

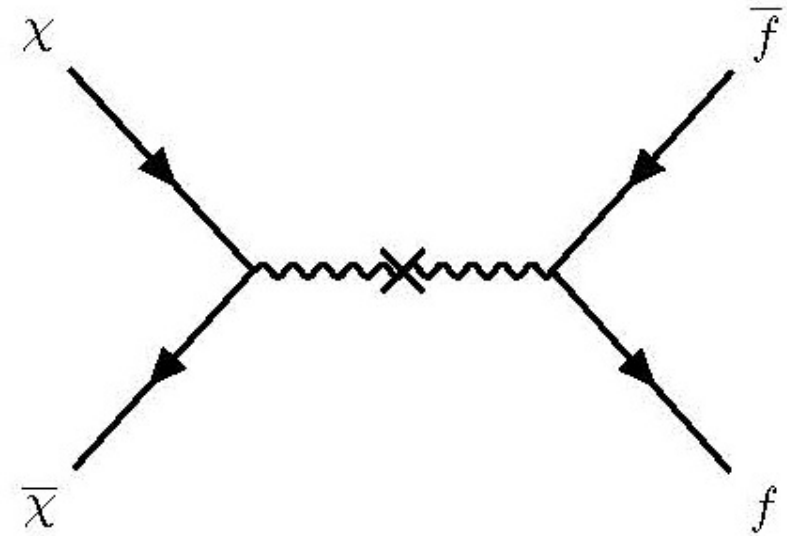
Portal Matter Model Building for g-2

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BASED ON 2211.09918 WITH L. L. EVERETT, S.T. EU, AND R.
XIMENES

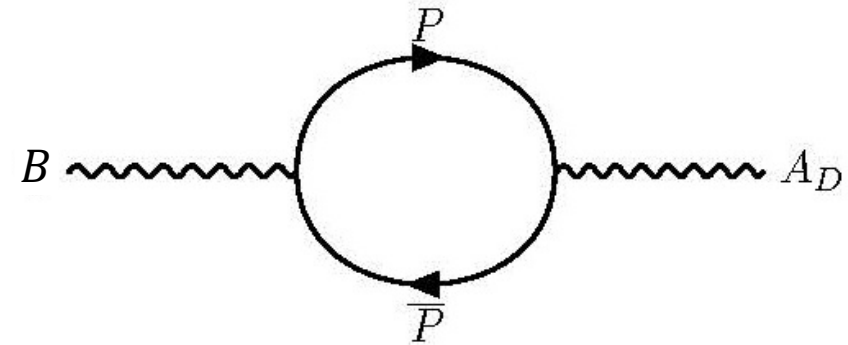
Vector Portal/Kinetic Mixing DM

- One popular idea: Dark Matter (DM) is a particle χ that interacts only with a new “dark force” given by the gauge group $U(1)_D$.
- The Standard Model (SM) is entirely uncharged under $U(1)_D$.
- Interaction between the SM and DM is achieved via kinetic mixing: $\sim \frac{\epsilon}{2 c_W} B_{\mu\nu} A_D^{\mu\nu}$.
- So, SM now couples with strength $\sim \epsilon e Q$ to A_D .
- $m_{DM}, m_{A_D} \sim 0.1 - 1$ GeV, $\epsilon \sim 10^{-(3-5)}$ reproduces the correct relic abundance without running afoul of other experimental constraints.
- Portal Matter: How do we generate ϵ ?



Portal Matter: Origins of ϵ

- The minimal setup: DM and a dark photon A_D . The small parameter ϵ is added by hand.
- By asking where ϵ comes from, can we get a window into higher-energy physics?
- A natural source for $\epsilon \sim 10^{-(3-5)}$ would be *portal matter*: Heavy particles charged under both $U(1)_D$ and SM hypercharge.^{1,2} Call the particle P , for portal matter.
- To get *finite* and *calculable* ϵ , we need $\sum Q_{Y_i} Q_{D_i} = 0$ to eliminate dependence on the renormalization scale μ .



$$\epsilon \propto \sum_i Q_{Y_i} Q_{D_i} \log \left(\frac{m_i^2}{\mu^2} \right)$$

¹B. HOLDOM, PHYS. LETT. B **166**, 196 (1986)

²B. HOLDOM, PHYS. LETT. B **178**, 65 (1986)

