

Testing a key foundation of the concordance model

Roy Maartens



Workshop on Tensions in Cosmology

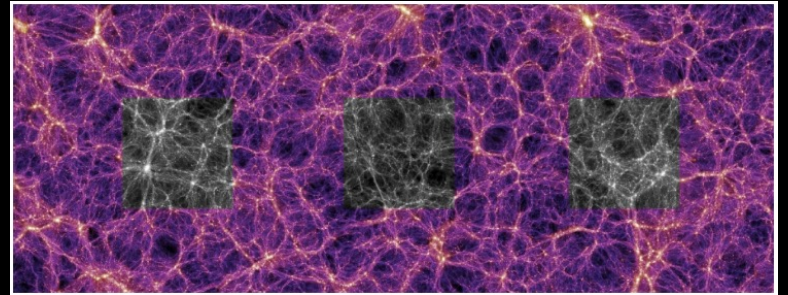
SEPTEMBER 7-12, 2022

The Cosmological Principle

The Universe is statistically isotropic and homogeneous

A critical foundation stone of LCDM.

The Universe is I+H on *average*,
on large enough scales.

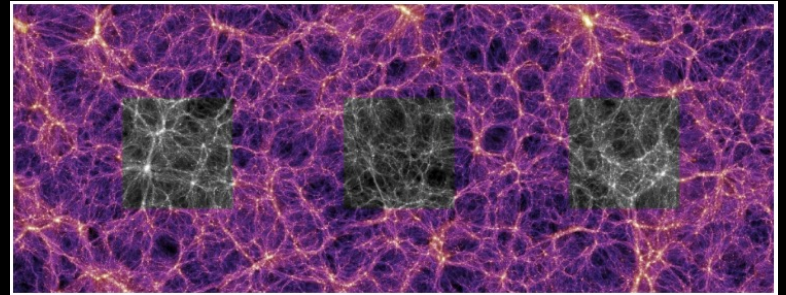


The Cosmological Principle

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on large enough scales.



Key point:

The CP implies a unique frame – or 4-velocity field u^a –
in which average isotropy and homogeneity holds

- All 'fundamental' observers u^a see isotropy + homogeneity.
- Any observer with 4-velocity different from u^a does not see I+H.

Testing the consistency of matter and radiation

For practical purposes – we assume that the CP holds and apply consistency tests to a perturbed FLRW model.

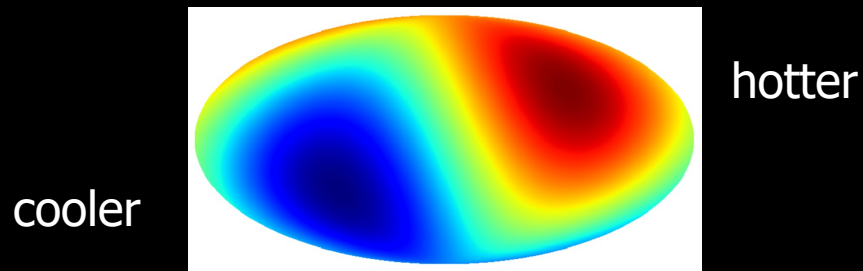
A key test:

isotropy in radiation and in matter should be consistent

Such a test was proposed by Ellis & Baldwin (1984).

Heliocentric observers are moving relative to the CMB rest-frame.

This generates a dipole in the CMB temperature –

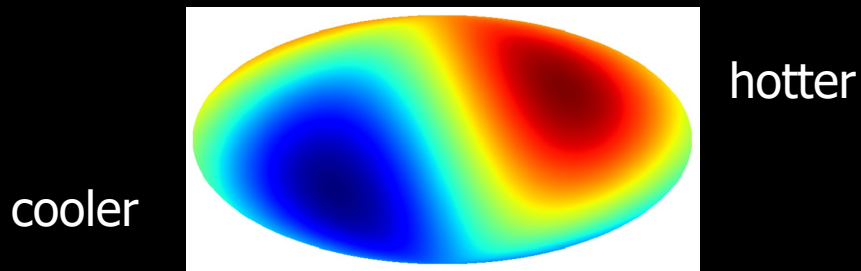


$$\begin{aligned} \tilde{T}(\tilde{\mathbf{n}}) &= T(\mathbf{n})(1 + \mathbf{n} \cdot \mathbf{v}_o) \quad v_o \approx 10^{-3} \\ \Rightarrow \tilde{\delta}_T(\mathbf{n}) &= \delta_T(\mathbf{n}) + \mathbf{n} \cdot \mathbf{v}_o \end{aligned}$$

at first order in perturbations.

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$$\Rightarrow \tilde{\delta}_T(\mathbf{n}) = \delta_T(\mathbf{n}) + \mathbf{n} \cdot \mathbf{v}_o$$

at first order in perturbations.

