

# Toward Cosmological Concordance with New Physics in the Dark Sector

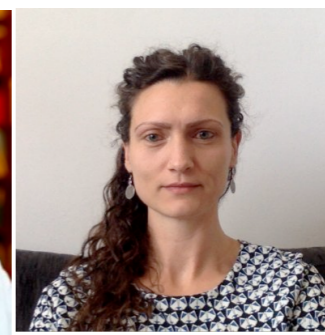
Colin Hill

Columbia University  
Flatiron Institute - Center for Computational Astrophysics

“Tensions in Cosmology” @ Corfu  
8 September 2022



2003.07355 w/ E. McDonough, M. Toomey, S. Alexander  
2006.11235 w/ EM,MT,SA + M. Ivanov, M. Simonovic, M. Zaldarriaga  
2109.04451 w/ ACT Collaboration  
2112.09128 w/ EM, W. Hu, M.-X. Lin, S. Zhou + work in prep.  
2112.10754 w/ A. La Posta, T. Louis, X. Garrido  
to appear w/ F. McCarthy + to appear with B. Bolliet



# Outline

- Early Dark Energy: ACT DR4 (+SPT-3G)
- —> Early Dark Sector
- Post-Recombination Reheating
- ~~Generalized Dark Matter~~ —> ~~Dark Radiation~~  
~~Conversion~~



# Early Dark Energy

Motivation: increase CMB-inferred  $H_0$

How does this work?

By decreasing the physical size of the sound horizon imprinted in the CMB

$$r_s^* = \int_0^{t_*} \frac{dt}{a(t)} c_s(t) = \int_{z_*}^{\infty} \frac{dz}{H(z)} c_s(z)$$

scale factor      sound speed

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Relevant ingredients in  $\Lambda$ CDM:  $\omega_b$ ,  $\omega_{\text{cdm}}$ ,  $\omega_v$ ,  $\omega_\gamma$

physical densities of  
baryons, CDM,  
neutrinos, photons

Angular sound horizon is (approx.) related to peak spacing:

$$\text{measured} \rightarrow \theta_s^* = \pi / \Delta \ell \longrightarrow D_A^* = r_s^* / \theta_s^* \longrightarrow H_0$$

$$D_A \sim 1/H_0$$

# Early Dark Energy

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$$r_s^* = \int_0^{t_*} \frac{dt}{a(t)} c_s(t) = \int_{z_*}^{\infty} \frac{dz}{H(z)} c_s(z)$$

Relevant ingredients in **EDE**:  $\omega_b$ ,  $\omega_m$ ,  $\omega_v$ ,  $\omega_\gamma$  + **EDE parameters**

Angular sound horizon is (approx.) related to peak spacing:

$$\theta_s^* = \pi / \Delta\ell \longrightarrow D_A^* = r_s^* / \theta_s^* \longrightarrow H_0$$

# Early Dark Energy

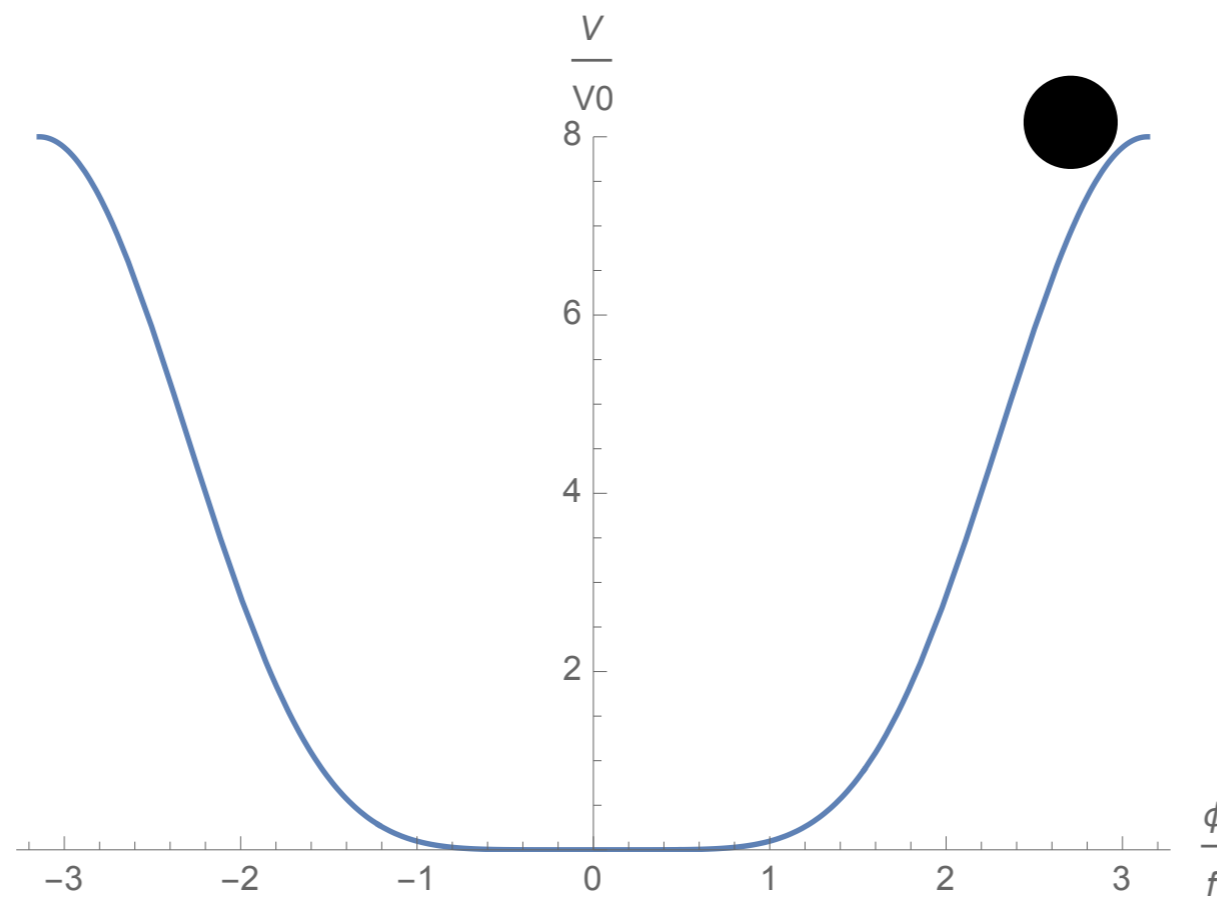
New component: (pseudo)-scalar field  $\phi$



# Early Dark Energy

New component: (pseudo)-scalar field  $\phi$

Idea: field initially frozen on its potential due to Hubble friction — acts as dark energy (equation of state  $w=-1$ )



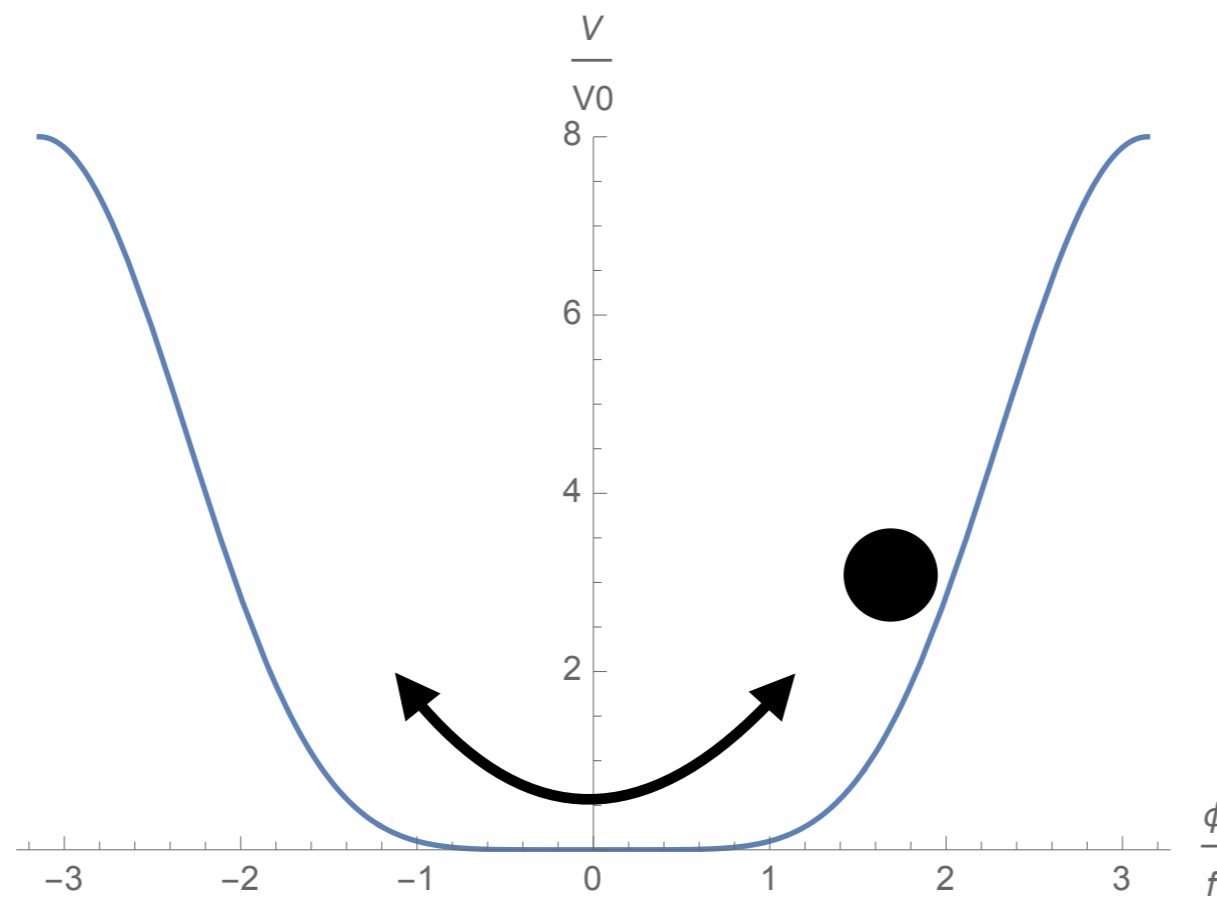
$$\ddot{\phi} + 3H\dot{\phi} + V'(\phi) = 0$$

$H \gg m$   
initially

# Early Dark Energy

New component: (pseudo)-scalar field  $\phi$

When  $H \sim m$  (field mass), it rolls down its potential and oscillates: effective  $w$  will depend on potential



For EDE, this must occur near  $\sim Z_{\text{CMB}}$



$$m \sim 10^{-27} \text{ eV}$$

e.g.,  $\phi(t) = \phi_i a^{-3/2} \cos(mt)$  if  $V(\phi) = m^2 \phi^2 / 2$

# Early Dark Energy

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Important: need late-time  $w>0$  so that EDE energy density contribution decays faster than matter

# Early Dark Energy

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Canonical EDE  
Potential:

$$V(\phi) = m^2 f^2 (1 - \cos(\phi/f))^n$$

Near minimum,  $V \sim \phi^{2n} \longrightarrow w_\phi = \frac{n-1}{n+1}$

$$m \sim 10^{-27} \text{ eV}$$

$$f \sim 10^{26-27} \text{ eV}$$

$n \geq 2$  (we fix to 3 throughout)

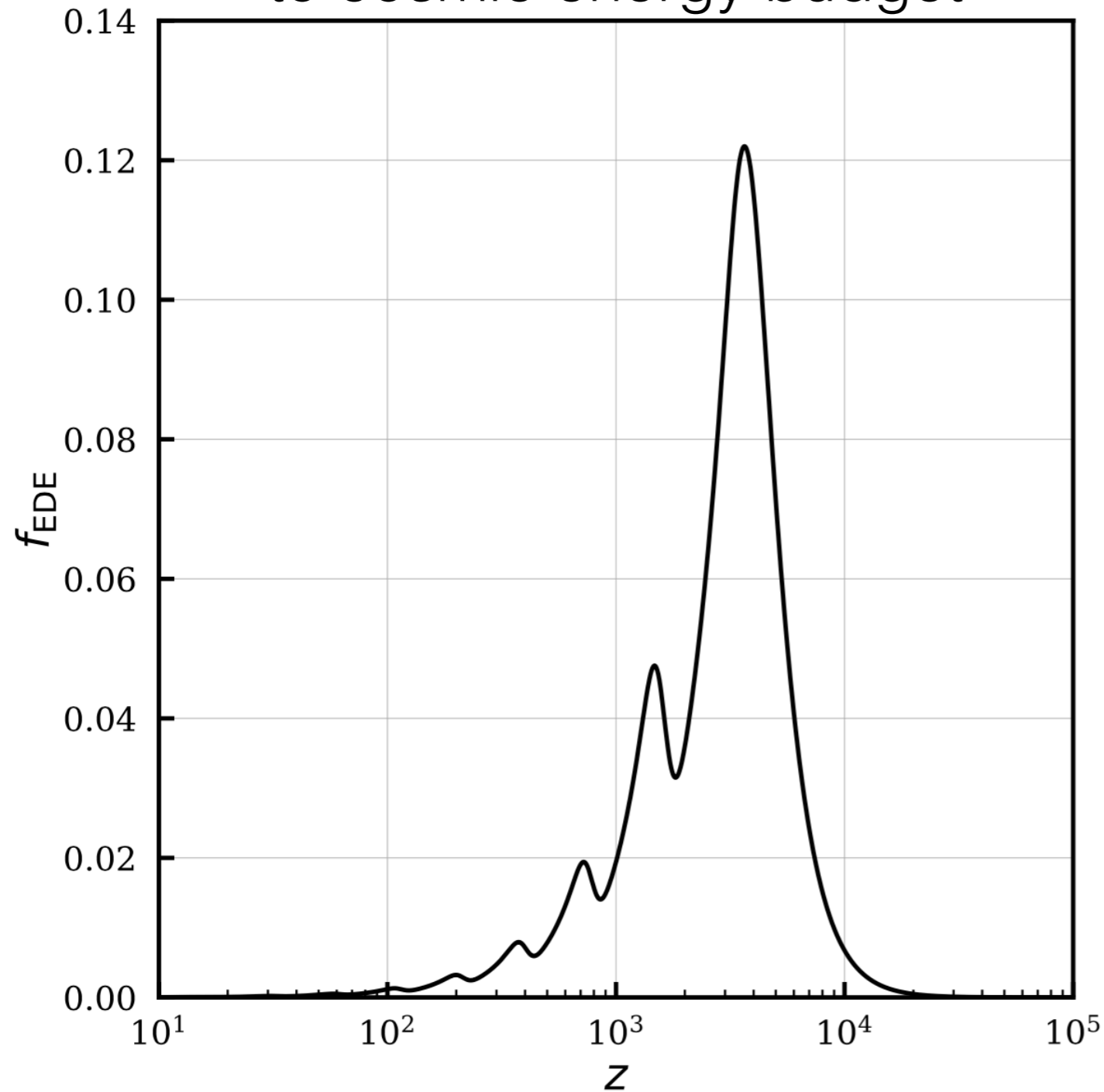


# Early Dark Energy

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## Parameterization

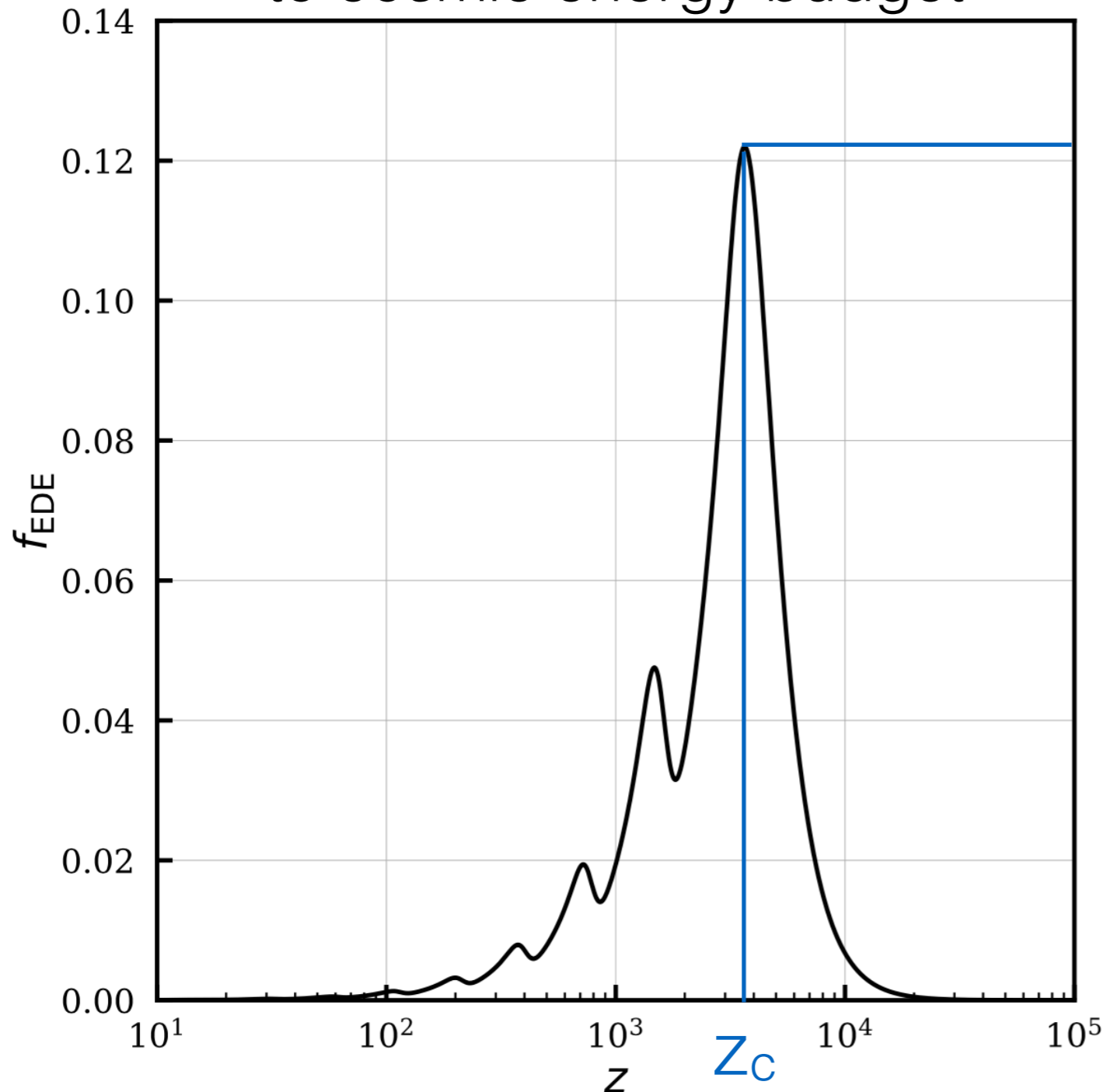
Fractional contribution of EDE  
to cosmic energy budget



# Early Dark Energy

## Parameterization

Fractional contribution of EDE  
to cosmic energy budget



Maximal contribution:  
 $f_{\text{EDE}}(z_c) \equiv (\rho_{\text{EDE}}/3M_{pl}^2 H^2)|_{z_c}$   
which occurs at redshift  $z_c$

Final parameter:  $\theta_i = \phi_i/f$   
(initial field displacement)

