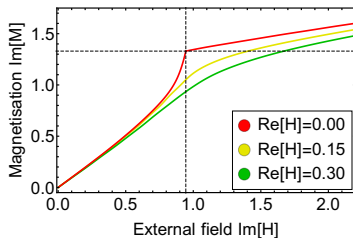
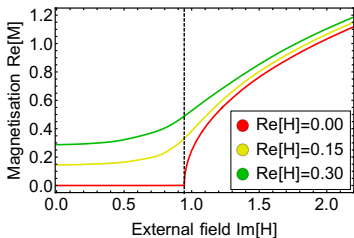


Towards locating the (real) critical endpoint

Friederike Ihssen

12/02 - CPOD 2022



Based on arXiv:2207.10057 in collaboration with J. M. Pawłowski



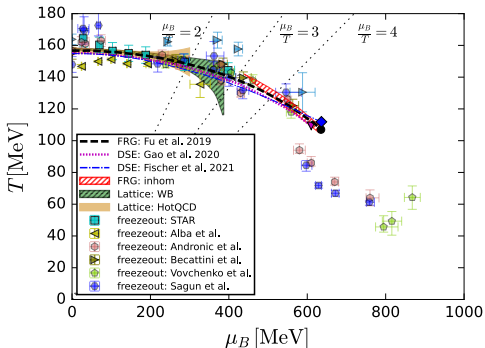
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Why are we interested in Lee-Yang Zeroes?

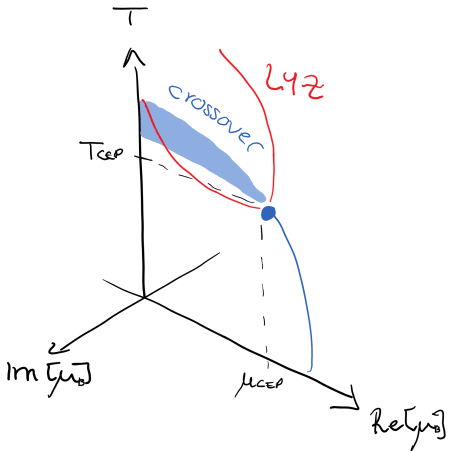


PhysRevD.101.054032: Fu, Pawłowski, Rennecke

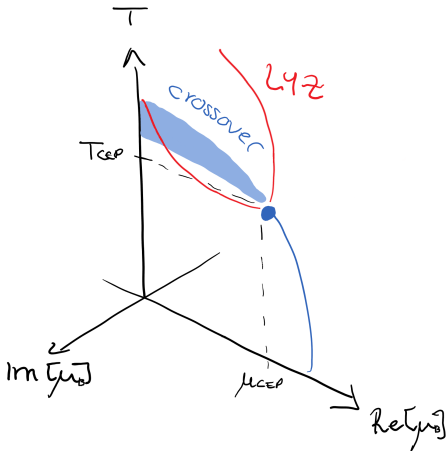
PhysLettB.2021.136584: Gao, Pawłowski

PhysRevD.104.054022: Gunkel, Fischer

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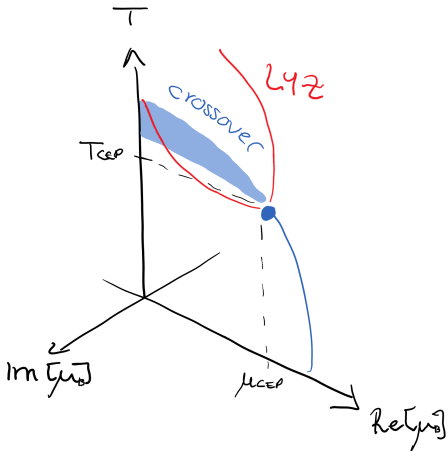
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- A zero-crossing of the generating functional/partition function \mathcal{Z}

