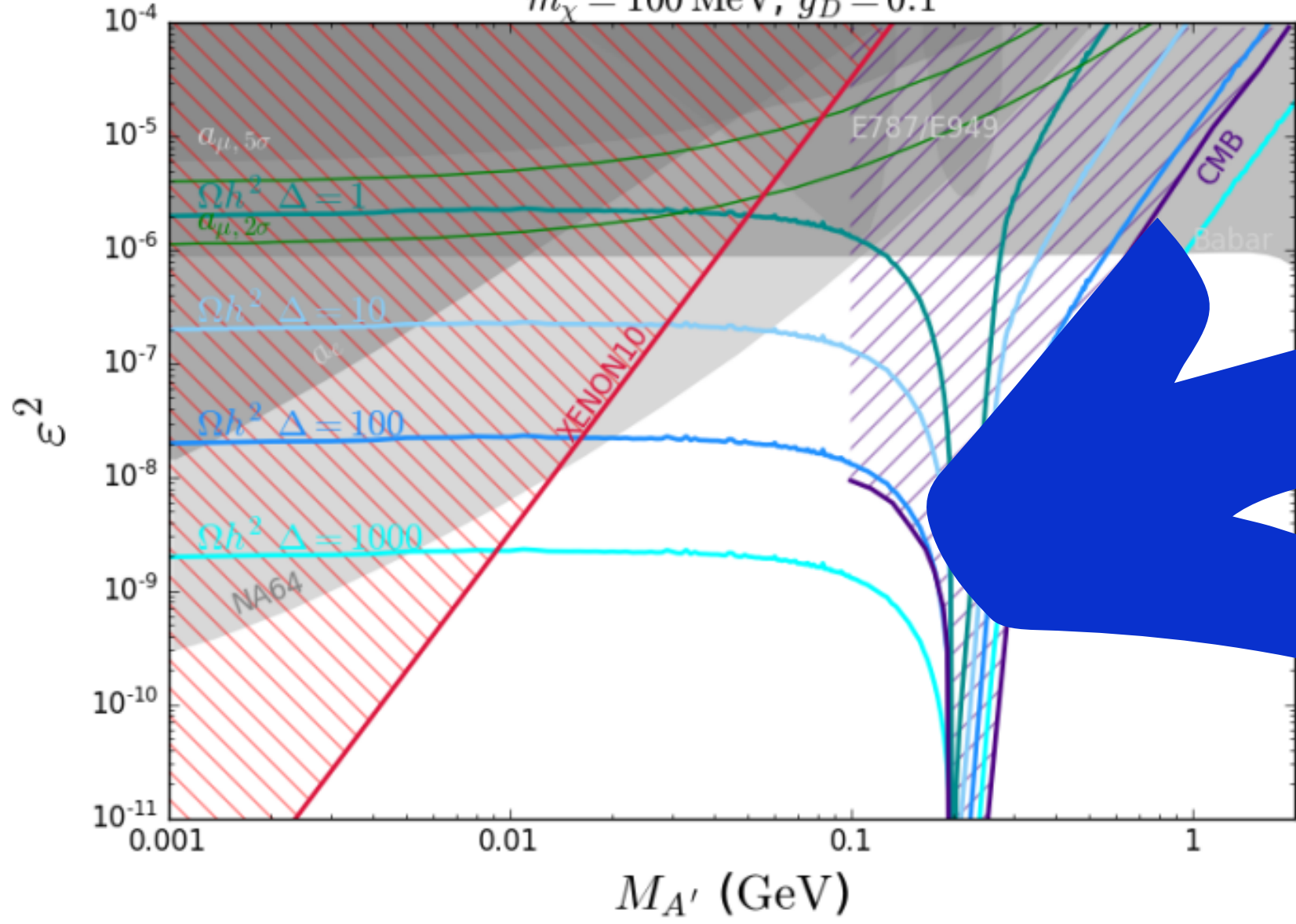
It has been excluded already by several experiments

Back to Standard Darkness

Taking a very constrained scenario: dark photon

$$\bar{\chi}(-g_D A' - m_\chi)\chi + \sum \bar{f}_i(-eq_{f_i}A - \epsilon eq_{f_i}A' - m_{f_i})f_i$$

$m_\chi = 100 \text{ MeV}, g_D = 0.1$



- Thermal relic
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- Direct detection
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- Accelerator
- +

Non-standard Cosmologies

Back to Standard Darkness

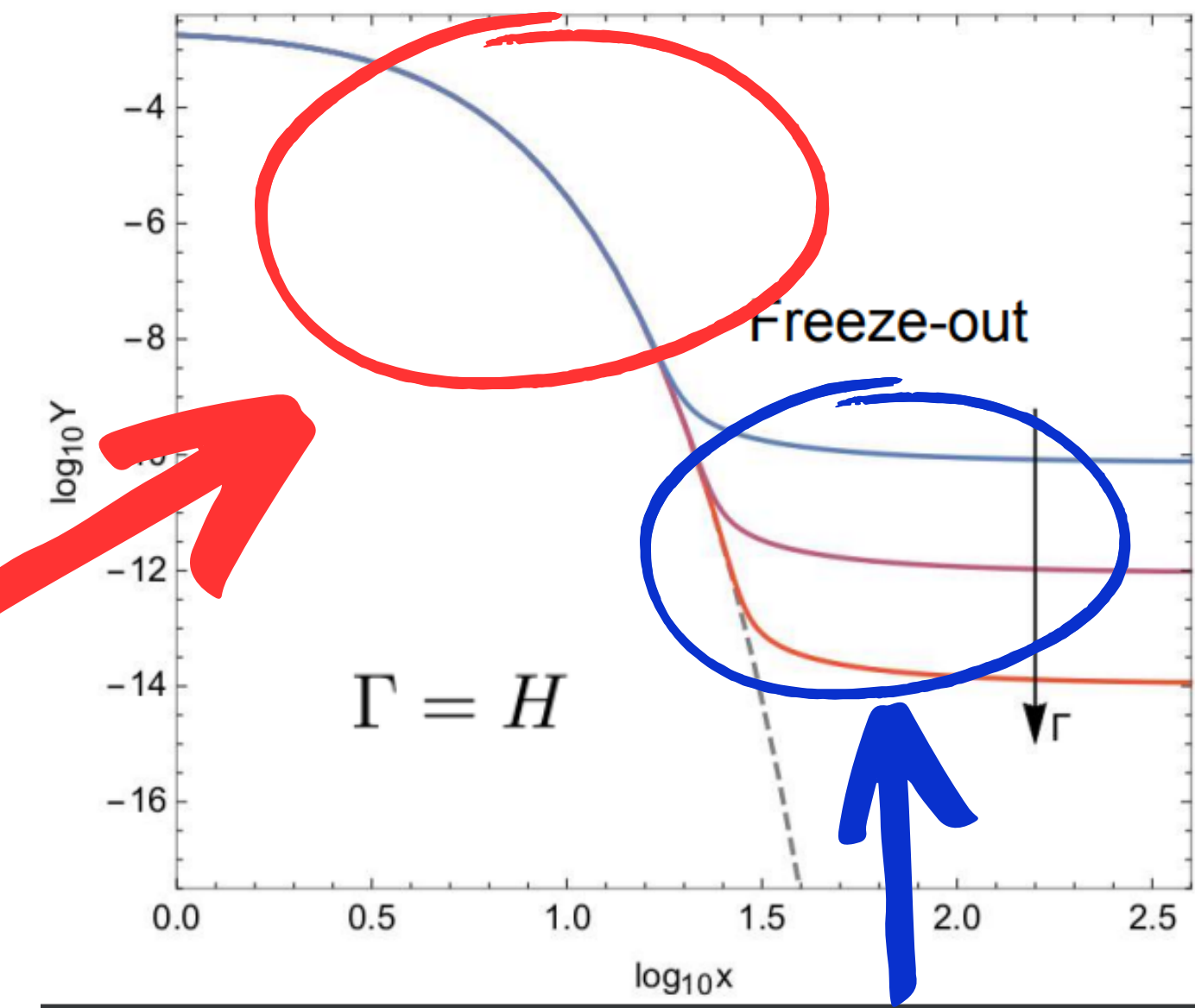
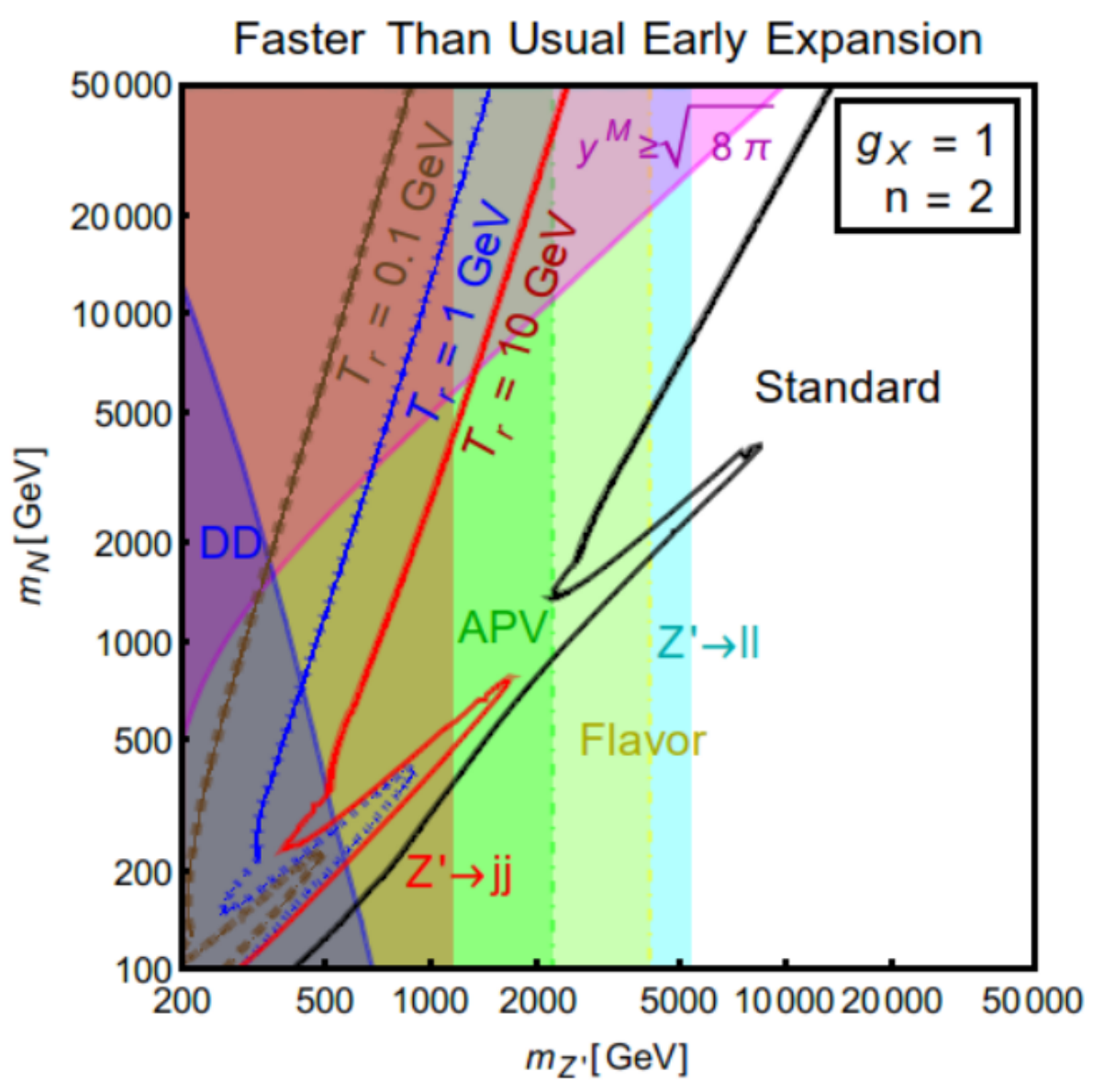
Taking a very constrained scenario: Z'

