

Summary of WG4: Mixing and mixing-related CP violation in B system: $\Delta m, \Delta\Gamma, \phi_s, \phi_1/\beta, \phi_2/\alpha, \phi_3/\gamma$

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1 Introduction

2 B -meson mixing

3 CP violation and mixing angles

4 Summary and Outlook

Introduction

- $B_s^0 - \bar{B}_s^0$ oscillations between flavor eigenstates $|B_s^0\rangle$ and $|\bar{B}_s^0\rangle$

$$i \frac{d}{dt} \begin{pmatrix} |B_s^0(t)\rangle \\ |\bar{B}_s^0(t)\rangle \end{pmatrix} = \left(\hat{M} - \frac{i}{2} \hat{\Gamma} \right) \begin{pmatrix} |B_s^0(t)\rangle \\ |\bar{B}_s^0(t)\rangle \end{pmatrix},$$

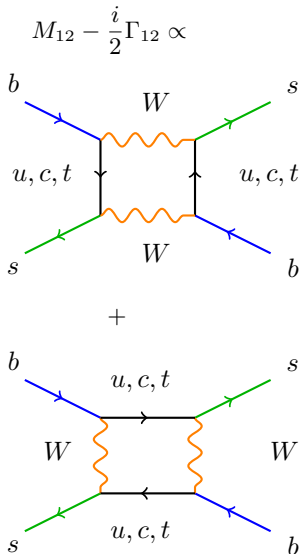
$$\hat{M} = \begin{pmatrix} M_{11} & M_{12} \\ M_{12}^* & M_{22} \end{pmatrix}, \quad \hat{\Gamma} = \begin{pmatrix} \Gamma_{11} & \Gamma_{12} \\ \Gamma_{12}^* & \Gamma_{22} \end{pmatrix}$$

- Diagonalize the matrices

$$|B_{s,L}\rangle = p |B_s^0\rangle + q |\bar{B}_s^0\rangle$$

$$|B_{s,H}\rangle = p |B_s^0\rangle - q |\bar{B}_s^0\rangle$$

- Mass eigenstates: $|B_{s,L}\rangle$ (lighter) and $|B_{s,H}\rangle$ (heavier)



- Physical observables depend on: $|M_{12}|$, $|\Gamma_{12}|$, ϕ_s
- Δm_s : $B_s^0 - \bar{B}_s^0$ oscillation frequency

$$\Delta M_s = M_H - M_L \approx 2|M_{12}|$$

t quark is dominant in SM, sensitivity to NP in the loops

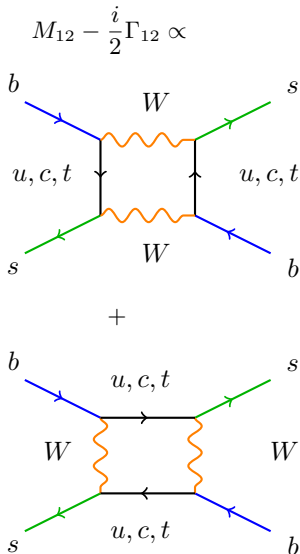
- $\Delta\Gamma_s$: $B_s^0 - \bar{B}_s^0$ width difference

$$\Delta\Gamma_s = \Gamma_L - \Gamma_H \approx 2|\Gamma_{12}| \cos \phi_s$$

only u and c contribute, precision probe of SM, little room for NP

- ϕ_s : CP-asymmetry in the mixing

$$a_{\text{fs}} = \text{Im} \left(\frac{\Gamma_{12}}{M_{12}} \right) = \left| \frac{\Gamma_{12}}{M_{12}} \right| \sin \phi_s$$



In this talk: constraints SM by studying the time-evolution of $B_{(s)}^0 - \bar{B}_{(s)}^0$:

- Recent LHCb & measurements of $\phi_s, \Delta m_s, \gamma$
- Recent LHCb & Belle II measurements of mixing-induced CPV in penguin decays
- Theory progress in understanding perturbative and nonperturbative aspects of B -mixing
- Theory determinations of γ and ϕ_s

B-meson mixing

Numerical results

$\Delta\Gamma$ to NNLO (Gerlach, Nierste, Shtabovenko, Steinhauser 2022)

$$\left. \frac{\Delta\Gamma_s}{\Delta M_s} \right|_{\text{pole}} = \left(3.79_{-0.58}^{+0.53}_{\text{scale}} \quad +0.09_{-0.19}_{\text{scale}, 1/m_b} \pm 0.11_{B\bar{B}_s} \pm 0.78_{1/m_b} \pm 0.05_{\text{input}} \right) \times 10^{-3}, \quad (33)$$

$$\left. \frac{\Delta\Gamma_s}{\Delta M_s} \right|_{\overline{\text{MS}}} = \left(4.33_{-0.44}^{+0.23}_{\text{scale}} \quad +0.09_{-0.19}_{\text{scale}, 1/m_b} \pm 0.12_{B\bar{B}_s} \pm 0.78_{1/m_b} \pm 0.05_{\text{input}} \right) \times 10^{-3}, \quad (34)$$

$$\left. \frac{\Delta\Gamma_s}{\Delta M_s} \right|_{\text{PS}} = \left(4.20_{-0.39}^{+0.36}_{\text{scale}} \quad +0.09_{-0.19}_{\text{scale}, 1/m_b} \pm 0.12_{B\bar{B}_s} \pm 0.78_{1/m_b} \pm 0.05_{\text{input}} \right) \times 10^{-3}. \quad (35)$$

