## A reanalysis of the SHOES data for $H_0$ : Effects of new degrees of freedom on the Hubble tension

#### Leandros Perivolaropoulos

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Slides of talk available at: https://cosmology.physics.uoi.gr/seminars/



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#### **Main Points**



A reanalysis of the latest SH0ES data for  $H_0$ : Effects of new degrees of freedom on the Hubble tension Leandros Perivolaropoulos (Ioannina U.), Foteini Skara (Ioannina U.) (Aug 23, 2022) e-Print: 2208.11169 [astro-ph.CO]

- 1. New physics at ultra low redshifts (z<0.02) can lead to a resolution of the Hubble tension.
- 2. New degrees of freedom can be introduced in the baseline SHOES analysis to detect such new physics signatures.
- If a change of the SnIa absolute luminosity M<sub>B</sub> is allowed at D<sub>c</sub>~50Mpc (new degree of freedom) in the SHOES analysis then the best fit value of H<sub>0</sub> drops from 73±1 km/(sec Mpc) to 67±4 km/(sec Mpc).
- 4. In the presence of the inverse distance ladder input on the SnIa absolute magnitude  $M_B$ , the transition model uncertainties drop dramatically and the extended model is strongly preferred over the SHOES baseline model of universal value of  $M_B$ .



 $H_0$  measurement using sound horizon standard ruler (inverse distance ladder):

Assumptions: P18ACDM E(z), Standard expansion before  $z_{rec}$ 

$$\theta_{s} = \frac{r_{s}}{D_{A}(z)} = \frac{H_{0} r_{s}}{\int_{0}^{z} \frac{dz}{E(z)}} \qquad r_{s} = \int_{0}^{t_{d}} c_{s} dt/a = \int_{0}^{a_{d}} c_{s} \frac{da}{a^{2}H(a)}$$
$$H_{0}^{P18} = 67.36 \pm 0.54 \text{ km s}^{-1} \text{Mpc}^{-1}$$



H<sub>0</sub> measurement using sound horizon standard ruler (inverse distance ladder):

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On the use of the local prior on the absolute magnitude of Type Ia supernovae in cosmological inference

David Camarena (Espirito Santo U.), Valerio Marra (Espirito Santo U. and Trieste Observ. and Trieste U.) (Jan 21, 2021) Published in: *Mon.Not.Roy.Astron.Soc.* 504 (2021) 5164-5171 • e-Print: 2101.08641 [astro-ph.CO]

Rapid transition of Geff at zt≈0.01 as a possible solution of the Hubble and growth tensions

Valerio Marra (Espirito Santo U. and Trieste Observ. and SISSA, Trieste and INFN, Trieste), Leandros Perivolaropoulos (Ioannina U.) (Feb 11, 2021)

Published in: Phys.Rev.D 104 (2021) 2, L021303 • e-Print: 2102.06012 [astro-ph.CO]

$$\theta_{s} = \frac{r_{s}}{D_{A}(z)} = \frac{H_{0} r_{s}}{\int_{0}^{z} \frac{dz}{E(z)}} \qquad r_{s} = \int_{0}^{t_{d}} c_{s} dt/a = \int_{0}^{a_{d}} c_{s} \frac{da}{a^{2}H(a)}$$

$$\mathbf{Z}_{rec} \qquad H_{0}^{P18} = 67.36 \pm 0.54 \text{ km s}^{-1} \text{Mpc}^{-1}$$

$$\mathcal{M} = M + 5log \frac{c/H_{0}}{Mpc} + 25 \qquad \mathcal{M}_{z>0.01} = 23.80 \pm 0.01$$

$$M_{pc}$$
  
 $M_{pc}$   
 $M$  tension.  
 $M_{z>0.01}^{P18} = -19.401 \pm 0.027$   $<$   $M_B^{R21} = -19.25 \pm 0.03$ 

 $\mathcal{M} = \mathcal{M}$ 



H<sub>0</sub> measurement using sound horizon standard ruler (inverse distance ladder):

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$$M_{z>0.01}^{P18} = -19.401 \pm 0.027 \quad \checkmark \quad M_B^{R21} = -19.25 \pm 0.03$$

M depends on  $G_{eff}$ .

 $H_0^{R21} = 73.04 \pm 1.04$ 



H<sub>0</sub> measurement using sound horizon standard ruler (inverse distance ladder):

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Rapid transition of Geff at  $zt \approx 0.01$  as a possible solution of the Hubble and growth tensions

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$$\mathcal{M}_{z>0.01} = 23.80 \pm 0.01$$
  
$$\mathcal{M}_{B}^{R21} = -19.25 \pm 0.03$$
  
$$\mathcal{M}_{B}^{R21} = -19.25 \pm 0.03$$
  
$$\mathcal{M}_{z>0.01} = M_{z<0.01}^{R20}$$

Here 
$$\theta_s = \frac{r_s}{D_A(z)} = \frac{H_0 r_s}{\int_0^z \frac{dz}{E(z)}}$$
  $r_s = \int_0^{t_d} c_s dt/a = \int_0^{a_d} c_s \frac{da}{a^2 H(a)}$   
 $R = M + 5log \frac{c/H_0}{Mpc} + 25$   $\mathcal{M}_{z>0.01} = 23.80 \pm 0.01$ 

M tension.  $M_{z>0.01}^{P18} = -19.401 \pm 0.027$   $\leq M_B^{R21} = -19.25 \pm 0.03$ 

M depends on 
$$G_{eff}$$
.  
 $H_0$  Tension  
 $H_0^{P18} = 67.36 \pm 0.54 \text{ km s}^{-1} \text{Mpc}^{-1}$ 



H<sub>0</sub> measurement using sound horizon standard ruler (inverse distance ladder):

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On the use of the local prior on the absolute magnitude of Type Ia supernovae in cosmological inference

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$$\mathcal{M}_{z>0.01} = 23.80 \pm 0.01$$

$$\mathcal{M}_{B}^{R21} = -19.25 \pm 0.03$$

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$$\mathcal{M}_{Z>0.01} = M_{z<0.01}^{R20}$$

$$\mathcal{M}_{z>0.01} = G_{\text{eff}}(z > 0.01)$$

$$\theta_{s} = \frac{r_{s}}{D_{A}(z)} = \frac{H_{0} r_{s}}{\int_{0}^{z} \frac{dz}{E(z)}} \qquad r_{s} = \int_{0}^{t_{d}} c_{s} dt/a = \int_{0}^{a_{d}} c_{s} \frac{da}{a^{2}H(a)}$$

$$\mathbf{z}_{rec} \qquad H_{0}^{P18} = 67.36 \pm 0.54 \text{ km s}^{-1} \text{Mpc}^{-1}$$

$$\mathcal{M} = M + 5 \log \frac{c/H_{0}}{Mpc} + 25 \qquad \mathcal{M}_{z>0.01} = 23.80 \pm 0.01$$

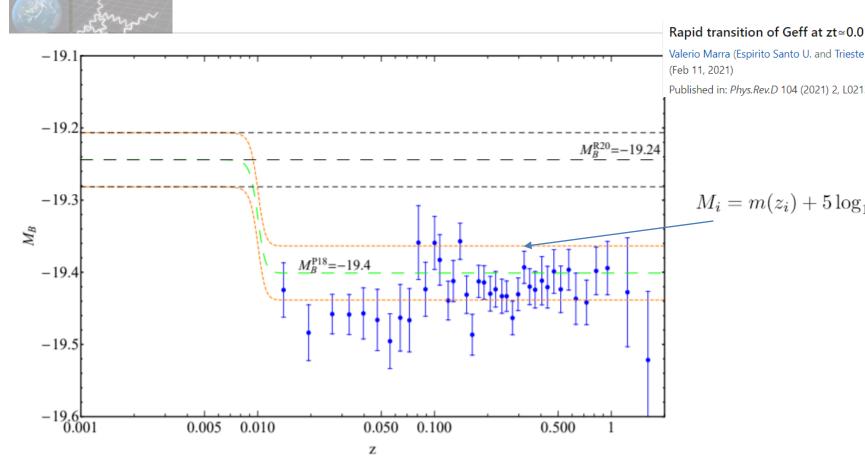
$$\mathbf{M} \text{ tension.}$$

$$M_{z>0.01}^{P18} = -19.401 \pm 0.027 \quad \checkmark \quad M_B^{R21} = -19.25 \pm 0.03$$

$$\begin{array}{c} \mbox{M depends on $G_{\rm eff}$}.\\ \mbox{H}_0^{R21} = 73.04 \pm 1.04 \end{tabular} > H_0^{\rm P18} = 67.36 \pm 0.54 \mbox{ km s}^{-1} \mbox{Mpc}^{-1} \end{array}$$

Assumption:  $G_{eff}(z<0.01)=G_{eff}(z>0.01)$ 

#### The M transition hypothesis



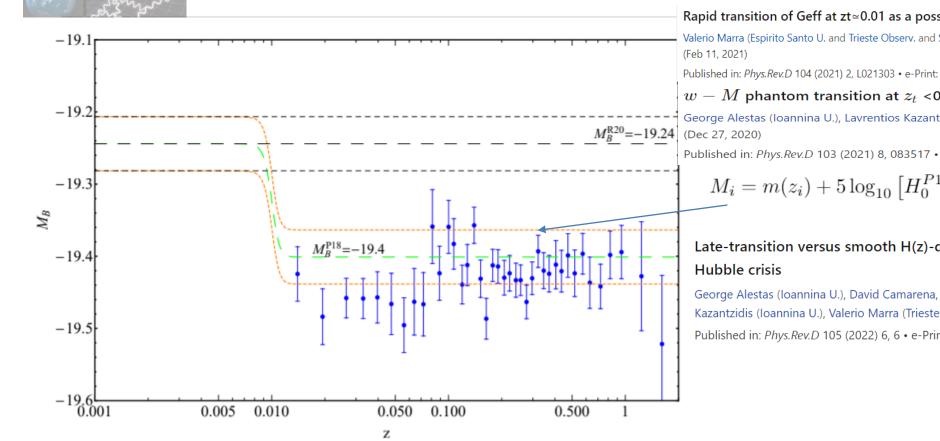
Rapid transition of Geff at zt≈0.01 as a possible solution of the Hubble and growth tensions

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$$M_i = m(z_i) + 5\log_{10} \left[ H_0^{P18} \cdot \text{Mpc}/c \right] - 5\log_{10}(D_L(z_i)) - 25$$

#### The M transition hypothesis



Rapid transition of Geff at zt≈0.01 as a possible solution of the Hubble and growth tensions

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w-M phantom transition at  $z_t$  <0.1 as a resolution of the Hubble tension George Alestas (Ioannina U.), Lavrentios Kazantzidis (Ioannina U.), Leandros Perivolaropoulos (Ioannina U.)

