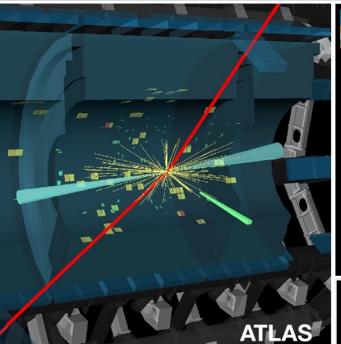
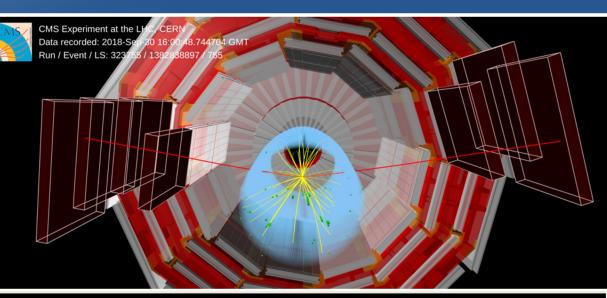


#### Rare processes and Searches for new Phenomena with **Higgs bosons**



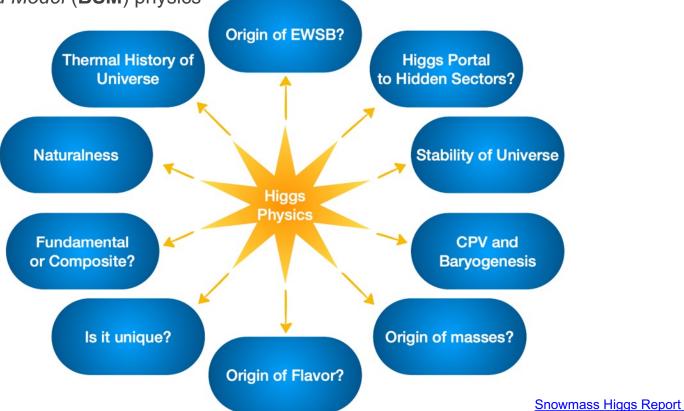


Bernd Stelzer for the ATLAS & CMS Collaboration Physics in Collision Conference, Arica, Oct 10<sup>th</sup>, 2023

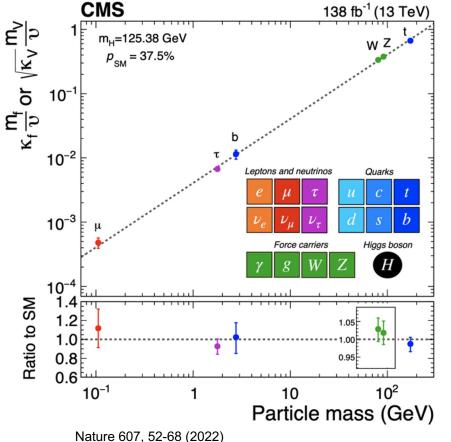


#### Introduction

The Higgs boson is **central** to **Standard Model** and *beyond Standard Model* (**BSM**) physics



### Introduction



CMS

- In the past 11 years, since the Higgs boson discovery, we have learned a lot about its properties, coupling to bosons and 3<sup>rd</sup> generation fermions (see previous talk)
- Large LHC datasets can uncover rare Higgs
   boson processes e.g. couplings to the 2<sup>nd</sup> generation of fermions, processes with destructive interference and processes that are not possible at tree level
  - Physics Beyond the SM (**BSM**) could greatly enhance their rate, motivating these analyses
  - The Higgs boson could also directly couple to
     BSM phenomena

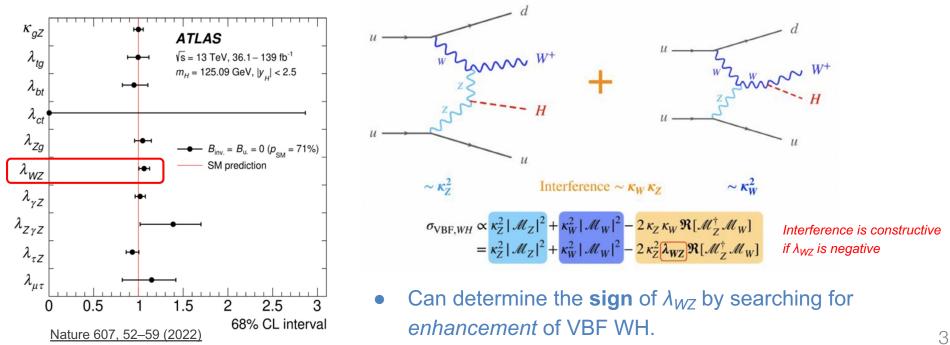


HIGGS boson

# 

### Rare Process, VBF HW

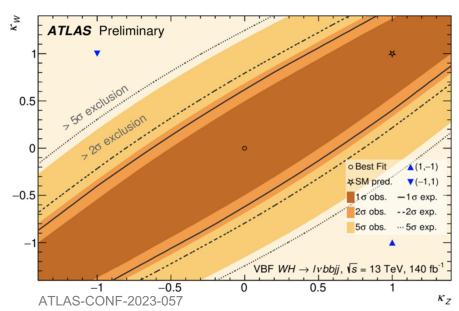
- Measuring the Higgs couplings to *W* and *Z* is essential for testing EWSB and custodial symmetry of the SM. Latest ATLAS result measure  $\lambda_{WZ} = \kappa_W / \kappa_Z = 1.06 \pm 0.06$
- Currently *little* sensitivity to the **sign** of  $\lambda_{WZ}$
- Destructive interference of VBF HW process leads to rare Higgs process in SM

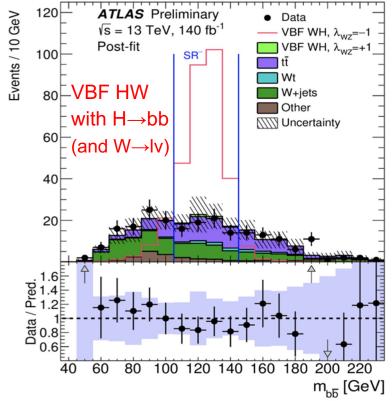


# 

## Rare Process, VBF HW

- Negative  $\lambda_{WZ}$  signal would be easily separable from background based on kinematics
- No excess is seen above SM prediction.
- Opposite sign couplings for  $\kappa_W$ ,  $\kappa_Z$  consistent with other Higgs measurements are **excluded** at > 8 $\sigma$ 
  - *W*, *Z* couplings to Higgs have same sign

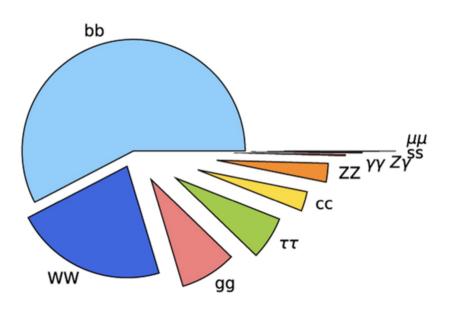




#### Higgs Rare Decays

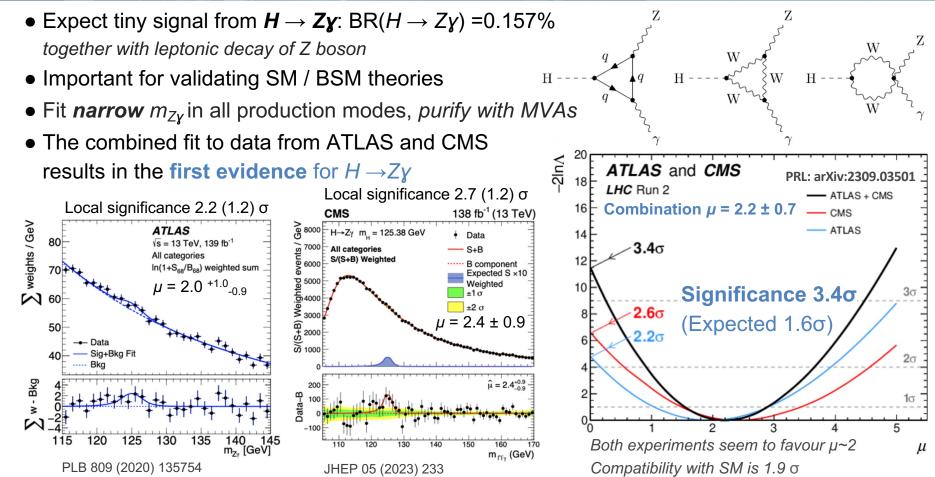
Decay channel	Branching	fraction	(%)
bb	57.63	$\pm 0.70$	
WW	22.00	$\pm 0.33$	
$\mathbf{g}\mathbf{g}$	8.15	$\pm 0.42$	
ττ	6.21	$\pm 0.09$	
сс	2.86	$\pm 0.09$	)
ZZ	2.71	$\pm 0.04$	
γγ	0.227	$\pm0.005$	
Zγ	0.157	$\pm 0.009$	)
SS	0.025	$\pm0.001$	
μμ	0.0216	$\pm 0.0004$	)

