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## From large to small systems



Pb—Pb









Phys. Lett. B 724 (2013) 213



pp





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Phys. Rev. Lett. 123, 212002 (2019)









## **Correlation functions**

### **Balance function**

Measure of balancing charges

$$R_2^{\alpha\beta} = \frac{\rho_2^{\alpha\beta}}{\rho_1^{\alpha}\rho_1^{\beta}} - 1$$

$$B^{\alpha\beta} = \frac{1}{2} \left\{ \rho_1^{\beta^-} \left[ R_2^{\alpha^+\beta^-} - R_2^{\alpha^-\beta^-} \right] + \rho_1^{\beta^+} \left[ R_2^{\alpha^-\beta^+} - R_2^{\alpha^+\beta^+} \right] \right\}$$



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- Measure of balancing charges
- $BF_{\pi\pi}$  narrow in central collision **coalescence**?









## **Balance function: Identified hadrons**



- Measure of balancing charges
- $\bullet$   $BF_{pp}$  wider than acceptance, no dependence on multiplicity • Early stage production?







## Balance function: Identified hadrons



• Pairing fractions not dependent on centrality, except  $I^{\pi\pi} \Rightarrow$  quantitative characterisation of the hadronisation of the QGP





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## **Balance function: Charged hadrons**



Narrowing of width of BF towards central collisions Radial flow, late particle production (coalescence) Similar behaviour in p—Pb collisions

VS





# VS

## **Balance function: Charged hadrons**



- Similar trends in p—Pb and Pb—Pb collisions



Consistent with the delayed hadronisation, support collectivity in p—Pb collisions <sup>•</sup> Models cannot describe data perfectly, though provide better description in  $\Delta \phi$ 



















### **G**<sub>2</sub>



 $\Delta \eta$  CI - different for all three systems • Pb—Pb - increase 24% - viscous effects of long-lived QGP with small  $\eta/s$ 









### $\mathbf{G}_2$



 $\left[ \frac{\int_{\Omega} p_{\mathrm{T},1} p_{\mathrm{T},2} \rho_2(p_1, p_2) \mathrm{d} p_{\mathrm{T},1} \mathrm{d} p_{\mathrm{T},2}}{\int_{\Omega} \rho_1(p_1) \mathrm{d} p_{\mathrm{T},1} \int_{\Omega} \rho_2(p_2) \mathrm{d} p_{\mathrm{T},2}} - \langle p_{\mathrm{T},1} \rangle \langle p_{\mathrm{T},2} \rangle \right]$  $G_2 = \frac{1}{\langle p_{\mathrm{T},1} \rangle \langle p_{\mathrm{T},2} \rangle} \, \, \mathbf{I}$ 

- $\Delta \eta$  CI - different for all three systems
  - Pb—Pb increase 24% viscous effects of long-lived QGP with small  $\eta/s$
  - pp slight decrease, p—Pb slight increase
    - Too small for viscous forces to equilibrate?
    - Different explanations?











### $\mathbf{G2}$



- $\Delta \eta$  CI different for all three systems
- Pb—Pb increase 24% viscous effects of long-lived QGP with small  $\eta/s$
- pp slight decrease, p—Pb slight increase
  - Too small for viscous forces to equilibrate?
  - Different explanations
    - Models without collective effects do not describe data











## Anisotropic flow

### **Initial spatial anisotropy**



### Final anisotropy in momentum space



medium (low  $p_{\rm T}$ )

- Reflects the conversion of the initial-state spatial anisotropy into final-state anisotropies in momentum space
- Anisotropy in distribution of final-state particles:

$$\frac{\mathrm{d}N}{\mathrm{d}\varphi} \propto 1 + 2\sum_{n=1}^{\infty} v_n \cos n(\varphi - \Psi_n)$$
$$v_n = \langle \cos(n(\varphi - \psi_n)) \rangle$$

- Initial conditions and transport properties of the created
  - Initial geometry affects energy loss of hard hadrons (high  $p_{\rm T}$ )







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- First measurement of Fourier coefficients up to  $p_{\rm T}$  = 200 GeV/c
- Compatible with previous measurements
  - Decreased uncertainties









