

NuDM-2022



# Non-standard neutrino interactions in light mediator models at reactor experiments

Sumit Ghosh  
Korea Institute for Advanced Study  
[ghosh@kias.re.kr](mailto:ghosh@kias.re.kr)

*w/ B. Dutta, T. Li, A. Thompson and A. Verma*

# MOTIVATIONS

- Neutrino mass and oscillations is a well known phenomenon. However, the neutrino mass generation mechanism remains unknown.
- So far, no experimental evidence for NSI of neutrinos, but if exist, directly indicate new physics beyond SM- well motivated from a phenomenological point of view.
- Recent interest in low-scale neutrino models containing new weakly coupled light mediating particles. Experiments such as GEMMA, Borexino, XENONnT constraint the parameter space.
- Ongoing/future reactor based neutrino experiments such as MINER, CONUS, CONNIE can investigate these low-scale models.

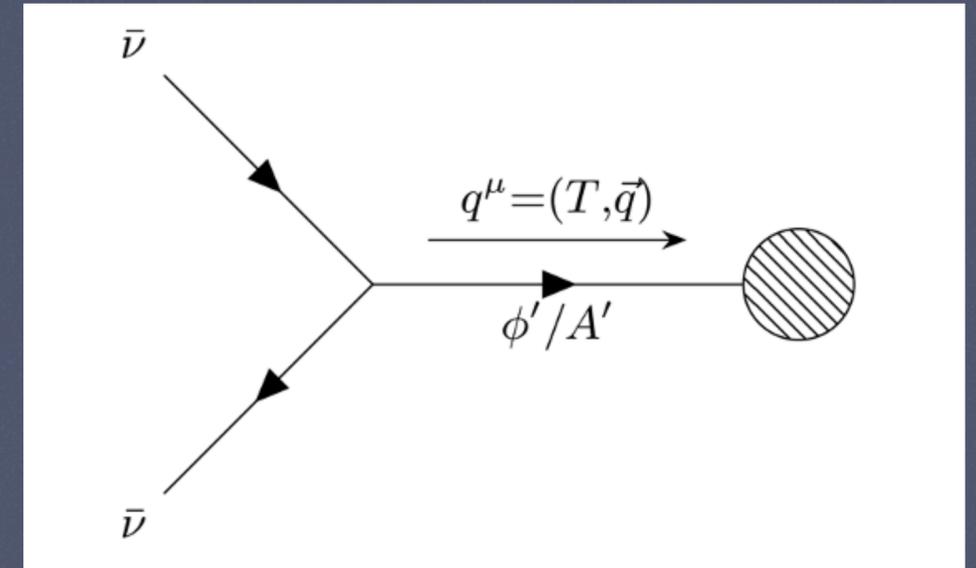
# LOW-SCALE NSI MODELS AT REACTOR EXPERIMENTS

- Non-standard neutrino-electron interactions induced by light scalar/vector mediator arise naturally in many low-scale models.
- Our strategy → investigate these low-scale models in reactor based neutrino experiment with low threshold Ge/Si detectors and find the prospect of probing/ruling out such models.

# NON-STANDARD NEUTRINO-ELECTRON SCATTERING

- Scalar NSI:  $\mathcal{L}_S \supset g_{e,\phi} \bar{e}e\phi + g_{\nu,\phi} \bar{\nu}\nu\phi$

$$\frac{d\sigma_e}{dT} - \frac{d\sigma_e^{\text{SM}}}{dT} = \frac{g_{\nu,\phi}^2 g_{e,\phi}^2 T m_e^2}{4\pi E_\nu^2 (2Tm_e + m_\phi^2)^2}$$



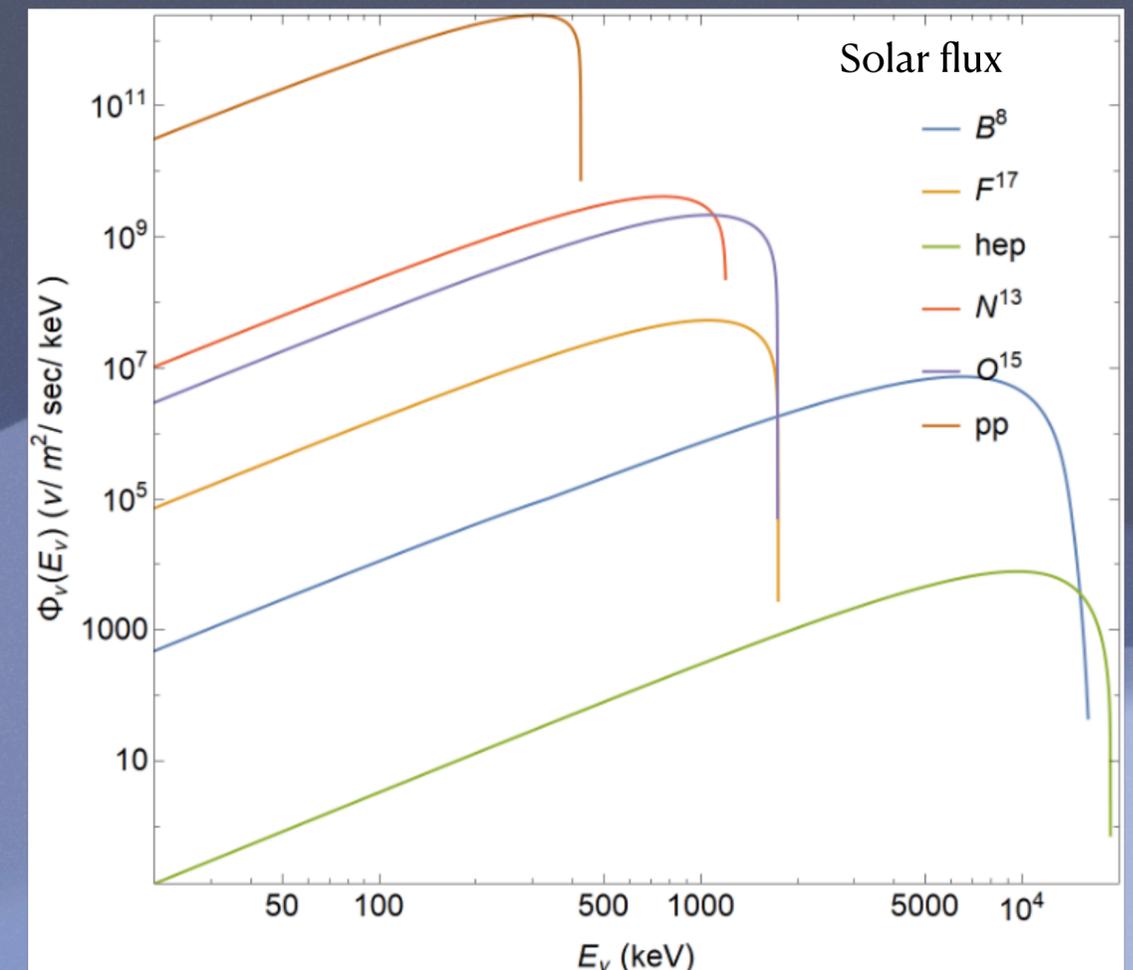
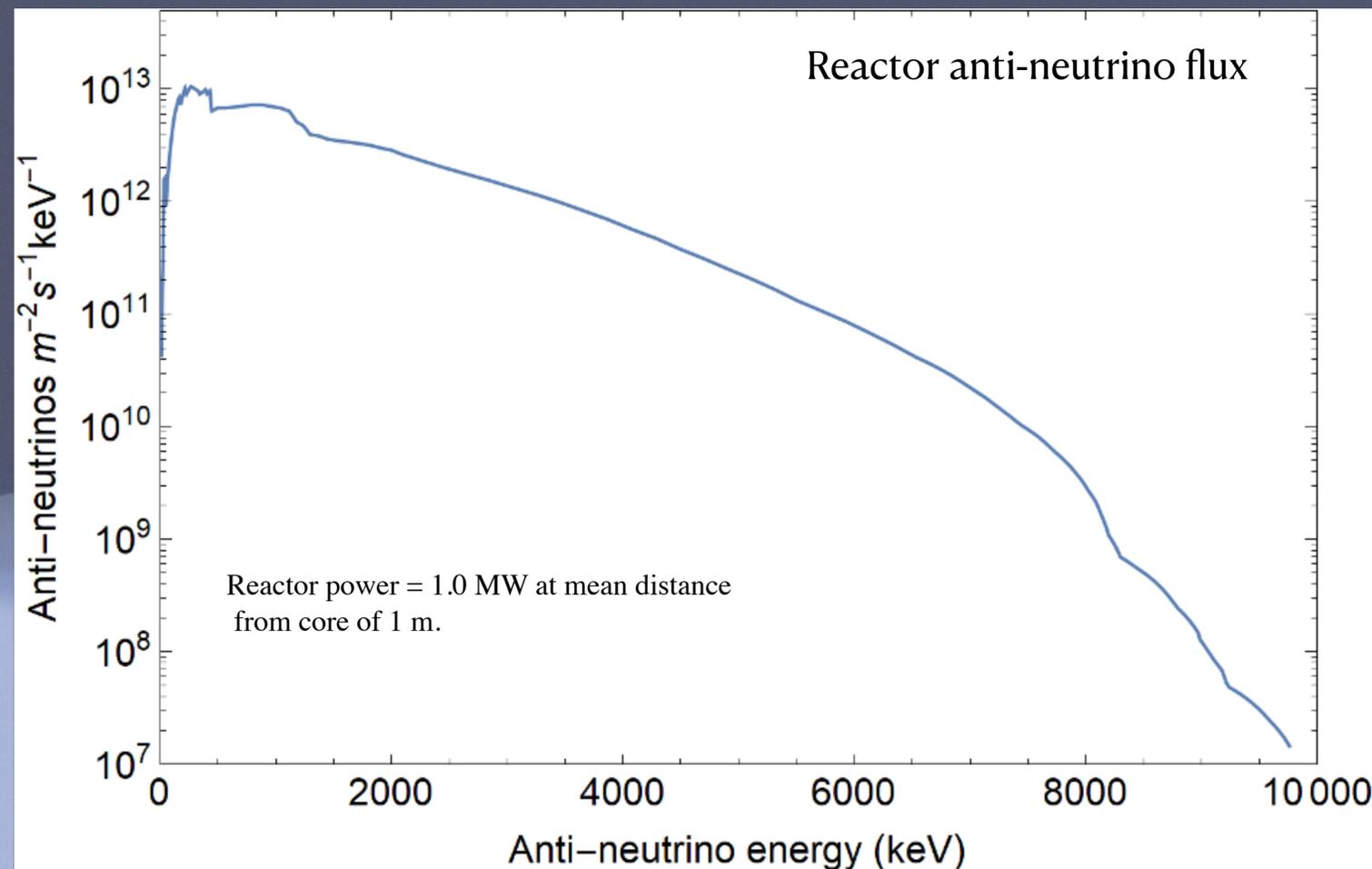
- Vector NSI:  $\mathcal{L}_V \supset g_{e,A'} \bar{e}\gamma^\rho e A'_\rho + g_{\nu,A'} \bar{\nu}_L \gamma^\rho \nu_L A'_\rho$

