

LHC constraints on monojet signatures from electroweakino DM and colored - superpartner decays

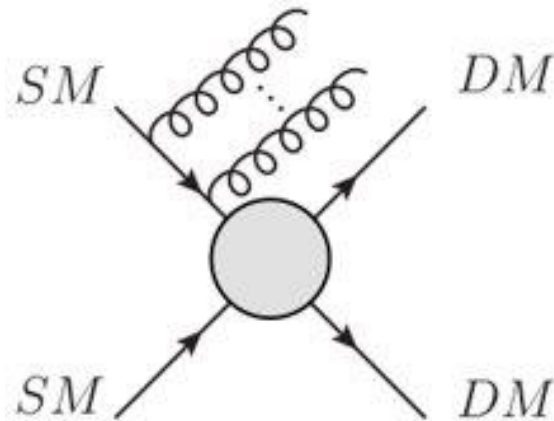
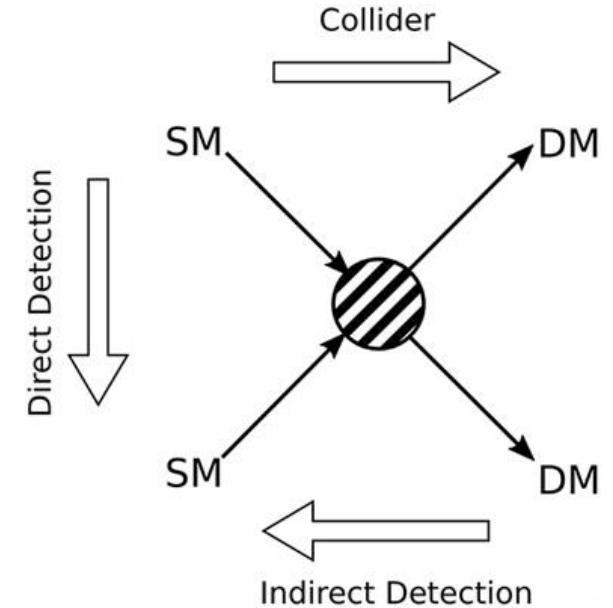
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[hep-ph]/2208.04342

7th LHC Reinterpretation Forum
CERN-December 2022



Supersymmetric DM

- DM existence strongly suggested by cosmological data.
- **Weakly interacting massive particle** fits well the picture.
- No direct detection and/or collider signal thus far.
- Supersymmetric partners of gauge and Higgs bosons are strong candidates but remain elusive.



Hadron colliders may produce DM particles in pairs, **associated with a few high p_T jets** originating from initial state QCD radiation

Two SUSY scenarios to re-visit:

- I. Light **higgsinos/gauginos** and rest of the spectrum decoupled. → Jets+MET signal originated from **Initial State Radiation**.
- II. Light **squarks/gluinos**. → Jets+MET signal originated from the **decay of coloured sparticles**.

Quick summary of chargino-neutralino sector:

$$\mathcal{L}_{\tilde{\chi}} = \overline{\tilde{\chi}_i^-} (\not{p} \delta_{ij} - P_L (U^* \textcolor{blue}{X} V^\dagger)_{ij} - P_R (V \textcolor{blue}{X}^\dagger U^T)_{ij}) \tilde{\chi}_j^- \\ + \frac{1}{2} \overline{\tilde{\chi}_i^0} (\not{p} \delta_{ij} - P_L (N^* \textcolor{green}{Y} N^\dagger)_{ij} - P_R (N \textcolor{green}{Y}^\dagger N^T)_{ij}) \tilde{\chi}_j^0$$

$$\textcolor{blue}{X} = \begin{pmatrix} \textcolor{blue}{M}_2 & \sqrt{2} M_W \sin \beta \\ \sqrt{2} M_W \cos \beta & \textcolor{red}{\mu} \end{pmatrix} \quad \text{diagonalised via} \quad \mathbf{M}_{\tilde{\chi}^+} = U^* \textcolor{blue}{X} V^\dagger$$

$$\textcolor{green}{Y} = \begin{pmatrix} \textcolor{blue}{M}_1 & 0 & -M_Z c_\beta s_W & M_Z s_\beta s_W \\ 0 & \textcolor{blue}{M}_2 & M_Z c_\beta c_W & -M_Z s_\beta c_W \\ -M_Z c_\beta s_W & M_Z c_\beta c_W & 0 & -\textcolor{red}{\mu} \\ M_Z s_\beta s_W & -M_Z s_\beta c_W & -\textcolor{red}{\mu} & 0 \end{pmatrix} \quad \text{diagonalised via} \quad \mathbf{M}_{\tilde{\chi}^0} = N^* \textcolor{green}{Y} N^\dagger$$

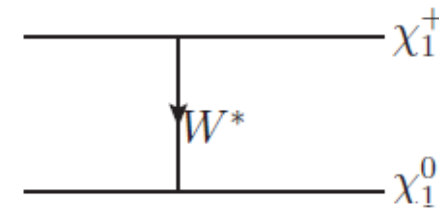
Small mixing expected if there is a hierarchy between M_1 , M_2 and μ and/or particles much heavier than the EW scale.

We concentrate in three cases:

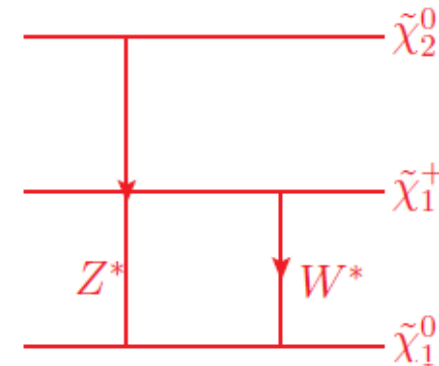
bingo-wino: almost mass degenerate winos
and bino LSP



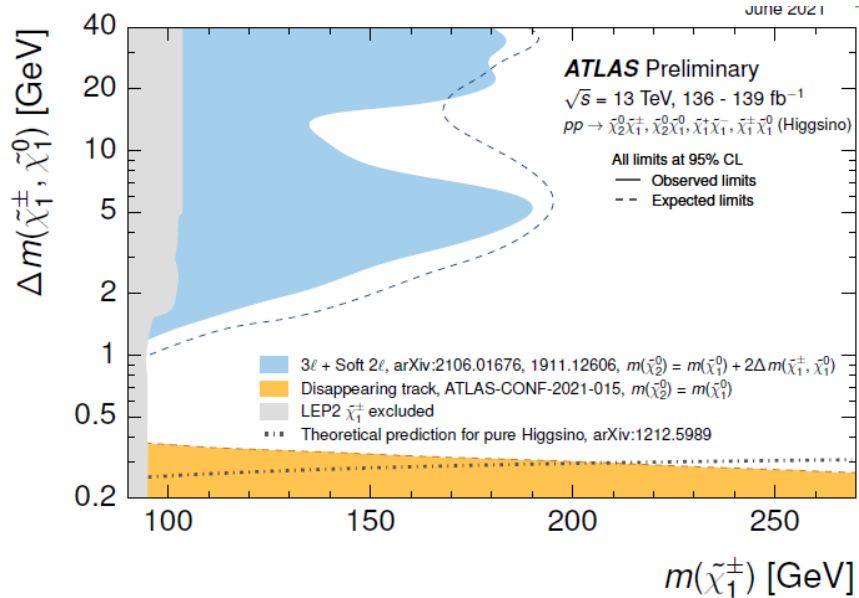
wino LSP: $M_2 \ll M_1, \mu$, two
quasi-degenerate states: χ_1^0, χ_1^\pm



higgsino LSP, $\mu \ll M_1, M_2$, three
quasi-degenerate states: $\tilde{\chi}_1^0, \tilde{\chi}_1^\pm, \tilde{\chi}_2^0$



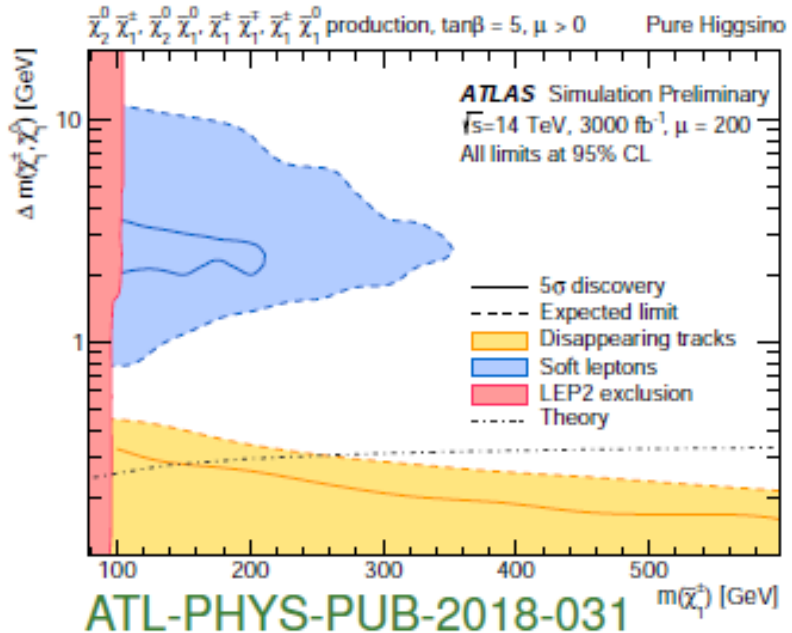
At the LHC this scenarios has been constrained focusing on:



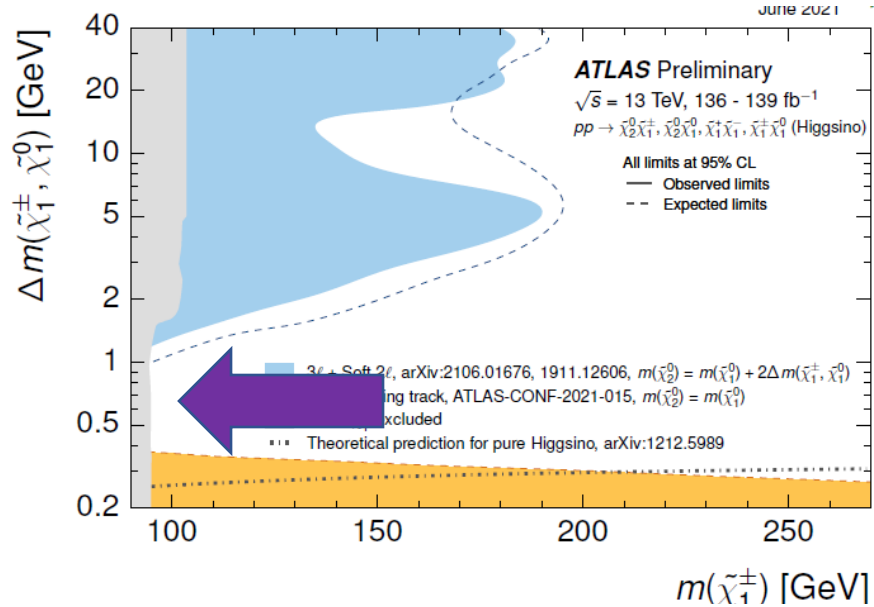
- Disappearing tracks: for sufficiently small mass gap, heavier states are long-lived.

