



**Kavli Institute for Astronomy and Astrophysics
Peking University**

北京大学科维理天文与天体物理研究所



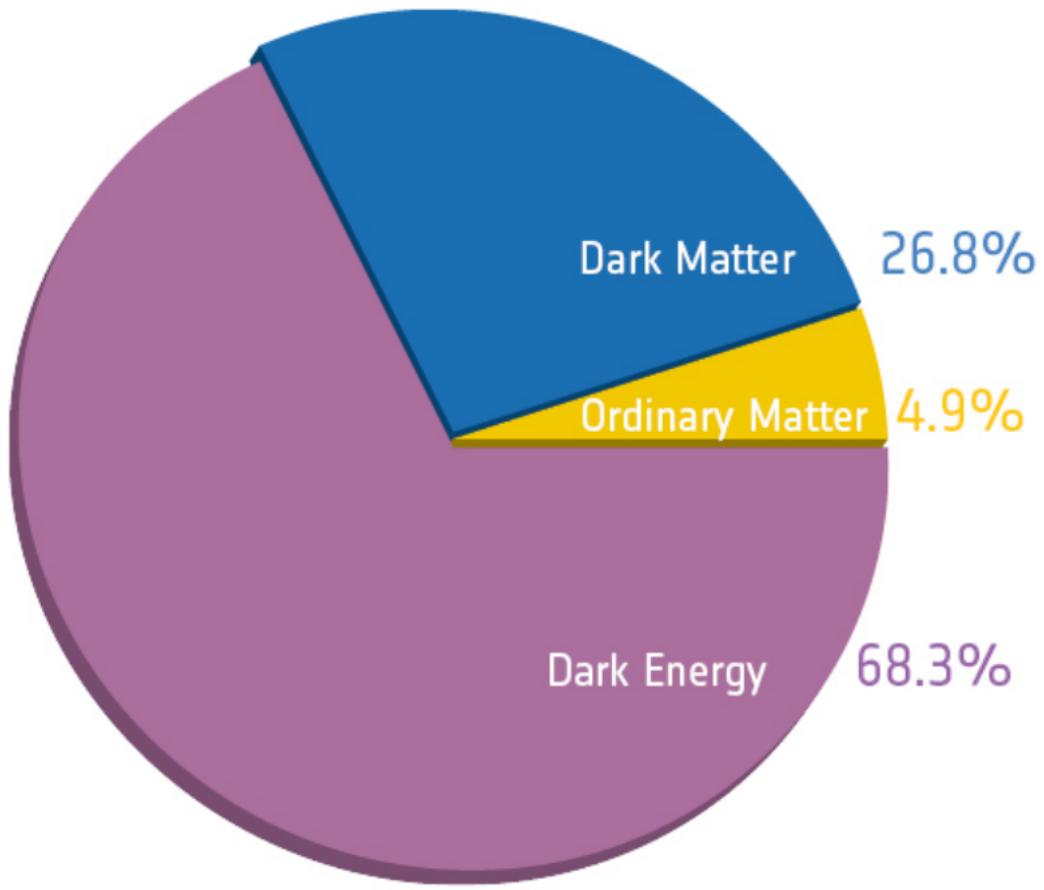
Intermediate Mass-Ratio Inspirals with Dark Matter Minispikes

Dicong Liang

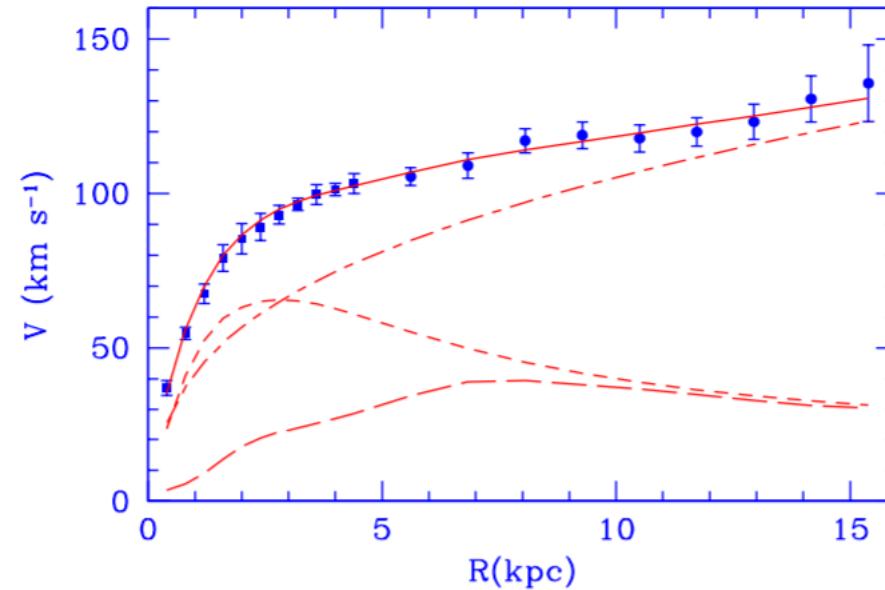
05/03/2023 Lisbon

Collaborators: Ning Dai, Yungui Gong and Tong Jiang
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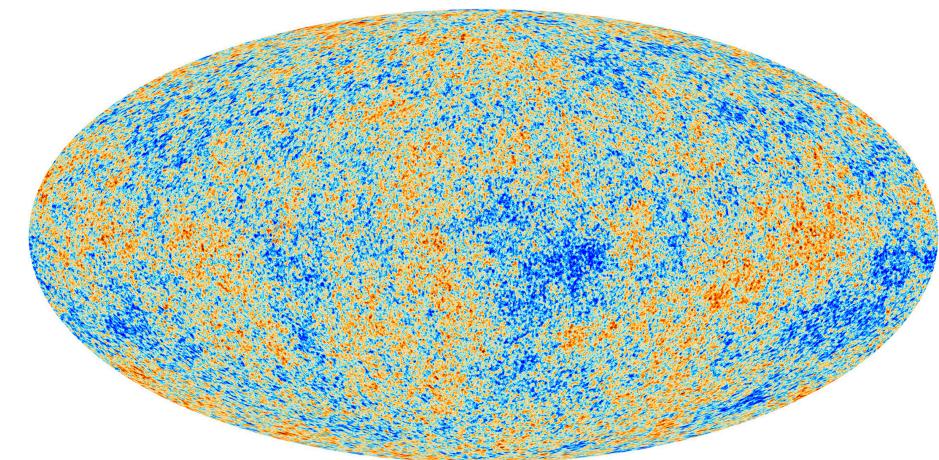
Dark Matter



