



Nuclear Science
Computing Center at CCNU



QCD phase transition from lattice QCD

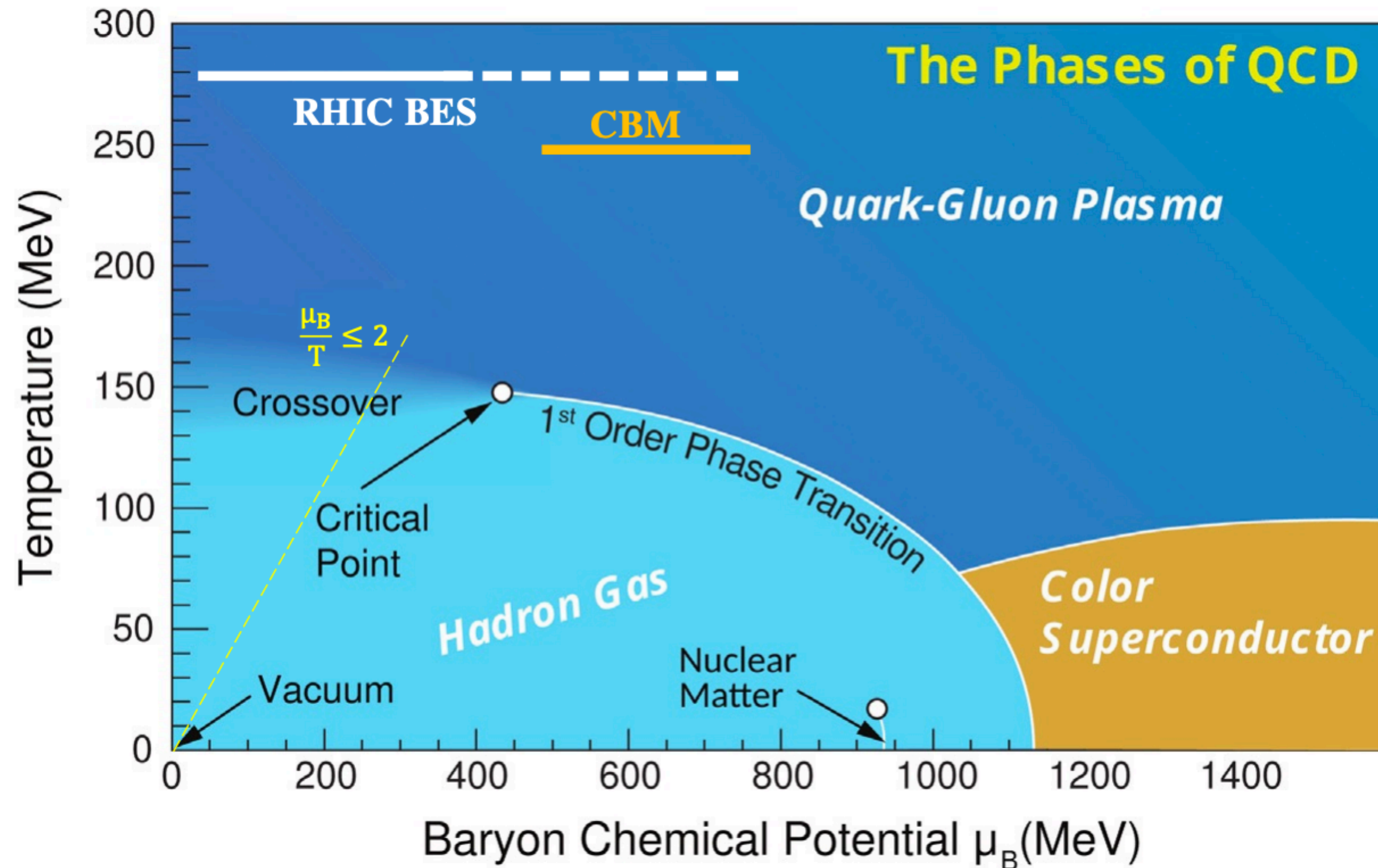
small quark mass & strong magnetic field

Heng-Tong Ding
Central China Normal University

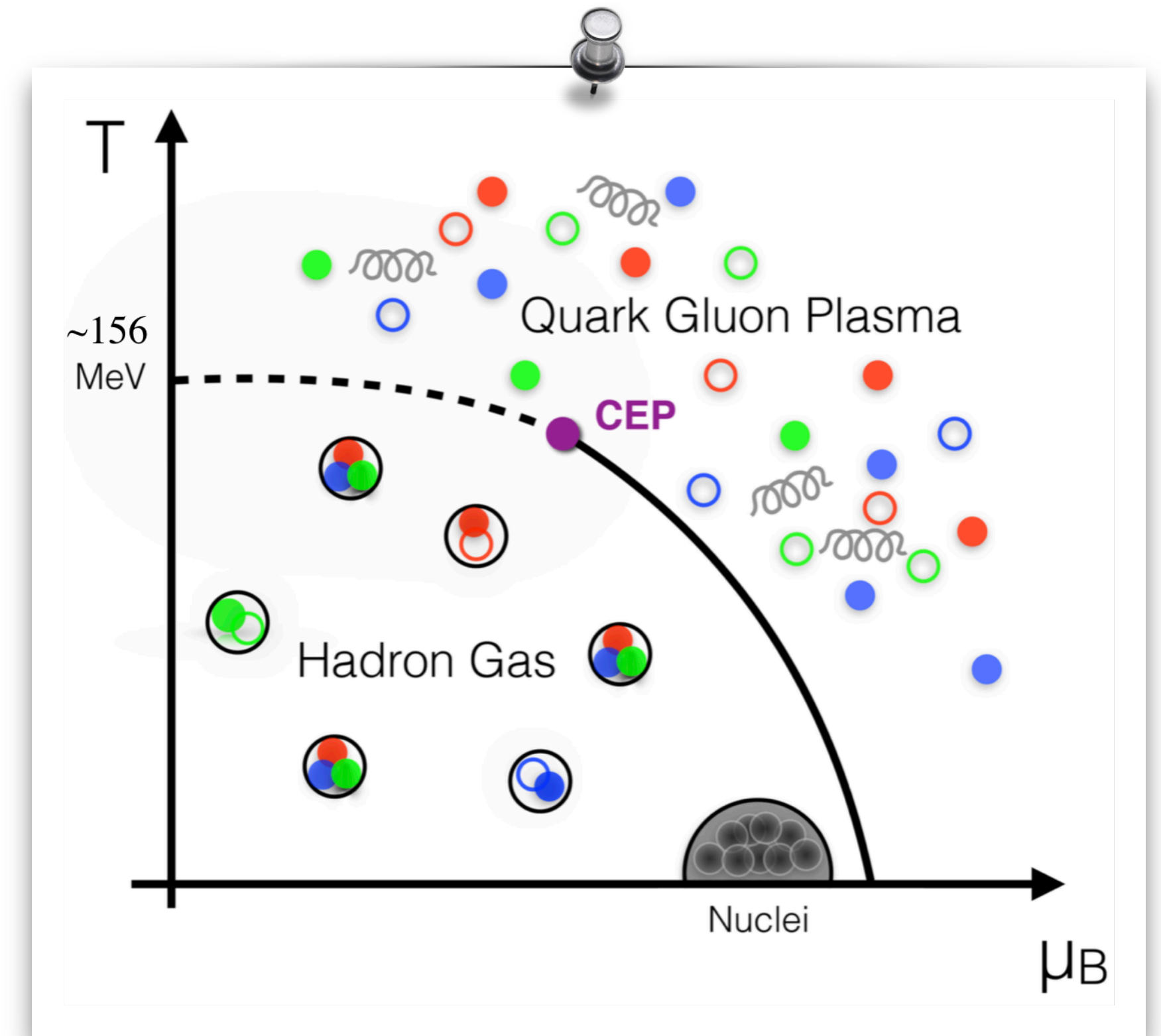
Online workshop on Critical Point and Onset of Deconfinement (CPOD2022)

28 Nov. - 2 Dec., 2022

Search for criticality



Almaalol et al., arXiv:2209.05009



HTD, F. Karsch, S. Mukherjee, arXiv:1504.05274

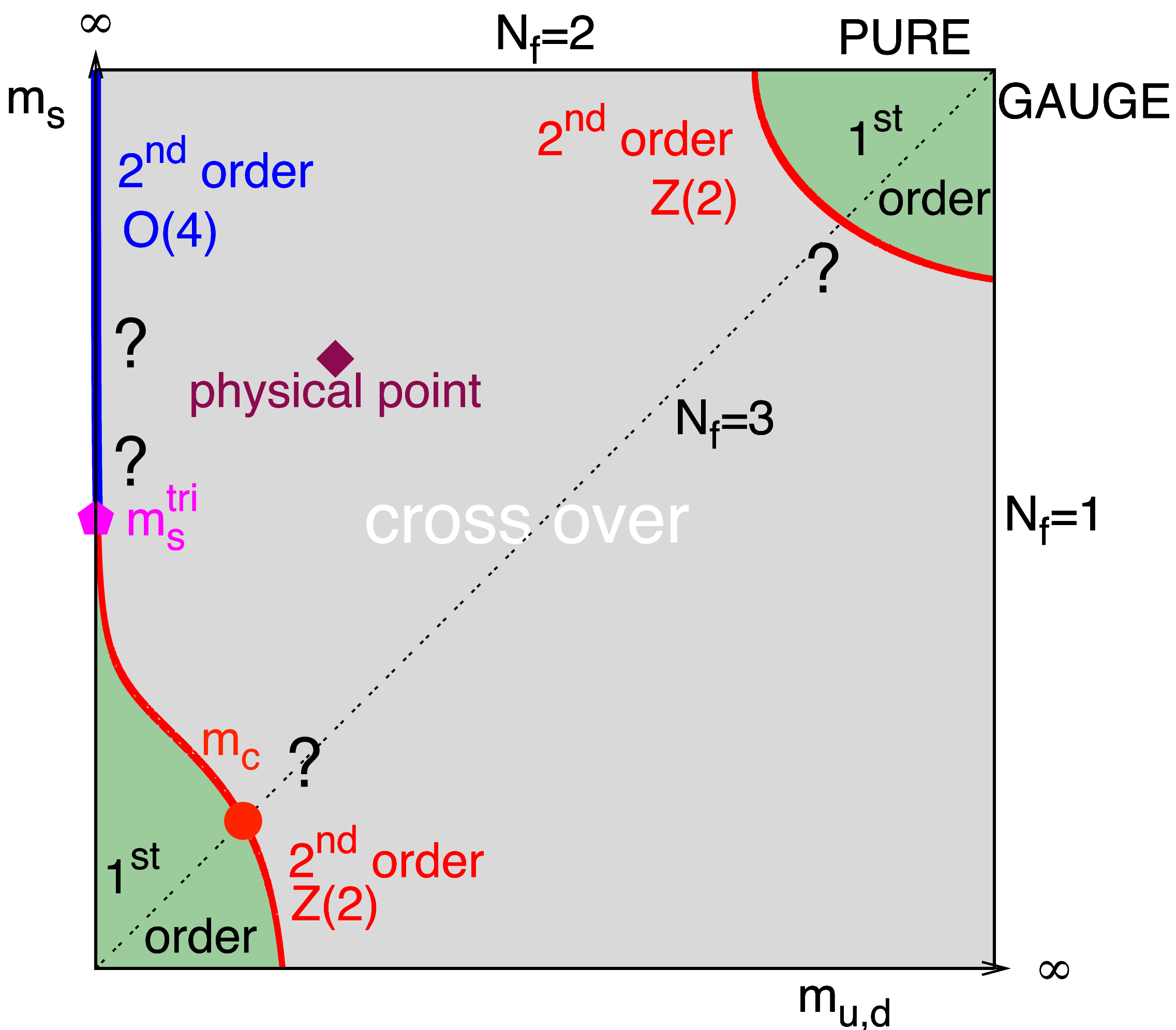
Sign Problem at $\mu_B \neq 0$

Taylor Expansion: **Jishnu Goswami [Tue]**

Imaginary μ_B : **David Clark [Tue]**

QCD criticality at $\mu_B=0$: relevance to CEP?

Columbia plot:
QCD phase diagram in quark mass plane



RG arguments: Pisarski & Wilczek, PRD29 (1984) 338

- $m_q=0$ or ∞ with $N_f=3$: a first order phase transition

- Critical lines of 2nd order transition
 - $N_f=2$: $O(4)$ universality class
 - $N_f=3$: $Z(2)$ universality class

K. Rajagopal & F. Wilczek, NPB 399 (1993) 395
Gavin, Gocksch & Pisarski, PRD 49 (1994) 3079
F. Wilczek, IJMPA 7(1992) 3911

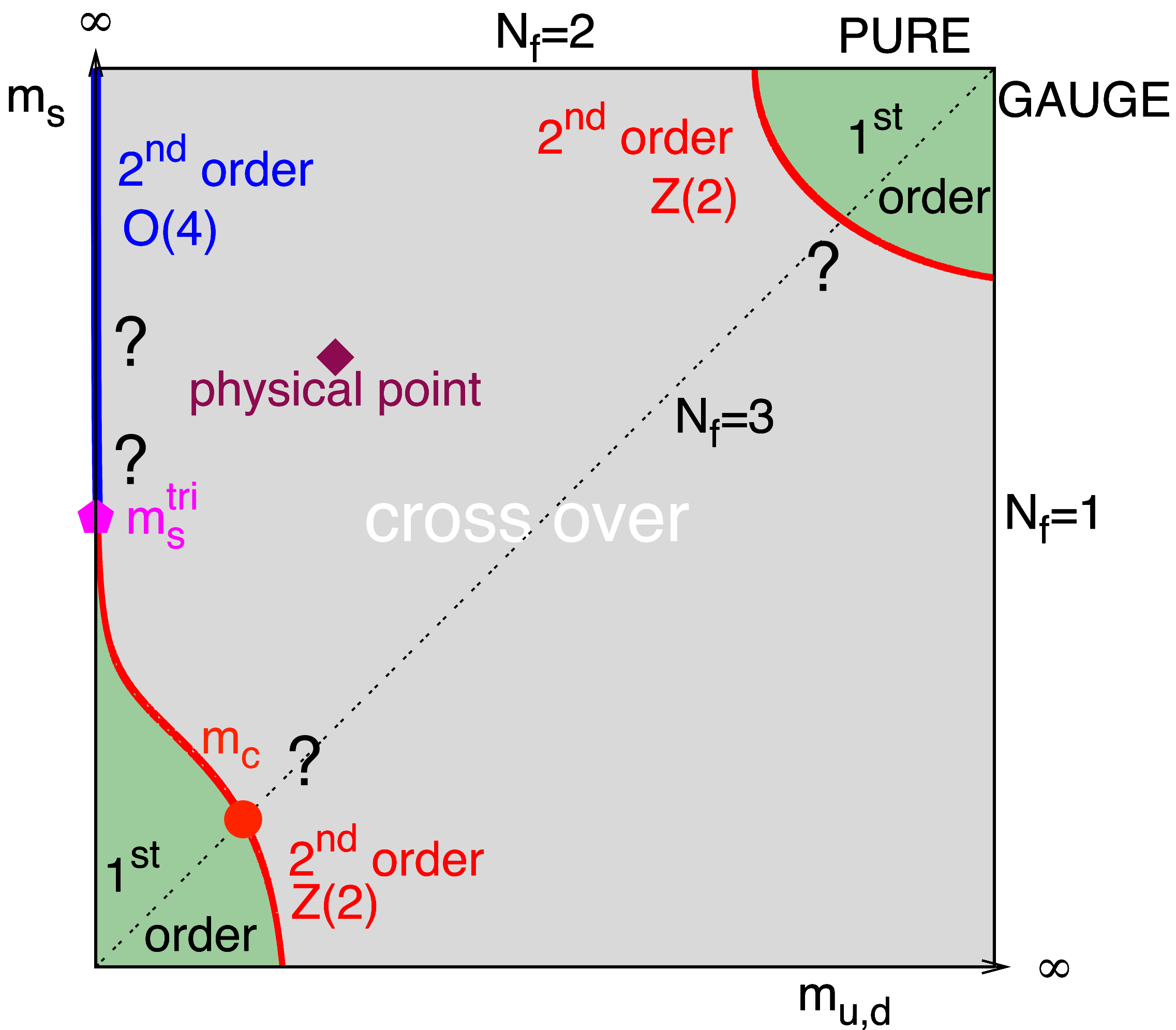
- Axial $U(1)$ anomaly in $N_f=2$ QCD
 - If manifested at T_c : 2nd order $O(4)$
 - If not: 1st order or 2nd order ($U(2)_L \otimes U(2)_R / U(2)_V$)

Butti, Pelissetto and Vicari, JHEP 08 (2003) 029
Pelissetto & Vicari, PRD 88 (2013) 105018
Grahl, PRD 90 (2014) 117904

1st order deconfinement phase transition region

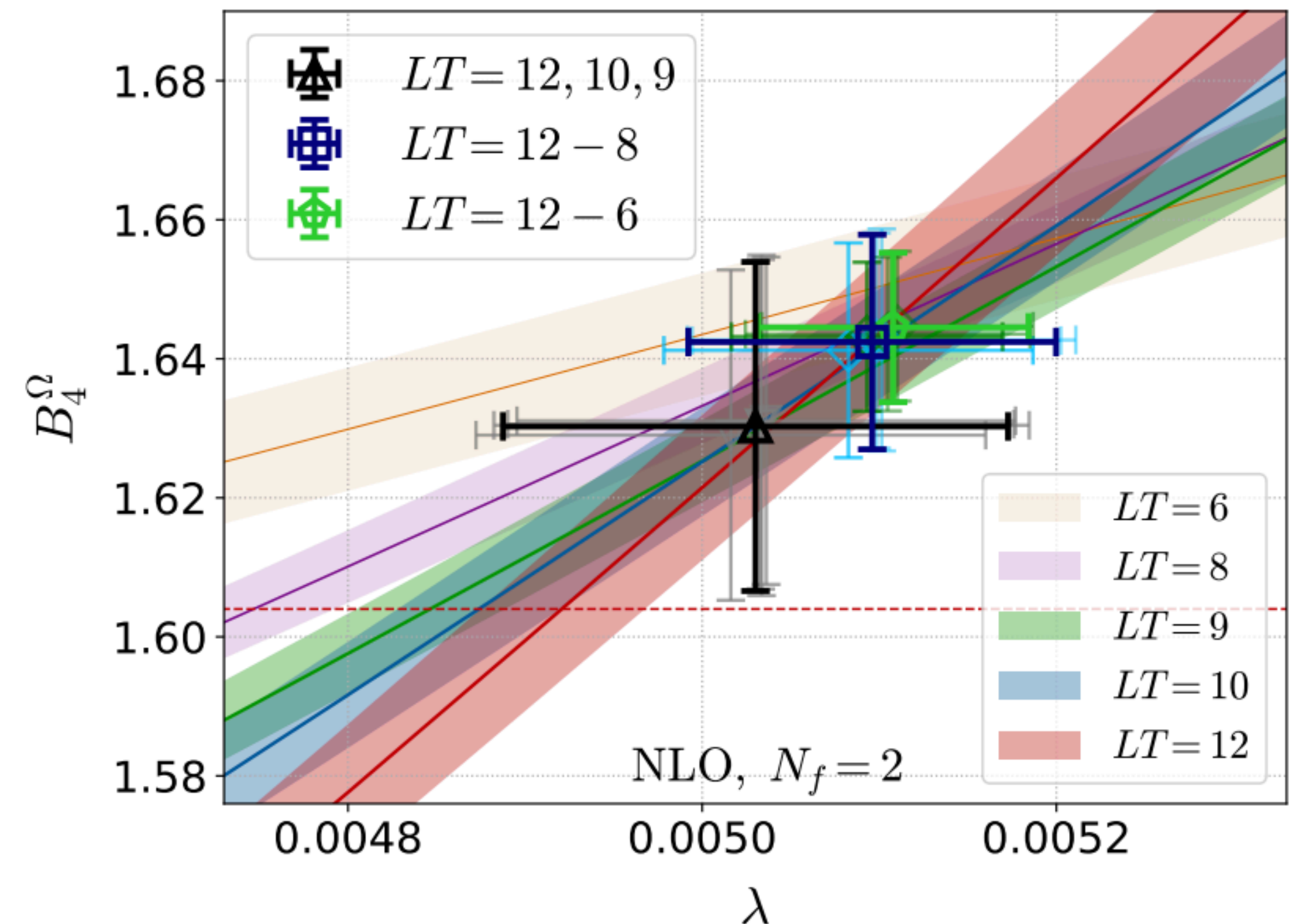
See Masakiyo Kitazawa's talk on Thu

Columbia plot:
QCD phase diagram in quark mass plane



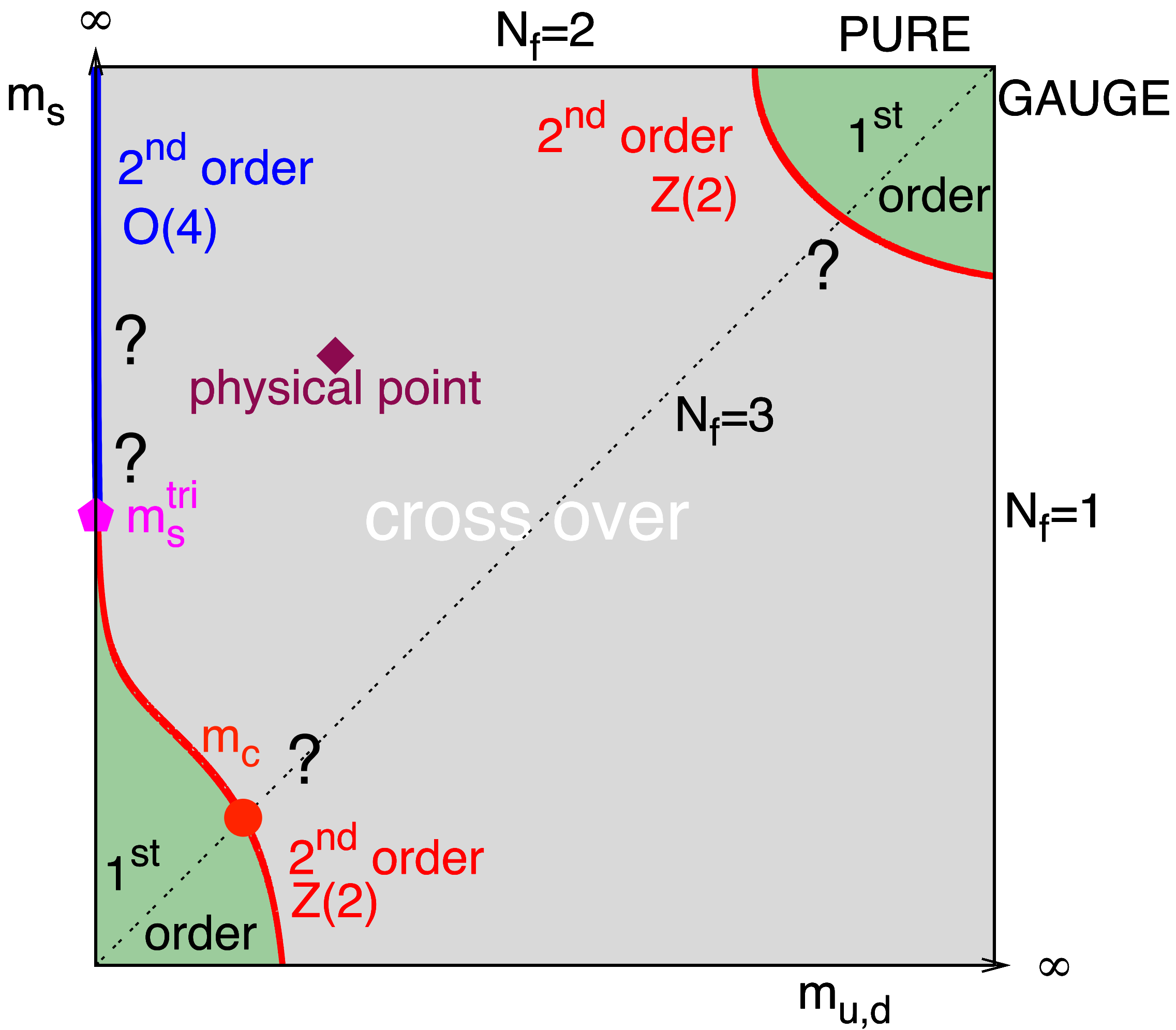
Upper right corner of the Columbia plot:

Observed a 2nd order $Z(2)$ transition in $N_f=2$ QCD with sufficiently heavy quarks

