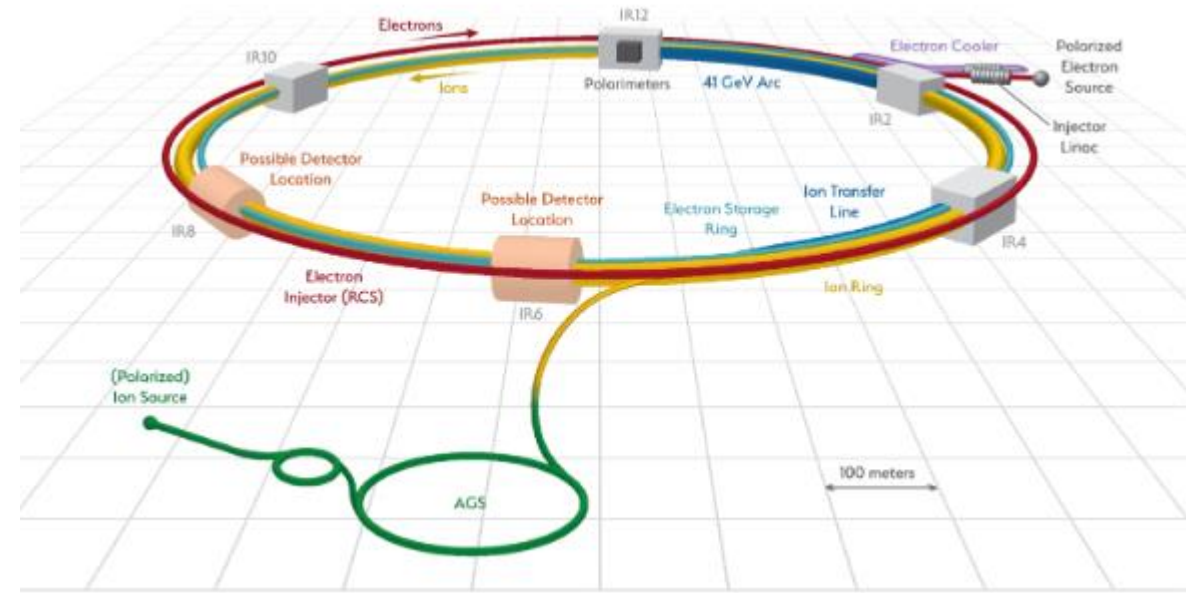
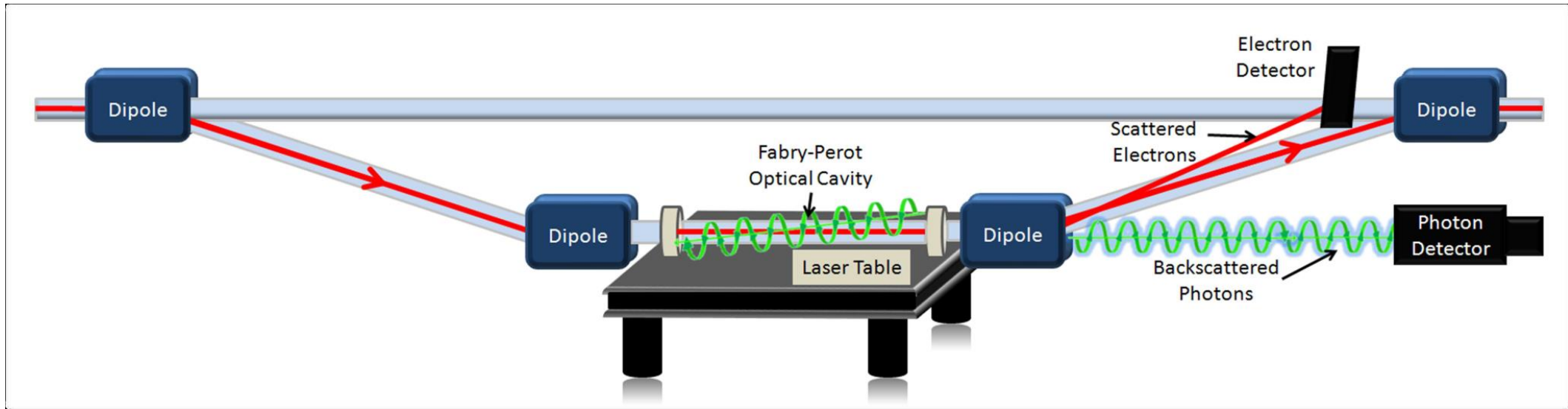


# Compton lasers at JLab and the EIC

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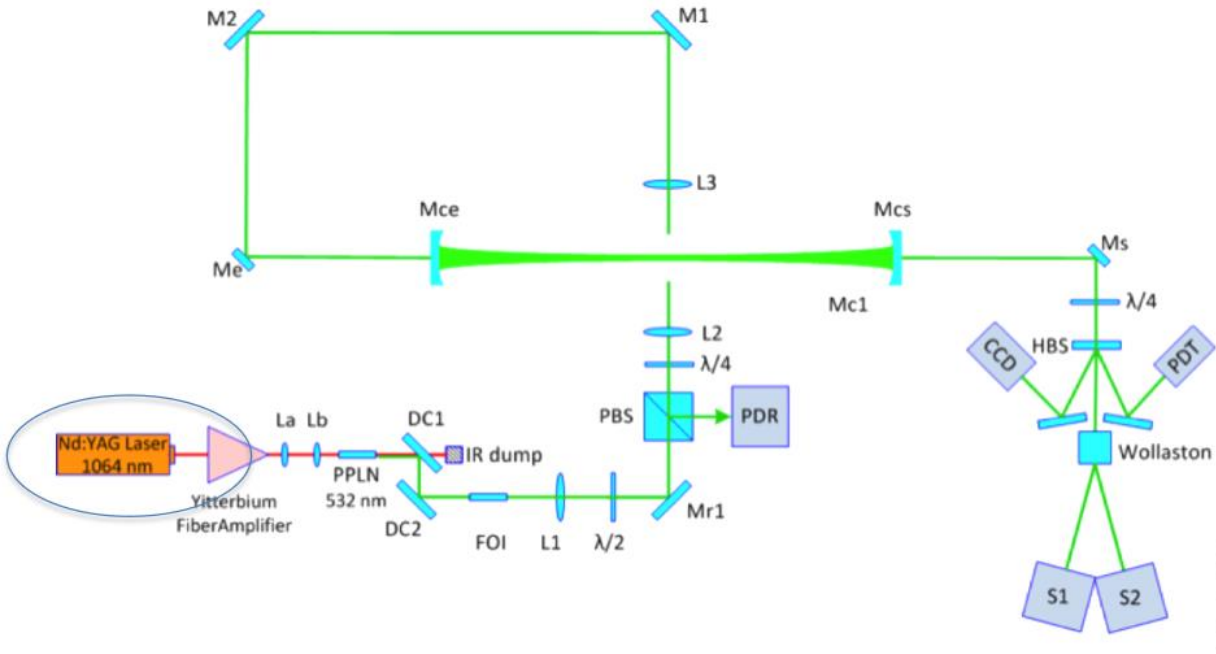


