



Imperial College
London

Data-driven determination of charm loop effects in $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ decays

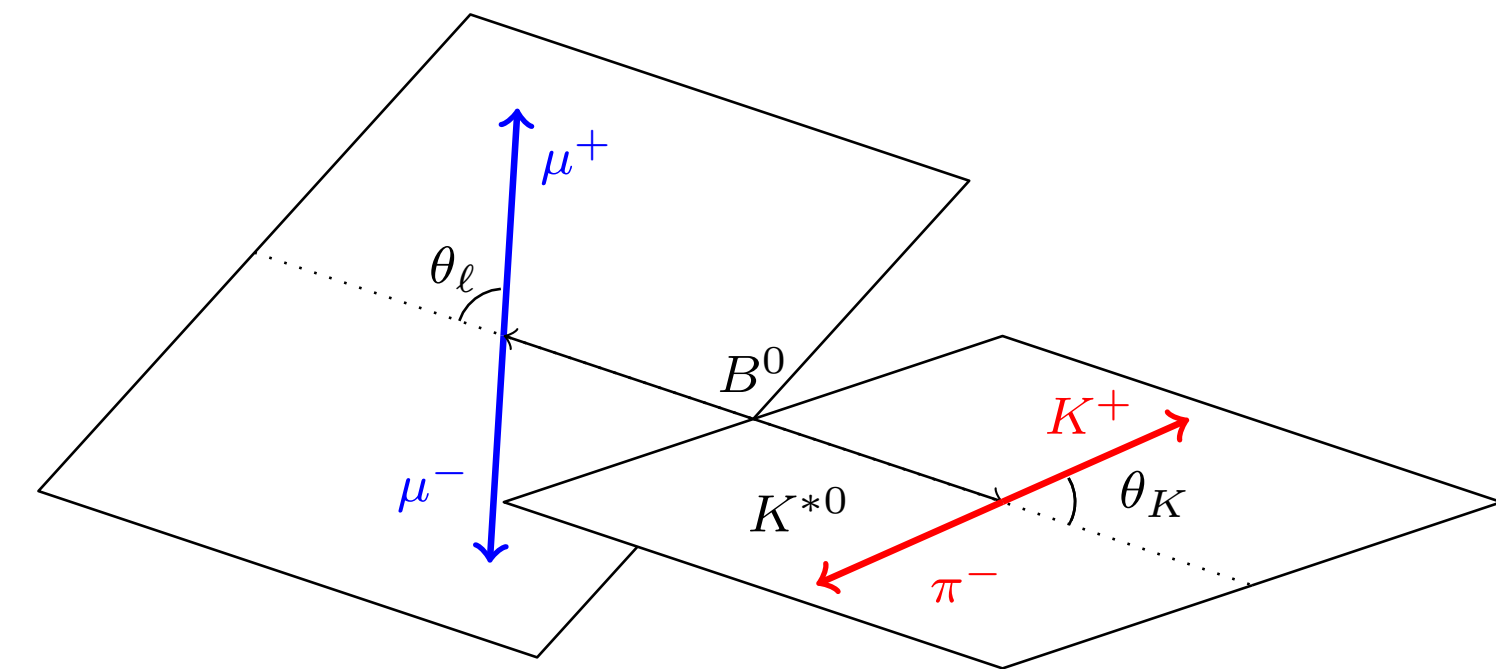
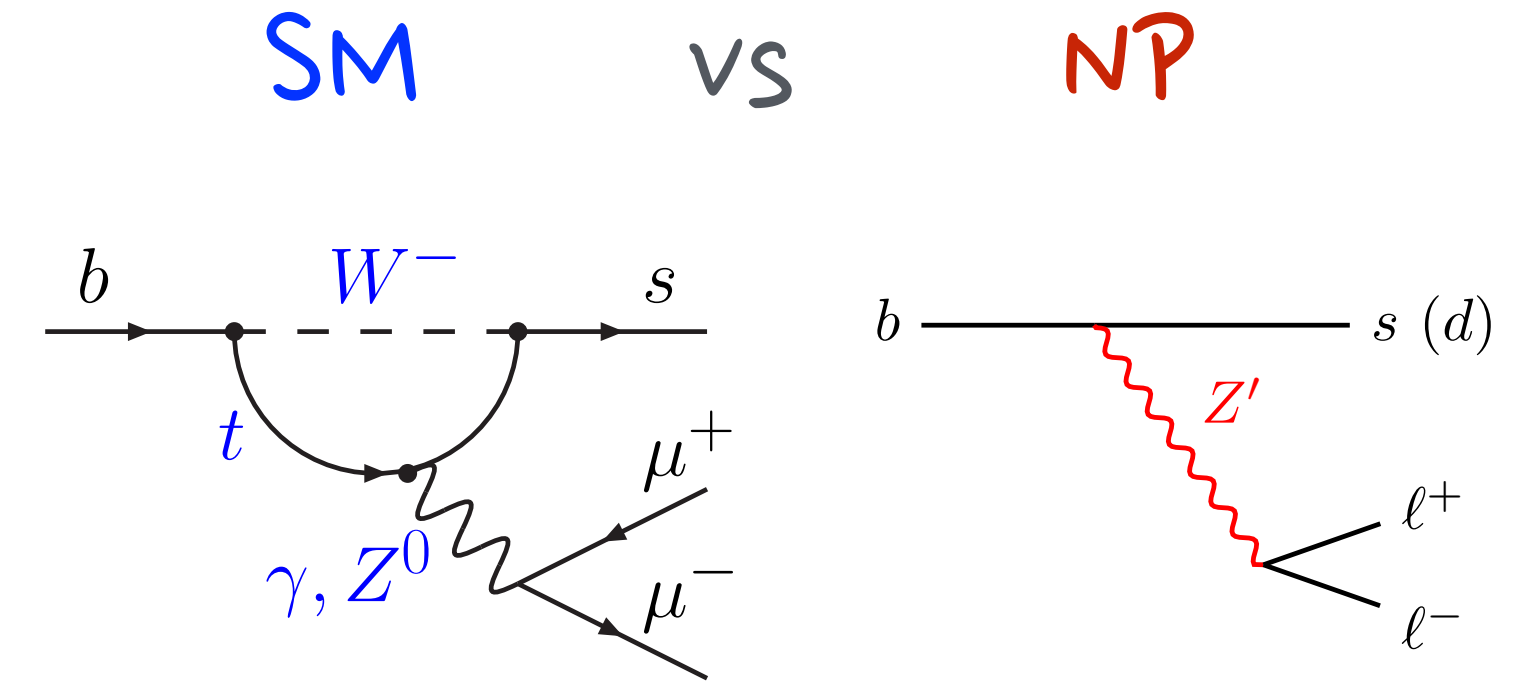
Andrea Mauri

On behalf of the LHCb collaboration

CKM 2023, Santiago de Compostela, 18-22 September 2023

Motivation

- $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ mediated by FCNC
 - ▶ suppressed in the SM \rightarrow sensitive to NP
 - ▶ rich angular structure
 - ▶ NP can alter branching ratio and ang. distributions

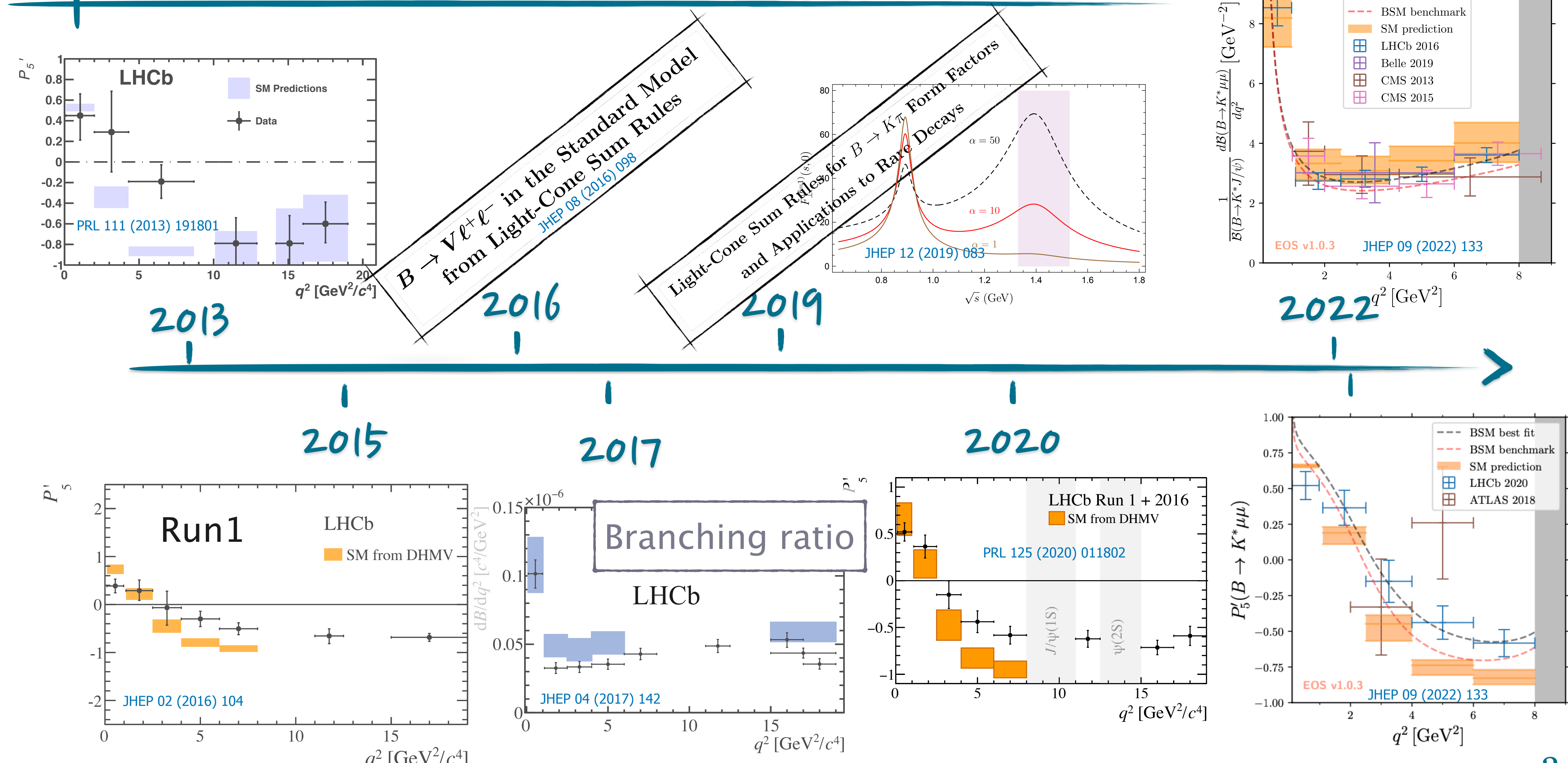


$$\frac{d^4\Gamma[B^0 \rightarrow K^{*0} \mu^+ \mu^-]}{dq^2 d\vec{\Omega}} = \frac{9}{32\pi} \sum_i I_i(q^2) f_i(\vec{\Omega})$$

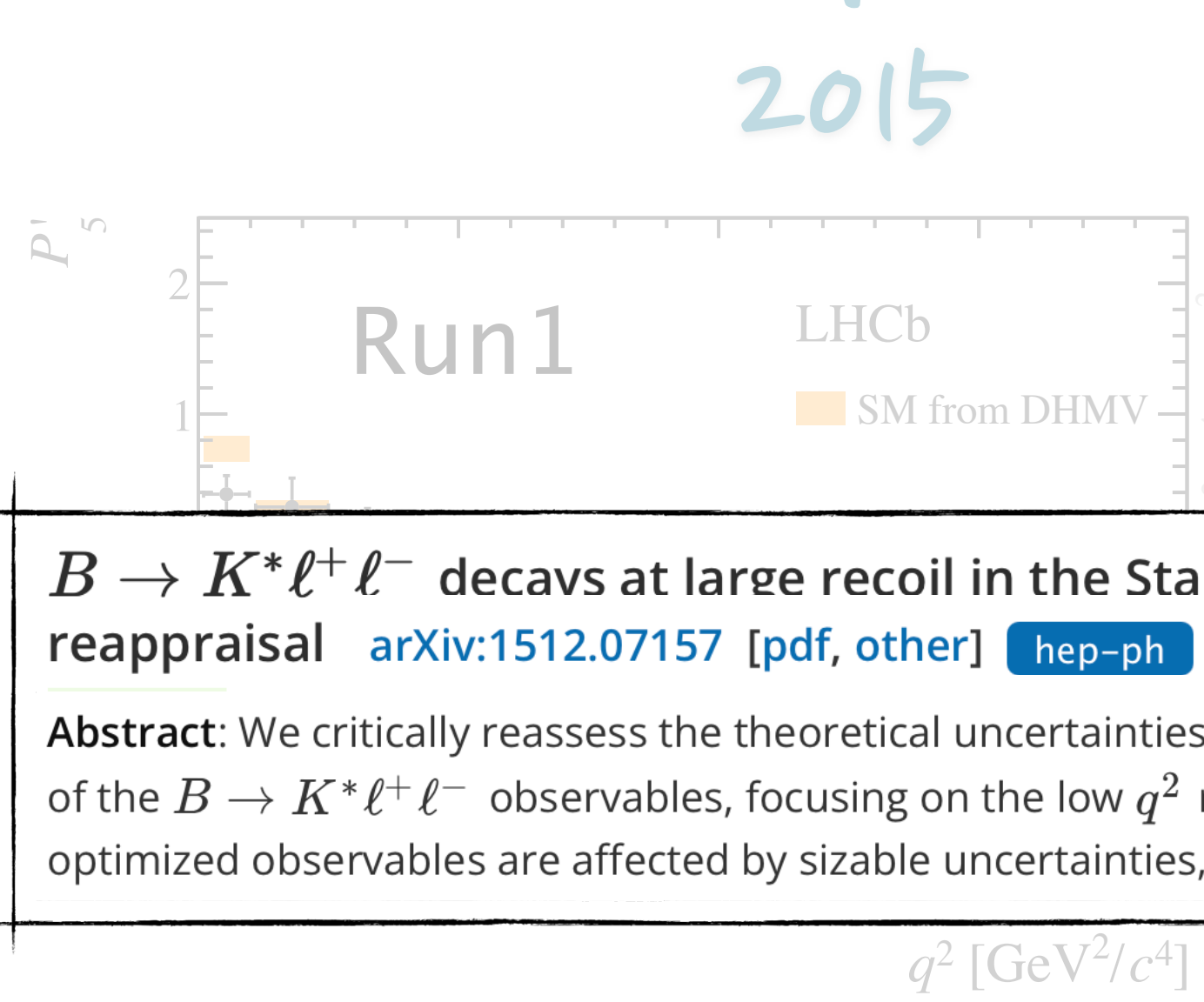
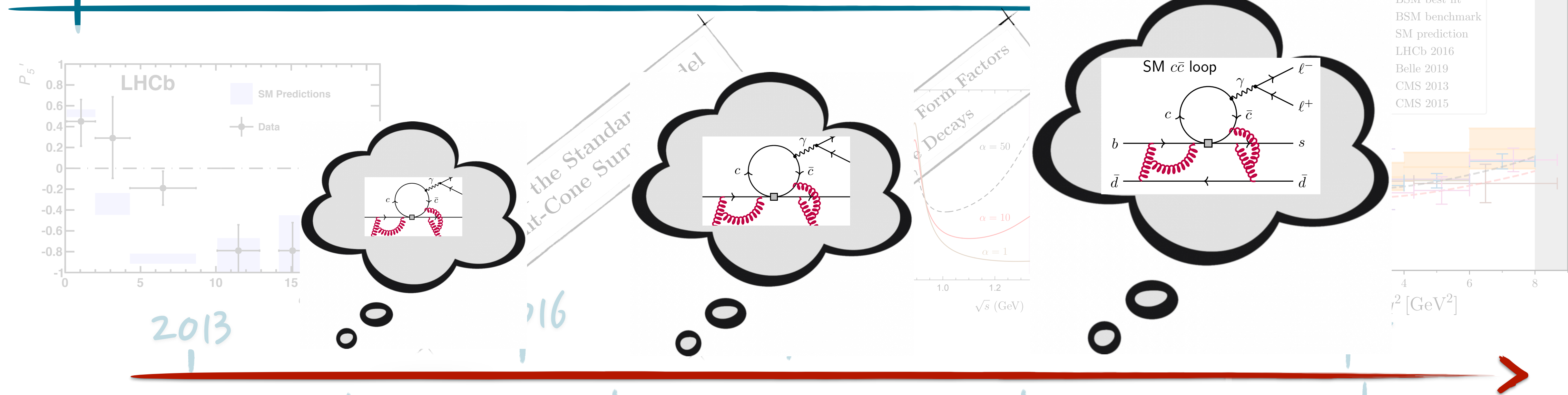
famous P'_5
observable

$$= \frac{9}{32\pi} \left[I_{1s} \sin^2 \theta_K + I_{1c} \cos^2 \theta_K + I_{2s} \sin^2 \theta_K \cos 2\theta_\ell + I_{2c} \cos^2 \theta_K \cos 2\theta_\ell + I_3 \sin^2 \theta_K \sin^2 \theta_\ell \cos 2\phi + I_4 \sin 2\theta_K \sin 2\theta_\ell \cos \phi + I_5 \sin 2\theta_K \sin \theta_\ell \cos \phi + I_6 \sin^2 \theta_K \cos \theta_\ell + I_7 \sin 2\theta_K \sin \theta_\ell \sin \phi + I_8 \sin 2\theta_K \sin 2\theta_\ell \sin \phi + I_9 \sin^2 \theta_K \sin^2 \theta_\ell \sin 2\phi \right],$$

Why we are here...a bit of history

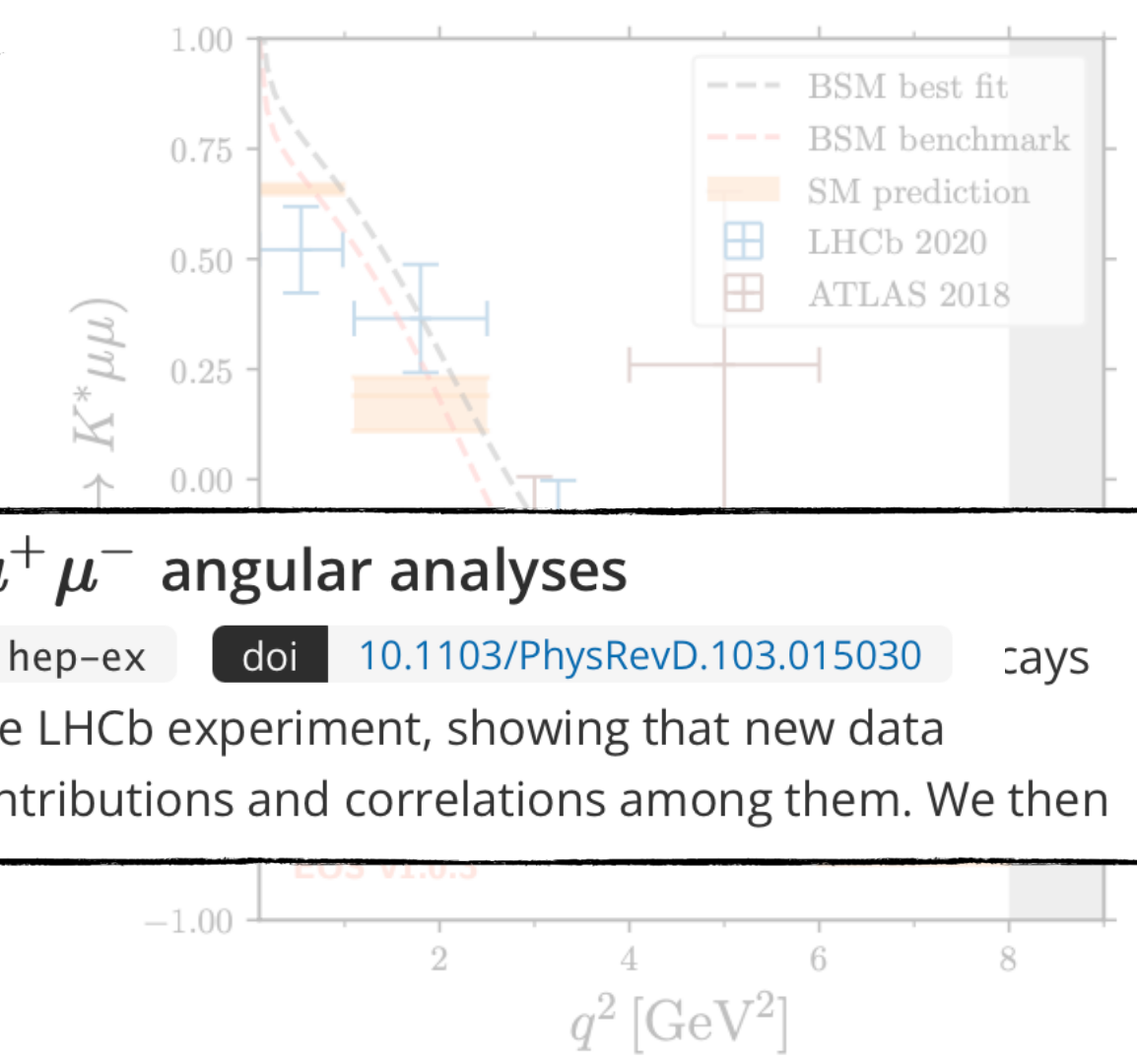


Why we are here...a bit of history



Hadronic uncertainties in the $B \rightarrow K^* \mu^+ \mu^-$ decay [arXiv:1809.03789](https://arxiv.org/abs/1809.03789)

Abstract: Motivated by the persisting 'anomaly' in the measurement of P_5' , we review hadronic uncertainties entering the angular observables of the decay $\bar{B} \rightarrow \bar{K}^* \mu^+ \mu^-$. We argue that hadronic uncertainties could account for the present measurements. We discuss how to



$B \rightarrow K^* \ell^+ \ell^-$ decays at large recoil in the Standard Model: a theoretical reappraisal [arXiv:1512.07157](https://arxiv.org/abs/1512.07157) [pdf, other] [hep-ph](https://arxiv.org/abs/1512.07157) doi [10.1007/JHEP06\(2016\)116](https://arxiv.org/abs/1512.07157)

