

Achievements and open questions WP1

Speakers:

Eliana Gianfelice, Ivan Koop (BINP), Dr Tatiana Pieloni (EPF Lausanne)

Ivan Koop

EPOL-2022, September 23, 2022

WP1: 3162/1-K01

Convener: Eliana Gianfelice



15:30

CEPC Polarization Simulations

🕒 30m

Speaker: Mr Zhe Duan (IHEP)

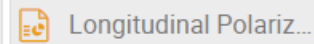
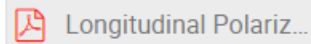


16:00

Longitudinal Polarization in CEPC

🕒 30m

Speaker: Tao Chen



16:30

FCCEe Polarization studies with BMAD

🕒 20m

Speaker: Yi Wu (EPFL - Ecole Polytechnique Federale Lausanne (CH))



16:50

Break

🕒 10m

17:00

FCCEe Polarization SITROS simulations

🕒 30m

Speaker: Eliana Gianfelice

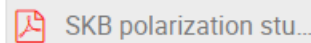


17:30

Spin Tracking Studies for the SuperKEKB Polarization Upgrade

🕒 30m

Speaker: Yuhao Peng



Zhe Duan: Overall progress in CEPC polarization studies

- Study of the radiative depolarization effects in CEPC collider rings [1] (this talk)
 - Spin tracking simulations for CEPC CDR lattice
 - Comparison between simulations with theories
- Longitudinally polarized colliding beams (Tao Chen's talk)
 - Polarization maintenance via the “spin resonance free” feature of the CEPC booster lattice [2,3]
 - Spin rotator design at CEPC-Z energy [4]
- Resonant depolarization (Sep 29 WP1 talk)
 - The option to prepare polarized e⁺/e⁻ bunches from the injector
- Compton polarimeter via scattered electron distribution [5]

[1] W. H. Xia, Z. Duan, Y. W. Wang, B. Wang, J. Gao, arXiv:2204.12718v1 [physics.acc-ph]

[2] V. Ranjbar, et al., PRAB 21, 111003 (2018). [3] Z. Duan, presentation at eeFACT 2022.

[4] W. Xia et al., RDTM (2022) doi: 10.1007/s41605-022-00344-2

[5] S. H. Chen et al., JINST 17, P08005, (2022)

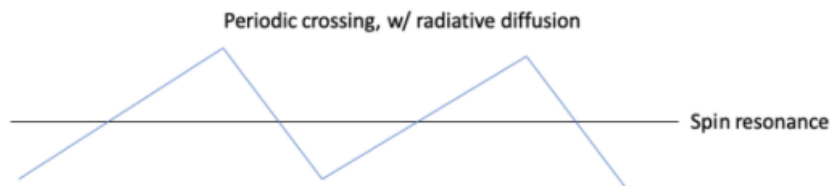
Correlated and uncorrelated regime of spin resonance crossing

- Follow the “dynamical picture” [1] that the instantaneous spin precession rate ν_s is dependent on the instantaneous energy deviation δ , underlying spin resonances could be crossed as a result of synchrotron oscillations
- The following two regimes of spin resonance crossing were also proposed in [1]

- Correlated regime: $\kappa = \frac{\nu_0^2 \lambda_p}{\nu_z^3} \ll 1$

- Non-resonant spin diffusion & perturbative treatment of $\frac{\partial \hat{n}}{\partial \delta}$ applies

$$\frac{\tau_p}{\tau_d} \approx \frac{11}{18} \sum_{k=n-l}^{n+l} \sum_{m=-\infty}^{\infty} \left(\frac{\nu_0^2 |\tilde{\omega}_k|^2 e^{-\sigma^2} I_m(\sigma^2)}{[(\nu_0 - k - m\nu_z)^2 - \nu_z^2]^2} + \frac{(\nu_0 - k)^2 |\tilde{\lambda}_k|^2 e^{-\sigma^2} I_m(\sigma^2)}{[(\nu_0 - k - m\nu_z)^2 - \nu_z^2]^2} \right)$$



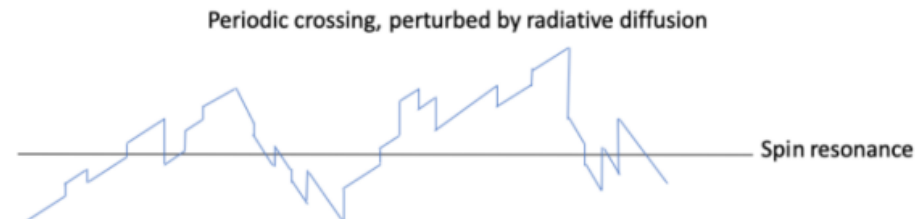
- Uncorrelated regime:

$$\kappa = \frac{\nu_0^2 \lambda_p}{\nu_z^3} \ll 1 \text{ is violated and } \frac{\nu_0 \sigma_\delta}{\nu_z} \gg 1$$

- Resonant spin diffusion

$$\lambda_d = \pi \sum_k \langle |\tilde{\epsilon}_k|^2 \delta(\nu - k) \rangle$$

$$\frac{\tau_p}{\tau_d} \approx \frac{\sqrt{\pi/2}}{\lambda_p} \sum_{k=n-l}^{n+l} \frac{|\tilde{\omega}_k|^2}{\sigma_0} \exp\left[-\frac{(\nu_0 - k)^2}{2\sigma_0^2}\right]$$

