# **High-repetition-rate** beam drivers and plasma sources

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# The desire for high-energy particle accelerators

#### **Radio-frequency cavity**



#### Charge-density wave in a plasma





To first order the particle energy at an accelerator facility defines its discovery reach

> Free-electron lasers: energy → wavelength

> *High-energy physics*: centre-of-mass energy

> Therefore higher energy is often desired but conventional technology is limited in accelerating gradient due to electrical breakdown

> Plasma-wakefield accelerators offer a route to higher energies with smaller facilities due to O(GV/m) gradients



# **Our customers: High-energy-physics (and photon-science) needs**

> Let's say that we could produce 500 GeV beams tomorrow, what else would we need to do?

> The luminosity demands that certain properties be maximised/minimised (similar demands for integrated brightness at FELs):





# **Our customers: High-energy-physics (and photon-science) needs**

> Let's say that we could produce 500 GeV beams tomorrow, what else would we need to do?

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#### **Selected results:**

Litos et al., "High-efficiency acceleration of an electron beam in a plasma wakefield accelerator", **Nature** (2014)

Wu et al., "High-throughput injection-acceleration of electron bunches from a linear accelerator to a laser wakefield accelerator", Nat. Phys. (2021)

Lindstrøm et al., "Energy-spread preservation and high efficiency in a plasma-wakefield accelerator", **Phys. Rev. Lett.** (2021)

Pompili et al., "Energy spread minimisation in a beam-driven plasma wakefield accelerator", Nat. Phys. (2021)

Lindstrøm et al., "Preservation of beam quality in a plasma-wakefield accelerator", **under review** (2023)

![](_page_3_Picture_13.jpeg)

![](_page_3_Figure_14.jpeg)

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# High-average-power requirements for linear colliders

![](_page_4_Figure_1.jpeg)

> We are many orders of magnitude away from where we need to be in terms of bunches per second and average power...

>... but why?

# High-average-power requirements for linear colliders

![](_page_5_Figure_1.jpeg)

![](_page_5_Figure_3.jpeg)

>... but why?

- >Energy → scalable staging to high energy remains an open challenge (see Carl L's talk)
- > Other research priorities  $\rightarrow$  solving other open challenges in the field applicable to low rep. rate
- > Other application goals → many facilities in Europe are motivated by application to photon science
- > Unknown limits → the physics effects that may limit/permit high rep. rate are currently unknown/ undefined

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# **Defining the repetition-rate upper limit**

#### What defines the minimum inter-bunch separation in metallic cavities?

- E.g. X-band (~12 GHz) normal-conducting accelerating cavities
- > Minimum possible separation is ~ 80 ps.
- > Long-range transverse wakefields induced in the metallic cavitie they lead to emittance blow-up.
- > Actual separation set at **0.5 ns** i.e. 2 GHz.

![](_page_6_Figure_7.jpeg)

> Long-range transverse wakefields induced in the metallic cavities from an acceleration event live longer than this and must be avoided as

![](_page_6_Figure_9.jpeg)

### Equivalent effect in plasma accelerators is long-term plasma motion

- achievable repetition rate

![](_page_7_Figure_4.jpeg)

![](_page_7_Picture_6.jpeg)

![](_page_7_Figure_7.jpeg)

![](_page_7_Picture_8.jpeg)

![](_page_8_Picture_0.jpeg)

![](_page_8_Figure_1.jpeg)

R. D'Arcy et al., Nature 603, 58-62 (2022)

point of recovery

# ns (for experimental settings)

> Recovery time defined as the separation at which all three experimental signals are consistent with zero to within experimental uncertainties

- > All residuals consistent with zero at ~63 ns\*
- > Equivalent to a repetition-rate upper limit of O(10 MHz)\*\*

\*for working point in argon plasma of density ~1E16 cm \*\*if CW operation is permitted by other physics effects/technical limits

![](_page_8_Picture_11.jpeg)

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# High-average-power requirements for linear colliders

![](_page_9_Figure_1.jpeg)

> We are many orders of magnitude away from where we need to be in terms of bunches per second and average power...

> ... but why?