CMS/



- Some rare processes also suffer from huge backgrounds, e.g.  $H \rightarrow cc$  or ss
- Others searches depend on small BR of boson decay, e.g.  $Z \rightarrow ee$ ,  $\mu\mu$  (BR=3.4%)

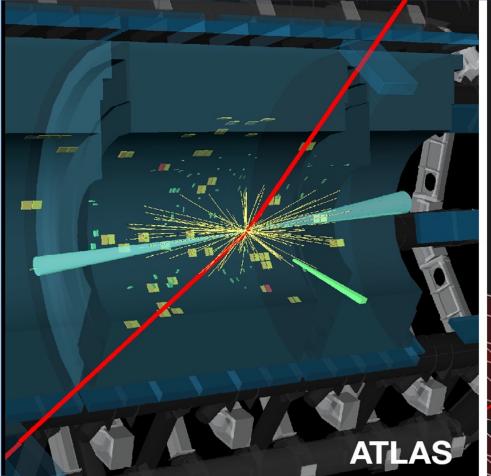
# $H \rightarrow Z\gamma$ Decays

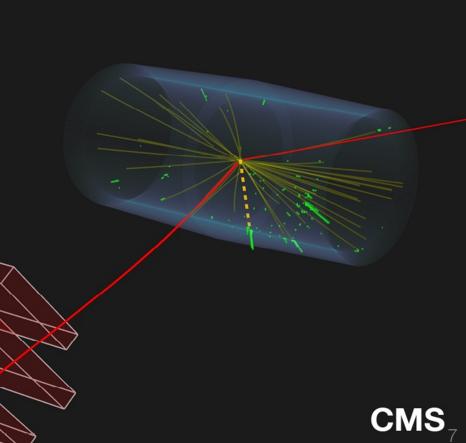


CMS



# $H \rightarrow Z\gamma Decays$





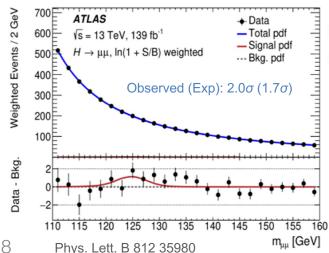


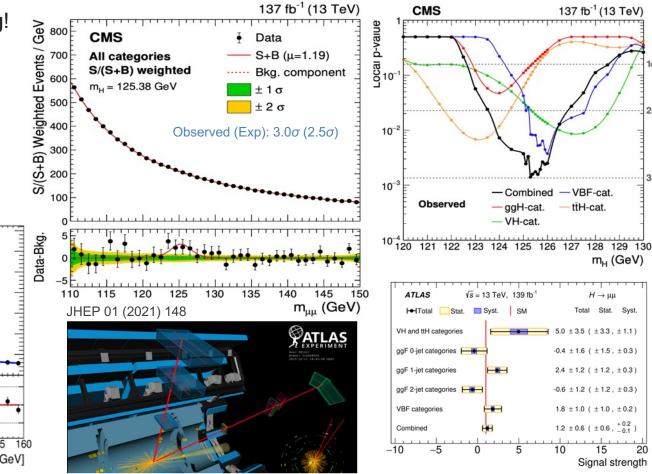
 $2\sigma$ 

•  $H \rightarrow \mu \mu$  is very challenging!  $BR(H \rightarrow \mu\mu) = 0.0216\%$ 

CMS

- Search for *narrow* m<sub>uu</sub> in all production modes, increase purity with MVAs
- CMS first evidence with a  $3\sigma$  signal significance!



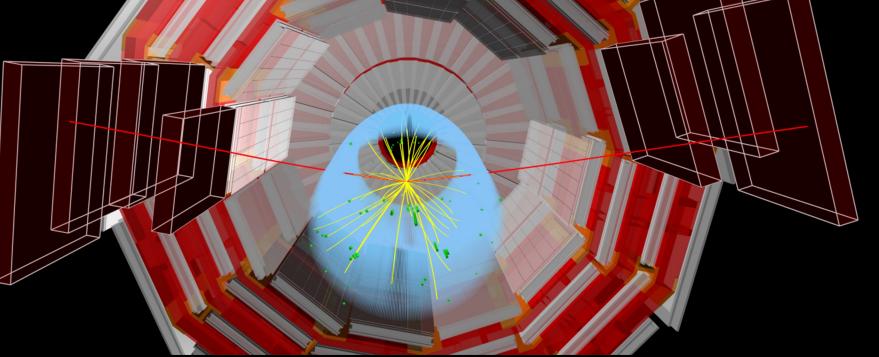






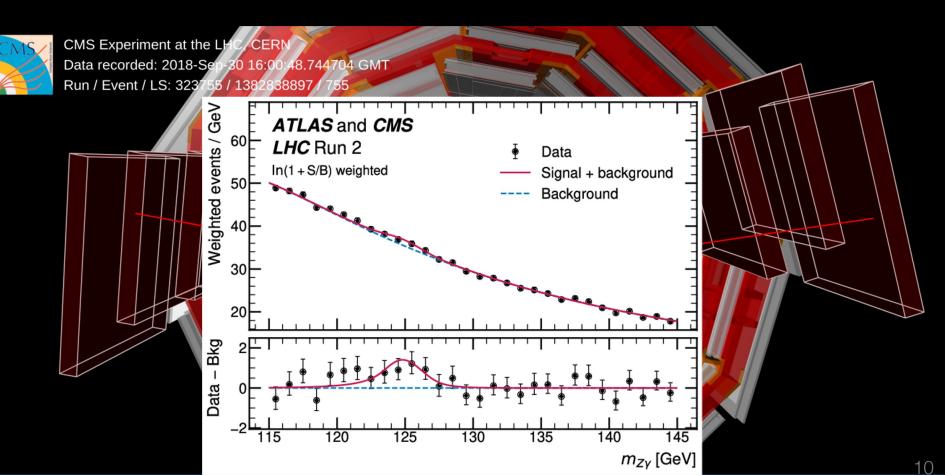


CMS Experiment at the LHC, CERN Data recorded: 2018-Sep-30 16:00:48.744704 GMT Run / Event / LS: 323755 / 1382838897 / 755







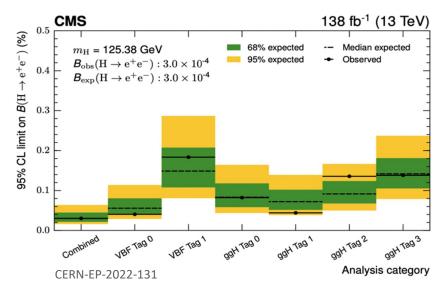


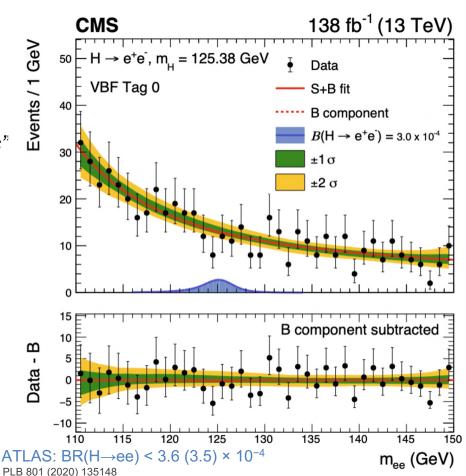


 Higgs decay to electron pairs has tiny BR(H → ee) = 0.0000005%

CMS

- Could be enhanced in BSM scenarios
- Search for narrow m<sub>ee</sub> resonance in all production modes and BDT purified "tag regions"
- BR(H $\rightarrow$ ee) < 3.0 × 10<sup>-4</sup> at 95%CL (exp and obs)



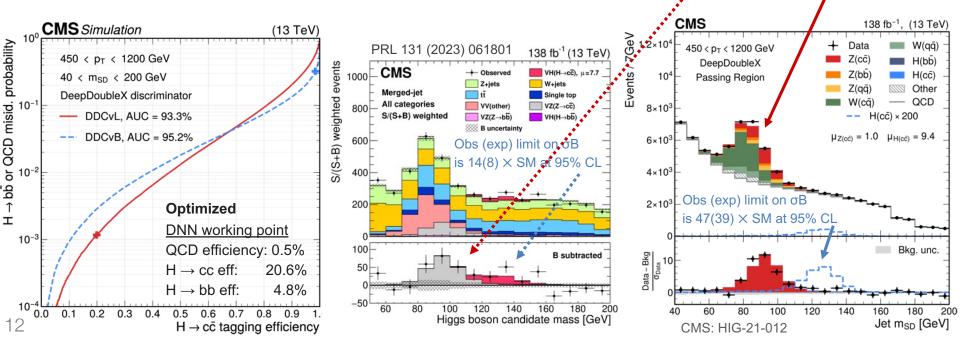




- Higgs→cc (charm) is a very challenging signature
  - Small branching ratio of 2.9% and difficulty of *c*-quark tagging given huge background rate
  - Search for ZH ( $H \rightarrow cc$ ) **and** boosted  $H \rightarrow cc$  events, reconstructed as a single large-radius jet
  - Use of  $Z \rightarrow cc$  data to validate the method yields a first observation of  $Z \rightarrow cc$  and  $Z \rightarrow cc$  + high  $p_T$  jets!

