Quantum Information with Top Quarks Y. Afik, JRMdN, EPJ Plus 136, 907 (2021) Y. Afik, JRMdN, Quantum 6, 820 (2022) Y. Afik, JRMdN, PRL 130, 221801 (2023)

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Motivation

- Standard Model is a Relativistic Quantum Field Theory = Special Relativity + Quantum Mechanics.
- Quantum Mechanics can be tested via Standard Model.
- Implementation of canonical techniques of Quantum Information \rightarrow Quantum Information Theory at High-Energy Colliders.
- Highest-energy study at the frontier of the known Physics!
- Interest: Genuinely relativistic environment, exotic interactions and symmetries, fundamental nature...



Part I: Quantum Information Theory

Quantum vs. Classical

- Quantum Mechanics: Particles are in superposition of states → Probabilistic description of measurements.
- Classical Mechanics can also describe random outputs using classical probability distributions (noise, experimental variations...).
- Something genuinely quantum? Yes: Wave nature of quantum mechanics!

- Quantum Correlations=Correlations not accounted by classical theories.
 - Quantum Discord
 - Entanglement
 - Steering
 - Bell nonlocality



Quantum State

- Quantum descriptions:
 - **Pure state** \rightarrow Wave function \rightarrow <u>Coherent</u> mixture of quantum states $\rightarrow |\Psi\rangle = \sum_{n} \alpha_{n} \cdot |\phi_{n}\rangle$, α_{n} are <u>amplitudes</u>
 - **Mixed state** \rightarrow Density matrix \rightarrow <u>Incoherent</u> mixture of quantum states $\rightarrow \rho = \sum_{n} p_n \cdot |\phi_n\rangle \langle \phi_n|$, p_n are probabilities
- Density matrix: Most general quantum state.
- Classical descriptions accounted by density matrices.



Qubits

- Qubit: Two-level quantum system $\left|\uparrow\right\rangle,\left|\downarrow\right\rangle\rightarrow$ Most simple quantum system.
- General density matrix (2×2) for 1 qubit \rightarrow 3 parameters B_i :

$$\rho = \frac{1 + \sum_i B_i \sigma^i}{2}$$

• Two qubits \rightarrow Most simple example of quantum correlations.

• General density matrix (4 × 4) for 2 qubits \rightarrow 15 parameters B_i^{\pm}, C_{ij}

$$\rho = \frac{1 + \sum_{i} \left(B_{i}^{+} \sigma^{i} \otimes 1 + B_{i}^{-} 1 \otimes \sigma^{i} \right) + \sum_{i,j} C_{ij} \sigma^{i} \otimes \sigma^{j}}{4}$$



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Quantum Discord

• Classically, two equivalent expressions for mutual information of bipartite system A and B (Alice and Bob):

$$I(A, B) = H(A) + H(B) - H(A, B) = H(A) - H(A|B)$$

$$H(A, B) = -\sum_{x,y} p(x, y) \log_2 p(x, y)$$

$$H(A|B) = \sum_{y} p(y)H(A|B = y)$$

• Quantum mechanics can introduce a "discord" between both expressions: Ollivier, Zurek PRL 88, 017901 (2001)

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$$\mathcal{D}(A,B) \equiv H(B) - H(A,B) + H(A|B) \neq 0$$

- Most basic form of quantum correlations!
- Quantum Discord is asymmetric $\mathcal{D}(A, B) \neq \mathcal{D}(B, A)$

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Quantum Discord: Two qubits

• How do we translate classical into quantum?