# JLab Compton Layout

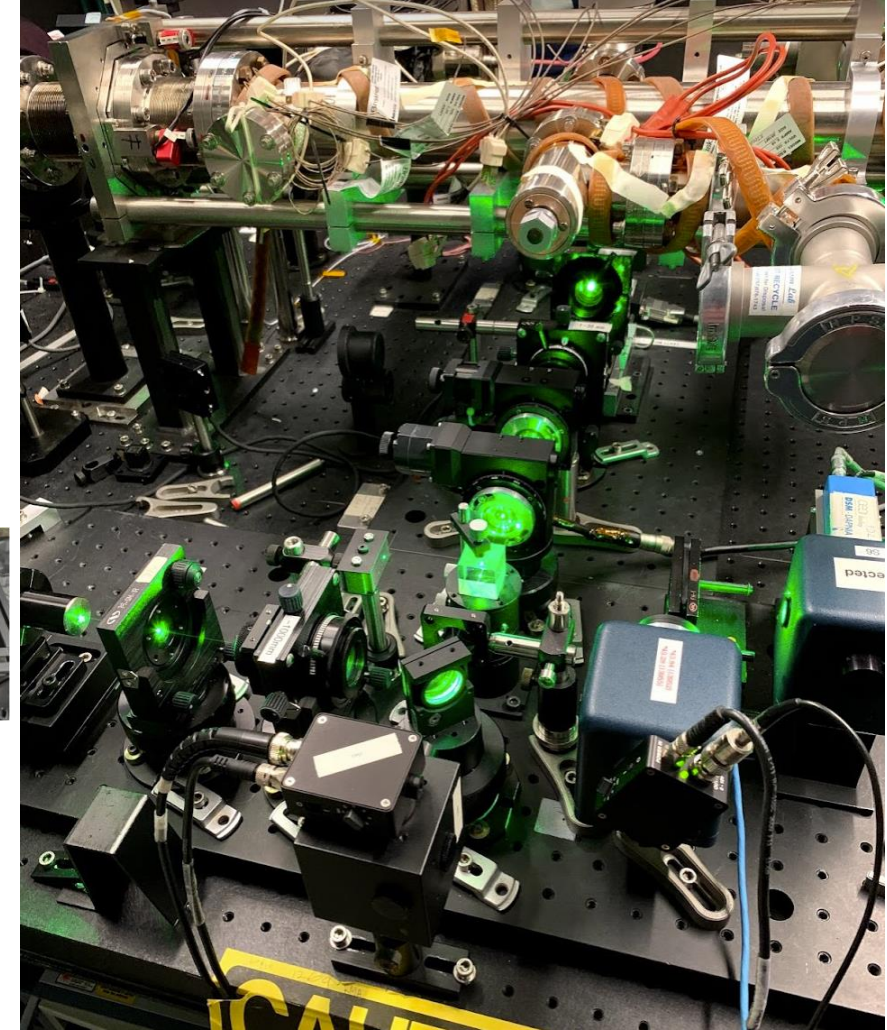


- The Compton setup at Jefferson Lab makes use of a chicane with 4 dipoles to steer the electron beam onto a laser table where interactions with a high power laser beams can occur
- The backscattered photon energy is measured in a calorimeter (crystal depends on the energy of the photons) and the scattered electron is tracked before the last dipole in the chicane to again infer energy loss
- System could in principle be used redundantly but more frequently we ended up using one of the two for the main measurement and used the other for a cross-check