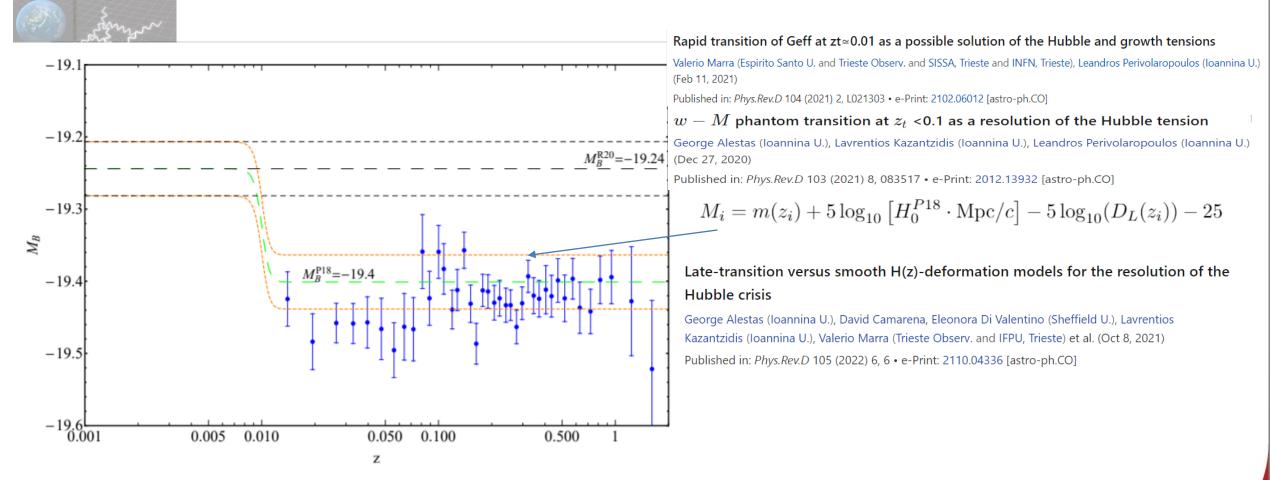
Published in: Phys.Rev.D 103 (2021) 8, 083517 • e-Print: 2012.13932 [astro-ph.CO]

 $M_i = m(z_i) + 5\log_{10} \left[ H_0^{P18} \cdot \text{Mpc}/c \right] - 5\log_{10}(D_L(z_i)) - 25$ 

#### Late-transition versus smooth H(z)-deformation models for the resolution of the

George Alestas (Joannina U.), David Camarena, Eleonora Di Valentino (Sheffield U.), Lavrentios Kazantzidis (Ioannina U.), Valerio Marra (Trieste Observ. and IFPU, Trieste) et al. (Oct 8, 2021) Published in: Phys.Rev.D 105 (2022) 6, 6 • e-Print: 2110.04336 [astro-ph.CO]

## The M transition hypothesis



A fundamental physics transition induces a transition of M (absolute magnitude or luminosity) at z<0.01.

#### Resolves M tension and Hubble tension.

Can potentially also resolve growth tension if the transition is connected with weaker gravity at z>z

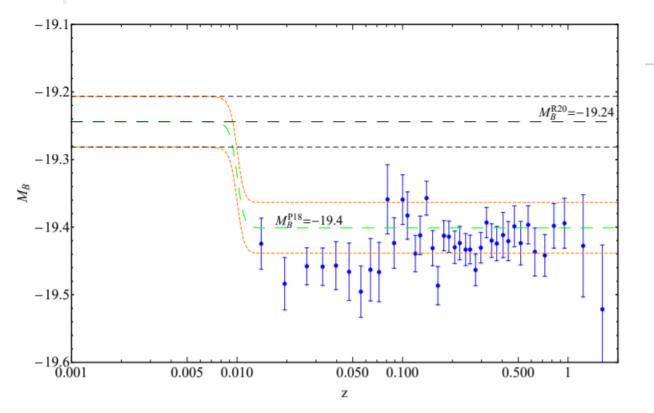
## Hints for the transition in data



#### Intrinsic tension in the supernova sector of the local Hubble constant measurement and its implications

Radosław Wojtak, Jens Hjorth (Jun 16, 2022) e-Print: 2206.08160 [astro-ph.CO] A reanalysis of the latest SH0ES data for  $H_0$ : Effects of new degrees of freedom on the Hubble tension Leandros Perivolaropoulos (Ioannina U.), Foteini Skara (Ioannina U.) (Aug 23, 2022) e-Print: 2208.11169 [astro-ph.CO]

#### **Gravitational Transition**



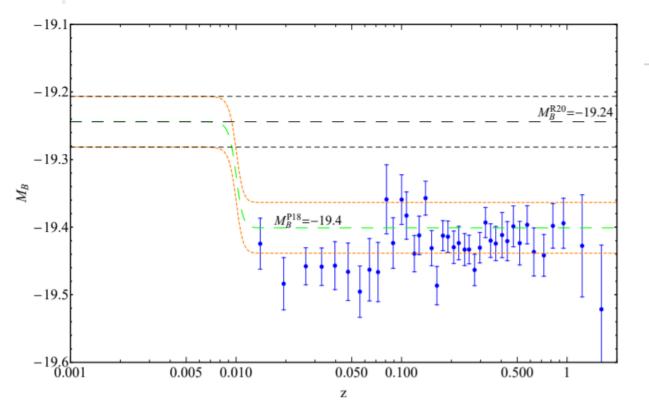
SnIa luminosities in the context of a Planck/ACDM background

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#### **Gravitational Transition**

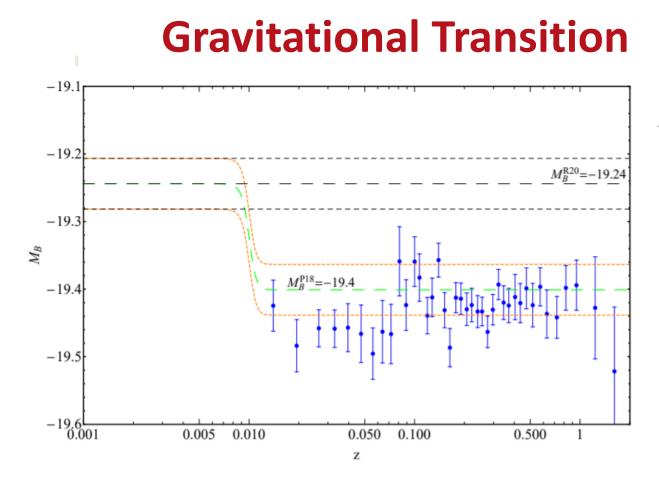


# A 10% transition of $G_{eff}$ is required for the reproduction of the required $\Delta M \sim 0.2$ for a pure Planck/ $\Lambda CDM$ background.

Rapid transition of Geff at zt≃0.01 as a possible solution of the Hubble and growth tensions Valerio Marra (Espirito Santo U. and Trieste Observ. and SISSA, Trieste and INFN, Trieste), Leandros Perivolaropoulos (Ioannina U.) (Feb 11, 2021)

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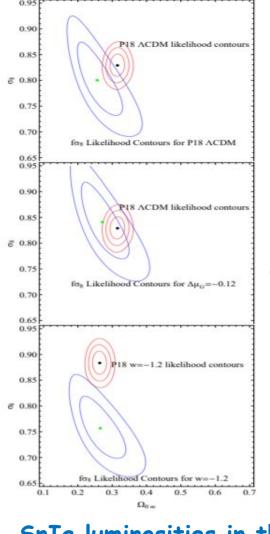
#### SnIa luminosities in the context of a Planck/ACDM background



A 10% transition of  $G_{eff}$  is required for the reproduction of the required  $\Delta M \sim 0.2$  for a pure Planck/ $\Lambda CDM$  background.

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The reduced value of  $G_{eff}$  leads to a higher  $\sigma_8$  value thus resolving the growth tension

Better fit to BAO. Growth tension resolved. M problem resolved

SnIa luminosities in the context of a Planck/ACDM background

#### The new SH0ES measurement of H<sub>0</sub>: The distance ladder in practice

Use the following system of 3492 equations fit for 47 unknown parameters

jth Cepheid in ith galaxy

 $m_{H,i,j}^{W} = \mu_{i} + M_{H}^{W} + b_{W}[P]_{i,j} + Z_{W}[O/H]_{i,j}$ 

 $m_{B,i} = \mu_i + M_B$ 

 $m_{B,i} - 5\log D_L(z_i) - 25 = M_B - 5\log H_0$ 

SnIa calibration

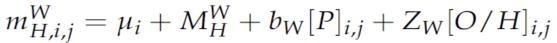
Cepheid calibration

 $M_B - 5 \log H_0$  m=µ(H<sub>0</sub>)+M<sub>B</sub> ->Hubble flow SnIa

#### The new SH0ES measurement of H<sub>0</sub>: The distance ladder in practice

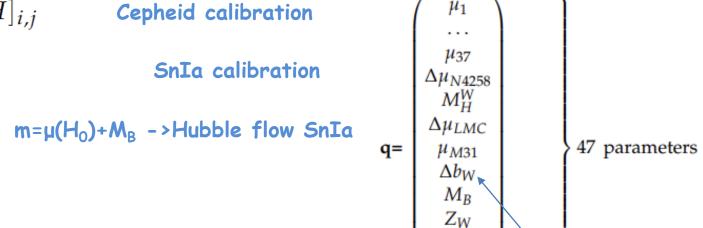
Use the following system of 3492 equations fit for 47 unknown parameters