Quantum Discord: Two qubits

- How do we translate classical into quantum?
- Shannon entropy \rightarrow Von Neumann entropy ($p_n \ge 0$, ρ eigenvalues)

$$\begin{array}{lll} H(A,B) & \to & H(\rho) = -\sum_n p_n \log_2 p_n \\ H(A) & \to & H(\rho_A), \ H(B) \to H(\rho_B), \ \rho_{A,B} = \operatorname{Tr}_{B,A}\rho \end{array}$$

Conditional probability → Conditional state ρ_{A|B} = One-qubit state after Bob's spin measurement along n̂:

$$H(A|B) = p_{\hat{\mathbf{n}}}H(\rho_{\hat{\mathbf{n}}}) + p_{-\hat{\mathbf{n}}}H(\rho_{-\hat{\mathbf{n}}})$$
$$\rho_{\hat{\mathbf{n}}} = \frac{\prod_{\hat{\mathbf{n}}}^{B}\rho\prod_{\hat{\mathbf{n}}}^{B}}{p_{\hat{\mathbf{n}}}} = \frac{1 + \mathbf{B}_{\hat{\mathbf{n}}}^{+} \cdot \sigma}{2}, \ \mathbf{B}_{\hat{\mathbf{n}}}^{+} = \frac{\mathbf{B}^{+} + \mathbf{C} \cdot \hat{\mathbf{n}}}{1 + \hat{\mathbf{n}} \cdot \mathbf{B}^{-}}, \ \rho_{\hat{\mathbf{n}}} = \frac{1 + \hat{\mathbf{n}} \cdot \mathbf{B}^{-}}{2}$$

● Genuine quantumness → Minimization over all spin directions:

$$\mathcal{D}(A,B) = H(\rho_B) - H(\rho) + \min_{\hat{\mathbf{n}}} p_{\hat{\mathbf{n}}} H(\rho_{\hat{\mathbf{n}}}) + p_{-\hat{\mathbf{n}}} H(\rho_{-\hat{\mathbf{n}}}) \neq 0$$

Entanglement

- Entanglement: Most genuine feature of Quantum Mechanics. Key resource for quantum technologies.
- Separability: $\rho = \sum_{n} p_{n} \rho_{n}^{a} \otimes \rho_{n}^{b}, \sum_{n} p_{n} = 1, \ p_{n} \ge 0$
- Classically correlated state in $\mathcal{H} \rightarrow$ Separable.
- Non-separability=Entanglement → Non-classical state.



Separable



Non-Separable

R. F. Werner, PRA 40, 4277 (1989)

Entanglement: Two qubits

• Two qubits: Separability=Positive P-representation $P(\mathbf{n}_A, \mathbf{n}_B) \ge 0$:

$$\rho = \int \mathrm{d}\Omega_{A} \mathrm{d}\Omega_{B} P(\mathbf{n}_{A}, \mathbf{n}_{B}) |\mathbf{n}_{A}\mathbf{n}_{B}\rangle \langle \mathbf{n}_{A}\mathbf{n}_{B}|, \ \int \mathrm{d}\Omega_{A} \mathrm{d}\Omega_{B} P(\mathbf{n}_{A}, \mathbf{n}_{B}) = 1$$

- Classical spins pointing at directions **n**_A, **n**_B!
- Separability=Purely classical correlations

$$C_{ij} = \langle \sigma^i \otimes \sigma^j \rangle = \int \mathrm{d}\Omega_A \mathrm{d}\Omega_B \ P(\mathbf{n}_A, \mathbf{n}_B) n_A^i n_B^j$$

● Entanglement=NO probability distribution → Genuine non-classical!



Steering: Two qubits

- Measurements of Bob can "steer" quantum state of Alice.
- Steering: Original conception of Schrödinger of EPR paradox → Only well-defined in 2007! (Wiseman, Jones, Doherty, PRL 98, 140402 (2007))
- Alice post-measurement state described by local-hidden states:

$$\tilde{\rho}_{\hat{\mathbf{n}}} = \Pi_{\hat{\mathbf{n}}}^{B} \rho \Pi_{\hat{\mathbf{n}}}^{B} = \int d\lambda \ p(1|\hat{\mathbf{n}}\lambda) p(\lambda) \rho_{B}(\lambda)$$

• If not, quantum state is steerable.



