



20 years of a holographic formula

based on work over the years with H. Dorn, R. Aros, F. Bugini and S. Acevedo

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Aim of the talk

Report a fruitful interplay **AdS/ CFT** correspondence , Conformal
Geometry/Spectral Theory

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- **Holographic formula**: a nostalgic overview

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- **Holographic formula**: a nostalgic overview
- **Gravity as a tool**: one-loop partition functions and **trace/Weyl/conformal anomalies**

Plan

1. Maldacena's conjecture
2. Holographic Weyl anomaly
3. A subleading $O(1)$ result
4. The holographic formula
5. Testing and tweaking
6. Outlook

Maldacena's conjecture

Realization of two deeply-rooted ideas in physics:

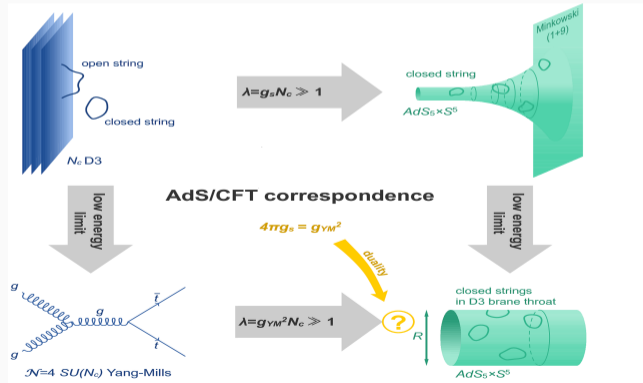
the holographic principle [G. 't Hooft / L. Susskind]

the string of the large-N gauge theory [G. 't Hooft]

Realization of two deeply-rooted ideas in physics:

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String/M-theory: partition function

- $AdS_{d+1} \times X$
- prescribed asymptotics at the conformal infinity

CFT_d : generating functional

- at conformal boundary
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- Most of the initial tests: **class. SUGRA**/ **leading large-N** regimes
- **weak/strong** duality: comparison with perturbative gauge regime only for protected quantities

Holographic Weyl anomaly

Bulk

- reconstruct a Poincare-Einstein metric from a given conformal infty
- $\frac{1}{G_5}(R - \frac{1}{2}R^2) Vol_g$, but in finite volume

? IR-UV connection [Susskind+Witten '98]

Boundary

- 1-loop effective potential (UV)
- QFT in curved spacetime: proper-time, heat-kernel, etc

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? Math peeking around the corner: Q-curvature \int volume renormalization of asympt. hyperbolic manifolds g = ratio of determinants of conf. inv. Laplacians / gen. Polyakov \int g [Branson, Fefferman+Graham, Graham+Zworski, etc]

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$$AdS_5 \times S^5: \quad \hbar T_{ren} \quad i = \frac{c}{8} Ric^2 - \frac{1}{3} R^2 \quad c = \frac{N^2 - 1}{4}$$

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A subleading $O(1)$ result

Beyond classical SUGRA:

$$l_p^4 = L^4 \quad 1 = N$$

- ? $O(N)$: tree-level after inclusion of open or unoriented closed strings
- ? $O(1)$ taking $N^2 \rightarrow N^2 - 1$: loop in SUGRA, but needs whole KK-towers and SUSY [Bilal+Chu' 99, Mansfield+Nozland+Ueno' 02]

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Bulk

- scalar $\frac{d^2}{4} < m^2 < \frac{d^2}{4} + 1$, two AdS-inv. quantizations

[Breitenlohner+Freedman' 82]

- generalized boundary condition \mathcal{F} : the only two conformal inv. choices $\mathcal{F} = 0; 1$

Boundary

- $\mathcal{F} = -$ CFT: same hologram (but different asymptotics) [Klebanov+Witten' 99]
- end points of RG-flow triggered by a relevant double-trace deformation $f \sim O^2$ of the $\mathcal{F} = -$ CFT.

- Background solution has $\epsilon = 0$) no effect on the classical SUGRA partition function
- But two AdS-invariant propagators G) quantum fluctuations of ϕ are sensitive to the boundary conditions (**Casimir effect**)

$$Z_{grav} = Z_{grav}^{class} \det (\square + m^2)^{-1/2}$$

- The ratio Z_{grav}^+ / Z_{grav} only contains the IR-divergence of the infinite AdS volume.

$$e^{-(V^+ - V) Vol(AdS)}$$

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AdS prediction: $O(1)$ correction to the holographic anomaly, polynomial in

Can this correction be reproduced on the boundary? **YES!!!**

- exploit the RG- flow picture: $f \sim O^2$
- Hubbard-Stratonovich transf. (auxiliary field trick)

$$Z = \int D\phi \int d\psi e^{-\int \psi^2 - \int \phi^2 - \int \psi \phi} = \int d\psi e^{-\int \psi^2 - \int \psi^2} = \int d\psi e^{-2\int \psi^2}$$

- large-N factorization

$$\int d\psi e^{-2\int \psi^2} = \int d\psi e^{-\int \psi^2} \int d\psi e^{-\int \psi^2}$$

- uct. det. of the auxiliary field: $\int d\psi e^{-\int \psi^2}$ as $f \sim 1$

$$Z = Z [\det(\dots)]^{-1/2}$$

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CFT confirmation: $O(1)$ correction to the trace anomaly ($d=2,4,6,8$)

The holographic formula

1. Bulk: non-zero $V^+ - V$ for both even and odd d , but for odd d only via numerics [Hartman+Rastelli '06]
2. Confusion: for odd d , no anomaly in CFT vs. nonzero effective potential in the bulk
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$$\text{AdS/CFT } \left(\frac{\det_+ (\dots + m^2)}{\det (\dots + m^2)} \right) \stackrel{?}{=} \det h_{\mathcal{O}} \mathcal{O}_i$$

Testing and tweaking

Unknown to our contemporary conformal geometers? GJMS ops.

