Observational Tensions in Kinetically Coupled Quintessence

Based on arXiv:2207.13682

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In collaboration with Bruno J. Barros, Vasco M. C. Ferreira and Noemi Frusciante

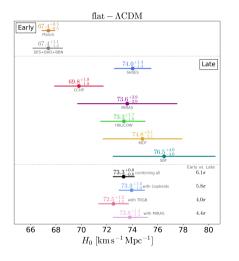




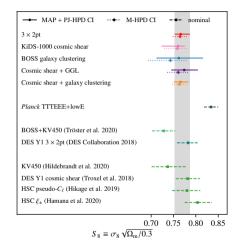




Cosmological Tensions



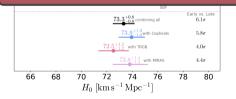
[Verde, Treu, Riess: Nature Astron. 3 891 (2019)]



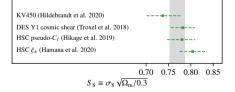
[C. Heymans et al.: Astron.Astrophys. 646 (2021)]



Interaction between Dark Energy and Dark Matter







[C. Heymans et al.: Astron.Astrophys. 646 (2021)]

Kinetic Model

Interacting Dark Energy with Kinetic Coupling

$$S = \int \mathrm{d}^4 x \sqrt{-g} \left[\frac{\mathsf{M}_{\mathrm{Pl}}^2}{2} R + X - V(\phi) + f(\phi, X) \tilde{\mathcal{L}}_c + \mathcal{L}_{\mathrm{SM}}(\psi_i, g_{\mu\nu}) \right]$$

[Barros: Phys. Rev. D 99, 064051 (2019)]

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- Function $f(\phi, X)$ how the field couples to the matter species
- Purely kinetic power-law coupling given by $f(X) = \left(M_{\text{Pl}}^{-4} X\right)^{\alpha}$ with $V(\phi) = V_0 e^{-\lambda \phi/M_{\text{Pl}}}$
- Equivalent to a conformally coupled theory with $\tilde{g}_{\mu\nu} = f^2(X)g_{\mu\nu}$

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In an <u>FLRW Universe</u> this gives

$$3M_{Pl}^{2}\mathcal{H}^{2} = a^{2}(\rho_{c} + \rho_{b} + \rho_{r} + \rho_{\phi})$$

$$\phi'' + 2\mathcal{H}\phi' + a^{2}V_{,\phi} = a^{2}Q$$

$$\rho_{c}' + 3\mathcal{H}\rho_{c} = -Q\phi', \quad \rho_{\phi}' + 3\mathcal{H}(1 + w_{\phi})\rho_{\phi} = Q\phi'$$

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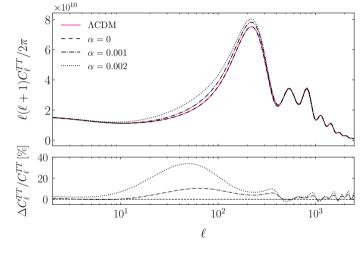
In an FLRW Universe this gives

$$\begin{cases} 3M_{\rm Pl}^2 \mathcal{H}^2 = a^2(\rho_c + \rho_b + \rho_r + \rho_{\phi}) \\ \phi'' + 2\mathcal{H}\phi' + a^2 V_{,\phi} = a^2 Q \\ \rho'_c + 3\mathcal{H}\rho_c = -Q\phi', \quad \rho'_{\phi} + 3\mathcal{H}(1 + w_{\phi})\rho_{\phi} = Q\phi' \end{cases}$$

Coupling function regulated by parameter α

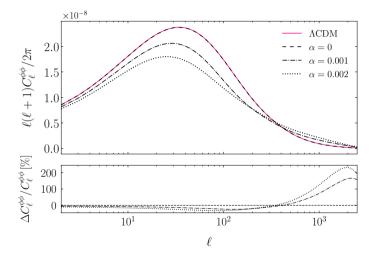
$$Q = 2\alpha \rho_c \frac{3\mathcal{H}\phi' + a^2 V_{,\phi}}{2\alpha a^2 \rho_c + (1+2\alpha) \phi'^2}$$

Cosmological observables

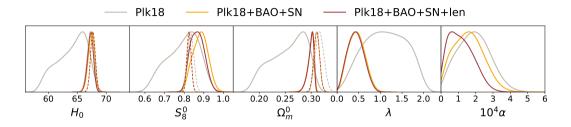


• General enhancement of TT power spectrum (ISW effect) - degeneracy between H_0 and α

Cosmological observables



• Overall suppression of the lensing power spectrum with increasing α



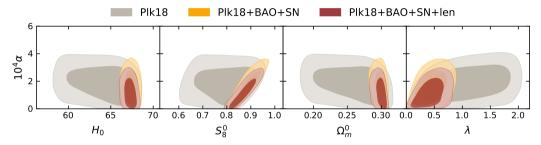
• Lower mean values for H_0 and Ω_m^0 in the Kinetic model

• α constrained to be of order 10⁻⁴ for all data combinations (all compatible at 68% C.L.)

