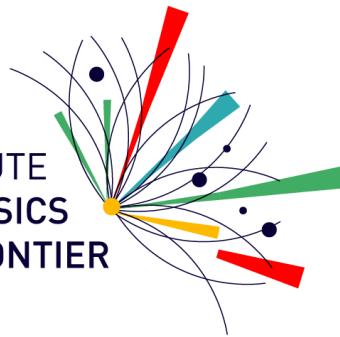


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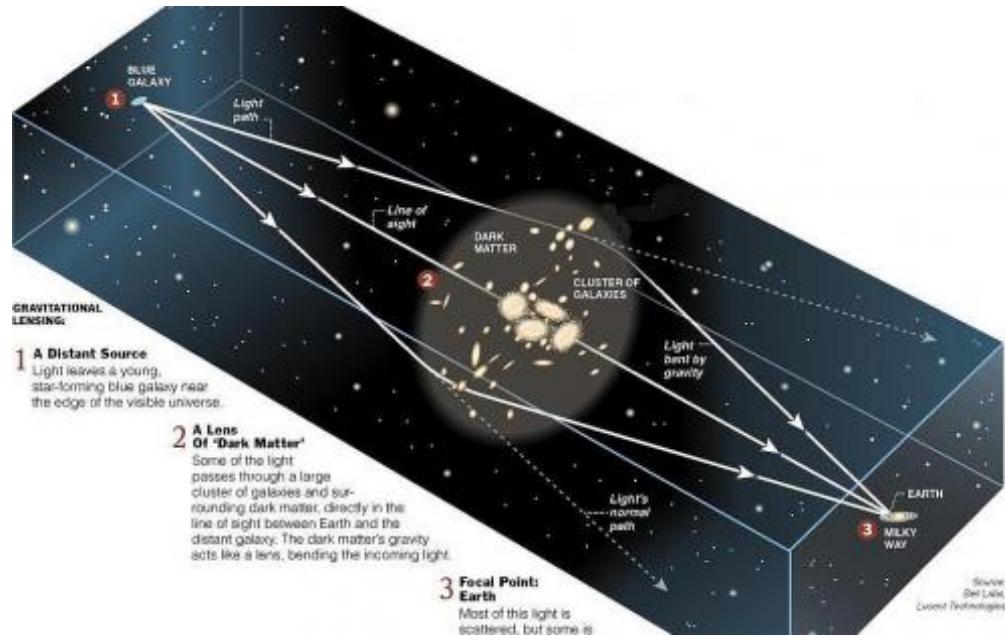
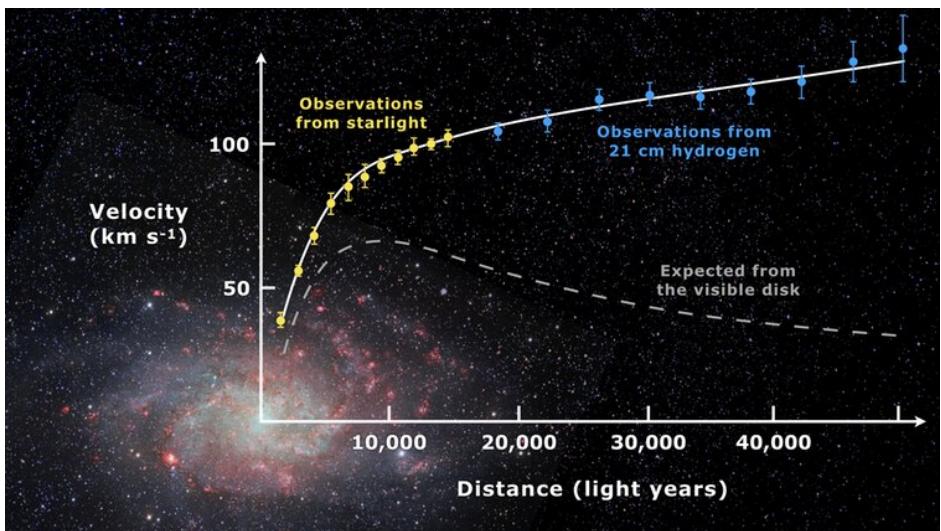
Particles and Space

Marco Aurelio Díaz

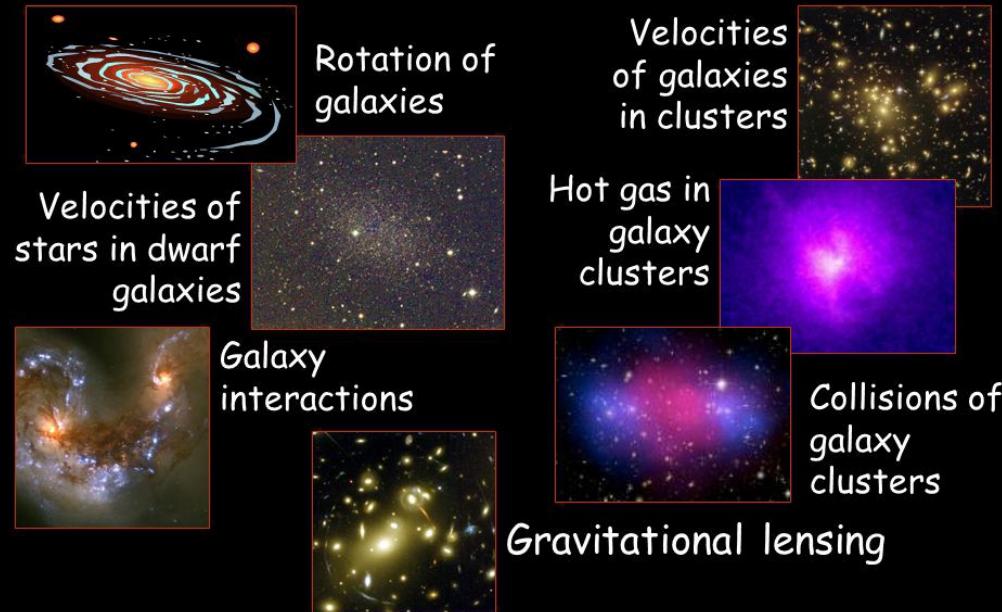
PUC



Dark Matter

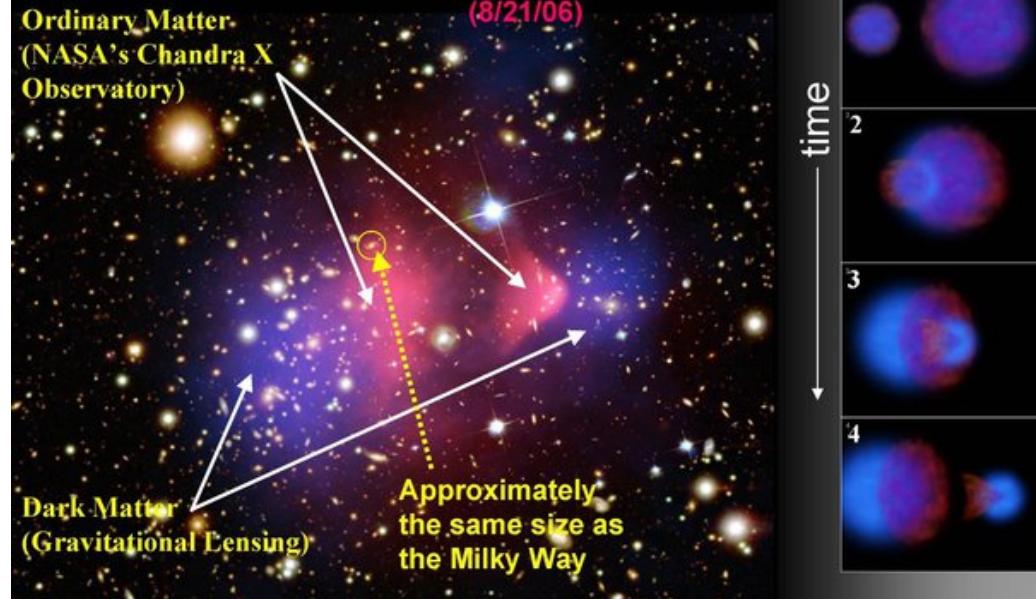


Evidence for Dark Matter



Cosmic Collision of 2 Galaxy Clusters splitting normal matter and dark matter apart