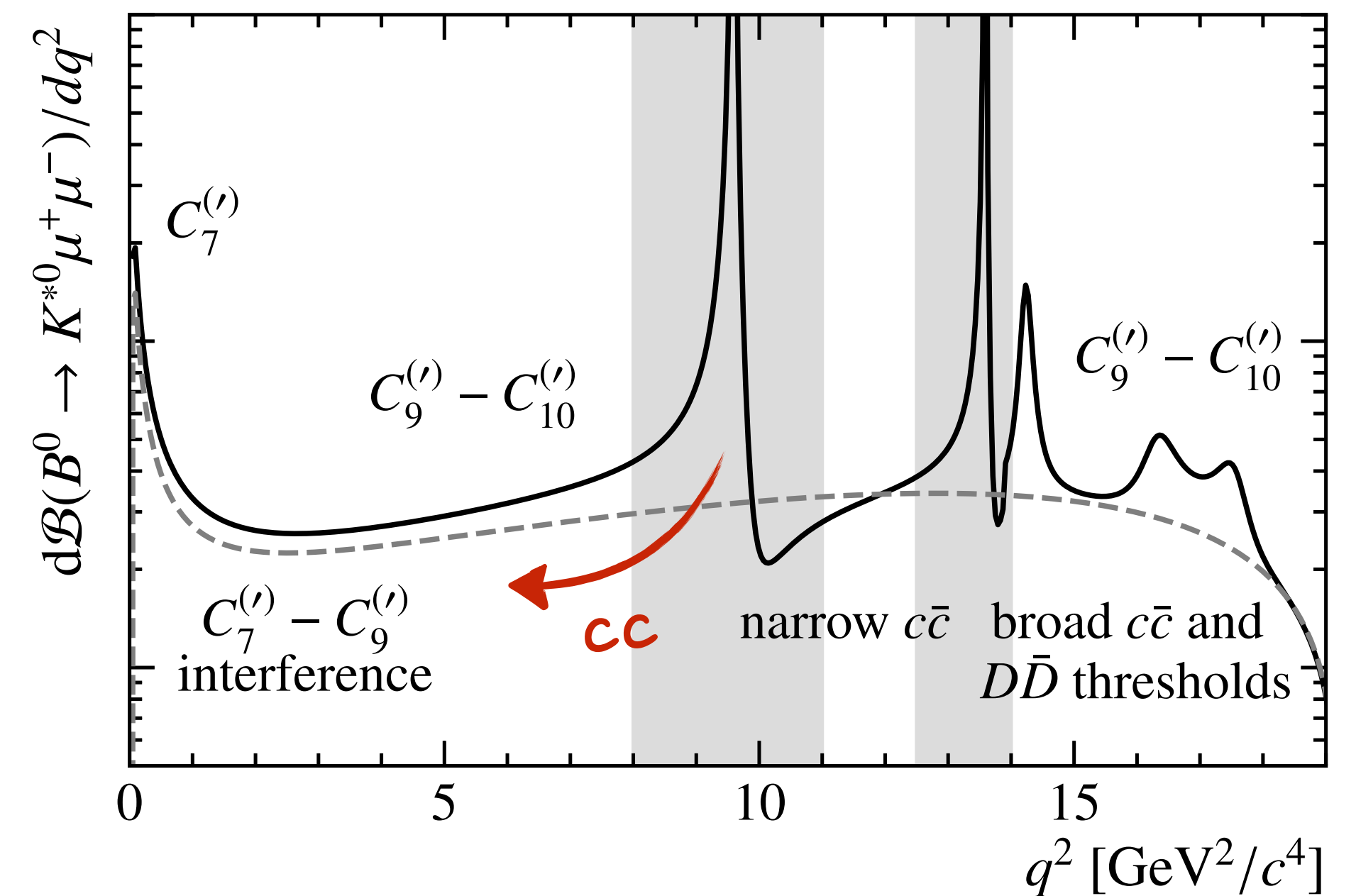
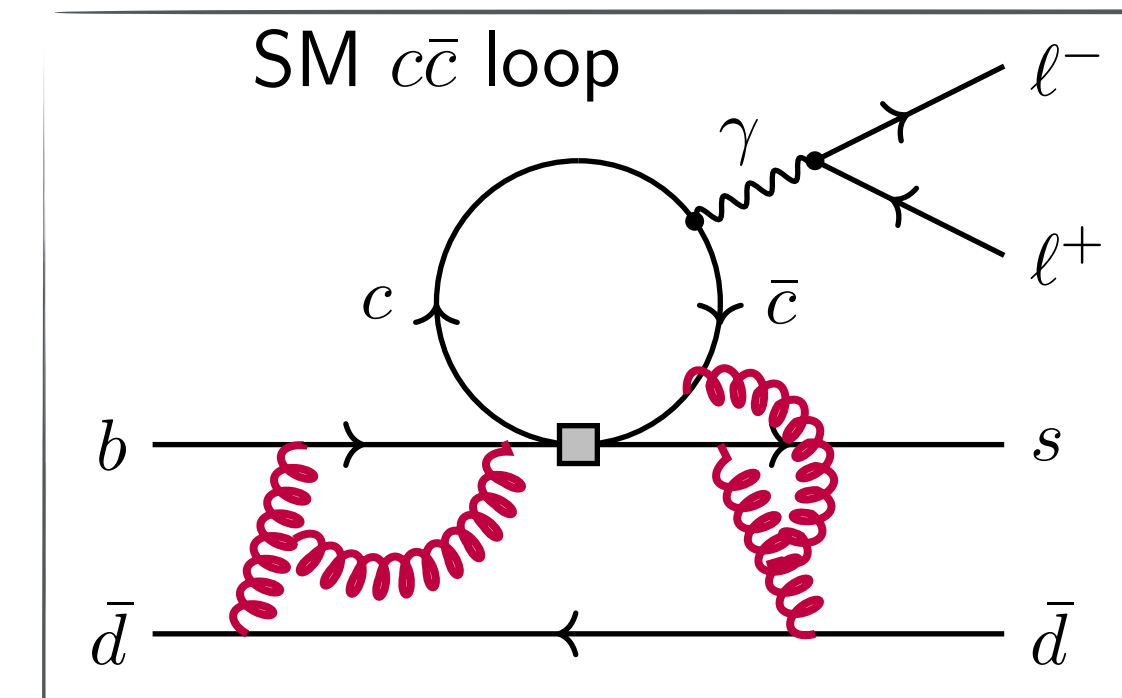
Abstract: We critically reassess the theoretical uncertainties in the Standard Model calculation of the $B \rightarrow K^* \ell^+ \ell^-$ observables, focusing on the low q^2 region. We point out that even optimized observables are affected by sizable uncertainties, since hadronic contributions

Lessons from the $B^{0,+} \rightarrow K^{*0,+} \mu^+ \mu^-$ angular analyses [arXiv:2011.01212](https://arxiv.org/abs/2011.01212) [pdf, other] [hep-ph](https://arxiv.org/abs/2011.01212) [hep-ex](https://arxiv.org/abs/2011.01212) doi [10.1103/PhysRevD.103.015030](https://arxiv.org/abs/2011.01212) says

in light of the recent measurements from the LHCb experiment, showing that new data strengthen the need for sizable hadronic contributions and correlations among them. We then

Hadronic uncertainties...

- Non-local hadronic contribution “*charm-loop*”
 - ▶ Difficult to calculate reliably from first principles
 - ▶ Can mimic NP
 - ▶ Can we access it from data?



- ❖ Important developments on the theory side (see talk by Meril...)
- ❖ SM prediction available @ $q^2 < 0$

Unbinned amplitude analysis

- Perform q^2 unbinned amplitude analysis
 - ▶ model *local* vs *non-local* contributions

non-local hadronic
matrix elements
“charm-loop”

$$A_{\lambda}^{L,R} = \mathcal{N}_{\lambda} \left\{ \left[\underbrace{(C_9 \pm C'_9)}_{\text{Wilson coeff.}} \mp \underbrace{(C_{10} \pm C'_{10})}_{\text{Wilson coeff.}} \right] \underbrace{\mathcal{F}_{\lambda}(q^2)}_{\text{Form Factors}} + \frac{2m_b M_B}{q^2} \left[\underbrace{(C_7 \pm C'_7)}_{\text{Wilson coeff.}} \underbrace{\mathcal{F}_{\lambda}^T(q^2)}_{\text{Form Factors}} - 16\pi^2 \frac{M_B}{m_b} \overline{\mathcal{H}_{\lambda}(q^2)} \right] \right\}$$

$\lambda = \perp, \parallel, 0$

polynomial expansion

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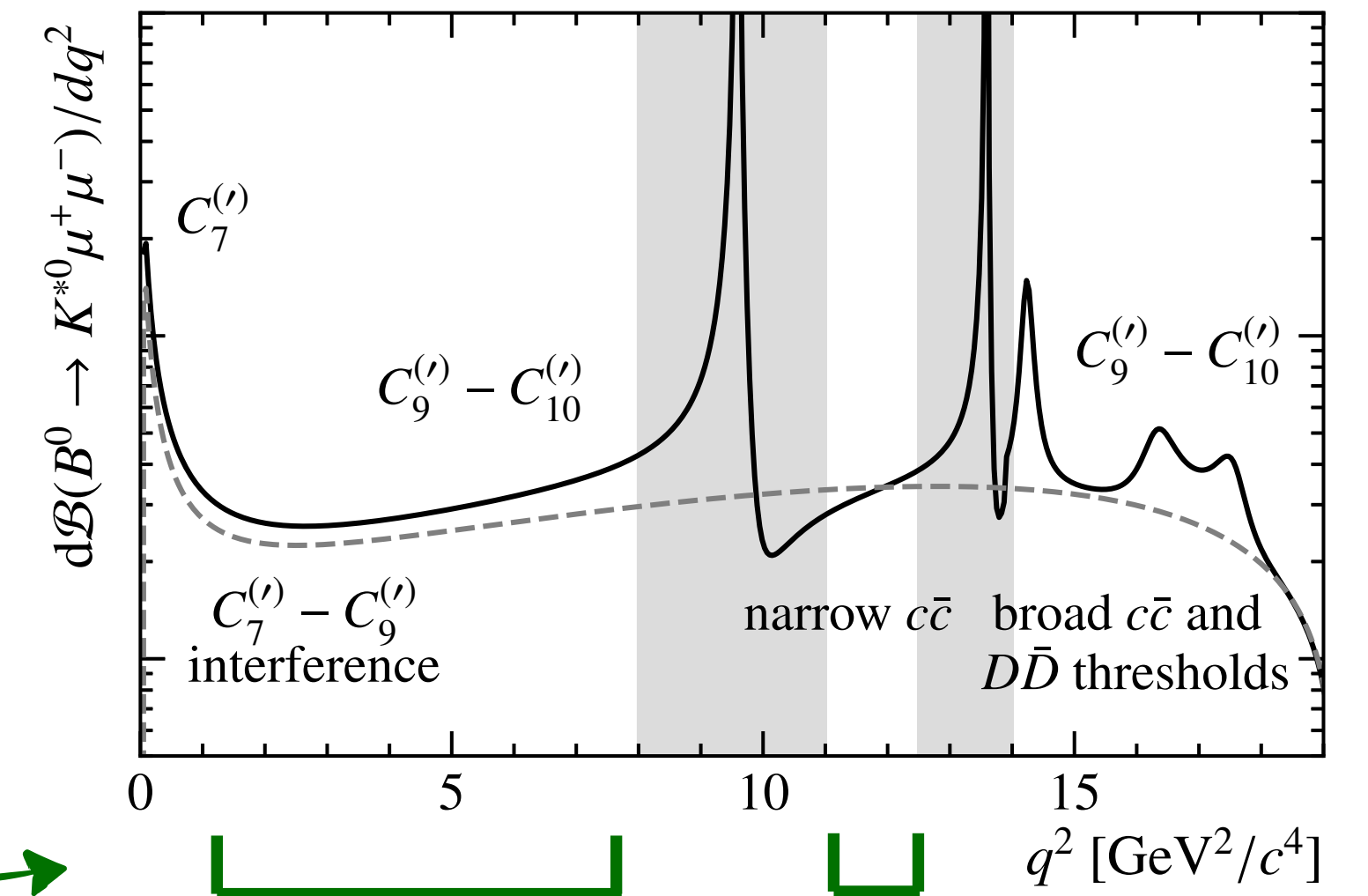
- ▶ **Fit 5-D differential decay rate!**

↳ $q^2, m_{K\pi}^2, \cos \theta_{\ell}, \theta_K, \phi$

$$\mathcal{H}_{\lambda}(z) = \frac{1 - z z_{J/\psi}^*}{z - z_{J/\psi}} \frac{1 - z z_{\psi(2S)}^*}{z - z_{\psi(2S)}} \times \dots \times \sum_n \alpha_{\lambda,n} z^n$$

Analysis overview

- Same dataset of previous LHCb $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ binned angular analysis (Run-I + 2016) 4.7 fb^{-1}
 - ▶ two q^2 regions: $[1.1, 8.0]$ & $[11, 12.5]$ GeV
- Six-dimensional fit
 - ▶ differential decay rate + invariant B mass to separate signal from combinatorial background
- Large number of signal parameters
 - ▶ Wilson coefficients: $C_9, C_{10}, C'_9, C'_{10}$ [floated] + C_7, C'_7 [fixed to SM]
 - ▶ local FF : [constrained to LCSR + latticeQCD] JHEP 01 (2019) 150, PoS LATTICE2014 (2015) 372
 - ▶ non-local hadronic parameters $\mathcal{H}_\lambda(q^2)$ (see next slides)
 - ▶ S-wave (FFs + relative magnitude&phase)



Non-local contributions

- Combining theory & experimental information to constrain charm-loop parameters

Theory information

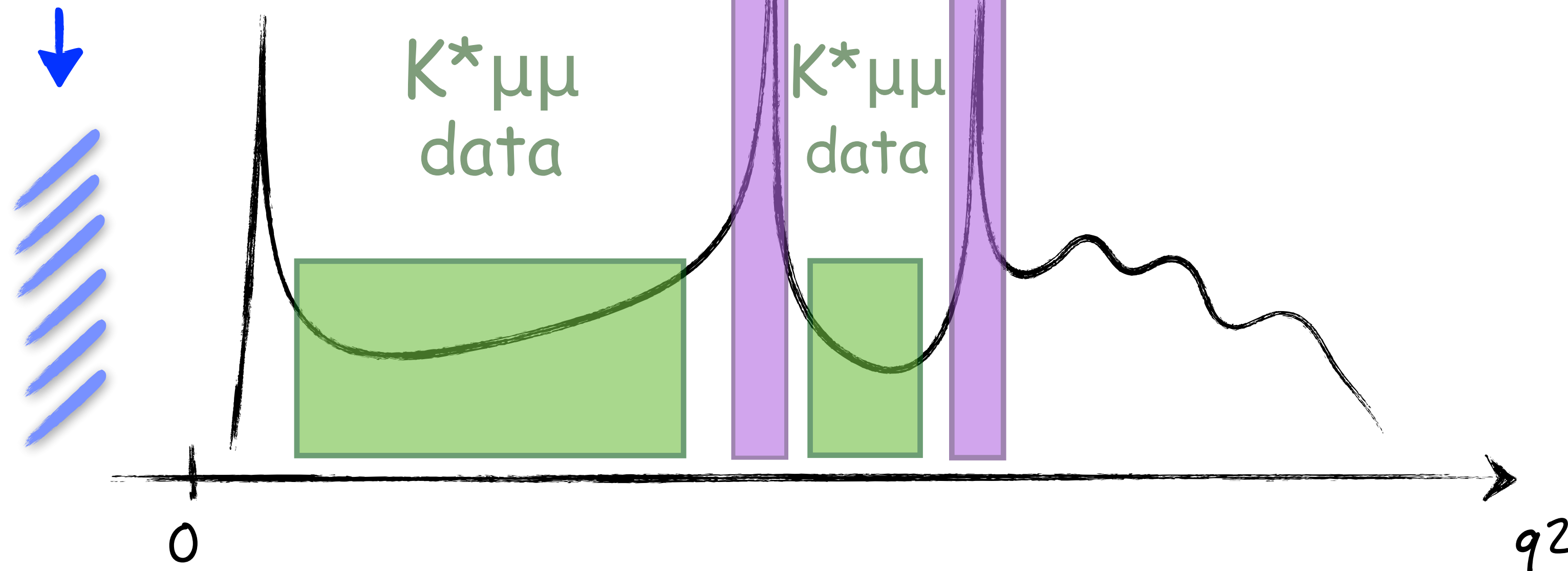
Value of charm-loop at $q^2 < 0$
► reliable for $q^2 \ll 4m_c^2$

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Experimental measurements

Branching ratio, polarization fraction and phase difference from $B^0 \rightarrow \psi_n K^{*0}$

PRD 76 031102(R) (2007)
PRD 88 052002 (2013)
PRD 88 074026 (2013)
PRD 90 112009 (2014)



Non-local contributions

Combining

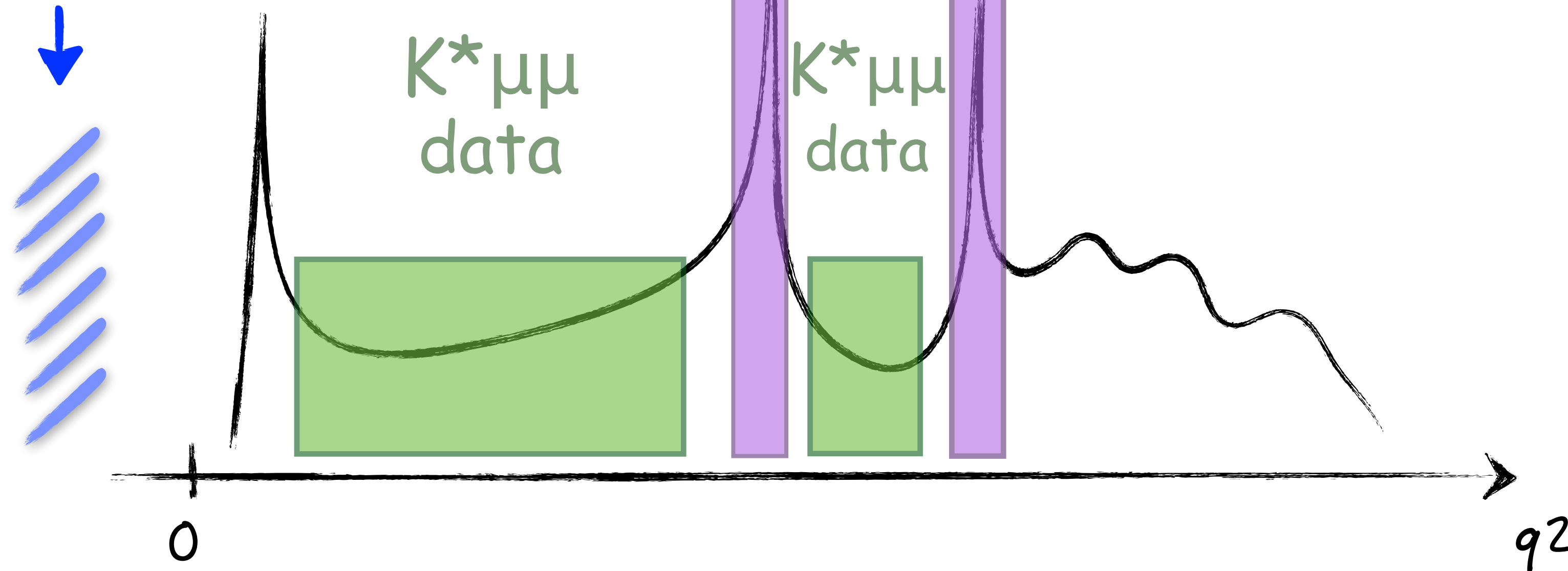
- $q^2 < 0$ *prior*: include theory points @ $q^2 < 0$
- $q^2 > 0$ *only*: exclude theory points @ $q^2 < 0$

Theory info

Value of charm-loop at $q^2 < 0$
 ► reliable for $q^2 \ll 4m_c^2$

phase difference from $B^0 \rightarrow \psi_n K^{*0}$

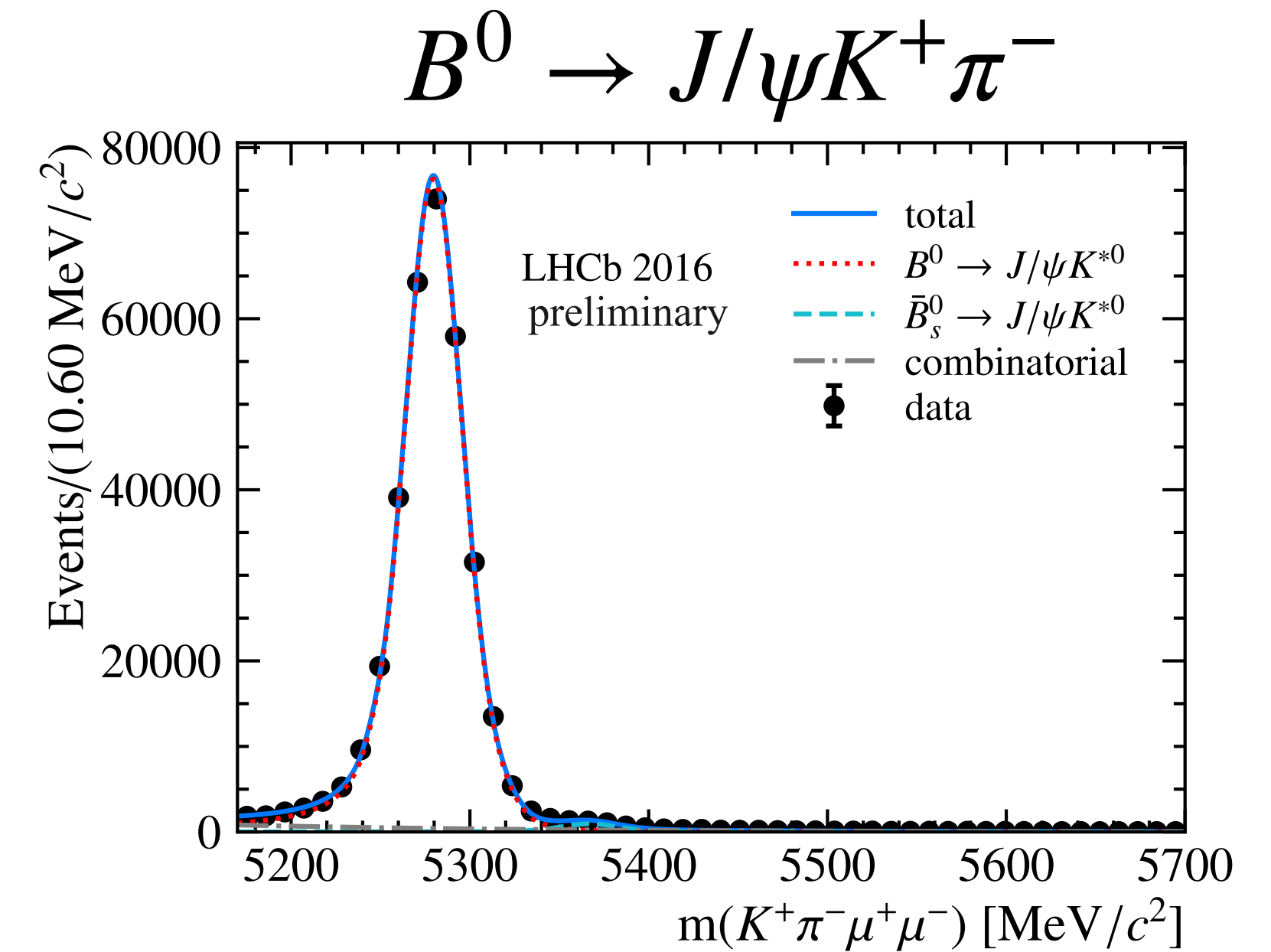
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PRD 76 031102(R) (2007)
 PRD 88 052002 (2013)
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Branching ratio constraint

- Angular information can only access the relative size of Wilson coefficients
 - ▶ Scale of Wilson coeff. set by branching ratio
- Include branching ratio measurement in the analysis
 - ▶ Normalised to $B^0 \rightarrow J/\psi K^+ \pi^-$ decays