$$\frac{d\sigma_e}{dT} - \frac{d\sigma_e^{\text{SM}}}{dT} = \frac{\sqrt{2} G_F m_e g_\nu g_{\nu,A'} g_{e,A'}}{\pi (2m_e T + m_{A'}^2)} + \frac{m_e g_{\nu,A'}^2 g_{e,A'}^2}{2\pi (2m_e T + m_{A'}^2)^2}$$

$$T_{\text{max}} = \frac{2E_\nu^2}{(m_N + 2E_\nu)}$$

# REACTOR VS SOLAR NEUTRINO FLUX

- MW reactor has similar energy profile to the solar neutrino flux with characteristic energy  $\leq 1$  MeV.
- Reactor antineutrino flux gets peaked around 200-300 keV and gives at least one order of magnitude more flux compared to solar flux at these keV energy scale.



# NEUTRINO-ELECTRON SCATTERING CROSS-SECTION

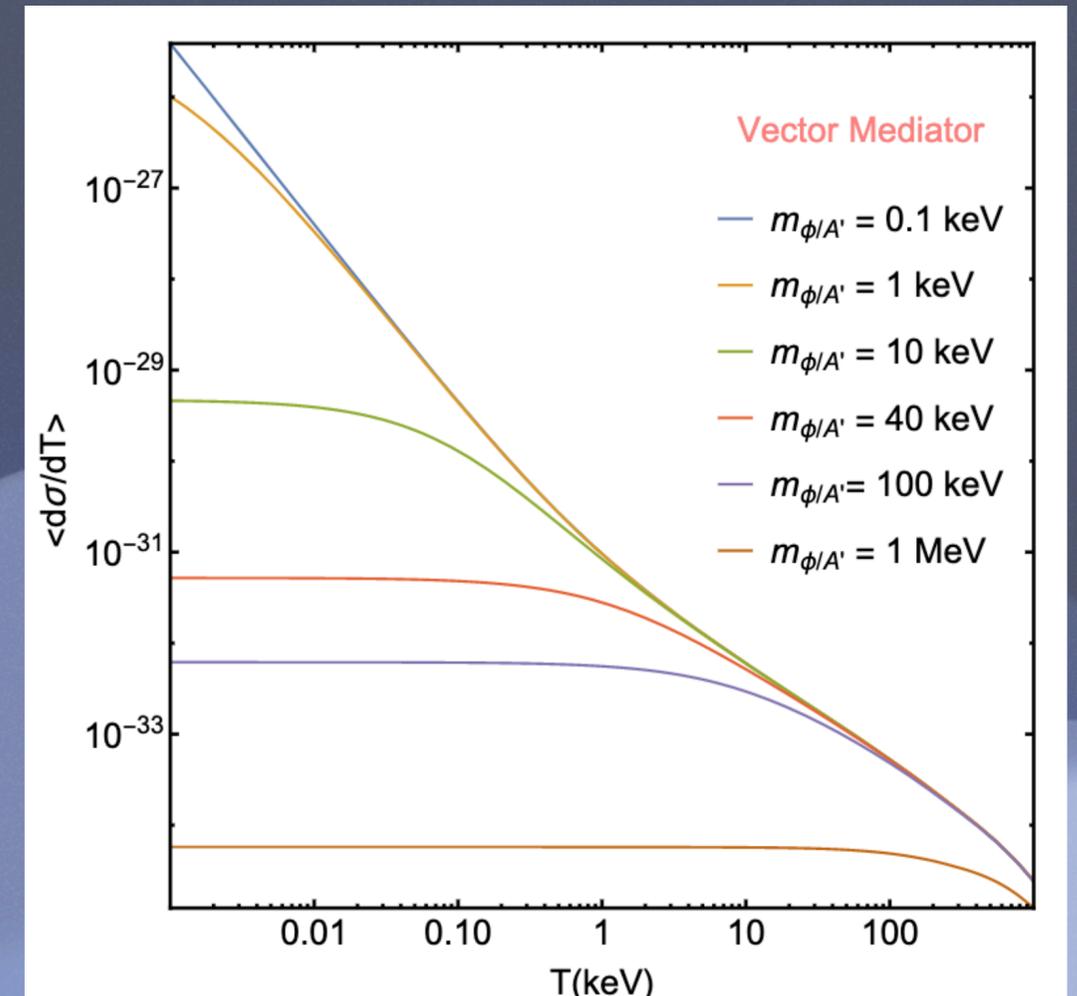
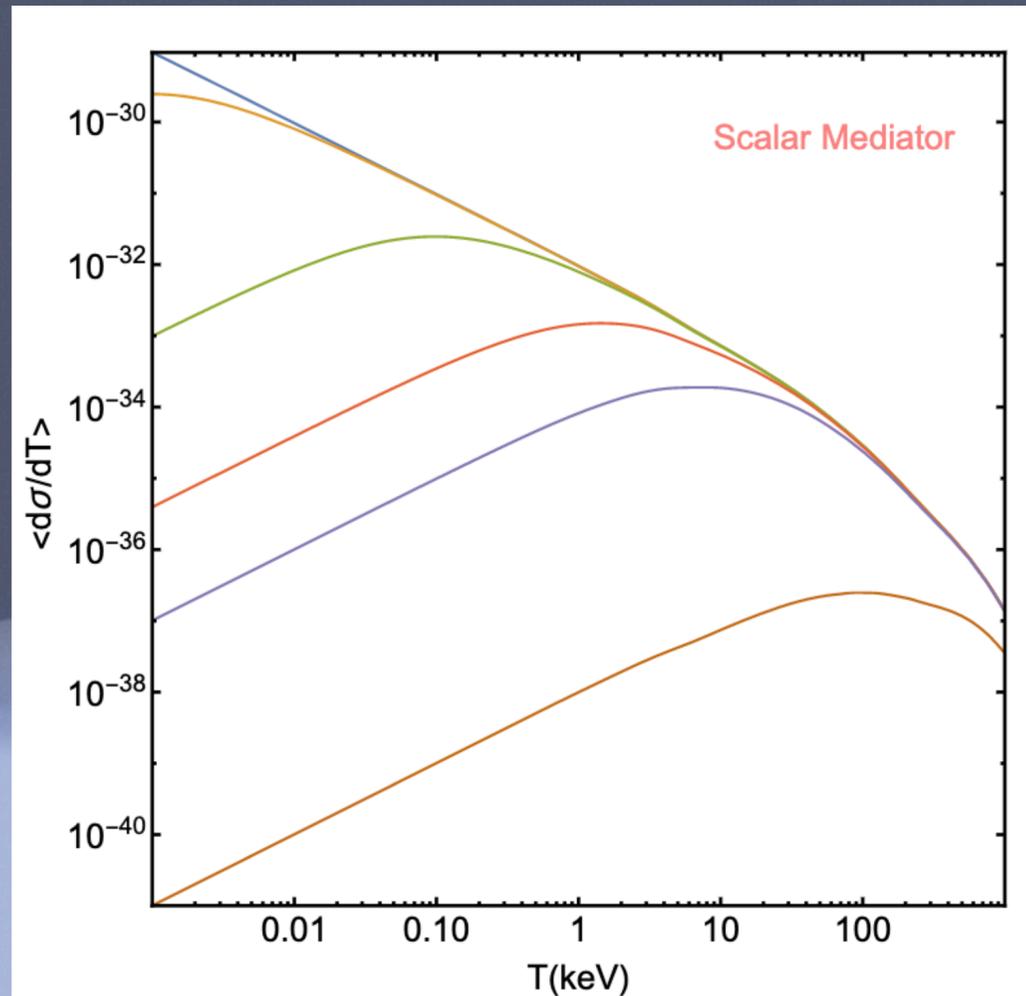
$$\left\langle \frac{d\sigma}{dT} \right\rangle = \frac{\int dE_\nu \Phi(E_\nu) \frac{d\sigma}{dT}}{\int dE_\nu \Phi(E_\nu)}$$