Planck data: prefers the higher mean value of α - degeneracies

BAO and SN data: slight decrease of the *best-fit* value of α

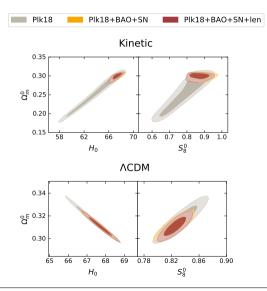
CMB lensing data: even lower central value of α - lensing excess



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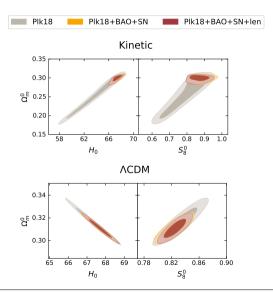
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- CMB lensing data: even lower central value of α lensing excess

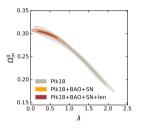
Cosmological Tensions



- **Opposite correlation** between H_0 and Ω_m^0 (positive) counterbalance the enhancement of the TT power spectrum for $\alpha \neq 0$
- Positive correlation between S_8^0 and Ω_m^0
- Planck data: $S_8^0 = 0.793^{+0.11}_{-0.064}$ for the kinetic model and $S_8^0 = 0.833 \pm 0.016$ for the standard scenario apparent easing of the S_8 tension
- Degeneracy broken by BAO and SN data -Hubble tension is also still present for all combinations

Cosmological Tensions





- Inclusion of BAO and SN data leads to narrower constraints on Ω_m^0 , and consequently on H_0 , S_8^0 and λ as well
- Anti-correlation between Ω_m^0 and λ : $\Omega_{\phi}^0 \approx 1 - \Omega_m^0$ and $w_{\phi} \approx 1 - 2V/3H^2$ cosmological constant like behaviour today pushes λ towards smaller values

	Plk18	Plk18 + BAO + SN	Plk18 + BAO + SN + len
$\Delta \chi^2_{ m eff}$	-0.9	0.7	1.0
ΔDIC	-0.316	0.7	1.62

(e) $\Delta \chi^2_{eff}$ to assess the favoured model and DIC to quantify the preference

Better fit to the data for the kinetic model for the Planck data, but not overly significant

Slight preference for ΛCDM for other data combinations

BAO and SN data change the fit to the TT likelihood and the CMB lensing data shows an excess of power - suppression for lpha
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Conclusions

- Interaction in the dark sector through a purely kinetic coupling function $f(X) = (M_{Pl}^{-4} X)^{\alpha}$
- Cosmological constraints on the parameters of the theory using CMB, CMB lensing, BAO and SN data
- The S_8 tension is apparently alleviated, while the H_0 tension is still present for all combinations
- The parameter α is consistently constrained to be of the order 10^{-4} but non-vanishing at 68% C.L.
- No clearly statistically favoured model between ΛCDM and the kinetic model
- It would still be of interest to consider the kinetic models for further investigations



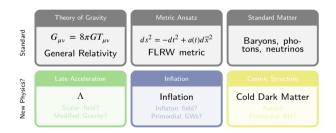
[Illustration Credits: Inês Viegas Oliveira]

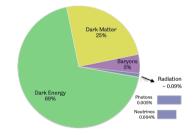
Thank you! Do you have any questions?

