

QCD equation of state at finite density with a critical point from an alternative expansion scheme

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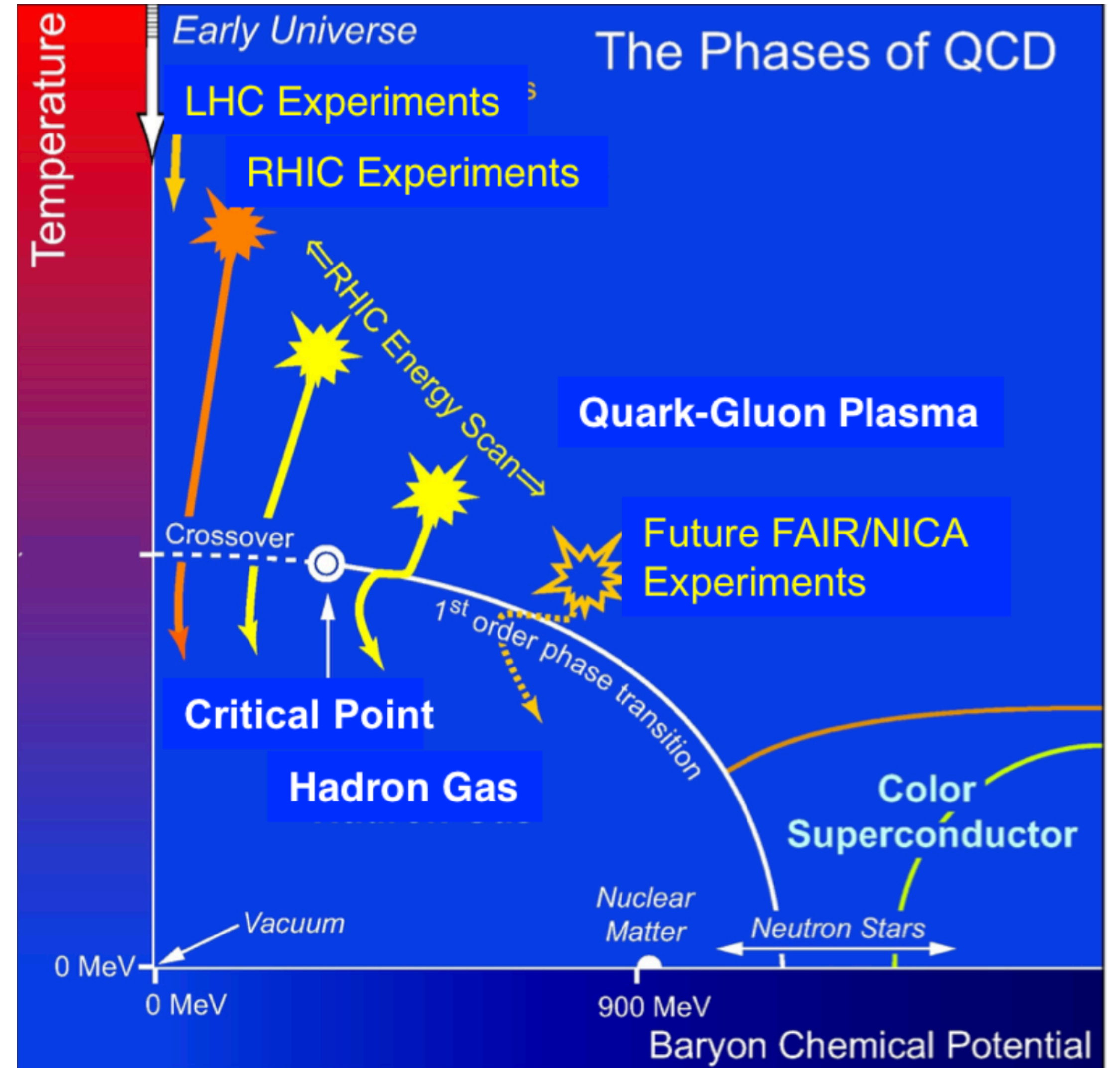
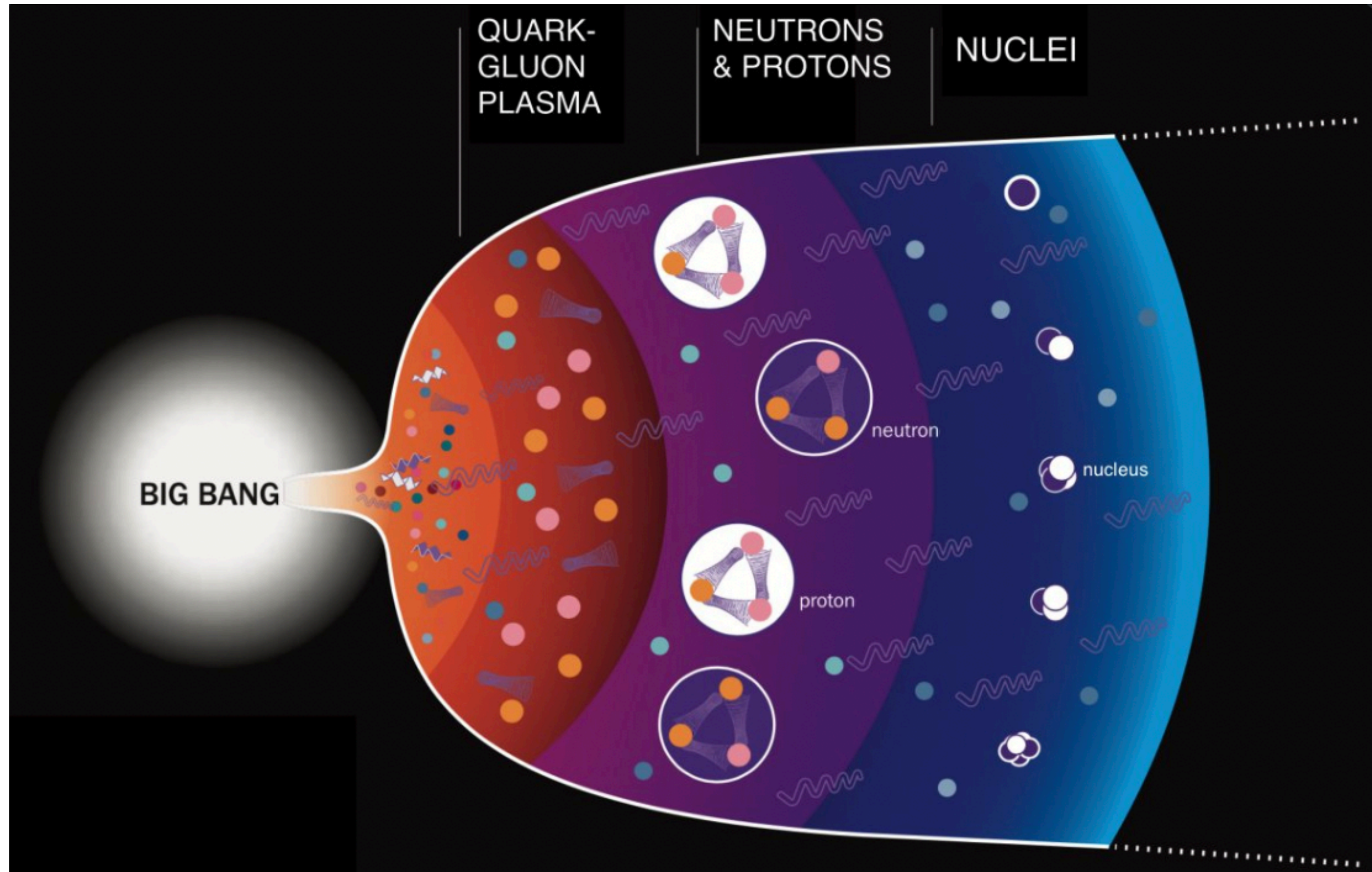
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Introduction

QCD Phase Diagram

Different phases of strongly interacting matter in equilibrium are captured in the QCD phase diagram

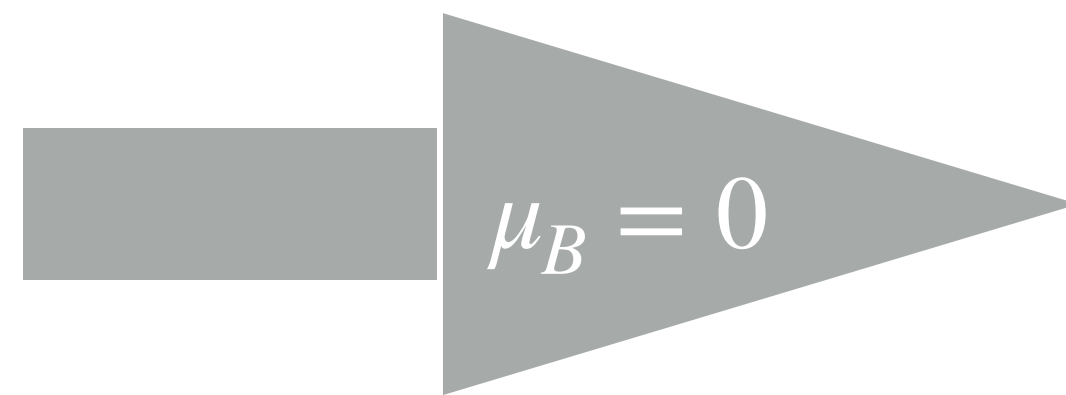


Lattice QCD results

Fermi Sign Problem!

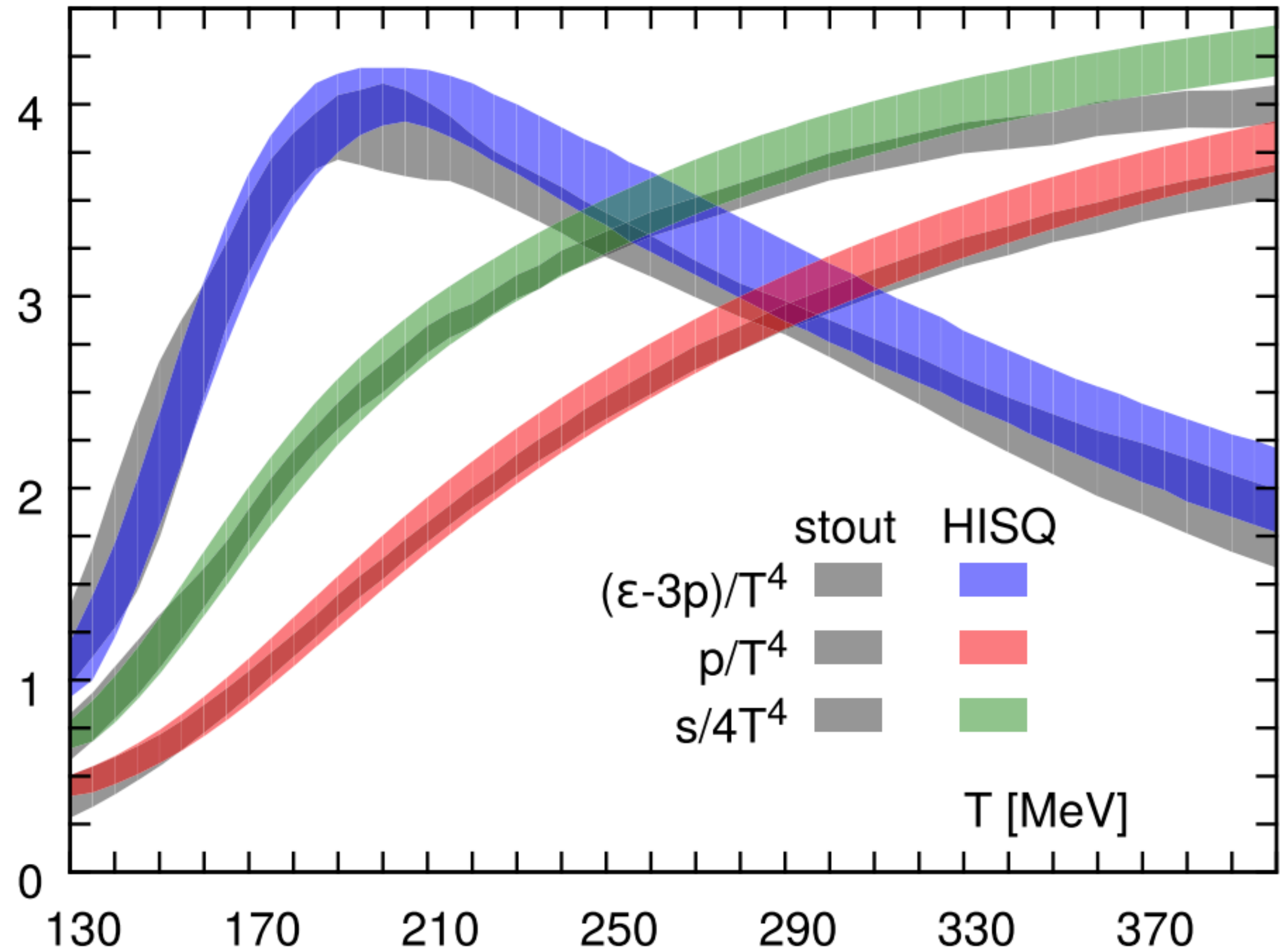
[Wuppertal-Budapest (stout) and HotQCD (HISQ) collaborations]

- Simulation at finite μ_B is Impossible



Way out!

- Simulating Imaginary $\mu_B \rightarrow$ Analytical Continuation
- Expansion around $\mu_B = 0$



[S. Borsanyi et al Phys. Lett. B 730 (2014) 99]

[A. Bazavov et al., PhysRevD.90.094503(2014)]

Part 1: Lattice EoS: Taylor Expansion

Lattice EoS at Finite Density

Taylor Expansion around $\mu_B = 0$

Limitations

Most Straightforward

$$\frac{P(T, \mu_B)}{T^4} = \sum_{n=0}^{\infty} \frac{1}{n!} \chi_n(T, \mu_B = 0) \left(\frac{\mu_B}{T} \right)^n$$

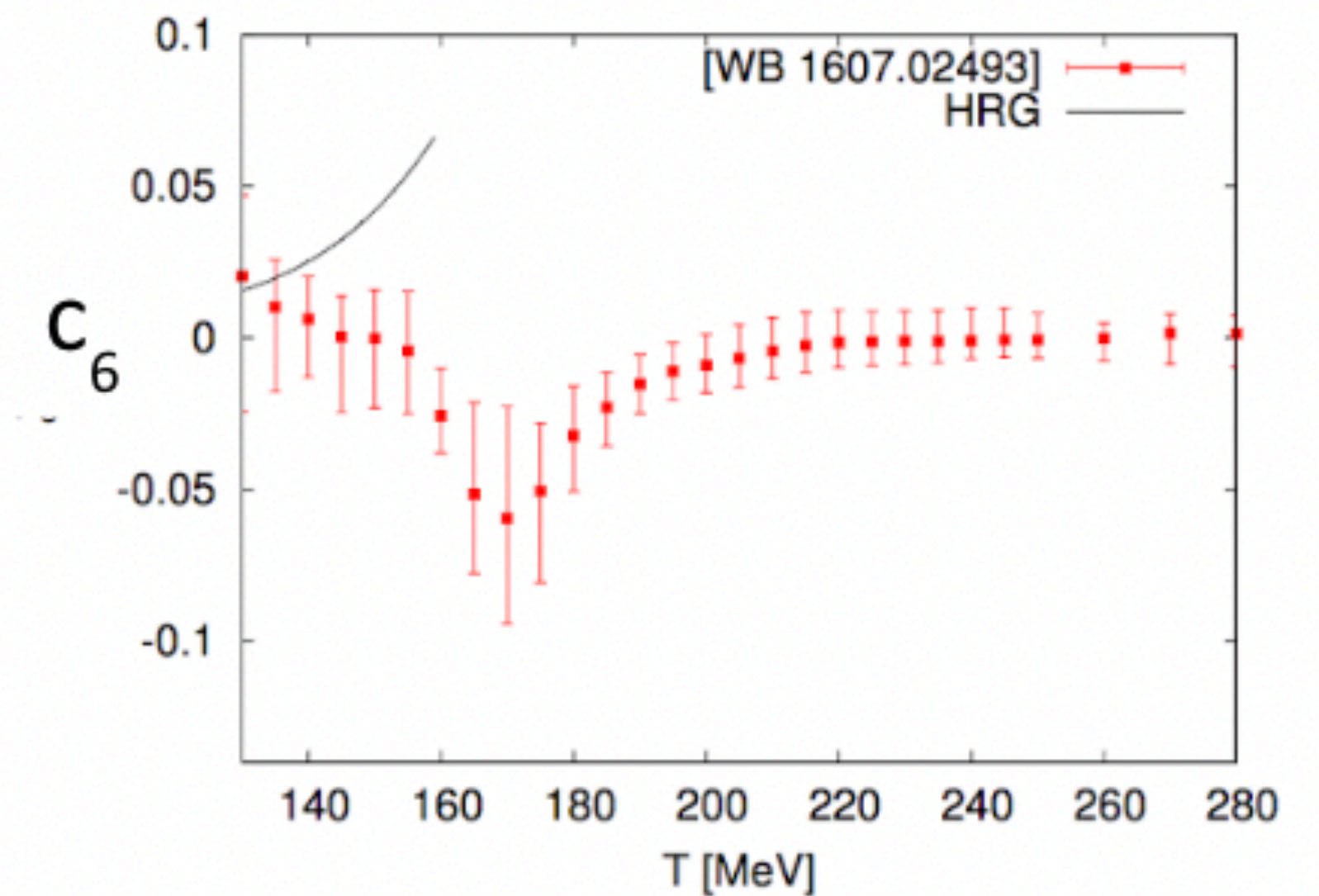
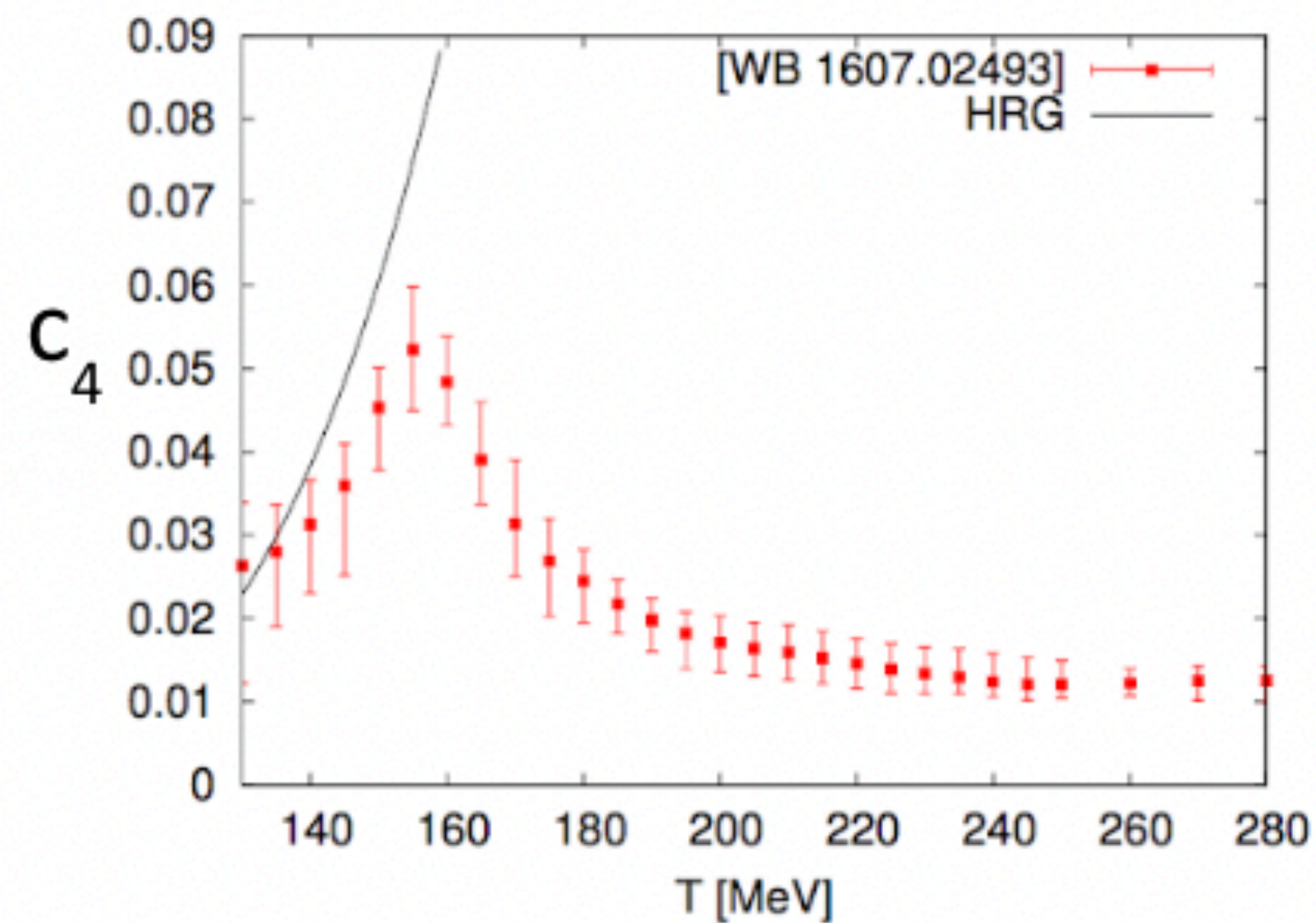
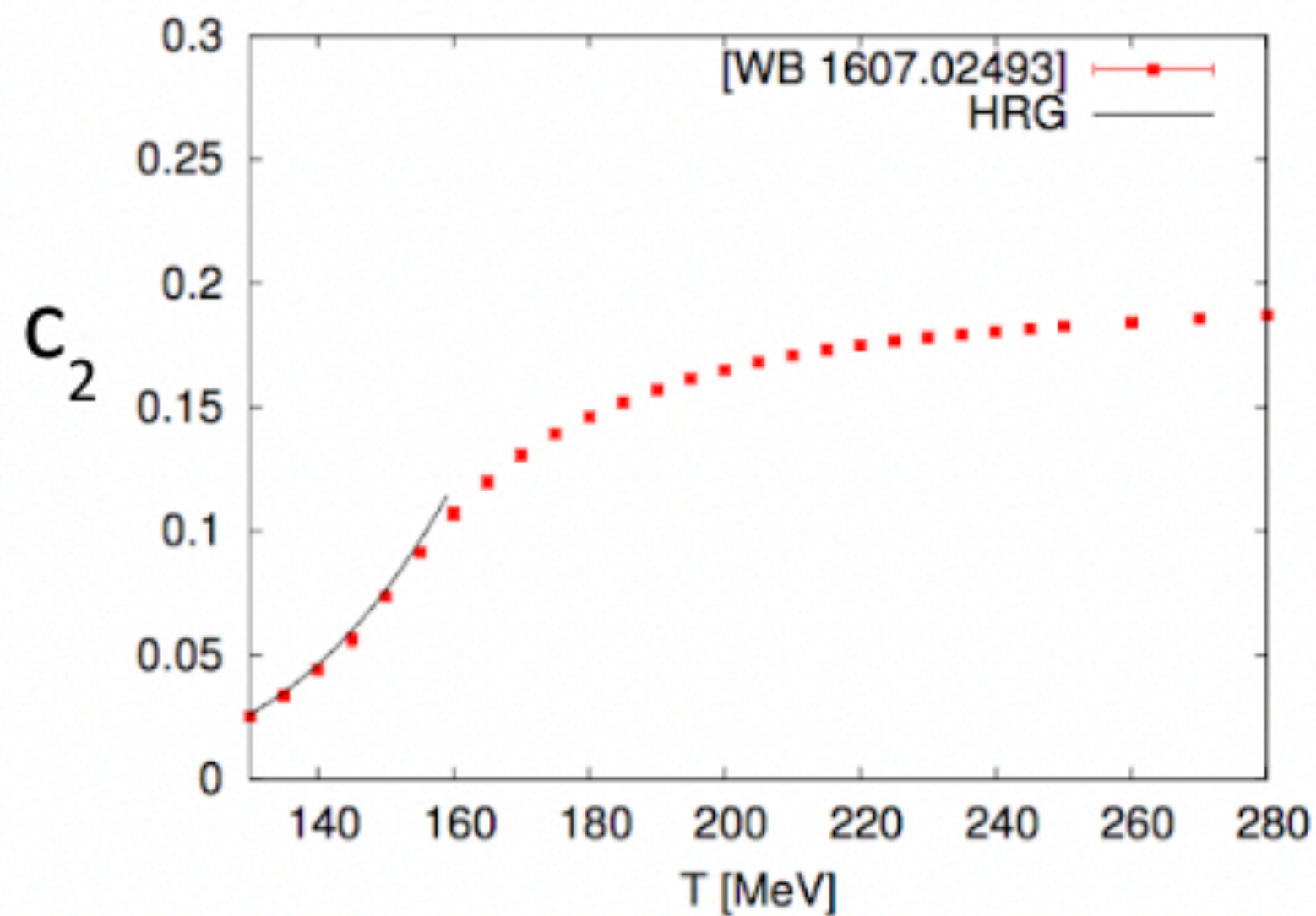
- Currently limited to $\frac{\mu_B}{T} \leq 2.5$ despite great Computational Power
- Adding one more Higher-Order term does not help in convergence
- Taylor expansion is carried out at $T = \text{constant}$ and doesn't cope well with μ_B -dependent transition temperature