$$+ \frac{1}{4} g_X (N_{1R} \gamma^\mu \gamma_5 N_{1R}) Z'_\mu - \frac{g_X}{2} Q_{X_f} (\bar{\psi}_f \gamma^\mu \psi_f) Z'_\mu$$

Was the universe expanding too fast?

$$H(T) \approx \frac{\pi}{3} \sqrt{\frac{g_\star}{10}} \frac{T^2}{M_{Pl}} \left(\frac{T}{T_r}\right)^{n/2}$$

A way to make under-abundant dark matter consistent with the data

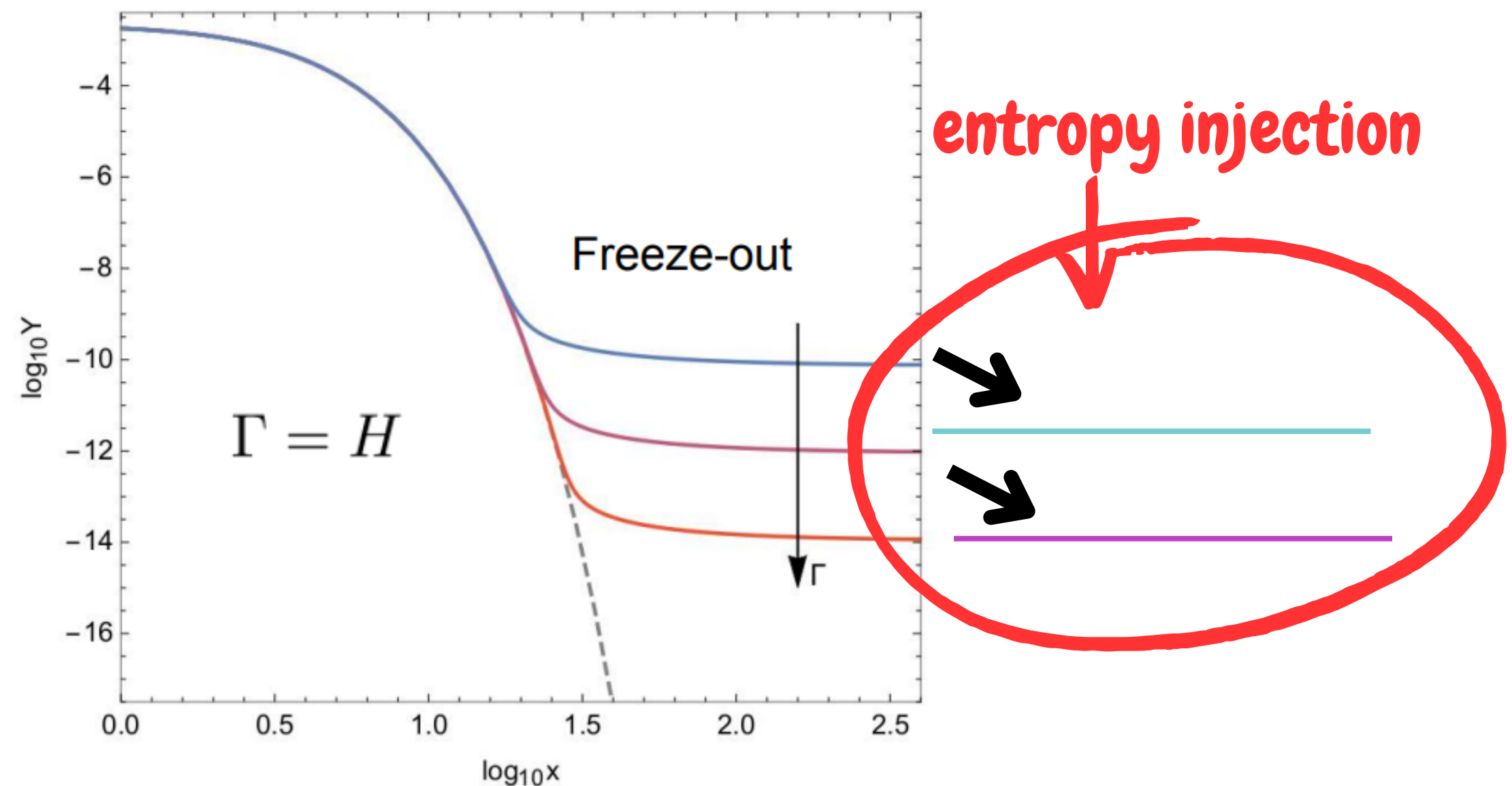


$$H_R(T) = \frac{\pi}{3} \sqrt{\frac{g_\star}{10}} \frac{T^2}{M_{Pl}}$$

Back to Standard Darkness

Taking a very constrained scenario: Z'

$$+ \frac{1}{4} g_X (N_{1R} \gamma^\mu \gamma_5 N_{1R}) Z'_\mu - \frac{g_X}{2} Q_{X_f} (\bar{\psi}_f \gamma^\mu \psi_f) Z'_\mu$$