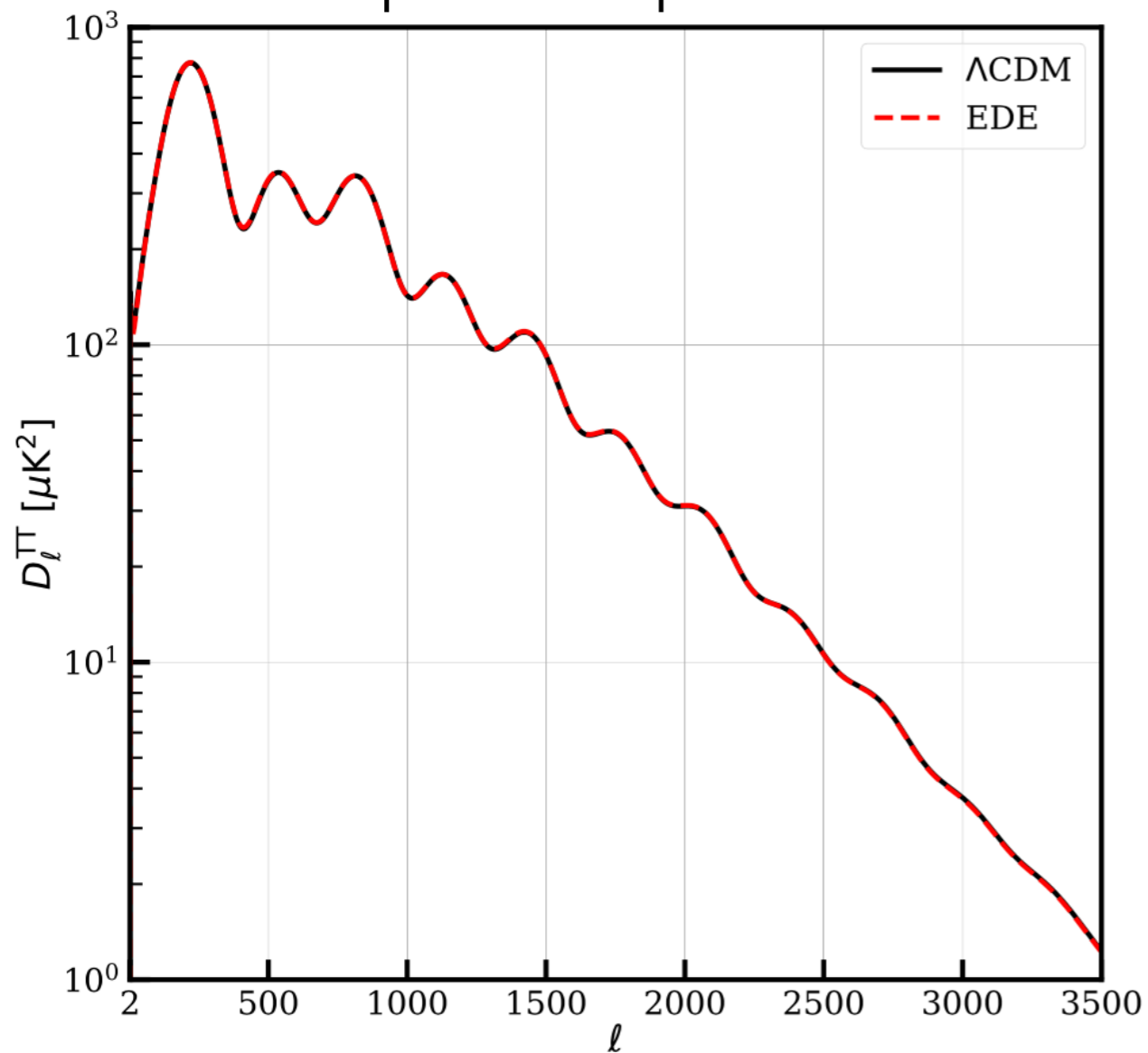
➔  $\{f_{\text{EDE}}, z_c, \theta_i\}$

N.B.: highly non-linear  
relation to physical scalar  
field parameters

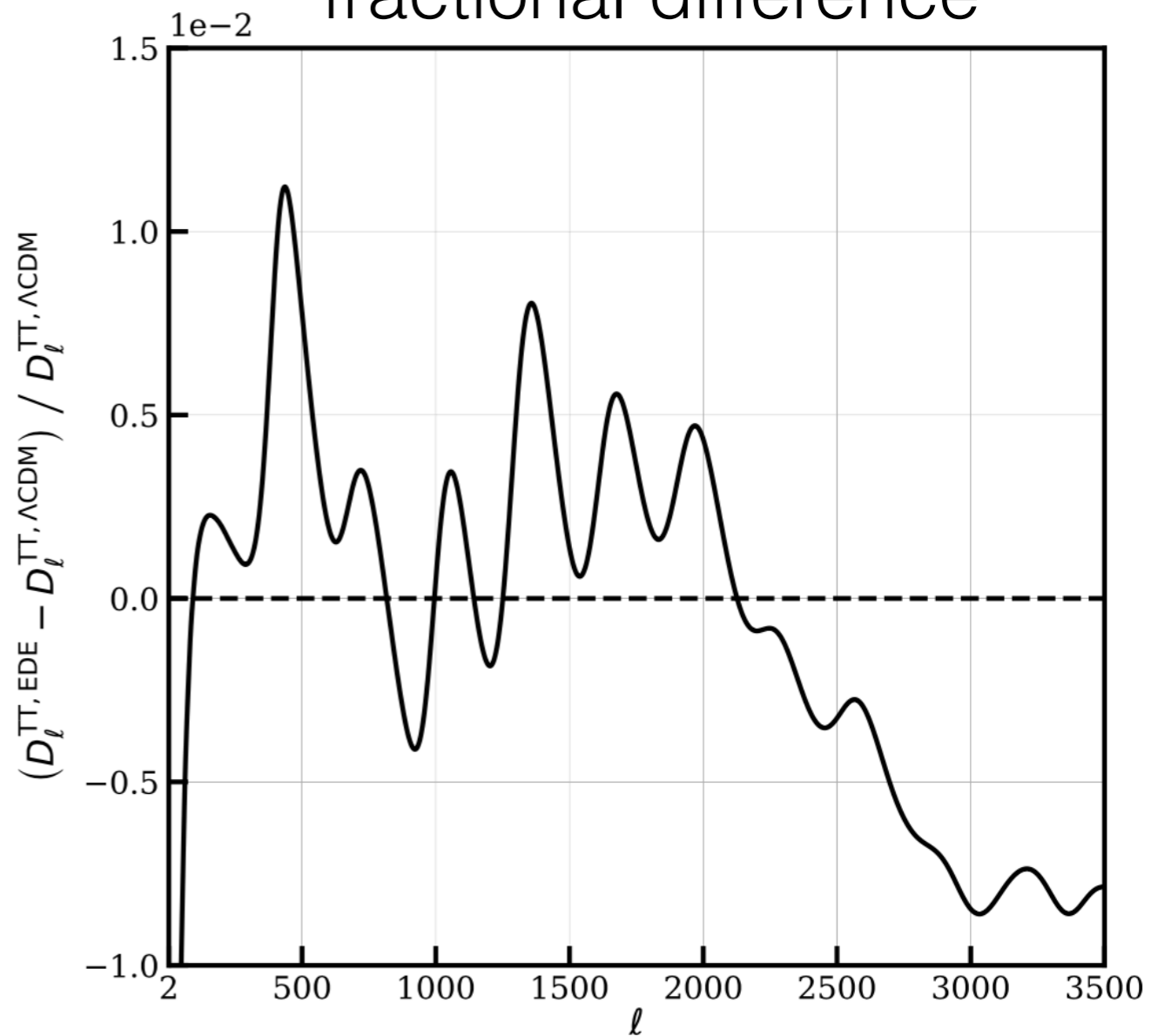
# Early Dark Energy

It maintains a good fit to CMB power spectrum data with higher  $H_0$

TT power spectrum



fractional difference



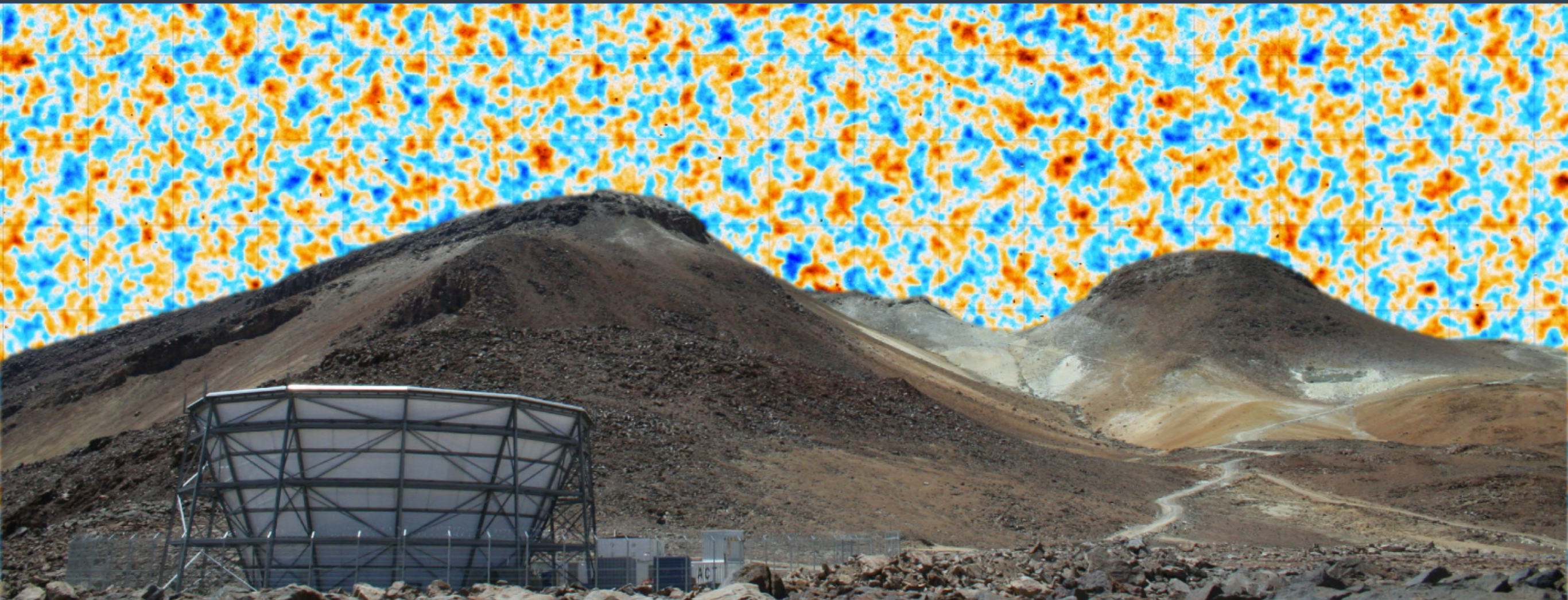
$\Lambda$ CDM model here has  $H_0 = 68.21$  km/s/Mpc

EDE model here has  $H_0 = 72.19$  km/s/Mpc

caused by  
decrease in  $r_s$



# The Atacama Cosmology Telescope

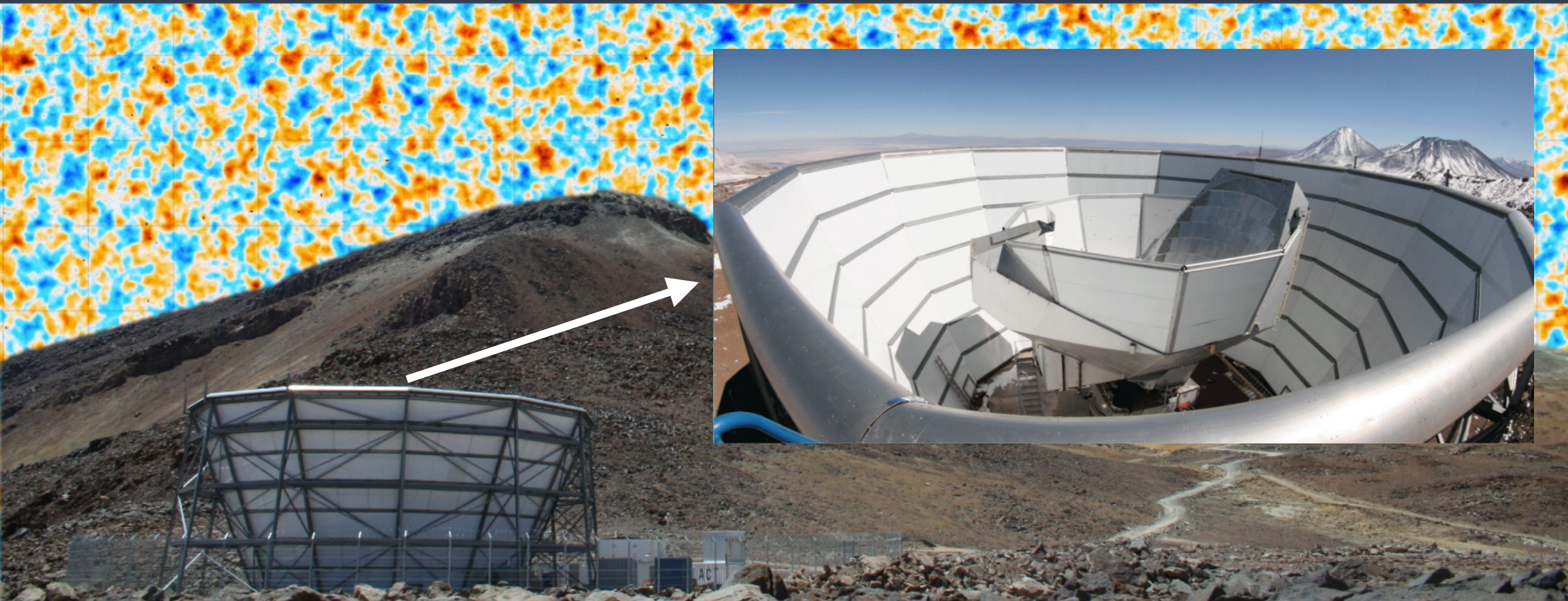


wide-area (~half-sky) multifrequency CMB survey  
observations: 2008-2022 (with some gaps for upgrades)





# The Atacama Cosmology Telescope

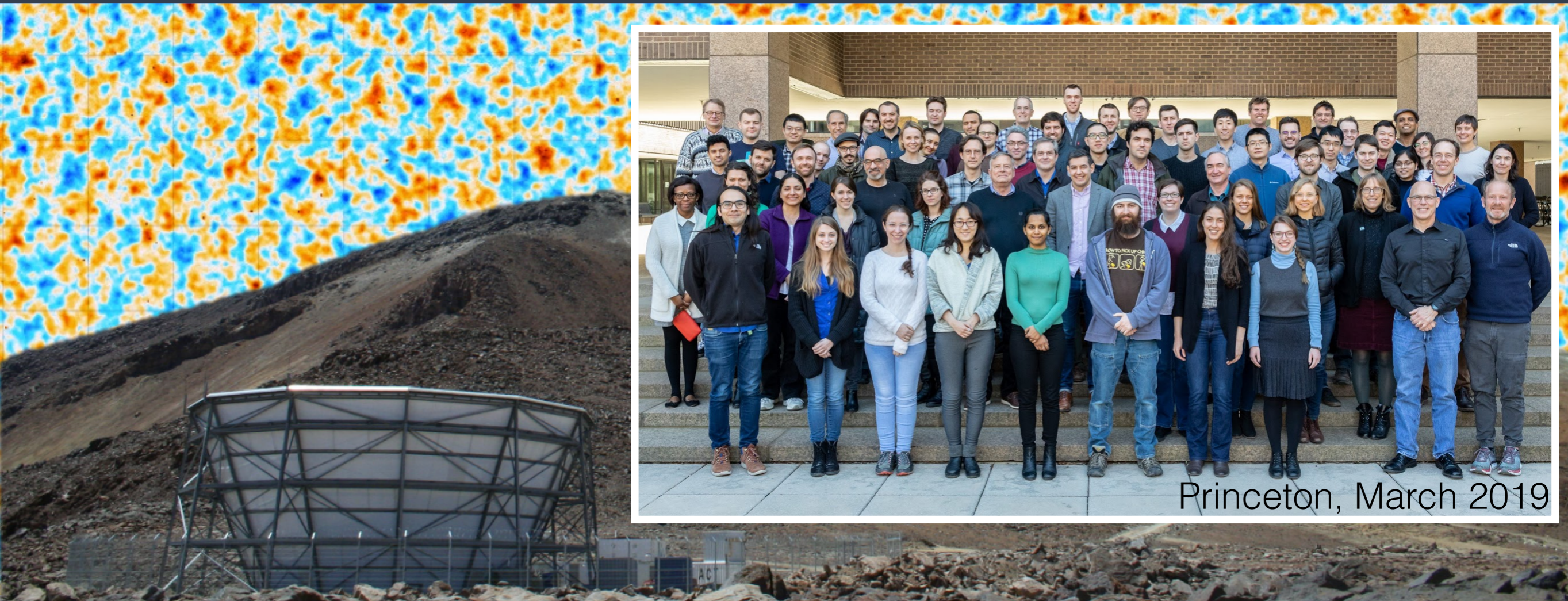


wide-area (~half-sky) multifrequency CMB survey observations: 2008-2022 (with some gaps for upgrades)



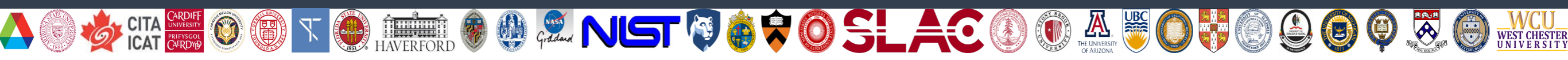


# The Atacama Cosmology Telescope



Princeton, March 2019

wide-area (~half-sky) multifrequency CMB survey observations: 2008-2022 (with some gaps for upgrades)

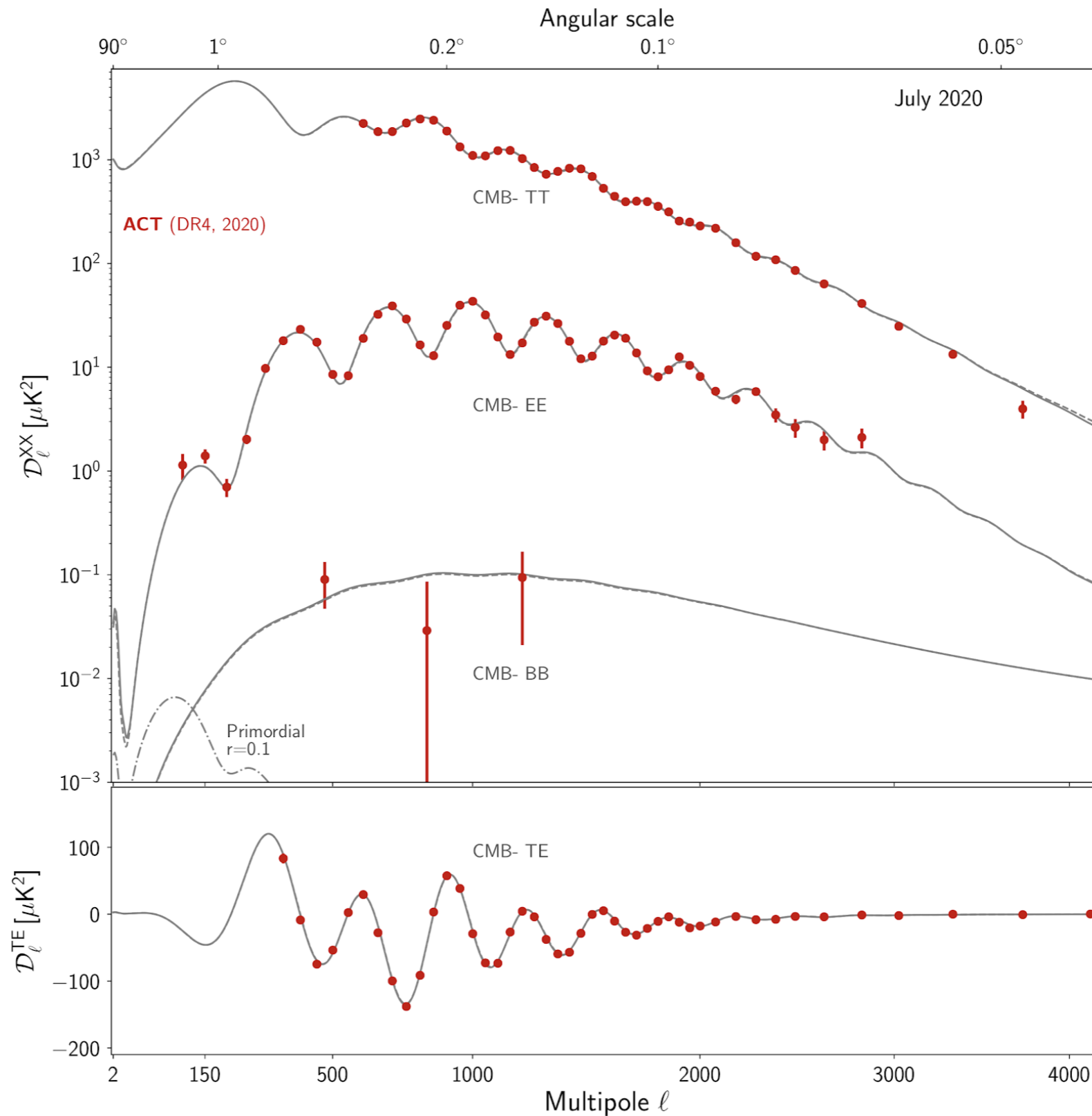




# ACT Data Release 4

## Foreground-marginalized CMB power spectra

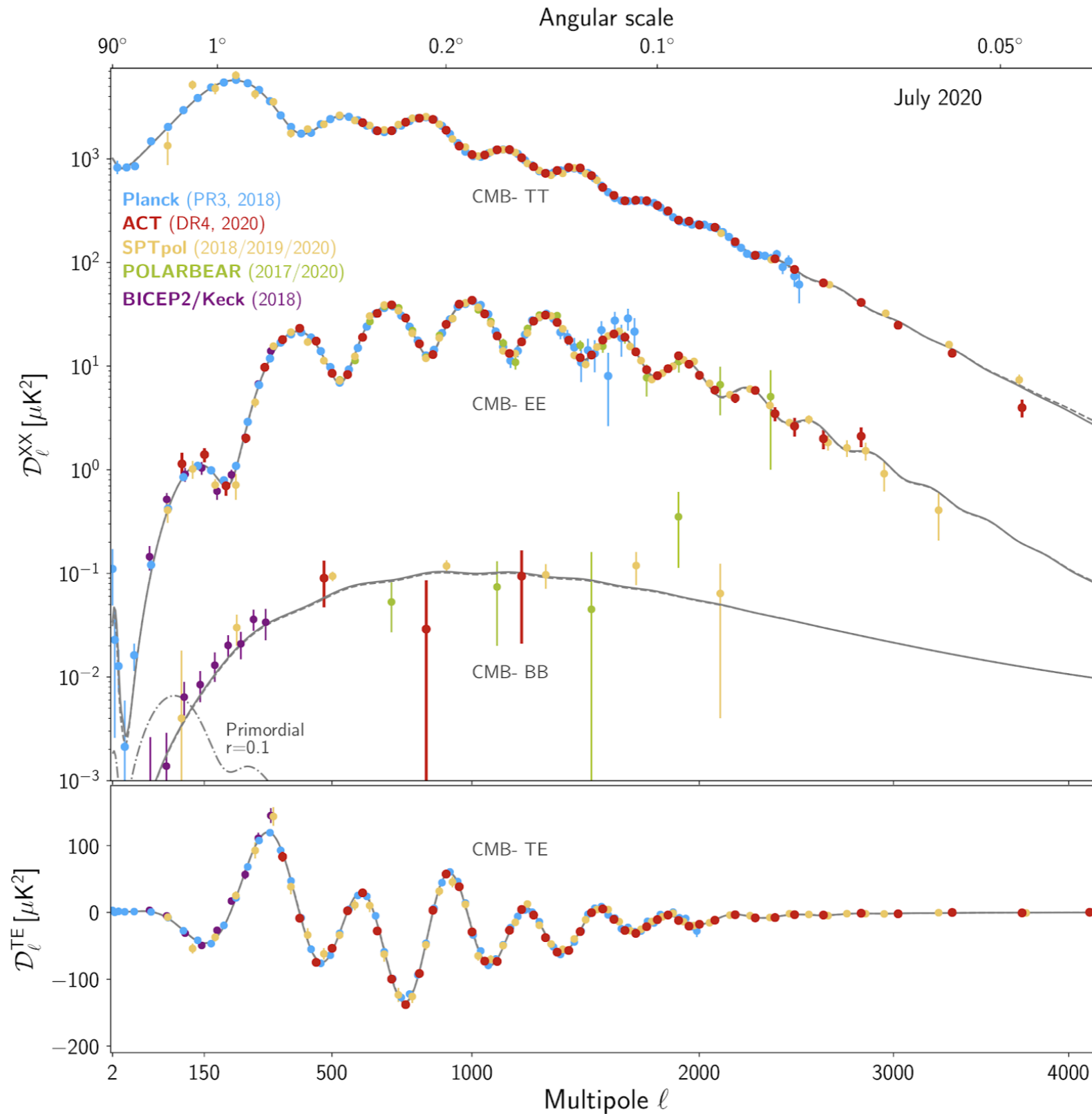
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# ACT Data Release 4

## Foreground-marginalized CMB power spectra

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# Constraints on Early Dark Energy



**JCH**, McDonough, Toomey, Alexander (2020, PRD Editors' Suggestion)

Ivanov, McDonough, **JCH**, Simonovic, Toomey, Alexander, Zaldarriaga (2020)

**JCH**, Calabrese, et al. [ACT Collaboration] (2021)

La Posta, Louis, Garrido, **JCH** (2021)





# ACT DR4 EDE Analysis

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**The Atacama Cosmology Telescope: Constraints on Pre-Recombination Early Dark Energy**

# ACT DR4 EDE Analysis

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## Motivation

- How robust are CMB-derived EDE constraints to the choice of CMB data set?
- What do we find if we replace Planck with ACT or ACT+WMAP?
- ACT and Planck are consistent at  $2.5\sigma$  in  $\Lambda$ CDM (with consistent  $H_0 \sim 67-68$  km/s/Mpc) — what about in EDE?
- N.B. we do not try to assess global concordance of any model w.r.t. all cosmological data in this analysis
- Data sets: ACT, WMAP, Planck, BAO, Planck CMB lensing

 Planck TT ( $l < 650$ )

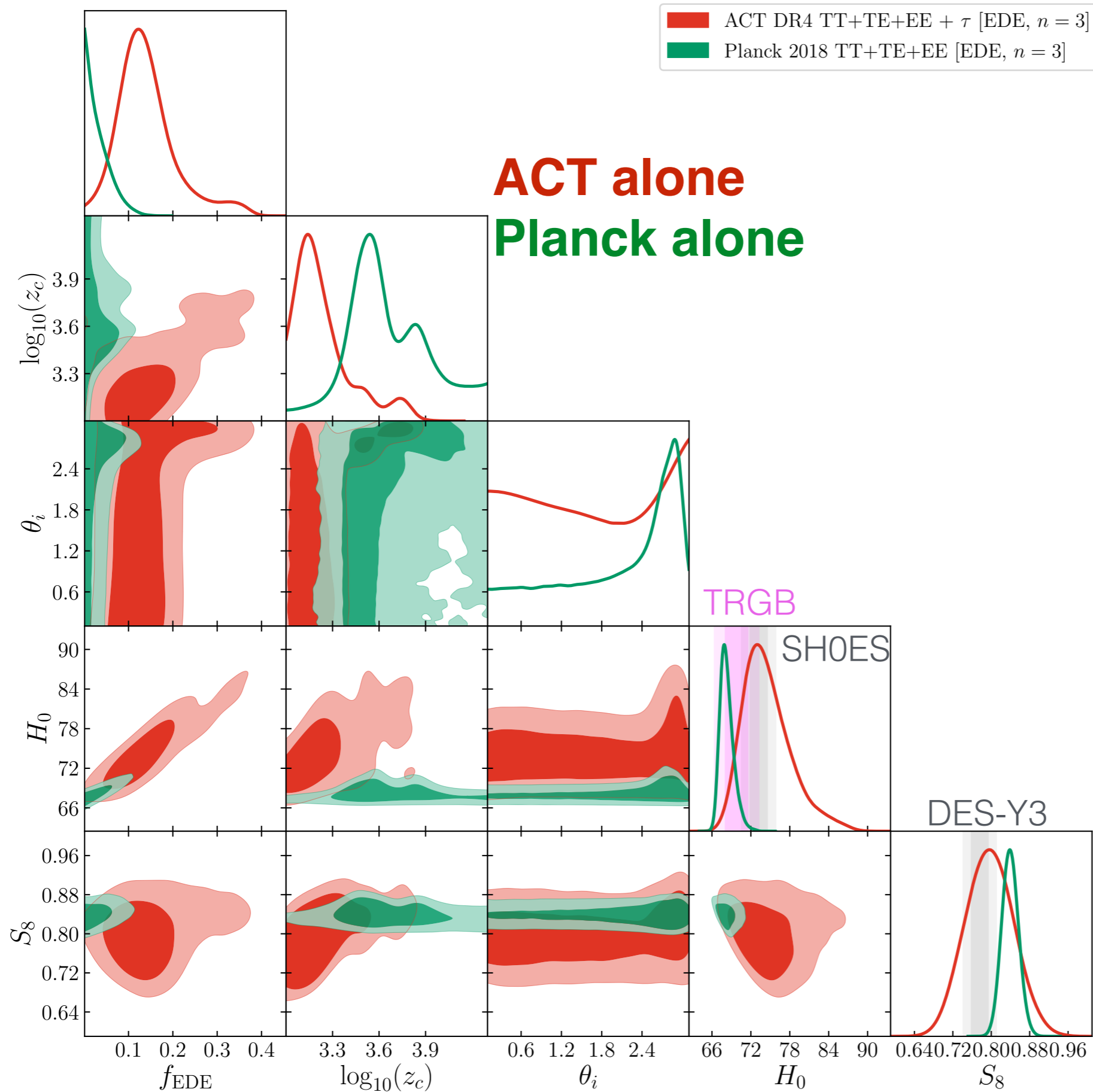
[JCH et al. \(2021\)](#)

See also [Poulin et al. \(2021\)](#)

Pipeline: **CLASS-EDE** (JCH+) + **Cobaya** (Torrado & Lewis)