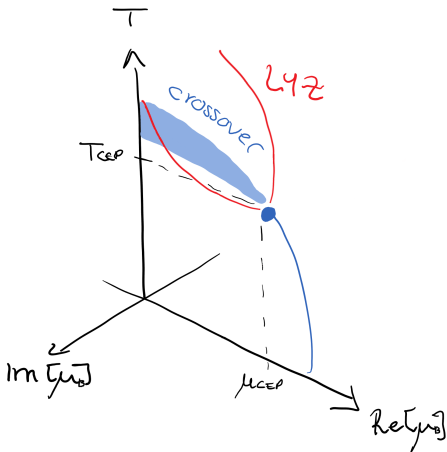
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- A zero-crossing of the generating functional/partition function \mathcal{Z}
- A divergence/ discontinuity in the free energy $F = -\beta^{-1} \log(\mathcal{Z})$.

⇒ How are complex LYZs connected to the (real) critical endpoint?

Complex phase transitions?

Partition function :

$$\mathcal{Z} = \sum_i \exp(-\beta E_i).$$

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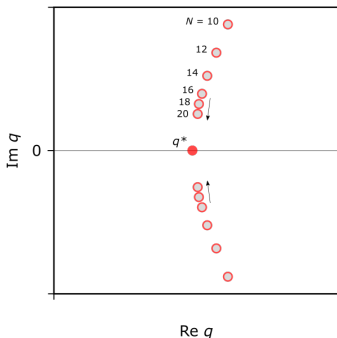
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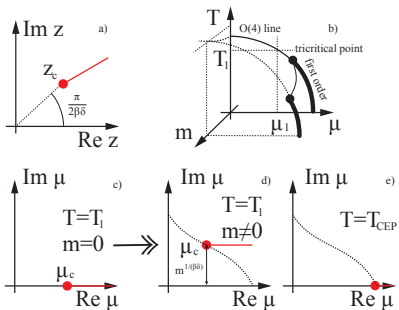
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⇒ Predict real phase transition!

Lee Yang edge singularities in QCD

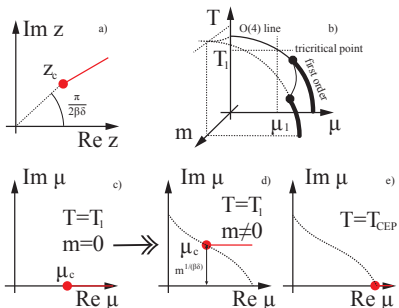
- Simple LEFT for meson interactions: O(4) theory



PhysRevLett.125.191602: Connelly, Johnson,
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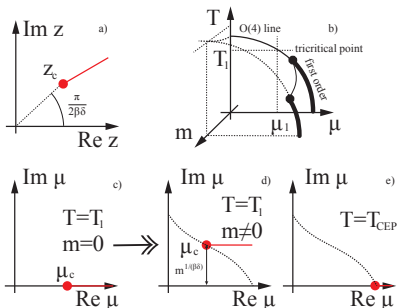
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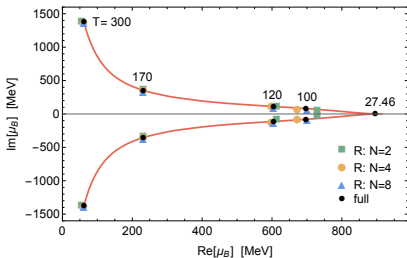
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Lee Yang edge singularities in QCD

- Simple LEFT for meson interactions: $O(4)$ theory
- At $m_\pi \rightarrow 0$ we find a second order phase transition at high T .
- As we go to physical pion masses the LY travels through the complex plane for increasing $\text{Re } \mu$.