[Borsányi, S. et al. PRL (2021)]

$$c_n(T) = \frac{\chi_n^B(T, \mu_B = 0)}{n!} = \frac{1}{n!} \left(\frac{\partial}{\partial(\mu_B/T)} \right)^n (P/T^4) \Big|_{\mu_B=0}$$

- Re-summation with Pade' Approximation can reach $\frac{\mu_B}{T} = 3.0$

.D. Bollweg et al., arXiv:2212.09043v1 [hep-lat] (2022)



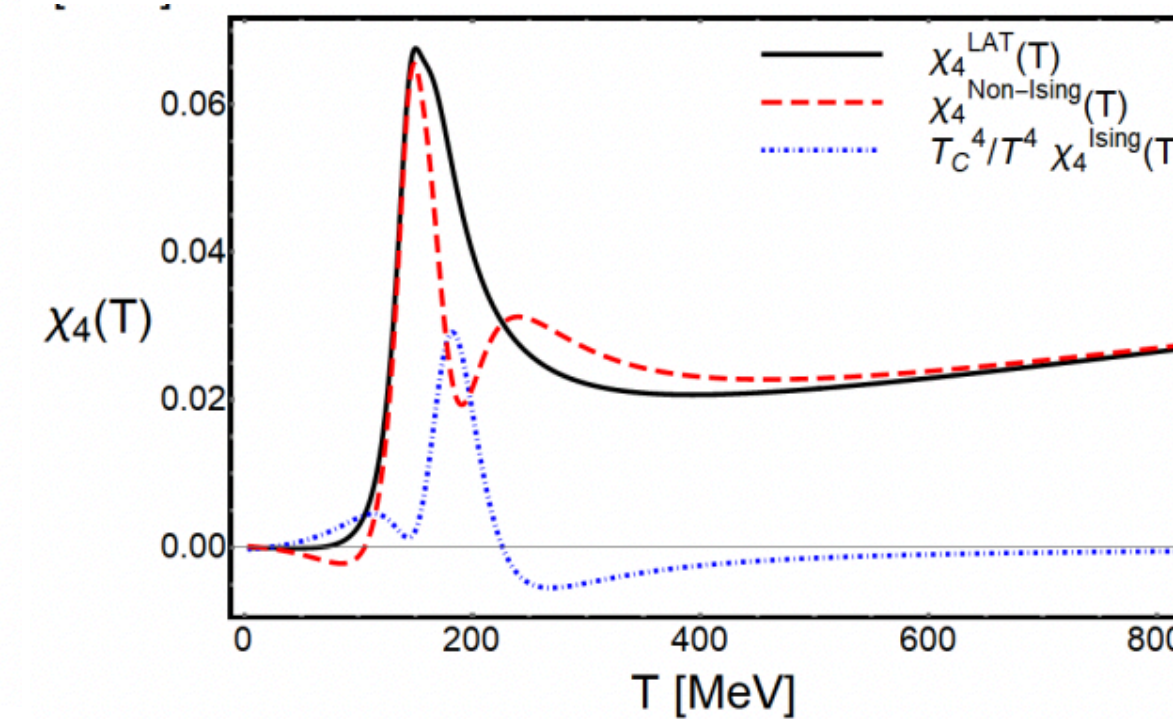
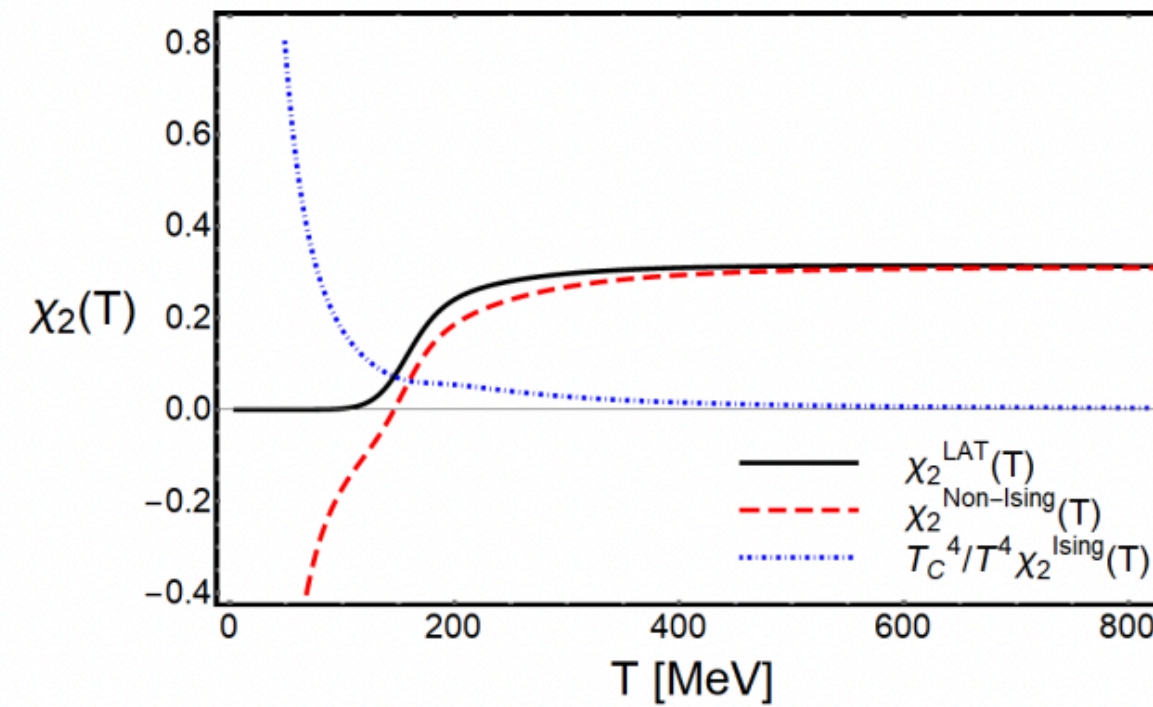
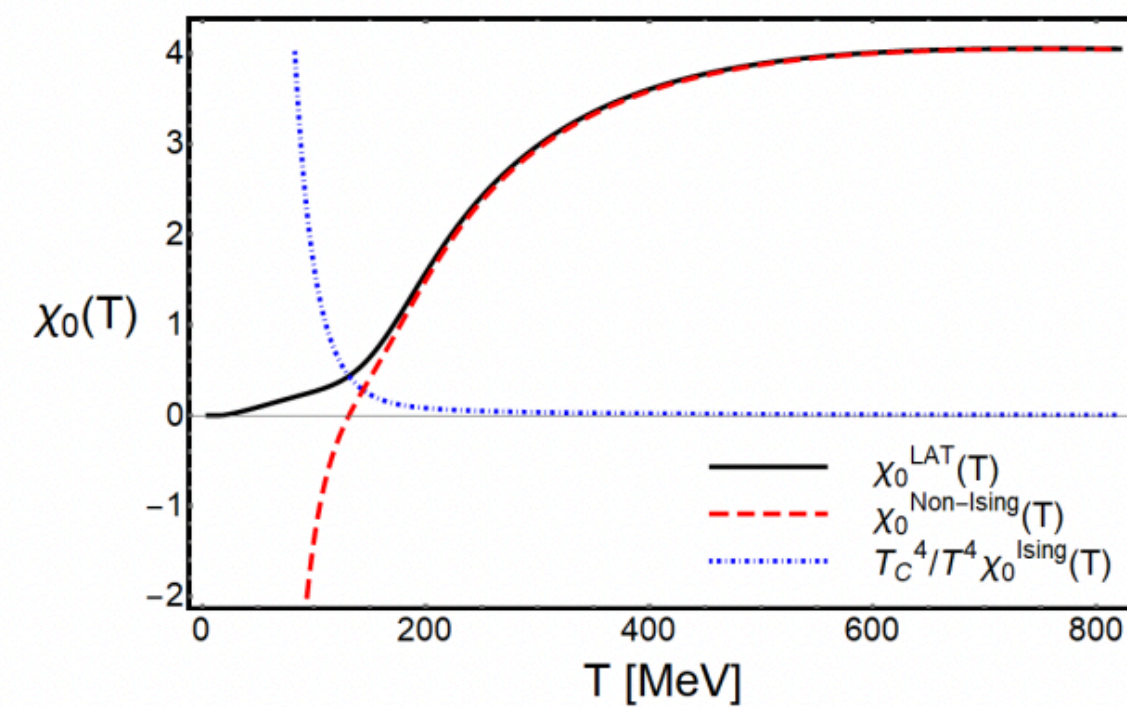
Wuppertal-Budapest

[Borsányi, S. et al. Journal of high energy physics (2018)(10), 1-29]

Taylor EoS with Critical point

$$P(T, \mu_B) = T^4 \sum_{n=0}^2 \frac{1}{(2n)!} \chi_{2n}^{Non-Ising}(T) \left(\frac{\mu_B}{T} \right)^{2n} + T_C^4 P_{symm}^{Ising}(T, \mu_B)$$

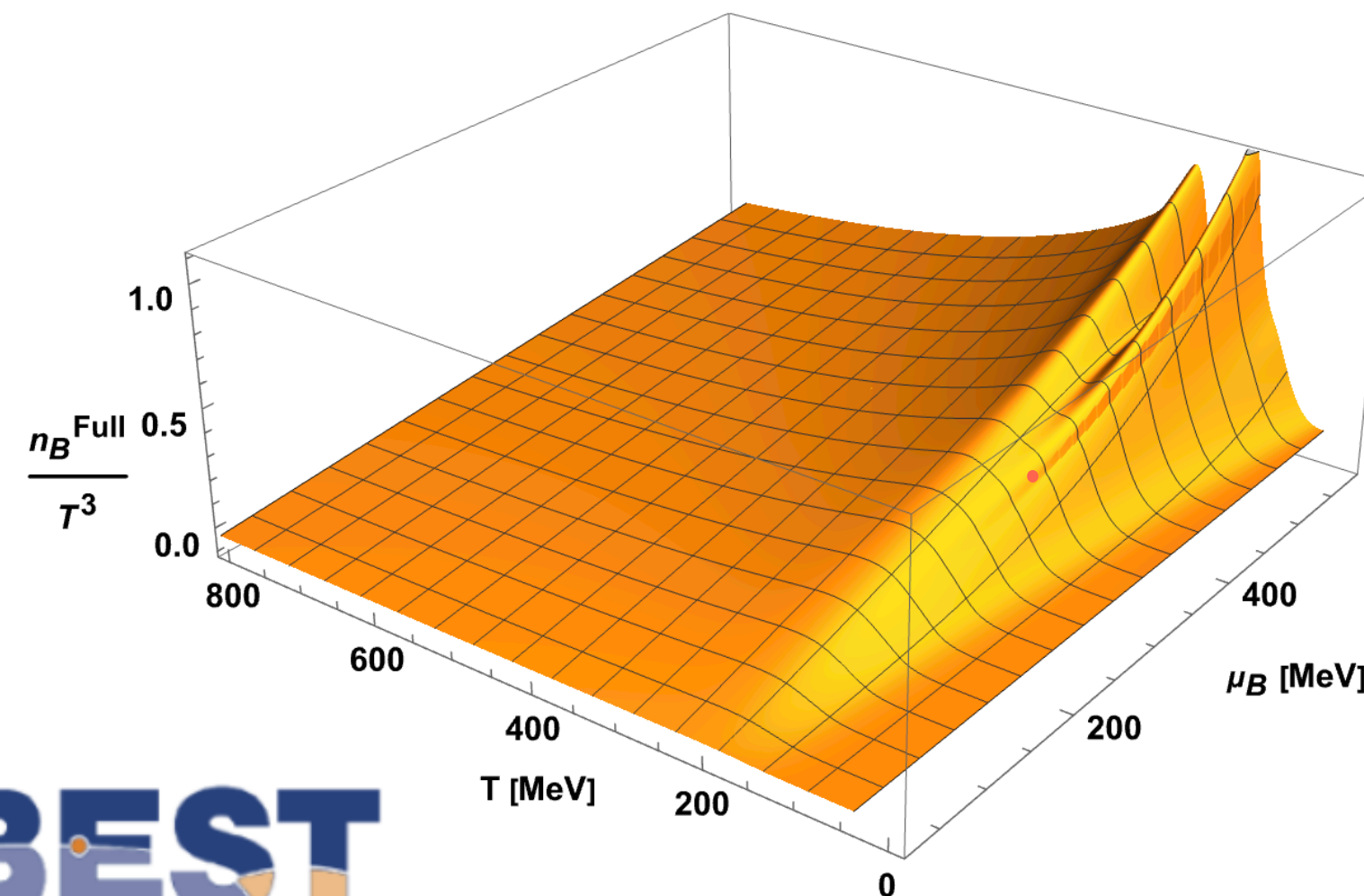
$$\chi_n^{Lat}(T) = \chi_n^{Non-Ising}(T) + \frac{T_C^4}{T^4} \chi_n^{Ising}(T)$$