— Another Clear Evidence of Dark Matter —



Neutrinos

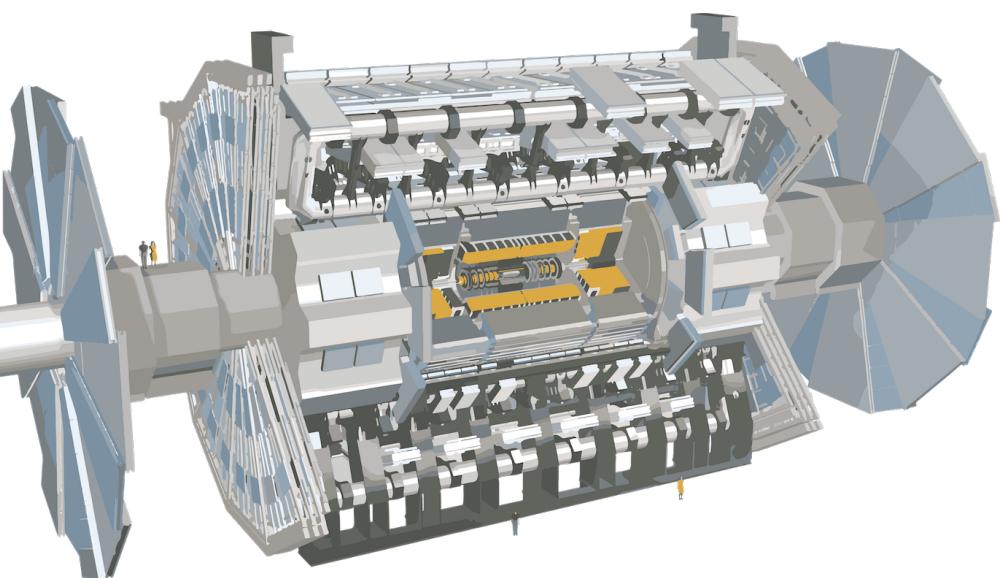
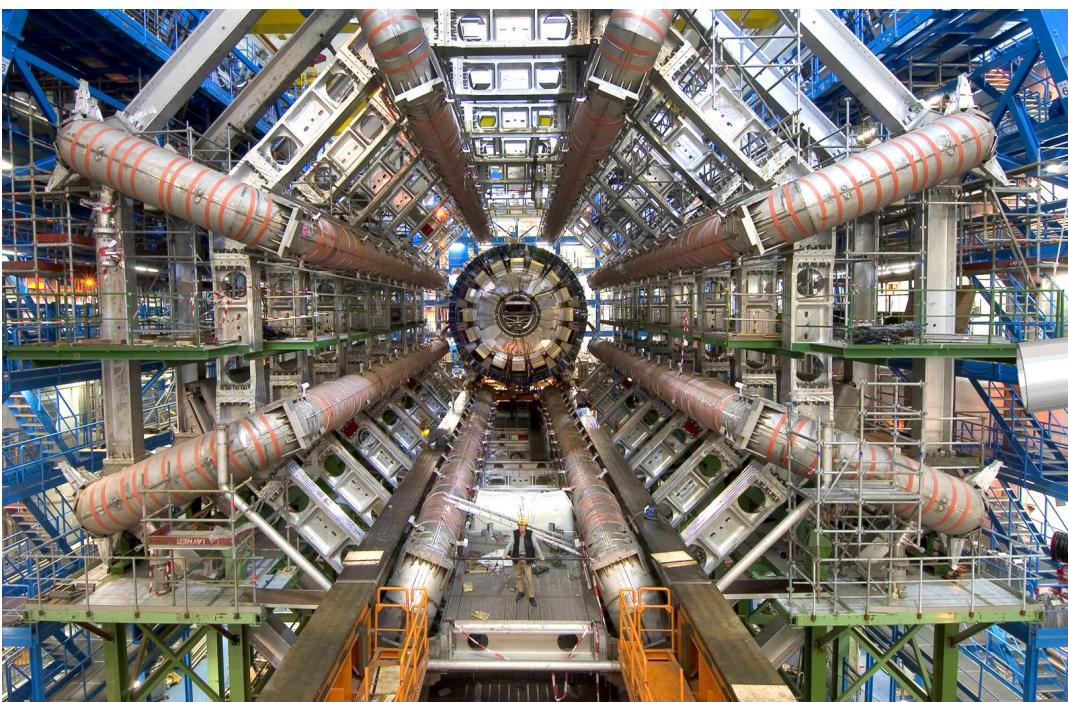
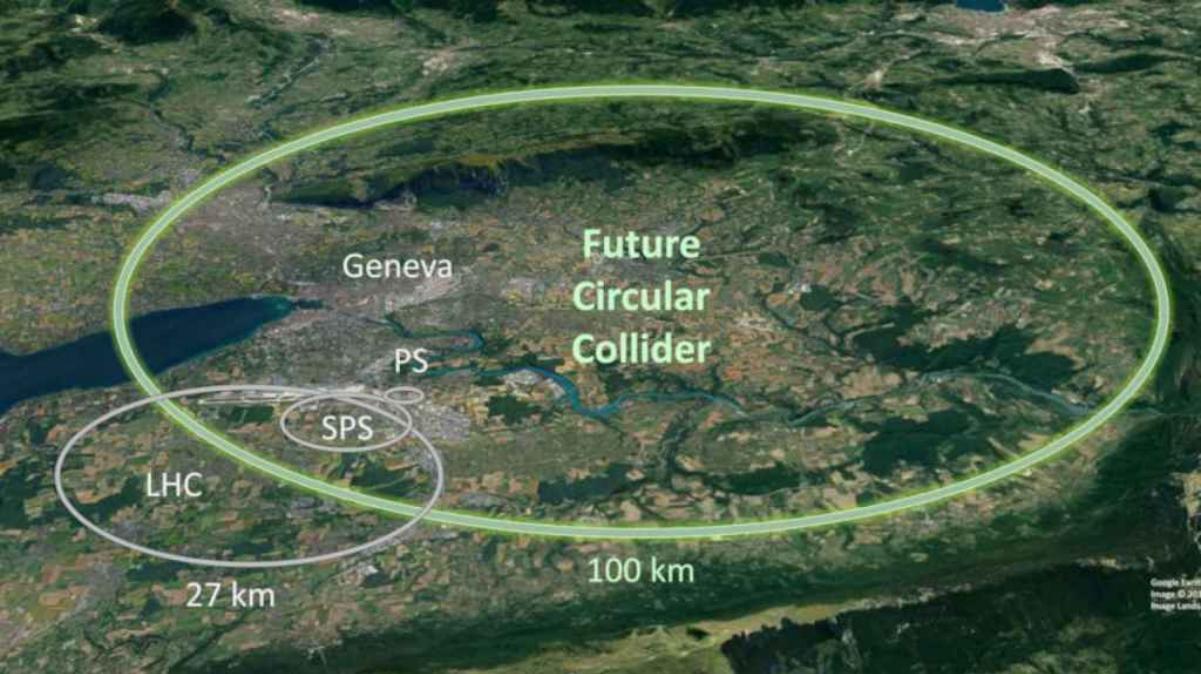
parameter	best fit $\pm 1\sigma$	2σ range	3σ range
$\Delta m_{21}^2 [10^{-5}\text{eV}^2]$	$7.50^{+0.22}_{-0.20}$	7.12–7.93	6.94–8.14
$ \Delta m_{31}^2 [10^{-3}\text{eV}^2]$ (NO)	$2.55^{+0.02}_{-0.03}$	2.49–2.60	2.47–2.63
$ \Delta m_{31}^2 [10^{-3}\text{eV}^2]$ (IO)	$2.45^{+0.02}_{-0.03}$	2.39–2.50	2.37–2.53
$\sin^2 \theta_{12}/10^{-1}$	3.18 ± 0.16	2.86–3.52	2.71–3.69
$\theta_{12}/^\circ$	34.3 ± 1.0	32.3–36.4	31.4–37.4
$\sin^2 \theta_{23}/10^{-1}$ (NO)	5.74 ± 0.14	5.41–5.99	4.34–6.10
$\theta_{23}/^\circ$ (NO)	49.26 ± 0.79	47.37–50.71	41.20–51.33
$\sin^2 \theta_{23}/10^{-1}$ (IO)	$5.78^{+0.10}_{-0.17}$	5.41–5.98	4.33–6.08
$\theta_{23}/^\circ$ (IO)	$49.46^{+0.60}_{-0.97}$	47.35–50.67	41.16–51.25
$\sin^2 \theta_{13}/10^{-2}$ (NO)	$2.200^{+0.069}_{-0.062}$	2.069–2.337	2.000–2.405
$\theta_{13}/^\circ$ (NO)	$8.53^{+0.13}_{-0.12}$	8.27–8.79	8.13–8.92
$\sin^2 \theta_{13}/10^{-2}$ (IO)	$2.225^{+0.064}_{-0.070}$	2.086–2.356	2.018–2.424
$\theta_{13}/^\circ$ (IO)	$8.58^{+0.12}_{-0.14}$	8.30–8.83	8.17–8.96
δ/π (NO)	$1.08^{+0.13}_{-0.12}$	0.84–1.42	0.71–1.99
$\delta/^\circ$ (NO)	194^{+24}_{-22}	152–255	128–359
δ/π (IO)	$1.58^{+0.15}_{-0.16}$	1.26–1.85	1.11–1.96
$\delta/^\circ$ (IO)	284^{+26}_{-28}	226–332	200–353

From:

JHEP 02, 071 (2021),
 P.F. de Salas, D.V. Forero,
 S. Gariazzo,
 P. Martínez-Miravé,
 O. Mena, C.A. Ternes,
 M. Tórtola, J.W.F. Valle.



Colliders



Particles from Space



Our Sun (one of a billion stars in the Milky Way), and our solar system is located NOT in the center, and NOT at the edge, but about 2/3 out on a spiral arm.



Scotogenic Model

Due to E. Ma, *Phys. Rev D* **73**, 077301 (2006).

1.- Neutrino Masses (1 loop \Rightarrow small)

2.- Dark Matter (DM)

	Standard Model			Fermions	Scalar
	L	e	ϕ	N	η
$SU(2)_L$	2	1	2	1	2
Y	-1	-2	1	0	1
\mathbb{Z}_2	+	+	+	-	-
l	1	1	0	1	0

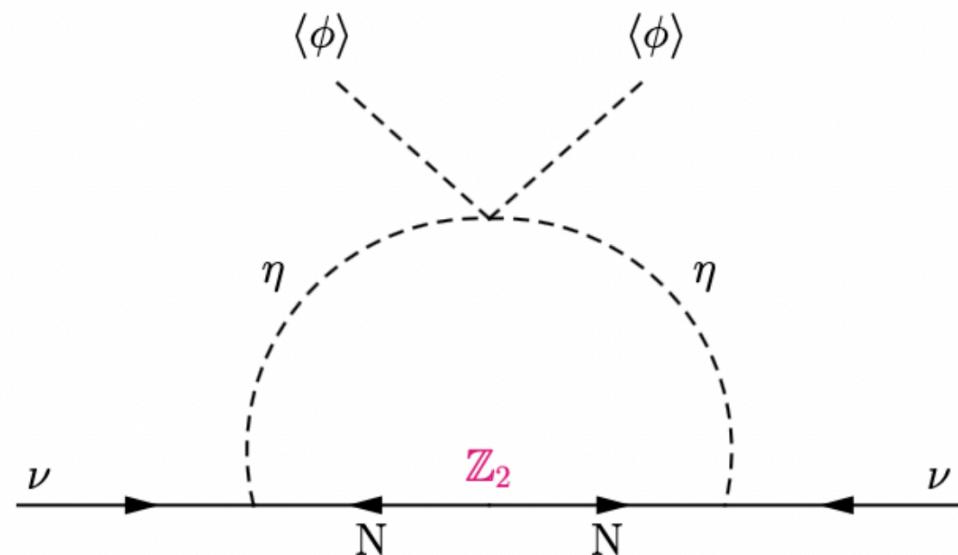
$$\mathcal{L} = - Y_N^{\alpha\beta} \bar{N}_\alpha \eta^\dagger L_\beta - \frac{1}{2} \bar{N}^\alpha M_{\alpha\beta} N^{\beta c} + \dots$$

$$m_\phi^2 = 2\lambda_1 v^2$$

$$m_{\eta^\pm}^2 = m_\eta^2 + \lambda_3 \frac{v^2}{2}$$

$$m_{\eta_R}^2 = m_\eta^2 + (\lambda_3 + \lambda_4 + \lambda_5) \frac{v^2}{2}$$

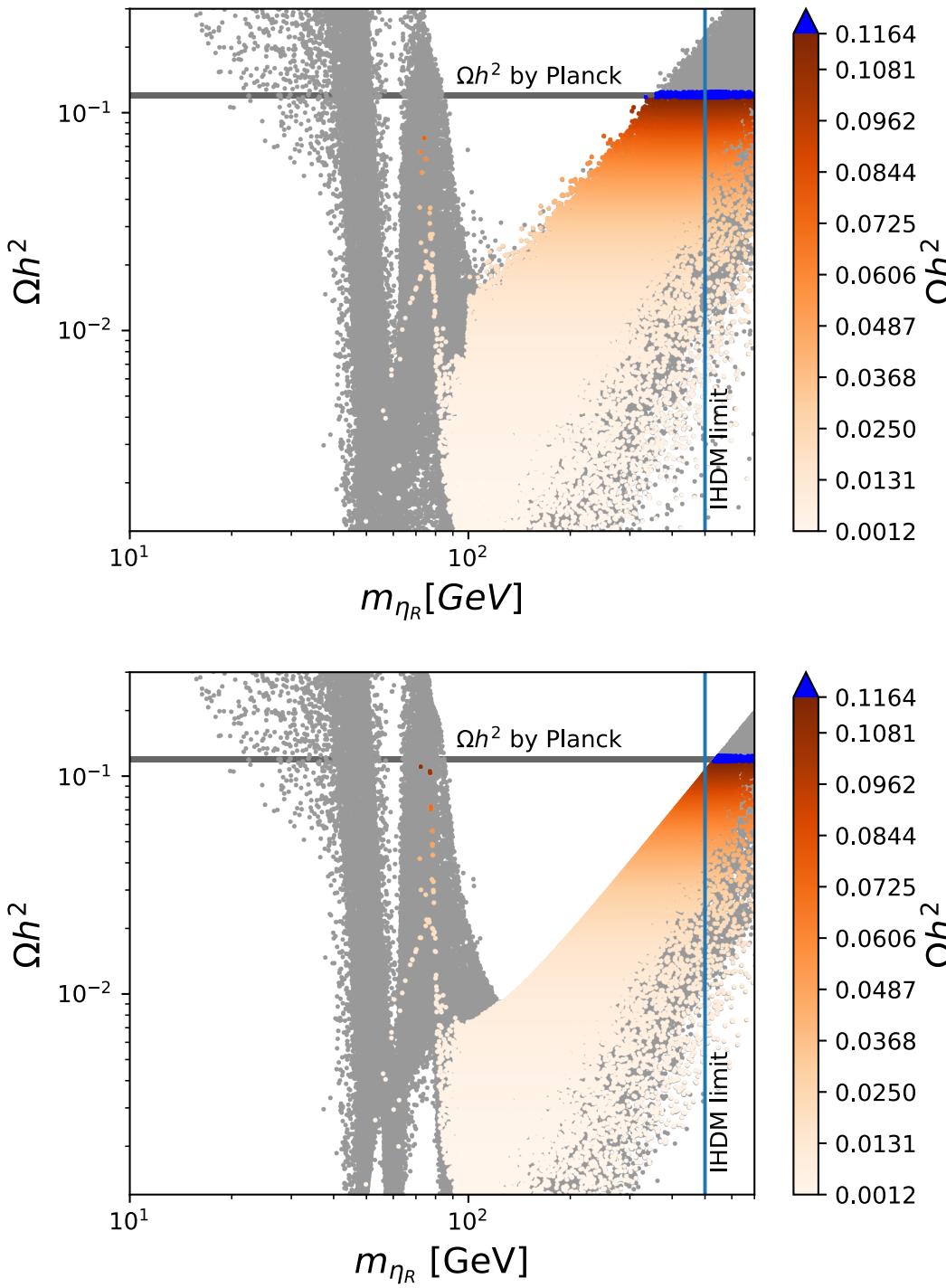
$$m_{\eta_L}^2 = m_\eta^2 + (\lambda_3 + \lambda_4 - \lambda_5) \frac{v^2}{2}$$