- Soft leptons: For a mass difference $\gtrsim 1 \text{ GeV}$ look for soft decay products.

- Long-standing limit at $\sim 100 \text{ GeV}$ from LEP

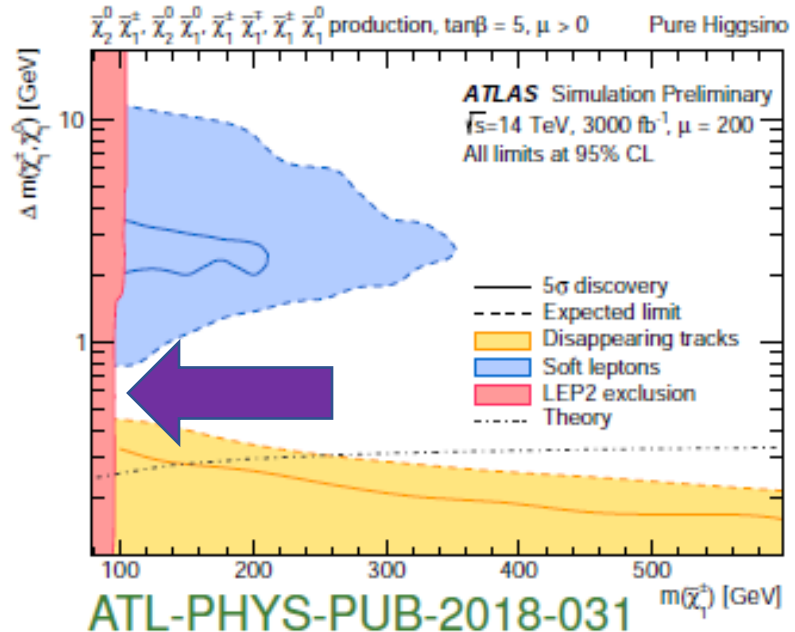


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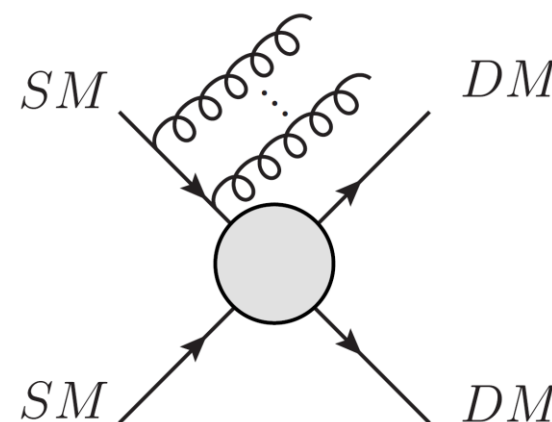


- Long-standing limit at ~ 100 GeV from LEP.

GAP between the two LHC searches. Use Mono/few-jets searches on this region.

Monojet searches from ATLAS and CMS are not sensitive (yet) to *electroweakino* DM. More than one jet emitted is possible, thus *more-than-one-jet* searches may be used also.

- We recast with CheckMATE a general search for squarks and gluinos, [arXiv:2010.14293](#), in total 70 signal regions.
- Basic (preselection) signal requirements:
 - no electrons or muons.
 - 2–6 jets
 - large missing energy > 300 GeV
 - hard leading jet $p_T > 200$ GeV
 - large effective mass > 800 GeV



- Some overlap of the final states with “mono”-jet.
- We focus on bins with the largest sensitivity (originally intended for squark pair production):

$$2\text{--}3 \text{ jets, } p_T^{\text{jet1}}, p_T^{\text{jet2}} > 250 \text{ GeV}$$

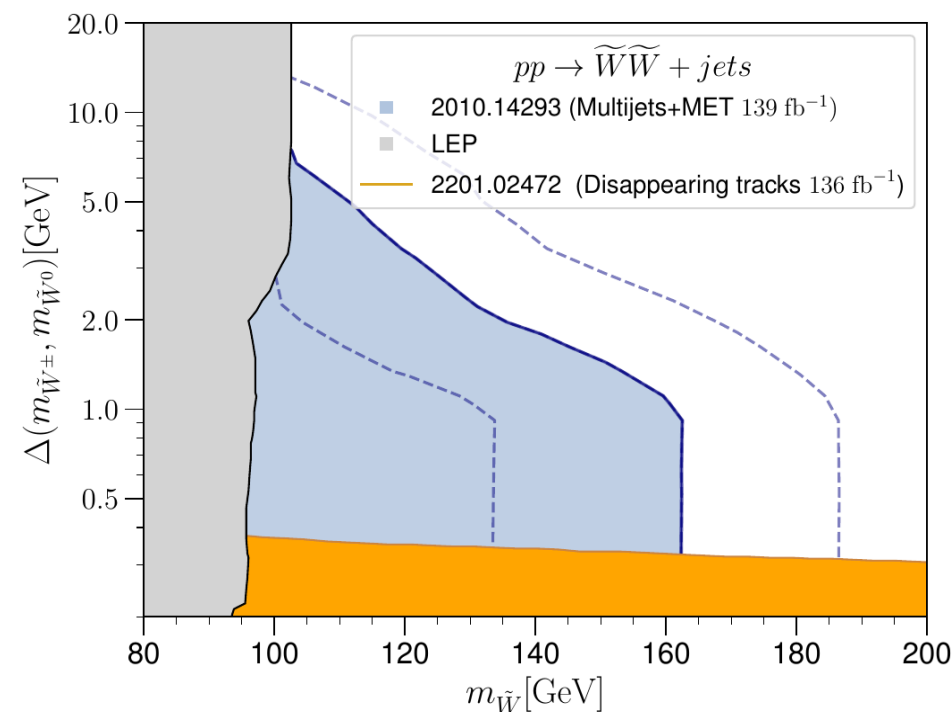
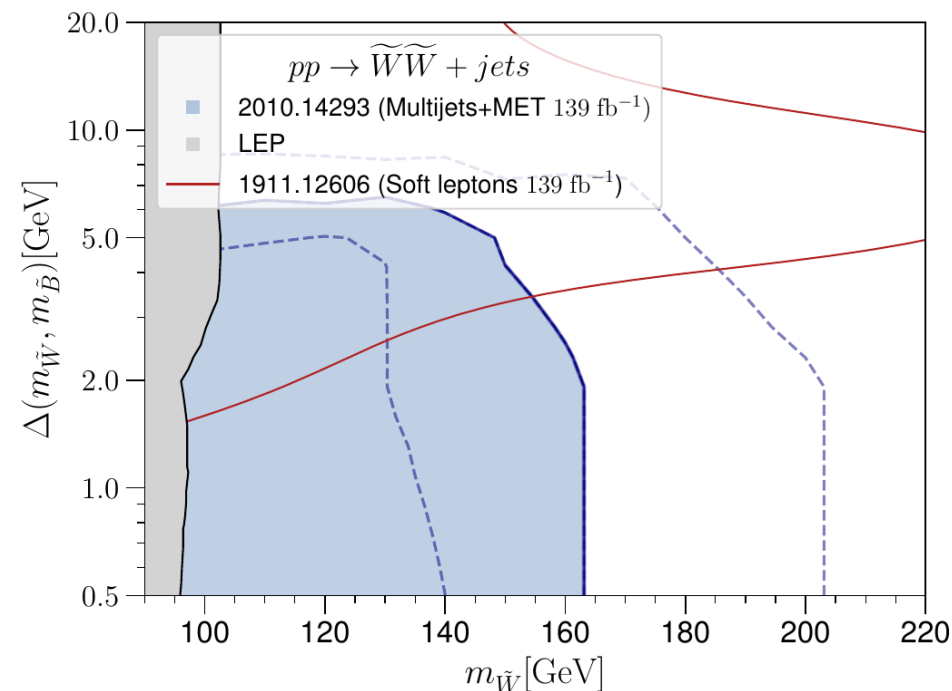
$$\text{effective mass} > 1600 \text{ GeV}$$

$$E_T^{\text{miss}} / \sqrt{H_T} > 16 \sqrt{\text{GeV}}$$

- Multi-bin fit using Histfitter.