(Credit: ESA and the Planck Collaboration)



Rotation Curve of M33
(E. Corbelli & P. Salucci, 2000)



Cosmic Microwave Background
(Credit: ESA and the Planck Collaboration)

Dark Matter Minispikes around an Intermediate Mass Black Hole

density profile

$$\rho_{\text{DM}}(r) = \begin{cases} \rho_{\text{sp}} \left(\frac{r_{\text{sp}}}{r} \right)^{\alpha}, & r_{\text{min}} \leq r \leq r_{\text{sp}} \\ 0, & r \leq r_{\text{min}} \end{cases}$$

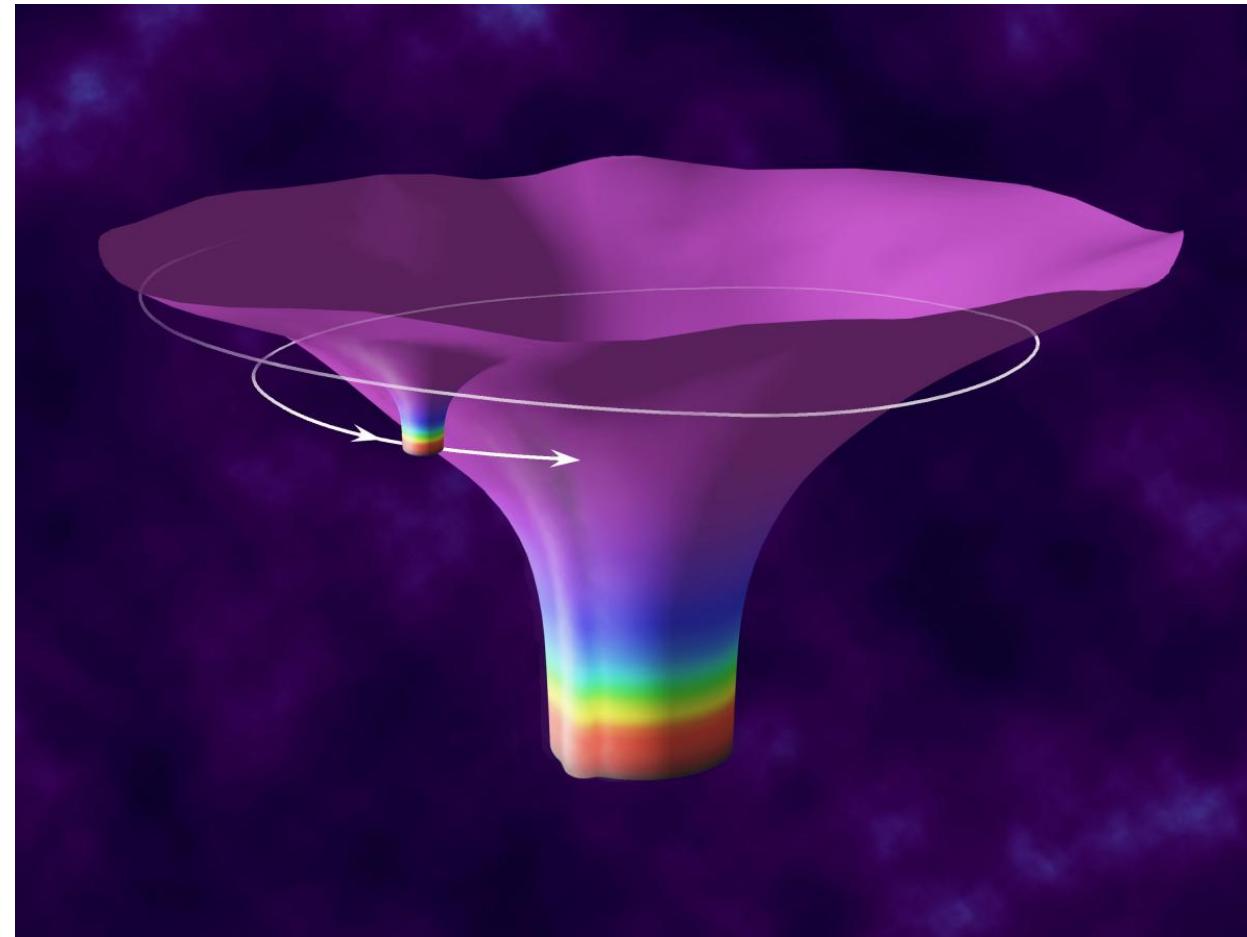
$$M = 10^3 M_{\odot} \quad \begin{aligned} r_{\text{min}} &= r_{\text{ISCO}} = 3R_s & r_{\text{sp}} &= 0.54 \text{ pc} \\ \rho_{\text{sp}} &= 226 M_{\odot}/\text{pc}^3 & 2.25 \leq \alpha &\leq 2.5 \end{aligned}$$

H. Zhao & J. Silk, 2005

K. Eda, et al. 2013

Intermediate Mass-Ratio Inspiral ($M/\mu = 10^2 \sim 10^4$)