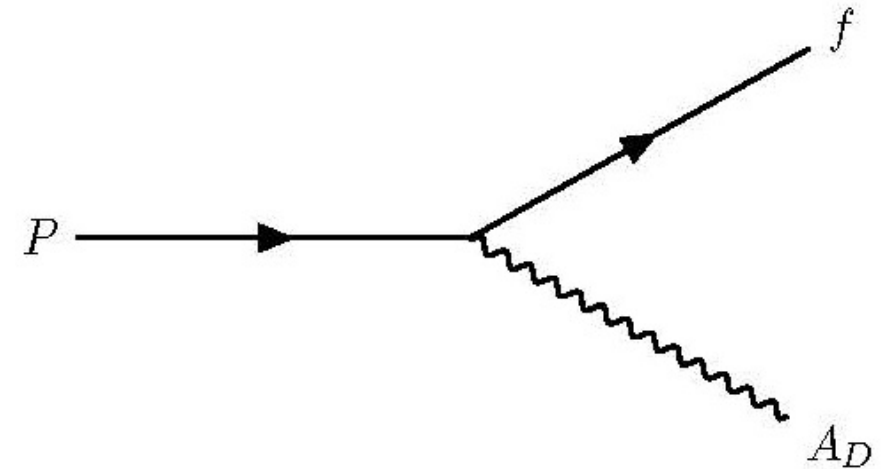
Portal Matter Theory 101

- What can we say about the portal matter? For this talk, we'll stick to fermionic portal matter.
 - 1) Precision electroweak constraints, $H \rightarrow gg, \gamma\gamma$ branching ratio $\rightarrow P$ is **vector-like**.
 - 2) If it's light enough to be thermally produced after reheating, it needs to decay before BBN
- How to get it to decay? Decay has to violate $U(1)_D$, and therefore should come from coupling to the scalar(s) that break that symmetry.
- Simplest path: Portal Matter has the same SM quantum numbers as some SM fermion and mixes with it (easy)



Portal Matter Phenomenology 101

- So we're looking for a portal matter ensemble that satisfies:
 - 1) All P are vector-like copies of SM fermions, but with $U(1)_D$ charge.
 - 2) Ensemble satisfies $\sum Q_{Y_i} Q_{D_i} = 0$
- **Highly suppressed** decay via regular vector-like fermion decay channels, e.g. $P \rightarrow f Z$, instead dominantly decays via $P \rightarrow f A_D, h_D$.
- Collider signature is then a **highly boosted** SM fermion (or jet) plus A_D or h_D .
- A_D, h_D signatures are model-dependent: May decay invisibly (to dark matter) or visibly (to light SM fermions)
- Leads to atypical signatures (e.g., displaced lepton-jets, slepton-like signatures).^{1,2}



¹T. G. RIZZO, PHYS. REV D **99**, NO.11, 115024 (2019) [ARXIV:1810.07531 [HEP-PH]]

²T. D. REUTER AND T. G. RIZZO, PHYS. REV. D **101**, NO.1, 015014 (2020) [ARXIV:1909.09160 [HEP-PH]]