 $m_{B,i} = \mu_i + M_B$ 

$$m_{B,i} - 5\log D_L(z_i) - 25 = M_B - 5\log H_0$$



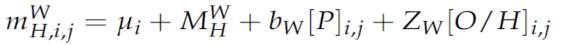
$$b_W = b_W^0 + \Delta b_W \equiv -3.286 + \Delta b_W$$

X  $\Delta zp$  $5 \log H_0$ 

#### The new SH0ES measurement of H<sub>0</sub>: The distance ladder in practice

Use the following system of 3492 equations fit for 47 unknown parameters





 $m_{B,i} = \mu_i + M_B$ 

$$m_{B,i} - 5\log D_L(z_i) - 25 = M_B - 5\log H_0$$

$$m=\mu(H_0)+M_B$$
 ->Hubble flow SnIa

SnIa calibration

Cepheid calibration

 $\mu_{37} \\ \Delta \mu_{N4258}$  $\Delta \mu_{LMC}$ 47 parameters  $\mu_{M31}$  $\Delta b_W$  $M_B$  $Z_W$ Х  $\Delta zp$  $5 \log H_0$  $b_W = b_W^0 + \Delta b_W \equiv -3.286 + \Delta b_W$ 

Express the system as linear vector transformation

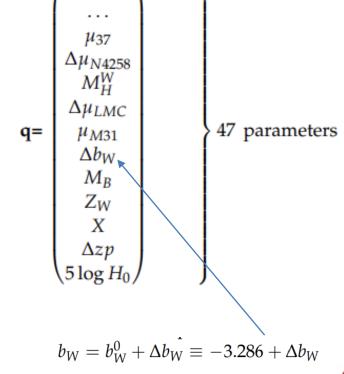
 $\mathbf{Y} = \mathbf{L}\mathbf{q}$  Minimize  $\chi^2$ :  $\chi^2 = (\mathbf{Y} - \mathbf{L}\mathbf{q})^{\mathrm{T}}\mathbf{C}^{-1}(\mathbf{Y} - \mathbf{L}\mathbf{q})$ 

#### The new SH0ES measurement of $H_0$ : The distance ladder in practice

Use the following system of 3492 equations fit for 47 unknown parameters jth Cepheid in ith galaxy  $m_{H,i,i}^{W} = \mu_{i} + M_{H}^{W} + b_{W}[P]_{i,j} + Z_{W}[O/H]_{i,j}$ Cepheid calibration  $m_{B,i}=\mu_i+M_B$   $\blacksquare$   $M_B^{R21}=-19.25\pm0.03$  SnIa calibration  $\Delta \mu_{N4258}$  $\Delta \mu_{LMC}$ m=µ(H<sub>0</sub>)+M<sub>B</sub> ->Hubble flow SnIa  $m_{B,i} - 5\log D_L(z_i) - 25 = M_B - 5\log H_0$ q=  $\mu_{M31}$  $\Delta b_W$  $M_B$  $H_0 = 73.04 \pm 1.04 \, km \, s^{-1} \, Mpc^{-1}$  $Z_W$ 

Express the system as linear vector transformation

 $\mathbf{Y} = \mathbf{L}\mathbf{q}$ Minimize  $\chi^2$ :  $\chi^2 = (Y - Lq)^T C^{-1} (Y - Lq)$ 



Y =	•	<b>The matrix equation</b> $\chi^2$ : $\chi^2 = (Y - Lq)^T C^{-1} (Y - Lq)$	Anchor constraints:	$M_{H}^{W} = Z_{W} = X$	$= -5.803 \pm 0.082$ = -5.903 \pm 0.025 = -0.21 \pm 0.12 = 0 \pm 0.00003 = 0 \pm 0.1
/-/)	$+ b_W[P]_{i,j} + Z_W[O/H]$	$[m_{B,i}]_{i,j}$ $m_{B,i} = \mu_i + M_B$ $m_{B,i} - 5\log D_L(z_i) - $	$25 = M_B - 5\log H_0$	$\Delta \mu_{N4258} =$	$= 0 \pm 10 \\ = 0 \pm 0.03 \\ = 0 \pm 0.026$
$\mathbf{Y} = \begin{pmatrix} \tilde{m}_{H,1}^{W} \\ & \cdots \\ \tilde{m}_{H,2150}^{W} \\ \bar{m}_{H,N4258,1}^{W} - \mu_{0,N4258} \\ & \cdots \\ \tilde{m}_{H,N4258,443}^{W} - \mu_{0,N4258} \\ & \tilde{m}_{H,N4258,443}^{W} - \mu_{0,N4258} \\ & \tilde{m}_{H,M31,1}^{W} \\ & \cdots \\ \tilde{m}_{H,MC,ground,1}^{W} - \mu_{0,LMC} \\ & \cdots \\ \tilde{m}_{H,LMC,ground,270}^{W} - \mu_{0,LMC} \\ & \tilde{m}_{H,SMC,ground,1}^{W} - \mu_{0,SMC} \\ & \cdots \\ \tilde{m}_{H,SMC,ground,143}^{W} - \mu_{0,SMC} \\ & \cdots \\ \tilde{m}_{H,SMC,ground,143}^{W} - \mu_{0,SMC} \\ & \cdots \\ \tilde{m}_{H,SMC,ground,143}^{W} - \mu_{0,SMC} \\ & \cdots \\ \tilde{m}_{H,LMC,HST,1}^{W} - \mu_{0,LMC} \\ & \cdots \\ \tilde{m}_{H,LMC,HST,69}^{W} - \mu_{0,LMC} \\ & \tilde{m}_{B,1}^{U} \\ & \cdots \\ \tilde{m}_{B,277}^{W} \\ \hline -5.803 \left( M_{H,Gaia}^{W} \right) \\ & -0.21 \left( Z_{W,Gaia} \right) \\ & 0 \left( \Delta zp \right) \\ & 0 \left( \Delta zp \right) \\ & 0 \left( \Delta \mu_{M258} \right) \\ & 0 \left( \Delta \mu_{LMC} \right) \\ & m_{B,1}^{0} - 5 \log[c_{21}(\cdots)] - 25 \\ & \cdots \\ m_{B,277}^{0} - 5 \log[c_{2277}(\cdots)] - 25 \end{pmatrix}$	<pre> } 2150 Cepheids in 37 SnIa hosts  980 Cepheids in non SnIa hosts  P80 Cepheids in non SnIa hosts L = } 77 SnIa in Cepheid hosts 8 External constraints 277 SnIa in Hubble flow </pre>	$= \left(\begin{array}{cccccccccccccccccccccccccccccccccccc$	Cepheids in non SnIa hosts <b>q=</b> 11a in Cepheid hosts	$ \mu_{1} \\ \mu_{37} \\ \mu_{N4258} \\ M_{H}^{W} \\ \mu_{LMC} \\ \mu_{M31} \\ \Delta b_{W} \\ M_{B} \\ Z_{W} \\ X \\ \Delta zp \\ \log H_{0} $	47 parameters

$m_{H,i,j}^W = \mu_i + M_H^W$	ELQ Minimize	<b>The matrix equation</b> $\chi^{2}:  \chi^{2} = (\Upsilon - Lq)^{T}C^{-1}(\Upsilon - Lq)$ $K^{2}:  \chi^{2} = (\chi - Lq)^{T}C^{-1}(\chi - Lq)$ $m_{B,i} = \mu_{i} + M_{B} \qquad m_{B,i} - 5 \log D_{L}(z_{i}) - 25 = M_{B} - 5$	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$
$ \begin{pmatrix} \bar{m}_{H,1}^W \\ \dots \\ \bar{m}_{H,2150}^W \\ \bar{m}_{H,N4258,1}^W - \mu_{0,N4258} \\ \dots \\ \bar{m}_{H,N4258,443}^W - \mu_{0,N4258} \\ \end{pmatrix} $	2150 Cepheids in 37 Sn Ia hosts	$\left(\begin{array}{cccccccccccccccccccccccccccccccccccc$	$(\mu_1)$
$\bar{m}_{H,M31,1}^{W}$  $\bar{m}_{H,M31,55}^{W}$ $\bar{m}_{H,LMC,ground,1}^{W} - \mu_{0,LMC}$  $\bar{m}_{H,LMC,ground,270}^{W} - \mu_{0,LMC}$ $\bar{m}_{H,SMC,ground,1}^{W} - \mu_{0,SMC}$	> 980 Cepheids in non SnIa hosts	$ \begin{bmatrix} 0 & \dots & 0 & 1 & 1 & 0 & 0 & [P]_{N4258,443} & 0 & [O/H]_{N4258,443} & 0 & 0 & 0 \\ 0 & \dots & 0 & 0 & 1 & 0 & 1 & [P]_{M31,1} & 0 & [O/H]_{M31,1} & 0 & 0 & 0 \\ \dots & \dots \\ 0 & \dots & 0 & 0 & 1 & 0 & 1 & [P]_{M31,55} & 0 & [O/H]_{M31,55} & 0 & 0 & 0 \\ 0 & \dots & 0 & 0 & 1 & 1 & 0 & [P]_{LMC,ground,1} & 0 & [O/H]_{LMC,ground,1} & 0 & 1 & 0 \\ \dots & \dots \\ 0 & \dots & 0 & 0 & 1 & 1 & 0 & [P]_{LMC,ground,270} & 0 & [O/H]_{LMC,ground,270} & 0 & 1 & 0 \\ \dots & \dots$	$ \begin{array}{c} \mu_{37} \\ \Delta \mu_{N4258} \\ M_H^W \\ \Delta \mu_{LMC} \end{array} $
$\mathbf{Y} = \begin{bmatrix} \bar{m}_{H,SMC,ground,143}^{W} - \mu_{0,SMC} \\ \bar{m}_{H,LMC,HST,1}^{W} - \mu_{0,LMC} \\ & \\ & \\ & \\ & \\ & \\ & \\ & \\ & \\ & \\ $	<pre> L = 77 SnIa in Cepheid hosts </pre>		$ \begin{array}{c c} \mathbf{q} = & \mu_{M31} \\ & \Delta b_W \\ & M_B \\ & Z_W \\ & X \end{array} $ 47 parameters
$\begin{array}{c} -0.21 \ (Z_{W,Gaia}) \\ 0 \ (X) \\ 0 \ (\Delta z p) \\ 0 \ (\Delta b_W) \\ 0 \ (\Delta \mu_{\rm N4258}) \end{array}$	8 External constraints	0        0       0       0       0       1       0       0       0       0         0        0       0       0       0       0       1       0 </td <td><math>\begin{pmatrix} \Delta zp \\ 5 \log H_0 \end{pmatrix}</math></td>	$\begin{pmatrix} \Delta zp \\ 5 \log H_0 \end{pmatrix}$
$ \begin{array}{c} 0 \ (\Delta \mu_{LMC}) \\ \hline m_{B,1}^0 - 5 \log[cz_1()] - 25 \\ \hline \dots \\ m_{B,277}^0 - 5 \log[cz_{277}()] - 25 \end{array} $	} 277 SnIa in Hubble flow	$\left(\begin{array}{cccccccccccccccccccccccccccccccccccc$	