- **Energy**  $\rightarrow$  scalable staging to high energy remains an open challenge (see Carl L's talk)
- > Other research priorities  $\rightarrow$  solving other open challenges in the field applicable to low rep. rate
- > Other application goals  $\rightarrow$  many facilities in Europe are motivated by application to photon science
- **Some unknown limits**  $\rightarrow$  many of the physics effects that may limit/permit high rep. rate are currently unknown/undefined

#### How do we define the rest?

#### Two fundamental components of a beam-driven plasma accelerator

#### Radio-frequency-cavity front end

![](_page_10_Picture_2.jpeg)

> Can be either warm or superconducting  $\rightarrow$  there are benefits to each

	FEL	Collider	Current
Bunches per second	10 <sup>1</sup> - 10 <sup>6</sup>	10 <sup>4</sup> - 10 <sup>5</sup>	<b>10<sup>1</sup> - 10</b> <sup>6</sup>
Avg. beam power (W)	10 <sup>1</sup> - 10 <sup>5</sup>	106	<b>10<sup>1</sup> - 10</b> <sup>6</sup>

#### **Plasma-accelerator stage**

![](_page_10_Picture_7.jpeg)

> Discharge-capillary plasma stages have been a workhorse in the field for the last ~20 years\*

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\*ignoring laser- and beam-generated sources in the interest of brevity Page 11

![](_page_10_Picture_11.jpeg)

![](_page_10_Picture_12.jpeg)

![](_page_10_Picture_13.jpeg)

# Bunch-train patterns: the path to O(10,000) bunches per second

- Operation in a *pulsed mode* enables higher accelerating gradients and higher efficiencies

![](_page_11_Figure_4.jpeg)

Radio-frequency cavities are capable of operating in *continuous-wave* (CW) *mode*  $\rightarrow$  different to 'CW' operation in lasers

However, accelerating gradient is limited due to inefficiencies/electrical breakdown  $\rightarrow$  larger driver complexes required

	CLIC example
<b>Toperty #1</b> in the separation $\Delta_b$	dissipation of long-range transverse wakefields
<b>operty #2</b> -train length <i>n</i> b	balance of RF pulse length, and acceleration field, and electrical breakdowns
<b>toperty #3</b> alse separation $\Delta_t$	dissipation of the cumulative heating fro each bunch train

![](_page_11_Picture_11.jpeg)

![](_page_11_Picture_12.jpeg)

![](_page_11_Picture_13.jpeg)

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# Bunch-train patterns: the path to O(10,000) bunches per second

How is the bunch pattern defined in conventional accelerators?

![](_page_12_Picture_2.jpeg)

warm radio-frequency accelerating cavities

![](_page_12_Figure_4.jpeg)

![](_page_12_Figure_6.jpeg)

**Property #1** 

Inter-bunch separation  $\Delta_b$ 

#### **Property #2**

Bunch train length  $n_b$ 

**Property #3** 

Macro-pulse separation  $\Delta_t$ 

![](_page_12_Picture_14.jpeg)

![](_page_12_Picture_15.jpeg)

![](_page_12_Picture_16.jpeg)

![](_page_12_Picture_17.jpeg)

# **Bunch-train-pattern comparison for different technologies**

	Plasma accelerator		clc
Inter-bunch separation	<i>O</i> (100 ns)	554 ns	0.5 ns
Bunch-train length	???	726 µs	156 ns
Macro-pulse separation	???	100 ms	20 ms
Max. # of bunches per second	???	13120	15600

> Seemingly compatible with ILC-type superconducting RF... but not with CLIC-type warm RF

![](_page_13_Picture_5.jpeg)

## **Bunch-train-pattern comparison for different technologies**

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Max. # of bunches per second	???	13120	15600

- > Example of a staged facility: ~100 m in total length
- > Path-length difference between stages: ~2m = ~6 ns → normal-conducting technology!
- > Total # of stages (in this example) i.e. # of drive bunches per accelerating bunch: 16
- > Inter-bunch separation: 16 x 6 ns = ~100 ns

![](_page_14_Figure_7.jpeg)

> Don't exclude any one type of technology just yet!