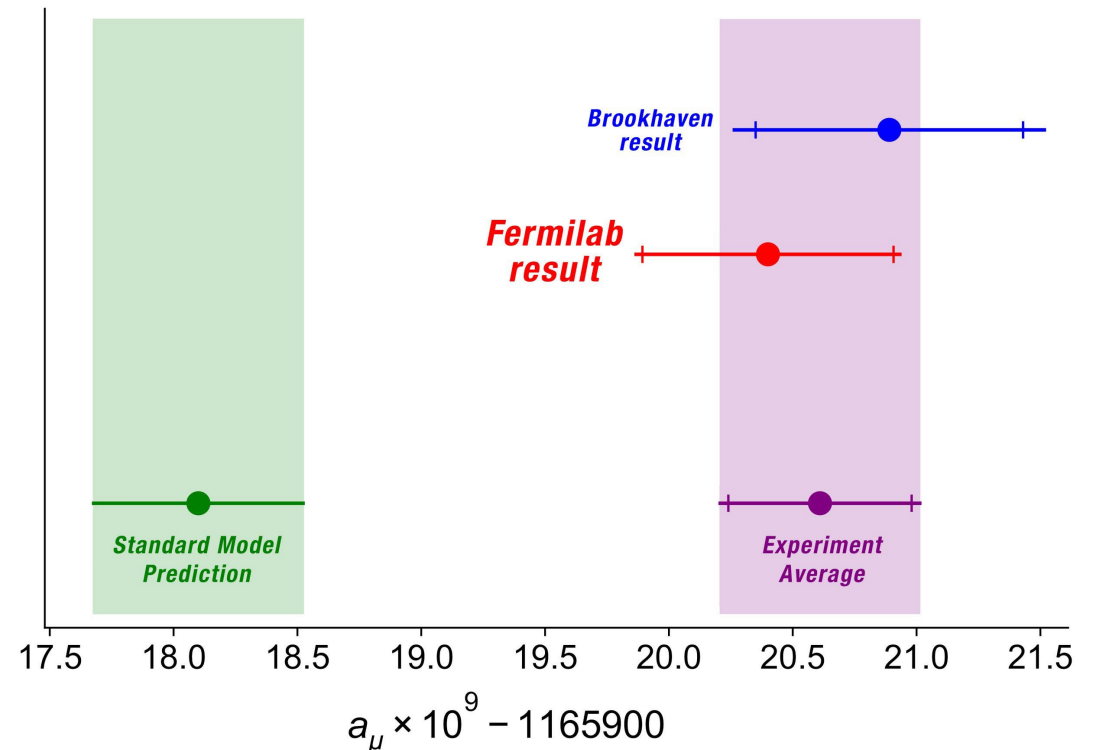
Muon $g - 2$

- An obvious question: Can portal matter address other mysteries in physics? Let's focus on one: **Muon anomalous magnetic moment**

- 4.2 σ discrepancy between SM expectation and experiment^{1,2,3}:

$$\Delta a_\mu = a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = (251 \pm 59) \times 10^{-11}$$

- Why? Could be hadronic vacuum polarization is wrong⁴
- Could be **new physics**



¹B. ABI *ET AL.* (MUON G-2), *PHYS. REV. LETT.* **126**, 141801 (2021), ARXIV:2104.03281 [HEP-EX]

²G. W. BENNETT *ET AL.* (MUON G-2), *PHYS. REV. D* **73**, 072003 (2006), ARXIV:HEP-EX/0602035

³T. AOYAMA *ET AL.*, *PHYS. REPT.* **887**, 1 (2020), ARXIV:2006.04822 [HEP-PH]

⁴S. BORSANYI *ET AL.*, *NATURE* **593**, 51 (2021), ARXIV:2002.12347 [HEP-LAT]

The BSM $g - 2$ Cookbook



- Three habits of highly effective $g - 2$ models:
 - 1) Couple something new to the muon, or you won't change its magnetic moment.
 - 2) New physics should have a scale $< O(10 \text{ TeV})$, or your correction will be too small.
 - 3) New physics couples to *both* chiralities μ_L and μ_R .
- Muon-mixed portal matter can do (1) and (2).
- Doesn't do so well at (3)– Only chirality violation is from m_μ , and falls short by an order of magnitude

A Minimal Workable Framework

Field	$SU(2)_L \times U(1)_Y$	Q_D
$\mathbf{l}_L = (v_L^\mu, \mu_L)^T$	$\left(2, -\frac{1}{2}\right)$	0
μ_R	$(1, -1)$	0
$\mathbf{L}_{L,R}^\pm = (N_{L,R}^\pm, L_{L,R}^\pm)^T$	$\left(2, -\frac{1}{2}\right)$	± 1
$E_{L,R}^\pm$	$(1, -1)$	± 1
$S = v_S + (h_D + i \sigma_D)/\sqrt{2}$	$(1, 0)$	+1

$$(\bar{\mu}_L \quad \bar{L}_L^+ \quad \bar{E}_L^+ \quad \dots) \begin{pmatrix} m_\mu & y_{SL}^+ v_S & 0 & & \\ 0 & M_L^+ & e^{i\phi_{LE}^+} \frac{y_{LE}^+}{y_\mu} m_\mu & \vdots & \\ y_{SE}^+ v_S & e^{i\phi_{EL}^+} \frac{y_{EL}^+}{y_\mu} m_\mu & M_E^+ & & \\ & & & \dots & \\ & & & & \ddots \end{pmatrix} \begin{pmatrix} \mu_R \\ L_R^+ \\ E_R^+ \\ \vdots \end{pmatrix}$$