3D Diagram

Baryon density

$$\frac{n_B(T, \mu_B)}{T^3} = \frac{\partial(P/T^4)}{\partial(\mu_B/T)}$$



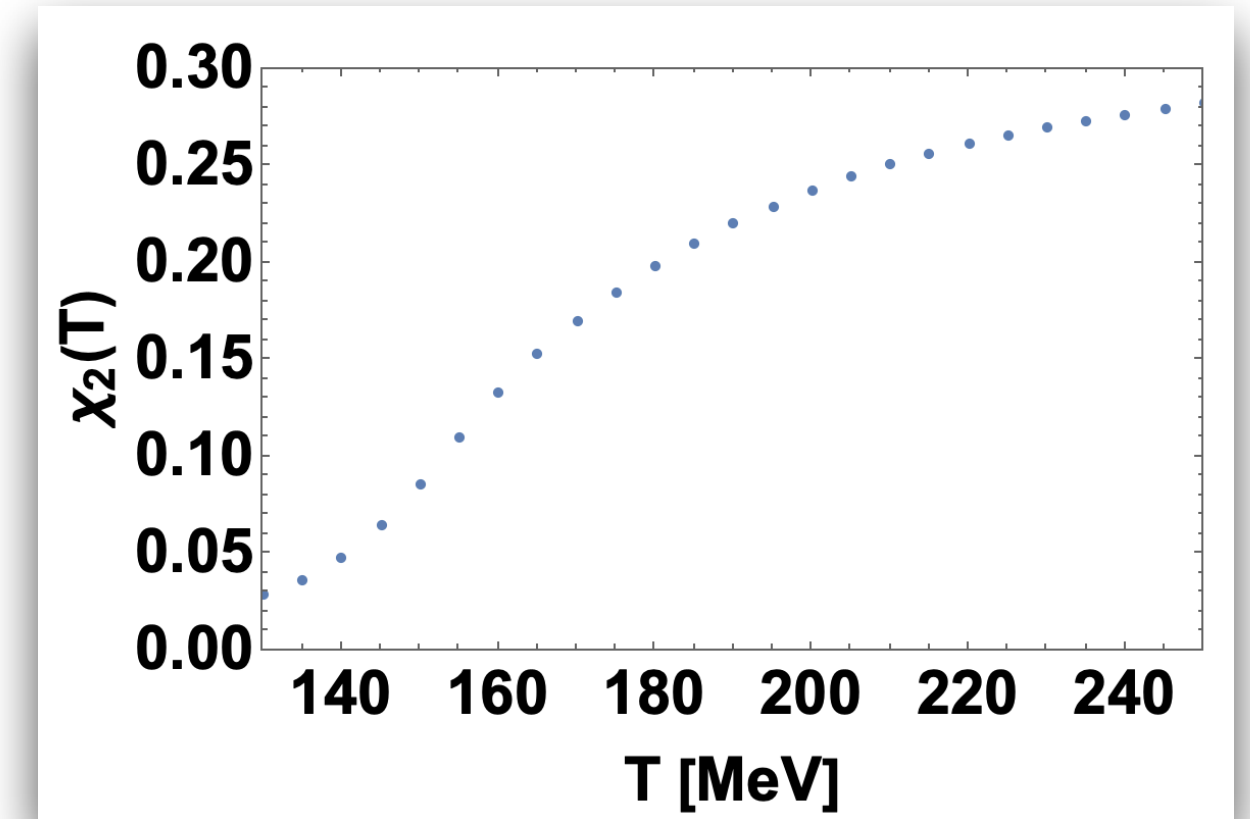
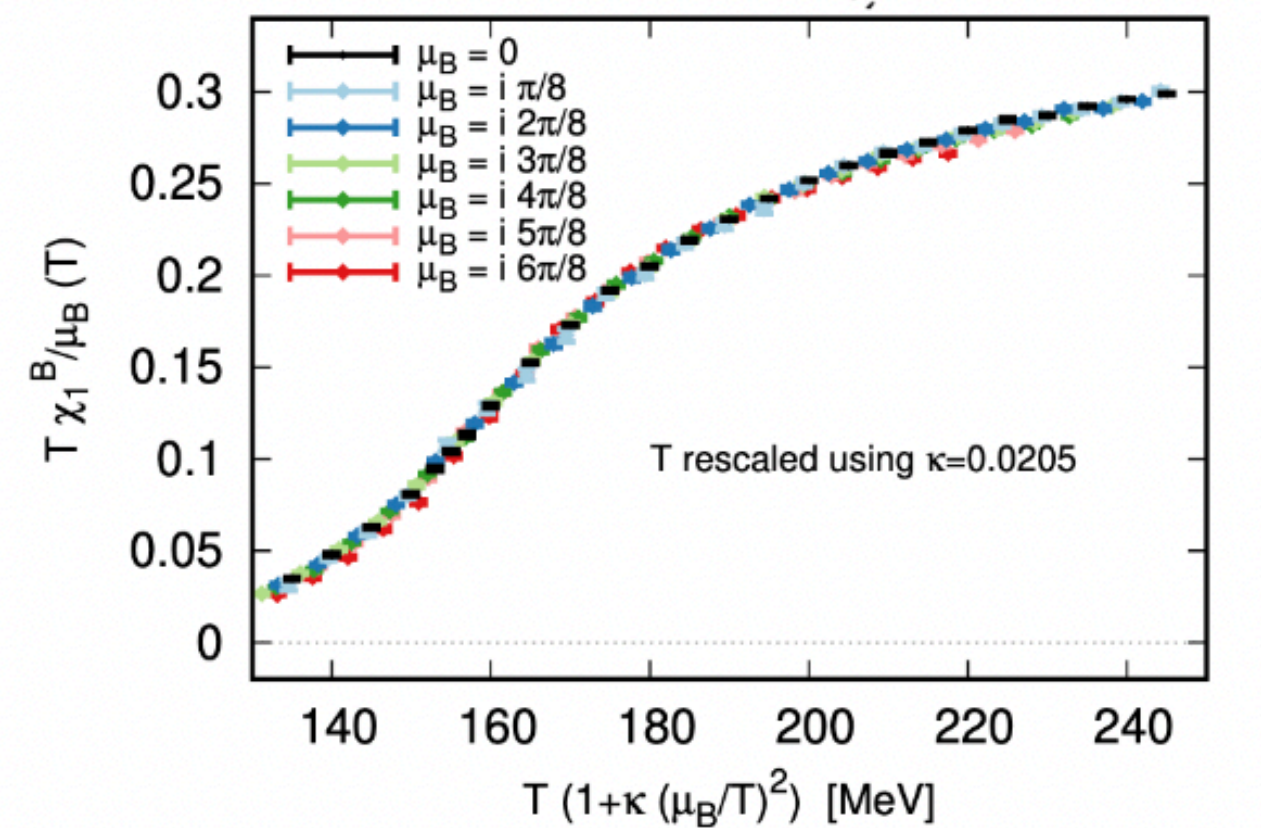
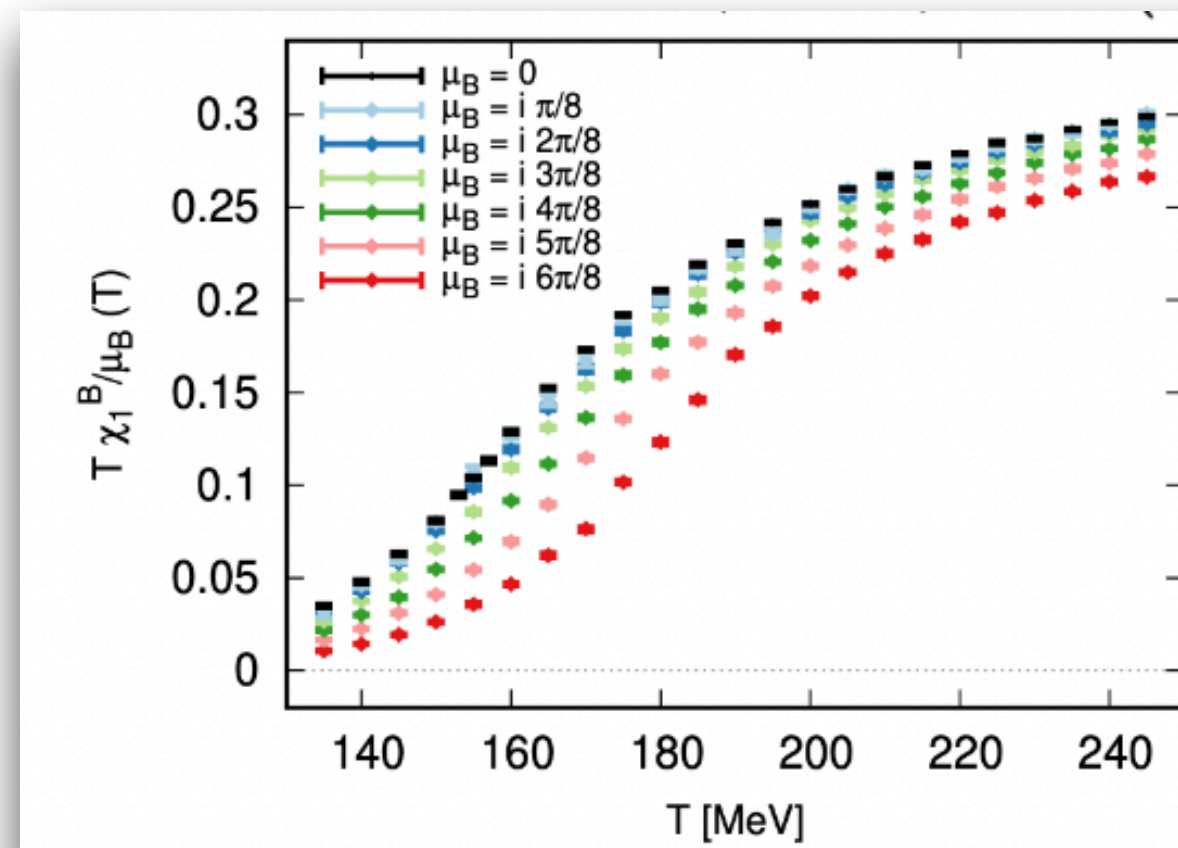
Critical Point at $\mu_{BC} = 350$ [MeV]

● Thermodynamic Observables at $\mu_B \geq 450$ MeV show unphysical behavior

Part 2: Lattice EoS: Alternative Expansion Scheme

Alternative Expansion Scheme: EoS

Simulating at Imaginary μ_B



[Borsányi, S. et al. PRL (2021)]

$$T \frac{\chi_1^B(T, \mu_B)}{\mu_B} = \chi_2^B(T', 0)$$

$$\chi_n^B(T, \mu_B) = \frac{1}{n!} \left(\frac{\partial}{\partial (\mu_B/T)} \right)^n (P/T^4)$$

$$T[T, \mu_B] = T \left[1 + \kappa_2^{BB}(T) \left(\frac{\mu_B}{T} \right)^2 + \kappa_4^{BB}(T) \left(\frac{\mu_B}{T} \right)^4 + \mathcal{O} \left(\frac{\mu_B}{T} \right)^6 \right]$$

- μ_B dependence is captured in T-rescaling.

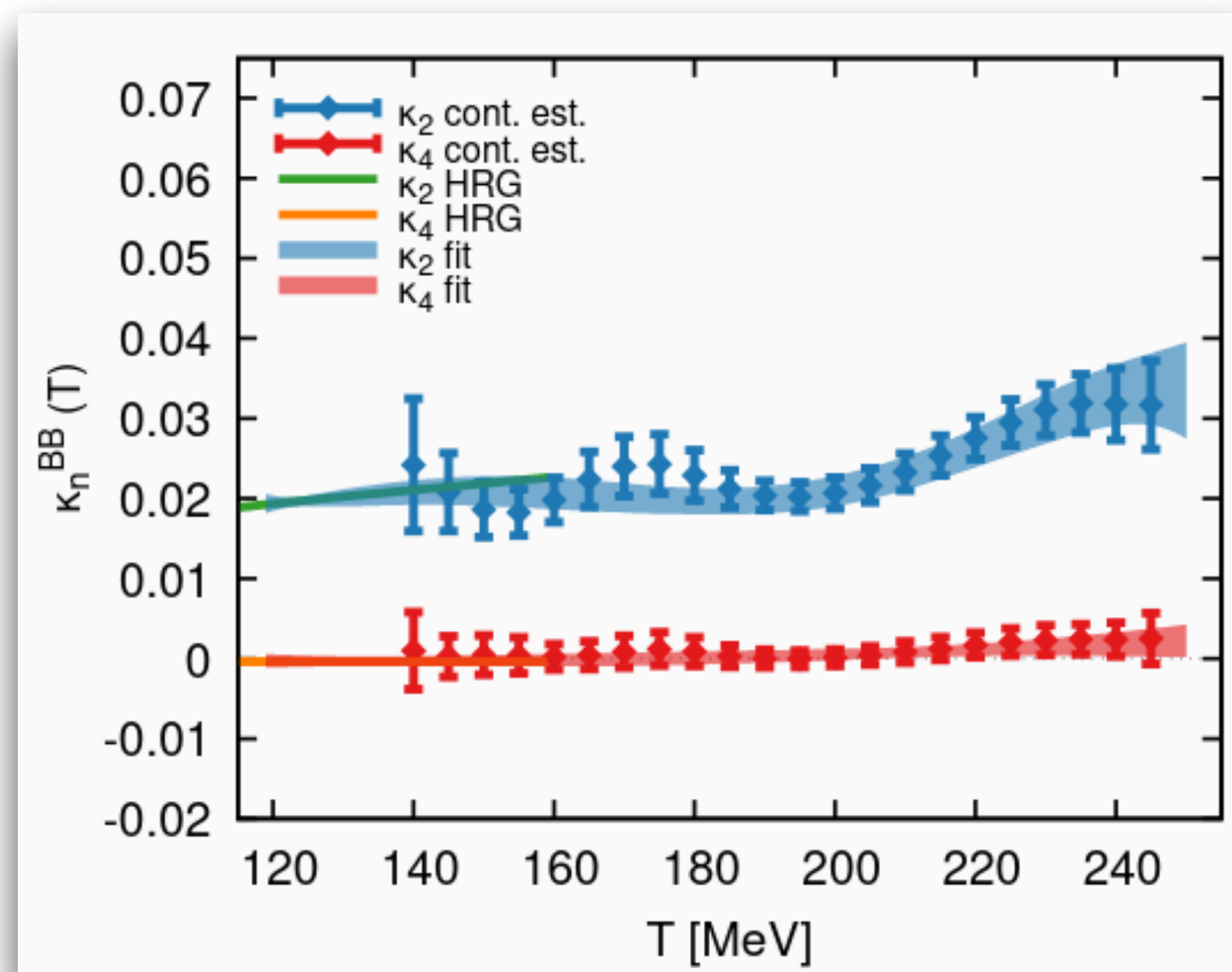
- Trusted up to $\frac{\mu_B}{T} = 3.5$

[Borsányi, S. et al. PRL (2021)]

Alternative expansion scheme

Comparing Taylor expansion and Alternative expansion

- $\kappa_2^{BB}(T) = \frac{1}{6T} \frac{\chi_4^B(T)}{\left(\frac{\partial \chi_2^B(T)}{\partial T}\right)}$
- $\kappa_4^{BB}(T) = \frac{1}{360\chi_2^B(T)^3} \left(3\chi_2'^{B2} \chi_6^B(T) - 5\chi_2^B(T)'' \chi_4^B(T)^2\right)$



[Borsányi, S. et al. PRL (2021)]

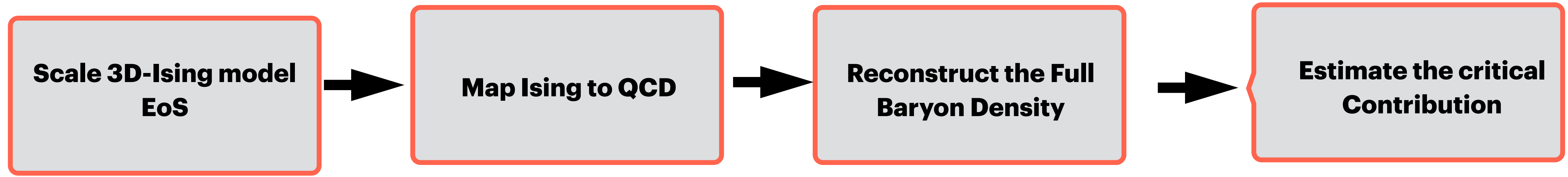
Pros

- $\kappa_2(T)$ is fairly constant over a large T-Range
- There is a separation of scale between $\kappa_2(T)$ and $\kappa_4(T)$
- $\kappa_4(T)$ is almost zero → faster convergence
- A good agreement with HRG results at Low Temperature

Part 3: Putting Critical Point into Alternative Expansion: EoS

EoS with a Critical point

Strategy



Scale 3D-Ising Model EoS

Close to the critical point, we define a parametrization for Magnetization M , Magnetic field h , and reduced temperature

QCD Critical point is in the 3D-Ising model Universality class

$$M = M_0 R^\beta \theta$$

$$h = h_0 R^{\beta\delta} \tilde{h}(\theta)$$

$$r = R(1 - \theta^2)$$

$$(R, \theta) \longmapsto (r, h)$$

$$r = \frac{T - T_C}{T_C}$$

- M_0, h_0 are normalization constants
- $\beta \approx 0.326$ and $\delta = 4.8$ are the 3D Ising model critical exponents
- $\tilde{h}(\theta) = (\theta + a\theta^3 + b\theta^5)$ with $a = -0.76201, b = 0.00804$
- The parameters take on the values $R \geq 0, |\theta| \leq \theta_0 \approx 1.154, \theta_0$ being the first non-trivial zero of $\tilde{h}(\theta)$

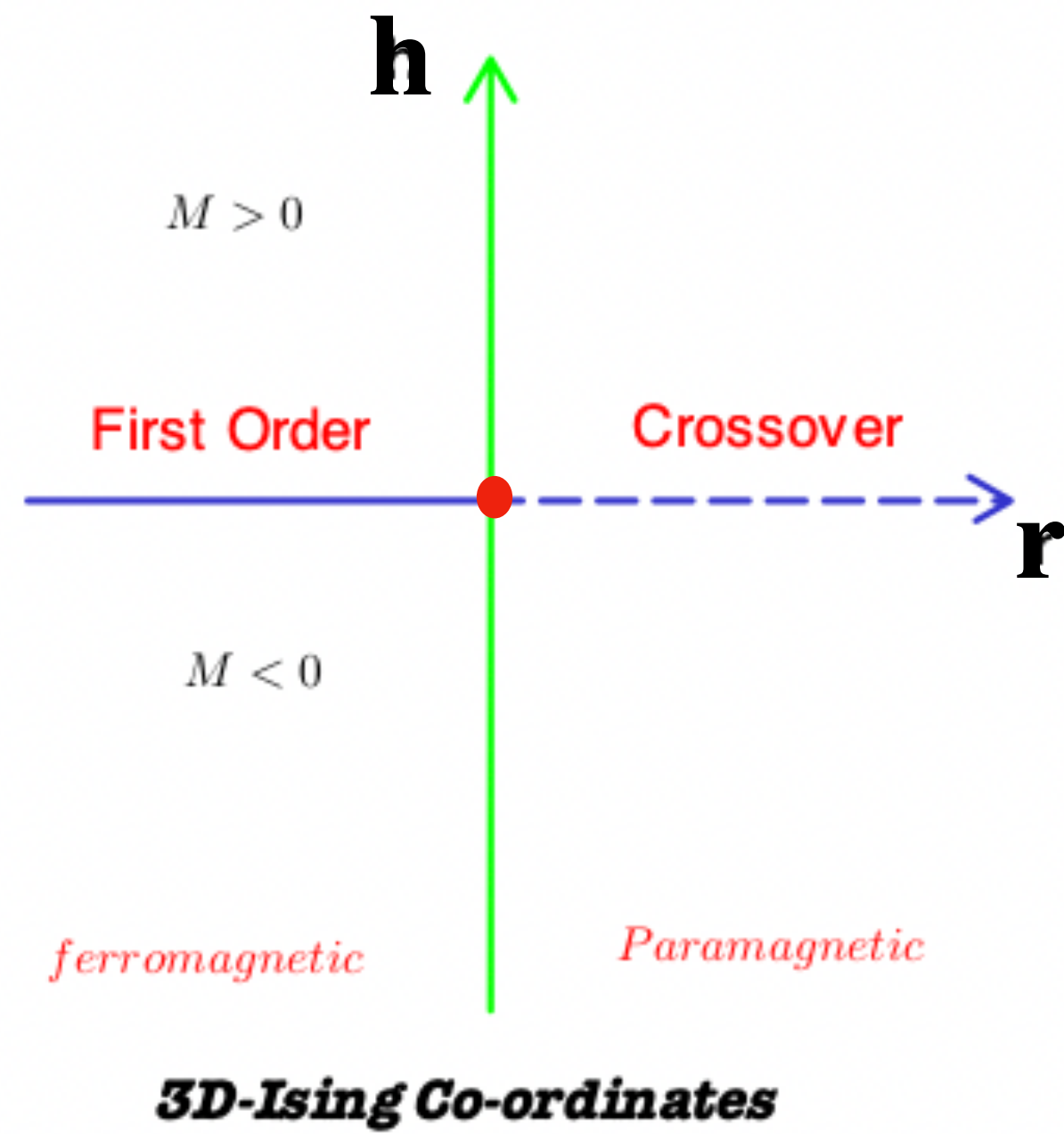
[Parotto et al PhysRevC.101.034901(2020)]

