

Measurement of the CP-violating phase ϕ^s with CMS: present and future

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Outline

- **1. Introduction What? Why?**
- **2. Measuring ϕ^s with the CMS detector Who? Where?**
- **3. Analysis and results How?**
- **4. Future prospects What now?**
- **5. Conclusions**

Motivations

- Decays of B_s mesons allow to study the time-dependent CP **violation generated by the interference between direct decays and flavour mixing**
	- **○** CPV in the interference is possible even if no CPV in decay and mixing
	- **○** *Golden channel:* B_s → J/ψ φ(1020) → μ⁺μ⁻ K⁺K⁻
- **● The weak phase ϕ^s is the main CPV observable**
	- Precisely predicted by the SM to be **ϕ^s ≈ -2βs ≈ -37 ± 1 mrad**, where $\bm{\beta}_{\rm s}$ is one of the angles of the $\bm{\mathsf{B}}_{\rm s}$ unitary triangle (determined very accurately by CKM global fits) [\[CKMfitter](http://ckmfitter.in2p3.fr/www/results/plots_spring21/num/ckmEval_results_spring21.html), [UTfit\]](https://arxiv.org/abs/2212.03894)
- New physics can change the value of $φ_$ up to ~100% via new particles contributing to the flavour oscillations [\[RMP88\(2016\)045002](https://doi.org/10.1103/RevModPhys.88.045002)]

State of the art (w. latest preliminary results from LHCb)

- Measurement **statistically limited** ➜ long-term commitment by multiple experimental collaborations
- Very active theoretical community (NP limits, penguin pollutions, predictions, …)
- Precision on ϕ_s close to 3 s.d. sensitivity for CPV in decay/mixing interference
	- $σ^{WA}$ ($φ_≤$) ≈ 15 mrad (40% relative uncertainty)

Measurement ingredients

● Time-dependent flavour asymmetry

- **● Essential ingredients**
	- Time-dependent **angular analysis** to separate the different CP eigenstate
	- Excellent **time resolution** and **flavour tagging** to see the $\mathsf{B}_{_\mathsf{S}}$ flavour oscillations (T $\scriptstyle\mathtt{\sim}$ 350 fs)
	- Time and angular **efficiencies**

$$
\text{sensitivity}\propto\sqrt{\frac{\epsilon_{\text{tag}}\mathcal{D}_{\text{tag}}^2N_{\text{sig}}}{2}}\sqrt{\frac{N_{\text{sig}}}{N_{\text{sig}}+N_{\text{bkg}}}}\,e^{-\frac{\sigma_t^2\Delta m_s^2}{2}}\Bigg|
$$

Why a CMS measurement?

CMS is a general-purpose detector well suited for studying $B_s \rightarrow J/\psi \phi(1020) \rightarrow \mu^+\mu^- K^+K^-$

- **● Silicon tracking system**
	- Excellent decay time resolution (σ_t ~ 60 fs)
	- **○** Large pseudorapidity range up to |η| = 2.5
- **● Superconducting solenoid**
	- **○** High momentum resolution for charged tracks
- **● Muon system**
	- **○** High efficiency in triggering/reconstructing J/ψ → $\mu^+\mu^-$
	- **σ**(p_T)/p_T ~ O(1%)
- **Enormous amount of data** collected at $\sqrt{s} = 13$ TeV
	- \circ O(1M) of B_s → J/ψ φ(1020) → μ⁺μ⁻ K⁺K⁻ candidates

$$
\circ \qquad P_{tag} \sim 7?
$$

Latest CMS results overview

- **Reference:** [Phys. Lett. B 816 \(2021\) 136188](https://doi.org/10.1016/j.physletb.2021.136188)
- **Dataset:** 2017-2018 ($L_{int} = 96$ fb⁻¹)
- **Trigger**: $J/\psi \rightarrow \mu^+\mu^-$ candidate plus an additional muon
- **Decay length cut:** >70 μm (to reduce prompt bkg.)
- **•** $m(K*K^{-})$ interval: $m(\phi(1020)) \pm 10$ MeV
- **Number of signal candidates:** 48500 ± 250
- **Flavour tagging: opposite-side muon**
	- \circ $\varepsilon_{\text{tag}} \approx 50\%$, $\mathsf{D}_{\text{tag}} \approx 0.2$, $\mathsf{P}_{\text{tag}} \approx 10\%$