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Steering: Two qubits

• Alice post-measurement state: same as for quantum discord.

$$\rho_{\hat{\mathbf{n}}} = \frac{\tilde{\rho}_{\hat{\mathbf{n}}}}{\mathrm{Tr}\tilde{\rho}_{\hat{\mathbf{n}}}} = \frac{1 + \mathbf{B}_{\hat{\mathbf{n}}}^+ \cdot \sigma}{2}, \ \mathbf{B}_{\hat{\mathbf{n}}}^+ = \frac{\mathbf{B}^+ + \mathbf{C} \cdot \hat{\mathbf{n}}}{1 + \hat{\mathbf{n}} \cdot \mathbf{B}^-}$$

- Set of conditional polarizations $B_{\hat{n}}^+$ describes an ellipsoid.
- Steering ellipsoid: Fundamental QI object, containing all information about the system.
- Similar for Bob \rightarrow Steering: also asymmetric between Alice and Bob.



Jevtic, Pusey, Jennings, Rudolph PRL 113, 020402 (2014)

Bell inequality: Two qubits

• Local realism: Joint Alice and Bob measurements M_A , M_B accounted by local hidden-variable model

$$p(a,b|M_AM_B) = \int d\lambda \ p(a|M_A\lambda)p(b|M_B\lambda)p(\lambda)$$

• Local realism holds if Bell inequality is satisfied. Two qubits \rightarrow CHSH inequality ($\mathbf{a}_i, \mathbf{b}_i$ spin axes of measurements M_A, M_B)

$$|\mathbf{a}_{1}^{\mathrm{T}}\mathbf{C}(\mathbf{b}_{1}-\mathbf{b}_{2})+\mathbf{a}_{2}^{\mathrm{T}}\mathbf{C}(\mathbf{b}_{1}+\mathbf{b}_{2})|\leq 2$$

Stronger condition than entanglement → "Spooky action at distance"



Hierarchy of Quantum Correlations

- Steering and Discord can be asymmetric between Alice and Bob.
- Bell Nonlocality and Entanglement are always symmetric.
- Quantum Hierarchy:

 $\textit{Bell Nonlocality} \subset \textit{Steering} \subset \textit{Entanglement} \subset \textit{Discord}$



Quantum Tomography: Two qubits

- **Quantum Tomography**: Reconstruction of quantum state from measurement of a set of observables.
- Quantum tomography \rightarrow Measurement of ALL quantum correlations.
- Most general density matrices for 1, 2 qubits:

$$\rho = \frac{1 + \sum_{i} B_{i} \sigma^{i}}{2}, \ \rho = \frac{1 + \sum_{i} \left(B_{i}^{+} \sigma^{i} + B_{i}^{-} \bar{\sigma}^{i}\right) + \sum_{i,j} C_{ij} \sigma^{i} \bar{\sigma}^{j}}{4}$$

- One-qubit quantum tomography=Measurement of 3 parameters, polarization vector **B**: $B_i = \langle \sigma^i \rangle$
- Two-qubit quantum tomography=Measurement of 15 parameters, polarization vectors B[±] and correlation matrix C:

$$B_{i}^{+} = \langle \sigma^{i} \rangle \,, \,\, B_{i}^{-} = \langle \bar{\sigma}^{i} \rangle \,, \,\, C_{ij} = \langle \sigma^{i} \bar{\sigma}^{j} \rangle$$



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Part II: Top Quark Physics

Who Top Quarks?

- Top quark is the most massive fundamental particle known to exist $(m_t c^2 \approx 173 \text{ GeV}).$
- First discovered by the D0 and CDF collaborations at the Tevatron in 1995.
- Top quarks produced in top-antitop $(t\bar{t})$ pairs through QCD or Electroweak processes.



Why Top Quarks?

- Large Width $\Gamma_t \sim 1 \text{ GeV} \rightarrow \text{Very short lifetime } \tau = 1/\Gamma_t \sim 10^{-25} \text{s}$
- Tops decay before
 - **)** Hadronisation $\sim 10^{-23} \mathrm{s.}$
 - Spin-decorrelation $\sim 10^{-21} {
 m s.}$
- \rightarrow NO DECOHERENCE OR RANDOMIZATION!
- Rotational invariance in $t\bar{t}$ rest frames $\rightarrow t\bar{t}$ spins measured from decay products.
- Measurements by D0 and CDF (Tevatron), ATLAS and CMS (LHC)
 → Well-established technique!