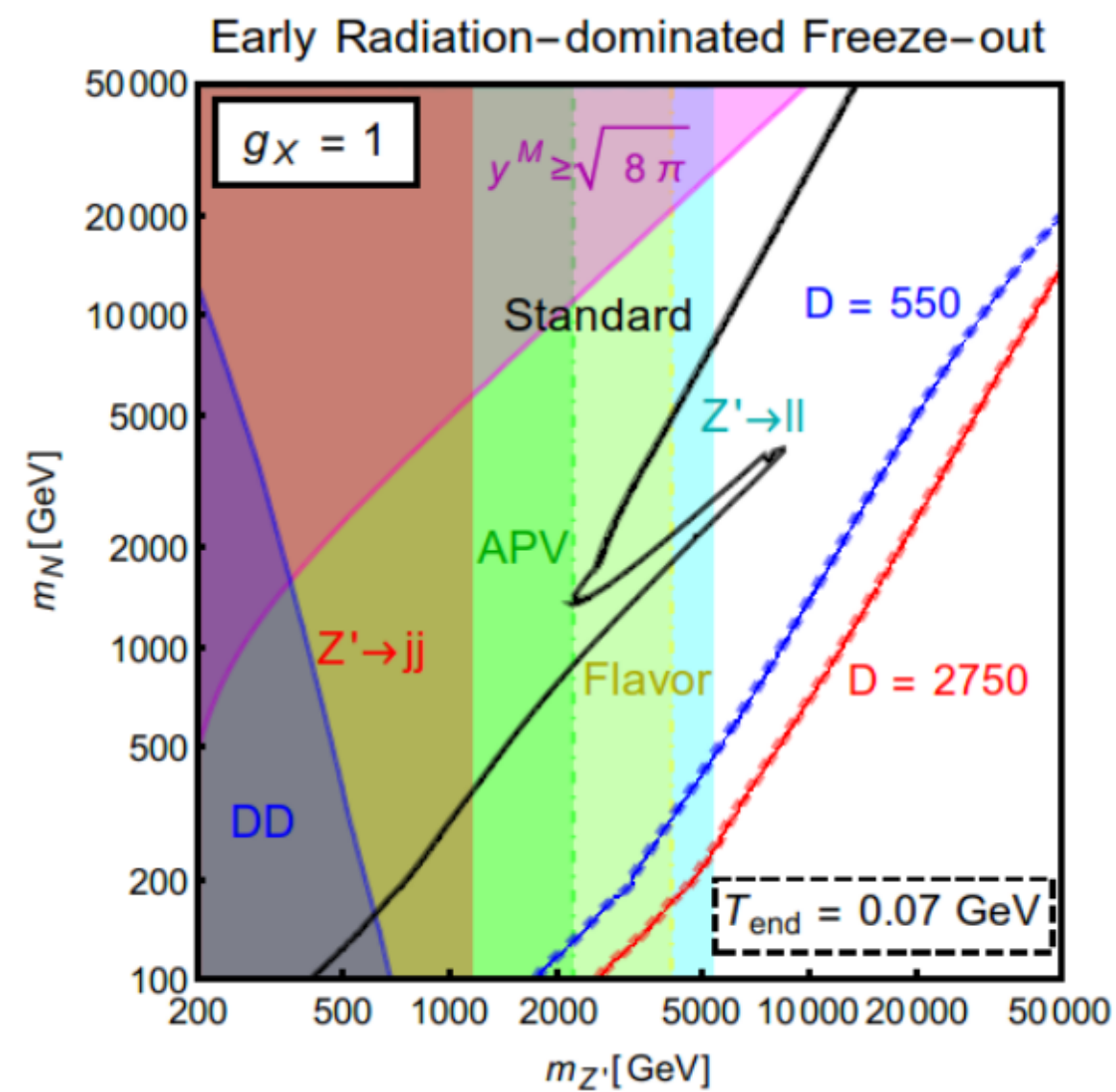


Back to Standard Darkness

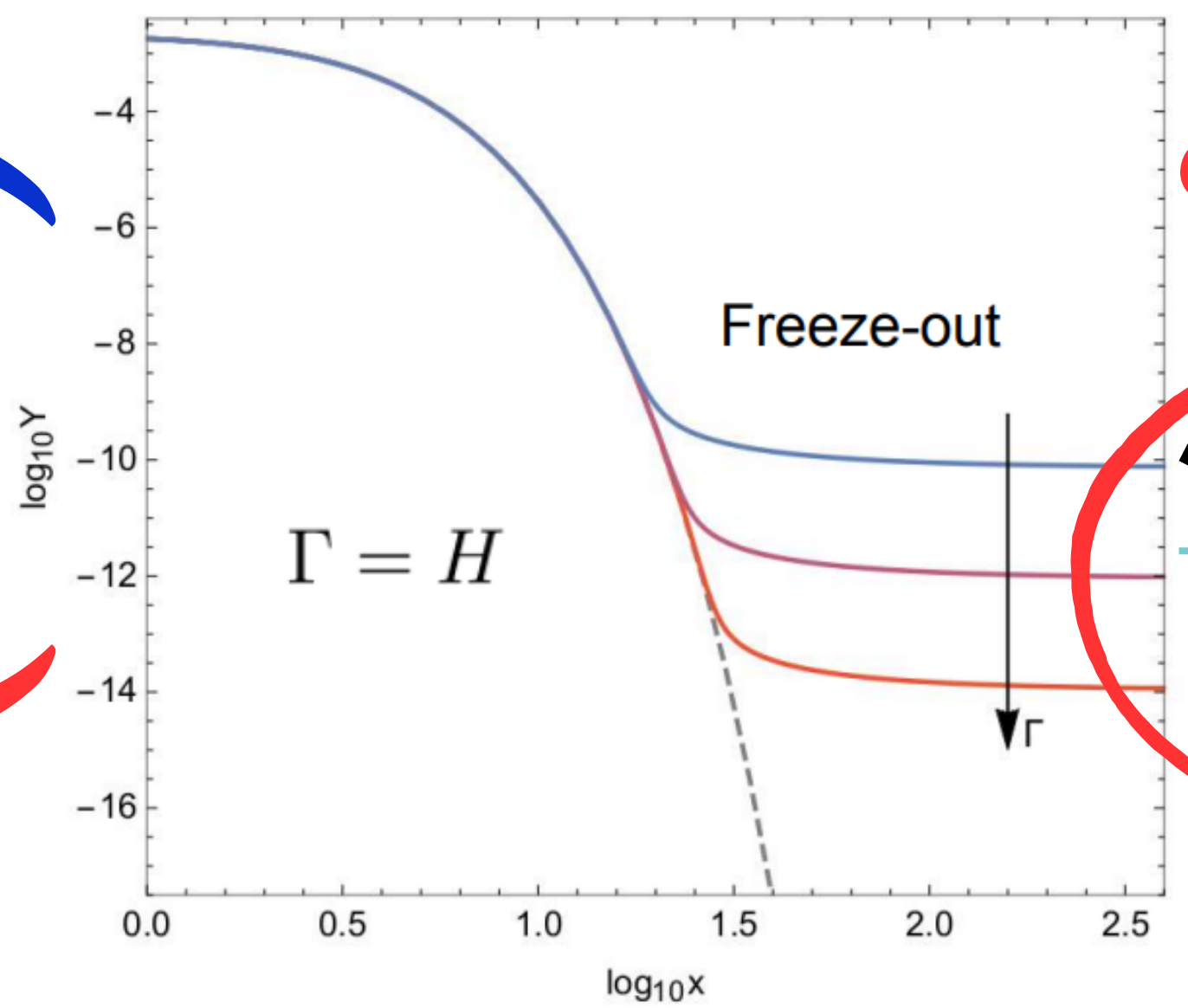
Taking a very constrained scenario: Z'

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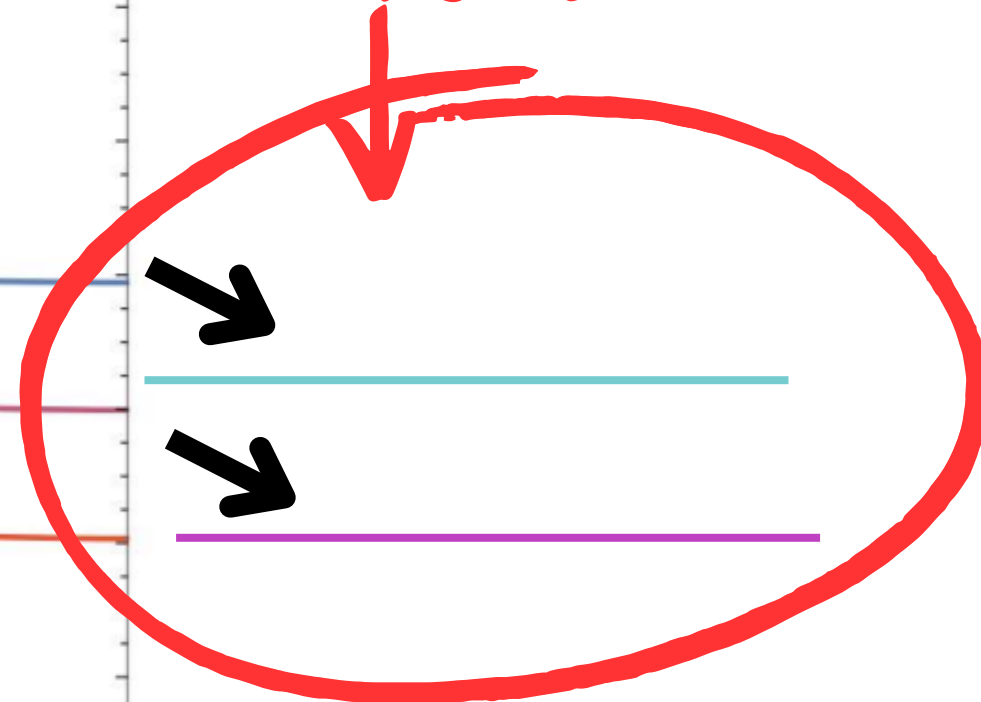
PROGRESS IS NEEDED HERE:
 Scalar dominates the energy density and decays into radiation