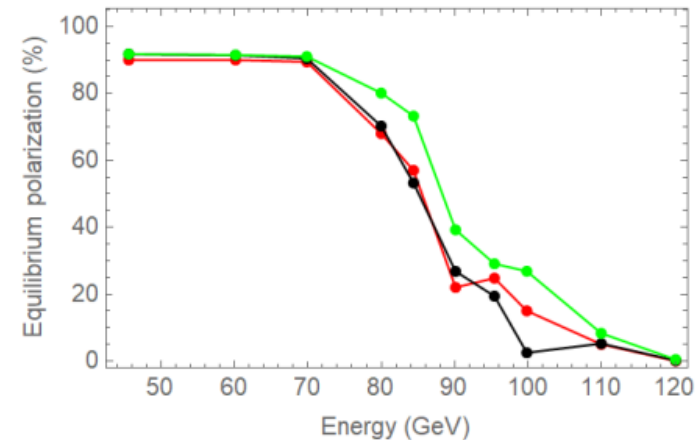


Case study: dependence on beam energy

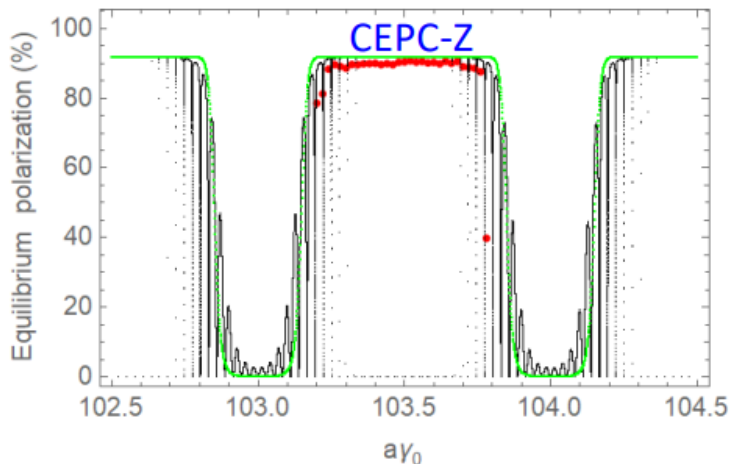
- Increasing beam energy lead to larger σ_δ , modulation index σ and correlation index κ

TABLE IV. The CEPC lattice parameters. (* indicates the planned operation energies in the CEPC CDR.)

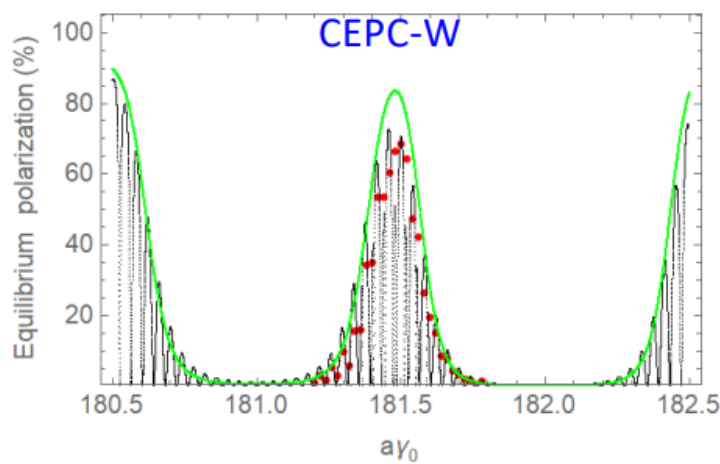
Beam energy (GeV)	$a\gamma_0$	$ \tilde{\omega}_k ^2 (\times 10^{-9})$	$ \tilde{\lambda}_k ^2 (\times 10^{-5})$	$\sigma_\delta (\times 10^{-4})$	ν_z	τ_p (h)	κ	σ
45.6*	103.5	$ \tilde{\omega}_{103} ^2 = 2.7$	$ \tilde{\lambda}_{103} ^2 = 2.2645$	3.77	0.028	252.72	0.03	1.39
		$ \tilde{\omega}_{104} ^2 = 2.8$	$ \tilde{\lambda}_{104} ^2 = 1.7166$					
60.1	136.5	$ \tilde{\omega}_{136} ^2 = 3.7$	$ \tilde{\lambda}_{136} ^2 = 0.8178$	4.96	0.028	63.34	0.20	2.42
		$ \tilde{\omega}_{137} ^2 = 14.7$	$ \tilde{\lambda}_{137} ^2 = 5.5717$					
69.8	158.5	$ \tilde{\omega}_{158} ^2 = 6.6$	$ \tilde{\lambda}_{158} ^2 = 0.9574$	5.77	0.0324	30.00	0.36	2.82
		$ \tilde{\omega}_{159} ^2 = 26.4$	$ \tilde{\lambda}_{159} ^2 = 4.4585$					
80.0*	181.5	$ \tilde{\omega}_{181} ^2 = 14.4$	$ \tilde{\lambda}_{181} ^2 = 4.9118$	6.61	0.0395	15.24	0.52	3.04
		$ \tilde{\omega}_{182} ^2 = 53.3$	$ \tilde{\lambda}_{182} ^2 = 15.6433$					
84.4	191.5	$ \tilde{\omega}_{191} ^2 = 16.3$	$ \tilde{\lambda}_{191} ^2 = 15.5332$	6.97	0.0425	11.65	0.61	3.14
		$ \tilde{\omega}_{192} ^2 = 19.8$	$ \tilde{\lambda}_{192} ^2 = 1.0088$					
90.1	204.5	$ \tilde{\omega}_{204} ^2 = 19.1$	$ \tilde{\lambda}_{204} ^2 = 3.7786$	7.43	0.0467	8.39	0.72	3.25
		$ \tilde{\omega}_{205} ^2 = 43.8$	$ \tilde{\lambda}_{205} ^2 = 0.6403$					
95.4	216.5	$ \tilde{\omega}_{216} ^2 = 15.0$	$ \tilde{\lambda}_{216} ^2 = 6.1547$	7.88	0.0515	6.31	0.80	3.31
		$ \tilde{\omega}_{217} ^2 = 34.8$	$ \tilde{\lambda}_{217} ^2 = 1.2292$					
99.8	226.5	$ \tilde{\omega}_{226} ^2 = 10.5$	$ \tilde{\lambda}_{226} ^2 = 35.2604$	8.24	0.0550	5.03	0.90	3.39
		$ \tilde{\omega}_{227} ^2 = 27.4$	$ \tilde{\lambda}_{227} ^2 = 4.9787$					
109.9	249.5	$ \tilde{\omega}_{249} ^2 = 56.9$	$ \tilde{\lambda}_{249} ^2 = 26.9851$	9.08	0.0585	3.10	1.48	3.87
		$ \tilde{\omega}_{250} ^2 = 41.3$	$ \tilde{\lambda}_{250} ^2 = 46.9963$					
120.1*	272.5	$ \tilde{\omega}_{271} ^2 = 770.4$	$ \tilde{\lambda}_{271} ^2 = 41.6290$	9.90	0.0650	2.03	1.95	4.15
		$ \tilde{\omega}_{272} ^2 = 95.8$	$ \tilde{\lambda}_{272} ^2 = 12.0825$					
		$ \tilde{\omega}_{273} ^2 = 1684.2$	$ \tilde{\lambda}_{273} ^2 = 361.7036$					