$$M(10^3 M_{\odot}) + \mu(10 M_{\odot})$$



Four Major Factors

- Gravity from DM
- Dynamic Friction
- Accretion
- Radiation Reaction

Osculating Orbit

$$\mathbf{a} = -\frac{Gm}{r^2} \mathbf{n} + \mathbf{f}$$

perturbative force

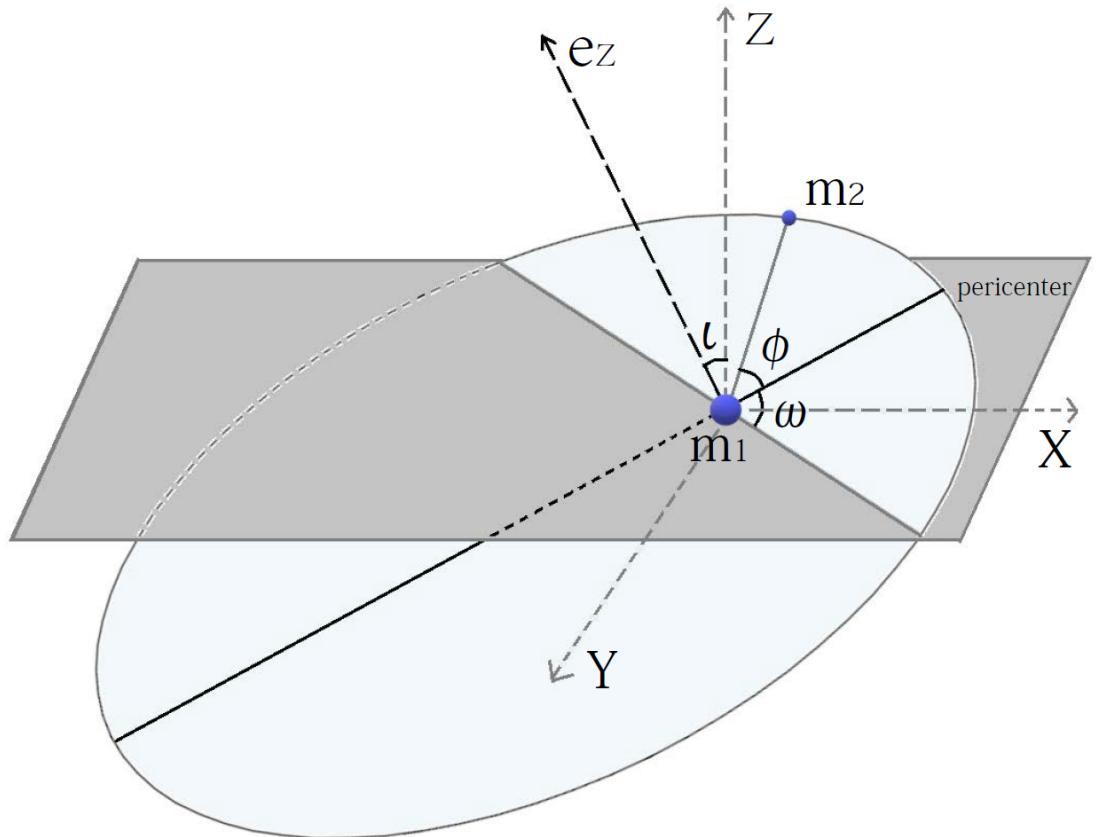
$$\mathbf{f} = \mathcal{R} \mathbf{n} + \mathcal{S} \mathbf{k} + \mathcal{W} \mathbf{e}_z$$

$$\frac{dp}{dt} = 2\sqrt{\frac{p^3}{Gm}} \frac{1}{1+e\cos\phi} \mathcal{S}$$

$$\frac{de}{dt} = \sqrt{\frac{p}{Gm}} \left[\sin\phi \mathcal{R} + \frac{2\cos\phi + e(1+\cos^2\phi)}{1+e\cos\phi} \mathcal{S} \right]$$

$$\frac{d\omega}{dt} = \frac{1}{e} \sqrt{\frac{p}{Gm}} \left[-\cos\phi \mathcal{R} + \frac{2+e\cos\phi}{1+e\cos\phi} \mathcal{S} - e \cot\iota \frac{\sin(\omega+\phi)}{1+e\cos\phi} \mathcal{W} \right]$$

$$\frac{d\phi}{dt} = \sqrt{\frac{Gm}{p^3}} (1+e\cos\phi)^2 + \frac{1}{e} \sqrt{\frac{p}{Gm}} \left[\cos\phi \mathcal{R} - \frac{2+e\cos\phi}{1+e\cos\phi} \sin\phi \mathcal{S} \right]$$



Gravity of DM minispikes

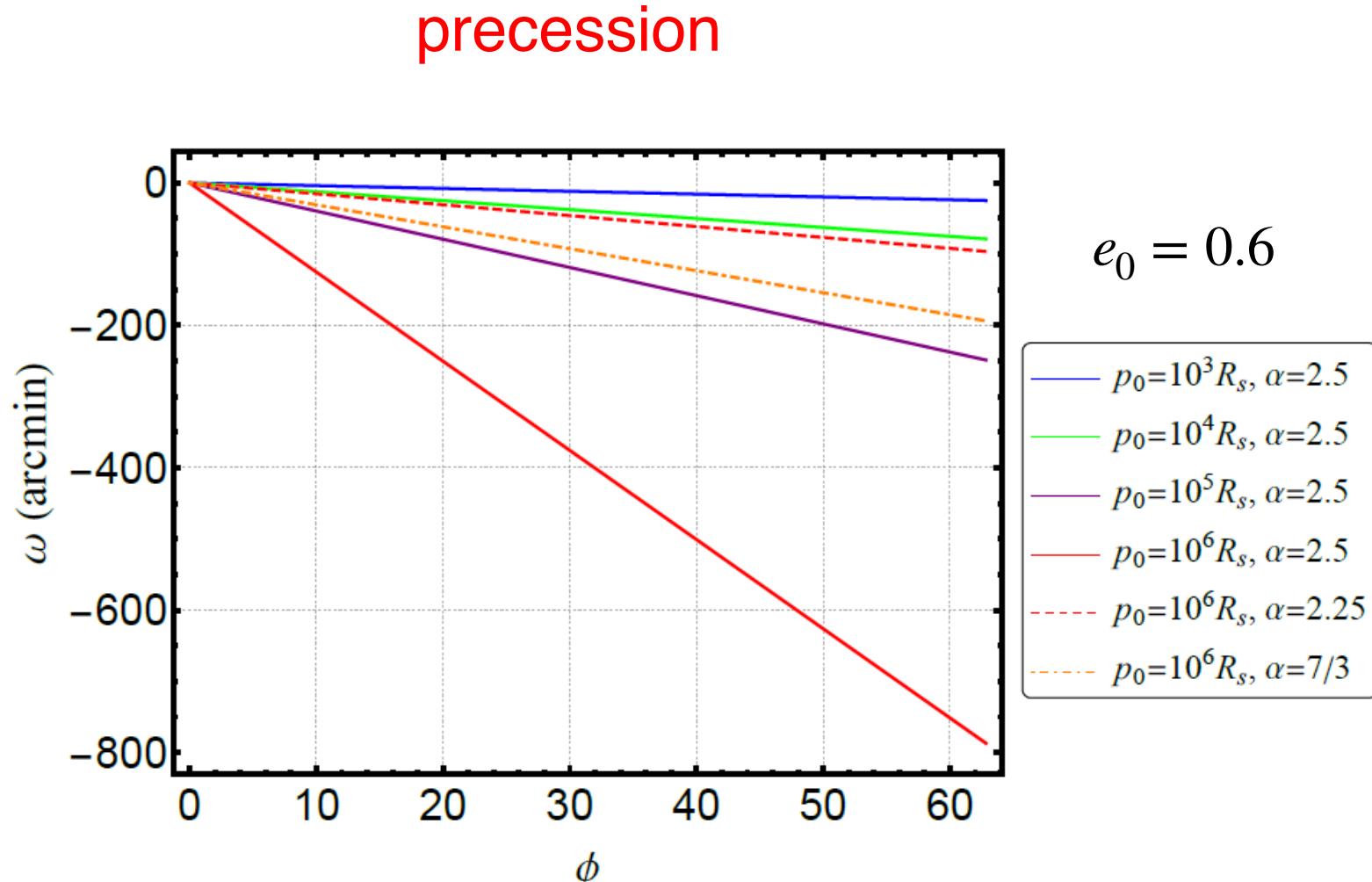
$$\mathbf{f}_G = - \frac{4\pi G \rho_{\text{sp}} r_{\text{sp}}^\alpha}{(3 - \alpha) r^{\alpha-1}} \mathbf{n}$$