If the Universe is isotropic about us on average, then

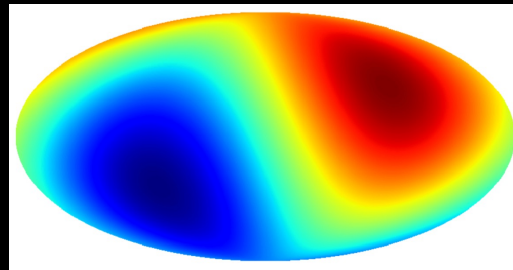
galaxy rest-frame = CMB rest-frame

$$\mathbf{v}_o|_{\text{gal}} = \mathbf{v}_o|_{\text{CMB}} \quad (\text{magnitude} + \text{direction})$$

– a critical test of the Cosmological Principle (Ellis & Baldwin 1984)

In other words, the same dipole should be seen in number counts:
highest counts in the direction \mathbf{v}_o , lowest counts in direction $-\mathbf{v}_o$

less galaxies



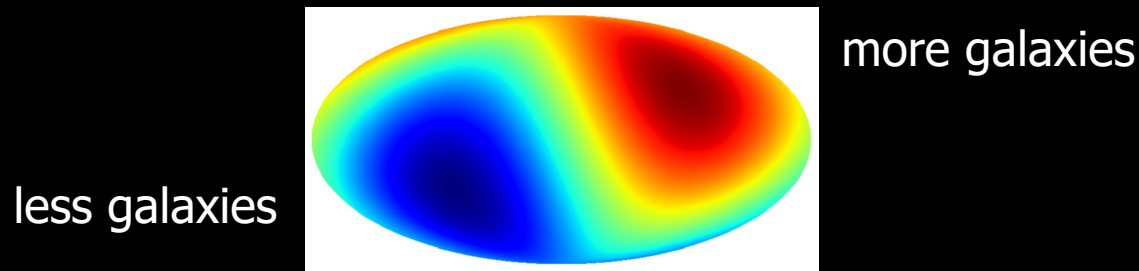
more galaxies

$$\tilde{\delta}_N(\mathbf{n}) = \delta_N(\mathbf{n}) + [2 + x(1 + \alpha)] \mathbf{n} \cdot \mathbf{v}_o \quad x, \alpha = \text{constant}$$

Ellis-Baldwin

Need surveys with large sky area and high numbers.

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Ellis-Baldwin

Need surveys with large sky area and high numbers.

Early tests with NVSS survey (JVLA telescope), e.g.,

Blake & Wall 2002; Singal 2011; Gibelyou & Huterer 2012;
Rubart & Schwarz 2013; Tiwari & Jain 2015; Colin et al 2017

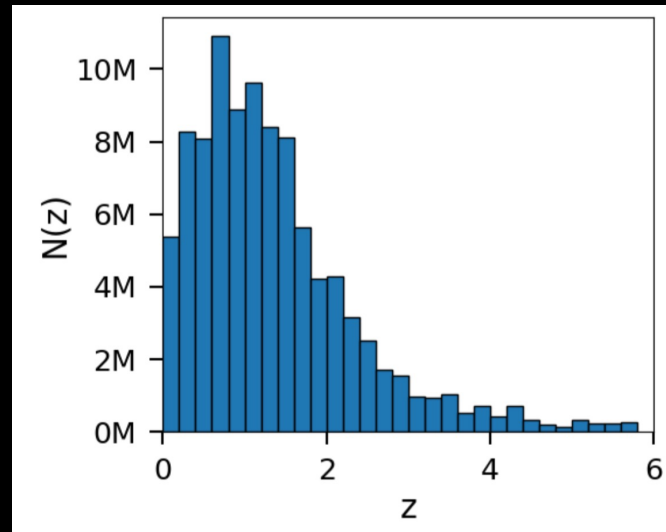
generally found consistency, with dipole magnitude $>$ theory.

A key issue is systematics on ultra-large scales – very difficult.

Also – to measure the LSS dipole, we must remove low redshift sources to avoid **nonlinear contamination** of the dipole.

We need a sample with many high- z sources – like SKA:

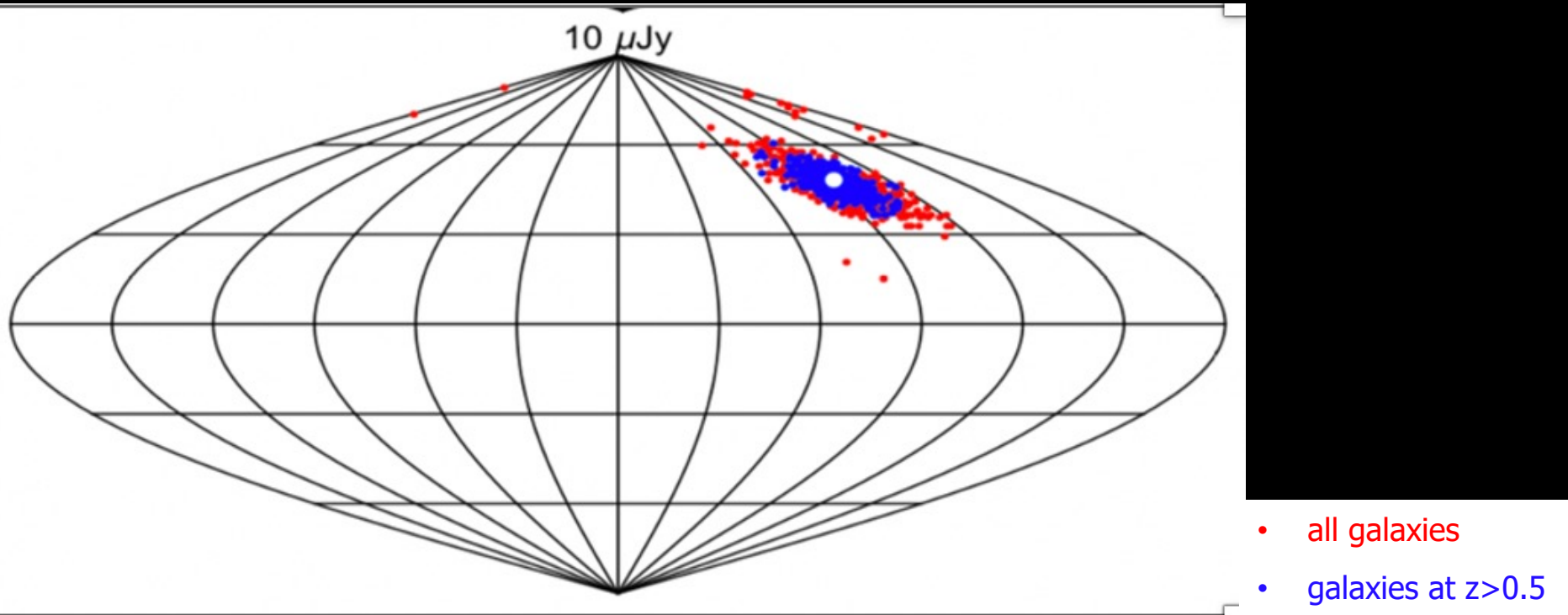
Simulations



(SKA Redbook, 2018)

To exclude low redshift radio galaxies – we can use redshift information from galaxy spectro- z and photo- z surveys.

Simulations – the dipole for SKA is well above the noise:



(Bengaly, Siewert, Schwarz, RM 2018)

Removing galaxies at $z < 0.5$ significantly improves the measurement.

SKA could measure the dipole with 5 – 10% error, giving a robust test of the Cosmological Principle.

Redshift-dependent dipole in galaxy redshift surveys

Radio continuum surveys detect galaxies by their radio emission – with no redshifts.

Number counts are projected on the 2D sky.

Redshift surveys in 3D lead to z -dependent dipole magnitude.

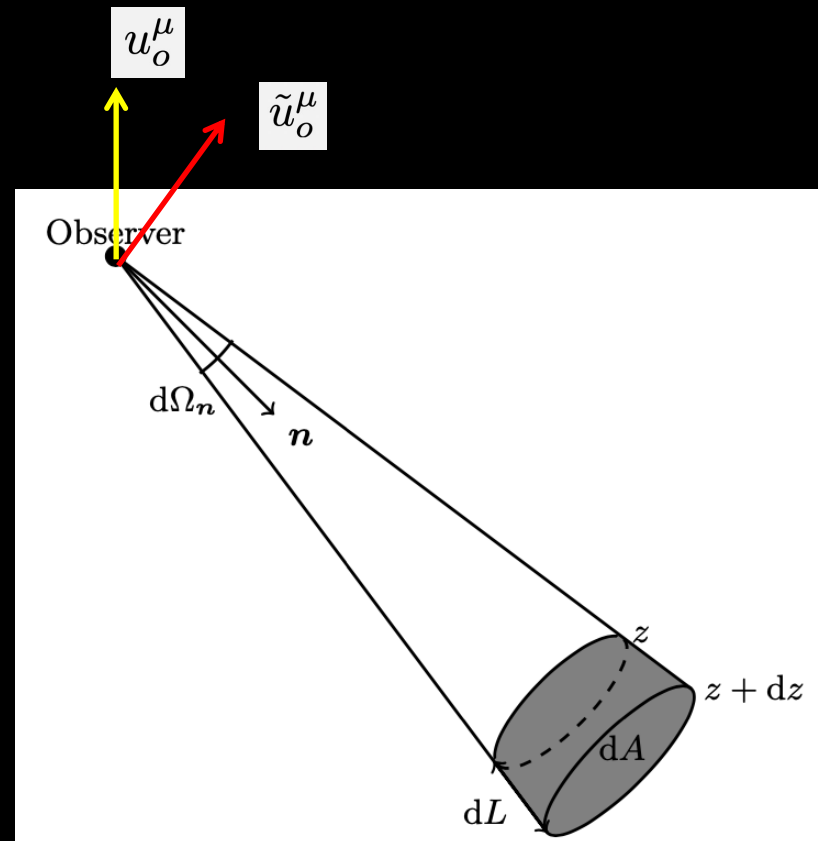
Redshift-dependent dipole in galaxy redshift surveys

Boosted observer 4-velocity:

$$\tilde{u}_o^\mu = \gamma(v_o) [u_o^\mu + v_o^\mu] = u_o^\mu + v_o^\mu + O(v_o^2) \quad \text{where} \quad u_o^\mu v_{o\mu} = 0$$

boosted
(heliocentric)

CMB rest-frame



The boosted observer measures redshifts and directions:

$$1 + \tilde{z} = (1 + z)(1 - \mathbf{n} \cdot \mathbf{v}_o) \quad \textit{Doppler boost}$$

$$\tilde{\mathbf{n}} = (1 - \mathbf{n} \cdot \mathbf{v}_o)\mathbf{n} + \mathbf{v}_o \quad \textit{aberration}$$

