

Excess of soft photons in inelastic hadronic collisions: a long lasting puzzle

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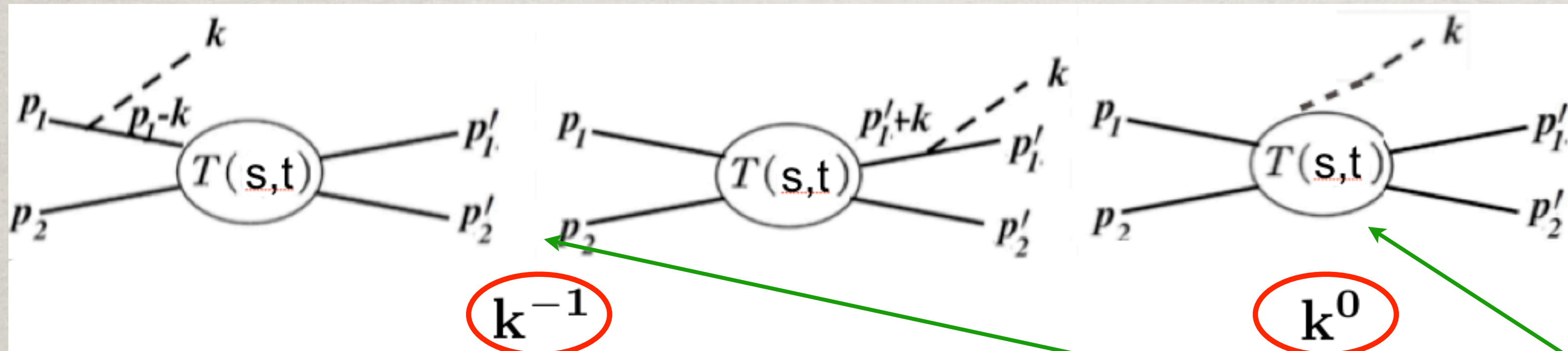
In collaboration with
Irina Potashnikova,
Ivan Schmidt

B. Kopeliovich, PIC 2023

The Low theorem, revisited

$$h_1 + h_2 \rightarrow h'_1 + \gamma + h'_2$$

F. Low, Phys. Rev. 110, 974 (1958)



The amplitude $\mathbf{M} = \mathbf{e}_\mu \mathbf{M}_\mu$ gets contributions from external and internal radiation.

$$\mathbf{M}_\mu = \mathbf{M}_\mu^{\text{ext}} + \mathbf{M}_\mu^{\text{int}}$$

$$\mathbf{M}_\mu^{\text{ext}} = \left(\frac{\mathbf{p}'_{1\mu}}{\mathbf{p}'_1 \mathbf{k}} - \frac{\mathbf{p}_{1\mu}}{\mathbf{p}_1 \mathbf{k}} \right) \mathbf{T}(s, t), \text{ is infrared divergent}$$

The two terms are related by gauge invariance $\mathbf{k}_\mu \mathbf{M}_\mu = 0$

Therefore \mathbf{M}^{int} is not divergent at $\mathbf{k} \rightarrow 0$, i.e. is **suppressed** in comparison with $\mathbf{M}_\mu^{\text{ext}}$

Landau-Pomeranchuk principle

L.Landau, I.Pomeranchuk,
Dokl. Akad. Nauk 92, 535(1953)

The Low theorem can also be treated as a formal proof of the Landau-Pomeranchuk principle, which states that any variation of the electric current within a short distance do not affect the spectrum of radiation with much longer coherence length,

$$I_c^\gamma = \frac{2E_{h1} x_1 (1 - x_1)}{k_T^2 + x_1 m_h^2} \quad x_1 = \frac{k_+^\gamma}{p_+^{h1}}$$

Low: "the distance a particle can move with energy imbalance $\Delta E=k$ "

Important is to keep the incoming ($l < -l_c$) and outgoing ($l > l_c \gamma$) currents unaffected by the current variations on shorter length scales. This means that only extrinsic radiation from initial and final hadrons $h1, h'1$ matters.

I_c^γ and E_{h1} are not Lorentz invariant, must be taken within the same reference frame

Fock state representation

The Low process $h_1 + h_2 \rightarrow h'_1 + \gamma + h'_2$ is diffractive excitation $h_1 \rightarrow h'_1 + \gamma$

The Good-Walker picture of diffraction:

M.L.Good & W.D.Walker,
Phys.Rev. 120, 1857(1960)

$$|h\rangle = C_0|h\rangle_0 + C_1|h\gamma\rangle_0 + \dots$$

The Fock components $|i\rangle$ are eigenstates of interaction, i.e. $\hat{f}|i\rangle = f_i|i\rangle$

Diffractive excitation occurs only due to diversity of the elastic eigenamplitudes, otherwise the incoming wave packet remains unchanged.

$$\langle h\gamma|\hat{f}|h\rangle = C_1^* C_0 (f_{h\gamma} - f_h)$$

Differently from the Feynman diagrams, one cannot say whether the photon is radiated before or after the interaction. The radiation amplitude is a linear combination of elastic eigenamplitudes.