For scalar mediator:

$$g_e = g_\nu = 5 \times 10^{-7}$$

For vector mediator:

$$g_e = g_\nu = 1.5 \times 10^{-7}$$



# NEUTRINO-ELECTRON SCATTERING RATE IN DETECTORS

- For incident neutrino energy less than the typical binding momentum of electron  $\sim Zm_e\alpha$  in a hydrogen like atom, the FEA is not accurate. This would correspond to  $\mathcal{O}(100)$  keV, typical neutrino energy from a MW reactor flux.
- For getting a better understanding of detector responses at low recoil one must include a proper treatment of many electron dynamics in atomic ionization that can be encapsulated in atomic ionization form factor.
- Si/Ge are crystal target with band structure with an energy gap between occupied valance band and unoccupied conduction band. At ultra-low energy transfers most of the transitions occurs from valance to conduction band. This is accounted for in the form of a crystal form factor.

# NEUTRINO-ELECTRON SCATTERING RATE IN DETECTORS

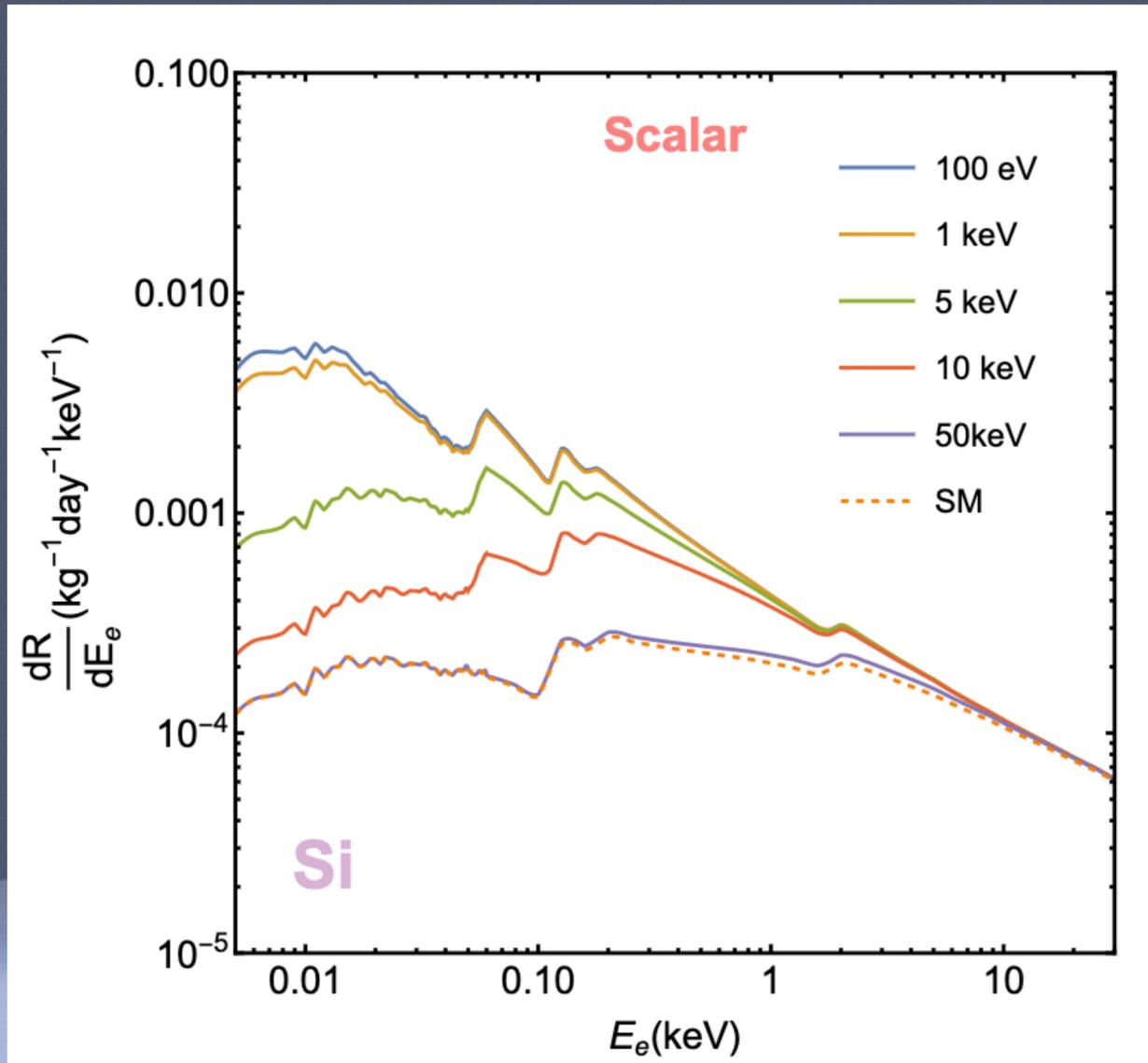
- The total rate is sum of the rates obtained for
  - atomic states to free states transition for deposited energy  $E_e > 60$  eV using atomic form factor from DarkARC code.

$$\frac{d\mathcal{R}_a}{d \ln E_e} = \frac{N_T}{4} \int dE_\nu \Phi(E_\nu) \int dq \left( \frac{d\sigma}{dq} \right) \left| f_{\text{ion}}^{n,l}(q, E_e) \right|^2$$

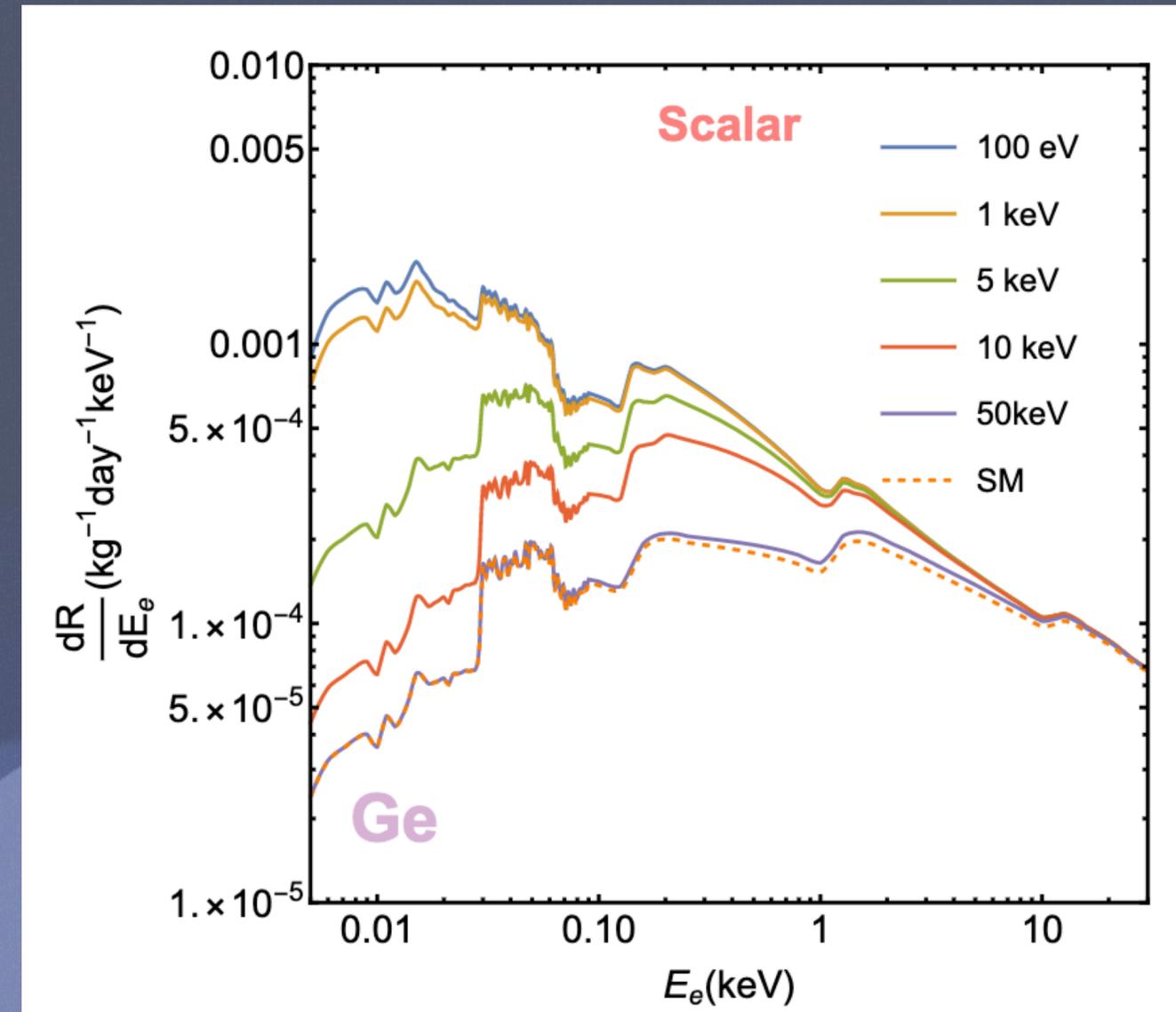
- valance to conduction band transition for deposited energy  $E_e \leq 60$  eV using crystal form factor from EXCEED-DM code.