Total number of particles is conserved:

$$\tilde{N} d\tilde{z} d\tilde{\Omega}_{\tilde{\mathbf{n}}} = N dz d\Omega_{\mathbf{n}}$$

Then the observed number per redshift per solid angle is

$$\tilde{N}(\tilde{z}, \tilde{\mathbf{n}}) = N(z, \mathbf{n}) [1 + 3\mathbf{n} \cdot \mathbf{v}_o]$$

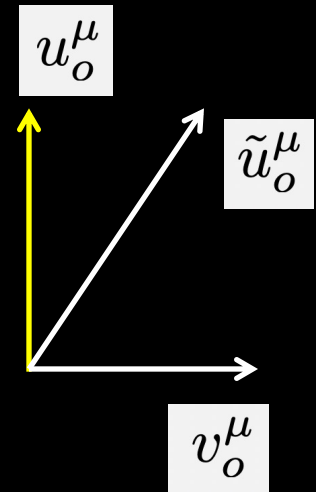
But – we must account for redshift and luminosity perturbations.

This generates a dipole in the observed number density contrast:

$$\delta_{\tilde{N}} = \delta_N + \delta_N^{\mathbf{v}_o} \quad \text{where} \quad \delta_N^{\mathbf{v}_o} = \mathcal{D}(z, m_*) \mathbf{n} \cdot \hat{\mathbf{v}}_o$$

For galaxies (RM, Clarkson, Chen 2017)

$$\mathcal{D}_{\text{gal}} = \left[3 + \frac{\dot{H}}{H^2} + (2 - 5s) \frac{(1+z)}{rH} - b_e \right] v_o$$



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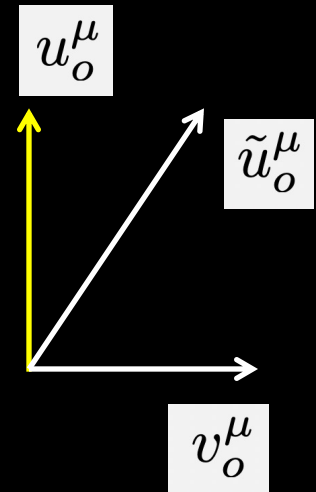
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evolution
bias

$$b_e(\bar{z}, m_*) = - \frac{\partial \ln [(1 + \bar{z})^{-3} \bar{\mathcal{N}}_s(\bar{z}, m < m_*)]}{\partial \ln(1 + \bar{z})}$$

magnification
bias

$$s(\bar{z}, m_*) = \frac{\partial \log \bar{\mathcal{N}}_s(\bar{z}, m < m_*)}{\partial m_*}$$



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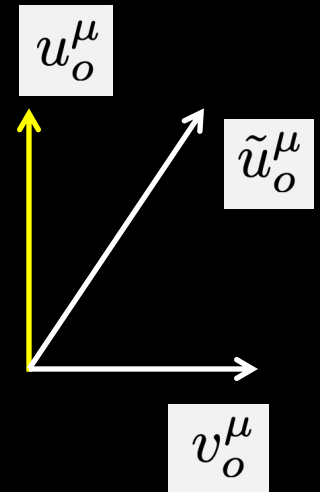
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luminosity
function

$$\bar{\mathcal{N}}_s(\bar{z}, m < m_*) = \frac{2}{5} \ln 10 \int_{-\infty}^{m_*} dm \bar{n}_s(\bar{z}, m) \quad \text{proper no. density (at source)}$$

where

$$\bar{N} = \left(\frac{r^2}{H} \right) (1+z)^{-3} \bar{\mathcal{N}}_s \quad \text{observed density related proper density}$$



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NB This is first-order – does not include **nonlinearities** that arise at low redshift

The projected 2D dipole

The 3D galaxy redshift dipole :

$$\mathcal{D}_{\text{gal}} = \left[3 + \frac{\dot{H}}{H^2} + (2 - 5s) \frac{(1+z)}{rH} - b_e \right] v_o$$

Do we recover the Ellis-Baldwin result for a 2D dipole

$$\langle \mathcal{D}_{\text{gal}} \rangle_{\text{EB}} = \left[2 + x(1 + \alpha) \right] v_o \quad \text{where } x, \alpha \text{ constant}$$

from our 3D expression?

Projection onto the 2D sky (Nadolny et al 2021)

$$\frac{\langle \mathcal{D}_{\text{gal}} \rangle (< m_*)}{v_o} = \frac{1}{\bar{N}_{\Omega} (< m_*)} \int_0^{\infty} dz \mathcal{D}_{\text{gal}}(z, < m_*) \bar{N}(z, < m_*)$$
$$\bar{N}_{\Omega} (< m_*) = \int_0^{\infty} dz \bar{N}(z, < m_*)$$

This leads to

$$\frac{\langle \mathcal{D}_{\text{gal}} \rangle}{v_o} = 2 + \frac{1}{\bar{N}_{\Omega}} \int_0^{\infty} dz x(1 + \alpha) \bar{N}$$

where

$$x = 2.5s$$

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- The original Ellis-Baldwin formula is only recovered **if we assume x and α are constant**
- Using $x(z)$ and $\alpha(z)$ in the EB formula is **incorrect**
- x and α should be determined from the luminosity function

The simplified Ellis-Baldwin model

$$\langle \mathcal{D}_{\text{gal}} \rangle_{\text{EB}} = \left[2 + x(1 + \alpha) \right] v_o$$

means that **in principle**, the apparent excess dipole magnitude

$$v_{o,\text{gal}} > v_{o,\text{cmb}}$$

could be due to the implicit approximation of constant magnification bias and spectral index (Dalang & Bonvin 2021).