Illustration Credits: Inês Viegas Oliveira (ivoliveira.com)

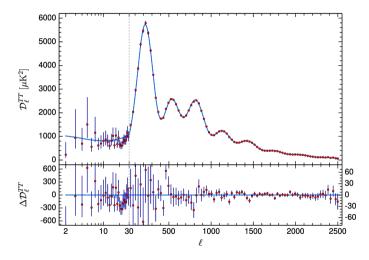
Standard Model of Cosmology

- Overwhelming observational evidence for the accelerated expansion of the Universe (CMB, SNe, BAO, ...)
- ACDM model is the standard model of Cosmology





[Ezquiaga and Zumalacárregui: Front.Astron.Space Sci. 5 (2018), 44]

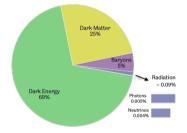


[Aghanim et al.: Astron.Astrophys. 641 (2020) A6]

Standard Λ CDM Model

- Overwhelming observational evidence for the accelerated expansion of the Universe (CMB, SNe, BAO, ...)
- Standard model of Cosmology based on 6 ingredients
- Theoretical simplicity of the unknown components and fewer number of free parameters
- ΛCDM correctly predicts the cosmic background radiation and the large-scale distribution of galaxies
- However: conceptual problems and observational tensions

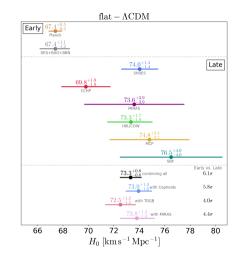
density remains constant with the expansion and has negative pressure ($w = p_{\lambda}/\rho_{\Lambda} = -1$)



Hubble Tension

- 4.4 σ tension between *Planck* 2018 and SH0ES:
 - CMB (Planck): $H_0 = 67.4 \pm 0.5 \text{ km/s/Mpc}$
 - SNe (SH0ES): $H_0 = 74.0 \pm 1.4 \text{ km/s/Mpc}$
- The *Planck* 2018 results are a grand confirmation of the Λ CDM model but they are model dependent
- Unlikely that the discrepancies could be explained by a single systematic error
- The magnitude and persistence hints at standard model flaws

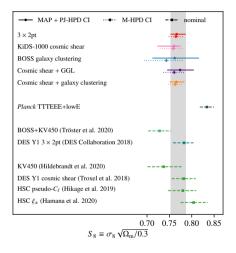
[Di Valentino et al.: arXiv:2008.11284]



[Verde, Treu, Riess: Nature Astron. 3 891 (2019)]

S_8 Tension

- S_8 Tension: discrepancy between CMB data and weak lensing and redshift surveys on the combined value of Ω_m^0 and σ_8^0 , expressed as $S_8^0 = \sigma_8 \sqrt{\Omega_m/0.3}$
- 3σ tension between *Planck* 2018 CMB data and KiDS-1000 combination of Cosmic Shear and Galaxy Clustering:
 - CMB (Planck): $S_8^0 = 0.834 \pm 0.016$
 - **CS+GC** (KiDS-1000): $S_8^0 = 0.766^{+0.020}_{-0.014}$
- Could be related to the excess of lensing measured by Planck, mimicking a larger S_8^0
- $\ensuremath{{ \bullet }}$ Correlation between the H_0 and S_8 tensions conjoined analysis
- Formulate extensions to the standard cosmological framework and test against the relevant constraints [Di Valentino et al.: arXiv:2008.11285]



[C. Heymans et al.: Astron.Astrophys. 646 (2021)]

"Quintessence" (ϕ) - dynamical scalar field that evolves in space and time, as opposed to Λ

New forces between DE and "normal matter" are heavily constrained by observations (e.g. solar system tests)

No fundamental principle which forbids interactions between the dark species

Modification to the conservation equations in FLRW Univese:

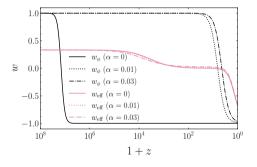
 $\dot{\rho}_{DM} + 3H\rho_{DM} = Q$ $\dot{\rho}_{\phi} + 3H (1 + w_{\phi}) \rho_{\phi} = -Q$

Changes in the background and linear perturbations evolution could naturally address the H_0 and S_8 tensions

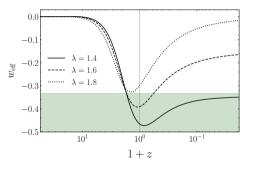
But how to achieve the phenomenological coupling?