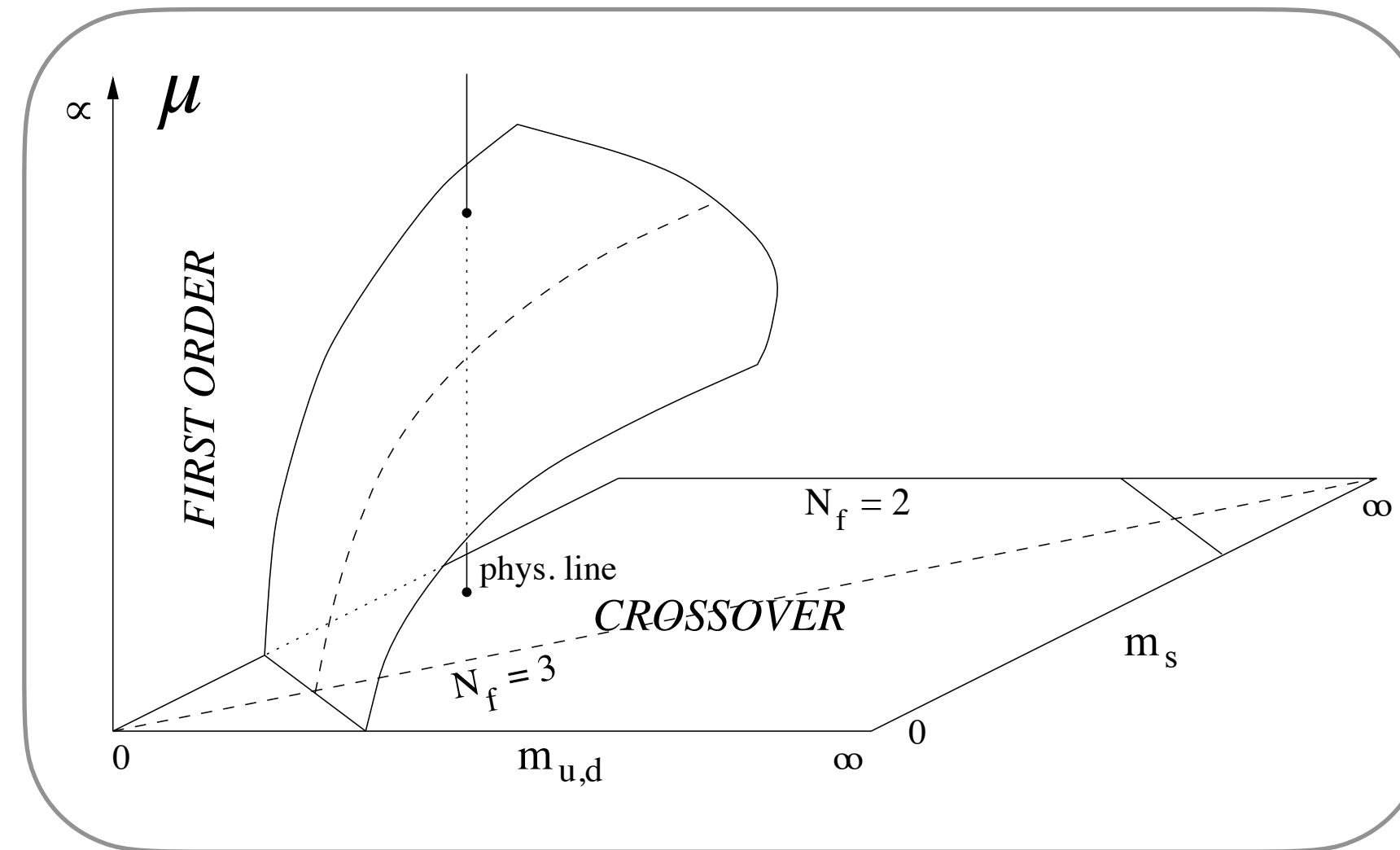


1st order chiral phase transition region

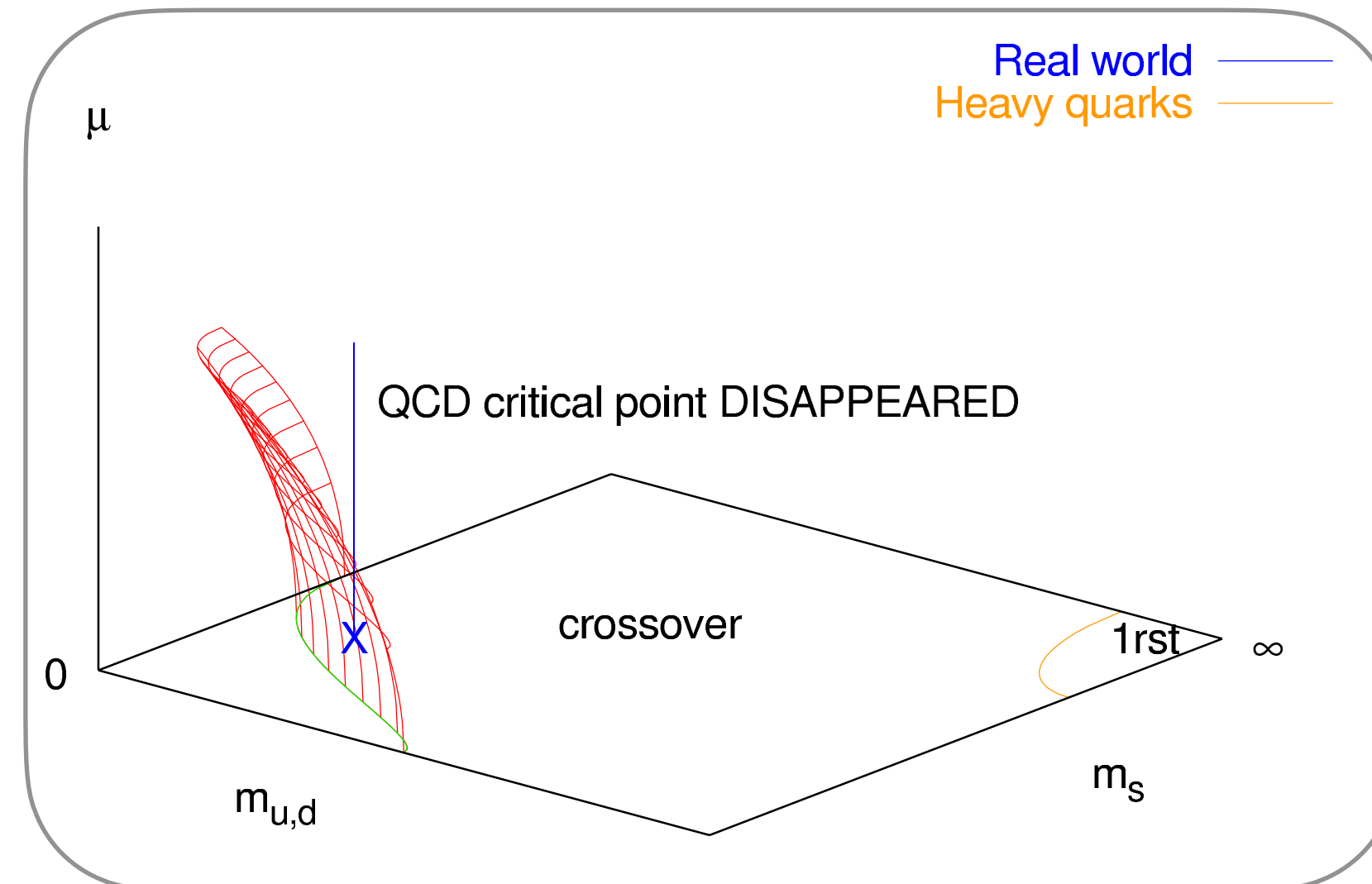
Columbia plot:
QCD phase diagram in quark mass plane



Bottom left corner of the Columbia plot:



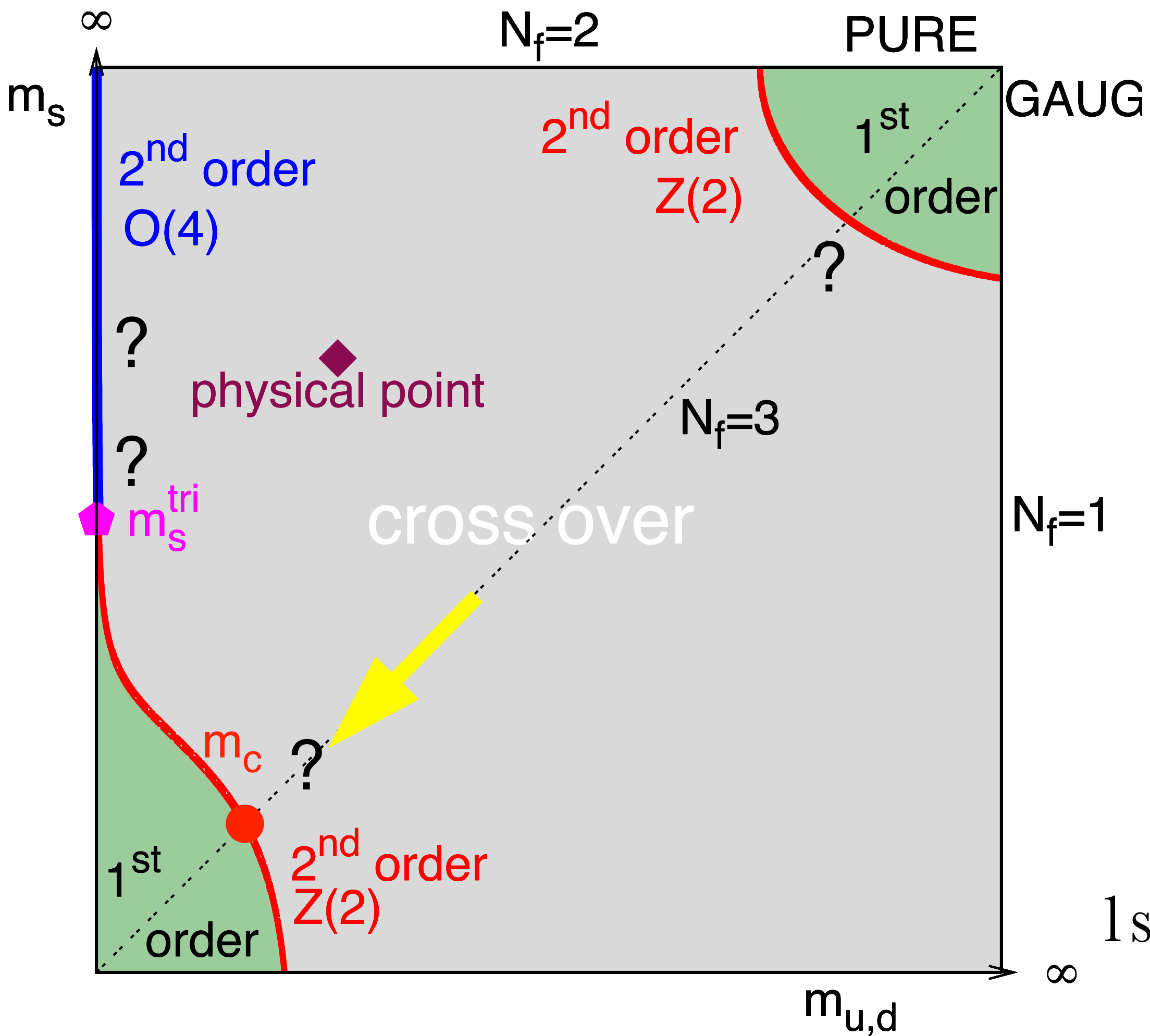
Karsch et al., '03,
Nakamura et al., 15'



de Forcrand & Philipsen, '07

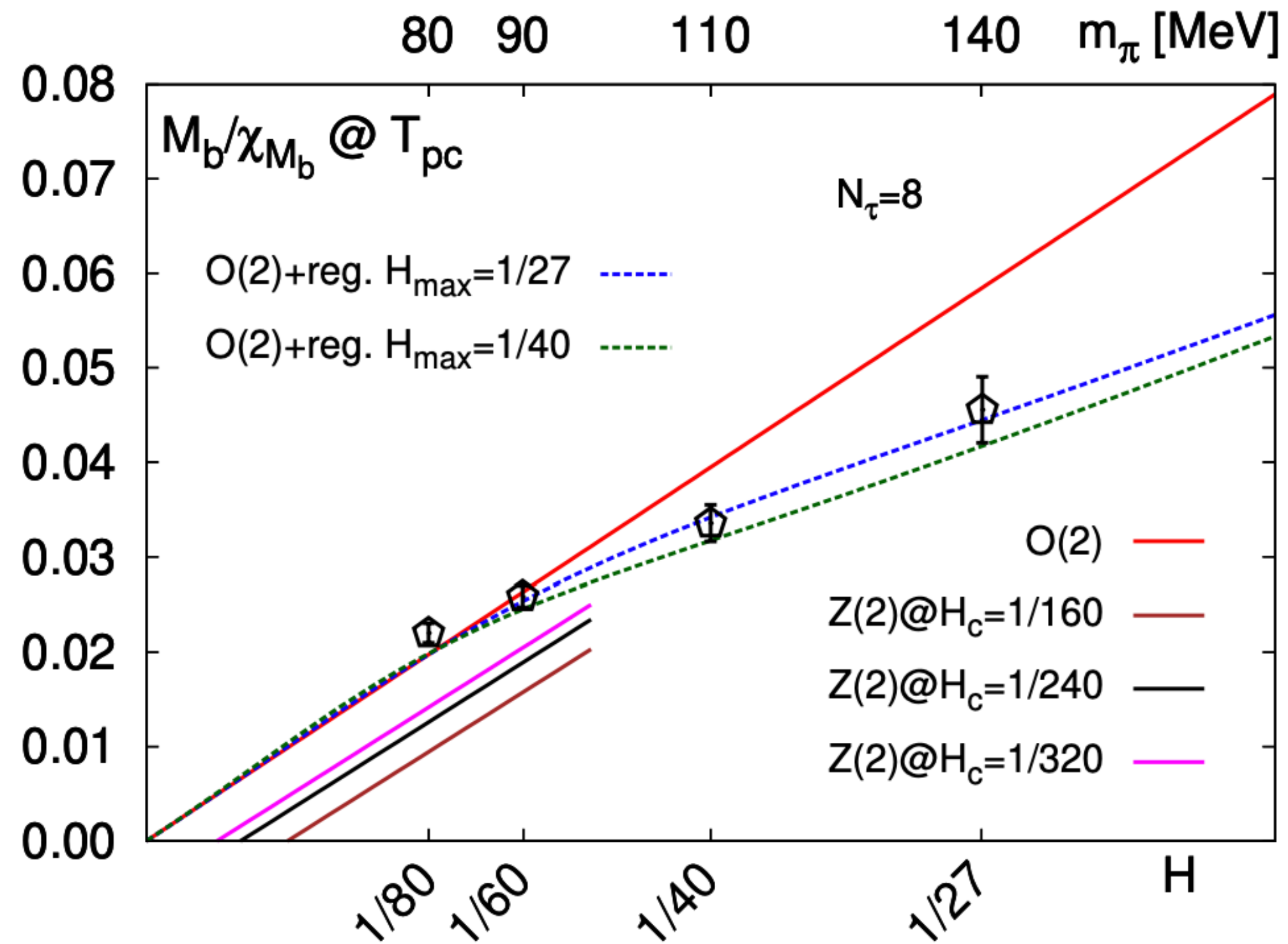
1st order chiral phase transition region

Columbia plot:
QCD phase diagram in quark mass plane



Bottom left corner of the Columbia plot:

HISQ fermions in $N_f=3$ QCD on $N_t=8$ lattices

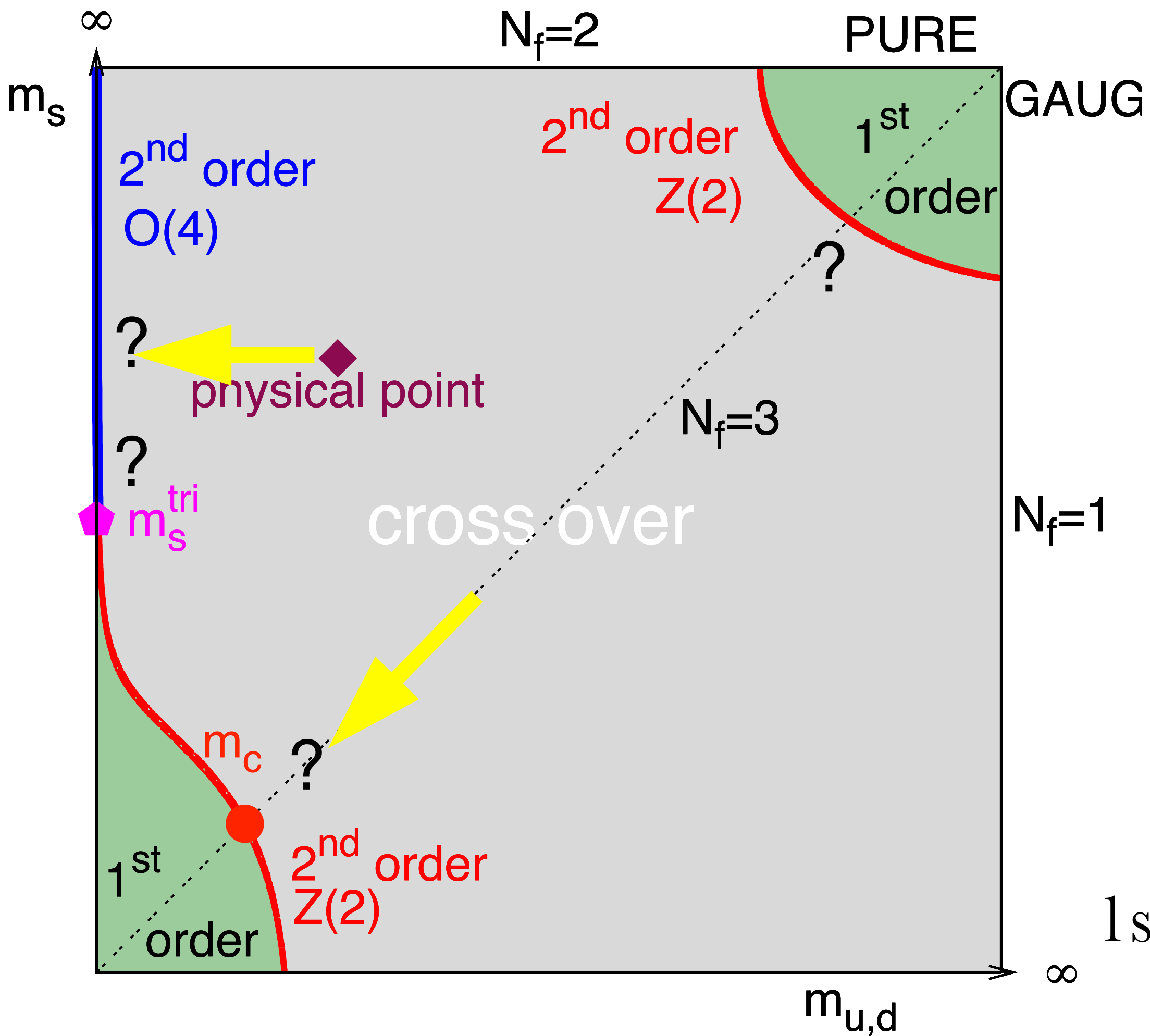


$$\frac{M}{\chi_M} \propto (H - H_c) \frac{f_G}{f_\chi}$$

1st order chiral phase transition region seems to be vanishing

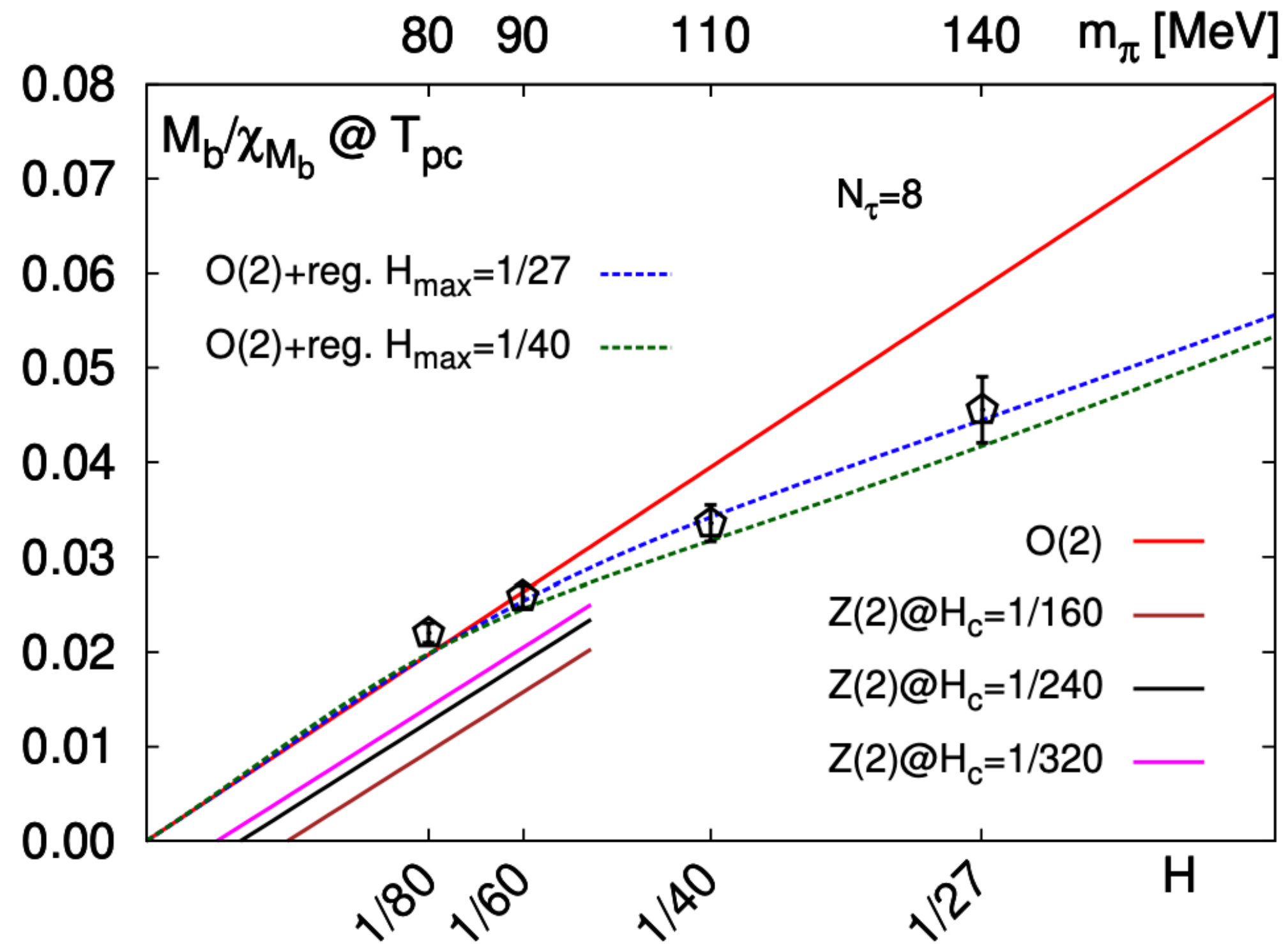
1st order chiral phase transition region

Columbia plot:
QCD phase diagram in quark mass plane



Bottom left corner of the Columbia plot:

HISQ fermions in $N_f=3$ QCD on $N_t=8$ lattices

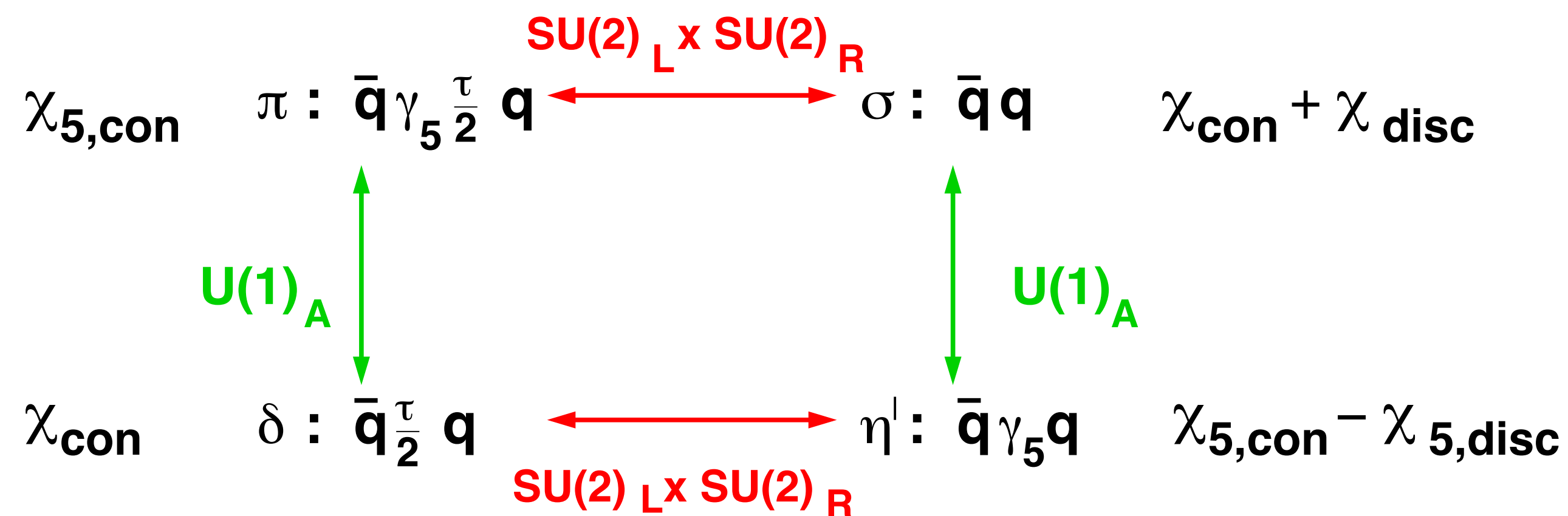


$$\frac{M}{\chi_M} \propto (H - H_c) \frac{f_G}{f_\chi}$$

1st order chiral phase transition region seems to be vanishing

Signatures of symmetry restorations

• \mathcal{S} Susceptibilities defined as integrated two point correlation functions of the local operators, e.g. $\chi_\pi = \int d^4x \langle \pi^i(x) \pi^i(0) \rangle$ with $\pi^i(x) = i\bar{\psi}_l(x) \gamma_5 \tau^i \psi_l(x)$