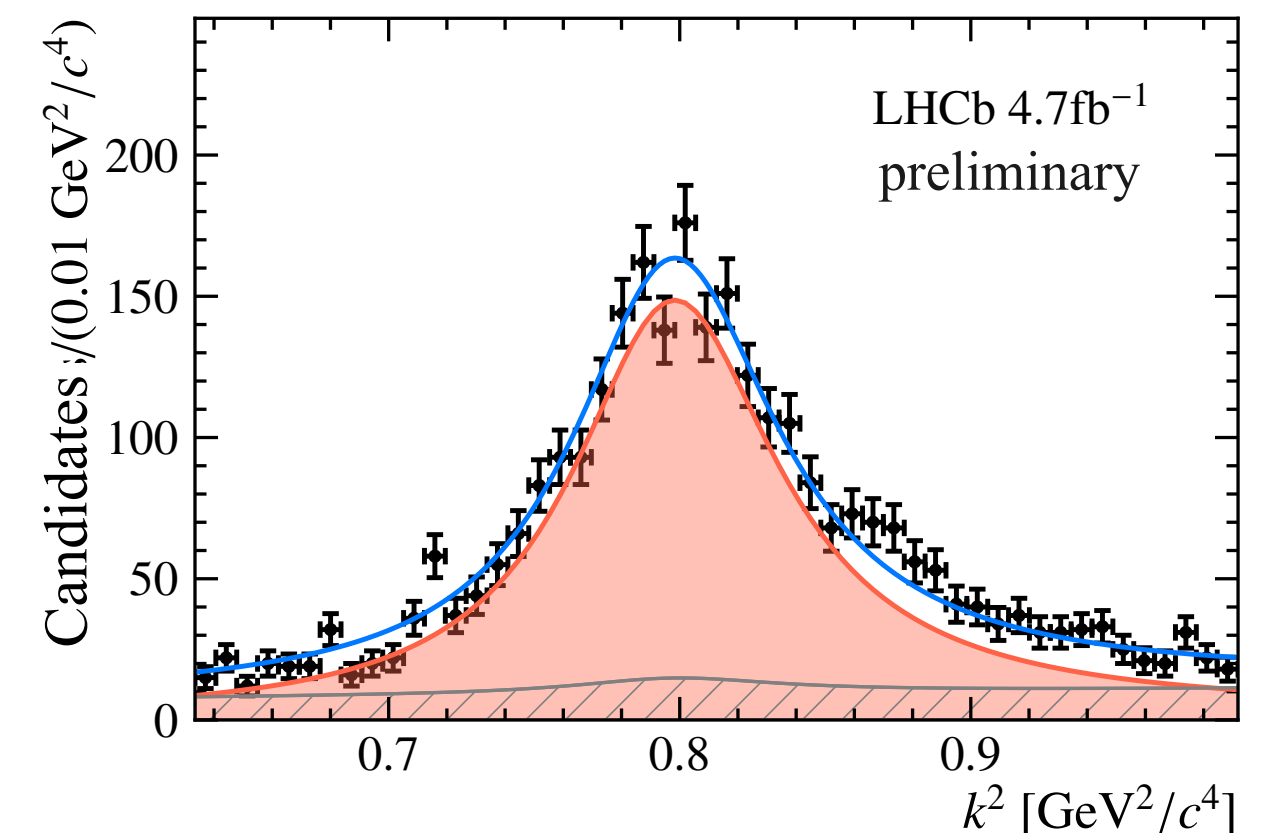
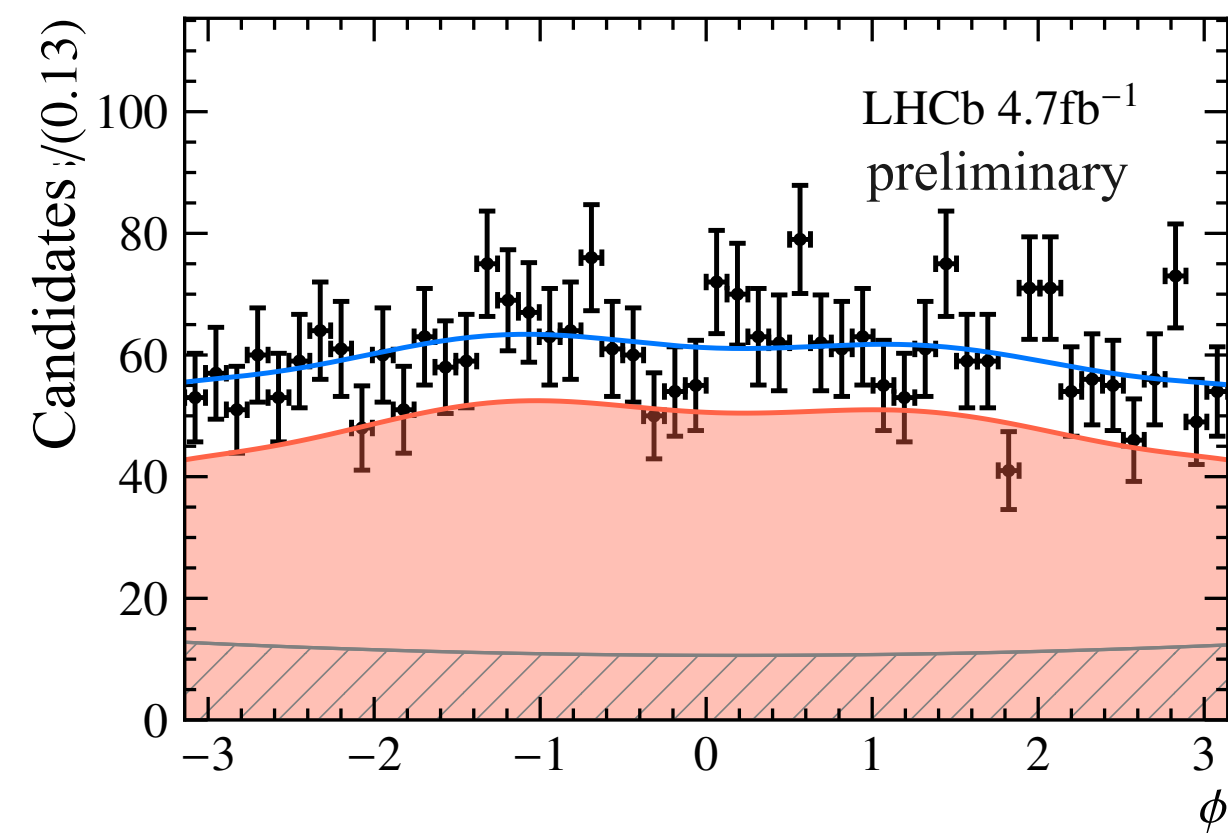
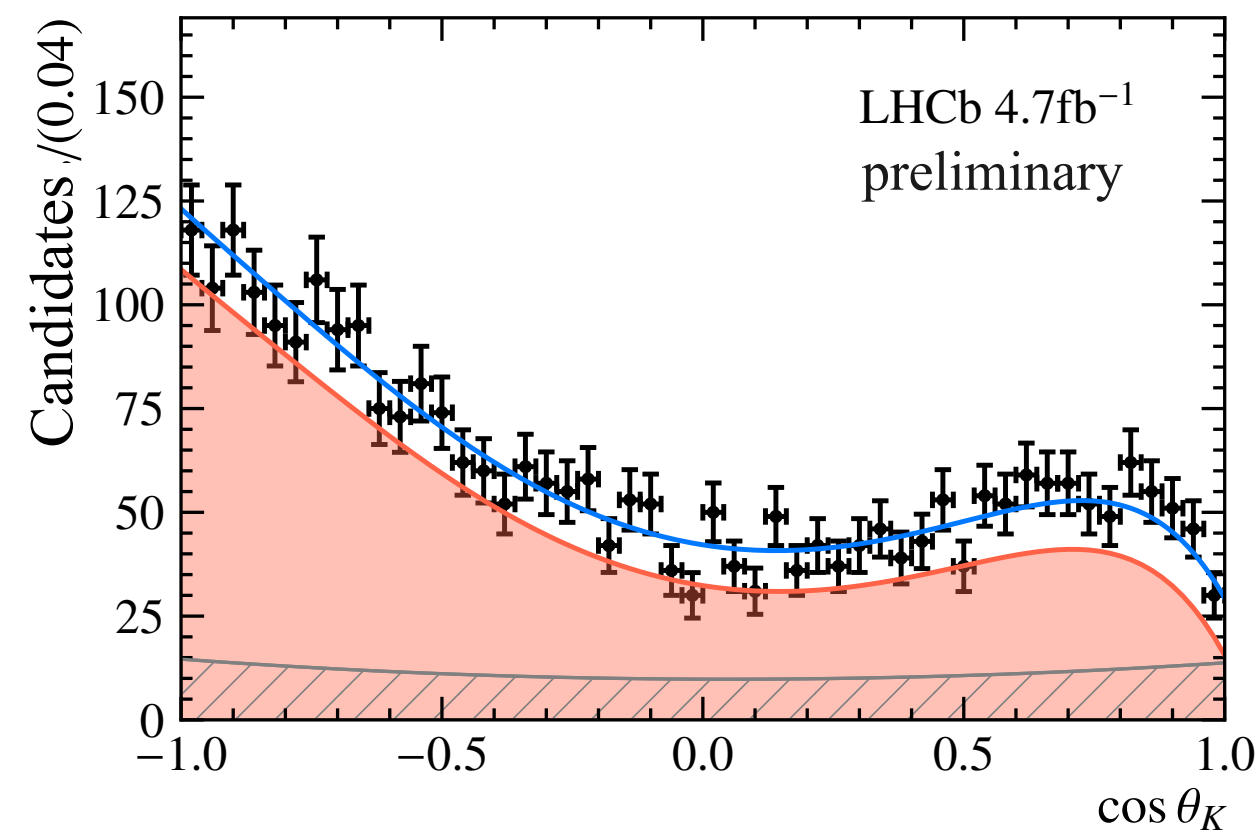
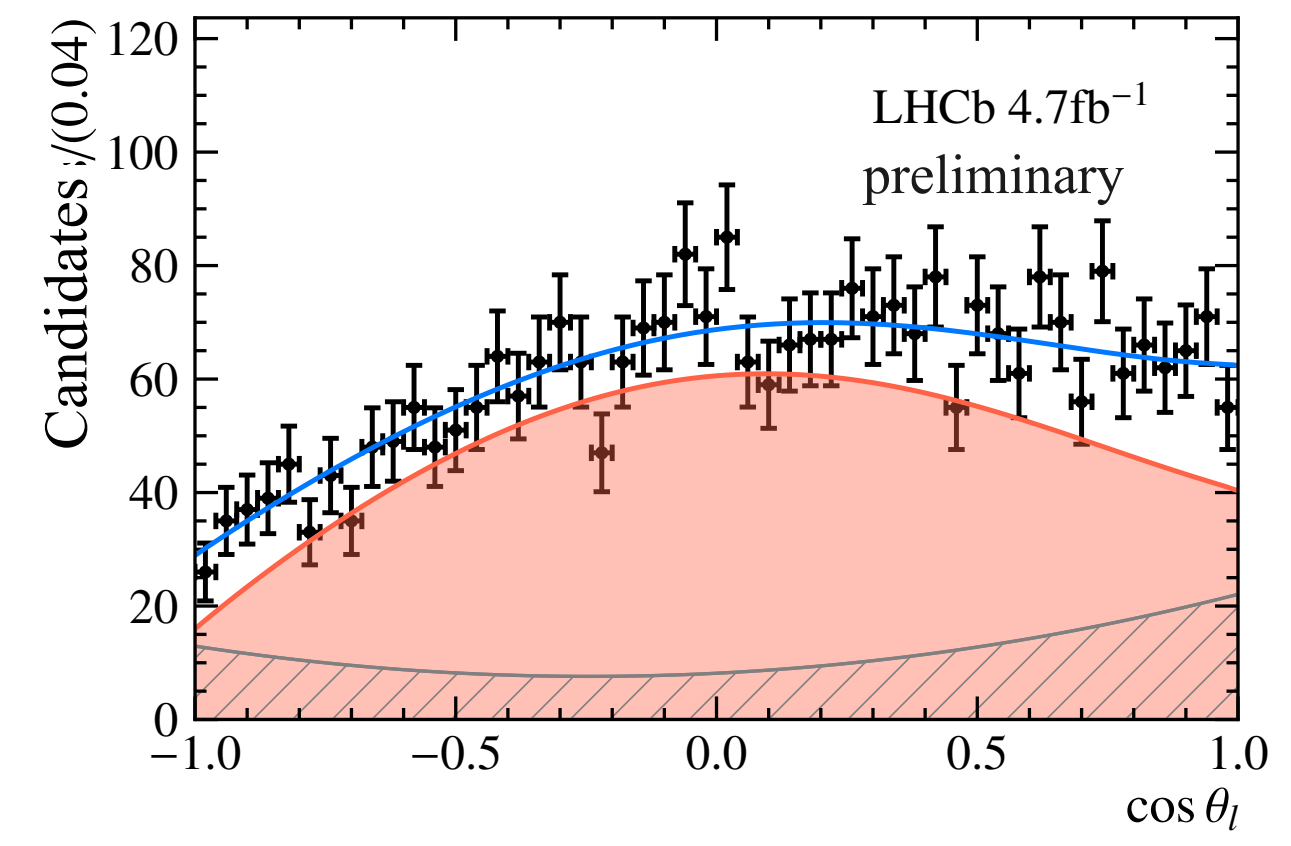
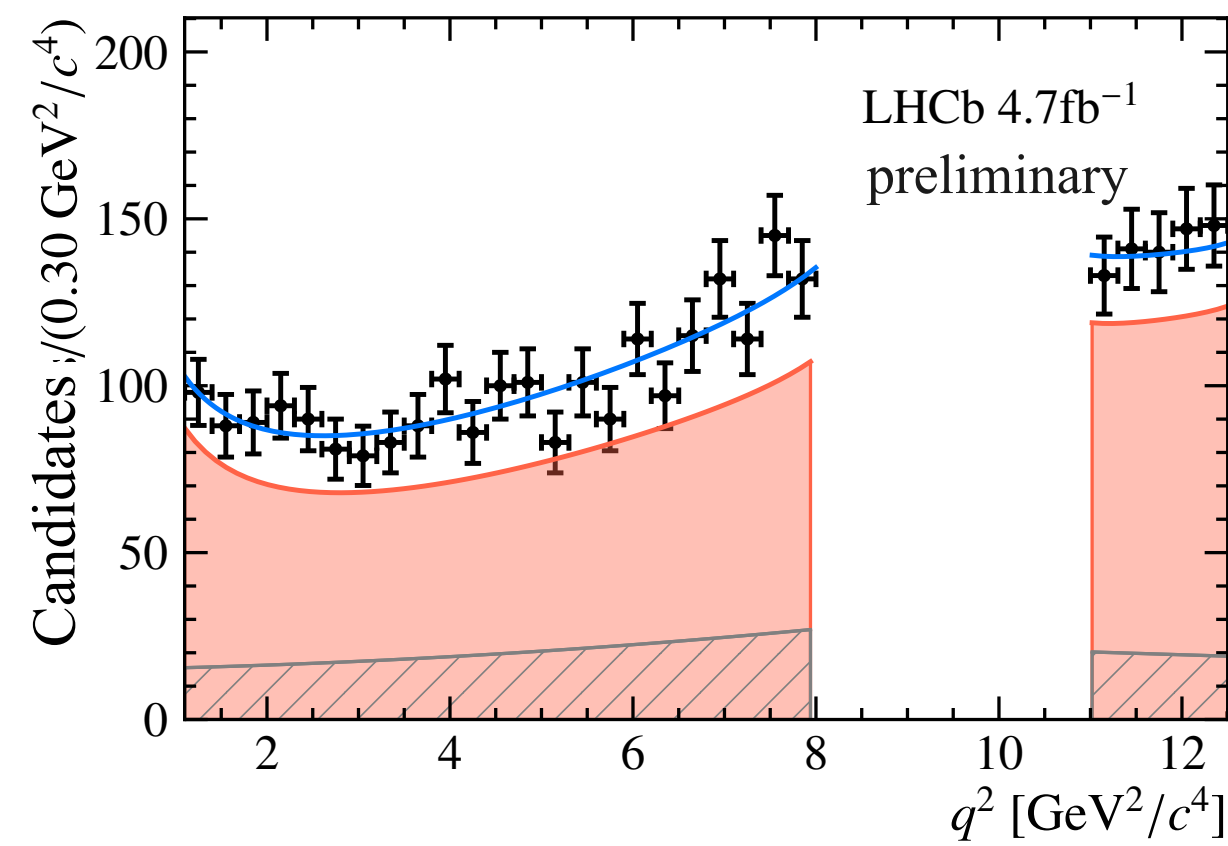
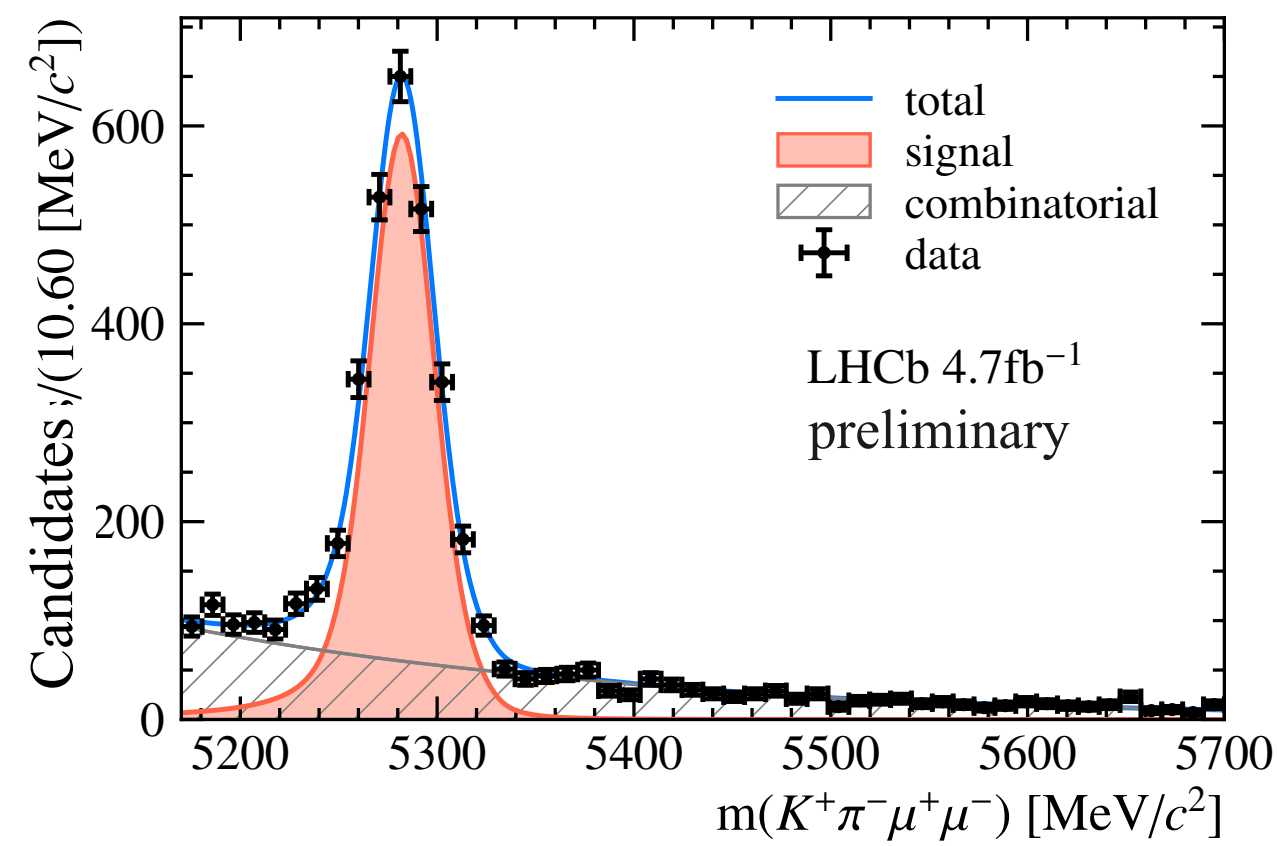
$$N_{sig} = N_{J/\psi K\pi} \times \frac{\mathcal{B}(B^0 \rightarrow K^{*0} \mu^+ \mu^-) \times \frac{2}{3}}{\mathcal{B}(B^0 \rightarrow J/\psi K^+ \pi^-) \times f^{J/\psi K\pi} \times \mathcal{B}(J/\psi \rightarrow \mu^+ \mu^-)} \times R_\epsilon,$$

- ▶ Largest systematic uncertainties:

- ❖ $\mathcal{B}(B^0 \rightarrow J/\psi K^+ \pi^-) = (1.15 \pm 0.01 \pm 0.05) \cdot 10^{-3}$ PRD 90 (2014) 112009
- ❖ $f_{\pm 100\text{MeV}}^{B^0 \rightarrow J/\psi K\pi} = 0.644 \pm 0.010$