$$\Delta p = 0 \quad \text{conservative}$$

$$\Delta e = 0$$

$$\Delta \omega_{\text{DM}} = \frac{4\pi p^{3-\alpha} \rho_{\text{sp}} r_{\text{sp}}^\alpha}{(3 - \alpha) M_{\text{eff}}} W_{\text{DM}}(e)$$

$$W_{\text{DM}}(e) = \int_0^{2\pi} \cos \phi (1 + e \cos \phi)^{\alpha-3} e^{-1} d\phi$$



Dynamic Friction

V. Cardoso, et al. 2020

$$\mathbf{f}_{\text{DF}} = -\frac{4\pi G^2 \mu \rho_{\text{DM}} I_\nu}{v^3} \mathbf{v}$$

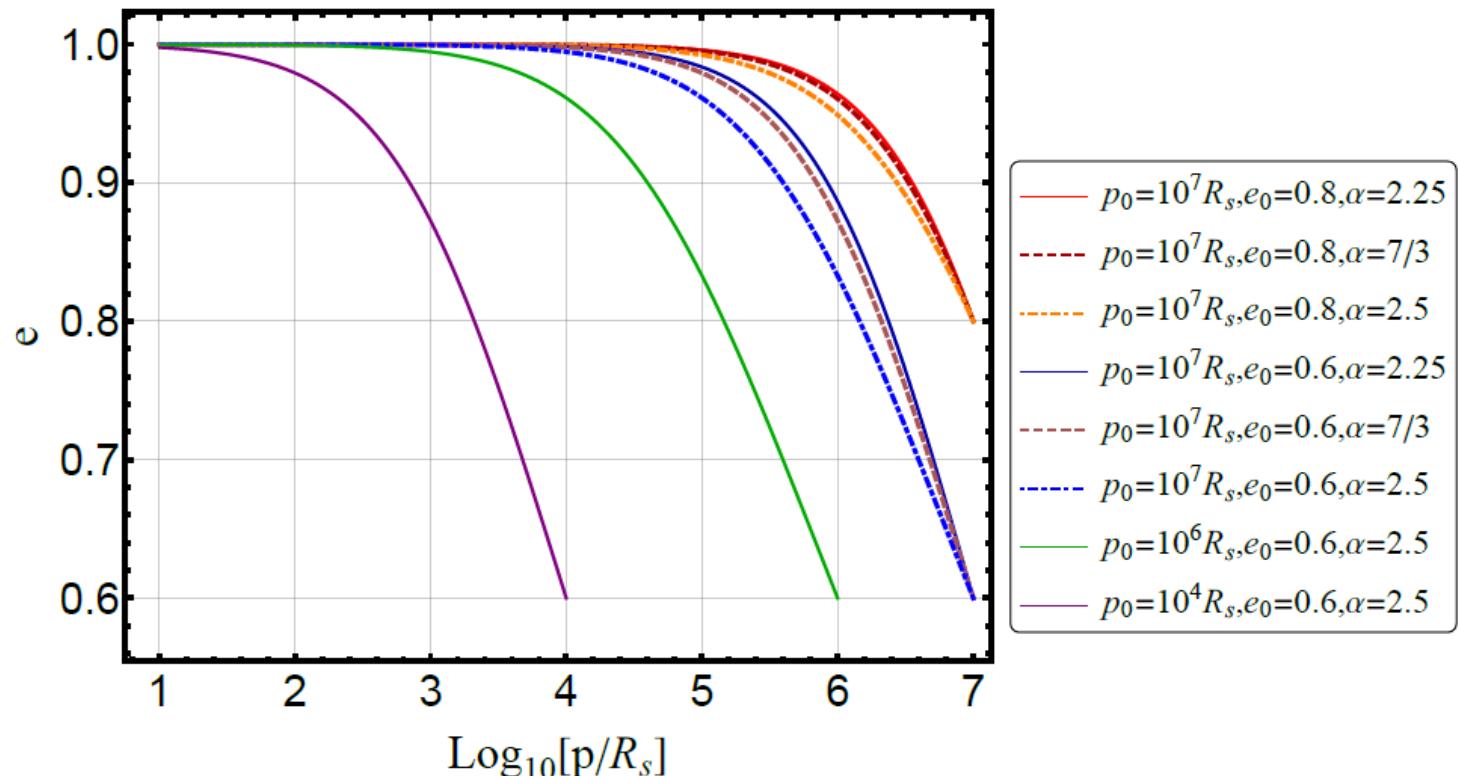
$$\left\langle \frac{dp}{d\phi} \right\rangle_{\text{DF}} = -\frac{4\mu \rho_{\text{sp}} r_{\text{sp}}^\alpha I_\nu}{M^2} p^{4-\alpha} g(e) < 0$$