# JLab Compton Layout

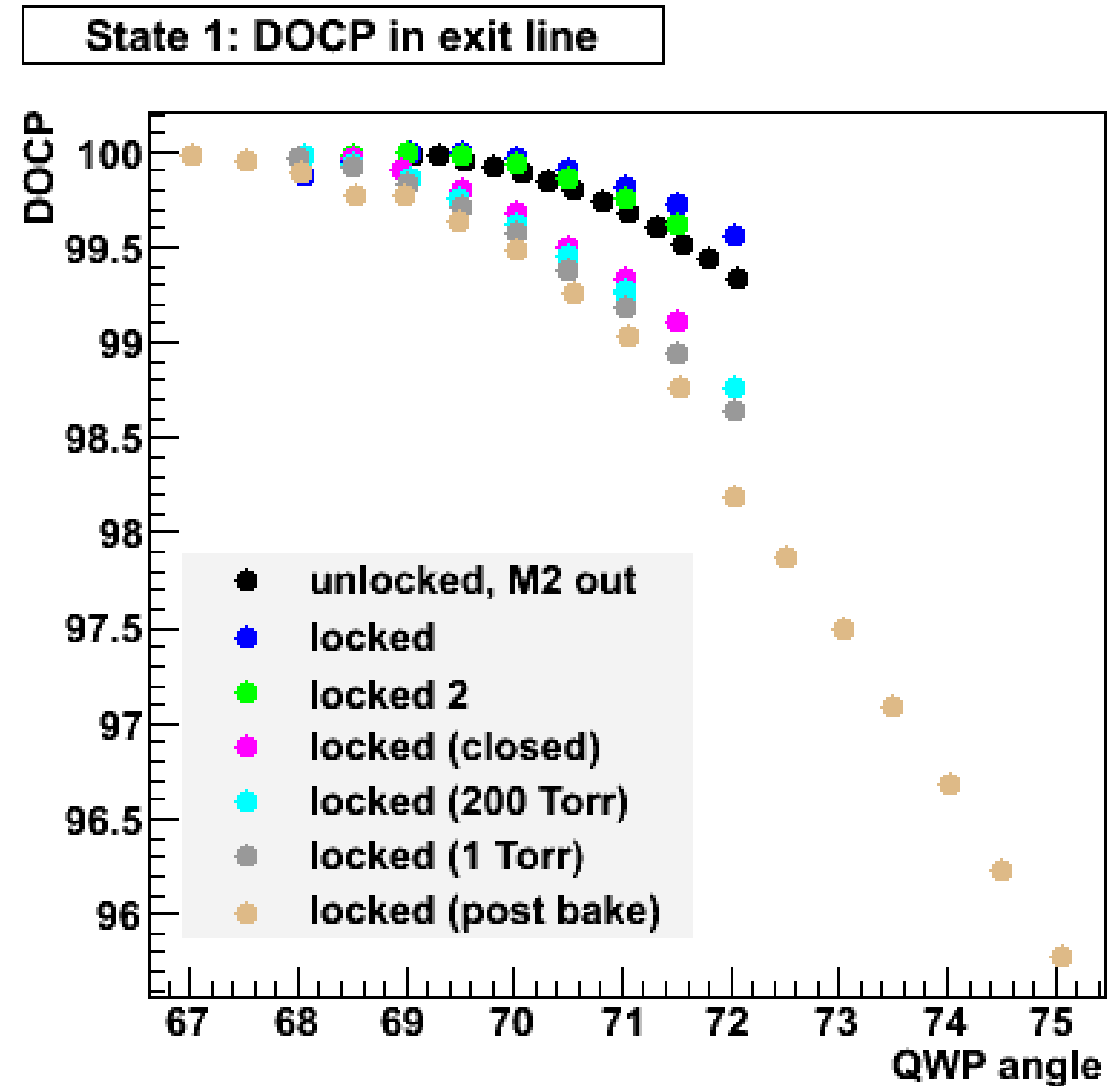


- The Jefferson Lab (Hall A) Compton laser is a CW high power system that enables measurements in integrating mode (multiple interactions per time unit)
- The high power is obtained through the use of Fabry-Perot high gain cavity ( $10^4$ ) cavity which typically leads to powers on the level of a few kW
- A critical aspect to reaching high precision (recent CREX results had 0.44% precision) is the ability to monitor and measure the degree of circular polarization (DOCP) inside the cavity



# DOCP through windows

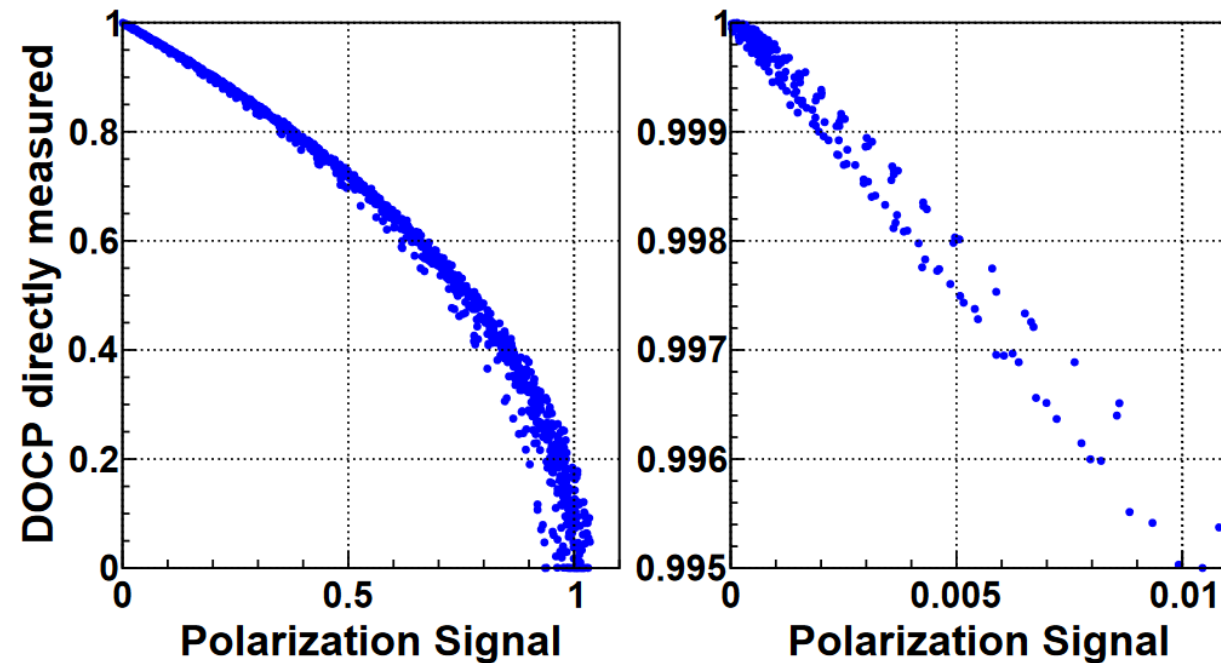
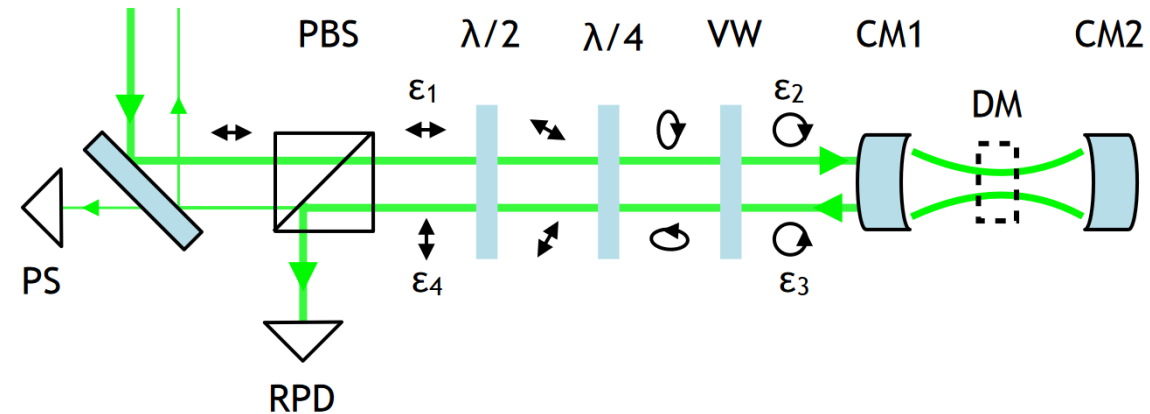
- Tests done with cavity at JLab showed that large differences in the degree of circular polarization can be obtained when straining the windows
- Typically the polarization is monitored through measurements of the transmitted laser light (after the IP)
- The “transfer function” can be measured on the bench but variations (such as tightening bolts or pulling vacuum) change the function making it unusable for the actual data taking