Overall result:

$$\Delta\Gamma^{th} = (0.076 \pm 0.017) \text{ps}^{-1} \quad (36)$$

Motivation

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Pascal Reeck – NNLO QCD corrections to $\Delta\Gamma_{(s)}$

Calculation

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Results

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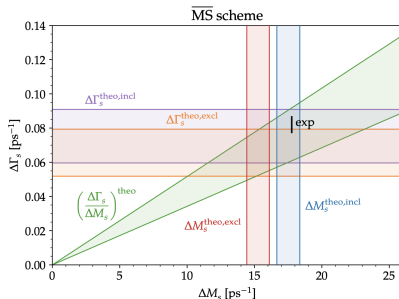
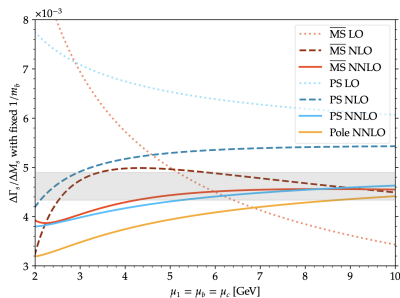
References

Santiago de Compostela, CKM 2023

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Source: NNLO QCD corrections to $\Delta\Gamma_{(s)}$ in the $B_{(s)} - \bar{B}_{(s)}$ system by **Pascal Reeck**

Visualisation of results



Motivation

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Pascal Reeck – NNLO QCD corrections to $\Delta\Gamma_{(s)}$

Calculation

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Results

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References

Santiago de Compostela, CKM 2023

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Source: NNLO QCD corrections to $\Delta\Gamma_{(s)}$ in the $B_{(s)} - \bar{B}_{(s)}$ system by **Pascal Reeck**

Outlook

Contribution	(Gerlach, Nierste, Shtabovenko, Steinhauser 2022)	WIP (Chen, Nierste, Reeck, Shtabovenko, Steinhauser)
$P_{1,2} \times P_{3-6}$	2 loops, $\mathcal{O}(z)$	3 loops, $\mathcal{O}(z^{10})$
$P_{1,2} \times P_8$	2 loops, $\mathcal{O}(z)$	3 loops, $\mathcal{O}(z^{10})$
$P_{3-6} \times P_{3-6}$	2 loops, $\mathcal{O}(z)$	3 loops, $\mathcal{O}(z^{10})$
$P_{3-6} \times P_8$	2 loops, $\mathcal{O}(z)$	3 loops, $\mathcal{O}(z^{10})$
$P_8 \times P_8$	2 loops, $\mathcal{O}(z)$	3 loops, $\mathcal{O}(z^{10})$
$P_{1,2} \times P_{1,2}$	3 loops, $\mathcal{O}(z)$	3 loops, $\mathcal{O}(z^{10})$

Motivation

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Pascal Reeck – NNLO QCD corrections to $\Delta\Gamma_{(s)}$

Calculation

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Results

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References

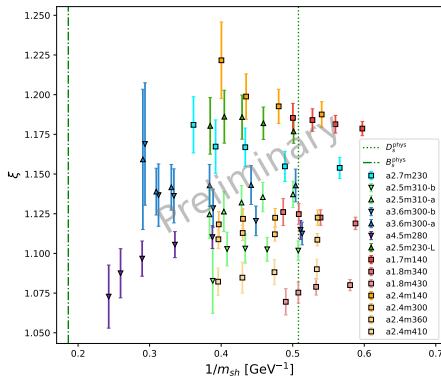
Santiago de Compostela, CKM 2023

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Source: NNLO QCD corrections to $\Delta\Gamma_{(s)}$ in the $B_{(s)} - \bar{B}_{(s)}$ system by **Pascal Reeck**

MIXING RATIOS ξ

- update of RBC/UKQCD work
[Boyle et al., arxiv 1812.08791]
- includes JLQCD ensembles
- completely new, fully correlated fitting strategy
- cancellation of renormalisation constants
- relatively flat $1/m_{sh}$ dependence with improved reach towards m_b^{phys}
- we are currently investigating various global fits on the data