#### Analytic minimization of $\chi^2$

$$\mathbf{Y} = \mathbf{L}\mathbf{q} \quad \text{Minimize } \chi^2: \quad \chi^2 = (\mathbf{Y} - \mathbf{L}\mathbf{q})^{\mathsf{T}}\mathbf{C}^{-1}(\mathbf{Y} - \mathbf{L}\mathbf{q})$$
Best fit parameter values:  

$$\chi^2 = (\mathbf{Y} - \mathbf{L}\mathbf{q})^{\mathsf{T}}\mathbf{C}^{-1}(\mathbf{Y} - \mathbf{L}\mathbf{q}) = \mathbf{q}^{\mathsf{T}}\mathbf{L}^{\mathsf{T}}\mathbf{C}^{-1}\mathbf{L}\mathbf{q} - 2\mathbf{q}^{\mathsf{T}}\mathbf{L}^{\mathsf{T}}\mathbf{C}^{-1}\mathbf{Y} + \mathbf{Y}^{\mathsf{T}}\mathbf{C}^{-1}\mathbf{Y}$$

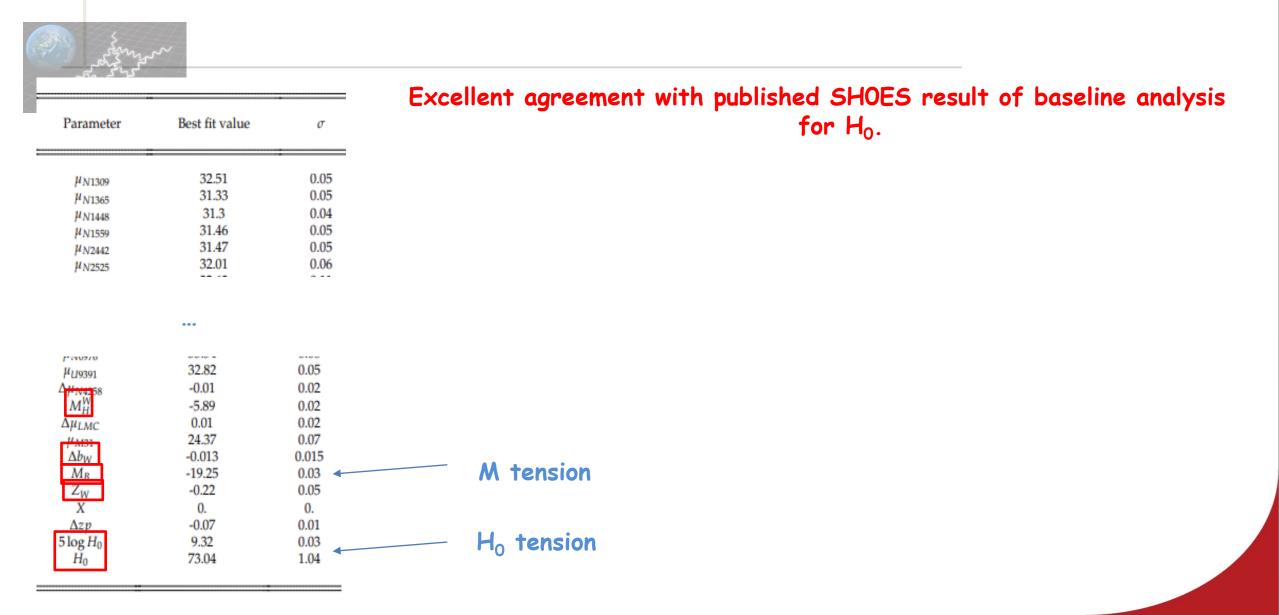
$$\frac{\partial\chi^2}{\partial\mathbf{q}}\Big|_{\mathbf{q}_{best}} = 0 => 2\mathbf{L}^{\mathsf{T}}\mathbf{C}^{-1}\mathbf{L}\mathbf{q}_{best} - 2\mathbf{L}^{\mathsf{T}}\mathbf{C}^{-1}\mathbf{Y} = 0$$

$$\mathbf{q}_{best} = (\mathbf{L}^{\mathsf{T}}\mathbf{C}^{-1}\mathbf{L})^{-1}\mathbf{L}^{\mathsf{T}}\mathbf{C}^{-1}\mathbf{Y}$$

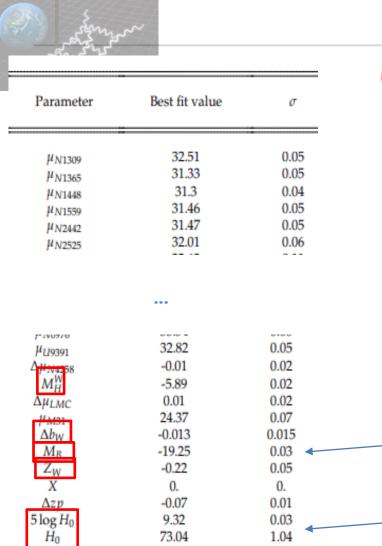
1 $\sigma$  errors of best fit parameters are diagonal elements of transformed covariance matrix:

$$\Sigma_{kl} = \sum_{i} \sum_{j} \left[ \frac{\partial q_{best,k}}{\partial Y_i} \right] C_{ij} \left[ \frac{\partial q_{best,l}}{\partial Y_j} \right] \qquad \Sigma = (\mathbf{L}^{T} \mathbf{C}^{-1} \mathbf{L})^{-1}$$

#### **Best fit parameter values**



#### **Best fit parameter values**



Excellent agreement with published SH0ES result of baseline analysis for  $H_0$ .

Main questions:

1. Self-consistency test: Are the values of the Cepheid modeling parameters  $b_{Wi}$  and  $Z_{Wi}$  obtained from each host i consistent with a universal value for each parameter?

2. Extension of baseline analysis with new degrees of freedom: How can new degrees of freedom and/or constraints be included in a generalized analysis?

**M** tension

 $H_0$  tension

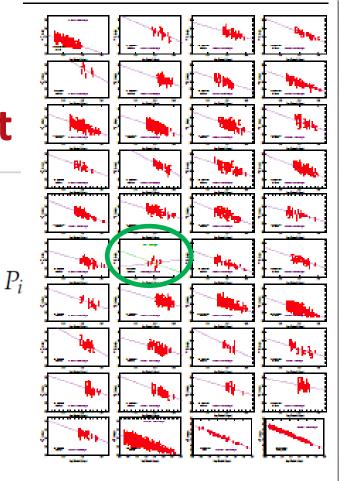
Best fit  $m_H^W$ -logP slopes:

#### Self consistency test I Modeling parameter b<sub>wi</sub> in each host



Find best fit b<sub>Wi</sub> slope and intercept s<sub>i</sub> in each Cepheid host i (37 SnIa+Cepheid + 3 pure Cepheid hosts)

$$m_{H,i,j}^{W} = \mu_i + M_H^{W} + b_W[P]_{i,j} + Z_W[O/H]_{i,j} \qquad m_{H,i,j}^{W} = s_i + b_{W,i} \log p_{i,j}$$

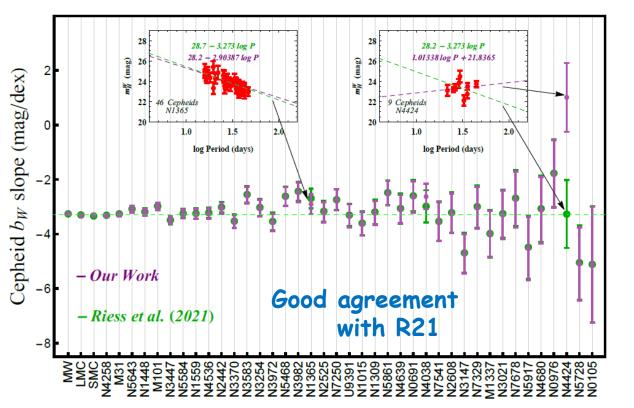


Best fit m<sub>H</sub><sup>W</sup>-logP slopes:

#### Self consistency test I **Modeling parameter b<sub>wi</sub> in each host**

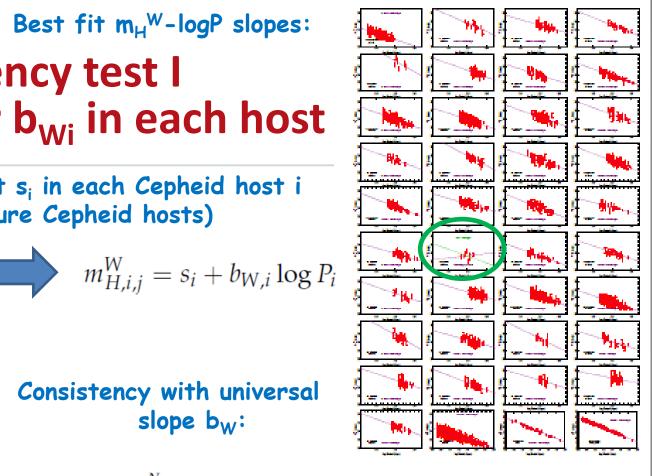
Find best fit b<sub>wi</sub> slope and intercept s<sub>i</sub> in each Cepheid host i (37 SnIa+Cepheid + 3 pure Cepheid hosts)

$$m_{H,i,j}^W = \mu_i + M_H^W + b_W [P]_{i,j} + Z_W [O/H]_{i,j}$$



Consistency with universal slope b<sub>w</sub>:

$$\chi^2(b_W) = \sum_{i=1}^N \frac{(b_{W,i} - b_W)^2}{\sigma_{b_{W,i}}^2} \quad \blacksquare$$



$$\frac{\chi^2_{bW,min}}{dof} = 1.55$$

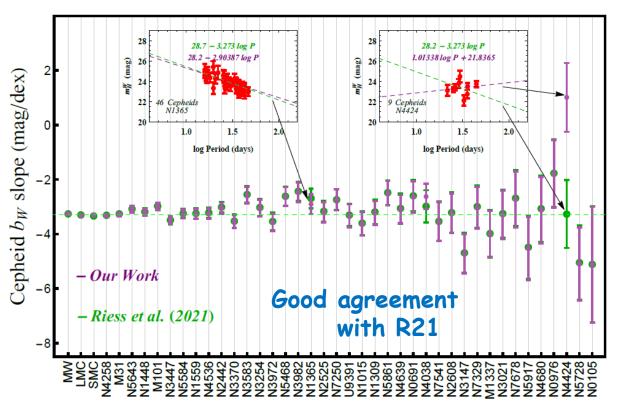
Best fit  $m_H^W$ -logP slopes:

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#### Self consistency test I Modeling parameter b<sub>wi</sub> in each host

Find best fit b<sub>Wi</sub> slope and intercept s<sub>i</sub> in each Cepheid host i (37 SnIa+Cepheid + 3 pure Cepheid hosts)

$$m_{H,i,j}^W = \mu_i + M_H^W + b_W[P]_{i,j} + Z_W[O/H]_{i,j}$$
  $m_{H,i,j}^W = s_i + b_{W,i}\log P_i$ 



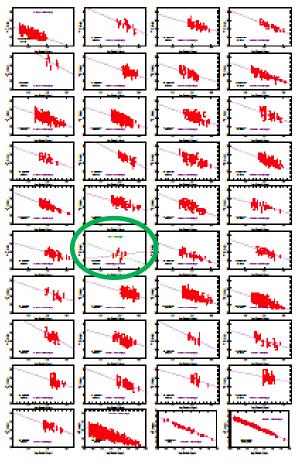
Consistency with universal slope b<sub>W</sub>:

$$\chi^2(b_W) = \sum_{i=1}^N \frac{(b_{W,i} - b_W)^2}{\sigma_{b_{W,i}}^2} \quad \blacksquare$$

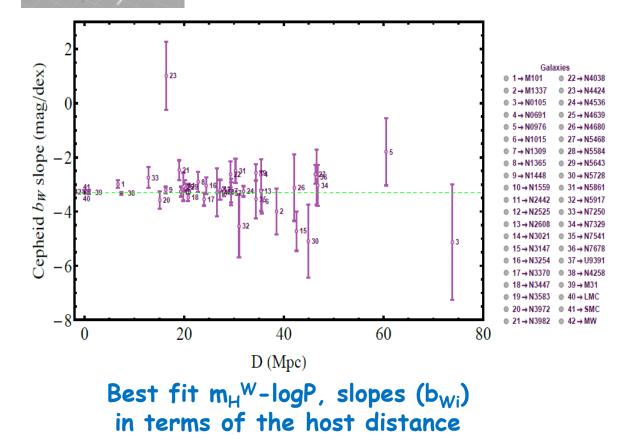
$$\frac{\chi^2_{bW,min}}{dof} = 1.55$$

$$\frac{\chi^2_{bW,min}}{dof} \simeq 1$$
  $\sigma_{b,scat} \simeq 0.18$ 

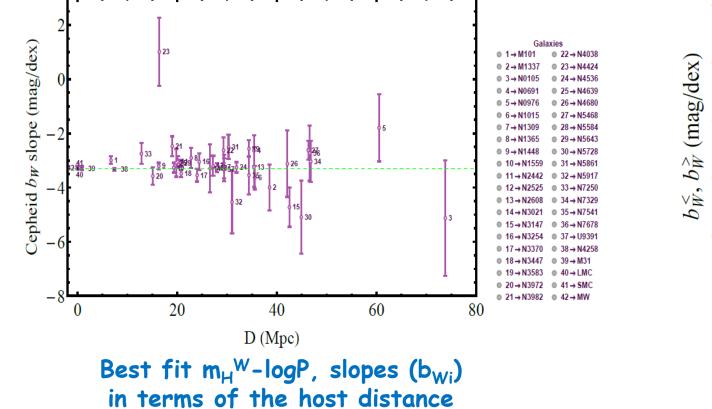
Allow for additional uncertainty  $\sigma_{\text{scat}}$ :  $\chi^2(b_W) = \sum_{i=1}^N \frac{(b_{W,i} - b_W)^2}{\sigma_{b_{W,i}}^2 + \sigma_{b,scat}^2}$ 

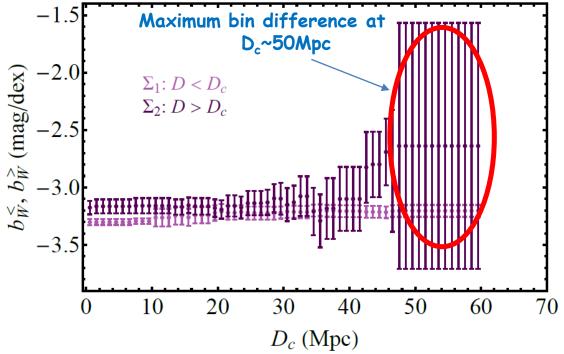


#### Modeling parameter b<sub>wi</sub> in each host i: Consistency of distance bins



#### Modeling parameter b<sub>wi</sub> in each host i: Consistency of distance bins





Consistency between high and low distance bins split at distance  $D_c$ 

#### Self consistency test II Modeling parameter Z<sub>Wi</sub> in each host i

Find best fit Z<sub>Wi</sub> slope and intercept s<sub>i</sub> in each Cepheid host i (37 SnIa+Cepheid + 2 pure Cepheid hosts)

$$m_{H,i,j}^{W} = \mu_{i} + M_{H}^{W} + b_{W}[P]_{i,j} + Z_{W}[O/H]_{i,j} \qquad m_{H,i,j}^{W} = s_{i} + Z_{W,i}[O/H]_{i,j}$$

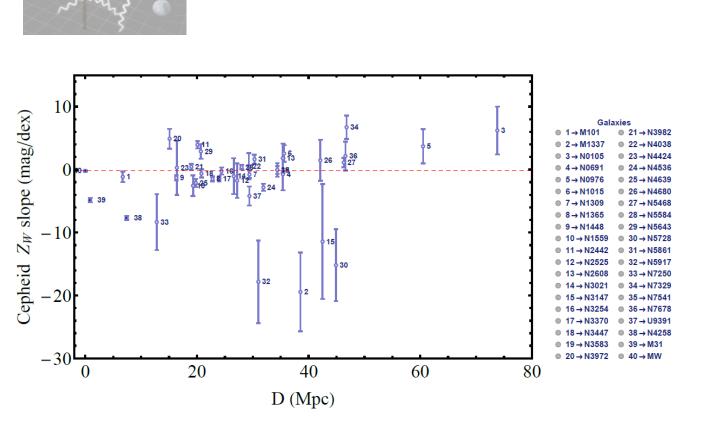
#### Consistency with universal metallicity slope $Z_W$ :

Best fit  $m_H^W$ -[O/H] slopes

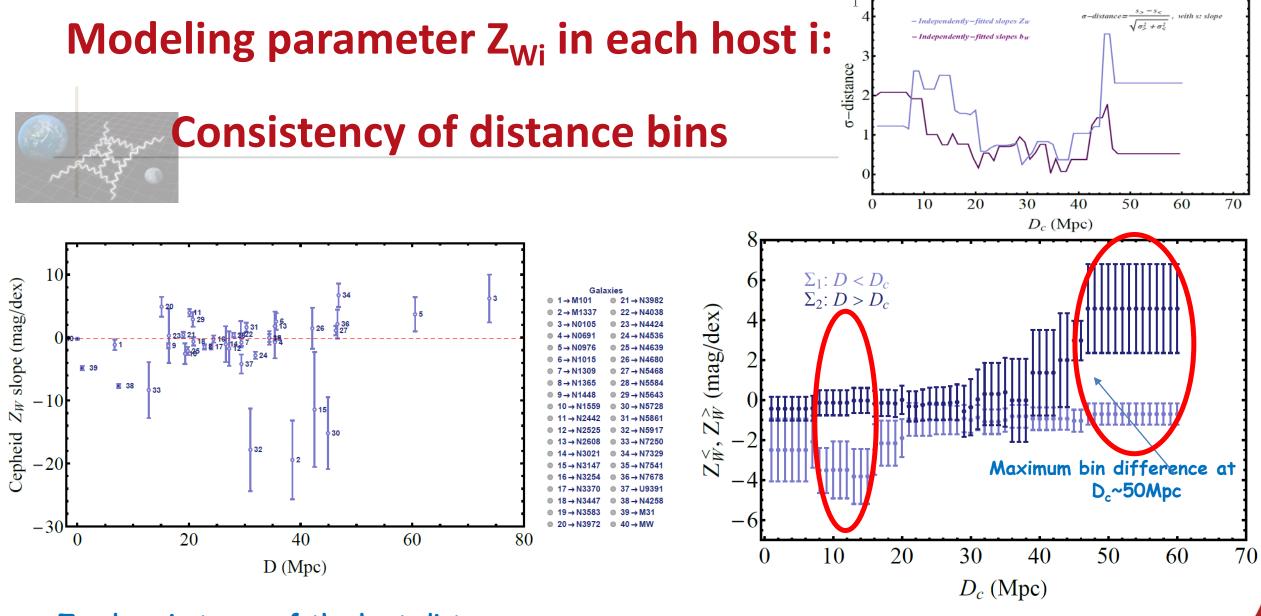
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#### Modeling parameter Z<sub>Wi</sub> in each host i:

**Consistency of distance bins** 



#### $Z_{wi}$ slope in terms of the host distance



#### $Z_{wi}$ slope in terms of the host distance

Consistency between high and low distance bins split at distance D<sub>c</sub>

#### Generalizing the baseline SH0ES modeling analysis: New degrees of freedom



Allow for a change (transition) of the modeling parameters  $M_W$ ,  $b_W$ ,  $Z_W$ ,  $M_B$  at a given distance  $D_c$  (cosmic time  $t_c$ ).