 $Higgs \rightarrow cc$ 

• Use of Deep Neural Network **charm-tagging** technique essential!



#### $H \rightarrow Meson + \gamma$

Direct

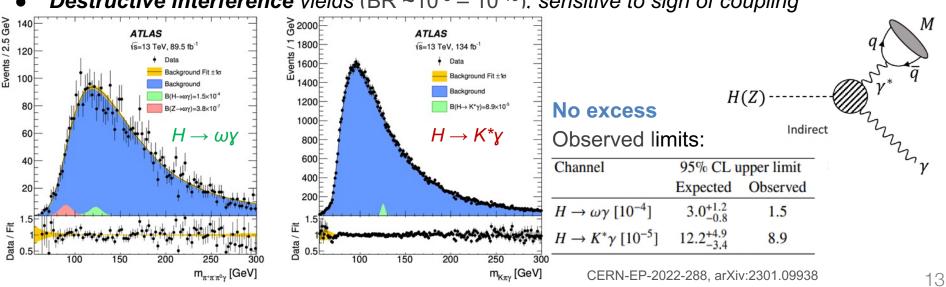
H(Z)

М

- Probe Higgs boson coupling to  $1^{st}$  and  $2^{nd}$  generation fermions in  $H \rightarrow Meson + \chi$
- Test Flavour conserving coupling to u and d quarks  $(H \rightarrow \omega \gamma)$ **Flavour violating** coupling to d and s quarks  $(H \rightarrow K^* y)$
- Standard Model prediction are driven by two contributions:
  - **Direct interaction:** Scales with Yukawa coupling 0
  - **Indirect interaction**:  $H \rightarrow \gamma^* \rightarrow M\gamma$  (M=meson) 0

CMS

**Destructive interference** vields (BR  $\sim 10^{-5} - 10^{-10}$ ). sensitive to sign of coupling

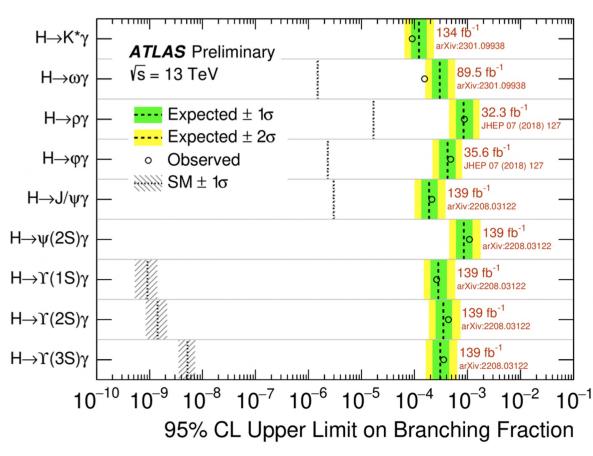


#### $H \rightarrow Meson + \gamma$

• Summary of searches using full Run-2 datasets:

CMS/

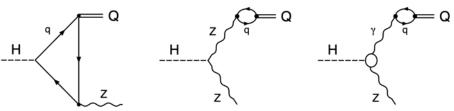
- b/c mesons:
   J/ψ, ψ(2S), Y(1S, 2S, 3S)
- Light/strange mesons:  $K^*$  and  $\omega$
- No significant excess, but many stringent limits!
- All results are statistically limited!
- Will make excellent use of larger Run-3 and HL-LHC datasets



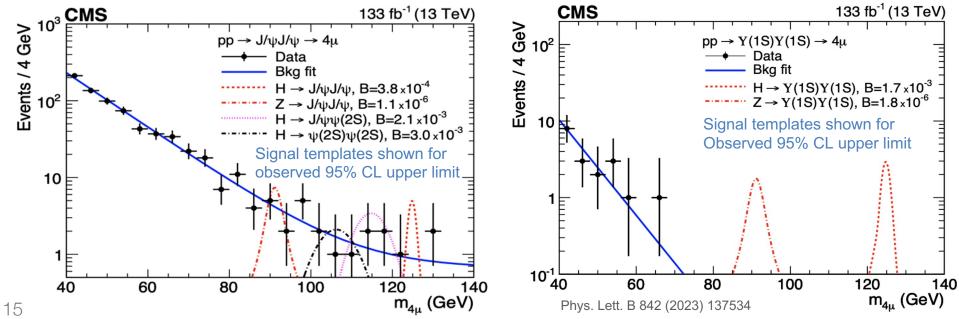
#### Higgs → Quarkonia

- Higgs → vector Quarkonia can be searched for in experimentally clean four-muon final state
- Tiny expected SM branching fractions ~10<sup>-9</sup> several orders of magnitude *below* sensitivity
- An observation would indicate BSM

CMS



#### No evidence for anomalously large BRs



#### $-LFV Higgs \rightarrow e\tau, \mu\tau$

W

 $\nu_{\ell'}$ 

95% CL

68% CL

 $Y_{\ell\tau}$ 

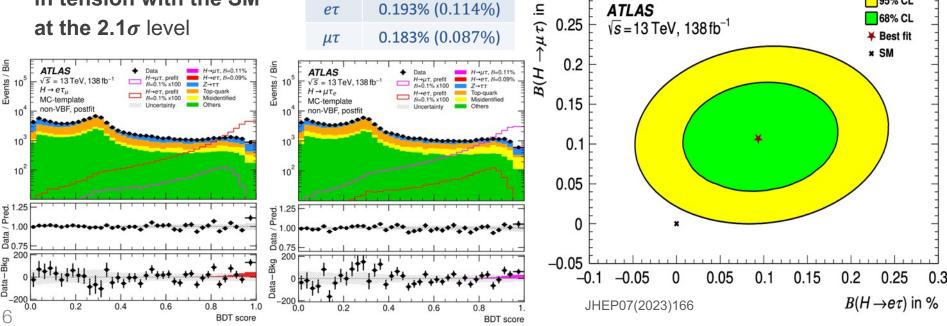
H

- Search for Lepton Flavour Violating (LFV) decay of Higgs boson, with  $H \rightarrow e\tau$  and  $H \rightarrow \mu\tau$  decay modes
- No significant excess is observed and limits are obtained:

ετ

• The simultaneous fit is in tension with the SM at the 2.1 $\sigma$  level

CMS,



Obs. (Exp.) 95% Upper

0.193% (0.114%)

0.3

0.25

ATLAS

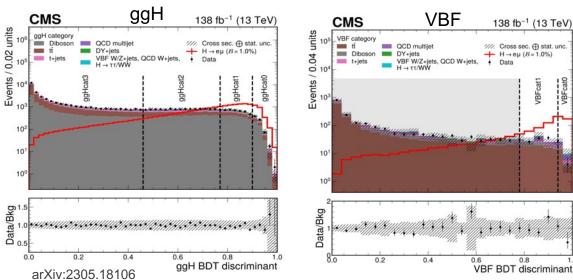
 $\sqrt{s} = 13 \text{ TeV}, 138 \text{ fb}^{-1}$ 

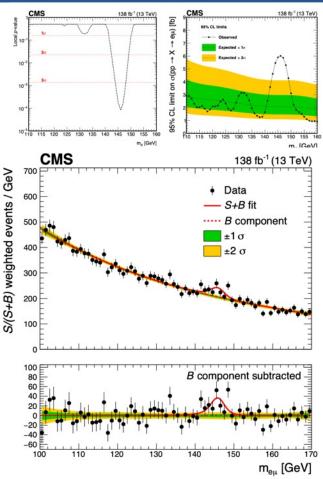
%



# LFV Higgs $\rightarrow e\mu$

- Search for LFV Decay  $H \rightarrow e\mu$
- Perform fit to m<sub>eµ</sub> distribution across production categories (ggH and VBF) purified by BDT categories
- BR(H→eµ) < 4.4 (4.7) × 10<sup>-5</sup> at 95%CL observed (exp)
- Scan for BSM Higgs boson: Modest excess in m<sub>eµ</sub> distribution at 146 GeV: 3.8σ (local), 2.8σ (global) significance.