$$\frac{d\mathcal{R}_{v \rightarrow c}}{d \ln E_e} = N_{\text{cell}} \int dE_\nu \Phi(E_\nu) \int dq \frac{d\sigma_e}{dq} \left| f_{v \rightarrow c}(q, E_e) \right|^2$$

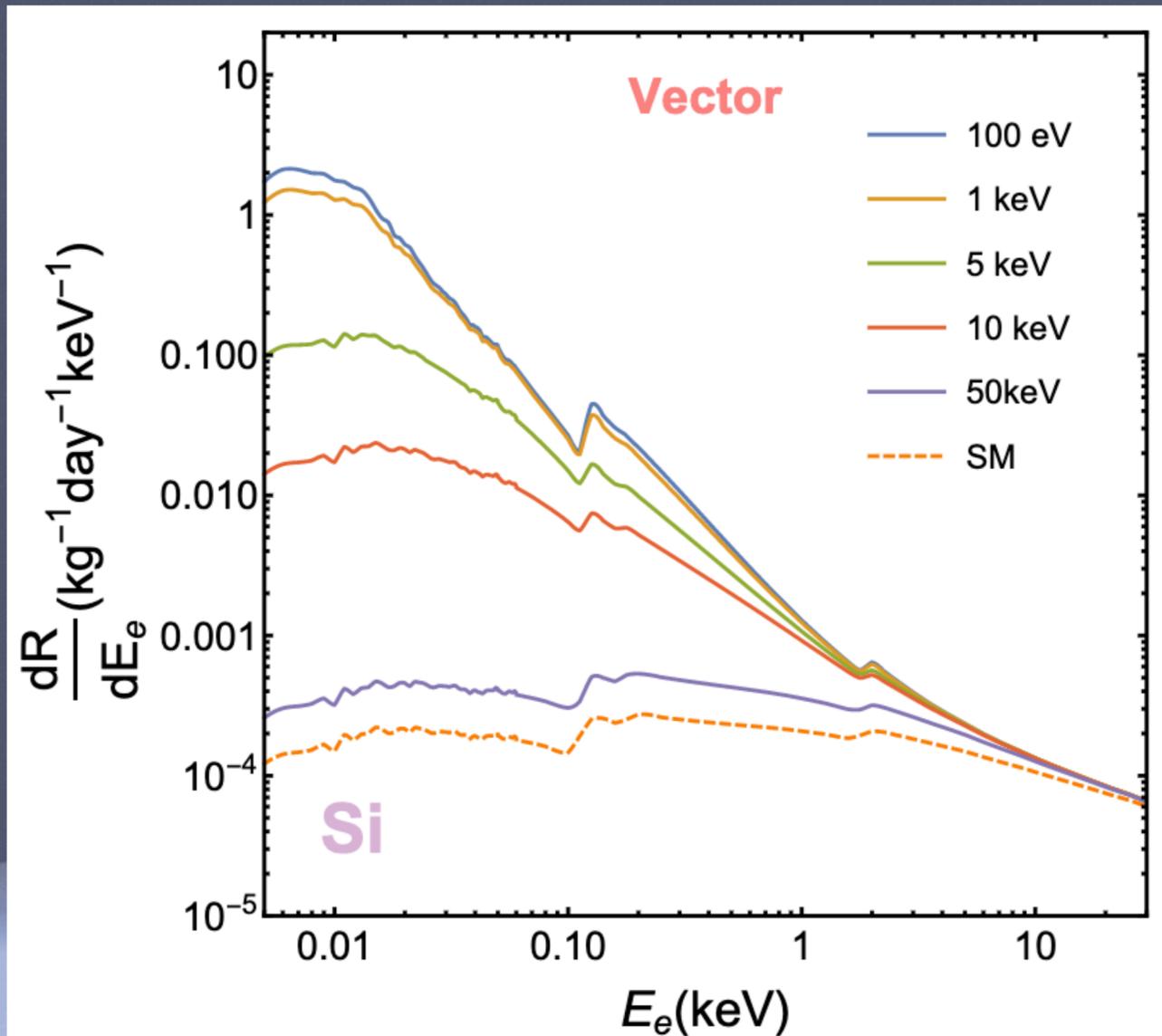
# RECOIL SPECTRA: SCALAR MEDIATOR



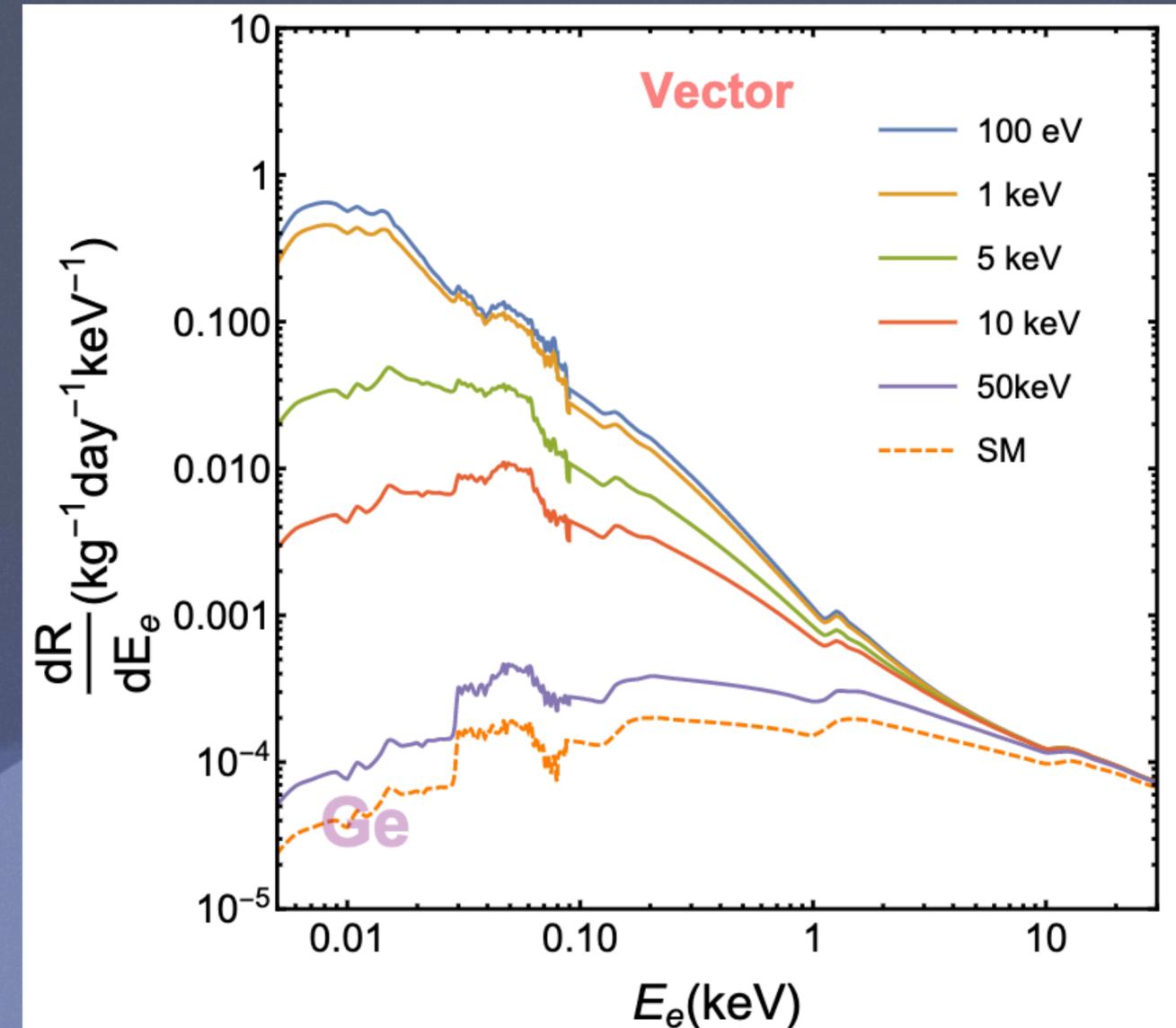
For scalar mediator:  
 $g_e = g_\nu = 5 \times 10^{-7}$