[Nonaka et al Physical Review C, 71(4), 044904.(2005)]

[Guida et al Nuclear Physics B, 489(3), 626-652.(1997)]

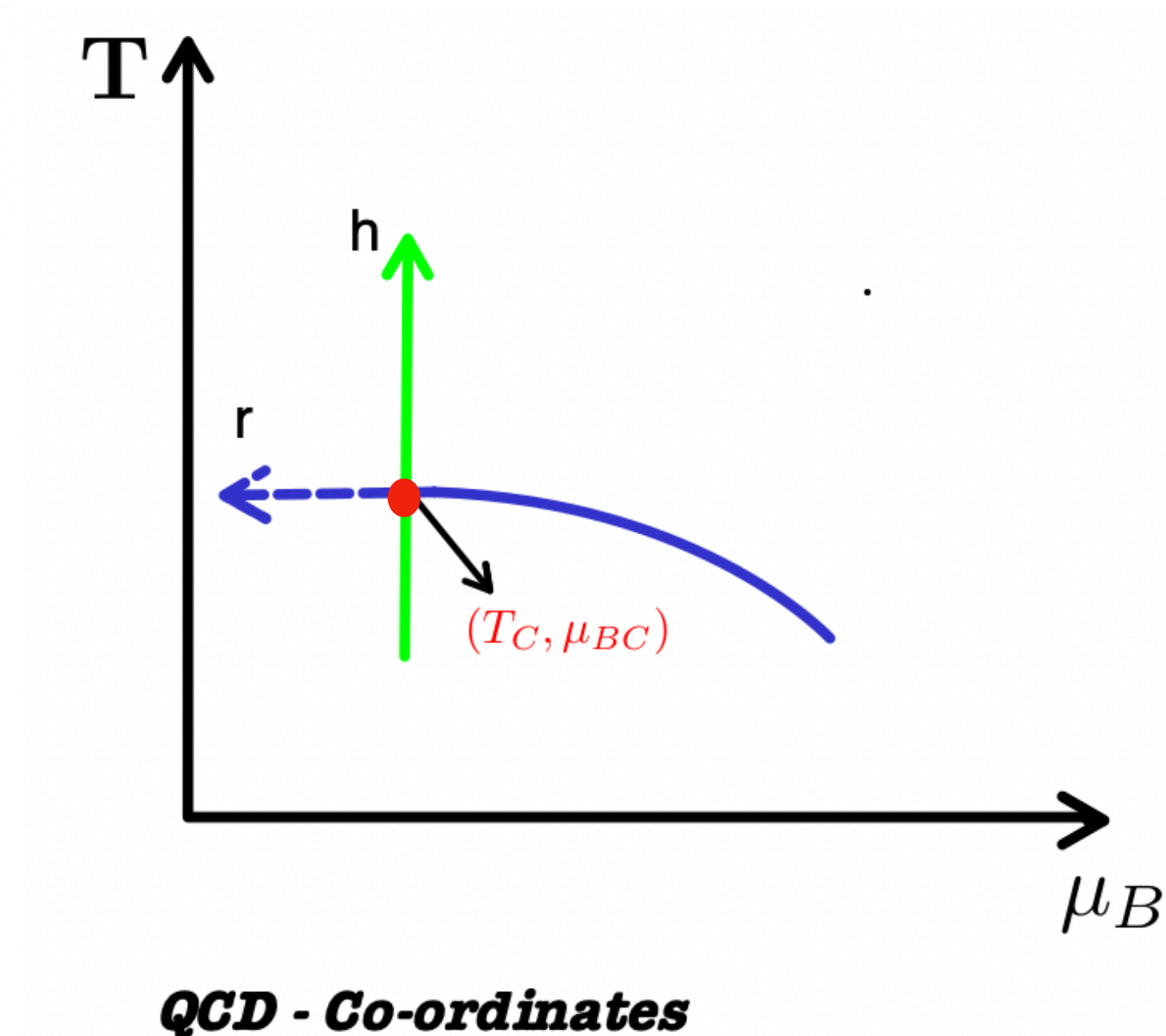
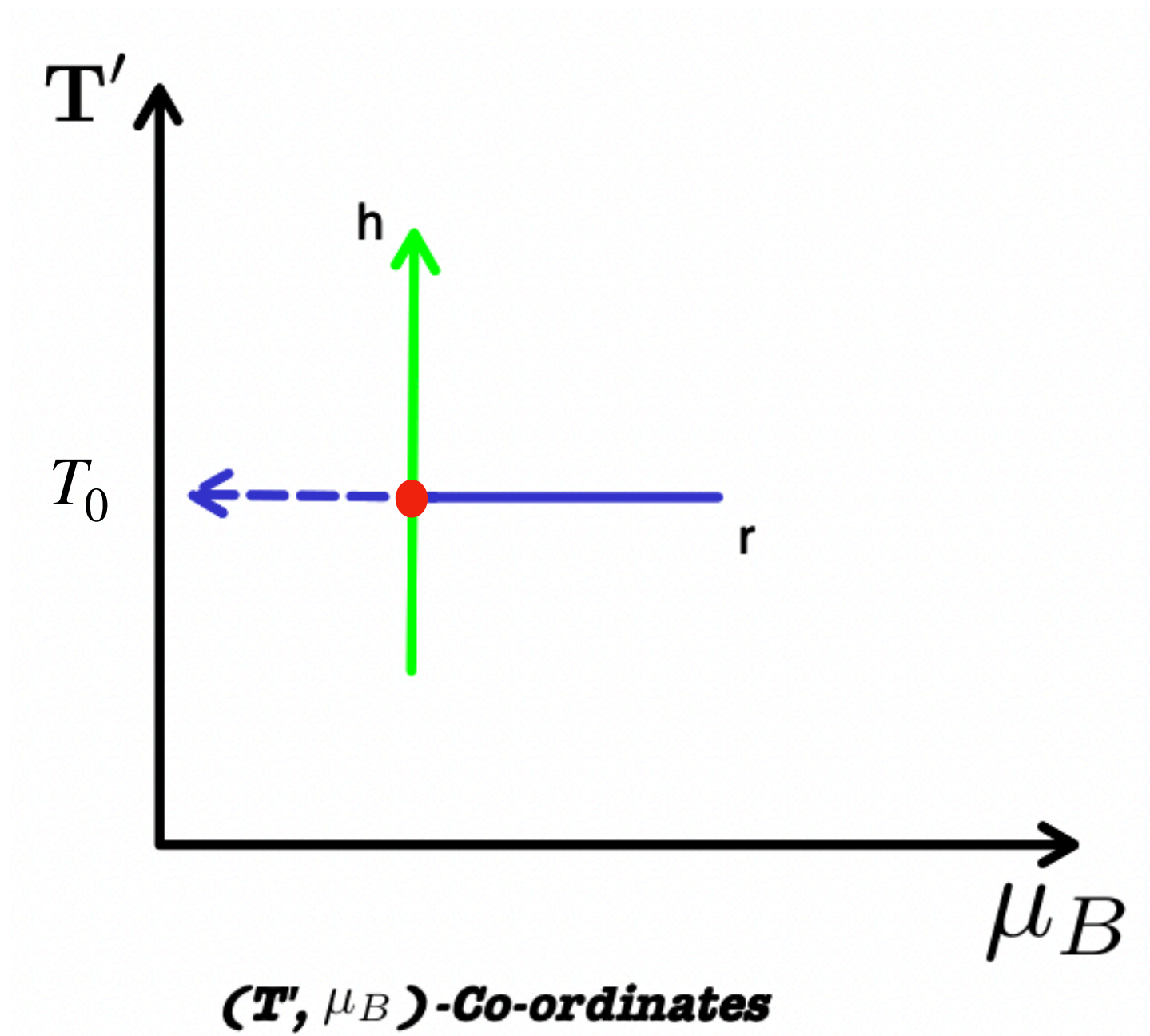
EoS with a Critical point

3D Ising to QCD Mapping



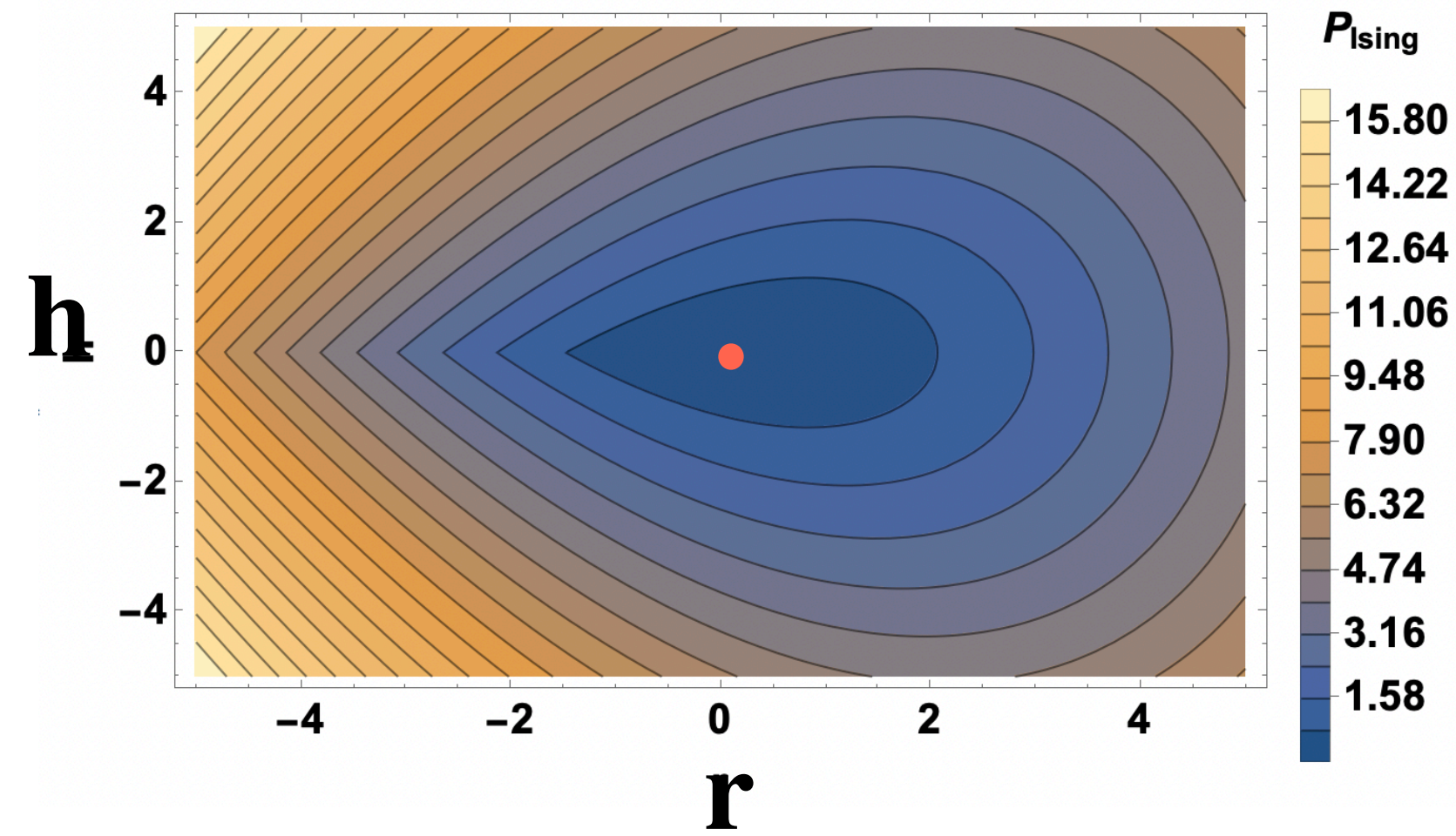
$$\frac{T' - T_0}{T_0} = W h \sin \alpha_{12}$$

$$\frac{\mu_B - \mu_{BC}}{T_0} = W (-r\rho - h \cos \alpha_{12})$$



$$T' = T \left[1 + \left(\frac{\mu_B}{T} \right)^2 \kappa_2^{BB}(T) + \mathcal{O} \left(\frac{\mu_B}{T} \right)^4 \right]$$

Ising Pressure

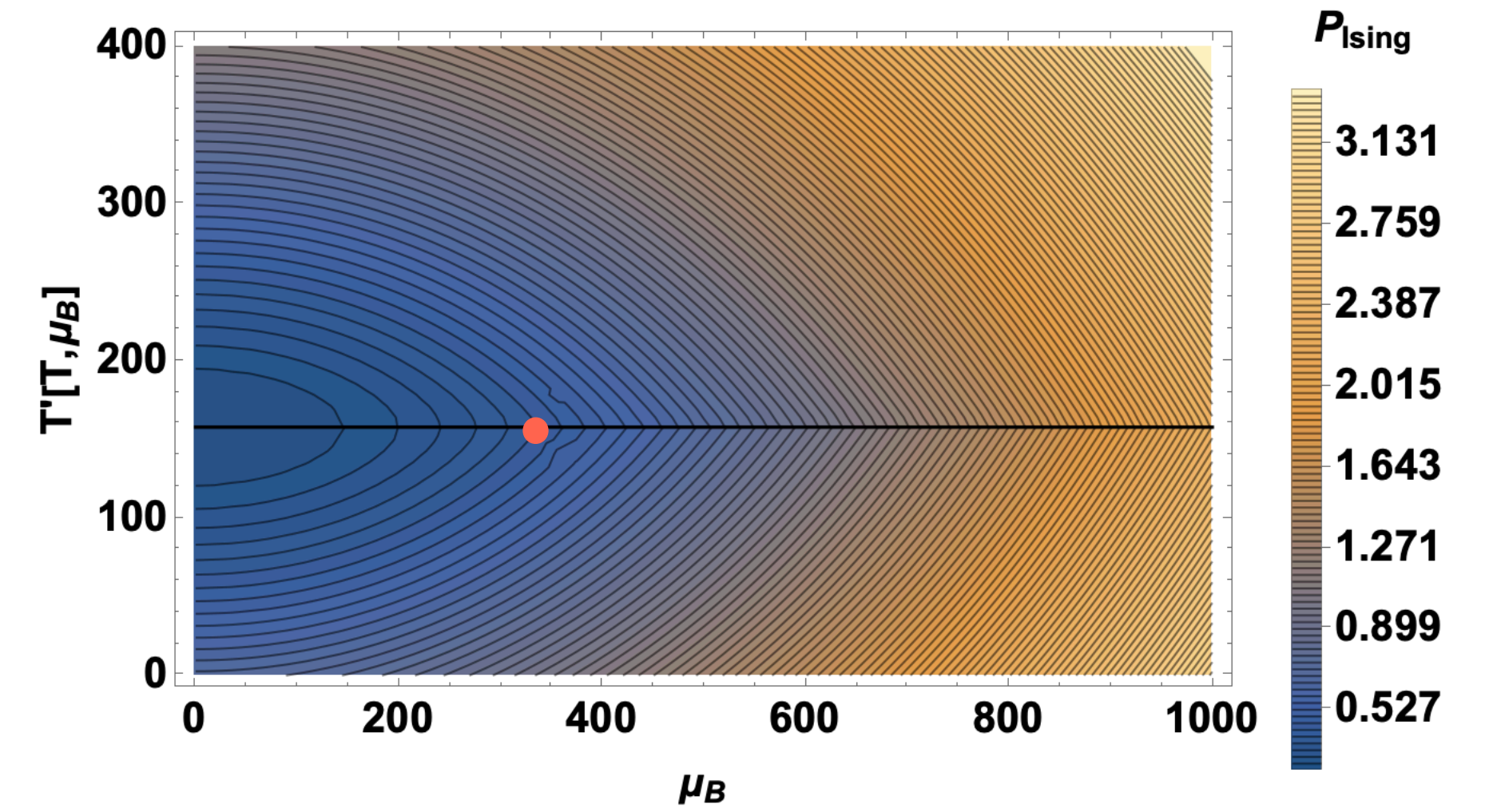


Ising Coordinates

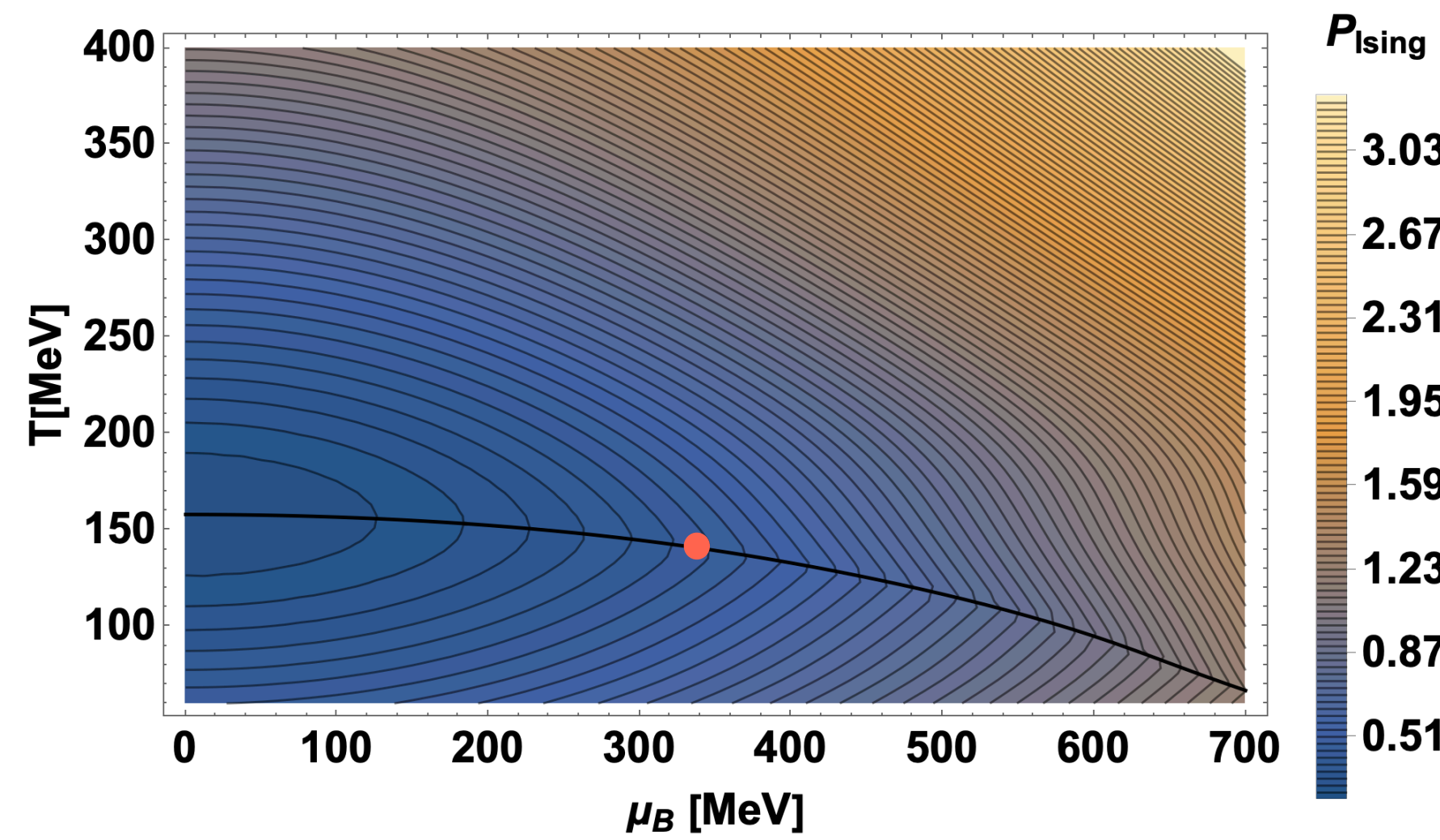
Parameters

$$w = 1, \rho = 2, \mu_{BC} = 350 \text{ MeV}, T_0 = 158 \text{ [MeV]}$$

$$T_C(\mu_B) = T_0 \left[1 - \kappa_2(T_0) \left(\frac{\mu_B}{T_0} \right)^2 \right]$$