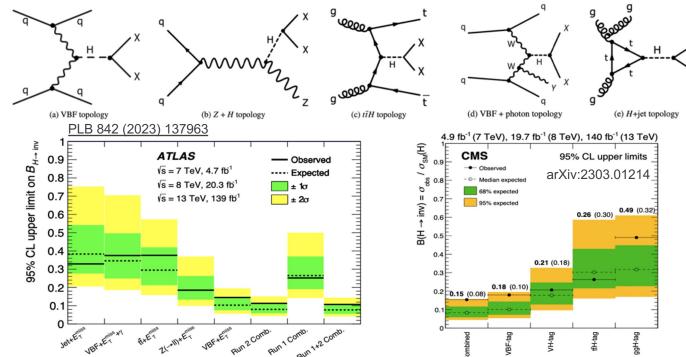


# Higgs to Invisible

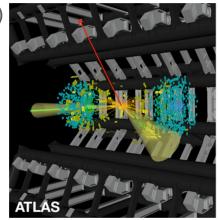
• Higgs-portal models to Dark Sector (e.g. Wilczek, Patt)

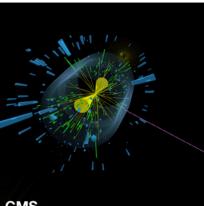
CMS,

- Within the SM, Higgs  $\rightarrow$  invisible decay is only H  $\rightarrow$  ZZ\*  $\rightarrow$  4v (BR ~ 0.1%)
- Searches for invisible Higgs decays, leading to missing p<sub>T</sub> + X



VBF + "missing p<sub>T</sub>"





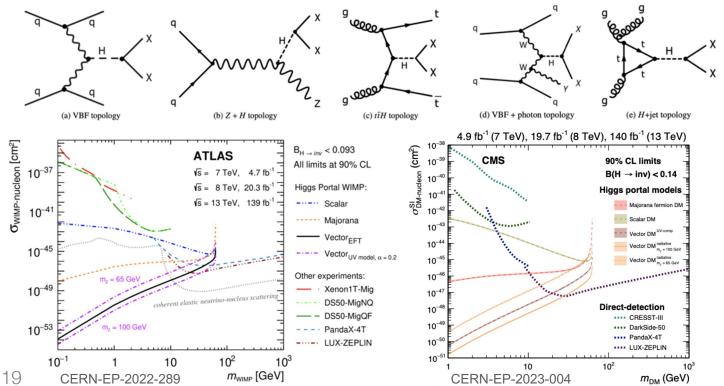
18 **Combined** ggH, VBF, VH, ttH; **CMS**: Br(H→inv) < 0.15 (0.08); **ATLAS**: Br(H→inv) < 0.107 (0.077) at 95% CL **CMS** 

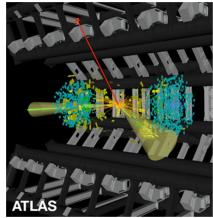
# Higgs to Invisible

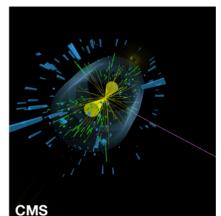
• Higgs-portal models to Dark Sector (e.g. Wilczek, Patt)

CMS

- Within the SM, **Higgs**  $\rightarrow$  invisible decay is only H  $\rightarrow$  ZZ\*  $\rightarrow$  4v (BR ~ 0.1%)
- Searches for invisible Higgs decays, leading to missing p<sub>T</sub> + X





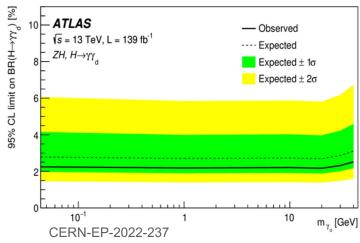


## Dark Photon, $H \rightarrow \chi \chi_d$

- Search for Higgs boson decaying into a (mixed) photon and a dark photon  $(\gamma_d)$
- Search in *ZH* production mode:

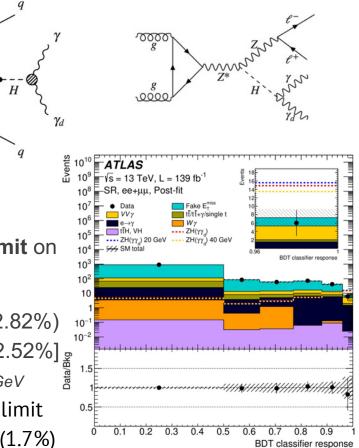
CMS.

- Leptonic Z  $\rightarrow$ II decay for clean signature
- BDT to enhance purity in signal region
- Search in *VBF* production mode:
  - Scan Higgs mass  $m_{H}$  instead of  $m\gamma_{d}$
- No excess of events above the SM expectation



Observed (exp) upper limit on BR(H  $\rightarrow \gamma \gamma_d$ ) at 95% CL:

- Massless γ<sub>d</sub>. \_\_\_\_
   Massive γ<sub>d</sub>: [2.19% 2.52%] <sup>m</sup>/<sub>g</sub>
   14 401 GeV VBF (m<sub>H</sub>=125 GeV) upper limit on BR(H  $\rightarrow \gamma \gamma_d$ ) of 1.8% (1.7%)



#### H→BSM Scalars

SM Higgs boson measurements are not the only part of LHC Higgs physics program

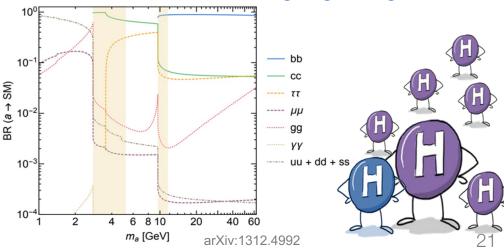
- Many BSM theories predict extended Higgs sector with additional scalars
- Allowed branching ratio to exotic decays O(10%)
- 2HDM is the simplest extension of the SM with:  $H^{\pm}$ , A(CP-odd), H and h (CP-even)
- Includes axion like particles (solution to strong CP problem Peccei–Quinn)
- Other extensions are considered e.g. MSSM, NMSSM, TRSM, 2HDM+S, g2HDM, ...

#### Example: 2HDM+S: 2 Higgs doublets models extended with one gauge singlet S

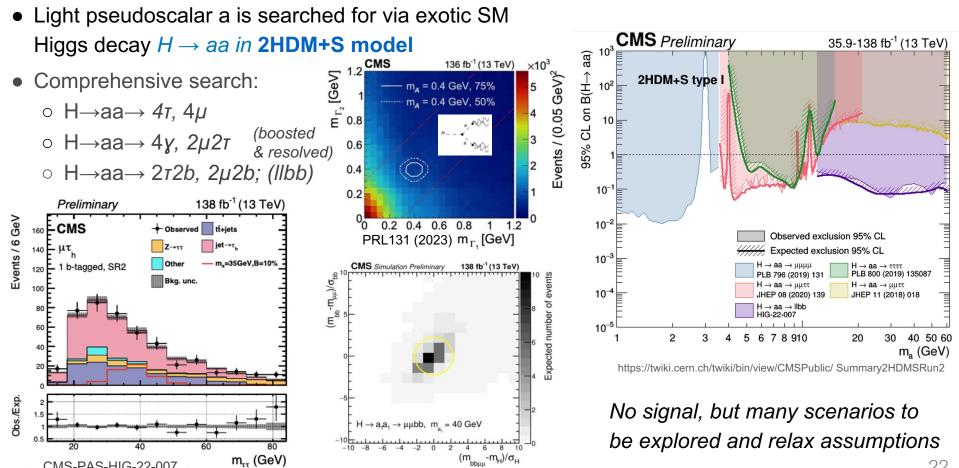
- $\circ~$  Predicts 7 physical Higgs states:
  - 3 Neutral scalars: h1 (SM like) h2, h3
  - 2 Charged Higgs: H<sup>±</sup>

CMS

- 2 pseudoscalars Higgs: A, a
- $\circ~$  The branching ratios depend on
  - 4 Model types (I-IV) of fermion couplings
  - $m_a$  (mass of the pseudoscalar)
  - $\tan \beta$  (ratio of the VEVs of the two doublets)



Higgs in 2HDM+S



CMS.