# ACT DR4 EDE Results

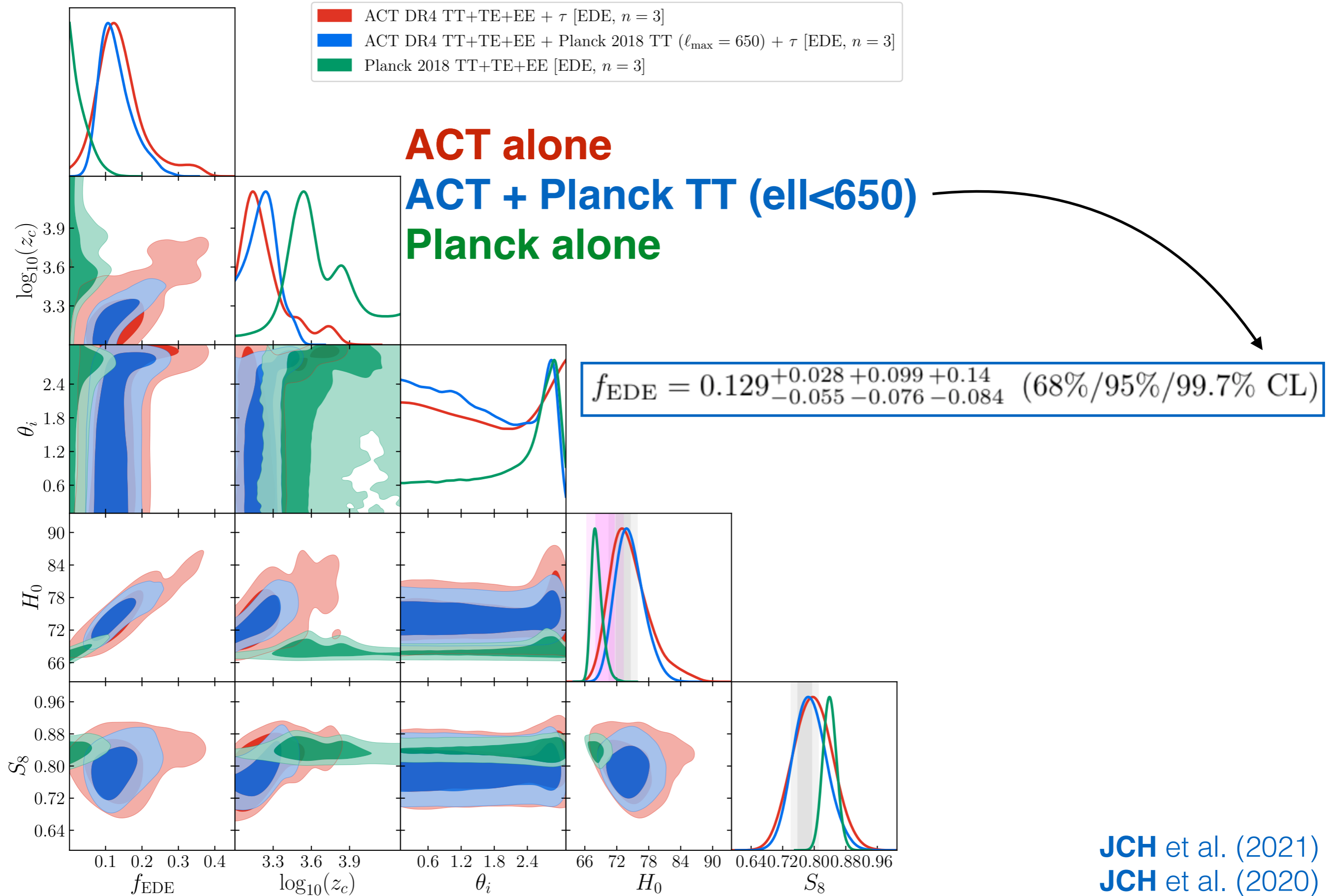
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JCH et al. (2021)  
JCH et al. (2020)

# ACT DR4 EDE Results

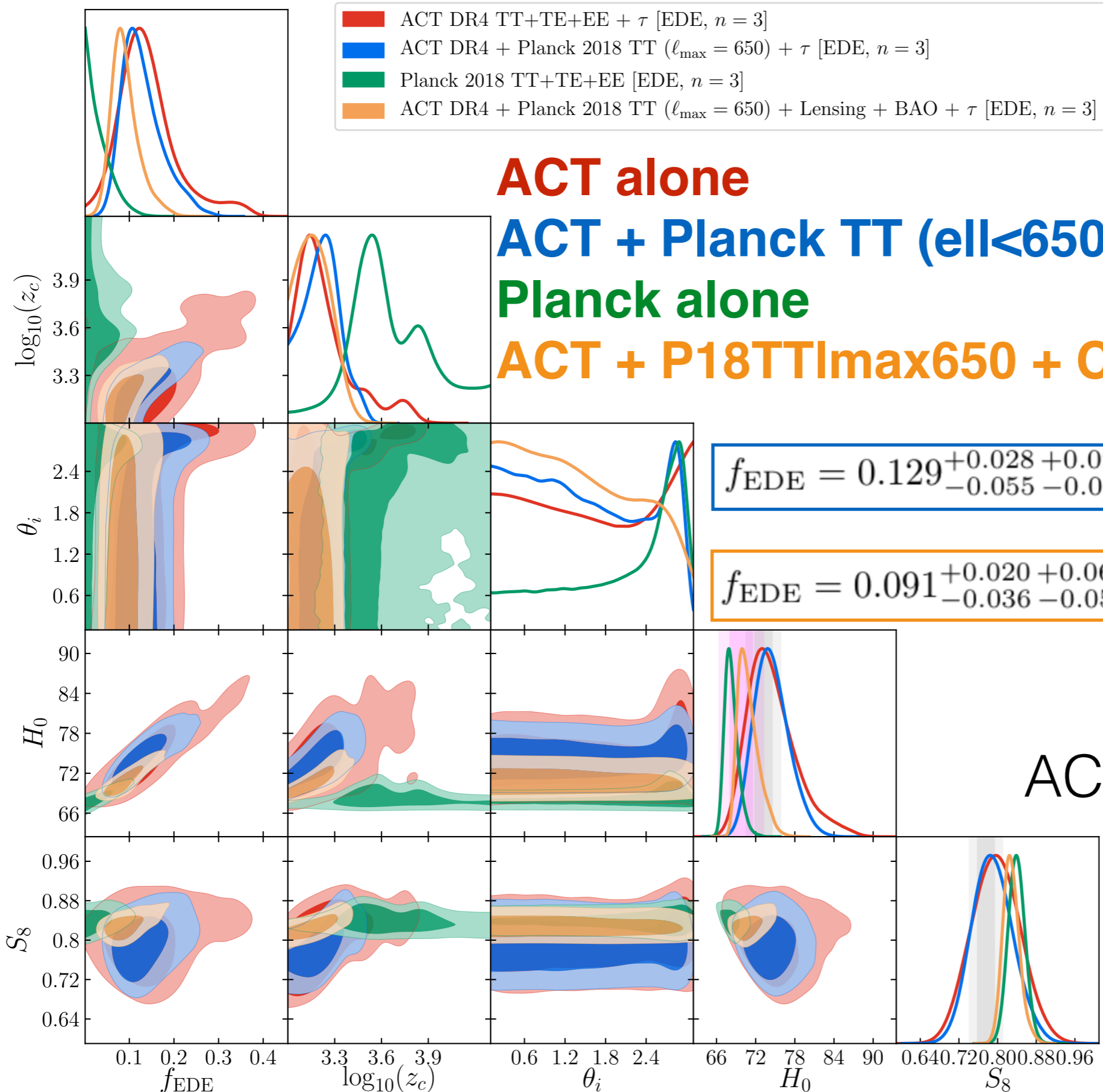
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**JCH** et al. (2021)  
**JCH** et al. (2020)

# ACT DR4 EDE Results

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**ACT alone**

**ACT + Planck TT ( $\ell < 650$ )**

**Planck alone**

**ACT + P18TTImax650 + CMB Lensing + BAO**

$$f_{\text{EDE}} = 0.129^{+0.028}_{-0.055} \quad +0.099 \quad +0.14 \quad (68\%/95\%/99.7\% \text{ CL})$$

$$f_{\text{EDE}} = 0.091^{+0.020}_{-0.036} \quad +0.069 \quad +0.11 \quad (68\%/95\%/99.7\% \text{ CL})$$

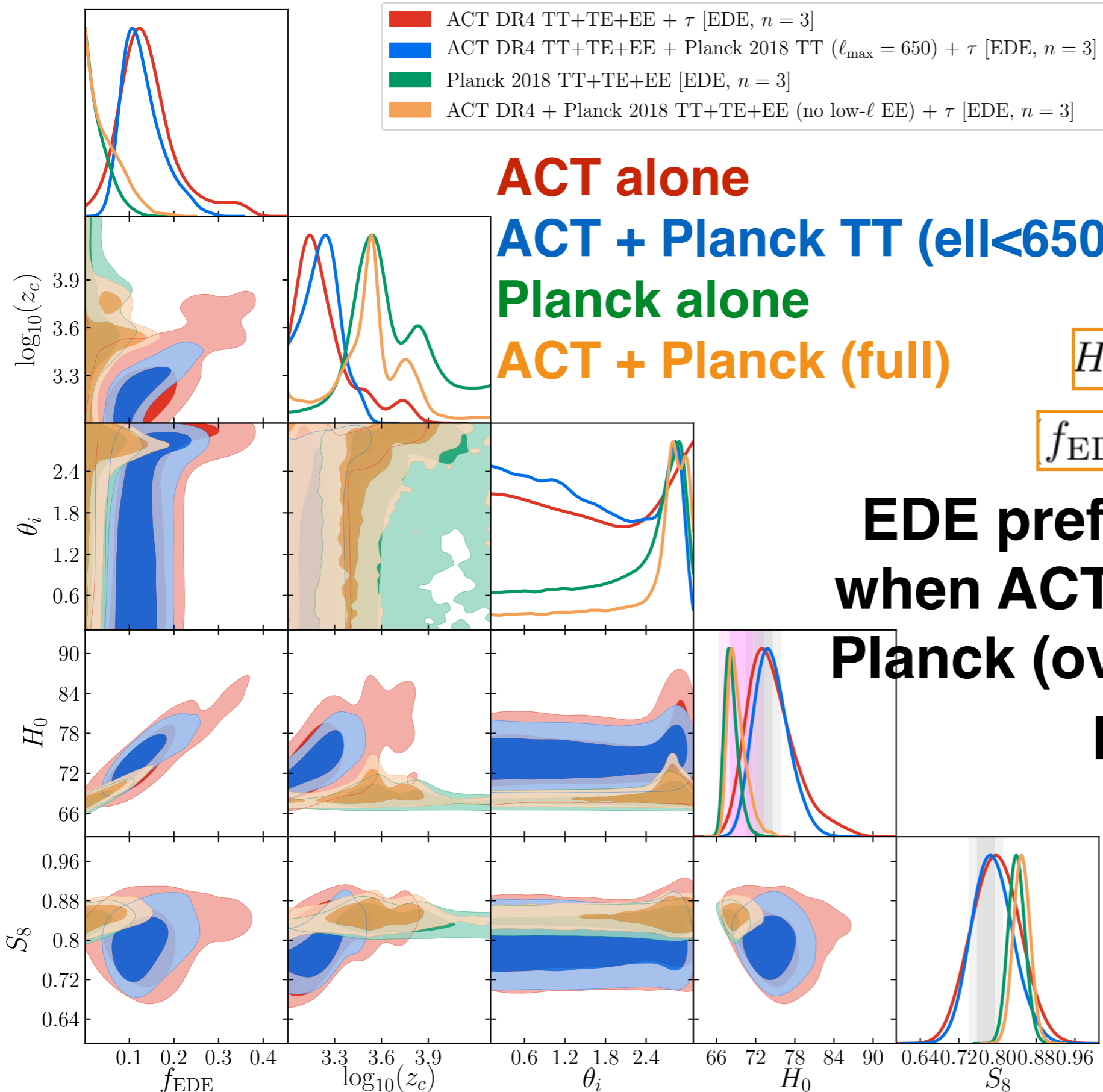
$$H_0 = 70.9^{+1.0}_{-2.0} \text{ km/s/Mpc}$$

ACT drives preference for non-zero  $f_{\text{EDE}}$  ( $>99.7\%$  CL in joint fits)



# ACT DR4 EDE Results

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- ACT DR4 TT+TE+EE +  $\tau$  [EDE,  $n = 3$ ]
- ACT DR4 TT+TE+EE + Planck 2018 TT ( $\ell_{\max} = 650$ ) +  $\tau$  [EDE,  $n = 3$ ]
- Planck 2018 TT+TE+EE [EDE,  $n = 3$ ]
- ACT DR4 + Planck 2018 TT+TE+EE (no low- $\ell$  EE) +  $\tau$  [EDE,  $n = 3$ ]

**ACT alone**  
**ACT + Planck TT ( $\ell < 650$ )**  
**Planck alone**  
**ACT + Planck (full)**

$$H_0 = 69.17^{+0.83}_{-1.70} \text{ km/s/Mpc}$$

$$f_{\text{EDE}} < 0.124 \text{ at } 95\% \text{ CL}$$

**EDE preference goes away when ACT is combined with Planck (overall constraining power still Planck-dominated)**

# Origin of ACT EDE Preference

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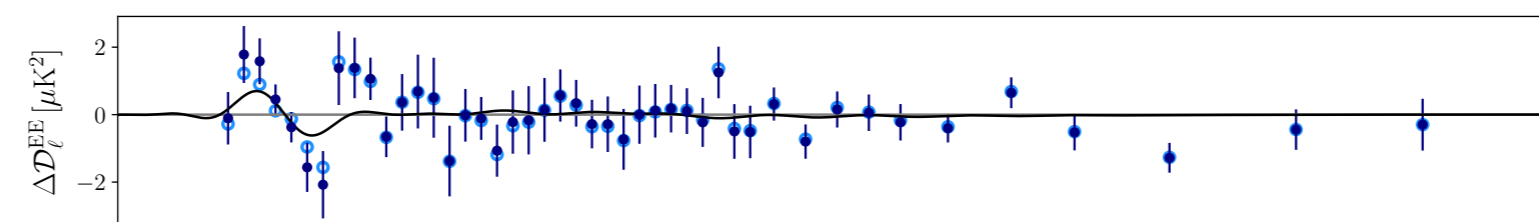
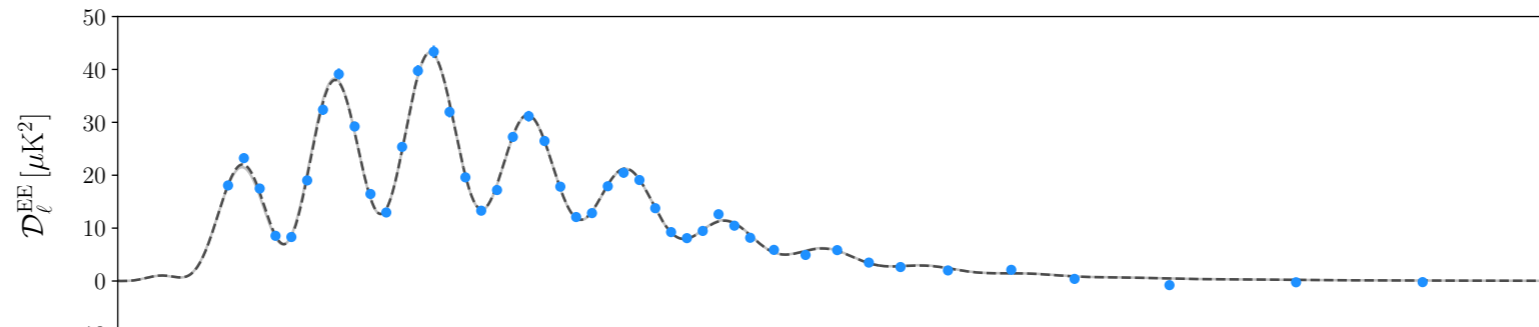
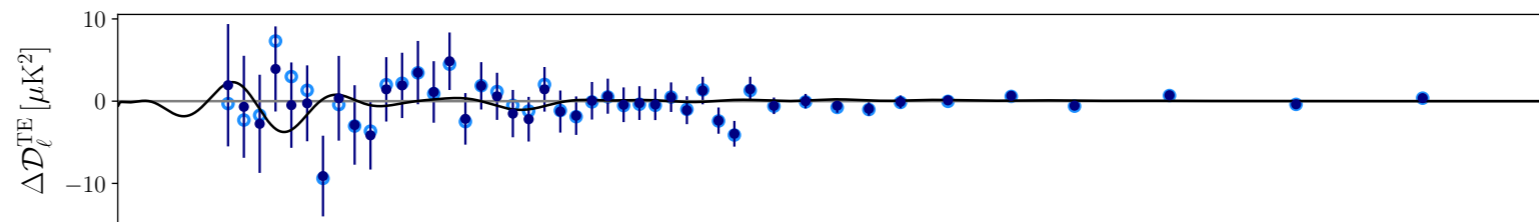
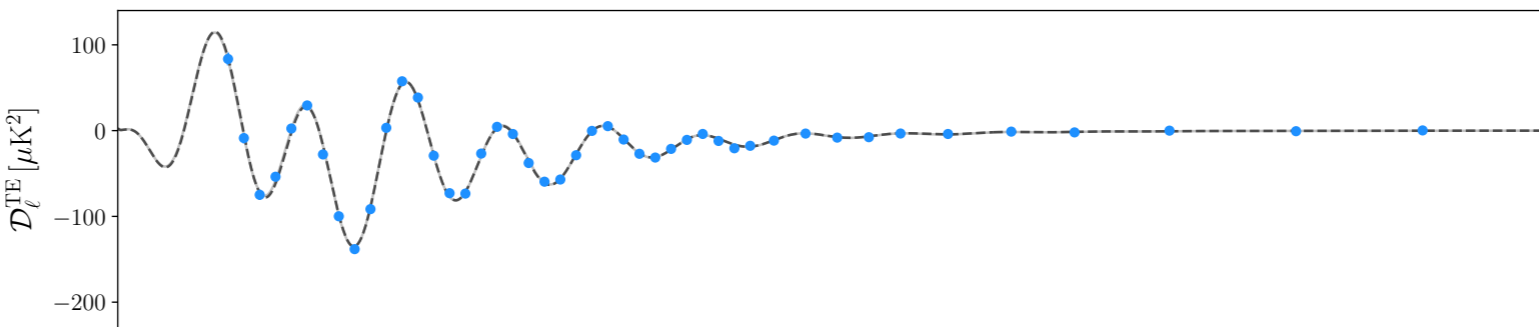
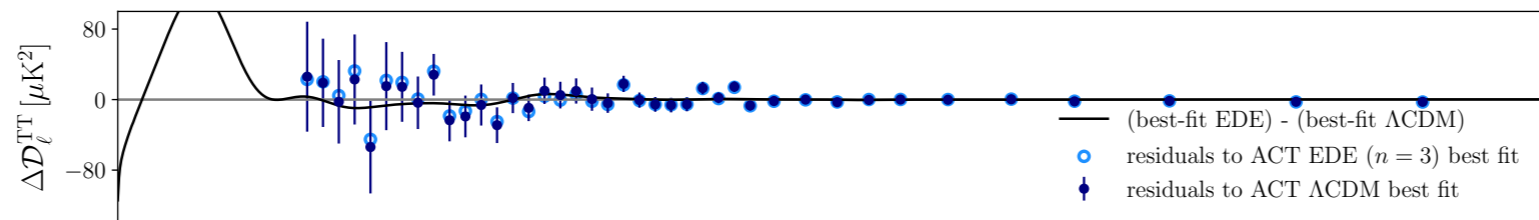
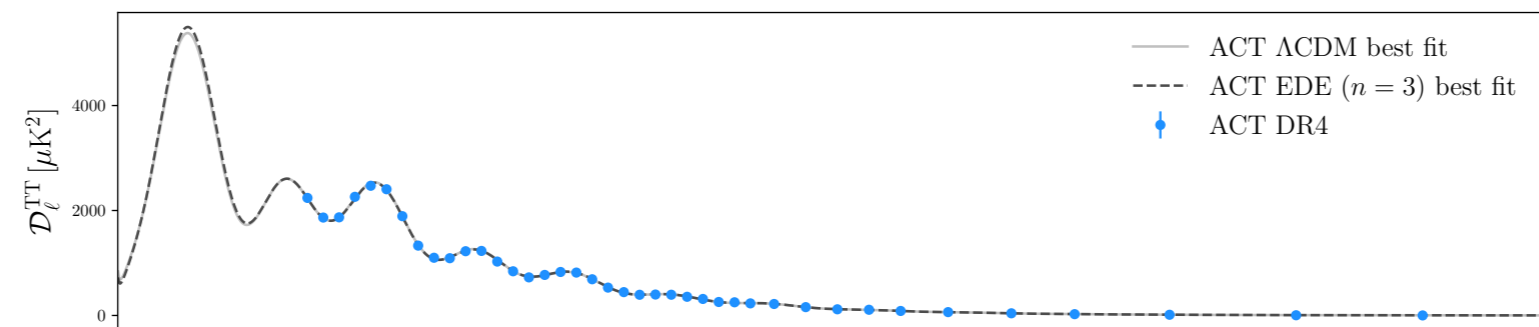
# Origin of ACT EDE Preference

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LCDM  
residuals

EDE  
residuals



$\ell$

# Origin of ACT EDE Preference

Colin Hill

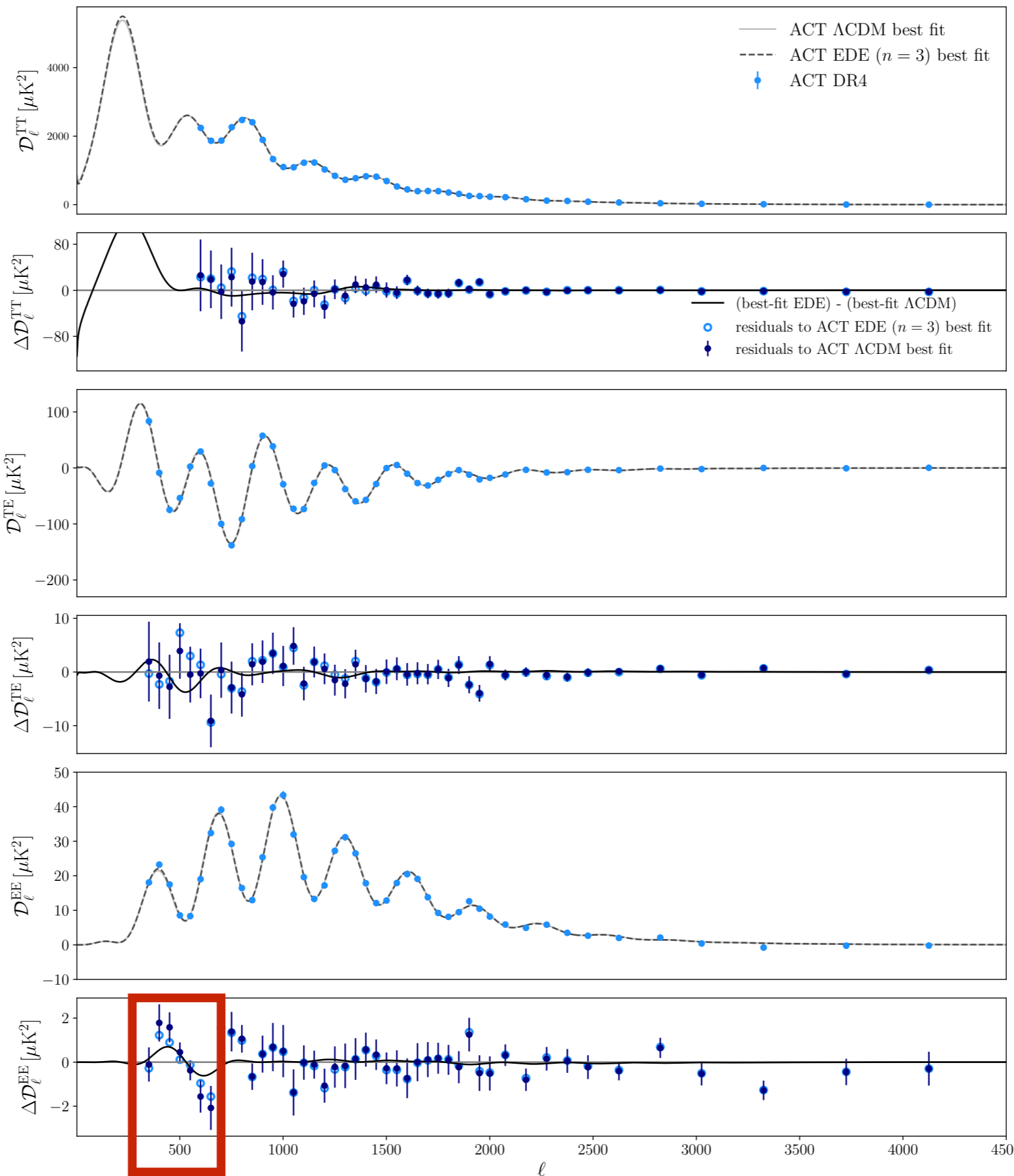
Columbia/CCA

LCDM  
residuals

EDE  
residuals

lowest  $\ell$   
bins in EE  
drive the  
preference

JCH et al. (2021)