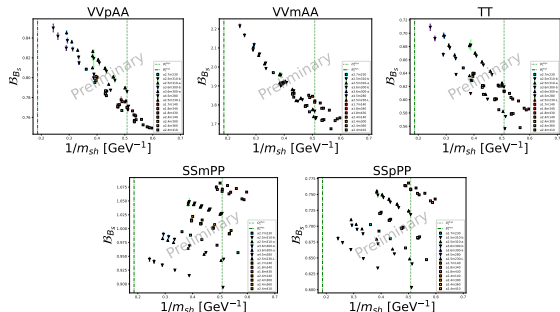


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Source: Update on $SU(3)$ -breaking ratios and bag parameters for $B_{(s)}$ mesons by **Felix Erben**

BAG PARAMETER \mathcal{B}_{hs} - ALL 5 OPERATORS

- heavy-strange bag parameters, renormalised at mass scale μ
- O_1, O_2 : mild α^2 dependence
- O_3, O_4 : strong α^2 dependence
- O_5 : medium α^2 dependence and curvature in $1/m_{sh}$
- very similar for heavy-light sector



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Source: Update on $SU(3)$ -breaking ratios and bag parameters for $B_{(s)}$ mesons by **Felix Erben**

CP violation and mixing angles

$B^0 \rightarrow J/\psi K^+ K^-$ at CMS Alberto Bragagnolo

CMS performed time-dependent analysis of

$B^0 \rightarrow J/\psi(\nu^+ \mu^-) K^+ K^-$ using:

- 2017-2018 data ($L = 96 \text{ fb}^{-1}$)
 $\Rightarrow 48500 \pm 250$ signal candidates
- Opposit-side muon tag
 \Rightarrow improvement in Run 3 using SS, electron, & jet-tagging

First measurement of Δm_s at CMS:

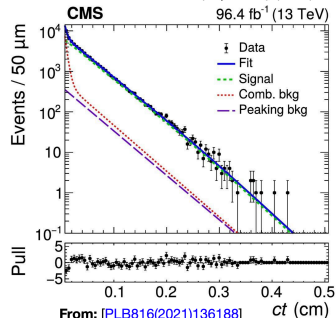
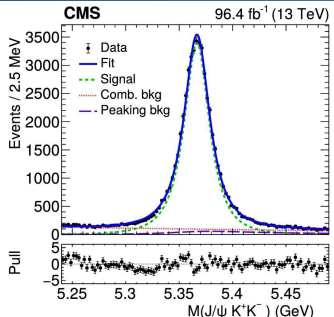
$$\Delta m_s = 17.51_{0.09}^{+0.10} \pm 0.03 \text{ ps}^{-1}$$

Measurement of ϕ_s and $\Delta\Gamma_s$ combined with previous analysis at 8 TeV:

$$\phi_s = 0.021 \pm 0.044 \pm 0.010 \text{ rad}$$

$$\Delta\Gamma_s = 0.1032 \pm 0.0095 \pm 0.0048 \text{ ps}^{-1}$$

- Compatible with SM expectation
- no evidence for CPV



From: [\[PLB816\(2021\)136188\]](#)

Session summary

$B^0 \rightarrow J/\psi K^+ K^-$ at LHCb Melissa Cruz Torres

LHCb recently updated time-dependent analysis of

$B^0 \rightarrow J/\psi(\mu^+ \mu^-) K^+ K^-$ using:

- Full Run 2 dataset
 $\Rightarrow \sim 349000$ signal candidates
- Combination of various OS and SS taggers
 tagging efficiency comparable as Run 1

Result:

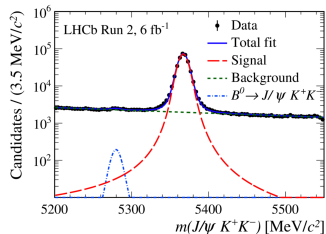
$$\begin{aligned} \bullet \phi_s &= -0.039 \pm 0.022 \pm 0.006 \text{ rad} & \bullet \Gamma_s - \Gamma_d &= -0.0056^{+0.0013}_{-0.0015} \pm 0.0014 \\ \bullet |\lambda| &= 1.001 \pm 0.011 \pm 0.005 & \bullet \Delta\Gamma_s &= 0.0845 \pm 0.0044 \pm 0.0024 \end{aligned}$$

Combined with Run 1 data:

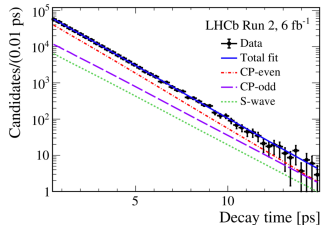
$$\phi_s = -0.044 \pm 0.020 \text{ rad}$$

$$|\lambda| = 0.990 \pm 0.010$$

Good compatibility with SM and no evidence for CPV



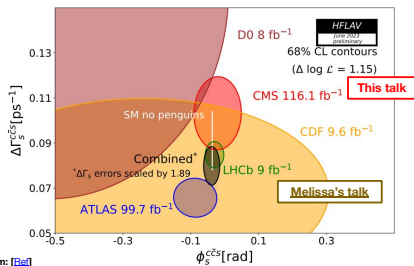
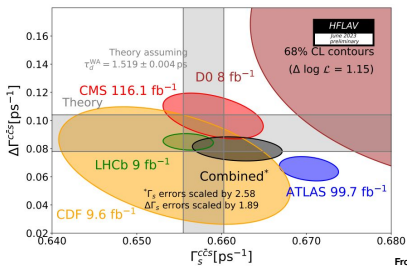
(arXiv: 2308.01468)