Fit: unbinned multidimensional extended maximum-likelihood

- **• Input observables:** m_{Bs} , ct, σ_{ct} , θ_{T} , ψ_{T} , φ_{T} , ξ_{tag} , ω_{tag}
- **Fitted parameters**
	- o CPV observables: ϕ_{s} , |λ|
	- \circ *B system properties:* ΔΓ , Γ , Δm
	- \overline{S} *Decay polarization*: $|A_0|^2$, $|A_1|^2$, $|A_5|^2$, δ_∥, δ_⊥, δ_{s⊥}
- **Bkg sources:** combinatorial, B^0 → J/ψ K^{*0} → μ⁺μ⁻ K⁺π⁻

OS-muon tagging

- **1. OS-muon selection** (very loose)
	- \circ p_T > 2 GeV, $|\eta|$ < 2.5, IP_z(μ, PV) < 1 cm, ΔR_{η,φ}(B_s) > 0.4
- **2. Tagging decision** (assuming $b \rightarrow \mu^{-} X$)
	- \circ $\mu \rightarrow \text{OS b} \rightarrow \text{signal } \overline{\text{b}} \text{ (B_s)}$
	- \circ μ^+ → OS \overline{b} → signal b (\overline{B}_s)
- **3. Mistag probability evaluation**
	- \circ Calibrated DNN trained on B_s MC and fine-tuned on self-tagging $B^+ \rightarrow J/\psi K^+$ data
	- Trained to discriminate *right* tags from *wrong* ones
	- **O** The output score s_{DNN} can be interpreted as a **probability with the DNN trained to reproduce**

Prob(right tag) = s_{DNN} **= 1-** ω_{evt}

Dilution sources: fakes, pileup, cascade decays, mixing of the OS-b

Systematic uncertainties

Leading systematic uncertainties

- \bullet ϕ_{s} \rightarrow model bias
- \bullet $\Delta\Gamma_{\rm s}$ and $\Gamma_{\rm s}$ \rightarrow lifetime efficiency
	- \rightarrow lifetime uncertainty
- \bullet Δm_s \bullet | λ | \rightarrow angular efficiency

Results

● Good agreement with SM predictions

$$
\begin{array}{lll}\n\circ & \varphi_s^{\text{SM}} & = -37 \pm 1 \text{ mrad} \\
\circ & \Delta \Gamma_s^{\text{SM}} & = 0.091 \pm 0.013 \text{ ps}^{-1} \text{ [Lenz & Tetlalmatzi-Xolocotz]} \\
\end{array}
$$

- \circ $|\lambda|^{SM}$ = 1 (no direct CPV)
- Δm_sSM =18.77 ± 0.86 ħps⁻¹ [<u>Lenz & Tetlalmatzi-Xolocotzi</u>]
- **• First measurement by CMS of Δm_s and |λ|**

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Combination with Run1

- **● The results of this analysis are combined with** the ones obtained by CMS at $\sqrt{s} = 8$ TeV^a
	- ϕ _s = -21 ± 44 (stat) ± 10 (syst) mrad
	- \circ ΔΓ_s = 0.1032 ± 0.0095 (stat) ± 0.0048 (syst) ps^{-1}
- **● Results in agreement with the SM predictions**
- The new trigger strategy, which trades the number of events for tagging power, pays off for φ_s while does not improve ΔΓ_s, which sensitivity is driven mainly by statistics

From: [[PLB816\(2021\)136188\]](https://doi.org/10.1016/j.physletb.2021.136188)

a [[PLB757\(2016\)9](http://dx.doi.org/10.1016/j.physletb.2016.03.046)]