For example if  $b_W$  was allowed to change, the Cepheid modeling would have to change as:

$$m_{H,i,j}^{W} = \mu_{i} + M_{H}^{W} + b_{W}[P]_{i,j} + Z_{W}[O/H]_{i,j}$$

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For example if  $b_W$  was allowed to change, the Cepheid modeling would have to change as:

 $m_{H,i,j}^{W} = \mu_{i} + M_{H}^{W} + b_{W}[P]_{i,j} + Z_{W}[O/H]_{i,j}$  $m_{H,i,j}^{W}(D) = \mu_{i} + M_{H}^{W} + b_{W}^{\geq}\Theta(D - D_{c})[P]_{i,j} + b_{W}^{\leq}\Theta(D_{c} - D)[P]_{i,j} + Z_{W}[O/H]_{i,j}$ 

The new matrix equation Y=L q would have the same data/constraints Y (labeled with their distance) the same covariance matrix C but different model matrix L and parameter matrix q.

#### The new matrix equation

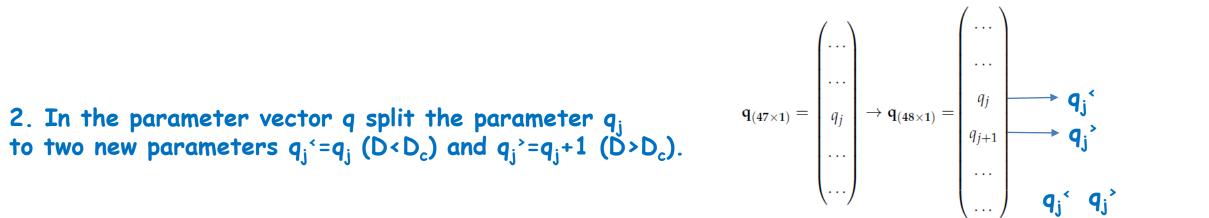
Allow for a change (transition) of the modeling parameter  $q_j$  at a given distance  $D_c$  (cosmic time  $t_c$ ).

1. Assign a distance  $D_i$  to each entry 1<i<3492 of the data/constraint vector Y.

#### The new matrix equation

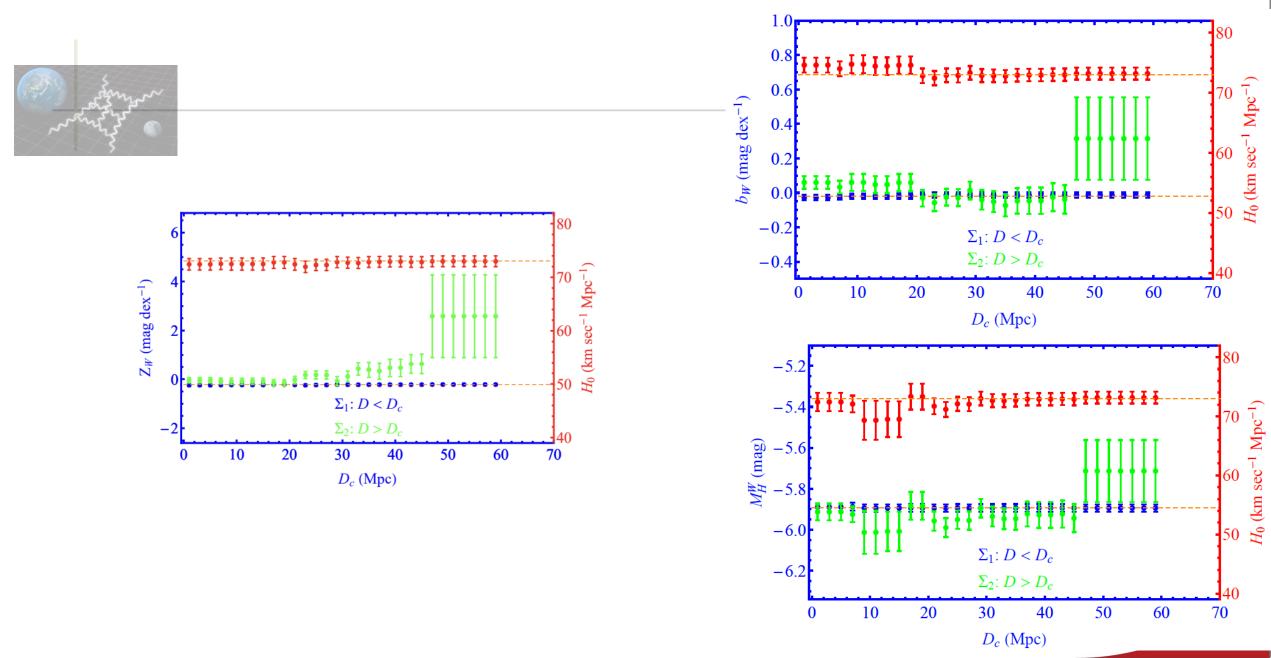
Allow for a change (transition) of the modeling parameter  $q_i$  at a given distance  $D_c$  (cosmic time t<sub>c</sub>).

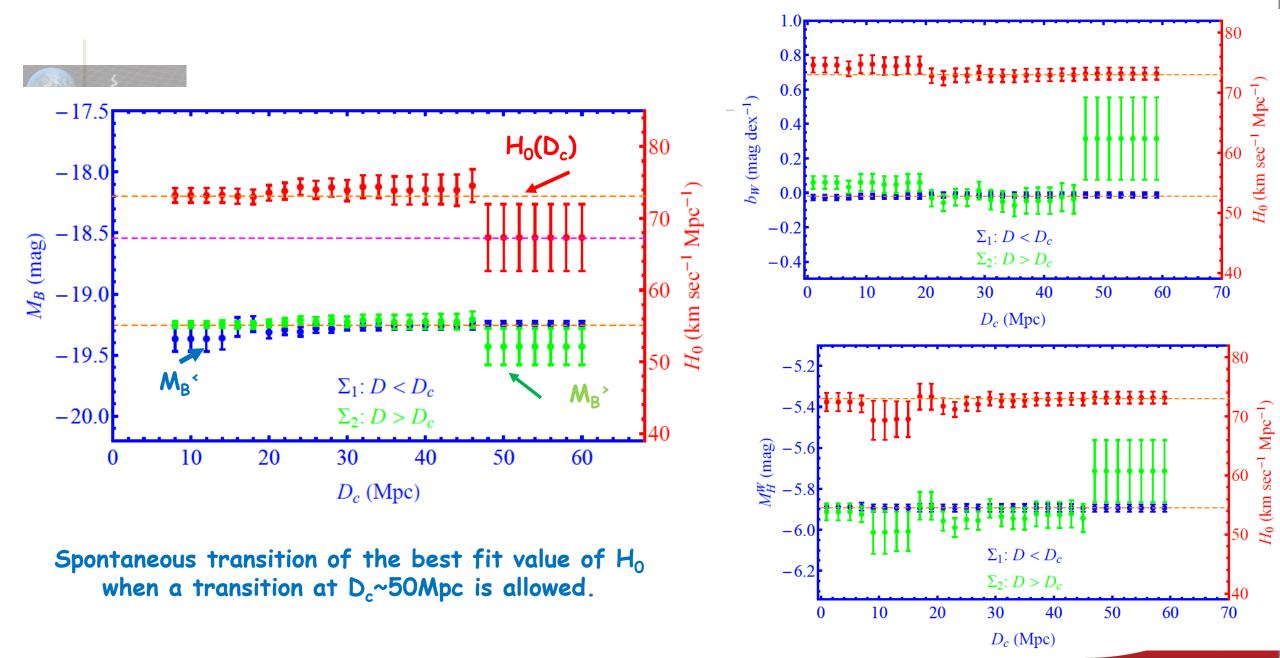
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#### The new matrix equation

Allow for a change (transition) of the modeling parameter  $q_i$  at a given distance  $D_c$  (cosmic time  $t_{c}$ ). 1. Assign a distance  $D_i$  to each entry 1<i<3492 of the data/constraint vector Y. 2. In the parameter vector q split the parameter  $q_{j}$ to two new parameters  $q_{j} \leq q_{j}$  (D<D<sub>c</sub>) and  $q_{j} \geq q_{j} + 1$  (D>D<sub>c</sub>). 3. In the modeling matrix L split column j to two new columns corresponding to D<D<sub>c</sub> (zero entries for rows with D>D<sub>c</sub> in Y) and D>D<sub>c</sub> (zero entries for rows with D<D<sub>c</sub> in Y).  $L_{3492,477} = \begin{pmatrix} \cdots & L_{1,j} & \cdots \\ \cdots & L_{3491,j} & \cdots \\ \cdots & L_{3491,j} & \cdots \\ \cdots & L_{3491,j} & \cdots \\ \cdots & L_{3492,477} = \begin{pmatrix} \cdots & L_{1,j} & \cdots \\ \cdots & L_{3491,j} & \cdots \\ \cdots & L_{3492,477} = \begin{pmatrix} \cdots & L_{1,j} & \cdots \\ \cdots & L_{3491,j} & \cdots \\ \cdots & \dots & \dots \\ D_{Y_{362}} > D_{c} & \cdots \\ \cdots & D_{Y_{362}} > D_{c}$ 





