

# Classical and Statistical Physics version of Proton's pressure distribution

Sabyasachi Ghosh  
(Assitant Professor,  
Physics Department, IIT Bhilai)



Cho Win Aung  
(Asean PhD, IIT  
Bhilai)



Thandar Zaw Win  
(Asean PhD, IIT  
Bhilai)



Jabed Umar  
(BS-MS, NISER)



Deependra Singh  
(BS-MS, NISER)



Motivation came from .....

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Letter | Published: 16 May 2018

**The pressure distribution inside the proton**

V. D. Burkert ✉, L. Elouadrhiri & F. X. Girod

Nature 557, 396–399 (2018) | Cite this article

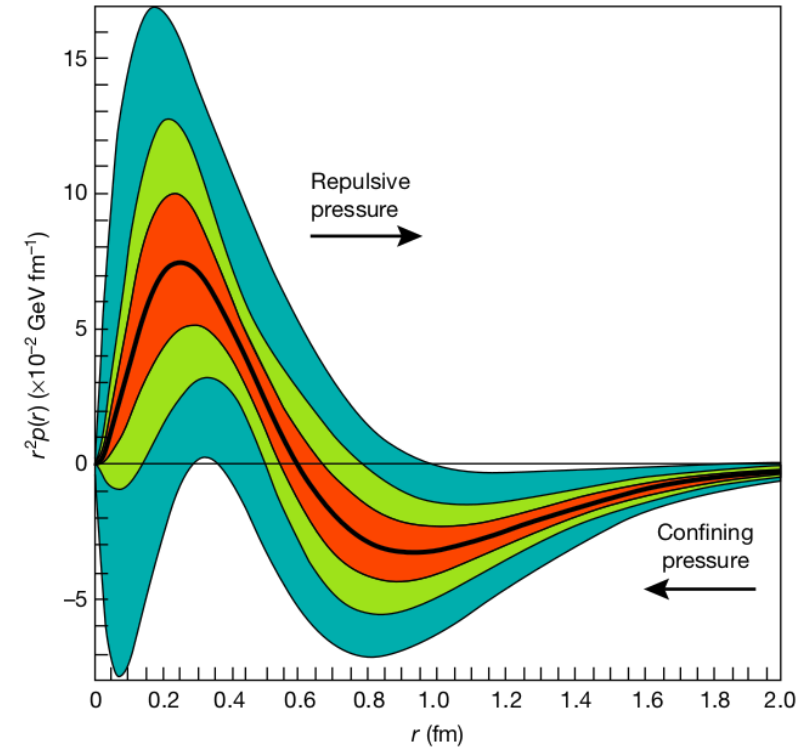
Thanks to....

