Slow Beam Degradation due to Electron Clouds at the LHC: Observations and Modelling

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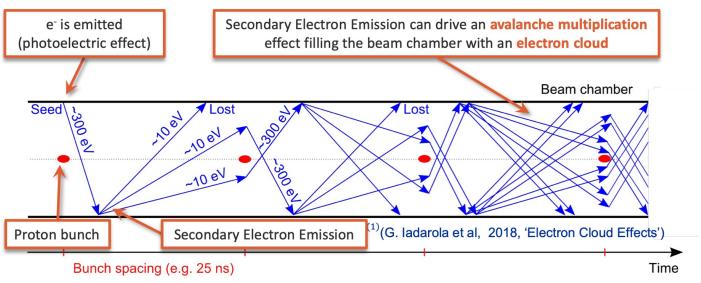
Outline

- 1. Introduction to electron clouds
- 2. Observations in LHC during Run 2
- 3. Progress in simulations
 - a) Description of e-cloud interaction
 - b) Simulation results

Outline

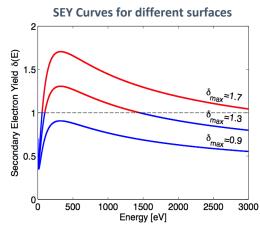
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Electron clouds

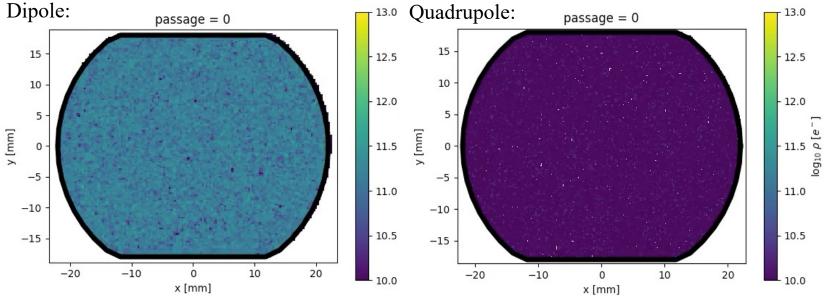


- Electrons are introduced into the beam chamber (residual gas ionization / synchr. rad. + photoelectric effect)
- 2. Electrons are accelerated by passing bunches and impact on beam chamber.
 - Depending on energy of electron and Secondary Emission Yield of surface, electrons can be emitted.
 SEY Curves for different surfaces

If conditions allow, **electrons multiply exponentially!**

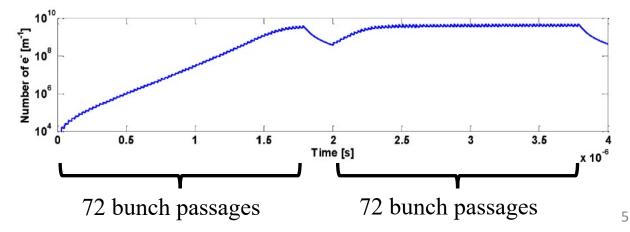


Electron clouds

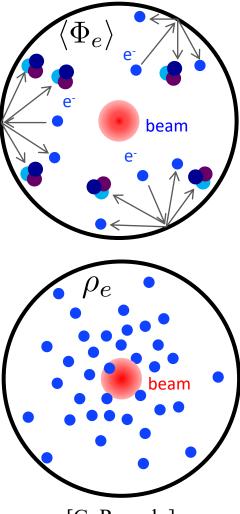


• Electrons multiply until a saturation is reached.

- Number of electrons quickly decays when bunches are not passing.
- Magnetic fields strongly affect the e-cloud.



Electron cloud effects



[G. Rumolo]

The electron flux to the wall is responsible for

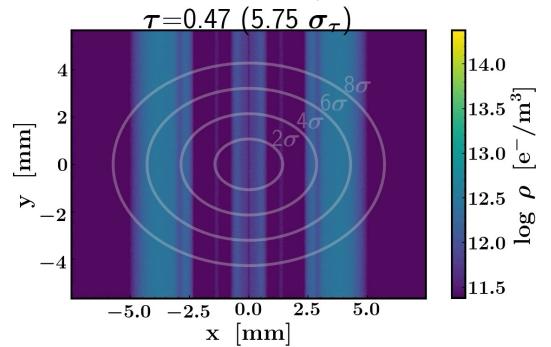
- Dynamic pressure rise
- Heat deposition
- Spurious signal for beam instrumentation

The electron density inside the chamber causes:

- Tune shift along the bunch train
- Synchronous phase shift along bunch train.
- Coherent beam instabilities (single and coupled bunch)
- **Incoherent effects** (beam lifetime degradation and slow emittance growth)

Electron cloud pinch

Incoherent electron cloud effects concern the motion of single particles under the influence of the non-linear forces induced by the electrons.

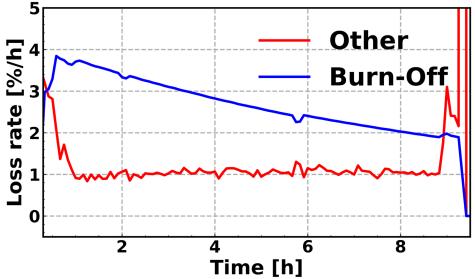


- Motion of electrons is very complex → Complex electron densities → complex induced forces.
- Protons from the beam are "moving" within these complex forces due to:
 - Betatron oscillations: up-down, left-right
 - Synchrotron oscillations: back-forth in "time"
- \rightarrow Increase of proton oscillation amplitude \rightarrow losses + emittance growth.

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LHC observations during a typical fill



Losses come from:

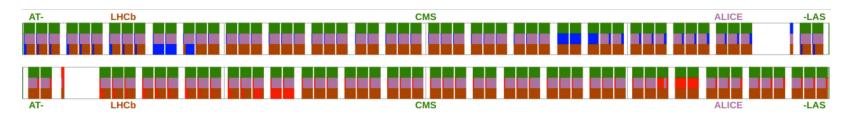
- Luminosity burn-off that decreases gradually.
- Continuous rate of additional losses.