The Fock component lifetime is the coherence time/length l_c^γ

Fock state expansion

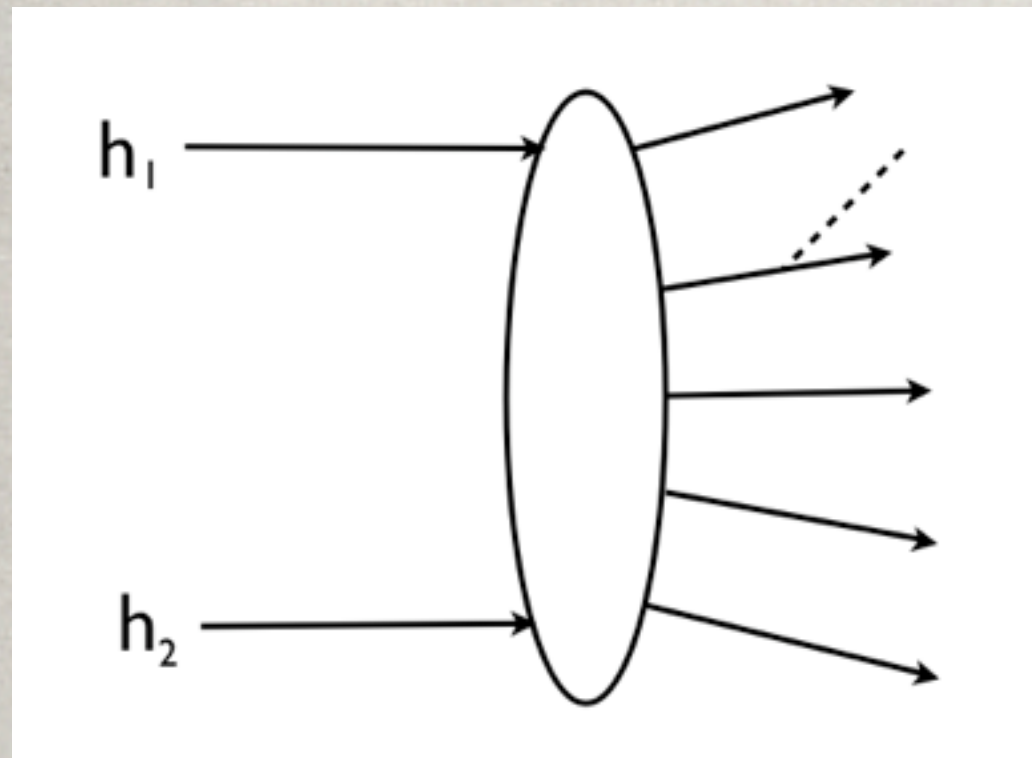
In both Fock components $|h\rangle$ and $|h\gamma\rangle$ only hadron h interacts, so at first glance $f_{h\gamma} = f_h$ and the diffractive amplitude vanishes. However, the impact parameters b are shifted, so the eigen-amplitudes $f_{h\gamma}(b) \neq f_h(b')$ cancel only in b -integrated amplitude, i.e. at $t=0$.

The Low amplitude $M_{\mu}^{\text{ext}} = \left(\frac{p'_{1\mu}}{p'_1 k} - \frac{p_{1\mu}}{p_1 k} \right) \mathbf{T}(s, t)$, indeed vanishes at $t \rightarrow 0$.

Photon radiation in inelastic collisions

The so-called **bremsstrahlung model (BM)** pretends to extend the Low theorem from the radiation in elastic scattering to inelastic collisions with multi-particle production

Photons are assumed to be radiated by participating charge particles, either the incoming, or outgoing.



$$M(\mathbf{p}_1, \mathbf{p}_2; \mathbf{p}_3 \dots \mathbf{p}_N \mathbf{k}) = M_0(\mathbf{p}_1, \mathbf{p}_2; \mathbf{p}_3 \dots \mathbf{p}_N) \left(\sum_i^{\text{charged}} \frac{\eta_i \mathbf{e}_i \mathbf{p}_i \cdot \boldsymbol{\epsilon}}{2\mathbf{p}_i \cdot \mathbf{k}} \right)$$

where $\eta_i = \pm 1$ for outgoing and incoming particles respectively, and M_0 is the amplitude of $2 \rightarrow N$ inelastic collision without radiation.

A.T. Goshaw et al., PRL 43(1979)1065

T. H. Burnett & N. M. Kroll, PRL 20(1968)86

• • •

The model turns out to contradict data significantly

A puzzling enhancement of low- k_T photons in high-energy hadronic and hadron-nucleus collisions in comparison with theoretical expectations has been repeatedly observed since 1984.