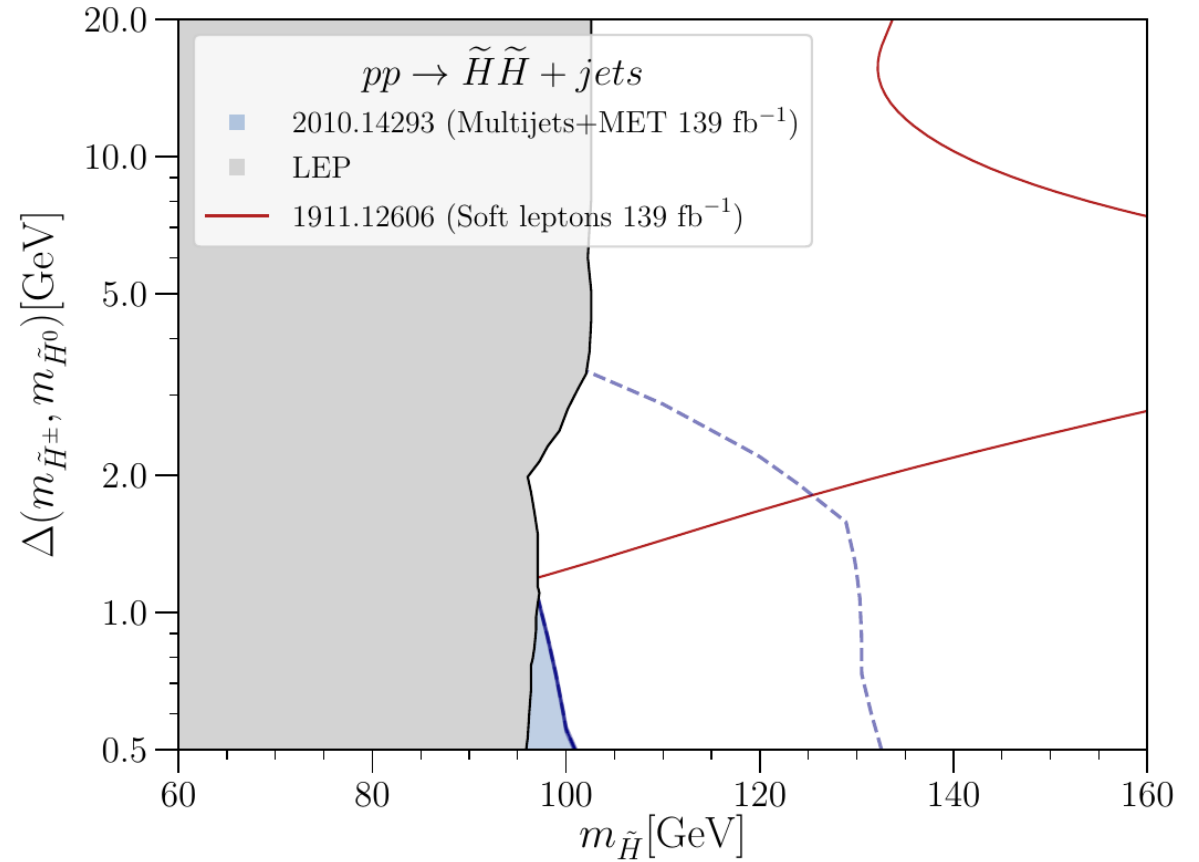
Limits over Wino

- bino-wino model
- $\widetilde{W}^\pm \rightarrow \widetilde{B}^0 W^*, \widetilde{W}^0 \rightarrow \widetilde{B}^0 Z^*$
- comparison with ATLAS exclusion (red line)
- wino model
- $\widetilde{W}^\pm \rightarrow \widetilde{W}^0 W^*$
- the new exclusion on top of LEP and long-lived charged wino limits



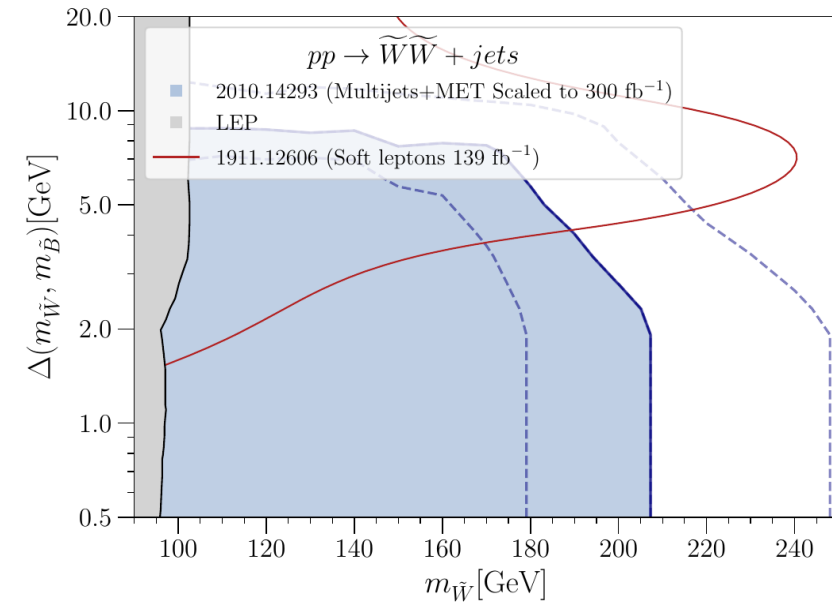
Limits over Higgsino

- higgsino model
- $pp \rightarrow \tilde{H}^\pm \tilde{H}_{1,2}^0, \tilde{H}^+ \tilde{H}^-, \tilde{H}_1^0 \tilde{H}_2^0$
- $\tilde{H}^\pm \rightarrow \tilde{H}_1^0 W^*, \tilde{H}_2^0 \rightarrow \tilde{H}_1^0 Z^*$