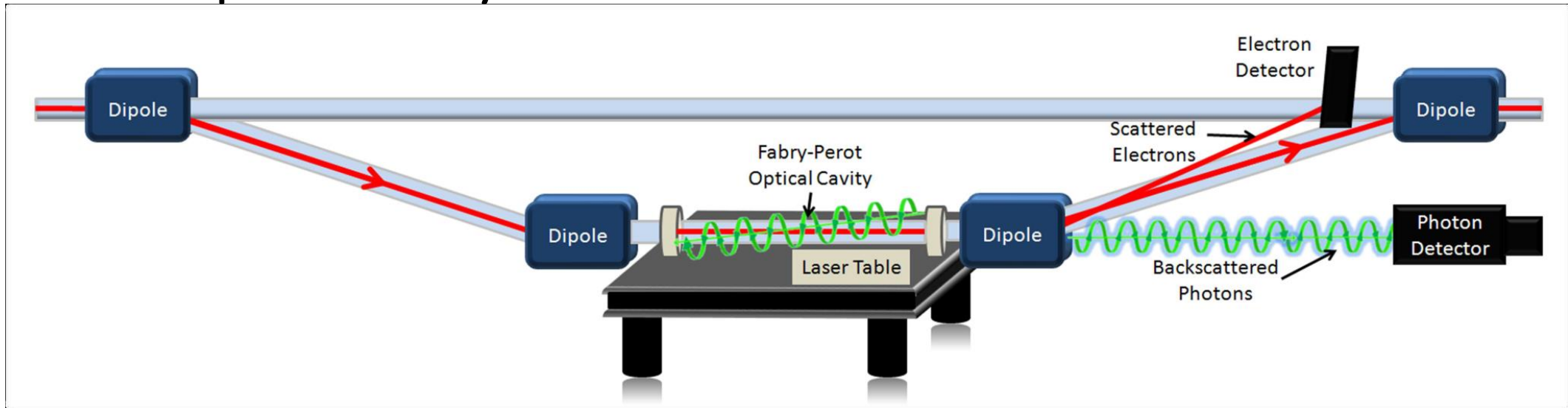


# JLab Compton polarimetry

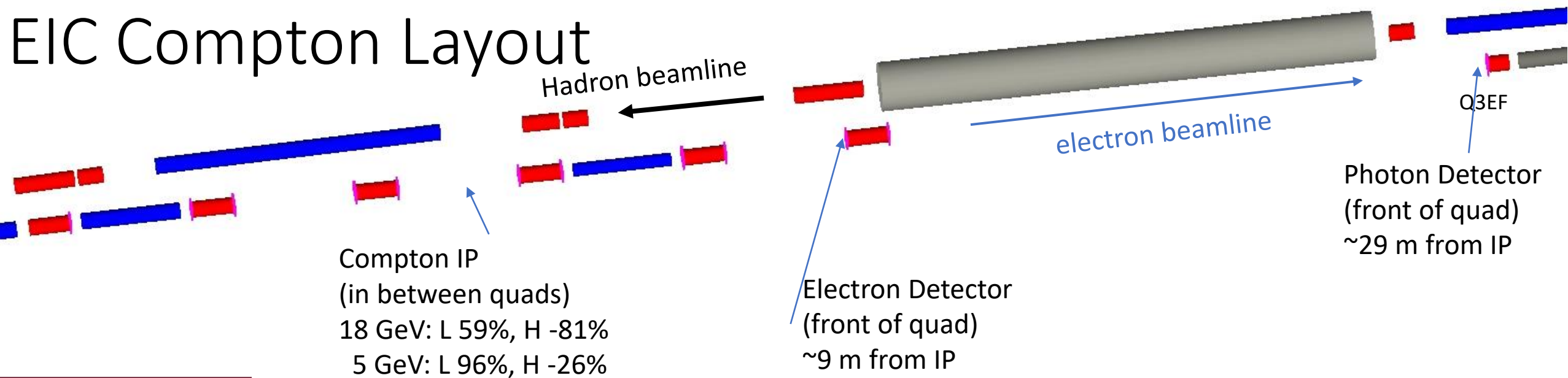
- In order to obtain circular polarization in the cavity one can use the information obtained from the back-reflected light
  - In this case it would be off of mirror M1
- Using the optical reversibility theorem one can relate the amount of light reaching “PS” to the degree of circular polarization inside the cavity
  - M. Dalton and D. Jones showed this to be true in a setup at JLab (previously observed and published by the LPOL group at HERA)
- By performing detailed scans of the half and quarter wave plates one can maximize the circular light at the IP and monitor it throughout the data taking