- DM: Real part of η (scalar DM)
- Look for long lived particles

$$\begin{aligned} V = & m_\phi^2 \phi^\dagger \phi + m_\eta^2 \eta^\dagger \eta + \lambda_1 (\phi^\dagger \phi)^2 + \lambda_2 (\eta^\dagger \eta)^2 \\ & + \lambda_3 (\phi^\dagger \phi)(\eta^\dagger \eta) + \lambda_4 (\phi^\dagger \eta)(\eta^\dagger \phi) \\ & + \frac{\lambda_5}{2} ((\phi^\dagger \eta) + (\eta^\dagger \phi))^2 \end{aligned}$$

Relic Density



Comparison of the relic density of DM in the Scotogenic model (top), with the Inert Higgs Doublet model (bottom).

The IHDM limit of about 500 GeV is reduced to about 300 GeV.

Parameter	Scanned range
λ_1	$[10^{-8}, 1]$
λ_2	$[10^{-8}, 1]$
λ_3	$\pm[10^{-8}, 1]$
λ_4	$\pm[10^{-8}, 1]$
λ_5	$\pm[10^{-8}, 1]$
m_η [GeV]	$[10, 1000]$
M_{N_1} [GeV]	$[50, 5000]$
M_{N_2} [GeV]	$[5 \times 10^3, 2 \times 10^6]$
M_{N_3} [GeV]	$[5 \times 10^3, 3.5 \times 10^6]$

Constraints.

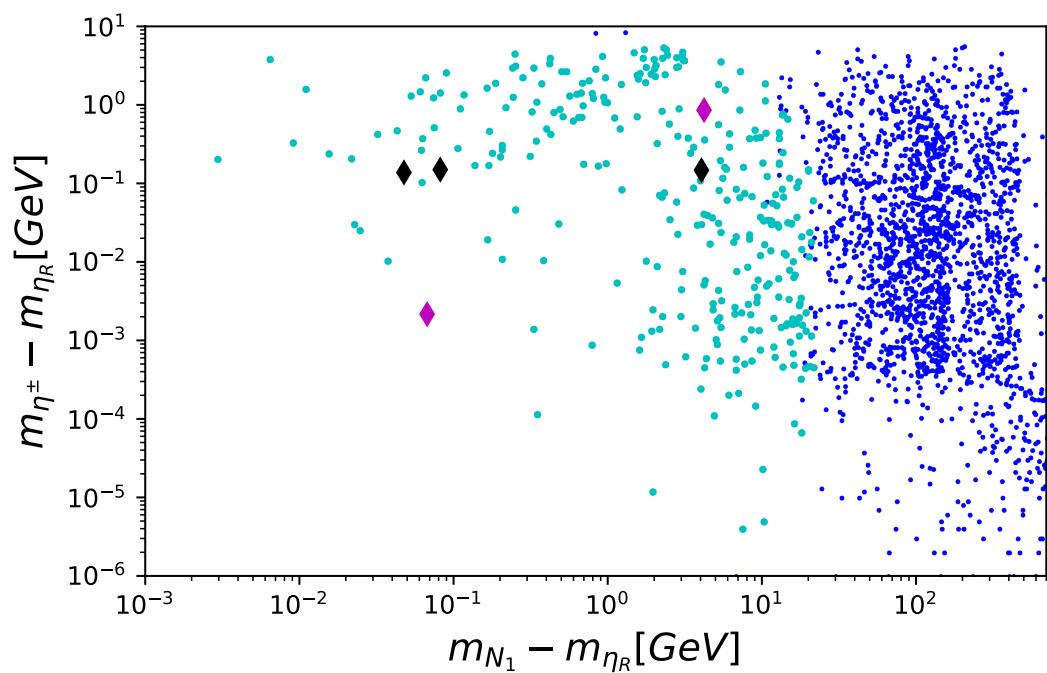
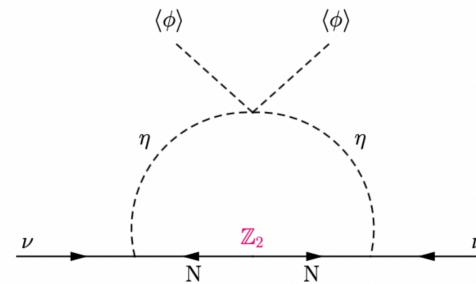
- Theoretical: η_R is the lightest (at tree level), potential bounded from below, perturbativity.
- Experimental: Neutrino masses and oscillations, electroweak precision, LFV, colliders, DM.

Neutrinos

Benchmark Points

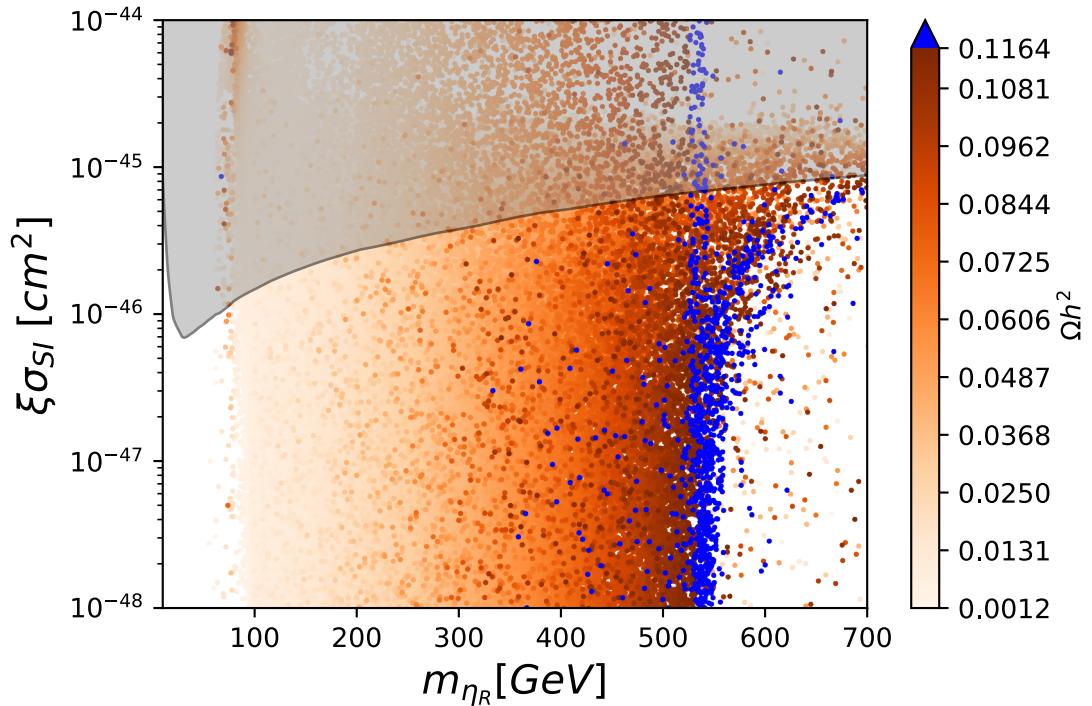
Parameter	B1	B2
λ_3	-2.809×10^{-4}	2.322×10^{-8}
λ_4	1.16×10^{-5}	-1.538×10^{-5}
λ_5	-2.511×10^{-2}	-2.878×10^{-5}
m_η^2 [GeV]	1.966×10^5	9.608×10^4
m_{η_R} [GeV]	442.535	309.961
m_{η_I} [GeV]	444.252	309.964
m_{η^\pm} [GeV]	443.394	309.964
m_{N_1} [GeV]	446.754	310.028
$c\tau_{N_1}$ [mm]	0.467	0.149
$\sigma(e^+e^- \rightarrow N_1 N_1)$ [fb]	9.89×10^{-20}	1.68×10^{-11}
Ωh^2	0.122	0.092

$$\mathcal{M}_{\alpha\beta}^\nu = \frac{m_{N_i}}{32\pi^2} Y_N^{\alpha i} Y_N^{\beta i} \left[\frac{m_{\eta_R}^2}{m_{\eta_R}^2 - m_{N_i}^2} \ln\left(\frac{m_{\eta_R}^2}{m_{N_i}^2}\right) - \frac{m_{\eta_I}^2}{m_{\eta_I}^2 - m_{N_i}^2} \ln\left(\frac{m_{\eta_I}^2}{m_{N_i}^2}\right) \right]$$



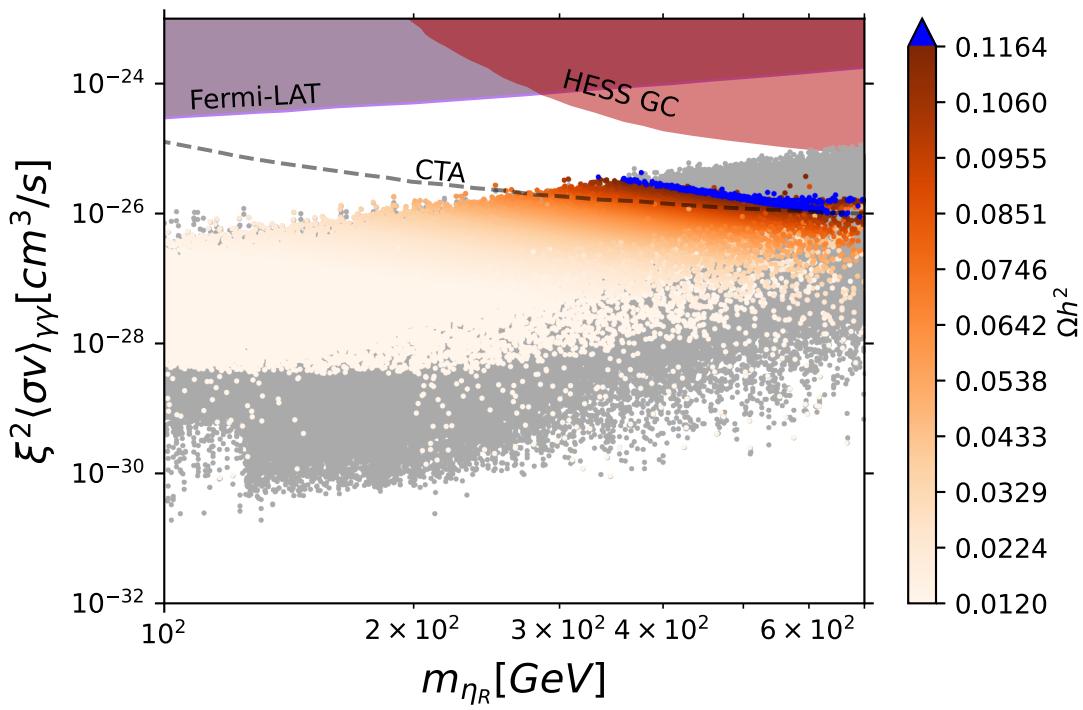
Parameter	B3	B4	B5
λ_3	-2.392×10^{-5}	3.305×10^{-6}	4.447×10^{-5}
λ_4	-6.923×10^{-7}	-1.46×10^{-3}	-3.293×10^{-6}
λ_5	-4.177×10^{-3}	-2.07×10^{-3}	-3.191×10^{-3}
m_η^2 [GeV]	1.851×10^5	1.276×10^5	1.234×10^5
m_{η_R} [GeV]	430.141	357.093	351.087
m_{η_I} [GeV]	430.435	357.269	351.362
m_{η^\pm} [GeV]	430.288	357.243	351.224
m_{N_1} [GeV]	434.197	357.175	351.134
$c\tau_{\eta^\pm}$ [mm]	16.859	14.587	28.412
$\sigma(pp \rightarrow \eta\eta j)$ [fb]	2.525	5.44	5.81
$N = \sigma \times BR \times \mathcal{L} \times \epsilon$	19.392	33.474	77.811
Ωh^2	0.121	0.121	0.119

Direct and Indirect Detection



-) Indirect searches.
-) DM annihilation cross section.
-) Grey points are excluded.
-) Blue points satisfy the total relic density.
-) Orange points satisfy it partially.