PhysRevLett.125.191602: Connelly, Johnson,
Mukherjee, Skokov



PhysRevD.105.014026: Mukherjee, Rennecke,
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A Functional Approach

The path-integral / generating functional:

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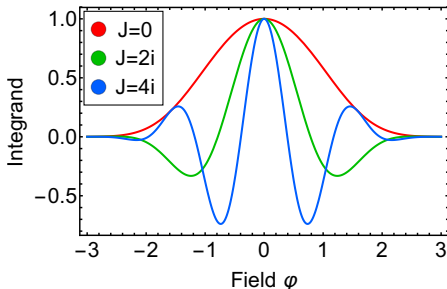
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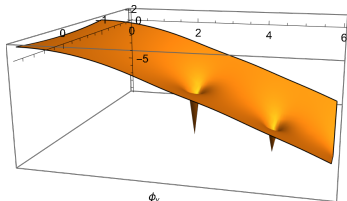
- $S[\varphi]$ classical action of the theory.
- Real function of a complex variable.
- Sample from oscillatory integrand (Lattice).



PoS LATTICE2021 (2022) 223: Attanasio,
Bauer, Kades, Pawłowski
arXiv 2203.01243: Pawłowski, Urban

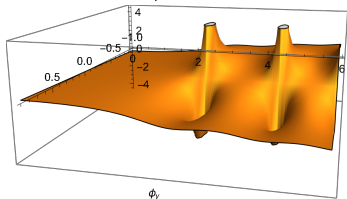
The $O(N)$ Model

$$S[\varphi] = \int_x \left\{ \frac{1}{2} \varphi(x) \left[-\partial_\mu^2 + m^2 \right] \varphi(x) + \frac{\lambda}{4!} \varphi(x)^4 \right\}$$

 ϕ_x

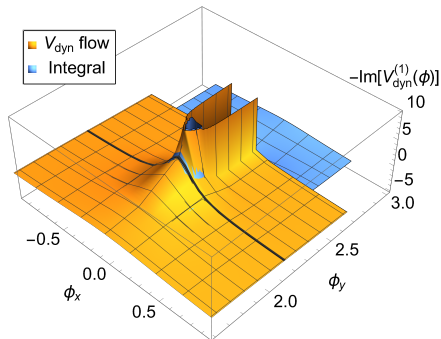
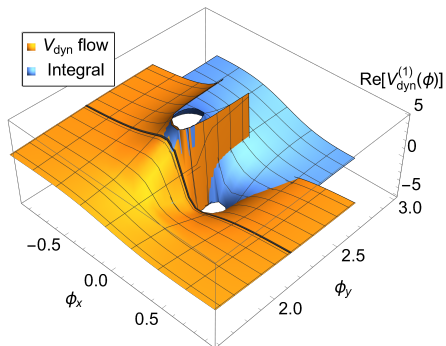
Test case:

- $O(N)$ with $d = 0, N = 1$
- Classical input parameters $m = 1, \lambda = 1$
- Direct evaluation of \mathcal{Z}