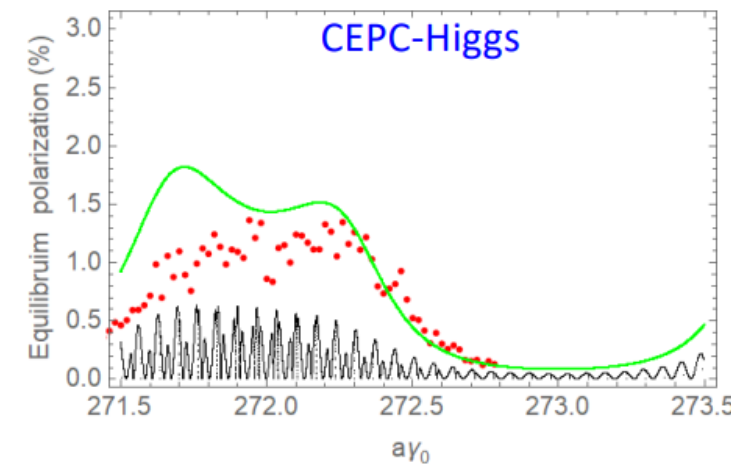
—● Monte Carlo —● Correlated regime —● Uncorrelated regime



● Monte Carlo ● Correlated regime ● Uncorrelated regime



● Monte Carlo ● Correlated regime ● Uncorrelated regime



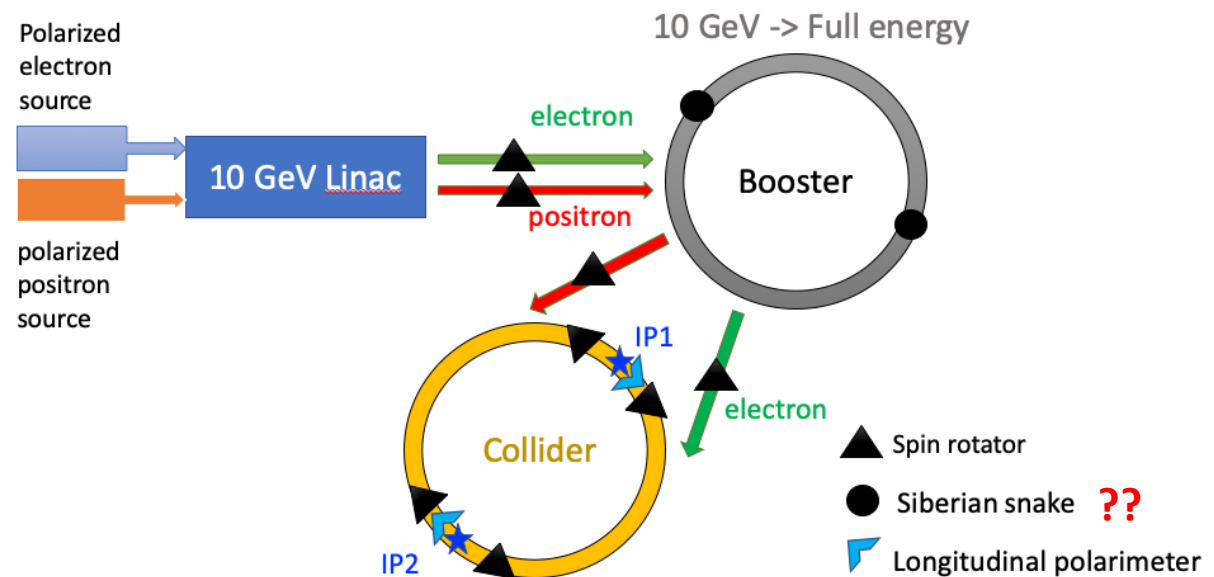
● Monte Carlo ● Correlated regime ● Uncorrelated regime

Summary

- We compared Monte-Carlo simulation of the radiative depolarization versus the two distinct theories that describe the influence of synchrotron oscillations & radiations at ultra-high beam energies.
- The comparison suggests a gradual evolution from the correlated regime to the uncorrelated regime, not clear at the moment. Work urgent is needed to clarify the theory. For example using the Bloch equation[1,2,3], that could merges into these theories at extremes.
- Generation of this study to more comprehensive lattice modeling and more error seeds is foreseen, for better understanding the radiative depolarization mechanisms and establishing correction methods to achieve a high beam polarization @ CEPC.

Tao Chen: Longitudinal polarization @ CEPC

- In the injector: preparation and maintenance of highly polarized e- (e+) beam(s).
 - Polarized source: polarized e- gun (specs defined), polarized e+ source (preliminary study)
 - Booster: polarization maintenance (underway)
 - Transfer lines: ensure the matching of polarization directions (to be studied)
- In the collider ring:
 - spin rotators - > longitudinal polarization[1] (done)
 - ensure $\tau_{DK} \gg \tau_b$, then $P_{avg} \approx P_{inj}$
 - Compton polarimeter[2] (under way)

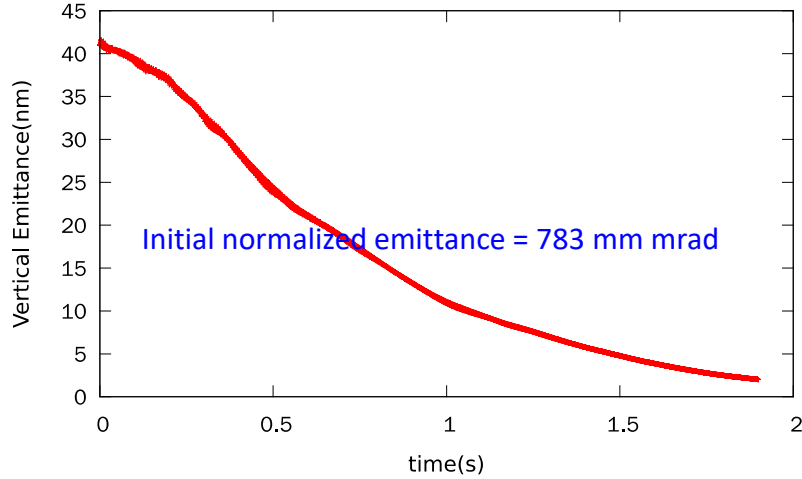


[1] W. H. Xia et al., RDTM (2022) doi: 10.1007/s41605-022-00344-2

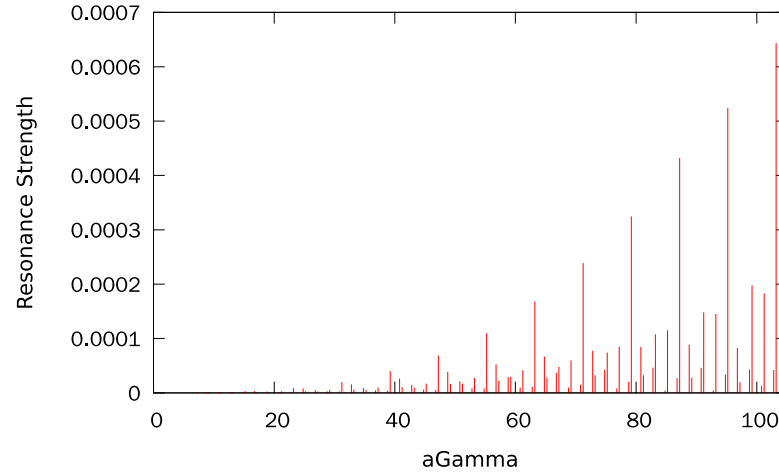
[2] S. H. Chen et al., JINST 17, P08005, (2022)