## QCD with Electron-Ion Collider (QEIC)

Jan 4 – 7, 2020

IIT Bombay

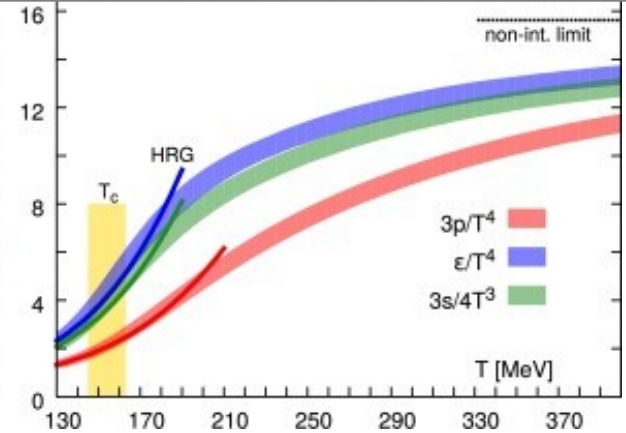
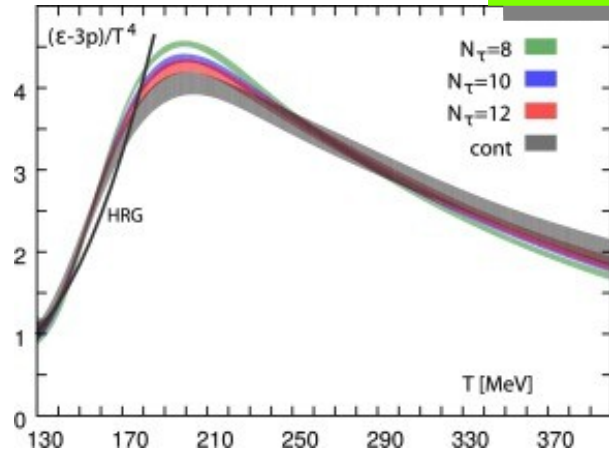
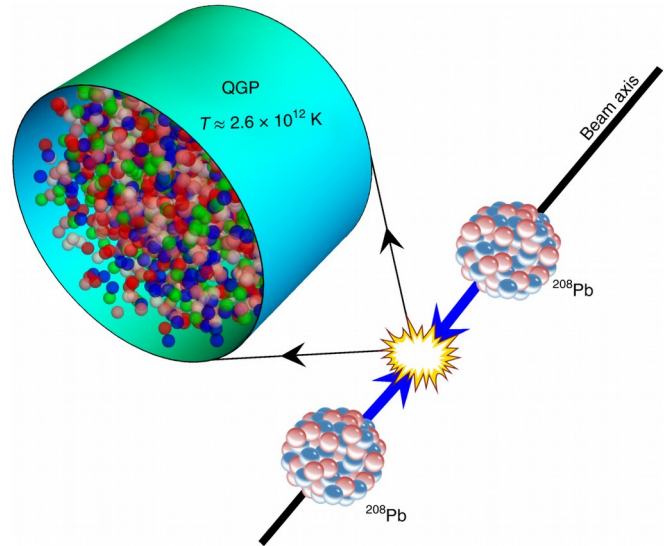
Asia/Kolkata timezone



**Fig. 1 | Radial pressure distribution in the proton.** The graph shows the pressure distribution  $r^2 p(r)$  that results from the interactions of the quarks in the proton versus the radial distance  $r$  from the centre of the proton. The thick black line corresponds to the pressure extracted from the D-term parameters fitted to published data<sup>22</sup> measured at 6 GeV. The corresponding estimated uncertainties are displayed as the light-green shaded area shown. The blue area represents the uncertainties from all the data that were available before the 6-GeV experiment, and the red shaded area shows projected results from future experiments at 12 GeV that will be performed with the upgraded experimental apparatus<sup>30</sup>. Uncertainties represent one standard deviation.