Dilution factors



entropy injection



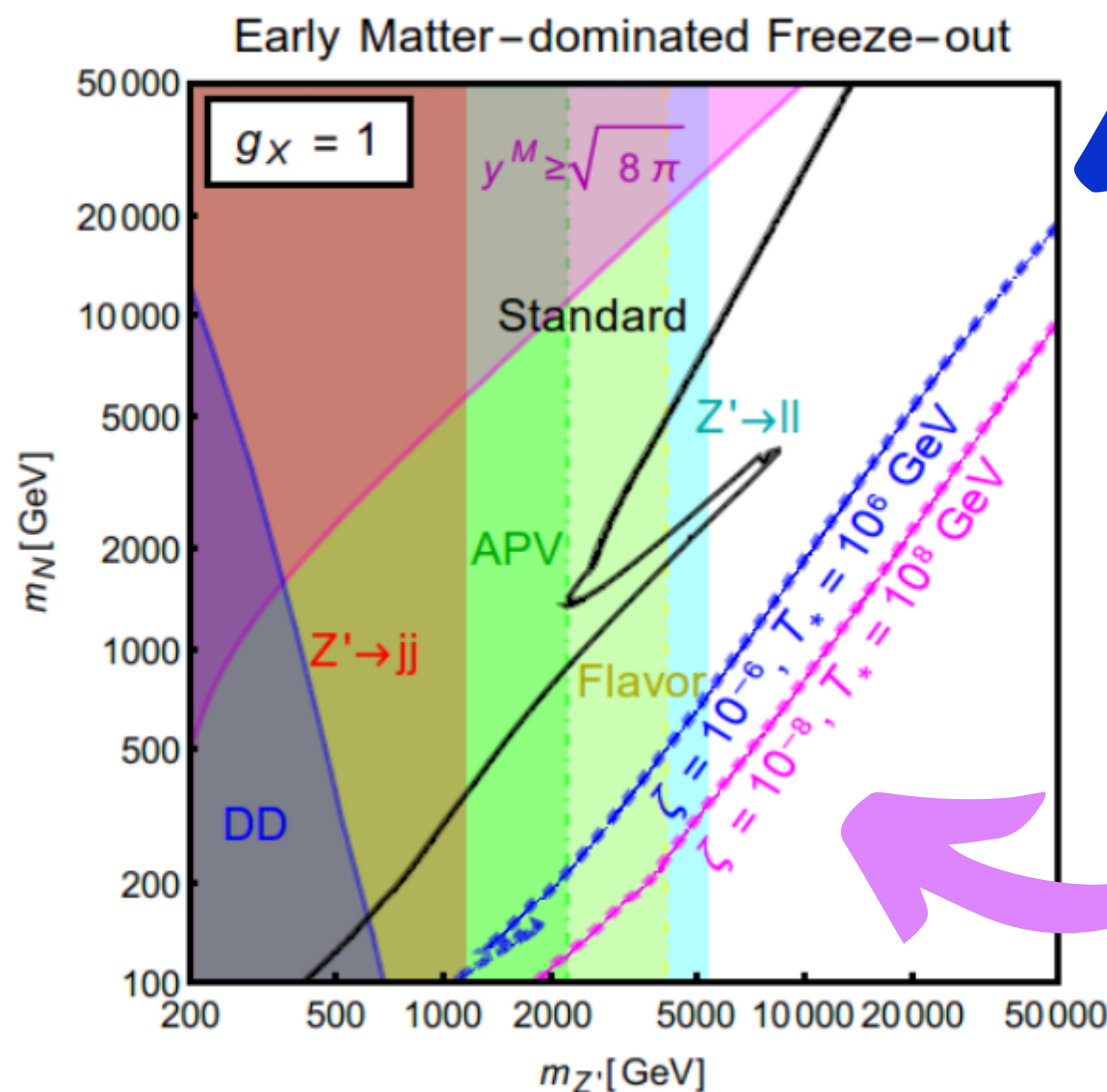
Back to Standard Darkness

Taking a very constrained scenario: Z'

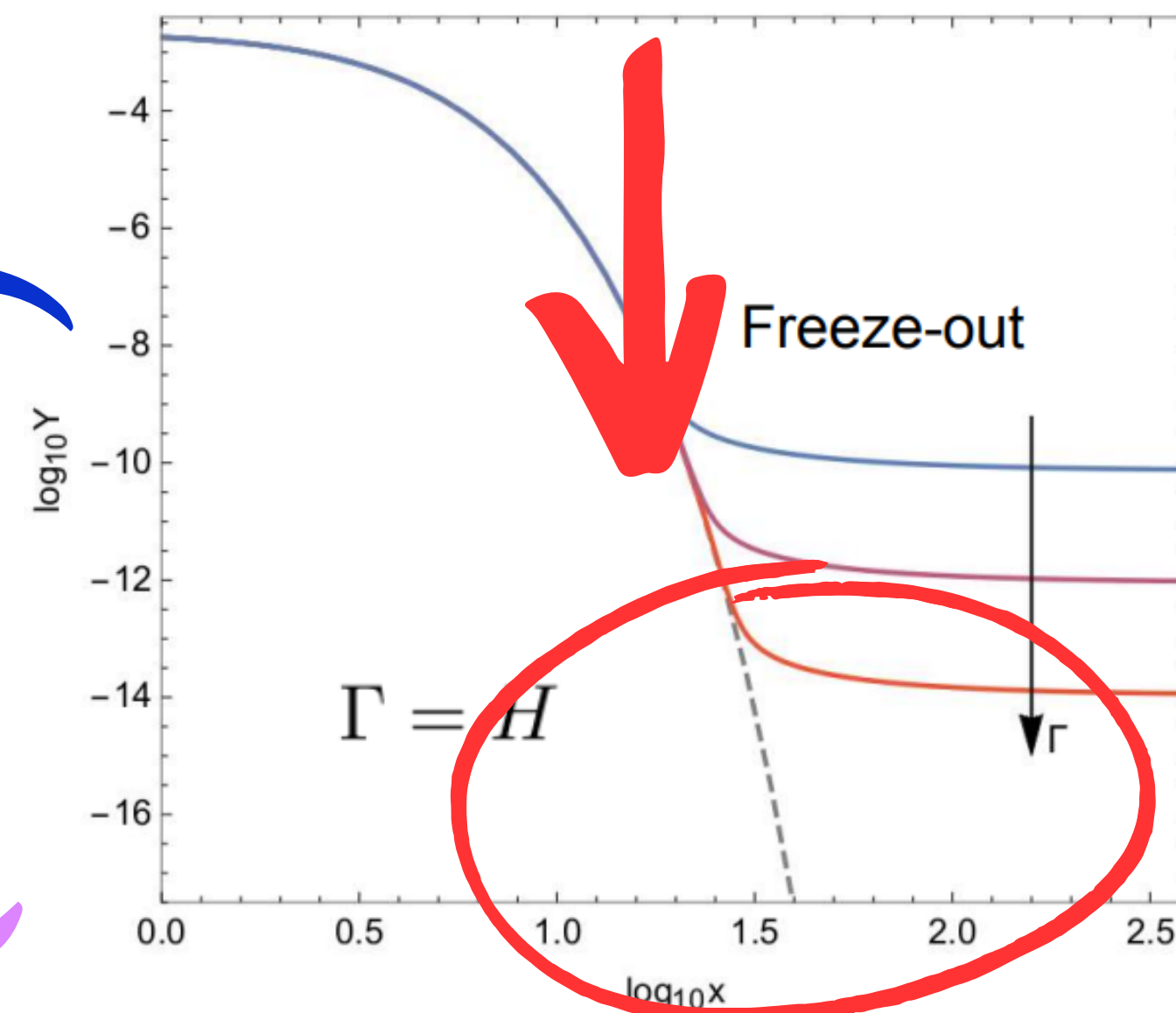
$$+ \frac{1}{4} g_X (N_{1R} \gamma^\mu \gamma_5 N_{1R}) Z'_\mu - \frac{g_X}{2} Q_{X_f} (\bar{\psi}_f \gamma^\mu \psi_f) Z'_\mu$$

$$\zeta = \frac{s(T_1)}{s(T_2)} \sim \frac{T_{end}}{T_\star}$$

$$T_{eq} \equiv T_\star \left(\frac{a_\star}{a(T_{eq})} \right)$$



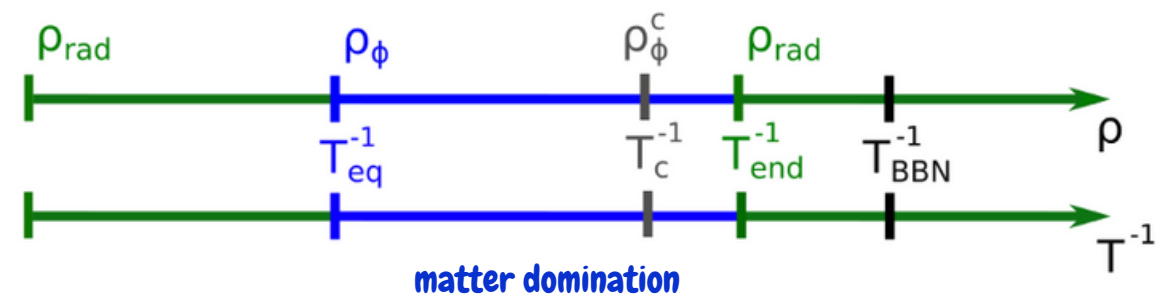
$H \propto T^{3/2}$ slower



Back to Standard Darkness

Taking a very constrained scenario: Z'

