

Quantum vacuum, a cosmic chameleon

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Introduction

Introduction

- ❖ Cosmological tensions are a motivation for new physics. Looking from different directions.
- ❖ Cosmological principle permits VED to be a function $\rho_{\text{vac}}(\chi(t))$, where $\chi(t)$ is a dynamical variable.
- ❖ VED faces vast theoretical problems, “the CC problem” regarding the $\sim m^4$ terms.

$$\rho_{\text{vac}}^{\text{obs}} / \rho_{\text{ZPE}} \sim \rho_{\text{vac}}^{\text{obs}} / m_e^4 \sim 10^{-34}. \quad (1)$$

and “the coincidence problem”,

$$\rho_{\text{vac}}^0 \sim \rho_{\text{m}}^0. \quad (2)$$

Running Vacuum Models

Running Vacuum Models

- Family of parametrizations of VED with long trajectory in literature. The Canonical RVM is

$$\rho_{\text{vac}}(H) = \frac{3}{8\pi G_N} (C_0 + \nu H^2), \quad (3)$$

where $|\nu| \ll 1$.

- Similar models have been explored.¹
- In particular, the RRVM was recently tested

$$\rho_{\text{vac}}(R) = \frac{3}{8\pi G_N} \left(c_0 + \frac{\nu}{12} R \right), \quad (4)$$

with $R = 12H^2 + 6\dot{H}$ is the Ricci scalar.

¹J. Solà Peracaula *et al* 1602.02103, 1703.08218, 1705.06723, 2102.12758.

Running Vacuum Models

- For type I models with threshold, $\rho_m \approx \rho_m^0 a^{-3+3\frac{\nu}{4}}$,

$$\rho_{\text{vac}}(a) = \begin{cases} \rho_{\text{vac}}^0 + \frac{\nu}{4}\rho_m^0(a^{-3} - 1) & , \quad a > a_{\text{th}}, \\ \rho_{\text{vac}}(a_{\text{th}}) & , \quad a < a_{\text{th}}. \end{cases} \quad (5)$$

Recent activation of the dynamics of DE.

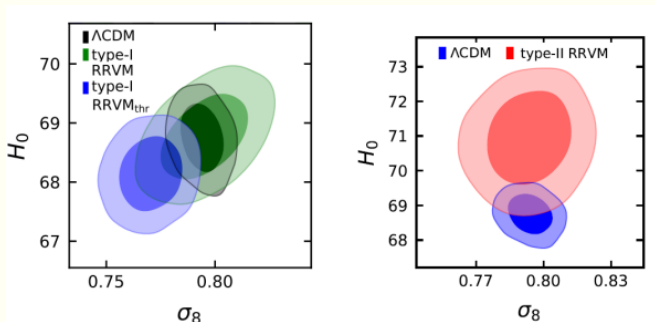
- For type II, $G(a) \sim 1 + \epsilon \ln a$, with $\epsilon \sim \mathcal{O}(\nu)$, and

$$\rho_{\text{vac}}(a) \approx \frac{3c_0}{8\pi G_N}(1 + 4\nu) + \nu\rho_m^0 a^{-3}, \quad (6)$$

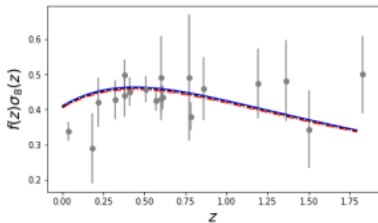
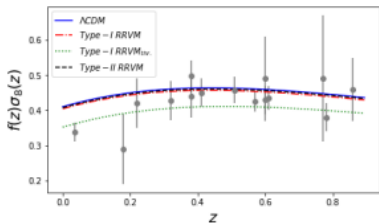
The initial value of G considered a free parameter.

Running Vacuum Models

Since $\dot{\rho}_{\text{vac}} \neq 0$ the perturbation equations are affected producing a departure from Λ CDM. Confronted against $SnIa + BAO + H(z) + f\sigma_8(z) + Planck$ data.



