

Probing gluon PDF at $x \rightarrow 0$ with Ultrapерipheral PbPb collisions at 5.02 TeV in CMS

André Ståhl

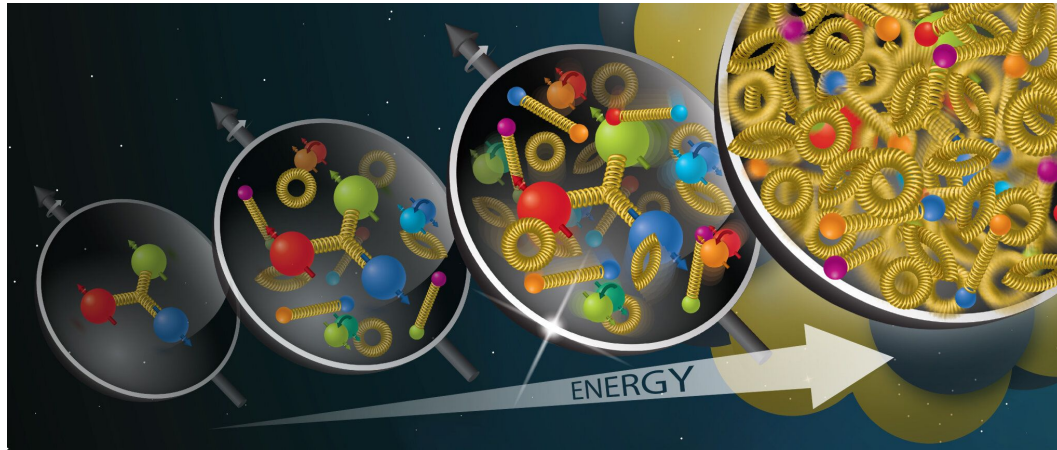
on behalf of the CMS Collaboration

European Organisation for Nuclear Research

LPCC CERN Seminar

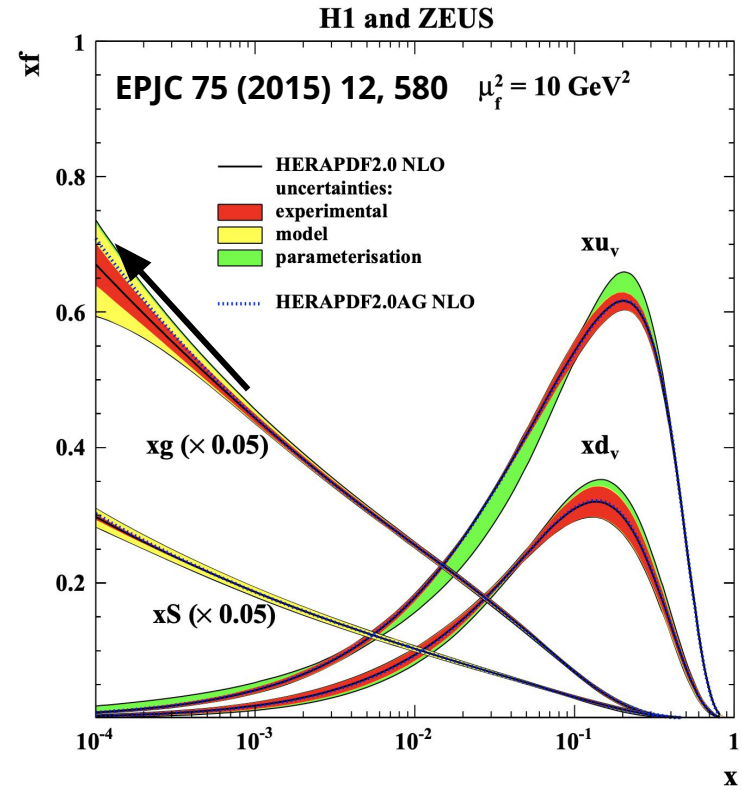


Understand the Glue that binds us all

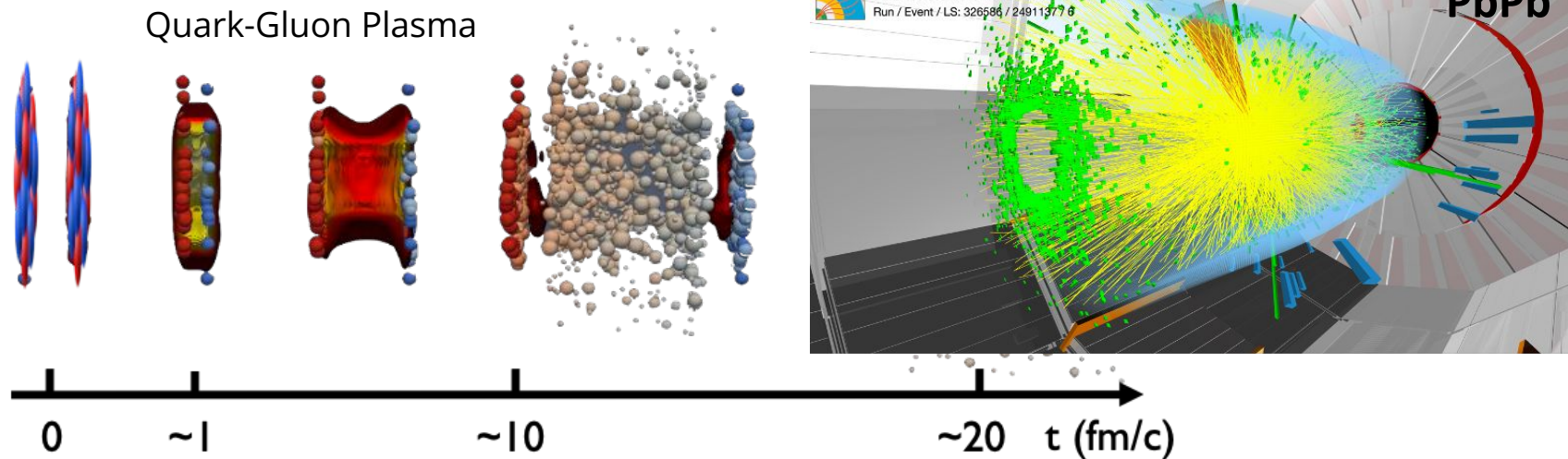


- DIS results show a rapid increase of gluon density towards small x , by gluon splittings.
- Indefinite growth at small- x violates unitarity \rightarrow mechanism beyond gluon splitting expected.

What is the fate of gluons at extreme densities toward the unitary limit?



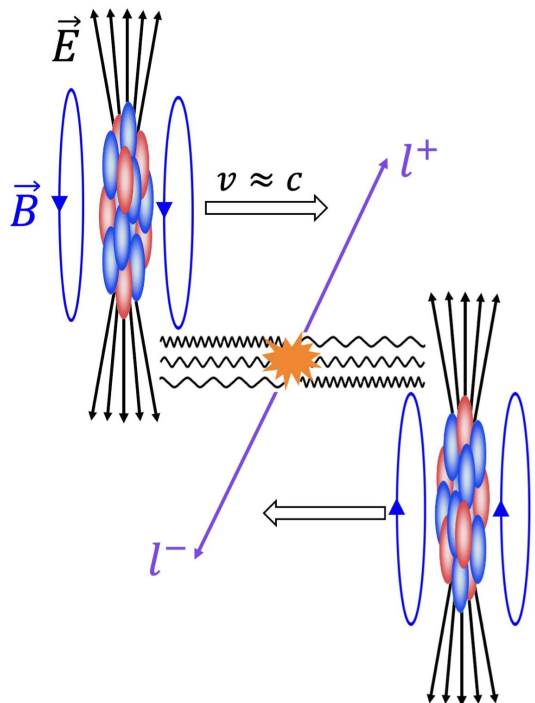
Heavy ion collisions and Hot Quark-Gluon Matter



A comprehensive program of investigating novel properties of quark-gluon plasma created at highest temperatures with the CMS experiment, using a multitude of probes in large (PbPb) and small (pp, pPb) systems

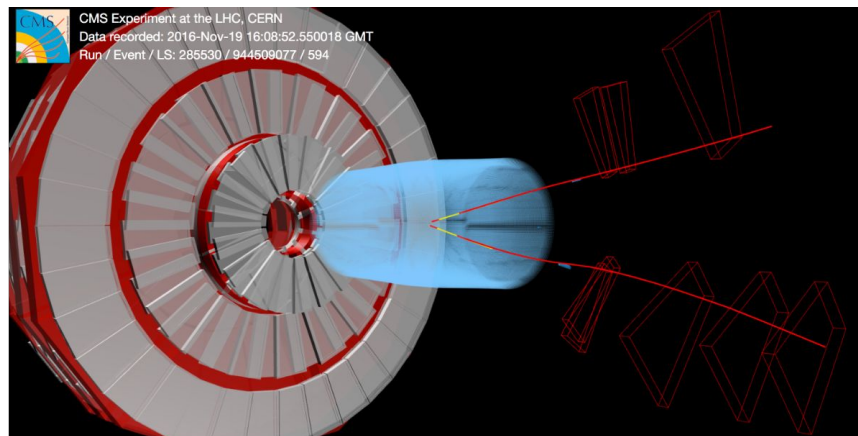
Ultra-peripheral nuclear collisions

When two ions “miss” each other, no QGP is created but,



- Strong EM fields generated by relativistic ions ($B \sim 10^{16}$ T).
- Lorentz contracted EM fields \rightarrow flux of quasi-real γ ($Q^2 < \hbar^2/R^2$). The photon flux $\propto Z^2$.
- Photon kinematics: $p_T < \hbar/R_A \sim 30$ MeV ($E_{\max} \sim 80$ GeV) at LHC.

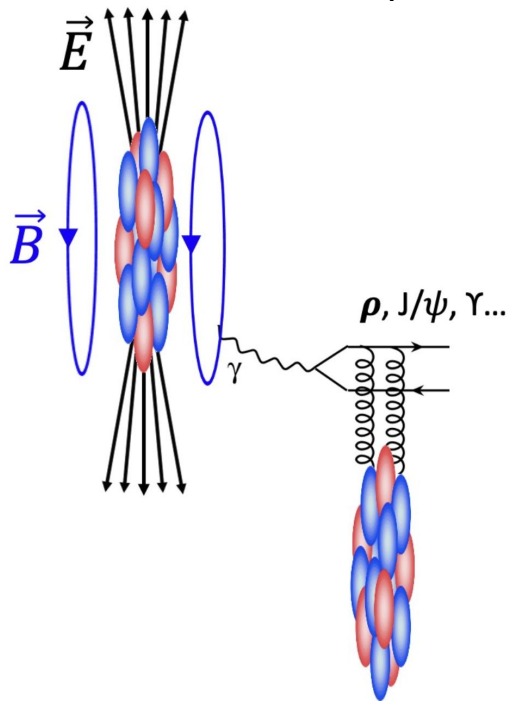
- light-light and light-Nucleus collider
- BSM searches (ALP, $g_{\tau-2}$).



Ultra-peripheral nuclear collisions: photon-nuclear interactions

Vector meson (e.g., J/ψ) photoproduction directly probes gluonic structure of nucleus and nucleon.

At LO in pQCD, cross section \sim photon flux \otimes $[xG(x)]^2$ (gluon PDFs)



Coherent production:

- Photon ($\hbar/k_L > 2R$) couples coherently to whole nucleus.
- Vector Meson (VM) $\langle p_T \rangle \sim 50$ MeV.
- Target nucleus usually remains intact.

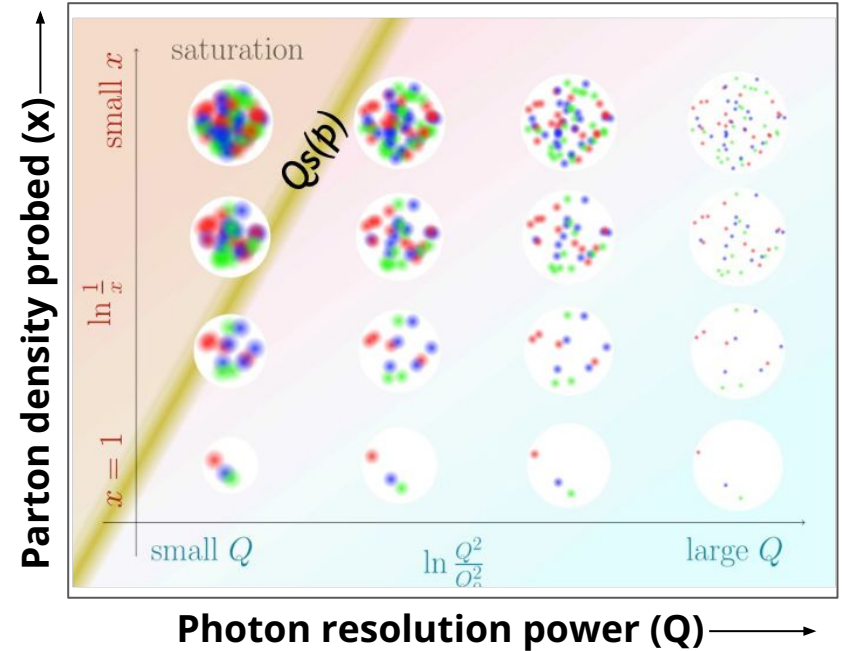
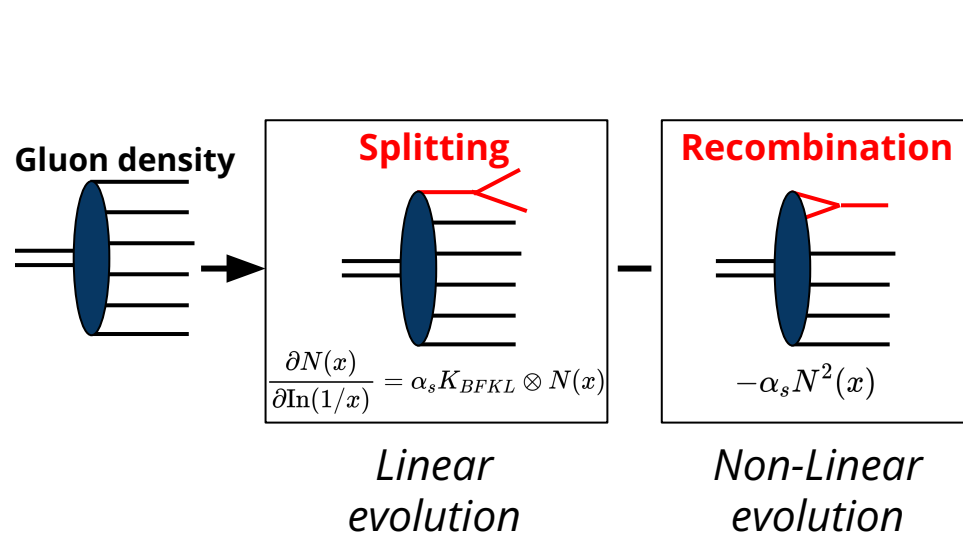
Incoherent production:

- Photon couples to part of nucleus.
- VM $\langle p_T \rangle \sim 500$ MeV.
- Target nucleus usually breaks.

Final state kinematics directly map to:

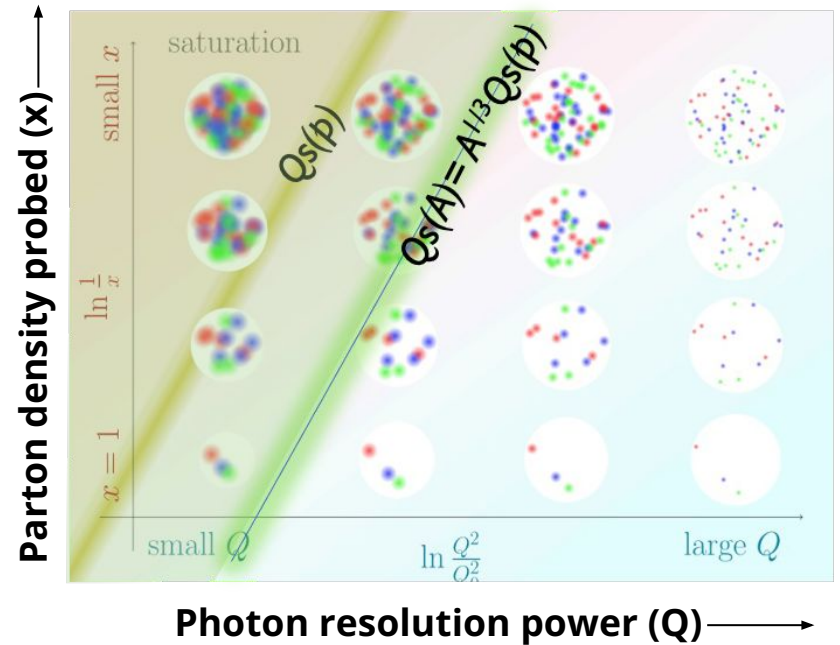
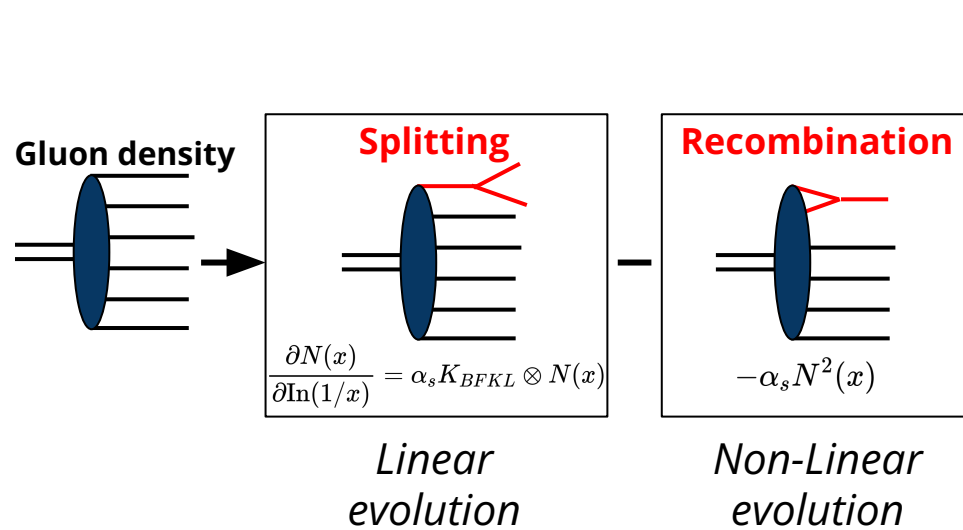
- Photon energy: $\omega = \frac{M_{VM}}{2} e^{\pm y}$
- **Bjorken-x** of gluons: $x = \frac{M_{VM}}{\sqrt{s_{NN}}} e^{\mp y}$

Gluon Saturation and Non-linear QCD regime



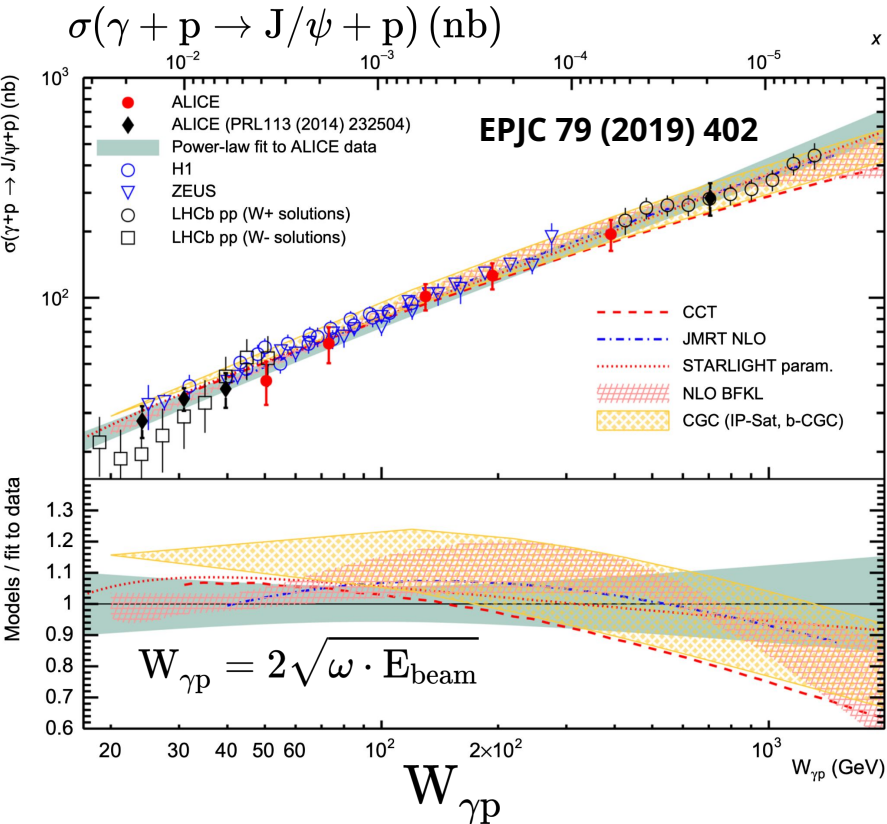
- A new regime of QCD: gluon saturation ($Q^2 < Q_s^2$) at gluon recom. = gluon splitting.