# Heavy Ion Collision (HIC) Direction.....

## Pressure of RHIC or LHC Matter



### 2.3.2 Pressure

The general expression

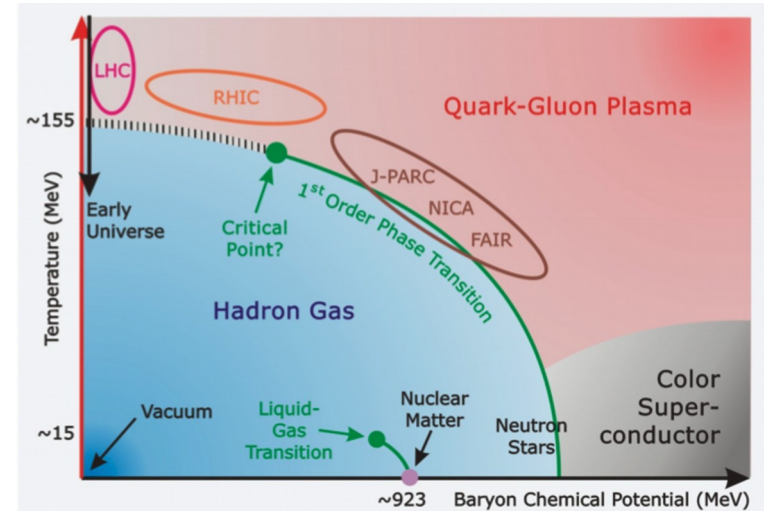
- Case:1  $T=0K$

$$P_{UR} = \frac{2}{\pi^2} T^4 f_4(A) \quad (37)$$

$$P_{UR} = \frac{1}{12\pi^2} \mu^4 \quad (38)$$

- Case:2  $\mu = 0$

$$P_{UR} = \frac{1.89}{\pi^2} T^4 \quad (39)$$



## Pressure (Microscopic expression) from Statistical Mechanics

$$\Phi(T, V, \mu) = -PV = -k_B T \ln Z_{GCE}$$

$$P = k_B T \frac{\partial}{\partial V} \ln Z_{GCE}$$

$$P = g \int \frac{d^3 p}{(2\pi)^3} \frac{pv}{3} f_0(\epsilon)$$

Massless  
case

$$\Phi = PV$$

$$= U - TS - \mu N$$

$$\Rightarrow d\Phi = dU - TdS - SdT - \mu dN - Nd\mu$$

$$dU = TdS - PdV + \mu dN$$

$$\Rightarrow d\Phi = -SdT - PdV - Nd\mu$$

### 2.3.2 Pressure

The general expression

$$P_{UR} = \frac{2}{\pi^2} T^4 f_4(A) \quad (37)$$

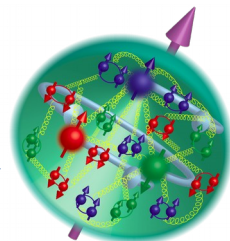
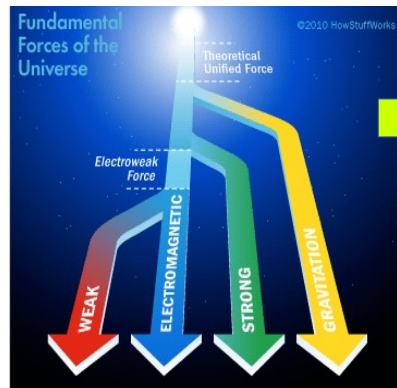
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# Pressure (Microscopic expression) in the direction of Form Factor



em:	$\partial_\mu J_{\text{em}}^\mu = 0$	$\langle N'   J_{\text{em}}^\mu   N \rangle$	$\rightarrow$	$Q = 1.602176487(40) \times 10^{-19} \text{C}$
				$\mu = 2.792847356(23) \mu_N$
weak:	PCAC	$\langle N'   J_{\text{weak}}^\mu   N \rangle$	$\rightarrow$	$g_A = 1.2694(28)$
				$g_p = 8.06(55)$
gravity:	$\partial_\mu T_{\text{grav}}^{\mu\nu} = 0$	$\langle N'   T_{\text{grav}}^{\mu\nu}   N \rangle$	$\rightarrow$	$m = 938.272013(23) \text{ MeV}/c^2$
				$J = \frac{1}{2}$
				$D = ?$

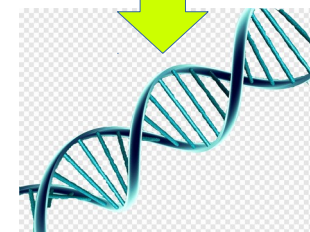
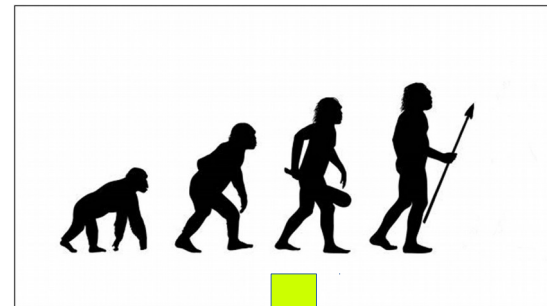