(μ_B, T') Coordinates



QCD Coordinates



Re-Constructing the Full Baryon Density

$$\frac{n_B^{full}(T, \mu_B)}{T^3} = \left(\frac{\mu_B}{T}\right) \chi_{2,Lattice}^B(T'_{full}(T, \mu_B), 0)$$

$$T'_{full}(T, \mu_B) = \underbrace{T'_{Lattice}(T, \mu_B)}_{\text{lowest order in } \left(\frac{\mu_B}{T}\right)} + \underbrace{T'_{crit}(T, \mu_B) - Taylor[T'_{crit}(T, \mu_B)]}_{\text{higher orders in } \left(\frac{\mu_B}{T}\right)}$$

Lattice Term
Ising Term

Introducing a Critical Point

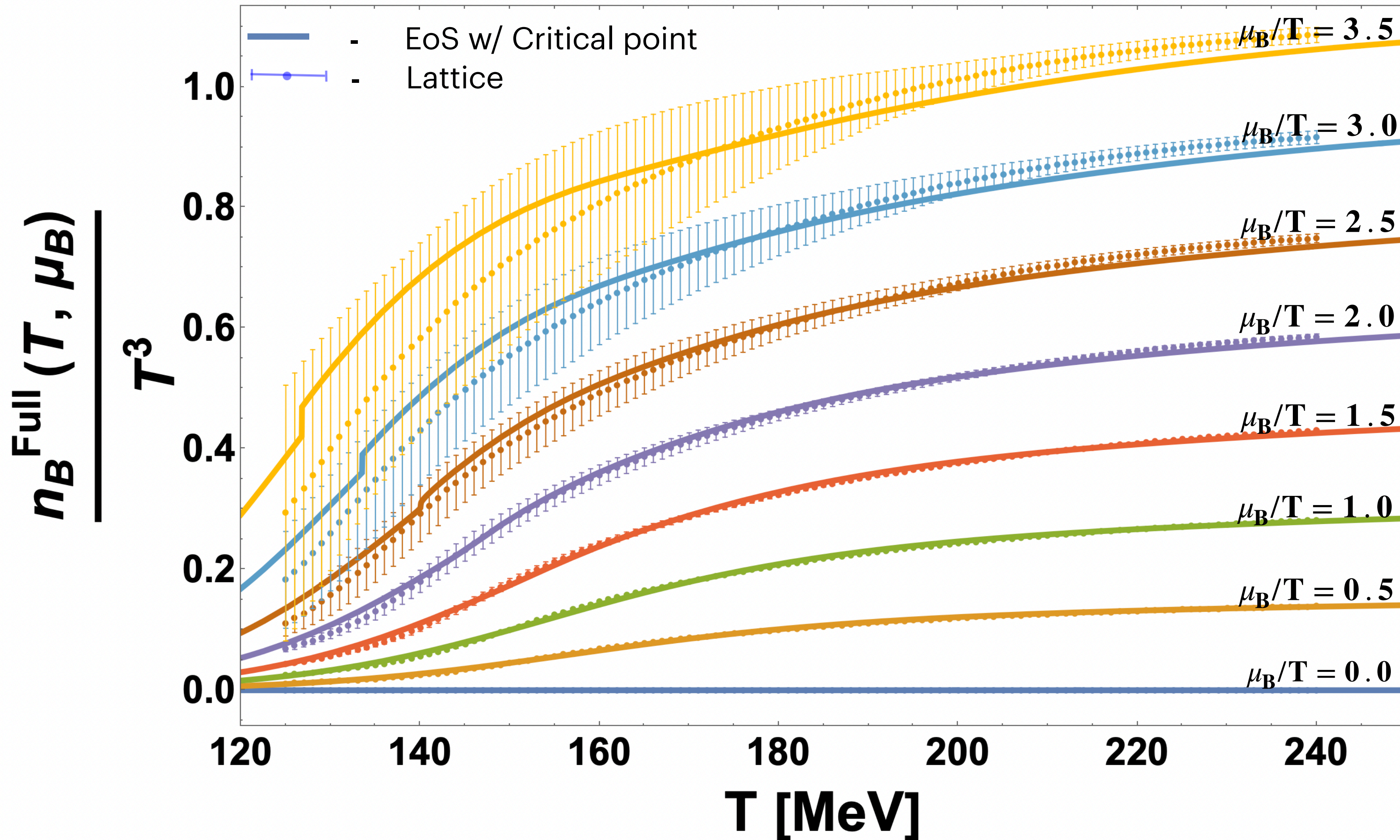
$$T'_{crit}(T, \mu_B) \approx T_0 + \left(\left. \frac{\partial \chi_{2,lattice}^B(T, 0)}{\partial T} \right|_{T=T_0} \right)^{-1} \frac{n_B^{crit}(T, \mu_B)}{T^3(\mu_B/T)} + \dots$$

$$\chi_1^{crit}(T, \mu_B) = \frac{n_B^{crit}(T, \mu_B)}{T^3} = \frac{\partial(P^{crit}(T, \mu_B)/T^4)}{\partial(\mu_B/T)}$$

Baryon density results

Full Baryon Density at a constant $\frac{\mu_B}{T}$ compared with Lattice

$\mu_B = 350$ [MeV], $\alpha_{12} = 90$, $\rho = 2$, $w = 2$

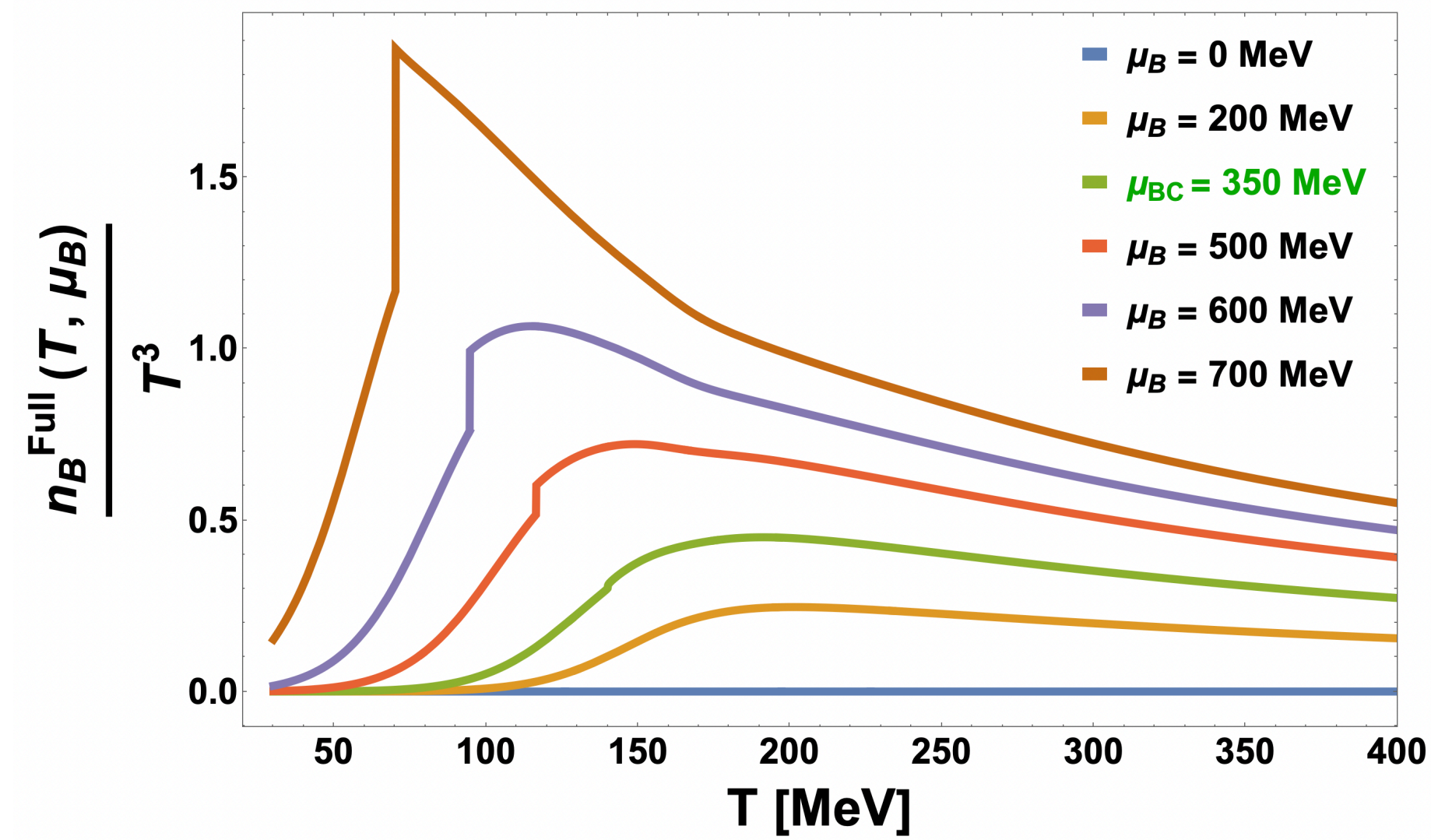


Baryon density results

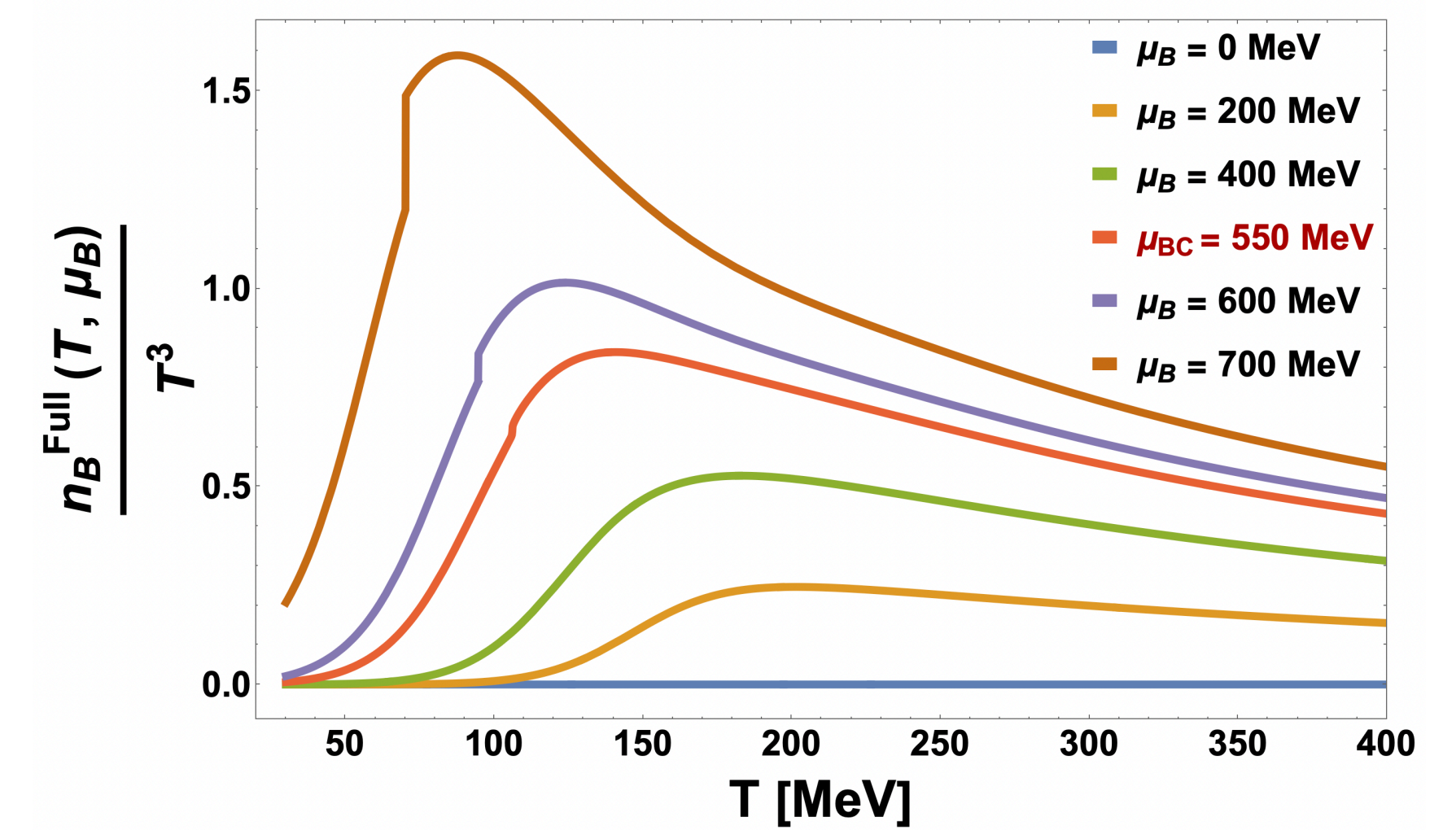
Baryon Density at a constant μ_B for different μ_{BC}