Gluon Saturation and Non-linear QCD regime



- A new regime of QCD: gluon saturation ($Q^2 < Q_s^2$) at gluon recom. = gluon splitting.
- **Saturation region is expected to be easier to be reached in nuclei: $Q_s \propto A^{1/3}$**

Search for gluon saturations in nucleons



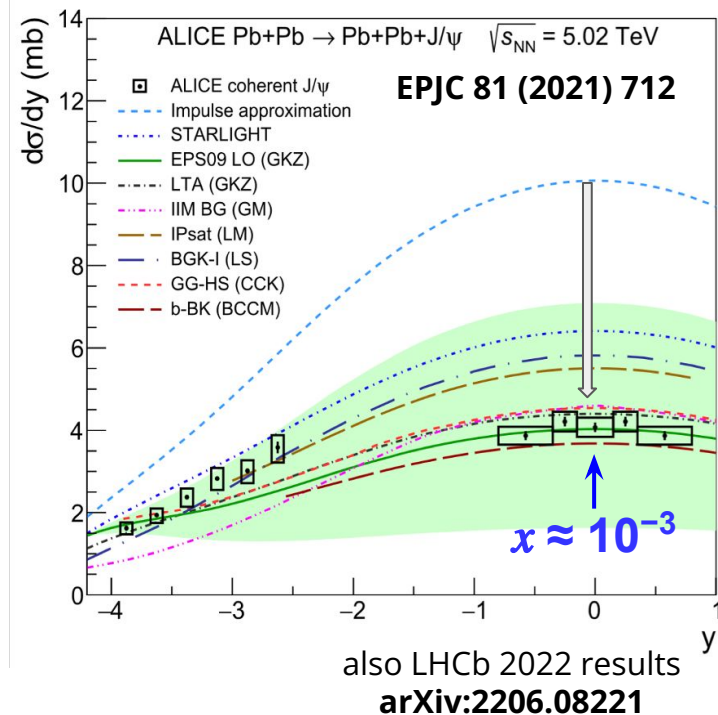
J/Ψ photoproduction from **photon-proton** interactions in ep, pPb and pp collisions

❖ Data follow a power-law trend, consistent with expectation from the rapidly increasing gluon density in a proton.

No clear indication of gluon saturation down to $x \sim 10^{-5}$ in a free nucleon.

Search for gluon saturations in heavy nuclei

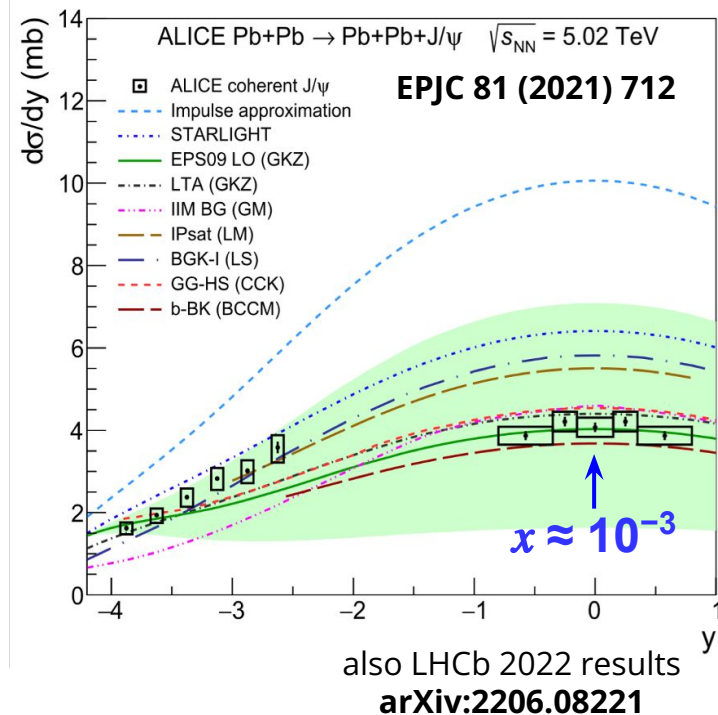
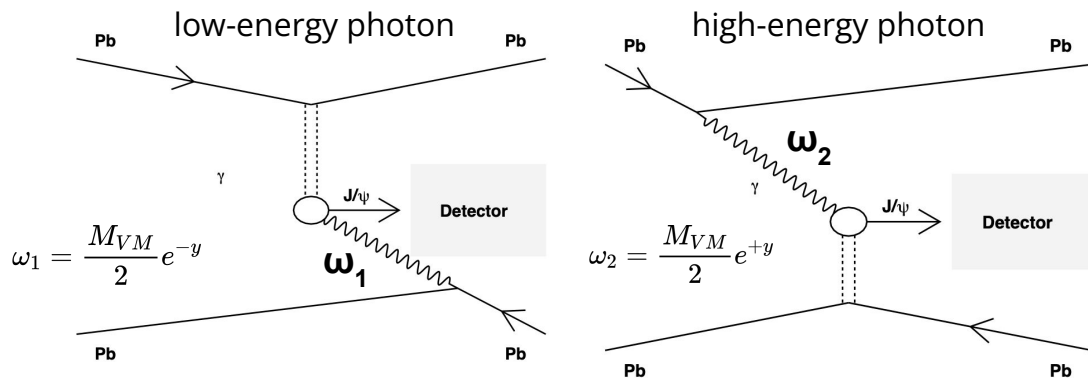
- Coherent Vector Meson production extensively measured at LHC.
- $\sigma(J/\Psi) \propto [xG(x)]^2 \rightarrow \sigma(J/\Psi) < \text{I.A.}$ (no nuclear effects) \rightarrow evidence for strong nuclear modification in heavy nuclei.
- No theory calculations (e.g., shadowing, saturation) can simultaneously predict mid- and forward rapidity data!?



Search for gluon saturations in heavy nuclei

Symmetric system: either ion can serve as the photon source or target nucleus

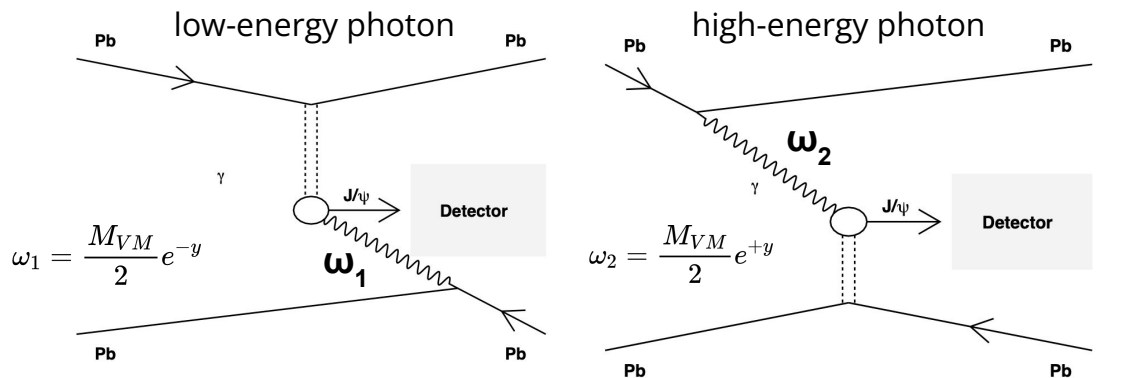
- A two-way ambiguity!



Search for gluon saturations in heavy nuclei

Symmetric system: either ion can serve as the photon source or target nucleus

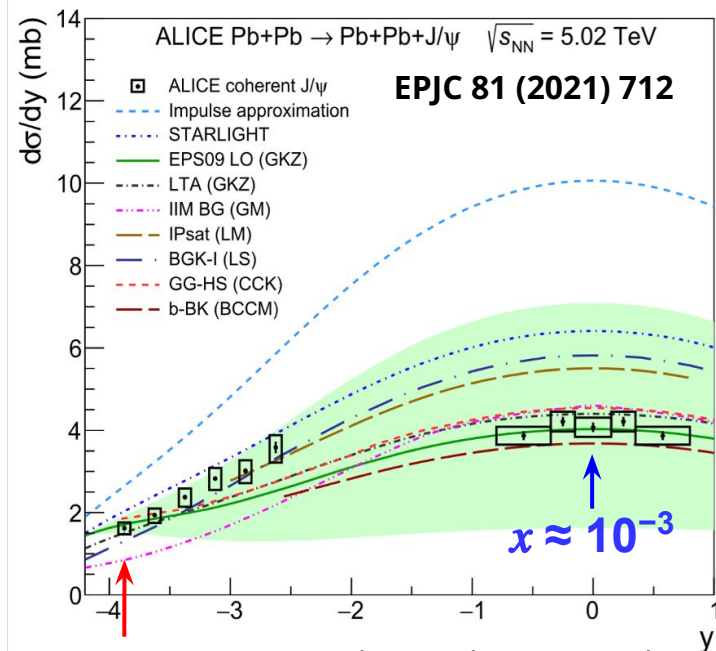
- A two-way ambiguity!



$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}}{dy} = N_{\gamma/A}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'}(\omega_1) + N_{\gamma/A}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'}(\omega_2)$$

The cross section at a given y consists of low- and high- x gluon contributions (except for $y=0$)

- **No easy access to $x \sim 10^{-5}$**



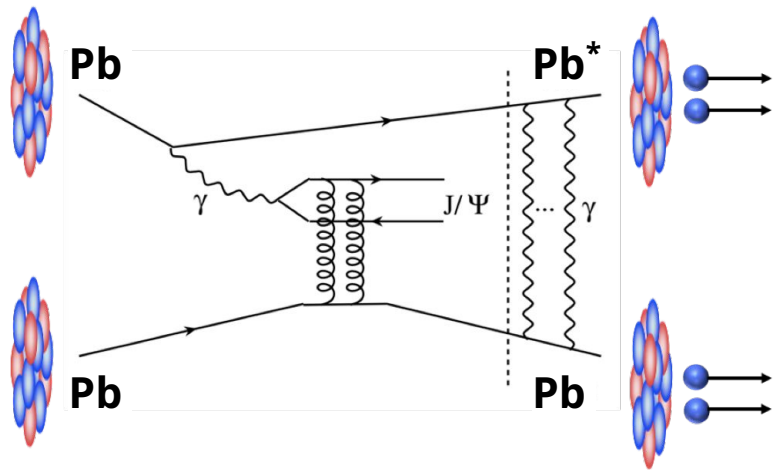
$x_1 \approx 10^{-5}$
or $x_2 \approx 10^{-2}$
(~95% high- x)

also LHCb 2022 results
arXiv:2206.08221

$$x_{1,2} = \frac{1}{\omega_{1,2}} \cdot \frac{M_{VM}^2}{2\sqrt{s_{NN}}}$$

A solution to the two-way ambiguity puzzle

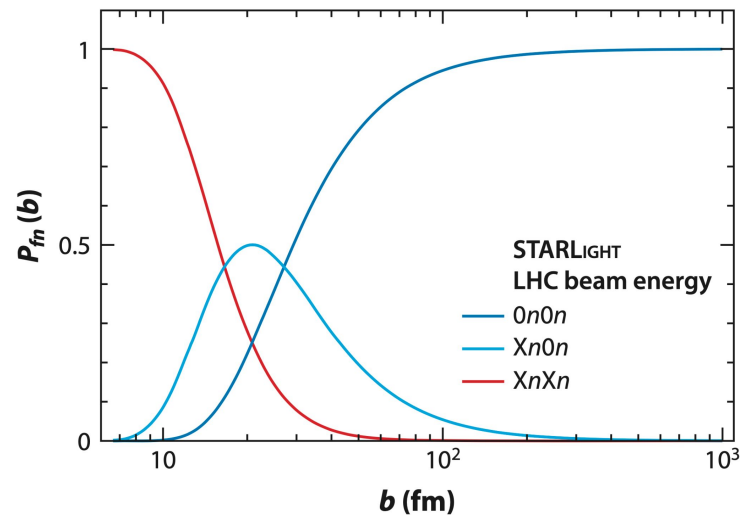
Control the impact parameter or “centrality” of UPCs via forward emitted neutrons



Nucleus excitation probability:

$$P_i(b) \propto 1/b^2$$

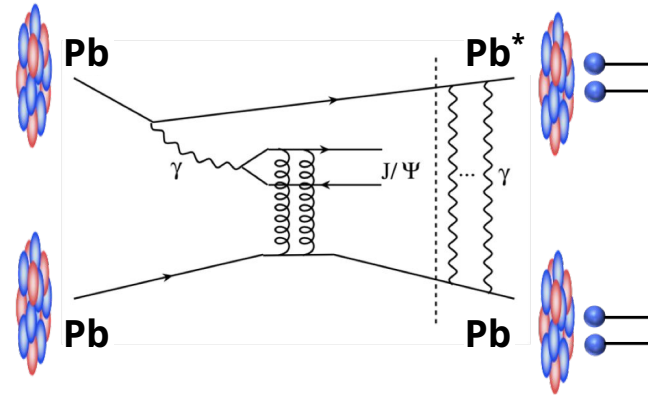
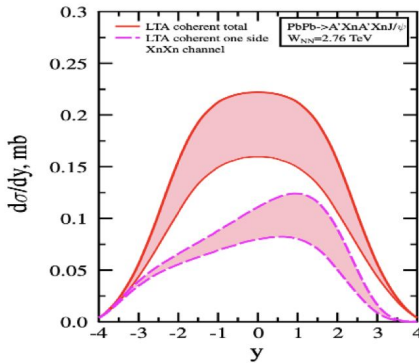
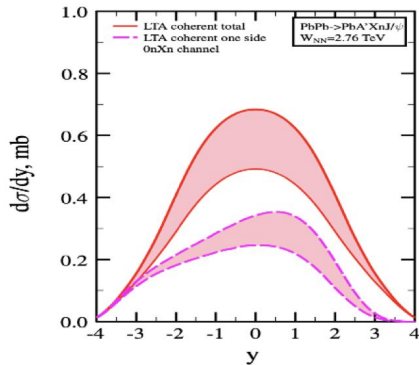
Klein & Steinberg,
Ann. Rev. Nucl. Part. Sci. **70** (2020) 323



- Analogous to centrality:
 - $b_{XnXn} < b_{0nXn} < b_{0n0n}$

A solution to the two-way ambiguity puzzle

Guzey et al., EPJC 74 (2014) 2942



$$w_1 = \frac{M_{VM}}{2} e^{-y}$$

$$w_2 = \frac{M_{VM}}{2} e^{+y}$$

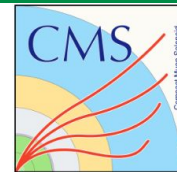
What is measured Photon flux from theory What we want to extract

$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{0nXn}}{dy} = N_{\gamma/A}^{0nXn}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'}(\omega_1) + N_{\gamma/A}^{0nXn}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'}(\omega_2)$$

$$\frac{d\sigma_{XnXn \rightarrow AA' J/\psi}^{XnXn}}{dy} = N_{\gamma/A}^{XnXn}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'}(\omega_1) + N_{\gamma/A}^{XnXn}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'}(\omega_2)$$

→ Solve for $\sigma_{\gamma A \rightarrow J/\psi A'}(\omega_1)$ and $\sigma_{\gamma A \rightarrow J/\psi A'}(\omega_2)$

Entering a new regime of small $x \sim 10^{-4}-10^{-5}$ in nuclei!



- **Introduction:**
 - General introduction to heavy ion physics and ultra-peripheral collisions.
 - State-of-the-art measurements and physics of interest.
- **Experimental setup:**
 - Compact Muon Solenoid detector.
 - Event selection and muon reconstruction.
 - Zero Degree Calorimeter and neutron multiplicity classification.
- **Controlling impact parameter dependence in UPC PbPb at 5.02 TeV:**
 - Neutron multiplicity dependence of $\mu^+\mu^-$ acoplanarity.
- **Coherent J/Ψ photoproduction in UPC PbPb at 5.02 TeV**
 - Results and discussion on new CMS result (CMS-PAS-HIN-22-002).
- **Future prospects**
- **Summary**

Compact Muon Solenoid Detector

SketchUpCMS

CMS DETECTOR

Total weight : 14,000 tonnes
Overall diameter : 15.0 m
Overall length : 28.7 m
Magnetic field : 3.8 T

STEEL RETURN YOKE
12,500 tonnes

SILICON TRACKERS

Pixel ($100 \times 150 \mu\text{m}^2$) $\sim 1.9 \text{ m}^2$ $\sim 124\text{M}$ channels
Microstrips ($80\text{--}180 \mu\text{m}$) $\sim 200 \text{ m}^2$ $\sim 9.6\text{M}$ channels

SUPERCONDUCTING SOLENOID
Niobium titanium coil carrying $\sim 18,000 \text{ A}$

MUON CHAMBERS

Barrel: 250 Drift Tube, 480 Resistive Plate Chambers
Endcaps: 540 Cathode Strip, 576 Resistive Plate Chambers

PRESHOWER

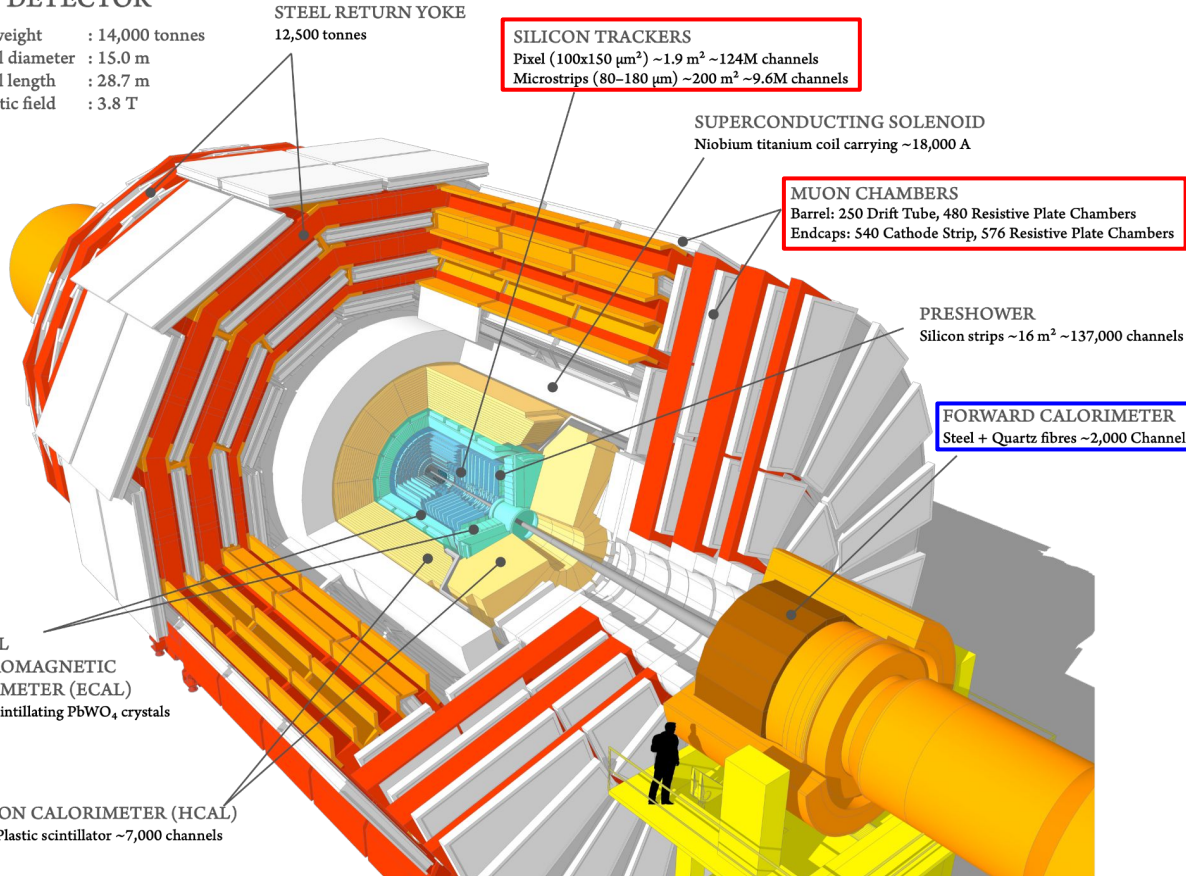
Silicon strips $\sim 16 \text{ m}^2$ $\sim 137,000$ channels

FORWARD CALORIMETER

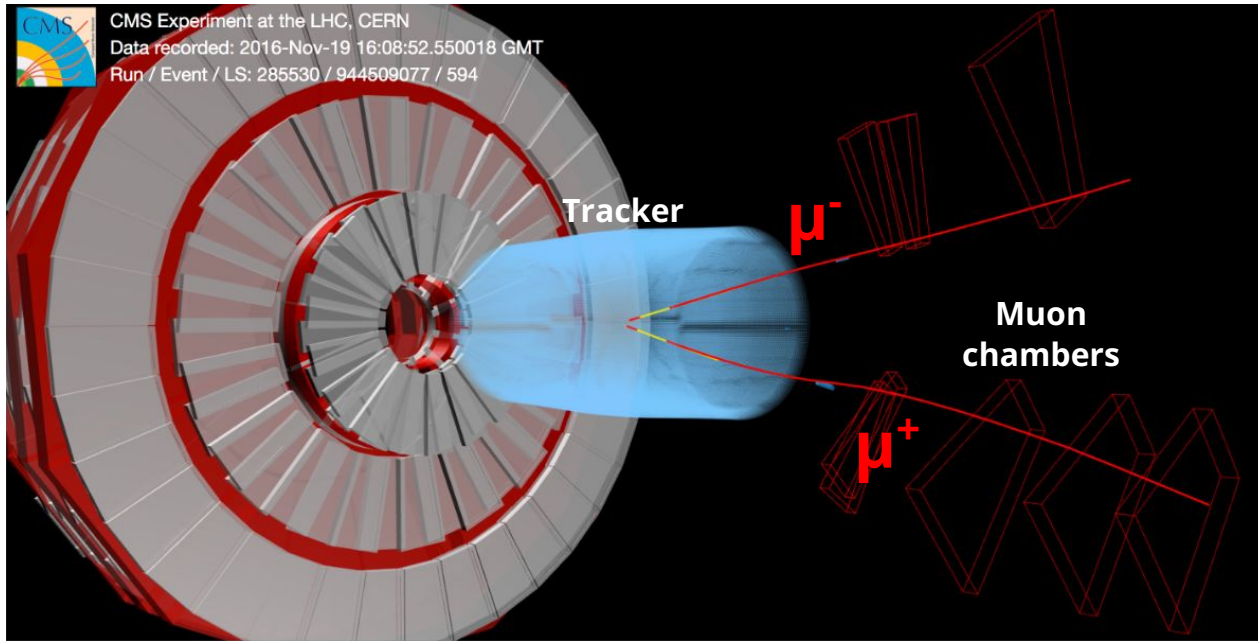
Steel + Quartz fibres $\sim 2,000$ Channels