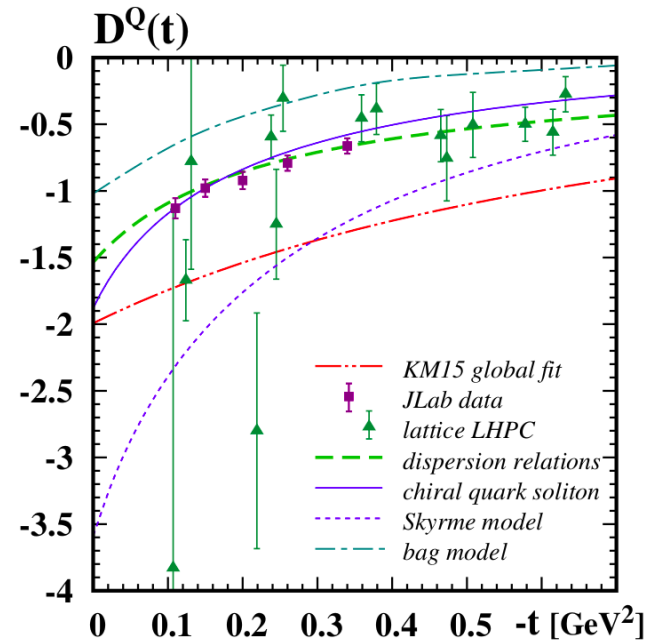
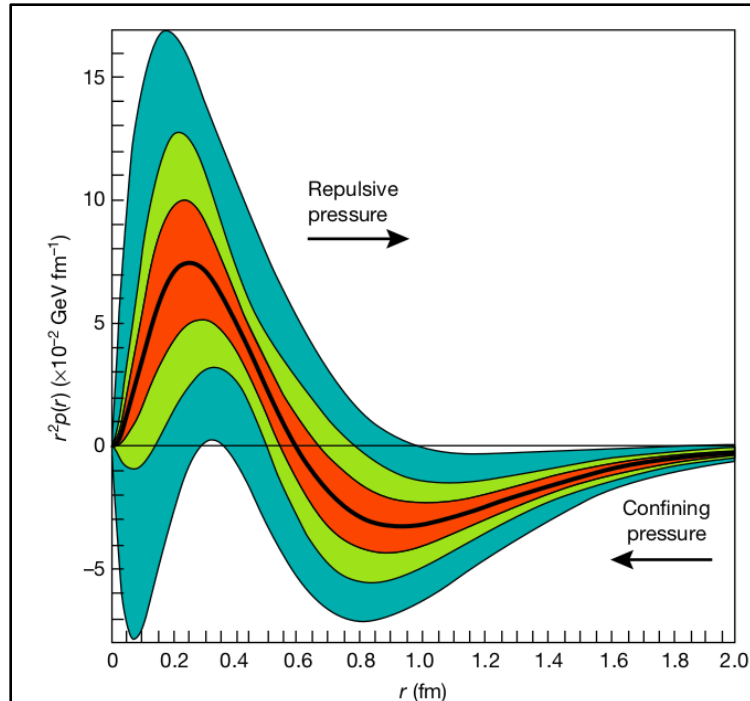
Table I. The global properties of the proton defined in terms of matrix elements of the conserved currents associated with respectively electromagnetic, weak, and gravitational interaction. Notice the weak currents include the partially conserved axial current, and  $g_A$  or  $g_p$  are strictly speaking defined in terms of transition matrix elements in the neutron  $\beta$ -decay or muon-capture. The values of the properties are from the particle data book [107] and [108] (for  $g_p$ ) except for the unknown  $D$ -term.