Experiment	Collision Energy	Photon k_T	Photon/Brem Ratio
$K^+ p$, CERN, WA27, BEBC (1984)	70 GeV/c	$k_T < 60$ MeV/c	4.0 ± 0.8
$K^+ p$, CERN, NA22, EHS (1993)	250 GeV/c	$k_T < 40$ MeV/c	6.4 ± 1.6
$\pi^+ p$, CERN, NA22, EHS (1997)	250 GeV/c	$k_T < 40$ MeV/c	6.9 ± 1.3
$\pi^- p$, CERN, WA83, OMEGA (1997)	280 GeV/c	$k_T < 10$ MeV/c	7.9 ± 1.4
$\pi^+ p$, CERN, WA91, OMEGA (2002)	280 GeV/c	$k_T < 20$ MeV/c	5.3 ± 0.9
pp , CERN, WA102, OMEGA (2002)	450 GeV/c	$k_T < 20$ MeV/c	4.1 ± 0.8

C.-Y. Wong, (2013)

BM "extension" of the Low theorem is not only unjustified, but in fact contradicts the Low theorem

First of all, the notion of incoming and outgoing particles is poorly defined. The process is characterized by two length scales:

(i) One scale is related to the coherence length of photon radiation (defined by Low as "the distance a particle can move with energy imbalance")

$$l_c^\gamma = \frac{2E_1 x_1 (1 - x_1)}{k_T^2 + x_1 m_h^2}$$

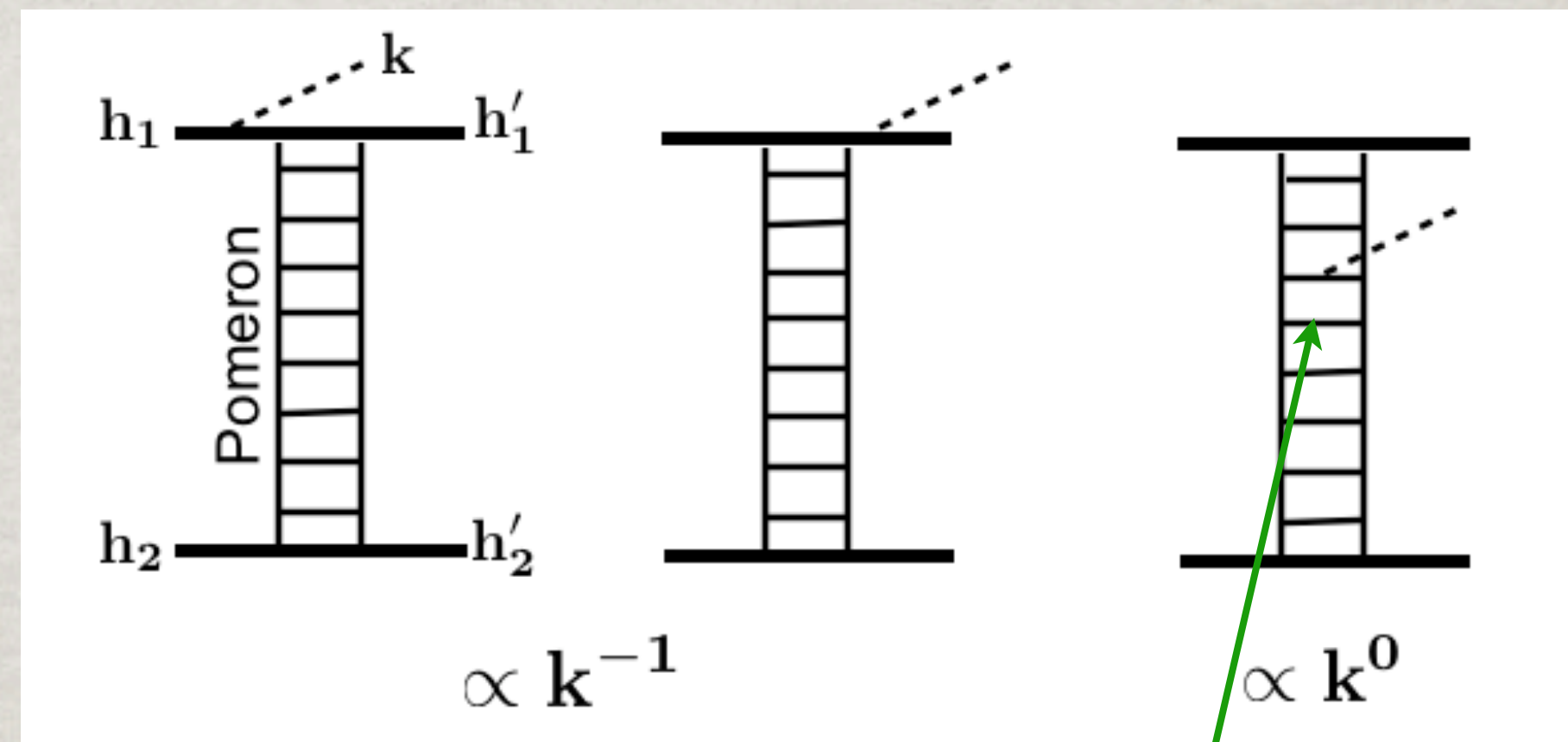
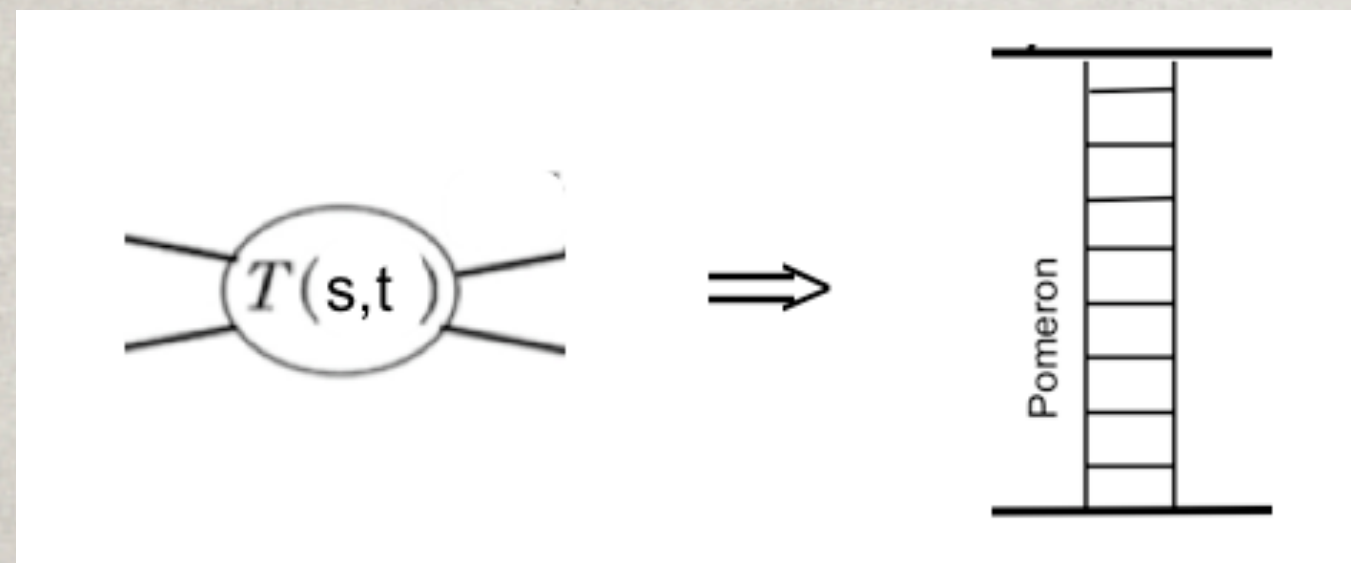
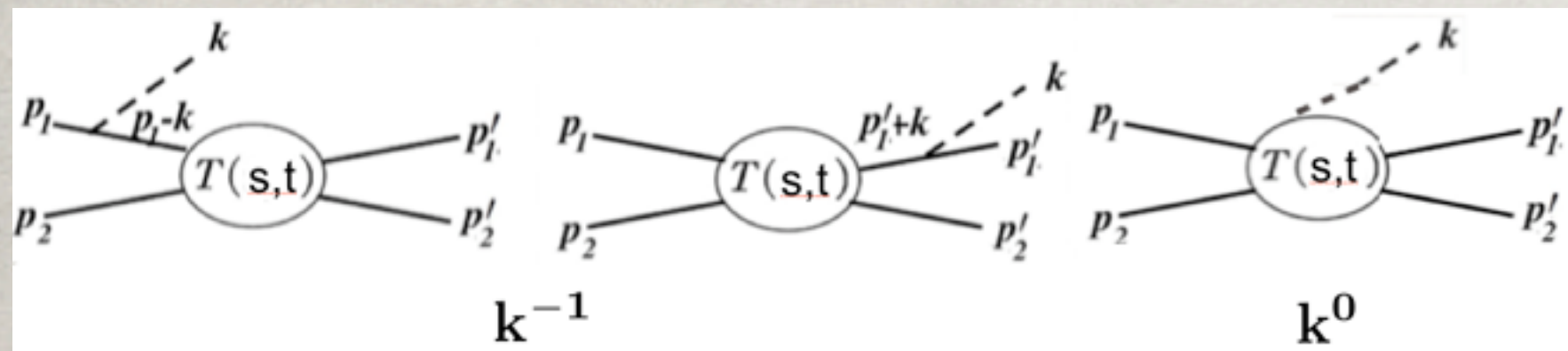
E_1 is the energy of h_1 in the rest frame of h_2 .
 x_1 is the fractional light-cone momentum of the photon