Top pair kinematics

• $t\bar{t}$ pair kinematically described by invariant mass $M_{t\bar{t}}$ and top direction \hat{k} in c.m. frame

$$egin{array}{rcl} k_t^{\mu} &=& (k_t^0, {f k}), k_{ar t}^{\mu} = (k_t^0, -{f k}) \ M_{tar t}^2 &\equiv& s_{tar t} \equiv (k_t + k_{ar t})^2 \end{array}$$

• Invariant mass is simply related to top c. m. velocity β

$$M_{t\bar{t}} = rac{2m_t}{\sqrt{1-\beta^2}}
ightarrow eta = 0
ightarrow M_{t\bar{t}} = 2m_t$$

• Threshold production: $M_{t\bar{t}} = 2m_t \approx 346 \text{ GeV}$



LO QCD Elementary Process

- Illustrative example: QCD analytical LO calculation.
 - Analytical results.
 - NLO corrections are small.
 - Building blocks of actual high-energy processes.
- Most elementary QCD processes:

$$q + \bar{q} \rightarrow t + \bar{t}, \ q = u, d...$$

 $g + g \rightarrow t + \bar{t}$

• Each initial state $I = q\bar{q}, gg$ gives rise to quantum state $\rho'(M_{t\bar{t}}, \hat{k})$



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LO QCD Realistic

- No free quarks or gluons → Hadrons: Bound states of quarks and gluons (partons).
- LHC, Tevatron: pp, $p\bar{p}$ collisions at high c.m. energies \sqrt{s} .

 $p + p \rightarrow \dots \rightarrow t + \overline{t}$ LHC $p + \overline{p} \rightarrow \dots \rightarrow t + \overline{t}$ Tevatron

• Quantum state depends now on c.m. energy \sqrt{s} :

$$\rho(M_{t\bar{t}},\hat{k}) = \sum_{I=q\bar{q},gg} w_I(M_{t\bar{t}},\sqrt{s})\rho'(M_{t\bar{t}},\hat{k})$$

• Total QCD process: *Incoherent* sum of elementary QCD processes with probability *w*₁.

• QCD Input: $w_I(M_{t\bar{t}}, \sqrt{s}), \rho^I(M_{t\bar{t}}, \hat{k}) \rightarrow QI$ Output: Textbook problem of *convex sum* of quantum states!





Part III: Quantum Tops

$t\bar{t}$ Quantum Correlations

- Quantum state $\rho(M_{t\bar{t}}, \hat{k})$: Function of scattering angle Θ and $M_{t\bar{t}}$.
- Two main regions of quantumness:
 - High- p_T for both $q\bar{q}$ and gg (spin triplet)
 - Threshold for gg (spin singlet).
- Colorbar: Discord.
- Solid, dashed-dotted, dashed: Boundaries of Entanglement, Steering, Bell Nonlocality → Hierarchy!
- a) $gg \rightarrow t\bar{t}$
- b) $q\bar{q} \rightarrow t\bar{t}$
- c) Run 2 LHC $\sqrt{s} = 13 \text{ TeV}$
- d) Tevatron $\sqrt{s} = 1.96 \text{ TeV}$



Total Quantum State

- Realistic measurement: Average over many different processes.
- Total quantum state: Events in window $[2m_t, M_{t\bar{t}}]$

$$\rho(M_{t\bar{t}}) \equiv \frac{1}{\sigma(M_{t\bar{t}})} \int_{2m_t}^{M_{t\bar{t}}} \mathrm{d}M \int \mathrm{d}\Omega \, \frac{\mathrm{d}\sigma}{\mathrm{d}M \mathrm{d}\Omega} \rho(M,\hat{k})$$

- Intuitively: Total quantum state=Sum of $t\bar{t}$ quantum states weighted with the differential cross-section.
- Rotational invariance around beam axis → Correlation matrix diagonal in beam basis C_{ij} = C_iδ_{ij}, C_x = C_y = C_⊥ → 2D dependence on C_⊥, C_z.
- Green: LHC. Orange: Tevatron.