TT

TE

EE

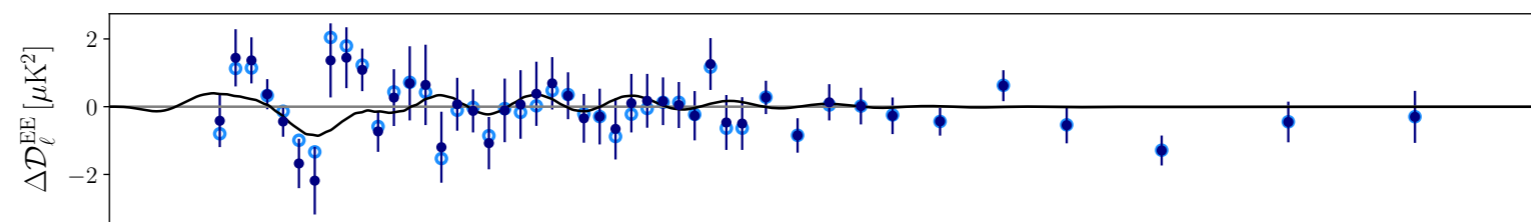
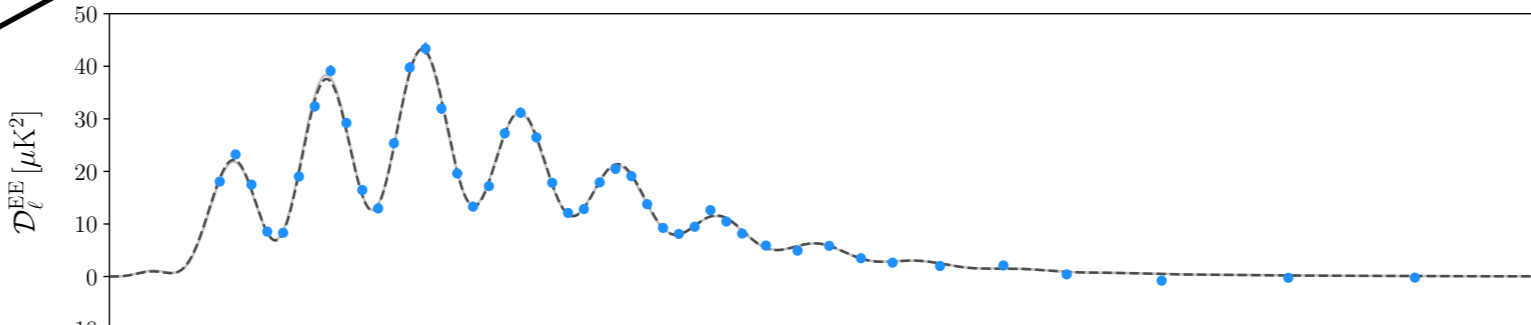
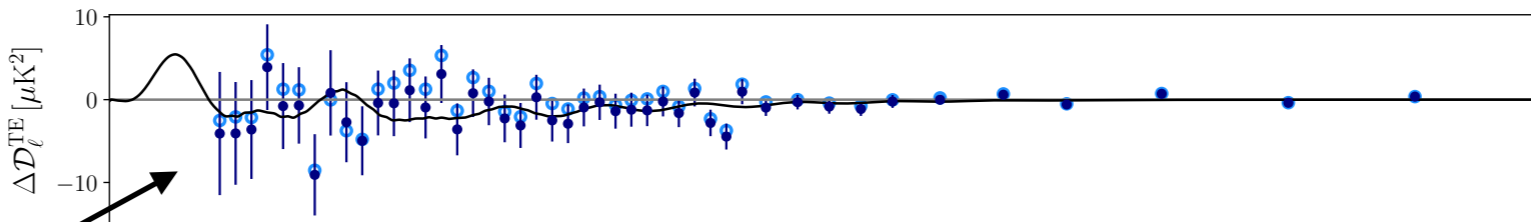
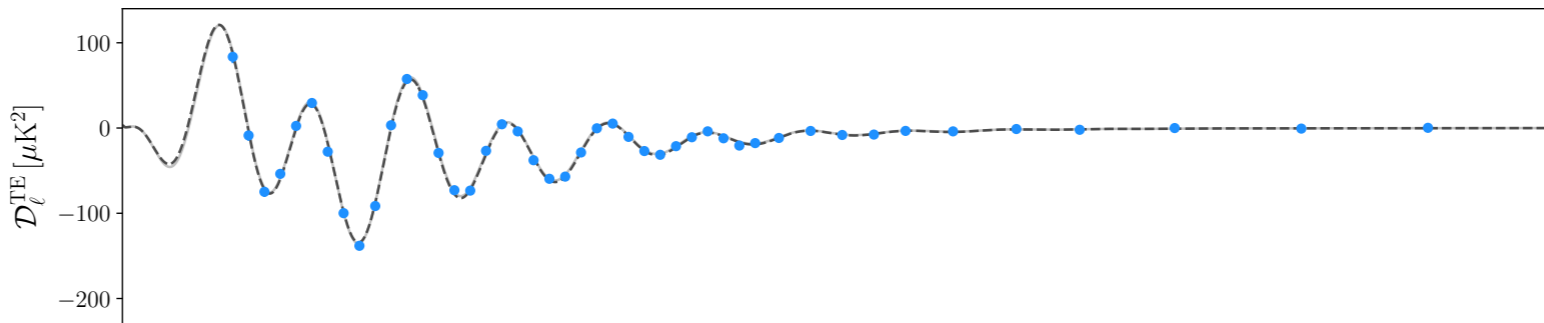
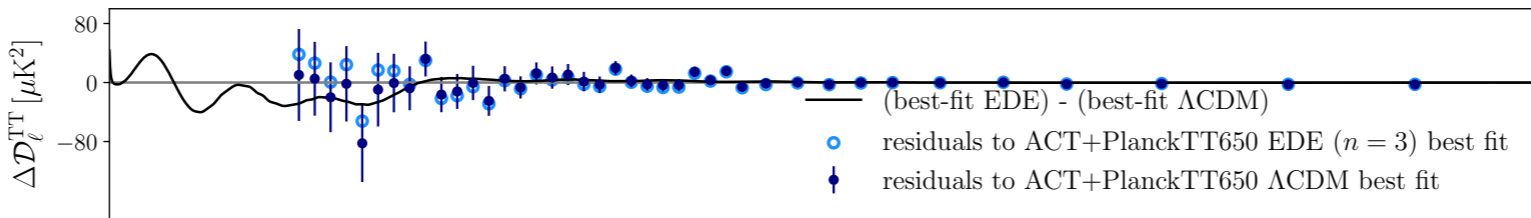
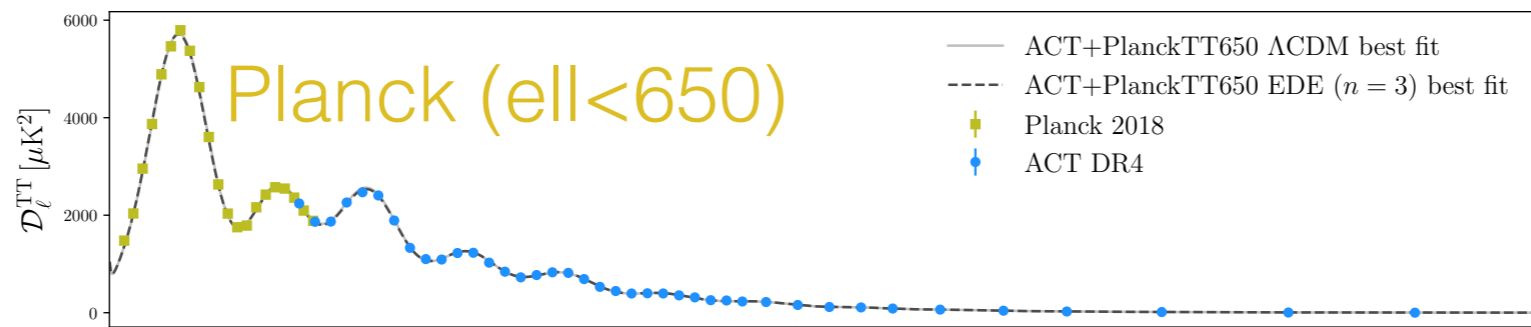
overall  
preference  
 $\sim 2.1\sigma$   
( $\Delta\chi^2 = -8.7$ )

# Origin of ACT EDE Preference

LCDM residuals

EDE residuals

TE plays an important role in driving EDE preference in joint fits



TT

TE

EE

overall preference  $\sim 3.2\sigma$   
 $(\Delta\chi^2 = -15.4)$



# Consistent with SPT-3G?

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Analysis using public SPT-3G TE/EE data

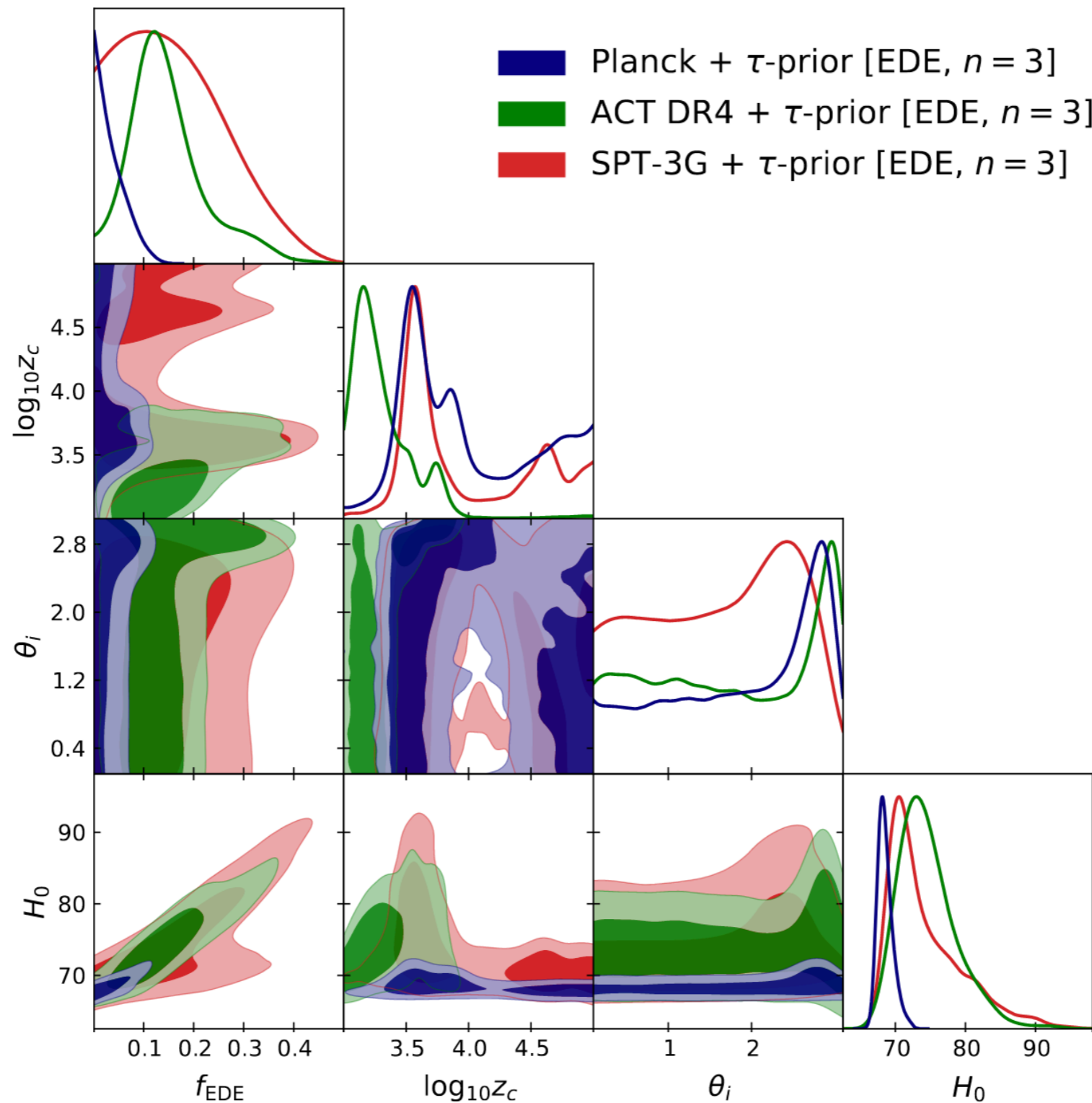
<https://arxiv.org/abs/2112.10754>; Phys. Rev. D 105, 083519

La Posta, Louis, Garrido, **JCH** (2021) ; SPT-3G data from Dutcher et al. (2021)

# Consistent with SPT-3G

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Analysis using public SPT-3G TE/EE data



# Consistent with SPT-3G

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Analysis using public SPT-3G TE/EE data

Inclusion of full Planck TT data still dominates overall constraining power and removes preference for non-zero EDE

Parameters	SPT-3G + Planck	SPT-3G + PlanckTT650 + ACT DR4	SPT-3G + Planck + ACT DR4
$f_{\text{EDE}}$	$< 0.088$	$0.121^{+0.040}_{-0.064}$	$< 0.107$
$H_0$ [km/s/Mpc]	$68.6^{+0.7}_{-1.1}$	$74.2^{+2.3}_{-3.0}$	$68.9^{+0.7}_{-1.6}$

↓  
2.6 $\sigma$  hint

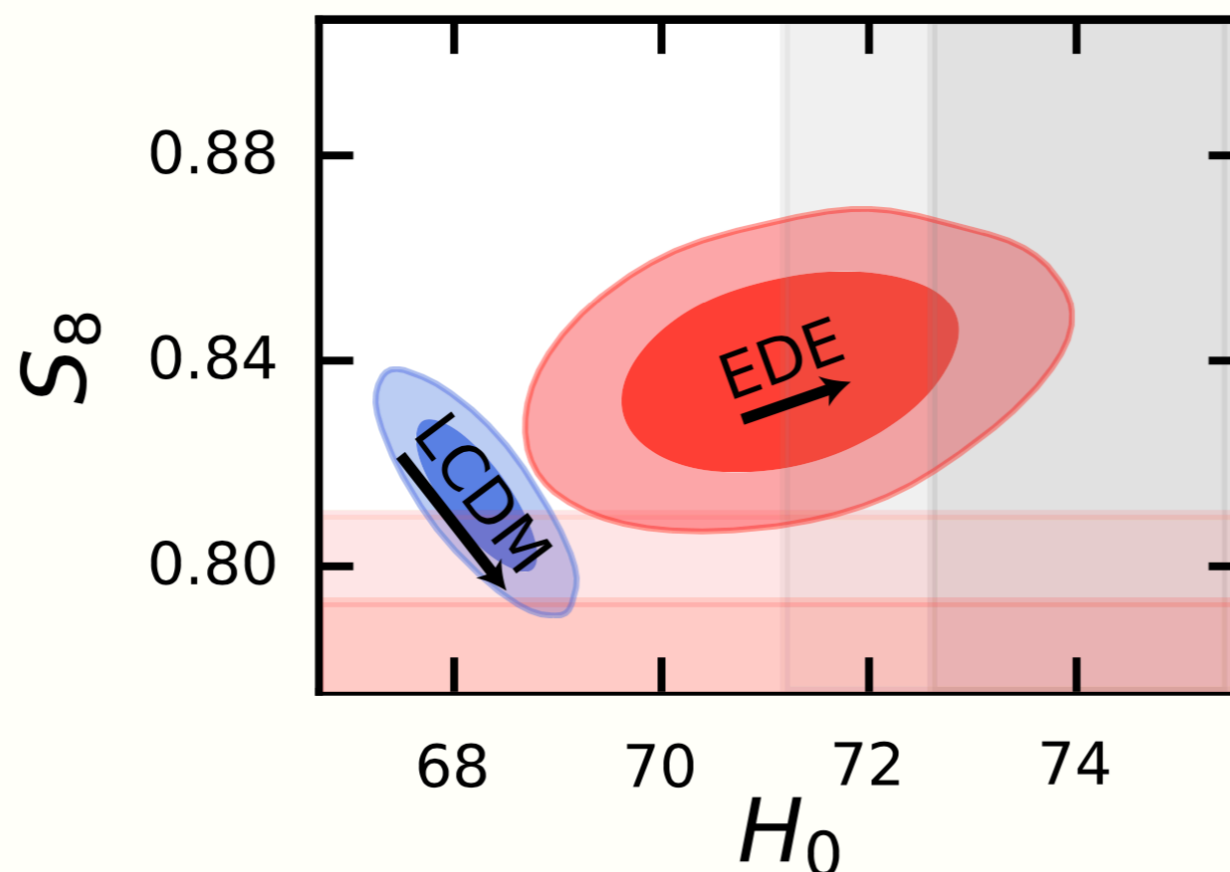
Upcoming data from ACT + SPT will be very interesting!

# EDE Puzzles & Problems

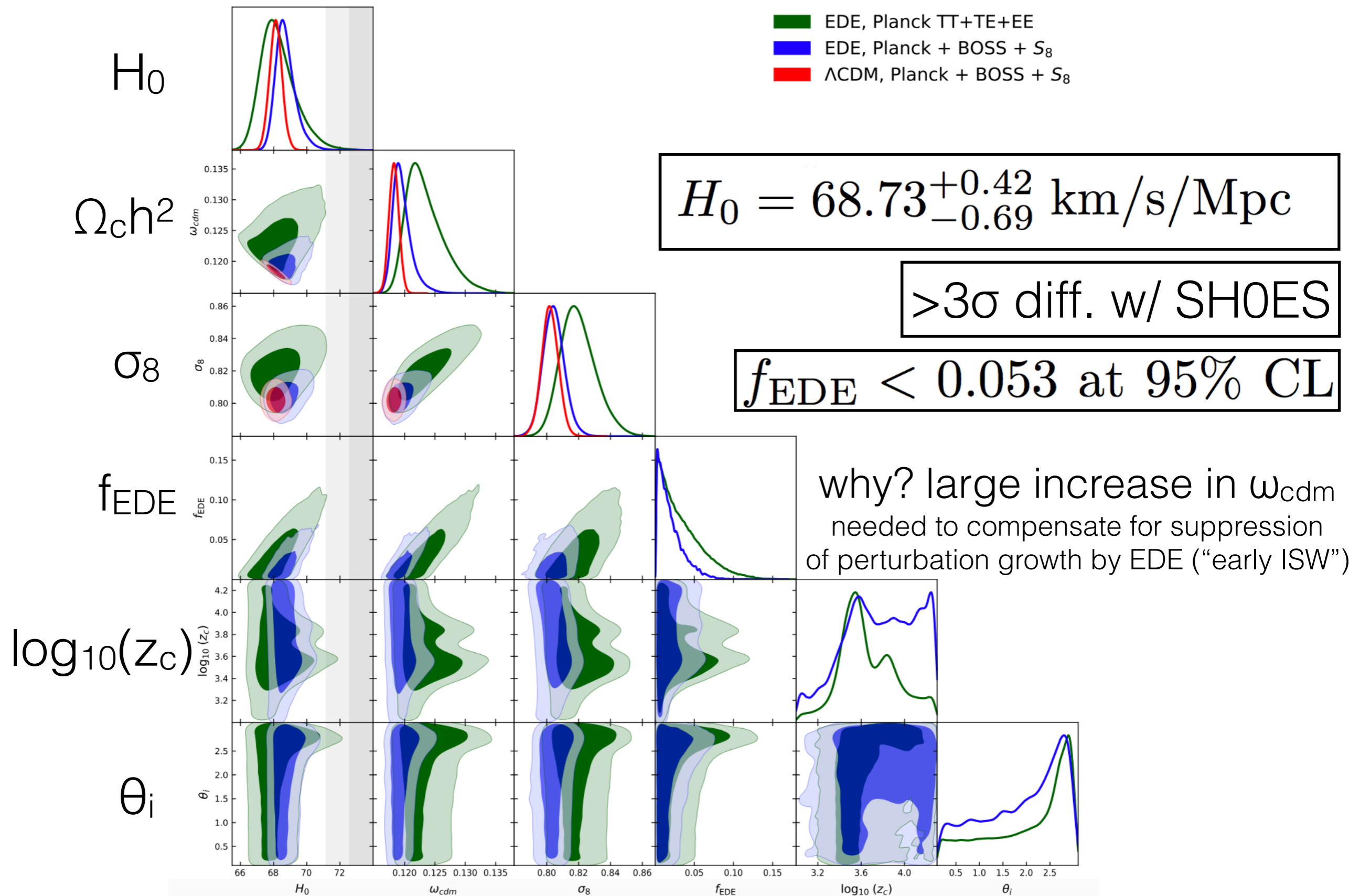


# EDE Puzzles & Problems

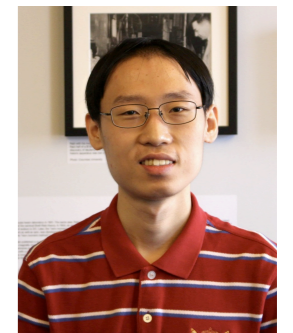
- Coincidence problem: why should these new dynamics appear near  $z_{\text{eq}}$ ? [ $\rightarrow V(\phi), V'(\phi)$ ]
- Initial conditions: axion-like field must start near top of cosine to fit Planck (e.g., Lin, Benevento, Hu, Raveri (2019)) [ $\rightarrow V''(\phi)$ ]
- Tension-trading:  $H_0$  is increased at the cost of adding significantly more dark matter, hence raising  $S_8$



# Planck + BOSS (EFT) + DES/HSC/KiDS ( $S_8$ )

Planck + BOSS (EFT) + DES/HSC/KiDS ( $S_8$ )

# Early Dark Sector



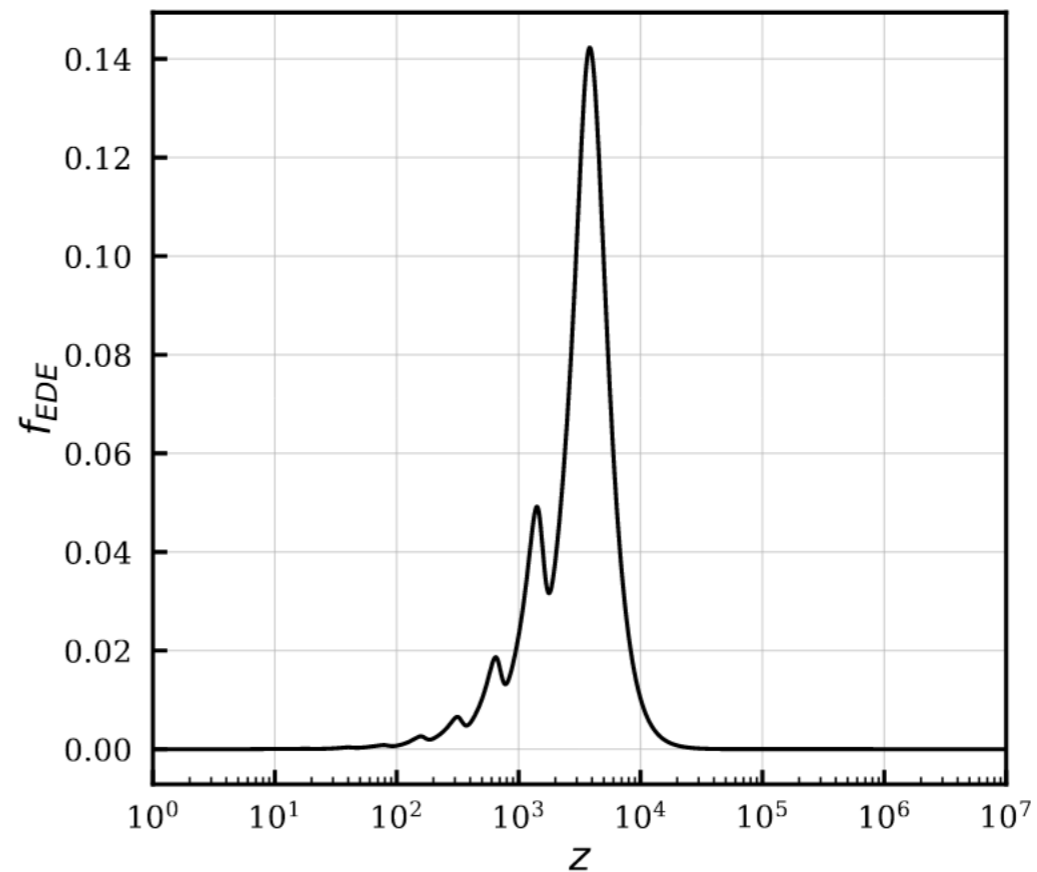
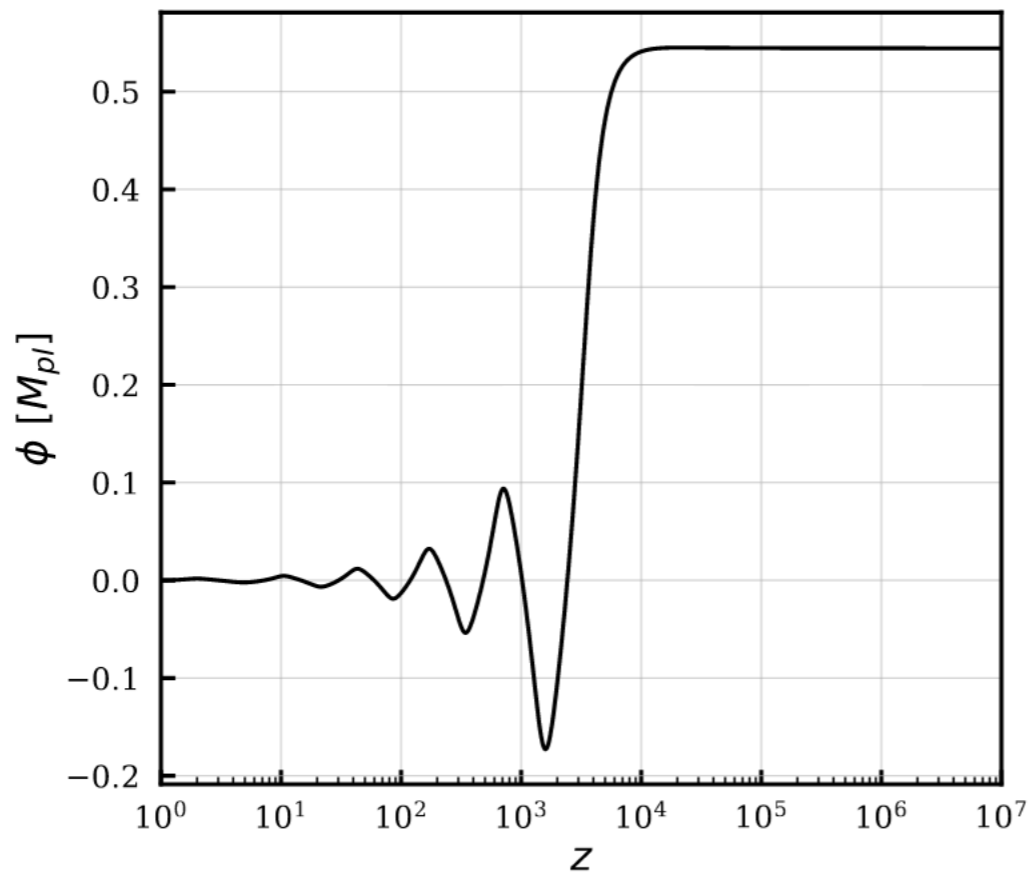
2112.04510 w/ Evan McDonough, Meng-Xiang Lin, Wayne Hu, Shengjia Zhou  
+ in prep. w/ Lin, McDonough, Hu