$\phi_s, \Gamma_s, \Delta\Gamma_s$: state of the art

State of the art (w. latest preliminary results from LHCb)

- Measurement **statistically limited** \rightarrow long-term commitment by multiple experimental collaborations
- Very active theoretical community (NP limits, penguin pollutions, predictions, ...)
- Precision on ϕ_s **close** to 3 s.d. sensitivity for CPV in decay/mixing interference
 - $\sigma^{\text{WA}}(\phi_s) \approx 15 \text{ mrad}$ (40% relative uncertainty)



Alberto Bragagnolo (UNIPD)

Measurement of the CPV phase ϕ_s

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Source: Measurement of the CP-violating phase ϕ_s with CMS: present and future by **Alberto Bragagnolo**

CP violation measurements in the penguin-mediated decay $B_s^0 \rightarrow \phi\phi$

Measured observables in the polarization-independent fit

arXiv:2304.06198

Parameter	Result
$\phi_s^{s\bar{s}s}$ [rad]	$-0.042 \pm 0.075 \pm 0.009$
$ \lambda $	$1.004 \pm 0.030 \pm 0.009$
$ A_0 ^2$	$0.384 \pm 0.007 \pm 0.003$
$ A_{\perp} ^2$	$0.310 \pm 0.006 \pm 0.003$
$\delta_{\parallel} - \delta_0$ [rad]	$2.463 \pm 0.029 \pm 0.009$
$\delta_{\perp} - \delta_0$ [rad]	$2.769 \pm 0.105 \pm 0.011$

The following parameters have been constrained to the measurements by LHCb collaboration

$$\Delta m_s = 17.766 \pm 0.006 \text{ ps}^{-1}$$

$$\Gamma_s = 0.657 \pm 0.002 \text{ ps}^{-1}$$

$$\Delta\Gamma_s = 0.078 \pm 0.006 \text{ ps}^{-1} \text{ with correlation coefficient of } -0.35$$

In combination with LHCb Run 1 measurements

$$\phi_s^{s\bar{s}s} = -0.074 \pm 0.069 \text{ rad and } |\lambda| = 1.009 \pm 0.07 \pm 0.030$$

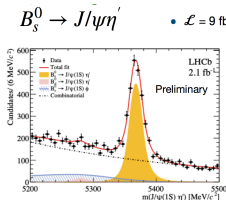
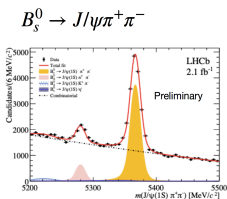
arXiv:2304.06198

This is the most precise measurement of CP violation in $B_s^0 \rightarrow \phi\phi$ to date

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Source: Measurement of the CP violating phase ϕ_s and $\phi_s^{s\bar{q}q}$ by **Melissa Cruz Torres (LHCb)**

$$B_s^0 \rightarrow J\psi\eta', B_s^0 \rightarrow J\psi\pi^+\pi^- \text{ at LHCb } (\Delta\Gamma_s)$$



• $\mathcal{L} = 9 \text{ fb}^{-1}$, Run 1 2011 + 2012 and Run 2 2015 to 2018 data,

LHCb-PAPER-2023-025

$\Delta\Gamma_s$ results and probability of χ^2

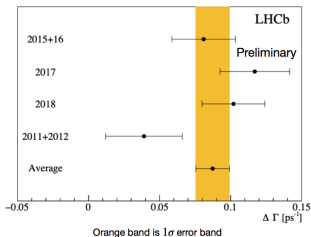
Dataset	$\Delta\Gamma_s$ [ps^{-1}]	$P(\chi^2)$
2011+12	0.039 ± 0.026	0.83
2015+16	0.081 ± 0.022	0.77
2017	0.117 ± 0.024	0.57
2018	0.102 ± 0.021	0.78

Summary

Using full pp-collision dataset between 2011 and 2018, $B_s^0 \rightarrow J\psi\eta'$ and $B_s^0 \rightarrow J\psi\pi^+\pi^-$, $\Delta\Gamma_s$ is measured to be

$$\Delta\Gamma_s = 0.087 \pm 0.012 \pm 0.009 \text{ ps}^{-1}$$

Comparison between the four data sets

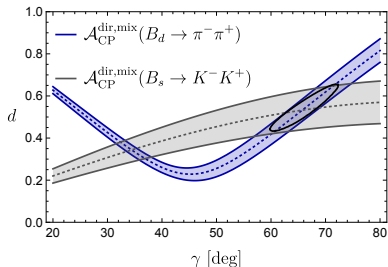


— This is the first $\Delta\Gamma_s$ measurement using the $B_s^0 \rightarrow J\psi\eta'$

