

# Anisotropic stars made of exotic matter in light of the complexity factor formalism

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# Outline

- **Introduction**
- **Relativistic stars in General Relativity**
- **Anisotropic matter: Complexity factor**
- **Results and Discussion**
- **Take home message**

# Introduction

- 1- Celestial bodies are not always made of isotropic fluid only!
- 2- Anisotropic stars can be made, at least theoretically, of exotic matter (as dark matter).
- 3- To study anisotropic stars, usually it is necessary to consider a concrete form of the anisotropic factor (problem!)
- 4- To bypass that, we will consider the complexity factor formalism to close the system and obtain the physical properties of such anisotropic stars.

# Relativistic stars in General Relativity

In absence of cosmological constant term, the classical Einstein field equations acquire the simple form:

$$G^{\nu}_{\mu} = 8\pi G T^{\nu}_{\mu},$$

The covariant energy-momentum tensor can be expressed in local Minkowski coordinates as

$$T^{\mu}_{\nu} = \{\rho, P_r, P_{\perp}, P_{\perp}\},$$

and the resulting field equations take the form:

# Relativistic stars in General Relativity

$$\rho = -\frac{1}{8\pi} \left[ -\frac{1}{r^2} + e^{-\lambda} \left( \frac{1}{r^2} - \frac{\lambda'}{r} \right) \right],$$

$$P_r = -\frac{1}{8\pi} \left[ \frac{1}{r^2} - e^{-\lambda} \left( \frac{1}{r^2} + \frac{\nu'}{r} \right) \right],$$

$$P_{\perp} = \frac{1}{32\pi} e^{-\lambda} \left( 2\nu'' + \nu'^2 - \lambda'\nu' + 2\frac{\nu' - \lambda'}{r} \right),$$

where the derivatives with respect to  $r$  are denoted by primes.

# Relativistic stars in General Relativity

We can combine the last equations to produce the generalized Tolman-Oppenheimer-Volkoff equation, i.e.

$$-\frac{1}{2}\nu' (\rho + P_r) - P_r' + \frac{2}{r} (P_\perp - P_r) = 0.$$

And removing the  $\nu'$  - dependence, we finally obtain the familiar equation

$$P_r' = -\frac{(m + 4\pi P_r r^3)}{r(r - 2m)} (\rho + P_r) + \frac{2}{r} (P_\perp - P_r),$$

# Relativistic stars in General Relativity

In what follow, we will rewrite the energy-momentum tensor in a more convenient way:

$$T_{\nu}^{\mu} = \rho u^{\mu} u_{\nu} - P h_{\nu}^{\mu} + \Pi_{\nu}^{\mu},$$

we set the four-velocity and the four acceleration as follow:

$$u^{\mu} = (e^{-\frac{\nu}{2}}, 0, 0, 0), \quad a^{\alpha} = u^{\alpha}_{;\beta} u^{\beta},$$

Where the unique non--vanishing component of the acceleration is

$$a_1 = -\nu' / 2,$$

# Relativistic stars in General Relativity

Notice that the set  $\{\Pi_{\nu}^{\mu}, \Pi, h_{\nu}^{\mu}, s^{\mu}, P\}$  is taken according to

$$\Pi_{\nu}^{\mu} = \Pi \left( s^{\mu} s_{\nu} + \frac{1}{3} h_{\nu}^{\mu} \right),$$

$$\Pi = P_{\perp} - P_r,$$

$$h_{\nu}^{\mu} = \delta_{\nu}^{\mu} - u^{\mu} u_{\nu},$$

$$s^{\mu} = (0, e^{-\frac{\lambda}{2}}, 0, 0),$$

$$P \equiv \frac{1}{3} \left( P_r + 2P_{\perp} \right),$$



# Relativistic stars in General Relativity

with the properties

$$s^\mu u_\mu = 0, \quad s^\mu s_\mu = -1.$$

The problem should be supplemented using certain boundary conditions on the surface

$$e^{\nu_\Sigma} = 1 - \frac{2M}{R},$$

$$e^{-\lambda_\Sigma} = 1 - \frac{2M}{R},$$

$$[P_r]_\Sigma = 0.$$

# Anisotropic matter: Complexity factor

The complexity factor describes the influence of the local anisotropy of pressure and density inhomogeneity on the Tolman mass