CRYSTAL
ELECTROMAGNETIC
CALORIMETER (ECAL)
 $\sim 76,000$ scintillating PbWO_4 crystals

HADRON CALORIMETER (HCAL)
Brass + Plastic scintillator $\sim 7,000$ channels



Event selection

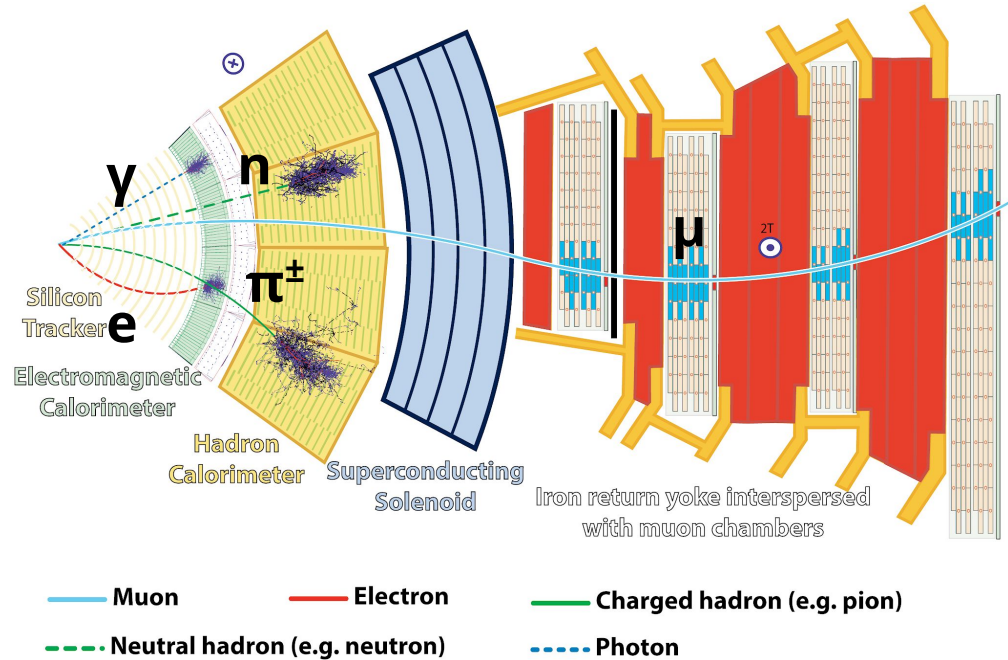
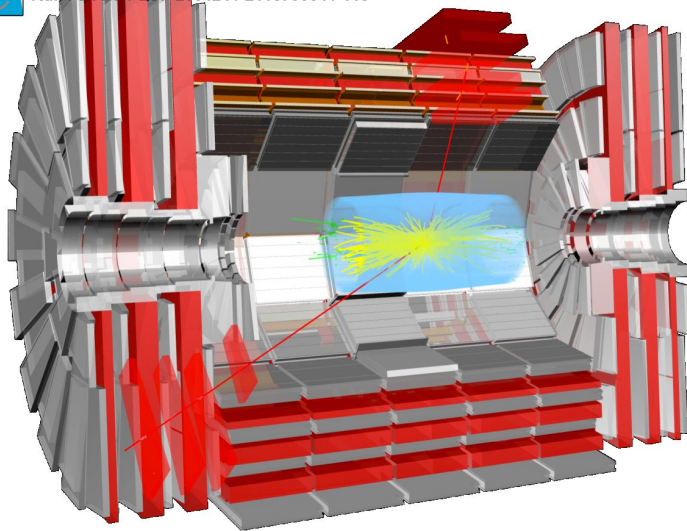


- UPC collisions produce few particles:
 - Require low energy measured in HF to suppress hadronic collisions.
- Select events with exactly two reconstructed tracks identified as muons.

Muon reconstruction

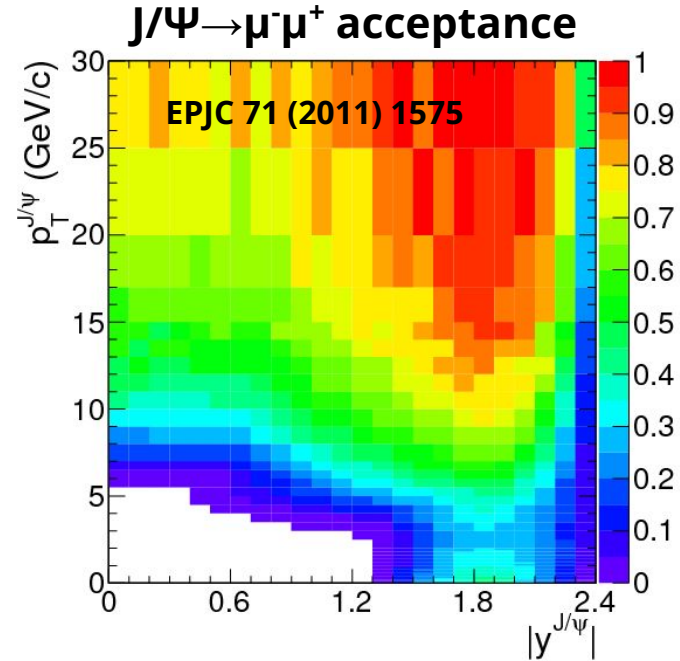
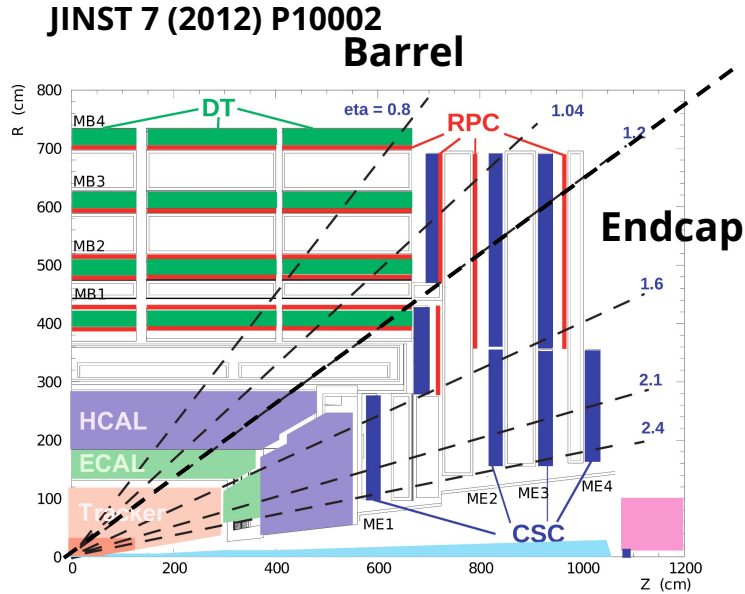


CMS Experiment at the LHC, CERN
Data recorded: 2015-Oct-30 19:23:54.631552 GMT
Run / Event / LS: / 260424 / 211873064 / 115



- Tracker and muon detectors used to reconstruct/identify muons.

Muon reconstruction



- Tracker and muon detectors used to reconstruct/identify muons.
- CMS able to reconstruct muons down to muon $p_T \sim 1$ GeV in forward region.
 - Key to identify J/Ψ mesons down to 0 GeV.

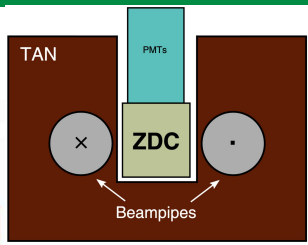
Zero Degree Calorimeter

ZDC2

POINT 5

140 m

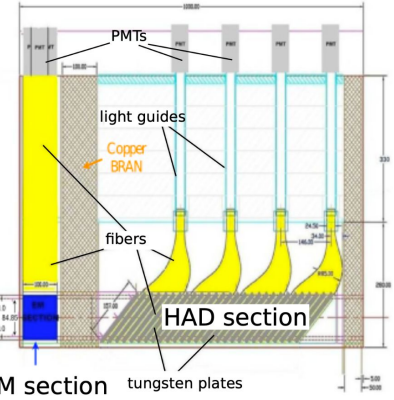
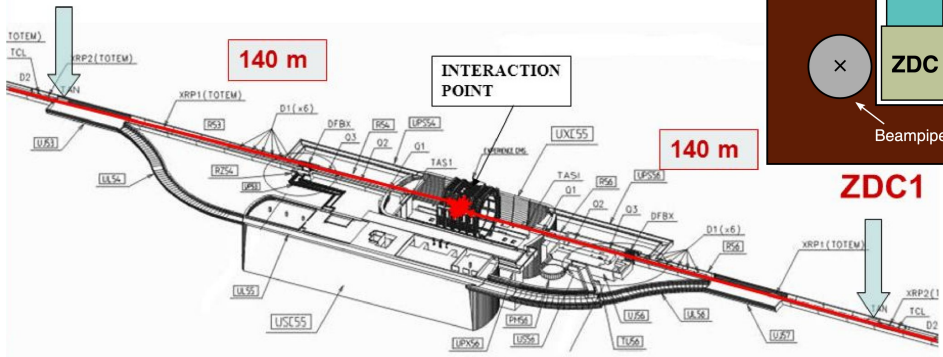
140 m



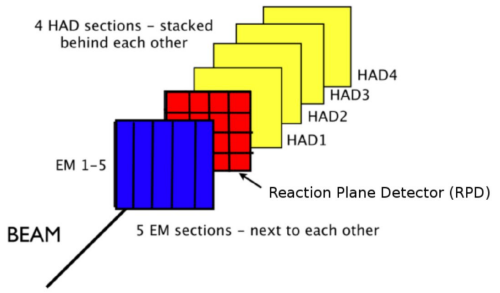
ZDC1

- Located ~ 140 m from IP5, between the two beam pipes.
- Measures forward neutral particles at $|\eta| > 8.5$.
- Used for neutron detection.

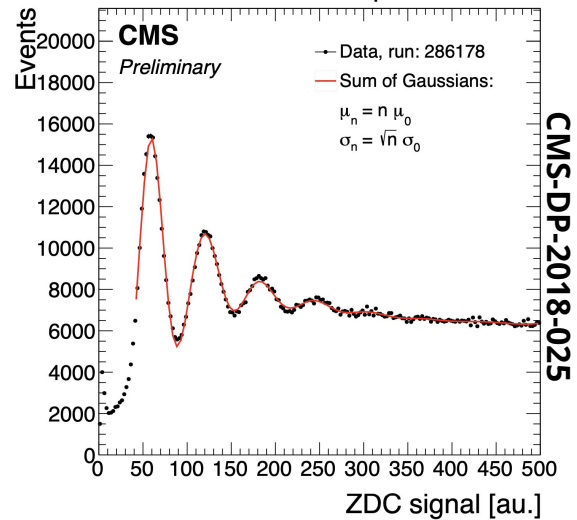
Universe 5 (2019) 210



ZDC Layout

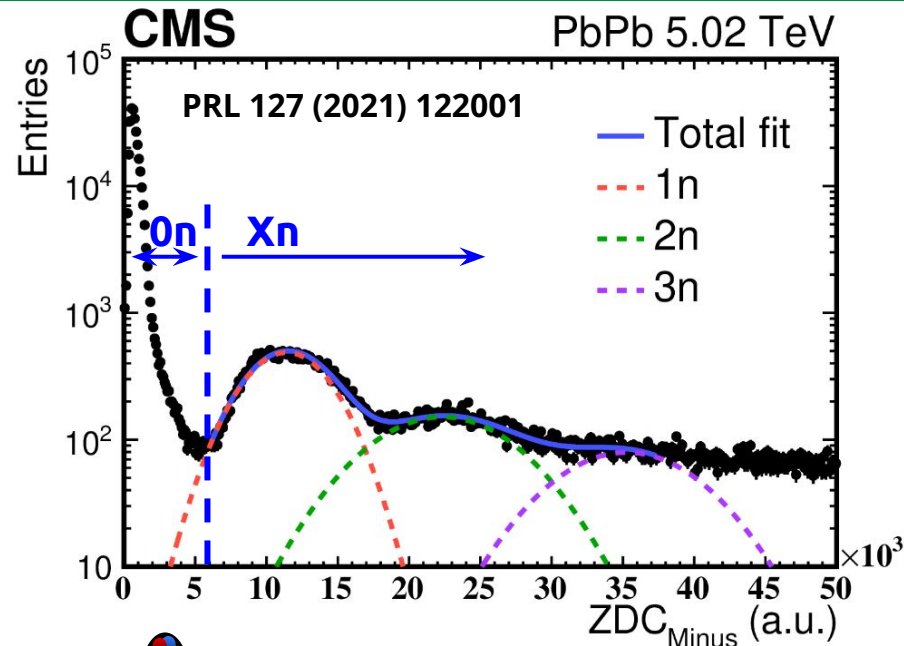
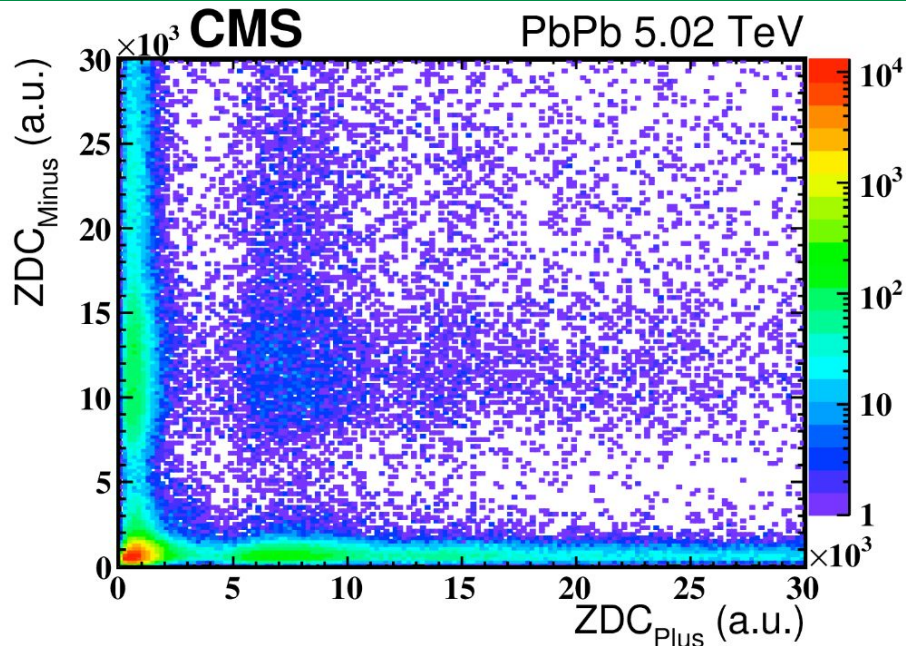


pPb 8.16 TeV

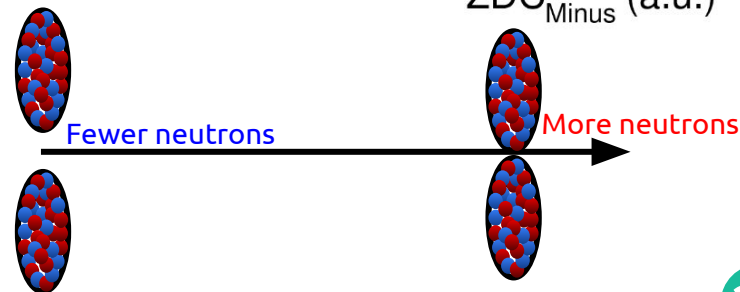


- Segmented in 1 RPD, 5 EM and 4 hadronic channels.

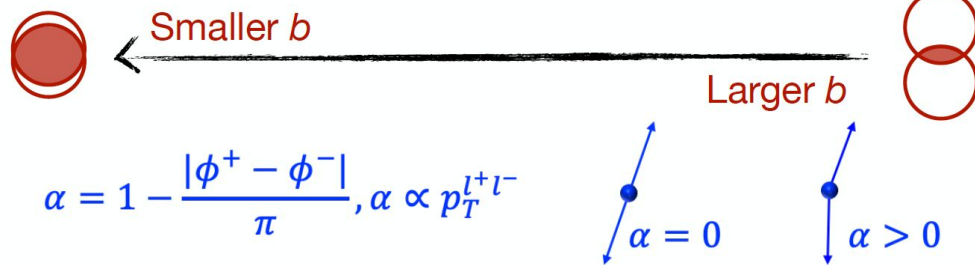
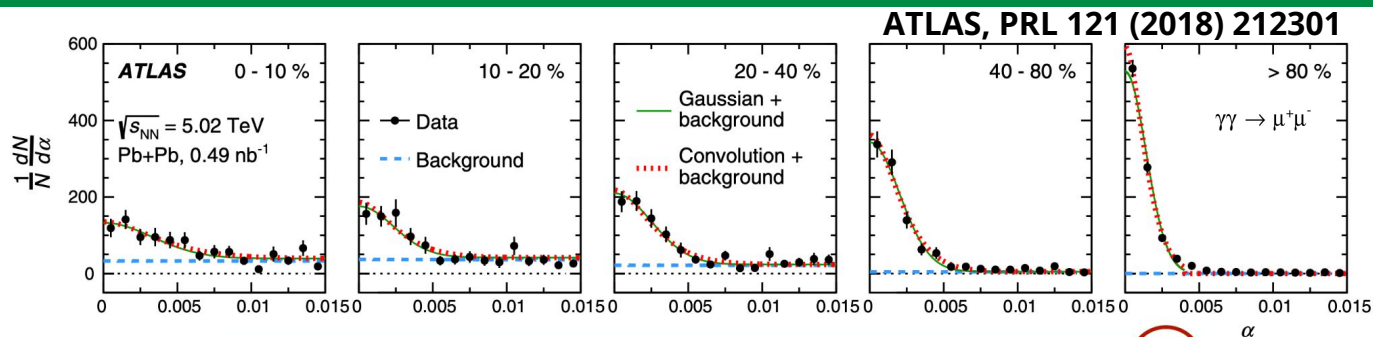
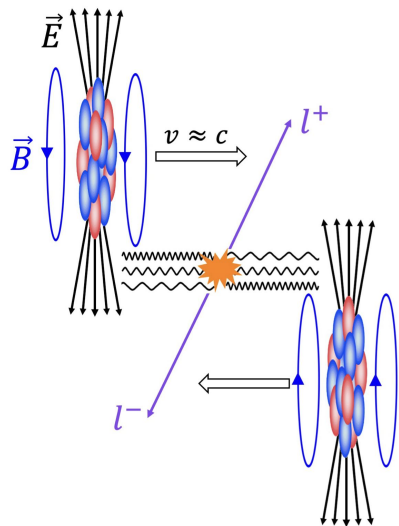
Event classification via neutron multiplicity



- Straight cut to disentangle neutrons:
 - $0n0n$, $0nXn$, $XnXn$ ($X: \geq 1$)



Strong-field QED processes with impact parameter dependence



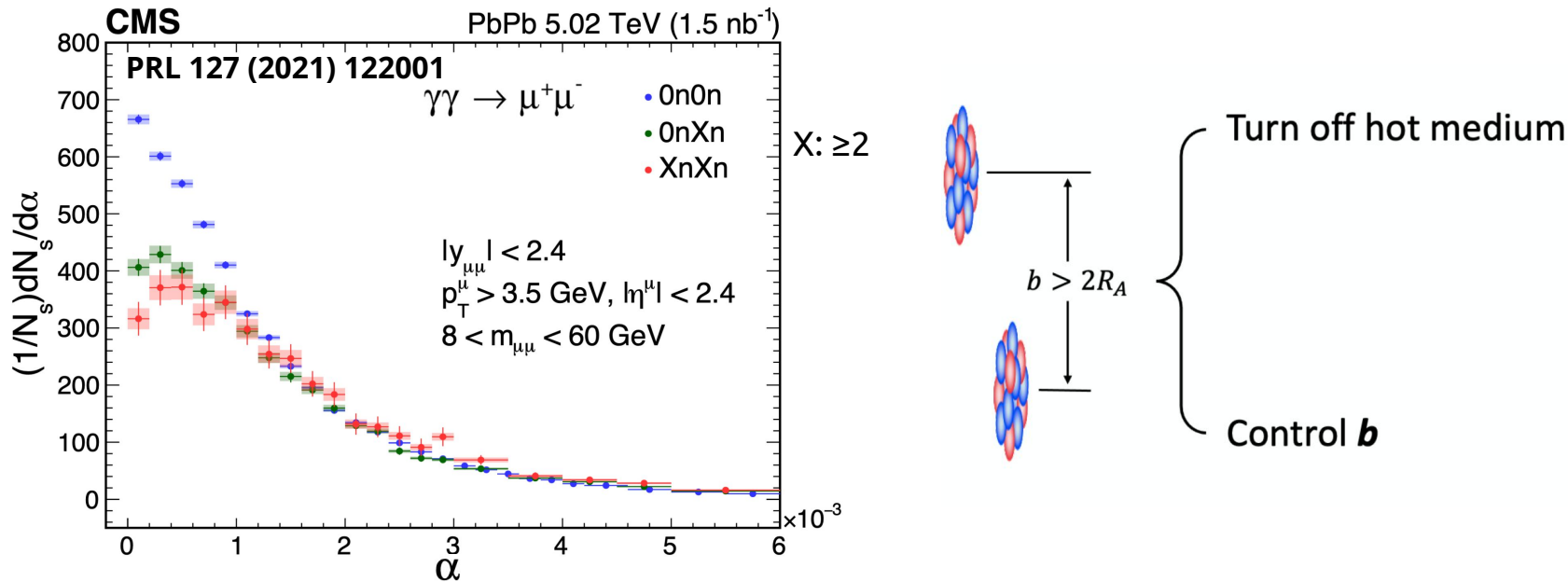
Back-to-back correlation becomes weaker towards central events

Final-State Effect?