-) Direct searches.
-) η_R —nucleon spin independent elastic scattering cross section.
-) Grey: XENON1T.
-) Blue points: satisfy the totality of the relic density (300 to 700 GeV).
-) Orange points satisfy it partially.

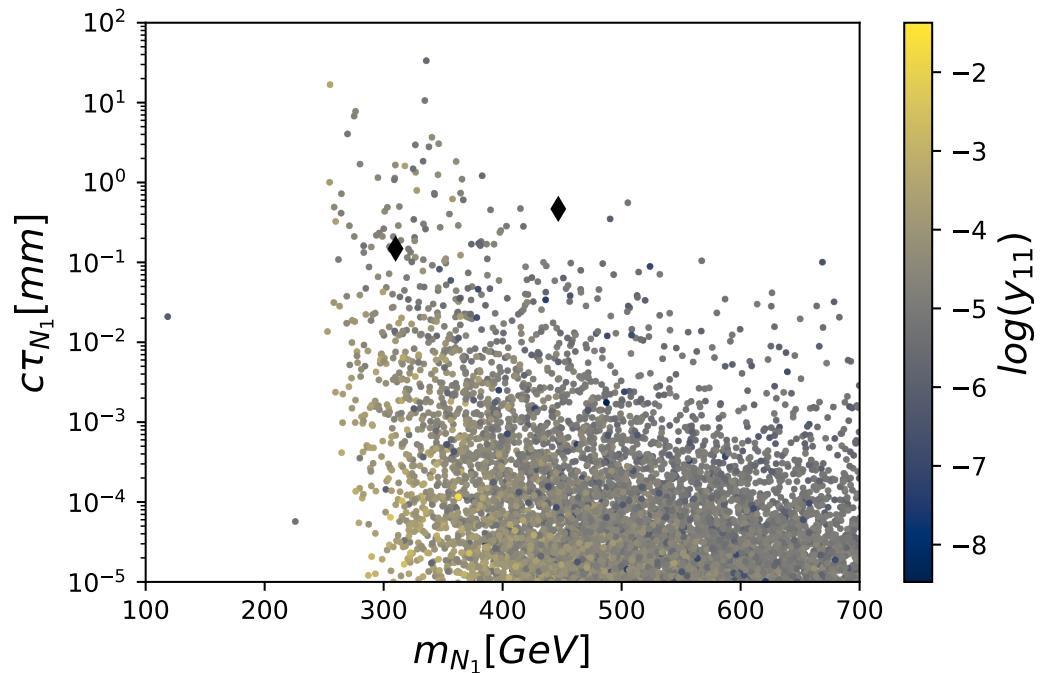
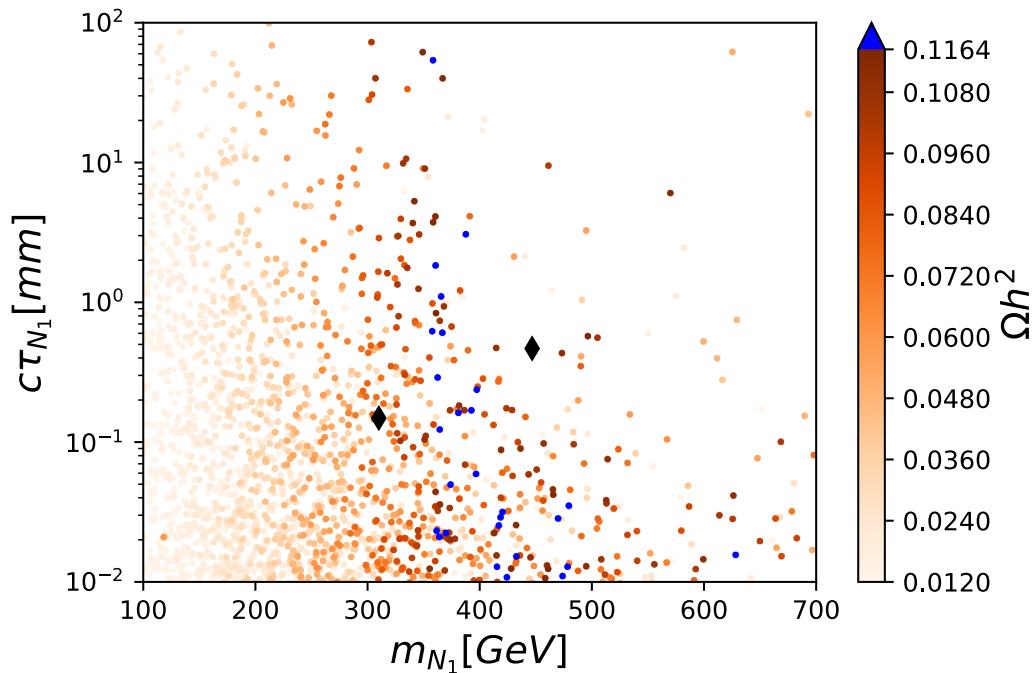


Long Lived N_1

$$\Gamma(N_k \rightarrow \ell_j^\pm \eta^\mp) = \frac{|Y_{jk}|^2}{32\pi} \frac{m_{N_k}^2 + m_{\ell_j}^2 - m_{\eta^\pm}^2}{m_{N_k}^3} \sqrt{(m_{N_k}^2 - m_{\ell_j}^2 - m_{\eta^\pm}^2)^2 - 4m_{\ell_j}^2 m_{\eta^\pm}^2}$$

$$\Gamma(N_k \rightarrow \nu \eta_\alpha) = \sum_j \frac{|Y_{jk}|^2}{32\pi} \frac{(m_{N_k}^2 - m_{\eta_\alpha}^2)^2}{m_{N_k}^3}$$

N_1 can be long-lived, but its production cross section is small.

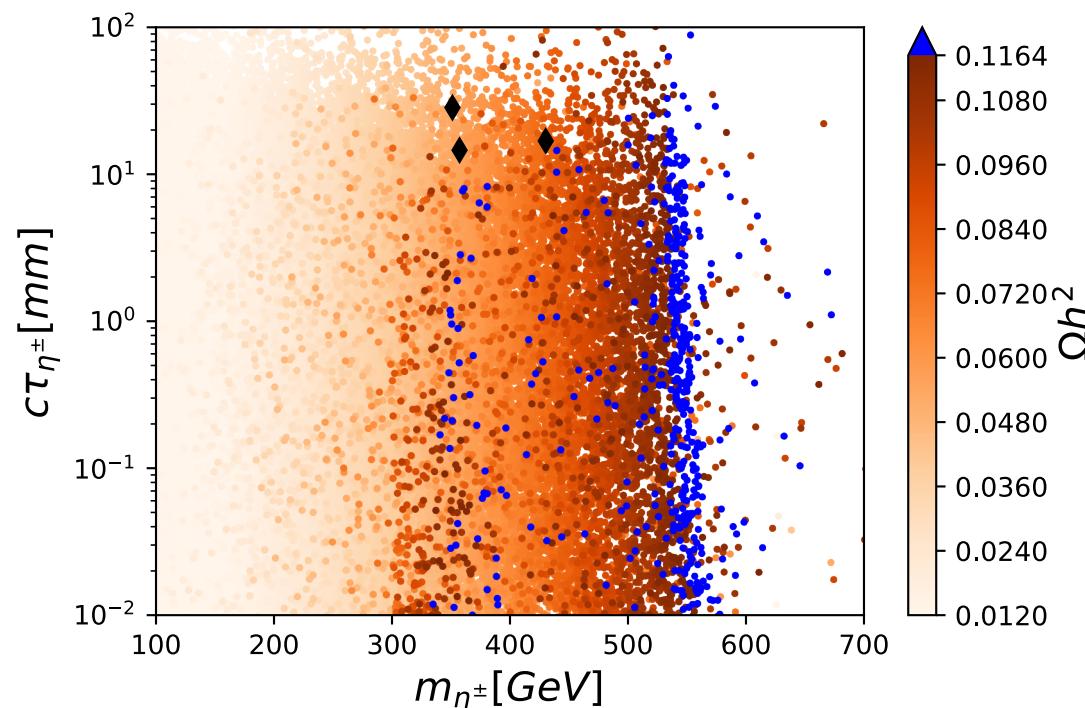


Long Lived η^\pm

$$\Gamma(\eta^\pm \rightarrow N_1 \ell^\pm) = \frac{Y_{\ell 1}^2 (m_{\eta^\pm}^2 - (m_{N_1} + m_\ell)^2)}{8\pi m_{\eta^\pm}} \sqrt{1 - \left(\frac{m_{N_1} - m_\ell}{m_{\eta^\pm}}\right)^2} \sqrt{1 - \left(\frac{m_{N_1} + m_\ell}{m_{\eta^\pm}}\right)^2}$$

$$\Gamma(\eta^\pm \rightarrow \eta_R \pi^\pm) = \frac{f_\pi^2 g^4 (m_{\eta^\pm}^2 - m_{\eta_R}^2)}{512\pi m_W^4 m_{\eta^\pm}} \sqrt{1 - \left(\frac{m_{\eta_R} - m_\pi}{m_{\eta^\pm}}\right)^2} \sqrt{1 - \left(\frac{m_{\eta_R} + m_\pi}{m_{\eta^\pm}}\right)^2}$$

η^\pm can be long-lived, and its production cross section is large.