$$m_T = (m_T)_\Sigma \left( \frac{r}{r_\Sigma} \right)^3 + r^3 \int_r^{r_\Sigma} \frac{e^{(\nu+\lambda)/2}}{\tilde{r}} Y_{TF} d\tilde{r},$$

where

$$Y_{TF} = 8\pi\Pi - \frac{4\pi}{r^3} \int_0^r \tilde{r}^3 \rho' d\tilde{r},$$

# Anisotropic matter: Complexity factor

To define the complexity factor, we need to perform the orthogonal decomposition of the Riemann tensor

$$Y_{\alpha\beta} = R_{\alpha\gamma\beta\delta}u^\gamma u^\delta,$$

$$Z_{\alpha\beta} = R^*_{\alpha\gamma\beta\delta}u^\gamma u^\delta = \frac{1}{2}\eta_{\alpha\gamma\epsilon\mu}R^{\epsilon\mu}{}_{\beta\delta}u^\gamma u^\delta,$$

$$X_{\alpha\beta} = R^*_{\alpha\gamma\beta\delta}u^\gamma u^\delta = \frac{1}{2}\eta_{\alpha\gamma}{}^{\epsilon\mu}R^*_{\epsilon\mu\beta\delta}u^\gamma u^\delta,$$

where

$$R^*_{\alpha\beta\gamma\delta} = \frac{1}{2}\eta_{\epsilon\mu\gamma\delta}R_{\alpha\beta}{}^{\epsilon\mu},$$

# Anisotropic matter: Complexity factor

We rewrite last expressions in term of the physical variables as follow

$$Y_{\alpha\beta} = \frac{4\pi}{3}(\rho + 3P)h_{\alpha\beta} + 4\pi\Pi_{\alpha\beta} + E_{\alpha\beta},$$

$$Z_{\alpha\beta} = 0,$$

$$X_{\alpha\beta} = \frac{8\pi}{3}\rho h_{\alpha\beta} + 4\pi\Pi_{\alpha\beta} - E_{\alpha\beta}.$$

where

$$E_{\alpha\beta} = C_{\alpha\gamma\beta\delta}u^\gamma u^\delta,$$

# Anisotropic matter: Complexity factor

We rewrite last expressions in term of the physical variables as follow

$$E_{\alpha\beta} = E \left( s_{\alpha} s_{\beta} + \frac{1}{3} h_{\alpha\beta} \right),$$

$$E = -\frac{e^{-\lambda}}{4} \left[ \nu'' + \frac{\nu'^2 - \lambda' \nu'}{2} - \frac{\nu' - \lambda'}{r} + \frac{2(1 - e^{\lambda})}{r^2} \right],$$

The following properties must be satisfied:

$$E^{\alpha}_{\alpha} = 0, \quad E_{\alpha\gamma} = E_{(\alpha\gamma)}, \quad E_{\alpha\gamma} u^{\gamma} = 0,$$

# Anisotropic matter: Complexity factor

Finally, last tensors can be represented in term of alternative scalar functions, i.e.,

$$X_T = 8\pi\rho,$$

$$X_{TF} = \frac{4\pi}{r^3} \int_0^r \tilde{r}^3 \rho' d\tilde{r},$$

$$Y_T = 4\pi(\rho + 3P_r - 2\Pi),$$

$$Y_{TF} = 8\pi\Pi - \frac{4\pi}{r^3} \int_0^r \tilde{r}^3 \rho' d\tilde{r},$$

# Anisotropic matter: Complexity factor

Demanding a vanishing complexity factor, we can obtain the anisotropic factor by solving

$$\Pi(r) = \frac{1}{2r^3} \int_0^r \tilde{r}^3 \rho'(\tilde{r}) d\tilde{r}.$$

Given that a profound comprehension of the idea of complexity is still under construction, the connection between any equation of state and the definition of complexity factor is still missing.

# Anisotropic matter: Complexity factor

In what follow, we will take a Chaplyin equation-of-state to close the system

$$P_r = A^2 \rho(r) - B^2 \rho(r)^{-1},$$

and with help of a vanishing complexity factor, we solve the problem numerically.

Also, for comparison, we will check our solution vs the more conventional approach (i.e., assuming a concrete form of the anisotropic factor).