Future prospects: precision measurement

CMS is currently working on a precision measurement of ϕ^s with the Run2 dataset by using all available triggers (see SL7)

- **Statistics:** expected to increase the number of signal candidates by a factor of **8~10**
- **Flavour tagging**: muon, electron, jet and same-side (first implementation without hadronic PID)
	- \circ Large enhancement of the effective statistics N_{Bs} x P_{tag}
- **Methodology:** various refinements to deal with the peculiarities of the new dataset
	- Efficiency modelization, background estimation, lifetime resolution, simulation corrections, …
	- *○ Not just a simple statistical scaling!*

- **● Large improvements are expected for all physics parameters**
	- ⊙ Reminder: *sensitivity*(φ_s, Δm_s) ∝ √(P_{tag}N_{Bs}) and *sensitivity*(ΔΓ_s, Γ_s) ∝ √(N_{Bs})
- This measurement will be the benchmark of several new analysis techniques, laying the foundations for future CMS works in the field CP violation

Conclusions

- **● The CPV phase ϕ^s and the decay width difference ∆Γs have been measured using 48 500** B_{s} **→** J/ψ ϕ(1020) signal candidates collected at √s = 13 TeV, corresponding to L_{int} = 96.4 fb⁻¹
- Events are selected using a trigger that requires an additional muon, which is exploited to infer the flavour of the B_s meson at production time, achieving P $_{\rm tag}$ ≈ 10% with small associated systematic uncertainties
- \bullet Results from this measurement are combined with those obtained at $\sqrt{s} = 8$ TeV, yielding

ϕ^s = −21 ± 44 (stat) ± 10 (syst) mrad

∆Γ s = 0.1032 ± 0.0095 (stat) ± 0.0048 (syst) ps-1

- **Results are found to be consistent with the Standard Model predictions, allowing to further** constrain possible contributions from new physics in the B_s meson decay and mixing
- \bullet With the increase in statistics and the development of new techniques, the future for $\phi_{\rm s}$ at LHC looks promising and challenging
- **● CMS is actively working to release an update of the Run2 measurement, adding new data sets and tagging strategies**
	- **○ Stay tuned in the next conference seasons!**

Thanks for the attention

Unitary triangles

. The unitary condition of the CKM matrix leads to the following set of constrains

 $\boxed{\sum_i |V_{ik}|^2 = \sum_k |V_{ik}|^2 = 1}$ \Longrightarrow weak universality

 $\boxed{\sum_i V_{ij}V_{ik}^* = 0}$ \implies six triangles in the complex plane ("unitary triangles")

• Of particular interest for this work is the so-called "B_s unitary triangle":

$$
V_{us}V_{ub}^\ast+V_{cs}V_{cb}^\ast+V_{ts}V_{tb}^\ast=0
$$

$$
\text{with angles:} \quad \alpha_{\text{S}} = \text{arg} \left(-\frac{V_{ts} V_{tb}^*}{V_{us} V_{ub}^*} \right), \quad \beta_{\text{S}} = \text{arg} \left(-\frac{V_{ts} V_{tb}^*}{V_{cs} V_{cb}^*} \right), \quad \gamma_{\text{S}} = \text{arg} \left(-\frac{V_{us} V_{ub}^*}{V_{cs} V_{cb}^*} \right).
$$

meson mixing B_{s}

- \cdot B_s mesons are subject to flavour mixing, that is oscillations between their C-conjugate states before decay
- · The light and heavy mass eigenstates are described by a superposition of flavour states, as

$$
\left|B^{L,H}_S\right\rangle=p\left|B^0_S\right\rangle\pm q\left|\overline{B}^0_S\right\rangle\quad\text{ with }|q|^2+|p|^2=1
$$

 \cdot The B_s system is characterized by the parameters

$$
\begin{aligned} m_s &\equiv \frac{m_H+m_L}{2}, \quad \Gamma_s \equiv \frac{\Gamma_H+\Gamma_L}{2} \\ \Delta m_s &\equiv m_H-m_L, \quad \Delta \Gamma_s \equiv \Gamma_L-\Gamma_H \end{aligned}
$$