Luminosity (ATLAS + CMS)

$$\left(-\frac{dI}{dt}\right)_{\text{other}} = \left(-\frac{dI}{dt}\right)_{\text{total}} - \sigma_{\text{inel.}} \cdot \mathcal{L}$$

Total loss rate (Fast Beam Current Transformer)

Filling scheme

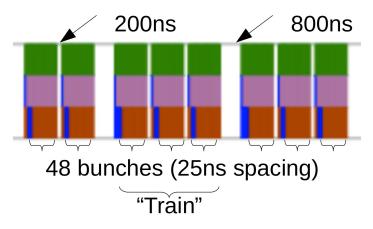


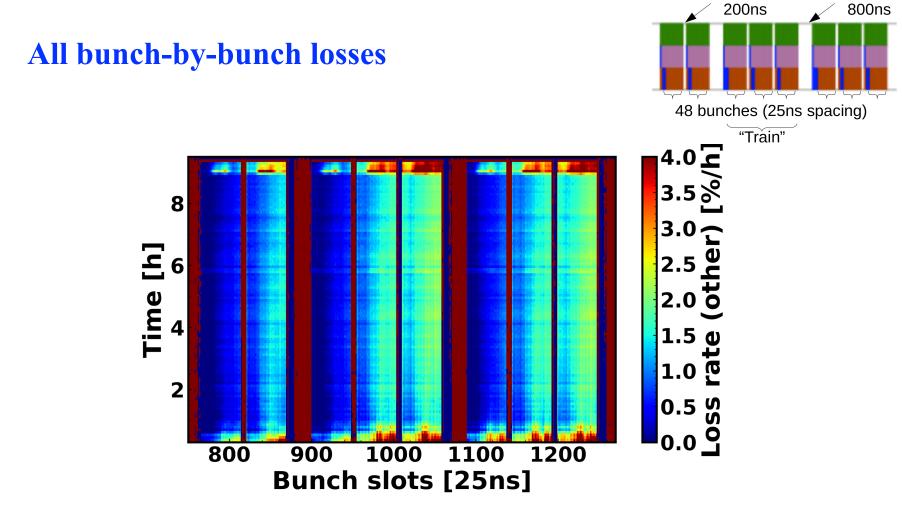
Standard 2018 Physics filling scheme (2556 bunches) [lpc.web.cern.ch]

Beam is composed of repeating patterns (trains):

- 2x48 bunches,
- 3x48 bunches.

Magnification:





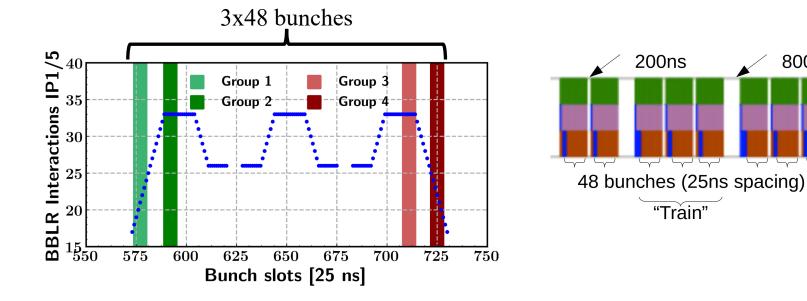
Global picture: Fairly constant loss rate (Corrected for burn-off).

• Grows from head to tail of each train

Number of BBLR interactions

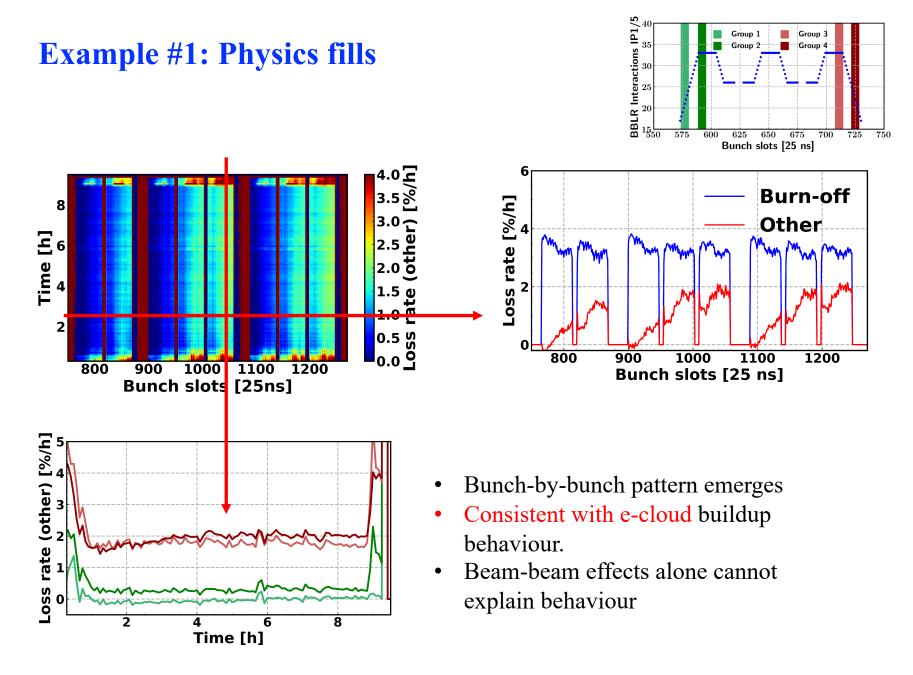
The only other non-linearities that depend on the bunch position are the Beam-Beam Long-Range (BBLR) interactions.

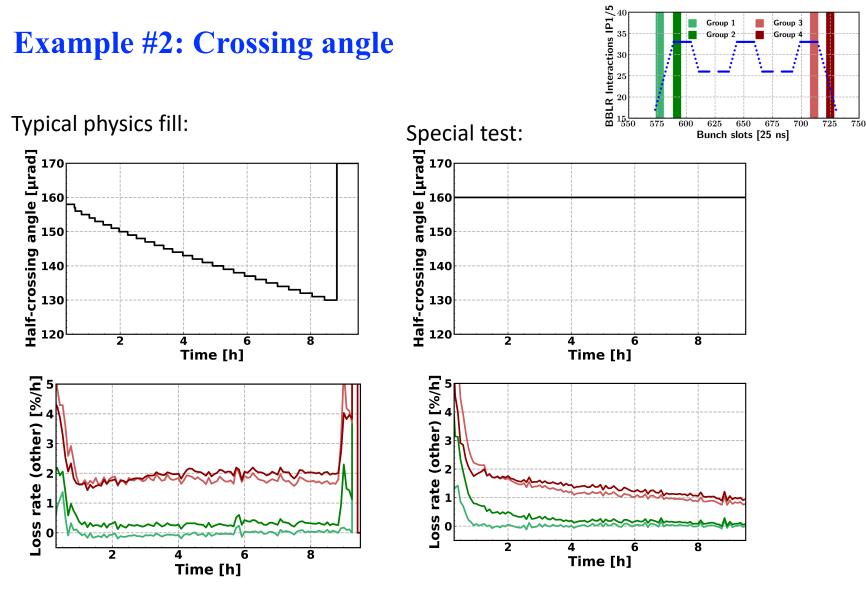
Number of long-range encounters changes for each bunch in the filling scheme.