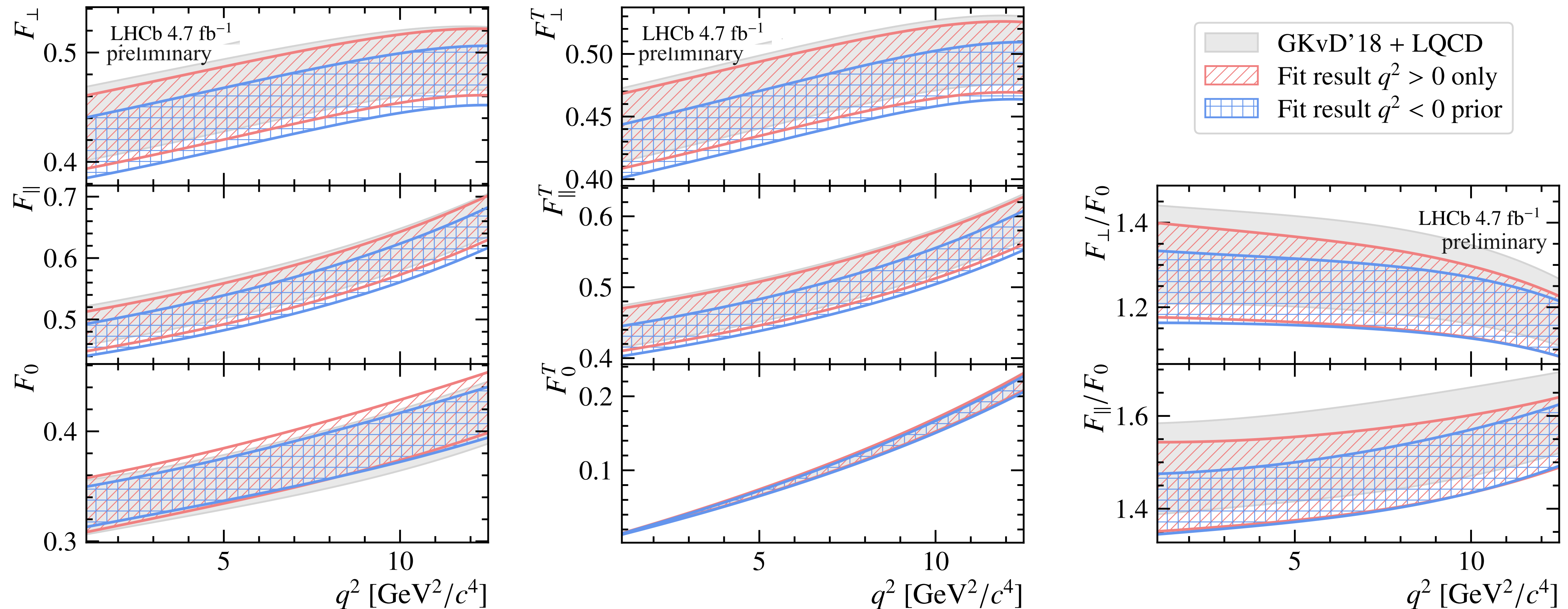
Fit projections

- Simultaneous fit to Run1 & 2016, $q^2 \in [1.1, 8.0] \text{ \& } [11, 12.5]$



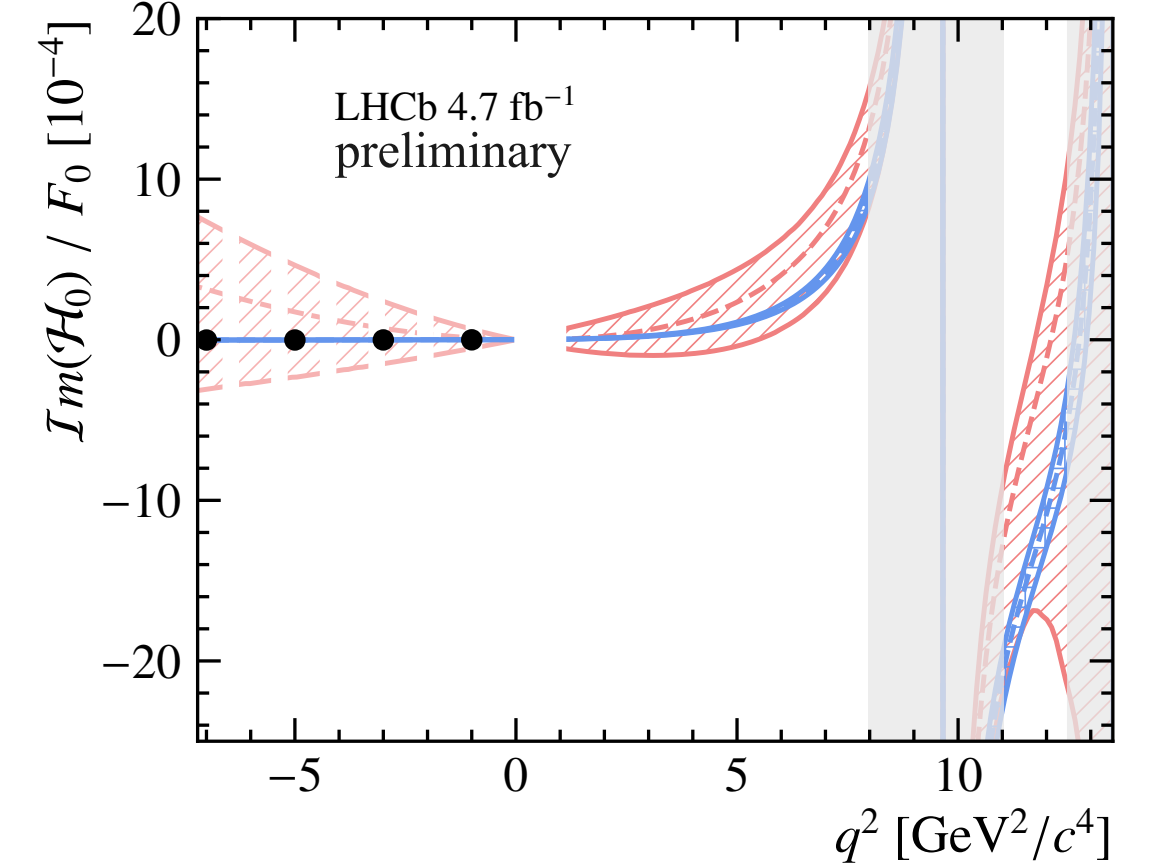
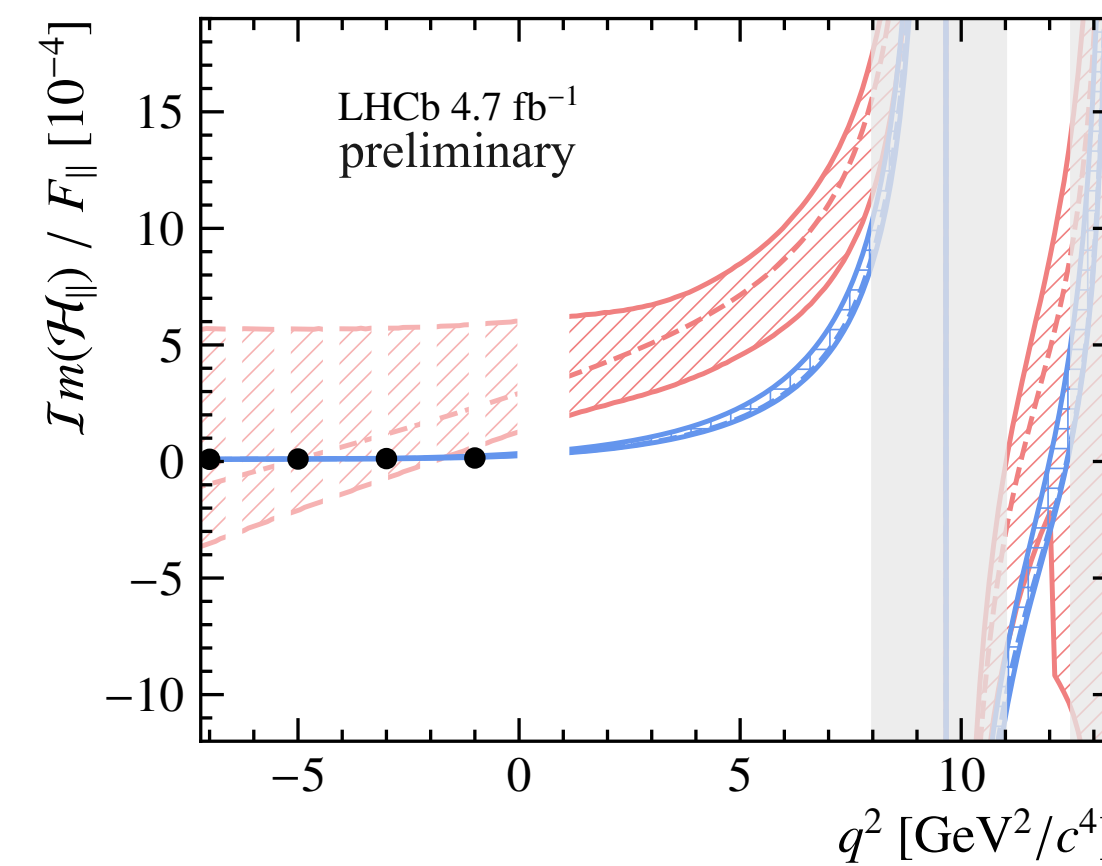
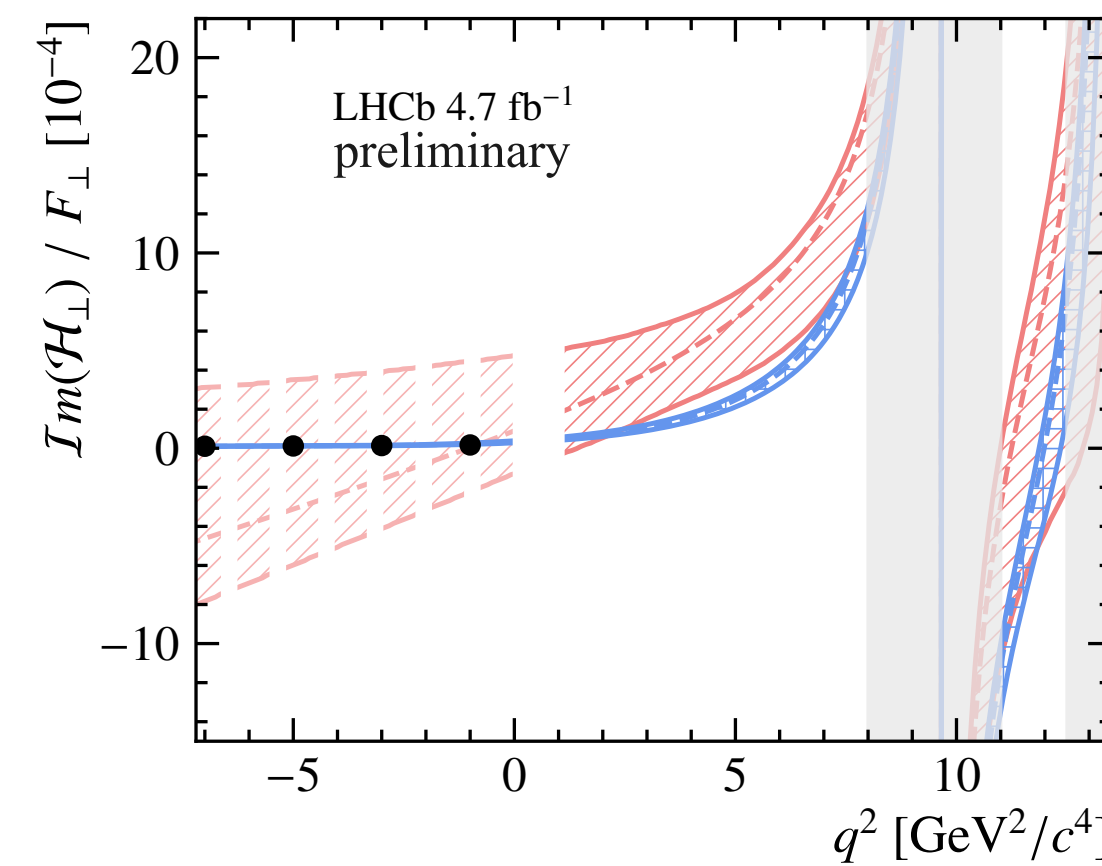
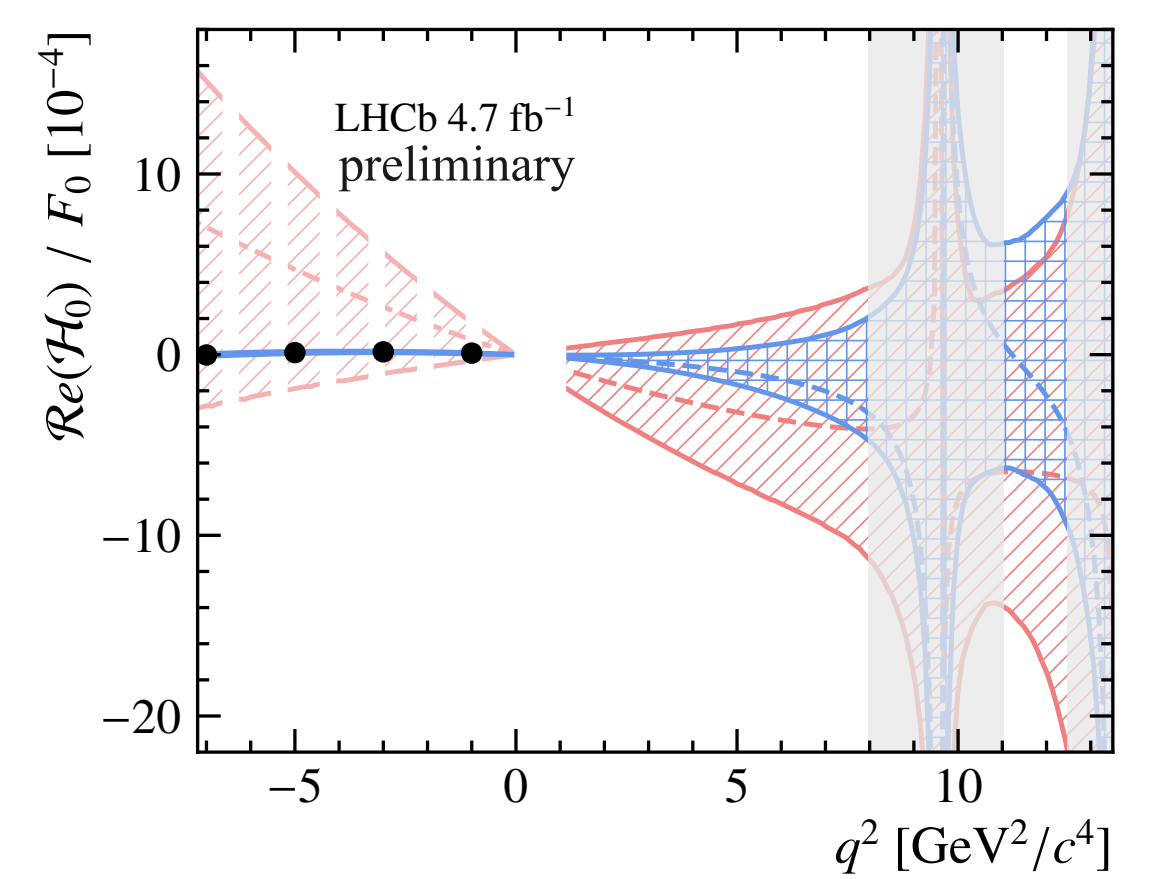
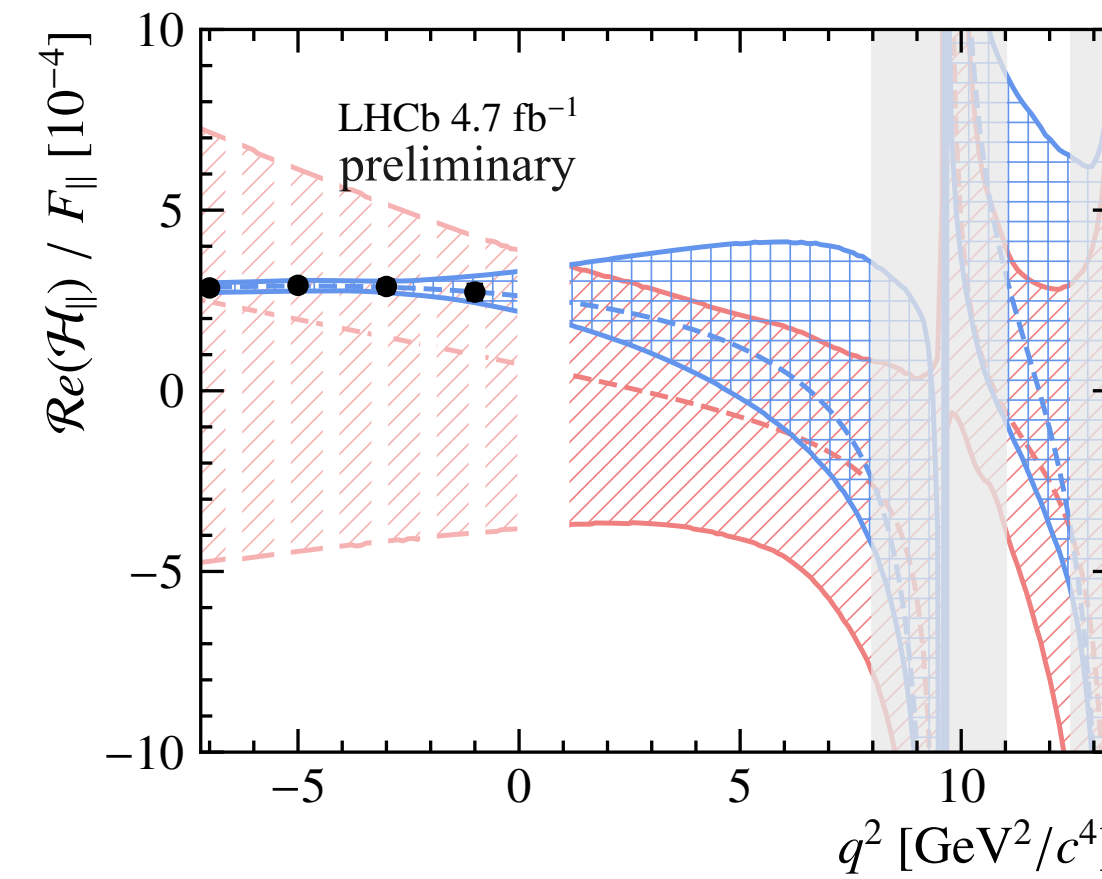
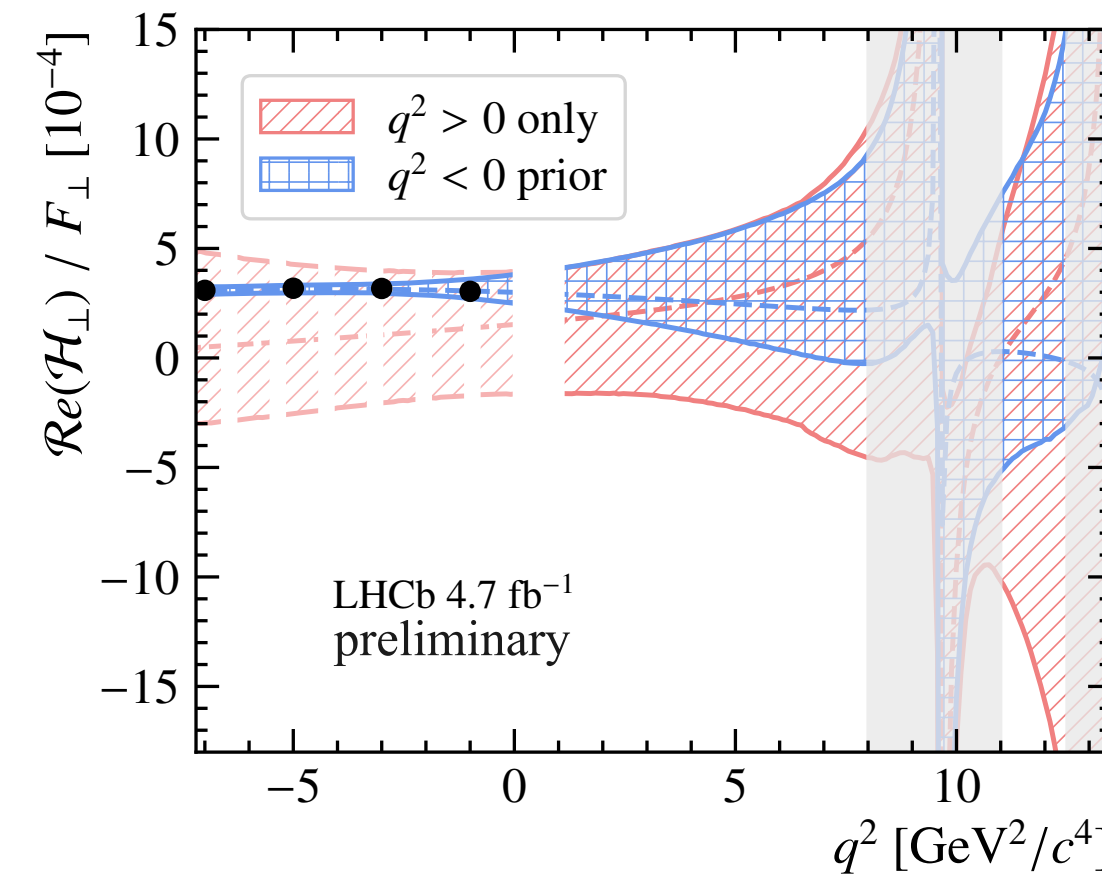
Form factor results

- Dominant uncertainty in $b \rightarrow s\ell\ell$ SM branching ratio prediction
- Fit results are found to require small adjustment in $\mathcal{F}_{\perp,\parallel}/\mathcal{F}_0$ ratio



Non-local hadronic results

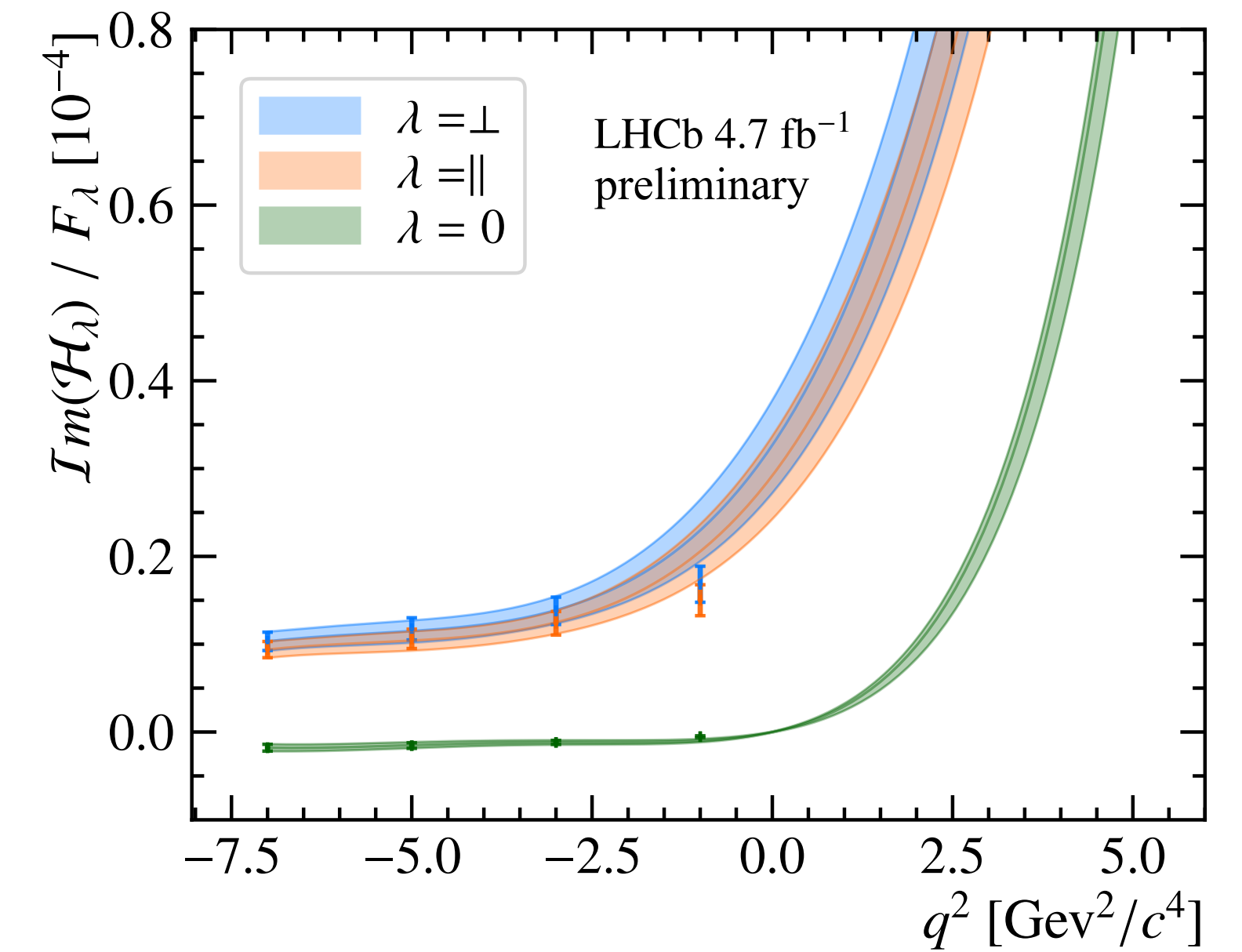
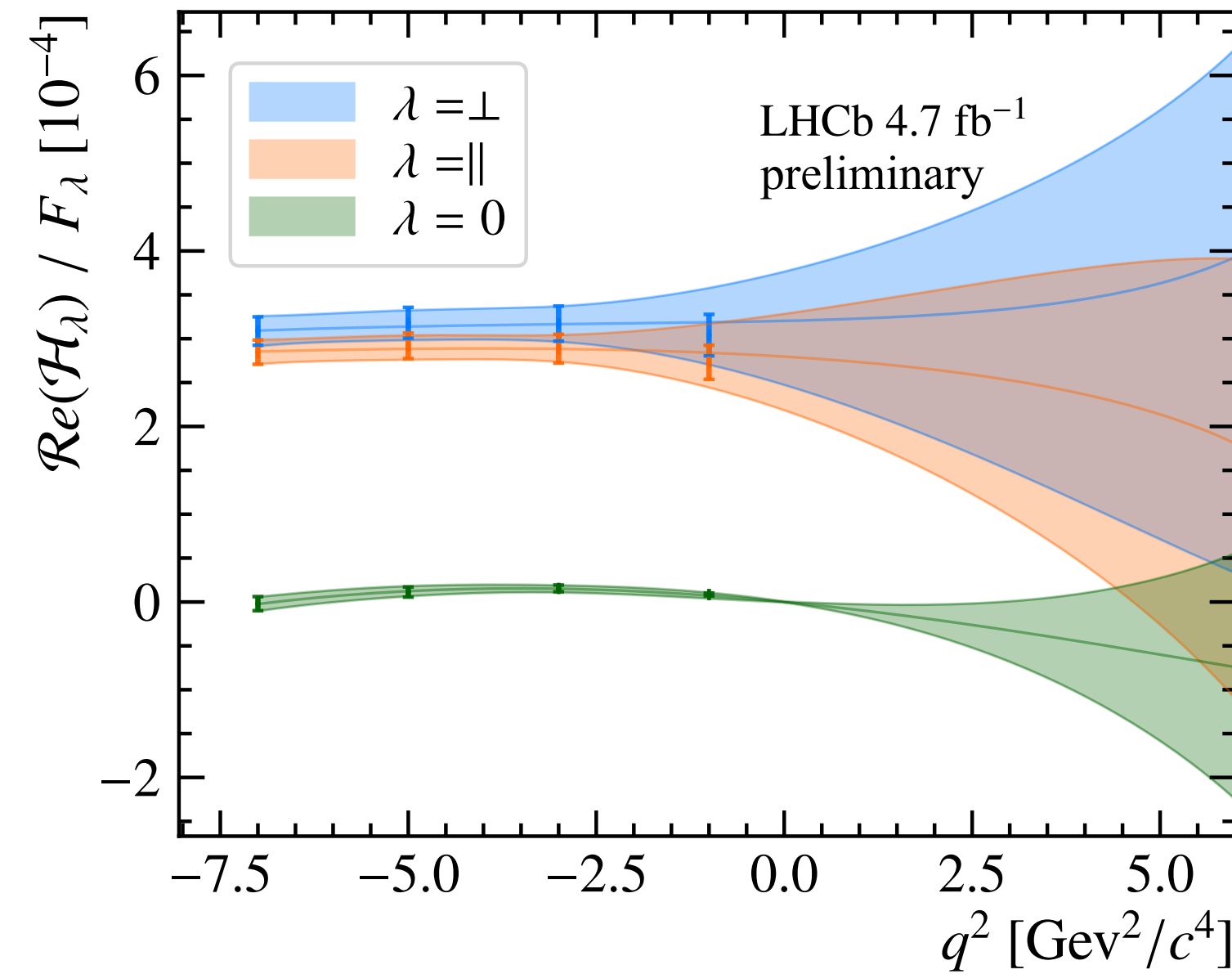
- Good overall agreement between the two configurations
- ▶ Small discrepancy in $Im \cdot \mathcal{H}_{\parallel}(q^2)$



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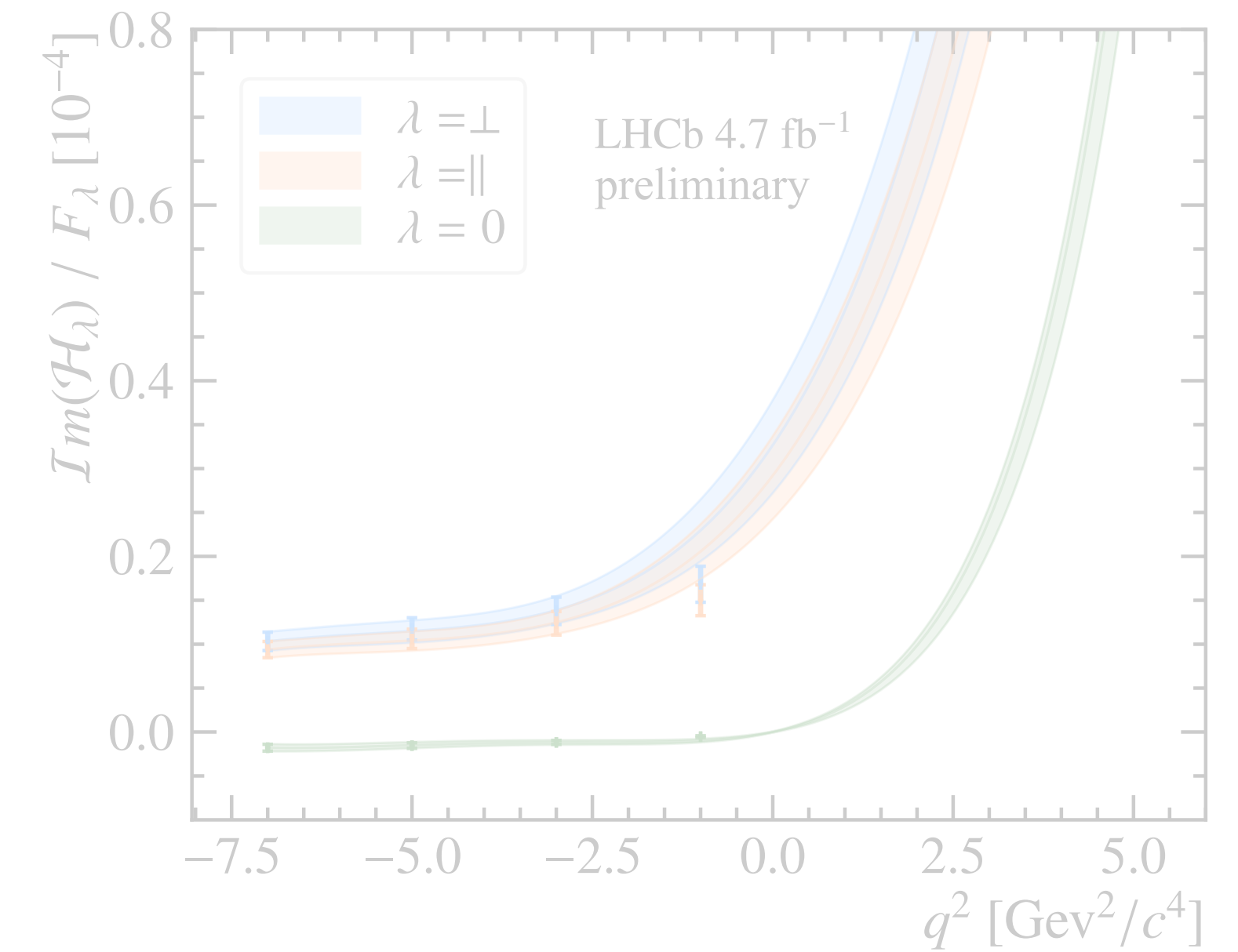
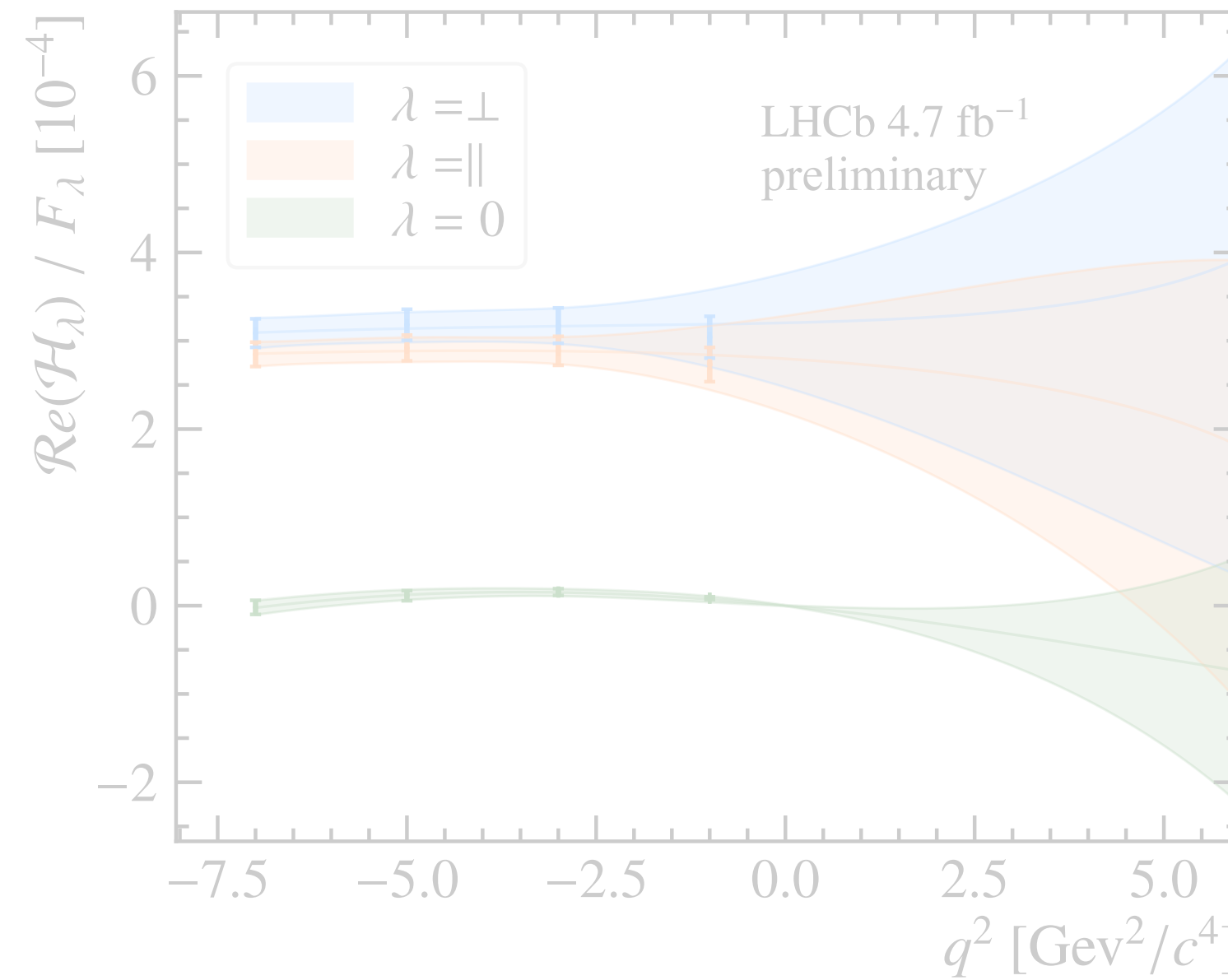