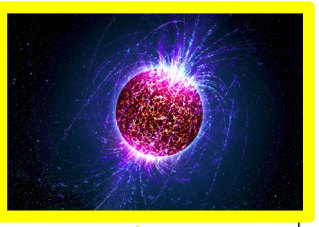
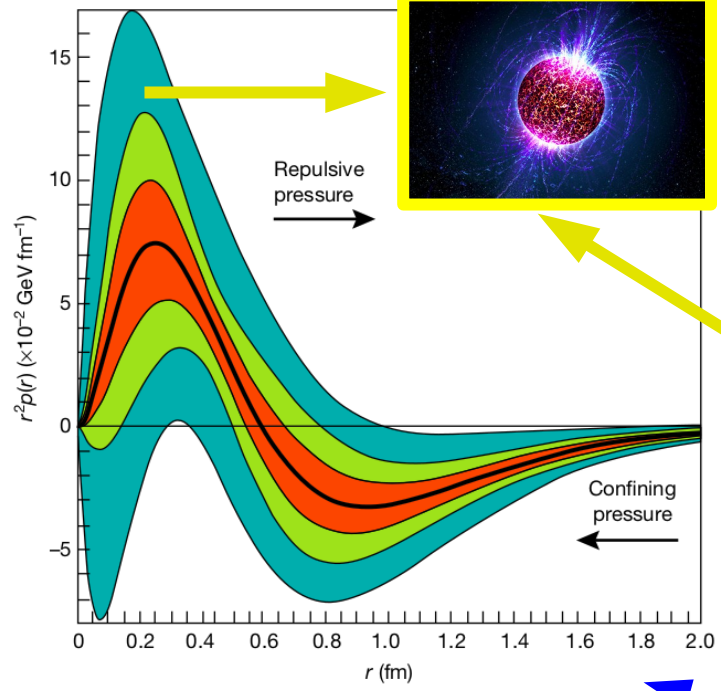
# Pressure (Microscopic expression) in the direction of Form Factor

## Energy Momentum Tensor

$$T^{ij}(\mathbf{r}) = \left( \frac{r^i r^j}{r^2} - \frac{1}{3} \delta^{ij} \right) s(r) + \delta^{ij} p(r).$$

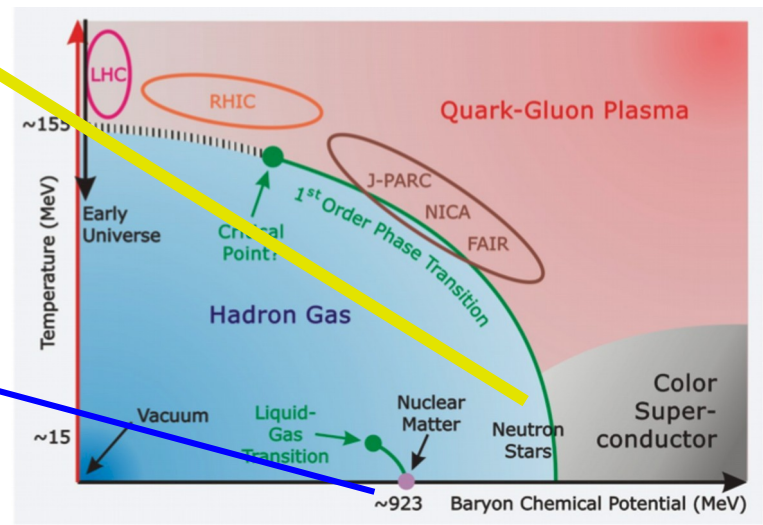
$$s(r) = -\frac{1}{4m} r \frac{d}{dr} \frac{1}{r} \frac{d}{dr} \tilde{D}(r), \quad p(r) = \frac{1}{6m} \frac{1}{r^2} \frac{d}{dr} r^2 \frac{d}{dr} \tilde{D}(r), \quad \tilde{D}(r) = \int \frac{d^3 \Delta}{(2\pi)^3} e^{-i\Delta \mathbf{r}} D(-\Delta^2).$$