Restoration of $SU(2)_L \times SU(2)_R$:

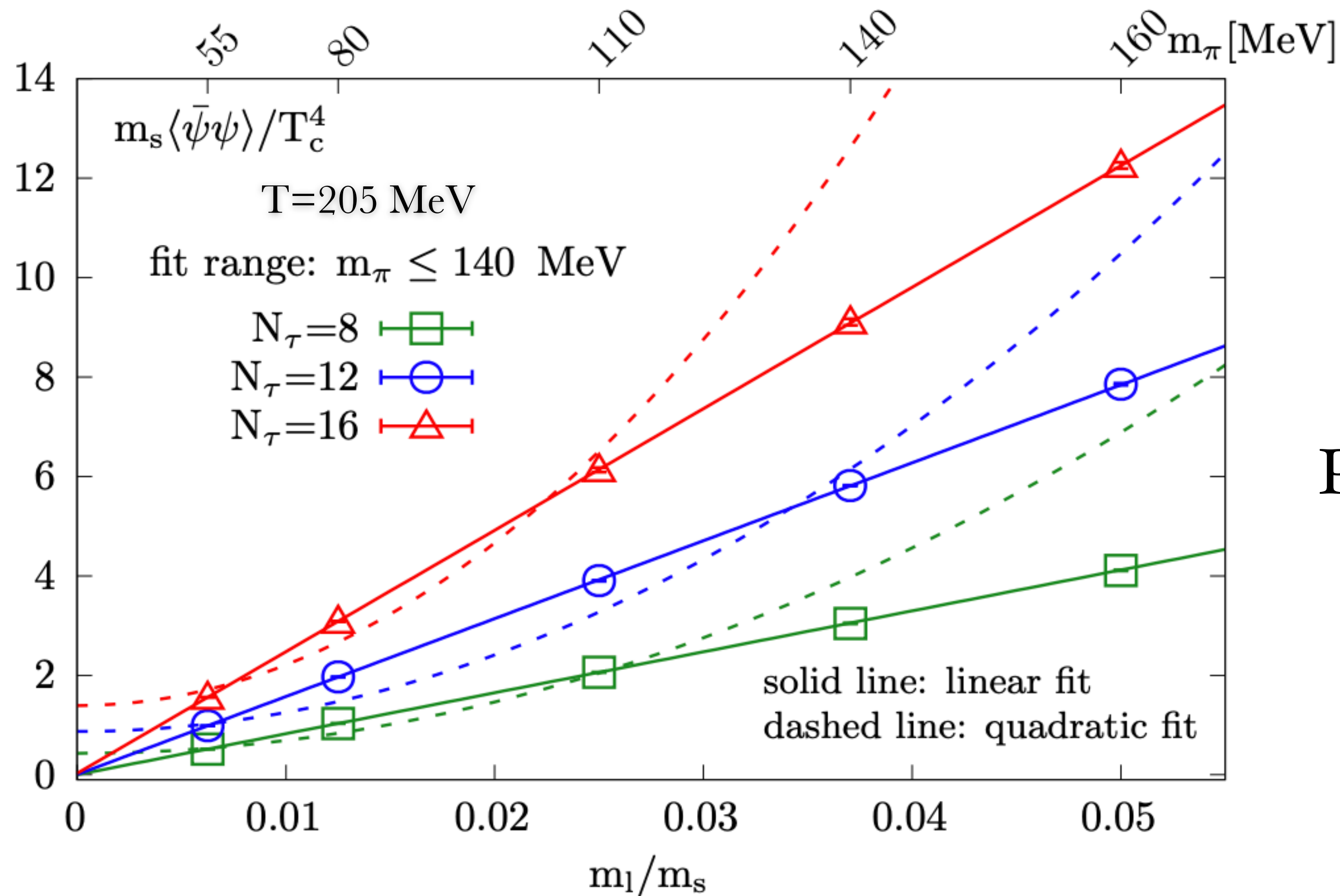
$$\begin{aligned} \chi_\pi - \chi_\sigma &= 0 \\ \chi_\delta - \chi_\eta &= 0 \end{aligned} \quad \Rightarrow \quad \chi_\pi - \chi_\delta = \chi_{\text{disc}} = \chi_{5,\text{disc}}$$

$$\chi_{\text{disc}} = \frac{T}{V} \int d^4x \left\langle \left[\bar{\psi}(x) \psi(x) - \langle \bar{\psi}(x) \psi(x) \rangle \right]^2 \right\rangle$$

Effective restoration of $U(1)_A$:

$$\begin{aligned} \chi_\pi - \chi_\delta &= 0 \\ \chi_\sigma - \chi_\eta &= 0 \end{aligned} \quad \Rightarrow \quad \chi_\pi - \chi_\delta = \chi_{\text{disc}} = \chi_{5,\text{disc}} = 0$$

SU(2)xSU(2) symmetry restoration at T=205 MeV



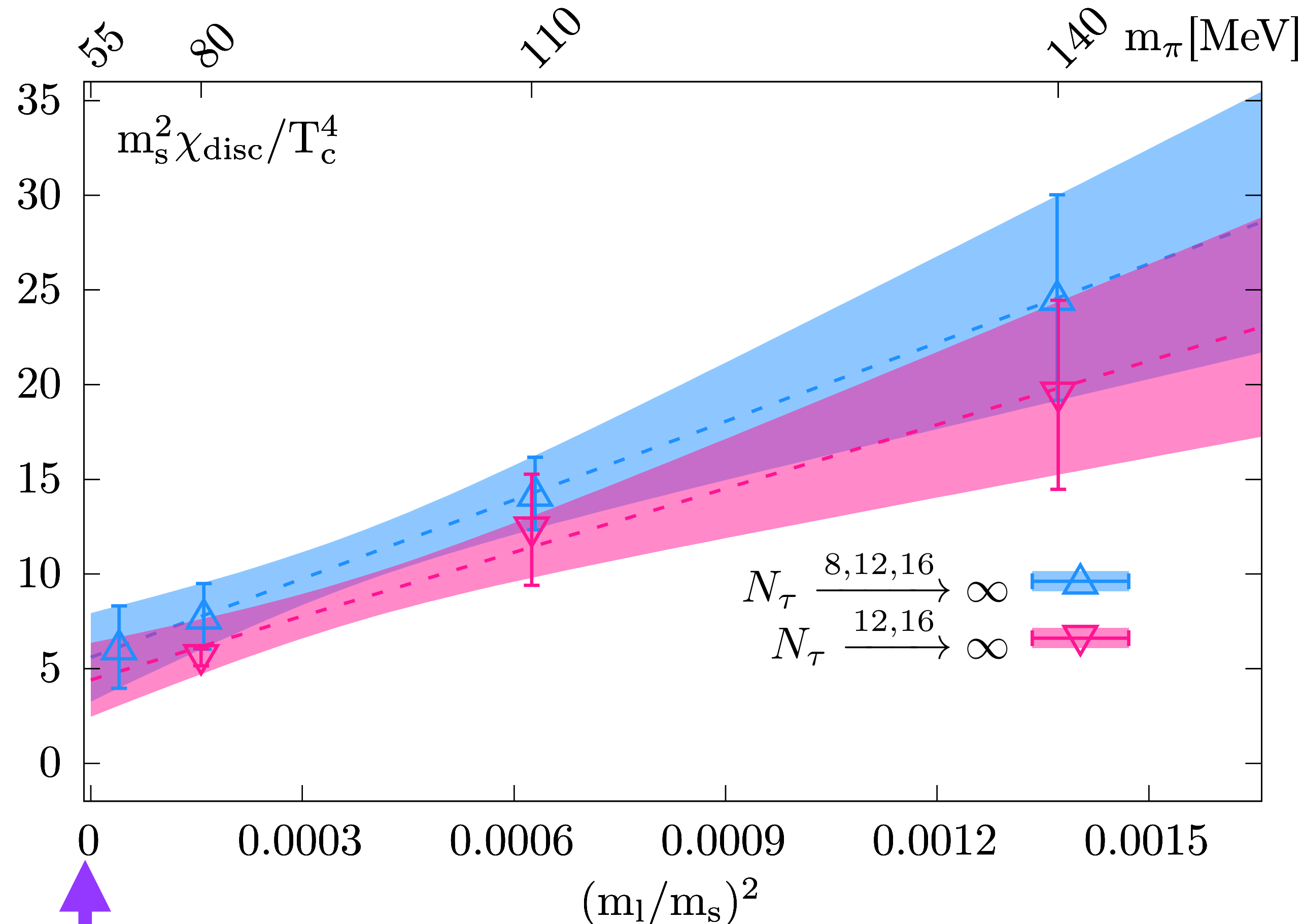
In the chiral symmetric phase
 Z(2) subgroup of SU(2)xSU(2) sym.

Partition function: even function of m

$$\langle \bar{\psi}\psi \rangle \propto m \text{ as } m \rightarrow 0$$

$$\chi_{disc} \propto m^2 \text{ as } m \rightarrow 0$$

Continuum and chiral extrapolations with $m_\pi \leq 140$ MeV data at $T \approx 205$ MeV in $N_f=2+1$ QCD



Joint fit: simultaneous fits

Continuum: $c_0 + c_1/N_\tau^2 + c_2/N_\tau^4$
Chiral: quadratic in quark mass

Value at $N_\tau \rightarrow \infty$ and $m \rightarrow 0$:
 5.6 ± 2.3

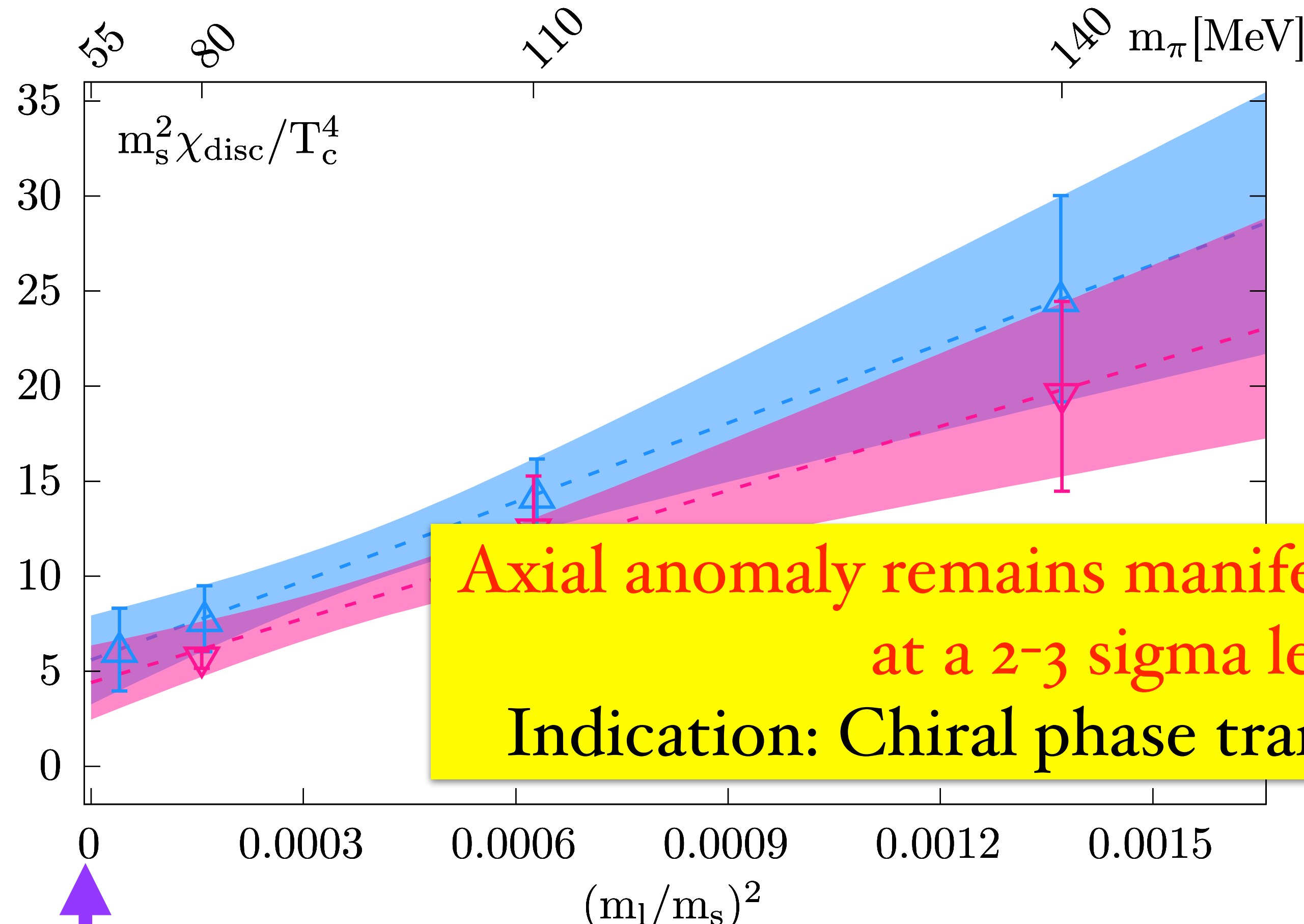
Sequential fit: first continuum and then chiral extrapo.

Continuum: quadratic in $1/N_\tau$ with $N_\tau=12$ & 16 data
Chiral: quadratic in quark mass

Value at $N_\tau \rightarrow \infty$ and $m \rightarrow 0$:
 4.4 ± 1.9

chiral limit

Continuum and chiral extrapolations with $m_\pi \leq 140$ MeV data at $T \approx 205$ MeV in $N_f=2+1$ QCD



Joint fit: simultaneous fits

Continuum: $c_0 + c_1/N_\tau^2 + c_2/N_\tau^4$
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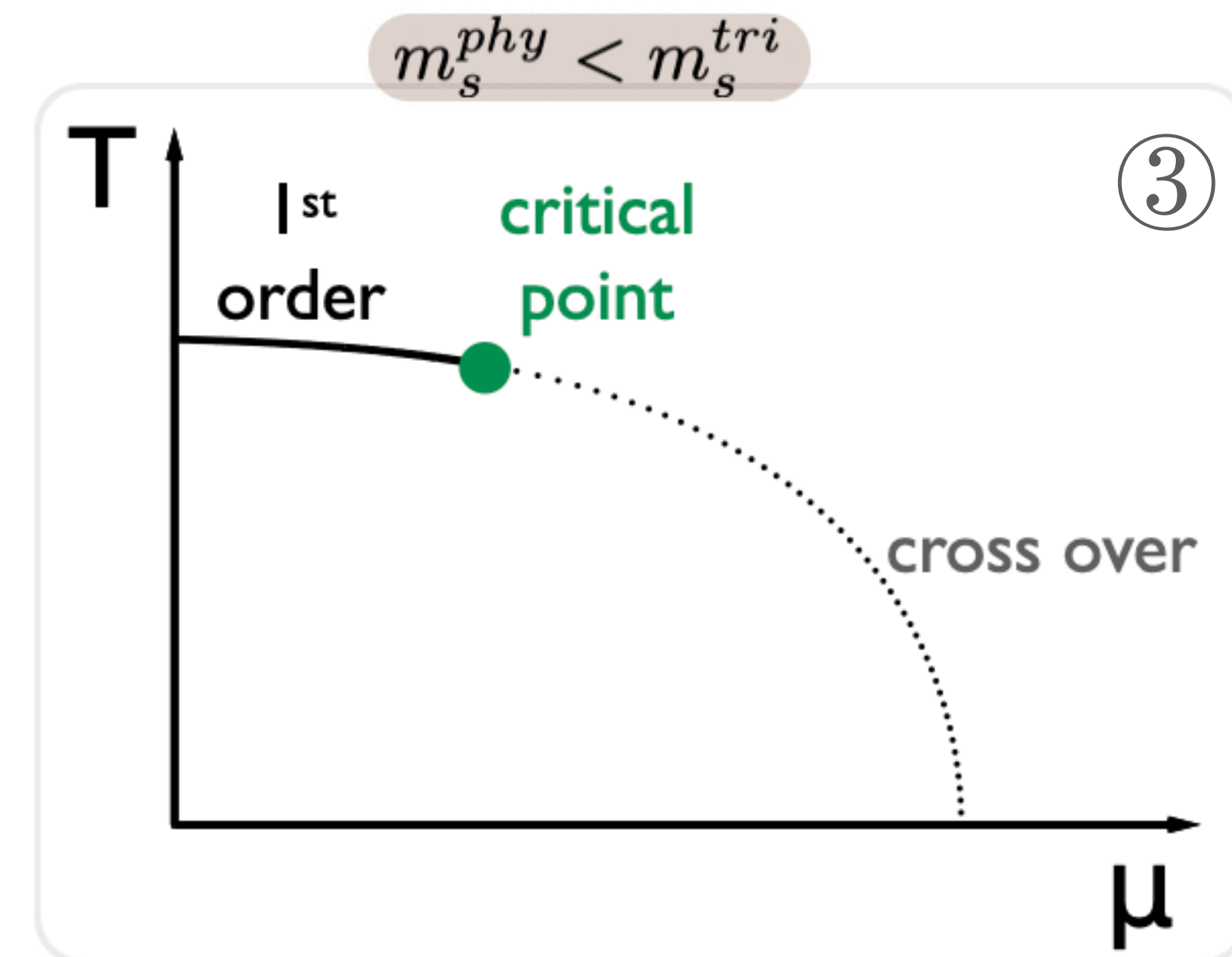
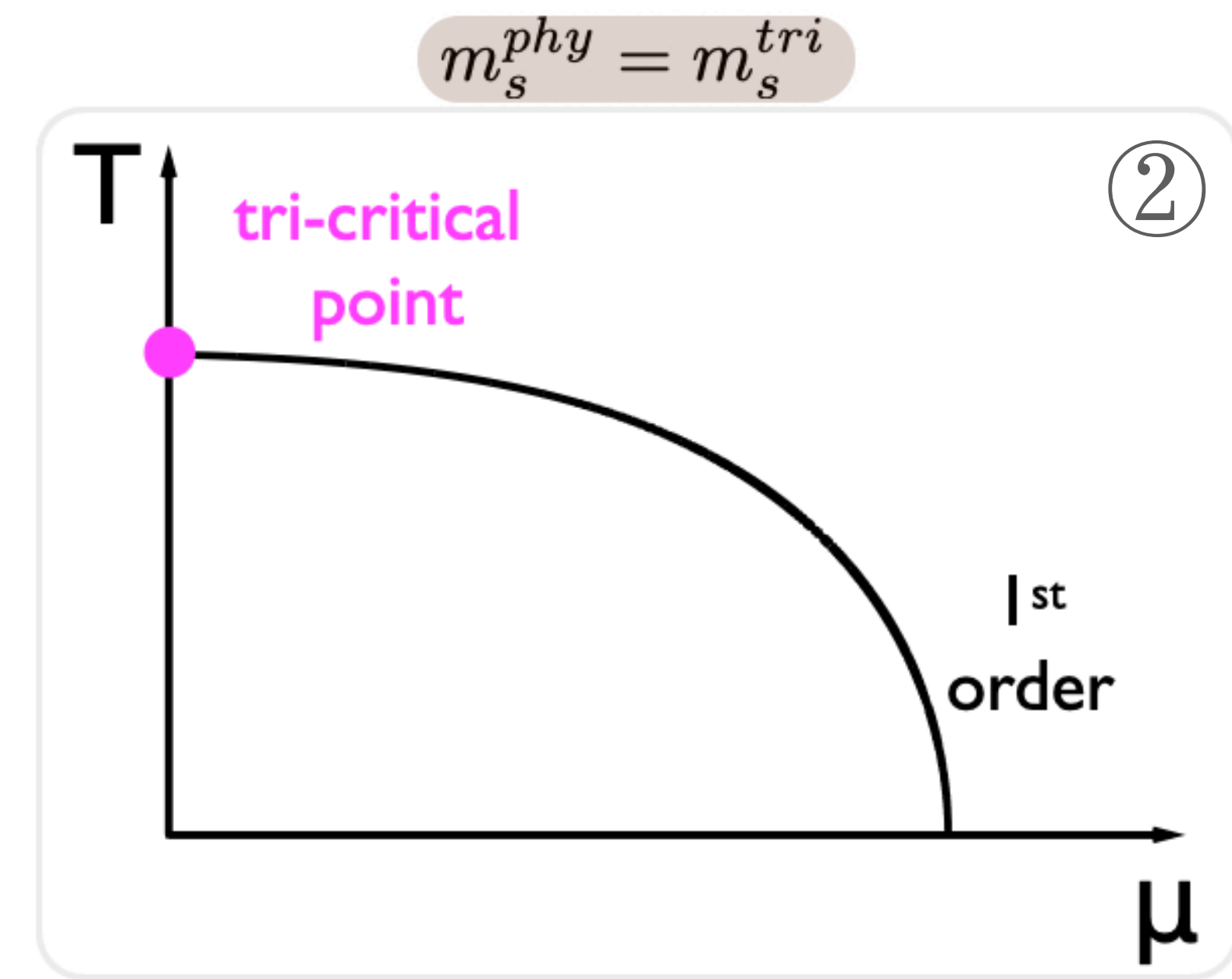
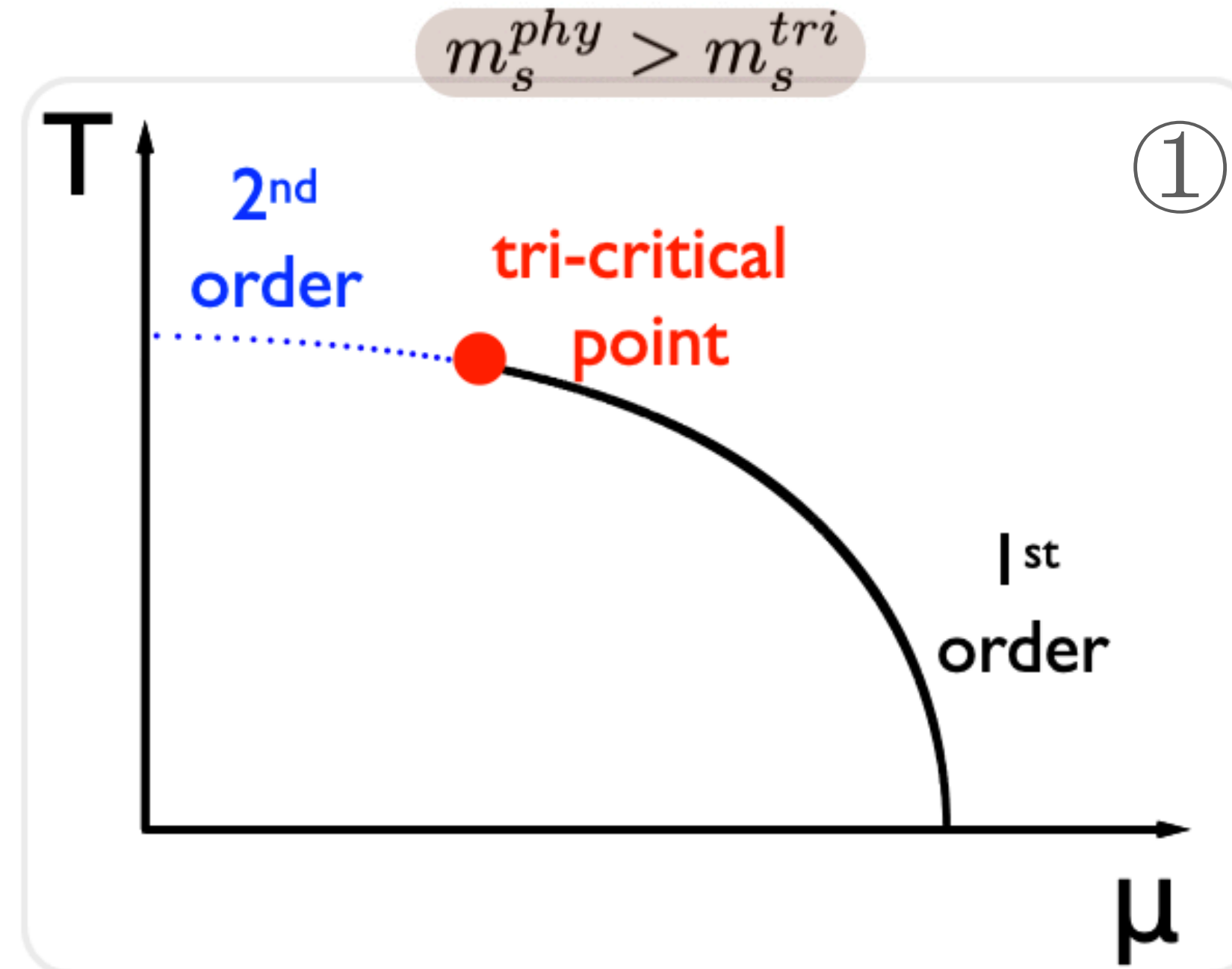
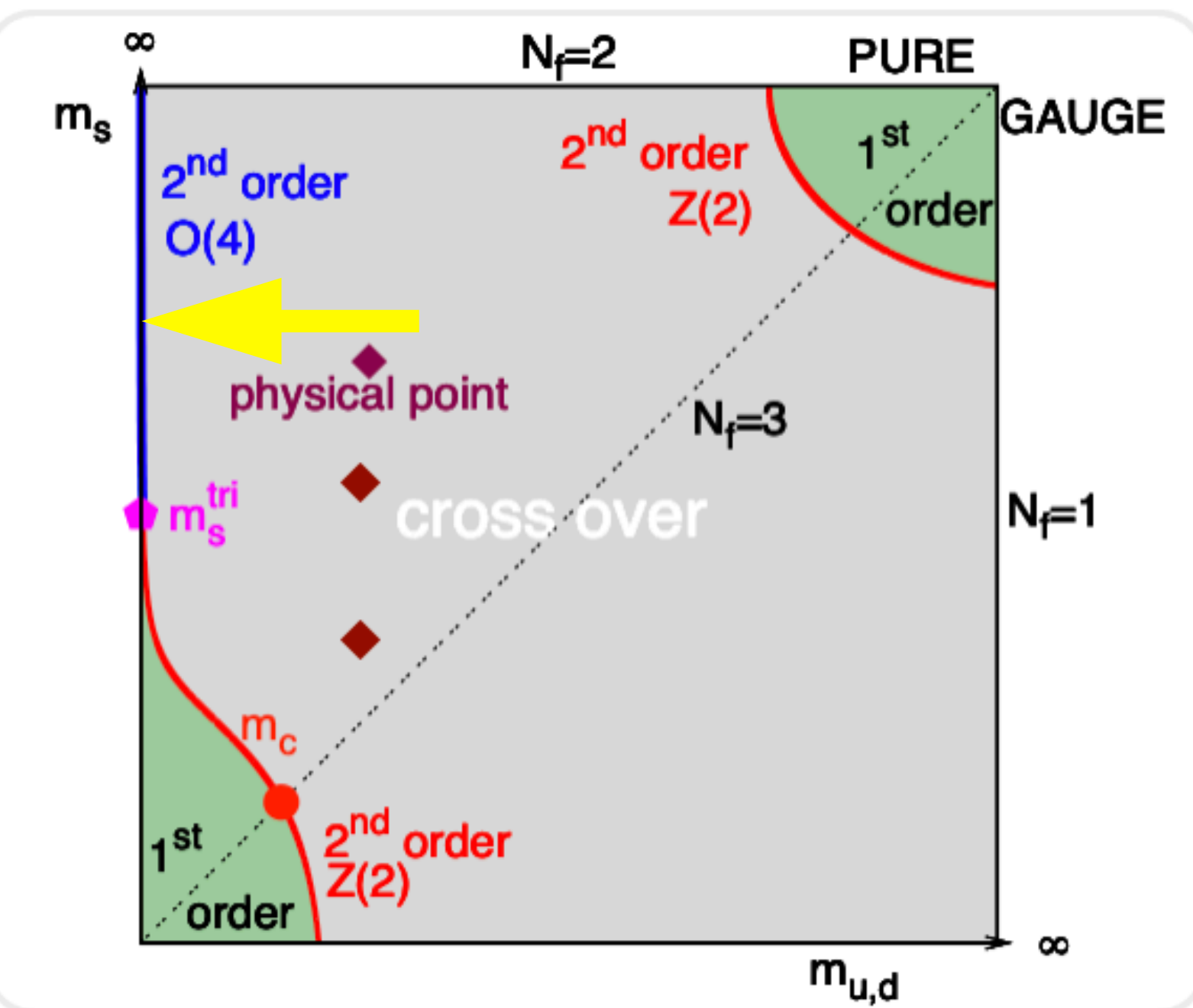
Axial anomaly remains manifested in the $U(1)_A$ measure at a 2-3 sigma level at $1.6T_c$
Indication: Chiral phase transition is 2nd order $O(4)$