- $\mathcal{H}_\lambda(q^2)$ constrained from 3 theory points $q^2 \in [-7, -5, -3] \text{ GeV}^2$
- Result can be tested at $q^2 = -1 \text{ GeV}^2$
 - ▶ Imaginary part tends to rise more rapidly than predictions

Non-local hadronic results

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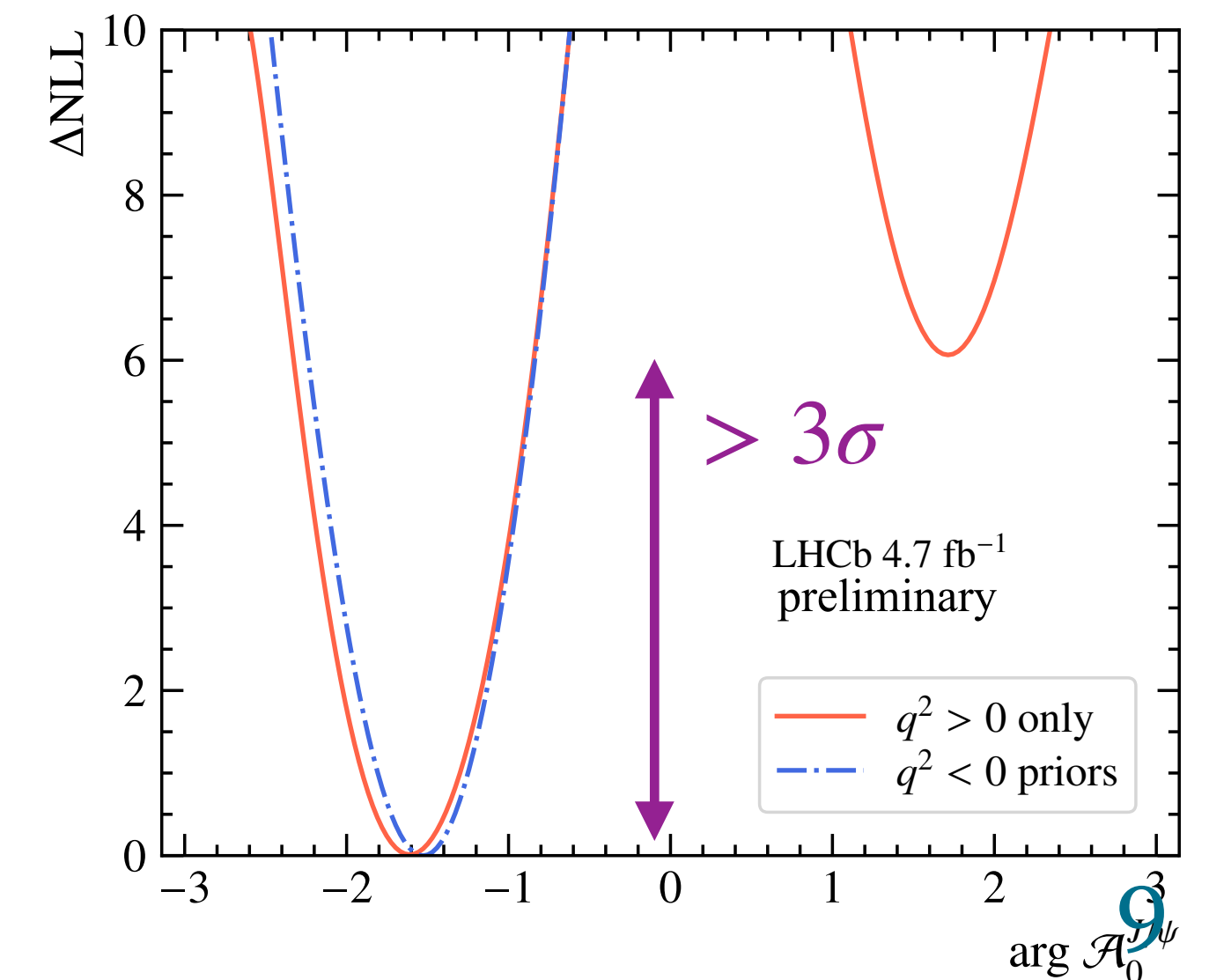
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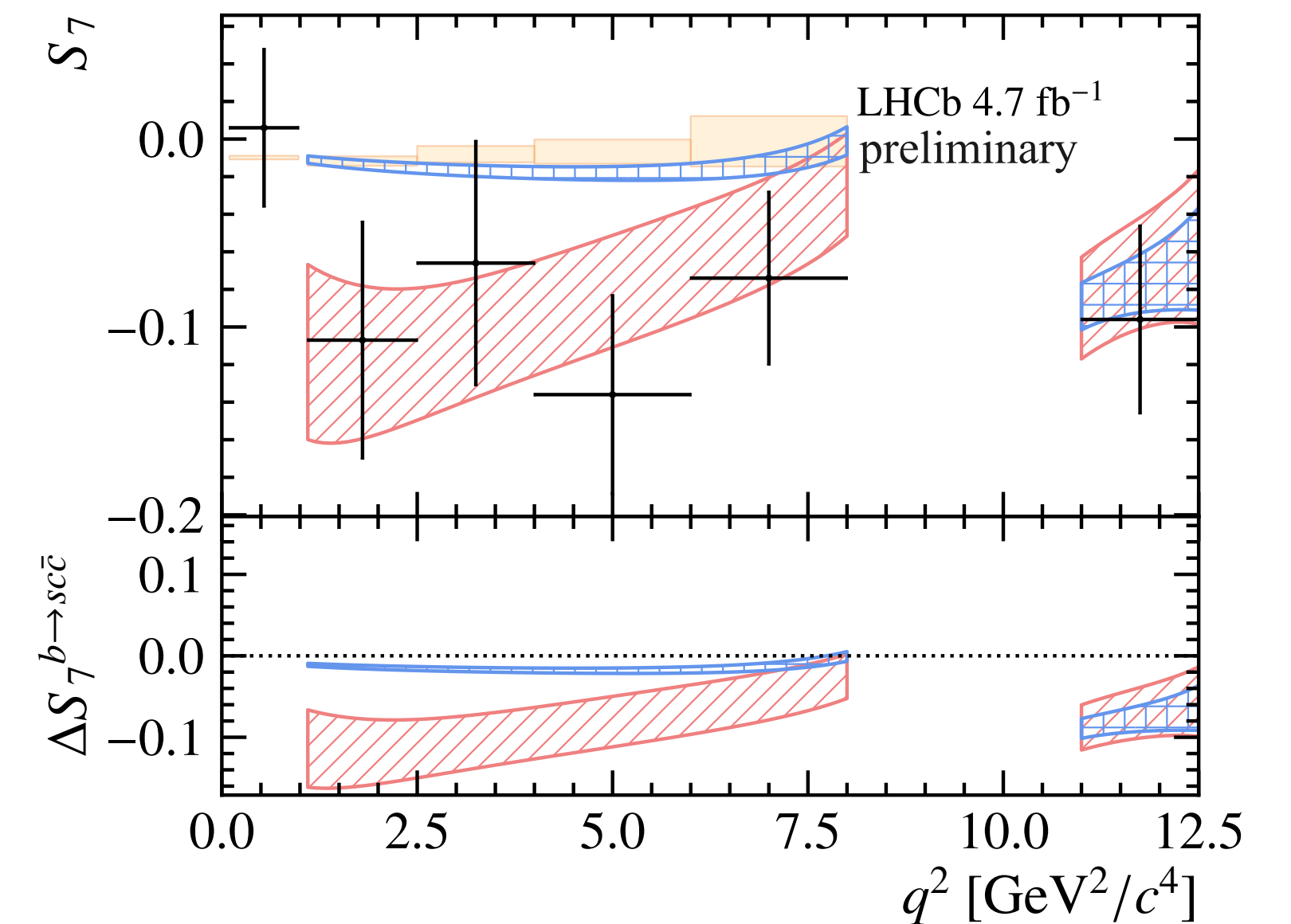
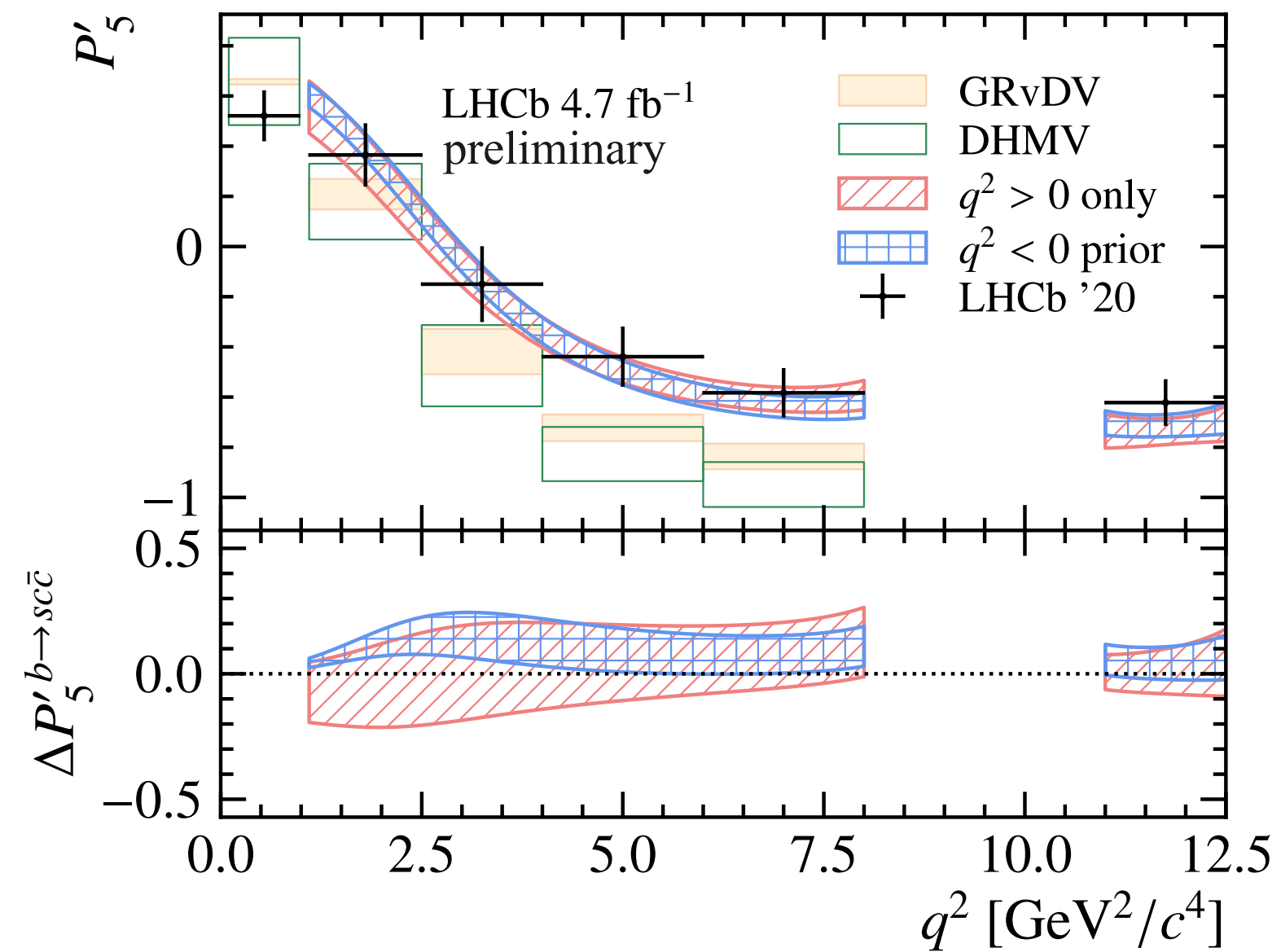
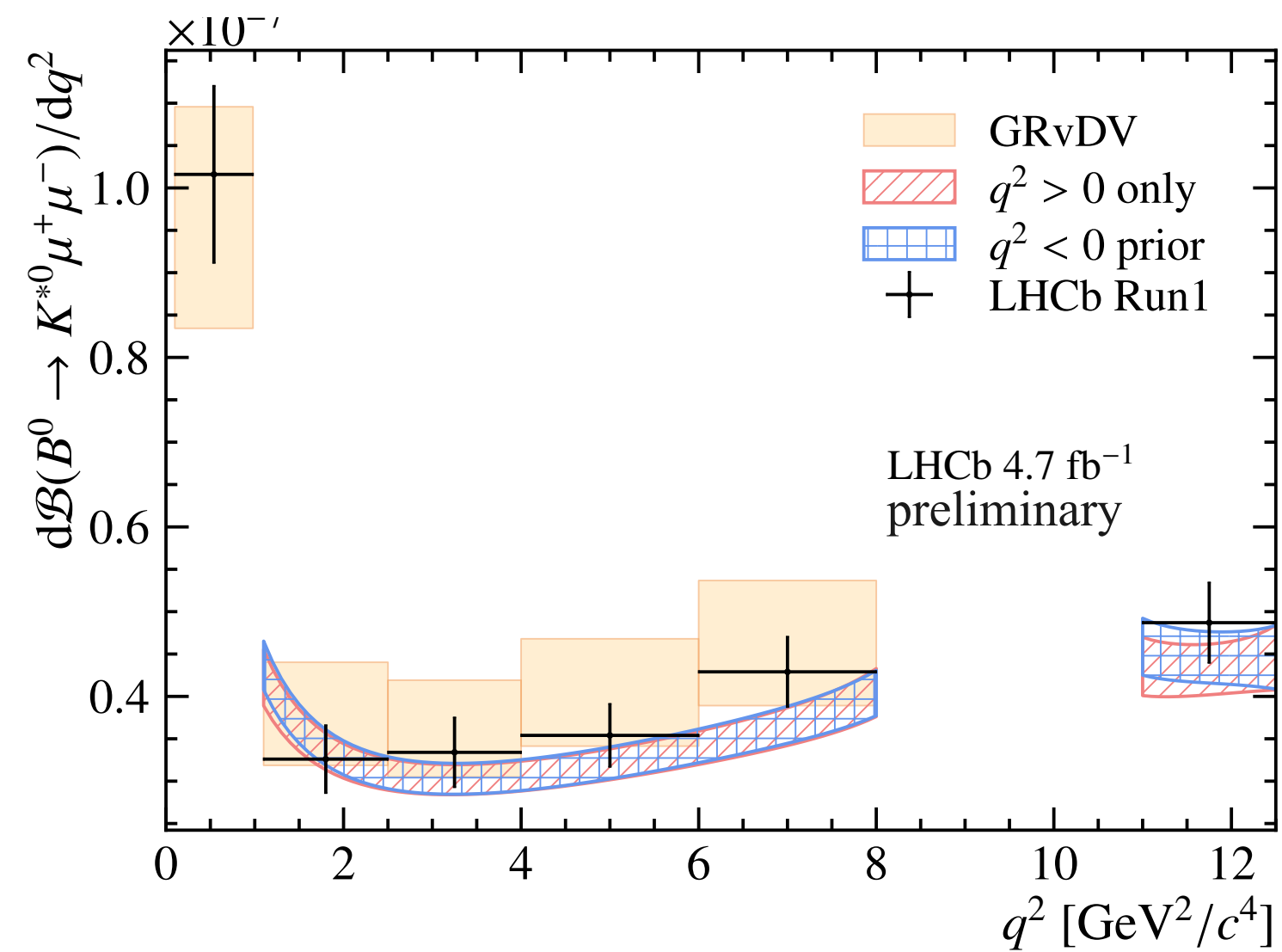
- Can measure phase difference between rare mode and Jpsi

- ▶ $\arg A_0^{J/\psi} = \begin{matrix} -1.55^{+0.22}_{-0.18} & [q^2 < 0] \\ -1.61^{+0.22}_{-0.20} & [q^2 > 0] \end{matrix}$ Compatible with what measured in $B^+ \rightarrow K^+ \mu^+ \mu^-$ decays