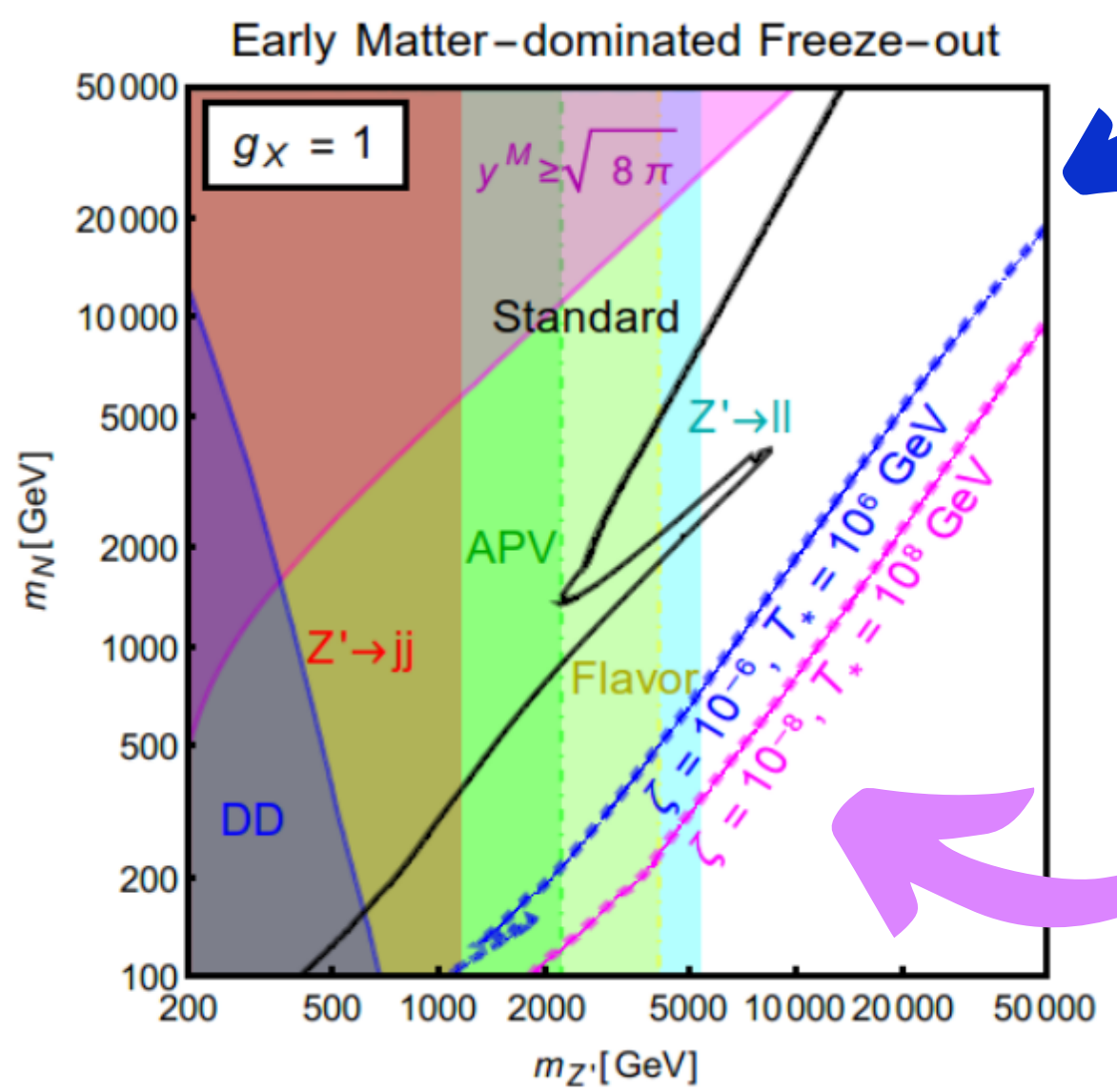
$$+ \frac{1}{4} g_X (N_{1R} \gamma^\mu \gamma_5 N_{1R}) Z'_\mu - \frac{g_X}{2} Q_{X_f} (\bar{\psi}_f \gamma^\mu \psi_f) Z'_\mu$$



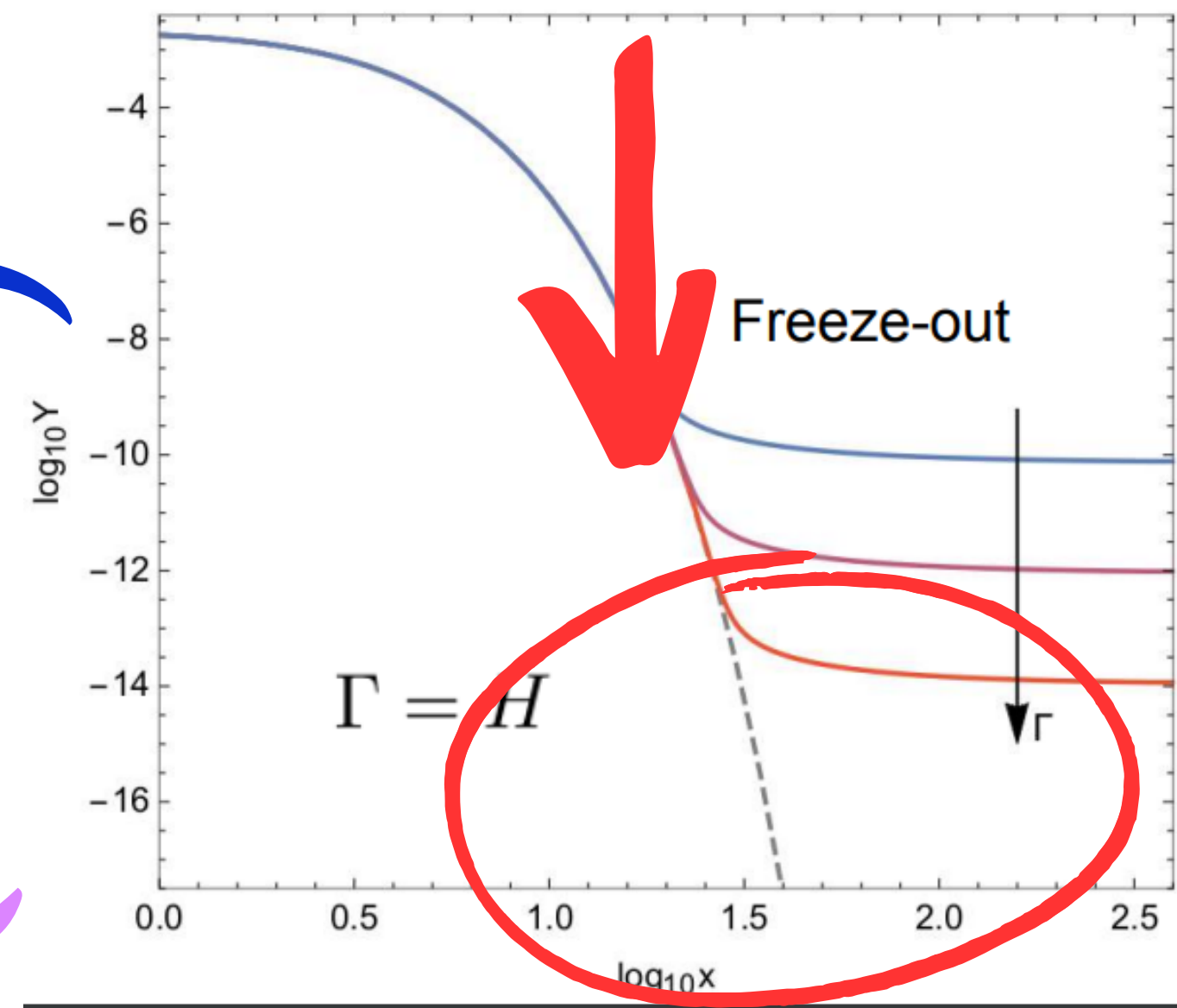
The universe is radiation-dominated at $T=T_*$

$$T_{end} \equiv \left[\frac{90 M_{Pl}^2}{\pi^2 g_*(T_{end})} \right]^{1/4} \Gamma_\phi^{1/2}$$

