FAKULTÄT FÜR INFORMATIK UND ELEKTROTECHNIK UNIVERSITÄT ROSTOCK


SIMULATION PROGRESS AND PLANS AT ROSTOCK/DESY

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## Overview

## Algorithms for 3D space charge calculations

- Multigrid based Poisson solver
- Software package MOEVE
- Simulations with GPT and ASTRA


## Particle tracking and e-cloud simulations

- Poisson solver for beam pipes with elliptical cross section
- Implementation of the tracking algorithm
- Plans for e-cloud simulations


## Algorithms for 3D Space Charge Calculations

Development of a 3D Poisson solver
(Gisela, since 1999)

- Particle mesh method
- Poisson solver based on multigrid
- For non-equidistant tensor product meshes
- Software package MOEVE 2.0
- Part of the tracking code GPT 2.7 (General Particle Tracer, Pulsar Physics)
- Part of ASTRA (test phase)
- Plan: (self-) adaptive multigrid



## Features of MOEVE

$-\Delta \varphi=\frac{\rho}{\varepsilon_{0}} \quad$ in $\Omega \subset \mathbb{R}^{3} \quad \begin{gathered}\text { Finite difference } \\ \text { discretization }\end{gathered} \quad A x=b$
iterative solvers:

- multigrid (MG)
- multigrid pre-conditioned conjugate gradients
- pre-conditioned conjugate gradients (Jacobi) (PCG)
- successive over relaxation (SOR)
- step size: non-equidistant
- numerical effort: $O(M) \quad M$ : total number of grid points
boundary conditions
- free space boundary
- perfect conducting rectangular box
- perfect conducting beam pipe with elliptical cross section


## Why Multigrid?

- Other Poisson solvers are much easier to implement
- They slow down considerably on nonequidistant meshes


Discretization of a spherical bunch


## Simultions with GPT

## COMPUMAG 2003



## Simulations with ASTRA

## EPAC 2006

- Gaussian particle distribution:

$$
\sigma_{x}=\sigma_{y}=0.75 \mathrm{~mm}, \sigma_{z}=1.0 \mathrm{~mm}
$$

-10,000 macro particles

- charge: -1 nC
- energy: 2 MeV
- tracking distance: 3 m
- quadrupol at $\mathrm{z}=1.2 \mathrm{~m}$
- number of steps (Poisson solver): N=32



## Beam Pipes with Elliptical Cross Section

Development of a Poisson solver for beam pipes and development of a tracking procedure (Aleksandar, since 2005)
-3D Space charge fields of bunches in a beam pipe of elliptical cross section, part of MOEVE 2.0

- iterative solvers: BiCG, BiCGSTAB
- step size: non-equidistant

Poisson equation:

$$
\Gamma=[-a, a] \times[-b, b] \times[-c, c]
$$



$$
\begin{aligned}
-\Delta \varphi & =\frac{\varrho}{\varepsilon_{0}} & & \text { in } \Omega \subset \mathbb{R}^{3}, \\
\varphi & =0 & & \text { on } \partial \Omega_{1}, \\
\frac{\partial \varphi}{\partial n}+\frac{1}{r} \varphi & =0 & & \text { on } \partial \Omega_{2},
\end{aligned}
$$



## Space Charge Simulation Results

Motivated from the paper "Simulation of transverse single bunch instabilities and emittance growth caused by electron cloud in LHC and SPS".
E. Benedetto, D. Schulte, F. Zimmermann, CERN, Switzerland, G. Rumolo, GSI, Germany

- construction of a fast cosine transformation for the simulation of conducting boundaries

Numerical example:

- spherical bunch
- $r \ll a, b$

( $r$ : radius of the bunch, $a, b$ : the half axis of the elliptical beam pipe)
- uniformly distributed charge of $\mathbf{1 n C}$
- bunch is located the centre of the beam pipe


## Space Charge Simulation Results

## sunsey

Electric field $E_{x}$ along the $x$-axis of a
square
$a=b$

rectangular box
$a=1.5 b$

elliptic b.c. beam pipe with elliptical cross section w/o.b.c. open b.c. on a rectangular w.b.c. conducting b.c. on a rectangular


## Space Charge Simulation Results

Electric field $E_{x}$ along $y= \pm b / 2$ of a

## square

$a=b$

rectangular box
$a=1.5 b$

elliptic b.c. beam pipe with elliptical cross section w/o.b.c. open b.c. on a rectangular w.b.c. conducting b.c. on a rectangular


## Space Charge Simulation Results

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## Particle Tracking

## Tracking Algorithm

- initial macro particle distribution
- deposit charges on the mesh nodes
- space charge computation in the center of mass system (3D Poisson solvers from MOEVE 2.0)
- interpolation of the fields for each macro particle in the laboratory frame
- time integration of the Newton-Lorentz equation for each macro particle (leap frog scheme)


## Time Integration of Particle Equations *

$$
p(t)=\gamma(t) v(t) \quad m=m_{0} \gamma \quad \gamma=\frac{1}{\sqrt{1-\frac{v^{2}}{c^{2}}}}
$$