Continuation: Δ ! $d=2$ $k (k \geq N)$) think of Ξ as inverse of k -th GJMS
[Graham+Jenne+Mason+Sparling' 91]

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Unknown to our contemporary conformal geometers? GJMS ops.

Continuation: Δ^{-k} ($d=2k$, $k \in \mathbb{N}$) think of Ξ as inverse of k -th GJMS
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d odd

- Analogous result for $d + 1$ even, for a generalized notion of determinant of GJMS [Guillarmou' 05]

d even

- “The delicate case of $d + 1$ odd where things do not renormalize correctly”, is still to be understood!

We anticipated that a proper treatment in $d = \text{even}$ should unveil the **Weyl anomaly**.

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a			
	$k = 1$	$k = 2$	$k = 3$
$n = 2$	2 ✗	-	-
$n = 4$	4=3 ✗	112=3 ✗	-
$n = 6$	10=3 ✗	64=3	738

- A compact formula in terms of Plancherel measure [DD' 08, Dowker' 10]
- The very same numbers found in recent years: log-term of the entanglement entropy for a massless free scalar through an even-sphere [Casi ni +Huerta' 10]

Maybe more important, the two chief roles of the Q-curvature are connected for the first time: **Branson vs. Fefferman+Graham**.

(i) regularized volume in the ambient construction [Graham+Zworski ' 01]

(ii) Polyakov f-Ia for GJMS [Branson' 93] ($d = 2; 4; 6$ and **conjecturally for all even**)

Holographically induced Polyakov formulas:

$$\log \frac{\det P_{2k}^d}{\det P_{2k}} = a \int_M w(\hat{Q} + Q) + \dots$$

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- ? From renormalized volume and its conformal variation under $\mathfrak{g} = e^{2w}g$ [Chang+Qing+Yang' 05], or alternatively, induced action for the conformal mode [Carlip' 05, Aros+Romo+Zamorano' 06], one gets the same structure!

Quotients $X = n \text{ AdS}_{n+1}$ (e.g. thermal AdS and BTZ bh)

[DD' 08, Aros+DD' 09]

Poincare patch

$$ds^2 = \frac{dz^2 + dx^2}{z^2} \text{ identification } (z; \mathbf{x}) \sim e^l(z; A\mathbf{x})$$

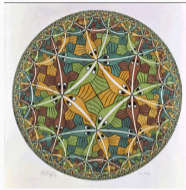
Gravity

- method of images
- in a nutshell, (Patterson-)Selberg zeta function Z [Patterson' 89]

CFT

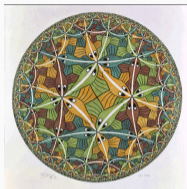
- thermal correlator $\langle hhO O \rangle$ at $T = 1/l; l$
- read off from scattering in X [Perry+Williams' 03]
- stationary Schrödinger in a Poschl-Teller barrier

Quotients



$$\frac{Z_{\Gamma}(n)}{Z_{\Gamma}(\)}^2 \exp(A_n \checkmark)$$

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In particular, for 'resonant' values of the scaling dimension $\neq 2$ in two dimensions:
the celebrated determinant of the Laplacian on the torus
[Ray+Singer '73, Polchinski '86]

Higher spins

Bulk AdS_{odd} [Giombi + Klebanov + Pufu + Safdi + Taronna '13]

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Boundary Einstein [Tseytlin' 13]

- Heat-kernel confirmation of the type-A trace anomaly coefficient
- Heat-kernel derivation of the **type-B trace anomaly** coefficient(s): remained holographically unaccounted for

Higher spins: building blocks

[Acevedo+Aros+Bugini +DD' 17, Aros+Bugini +DD' 19' 21' 22]

$$\frac{\det_{;TT} \left(n \hat{L}^{(s)} + s(n + s - 2) \frac{n^2}{4} + k^2 \right)}{\det_{+;TT} \left(n \hat{L}^{(s)} + s(n + s - 2) \frac{n^2}{4} + k^2 \right)} = \det_{TT} P_{2k}^{(s)} \cdots \det P_{2k}^{(0)}$$

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Three key ingredients:

- Simple holographic recipe to read off type-B anomaly [Bugini+DD'16].
- Extrapolation of b_6 heat coefficient, recently computed by [Liu+McPeak' 18]
- WKB-exactness of the heat-kernel for transverse-traceless totally symmetric rank- s tensors.

Outlook

Summary

- "It is very likely that the holographic formula is right"

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MANY THANKS FOR YOUR ATTENTION

¡GRACIAS!