

# Cosmology under the fractional calculus approach: a possible $H_0$ tension resolution?

Tensions in Cosmology

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# Cosmology under the fractional calculus approach

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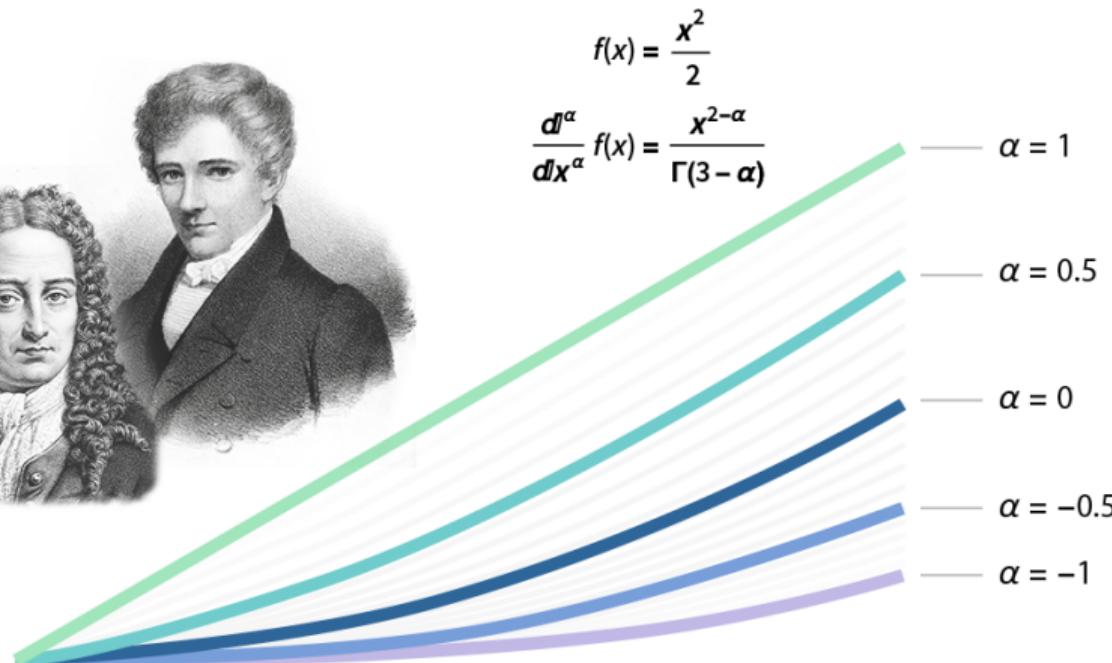
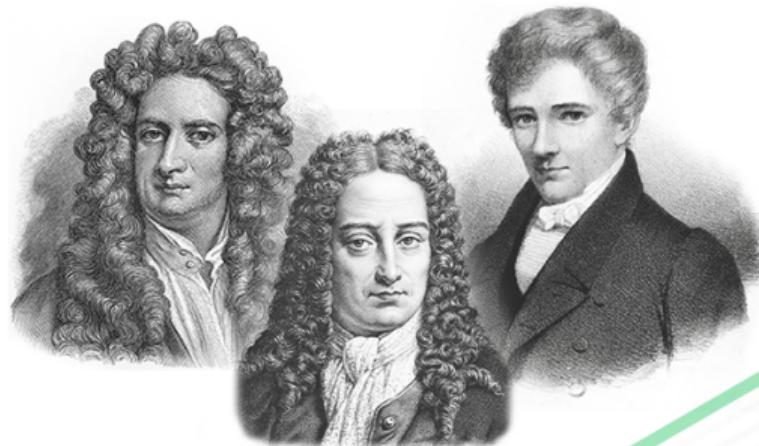
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- ① Fractional differentiation has drawn increasing attention in the study of so-called "anomalous" social and physical behaviours, where the scaling power law of fractional order appears universal as an empirical description of such complex phenomena.
- ② In the standard mathematical models of integer-order derivatives, including nonlinear models, do not work adequately in many cases where power law is clearly observed.
- ③ To accurately reflect the non-local, frequency- and history-dependent properties of power law phenomena, alternative modelling tools have to be introduced, such as fractional calculus.
- ④ Research in fractional differentiation is inherently multi-disciplinary and has its application across various disciplines.