![](_page_14_Figure_10.jpeg)

# **Bunch-train-pattern comparison for different technologies**

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![](_page_15_Figure_3.jpeg)

![](_page_15_Figure_5.jpeg)

![](_page_15_Figure_6.jpeg)

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# **Building up plasma-accelerator stages to high repetition rate**

Radio-frequency-cavity front end

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**Requirement:** increase the bunch-train length

**Result:** the desired bunches per second

**Requirement:** manage the increased average power -

**Plasma-accelerator stage** 

![](_page_16_Picture_8.jpeg)

		FEL	Collider	Curre
	Inter-bunch sep. (µs)	10 <sup>-1</sup> - 10 <sup>3</sup>	10 <sup>-1</sup> - 10 <sup>3</sup>	<b>10</b> -1
┥	Bunch-train length (#)	10 <sup>1</sup> - 10 <sup>3</sup>	10 <sup>1</sup> - 10 <sup>3</sup>	<b>10</b> º
Ì	Macro-pulse rate (Hz)	10 <sup>1</sup> - 10 <sup>2</sup>	10 <sup>1</sup> - 10 <sup>2</sup>	<b>10</b> <sup>1</sup>
┥	Bunches per second	10 <sup>1</sup> - 10 <sup>6</sup>	10 <sup>4</sup> - 10 <sup>5</sup>	<b>10</b> <sup>1</sup>
┥	Avg. beam power (W)	10 <sup>1</sup> - 10 <sup>5</sup>	106	<b>10</b> <sup>1</sup>

![](_page_16_Picture_11.jpeg)

![](_page_16_Picture_12.jpeg)

# **Returning to the comparison with conventional accelerators**

![](_page_17_Figure_1.jpeg)

ators* Plasma accelerators
----------------------------

dissipation of long-term plasma motion → O(100 ns)

???

???

\*CLIC example

![](_page_17_Picture_11.jpeg)

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# Returning to the comparison with conventional accelerators

![](_page_18_Figure_1.jpeg)

**Plasma accelerators** 

dissipation of long-term plasma motion → O(100 ns)

goal: similar plasma properties for each acceleration event

goal: plasma source capable of withstanding large heat loads

\*CLIC example

![](_page_18_Picture_10.jpeg)

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**Challenge:** Plasma-electron density decays exponentially on the µs timescale due to **expulsion** and **recombination** 

Image credit: Kyrre Ness Sjøbæk

![](_page_19_Picture_3.jpeg)

#### Plasma expulsion from a capillary-discharge plasma source

![](_page_19_Figure_6.jpeg)

Plasma-density decay in a open-ended discharge-capillary plasma stage

![](_page_19_Picture_9.jpeg)

**Challenge:** Plasma-electron density decays exponentially on the µs timescale due to **expulsion** and **recombination** 

![](_page_20_Figure_2.jpeg)

HALHF: Foster, D'Arcy, & Lindstrøm https://arxiv.org/abs/2303.10150

decreases the density

![](_page_20_Picture_6.jpeg)

![](_page_20_Picture_7.jpeg)

> Challenge: Cumulative heating of the plasma from inefficiencies in the system may modify the wakefield properties

C.A. Lindstrøm et al., Phys. Rev. Lett. **126**, 014801 (2021)

![](_page_21_Figure_3.jpeg)

Shot-to-shot stability of beams at FLASHForward

![](_page_21_Figure_6.jpeg)

Power-transfer diagram in a plasma accelerator

![](_page_22_Figure_2.jpeg)

Plasma evolution as a result of energy deposited on axis by plasma acceleration

**Challenge:** Cumulative heating of the plasma from inefficiencies in the system may modify the wakefield properties

![](_page_22_Figure_5.jpeg)

K. V. Lotov, PRSTAB 6, 061301 (2003)

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#### Simulations of how plasma background temperature modifies the plasma-wakefield properties

**Challenge:** Cumulative heating of the plasma will lead to **cumulative heating of the plasma stage** 

![](_page_23_Figure_2.jpeg)

Power-transfer diagram in a plasma accelerator

R. Zgadzaj *et al.*, Nat. Commun. **11**, 4753 (2020)

![](_page_23_Figure_7.jpeg)

#### Simulated energy-transport channels after driving a wake

![](_page_23_Picture_10.jpeg)

**Challenge:** Cumulative heating of the plasma will lead to **cumulative heating of the plasma stage** 

Image credit: Anthony Gonsalves

![](_page_24_Picture_3.jpeg)

#### **Grown-diamond capillary-discharge waveguides at LBNL**

![](_page_24_Figure_7.jpeg)

#### Heat-transfer simulation of a liquid-cooled plasma source

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#### **Challenge: Cooling requirements may be beyond what is achievable** with near-future technology

![](_page_25_Figure_2.jpeg)