Continuum and chiral extrapo.
quadratic in $1/N_\tau$ with
 $N_\tau=12$ & 16 data
Chiral: quadratic in quark mass

Value at $N_\tau \rightarrow \infty$ and $m \rightarrow 0$:
 4.4 ± 1.9

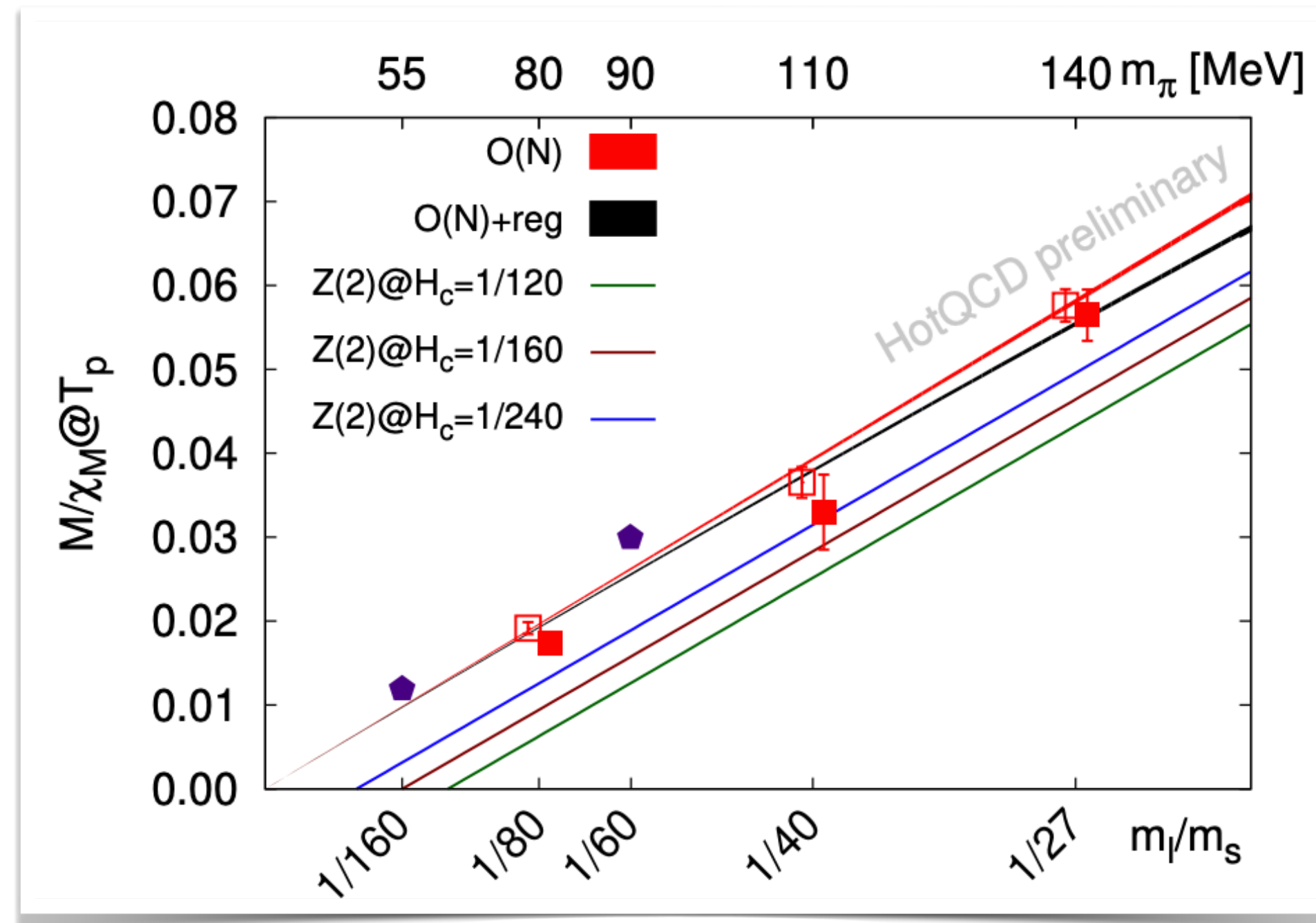
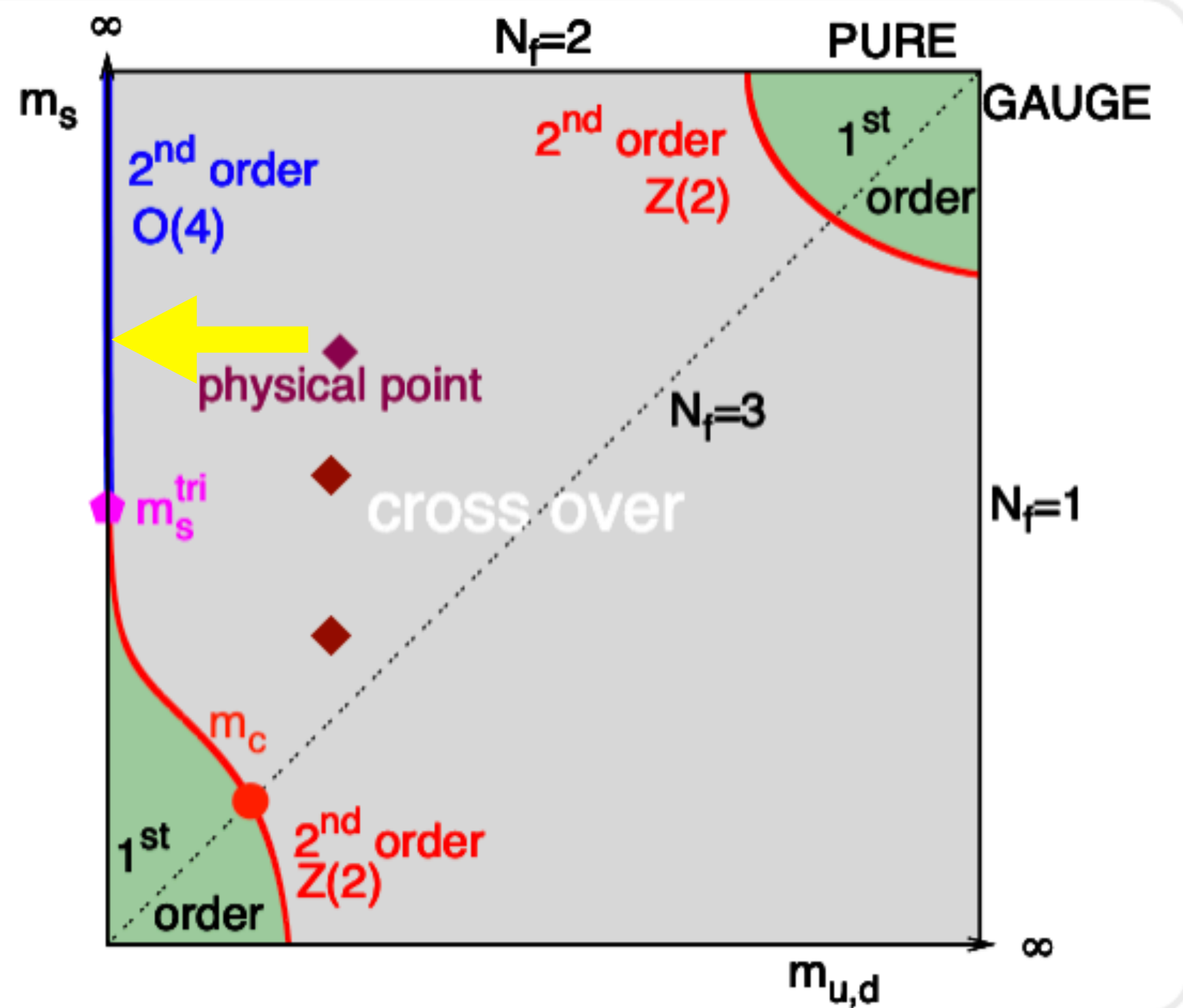
HTD, S.-T. Li, A. Tomiya, S. Mukherjee, X.-D. Wang, Y. Zhang
PRL126(2021)082001

Scenarios of QCD chiral phase transition at $m_l=0$



chiral phase transition in $N_f=2+1$ QCD

Ratio of order parameter and its sus.



$$\frac{M}{\chi_M} \propto \left(\frac{m_l}{m_s} - H_c \right) \frac{f_G}{f_\chi}$$

$H_c \approx 0$

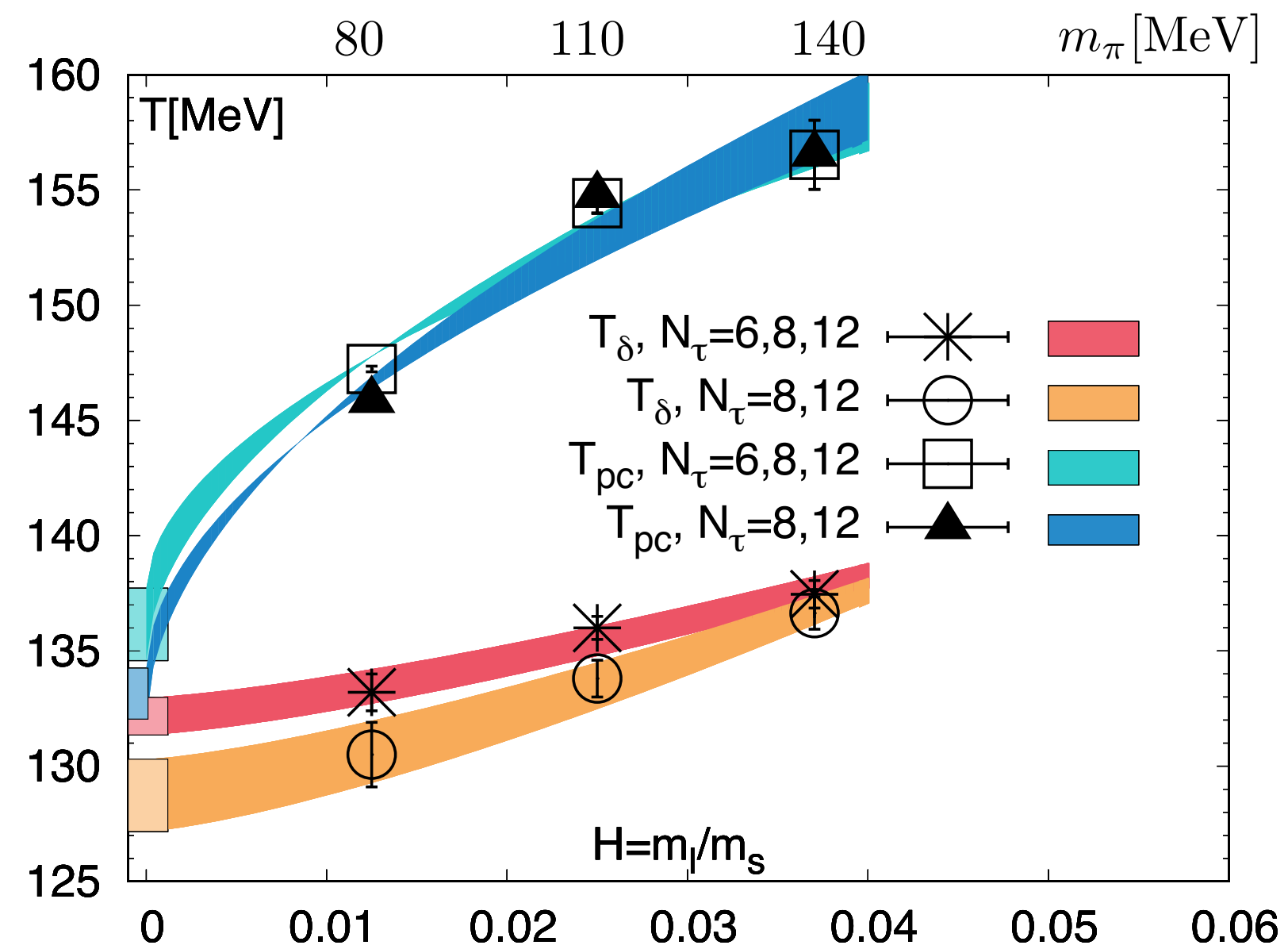
O. Kaczmarek et al., Acta Phys. Pol. B Proc. Suppl. 14 (2021) 291

$m_s^{phy} > m_s^{tri}$: favors a 2nd order $O(4)$ but not $Z(2)$ chiral phase transition

chiral phase transition temperature in $N_f=2+1$ QCD

$$\text{O(4) scaling fit: } T_X(H, L) = T_c^0 \left(1 + \left(\frac{z_X(z_L)}{z_0} \right) H^{1/\beta\delta} \right) + c_X H^{1-1/\delta+1/\beta\delta}$$

HISQ fermions



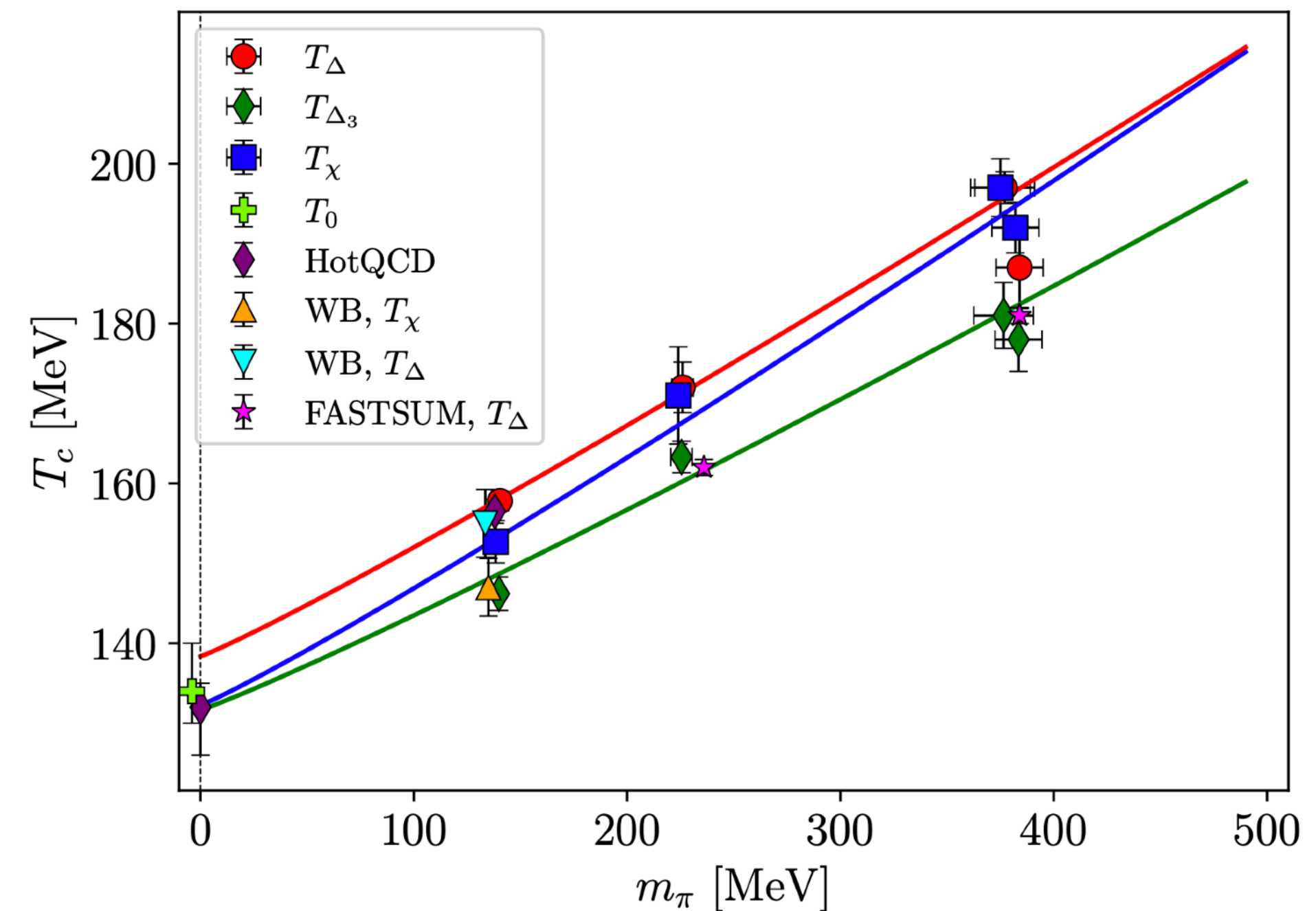
HTD et al.[HotQCD], Phys.Rev. Lett. 123 (2019) 062002,

S.-T. Li et al., Nucl.Phys.Rev. 37 (2020) 3, 674,

O. Kaczmarek et al., Acta Phys. Pol. B Proc. Suppl. 14 (2021) 291

$$T_c^0 = 132_{-6}^{+3} \text{ MeV}$$

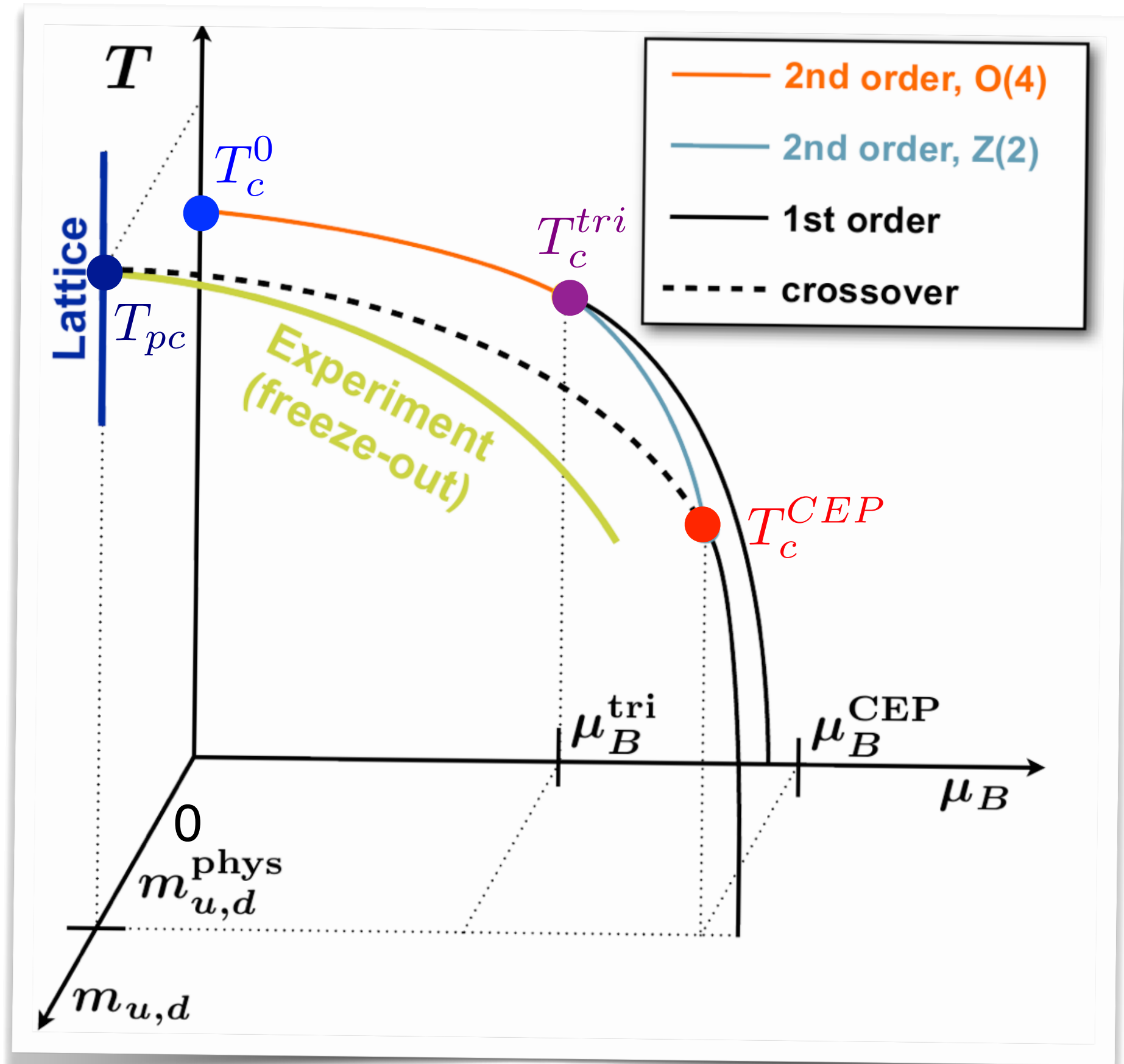
Twisted Wilson fermions



A.Y.Kotov, M.P.Lombardo and A.Trunin, Phys. Lett. B 823 (2021)136749

$$T_c^0 = 134_{-4}^{+6} \text{ MeV}$$

Recap: relevance of criticality at $\mu_B=0$ to CEP



T_{pc} : ≈ 156 MeV, chiral crossover T at $\mu_B=0$

Bazavov et al., [HotQCD] Phys. Lett. B795 (2019) 15
Borsanyi et al., Phys.Rev.Lett. 125 (2020) 052001

T_c^{CEP} : transition T at the critical end point

T_c^0 : ≈ 132 MeV chiral phase transition T at $m_q=0$ and $\mu_B=0$

HTD et al.[HotQCD], PRL123 (2019) 062002
A.Y.Kotov et al., PLB 823 (2021)136749

T_c^{tri} : transition T at the tri-critical point

Random Matrix Model & NJL suggests:

$$T_c^{tri} - T_c^{CEP}(m_q) \propto m_q^{2/5}$$

Y. Hatta & T. Ikeda, PRD67 (2003) 014028
M. A. Halasz et al, PRD 58 (1998) 096007
M. Buballa, S. Carignano, PLB791(2019)361