Initial-State Effect?

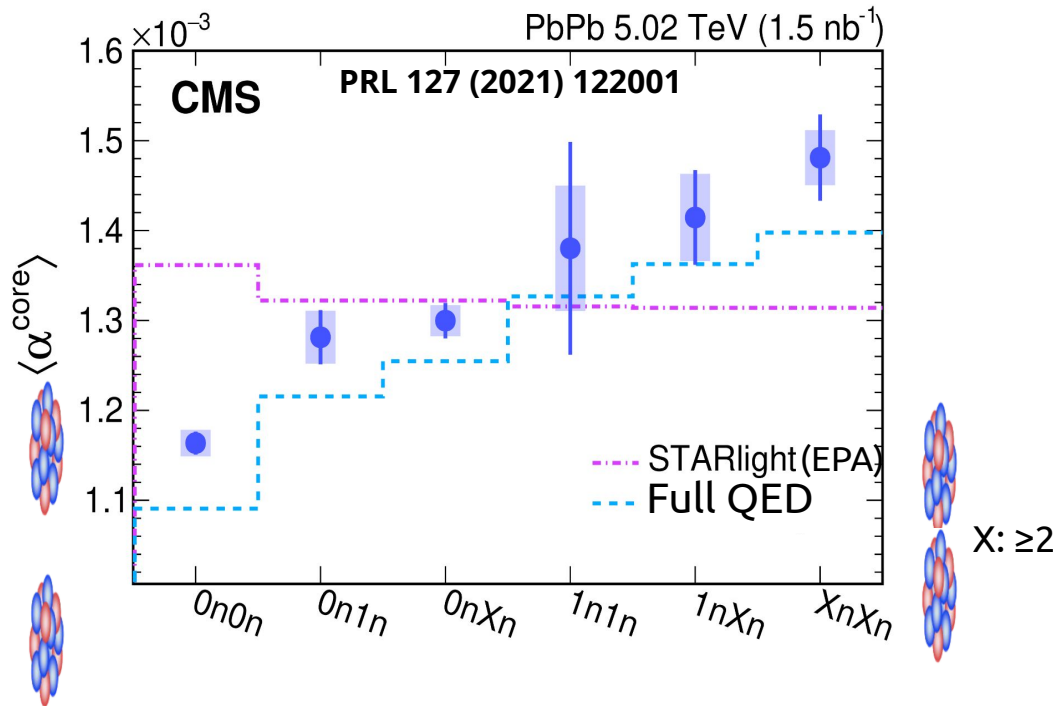
- EM rescatterings by the QGP medium?
- OR
- Initial photon p_T varies with b ?

α spectrum vs. neutron multiplicity class



Significant broadening of back-to-back correlations toward smaller b still observed when final effects are turned off, which must be related to the initial photon p_T

$\langle \alpha^{\text{core}} \rangle$ vs. neutron multiplicity class

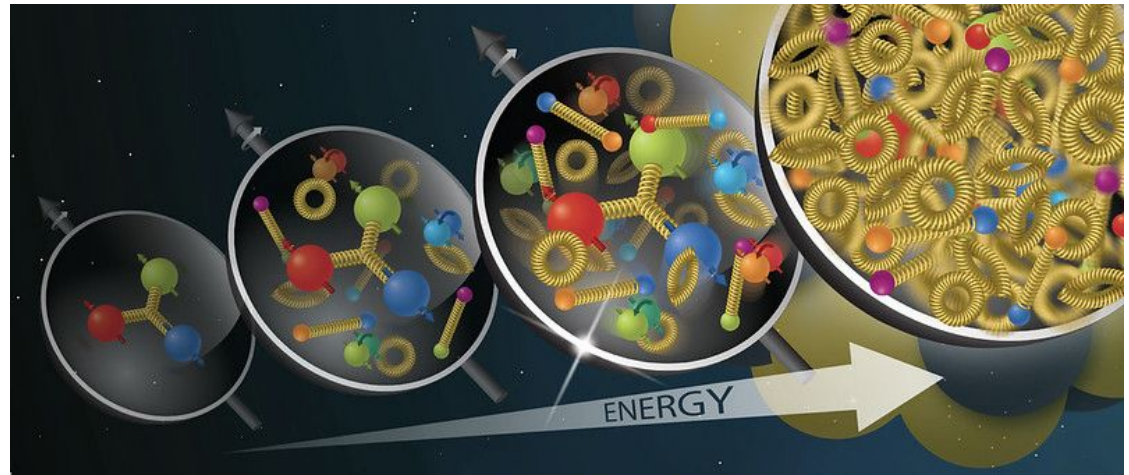
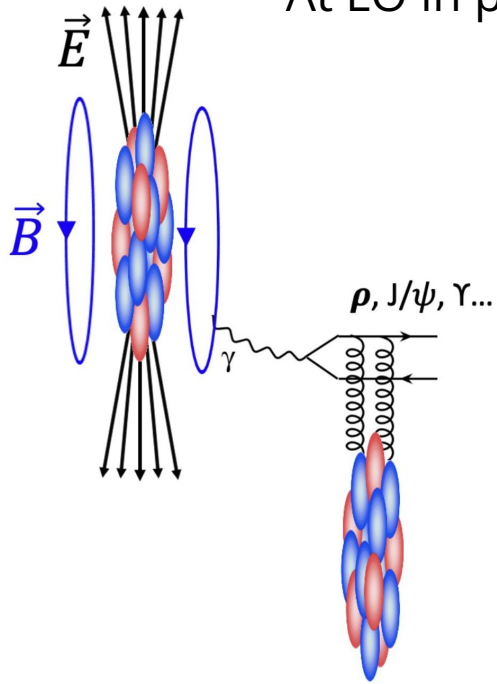


Strong (5.7σ) neutron multiplicity dependence of $\langle \alpha^{\text{core}} \rangle \propto \gamma p_T$

- b dependence of initial photon p_T , not captured by STARLight
- **Described by a leading order QED calculation with b dependence.**

Imaging the gluonic structure of nuclei

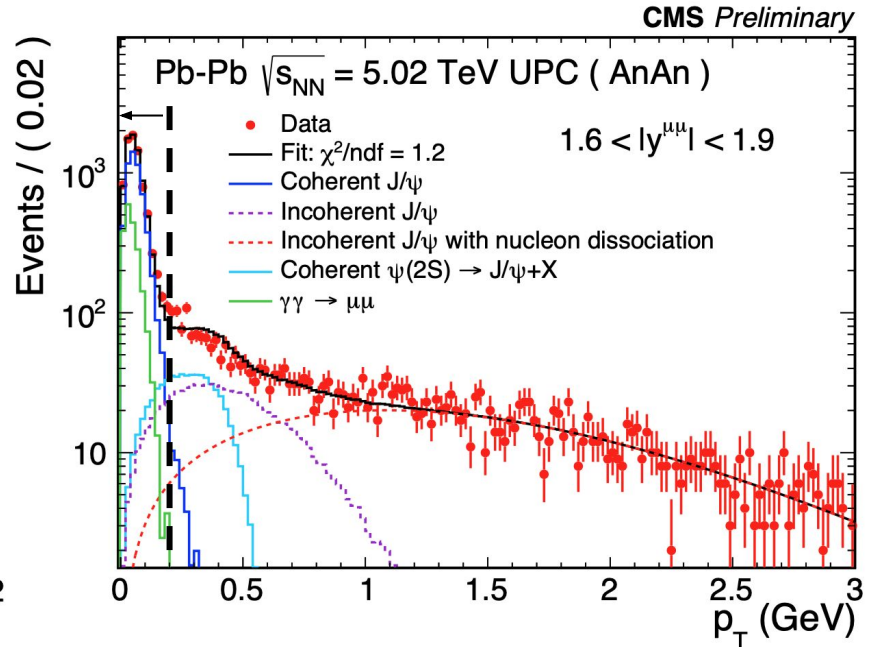
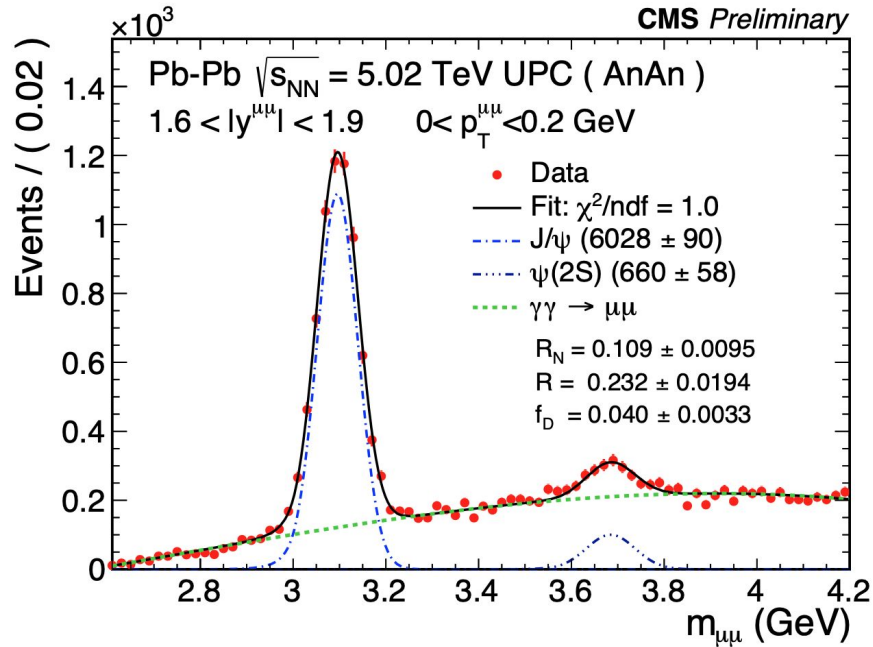
At LO in pQCD, cross section \sim photon flux $\otimes [xG(x)]^2$ (gluon PDFs)



Smaller x

Signal extraction

CMS-PAS-HIN-22-002



Signal yields are extracted by fitting the mass and transverse momentum spectra.

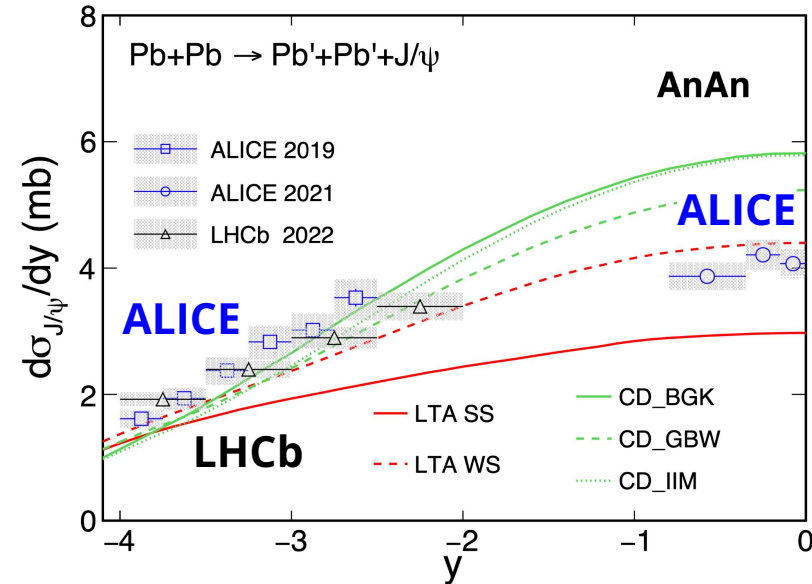
AnAn: All possible neutron emissions

Coherent J/ψ in AnAn

ALICE, EPJC 81 (2021) 712

LHCb, arXiv:2206.08221

$$\frac{d\sigma_{J/\psi}^{coh}}{dy} = \frac{N(J/\psi)}{(1 + f_I + f_D) \cdot \epsilon(J/\psi) \cdot Acc(J/\psi) \cdot BR(J/\psi \rightarrow \mu\mu) \cdot L_{int} \cdot \Delta y}$$

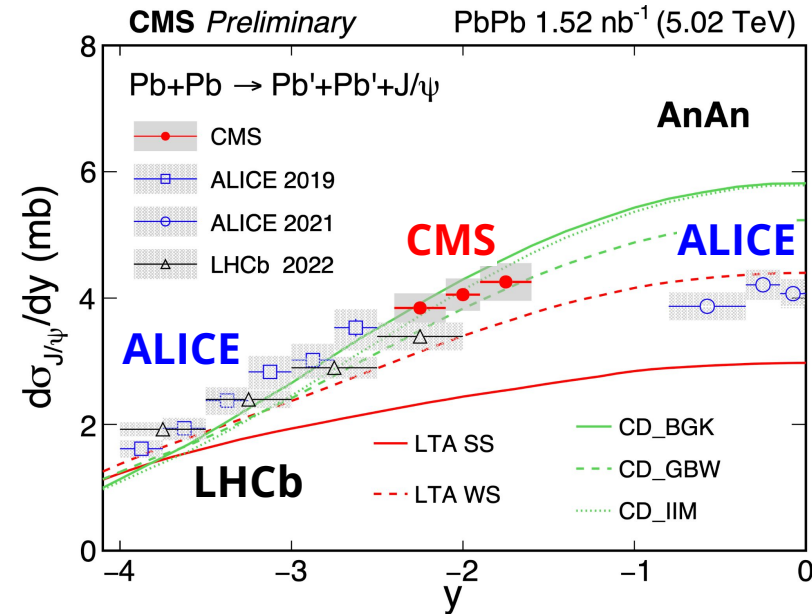


AnAn: All possible neutron emissions

- A tension between ALICE forward and mid rapidity data?
- LHCb data seems to better connect to ALICE mid rapidity data.

Coherent J/ψ in AnAn

CMS-PAS-HIN-22-002



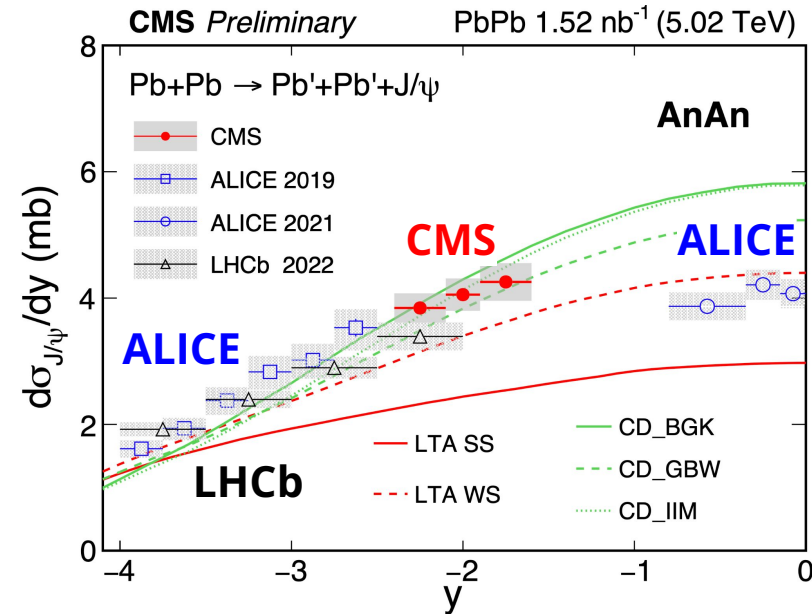
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- LHCb data seems to better connect to ALICE mid rapidity data.
- **CMS data cover a unique rapidity region.**

Coherent J/ψ in AnAn

CMS-PAS-HIN-22-002



AnAn: All possible neutron emissions

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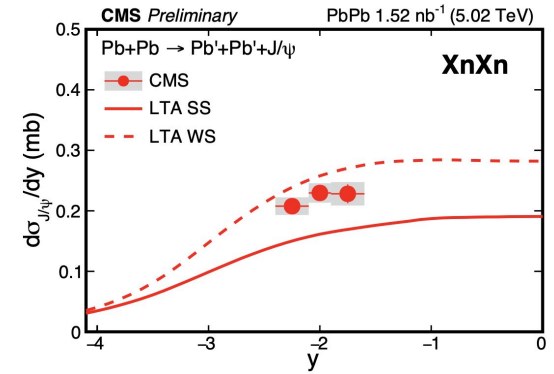
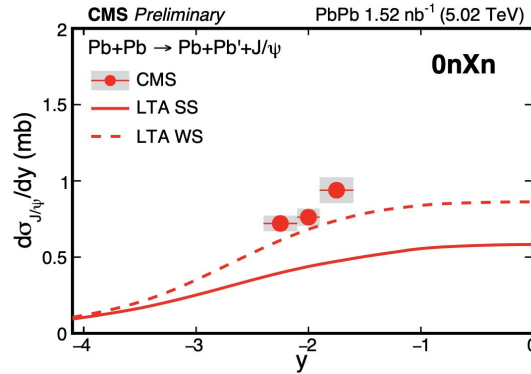
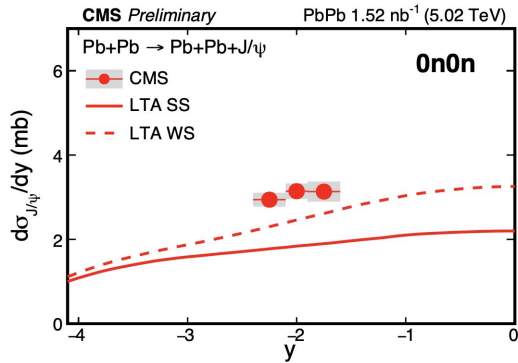
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- **CMS data cover a unique rapidity region.**

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- A deeper look at J/ψ production from γ+Pb at a given ω without the “two-way ambiguity” may tell more.

Coherent J/ψ in $0n0n$, $0nXn$, $XnXn$

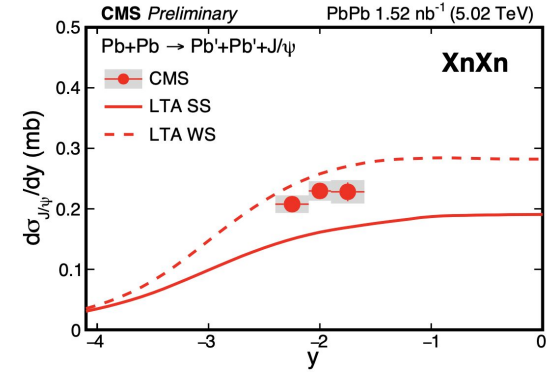
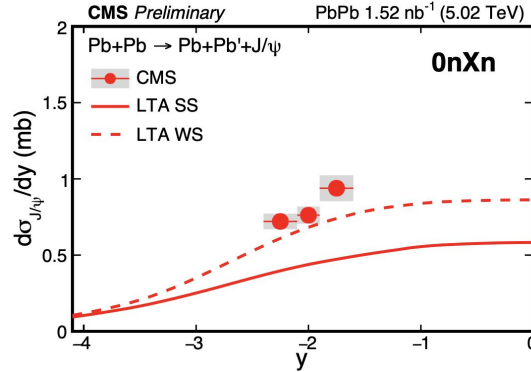
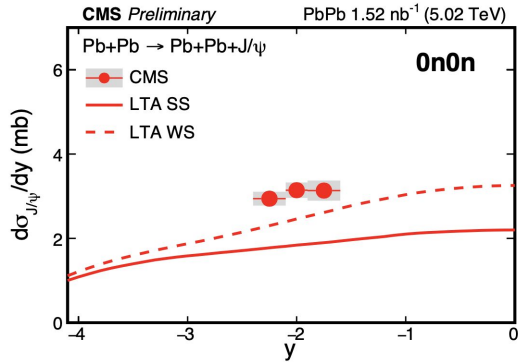
CMS-PAS-HIN-22-002



- Data in $0n0n$ and $0nXn$ are higher than Leading Twist Approximation (LTA) prediction.
- Data in $XnXn$ stay in between LTA weak suppression (WS) and strong suppression (SS) assumptions.