$$\left\langle \frac{de}{d\phi} \right\rangle_{\text{DF}} = -\frac{4\mu \rho_{\text{sp}} r_{\text{sp}}^\alpha I_\nu}{M^2} p^{3-\alpha} f(e) > 0$$

$$\left\langle \frac{d\omega}{d\phi} \right\rangle_{\text{DF}} = 0 \quad \text{dissipative}$$

$$g(e) = \int_0^{2\pi} \frac{d\phi}{(1 + 2e \cos \phi + e^2)^{3/2}(1 + e \cos \phi)^{2-\alpha}}$$

$$f(e) = \int_0^{2\pi} \frac{(\cos \phi + e)d\phi}{(1 + 2e \cos \phi + e^2)^{3/2}(1 + e \cos \phi)^{2-\alpha}}$$



Bondi-Hoyle Accretion

H. Bondi & F. Hoyle, 1944
C. Macedo, et al. 2013

$$\dot{\mu} = 4\pi G^2 \lambda \frac{\mu^2 \rho_{\text{DM}}}{(v^2 + c_s^2)^{3/2}}$$

$$\mathbf{f}_a \simeq - \frac{4\pi G^2 \mu \lambda \rho_{\text{DM}}}{v^3} \mathbf{v}$$

dissipative

similar to dynamic friction

Radiation Reaction

T. Damour & N. Deruelle, 1981

$$\mathbf{f}_{\text{GW}} = \frac{8}{5} \frac{G^2 M \mu}{c^5 r^3} \left[\left(3v^2 + \frac{17}{3} \frac{Gm}{r} \right) \dot{r} \mathbf{n} - \left(v^2 + 3 \frac{Gm}{r} \right) \mathbf{v} \right]$$

$$\left\langle \frac{dp}{d\phi} \right\rangle_{\text{GW}} = - \frac{8}{5} \eta \frac{(Gm)^{5/2}}{c^5 p^{3/2}} (8 + 7e^2) < 0$$

$$\left\langle \frac{de}{d\phi} \right\rangle_{\text{GW}} = - \frac{8}{5} \eta \frac{(Gm)^{5/2}}{c^5 p^{5/2}} \left(\frac{304}{24} e + \frac{121}{24} e^3 \right) < 0$$

$$\left\langle \frac{d\omega}{d\phi} \right\rangle_{\text{GW}} = 0$$

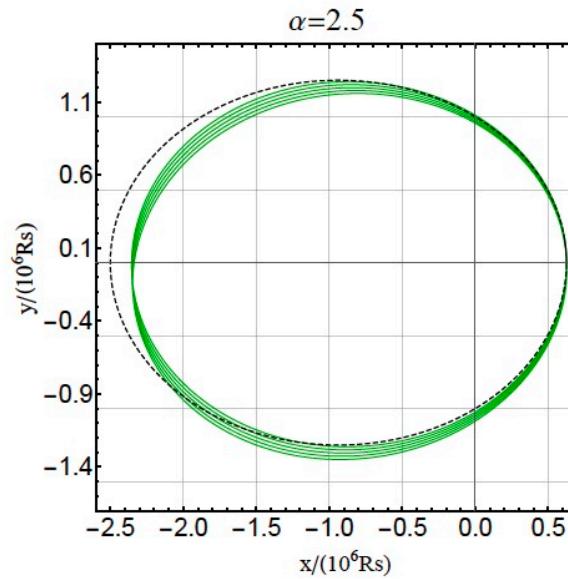
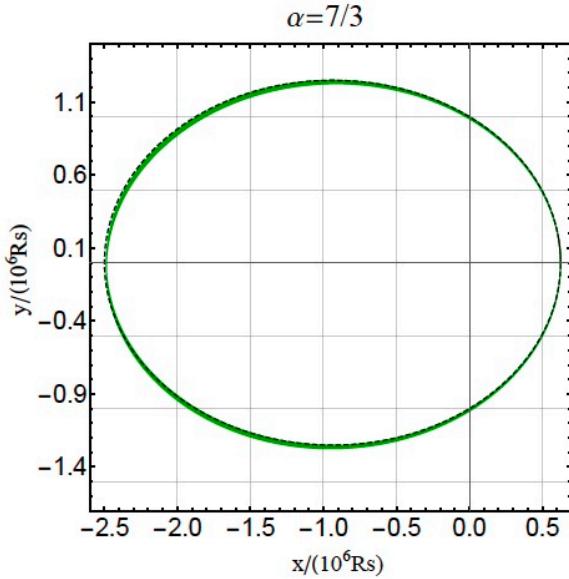
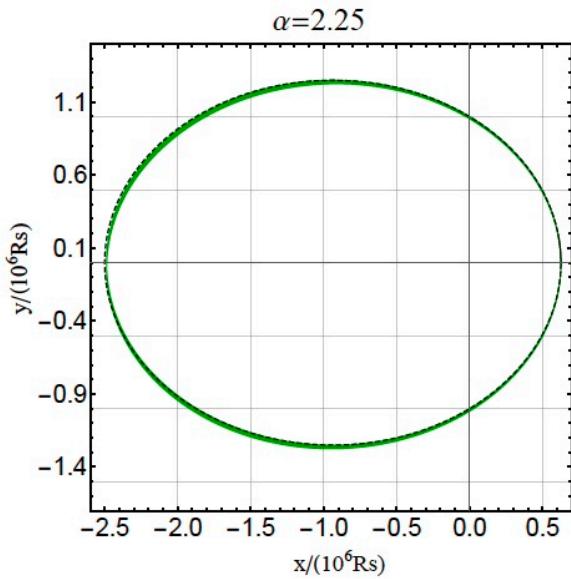
dissipative