Tao Chen: Simulation of polarization transmission to 45.6 GeV

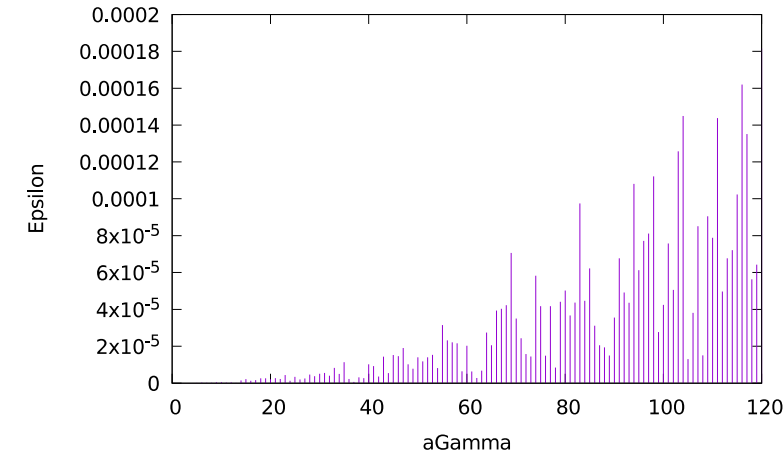
Evolution of vertical rms emittance



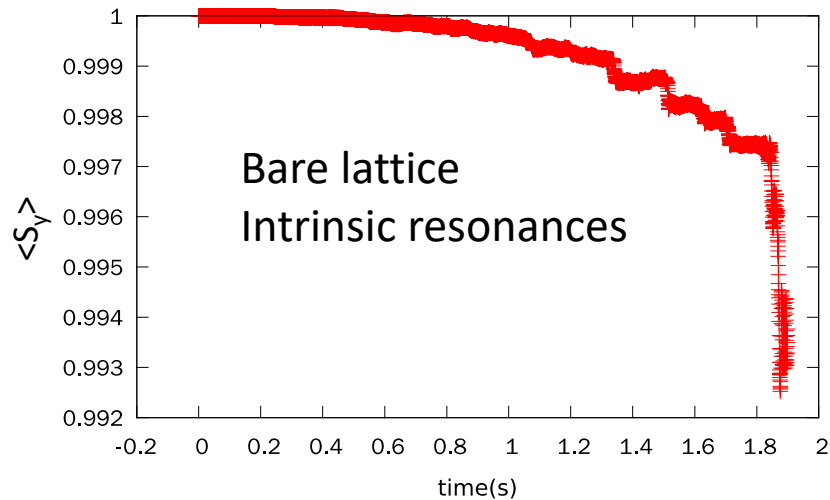
Intrinsic resonance strength scaled with emittance



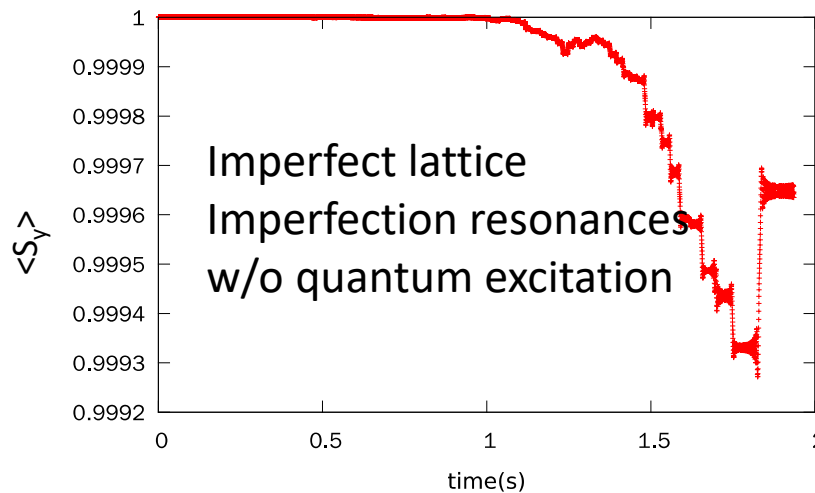
Imperfection resonance strength



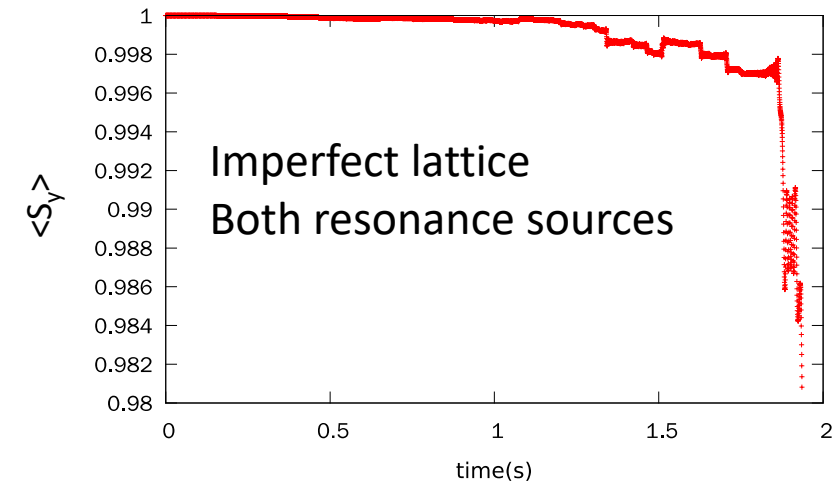
200 particles, initial vertical emittance 40nm