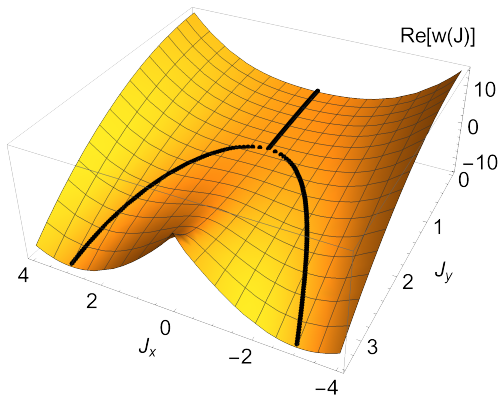
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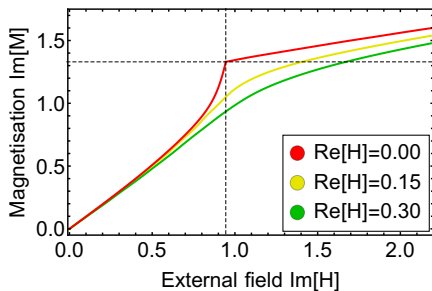
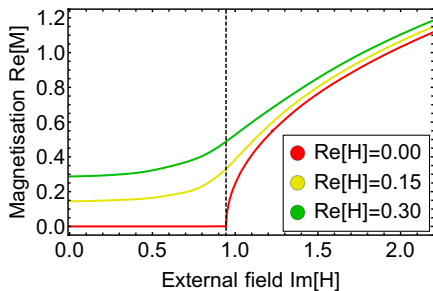
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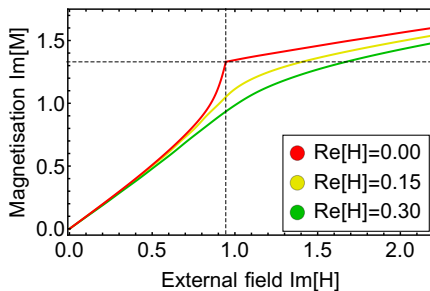
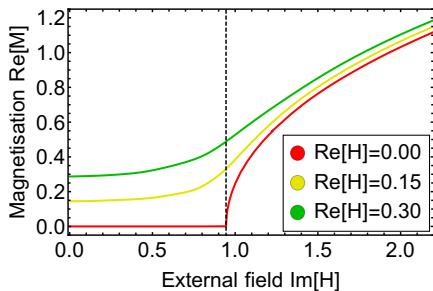
Magnetisation $M = \langle \phi \rangle$, at minimum of the effective potential $w[J]$.



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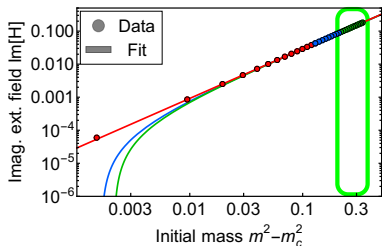
LY-location for $\text{Im}[H] = 0.994$ and initial condition $m^2 = 1$ at $\text{Im}[M] = 1.33$!

Recovering a Real Phase Transition

Scaling behaviour: $\frac{m^2 - m_c^2}{m_c^2} = \text{const } H^{1/\Delta}$

Recovering a Real Phase Transition

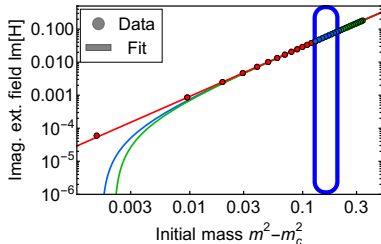
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How close do we need to go to extrapolate m_c^2 ?

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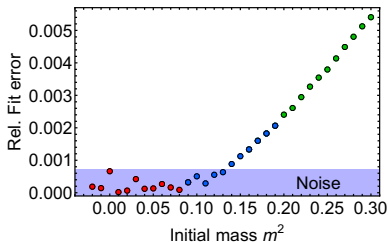
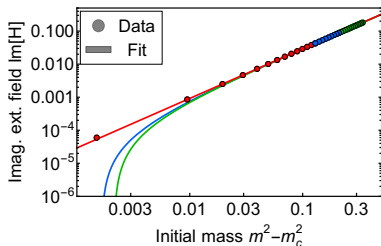
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Δ	m_c^2
1.4969(66)	-0.03943(17)

⇒ Last computed physical mass:

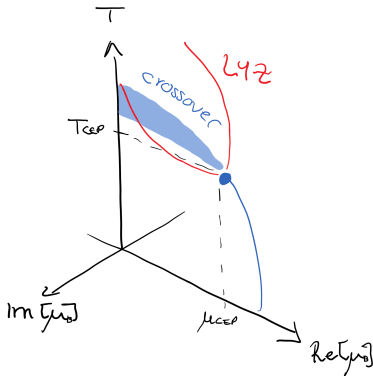
$$m_0^2(m^2 = -0.039) = 4.2 \cdot 10^{-4}$$

Summary & Outlook

- Investigated the convergence of functional flows in the complex plane using DG-methods.
- Retrieved location of the Lee-Yang singularity in $d = 0$ and $d = 4$.
- We extracted the critical mass m_c^2 of the real phase transition in $d = 4$!