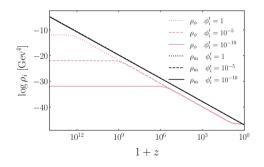
- EoS parameters w_φ and w_{eff}: transition towards accelerating state occurs later for larger values of α
- ${\ensuremath{\bullet}}$ For a stronger interaction, the field remains frozen in the scaling regime for longer, with $w_{\phi}=1$



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- EoS parameters w_φ and w_{eff}: transition towards accelerating state occurs later for larger values of α
- \circledast For a stronger interaction, the field remains frozen in the scaling regime for longer, with $w_{\phi}=1$
- Increasing the value of λ decreases the chances of achieving accelerated expansion even if just transient
- Increasing ϕ'_i leads to an earlier onset of the scaling regime only, setting the duration of the scaling no influence whatsoever on the dynamics



Linear Perturbations

Scalar perturbations in the conformal Newtonian gauge

$$ds^{2} = a^{2}(\tau) \left[-(1 + 2\Psi) d\tau^{2} + (1 - 2\Phi) \delta_{ij} dx^{i} dx^{j} \right]$$

Perturbed continuity and Euler equations for CDM ($\delta_c = \delta \rho_c / \rho_c$ and $\theta_c = \partial_i \partial^i v_c$)

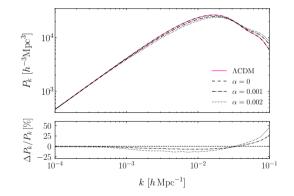
$$\begin{cases} \delta_c' + \theta_c - 3\Phi' = \frac{\varrho}{\rho_c} \left(\phi' \delta_c - \delta \phi' \right) - \frac{\phi'}{\rho_c} \delta Q \\ \theta_c' + \mathcal{H} \theta_c - k^2 \Psi = \frac{\varrho}{\rho_c} \left(\phi' \theta_c - k^2 \delta \phi \right) \end{cases}$$

These equations describe the clustering of matter that leads to structure formation

$$\delta\phi'' + 2\mathcal{H}\delta\phi' + \left(a^2V_{,\phi\phi} + k^2\right)\delta\phi - \left(\Psi' + 3\Phi'\right)\phi' + 2a^2\Psi V_{,\phi} = a^2\delta Q + 2a^2Q\Psi$$

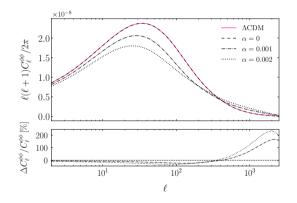
$$\delta Q = \frac{2\alpha \rho_c \left\{ -3\Phi' \phi' - \phi' \theta_c + \left[3\mathcal{H} \phi' + a^2 (V_{,\phi} - Q) \right] \delta_c + \left(2k^2 + a^2 V_{,\phi\phi} \right) \delta \phi - \left[3\mathcal{H} \phi' + 2a^2 (V_{,\phi} - Q) \right] \frac{\delta \phi'}{\phi'} + 2a^2 \Psi \left(Q - V_{,\phi} \right) \right\}}{2\alpha a^2 \rho_c + (1 + 2\alpha) \phi'^2}$$

- Matter power spectrum is significantly suppressed at intermediate scales, $10^{-3} < k < 3 \times 10^{-2}$ and enhanced at the smaller scales, $k > 3 \times 10^{-2}$
- The growth of the matter perturbations leads to a larger σ_8 for the kinetic model



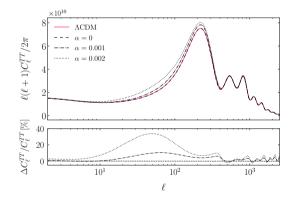
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Cosmological observables

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- () The growth of the matter perturbations leads to a larger σ_8 for the kinetic model
- Overall suppression of the lensing power spectrum itself with increasing α
- General enhancement of TT power spectrum with greater amplitude and broadness of the first peak - ISW effect



Given a data set d, we want to sample posteriors on the model parameters θ that maximise the likelihood:



Modified version of Einstein-Boltzmann code CLASS interfaced with the MontePython sampler [Blas, Lesgourgues, Tram: JCAP 1107 (2011) 034; Audren et al.: JCAP 1302 (2013) 001; Brinckmann, Lesgourgues: Phys. Dark Univ. 24 (2019) 100260]

Employ an MCMC sampling method and analyse results in GetDist [Lewis: arXiv:2008.11284]



The Λ CDM model is based on 6 free parameters:

- \odot the baryon and dark matter densities $\Omega_b h^2$ and $\Omega_c h^2$
- () the angular size of the sound horizon at decoupling θ_s
- () the reionisation redshift z_{reio}
- the spectral index n_s and the amplitude A_s of inflationary scalar perturbations