$T_c^0(\mu_B)$ decreases as μ_B up to NLO from LQCD

O. Kaczmarek et al., PRD83 (2011) 014504
P. Hegde & HTD, PoS LATTICE2015 (2016) 141
O. Kaczmarek et al., PoS LATTICE2021 (2022) 429

Indication

$$T_c^0 > T_c^{tri} > T_c^{CEP}$$

Microscopic manifestation of the symmetry restorations: Dirac Eigenvalues and their correlations C_n

$$\langle \bar{\psi}\psi \rangle = \int_0^\infty \frac{4m_l \rho}{\lambda^2 + m_l^2} d\lambda, \quad \lim_{m_l \rightarrow 0} \langle \bar{\psi}\psi \rangle = \pi \rho(0)$$

$$\text{E.g.: } \frac{V}{T} \frac{\partial \rho}{\partial m_l} = \int_0^\infty d\lambda_2 \frac{4m_l C_2(\lambda, \lambda_2; m_l)}{\lambda_2^2 + m_l^2}$$

$$\chi_1 \equiv \chi_{disc} = \int_0^\infty \frac{4m_l \partial \rho / \partial m_l}{\lambda^2 + m_l^2} d\lambda$$

$$\propto \langle (\bar{\psi}\psi - \langle \bar{\psi}\psi \rangle)^2 \rangle = \kappa_2(\bar{\psi}\psi)$$

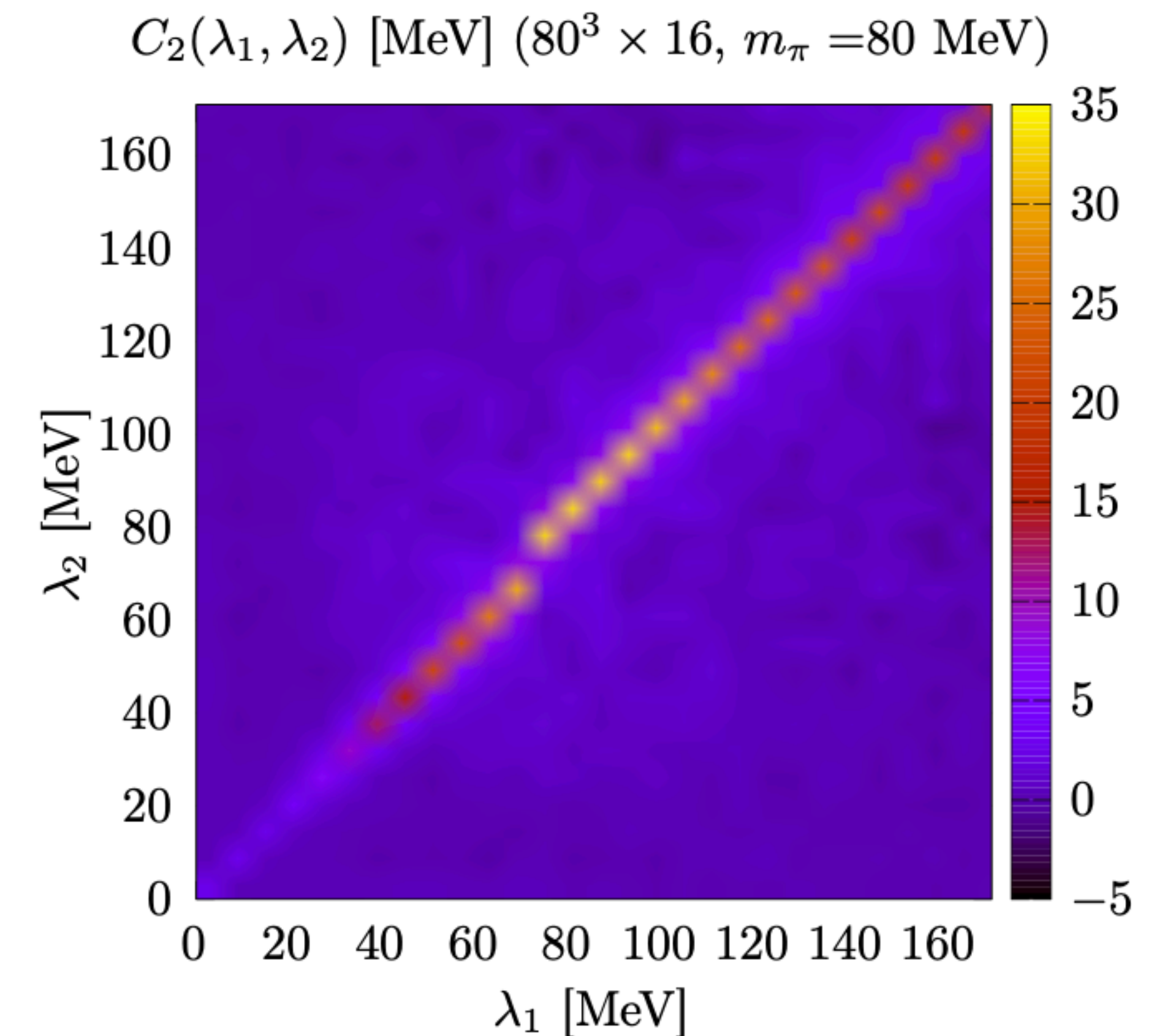
$$\chi_2 = \int_0^\infty \frac{4m_l \partial^2 \rho / \partial m_l^2}{\lambda^2 + m_l^2} d\lambda \propto \kappa_3(\bar{\psi}\psi) + \dots$$

... ..

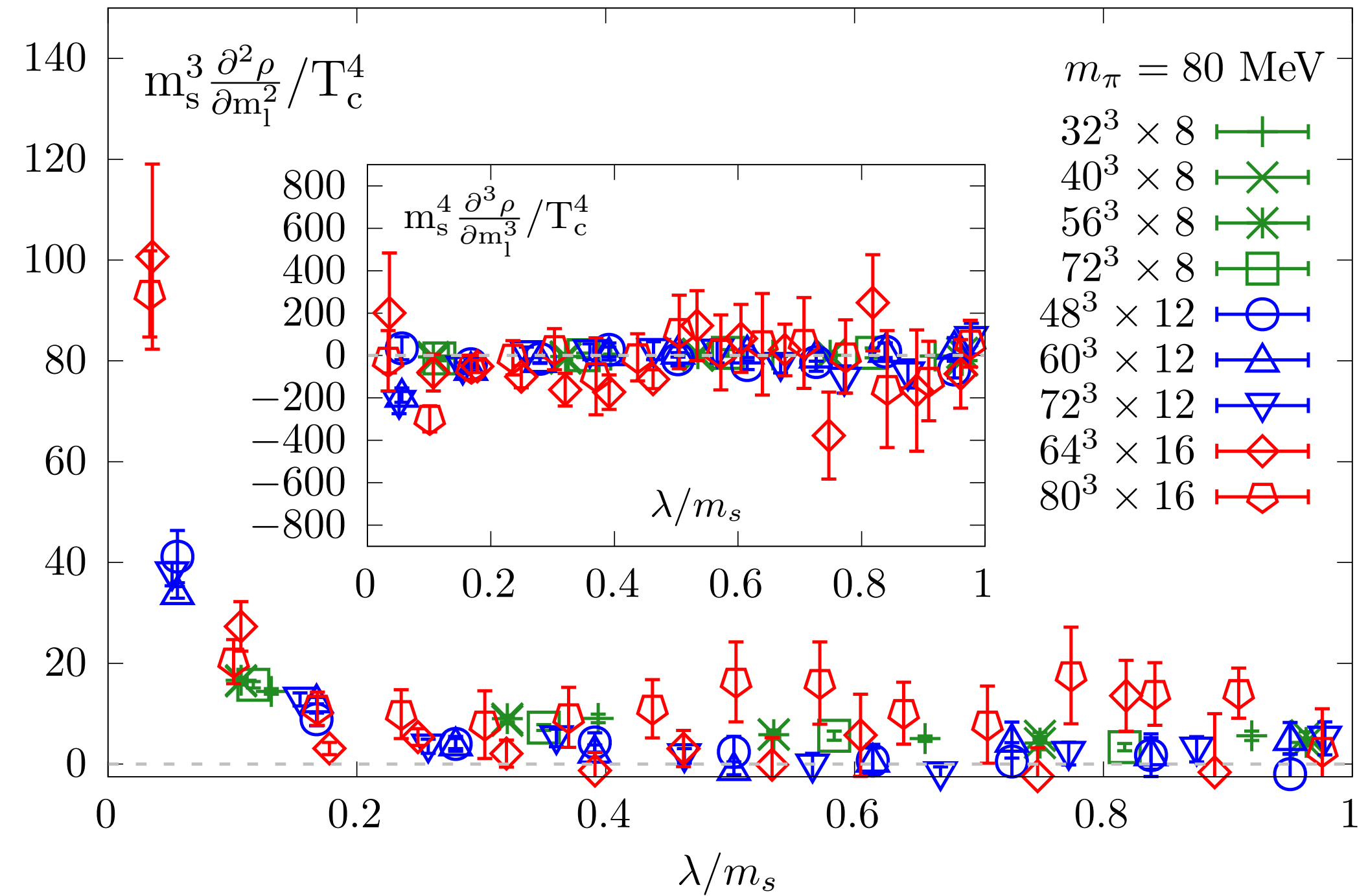
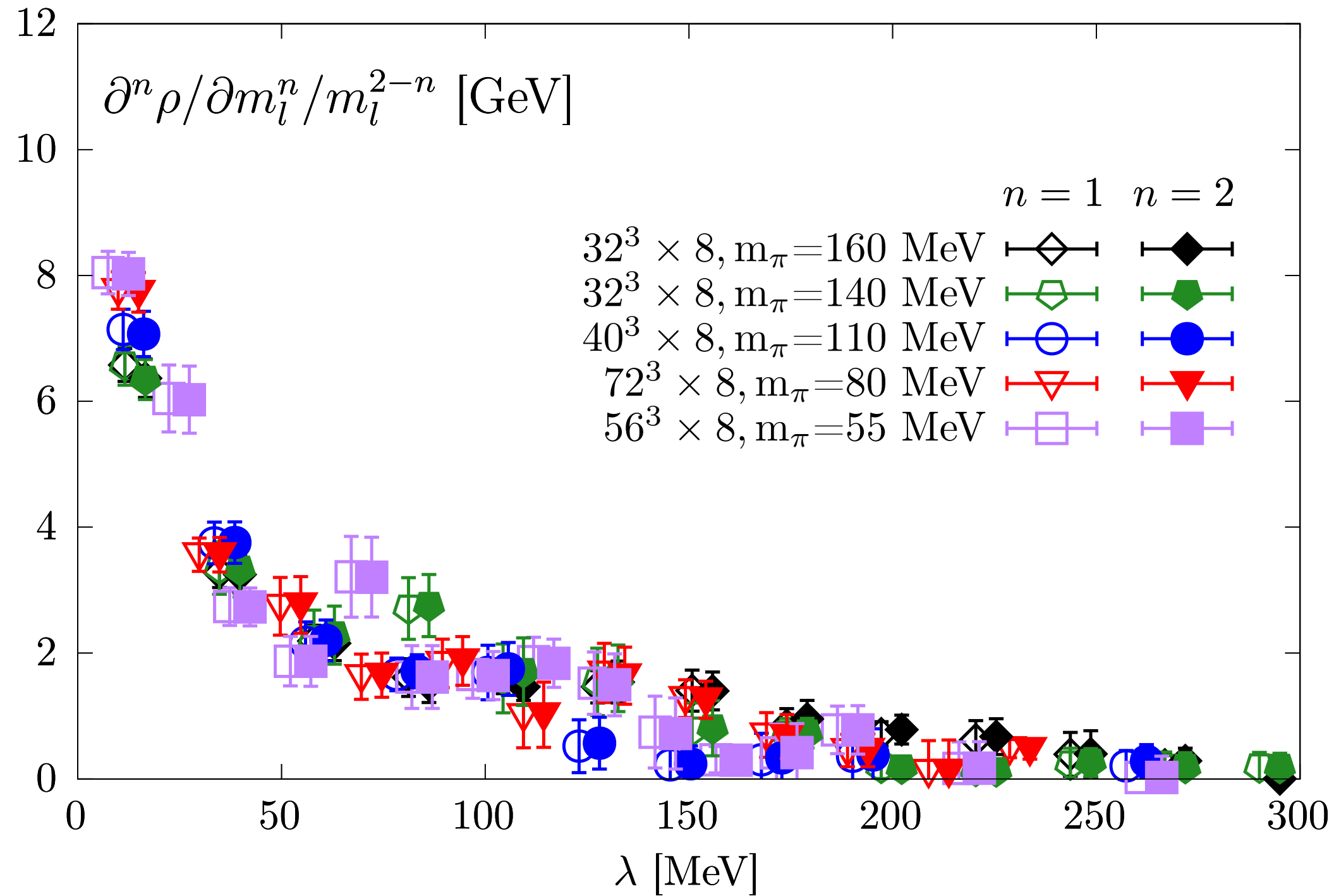
$$\chi_n = \int_0^\infty \frac{4m_l \partial^n \rho / \partial m_l^n}{\lambda^2 + m_l^2} d\lambda \propto \kappa_{n+1}(\bar{\psi}\psi) + \dots$$

$$\partial^n \rho(\lambda) / \partial m_l^n = f(C_{n+1}, C_n, \dots, C_2)$$

$$C_n = \kappa_n(\rho_U(\lambda))$$



1st, 2nd & 3rd quark mass derivative of ρ on $N_\tau=8$ lattices at $T \approx 205 \text{ MeV} \simeq 1.6T_c$



$$m_l^{-1} \partial \rho / \partial m_l \approx \partial^2 \rho / \partial m_l^2$$

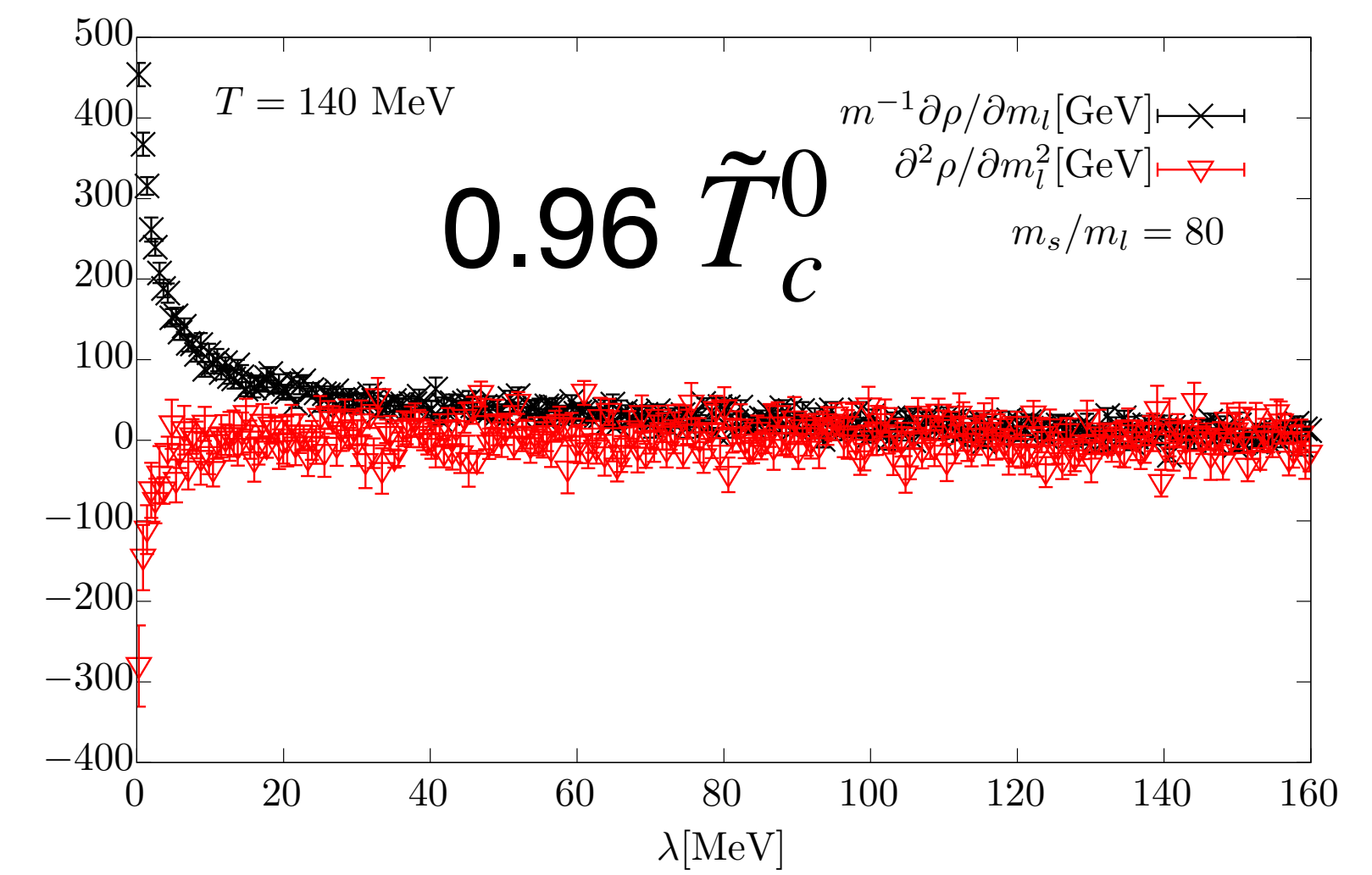
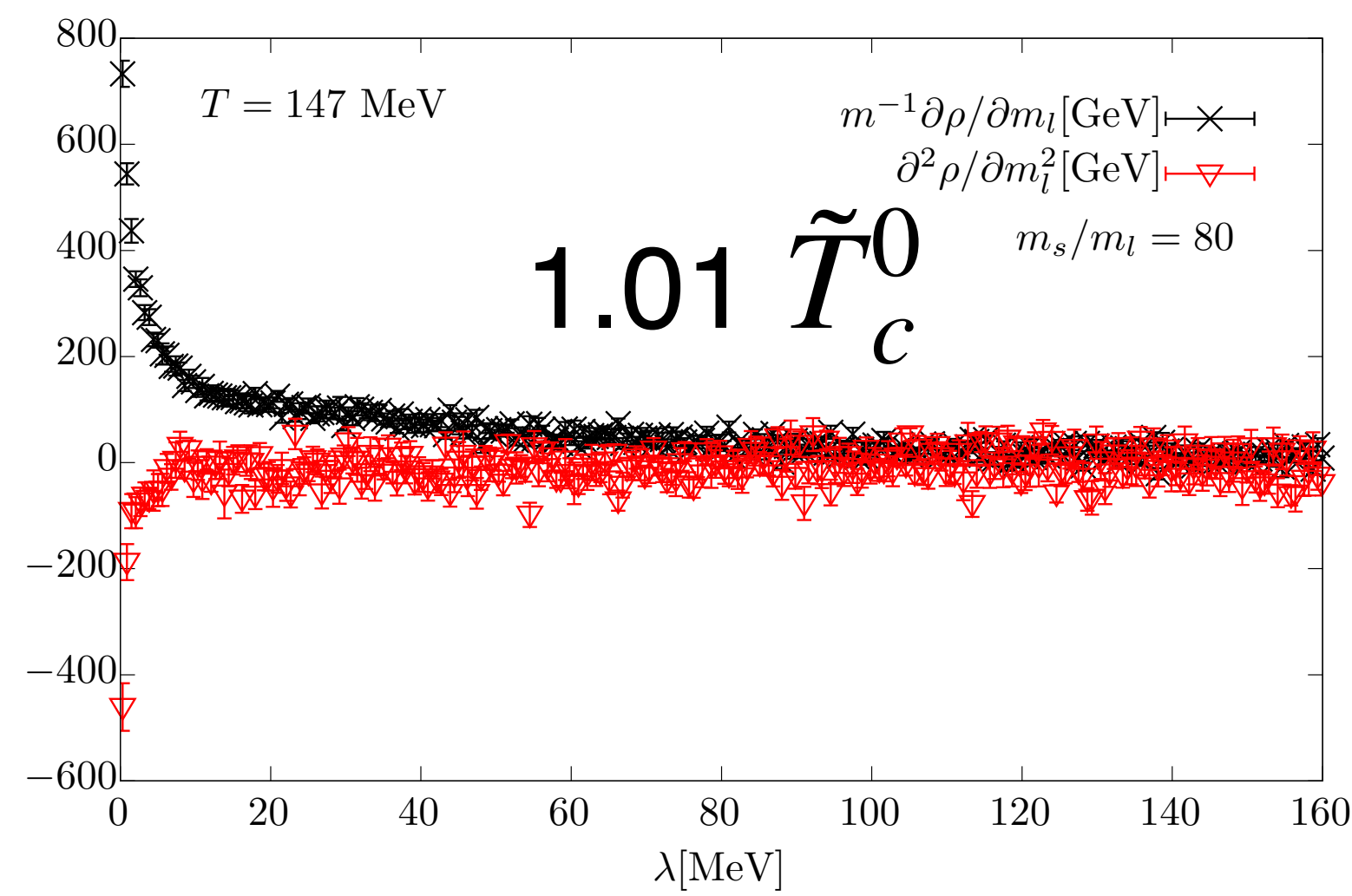
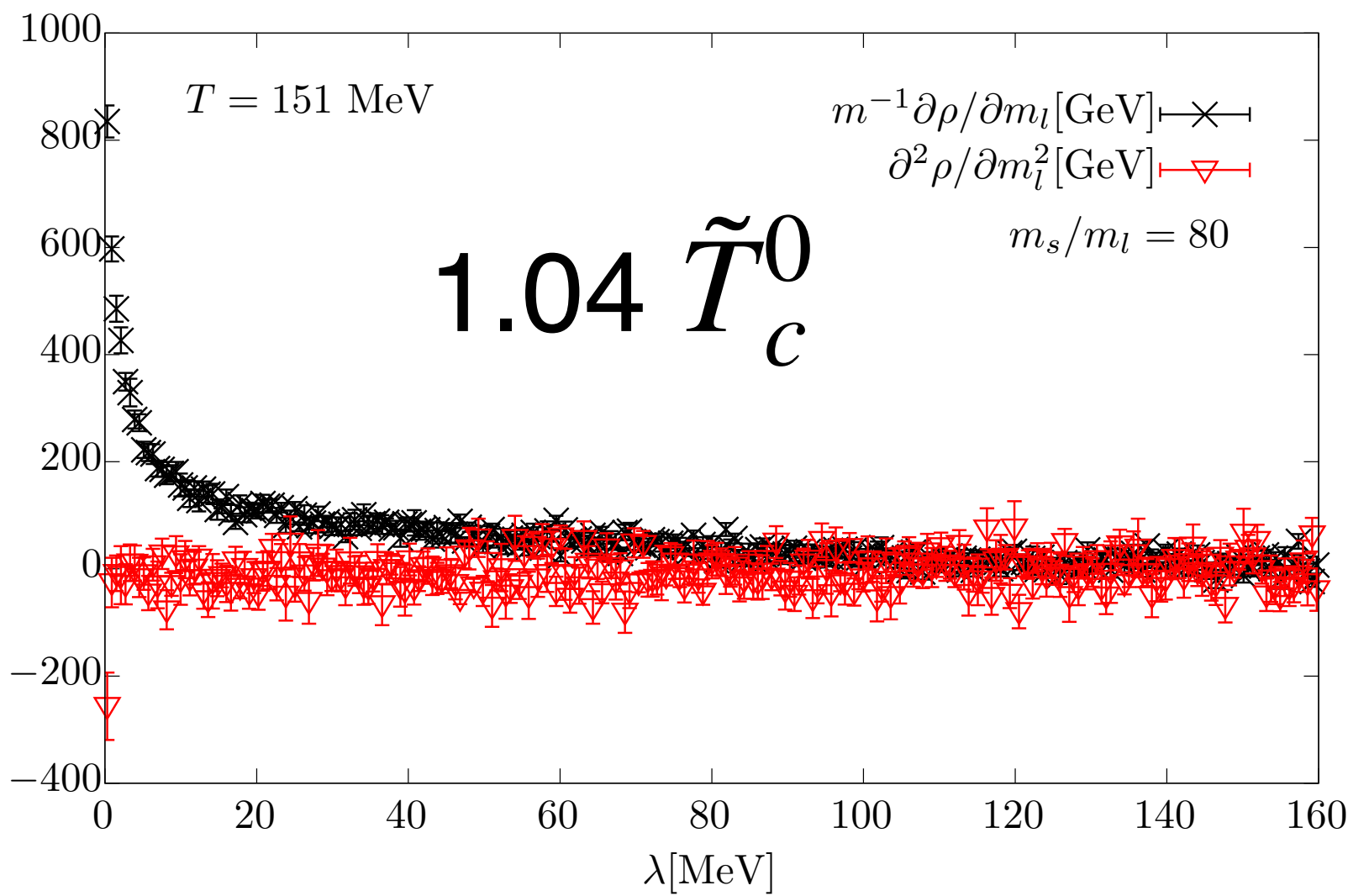
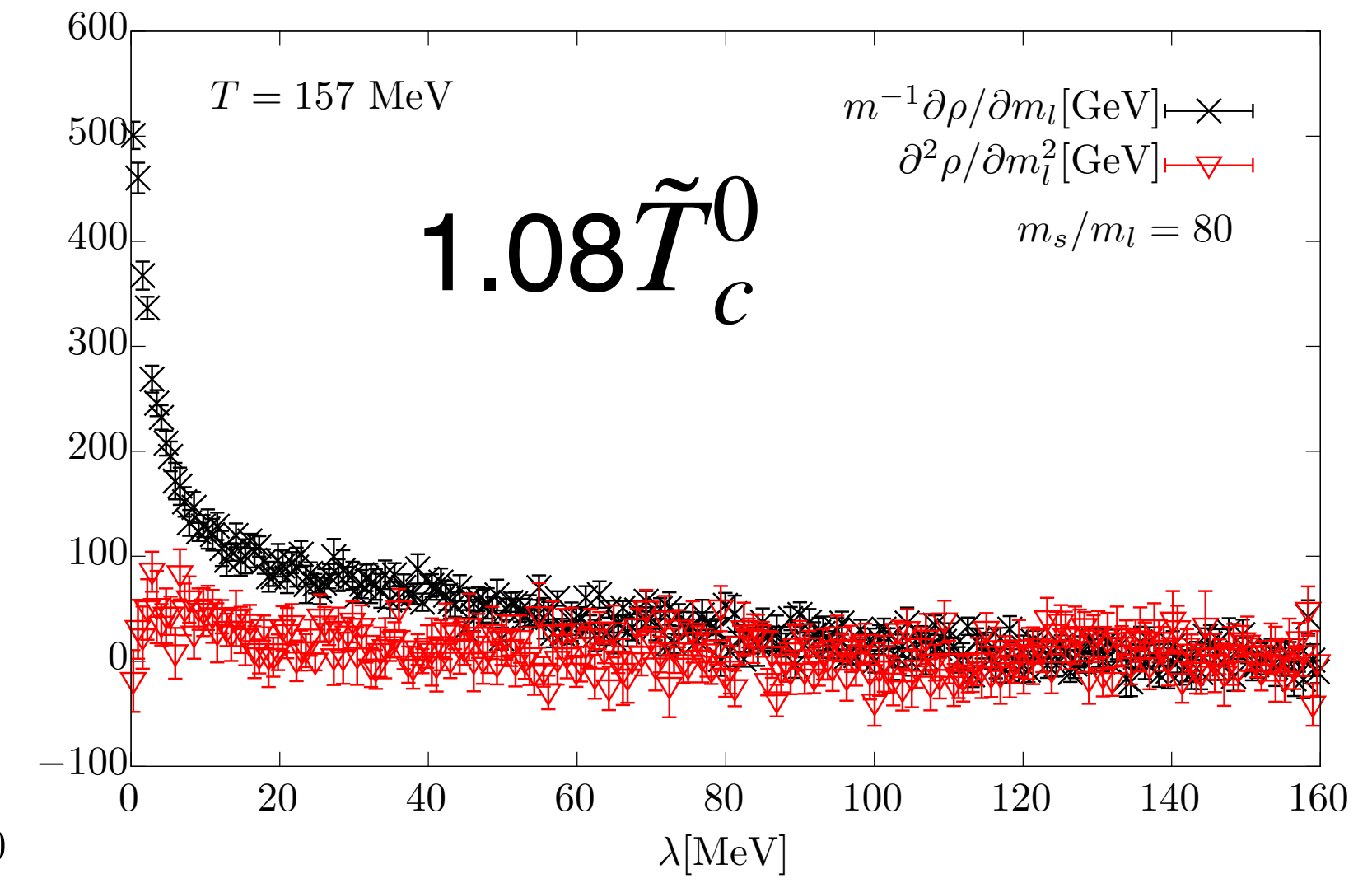
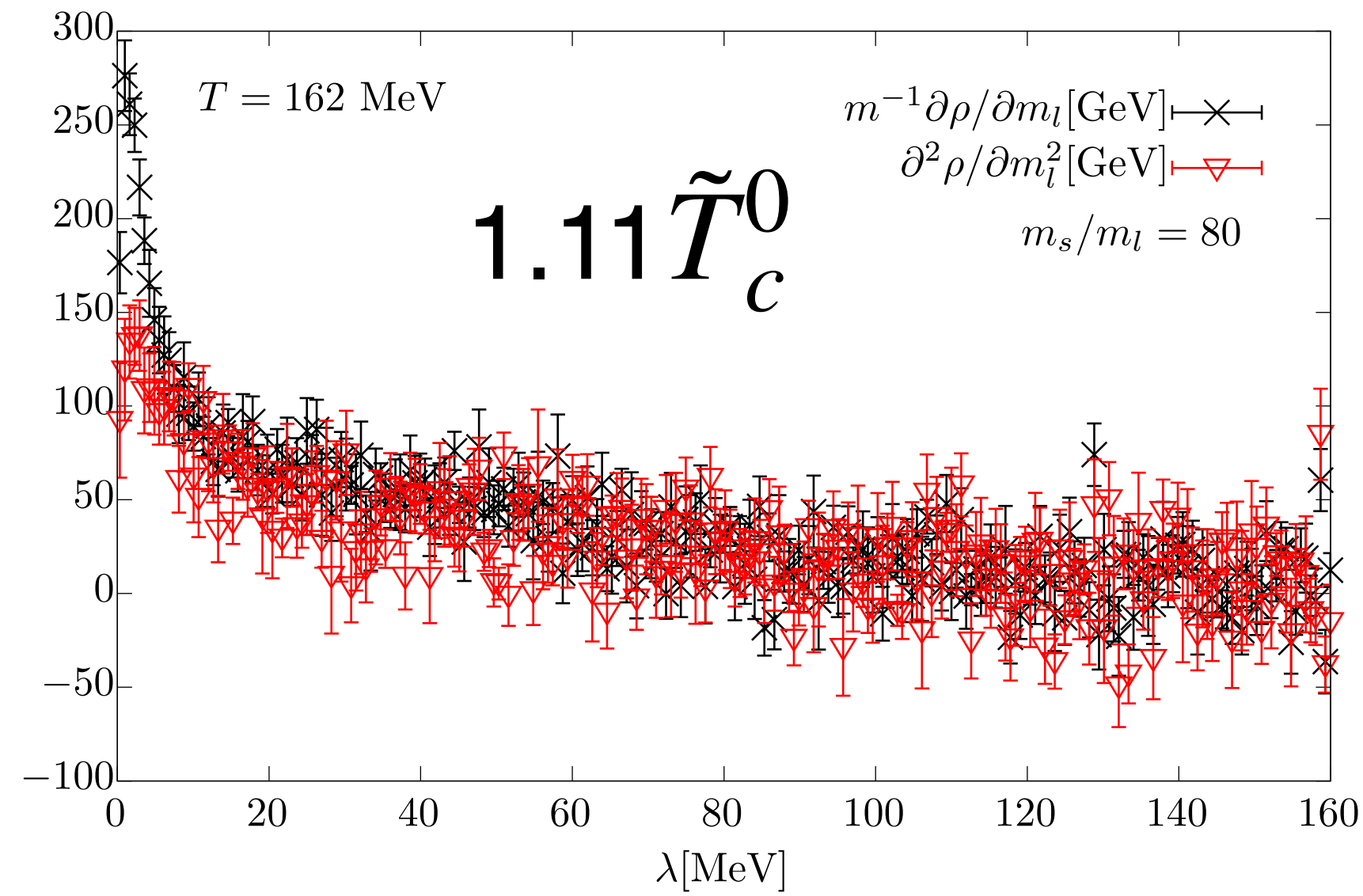
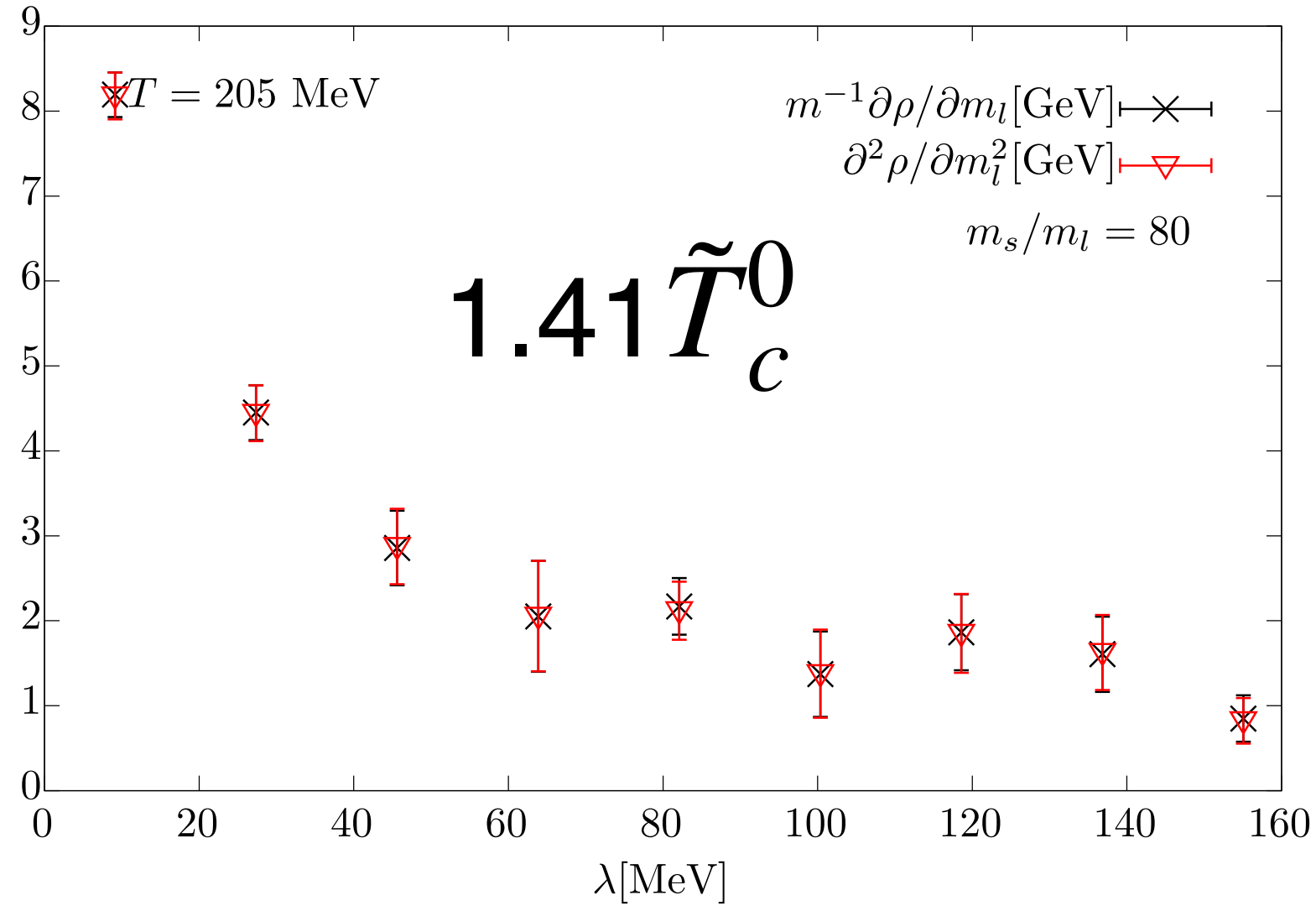
$$\partial^3 \rho / \partial m_l^3 \approx 0$$