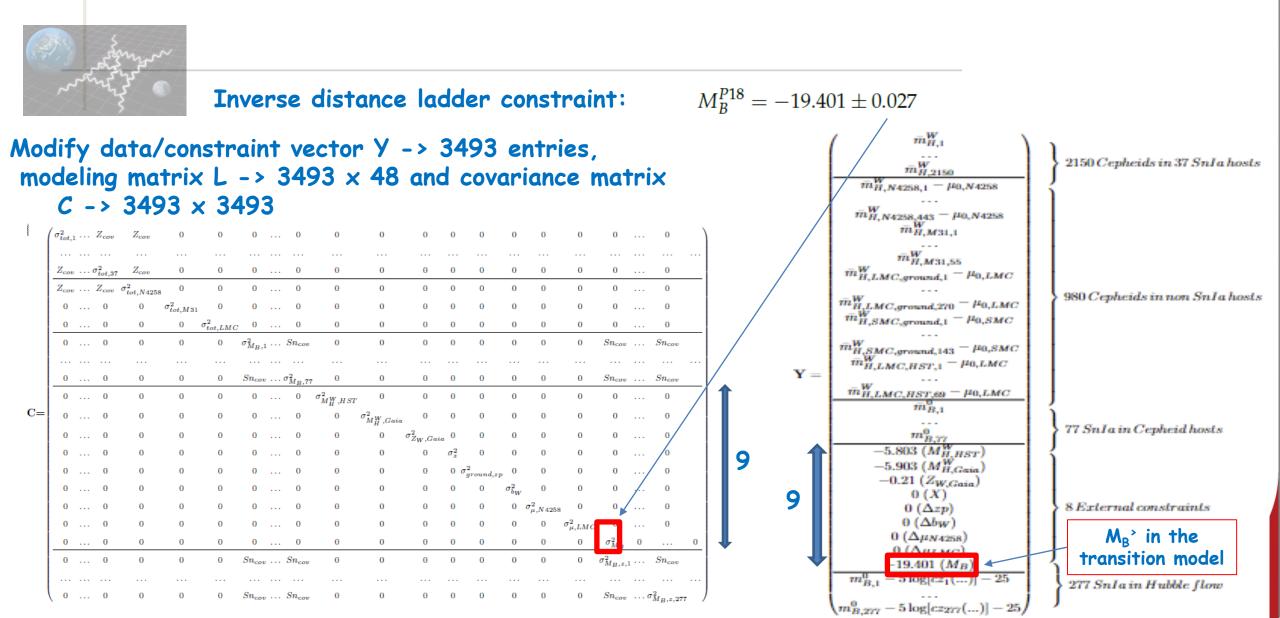
# **Including the inverse distance ladder constraint**



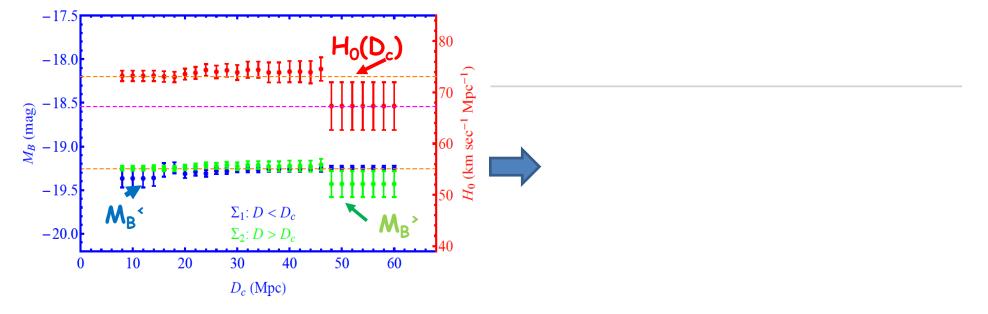
**Inverse distance ladder constraint:**  $M_B^{P18} = -19.401 \pm 0.027$ 

Modify data/constraint vector Y -> 3493 entries, modeling matrix L -> 3493 × 48 and covariance matrix C -> 3493 x 3493

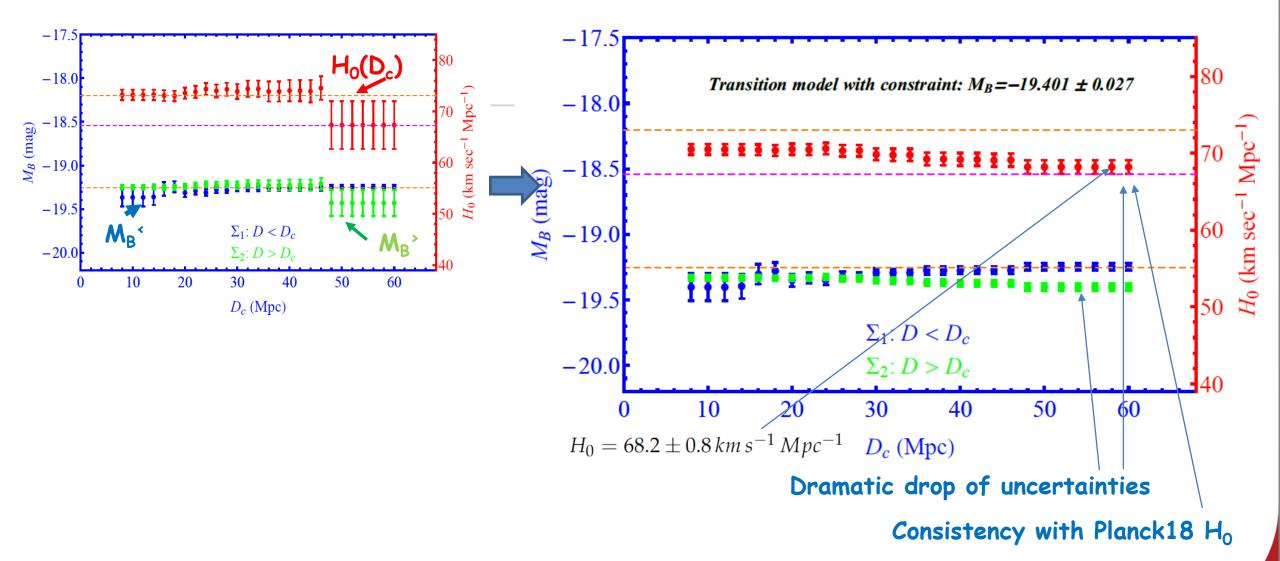
#### Including the inverse distance ladder constraint



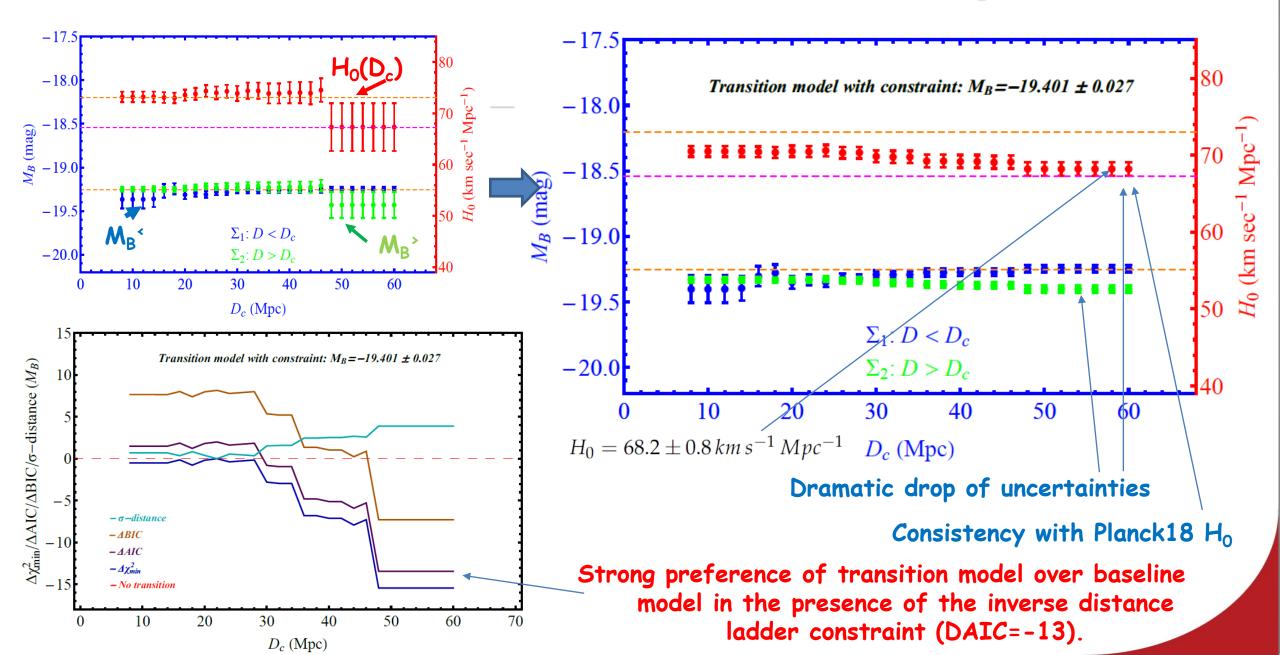
# **Results of the Generalized Analysis with M<sub>B</sub> constraint**