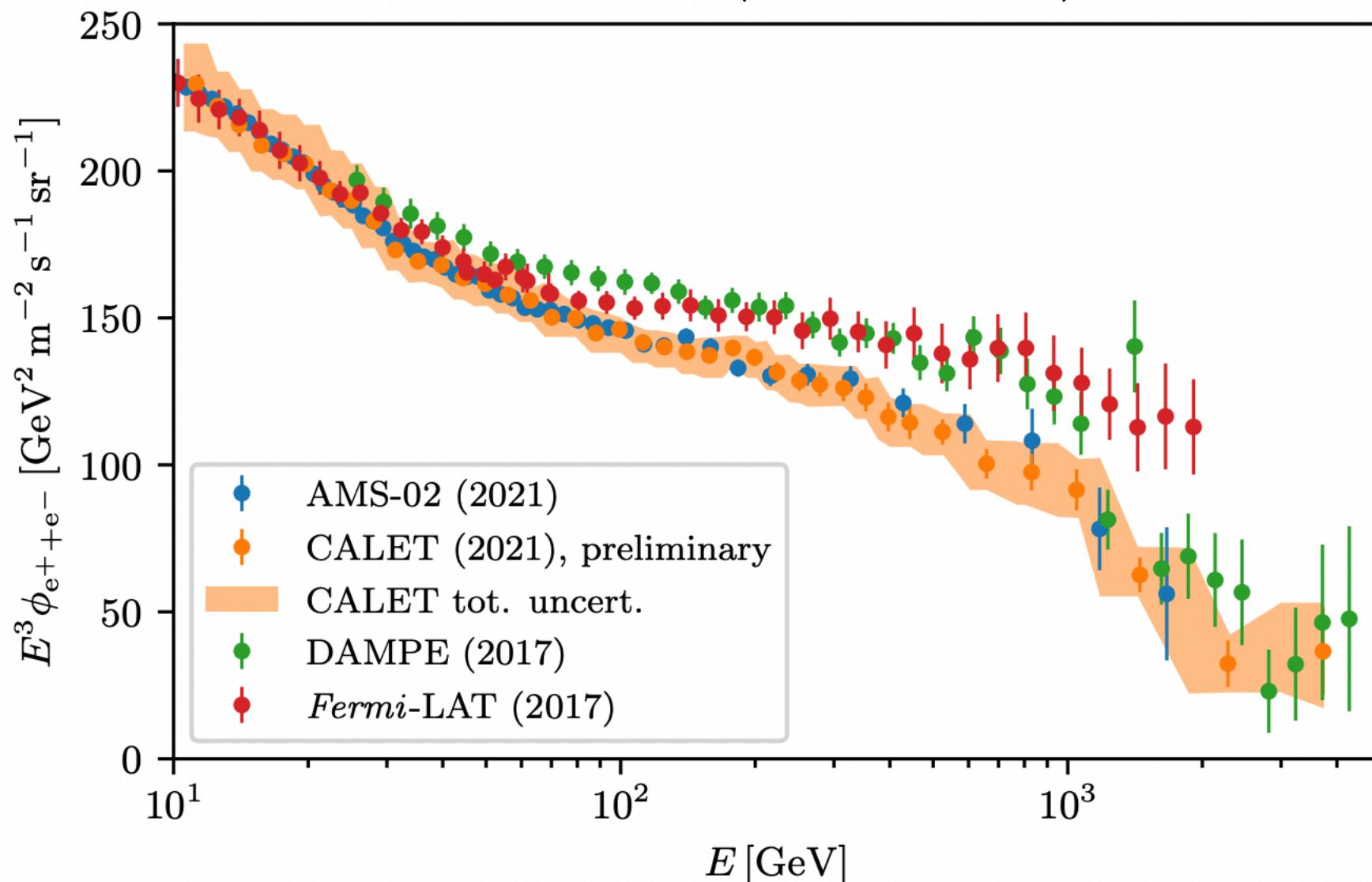
Particles from Space

All-electron Flux

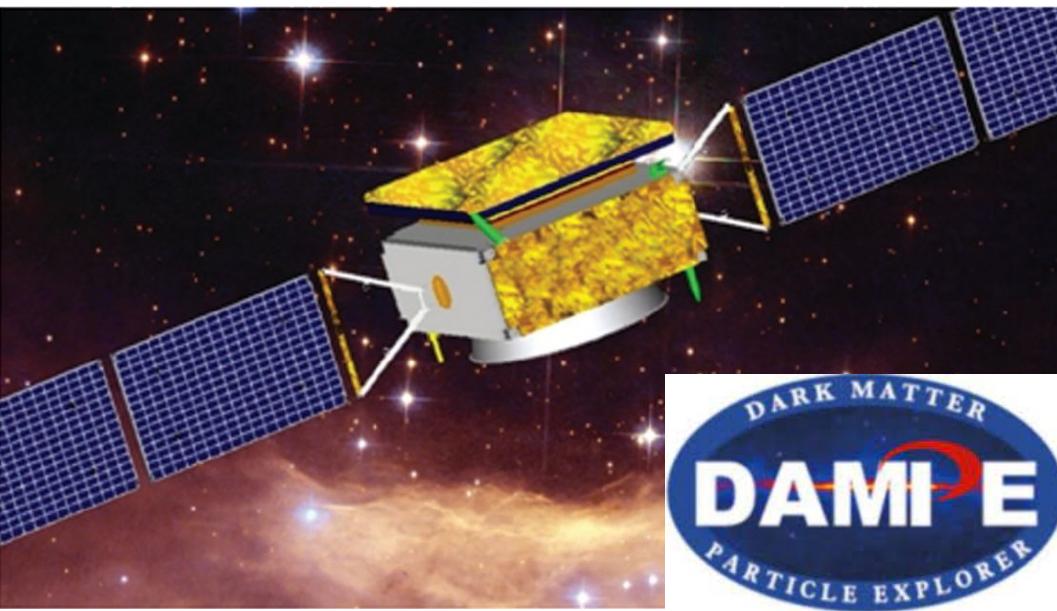
Between 10 and 30 GeV: $E^{-3.25}$

Between 30 GeV and cutoff: $E^{-3.1}$ (DAMPE, Fermi-LAT)

$E^{-3.15}$ (AMS-02, CALET)

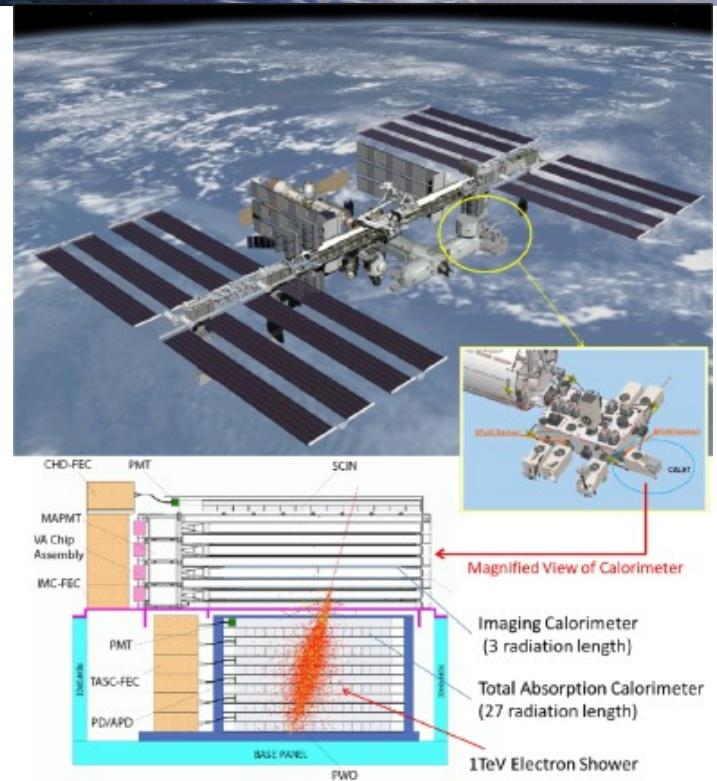
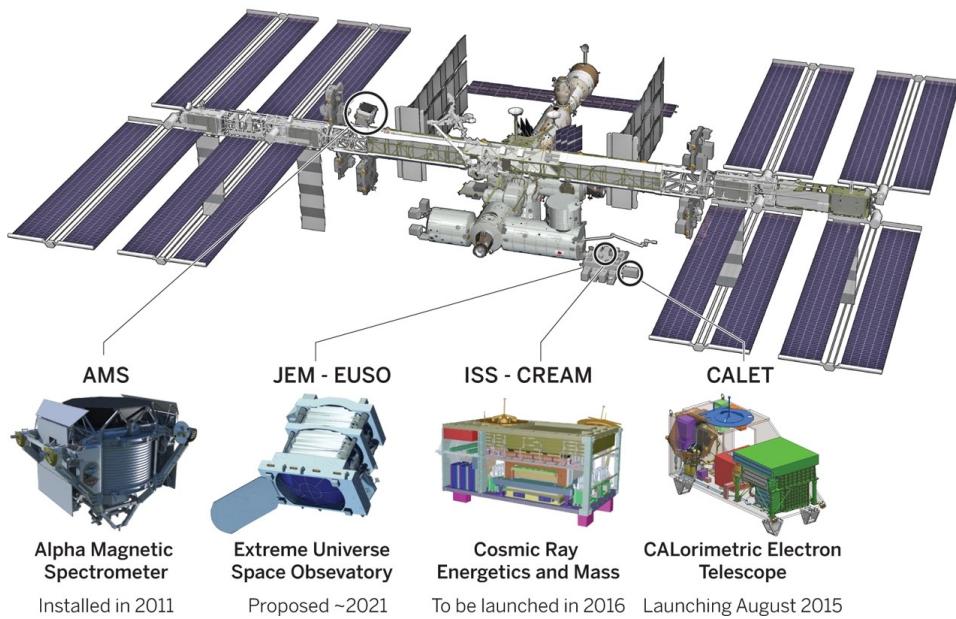


Satellites

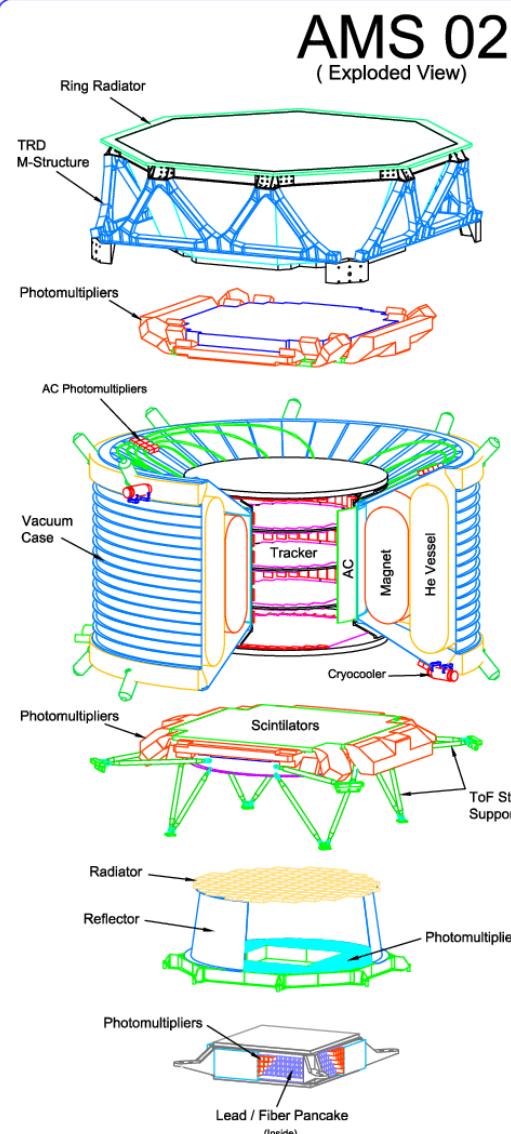


Cosmic ray detectors on the ISS

New experiments, perched outside Earth's atmosphere, promise to turn the International Space Station into a well-rounded platform for unlocking the secrets of supernovae and even dark matter.