# RECOIL SPECTRA: VECTOR MEDIATOR



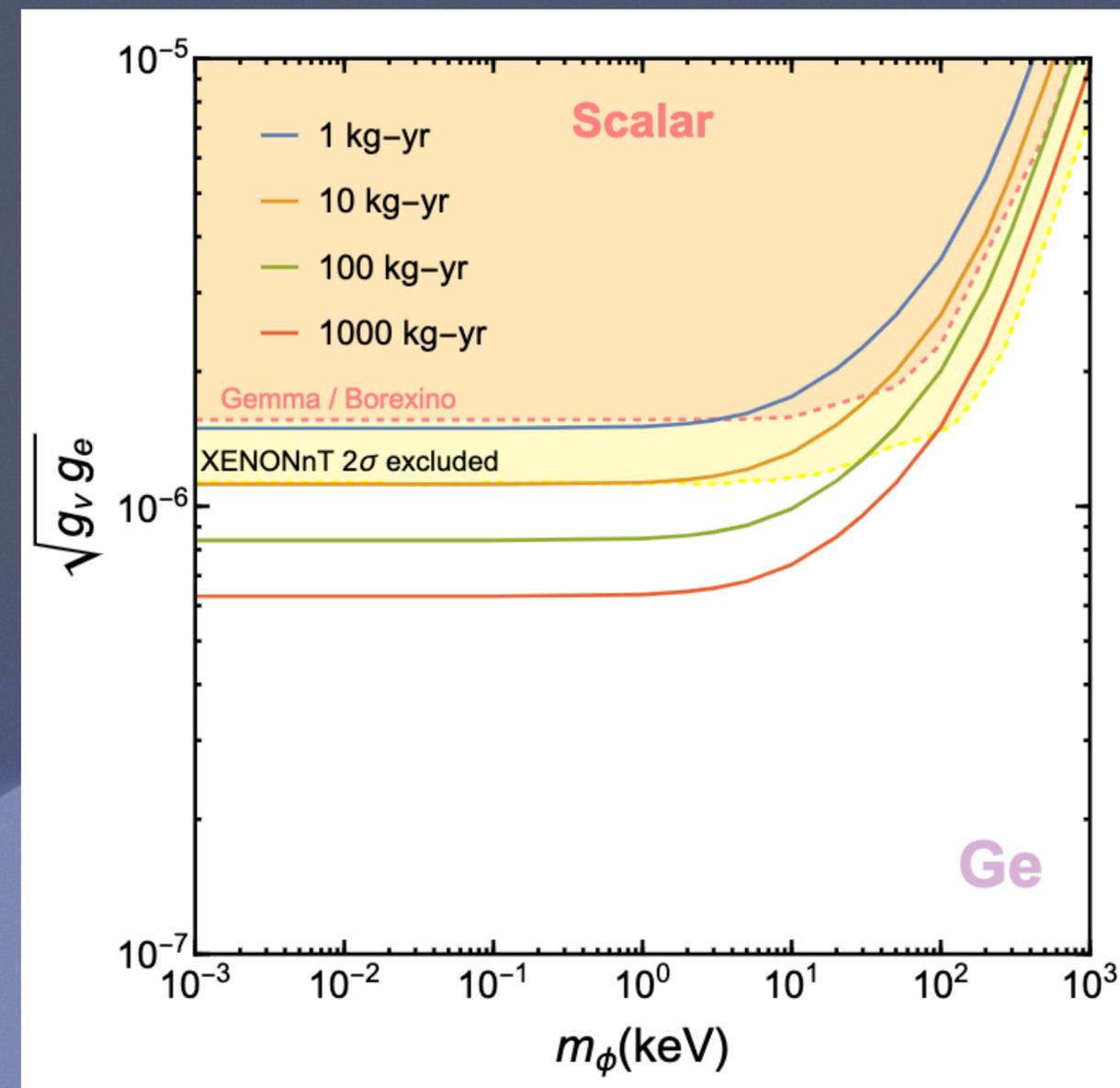
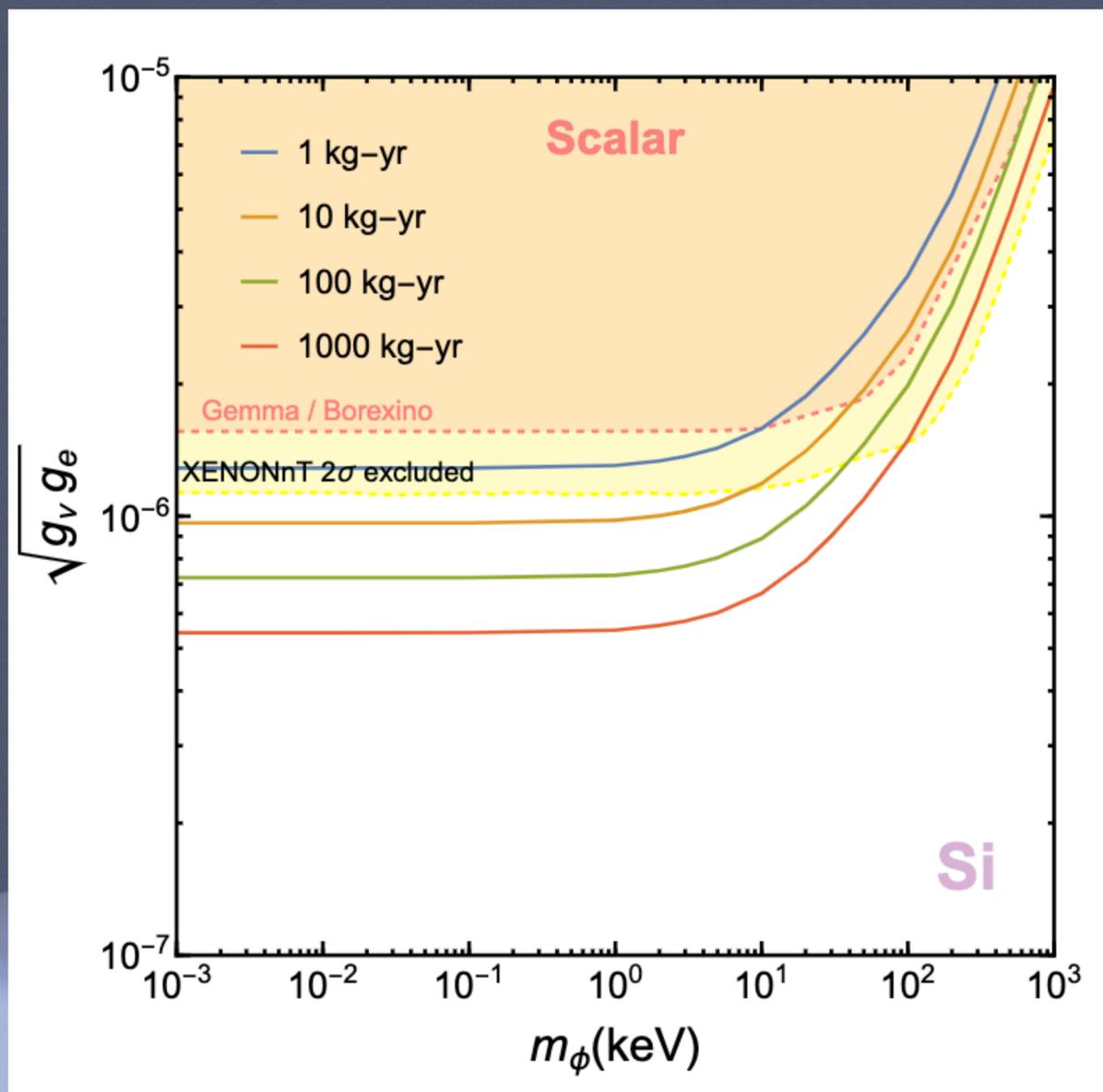
For vector mediator:  
 $g_e = g_\nu = 1.5 \times 10^{-7}$