Source: Measurement of the CP violating phase ϕ_s and $\phi_s^{sq\bar{q}}$ by **Melissa Cruz Torres (LHCb)**

CKM-angle γ from non-tree decays

Fleischer [1999,2007]; Fleischer, Kneijens [2011]; Fleischer, Malami, Jaarsma, KKV [2016]
 Cuichini, Franco, Mishima, Silvestrini [2012], Data from LHCb [2022] **Fleischer, Jaarsma, KKV [2211.08346]**



- **New!** First observation of CP violation in penguin dominated $B_s \rightarrow K^+ K^-$ LHCb 2022
- **New!** First determination of γ with only CP asymmetries
- $\gamma = (65_{-5}^{+7})^\circ$ **Fleischer, Jaarsma, KKV [2111.08346]**
- Agrees with tree determinations: $\gamma = (64.9 \pm 4.5)^\circ$ LHCb [2021] without B_s modes
- Limited by U -spin breaking corrections

Source: Theory determination of γ and ϕ_s from $B_{(s)} \rightarrow hh$ decays by **Keri Vos**

Determination of ϕ_s (II)

Fleischer, Jaarsma, KKV, JHEP 02 (2023) 081 [2211.08346]

Strategy II:

- Use ratio of branching ratios of $B_d \rightarrow \pi\pi$ and $B_s \rightarrow KK$ decays $K = 105.3 \pm 9.6$
- Non-factorisable U -spin-breaking contributions:

$$\xi_{\text{NF}}^a \equiv \left| \frac{1+r_P}{1+r'_P} \right| \left| \frac{1+x}{1+x'} \right| \left| \frac{a_{\text{NF}}^T}{a_{\text{NF}}^{T'}} \right| = 1.00 \pm 0.07,$$

- $\Delta\phi_{KK} = -(4.5 \pm 5.3)^\circ$
- With $\phi_s^{\text{eff}} = -(8.1 \pm 1.9)^\circ \rightarrow \phi_s = -(3.6 \pm 5.7)^\circ$

Remarkable agreement with $B_s^0 \rightarrow J/\psi\phi$ determination: $\phi_s = -(4.2 \pm 1.4)^\circ$

$$B^0 \rightarrow K_S^0 \pi^0 \gamma$$

- Radiative penguin, fully neutral final state
- Most precise result to date:

$$S = 0.00_{-0.26}^{+0.27} \pm 0.03$$

$$C = 0.10 \pm 0.13 \pm 0.03$$

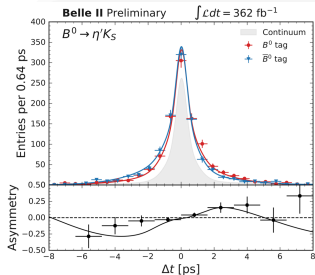
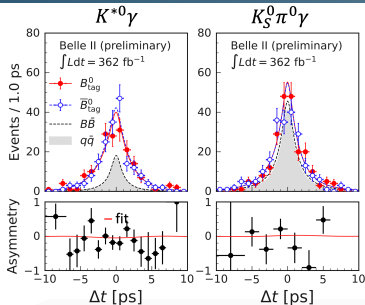
in K^{*0} resonance region

$$B^0 \rightarrow \eta' K_S^0$$

- Most abundant $b \rightarrow s$ had. penguin
- Precision approaching world best:

$$S = 0.67 \pm 0.10 \pm 0.04$$

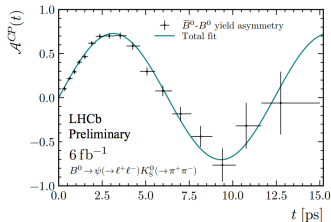
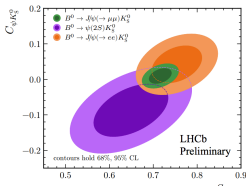
$$C = -0.19 \pm 0.08 \pm 0.03$$



Precision of TD analyses will increase w/ new flavor tagger: $\sim 37\%$ eff vs $\sim 29\%$ at Belle.