- Group 1: Few BBLR, reduced e-cloud effects
- Group 2: Max BBLR, reduced e-cloud effects
- Group 3: Max BBLR, stronger e-cloud effects
- **Group 4**: Few BBLR, stronger e-cloud effects

800ns

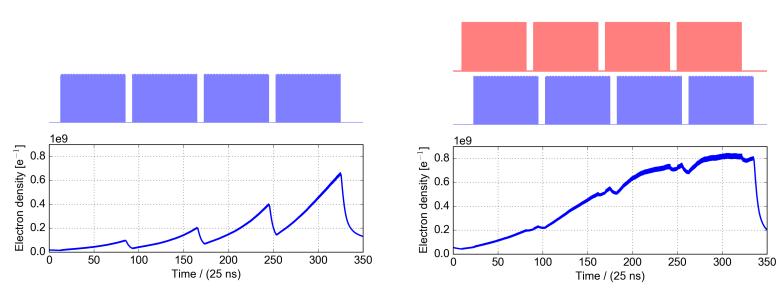




- A reduced crossing angle typically enhances BBLR interactions.
- In this case, it enhances the e-cloud pattern losses.

Example #3: Buildup simulations in Inner Triplet quadrupoles

One beam:



Two beams:

- One beam: In the small 200 ns between batches, the electron cloud decays significantly.
- Two beams: Beams are not synchronized and the e-cloud does not decay.

Example #3: Buildup simulations in Inner Triplet quadrupoles

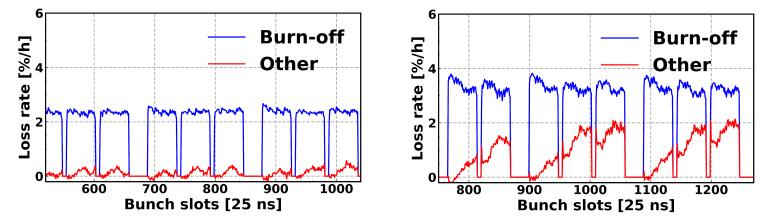
Two beams: 1e9 1e9 Electron density [e⁻¹] 0.8 0.6 0.4 0.2 0.0 0.0 250 300 350 í٥ 50 100 150 200 50 100 150 200 250 300 350 0 Time / (25 ns) Time / (25 ns) 6 6 **Burn-off Burn-off** Loss rate [%/h] Loss rate [%/h] N b Other Other 4 1300 1400 1500 1600 1700 800 900 1000 1100 1200 Bunch slots [25 ns] Bunch slots [25 ns]

The bunch-by-bunch pattern of the losses resembles the e-cloud buildup simulations of the Inner Triplet quadrupoles.

One beam:

Example #4: Measurements with different betatron functions

 $\beta^* = 65 \text{ cm}, \phi = 120 \mu \text{rad}$ Large ATS telescope¹ \rightarrow \rightarrow enhancement of arc beta functions $\beta^* = 30 \text{ cm}, \phi = 150 \mu \text{rad}$ Moderate ATS telescope



- Decreasing β in the inner triplet quadrupoles should reduce effect of the e-cloud in the inner triplet.
- Increasing β in arcs should enhances e-cloud effect: no significant losses.

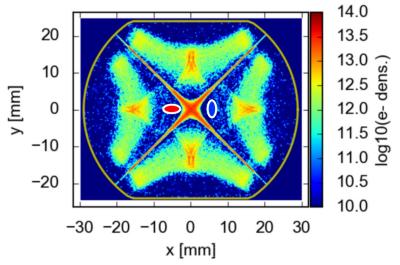
Summary - Observations

Electron cloud related losses are enhanced when:

- 1. reducing β^* (increasing β in IT)
- 2. reducing crossing angle (changes closed orbit in IT)
- 3. Two beams are present (enhanced buildup in IT) but not when:

but not when:

4. Increasing β in arcs

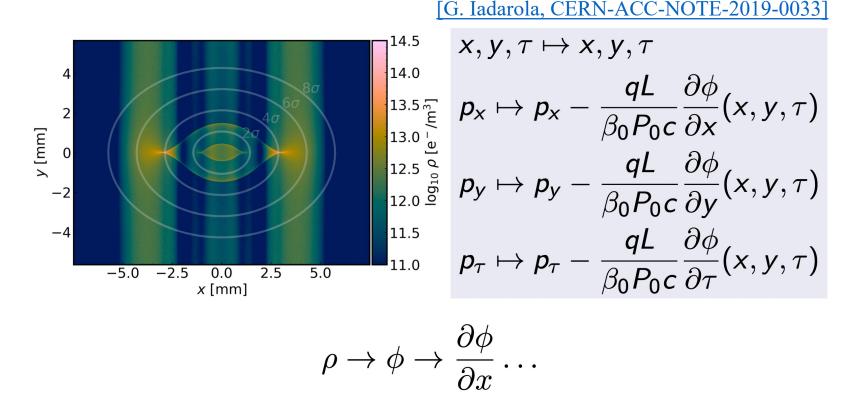


All observations point to the Inner Triplet Quadrupoles. Good news: HL-LHC Inner Triplet will have a-C coating to suppress e-cloud.

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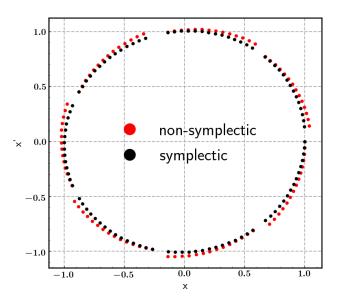
Introduction to simulations



- Complex time-dependent e-cloud density \rightarrow complex time-dependent forces
- Slow incoherent effects → e-cloud can be re-used = weak-strong approximaton (no self-consistency)
- But: e-cloud potential (PIC) is defined on a 3D grid. Needs to be interpolated.

Symplecticity

- Numerical methods in solving Hamiltonian systems can break the symplectic condition, making them less accurate at long timescales. (Millions of turns)
- Typically important to preserve symplecticity, even at the expense of accuracy.