$\alpha_{12} = 90, \rho = 2, w = 2$

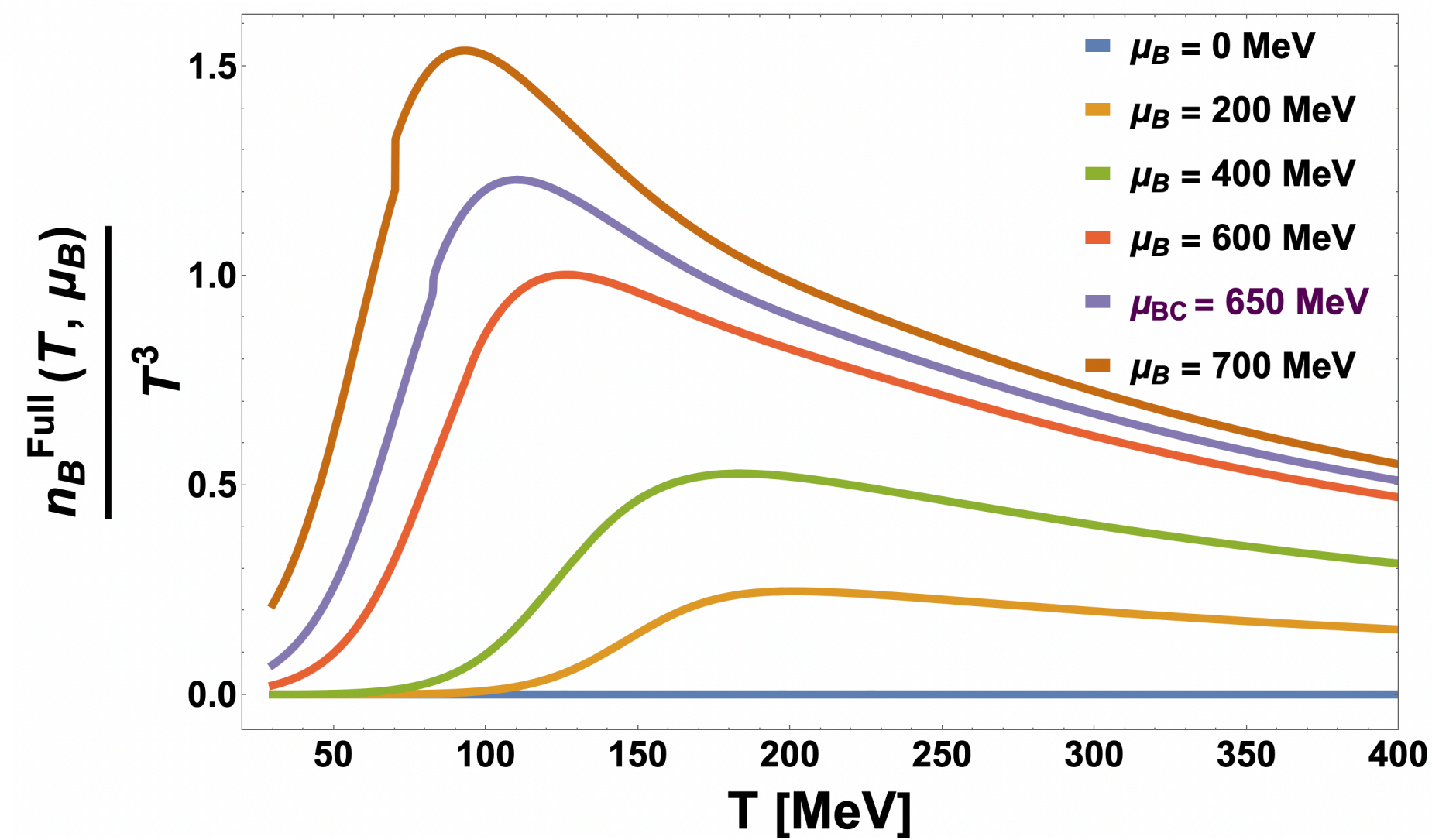
$\mu_{BC} = 350 \text{ MeV}$



$\mu_{BC} = 550 \text{ MeV}$



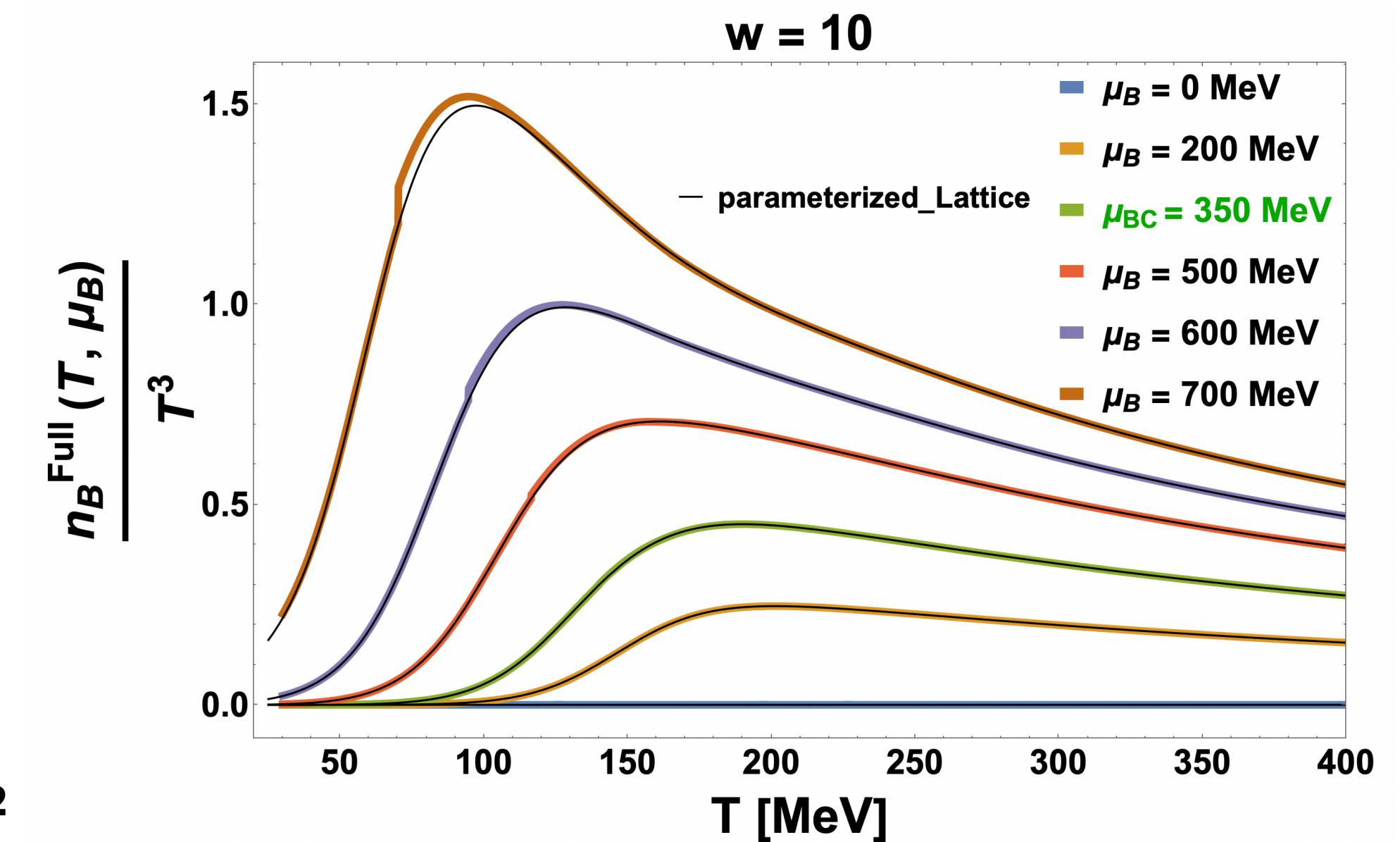
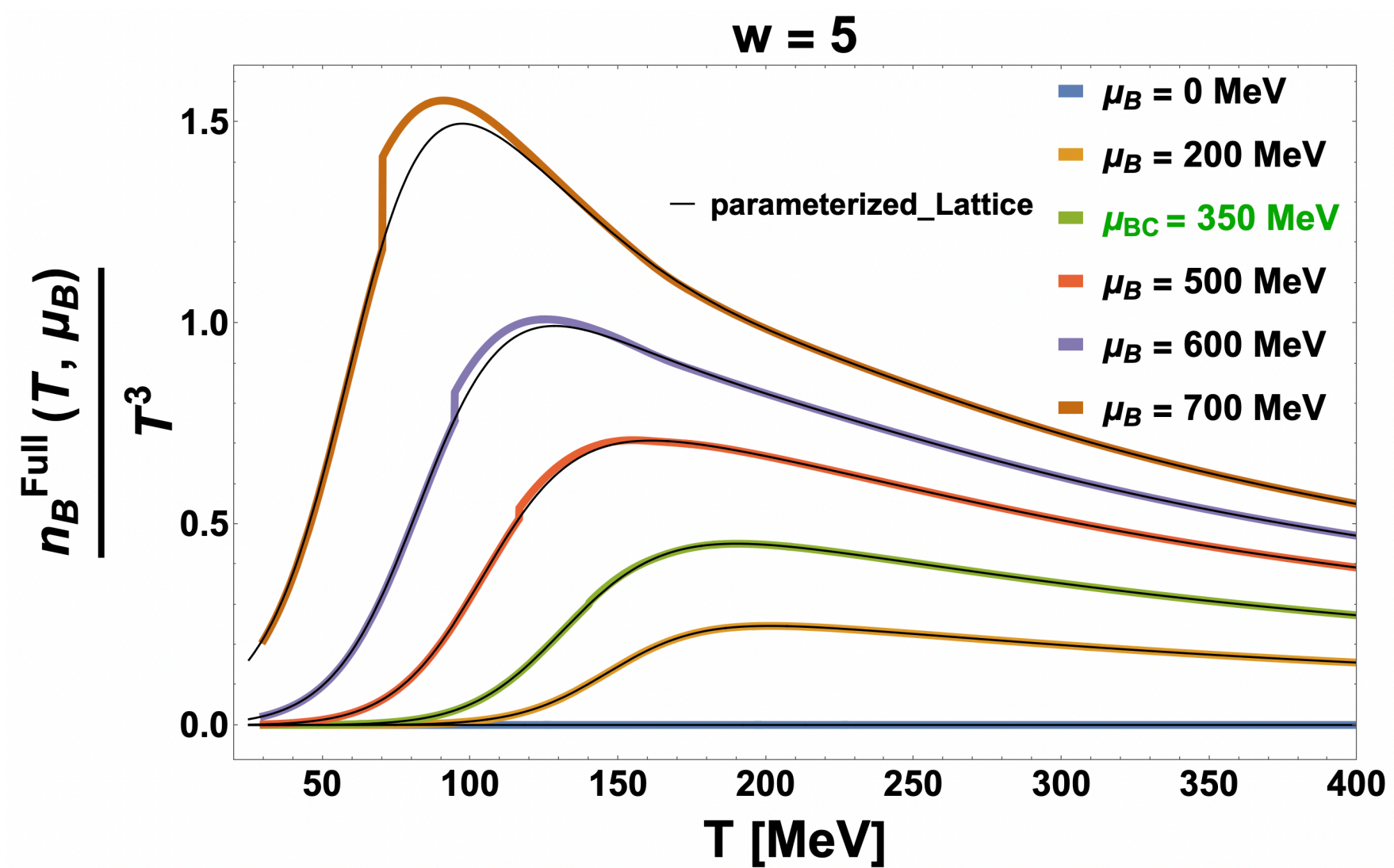
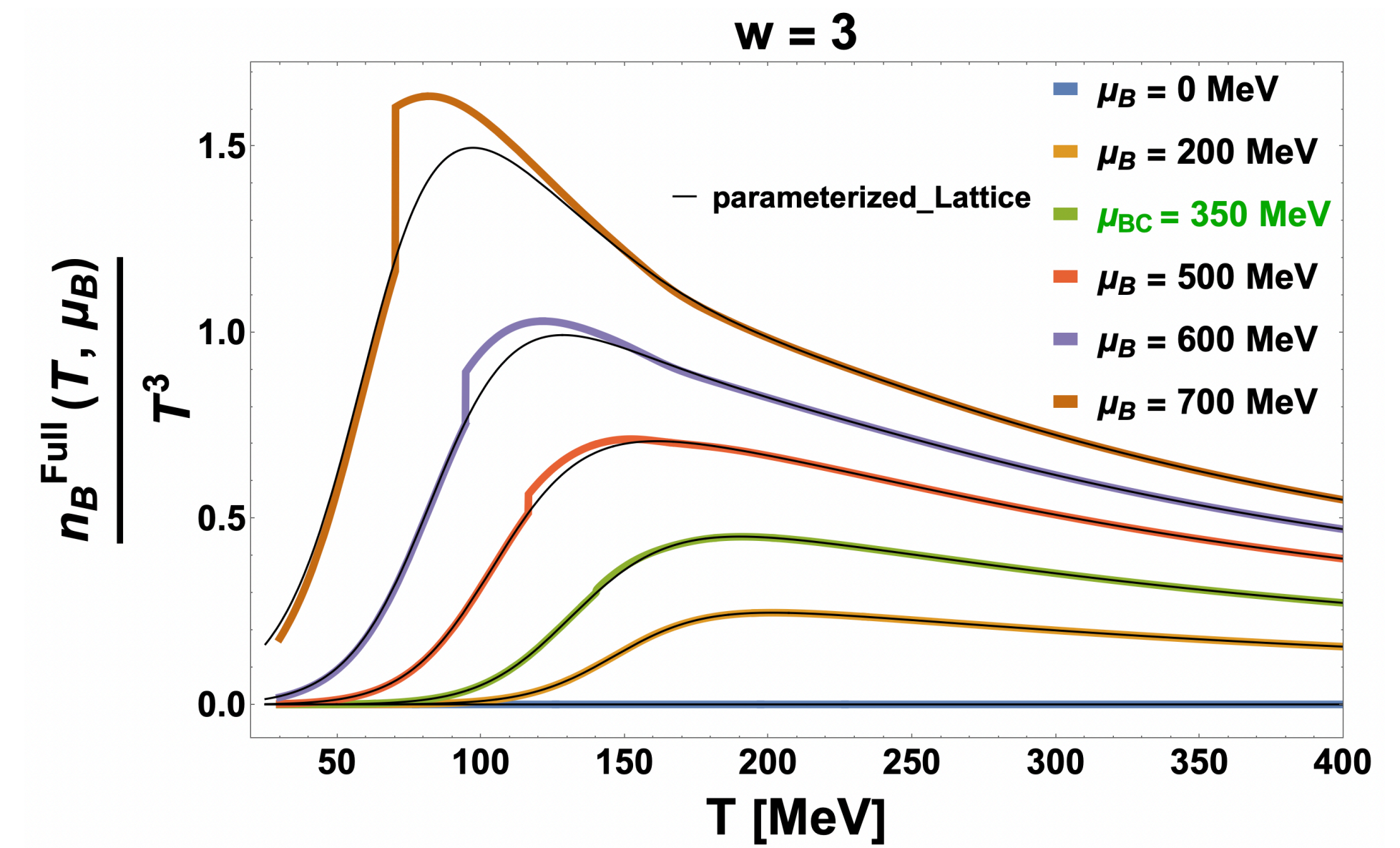
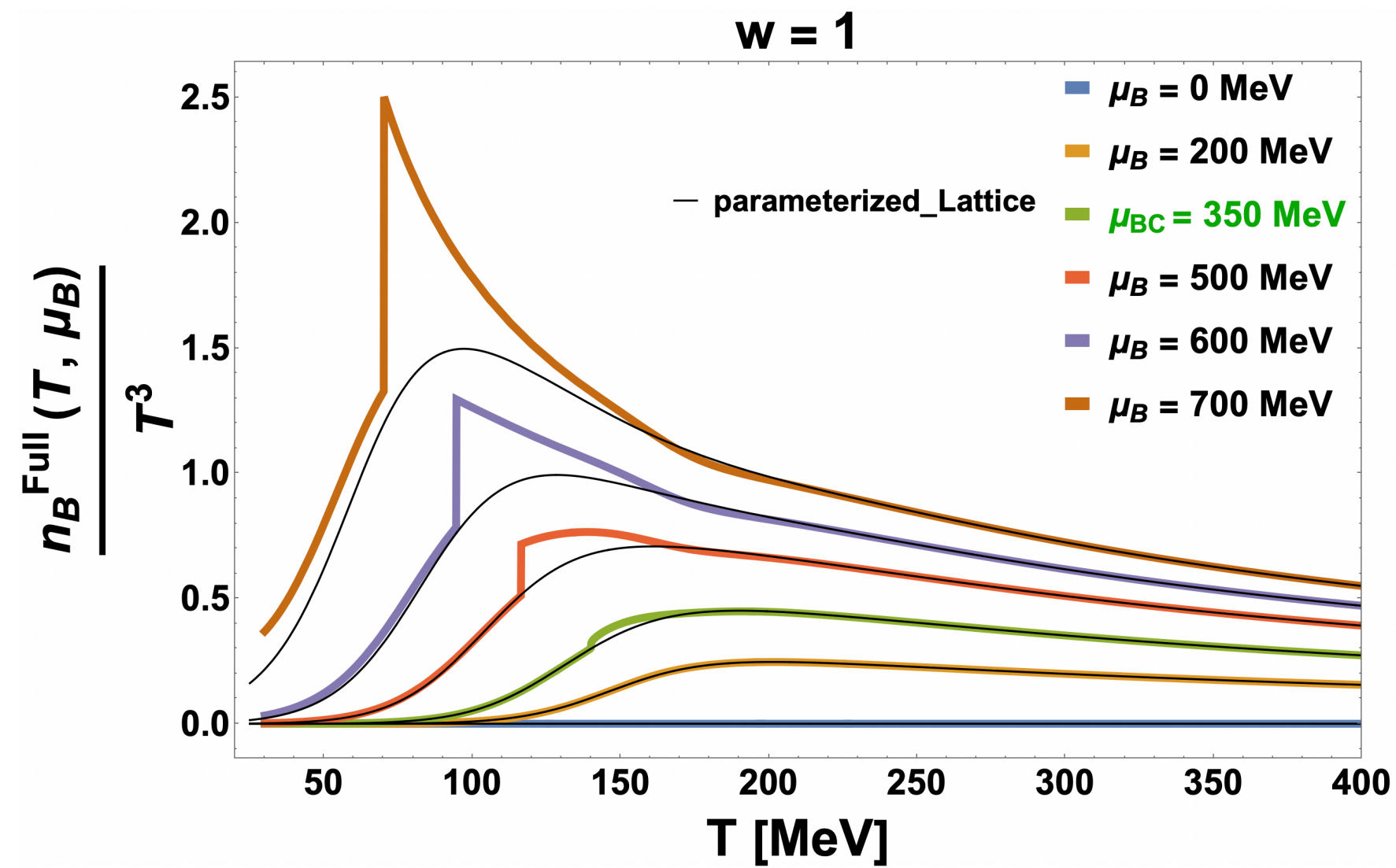
$\mu_{BC} = 650 \text{ MeV}$



Estimating the Critical contribution

Baryon Density at a constant μ_B for 1 to w=10

$\mu_{BC} = 350$ [MeV], $\alpha_{12} = 90$, $\rho = 2$



Thermodynamic Relations

$$\frac{P(T, \mu_B)}{T^4} = \frac{P(T, 0)}{T^4} + \int_0^{\mu_B} d\hat{\mu}'_B \frac{n_B(T, \hat{\mu}'_B)}{T^3}$$

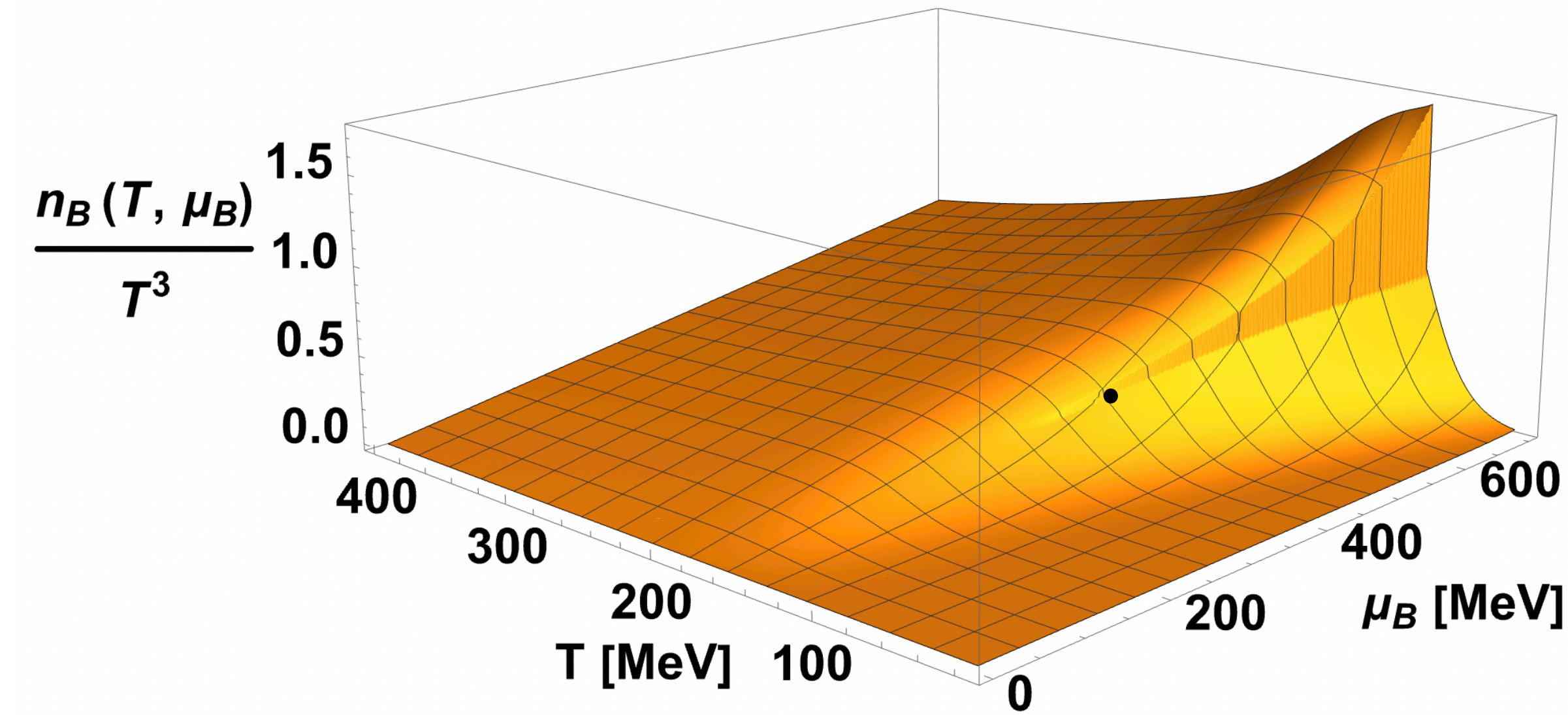
$$\frac{S(T, \mu_B)}{T^3} = \frac{1}{T^3} \left(\frac{\partial P}{\partial T} \right) \Big|_{\mu_B}$$

$$\frac{\epsilon(T, \mu_B)}{T^3} = \frac{S}{T^3} - \frac{P}{T^4} + \frac{\mu_B}{T} \frac{n_B}{T^3}$$

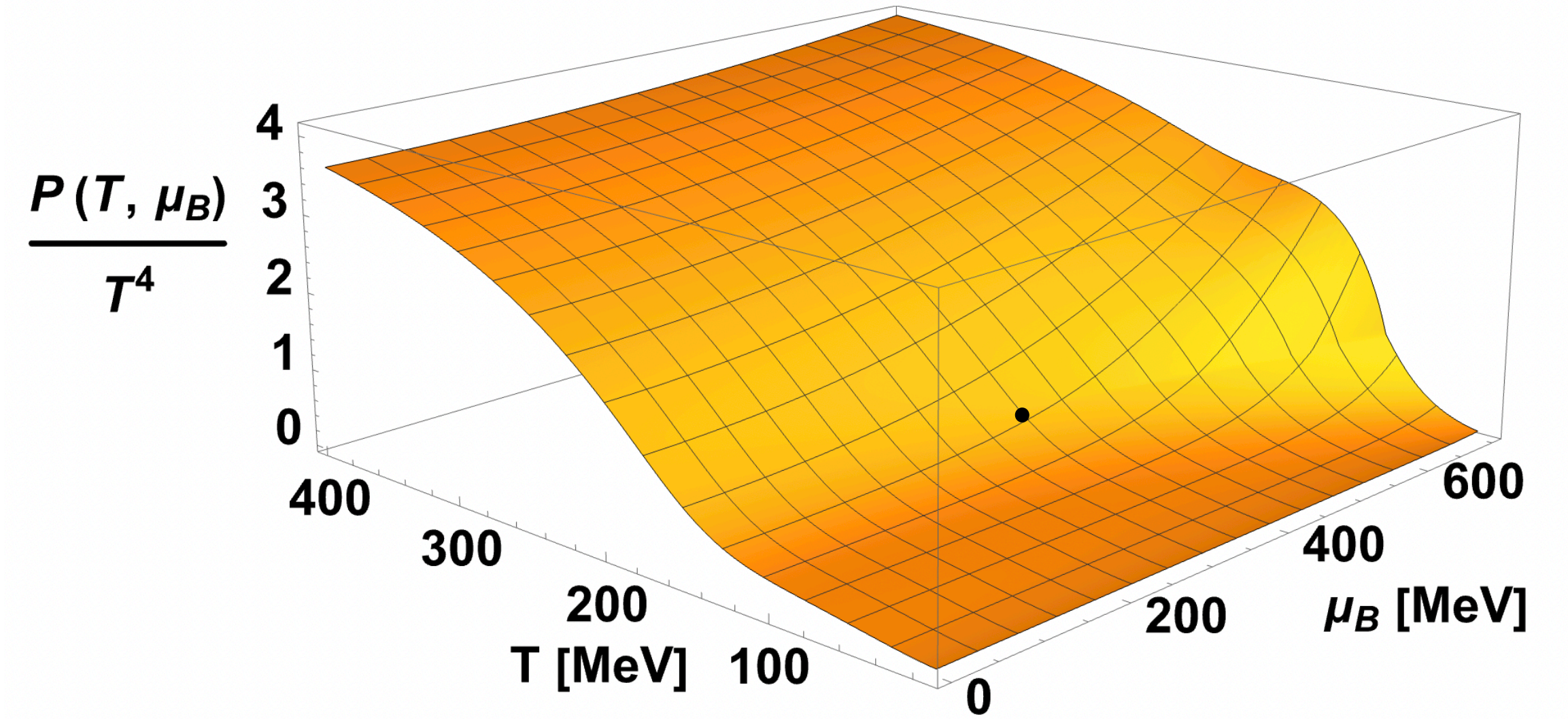
$$\frac{\chi_2(T, \mu_B)}{T^2} = T \left(\frac{\partial (n_B/T^3)}{\partial \mu_B} \right) \Big|_{\mu_B}$$