$\sin(2\beta)$ at LHCb

Mode	S	C	ρ	
$J/\psi(\rightarrow \mu^+\mu^-) K_S^0$	$0.716 \pm 0.015 \pm 0.007$	$+0.010 \pm 0.014 \pm 0.003$	0.446	~306k signal decays
$\psi(2S)(\rightarrow \mu^+\mu^-) K_S^0$	$0.649 \pm 0.053 \pm 0.018$	$-0.087 \pm 0.048 \pm 0.005$	0.503	~23k signal decays
$J/\psi(\rightarrow e^+e^-) K_S^0$	$0.754 \pm 0.037 \pm 0.008$	$+0.042 \pm 0.034 \pm 0.008$	0.374	~43k signal decays



Combination from simultaneous fit of the three decay modes:

$$S_{\psi K_S^0} = 0.717 \pm 0.013 \pm 0.008$$

$$C_{\psi K_S^0} = 0.008 \pm 0.012 \pm 0.003$$

$$\rho = 0.441$$

- To date, most precise single measurement of $\sin(2\beta)$
- Should improve precision of world average by $\sim 35\%$
- Still statistically limited

Source: Measurements of $\sin(2\beta)$ at LHCb by **Thomas Latham** - LHCb-PAPER-2023-013

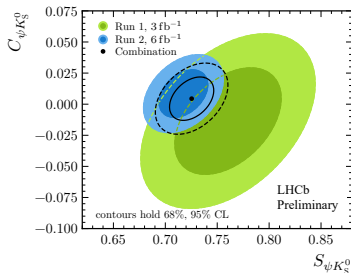
Combination with Run 1

- Combination with the previous Run 1 results
 - Assumes sources of systematic uncertainties from external parameters $\Delta m_d, \Delta \Gamma_d, A_P$ are fully correlated

$$S_{\psi K_S^0} = 0.724 \pm 0.014$$

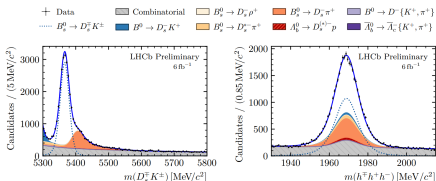
$$C_{\psi K_S^0} = 0.004 \pm 0.012$$

$$\rho = 0.40$$

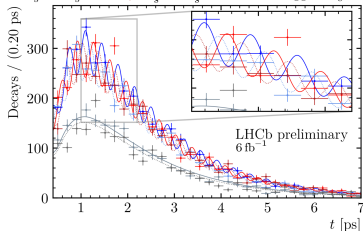


$B_s^0 \rightarrow D_s^\mp K^\pm$ - Run2

[4] LHCb-CONF-2023-004

• 20950 \pm 180 candidates

$B_s^0 \rightarrow D_s^- K^+$ $\bar{B}_s^0 \rightarrow D_s^- K^+$ $B_s^0 \rightarrow D_s^- K^+$ Untagged $D_s^- K^+$
 $B_s^0 \rightarrow D_s^+ K^-$ $\bar{B}_s^0 \rightarrow D_s^+ K^-$ $B_s^0 \rightarrow D_s^+ K^-$ Untagged $D_s^+ K^-$



- Significant CP violation in the interference
 $S_f \neq -S_{\bar{f}}$ at 8.8 σ

$$C_f = 0.791 \pm 0.061 \pm 0.022$$

$$A_f^{\Delta\Gamma} = 0.051 \pm 0.134 \pm 0.037$$

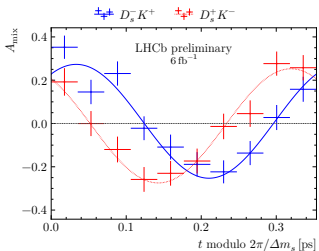
$$S_f = -0.571 \pm 0.084 \pm 0.023$$

$$A_{\bar{f}}^{\Delta\Gamma} = 0.303 \pm 0.125 \pm 0.036$$

$$S_{\bar{f}} = -0.503 \pm 0.084 \pm 0.025$$

- No systematic limitation expected in Run3

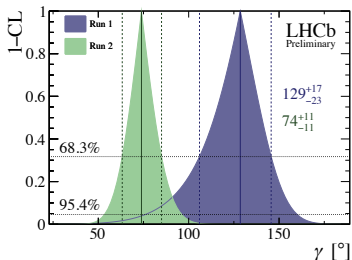
Source: Decay-time-dependent measurements of the CKM angle γ at LHCb by **Quentin Fühling**



- Extraction of physics parameters
 - External input [9]
 - $-2\beta_s = \phi_s = (-0.031 \pm 0.018) \text{ rad}$
- Run2 standalone result:

$$\gamma = (74 \pm 11)^\circ$$

$$\delta = (346.9 \pm 6.6)^\circ \quad r_{D,K} = 0.327 \pm 0.038$$



- Compatibility to Run1 [2] at 1.3σ
 - Driven by γ at 2σ and $\text{Re}[\lambda_f]$
 - $r_{D,K}$ and δ at 0.6σ each
- Updated machinery reproduces Run1 result [2]
- Combination in preparation