In the Kinetic Model scenario we also allow sampling of:

- ${\ensuremath{\, \bullet }}$ the effective coupling parameter α and the steepness of the potential λ
- \implies 2 additional parameters

The remaining cosmological parameters are either fixed to standard Planck 2018 values or derived from the main ones

Parameter	Prior
$\Omega_{ m b} h^2$	[0.005, 0.1]
$\Omega_{ m c} h^2$	[0.001, 0.99]
$100\theta_s$	[0.5, 10]
z _{reio}	[0., 20.]
n _s	[0.7, 1.3]
$\log\left(10^{10}A_{s} ight)$	[1.7, 5.0]

Parameter	Prior
α	[0, 1]
λ	[0,2]

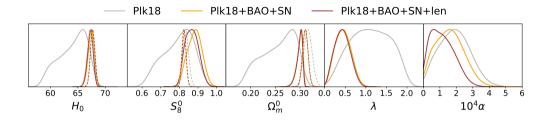
Data Sets

- Baseline data set is "Plk18": CMB Planck 2018 data for large angular scales ℓ = [2, 29] and a joint of TT, TE and EE likelihoods for the small angular scales [Aghanim et al.: Astron.Astrophys. 641 (2020) A5]
- "Plk18+BAO+SN": "Plk18" plus compilation of baryon acoustic oscillations (BAO) distance and expansion rate measurements and distance moduli measurements of type Ia Supernova (SN) data from Pantheon.

[Ross et. al: Mon. Not. Roy. Astron. Soc. 449 (2015) 835; Beutler et al.:
 Mon. Not. Roy. Astron. Soc. 464 (2017) 3409; Beutler et al.: Mon. Not. Roy.
 Astron. Soc. 416 (2011) 3017; Scolnic et. al: Astrophys. J. 859 (2018) 101]

• "PIk18+BAO+SN+len": "PIk18+BAO+SN" plus CMB lensing potential data from Planck 2018 [Aghanim et al.: Astron.Astrophys. 641 (2020) A8]

Cosmological Bounds

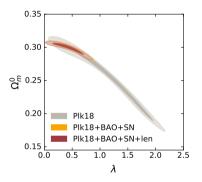


Kinetic Model			
Parameter	Plk18	Plk18+	Plk18+
- arameter		BAO+SN	BAO+SN+Ien
S_{8}^{0}	$0.793^{+0.11}_{-0.064}$	$0.875^{+0.037}_{-0.043}$	$0.863^{+0.030}_{-0.039}$
Ω_m^0	$0.257^{+0.045}_{-0.025}$	$0.2988^{+0.0072}_{-0.0036}$	$0.2982^{+0.0070}_{-0.0035}$
H_0	$64.0^{+3.3}_{-1.8}$	67.14 ± 0.62	$66.94^{+0.60}_{-0.54}$
λ	1.11 ± 0.48	$0.42^{+0.18}_{-0.21}$	$0.41^{+0.17}_{-0.22}$
$10^4 \alpha$	1.88 ± 0.95	$1.37^{+0.67}_{-1.0}$	$1.05^{+0.51}_{-0.87}$

ACDM Model				
Parameter	Plk18	Plk18+	Plk18+	
		BAO+SN	BAO+SN+len	
S_{8}^{0}	0.833 ± 0.016	$0.831^{+0.013}_{-0.015}$	0.834 ± 0.013	
$\Omega_{\rm m}^0$	0.3163 ± 0.0085	$0.3151^{+0.0060}_{-0.0075}$	0.3162 ± 0.0073	
H_0	67.31 ± 0.61	$67.39^{+0.53}_{-0.45}$	67.32 ± 0.53	

Cosmological Bounds

- \circledast Inclusion of BAO and SN data leads to narrower constraints on Ω^0_m , and consequently on $H_0,~S^0_8$ and λ as well
- Anti-correlation between Ω_m^0 and λ : $\Omega_{\phi}^0 \approx 1 \Omega_m^0$ and $w_{\phi} \approx 1 - 2V/3H^2$ - cosmological constant like behaviour today pushes λ towards smaller values
- Better fit to the data for the kinetic model considering the Planck data, although not overly significant
- Slight preference for Λ CDM for other data combinations inclusion of BAO and SN data changes the fit to the TT likelihood and the CMB lensing data shows an excess of power, while $\alpha \neq 0$ introduces an overall suppression
- There is no statistical evidence in support for either of the two models



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