# Early Dark Sector

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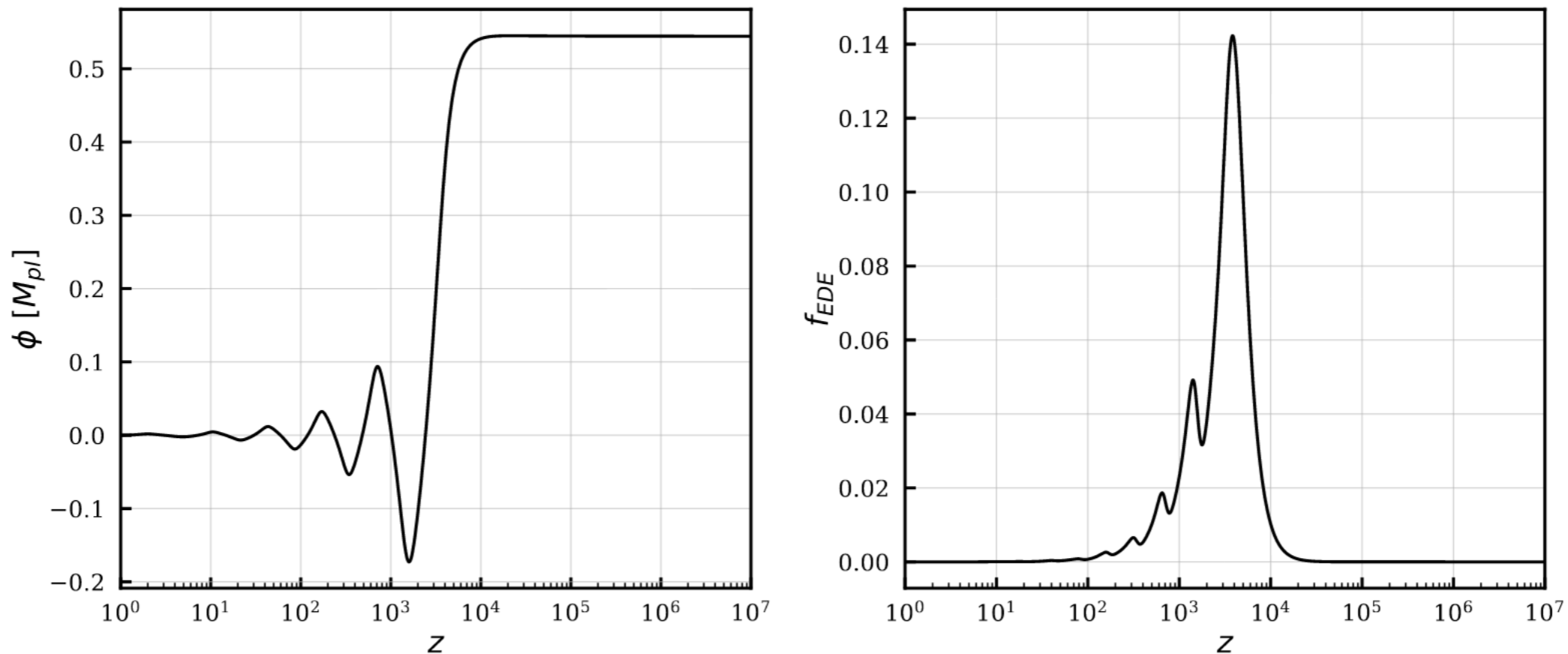
Theoretical motivation: in  $H_0$ -tension-resolving EDE models, the scalar  $\phi$  generically undergoes a Planckian field excursion



# Early Dark Sector

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Theoretical motivation: in  $H_0$ -tension-resolving EDE models, the scalar  $\phi$  generically undergoes a Planckian field excursion

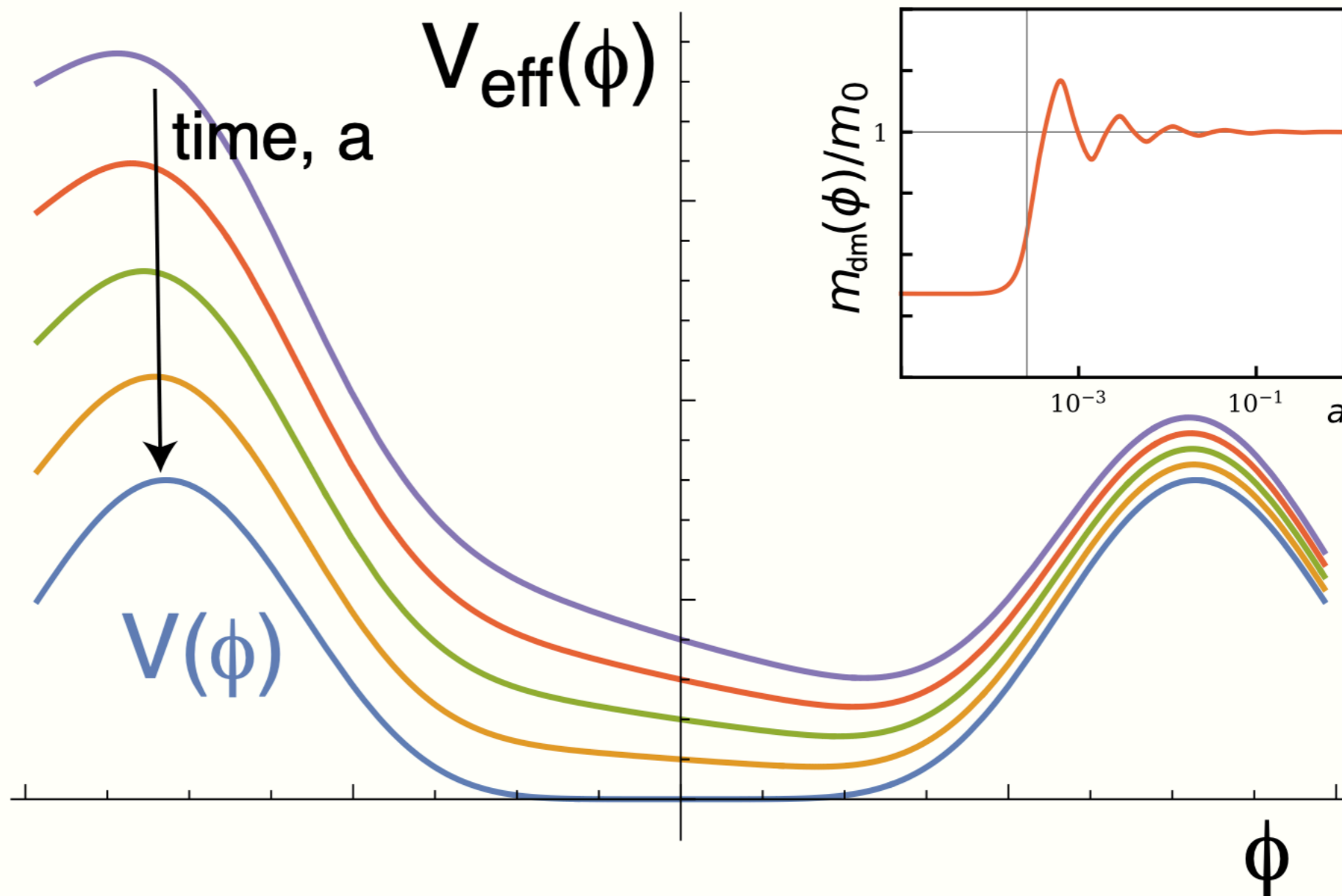


- Swampland distance conjecture (SDC) [Vafa, Ooguri]:  
breakdown of EFT that occurs at Planckian field excursions is encoded in an exponential sensitivity of the mass spectrum of the effective theory  $\rightarrow$  suggests  $m_{DM} \sim e^{\phi/M_{pl}}$

# Early Dark Sector

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- Field dependent dark matter mass:  $m_{\text{dm}}(\phi) = m_0 e^{c\phi}$
- Effective potential:  $V_{\text{eff}} = V_0 + m_{\text{dm}}(\phi)n_{\text{dm}}$



# Early Dark Sector

Effects on growth of structure (holding  $\Omega_c h^2$  fixed)

- $c < 0$  ( $c > 0$ ): DM mass lighter (heavier) at  $z > z_c$ , and thus matter-radiation equality occurs later (earlier), leading to less (more) growth by today



# Early Dark Sector

Effects on growth of structure (holding  $\Omega_c h^2$  fixed)

- $c < 0$  ( $c > 0$ ): DM mass lighter (heavier) at  $z > z_c$ , and thus matter-radiation equality occurs later (earlier), leading to less (more) growth by today
- On small scales,  $\phi$  mediates a fifth force that scales as  $\sim c^2$  and enhances growth

$$G_{\text{eff}} = G_N \left( 1 + \frac{2c^2 k^2}{k^2 + a^2 d^2 V / d\phi^2} \right)$$

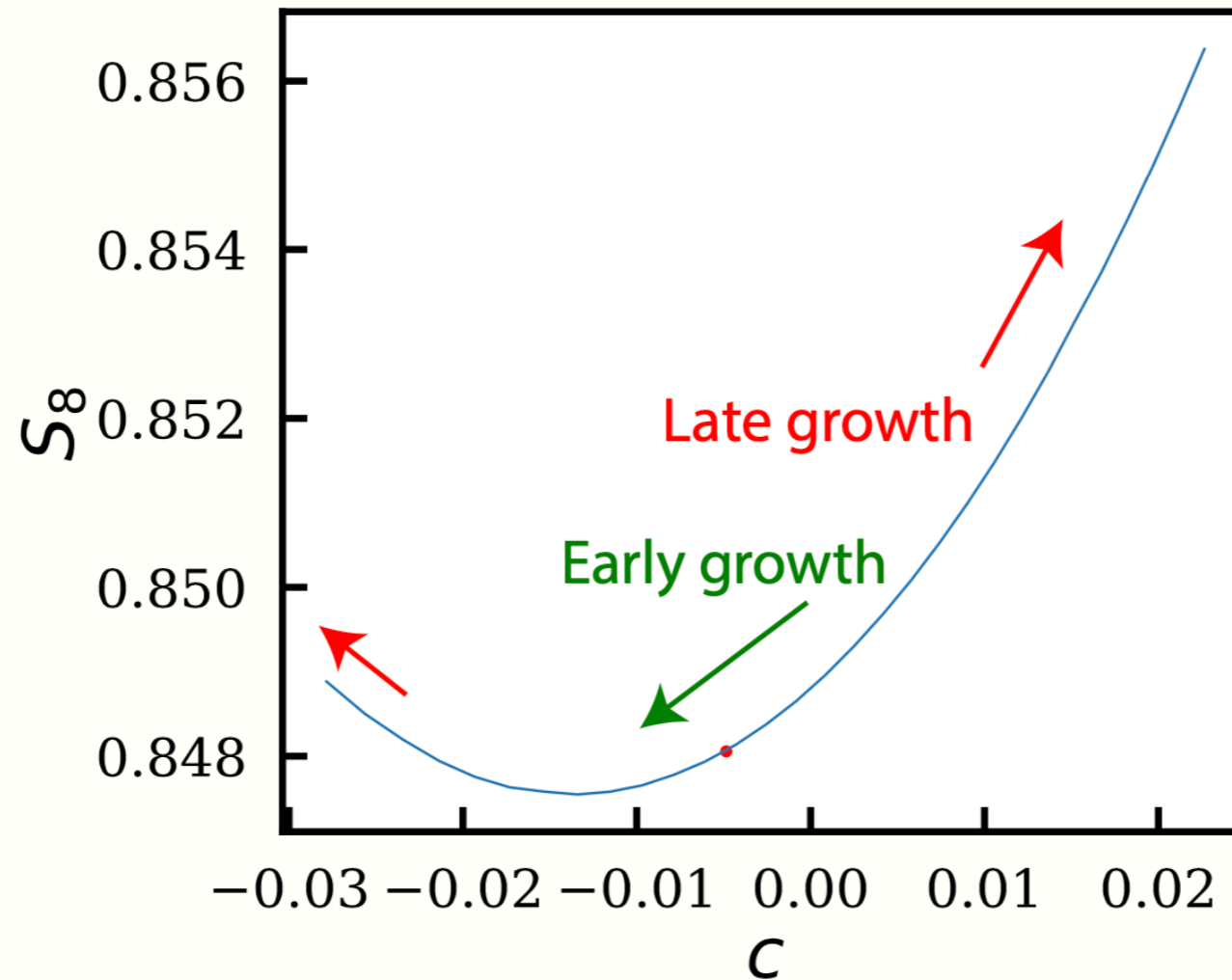
—————→  $G_{\text{eff}} = G_N (1 + 2c^2)$   
high-k  
limit

# Early Dark Sector

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## Competing $S_8$ effects

- Early: linear in  $c$  and  $c < 0$  reduces  $S_8$
- Late: quadratic in  $c$  and all  $c$  enhances  $S_8$

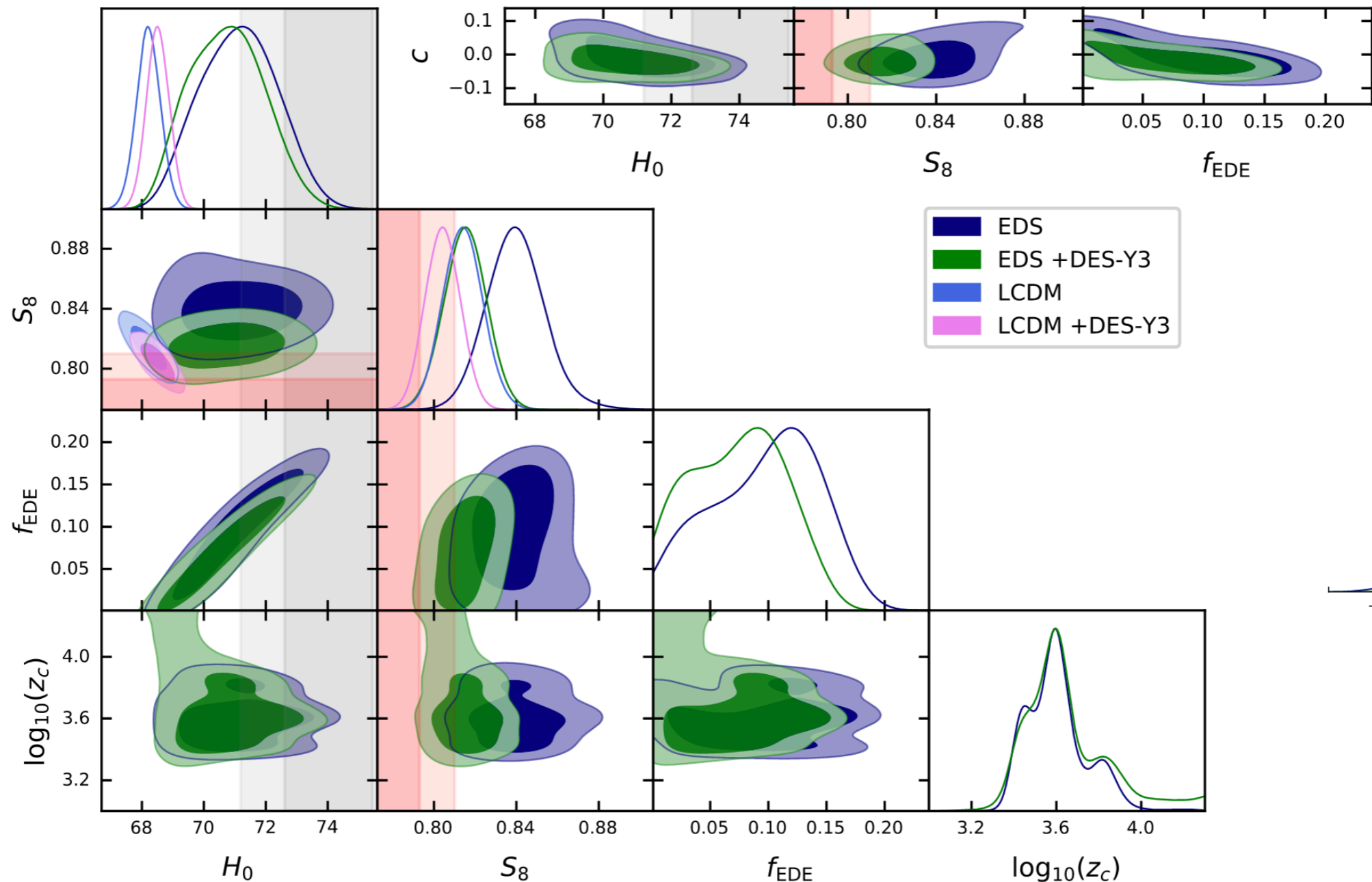


- Combination: only small ability to lower  $S_8$

# Early Dark Sector

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Additional low- $S_8$ , high- $H_0$  parameter space is now allowed:  
we've reduced the "tension of tensions", but fit does not substantially  
improve over  $\Lambda$ CDM



Data sets = Planck + BAO +  
SNIa + SH0ES (+DES)

# Early Dark Sector 2

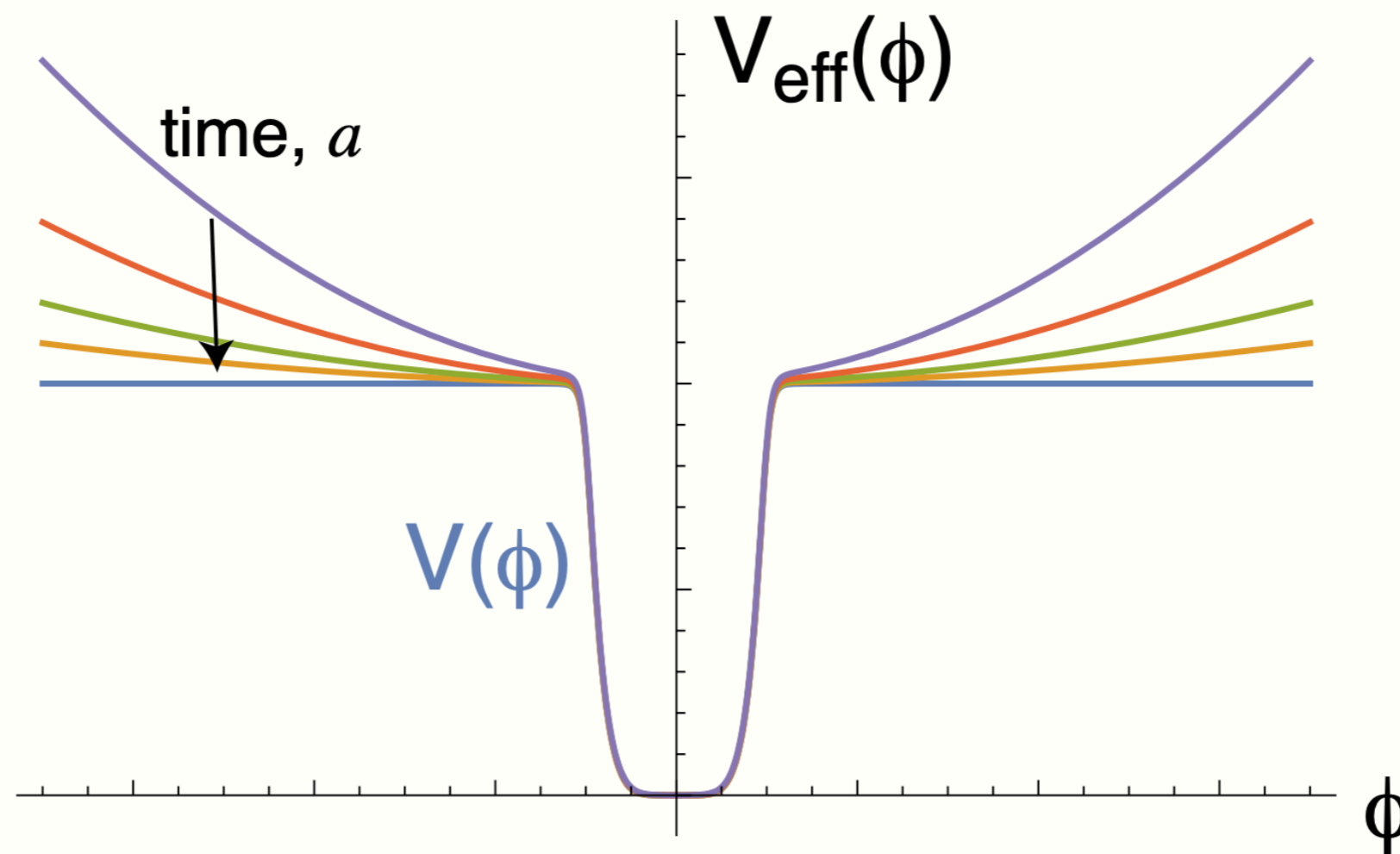
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Can we perhaps solve (at least) the EDE “coincidence” and initial-conditions problems?

# Early Dark Sector 2

## Toy solution

- Problem for acceptable  $\Delta m_{\text{dm}}/m_{\text{dm}}$  and generic initial conditions, slope of bare potential too high to trigger off coupling
- Flatten the bare potential into a plateau and change  $m(\phi)$ ,  $V_{\text{eff}} \propto \rho_{\text{dm}}$  and overcomes Hubble drag near equality