On the closed orbit

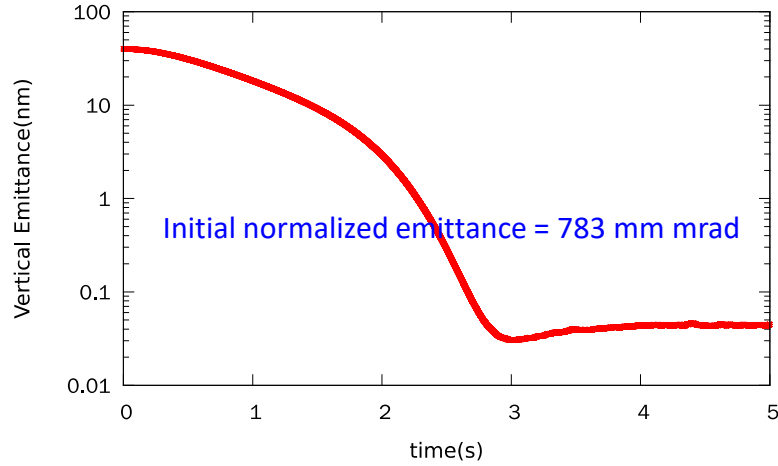


200 particles, initial vertical emittance 40nm

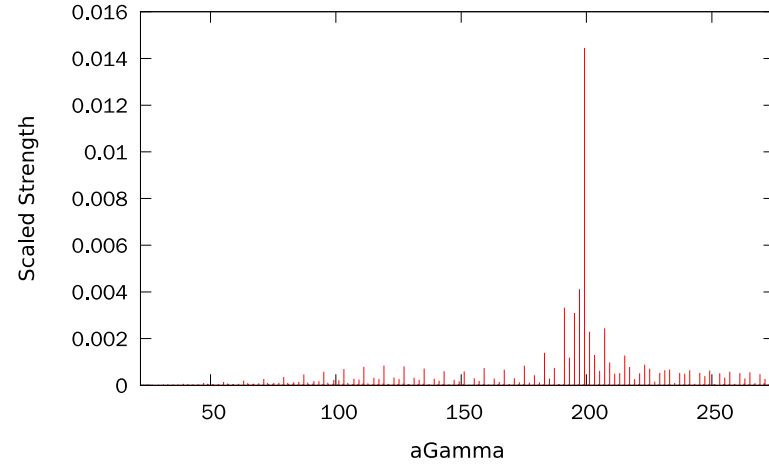


Tao Chen: Simulation of polarization transmission to 120 GeV

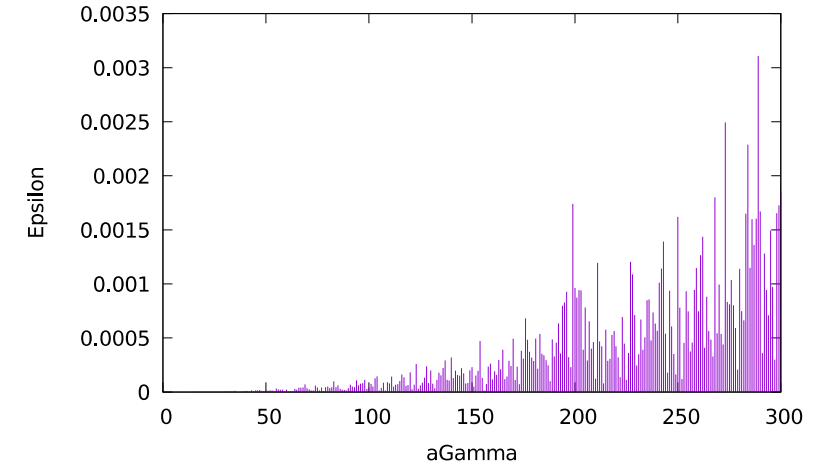
Evolution of vertical rms emittance



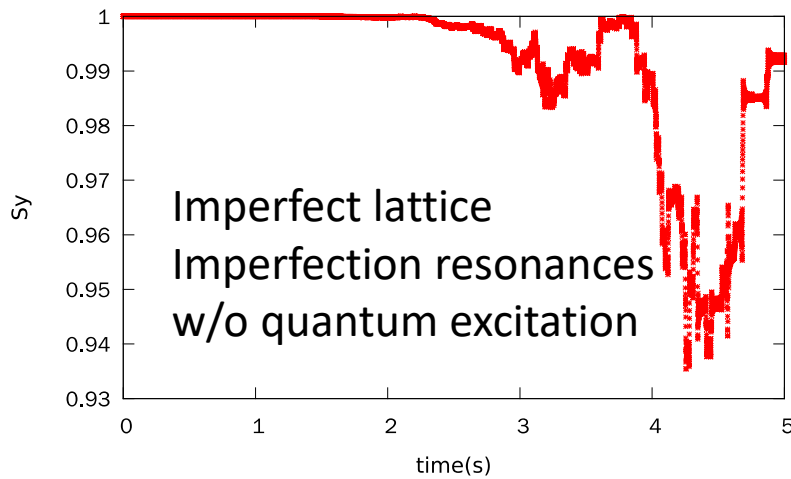
Intrinsic resonance strength scaled with emittance



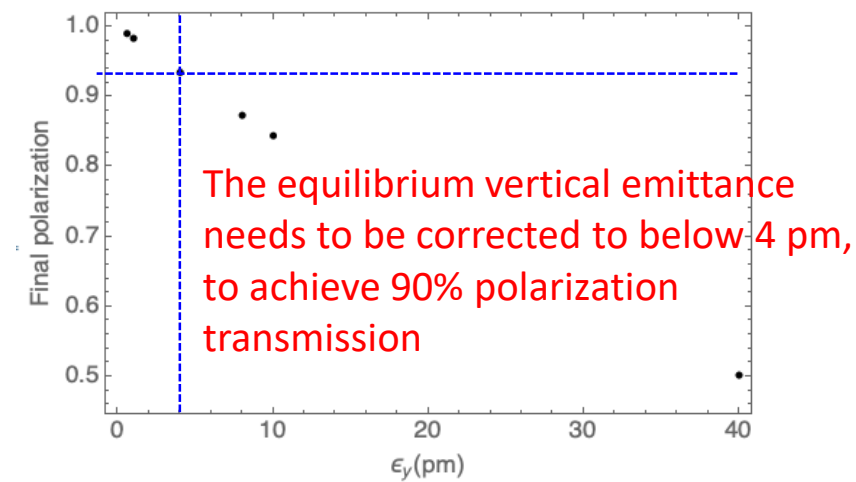
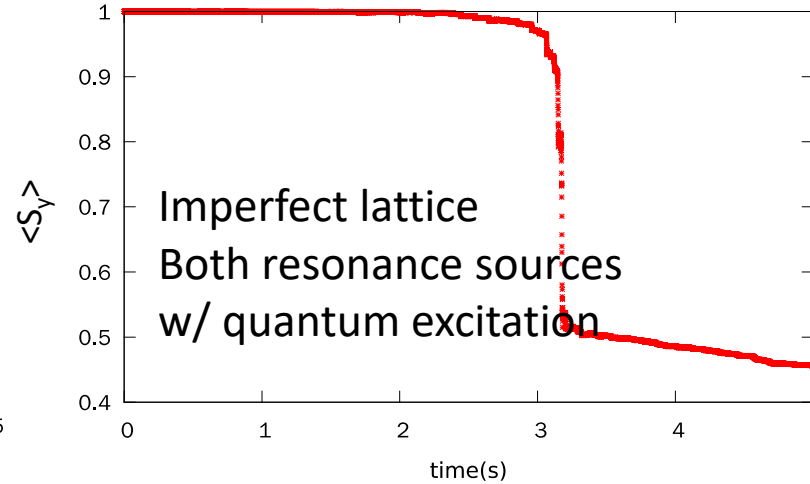
Imperfection resonance strength