• Interpolation scheme should guarantee symplecticity.

In our case, symplecticity:
$$\frac{\partial^2 \phi}{\partial x \partial y} = \frac{\partial^2 \phi}{\partial y \partial x}$$

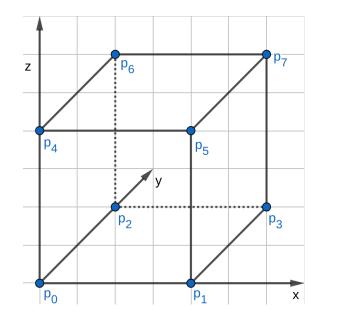
• Linear interpolation is not symplectic.

$$\begin{aligned} x, y, \tau &\mapsto x, y, \tau \\ p_x &\mapsto p_x - \frac{qL}{\beta_0 P_0 c} \frac{\partial \phi}{\partial x}(x, y, \tau) \\ p_y &\mapsto p_y - \frac{qL}{\beta_0 P_0 c} \frac{\partial \phi}{\partial y}(x, y, \tau) \\ p_\tau &\mapsto p_\tau - \frac{qL}{\beta_0 P_0 c} \frac{\partial \phi}{\partial \tau}(x, y, \tau) \end{aligned}$$

Tricubic interpolation

Given a regular 3D grid of any function f^{ijk} , we interpolate locally in a way that the following quantities are continuous globally.

$$\left\{f, \frac{\partial f}{\partial x}, \frac{\partial f}{\partial y}, \frac{\partial f}{\partial z}, \frac{\partial^2 f}{\partial x \partial y}, \frac{\partial^2 f}{\partial x \partial z}, \frac{\partial^2 f}{\partial y \partial z}\right\}$$



Lekien and Marsden^{*} proved that it is possible to meet this condition by using a tricubic interpolation scheme.

$$f(x, y, z) = \sum_{i=0}^{3} \sum_{j=0}^{3} \sum_{k=0}^{3} a_{ijk} x^{i} y^{j} z^{k}$$

The coefficients a_{ijk} change from cell to cell but required quantities stay continuous across the cells.

• Analytical derivatives for interaction.

Small digression

Demonstration of symplecticity violation/preservation

Consider the Hamiltonian:
$$H = rac{p_1^2}{2} + rac{p_2^2}{2} + \phi(q_1,q_2)$$
 with $\phi(q_1,q_2) = e^{q_1-q_2}$

These quantities are conserved: (along with the Hamiltonian)

$$J_{1} = (p_{1} - p_{2})^{2} + 4e^{q_{1} - q_{2}},$$

$$I_{1} = \frac{p_{1} - p_{2} + \sqrt{J_{1}}}{p_{1} - p_{2} - \sqrt{J_{1}}} \exp\left(\sqrt{J_{1}} \frac{q_{1} + q_{2}}{p_{1} + p_{2}}\right)$$

We can numerically solve the equations of motion with

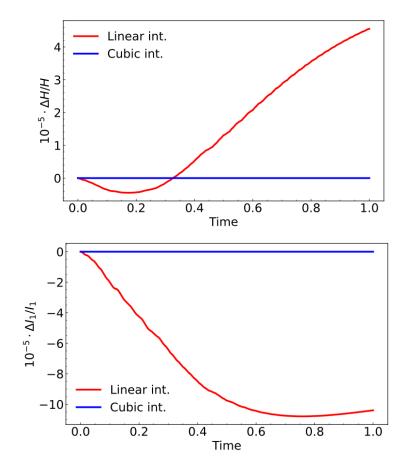
the method: $q_1^f = q_1^i + p_1 \cdot \Delta t$ $q_{1} = q_{1} + p_{1} \cdot \Delta t$ $q_{2}^{f} = q_{2}^{i} + p_{2} \cdot \Delta t$ $p_{1}^{f} = p_{1}^{i} - \begin{array}{c} \frac{\partial \phi}{\partial q_{1}}(q_{1}^{f}, q_{2}^{f}) \cdot \Delta t \\ \frac{\partial \phi}{\partial q_{2}}(q_{1}^{f}, q_{2}^{f}) \cdot \Delta t \end{array}$ • The potential is discreduled by the two intermetion methods are used. I. Linear interpolation 2. Cubic interpolation

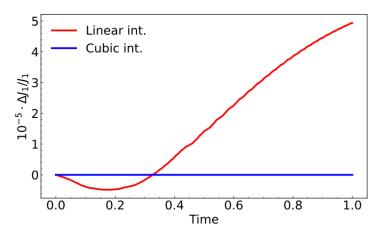
- The potential is discretized on a grid and the two interpolation

Demonstration of symplecticity violation/preservation

Non-symplectic method: Use (bi)linear interpolation on the derivatives of $\phi(q_1, q_2) = e^{q_1 - q_2}$.

Symplectic method: Use (bi)cubic interpolation on $\phi(q_1, q_2) = e^{q_1 - q_2}$.

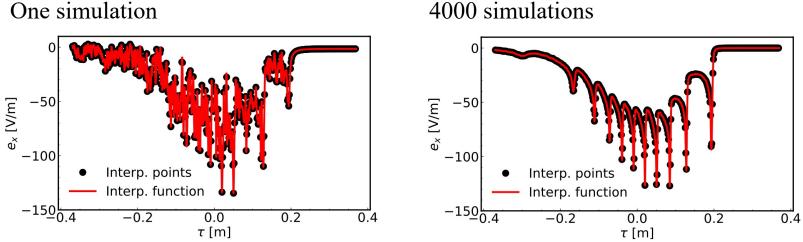




- The relative error on the integrals of motion does not grow with a symplectic method,
- While it grows for non-symplectic methods.

End of digression

Issue with PIC potential

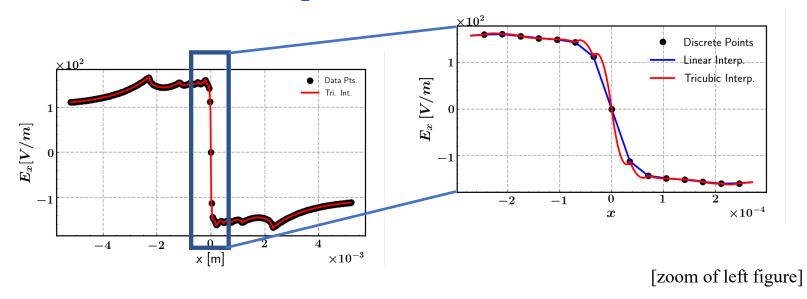


One simulation

- PIC simulation suffers from macroparticle noise. ٠
- Can be reduced by averaging many simulations. ٠

Averaging 4000 reveals the physical structures in the induced forces.