• Cross:
$$\beta = 0$$
; Circle: $\beta = 1$.



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Part IV: Experimental Analysis

Top pair Quantum Tomography

- $\rho(M_{t\bar{t}}) \rightarrow \text{Two qubit quantum state} \rightarrow \text{Quantum tomography} = \text{Measurement of spin polarizations and spin correlations.}$
- Spin polarizations \mathbf{B}^{\pm} and spin correlation matrix \mathbf{C} extracted from cross-section $\sigma_{\ell\bar{\ell}}$ of dileptonic decay

$$\frac{1}{\sigma_{\ell\bar{\ell}}} \frac{\mathrm{d}\sigma_{\ell\bar{\ell}}}{\mathrm{d}\Omega_{+}\mathrm{d}\Omega_{-}} = \frac{1}{(4\pi)^{2}} \left[1 + \mathbf{B}^{+} \cdot \hat{\ell}_{+} - \mathbf{B}^{-} \cdot \hat{\ell}_{-} - \hat{\ell}_{+} \cdot \mathbf{C} \cdot \hat{\ell}_{-} \right]$$

• $\hat{\ell}_{\pm}$: lepton directions in each top (antitop) rest frames.



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Entanglement in $t\bar{t}$ production at LHC $\sqrt{s} = 13$ TeV

- Entanglement witness: $W = D + 1/3 < 0, D \equiv \text{tr } \mathbf{C}/3 \rightarrow$ Entanglement only close to threshold.
- *D* directly measurable from decay cross-sections:

$$\frac{1}{\sigma}\frac{d\sigma}{d\cos\varphi}=\frac{1}{2}(1-D\cos\varphi)$$

- Entanglement detection from one single magnitude! → No need for Quantum Tomography!
- High-statistical significance!
- Entanglement also at high-p_T: Fabbrichesi, Floreanini, Panizzo, PRL 127, 161801 (2021), Severi, Boschi, Maltoni, Sioli, EPJC 82, 285 (2022)









Discord and Steering

Normalized dileptonic cross-section → Angular probability distribution:

$$p(\hat{\ell}_+, \hat{\ell}_-) = \frac{1}{\sigma_{\ell\bar{\ell}}} \frac{\mathrm{d}\sigma_{\ell\bar{\ell}}}{\mathrm{d}\Omega_+ \mathrm{d}\Omega_-} = \frac{1 + \mathbf{B}^+ \cdot \hat{\ell}_+ - \mathbf{B}^- \cdot \hat{\ell}_- - \hat{\ell}_+ \cdot \mathbf{C} \cdot \hat{\ell}_-}{(4\pi)^2}$$

• Direct one-qubit tomography of $\rho_{A,B}$, $\rho_{\hat{\mathbf{n}}}$ from Bloch vectors \mathbf{B}^{\pm} , $\mathbf{B}_{\hat{\mathbf{n}}}^{\pm}$:

$$p(\hat{\ell}_{\pm}) = \int \mathrm{d}\Omega_{\mp} \ p(\hat{\ell}_{\pm}, \hat{\ell}_{-}) = \frac{1 \pm \mathbf{B}^{\pm} \cdot \hat{\ell}_{\pm}}{4\pi}$$
$$p(\hat{\ell}_{\pm} | \hat{\ell}_{\mp} = \mp \hat{\mathbf{n}}) = \frac{p(\hat{\ell}_{\pm}, \hat{\ell}_{\mp} = \mp \hat{\mathbf{n}})}{p(\hat{\ell}_{\mp} = \mp \hat{\mathbf{n}})} = \frac{1 \pm \mathbf{B}_{\hat{\mathbf{n}}}^{+} \cdot \hat{\ell}_{\pm}}{4\pi}$$

- Actual discord \rightarrow Evaluated from minimization over $\hat{\mathbf{n}}$.
- Measurement of $\mathbf{B}_{\hat{\mathbf{n}}}^{\pm} \rightarrow \text{Reconstruction of } t, \bar{t} \text{ steering ellipsoids.}$
- Highly-challenging measurements in conventional setups → Natural implementation in colliders!