• For the B_s^0 system $|q/p| \simeq 1$ is observed¹, so that **the ratio q/p** can be expressed in terms of a complex phase:

$$
\boxed{\frac{q}{p} \equiv e^{-i\phi_M} \simeq \frac{V_{ts}V_{tb}^*}{V_{ts}^*V_{tb}}}
$$

¹World-average value: $|q/p| = 1.0003 \pm 0.0014$ [HFLAV]

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• The flavour eigenstates oscillate with a period of

$$
T=\frac{2\pi}{\Delta m_S}\sim 350\text{ fs}
$$

CPV in mesons

- Observable CP violation is generated by interference between amplitudes
- . Three different types of CP violation are possible
	- 1. "Direct" CPV in decays
		- Observed in kaons, B and D mesons¹
	- 2. "Indirect" CPV in mixing
		- Observed in K^0 oscillations²
	- 3. CPV in the **interference** of decays and mixing
		- Observed in K^0 and B^0 mesons³

$$
\left|\,\mathcal{P}(P\rightarrow f)\neq\mathcal{P}(\overline{P}\rightarrow\overline{f})\,\right|
$$

$$
\mathcal{P}(P^0\to \overline{P}{}^0)\neq \mathcal{P}(\overline{P}{}^0\to P^0)
$$

$$
\boxed{\mathcal{P}(P^0 \to f) \neq \mathcal{P}(P^0 \to \overline{P}^0 \to f)}
$$

• Defining A_f as the P \rightarrow f amplitude, CPV information is encoded in the rephasing-invariant complex parameter λ :

$$
\boxed{\lambda \equiv \frac{q}{p} \frac{\overline{A}_f}{A_f}} \begin{cases} \left| \overline{A}_{\overline{f}} / A_f \right| \neq 1 & \to \text{ direct CPU} \\ \left| q/p \right| \neq 1 & \to \text{ indirect CPU} \\ \left| \lambda \right| = 1, \text{ Im}(\lambda) \neq 0 & \to \text{ interference CPU} \end{cases}}
$$

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CPV in Bs ➜ **J/ψ ϕ(1020)** ➜ **μ⁺ μ - K+K -**

• No direct CPU:
$$
\left| \frac{A_{\overline{f}}}{A_{f}} \right| = \left| \frac{V_{cb}V_{cs}^{*}}{V_{cb}^{*}V_{cs}} \right| =
$$

- $\Big| \ = \left| \frac{V_{tb}^* V_{ts}}{V_{tb} V_{ts}^*} \right| \simeq 1$ $\frac{q}{p}$ · No CPV in mixing:
- CPV in the **interference** (neglecting penguin contributions)¹:

$$
\lambda = \underbrace{\eta_f \left(\frac{V_{cb} V_{cs}^*}{V_{cb}^* V_{cs}}\right)}_{=\overline{A}_f/A_f} \underbrace{\left(\frac{V_{tb}^* V_{ts}}{V_{tb} V_{ts}^*}\right)}_{=q/p} \stackrel{|\lambda|=1}{=} \eta_f \, e^{-i \phi_S}
$$

$$
\phi_s = -2 \arg \left(-\frac{V_{ts} V_{tb}^*}{V_{cs} V_{cb}^*}\right) = -2 \beta_s
$$

where η_f is the CP eigenvalue of the final state and β_s is one of the angles of the B_s^0 unitary triangle

 \cdot β _s can be determined very precisely with global CKM fits

¹ Penguin transitions are predicted to change the value of ϕ_S by about \sim 1 mrad, almost two order of magnitudes smaller than the current experimental sensitivity (\sim 30-50 mrad)