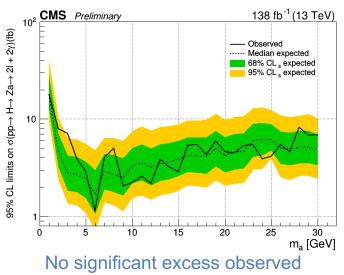
CMS-PAS-HIG-22-007

#### $H \rightarrow Za \rightarrow II\chi\chi$

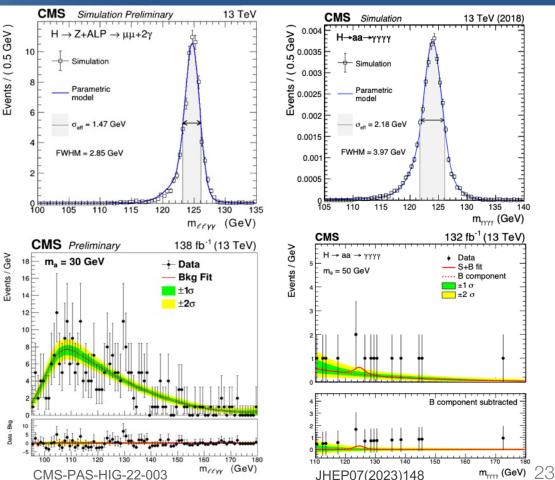
• Dedicated serach for  $H \rightarrow Za \rightarrow II\chi\chi$  and  $\chi\chi\chi\chi$ 

CMS

- Use high-resolution 4-body invariant mass m<sub>//yy</sub> and m<sub>yyyy</sub>
- Sensitivity optimized using **BDTs** for *ee* and  $\mu\mu$  channel



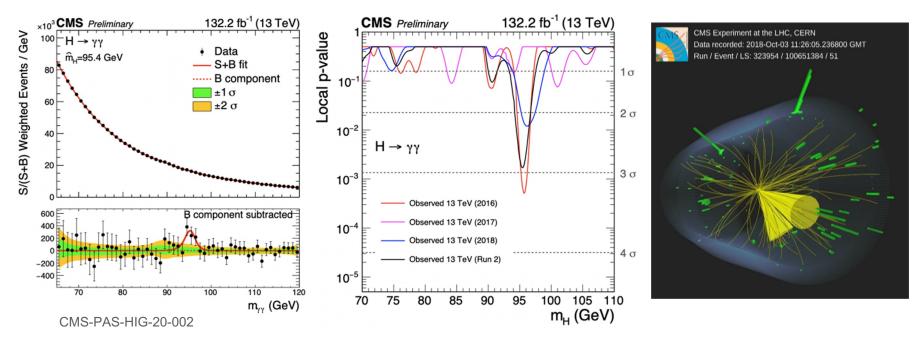
Events / GeV





## Low Mass Higgs Boson?

- CMS performed a diphoton resonance search for a standard model-like Higgs boson in the mass range between 70 and 110 GeV with full LHC Run 2 data.
- Modest excess at  $m_{yy} = 95.4 \text{ GeV}$  with 2.9 $\sigma$  local (1.3 $\sigma$  global) significance.
- Large "look elsewhere effect" in high-resolution di-photon channel



#### Low Mass Higgs Boson

Model independent

100

ATLAS Preliminary - Observed

90

 $\sqrt{s} = 13 \text{ TeV}, 140 \text{ fb}^{-1}$ 

80

10-

10-4

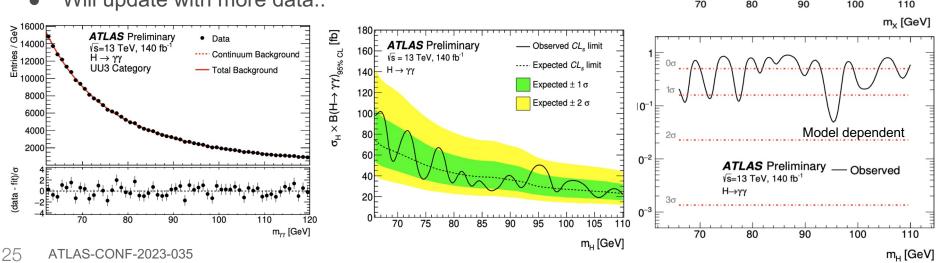
10-

70

- ATLAS performed recently a similar analysis in range between m<sub>vv</sub> = 66 and 110 GeV
- Perform two searches:

CMS.

- **Model independent** search for spin-0 particle (NWA) Ο (three categories based on photon conversion, UU, UC, CC)
- Model dependent search for light SM-like particle Ο (additionally use three BDT categories to purify sample)
- No significant excess observed in full Run-2 dataset
- Will update with more data...



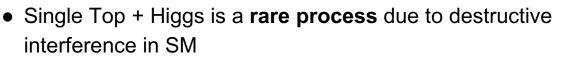
## Single Top + Higgs

**CMS** Preliminary

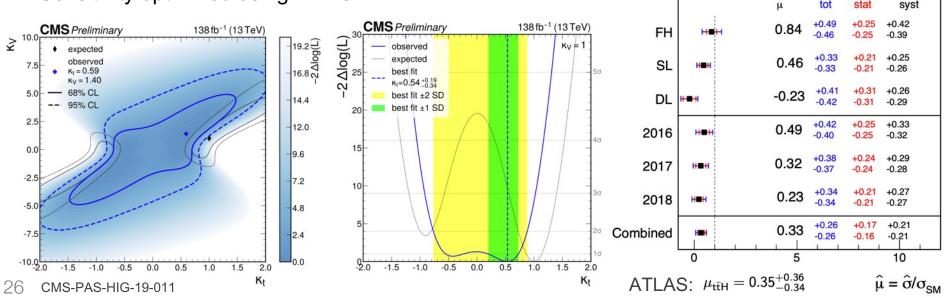
W

138 fb<sup>-1</sup> (13 TeV)





- Sensitive to absolute value of the top Yukawa coupling, the Higgs boson coupling to vector bosons and their **relative sign**
- New CMS result on search for tH production with  $H \rightarrow bb$
- Sensitivity optimized using ANNs





- Search for new couplings  $\rho_{tu}$ ,  $\rho_{tc}$  in Higgs decays in **same-sign tt final states**
- Possible explanations for the *electroweak baryogenesis* and muon *anomalous magnetic moment* in the generalized 2HDM model with FCNC H/A Higgs Yukawa couplings

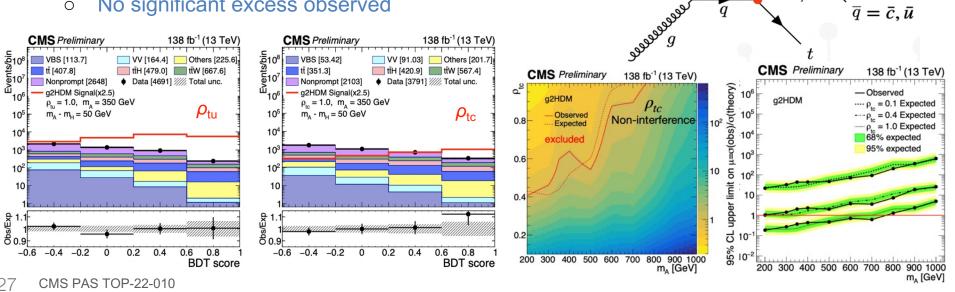
 $H/A \rightarrow tq$ 

tu,tc

H/A

 $\rho_{\rm tu,tc}$ 

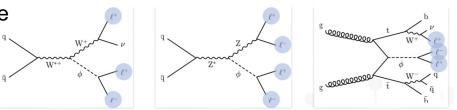
- Analysis optimized for top-pair final state (BDT)
  - Search same sign dilepton final states ( $e^{\pm}e^{\pm}$ ,  $\mu^{\pm}\mu^{\pm}$ ,  $e^{\pm}\mu^{\pm}$ ) Ο
  - No significant excess observed Ο

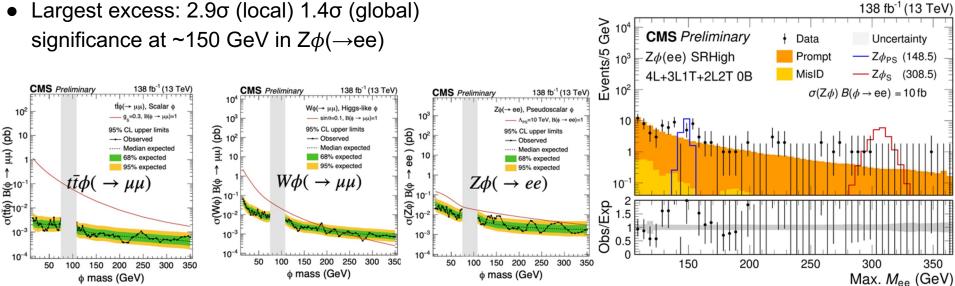






- First search for a **BSM spin-0 scalar**  $\phi$  resonance in association with bosons (V=W,Z) or tops (tt)
- Target multilepton final state with  $\phi \rightarrow dilepton$ Considering decays into all lepton flavors
- Largest excess:  $2.9\sigma$  (local)  $1.4\sigma$  (global) significance at ~150 GeV in  $Z\phi(\rightarrow ee)$