Retrieving the angular observables

- From the fit result we can reproduce the classic binned observables



- Lower BR compared to LHCb Run1 due to updated normalisation inputs

- Bottom: $c\bar{c}$ impact to P'_5

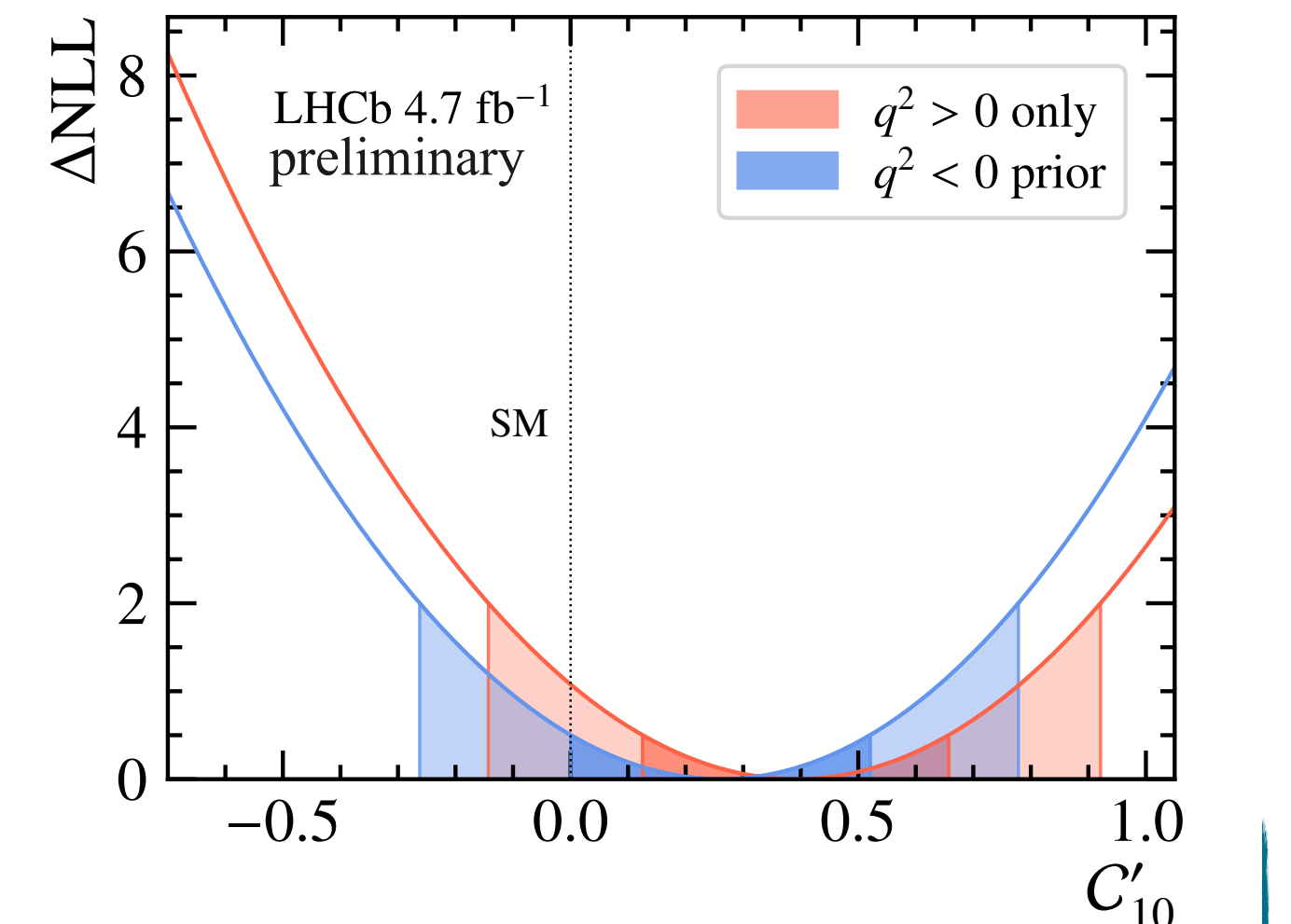
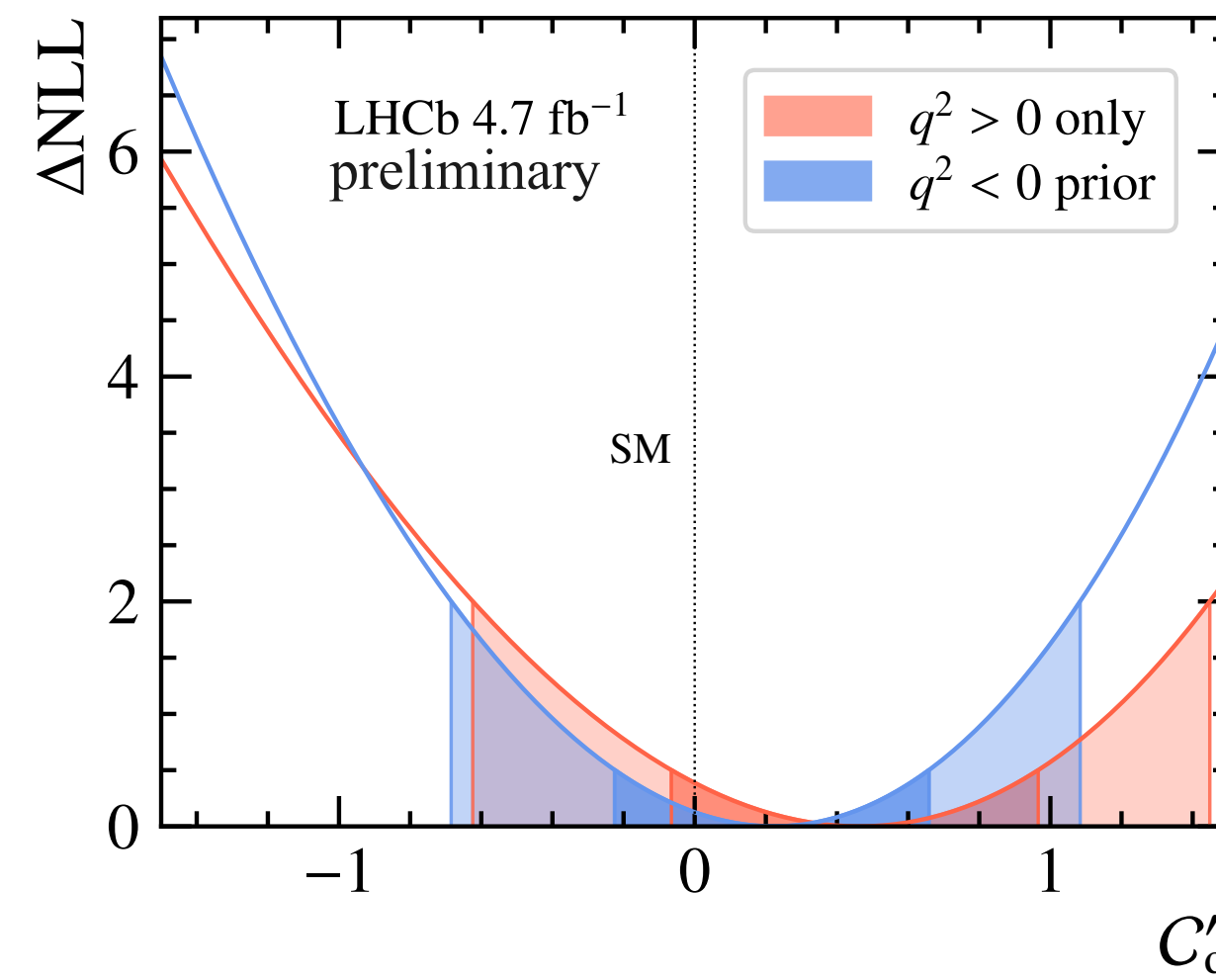
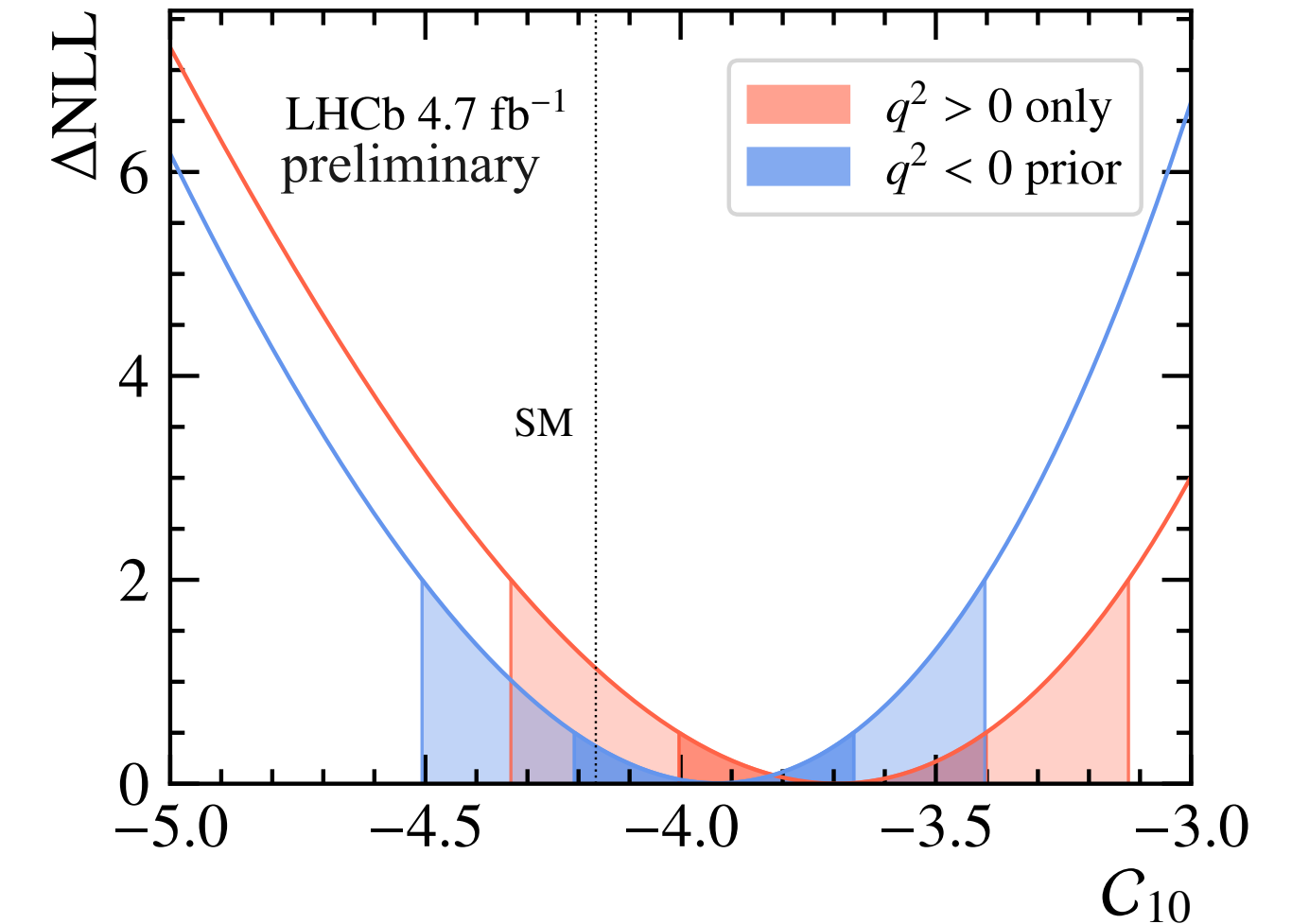
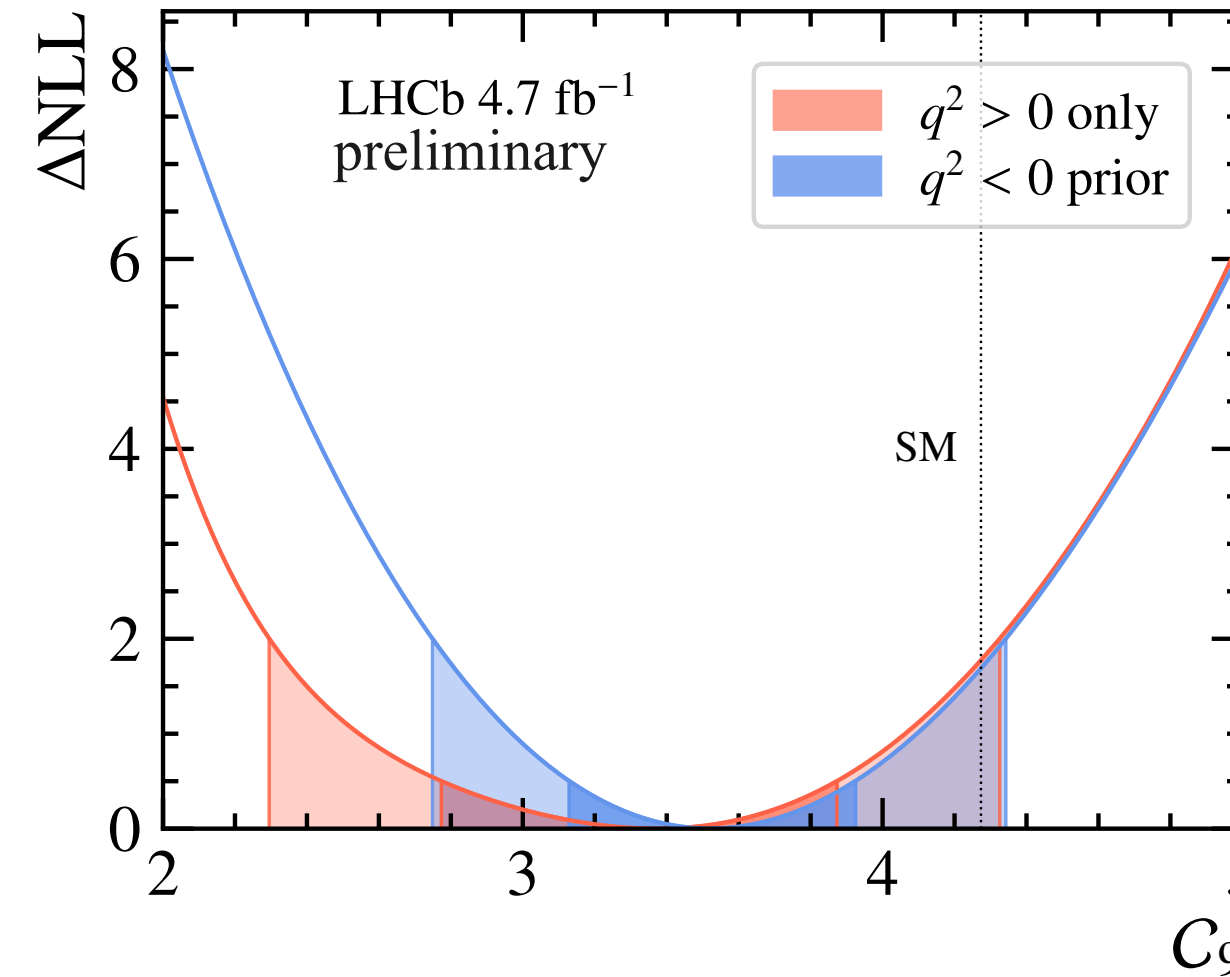
$$J_7 = \frac{3\sqrt{2}}{4} \beta_l \left[\text{Im}(A_0^L A_{\parallel}^{L*}) - (L \rightarrow R) \right]$$

Wilson coefficients

- Uncertainty obtained from neg. log-likelihood profile

	$q^2 > 0$ only	
	Fit result	deviation from SM
C_9	$-0.93^{+0.53}_{-0.57}$	1.9σ
C_{10}	$0.48^{+0.29}_{-0.31}$	1.5σ
C'_9	$0.48^{+0.49}_{-0.55}$	0.9σ
C'_{10}	$0.38^{+0.28}_{-0.25}$	1.5σ
	$q^2 < 0$ prior	
C_9	$-0.68^{+0.33}_{-0.46}$	1.8σ
C_{10}	$0.24^{+0.27}_{-0.28}$	0.9σ
C'_9	$0.26^{+0.40}_{-0.48}$	0.5σ
C'_{10}	$0.27^{+0.25}_{-0.27}$	1.0σ

- Global compatibility [4 d.o.f.] with SM 1.3 (1.4) σ

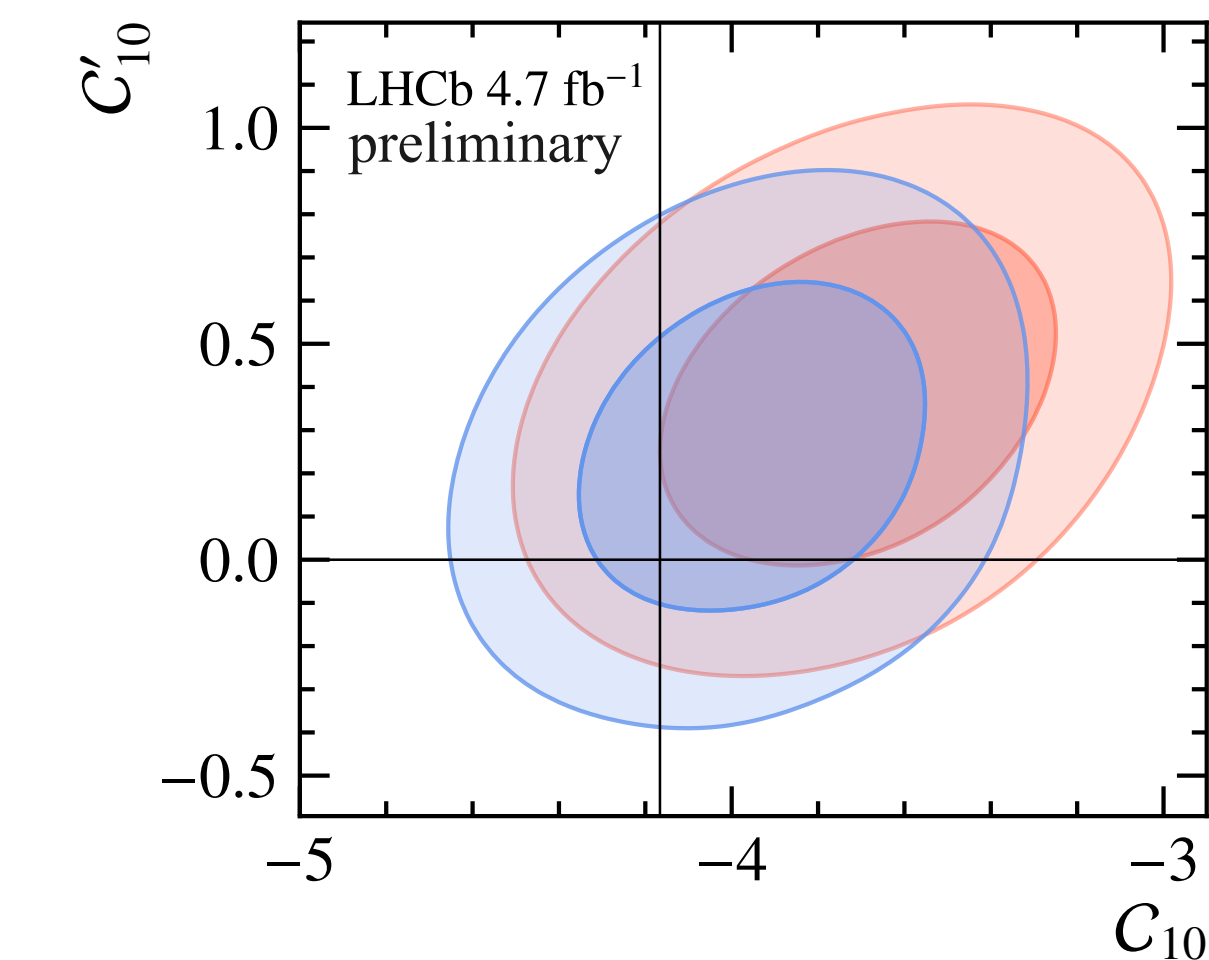
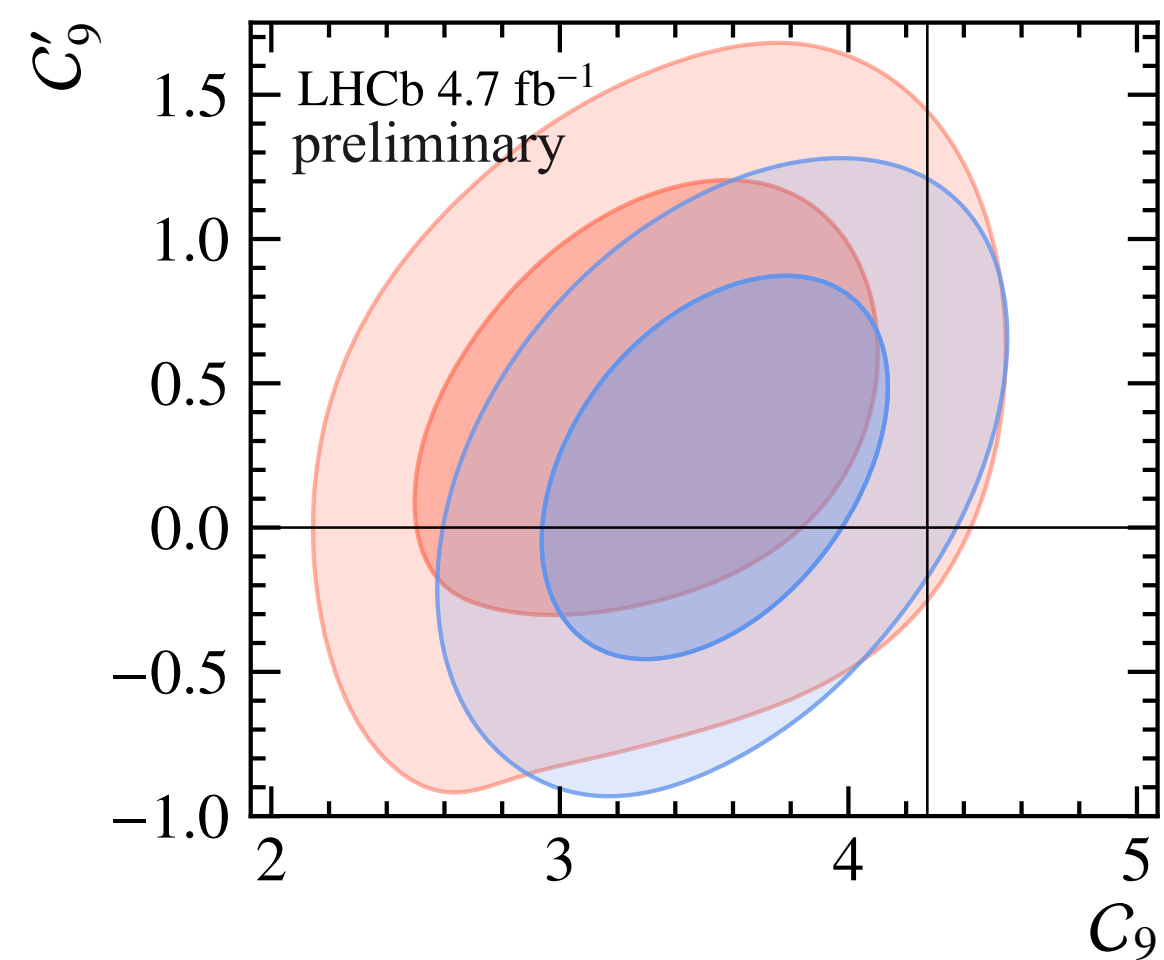
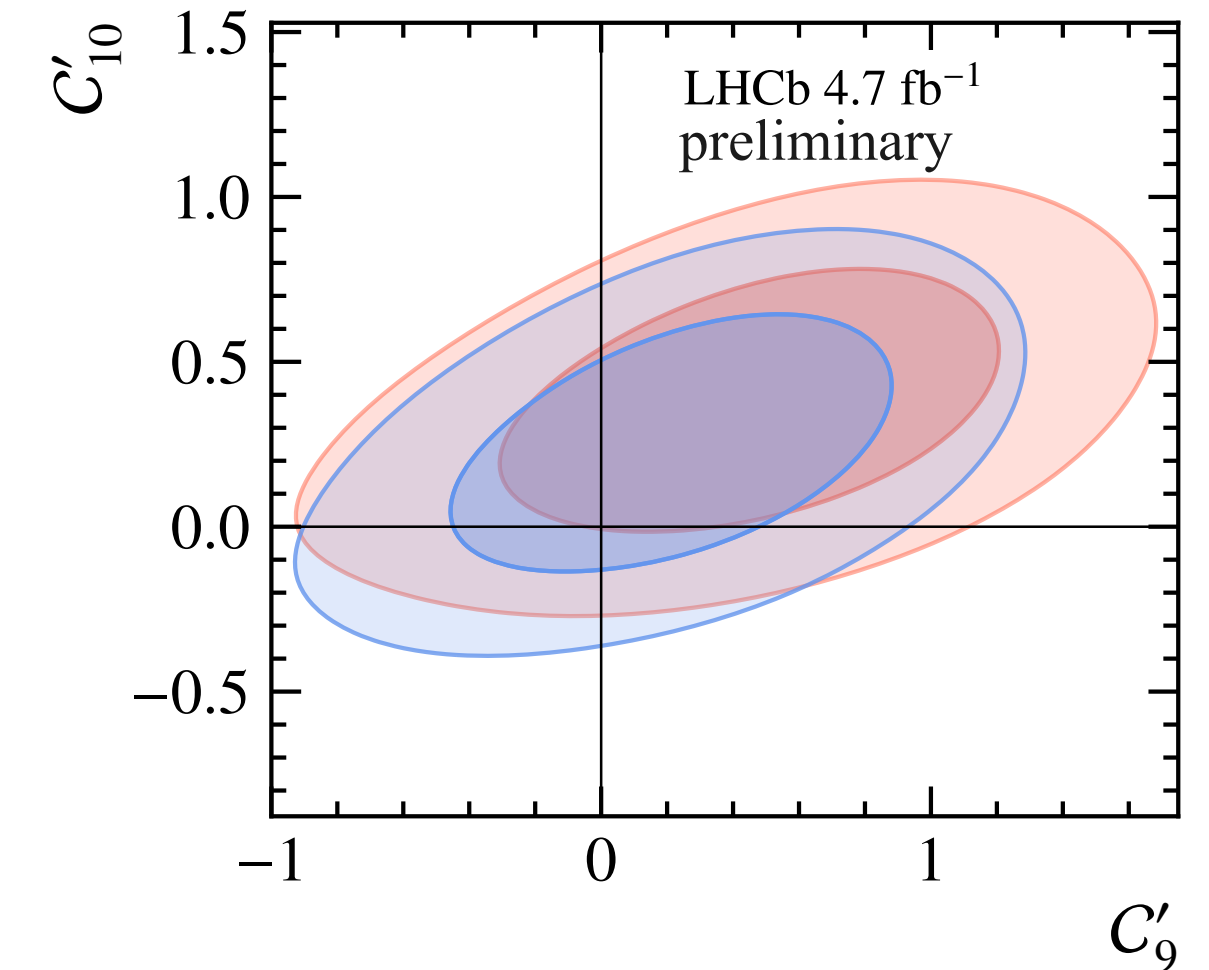
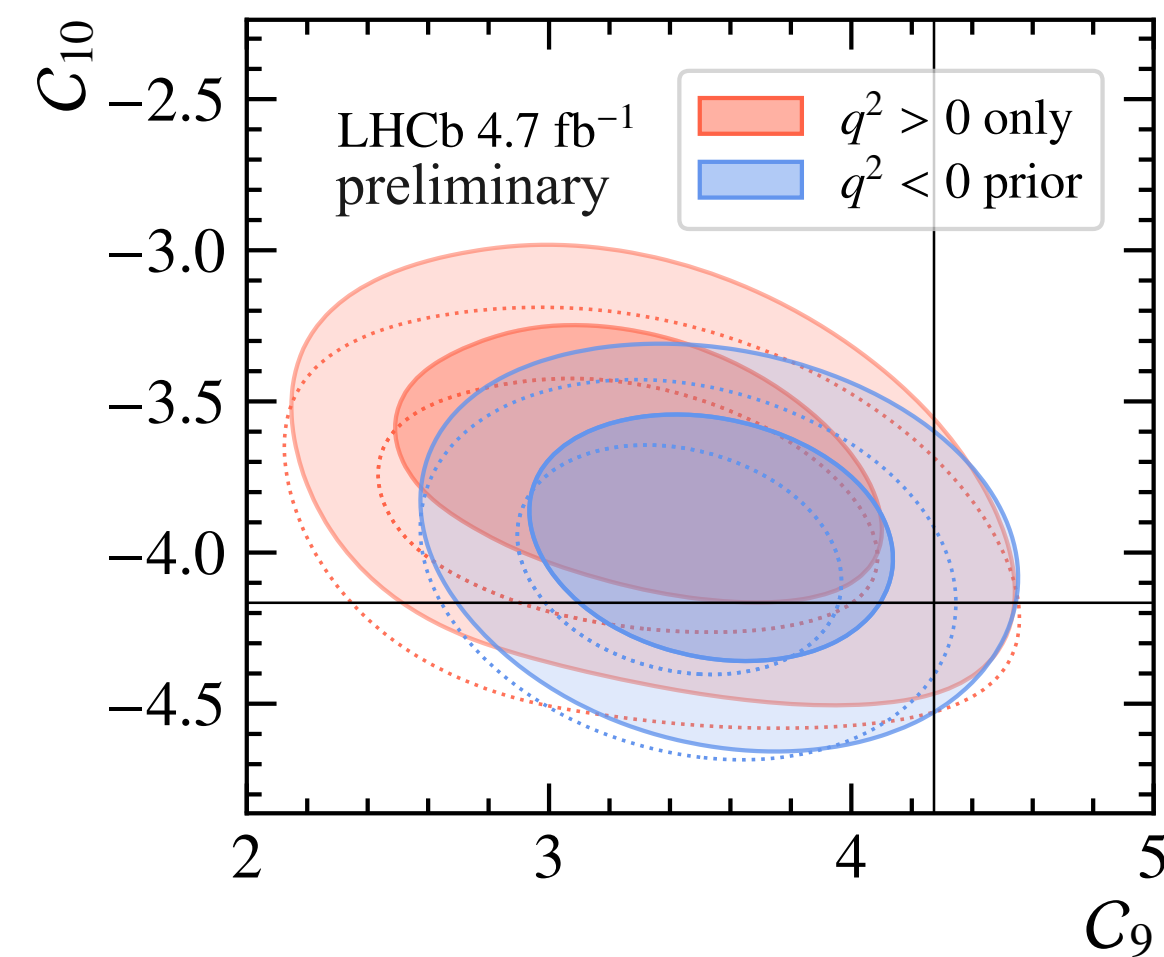


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Conclusion

- First unbinned amplitude analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ decays
- Complementary and more in dept set of information w.r.t. previous binned analyses
- Non-local hadronic determined from data under two assumptions
- Despite the extra freedom given by $c\bar{c}$ pars, fit still prefers to insert a shift in \mathcal{C}_9
- Tension reduced to $\sim 2\sigma$ in \mathcal{C}_9 and $\sim 1.4\sigma$ global

Backup

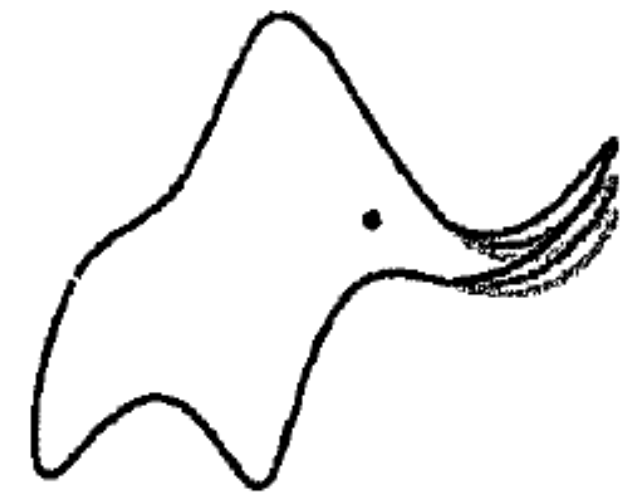
Results 68% 95% CL

	$q^2 > 0$ only			
	best fit value	68% CL	95% CL	deviation from SM
\mathcal{C}_9	3.34	[2.77, 3.87]	[2.30, 4.33]	1.9 σ
\mathcal{C}_{10}	-3.69	[-4.00, -3.40]	[-4.33, -3.12]	1.5 σ
\mathcal{C}'_9	0.48	[-0.07, 0.97]	[-0.62, 1.45]	0.9 σ
\mathcal{C}'_{10}	0.38	[0.13, 0.66]	[-0.14, 0.92]	1.5 σ
	$q^2 < 0$ prior			
\mathcal{C}_9	3.59	[3.13, 3.92]	[2.75, 4.34]	1.8 σ
\mathcal{C}_{10}	-3.93	[-4.21, -3.66]	[-4.51, -3.40]	0.9 σ
\mathcal{C}'_9	0.26	[-0.22, 0.66]	[-0.68, 1.08]	0.5 σ
\mathcal{C}'_{10}	0.27	[0.00, 0.52]	[-0.26, 0.78]	1.0 σ