Issue with cubic interpolator



- Close look reveals irregularities from Tricubic interpolation.
- Inaccuracies are correlated with **discontinuity of second derivative accross cells**.

$$\mathcal{E}_{x}^{ijk} = \left. \frac{\partial e_{x}^{int}}{\partial x} \right|_{x \to x_{i}^{+}} - \left. \frac{\partial e_{x}^{int}}{\partial x} \right|_{x \to x_{i}^{-}} = -2 \frac{\partial^{3} \phi}{\partial x^{3}} (x_{i}, y_{j}, \tau_{k}) \Delta x + O(\Delta x^{3})$$

Refinement of potential

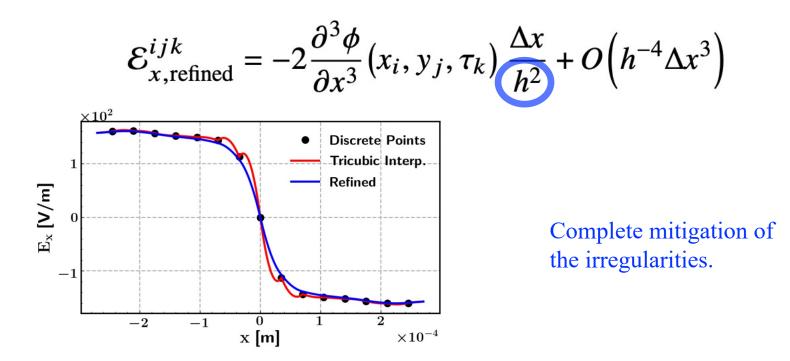
We found that we can treat our potential by:

- 1. Interpolate charge density on an auxilliary finer grid (by factor h).
- 2. Recalculate φ and derivatives in the finer grid.
- 3. Store recalculated φ and derivatives on original grid.

 $=\frac{\Delta x}{h}$

 $\Delta x_{\text{refined}}$

Minimal expense on memory and speed (performed during pre-processing) Proved analytically that error becomes:



Quick recap

- Analytical form of e-cloud kick.
- Used a high-order interpolation scheme (tri-cubic) to preserve symplecticity everywhere in phase space.
- Averaged multiple Particle-In-Cell e-cloud simulations to reduce macroparticle noise in the interpolated data.
- Solved Poisson's equation in a finer auxiliary grid (done only once) to improve performance of the interpolation scheme.

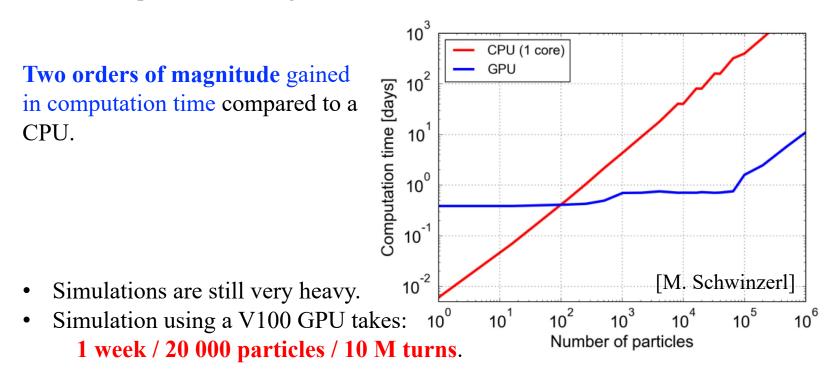
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Next:

- Direct tracking simulation results of the **incoherent effect of electron clouds in the main dipole and quadrupole magnets of the LHC at injection energy**.
- Interaction was implemented in SixTrackLib/XSuite software to use GPUs and including the **full lattice model** of the LHC.
- SixTrackLib/XSuite simulates beam particles through each element of the lattice using symplectic (non-linear) maps.

Graphics Processing Units

Realistic e-cloud simulation studies were made possible only through access to modern Graphics Processing Unit hardware.



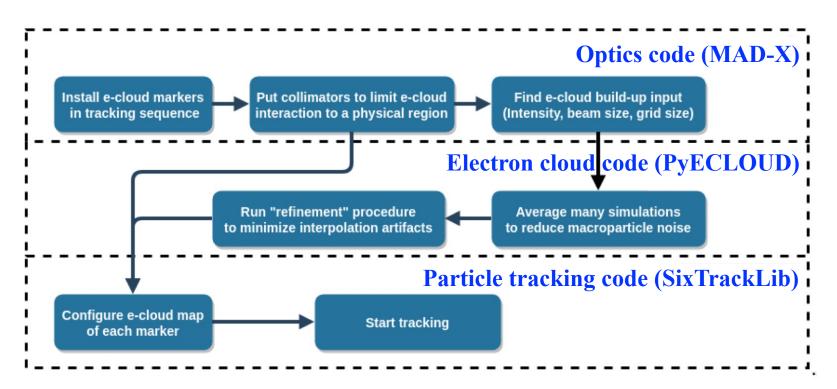
Following studies used several GPUs available from:

- CERN IT batch service
- INFN-CNAF in Bologna (through HL-LHC collaboration)

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General procedure for the simulation



Examples:

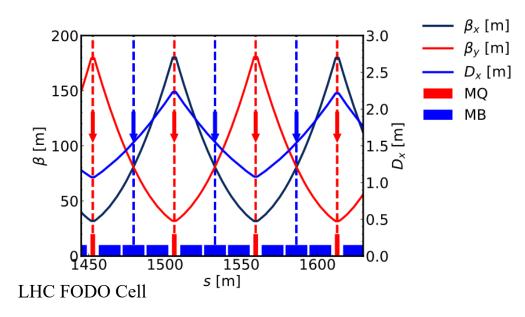
- Dynamic aperture tune scan to optimize accelerator's configuration.
- Dynamic aperture (SEY, bunch intensity): qualitative characterization of e-cloud effect, tolerances/intensity reach.
- Frequency Map Analysis for insight on the mechanism.
- Long-term simulations for estimations of losses and emittance growth.