# JLab Compton Layout



# EIC Compton Layout



# e-Polarimetry requirements for the EIC

## Fast

- At 18 GeV bunches will be replaced every 1-3 min
  - A full polarimetry measurement needs to happen in a shorter time span
- The amount of electrons per bunch is fairly small  $\sim 24$  nC
  - will need bright laser beam to obtain needed luminosity
- A fast polarimeter will allow for faster machine setup

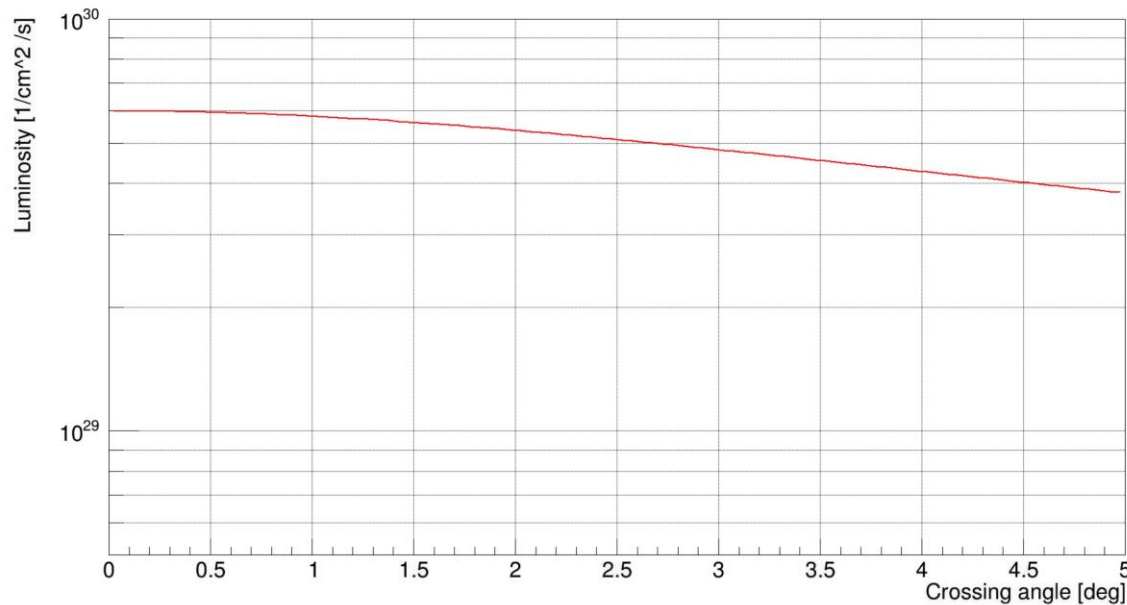
## Precise

- Distance between buckets is  $\sim 10$  ns (@5,10 GeV)
  - bunch by bunch measurement cannot be done with a CW laser without very fast detectors
- For systematic studies we would like to have the ability to either measure a single bunch ( $\sim 78$  kHz) or have interactions with all 1160 (260) bunches at 10 and 5 GeV (18 GeV)
- Backgrounds need to be under control
- Laser polarization needs to be known to a high degree

# Luminosity calculations

$$\mathcal{L} = f_0 N_1 N_2 \frac{\cos(\theta/2)}{2\pi} \frac{1}{\sqrt{(\sigma_{x,1}^2 + \sigma_{x,2}^2)}} \times \frac{1}{\sqrt{(\sigma_{y,1}^2 + \sigma_{y,2}^2) \cos^2(\theta/2) + (\sigma_{z,1}^2 + \sigma_{z,2}^2) \sin^2(\theta/2)}} \quad (1)$$

S. Verdu-Andres (CAD): <https://www.bnl.gov/isd/documents/95396.pdf>



- The dependence of the luminosity of crossing angle needs to take into account the transverse profile of the beam and the length of the pulse
- The estimation on the left is made for a single pulse
- For a 10W 100MHz pulsed laser with a 12ps pulse can provide about  $6 \cdot 10^5$  1/(barn\*s) of luminosity



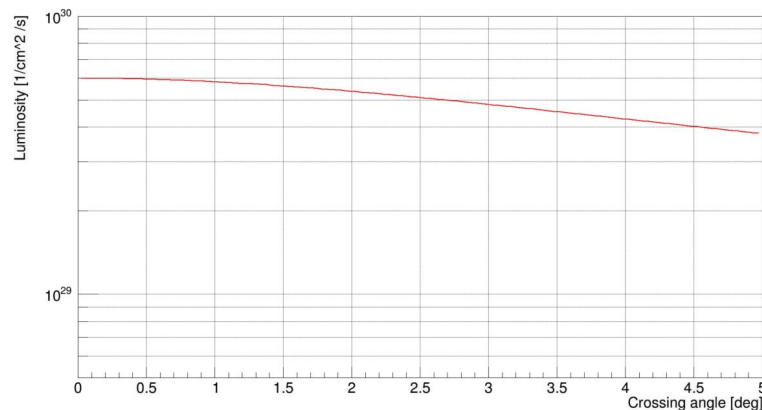
# Luminosity calculations

Configuration	Beam energy [GeV]	Unpol Xsec[barn]	A	A^2	L	1/t(1%)	t[s]	t[min]
laser:532nm, photon	18	0.432	0.072	5.18E-03	1.81E+05	2.93E-02	34	0.57
laser:532nm, electron	18	0.432	0.075	5.63E-03	1.81E+05	3.18E-02	31	0.52
laser:1064nm, photon	18	0.333	0.046	2.12E-03	2.35E+05	1.20E-02	84	1.39
laser:1064nm, electron	18	0.333	0.046	2.12E-03	2.35E+05	1.20E-02	84	1.39