CMS-PAS-EXO-21-018

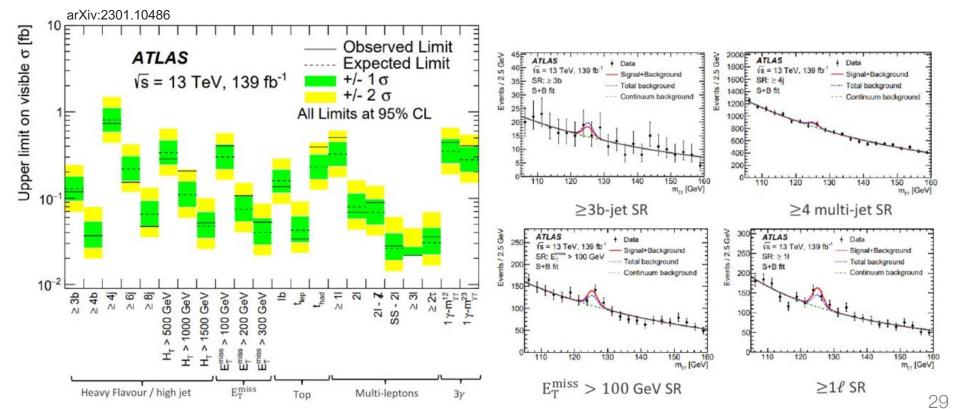
(qd) (ท่ท่ ←

# *Higgs* + *Anything*...

• Explores Higgs produced with additional objects to probe BSM in model independent way

CMS

 $\rightarrow$ I.e. Signature-based search for 125 GeV Higgs  $\rightarrow$  yy produced <u>with</u> additional objects





30

# Effective Field Theory (EFT)

- To maximize sensitivity to possible BSM effects in Higgs measurements, **ATLAS and CMS** target *well-defined kinematic regions*, split by production and decay mode The set of *σB* results are used in EFT interpretations
- Expand the SM Lagrangian with higher dimensional operators, to parameterize BSM effects

$$\mathscr{L}_{SMEFT} = \mathscr{L}_{SM} + \sum_{i} \frac{c_i^{d=6}}{\Lambda^2} \mathcal{O}^{d=6} + \sum_{i} \frac{c_i^{d=8}}{\Lambda^4} \mathcal{O}^{d=8} + \dots$$

Deviation from SM quantified by d=

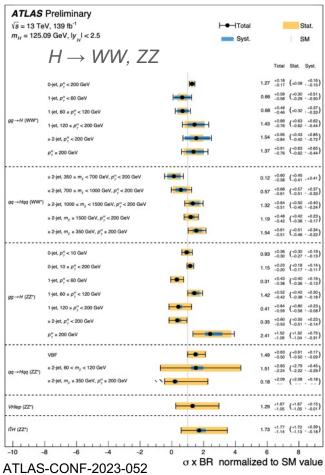
$$\sigma \propto \left| \mathcal{M}_{\rm SM} + \sum_{i} \frac{c_i^{(6)}}{\Lambda^2} \mathcal{M}_i^{(6)} + \sum_{i} \frac{c_i^{(8)}}{\Lambda^4} \mathcal{M}_i^{(8)} \right|^2$$
$$= \left| \mathcal{M}_{\rm SM} \right|^2 + \underbrace{\sum_{i} 2Re\left( \mathcal{M}_{\rm SM}^* \mathcal{M}_i^{(6)} \right) \frac{c_i^{(6)}}{\Lambda^2}}_{\text{Linear term (interference with SM)}} + \underbrace{\sum_{ij} 2Re\left( \mathcal{M}_i^{(6)*} \mathcal{M}_j^{(6)} \right) \frac{c_i^{(6)} c_j^{(6)}}{\Lambda^4}}_{\text{Quadratic term (purely BSM effect)}}$$

	$\begin{pmatrix} 1 \end{pmatrix}$
+O	$\left( \frac{1}{\Lambda^6} \right)$
	$\langle 1 \rangle$

	<sup>C</sup> H1.22	(11.1
uantified by d=6 Wilson coefficients		$(H^{\dagger}i$
$(2)$ $1^2$	$c_{He,11}$	$(H^{\dagger}i$
$\left[\sum_{i=1}^{n} c_i^{(8)} + c_i^{(8)}\right]^{-1}$	$C_{He,22}$	$(H^{\dagger}i$
$\sum \frac{i}{\Lambda^4} \mathcal{M}_i^{(\circ)}$	<i>c</i> <sub><i>He</i>,33</sub>	$(H^{\dagger}i$
	$c_{Hq}^{\scriptscriptstyle (1)}$	$(H^{\dagger}$
(6) $C_{i}^{(6)}$ $\sum_{a \in \mathbb{Z}} (x_{i}^{(6)*} + y_{i}^{(6)}) c_{i}^{(6)} c_{i}^{(6)}$	$c_{Hq}^{_{(3)}}$	$(H^{\dagger}i$
$_{\mathrm{M}}\mathcal{M}_{i}^{(6)}\left(\frac{c_{i}^{(6)}}{\Lambda^{2}}\right) + \sum 2Re\left(\mathcal{M}_{i}^{(6)*}\mathcal{M}_{j}^{(6)}\right)\frac{c_{i}^{(5)}c_{j}^{(5)}}{\Lambda^{4}}$	$c_{Hu}$	$(H^{\dagger}i$
	$c_{Hd}$	$(H^{\dagger}i$
ference with SM) Quadratic term (purely BSM effect)	$c_{HQ}^{\scriptscriptstyle (1)}$	$(H^{\dagger})$
	$c_{HQ}^{\scriptscriptstyle (3)}$	$(H^{\dagger}i)$
Neglect linear d=8 term, and	$c_{Ht}$	(H
model-dependent truncation	c <sub>Hb</sub>	$(H^{\dagger}$