E-cloud setup

E-cloud exists across the full length of the LHC beam pipe. **Different magnetic fields lead to completely different e-clouds.** Most significant contributors:

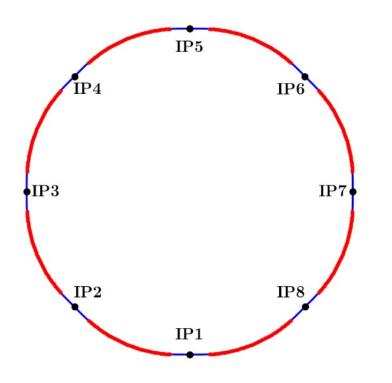
- 1. E-cloud in arc dipoles (MB) (66%)
- 2. E-cloud in arc quadrupoles (MQ) (7%)

We place one interaction for each three dipoles and each quadrupole.



- Betatron and dispersion functions stay the same between each cell.
- Approximate SEY as uniform everywhere. Large fluctuations in reality.
- Effect from saturated e-cloud.

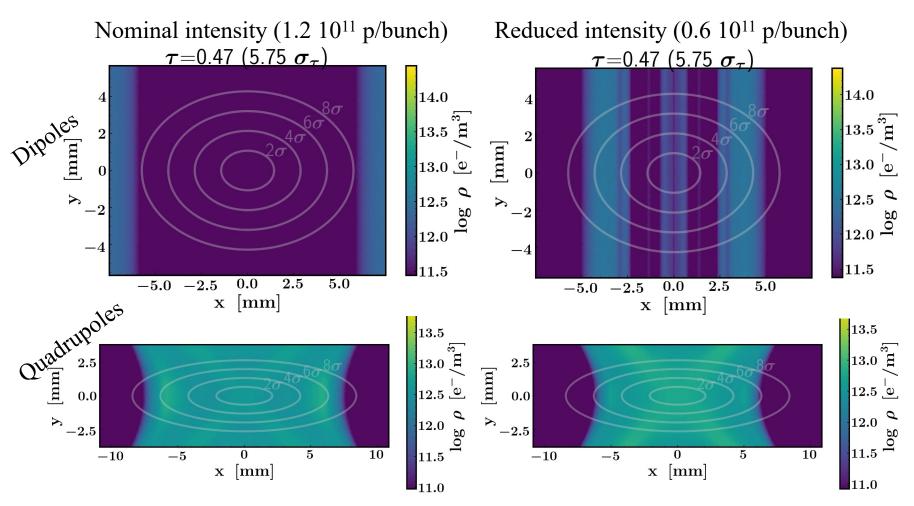
E-cloud setup



- One MB e-cloud per half-cell
- \rightarrow 46 interactions per arc
- \rightarrow 368 interactions.
- One MQ e-cloud per half-cell
 → 45 interactions per arc
 → 360 interactions.

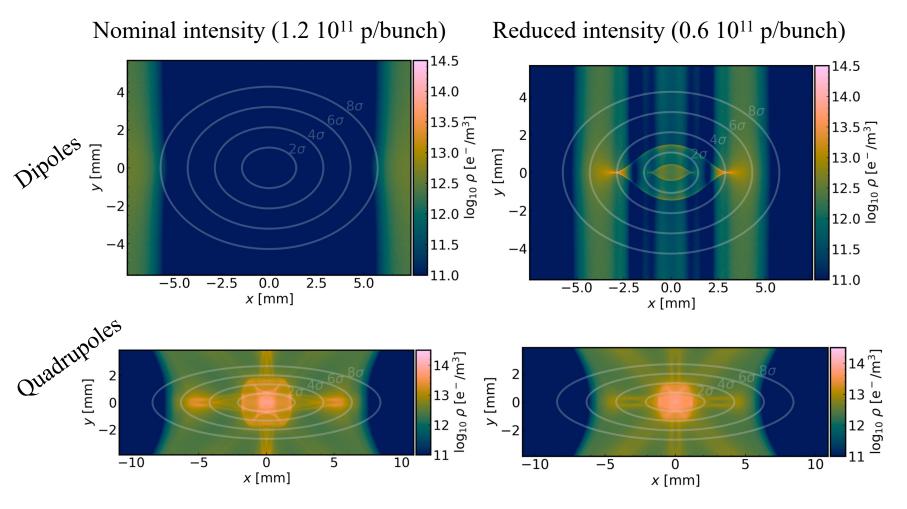
Tracking time per e-cloud type (~360 interactions) is about as much as rest of the lattice (11k tracking elements).

E-cloud setup



- Dipoles: Reduced bunch intensity leads to larger e⁻ density close to the beam.
- Quadrupoles: Small dependence on bunch intensity, large e⁻ densities close to beam.

E-cloud setup



- Dipoles: Reduced bunch intensity leads to larger e⁻ density close to the beam.
- Quadrupoles: Small dependence on bunch intensity, large e⁻ densities close to beam.

Simulation Parameters

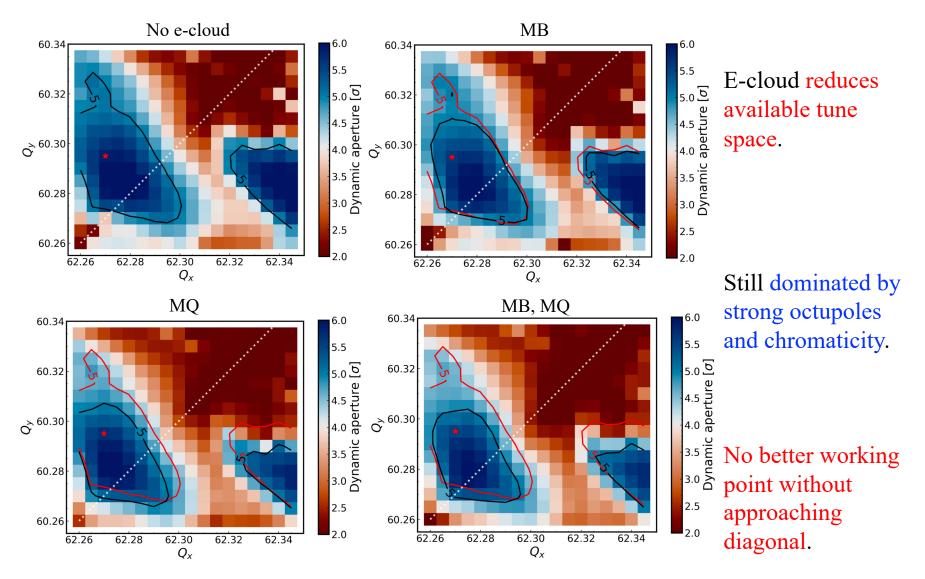
Typical LHC at injection, 2018

Bunch intensity : $1.20\ 10^{11}$ protons Energy : $450\ GeV$ Chromaticity : 15/15Octupole magnet's current : $40\ A$ Bunch spacing : $25\ ns$ Transverse norm emittances : $2\ \mu m/2\ \mu m$ R.M.S. bunch length : $0.09\ m$ Betatron tunes : 62.270/60.295RF voltage : $6\ MV$

The three primary collimators (TCP) in IR7 (as black absorbers) are included in the lattice at their typical configuration (5.7 "collimation" $\sigma \rightarrow 7.5$ beam σ).