- ① Fractional-order Systems and Controls: Fundamentals and Applications, Concepción A. Monje, YangQuan Chen, Blas M. Vinagre, Springer 2010.
- ② Stabilization and Control of Fractional Order Systems: A Sliding Mode Approach, Bandyopadhyay Bijnan, Shyam Kamal, Springer 2015
- ③ Advances in Robust Fractional Control, Fabrizio Padula, Antonio Visioli, Springer 2014.
- ④ Fractional Calculus: An Introduction for Physicists, Richard Herrmann, World Scientific, 2014.
- ⑤ Fractional Dynamics: Applications of Fractional Calculus to Dynamics of Fractional Calculus to Dynamics of Particles, Fields and Media, Vasily E. Tarasov, Springer 2010.
- ⑥ Fractional Dynamics: Recent Advances, Joseph Klafter, S. C. Lim, Ralf Metzler, World Scientific, 2012.
- ⑦ Advanced Methods in the Fractional Calculus of Variations, Agnieszka B. Malinowska, Tatiana Odzijewicz, Delfim F.M. Torres, Springer 2015.
- ⑧ Fractional Trigonometry: With Applications to Fractional Differential Equations and Science, Carl F. Lorenzo, Tom T. Hartley, Wiley 2016.

Cauchy's formula for the integral multiple of integer order  $\mu > 0$  [Uchaikin V. V., 2013,

Fractional derivatives for physicists and Engineers. Higher Education Press]

$${}_c I_x^\mu f(x) = \frac{1}{\Gamma(\mu)} \int_c^x f(t)(x-t)^{\mu-1} dt. \quad (1)$$

The Riemann-Liouville derivative (RLD) with  $\alpha \geq 0$  for  $f(x)$  is defined by

$$\begin{aligned} D_x^\alpha f(x) &\equiv \frac{d^n}{dx^n} I_x^{n-\alpha} f(x) \\ &= \Gamma(n-\alpha)^{-1} \frac{d^n}{dx^n} \int_c^x \frac{f(t)}{(x-t)^{\alpha-n+1}} dt, \end{aligned} \quad (2)$$

where  $n = [\alpha] + 1$  being  $\alpha \in (n-1, n)$ .



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# Some rules

$$D_x^\alpha [D_x^\beta f(x)] = D_x^{\alpha+\beta} f(x) - \sum_{j=1}^n D_x^{\beta-j} f(c+) \frac{(x-c)^{-\alpha-j}}{\Gamma(1-\alpha-j)}, \quad (3)$$

or in other words  $D_x^\alpha D_x^\beta f(x) \neq D_x^{\alpha+\beta} f(x)$ , if only not all derivatives  $D_x^{\beta-j} f(c+)$  at the beginning of the interval are equal to zero.

Leibniz rule reads as

$$D_x^\alpha [f(x)g(x)] = \sum_{k=0}^{\infty} \frac{\Gamma(\alpha+1)}{k! \Gamma(\alpha-k+1)} D_x^{\alpha-k} f(x) D_x^k g(x), \quad (4)$$

having the usual when  $\alpha = n$ .



# Caputo Approach

The Caputo definition is:

$$\begin{aligned} {}^C D_x^\alpha f(x) &\equiv {}_c I_x^{n-\alpha} D_x^n f(x) \\ &= \Gamma(n - \alpha)^{-1} \int_c^x \frac{\frac{d^n}{dt^n} f(t)}{(x - t)^{\alpha - n + 1}} dt \end{aligned} \quad (5)$$

where  $n = \max\{0, [\alpha]\}$ .

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The background cosmology is based on the Friedmann-Lemaître-Robertson-Walker (FLRW) metric for a flat Universe. The fractional effective action can be written in the form

$$S_{eff} = \Gamma(\alpha)^{-1} \int_0^t \left[ \frac{3}{8\pi G} \left( a^2 \ddot{a} + a \dot{a}^2 - a^2 \dot{a} \right) + a^3 \mathcal{L}_m \right] \times (t - \tau)^{\alpha-1} d\tau, \quad (6)$$

where  $\mathcal{L}_m$  is the matter Lagrangian,  $\alpha$  is the fractional constant parameter,  $t$  and  $\tau$  are the physical and proper time respectively and where the  $\Lambda$  is not considered [Shchigolev V. K., 2011, Commun. Theor. Phys., 56, 389]



# Modified Friedmann equations

The minimization of the fractional action (6) generates the modified Friedmann equation

$$H^2 + \frac{(1 - \alpha)}{t} H = \frac{8\pi G}{3} \sum_i \rho_i, \quad (7)$$

and the continuity equation is given by

$$\sum_i \left[ \dot{\rho}_i + 3 \left( H + \frac{1 - \alpha}{3t} \right) (\rho_i + p_i) \right] = 0. \quad (8)$$

When  $\alpha = 1$ , the standard cosmology is recovered without  $\Lambda$ .