On the closed orbit



5000 particles, initial vertical emittance 40nm



Tao Chen: Short summary on polarization maintenance in booster

Findings:

- A large ramping rate of spin precession frequency α , due to the large circumference
- Spin resonances are generally weak, due to the high periodicity & cancellation
- Depolarization is negligible, in the fast crossing regime $\frac{\epsilon}{\sqrt{\alpha}} \ll 0.1$, up to 45.6 GeV
- The strong intrinsic resonance at ~ 87 GeV leads to large depolarization, and hurts the polarization transmission up to 120 GeV, potential mitigations:
 - A new lattice with the first strong intrinsic resonance larger than 120 GeV
 - The above study used the lattice of CDR, In the new design of TDR the condition is satisfied.
 - Control the vertical equilibrium beam emittance to below ~ 4 pm (coupling $\sim 0.1\%$)
- Further research is needed on the tolerance of the highly efficient polarization transmission to the corrected closed-orbit amplitude

Ideas for longitudinal polarization at the Z/W/H/top factory

I.Koop

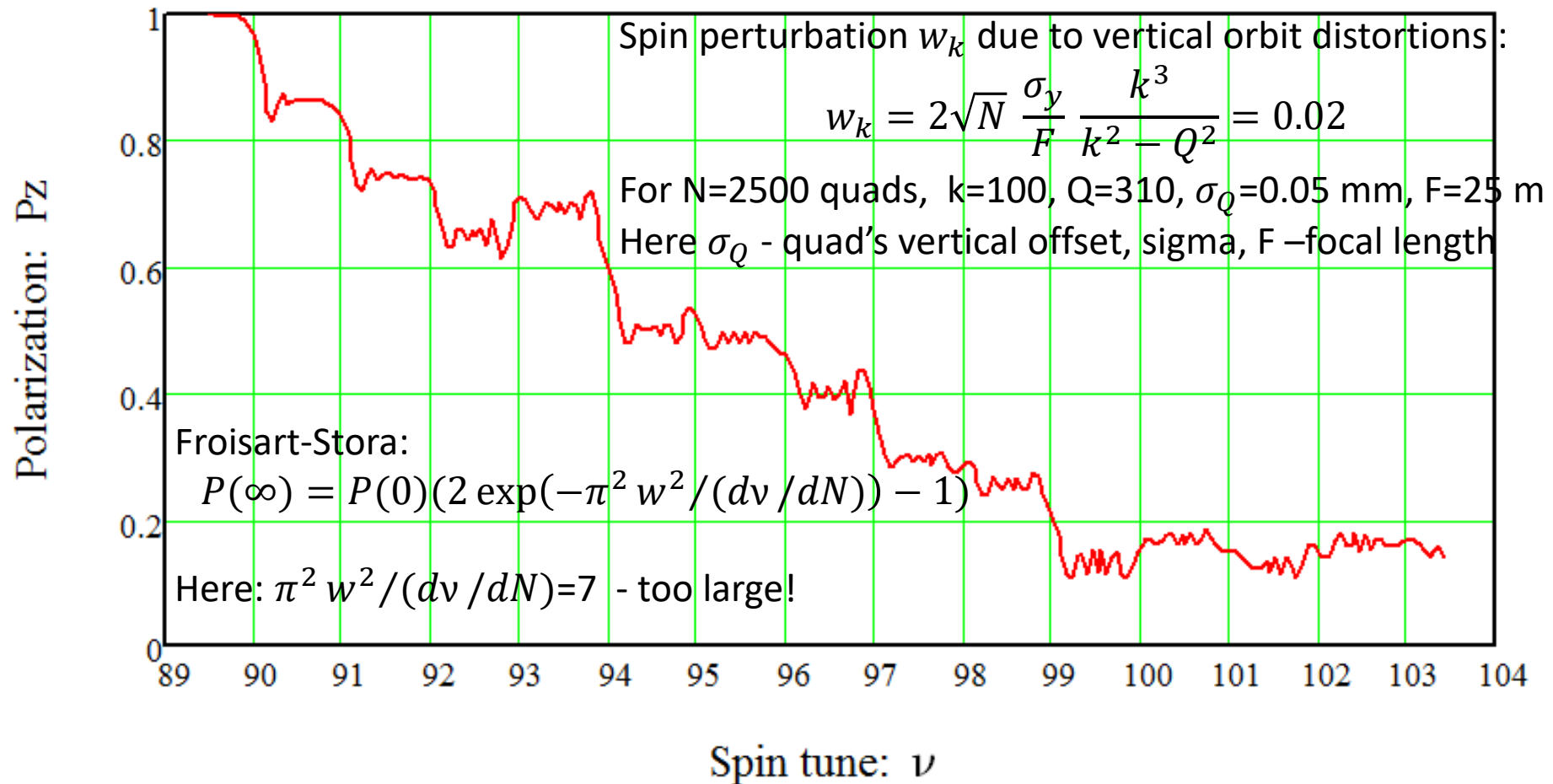
Budker Institute of Nuclear Physics,
Novosibirsk State University,
Novosibirsk State Technical University, 630090,
Novosibirsk, Russia

62nd ICFA Advanced Beam Dynamics Workshop
on High Luminosity Circular e^+e^- Colliders
(eeFACT2018)