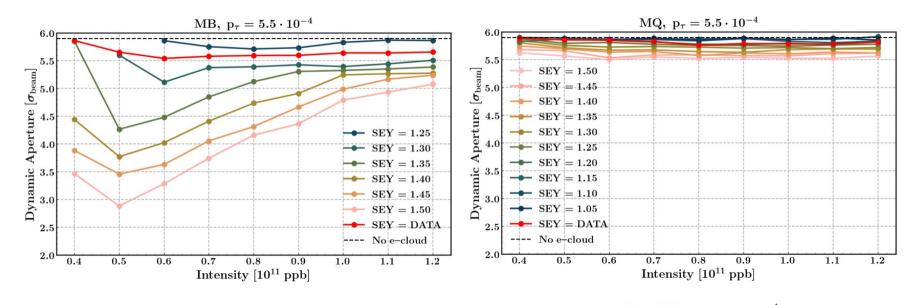
There is **no** uncorrected linear coupling, magnet field imperfections, magnet misalignments or beam-beam interactions in the lattice.

Tune scan

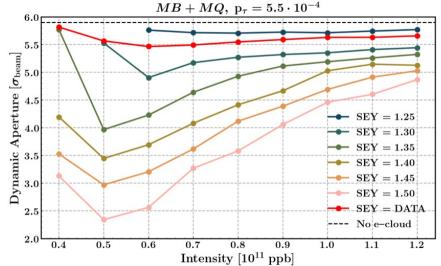


 $⁽SEY_{MB}: 1.3, SEY_{MQ}: 1.3)$

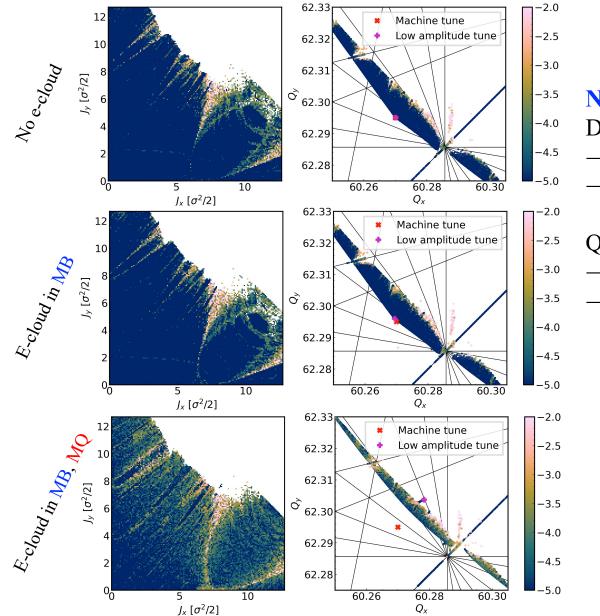
Secondary Emission Yield (SEY) - Intensity scan



- Larger Secondary electron Emission Yield (of beam pipe) →
 → stronger e-cloud → less DA
- Dipoles (MB): strong dependence with bunch intensity, correlated to edensity close to the beam.
- Quadrupoles (MQ): weak dependence with bunch intensity



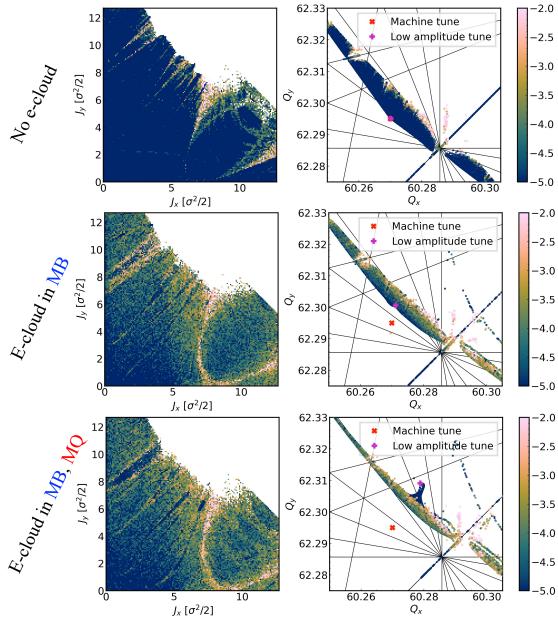
Frequency Map Analysis – Nominal intensity (1.2 10¹¹ p/b)



Nominal intensity

- Dipoles (MB):
 - \rightarrow tiny tune-shift
 - \rightarrow negligible effect
- ⁵ Quadrupoles (MQ):
- \rightarrow large tune-shift
- \rightarrow more resonances

Frequency Map Analysis – Reduced intensity (0.6 10¹¹ p/b)



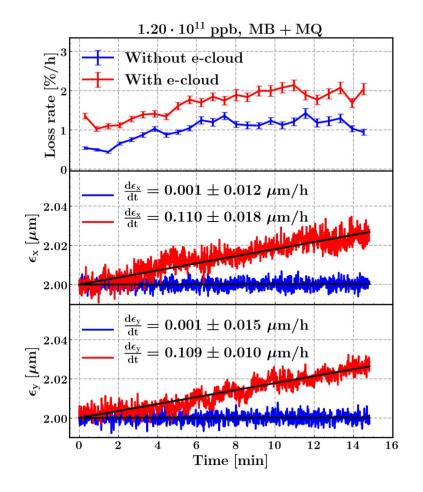
- **Reduced** intensity
- Dipoles (MB):
 - \rightarrow larger tune-shift
 - \rightarrow more resonances
- ⁵ Quadrupoles (MQ):
- \rightarrow large tune-shift
- \rightarrow more resonances

Reminder:

Significant e⁻ density appears on the beam location in dipole magnets for reduced intensity.

Particles are on-momentum. Work in progress to try identify synchro-betatron resonances.