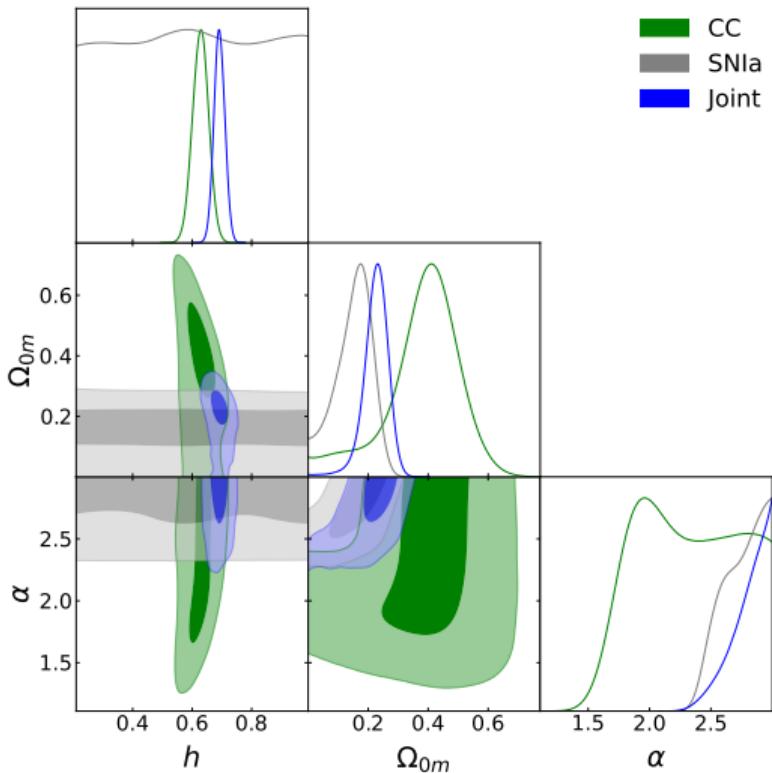
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Using the equation of state  $p_i = w_i \rho_i$ , where  $w_i \neq -1$  are constants,

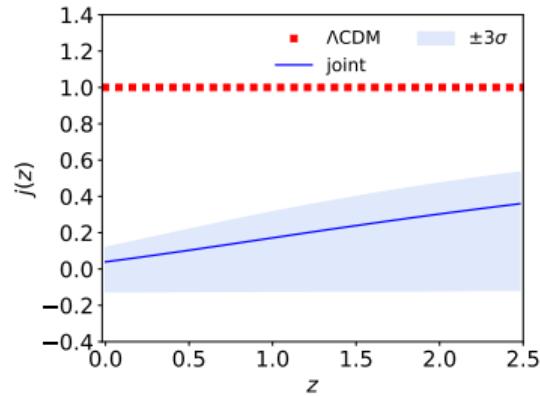
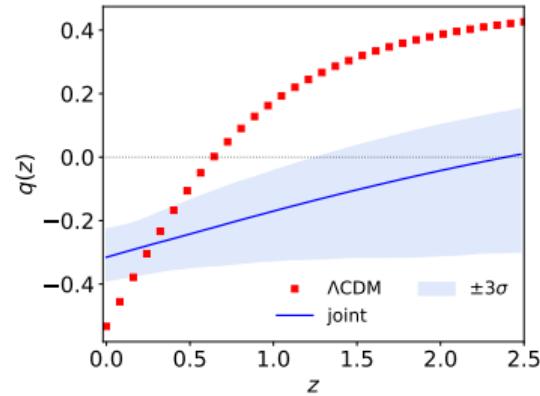
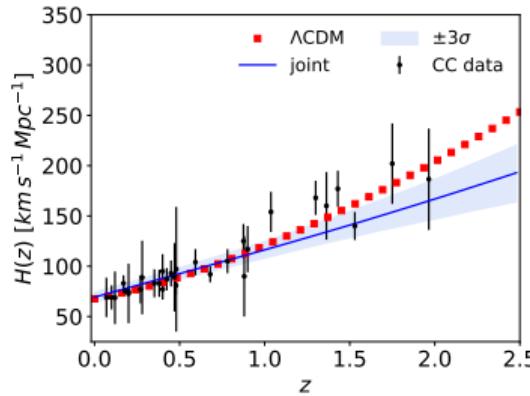
$$\begin{aligned} & \sum_i (1 + w_i) \rho_i \left[ \frac{\dot{\rho}_i}{(1 + w_i) \rho_i} + 3 \frac{\dot{a}}{a} + \frac{1 - \alpha}{t} \right] \\ &= \sum_i (1 + w_i) \rho_i \frac{d}{dt} \left[ \ln \left( \rho_i^{1/(1+w_i)} a^3 t^{1-\alpha} \right) \right]. \end{aligned} \quad (9)$$