Coherent J/ψ in $0n0n$, $0nXn$, $XnXn$ help to disentangle

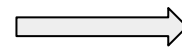
CMS-PAS-HIN-22-002



$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{0n0n}}{dy} = N_{\gamma/A}^{0n0n}(w_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'}(w_1) + N_{\gamma/A}^{0n0n}(w_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'}(w_2)$$

$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{0nXn}}{dy} = N_{\gamma/A}^{0nXn}(w_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'}(w_1) + N_{\gamma/A}^{0nXn}(w_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'}(w_2)$$

$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{XnXn}}{dy} = N_{\gamma/A}^{XnXn}(w_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'}(w_1) + N_{\gamma/A}^{XnXn}(w_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'}(w_2)$$



Low-energy γ

$$w_1 = \frac{M_{VM}}{2} e^{-y}$$

High-energy γ

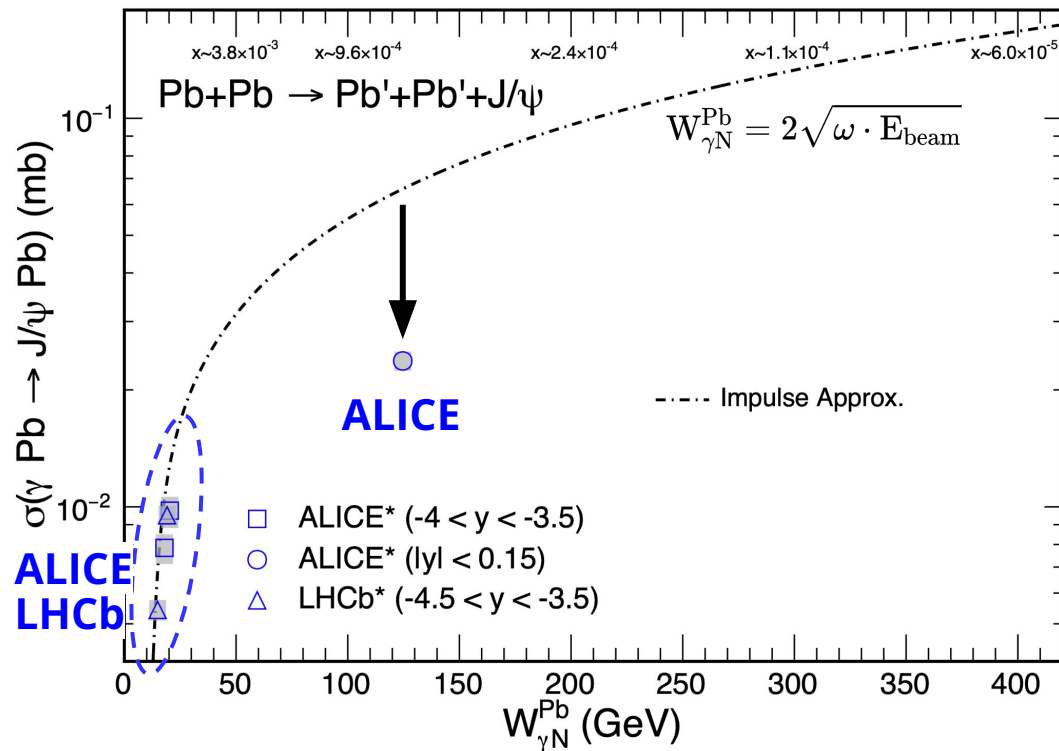
$$w_2 = \frac{M_{VM}}{2} e^{+y}$$

- Disentangle the low- and high- energy photon-nucleus contributions of a single γ +Pb.

Coherent J/ψ cross section of single γ +Pb vs. W

ALICE, EPJC 81 (2021) 712

LHCb, arXiv:2206.08221

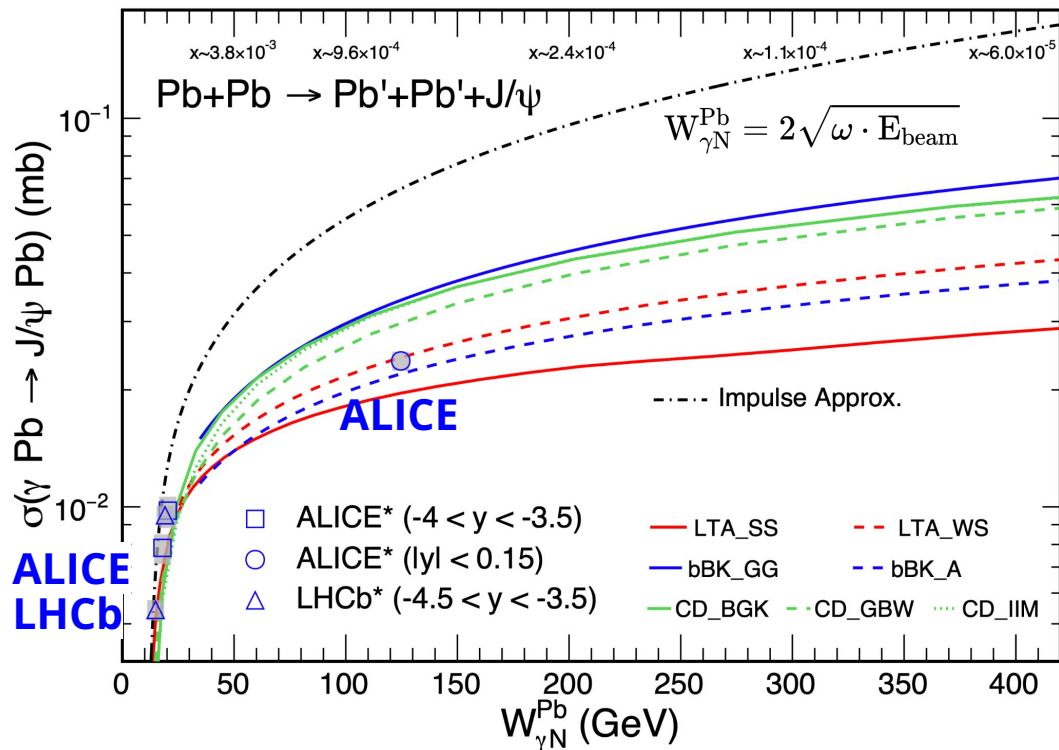


ALICE, LHCb vs. IA:

- Data is close to IA at low W .
- Data is significant lower than IA at $W \sim 125$ GeV.

Coherent J/ψ cross section of single γ +Pb vs. W

ALICE, EPJC 81 (2021) 712
LHCb, arXiv:2206.08221



ALICE, LHCb vs. IA:

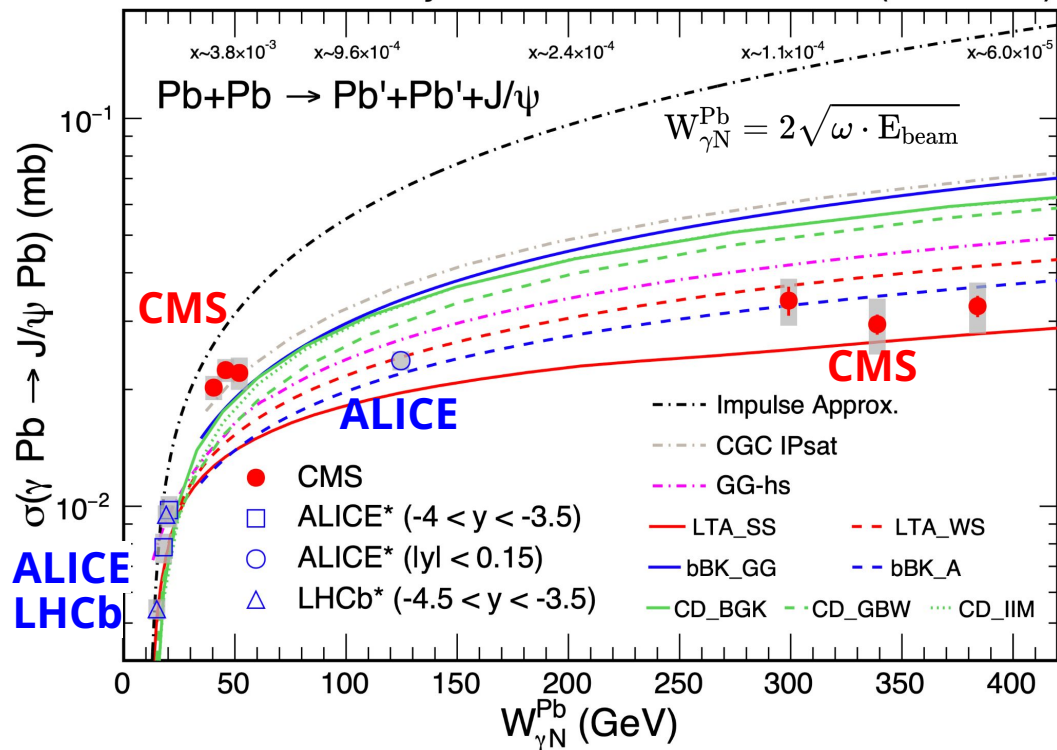
- Data is close to IA at low W .
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- Larger suppression towards higher W .

Coherent J/ψ cross section of single γ +Pb vs. W

CMS-PAS-HIN-22-002

CMS Preliminary

PbPb 1.52 nb⁻¹ (5.02 TeV)



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New data from CMS:

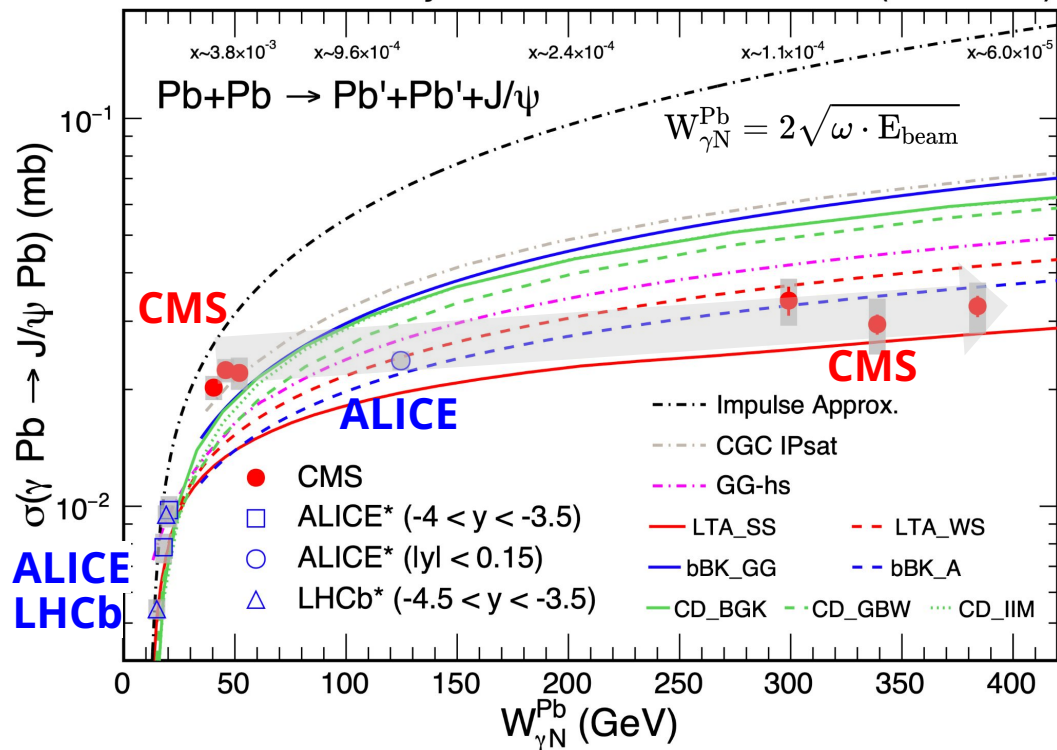
- Rapid increase at $W < 40$ GeV.

Coherent J/ψ cross section of single γ +Pb vs. W

CMS-PAS-HIN-22-002

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- Rapid increase at $W < 40$ GeV.
- Turn into a nearly flat trend for $W > 40$ GeV.

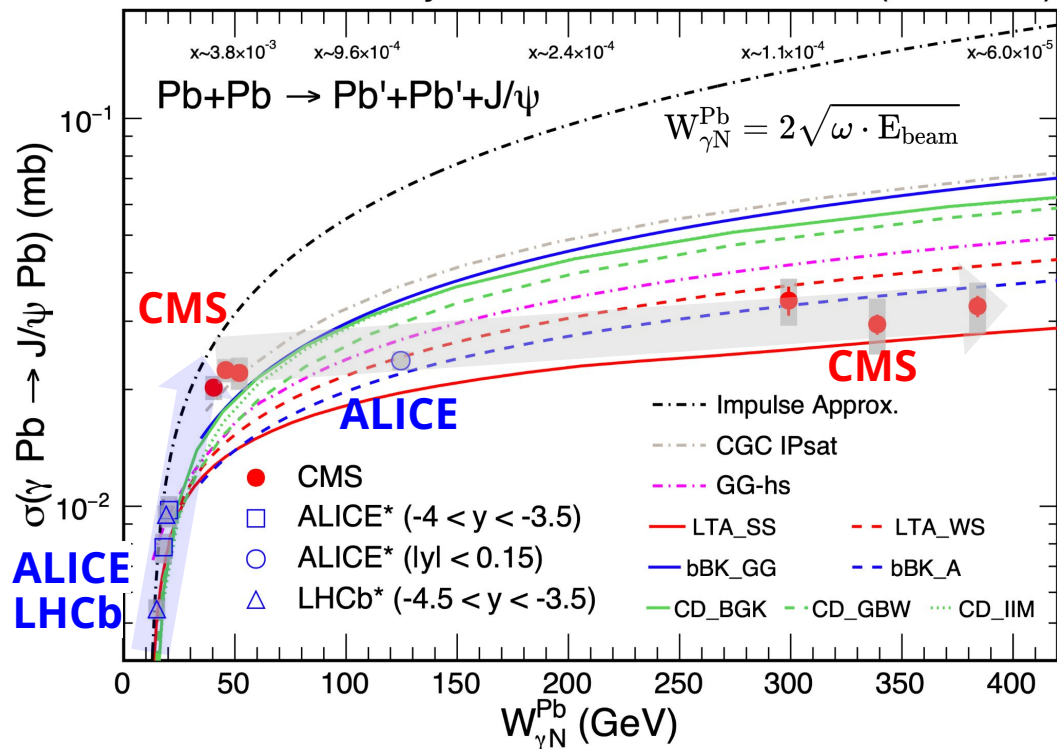
Slope = 2.98 ± 0.42 (stat.) ± 1.06 (syst.) $\times 10^{-5}$ mb/GeV

Coherent J/ψ cross section of single γ +Pb vs. W

CMS-PAS-HIN-22-002

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New data from **CMS**:

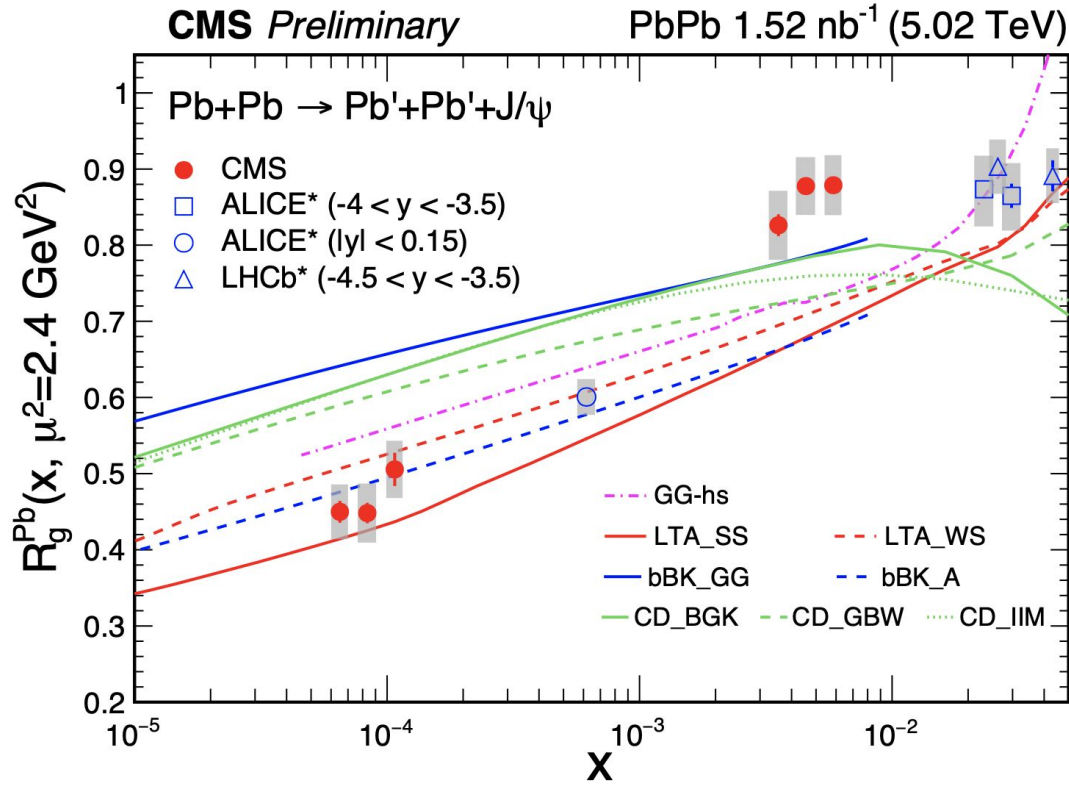
- Rapid increase at $W < 40$ GeV.
- Turn into a nearly flat trend for $W > 40$ GeV.

Slope = 2.98 ± 0.42 (stat.) ± 1.06 (syst.) $\times 10^{-5}$ mb/GeV

No models can describe the entire data distribution.

Nuclear gluon suppression factor

CMS-PAS-HIN-22-002



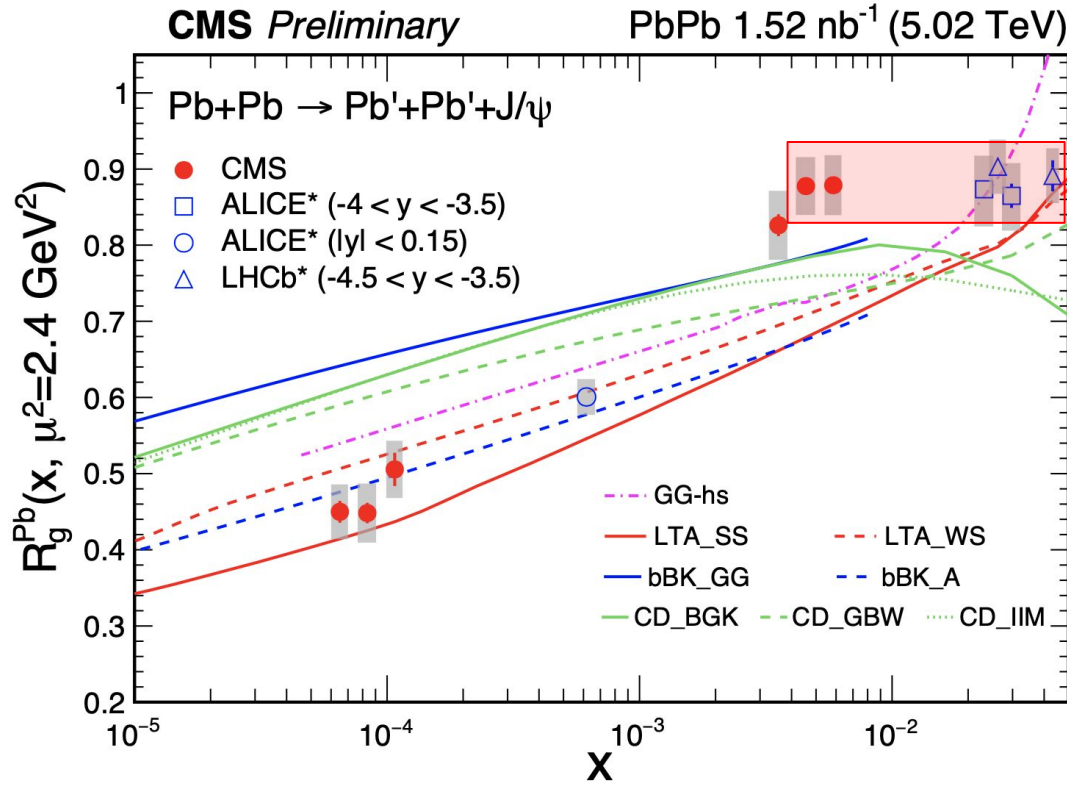
$$R_g^A = \left(\frac{\sigma_{\gamma A \rightarrow J/\psi A}^{exp}}{\sigma_{\gamma A \rightarrow J/\psi A}^{IA}} \right)^{1/2}$$

Impulse approx. (IA)
neglects all nuclear effects.

- R_g represents nuclear gluon suppression factor at LO.