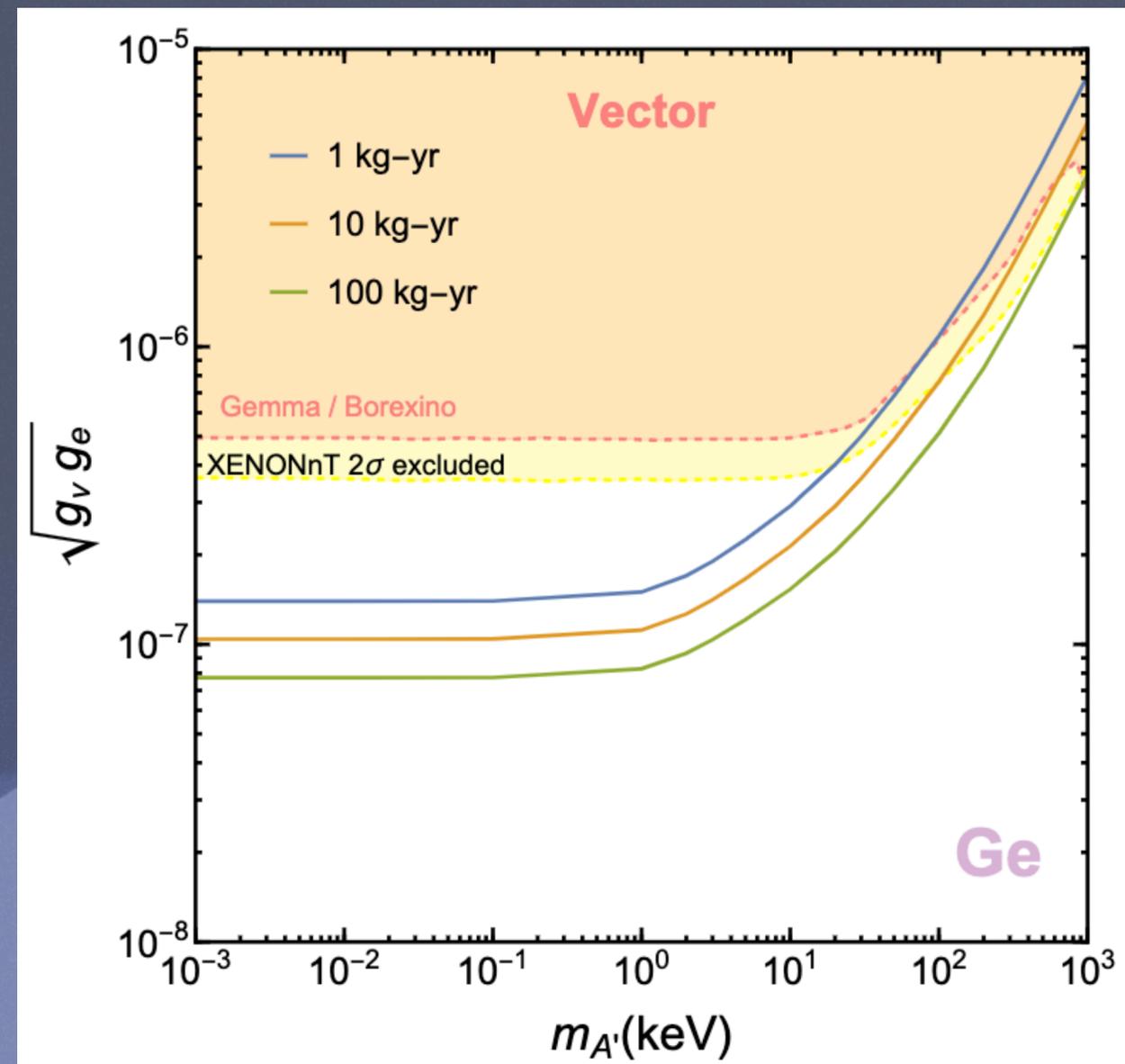
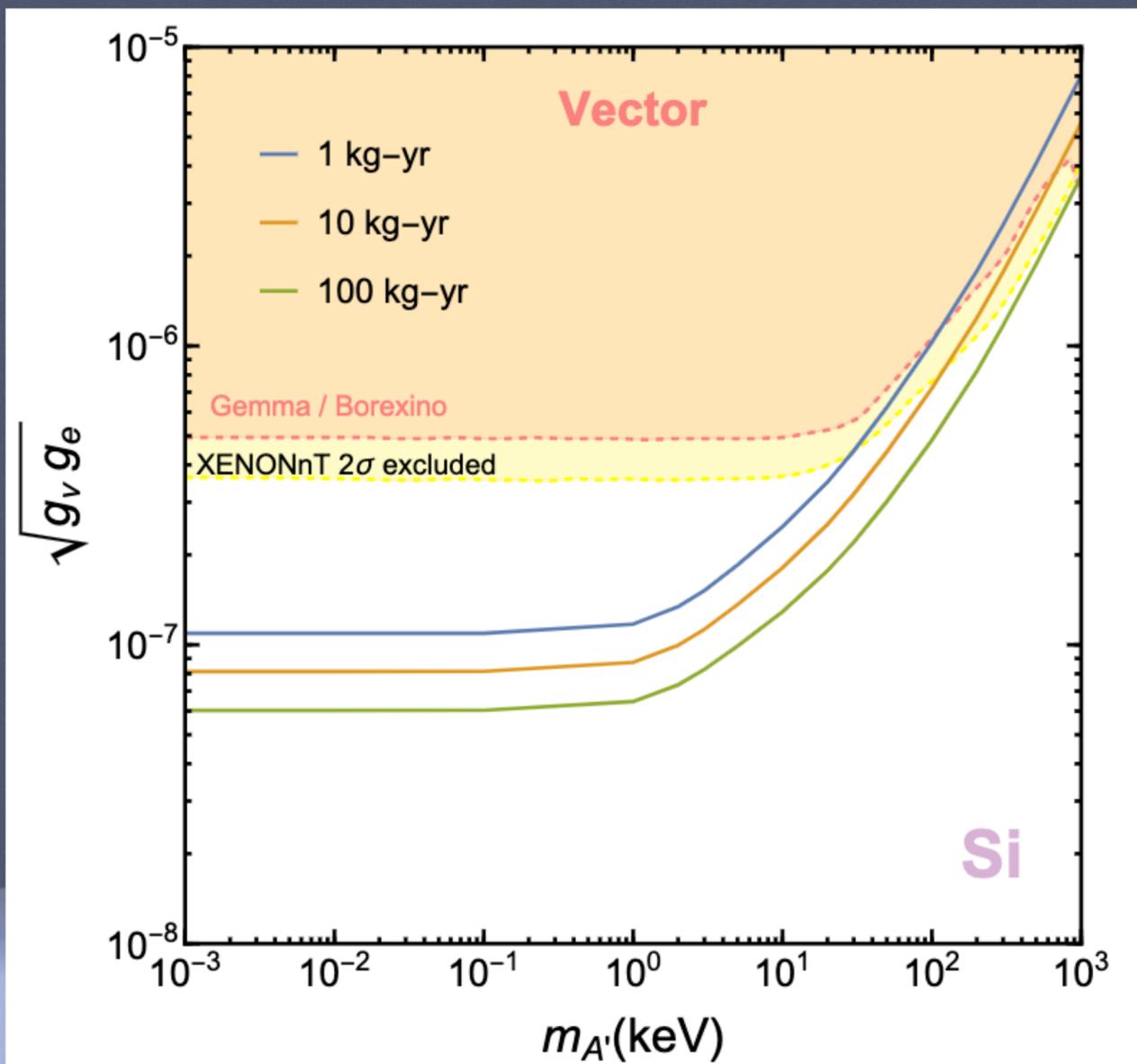
# RESULT: SENSITIVITY CURVE

- To obtain the sensitivity curve
  - we show projected sensitivity at 90% confidence level.
  - corresponding constraints from Borexino, GEMMA and XENONnT are shown.
  - we assume a constant background of 0.1 dru in addition to SM neutrino-electron background.
  - we also assume the detector is capable of discriminating between ER and NR signals down to the threshold of 5 eV.