Dilute instanton gas approximation at $1.6 T_c$: $\rho(\lambda \rightarrow 0, m_l \rightarrow 0) \propto m_l^2 \delta(\lambda)$

Dilute instanton gas approximation does not hold towards T_c

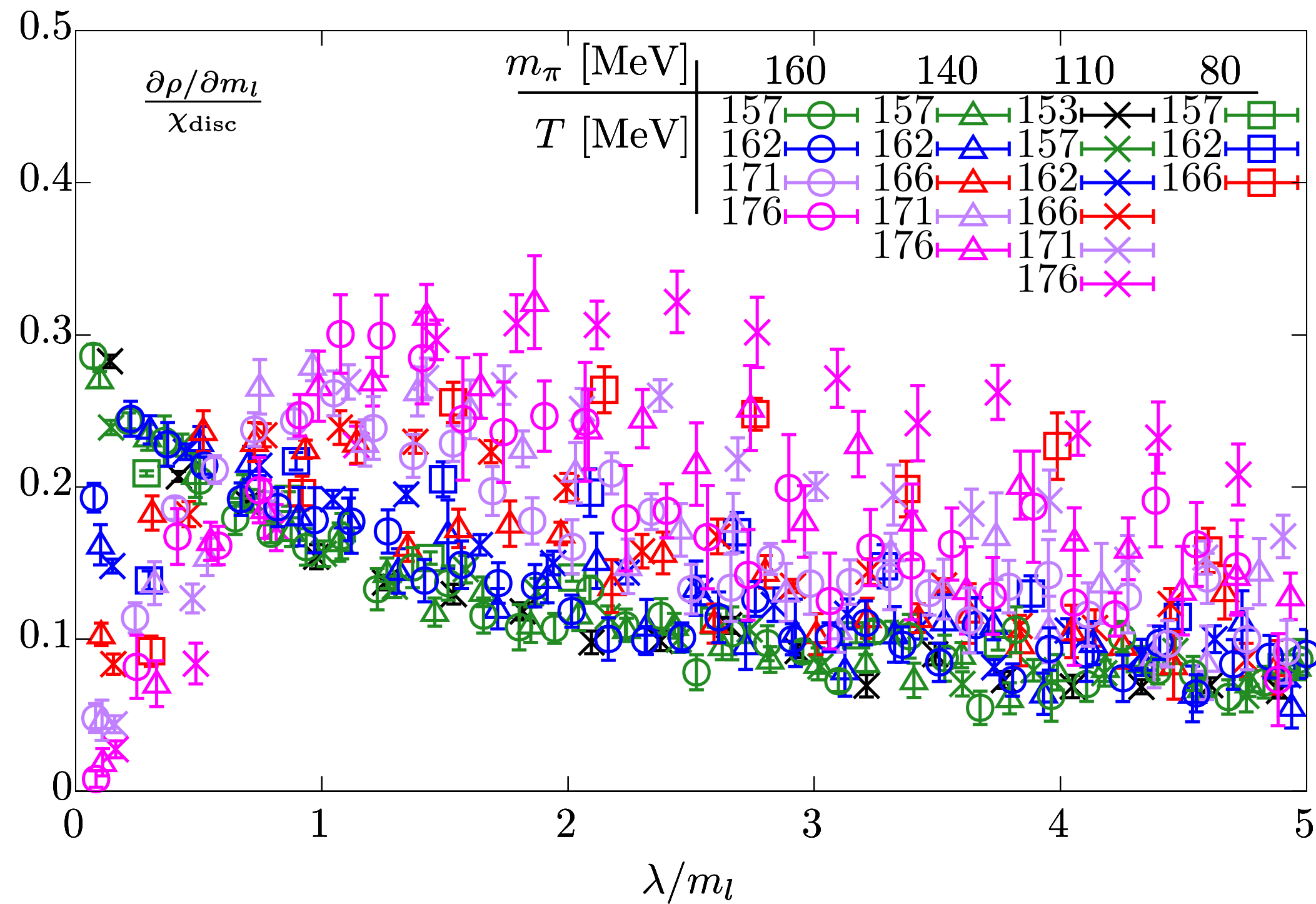
$$\partial^2 \rho / \partial m_l \neq m_l^{-1} \partial \rho / \partial m$$

$$\tilde{T}_c^0 \equiv T_c^0(N_\tau = 8) \approx 145.6 \text{ MeV}$$

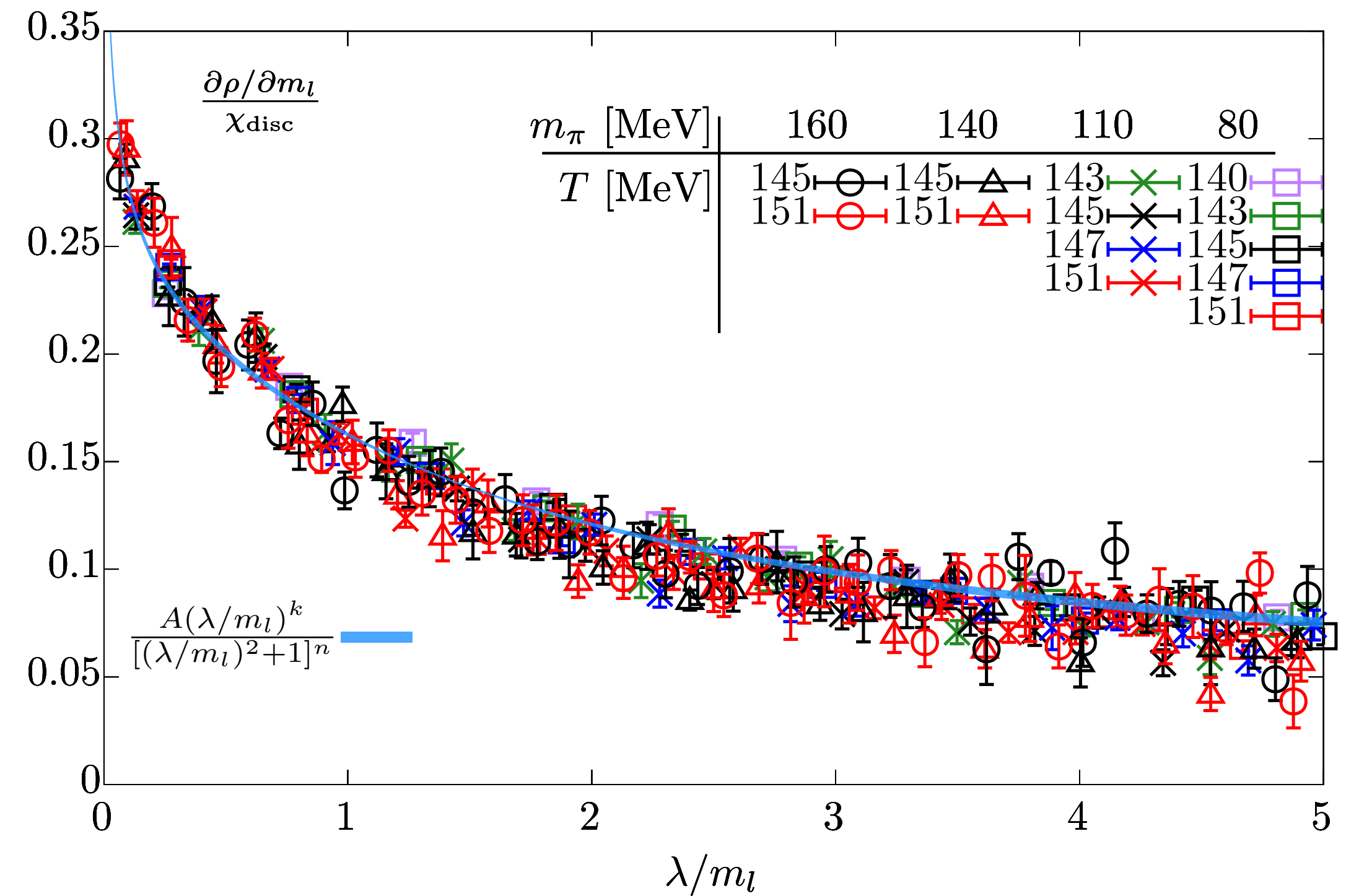


Scaling behavior in $\partial\rho/\partial m_l/\chi_{disc}$ v.s. λ/m_l near T_c

No scaling at $T \in [1.08, 1.2]\tilde{T}_c^0$



Scaling at $T \in [0.96, 1.04]\tilde{T}_c^0$

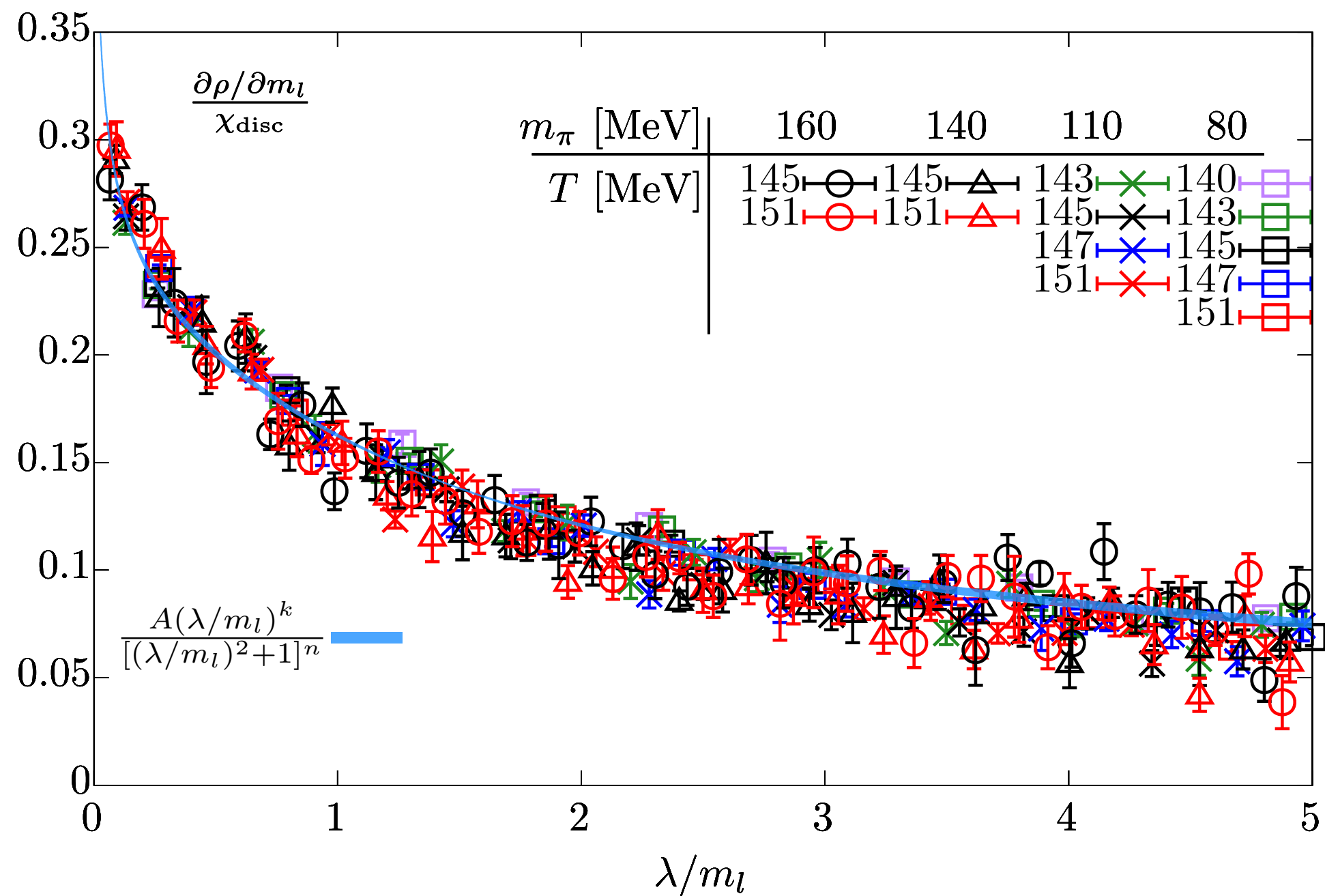


Wei-Ping Huang et al., work in progress,
HTD, Wei-Ping Huang, Min Lin et al., PoS LATTICE2021 (2022) 591