$$N_{Compton} = \frac{\mathcal{L} \cdot \sigma_{unpol}}{f_{beam}}$$

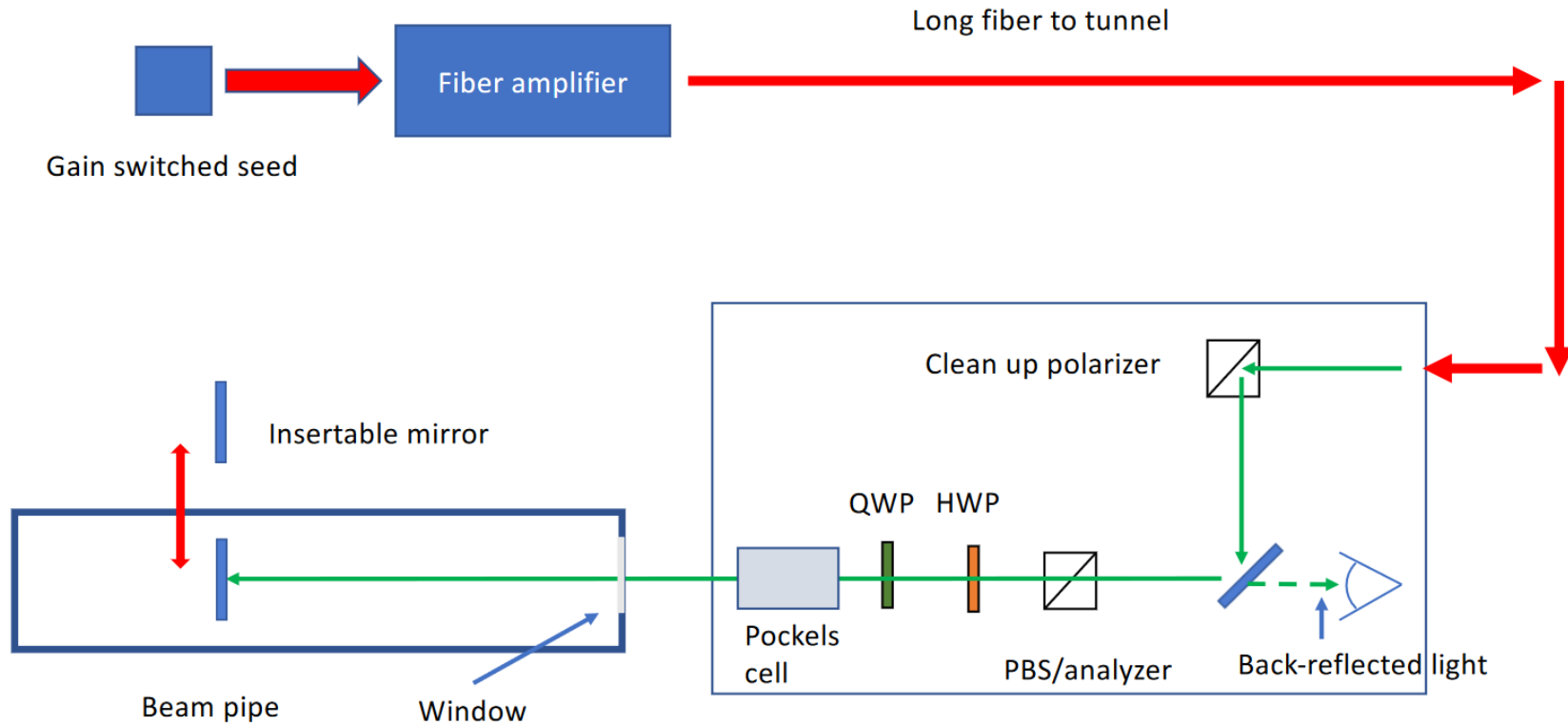
$$t_{meth} = \left( \mathcal{L} \sigma_{Compton} P_e^2 P_\gamma^2 \left( \frac{\Delta P_e}{P_e} \right)^2 A_{meth}^2 \right)^{-1}$$

G. Bardin, et al., Conceptual design report of a Compton polarimeter in cebaf hall a, JLab Internal note.



- Assuming one scattered particle per bunch would allow us to calculate the luminosity needed and a time estimate for how long it would take to reach a 1% statistical precision
- For all configurations envisioned for the EIC (5-18 GeV) the luminosity requirements are on the level of few 1/(barn\*s)
  - Comparing this to the estimate for the 10W laser proves that such a laser will be sufficient
- The times needed to the needed statistics for the signal are on the level 30s at 18 GeV
  - Lower energies are less of a concern due to the longer lived stores
  - This would allow for simultaneous measurement of all bunches