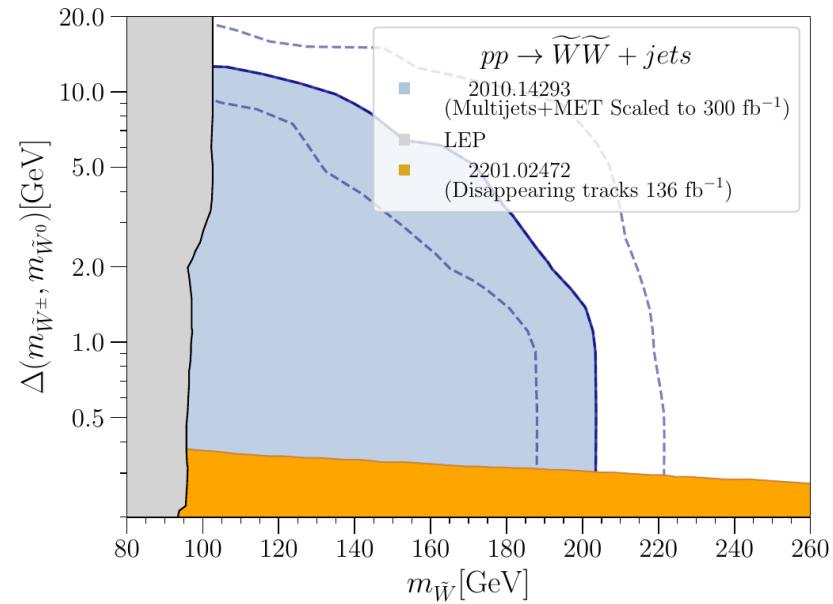


Prospects for Run 3

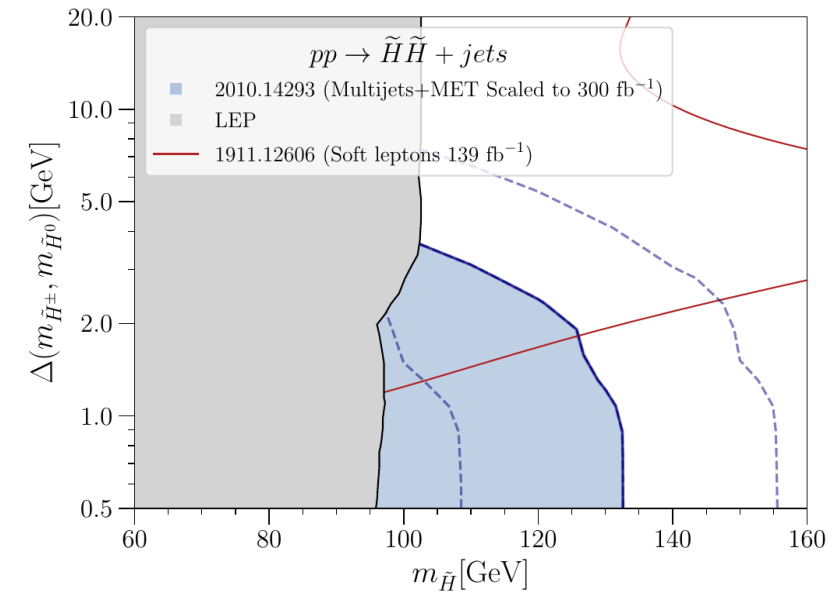
Wino-bino model



Wino model



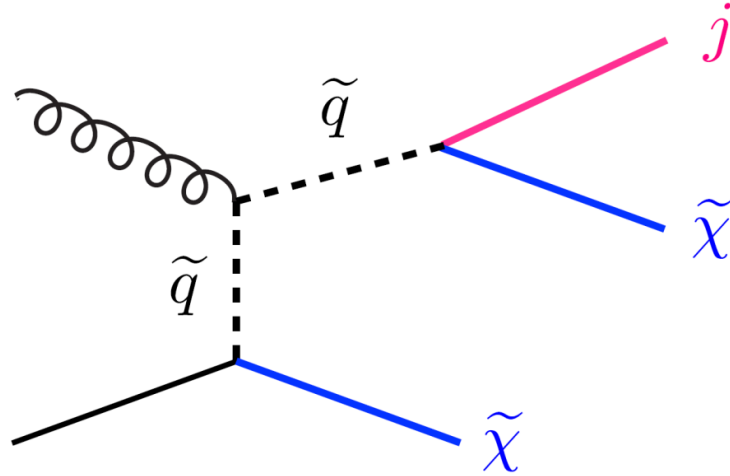
Higgsino



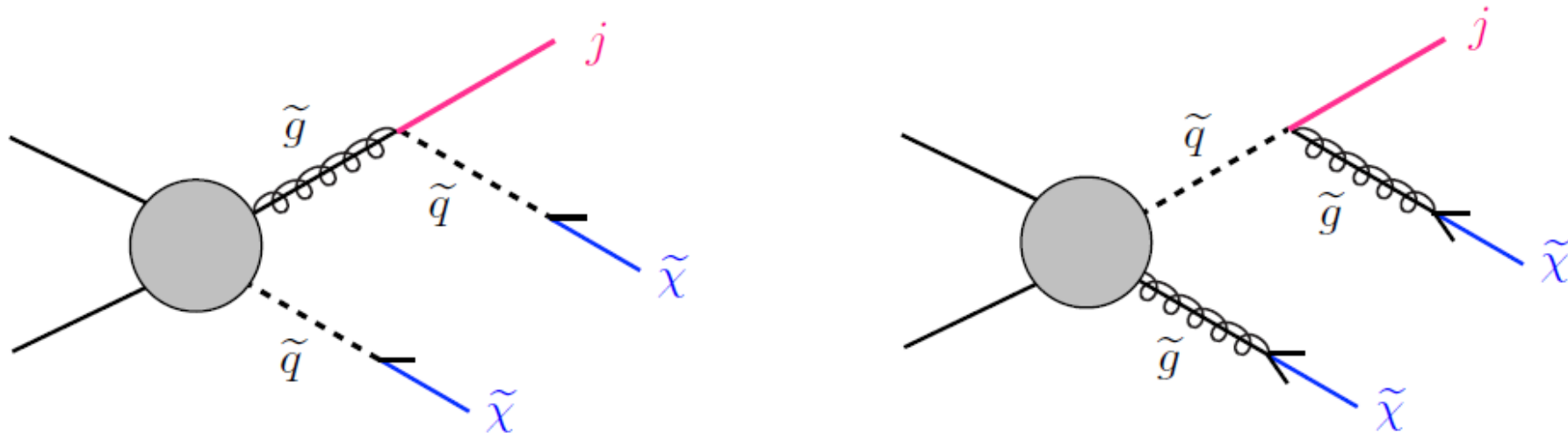
Two SUSY scenarios to re-visit:

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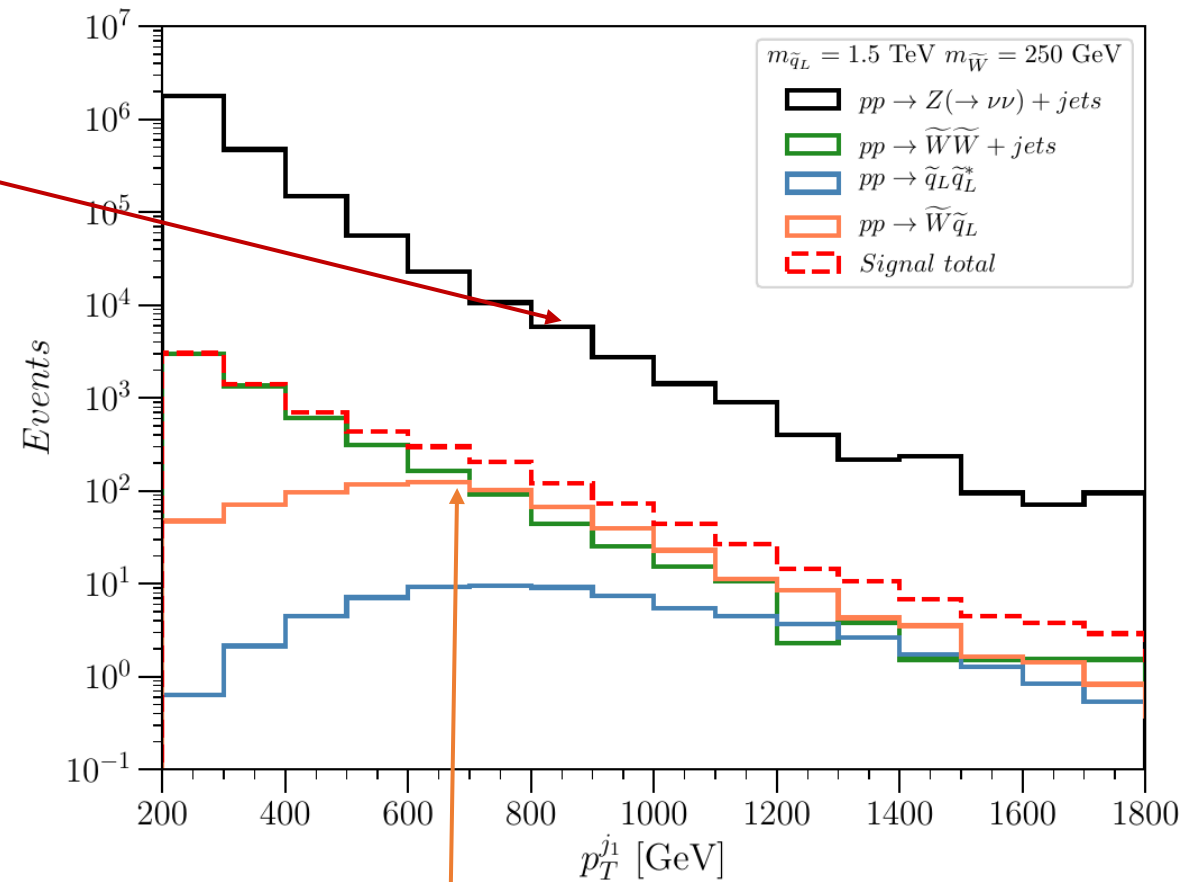
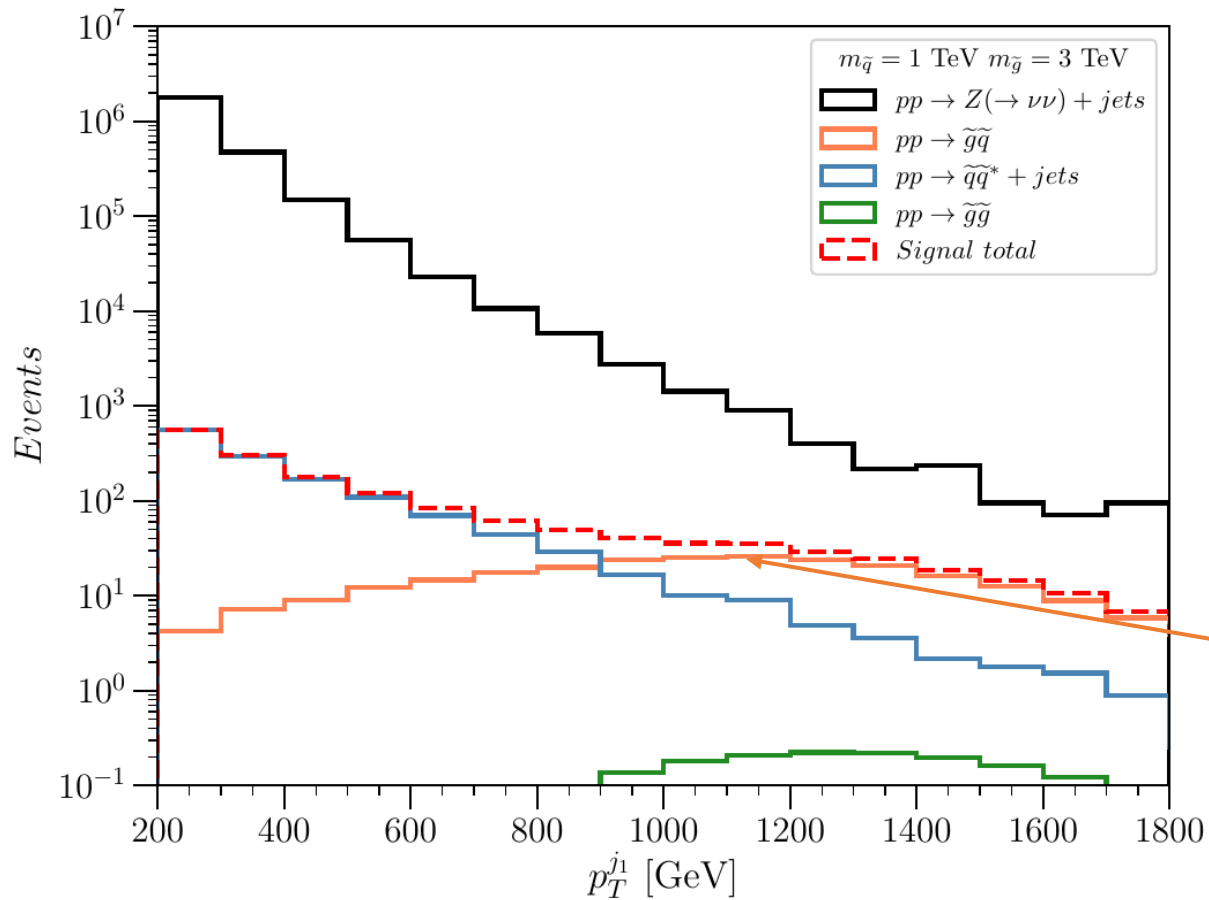
- I. **Large mass hierarchy between the squarks and the wino (LSP).**
The **associated squark-wino production** produces a monojet-like signature.



- II. **Lighter of squark/gluino is nearly degenerated with LSP.**
The **associated gluino-squark production** leads to a monojet signal.