Thermodynamic observables

Baryon Density



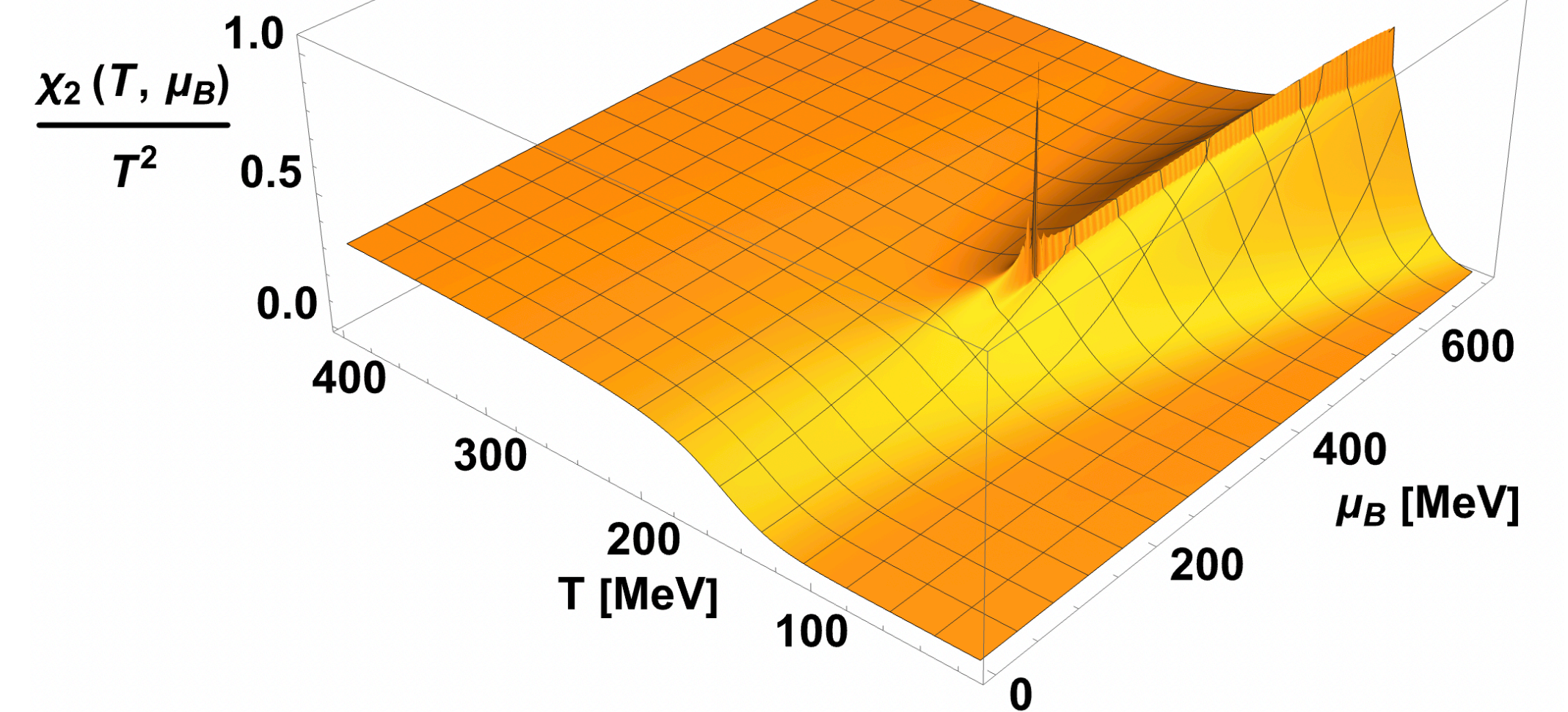
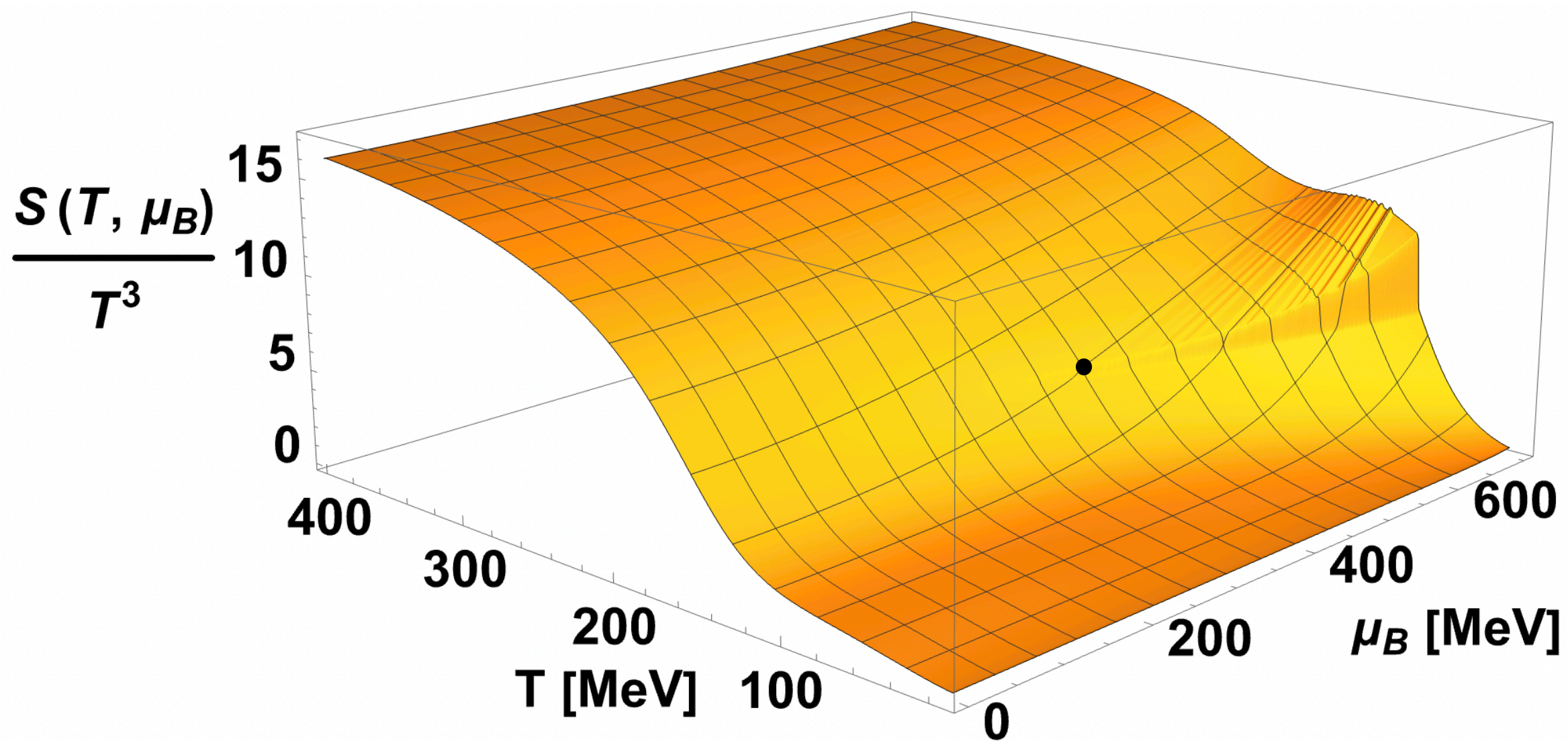
Pressure



• **Critical Point**

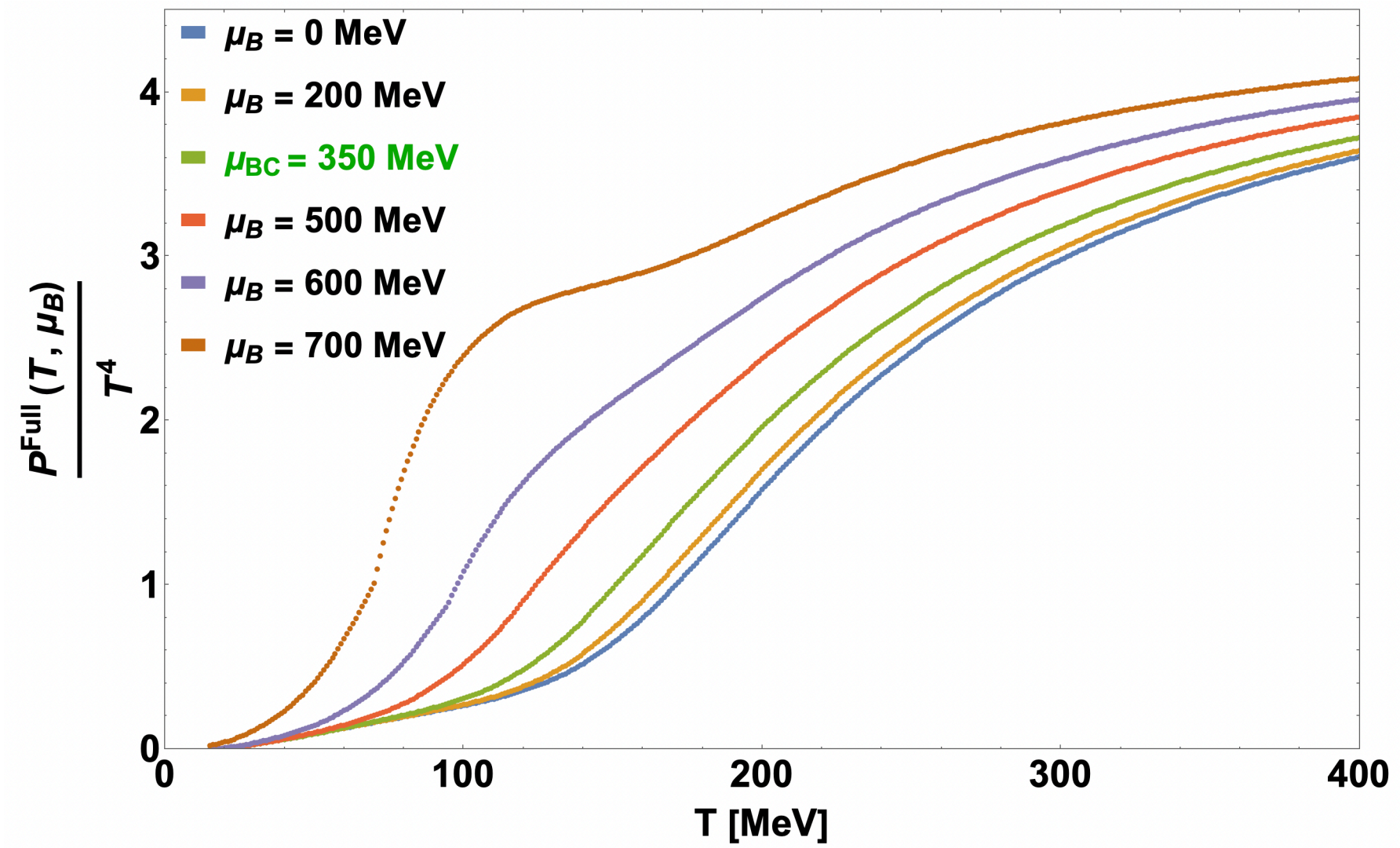
$$w = 2, \rho = 2 \text{ \& } \alpha_{12} = 90$$