$$\Omega_N h^2 = \zeta \Omega_N^{MD} h^2$$



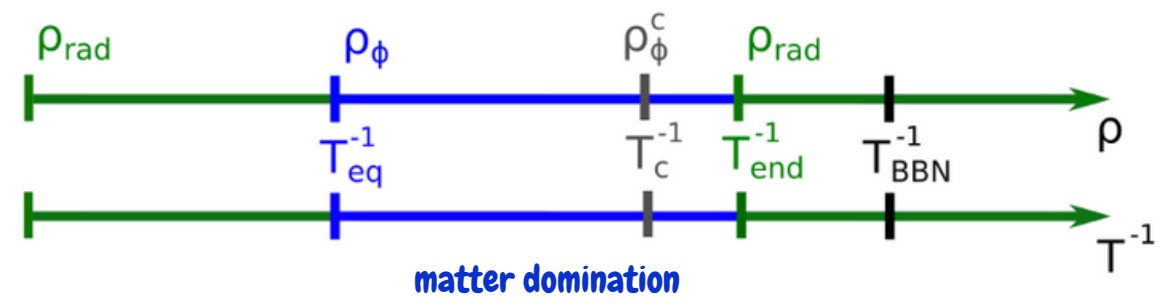
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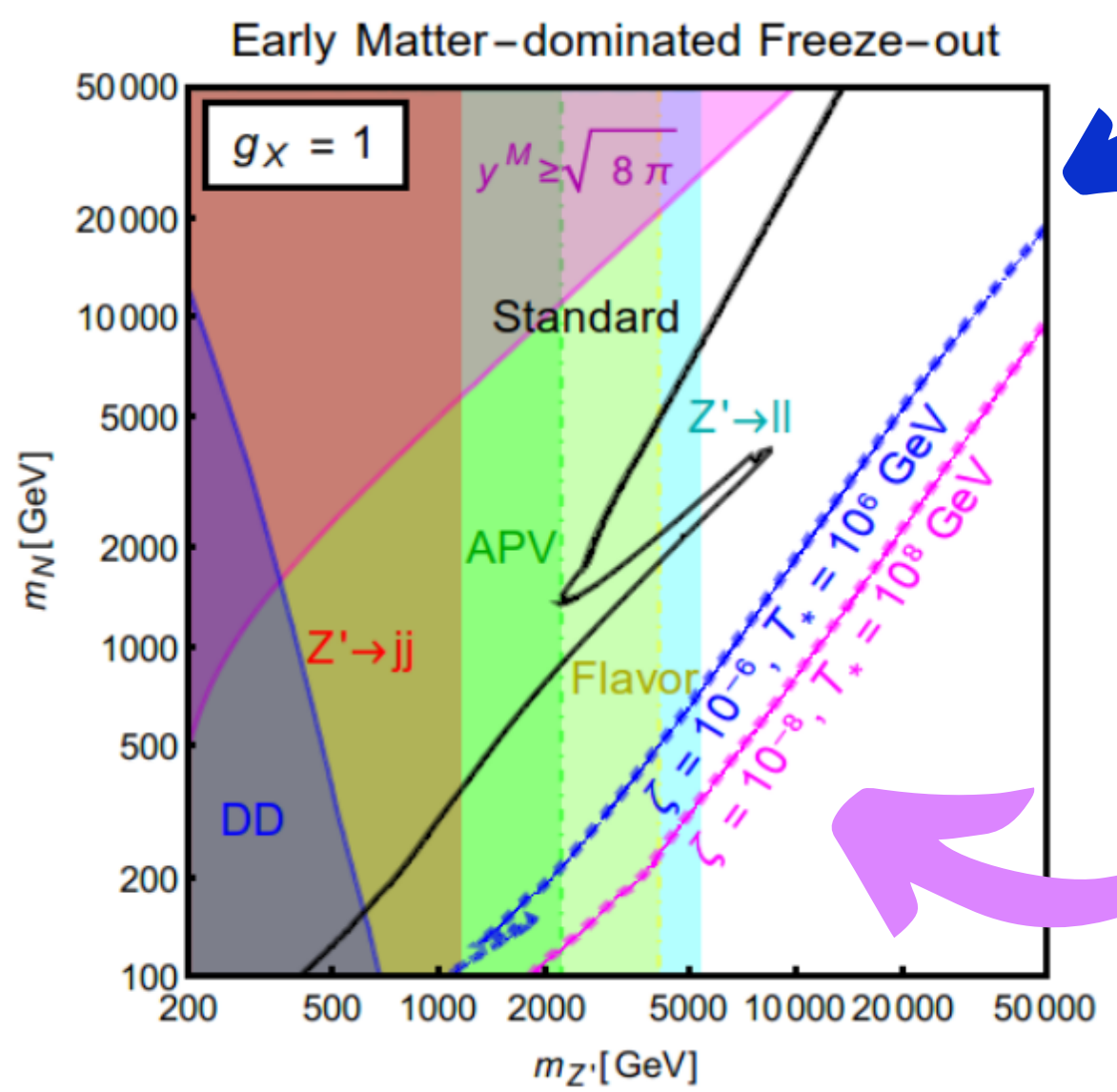
$$+ \frac{1}{4} g_X (N_{1R} \gamma^\mu \gamma_5 N_{1R}) Z'_\mu - \frac{g_X}{2} Q_{X_f} (\bar{\psi}_f \gamma^\mu \psi_f) Z'_\mu$$



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Non-standard Cosmologies

- Thermal relic
- +
- Direct detection
- +
- Indirect Detection
- +
- Colliders
- +

Neutron stars as laboratories

Rate at which energy is transferred

$$\dot{E} = \frac{E_{\chi}^R \dot{m}}{m_{\chi}} f$$

$\xrightarrow{\hspace{1.5cm}}$ Flux of DM
 $\xrightarrow{\hspace{1.5cm}}$ is the capture efficiency
 $\xrightarrow{\hspace{1.5cm}}$ DM mass

Heating

$$T_{NS} = \left[\frac{(\gamma^R - 1) b_{max}^2 v_{\chi} \rho_{\chi}}{\sigma_B R^2} \right]^{1/4} \left(1 - \frac{2GM_{\star}}{R} \right)^{\frac{1}{2}} f^{1/4}$$

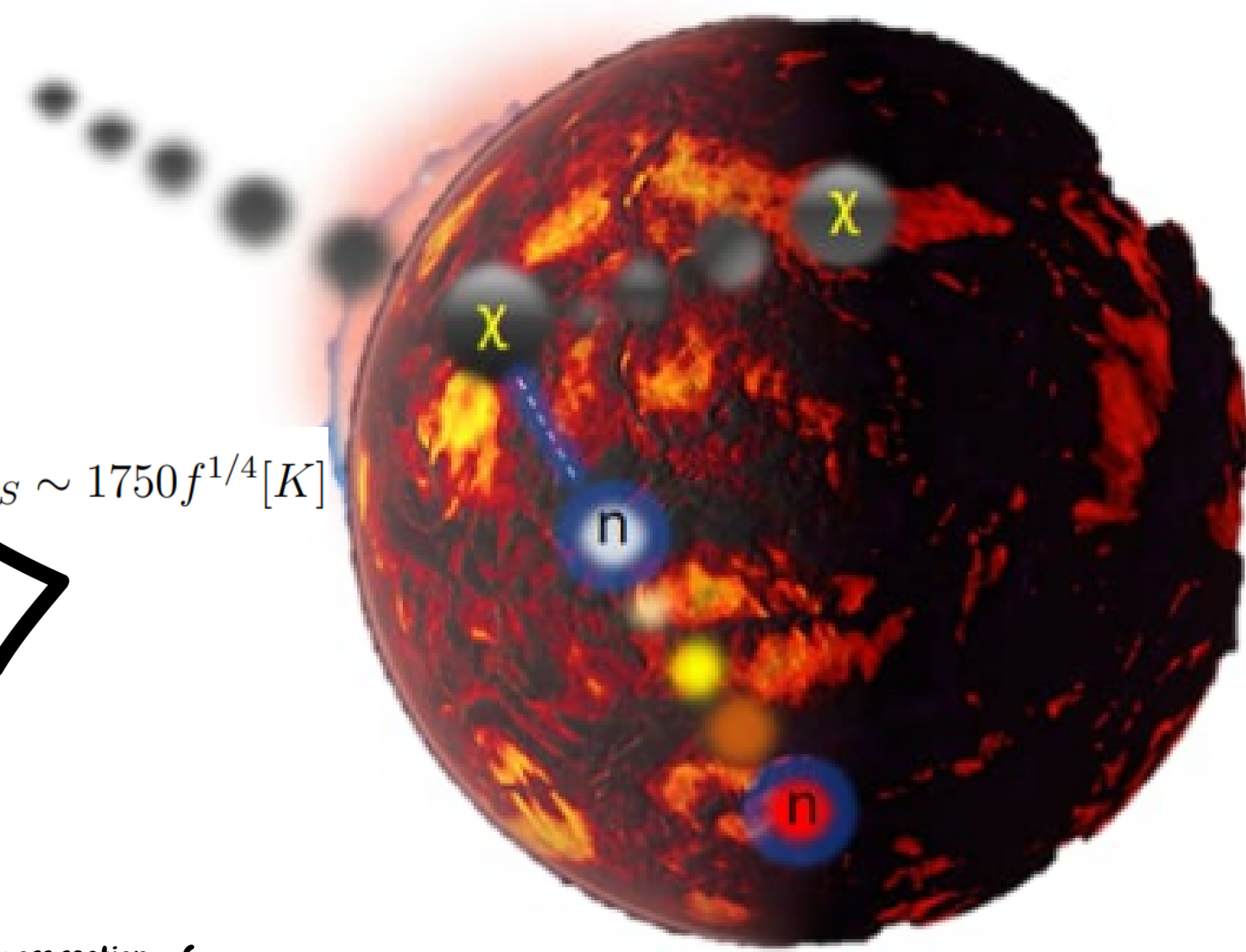
$T_{NS} \sim 1750 f^{1/4} [K]$

Astrophysics

Connection to Particle Physics

It depends on the relation between dark matter-nucleon scattering cross section (σ_n) and a saturation cross section, σ_S , above which all the transient dark matter is captured so that