Nuclear gluon suppression factor

CMS-PAS-HIN-22-002



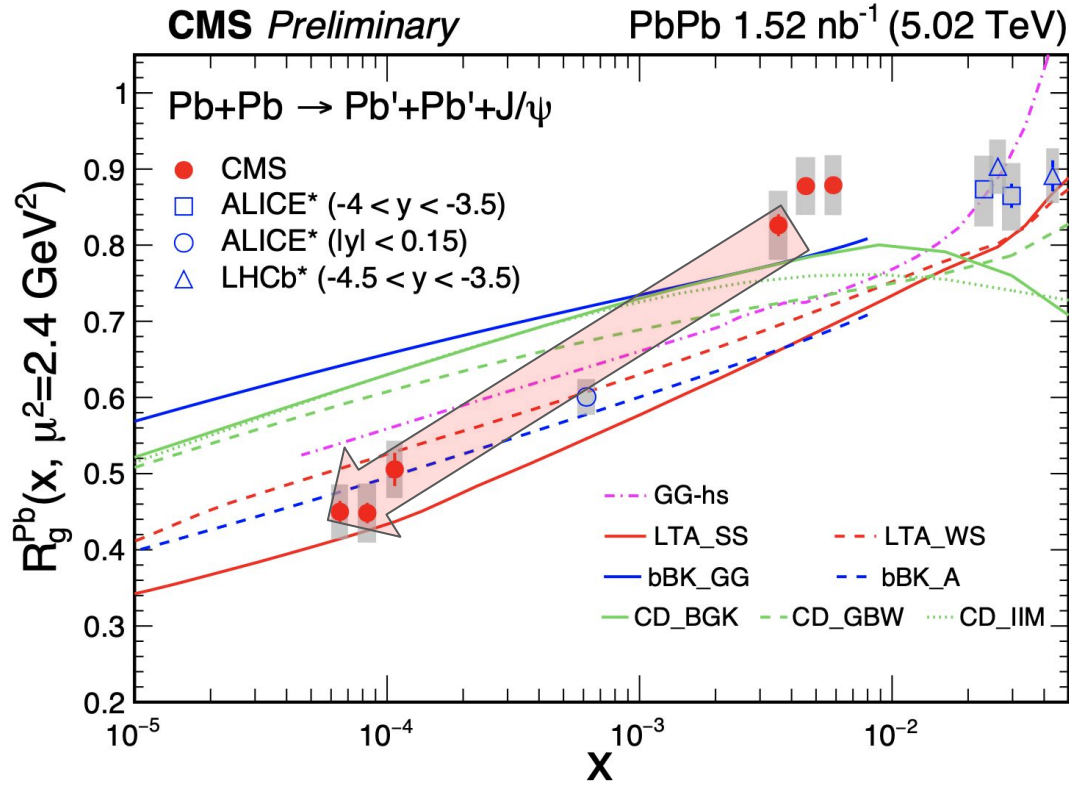
$$R_g^A = \left(\frac{\sigma_{\gamma A \rightarrow J/\psi A}^{\text{exp}}}{\sigma_{\gamma A \rightarrow J/\psi A}^{\text{IA}}} \right)^{1/2}$$

Impulse approx. (IA)
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- R_g represents nuclear gluon suppression factor at LO.
- $x \sim 10^{-3} - 10^{-2}$: flat trend.

Nuclear gluon suppression factor

CMS-PAS-HIN-22-002



$$R_g^A = \left(\frac{\sigma_{\gamma A \rightarrow J/\psi A}^{exp}}{\sigma_{\gamma A \rightarrow J/\psi A}^{IA}} \right)^{1/2}$$

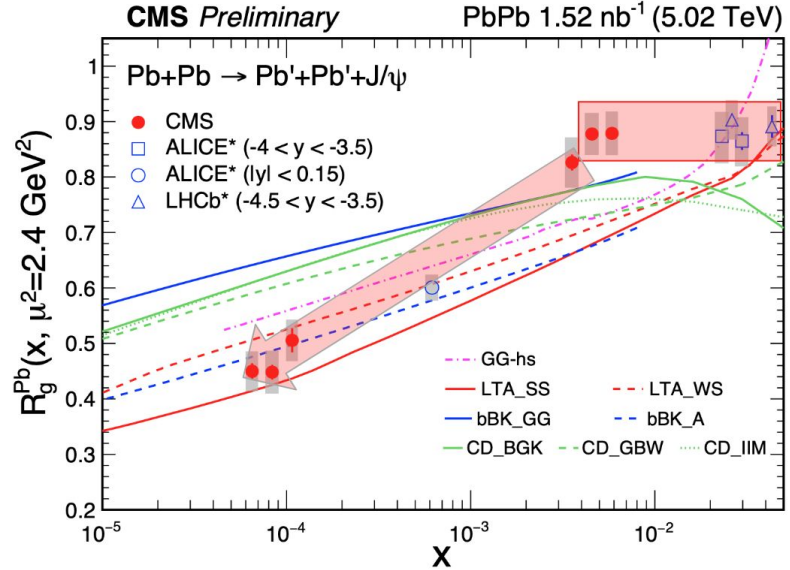
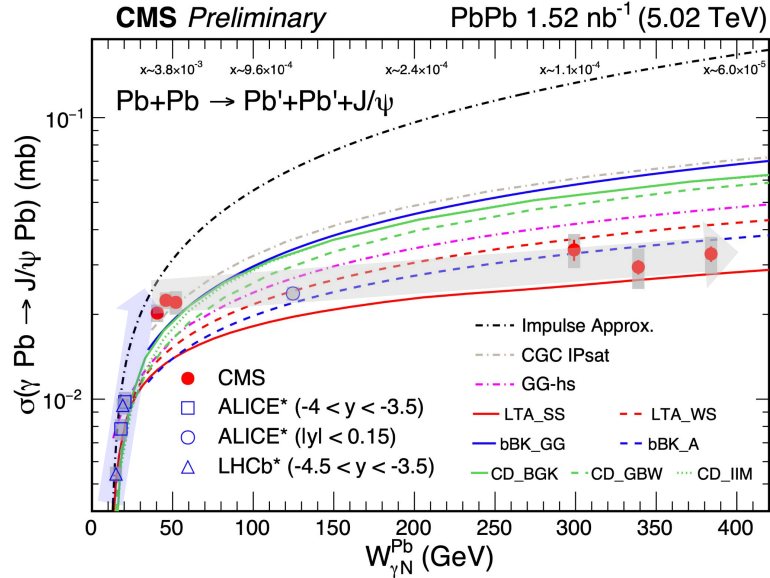
Impulse approx. (IA)
neglects all nuclear effects.

- R_g represents nuclear gluon suppression factor at LO.
- x ~ 10⁻³ - 10⁻²: flat trend.
- Quickly decrease towards lower x region.

Beyond model expectation

What we learnt

CMS-PAS-HIN-22-002

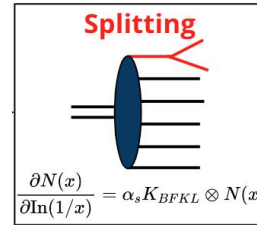
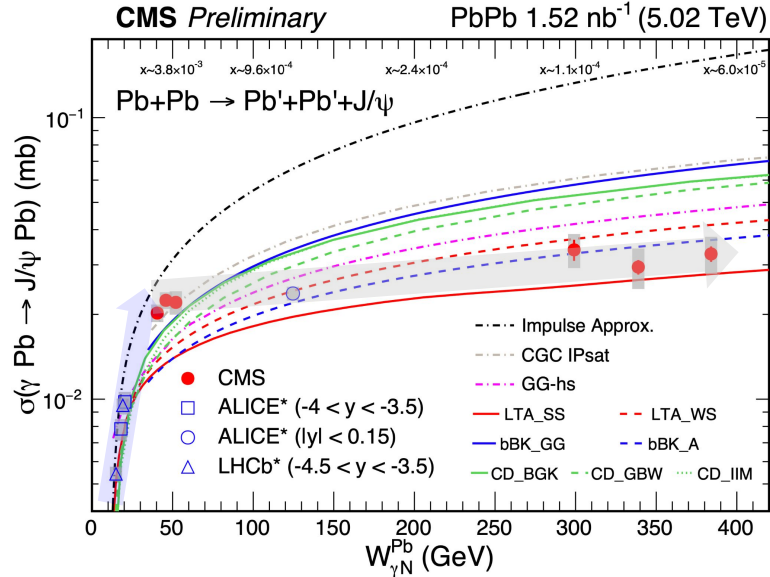


Coherent J/ Ψ production vs. W from single γ +Pb:

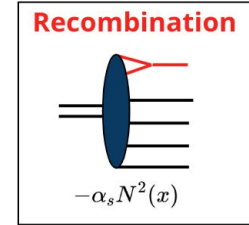
- At $W < 40$ GeV, quickly rise and the growth slope \sim free nucleons calculation.
- At $W > 40$ GeV, nearly flat with a much smaller increasing trend.
- CMS, ALICE, LHCb data follows the same trend.

What physics behind?

CMS-PAS-HIN-22-002



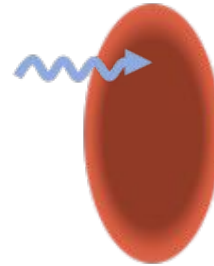
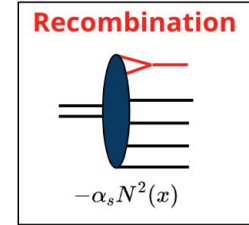
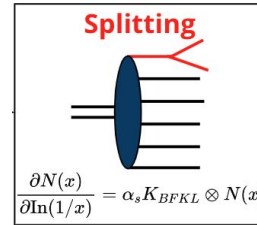
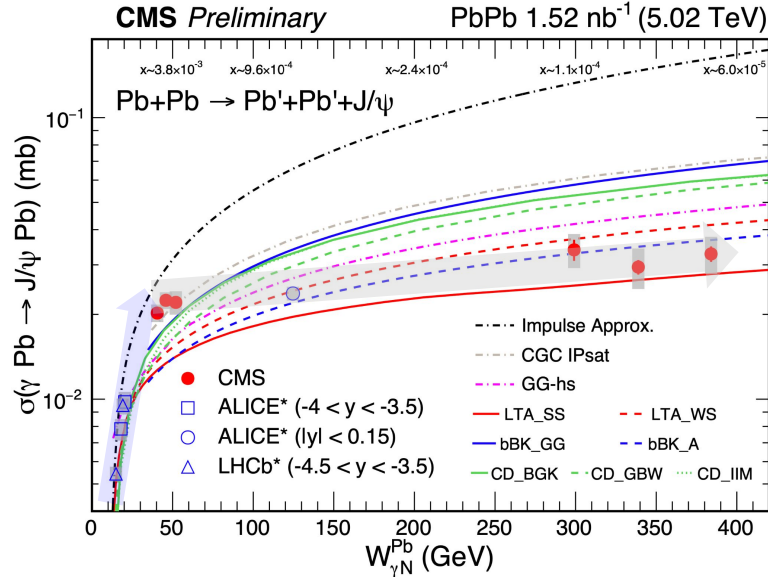
=



- σ stops rapid rising trend → splitting and recombination of gluons become equal
 - **Clear evidence for gluon saturation!!?**

What physics behind?

CMS-PAS-HIN-22-002



$$\hat{\sigma}_{\text{PQCD}}^{\text{inel}} \leq \hat{\sigma}_{\text{black}} = \pi R_{\text{target}}^2$$

- σ stops rapid rising trend → splitting and recombination of gluons become equal
 - **Clear evidence for gluon saturation!!?**

OR

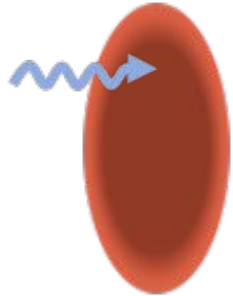
- Nucleus target becomes totally absorptive to incoming photons → **Black Disk Limit!!?**
 - **Nucleus becomes a black disk, internal structure is invisible.**

Another novel regime of QCD: Black Disk Limit

L. Frankfurt, V. Guzey, M. McDermott, M. Strikman **PRL 87 (2001)192301**

L. Frankfurt, M. Strikman, M. Zhalov, **PLB 537 (2002) 51**

In the *strong absorption scenario*, the interaction probability may reach the unitarity limit. The nucleus target becomes totally absorptive to incoming photons.



$$\hat{\sigma}_{\text{PQCD}}^{\text{inel}} \leq \hat{\sigma}_{\text{black}} = \pi R_{\text{target}}^2$$

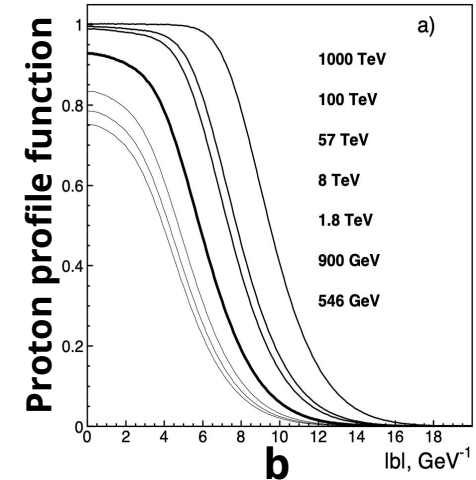
“Black Disk Limit (BDL)”

- opposite to the “color transparency”

... Inner structure disappears



energy increase

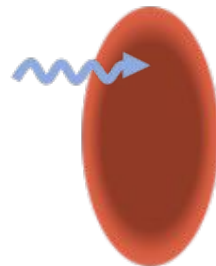
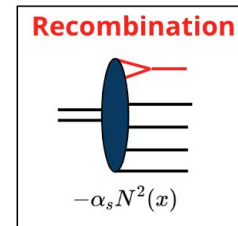
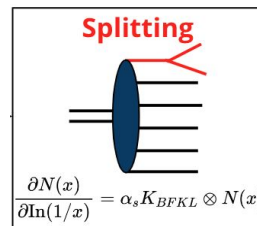
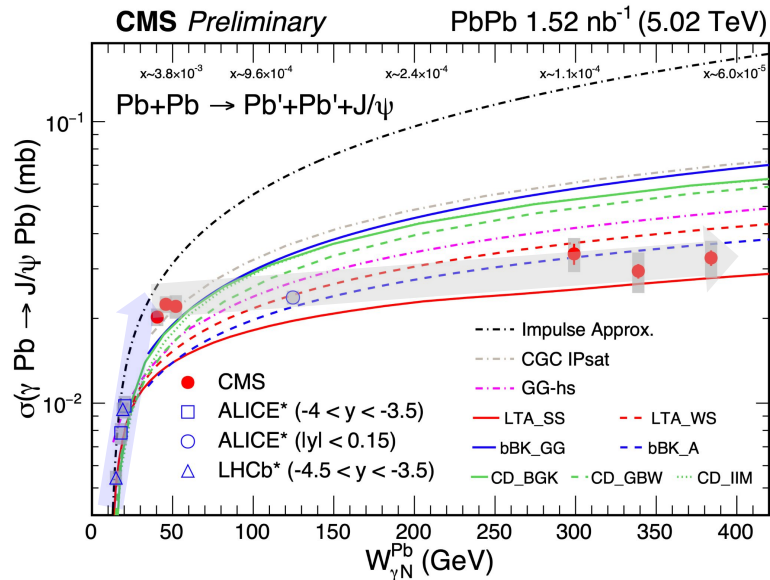


The BDL represents a novel regime at small x when the LO QCD and the notion of the parton distributions becomes inapplicable for describing hard processes .

- **New theoretical tools are needed in this regime!**

The slowly increasing trend at high W

CMS-PAS-HIN-22-002

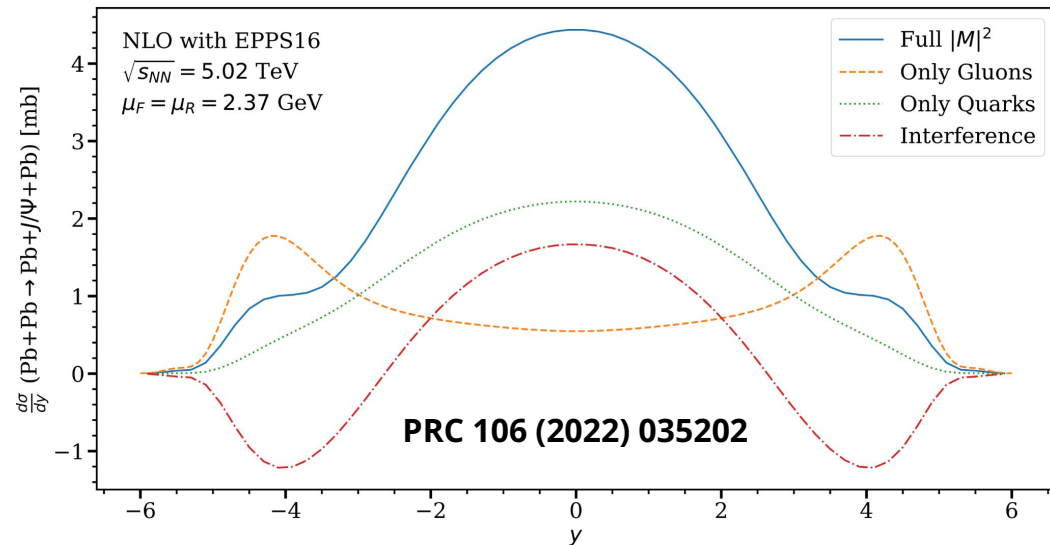


$$\hat{\sigma}_{\text{PQCD}}^{\text{inel}} \leq \hat{\sigma}_{\text{black}} = \pi R_{\text{target}}^2$$

- Periphery of nucleus may not be fully saturated or fully black at $W \sim 40$ GeV, but gradually turn to saturated or fully black with further increasing of the probing energy.

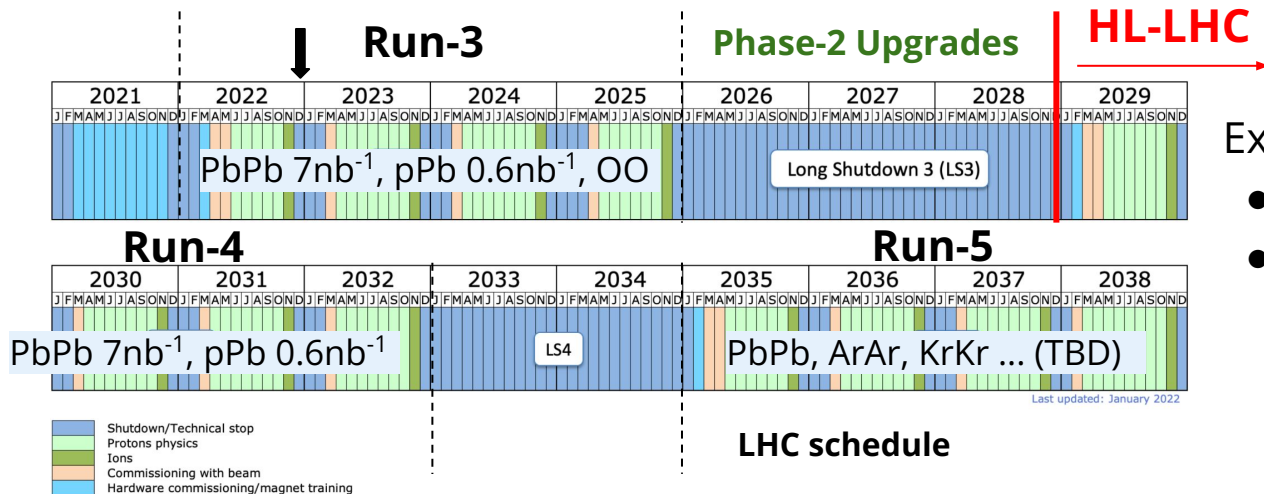
First NLO calculations on exclusive J/Ψ production

- First NLO pQCD calculations published recently (Eskola et al).
- Quark contributions at NLO + strong cancellations between LO and NLO gluons → *dominance of quark contribution at central rapidities.*



- Needs careful attention when interpreting the data.
- “ $\sigma \propto (\text{gluon PDF})^2$ ” not true at NLO.

Future opportunities



Exciting opportunities ahead by:

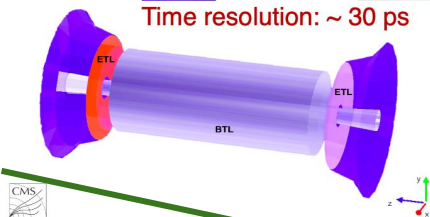
- Higher luminosities.
- A variety of ion species.

Future opportunities

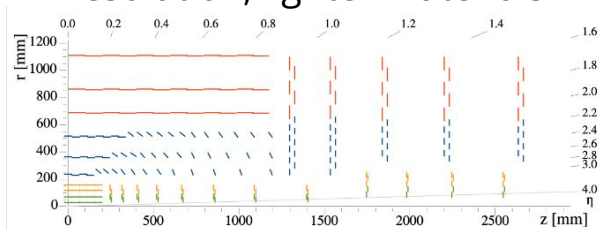
MIP Timing Detector for PID

- BTL: LYSO bars + SiPM readout:**
- The BTL readout is $\sim 4.4\sigma$
 - Inner radius: 1148 mm (40 mm thick)
 - Length: ± 2.0 m along z
 - Surface: ~ 30 m², 3324 channels
 - Fluence at 4 ab⁻¹: 2×10^{10} n_{om}/cm²
- ETL: Si with internal gain (LGAD):**
- On the CE: inner: 1.6×10^9 e⁻/s
 - Radius: 315 <math>R < 1200 mm
 - Position in z: ± 0.3 m (65 mm thick)
 - Surface: ~ 14 m², $\sim 8.9M$ channels
 - Fluence at 4 ab⁻¹: up to 2×10^{10} n_{om}/cm²

Time resolution: ~ 30 ps



Tracker with $|\eta| < 4$ and better resolution, lighter materials

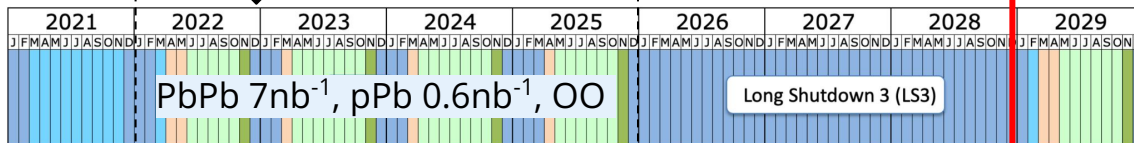


- Muon systems with $|\eta| < 2.8$
- Trigger and DAQ rate: ~ 10 x
-

Run-3

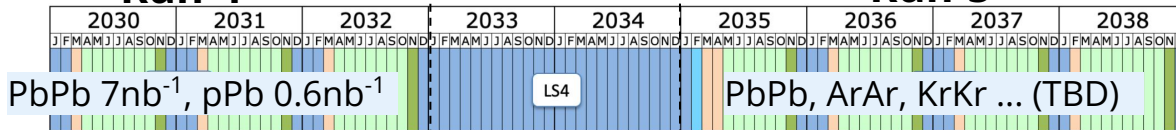
Phase-2 Upgrades

HL-LHC



Run-4

Run-5



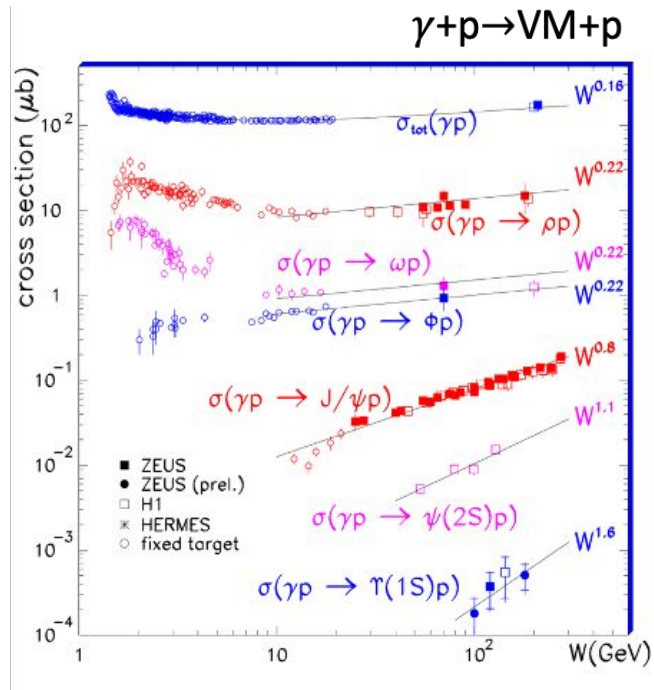
LHC schedule

- Shutdown/Technical stop
- Protons physics
- Ions
- Commissioning with beam
- Hardware commissioning/magnet training

Exciting opportunities ahead by:

- Higher luminosities.
- A variety of ion species.
- Upgrades enabled by new technologies!