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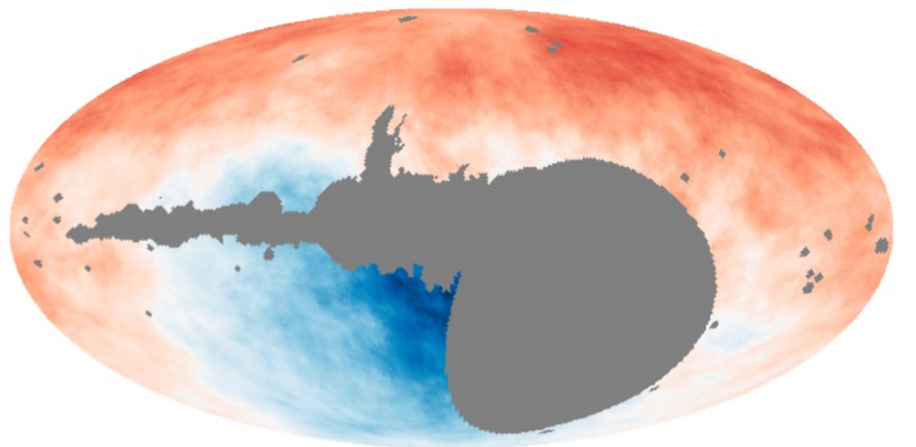
$$v_{o,\text{gal}} > v_{o,\text{cmb}}$$

could be due to the implicit approximation of constant magnification bias and spectral index (Dalang & Bonvin 2021).

In other words, it is possible that

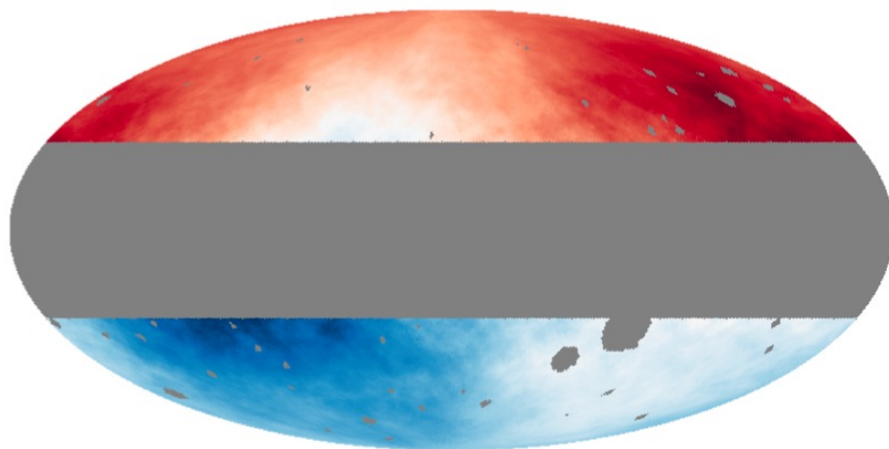
$$v_{o,\text{gal}} \approx v_{o,\text{cmb}} \quad \text{since} \quad \langle \mathcal{D}_{\text{gal}} \rangle_{\text{EB}} < \langle \mathcal{D}_{\text{gal}} \rangle_{\text{true}}$$

Measuring the dipole in a better 2D sample



16.6 source deg^{-2} 17.2

NVSS

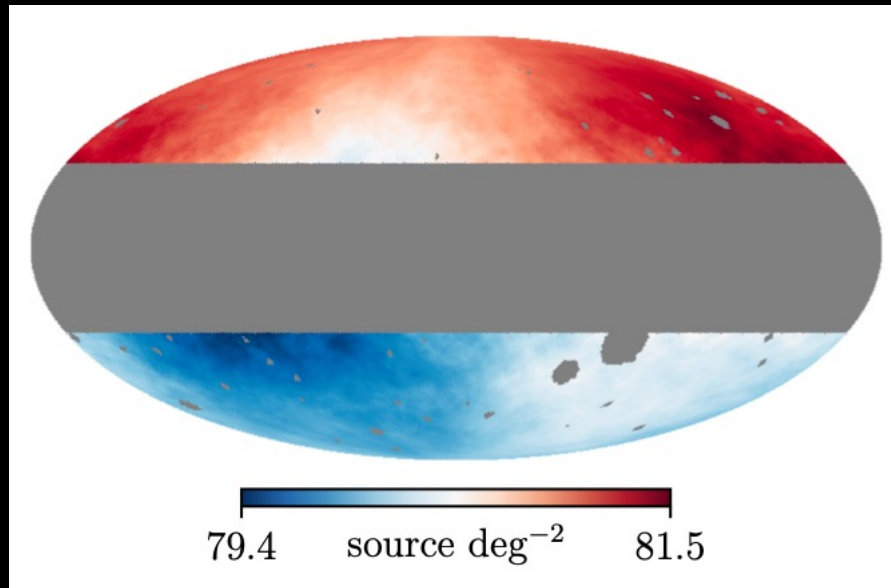


79.4 source deg^{-2} 81.5

CatWISE2020
(1.36M quasars, mid-IR)