Assuming separated conservation equations, setting  $a(t_U) = 1$ , where  $t_U$  is the age of the Universe, and denoting by  $\rho_{0i}$  the current value of energy density of the  $i$ -th species,

$$\rho_i(t) = \rho_{0i} a(t)^{-3(1+w_i)} (t/t_U)^{(\alpha-1)(1+w_i)}. \quad (10)$$



**Figure:** 2D likelihood contours at 68% and 99.7% CL, alongside the corresponding 1D posterior distribution of the free parameters



**Figure:** Left to right: reconstruction of the  $H(z)$ ,  $q(z)$ , and  $j(z)$ , in fractional cosmology represent the results of  $\Lambda$ CDM cosmology with  $h = 0.6766$  and  $\Omega_{m0} = 0.3111$  [Planck:2018]



Sample	$\chi^2_{\text{min}}$	$h$	$\Omega_{0m}$	$\alpha$
CC	16.14	$0.629^{+0.027}_{-0.027}$	$0.399^{+0.093}_{-0.122}$	$2.281^{+0.492}_{-0.433}$
SNIa	54.83	$0.599^{+0.275}_{-0.269}$	$0.160^{+0.050}_{-0.072}$	$2.771^{+0.161}_{-0.214}$
Joint	78.69	$0.692^{+0.019}_{-0.018}$	$0.228^{+0.035}_{-0.040}$	$2.839^{+0.117}_{-0.193}$

**Table:** Best-fit values and their 68% CL uncertainties for fractional cosmology (dust + radiation) with CC, SNIa and a Joint analysis.



## Defining the dimensionless variables

$$x_{1,2} = \frac{8\pi G t^2 \rho_{m,r}}{3} / \left( A + \frac{|1-\alpha|}{2} \right)^2, \quad A = tH, \quad \tau = \ln(t), \quad (11)$$

we produce a 2D dynamical system subject to the restriction

$$G(x_1, x_2, A) := A^2 + (1 - \alpha)A - (x_1 + x_2)(A + |1 - \alpha|/2)^2 = 0. \quad (12)$$

Label	$\Omega_m$	$\Omega_r$	$H$	$q$	Solution	$a(t) = (t/t_U)^A$
$P_1$	Indeterminate	Indeterminate	0	Indeterminate	Static universe	$a(t) = \text{constant}$
$P_2$	0	$\frac{7-4\alpha}{2\alpha+1}$	$\frac{2\alpha+1}{6t}$	$-\frac{2\alpha-5}{2\alpha+1}$	Power-law (decelerated if $\alpha < \frac{5}{2}$ ) Power-law (accelerated if $\alpha > \frac{5}{2}$ )	$a(t) = (t/t_U)^{(2\alpha+1)/6}$
$P_3$	$\frac{2(2-\alpha)}{\alpha+1}$	0	$\frac{\alpha+1}{3t}$	$-\frac{\alpha-2}{\alpha+1}$	Power-law (decelerated if $\alpha < 2$ ) Power-law (accelerated if $\alpha > 2$ )	$a(t) = (t/t_U)^{\frac{1+\alpha}{3}}$
$P_4$	0	0	$\frac{\alpha-1}{t}$	$-\frac{\alpha-2}{\alpha-1}$	Power-law (accelerated if $\alpha < 1$ or $\alpha > 2$ ) Power-law (decelerated if $1 < \alpha < 2$ )	$a(t) = (t/t_U)^{\alpha-1}$

Table: Equilibrium points of associated dynamical system.