# Results and Discussion

We will use the following numerical values:

Case # 1:

$$A = \sqrt{0.400}, \quad B = 0.230 \times 10^{-3} km^{-2}.$$

Case # 2:

$$A = \sqrt{0.425}, \quad B = 0.215 \times 10^{-3} km^{-2}.$$

Case # 3:

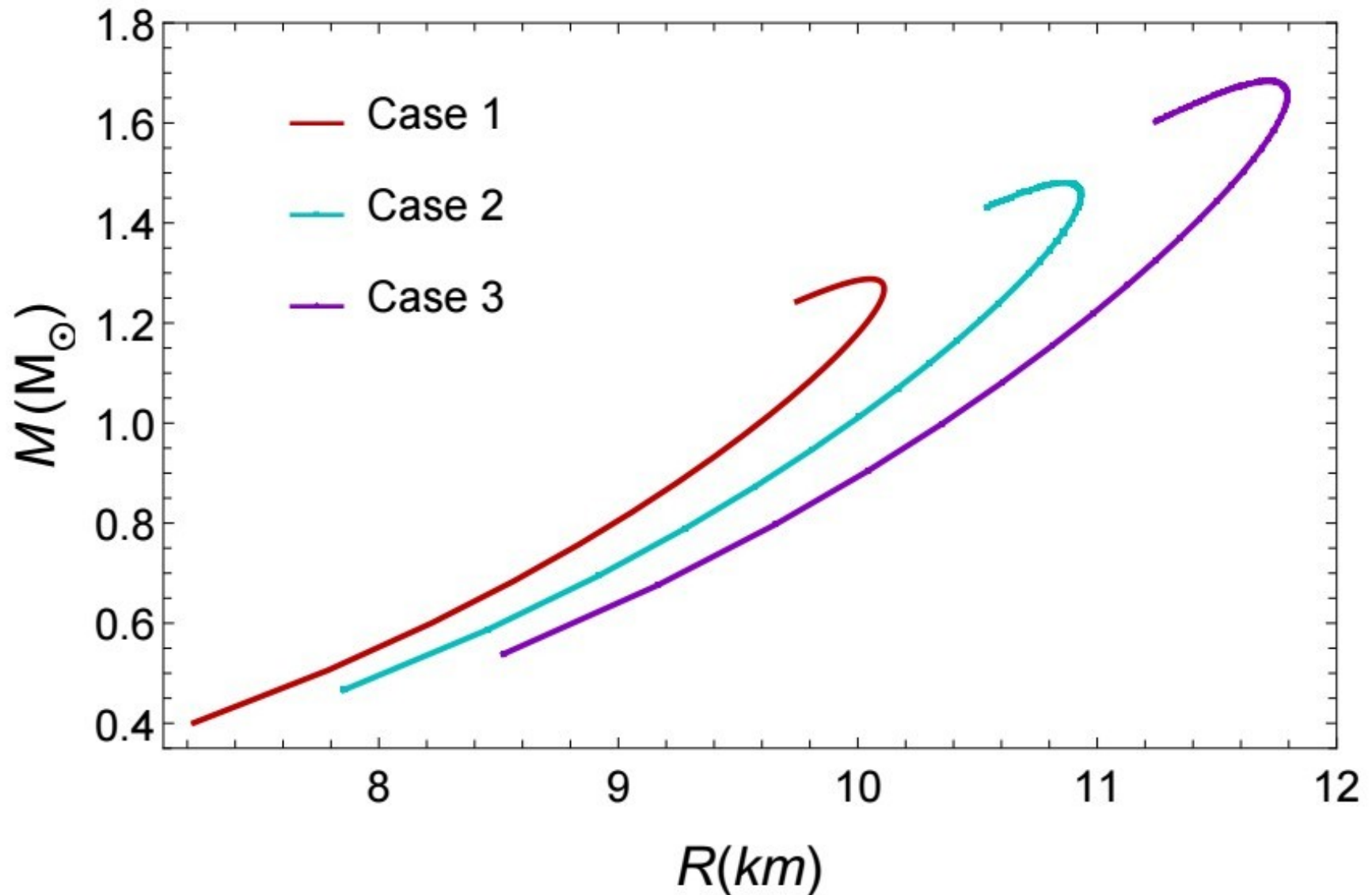
$$A = \sqrt{0.450}, \quad B = 0.200 \times 10^{-3} km^{-2}.$$