$$x_1 = k_+^\gamma / p_+^{h_1}$$

(ii) The range of strong interactions, is assumed to be short. In the target rest frame it is given by the inverse pion mass $l_h \sim 1 \text{ fm}$. The main condition of the Low theorem is $l_c^\gamma \gg l_h$

At high **energies** the Low process essentially simplifies. Elastic scattering dominates, so $h_1 = h'_1$ and $T(s,t)$ is the elastic amplitude, dominated by Pomeron exchange.

Besides, the energy dependence of $T(s,t)$ can be neglected, $\frac{\partial T(s,t)}{\partial s} \approx \frac{1}{s}$



M^{ext}

M^{int}

The key result of the Low theorem is suppression of internal radiation.

Unitarity relation

Optical theorem relates multi-particle production with the imaginary part of the forward elastic amplitude.

$$2 \operatorname{Im} \left[\text{Diagram} \right] = \left| \text{Diagram} \right|^2$$

Thus, radiation by produced charged particles, within the "bremsstrahlung model", corresponds to the strongly suppressed internal radiation term M^{int} of the Low process.

On the contrary to the claim that the "bremsstrahlung model" is an extended version of the Low theorem, in fact it strictly contradicts it.

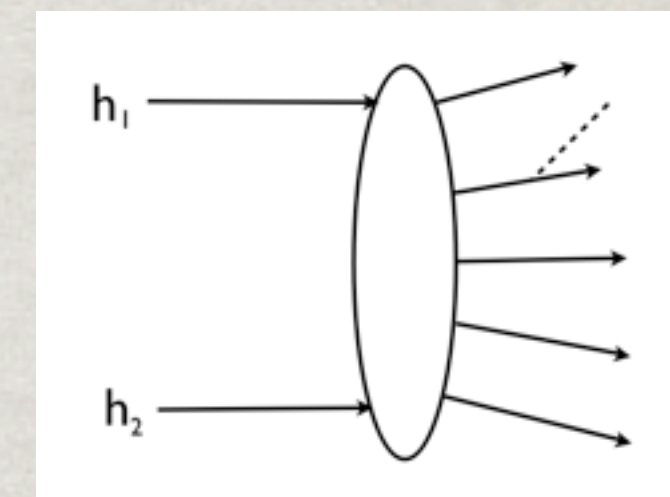
Charge screening

Specific feature of the Pomeron is electric neutrality. This means that electric charges of the colliding hadrons h_1, h_2 do not flow towards the central rapidities. A charge can be transported through large rapidity intervals only by Reggeons, having intercepts $\alpha(0) \approx 1/2$ and lower. Therefore production of a charged hadron with rapidity η is suppressed by $\exp[-(Y-\eta)/2]$, where $Y = \ln(\sqrt{s}/2)$. Thus, the electric charges of the colliding hadrons remain at maximal and minimal rapidities and photons cannot be radiated at the medium rapidities, as is incorrectly assumed in the BM.

Of course one could introduce quark loops in the gluonic ladder of the Pomeron, however photon radiation from quarks and antiquarks cancels due to charge screening.

This is another deficiency of BM.

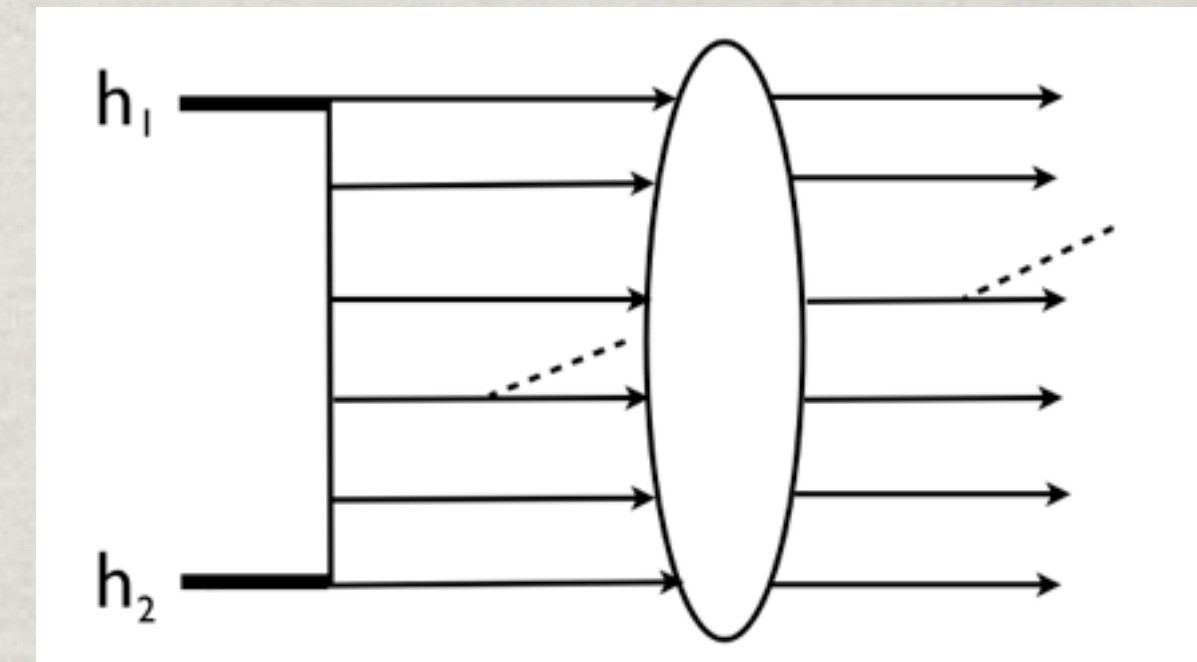
What is wrong with BM,
why Feynman rules cannot be applied ?



?

The propagator in coordinate representation acquires an infinitely long radiation length.

$$\sum_i^{\text{charged}} \frac{\eta_i \mathbf{e}_i \mathbf{p}_i \cdot \boldsymbol{\epsilon}}{2\mathbf{p}_i \cdot \mathbf{k}} \longrightarrow \frac{1}{p^2 - m_h^2} = \frac{x_1(1-x_1)}{k_T^2 + x_1 m_h^2} = \frac{l_c^\gamma}{2E_1}$$



Thus, the process cannot be treated as radiation either from the two incoming, or from N outgoing charges. All the produced partons pre-exist the interaction.