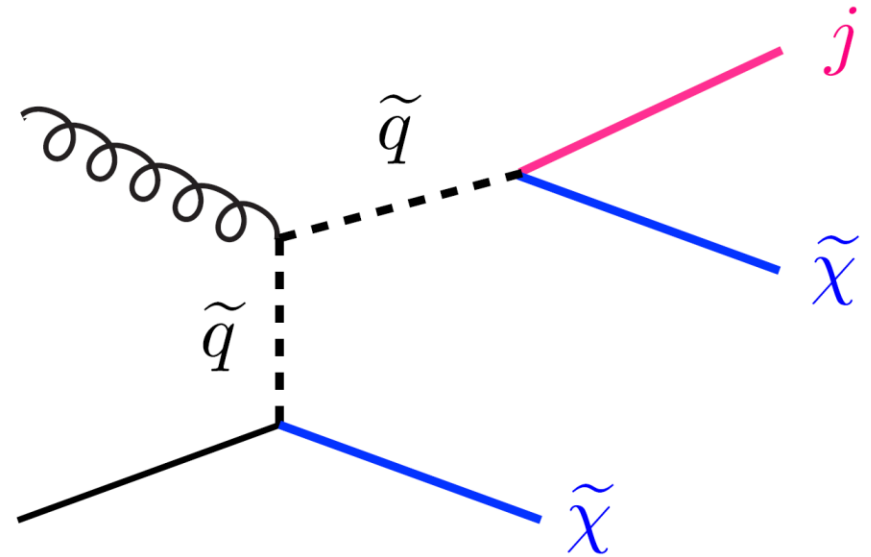
QCD radiation has a **monotonically falling spectrum** and the signal acceptance becomes low once **tight pT cuts** are imposed on the jets.



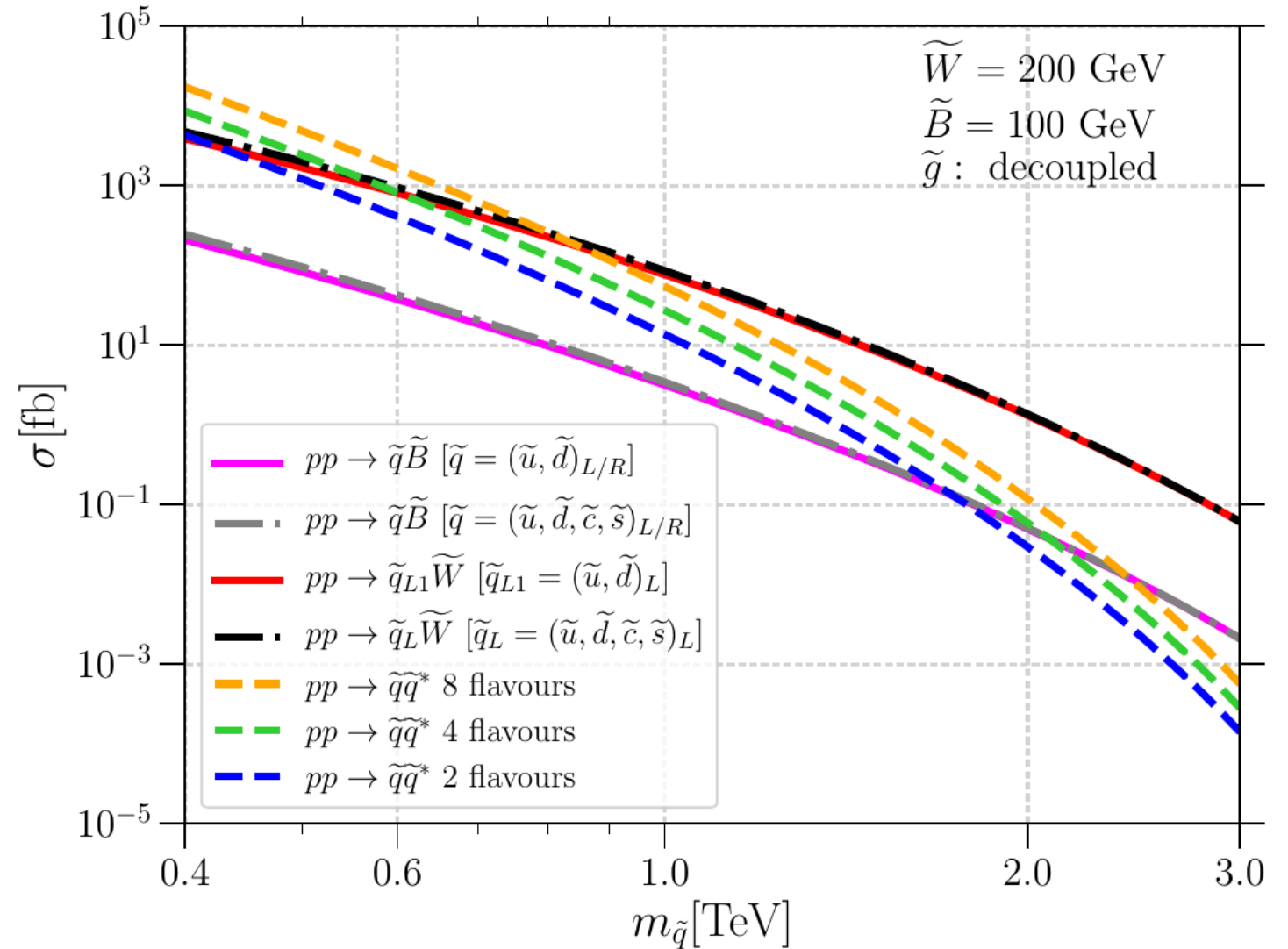
Unlike ISR, in these processes the q-jet has the **energy scale characterised by the mass differences** between the squark and neutralino in (I) and the squark and gluino in (II).

Mono-jets from squark-wino production

- An isolated energetic jet could result from the associated production of squarks together with electroweakinos.
- Specifically sensitive to 1st generation of squarks.
- Not relevant for higgsino-like electroweakinos.



- This signal is generally disregarded in front of strong production of squarks.
- For wino-like neutralino/chargino and **squark mass ~ 1 TeV**, the cross section is competitive with squark pair production ($m_W = 200$ GeV).
- For bino-like neutralino, the cross section becomes relevant for **squark mass > 2 TeV** ($m_B = 100$ GeV).



We use again the search for squarks and gluinos : [arXiv:2010.14293](https://arxiv.org/abs/2010.14293).

Focus on the regions with the largest sensitivity: MB-C-2 and MB-SSd-2

jets + E_T^{miss} (139 fb⁻¹) MB-C-2
e, μ veto
$E_T^{\text{miss}} > 300$
$p_T^{j1} > 600, \eta^{j1} < 2.8$
$p_T^{j2} > 50, \eta^{j2} < 2.8$
$N_j(p_T > 50, \eta < 2.8) \leq 3$
$\Delta\phi(\text{jet}, \mathbf{p}_T^{\text{miss}}) > 0.4$
$(E_T^{\text{miss}}/\sqrt{H_T}, m_{\text{eff}})$ binned

jets + E_T^{miss} (139 fb⁻¹) MB-SSd-2
e, μ veto
$E_T^{\text{miss}} > 300$
$p_T^{j1} > 250, \eta^{j1} < 2$
$p_T^{j2} > 250, \eta^{j2} < 2$
$N_j(p_T > 50, \eta < 2.8) \leq 3$
$\Delta\phi(\text{jet}, \mathbf{p}_T^{\text{miss}}) > 0.8$
$(E_T^{\text{miss}}/\sqrt{H_T}, m_{\text{eff}})$ binned

Improvement of the limit based on the combination of the signals:

- Squark pair production.

$$pp \rightarrow \tilde{q}\tilde{q}$$

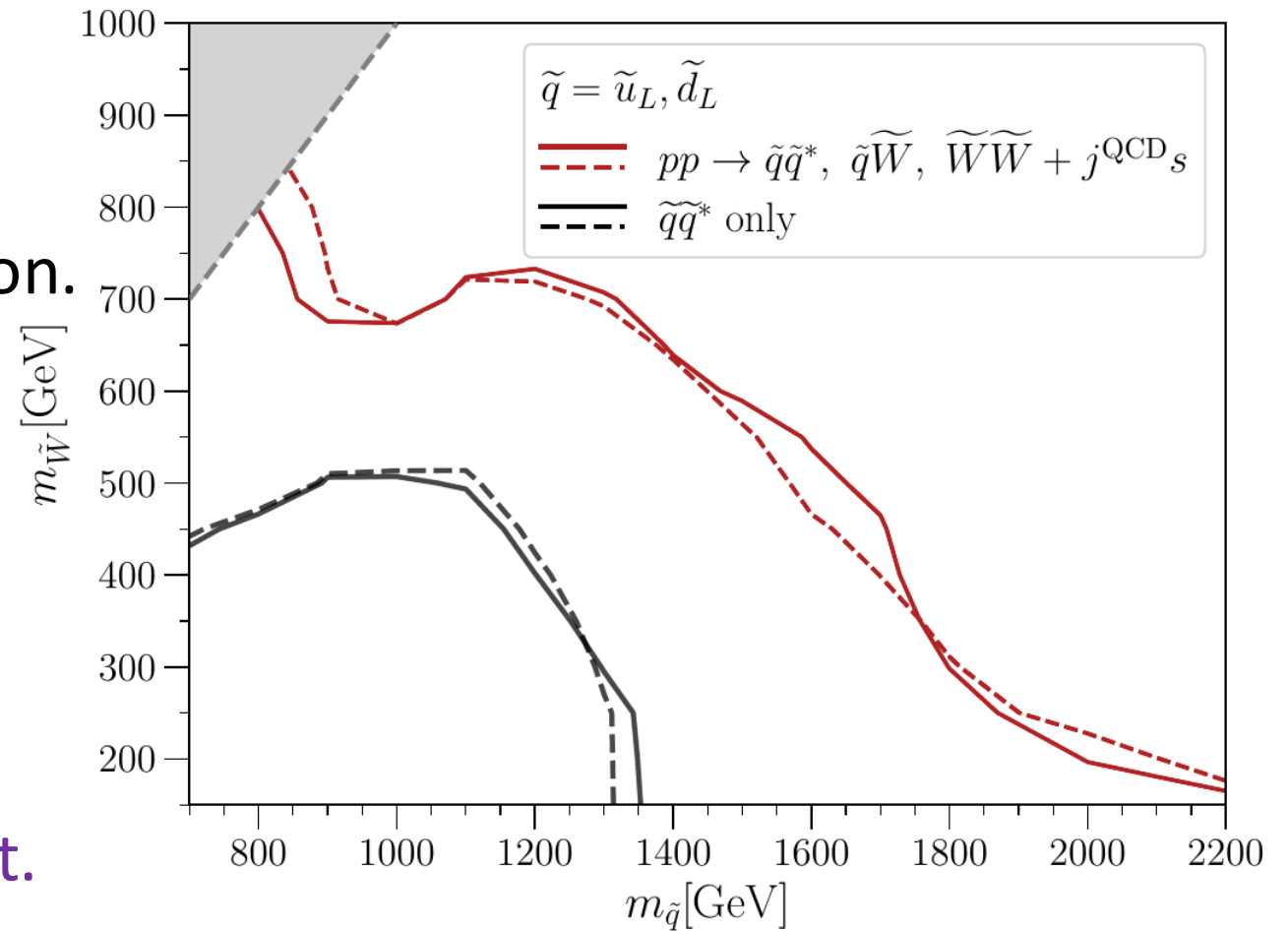
- Squark-wino associated production.