Choice of the z order

- Truncation order of polynomial expansion chosen based on NLL improvement observed in data
 - ▶ $\mathcal{H}_\lambda[z^2]$ for fit @ $q^2 > 0$
 - ▶ $\mathcal{H}_\lambda[z^4]$ for fit with $q^2 < 0$ prior

	$2\Delta \log \mathcal{L}$	
	$q^2 < 0$ prior	$q^2 > 0$ only
$\mathcal{H}_\lambda[z^3] - \mathcal{H}_\lambda[z^2]$		3.6
$\mathcal{H}_\lambda[z^4] - \mathcal{H}_\lambda[z^3]$	21.22	-
$\mathcal{H}_\lambda[z^5] - \mathcal{H}_\lambda[z^4]$	8.64	-

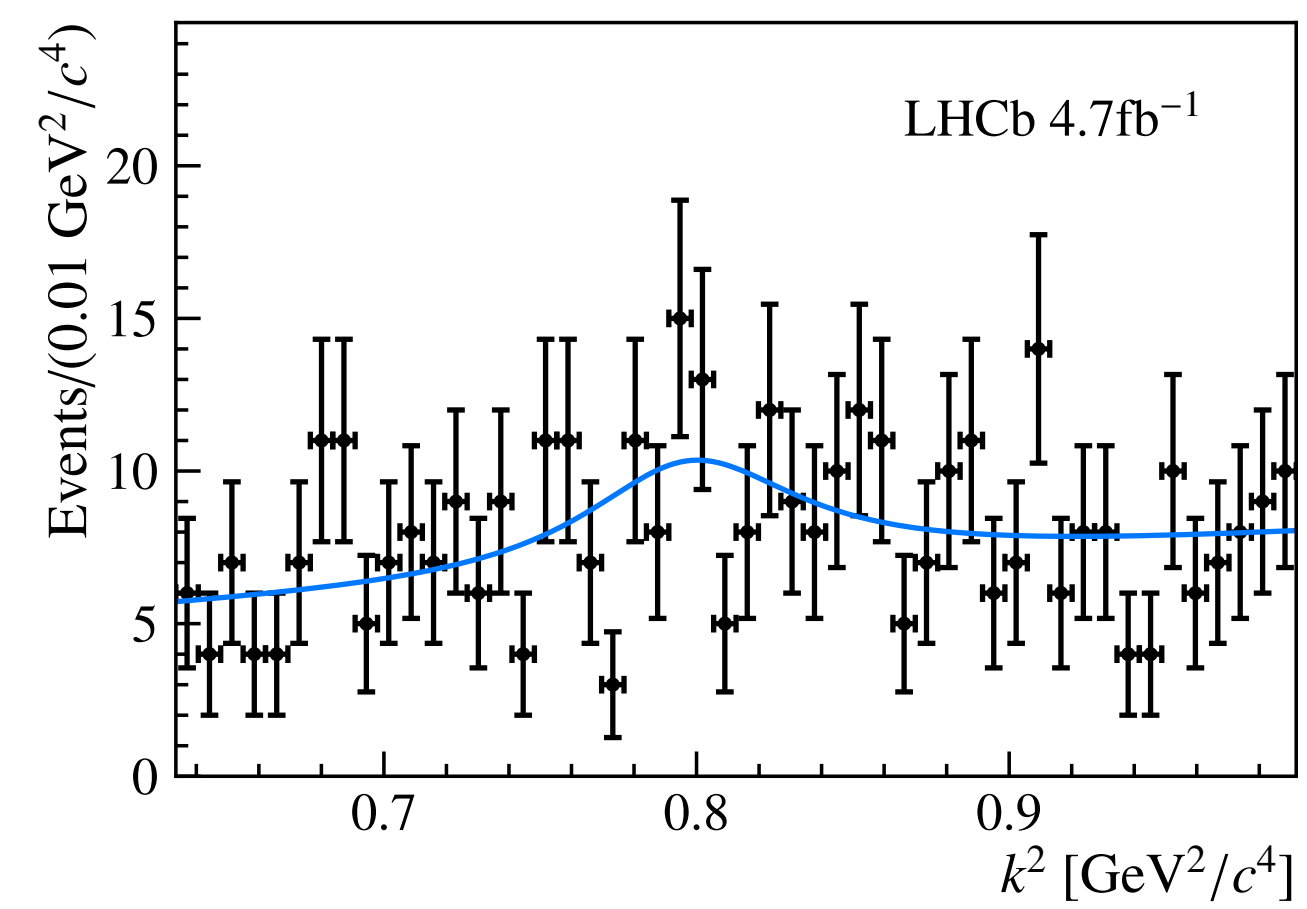
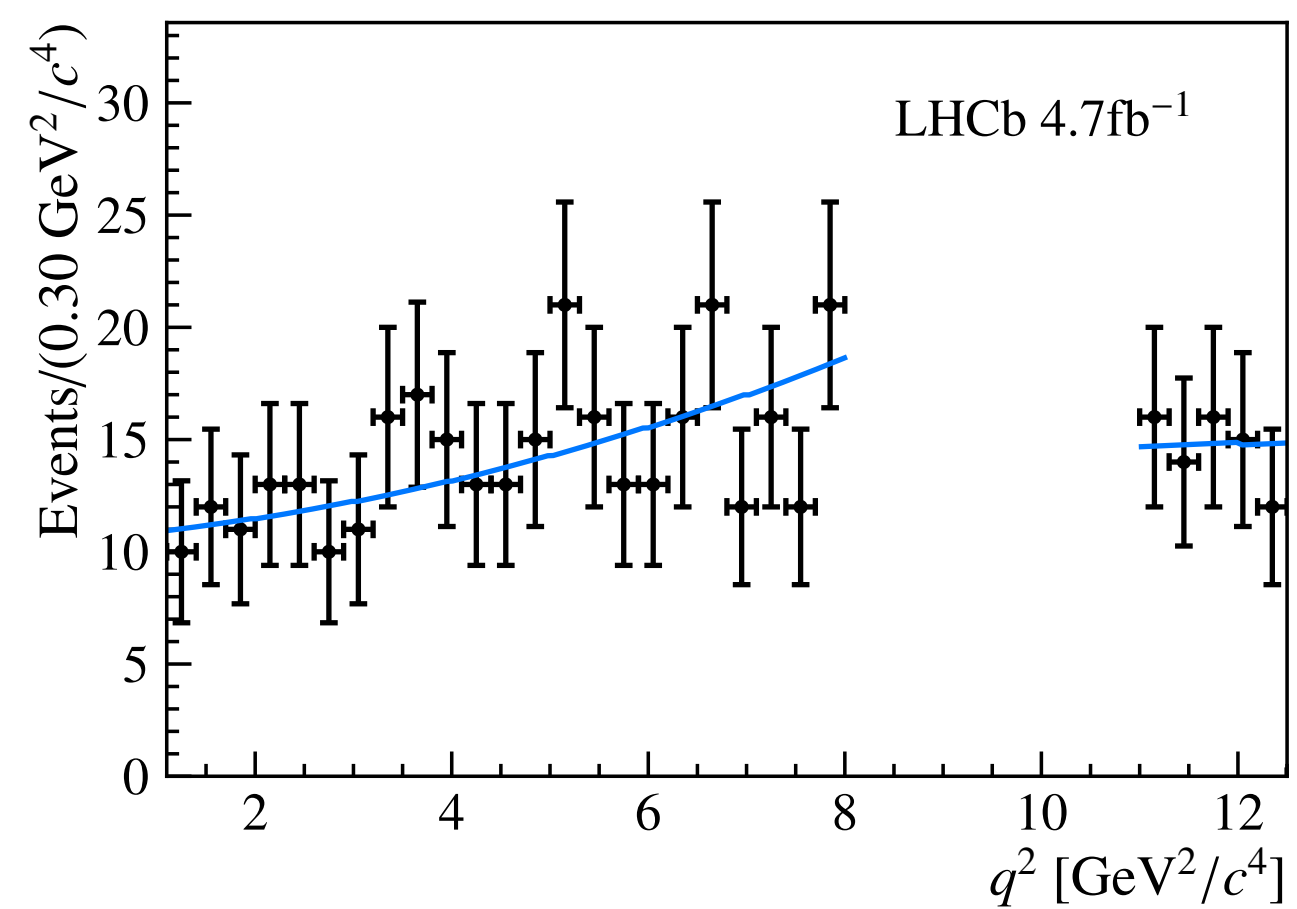
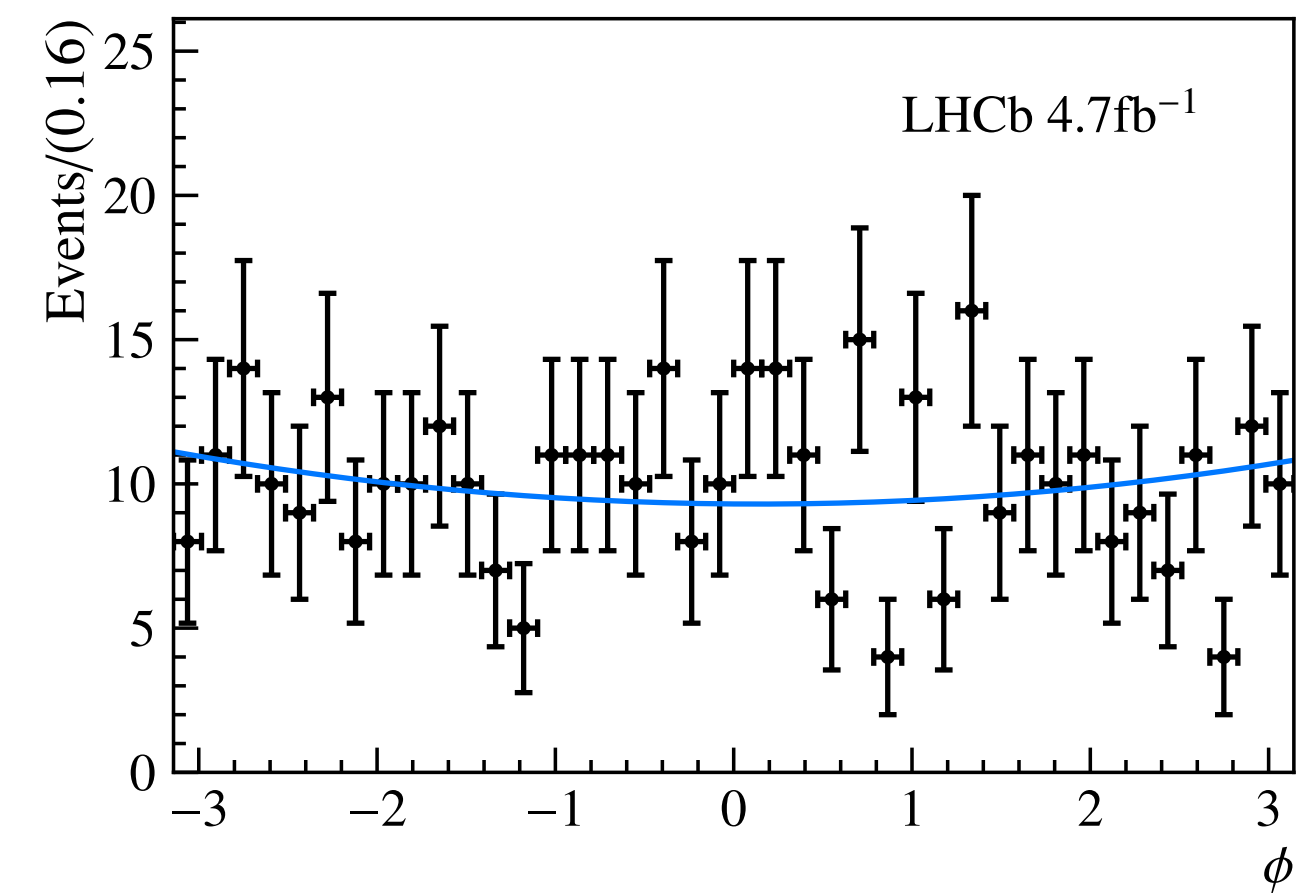
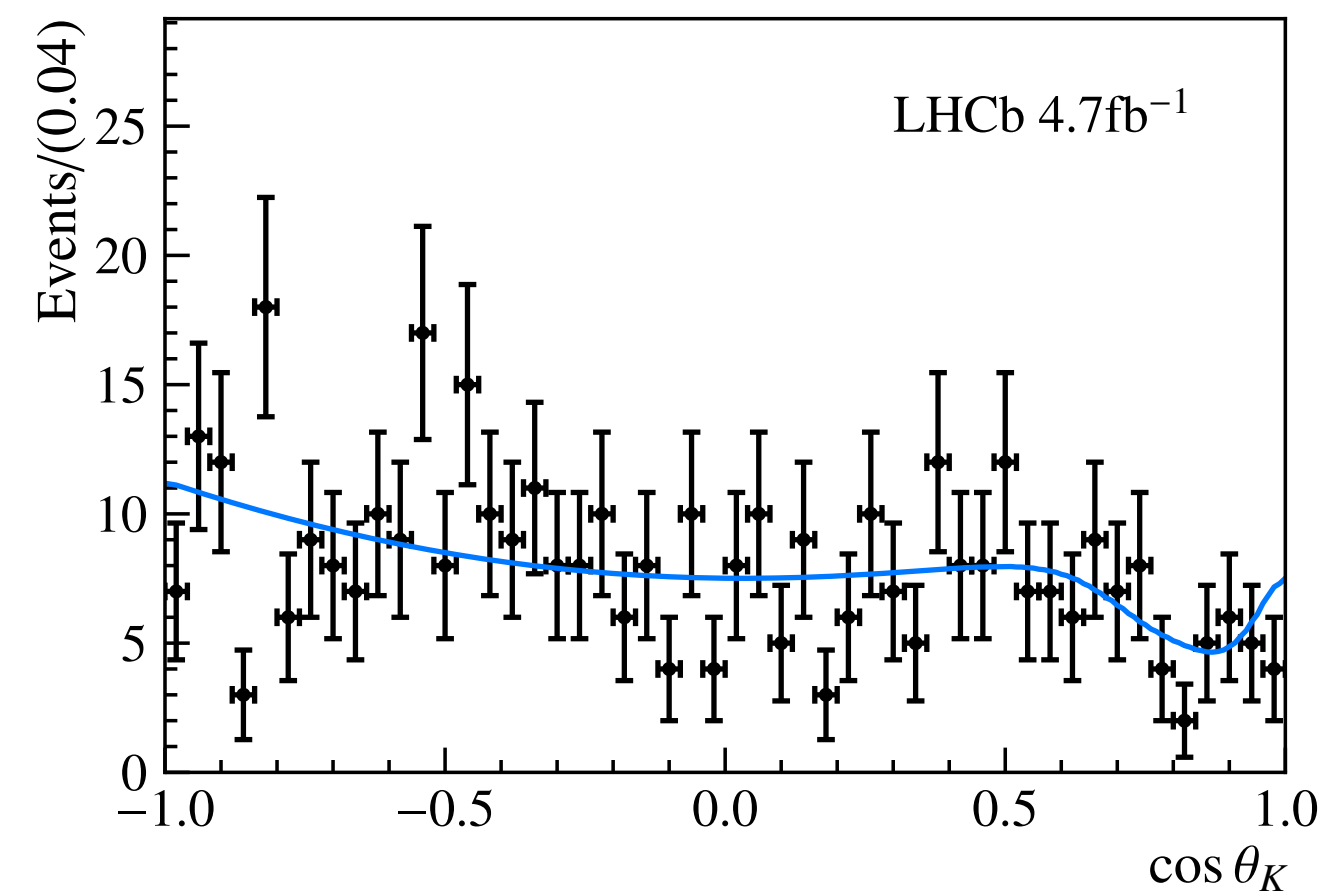
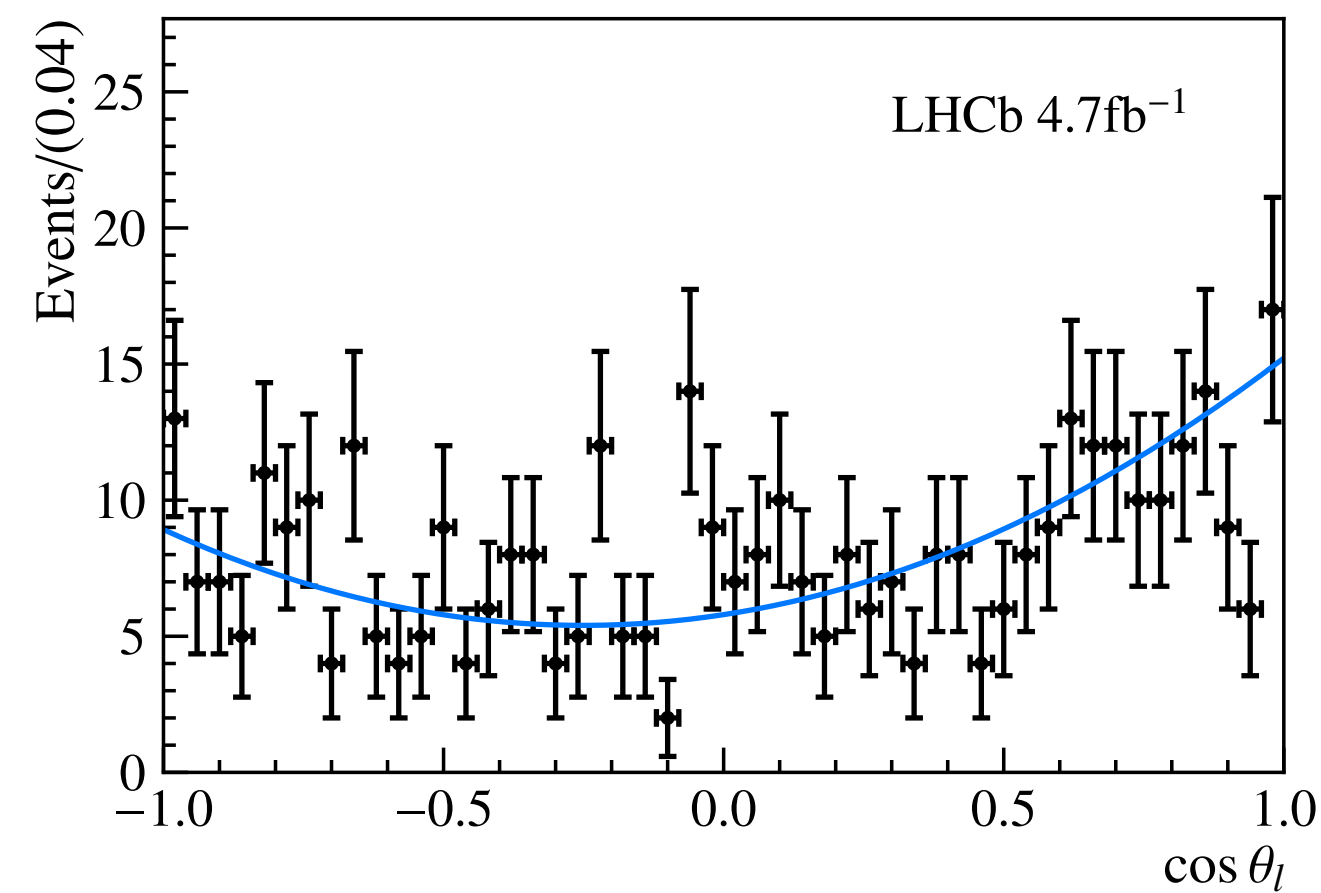


“With four parameters I can fit an elephant, and with five I can make him wiggle his trunk.” J. von Neumann