Non-zero  $v_2$  at high  $p_T$ 

Increasing towards semicentral collisions

• Hard hadrons influenced by initial geometry







## $v_n$ up to high $p_T$ in Pb—Pb



- Non-zero  $v_2$  at high  $p_T$ 
  - Increasing towards semicentral collisions
  - Hard hadrons influenced by initial geometry
- $v_3$  at high  $p_T$  compatible with zero
  - Hard hadrons not influenced by initial fluctuations









Fourier coefficients of dijets measured with two-particle correlations





## v<sub>n</sub> of dijets (Pb—Pb)



- Fourier coefficients of dijets measured with two-particle correlations
- Non-zero  $V_2$ :
  - More jets observed coplanar with event plane
    - $^{\circ}$  Less energy loss  $\rightarrow$  higher chance to pass selection criteria







- Fourier coefficients of dijets measured with two-particle correlations
- Non-zero  $V_2$ :
  - More jets observed coplanar with event plane
    - Less energy loss  $\rightarrow$  higher chance to pass selection criteria
- $v_3$ ,  $v_4$  compatible with zero:

• The fluctuations in the initial state **do not impact** the azimuthal distributions of dijets





## v<sub>2</sub> of particles in jets (p–Pb, Pb–Pb)





### • Non-zero jet $v_2$ in p—Pb and Pb—Pb collisions • Smaller magnitude than inclusive $V_2$

- No dependence on  $p_{T}^{assoc}$ 
  - At high  $p_{\rm T}$  similar magnitude as in Pb—Pb
- $v_2$  driven by the non-equilibrium anisotropic parton escape mechanism





## v2 of particles in jets (pp)



•  $h^{UE}$  separated by  $\Delta \eta > 1$  from jet with  $p_T > 15$  GeV/c • h<sup>J</sup> constituents of jets





• Study of influence of jets on inclusive  $v_2$  and jet  $v_2$  - origin of  $v_n$  in pp?



11

## v2 of particles in jets (pp)



No multiplicity dependence



• Presence of a jet with  $p_T > 15$  GeV/c does not influence the  $v_2$  of  $h^{UE}$ 

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11

## v2 of particles in jets (pp)



- - No multiplicity dependence
- $v_2$  of h<sup>J</sup> compatible with zero





• Presence of a jet with  $p_T > 15$  GeV/c does not influence the  $v_2$  of h<sup>UE</sup>

### • The inclusive $v_2$ is not driven by jet fragmentation, but rather by bulk The collective system is too small to influence jets - no energy loss

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11





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### Limit of collectivity in small systems



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### Long-range two-particle correlations measured up to the smallest multiplicities









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### Limit of collectivity in small systems



### Integrated ridge yield as a function of multiplicity calculated with great precision









• At  $N_{ch} > 30$  compatible with CMS measurement

## Limit of collectivity in small systems



Integrated ridge yield as a function of multiplicity calculated with great precision









• At  $N_{ch} > 30$  compatible with CMS measurement • At lower multiplicities - increased precision

## Limit of collectivity in small systems



Integrated ridge yield as a function of multiplicity calculated with great precision









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• Comparison with  $e^+e^-$  collisions •  $Y^{pp} > Y^{ee}$  for  $\langle N_{ch} \rangle \approx 15$  with  $3\sigma$ First quantitative comparison between pp and  $e^+e^-$  collisions New insight to processes contributing to the long-range ridge

### Limit of collectivity in small systems







### • Large systems

- Deconfined medium with small  $\eta/s$
- Late production of pions via coalescence, hint of early production of protons •  $v_2$  of jets and jet particles induced by path length dependent energy loss

### • Small systems

- Collectivity supported by narrowing of the peak width of BF and  $G_2$  correlation functions of low  $p_{\rm T}$  hadrons and non-zero  $v_2$ 
  - Viscous forces do not have time to equilibrate the system
- $v_2$  in pp collisions not driven and not influenced by jet fragmentation
- Significant ridge yield in pp down to  $\langle N_{ch} \rangle \approx 10$ , larger than in  $e^+e^-$





Thank you for your attention!





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### More results

## PID flow in Pb-Pb



First  $p_{\rm T}$ -differential  $v_2$  measurements using four-particle cumulants for identified particles

- In the intermediate  $p_T$  range,  $v_2$ {4}for baryons is larger than that for mesons by about 50%
- $F(v_2)$  an apparent splitting between baryons and mesons for centrality above
- $30\% \Rightarrow$  a significant role for final-state interactions in developing this observable



## Backup