- relativistic generalisation of the Newton-Lorentz equation:

$$
m_{0} \frac{d \vec{p}}{d t}=q \vec{E}+q\left(\frac{\vec{p}}{\gamma} \times \vec{B}\right)
$$

$$
\begin{gathered}
\frac{\vec{p}^{n+\frac{1}{2}}-\vec{p}^{n-\frac{1}{2}}}{\Delta t}=\frac{q}{m_{0}}\left[\vec{E}^{n}+\frac{\vec{p}^{n+\frac{1}{2}}-\vec{p}^{n-\frac{1}{2}}}{2 \gamma^{n}} \times \vec{B}^{n}\right] \\
\vec{p}^{n+\frac{1}{2}}=\vec{p}^{+}+\frac{q \vec{E}^{n} \Delta t}{2 m_{0}} \\
\vec{p}^{n-\frac{1}{2}}=\vec{p}^{-}-\frac{q \vec{E}^{n} \Delta t}{2 m_{0}}
\end{gathered}
$$

- Boris rotation for 3D fields:

$$
\frac{\vec{p}^{+}-\vec{p}^{-}}{\Delta t}=\frac{q}{2 \gamma^{n} m_{0}}\left(\vec{p}^{+}-\vec{p}^{-}\right) \times \vec{B}^{n}
$$

*Birdsall, C.K. and A.B. Langdon, Plasma Physics via Computer Simulation, McGraw-Hill, New York, 1985.

## First Tracking Simulations

- cylindrical bunch
- 50,000 macro particles
- $\sigma_{\mathrm{x}}=\sigma_{\mathrm{y}}=\sigma_{\mathrm{z}}=1.0 \mathrm{~mm}$
- uniform distribution
- charge -1 nC
- $\mathrm{E}_{\text {kin }}=5 \mathrm{MeV}$
- tracking time 110 ps
- time step 5 ps



## Plans: Interaction Beam - E-Cloud

- starting with uniform particle distribution for the e-cloud

- separate space charge computation for the relativistic positron beam and the e-cloud
- each time step a new part of the computational domain gets uniform particle distribution for the e-cloud


The End - ECL2

## Comparisson of the tracking with ASTRA*

Tracking with ASTRA (employing 3D space charge routine for elliptical domains and ASTRA's own FFT space charge solver )

- Gaussian bunch, 10,000 macro particles

$$
\sigma_{x}=\sigma_{y}=0.75 \mathrm{~mm} \sigma_{z}=1.0 \mathrm{~mm}
$$

- charge of -1nC, average energy of 2 MeV
- beam pipe has a diameter of 24 mm

Investigate the transverse electric field after a drift of $z=0.3 \mathrm{~m}$

[^0]
## Comparisson of the tracking with ASTRA

The transverse electric field after a drift $z=0.3 \mathrm{~m}$ in a circular beam pipe.


## Comparisson of the tracking with ASTRA

The same drift was simulated with the 3D FFT space charge routine of Astra


## Finite difference equation and the system of linear equations

$\Gamma$ is generally discretized in $N_{x}, N_{y}$ and $N_{z}$ non-equidistant steps:

$$
\begin{gathered}
h_{x, 0}, h_{x, 1}, \ldots, h_{x, N_{x}-1} \\
\tilde{h}_{x, i}=\left\{\begin{array}{cc}
\frac{h_{x, i-1}+h_{x, i}}{2}, & i=1, \ldots, N_{x}-1 \\
\frac{h_{x, i}}{2}, & i=0, N_{x}
\end{array}\right. \\
+\tilde{h}_{y, j} \tilde{h}_{z, k}\left(-\frac{1}{h_{x, i-1}} \varphi_{i-1, j, k}+\left(\frac{1}{h_{x, i-1}}+\frac{1}{h_{x, i}}\right) \varphi_{i, j, k}-\frac{1}{h_{x, i}} \varphi_{i+1, j, k}\right) \\
+\quad \tilde{h}_{x, i} \tilde{h}_{z, k}\left(-\frac{1}{h_{y, j-1}} \varphi_{i, j-1, k}+\left(\frac{1}{h_{y, j-1}}+\frac{1}{h_{y, j}}\right) \varphi_{i, j, k}-\frac{1}{h_{y, i}} \varphi_{i, j+1, k}\right) \\
+\quad \tilde{h}_{x, i} \tilde{h}_{y, j}\left(-\frac{1}{h_{z, k-1}} \varphi_{i, j, k-1}+\left(\frac{1}{h_{z, k-1}}+\frac{1}{h_{z, k}}\right) \varphi_{i, j, k}-\frac{1}{h_{z, k}} \varphi_{i, j, k+1}\right) \\
=\tilde{h}_{x, i} \tilde{h}_{y, j} \tilde{h}_{z, k} f_{i, j, k}
\end{gathered}
$$

## Finite difference equation and the system of linear equations

Boundary adapted 7-point star of grid points inside the ellipticall domain $\Omega$


Linear system :

$$
A u=b
$$

System matrix $\boldsymbol{A}$ is:

- block structured
- nonsymmetric
- positve definite

Use of BiCGSTAB to solve the system


Non-symmetric 2-D Shortley-Weller Star.

## Solver performance




Computation time of BiCG (left) and BiCGSTAB (right) for gradually increasing number of discretization points in each coordinate ( $N_{x}, N_{y}, N_{z}$ ).


[^0]:    * K. Flöttmann, "Astra", DESY, Hamburg www.desy.del ~mpyflo, 2000.