on coefficient	Operator	Wilson coefficient	Operator
c <sub>H</sub>	$(H^{\dagger}H)^3$	$c_{Qq}^{(1,1)}$	$(\bar{Q}\gamma_{\mu}Q)(\bar{q}\gamma^{\mu}q)$
$c_{H\square}$	$(H^{\dagger}H)\Box(H^{\dagger}H)$	$c_{Qq}^{(1,8)}$	$(\bar{Q}T^a\gamma_\mu Q)(\bar{q}T^a\gamma^\mu q)$
$c_G$	$f^{abc}G^{a u}_\mu G^{b ho}_ u G^{c\mu}_ ho$	$c_{Qq}^{(3,1)}$	$(\bar{Q}\sigma^i\gamma_\mu Q)(\bar{q}\sigma^i\gamma^\mu q)$
CW	$\epsilon^{IJK}W^{I\nu}_{\mu}W^{J\rho}_{\nu}W^{K\mu}_{\rho}$	(3,8)	$(\bar{Q}\sigma^i T^a \gamma_\mu Q)(\bar{q}\sigma^i T^a \gamma^\mu q)$
C <sub>HDD</sub>	$\left(H^{\dagger}D^{\mu}H\right)^{*}\left(H^{\dagger}D_{\mu}H\right)$	$c^{Qq}_{qq}$	$(\bar{q}\sigma^i\gamma_\mu q)(\bar{q}\sigma^i\gamma^\mu q)$
C <sub>HG</sub>	$H^{\dagger}HG^{A}_{\mu u}G^{A\mu u}$	$c_{qq}$ $c_{tu}^{(1)}$	
CHB	$H^{\dagger}HB_{\mu u}B^{\mu u}$		$(\bar{t}\gamma_{\mu}t)(\bar{u}\gamma^{\mu}u)$
C <sub>HW</sub>	$H^{\dagger}HW^{I}_{\mu u}W^{I\mu u}$	$c_{tu}^{(8)}$	$(\bar{t}T^a\gamma_\mu t)(\bar{u}T^a\gamma^\mu u)$
C <sub>HWB</sub>	$H^\dagger  au^I H W^I_{\mu u} B^{\mu u}$	$c_{td}^{(1)}$	$(\bar{t}\gamma_{\mu}t)(\bar{d}\gamma^{\mu}d)$
$c_{Hl,11}^{_{(1)}}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\overline{l}_{1}\gamma^{\mu}l_{1})$	$c_{td}^{(8)}$	$(\bar{t}T^a\gamma_\mu t)(\bar{d}T^a\gamma^\mu d)$
$c_{Hl,22}^{(1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{l}_{2}\gamma^{\mu}l_{2})$	$c_{Qu}^{(1)}$	$(\bar{Q}\gamma_{\mu}Q)(\bar{u}\gamma^{\mu}u)$
$c_{Hl,33}^{(1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{l}_{3}\gamma^{\mu}l_{3})$	$c_{Qu}^{(8)}$	$(\bar{Q}T^a\gamma_\mu Q)(\bar{u}T^a\gamma^\mu u)$
$c_{Hl,11}^{(3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\overline{l}_{1}\tau^{I}\gamma^{\mu}l_{1})$	$c_{Qd}^{(1)}$	$(\bar{Q}\gamma_{\mu}Q)(\bar{d}\gamma^{\mu}d)$
$c_{Hl,22}^{_{(3)}}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{l}_{2}\tau^{I}\gamma^{\mu}l_{2})$	$c_{Qd}^{(8)}$	$(\bar{Q}T^a\gamma_\mu Q)(\bar{d}T^a\gamma^\mu d)$
$c_{Hl,33}^{(3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{l}_{3}\tau^{I}\gamma^{\mu}l_{3})$	$c_{tq}^{(1)}$	$(\bar{q}\gamma_{\mu}q)(\bar{t}\gamma^{\mu}t)$
$c_{He,11}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{e}_{1}\gamma^{\mu}e_{1})$	$c_{tq}^{_{(8)}}$	$(\bar{q}T^a\gamma_\mu q)(\bar{t}T^a\gamma^\mu t)$
$c_{He,22}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{e}_{2}\gamma^{\mu}e_{2})$	C <sub>eH,22</sub>	$(H^{\dagger}H)(\bar{l}_{2}e_{2}H)$
<i>c<sub>He,33</sub></i>	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{e}_{3}\gamma^{\mu}e_{3})$	СеН,33	$(H^{\dagger}H)(\bar{l}_{3}e_{3}H)$
$c_{Hq}^{(1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{q}\gamma^{\mu}q)$	CuH	$(H^{\dagger}H)(\bar{q}Y_{u}^{\dagger}u\widetilde{H})$
$c_{Hq}^{_{(3)}}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{q}\tau^{I}\gamma^{\mu}q)$	$c_{tH}$	$(H^{\dagger}H)(\bar{Q}\widetilde{H}t)$
c <sub>Hu</sub>	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{u}_{p}\gamma^{\mu}u_{r})$	C <sub>bH</sub>	$(H^{\dagger}H)(ar{Q}Hb)$
$C_{Hd}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{d}_{p}\gamma^{\mu}d_{r})$	$c_{tG}$	$(\bar{Q}\sigma^{\mu\nu}T^{A}t)\widetilde{H}G^{A}_{\mu\nu}$
$c_{HQ}^{(1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{Q}\gamma^{\mu}Q)$	CtW	$(\bar{Q}\sigma^{\mu\nu}t)\tau^I \tilde{H} W^I_{\mu\nu}$
$c_{HQ}^{\scriptscriptstyle (3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{Q}\tau^{I}\gamma^{\mu}Q)$ $(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{t}\gamma^{\mu}t)$	$C_{tB}$	$(\bar{Q}\sigma^{\mu\nu}t)\widetilde{H}B_{\mu\nu}$
c <sub>Ht</sub> c <sub>Hb</sub>	$(H^{\dagger}i \overleftrightarrow{D}_{\mu}H)(t\gamma^{\mu}t) (H^{\dagger}i \overleftrightarrow{D}_{\mu}H)(\bar{b}\gamma^{\mu}b)$	c <sub>11,1221</sub>	$(\bar{l}_1\gamma_\mu l_2)(\bar{l}_2\gamma^\mu l_1)$
c <sub>Hb</sub>	$(H^{\dagger}i\widetilde{D}_{\mu}H)(\bar{b}\gamma^{\mu}b)$	<i>c</i> <sub><i>ll</i>,1221</sub>	$(\bar{l}_1\gamma_\mu l_2)(\bar{l}_2\gamma^\mu l_1)$





	Preliminary	H	Stat.
	eV, 139 fb <sup>-1</sup>	Syst.	SM
$m_{H} = 125$	.09 GeV, ly 1 < 2.5		SM
	$H \rightarrow \gamma \gamma, Z \gamma, t t$	H, tH	otal Stat. Syst.
	0-jet, p <sup>µ</sup> <sub>7</sub> < 10 GeV ►	0.66	0.27 (+0.24 , +0.12
	0-jet, 10 = $p_{\gamma}^{H}$ < 200 GeV	1.24	0.18 (+0.15 - +0.10 0.17 (+0.150.00
	1-jet, p <sup>µ</sup> <sub>γ</sub> < 60 GeV	1.16	0.39 (+0.36 , +0.5) 0.38 (+0.36 , -0.1)
	1-jet, 60 ≤ p <sup>H</sup> <sub>T</sub> < 120 GeV	1.14	0.40 (+0.33 . +0.20 0.36 (+0.330.10
	1-jet, 120 s p <sup>N</sup> <sub>7</sub> < 200 GeV	0.93	0.57 (+0.53 +0.20
gg →H (үү)	≥ 2-jet, m <sub>j</sub> < 350 GeV, p <sup>H</sup> <sub>2</sub> < 120 GeV	0.58	0.56 (+0.53 +0.11
	$\ge 2$ -jet, $m_j < 350$ GeV, $120 \le p_{\gamma}^{\prime\prime} < 200$ GeV	H 1.31	0.50 (+0.48 +0.15
	≥ 2-jet, m <sub>j</sub> ≥ 350 GeV, p <sup>H</sup> <sub>1</sub> < 200 GeV	1.09	0.95 (-0.91 -0.30
	200 s p <sup>v</sup> <sub>r</sub> < 300 GeV	1.56	0.45 (+0.41 +0.10
	300 ≤ p <sup>v</sup> <sub>r</sub> < 450 GeV	0.17	0.56 (+0.54 +0.54
	p <sup>+</sup> <sub>7</sub> ≥ 450 GeV		1.47 (*1.42 *0.4 1.18 (-1.15 * -0.2
	« 1-jet and VH-veto		0.96 (-0.90 +0.30
	> 2-jet, VH-had	0.21	0.74 (+0.72 +0.54
	$\ge 2$ -jet, 350 $\le m_j < 700$ GeV, $p_{\gamma}^{\vee} < 200$ GeV	1.28	0.80 (-0.61 +0.51
qq→Hqq (ŗỵ)	≥ 2-jet, 700 ≤ $m_j$ < 1000 GeV, $p_j^H$ < 200 GeV	1.47	0.84 ( 0.72 +0.43
	$\ge 2$ -jet, $m_j \ge 1000$ GeV, $p_{\gamma}^{N} < 200$ GeV	H 1.31	0.46 (+0.36 +0.21
	≥ 2-jet, 350 ≤ $m_g$ < 1000 GeV, $p_g^H$ ≥ 200 GeV		0.74 (+0.73 +0.13
	≥ 2-jet, m <sub>j</sub> ≥ 1000 GeV, p <sup>N</sup> <sub>7</sub> ≥ 200 GeV		0.67 (+0.61 +0.20 0.57 (-0.520.25
	p <sup>v</sup> <sub>2</sub> < 150 GeV	1.75	0.82 (+0.80 +0.50 0.73 (-0.720.00
qq→Hiv (yy)	p <sup>v</sup> <sub>7</sub> = 150 GeV		
		-0.64	0.88 (+0.87 ,+0.13
ggiqq→Hll/vv (	ry) p <sup>v</sup> <sub>7</sub> < 150 GeV		
	p, # 100 GeV		0.92 -0.91 + -0.16
	<i>p</i> <sup><i>H</i></sup> <sub>7</sub> < 60 GeV	0.83	0.82 (+0.81 +0.11
tîH ( <sub>111</sub> )	60 s p <sup>v</sup> <sub>7</sub> < 120 GeV ►		0.60 (+0.59 +0.00
	120 ≈ p <sup>N</sup> <sub>7</sub> < 200 GeV	0.65	0.64 (+0.63 +0.15
	200 « p <sup>N</sup> <sub>r</sub> < 300 GeV	1.23	0.81 (-0.80 +0.11
	<i>p</i> <sup>H</sup> <sub>7</sub> ≥ 300 GeV	1.17	0.96 (+0.95 +0.16 0.75 (-0.74 + -0.15
1H (17)	·····	2.06	4.13 (+3.94 +1.25 3.27 (-3.140.96
			0.97 (+0.88 -0.4
H(Z Y)		2.05	0.97 (+0.88 +0.4 0.93 (-0.870.3
	-6 -4 -2 0	2 4 6 8	10