## **Results of the Generalized Analysis with M<sub>B</sub> constraint**



## **Results of the Generalized Analysis with M<sub>B</sub> constraint**



	Model	$\chi^2_{min}$	$\chi^2_{red}$ <sup>a</sup>	ΔAIC	ΔBIC	$H_0$ [Km s <sup>-1</sup> Mpc <sup>-1</sup> ]	$M_B$ [mag]	$M_H^W$ [mag]	∆b <sub>W</sub> [mag/dex]	Z <sub>W</sub> [mag/dex]
	Baseline	3552.76	1.031	0	0	$73.043 \pm 1.007$	$-19.253 \pm 0.029$	$-5.894 \pm 0.018$	$-0.013 \pm 0.015$	$-0.217 \pm 0.045$
	Transition <sup>b</sup> $M_B$	3551.31	1.031	0.55	6.71	$67.326 \pm 4.647$	$-19.250 \pm 0.029$ $-19.430 \pm 0.150$ $1.2\sigma$	$-5.894 \pm 0.018$	$-0.013 \pm 0.015$	$-0.217 \pm 0.045$
	Transition <sup>b</sup> $M_H^W$	3551.31	1.031	0.55	6.71	$73.162 \pm 1.014$	$-19.250 \pm 0.029$	$-5.894 \pm 0.018$ $-5.713 \pm 0.151$ $1.2\sigma$	$-0.013 \pm 0.015$	$-0.217 \pm 0.045$
	Transition <sup>b</sup> $Z_W$	3549.99	1.030	-0.77	5.39	$72.981 \pm 1.007$	$-19.255 \pm 0.029$	$-5.894 \pm 0.018$	$-0.014 \pm 0.015$	$-0.217 \pm 0.045$ $2.588 \pm 1.686$ $1.7\sigma$
	Transition <sup><math>b</math></sup> $b_W$	3550.86	1.030	0.10	6.26	$73.173 \pm 1.013$	$-19.249 \pm 0.029$	$-5.894 \pm 0.018$	$\begin{array}{c} -0.013 \pm 0.015 \\ 0.315 \pm 0.239 \\ 1.4\sigma \end{array}$	$-0.217 \pm 0.045$
	Baseline+Constraint <sup>c</sup>	3566.78	1.035	0	0	$70.457 \pm 0.696$	$-19.332 \pm 0.020$	$-5.920 \pm 0.017$	$-0.026 \pm 0.015$	$-0.220 \pm 0.045$
	Transition <sup><math>b,c</math></sup> $M_B$ +Constraint	3551.34	1.031	-13.44	-7.27	$68.202 \pm 0.879$	$-19.249 \pm 0.029$ $-19.402 \pm 0.027$ $3.9\sigma$	$-5.893 \pm 0.018$	$-0.013 \pm 0.015$	$-0.217 \pm 0.045$

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#### **Main Points**



- 1. New physics at ultra low redshifts (z<0.02) can lead to a resolution of the Hubble tension.
- 2. New degrees of freedom can be introduced in the baseline SHOES analysis to detect such new physics signatures.
- If a change of the SnIa absolute luminosity M<sub>B</sub> is allowed at D<sub>c</sub>~50Mpc (new degree of freedom) in the SHOES analysis then the best fit value of H<sub>0</sub> drops from 73±1 km/(sec Mpc) to 67±4 km/(sec Mpc).
- 4. In the presence of the inverse distance ladder input on M<sub>B</sub> the transition model uncertainties drop dramatically and the extended model is strongly preferred over the SHOES baseline model of universal value of M<sub>B</sub>.

#### **Theoretical Model: Scalar Tensor Theory**



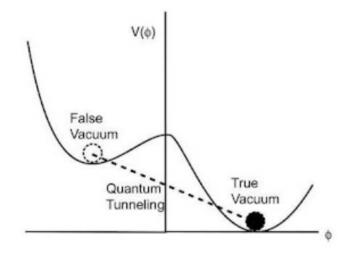
# $\text{Scalar Tensor Transition:} \quad S = \int d^4x \sqrt{|g|} \left[ \frac{1}{2} \xi \varphi^2 R - \frac{1}{2} (\partial \varphi)^2 - V(\varphi) + \mathcal{L}_m \right], \qquad 8\pi G_N = \xi^{-1} v^{-2}$

v: potential minimum

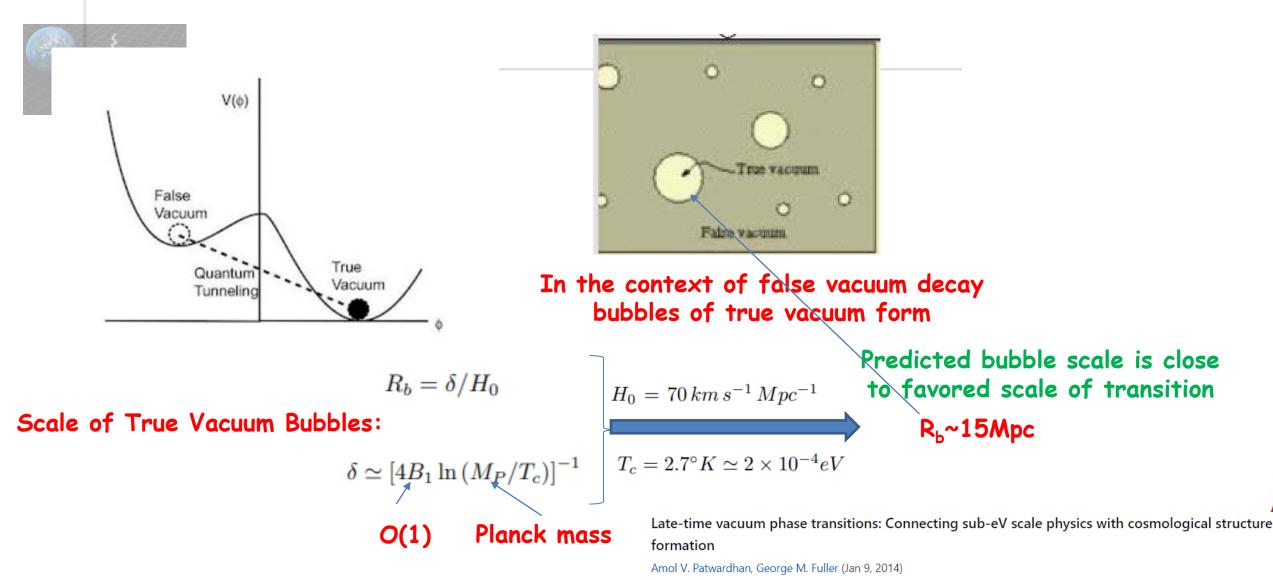
$$v = \frac{1}{\sqrt{8\pi G_N}} = M_{\rm Pl} \sim 10^{19} {\rm GeV},$$
 Cosmologic

Cosmological Constant:  $\Lambda = V(v)$ 

A phase transition (false vacuum decay) would induce a transition in the strength of gravity as well



#### **Generic Distance Scale**



Published in: Phys.Rev.D 90 (2014) 6, 063009 • e-Print: 1401.1923 [astro-ph.CO]

#### **Constraints**

**Table 1.** Solar system, astrophysical and cosmological constraints on the evolution of the gravitational constant. Methods with star (\*) constrain  $G_N$ , while the rest constrain  $G_{\text{eff}}$ . The latest and strongest constraints are shown for each method.

Method	$\left  \frac{\Delta G_{\rm eff}}{G_{\rm eff}} \right _{max}$	$\left \frac{\hat{G}_{\rm eff}}{G_{\rm eff}}\right _{max} (yr^{-1})$	Time Scale (Yr)	References
Lunar ranging		$1.47 imes10^{-13}$	24	[34]
Solar system		$4.6 imes10^{-14}$	50	[35,36]
Pulsar timing		$3.1  imes 10^{-12}$	1.5	[37]
Strong Lensing		$10^{-2}$	0.6	[38]
Orbits of binary pulsar		$1.0  imes 10^{-12}$	22	[39]
Ephemeris of Mercury		$4 imes 10^{-14}$	7	[40]
Exoplanetary motion		$10^{-6}$	4	[41]
Hubble diagram SnIa	0.1	$1  imes 10^{-11}$	${\sim}10^8$	[42]
Pulsating white-dwarfs		$1.8 imes10^{-10}$	0	[43]
Viking lander ranging		$4 imes 10^{-12}$	6	[44]
Helioseismology		$1.6 imes10^{-12}$	$4 imes 10^9$	[45]
Gravitational waves	8	$5 imes 10^{-8}$	$1.3 imes10^8$	[46]
Paleontology	0.1	$2 imes 10^{-11}$	$4 imes 10^9$	[47]
Globular clusters		$35  imes 10^{-12}$	${\sim}10^{10}$	[48]
Binary pulsar masses		$4.8 imes10^{-12}$	${\sim}10^{10}$	[49]
Gravitochemical heating		$4 imes 10^{-12}$	${\sim}10^8$	[50]
Strong lensing		$3 imes 10^{-1}$	${\sim}10^{10}$	[38]
Big Bang Nucleosynthesis *	0.05	$4.5 imes10^{-12}$	$1.4 imes10^{10}$	[30]
Anisotropies in CMB *	0.095	$1.75  imes 10^{-12}$	$1.4 imes10^{10}$	[51]

#### Hints for a Gravitational Transition in Tully-Fisher Data

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# Main Points / Conclusion

Viable early and late approaches to the Hubble tension appear to require the existence of an abrupt transition event either at  $t_{rec}$  or at present  $t_0$ .

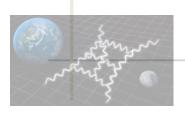
The late transition event may involve a sudden diming of the SnIa intrinsic luminosity occurring less than 150 million years ago (z<sub>t</sub><0.01).

Such a diming may be due to a sudden increase of the strength  $G_{eff}$  of gravitational interactions by about 10% at  $z_t < 0.01$ . This is a viable and testable conjecture.

There are hints for such a transition in the recent Cepheid+SnIa SHOES data for the measurement of HO with distance ladder methods.

Theoretical models supporting such an event may involve a false vacuum decay models where the true vacuum has a similar energy scale as the observed cosmological constant scale (0.002eV) and a decay rate Γ~H<sub>0</sub>.

## The distance ladder and the new SHOES data



Calibrating Cepheids+SnIa with geometric distance measurements: The distance ladder (distance  $\langle - \rangle$  absolute luminosity ladder m=µ+M(p))

- 1. Measure geometrically (parallax etc) the distance to nearby Cepheid variable stars and their properties (period, metallicity). Thus find Cepheid absolute luminosity in anchor galaxies (d<8Mpc).
- 2. Calibrate Cepheids: Connect absolute luminosity to their properties (period, metallicity).
- 3. Use calibrated Cepheids to measure distances to Cepheid+SnIa hosts (larger distances). Thus find SnIa absolute luminosities
- 4. Calbrate SnIa: Connect absolute luminosity to their properties (light curve stretch, color).
- 5. Use calibrated SnIa (known absolute luminosity) to measure distances in the Hubble flow. Thus, measure H<sub>0</sub>.

New SHOES data release: Significant increase of number of Cepheid calibrators (from 2416 to 3130) of SnIa calibrators in SnIa hosts (from 19 to 42) and of SnIa in Hubble flow (from 217 to 277).