Future opportunities



Various vector meson species in γPb as a function of a broad W range with neutron taggings

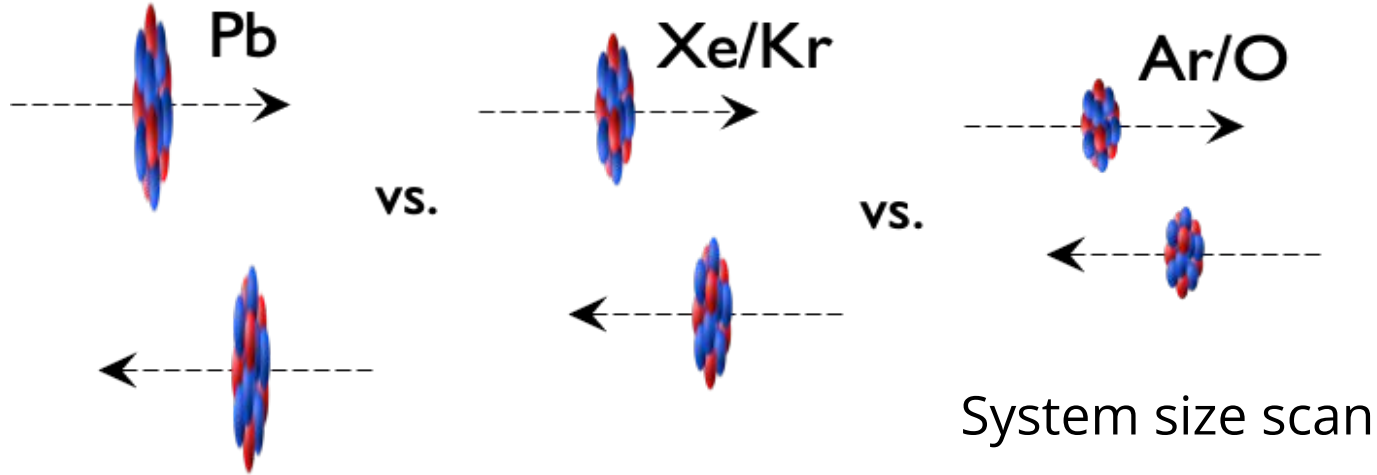
➤ e.g., control of dipole sizes and hard scales.

CERN yellow report, [arXiv:1812.06772](https://arxiv.org/abs/1812.06772)

Condition	Tot.	Central 1	Central 2	Forward 1	Forward 2
		Narrow	Wide	Narrow	Wide
Rapidity	-	$ y < 0.9$	$ y < 2.4$	$2.5 < y < 4.0$	$2 < y < 5$
$e/\pi/\mu$ pseudorapidity	-	$ \eta < 0.9$	$ \eta < 2.4$	$2.5 < \eta < 4.0$	$2 < \eta < 5$

PbPb $L_{\text{int}} = 13 \text{ nb}^{-1}$						
Meson	σ	All Total	Central 1 Total	Central 2 Total	Forward 1 Total 1	Forward 2 Total
$\rho \rightarrow \pi^+\pi^-$	5.2b	68 B	5.5 B	21B	4.9 B	13 B
$\rho' \rightarrow \pi^+\pi^-\pi^+\pi^-$	730 mb	9.5 B	210 M	2.5 B	190 M	1.2 B
$\phi \rightarrow K^+K^-$	0.22b	2.9 B	82 M	490 M	15 M	330 M
$J/\psi \rightarrow \mu^+\mu^-$	1.0 mb	14 M	1.1 M	5.7 M	600 K	1.6 M
$\psi(2S) \rightarrow \mu^+\mu^-$	30 μb	400 K	35 K	180 K	19 K	47 K
$Y(1S) \rightarrow \mu^+\mu^-$	2.0 μb	26 K	2.8 K	14 K	880	2.0 K

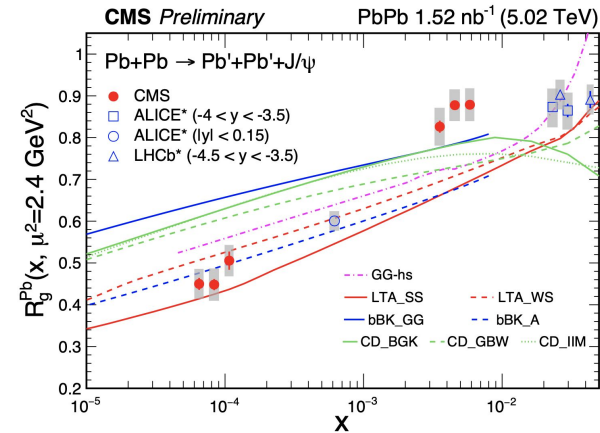
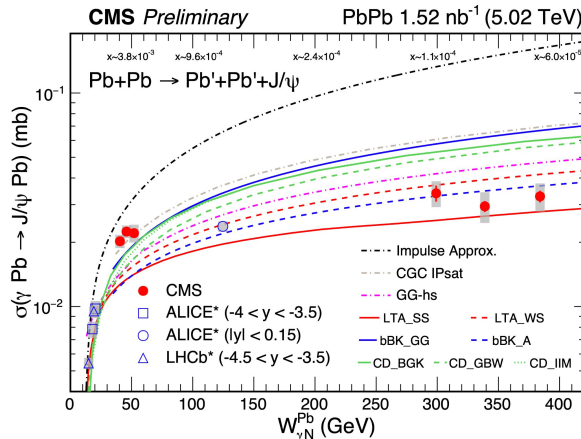
Future opportunities



- Variation of saturation scales in search for gluon saturation.
- When approaching the BDL:
 - Coh. cross section scales with $A^{2/3}$
 - Incoh. cross section strongly suppressed, internal substructure becomes invisible

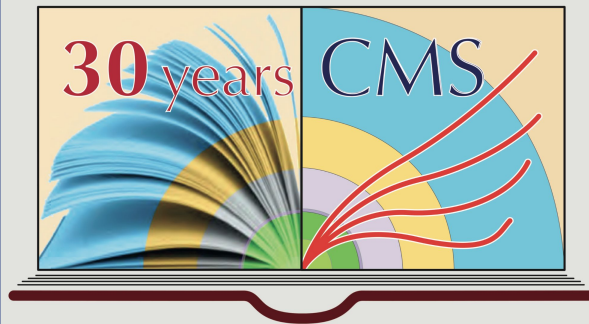
Summary

- For the first time, **disentangled the low and high y energy** contributions to coh. J/Ψ .
- CMS measured coh. J/Ψ at a **new unprecedentedly low- x gluon regime** (10^{-4} - 10^{-5}).
- **Flattening** of coh. $\sigma(J/\Psi)$ vs. W not predicted by state of the art models
 - **Gluon saturation?** or **black disk limit?** or other physic effects?
- HL-LHC including CMS Phase-2 upgrades will bring new exciting opportunities.



CMS-PAS-HIN-22-002

CMS 30 years anniversary



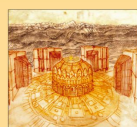
The Compact Muon Solenoid
Technical Proposal



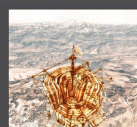
The Magnet System
Technical Design Report



The Hadron Calorimeter Project
Technical Design Report



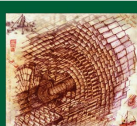
The Phase 2 Upgrade of the
CMS Endcap Calorimeter
Technical Design Report



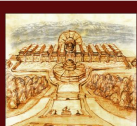
The Compact Muon Solenoid
Phase II Upgrade
Technical proposal



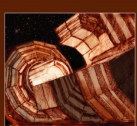
CMS TECHNICAL DESIGN REPORT
Muon Endcap Upgrade
GEM V1 - The Station 1 GEM Project



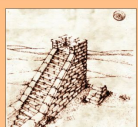
The Tracker Project
Technical Design Report



The Phase 3 Upgrade of the
CMS Beam Radiation Instrumentation
and High Level Trigger
Technical Design Report



The Phase 2 Upgrade of the
CMS Beam Radiation Instrumentation
and High Level Trigger
Technical Design Report



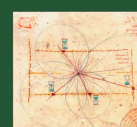
Upgrade of CMS detector
through 2020
Technical Proposal



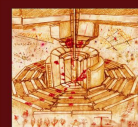
TECHNICAL DESIGN REPORT FOR
CMS-TOTEM
PRECISION PROTON SPECTROMETER



The Electromagnetic
Calorimeter Project
Technical Design Report



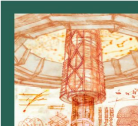
A MIP Timing Detector
for the CMS Phase-1 Upgrade
Technical Design Report



The Phase-2 Upgrade of the
CMS Level-1 Trigger
Technical Design Report



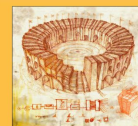
The Muon System Project
Technical Design Report



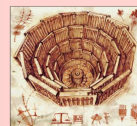
The Phase-2 Upgrade of the
CMS Tracker
Technical Design Report



CMS TECHNICAL DESIGN REPORT
FOR THE PIXEL DETECTOR UPGRADE



CMS TECHNICAL DESIGN REPORT
FOR THE PHASE 1 UPGRADE
OF THE HADRON CALORIMETER



CMS TECHNICAL DESIGN REPORT
FOR THE LEVEL-1 TRIGGER UPGRADE



Software and Detector
Performance
Physics Technical Design Report Vol 1



The Phase-2 Upgrade of the
CMS Muon Detectors
Technical Design Report



The Phase 2 Upgrade of the
CMS Barrel Calorimeters
Technical Design Report



Physics Performance
Physics Technical Design Report Vol II



The Phase-2 Upgrade of the
CMS Muon Detectors
Technical Design Report



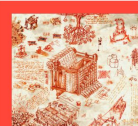
The Phase 2 Upgrade of the
CMS Barrel Calorimeters
Technical Design Report



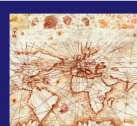
High Density QCD
with Heavy Ions
Physics Technical Design Report Addendum 1



The Phase 2 Upgrade of the
CMS Muon Detectors
Technical Design Report



The Phase 2 Upgrade of the
CMS Muon Detectors
Technical Design Report



The Computing thing
Technical Design Report

Thank you for your attention!

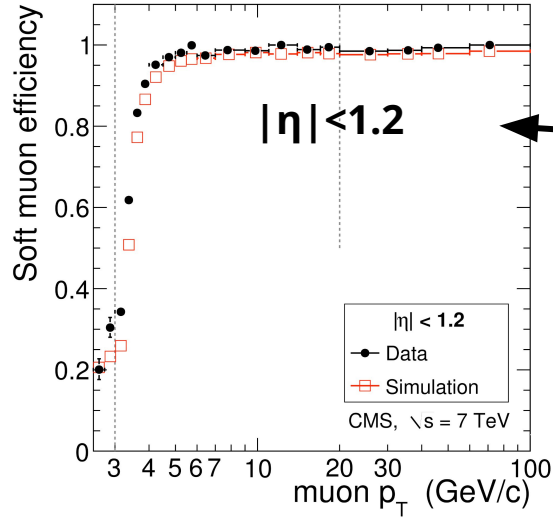


BACKUP

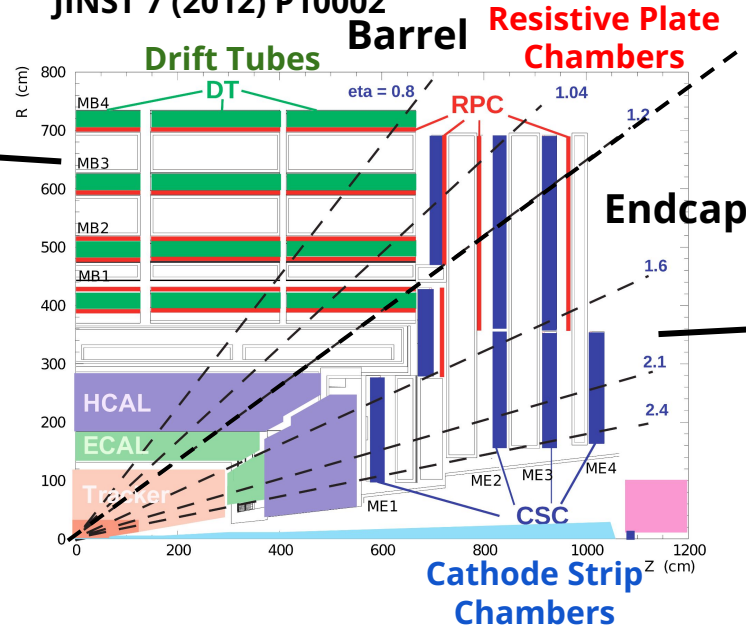


Muon reconstruction

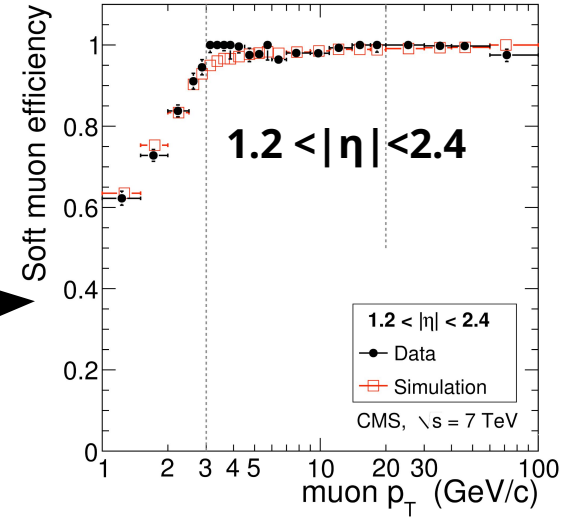
Muon efficiency



JINST 7 (2012) P10002

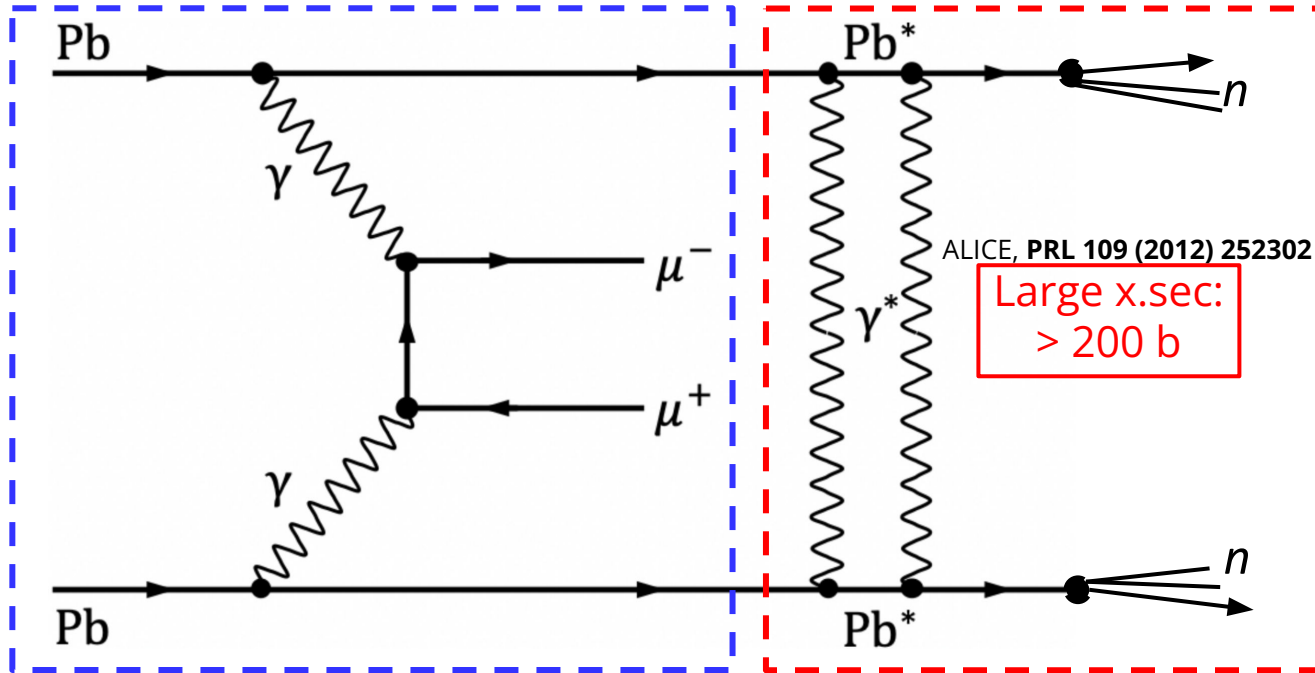


Muon efficiency



- Tracker and muon detectors used to reconstruct/identify muons.
- CMS able to reconstruct muons down to muon $p_T \sim 1$ GeV in forward region.

EM dissociative pileup correction



Impact of dissociative PU corrected by measuring neutron multiplicity in events without any activity in CMS tracker.