(Secret et al 2022)

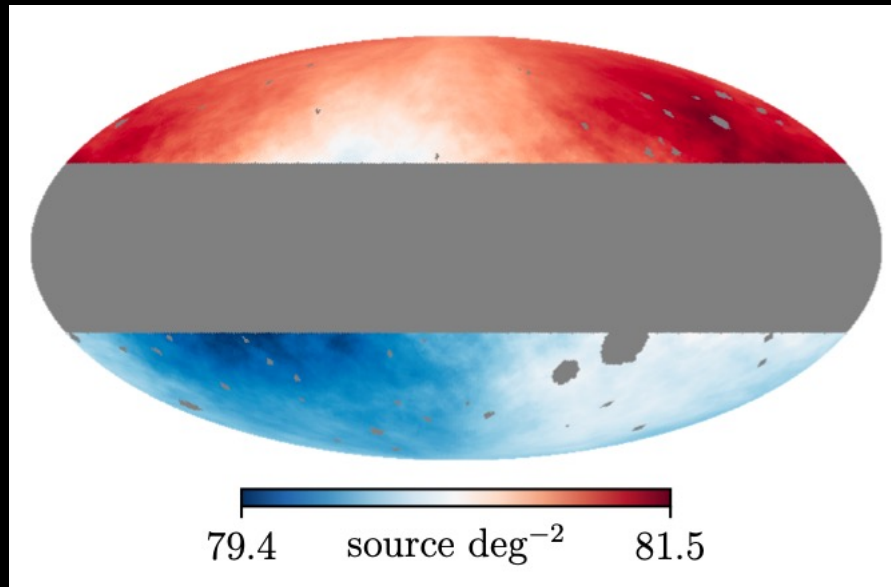
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CatWISE2020

Secrest et al 2021, 2022 use the Ellis-Baldwin formula to find that
the dipole magnitude is in tension with the CMB at $>4\sigma$

Measuring the dipole in a better 2D sample

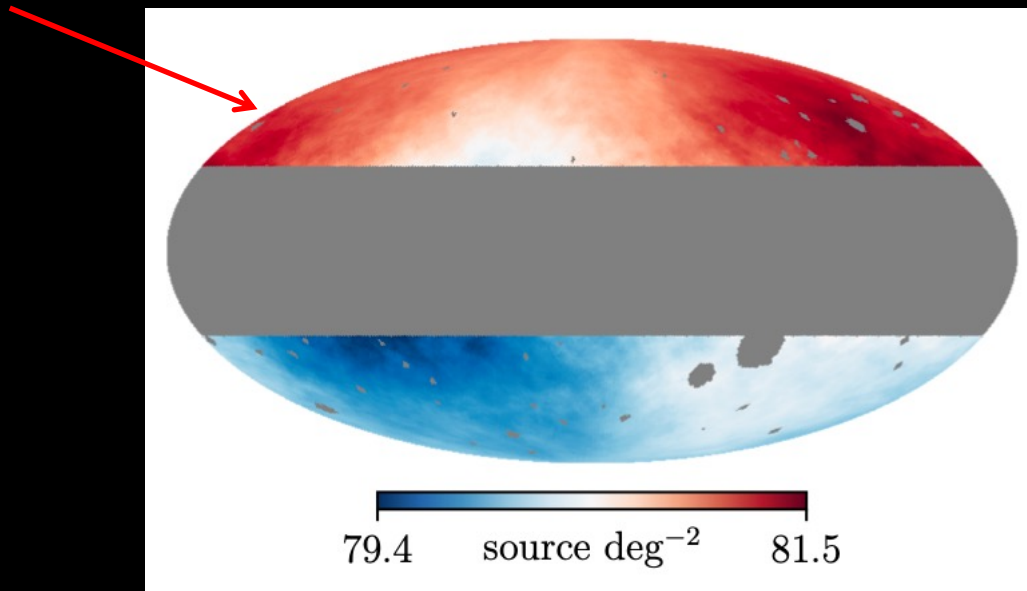


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Is this robust?