HKUST, Hongkong, 24-27 September, 2018

Koop: option-2: Fast acceleration in a booster ring

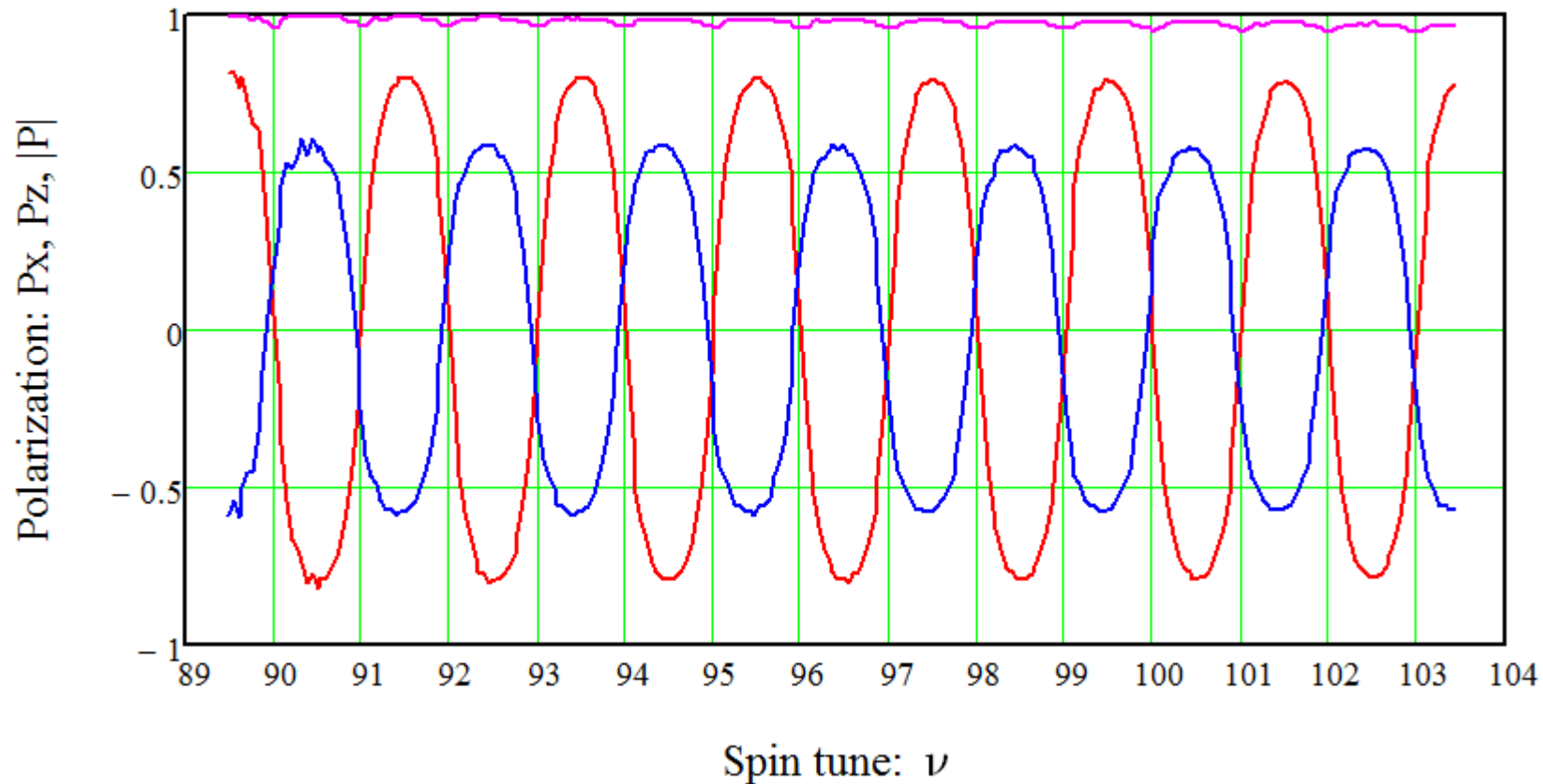
39-45 GeV, $Q_s=0.02$, $\sigma_\delta=0.0004$, $w=0.02$, $dv/dN=0.056$



As one can see, the presented here simulation shows very strong depolarization in the FCC-ee booster synchrotron during acceleration with the nominal ramp rate: 25 GeV/ 0.32 s. Besides, we can expect up to 3 times stronger harmonics due to statistical fluctuations.

Koop: option 3: Adiabatic crossing of integers near Z

45.6 GeV, $Q_s=0.02$, $\sigma_\delta=0.0004$, $w=0.2$, $d\nu/dN=0.056$



Beam polarization is well preserved with the use of single Partial Snake for the acceleration in the FCC-ee booster ring. Polarization loss is only 3%, energy ramp rate 25 GeV/0.32 s. Can use static solenoid with field integral **BL=200 T·m**. Then, at 20 GeV we will have $w=0.5$ (full snake!) and $w=0.22$ at 45.6 GeV. Quads of spin-rotator will ramp to keep $Q_{x,y} = const$.

Koop:

I should check – how closed orbit distortion amplitudes used in my simulations in a toy ring are relevant to a real situation in the fast ramping booster synchrotron...

That strongly affects the parameters of a needed partial snake for adiabatic crossing of integer resonances.

And, probably, fast crossing can work without any problems?
That is a simplest solution! EIC and CEPC rely on such approach!

Spin Polarization Simulations for the Future Circular Collider e^+e^- using BMAD

Yi Wu¹, Félix Carrier², Tatiana Pieloni¹

¹École Polytechnique Fédérale de Lausanne (EPFL)

²The European Organization for Nuclear Research (CERN)

Acknowledgments to Alain Blondel, Desmond Barber, David Sagan, Eliana Gianfelice-Wendt, Tessa Charles, Werner Herr, Léon Van Riesen-Haupt and all colleagues

The logo for EPFL (École Polytechnique Fédérale de Lausanne) consists of the letters 'EPFL' in a bold, red, sans-serif font.

- SITF, the linear spin simulation module in SITROS
- Underlying differences between two codes exist → check step by step