Running Vacuum Models



- ❖ The Type I model with threshold has a bigger impact on solving the σ_8 tension.
- ❖ The Type II model seems to be able to alleviate the H_0 tension without altering the σ_8 one.
- ❖ There were no strong theoretical grounds from QFT or quantum gravity supporting them.

Renormalization of Vacuum Energy

Renormalization of Vacuum Energy

- ❖ We start with EH action + ρ_Λ . For simplicity, in the matter sector only a scalar field ϕ ,

$$S = - \int dx^4 \sqrt{-g} \left(\frac{1}{2} g^{\mu\nu} \partial_\nu \phi \partial_\mu \phi + \frac{1}{2} (m^2 + \xi R) \phi^2 \right) \quad (7)$$

where ξ is a non-minimal coupling.

- ❖ With a flat FLRW background, $g_{\mu\nu} = a^2 \text{diag}(-1, 1, 1, 1)$.
- ❖ Splitting the field in a background part and in a fluctuating part

$$\phi(\tau, \mathbf{x}) = \phi(\tau) + \delta\phi(\tau, \mathbf{x}). \quad (8)$$

Renormalization of Vacuum Energy

- ❖ The Fourier decomposition in modes is

$$\delta\phi(\tau, \mathbf{x}) = \frac{1}{(2\pi)^{(3/2)} a} \int dk^3 \left[A_k e^{i\mathbf{k}\mathbf{x}} h_k(t) + A_k^\dagger e^{-i\mathbf{k}\mathbf{x}} h_k^*(t) \right], \quad (9)$$

with usual commutation relations.

- ❖ KG eq: $h_k'' + \Omega_k^2 h_k = 0$, $\Omega_k \equiv k^2 + a^2 m^2 + a^2 (\xi - 1/6) R$
- ❖ Traditional solution: WKB ansatz and recursive iterations,

$$h_k(\tau) \sim W_k^{-1/2}(\tau) e^{-i \int_1^\tau W_k(\tau_1) d\tau_1}, \quad (10)$$

such that $W_k^2 = \Omega_k^2 - \frac{1}{2} \frac{W_k''}{W_k} + \frac{3}{4} \left(\frac{W_k'}{W_k} \right)^2$

Renormalization of Vacuum Energy

- ❖ Background evolving slowly. Solution is organized in what we call adiabatic orders \equiv number of time derivatives.
- ❖ We can use conjecture the renormalized Zero-Point Energy:

$$\begin{aligned}\langle T_{00}^{\delta\phi} \rangle_{\text{ren}}(M) &\equiv \langle T_{00}^{\delta\phi} \rangle(m) - \langle T_{00}^{\delta\phi} \rangle^{(0-4)}(M) \\ &= \frac{a^2}{128\pi^2} \left(-M^4 + 4m^2M^2 - 3m^4 + 2m^4 \log \frac{m^2}{M^2} \right) \\ &\quad - \frac{3a^2H^2 \left(\xi - \frac{1}{6} \right)}{16\pi^2} \left(m^2 - M^2 - m^2 \log \frac{m^2}{M^2} \right) + \mathcal{O}\left(H^4, \frac{H^6}{m^2}, \dots\right)\end{aligned}\tag{11}$$

M is an arbitrary off-shell mass scale.

Renormalization of Vacuum Energy

- ❖ We define

$$\langle T_{\mu\nu}^{\text{vac}} \rangle(M) \equiv -\rho_{\Lambda}(M)g_{\mu\nu} + \langle T_{\mu\nu}^{\delta\phi} \rangle_{\text{ren}}(M) \quad (12)$$

In plain words “ $VED = ZPE + \Lambda$ ”,

$$\rho_{\text{vac}}(M) = \frac{\langle T_{00}^{\delta\phi} \rangle_{\text{ren}}(M)}{a^2} + \rho_{\Lambda}(M). \quad (13)$$

- ❖ Set $M = H$ and do the subtraction at two different scales,

$$\rho_{\text{vac}}(H) = \rho_{\text{vac}}(H_0) + \frac{3\nu_{\text{eff}}}{8\pi} (H^2 - H_0^2) m_{\text{Pl}}^2 + \mathcal{O}(H^4), \quad (14)$$

where $\nu_{\text{eff}} \equiv \frac{(\xi - \frac{1}{6})}{2\pi} \frac{m^2}{m_{\text{Pl}}^2} \ln \frac{m^2}{H_0^2}$ is expected to be small.