This explains the observed contradiction between the BM and Low/optical theorems.

As far as the bremsstrahlung model turns out to be incorrect, an alternative description of soft photon radiation is required.

BK, I.Potashnikova, M.Krelina, K.Reygers, [2212.03429](#)

The usual factorization relation for photon radiation

$$\frac{d^4\sigma}{dM^2 dx_F d^2k_T} = \frac{\alpha_{em}}{3\pi M^2} \frac{x_1}{x_1 + x_2} \int_{x_1}^1 \frac{d\alpha}{\alpha^2} \sum_f Z_f^2 \left\{ q_f \left(\frac{x_1}{\alpha} \right) + q_{\bar{f}} \left(\frac{x_1}{\alpha} \right) \right\} \frac{d\sigma(qN \rightarrow \gamma^* X)}{d \ln \alpha d^2k_T}$$

$$x_1 x_2 = M^2/s; \quad 2x_2 = \sqrt{x_F^2 + 4M^2/s} - x_F \quad \alpha = \mathbf{p}_+^\gamma / \mathbf{p}_+^q$$

cannot be used in the soft domain of massless photons and small k_T

For the quark distribution function we rely on the Quark-Gluon String model (QGSM)

A. Kaidalov, Phys.Lett. B116(1982)459

A. Kaidalov and M. Poghosyan, *Eur.Phys.J.C* 67 (2010) 397

or Dual Parton Model

A. Capella, U. Sukhatme, C-I Tan, J. Tran Thanh Van, Phys.Rept. 236(1994)225

A.Capella, A.Kaidalov, C.Merino, J.Tran Thanh Van, Phys.Lett.B 337 (1994) 358

which are nearly the same, both based on Regge phenomenology

The radiation cross section $\frac{d\sigma(qp \rightarrow \gamma p)}{d \ln \alpha d^2 k_T}$