# Results and Discussion

## 1- Complexity Factor Formalism

$$\Pi(r) = \frac{1}{2r^3} \int_0^r \tilde{r}^3 \rho'(\tilde{r}) d\tilde{r}.$$

# Results and Discussion



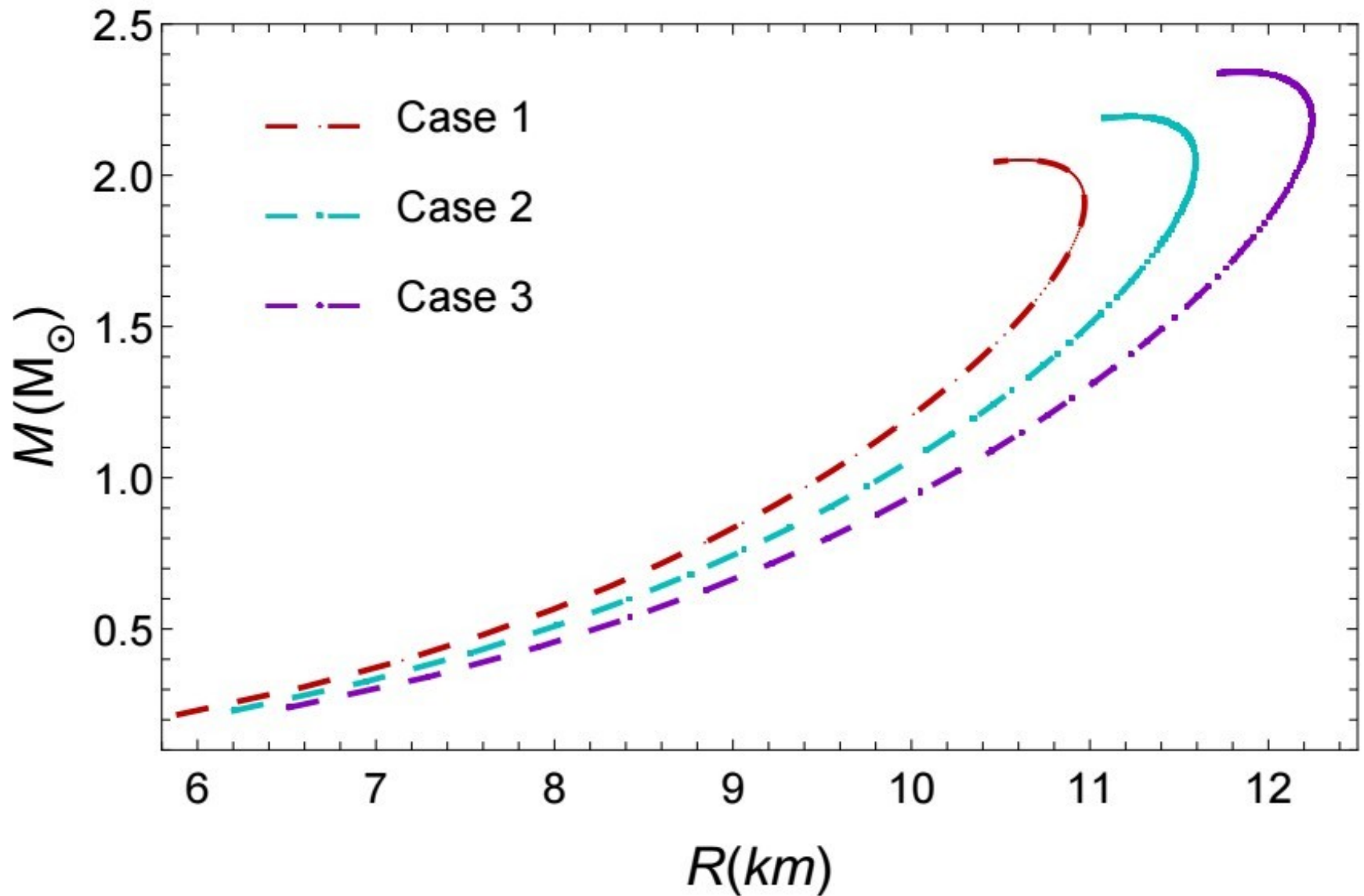
# Results and Discussion

## 2- Conventional Approach

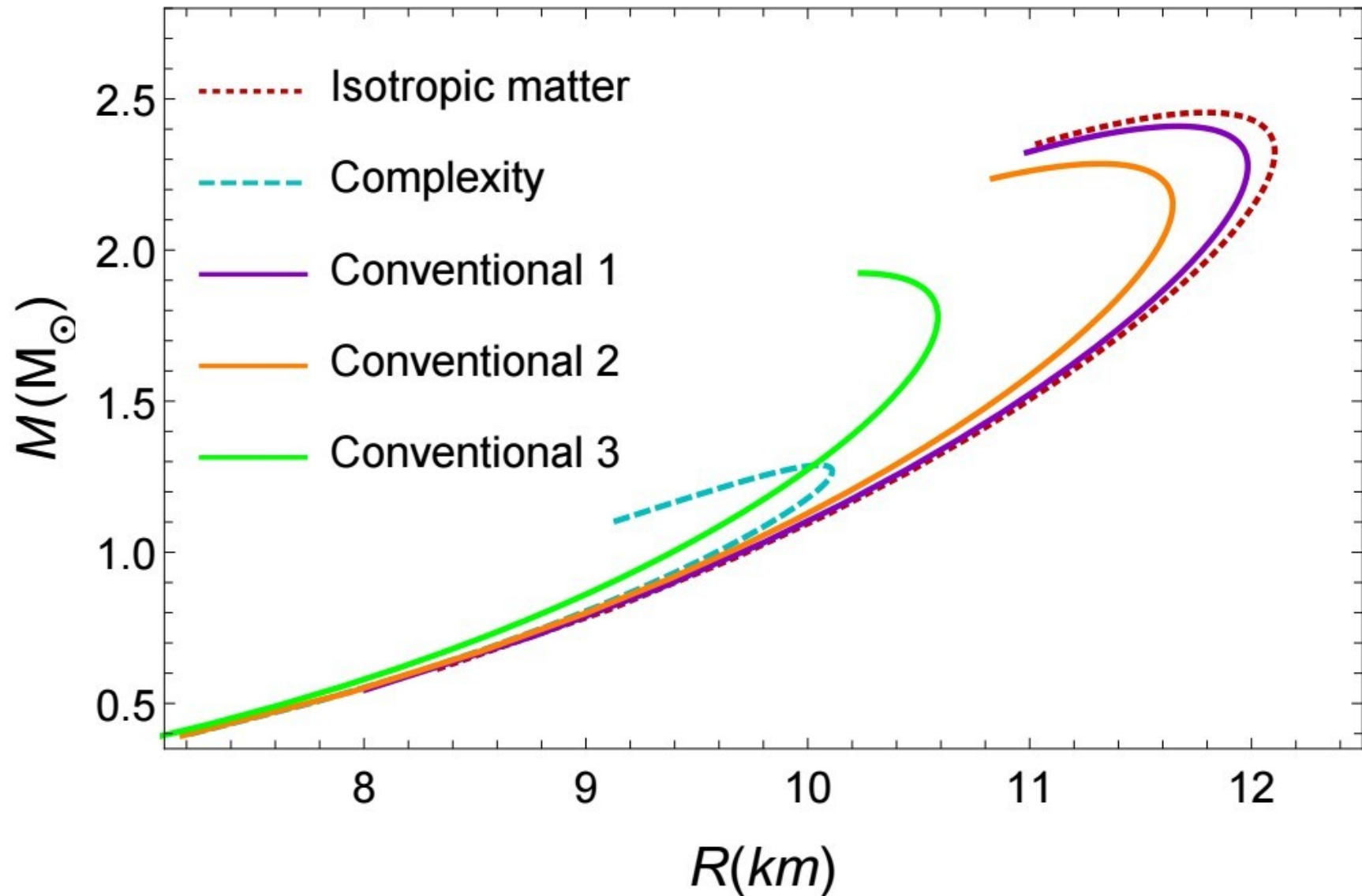
$$\Pi(r) = - \left( \frac{r}{a} \right)^2 \rho(r).$$

$$a \equiv a_{\text{large}} = 30 \text{ km.}$$

# Results and Discussion



# Results and Discussion



# Take home message

The complexity factor formalism predicts smaller and lighter objects compared to the conventional approach, despite both cases having a negative anisotropic factor.