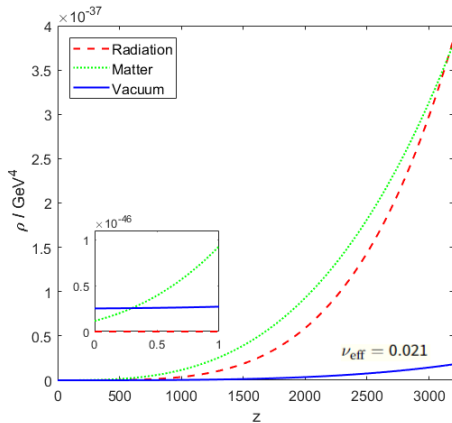
Chameleonic VED

Chameleonic VED

- ❖ Mild dynamics: Two values, at H_1, H_2 , of the vacuum energy in the late times are smoothly related through $\sim \nu_{\text{eff}} m_{\text{Pl}}^2 (H_1^2 - H_2^2)$.
- ❖ Running free from $\sim m^4$ terms.
- ❖ Gravitational constant is also shown to be running (logarithmically),

$$G(H) \approx \frac{G(H_0)}{1 - \nu_{\text{eff}} \frac{\ln H^2/H_0^2}{\ln m^2/H_0^2}}. \quad (15)$$

Chameleonic VED



Evolution of the energy densities for $\nu_{\text{eff}} = 0.02$, $\Omega_m^0 = 0.32$, $\Omega_r^0 = 0.0001$.

Chameleonic VED

- Analogously one may compute the associated vacuum pressure. One can infer the equation of state:

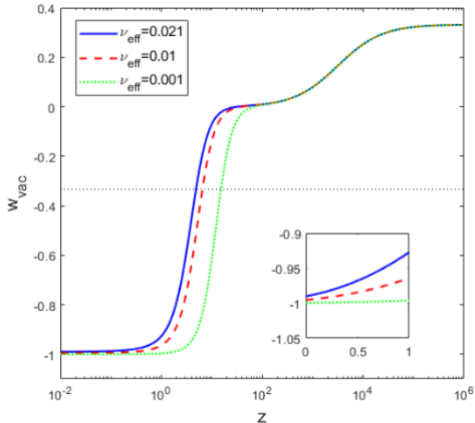
$$1 + w_{\text{vac}}(z) \approx \frac{\nu_{\text{eff}} (\Omega_m^0 (1+z)^3 + \frac{4}{3} \Omega_r^0 (1+z)^4)}{\Omega_{\text{vac}}^0 + \nu_{\text{eff}} [-1 + \Omega_m^0 (1+z)^3 + \Omega_r^0 (1+z)^4 + \Omega_{\text{vac}}^0]}. \quad (16)$$

- Vacuum Energy has a chameleonic EoS, mimicking the dominant component.
- It behaves as Dark Energy ($w_{\text{vac}} < -0.33$) for low redshifts,

$$w_{\text{vac}}(z) = -1 + \nu_{\text{eff}} \frac{\Omega_m^0}{\Omega_{\text{vac}}^0} (1+z)^3, \quad (17)$$

Eventually $w_{\text{vac}} = -1$ if $z \rightarrow -1$.

Chameleonic VED



The VED has a quintessence behaviour for $\nu_{\text{eff}} > 0$.

Conclusions and Future works

Conclusions and Future works

- ❖ We presented a computation of VED from QFT in curved spacetime.
- ❖ VED is mildly dynamical in the late universe. G is expected to vary logarithmically.
- ❖ The EoS for vacuum is not -1 along the whole story, mimics the dominant component and behaves as quintessence DE in the late universe.
- ❖ The higher powers in the adiabatic expansion may have an interesting role in the primeval era.
- ❖ All in all, the extra features of the VED may alleviate the σ_8 and H_0 tension, but should be contrasted with data.

THANKS!!



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