$$pp \rightarrow \tilde{q}_L \tilde{W}$$

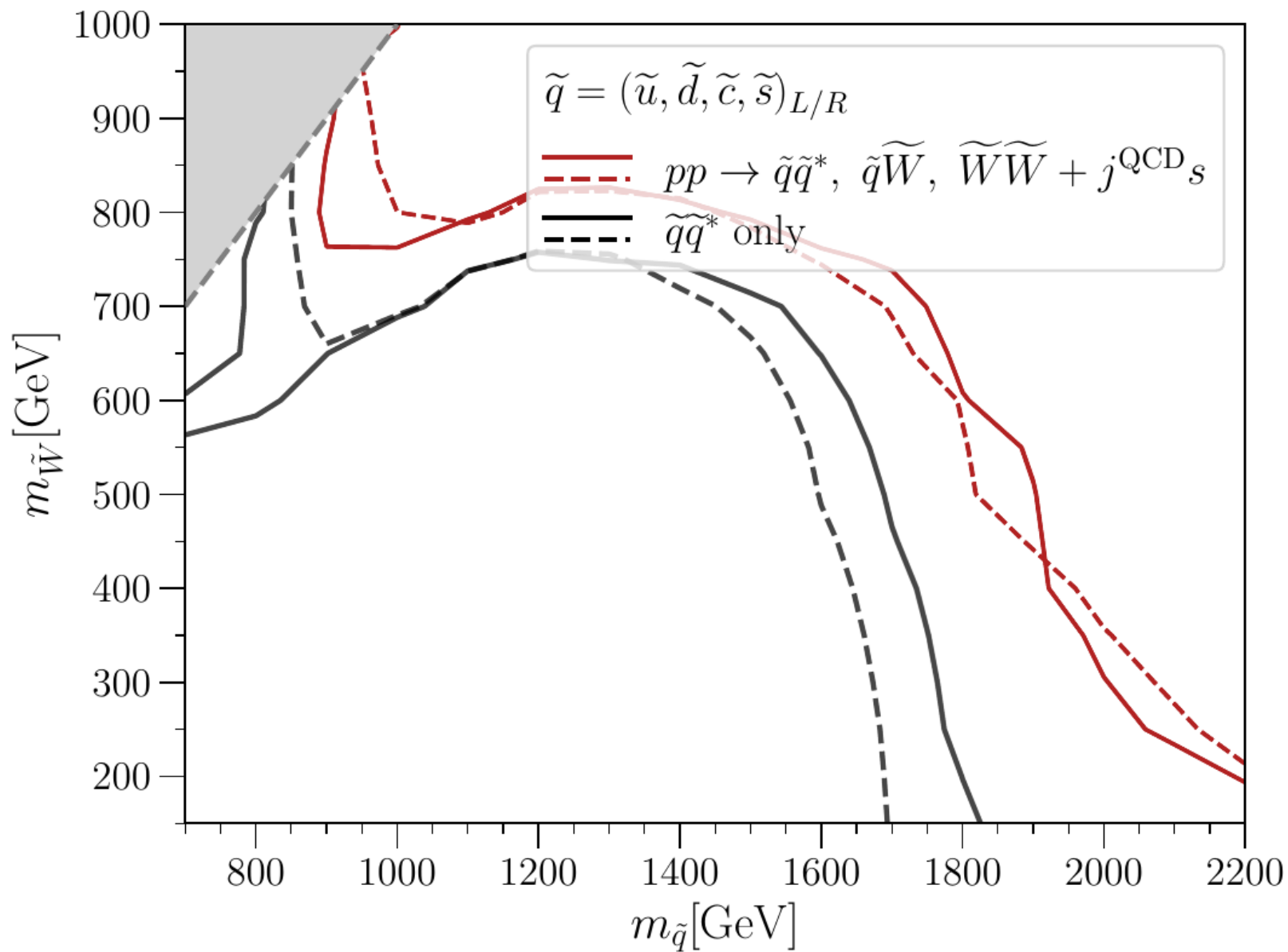
- Wino pair production + ISR jets.

$$pp \rightarrow \tilde{W}\tilde{W} + jets$$

Only 1st generation left squarks light.

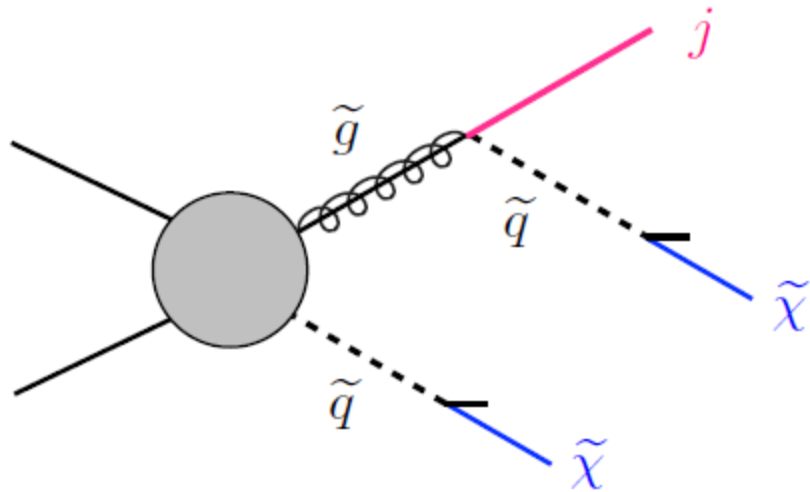


Also if squarks 8-fold degenerated.



Mono-jets from gluino-squark production

If Squark/gluino is degenerate with neutralino (**coannihilation!**) we have monojets from squark-gluino associated production.

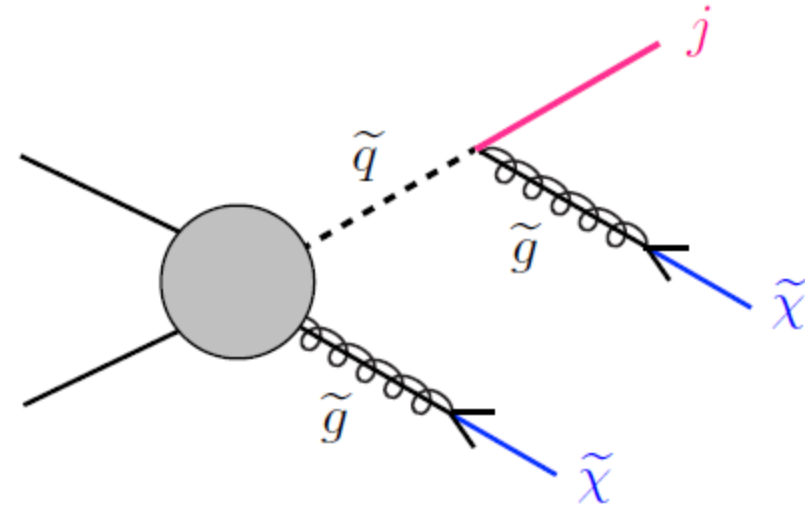


Case $m_{\tilde{g}} > m_{\tilde{q}}$:

$$pp \rightarrow \tilde{g}\tilde{q}, \quad \tilde{g} \rightarrow q\tilde{\chi} \quad \dots \text{ (a1)}$$

$$pp \rightarrow \tilde{g}\tilde{g}, \quad \tilde{g}\tilde{g} \rightarrow (q\tilde{q})(q\tilde{q}) \quad \dots \text{ (a2)}$$