Measuring the dipole in a better 2D sample



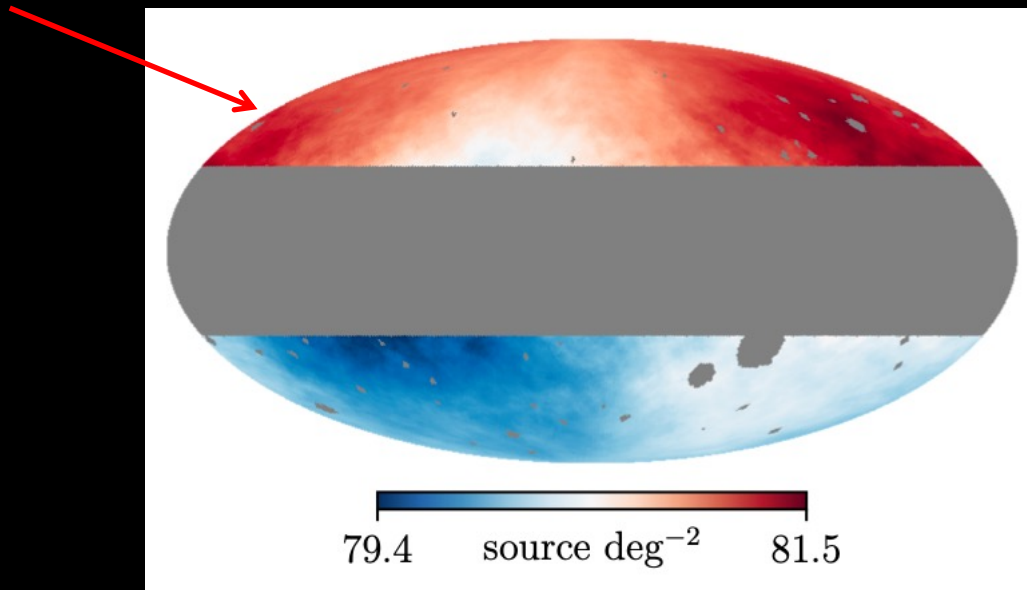
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Is this robust?

- There could be unaccounted for systematics on very large scales

Measuring the dipole in a better 2D sample



CatWISE2020

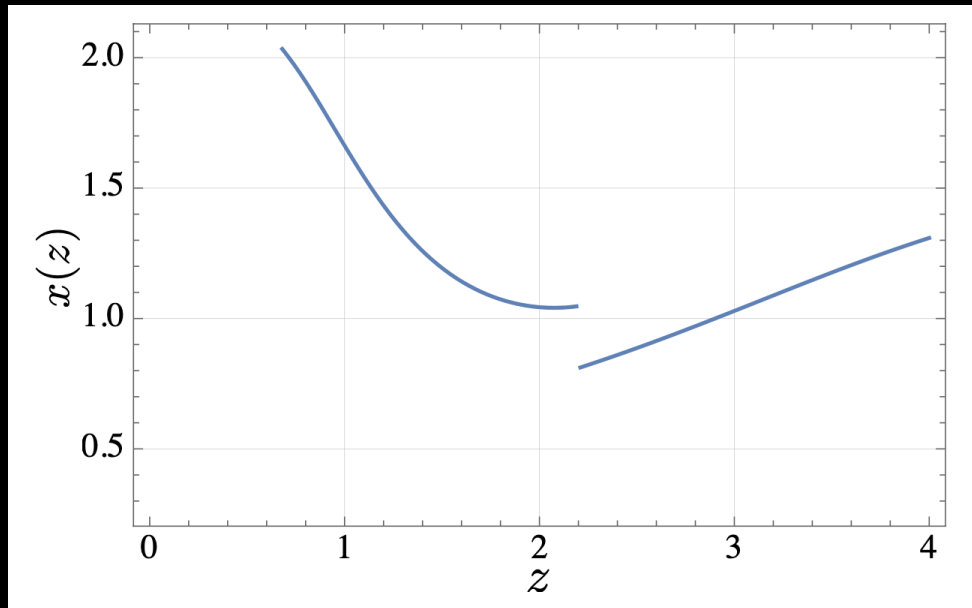
Secrest et al 2021, 2022 use the Ellis-Baldwin formula to find that
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Is this robust?

- There could be unaccounted for systematics on very large scales
- The Ellis-Baldwin formula could be a bad approximation

x and α should be determined by the luminosity function.

For example – the eBOSS quasar LF (Wang et al 2020) gives for $x(z)$:



(Dalang & Bonvin 2021)

Clearly x is not constant – and neither is α .

Extrapolating from eBOSS to CatWISE2020 –

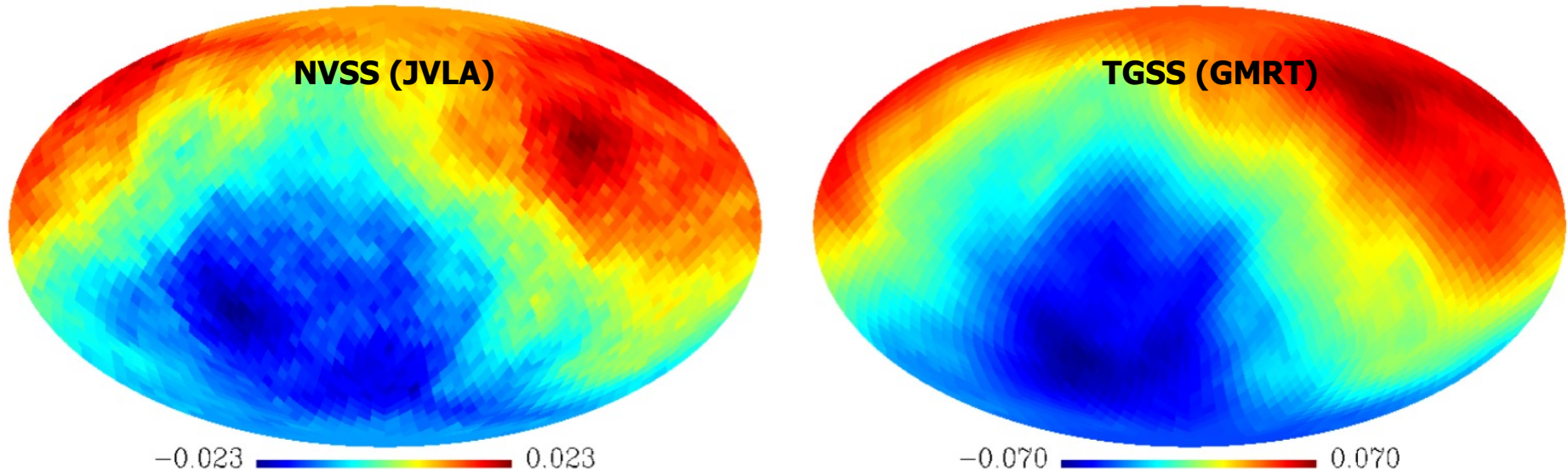
the tension with CMB can in principle be removed (Dalang & Bonvin 2021)