Summary & Outlook

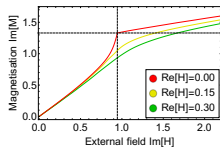
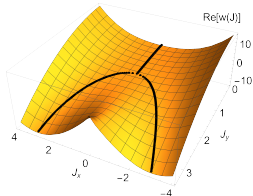
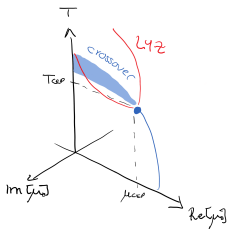
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⇒ Go towards QCD phase structure and complex μ !

Thank you for your attention!

Friederike Ihssen
ITP Heidelberg



A Functional Approach

The infrared regularized path-integral / generating functional:

$$\mathcal{Z}_k[J] = \int [d\varphi]_{\text{ren}, p^2 \geq k^2} \exp \{ - S[\varphi] + J \cdot \varphi \},$$

$$\int [d\varphi]_{\text{ren}, p^2 \geq k^2} = \int [d\varphi]_{\text{ren}} \exp \left\{ - \frac{1}{2} \varphi \cdot R_k \cdot \varphi \right\}$$

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- Suppression of IR momentum modes by k -dependent mass term.

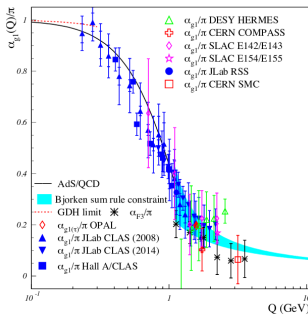
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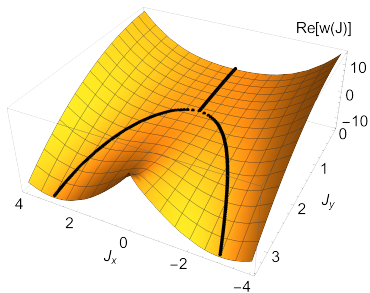
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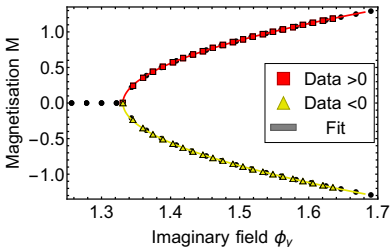
- Suppression of IR momentum modes by k -dependent mass term.
- Access to non-perturbative IR regime by stepwise integration of low energy momentum modes.



$$M(\phi_y, H) = \phi_{\text{EoM}}, \quad \text{with} \quad \partial_{\phi_x} w[J(\phi_x, \phi_y)] - H|_{\phi_x = \phi_{x, \text{EoM}}} = 0$$



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Scaling relations:

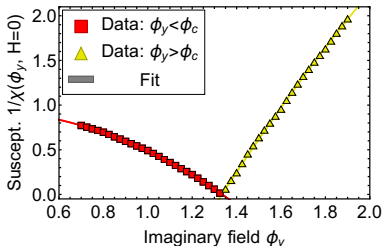
- Preserves symmetry:

$$\text{Re} [M(\phi_y, H = 0)] = B \left(\frac{\phi_y - \phi_c}{\phi_c} \right)^\beta$$

Mean Field: $\beta = 1/2$

Fit: 0.505(23)

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Scaling relations:

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- Breaks symmetry:

$$\text{Re} [M(\phi_y = \phi_c, H)] = B_c H^{1/\delta}$$

Mean Field: $|C_+| = 2|C_-|$

$$C_+ = -0.4369(11), \quad C_- = 0.2025(21)$$

$$R_\chi = \frac{C_+ B^{\delta-1}}{B_c^\delta} = 1.010(57)$$

- Susceptibility:

$$\chi(\phi_y, H = 0) = C_{+/-} \left(\frac{\phi_y - \phi_c}{\phi_c} + \text{subl.} \right)^{-\gamma}$$