is calculated within the **color dipole approach**

B.K., Hirscheegg 1995 hep-ph/9609385

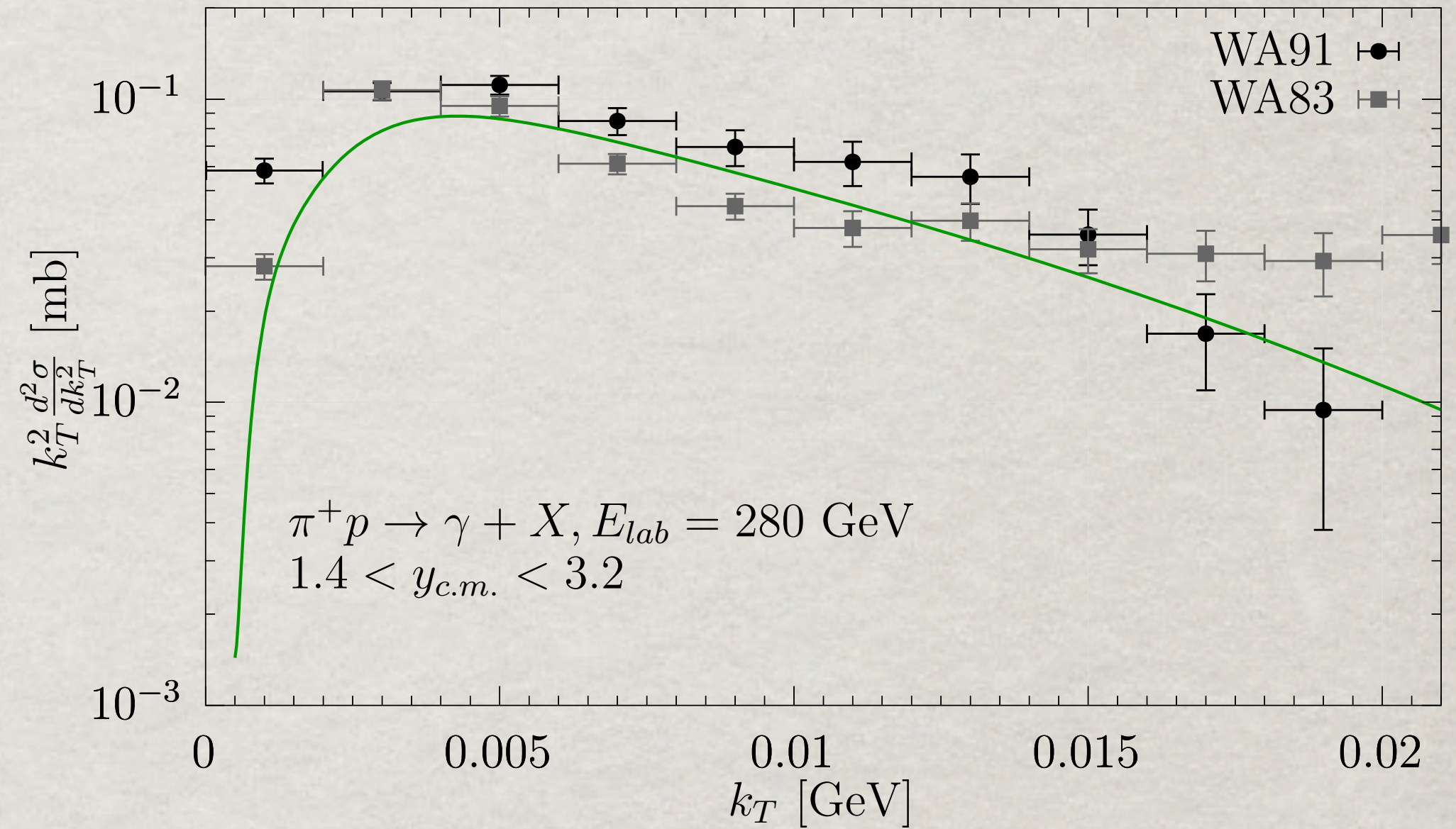
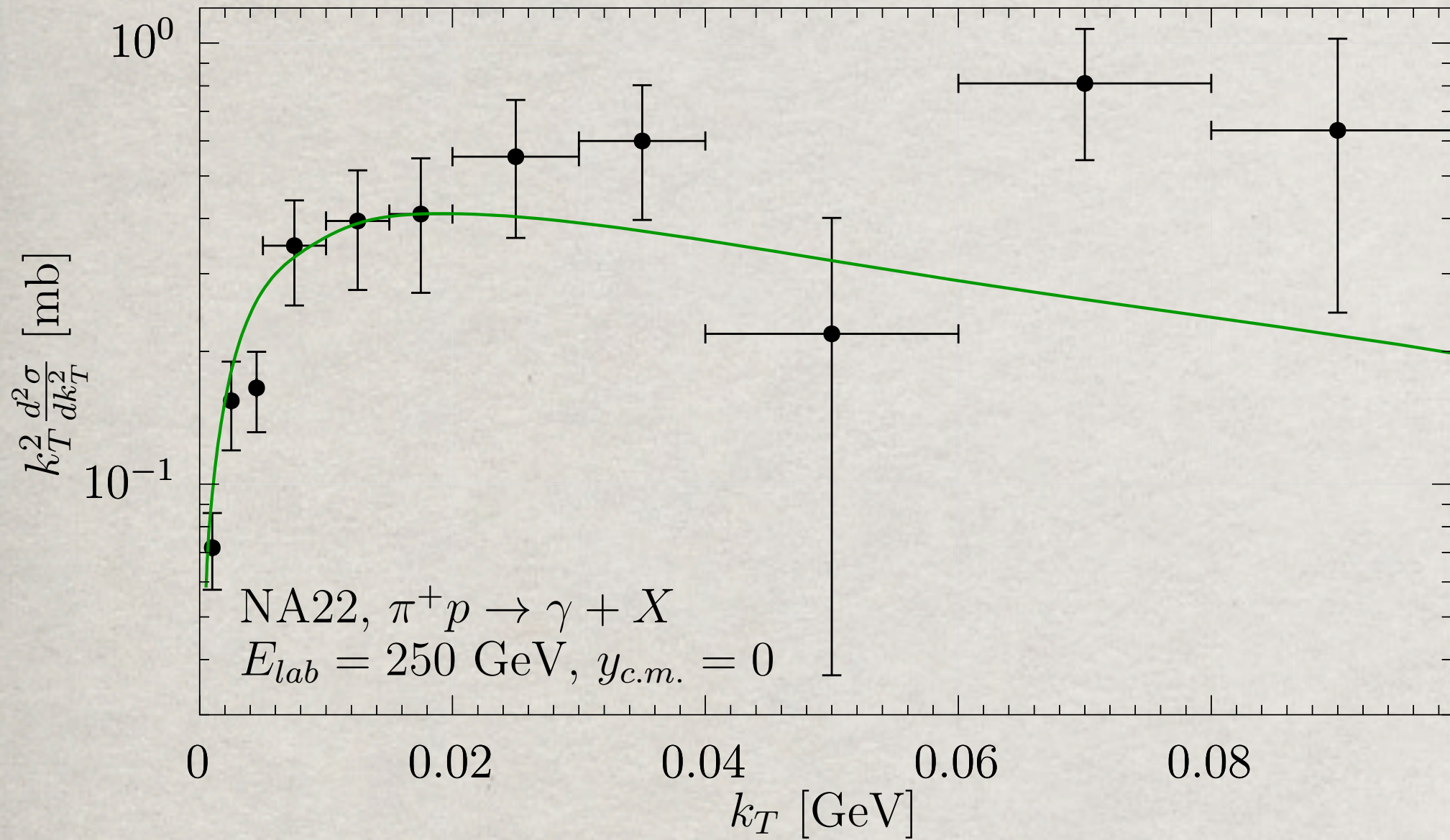
B.K., A.Schäfer, A.Tarasov,
Phys.Rev.C 59(1999)1609;
Phys.Rev.D 62 (2000) 054022