$$pp \rightarrow \tilde{q}\tilde{q} + \text{ISR} \quad \dots \text{ (a3)}$$



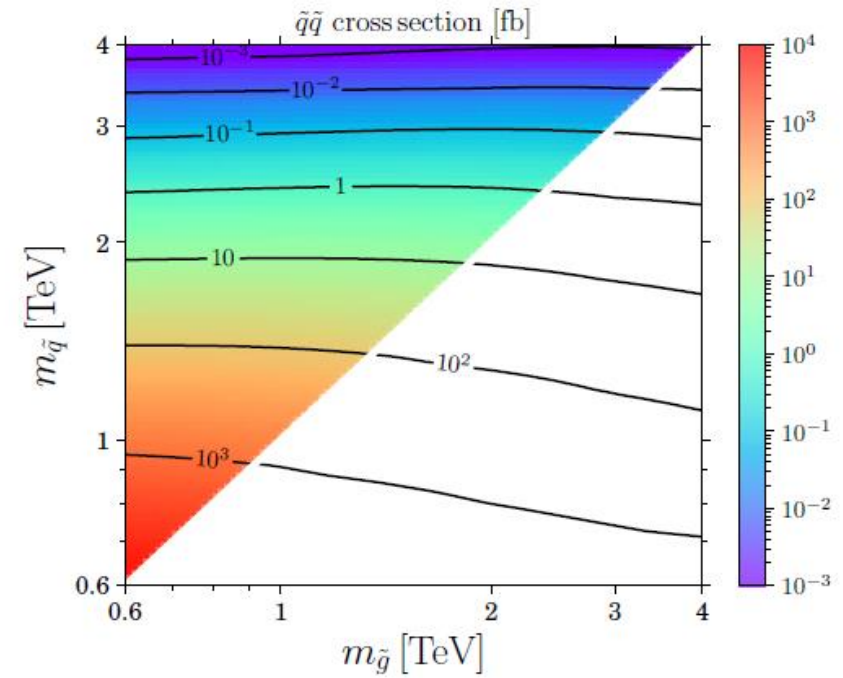
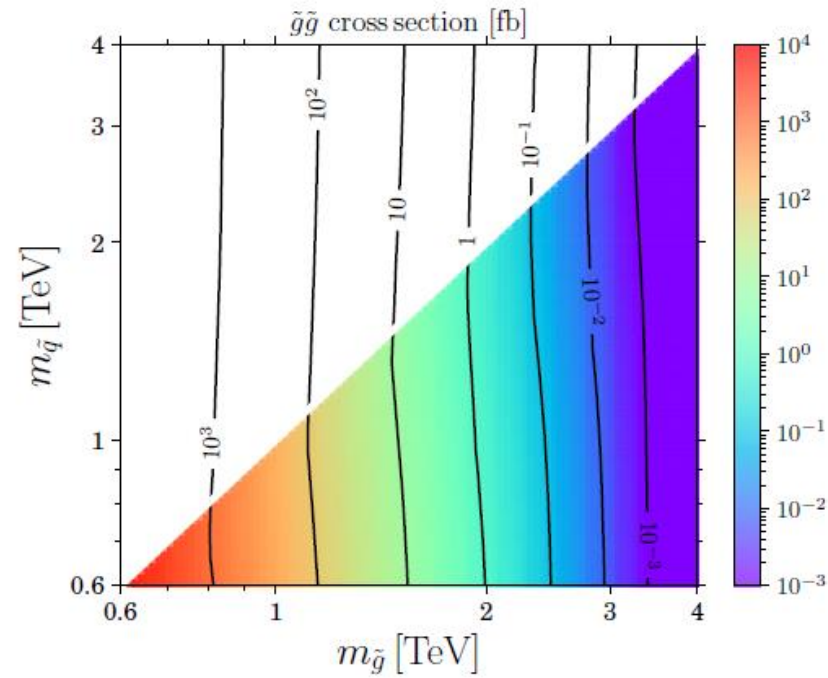
Case $m_{\tilde{q}} > m_{\tilde{g}}$:

$$pp \rightarrow \tilde{g}\tilde{q}, \quad \tilde{q} \rightarrow q\tilde{g} \quad \dots \text{ (b1)}$$

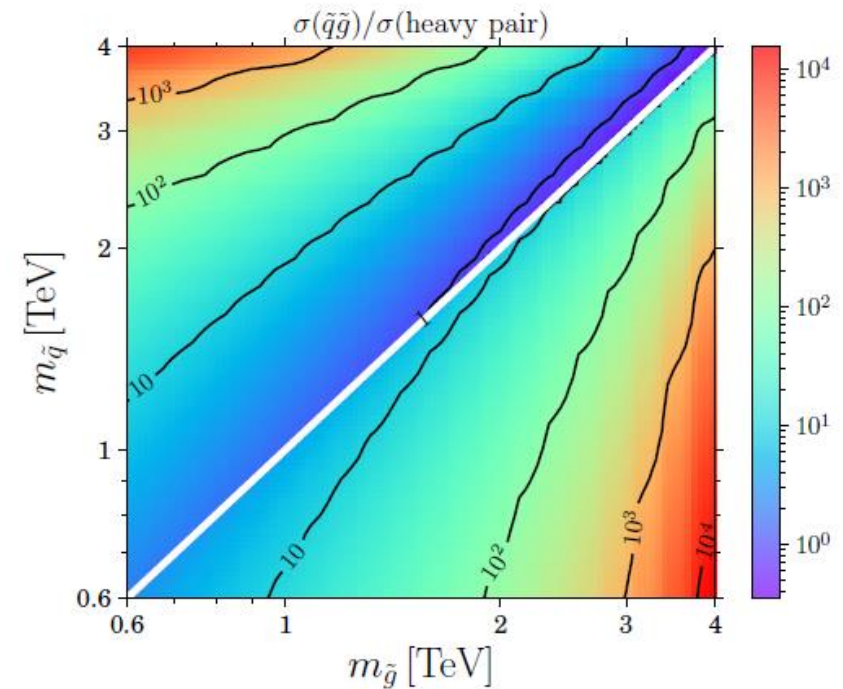
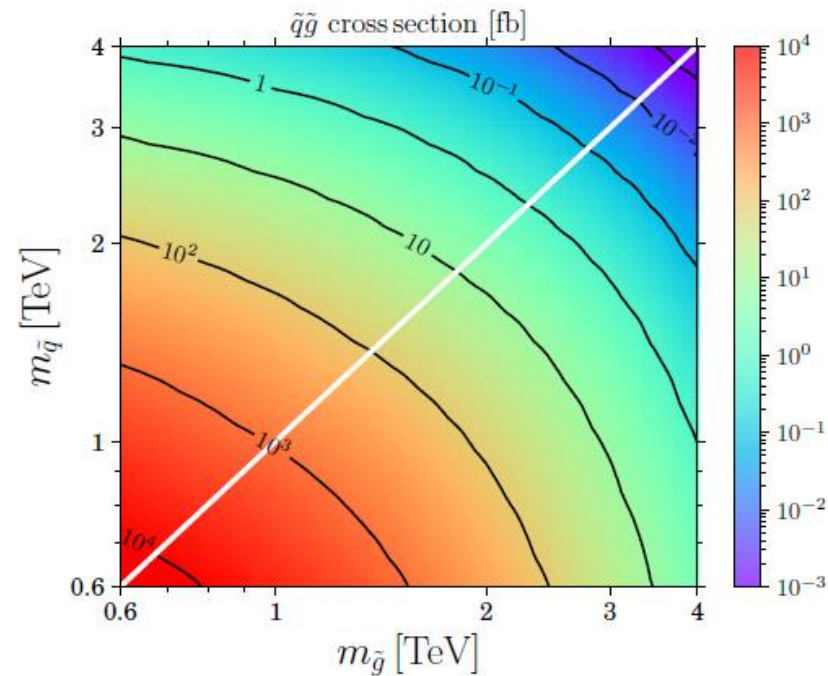
$$pp \rightarrow \tilde{q}\tilde{q}, \quad \tilde{q}\tilde{q} \rightarrow (q\tilde{g})(q\tilde{g}) \quad \dots \text{ (b2)}$$