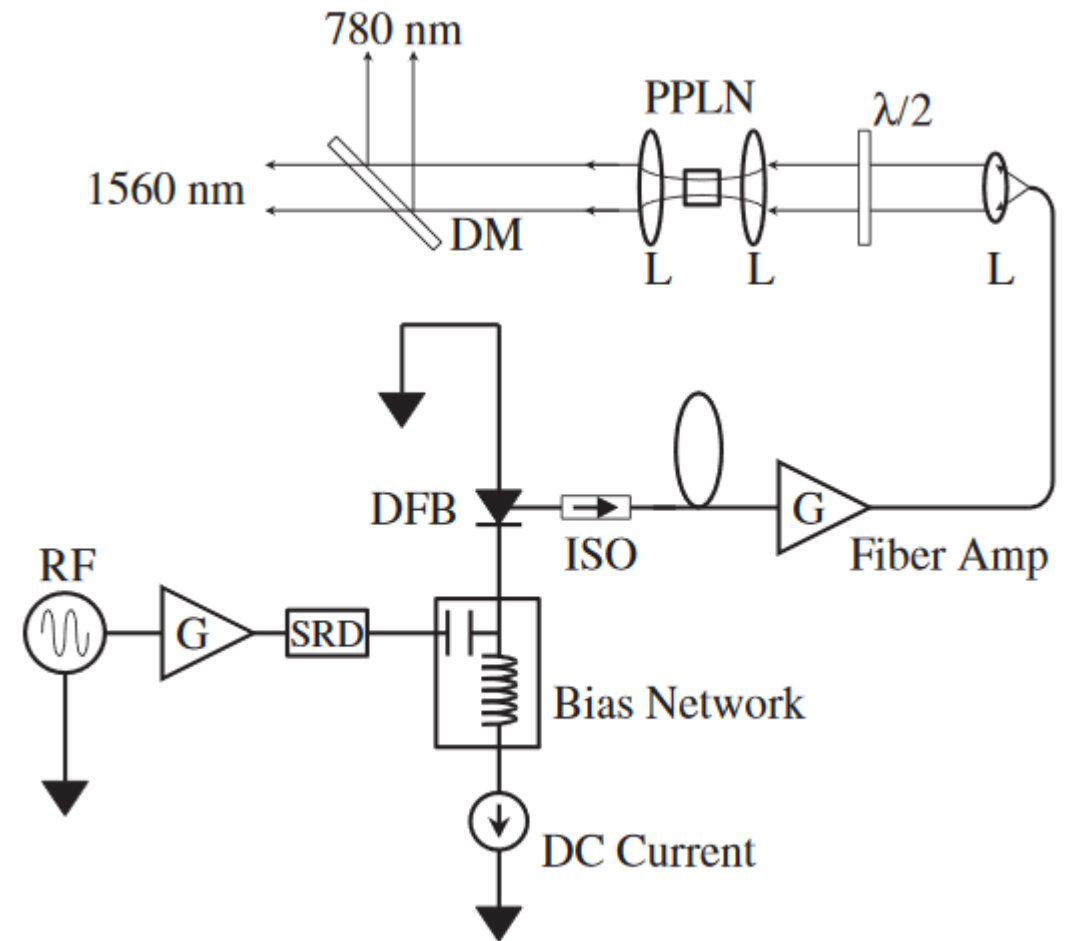
# Current design of EIC laser system



- The initial laser system design uses most of the design features highlighted in the previous Compton polarimeter implementations
  - As was before we need the laser system to be away from potential fatal radiation fields inside the tunnel (we plan to evaluate the use of high power laser fiber)
- The vacuum resident insertable mirror will be needed in order to be able to monitor the DOCP at the interaction point

# Gain switched seed

- The gain switched seed laser design developed at CEBAF for the injector satisfies all the requirements that we discussed so far
  - The RF lock allows us to synchronize to all or specific electron bunches
  - The pulse longitudinal width will be smaller than the electron bunch (allowing us to potentially measure the longitudinal polarization profile)
  - The PPLN or LBO crystal will allow us to frequency double the 1064nm light to 532
- The system has proven to be very reliable and has been adopted by other facilities (such as the Mainz Microtron)

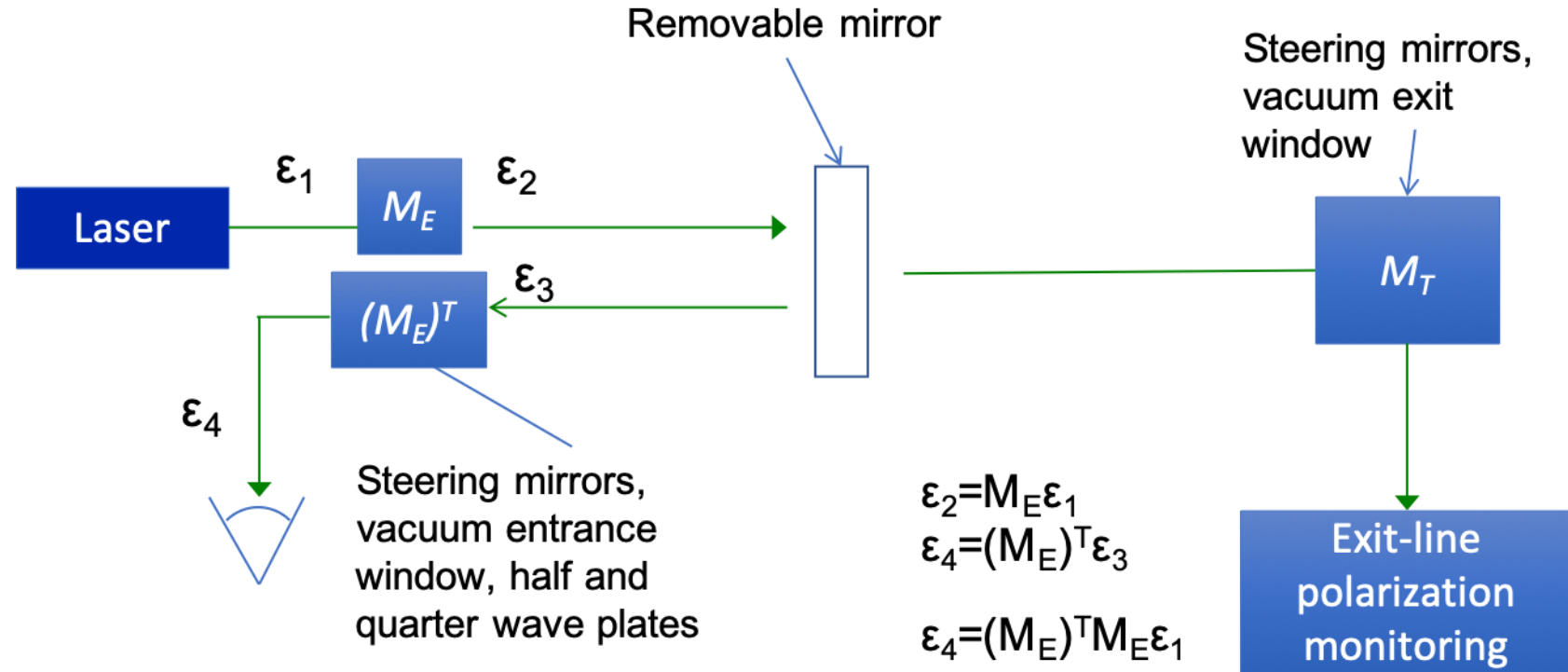
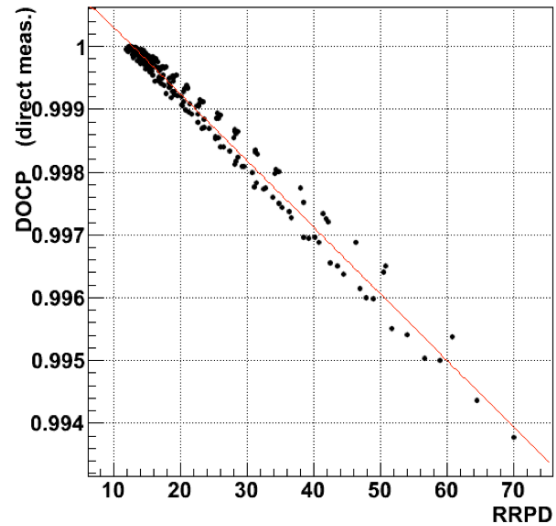