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Bell Test Loopholes in a Collider Experiment

- Loopholes: Experimental tests of Bell's inequality may not fulfill all hypotheses of Bell's theorem.
- Collider experiment:
 - Free-will loophole: Spin measurement directions should be free, independent from hidden-variables. → Not even single-detection events from Alice and Bob!
 - Detection loophole: Only a subset of events selected for measurement \rightarrow Bias!
- Quite natural: Colliders were not designed to test Bell's Inequality!



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New Physics Witnesses

- Approximate *CP*-invariance of Standard Model → C = C^T, B⁺ = B⁻ → Symmetric discord and steering!
- Therefore: Discord and/or Steering asymmetry → New Physics!
- New physics witnesses: Symmetry protected observables by SM, only non-zero for New Physics:
 - - Asymmetries in ellipsoid centers and/or semiaxes.
- No SM contribution to New Physics witnesses!



Conclusions and outlook

- Quantum Information theory ↔ High-Energy Physics. Interdisciplinary, huge potential and great interest!
- QI perspective:
 - Highest-energy observation of entanglement ever!
 - ② Genuinely relativistic, exotic symmetries and interactions, fundamental nature → Frontier of known Physics!
 - **3** Highly-demanding measurements naturally implemented at LHC.
- HEP perspective:
 - Quantum Tomography: Novel experimental tool.
 - QI techniques can inspire new approaches for searching New Physics:
 - Aoude, Madge, Maltoni, Mantani, PRD (2022).
 - Severi, Vryonidou, JHEP (2023).
 - Fabbrichesi, Floreanini, Gabrielli, EPJC (2023).
- Extension to e⁺e[−] colliders: Spin of initial state can be controlled!
 → Manipulation of qubits? Quantum gates?
- Adaptation to τ leptons, qutrits W^{\pm}, Z^{0} .

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 - 2 QI techniques can inspire new approaches for searching New Physics:
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 - Severi, Vryonidou, JHEP (2023).
 - Fabbrichesi, Floreanini, Gabrielli, EPJC (2023).
- The first measurements of entanglement between a pair of top-quarks are ongoing within ATLAS and CMS.
- The results are expected to be public soon stay tuned!

Thank You



Backup Slides

Backup

Quantum Discord: Two qubits

• Classical two-qubit state:

$$p = p_{++} |\hat{\mathbf{n}} \hat{\mathbf{n}}' \rangle \langle \hat{\mathbf{n}} \hat{\mathbf{n}}' | + p_{+-} |\hat{\mathbf{n}} - \hat{\mathbf{n}}' \rangle \langle \hat{\mathbf{n}} - \hat{\mathbf{n}}' | + p_{-+} |-\hat{\mathbf{n}} \hat{\mathbf{n}}' \rangle \langle -\hat{\mathbf{n}} \hat{\mathbf{n}}' | + p_{--} |-\hat{\mathbf{n}} - \hat{\mathbf{n}}' \rangle \langle -\hat{\mathbf{n}} - \hat{\mathbf{n}}' |$$

- Incoherent statistical mixture of $|\pm \hat{\mathbf{n}} \pm \hat{\mathbf{n}}' \rangle$ states \rightarrow Quantum mechanics if it was just some probabilistic theory
- Classical states ⇐⇒ Zero discord when measuring B along n̂' direction!

$$\mathcal{D}(A,B) = H(\rho_B) - H(\rho) + p_{\mathbf{\hat{n}}'}H(\rho_{\mathbf{\hat{n}}'}) + p_{-\mathbf{\hat{n}}'}H(\rho_{-\mathbf{\hat{n}}'}) = 0$$

Classical states included in separable states:

$$\rho = \int \mathrm{d}\Omega_A \mathrm{d}\Omega_B \, P(\mathbf{n}_A, \mathbf{n}_B) \left| \mathbf{n}_A \mathbf{n}_B \right\rangle \left\langle \mathbf{n}_A \mathbf{n}_B \right|$$

Entanglement → Quantum discord

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