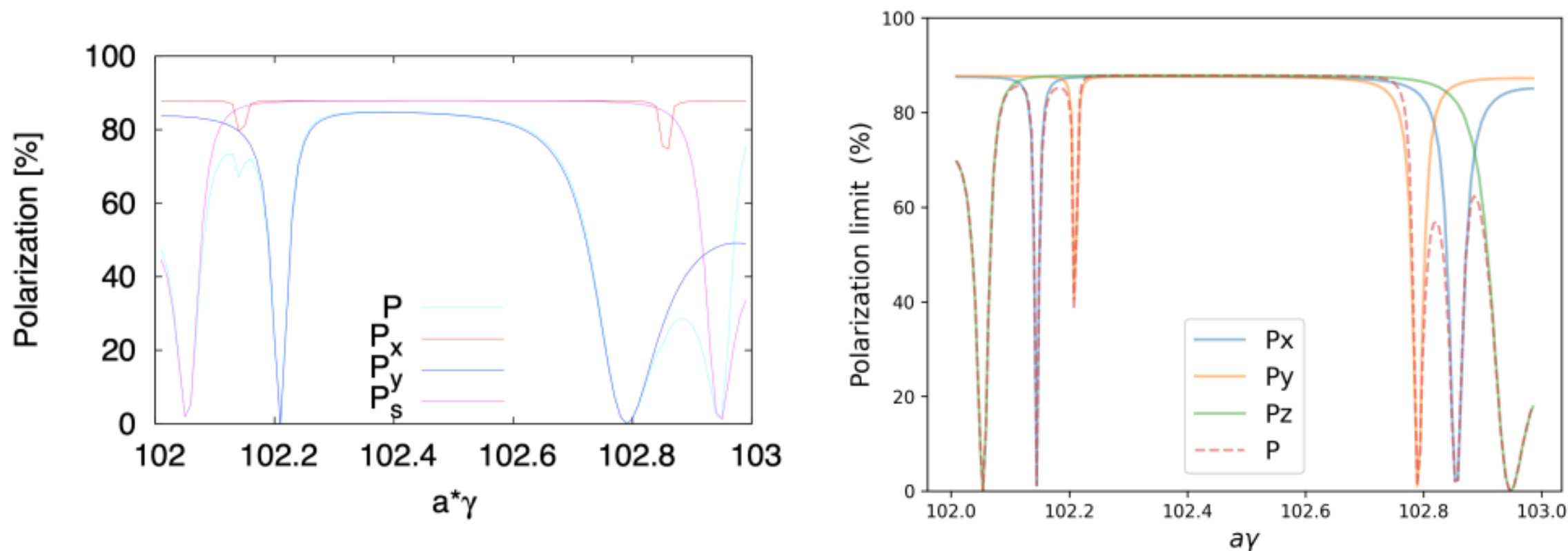


Figure: Energy scan using sequence version 213 seed 13 in SITF (left) and Tao (right)

FCCEe Polarization simulations with SITROS

Content:

- Introduction
- Polarization wigglers
- Simulations at 45 and 80 GeV
- The importance of damping in the 8x8 matrix
- Summary

Eliana GIANFELICE (Fermilab)
EPOL2022, September 22, 2022

Eliana's Summary

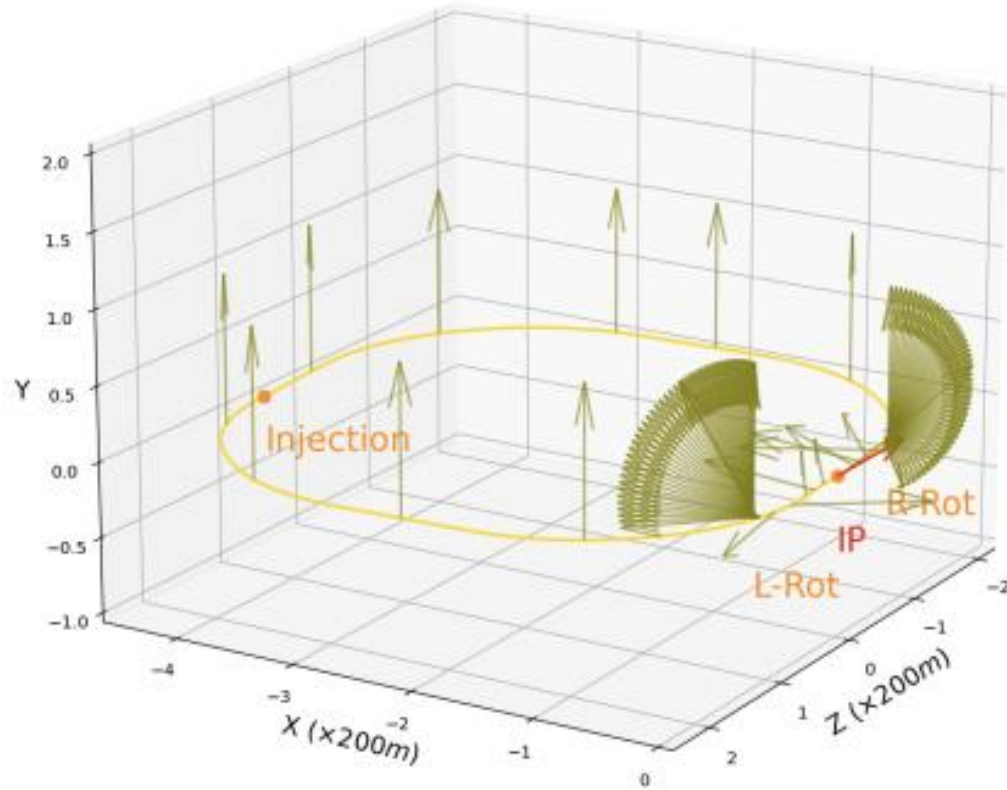
Due to the demanding IR optics design and the machine size, establishing a closed orbit and keeping a stable machine look challenging.

- Beam polarization is obtained “for free” through Sokolov-Ternov effect.
 - At 45 GeV wigglers are required to get $\tau_{10\%} \approx 2-3$ h.
They do not harm polarization.
- P_{∞} depends on how well is the machine aligned/corrected, requirements becoming stricter at high energy.
 - Extremely well corrected orbit/optics is required for a large chromatic machine with $\beta_y^* = 0.8 - 1$ mm as FCC-ee to work and meet required performance.
 - * This benefits also polarization.
- The puzzling small P_y , in particular at 80 GeV, has been likely understood.

Thanks!

Spin Studies with BMAD for a SuperKEKB Polarization Upgrade

1



Yuhao Peng

2022.09.22

- Sextupole pairs located at the Rotator tuning area are turned off because the phase difference between these identical pairs is no longer π (the condition to cancel out the non-linear effects)
- Adjust sextuples in 4 arc section (45 pairs) shown in the picture above to match the original Chromaticity

Ring Parameters Comparison after performing the closed-geometry optimization

29

Machine Parameter	Original Ring	Rot Installed
Tune Q_x	45.530994	45.530994
Tune Q_y	43.580709	43.580709
Chromaticity ξ_x	1.593508	1.593508
Chromaticity ξ_y	1.622865	1.622865
Damping partition J_x	1.000064	0.984216
Damping partition J_y	1.000002	1.005266
Emittance ε_x (m)	4.44061×10^{-9}	4.89628×10^{-9}
Emittance ε_y (m)	5.65367×10^{-13}	3.96631×10^{-12}

Tune and Chromaticity are matched to the original

Conclusion

Acceleration of polarized beams (both – e+ and e-) in a booster – most challenging and most attractive solution. Shall study seriously!

Simulation codes are developing, but need to be faster and become more friendly for users, as MADx and other accelerator codes.