But this does not take account of differences in selection criteria and redshift range.

This seems to be an open question for further investigation – which is critical for testing the Cosmological Principle.

Extra slides

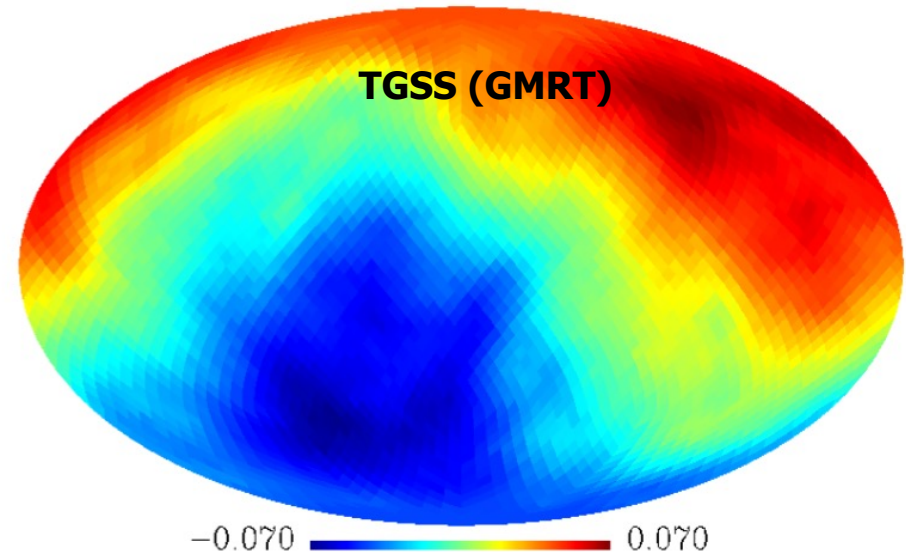
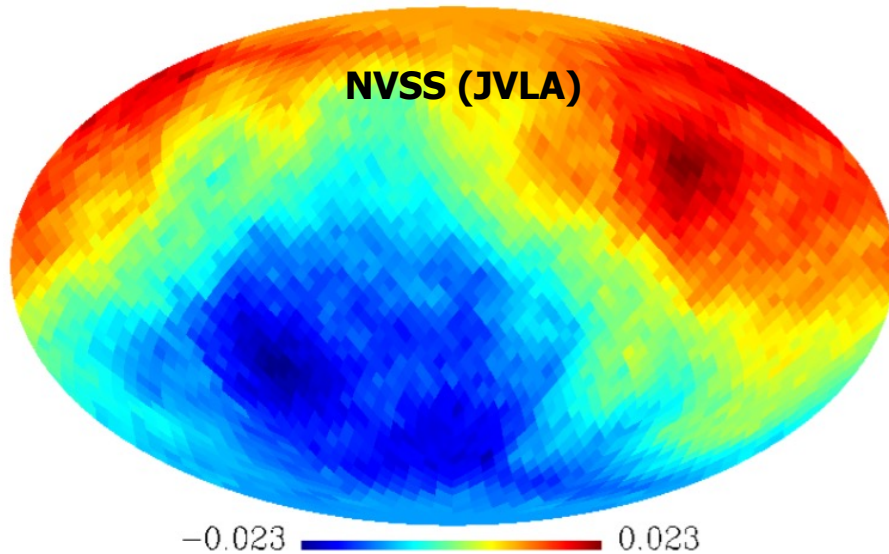
Including a newer radio continuum survey TGSS:



(Bengaly, RM, Santos 2017)

- Dipole direction is roughly consistent with CMB.
- Dipole magnitude in TGSS even larger – due to flux calibration systematics (Tiwari et al 2019; Secrest et al 2022).

Including a newer radio continuum survey TGSS:



(Bengaly, RM, Santos 2017)

Conclusions:

- Dipole noise is too large (not enough galaxies).
- The test is not robust with current radio continuum surveys.

Angular power spectrum dipole:

$$\tilde{C}_1 = C_1 + C_1^{v_o}$$

intrinsic \ll kinematic

$$C_1^{v_o}(z, z', m_*) = \frac{4\pi}{9} \mathcal{D}(z, m_*) \mathcal{D}(z', m_*)$$

SKA dipole
relative to
CMB dipole
for $z'=z$