# SENSITIVITY PLOT: SCALAR MEDIATOR



# SENSITIVITY PLOT: VECTOR MEDIATOR



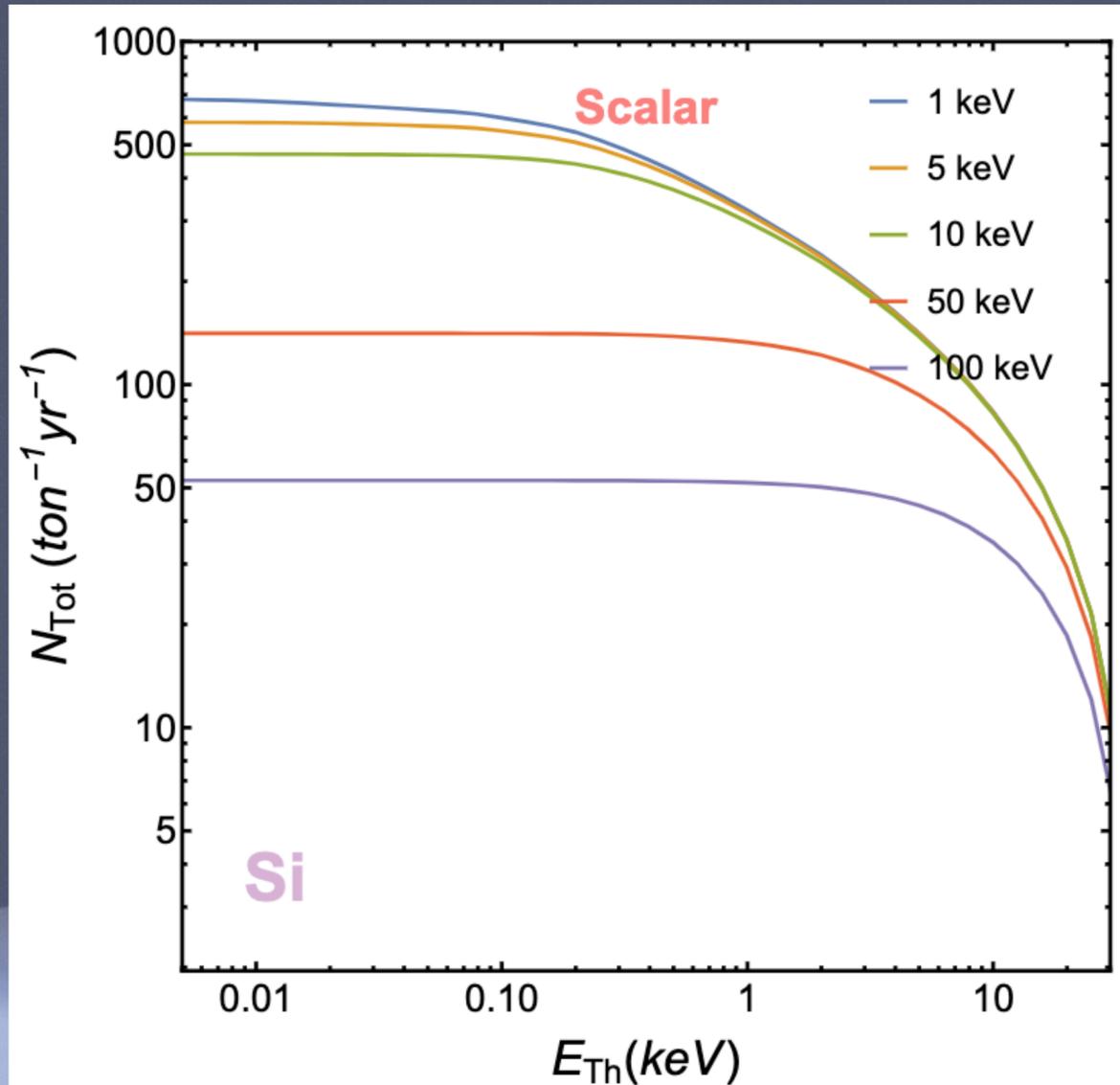
# NATURE OF THE CURVE

- For heavier mediator masses:
  - the coupling sensitivity decreases since for  $q \ll M_{\phi/A'}$  it is Fermi like contact interactions.
  - gives a line with a slope of 2 in the log-log space of the coupling vs mediator mass plot.
- For lighter mediator masses:
  - the mediator mass can be neglected in the propagator for  $q \gg M_{\phi/A'}$ .
  - sensitivity becomes flat in the log-log space of the coupling vs mediator mass plot.

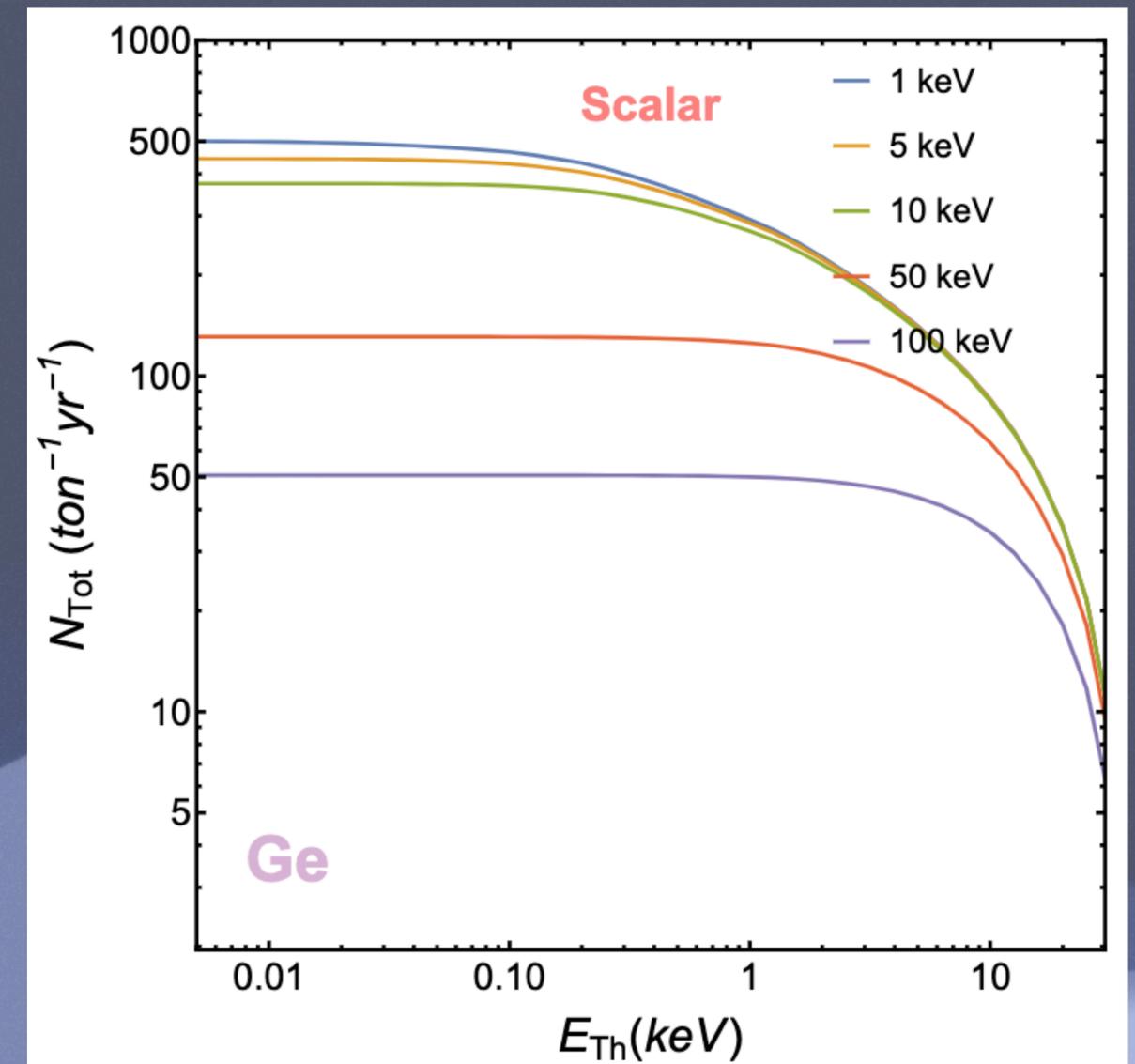
# RESULTS

- Using a 10 kg-yr Ge detector and MW reactor one can probe coupling values of  $\sqrt{g_e g_\nu} \sim 1 \times 10^{-6}$  and  $1 \times 10^{-7}$  for a scalar and vector mediator with masses below 1 keV respectively.
- The same exposure and setup for Si detector gives coupling reach of  $\sqrt{g_e g_\nu} \sim 9.5 \times 10^{-7}$  and  $8 \times 10^{-8}$  for scalar and vector mediator with masses below 1 keV respectively.