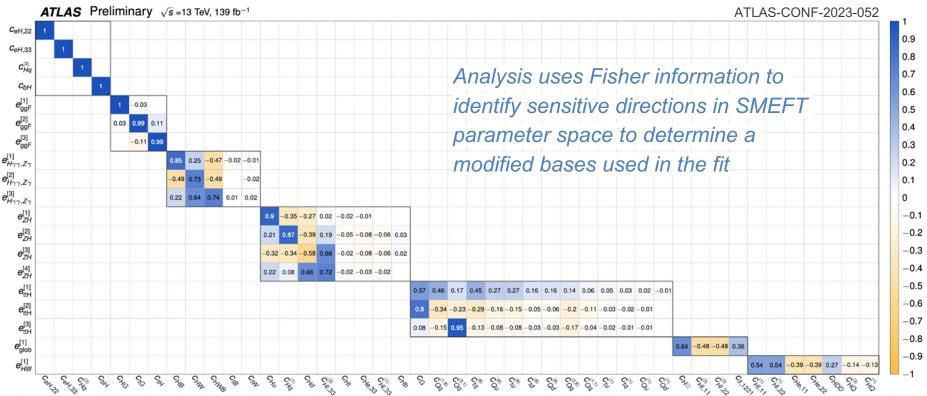
√s = 13 TeV m <sub>H</sub> = 125.09	, 139 fb <sup>-1</sup> 9 GeV, ly <sub>H</sub> I < 2.5		Syst.	Stat.
	$H \rightarrow \tau \tau$ , b	b, μμ		
				Total Stat. Syst.
	1-jet, 120 s $p_{\gamma}^{n}$ < 200 GeV	HH		.19 +0.68 (+0.41 +0.55
	$\ge$ 1-jet, $m_j < 350~{\rm GeV}, 0 \le p_{\gamma}^{\rm v} < 60~{\rm Ge}$		0	31 ± 0.94 (±0.56 , ±0.75
99 →H (++)	$\approx$ 2-jet, $m_g$ < 350 GeV, 120 $\propto p_{\gamma}^{\rm H}$ < 200	Gev H	0	.60 +0.87 -0.78 (+0.54 , +0.68
	a 2-jet, $m_j$ a 350 GeV, $\rho_{\gamma}^{\prime\prime}$ < 200 GeV	⊢ <b>−</b>	• · · · · · · · · · · · · · · · · · · ·	.55 +2.33 (+1.31 +1.92 -1.56 (-1.30 + -1.36
	$200 \le p_{\gamma}^{\prime\prime} < 300 \; {\rm GeV}$	H	1	.02 +0.55 (+0.31 +0.46 -0.41 (-0.30 -0.28
	ρ <sup>µ</sup> <sub>7</sub> ≥ 300 GeV		1	27 +0.77 (+0.46 +0.61 -0.45 + -0.31
	≥ 2-jet, 60 ≤ m, ≤ 120 GeV			.97 *0.66 (*0.55 *0.36 -0.63 (*0.55 *0.36
φα→Ηφα (ττ)	≥ 2-jet, m <sub>j</sub> ≥ 350 GeV	-		.80 +0.23 (-0.53 +-0.34 .80 +0.23 (+0.17 +0.15 -0.20 (+0.16 +-0.12
tĨH (m)		H <b></b> -I	1	24 +1.35 (+1.11 +0.77 -1.12 (-0.96 + -0.55
qq →Hqq (bb)			0	.98 +0.39 (=0.33 , +0.20 -0.38 (=0.33 , -0.18
	$150 \le \rho_{\gamma}^{\gamma} < 250 \text{ GeV}$	H <b>ee</b> H	0	79 +0.50 (+0.34 +0.37 -0.49 (-0.33 +0.36
qq →HIv (bb)	250 s $\rho_{\gamma}^{\gamma}<400~{\rm GeV}$		1	10 +0.41 (+0.35 +0.20 -0.38 (-0.34 + -0.18
	<i>p</i> <sup>v</sup> <sub>7</sub> ≥ 400 GeV	H	1	.50 +0.93 (+0.78 +0.51 -0.83 (-0.72 + -0.41
	$75 \le p_{\gamma}^{\gamma} < 150 \text{ GeV}$	H <b></b> -I	0	.90 +0.71 (±0.47 , +0.52
gg/qq →Hll/vv (bb)	$150 \le p_{\tau}^{\vee} < 250 \; \mathrm{GeV}$	H <b>H</b>	1	.13 -0.37 (+0.27 , +0.25
99r94	$250 \le p_{\tau}^{\vee} < 400 \; \mathrm{GeV}$		1	.01 -0.39 (+0.35 . +0.17
	$p_T^{\nu} \approx 400 \text{ GeV}$	H <b></b> H	0	29 +0.92 (+0.76 +0.53
	p <sup>H</sup> <sub>2</sub> < 120 GeV			
	$p_{\gamma}^{\nu} < 120 \text{ GeV}$ 120 s $p_{\gamma}^{\nu} < 200 \text{ GeV}$			10 +1.05 (+0.48 , +0.94 -0.99 (+0.48 , -0.87 0.22 +1.02 (+0.72 +0.73 -0.07 +0.75
tĨH (bb)	$200 \le p_{\gamma}^{H} < 300 \text{ GeV}$			98 +0.91 (+0.71 +0.57 -0.86 (-0.680.53
	$300 \le p_{\gamma}^{\prime\prime} < 450 \text{ GeV}$			0.23 +0.73 (+0.58 +0.45 -0.72 (-0.54 + -0.47
	p <sup>H</sup> <sub>7</sub> > 450 GeV ⊢		4	0.19 +1.48 (+1.06 +1.03 -1.40 (-0.911.06
gg →H, tĨH (μμ)			0	.54 × 0.85 (×0.83 , +0.22
qq→Hqq, VH (µµ)				23 +1.32 (+1.28 +0.31 -1.24 (-1.22 + -0.25
-8 -	-6 -4 -2	0 2	4 6	8
<i>gg →H</i> (bb)	450 ≤ p <sup>H</sup> <sub>7</sub> < 650 GeV	•••••		4.2 +6.4 (+5.0 +3.9 -5.0 +7.9 8.6 +14.7 (+10.3 +10.4 -12.6 (+10.2 +-7.4
1				8.6 -12.6 -10.2 -7.4

31

Data is currently insufficient to constrain all Wilson Coefficients simultaneously!

CMS

 $\rightarrow$  Motivate a modified basis reducing the number of free parameters in the fit (39 $\rightarrow$ 19 shown here)



ATLAS Preliminary ATLAS-CONF-2023-052  $\sqrt{s}$  =13 TeV, 139 fb<sup>-1</sup>,  $m_H$  = 125.09 GeV SMEFT  $\Lambda = 1$  TeV  $H \rightarrow \gamma\gamma$  $H \rightarrow Z\gamma$  $H \rightarrow WW^* \rightarrow l\nu l\nu$ Expected contribution production decay H  $\rightarrow 77^{*} \rightarrow 4$ 0.4  $H \rightarrow b\bar{b}$  $H \rightarrow \tau \tau$ The production and decay channels 0.2  $H \rightarrow \mu\mu$ agF contributing to the sensitivity VBF 0.8 WH 0.6 ZH 0.4 t tH tH 0.2 8 inclusive 0.32 10<sup>1</sup> Linear (obs.) Linear (exp.) ö) [TeV] Symmetrized uncertainty ( $\sigma$ ) 10<sup>0</sup> robed Scale (A/ 10-1 10-10 Linear (obs.) Parameter value scaled by symmetrized uncertainty ( $c'/\sigma$ )  $p_{SM} = 94.5\%$ Linear (exp.) The fit is performed with 19 free Best Fit 68 % CL ··· 95 % Cl parameters profiled simultaneously! No significant deviation from SM C200 0/1/ 0/1

CMS

33

# Summary and Outlook

- Rare Higgs boson processes offer fruitful ways to test the SM and search for BSM physics in a complementary way to the more *established* Higgs processes.
- Many results available with Run-2 dataset of ~140 fb<sup>-1</sup> at 13 TeV (not all could be shown)
   ATLAS: https://twiki.cern.ch/twiki/bin/view/AtlasPublic/HiggsPublicResults
  - o CMS: https://cms.cern/physics/cms-higgs-results
- In many cases, improved sensitivity leads to tighter constraints beyond simple luminosity scaling, thanks to advanced analysis techniques
- Possible BSM physics is being constrained in Effective Field Theory Interpretations using Higgs boson data from ATLAS and CMS
- Run 3 physics program has just started!

CMS/

- Much larger dataset will be available soon and beyond (HL-LHC)
- Rare Higgs processes promise to become even more interesting as the luminosity increases and observation comes within reach.
- New Physics could appear in the coming years...





# Backup



Demonstrate applicability of constraints using 2HDM benchmark