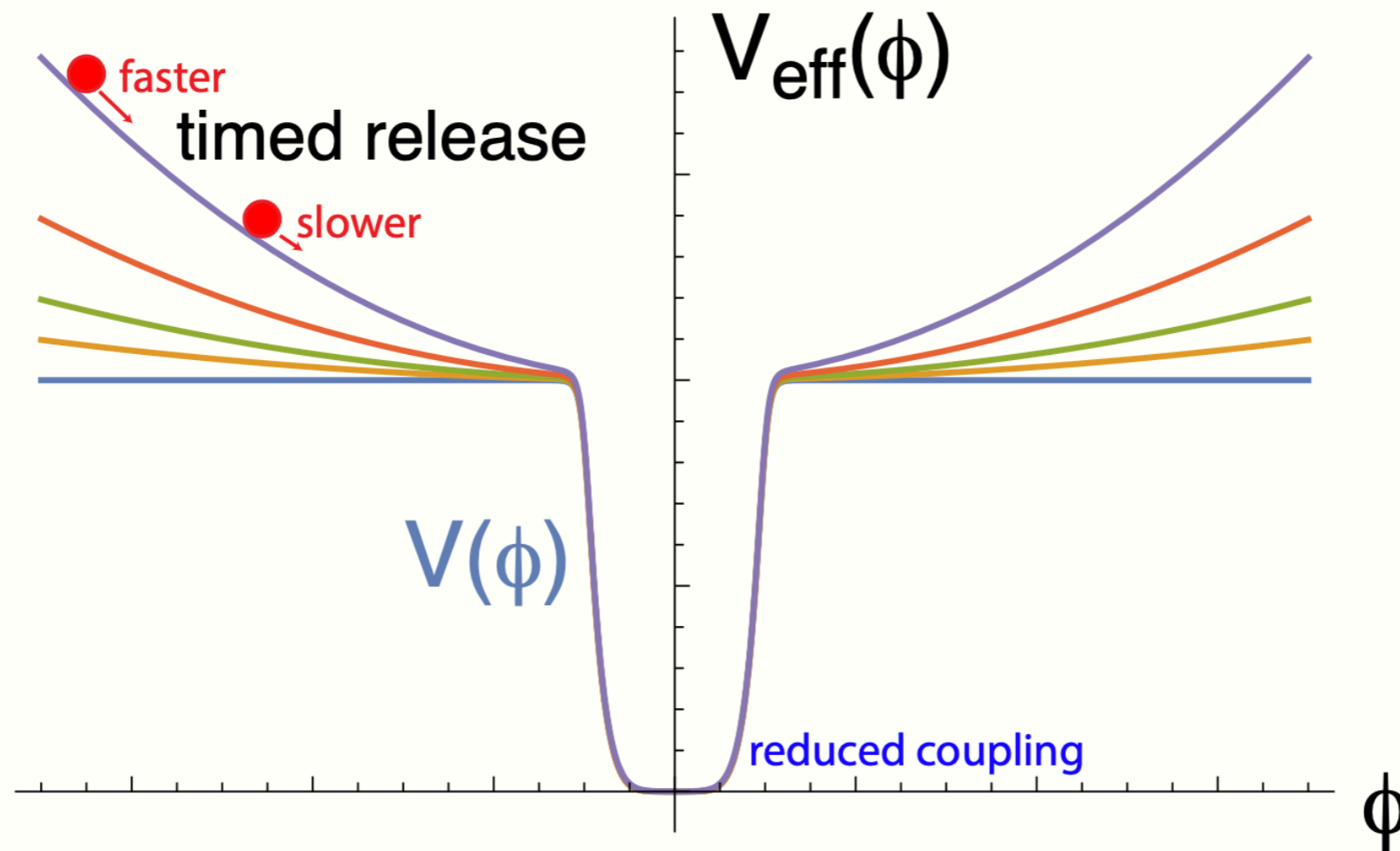


# Early Dark Sector 2

Colin Hill  
Columbia/CCA

## Toy solution

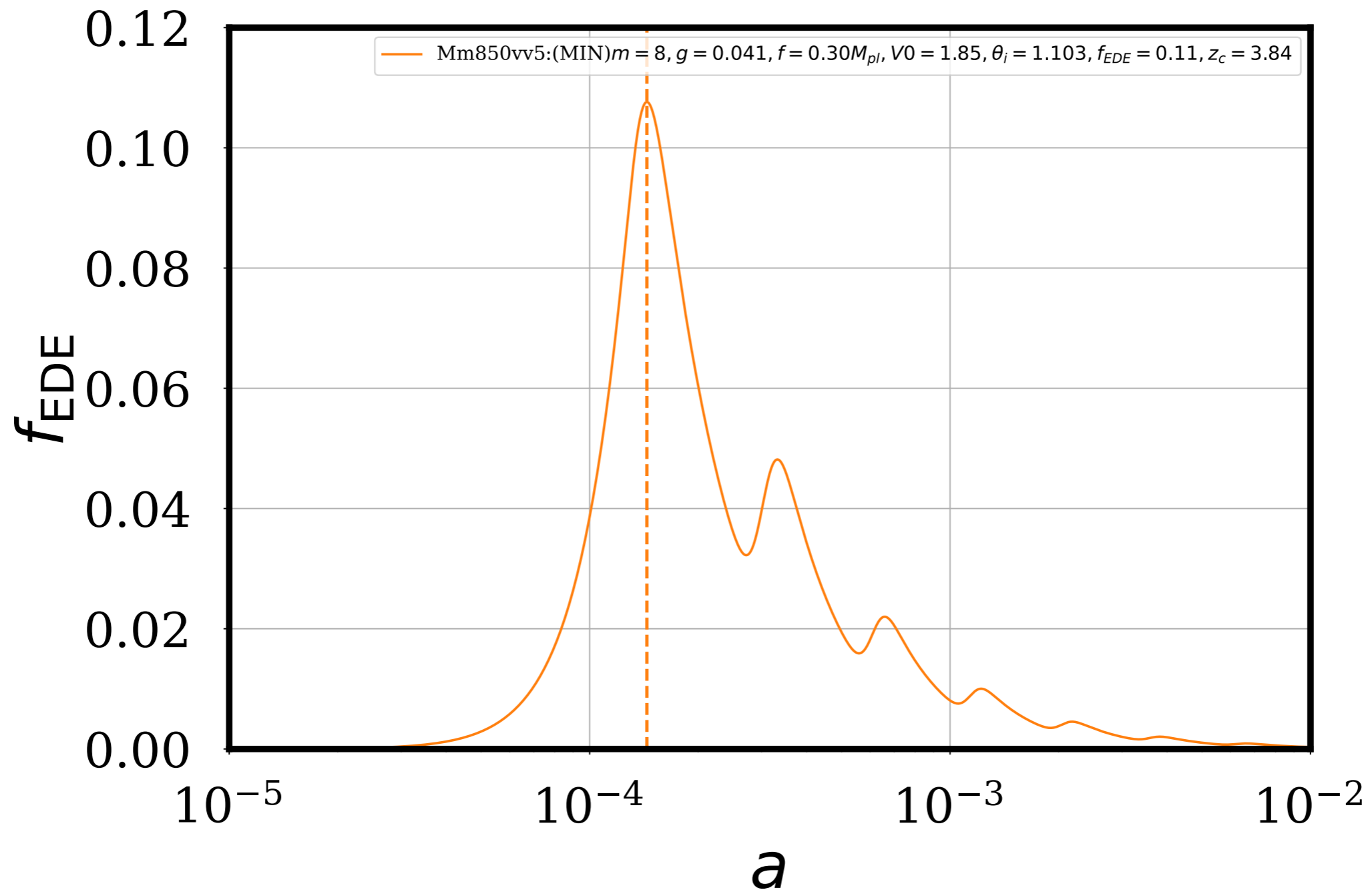
- Coincidence solved: field starts to roll because of equality
- Initial tuning solved: field will roll to edge of plateau from wide range of initial field positions
- Late growth solved:  $m(\phi) \propto 1 + g\phi^2$  suppresses 5th force  $\phi \rightarrow 0$



# Early Dark Sector 2

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Basic validation: can successfully lower  $r_s$ , raise  $H_0 \sim 70.5$  km/s/Mpc

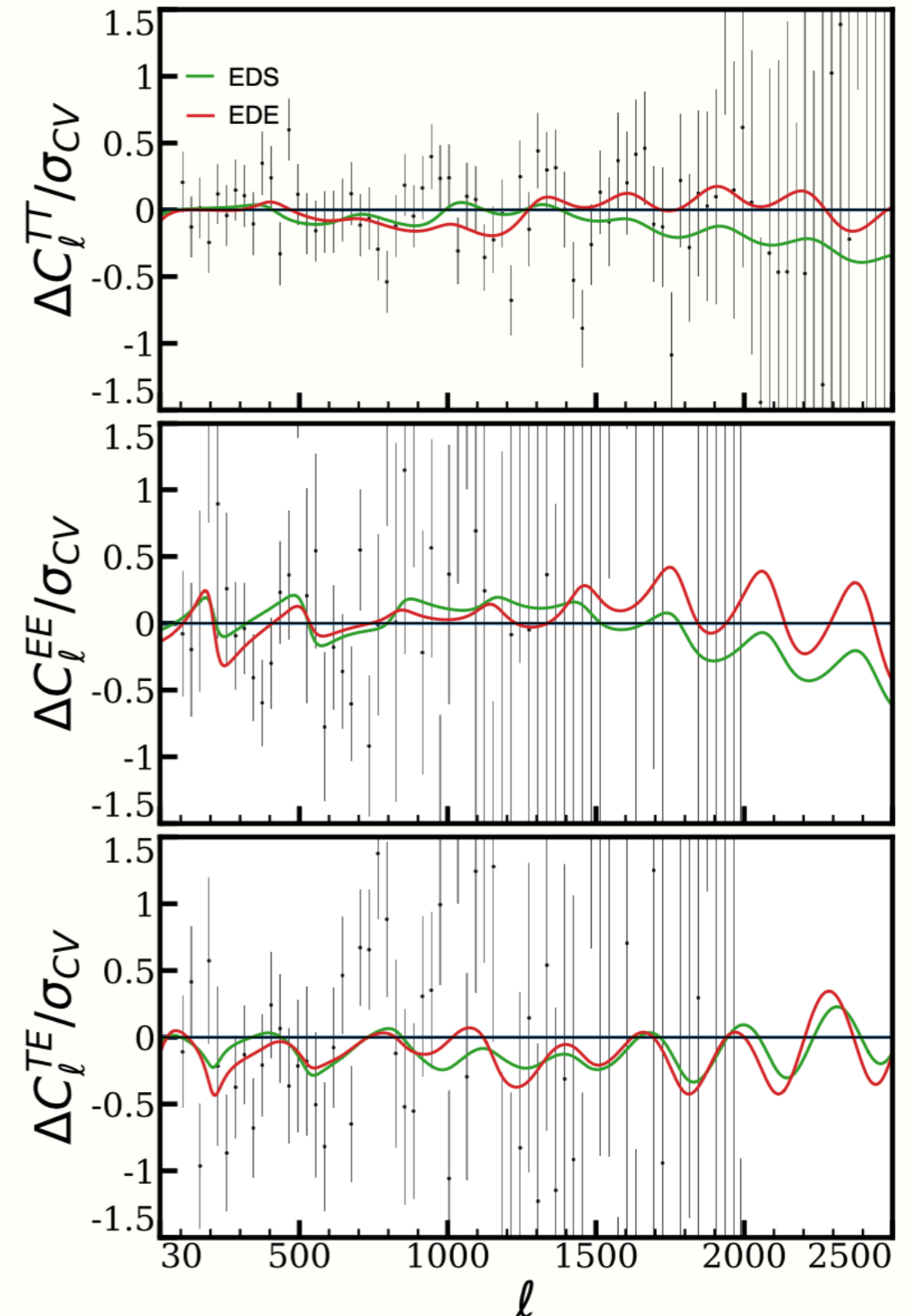


# Early Dark Sector 2

Colin Hill  
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However:

- Simple toy model achieves  $H_0 > 70$  (better fit than  $\Lambda$ CDM but not as good as EDE)

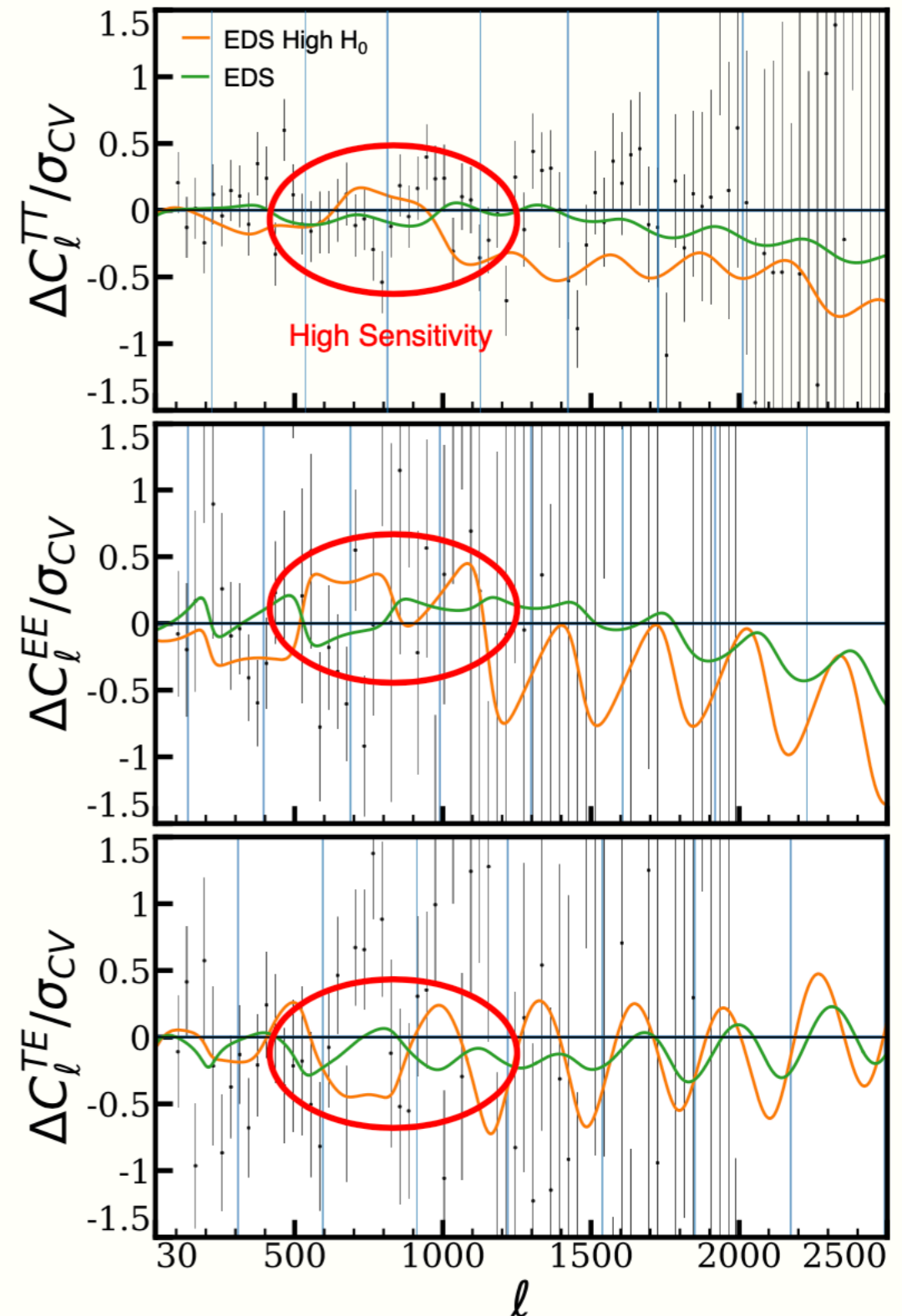


# Early Dark Sector 2

Colin Hill  
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However:

- Simple toy model achieves  $H_0 > 70$  (better fit than  $\Lambda$ CDM but not as good as EDE)
- Even larger  $H_0$  produces too large changes to gravitational driving for CMB acoustic modes around  $\ell \sim 500$  (see also [Lin, Hu, Raveri 2009.08974](#))
- Potentially related to  $A_L$  anomaly ( $TT$  peaks too smooth)
- Incidental or fundamental? rolling in the effective potential produces problematic field fluctuations



W.I.P.

Next: ACT DR6



# DR6 Forecasts

***ACT TT + TE + EE : precision cosmology beyond Planck***

	ACT DR4	ACT DR4 + WMAP	Planck	Planck + ACT DR6
$\sigma(H_0)$	1.5	1.1	0.5	0.4
$\sigma(n_s)$	0.015	0.006	0.004	0.003
$\sigma(N_{\text{eff}})$	0.4	0.3	0.2	0.1

**Large improvements in beyond- $\Lambda$ CDM parameters:  
~2x increase in sensitivity to new light relic particles**



***PRELIMINARY FORECAST***

**Upcoming ACT DR6 precision cosmology constraints will surpass those from Planck ( $H_0$ ,  $N_{\text{eff}}$ ,  $\Sigma m_\nu$ ,  $\sigma_8$ , + beyond- $\Lambda$ CDM models) — stay tuned!**

# Discovering EDE in the CMB

Colin Hill  
Columbia/CCA

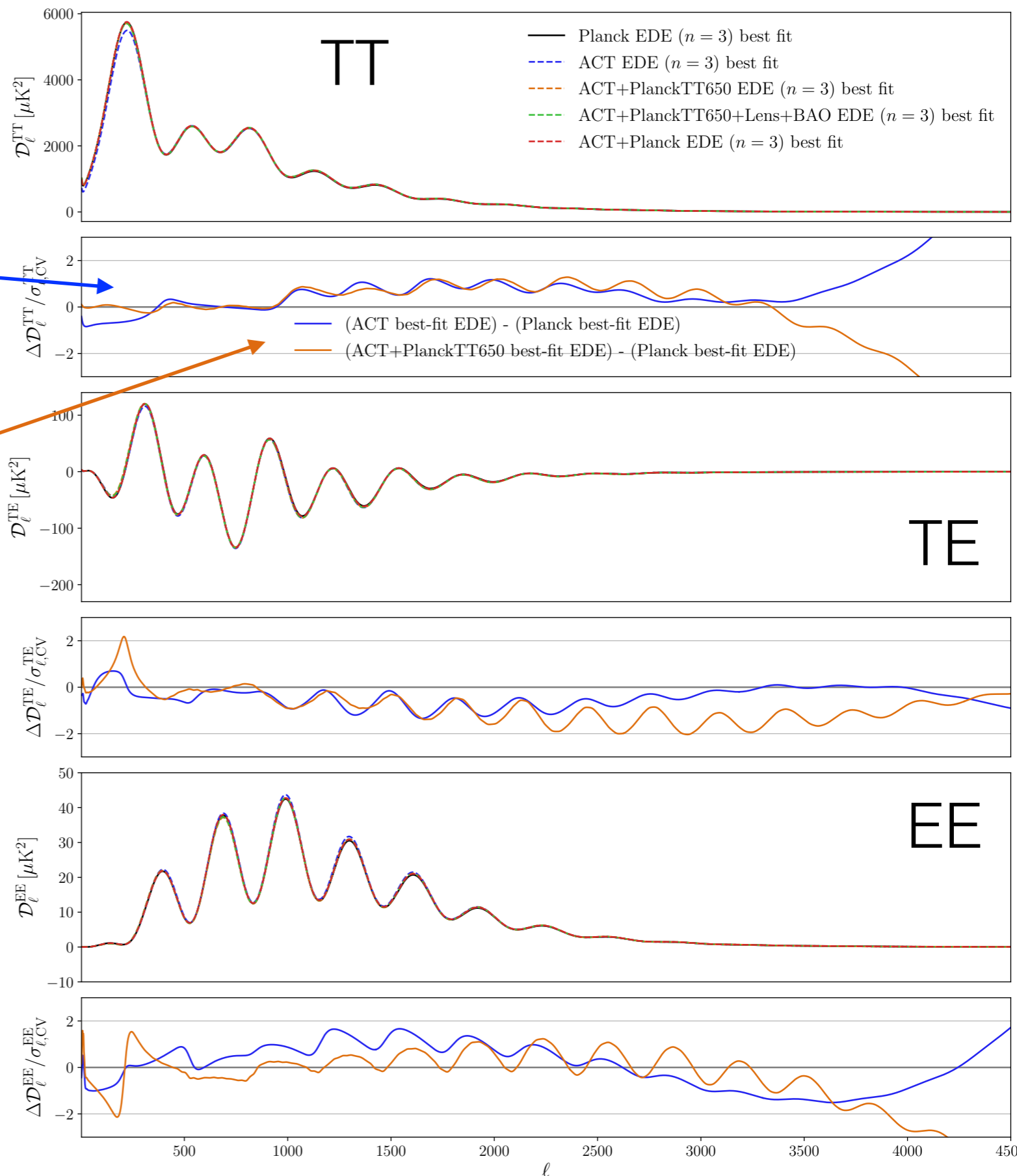
# Discovering EDE in the CMB

Colin Hill  
Columbia/CCA

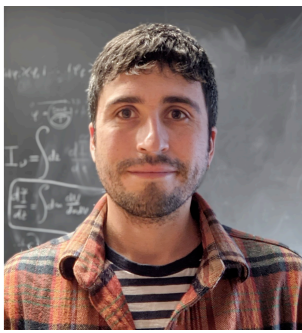
ACT best-fit EDE -  
Planck EDE

ACT+P18TT650 EDE -  
Planck EDE

Imminent potential  
discovery with upcoming  
ACT DR6 (~2023): the  
models shown  
here can be  
distinguished at  $\sim 20\sigma$



# Post-Recombination Reheating (PRR)



# $H_0$ and $T_{\text{CMB},0}$

*$H_0$  tension or  $T_0$  tension?*

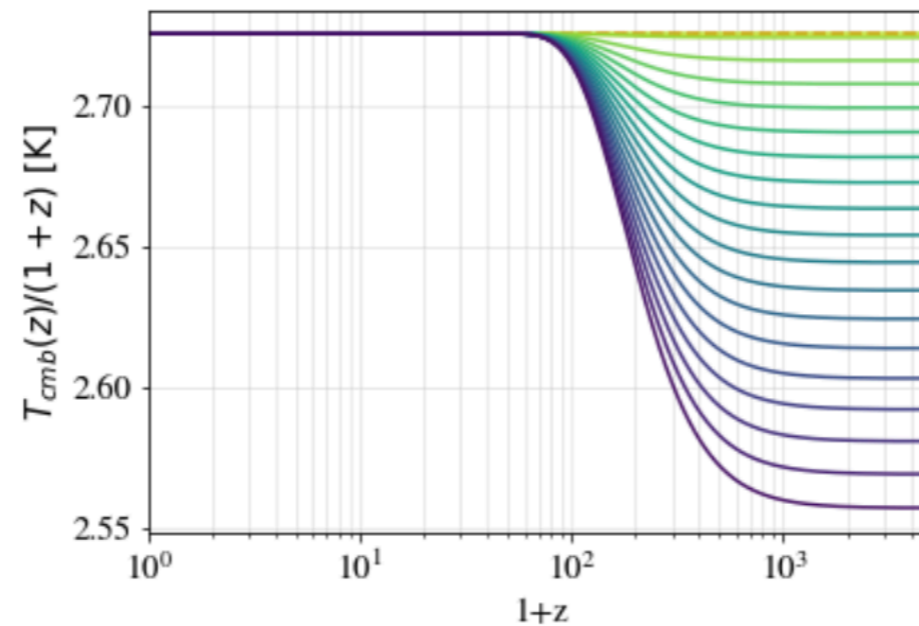
Geometric degeneracy: background and linear perturbation evolution depend (almost) only on parameter combination  $H_0 T_{\text{CMB},0}^{1.2}$

- Cosmology would be very different without COBE-FIRAS!
- Ignoring FIRAS, Planck+SH0ES can be fit with  $T_{\text{CMB},0} = 2.56 \pm 0.05$  K
- BAO breaks degeneracies: Planck + BAO yield  $T_{\text{CMB},0} = 2.71 \pm 0.02$  K
- FIRAS result:  $T_{\text{CMB},0} = 2.72548 \pm 0.00057$  K
- Can we build on this idea while maintaining agreement with FIRAS?



# Did the Universe (Slightly) Reheat after Recombination?

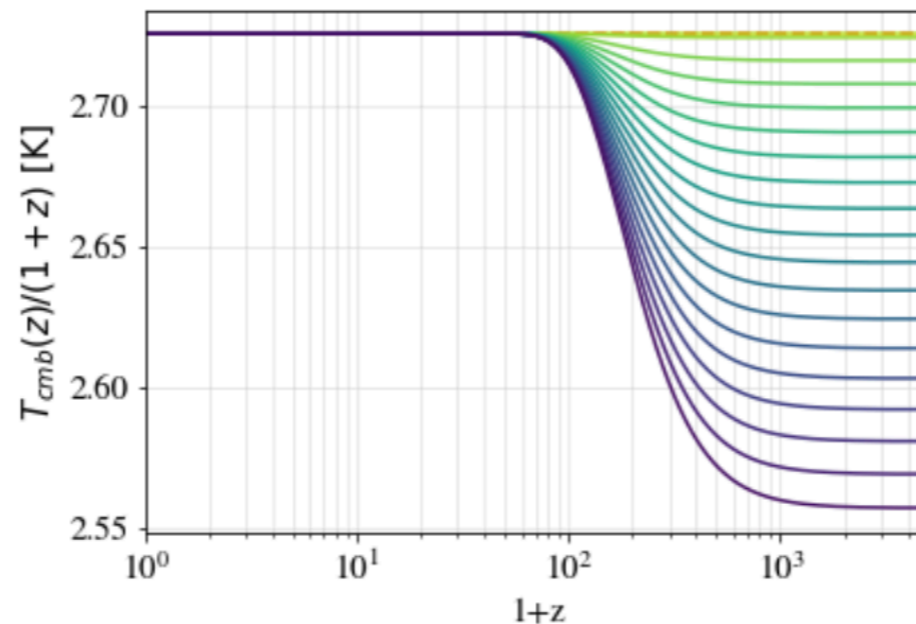
- Suppose  $T_{\text{CMB}}(z^*) < T_{\text{CMB}}(z^*)^{\Lambda\text{CDM}}$ , but a process injects energy at  $z < z^*$



Direct constraints  
on  $T_{\text{CMB}}(z)$  only  
exist at  $z < \sim 3$

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Direct constraints  
on  $T_{\text{CMB}}(z)$  only  
exist at  $z < \sim 3$

- A conspiracy of integrals leads to higher  $H_0$ : the sound horizon is not decreased (in fact it increases due to lower  $H(z)$  at early times)

$$r_s = \int_{z_*}^{\infty} \frac{dz}{H(z)} c_s(z)$$

- To keep  $\theta_s^*$  fixed,  $D_A^*$  must increase, but since  $z^*$  increases, one must increase  $H_0$  to compensate otherwise increased value of the integrand:

$$D_A^* = \int_0^{z^*} c dz / H(z)$$

- To keep  $k_{\text{eq}}$  fixed,  $\Omega_m$  must decrease, and hence  $S_8$  decreases

# Did the Universe (Slightly) Reheat after Recombination?

Concrete model: sub-component of CDM decays into photons after  $z^*$

- Background evolution similar to usual decaying DM- $\rightarrow$ DR:

$$\begin{aligned}\rho'_{\text{DCDM}} &= -3aH\rho_{\text{DCDM}} - a\Gamma\rho_{\text{DCDM}} \\ \rho'_{\gamma} &= -4aH\rho_{\gamma} + a\Gamma\rho_{\text{DCDM}},\end{aligned}$$

- New parameters:

$\Gamma$  = decay rate

$\omega_{\text{DCDM,ini}}$  = initial decaying CDM density

( $T_{\text{CMB,ini}}$  = initial CMB monopole temperature) — not really new

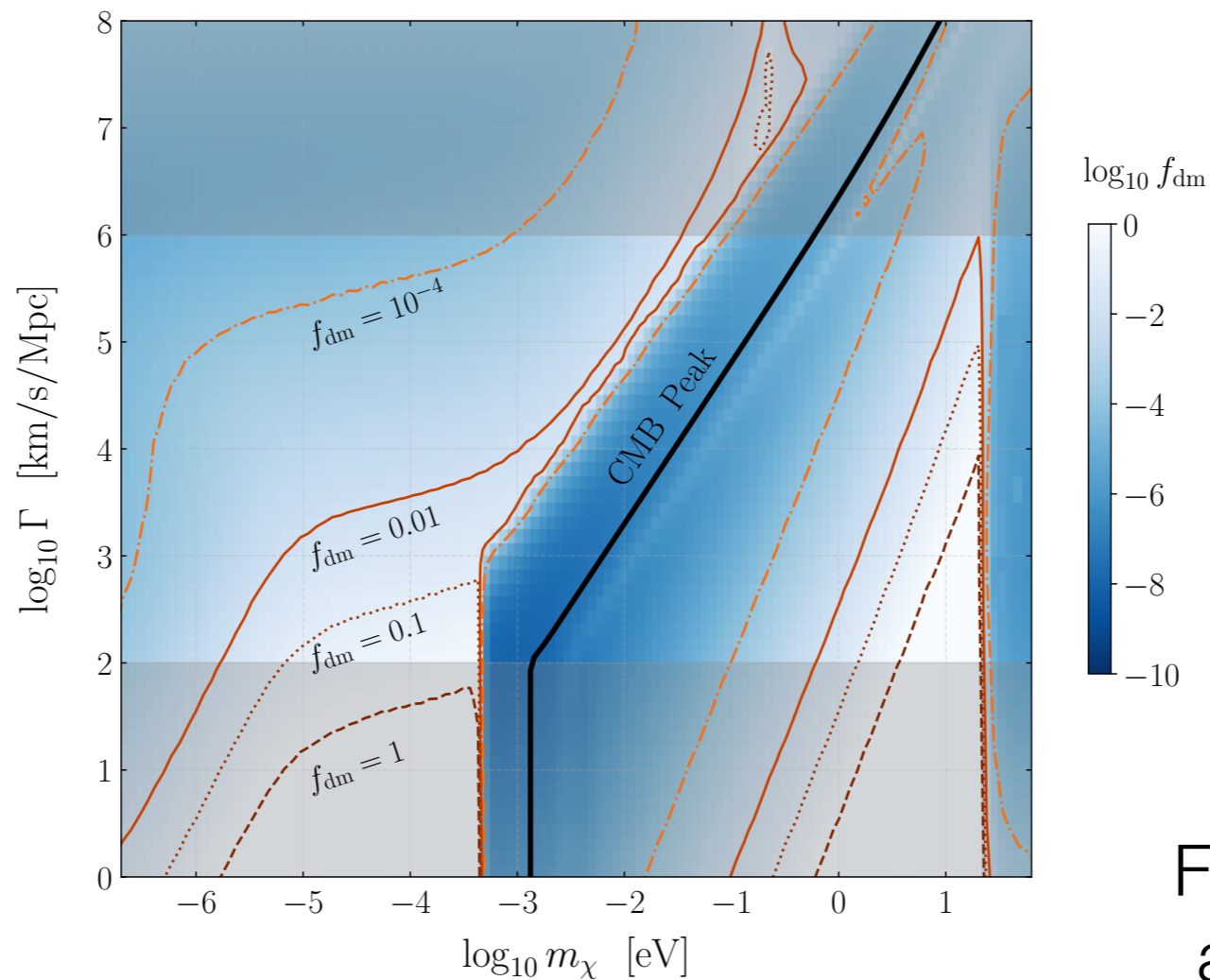
- Perturbation evolution equations for photons acquire new terms not present in  $\Lambda$ CDM or usual DCDM- $\rightarrow$ DR model

# Did the Universe (Slightly) Reheat after Recombination?

- Key feature: only a tiny amount of decaying CDM is needed to increase  $T_{\text{CMB}}$  by the necessary magnitude (e.g., for decay at  $z=22$ , only 0.02% of CDM decaying will increase  $T_{\text{CMB}}$  by 1%)

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- What about spectral distortion constraints?