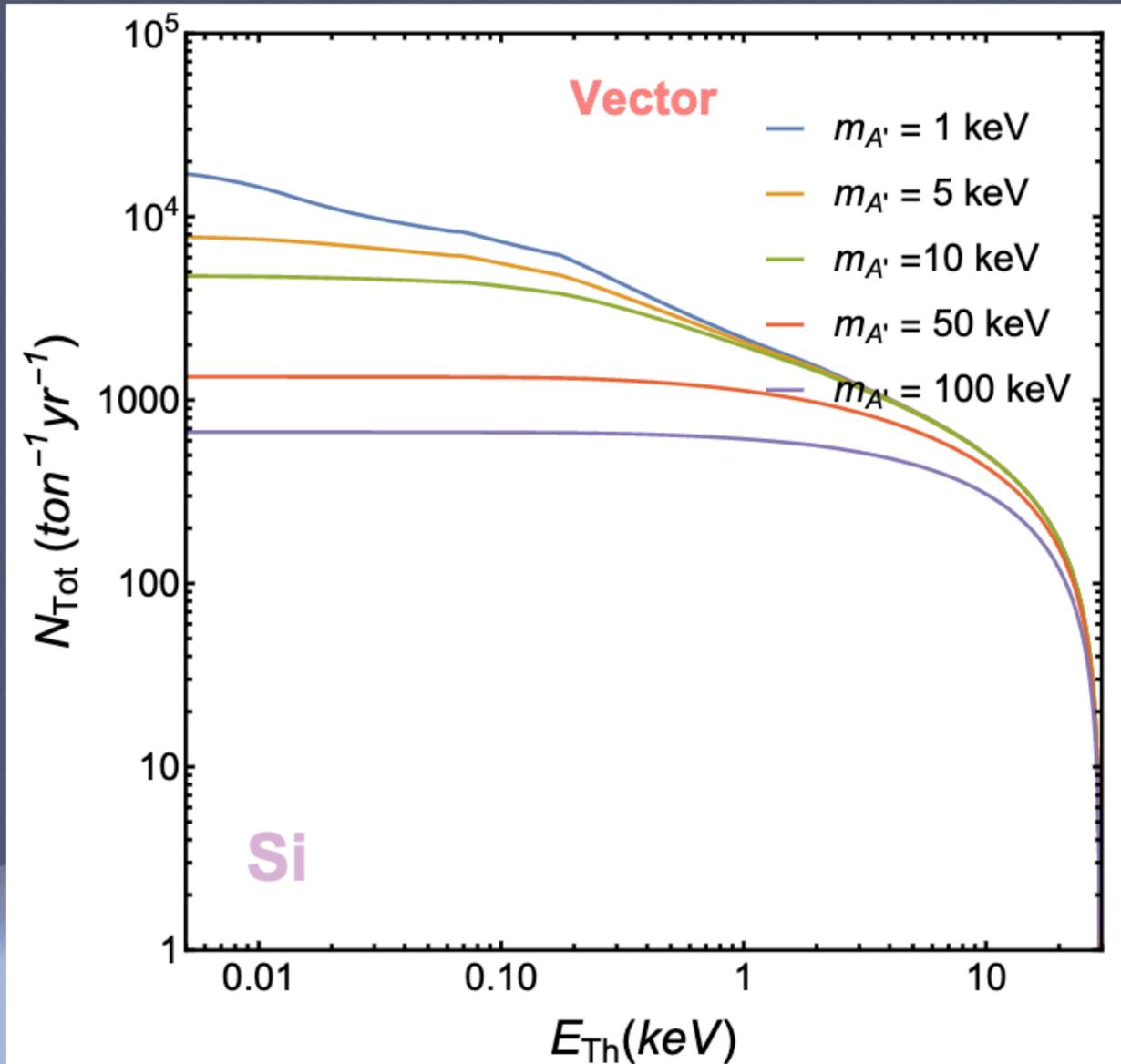
# INTEGRATED EVENTS: SCALAR MEDIATOR



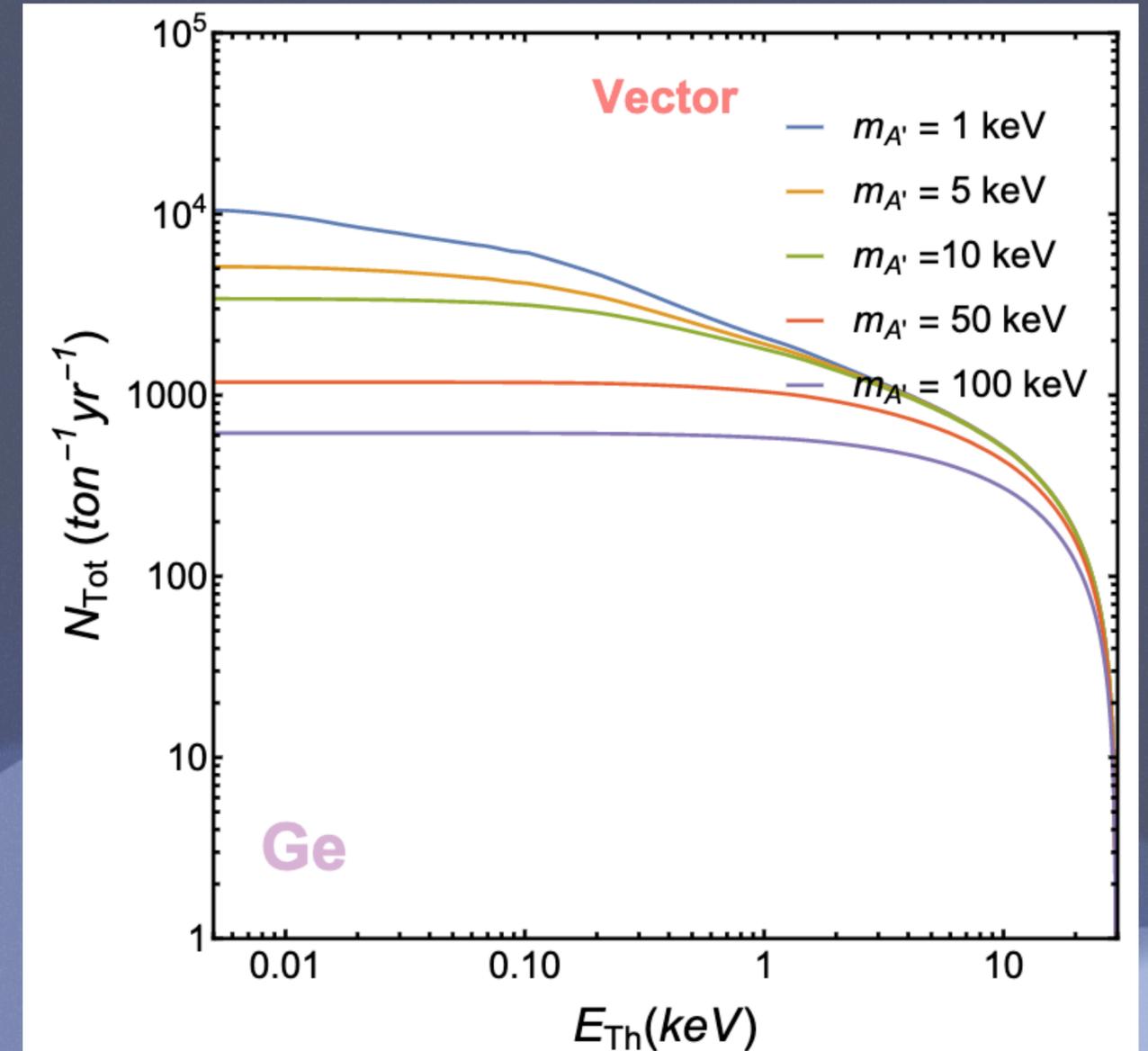
For scalar mediator:  
 $g_e = g_\nu = 5 \times 10^{-7}$



# INTEGRATED EVENTS: VECTOR MEDIATOR



For vector mediator:  
 $g_e = g_\nu = 1.5 \times 10^{-7}$



# LIGHT MEDIATOR MODELS: EXAMPLES

- Light scalar mediator models: we propose one

- by extending the SM scalar sector by

$$H_1 \sim (2, 1/2), \quad H_2 \sim (2, 1/2), \quad \phi \sim (1, 0), \quad \Delta \sim (3, 1).$$

- mixing in the scalar sector give rise to one  $\mathcal{O}(1)$  keV scalar particle.

- Yukawa term  $\bar{l}_{L_i}^c (y')_{ij} i\sigma_2 \Delta l'_{L_j}$  gives NSI of neutrinos .

- Light vector models: the extension of SM by an anomaly-free  $U(1)$  gauge group such as,  $U(1)_{L_i - l_j}$ ,  $U(1)_{B-L}$ ,  $U(1)_{T3R}$  etc.

# LIGHT MEDIATOR MODELS: CONSTRAINTS

- Neutrino-neutrino, neutrino-electron and electron-electron interactions induced by the new light mediator:
  - Borexino and GEMMA are quite sensitive to the modification of the scattering cross-section by NSI with low mass mediators.
  - different astrophysical processes such as energy loss of SN1987A, solar cooling process and cooling of star in globular clusters can put bounds. Can be evaded by Chameleon effect
  - contribution to  $\Delta N_{eff}$  can also put bounds.
  - vector NSI induces flavor dependent matter potential and scalar NSI gives a correction to neutrino mass matrix. Can affect oscillations and scattering experiment.

# CONCLUSIONS

- A well-motivated phenomenological approach to search for new physics, in neutrino sector, is that of NSI.
- At neutrino energy around 200 keV reactor anti-neutrino flux is one order of magnitude larger than the solar flux. Also atomic physics effects becomes important at this energy scale.
- With ER and NR discrimination and low background, the low threshold Si/Ge detector placed in reactor based experiment can probe large parameter space of light mediator induced NSI models.