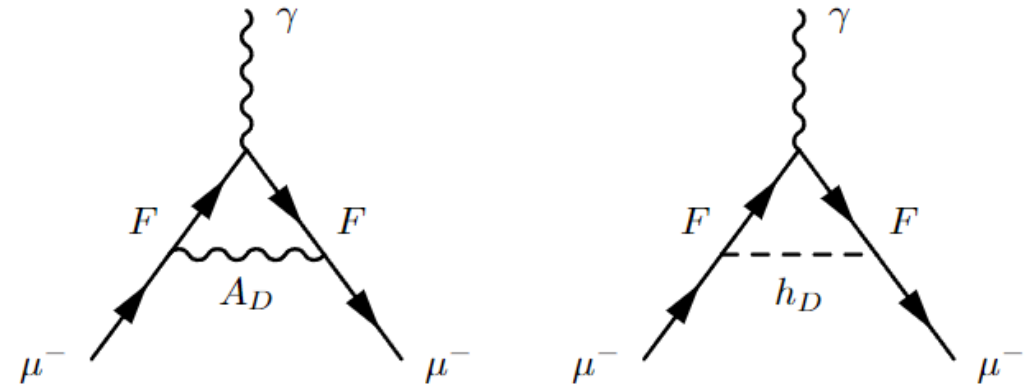
- There are 5 chirality-flipping masses:

$$m_\mu, e^{i\phi_{EL}^\pm} \frac{y_{EL}^\pm}{y_\mu} m_\mu, e^{i\phi_{LE}^\pm} \frac{y_{LE}^\pm}{y_\mu} m_\mu.^1$$

- Naively, $y_{EL,LE}^\pm$ can be $\gg y_\mu \sim 10^{-4}$
- In this basis, all Yukawas are positive, only complex phases are ϕ 's

$g - 2$ in the Minimal Model

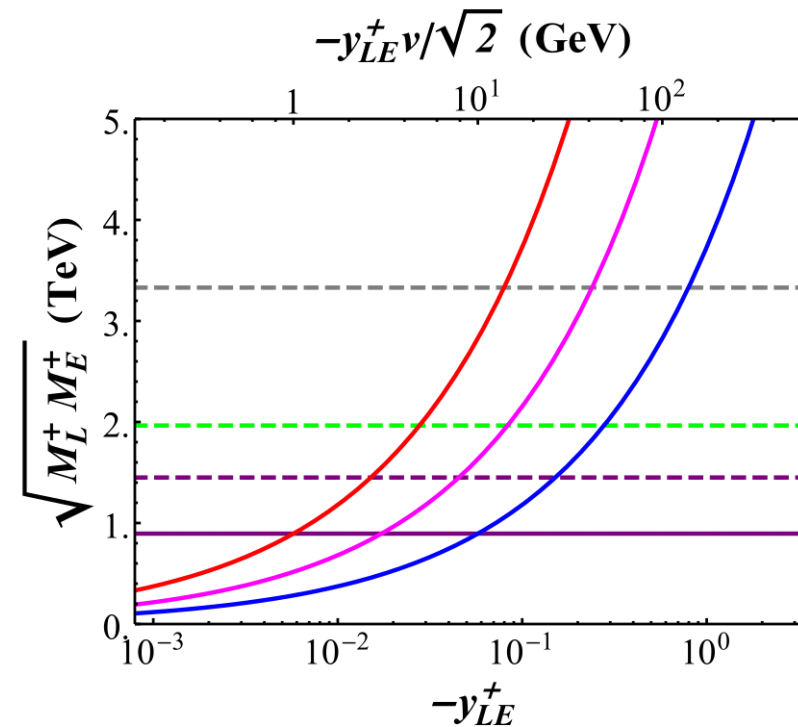
- Observed anomaly is entirely accounted for with $y_{LE}^{\pm}/y_{\mu} \sim O(10)$, $M_{L,E}^{\pm} \sim O(\text{TeV})$
- Light dark Higgs, dark photon means that there is *no* dependence on any parameters that directly enter the simplified dark matter model— this is entirely a portal matter effect.
- The physical phases ϕ_{LE}^{\pm} are critical to generating the correct sign correction here



$$\Delta a_{\mu} \approx -\frac{m_{\mu}^2}{16 \pi^2} \sum_{i=+,-} \frac{y_{SL}^i y_{SE}^i}{M_L^i M_E^i} \frac{y_{LE}^i}{y_{\mu}} \cos \phi_{LE}^i$$