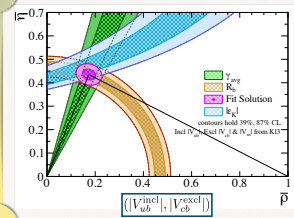
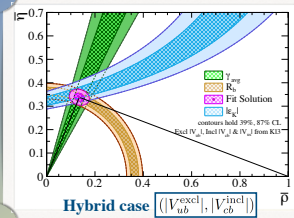
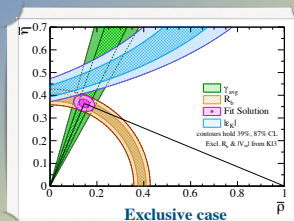
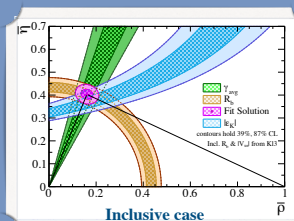
[4] LHCb-CONF-2023-004

Source: Decay-time-dependent measurements of the CKM angle γ at LHCb by **Quentin Fühling**

Unitarity Triangle Apex Determination

► Making a fit to γ and R_b

Incl	$\bar{\rho} = 0.161 \pm 0.025$,	$\bar{\eta} = 0.403 \pm 0.022$
Excl	$\bar{\rho} = 0.146 \pm 0.022$,	$\bar{\eta} = 0.364 \pm 0.018$
$(V_{ub}^{\text{excl}} , V_{cb}^{\text{incl}})$	$\bar{\rho} = 0.135 \pm 0.021$,	$\bar{\eta} = 0.338 \pm 0.017$
$(V_{ub}^{\text{incl}} , V_{cb}^{\text{excl}})$	$\bar{\rho} = 0.174 \pm 0.027$,	$\bar{\eta} = 0.435 \pm 0.023$



Source: New Physics in $B_q - \bar{B}_q$ mixing in connection with CKM angle γ by **Eleftheria Malami**

How can we determine the Unitarity Triangle Apex?

2) Utilising Mixing and R_b - without γ

- The UT side R_t is defined as:

$$R_t \equiv \left| \frac{V_{td}V_{tb}}{V_{cd}V_{cb}} \right| = \sqrt{(1-\bar{\rho})^2 + \bar{\eta}^2} = \frac{1}{\lambda} \left| \frac{V_{td}}{V_{ts}} \right| \left[1 - \frac{\lambda^2}{2} (1-2\bar{\rho}) \right] + \mathcal{O}(\lambda^4)$$

Incl	$\Delta m_d^{\text{SM}} = (0.511 \pm 0.040) \text{ ps}^{-1}$,	$\Delta m_s^{\text{SM}} = (17.23 \pm 0.87) \text{ ps}^{-1}$
Excl	$\Delta m_d^{\text{SM}} = (0.438 \pm 0.033) \text{ ps}^{-1}$,	$\Delta m_s^{\text{SM}} = (14.80 \pm 0.76) \text{ ps}^{-1}$
$(V_{ub}^{\text{excl}} , V_{cb}^{\text{incl}})$	$\Delta m_d^{\text{SM}} = (0.509 \pm 0.037) \text{ ps}^{-1}$,	$\Delta m_s^{\text{SM}} = (17.19 \pm 0.87) \text{ ps}^{-1}$
$(V_{ub}^{\text{incl}} , V_{cb}^{\text{excl}})$	$\Delta m_d^{\text{SM}} = (0.442 \pm 0.036) \text{ ps}^{-1}$,	$\Delta m_s^{\text{SM}} = (14.84 \pm 0.76) \text{ ps}^{-1}$

fit to the sides R_b and R_t

Incl, $K\bar{L}3$	$\bar{\rho} = 0.180 \pm 0.014$,	$\bar{\eta} = 0.395 \pm 0.020$
Excl, $K\bar{L}3$	$\bar{\rho} = 0.163 \pm 0.013$,	$\bar{\eta} = 0.357 \pm 0.017$
Hybrid, $K\bar{L}3$	$\bar{\rho} = 0.153 \pm 0.013$,	$\bar{\eta} = 0.330 \pm 0.016$

scenarios with γ are a factor 2 less precise than the scenarios without γ

assume Δm_s^{SM}
 Δm_d^{SM}

$$\left| \frac{V_{td}}{V_{ts}} \right| = \xi \sqrt{\frac{m_{B_s} \Delta m_s^{\text{SM}}}{m_{B_d} \Delta m_d^{\text{SM}}}}$$

lattice $\xi \equiv \frac{f_{B_s} \sqrt{\bar{B}_{B_s}}}{f_{B_d} \sqrt{\bar{B}_{B_d}}}$

FLAG(2021), arXiv:1907.01025
 $\xi = 1.212 \pm 0.016$

$$\left| \frac{V_{td}}{V_{ts}} \right| = 0.2063 \pm 0.0004 \pm 0.0027$$

→ due to lattice input

due to experiment

* UT apex determination through R_b and R_t is more precise

* R_t determined assuming SM Δm_d and Δm_s

↓
ignores possible NP in $B_q^0 - \bar{B}_q^0$ mixing
• NP will contaminate R_t determination