Entropy Density

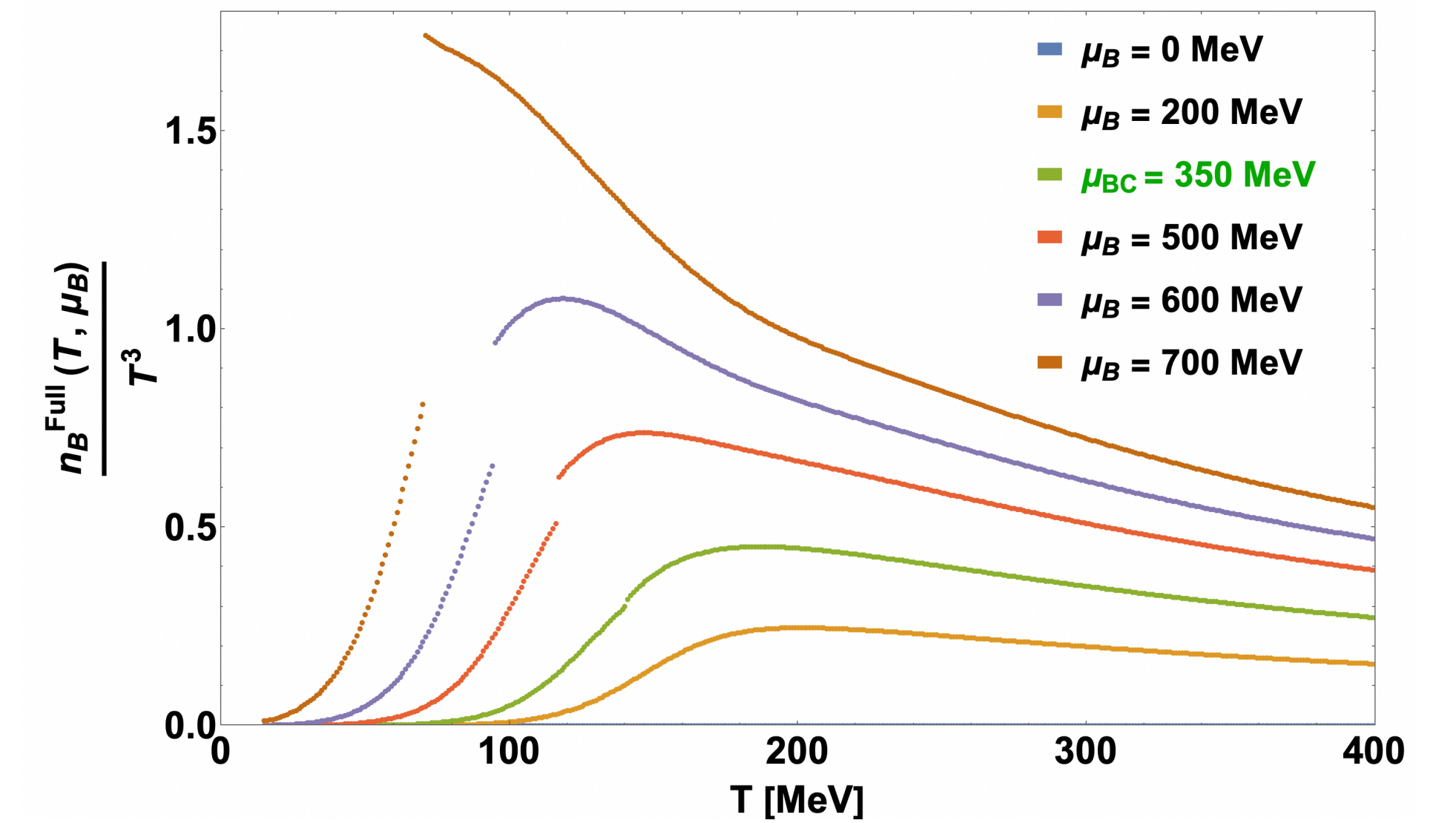


Thermodynamic observables

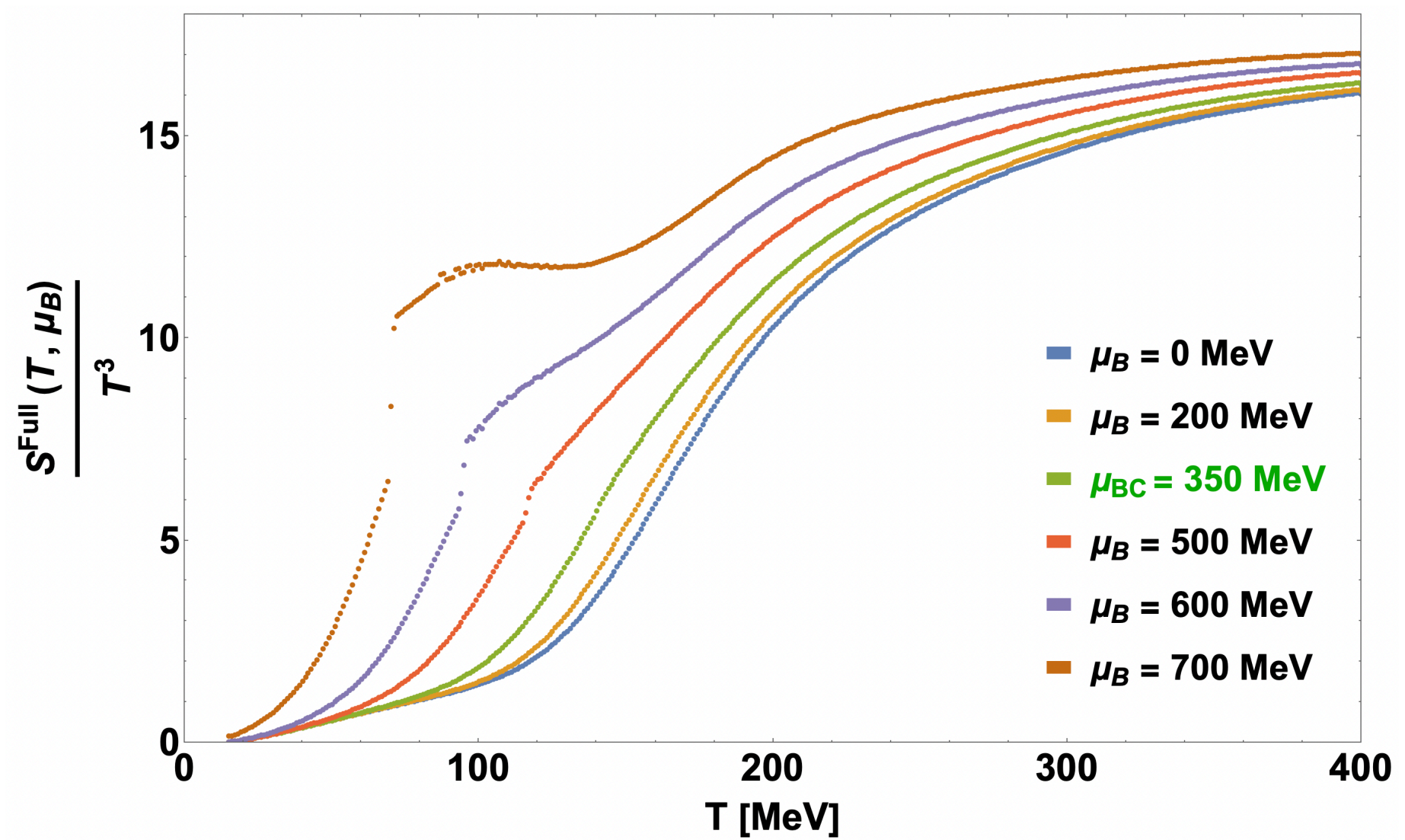
Pressure



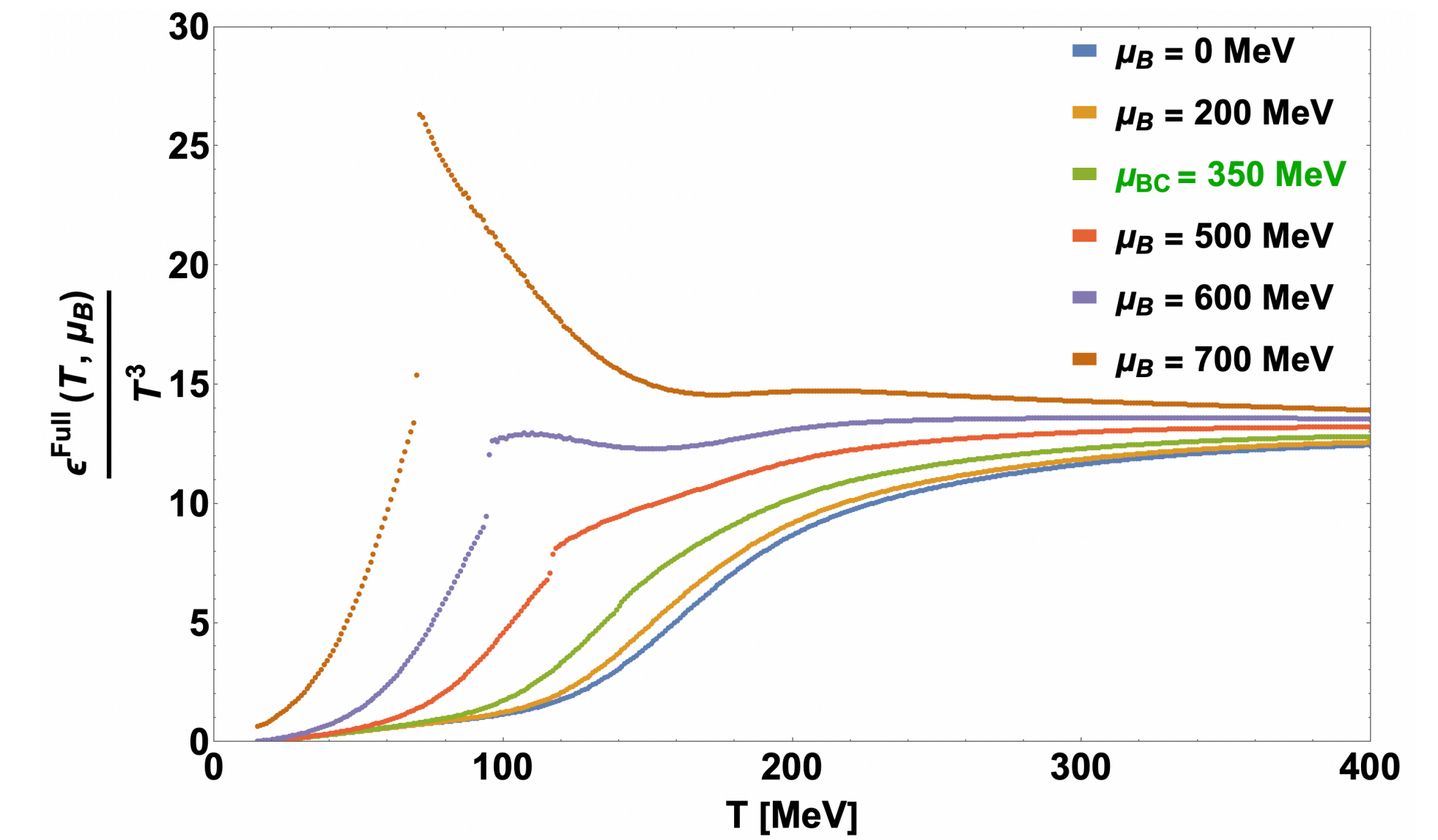
Baryon density



Entropy density



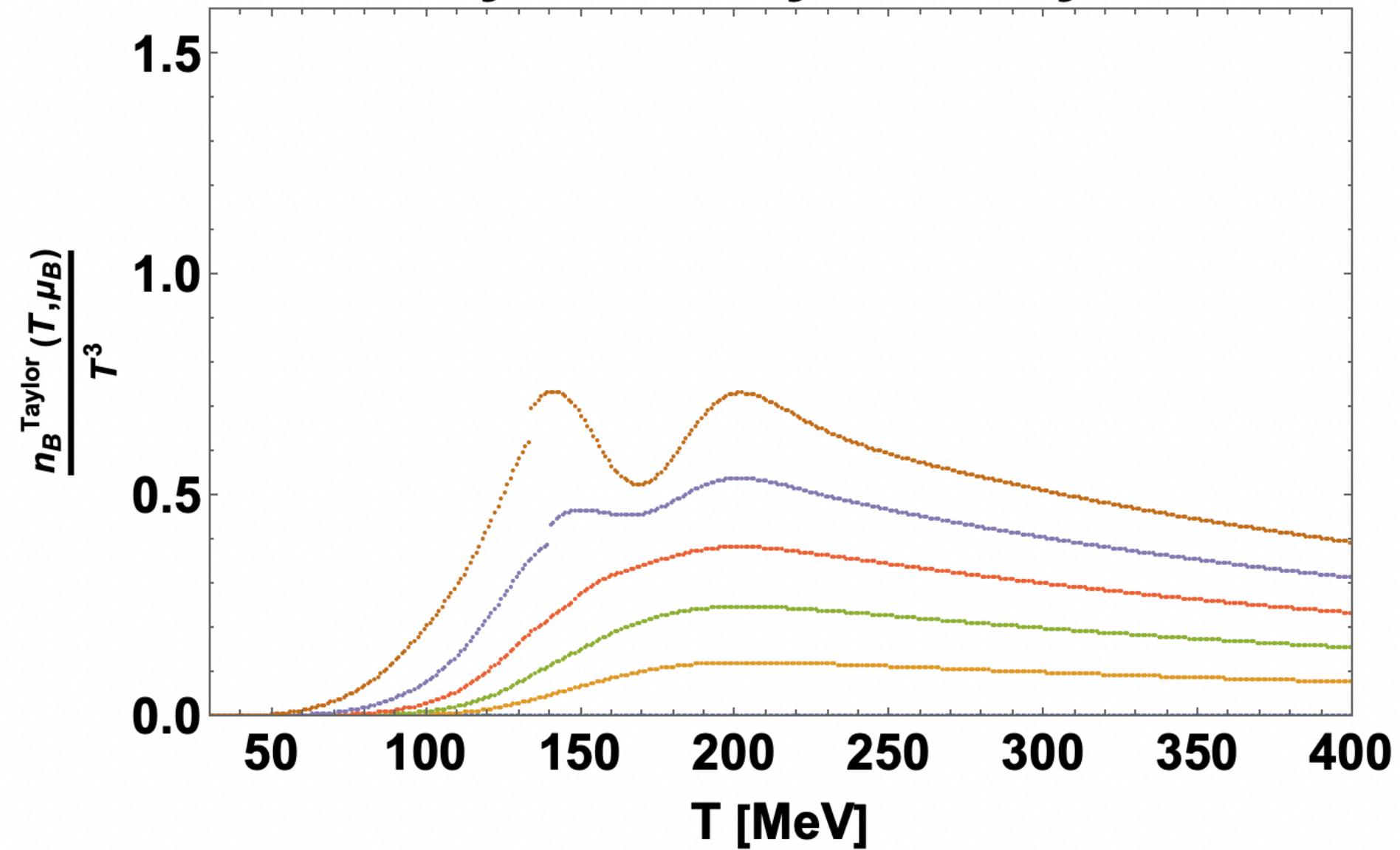
Energy density



Summary

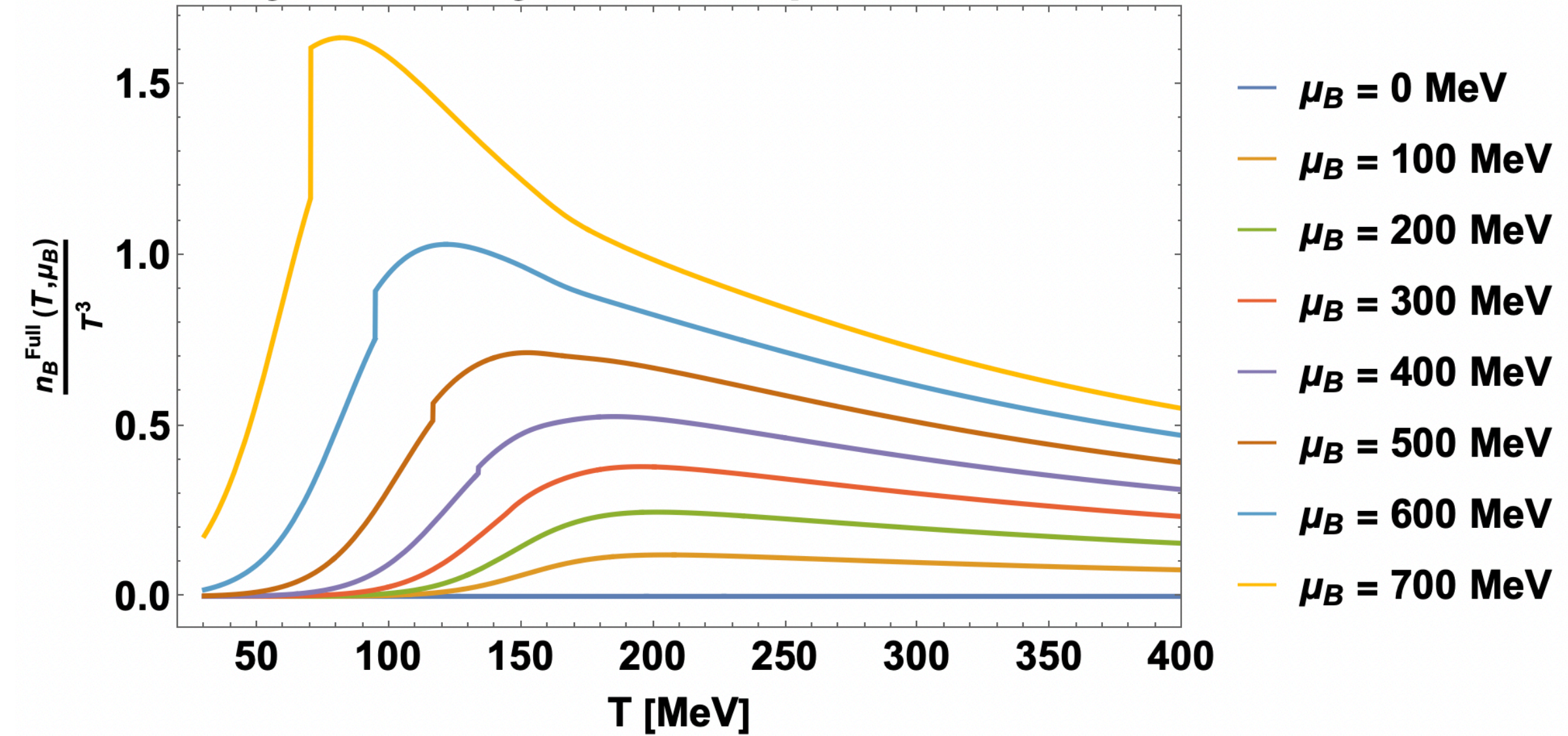
$\mu_B = 350$ [MeV], $\alpha_{12} = 90$, $w = 3$, $\rho = 2$

Baryon Density from Taylor



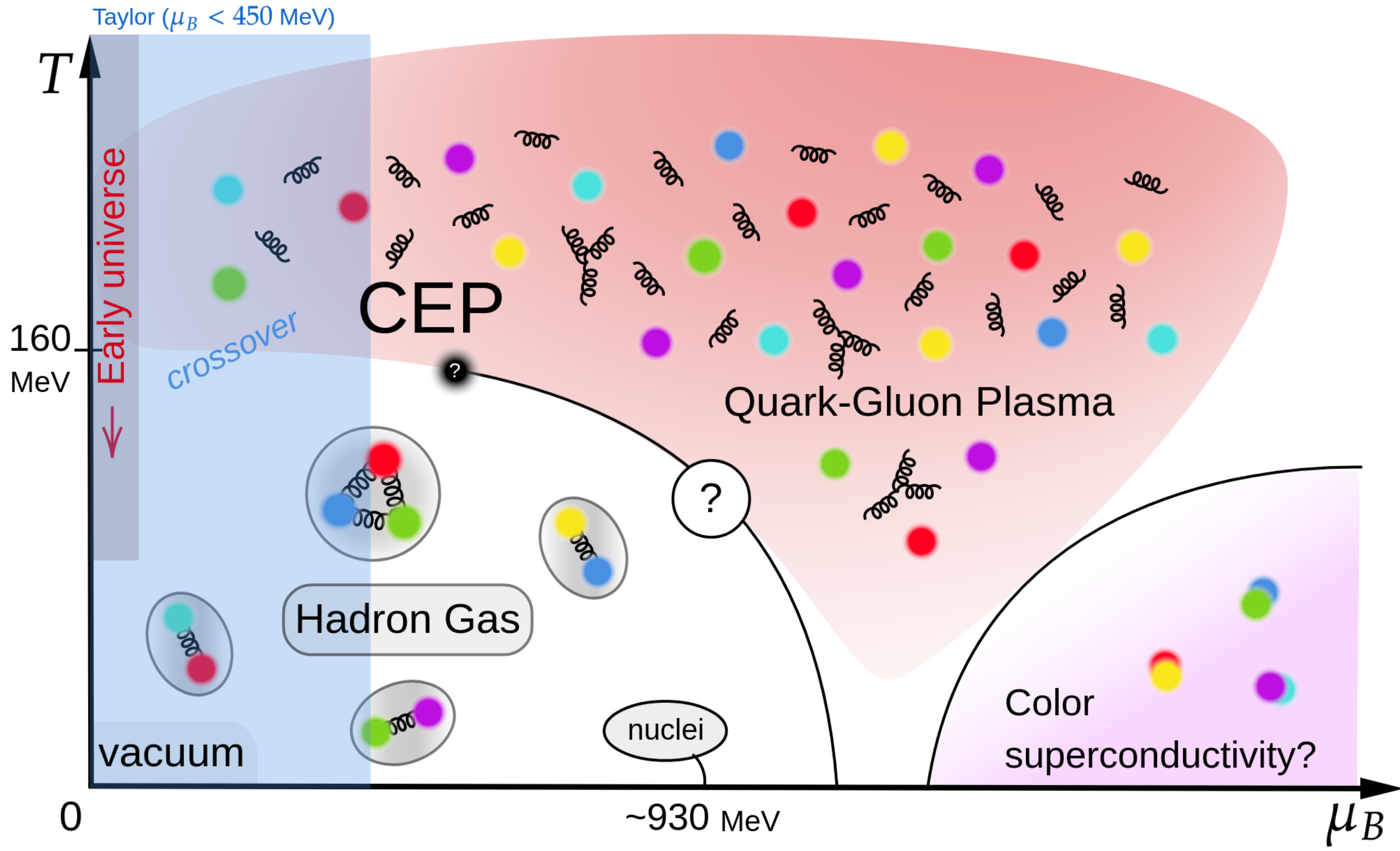
- $\mu_B = 0$ MeV
- $\mu_B = 100$ MeV
- $\mu_B = 200$ MeV
- $\mu_B = 300$ MeV
- $\mu_B = 400$ MeV
- $\mu_B = 500$ MeV

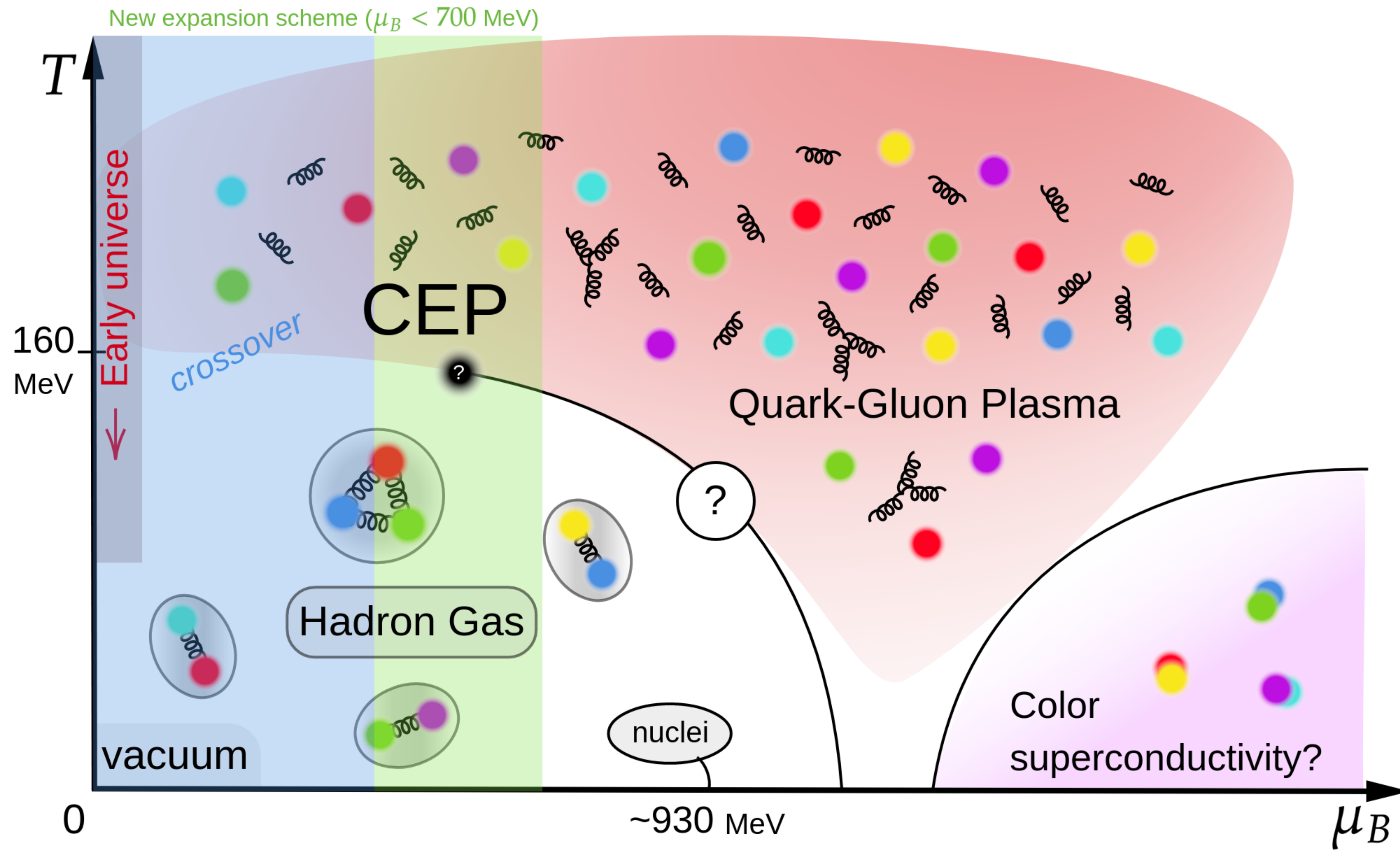
Baryon Density from T-Expansion Scheme



- $\mu_B = 0$ MeV
- $\mu_B = 100$ MeV
- $\mu_B = 200$ MeV
- $\mu_B = 300$ MeV
- $\mu_B = 400$ MeV
- $\mu_B = 500$ MeV
- $\mu_B = 600$ MeV
- $\mu_B = 700$ MeV

- **A more physical EoS that captures a large part of the Phase diagram is required.**
- **We provide a family of EoS with a correct Critical point up $\mu_B = 700$ MeV.**
- **Our EoS allows users to change parameters and compare with the data from the Experiment (Beam Energy Scan II)**





Future Work

- **Explore and constrain the Parameters space by requesting thermodynamics Stability and causality of our EoS**

$$P(T, \mu_B) > 0, \quad n_B(T, \mu_B) > 0, \quad s(T, \mu_B) > 0 \quad \epsilon(T, \mu_B) > 0 \quad c_s^2(T, \mu_B) > 0$$

$$c_s^2(T, \mu_B) < 1$$

- **Merge with Nuclear Matter EoS at low Temperatures**

Thank you for your attention!

Back up!

Energy Density