Long simulations (10M turns \rightarrow 15min beam time)



Incoherent effects in the LHC are typically very slow processes. Need to simulate long timescales. Recent advances (SixTrackLib/XSuite) allow the direct simulation of particle distributions with GPUs for such times.

In long term simulations we observe:

- small increase of losses
- horizontal emittance growth,
- vertical emittance growth, when e-clouds are included.

Experimental observations show emittance growth in the same order of magnitude. For quantitative comparisons we have planned dedicated MDs in Run 3.

Conclusion and Remarks

Observations:

• Electron cloud in the insertion region quadrupoles is significant. Reduces integrated luminosity.

Simulations:

- We can do particle tracking simulations with **arbitrarily complex e-clouds in arbitrarily complex lattices for millions of turns.**
- Simulated simplified scenario at injection energy. Interplay with non-linear magnetic imperfections expected.
- Simulations have reproduced the expected qualitative behavior.
- Very long simulation timescales (several minutes) are in reach. (Using GPUs)

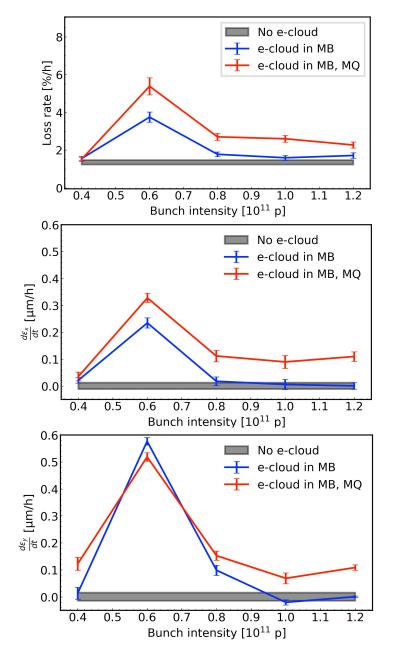
Outlook for the future:

- Comparison with experimental measurements needs specialized tests.
 → Soon to be carried out in the LHC.
- Simulate more complex scenario of collisions in LHC: Strong electron clouds in the Insertion Region quadrupoles + beam-beam effects.

Thank you for your attention! Konstantinos Paraschou

Backup slides

Long simulations (10M turns \rightarrow 15min beam time)



MB (Dipoles):

- Losses stronger at reduced intensity.
- Emittance growth only at reduced intensity.
- Vertical growth larger than horizontal.

MQ (Quadrupoles):

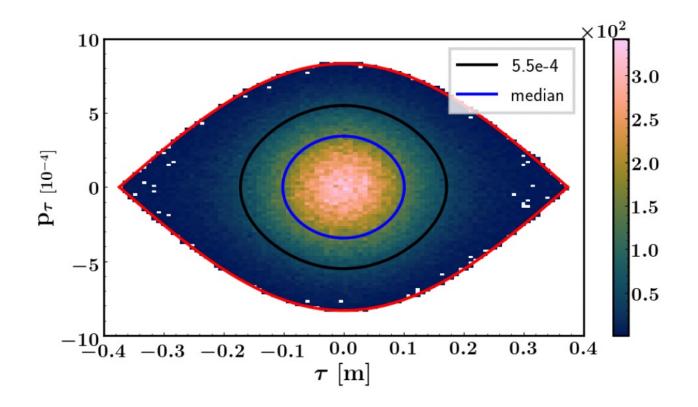
- Losses across all intensities.
- Emittance growth at all intensities.
- Similar growths in both horizontal and vertical.

Effects strongly correlated with the edensity close to the beam.

Reminder:

- MB show large densities around the beam for reduced intensities,
- MQ for all intensities.

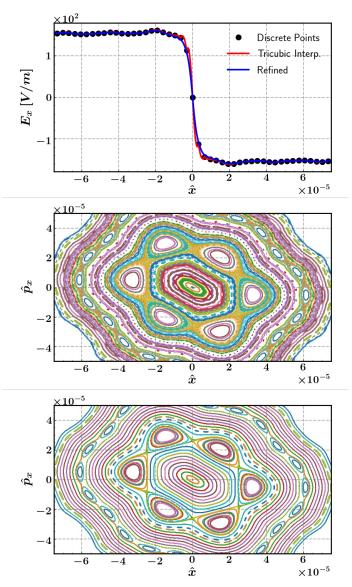
The RF bucket



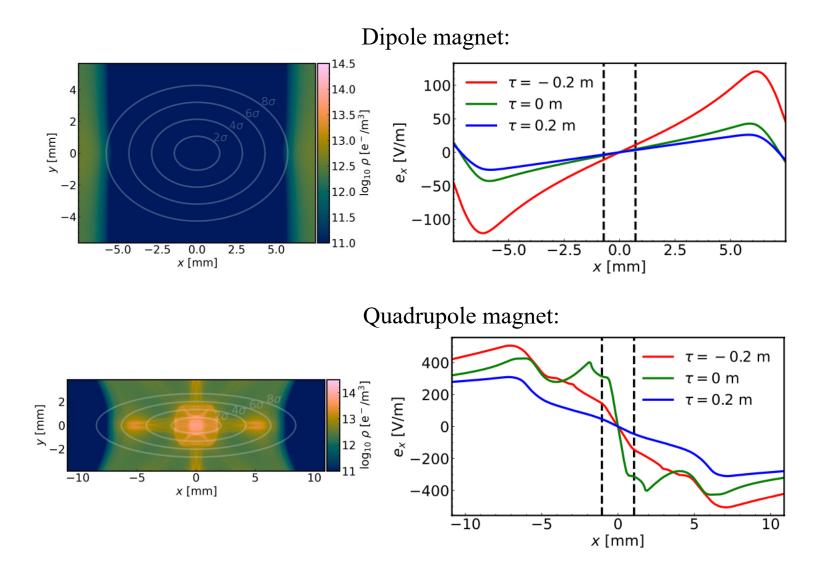
- DA simulations done for off-momentum particles ($p_{\tau} = 5.5 \ 10^{-4}$).
- FMA simulations done for on-momentum particles ($p_{\tau} = 0$).
- Long-term tracking simulations with particles across the full bucket.
- Work in progress: FMA with off-momentum particles.

Impact of tricubic interpolation irregularities

- Simple tracking of linear 2D phase space rotation and an e-cloud symplectic kick.
- Very important to minimize irregularities.
- By reducing them, there is significant impact on the particle motion.



Induced forces



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