* To determine NP in $B_q^0 - \bar{B}_q^0$ mixing in a general scenario: UT apex determination through R_b and γ

Source: New Physics in $B_q - \bar{B}_q$ mixing in connection with CKM angle γ by Eleftheria Malami

Introducing NP Parameters

$$\Delta m_q = \Delta m_q^{\text{SM}} (1 + \kappa_q e^{i\sigma_q})$$

$$\phi_q = \phi_q^{\text{SM}} + \phi_q^{\text{NP}} = \phi_q^{\text{SM}} + \arg(1 + \kappa_q e^{i\sigma_q})$$

Model independent parametrization

size of the NP effects is described by κ_q

σ_q is a complex phase for additional CP-violating effects

We explore 2 different NP scenarios

Scenario I → most general case

utilise UT apex determination for the SM predictions of Δm_q and ϕ_q

NP parameters (κ_d, σ_d) and (κ_s, σ_s) independently from each other

Scenario II

we consider NP contributions are equal in the B_d and the B_s systems

FUNP

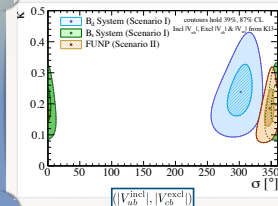
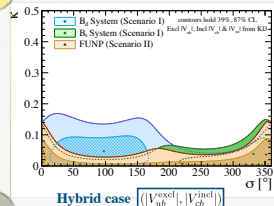
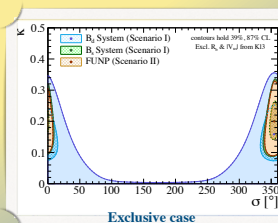
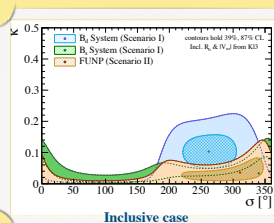
Comparing FUNP with Scenario I
test of FUNP assumption
Impact of the assumptions on the constraints on parameters space of NP in mixing

UT apex determination that only relies on R_b and mixing parameters

without information on γ

NP in γ will not affect the results

Source: New Physics in $B_q - \bar{B}_q$ mixing in connection with CKM angle γ by **Eleftheria Malami**

Comparison between Scenarios I and II for κ_q and σ_q 

Source: New Physics in $B_q - \bar{B}_q$ mixing in connection with CKM angle γ by **Eleftheria Malami**

Determining NP in $B_s^0 \rightarrow \mu^+ \mu^-$

arXiv:hep-ph/0303060
arXiv:2104.09521
arXiv:2109.11032

$$\mathcal{R}_{s\mu} \equiv \left| \frac{\tilde{\mathcal{B}}(B_s \rightarrow \mu^+ \mu^-)}{\Delta m_s} \right|$$

NP can modify its branching ratio
(Pseudo)-Scalar $B_s^0 - \bar{B}_s^0$ mixing

CKM elements drop out in the SM ratio

Including NP effects in both $B(B_s \rightarrow \mu^+ \mu^-)$ and Δm_s we get the generalised expression

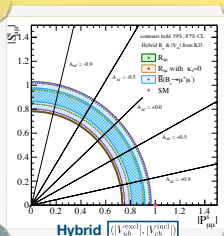
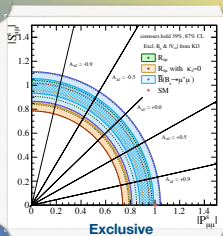
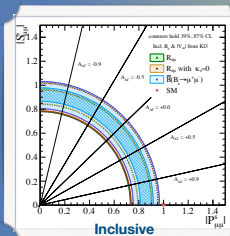
$$\mathcal{R}_{s\mu} = (1.60 \pm 0.19) \times 10^{-10}$$

Comparing with the SM, we obtain extra contours

$$\mathcal{R}_{s\mu}^{\text{SM}} = \frac{\tau_{B_s}}{1 - y_s} \frac{3G_F^2 m_W^2 \sin^4 \theta_W}{4\pi^3} \frac{|C_{10}^{\text{SM}}|^2}{S_0(x_t) \eta_{2B} \hat{B}_{B_s}} m_\mu^2 \sqrt{1 - 4 \frac{m_\mu^2}{m_{B_s}^2}}$$

$$\mathcal{R}_{s\mu}^{\text{SM}} = (2.22 \pm 0.10) \times 10^{-10}$$

introduces a dependence on the CKM matrix elements through the NP parameters (κ_s, σ_s)



We can minimise this dependence, creating the following ratio R_{sp}

Source: New Physics in $B_q - \bar{B}_q$ mixing in connection with CKM angle γ by **Eleftheria Malami**

Summary and Outlook

Summary

- Many exciting new experimental results utilizing large datasets
- Increased precision due to higher statistics
- Theory doing its best to keep up