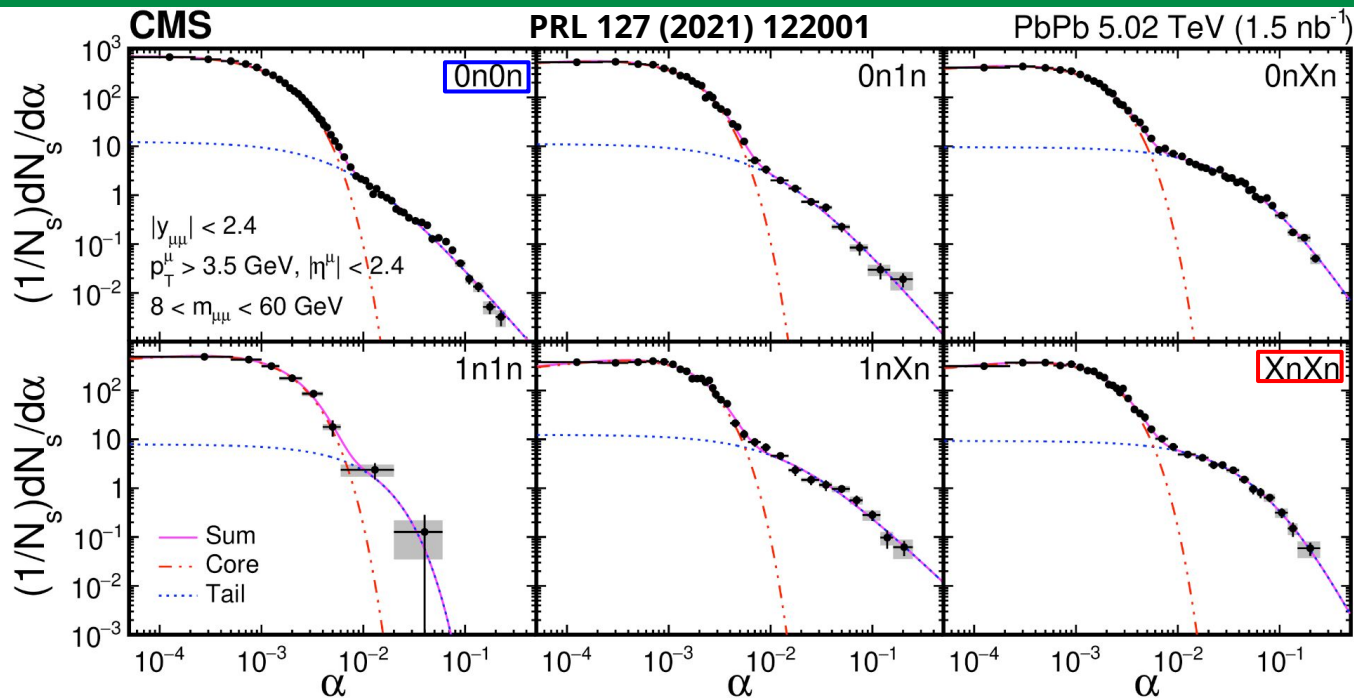
$\gamma\gamma \rightarrow \mu^+\mu^-$ with/without neutron emitting

different collisions

EM dissociation without any $\gamma\gamma \rightarrow \mu^+\mu^-$

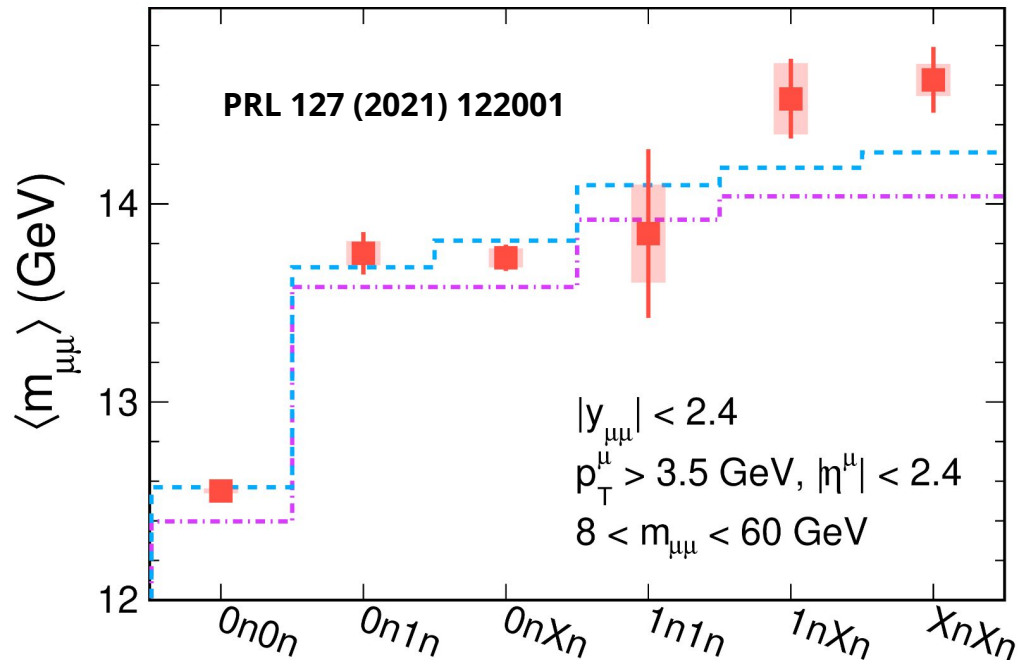
$\gamma\gamma \rightarrow \mu^+\mu^-$ with neutron multiplicity migration

α spectrum vs. neutron multiplicity



- **0n0n (fewer neutrons) → XnXn (more neutrons)**
 - Tail contribution becomes larger.
 - Seems has depletion in the very small α .

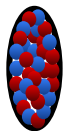
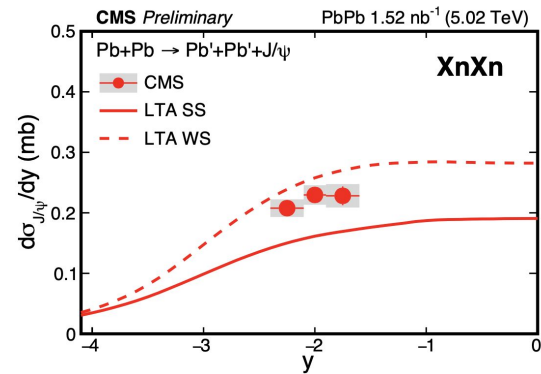
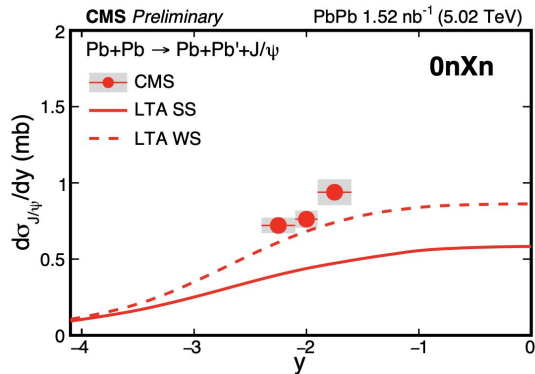
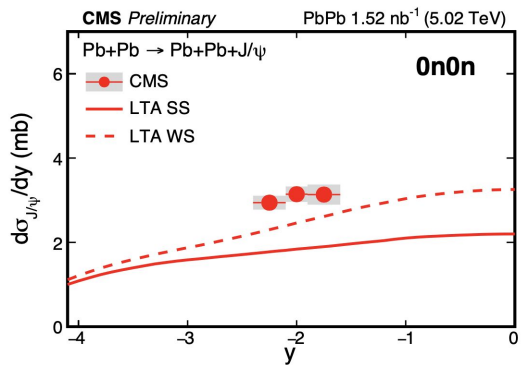
$\langle m_{\mu\mu} \rangle$ vs. neutron multiplicity



- Strong neutron multiplicity dependence of $\langle m_{\mu\mu} \rangle$
 - Deviation from constant $\gg 5\sigma$
 - b dependence of initial photon energy.

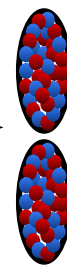
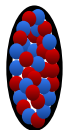
Coherent Jpsi in 0n0n, 0nXn, XnXn

CMS-PAS-HIN-22-002



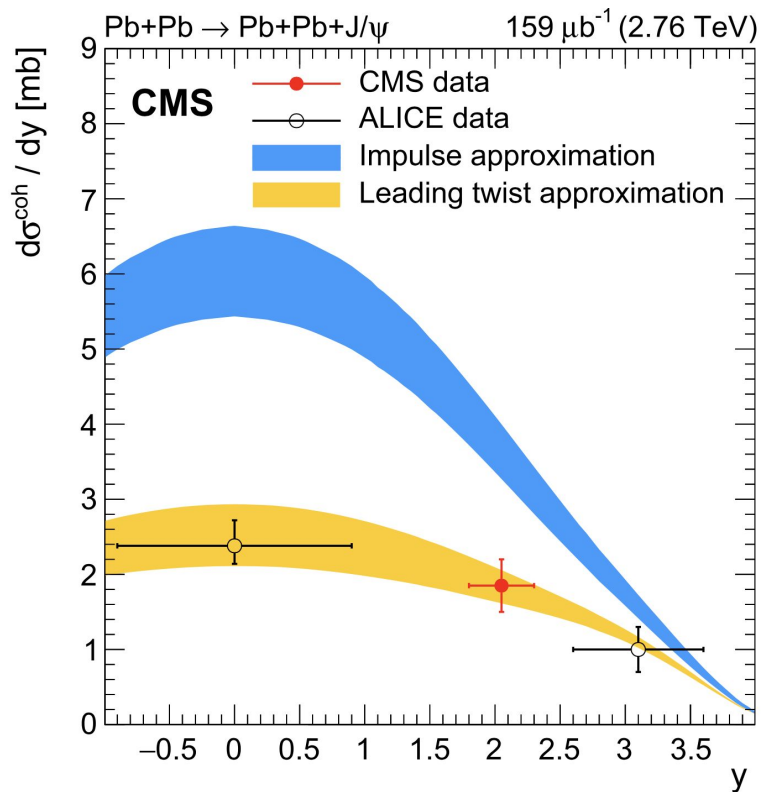
Fewer neutrons

More neutrons



Coherent Jpsi in Run 1

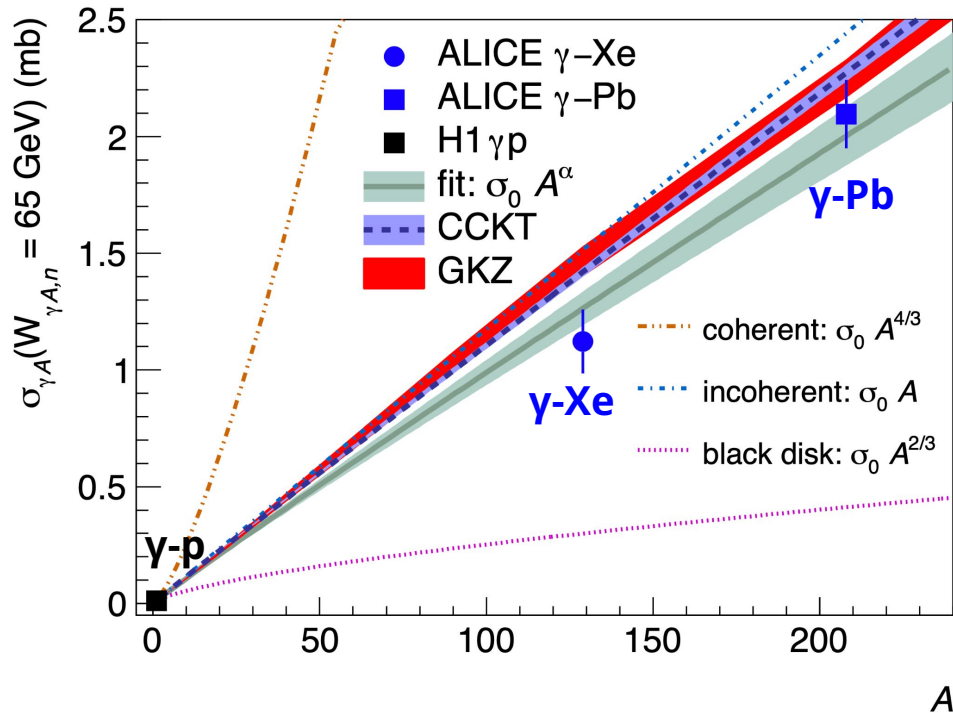
PLB 772 (2017) 489



- Run 1 data from CMS and ALICE well consistent with LTA model calculations
- Large uncertainties and wide y bins

ALICE UPC ρ vs system size

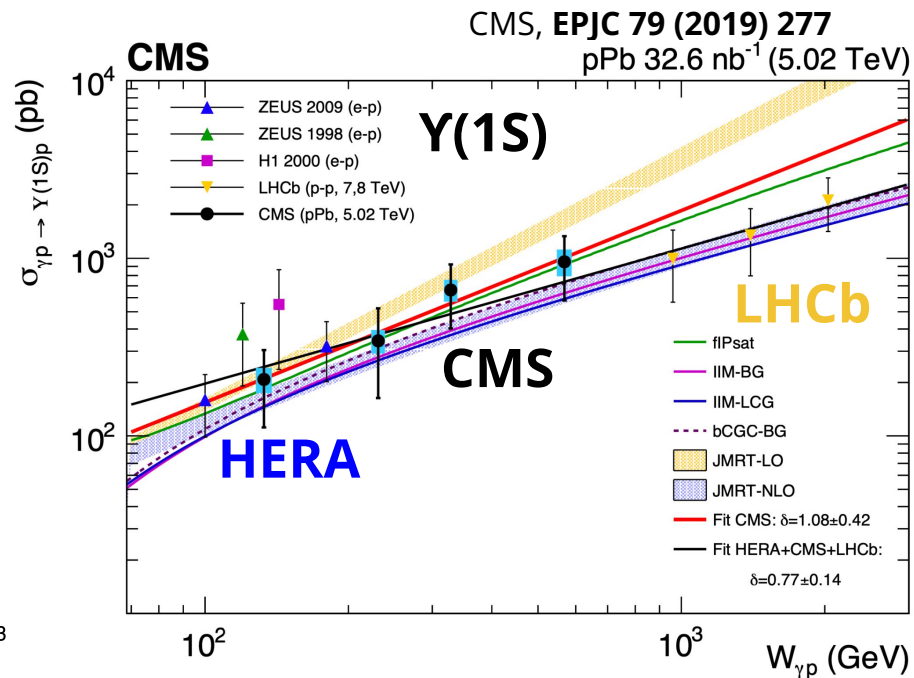
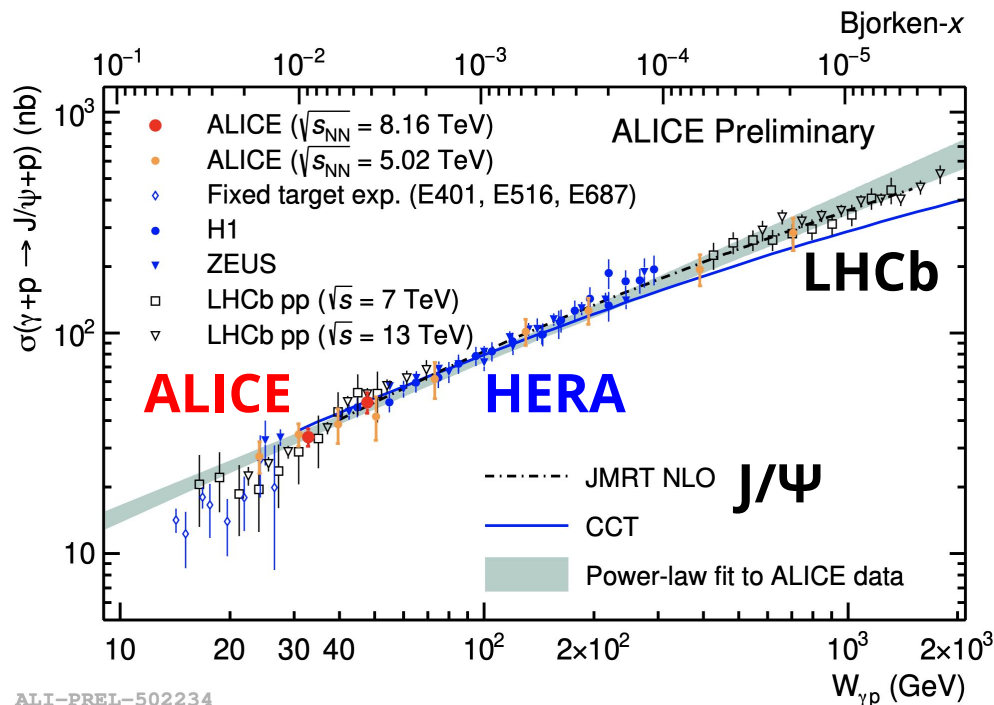
ALICE, PLB 820 (2021) 136481



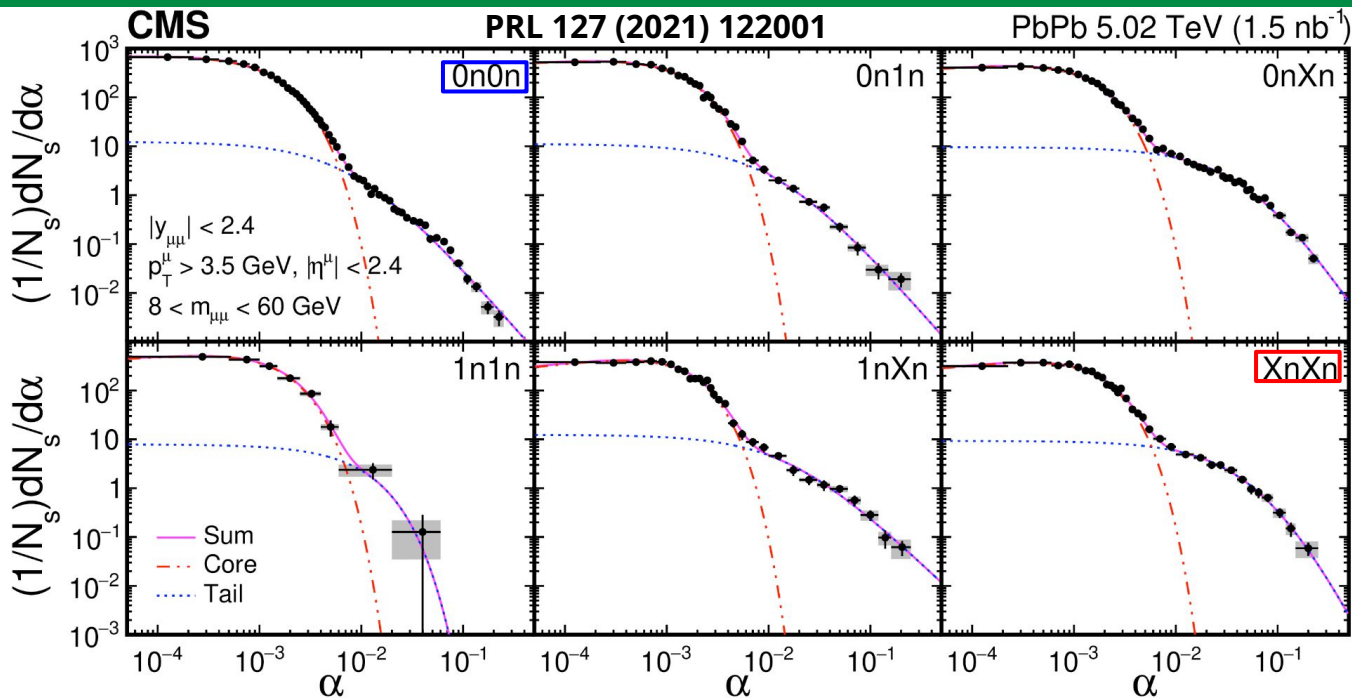
If J/Ψ -nucleus approaches BDL, why ρ -Nucleus does not?

- With A decrease, it is harder to reach BDL \rightarrow the direct $A^{2/3}$ cannot scale to small A .
- Relation of dipole size vs. M in seen by nucleus is different to what seen by nucleon?

Quarkonium photoproduction in γ -p

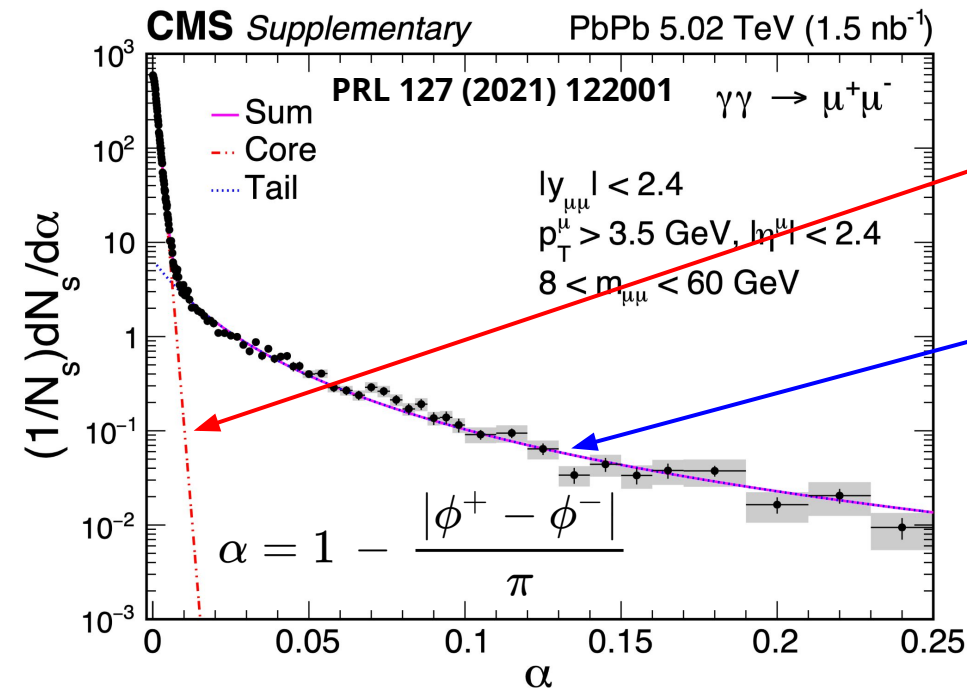


α spectrum vs. neutron multiplicity class



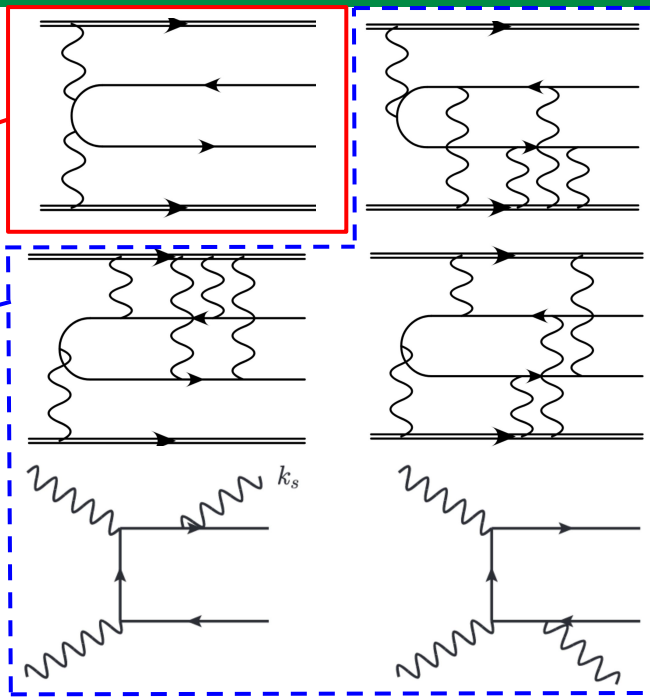
- **0n0n (fewer neutrons) → XnXn (more neutrons)**
 - Tail contribution becomes larger.
 - Seems has depletion in the very small α .

Decomposing LO and HO in α spectrum



Leading Order

Higher Order



Zero Degree Calorimeter