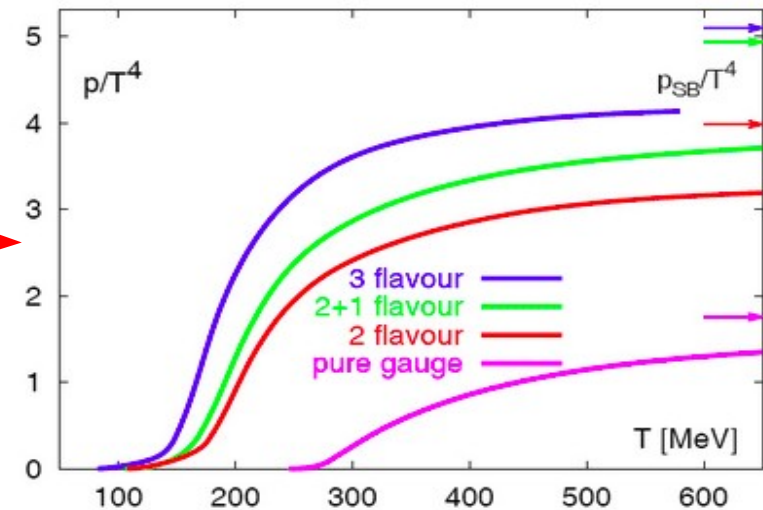
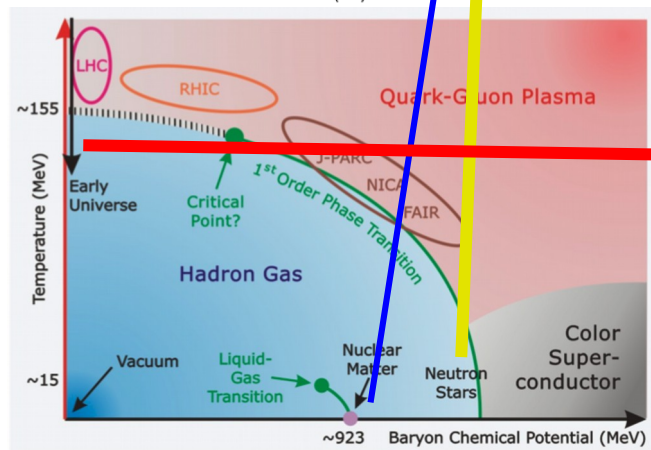
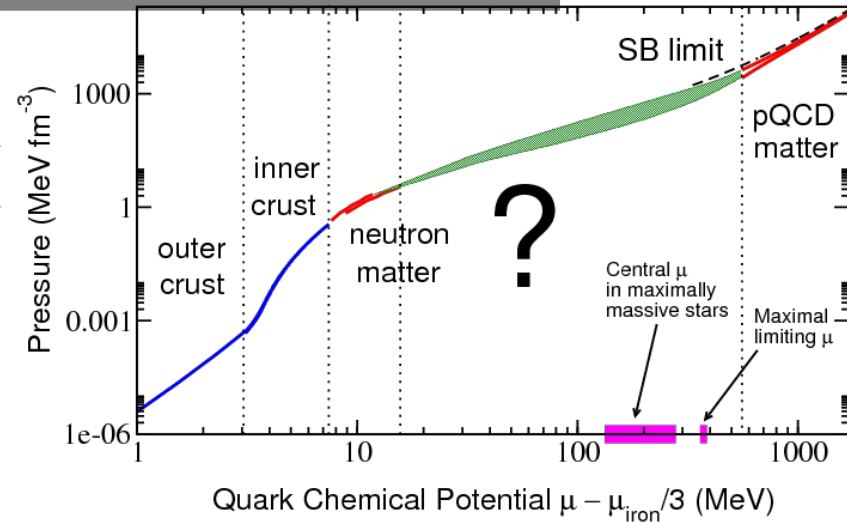
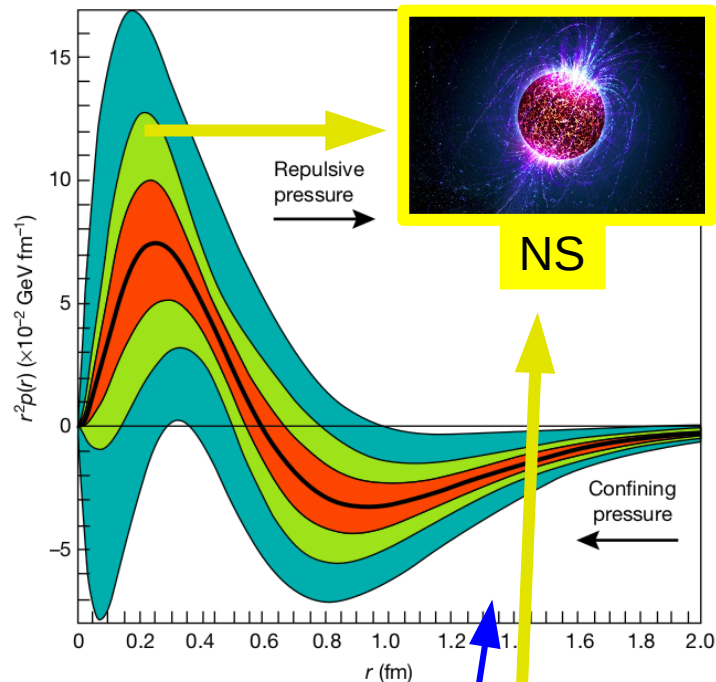