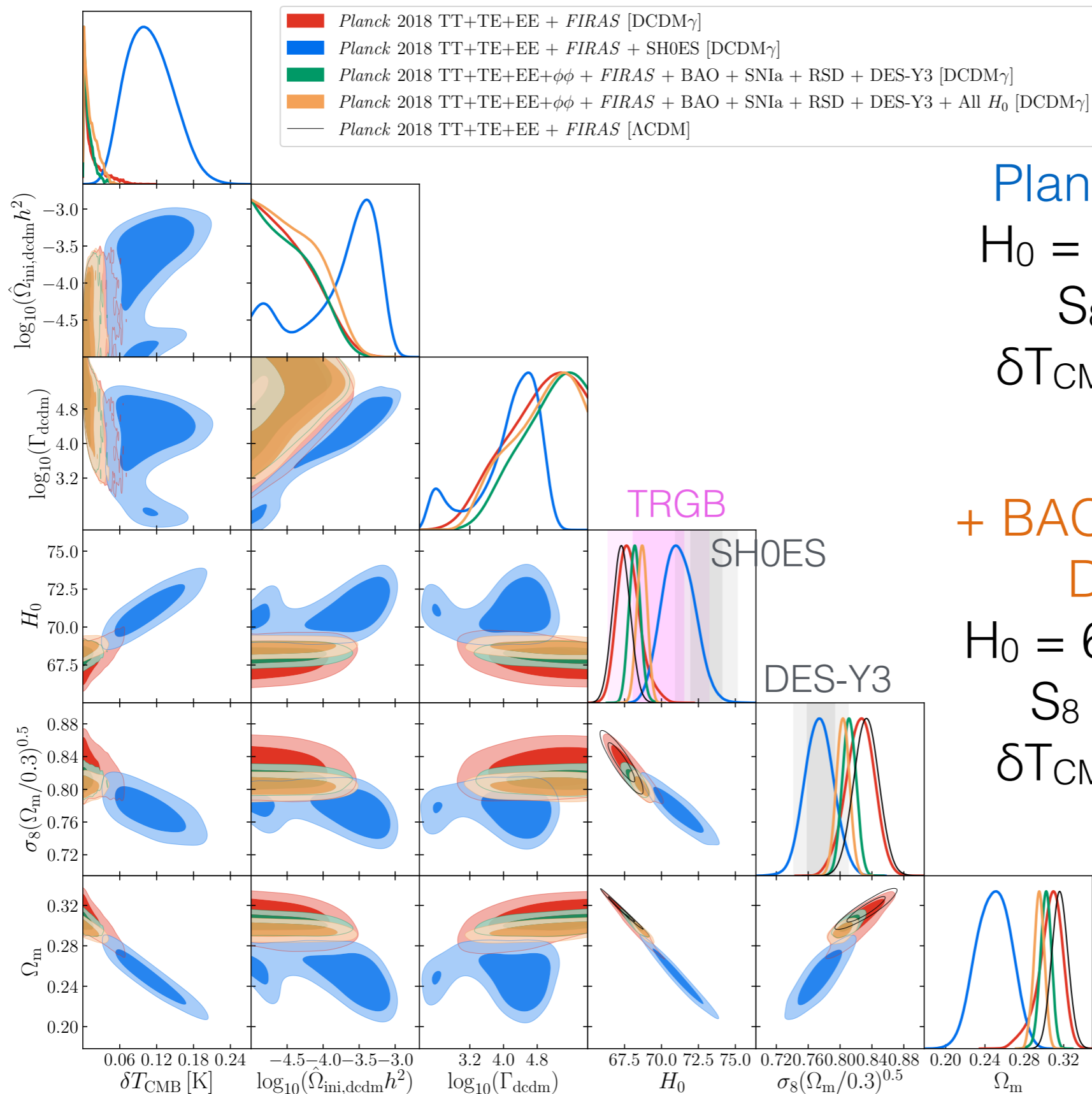


$f_{\text{DM}}$  = fraction  
of DM that  
decays

For simplicity here, we  
assume that injected  
photons are thermal



# PRR Analysis



Planck + FIRAS + SH0ES:

$$H_0 = 71.2 \pm 1.1 \text{ km/s/Mpc}$$

$$S_8 = 0.774 \pm 0.018$$

$$\delta T_{\text{CMB}} = 0.109^{+0.033}_{-0.044} \text{ K}$$

+ BAO + SNIa +  $\phi\phi$  + RSD +  
DES + MCP + SBF:

$$H_0 = 68.7 \pm 0.35 \text{ km/s/Mpc}$$

$$S_8 = 0.8035 \pm 0.0081$$

$$\delta T_{\text{CMB}} = 0.013^{+0.004}_{-0.013} \text{ K}$$

$\chi^2_{\text{Planck}}$  is as  
good as that of  
 $\Lambda$ CDM fit to  
Planck alone

# PRR Analysis

## Takeaways

- Main obstruction to success of the model: conflict with  $\Omega_m$  constraint from BAO and SNIa
- One tweak: allow  $N_{\text{eff}}$  to 'restore' early-universe radiation density back to its normal value, and mitigate decrease in  $\Omega_m$  (thus giving up  $S_8$  fix)
- Such a model fits data better than  $\Lambda\text{CDM}+N_{\text{eff}}$ , but not dramatically
- Key points:
  - There are large swathes of cosmic history where (semi-)dramatic changes to the model could still lurk
  - Seemingly small changes ( $\delta T_{\text{CMB}} \sim 10\text{-}100$  mK) can have big effects
  - We need to measure CMB spectral distortions much better than FIRAS! Strong motivation for PIXIE, FOSSIL, BISOU +++



# Take-Home Messages

- 1) ACT and Planck prefer somewhat different EDE model parameters, with ACT yielding higher  $H_0$
- 2) Early dark sector may help w/  $S_8$ , coincidence, ICs of EDE
- 3) Small post-recombination reheating is allowed by data and moves  $H_0$  and  $S_8$  in the 'right direction'
- 4) Early(ish)-universe  $H_0 / S_8$  resolutions generically predict clear deviations from  $\Lambda$ CDM in the CMB — imminently testable with ACT DR6



# Bonus

# ACT DR4 Cosmology

Colin Hill

Columbia/CCA

ACT: completely independent check of WMAP and Planck results

## $\Lambda$ CDM Parameter Refresher:

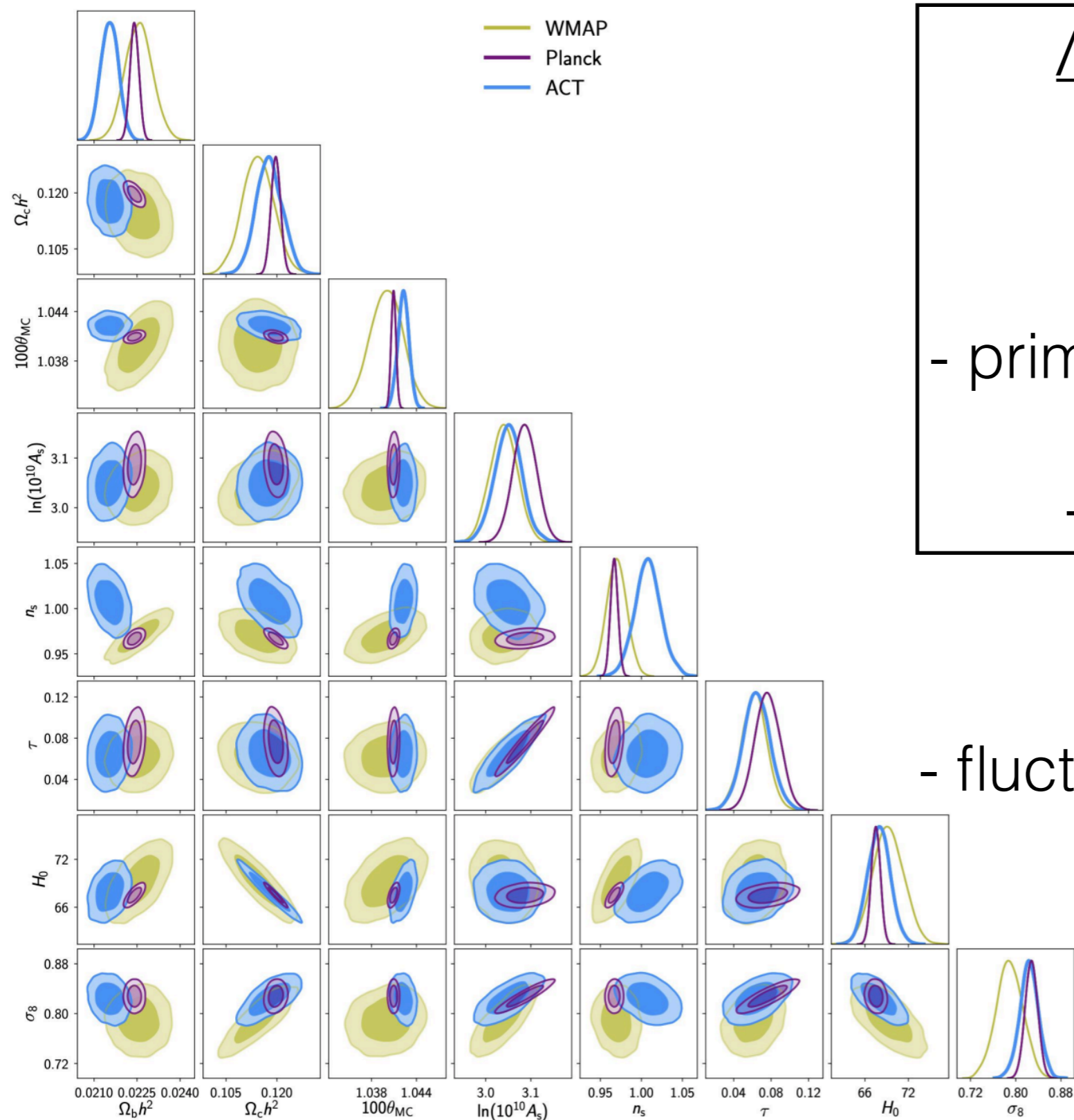
- baryon density
- cold dark matter density
- angular acoustic scale
- primordial fluctuation amplitude
  - primordial spectral index
  - optical depth (reionization)

## Derived Parameters:

- Hubble constant ( $H_0$ )
- fluctuation amplitude at  $z=0$  ( $\sigma_8$ )

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ACT: completely independent check of WMAP and Planck results



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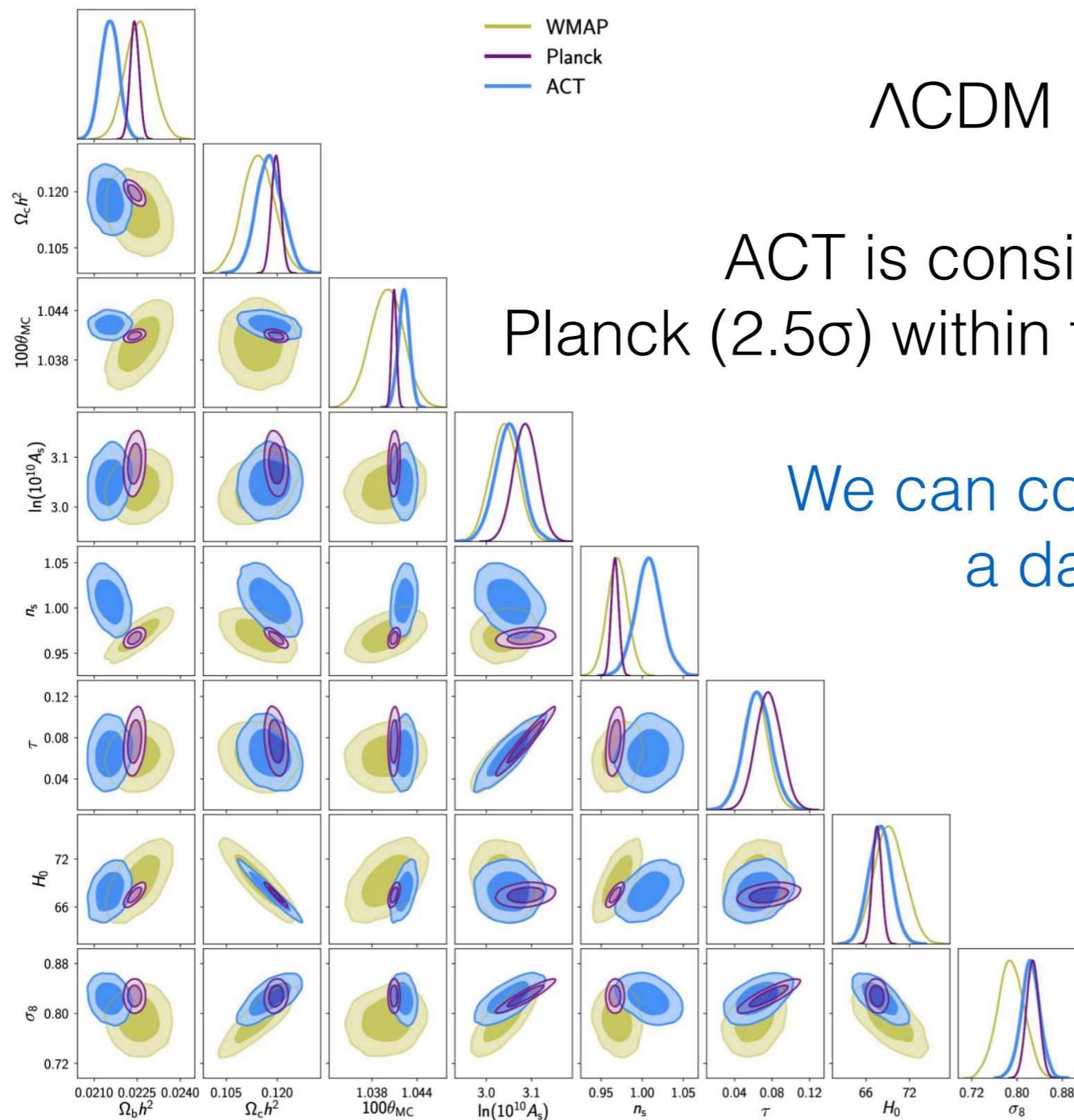
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ACT: completely independent check of WMAP and Planck results



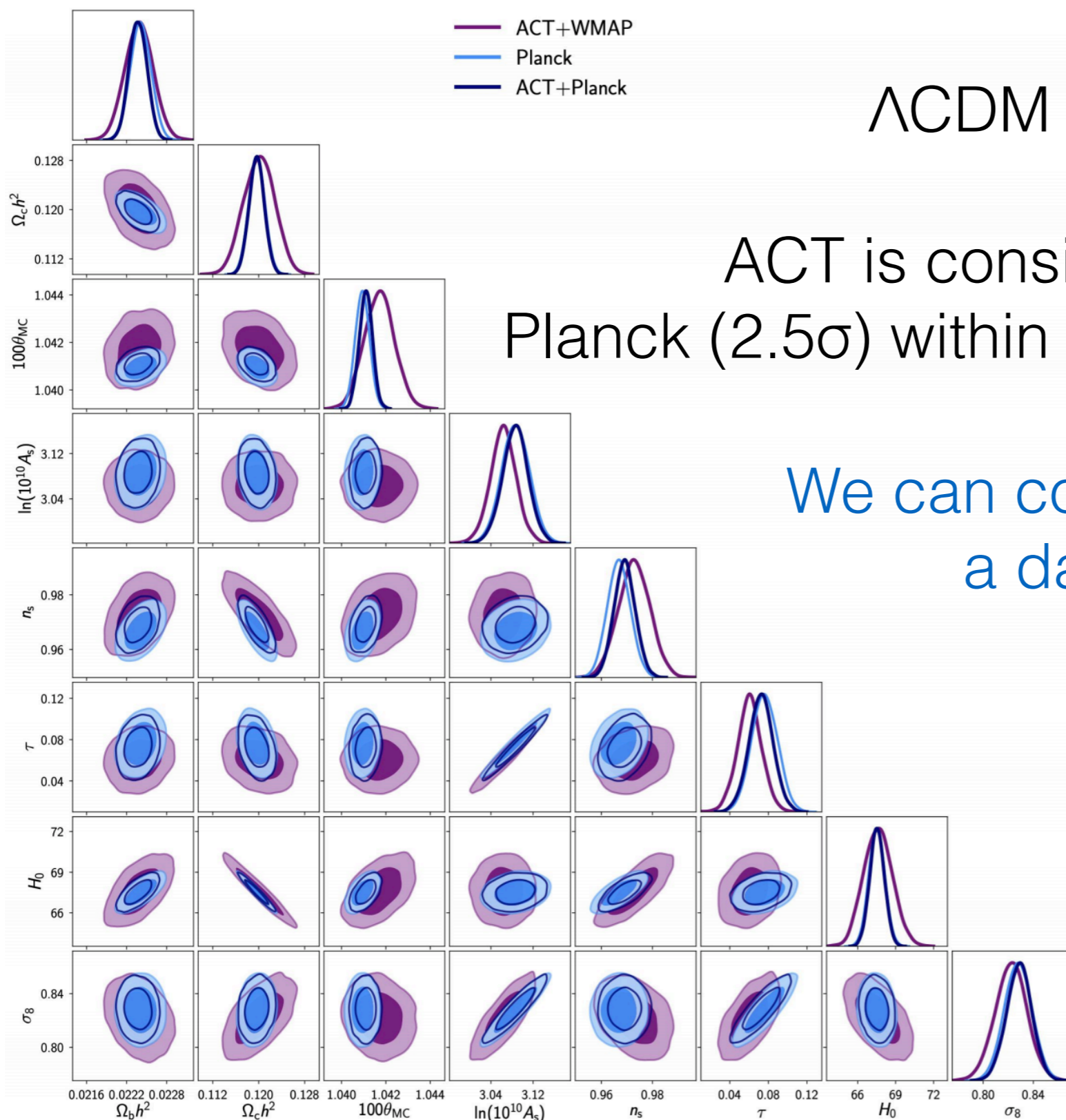
$\Lambda$ CDM is a good fit to the ACT data

ACT is consistent with WMAP ( $2.4\sigma$ ) and Planck ( $2.5\sigma$ ) within the  $\Lambda$ CDM parameter space

We can combine ACT+WMAP to form a data set with statistical power approaching Planck

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ACT(+WMAP): completely independent check of Planck results



$\Lambda$ CDM is a good fit to the ACT data

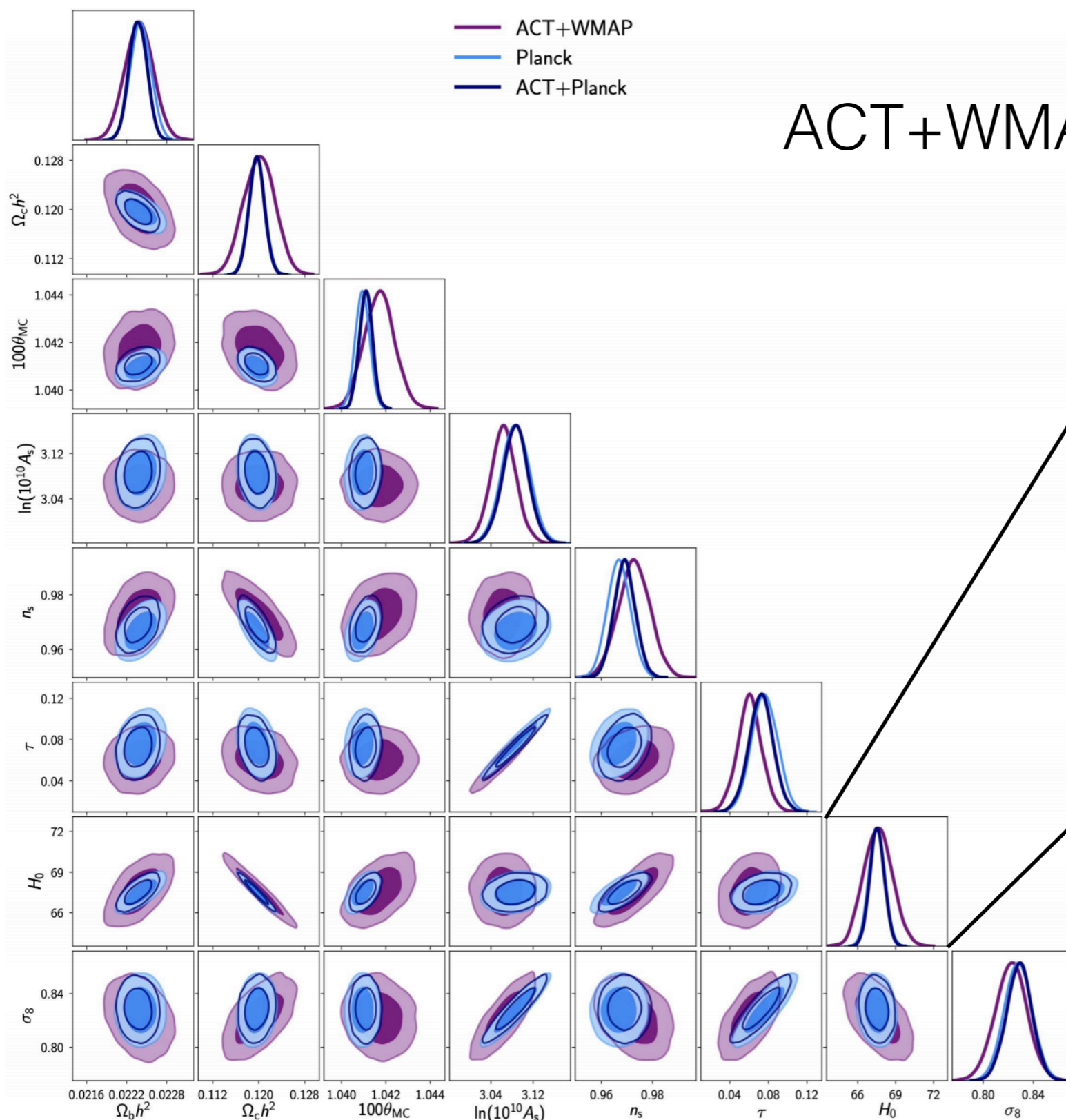
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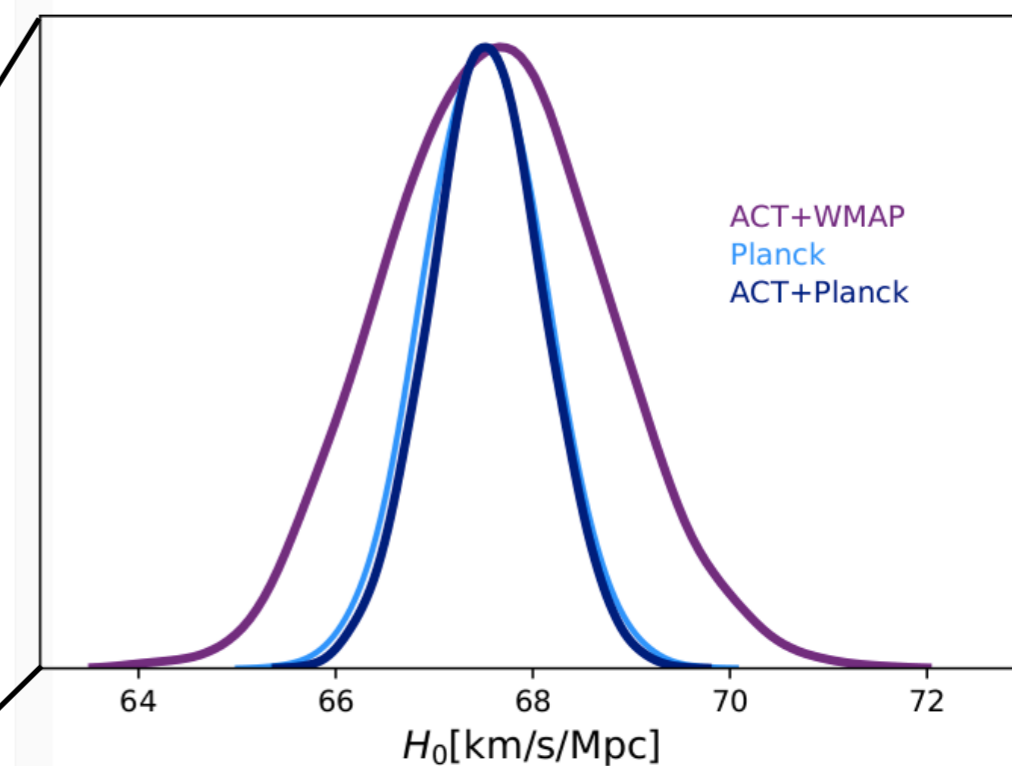
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ACT(+WMAP): completely independent check of Planck results