# Results

- ①  $\alpha$ , is estimated for each dataset, and in particular, we have  $\alpha = 2.839_{-0.193}^{+0.117}$  for the joint analysis, allowing an accelerated Universe.
- ② We recover traditional calculus when  $\alpha = 1$ ; however, in the region  $0 < \alpha < 1$ , obtaining an accelerated physical Universe at late stages is not feasible. There we obtain an accelerated power-law solution corresponding to negative values for the age of the Universe.
- ③ Hence, one way to avoid this problem is under the introduction of  $\Lambda$ , which will act as a cosmological constant. However, we are in a loop because the idea explains the Universe's acceleration through the  $\alpha$  term, which contributes to the fractional calculus theory.
- ④ The other way is to consider  $\alpha > 2$ , and then we get an accelerated physical Universe at late stages.



# Results

- ① On the other hand, the age of the Universe is estimated for each dataset,  $t_U/\text{Gyrs} = 33.633^{+14.745}_{-15.095}$  (CC),  $33.837^{+27.833}_{-10.788}$  (SNIa) and  $33.617^{+3.411}_{-4.511}$  (joint).
- ② For the Joint value, we obtain around 2.4 times larger than the age of the Universe expected under the standard paradigm, which is also in disagreement with the value obtained with globular clusters,  $t_U = 13.5^{+0.16}_{-0.14} \pm 0.23$  [Valcin D., Jimenez R., Verde L., Bernal J. L., Wandelt B. D., 2021, Journal of Cosmology and Astroparticle Physics, 2021, 017].
- ③ The term  $(1 - \alpha)H/t$ , which acts as an extra source of mass leading to a closed Universe, could be the origin of this older Universe. In closed scenarios, the Universe becomes older than the standard prediction [Di Valentino E., et al., 2021b, Astropart. Phys., 131, 102607].



# Results

- ① Figure 3 displays the reconstruction of the  $H_0(z)$  diagnostic [Krishnan C., Ó Colgáin E., Sheikh-Jabbari M., Yang T., 2021, Phys. Rev. D, 103] for the fractional cosmology and its error band at  $3\sigma$  CL.
- ② Although the path (solid line) for the fractional cosmology is consistent within  $3\sigma$  with the CMB Planck value [Aghanim N., et al., 2018, Planck 2018 results. VI. Cosmological parameters (arXiv:1807.06209)] for  $z \lesssim 1.5$ , we can observe that it presents a trend to the  $H_0$  value obtained by SH0ES [Riess A. G., Casertano S., Yuan W., Macri L. M., Scolnic D., 2019, Astrophys. J., 876, 85] for the present time.
- ③ Nevertheless, the  $H_0$  value for  $1.5 < z < 2.5$  is lower than the Planck value, suggesting a tension in this value.



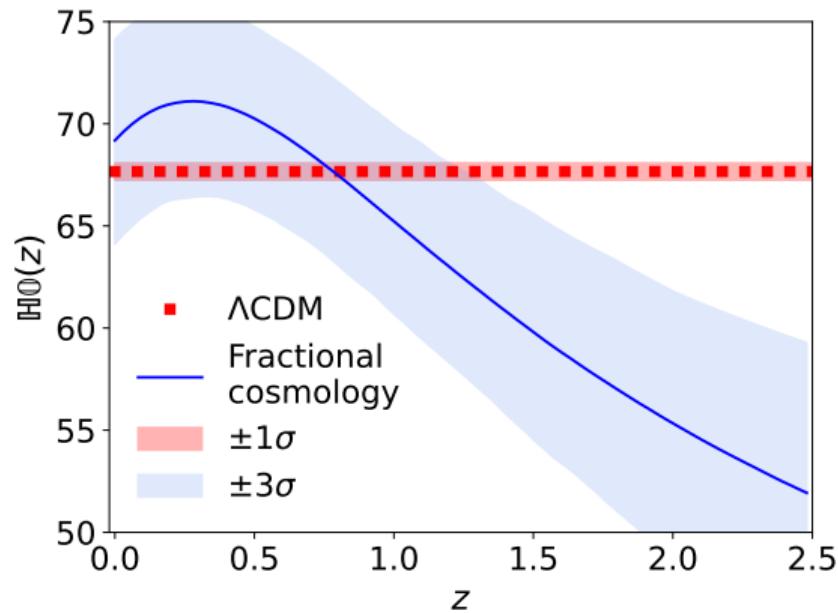


Figure:  $H_0(z)$  diagnostic for fractional cosmology and its comparison against  $\Lambda\text{CDM}$  model.

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# Thank you for your attention!