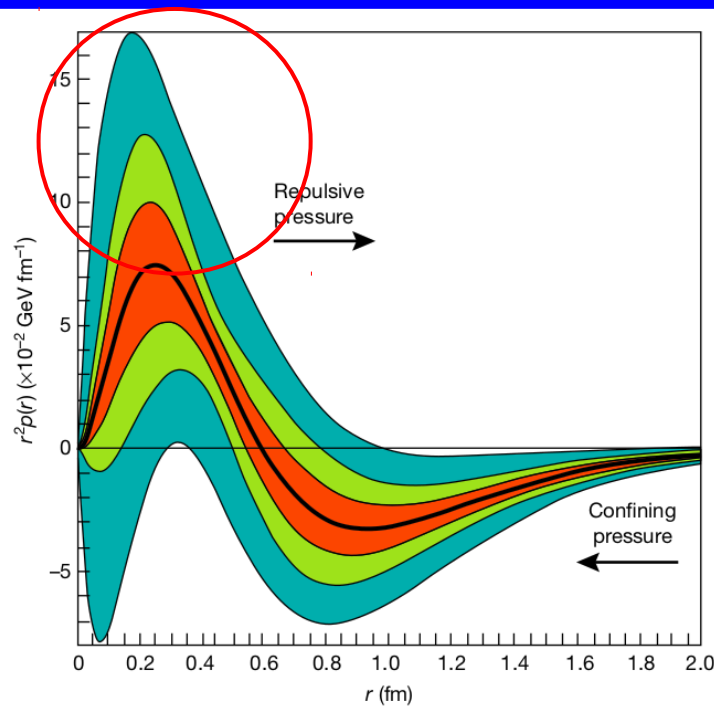
The average peak pressure near the centre is about  $10^{35}$  pascals, which exceeds the pressure estimated for the most densely packed known objects in the Universe, neutron stars.

$$1 \text{ GeV}^4 \rightarrow 2.0852 \times 10^{37} \text{ J m}^{-3}$$



$$1 \text{ GeV}^4 \rightarrow 2.0852 \times 10^{37} \text{ J m}^{-3}$$





$$P(r)V = \frac{g}{h^3} \int_V \int_0^\infty \left( \frac{pv}{3} \right) f(\epsilon_i) d\tau d^3p$$

$$P(r) = \frac{4\pi(2s+1)}{h^3} \int_0^{p_f(r)} \left( \frac{pv}{3} \right) p^2 dp$$

where  $g = 2s + 1$  is the degeneracy,  $p_f(r)$  is the fermi momentum obtained from the relation  $\epsilon_f(\mu) = \sqrt{p_f^2 c^2 + m^2 c^4}$  and  $d\tau$  is an infinitesimal volume element.

Velocity  $v$  is given by:

$$v = \frac{pc^2}{\sqrt{p^2 c^2 + m^2 c^4}}$$

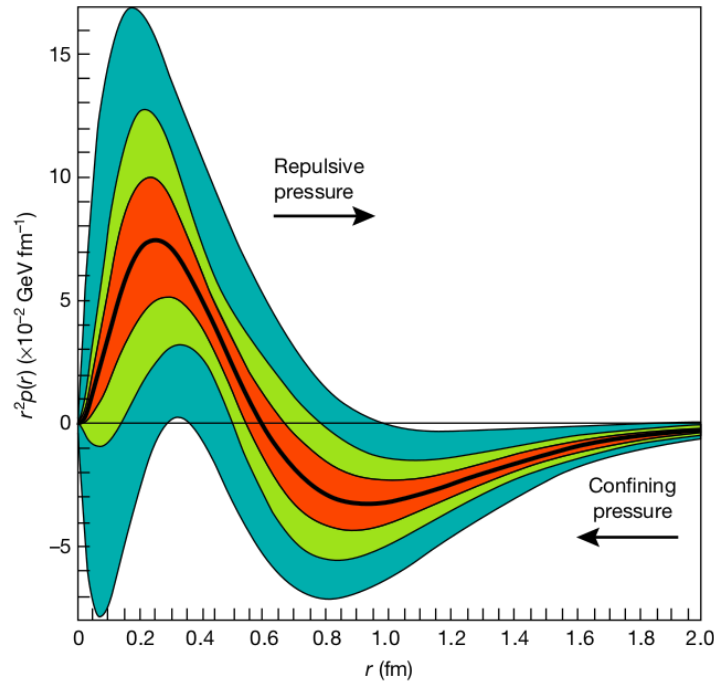


$$p(r)r^2 = 25.93 \times 10^{-2} (4 - 6.325r)r \exp(-6.325r/2).$$

Pressure inside Proton  
from the angle of  
Statistical Mechanics

$$P(r) = 1.047 \left( (2p_f^3 - 3p_f m^2) \sqrt{p_f^2 + m^2} - 3m^4 \left( \frac{p_f}{m} \right) \right)$$

# Pressure inside Proton from the angle of Classical Mechanics



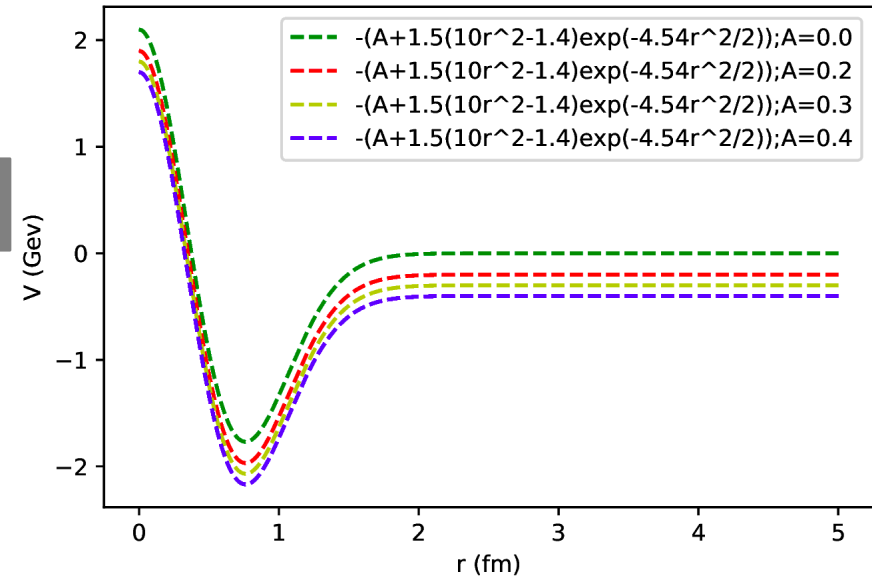
$$\tilde{V}(r, \theta, \phi) = -(A - B (br^2 + d) \exp(-fr^2))$$

$$\frac{\partial V(r)}{\partial r} = 2Br \cdot (bfr^2 + df - b) e^{-fr^2}$$

$$p(r)r^2 = 25.93 \times 10^{-2} (4 - 6.325r) r \exp(-6.325r/2)$$



potential inside the proton



# Summary

- We have attempted to estimate degenerate quark distribution inside Proton (motivated from the fact – peak pressure can denser than neutron star environment).
- Starting with this simple degenerate quark fluid motion, our future aim to project the formalism towards the standard Wigner function approach and also towards Kubo formalism approach with an expectation of non-zero dissipation!
- We have sketch the mechanical potential profile for the pressure distribution inside proton, which might be interesting to compare with our known form of potential like Yukawa, Coulomb, string type potentials.
- Our statistical and Classical mechanical version attempt for the pressure profile may be considered as a beginning steps from the understanding of heavy ion collision (HIC) direction, which may be eventually matured via mapping between EIC and HIC tools.