$$f = \min(\sigma_{\chi n} / \sigma_S, 1)$$



Bell, Busoni, Robles, Virgator, arxiv: 2004.14888

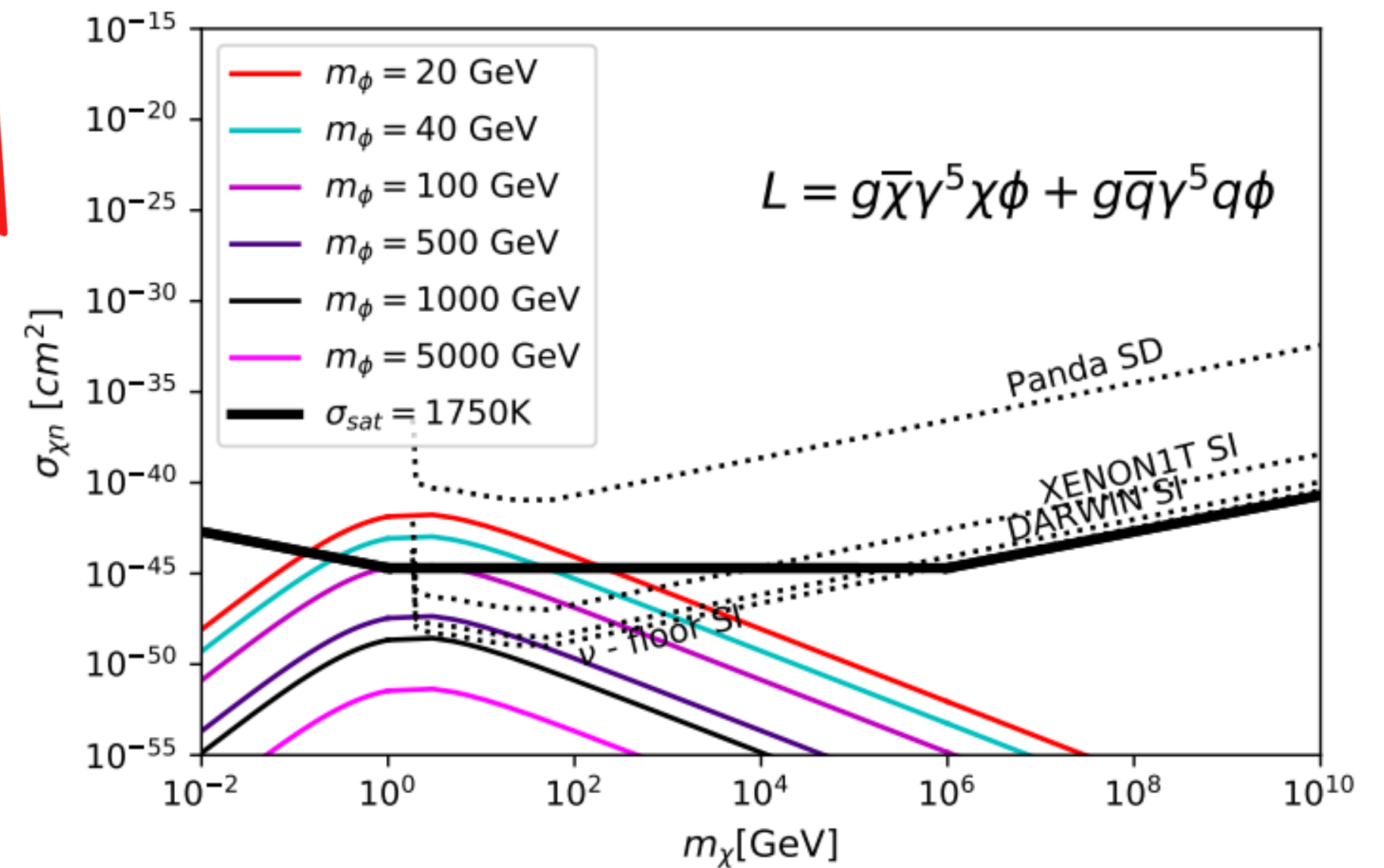




Neutron stars with temperatures around 2000 K located 10 pc away could be detected at the JWST. Owing to the compact ($R \sim 10$ km) size of NS, they will appear as unresolved, extremely faint point sources, even for a cutting-edge facility such as the JWST. The Near-Infrared Camera has the potential to discover such neutron stars with 24h of exposure.

SCATTERING CROSS SECTION

Fermion DM + pseudoscalar Mediator

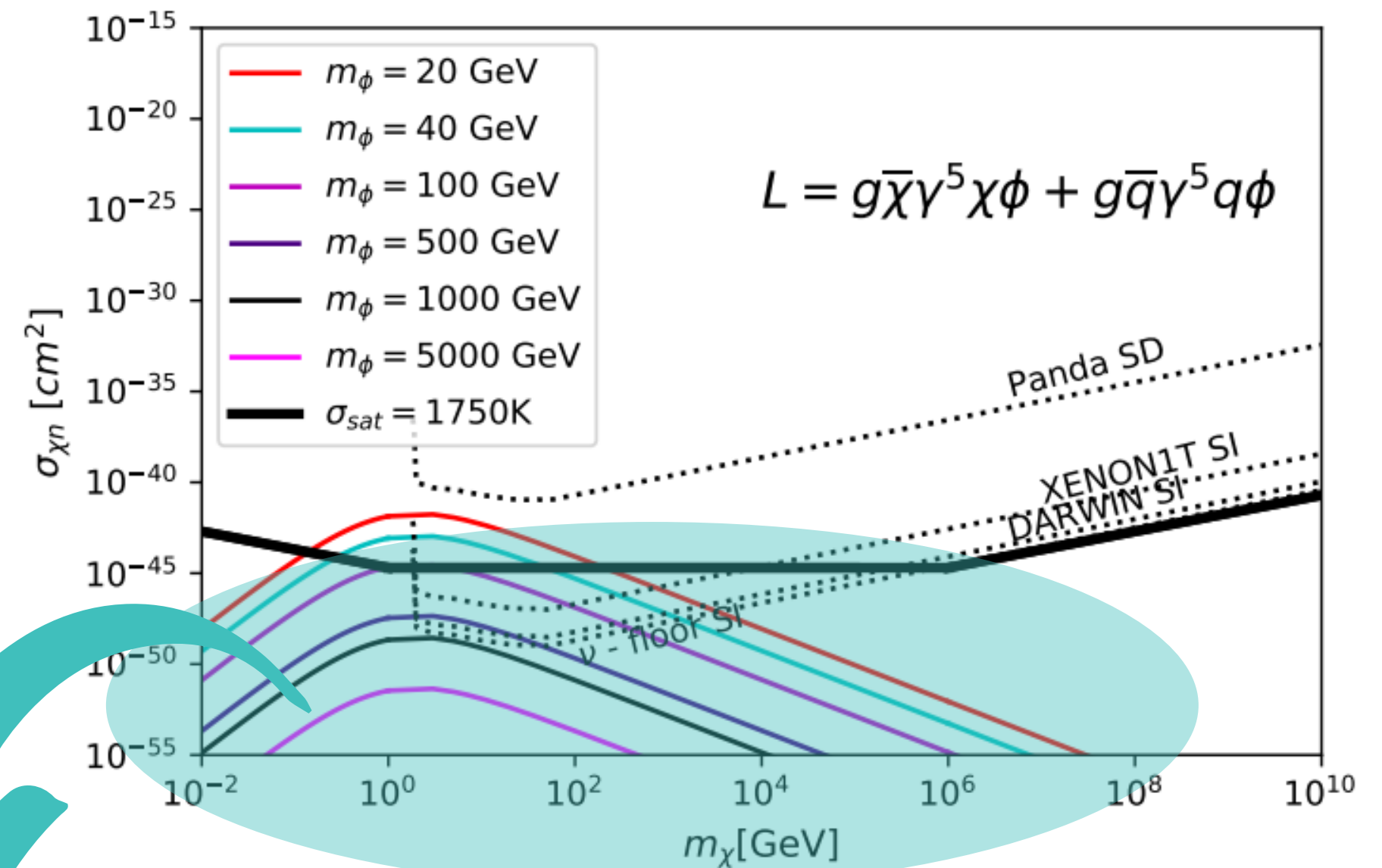
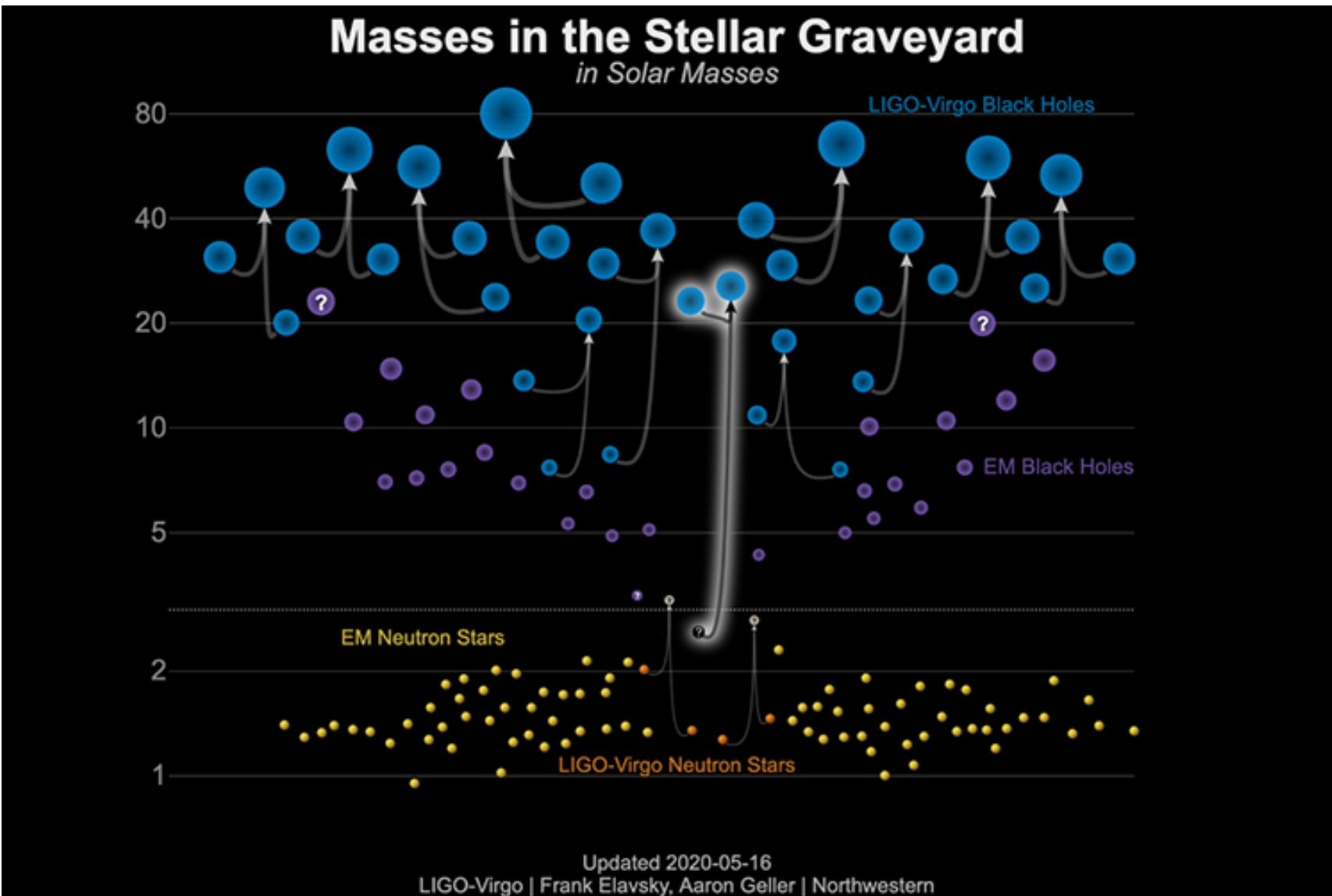


MASS

Maity, FSQ ariv: 2104.02700

Collective effect is needed

Fermion DM + pseudoscalar Mediator



Neutron stars will constitute an important probe!