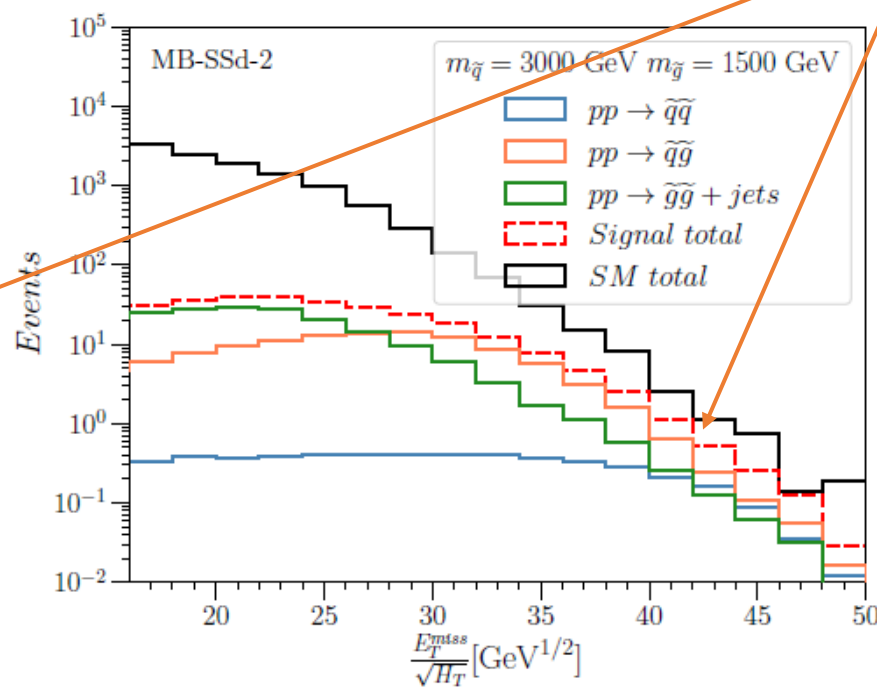
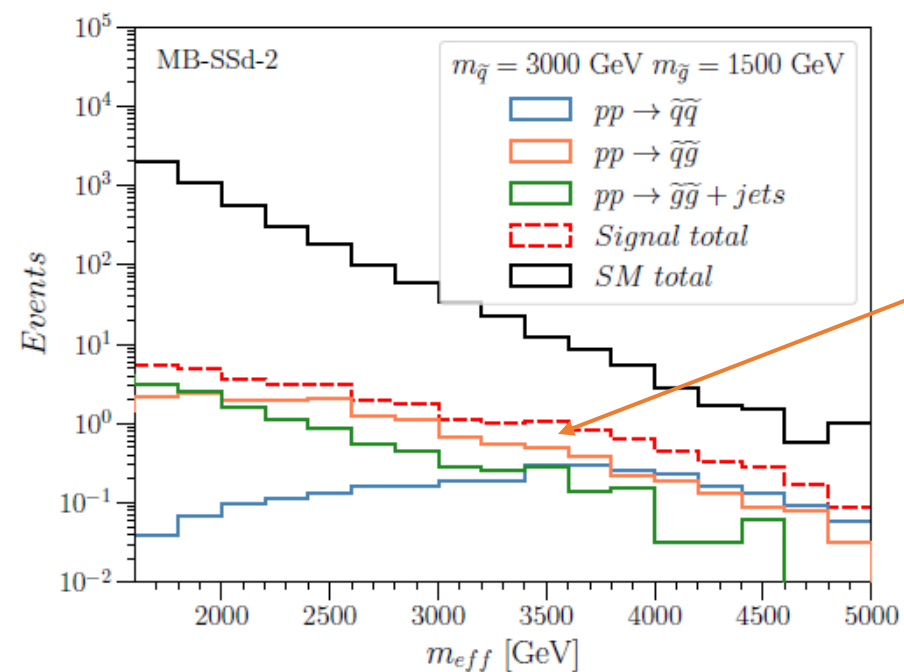
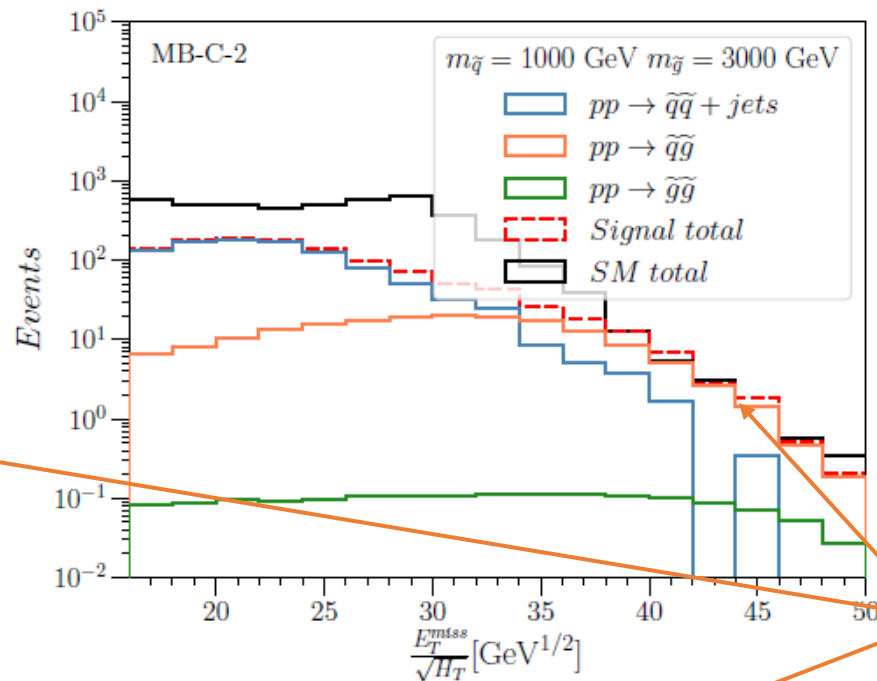
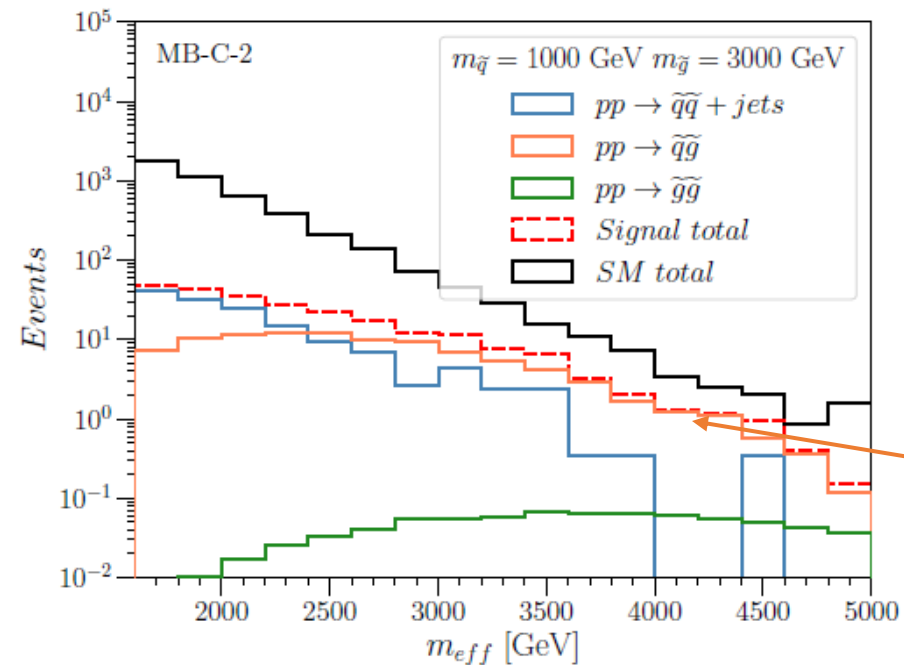
$$pp \rightarrow \tilde{g}\tilde{g} + \text{ISR} \quad \dots \text{ (b3)}$$

The cross section of the **associated production** is almost always larger than that of the (a2) and (b2) processes



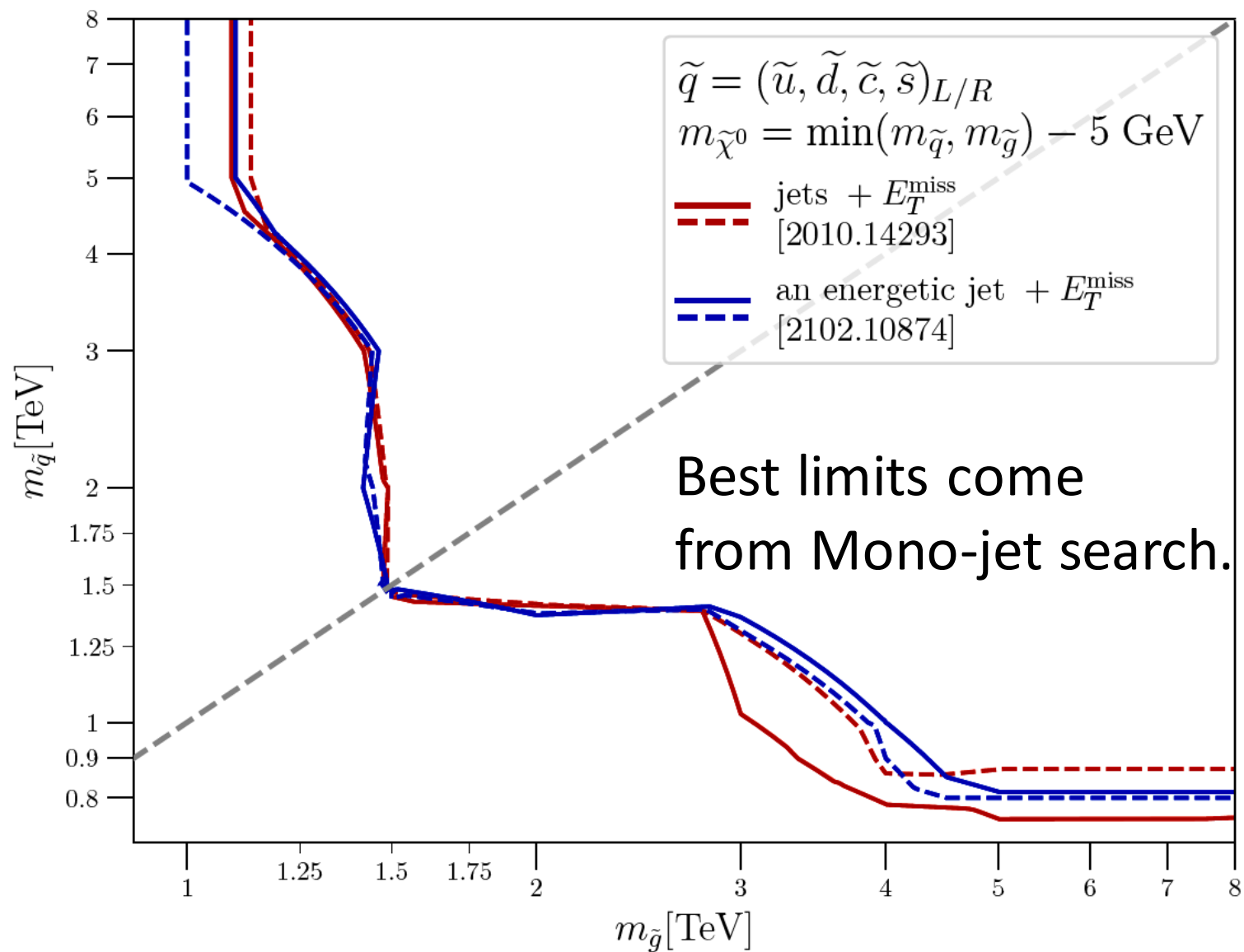
Relative rate of the associated production is enhanced particularly in the **hierarchical mass regions**.





The gluino-squark associated production dominates the signal distributions for higher values of the discriminant variables. $E_T^{miss} / \sqrt{H_T}$ and m_{eff} .

We compare the limits obtained from recasted Monojet ATLAS (blue) search vs recasted gluino/squark ATLAS search (red).



Summary

- Initial state radiation can give a handle on challenging bits of LSP parameter space.
- Squark search outperforms dedicated monojet analysis for electroweakinos. New constraints closing the gap in (model independent) wino exclusion. Higgsinos more difficult but with some promise. HL prospects to be seen.
- Two distinct SUSY scenarios where a single high p_T jet originates from SUSY particle pair production, to which mono- and di-jet event selections are particularly sensitive:
 1. A squark-wino simplified model \rightarrow non negligible contribution of the associated squark-wino production. Negligible in the squark-bino model with current luminosity but need to be reconsidered for the HL-LHC.
 2. A gluino-squark simplified model with nearly degenerate bino LSP. \rightarrow We derived current exclusion limit on the gluino-squark mass plane



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Understanding the Early Universe:
interplay of theory and collider experiments

Joint research project between the University of Warsaw & University of Bergen