The B s ➜ **J/ψ ϕ(1020)** ➜ **μ⁺ μ - K+K - decay**

- **● B s** ➜ **J/ψ ϕ(1020)** ➜ **μ⁺ μ[−] K+K − is the golden channel for measuring ϕ^s**
	- 1. The final state can be reconstructed with high S/B ratio
	- 2. $J/\psi \rightarrow \mu^+\mu^-$ is easy to trigger
	- 3. Only one CP-violating phase ("golden mode") if neglecting penguin contributions
	- 4. SM predicts no direct CPV
- **● The final state is a mixture of CP-even and CP-odd eigenstates**
	- Spin-0 pseudo-scalar meson (B_s) decaying into two spin-1 vector mesons (J/ψ φ(1020))
	- o The CP eigenvalue of the final state depends on the value of the orbital momentum, as η_f = (−1)^l
- **● The B s** ➜ **J/ψ ϕ(1020)** ➜ **μ⁺ μ[−] K+K −decay amplitude can be decomposed into three polarization states**
	- **A**₀ : l = 0 → CP-even
	- **A**_⊥ : l = **1 → CP-odd**
	- A_{| :} | = 2 → CP-even
- **● Additional contribution ("S-wave") from non-resonant**
	- **B s** ➜ **J/ψ K+K − and B s** ➜ **J/ψ f⁰ (980) is assumed**

$$
\circ \quad A_{\rm S}: I = 0 \rightarrow CP\text{-odd}
$$

Decay rate model

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Angular efficiency

- Detector acceptance and event selection lead to non-uniform angular efficiency
- 3D angular efficiency functions are evaluated in bins of $\cos\theta_T$, $\cos\psi_T$ and φ_T , separately for 2017 and 2018, using simulated samples
- The efficiency functions are obtained with a projection of the 3D angular efficiency histograms on an orthogonal basis:
	- 1. Construct efficiency histograms
		- Numerator: 3D angular RECO histograms from $\Delta\Gamma_s = 0$ MC samples
		- Denominator: 3D angular GEN histograms from GEN only sample
		- Binning: 70 bins for $\cos \theta_T$ and $\cos \psi_T$, and 30 for φ_T
	- 2. Project on Legendre orthogonal basis

$$
b_{l,k,m}(\Theta) = P_l^m(\cos\theta_T)\cdot P_k^m(\cos\psi_T)\cdot \begin{cases} \sin(m\,\varphi_T) & \text{ if } m < 0 \\ \cos(m\,\varphi_T) & \text{ if } m > 0 \\ 1/2 & \text{ if } m = 0 \end{cases}
$$

- up to order 6
- 3. Construct angular efficiency as

$$
\epsilon(\Theta) = \sum_{l,k,m} c_{l,k,m} \cdot b_{l,k,m}(\Theta)
$$

 \cdot C_{LK, m} are the projection coefficients

Measurement of the CPV phase ϕ_s

DNNs as probability estimators

- \cdot A DNN can be (naturally) engineered to predict the probability of a given input example to belong to one of the classes for which the network has been trained
	- . e.g.: what is the probability that the flavour inference for a given event is correct?
- To this end, the softmax activation function can be used in the last layer to normalize the output of the network to a probability distribution consisting of K probabilities in the interval [0, 1] that add up to 1:¹

$$
\boxed{\sigma(\vec{z})_i = \frac{e^{z_i}}{\sum_{j=1}^K e^{z_j}} \quad (i = 1, \ldots, K)}
$$

- A DNN is called "calibrated" if the output probability of the predicted class reflects its true posterior probability (the softmax function does not ensure calibrated networks)
- To improve probability calibration, the cross-entropy loss function can be used in the training process

$$
\boxed{\mathcal{L}_{CE} = -\sum_{i=1}^K t_i \log(s_i)}
$$
 where $\begin{array}{l} s_i = \text{network scores} \\ t_i = \text{one-hot encoded truth labels}^2 \end{array}$

• \mathcal{L}_{CF} can be interpreted as the negative log-likelihood for the conditional probability $P(\vec{t} | \vec{s})$

1Based on the Luce's choice axiom. Ref: R.D. Luce, "Individual choice behavior: a theoretical analysis", Wiley, New York (1959)