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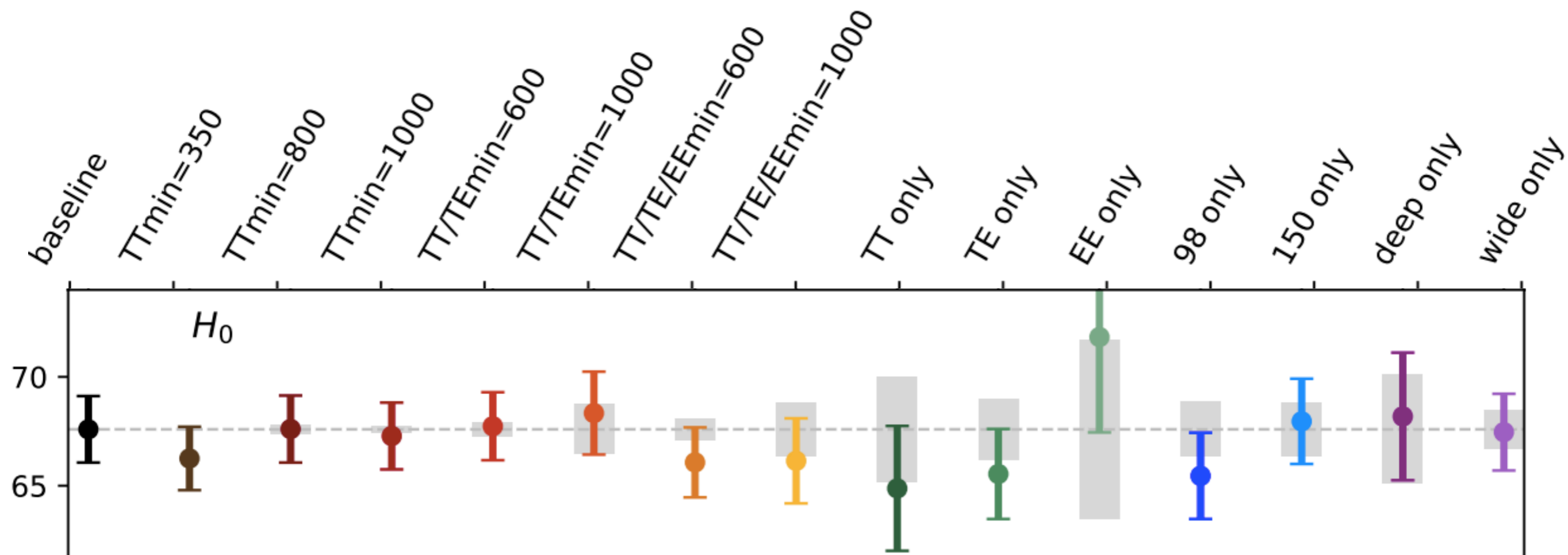


$H_0 = 67.6 \pm 1.1$  km/s/Mpc ACT+WMAP  
 $H_0 = 67.9 \pm 1.5$  km/s/Mpc ACT  
 $H_0 = 67.5 \pm 0.6$  km/s/Mpc Planck

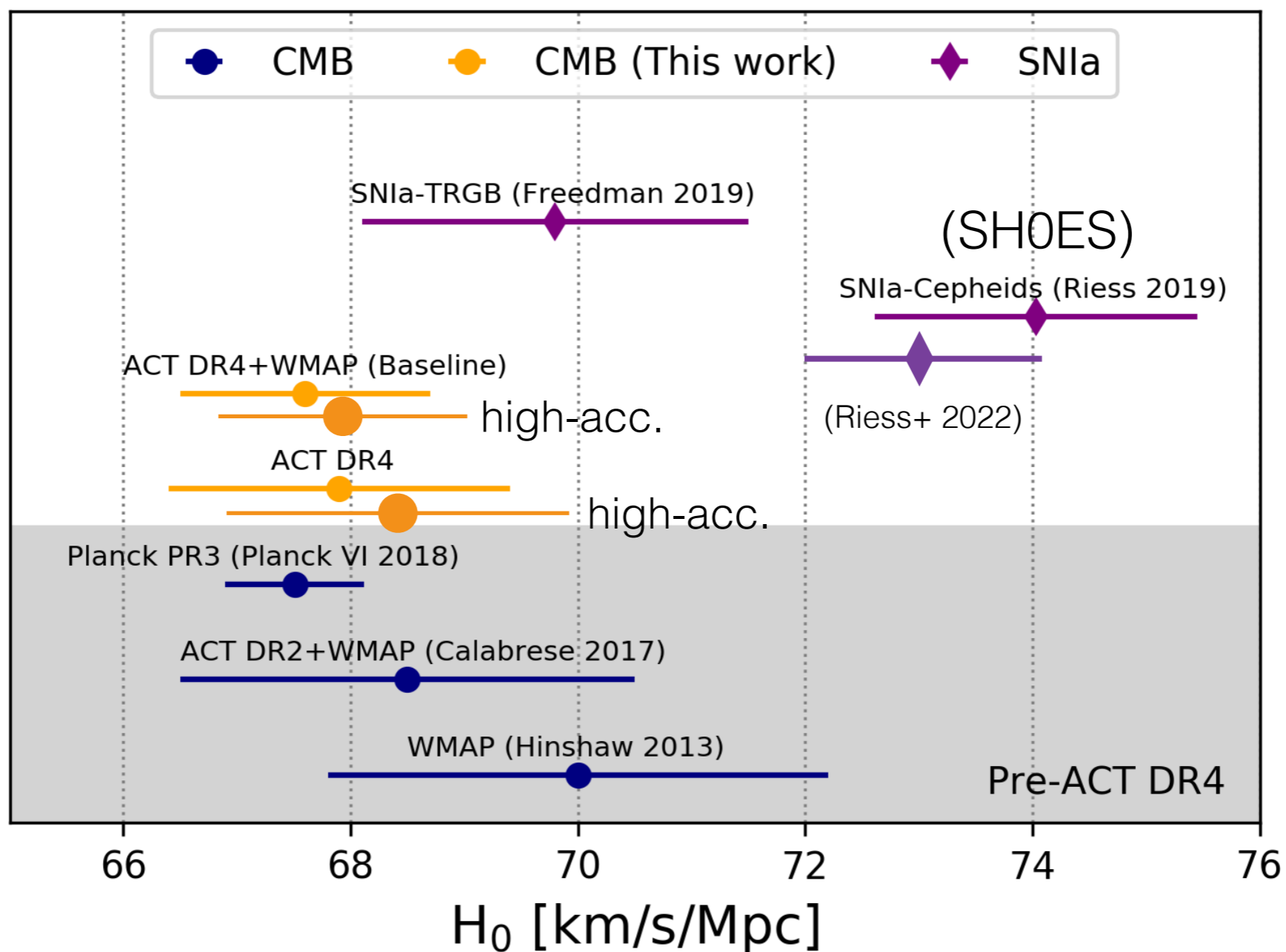


# ACT DR4 Cosmology

$H_0$  results are stable for a wide range of analysis choices



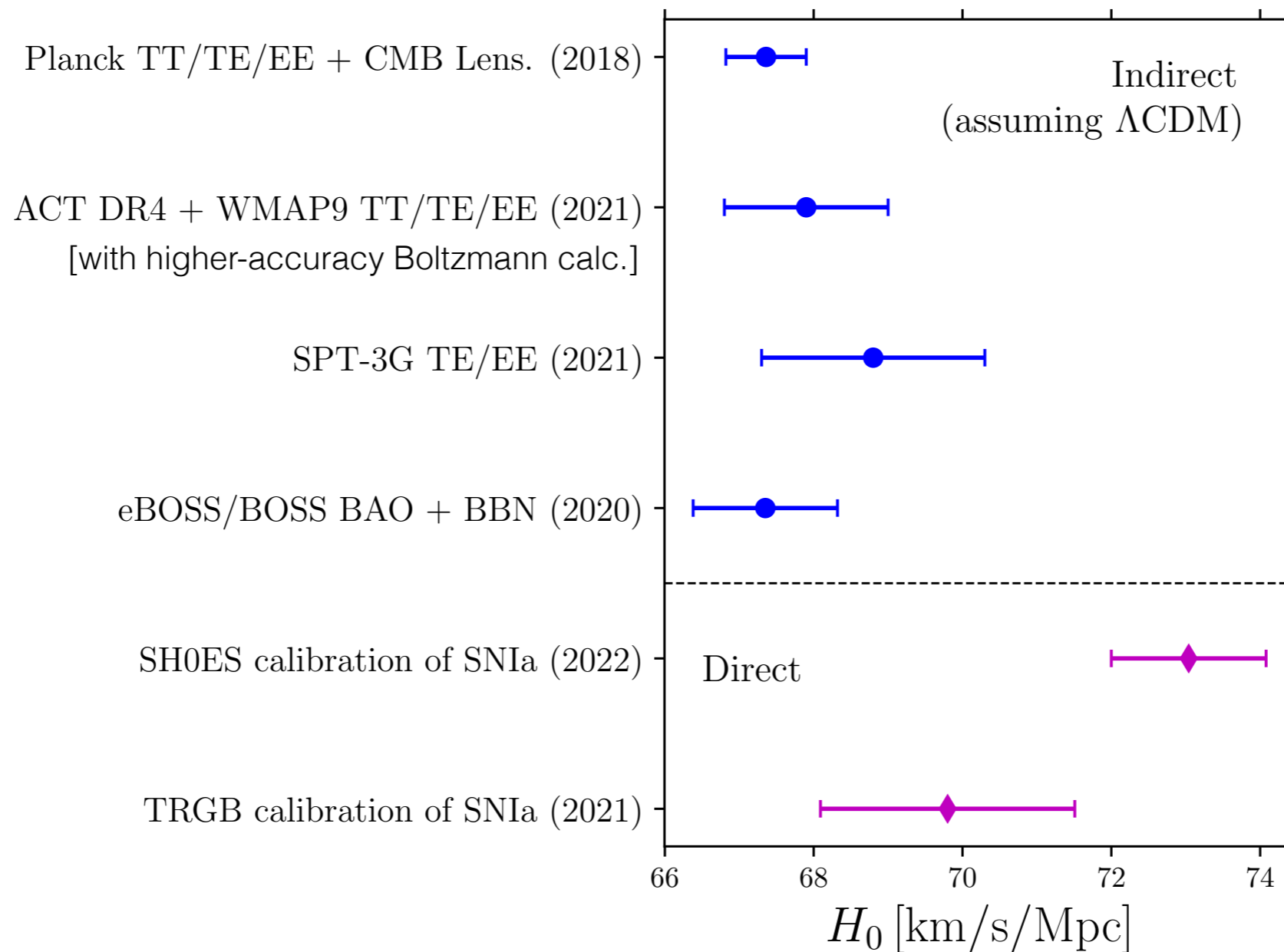
# ACT DR4 Cosmology



~3.4 $\sigma$  difference between ACT+WMAP (high-acc.,  $\Lambda$ CDM) and Cepheid-calibrated SNIa (SH0ES 2022)

Agreement within  $\sim 1\sigma$  between ACT+WMAP and TRGB-calibrated SNIa

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Agreement within  $\sim 1\sigma$  between ACT+WMAP and TRGB-calibrated SNIa

[Aiola et al. \(2020\)](#); [JCH et al. \(2021\)](#); see also [McCarthy, JCH, Madhavacheril \(2021\)](#)



# Did the Universe (Slightly) Reheat after Recombination?

## Linear perturbation theory

$$\delta'_{\text{DCDM}} = -\theta_{\text{DCDM}} - m_{\text{cont}} - a\Gamma m_{\psi} \quad (3)$$

$$\theta'_{\text{DCDM}} = -\frac{a'}{a}\theta_{\text{DCDM}} + k^2 m_{\psi} \quad (4)$$

$$\delta'_{\gamma} = -\frac{4}{3}\theta_{\gamma} - \frac{4}{3}m_{\text{cont}} + a\Gamma \frac{\rho_{\text{DCDM}}}{\rho_{\gamma}} (\delta_{\text{DCDM}} - \delta_{\gamma} + m_{\psi}) \quad (5)$$

$$\theta'_{\gamma} = k^2 \left( \frac{1}{4}\delta_{\gamma} - \sigma_{\gamma} \right) + k^2 m_{\psi} + an_e \sigma_{\text{T}} (\theta_b - \theta_{\gamma}) - \frac{3}{4}a\Gamma \frac{\rho_{\text{DCDM}}}{\rho_{\gamma}} \left( \frac{4}{3}\theta_{\gamma} - \theta_{\text{DCDM}} \right) \quad (6)$$

$$F'_{\gamma,2} = 2\sigma'_{\gamma} = \frac{8}{15}\theta_{\gamma} - \frac{3k}{5}F_{\gamma,3} + \frac{8}{15}m_{\text{shear}} - \frac{9}{5}an_e \sigma_{\text{T}} \sigma_{\gamma} + \frac{1}{10}an_e \sigma_{\text{T}} (G_{\gamma,0} + G_{\gamma,2}) - 2\sigma_{\gamma} a\Gamma \frac{\rho_{\text{DCDM}}}{\rho_{\gamma}} \quad (7)$$

$$F'_{\gamma,\ell} = \frac{k}{2\ell+1} [\ell F_{\gamma,\ell-1} - (\ell+1) F_{\gamma,\ell+1}] - a\Gamma F_{\gamma,\ell} \frac{\rho_{\text{DCDM}}}{\rho_{\gamma}} \quad (8)$$

All terms containing  $\Gamma$  in photon perturbation equations are new

Gauge	Synchronous	Newtonian
$m_{\text{cont}}$	$h'/2$	$-3\phi'$
$m_{\psi}$	0	$\psi$
$m_{\text{shear}}$	$(h' + 6\eta')/2$	0

# Did the Universe (Slightly) Reheat after Recombination?

“Hat” variables: all cosmological quantities at a given  $T_{\text{CMB}}$  depend on quantities proportional to baryon-to-photon and dark matter-to-photon number ratios

$$\begin{aligned}\hat{\omega}_b &\equiv \omega_b \left( \frac{T_{\text{CMB,ini}}}{T_{\text{FIRAS}}} \right)^{-3} \\ \hat{\omega}_c &\equiv \omega_c \left( \frac{T_{\text{CMB,ini}}}{T_{\text{FIRAS}}} \right)^{-3} \\ \hat{\omega}_{\text{DCDM,ini}} &\equiv \omega_{\text{DCDM,ini}} \left( \frac{T_{\text{CMB,ini}}}{T_{\text{FIRAS}}} \right)^{-3} \\ \hat{A}_s &\equiv A_s \left( \frac{T_{\text{CMB,ini}}}{T_{\text{FIRAS}}} \right)^{n_s-1}.\end{aligned}$$

BBN abundances depend only on  $\hat{\omega}_b$

Consistency of this parameter with Planck thus maintains (approximate) consistency with BBN [modulo varying  $N_{\text{eff}}$ , etc.]

# Did the Universe (Slightly) Reheat after Recombination?

Constraints on light particles coupled to photons from, e.g., white dwarf lifetimes can be evaded by formulating the model using an excited dark matter state

Consider DM with non-zero dipole moment, coupled to SM sector through a kinetically mixed massive dark photon (DP)

The DP allows for transitions between ground and excited states of the DM

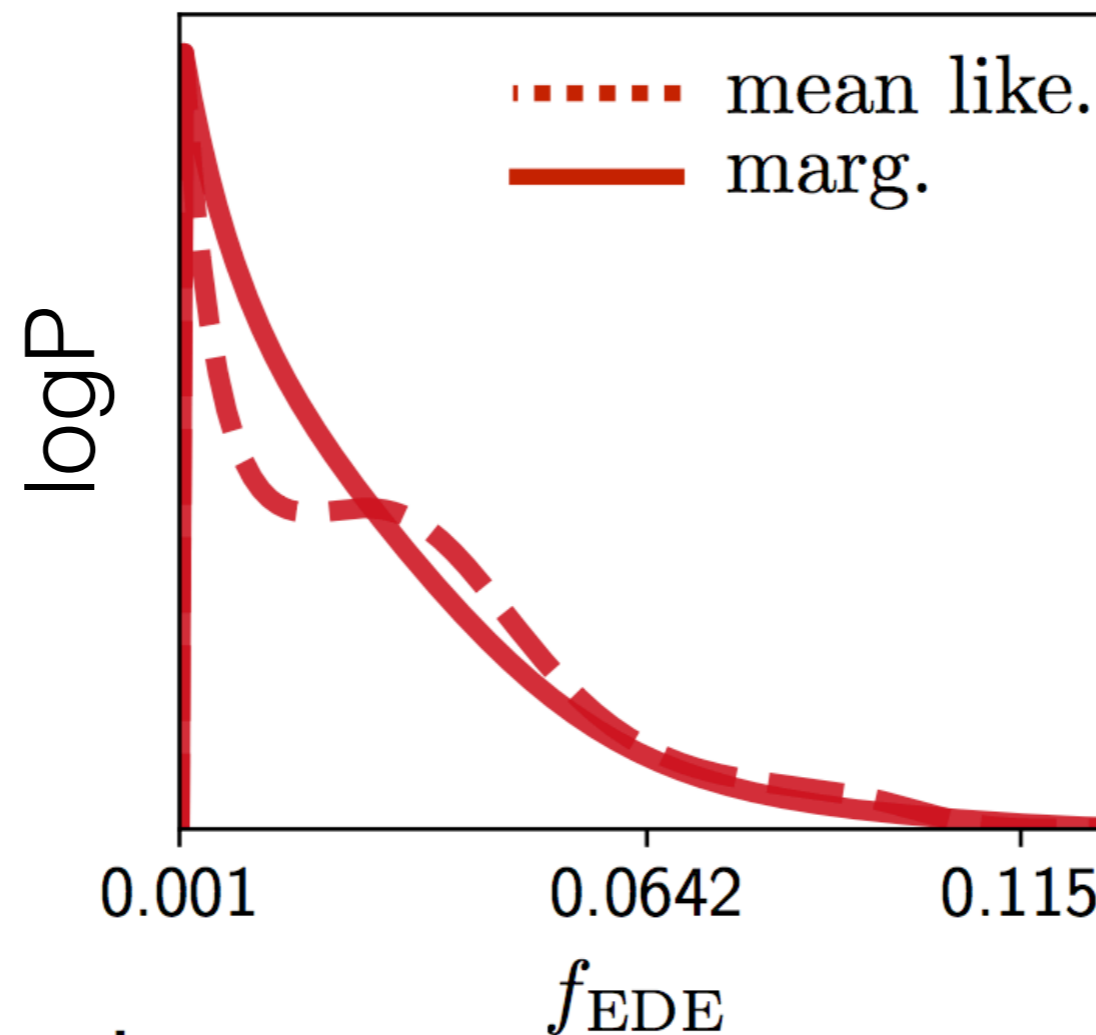
If the energy splitting of the states is  $\sim 0.1$  eV, the excited state is metastable for  $\sim 1-10$  Myr as we need in the PRR scenario

If  $m_{\text{DP}} > \text{MeV}$ , it is not produced in stars or supernovae

In general our results suggest that cosmological implications of such scenarios should be considered

# Prior Volume Effects?

non-preference for  $H_0$ -resolving EDE in Planck is robust in either frequentist or Bayesian methodology

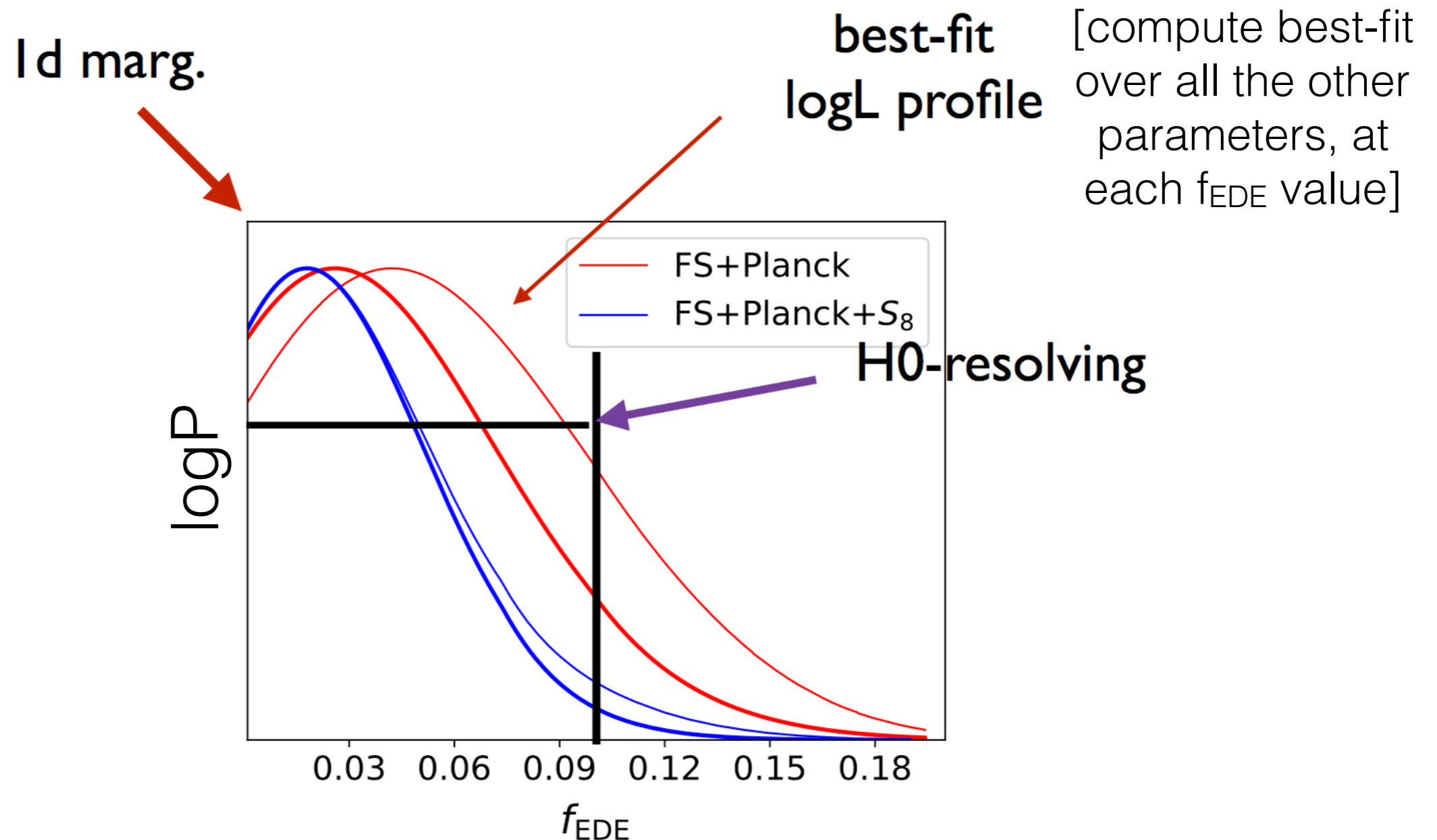


mean like. =  
[compute average  
log(likelihood),  
averaging over  
all the other  
parameters, at  
each  $f_{\text{EDE}}$  value]

**FS+Planck only**

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