Alpha Magnetic Spectrometer (AMS-02)



TRD:
Transition
Radiation
Detector

ToF: (s1,s2)
Time of Flight
Detector

TR:
Silicon Tracker

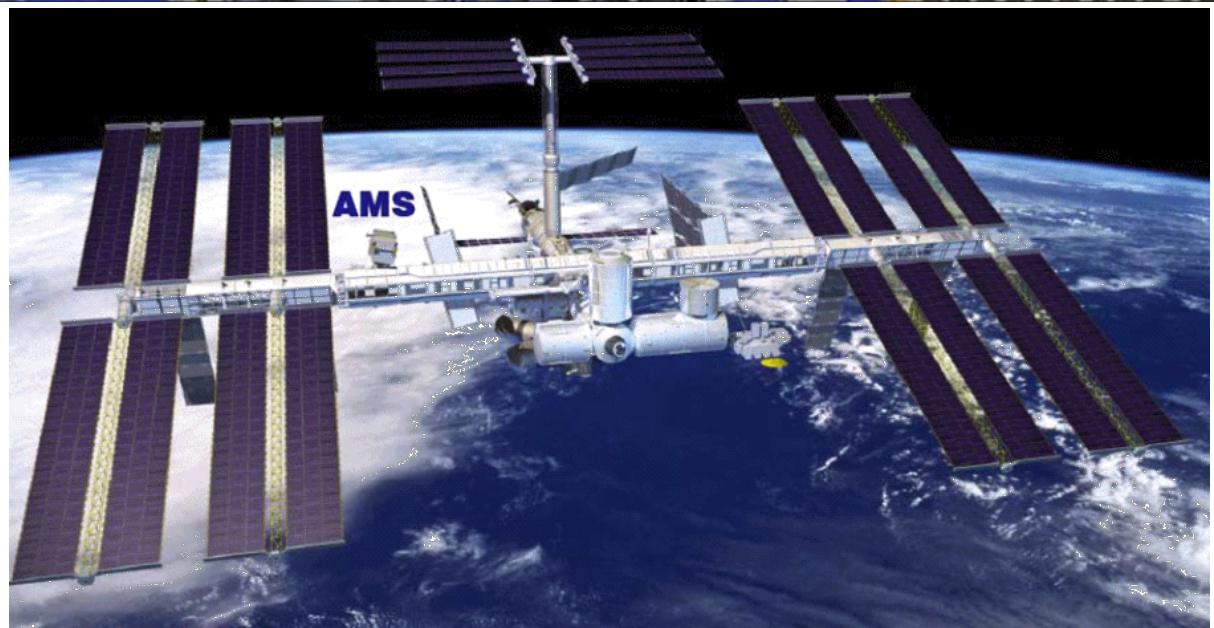
AC:
Anticoincidence
Counter

MG:
Magnet

ToF: (s3,s4)
Time of Flight
Detector

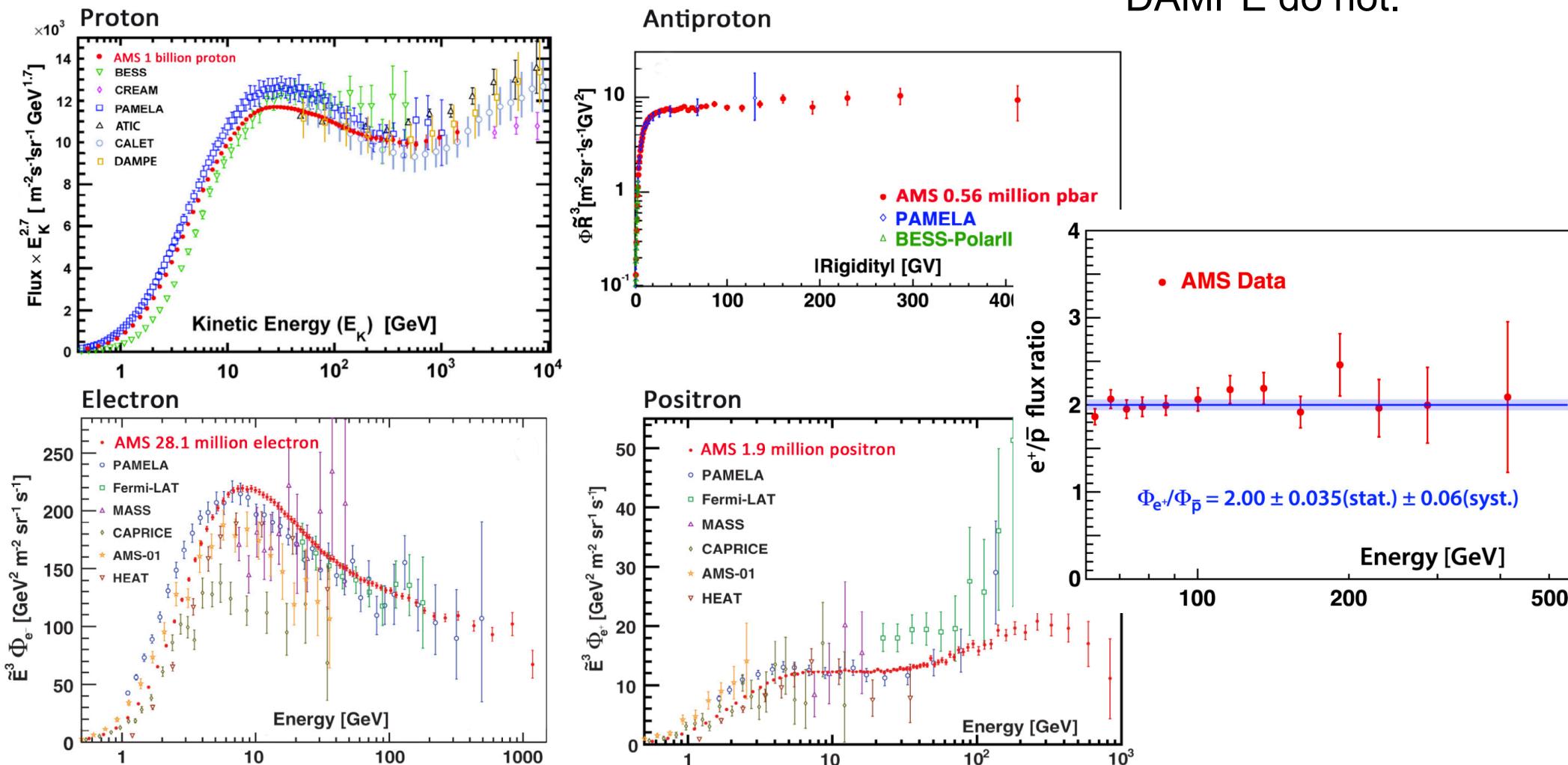
RICH:
Ring image
Cherenkov Counter

EMC:
Electromagnetic
Calorimeter



Latest results from AMS-02

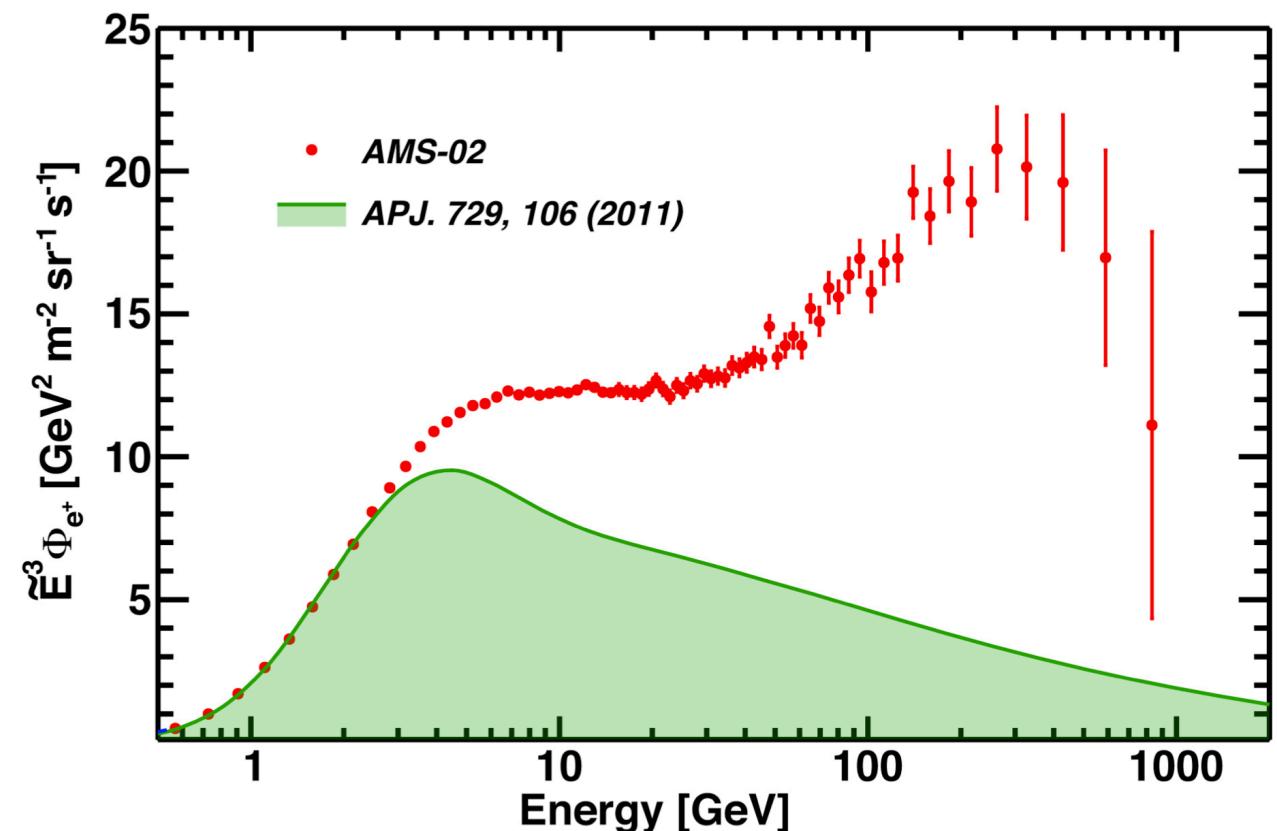
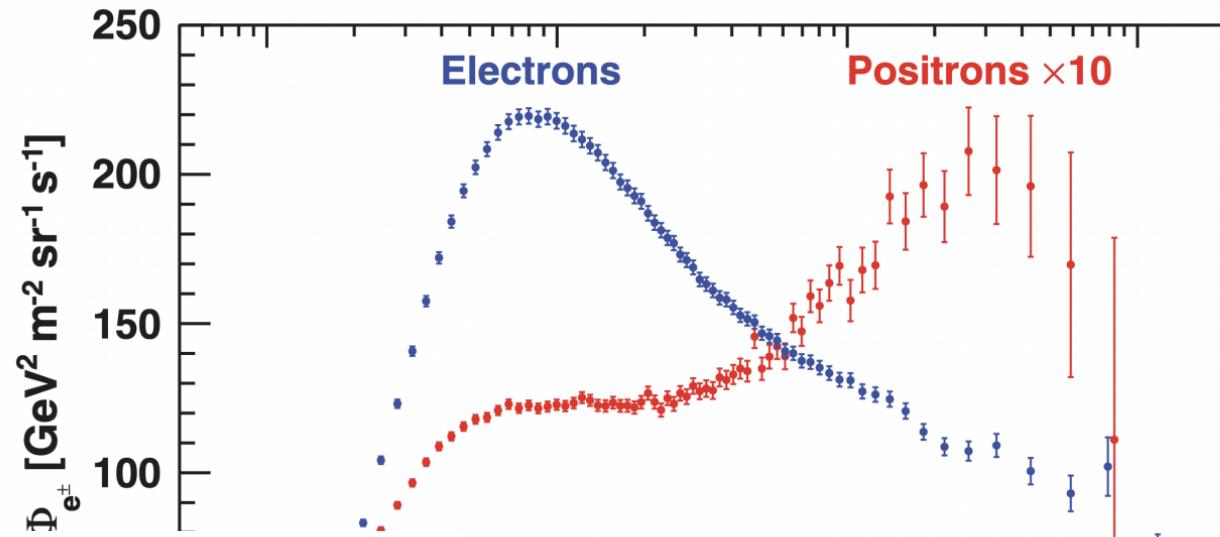
AMS-02 can differentiate between electrons and positrons, but CALET and DAMPE do not.



Significant excess starting from 25.2 ± 1.8 GeV, with a sharp drop above 284^{+91}_{-64} GeV [Z. Weng, PoS 045, ICHEP2020 (2020)]. Antiproton flux has a similar behavior.

Latest Results on Electrons and Positrons from AMS

Rise on positron flux persists with time [Z. Weng, PoS 045, ICHEP2020 (2020)].



Positron flux measured by AMS-02 [J. Bedugo, PoS 016, ICRC2021 (2021)] compared to GALPROP prediction for cosmic rays.