 2 Every entry of this equal to 0 but the one corresponding to the true class, which is equal to 1

Measurement of the CPV phase ϕ

Deep Neural Network for flavour tagging

· Training features

- · Muon variables
	- p_T , η , IP_{xy}, $\sigma_{IP_{xy}}$, IP_z, σ_{IP_z} , $\Delta R(B_s^0)$, ...
- · Surrounding activity variables ("muon cone")
	- Constructed from tracks around the muon direction
	- ISO_{μ} , Qcone, P_{T.cone}, E_{μ}/E_{cone}, ...
- · Architecture: fully connected
	- 3 layers of 200 neurons
	- Rectified Linear Unit activation: $f(x) = max(0, x)$
	- 40% dropout probability
		- Probability of temporarily removing a neuron in each training iteration, used to reduce overtraining
- Output: softmax
- Loss: binary cross-entropy
- · Optimizer: Adam
	- Adaptive optimization algorithm specifically designed for training DNNs

Measurement of the CPV phase ϕ_s

Fit model

$$
P = \frac{N_{sgn}}{N_{tot}} P_{sgn} + \frac{N_{bkg}}{N_{tot}} P_{bkg}
$$

 $\mathsf{P}_{\mathsf{sgn}} = \epsilon(\mathsf{ct}) \, \epsilon(\Theta) \, \left[\mathsf{f}(\Theta, \mathsf{ct} \, | \, \alpha) \otimes \mathsf{G}(\mathsf{ct}, \sigma_{\mathsf{ct}}) \right] \, \mathsf{P}_{\mathsf{sgn}}(\mathsf{m}_{\mathsf{B}^0_\epsilon}) \, \mathsf{P}_{\mathsf{sgn}}(\sigma_{\mathsf{ct}}) \, \mathsf{P}_{\mathsf{sgn}}(\xi_{\mathsf{tag}})$

- \cdot ϵ (ct) ϵ (Θ): efficiency functions
- \cdot f(Θ , ct | α): differential decay rate pdf
- $G(ct, \sigma_{ct})$: Gaussian resolution function

 $\cdot \Theta = (\cos \theta_{\text{T}}, \cos \psi_{\text{T}}, \varphi_{\text{T}})$

- $P_{sgn}(m_{B2})$: mass pdf
- $P_{\text{sgn}}(\sigma_{\text{ct}})$: proper decay length uncertainty pdf
- $P_{sgn}(\xi_{tag})$: tag decision pdf

 $\alpha = (\phi_s, \Gamma_s, \Delta\Gamma_s, \Delta\mathfrak{m}_s, |\lambda|, A_0, A_\perp, A_s, \delta_\parallel, \delta_\perp, \delta_{s\perp})$

 $P_{bkg} = P_{bkg}(\cos\theta_T, \varphi_T) P_{bkg}(\cos\psi_T) P_{bkg}(ct) P_{bkg}(m_{B^0}) P_{bkg}(\sigma_{ct}) P_{bkg}(\xi_{tag})$

- $P_{bkg}(\cos\theta_T, \varphi_T)$, $P_{bkg}(\cos\psi_T)$, $P_{bkg}(\text{ct})$: background angular and proper decay length pdfs
- \cdot P_{bkg} contains a dedicated term to model the **peaking background** from $B^0 \rightarrow J/\psi K^*(892)^0 \rightarrow \mu^+\mu^- K^+\pi^-$ where the pion is misidentified as a kaon
	- The peaking background from $\Lambda_b^0 \to J/\psi K^- p \to \mu^+ \mu^- K^- p$ is estimated to be negligible

Fit results

Correlations in the 13 TeV results

Table 1: Statistical correlation matrix between the physics parameters as obtained from the ML fit to the 13 TeV data.

Full combination results

Correlation in the combination

Table 3: Correlations between the physics parameters as obtained from the combination between the CMS 8 TeV and 13 TeV results. Correlations are both statistical and systematic.