B.K., J.Raufeisen and A.Tarasov, Phys.Lett. B503(2001)91

B.K., A.Rezaeian, H-J.Pirner, I.Schmidt, Phys.Lett. B653(2007)210

B.K., E.Levin, A.Rezaeian, I.Schmidt, Phys.Lett. B675(2009)190

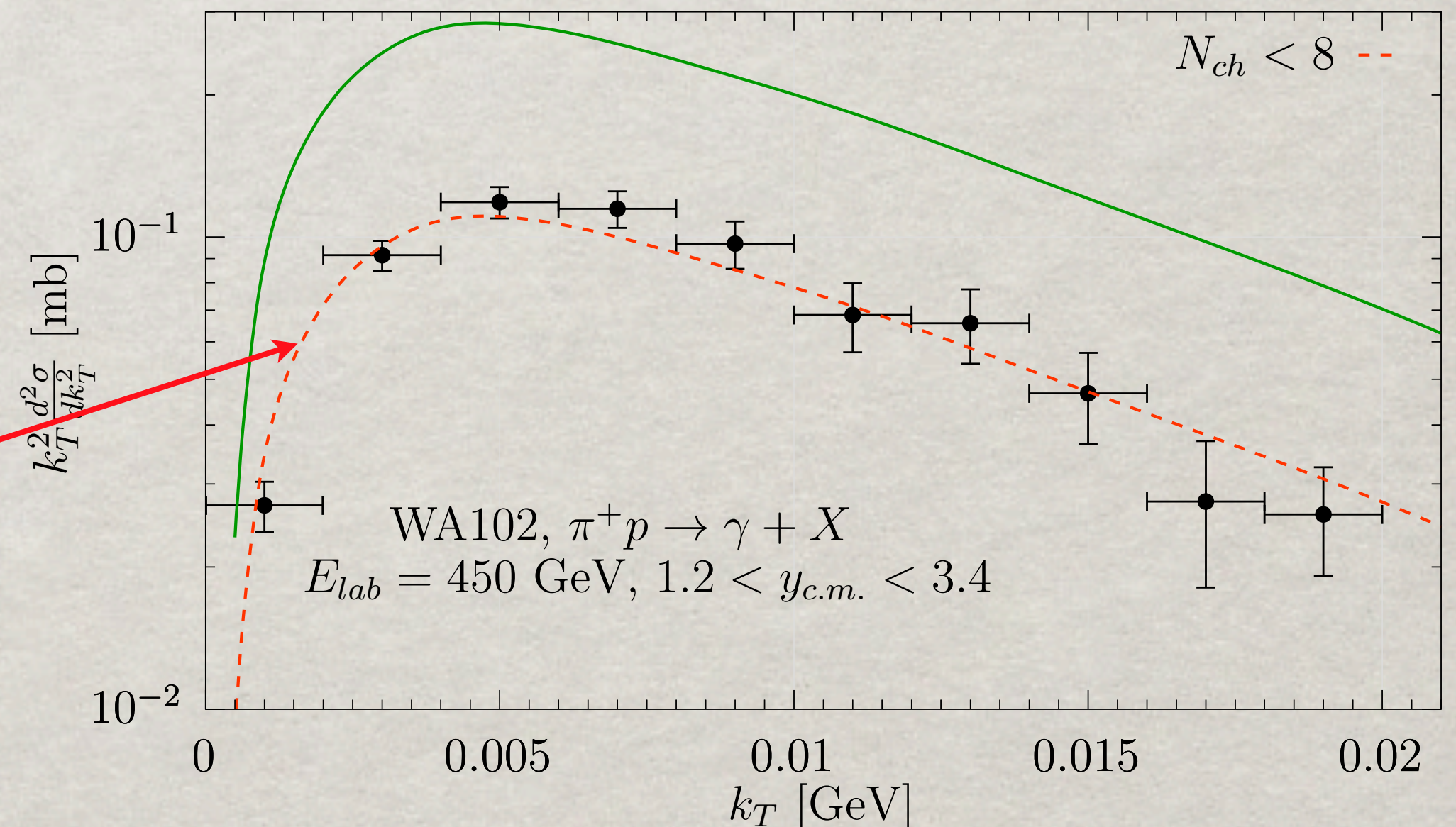
Data



The primordial q_T distribution of the quarks was taken into account: $\tilde{k}'_T = \tilde{k}_T - \alpha \tilde{q}_T$

WA102 had a cut for number of charged tracks $N_{ch} < 8$

That leads to the suppression factor 0.39 (QGSM)



Conclusions

The observed enhancement of low- k_T photons in comparison with incorrect calculations, should not be treated as a puzzle.

Low's paper considered a large rapidity gap process of diffractive photon radiation $h \rightarrow h + \gamma$, which has little to do with multiple production.

Radiation of small- k_T photons takes long time $l^\gamma \propto 1/k_T$, so they cannot appear momentarily from the interaction blob, as is assumed in the Bremsstrahlung Model.

Due to electromagnetic neutrality of the Pomeron photon radiation at the mid-rapidities cancels. This interference effect is missed in BM.

Calculations based on QGSM and color dipole description of photon radiation well agree with data in the whole range of k_T .