Supersymmetry

1.- R-Parity (Rp) Conserving:

a.- Dark Matter (DM) candidate: neutralino

b.- Superpotential:

$$W_{MSSM} = (\mu \hat{H}_u^a \hat{H}_d^b + h_{ij}^e \hat{H}_d^a \hat{H}_u^b + h_{ij}^d \hat{H}_d^a \hat{Q}_i^b \hat{D}_j - h_{ij}^u \hat{H}_u^a \hat{Q}_i^b \hat{U}_j) \varepsilon_{ab}$$

2.- R-Parity Violating:

a.- Dark Matter candidate: gravitino

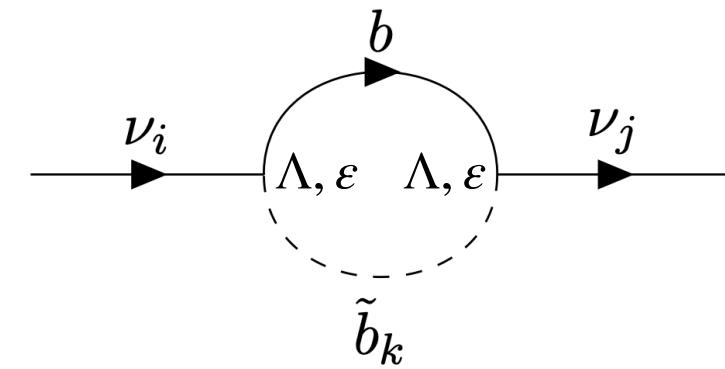
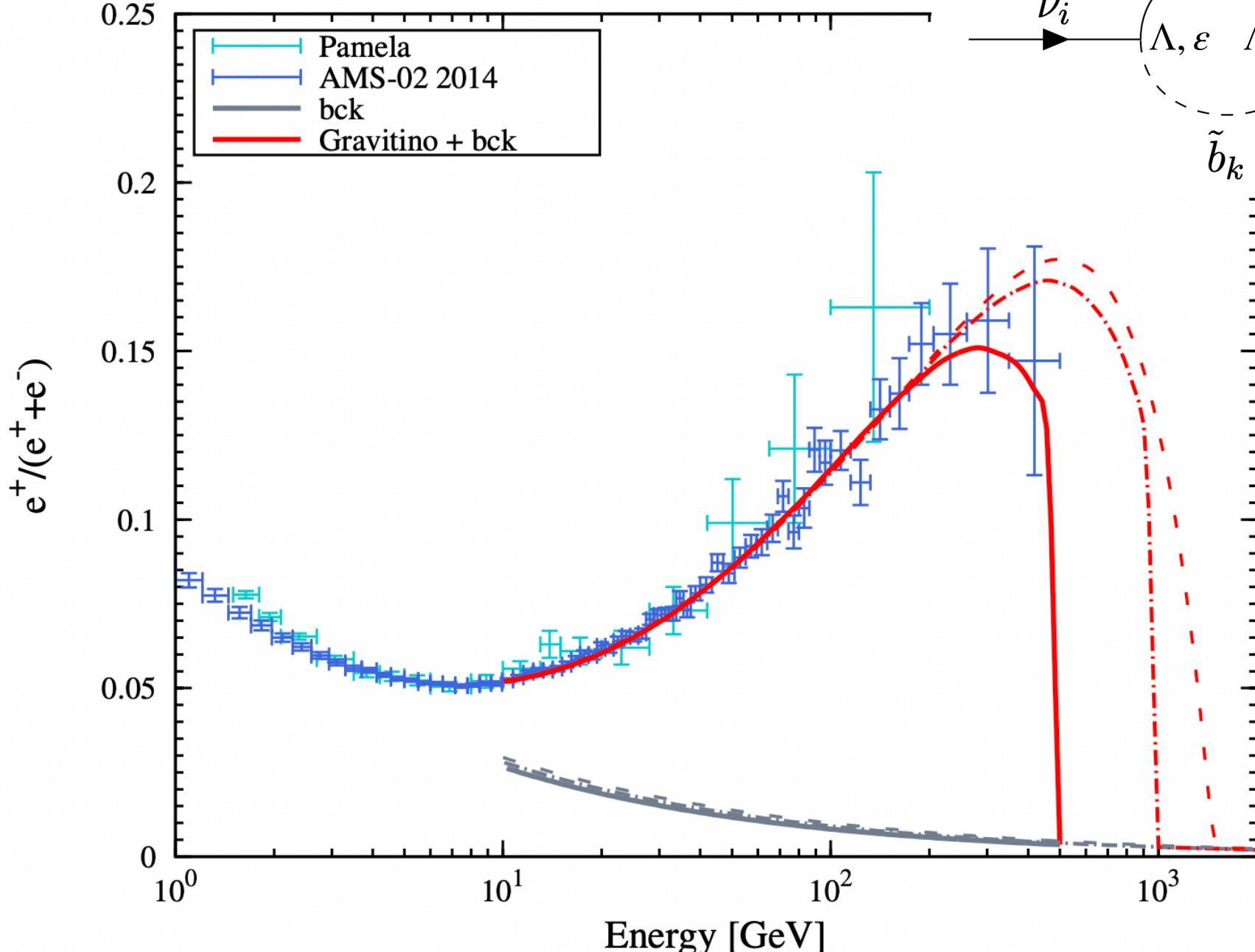
b.- Superpotential:

$$\begin{aligned} W_{MSSM+RpV} = & W_{MSSM} + (\epsilon_i \hat{H}_u^a \hat{L}_i^b \\ & + \frac{1}{2} \lambda_{ijk} \hat{L}_i^a \hat{L}_j^b \hat{R}_k + \lambda'_{ijk} \hat{L}_i^a \hat{Q}_j^b \hat{D}_k) \varepsilon_{ab} + \frac{1}{2} \lambda''_{ijk} \hat{U}_i \hat{D}_j \hat{D}_k \end{aligned}$$

3.- The gravitino interaction is governed by:

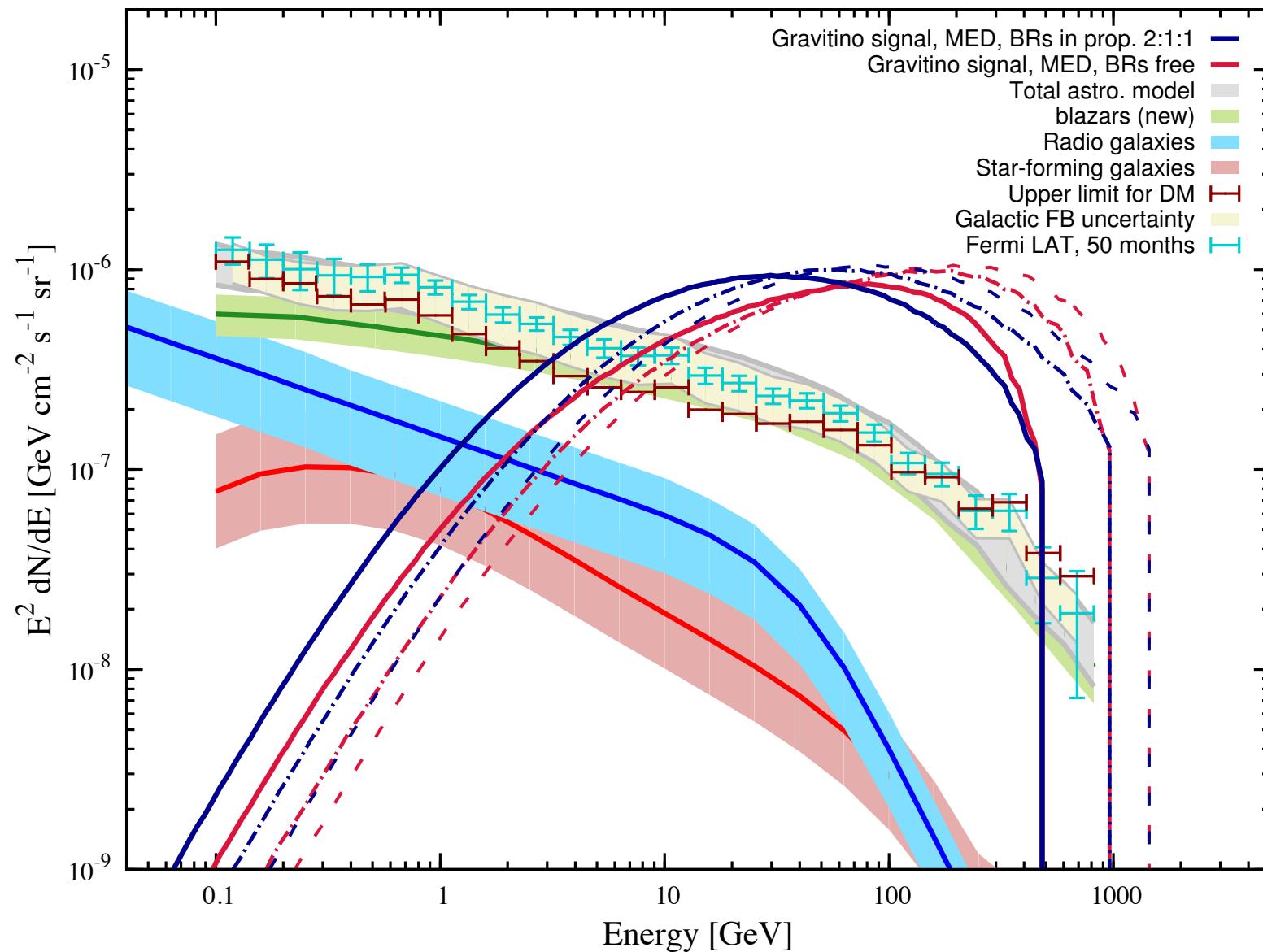
$$\mathcal{L} = -\frac{1}{\sqrt{2} M_P} \bar{F} \gamma^\mu \gamma^\nu \partial_\nu \tilde{F} \tilde{G}_\mu$$

BRpV Only

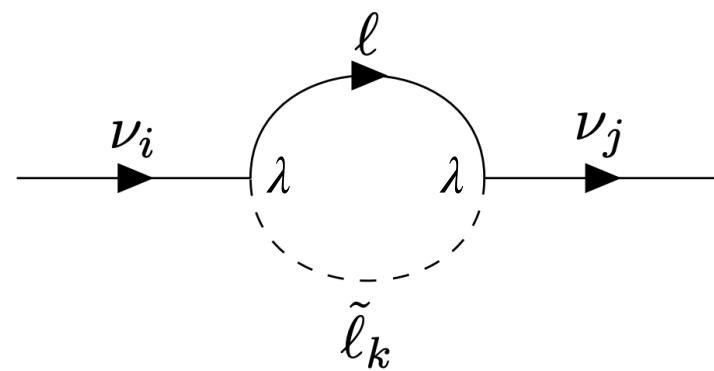


Based on *Phys. Dark Univ.* **11**, 1 (2016), in collaboration with E. Carquin, G. Gómez-Vargas, B. Panes, and N. Viaux.

Explains the rise and drop of the positrons flux, but emits too many photons. That is, the galaxy would be too bright.