Large Orbital Distance

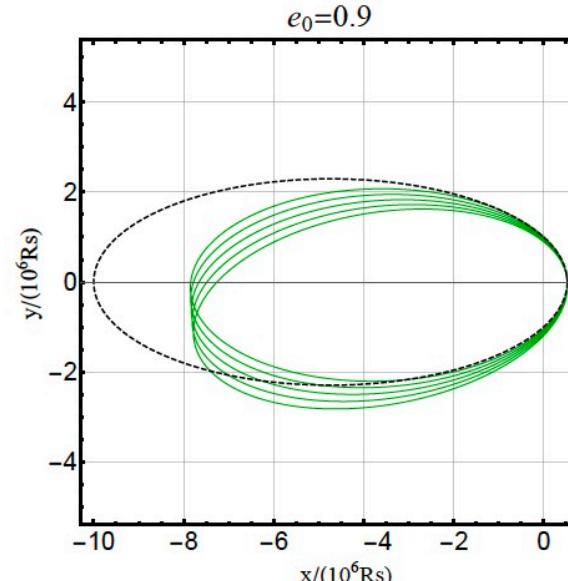
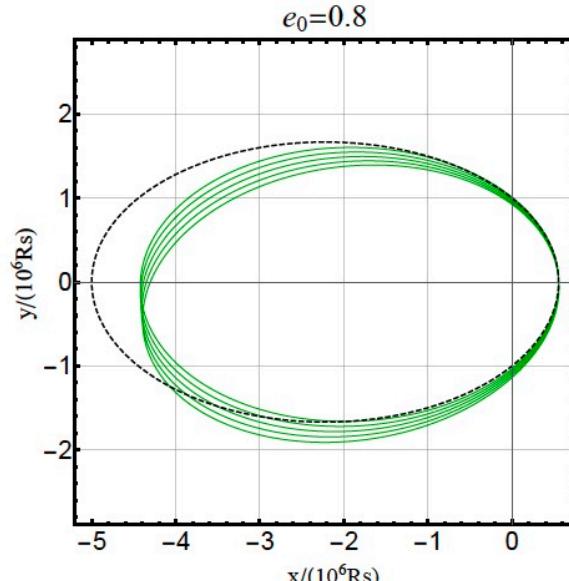
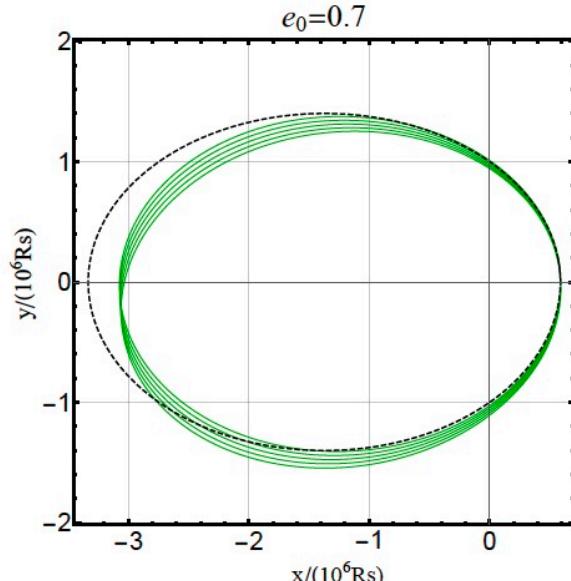
$$p_0 = 10^6 R_s$$

period~100 years

$$e_0 = 0.6$$



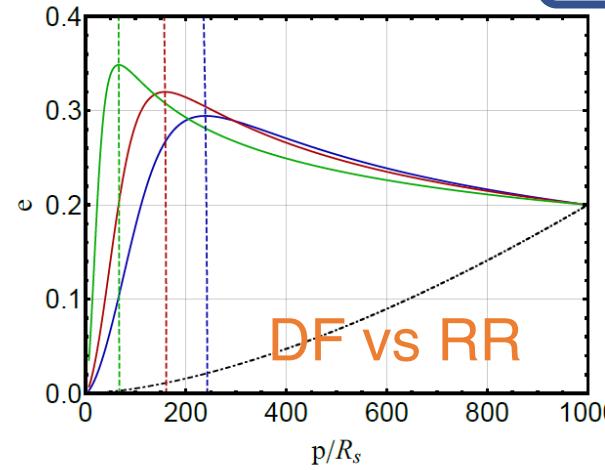
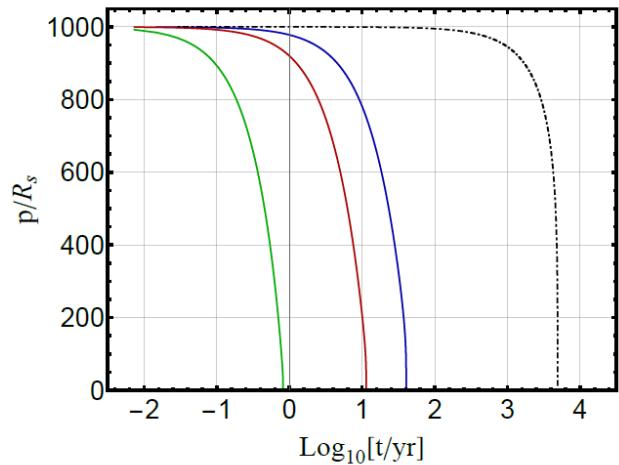
$$\alpha = 2.5$$



Small Orbital Distance

$$p_0 = 10^3 R_s$$

orbital evolution



$$p : 10^3 R_s \rightarrow 10 R_s$$

evolution time (yr)

e	No DM	$\alpha = 2.25$	$\alpha = 7/3$	$\alpha = 2.5$
0	4829	41.0	11.5	0.813
0.2	4901	40.4	11.4	0.815
0.4	5178	38.6	11.1	0.826
0.6	5928	35.6	10.5	0.848
0.8	8354	30.3	9.5	0.879
0.9	12625	25.5	8.4	0.898