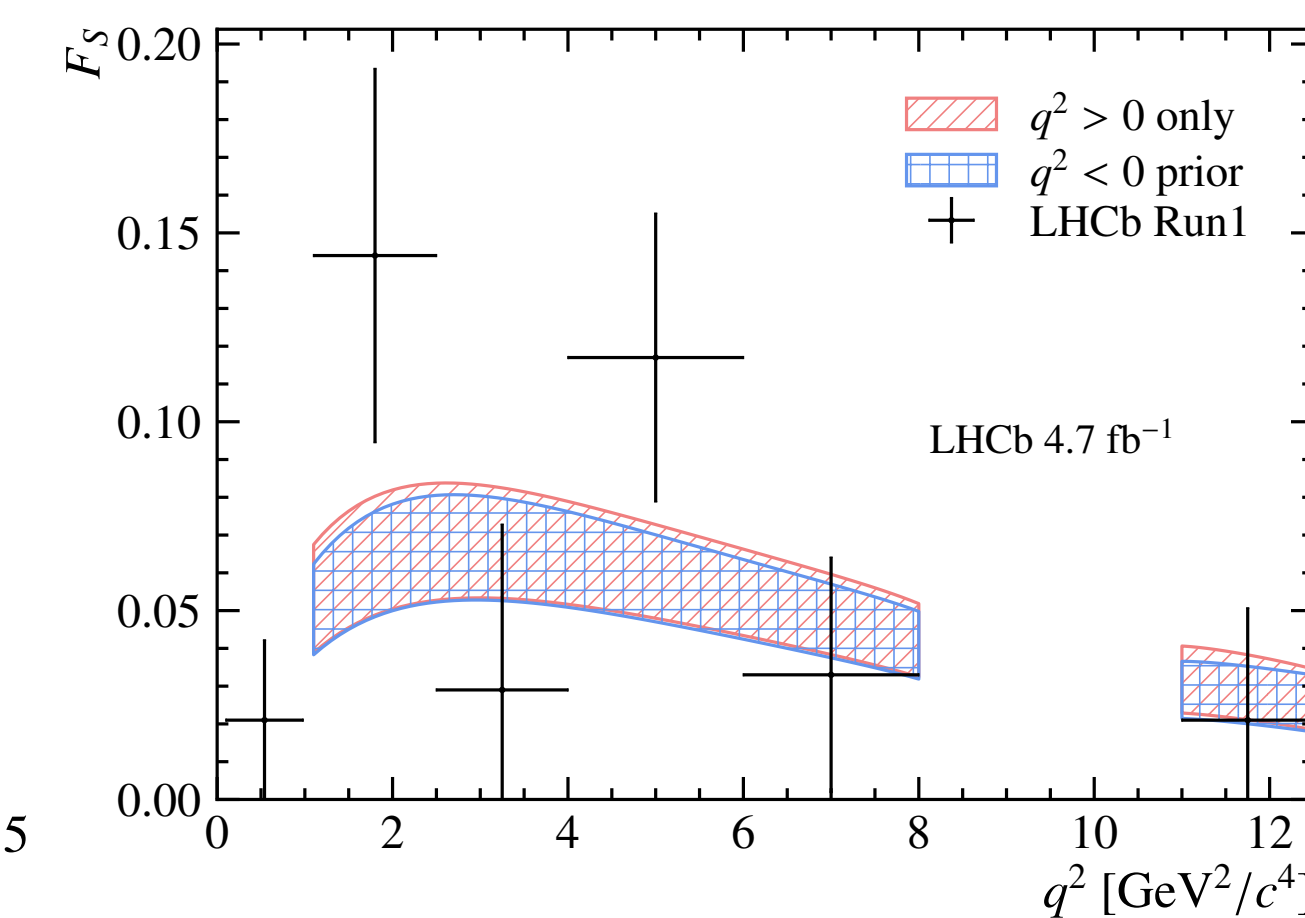
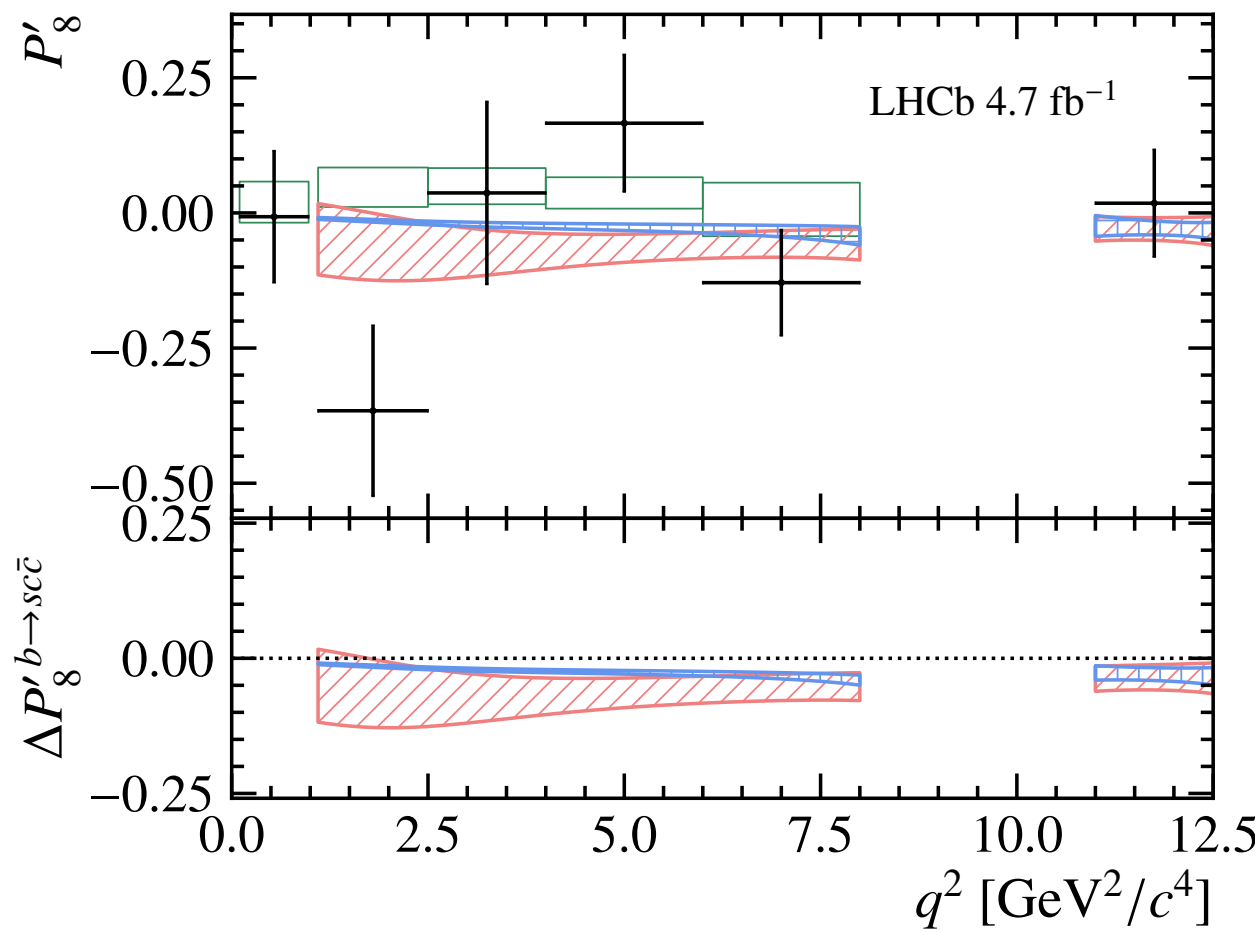
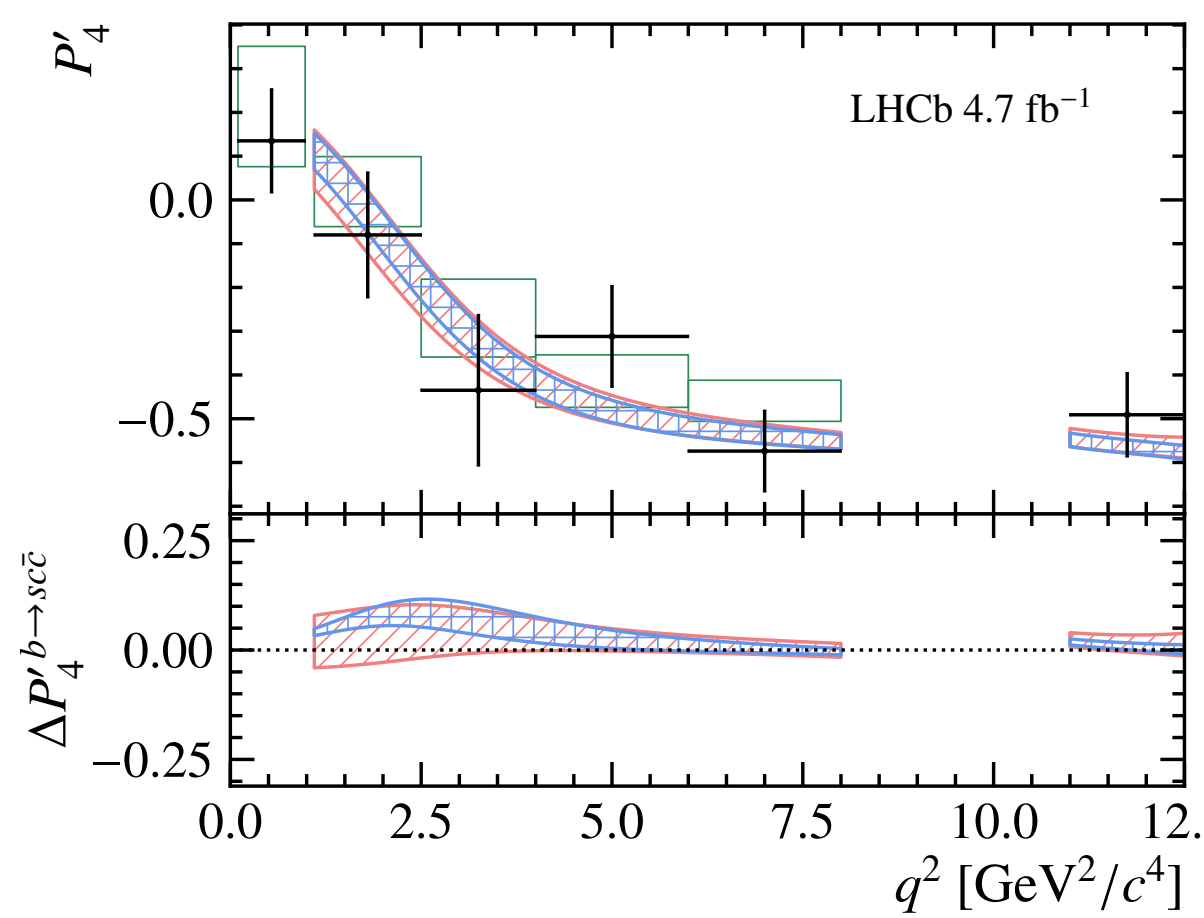
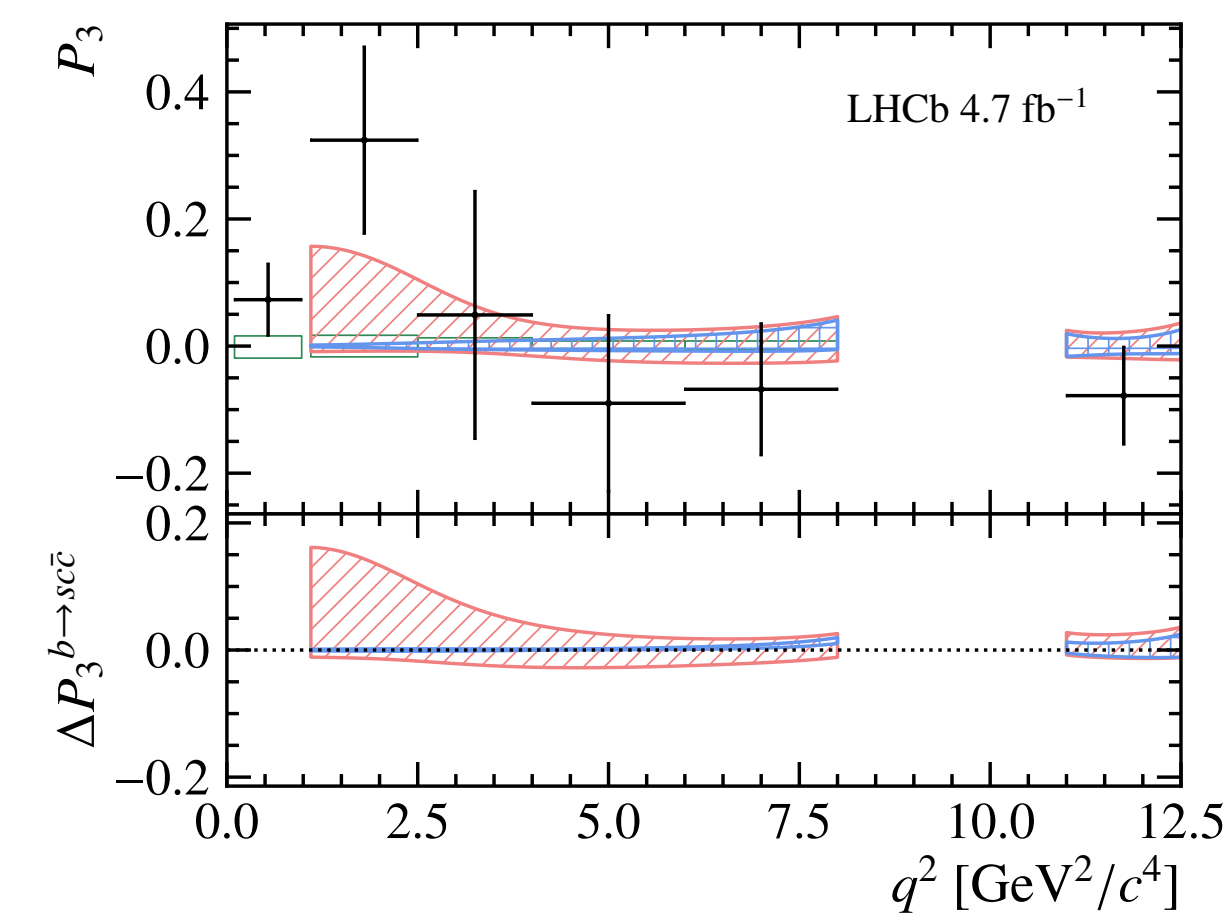
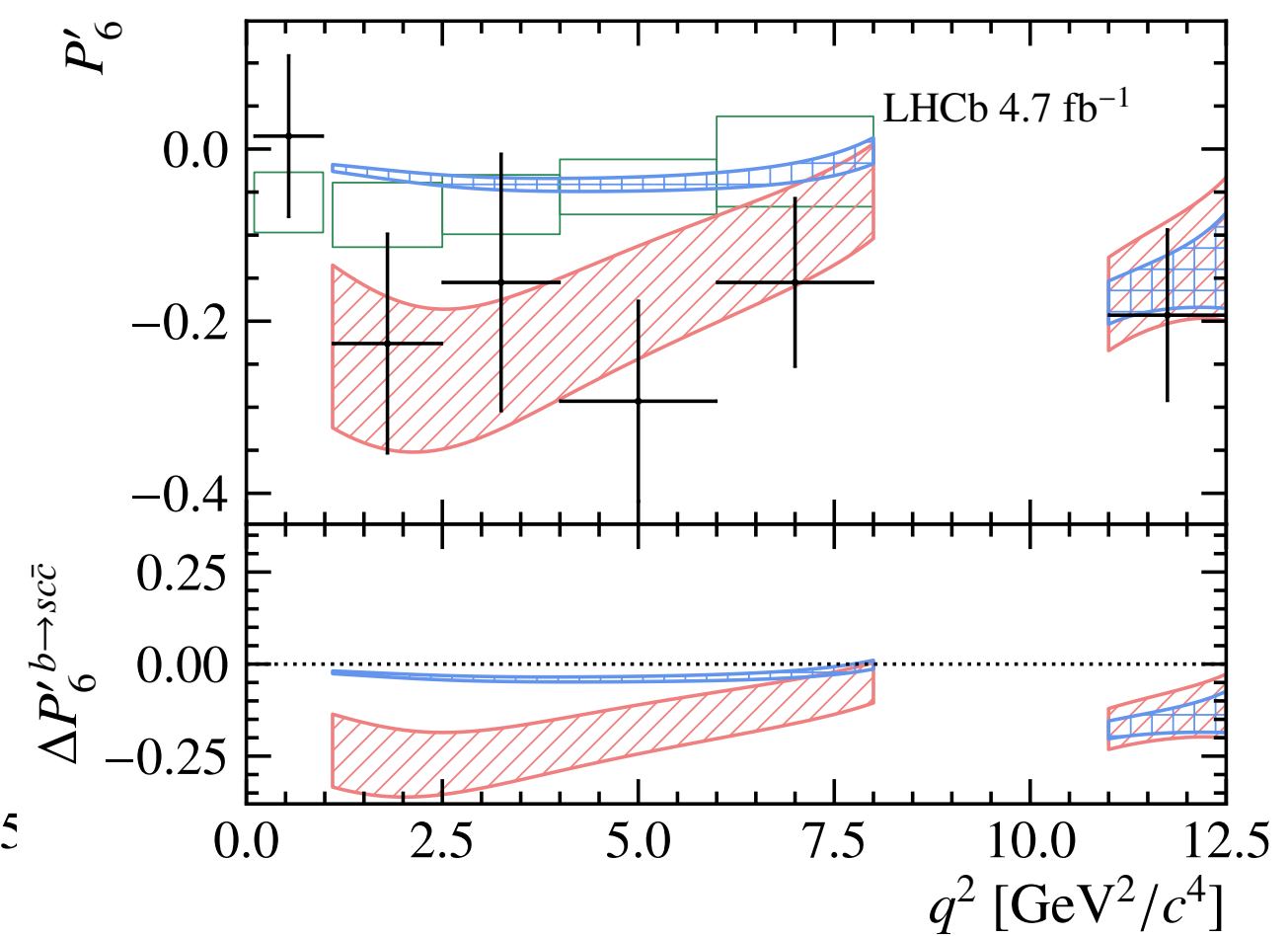
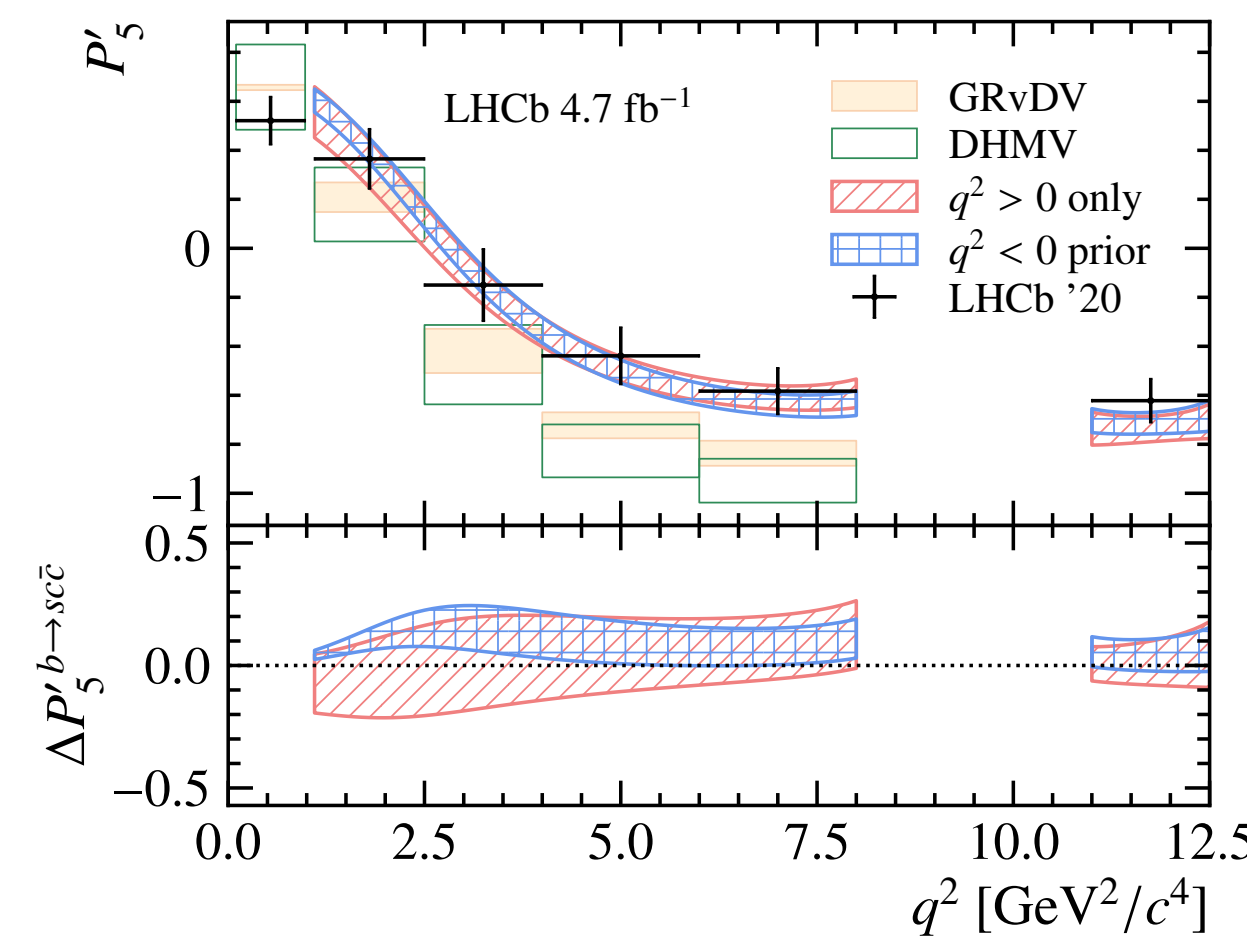
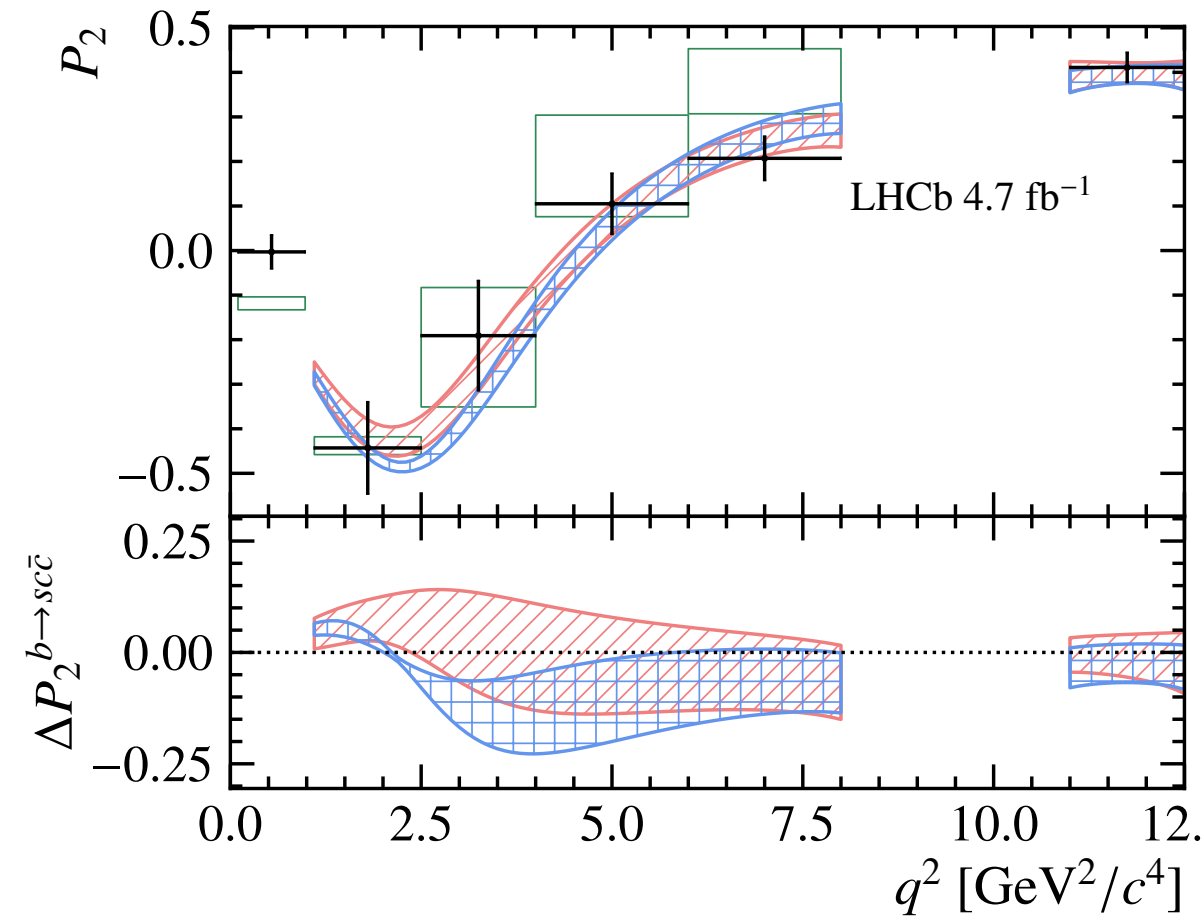
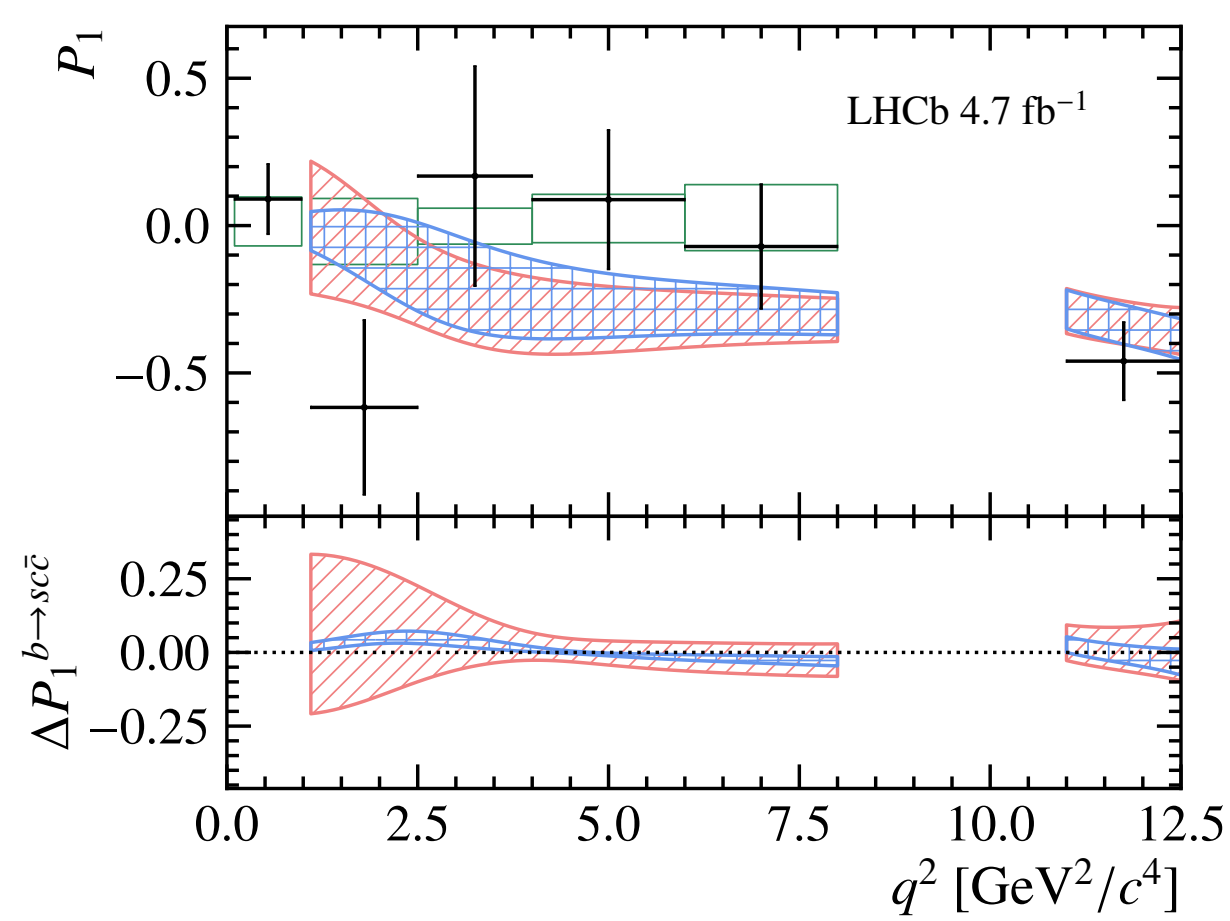
Systematics

	\mathcal{C}_9	\mathcal{C}_{10}	\mathcal{C}'_9	\mathcal{C}'_{10}
Amplitude model				
S-wave form factors	< 0.01	< 0.01	< 0.01	< 0.01
S-wave non-local hadronic	0.02	0.02	0.14	0.04
S-wave k^2 model	< 0.01	< 0.01	0.05	0.03
Subtotal	0.02	0.02	0.15	0.05
External inputs on BR				
$\mathcal{B}(B^0 \rightarrow J/\psi K^+ \pi^-)$	0.05	0.08	0.02	0.01
$f_{\pm 100\text{MeV}}^{B^0 \rightarrow J/\psi K\pi}$	0.03	0.03	0.01	< 0.01
Others	0.03	0.04	0.03	0.01
Subtotal	0.07	0.09	0.04	0.01
Background model				
Chebyshev polynomial order	0.01	0.01	0.01	< 0.01
Combinatorial shape in k^2	0.02	< 0.01	0.02	< 0.01
Background factorisation	0.01	0.01	0.01	0.01
Peaking background	0.01	< 0.01	0.02	0.01
Subtotal	0.03	0.02	0.03	0.01
Experimental effects				
Acceptance parametrisation	< 0.01	< 0.01	< 0.01	< 0.01
Statistical uncertainty on acceptance	0.02	< 0.01	0.02	< 0.01
Subtotal	0.02	< 0.01	0.02	< 0.01
Total systematic uncertainty	0.08	0.10	0.16	0.05

Upper mass projections



Ang. obs (P-basis)



Ang. obs (S-basis)