#### **Parameters**:

 $dP_{cool}/ds$  : Required cooling rate

 $\Delta E_d$ : Driver energy loss

 $Q_d$ : Driver charge

 $Q_w$ : Witness charge

f: Collision frequency

 $L_{\rm s}$ : Plasma length

 $\eta_d$ : Efficiency from driver to wake

 $\eta_w$ : Efficiency from wake to witness

- $E_{7}$ : Accelerating gradient
- T: Transformer ratio

#### **Challenge: Cooling requirements may be beyond what is achievable** with near-future technology

![](_page_26_Figure_2.jpeg)

gradient is directly limited by the achievable cooling rate

#### > ... but are we limited in practice?

$$\frac{-\eta_w}{2} \approx E_z \cdot Q_w f \cdot (\frac{1}{\eta_w} - 1)$$
  
'space efficiency' 'energy efficiency'  
'physics efficiency'

**Conclusion:** For a certain particle flux (luminosity) and a certain energy-transfer efficiency (sustainability), the acceleration

![](_page_26_Picture_10.jpeg)

#### Challenge: Cooling requirements may be beyond what is achievable with near-future technology

![](_page_27_Figure_2.jpeg)

#### example for average power (per second): 10 GV/m x 1 nC x 10,000 s<sup>-1</sup> x (1/0.6 - 1) = 67 kW/m

- ... but the average power in a MHz bunch train will be **x100** the average power over a second at 10 kHz
- $> \dots$  and it may not be possible to 'manage' rapid temperature increases/stresses on the  $\mu$ s timescale

$$\frac{-\eta_{w}}{2} \approx E_{z} \cdot Q_{w} f \cdot (\frac{1}{\eta_{w}} - 1)$$
  
'space efficiency' 'energy efficiency'  
'physics efficiency'

> CLIC is expected to be able to 'manage'  $\sim$ 20 kW/m  $\rightarrow$  in the right ballpark but likely using very different cooling schemes

![](_page_27_Picture_14.jpeg)

#### **Challenge:** Cooling requirements may be beyond what is achievable with near-future technology

![](_page_28_Figure_2.jpeg)

- > If we can't boost the cooling rates for plasma stages over those of CLIC, where do we compromise?
- Inter-stage optics dominate the plasma-accelerator length at high energies  $\rightarrow$  lower gradients and longer stages?
- does this all necessitate operating the conventional linac at CW? Or
- > **Caveat**: This assumes that all the power makes it to the wall of the plasma stage
- > ...
- This challenge cannot be tackled in isolation  $\rightarrow$  iteration loop with attempts to solve the other challenges

$$\frac{-\eta_{w}}{2} \approx E_{z} \cdot Q_{w} f \cdot (\frac{1}{\eta_{w}} - 1)$$

$$\frac{1}{2} \cdot Q_{w} f \cdot (\frac{1}{\eta_{w}} - 1)$$

but does it? Expulsion of power with expulsion of plasma? Do unknown energy-transport channels help us? etc. etc.

![](_page_28_Picture_17.jpeg)

#### Plasma-wakefield accelerators at high repetition rates Summary and outlook

- The recovery time of a plasma-wakefield accelerator indicates compatibility with radiofrequency bunch-train patterns → a great first step... but still just a first step
- > The big challenge now is bridging the up-to-five order-of-magnitude gap from state-of-the-art to what is required
- Many outstanding scientific and technical goals to be reached with an emphasis on simulation tools, driver development, and plasma-source technology
- Schemes discussed here do not utilise high-power lasers but a future linear collider based on novel-accelerator technology will likely be based on both types of technologies e.g. laser drivers and discharge-based stages
- A coordinated international effort from both the LWFA and PWFA communities will be required to solve all the problems in the next decade

![](_page_29_Picture_6.jpeg)

# Plasma-wakefield accelerators at high repetition rates

#### **Open questions**

#### Beam drivers:

- The necessary beam drivers seemingly exist so it would be sensible to leverage 'shovel ready' designs if possible... But can we fully utilise them in their current/planned bunch-train format?
- If not, do we need to push the conventional community to reimagine how they operate their machines to best conform to our needs?
- > Bunch-train patterns have been the focus here but are the bunch parameters at this repetition rate sensible?

#### > Plasma stages:

- Can similar plasma properties be reproduced at MHz to enable acceleration of high-rep.-rate bunch trains?
   How hot does the plasma get due to MHz plasma acceleration? And will it substantially modify the plasma
- How hot does the plasma get due to MHz plasm properties?
- Can kW-MW levels of average power left in the plasma stage be managed?
- Do the plasma stages mandate CW operation of the linac? And if so, are CW discharge-generated plasma stages possible?

![](_page_30_Picture_11.jpeg)