Constraints on muon spin force from co-magnetometer experiments

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Spin-mass coupling

• Spin-mass coupling mediated by light CP-violating scalar:

$$\mathscr{L} = g_s \phi \,\bar{\psi}_1 \psi_1 + \frac{g_p}{m} (\partial_\alpha \phi) \,\bar{\psi}_2 \gamma^\alpha \gamma_5 \psi_2 \qquad \qquad V(r) = \frac{g_s g_p}{8\pi m} \left(\frac{m_\phi}{r} + \frac{1}{r^2}\right) e^{-m_\phi r} \,\vec{\sigma}_2 \cdot \hat{r}$$

• Acts as "pseudo" magnetic field:



Muon spin-mass coupling

• Muon spin-mass coupling: additional spin rotation to storage ring experiment.

$$\mathscr{L} = g_s \phi \bar{N} N + \frac{g_p}{m_\mu} \partial_\alpha \phi \,\bar{\mu} \gamma^\alpha \gamma_5 \mu$$

$$\overrightarrow{\Omega} = \frac{e}{m} \left[a \vec{B} - \left(a - \frac{1}{\gamma^2 - 1} \right) \vec{\beta} \times \vec{E} + \frac{1}{\gamma} \frac{m}{e} \vec{b} \right], \quad b = \frac{g_p |\vec{\nabla} \phi|}{m_\mu}.$$

• Muon g-2 may be explained if $b = \gamma \frac{eB}{m_{\mu}} \Delta a_{\mu} \sim 120 \,\text{Hz}$. [Davoudiasl+ 22; Agrawal+ 22]

(Effect would be bigger at J-PARC due to smaller γ .)

A way of constraining this coupling in the laboratory?

Muons at loop level

• By closing muon loops, several different spin-mass couplings are generated.



• These couplings probed by the co-magnetometer experiments.

Co-magnetometer experiments

• Co-magnetometer technique is key to suppress noise from the magnetic field.

$$\mathcal{\mathcal{H}} = \overrightarrow{\sigma}_{N} \cdot \left(\mu_{N} \overrightarrow{B} + \chi_{N} \overrightarrow{g} \right) \qquad \overrightarrow{\Omega}_{N_{1}} - \frac{\mu_{N_{1}}}{\mu_{N_{2}}} \overrightarrow{\Omega}_{N_{2}} = \overrightarrow{g} \left(\chi_{N_{1}} - \frac{\mu_{N_{1}}}{\mu_{N_{2}}} \chi_{N_{2}} \right).$$

(Evaluation of χ_N requires the nuclear physics input.)

• Recent 129 Xe - 131 Xe result: [Zhang+ 23]

$$|(\chi_{129} + 1.1245 \chi_{131}) \overrightarrow{g}| < 8.5 \times 10^{-23} \,\mathrm{eV}.$$

• If we use the Schmidt model to evaluate χ_N , we obtain

(The lattice result [FLAG Review 2021] used to go from quark to nucleon spin.)

$$\begin{cases} \rho_N \vec{\nabla} \phi \cdot \vec{B}_N \\ \vec{\nabla} \phi \cdot \vec{\sigma}_N \end{cases} \qquad b < 0.44 \text{ Hz}, \\ b < 0.18 \text{ Hz}. \end{cases}$$

This is to be compared with $b \simeq 120 \,\text{Hz}$ for muon g-2.

Caveat on nuclear physics

• The Schmidt model does not reproduce the nuclear magnetic dipole moments.

(still usually applied to extract constraints in the spin-mass coupling experiments.)

$$\begin{cases} \mu_{129}^{\text{(Schmidt)}} = \mu_n, & \mu_{129}^{\text{(exp)}} = 0.41 \,\mu_n, \\ \mu_{131}^{\text{(Schmidt)}} = -0.6 \,\mu_n, & \mu_{131}^{\text{(exp)}} = -0.36 \,\mu_n. \end{cases}$$

• Different nuclear computations give different results:

[Klos+ 13] $\rightarrow b < 1.2$ Hz, [Ressell, Dean 97] (Bonn A) $\rightarrow b < 0.56$ Hz,

[Ressell, Dean 97] (Nijmegen II) $\rightarrow b < 0.49 \,\text{Hz}$, (from nucleon spin operator.)

• Even a possibility of perfect cancellation, $\chi_{129} = \frac{\mu_{129}}{\mu_{131}}\chi_{131}$, not fully excluded.

More precise input from nuclear physics desired.

Summary

• Muon spin-mass coupling motivated by muon g-2 anomaly.



• Muon loops induces operators probed by co-magnetometer experiments.



• Spin-mass coupling explaining muon g-2 anomaly is disfavored.

But

• Nuclear uncertainty exists, more precise nuclear physic input desired.