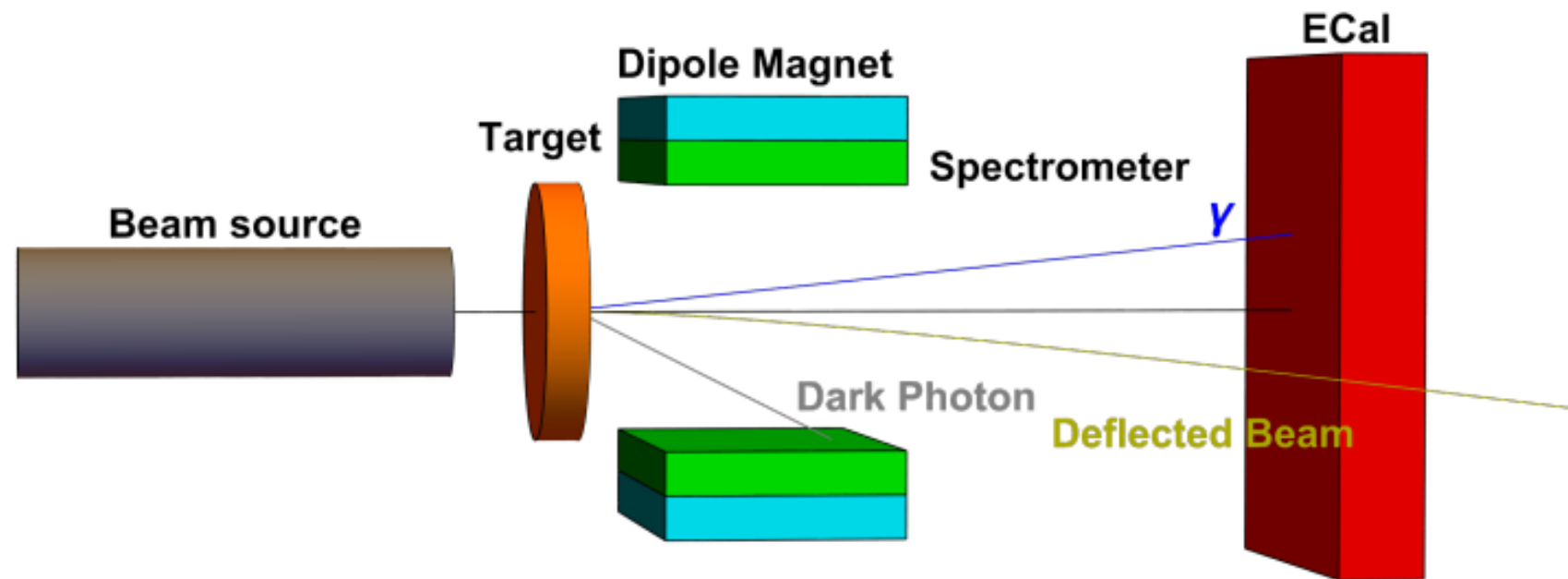


Search for Dark Sector by Repurposing the UVX Brazilian Synchrotron

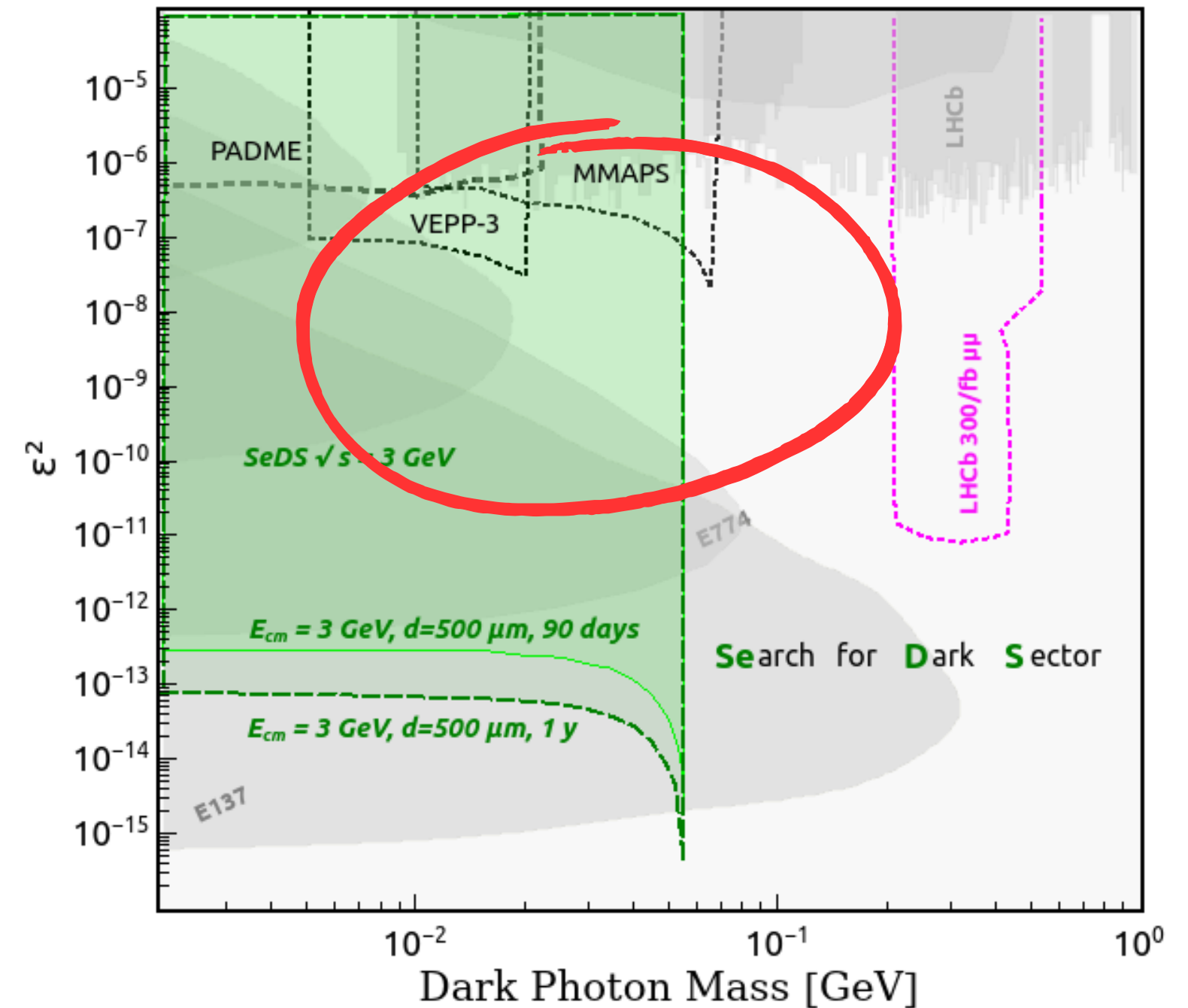
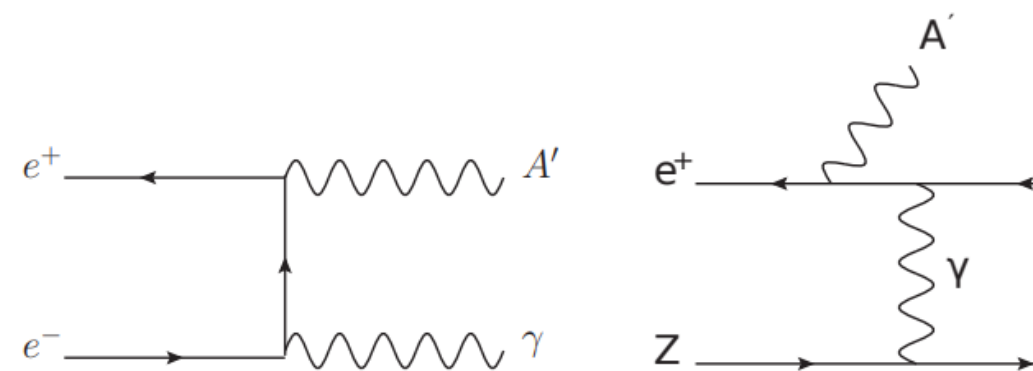
#1

L. Duarte (IIP, Brazil), L. Lin (LNLS, Campinas), M. Lindner (Heidelberg, Max Planck Inst.), V. Kozhuharov (Sofiya U. and INFN, Italy), S.V. Kuleshov (Andres Bello Natl. U. and Unlisted, CL) et al. (Jun 10, 2022)

e-Print: 2206.05305 [hep-ph]



**e^+e^- collisions
competitive probe
In Latin-America**



Thank you

Without interdisciplinary searches one cannot claim the discovery of dark matter particles. Alternatives models for particle dark matter are welcome but without a strong experimental program involved discoveries are not expected