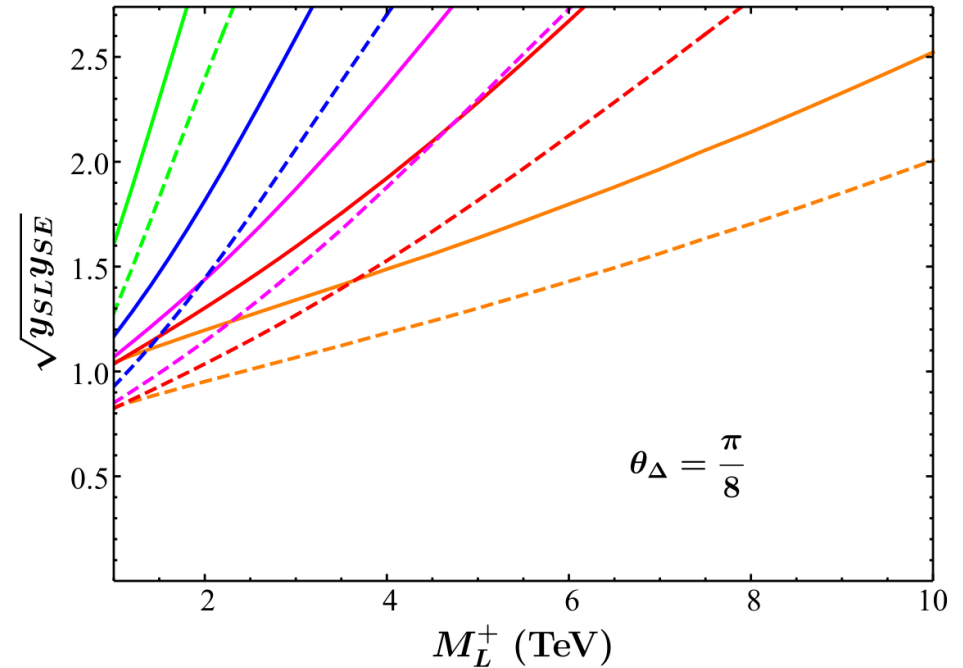
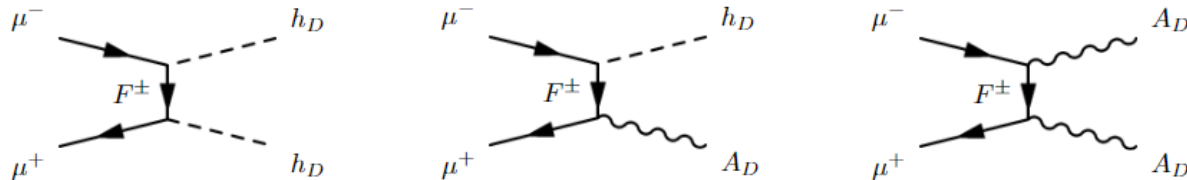
Portal Matter Collider Production

- First step in constraining the model: Limits on portal matter masses.
- On the right, we have the $\sqrt{M_L^+ M_E^+}$ needed to recreate Δa_μ , plus LHC constraints on M_E^+ , as well as HL-LHC, HE-LHC, hh-FCC¹
- $y_{LE}^\pm, y_{SL,SE}^\pm$ can only be constrained (at LHC) by perturbative unitarity.
- Multi-TeV muon collider can do better, probing portal matter masses up to $\lesssim \sqrt{s}/2$ and constraining $y_{SL,SE}^\pm$



Portal Matter Monophoton Signal

- With a muon collider, we could constrain $\Delta a_\mu \propto y_{SL}^\pm y_{SE}^\pm$ through monophoton events
- On the right, estimated discovery (solid) and exclusion (dashed) reach for c.o.m. energies of 3, 6, 10, 14, 30 TeV
- Feasible* to probe region $y_{SL,SE}^\pm \sim O(1)$, which is exactly what we need!



Conclusions

- Portal matter provides guidance to BSM model building, with ready constraints from searches for distinctive vector-like fermions
- Example: Muon $g - 2$ model building led us to a new parameter space
- In the muon $g - 2$ model, parameter space that explains the anomaly can be readily probed in a multi-TeV muon collider

Thank You!