TRpV Only



$$\Phi_{e^-} = C_e \left(\frac{E}{1 \text{ GeV}} \right)^{-\gamma_e}$$

$$\Phi_{e^+} = C_p \left(\frac{E}{1 \text{ GeV}} \right)^{-\gamma_p}$$

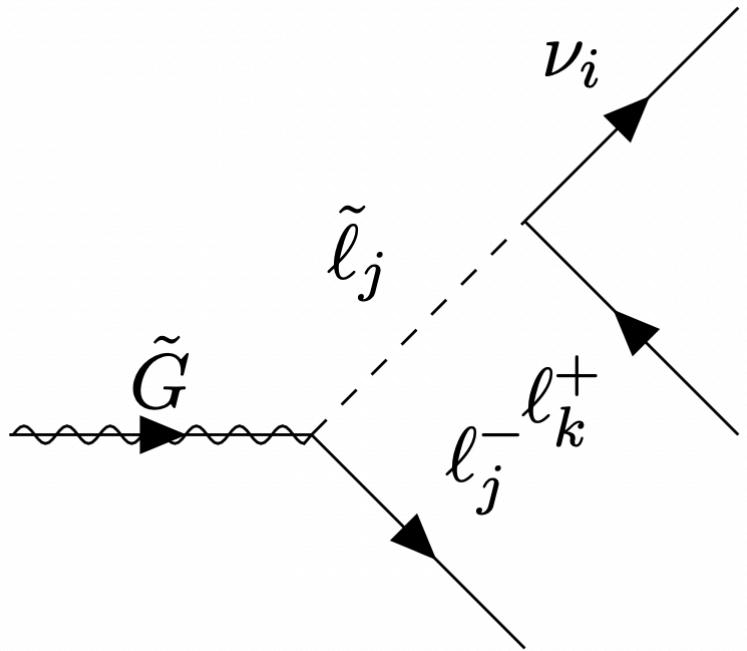
$$\alpha_1 = B(\tilde{G} \rightarrow e^- e^+ \nu) + B(\tilde{G} \rightarrow e^- \mu^+ \nu) + B(\tilde{G} \rightarrow e^- \tau^+ \nu)$$

$$\alpha_2 = B(\tilde{G} \rightarrow \mu^- e^+ \nu) + B(\tilde{G} \rightarrow \mu^- \mu^+ \nu) + B(\tilde{G} \rightarrow \mu^- \tau^+ \nu)$$

$$\alpha_3 = B(\tilde{G} \rightarrow \tau^- e^+ \nu) + B(\tilde{G} \rightarrow \tau^- \mu^+ \nu) + B(\tilde{G} \rightarrow \tau^- \tau^+ \nu)$$

Parameter	Case 1	Case 2	Case 3	Case 4
C_p [1/GeV cm ² s str]	14.90	14.74	14.93	14.37
γ_p	3.11	3.10	3.11	3.09
C_e [1/GeV cm ² s str]	426.10	421.77	422.08	422.67
γ_e	3.27	3.27	3.27	3.27
m_G [GeV]	1281	2274	3604	3751
τ_G [10 ²⁶ s]	4.61	3.59	2.27	2.29
$\alpha_1 : e^- l^+ \nu$	0.43	0.06	0.03	0.32
$\alpha_2 : \mu^- l^+ \nu$	0.03	0.36	0.15	0
$\alpha_3 : \tau^- l^+ \nu$	0.54	0.58	0.82	0.68

Gravitino decay



$$\tilde{m} \equiv m_{\tilde{\nu}_i} = m_{\tilde{\ell}_{i,1}} = m_{\tilde{\ell}_{i,2}}$$

$$\Gamma(\tilde{G} \rightarrow \nu_i e_j \bar{e}_k) \approx \frac{\lambda_{ijk}^2}{3\pi^3 2^{11}} \frac{m_{\tilde{G}}^7}{M_P^2 \tilde{m}^4}$$

$$\tau_{\tilde{G}} = 4 \times 10^{26} \text{ s} \left(\frac{1}{\lambda_{ijk} \lambda_{ijk}} \right) \left(\frac{\tilde{m}}{10^8 \text{ GeV}} \right)^4 \left(\frac{2 \text{ TeV}}{m_{\tilde{G}}} \right)^7$$

Neutrinos

$$M_{ij}^{\nu(1)} \approx \frac{1}{16\pi^2} \sum_{gr} s_{\tilde{\ell}} c_{\tilde{\ell}} (\lambda_{igr} \lambda_{jrg} + \lambda_{jgr} \lambda_{irg}) m_g \ln \frac{m_{\tilde{\ell}_{r,2}}^2}{m_{\tilde{\ell}_{r,1}}^2}$$

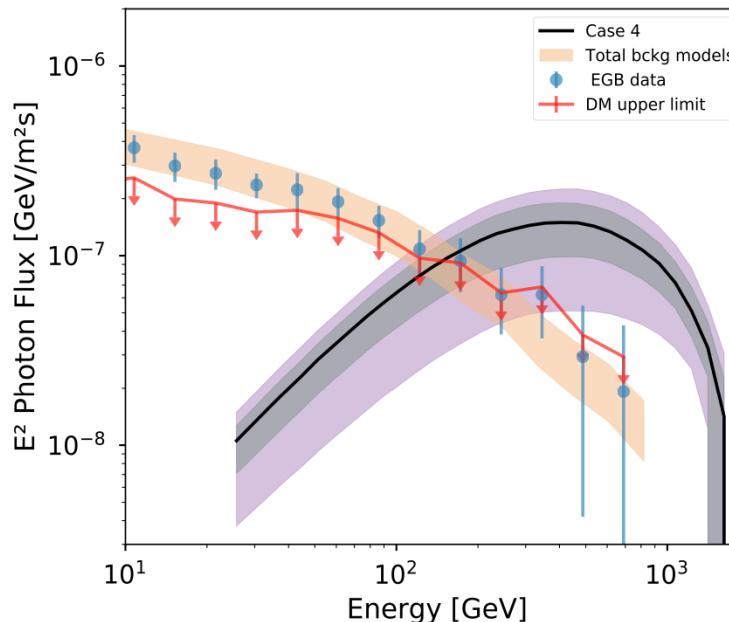
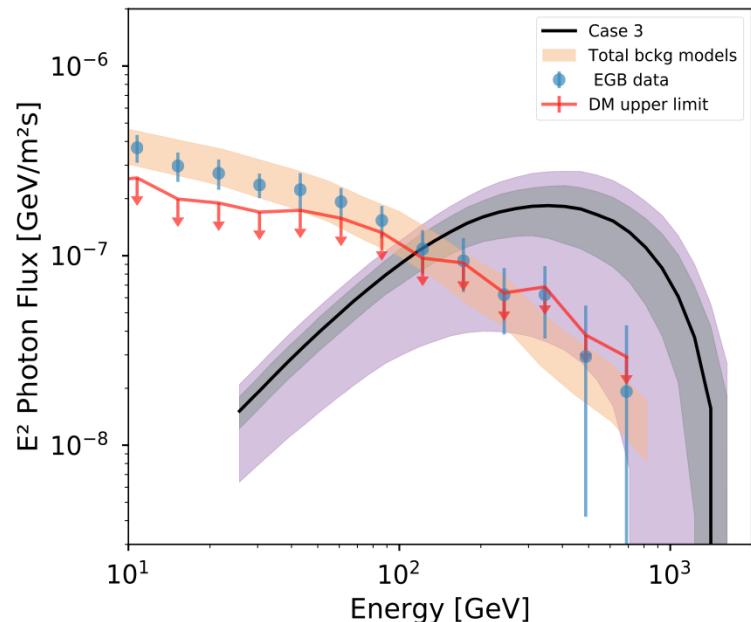
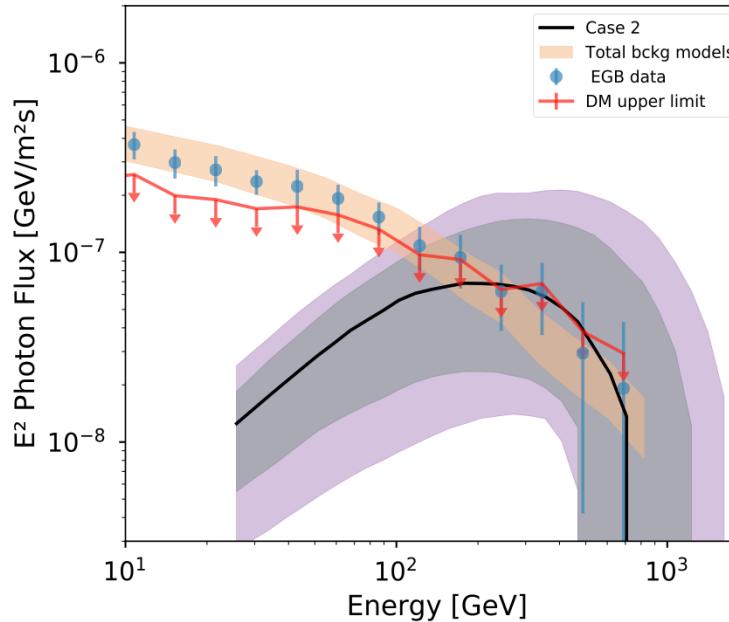
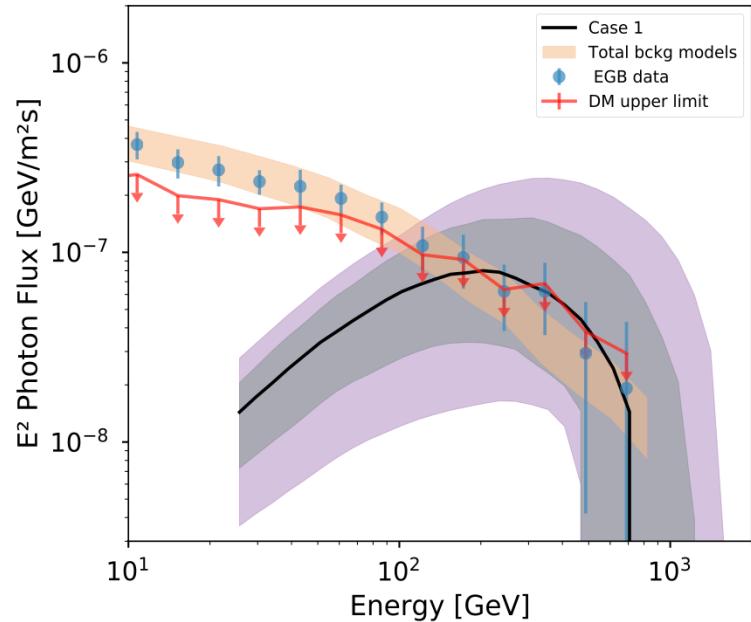
$$M_{ij}^{\nu(1)} \approx \frac{1}{8\pi^2} \lambda_{i23} \lambda_{j32} \frac{m_\mu m_\tau A_\tau}{\tilde{m}^2}$$

$$\approx 2 \times 10^{-2} \text{ eV} \lambda_{i23} \lambda_{j32} \left(\frac{10^8 \text{ GeV}}{\tilde{m}} \right)$$

$$\approx 2 \times 10^{-2} \text{ eV} (\lambda_{i23} \lambda_{j32})^{3/4} \left(\frac{4 \times 10^{26} \text{ s}}{\tau_{\tilde{G}}} \right)^{1/4} \left(\frac{2 \text{ TeV}}{m_{\tilde{G}}} \right)^{7/4}$$

Photon Flux

Pure TRpV can explain AMS-02



Another effect:
DAMPE and CALET
can measure all
electron flux to higher
energies.

There is tension be-
tween this explanation
and the photon flux
when we add CALET
and/or DAMPE.

Conclusions

- 1.- In the original Scotogenic model masses of dark matter can be smaller than 500 GeV, as noticed before (in contrast in the IHDM it cannot).
- 2.- In this model, charged and neutral new particles can be long lived.
- 3.- The positron fraction in cosmic rays rises with energy. A decaying Gravitino with mass of a few TeV can explain it. In this case, pure BRpV produce in too many photons.
- 4.- Pure TRpV produce less photons than pure BRpV, and can explain the rise of positron fraction measured by AMS-02. But there is tension between the measurement of different satellites. In addition, neutrino masses and mixing angles should be evaluated.
- 5.- Cosmic rays provide complementary information, thus satellites are an important tool. Particles can reach higher energies, although we loose control.
- 6.- To establish a model beyond the SM we should use a combination of results from different experiments: neutrino, colliders, and satellites.
- 7.- Any possible discrepancy between measurements coming from different satellites might be important and should be resolved.