$$\partial\rho/\partial m_l = f(\lambda/m_l) \times \chi_{disc}$$

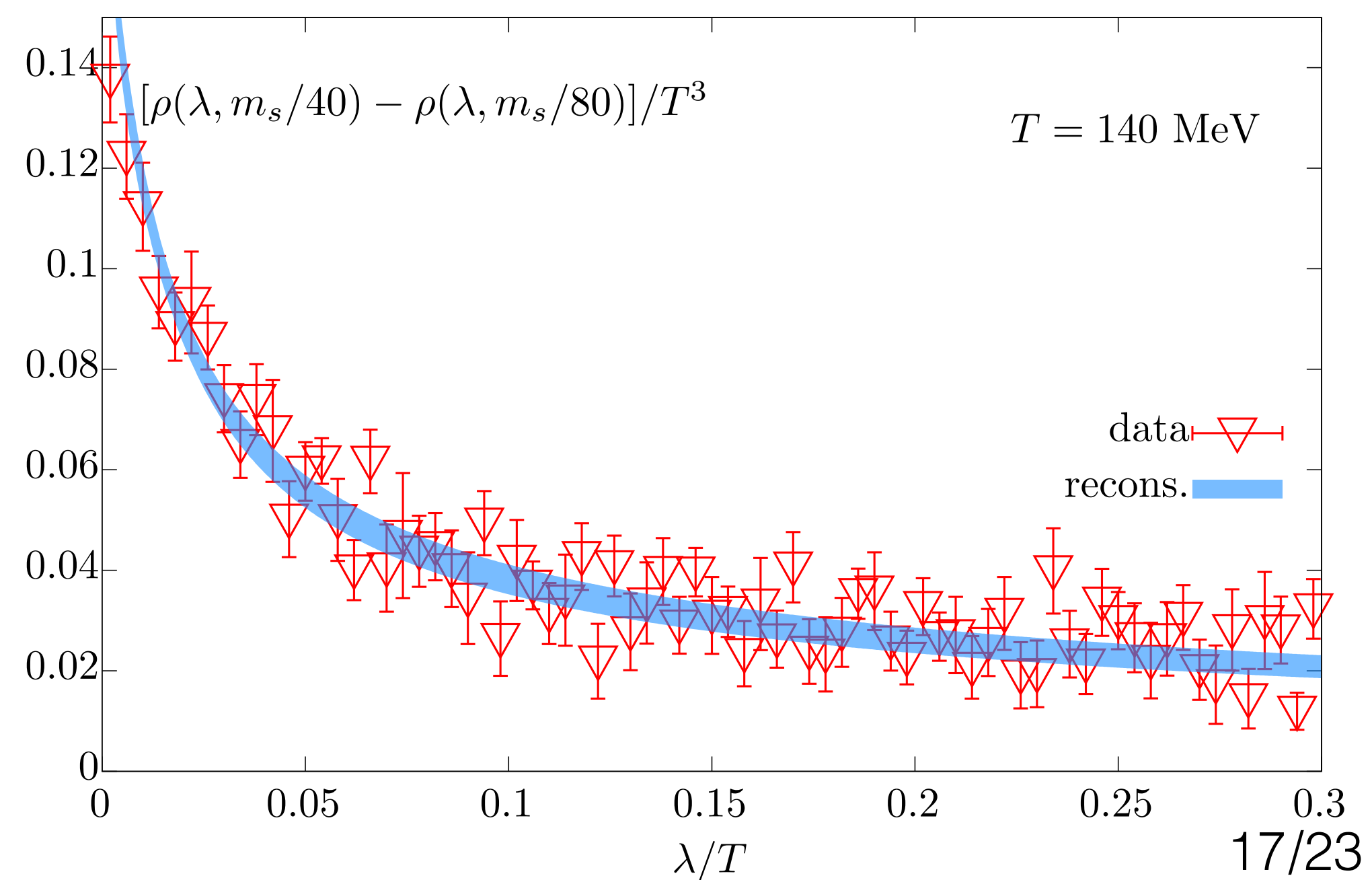
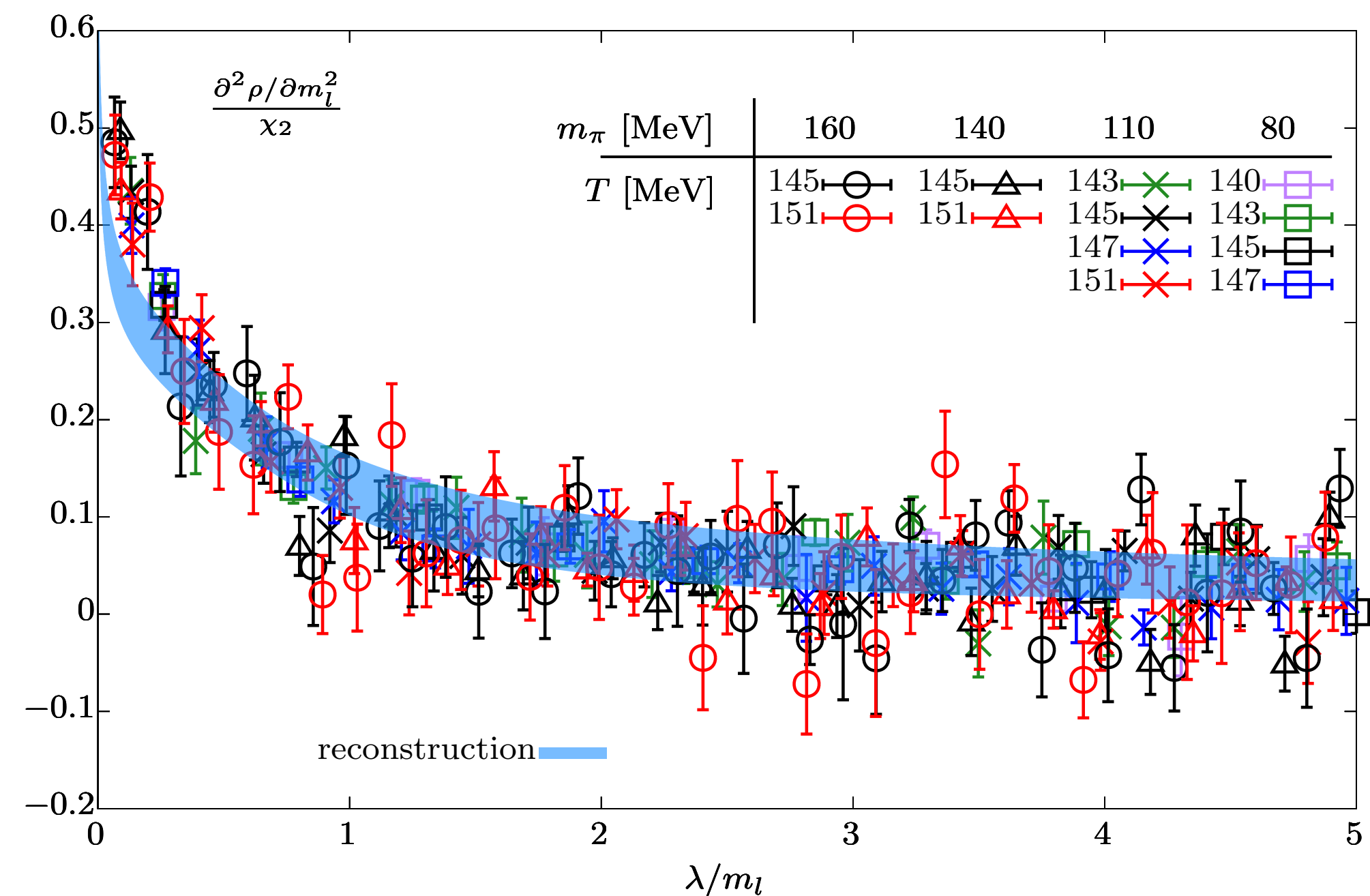
Criticality of QCD in correlated Dirac Eigenvalues

$$\partial\rho/\partial m_l = f(\lambda/m_l) \times \chi_{disc}$$



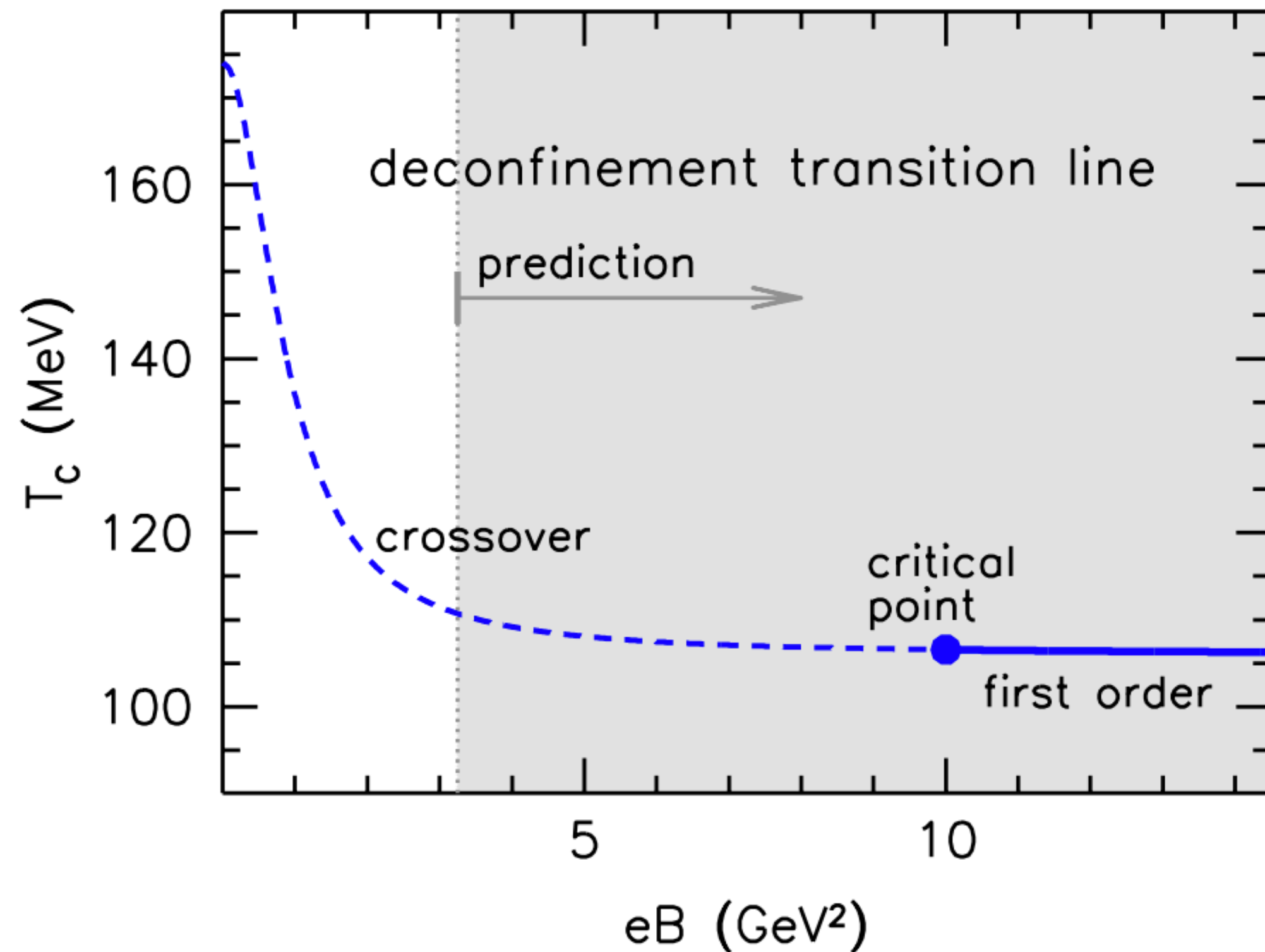
Wei-Ping Huang et al., work in progress

Reproduction of
 $\partial^2\rho/\partial m_l^2/\chi_2$ & ρ
without fits!



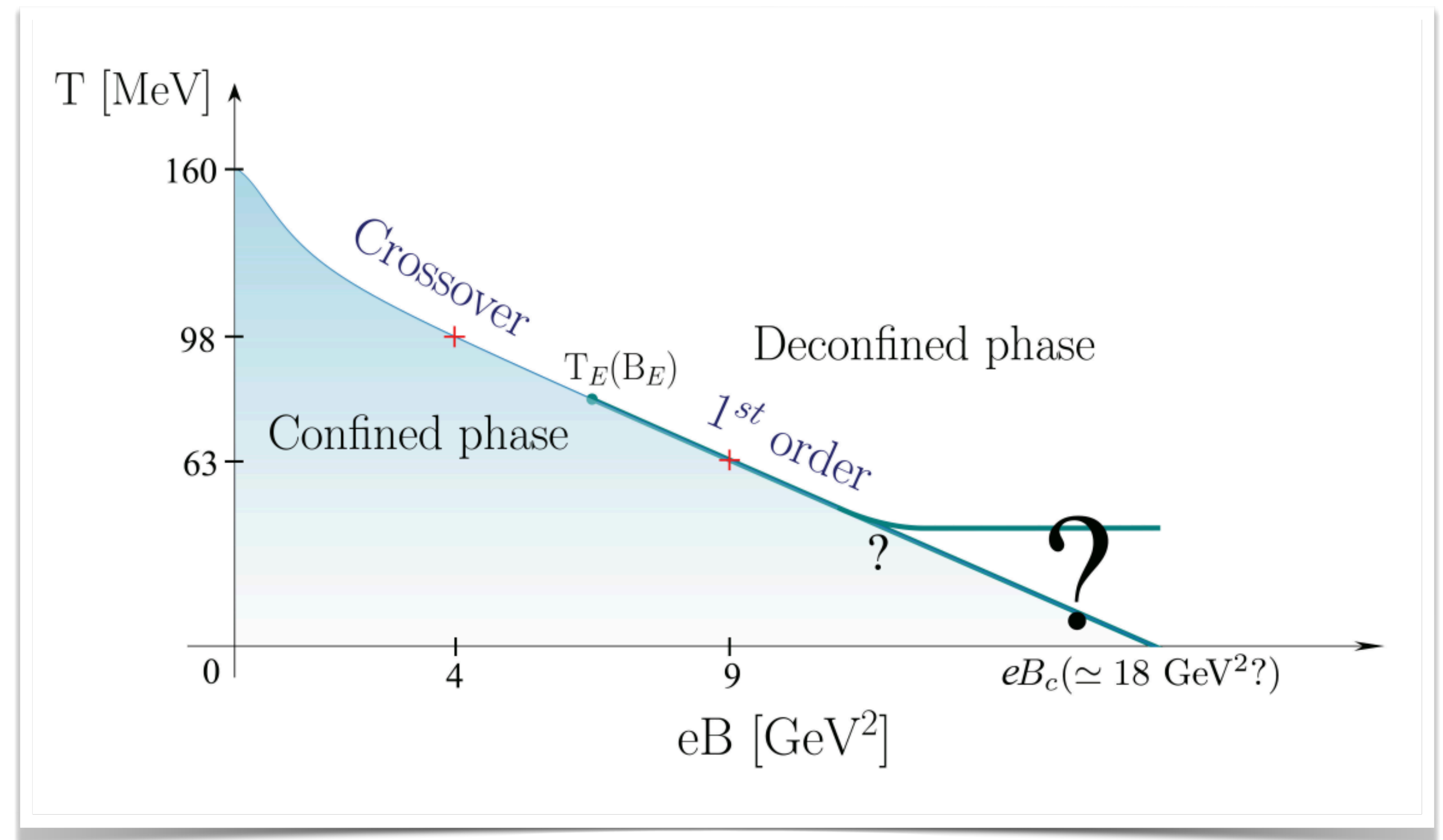
A plausible CEP in T-eB plane

Prediction of
a Critical End Point in the T-eB plane



G. Endrodi, JHEP 07 (2015) 173

Observation of a 1st order phase transition at $eB=9 \text{ GeV}^2$



M. D'Elia et al., Phys.Rev.D 105 (2022) 3, 034511

QCD transition in strong magnetic fields

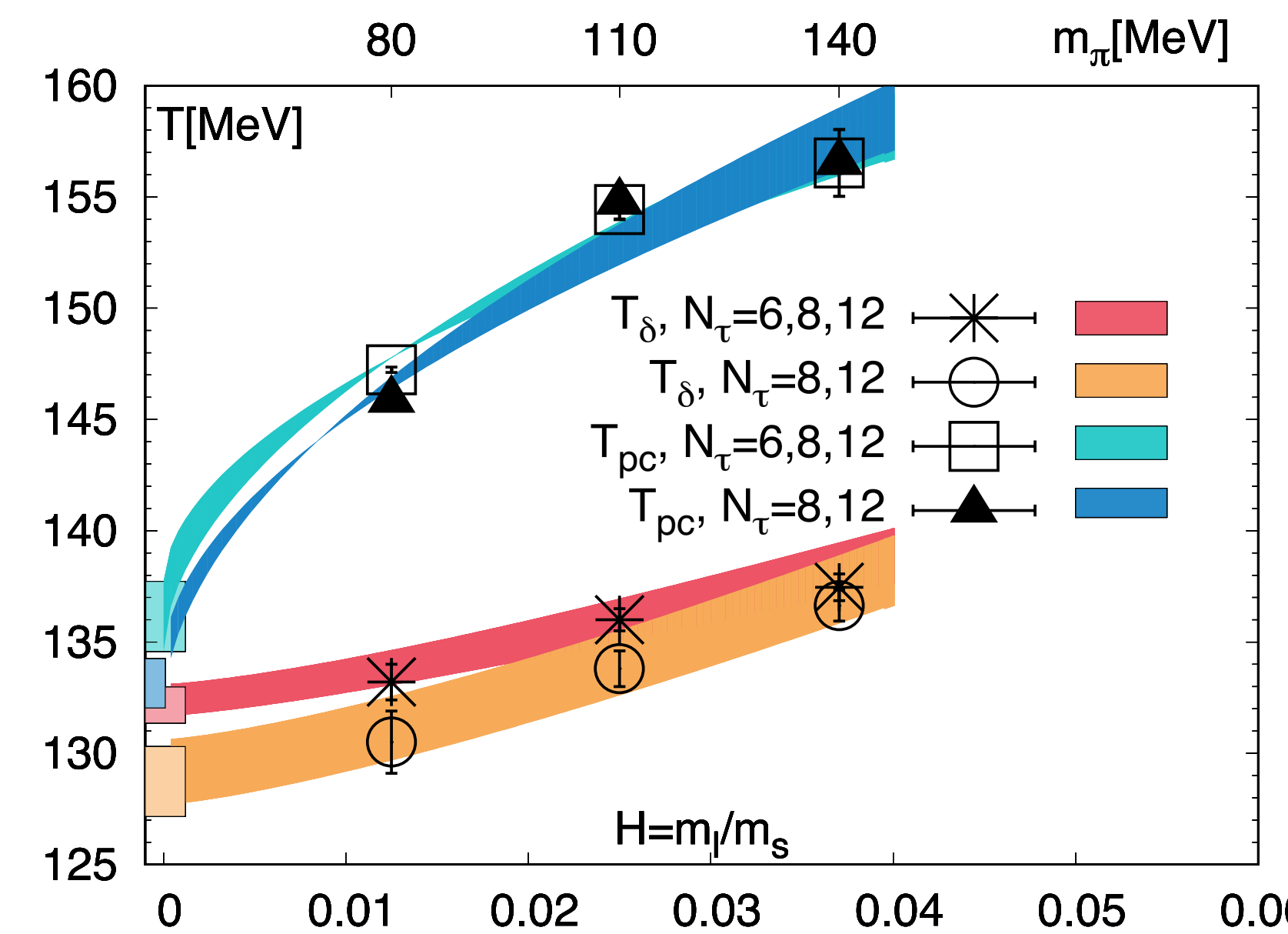
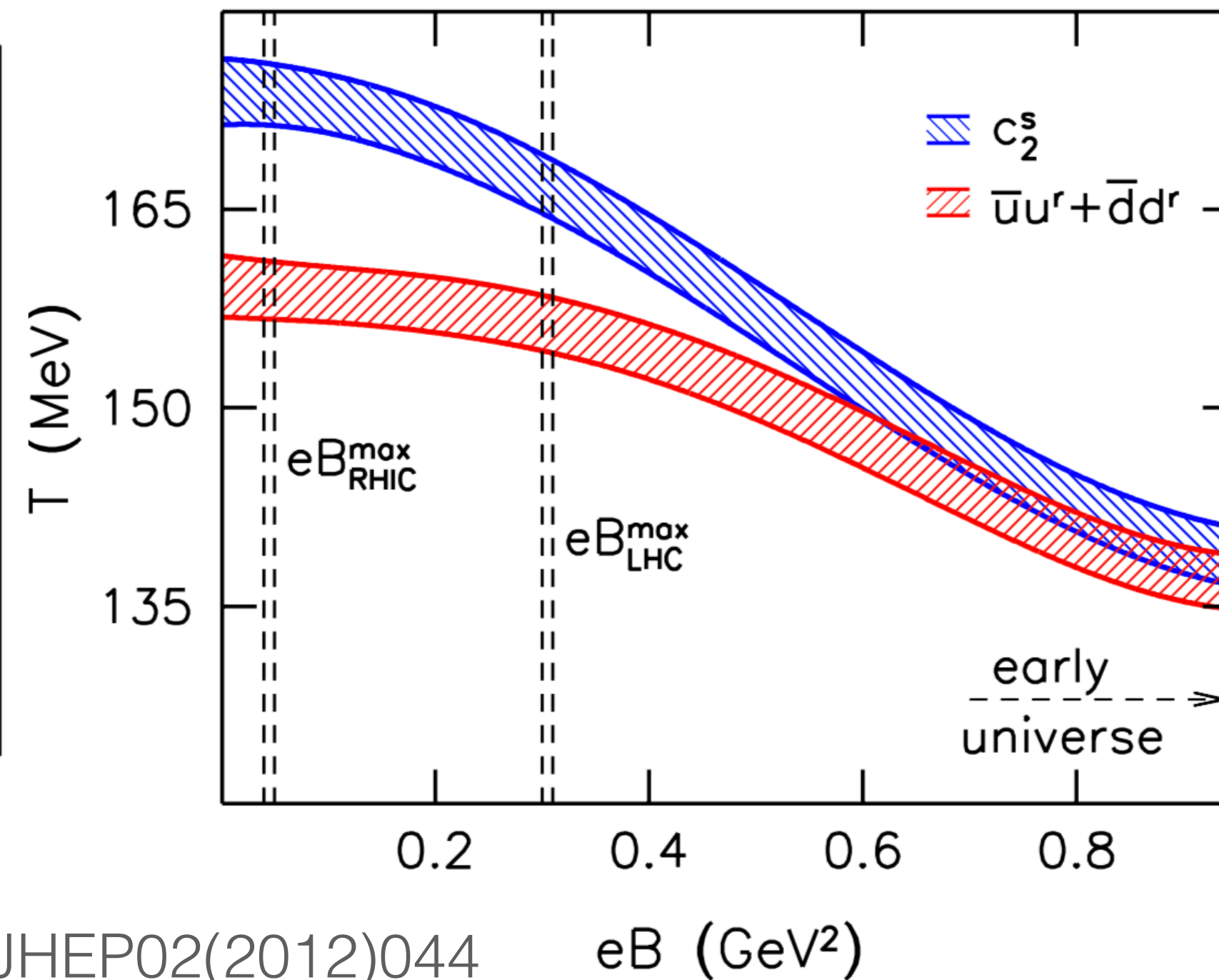
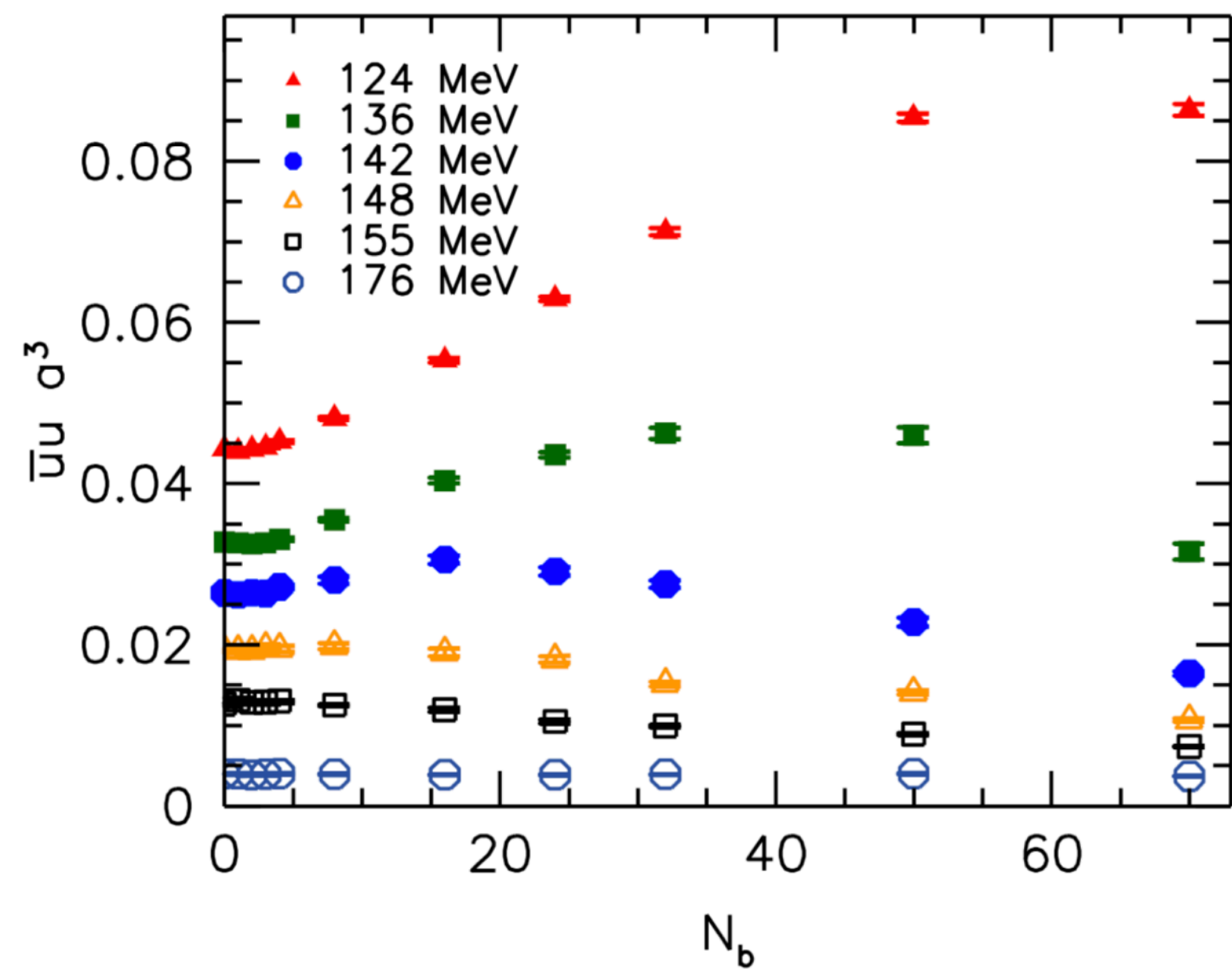
Inverse magnetic catalyses and reduction of T_{pc}

Continuum extrapolated lattice QCD results with physical pion mass

Inverse magnetic catalysis (IMC)

$eB \uparrow \quad T_{pc} \downarrow$

$eB=0, N_f=2+1$ QCD



Bali et al., JHEP02(2012)044

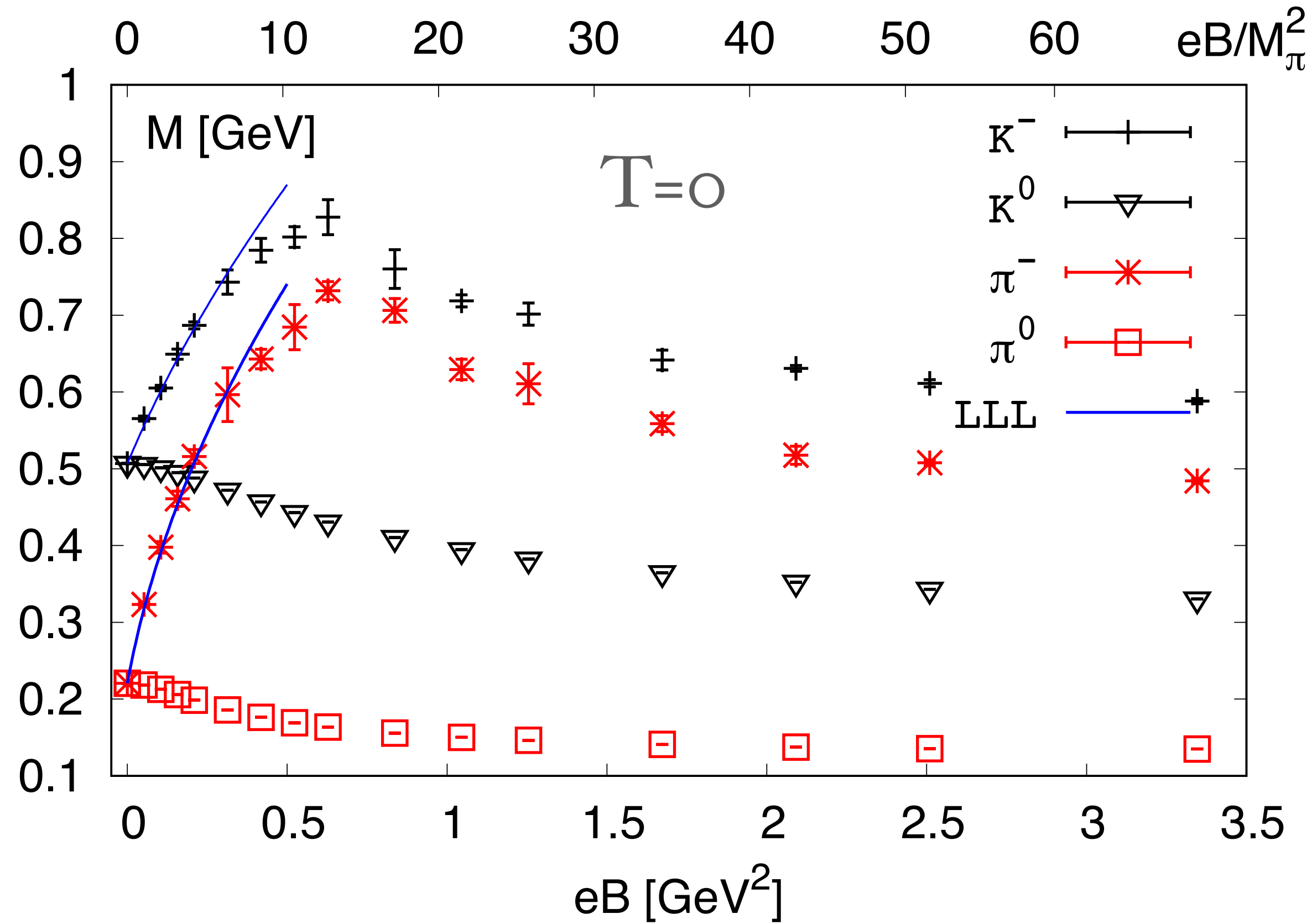
HTD, P. Hegde, O. Kaczmarek et al. [HotQCD],
Phys. Rev. Lett. 123 (2019) 062002
HTD, arXiv:2002.11957

Reduction of T_{pc} always associated with IMC? Not necessarily!