Phys. Rev. ST Accel. Beams **9**, 063501 (2006)

<https://journals.aps.org/prab/pdf/10.1103/PhysRevSTAB.9.063501>

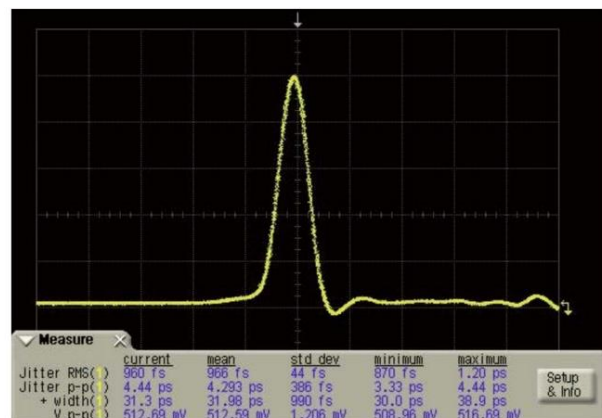
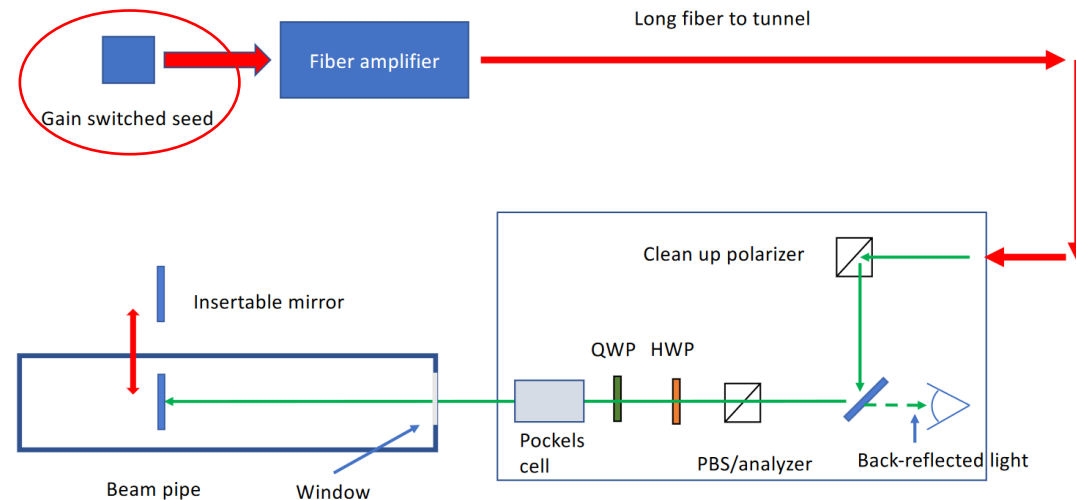
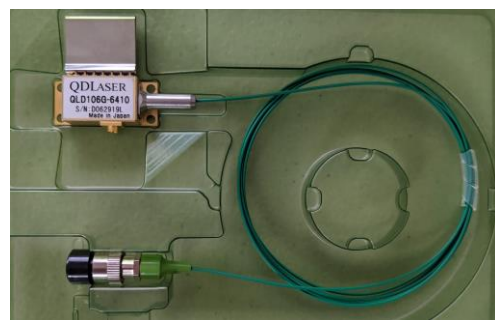
# Current design of EIC laser system

DOCP vs reflected power



- The polarization setup for the EIC Compton will follow the same logical reasoning as the Jefferson Lab measurements

# Work start



- We have started to
- QLD106G-6410 provides 1064nm short pulses (~20ps)
  - Average power 0.1mW @100MHz (-10dBm)
- Electrical pulse generator: 30-230ps pulses, for 1MHz to 1GHz
  - The plan is to have fixed pulse energy at the highest rep rate and use a pulse picker (such as an fiber coupled EOM) to get us the lower frequencies
- A pre-amplifier (YDFA100P) with an output of ~12dBm (getting us to ~16mW) is going to be needed for characterization and input into the fiber amplifier



# Conclusions and further questions

- Due to the specifics of the machine at the EIC we cannot make use of the excellent system that has been employed at JLab over the past decade
- A pulsed laser with an average power of 10W will allow for enough luminosity to be able to have (on average) one interaction per collision
- We plan to have the radiation sensitive electronics (such as amplifiers and seed laser) away from the tunnel to increase longevity of the system
- The DOCP monitoring throughout data taking should be possible by making use of the same scheme developed at JLab
- Questions:
  - How will the laser beam parameters (such as polarization) change as they pass through the very long fiber?

# Backup