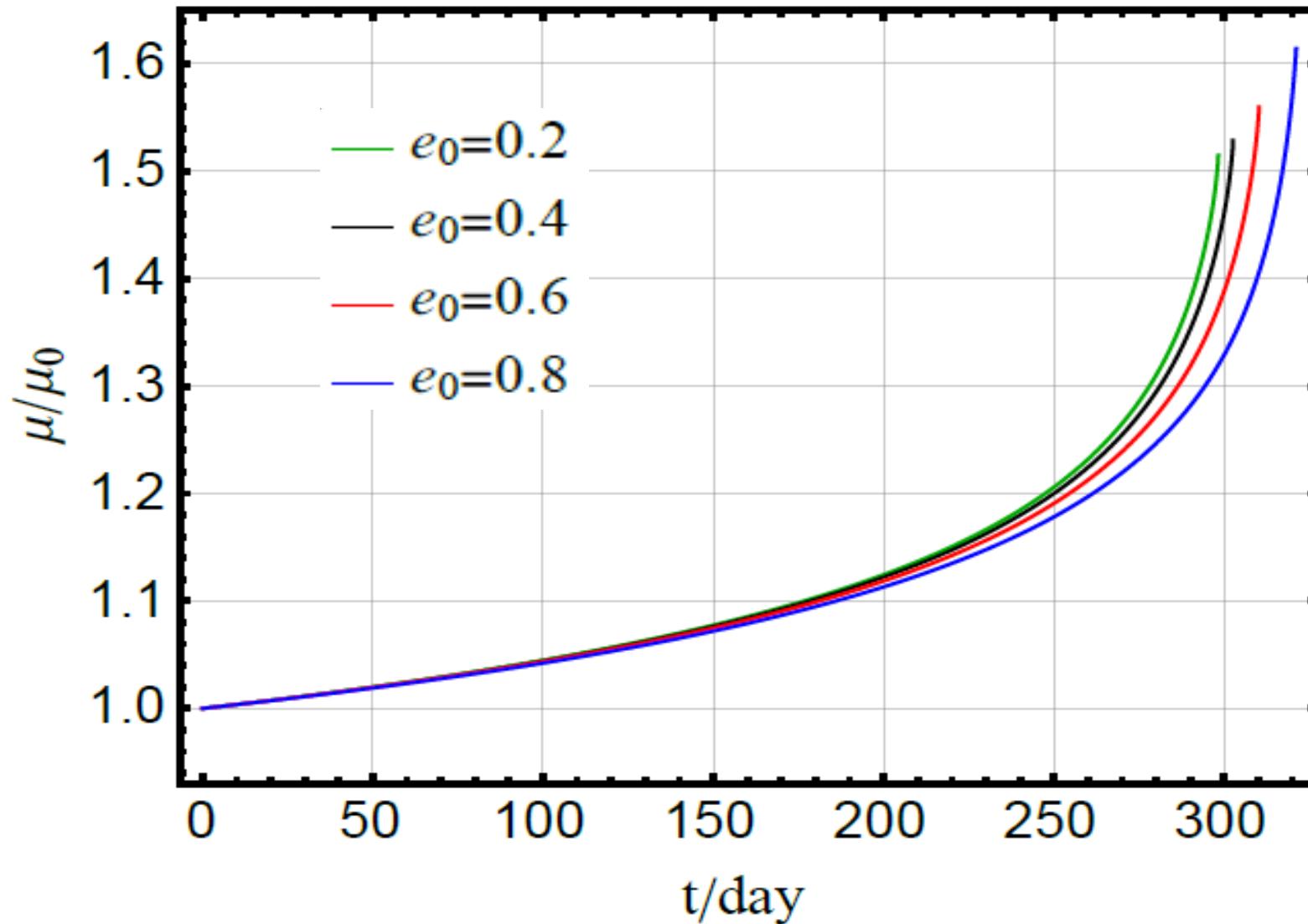
acceleration

Small Orbital Distance

$$p_0 = 10^3 R_s$$

accretion

$$\alpha = 2.5$$

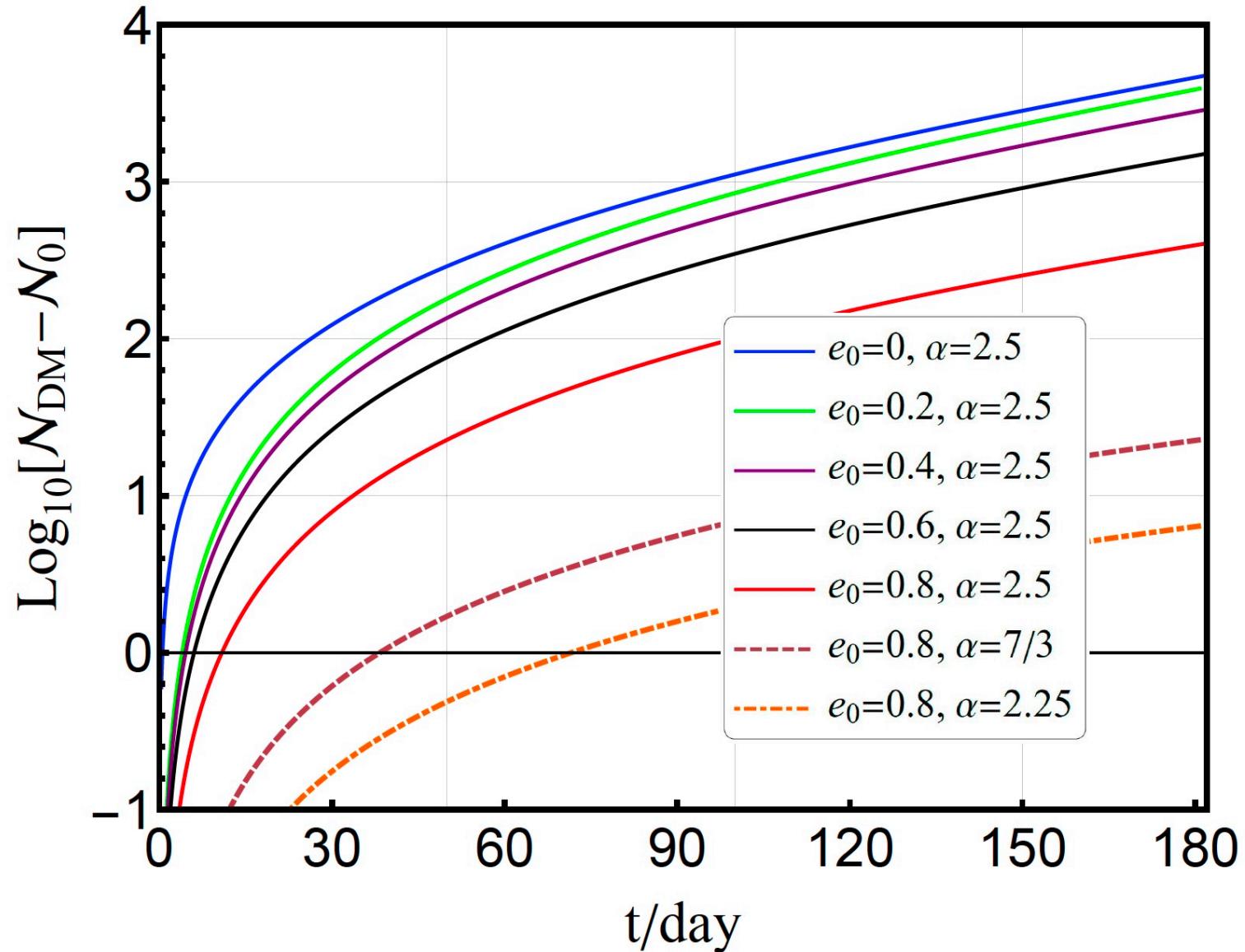


Small Orbital Distance

orbital cycles

$$\mathcal{N} = \frac{1}{2\pi} \int_{t_i}^{t_f} f(t) dt$$

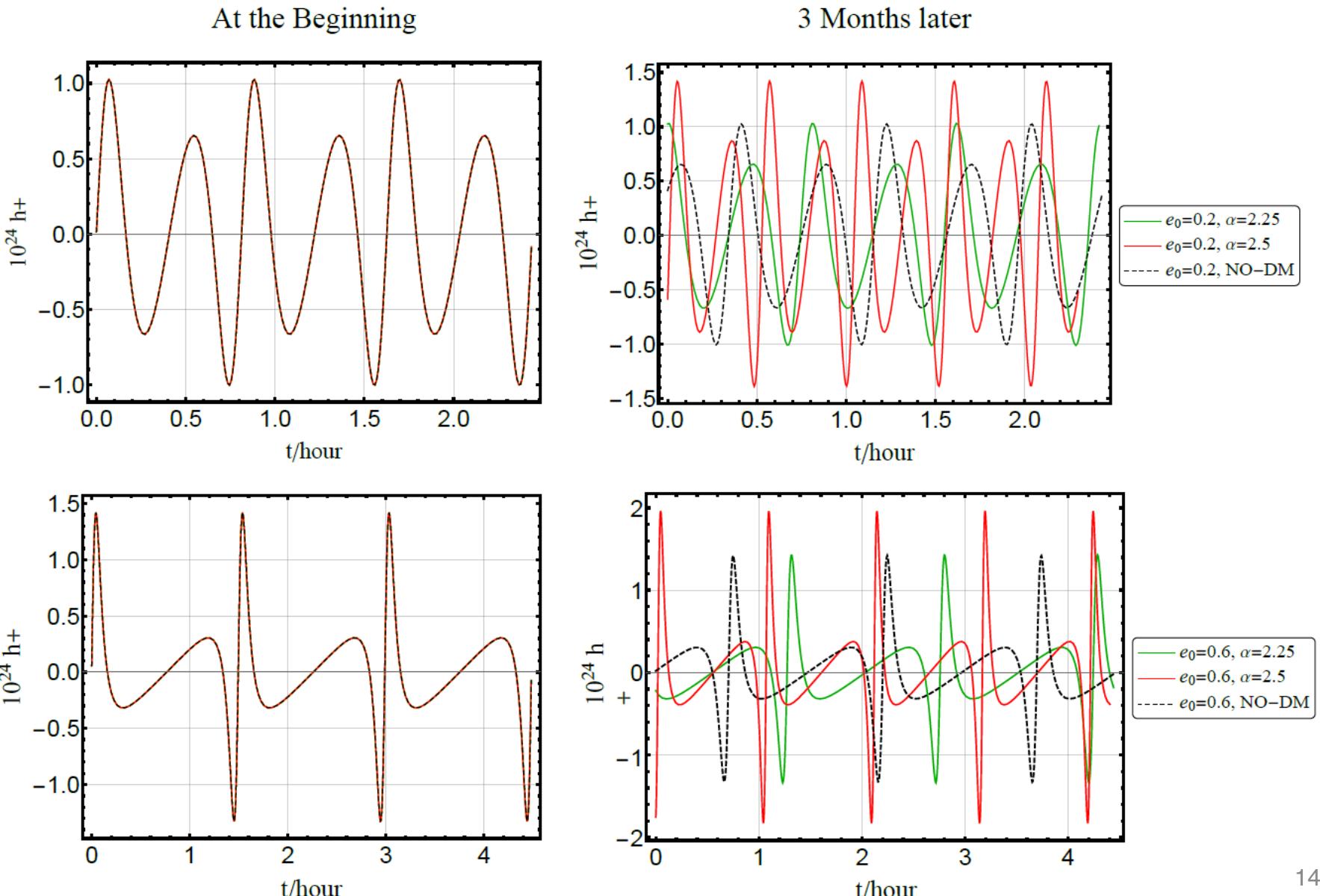
$$p_0 = 10^3 R_s$$



Small Orbital Distance

$$p_0 = 10^3 R_S$$

GW waveform



Small Orbital Distance

one-year evolution before $10R_S$

$\alpha = 7/3$

$D_L = 100\text{Mpc}$

e_0	SNR_0	SNR_D	Mismatch	D_{\max}
0.2	34.13	42.97	0.99992	358.1
0.4	34.19	47.74	0.99944	397.8
0.6	34.44	36.71	0.99965	305.9

N. Dai, et al. 2022

$$\text{Mismatch} = 1 - \max_{t_c, \phi_c} \frac{(\tilde{h}_0, \tilde{h}_{\text{DM}})}{\sqrt{(\tilde{h}_0, \tilde{h}_0)(\tilde{h}_{\text{DM}}, \tilde{h}_{\text{DM}})}}$$

Summary

- ◆ The existence of dark matter significantly affects the evolution of IMRIs
- ◆ How to model these environmental effects will be essential in the future

Thank you for listening