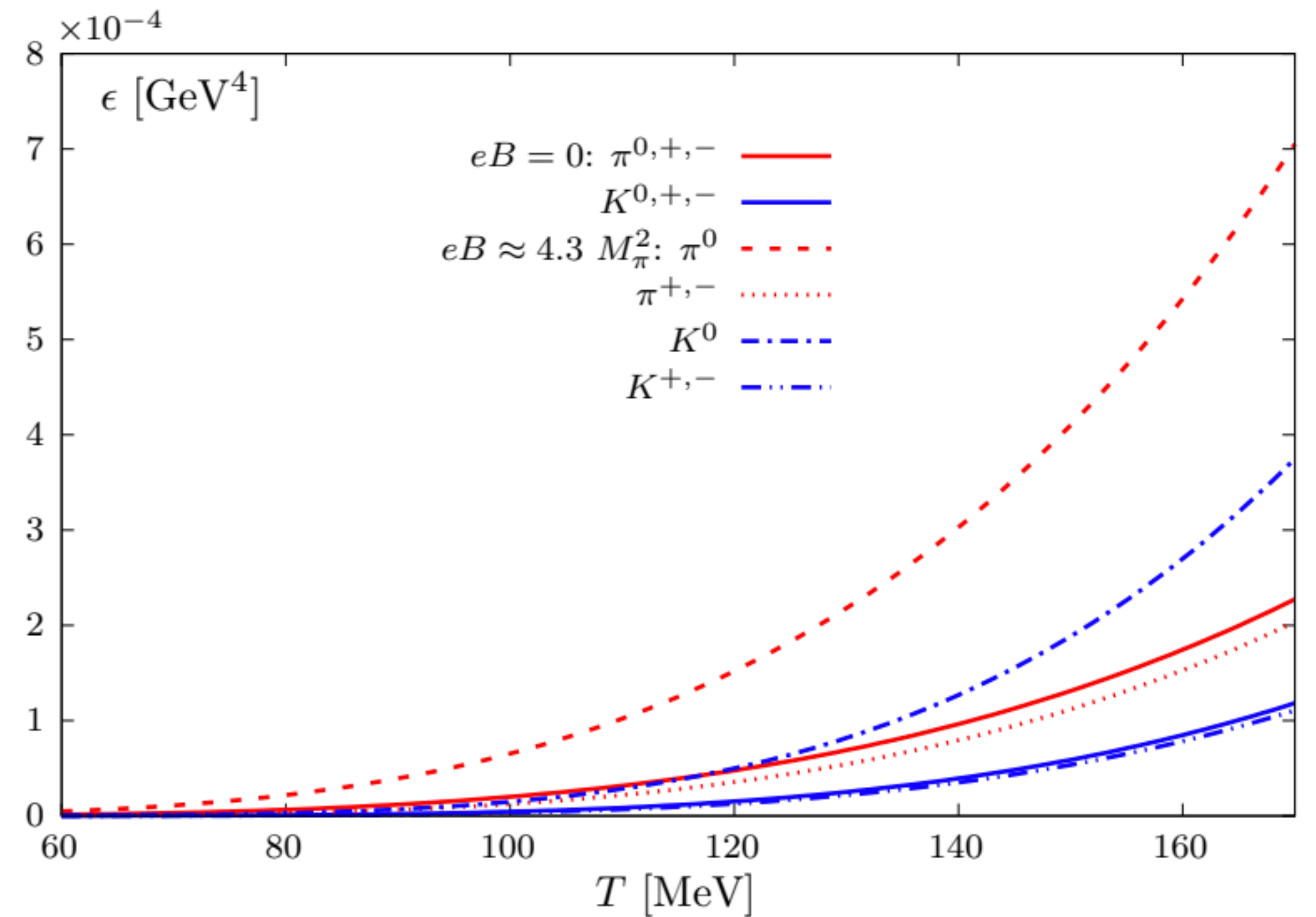
Role of hadrons?

Masses of $\pi^{0,\pm}$ and $K^{0,\pm}$ and energy density

$N_f=2+1$ QCD, $M_\pi(eB=0) \approx 220$ MeV,
 $32^3 \times 96$ lattices with $a^{-1} \approx 1.7$ GeV and HISQ action



Energy density in Hadron resonance gas model



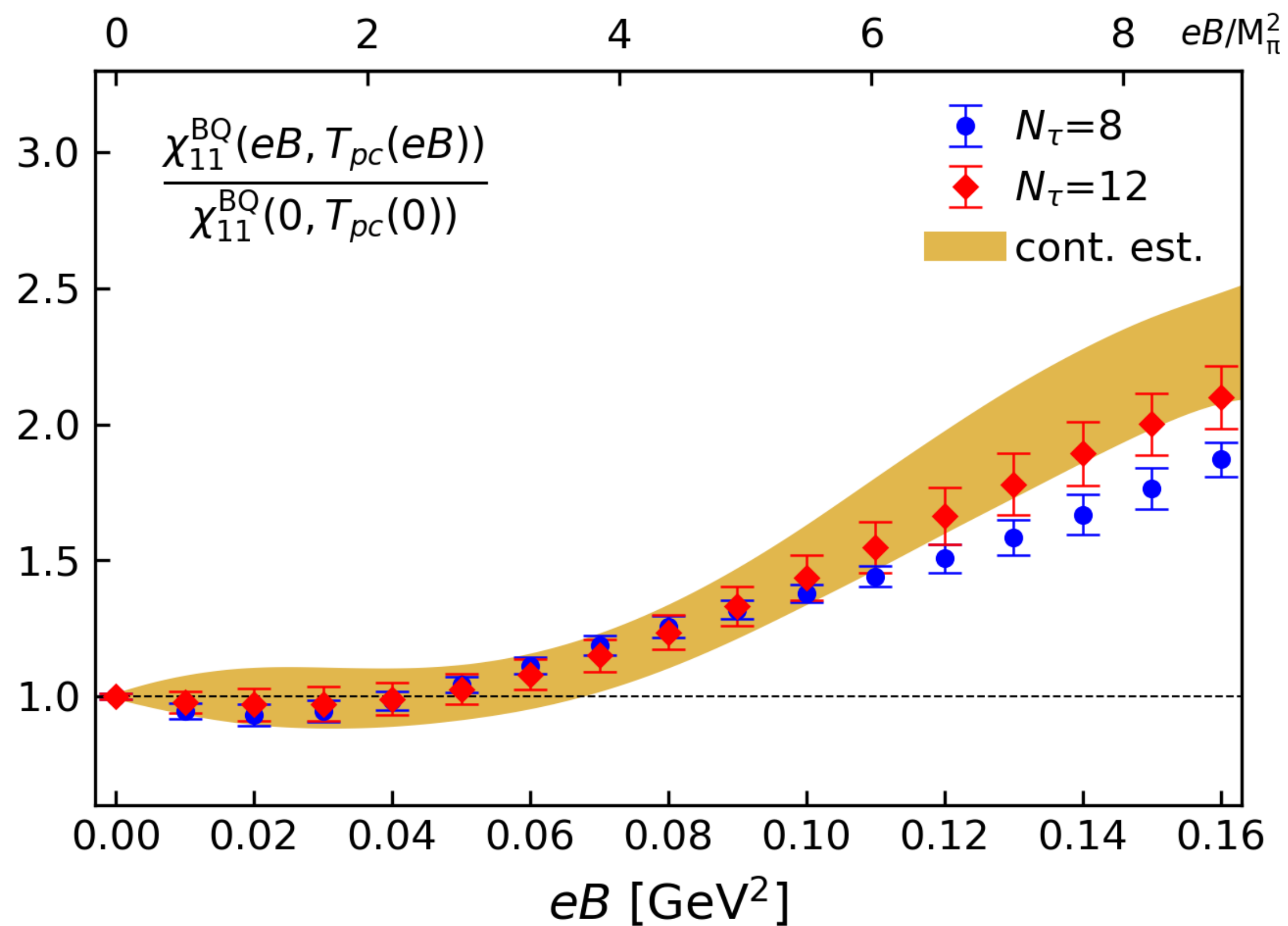
HTD, S.-T. Li, A. Tomiya, X.-D. Wang, Y. Zhang, PRD 126 (2021) 082001

See quenched LQCD results in Bali et al., PRD 97 (2018) 034505,
 Lushevskaya et al., NPB 898 (2015) 627

HTD, S.-T. Li, Q. Shi, A. Tomiya, X.-D. Wang, Y. Zhang, arXiv: 2011.04870

Ratio $X(eB)/X(eB=0)$ for 2nd order off-diagonal fluctuations

$N_f=2+1$ QCD, $M_\pi(eB=0) \approx 135$ MeV, $T_{pc}(eB=0) \approx 157$ MeV, $32^3 \times 8$ and $48^3 \times 12$ lattices with HISQ action



$X(eB)/X(eB=0)$: Rcp like observable

At $eB \simeq M_\pi^2$: deviation from unity is mild

At $eB \simeq 8M_\pi^2$: ~ 2 !

Note: $T_{pc}(eB \simeq 10M_\pi^2)/T_{pc}(eB=0) \sim 99\%$



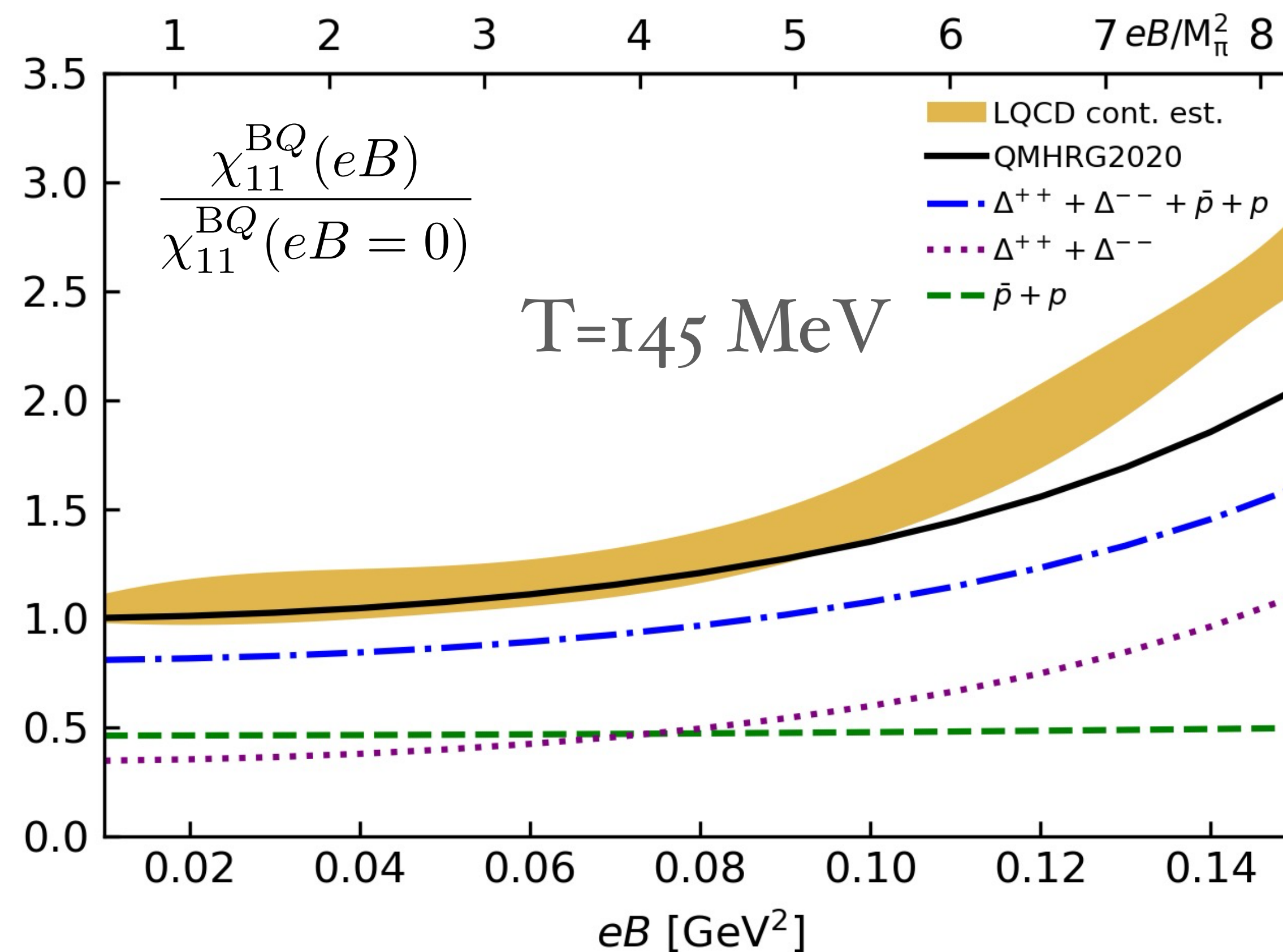
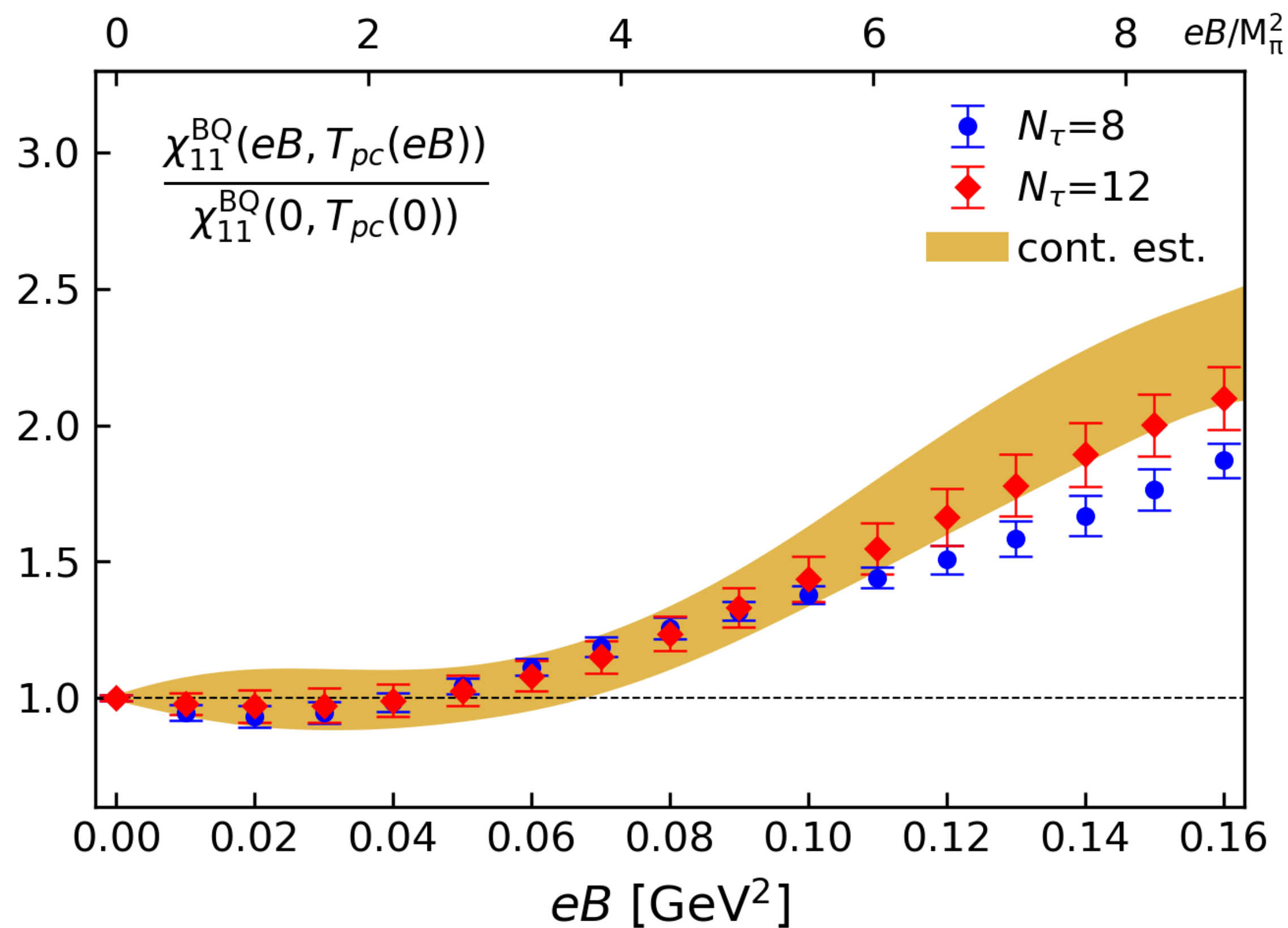
 Central Collisions Peripheral Collisions

 Smaller eB Larger eB

HTD, S.-T. Li, J.-H. Liu and X.-D. Wang, QM2022, arXiv:2208.07285

Ratio $\chi(eB)/\chi(eB=0)$ for 2nd order off-diagonal fluctuations

$N_f=2+1$ QCD, $M_\pi(eB=0) \approx 135$ MeV, $T_{pc}(eB=0) \approx 157$ MeV, $32^3 \times 8$ and $48^3 \times 12$ lattices with HISQ action



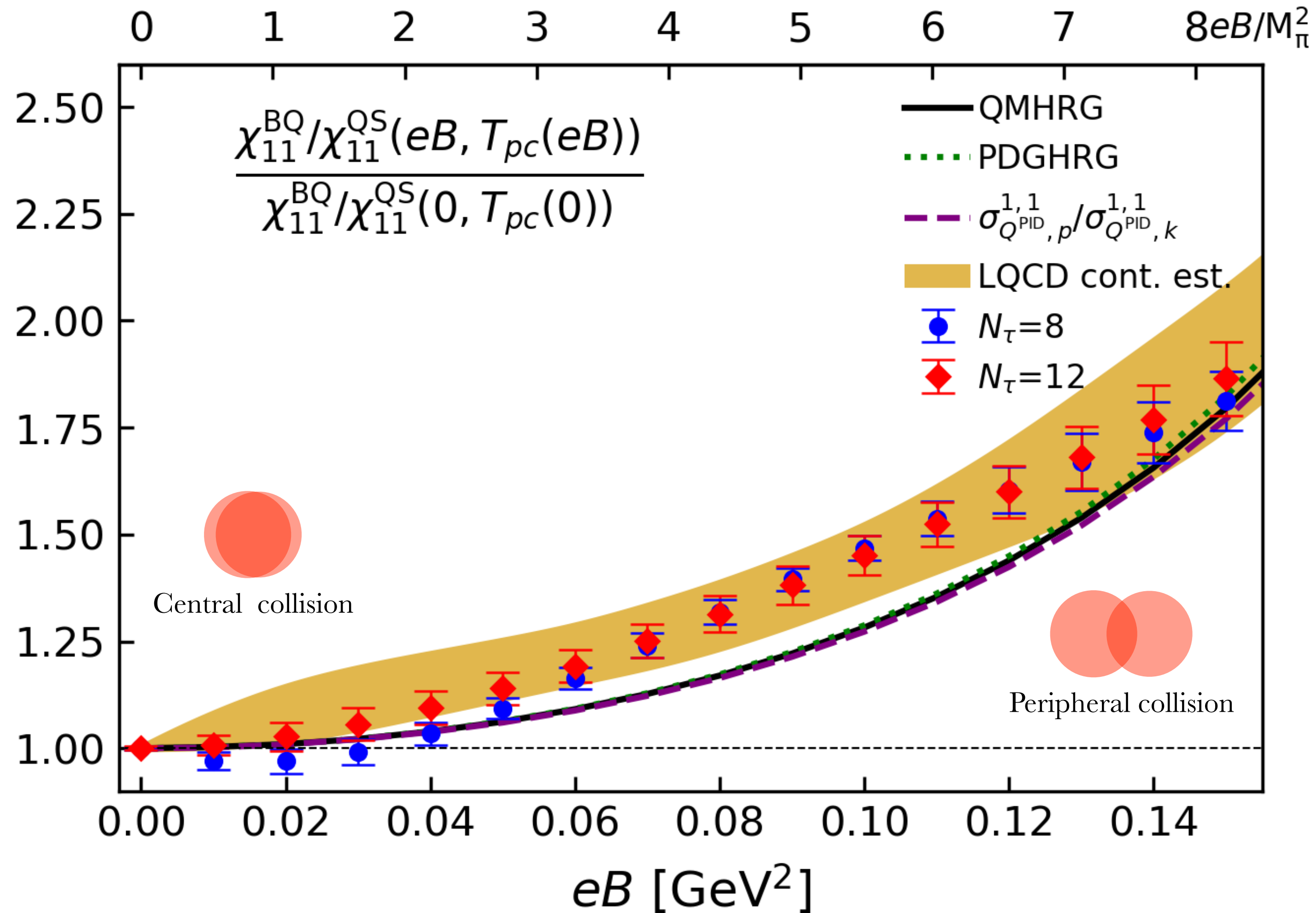
Central Collisions
 Smaller eB

Peripheral Collisions
 Larger eB

HTD, S.-T. Li, J.-H. Liu and X.-D. Wang, QM2022, arXiv:2208.07285

QCD benchmarks for the manifestation of eB in conserved charge fluctuations

See Jun-Hong Liu's talk on Fri.



Summary & Outlook

📌 Most relevant criticality at $\mu_B = 0$ to thermodynamics at CEP:

2nd $O(4)$ phase transition $\implies T_c^{CEP} < T_c^0 \approx 132\text{MeV}$

📌 Criticality in Dirac eigenvalue correlation: Microscopic manifestation of the criticality in a narrow T window

📌 QCD benchmarks for 2nd off-diagonal fluctuations in a background magnetic field: possibility to detect the existence of a magnetic field in HIC

▶ Search for criticality in the T - eB plane